



2023 LATTICE

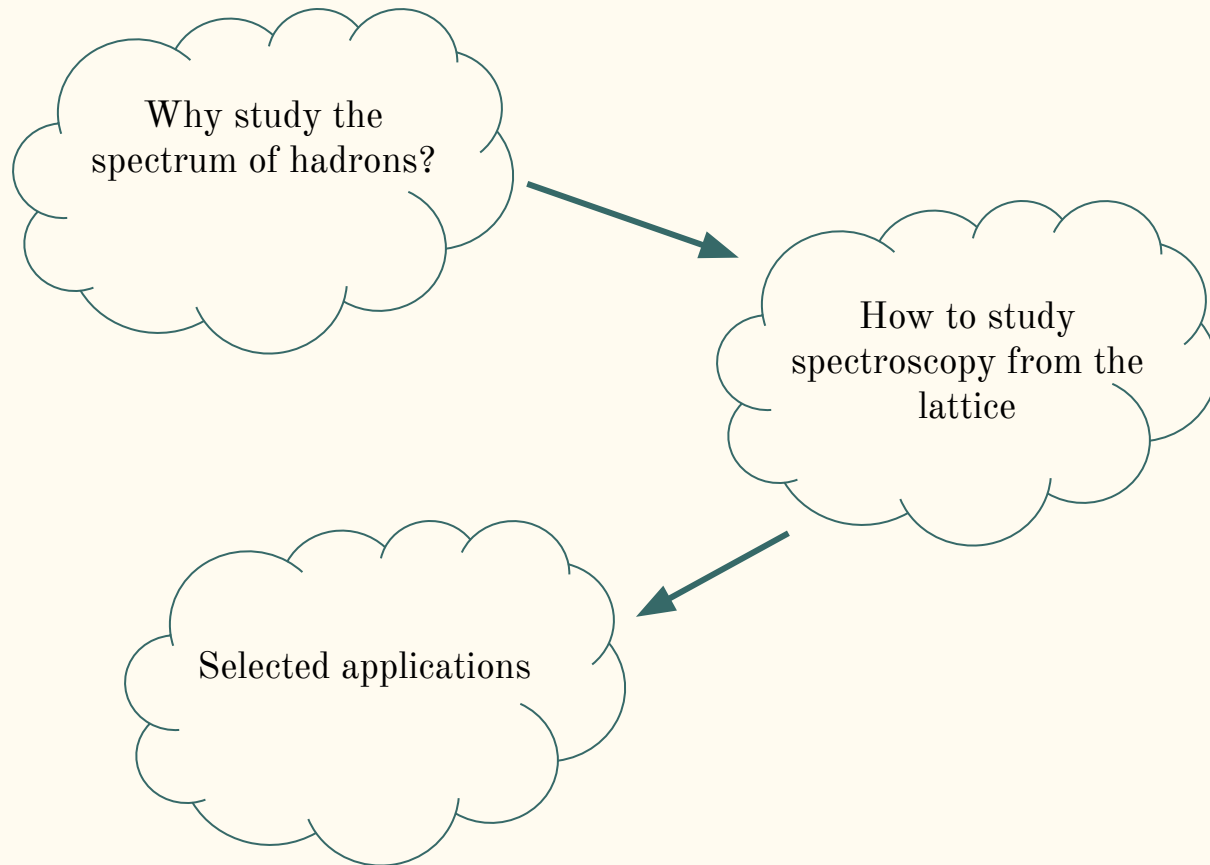
Hadron spectroscopy and few-body dynamics

Andrew Hanlon
Brookhaven National Laboratory

The 40th Lattice Conference
Fermilab
Aug. 3, 2023

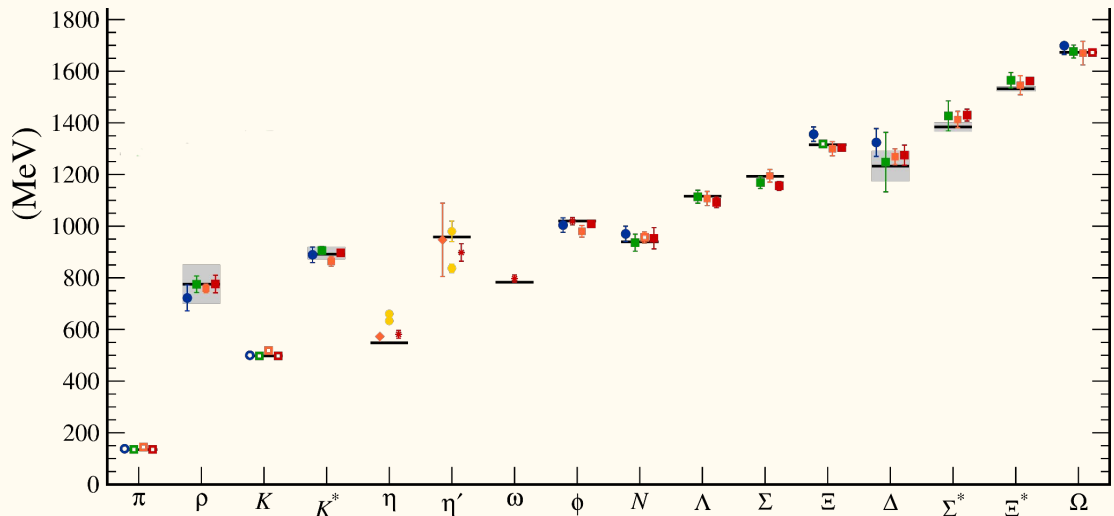


Outline



Spectroscopy from the lattice

Spectrum of QCD-‘stable’ states agrees well with experiment



[A. Kronfeld 1203.1204]

However:

- Bound states near threshold can have significant finite-volume effects
- Further, most hadrons are *resonances*, i.e. unstable states

How to proceed?:

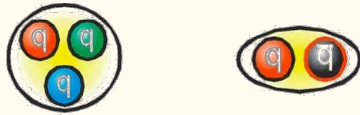
- Robust extraction of resonance information from multi-hadron states on the lattice can be (and has been) achieved!

[Lüscher '86, '91; generalizations]

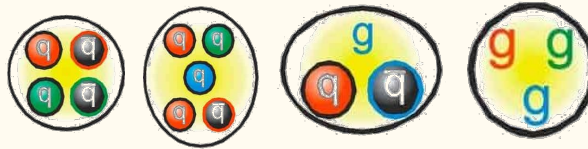
The limits of the quark model

The quark model was mostly successful at describing a wide range of the observed resonances, but is clearly incomplete

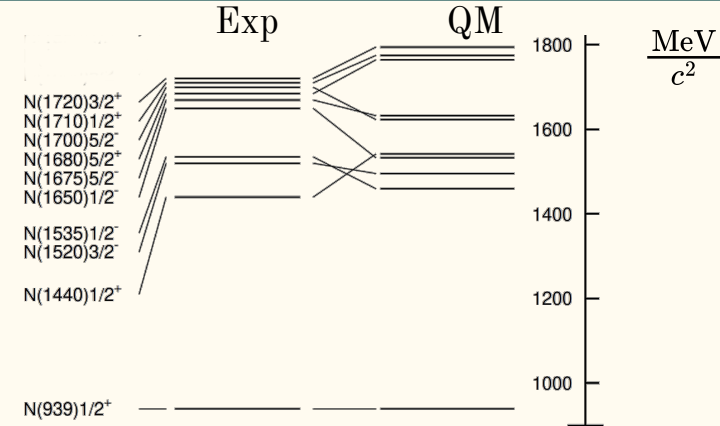
- It was puzzling that up until ~ 2000 , all hadrons could be described by conventional baryons/mesons



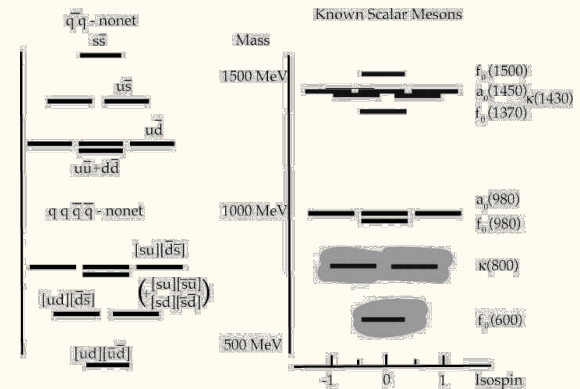
- Exotic hadrons possible



- However, many questions
 - why is the Roper, $N(1440) 1/2^+$, so light?
 - why is the $\Lambda(1405)$ lighter than its nucleon counterpart, i.e. the $N(1535) 1/2^-$?
 - is the nonet of light scalar mesons better described as tetraquarks?
 - and lots more...



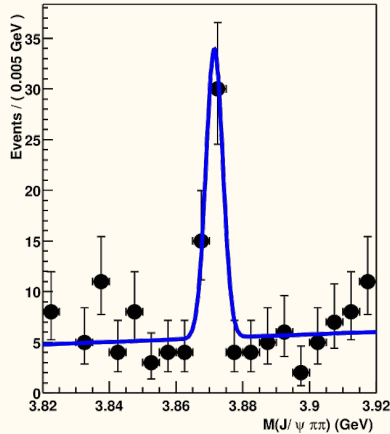
[Particle Data Group, *PTEP* 2022 083C01]



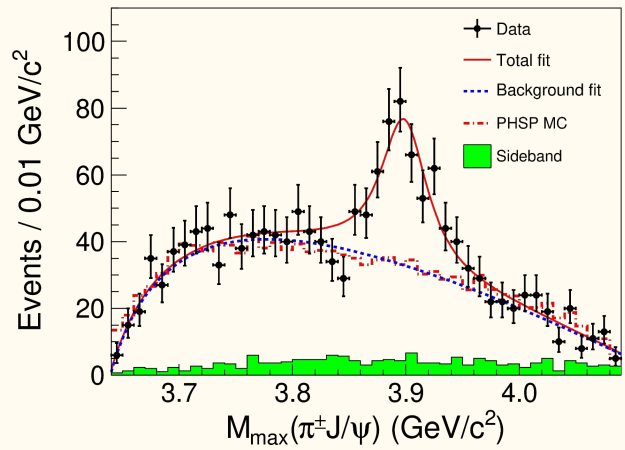
[R.L. Jaffe hep-ph/0409065]

Exotic candidates

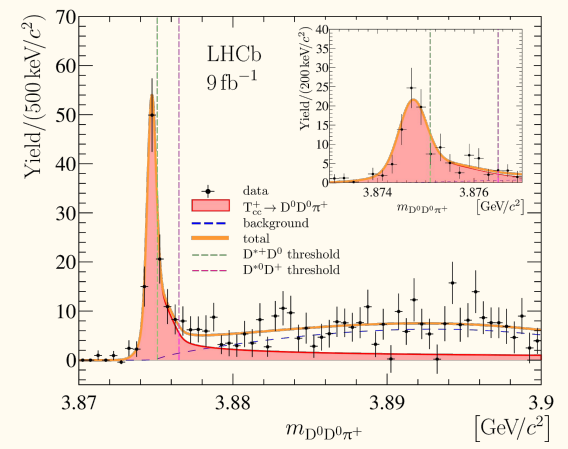
Since the early 2000s, several resonances that do not fit in the quark model have appeared



[Belle hep-ex/0309032]



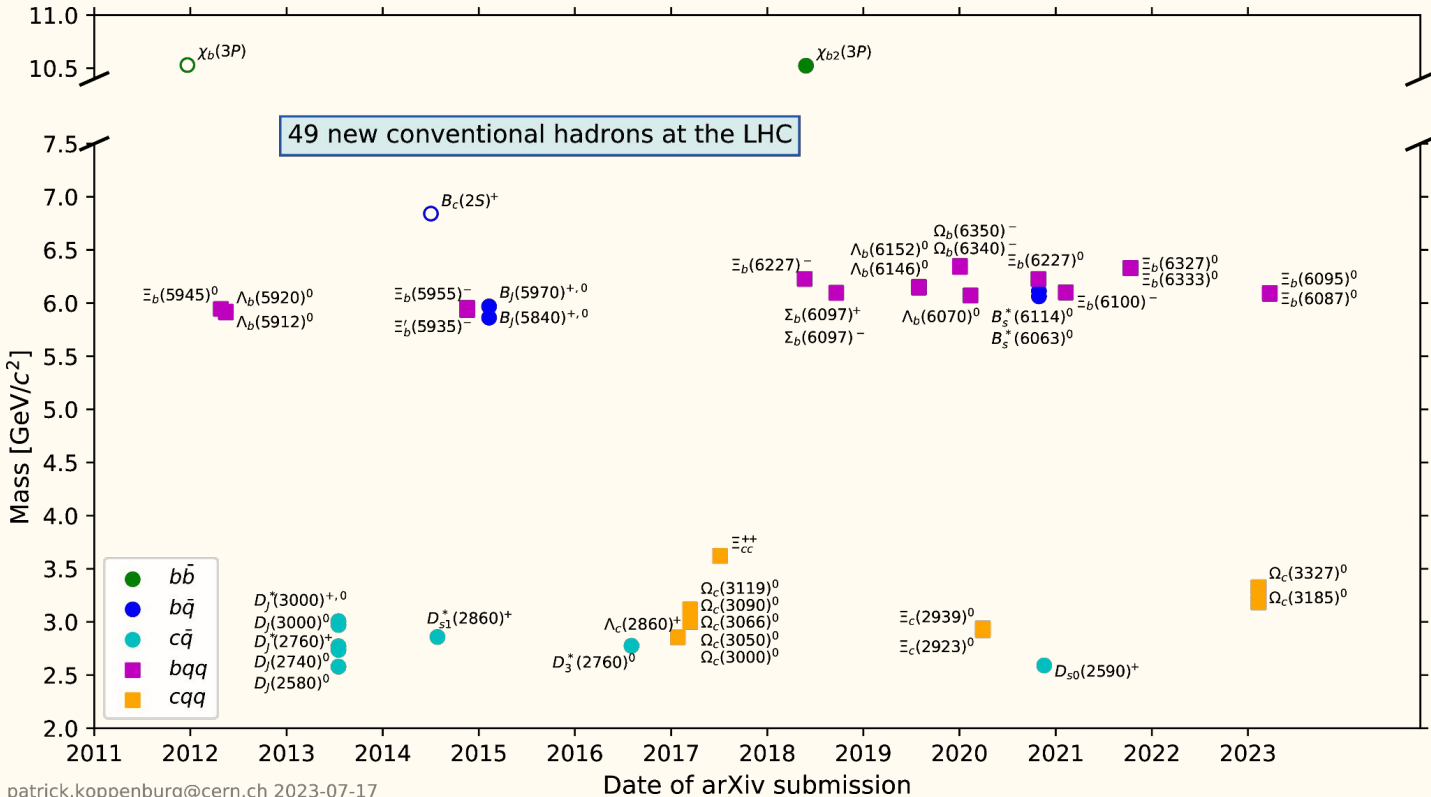
[BESIII 1303.5949]



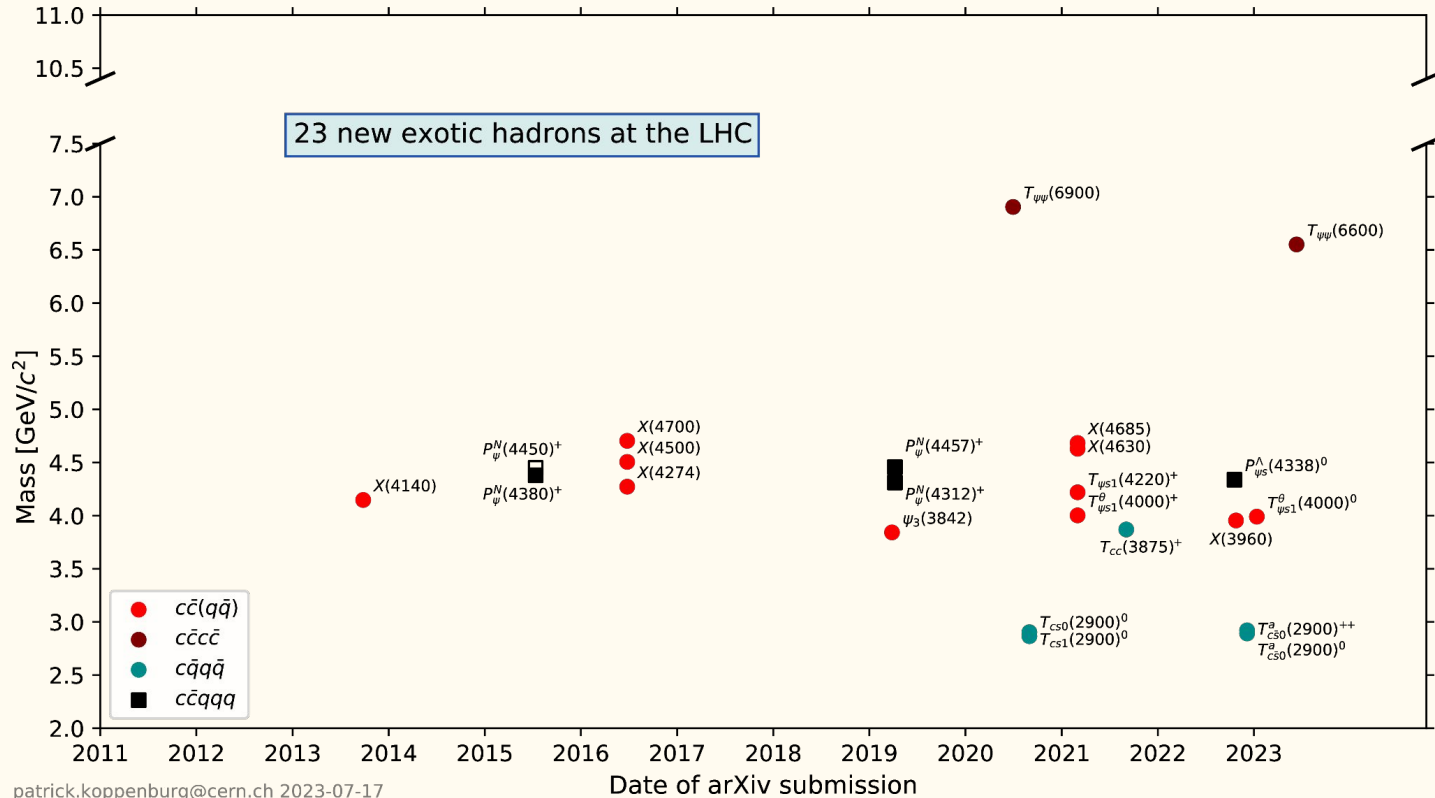
[LHCb 2109.01056]

- $X(3872)$ - could be the $\chi_{c1}(2P)$, but very different from expected behavior
- $Z_c(3900)^+$ - charged charmonium-like state, must be more than $c\bar{c}$
- $T_{cc}^+(3875)^+$ - two charm and overall integer spin, must be exotic to be a color singlet

New hadrons at the LHC



New hadrons at the LHC

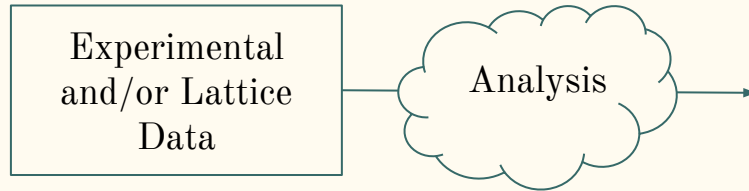


Extracting resonances from data

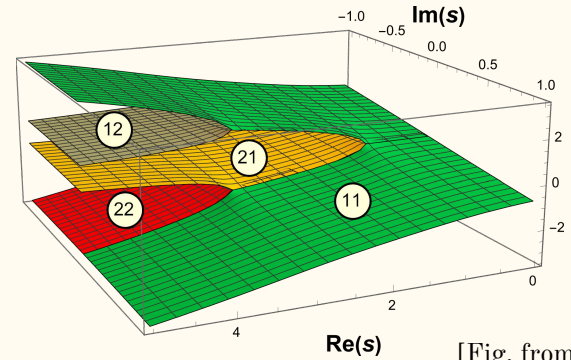
- Resonances appear as poles *off* the real-axis on the *unphysical* Riemann sheets of T_{fi}

$$S_{fi} = \delta_{fi} + i(2\pi)^4 \delta^4(p^f - p^i) T_{fi}$$

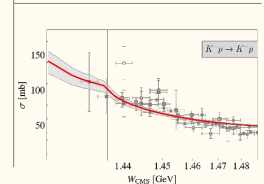
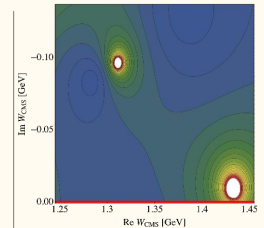
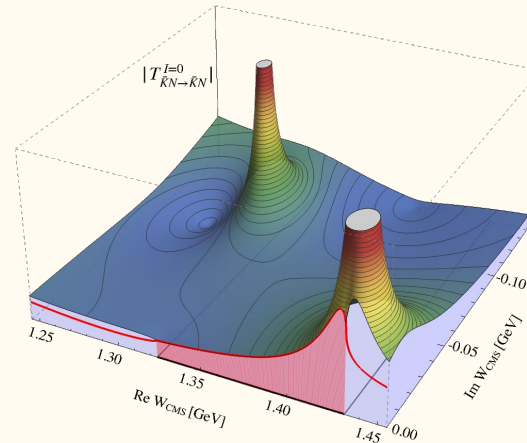
- Each scattering threshold doubles the number of Riemann sheets in which the scattering amplitude lives
- Must infer poles from data on real axis



Goal: Reliably determine resonance properties by robustly extracting pole position and residue!



[Fig. from 2212.07856]



[M. Mai 2010.00056]

Lüscher two-particle formalism

Compact formula for quantization condition

$$\det \left[F(E_2, \mathbf{P}, L)^{-1} + \mathcal{K}_2(E_2^*) \right] = 0$$

E - finite-volume energies

\mathcal{K}_2 - 2-to-2 K-matrix

F - known geometric function

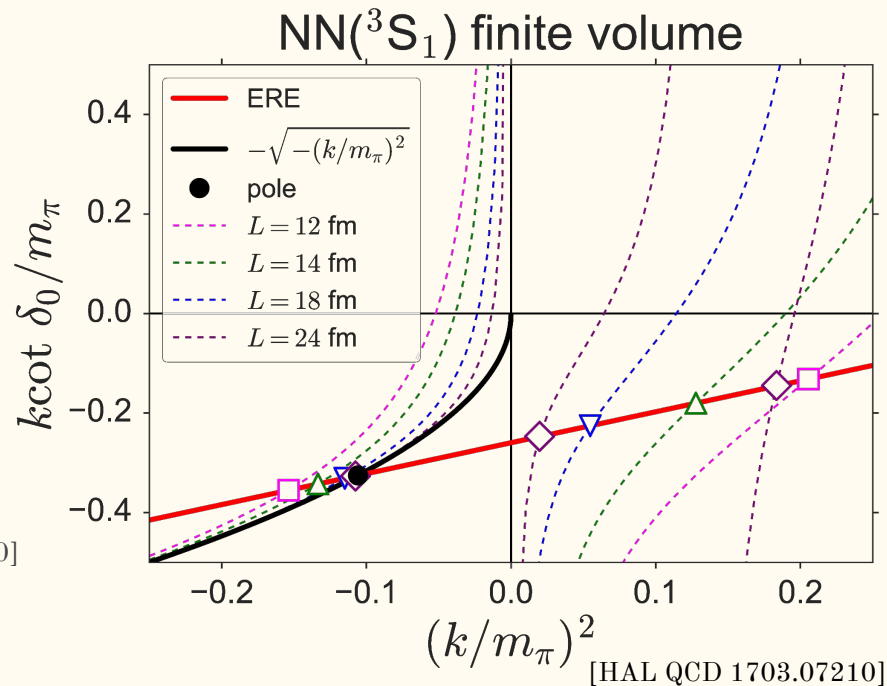
Caveats:

- truncated at some max ℓ
- only valid above left-hand cut and below 3 (or 4) particle threshold
- assumes continuum energies
- ignores exponentially small contributions

Alternatives:

- HAL QCD potential method [S. Aoki & T. Doi 2003.10730]
- Spectral functions from Euclidean correlators [J. Bulava & M. Hansen 1903.11735]

Truncated at $\ell_{\max} = 0 \Rightarrow$ one-to-one mapping



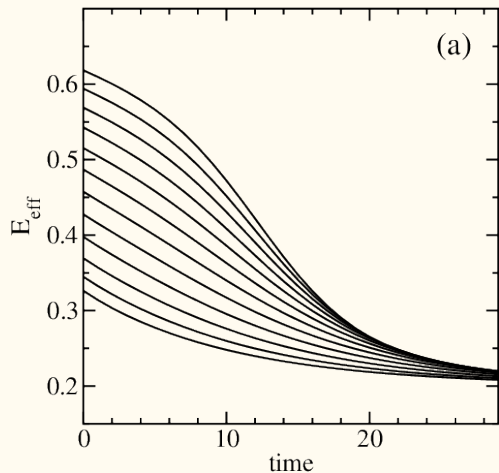
Finite-volume spectrum from correlator matrix

[Lüscher & Wolff *Nucl.Phys.B* 339, 222 (1990)]
[ALPHA 0902.1265]

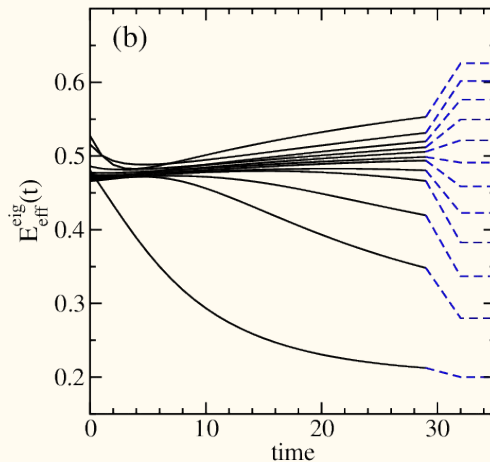
$$C_{ij}(t) = \langle \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) \rangle = \sum_{n=0}^{\infty} Z_i^{(n)} Z_j^{(n)*} e^{-E_n t}, \quad Z_j^{(n)} = \langle 0 | \mathcal{O}_j | n \rangle$$

Toy Model:

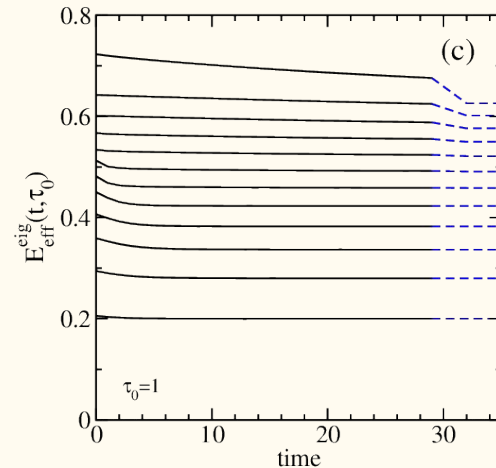
$$E_n = E_{n-1} + \frac{0.08}{\sqrt{n}}, \quad n = 1, 2, \dots, 199, \quad E_0 = 0.20, \quad Z_j^{(n)} = \frac{(-1)^{j+n}}{1 + 0.05(j-n)^2}$$



Diagonal elements of $C(t)$



Eigenvalues of $C(t)$



Generalized eigenvalues of $C(t)$

The $\Lambda(1405)$ - Two poles or one?

An $I=0, J^P=1/2^-$ resonance, with possible lower mass partner $\Lambda(1380)$ not predicted by quark model

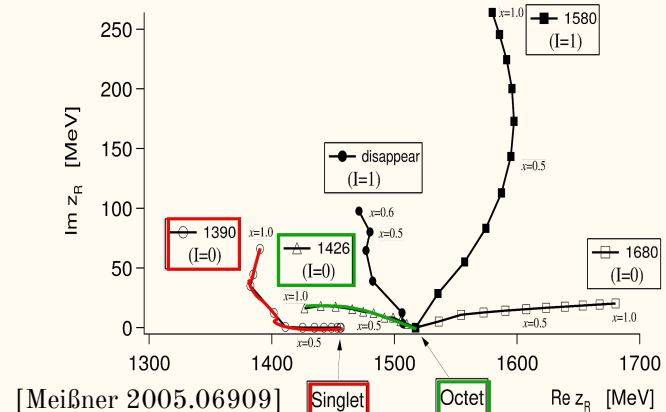
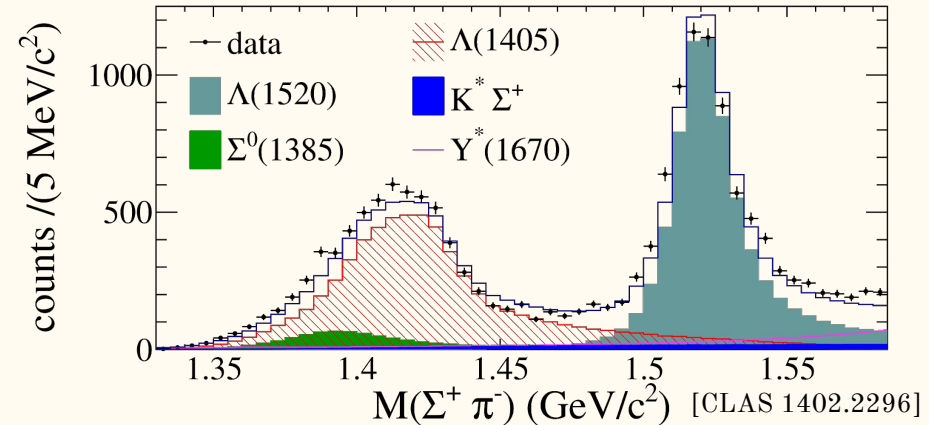
Experiment tends to be mixed regarding this

- CLAS, BGOOD, ALICE, GlueX suggest two poles
- J-PARC consistent with one pole
- Global analysis says one
[Anisovich et al, EPJA 2020]

Unitarized Chiral EFT predicts two poles

Reference	high-mass pole [MeV]	low-mass pole [MeV]
arXiv:1109.3005 & 1201.6549	$1424_{-23}^{+7} - i26_{-14}^{+3}$	$1381_{-6}^{+18} - i81_{-8}^{+19}$
arXiv:1210.3485	$1421_{-2}^{+3} - i19_{-5}^{+8}$	$1388_{-9}^{+9} - i114_{-25}^{+24}$
arXiv:1411.7884, sol. #2	$1434_{-2}^{+2} - i10_{-1}^{+2}$	$1330_{-5}^{+4} - i56_{-11}^{+17}$
arXiv:1411.7884, sol. #4	$1429_{-7}^{+8} - i12_{-3}^{+2}$	$1325_{-15}^{+15} - i90_{-18}^{+12}$

[Particle Data Group, *Chpt.* 83, PTEP 2022 083C01]



What does QCD on the lattice say?

Recent determination of the $I=0$ coupled channel $\pi\Sigma$ - $\bar{K}N$ scattering amplitudes with $m_\pi \sim 200$ MeV

Several parameterizations considered, with effective-range expansion as preferred fit

$$\frac{E_{\text{cm}}}{m_\pi} \tilde{K}_{ij} = A_{ij} + B_{ij} \Delta_{\pi\Sigma}$$

with $B_{00} = B_{11} = 0$, and where $\Delta_{\pi\Sigma} = \frac{E_{\text{cm}}^2 - (m_\pi + m_\Sigma)^2}{(m_\pi + m_\Sigma)^2}$

Final estimate for pole positions consistent with results from UChPT

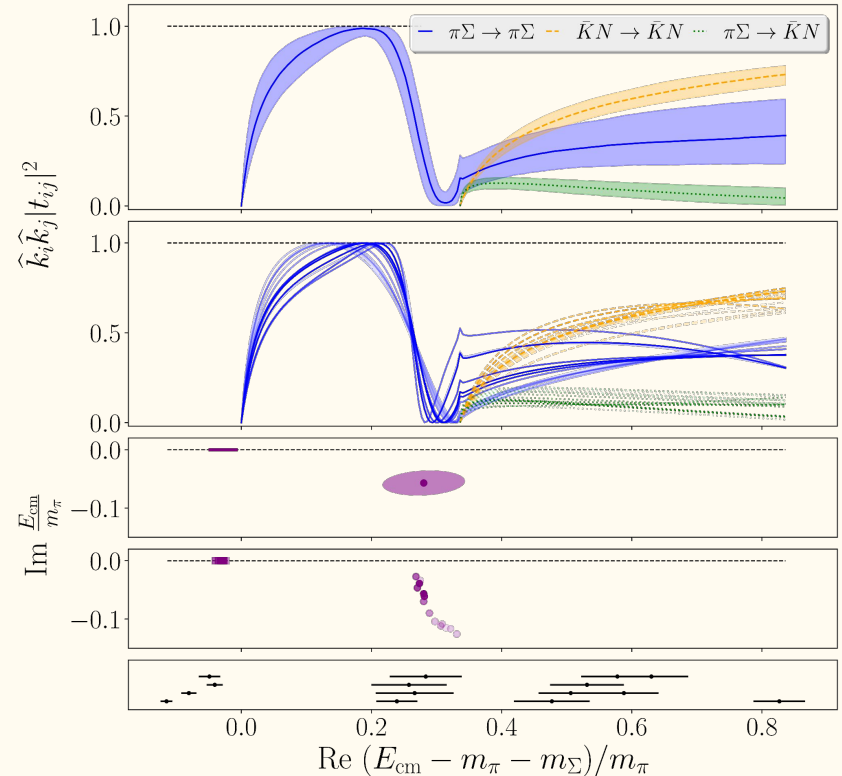
$$E_1 = 1392(9)_{\text{stat}}(2)_{\text{model}}(16)_a \text{ MeV},$$

$$E_2 = [1455(13)_{\text{stat}}(2)_{\text{model}}(17)_a$$

$$- i \times 11.5(4.4)_{\text{stat}}(4)_{\text{model}}(0.1)_a] \text{ MeV}.$$

See talks by:

- Fernando Romero-López (this work), Jul. 31 at 5:00 pm
- Kotaro Murakami (HAL QCD), Jul. 31 at 5:20 pm



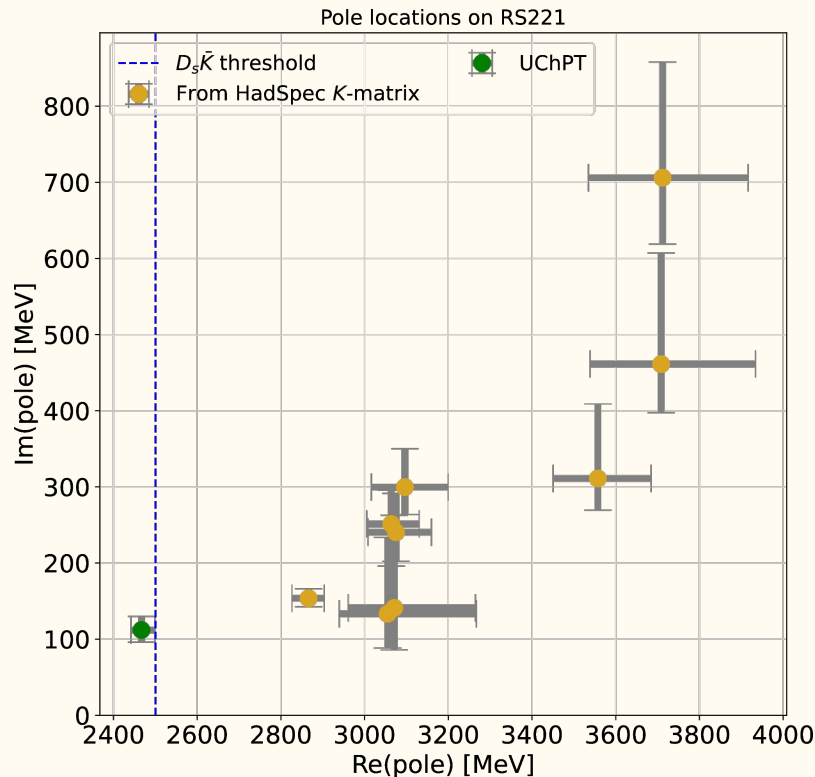
[J. Bulava, ..., **ADH** et al. 2307.10413, 2307.13471]

Two-pole structure for other resonances?

Pole extractions can be unstable if located on sheets not directly connected to the physical one

Lattice results for coupled $D\pi$ - $D\eta$ - $D_s\bar{K}$ showed strong dependence for second $D_0^*(2300)$ resonance pole on the scattering amplitude parameterization, thus was considered ‘unreliable’

Poles with a small residue close to threshold look similar to poles with a large residue far from threshold



[A. Asokan *et al.* 2212.07856]

[Data: HadSpec 1607.07093]

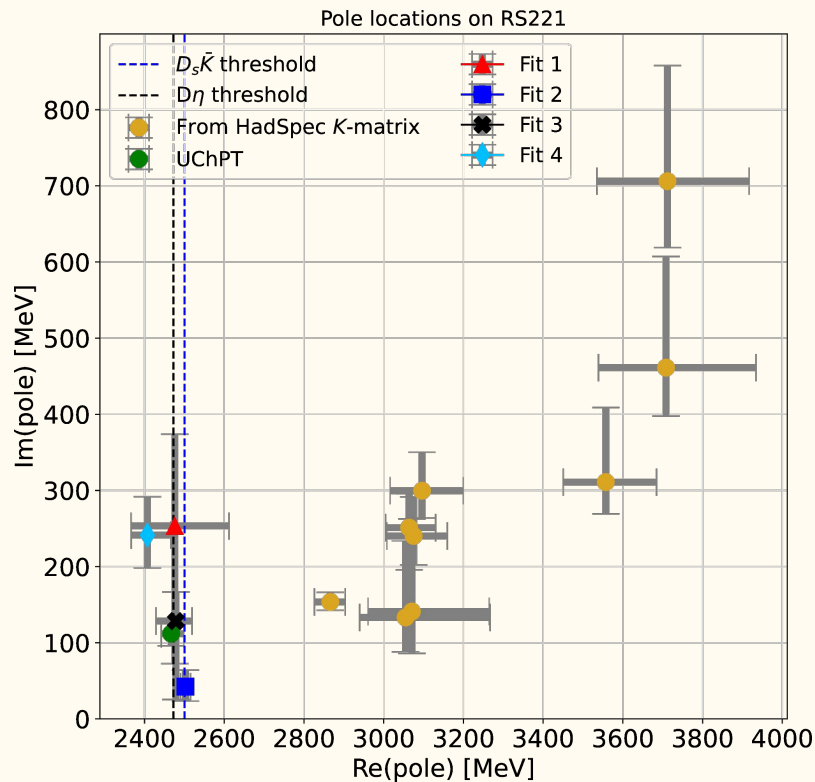
Two-pole structure for other resonances?

Pole extractions can be unstable if located on sheets not directly connected to the physical one

Lattice results for coupled $D\pi$ - $D\eta$ - $D_s\bar{K}$ showed strong dependence for second $D_0^*(2300)$ resonance pole on the scattering amplitude parameterization, thus was considered ‘unreliable’

Poles with a small residue close to threshold look similar to poles with a large residue far from threshold

Using constraints from ChPT for the parameterizations, stable pole positions could be extracted!



[A. Asokan *et al.* 2212.07856]

[Data: HadSpec 1607.07093]

General structure of the scattering amplitude

In cases where the resonance pole lies deep in the complex plane, the full structure should be taken into account

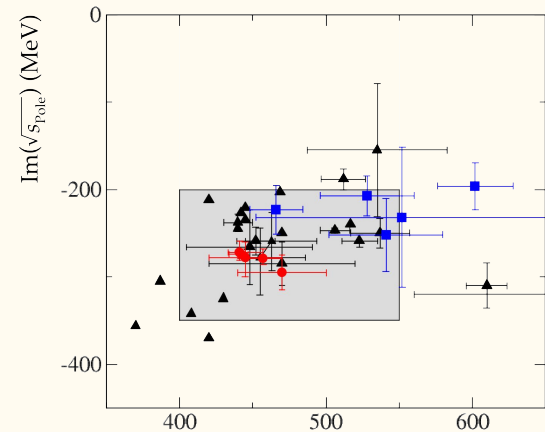
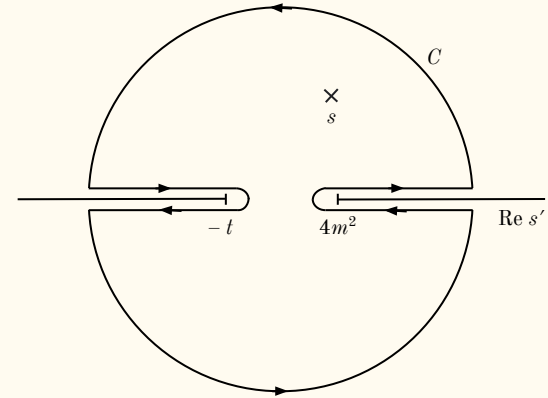
- Unitarity ($\mathcal{S}\mathcal{S}^\dagger = \mathbf{1}$, K -matrix enforces this)
- Crossing (relation between scattering channels)
- Analyticity in Mandelstam variables (from causality)

Analytic properties can be imposed via dispersion relations

$$T(s, t, u) = \underbrace{\frac{1}{\pi} \int_{4m^2}^{\infty} ds' \frac{\text{Im} T(s', t, u)}{s' - (s + i\varepsilon)}}_{\text{threshold cut}} + \underbrace{\frac{1}{\pi} \int_{-\infty}^{-t} ds' \frac{\text{Im} T(s', t, u)}{s' - (s + i\varepsilon)}}_{\text{left-hand cut}}$$

*May require subtractions in order to ignore contour at ∞

Leads to much more stable pole determination! (red points)



σ resonance from lattice with dispersion relations

Dispersion relation techniques can also be used with lattice data

In general, the partial-wave projected amplitudes

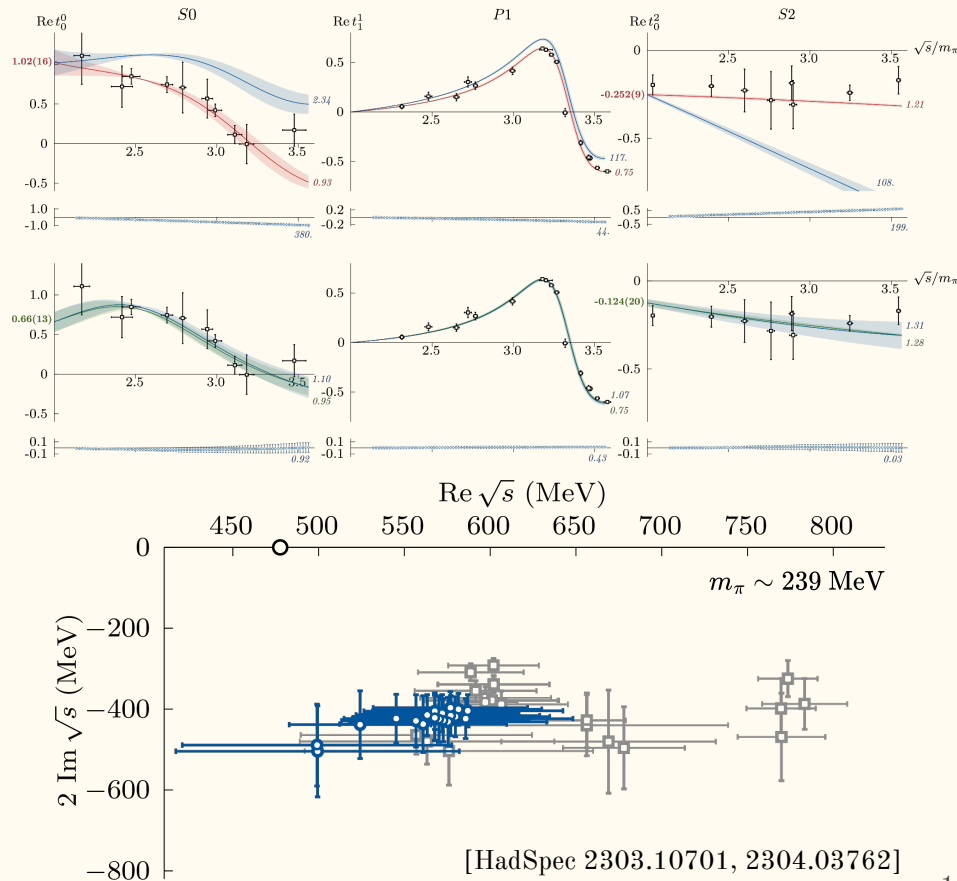
$$t_\ell(s) = \frac{1}{64\pi} \int d\cos\theta_s T(s, t, u) P_\ell(\cos\theta_s)$$

constrained from lattice data do not respect crossing symmetry. However, those that do will respect the dispersion relation

$$\tilde{t}_\ell^I(s) = \tau_\ell^I(s) + \sum_{I', \ell'} \int_{4m_\pi^2}^{\infty} ds' K_{\ell\ell'}^{II'}(s', s) \text{Im} t_{\ell'}^{I'}(s'),$$

i.e. they will result in $\tilde{t}_\ell^I(s) = t_\ell^I(s)$

Again leads to more stable pole positions!



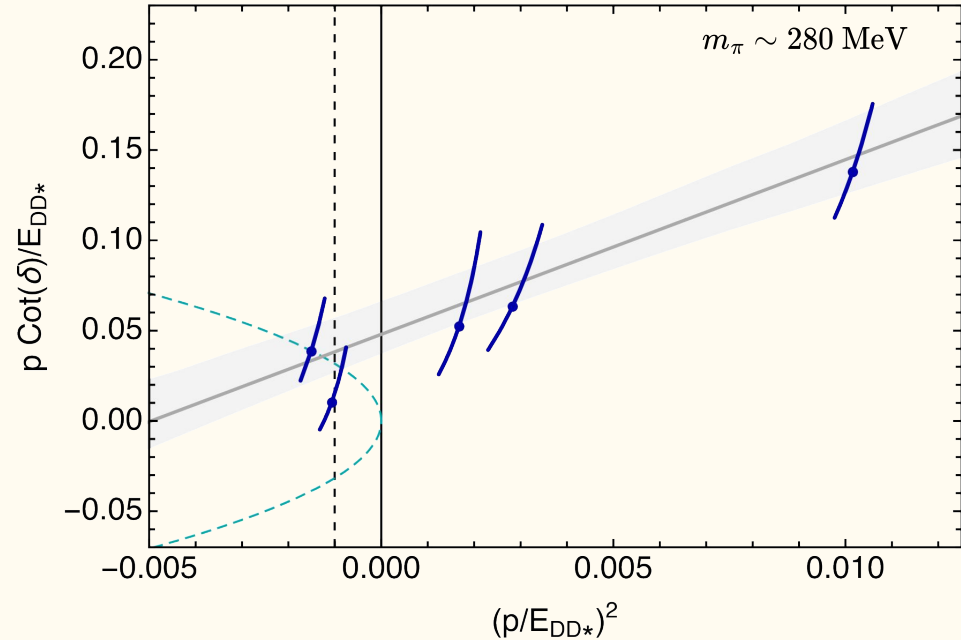
Speaking of left-hand cuts...

The T_{cc}^+ has been studied from the lattice using pion masses in which the D^* -meson is stable

Lüscher quantization condition not valid on the left-hand cut

- For DD^* system, this cut can lie close to threshold, spoiling the conclusions drawn

Work towards extending the quantization condition is underway [Raposo and Hansen 2301.03981]



[M. Padmanath & S. Prelovsek 2202.10110]

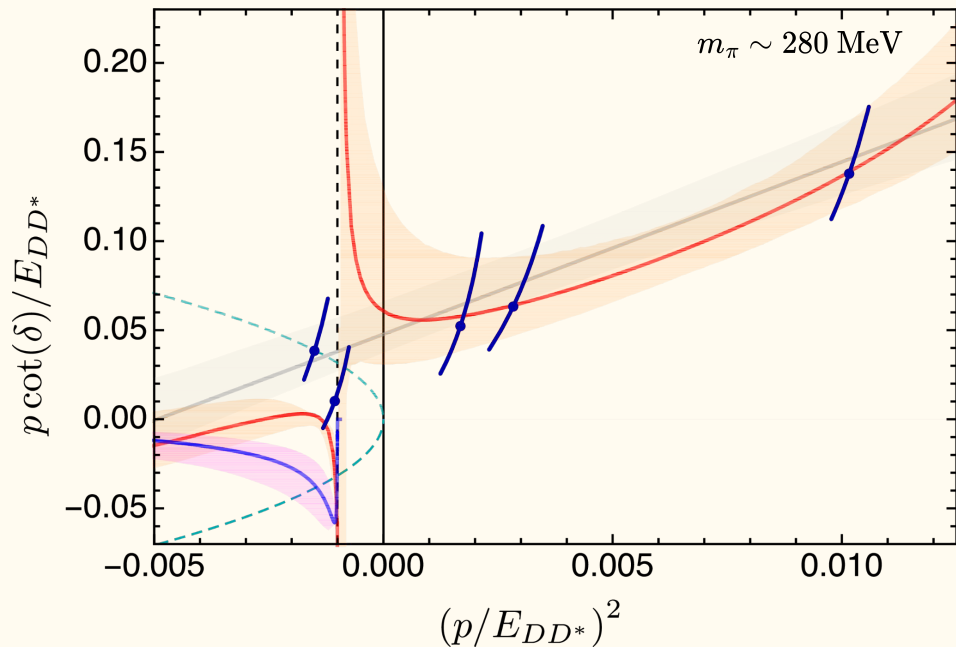
Speaking of left-hand cuts...

The T_{cc}^+ has been studied from the lattice using pion masses in which the D^* -meson is stable

Lüscher quantization condition not valid on the left-hand cut

- For DD^* system, this cut can lie close to threshold, spoiling the conclusions drawn

Work towards extending the quantization condition is underway [Raposo and Hansen 2301.03981]



[M.-L. Du *et al.* 2303.09441]

[Data: M. Padmanath & S. Prelovsek 2202.10110]

Speaking of left-hand cuts...

The T_{cc}^+ has been studied from the lattice using pion masses in which the D^* -meson is stable

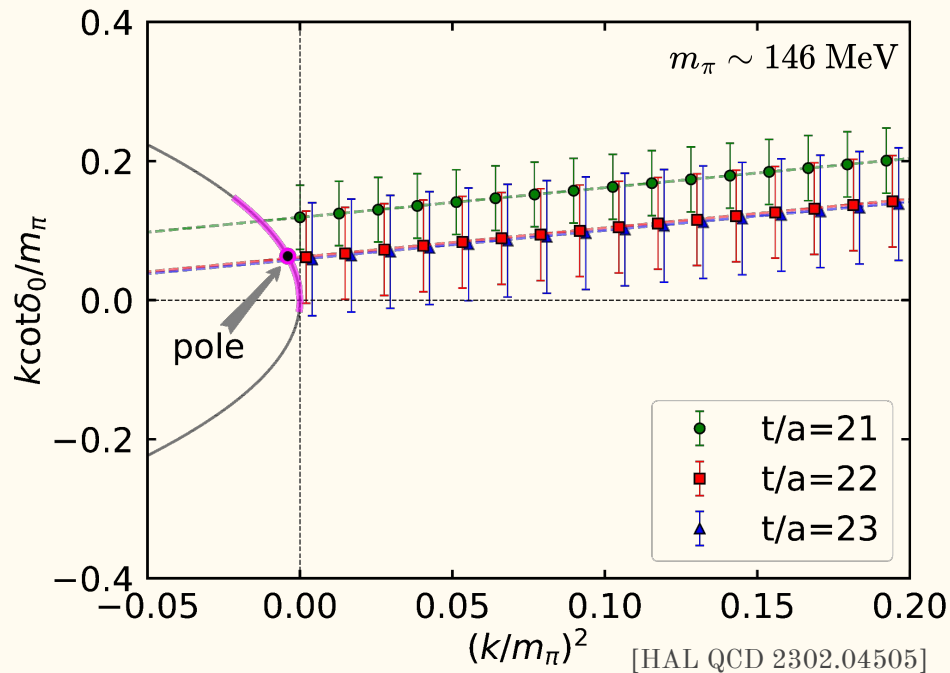
Lüscher quantization condition not valid on the left-hand cut

- For DD^* system, this cut can lie close to threshold, spoiling the conclusions drawn

Work towards extending the quantization condition is underway [Raposo and Hansen 2301.03981]

HAL QCD results just above left-hand cut

Lattice results at $m_\pi \sim 350$ MeV find attractive interaction in $I=0$, only one energy near threshold [CLQCD 2206.06185]

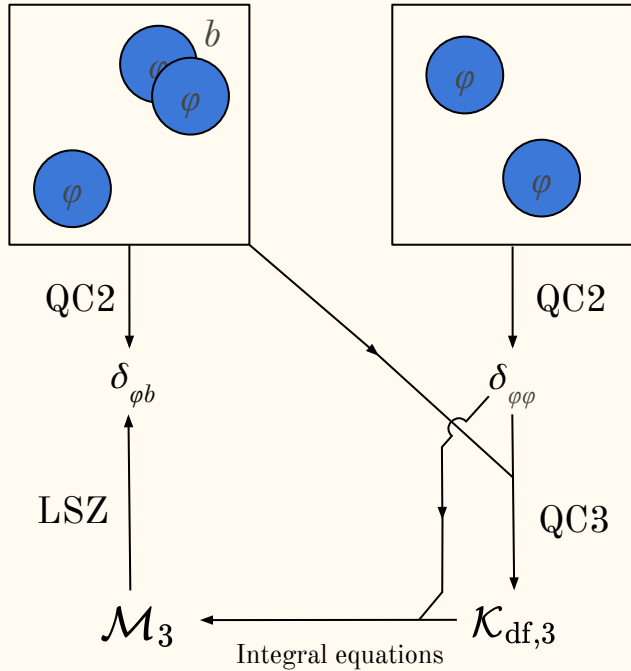


See talks by:

- Jeremy Green (DD^* with distillation), Jul. 31 at 1:30 pm
- Emmanuel Ortiz Pacheco (diquark-antidiquark operators), Jul. 31 at 1:50 pm
- Sinya Aoki (T_{cc} with HAL QCD method), Jul. 31 at 2:10 pm
- Andre Baiao Raposo (Scattering on left-hand cut), Aug. 3 at 1:30 pm

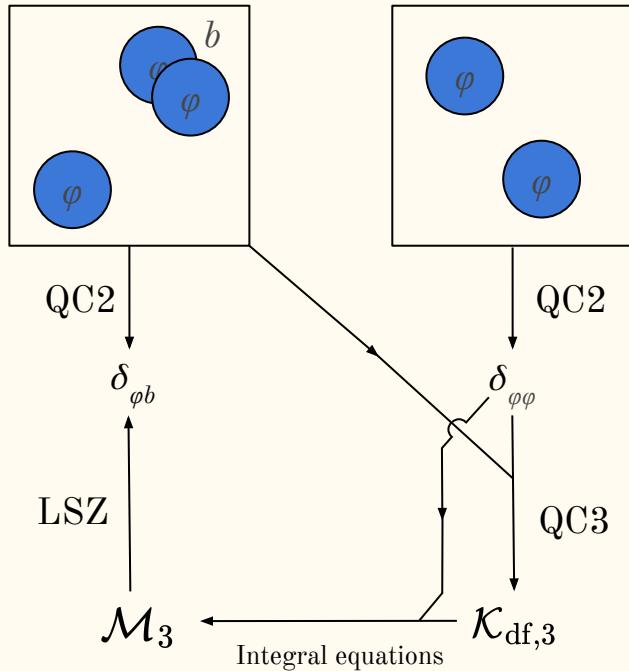
Three-particle formalism saves the day

Consider scattering of a bound state b (of two φ particles) and single particle φ

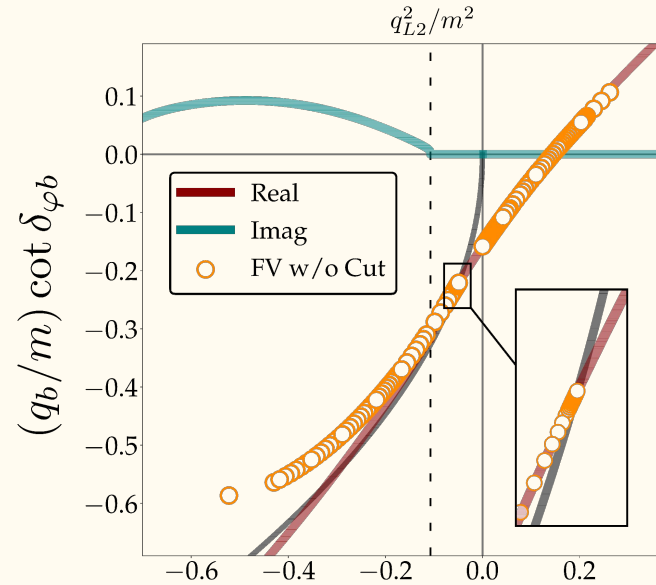


Three-particle formalism saves the day

Consider scattering of a bound state b (of two φ particles) and single particle φ



Breakdown of equivalence between methods beyond left-hand cut demonstrated recently



[Dawid, Islam, Briceño 2303.04394]

[Orange data: F. Romero-López *et al.* 1908.02411]

See talks by:

- Md Habib E Islam (Breakdown from left-hand cut), Aug. 3 at 1:50 pm
- Steve Sharpe (Solving left-hand cut for the T_{cc}), Aug. 3 at 2:10 pm

Three-particle Quantization Conditions

- Most QCD resonance decays involve three or more particles
 $\omega(782) \rightarrow \pi\pi\pi$, $a_1(1260) \rightarrow \pi\pi\pi$, $N(1440) \rightarrow N\pi\pi$
- Many recent developments on the theoretical side (and their applications)
- Three competing formalisms to interpret three-particle finite-volume energies
 - Relativistic Field Theory (RFT) approach [Hansen & Sharpe 1408.5933, 1504.04248, ...]
 - Finite-volume unitarity (FVU) approach [Mai, Döring 1709.08222]
 - Non-relativistic effective field theory (NREFT) [Hammer, Pang, Rusetsky 1706.07700, 1707.02176]
 - Basis for recent formalism studying the Roper [Severt, Mai, Meißner 2212.02171]

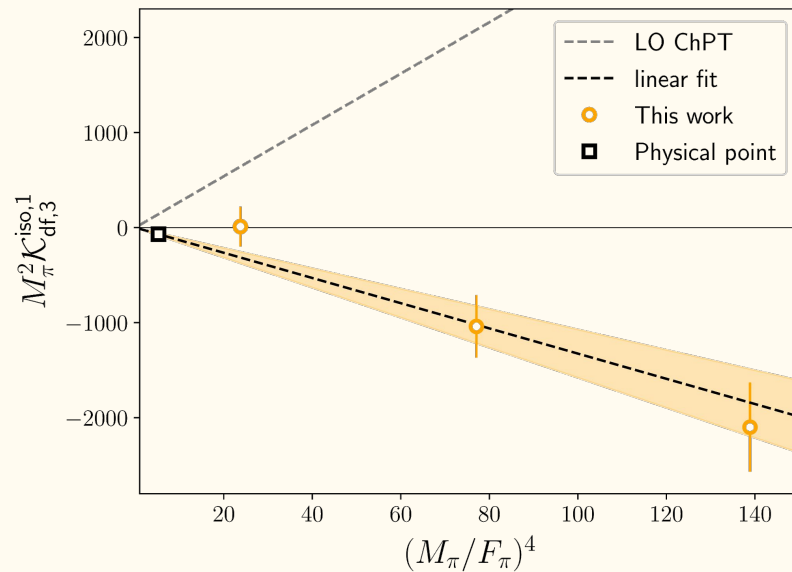
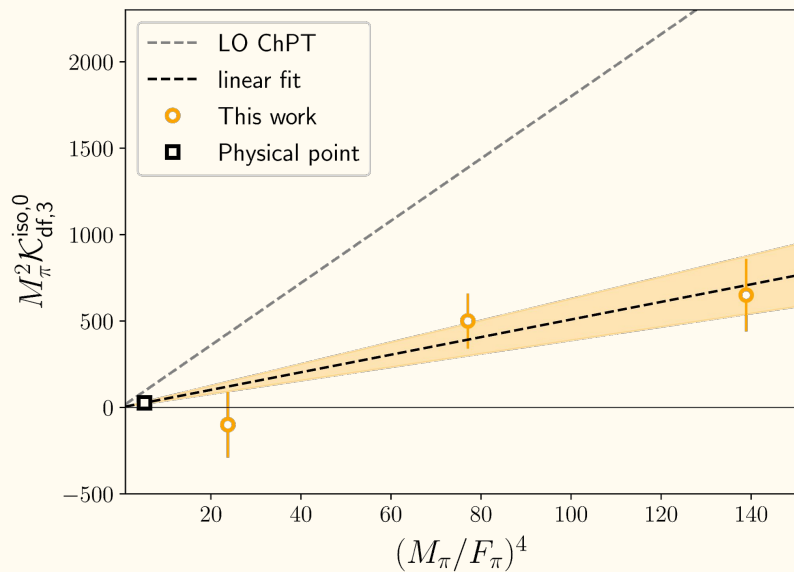
RFT QC

$$\det \left[F_3(E, \mathbf{P}, L)^{-1} + \mathcal{K}_{\text{df},3}(E^*) \right] = 0$$

- Equation in $k\ell m$ basis (spectator - dimer)
- F_3 contains both kinematic functions and the two-particle K-matrix
- $\mathcal{K}_{\text{df},3}$ is a real, analytic, infinite-volume quantity but is scheme-dependent
- Must solve integral equation to obtain three-particle scattering amplitude

Three-pion K -matrix at NLO in ChPT

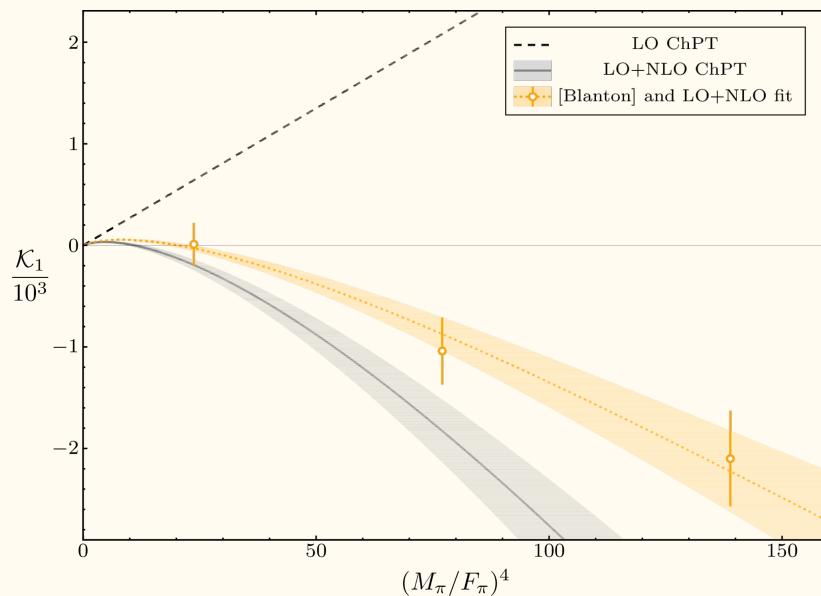
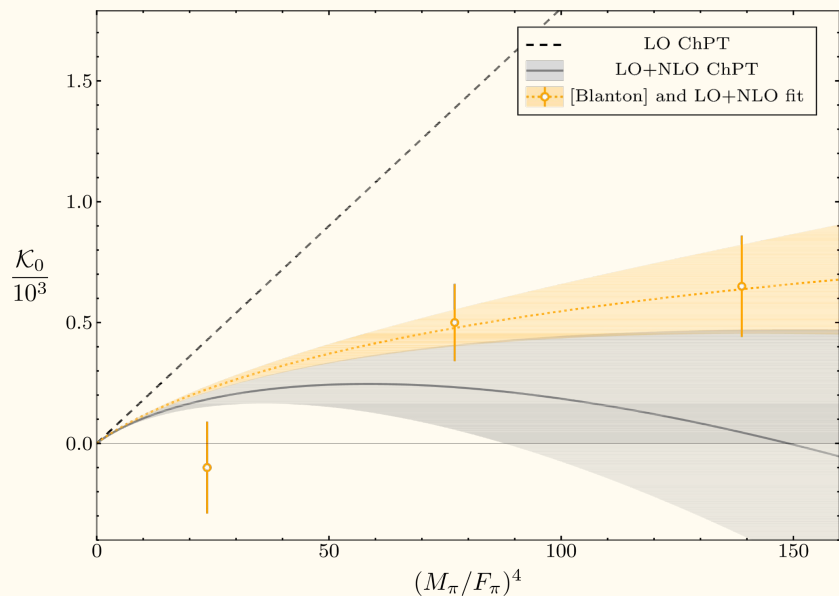
Initial worry regarding comparison to LO ChPT [T. Blanton, **ADH** *et al.* 2106.05590]



Three-pion K -matrix at NLO in ChPT

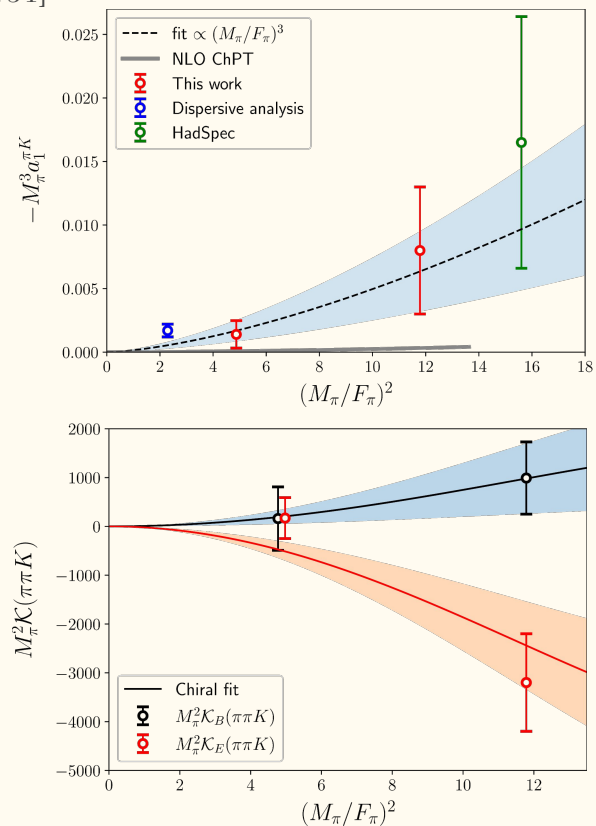
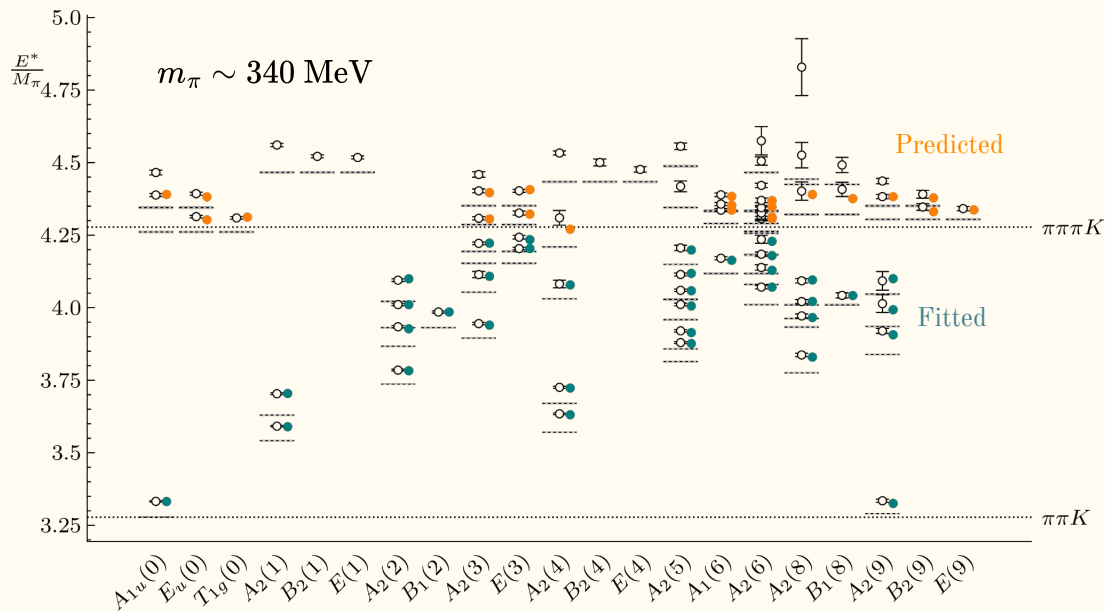
Initial worry regarding comparison to LO ChPT [T. Blanton, **ADH** *et al.* 2106.05590]

Significant NLO correction gives better agreement with lattice! [J. Baeza-Ballesteros *et al.* 2303.13206]



Mixed flavor three-hadron systems

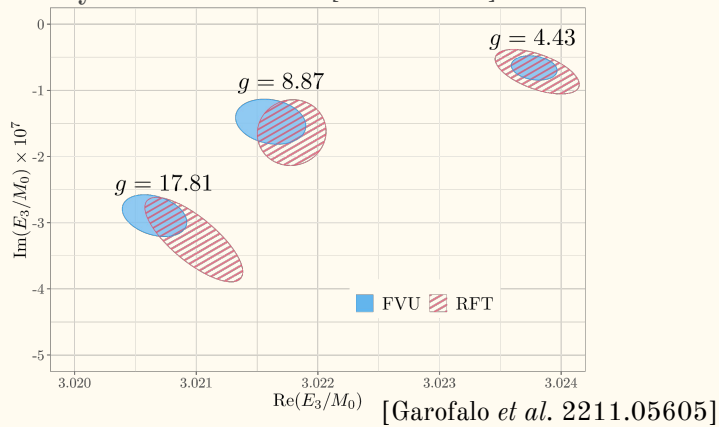
- Formalism for “2+1” flavor systems recently developed [T. Blanton, *et al.* 2111.12734]
- Application to $\pi\pi K$ and $KK\pi$ systems [Z. Draper, **ADH** *et al.* 2302.13587]



Present and future for three particles on a lattice

Recent progress

- RFT/FVU formalisms equivalent [1905.12007, 2007.16190, 2208.10587]
- Non-maximal isospin [2003.10974]
- Non-degenerate mesons [2106.05590]
- Integral equations with resonances [2211.05605]
- Non-zero spin [2303.10219, 2304.13635]
- Analytic continuations [2303.04394]

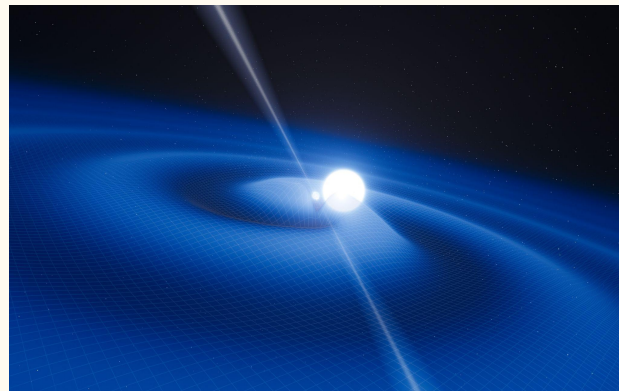


See talks by:

- Sebastian Dawid (Analytic continuations), Aug. 3 at 2:50 pm
- Zachary Draper (3 spinning particles), Aug. 3 at 3:10 pm

On the horizon?

- Formalism for proper extraction of T_{cc}^+
- Formalism for Roper resonance
- Lattice data with three-neutron formalism

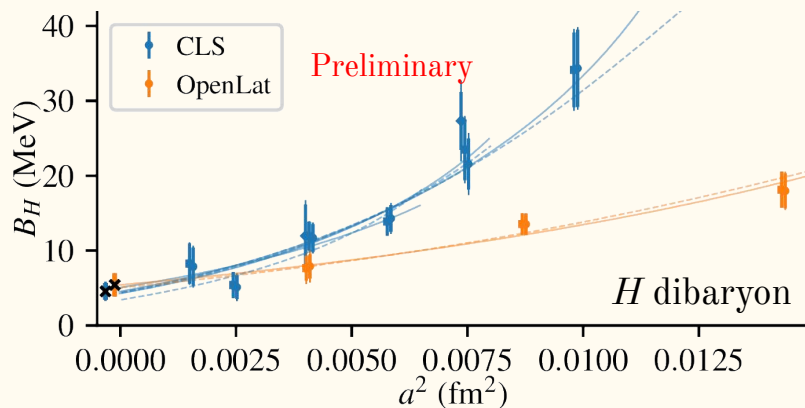


[Artist's impression of the pulsar PSR J0348+0432 and its white dwarf companion, Credit: ESO/L. Calçada, <https://www.eso.org/public/images/eso1319c/>]

But we should not get too ahead of ourselves regarding three nucleons before...

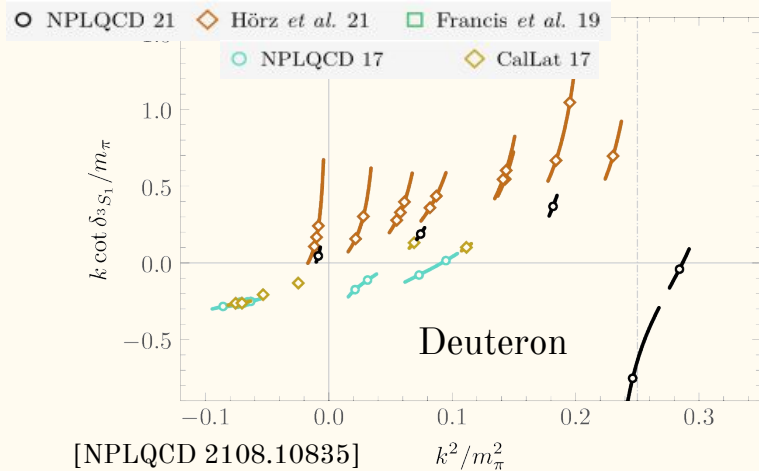
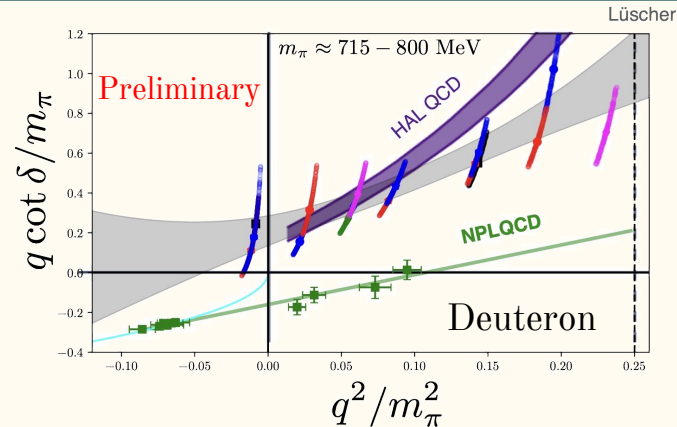
Two-baryon calculations

- NN results in continuum see only virtual bound state [Mainz]
- GEVP results see no bound state, asymmetric correlators do [NPLQCD]
- Lüscher and HAL QCD method agree on same ensemble [CoSMoN]
- Universality of binding energy observed from two actions [BaSc: Mainz+CoSMoN]



See talks by:

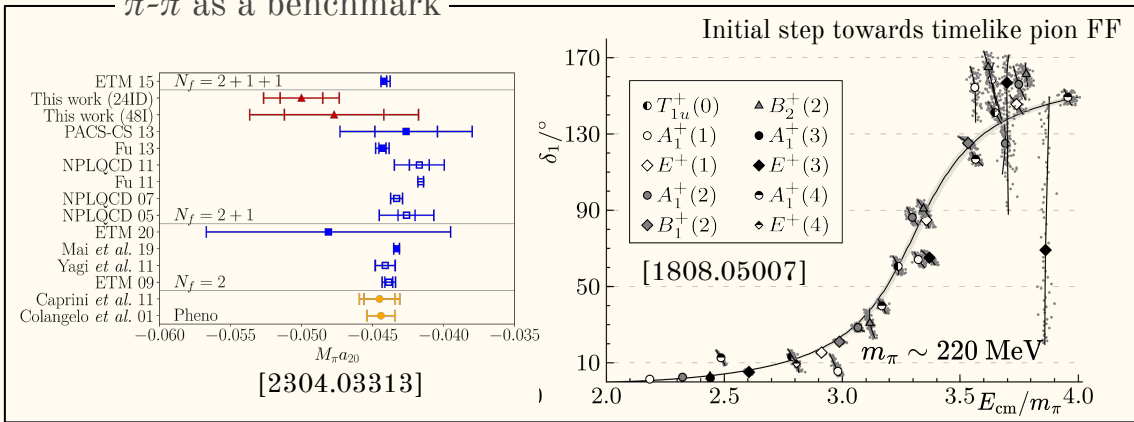
- Takahiro Doi (Nucleon-hyperon, HAL QCD), Aug. 3 at 5:00 pm
- André Walker-Loud (NN, HAL QCD vs Lüscher), Aug. 3 at 5:20 pm
- Phiala Shanahan (NN, $m_\pi \sim 800, 170$ MeV), Aug. 3 at 5:40 pm
- Anthony Francis (SWF from OpenLat), Aug. 4 at 9:00 am



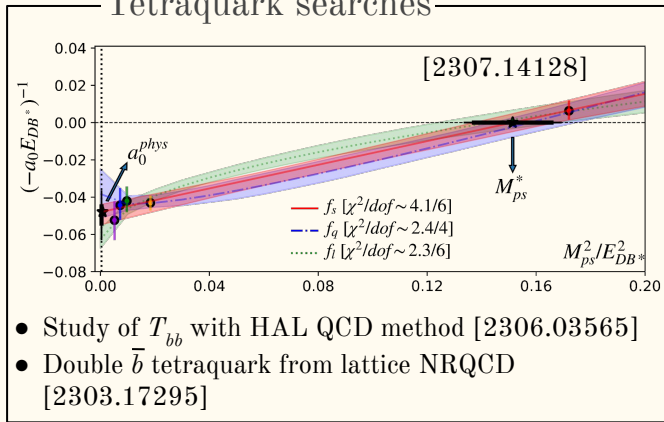
[NPLQCD 2108.10835]

Incomplete selection of applications of formalism

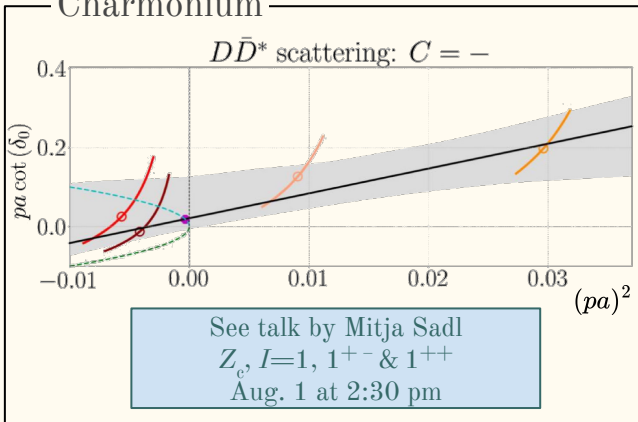
$\pi\text{-}\pi$ as a benchmark



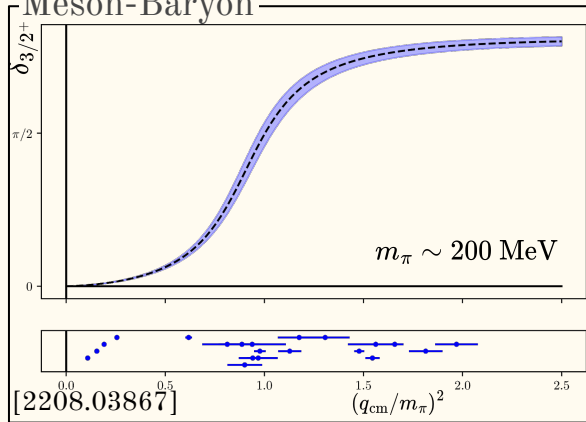
Tetraquark searches



Charmonium



Meson-Baryon



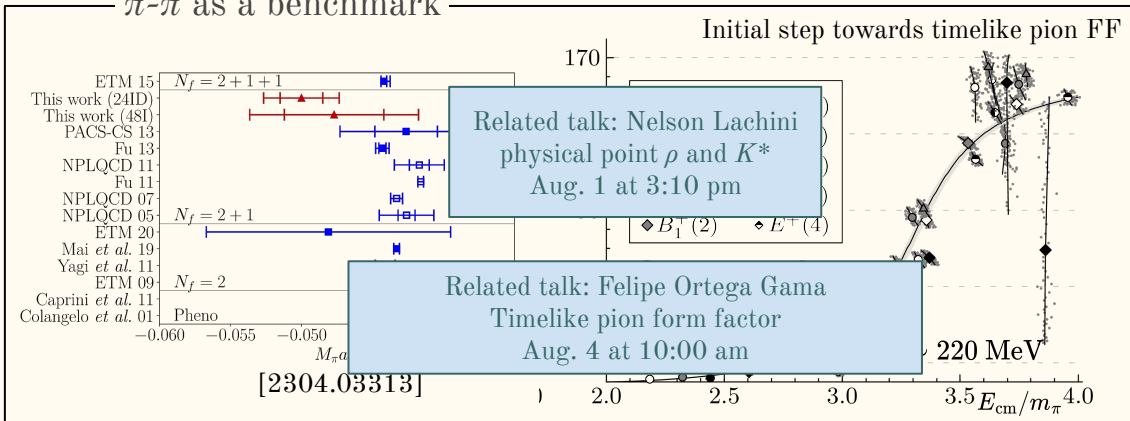
Studied with the HAL QCD method as well
 [2210.05395]

	$m_\pi(\text{MeV})$	$M_\pi a_0^{3/2}$
Alexandrou <i>et al.</i> [2307.12846]	139	-0.13(4)
Bulava <i>et al.</i> [2208.03867]	200	-0.2735(81)
ChPT [1003.4444]	140	-0.0785(32)
Unitarized ChPT [1012.2233]	140	-0.0894(17)
Phenomenology	140	-0.101(4)

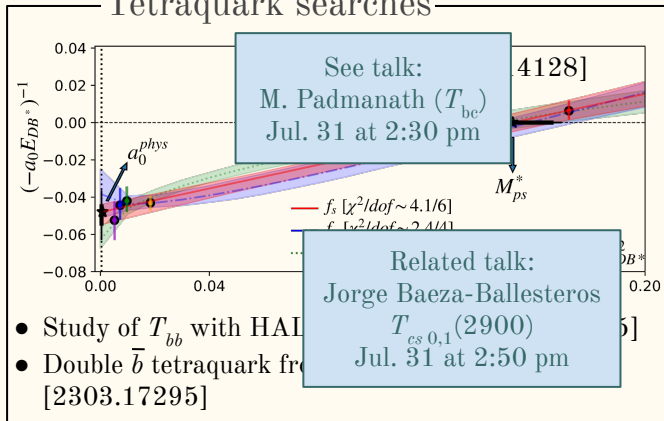
Initial step towards $N \rightarrow \Lambda$ transition FF
 [2211.12278]

Incomplete selection of applications of formalism

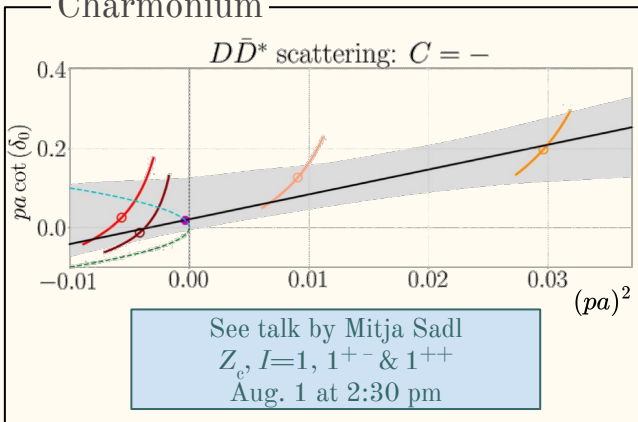
π - π as a benchmark



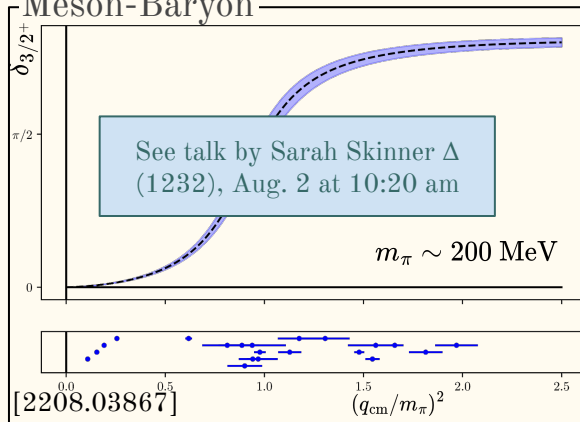
Tetraquark searches



Charmonium



Meson-Baryon



Studied with the HAL QCD method as well
[2210.05395]

Related talk: Anthony Grebe
 $N\pi$ scattering, Aug. 2 at 10:00 am

Unitarized ChPT [1012.2233]	140	-0.0894(17)
Phenomenology	140	-0.101(4)

Initial step towards $N \rightarrow \Lambda$ transition FF
[2211.12278]

Conclusions and Outlooks

- Robust extractions of pole positions from both lattice and experimental data
- Applications of two-hadron interactions involving baryons becoming more prevalent
- Left-hand cuts beginning to be dealt with
- Progress for three particle systems continues

Exciting time for spectroscopy!

- Dispersion relations for other broad resonances deep in the complex plane
- Control over systematics becoming the norm
 - Excited-state contamination and discretization errors
- Three-particle formalism opening up possibilities to study new resonances
 - Roper, T_{cc}^+ , and many others..

Stay tuned!

Thank you
for your
attention!

