

# Towards charm physics with stabilised Wilson Fermions

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# Motivation

## Charm physics with Stabilised Wilson Fermions

- ▶ first phenomenological studies with Stabilised Wilson Fermions
- ▶ planned: charm physics
  - ▶ meson spectrum
  - ▶ leptonic decays
  - ▶ semi-leptonic decays
  - ▶ baryons
- ▶ ToDo:
  - ▶ Step 1: fix the charm mass  $m_c$ 
    - ▶ investigate relative cutoff effects from different fixings (LCPs)
    - ▶ via charmed mesons
    - ▶ via RGI quark mass ratio
  - ▶ Step 2: calculations @ fixed charm mass

## Reminder: Stabilised Wilson Fermions

- ▶ even-odd preconditioning  $\Rightarrow$  exp. variant of  $O(a)$  improved Wilson fermions [Francis et al., 2019, 2022]

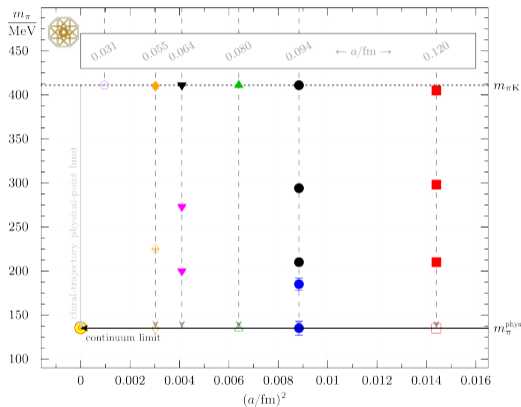
$$D_{ee} + D_{oo} = (m_0 + 4) + c_{\text{sw}} \frac{i}{4} \sigma_{\mu\nu} F_{\mu\nu} \rightarrow (m_0 + 4) \exp \left( \frac{c_{\text{sw}}}{(m_0 + 4)} \cdot \frac{i}{4} \sigma_{\mu\nu} F_{\mu\nu} \right)$$

- ▶ includes SMD algorithm and other stabilising measures
- ▶ potentially smaller cut-off effects
- ▶ could be beneficial for heavy quarks

# Our setup

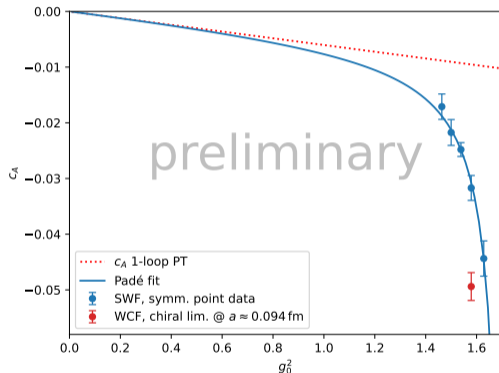
- ▶ Stabilised Wilson Fermions
- ▶ openLAT initiative<sup>a</sup> ensembles at  $N_f = 3$  flavour-sym. point
  - ↳ A. Francis, 4.8., 09:00
  - ▶  $\kappa_{\text{crit}}$  not yet known
  - ▶ expect no impact on charm observables
- ▶ periodic boundary conditions
- ▶ four lattice spacings
  - $a \approx 0.120, 0.095, 0.080, 0.064$  fm

<sup>a</sup><https://openlat1.gitlab.io>



# Our tools

- ▶  $O(a)$  improvement of axial-vector current
 
$$(A_\mu^a)_I(x) = A_\mu^a(x) + c_A \partial_\mu P^a(x)$$
  - ▶ done for SWF in the SF [Fritzsch et al., 2022] [Della Morte et al., 2005; Bulava et al., 2015]
  - ▶ also directly at sym. point
- ▶ software stack: Hadrons [Portelli et al., 2023]/GRID [Boyle et al., 2015]
  - ▶ Notation: mass parameter  $M_0 = \frac{1}{2\kappa} - 4$
- ▶ Gamma-method [Wolff, 2003] (pyerrors [Joswig et al., 2022])



# The renormalised quark mass

- ▶  $O(a)$  improved renormalised quark mass

$$m_{i,R} = Z_m \left\{ \left[ m_{q,i} + (r_m - 1) \frac{\text{Tr} [M_q]}{N_f} \right] + aB_i \right\} + O(a^2)$$

$$B_i = b_m m_{q,i}^2 + \bar{b}_m m_{q,i} \text{Tr} [M_q] + (r_m d_m - b_m) \frac{\text{Tr} [M_q^2]}{N_f} + (r_m \bar{d}_m - \bar{b}_m) \frac{\text{Tr} [M_q]^2}{N_f}$$

- ▶ ...from the renorm. PCAC relation we also have

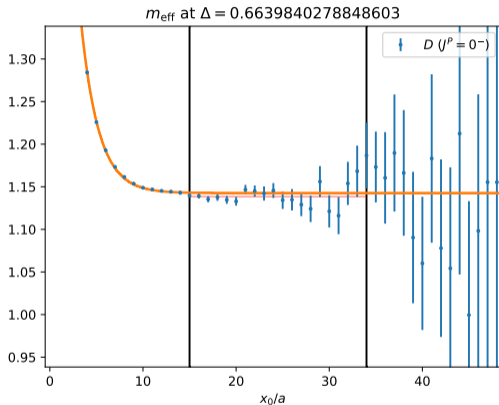
$$\frac{m_{i,R} + m_{j,R}}{2} = \frac{Z_A}{Z_P} m_{ij} \left[ 1 + (b_A - b_P) a m_{q,ij} + (\bar{b}_A - \bar{b}_P) a \text{Tr} [M_q] \right] + O(a^2)$$

## Two strategies

- ▶ fix  $m_c$  via meson mass
  - ▶ here: only via PS channel
    - ▶ stable plateaus
  - ▶  $D$  and  $\eta_c$
- ▶ fix  $m_c$  via RGI quark mass [de Divitiis et al., 2019] (for now: quark mass ratio)
  - ▶ determine estimators for  $(b_A - b_P)$ ,  $\frac{Z_P Z_m}{Z_A}$ ,  $b_m$
  - ▶ employ massive renormalisation scheme
  - ▶ tune e.g.  $\frac{m_c}{m_{\text{sym}}} = 32.89$ , where  $m_{\text{sym}}$  fixed through hadronic scheme at flavour sym. point

# Input meson masses (2 heavy LCP)

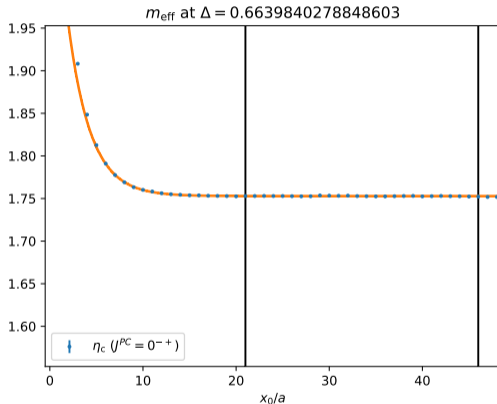
- ▶ eff. masses
- ▶ two-state fit  $m_{\text{eff}}(x_0) = m_0 + c_1 e^{-\Delta_1 x_0} + \dots$
- ▶ plateau start at  $\sigma_{\text{stat}}/4 > c_1 e^{-\Delta_1 x_0}$
- ▶  $\Delta_h = M_{0,h} - M_{0,\text{sym}}$
- ▶ end at  $\sigma_{\text{stat}} > 3\%$
- ▶ fix via
 
$$m_{\bar{D}} = \frac{1}{3}(2m_D + m_{D_s}) = 1866.10(45) \text{ MeV}$$
- ▶ find  $m_{\bar{D}}(\Delta)|_{\Delta_h} = m_{\bar{D}}^{\text{phys}}$





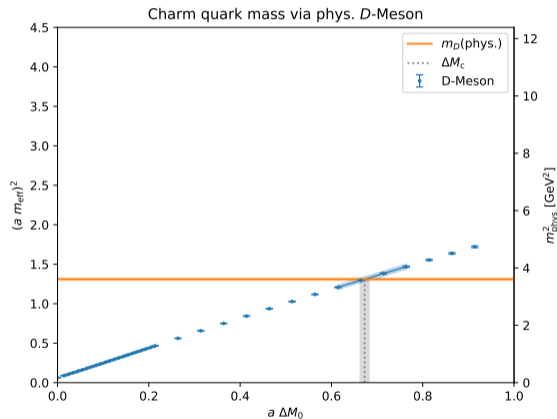
# Input meson masses (2 heavy LCP)

- ▶ eff. masses
- ▶ two-state fit  $m_{\text{eff}}(x_0) = m_0 + c_1 e^{-\Delta_1 x_0} + \dots$
- ▶ plateau start at  $\sigma_{\text{stat}}/4 > c_1 e^{-\Delta_1 x_0}$
- ▶  $\Delta_h = M_{0,h} - M_{0,\text{sym}}$
- ▶ fix via  $m_{\eta_c} = 2983.90(40) \text{ MeV}$
- ▶ find  $m_{\eta_c}(\Delta)|_{\Delta_h} = m_{\eta_c}^{\text{phys}}$

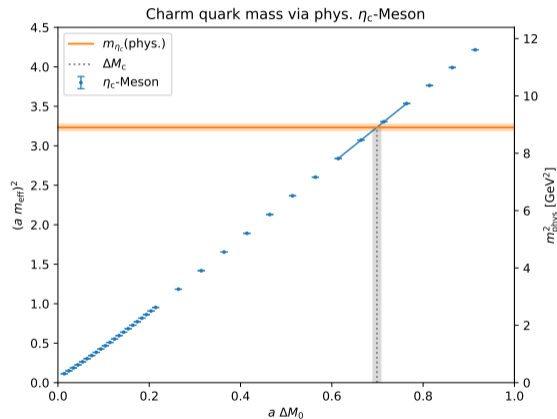


# Input meson masses (2 heavy LCP)

► find mass parameter  $M_{0,c}$  [D]



► find mass parameter  $M_{0,c}$  [ $\eta_c$ ]



# The Quark Mass Ratio

- ▶ fix  $M_{0,c}$  by **Quark Mass Ratio**  $\rightarrow M_{0,c}$  [QMR]
- ▶ ideally: fix by renormalised quark mass (need  $Z_A$ ,  $Z_P$ )
- ▶ do not know critical point  $\Rightarrow m_q$  not yet determined for use in

$$R_1(\Delta_h, \Delta_l) \equiv \frac{m_q(\Delta_h)Z(\Delta_h)[1 + b_m(\Delta_h)am_q(\Delta_h)]}{m_q(\Delta_l)Z(\Delta_l)[1 + b_m(\Delta_l)am_q(\Delta_l)]} + \mathcal{O}(a^2)$$

- ▶ use PCAC masses instead [de Divitiis et al., 2019; Fritzsche et al., 2023]

$$m_{ij} = Z \left[ m_{q,ij} + (r_m - 1) \frac{\text{Tr}[M_q]}{N_f} + aB_{ij} \right] + \mathcal{O}(a^2)$$

with  $Z = \frac{Z_P Z_m}{Z_A}$ ,  $\text{Tr}[M]$  fixed by  $\phi_4$  (see CLS, [Bruno et al., 2015])

# The Quark Mass Ratio

- ▶  $Z, b_m, b_A - b_P$  carry explicit  $\text{Tr}[M_q]$  dependence

$$\tilde{R}_2(\Delta_h, \Delta_l) \equiv \frac{m(\Delta_h)[1 + \tilde{b}_{AP}(\Delta_h)am(\Delta_h)]}{m(\Delta_l)[1 + \tilde{b}_{AP}(\Delta_l)am(\Delta_l)]} + \mathcal{O}(a^2) .$$

- ▶ with  $\tilde{b}_{AP} = \frac{b_A - b_P}{Z}$
- ▶ parametrize  $m(\Delta)$  as a polynomial
- ▶ tune to/compare in continuum to

$$\tilde{R}_2(\Delta_c, \Delta_{\text{sym}}) \equiv \tilde{R}_2(\Delta_h, \Delta_l) |_{\text{phys}} = \frac{m_c}{m_{\text{sym}}} \approx 32.89$$

- ▶ need estimators for  $b_m, b_{AP}, Z$

## Renormalisation constants

- ▶ define estimators for the renormalisation constants [Guagnelli et al., 2001; de Divitiis et al., 2019]

$$R_m = \frac{2(m_{12} - m_{33})}{(m_{22} - m_{11})a\Delta} \xrightarrow{\Delta, m_q \rightarrow 0} b_m$$

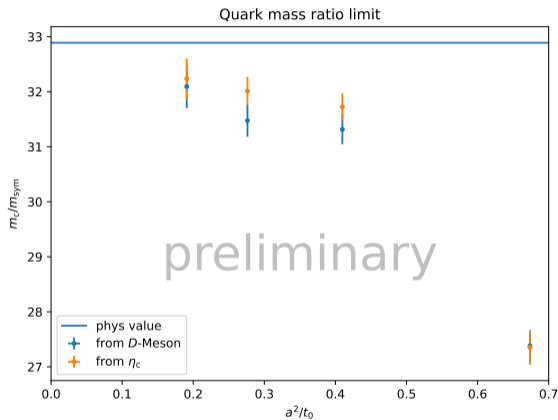
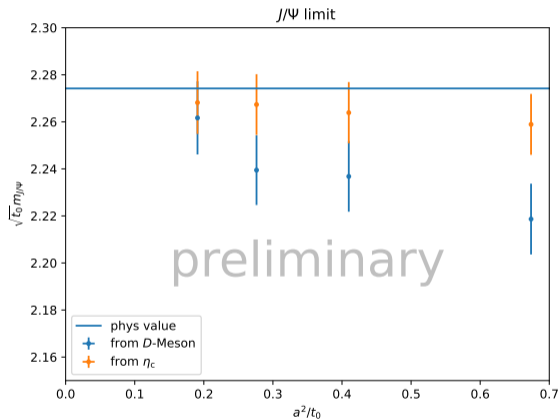
$$R_{AP} = \frac{2m_{12} - m_{11} - m_{22}}{(m_{22} - m_{11})a\Delta} \xrightarrow{\Delta, m_q \rightarrow 0} (b_A - b_P)$$

$$R_Z = \frac{m_{22} - m_{11}}{2\Delta} + (R_{AP} - R_m)(am_{11} + am_{22}) \xrightarrow{\Delta, m_q \rightarrow 0} Z$$

with  $m_{33} = m_{22}(\Delta/2)$

- ▶ evaluate with polynomial interpolation in  $\Delta$

# Preliminary results



# Outlook

- ▶ multistate fit for some meson masses
- ▶ spin averaging for charmed mesons
- ▶ fix charm quark mass via QMR and RGI quark mass directly
- ▶ measurements of phys. observables at fixed  $m_c$

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