

Trace anomaly form factors of the pion and the nucleon from lattice QCD

Bigeng Wang

Collaborators: Fangcheng He, Gen Wang, Yi-Bo Yang, Jian Liang, Terrence Draper,
Keh-Fei Liu
(χ QCD collaboration)

Department of Physics and Astronomy
University of Kentucky

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- Energy momentum tensor (EMT)


$$T_{\mu\nu} = \frac{1}{4}\bar{\psi}\gamma_{(\mu}\overleftrightarrow{D}_{\nu)}\psi + G_{\mu\alpha}G_{\nu\alpha} - \frac{1}{4}\delta_{\mu\nu}G^2 \quad (1)$$

- Pion mass can be obtained from the trace of the EMT:

$$m_\pi = \underbrace{\frac{\langle\pi|\int d^3\vec{x}[\frac{\beta(g)}{2g}G^2 + \sum_f\gamma_m(g)m_f\bar{\psi}_f\psi_f]|\pi\rangle}{\langle\pi|\pi\rangle}}_{\text{conformal symmetry breaking} \leftrightarrow \text{trace anomaly ME}} + \underbrace{\frac{\langle\pi|\int d^3\vec{x}\sum_f m_f\bar{\psi}_f\psi_f|\pi\rangle}{\langle\pi|\pi\rangle}}_{\sigma \text{ term, } \frac{1}{2}m_\pi \propto \sqrt{m_q}^{-1}} \quad (2)$$

- 1st order in the **chiral symmetry breaking**:

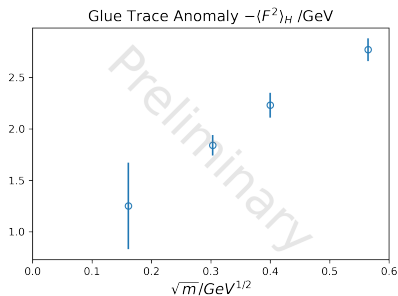
$$m_\pi \propto \sqrt{m_q}, \quad \text{for } m_q = m_u = m_d \quad (3)$$

¹Based on the Gellmann-Oakes-Renner relation and the Feynman-Hellman theorem. 

Physics Motivation: Pion Mass Puzzle

- **Does** the trace anomaly matrix element keep itself proportional to $\sqrt{m_q}$ as $m_q \rightarrow 0$?
F. He, P. Sun and Y.B. Yang (χ QCD) (PRD 2021, 2101.04942)

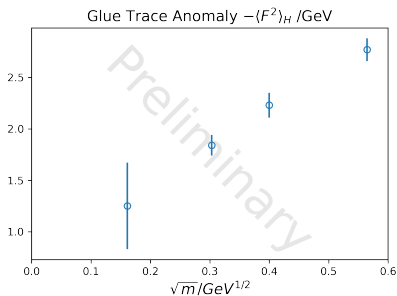
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This is a suggestion that the **conformal** (scale) symmetry breaking in the **pion** is linked to the **chiral** symmetry breaking. K.F. Liu arXiv:2302.11600

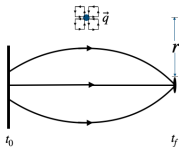
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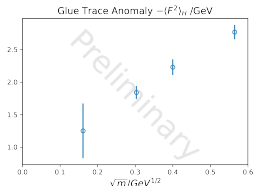
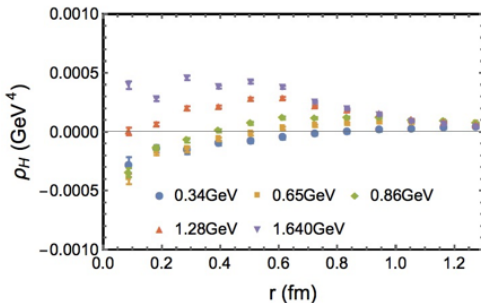
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Calculate a density function

$\rho_H(r)$:



$$\int \rho_H(r) r^2 dr \sim \frac{\beta(g)}{2g} G^2 \quad (5)$$



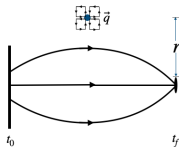
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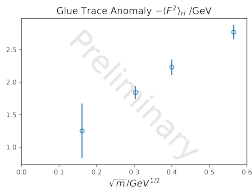
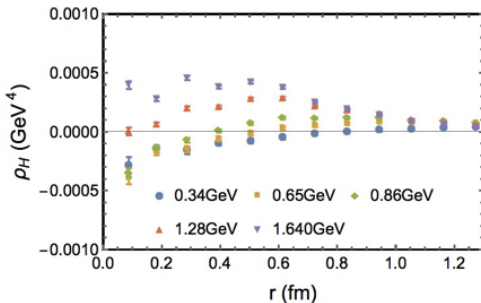
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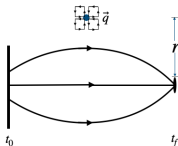
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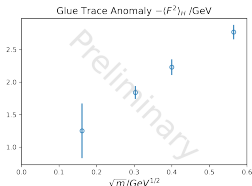
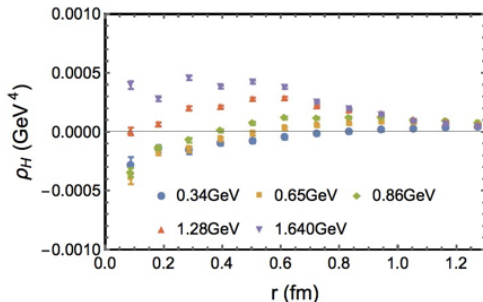
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- Will the **form factors** changes sign as well? (if they are connected by some sort of Fourier transform)

- 1 Normalization convention:

$$1 = \int \frac{d^3p}{(2\pi)^3} |p\rangle \frac{m}{E_p} \langle p|, \quad |p\rangle = \sqrt{\frac{E_p}{m}} a_p^+ |\Omega\rangle \quad (6)$$

- 2 Define a **dimensionless** trace anomaly form factor $G_H(Q^2)$, where $Q^2 = -(P' - P)^2$:
- for spin- $\frac{1}{2}$ particle like proton:

$$\langle P' | T_\mu^\mu | P \rangle = m_N G_N(Q^2) \bar{u}(P') u(P), \quad (7)$$

- for spin-0 particle like pion:

$$\langle P' | T_\mu^\mu | P \rangle = m_\pi G_\pi(Q^2). \quad (8)$$

- 3 $G_H(Q^2 = 0)$ is the contribution to the total mass of hadron H .

Renormalization of the trace anomaly form factors

The trace anomaly terms need renormalization:

$$\langle \pi | \int d^3 \vec{x} \gamma \left[\frac{\beta(g)}{2g} G^2 + \sum_f \gamma_m(g) m_f \bar{\psi}_f \psi_f \right] | \pi \rangle \quad (9)$$

Renormalization method F. He, P. Sun and Y.B. Yang (χ QCD) (PRD 2021, 2101.04942) :

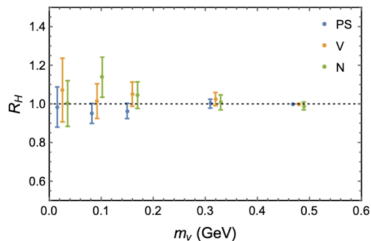
- $\frac{\beta(g)}{2g}$ and $\gamma_m(g)$ are independent of the hadron state
- Verify the assumption: sum rule satisfied for other masses

- Solve the mass sum-rule equations for pseudo-scalar(π) and vector meson(ρ) at one mass $m_v a = 0.3$:

$$M_{\text{PS}} - (1 + \gamma_m) \langle H_m \rangle_{\text{PS}} - \frac{\beta(g)}{2g} \langle F^2 \rangle_{\text{PS}} = 0, \quad (10)$$

$$M_{\text{V}} - (1 + \gamma_m) \langle H_m \rangle_{\text{V}} - \frac{\beta(g)}{2g} \langle F^2 \rangle_{\text{V}} = 0, \quad (11)$$

and obtain the bare γ_m and $\frac{\beta(g)}{2g}$.



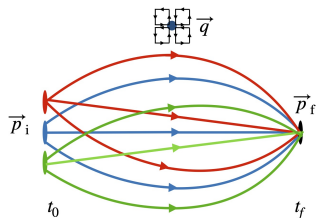
$$R_H = \left[(1 + \gamma_m) \langle H_m \rangle_H + \frac{\beta(g)}{2g} \langle F^2 \rangle_H \right] / m_H \quad (12)$$

Calculation of **glue** trace anomaly FF using two- and three-point correlators, with **grid source** and **low-mode substitution (LMS)**

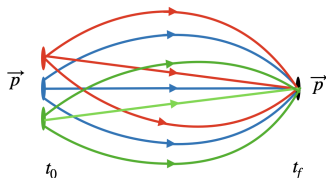
- For proton, the σ term is small, and glue trace anomaly dominates
- For pion, trace anomaly $\sim \frac{1}{2}m_\pi$ and γ_m is not very large.

→ calculate **glue** trace anomaly form factors

Three-point correlators:



Two-point correlators:



Need large statistics for various momentum transfer values:

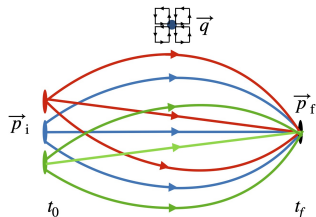
- Add source and sink momenta using phase factors (**no need for extra inversions**)

$$C_N^{\mathcal{G}}(\vec{p}_i, \vec{p}_f, t) = \frac{1}{n} \sum_i^n e^{i\vec{p}_i \cdot \vec{w}_i} C_N^{\mathcal{G}, \vec{w}_i}(\vec{p}_f, t) + C_N^{\mathcal{G}, H}(\vec{p}_i, \vec{p}_f, t) \quad (13)$$

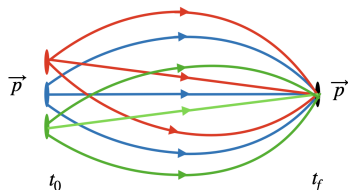
G. Wang, Y.-B. Yang, J. Liang, T. Draper, and K.-F. Liu, Phys. Rev. D **106**, 014512 (2022) [arXiv:2108.08111]

Calculation of **glue** trace anomaly FF using two- and three-point correlators, with **grid source** and **low-mode substitution (LMS)**

Three-point correlators:



Two-point correlators:



Taking ratios of the 3pt and 2pt (using pion as an example):

$$R_{\pi}^{S_i S_f}(t, \tau; \vec{p}_i, \vec{p}_f) = \frac{C_{\pi, 3pt}^{S_i S_f}(t, \tau; \vec{p}_i, \vec{p}_f)}{C_{\pi, 2pt}^{S_f}(t; \vec{p}_f)} \sqrt{\frac{C_{\pi, 2pt}^{S_i}(t - \tau; \vec{p}_i) C_{\pi, 2pt}^{S_f}(t; \vec{p}_f) C_{\pi, 2pt}^{S_f}(\tau; \vec{p}_f)}{C_{\pi, 2pt}^{S_f}(t - \tau; \vec{p}_f) C_{\pi, 2pt}^{S_i}(t; \vec{p}_i) C_{\pi, 2pt}^{S_i}(\tau; \vec{p}_i)}}$$

$$\xrightarrow{t \gg \tau \gg 0} \frac{m_{\pi}^2}{\sqrt{E_{\pi, \vec{p}_i} E_{\pi, \vec{p}_f}}} G_{\pi}(Q^2)$$

$$+ C'_1 e^{-\Delta E_i^1 \tau} + C'_2 e^{-\Delta E_f^1 (t - \tau)} + C'_3 e^{-\Delta E_i^1 \tau - \Delta E_f^1 (t - \tau)}$$

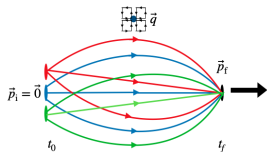
- Overlap fermions on DWF at near-physical pion mass:

Ensemble	$L^3 \times T$	a (fm)	L (fm)	m_π (MeV)	N_{conf}	N_{src}
24l	$24^3 \times 64$	0.1105(3)	2.65	340	783	64×2

Three different momentum transfer scenarios, up to $\sim O(100)$ source-sink momentum combinations (with same Q^2 averaged)

- source at rest:

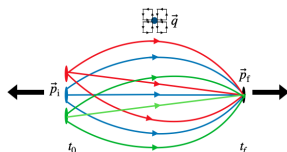
$$|\vec{p}_i| = 0 \text{ with } \vec{q} = \vec{p}_f$$



small Q^2

- back-to-back:

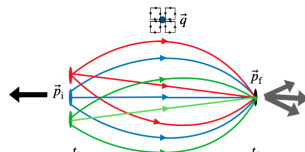
$$\vec{p}_f = -\vec{p}_i \text{ with } \vec{q} = 2\vec{p}_f$$



large Q^2

- near-back-to-back:

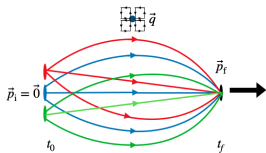
$$\vec{p}_f \neq -\vec{p}_i, \vec{p}_f \text{ \& } -\vec{p}_i \sim \vec{q}/2$$



other nearby Q^2

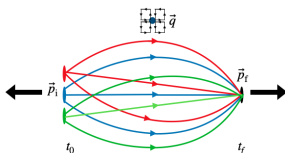
- source at rest:

$$|\vec{p}_i| = 0 \text{ with } \vec{q} = \vec{p}_f$$



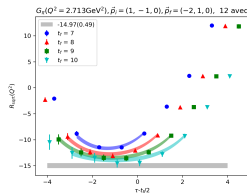
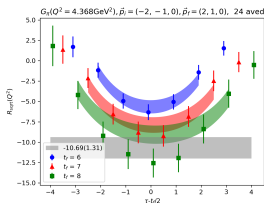
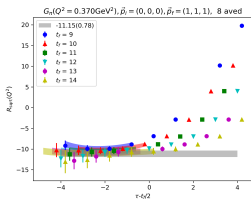
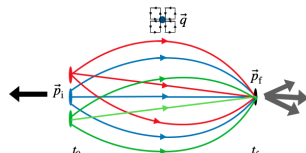
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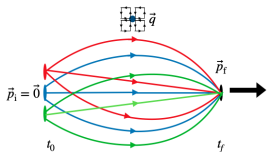
- near-back-to-back:

$$\vec{p}_f \neq -\vec{p}_i, \vec{p}_f \ \& \ -\vec{p}_i \sim \vec{q}/2$$

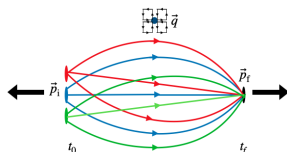


Results for the pion **preliminary**

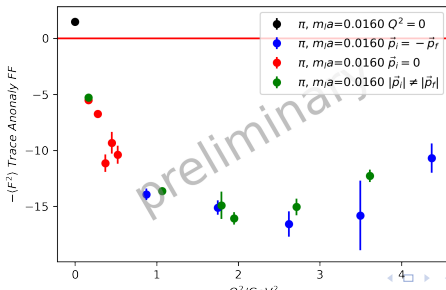
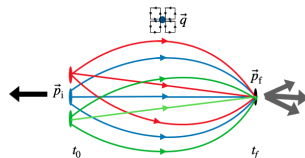
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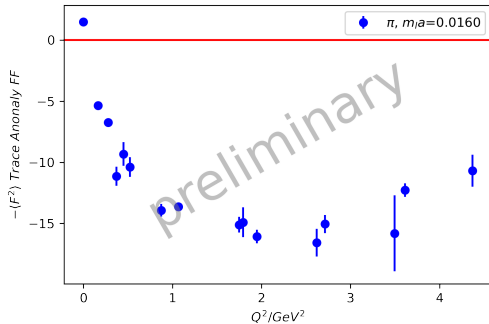
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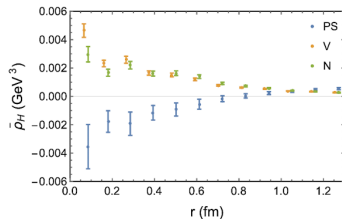
Current work: form factors



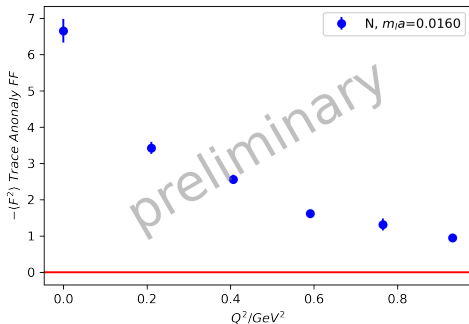
- **positive** at $Q^2 = 0 \text{ GeV}^2$ (contribution to the pion mass from glue)
- **sign change** of glue trace anomaly form factors for **pion**, consistent with the density function.
- form factor calculated up to $Q^2 \sim 4.3 \text{ GeV}^2$

Previous results: density functions

F. He, P. Sun and Y.B. Yang (χ QCD)
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Current work: form factors

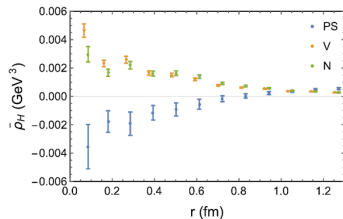


- **positive** at $Q^2 = 0 \text{ GeV}^2$ (contribution to the proton mass from glue)
- **NO sign change**, monotonically decreasing, consistent with the density function.

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1 pion mass puzzle (motivation):

- trace anomaly matrix element is proportional to $\sqrt{m_q}$ as $m_q \rightarrow 0$.
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Outlook

- Extract radius of trace anomaly form factors for π , ρ , and N .
- We expect the calculation on the 48l ensemble will give a prediction of the trace anomaly form factors **at physical pion mass**.

Thanks for your attention!