

# Neutrino Detectors

*(Part I)*

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# About Myself

- I am an Experimental Particle Physicist
- Ph.D. (Wayne State University)
  - CMS experiment at CERN, LHC
- Postdoc (Kansas State University)
  - *Switched to neutrinos as a postdoc — best decision of my life!*
  - *Spent most of my postdoctoral career on MicroBooNE*
- Faculty (University of Tennessee, Knoxville, 2016-19)
- Senior Scientist at Los Alamos National Lab since 2019
- Experimental collaborations
  - Short and long-baseline oscillation experiments (MicroBooNE, SBND, DUNE)
  - Neutron cross section experiments (ARTIE, MarEX)
- Involvement
  - *Oscillations, argon cross sections, detector physics & calibration analyses*
  - *Detector R&D, calibration & cryogenic instrumentation, slow controls*

*My first trip to Fermilab  
as a graduate student  
(2009)*



*Constructing MicroBooNE TPC*



# Neutrino Detector Lectures

## We have two 1-hour lectures

- Lecture 1: Neutrino Detection Basics
- Lecture 2: Detector Technologies

## Disclaimer:

- There is no way I can cover everything...plus I have my own biases
- My goal with these lectures is to cover the general principles of neutrino detection, their interactions and experimental signatures along with an overview of technologies.
- Biggest challenge in preparing these lectures: *getting the CaPiTaLiZatioN right!*



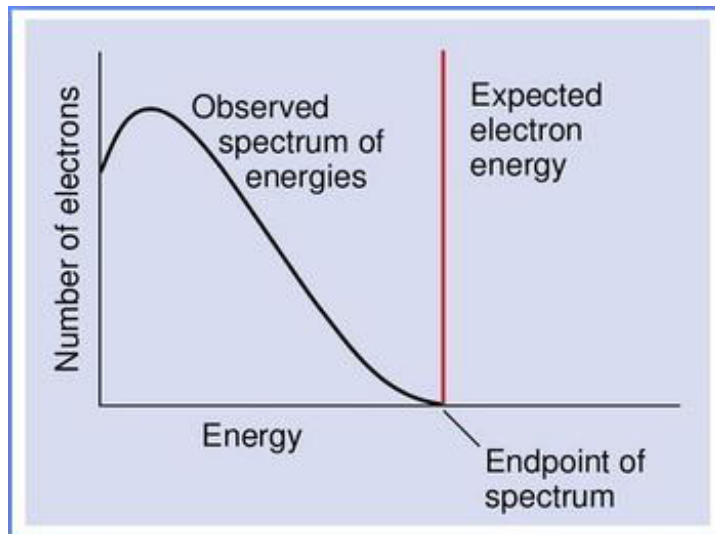
# Lecture 1: Outline

- Lecture 1 (Neutrino Detection Basics)
  - Brief Intro to Neutrinos
  - Neutrino Sources
  - Neutrino Detection Challenges
  - Neutrino Detector Goals
  - Neutrino Interactions
  - Passage of Particles in Matter

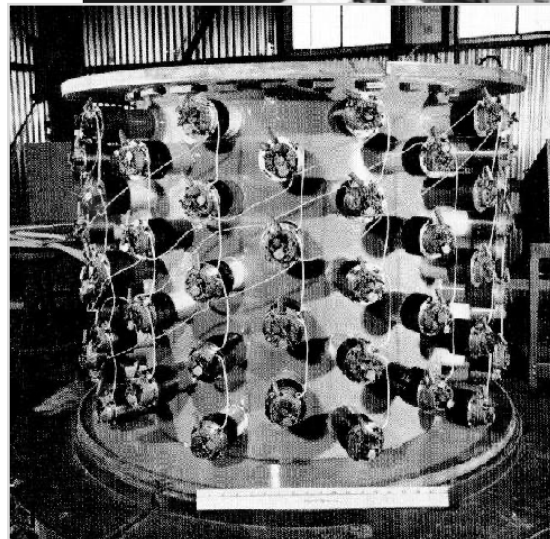
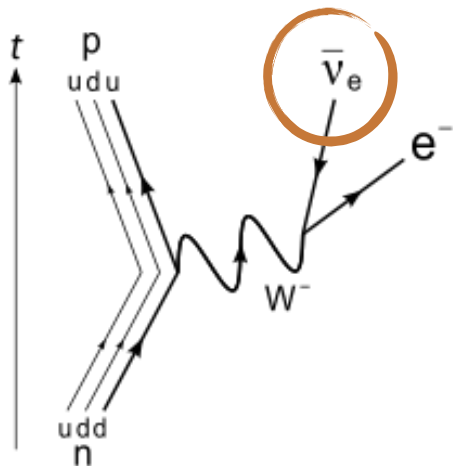


# Project Poltergeist

A desperate remedy by Pauli to explain the observed energies in  $\beta$  decay (1931)



Nuclear  $\beta$  decay



1956: F. Reines and C. Cowen from LANL detected the first neutrino from a nuclear reactor (*Nobel Prize 1995*)



June 5, 1998

## Mass Found in Elusive Particle; Universe May Never Be the Same

### Discovery on Neutrino Rattles Basic Theory About All Matter

By MALCOLM W. BROWNE

TAKAYAMA, Japan, June 5 — In what colleagues hailed as a historic landmark, 120 physicists from 23 research institutions in Japan and the United States announced today that they had found the existence of mass in a notoriously elusive subatomic particle called the neutrino.

The neutrino, a particle that carries no electric charge, is so light that it was assumed for many years to have no mass at all. After today's announcement, cosmologists will have to confront the possibility that much of the mass of the universe is in the form of neutrinos. The discovery will also compel scientists to revise a highly successful theory of the composition of matter known as the Standard Model.

Word of the discovery had drawn some 300 physicists here to discuss neutrino research. Among other things, they said, the finding of neutrino mass might affect theories about the formation and evolution of galaxies and the ultimate fate of the universe. If neutrinos have sufficient mass, their presence throughout the universe would increase the overall mass of the universe, possibly slowing its present expansion.

Others said the newly detected but as yet unmeasured mass of the neutrino must be too small to cause cosmological effects. But whatever the case, there was general agreement here that the discovery will have far-reaching consequences for the investigation of the nature of matter.

Speaking for the collaboration of scientists who discovered the existence of neutrino mass using a huge underground detector called Super-Kamiokande, Dr. Takaaki Kajita of the Institute for Cosmic Ray Research of Tokyo University said that all explanations for the data collect-

#### Detecting Neutrinos



Neutrinos pass through the Earth's surface to a tank filled with 12.5 million gallons of ultra-pure water ...

... and collide with other particles ...

... producing a cone-shaped flash of light.



LIGHT AMPLIFIER

The light is recorded by 11,200 20-inch light amplifiers that cover the inside of the tank.

#### And Detecting Their Mass

By analyzing physicists' neutrinos have their journey form, they m

Source: Univers

ed by the de  
ence of neutr  
sentially rule

Dr. Yoji T  
coalition and  
mioka Neutri  
the underground  
30 miles north  
Alps, acknowle

announcement was "very strong," but said, "We have investigated all

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# Neutrinos Oscillate and so they have mass! (albeit very tiny)

Until as recently as 1998, neutrinos were considered to be *massless*

We have detected  
oscillations from

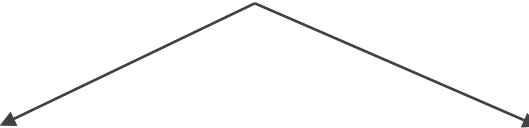
Atmospheric  
Solar  
Accelerator  
Reactor





# There is still a lot we don't know

SM = Standard Model



## Within SM 3-flavor mixing

- Absolute mass of neutrinos?
- Neutrinos Majorana or Dirac?  
**Are neutrinos their own anti-particles?**
- Precision Measurement of neutrino mixing parameters?
- Which neutrino is the lightest?
- CP violation in the neutrino sector? **Matter-antimatter puzzle**

## Beyond SM 3-flavor mixing

- Absolute mass of neutrinos?  
Are there more than 3 neutrinos? **Candidates for dark matter?**
- Other New physics e.g. non-standard interactions

We have been studying neutrinos for about a century now and there is still a lot we don't know. The story of neutrinos is far from being complete

# There is still a lot we don't know

Direct Mass  
Measurement  
Experiments

Within SM 3-flavor mixing

- Absolute mass of neutrinos?
- Neutrinos Majorana or Dirac?
- Precision Measurement of mixing parameters?
- Neutrino mass ordering?
- Is  $\theta_{23}$  maximal mixing?
- CP violation in the neutrino sector?

Long-Baseline  
Neutrino  
Oscillation  
Experiments

Neutrinoless  
Double Beta  
Decay Expts.

Beyond SM 3-flavor mixing

- Absolute mass of neutrinos? Are there more than 3 neutrinos?
- Other BSM physics e.g. non-standard interactions

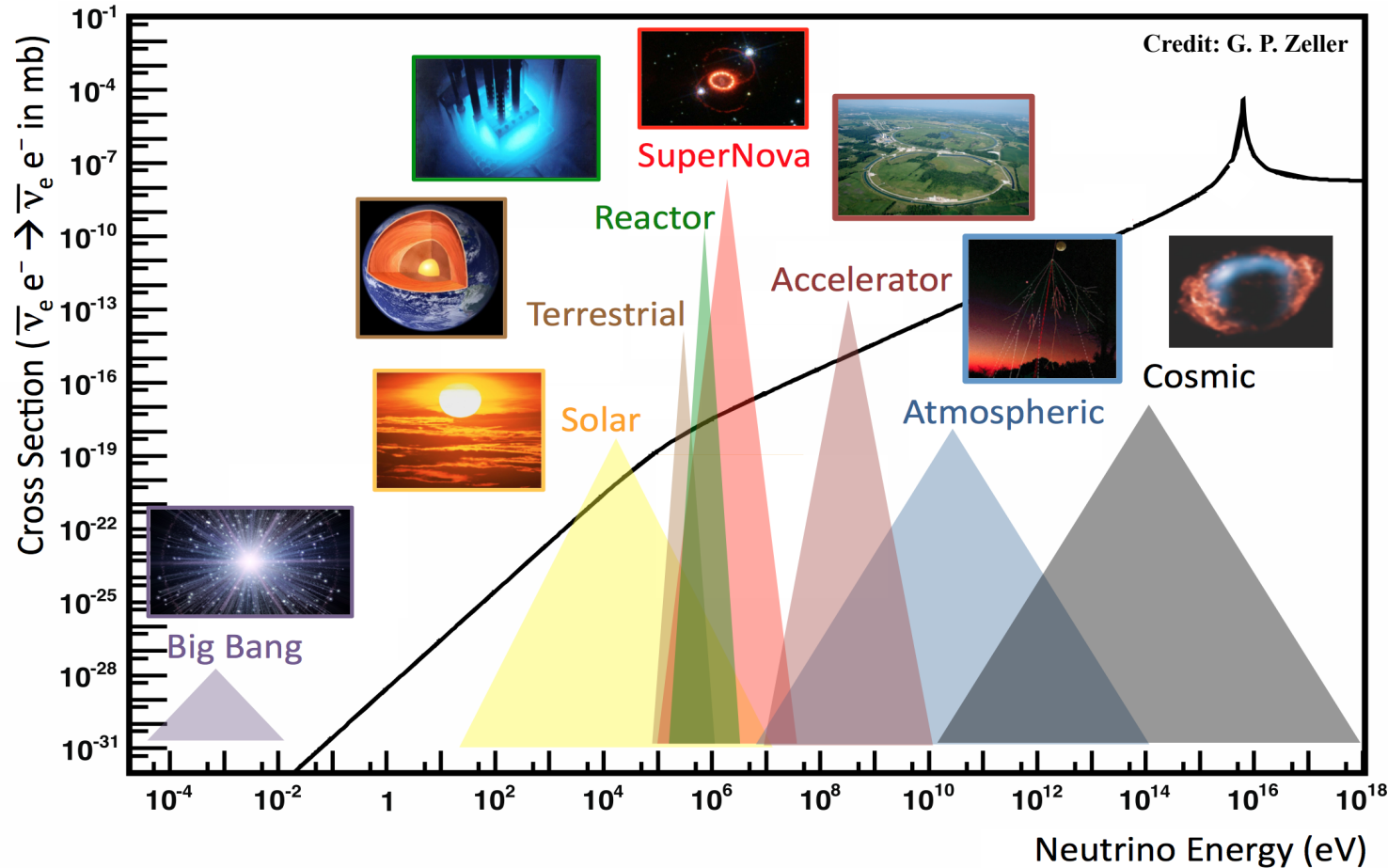
Short-Baseline  
Neutrino  
Experiments

Short- and  
Long-Baseline  
Experiments

Neutrino Physics is a  
very active field and will be for  
the next few decades!



# Neutrino Sources



Neutrinos span  
multiple frontiers!

Particle Physics

AstroPhysics

Cosmology

High energy Astro-  
particle physics

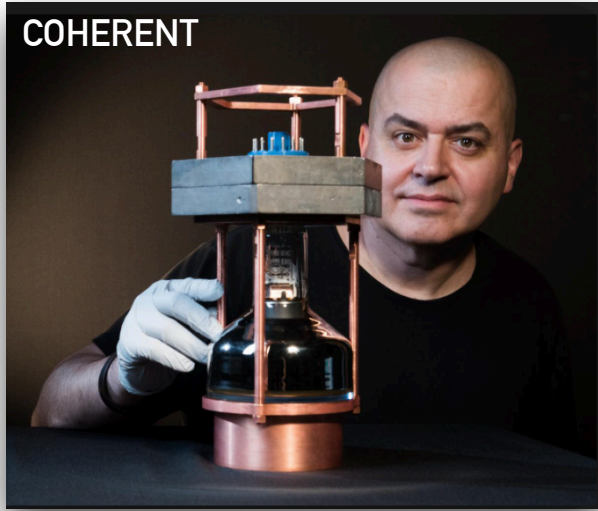
Nuclear physics

- Overwhelming number of sources, wide range of energies
- Need wide spectrum of experiments and technologies!

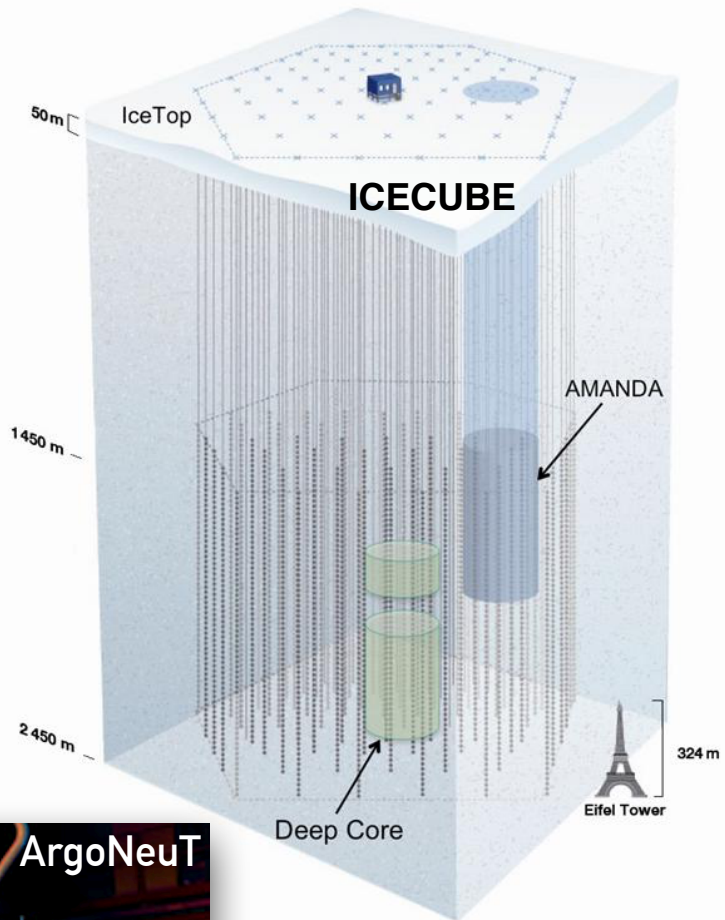
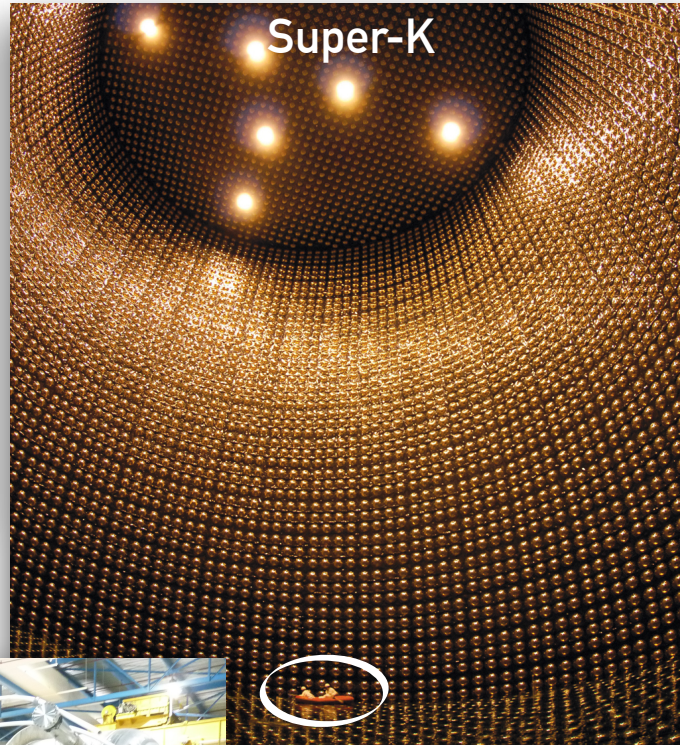


# Neutrino Detectors at All Scales

COHERENT



Super-K



MicroBooNE



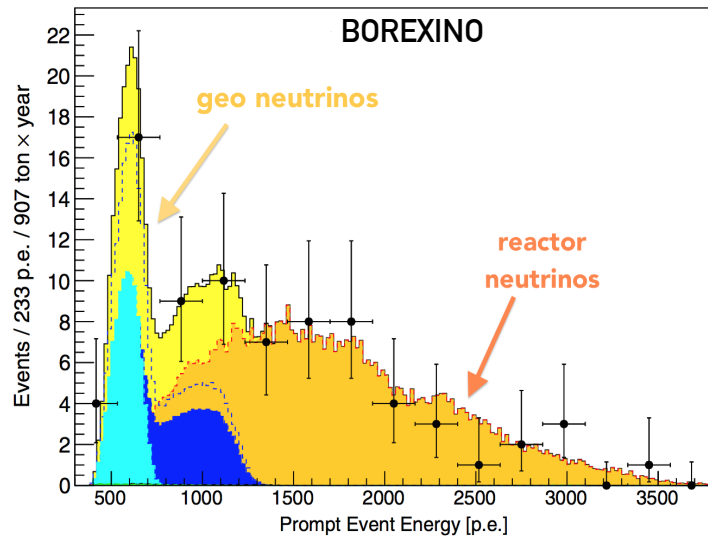
ArgoNeuT



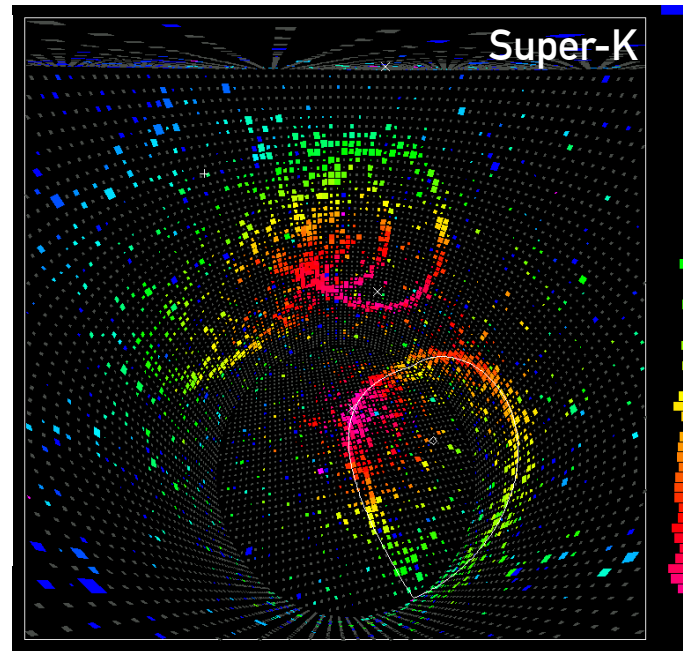
**Technologies  
advances at every  
turn have enabled and  
continue to enable  
neutrino discoveries**



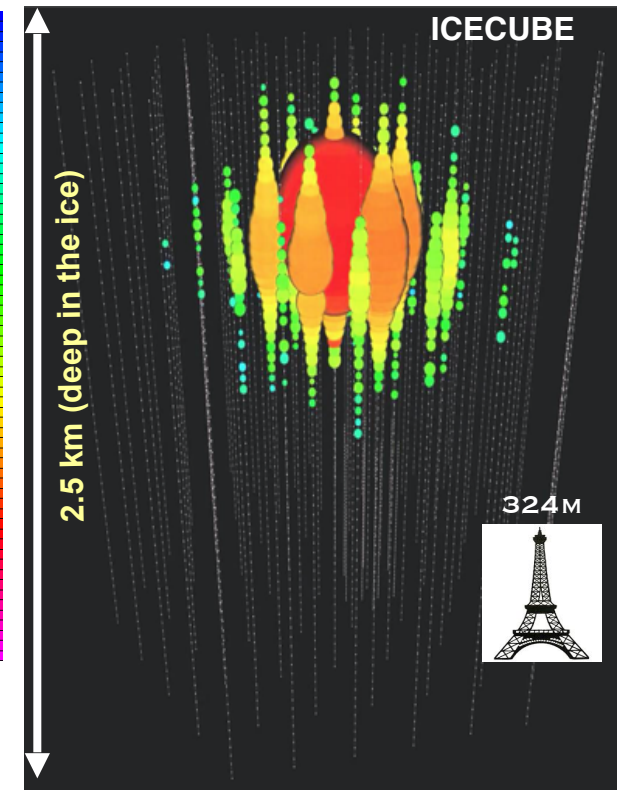
# Visualizing Neutrinos



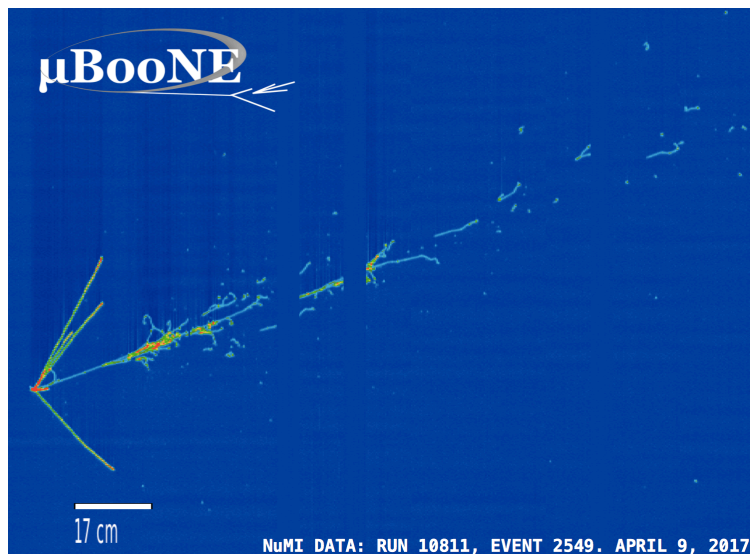
MeV-scale neutrino



A few-100 MeV neutrino



A 2 PeV scale astro physical event in the detector



MeV ————— We have observed neutrinos at wide range of energies —————> PeV

# Detecting Neutrinos is Challenging

- They are invisible (no charge)
- They are extremely weakly interacting

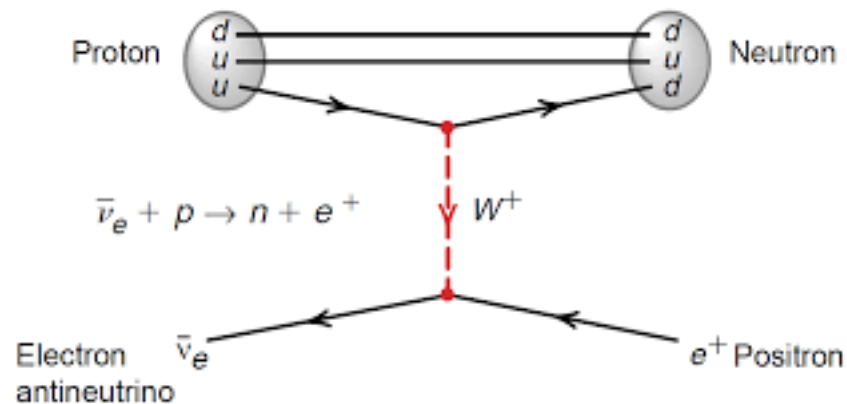


# Detecting Neutrinos is Challenging

- They are invisible (no charge)
- They are *extremely weakly* interacting
- In other words, they have very small interaction cross sections
- MeV-scale neutrino (typical energy of a neutrino emitted from sun or a nuclear reactor) has a cross section,  $\sigma \sim 10^{-44} \text{ cm}^2$  — tiny!

Inverse Beta Decay

$$\bar{\nu}_e + p \rightarrow e^+ + n$$



$$\sigma_{\bar{\nu}p} \approx 5 \times 10^{-44} \text{ cm}^2 \quad \text{for} \quad (E_{\bar{\nu}} \sim 2 \text{ MeV})$$

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- GeV-scale neutrino (typical energy of a neutrino from a particle accelerator) has a cross section,  $\sigma \sim 10^{-40} \text{ cm}^2$  — still tiny!
- Mean free path of a neutrino in lead
  - MeV-scale neutrino:  $d_{\text{lead}} \sim 10^{16} \text{ m}$  (over a light year of lead!)
  - GeV-scale neutrino:  $d_{\text{lead}} \sim 10^{12} \text{ m}$  (still almost a trillion miles of lead!)
- What about a GeV-scale proton?  $\sigma \sim 10^{-25} \text{ cm}^2$ 
  - GeV-scale proton:  $d_{\text{lead}} \sim 10 \text{ cm!}$

$$d_{\text{lead}} = \frac{1.66 \times 10^{-27} \text{ kg}}{(\sigma_{\text{v-N}} \text{ m}^2)(11400 \text{ kg/m}^3)}$$

atomic mass unit

v-N cross-section

density of lead

# Detecting Neutrinos 101

- *Basic Strategy*
  - Produce them in large quantities in a well defined area
  - Put something **very Dense**, **very BIG** and **very Sensitive** for neutrinos to interact
- *In other words*
  - High intense beams (typically kW beams, now moving to MW)
  - Large neutrino fluxes
  - Long exposure time
  - Dense targets (e.g. Argon)
  - Large target mass (tens of meters, hundreds to multi-kiloton-scale)
  - Low background (place them underground; design for maximum signal sensitivity; efficient background tagging etc.)

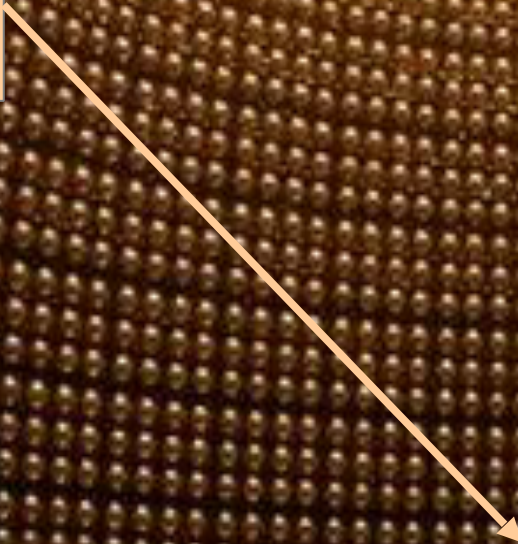


# The Super Kamiokande Experiment (Japan)

Dimensions:  
41 m height  
30m diameter tank  
50,000 tons of water  
11,000 PMTs

To study solar and atmospheric neutrinos  
**(1000 m underground)**  
A water Cherenkov detector

Researchers sitting  
in a boat  
inside the detector  
*How cool is that?*

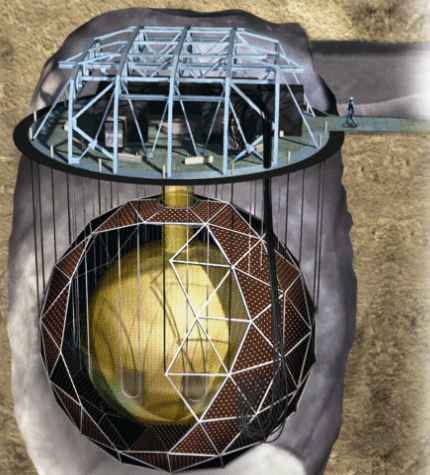
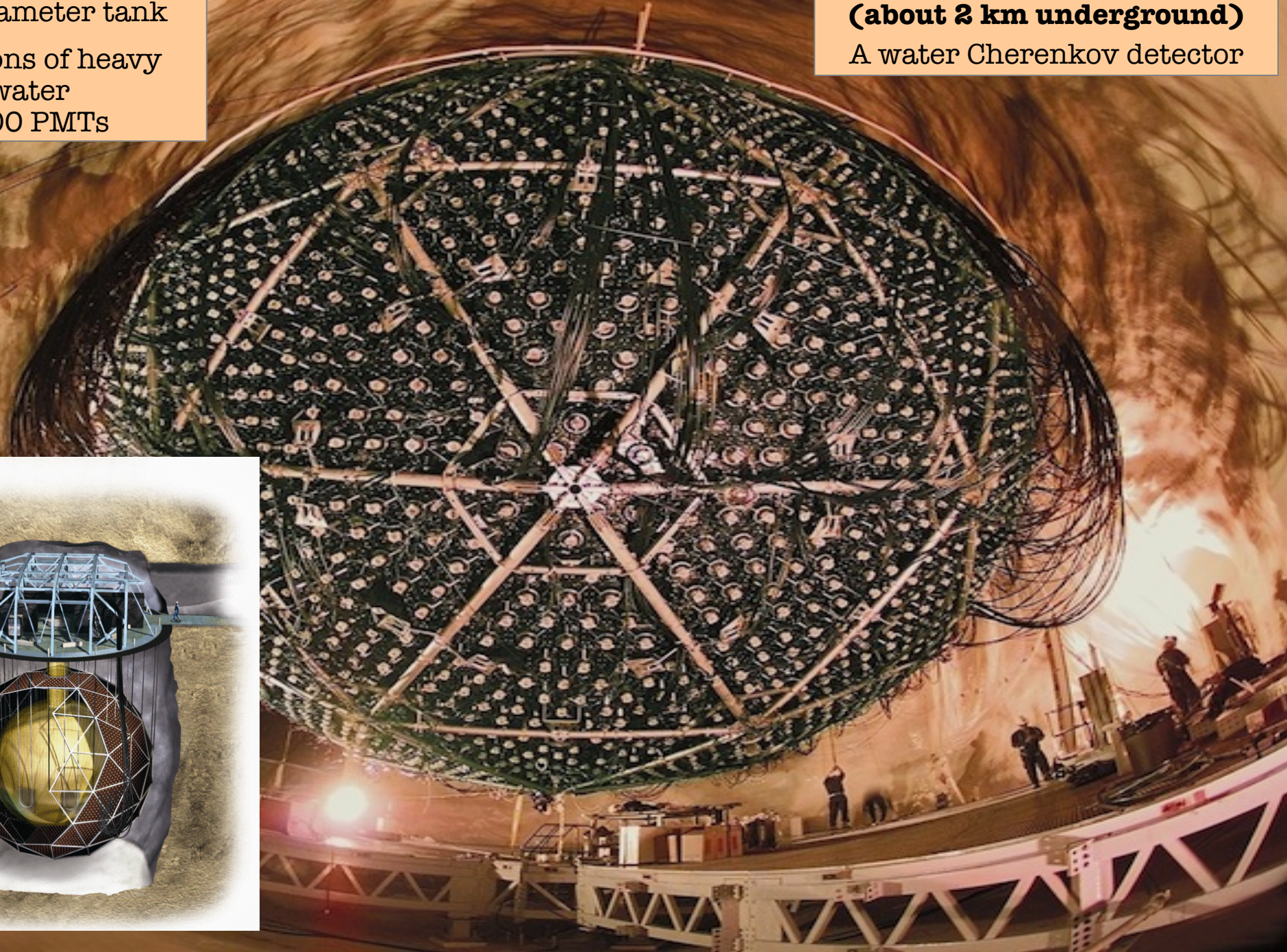




# The Sudbury Neutrino Observatory (Canada)

Dimensions:  
12m diameter tank  
1000 tons of heavy  
water  
9000 PMTs

To study solar neutrinos  
**(about 2 km underground)**  
A water Cherenkov detector





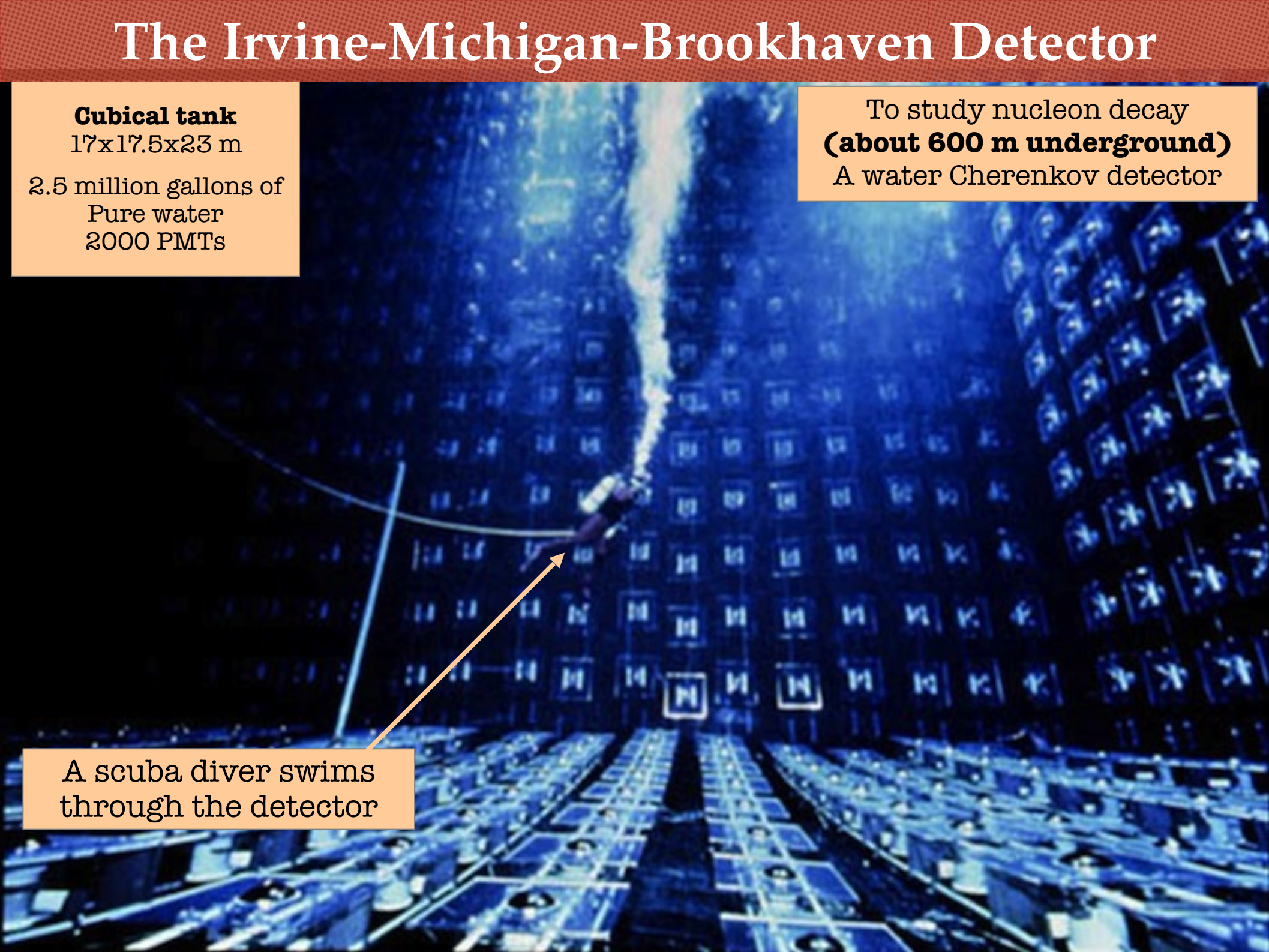
# The Irvine-Michigan-Brookhaven Detector

## Cubical tank

17x17.5x23 m

2.5 million gallons of  
Pure water  
2000 PMTs

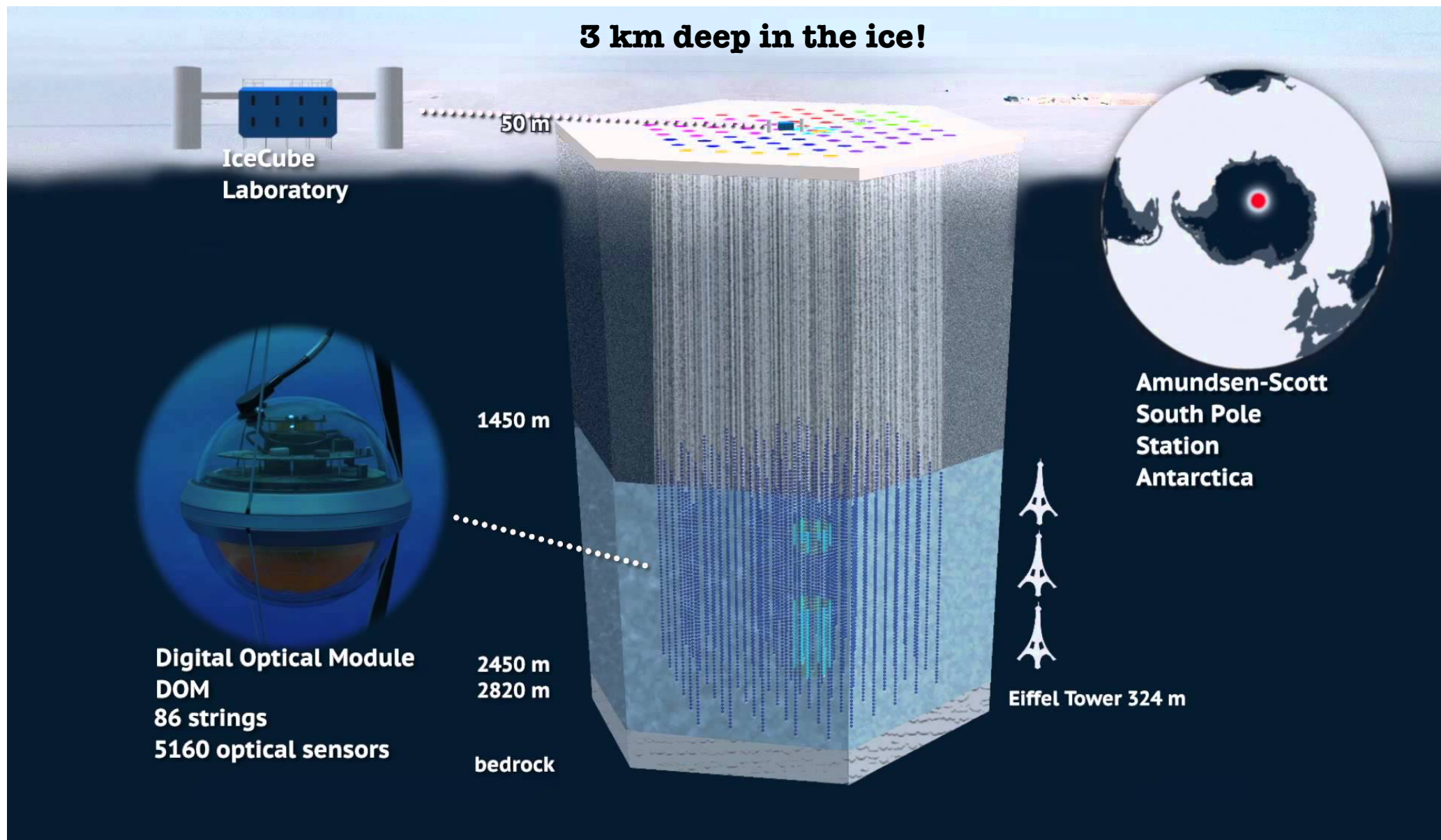
To study nucleon decay  
**(about 600 m underground)**  
A water Cherenkov detector

A photograph of the interior of the Irvine-Michigan-Brookhaven Detector. The image shows a vast, dark blue space filled with a grid of photomultiplier tubes (PMTs) arranged in a cubic pattern. A scuba diver is visible in the center, swimming through the water. A bright, vertical streak of light is visible in the background. An orange arrow points from the text box to the diver.

A scuba diver swims  
through the detector



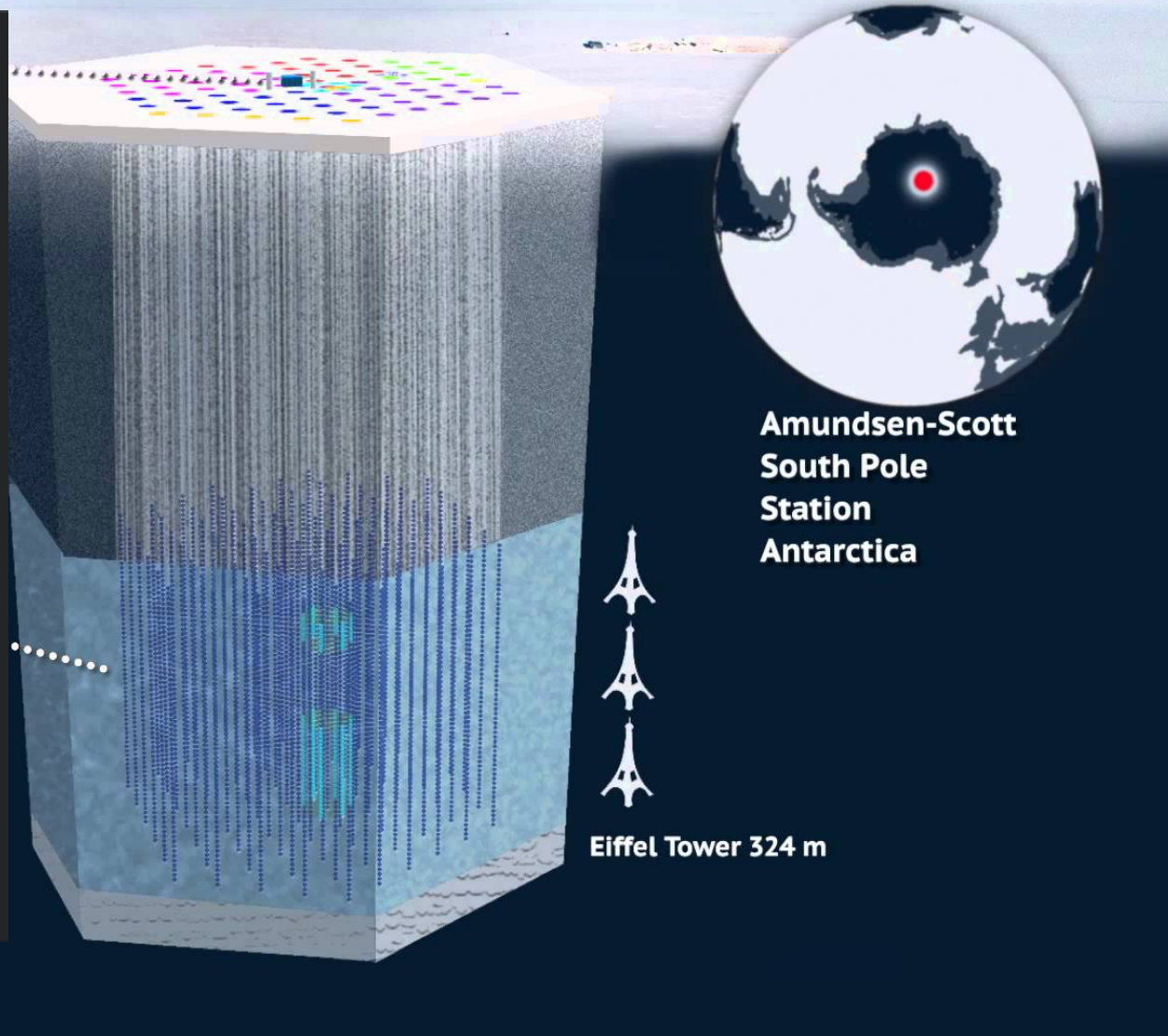
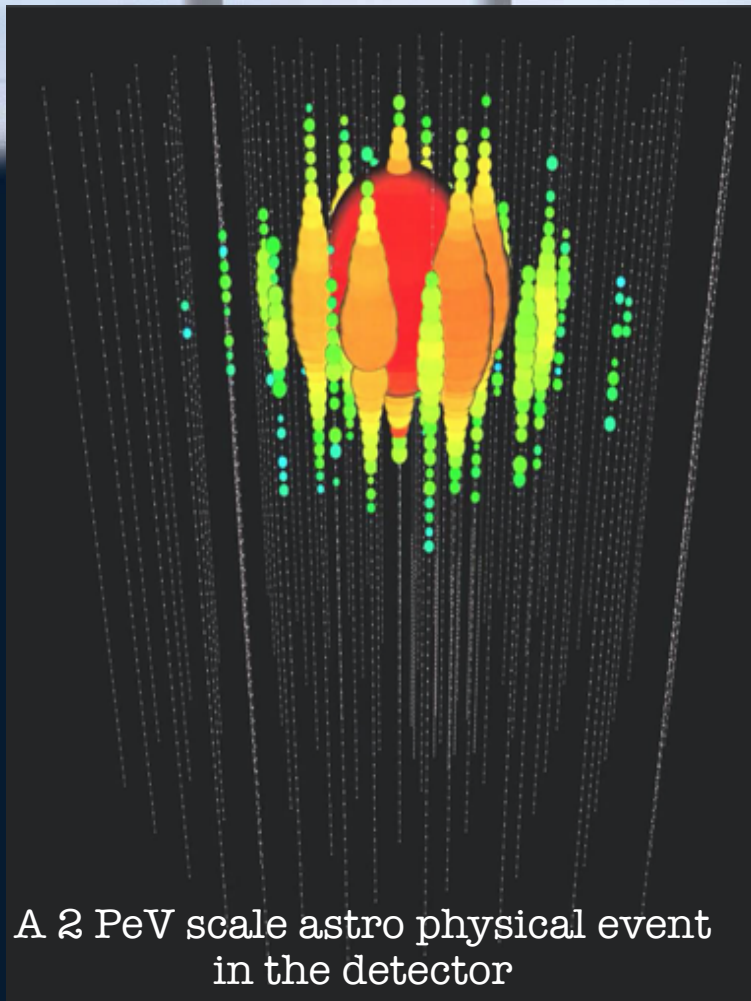
# The IceCUBE Experiment (South Pole)



searches for neutrinos from the most violent astrophysical sources: exploding stars, gamma-ray bursts, black holes and neutron stars.

# The IceCUBE Experiment (South Pole)

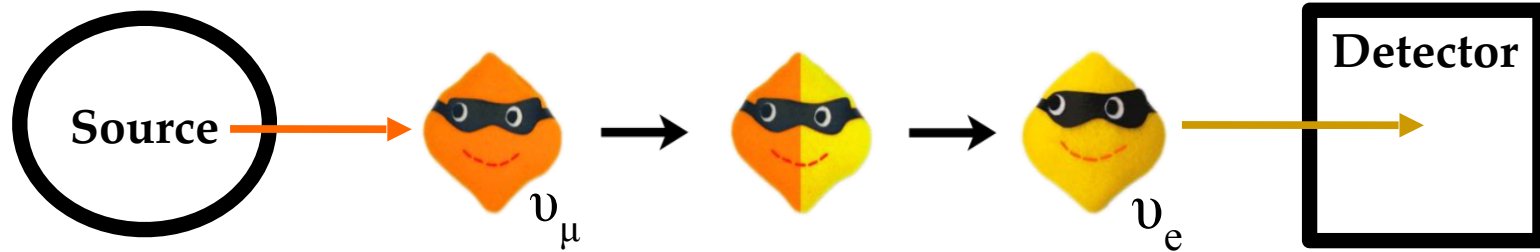
**3 km deep in the ice!**



searches for neutrinos from the most violent astrophysical sources: exploding stars, gamma-ray bursts, black holes and neutron stars.



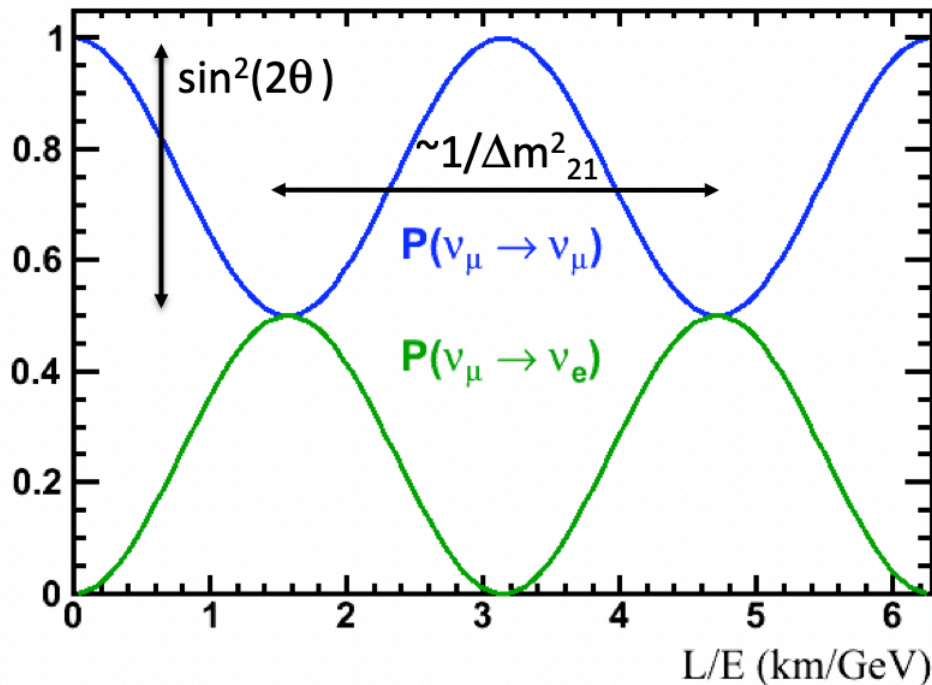
# 2-Flavor Oscillations



Oscillation  
Probability

$$P(\nu_\alpha \rightarrow \nu_\beta) = \sin^2(2\theta) \sin^2\left(1.27 \Delta m^2 \frac{L(\text{km})}{E(\text{GeV})}\right)$$

$$\Delta m_{21}^2 = m_2^2 - m_1^2$$



Experimental parameters:  $L, E$

Parameter of nature:  $\Delta m^2, \sin^2 2\theta$

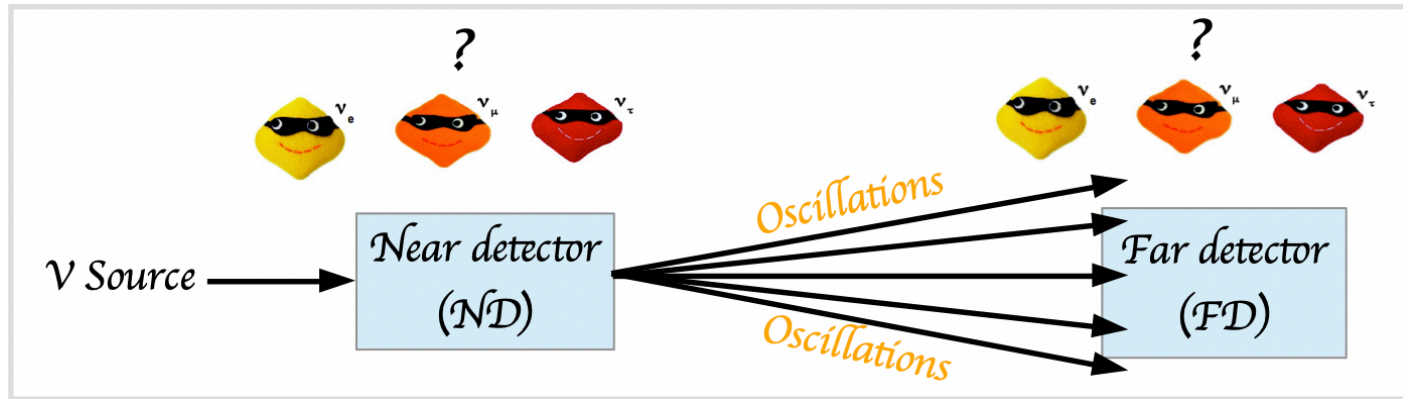
**Long-baseline:**  $L \sim 1000 \text{ km}$

**Short-baseline:**  $L \sim 1 \text{ km}$

# A Typical Oscillation Experiment

## A Typical Oscillation Experiment

*Oscillation experiments are basically counting experiments*



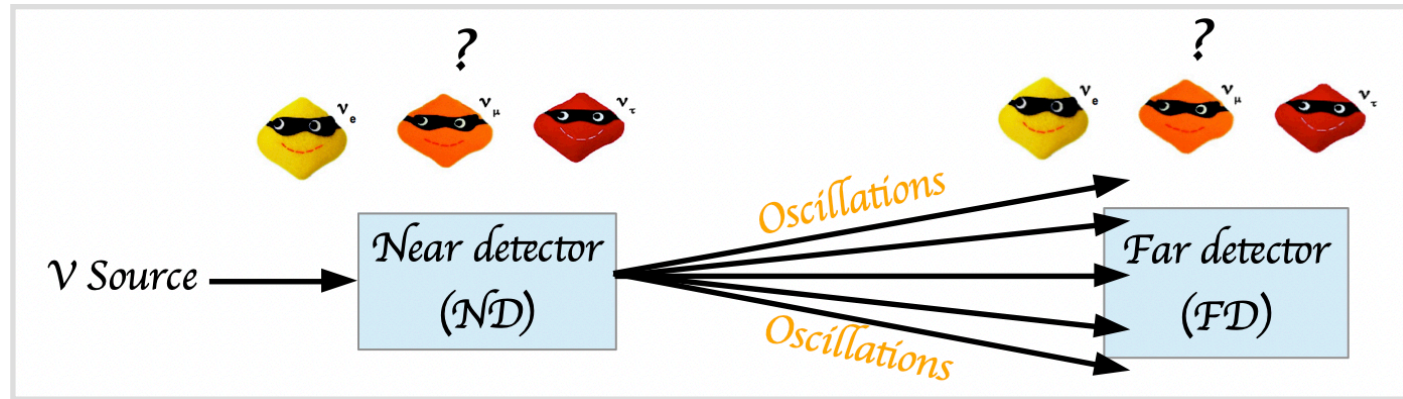
- Start with an intense source of neutrinos (e.g.  $\nu_\mu$ )
- Build a near detector and a far detector with distance optimized for oscillations to occur
- Measure unoscillated flavor and energy spectrum at  $L \sim 0$  (near detector)
- Measure oscillated flavor and energy spectrum again at  $L \sim \text{oscillation maximum}$  (far detector)
- Compare



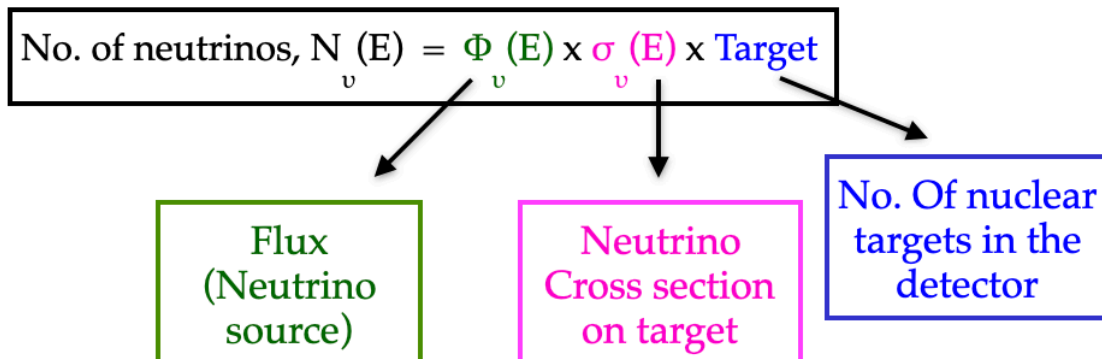
# A Typical Oscillation Experiment

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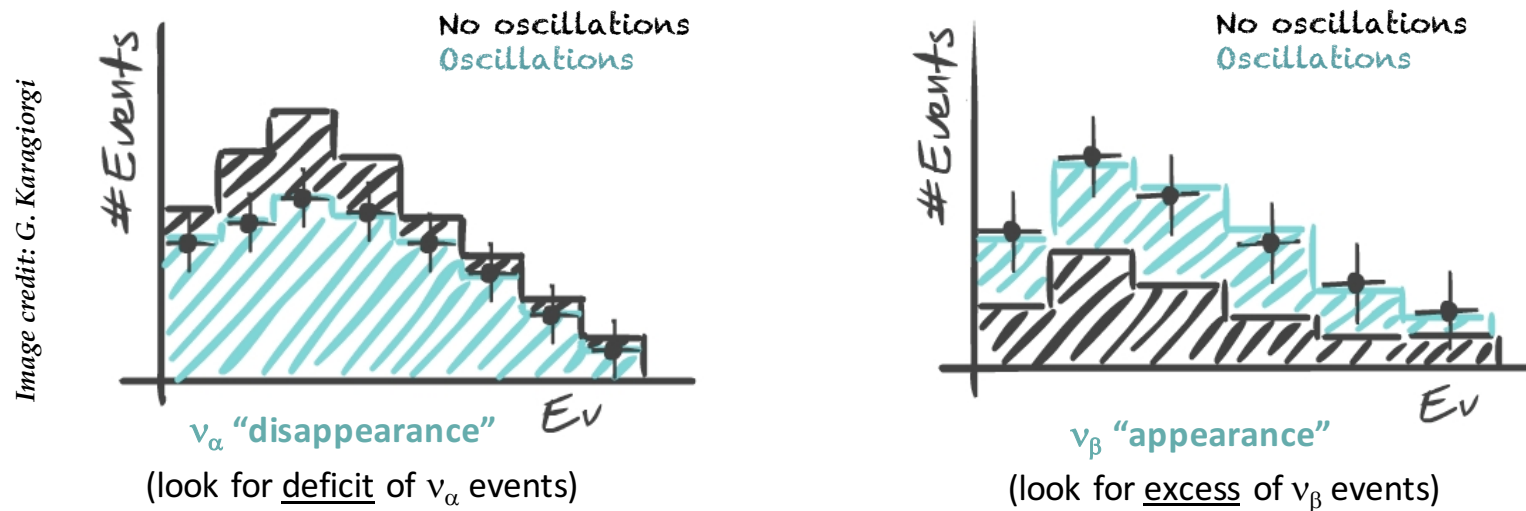
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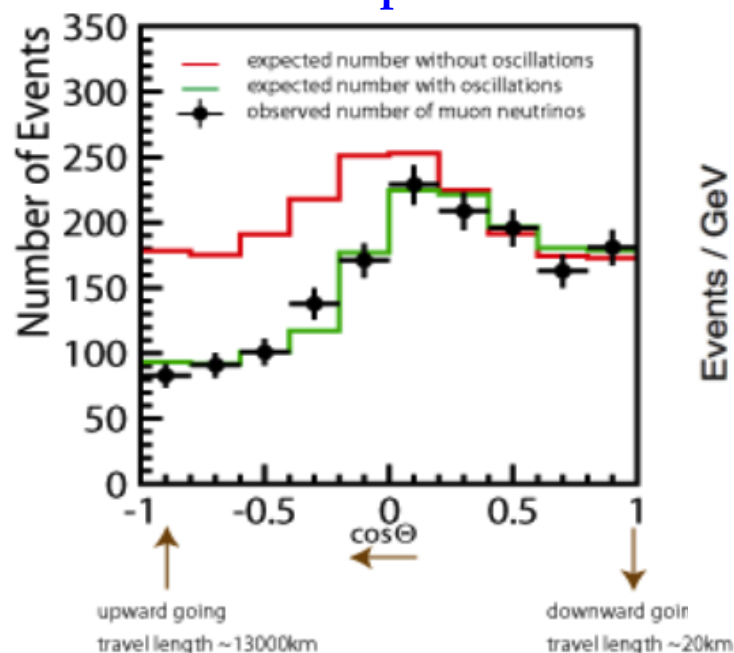


Can perform "Appearance" or "Disappearance" measurements

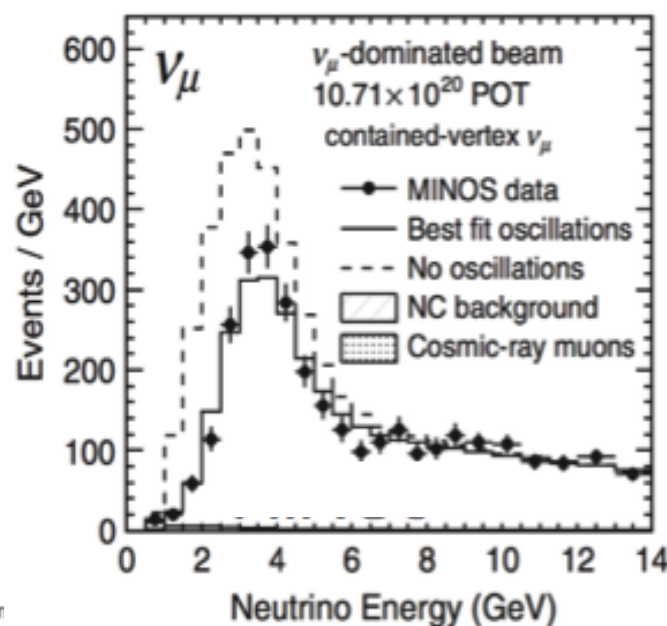


# Oscillation Experiments

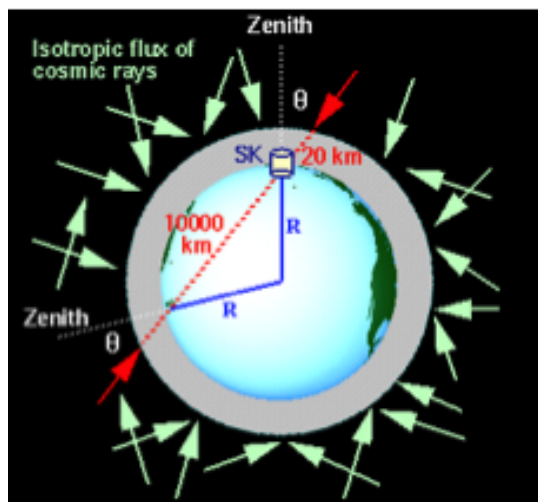
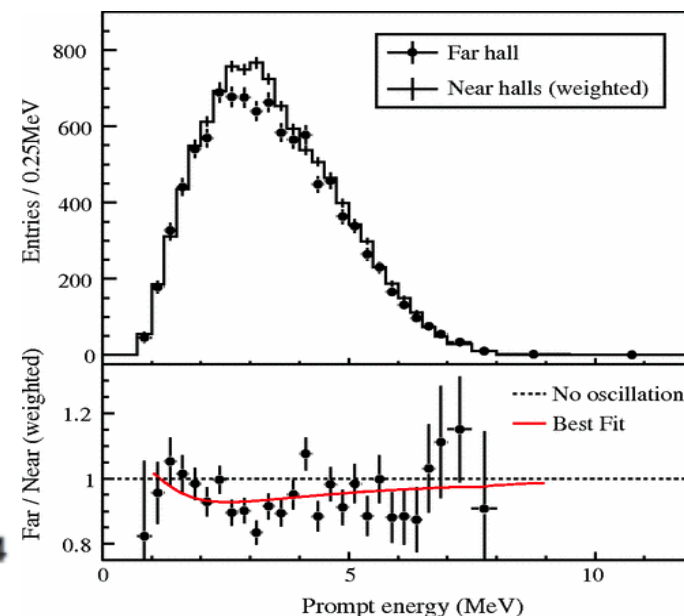
**SUPER-K**  
**Atmospheric**



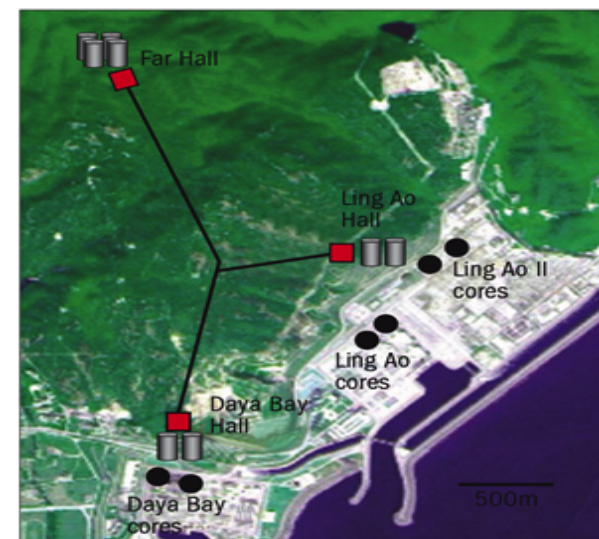
**MINOS**  
**Accelerator**



**DAYABAY**  
**Reactor source**



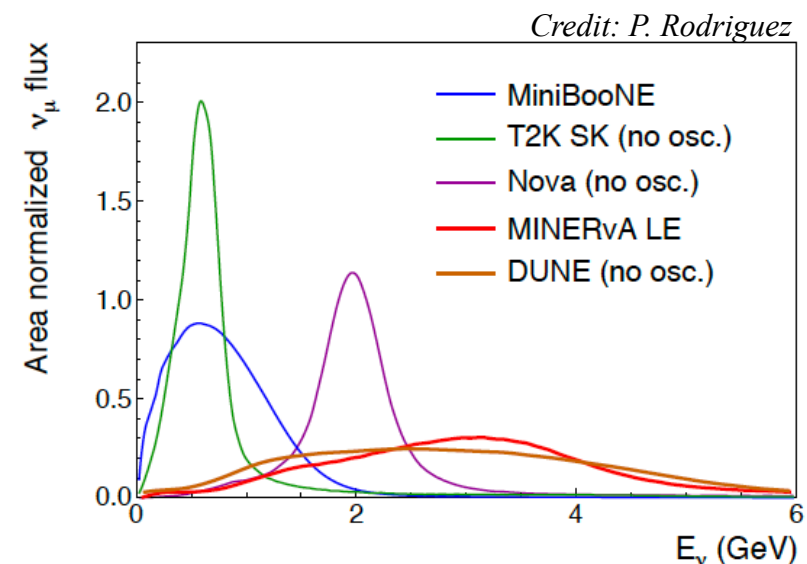
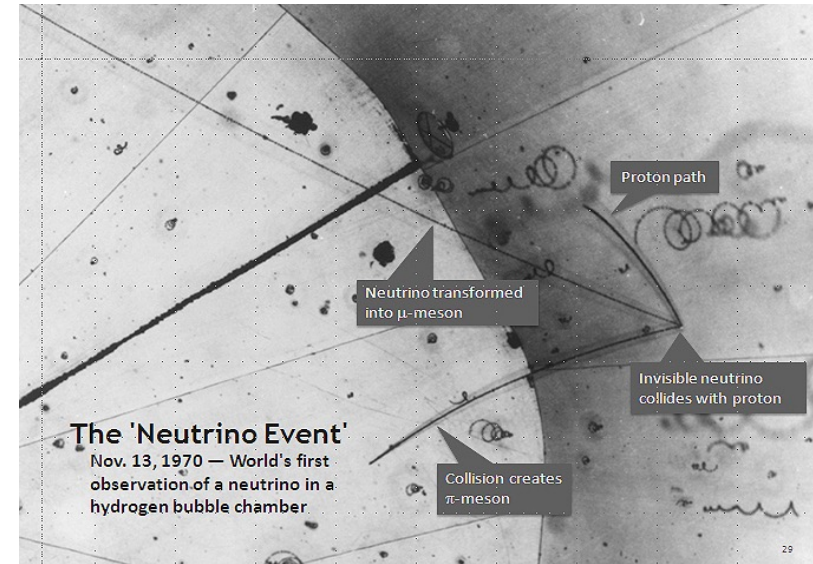
**NuMI**  
**Fermilab**



# Neutrino Detection Goals

This depends on the experiment but typically

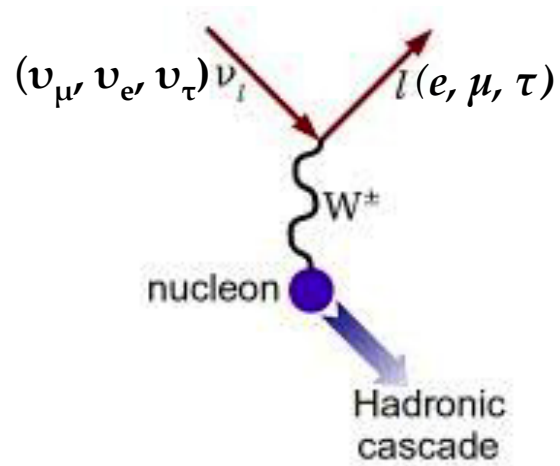
- **Identify the flavor of the neutrino**
  - Only indirect detection *via* particles produced in neutrino interactions
  - Need to know the reaction channel
    - ✓ Charged Current vs Neutral Current
    - ✓ Various interaction modes within each reaction channel
- **Measure the  $E_\nu$  as accurately as possible**
  - Not easy since neutrino sources are not always monochromatic
- **Neutrino or Anti-neutrino?**
  - differentiate this e.g. oscillation experiments aiming to measure Charge-Parity Violation
  - Charged Current interactions can provide handles



# Neutrino Flavor Tagging

**Charged Current (CC)**  
( $W^{+/-}$  exchange)

Neutrino-nucleon  
scattering



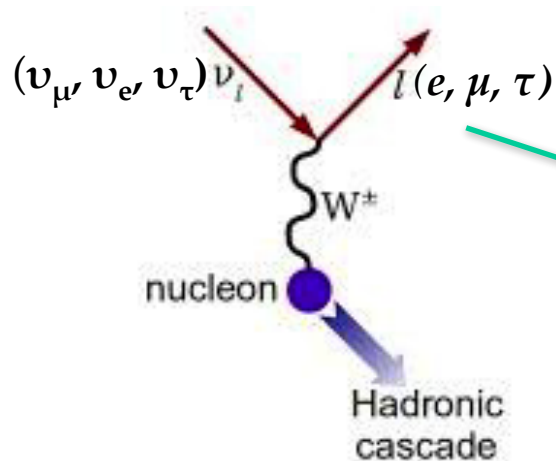
- Outgoing lepton determines the  $\nu$  flavor
- Outgoing hadrons: protons, neutrons, pions
- Typically, your signal event



# Neutrino Flavor Tagging

**Charged Current (CC)**  
( $W^{+/-}$  exchange)

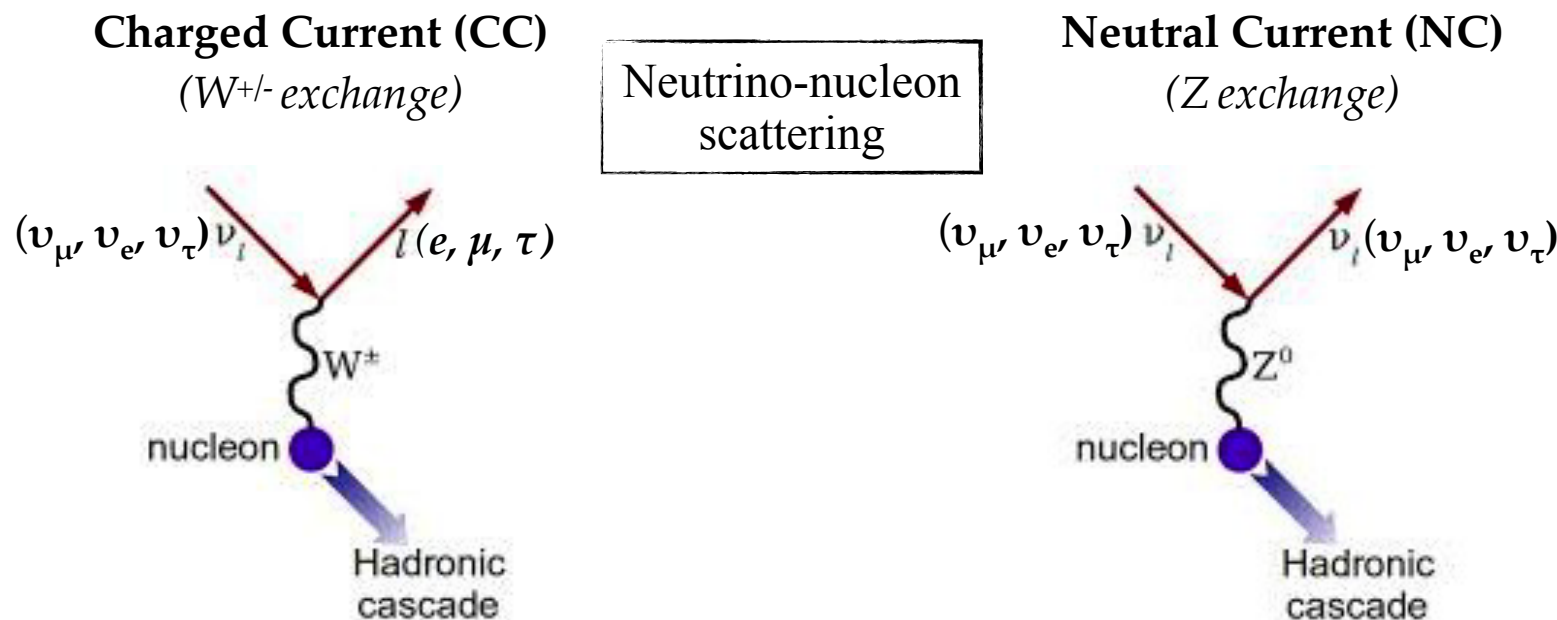
Neutrino-nucleon  
scattering



- Production of a lepton requires minimum energy. Thresholds
  - $E_\nu \sim 500 \text{ keV}$  (for an electron)
  - $E_\nu \sim 120 \text{ MeV}$  (for a muon)
  - $E_\nu \sim 3.5 \text{ GeV}$  (for a tau)

- Outgoing lepton determines the  $\nu$  flavor
- Outgoing hadrons: protons, neutrons, pions
- Typically, your signal event

# Neutrino Flavor Tagging



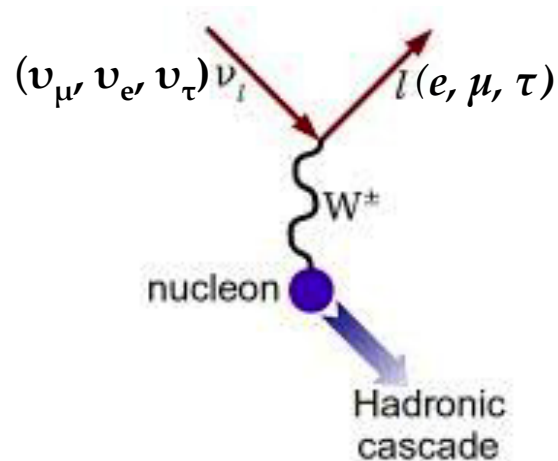
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- No outgoing lepton to tag the  $\nu$  type
- Can only see hadrons in the final state
- Typically, your background event (e.g. in “appearance oscillation” measurements)

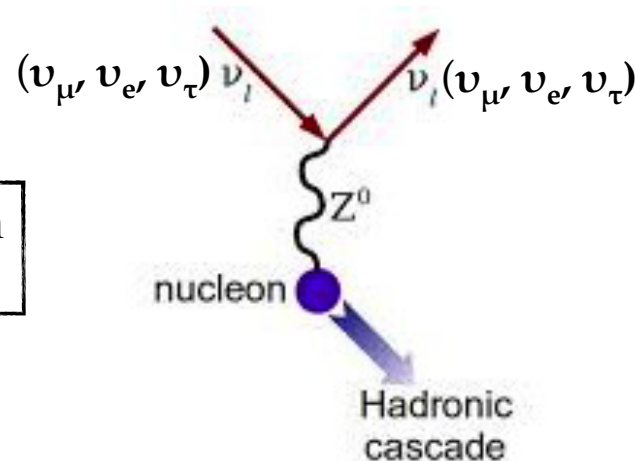


# Neutrino Flavor Tagging

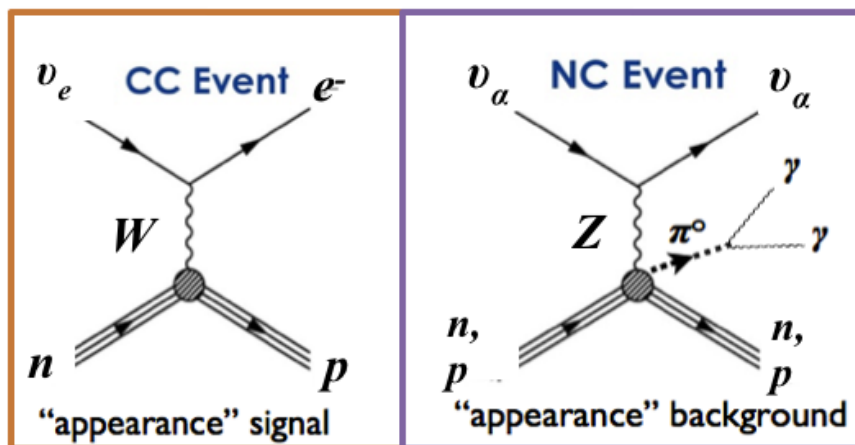
**Charged Current (CC)**  
( $W^{+/-}$  exchange)



**Neutral Current (NC)**  
( $Z$  exchange)



Neutrino-nucleon  
scattering

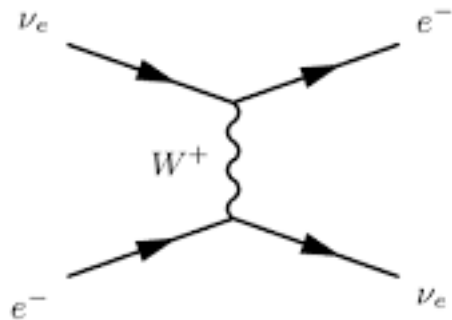


- No outgoing lepton to tag the  $\nu$  type
- Can only see hadrons in the final state
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# Neutrino Flavor Tagging

## Charged Current (CC)

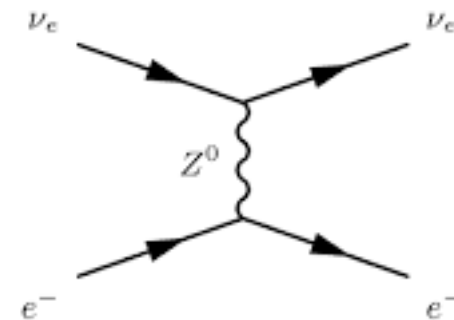
( $W^{+/-}$  exchange)



Neutrino-electron  
scattering

## Neutral Current (NC)

( $Z$  exchange)



- Neutrino scattering off of an electron
- Signal is a single final state electron



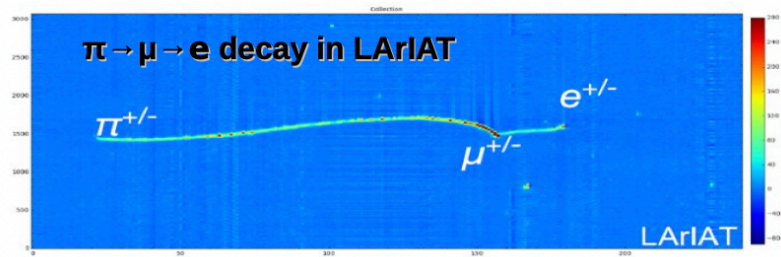
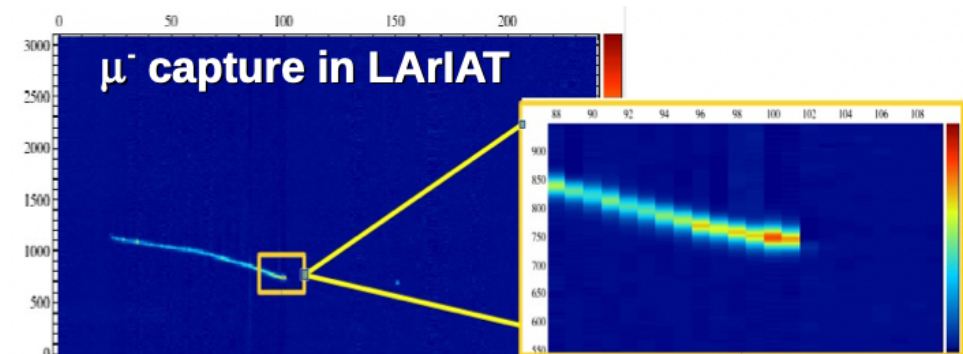
# Neutrino vs Anti-Neutrino Tagging

- Key for many experimental searches such as oscillation experiments looking to measure charge-parity violation in the neutrino sector
- Magnetic field is ideal for charge sign determination, however,
  - Neutrino detectors are typically huge
  - High volume magnetic field is hard
  - Expensive
  - Impacts other detector elements e.g. electronics



# Neutrino vs Anti-Neutrino Tagging

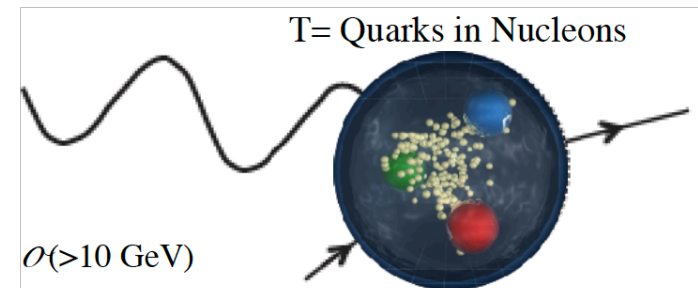
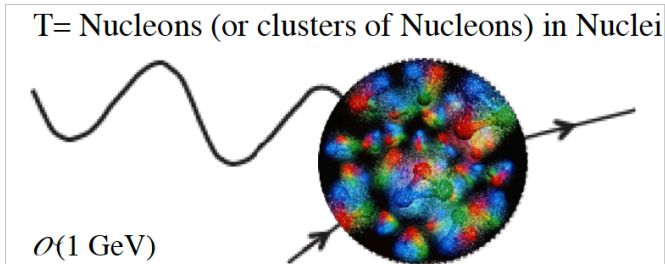
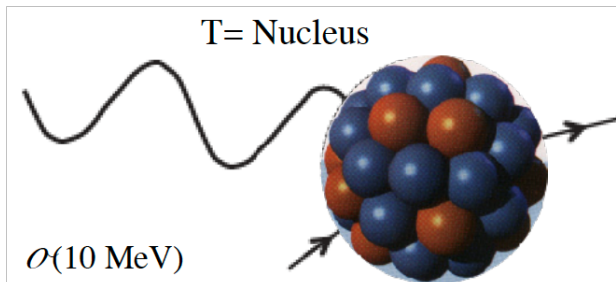
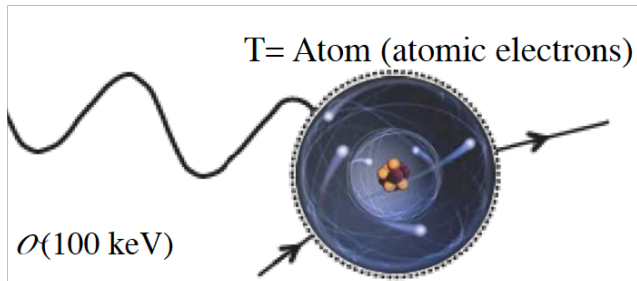
- Key for many experimental searches such as oscillation experiments looking to measure charge-parity violation in the neutrino sector
- Magnetic field is ideal for charge sign determination, however,
  - Neutrino detectors are typically huge
  - High volume magnetic field is hard
  - Expensive
  - Impacts other detector elements e.g. electronics
- One can use topology (e.g. decay vs capture) for particle sign identification in the absence of a magnetic field
- NC interactions cannot distinguish
- But, CC can distinguish  $b/n \nu$  and anti- $\nu$  using
  - opposite lepton charge
  - different final state hadrons
  - Muons from hadron decays
  - Requires good final state reconstruction





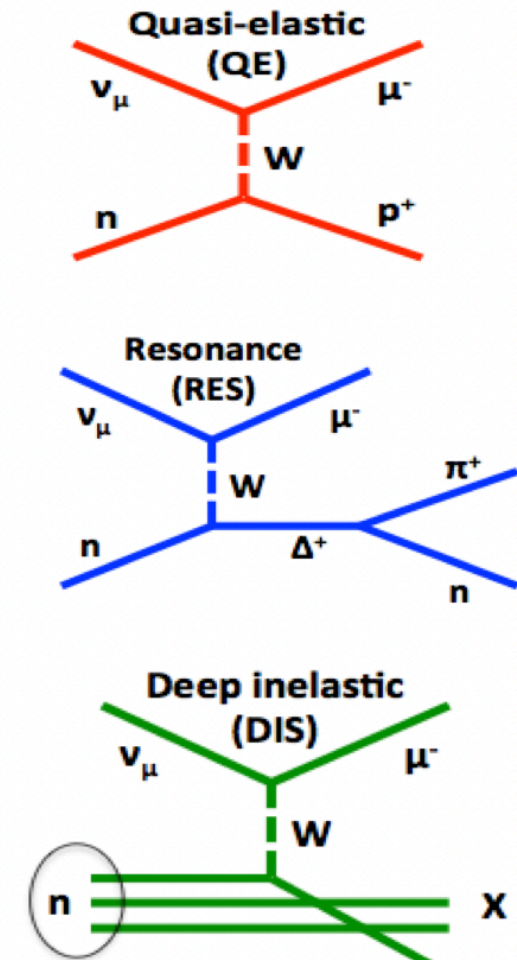
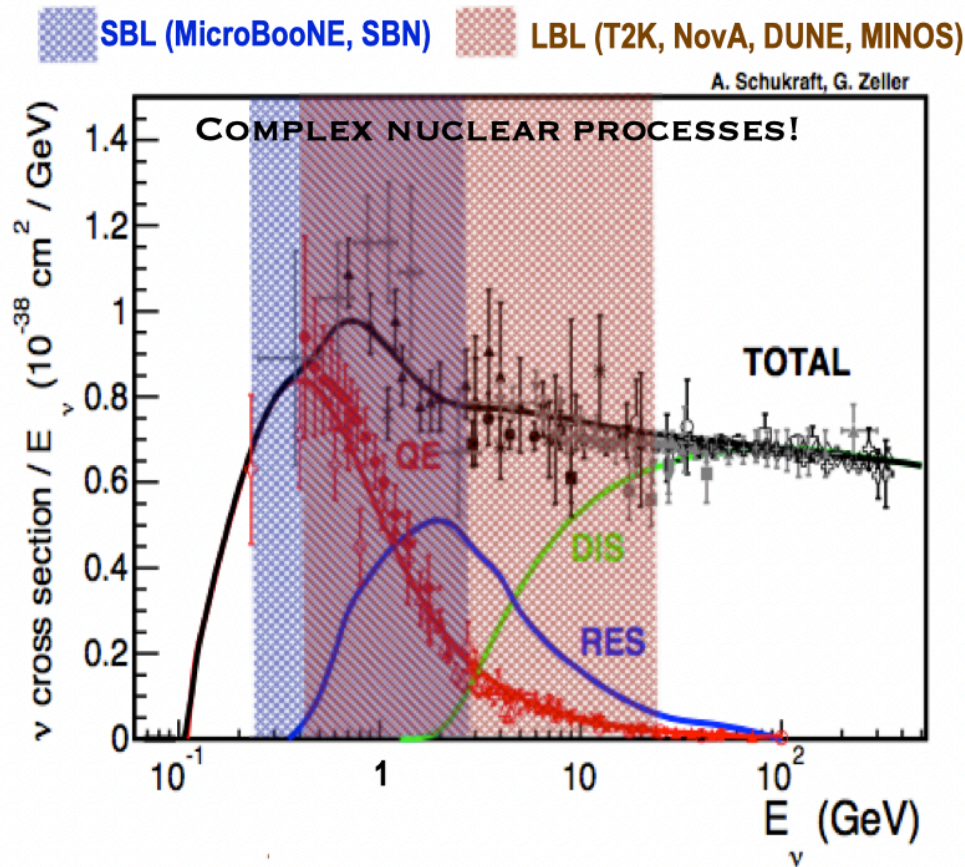
# Neutrino Interactions are Complex

Neutrinos probe matter from its Atomic structure to quark structure depending on the energy of the incoming neutrino



# Neutrino Interactions are Complex

Current and future neutrino oscillation experiments focus in the few GeV range



- Higher energies are more messy due to superposition of different channels



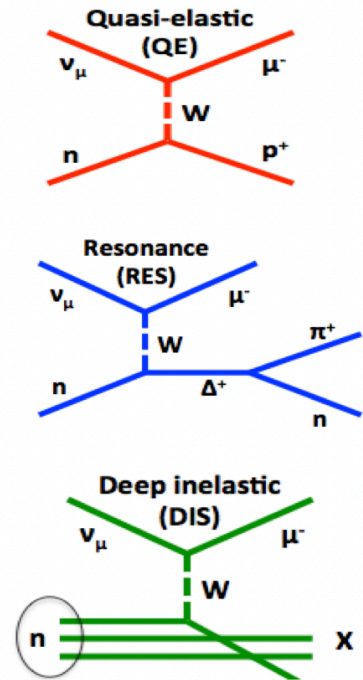
# Neutrino Energy Reconstruction

- Neutrino energy from charged lepton kinematics for CCQE

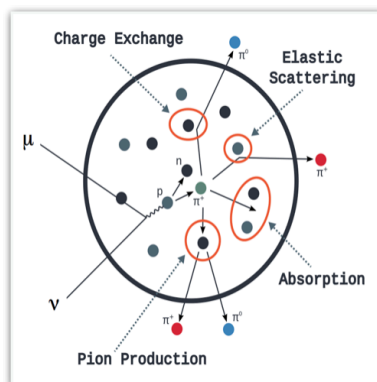
$$E_\nu = \frac{m_p^2 - m_{n'}^2 - m_\ell^2 + 2m_{n'} E_\ell}{2(m_{n'} - E_\ell + p_\ell \cos \theta_{\text{beam}})}$$

(2 body kinematics; assumes the target nucleon is at rest)

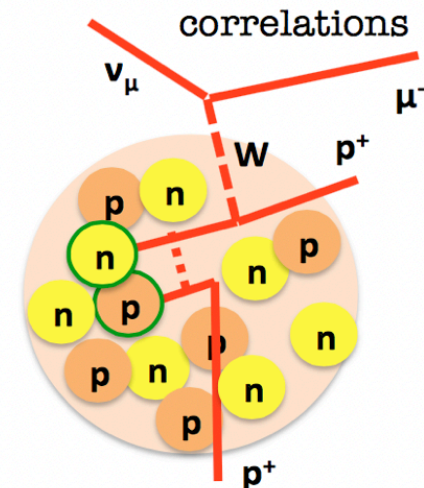
- More complicated final states for RES and DIS channels
- Both lepton and hadron kinematics important for an accurate measurement for all reaction channels
- Modern experiments use denser targets (e.g. Argon) making this picture even more complex — **thorough understanding of neutrino-nucleus interaction theory is key**



Final State Interactions

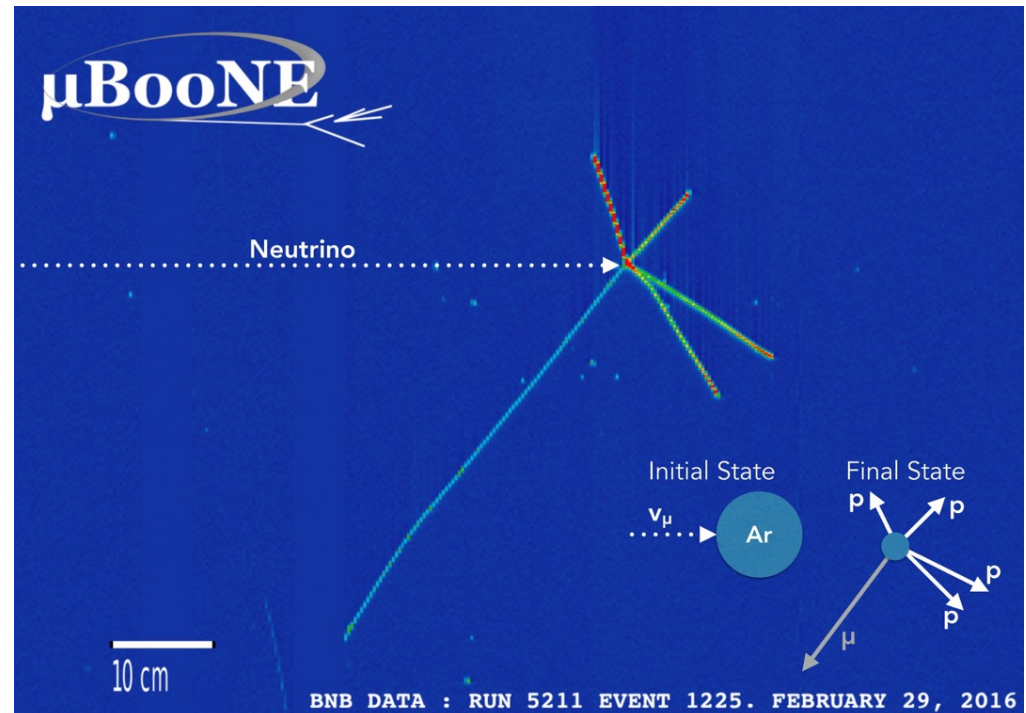
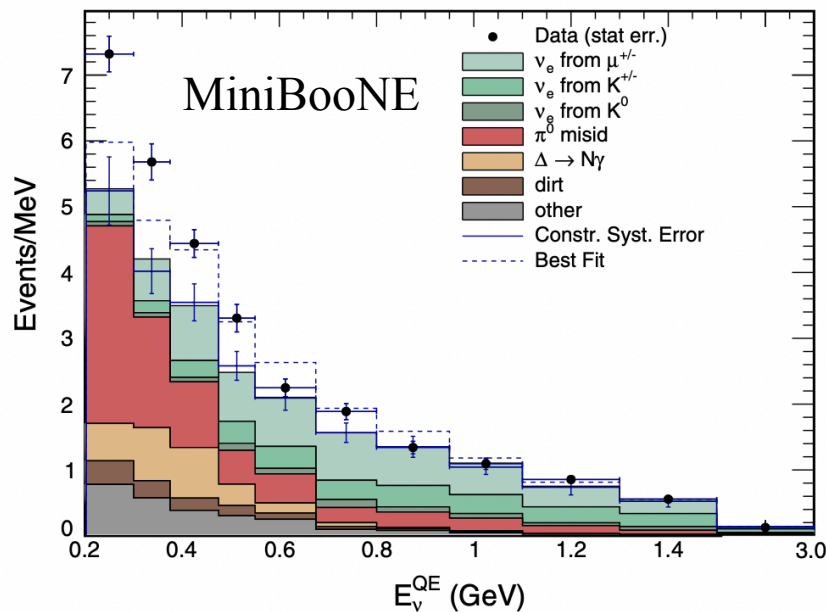


Nucleon-nucleon correlations



# Neutrino Detector Goals

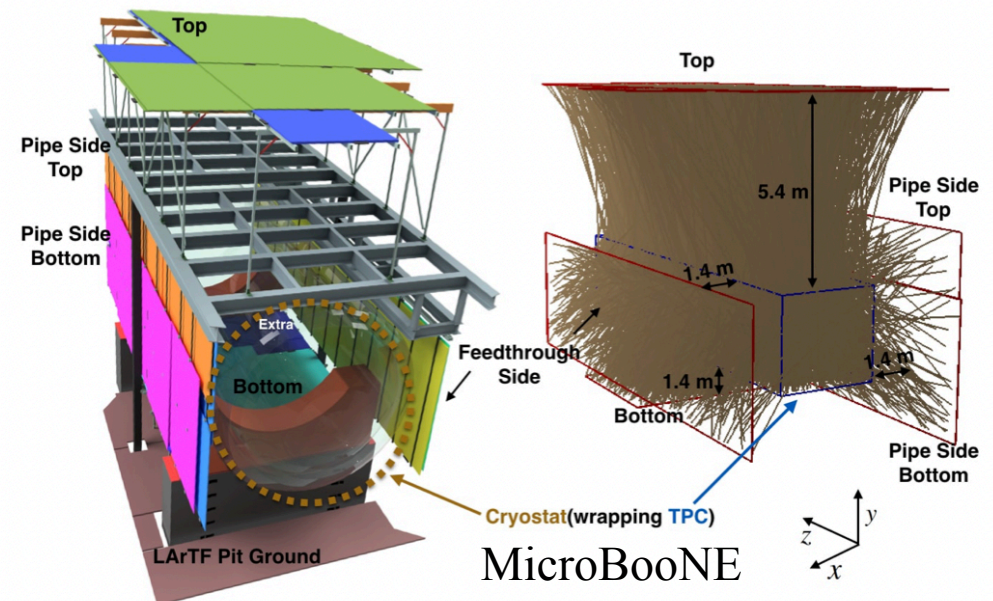
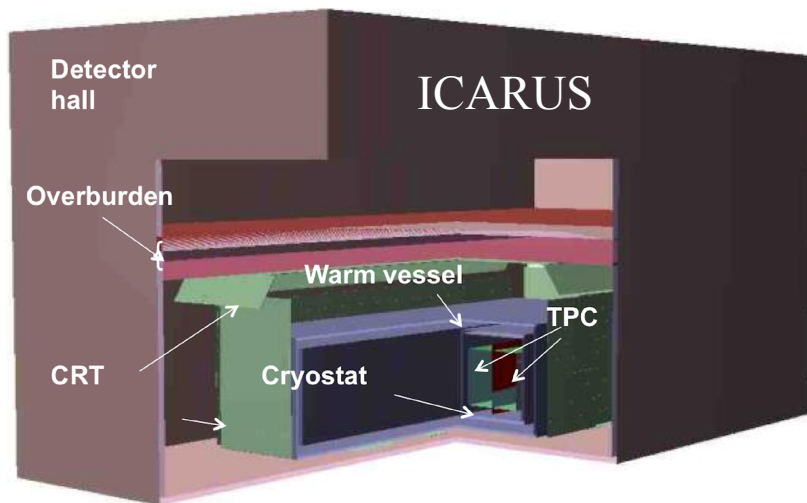
- Neutrino detectors need to work over a broad energy range (from MeV to PeV)
- They should
  - detect leptons and hadrons (protons, pions etc.)
  - distinguish electrons from photons (key for  $\nu_\mu \rightarrow \nu_e$  Appearance experiments)
  - reduce backgrounds and measure them when necessary





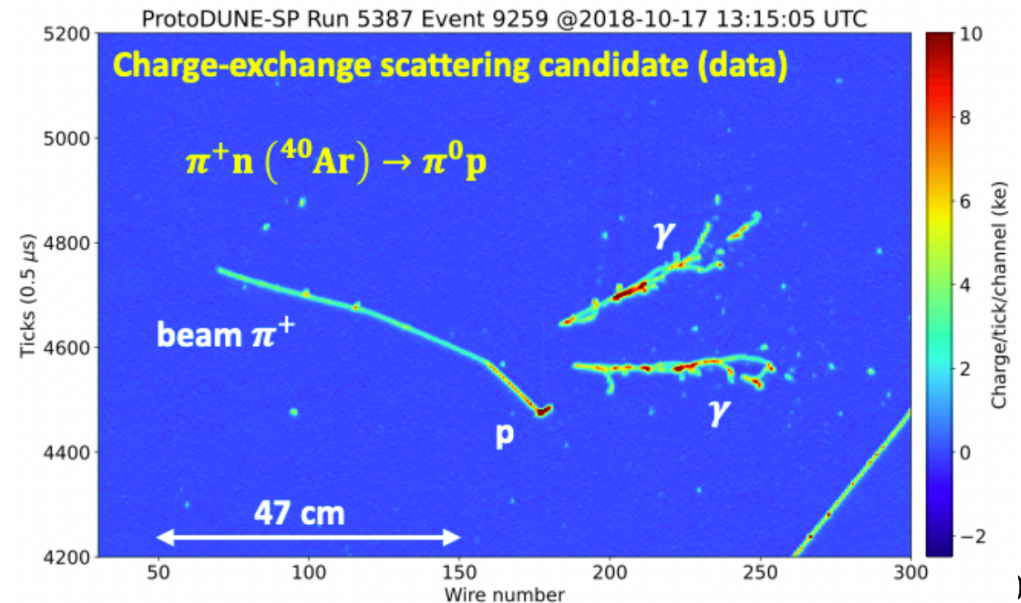
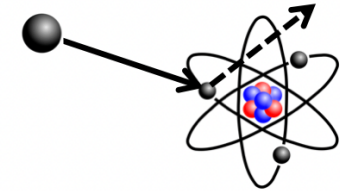
# About Backgrounds

- Backgrounds vary with neutrino energy
  - backgrounds for MeV-scale are not the same as GeV-scale neutrinos)
- Cosmic rays are a worrisome background, many ways to handle this
  - place your detector underground when possible
  - Take beam-off runs to measure cosmic ray background
  - If on surface, implement a cosmic veto/tagger system and/or shielding/overburden
- Reactors produce copious amounts of low-energy ( $< 10$  MeV) neutrinos
  - place your detectors far away from reactors
- Low-energy backgrounds are a concern for many experiments (e.g. reactor, solar / atmospheric, geo-neutrino etc.)



# Particle Interactions in Matter

- As particles move through matter, many things can happen
  - Ionization: strip electrons off of atoms in the medium
  - Scintillation: excite atoms and produce scintillation light
  - Cherenkov radiation
  - Decay into other particles
  - Produce new particles
- Many processes can occur that can result in energy loss of the particle
  - Common energy loss process is inelastic collisions with atomic electrons (Ionization)
  - Elastic scattering from nuclei
  - Atomic excitations
  - Hadronic interactions
  - Compton scattering
  - Bremsstrahlung
  - Pair production
  - Photoelectric effect...and so on





# Particle Interactions in Matter

- As particles move through matter, many things can happen

- Ionization: send electrons off atoms (the workhorse)

**Read PDG Review on this (available online)**

- Scintillation: excite atoms and produce scintillation light

- Cherenkov radiation



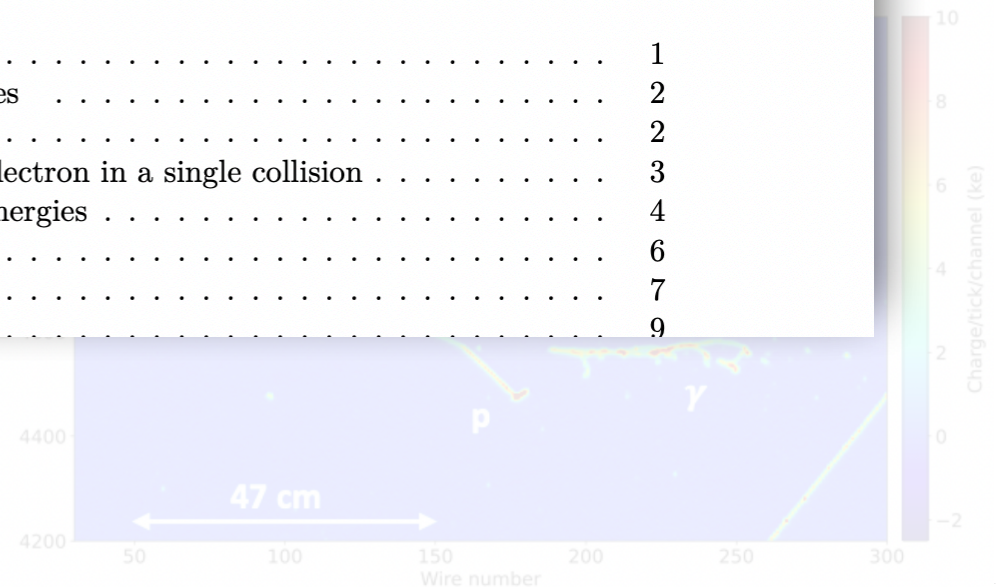
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## *34. Passage of Particles Through Matter*

### 34. Passage of Particles Through Matter

Revised August 2021 by D.E. Groom (LBNL) and S.R. Klein (NSD LBNL; UC Berkeley).

34.1	Notation . . . . .	1
34.2	Electronic energy loss by heavy particles . . . . .	2
34.2.1	Moments and cross sections . . . . .	2
34.2.2	Maximum energy transfer to an electron in a single collision . . . . .	3
34.2.3	Stopping power at intermediate energies . . . . .	4
34.2.4	Mean excitation energy . . . . .	6
34.2.5	Density effect . . . . .	7
34.2.6	Energy loss at low energies . . . . .	9



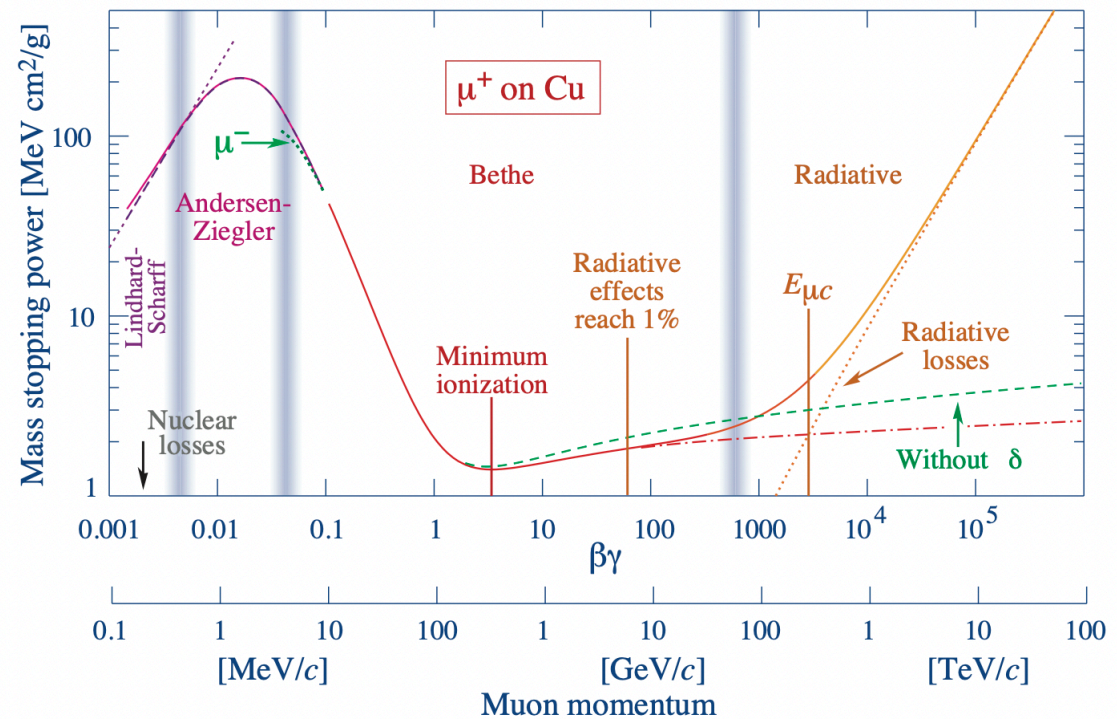
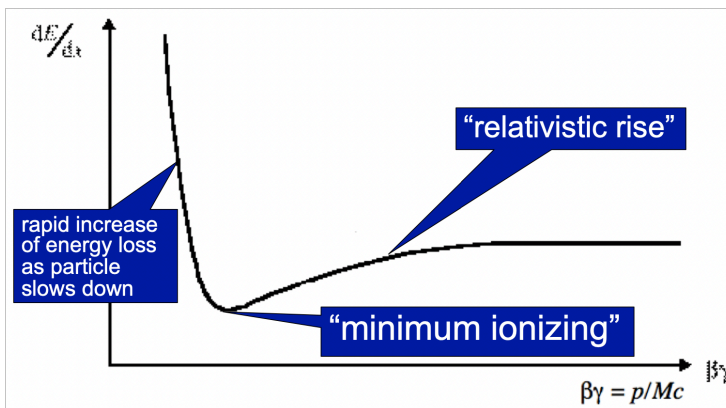
# Energy Loss of Particles (Heavy Charged Particles)

- Mean rate of energy loss for heavy (heavier than  $e^{+/-}$ ) charged particles
  - muon, pion, proton, kaon etc.

- Good to a few % in the MeV to GeV energy range and for intermediate Z materials

$$0.1 \lesssim \beta\gamma \lesssim 1000$$

Basic shape of the curve



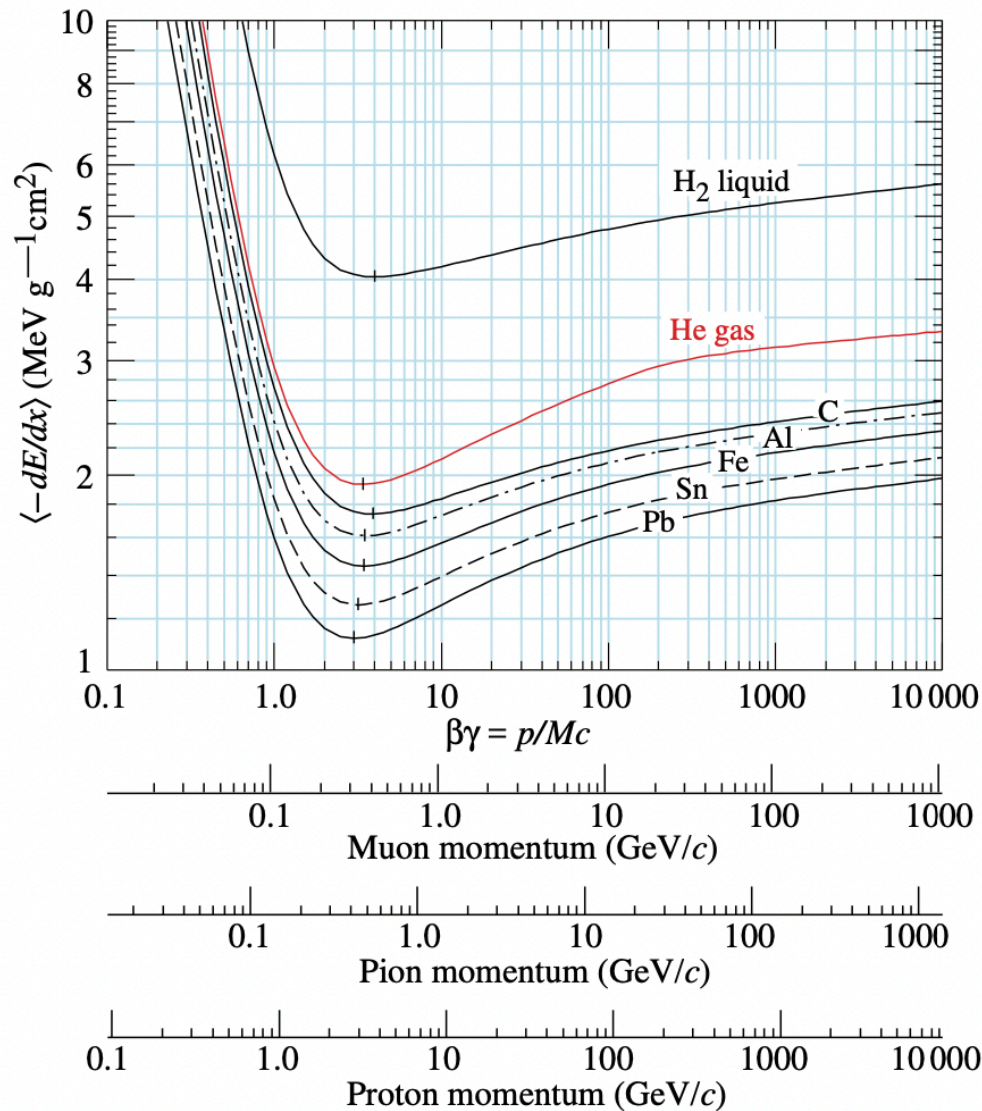
$$\left\langle -\frac{dE}{dx} \right\rangle = K z^2 \frac{Z}{A} \frac{1}{\beta^2} \left[ \frac{1}{2} \ln \frac{2m_e c^2 \beta^2 \gamma^2 W_{\max}}{I^2} - \beta^2 - \frac{\delta(\beta\gamma)}{2} \right]$$

Bethe-Bloch Equation



# Energy Loss of Particles (Heavy Charged Particles)

- Mean rate of energy loss for different materials and different incident particles

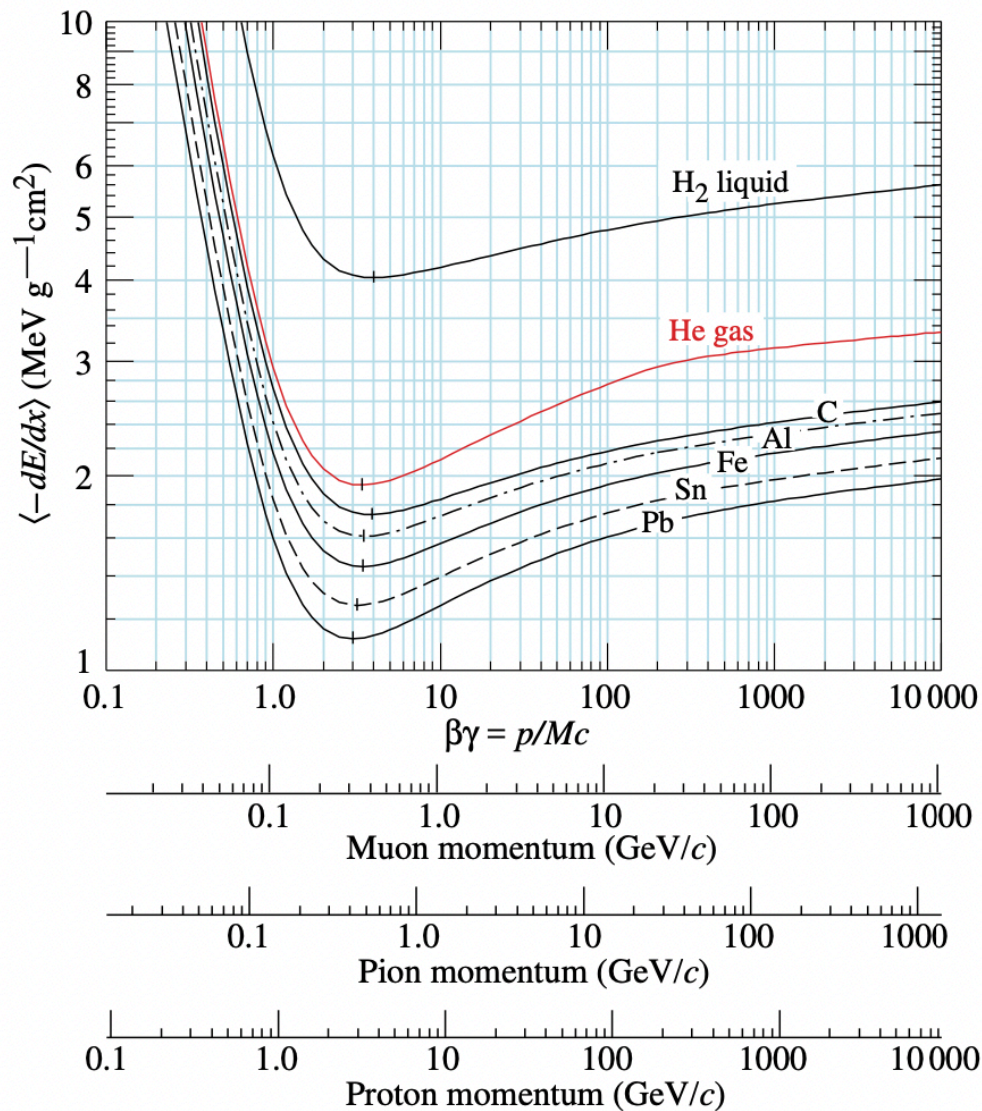


Material	Minimum Ionizing $dE/dx$ (MeV/cm)
Liquid Argon	2.1
Water	2.0
Steel	11.4
Scintillator (CH)	1.9
Lead	12.7

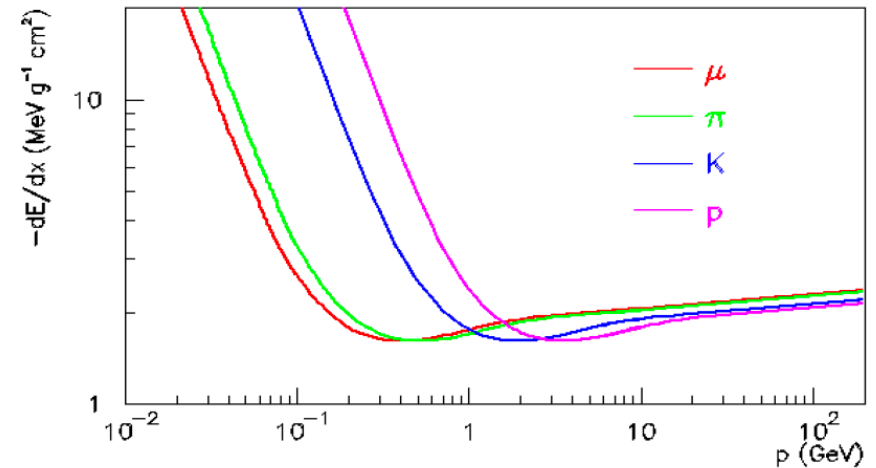
- Stopping power is an important number as it tells you how far a particle can travel in the detector medium
- Important for determining how big your detector needs to be

# Energy Loss of Particles (Heavy Charged Particles)

- Mean rate of energy loss for different materials and different incident particles



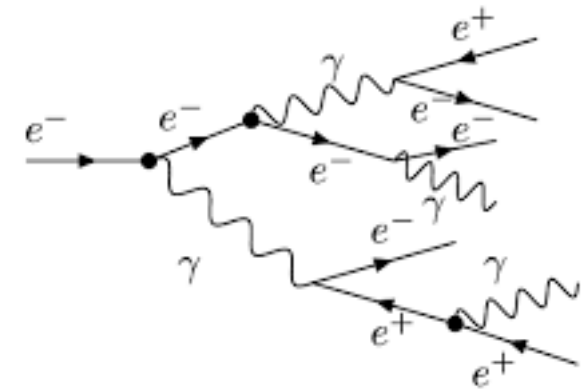
If you know  $dE/dx$  and  $p$ , you can do particle ID — a common technique used in experiments





# Energy Loss of Particles (Electrons)

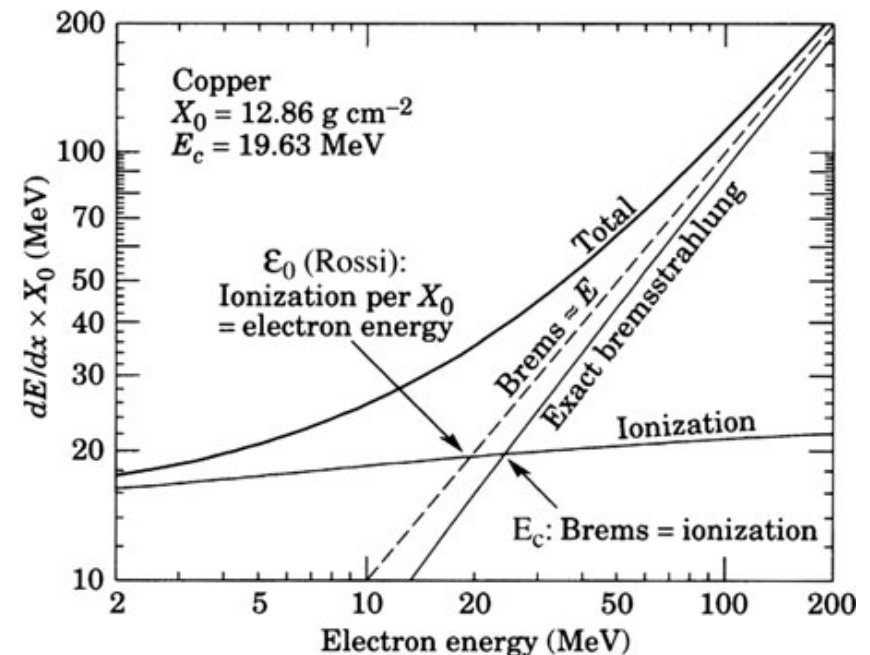
- Electrons act differently in a detector medium
- Electrons above the critical energy ( $E_c$ ) lose energy via both radiation (Bremsstrahlung) and collisional processes resulting in an electromagnetic shower (or spray) in the detector
  - Electrons above  $E_c$  create photons via Bremsstrahlung which then produce  $e^+e^-$  pairs which then go on to produce more photons until energies drop below  $E_c$
- Critical energy depends on the medium (typically few tens of MeV) and is defined as the cross over point where Bremsstrahlung energy loss > ionization loss



Bethe-Heitler approximation

$$E_c \approx \frac{1600m_e c^2}{Z}$$

Material	$E_c$ (MeV)
Lead	9.51
Aluminum	51.0
Iron	27.4
Copper	24.8
Water	92



# Energy Loss of Particles (Electrons)

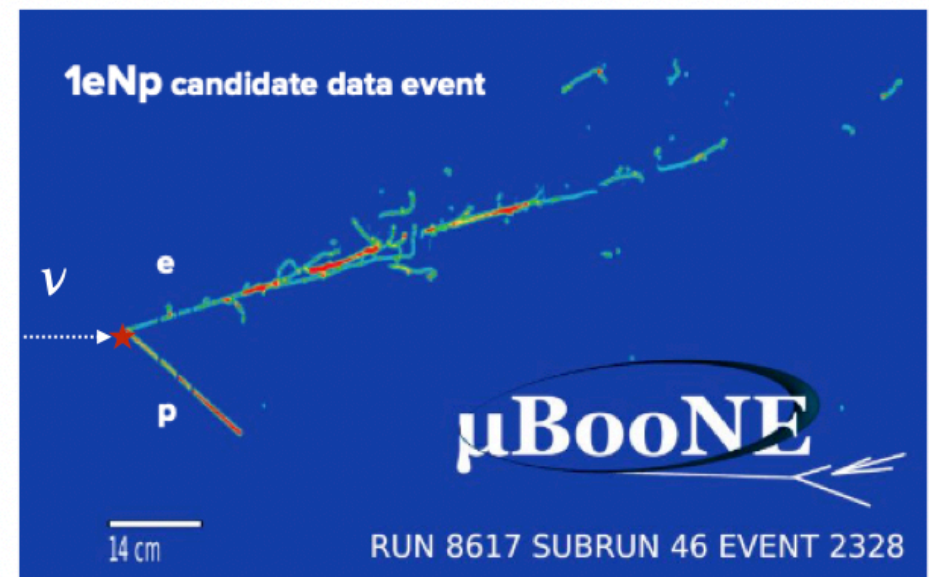
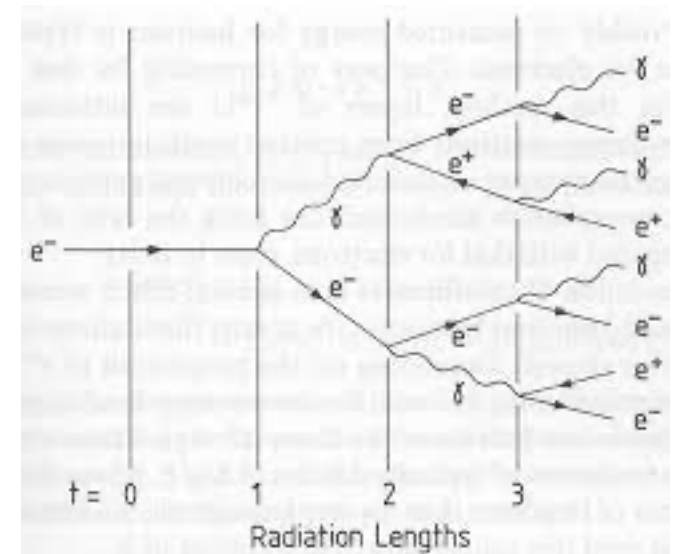
- Radiation length ( $X_0$ ) is an important parameter for electromagnetic showers
- It is defined as the distance over which electrons lose  $1/e$  of their energy by radiation i.e., at every  $X_0$  an electron would Bremsstrahlung creating a photon

$$X_0 = \frac{716.4A}{Z(Z+1)\ln(287/\sqrt{Z})} \left[ \frac{\text{g}}{\text{cm}^2} \right]$$

- Distance over which photons pair produce is  $(9/7) X_0$
- Electromagnetic shower size estimate after “t” radiation lengths

$$t_{\text{max}} = \frac{\ln(E_0/E_c)}{\ln 2}$$

Material	$X_0$ (cm)
Liquid Argon	14
Water	37
Steel	1.76
Scintillator (CH)	42
Lead	0.56

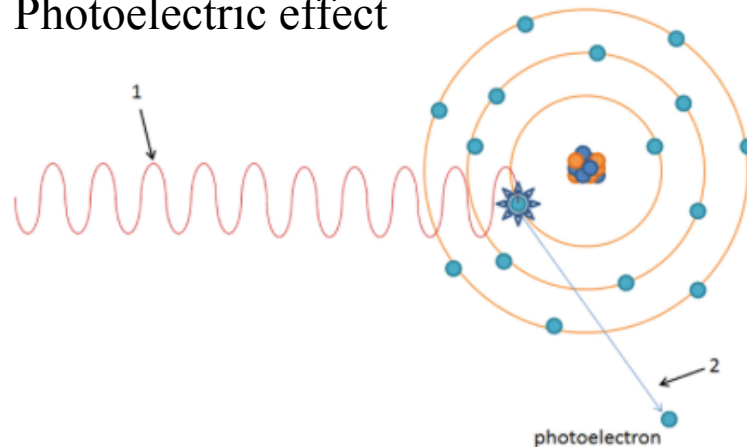




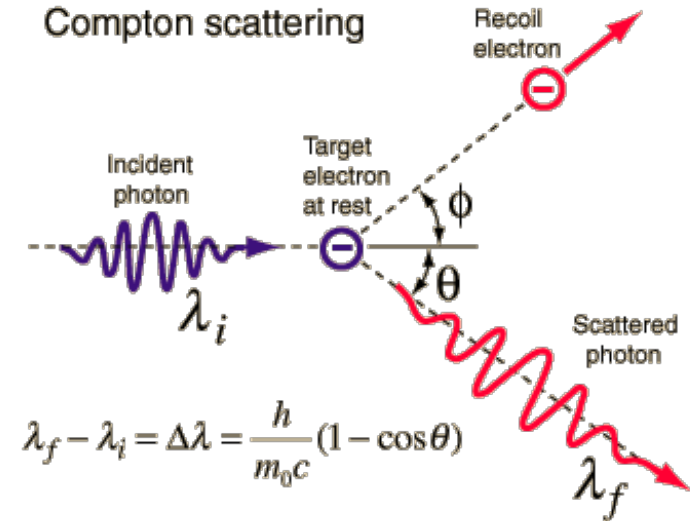
# Energy Loss of Particles (Photons)

- Photon primarily lose energies through 4 ways

Photoelectric effect



Compton scattering



Pair production

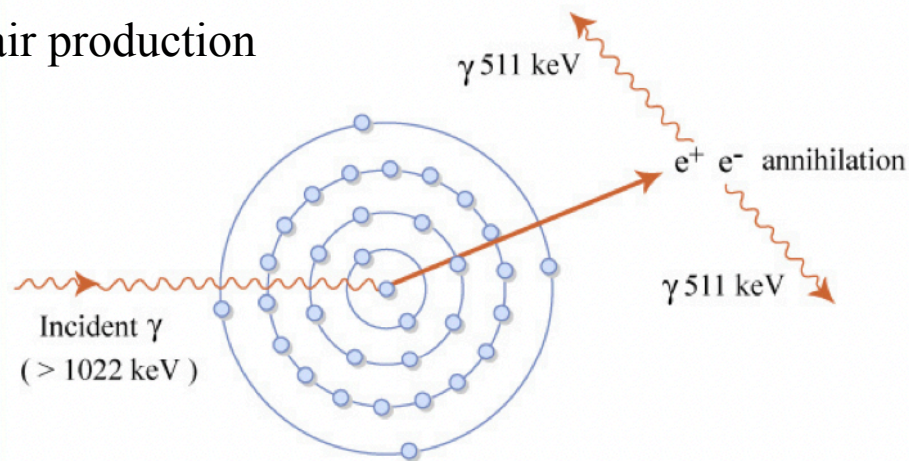
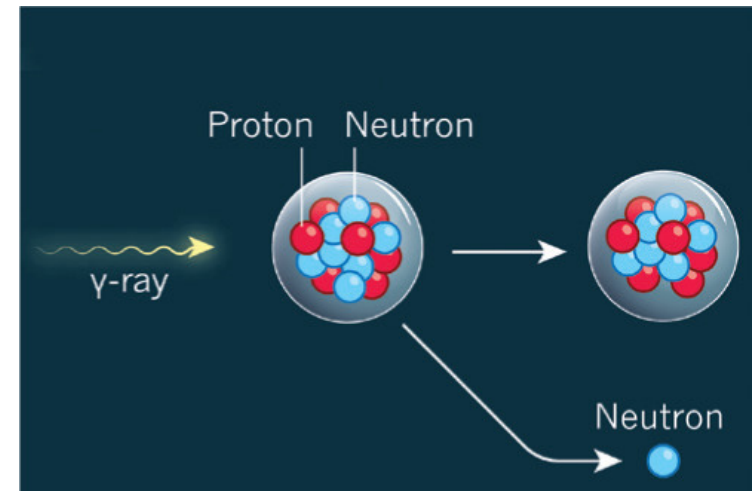


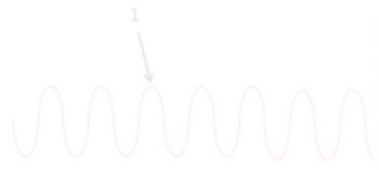
Photo-nuclear effect



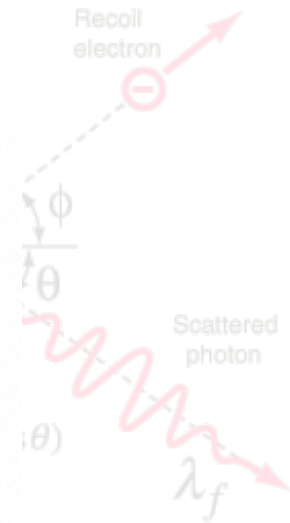
# Energy Loss of Particles (Photons)

- Photon primarily lose energies through 4 ways

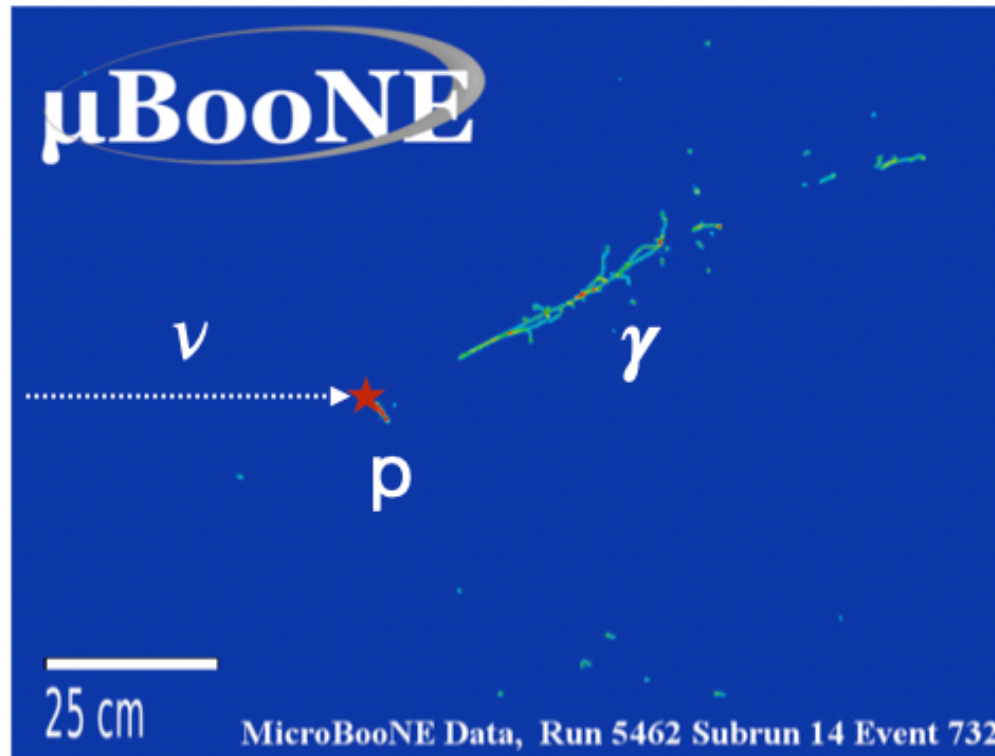
Photoelectric effect



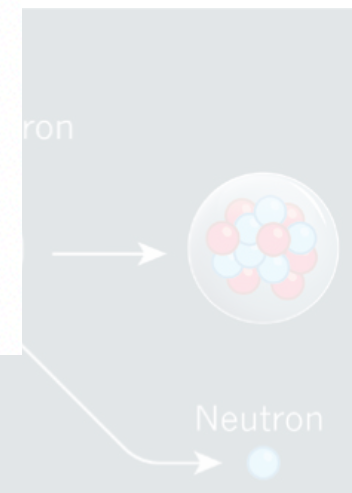
Compton scattering



Pair production



Nuclear effect

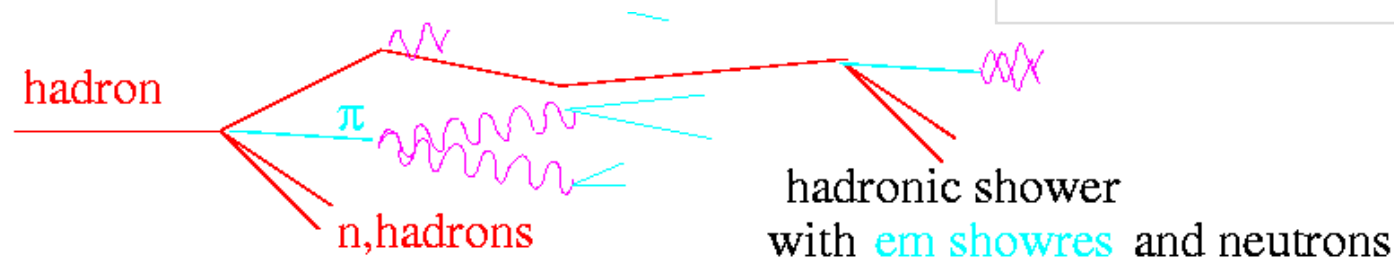
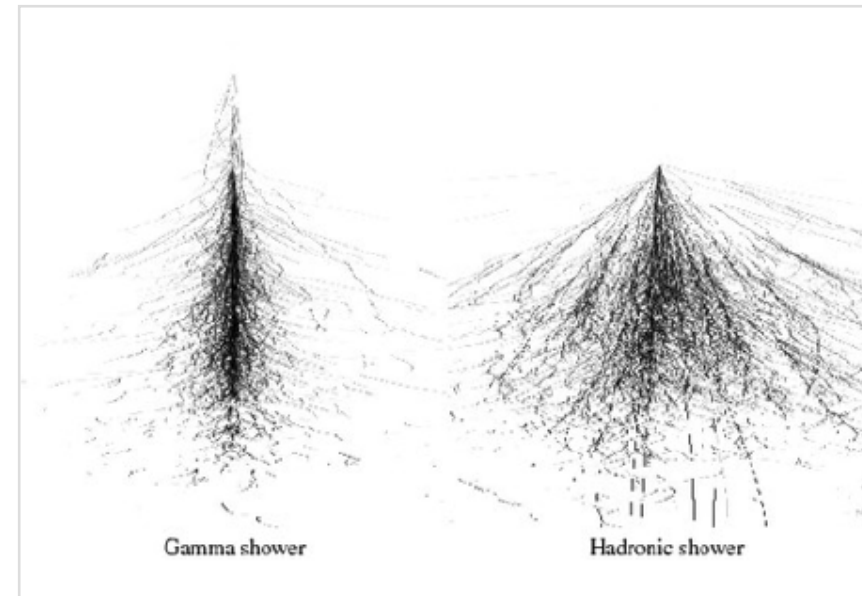




# Energy Loss of Particles (Hadrons)

- For hadrons, interaction ( $\lambda$ ) length defines the distance they travel before undergoing a strong nuclear interaction
- Neutron, for example, interacts via the weak force and undergoes energy loss through various mechanisms
  - Elastic scattering, inelastic scattering, neutron capture, hadronic showers etc.
  - Elastic scattering on nuclei is the main mechanism of energy loss for neutrons
  - Neutron capture important for low energy neutrino experiments
- Hadron interaction lengths are longer than radiation lengths
- At higher energies, hadronic showers initiated by e.g. proton, neutron become relevant
- Hadronic showers are broader than EM showers

Material	X <sub>0</sub> (cm)	$\lambda_{\text{INT}}(\text{cm})$
Liquid Argon	14	83.5
Water	37	83.6
Steel	1.76	17
Scintillator (CH)	42	~80
Lead	0.56	17





Thank you, we will continue tomorrow!

