# **Caveats concerning the G4-DF simulation**

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1

- **● G4 receives a set of wavelength-dependent transmitance curves**
- **● Each curve is defined for a certain angle of incidence (AOI)**
- We endow a surface with the provided transmitance data
- When a photon reaches such surface, G4 performs a 2D-interpolation (wavelength and AOI) and comes up with a transmitance value for such event
- The photon is transmitted or specularly-reflected up to the interpolated transmitance



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If reflection takes place, the photon is specularly reflected If transmission takes place, **no snell refraction is applied!**



From Geant4 user's guide > Tracking and physics > Physics processes > Optical photon processes > Boundary process: *" As expressed in Maxwell's equations, Fresnel reflection and refraction are intertwined through their relative probabilities of occurrence. Therefore neither of these processes, nor total internal reflection, are viewed as individual processes deserving separate class implementation. "*

Fresnel reflection/refraction are regarded as processes which are ruled-by and inseparable-from their relative probabilities. Thus, setting a custom transmission probability destroys such processes: now just reflection and "raw" transmission are considered.



#### **DF previous considerations**



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#### **DF computational model**





For a realistic simulation of the DF, refraction in every surface must be simulated, but as we have seen, refraction cannot be separated from reflection

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So, if we measure certain TCs in the lab, and assign them to the substrate->MLS interface, these interfaces will contribute with an extra reflection probability.









 $I_m(\lambda, \theta_{sm}) = I_o(\lambda) \cdot FT^{n_{sm}}_{n_{su}}(\theta_{sm}) \cdot X(\lambda, \theta_{sm})$ 

**First caveat First caveat But, since we want to simulate the DF** in medias that are different from that of the lab, we are interested in decoupling the transmission in the last interface, which depends on the SM r. index



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 $I_m(\lambda, \theta_{sm}) = I_o(\lambda) \cdot FT_{n_{su}}^{n_{sm}}(\theta_{sm}) \cdot {ITC}(\lambda, \theta_{sm}) \cdot FT_{n_{sm}}^{n_{DF}}(\alpha(\theta_{sm}))$ 

$$
ITC(\lambda, \theta_{sm}) = \frac{T_m(\lambda, \theta_{sm})}{FT_{n_{su}}^{n_{sm}}(\theta_{sm}) \cdot FT_{n_{sm}}^{n_{DF}}(\alpha(\theta_{sm}))}
$$
 If we knew everything here, we could compute the 'intrinsic' transmitance curve of the dichroic filter as shown above 
$$
I_m(\lambda, \theta_{sm}) = I_o(\lambda) \cdot FT_{n_{su}}^{n_{sm}}(\theta_{sm}) \cdot ITC(\lambda, \theta_{sm}) \cdot FT_{n_{sm}}^{n_{DF}}(\alpha(\theta_{sm}))
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$$

The truth is we know everything except for **this**, but we can overcome this if we assume the following conceptual model:

# **DF conceptual model**



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 $\rightarrow$  Substitute the MLS by an homogeneous volume with an effective refractive index



Within this model, to explain the dependence of TCs with the surrounding media, our current hypothesis is that **what determines the TC is the angle of refraction (AOR) within the ~MLS.**

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# **Model backups**

There are **experimental** measurements and Transfer-matrix-method simulations backing up this hypothesis:

**ZOOM** 









25

# **Model backups**

There are experimental measurements and Transfer-matrix-method **simulations** backing up this hypothesis:



$$
ITC(\lambda, \theta_{sm}) = \frac{T_m(\lambda, \theta_{sm})}{FT_{nsu}^{n_{sm}}(\theta_{sm}) \cdot FT_{nsm}^{\overline{n_{DF}}(\alpha(\theta_{sm}))}}
$$
  
\nWithin this approximation, we know what the refractive index of the dichroic filter is (the effective r. index of the MLS) and we can compute the angle of the light that propagates within the ~MLS using Snell's law  
\n
$$
ITC(\lambda, \theta_{sm}) = \frac{T_m(\lambda, \theta_{sm})}{FT_{nsu}^{n_{sm}}(\theta_{sm}) \cdot FT_{nsm}^{n_{eff}}(S(\theta_{sm}))}
$$



# **DF computational model**



#### **Second caveat**



To simulate a realistic angular distribution of the photons that have entered the XA, we need to set the ~MLS volume refractive index to that of the substrate

**Previous considerations on parallel-faces stackings**



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Indeed, **we only need to know the pair (theta\_i, n\_i) for some**   $n_i$ sin $\theta_i$  = cte  $\forall i$ **layer**, to know the light propagation direction in any other layer

#### **Second caveat**



Since no refraction is simulated in the ITC interface, the angular distribution of photons within the ~MLS volume is that of the photons in the substrate volume (this is a simulation artifact). To prevent this artifact from spoiling a realistic simulation of the angular distribution of photons in LAr, the refraction in the ~MLS->LAr interface must account for the substrate refractive index.

**In G4, the TC is applied at some boundary**, so there's actually no volume where an AOR is defined. I.e. when a photon reaches the dichroic boundary, G4 won't have access to any AOR information: **The only information that G4 counts on to decide which TC to apply is the AOI**

Assuming the AOR model, a way to proceed is the following one:



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Assuming the AOR model, a way to proceed is the following one:



1) A photon impinges on the dichroic boundary with a certain AOI theta, from a media with refractive index n\_1

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2) Ask yourself, what would the AOR be if there was actually a volume with a refractive index equal to the MLS effective refractive index?

I.e. compute "forward" snell's law

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3) Ask yourself, what AOI in air would give rise to such AOR?

I.e. compute "backward" snell's law

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Assuming the AOR model, a way to proceed is the following one:



40

 $-75°$ 

 $20<sup>5</sup>$  $25°$ 

 $30^\circ$  $35<sup>o</sup>$ 40° 45º 50° 55° 60° 65º 70<sup>o</sup>

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As a result, we redefine the angles of the TCs we measured in air to overcome the fact that no MLS refraction happens in the simulation



Assume we have measured a DF in air from 0° to 75° AOI, and assume we are modelling the DF using the first approach, where photons impinge over the dichroic boundary from a fused silica (FS) media (n≃1.47). Then, this method assigns the 0º TC curve in air to the 0º TC curve in FS, and **the 75º TC curve in air to the 41º TC curve in FS**



**\***The information loss is actually the same in both cases, if we assume that for both cases, the incoming light was emitted by PTP with a fixed AD

If we want to simulate a PTP (n≃1.65) layer right above the dichroic boundary, the information loss is even worse**\***



Assuming  $n_{\text{eff}} \simeq 1.68$ 

The underlying cause is that sweeping 0-75º AOI in air does not let us explore the full range of AOR that refraction from PTP gives:



The most immediate 'solution' is to extrapolate the nearest (angle–wise) TC to every unknown value:



Since the bigger the AOI, the worse DF performance, extrapolating the biggest-AOI available TC to bigger angles sets an upper bound to the DF performance for AOIs whose TC we have not measured yet.

Measuring the TCs within water would widen our AOR scope

![](_page_45_Figure_2.jpeg)

Assuming  $n_{\text{eff}}\simeq 1.68$