

# (Re)-Assessing the Role of Meson-Exchange-Currents in Lepton-Nucleus Interactions

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# OUTLINE

- ★ Preamble
- ★ Qualitative features of MEC contributions emerging from early studies of electron-nucleus scattering
- ★ Studies based on more realistic models of electron-nucleus scattering
  - ▶ Superscaling
  - ▶ Factorisation
- ★ Specific issues associated with the description of neutrino-nucleus scattering
  - ▶ Flux average
  - ▶ Contribution of axial current
- ★ Results of recent studies
- ★ Summary & Outlook

## PREAMBLE

- ★ Nuclear Many-Body Theory is based on the tenet—strongly supported by low-energy nuclear phenomenology—that nucleon dynamics below pion production threshold can be described by the Hamiltonian

$$H_A = \sum_{i=1}^A \frac{\mathbf{p}_i^2}{2m} + \sum_{j>i=1}^A v_{ij} + \dots ,$$

and the associated electro-weak current operator

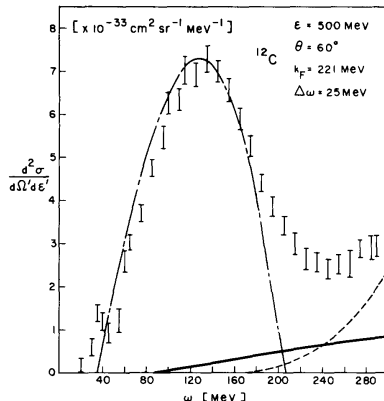
$$J_A(q) = \sum_{i=1}^A j_i + \sum_{j>i=1}^A j_{ij} + \dots .$$

- ★ In the late 1970s, significant two-nucleon meson-exchange current (MEC) contributions, leading to the excitation of 2p2h final states, were advocated to explain inclusive electron scattering cross sections in the *dip* region, between the quasi-elastic and  $\Delta$ -production peaks, corresponding to  $\omega \approx Q^2/m$

## EARLY ELECTRON SCATTERING STUDY

- ★ The amount of strength needed to describe the data depends on the treatment of the one-nucleon current contributions

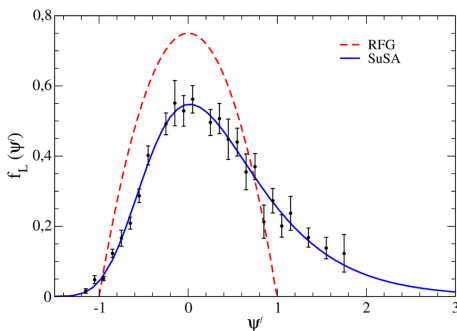
- ★ Calculations by  
T.W. Donnelly  
*et al.*, PLB **76**, 393 (1978).  
RFGM + MEC
- ★ Data from E.J. Moniz  
*et al.*, PRL **26**, 445 (1971)



- ★ Most studies only considered transitions to 2p2h final states, neglecting interference between 1p1h and 2p2h amplitudes altogether

# ENTER NUCLEAR DYNAMICS

- ★ Taking into account the effects of nuclear dynamics not included in the RFGM leads to the appearance of sizeable asymmetric tails, originating from both initial state dynamics—primarily Short-Range Correlations (SRC)—and Final State Interactions (FSI)



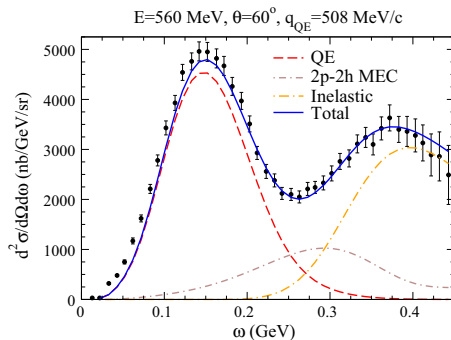
- ★ J.E. Amaro *et al.*,  
Eur. Phys. J. Spec. Top.  
**230**, 4321 (2021)
- ★ Data from J. Jourdan,  
NPA **603**, 17 (1996)

- ★ One- and two-nucleon current contributions should be treated consistently, within a unified realistic model

## MORE ADVANCED MODELS: SuSA

- ★ In the phenomenological superscaling model (SuSA) the single-nucleon knock out contribution to the cross section is obtained from electron scattering data
- ★ MEC contributions computed within the RFGM, including only transitions to 2p2h final states

- ★ J.E. Amaro *et al.*,  
Eur. Phys. J. Spec. Top.  
**230**, 4321 (2021)
- ★ Carbon data from  
P. Barreau *et al.*, NPA **402**,  
515 (1983)

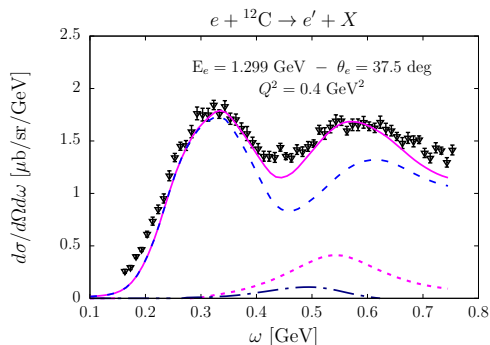


- ★ The SuSA model is inherently unable to take into account interference between 1p1h and 2p2h amplitudes

## MORE ADVANCED MODELS: FACTORISATION

- ★ Factorisation of the nuclear cross section allows to treat one- and two-body current contributions within a consistent framework, using spectral functions obtained from a state-of-the-art microscopic model of nuclear dynamics and fully relativistic current operators

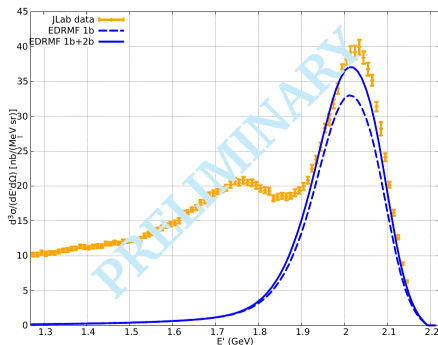
- ★ N. Rocco *et al.*, PRL **116**, 192501 (2016)
- ★ Carbon data from R. Sealock *et al.*, PRL **62**, 1350 (1989)



- ★ Transitions to 1p1h final state induced by the two-nucleon current are neglected

## RECENT STUDIES: ED-RMF MODEL

- ★ Relativistic Mean-Field Model, extended to consistently describe FSI and corrected to account for ground-state correlations;  
T. Franco-Munoz *et al.* [arXiv:2203.09996 \[nucl-th\]](#)
- ★ Fully relativistic current operators, transitions to  $1p1h$  final states and interference contributions included
- ★ Figure courtesy of Raul Gonz  les-Jim  nez
- ★ Data from H. Dai *et al.* PRC **99**, 054608 (2019)
- ★ The two-body current contribution, yielding  $\sim 20\%$  of the full cross section, peaks at  $\omega \approx Q^2/2m$

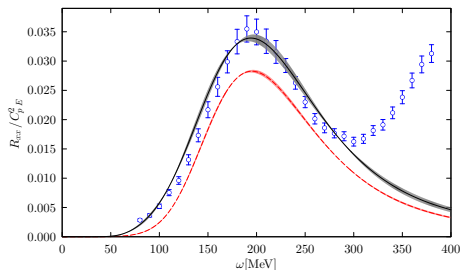


- ★ Distinct energy dependence of MEC contributions



# HINTS FROM GFMC CALCULATIONS

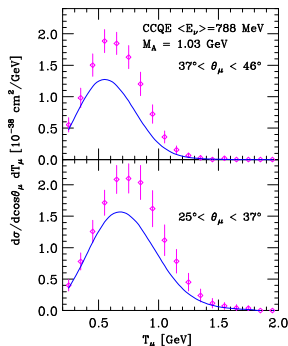
- ★ The Green Function Monte Carlo (GFMC) technique allows one to perform *ab initio* calculations of the nuclear response in the non relativistic regime. All one- and two-body current contributions, including interference, consistently taken into account
- ★ Lovato *et al.*, PRC **91**, 062501 (2015); red: one-body current only; blue: full calculation
- ★ Transverse response of Carbon at  $q = 600$  MeV from J. Carlson *et al.* PRC **65**, 024002 (2002);



- ★ Transitions to 1p1h final states and interference appear to play a critical role

# NEUTRINO-NUCLEUS INTERACTIONS

- ★ State-of-the-art models of electron-nucleus scattering largely fail to explain the measured neutrino cross sections in terms of single-nucleon knock out induced by the one-body current
- ★ Benhar *et al.* PRL **105**, 132301 (2010)
- ★ MiniBooNE data, PRD **81**, 092005 (2010)

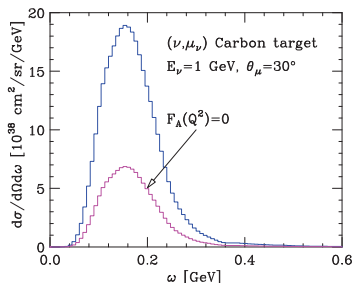


Processes involving MEC—somewhat misleadingly referred to as 2p2h contributions—have been advocated as the main source of the missing strength

# NEUTRINO-NUCLEUS INTERACTIONS

- ★ Compared to the electron-nucleus cross sections, the measured neutrino-nucleus cross sections involve two important differences:
  - ▶ the average over a broad neutrino flux, which severely hampers a clear-cut identification of different reaction mechanisms
  - ▶ a large contribution of the axial-vector current

- ★ Jen *et al.* PRD **90**, 093004 (2014); dipole fit with  $M_A = 1.03$  GeV



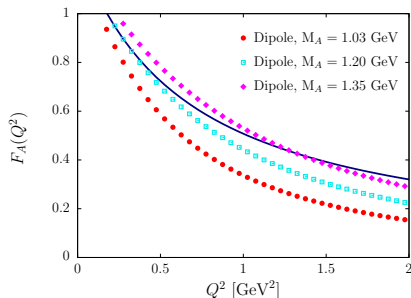
- ★ The size of the contribution from two-nucleon currents is largely driven by the uncertainty on the axial structure of the nucleon

## $Q^2$ -DEPENDENCE OF THE AXIAL FORM FACTOR

- ★ The available data—from, e.g., MiniBooNE [PRD **81**, 092005 (2010)] and T2K [PRD **92**, 112003 (2015)] can be explained by significantly increasing the nucleon axial mass from its *canonical* value  $M_A = 1.03 \text{ GeV}$ . However, in the absence of a convincing motivation, this prescription appears to be largely arbitrary

- Comparison between the results of Park et al. [PRD **105**, 054505 (2022)] (full line) obtained from lattice QCD, and the dipole parametrisation

$$F_A(Q^2) = g_A (1 + Q^2/M_A^2)^{-1}$$

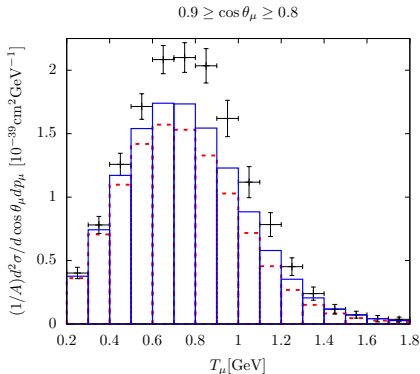


- ★ At  $Q^2 \lesssim 0.5 \text{ GeV}^2$  the dipole fit with  $M_A = 1.2 \text{ GeV}$  is remarkably close to the results of Park *et al*,

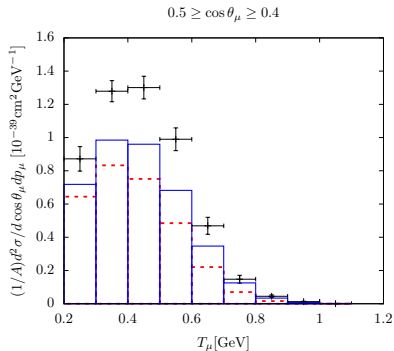
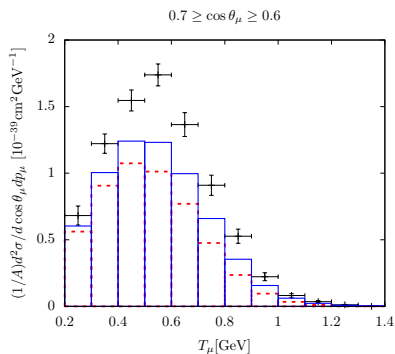
## COMPARISON TO MINIBOONE DATA (FROM OB'S TALK AT NUINT22)

- ★ Replacing the  $M_A = 1.03$  MeV dipole parametrisation with the lattice QCD axial form factor of Parks *et al.* leads to a  $\sim 10 - 15\%$  enhancement of the single-nucleon knock out cross section, entailing a corresponding reduction of the missing strength

- ★ Theoretical calculations carried out using a realistic model of the carbon spectral function [same as in PRL **105**, 132301 (2010)]

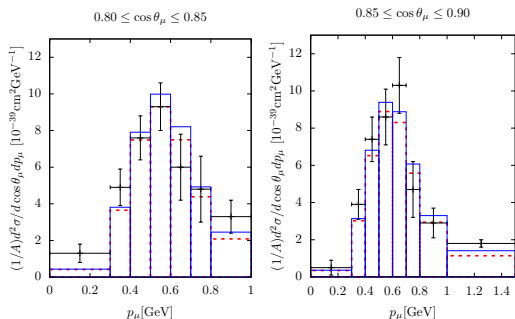


★ Similar pattern observed at all muon emission angles



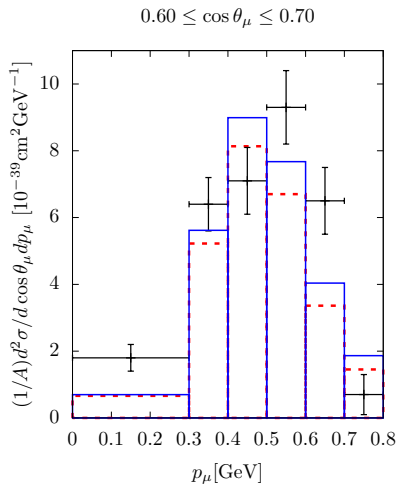
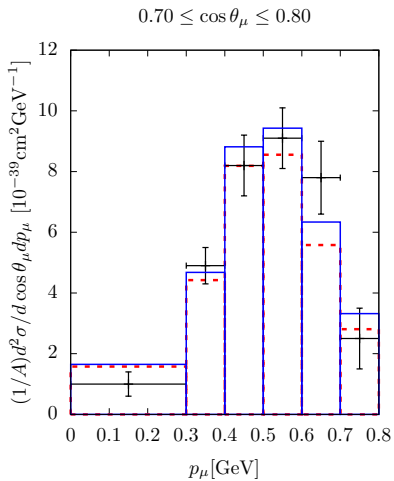
## COMPARISON TO T2K DATA (FROM OB'S TALK AT NUINT22)

- ★ A comparison to T2K CCQE data [K. Abe et al., PRD 93, 112012 (2016)] suggests in this instance there is less room for contributions other than single-nucleon knock out



- ★ This observation is consistent with the results of the analysis of T2K data based on the dipole parametrisation of the axial form factor, yielding  $M_A = 1.26\text{ GeV}$  (to be compared with  $M_A = 1.35\text{ GeV}$  reported by MiniBooNE)

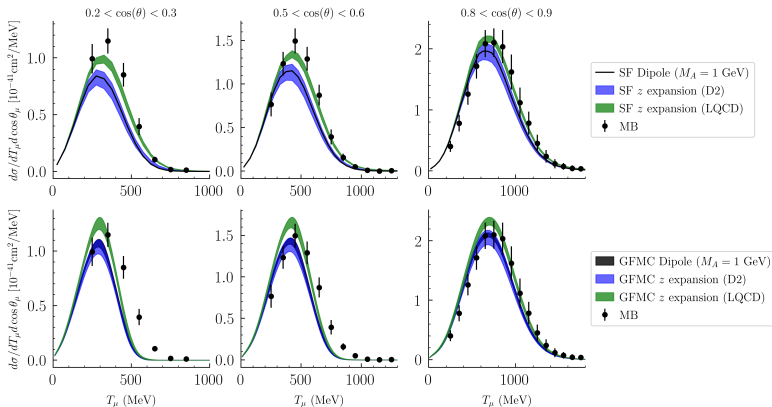
- ★ Similar pattern observed at all muon emission angles





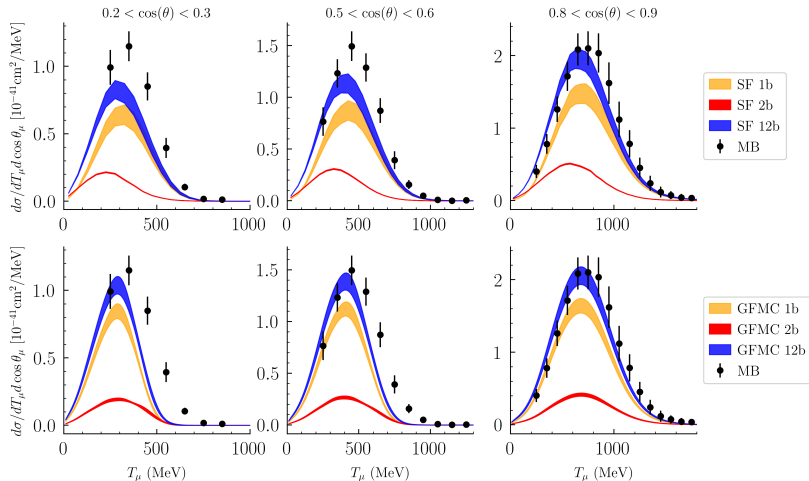
# SIMILAR ANALYSIS BY SIMONS *et al.* ([ARXIV:2210.02455](https://arxiv.org/abs/2210.02455) [HEP-PH])

- ★ MiniBooNE data analysed using the GFMC and SF formalisms and different prescriptions for the axial form factor

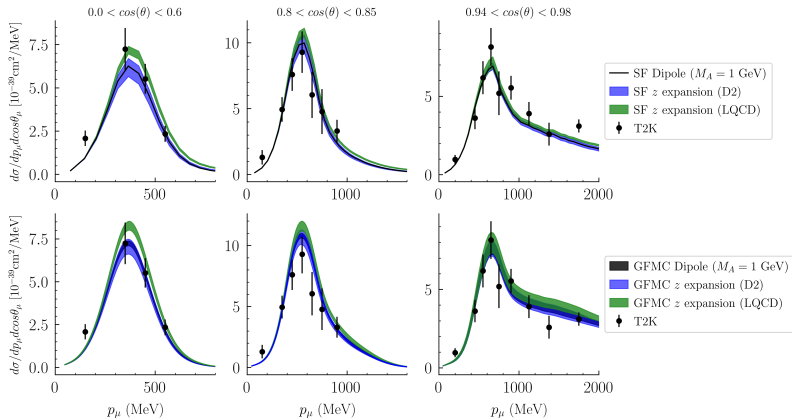


# ONE- AND TWO-NUCLEON CURRENT CONTRIBUTIONS

★ SF results do not include interference



# COMPARISON TO T2K DATA



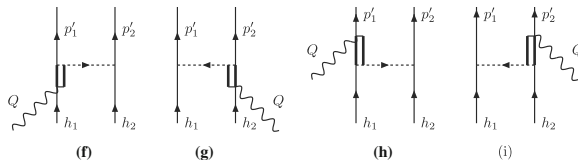
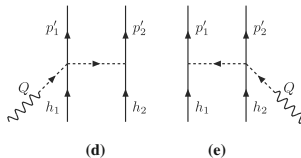
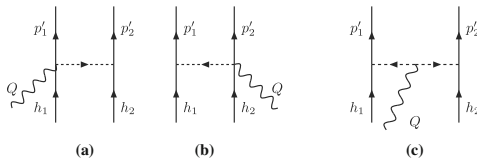
## SUMMARY & OUTLOOK

- ★ The measured neutrino-nucleus cross sections exposed specific features of MEC contributions which had not, or could not, been studied by electron scattering experiments
- ★ One- and two-nucleon contributions to the nuclear cross sections must be treated *consistently* within a *realistic* model of nuclear dynamics, including *interference* and using *fully relativistic currents*
- ★ Recent results suggest that the role of interference contributions involving 1p1h transitions induced by MEC, which received little or no attention in electron scattering studies, may, in fact, be important
- ★ While, in general, MC simulation are inherently unable to include interference effects, the distinct energy dependence emerging in the MEC 1p1h sector naturally lends itself to a simplified treatment based on, yes, a modification of the nucleon form factors.

Thank you!

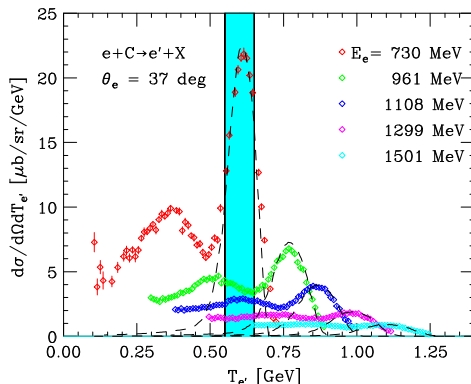
## Backup slides

# PROCESSES INVOLVING TWO-NUCLEON CURRENTS



## THE TROUBLE WITH FLUX AVERAGE

- ★ In neutrino-nucleus interactions, e.g. ,  $\nu_\mu + A \rightarrow \mu^- + X$ , the beam energy is unknown, and so is the energy transfer . As a consequence, different reaction mechanisms contribute to the cross section at fixed muon energy and emission angle
- ★ This problem clearly emerges from the analysis of electron-scattering data corresponding to different beam energies

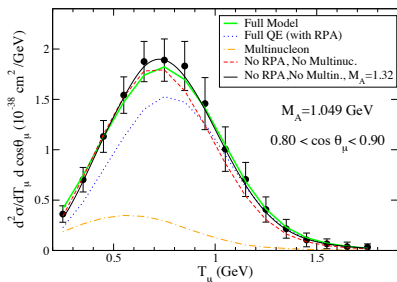




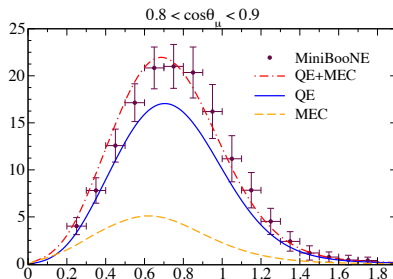
# IDENTIFICATION OF DIFFERENT MECHANISMS

- ★ While involving somewhat different assumptions, several models agree in predicting that the MiniBooNE data can be explained taking into account the contribution of processes involving two-nucleon currents (MEC), associated with **2p2h** final states

## ► Valencia model: Nieves *et al*



## ► Superscaling: Megias *et al*



- ★ Assessing the role of the **2p2h** sector requires an accurate description of the dominant single-nucleon knock out process