

New Physics models at the precision & intensity frontiers

Stefania Gori
UC Santa Cruz



Beyond the SM from Colliders to the Early Universe workshop
Carena-Wagner fest

Chicago,
May 30, 2023

What's this talk about

- * Physics and memories with Carlos and Marcela since 2009
- * Higgs physics and the connection to dark sectors (precision physics)
- * Dark sector physics at high intensity experiments: DarkQuest

Meeting Carlos and Marcela in 2009

I first met Carlos at the **SUSY** conference in Boston, June 2009
(my first time in the US!).

My first impression: how many questions does he ask during talks??

Then I visited Chicago and met Marcela in November 2009

Then another **SUSY** conference,
Bonn, 2010



I moved to Chicago in September 2010: the start of my “US adventure”:

Professionally incredibly stimulating...

and a lot of fun! These guys know how to party!

Fun memories over the years: Aspen



2018



Partying

Eating/drinking

... and...

Fun memories over the years: Aspen



Partying



| Large Formals | |
|---|------------|
| Vintage Champagne | |
| Veuve Clicquot Cuvée Privée 1982 | \$2,100.00 |
| Veuve Clicquot Cuvée Privée Rosé 1979 | \$2,300.00 |
| Dom Pérignon P2 Vintage 2000 | \$3,300.00 |
| Dom Pérignon P2 Rosé Vintage 1995 | \$1,025.00 |
| Dom Ruart Rosé 2004 | |
| Nebuchadnezzar 15 Liter | |
| Armand de Brignac Ace of Spades Rosé | \$150,000 |
| Armand de Brignac Ace of Spades Brut | \$120,000 |
| Jeroboam 3 Liter | |
| Dom Pérignon Rosé | \$10,500 |
| Dom Pérignon Brut | \$8,000 |
| Perrier-Jouët Belle Époque Rosé | \$4,000 |
| Perrier-Jouët Belle Époque Le Fleur Brut | \$3,500 |
| Veuve Clicquot Yellow Label Brut | \$2,000 |
| Magnum 1.5 Liter | |
| Louis Roederer Cristal Rosé | \$5,000 |
| Perrier-Jouët Belle Époque Le Fleur Rosé | \$5,000 |
| Dom Luminos Rosé | \$3,600 |
| Dom Luminos Brut | \$3,200 |
| Veuve Clicquot Pink Label Rosé | \$2,200 |
| Perrier-Jouët Belle Époque Brut | \$2,000 |
| Veuve Clicquot Yellow Label Brut | \$1,400 |
| RESERVE WATER \$10 Each, or Six for \$50 | |
| Red Bull or Bottled Water \$8 Each, or 5 for \$25 | |
| Bottle Prices Do Not Include Sales Tax. | |
| 20% Service Charge Will Be Added. | |

Eating/drinking



... and...

Fun memories over the years: Aspen



2016

No Carlos here.

We are keeping up the tradition!

Partying

Eating/drinking

... and volleyball

Fun memories over the years: Aspen



2016

No Carlos here.

We are keeping up the tradition!



Papers with Carlos and Marcela

Search for Higgs Bosons in SUSY Cascade Decays and Neutralino Dark Matter,
SG, P. Schwaller, **C.E.M. Wagner**, Phys. Rev. D 83 (2011) 115022

W plus two Jets from a Quasi-Inert Higgs Doublet,
Q.H. Cao, **M. Carena**, SG, A. Menon, P. Schwaller, **C.E.M. Wagner**, L.T. Wang, JHEP 1108 (2011) 002

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W. Altmannshofer, **M. Carena**, SG, A. de la Puente, Phys. Rev. D 84 (2011) 095027

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LHC Discovery Potential for Non-Standard Higgs Bosons in the 3b Channel,
M. Carena, SG, A. Juste, A. Menon, **C.E.M. Wagner**, L.T. Wang, JHEP 1207 (2012) 091

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M. Carena, SG, N.R. Shah, **C.E.M. Wagner**, L.T. Wang, JHEP 1207 (2012) 175

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Find the intruder

W+jets

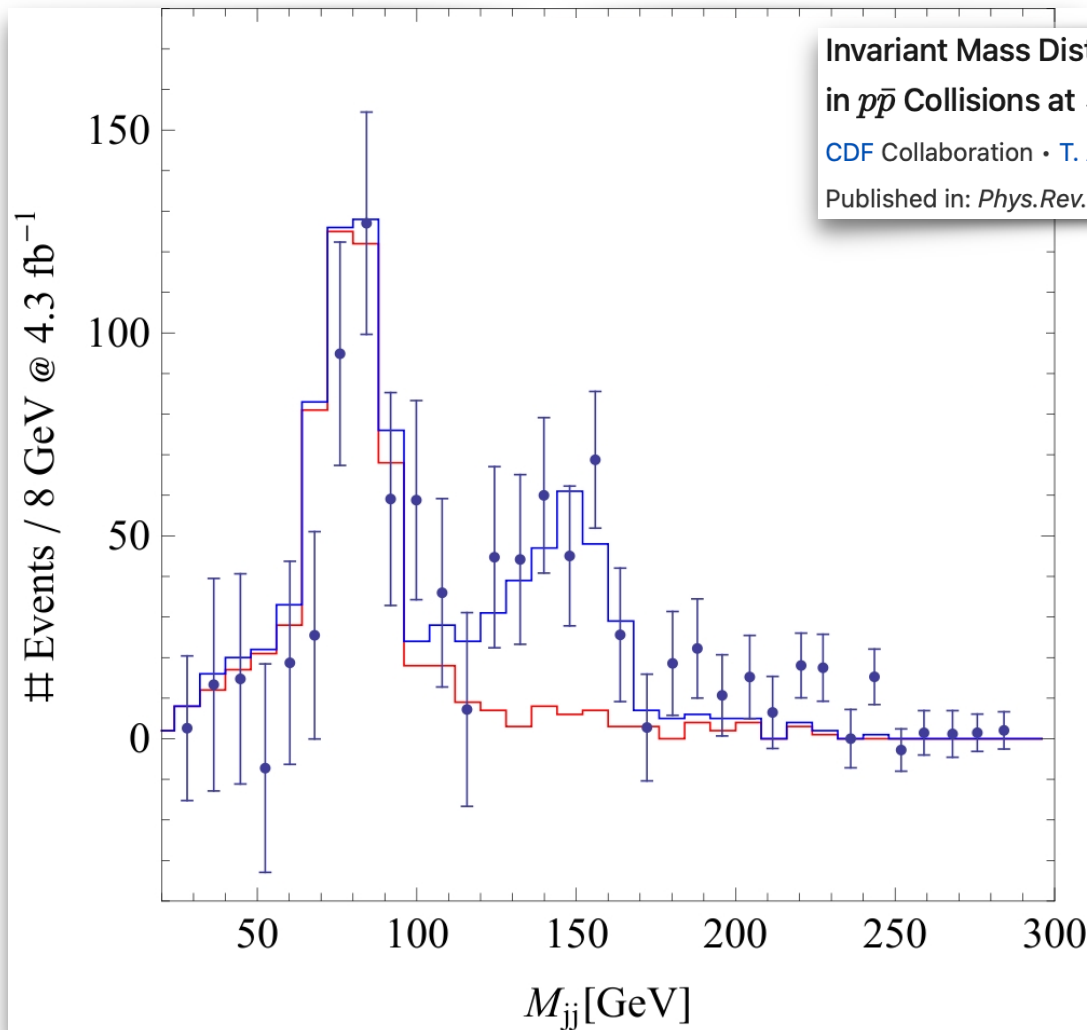
W plus two Jets from a Quasi-Inert Higgs Doublet,

Q.H. Cao, **M. Carena**, SG, A. Menon, P. Schwaller, **C.E.M. Wagner**, L.T. Wang, JHEP 1108 (2011) 002

Invariant Mass Distribution of Jet Pairs Produced in Association with a W boson in $p\bar{p}$ Collisions at $\sqrt{s} = 1.96$ TeV

CDF Collaboration • [T. Aaltonen](#) (Helsinki Inst. of Phys.) et al. (Apr, 2011)

Published in: *Phys.Rev.Lett.* 106 (2011) 171801 • e-Print: [1104.0699](#) [hep-ex]



3.2 σ excess at CDF

our solution:

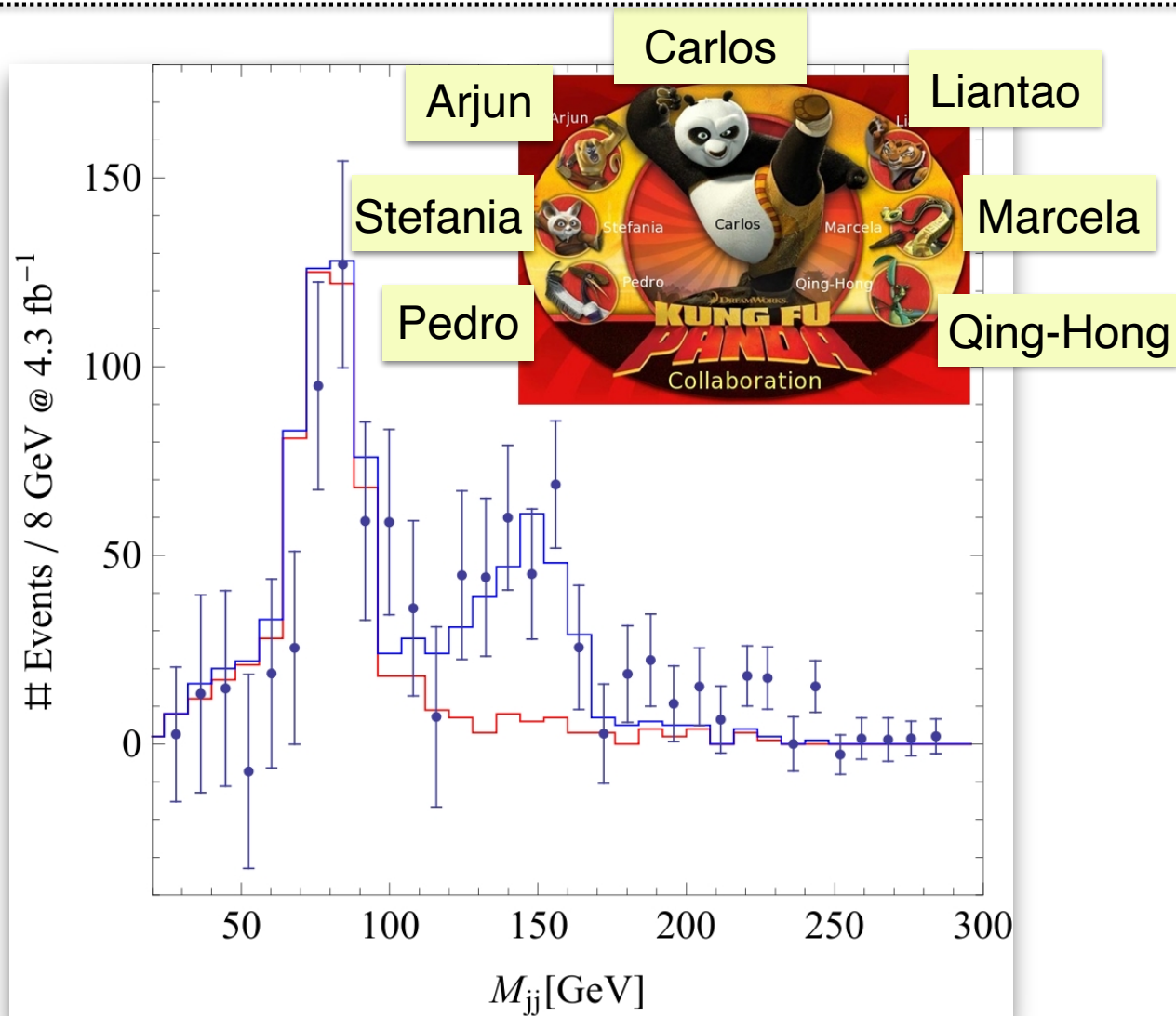
$$p\bar{p} \rightarrow H^\pm \rightarrow W^\pm H^0 / W^\pm A^0 \rightarrow \ell^\pm \nu jj$$

W+jets and the Panda Collaboration



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Q.H. Cao, **M. Carena**, SG, A. Menon, P. Schwaller, **C.E.M. Wagner**, L.T. Wang, JHEP 1108 (2011) 002

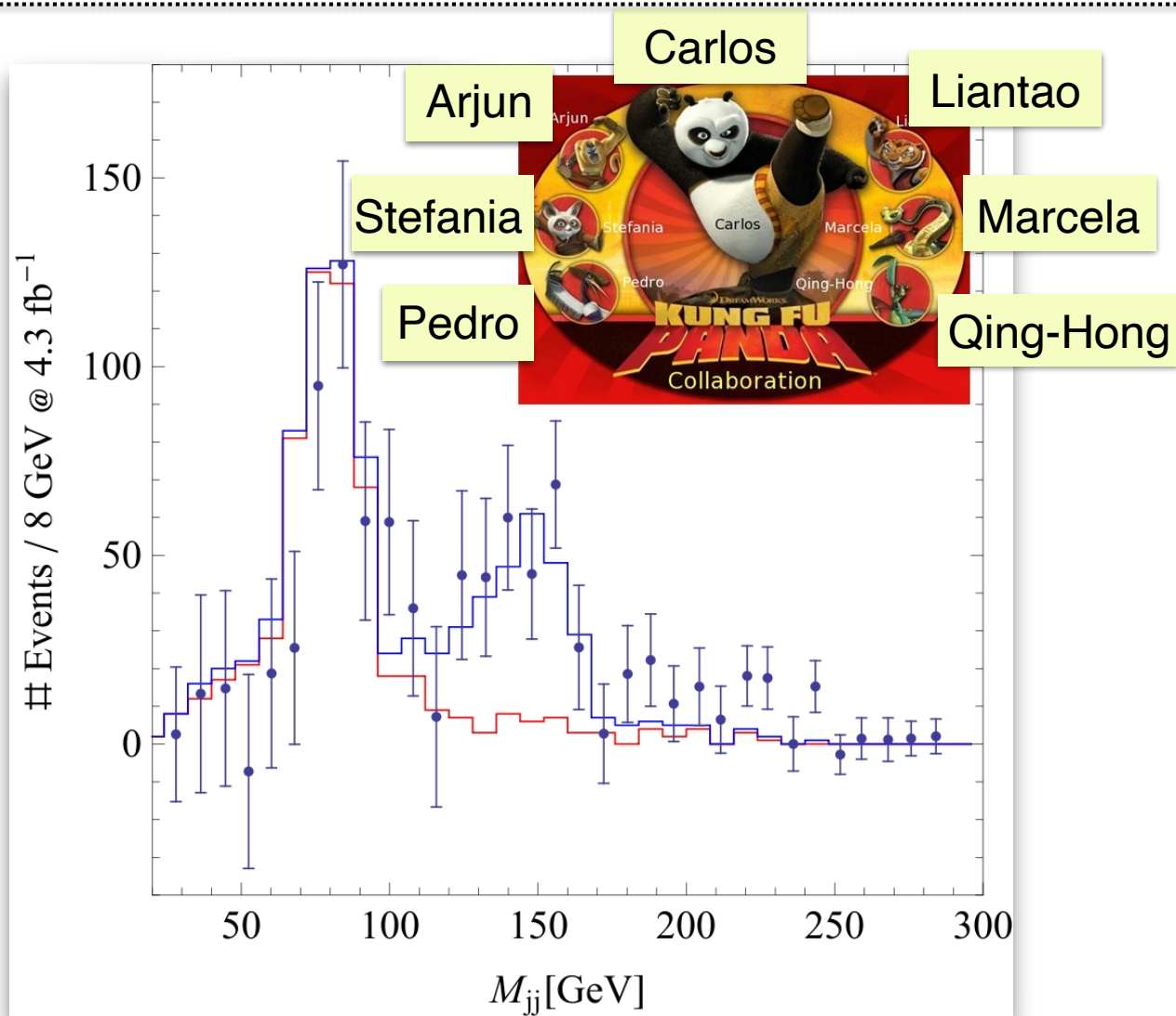


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Higgs pheno before the Higgs discovery

Signals of CP Violation Beyond the MSSM in Higgs and Flavor Physics

Wolfgang Altmannshofer (Fermilab), Marcela Carena (Fermilab and Chicago U., EFI), Stefania Gori (Chicago U., EFI), Alejandro de la Puente (Fermilab and Notre Dame U.) (Jul, 2011)

Published in: *Phys.Rev.D* 84 (2011) 095027 • e-Print: [1107.3814](#) [hep-ph]

MSSM+additional $1/M$ suppressed effective operators that carry CP violation.

Higgs and flavor phenomenology for Tevatron and LHC.

Search for Higgs Bosons in SUSY Cascade Decays and Neutralino Dark Matter

Stefania Gori (Chicago U.), Pedro Schwaller (Illinois U., Chicago and Argonne), Carlos E.M. Wagner (Chicago U. and Argonne and Chicago U., EFI) (Mar, 2011)

Published in: *Phys.Rev.D* 83 (2011) 115022 • e-Print: [1103.4138](#) [hep-ph]

“Due to the large QCD background, searches for such a Higgs, decaying into a pair of bottom quarks, is very challenging at the LHC”.

Study of boosted Higgs production from neutralino and chargino decay and compatibility with DM scenario.

Dreaming of the Higgs

A 125 GeV SM-like Higgs in the MSSM and the gamma-gamma Rate,

M. Carena, SG, N.R. Shah, C.E.M. Wagner, JHEP 1203 (2012) 014

arXiv: 1112.3336

This paper was after the first LHC hints for a Higgs at $\sim(124 - 126)$ GeV with possibly an enhanced gamma gamma rate:

ATLAS-CONF-2011-163

global significance of 2.3σ

local significance of 3.6σ

$H \rightarrow \gamma\gamma$

$H \rightarrow ZZ \rightarrow 4l$

$H \rightarrow WW^* \rightarrow l\nu l\nu$

CMS-PAS-HIG-11-032

global significance of 1.9σ

local significance of 2.6σ

$H \rightarrow \gamma\gamma$

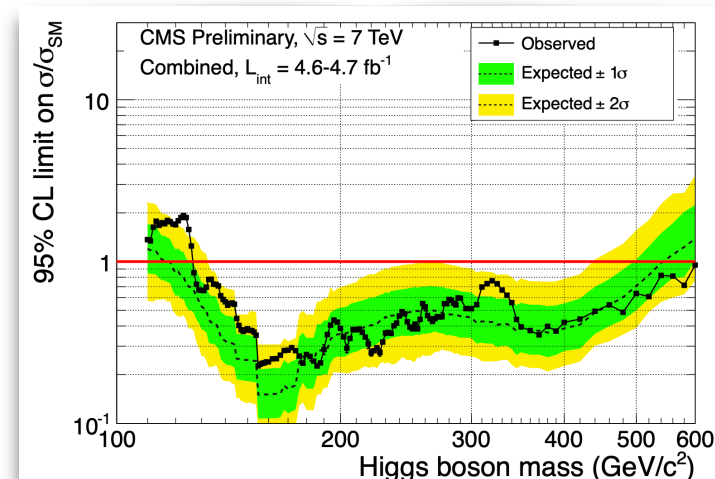
$H \rightarrow ZZ \rightarrow 4l$

$H \rightarrow WW^* \rightarrow l\nu l\nu$

$H \rightarrow b\bar{b}$

$H \rightarrow \tau\tau$

4.9 fb⁻¹
at 7 TeV



(The enhancement of the di-photon rate persisted after the Higgs discovery)

Dreaming of the Higgs

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1. In the MSSM, stops have to lift the mass of the Higgs by ~ 35 GeV.

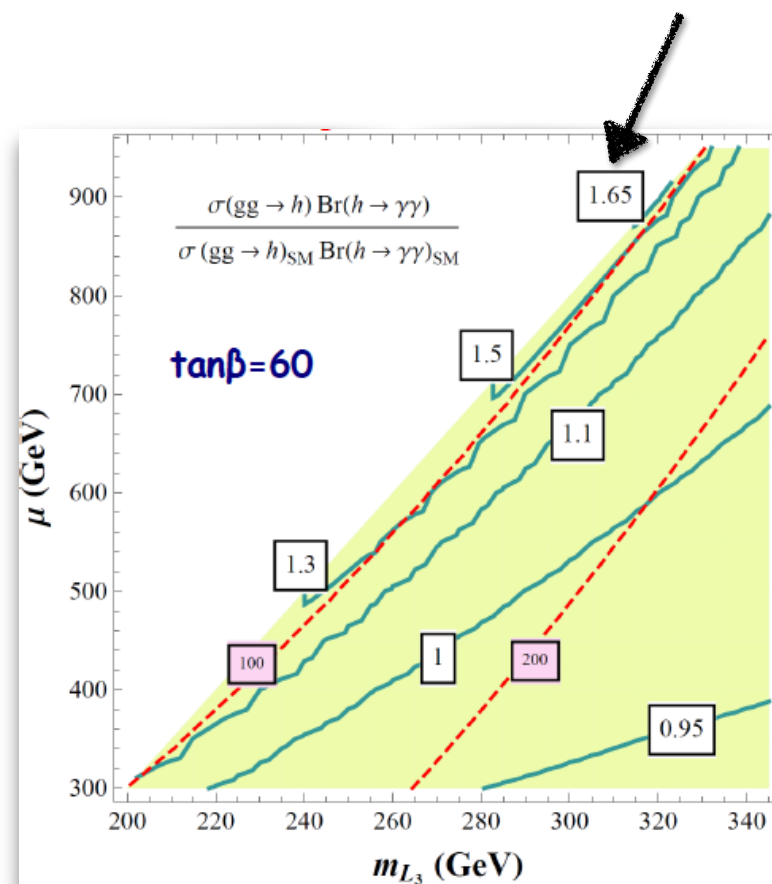
1.1 stops in the several hundred GeV mass range with somewhat large mixing, or **1.2** a large hierarchy between the two stop masses in the case that one of the two stops is light.

2. Light and heavily mixed staus can enhance the di-photon rate

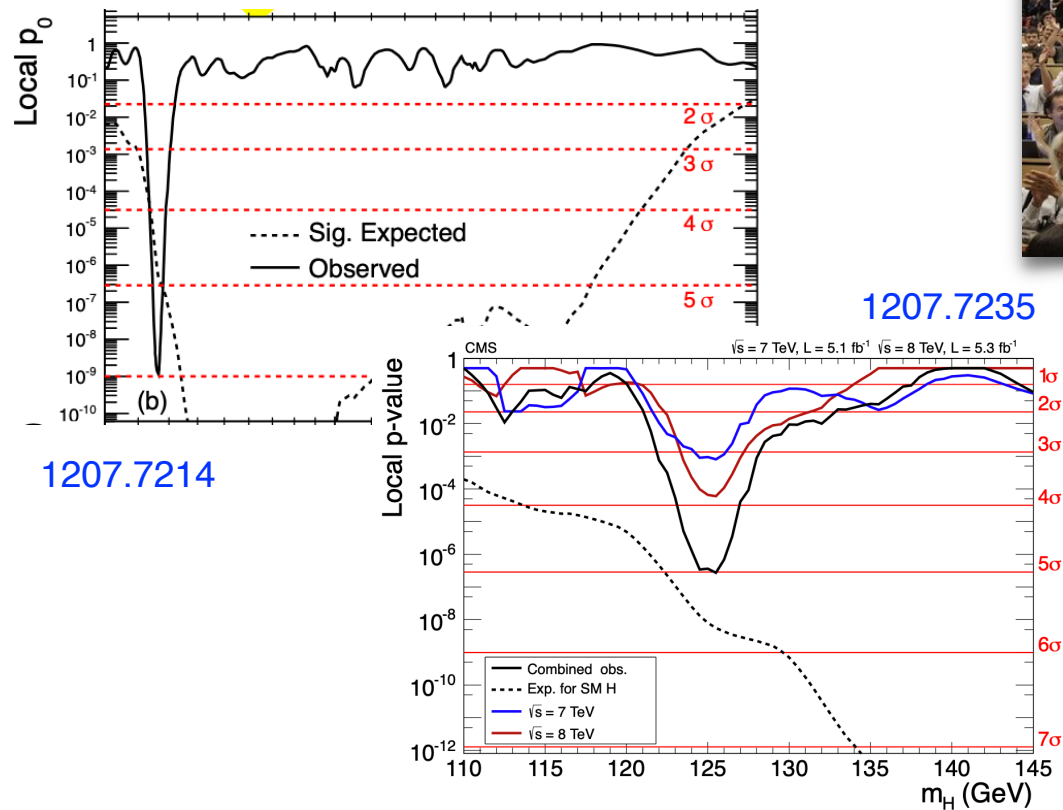
$$\mathcal{M}_{\tilde{\tau}}^2 \simeq \begin{pmatrix} m_{L3}^2 + m_{\tilde{\tau}}^2 + D_L^\tau & m_{\tilde{\tau}}(A_{\tilde{\tau}} - \mu \tan \beta) \\ m_{\tilde{\tau}}(A_{\tilde{\tau}} - \mu \tan \beta) & m_{E3}^2 + m_{\tilde{\tau}}^2 + D_R^\tau \end{pmatrix}$$

$$\Delta A_{\gamma\gamma} \propto - \frac{(\mu \tan \beta)^2 m_{\tilde{\tau}}^2}{m_{L3}^2 m_{E3}^2 - m_{\tilde{\tau}}^2 (\mu \tan \beta)^2}$$

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Discovery!



July 4, 2012
starting at 2am local time in Chicago



Discovery party



Discovery party



The reaction to the Higgs discovery

The year of the Higgs discovery was the most exciting year

Light Stau Phenomenology and the Higgs gamma-gamma Rate,
M. Carena, SG, N.R. Shah, **C.E.M. Wagner**, L.T. Wang, JHEP 1207 (2012) 175

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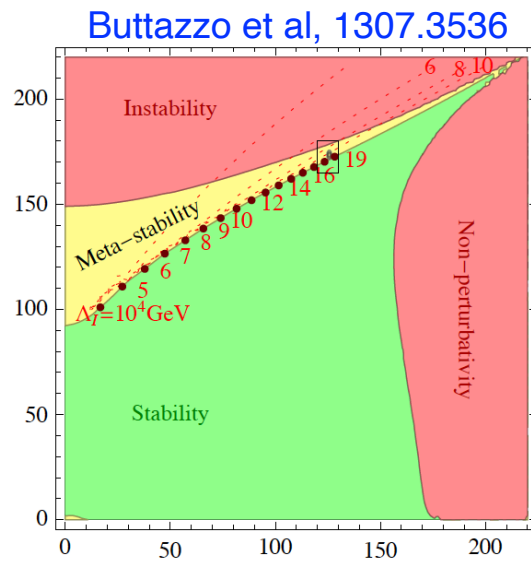
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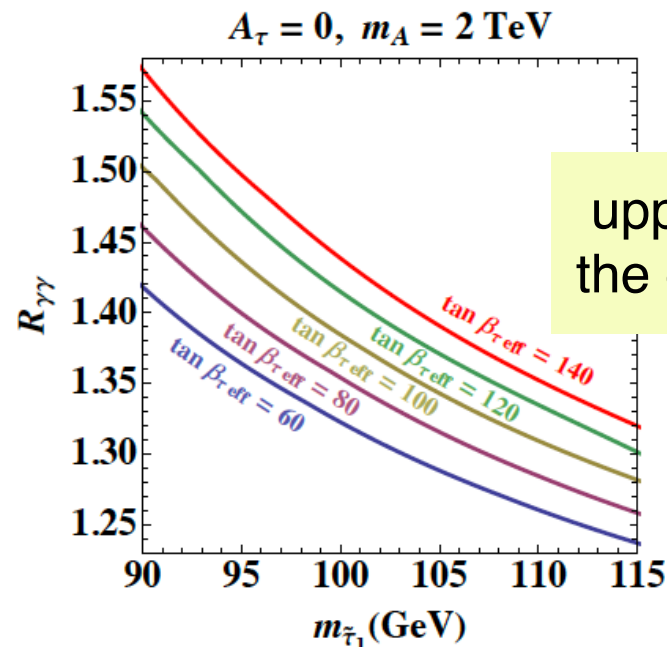
Vacuum Stability and Higgs Diphoton Decays in the MSSM,
M. Carena, SG, I. Low, N.R. Shah, C.E.M. Wagner, JHEP 1302 (2013) 114

The SM has a metastable minimum



Light staus can destabilize the EW minimum.
Appearance of **charged breaking minima**

$$V \supset -2y_\tau \mu \tilde{L} \tilde{\tau} \phi_u + \tilde{L}^2 \tilde{\tau}^2 \left(y_\tau^2 - \frac{g_1^2}{2} \right)$$



$$y_\tau \mu \sim \sqrt{2} \frac{m_\tau}{v} \mu \tan \beta$$

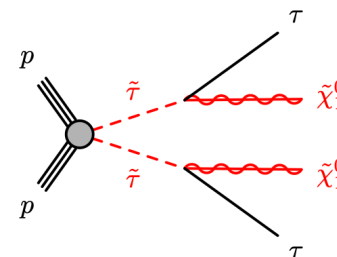
upper bound on
the di-photon rate

Staus in 2023: still pretty hidden

Quoted proposed scenario: $m_{L3} = m_{e3} = 280$ GeV, $\tan \beta = 60$, $\mu = 650$ GeV and $M_1 = 35$ GeV, giving a light stau ~ 95 GeV, a very light LSP, $m_{\tilde{\chi}_1^0} \sim 35$ GeV and a light sneutrino, ~ 270 GeV.

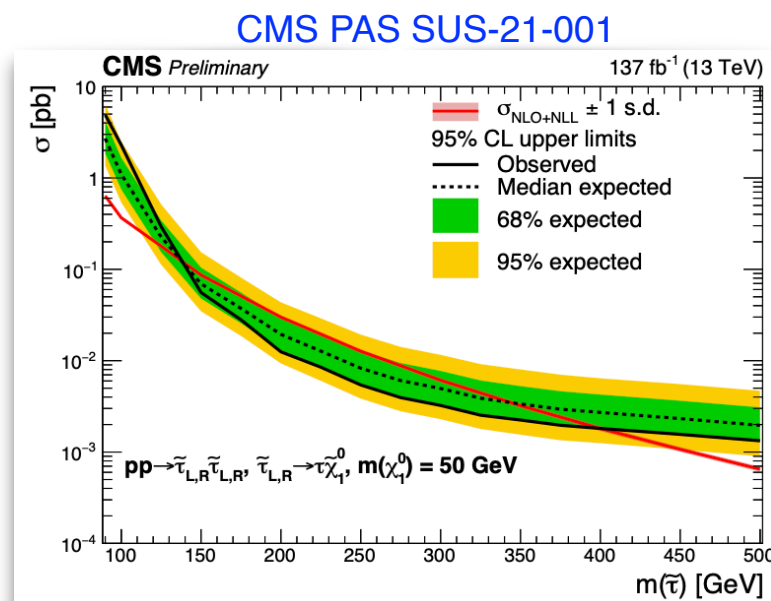
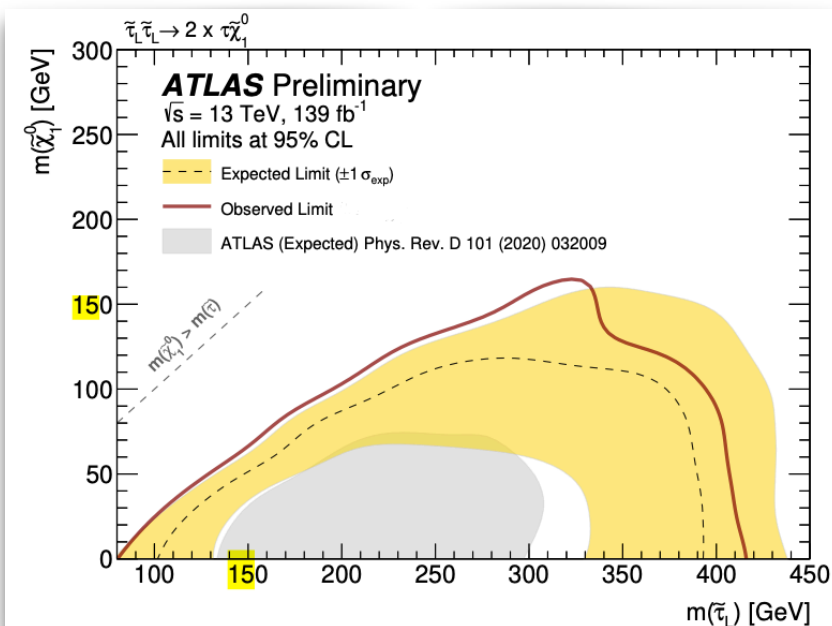
1205.5842

| | Signature | 8 TeV LHC (fb) | 14 TeV LHC (fb) |
|--|-----------------------------|----------------|-----------------|
| $pp \rightarrow \tilde{\tau}_1 \tilde{\tau}_1$ | $2\tau, \cancel{E}_T$ | 55.3 | 124.6 |
| $pp \rightarrow \tilde{\tau}_1 \tilde{\tau}_2$ | $2\tau, Z, \cancel{E}_T$ | 1.0 | 3.2 |
| $pp \rightarrow \tilde{\tau}_2 \tilde{\tau}_2$ | $2\tau, 2Z, \cancel{E}_T$ | 0.15 | 0.6 |
| $pp \rightarrow \tilde{\tau}_1 \tilde{\nu}_\tau$ | $2\tau, W, \cancel{E}_T$ | 14.3 | 38.8 |
| $pp \rightarrow \tilde{\tau}_2 \tilde{\nu}_\tau$ | $2\tau, W, Z, \cancel{E}_T$ | 0.9 | 3.1 |
| $pp \rightarrow \tilde{\nu}_\tau \tilde{\nu}_\tau$ | $2\tau, 2W, \cancel{E}_T$ | 1.6 | 5.3 |



$$pp \rightarrow \tilde{\tau}_1 \tilde{\nu}_\tau \rightarrow \tilde{\tau}_1 (W \tilde{\tau}_1) \rightarrow \tau \chi_1 W \tau \chi_1$$

Our signature has not been looked for yet by ATLAS and CMS





Future Higgs discoveries and the Higgs portal

“The Higgs is SM-like” but...

We need to understand if the Higgs

- * interacts with the 2nd generation



Future Higgs discovery
(for Run III)

- * interacts with itself

- * is CP violating

- * interacts with DM/a dark sector

- * interacts with new Higgs bosons

- * ...

Future Higgs discoveries and the Higgs portal



“The Higgs is SM-like” but...

We need to understand if the Higgs

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Future Higgs discovery
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- * ...



IH^2 is the lowest dimensional gauge and Lorentz invariant combination of SM fields that we can write down



The Higgs easily couples
to any new physics

One of the only 3 possible renormalizable operators that connect the SM to gauge singlet particles (aka **dark sector** particles): $IH^2 IS^2$

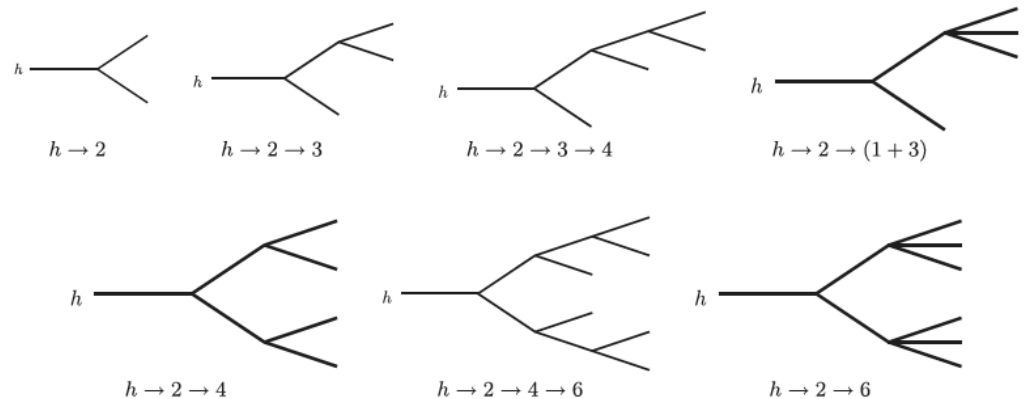
Higgs portal

Higgs exotic decays

Aspen (summer 2012) was also the place where several colleagues (many from the Chicago area) and my self started the effort studying Higgs exotic decays

Exotic Decays of the 125 GeV Higgs Boson
1312.4992,

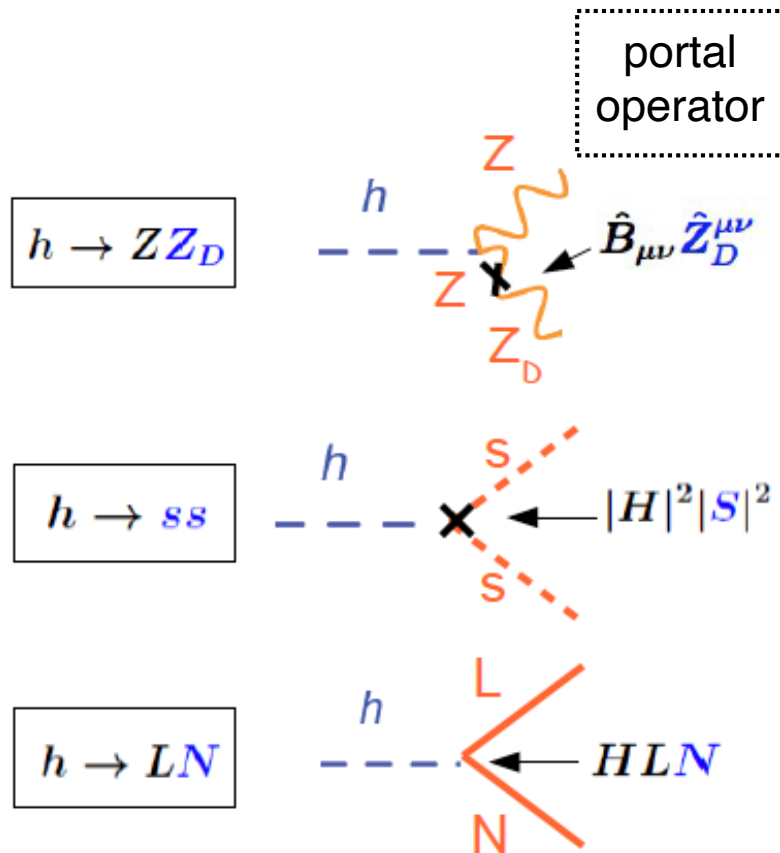
D. Curtin, R. Essig, SG, P. Jaiswal, A. Katz, T. Liu, Z. Liu,
D. McKeen, J. Shelton, M. Strassler, Z. Surujon,
B. Tweedie, Y-M. Zhong



Some inspiration from models like “Dark higgs”,
Draper, Wagner, Wang, Zhang, 1009.3963
(NMSSM with an approximate PQ symmetry)

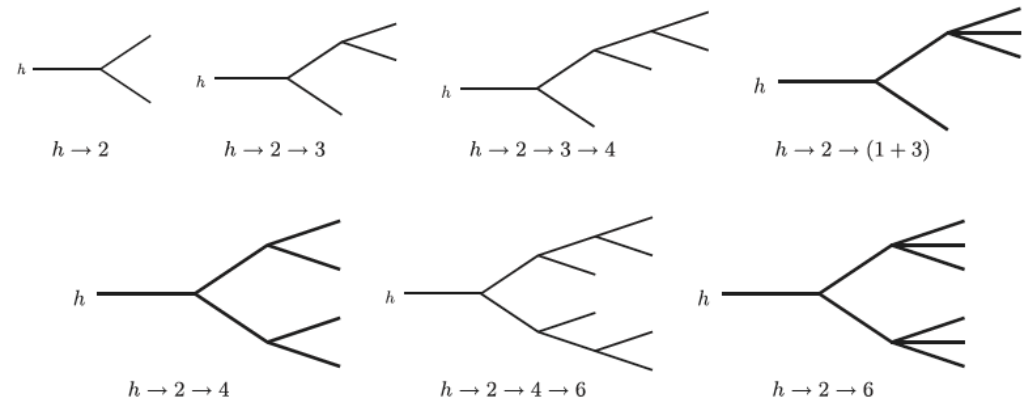
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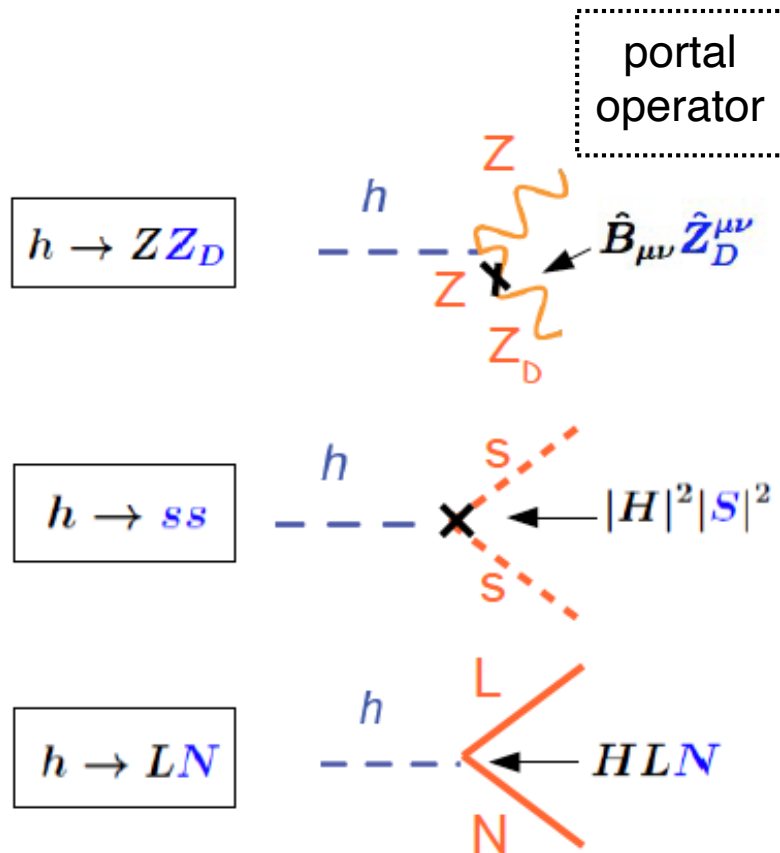
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The LHC can be thought as
a high-intensity machine

- * Now 7 million Higgs
- * HL-LHC 170 million Higgs

This is particularly relevant for
rare processes like Higgs exotic decays

Higgs exotic decays: one of the
best ways to probe
light new particles at the LHC

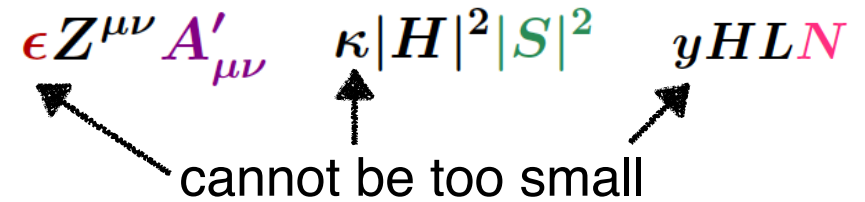
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(NMSSM with an approximate PQ symmetry)

Lower bound on the interaction strength

If we ask for a thermal freeze-out scenario:

$$\epsilon Z^{\mu\nu} A'_{\mu\nu} \quad \kappa |H|^2 |S|^2 \quad y H L N$$

cannot be too small

The diagram shows three interaction terms: $\epsilon Z^{\mu\nu} A'_{\mu\nu}$, $\kappa |H|^2 |S|^2$, and $y H L N$. Three arrows originate from the text "cannot be too small" and point to each of these terms. The first arrow points to ϵ , the second to κ , and the third to y .

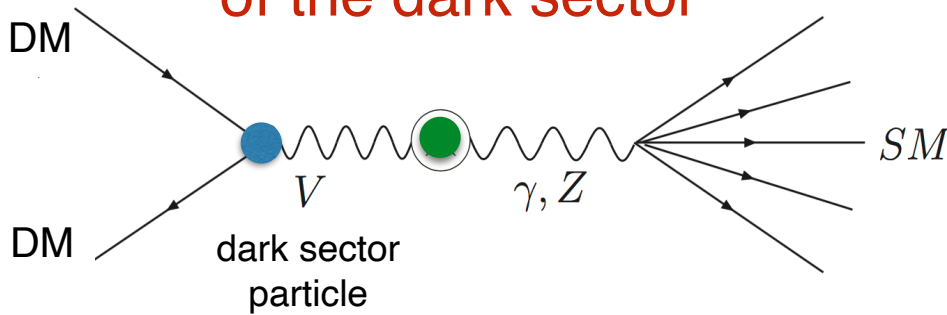
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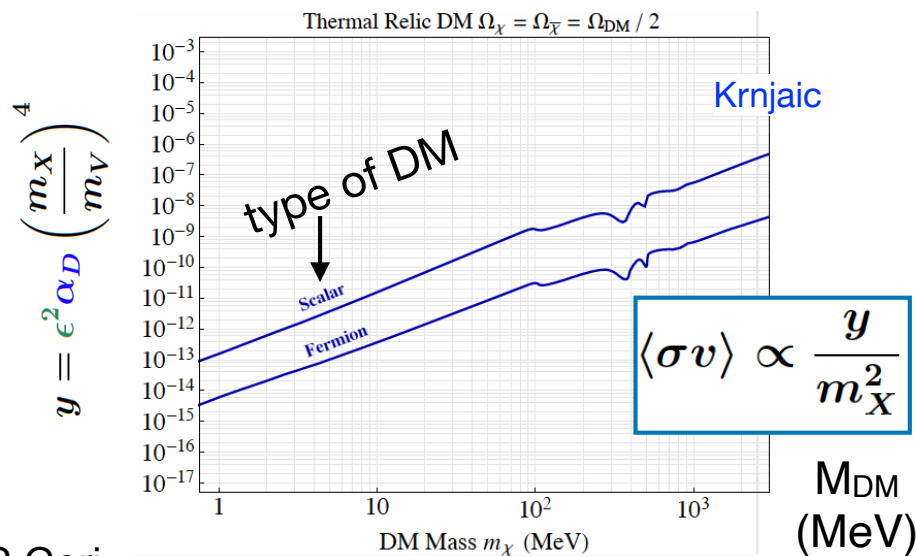
$$\epsilon Z^{\mu\nu} A'_{\mu\nu} \quad \kappa |H|^2 |S|^2 \quad y H L N$$

cannot be too small

1. DM is the lightest state of the dark sector



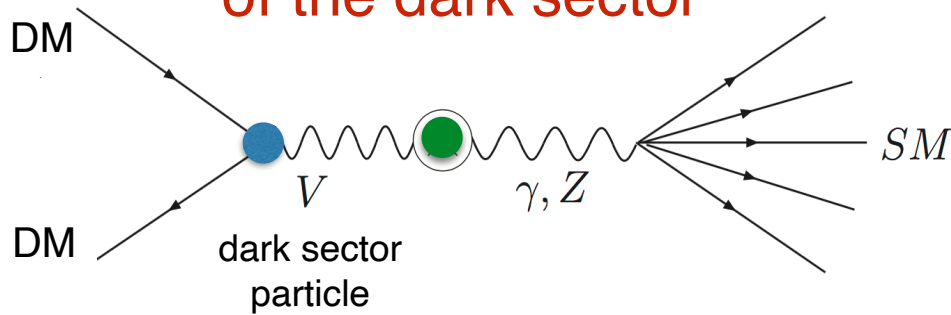
Relic abundance regulated by ●, ●



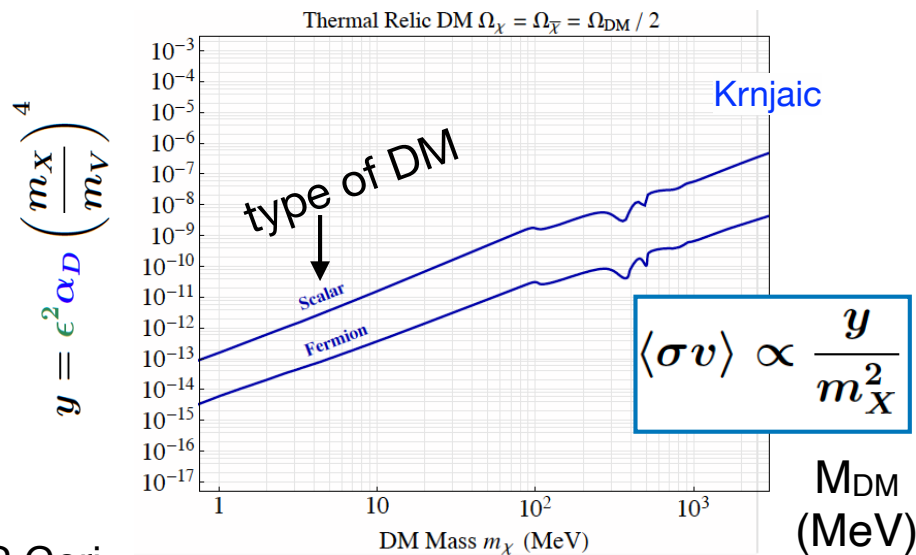
Lower bound on the interaction strength

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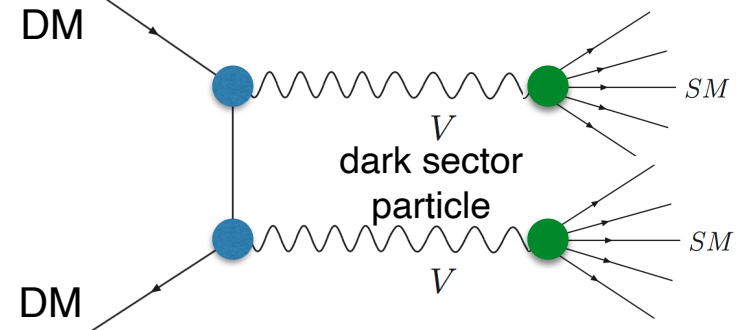


Relic abundance regulated by ●, ●

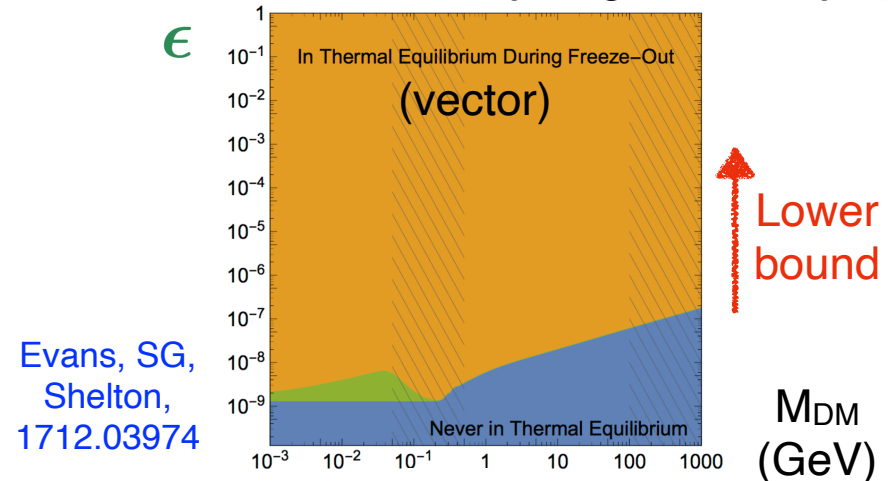


S.Gori

2. One (or more) particles of the dark sector are lighter than DM ("secluded" case)



Thermalization mainly regulated by ●

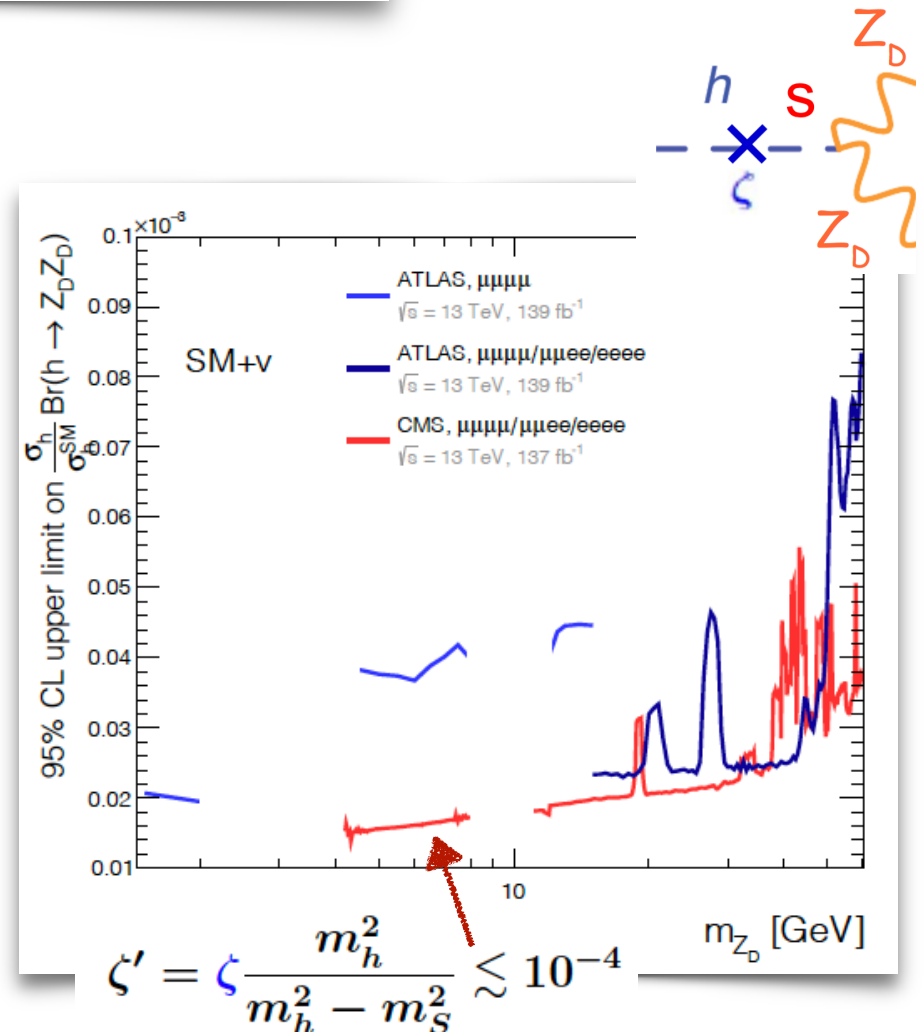
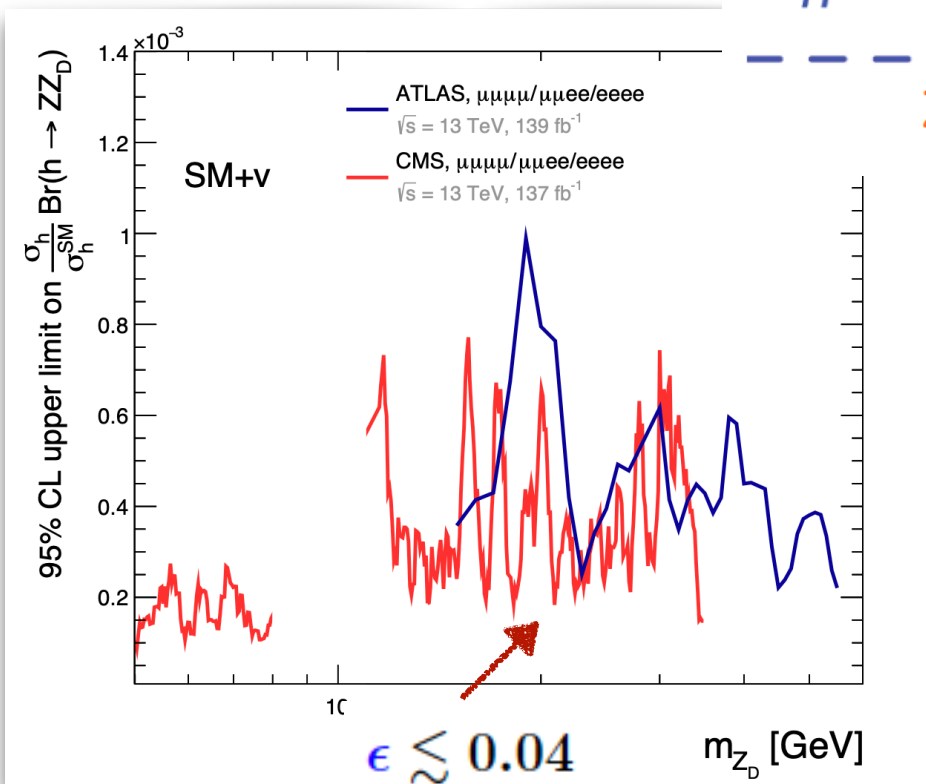


Dark photons and scalars from Higgs exotic decays

Curtin, Essig, SG,
Shelton, 1412.0018

$$\mathcal{L} \supset \frac{\epsilon}{2 \cos \theta} \hat{V}_{\mu\nu} \hat{B}^{\mu\nu} + \frac{1}{8} \langle S \rangle^2 g_D^2 (\hat{V}_\mu)^2 + \zeta |S|^2 |H|^2$$

Cepeda, SG, Martínez Outschoorn,
Shelton, 2111.12751



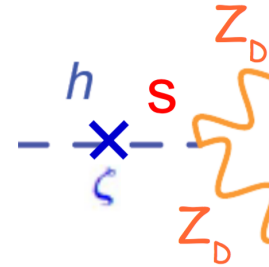
Prompt decays of the dark photon.

Dark photons and scalars from Higgs exotic decays

Curtin, Essig, SG,
Shelton, 1412.0018

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Cepeda, SG, Martínez Outschoorn,
Shelton, 2111.12751



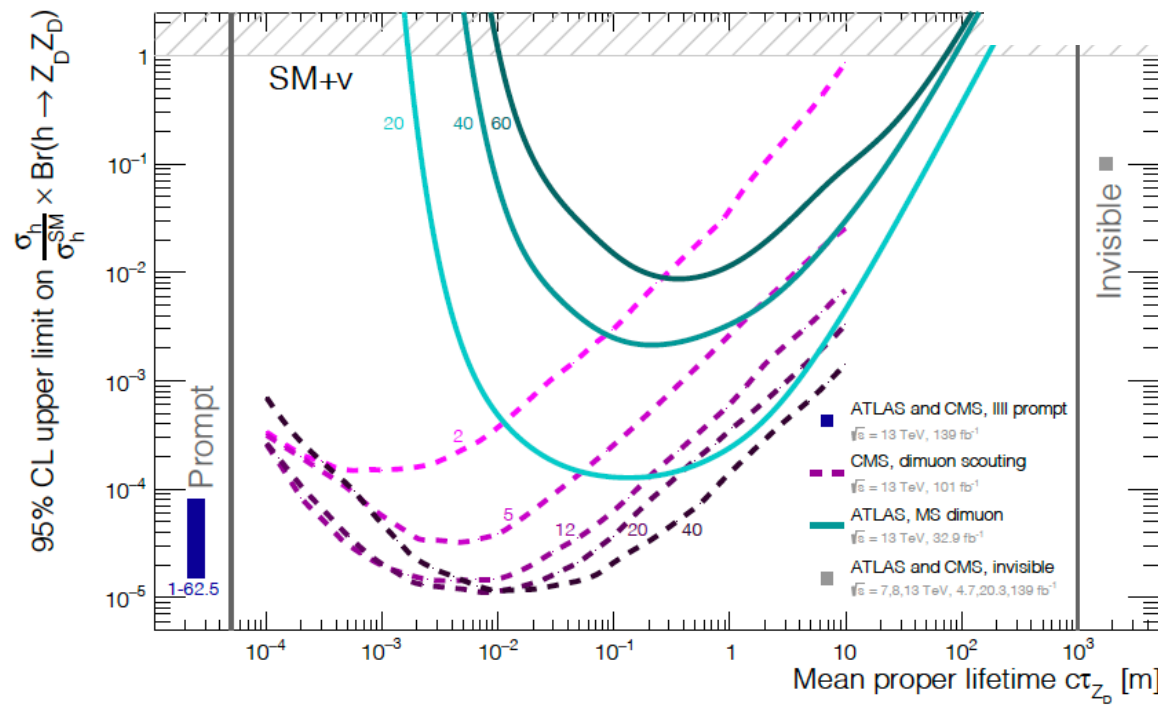
* Branching ratio depends on

$$\zeta' = \zeta \frac{m_h^2}{m_h^2 - m_S^2}$$

* Life time depends on ϵ .
At low mass, calculation done
with the $R(s)$ ratio:

$$\Gamma_{Z_D} = R_{Z_D} \Gamma(Z_D \rightarrow \mu^+ \mu^-) + \sum_{f=e,\mu,\tau,\nu_i} \Gamma(Z_D \rightarrow f \bar{f})$$

$$R_{Z_D}(m_{Z_D}) = R(s) \equiv \frac{\sigma(e^+ e^- \rightarrow \text{hadrons})}{\sigma(e^+ e^- \rightarrow \mu^+ \mu^-)}$$



Displaced decays of the dark photon.

Dark sectors at high intensity experiments

The **several portals** will lead to the **production of dark particles** at high intensity & high energy experiments

$$\epsilon Z^{\mu\nu} A'_{\mu\nu} \quad \kappa |H|^2 |S|^2$$

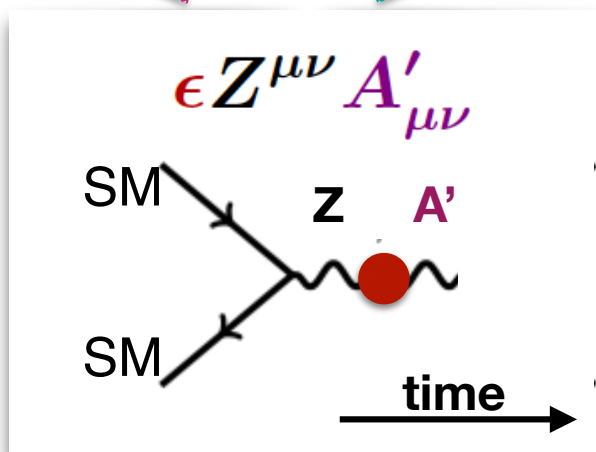
$$y H L N$$

(Typically) high energy

Colliding beam experiments



- * B-factories (Belle-II)
e⁺e⁻ collider
- * The LHC (pp collider)



(Typically) high intensity

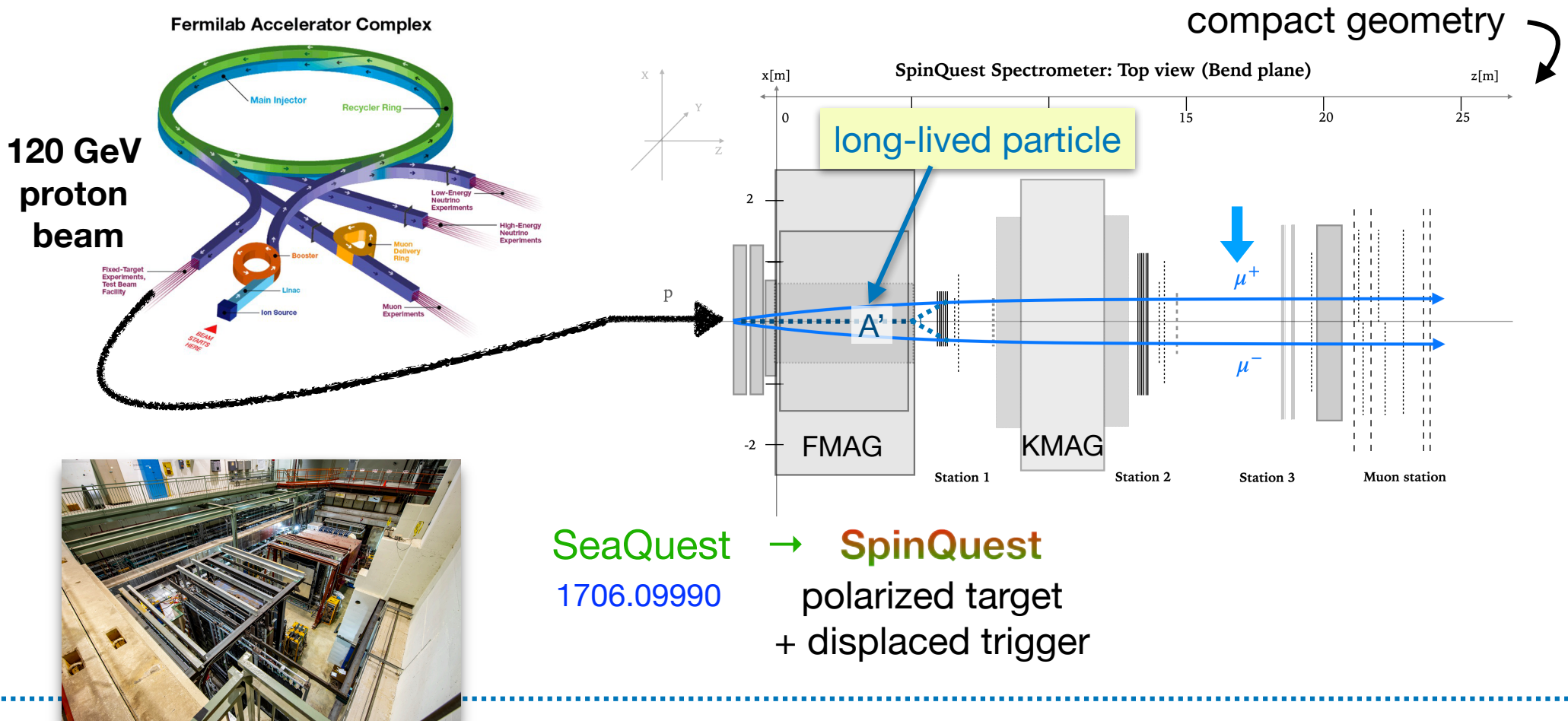
Fixed target experiments



- * Kaon exp. (NA62, KOTO)
- * proton beam dump exp.
- * electron beam dump exp.
- * electron fixed target exp.
- * neutrino exp.
- * light meson (e.g. pion) exp.

Proposing new experiments and searches to achieve a broad program of probes of the dark sector. Simplified model strategy vs. complete models

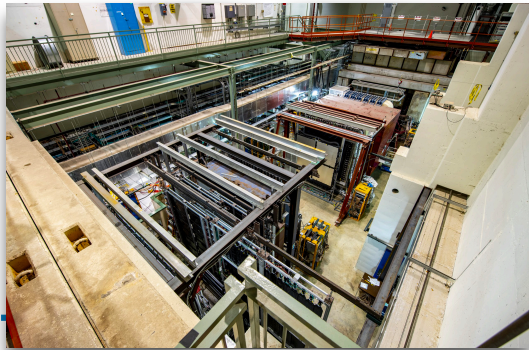
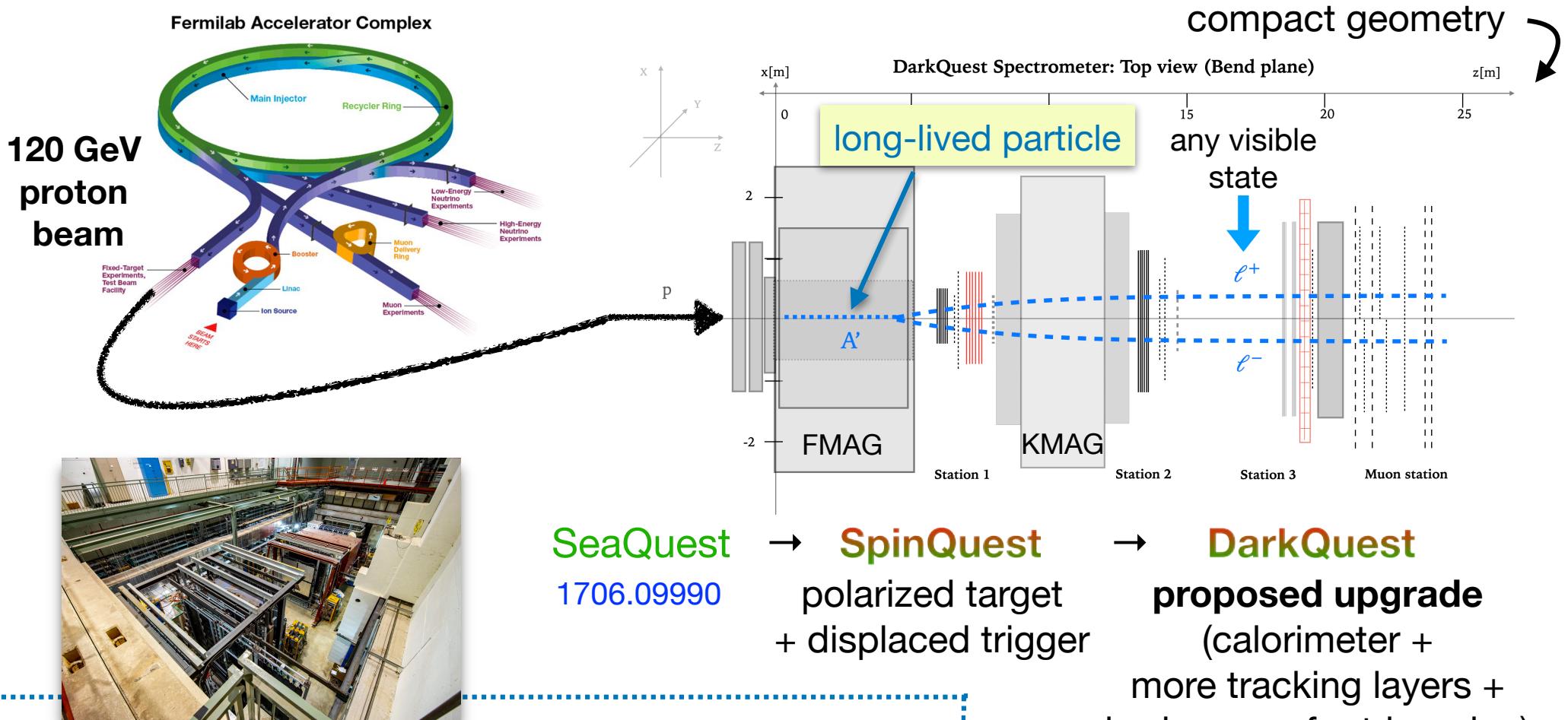
SpinQuest and DarkQuest



Nuclear physics: Measuring the Drell-Yan muon process for studies of the proton structure

Particle Physics: Visible dark sector searches (muons)

SpinQuest and DarkQuest



SeaQuest
1706.09990

→ **SpinQuest**
polarized target
+ displaced trigger

→ **DarkQuest**
proposed upgrade
(calorimeter +
more tracking layers +
hodoscope for triggering)

Nuclear physics: Measuring the Drell-Yan muon process for studies of the proton structure

Particle Physics: Visible dark sector searches
(any visible)

Initial proposal:

Gardner, Holt, Tadeipalli, 1509.00050

Berlin, SG, Schuster, Toro, 1804.00661

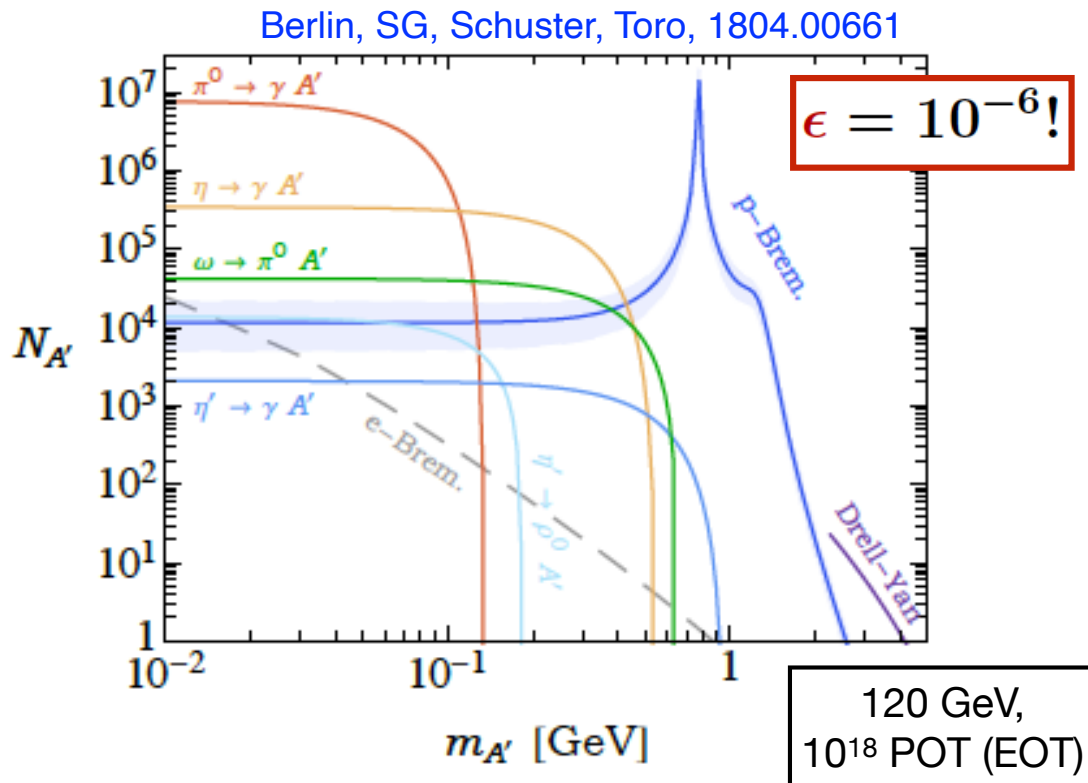
Snowmass white paper: 2203.08322

Why a high energy compact proton beam dump?

1. Large production rates of dark particles

In the case of a dark photon:

$$\mathcal{L} \supset \frac{\epsilon}{2 \cos \theta} \hat{V}_{\mu\nu} \hat{B}^{\mu\nu}$$

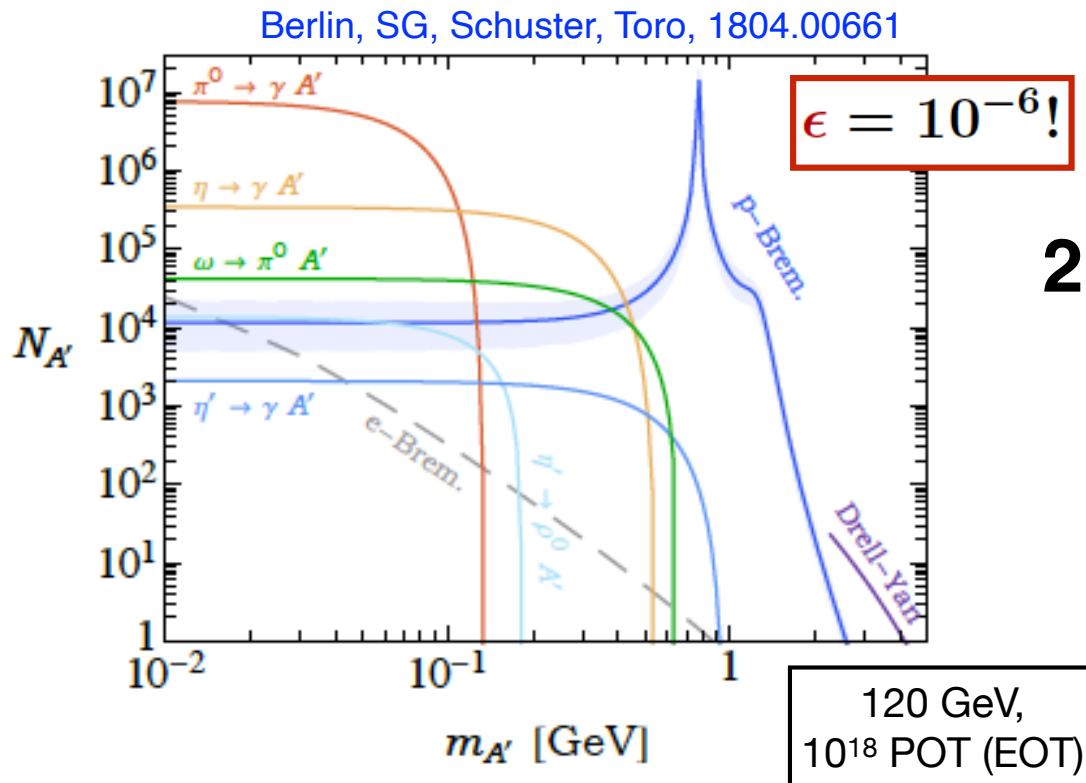


Why a high energy compact proton beam dump?

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1. A lot of light mesons:

$\sim 10^{19}$ pions

past pion factories: $\sim 10^{11}$ pi

2. Larger Bremsstrahlung production than at electron fixed target experiments:

proton: $\sigma \sim \alpha_{\text{em}} \epsilon^2 \times \sigma_{pp}$

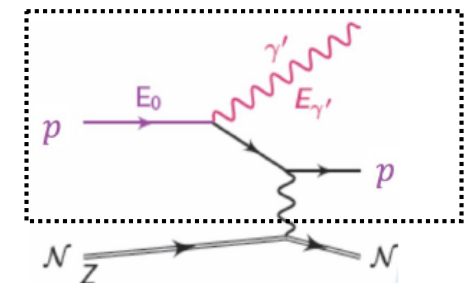
electron: $\sigma \sim \frac{\alpha_{\text{em}}^3 \epsilon^2}{m_{A'}^2} Z^2$

Blümlein et al, 1311.3870.

deNiverville et al, 1609.01770

Uncertainties in the calculation of the Bremsstrahlung
Generalized Williams-Weizsäcker approximation

$$E_p, E_{A'}, E_p - E_{A'} \gg m_p, m_{A'}, |p_T^{A'}|$$



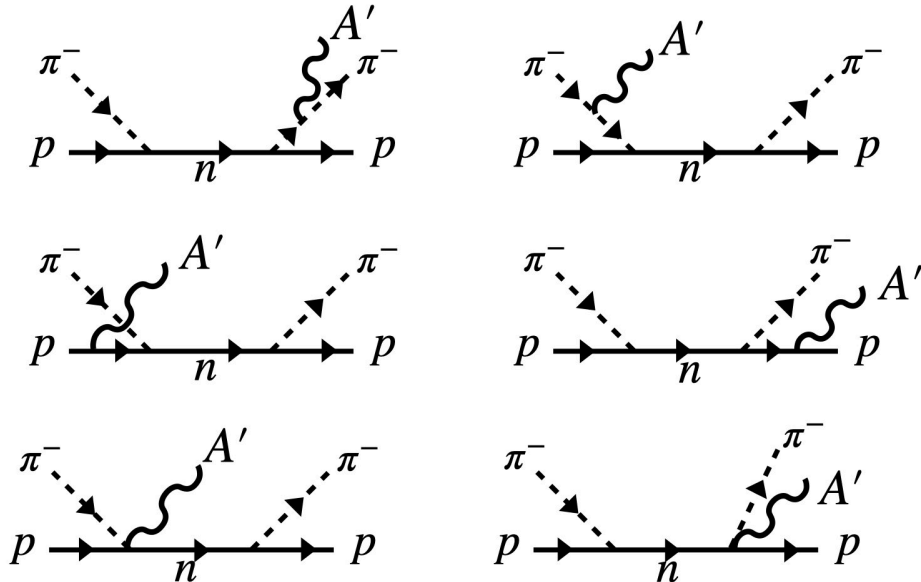
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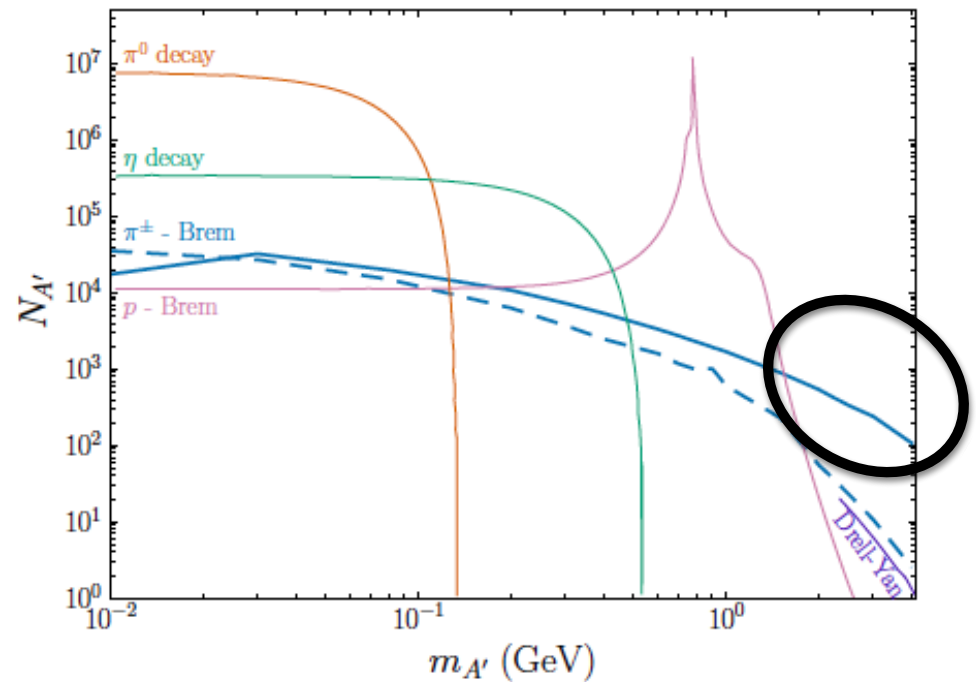
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Pion Bremsstrahlung



Secondary pion beam



Curtin, Kahn, Nguyen,
appearing soon

Why a high energy compact proton beam dump?

1. Large production rates of dark particles

Large production of heavier mesons as well

Batell, Evans, SG, Rai, 2008.08108

| <i>K</i> mesons* | | <i>D</i> mesons | | <i>B</i> mesons | | Leptons | |
|------------------|---------------------------|------------------|---------------------------|------------------|------------------------|------------------|---------------------------|
| K^\pm | $\sim 1.8 \times 10^{15}$ | D^\pm | $\sim 6.8 \times 10^{14}$ | B^\pm | $\sim 5.3 \times 10^7$ | τ^\pm | $\sim 4.7 \times 10^{10}$ |
| K_L^0 | $\sim 2.2 \times 10^{14}$ | D_s^\pm | $\sim 2.0 \times 10^{13}$ | B_d, \bar{B}_d | $\sim 5.3 \times 10^7$ | $\tau_{D_s}^\pm$ | $\sim 1.1 \times 10^{12}$ |
| K_S^0 | $\sim 1.2 \times 10^{17}$ | D^0, \bar{D}^0 | $\sim 1.3 \times 10^{14}$ | | | | |

number that decay before
one nuclear interaction length

Comparison with meson factories:
NA62, KOTO $\sim 10^{13}$ Kaons
Belle II (ultimate) 10^{10} B mesons

Why a high energy compact proton beam dump?

1. Large production rates of dark particles

$$\mathcal{L} \supset \zeta |S|^2 |H|^2$$

In the case of a dark scalar:

Large production of heavier mesons as well

Batell, Evans, SG, Rai, 2008.08108

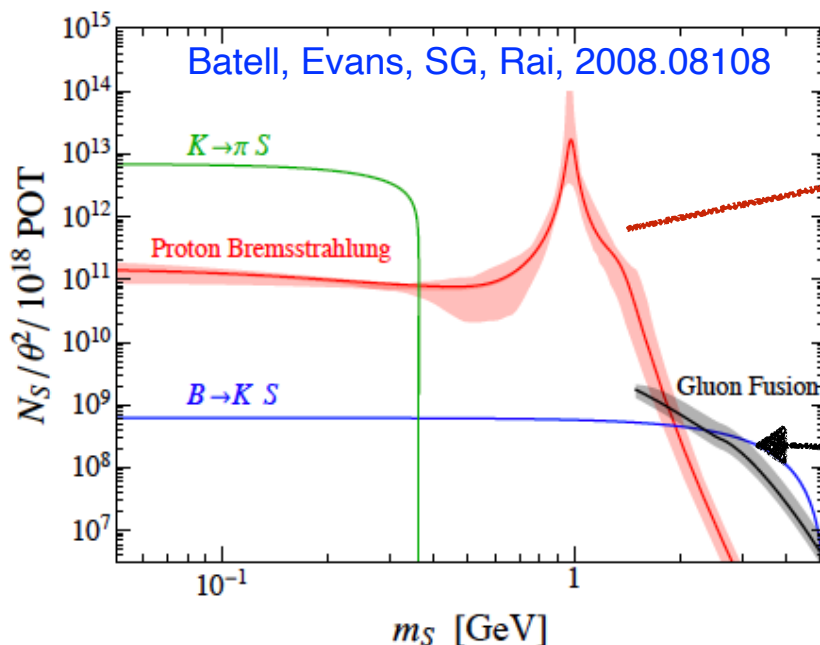
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number that decay before
one nuclear interaction length

Comparison with meson factories:

NA62, KOTO $\sim 10^{13}$ Kaons

Belle II (ultimate) 10^{10} B mesons



Williams-Weizsäcker approximation:

$$\frac{d\sigma_{\text{brem}}}{dz dp_T^2} \approx \sigma_{pp}(s') P_{p \rightarrow pS}(z, p_T^2).$$

$$P_{p \rightarrow pS}(z, p_T^2) \approx |F_S(m_S^2)|^2 \frac{g_{SNN}^2 \theta^2}{8\pi^2} \frac{z [m_p^2 (2-z)^2 + p_T^2]}{[m_p^2 z^2 + m_S^2 (1-z) + p_T^2]^2}$$

leading order calculation.
O(1) corrections expected

Why a high energy compact proton beam dump?

1. Large production rates of dark particles

$$\mathcal{L} \supset \zeta |S|^2 |H|^2$$

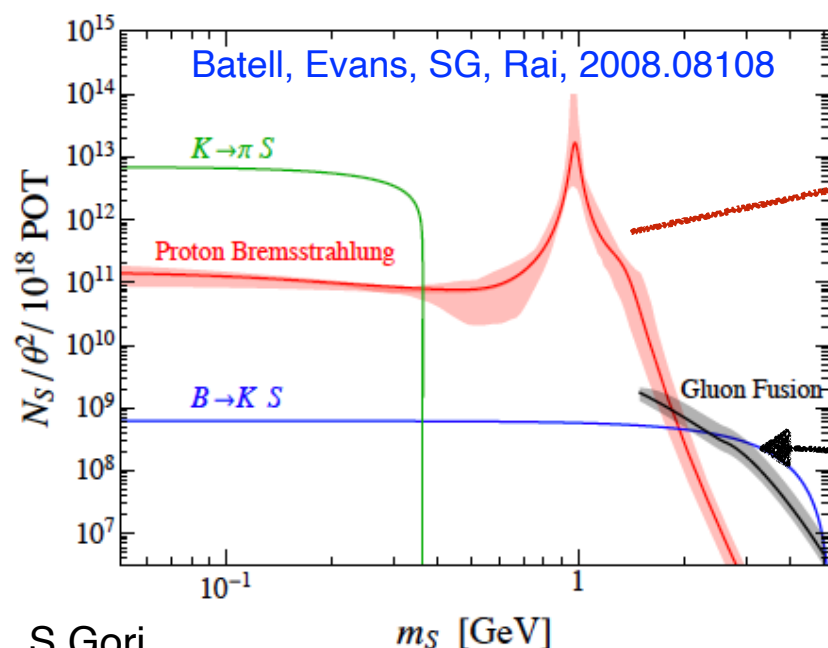
In the case of a dark scalar:

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Batell, Evans, SG, Rai, 2008.08108

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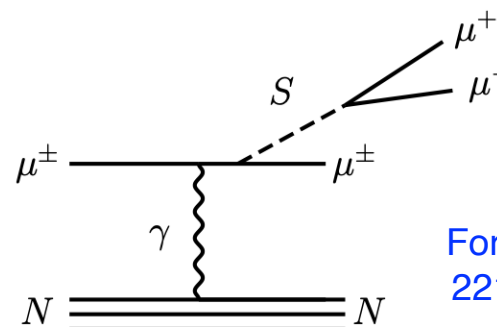
number that decay before
one nuclear interaction length



Comparison with meson factories:

Note: **very intense secondary muon beam**
 Reach on muon-coupled dark sectors

example:



Forbes et al,
2212.00033

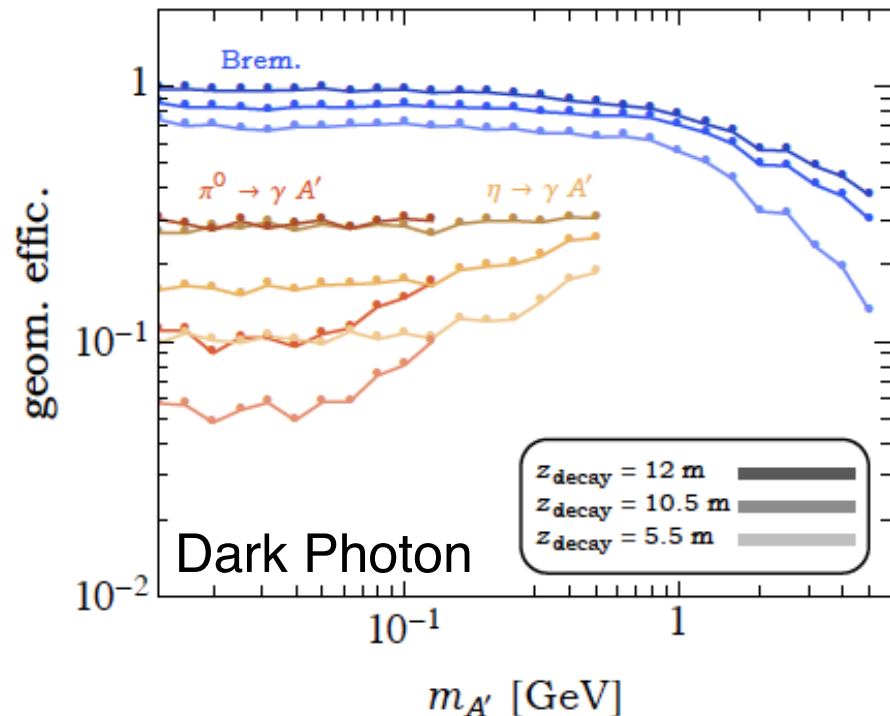
leading order calculation.

O(1) corrections expected

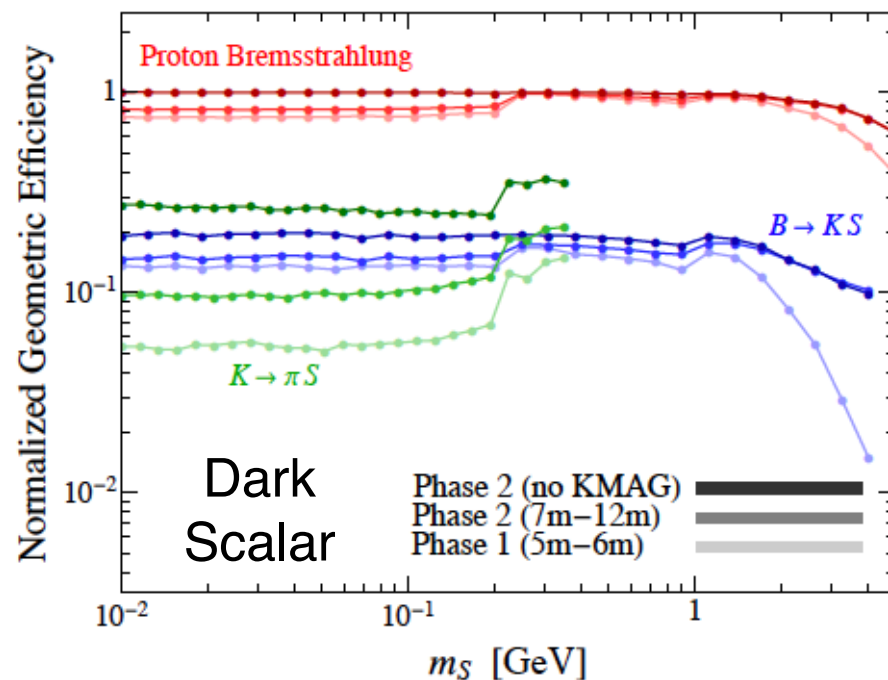
Why a high energy compact proton beam dump?

2. High geometric acceptance

Berlin, SG, Schuster, Toro, 1804.00661

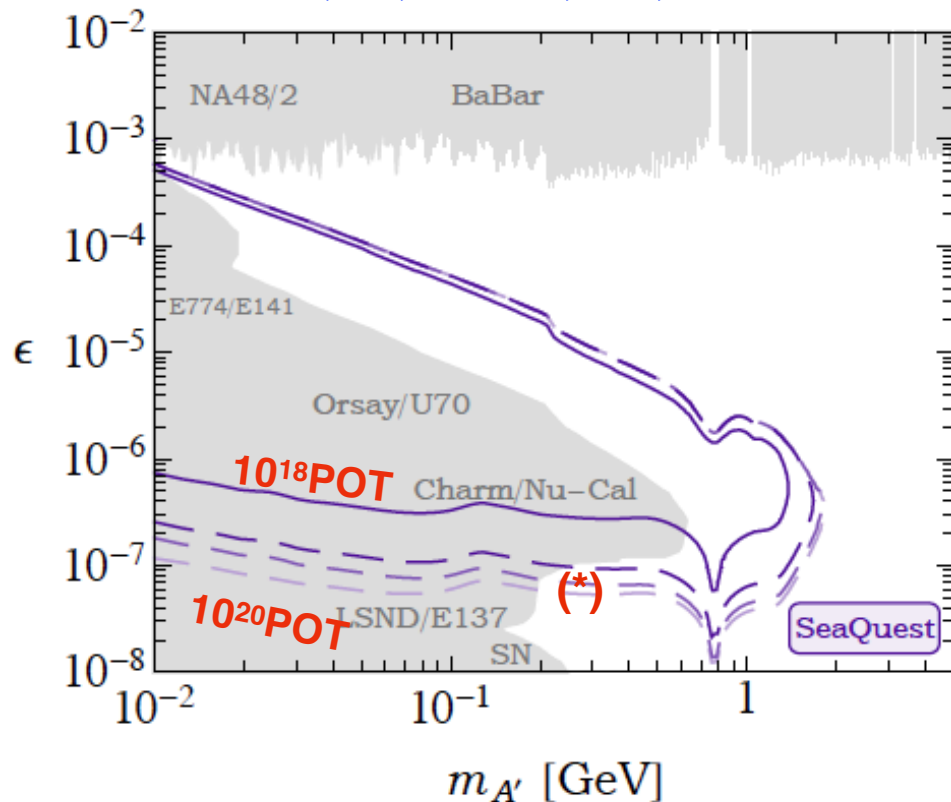


Batell, Evans, SG, Rai, 2008.08108



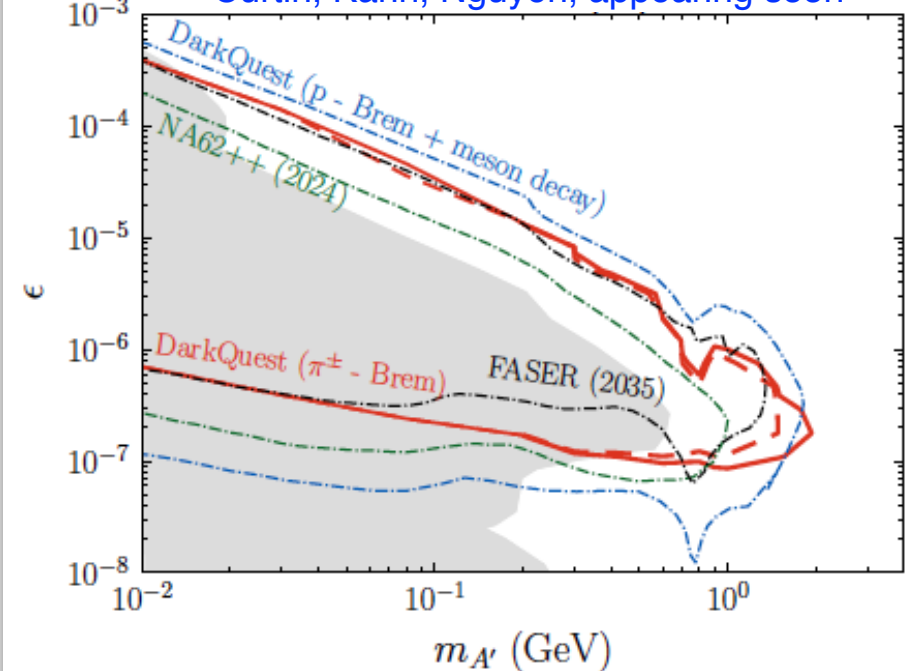
The reach on a minimal dark photon

Berlin, SG, Schuster, Toro, 1804.00661

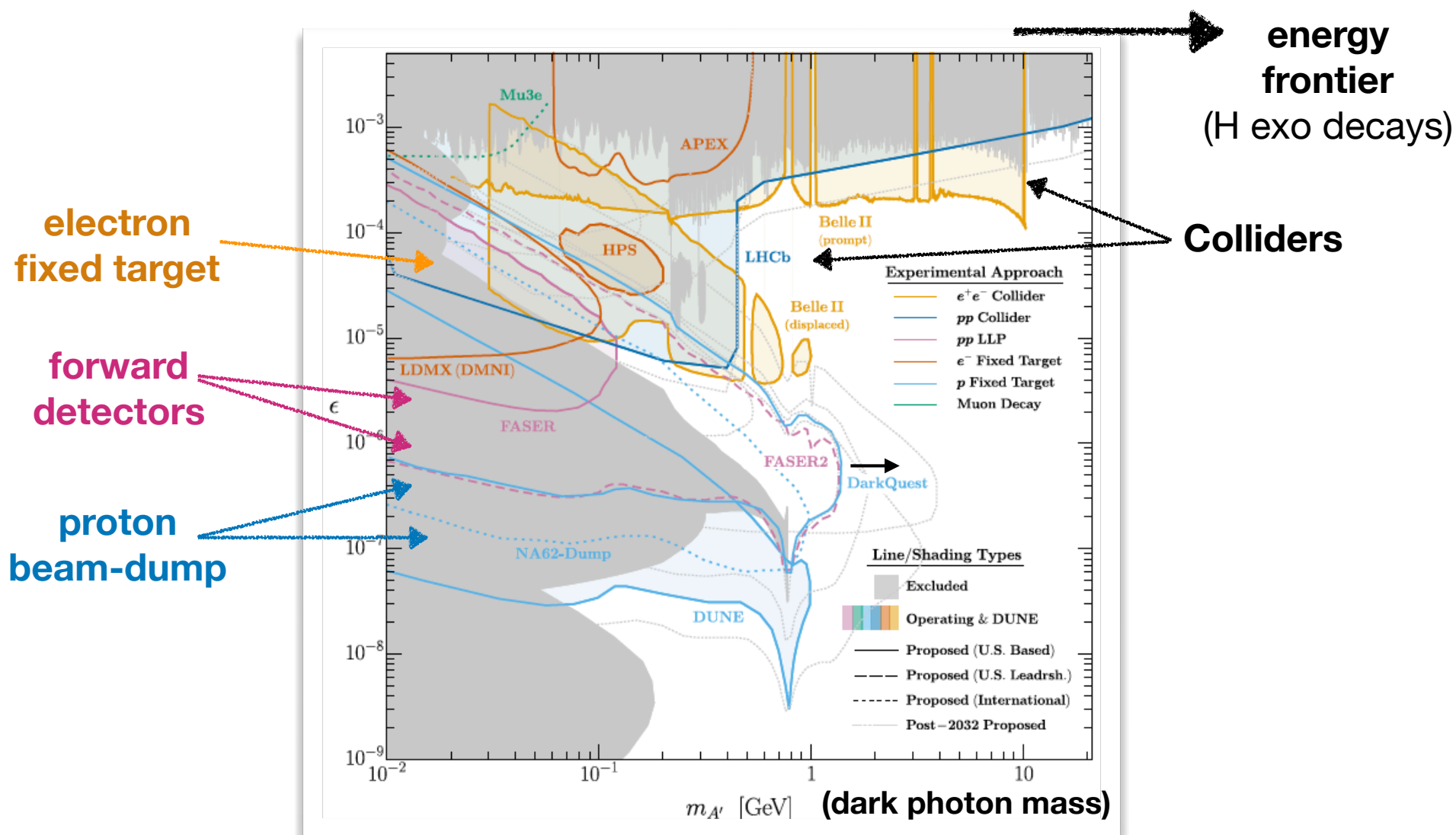


(*) decay regions: (5-6)m, (5-9)m, (5-12)m

Curtin, Kahn, Nguyen, appearing soon



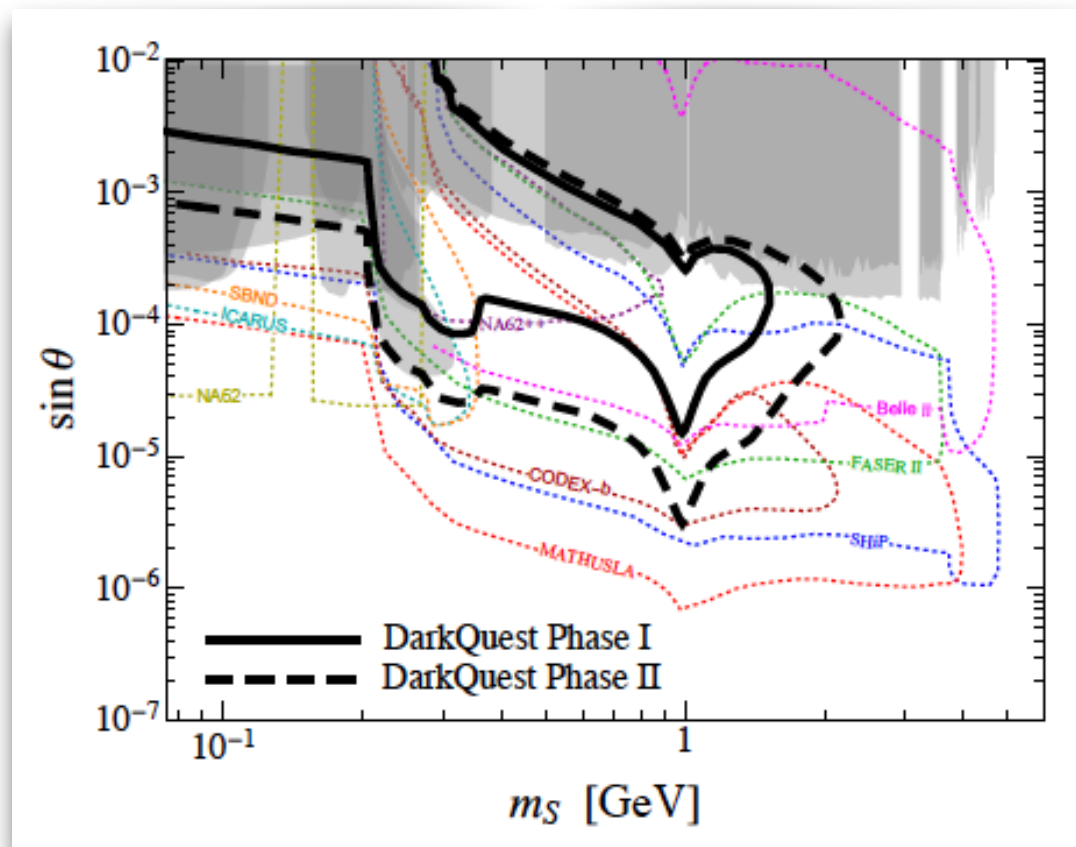
Putting it in a broader context



Batell, Blinov, Hearty, McGehee, 2207.06905,
Snowmass white paper

The reach on a minimal dark scalar

Batell, Evans, SG, Rai, 2008.08108



Large uncertainties in the calculations of branching ratios and life time of the scalar

Winkler, 1809.01876

- chiral perturbation theory
- perturbative spectator
- dispersion analyses

↓ mass

Many searches can be performed

The main strength of the experiment is on dark photon and axion-mediated models + models with new muon-coupled particles.

Several studies have been performed, including less minimal models for freeze-out Dark Matter.

| Signature | Model |
|------------------------------------|---|
| e^+e^- | dark photon dark Higgs leptophilic scalar* |
| $e^+e^-e^+e^-$ | Higgsed dark photon |
| $e^\pm\pi^\mp, e^\pm K^\mp, \dots$ | sterile neutrino |
| $e^+e^- + \text{MET}$ | inelastic dark matter strongly interacting dark matter hidden valleys |
| $\pi^+\pi^-, K^+K^-, \dots$ | dark Higgs* |
| $\gamma\gamma$ | axion-like particle* |

Many searches can be performed

The main strength of the experiment is on dark photon and axion-mediated models + models with new muon-coupled particles.

Several studies have been performed, including less minimal models for freeze-out Dark Matter.

* Two examples:

Inelastic DM models (IDM): $\chi_1 \chi_2$ co-annihilation,

$$A' \rightarrow \chi_1 \chi_2, \chi_2 \rightarrow \chi_1 e^+ e^-$$

long-lived since small mass splitting are needed for the co-annihilation to be efficient

$$\Gamma(\chi_2 \rightarrow \chi_1 e^+ e^-) \simeq \frac{4\epsilon^2 \alpha_{\text{em}} \alpha_D \Delta^5 m_1^5}{15\pi m_{A'}^4}$$

* Strongly interacting DM models (SIMP): $3\pi_D \rightarrow 2\pi_D$ (QCD-like models)

$$A' \rightarrow \pi_D V_D, V_D \rightarrow \pi_D e^+ e^-$$

$$A' \rightarrow \pi_D V_D, V_D \rightarrow e^+ e^-$$

Berlin, Blinov, SG,
Schuster, Toro, 1801.05805

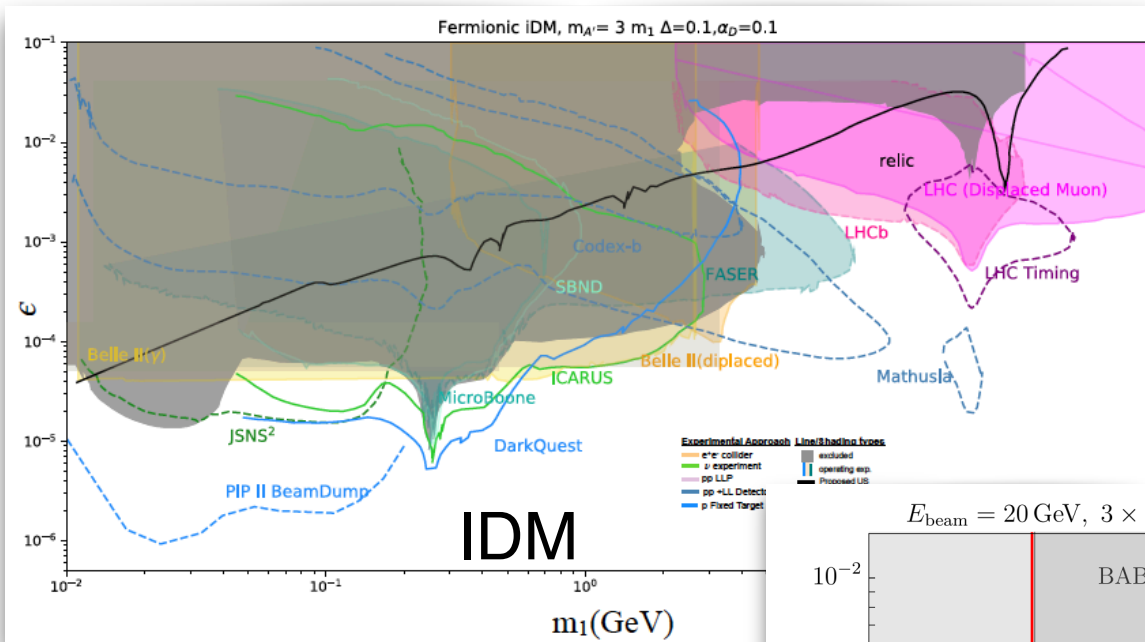
long-lived

Width is $\epsilon^2 \alpha_{\text{em}}$
suppressed

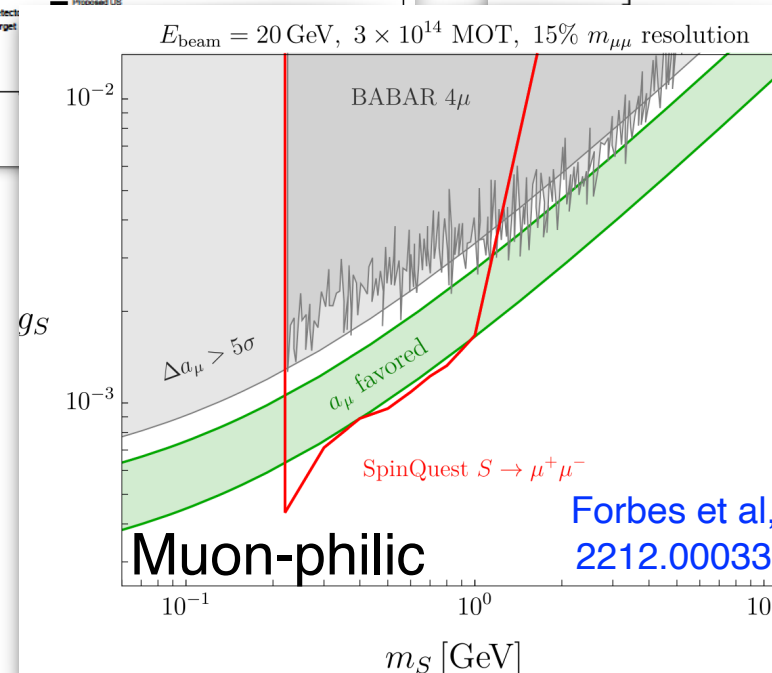
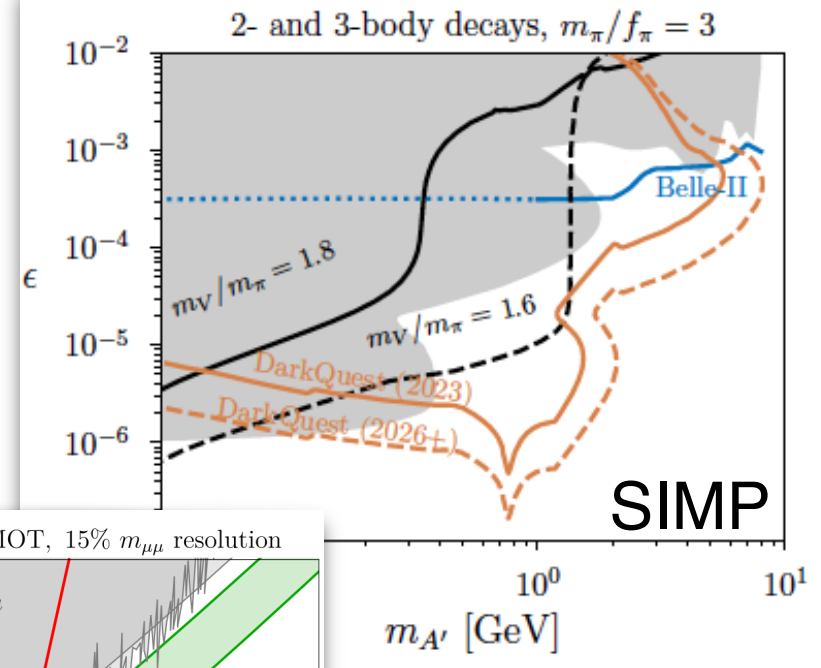
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| $e^+ e^- + \text{MET}$ | inelastic dark matter strongly interacting dark matter hidden valleys |
| $\pi^+ \pi^-, K^+ K^-, \dots$ | dark Higgs* |
| $\gamma\gamma$ | axion-like particle* |

The reach on DM /anomaly motivated models

Harris, Schuster, Zupan, 2207.08990



Apyan et al, 2203.08322



Forbes et al,
2212.00033

List of experimental studies

In the past two years a lot of progress has been made:

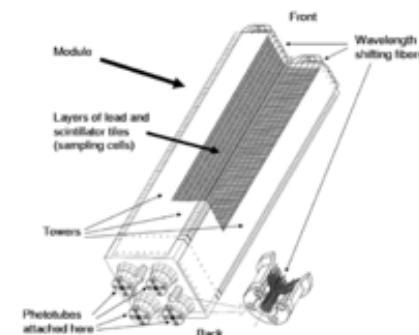
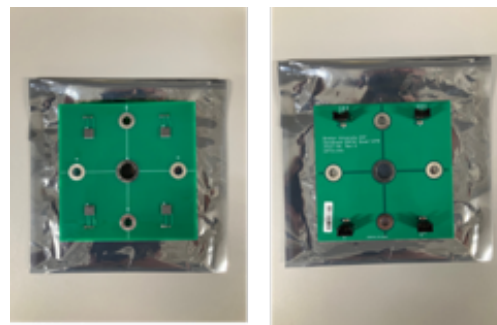
* Detector:

- o EMCal integration into the SpinQuest spectrometer
- o Extra Tracking layer integration into the SpinQuest spectrometer

* GEANT4 - based simulations:

- o EMCal simulations
- o Triggering
- o Tracking & vertexing
- o ParticleID: tracking + calorimeter information

Custom 4-ch SiPM Board



EMCal Test Stand at BU



Strong connections
with the SpinQuest
collaboration

Thank you, Carlos and Marcela!



SpinQuest/DarkQuest status and timeline

