



Rescattering in cascade models with RDWIA inputs

Alexis Nikolakopoulos

Workshop on neutrino event generators

Fermilab 15 March 2023

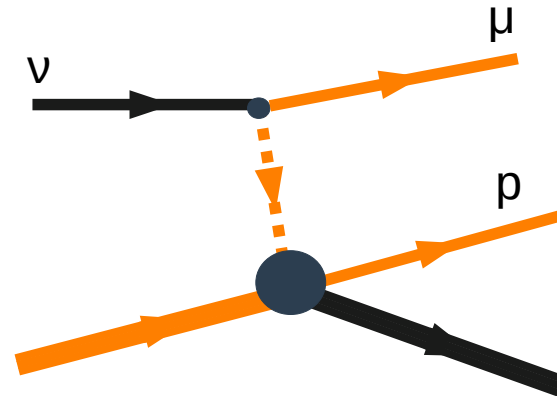
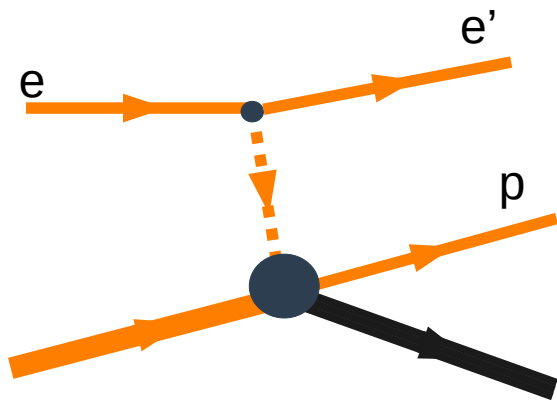
Outline

- 1 How FSI in exclusive one-nucleon knockout is described:
Optical potential approach
- 2 Comparison of the optical potential and the NEUT cascade
- 3 What is the input to the cascade ? :
How does the generator decide on hadrons from inclusive cross section?

Electron and neutrino scattering: Similar in theory, not in experiment

$$\mathcal{M} = j_{lep,\mu} \langle \Psi_f | \mathcal{O}^\mu | \Psi_i \rangle$$

$\langle \Psi_f |$ $\langle \Psi_i |$ \longrightarrow Strong interaction: nuclear many-body states
The 'same physics' in $A(e,e' p) X$ and $A(\nu, \mu p) X$



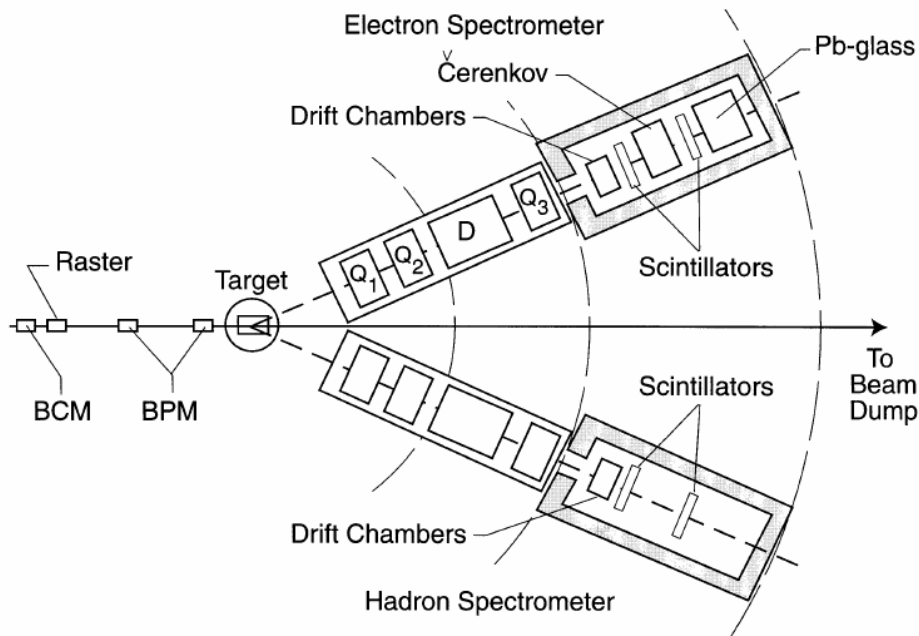
Experimental conditions are different!

$A(e,e' p) X \longrightarrow$ (semi)-exclusive

$A(\nu, \mu p) X \longrightarrow$ (semi)-inclusive

Electron and neutrino scattering: Similar in theory not in experiment

$\langle \Psi_f |$ $\langle \Psi_i |$ \longrightarrow Strong interaction: nuclear many-body states
The 'same theory' in $A(e,e' p) B$ and $A(\nu, \mu p) B$



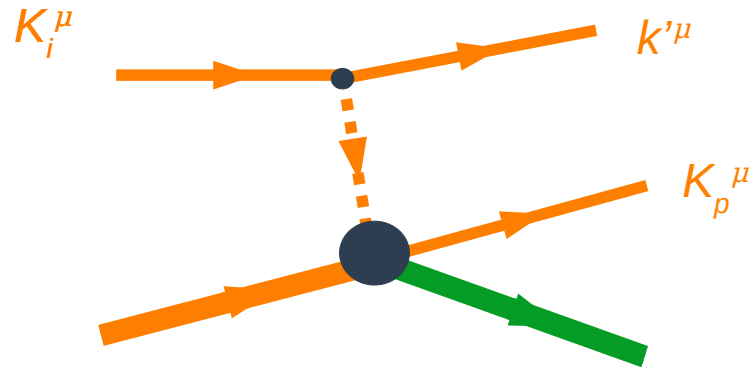
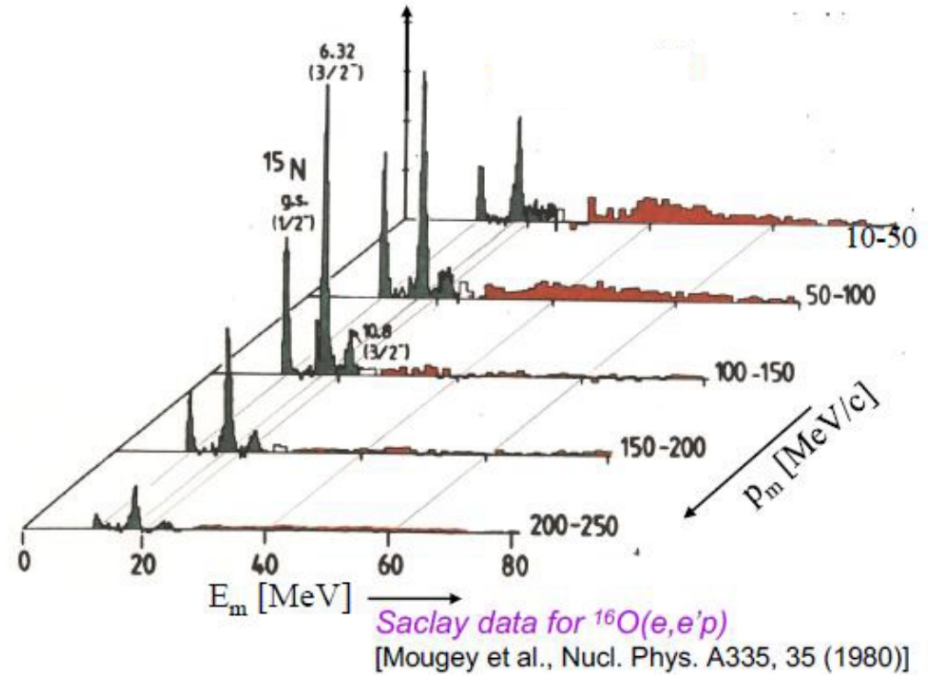
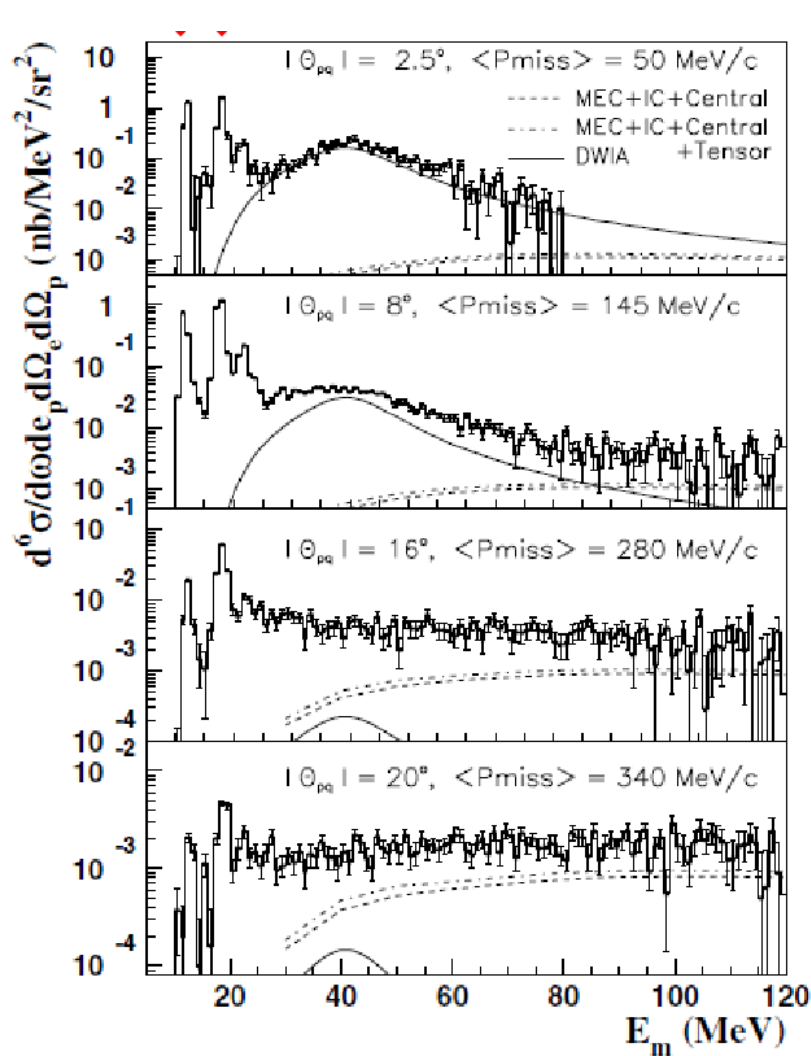
Very different experiments

- Fixed kinematics of e' , p
- Known incoming energy!

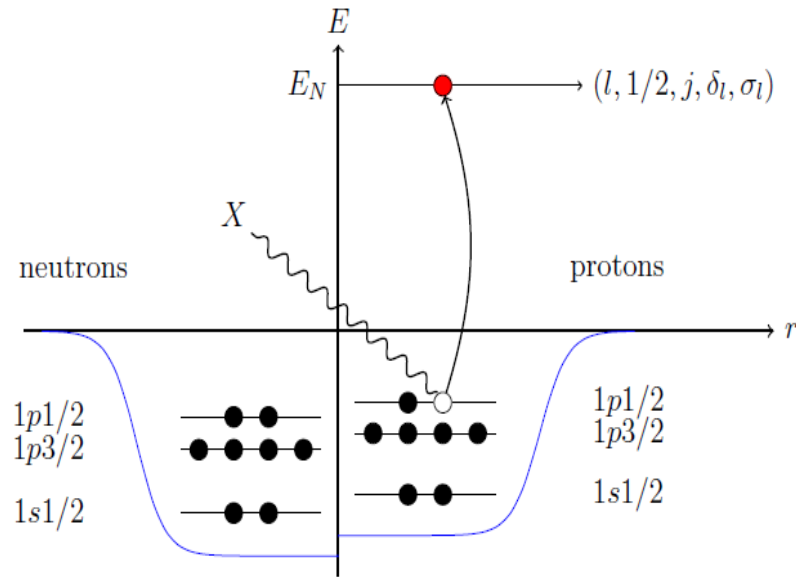
Experimental conditions determine
the theoretical tools for $\langle \Psi_f |$

JLAB Hall A [PRC 42 38]

Exclusive electron scattering: Missing energy distributions



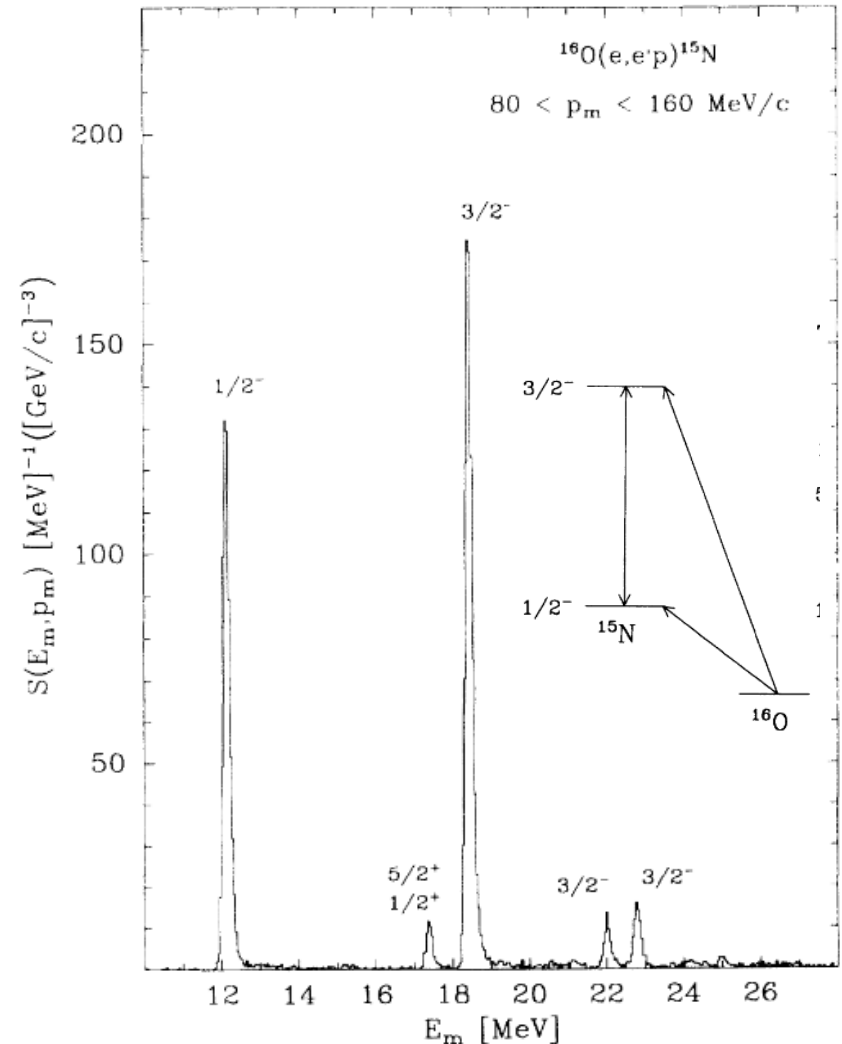
Exclusive electron scattering



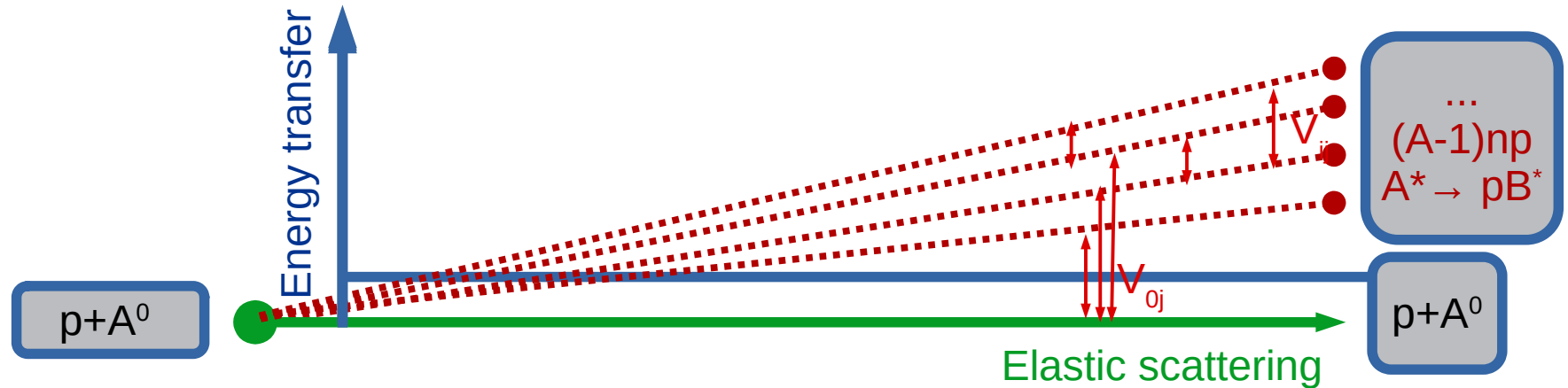
$$\mathcal{M} = j_{lep,\mu} \langle \Psi_f | \mathcal{O}^\mu | \Psi_i \rangle$$

$$\langle \Psi_f | = \langle \phi_p | \langle {}^{15}\text{N}^* |$$

Pure 1-proton knockout from a nuclear shell
How do we describe this state ?



Optical potential approach



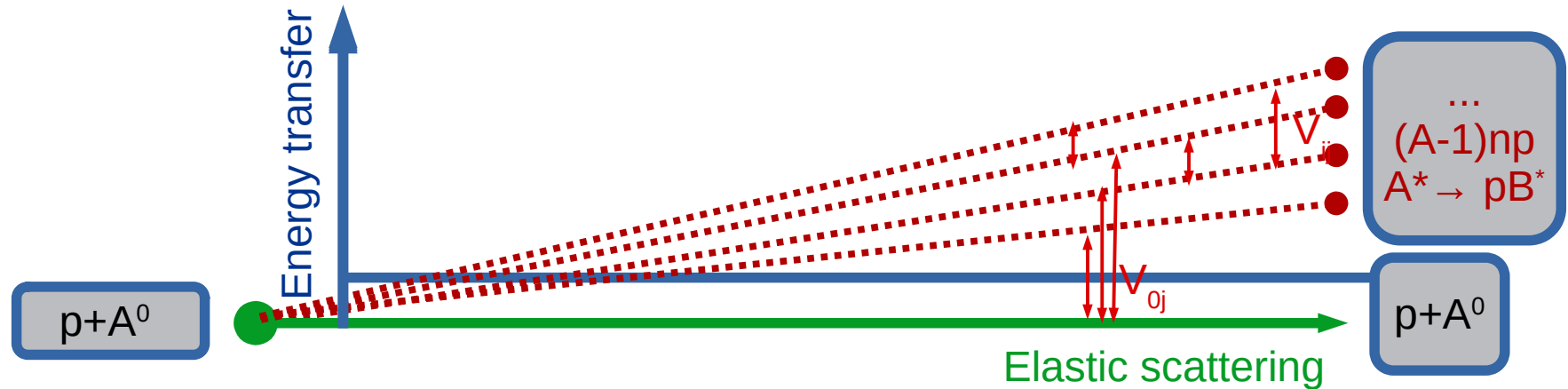
Coupled channels problem, separate out **elastic channel**

$$(H_{00} - E)|\phi_0\rangle = -V_{0j}|\phi_j\rangle \quad (j > 0) \quad \bullet$$

$$(H_{ij} - E)|\phi_j\rangle = -V_{j0}|\phi_0\rangle \quad (i, j > 0) \quad \bullet$$

$$H_{ij} = H^{free}\delta_{ij} + V_{ij}^{nA} + \epsilon_i\delta_{ij}$$

Optical potential approach



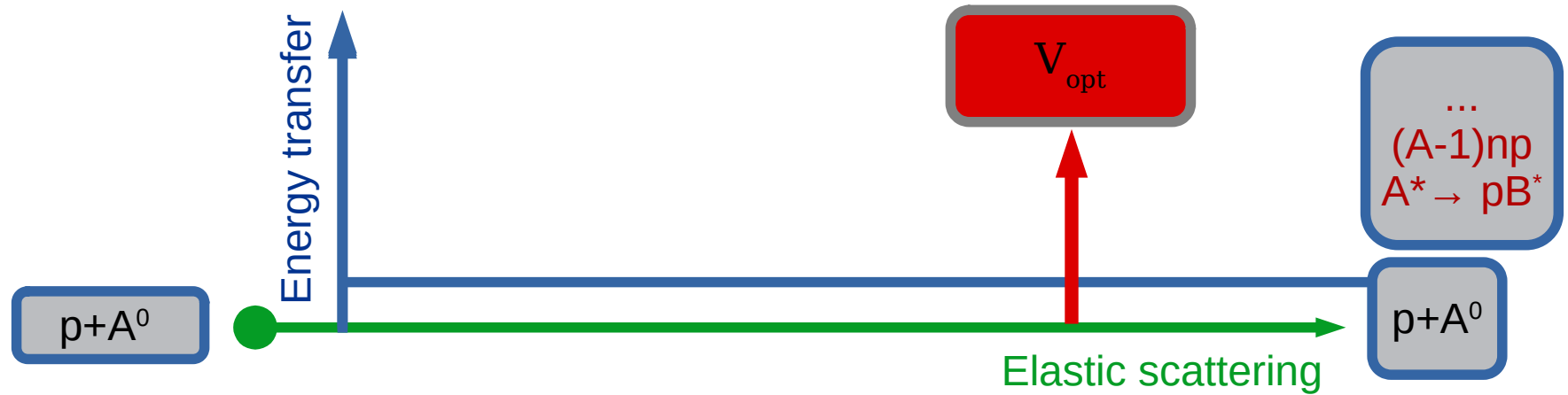
$$(H_{00} - E)|\phi_0\rangle = -V_{0j}|\phi_j\rangle \quad (j > 0) \quad \bullet$$

$$(H_{ij} - E)|\phi_j\rangle = -V_{j0}|\phi_0\rangle \quad (i, j > 0) \quad \bullet$$

$$\left[H^{free} + V_{00}^{nA} + V_{0j} \frac{1}{E - H_{ij} + i\eta} V_{j0} - E \right] |\phi_0\rangle \quad \bullet$$

Coupled channels problem \rightarrow Effective one-body problem as a formal solution

Optical potential approach



$$\left[H^{free} + V_{00}^{nA} + V_{0j} \frac{1}{E - H_{ij} + i\eta} V_{j0} - E \right] |\phi_0\rangle \bullet$$

$$\approx \left[H^{free} + \mathcal{V}^{opt} - E \right] |\phi_0\rangle$$

Coupled channels problem \rightarrow Effective one-body problem with optical potential

Global Dirac phenomenology for proton-nucleus elastic scattering

E. D. Cooper, S. Hama, and B. C. Clark

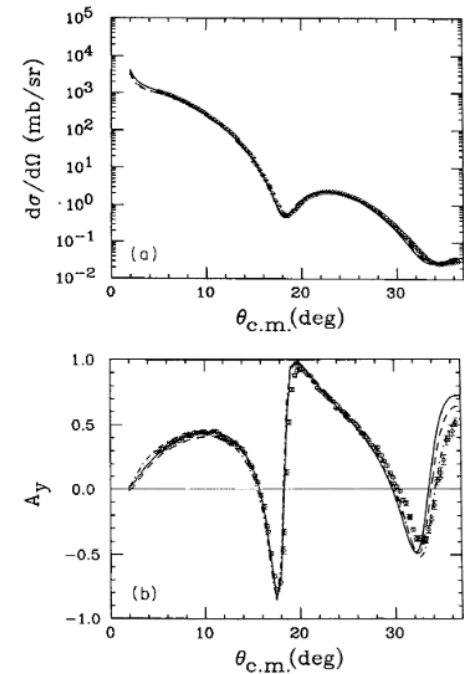
Department of Physics, The Ohio State University, Columbus, Ohio 43210

R. L. Mercer

IBM Thomas J. Watson Research Center, Yorktown Heights, New York 10598

(Received 31 August 1992)

Target	T_p (MeV)	σ_R (mb)				Reference
		EDAI-fit	EDAD-fit			
			fit 1	fit 2	fit 3	
^{12}C	29.00	420.2	435.5	433.1	422.7	[6]
	30.30	415.9	429.0	425.6	414.2	[7]
	49.00	358.8	363.0	348.4	327.7	[6]
	49.48	357.4	361.8	347.0	326.1	[8]
	61.40	323.3	335.6	317.0	294.8	[9]
	65.00	313.5	329.0	309.7	287.4	[10]
	122.00	202.2	269.0	254.4	230.5	[11]
	160.00	177.8	252.3	246.4	215.2	[11]
	200.00	177.6	243.0	243.9	205.0	[11–13]
	300.00	201.1	233.0	235.4	194.9	[14]
	398.00	215.8	227.4	218.6	199.1	[15]
	494.00	227.2	223.7	203.0	211.6	[16]
	797.50	238.4	235.3	209.9	250.0	[17,18]
	1040.00	198.6	259.4	243.8	232.2	[19,20]



Fit to elastic proton-nucleus scattering data

Exclusive electron scattering with Optical potential: 'standard approach'

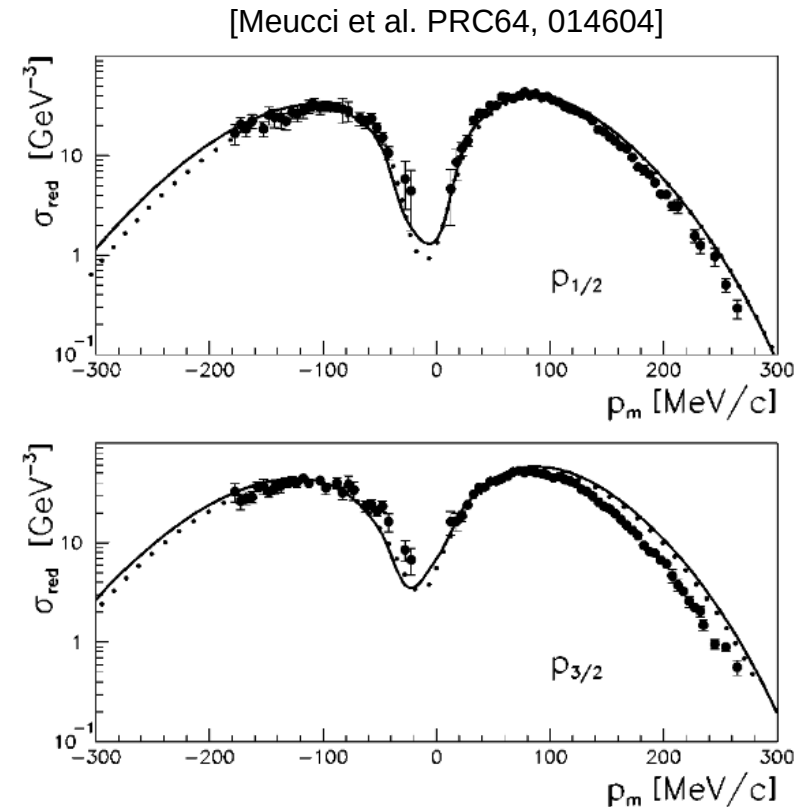
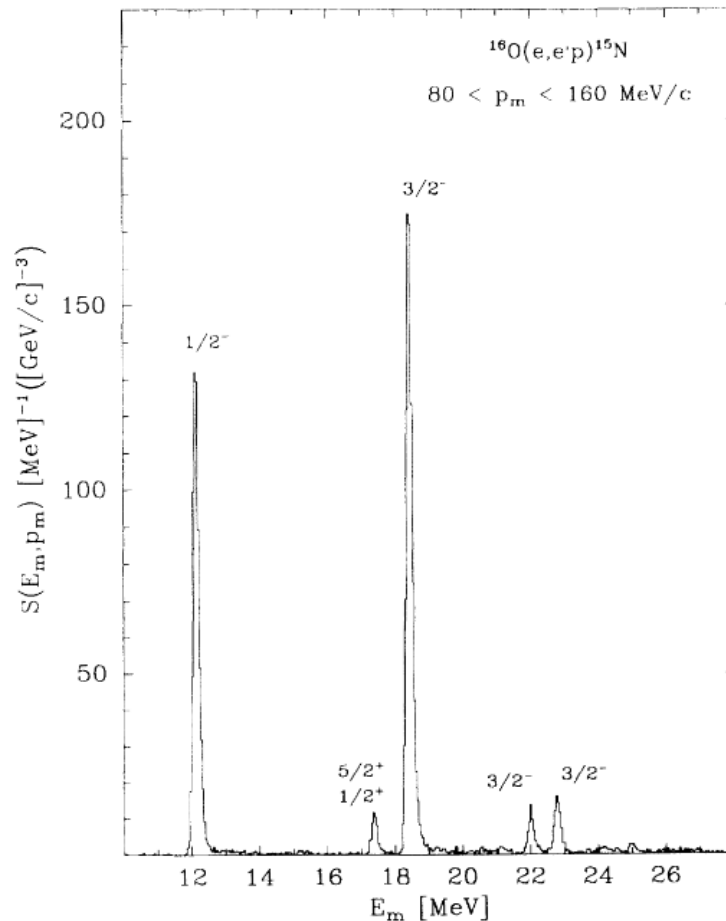
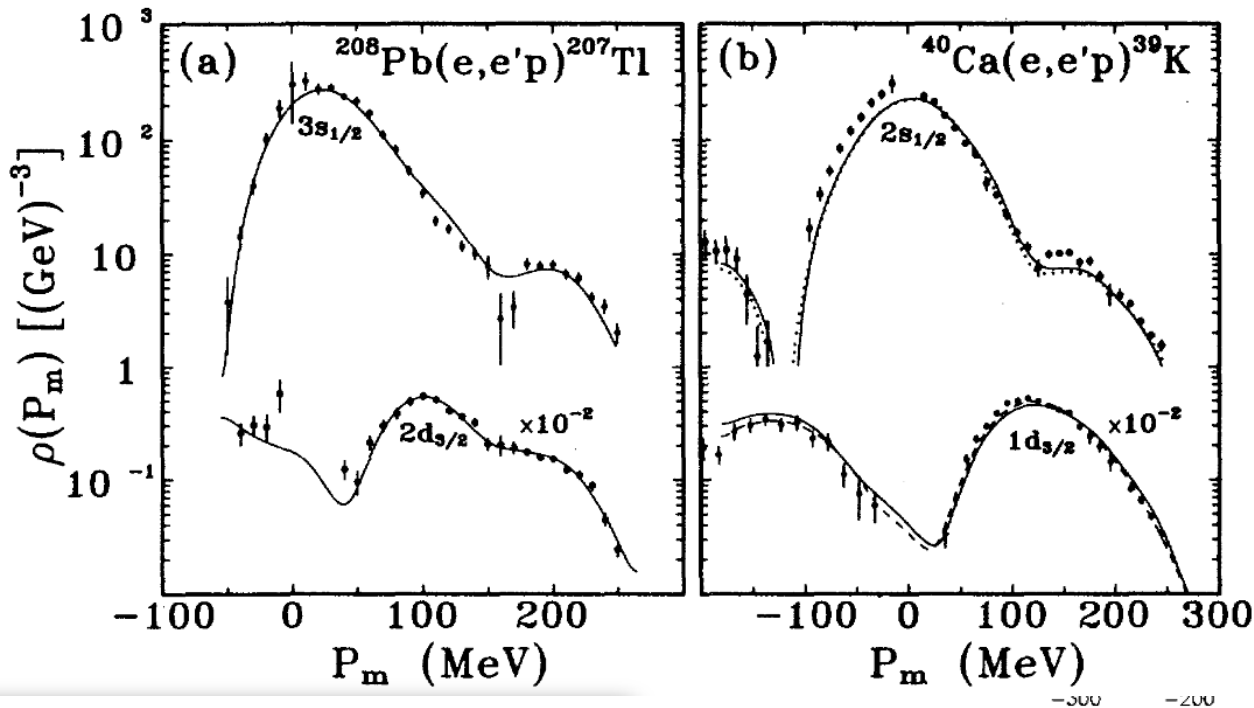


FIG. 11. The reduced cross section (σ_{red}) of the $^{16}\text{O}(e, e'p)$ reaction as a function of the recoil momentum p_m for the transitions to the $1/2^-$ ground state and to the $3/2^-$ excited state of ^{15}N , in

Exclusive electron scattering with Optical potential: 'standard approach'

[Udias et al. PRC48, 2731]



[Meucci et al. PRC64, 014604]

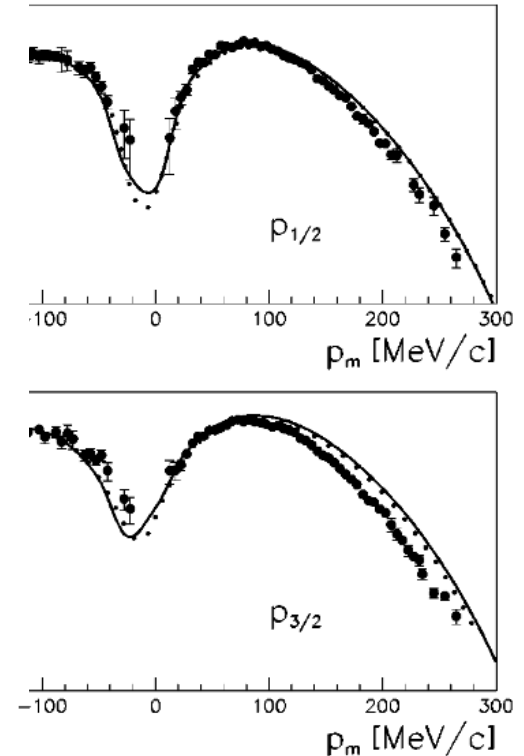


FIG. 11. The reduced cross section (σ_{red}) of the $^{16}\text{O}(e,e'p)$ reaction as a function of the recoil momentum p_m for the transitions to the $1/2^-$ ground state and to the $3/2^-$ excited state of ^{15}N , in

Benchmarking intranuclear cascade models for neutrino scattering with relativistic optical potentials

A. Nikolakopoulos^{1,2,*}, R. González-Jiménez³, N. Jachowicz¹, K. Niewczas^{1,4}, F. Sánchez⁵, and J. M. Udías³

1. Input to NEUT INC

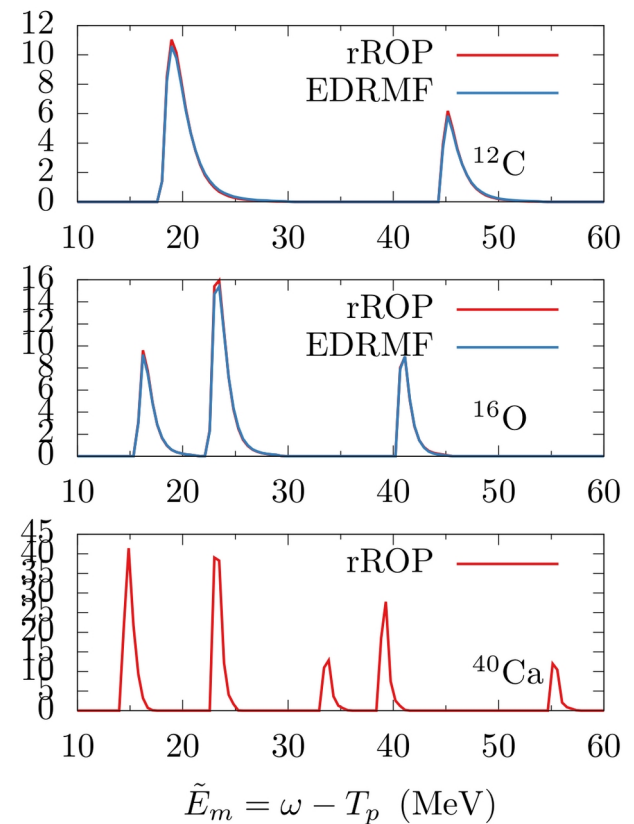
Events according to **unfactorized**
five-fold differential nucleon knockout
cross section

2. Kinematic cuts on NEUT result

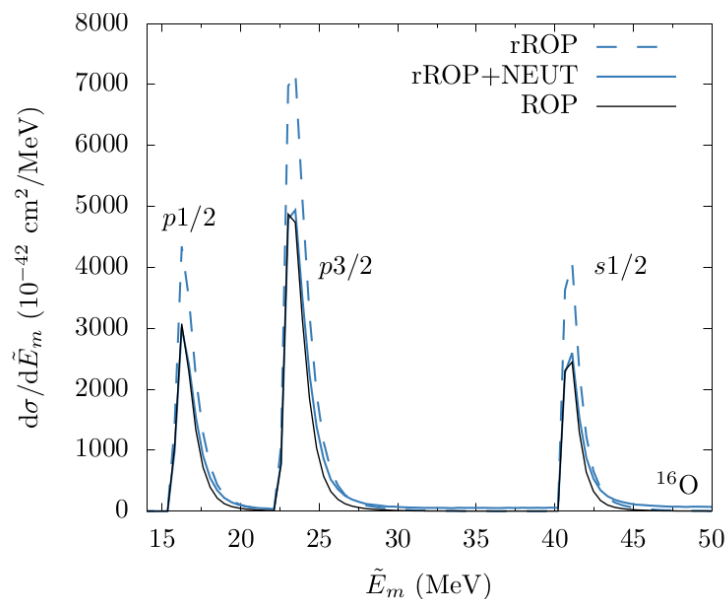
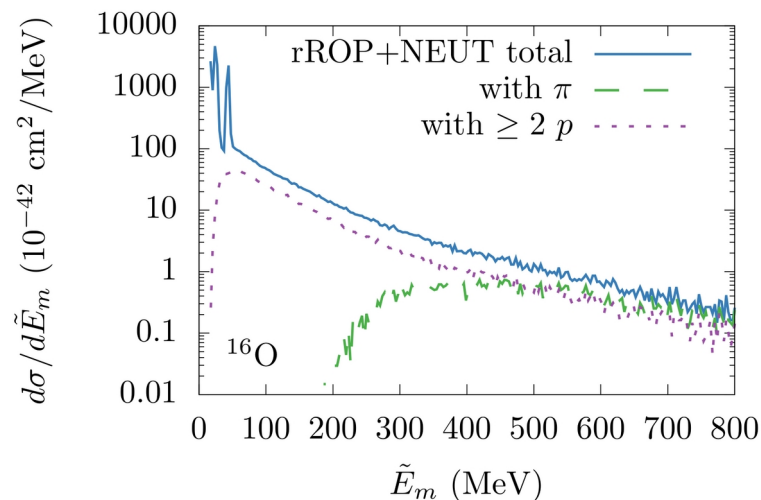
Select events with E_m in shell-model

Is an 'exclusive' signal like in (e,e'p)

Cascade can be compared to ROP



Cascade model with rROP inputs

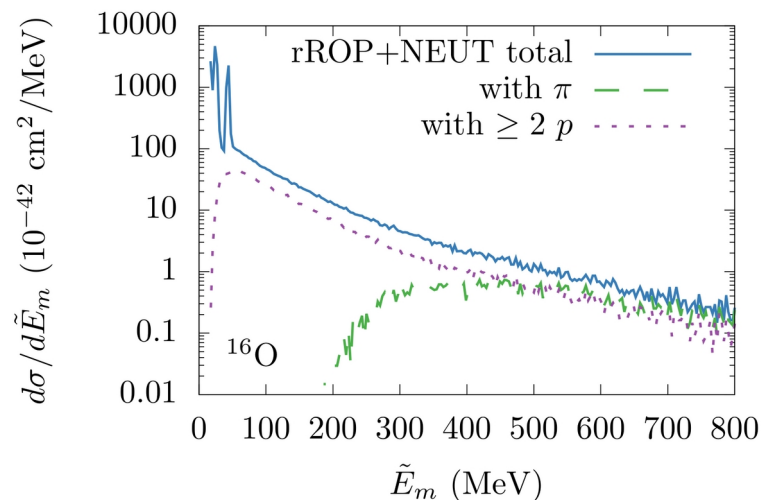


$$\tilde{E}_m = E_i - E_l - T_p$$

T2K flux-folded calculations

- Cascade moves strength from shell model peaks to larger E_m
=> Rescattering into different final states
- Strength of shell model peaks agrees with ROP predictions
- A nice consistent picture emerges between inclusive and exclusive results

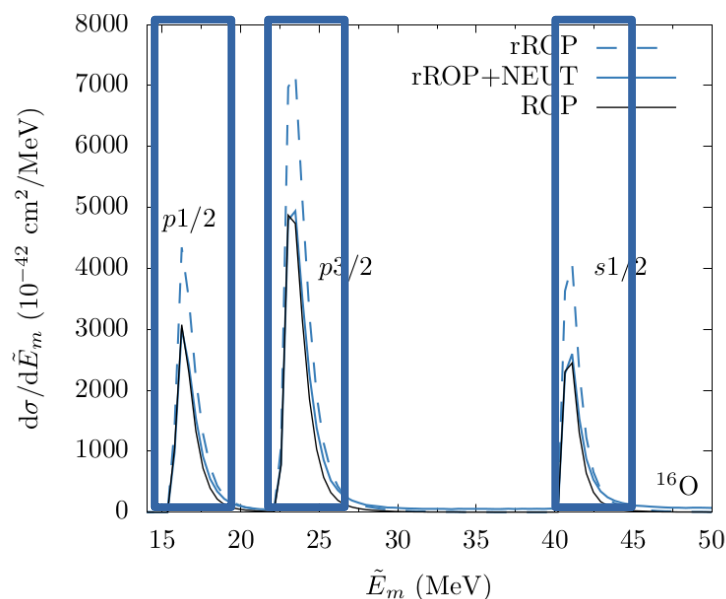
Cascade model with rROP inputs



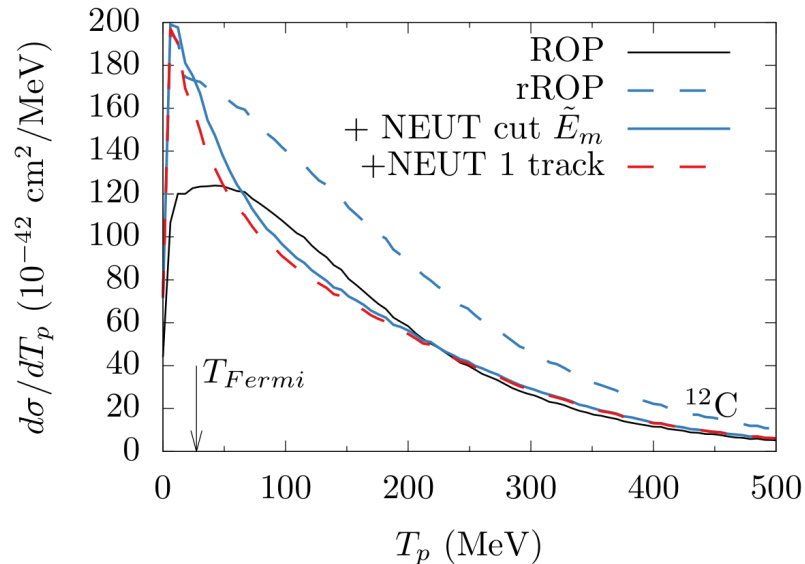
$$\tilde{E}_m = E_i - E_l - T_p$$

T2K flux-folded calculations

- Cascade moves strength from shell model peaks to larger E_m
=> Rescattering into different final states
- Strength of shell model peaks agrees with ROP predictions
- Make kinematic cuts like in (e,e'p)
=> Remove the rescattering
=> exclusive conditions

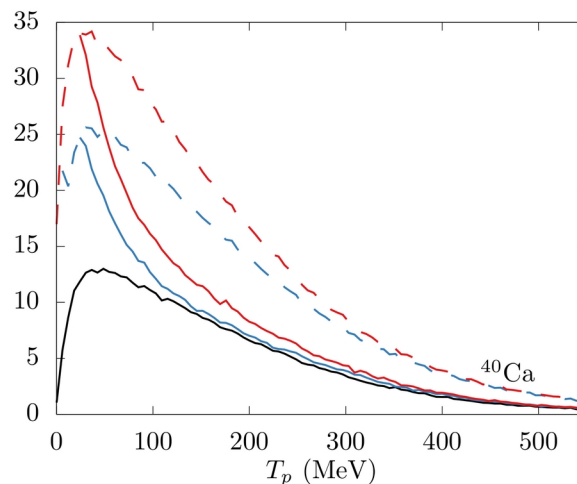
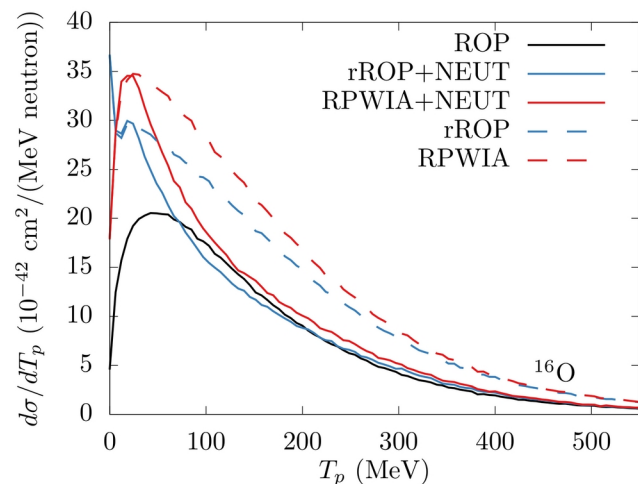
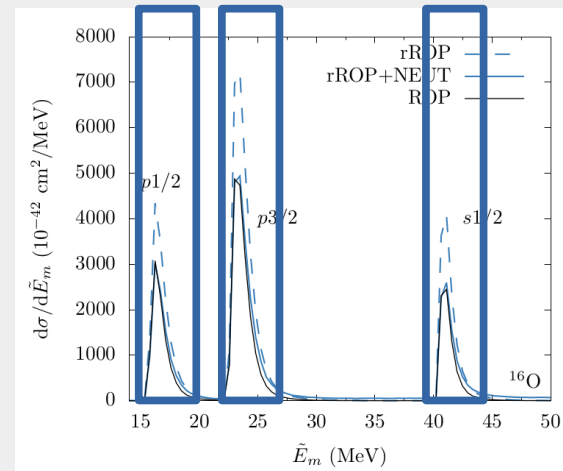


Cascade model with rROP inputs



T2K flux-folded calculations

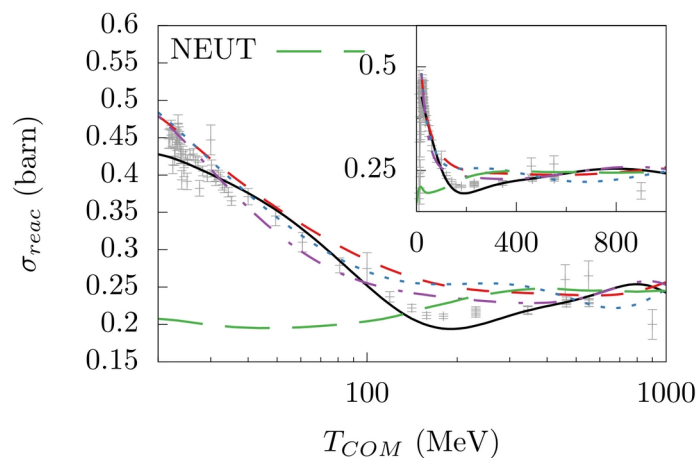
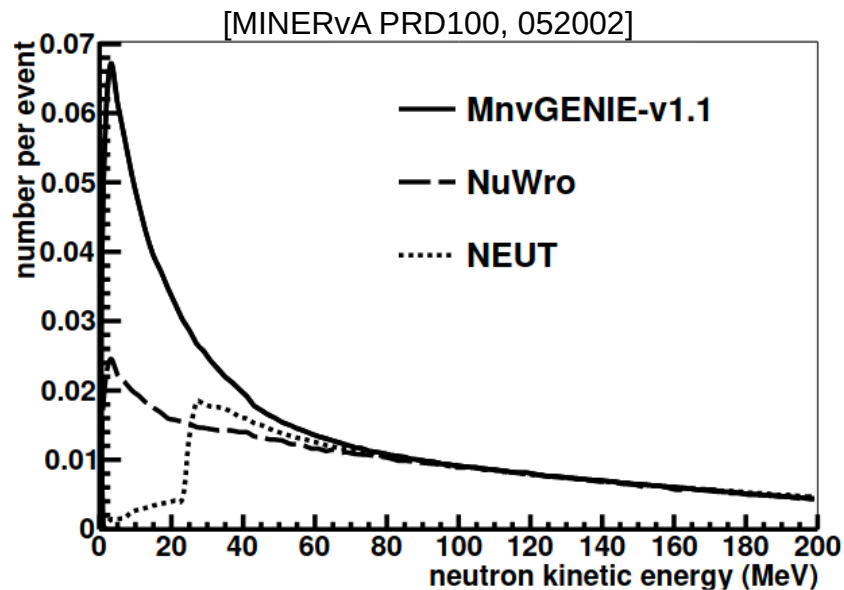
With missing energy cuts



The INC = ROP for high T_p

For $T_p < 100$
differences of up to 100% !

Discrepancies at low- T_N

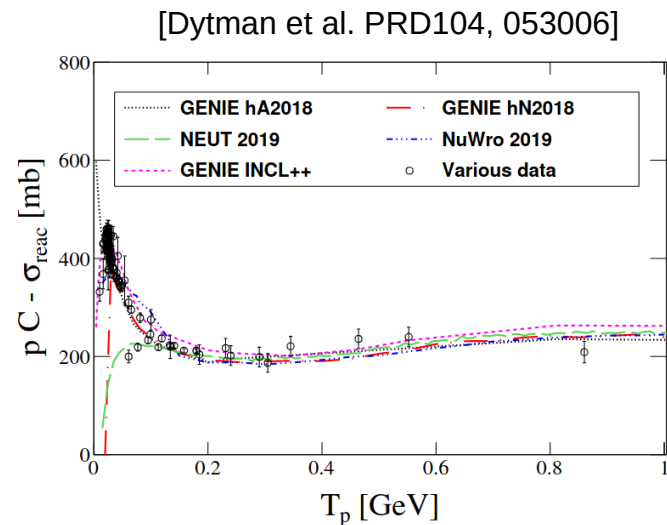


T2K flux-folded calculations

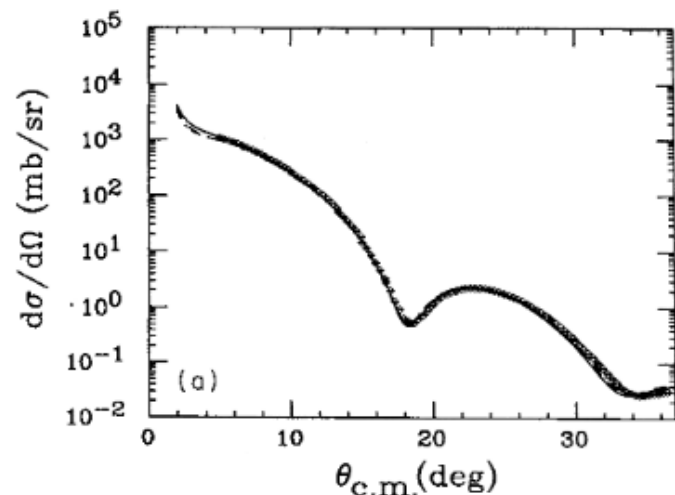
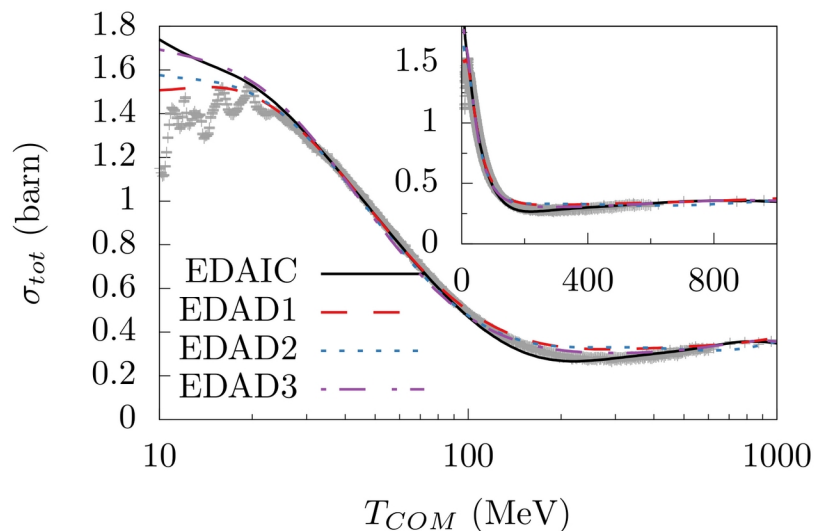
The INC = ROP for high T_p

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Constraints for cascades:
Reaction cross section



What the INC can learn from optical potentials



T2K flux-folded calculations

The INC = ROP for high T_p

For $T_p < 100$
differences of up to 100% !

Constraints for cascades:
Reaction cross section

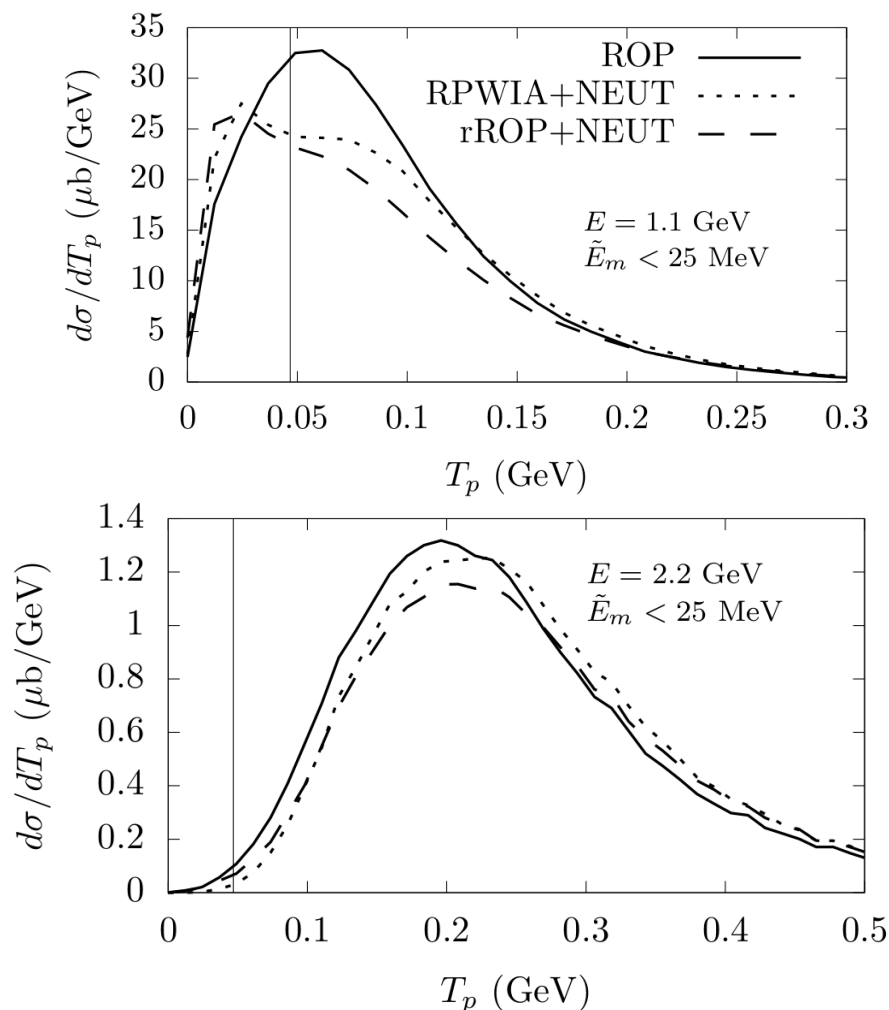
This study:

- Brings the constraints from the analysis of *elastic* and total cross sections
- Gives a **quantum-mechanical** benchmark in the low- T_N region

What the INC can learn from optical potentials

$$\tilde{E}_m < 25 \text{ MeV}$$

→ No 2p2h, RES, inelastic FSI



This study:

- Brings the constraints from the analysis of *elastic* and total cross sections
- Gives a **quantum-mechanical** benchmark in the low- T_N region

Ultimately we should test with data

- Results for e4nu kinematics with a strict E_m cut

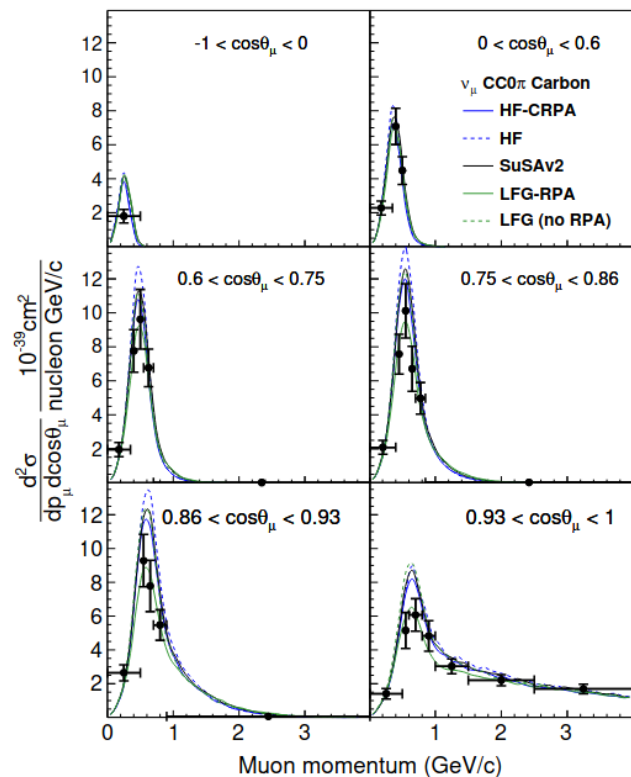
From inclusive to semi-inclusive one-nucleon knockout in neutrino event generators

Alexis Nikolakopoulos^{1,*}, Steven Gardiner¹, Afroditi Papadopoulou², Stephen Dolan³ and Raúl González-Jiménez⁴

[arXiv:2302.12182]

The question:

How do event generators generate hadrons from the **inclusive** cross section ?



Implementation of the SuSAv2-MEC 1p1h and 2p2h models in GENIE and analysis of nuclear effects in T2K measurements

S. Dolan,^{1,2,3} G.D. Megias,^{1,2,4} and S. Bolognesi²

Implementation of the CRPA model in the GENIE event generator and analysis of nuclear effects in low-energy transfer neutrino-nucleus interactions

S. Dolan,¹ A. Nikolakopoulos,² O. Page,³ S. Gardiner,⁴ N. Jachowicz,² and V. Pandey⁵

Implementation of SuSAv2, HF-CRPA, LFG+RPA (Valencia), SuSA-MEC, ... provide the **inclusive** cross section

=> The information on outgoing nucleon is not available

(Similar issues with MEC responses, Resonance production, ...)

From inclusive to semi-inclusive one-nucleon knockout in neutrino event generators

Alexis Nikolakopoulos ^{1,*}, Steven Gardiner ¹, Afroditi Papadopoulou ², Stephen Dolan ³ and Raúl González-Jiménez ⁴

[arXiv:2302.12182]

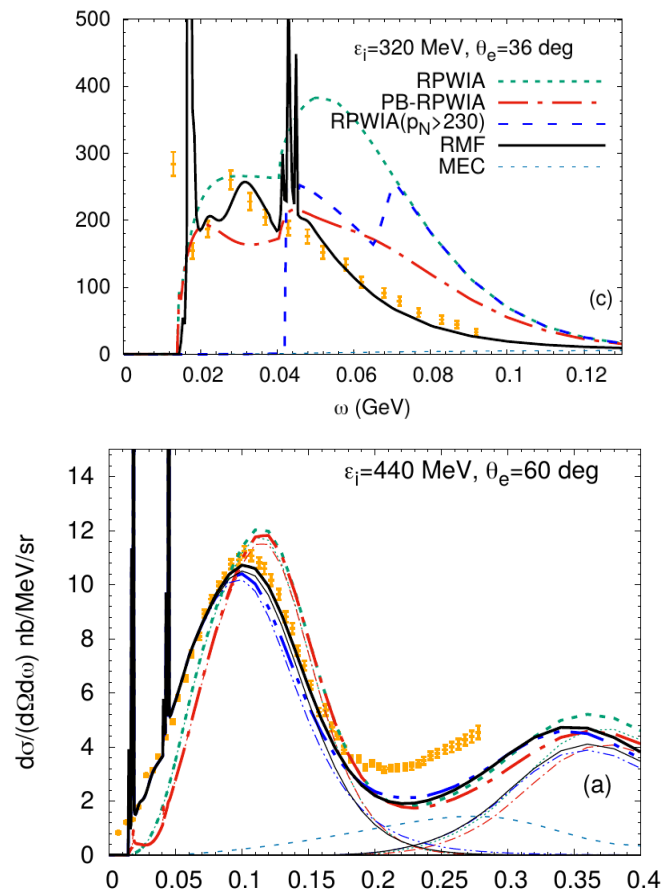
$$P(E_l, \theta_l, T_N, \Omega_N) = \sum_{M_B} \frac{d^4\sigma(E_e, M_B)}{dE_{e'} d\cos\theta_{e'} d\Omega_N}.$$

We generate events for (e,e'p) in RDWIA with **real potential**

- Full consistent description of exclusive kinematics 1e1p

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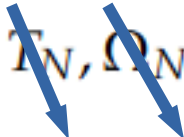
[Phys. Rev. C 100, 045501 (2019)]

We generate events for (e,e'p) in RDWIA with **real potential**

- Full consistent description of exclusive kinematics 1e1p
- Integrate over the proton \rightarrow get the correct inclusive cross section (=includes 'elastic' FSI!)

From inclusive to semi-inclusive one-nucleon knockout in neutrino event generators

Alexis Nikolakopoulos ^{1,*}, Steven Gardiner ¹, Afroditi Papadopoulou ², Stephen Dolan ³ and Raúl González-Jiménez ⁴

$$P(E_l, \theta_l, T_N, \Omega_N)$$


Replace by
'factorized approach'

We get the GENIE version
based on **the same** inclusive
cross section!

We generate events for (e,e'p) in
RDWIA with **real potential**

- Full consistent description of exclusive kinematics 1e1p
- Integrate over the proton → get the correct inclusive cross section (=includes 'elastic' FSI!)
- For every event we **replace** the nucleon kinematics by the GENIE prediction (SuSAv2 implementation)

Nucleon kinematics from inclusive cross section in GENIE

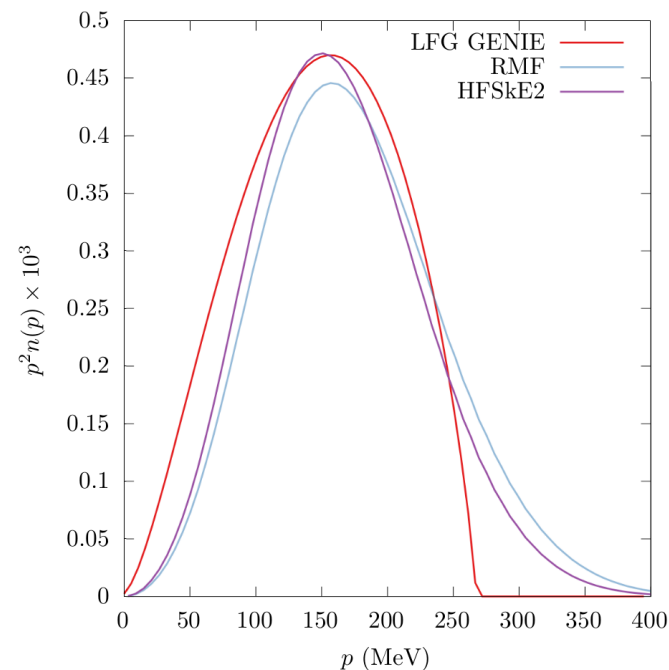
$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

1. Draw initial nucleon \mathbf{p}_m from $p^2 n(p)$ (e.g. LFG)

!! 2. Compute $E_m^2 = p_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$



Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

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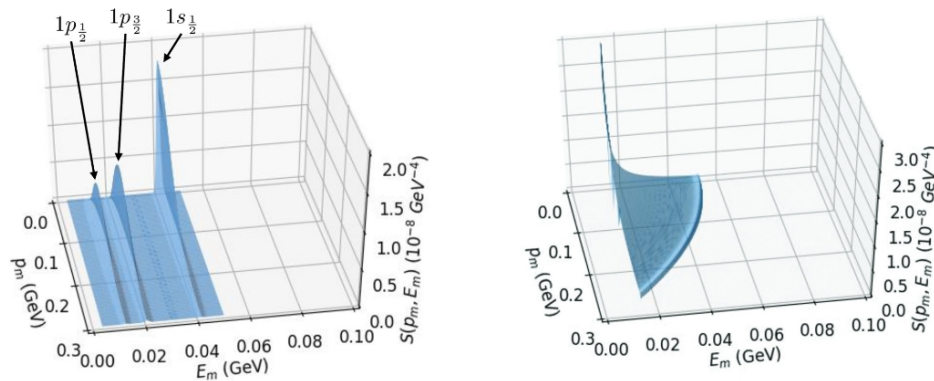
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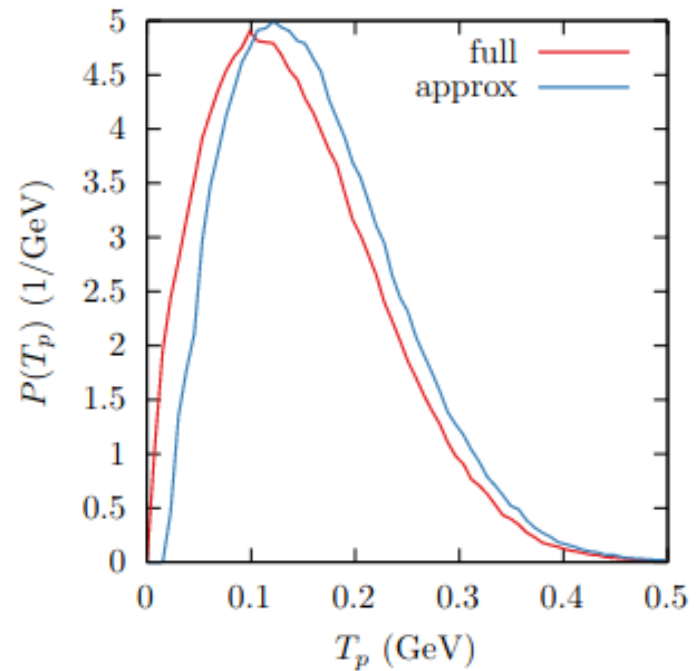
3. $E_N = E_m + \omega - E_b(q)$

Get a shift in total energy spectrum →

LFG $E_m^2 = \mathbf{p}_m^2 + M_N^2$ is not realistic!



[V. Orden & Donnelly PRC 100 044620]



E4nu kinematics $E = 1.15$ GeV

Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information \rightarrow Need to generate it in GENIE

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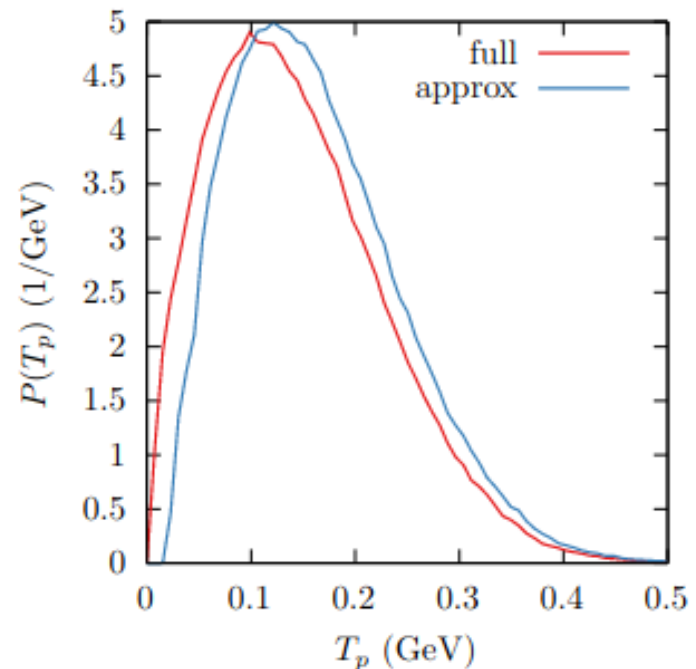
!! 2. Compute $E_m^2 = \mathbf{p}_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$

4. $k_N^2 = E_N^2 - M_N^2$

!! $|\mathbf{p}_m + \mathbf{q}| \neq k_N = \sqrt{E_N^2 - M_N^2}$

$$\rightarrow \mathbf{k}_N = \frac{k_N}{|\mathbf{p}_m + \mathbf{q}|} (\mathbf{p}_m + \mathbf{q})$$



Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE

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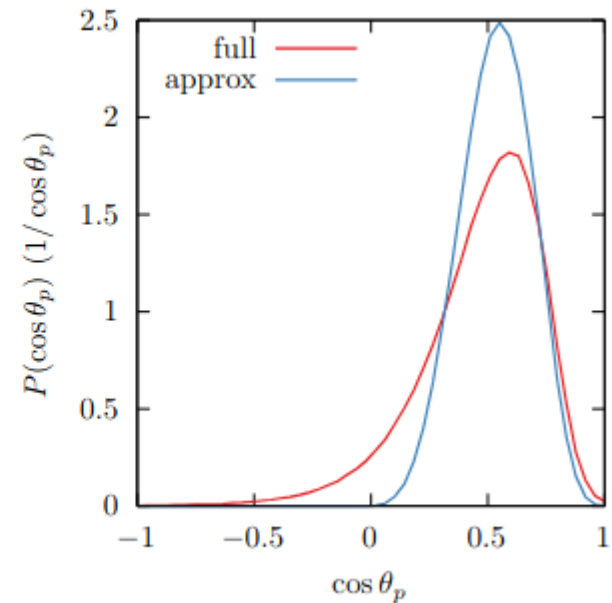
!! 2. Compute $E_m^2 = \mathbf{p}_m^2 + M_N^2$

3. $E_N = E_m + \omega - E_b(q)$

4. $k_N^2 = E_N^2 - M_N^2$

!! $|\mathbf{p}_m + \mathbf{q}| \neq k_N = \sqrt{E_N^2 - M_N^2}$

→ $\mathbf{k}_N = \frac{k_N}{|\mathbf{p}_m + \mathbf{q}|} (\mathbf{p}_m + \mathbf{q})$

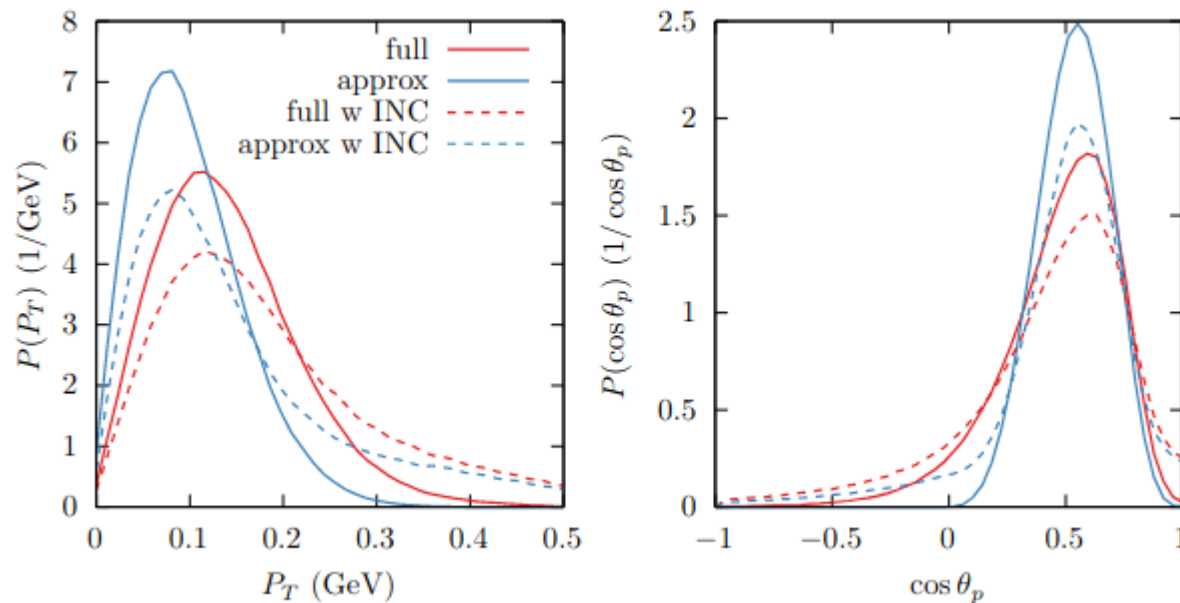


Serious differences in angular distributions!

Nucleon kinematics from inclusive cross section in GENIE

$$\frac{d\sigma(E_\nu)}{dE_l d\cos\theta_l} = G^2 \frac{k_l}{E_\nu} L_{\mu\nu} \int d\Omega_N \sum_{n,\kappa} H_{n,\kappa}^{\mu\nu}(\omega, q, \Omega_N, E_{n,\kappa})$$

Lost nucleon information → Need to generate it in GENIE



Results for e4nu kinematics $E=1.159$ including the GENIE cascade!

Shape differences biggest in P_T and angular distributions

Conclusions

- The ROP and INC approaches use nucleon-nucleus scattering to constrain FSI, in different ways.
- A consistent comparison of the NEUT INC and optical potential shows that there is quantitative agreement at large kinetic energies. For small kinetic energy the differences are up to 100% !!
- The ROP should be more reliable in this comparison, but the true answer is unknown! Should measure $(e,e'p)$ over large phase space with cut on missing energy
- Results of the generator will depend crucially on the input to the INC!
- Current implementations (necessarily) use unrealistic approximations
- **Unfactorized Events** for flux-averaged signals over the whole phase space can be generated combined with INC provides inclusive and exclusive CS
→ **You can use these for validation/error estimation/... of your own INC/simulation/...**