

Detecting Lee Yang/Fisher singularities by multi-point Padé

Francesco Di Renzo (University of Parma and INFN)

LATTICE 2023

Fermilab, 07/31/2023

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In collaboration with P. Dimopoulos (Parma), K. Zambello (Pisa),
G. Nicotra, C. Schmidt, S. Singh (Bielefeld), J. Goswami (RIKEN), D. Clarke (Utah)

see also talks by C. Schmidt and D. Clarke on Friday



**UNIVERSITÀ
DI PARMA**



Agenda

- An invitation (multi-point Padè)
- What we have been doing in Lattice QCD
- What we could do for the 2d Ising model: finite size scaling
- Can we do the same for QCD?

An invitation

A few words on multi-point PADÈ

Suppose you know the **values** of a **function** (and of its derivatives) at a number of points

$$\dots, f(z_k), f'(z_k), \dots, f^{(s-1)}(z_k), \dots \quad k = 1 \dots N$$

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If you want to **approximate the function with a rational function**

$$R_n^m(z) = \frac{P_m(z)}{\tilde{Q}_n(z)} = \frac{P_m(z)}{1 + Q_n(z)} = \frac{\sum_{i=0}^m a_i z^i}{1 + \sum_{j=1}^n b_j z^j}$$

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This is the starting point for a **multi-point Padè approximation**: solve the linear system

$$\begin{aligned} & \dots \\ & P_m(z_k) - f(z_k)Q_n(z_k) = f(z_k) \\ & P'_m(z_k) - f'(z_k)Q_n(z_k) - f(z_k)Q'_n(z_k) = f'(z_k) \\ & \dots \end{aligned}$$

from which we want to get the unknown

$$\{a_i \mid i = 0 \dots m\} \quad \{b_j \mid j = 1 \dots n\} \quad n + m + 1 = N s$$

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In principle you should require that $\nexists z_0 : P_m(z_0) = \tilde{Q}_n(z_0) = 0$

... but we will need to give up with this ...

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Yes! LATTICE QCD at IMAGINARY values of the baryonic chemical potential

PHYSICAL REVIEW D **105**, 034513 (2022)

Contribution to understanding the phase structure of strong interaction matter: Lee-Yang edge singularities from lattice QCD

P. Dimopoulos¹, L. Dini², F. Di Renzo¹, J. Goswami², G. Nicotra², C. Schmidt²,
S. Singh^{1,*}, K. Zambello¹ and F. Ziesché²

... where we computed and “multi-point Padè approximated”

$$\begin{aligned}\chi_n^B(T, V, \mu_B) &= \left(\frac{\partial}{\partial \hat{\mu}_B} \right)^n \frac{\ln Z(T, V, \mu_l, \mu_s)}{VT^3} \\ &= \left(\frac{1}{3} \frac{\partial}{\partial \hat{\mu}_l} + \frac{1}{3} \frac{\partial}{\partial \hat{\mu}_s} \right)^n \frac{\ln Z(T, V, \mu_l, \mu_s)}{VT^3}\end{aligned}$$

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Why a rational approximation instead of a polynomial? Because you have **POLES** that can mimic the **SINGULARITIES** of your function! (*at least the nearest ones ...*)

Any useful ...?

Yes! LATTICE QCD at **IMAGINARY** values of the **baryonic chemical potential**
... a natural analytic continuation to real chemical potential!

PHYSICAL REVIEW D **105**, 034513 (2022)

... and not only that: **singularities!**

Contribution to understanding the phase structure of strong interaction matter: Lee-Yang edge singularities from lattice QCD

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We present a calculation of the net baryon number density as a function of imaginary baryon number chemical potential, obtained with highly improved staggered quarks at temporal lattice extent of $N_\tau = 4, 6$. We construct various rational function approximations of the lattice data and discuss how poles in the complex plane can be determined from them. We compare our results of the singularities in the chemical potential plane to the theoretically expected positions of the Lee-Yang edge singularity in the vicinity of the Roberge-Weiss and chiral phase transitions. We find a temperature scaling that is in accordance with the expected power law behavior.

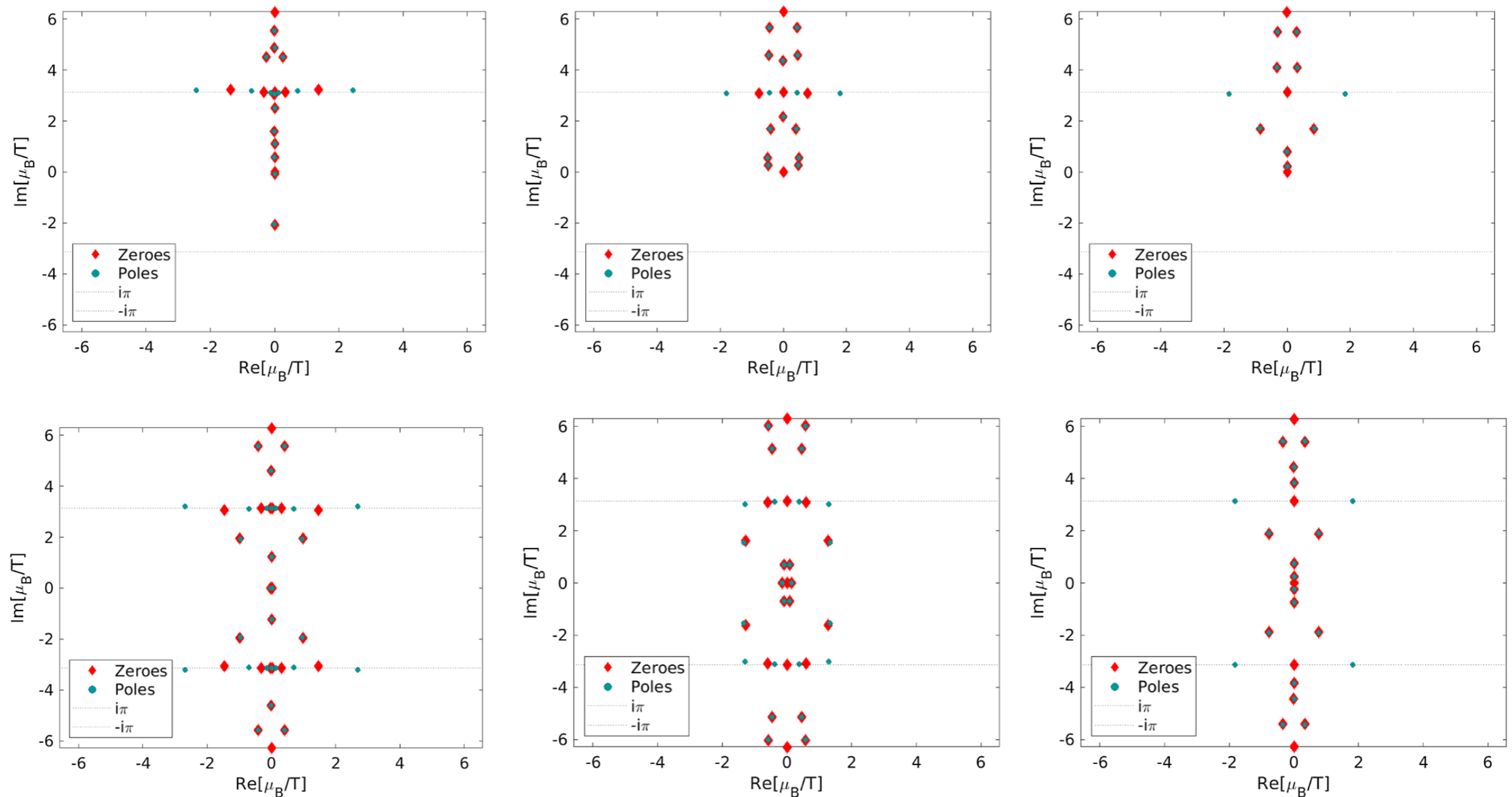


FIG. 5. Singularity structure in the $\hat{\mu}_B$ plane for three different temperatures (from left to right $T = 201.4, 186.3, 167.4$). Upper row: *Ansatz (15)*; lower row: *Ansatz (20)*.

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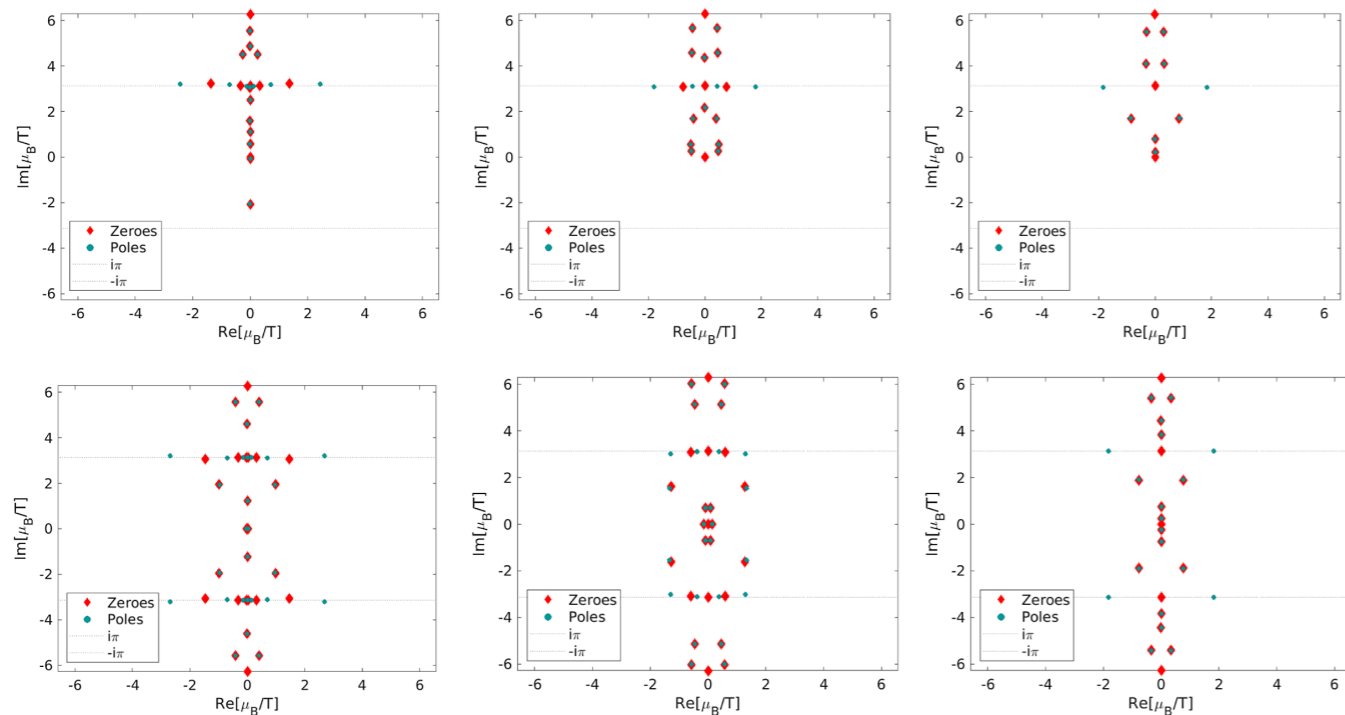


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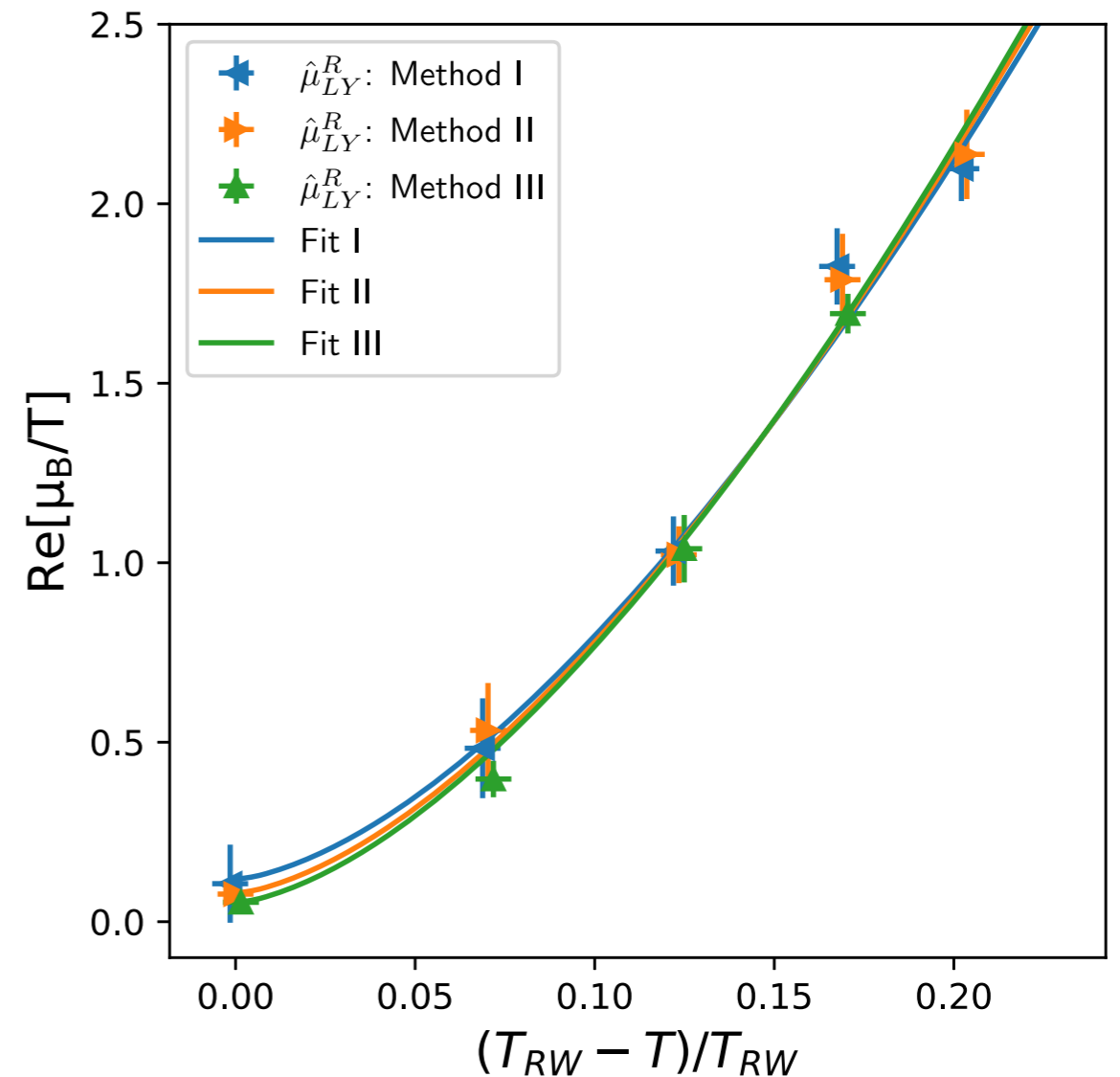


FIG. 9. Scaling fit to the Lee-Yang edge singularities in the vicinity of the Roberge-Weiss transition to the *Ansatz* (22). Shown are three distinct data sets for the real parts of the $\hat{\mu}_B$ (imaginary parts of h) as a function of the reduced temperature $(T_{RW} - T)/T_{RW}$, as obtained from methods I–III.

Can we do *better* than this ?!?

As reported at Lattice 2022,
we tested our approach on the
2D ISING model


ISING model

$$U(\sigma) = -J \sum_{\langle i,j \rangle} \sigma_i \sigma_j - h \sum_i \sigma_i$$

$$Z(\beta, h) = Z(0, h) e^{\beta c} \prod_k (1 - \beta/\beta_k)$$

$$\langle\langle U^n \rangle\rangle = (-1)^{(n-1)} \sum_k \frac{(n-1)!}{(\beta_k - \beta)^n}, \quad n > 1$$

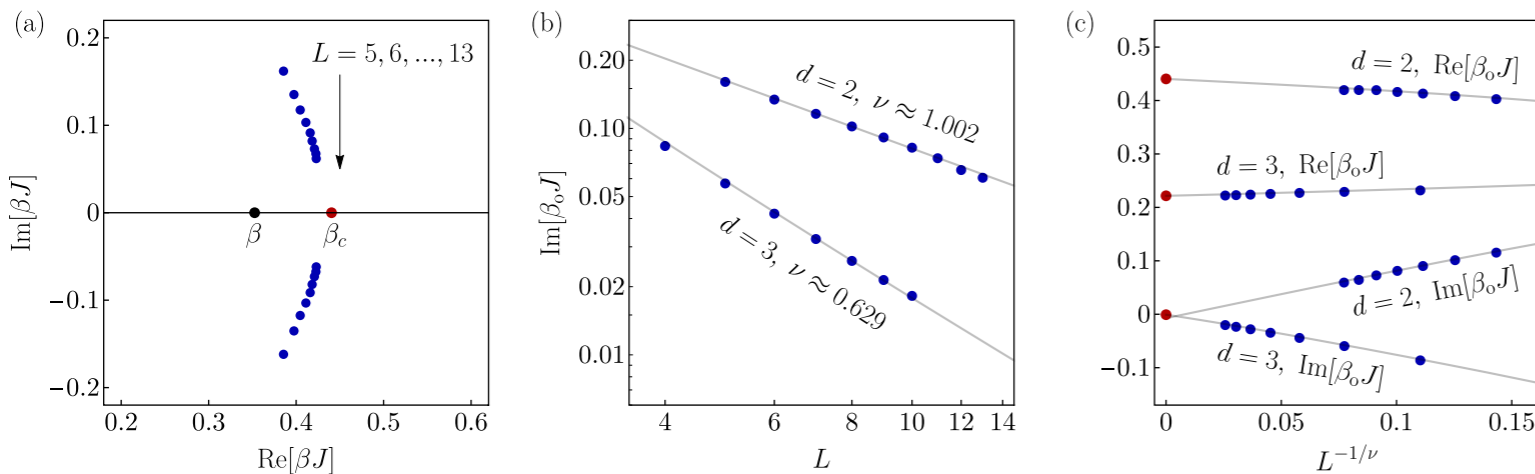
Determination of universal critical exponents using Lee-Yang theory

Aydin Deger and Christian Flindt 

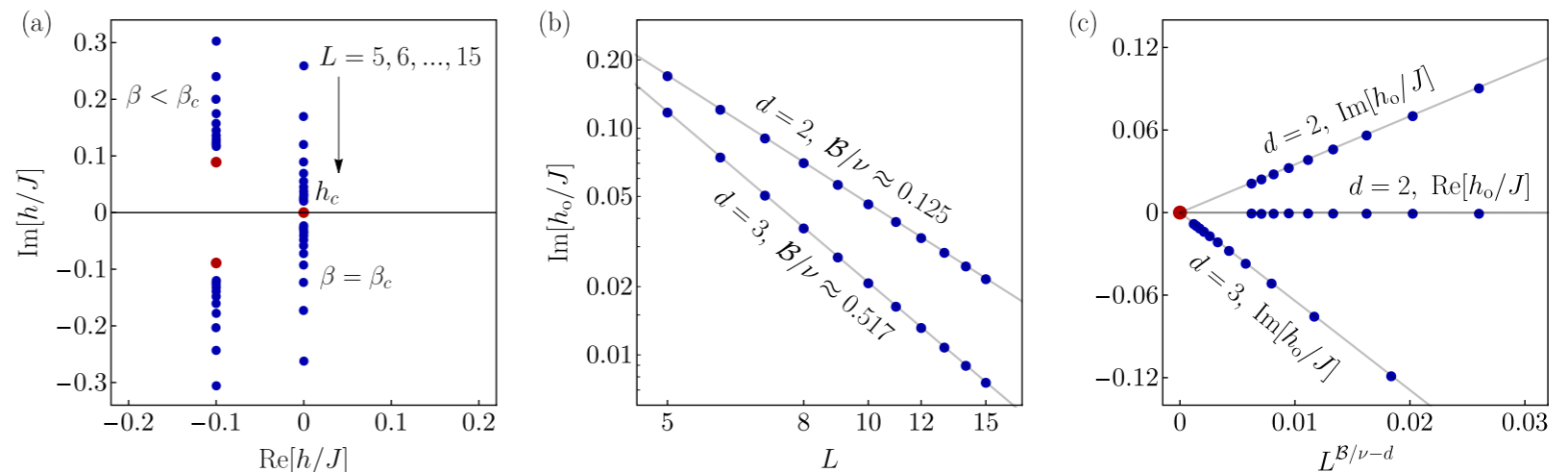
Department of Applied Physics, Aalto University, 00076 Aalto, Finland

Zeros of the partition function determined via computations of **cumulants** (derivatives of $\log(Z)$ with respect to inverse temperature)

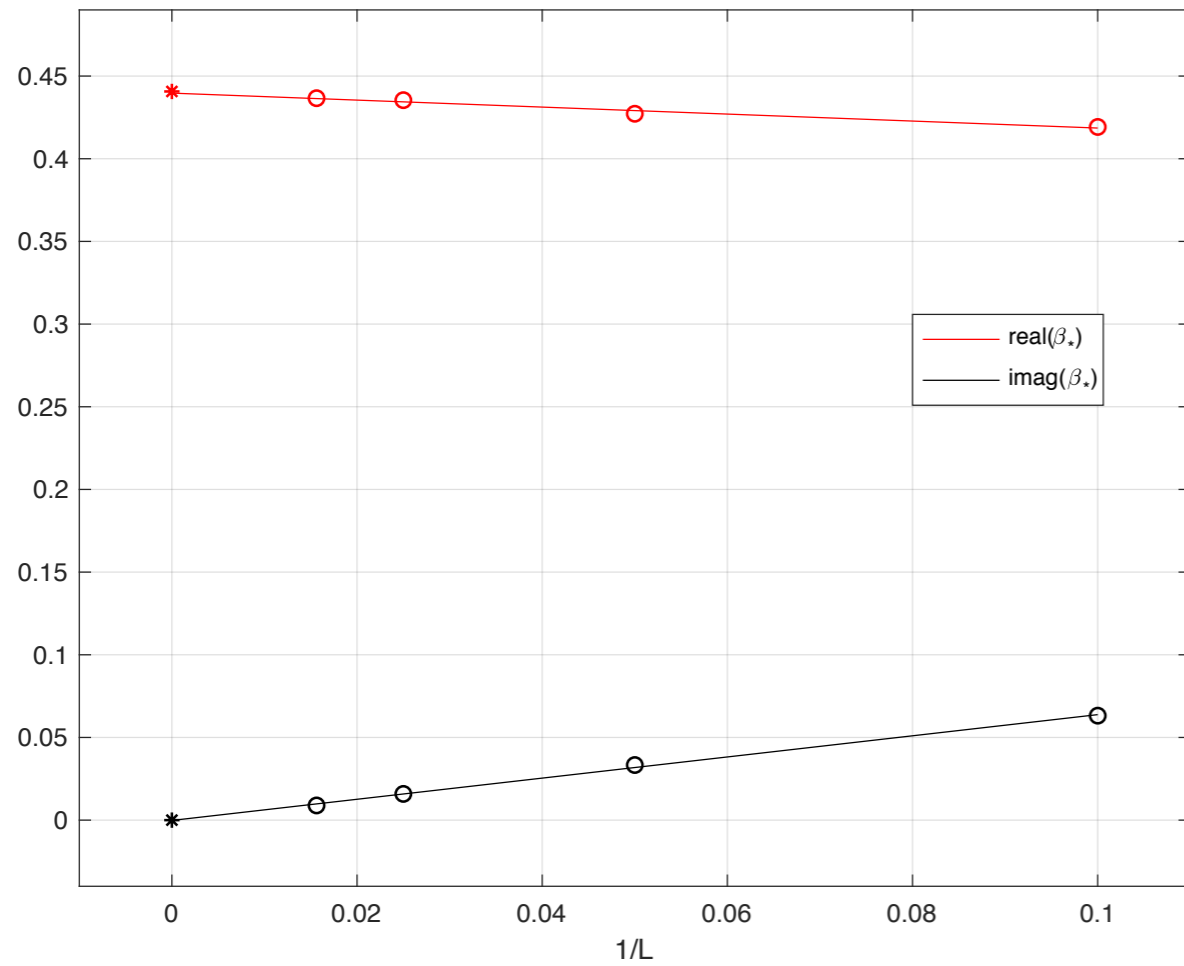
Scaling relations are supposed to describe the approach of leading zeros to critical inverse temperature.



Dealing instead with **leading zeros from magnetisation cumulants** (now derive with respect to magnetic field)

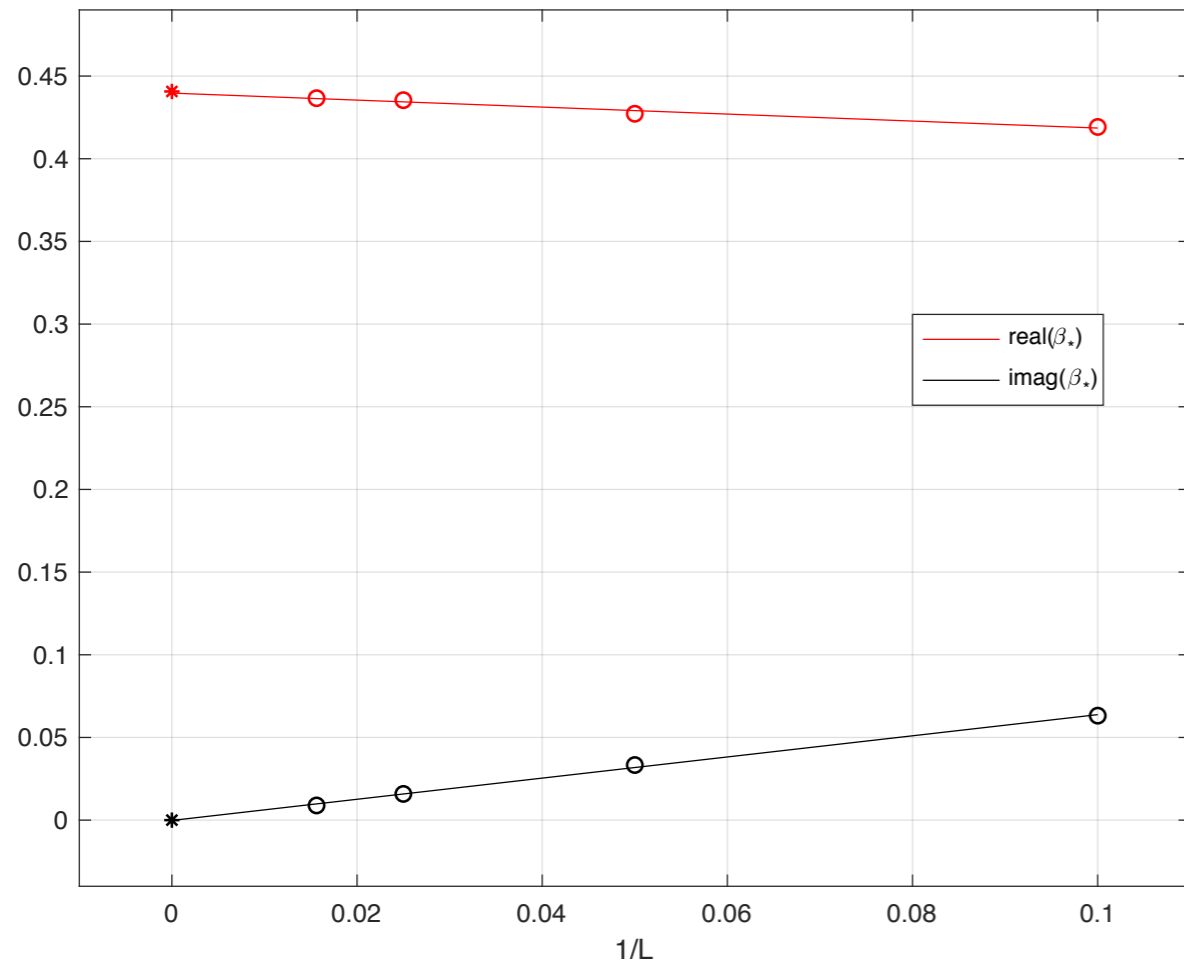


The same scaling plots obtained with our multi-point PADÈ method

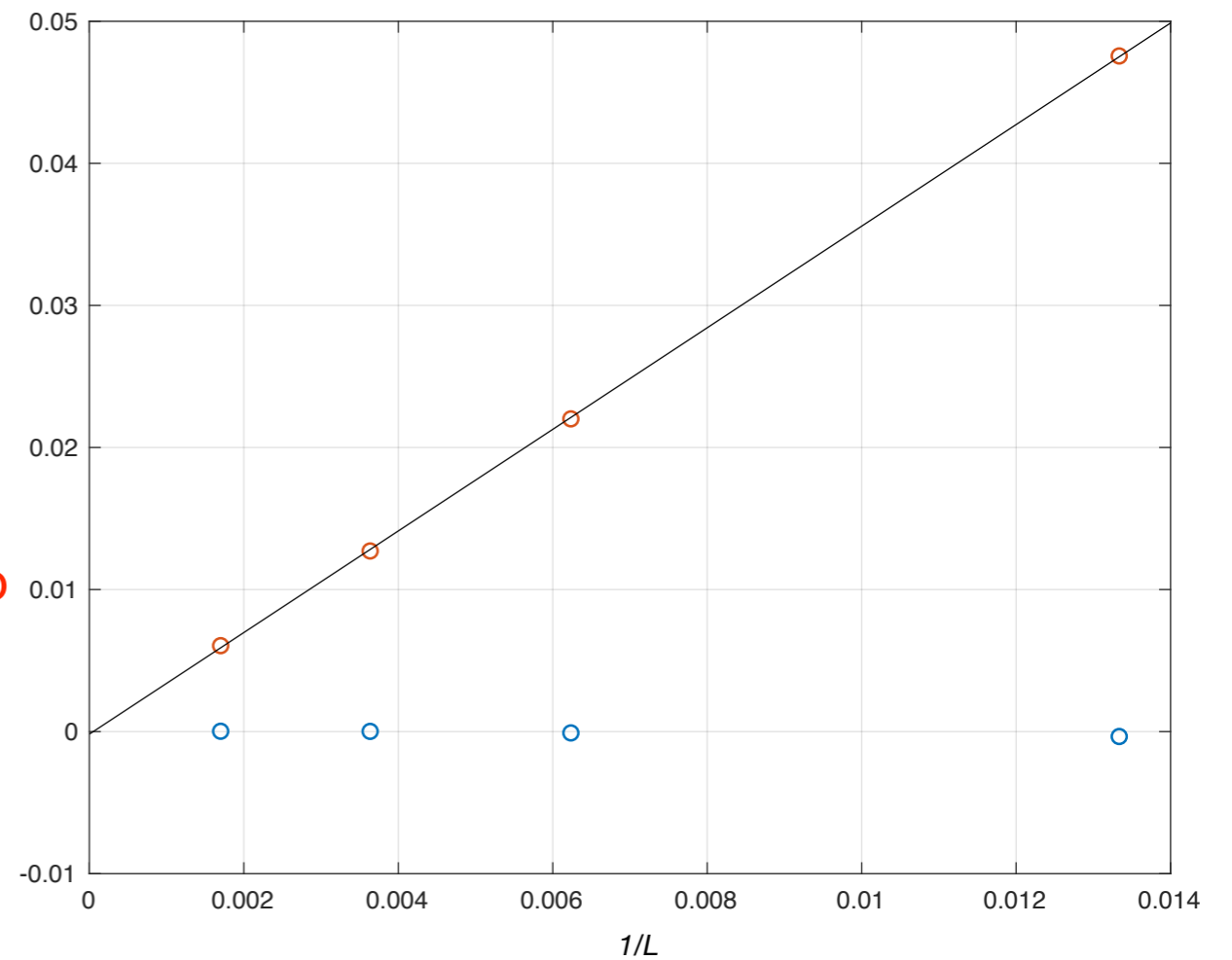


Real and imaginary parts of the **leading (Fisher) zero**
as a function of (inverse) **lattice size**
*(from the computation of specific heat at different
temperatures and lattice sizes)*

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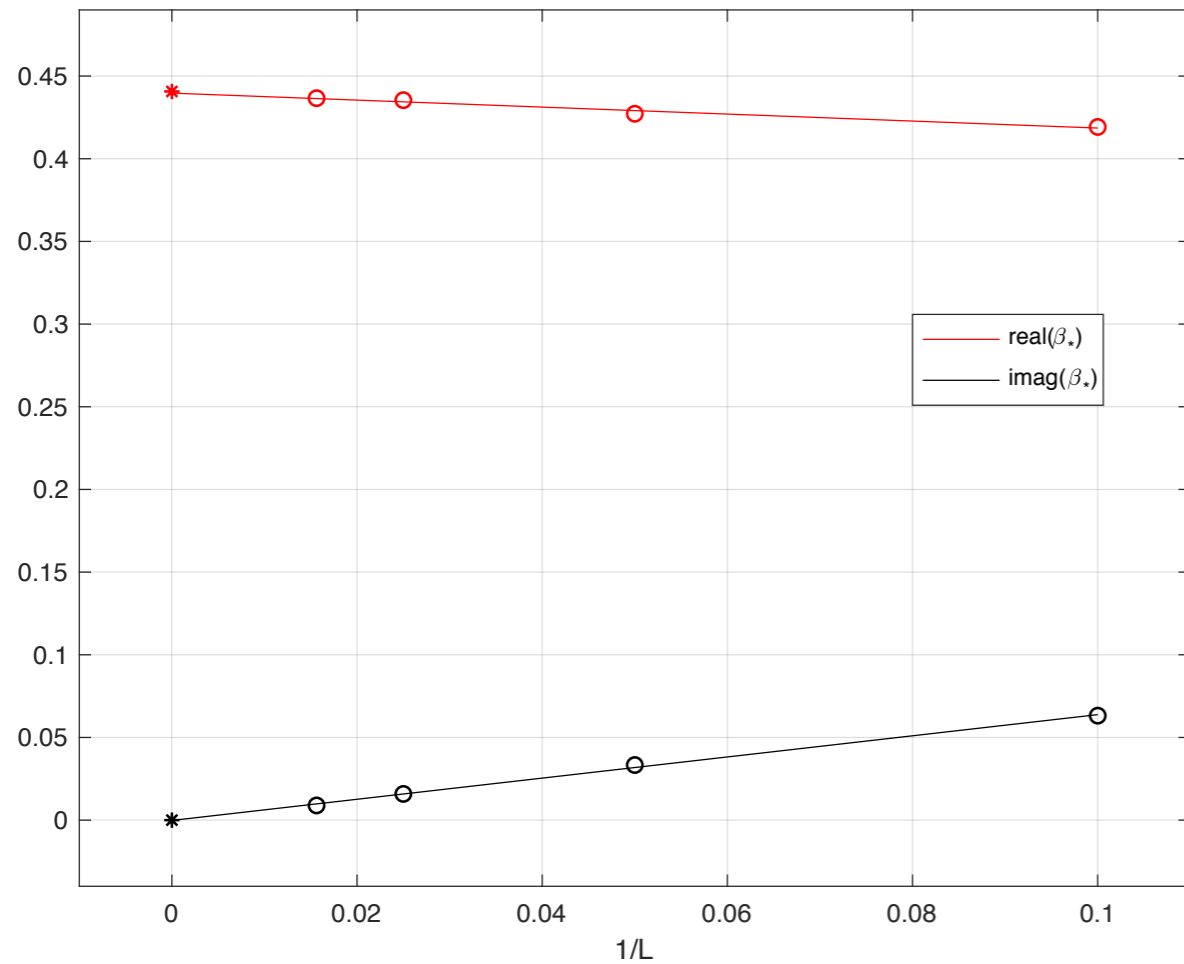


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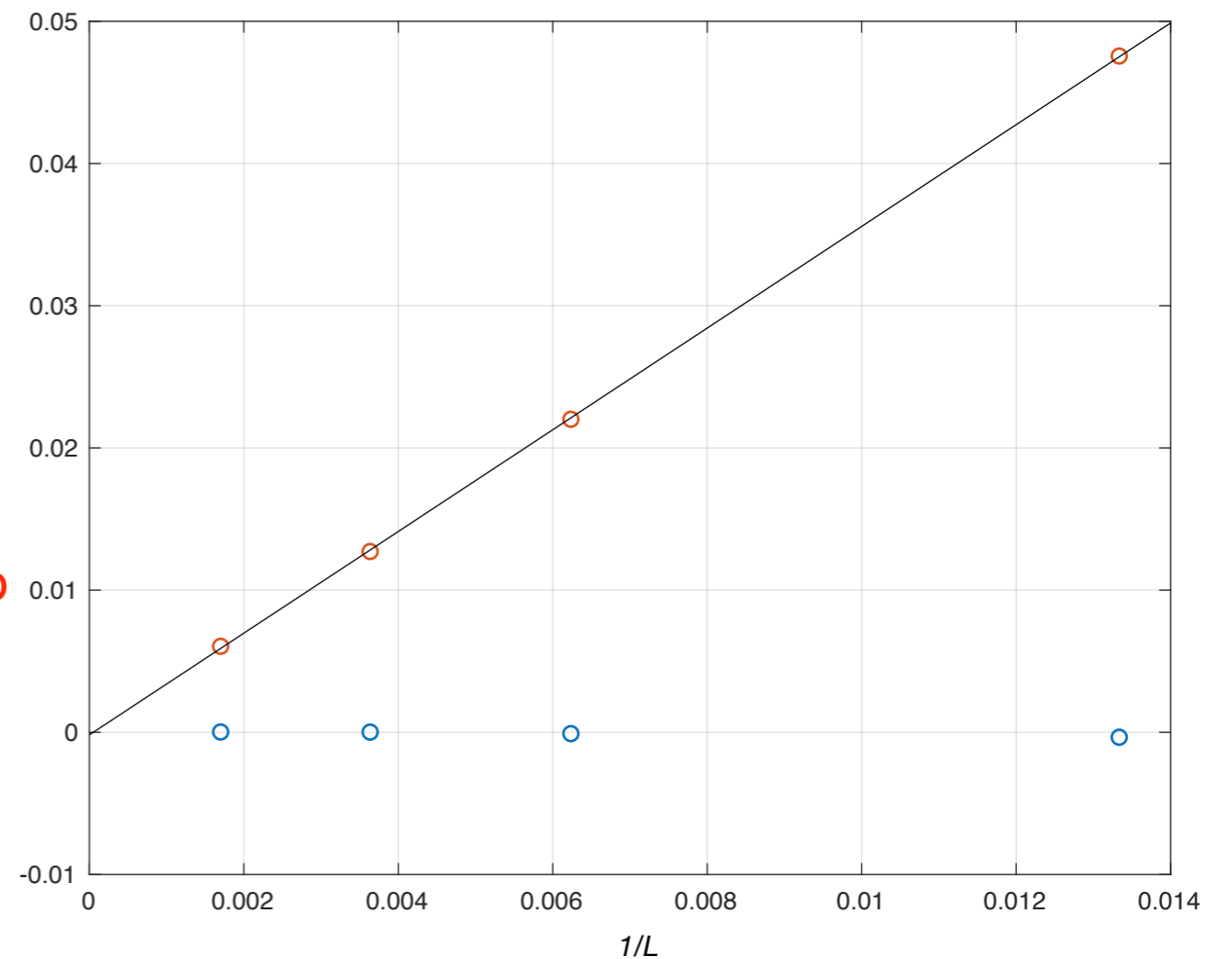


Real and imaginary parts of the **leading (Lee Yang) zero** as a function of (inverse) **lattice size**
(from the computation of magnetisation at different magnetic fields and lattice sizes, sitting at β_c)

The same scaling plots obtained with our multi-point PADÈ method



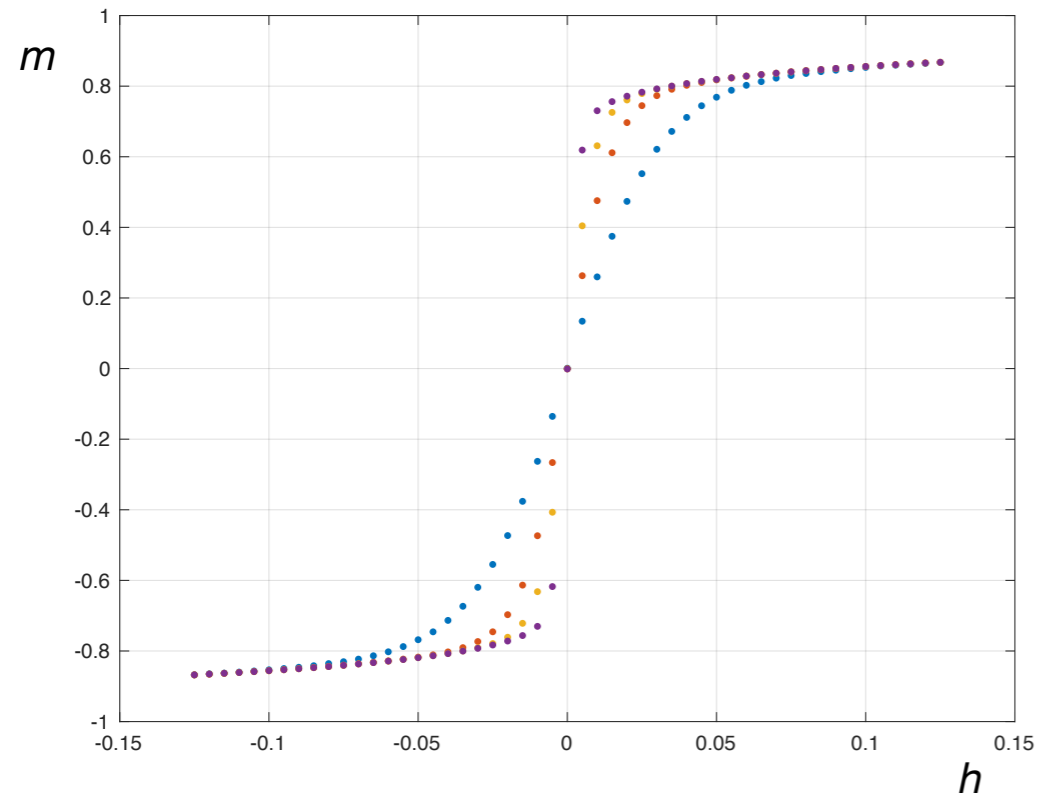
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(from the computation of magnetisation at different magnetic fields and lattice sizes, sitting at β_c)

Please notice: **here we are following the right order!** (i.e. we first get the critical temperature)

Let us focus on the second plot (i.e. we look at m as a function of h , sitting at the critical temperature)

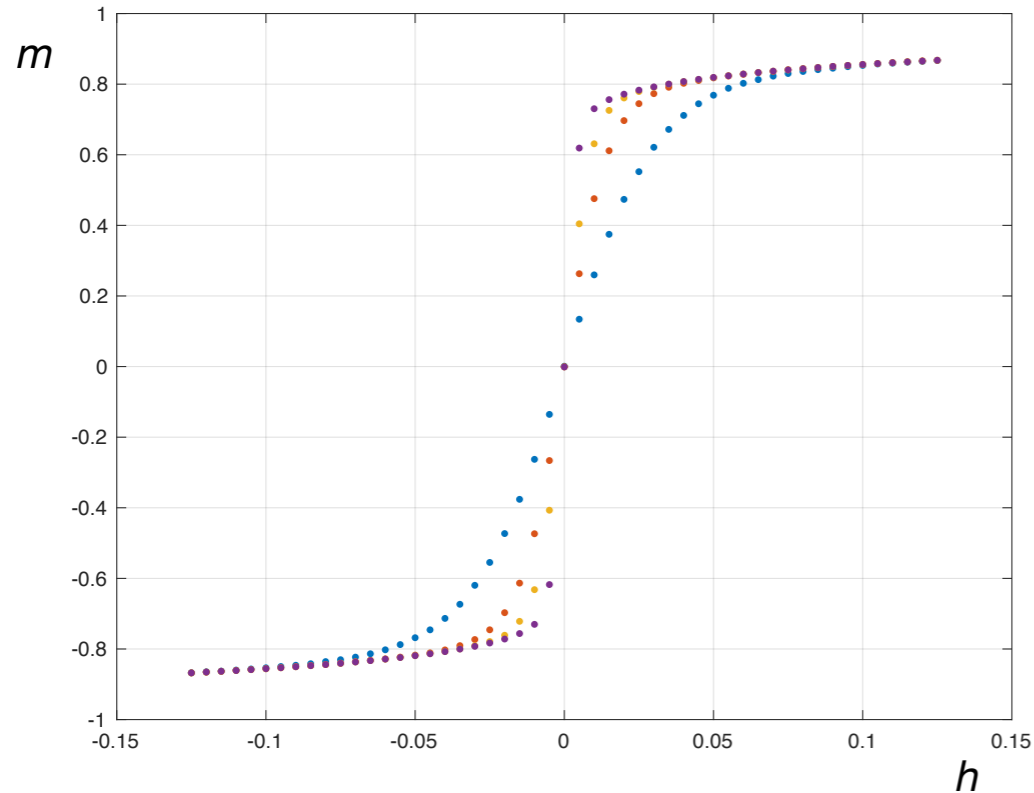


We compute **magnetisation**
at different magnetic **fields** and lattice **sizes**
and obtain **Padè** approximants

i.e.

$$f(z) \mapsto m^{(L)}(h) \quad R_m^n(z) \mapsto R_m^{n(L)}(h)$$

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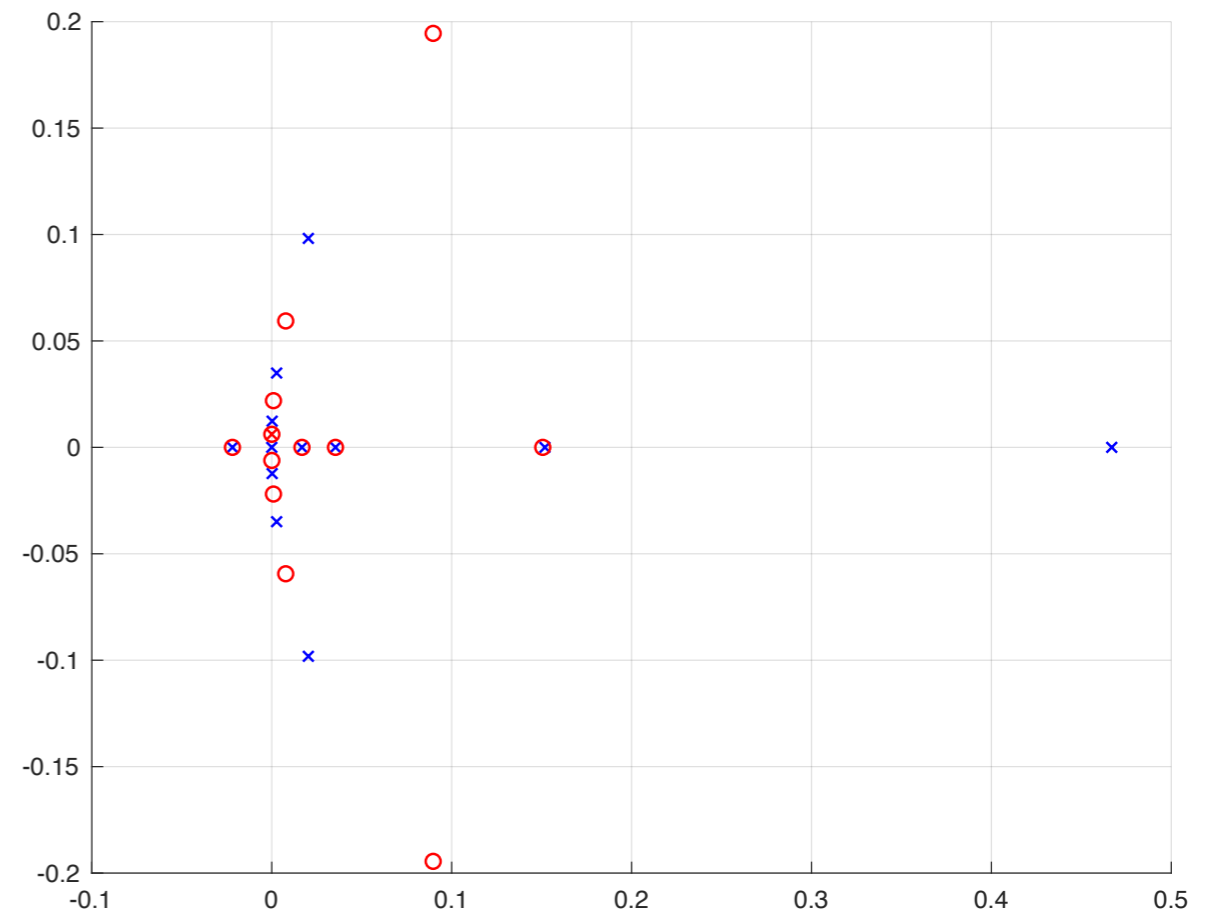
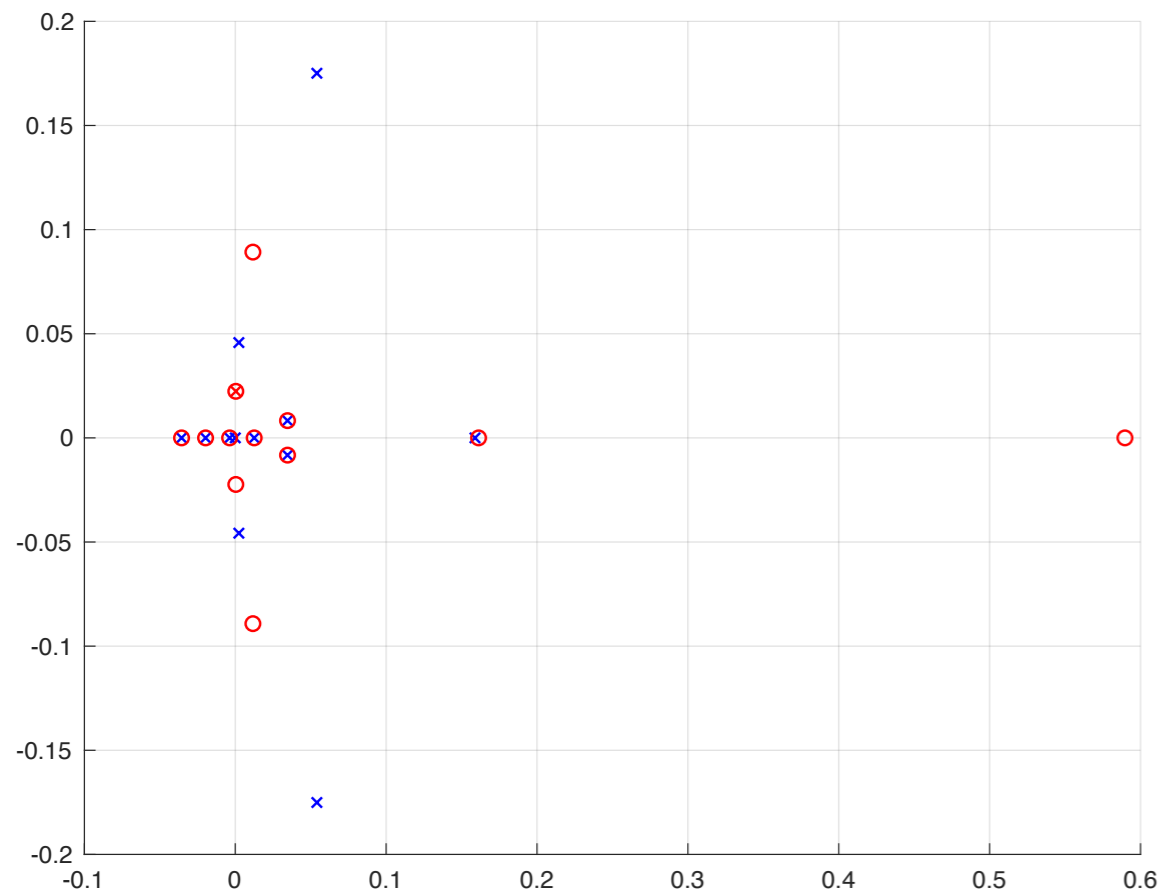
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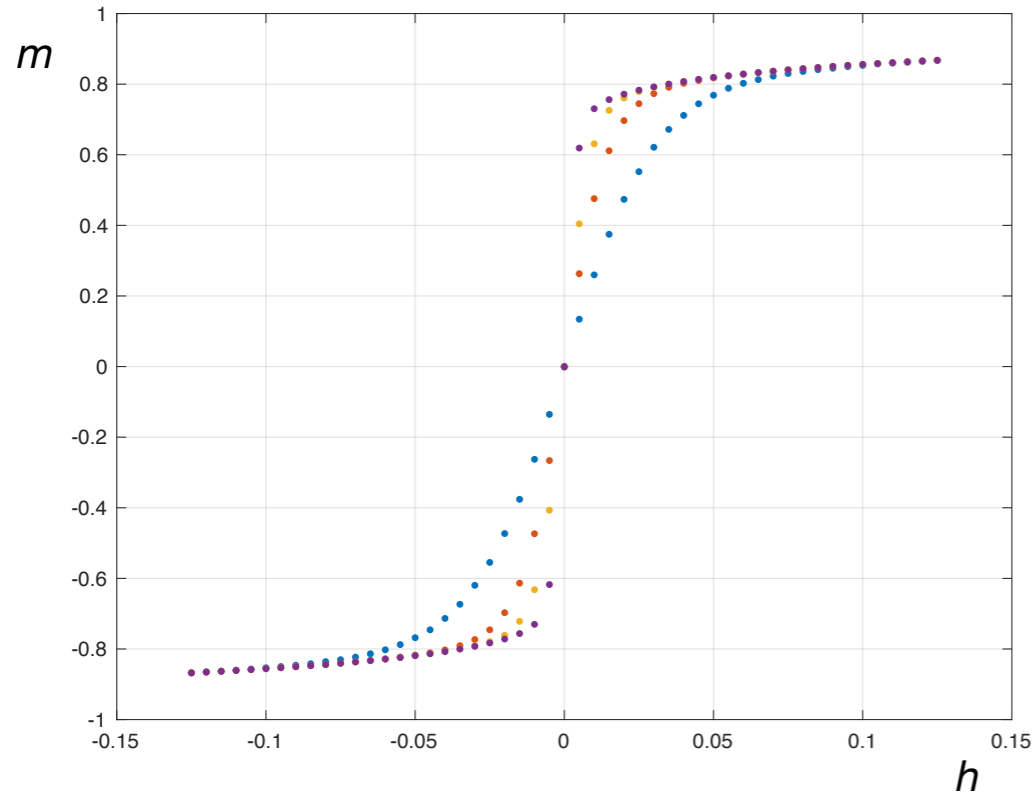
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We determine the **leading pole** (red circle) from the **Padè** approximants

(figures are in **complex h plane**)



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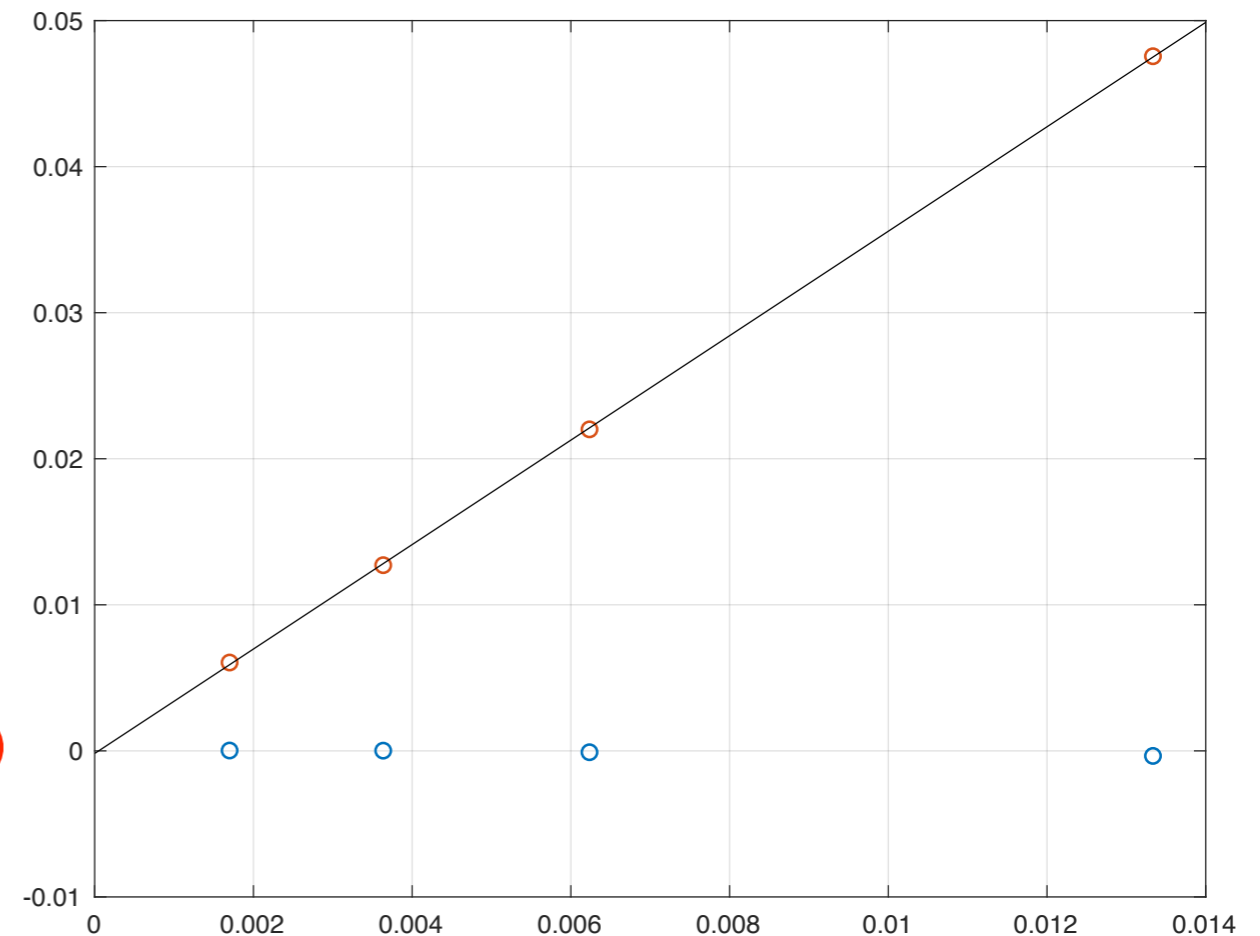
We determine the **leading pole** (**red circle**) from the **Padè** approximants

And finally we inspect the **scaling** of **leading poles** (**red circles**)

(figure again in **complex h plane**)

$$h_0 \sim L^{-\frac{\beta\delta}{\nu}}$$

Combination of the relevant **critical exponents** got to a few **per mille**



Can we repeat this
for (Lattice) QCD?

Who is who ...

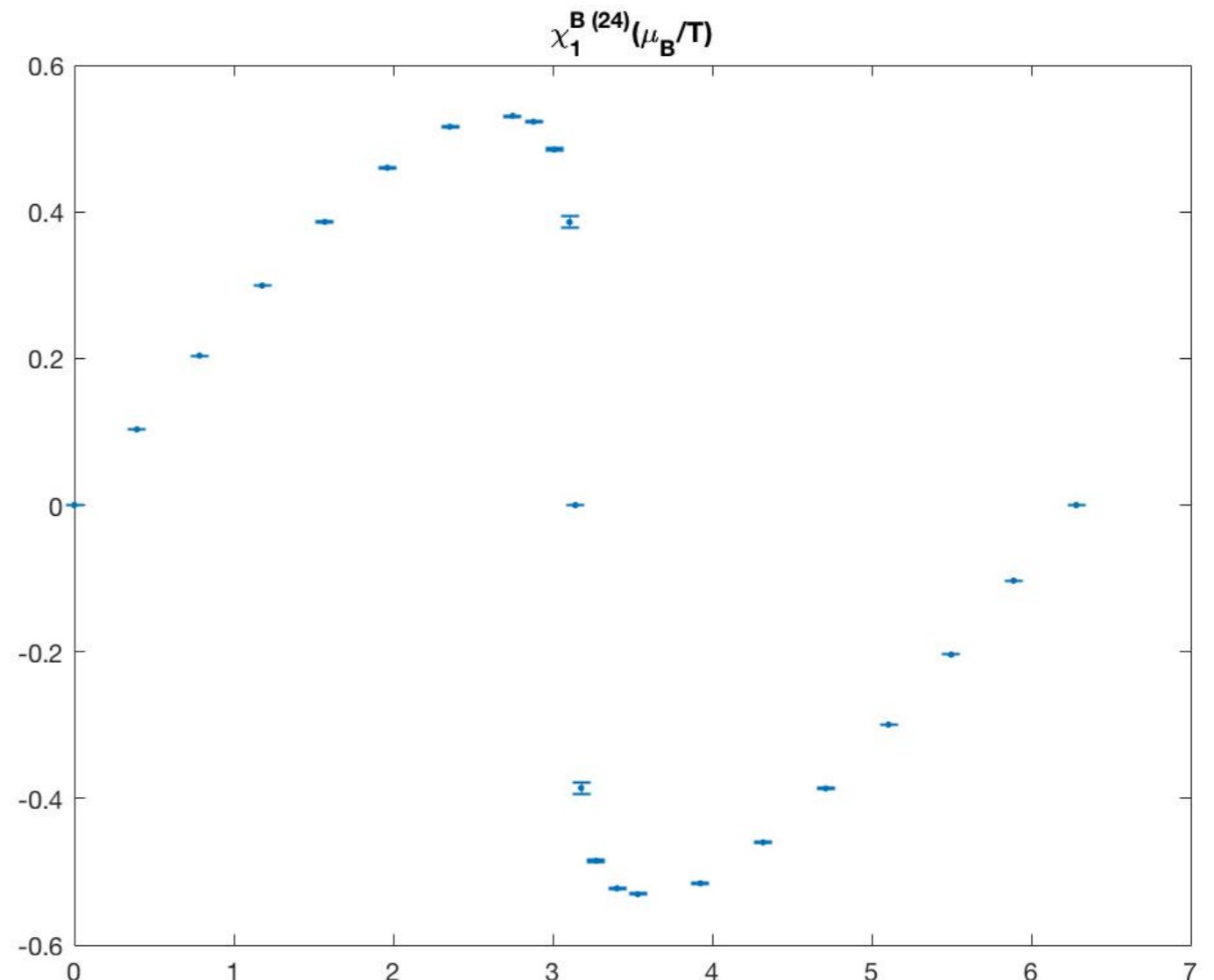
We now move to the setting of QCD with **baryonic chemical potential**

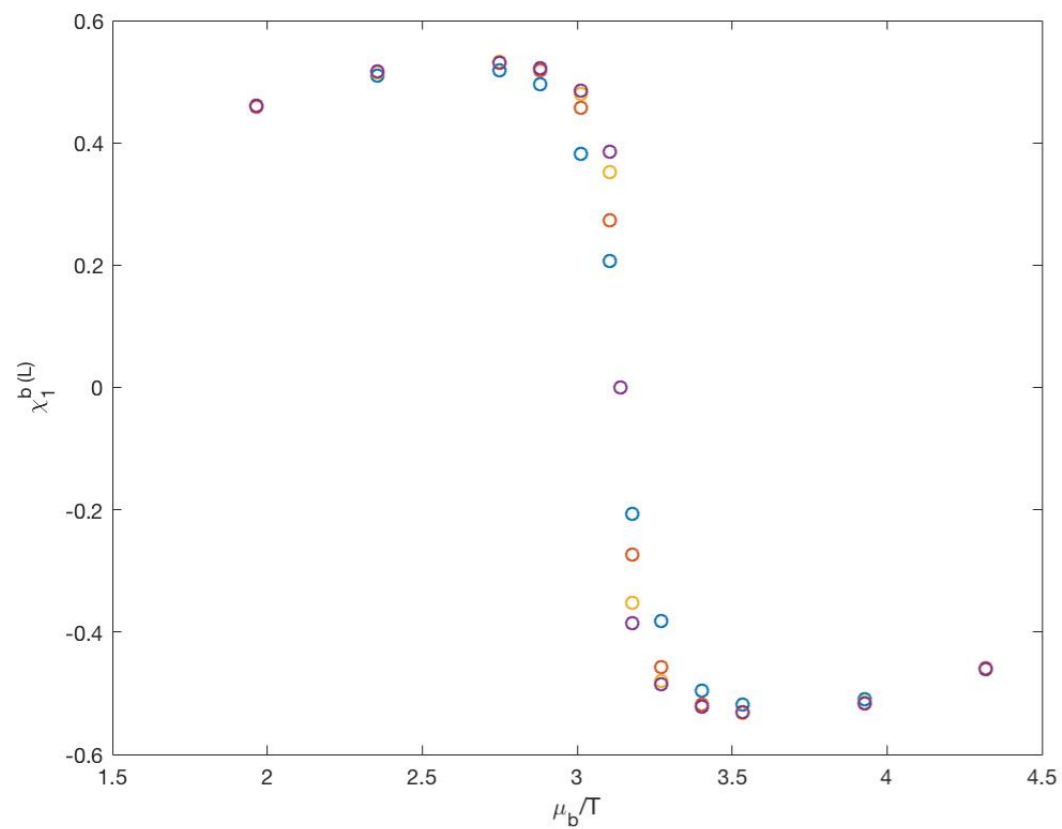
$$h \mapsto \hat{\mu}_B = \frac{\mu_B}{T} \quad m^{(L)}(h) \mapsto \chi_1^{B(L)}(\hat{\mu}_B) = \frac{\partial}{\partial \hat{\mu}_B} \frac{\ln Z}{VT^3}$$

which we probe at $T \sim 200 \text{ MeV}$ on a **coarse** lattice $N_\tau = 4$ with $N_\sigma = L = 12, 16, 20, 24$ at **imaginary baryonic chemical potential**, looking for the **Roberge Weiss transition**.

$$\hat{\mu}_{B \text{ cr}} = i\pi$$

and we have to look for the scaling of $\hat{\mu}_{B0}^{(R)}$

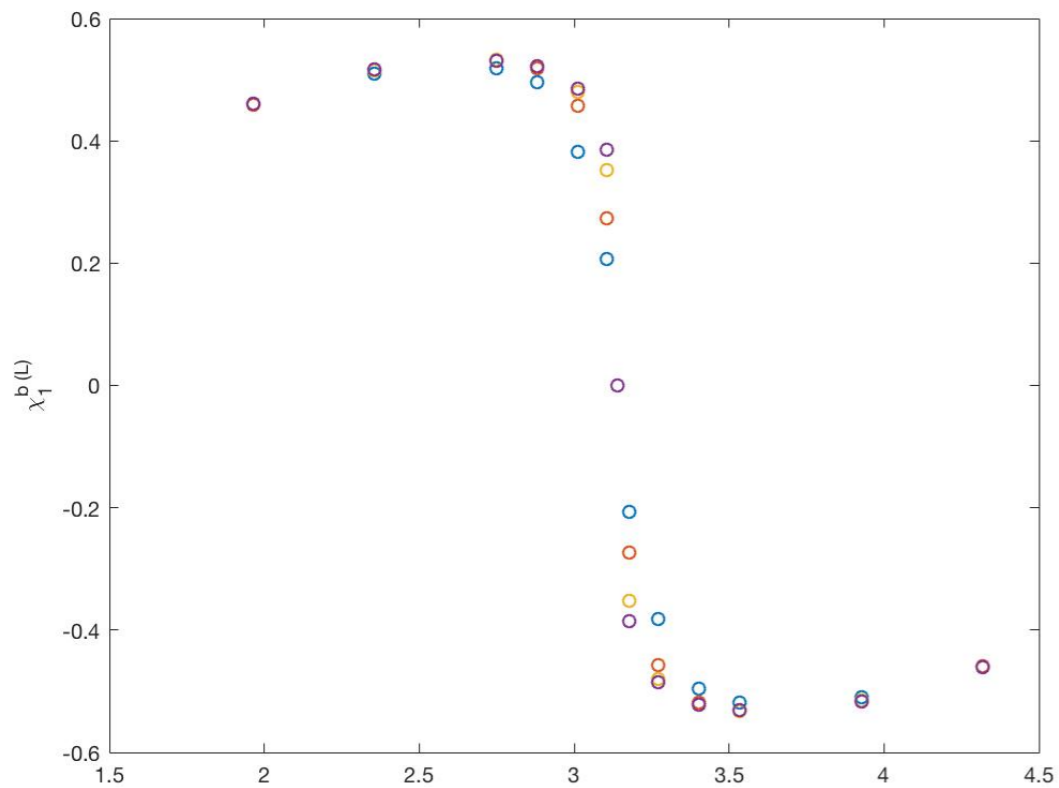




We compute **number density**
at different values of **imaginary chemical**
potential and **lattice size**
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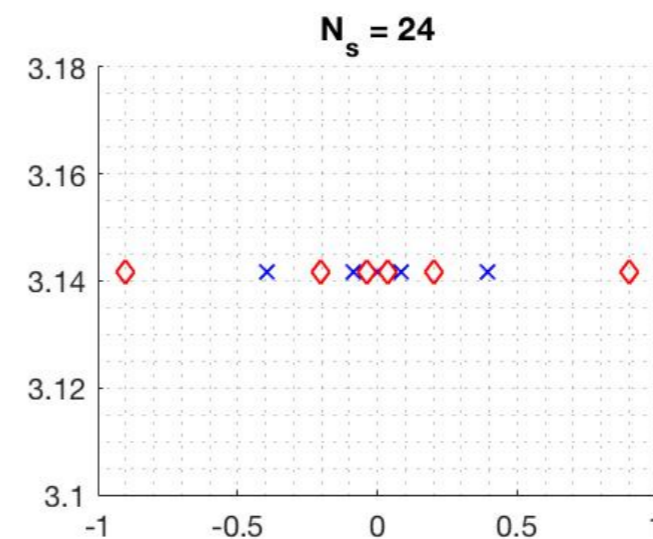
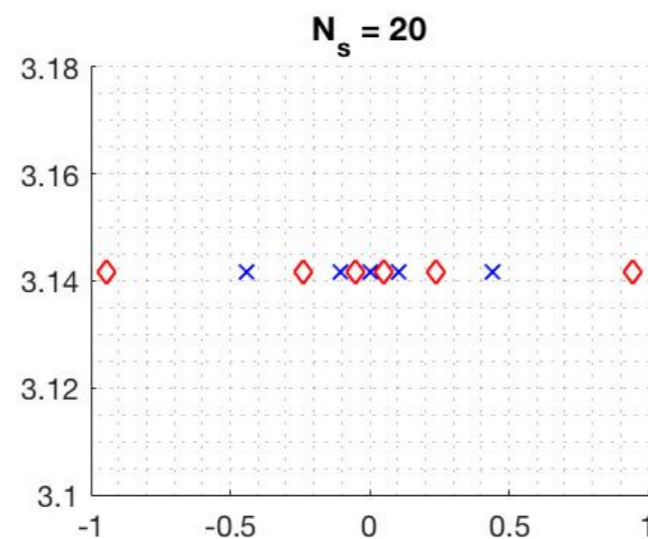
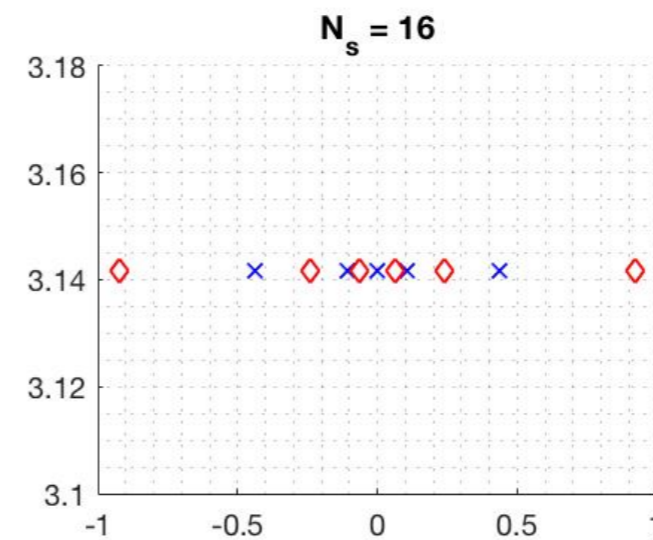
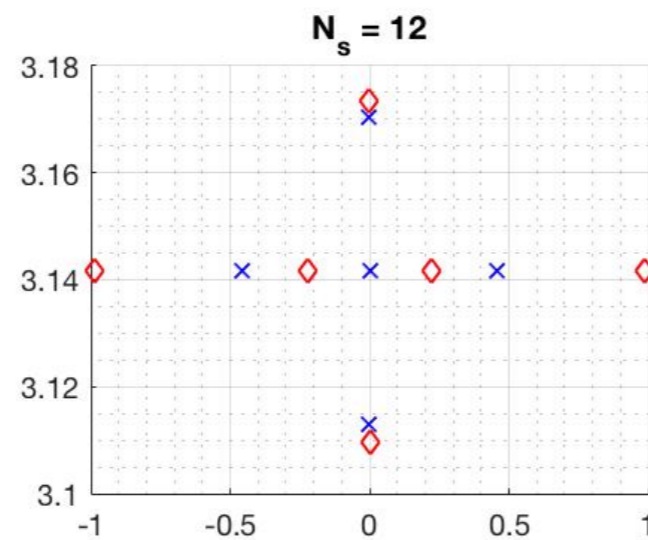
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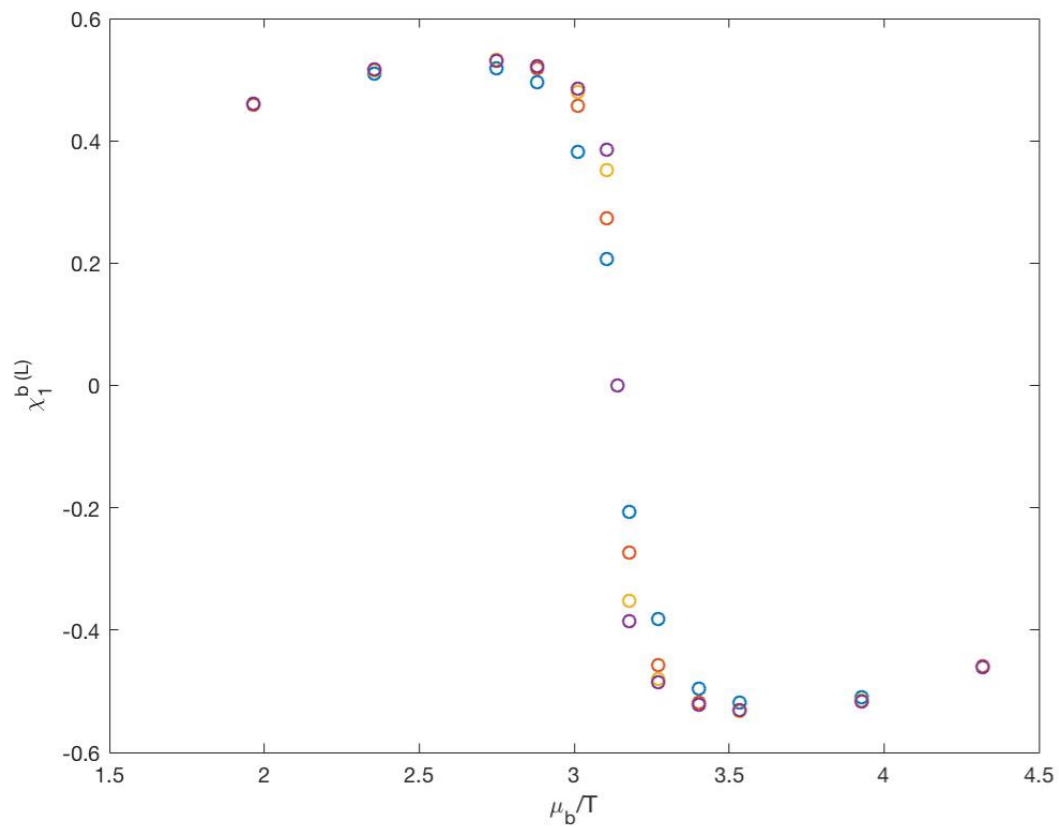
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$$f(z) \mapsto \chi_1^{B(L)}(\hat{\mu}_B) \quad R_m^n(z) \mapsto R_m^n(L)(\hat{\mu}_B)$$

We determine the **leading pole** (red symbol) from the **Padè approximants**

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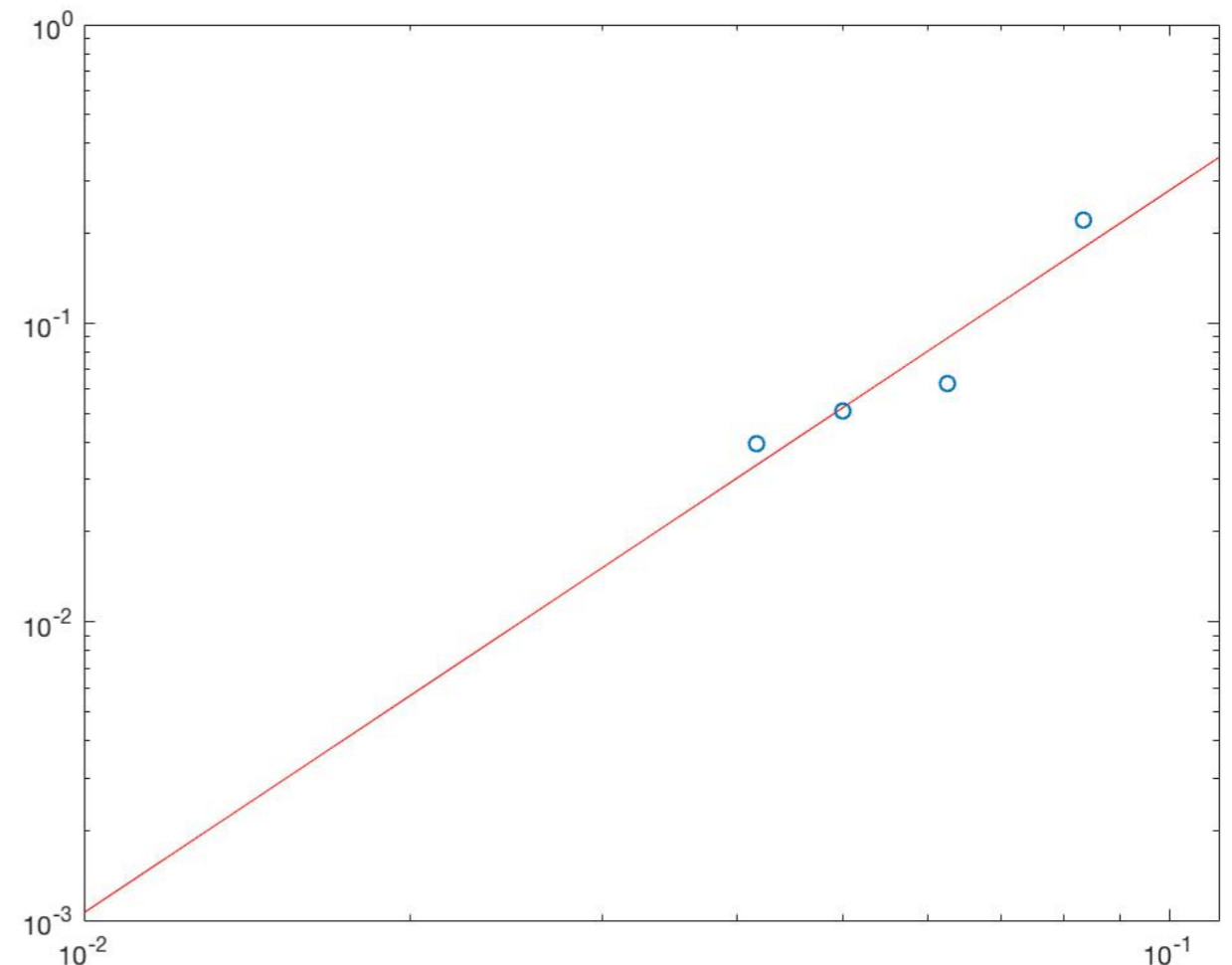
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And finally we inspect the **scaling** of **leading poles** (red circles)

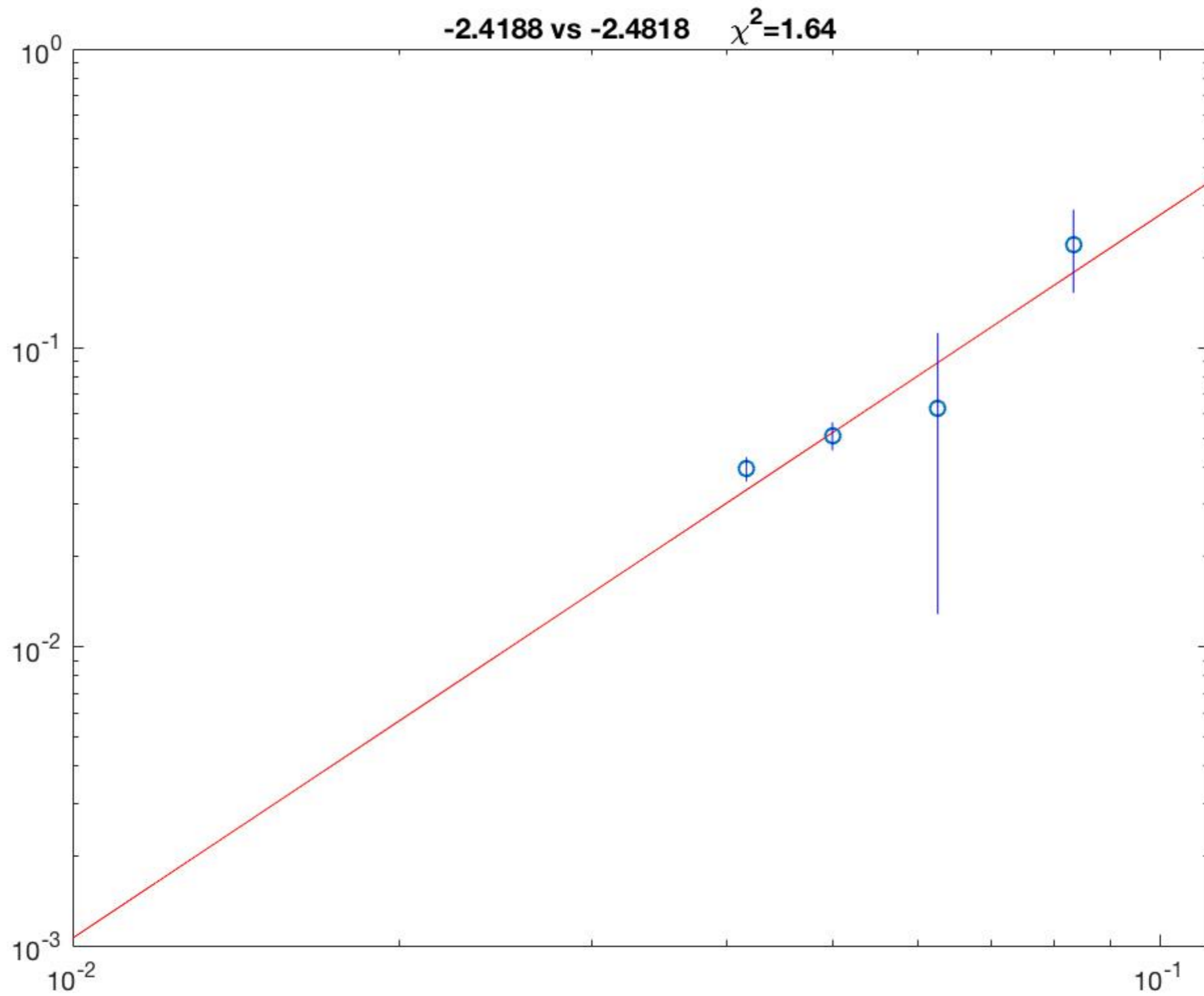
(figure again in **complex chem.pot. plane**)

$$\hat{\mu}_{B0}^{(R)} \sim L^{-\frac{\beta\delta}{\nu}}$$

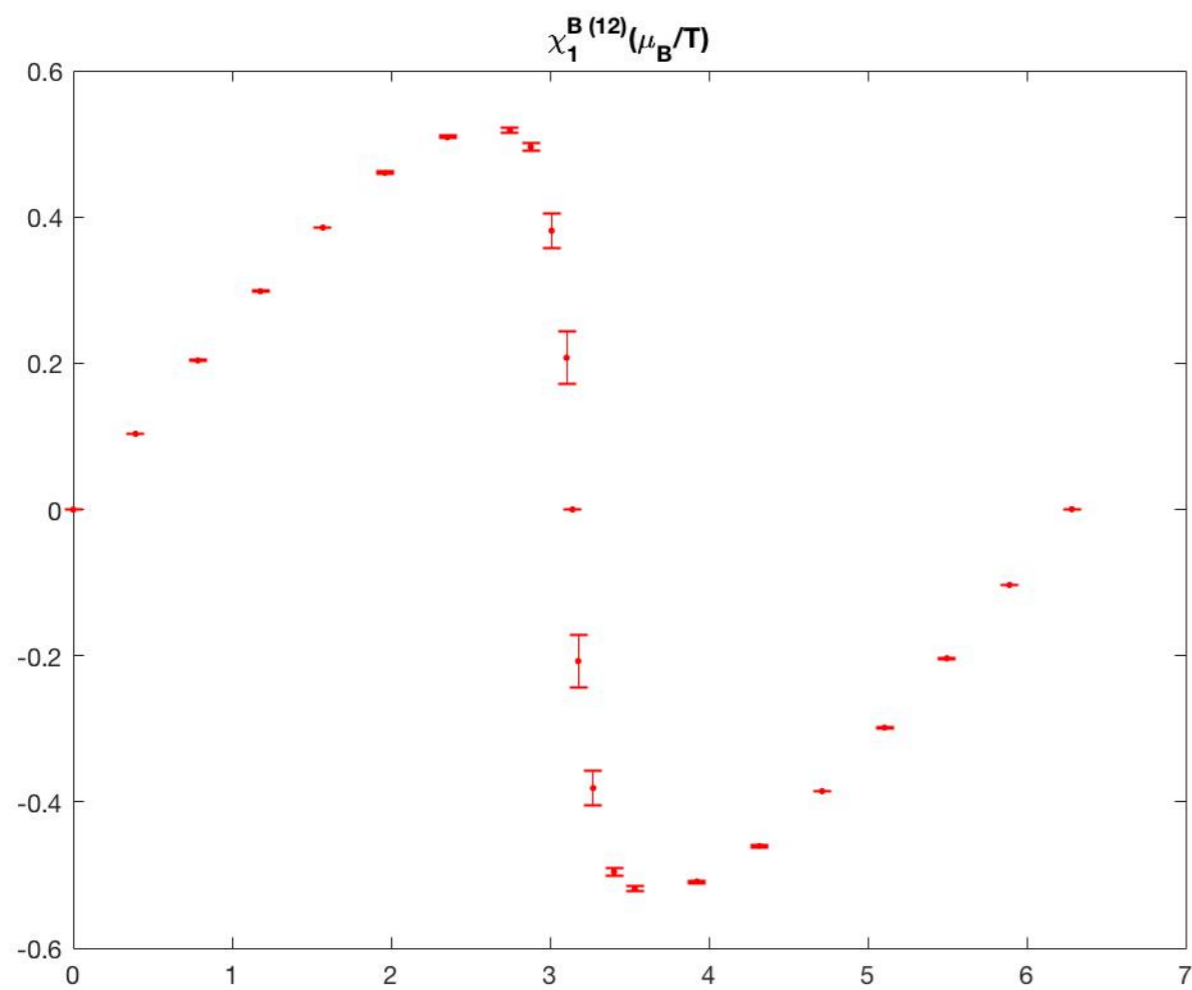
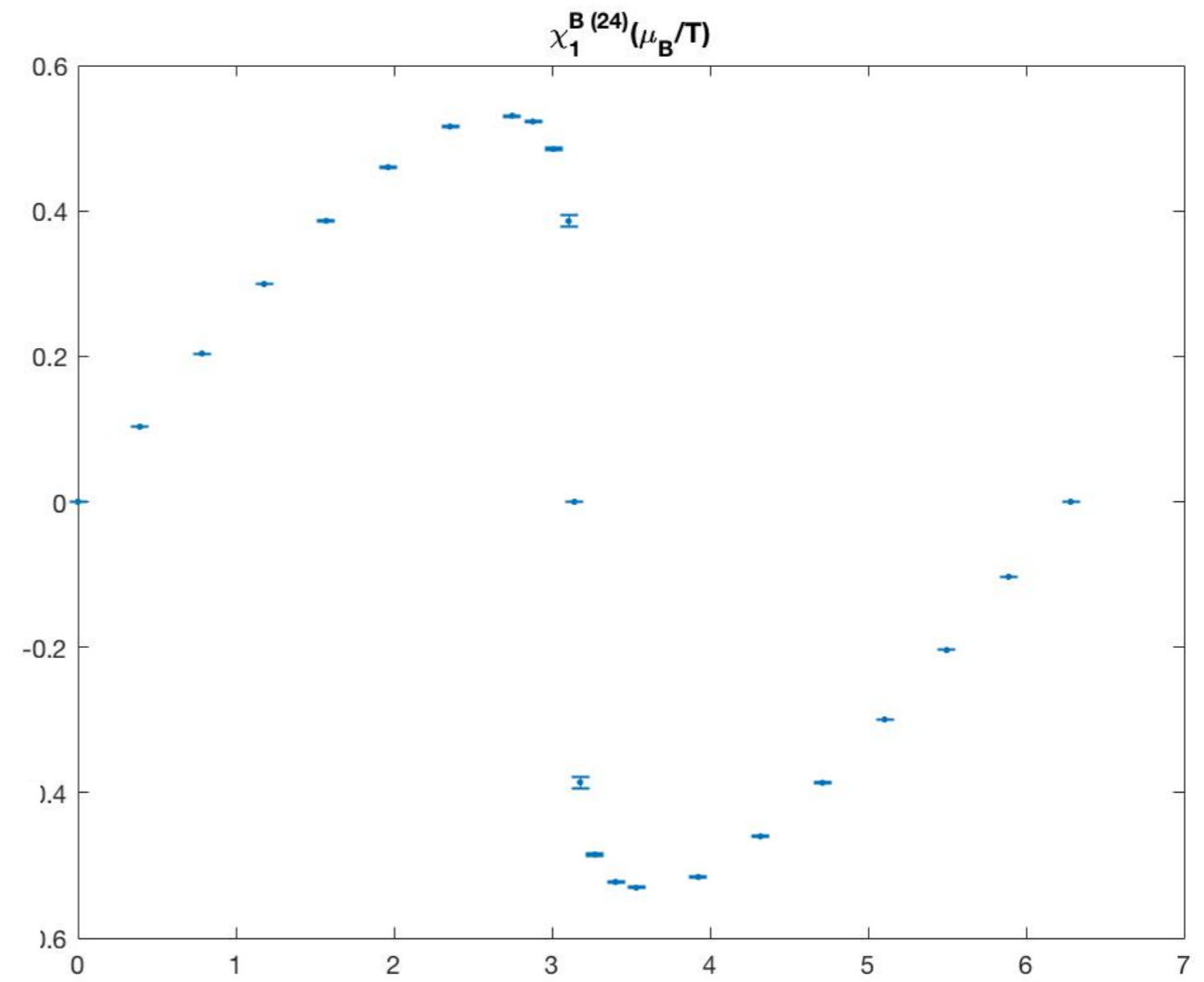
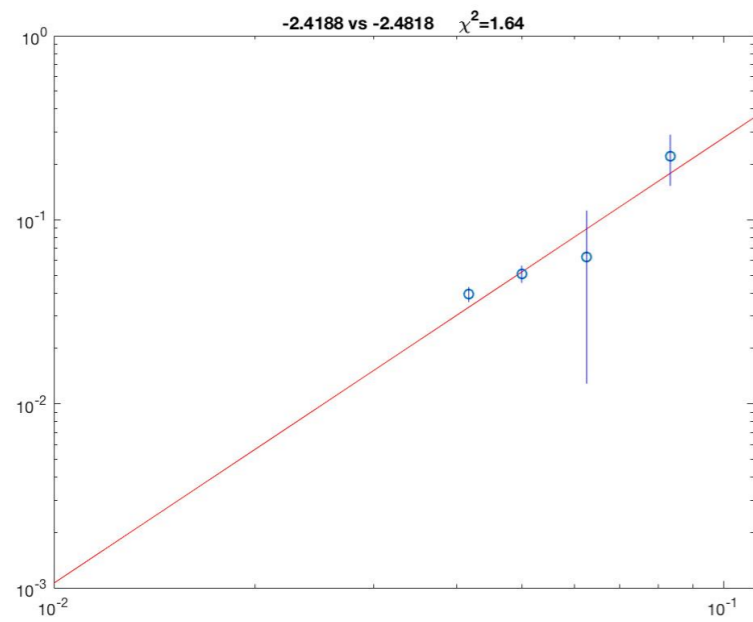
As for the combination of the relevant **critical exponents** we get **2.4188** vs **2.4818**



You have to cheat honestly (C. Michael, private communication)



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We need to enlarge our **statistics** ...

Something still preliminary...

(Also, again,) in view of inspecting finite size scaling, we can play the other game

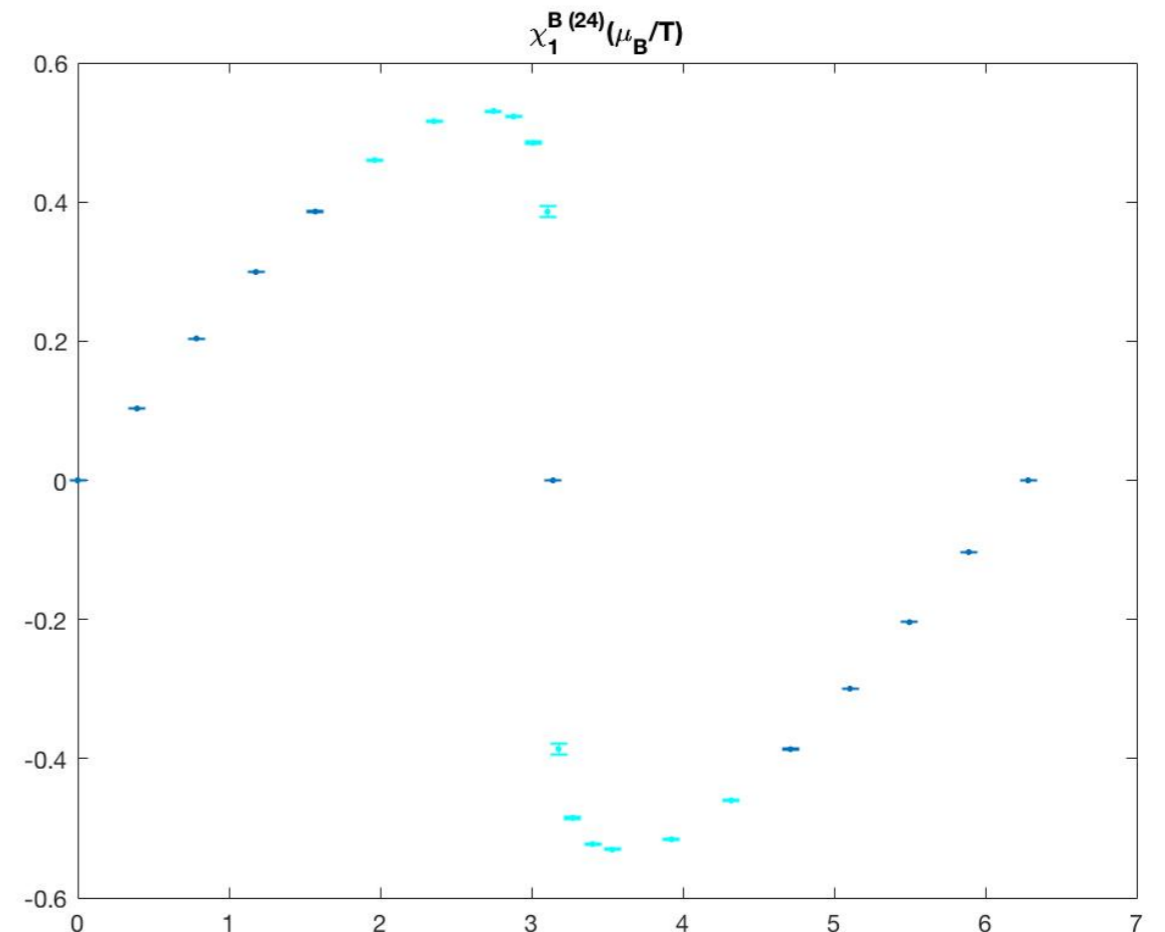
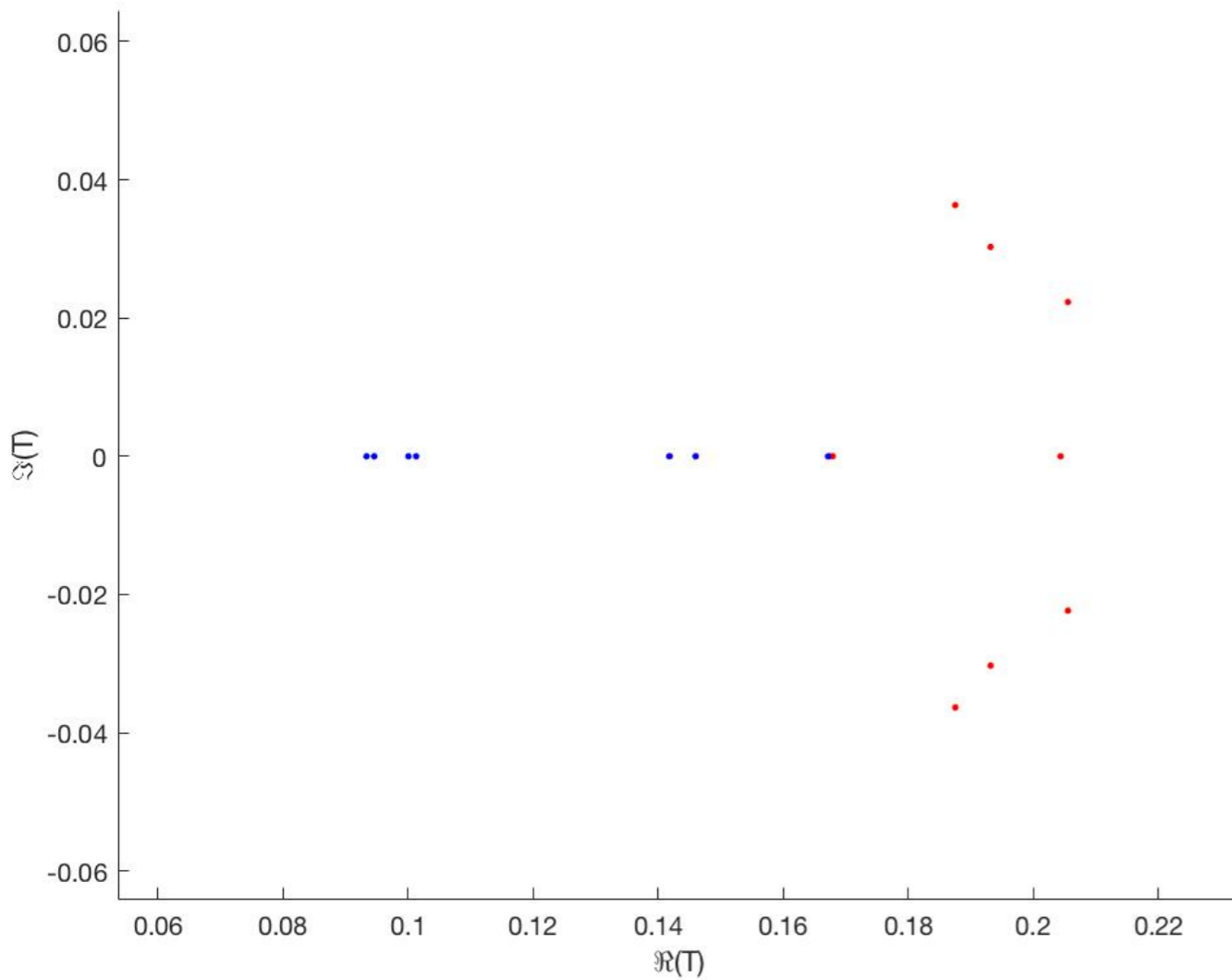
$$f(z; -) \mapsto \chi_1^{B(L)}(T; \mu_B/T) \quad R_m^n(z; -) \mapsto R_m^{n(L)}(T; \mu_B/T)$$

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Beware! Here we are *moving* in **complex T plane**, at **different values of the imaginary chemical potential**, but (at the moment) only **one value (L=24) of lattice size**...

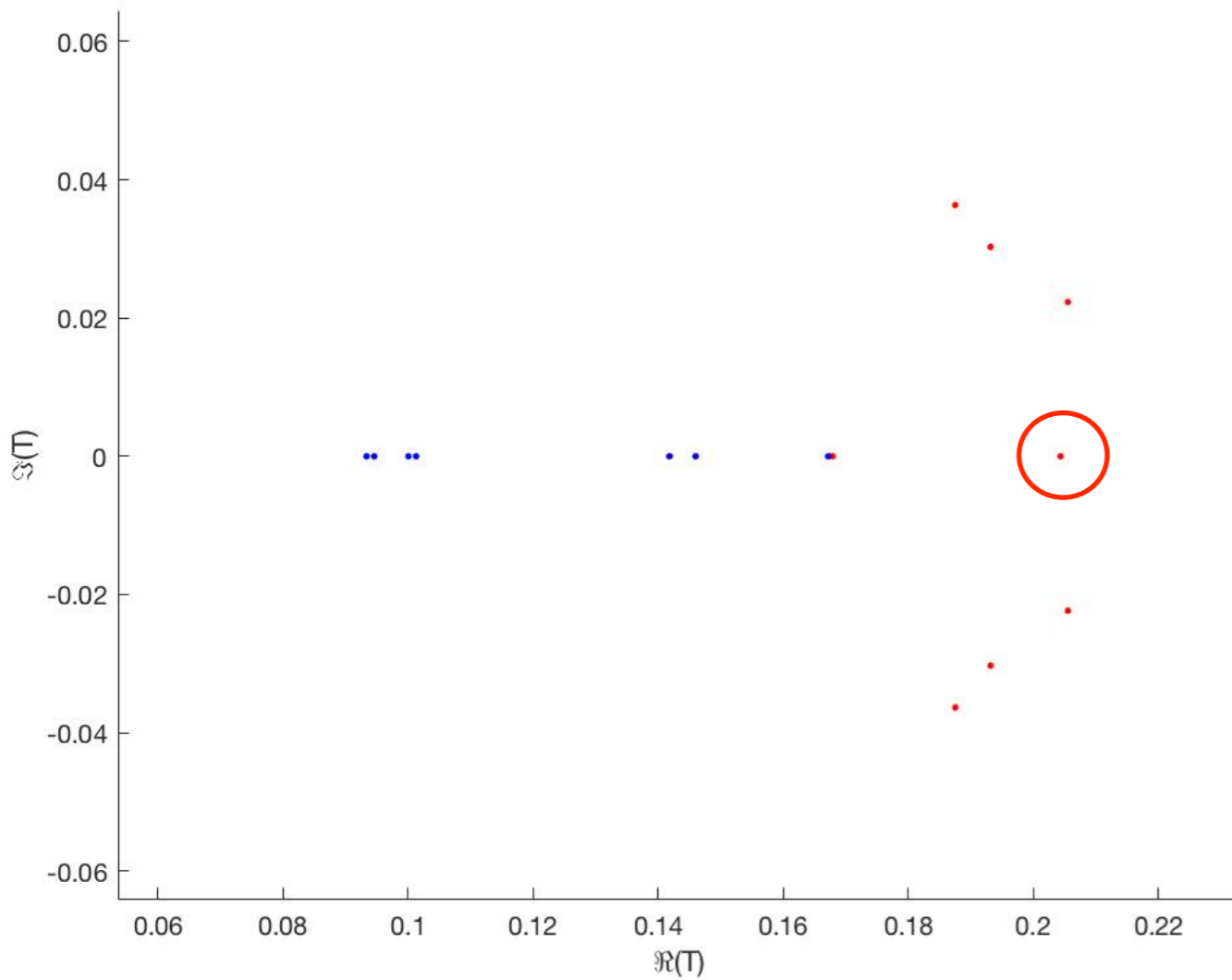


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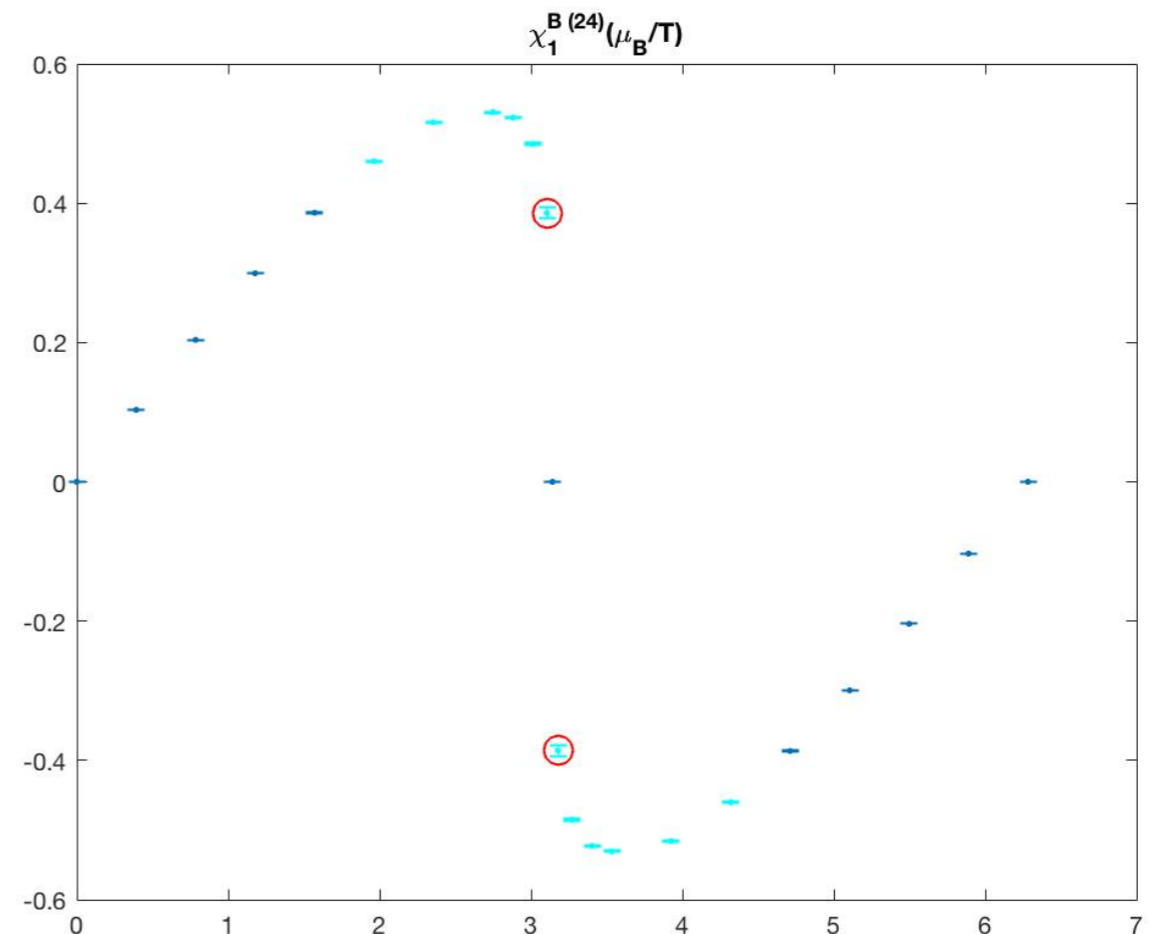
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... making sense ...



CONCLUSIONS

1. The program of (multi-point) Padè analysis in the complex baryonic chemical potential plane could provide interesting informations on Lee Yang edge singularities in QCD. RW seems solid (**TALK BY Christian Schmidt on Fri**); we are trying to better understand *other* transitions... The Holy Grail (needless to say) is the critical point... **MORE ON THIS IN THE TALK BY DAVID CLARKE (Fri)!**
2. Having gained more confidence in the method (from Ising 2d) we now think we can inspect finite size scaling (also) in LQCD.
3. There is much to do for Padé analysis in the complex temperature plane. Results started making sense...