Caveats concerning the G4-DF simulation

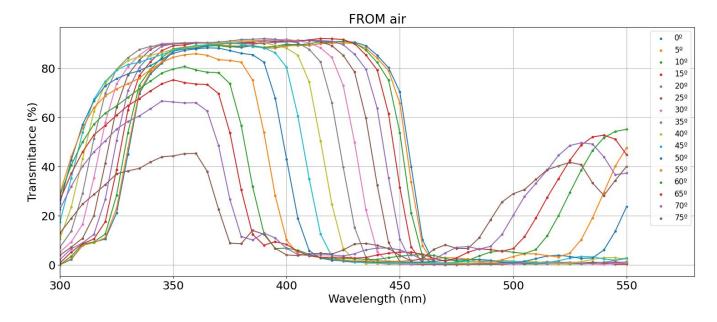
Julio Ureña, Justo Martín-Albo, Anselmo Cervera

Photon Collectors WG - 7 February 2023

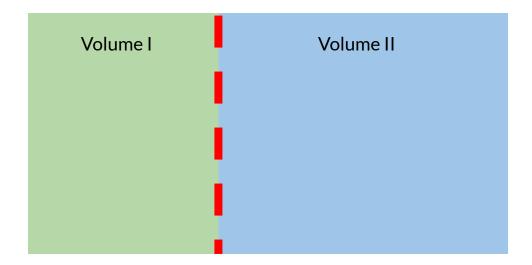


1

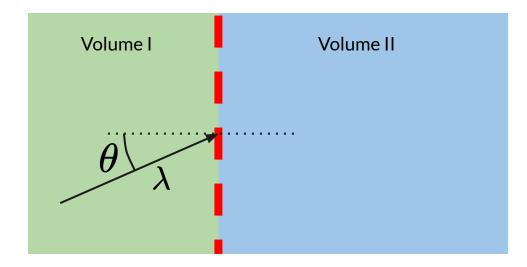
- G4 receives a set of wavelength-dependent transmitance curves
- Each curve is defined for a certain angle of incidence (AOI)
- We endow a surface with the provided transmitance data
- When a photon reaches such surface, G4 performs a 2D-interpolation (wavelength and AOI) and comes up with a transmitance value for such event
- The photon is transmitted or specularly-reflected up to the interpolated transmitance



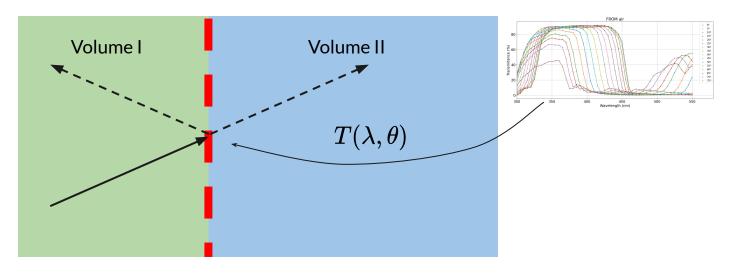
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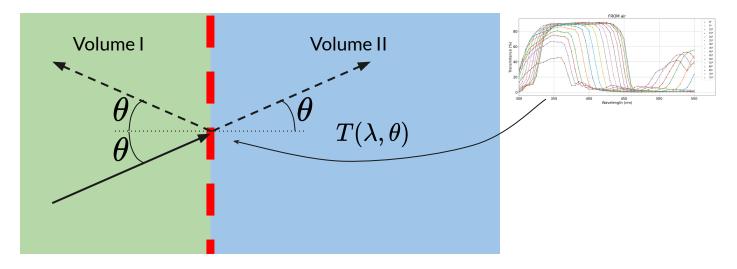
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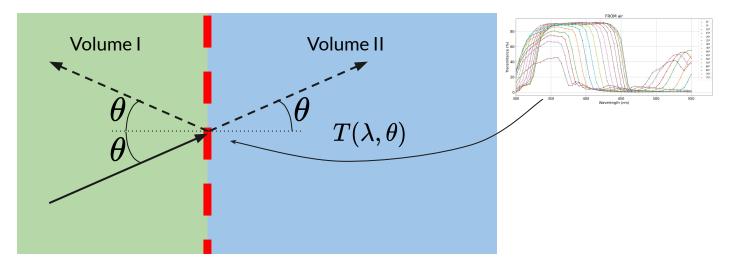


If reflection takes place, the photon is specularly reflected If transmission takes place, **no snell refraction is applied!**

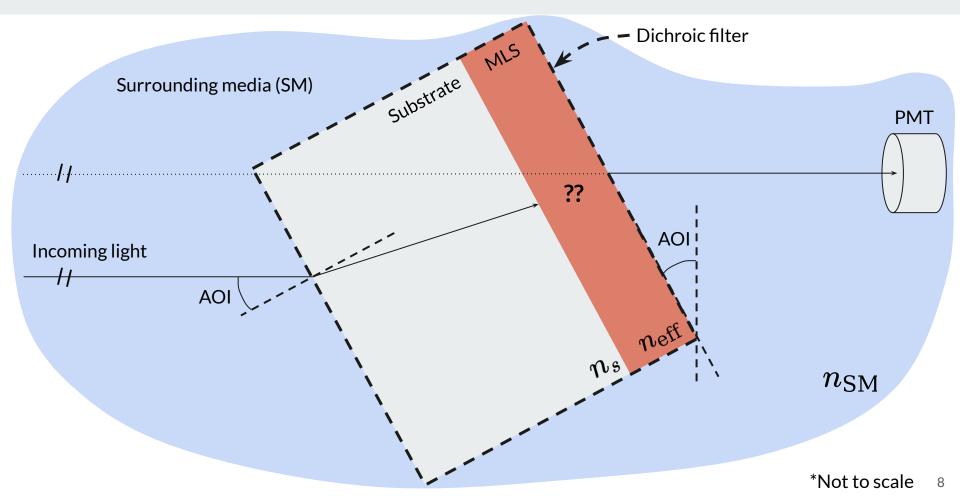


From Geant4 user's guide > Tracking and physics > Physics processes > Optical photon processes > Boundary process:
"As expressed in Maxwell's equations, Fresnel reflection and refraction are intertwined through their relative probabilities of occurrence. Therefore neither of these processes, nor total internal reflection, are viewed as individual processes deserving separate class implementation."

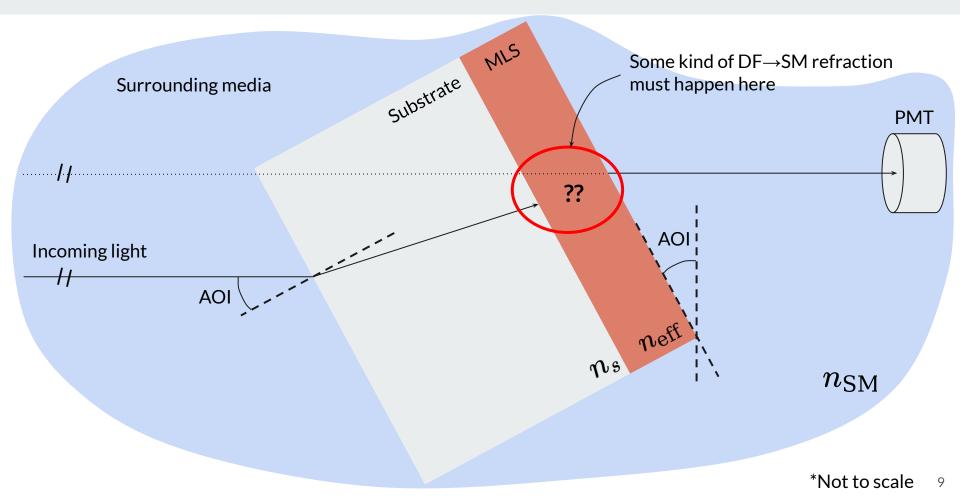
Fresnel reflection/refraction are regarded as processes which are ruled-by and inseparable-from their relative probabilities. Thus, setting a custom transmission probability destroys such processes: now just reflection and "raw" transmission are considered.



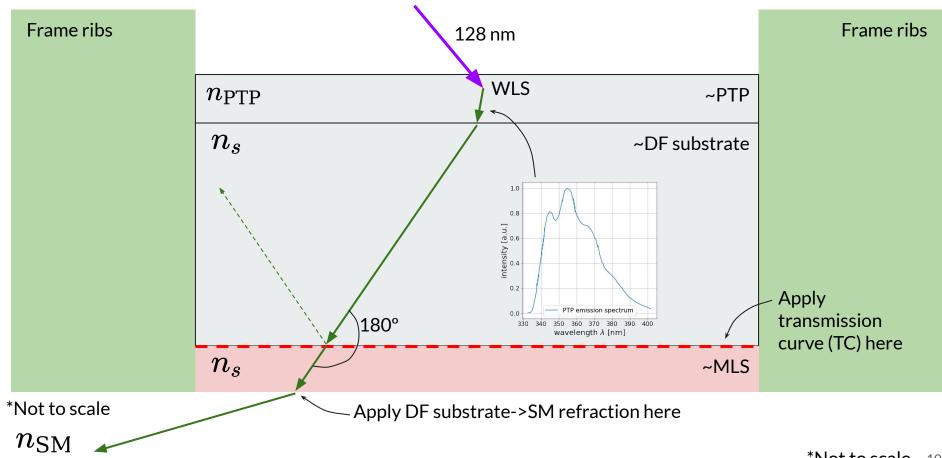
DF previous considerations



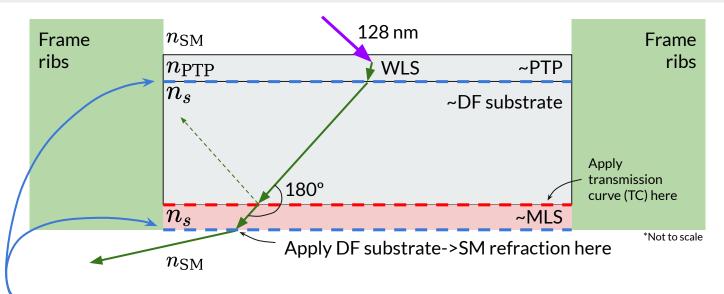
DF previous considerations



DF computational model

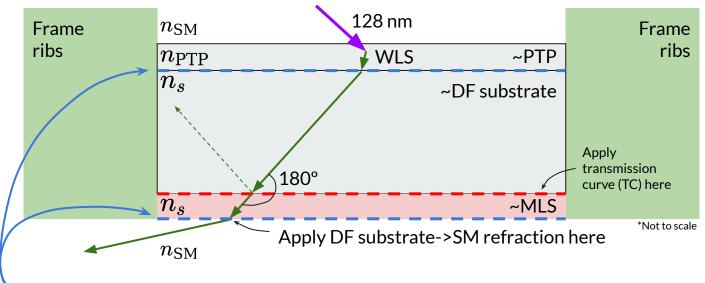


*Not to scale 10

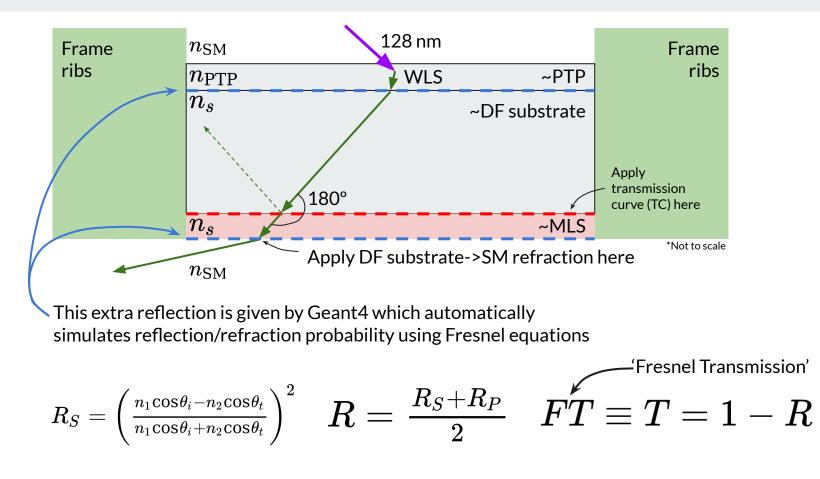


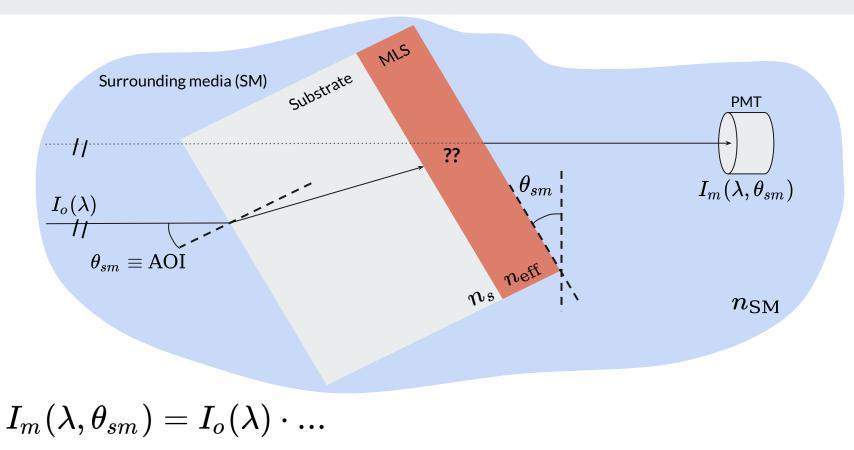
For a realistic simulation of the DF, refraction in every surface must be simulated, but as we have seen, refraction cannot be separated from reflection

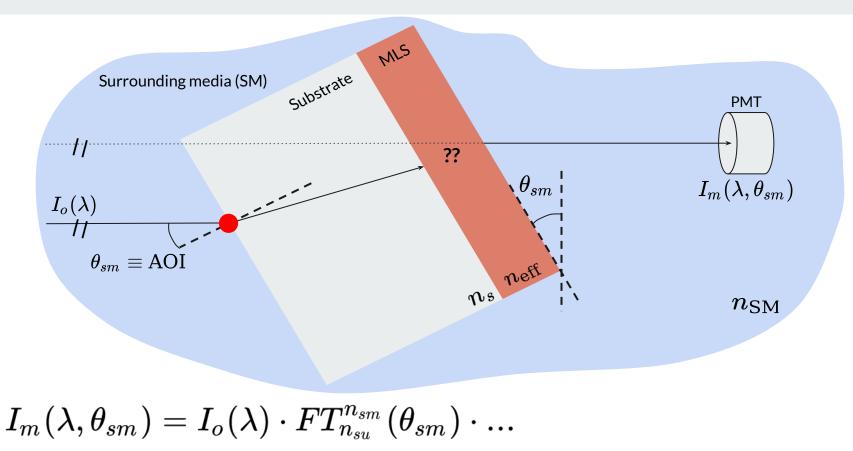
From Geant4 user's guide > Tracking and physics > Physics processes > Optical photon processes > Boundary process: "As expressed in Maxwell's equations, Fresnel reflection and refraction are intertwined through their relative probabilities of occurrence. Therefore neither of these processes, nor total internal reflection, are viewed as individual processes deserving separate class implementation."

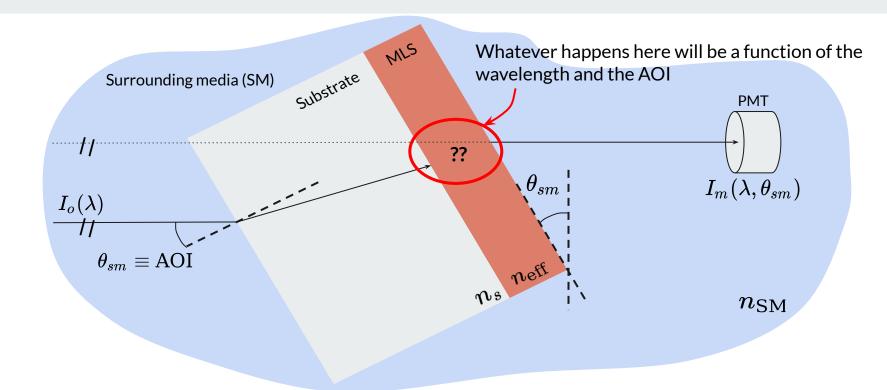


So, if we measure certain TCs in the lab, and assign them to the substrate->MLS interface, these interfaces will contribute with an extra reflection probability.



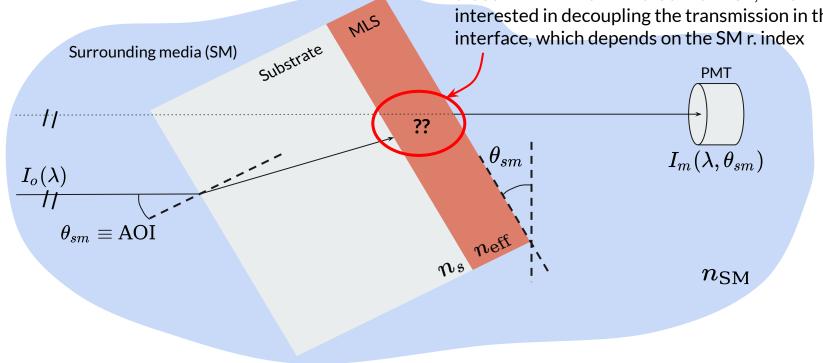




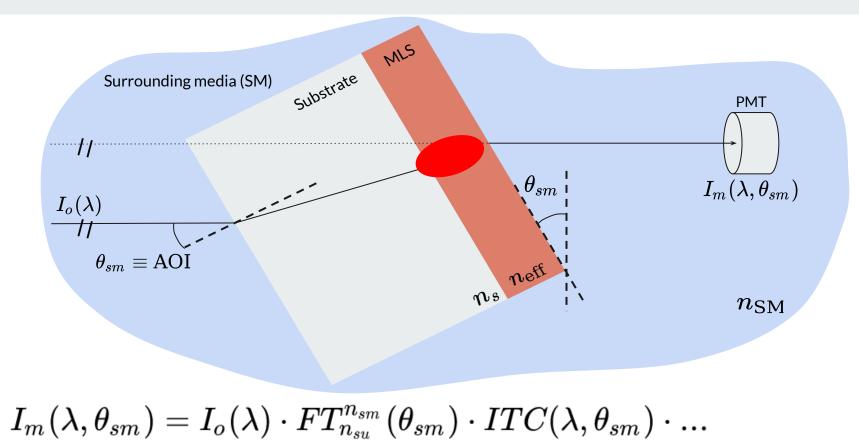


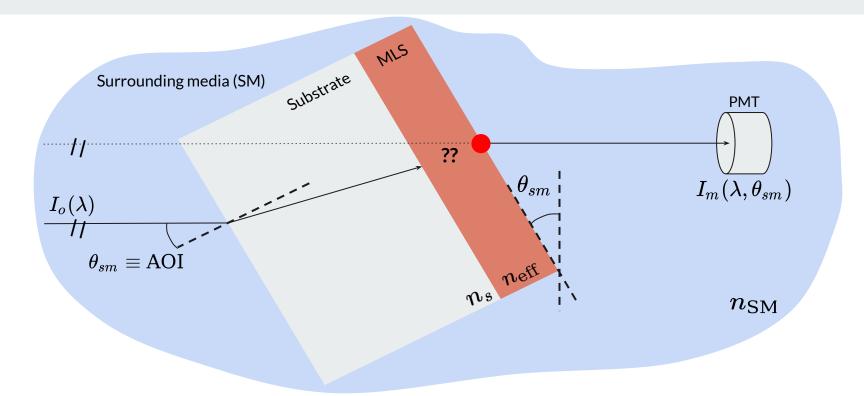
$I_m(\lambda, heta_{sm}) = I_o(\lambda) \cdot FT^{n_{sm}}_{n_{su}}(heta_{sm}) \cdot X(\lambda, heta_{sm})$

But, since we want to simulate the DF in medias that are different from that of the lab, we are interested in decoupling the transmission in the last interface, which depends on the SM r. index



 $I_m(\lambda, heta_{sm}) = I_o(\lambda) \cdot FT^{n_{sm}}_{n_{su}}(heta_{sm}) \cdot X(\lambda, heta_{sm})$





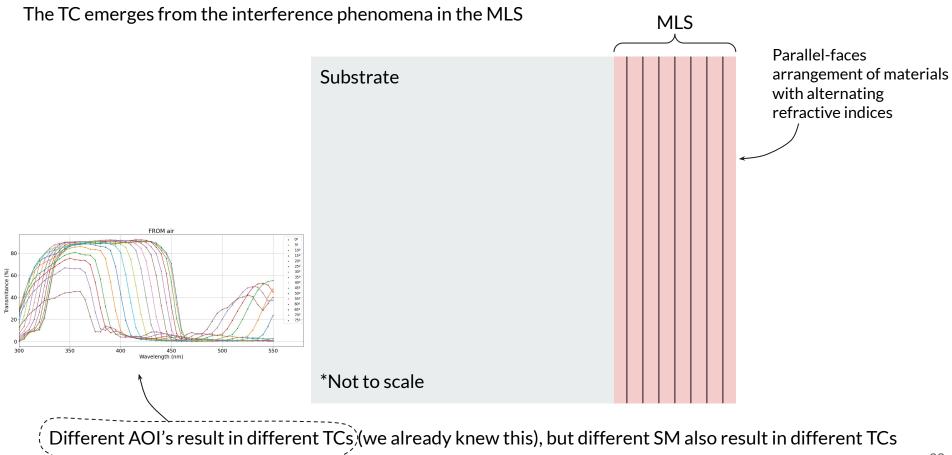
 $I_m(\lambda, heta_{sm}) = I_o(\lambda) \cdot FT^{n_{sm}}_{n_{su}}(heta_{sm}) \cdot ITC(\lambda, heta_{sm}) \cdot FT^{n_{DF}}_{n_{sm}}(lpha(heta_{sm}))$

$$ITC(\lambda, \theta_{sm}) = \frac{T_m(\lambda, \theta_{sm})}{FT_{n_{su}}^{n_{sm}}(\theta_{sm}) \cdot FT_{n_{sm}}^{n_{DF}}(\alpha(\theta_{sm}))}$$
If we knew everything here, we could compute the 'intrinsic' transmitance curve of the dichroic filter as shown above
$$I_m(\lambda, \theta_{sm}) = I_o(\lambda) \cdot FT_{n_{su}}^{n_{sm}}(\theta_{sm}) \cdot ITC(\lambda, \theta_{sm}) \cdot FT_{n_{sm}}^{n_{DF}}(\alpha(\theta_{sm}))$$

$$ITC(\lambda, heta_{sm}) = rac{T_m(\lambda, heta_{sm})}{FT_{n_{su}}^{n_{sm}}(heta_{sm})\cdot FT_{n_{sm}}^{n_{DF}}(lpha(heta_{sm})))}$$

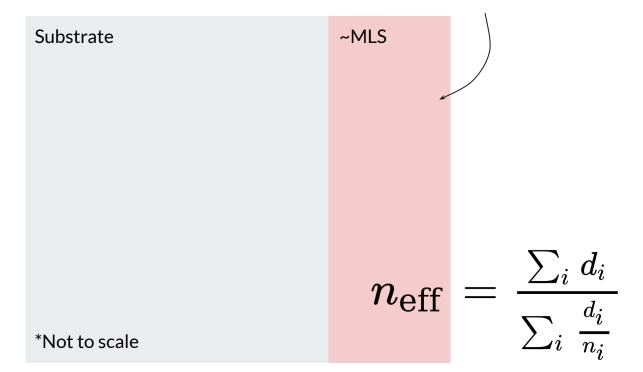
The truth is we know everything except for this, but we can overcome this if we assume the following conceptual model:

DF conceptual model



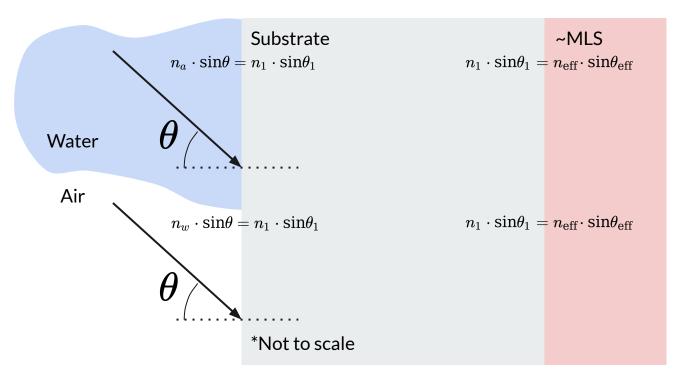
DF conceptual model

 \rightarrow Substitute the MLS by an homogeneous volume with an effective refractive index



Within this model, to explain the dependence of TCs with the surrounding media, our current hypothesis is that **what** determines the TC is the angle of refraction (AOR) within the ~MLS.

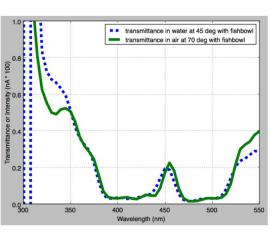
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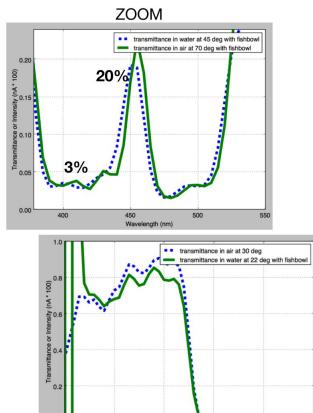


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Model backups

There are **<u>experimental</u>** measurements and Transfer-matrix-method simulations backing up this hypothesis:





400

Wavelength (nm)

450

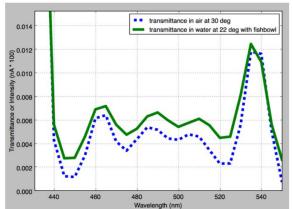
500

550

0.0

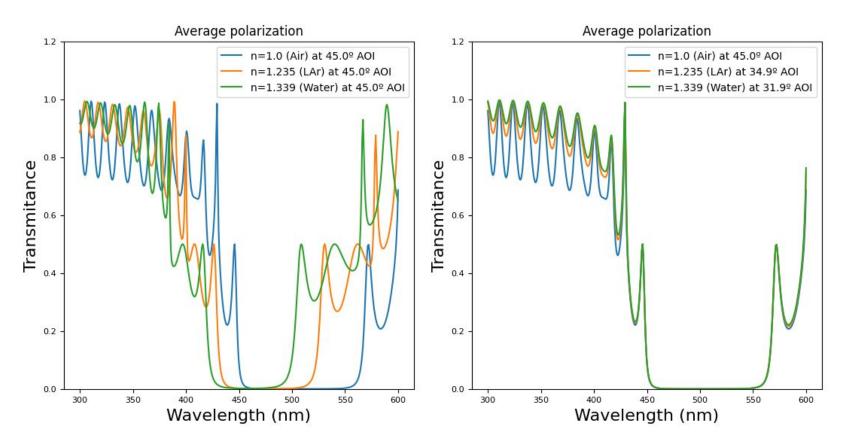
350





Model backups

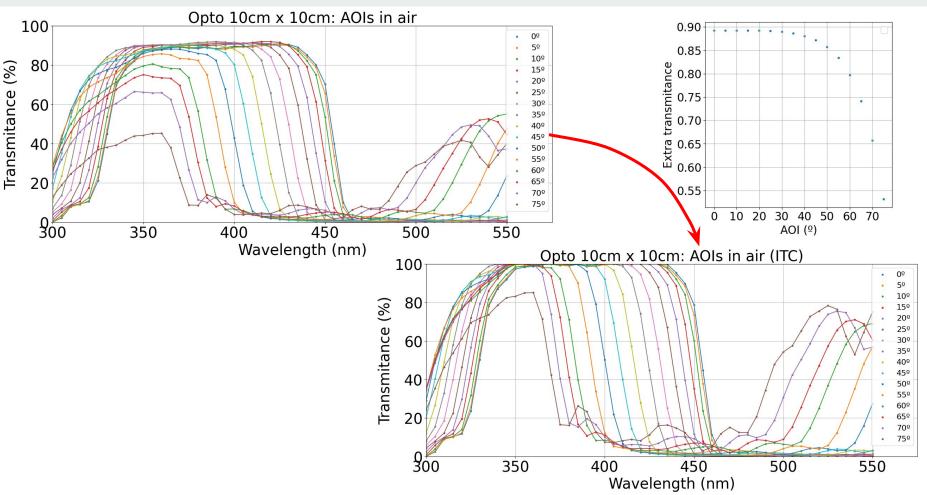
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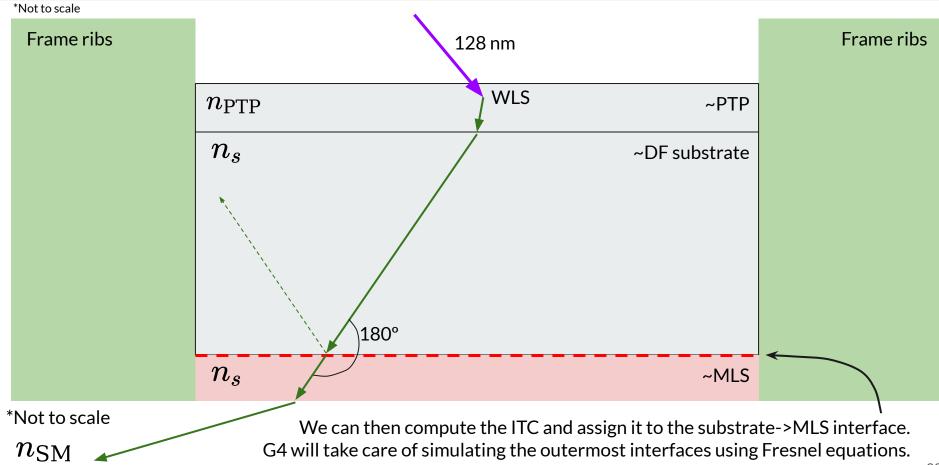
$$ITC(\lambda, heta_{sm}) = rac{T_m(\lambda, heta_{sm})}{FT_{n_{su}}^{n_{sm}}(heta_{sm})\cdot FT_{n_{sm}}^{n_{DF}}(lpha(heta_{sm})))}$$

Within this approximation, we know what the refractive index of the dichroic filter is (the effective r. index of the MLS) and we can compute the angle of the light that propagates within the ~MLS using Snell's law

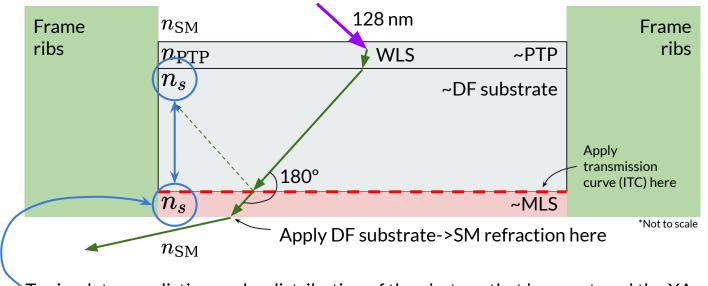
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ight)}
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DF computational model

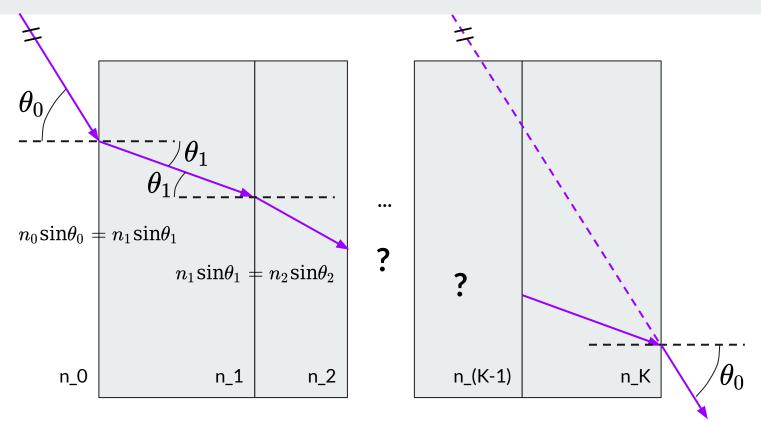


Second caveat

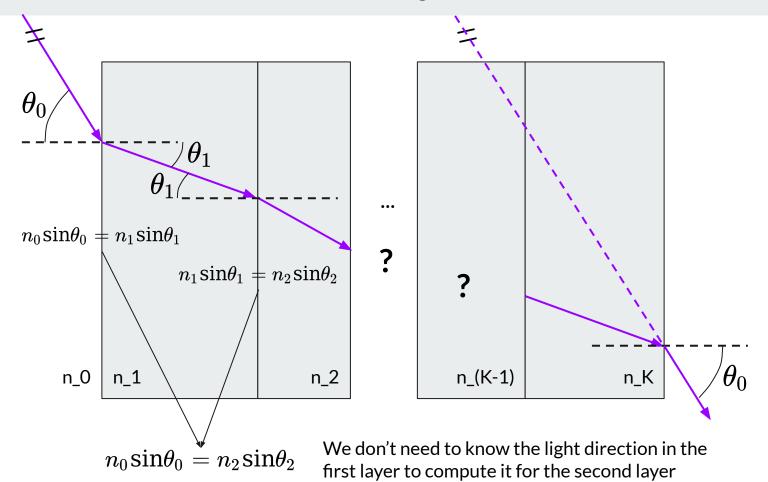


To simulate a realistic angular distribution of the photons that have entered the XA, we need to set the ~MLS volume refractive index to that of the substrate

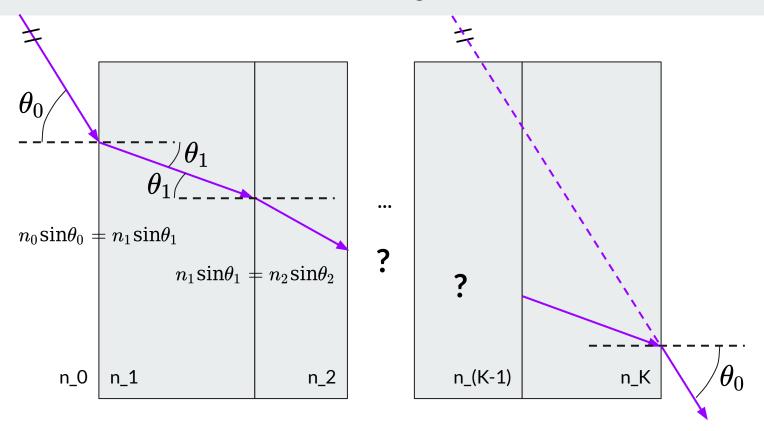
Previous considerations on parallel-faces stackings



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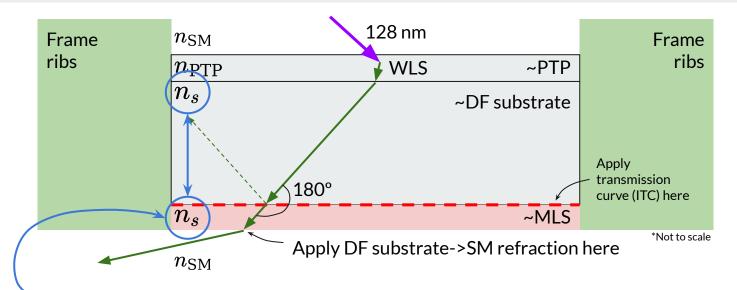


Previous considerations on parallel-faces stackings



 $n_i \sin \theta_i = \text{cte} \ \forall i$ Indeed, we only need to know the pair (theta_i, n_i) for some layer, to know the light propagation direction in any other layer

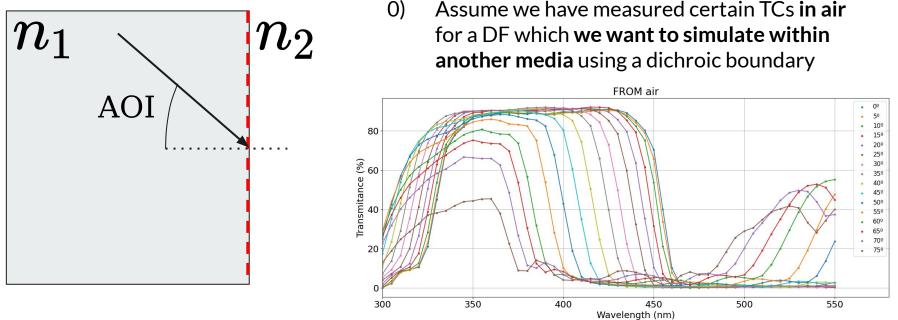
Second caveat



Since no refraction is simulated in the ITC interface, the angular distribution of photons within the ~MLS volume is that of the photons in the substrate volume (this is a simulation artifact). To prevent this artifact from spoiling a realistic simulation of the angular distribution of photons in LAr, the refraction in the ~MLS->LAr interface must account for the substrate refractive index.

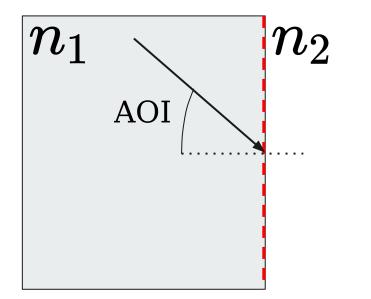
In G4, the TC is applied at some boundary, so there's actually no volume where an AOR is defined. I.e. when a photon reaches the dichroic boundary, G4 won't have access to any AOR information: The only information that G4 counts on to decide which TC to apply is the AOI

Assuming the AOR model, a way to proceed is the following one:



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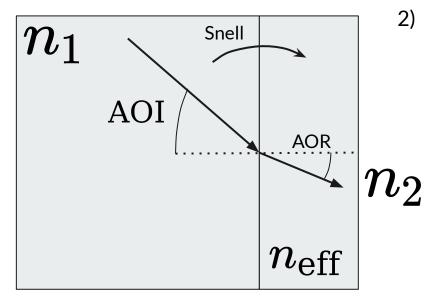
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1) A photon impinges on the dichroic boundary with a certain AOI theta, from a media with refractive index n_1

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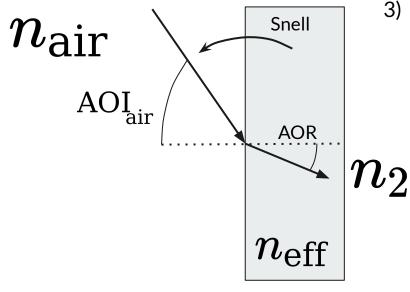


Ask yourself, what would the AOR be if there was actually a volume with a refractive index equal to the MLS effective refractive index?

I.e. compute "forward" snell's law

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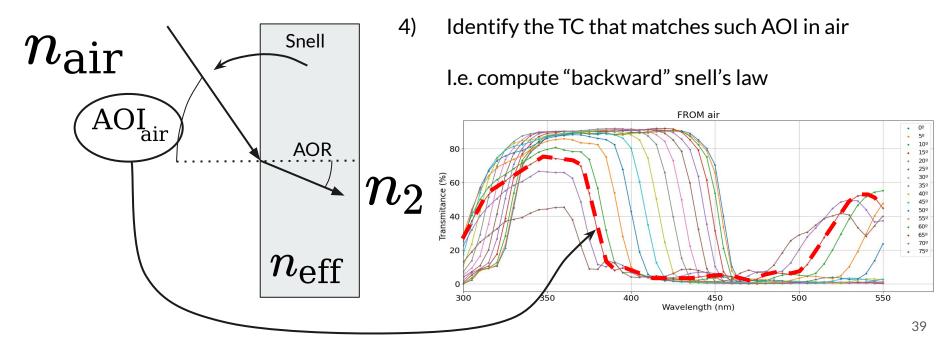


Ask yourself, what AOI in air would give rise to such AOR?

I.e. compute "backward" snell's law

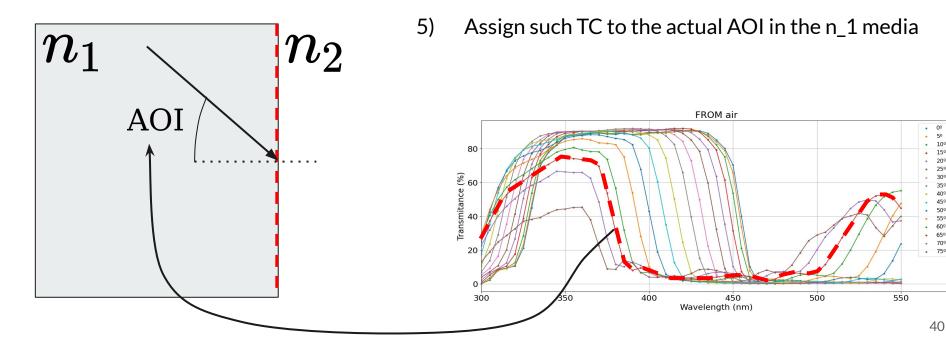
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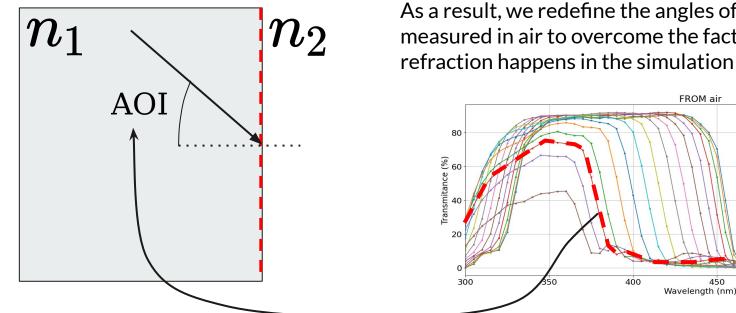
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30 35 40º 45º 50º 55º 60º 65º 70⁰ 75º

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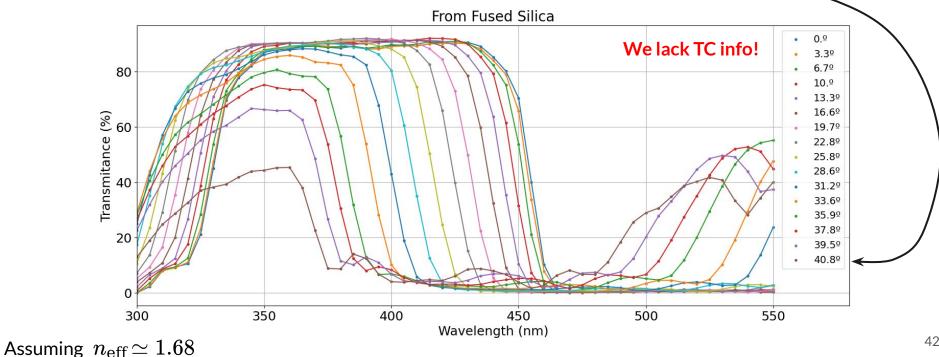
As a result, we redefine the angles of the TCs we measured in air to overcome the fact that no MLS refraction happens in the simulation



500

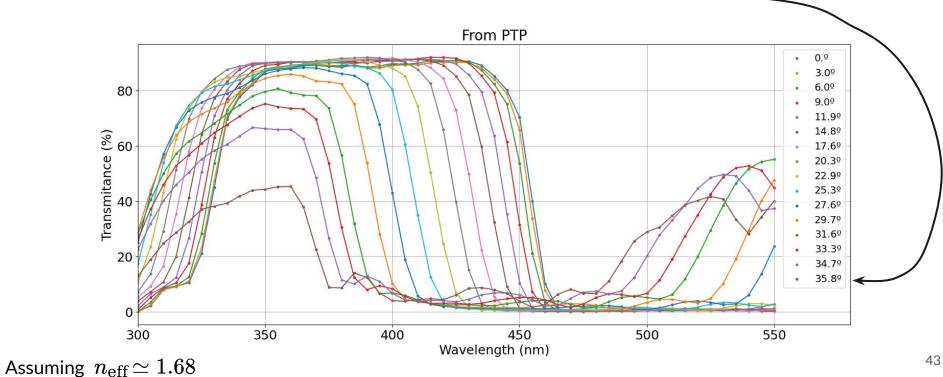
30 35 45º 50º 55º 60º 65º 70⁰

Assume we have measured a DF in air from 0° to 75° AOI, and assume we are modelling the DF using the first approach, where photons impinge over the dichroic boundary from a fused silica (FS) media ($n \approx 1.47$). Then, this method assigns the 0° TC curve in air to the 0° TC curve in FS, and the 75° TC curve in air to the 41° TC curve in FS

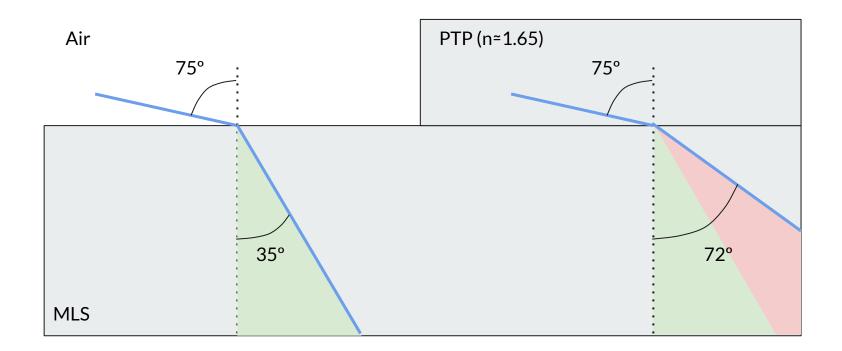


*The information loss is actually the same in both cases, if we assume that for both cases, the incoming light was emitted by PTP with a fixed AD

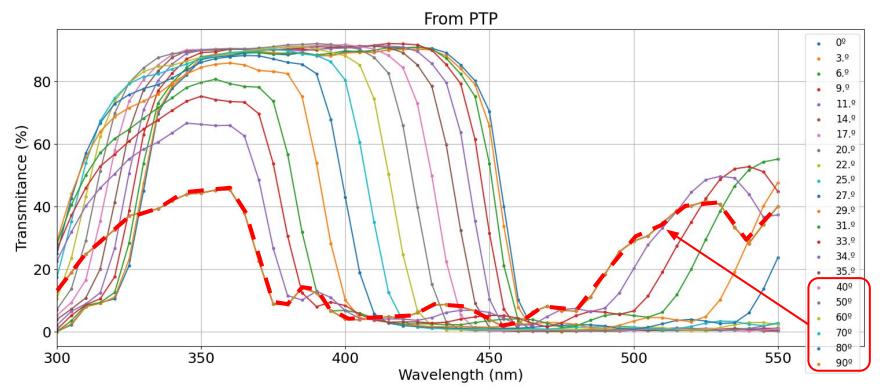
If we want to simulate a PTP (n~1.65) layer right above the dichroic boundary, the information loss is even worse*



The underlying cause is that sweeping 0-75° AOI in air does not let us explore the full range of AOR that refraction from PTP gives:



The most immediate 'solution' is to extrapolate the nearest (angle-wise) TC to every unknown value:



Since the bigger the AOI, the worse DF performance, extrapolating the biggest-AOI available TC to bigger angles sets an upper bound to the DF performance for AOIs whose TC we have not measured yet.

Measuring the TCs within water would widen our AOR scope

