

# Searching for BSM in the DUNE ND

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**DUNE-ND Coll. meeting  
IC London**

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Horizon2020

# Evidence beyond the SM

There is evidence that the Standard Model is incomplete.

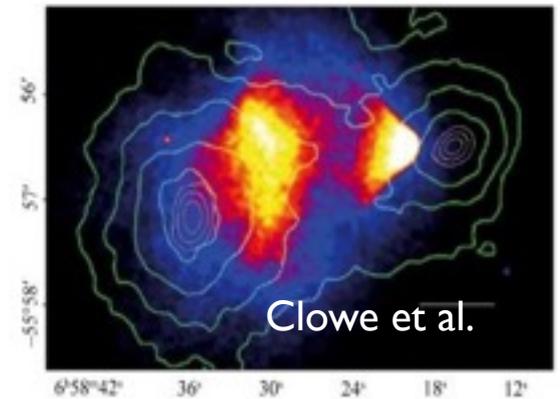
*Neutrino masses*



*Dark Matter*



*Baryon asymmetry*



*What is the new physics  
scale?*

*Are there new:  
symmetries?  
particles?  
interactions?*

# New physics scale: High Energy frontier and above

eV

keV

MeV

GeV

TeV

Intermediate scale

GUT scale

*At TeV SUSY MSSM, split SUSY, Composite Higgs, Technicolor....*

LHC new particle searches

Leptogenesis/EW baryogenesis

LLP searches

WIMPs

CLFV

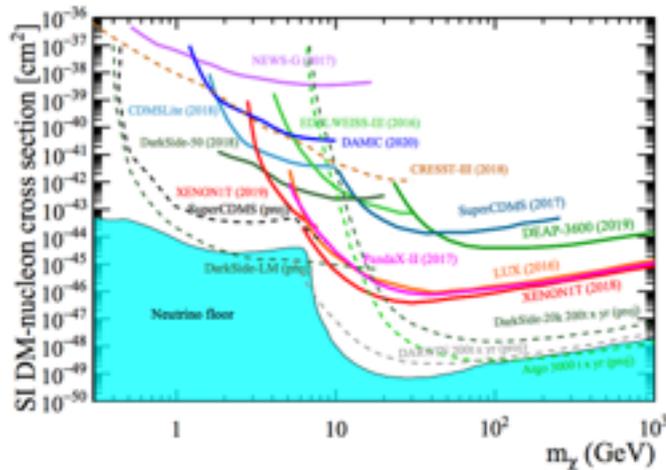
*GUT theories*

Proton decay

Neutrino masses via see-saw Type I

Leptogenesis

GW



Despite intense searches in colliders, flavour and DM exp,  
**no hints of TeV new physics have been found.**

# Going low in energy: Dark sectors

A change of paradigm might be needed: new physics may be light but hidden because too weakly interacting (dark or hidden sectors).



*Low E See-saw models, NuMSM, extended see-saw...*

Sterile nu  
oscillations

light DM  
Leptogenesis

HNL searches: peak, kinks, decays,

Recently, a strong theoretical effort has been done together with a blooming of experimental opportunities at the intensity frontier. DUNE ND inscribes in this thriving context.

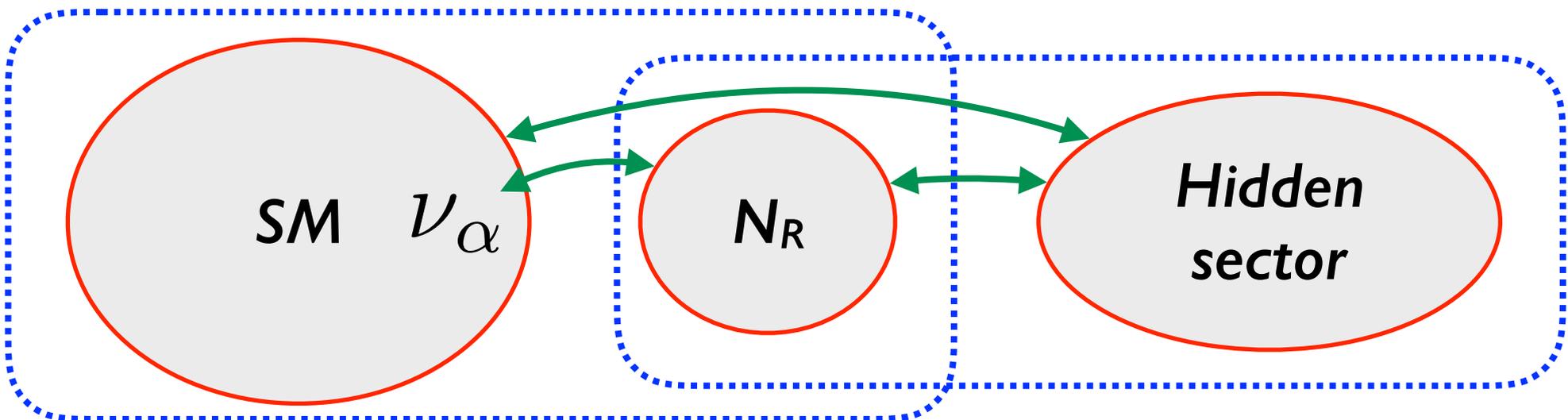
See e.g. Artuso et al., 2210.04765, S. Gori et al, 2209.04671, FIPs 2023 report, 2305.01715

The dark sector can interact with SM via portals:

- the kinetic mixing portal (dark photons);
- the scalar portal (dark scalar);
- the neutrino portal (heavy neutral leptons);
- the axion portal (axion);
- millicharged particles.

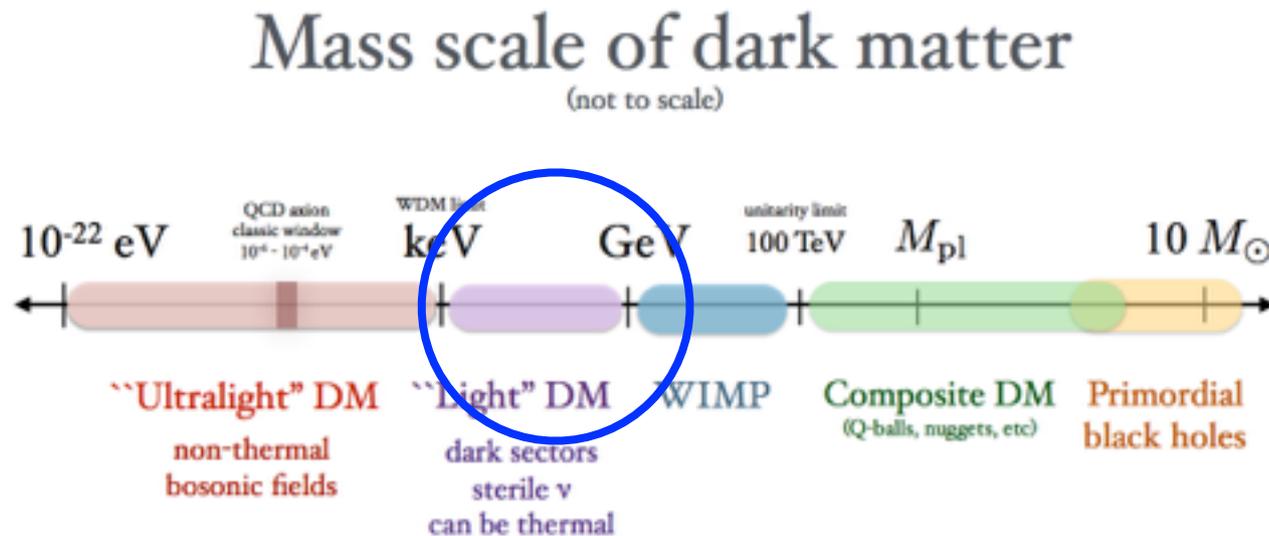
See J.Yoo's talk

Portal	Coupling
Dark Photon, $A'$	$-\frac{\epsilon}{2\cos\theta_W} F'_{\mu\nu} B^{\mu\nu}$
Axion-like particles, $a$	$\frac{a}{f_a} F_{\mu\nu} \tilde{F}^{\mu\nu}, \frac{a}{f_a} G_{i,\mu\nu} \tilde{G}_i^{\mu\nu}, \frac{\partial_\mu a}{f_a} \bar{\psi} \gamma^\mu \gamma^5 \psi$
Dark Higgs, $S$	$(\mu S + \lambda_{HS} S^2) H^\dagger H$
Heavy Neutral Lepton, $N$	$y_N L H N$
millicharged particle, $\chi$	$\epsilon A^\mu \bar{\chi} \gamma_\mu \chi$



# Notable example I: Vector portal and sub-GeV DM

Light gauge extensions of the SM can involve secluded or (B-)leptonic  $U(1)'$ . The associated gauge bosons are called dark photons or  $Z'$ .



T. Lin, 1904.07915

The DS can also contain a light DM particle, whose annihilations can be mediated by the dark photon (or other particle).

## Notable example 2: MeV-GeV see-saw

Sterile neutrinos: hypothetical neutral fermionic singlets of the Standard Model. Generically they mix with the light neutrinos and the mass states are HNL.

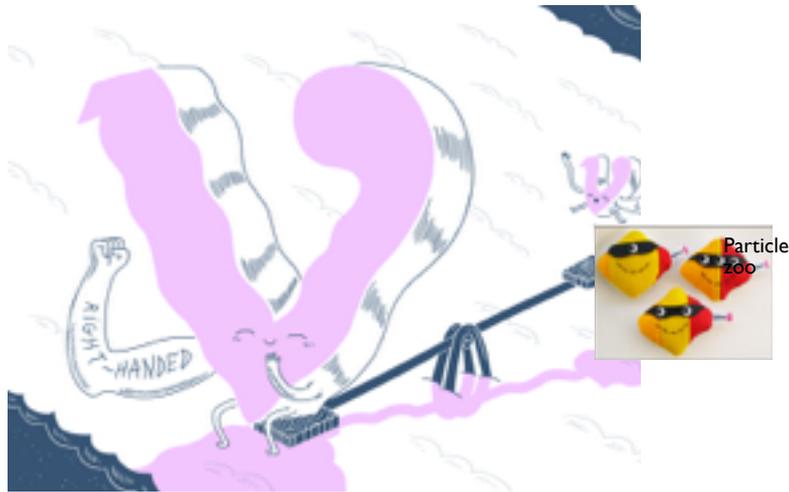
$$\text{Flavour state} \begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \\ \nu_s \end{pmatrix} = U_{4 \times 4} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \\ \nu_4 \end{pmatrix}$$

**Massive state**  
**Nearly-sterile**  
**neutrino, commonly**  
**called sterile neutrino**

$$\mathcal{L} = \dots + \bar{\ell}_L U_{\ell 4} \gamma_\mu \nu_{4,L} W^\mu + \text{NC} + \text{h.c.}$$

HNLs can interact with the SM via mixing and hence be produced and decay into SM particles.

HNLs can be inscribed in see-saw models (type I and extensions, e.g. inverse, linear and extended see saw) and can give raise to neutrino masses.



Symmetry magazine

- Introduce a right handed neutrino **N**
- Couples to the Higgs and has a Majorana mass

$$\mathcal{L} = -Y_\nu \bar{N} L \cdot H - 1/2 \bar{N}^c M_R N$$

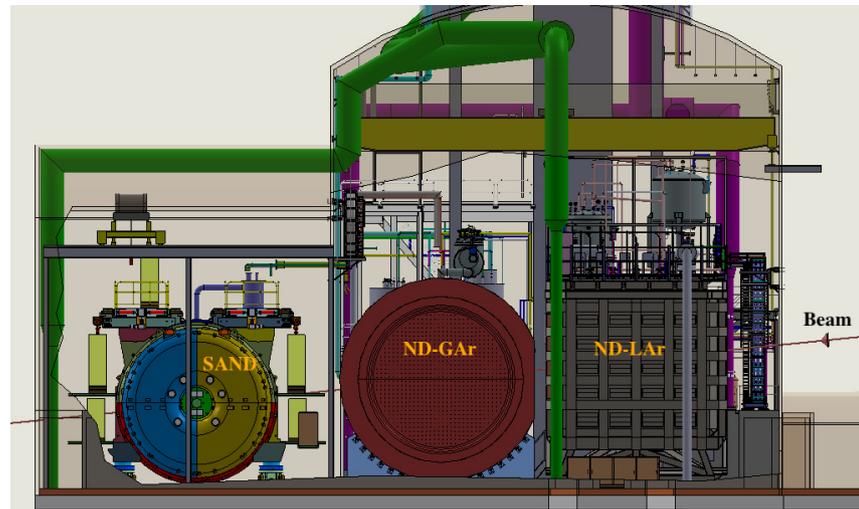
$$\begin{pmatrix} 0 & m_D \\ m_D^T & M_N \end{pmatrix} \rightarrow m_\nu = \frac{Y_\nu^2 v_H^2}{M_N} \sim \frac{1 \text{ GeV}^2}{10^{10} \text{ GeV}} \sim 0.1 \text{ eV}$$

GeV scale See-saw models can also explain the baryon asymmetry via leptogenesis.

# Searching for BSM in DUNE ND

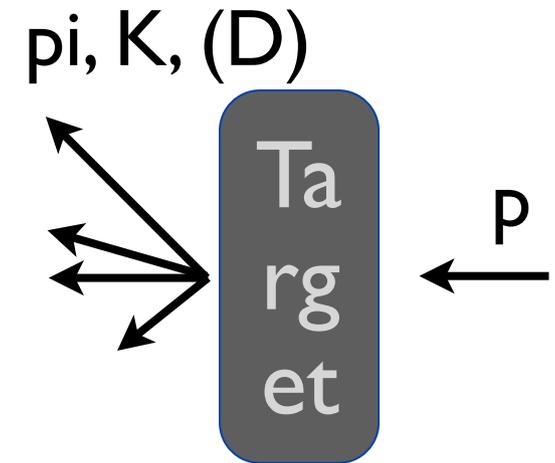
To search directly for BSM in the DUNE ND, one can employ two main exp strategies:

1. a-la beam dump experiments.
2. production (e.g. via NC) and detection in the same detector due to neutrino or other particle (e.g. DM) beam.

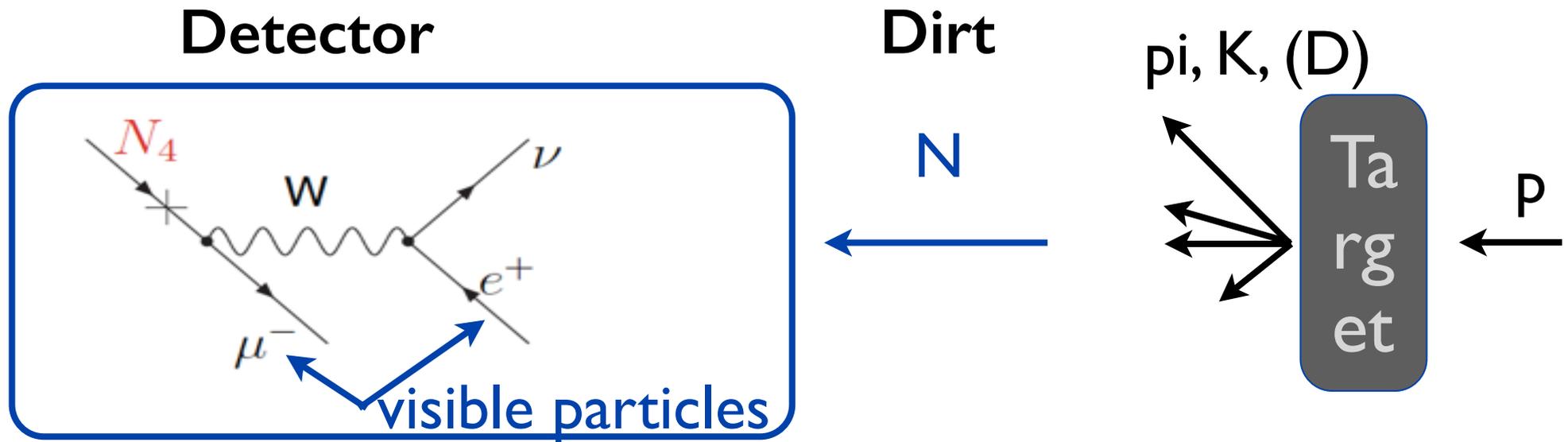


I will refer to the talk presented by M. Bishai yesterday on DUNE ND Phase I and II.

# “A la beam dump” experiment



# “A la beam dump” experiment

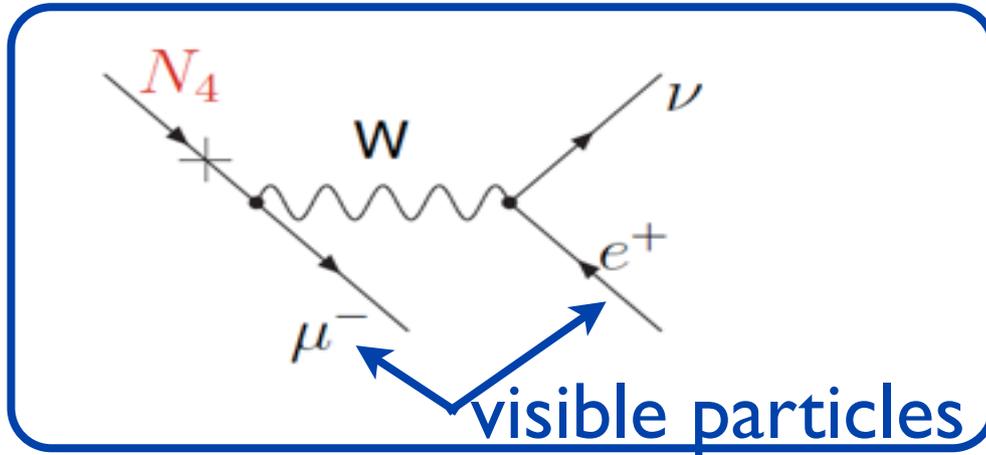


A part from the exp in the 90', in particular PSI91, neutrino accelerator experiments (see recent analysis by T2K) and NA62 can search in this mode.

The typical candidate is a **heavy neutral lepton** or sterile neutrino with  $\sim 100$  MeV mass.

# “A la beam dump” experiment

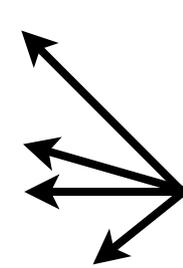
Detector



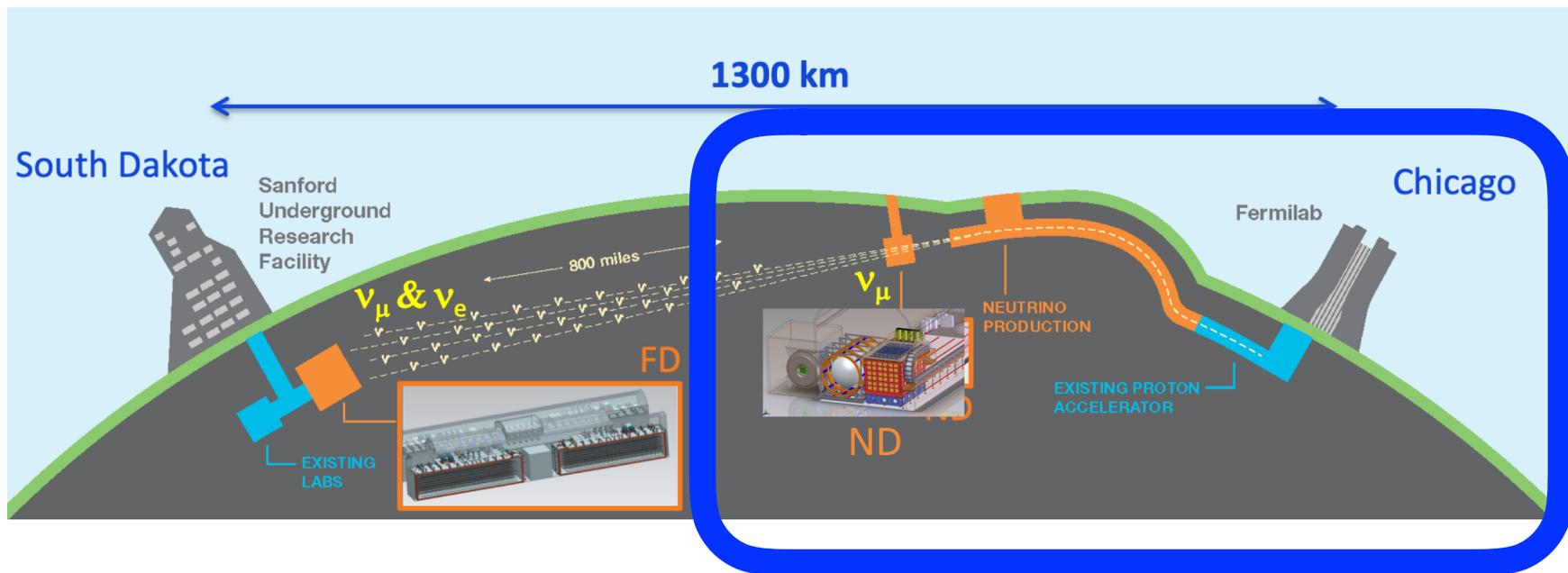
Dirt



$\pi, K, (D)$

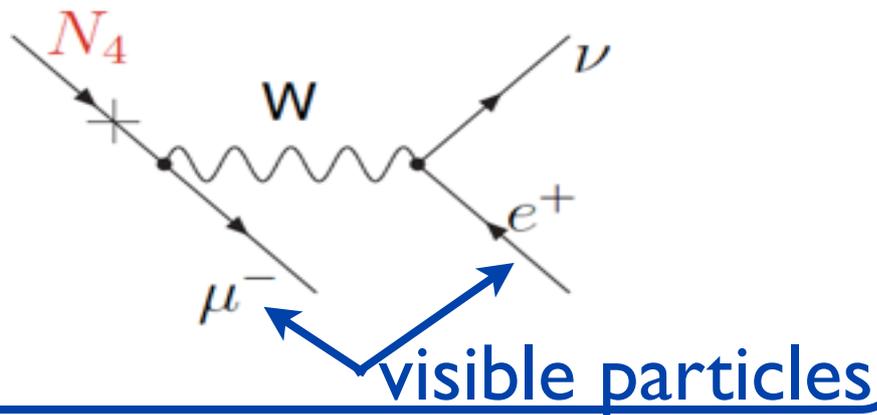


Target



# “A la beam dump” experiment

## Detector



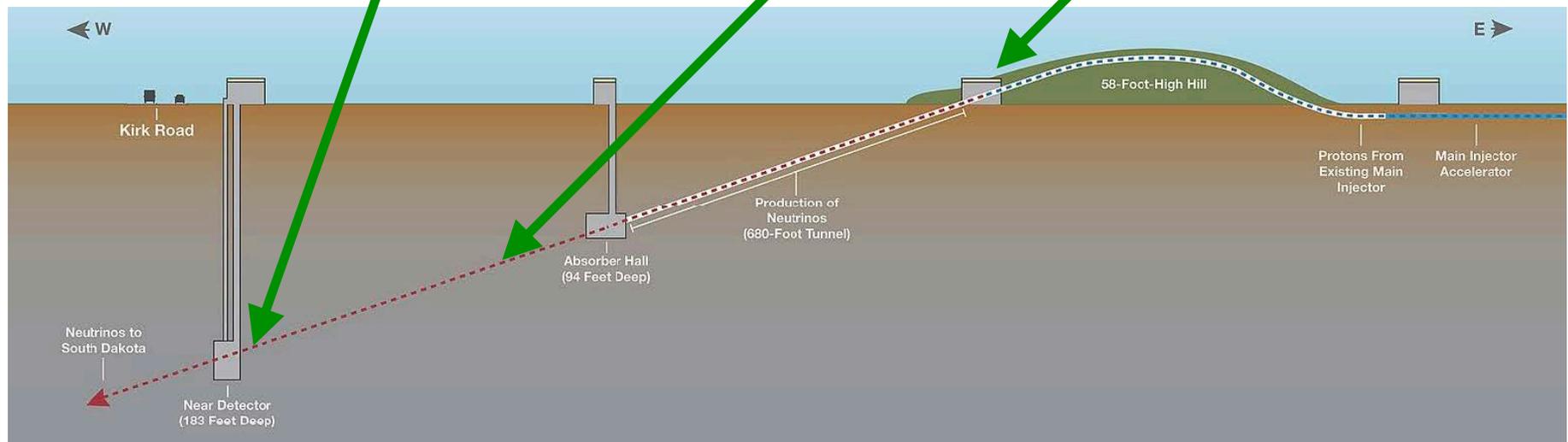
## Dirt

$N$

$\pi, K, (D)$

Target

$P$



## The key characteristic:

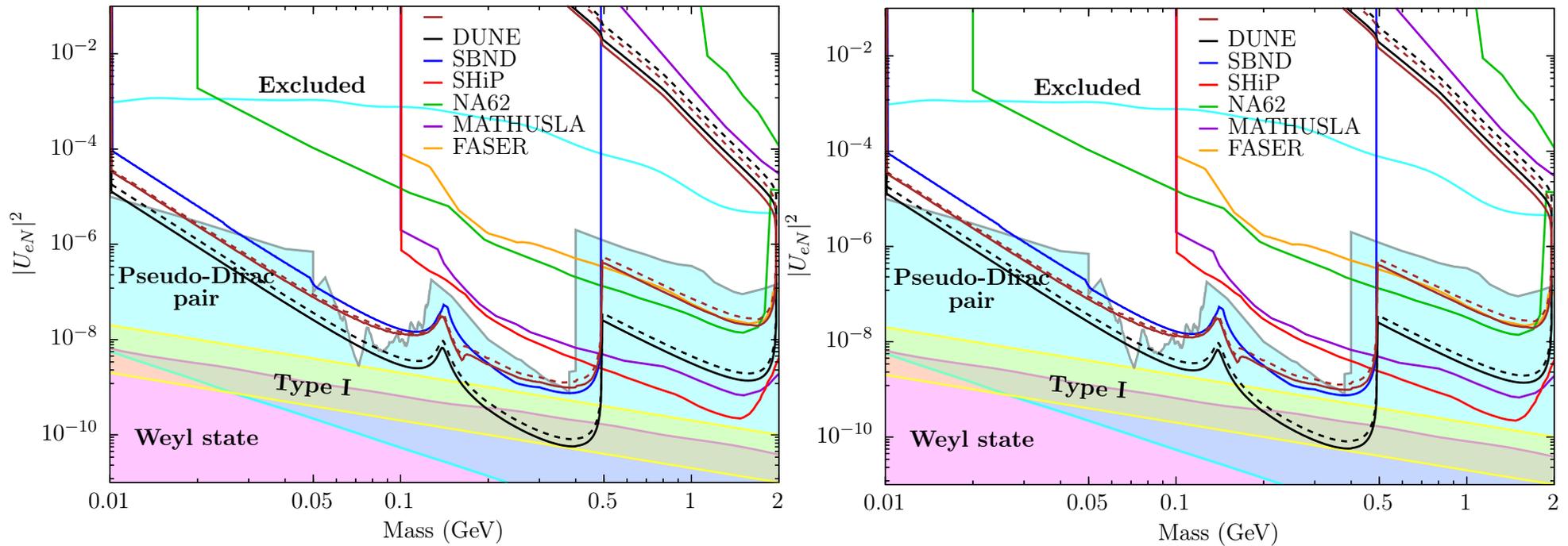
- high proton intensity;
- large volume (not large mass) with excellent reconstruction capabilities (typical backgrounds are gammas,  $\pi^0$ ,  $C\mu^+ l \pi$ );
- (distance from target).

	PS191	DUNE ND	SBND	NA62	SHiP
Baseline	128 m	574 m	110 m	220 m	60 m
Volume	216 m <sup>3</sup>	150 m <sup>3</sup>	80 m <sup>3</sup>	750 m <sup>3</sup>	590 m <sup>3</sup>
Energy	19.2 GeV	80 GeV	8 GeV	400 GeV	400 GeV
POT	$0.86 \times 10^{19}$	$1.32 \times 10^{22}$	$6.6 \times 10^{20}$	$3 \times 10^{18}$	$2 \times 10^{20}$
Exposure	1.0	220.9	16.4	8.5	5820

Ballett, Boschi, SP, I905.00284

**DUNE ND has an exposure few 100 x PS191!**

# Experimental bounds: decays.

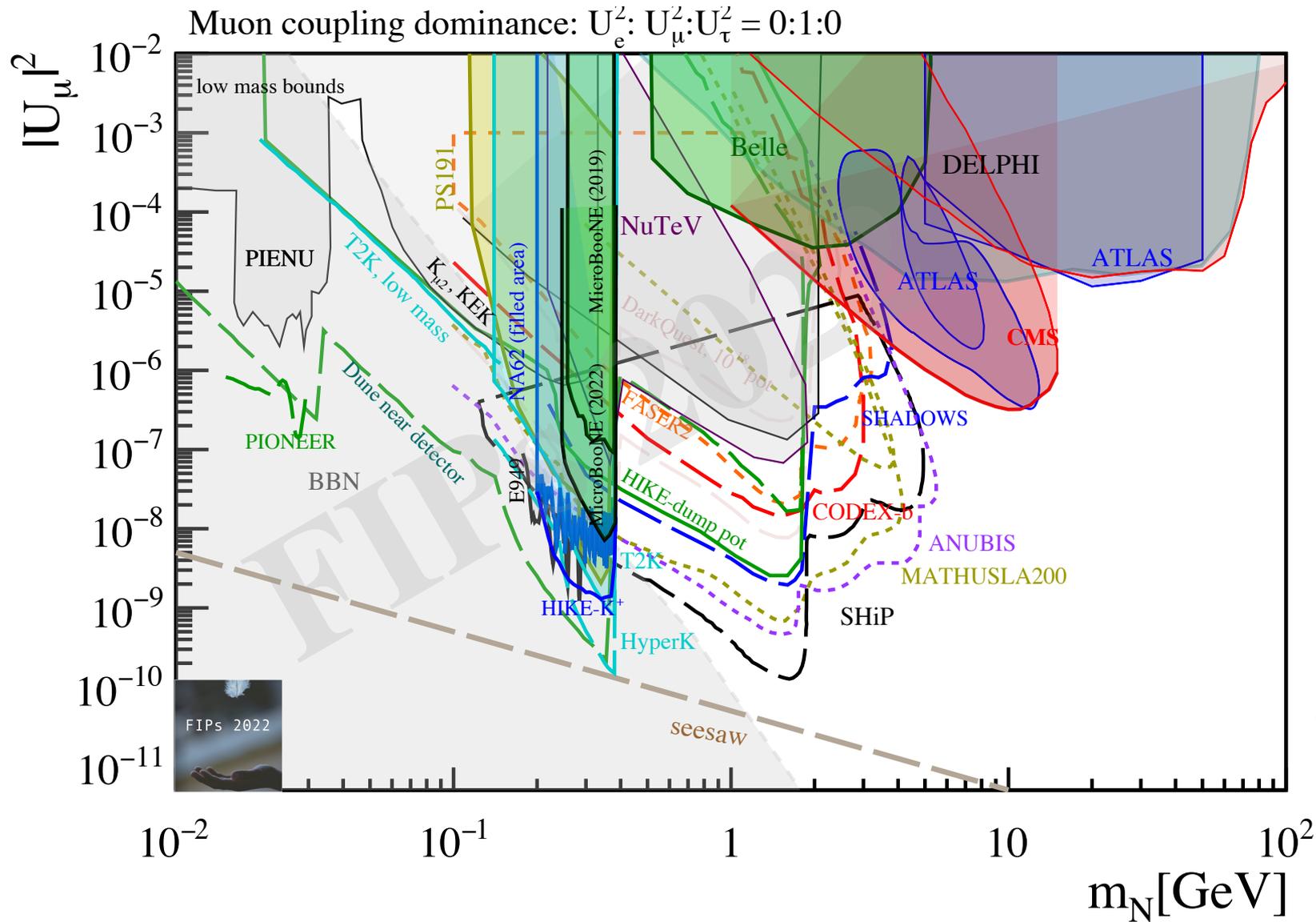


Ballett, Boschi, SP, I 905.00284

**DUNE ND is excellently suited for this type of search and has access to GeV masses.**

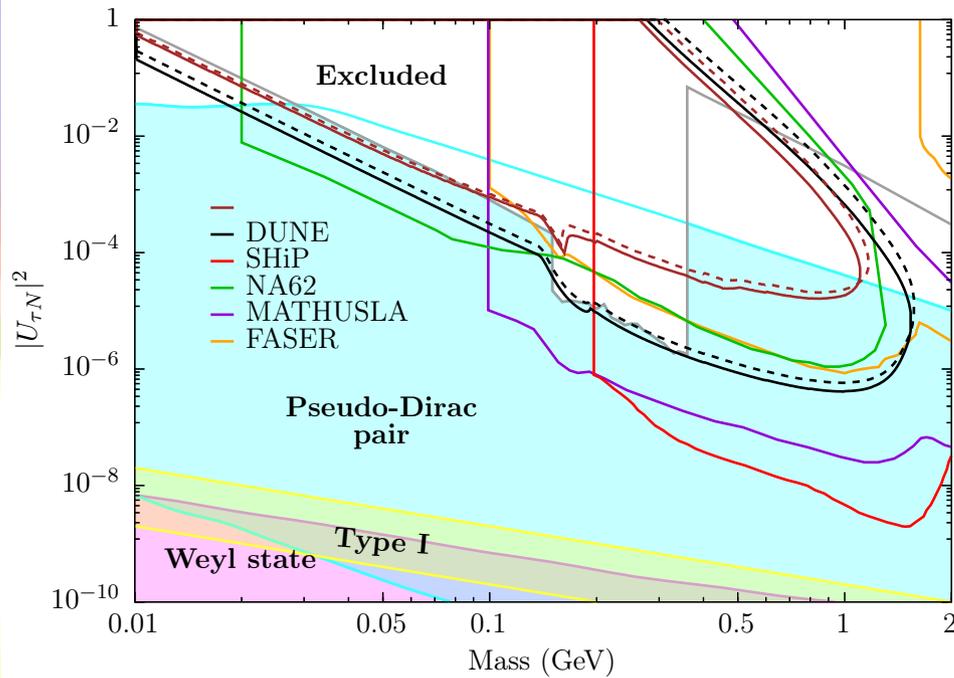
See K. Kelly's and F. Deppisch's talks

# Experimental bounds: decays.

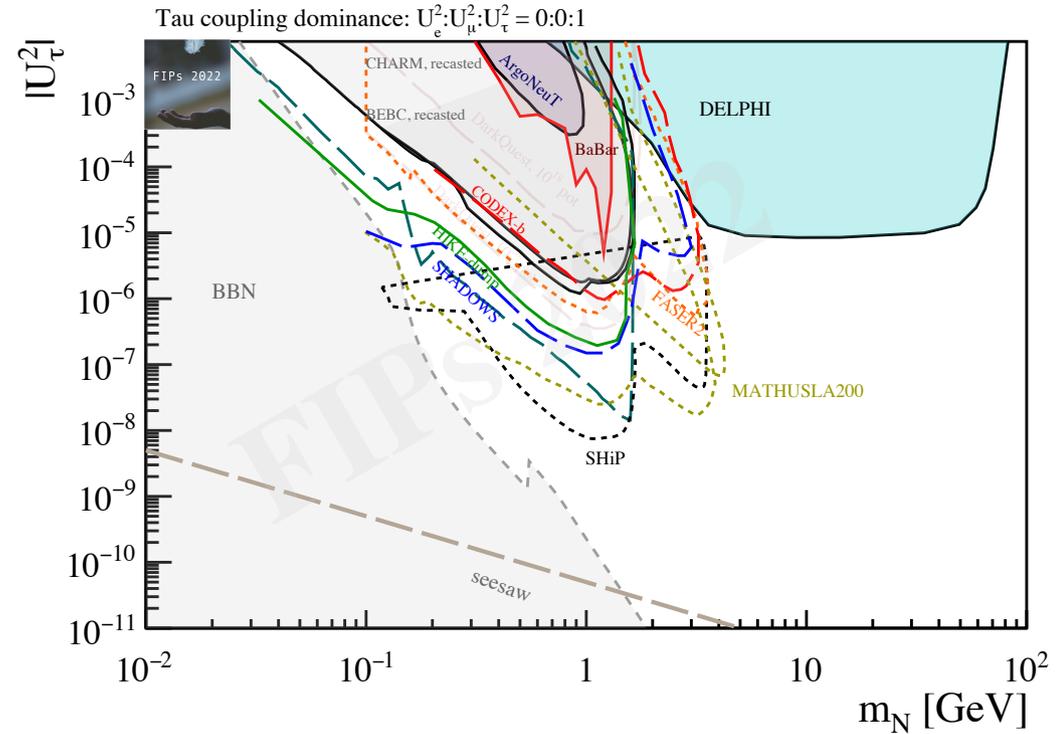


Note: these studies focus on one flavour mixing at a time, but this is not compatible with see-saw type I.

# Experimental bounds: tau mixing.



Ballett, Boschi, SP, 1905.00284

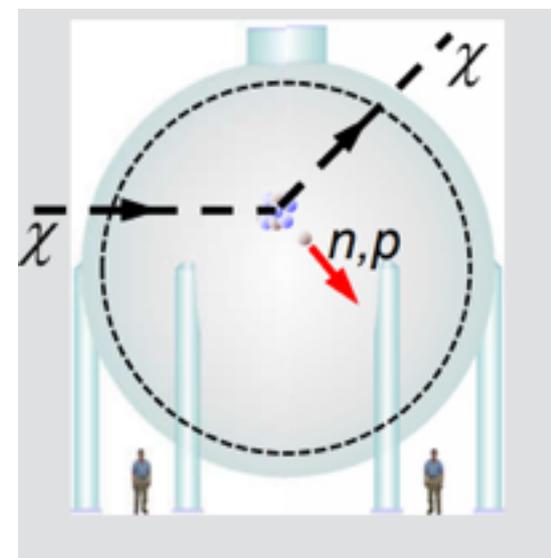
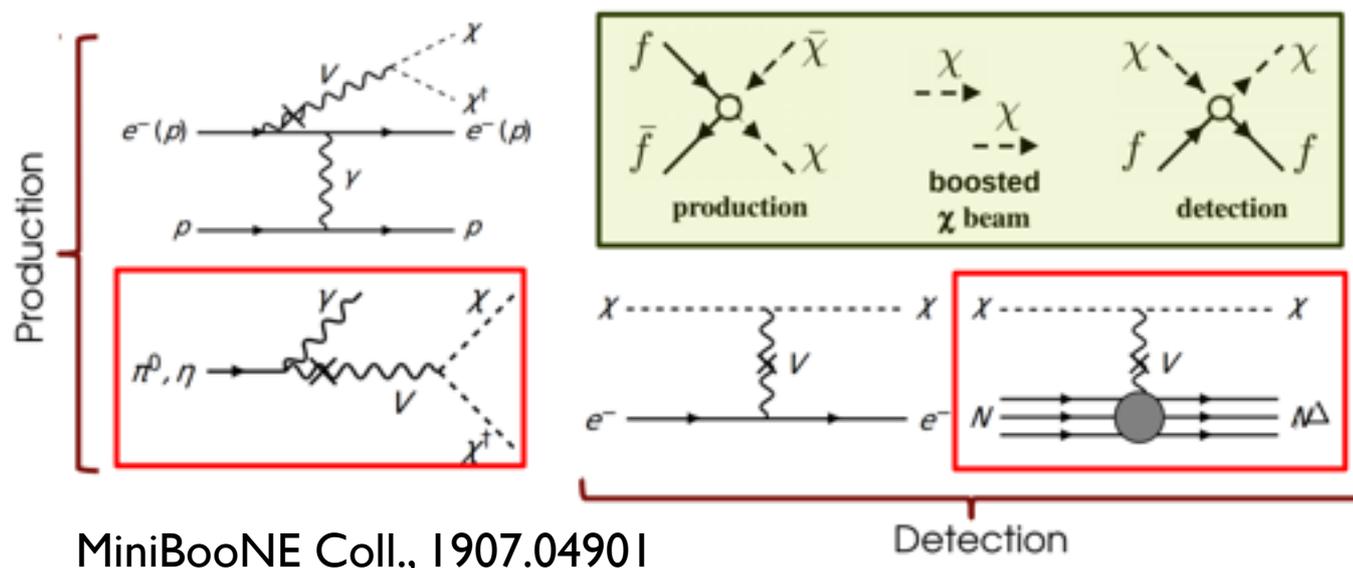


FIPs 2022 report, 2305.01715

Thanks to the  $D_s$  production, DUNE ND will have also sensitivity to the mixing with tau neutrinos.

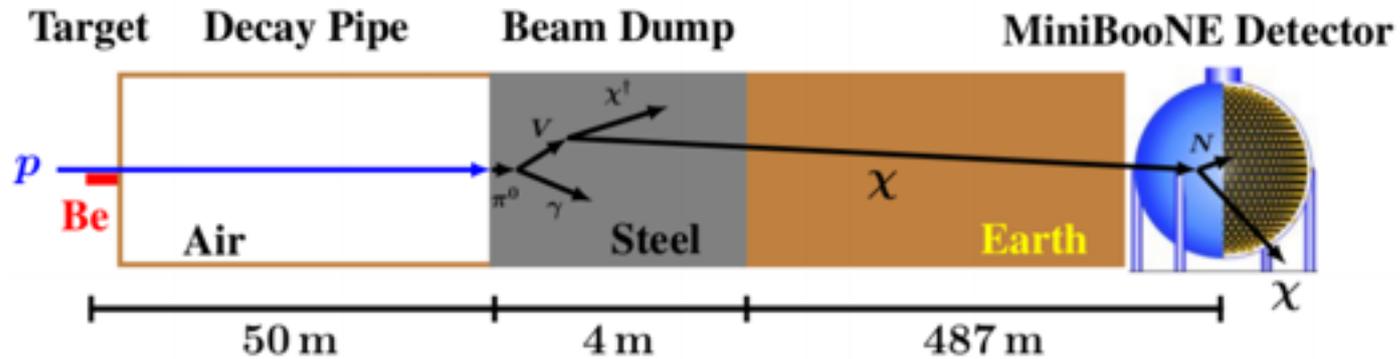
# Neutrino or dark sector scattering

If dark sector particles (e.g. LDM) interacts with the SM, it can be produced in beam dump experiments (directly or via subsequent DS particle decays).



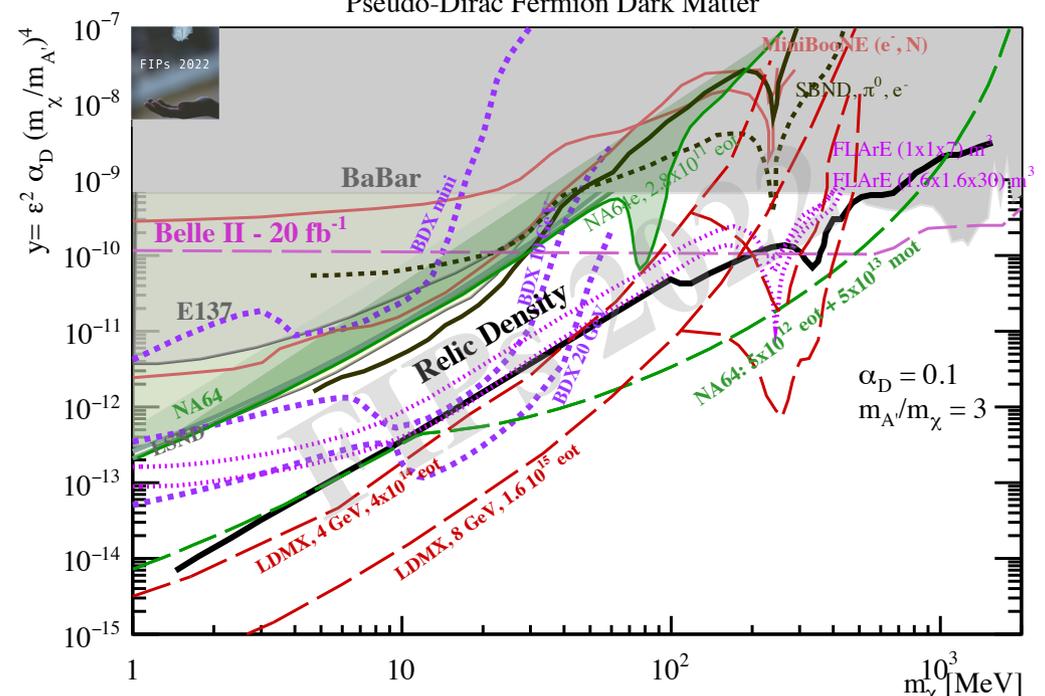
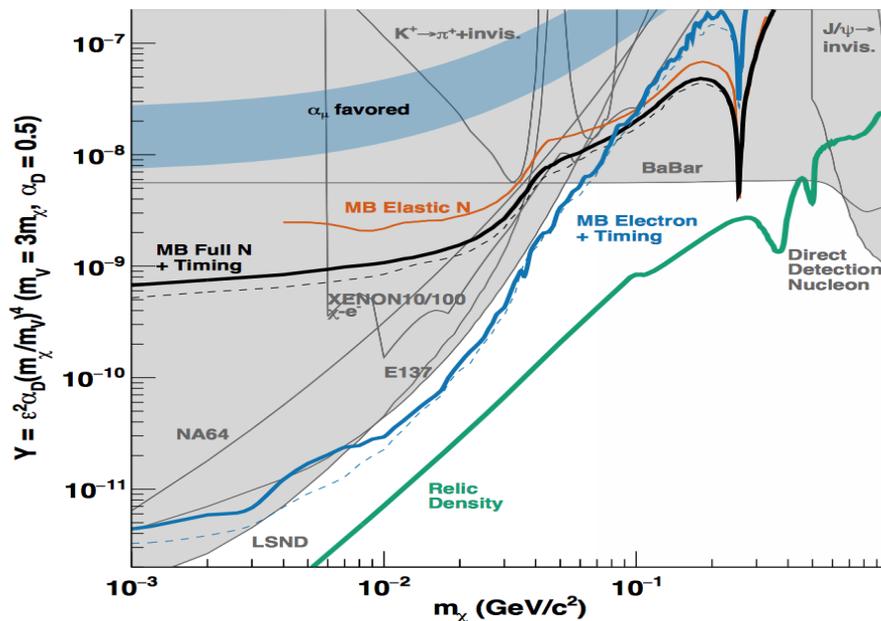
Detection proceeds via long lived Dark Sector particle scattering in the detector, e.g. nucleon elastic scattering with recoil energy  $\sim 100$  MeV (similar to neutrino-nucleon NC elastic scattering).

# MiniBooNE has done a search both with nucleon and electron recoils.



MiniBooNE Coll., PRL 118 (2017)

Pseudo-Dirac Fermion Dark Matter



MiniBooNE DM Collaboration, 1807.06137

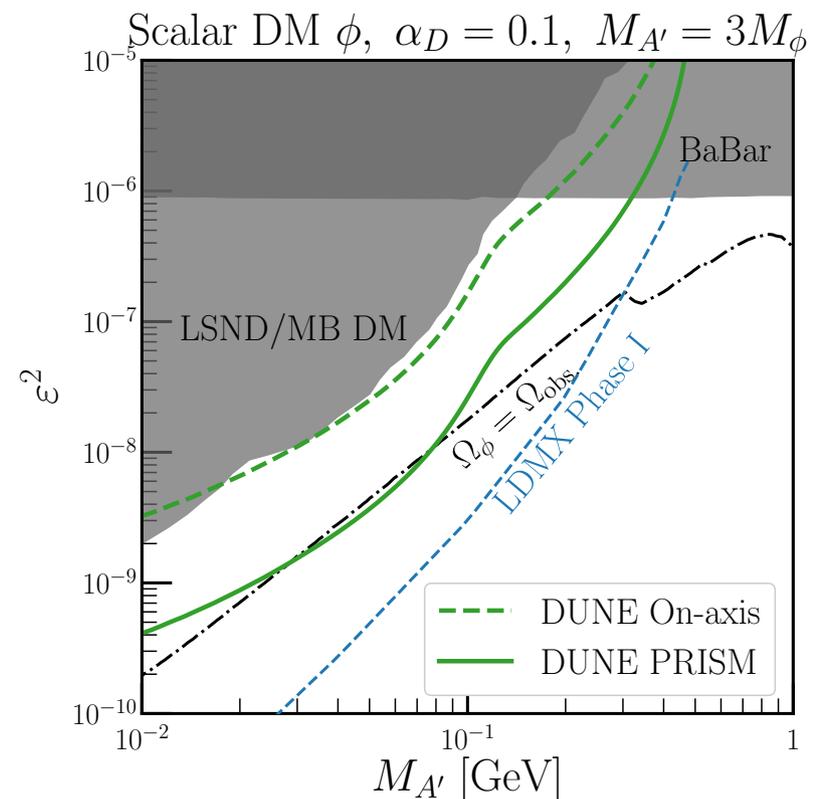
FIPs 2022 report, 2305.01715

@Silvia Pascoli

## Key characteristics:

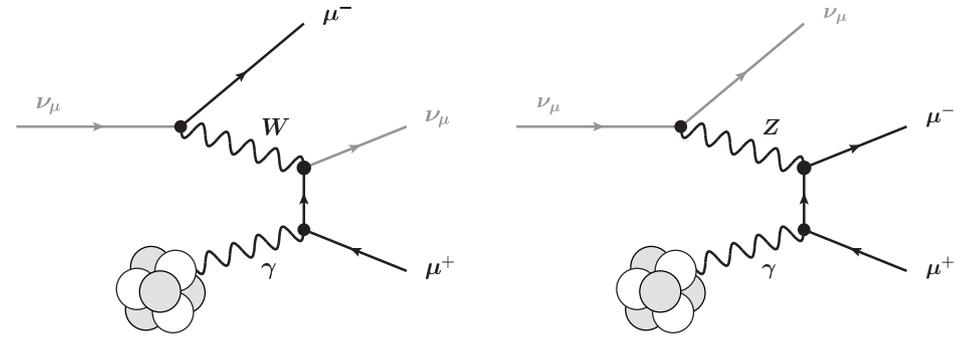
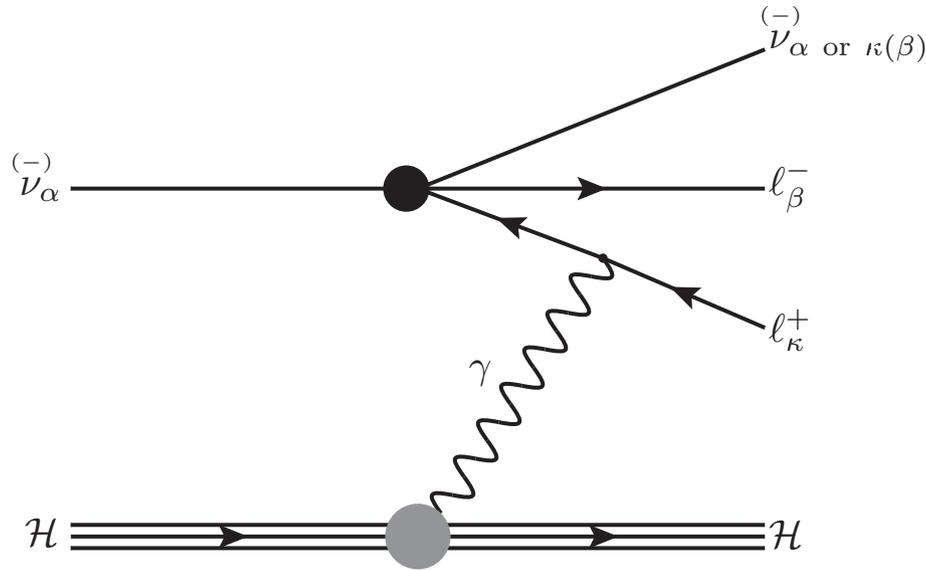
- **beam intensity**
- **detector mass (needed for scattering)**
- **position on/off-axis for bkgd reduction (DUNE Prism) due to different DM flux.**

Dominant background is typically due to neutrino interactions (e.g. QE NC events, nu scattering off electrons and QE CC electron scattering events).



De Romeri, Kelly, Machado, 1903.10505. See also, Breitbach et al., 2102.03383, Brdar et al., 2206.06380

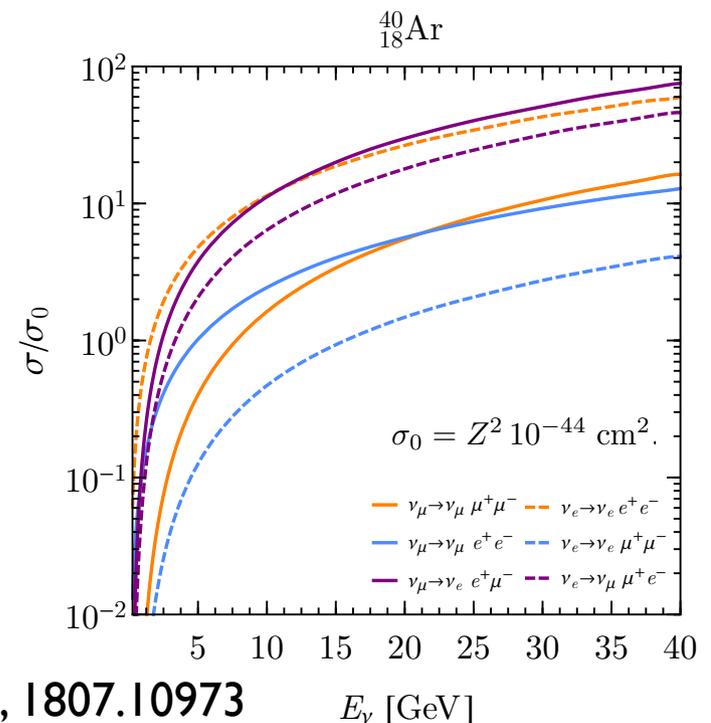
# Trident



Altmannshofer et al., 1902.06765

Trident processes are expected in the SM but are very suppressed.

$$\frac{\sigma(\nu_\mu \rightarrow \nu_\mu \mu^+ \mu^-)_{\text{exp}}}{\sigma(\nu_\mu \rightarrow \nu_\mu \mu^+ \mu^-)_{\text{SM}}} = \begin{cases} 1.58 \pm 0.64 & (\text{CHARM-II}) \\ 0.82 \pm 0.28 & (\text{CCFR}) \\ 0.72^{+1.73}_{-0.72} & (\text{NuTeV}) \end{cases}$$



Ballett et al., 1807.10973

$E_\nu$  [GeV]

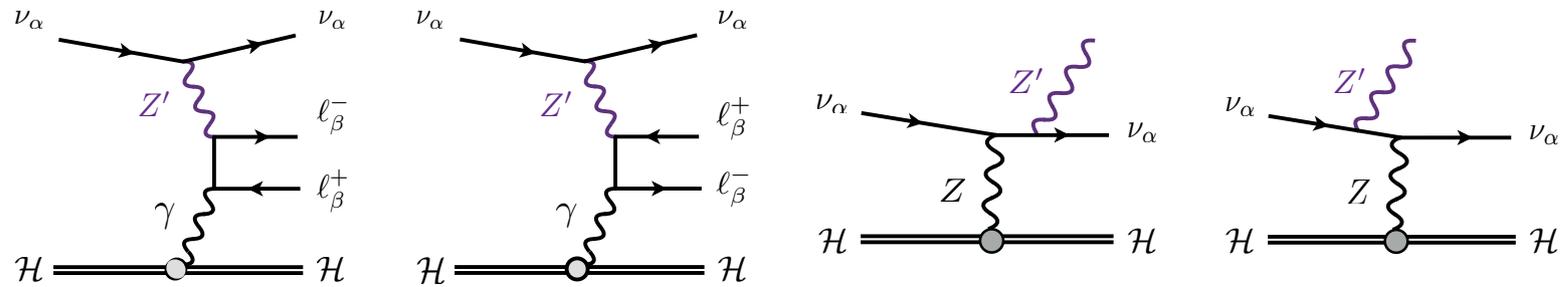
# Neutrino accelerator experiments offer new opportunities to discover them.

Experiment	Baseline (m)	Total Exposure (POT)	Fiducial Mass (t)	$E_\nu$ (GeV)
SBND	110	$6.6 \times 10^{20}$	112	0 – 3
$\mu$ BooNE	470	$1.32 \times 10^{21}$	89	0 – 3
ICARUS	600	$6.6 \times 10^{20}$	476	0 – 3
DUNE	574	$12.81 (12.81) \times 10^{21}$	50	0 – 40
$\nu$ STORM	50	$10^{21}$	100	0 – 6

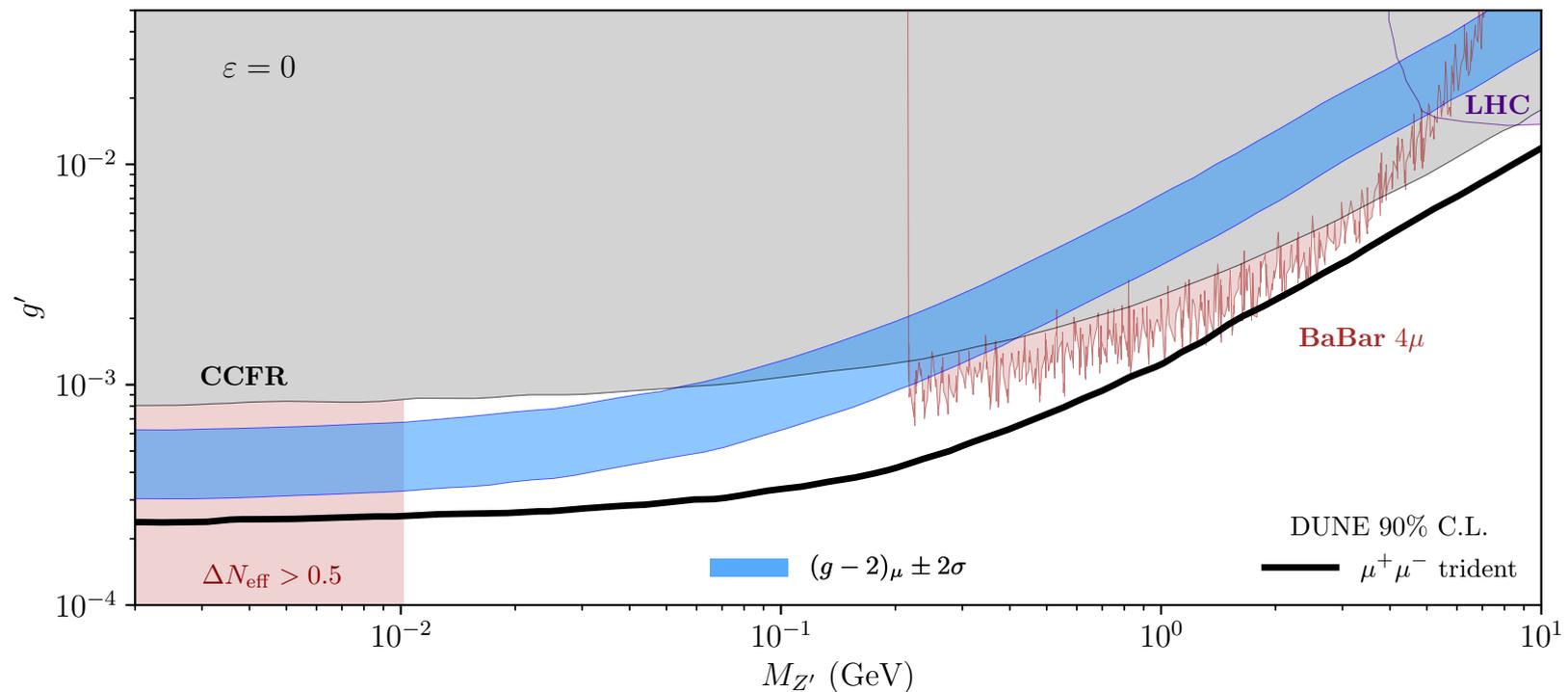
It is a very rare process with a cross section which increases with energy. DUNE is well suited for this search.

$\nu_\mu \rightarrow \nu_\mu e^+ e^-$	6	0.4	0.7	913 (58)	73
	0.7	0.04	0.07	128 (9)	9
$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu e^- e^+$	0.2	0.01	0.02	34 (695)	9
	0.03	0.001	0.002	5 (95)	1
$\nu_e \rightarrow \nu_e e^- e^+$	0.2	0.01	0.02	50 (13)	32
	0.02	0.001	0.002	8 (2)	4
$\bar{\nu}_e \rightarrow \bar{\nu}_e e^+ e^-$	0.02	0.001	0.002	10 (34)	–
	0.003	0.0001	0.0002	2 (5)	–
Total $e^+ e^-$	6	0.4	0.7	1007 (800)	114
	0.7	0.0	0.1	143 (111)	14
$\nu_\mu \rightarrow \nu_\mu \mu^+ \mu^-$	0.4	0.03	0.04	271 (32)	9
	0.4	0.03	0.04	186 (19)	8
$\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu \mu^- \mu^+$	0.01	0.001	0.001	14 (177)	2
	0.01	0.0009	0.001	9 (127)	1
$\nu_e \rightarrow \nu_e \mu^+ \mu^-$	0.002	0.0001	0.0001	1 (0.5)	0.4
	0.001	0.0001	0.0001	0.7 (0.2)	0.3
$\bar{\nu}_e \rightarrow \bar{\nu}_e \mu^+ \mu^-$	0.0002	0.0000	0.0000	0.3 (0.9)	–
	0.0001	0.0000	0.0000	0.2 (0.5)	–
Total $\mu^+ \mu^-$	0.4	0.0	0.0	286 (210)	11
	0.4	0.0	0.0	196 (147)	9

# Trident events can be used to search for light vector boson mediators ( $Z'$ ).

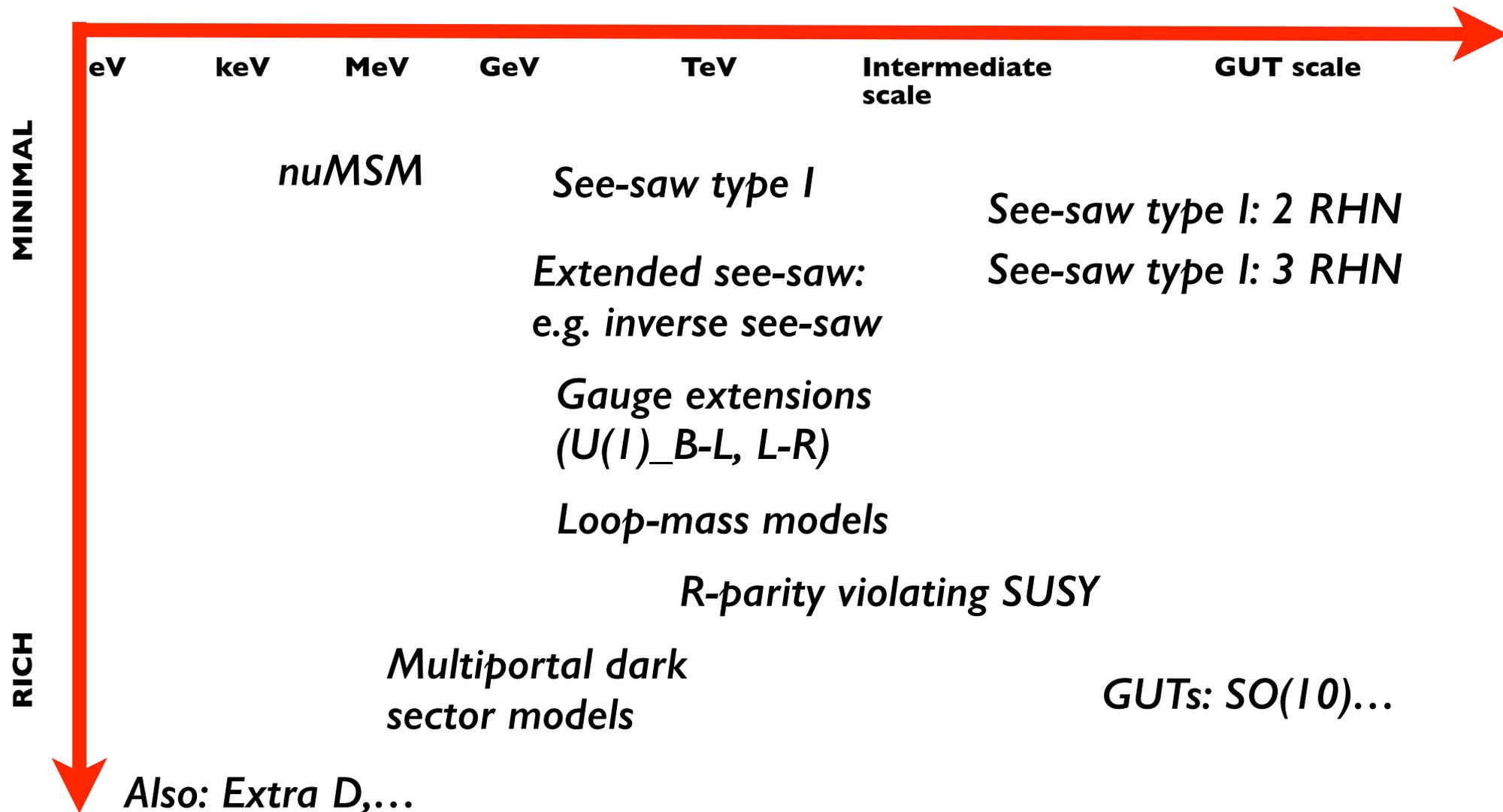


$L_\mu - L_\tau$ , DUNE ND, 75 tonnes, 5 y  $\nu$ -mode + 5 y  $\bar{\nu}$ -mode, 120 GeV  $p^+$ ,  $\sigma_{\text{norm}} = 5\%$



Ballett et al., 1902.08579, see also Altmannshofer et al., 1902.06765

# From minimality to richness

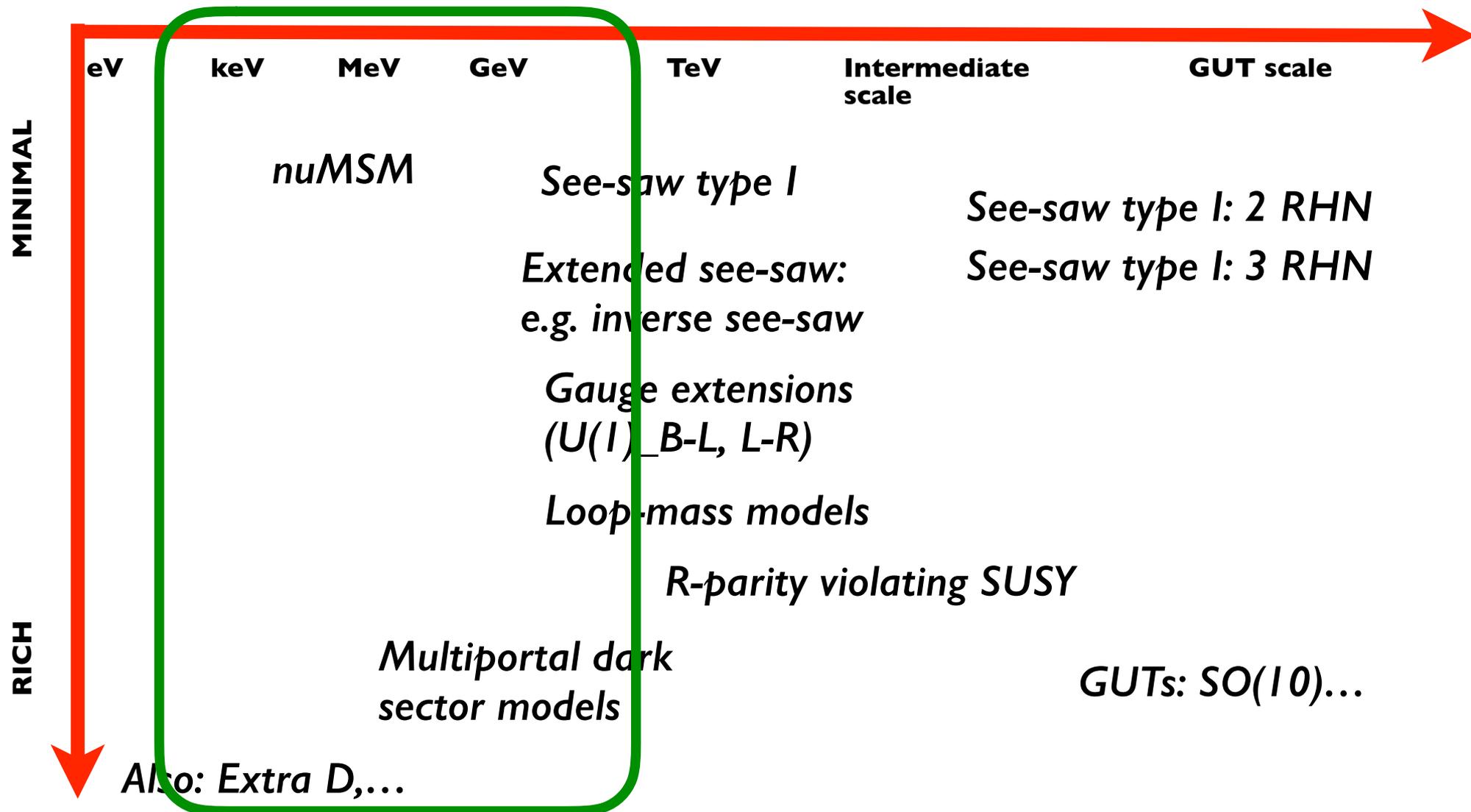


Two contrasting approaches can be taken:

Minimality: the fewest ingredients -> predictivity

Richness (theory-motivated): connections, new signatures

# From minimality to richness



Two contrasting approaches can be taken:

**Minimality: the fewest ingredients -> predictivity**

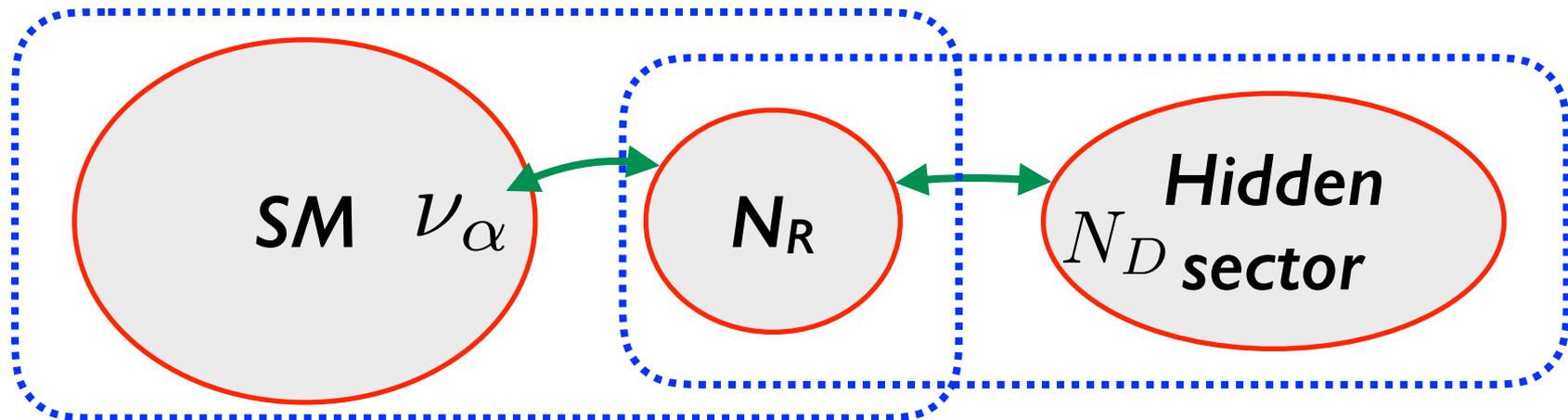
**Richness (theory-motivated): connections, new signatures**

HNLs and other DS particles could be part of a new low energy sector that contains several new states (neutral fermions, gauge bosons, scalars, DM)... This is the case for the Standard Model!

Example: **Three portal model**

***neutrino portal***

$$\bar{L} \cdot H N_R \quad (+ \dots \overline{N_R} N_S)$$

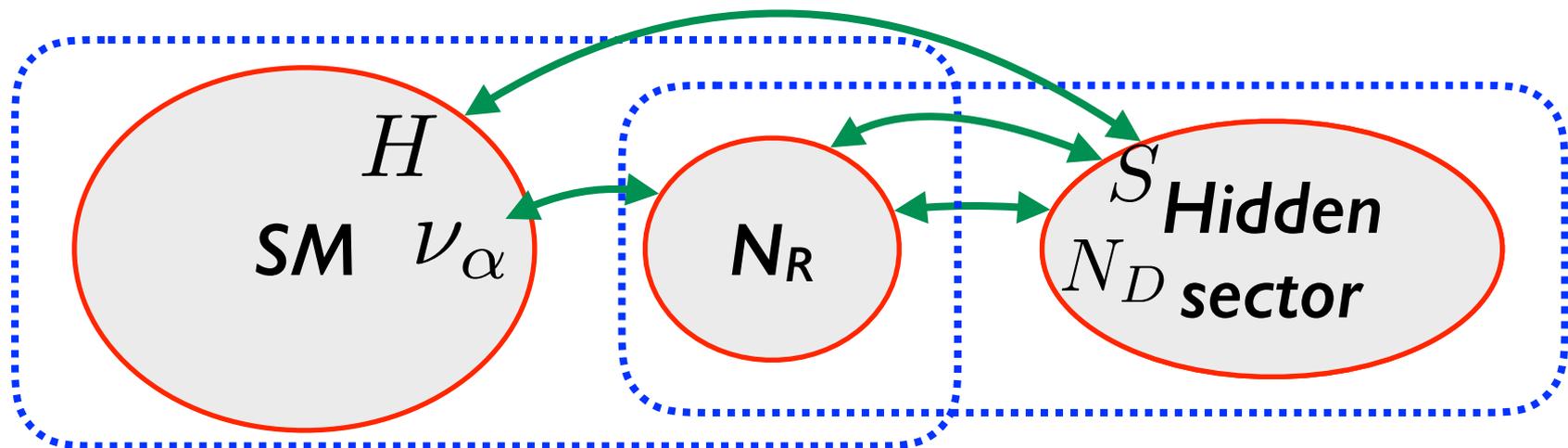


HNLs and other DS particles could be part of a new low energy sectors which contain several new states (neutral fermions, gauge bosons, scalars, DM)... This is the case for the Standard Model!

Example: **Three portal model**

**neutrino portal, Higgs portal**

$$\bar{L} \cdot H N_R \quad (+ \dots \bar{N}_R N_S) \quad \lambda_{\phi H} \phi^\dagger \phi H^\dagger H$$



HNLs and other DS particles could be part of a new low energy sectors which contain several new states (neutral fermions, gauge bosons, scalars, DM)... This is the case for the Standard Model!

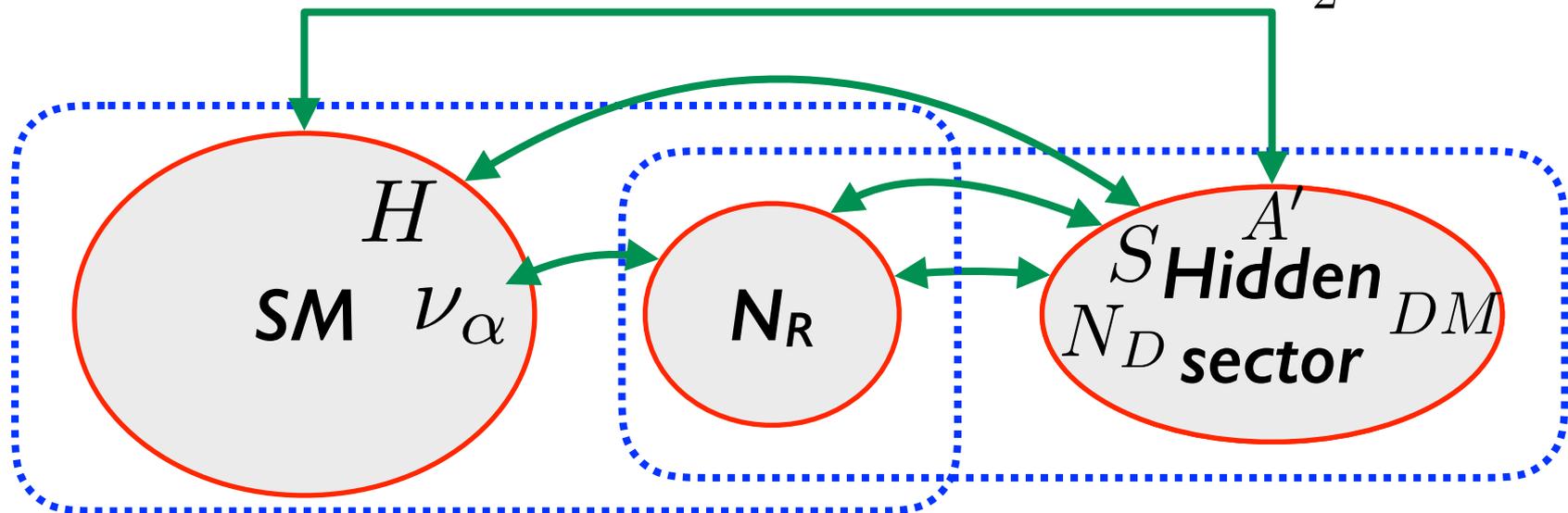
## Example: **Three portal model**

**neutrino portal, Higgs portal, vector portal**

$$\bar{L} \cdot H N_R \quad (+ \dots \bar{N}_R N_S)$$

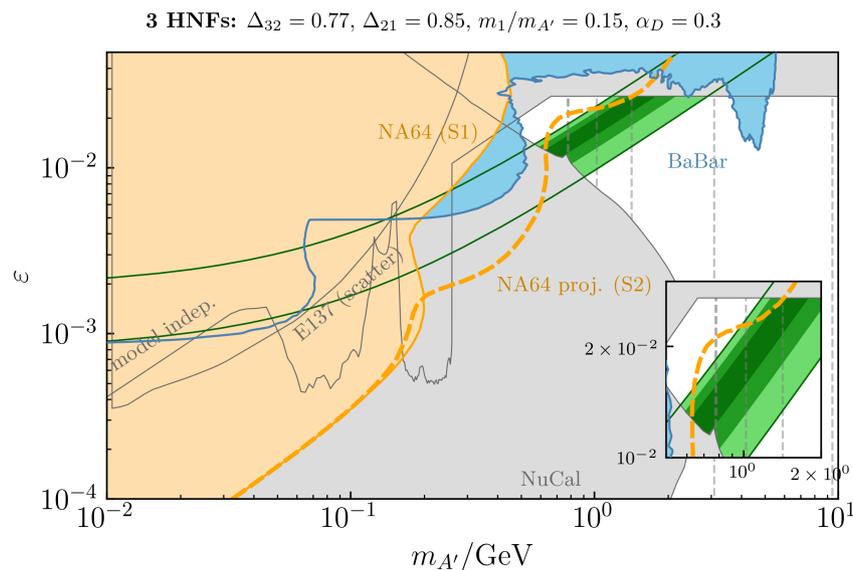
$$\lambda_{\phi H} \phi^\dagger \phi H^\dagger H$$

$$\frac{\sin \chi}{2} X_{\mu\nu} B^{\mu\nu}$$



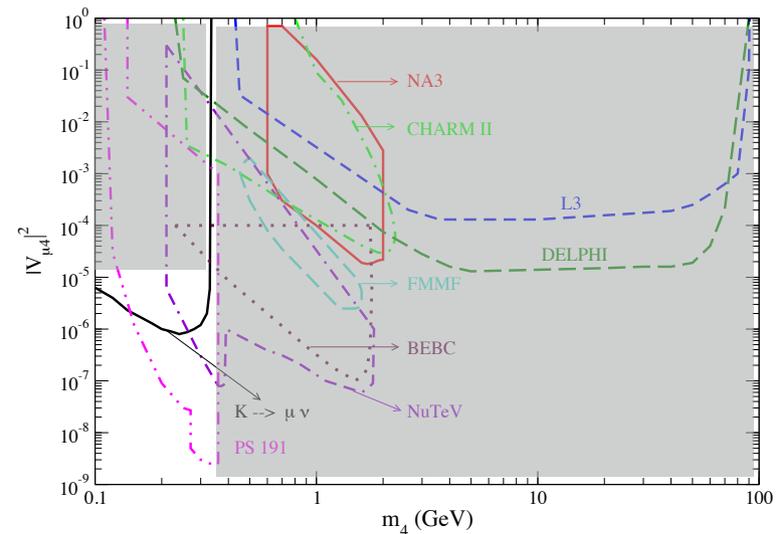
This type of models have a distinct phenomenology, very different from minimal models.

- They can evade standard bounds as the DS particles can decay fast and semivisibly.
- Typical signatures: semivisible decays into multileptons and missing energy with decay chains.



A. Abdullahi et al., 2302.05410

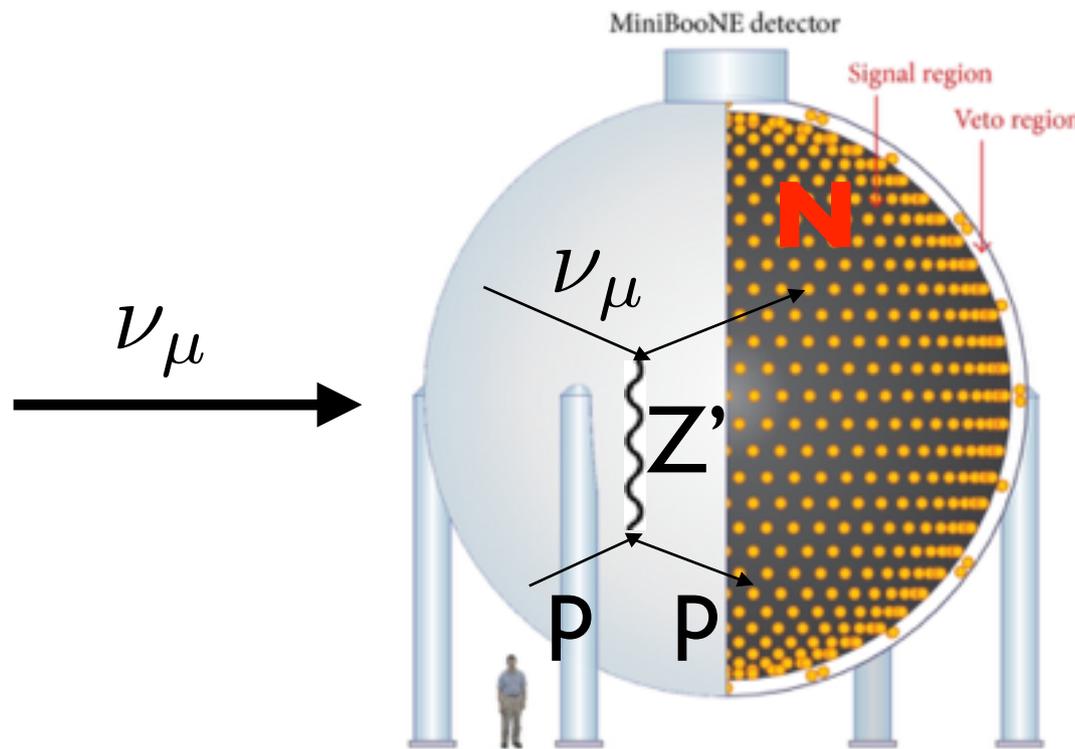
Collider bounds  
can be weakened.



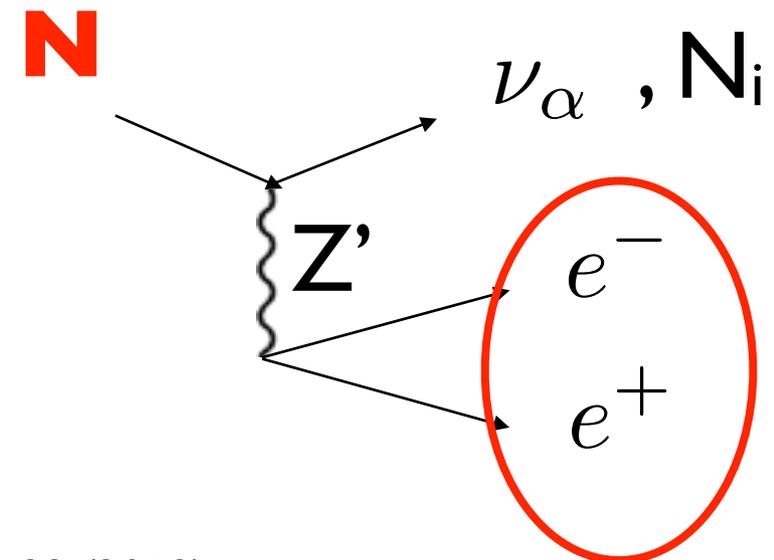
HNL decay bounds need  
to be recomputed.

# A MiniBooNE low-E excess explanation

A viable explanation of the MiniBooNE low-E excess is provided by the **up-scattering of an HNL  $N$**  in the detector and **its decay into  $ee \nu$** .



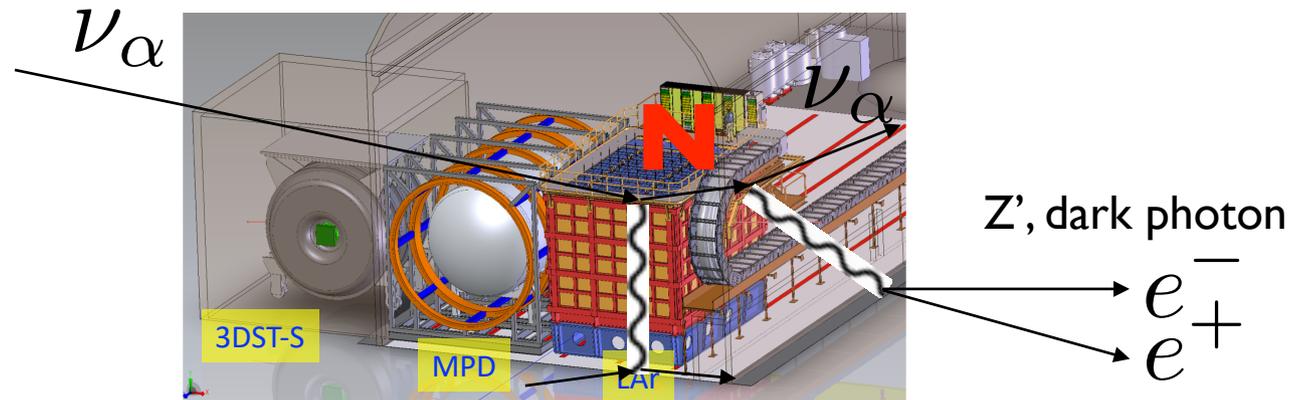
It builds on a decay explanation of MiniBooNE by S. Gnineko, PRL 103 (2009). A similar analysis appeared at the same time but with light  $Z'$  by E. Bertuzzo et al., PRL 121 (2018).



P. Ballett, S. Pascoli, M. Ross-Lonergan, PRD 99 (2019)

# Unique signatures and future tests

These models has key signatures which can be tested.



One can expect **displaced vertices**, decay chains and unique HNL and dark photon phenomenology (typically, semivisible decays):

- MicroBooNE, T2K ND, DUNE-ND;
- NA62&SHADOWS;
- Nu@LHC programme;
- NA64;
- BelleII and BESIII.

## Weak Mixing Angle

$$\sin^2 \theta_W = \frac{g'^2}{g^2 + g'^2}$$

The weak mixing angle is a key parameter in the SM. It can be measured in electron scattering.

Agarwalla, Huber, 1005.1254

## Sterile neutrinos

Sterile neutrinos with eV masses could be searched for at the DUNE ND with a different L/E dependence w.r.t. SBN.

See J. Penedo's talk

## Non-unitarity and 0-distance NSIs

Non-unitarity effects or NSIs could lead to specific appearance and disappearance signatures.

See M. Bishai's comments in talk

# Conclusions

**BSM models with light sectors** (eg. heavy neutral fermions, dark photons, light DM...) are very interesting extension of the SM. They can explain neutrino masses and can be connected to the other compelling observational evidences of BSM (DM and the baryon asymmetry).

**DUNE-ND has unique opportunities** to search for them thanks to high flux, very capable detector complex (and DUNE-PRISM):

- **a-la beam dump experiment;**
- **scattering searches;**
- **other searches (sterile neutrinos, NSI...).**

## Models with enhanced mixing

The mixing-mass relation can be avoided if neutrino masses are suppressed. A typical example are the **inverse see-saw and extended see-saw models**, in which two sterile neutrinos are introduced.

$$\mathcal{L} = Y\bar{L} \cdot H N_1 + Y_2\bar{L} \cdot H N_2^c + \Lambda\bar{N}_1 N_2 + \mu' N_1^T C N_1 + \mu N_2^T C N_2$$

The neutrino mass matrix is

$$\begin{pmatrix} 0 & Yv & Y_2v \\ Yv & \mu' & \Lambda \\ Y_2v & \Lambda & \mu \end{pmatrix}$$

Neutrino masses are suppressed by the (small) lepton number parameters:

$$m_{tree} \simeq -m_D^T M^{-1} m_D \simeq \frac{v^2}{2(\Lambda^2 - \mu'\mu)} (\mu Y_1^T Y_1 + \epsilon^2 \mu' Y_2^T Y_2 - \Lambda \epsilon (Y_2^T Y_1 + Y_1^T Y_2))$$

## Two limits:

- **Inverse see-saw:**  $\Lambda \gg \mu, Y_2 v, \mu'$  Gavela et al., 0906.1461; Ibarra, Molinaro, Petcov, 1103.6217

## Two quasi-Dirac neutrinos with large mixing:

$$m_4 \approx -m_5 \approx \tilde{M}_1 \approx -\tilde{M}_2 \approx \Lambda, \quad U_{e4} \approx U_{e5} \approx Y_{1e} v / 2\Lambda,$$
$$\Delta\tilde{M} \equiv |\tilde{M}_2| - |\tilde{M}_1| \approx \mu',$$

- **Extended see-saw:**  $\mu' \gg \Lambda, \mu$  Kang, Kim, 2007; Majee et al., 2008; Mitra, Senjanovic, Vissani, 1108.0004

## One light and one heavy sterile neutrino:

$$m_4 \approx \tilde{M}_1 \approx -\Lambda^2 / \mu', \quad U_{e4} \approx Y_{1e} v / \sqrt{2}\Lambda$$
$$m_5 \approx \tilde{M}_2 \approx \mu', \quad U_{e5} \approx Y_{1e} v / \sqrt{2}\mu'$$

Other models with enhanced sterile neutrino production require new interactions (e.g.  $Z'$ , see-saw type III...)