

Simulating Z₂ lattice gauge theory on a quantum computer

Charles et al. - 2305.02361 [hep-lat]

Hank Lamm

August 3, 2023

This work is an outcome of the QCIPU program @ Fermilab

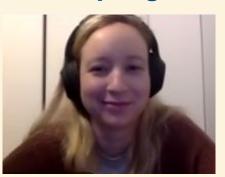


Ruth Van de Water Fermilab

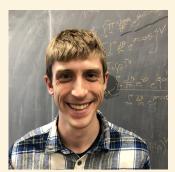


Clement Charles

→ Grad Student @ Maryland



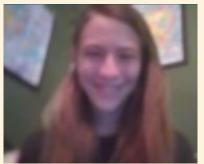
Elizabeth Hardt Argonne



Michael Wagman Fermilab



Florian Herren Fermilab



Sara Starecheski Undergrad @ UIUC



Norman Hogan Grad Student @ NCSU



Erik Gustafson NASA

Quantum Computing Internship For Physics Undergraduates 2020 - Present

3-week summer school for 17 students + year-long internship for 4-5 students with goal to develop diverse

quantum workforce with skills needed to succeed in academia and industry

Young field provides opportunity to build inclusive community

- · Students paid competitive hourly wage
 - -Essential to enable participation by students from all socioeconomic backgrounds.
- Topical lectures by experts in the field
 - -Quantum physics & math, quantum algorithms, error mitigation & correction, quantum hardware. Self-contained and *accessible to all preparation levels*.
- Pair programming on quantum simulators & real devices
 - -Computational exercises in Python + Qiskit on classical and quantum algorithms. Final project simulating **1+1d lattice gauge theory** on real devices.
- Panels and informal discussions on career opportunities
 - Panelists from both academia and industry. Information about applying to and paying for graduate school especially important for first-generation college students.
- Year-long interns perform publishable research

PHYSICAL REVIEW D 106, 114501 (2022)

Primitive quantum gates for an SU(2) discrete subgroup: Binary tetrahedral

Erik J. Gustafson[®], ^{1,*} Henry Lamm[®], ^{1,†} Felicity Lovelace, ^{2,‡} and Damian Musk[®], ^{3,†} Fermi National Accelerator Laboratory, Batavia, Illinois 60510, USA

²Department of Physics, University of Illinois at Chicago, Chicago, Illinois 60607, USA

³Stanford University Online High School, Redwood City, California 94063, USA

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Simulating \mathbb{Z}_2 lattice gauge theory on a quantum computer

Clement Charles, ^{1,2} Erik J. Gustafson, ^{3,4,5} Elizabeth Hardt, ^{6,7} Florian Herren, ³ Norman Hogan, ³ Henry Lamm, ³ Sara Starecheski, ^{3,10} Ruth S. Van de Water, ³ and Michael L. Wagman ³ Department of Physics, The University of the West Indies, St. Augustine Campus, Frinidad & Tobago ² Physics Division, Lawrence Berkeley National Laboratory, Berkeley, CA 9,1720, USA ³ Fermi National Accelentor Laboratory, Berkeley, CA 9,1720, USA ⁴ Quantum Artificial Intelligence Laboratory (QuAIL), NASA Ames Research Center, Moffett Field, CA, 9,4035, USA ⁵ USRA Research Institute for Advanced Computer Science (RIACS), Mountain View, CA, 9,4043, USA ⁵ Department of Physics, University of Hinois at Chicago, Chicago, Rilmos 60007, USA

Department of Physics, University of Illinois at Chicago, Chicago, Illinois 60907, USA
 Advanced Photon Source, Argonne National Laboratory, Argonne, Illinois 60939, USA
 Department of Physics, North Carolina State University, Raleigh, North Carolina 27695, USA
 Department of Physics, Sarah Lawrence College, Bronzville, NY 10708, USA
 Department of Physics, University of Illinois at Urbana-Chanapign, Urbana, IL 61801, USA

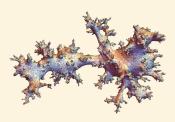
Department of Physics, University of Illinois at Urbana-Cl (Dated: May 5, 2023)

Quantum Computing for Particle Physics, it's a need

- The world is quantum, and we are lucky anything is amenable to classical computers
 - Large-scale quantum computers can tackle computations in HEP otherwise inaccessible
 - This opens up new frontiers & extends the reach of LHC, LIGO, EIC & DUNE

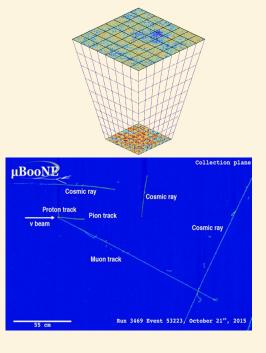


- Ab initio cross sections for colliders and neutrino experiments
- Cosmic inflation and the evolution of matter asymmetry in the early universe
- Explorations of BSM, supersymmetry, and quantum gravity
- Hadronization and Hydrodynamics in Heavy-Ion collisions



While broad, these topics often are formulated as **lattice field theories**

Quantum Simulation for High-Energy Physics Bauer, Davoudi et al. - PRX Quantum 4 (2023) 2, 027001 Wonderful survey of physics questions, methods, and outstanding problems in field



Premature optimization is the root of all evil

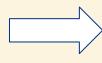
2 is the smallest and only even prime number

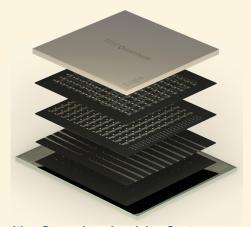
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1+1d Z₂ lattice gauge theory on quantum computer

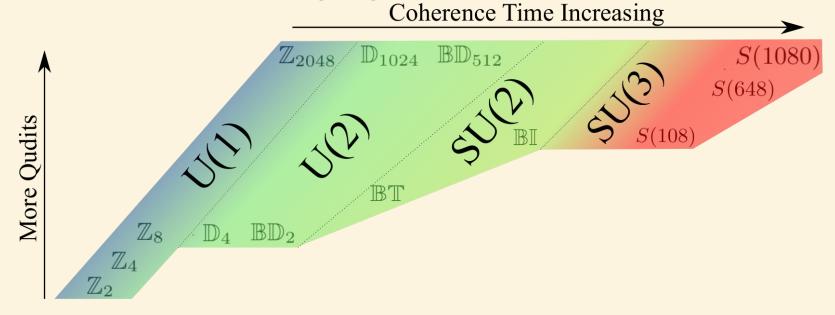






Experiments with a Gauge-Invariant Ising System Michael Creutz, Laurence Jacobs, and Claudio Rebbi *Phys. Rev. Lett. 42, 1390 (1979)* Early simulation of 3+1d 8^4 Euclidean $\mathbf{Z_2}$ pure gauge theory

The ladder of discrete gauge theories in HEP calculations



Gluon Field Digitization for Quantum Computers NuQS collaboration - *Phys.Rev.D 100 (2019) 11, 114501* Demonstrated that S(1080) approximates certain 3+1d SU(3) observables

The ladder of discrete gauge theories in HEP calculations

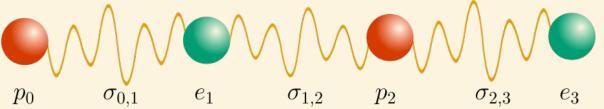
Coherence Time Increasing \mathbb{D}_{1024} More Qudits

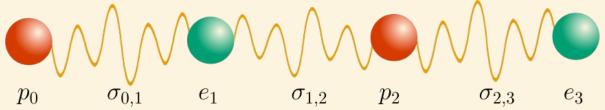
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Digitising SU(2) gauge fields and the freezing transition Hartung et al. - Eur.Phys.J.C 82 (2022) 3, 237 Understanding the scaling of freezing transitions with approximations

$$\beta_{f,U(1)} = \frac{\log(1+\sqrt{2})}{1-\cos\left(\frac{2\pi}{N}\right)} \approx \kappa_2 N^2, \text{ which extends to } \beta_{f,SU(N_c)} \approx \kappa N^{\frac{N_c^2-1}{2}}$$
 But whereas \mathbb{Z}_N can be **taken to** ∞ , **limited** number for $SU(N_c)$

 $\beta \propto \frac{1}{\log(a)} \implies a_f \propto e^{-\beta_f}$

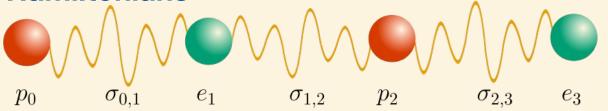




Kogut-Susskind Hamiltonian [with O(a²) errors]

$$H = \sum_{n=1}^{N_s - 1} \left[\frac{1}{2} \sigma_{n,n+1}^x + \frac{\eta}{2} (\bar{\psi}_n \sigma_{n,n+1}^z \psi_{n+1} + h.c.) \right] + m_0 \sum_{n=1}^{N_s} (-1)^n \bar{\psi}_n \psi_n,$$

Hamiltonian Formulation of Wilson's Lattice Gauge Theories Kogut & Susskind *Phys.Rev.D 11 (1975) 395-408* Formulated O(a²) lattice Hamiltonian for LGT with staggered matter



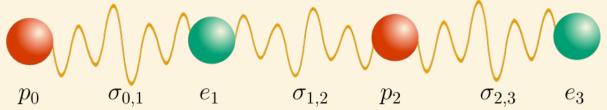
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Qubit Hamiltonian via Jordan-Wigner

$$\begin{split} H = & \frac{1}{2} \sum_{n=0}^{N_s - 1} \sigma_{n,n+1}^x - \frac{m_0}{2} \sum_{n=0}^{N_s - 1} (-1)^n Z_n \\ & + \frac{\eta}{4} \sum_{n=0}^{N_s - 2} (X_n X_{n+1} + Y_n Y_{n+1}) \sigma_{n,n+1}^z. \end{split}$$



Kogut-Susskind Hamiltonian [with O(a2) errors]

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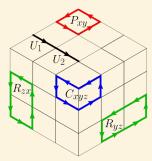
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Alway remember: lattice Hamiltonian is a choice

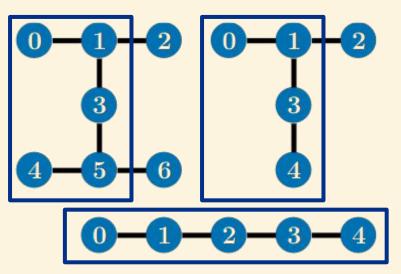
Improved Hamiltonians for Quantum Simulation of Gauge Theories Carena, Lamm, Li, Liu *PRL 129 (2022) 5*Developed quantum circuits for O(a⁴) pure-gauge Hamiltonian

Quantum Simulation of Lattice QCD with Improved Hamiltonians Ciavarella 2307.05593 [hep-lat]
Formulated Hamiltonian with reduced truncation errors

Improved Fermion Hamiltonians for quantum simulations Gustafson & Van de Water - in prep (Talk @ 4:20 PM on Thurs.) Formulating Hamiltonians for ASQTAD fermions



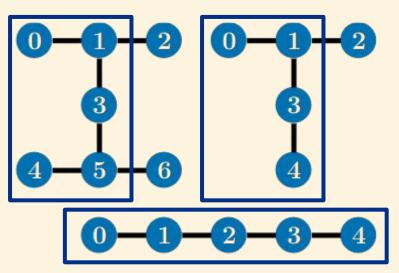
Hamiltonian Gates for Trotterization with restricted connectivity



Restricting to longest linear graph

$$\text{for heavy-polygon with p sides: } \frac{N_p-2}{N_p-1} \leq 86\% \, \frac{\text{BAD!}}{}$$

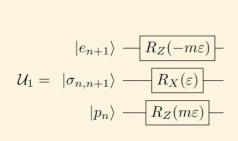
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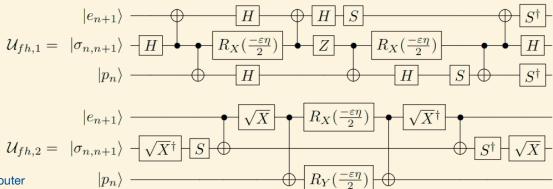


Restricting to longest linear graph

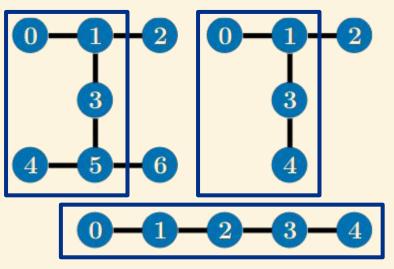
 $ext{for heavy-polygon with } p ext{ sides: } rac{N_p-2}{N_p-1} \leq 86\% ext{ BAD!}$

 ${ t ibm_nairobi} o 43\% ext{ WORSE!}$





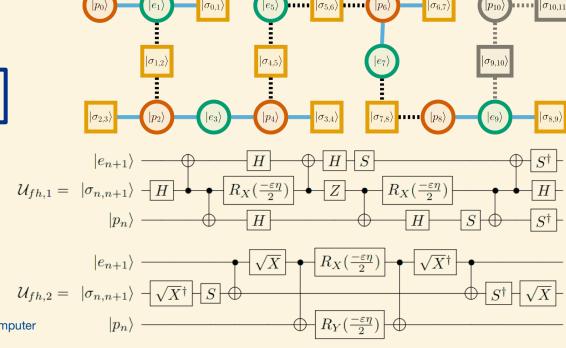
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8/2/2022 Hank Lamm | Simulating Z₂ on Quantum Computer

 $\mathcal{U}_1 = |\sigma_{n,n+1}\rangle \longrightarrow R_X(\varepsilon)$

 $|e_{n+1}\rangle - R_Z(-m\varepsilon)$

 $|p_n\rangle \longrightarrow R_Z(m\varepsilon)$

 Want to measure a correlator after preparing in a superposition of vacuum and "particle" state

$$C(t) = \langle \phi(N_s) | \mathcal{U}^{\dagger}(t) \ O \ \mathcal{U}(t) | \phi(N_s) \rangle$$

= \cos(Mt) + \dots,

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= \cos(Mt) + \dots,

• Trotteriztion introduces discretization errors into correlator, and thus scale setting M $\mathfrak{C}(t/\varepsilon) = \left\langle \phi(N_s) \middle| \mathcal{U}^\dagger(t/\varepsilon)^{N_t} \; O \; \mathcal{U}(t/\varepsilon)^{N_t} \; \middle| \phi(N_s) \right\rangle$

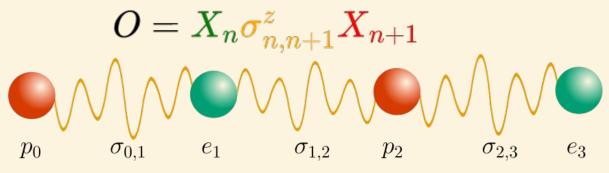
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"Particle" state and operator insertion are given by "meson" excitation operator



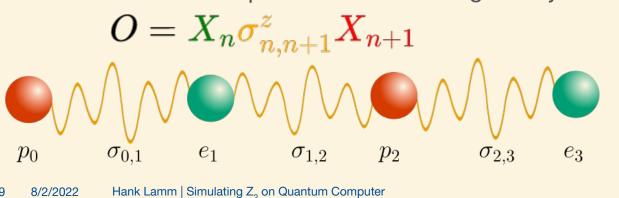
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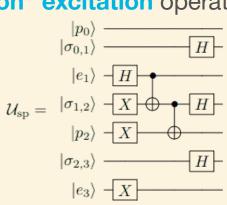
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Trotteriztion introduces discretization errors $\mathfrak{C}(t/\varepsilon) = \langle \phi(N_s) | \mathcal{U}^{\dagger}(t/\varepsilon)^{N_t} O \mathcal{U}(t/\varepsilon)^{N_t} | \phi(N_s) \rangle$ into correlator, and thus scale setting M

"Particle" state and operator insertion are given by "meson" excitation operator

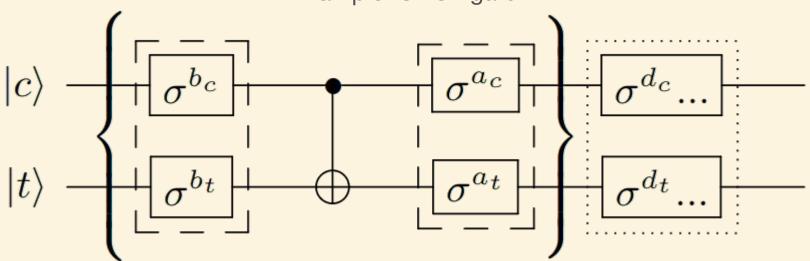




First error mitigation: Pauli Twirling

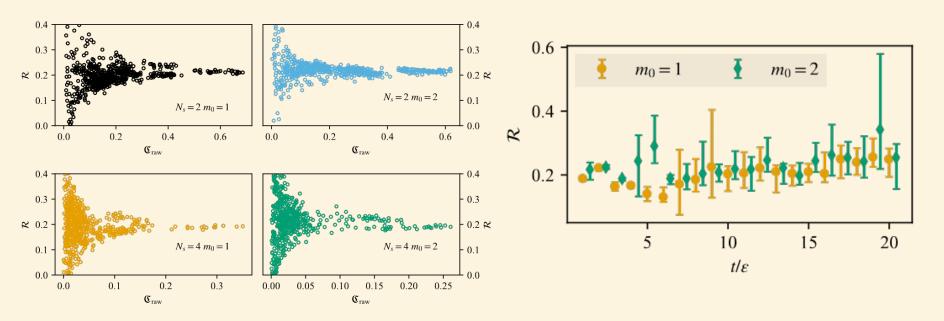
Converts coherent errors to stochastic ones





Simulating one-dimensional quantum chromodynamics on a quantum computer: Real-time evolutions of tetra- and pentaquarks
Y. Y. Atas et al., 2207.03473 [quant-ph]
Early example of use in 1+1d SU(3) lattice simulation

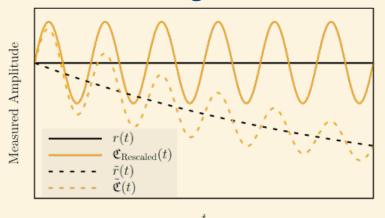
Second error mitigation: Readout error mitigation



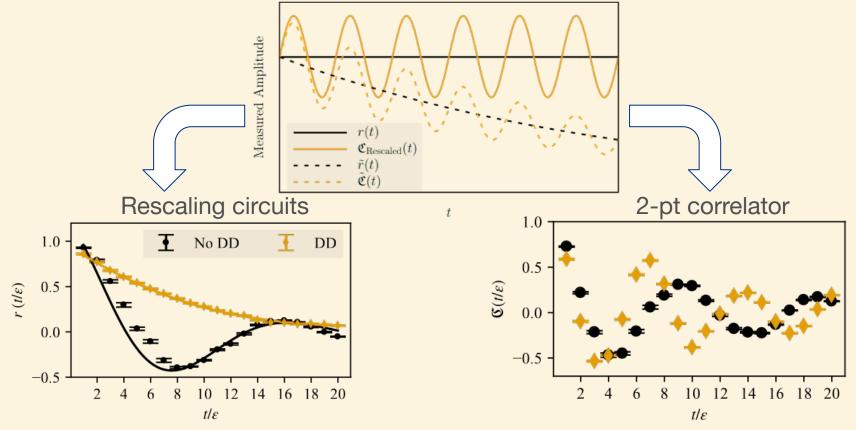
Error arising from readout typically ~20% albeit dependent on correlator value

Genuine 12-qubit entanglement on a superconducting quantum processor M. Gong et al. Phys. Rev. Lett. 122, 110501 (2019)
Good example of inversion matrix method for readout error mitigation

Third error mitigation: Rescaling



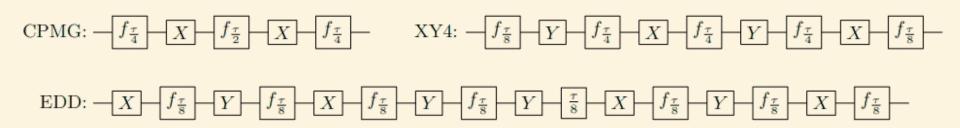
Third error mitigation: Rescaling



Self-mitigating Trotter circuits for SU(2) lattice gauge theory on a quantum computer S. A Rahman, R. Lewis, E. Mendicelli, and S. Powell, Phys. Rev. D 106 First demonstration of rescaling in a lattice simulation

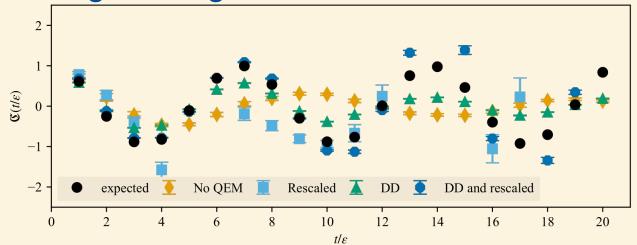
Fourth error mitigation: Dynamical Decoupling

Perform gates operators on spectators to prevent noise

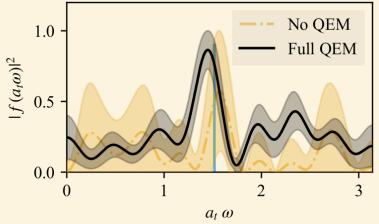


For our lattice simulations on IBM devices, we found XY4 the best balance

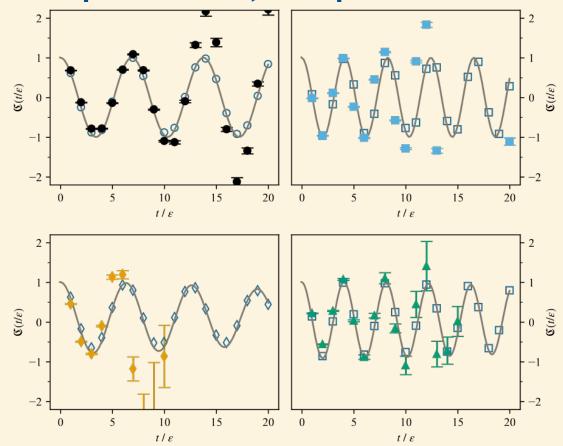
Putting it all together



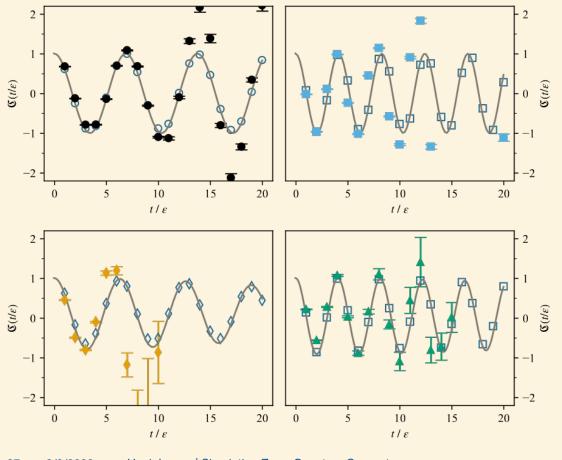
Error mitigation allows up to 6x longer evolution

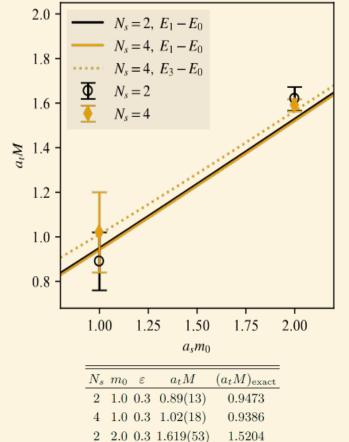


Multiple volumes, multiple masses



Multiple volumes, multiple masses



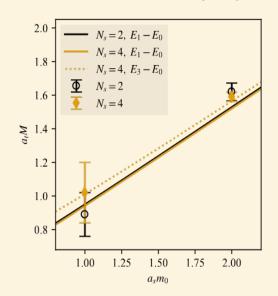


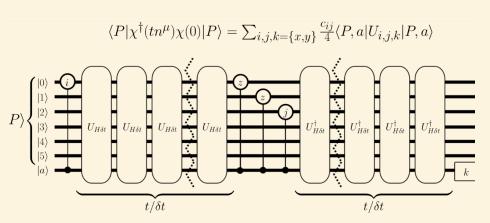
4 2.0 0.3 1.591(27)

1.5168

Endgame

- The road to quantum practicality in HEP will be long and winding
- Error mitigation provides small, but non-trivial extensions in evolution time
- Scale setting directly on the quantum computer is possible
- Current QCIPU project is computing hadronic tensor in this theory





Parton Physics on Quantum Computers
Lamm, Lawrence, Yamauchi - *Phys.Rev.Res* 2 (2020) 1, 013272
Formulation of Practical HEP Quantum Advantage Problem