

## Scalar Content of Nucleon with the Gradient Flow and Machine Learning

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## Calculation of Hadronic Matrix Elements that Couple to Dark Matter

- A weakly interacting massive particle (WIMP) is a very popular dark matter (DM) candidate, with various ongoing experiments around the world provide rather severe constraints for the parameters of the models.
- A WIMP type of DM particles, due to its assumed large mass, produces a Higgs boson that couples to the various quark flavor scalar density operators taken between nucleon states.
- At zero momentum transfer, the cross section for spin independent elastic WIMP-nucleon ( $\chi N$ ) scattering reads

$$
\begin{equation*}
\sigma_{\chi N} \sim\left|\sum_{f} G_{f}\left(m_{\chi}^{2}\right) f_{T_{f}}\right|^{2} \quad \text { with } \quad f_{T_{f}}=\frac{m_{f}}{m_{N}}\langle N| \bar{q}_{f} q_{f}|N\rangle=\frac{m_{f}}{m_{N}} g_{S}^{f} . \tag{1}
\end{equation*}
$$

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## Lattice Calculation

- Lattice QCD determination of the nucleon matrix elements from first principles.
- Direct calculation is numerically very challenging (disconnected diagrams).
- Gradient flow helps to reduce the statistical error. ${ }^{1}$
- The matrix element we are interested in

$$
\begin{align*}
& g_{S}^{\ell}=\frac{G_{\pi}}{G_{\pi, t}} \cdot\left[\left\langle\mathcal{N} S^{\ell}(t) \overline{\mathcal{N}}\right\rangle^{c o n}+\left\langle\mathcal{N} S^{\ell}(t) \overline{\mathcal{N}}\right\rangle^{d i s c}-\left\langle S^{\ell}(t)\right\rangle\langle\mathcal{N} \overline{\mathcal{N}}\rangle\right] .  \tag{2}\\
& g_{S}^{s}=\frac{G_{\pi}}{G_{\pi, t}} \cdot\left[\left\langle\mathcal{N} S^{s}(t) \overline{\mathcal{N}}\right\rangle^{d i s c}-\left\langle S^{s}(t)\right\rangle\langle\mathcal{N} \overline{\mathcal{N}}\rangle\right] . \tag{3}
\end{align*}
$$

[^0]
## Numerical Strategy

- We compute the quark disconnected diagrams stochastically using the Hutchinson trace method.

$$
\begin{equation*}
\operatorname{Tr}[A] \approx \frac{1}{N_{\xi}} \sum_{\xi}^{N_{\xi}} z_{\xi}^{\dagger} A z_{\xi} \tag{4}
\end{equation*}
$$

- Here $z_{\xi}$ is one of $\{1,-1, i,-i\}$ as in a $Z_{4}$ random noise source.
- We apply the gradient flow and we define $A=K(t) S K^{\dagger}$, where $K^{\dagger}$ is the adjoint kernel which flows the source of the propagator $S$, while $K$ flows the propagator sink.
- Hence, $K S K^{\dagger}=S(t, t)$.

$$
\begin{equation*}
\operatorname{Tr}[S(t)] \approx \frac{1}{N_{\xi}} \sum_{\xi}^{N_{\xi}} z_{\xi}^{\dagger} K(t, 0) S K^{\dagger}(0, t) z_{\xi} . \tag{5}
\end{equation*}
$$

- The gradient flow has the added benefit of a noise reduction effect.


## Signal to Noise Ratio Study

- We combine the disconnected diagram calculation with their respective nucleon-nucleon 2-pt function.
- We study the signal to noise ratio as a function of the number of stochastic sources $N_{\xi}$ in $N_{s r c}$ sources stochastic locations for the nucleon 2pt correlator at $t / a^{2}=1$.



As a consequence, we have determined that with $N_{\xi} \simeq 10$ and $N_{s r c} \simeq 20$, the signal to noise ratio is saturated.

## Quark Condensate $S^{f}(t)=\bar{\chi}^{f}(t, x) \chi^{f}(t, x)$

- Light and strange quark condensates at $t / a^{2}=0$.

- Euclidean time independent quantity.
- A huge number of inversions is required: $N_{\text {src }} \times N_{\text {conf }} \times N_{\text {flow }} \times N_{\text {flavors }}$
- Since they are strongly correlated, we developed a ML method which uses subset of the ensemble and predict the remains.


## Correlations in the Training Set




This correlation provides a non-trivial mapping from $S^{s}\left(t / a^{2}=0\right)$ to $S^{\ell}\left(t / a^{2}=0\right)$.

## Lattice Ensemble and ML Sets

| Ens | $\beta$ | $\kappa_{l}$ | $\kappa_{s}$ | V | $m_{\pi}[\mathrm{MeV}]$ | $m_{N}[\mathrm{GeV}]$ | $N_{\text {conf }}$ | $a[\mathrm{fm}]$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $M_{1}$ | 1.90 | 0.13700 | 0.1364 | $32^{3} \times 64$ | $699.0(3)$ | $1.585(2)$ | 399 | $0.0907(13)$ |

- light quark condensate $S^{\ell}(t)$

| $i_{\text {src }}$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $t / a^{2}$ |  |  |  |  |  |  |  |  |  |  |
| 0.0 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 0.5 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 1.0 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |

- strange quark condensate $S^{s}(t)$

| $i_{\text {src }}$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 0.0 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |

- $i_{\text {src }}=1$ are used for training set.
- test ML method for two cases:
- Between quark flavors: $S^{s}\left(t / a^{2}=0\right)$ and $S^{\ell}\left(t / a^{2}=0\right)$
- For different flow time: $S^{\ell}\left(t / a^{2}=0.5\right)$ and $S^{\ell}\left(t / a^{2}=1\right)$


## Decision Tree Results

- We used Decision Trees with infinite depth for the regression problem, flexible and easy to train (using scikit-learn)
- We make use of translation invariance in order to increase the sample of training data in the ML model.
- The number of training points : $N_{\text {src }} \times N_{\text {conf }}=50 \times 64=3200$




## Quark Condensate Results



- Deviation: $\Delta S=\frac{\left|S^{\text {direct }}-S^{\mathrm{DT}}\right|}{\sqrt{\left(\delta S^{\text {direct }}\right)^{2}+\left(\delta S^{\mathrm{DT}}\right)^{2}}}$


## Combining with Nucleon 2-point Function

$$
\begin{align*}
& R\left(y_{4}, x_{4}, t\right)  \tag{6}\\
& \equiv \frac{\left\langle N\left(y_{4}, 0\right)\right| q q^{f}\left(x_{4}, t\right)|\bar{N}(0,0)\rangle-\left\langle q q^{f}\left(x_{4}, t\right)\right\rangle\left\langle N\left(y_{4}, 0\right) \bar{N}(0,0)\right\rangle}{\left\langle N\left(y_{4}, 0\right) \bar{N}(0,0)\right\rangle} \frac{G_{\pi}(0)}{G_{\pi}(t)}  \tag{7}\\
& =g_{S}^{f}+O\left(e^{-\Delta E\left|x_{4}\right|}\right)+O\left(e^{-\Delta E\left|y_{4}-x_{4}\right|}\right) \tag{8}
\end{align*}
$$




## Comparison with Direct Calculation






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## Computing time saving

Calculation of light quark condensate.

| $i_{\text {sro }}$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 | 10 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $t / a^{2}$ | 1 | 2 |  |  |  |  |  |  |  |  |
| 0.0 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 0.1 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 0.2 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 0.3 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 0.4 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 0.5 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 0.6 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 0.7 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |
| 0.8 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 0.9 | 50 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 | 0 |
| 1.0 | 50 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 | 399 |

- full calculation : $399 \times 9 \times 11 \times 2=79002$ data points
- Traning and input set : $50 \times 11 \times 2+399 \times 9=4691$ confs
- $5.94 \%$ of full calculation is needed for both light and strange quark condensates.


## Results of $\sigma_{f}(t)$ (Preliminary)

- $\sigma_{f}=m_{f} g_{S}^{f}$
- Input is $g_{S}^{\ell}(t=0.5)$ with 9 stochastic sources.




## Light quark connected diagram




## Adding light quark connected and disconnected diagrams




## Summary and Outlook

- ML can help us reduce the computational cost of computing observables in case of strong correlations
- We computed $\sigma_{s}$ terms from disconnected diagram calculations and $\sigma_{\pi N}$ from connected and disconnected contributions.
- Using a machine learning method we achieves similar precision using about $5 \%$ of the computing resources compare to direct calculations.
- Need to analyze other ensembles for the continuum limit
- Other observables (vector, Dslash, tensor) currents intead of scalar are available with the same machine learning method.

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[^0]:    ${ }^{1}$ See DOI:10.22323/1.214.0251 and A. Shidler talk from Monday 31/7 for more details.

