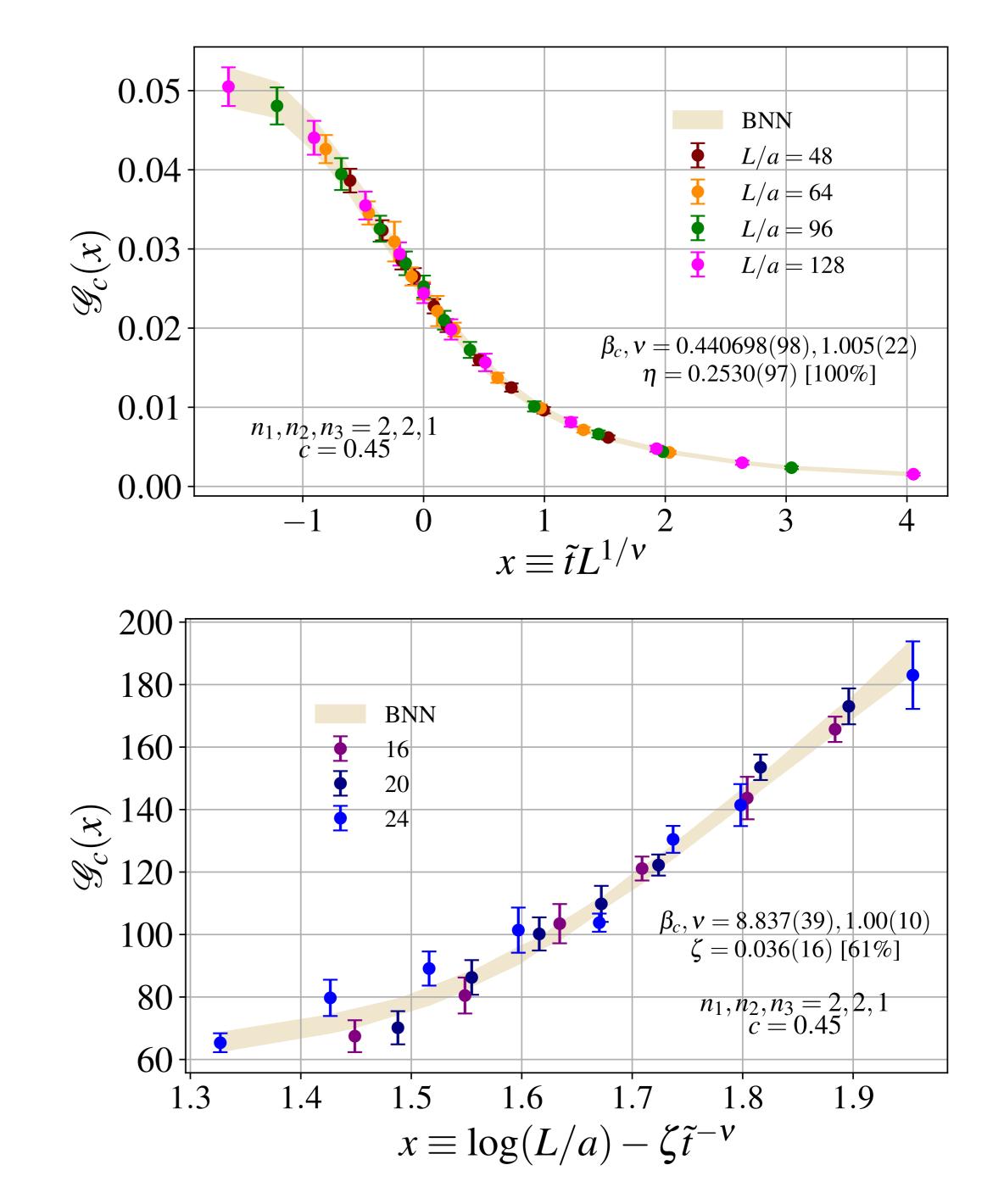
# **Constrained curve fitting with Bayesian neural networks Curtis Taylor Peterson & Anna Hasenfratz** University of Colorado Boulder

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#### Introduction

Common to many analysis pipelines in lattice field theory is the need to fit data to a model that is determined only partially by a finite number of model parameters. Familiar examples include analyses of finite-size scaling (FSS) and ground state spectroscopy. Motivated by promising results in condensed matter physics [1], we conjecture that the expressivity of neural networks makes them good candidates for parameterizing the unknown component of such models. We test this conjecture by performing a curve collapse analysis of the 2nd-order finite-temperature phase transition of the 2D Ising model and the zero-temperature phase transition of the 4D massless  $N_f/Nc = 8/3$  gauge-fermion system.

#### Results





#### Network architecture & training

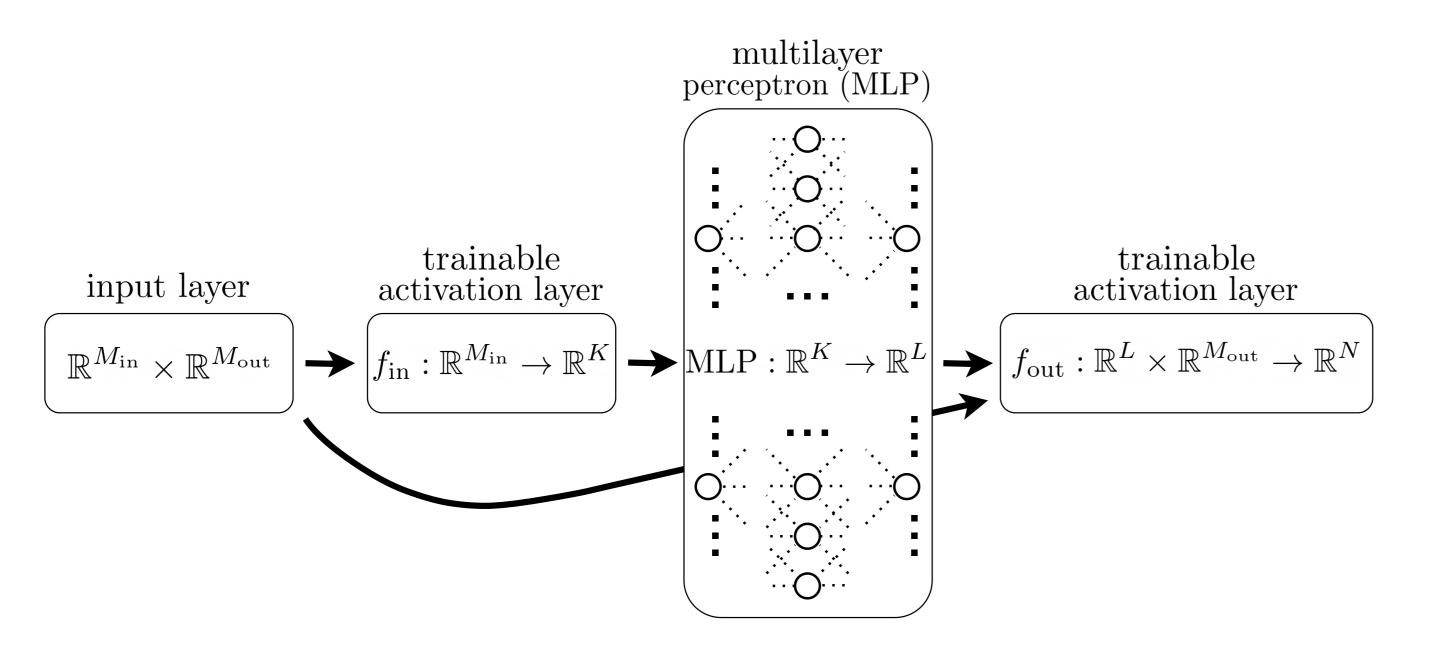


Figure 1: Neural network architecture.

#### Architecture

- Full model  $\mathcal{M}_{\Theta}$  maps input features in  $\mathbb{R}^{M_{\text{in}}} \times \mathbb{R}^{M_{\text{out}}}$  to output features in  $\mathbb{R}^{N}$  given the model parameters  $\{\Theta\} = \{\theta_{\text{in}}, \theta_{\text{MLP}}, \theta_{\text{out}}\}$ ; see Fig. 1.
- $-\{\theta_{MLP}\}\$ denote learnable parameters of a multilayer perception (MLP) that pa-

Figure 2: (Top panel) Curve collapse for the 2D Ising model using a 3-layer network with  $n_1 = n_2 = 2$  nodes in the first and second layers and  $n_3 = 1$  nodes in the third (output) layer. (Bottom panel) Curve collapse for the 4D massless  $N_f/N_c = 8/3$  gauge-fermion system with  $n_1, n_2, n_3 = 2, 2, 1$  nodes in each layer.

- rameterize the a priori unknown component of the model, such as the scaling function in an FSS analysis or the infinite sum of excited-state exponentials in an analysis of hadron spectroscopy.
- $-\{\theta_{in}\}\$  and  $\{\theta_{out}\}\$  denote additional learnable parameters of the model, such as critical exponents or ground-state energies/amplitudes, that enter the model function through explicitly-defined activation functions that act upon a combination of input features of the full network and output features of the MLP.

### Training

• Training is performed by least-squares minimization of the augmented  $\chi^2$  [2]

 $\chi^2_{\mathrm{aug.}} = \chi^2_{\mathrm{data}} + \chi^2_{\mathrm{priors}},$ 

where  $\chi^2_{data}$  and  $\chi^2_{priors}$  include information from the covariance of the data & priors, respectively.

- Priors on the networks weights in  $\{\theta_{MLP}\}$  provide a lever arm for controlling overfitting; equivalent to L2 regularization in the ML literature
- Appropriate priors can be estimated using the *empirical Bayes method*
- We optimize  $\chi^2_{aug.}$  using the "basin hopping" algorithm championed in Ref. [3]:

1. Random perturbation of coordinates  $(\Theta)$ 

2. Local minimization (trust region reflective)

3. Metropolis accept/reject

## **Conclusions & forthcoming research**

- Our FSS results are consistent with the published literature
- -Prediction for Ising critical exponents consistent with exact prediction:  $\beta_c \approx 0.4407$  and  $\nu, \eta = 1, 1/4$ .
- $-N_f = 8$  test rules out 1st-order scaling and appears to have a very slight preference for BKT.
- Even a relatively small neural network is sufficiently expressive to accurately represent the scaling function for the systems that we have tested
- -Bayesian priors provide level arm to control overfitting
- We are currently testing this method on the XY model
- •We will begin testing the application of this method to ground state spectroscopy

#### References

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#### **Gradient flow coupling & FSS**

The gradient flow coupling at smearing radius  $\sqrt{8t} = cL$  defined by

 $g_c^2(\beta,L) \sim t^2 \langle E(t) \rangle \big|_{8t=(cL)^2}$ 

can be utilized as a scaling variable for use in finite-size scaling [4]. • At a 1st/2nd-order phase transition,  $g_c^2(\beta, L)$  is expected to scale as

 $g_c^2(eta,L) \sim L^\eta \mathcal{G}_c( ilde{t}L^{1/
u})$ 

with  $\nu = 1/d$  signaling first order transition. • At a Berezinskii–Kosterlitz–Thouless (BKT) transition the scaling is  $g_c^2(\beta, L) \sim L^{\gamma/\nu} \mathcal{G}_c(L \exp(-\zeta \tilde{t}^{-\nu})),$ 

where  $\tilde{t} \equiv \beta / \beta_c - 1$ .

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