

Pion and kaon generalized form factors from lattice QCD

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Outline

- Introduction
- Theoretical setup – Calculation parameters
- Pion generalized form factors A_{20} , B_{20}
- Kaon generalized form factors A_{20} , B_{20}
- SU(3) flavor symmetry breaking
- Summary

Updated results for scalar, vector, tensor form factors @
APS GHP 2023

<https://indico.jlab.org/event/667/contributions/12289/>

PHYSICAL REVIEW D **103**, 014508 (2021)

Mellin moments $\langle x \rangle$ and $\langle x^2 \rangle$ for the pion and kaon from lattice QCD

Constantia Alexandrou,^{1,2} Simone Bacchio,^{1,2} Ian Cloët,³ Martha Constantinou^{Ⓞ,4},
Kyriakos Hadjiyiannakou,^{1,2} Giannis Koutsou,² and Colin Lauer^{Ⓞ,3,4}

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PHYSICAL REVIEW D **104**, 054504 (2021)

**Pion and kaon $\langle x^3 \rangle$ from lattice QCD and PDF reconstruction
from Mellin moments**

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**Scalar, vector, and tensor form factors for the pion
and kaon from lattice QCD**

Constantia Alexandrou,^{1,2} Simone Bacchio,² Ian Cloët,³ Martha Constantinou^{Ⓞ,4} Joseph Delmar,⁴
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(ETM Collaboration)



Motivation

- Understanding the structure of pion and kaon is as important for describing QCD dynamics as the proton (pions & kaons arise from dynamical chiral symmetry breaking)
- Important for studying SU(3) flavor symmetry breaking
- Pion and kaon less studied than proton:
Most experimental, theoretical, and lattice studies on pion structure regard the vector form factor
- Increased interest in comparing the pion and proton structure (e.g., simultaneous extraction of PDFs and TMDs)
[JAM collaboration: P. Barry et al., arXiv:2302.01192]
- Meson structure under experimental study at JLab 12 GeV (e.g., Measurement of the Charged Pion Form Factor in JLab Hall C)
- Pion and kaon structure will be studied at the Electron-Ion Collider (EIC)
[EIC Yellow Report, arXiv:2103.05419; Aguilar et al., EPJA 55, 190 (2019)]

Important to complement experimental effort from first-principle calculations

Theoretical Setup

- Form factors and generalizations obtained from matrix elements of local operators

$$\mathcal{O}^{\{\mu\nu\}} \equiv \bar{\psi} \left[\frac{1}{2} (\gamma^\mu \vec{D}^\nu + \gamma^\nu \vec{D}^\mu) - \frac{1}{4} \left(\sum_{\rho=1}^4 \delta_{\mu\nu} \gamma^\rho \vec{D}^\rho \right) \right] \psi \quad \mathcal{O}^{\{\mu\nu\rho\}} \equiv \bar{\psi} \gamma^{\{\mu} D^\nu D^\rho \} \psi \quad \mathcal{O}^{\{\mu\nu\rho\sigma\}} \equiv \bar{\psi} \gamma^{\{\mu} D^\nu D^\rho D^\sigma \} \psi$$

- Matrix elements decompose to form factors and generalized form factors [P. Hagler, Phys. Rept. 490, 49 (2010)]

$$\left\langle M(p') | \mathcal{O}^{\{\mu\nu\}} | M(p) \right\rangle = C \left[2P^{\{\mu} P^{\nu\}} A_{20}(Q^2) + 2\Delta^{\{\mu} \Delta^{\nu\}} B_{20}(Q^2) \right]$$

$$B_{20} = A_{22}$$

$$\left\langle M(p') | \mathcal{O}^{\{\mu\nu\rho\}} | M(p) \right\rangle = C \left[2iP^{\{\mu} P^\nu P^\rho \} A_{30}(Q^2) + 2i\Delta^{\{\mu} \Delta^\nu P^\rho \} B_{30}(Q^2) \right]$$

$$B_{30} = A_{32}$$

$$\left\langle M(p') | \mathcal{O}^{\{\mu\nu\rho\sigma\}} | M(p) \right\rangle = C \left[-2P^{\{\mu} P^\nu P^\rho P^\sigma \} A_{40}(Q^2) - 2\Delta^{\{\mu} \Delta^\nu P^\rho P^\sigma \} B_{40}(Q^2) - 2\Delta^{\{\mu} \Delta^\nu \Delta^\rho \Delta^\sigma \} C_{40}(Q^2) \right]$$

$$B_{40} = A_{42}$$

$$C_{40} = A_{44}$$

- Form factors are frame independent
- Two methods to extract ground-state contribution
 - fit plateau region to a constant value
 - perform a two-state fit on three-point function
- Non-perturbative renormalization ($\overline{\text{MS}}$ at 2 GeV)

Lattice Details

- Two ensembles of twisted-clover fermions and Iwasaki improved gluons [\[C. Alexandrou et al., PRD 104, 074520 \(2021\)\]](#)

Ensemble	β	$a(\text{fm})$	$L^3 \times T$	N_f	$m_\pi(\text{MeV})$	$L(\text{fm})$
cA211.30.32	1.726	0.094	$32^3 \times 64$	2 + 1 + 1	265	3.0
cB211.25.48	1.778	0.079	$48^3 \times 96$	2 + 1 + 1	250	3.79

- Calculation of connected contributions

- Use of different kinematic frames:

-- Rest frame: desirable for charges and usual form factors

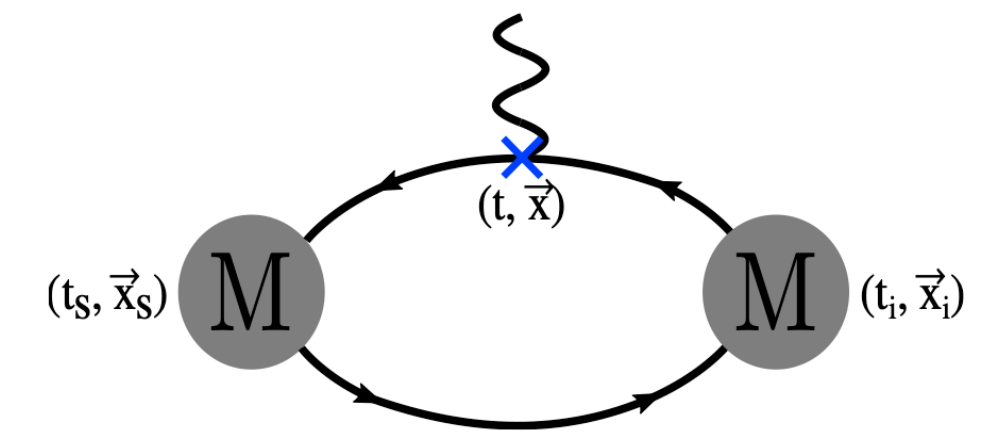
-- Boosted frame: necessary for higher Mellin moments of GPDs without power-divergent mixing

- Advantage of boosted frame: access to a denser range of $-t$ $(-t = \vec{q} - (E(p') - E(p))^2)$

- Boosted frame of choice: fixed final momentum

(use of sequential method: matrix element for all values of \vec{q} are obtained at once)

$$\vec{p}' = \frac{2\pi}{L}(\pm 1, \pm 1, \pm 1), \quad \vec{p} = \vec{p}' - \vec{q} \quad \vec{q} = \frac{2\pi}{L}(\pm n_x, \pm n_y, \pm n_z) \quad n_x, n_y, n_z \in [0, 8]$$

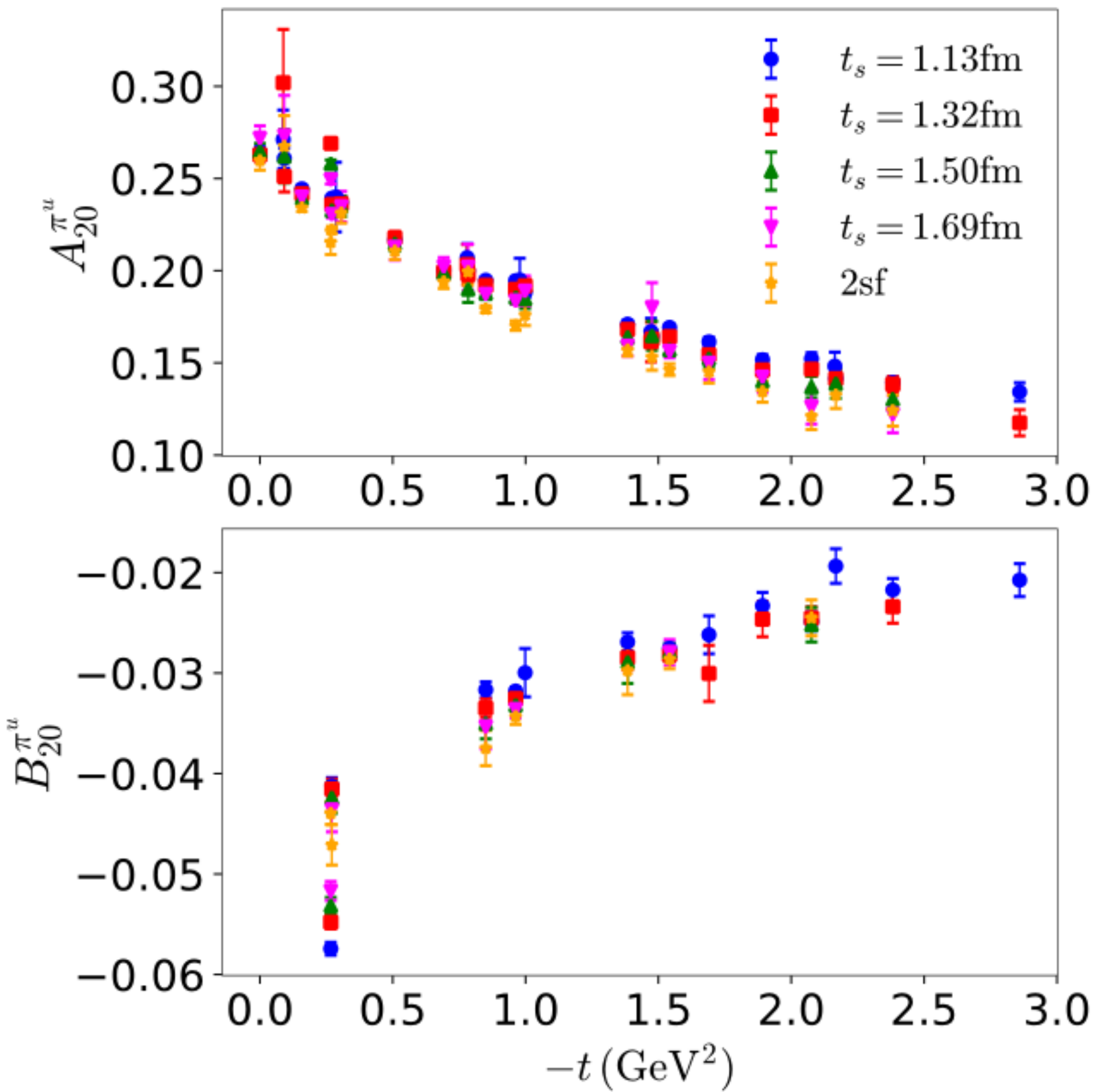


Statistics

Ensemble	Frame	$\vec{p}(2\pi/L)$	t_s/a	t_s (fm)	confs	src pos.	Total
cA211.30.32	R	(0, 0, 0)	12, 14, 16, 18, 20, 24	1.13, 1.32, 1.50, 1.69, 1.88, 2.256	122	16	1,952
cA211.30.32	B	($\pm 1, \pm 1, \pm 1$)	12, 14, 16, 18	1.13, 1.32, 1.50, 1.69	122	136	132,736
cB211.25.48	B	($\pm 1, \pm 1, \pm 1$)	14, 16, 18, 20	1.11, 1.26, 1.42, 1.58	76	8	4,864

Pion Generalized Form Factors

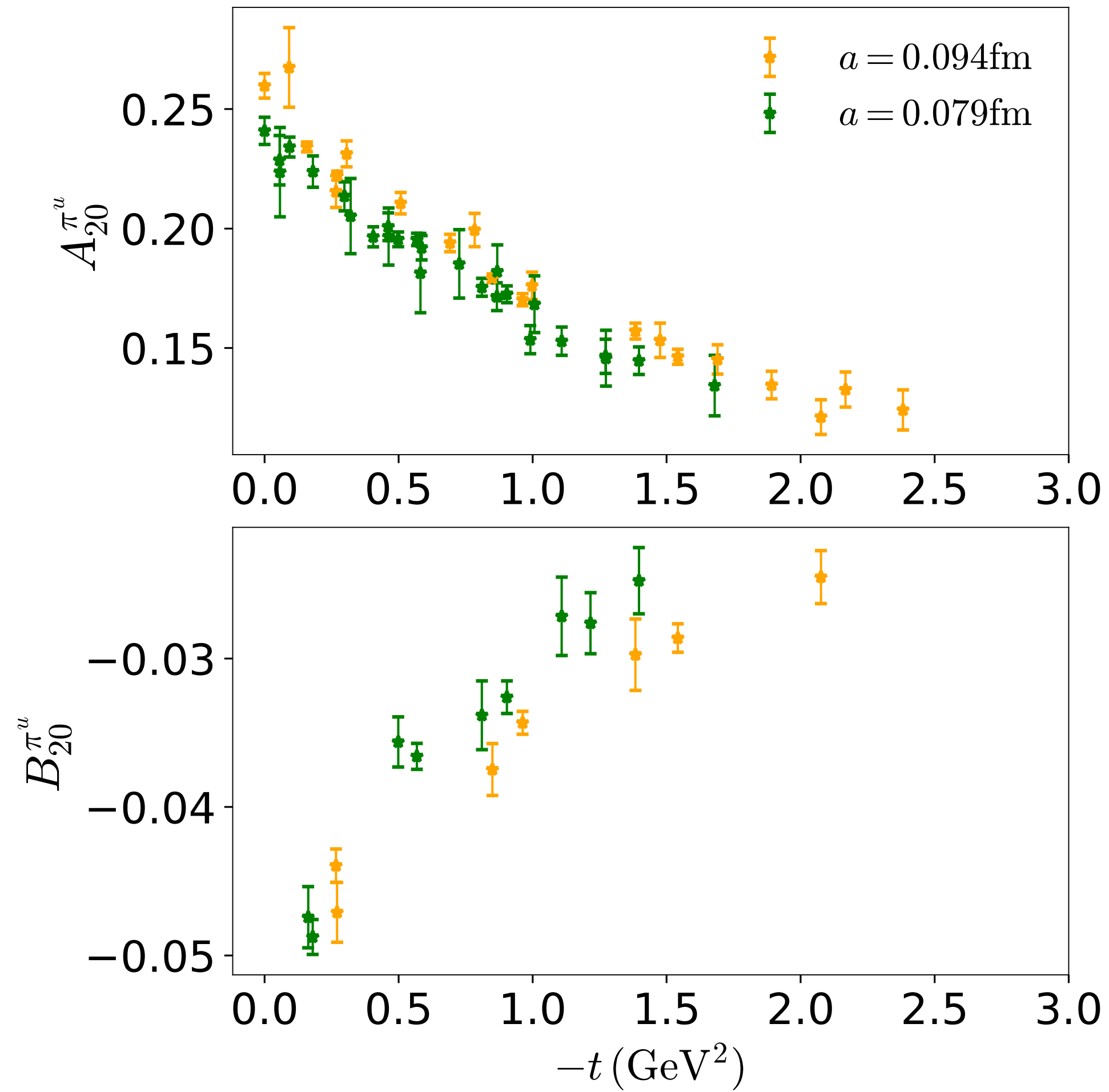
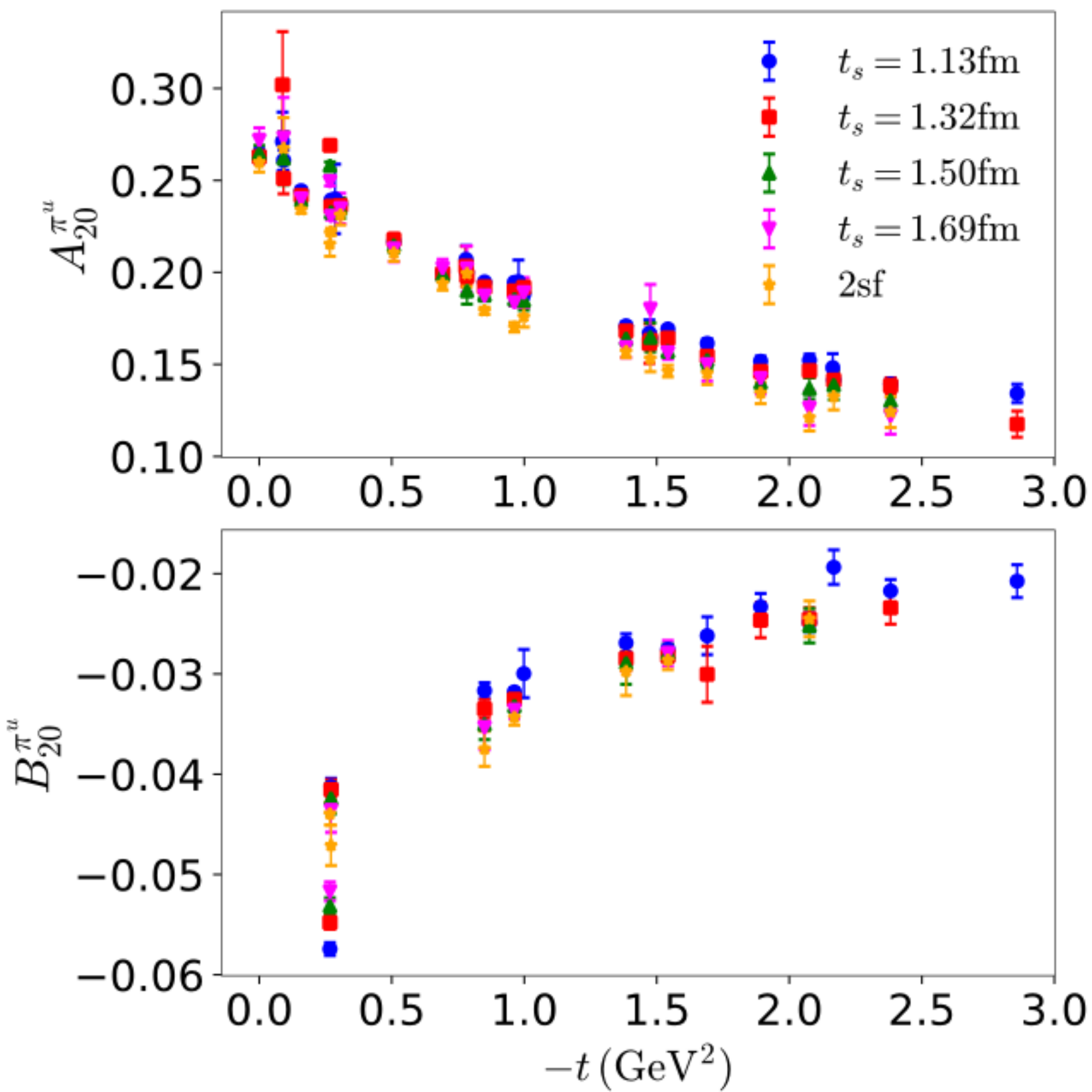
Pion Generalized Form Factors



- Increase of statistical error is not linear with \vec{q} ($-t = \vec{q}^2 - (E(p') - E(p))^2$)

→ a careful analysis required to select data points with controlled uncertainties

Pion Generalized Form Factors



Comparison of
2-state fits:

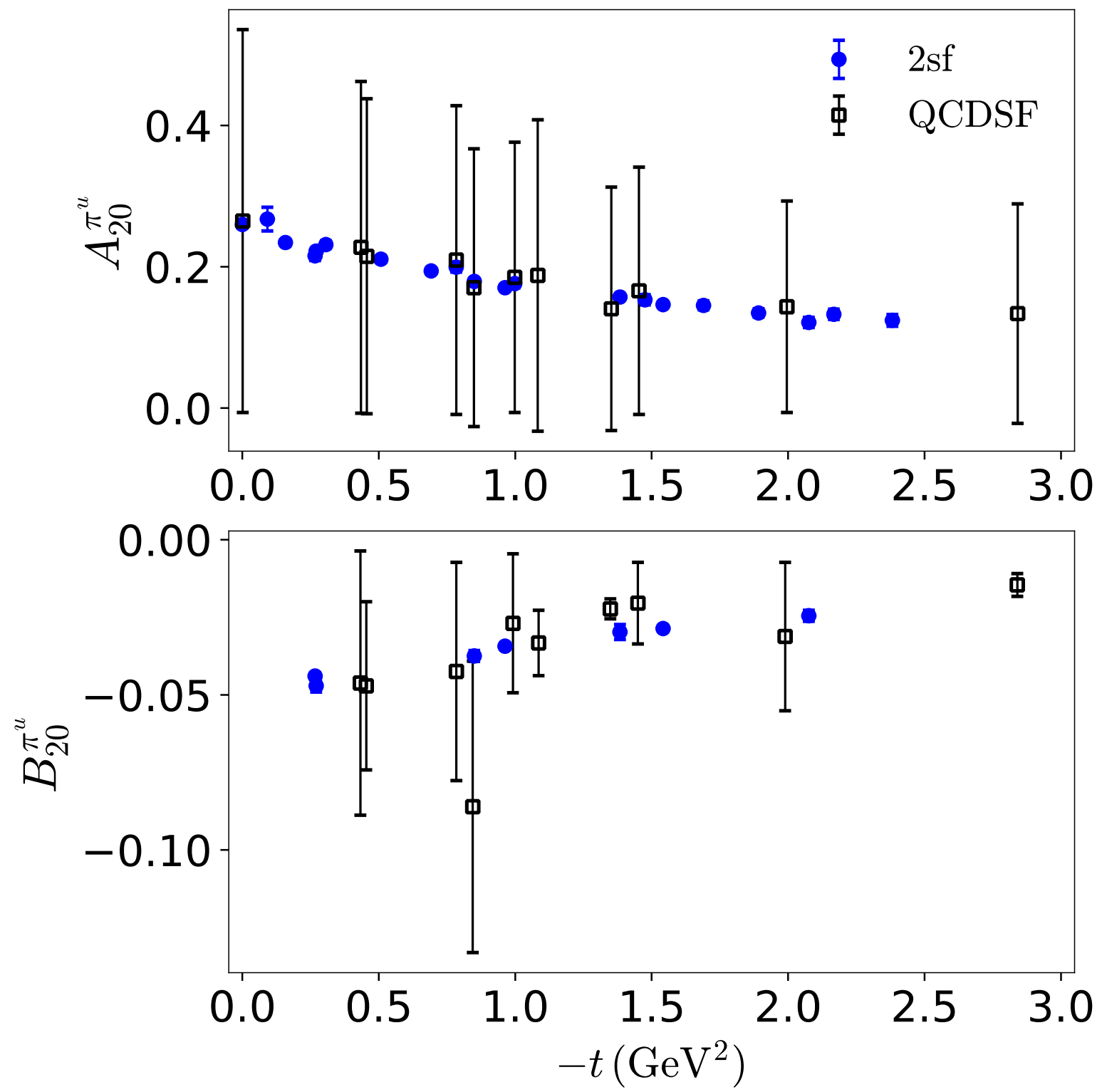
- Systematic effects more prominent in B_{20}
- Discretization effects less significant than scalar, vector, tensor FFs

- Increase of statistical error is not linear with \vec{q} ($-t = \vec{q}^2 - (E(p') - E(p))^2$)

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Pion Generalized Form Factors

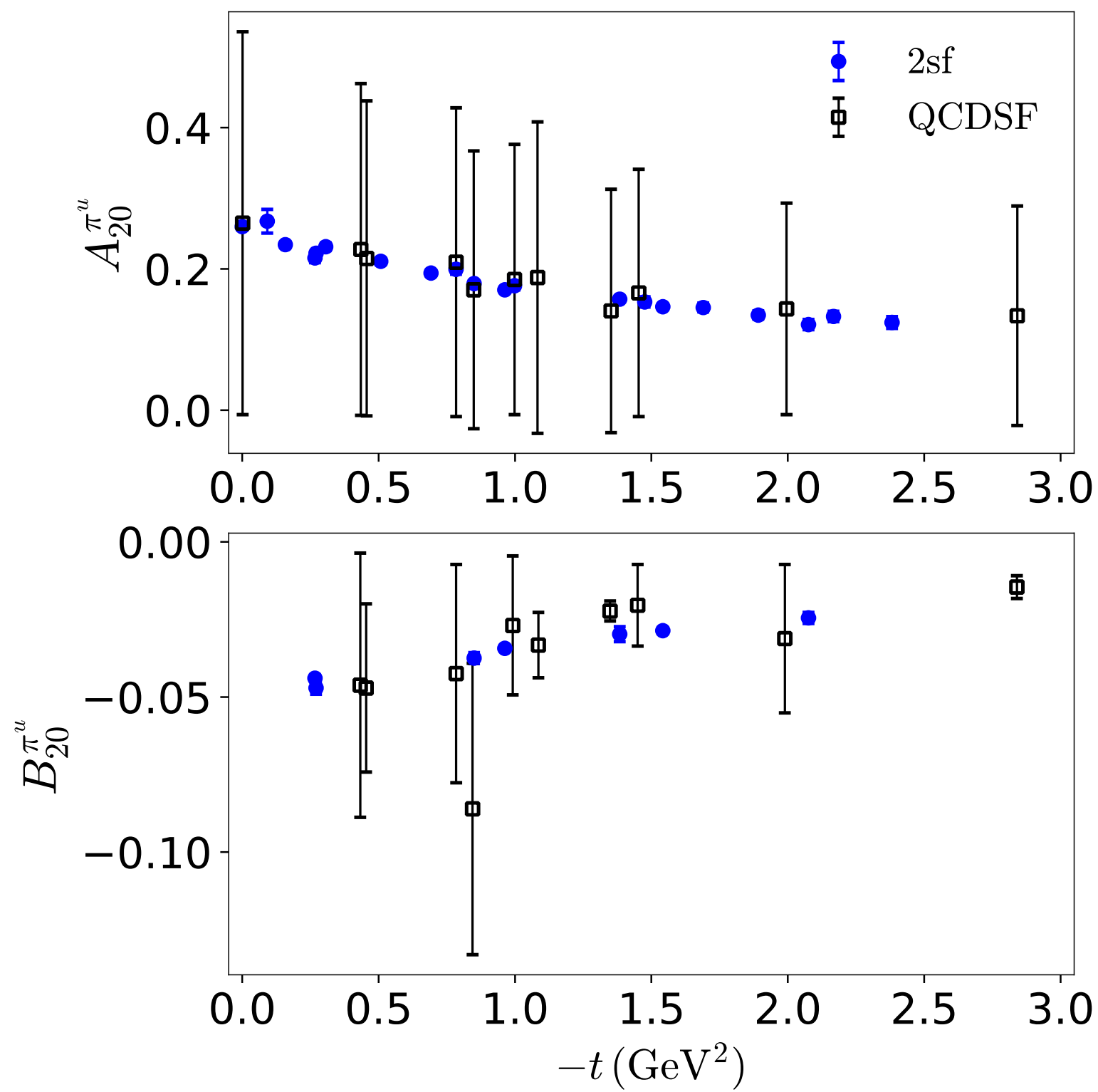
Pion Generalized Form Factors



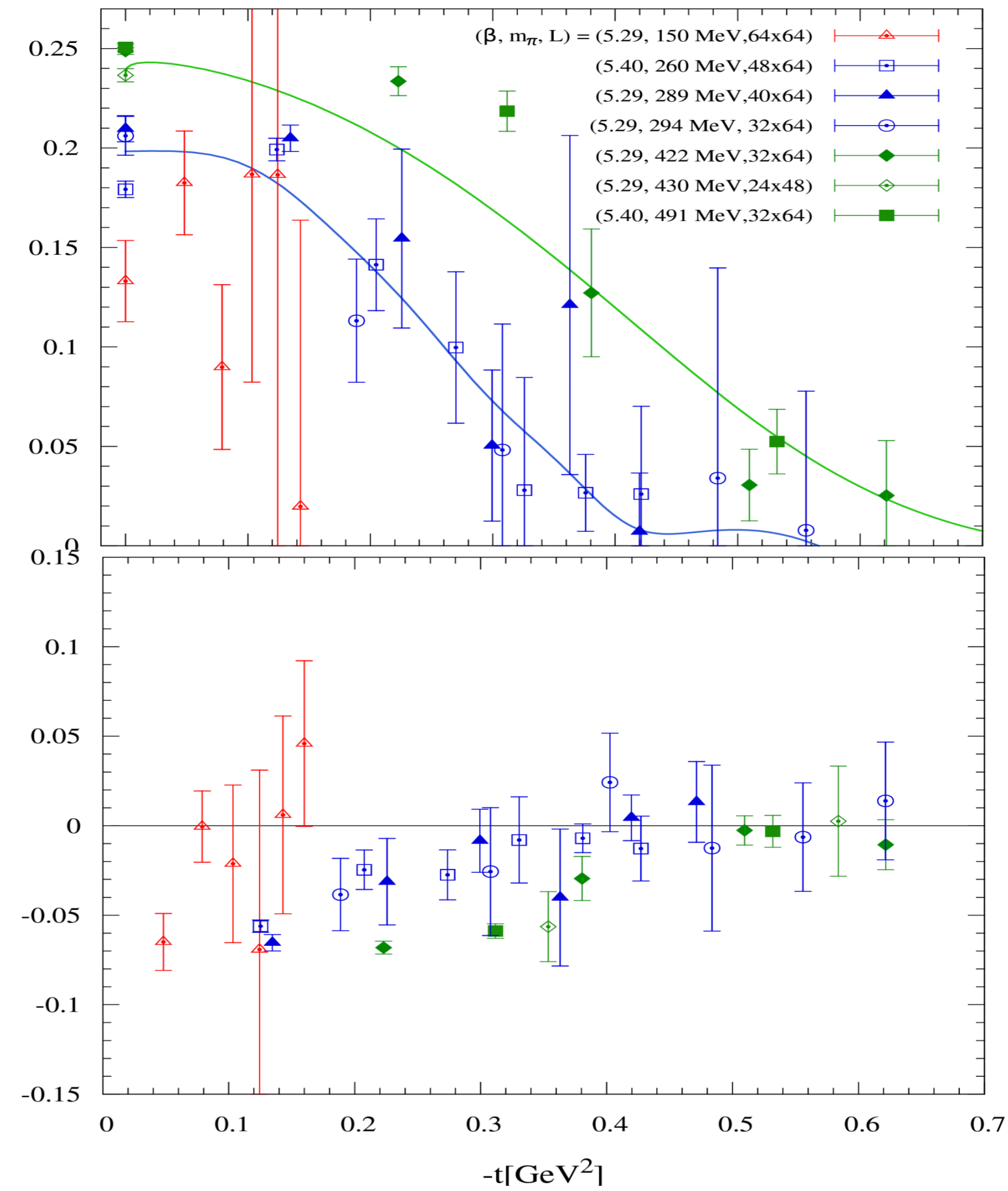
- Results fully compatible with QCDSF [PoS(LAT2005)360]

Pion Generalized Form Factors

Comparison with other studies



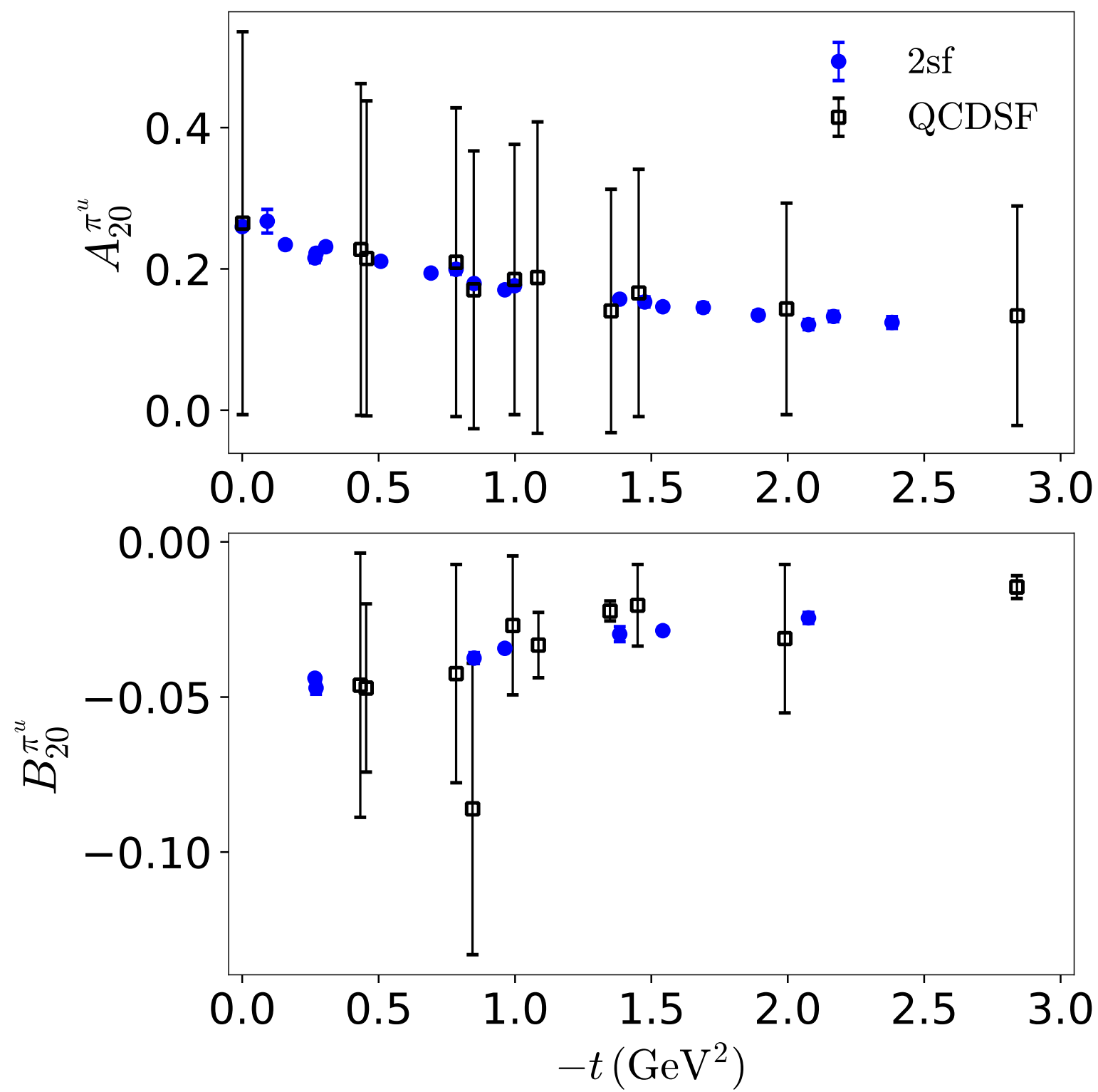
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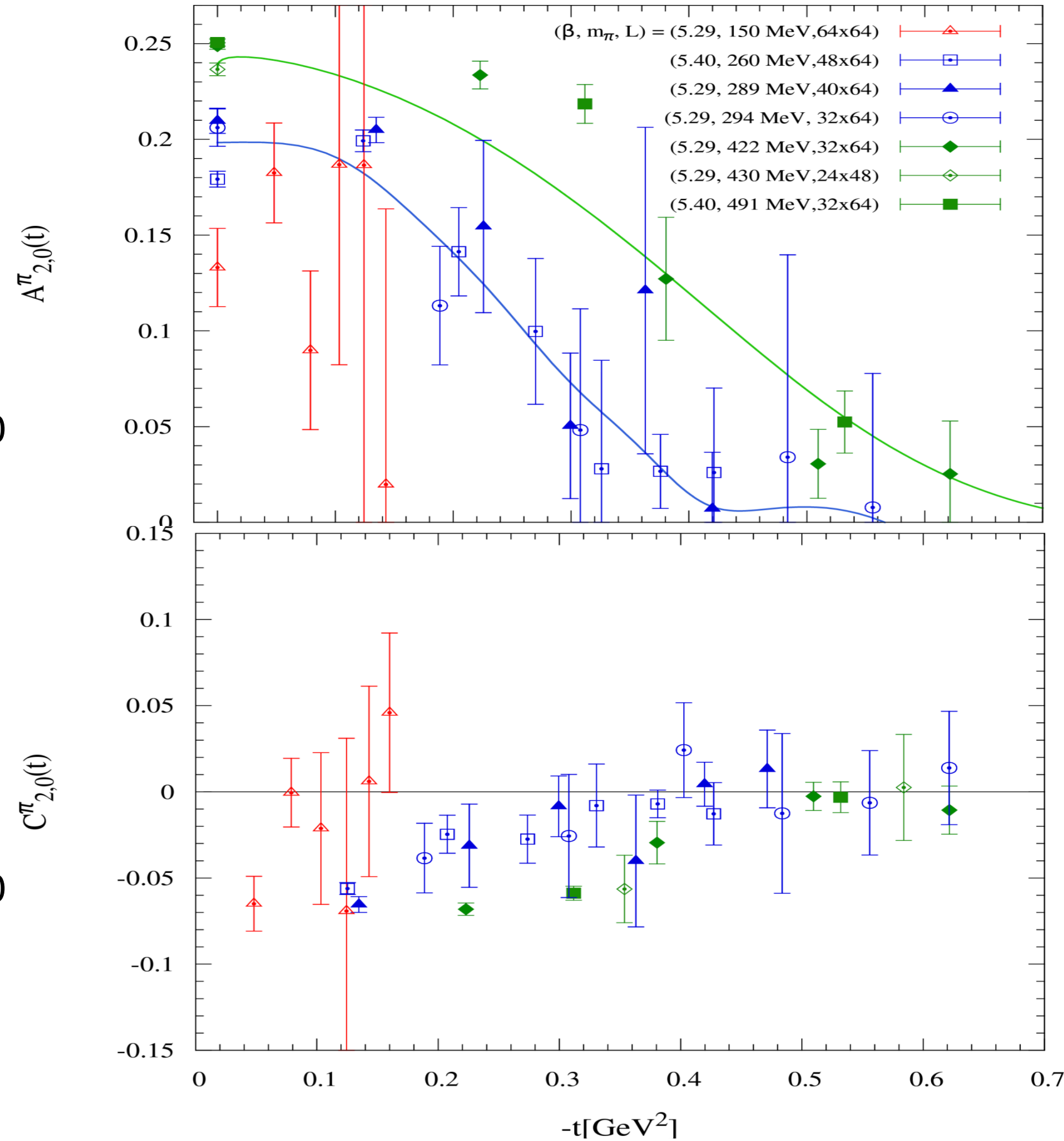
- RQCD demonstrates large systematic uncertainties [PoS LAT2013 (2014) 447]

Pion Generalized Form Factors

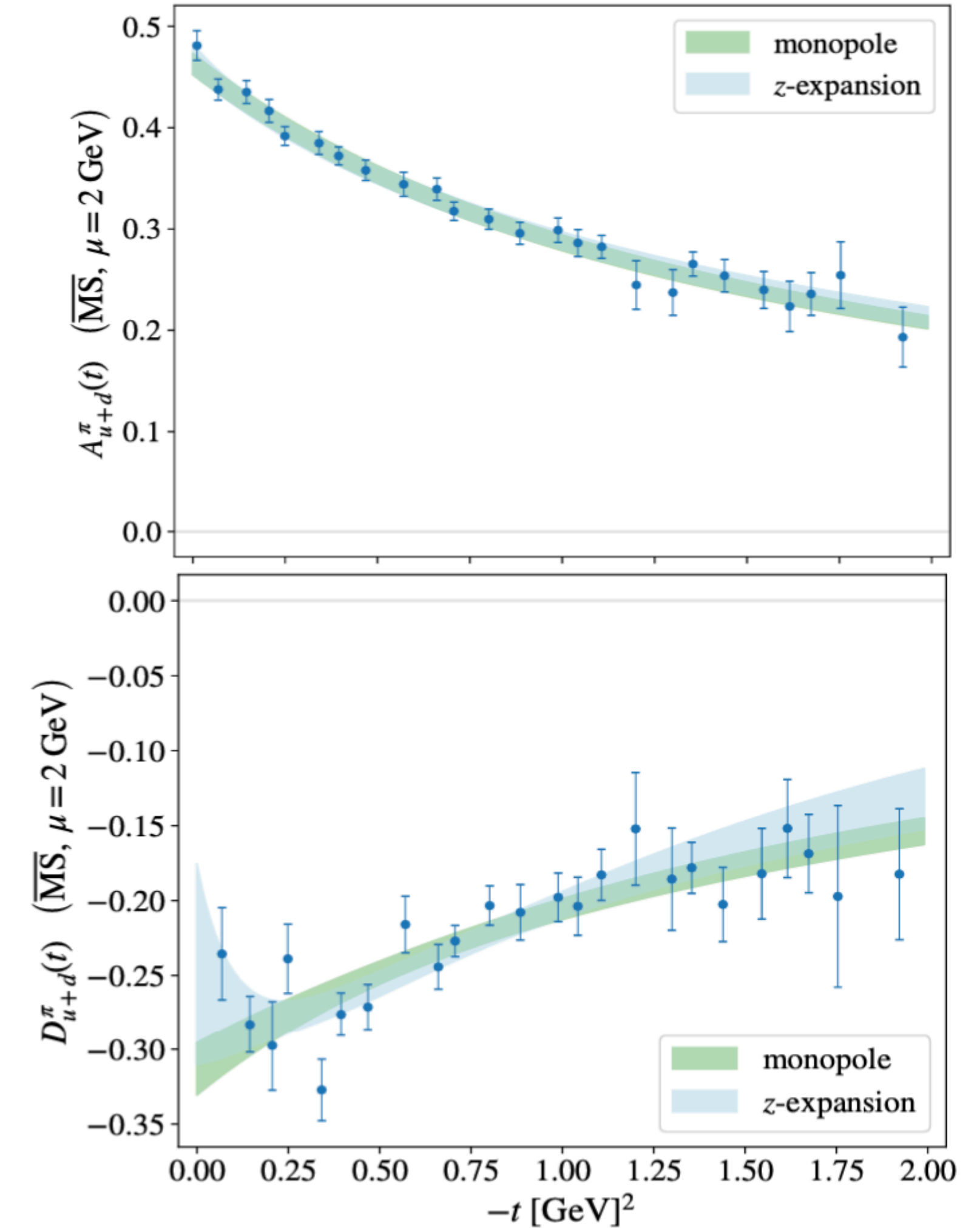
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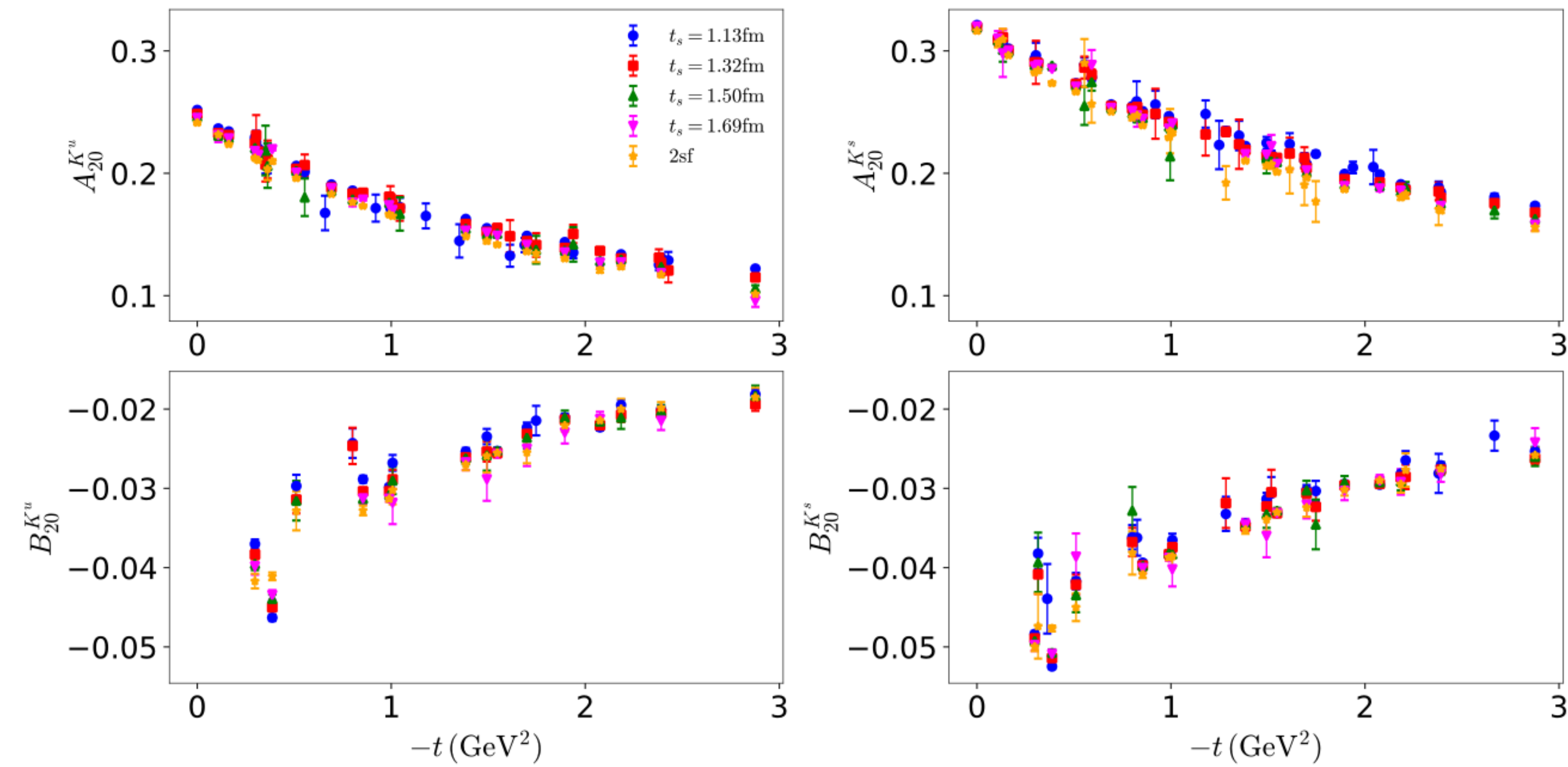
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- Work by MIT (arXiv:2307.11707) includes disconnected contributions and elimination of mixing with gluon contributions (see talk by D. Pefkou)

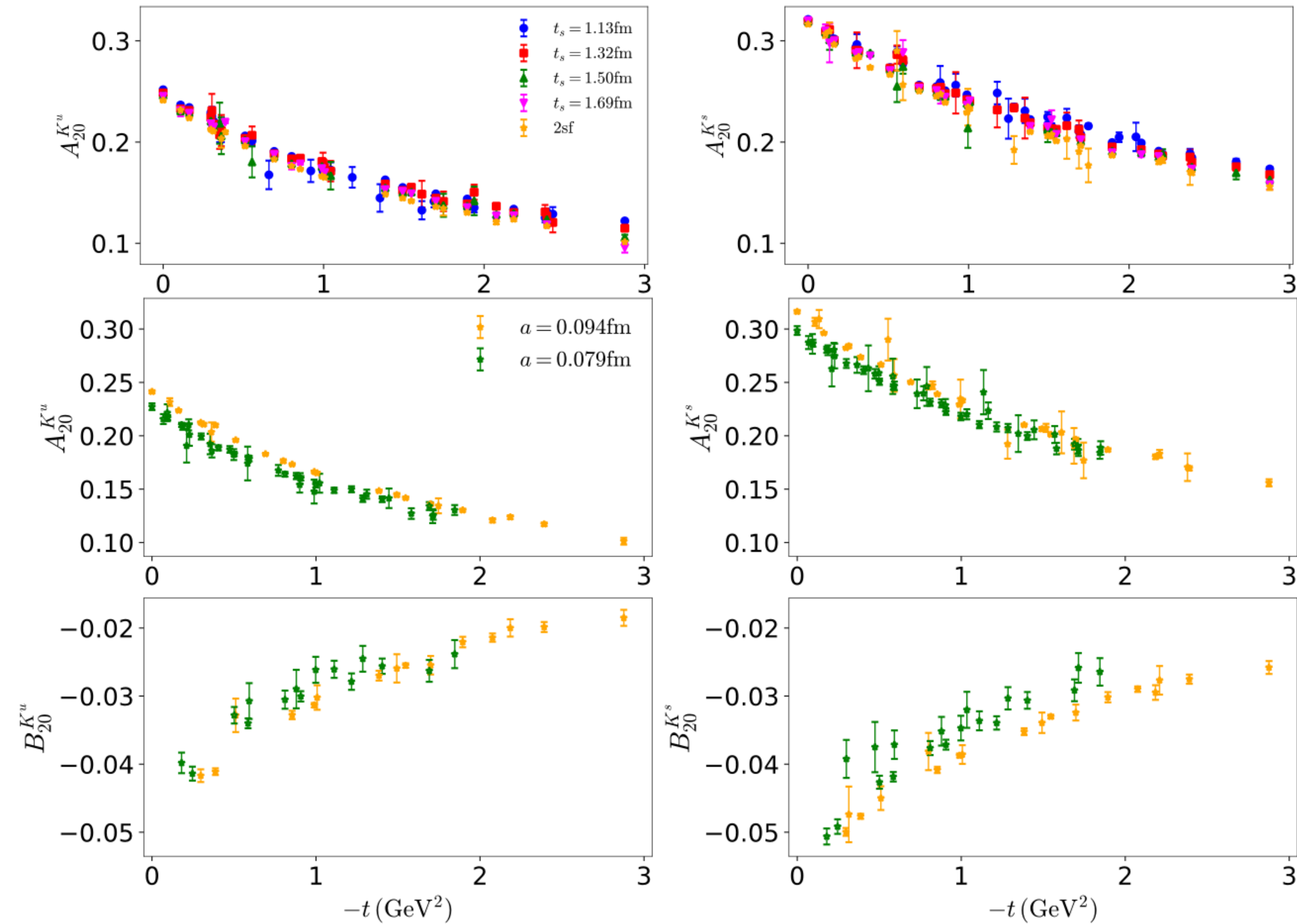
Kaon Generalized Form Factors

Kaon Generalized Form Factors



- Less statistical noise compared to pion
- Dense range of $-t$ values
- Strange contributions dominant in magnitude
- Small excited state effects (typically within errors)

Kaon Generalized Form Factors



- Less statistical noise compared to pion
- Dense range of $-t$ values
- Strange contributions dominant in magnitude
- Small excited state effects (typically within errors)
- Accuracy of data reveals systematic effects in both A_{20} and B_{20}

Pion Parameterization

- Generalized form factors provide information on QCD dynamics (eg. SU(3) flavor symmetry effects)

We parameterize the $-t$ dependence using a monopole (n=1) and an n-pole fit $F_{\Gamma}(Q^2) = \frac{F_{\Gamma}(0)}{\left(1 + \frac{Q^2}{M_{\Gamma}^2}\right)^n}$

- Other fit functions to be explored (e.g., z-expansion)

- Examined effect of $-t$ range included in fit ($\sim 1 \text{ GeV}^2$, $\sim 3 \text{ GeV}^2$)

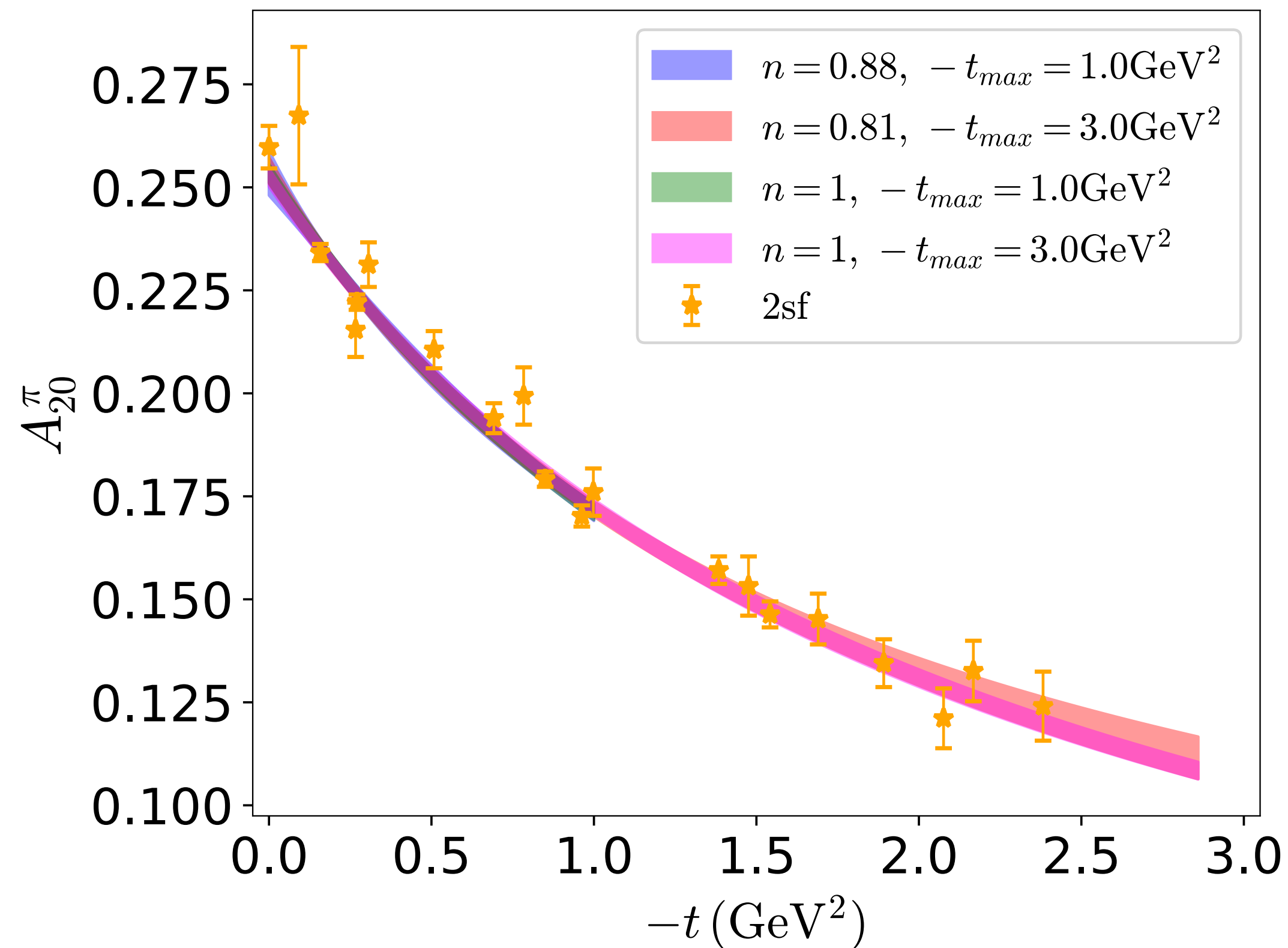
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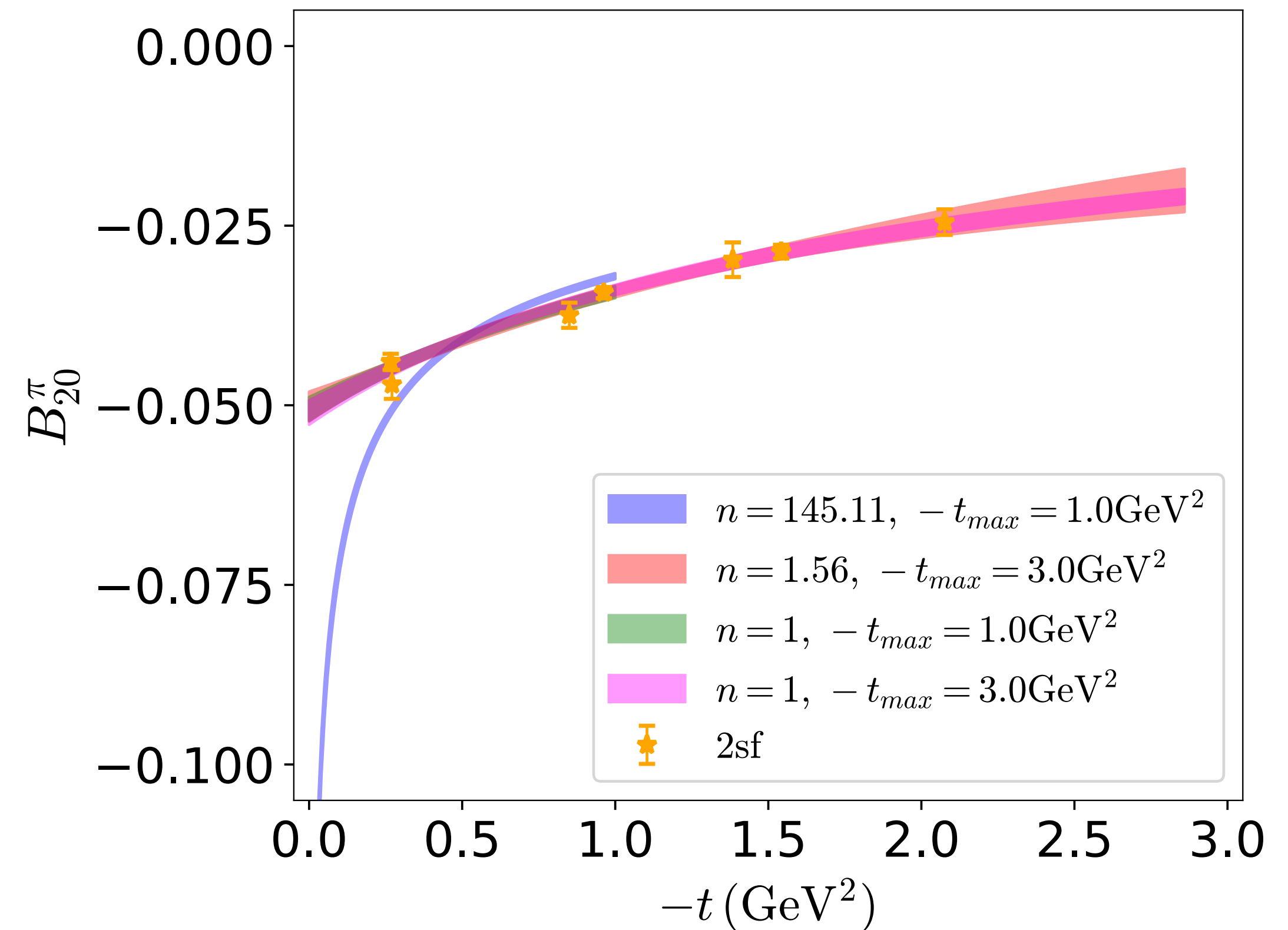
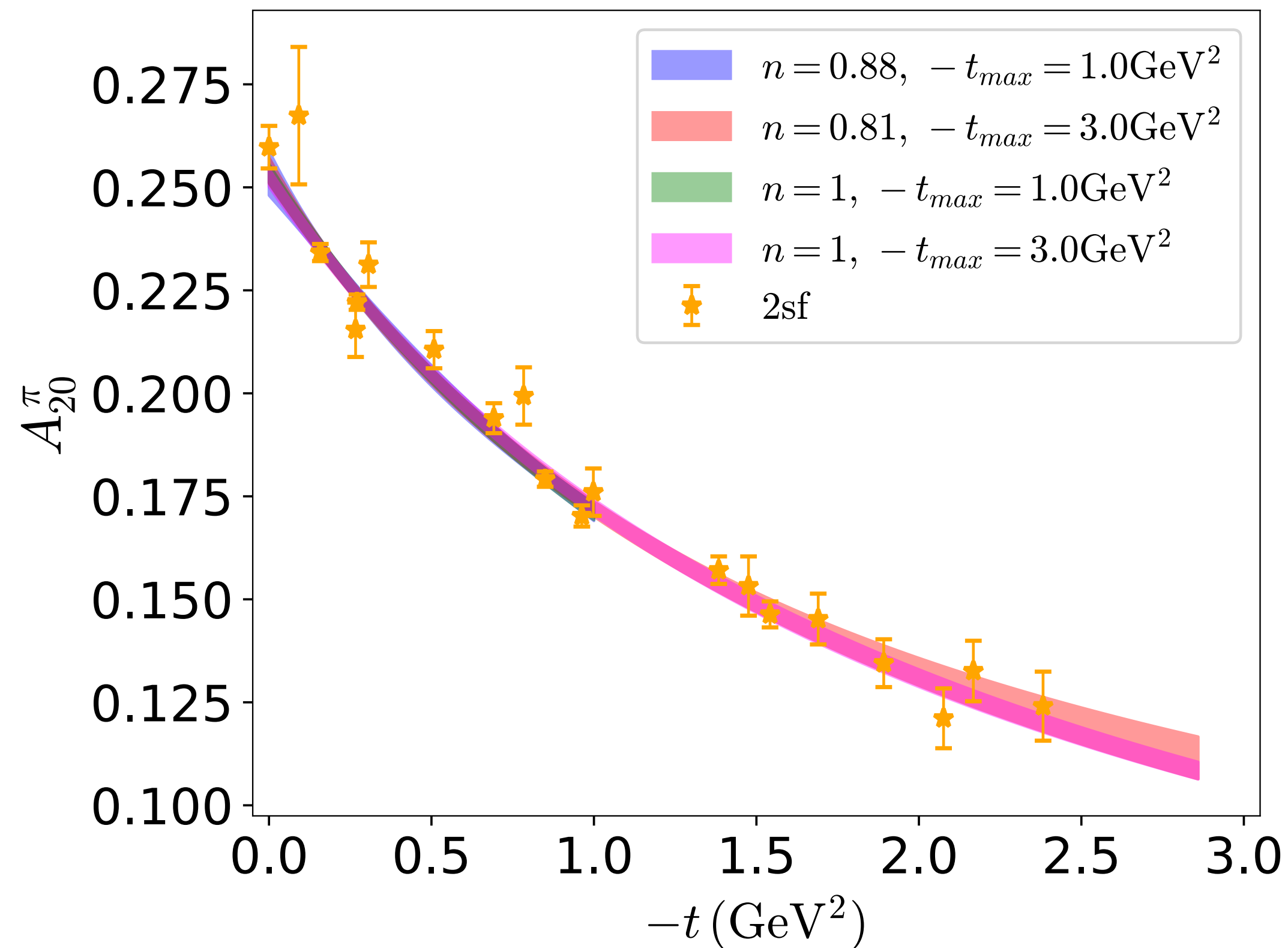
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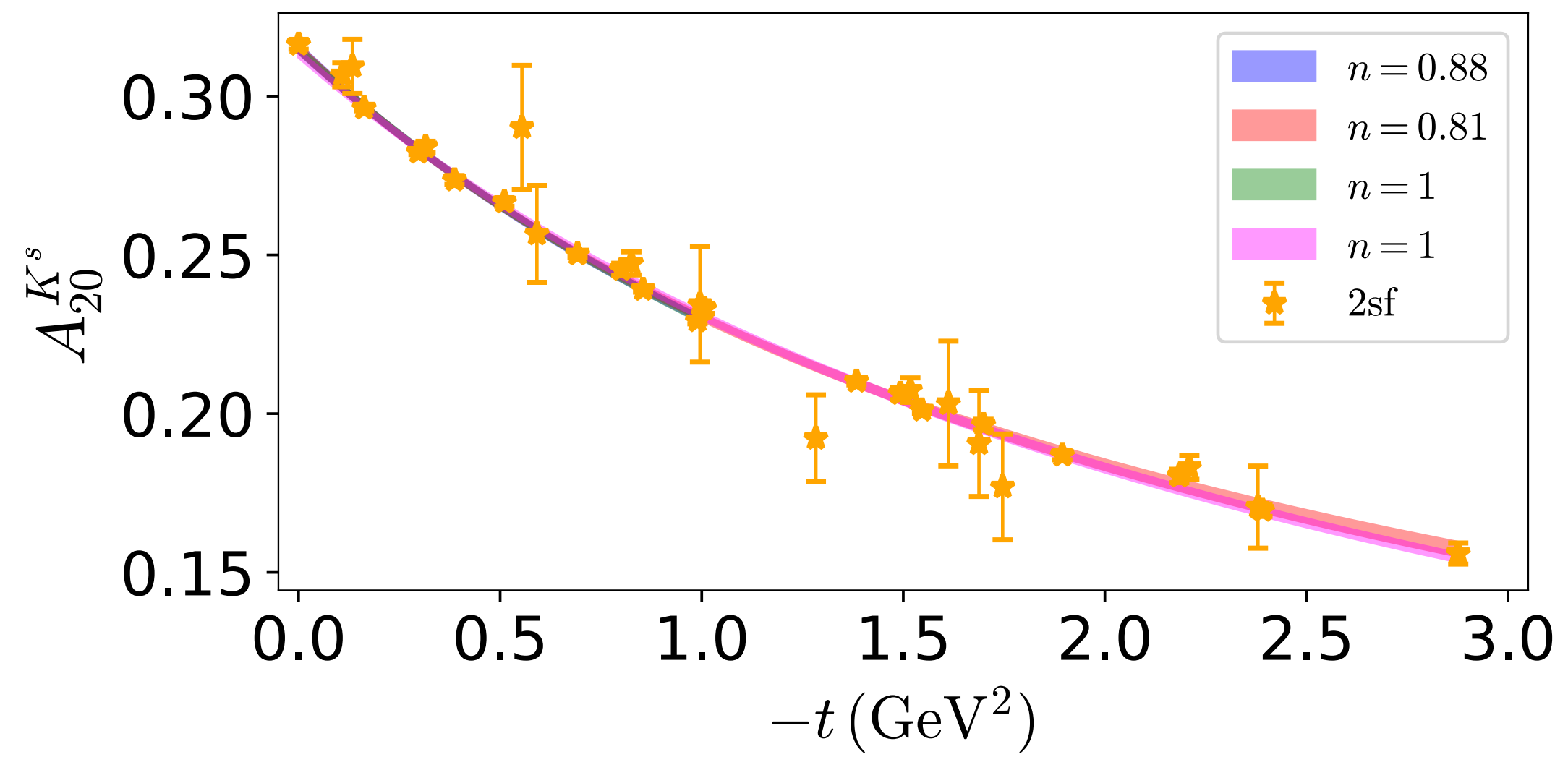
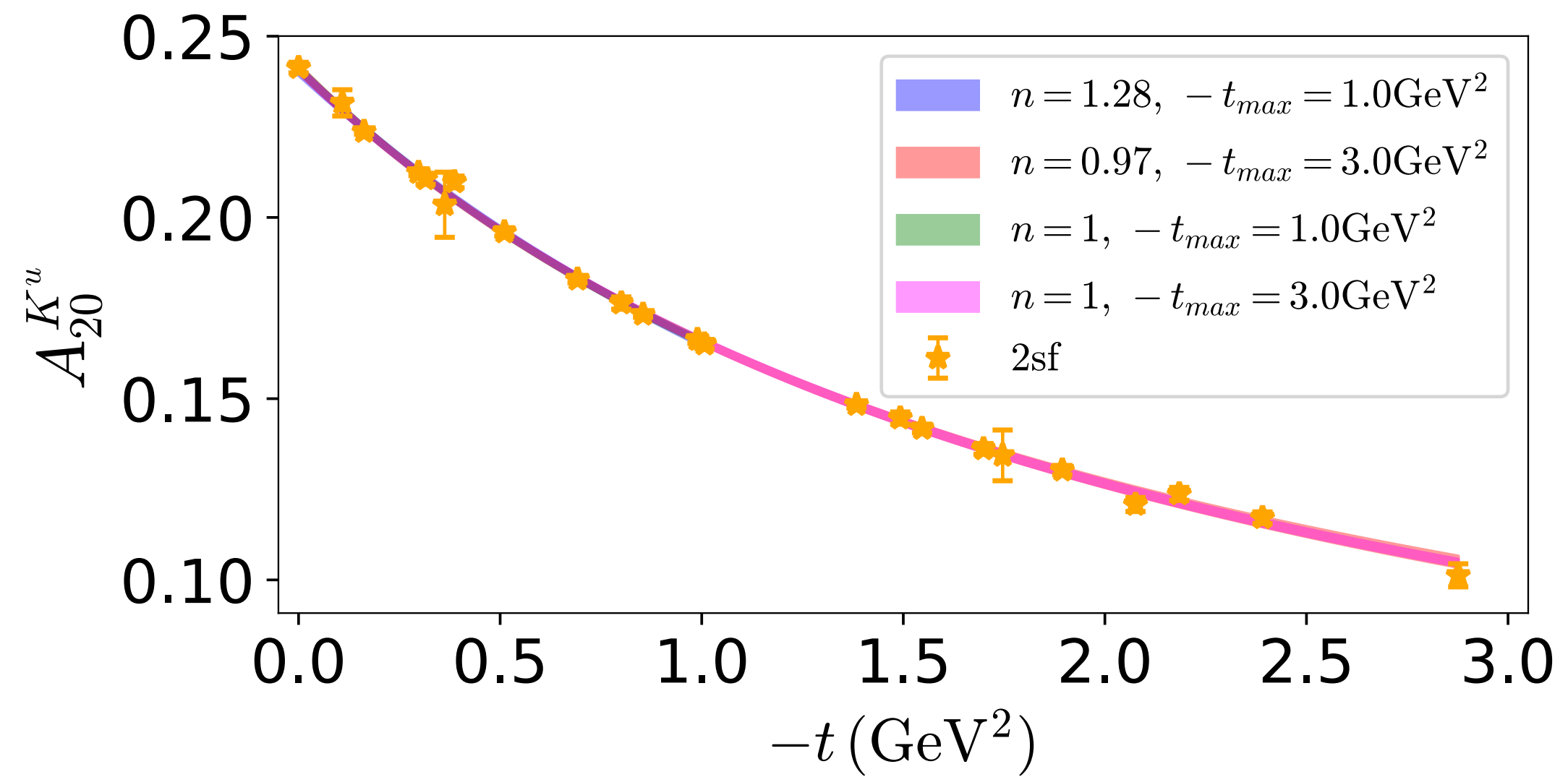
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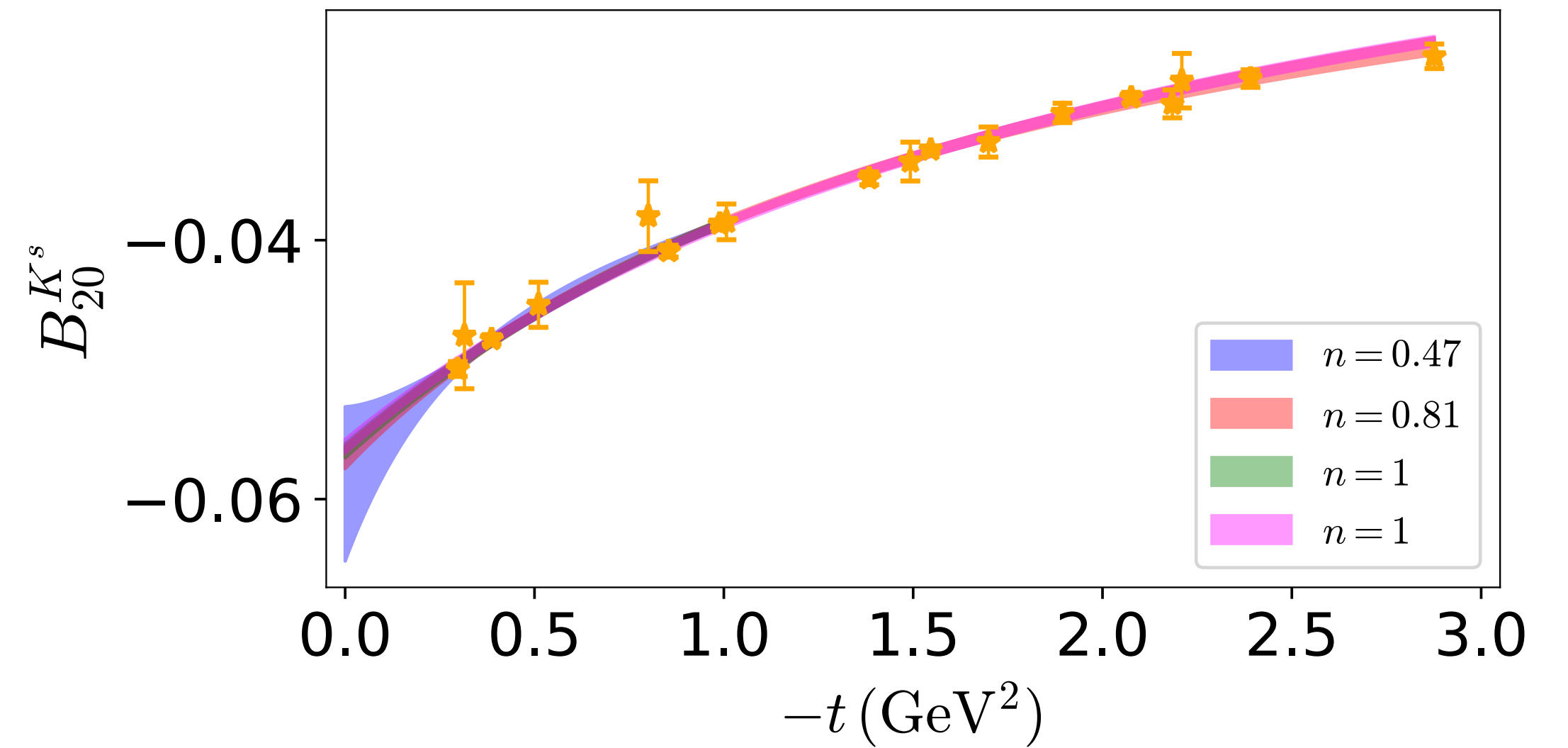
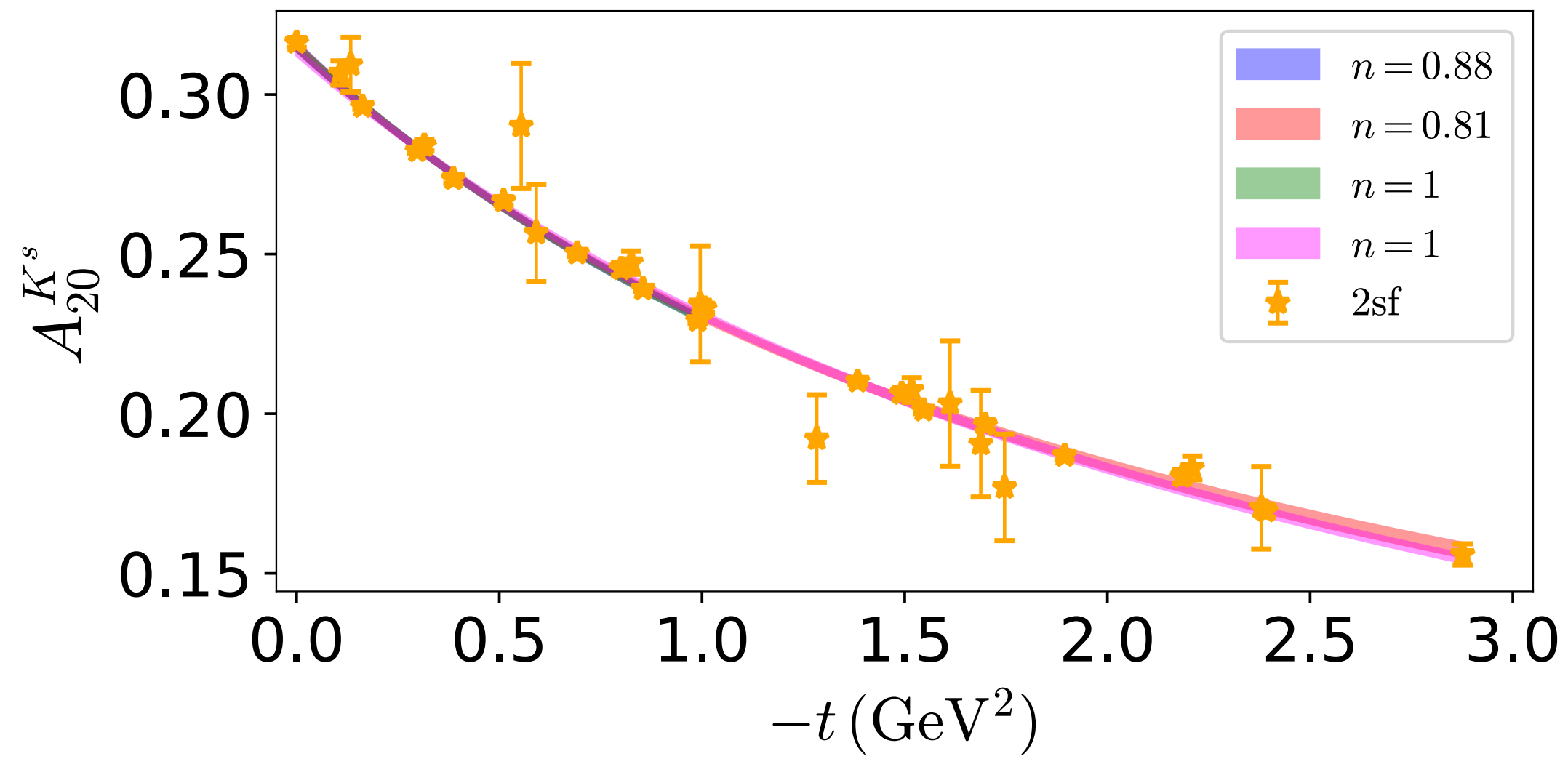
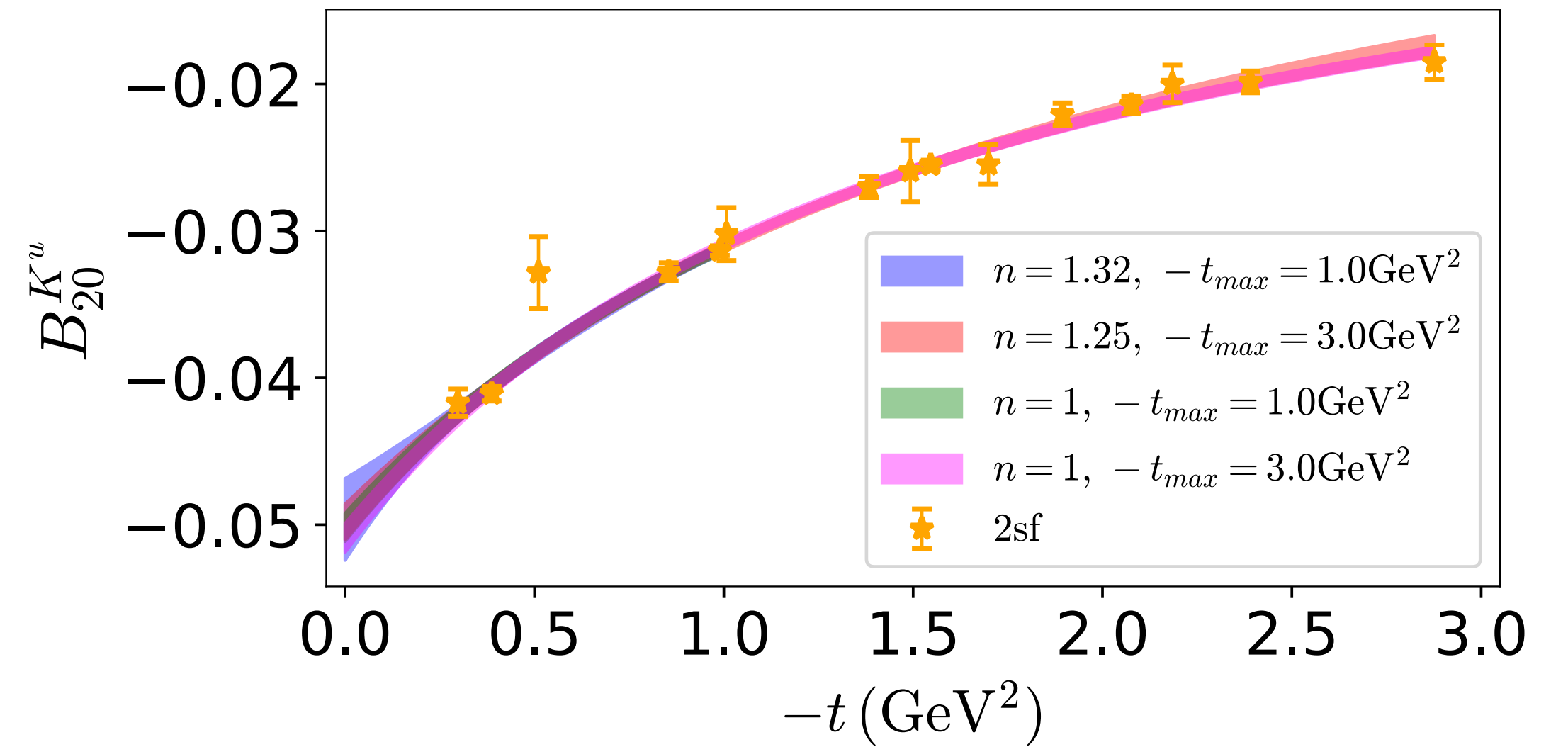
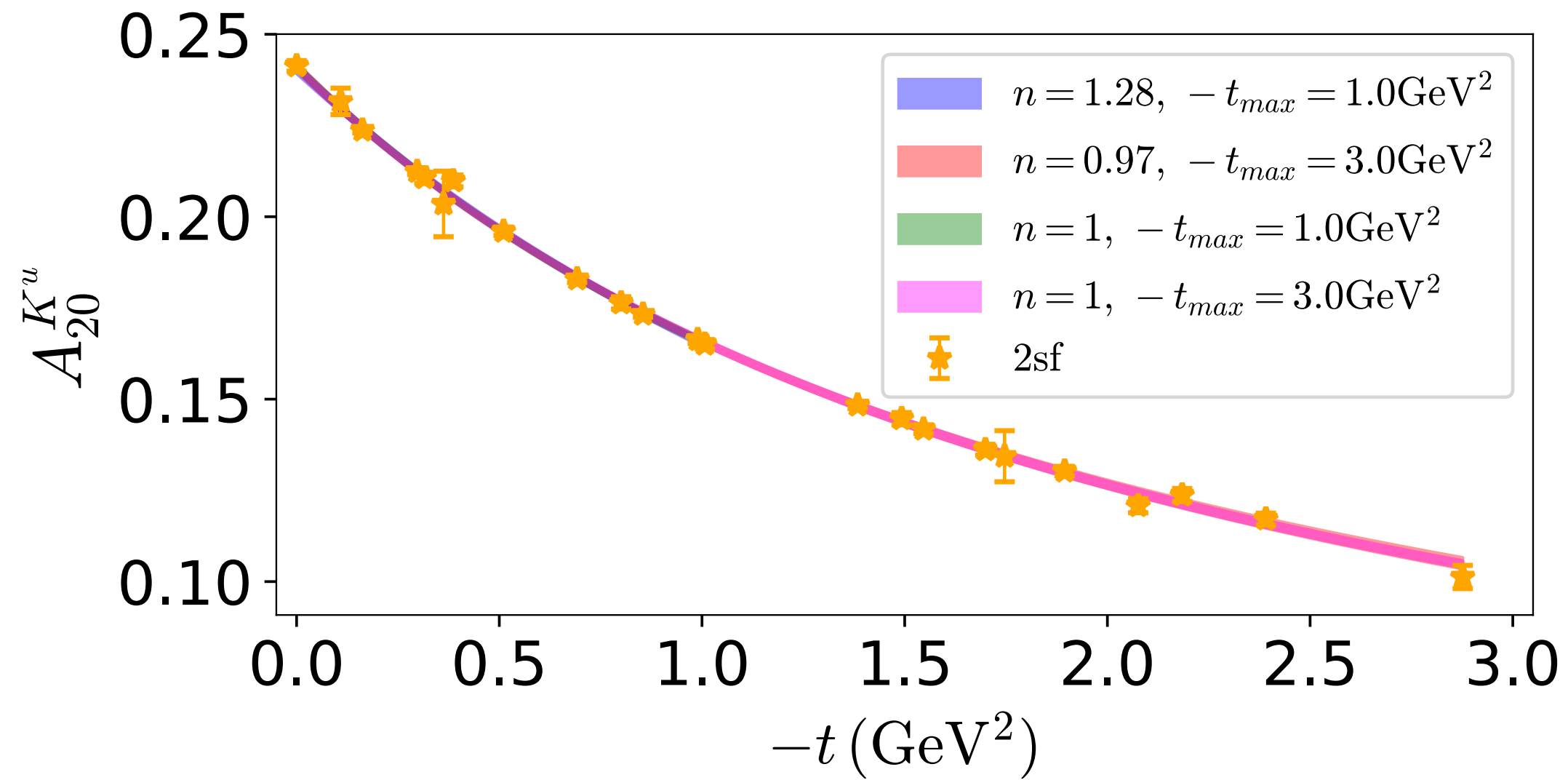


Kaon Parameterization

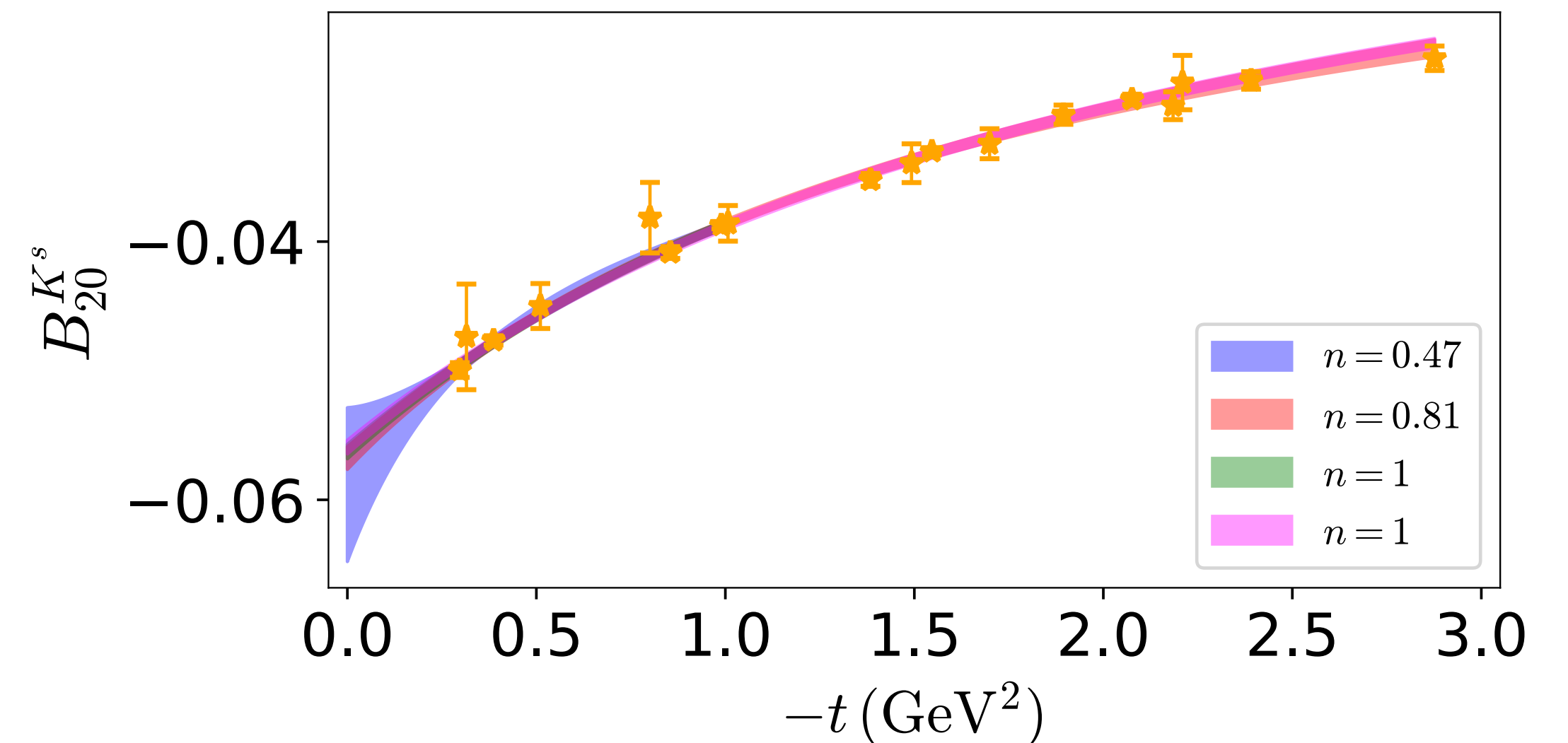
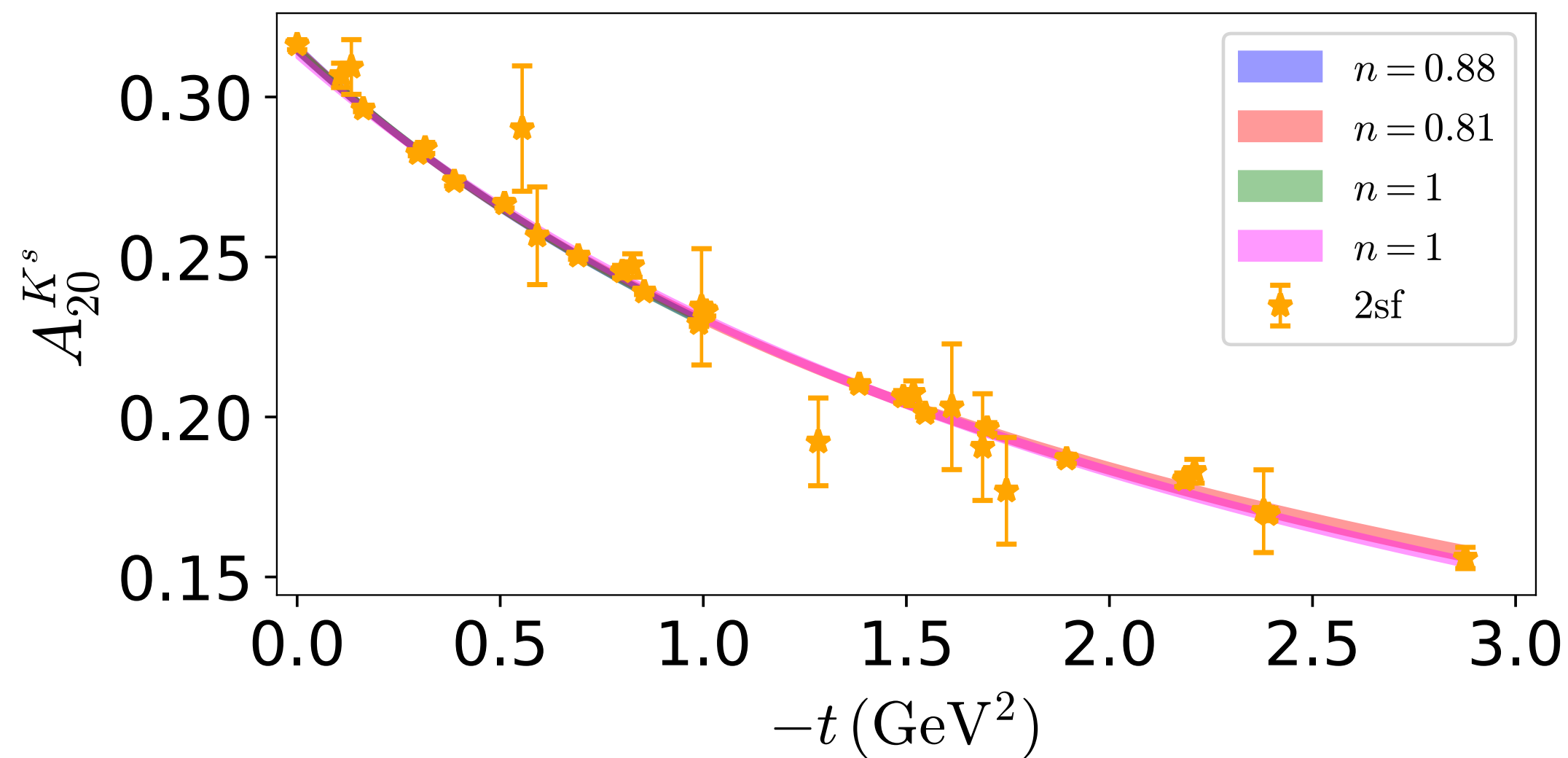
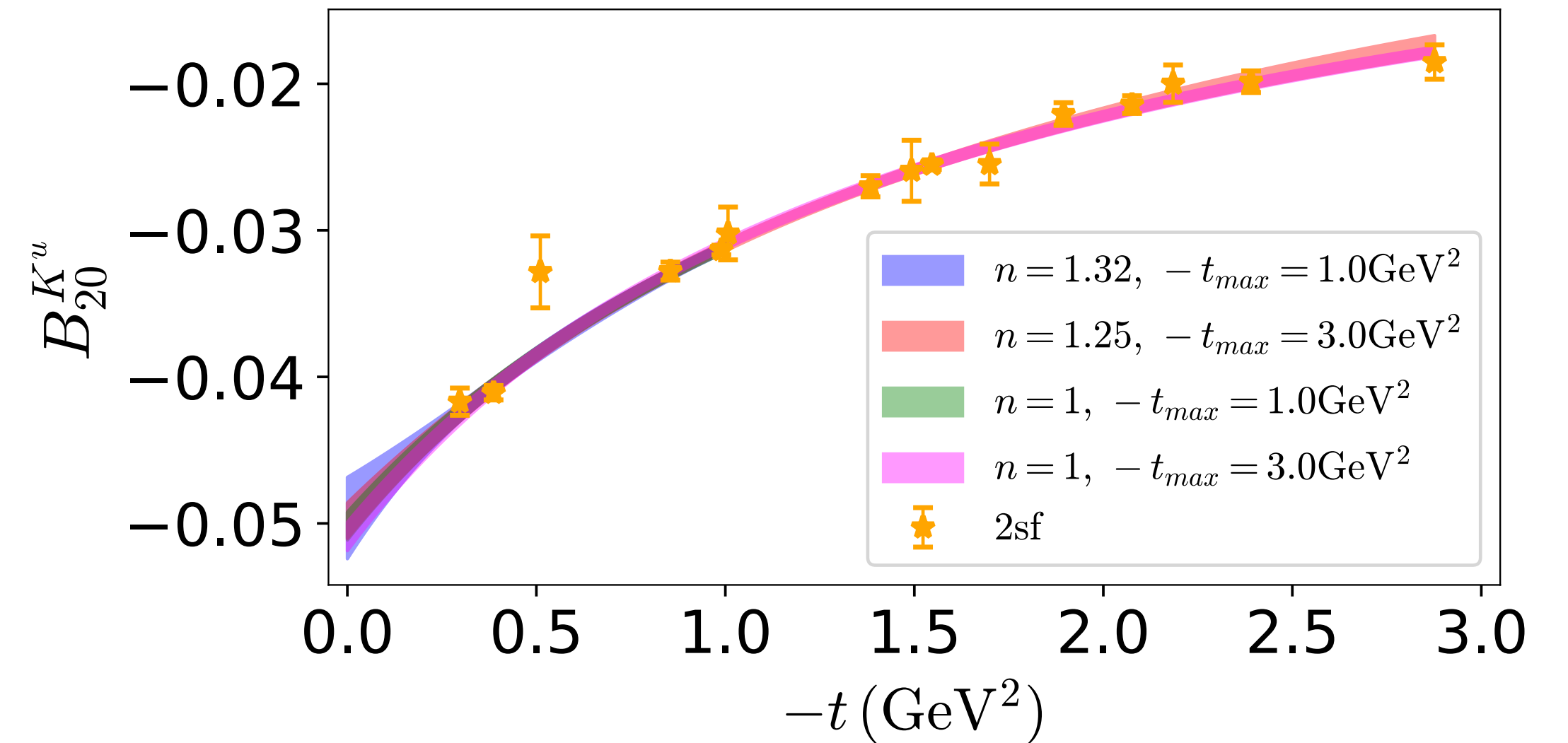
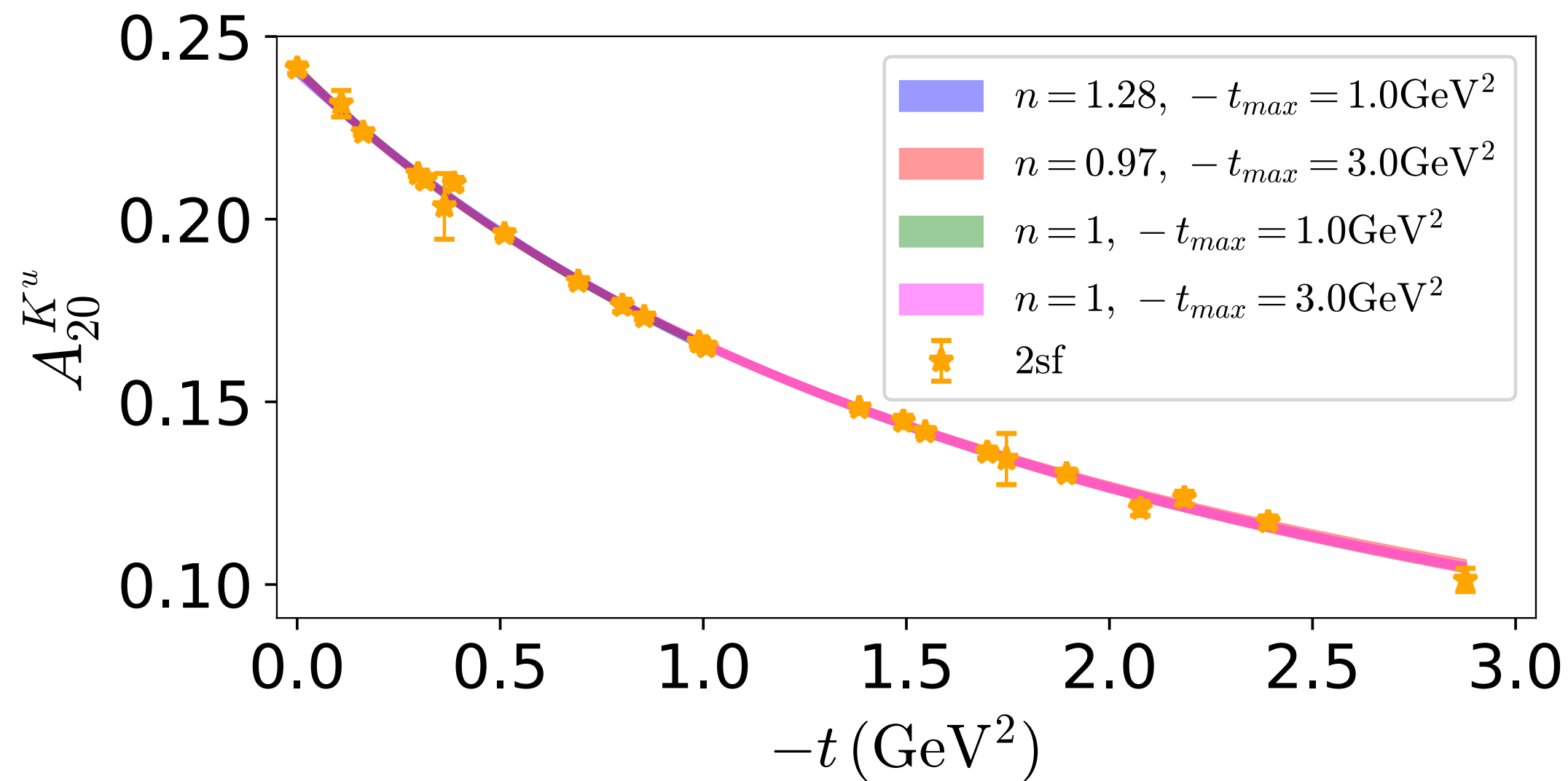
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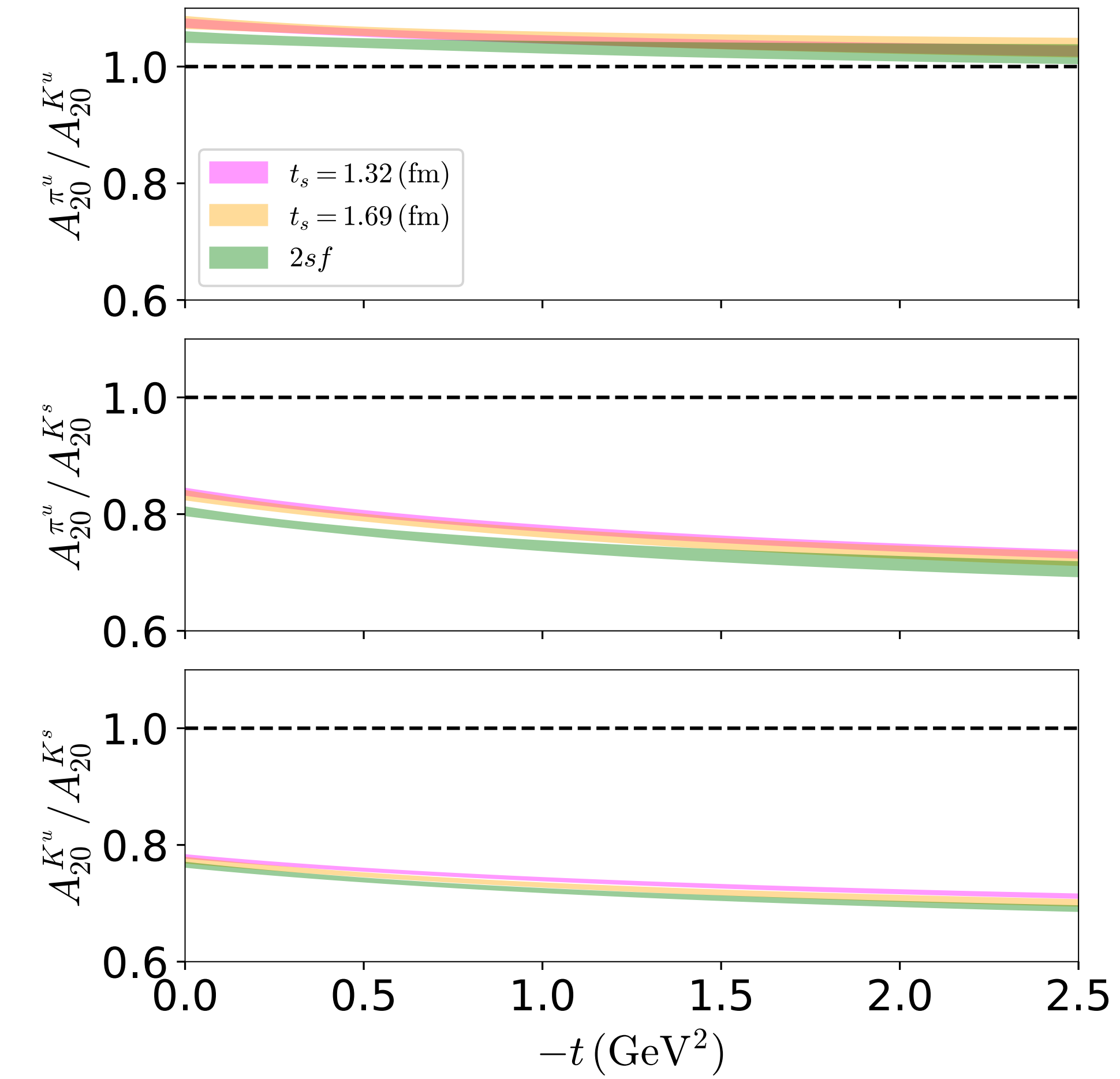
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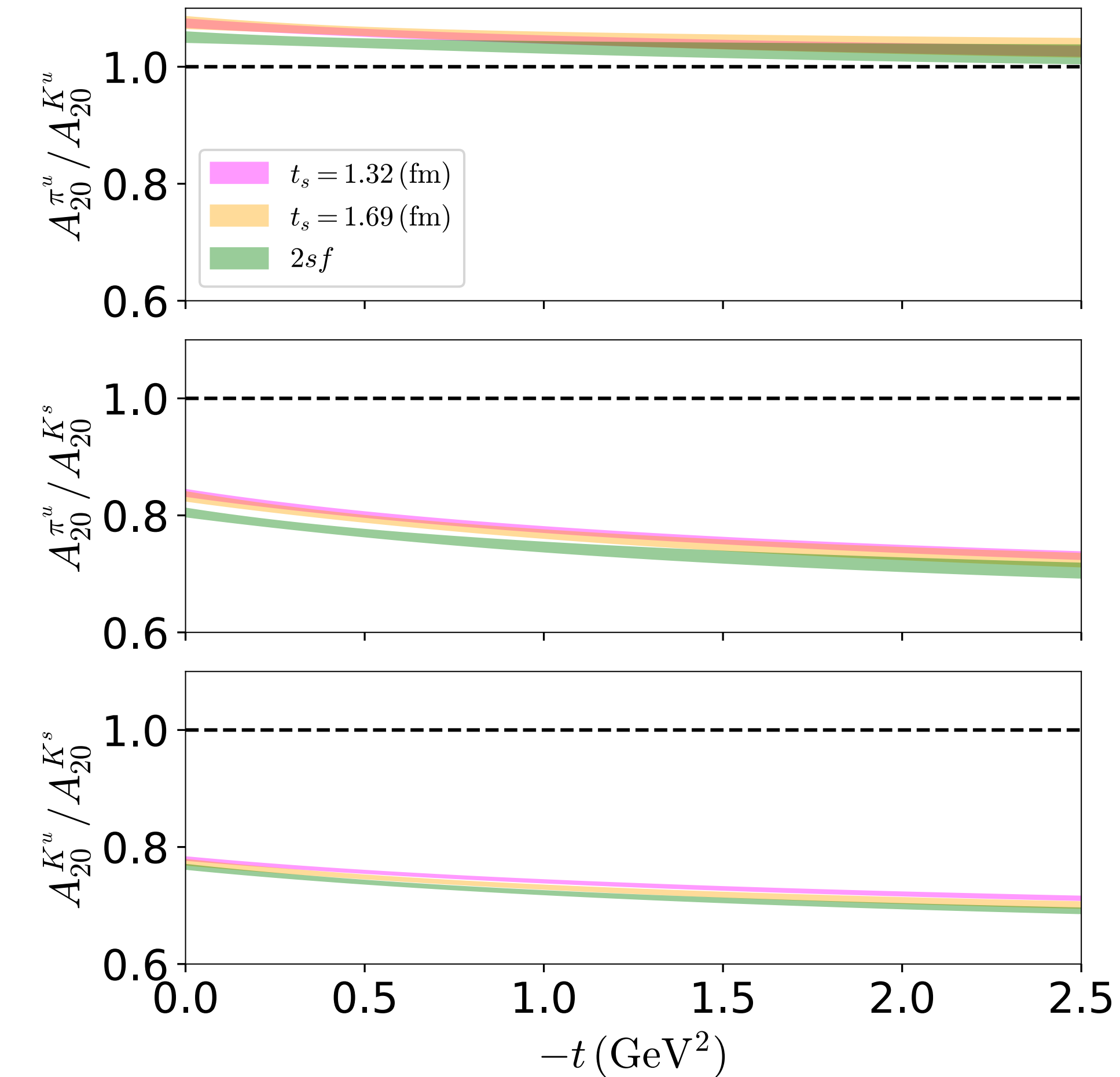
- All parametrizations describe the data very well
- Better stability in the fit for B_{20} in both flavors for the kaon compared to the pion

SU(3) Flavor Symmetry Breaking

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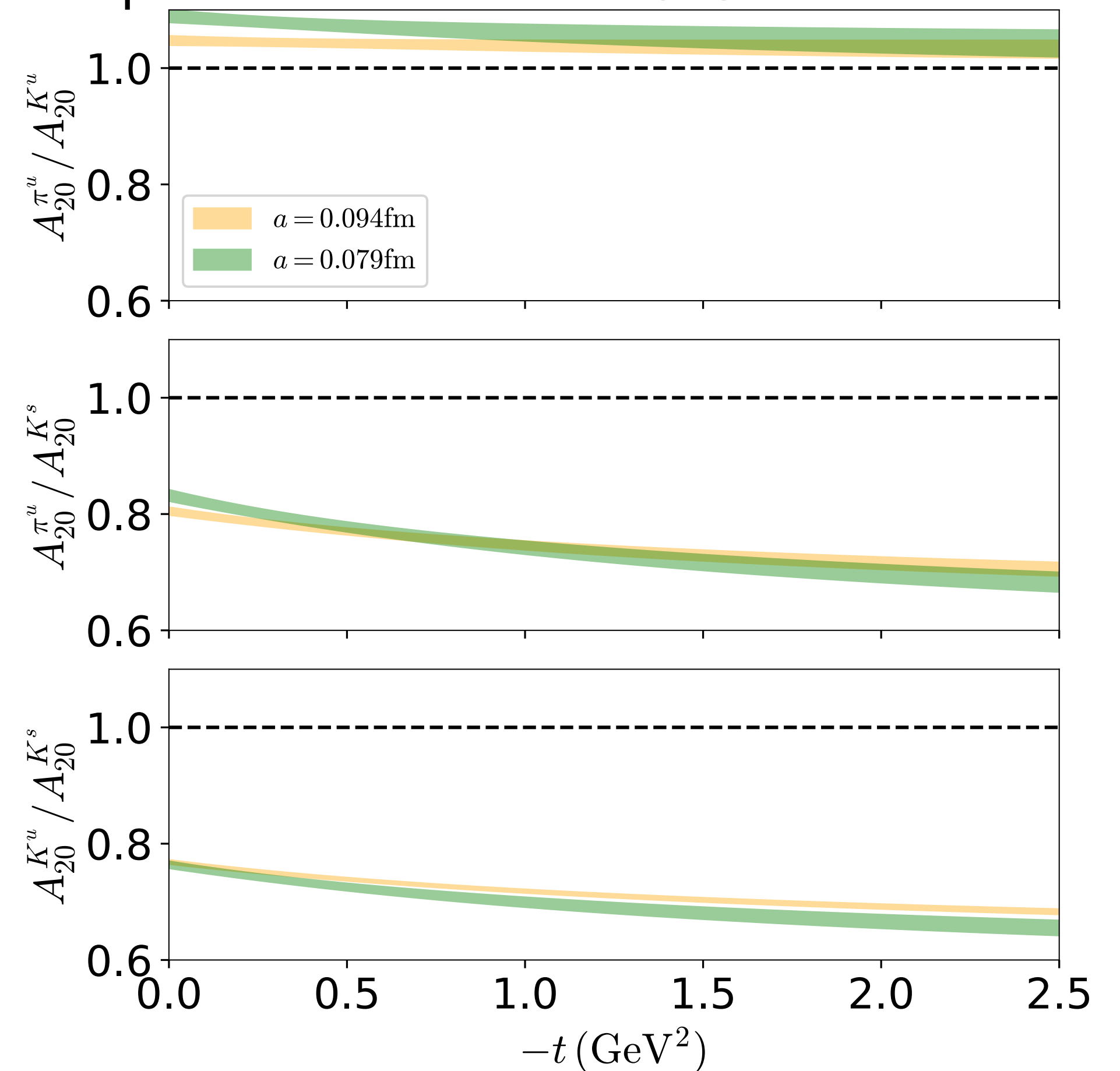
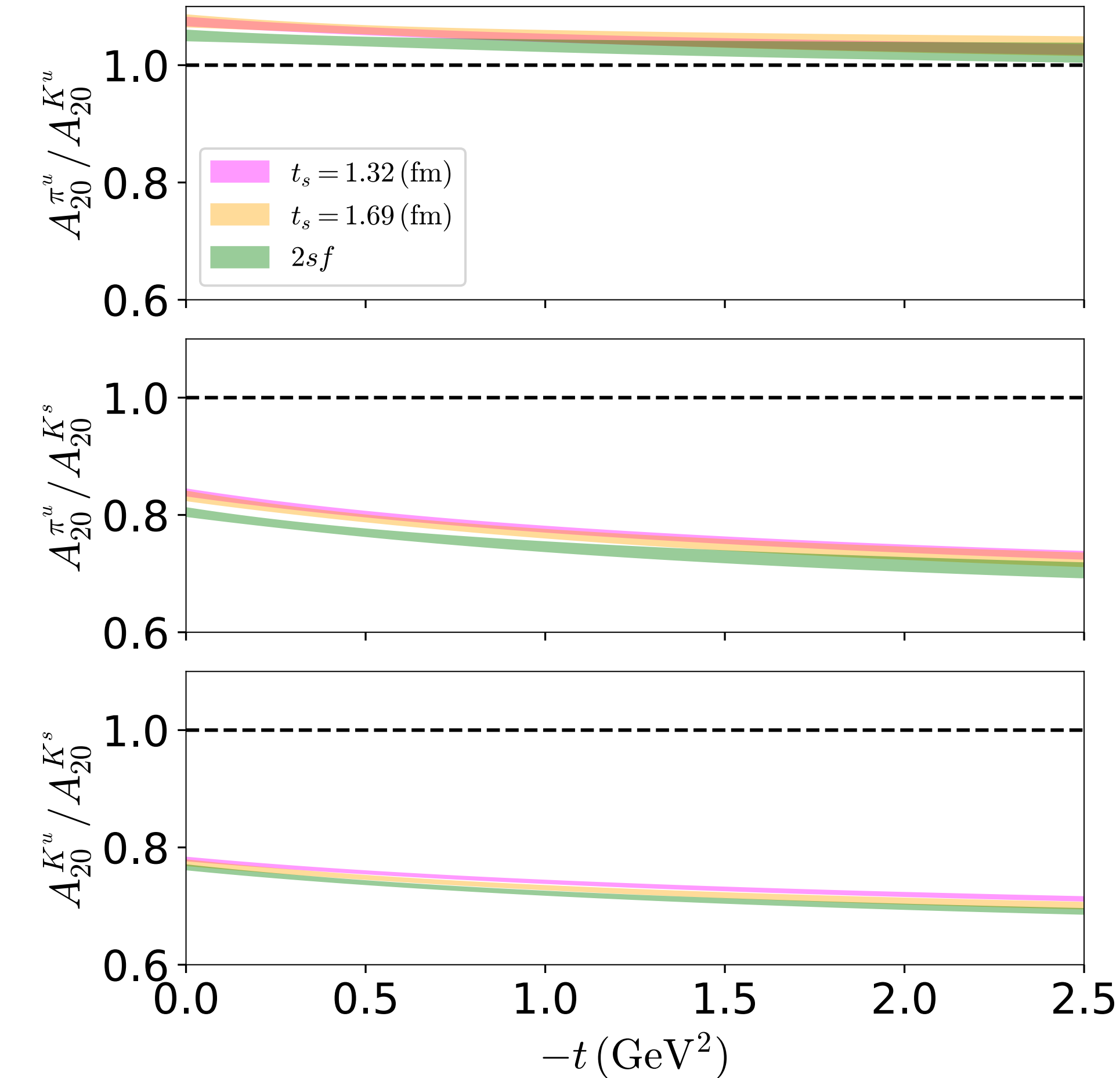
SU(3) Flavor Symmetry Breaking



- Plots show ratios of parametrized data (monopole fit)
- Up quark equivalent contribution in both particles (about 5% - 10% higher in pion)
- SU(3) breaking effects between up and strange quarks at 20% for forward limit and up to 30% for $-t = 2.5 \text{ GeV}^2$

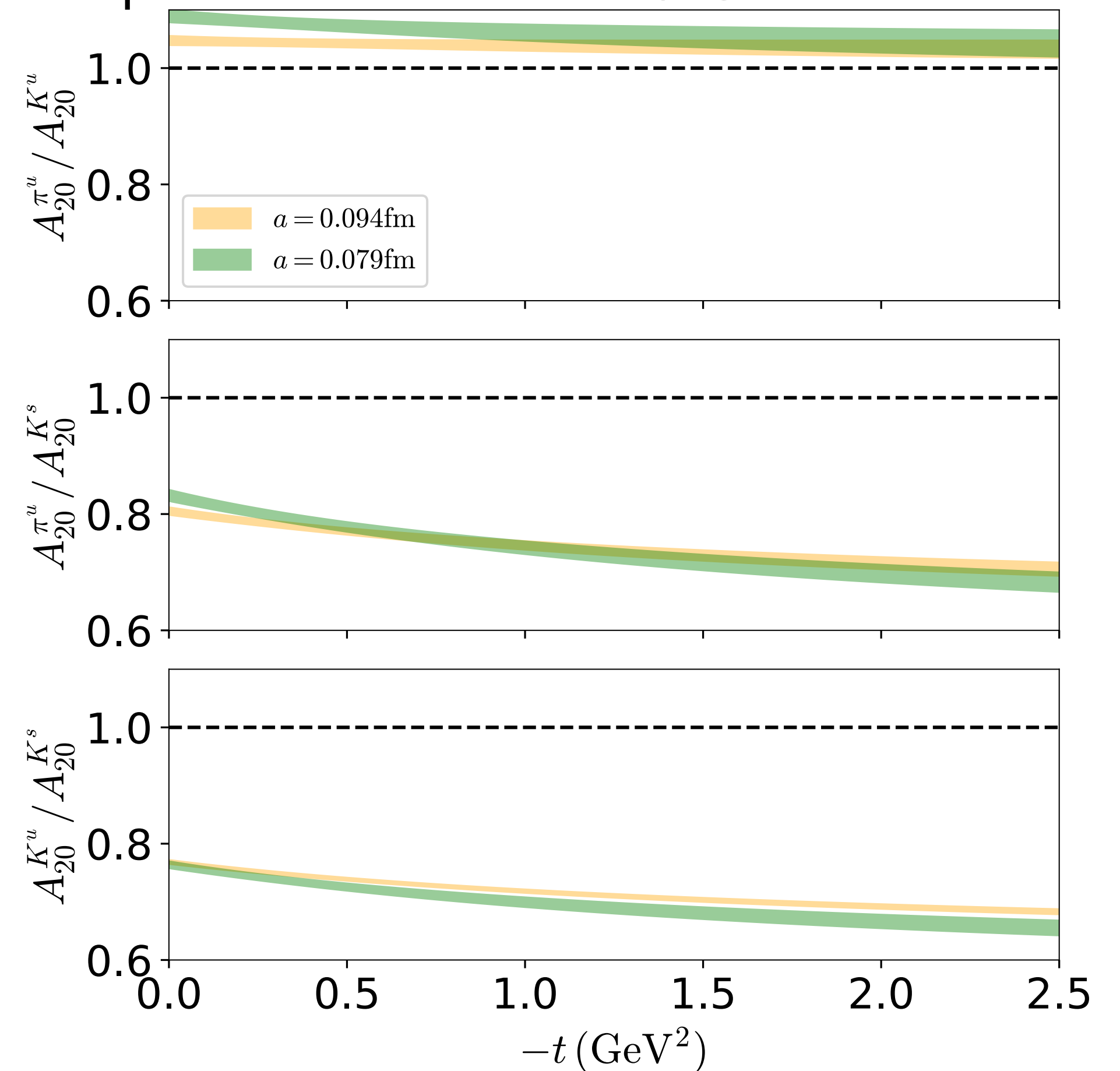
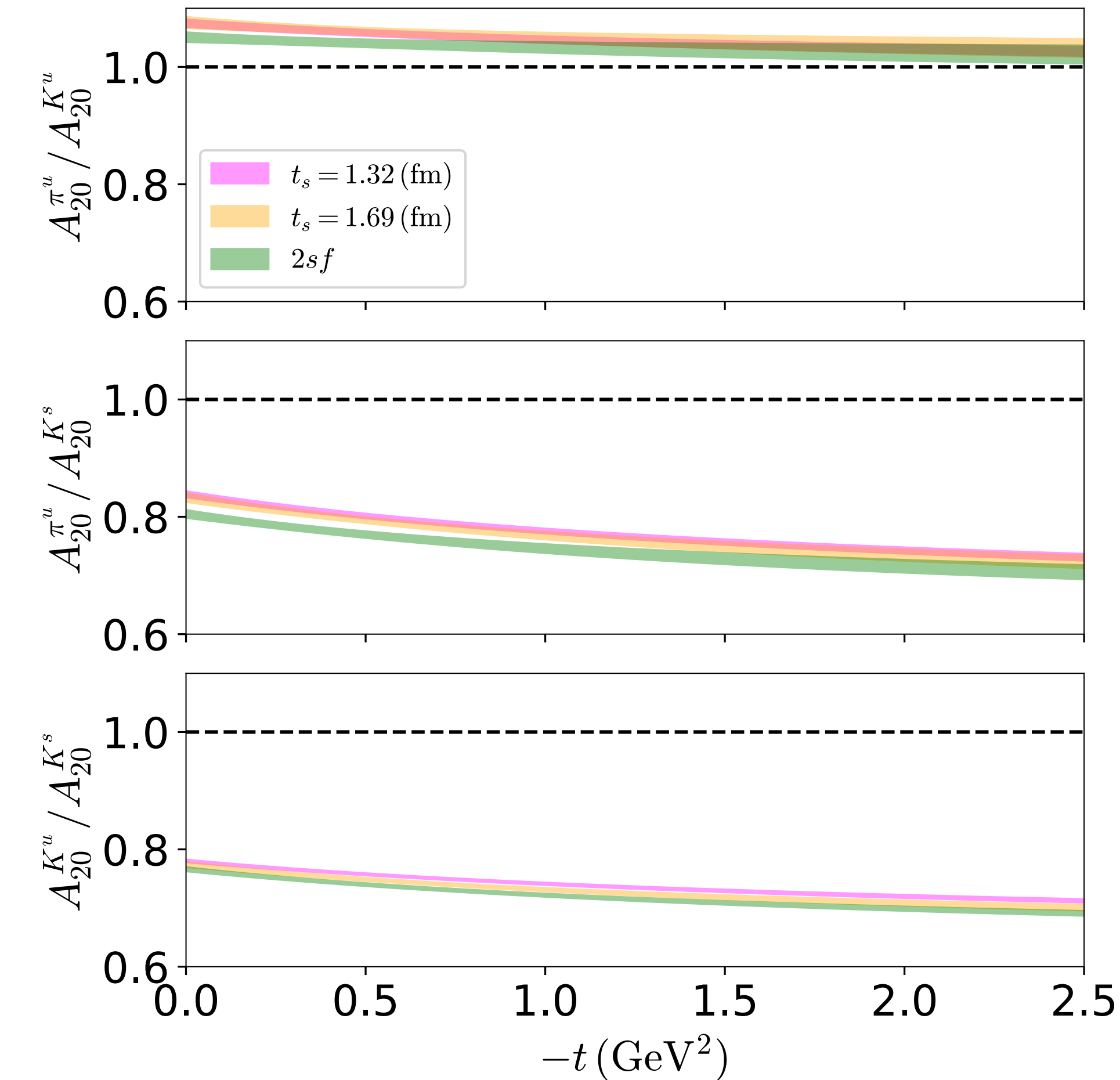
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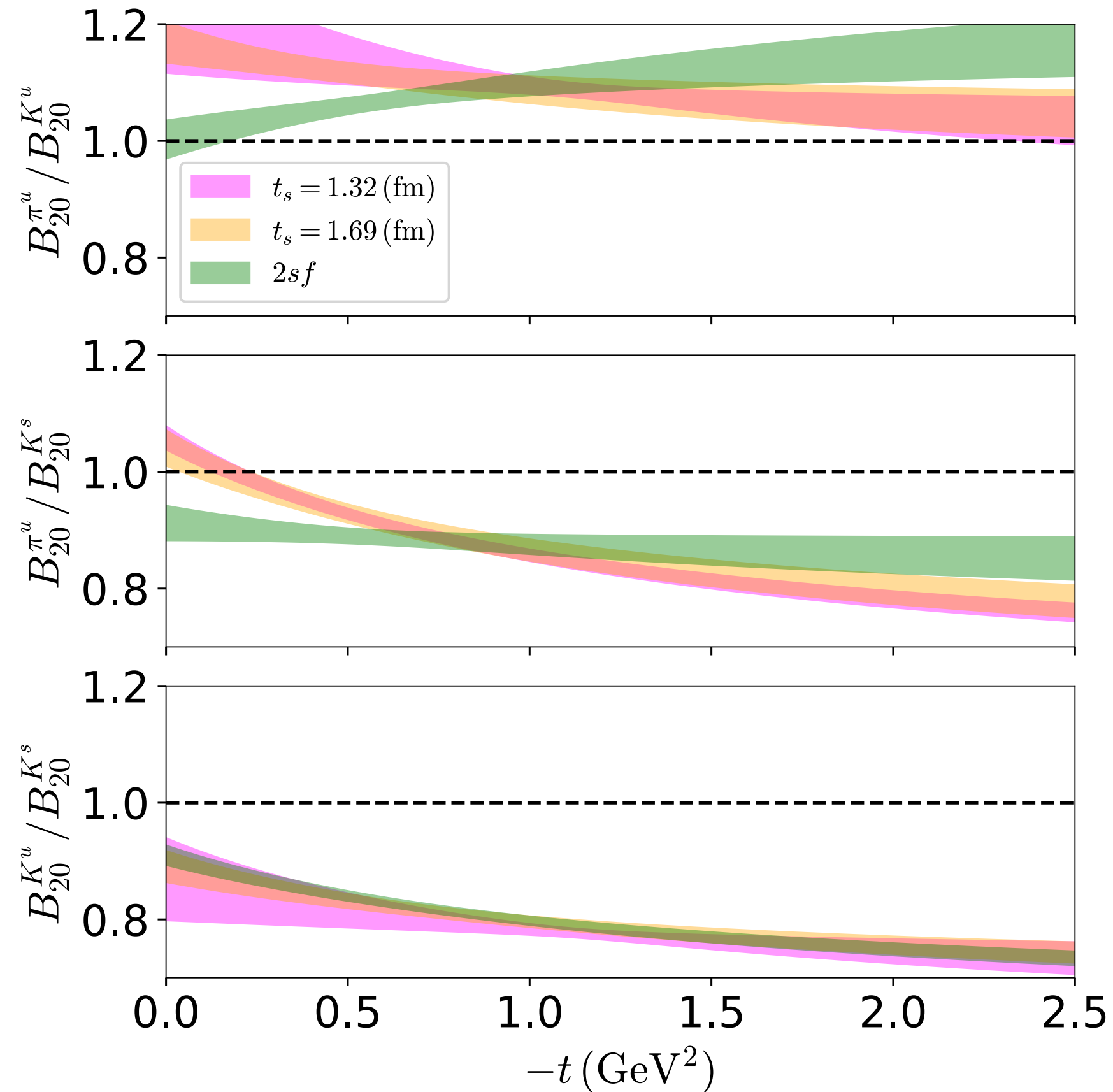
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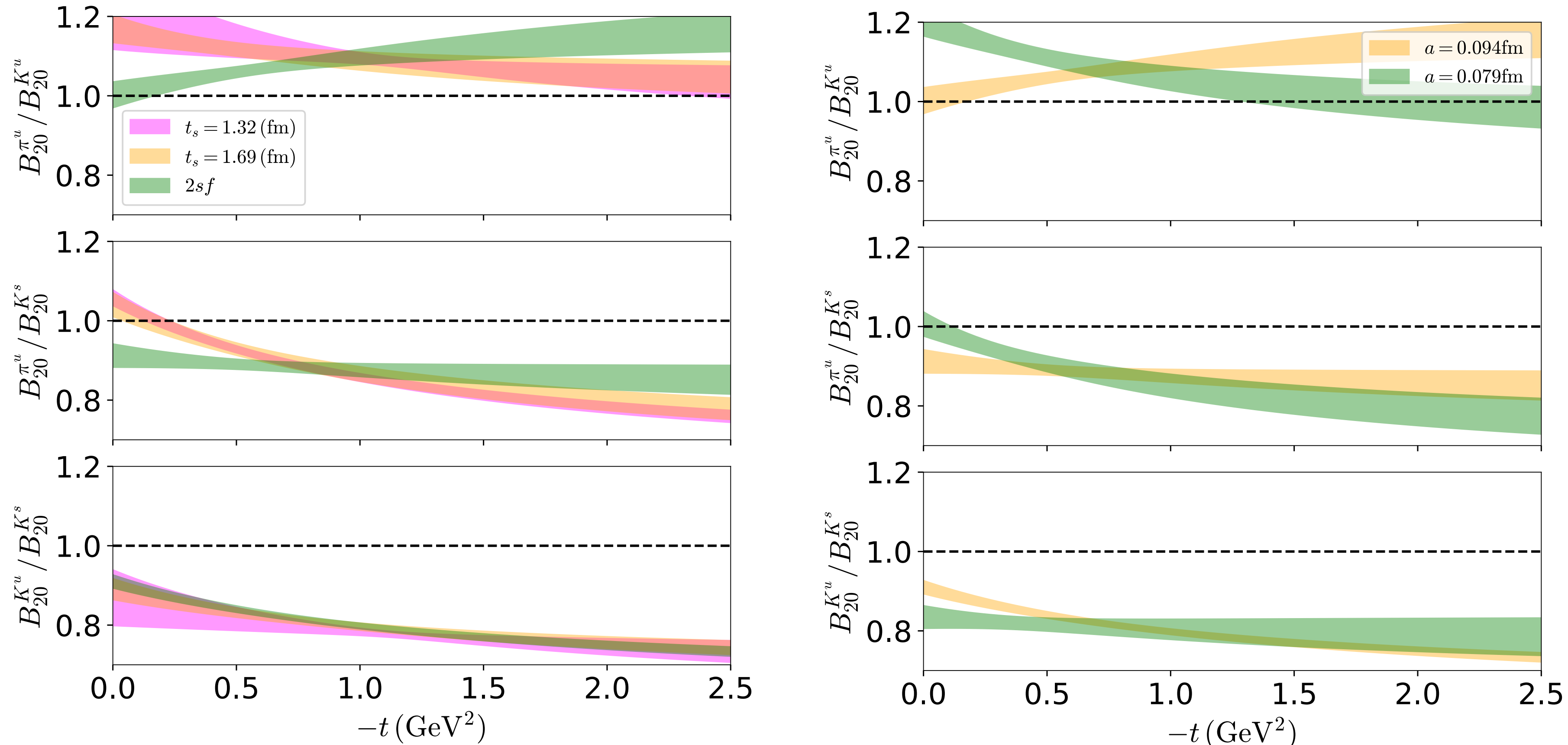
- Similar behavior between ensembles and small differences in certain regions of $-t$
- $a = 0.079$ fm has larger errors (lower statistics) but still under control

SU(3) Flavor Symmetry Breaking



- Ratio affected by small values of B_{20} and large uncertainties (particularly for the comparison of the role of the up-quark in the two particles)
- Excited-states effects more prominent than A_{20}
- 2-state fit: - small SU(3) flavor symmetry breaking between up and strange (~10%)
- Up quark contribution increases in pion as $-t$ increases (up to ~ 20% effect)

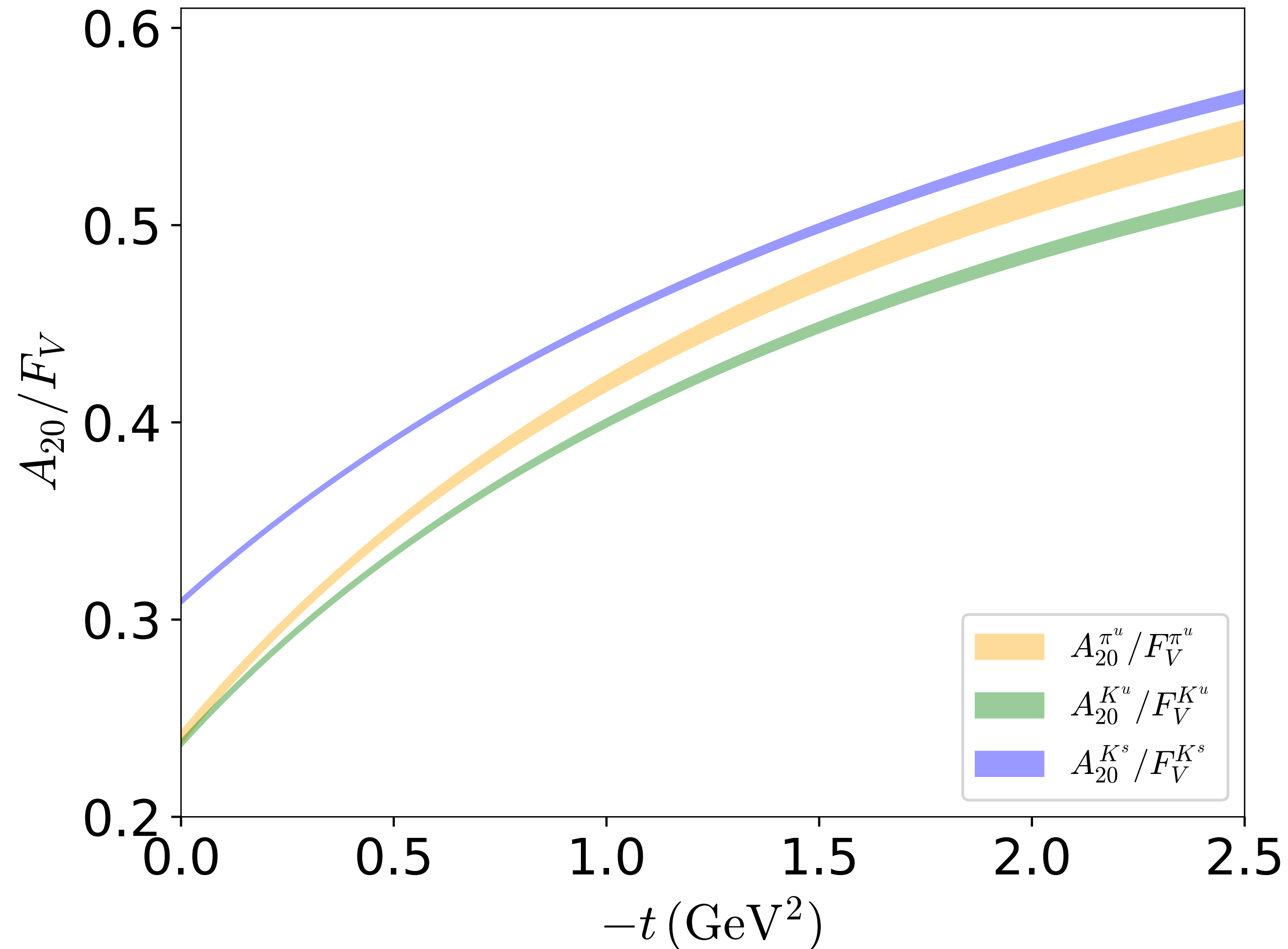
SU(3) Flavor Symmetry Breaking



- Possible discretization and/or volume effects observed

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Comparison of different order of Mellin moments



- Higher Mellin moments receive support at higher values of x and are expected to suppress in value compared to lower Mellin moments.
- Ratio at $-t = 0$ equivalent to $A_{20}(0)$
- Slope of A_{20} different than of F_V (similar behavior for both particles/flavors)
- Values of A_{20} becomes about 50% of F_V at $-t = 2.5 \text{ GeV}^2$

Mellin Moments of PDFs (Preliminary B-ensemble)

- We calculate the Mellin moments of PDFs up to $\langle x^3 \rangle$ with non-perturbative renormalization using operators that avoid mixing
- Pion Mellin moments $\langle x^n \rangle$:

$$\langle x \rangle_A^\pi = 0.273(07) \quad \langle x^2 \rangle_A^\pi = 0.112(06) \quad \langle x^3 \rangle_A^\pi = 0.036(13)$$

$$\langle x \rangle_B^\pi = 0.227(14) \quad \langle x^2 \rangle_B^\pi = 0.090(12) \quad \langle x^3 \rangle_B^\pi : \text{too noisy}$$

- Kaon Mellin moments $\langle x^n \rangle$:

$$\langle x \rangle_A^{K^u} = 0.243(5) \quad \langle x^2 \rangle_A^{K^u} = 0.096(2) \quad \langle x^3 \rangle_A^{K^u} = 0.042(6)$$

$$\langle x \rangle_B^{K^u} = 0.234(6) \quad \langle x^2 \rangle_B^{K^u} = 0.089(4) \quad \langle x^3 \rangle_B^{K^u} = 0.060(11)$$

$$\langle x \rangle_A^{K^s} = 0.307(7) \quad \langle x^2 \rangle_A^{K^s} = 0.129(4) \quad \langle x^3 \rangle_A^{K^s} = 0.060(12)$$

$$\langle x \rangle_B^{K^s} = 0.320(4) \quad \langle x^2 \rangle_B^{K^s} = 0.138(2) \quad \langle x^3 \rangle_B^{K^s} = 0.082(4)$$

Some differences observed outside statistical errors that lead to indication for systematic uncertainties

Summary

- Calculation of pion and kaon generalized form factors fully accessible from a boosted kinematic frame up to 3 GeV^2
- Lattice data reveal systematic effects attributed to a combination of finite- a and finite volume
- Parametrization of A_{20} and B_{20} favors a monopole fit
- SU(3) flavor symmetry breaking observed in u-quark (both particles) vs s-quark (up to 20-30%)
- SU(3) flavor symmetry breaking effects different between A_{20} and B_{20}
- Ratio of A_{20} and the vector FF shows up to 50% suppression in value

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Future Work

- Kinematic setup designed to access generalized form factors up to 3-derivative operator
- Higher statistics required to go beyond A_{20} and B_{20}
- Extraction of second Mellin moment of the quark probability density in impact parameter space, ρ , with addition of tensor GFFs
- Addition of two more ensembles with finer lattice spacing for continuum limit

Acknowledgements

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- Work supported by U.S. Department of Energy Early Career Award (Grant No. DE-SC0020405)
- This work used resources provided by ACCESS (Advanced Cyberinfrastructure Coordination Ecosystem: Services & Support, formerly XSEDE), funded by the National Science Foundation (NSF) through award ACI-1540931