

# Beyond Generalized Eigenvalues

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## Notation

$C_{ab}(t)$	scalar correlation function $\in \mathbb{C}^{1 \times 1}$
$t$	integer-spaced Euclidean time separations
$\mathbf{C}(t)$	matrix correlation function $\in \mathbb{C}^{N \times N}$
$N$	dimension of operator basis
$a, b, n$	operator indices $\in [1, N]$
$L$	exact polynomial order
$K$	assumed polynomial order
$\ell$	index $\in [0, L-1]$ or $[0, K-1]$
$k$	index of states contributing to $\mathbf{C}(t)$
$\lambda_k$	$\exp(-E_k)$ energy of state contributing to $\mathbf{C}(t)$
$z_{ak}$	$\langle 0   \mathcal{O}_a   k \rangle$ amplitude of state contributing to $\mathbf{C}(t)$

## Correlation Function

$$\begin{aligned} C_{ab}(t) &= \sum_k \langle 0 | \mathcal{O}_a | k \rangle \exp(-E_k t) \langle k | \mathcal{O}_b^\dagger | 0 \rangle \\ &= z_{ak} \lambda_k^t z_{bk}^* = \mathbf{Z} \Lambda^t \mathbf{Z}^\dagger \end{aligned}$$

We assume operator basis constructed so  $\mathbf{C}(t)$  is Hermitian

## Simple Effective Energy

$$\begin{aligned} C(t) &= [\lambda_{\text{eff}}(t, t_0)]^{t-t_0} C(t_0) \quad (t > t_0) \\ \Rightarrow [\lambda_{\text{eff}}(t, t_0)]^{t-t_0} - [C(t_0)]^{-1} C(t) &= 0 \\ \Rightarrow E_{\text{eff}}(t, t_0) &= -\frac{1}{t-t_0} \log \frac{C(t)}{C(t_0)} \end{aligned}$$

## Generalized Eigenvalues

Extends simple effective mass to matrix-valued correlation functions

$$\begin{aligned} \det \{ \mathbf{C}(t) - [\lambda_{\text{eff},k}(t, t_0)]^{t-t_0} \mathbf{C}(t_0) \} &= 0 \\ \Rightarrow \det \{ [\lambda_{\text{eff},n}(t, t_0)]^{t-t_0} \mathbf{I} - [\mathbf{C}(t_0)]^{-1} \mathbf{C}(t) \} &= 0 \end{aligned}$$

Reduces to finding eigenvalues provided  $\mathbf{C}(t_0)$  invertable.

Extending effective mass to matrices increases solution space  $1 \rightarrow N$ .

## Prony's Method (1795)

Linear Prediction

$$\begin{aligned} \mathcal{P}_K(\lambda) &= \prod_{k=1}^K (\lambda - \lambda_k) = \lambda^K + \sum_{\ell=0}^{K-1} p_\ell \lambda^\ell \\ \mathcal{P}_K(\lambda_k) &= 0 \\ \Rightarrow \lambda_k^K &= -\sum_{\ell=0}^{K-1} p_\ell \lambda_k^\ell \end{aligned}$$

Assuming only  $K$  states in  $\mathbf{C}(t)$

$$C(t) = -\sum_{\ell=0}^{K-1} p_\ell C(t-K+\ell), \quad t \geq K$$

## Prony's Method (1795) cont.

Solve for  $p_\ell$ .  $\mathbf{C}_t \equiv \mathbf{C}(t-t_0)$

$$\begin{bmatrix} C_K \\ C_{K+1} \\ \vdots \\ C_{2K-1} \end{bmatrix} = - \begin{bmatrix} C_0 & \cdots & C_{K-1} \\ C_1 & \cdots & C_K \\ \vdots & \ddots & \vdots \\ C_{K-1} & \cdots & C_{2K-2} \end{bmatrix} \begin{bmatrix} p_0 \\ p_1 \\ \vdots \\ p_{K-1} \end{bmatrix}$$

Find roots of polynomial  $\lambda_k$ . Alternatively, find eigenvalues of companion matrix.

$$\mathbf{C} = \begin{bmatrix} 0 & 0 & \cdots & 0 & -p_0 \\ 1 & 0 & \cdots & 0 & -p_1 \\ 0 & 1 & \ddots & \vdots & -p_2 \\ \vdots & \ddots & \ddots & 0 & \vdots \\ 0 & \cdots & 0 & 1 & -p_{K-1} \end{bmatrix}$$

After  $\lambda_k$  are known, amplitudes can be determined, just like simple case.

Extending effective mass from 2 times to  $2K$  times increased solution size  $1 \rightarrow K$ .

## Goal of This Work

Use  $N \times N$  correlation matrices at  $2 \cdot K$  equally-spaced times to find  $K \cdot N$  effective energy solutions.

## Block Prony Method (BPM)

Block Companion Matrix

$$\mathbf{C} = \begin{bmatrix} 0 & 0 & \cdots & 0 & -\mathbf{P}_0 \\ \mathbf{I} & 0 & \cdots & 0 & -\mathbf{P}_1 \\ 0 & \mathbf{I} & \ddots & \vdots & -\mathbf{P}_2 \\ \vdots & \ddots & \ddots & 0 & \vdots \\ 0 & \cdots & 0 & \mathbf{I} & -\mathbf{P}_{K-1} \end{bmatrix}$$

Assume  $K \cdot N$  non-zero, non-degenerate eigenvalues  $\lambda_k$ . Let  $\mathbf{v}_{k\ell}$  be size  $N$  blocks of  $k$ th eigenvector of  $\mathbf{C}$ , satisfying

$$\begin{aligned} \lambda_k \mathbf{v}_{k,0} &= -\mathbf{P}_0 \mathbf{v}_{k,K-1} \\ \lambda_k \mathbf{v}_{k,\ell} &= \mathbf{v}_{k,\ell-1} - \mathbf{P}_\ell \mathbf{v}_{k,K-1} \end{aligned}$$

After substitution

$$\lambda_k^K \mathbf{v}_{k,K-1} = \left( -\sum_{\ell=0}^{K-1} \lambda_k^\ell \mathbf{P}_\ell \right) \mathbf{v}_{k,K-1}$$

## Block Prony Method (BPM) cont

Let  $\mathbf{w}_\ell$  be size  $N$  blocks of any normalizable vector  $\mathbf{w}$  in eigenspace of  $\mathbf{C}$

$$\mathbf{C}_t \mathbf{w}_{K-1} = -\sum_{\ell=0}^{K-1} \mathbf{C}_{t-K+\ell} \mathbf{P}_\ell \mathbf{w}_{K-1}, \quad t \geq K$$

Solve for  $\mathbf{P}_\ell$ .  $\mathbf{C}_t \equiv \mathbf{C}(t-t_0)$

$$\begin{bmatrix} C_K \\ C_{K+1} \\ \vdots \\ C_{2K-1} \end{bmatrix} = - \begin{bmatrix} C_0 & \cdots & C_{K-1} \\ C_1 & \cdots & C_K \\ \vdots & \ddots & \vdots \\ C_{K-1} & \cdots & C_{2K-2} \end{bmatrix} \begin{bmatrix} \mathbf{P}_0 \\ \mathbf{P}_1 \\ \vdots \\ \mathbf{P}_{K-1} \end{bmatrix}$$

Then  $K \cdot N$  eigenvalues of block companion  $\mathbf{C}$  matrix are effective energies.

## Effective Amplitudes

Once  $\lambda_k$  are known, challenge is to solve bilinear system of equations (BLS)

$$C_{ab}(t) = \sum_{k=1}^{KN} z_{ak} \lambda_k^t z_{bk}^*$$

No general method to solve BLS is known. See [Johnson et al.](#)

Equal number of equations and unknowns.

$$C_{ab}(t) : 2K \frac{N(N+1)}{2}, \quad \lambda_k : KN, \quad z_{ak} : N(KN)$$

Exact solution should be possible up to signs:  $z_{ak} \rightarrow -z_{ak}$  for any  $k$ .

For now, we recommend searching for zero residual solutions using non-linear least squares minimization.

## References for Further Reading

[Fleming, Cohen, Lin and Pereyra, arXiv:0903.2314 \[hep-lat\]](#)

[Gohberg, Lancaster and Rodman, Matrix Polynomials, SIAM 2009.](#)

[Johnson, Šmigoc and Yang, arXiv:1303.4988 \[math.RA\]](#)

[Cushman and Fleming, arXiv:1912.08205 \[hep-lat\]](#)