

# FROM HIGGS TO CHARM

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Standard Model @ the LHC

July 12, 2023



TH, X. Wang, arXiv:1704.00790; TH, B. Nachman, X. Wang, arXiv:1812.06992;  
B. Carlson, TH, S.C.I. Leung, arXiv:2105.08738; TH, Y. Ma, et al. arXiv:2202.08273

# MASSES IN THE STANDARD MODEL

The good, the bad, and the ugly

(1).  $M_{W,Z}$ : the good!

Prediction:  $M_W, M_Z = g v/2$  ;  $\delta m_w^2 \sim m_w^2 \ln(\Lambda/m_w)$

BSM: all calculable and predictable  $\rightarrow$  e.g. precision  $M_W$

(2).  $m_H$ : the bad!  $m_H = \sqrt{2} \mu = (2\lambda)^{1/2} v = 125 \text{ GeV}$

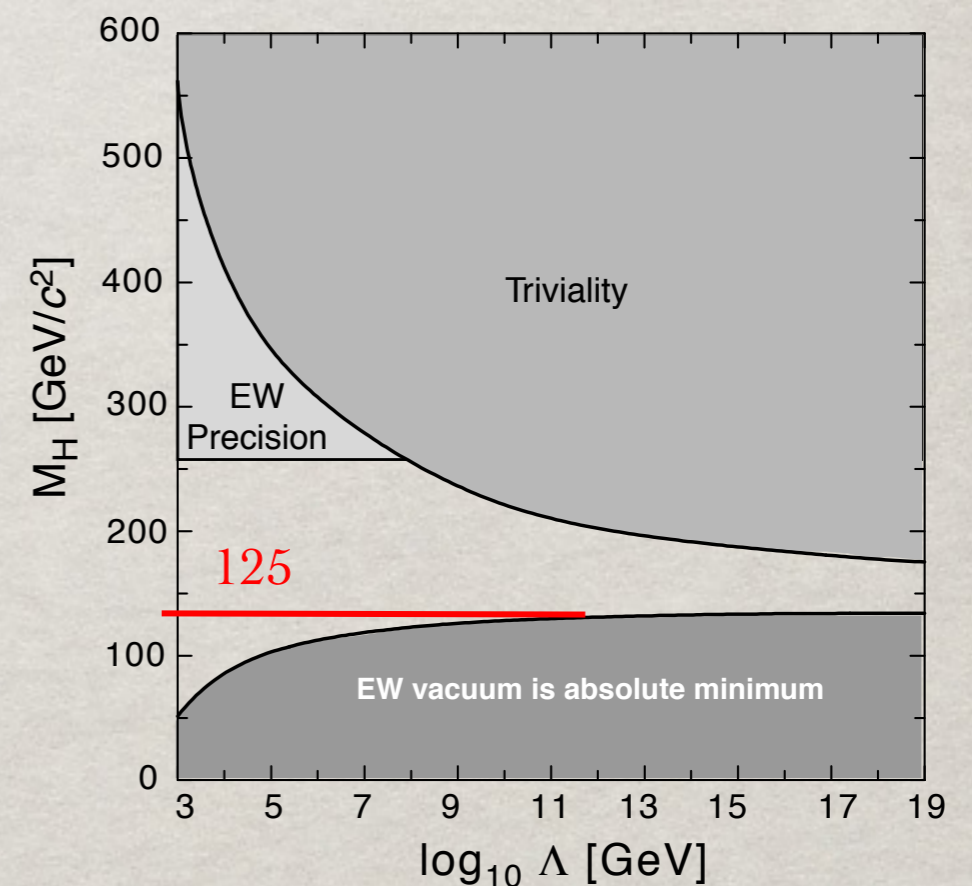
No prediction on  $m_H$

Quadratic corrections from the new physics scale:

$$\delta m_H^2 \propto -\frac{k^2}{4\pi^2} \Lambda^2$$

$\rightarrow$  “Naturalness” or “hierarchy puzzle” ?

Note: the quadratic mass corrections are NOT experimentally observable!



# (3). mf: the ugly!

☺ • Couplings are fixed by the masses:

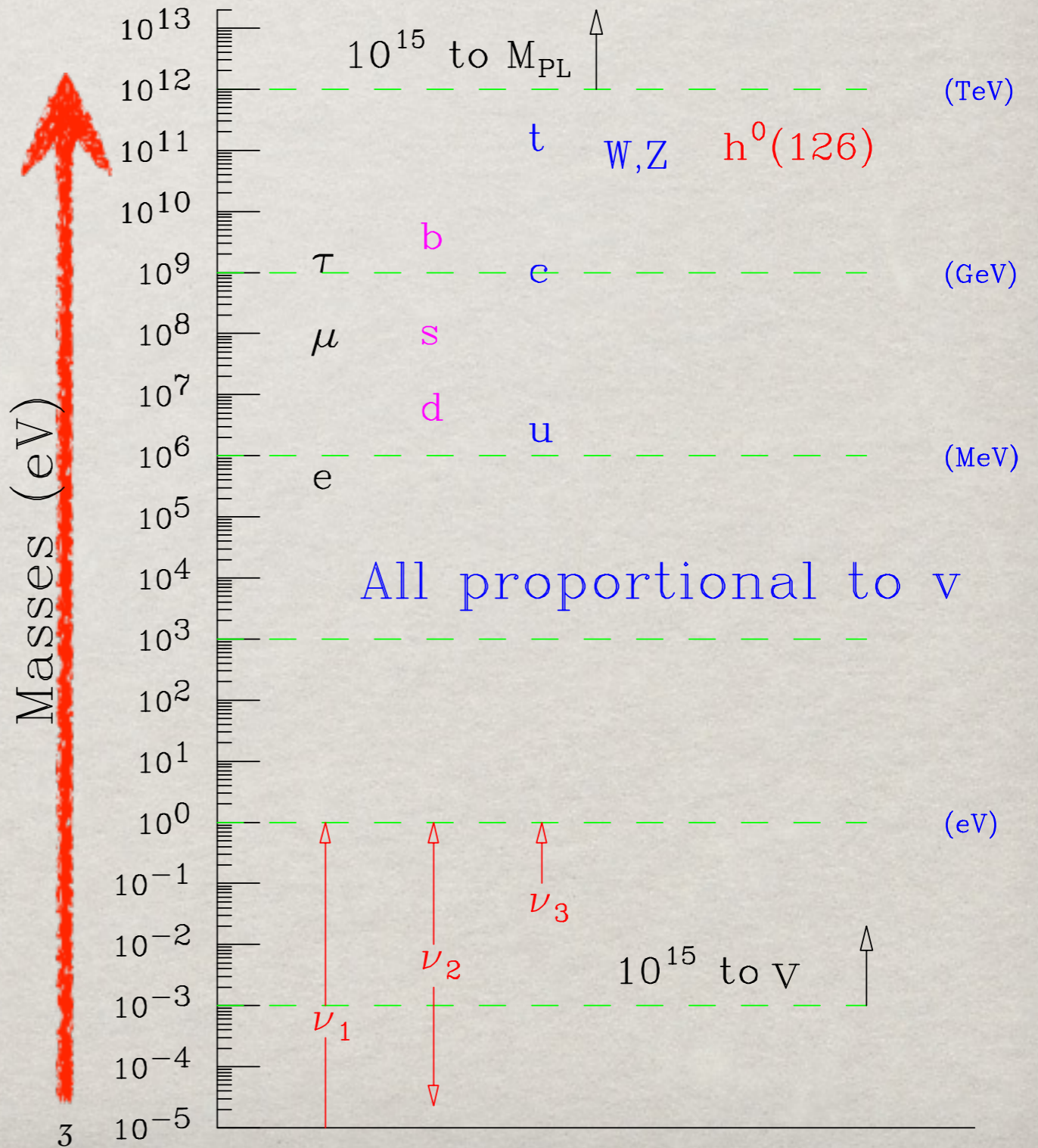
$$\mathcal{L}_Y \sim \sum_f m_f \bar{f} f (1 + H/v)$$

$$\delta m_f \sim m_f \ln(\Lambda/m_w) \quad (\text{chiral symm})$$



- Particle mass hierarchy
- Patterns of quark, neutrino mixings
- New CP-violation sources?
- Tiny neutrino masses!

**Higgs Yukawa couplings as the pivot for all !**



# YUKAWA COUPLINGS - THEORY:

## (1). Generate flavor hierarchies

- **Horizontal flavor symmetry:** Froggatt-Nielsen mechanism  
SM fermions charged  $[q_i, u_i, d_i]$  under  $U(1)_{FN}$  symmetry  
broken by  $\langle \phi \rangle / M \sim 0.2$

Froggatt & Nielsen (1979)

$$(Y_u)_{ij} \sim \left( \frac{\langle \phi \rangle}{M} \right)^{[q_i] - [u_j]}, \quad (Y_d)_{ij} \sim \left( \frac{\langle \phi \rangle}{M} \right)^{[q_i] - [d_j]}$$

J. Zupan & W. Altmannshofer,  
2203.07726

- **Warped extra-dimension:**  
Yukawa couplings determined by  
the overlapping with the Higgs brane.  
→ dual to (partial) composite model.

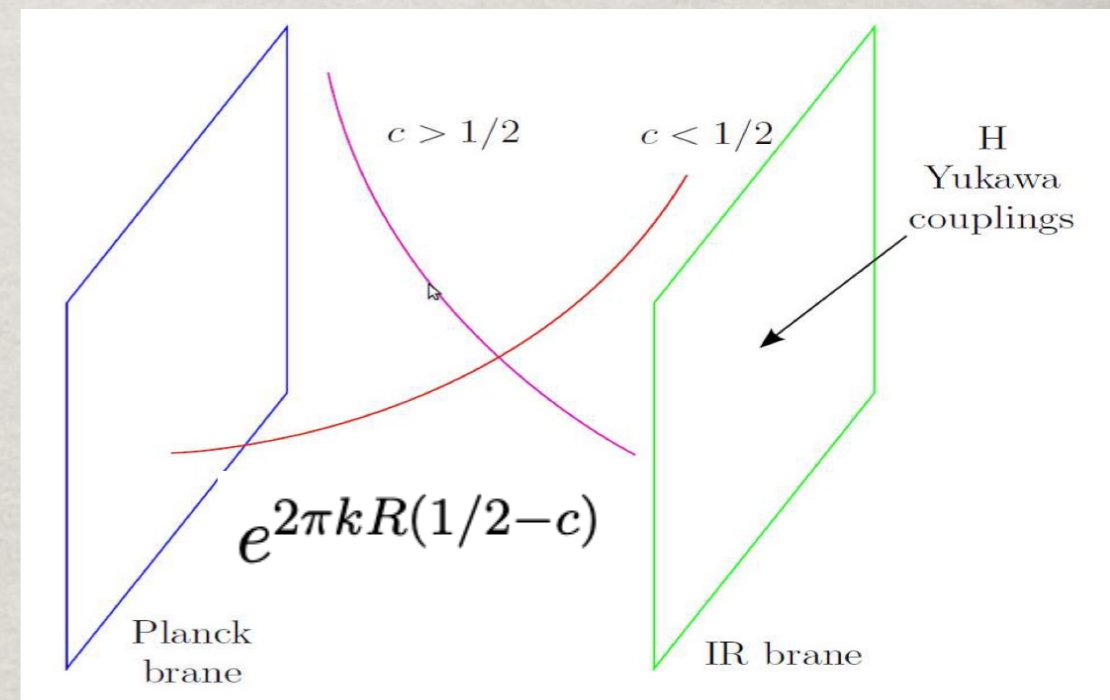
Randall & Sundrum (1999); Huber & Shafi (2001);  
Agashe et al. (2005)

- **Radiative generation** of  $m_f$ :

The 3<sup>rd</sup> generation @ tree-level

Light generations by new particle loops  $\sim 1/16\pi^2 \sim 10^{-2}$ .

S. Weinberg (1972)



## (2). The Higgs sector extension

- 2HDM (MSSM): well-motivated  
( $\tan\beta = v_2/v_1$ ;  $\alpha$  the neutral Higgs mixing)

**Talks in HPNP2023:**  
Pilaftsis; Haber; Takeuchi;  
Yagyu; Ferreira; De Curtis;  
Song; Dey; Kartayama;  
Heinemeyer; Muta; Ivanov;  
Sakurai ... ..

	Tree-level Normalized Higgs couplings			
	$\kappa_h^u$	$\kappa_h^d$	$\kappa_h^e$	$\kappa_h^V$
Type-I	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\sin(\beta - \alpha)$
Type-II	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\sin(\beta - \alpha)$
Type-L	$\frac{\cos\alpha}{\sin\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\sin(\beta - \alpha)$
Type-F	$\frac{\cos\alpha}{\sin\beta}$	$-\frac{\sin\alpha}{\cos\beta}$	$\frac{\cos\alpha}{\sin\beta}$	$\sin(\beta - \alpha)$

Decoupling/  
Alignment limit:

$$\kappa's \rightarrow 1$$

H. Haber & Y. Nir (1990)

- Plus a singlet  $S$  (NMSSM):  
more mixing & flavor physics, connect to dark sector

- Add a triplet  $\Phi$  (Type-II seesaw):  
 $\phi^{\pm\pm}, \phi^{\pm}, \phi^0$ ; connect to neutrino Majorana mass

For a review, see, i.e., G.C. Branco, M. Sher et al., arXiv:1106.0034 ...

Along with searching for BSM new physics,  
 precision Higgs measurement is a must!  
 → Parameterize Higgs couplings in a simple way

## SMEFT

SM Effective Field Theory: a linear representation

$$\text{SM-like Higgs } \Phi = \frac{1}{\sqrt{2}} \begin{pmatrix} \sqrt{2}\phi^+ \\ v + h + i\phi^0 \end{pmatrix}$$

$$\mathcal{L}_Y \sim \sum_{n=0} \frac{Y_{ij}^n}{\Lambda^{2n}} (\Phi^\dagger \Phi)^n \bar{L}_{iL} \Phi e_{jR} \rightarrow m_f = \frac{v}{\sqrt{2}} \sum_{n=0} Y_n^f \frac{v^{2n}}{\Lambda^{2n}}$$

Yukawa coupling deviates from the mass relation!

At the dim-6 leading order:

$$\rightarrow \delta\kappa_f \sim Y_1 \frac{v^2}{\Lambda^2} \sim O(\text{a few}\%) \text{ for } \Lambda \sim 2 \text{ TeV!}$$

This is the immediate target @ LHC!

# HEFT

Higgs Effective Field Theory: a non-linear representation

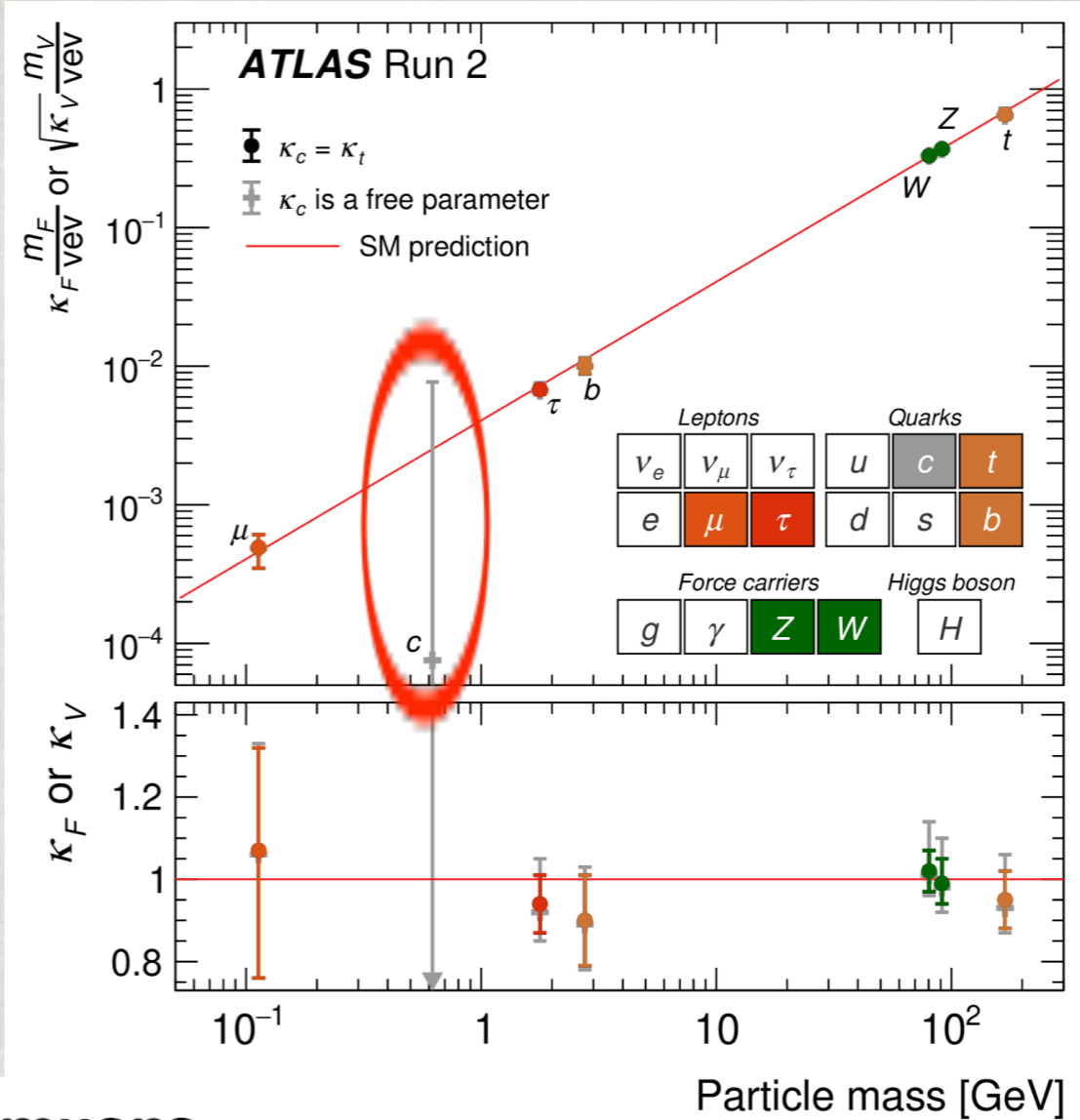
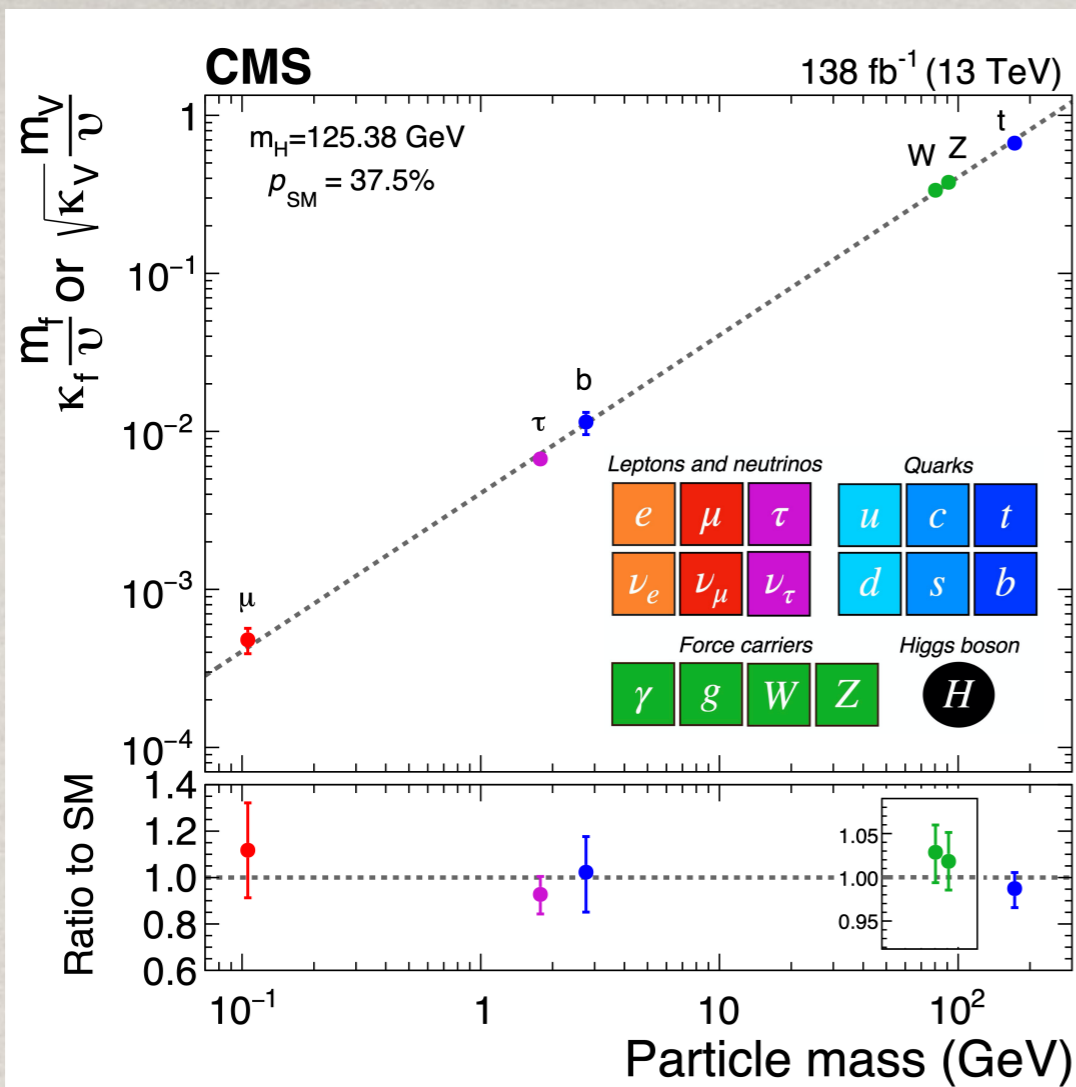
$$U = e^{i\phi^a \tau_a / v} \quad \text{with} \quad \phi^a \tau_a = \sqrt{2} \begin{pmatrix} \frac{\phi^0}{\sqrt{2}} & \phi^+ \\ \phi^- & -\frac{\phi^0}{\sqrt{2}} \end{pmatrix}$$
$$L_Y \sim -\frac{v}{2\sqrt{2}} \left[ \sum_{n \geq 0} y_n \left( \frac{H}{v} \right)^n (\bar{\nu}_L, \bar{\mu}_L) U (1 - \tau_3) \begin{pmatrix} \nu_R \\ \mu_R \end{pmatrix} + \text{h.c.} \right]$$

$$Y_f(H) = \frac{\sqrt{2}m_f}{v} + \sum_{n=1} y_{fn} \left( \frac{H}{v} \right)^n$$

- The scale for new dynamics is at  $\Lambda \sim 4\pi v$   
→ close by! The deviation can be sizable:

$$\rightarrow \delta\kappa_f \sim Y_1 \frac{H}{v} \sim O(1)$$

- Multiple  $W_L$ /Higgs may be enhanced.



## Evidence for Higgs boson decay to a pair of muons

ABSTRACT: Evidence for Higgs boson decay to a pair of muons is presented. This result combines searches in four exclusive categories targeting the production of the Higgs boson via gluon fusion, via vector boson fusion, in association with a vector boson, or in association with a top quark-antiquark pair. The analysis is performed using collision data at  $\sqrt{s} = 13 \text{ TeV}$ , corresponding to an integrated luminosity of  $37.1 \text{ fb}^{-1}$  recorded by the CMS experiment at the CERN LHC. An excess of events above the Standard Model expectation is observed in data with a significance of  $3.0 \sigma$ .

CMS coll., arXiv:2009.04363

## $h \rightarrow \mu^+ \mu^-$ via gluon fusion at the LHC

Tao Han and Bob McElrath

Department of Physics, University of Wisconsin-Madison, WI 53706

### Abstract

We study the observability of the  $h \rightarrow \mu^+ \mu^-$  decay in the Standard Model and the MSSM at the LHC. The observation of the  $h\mu\mu$  coupling is important to determine whether the Higgs particle that generates mass for the weak bosons is also responsible for mass generation of the second generation of fermions. We find that the signal via the gluon fusion channel is comparable to that from the weak-boson fusion. By combining these two channels, observing  $h \rightarrow \mu^+ \mu^-$  is feasible at the LHC with a delivered luminosity of  $300 \text{ fb}^{-1}$  at  $3\sigma$  statistical significance for  $110 \text{ GeV} < m_h < 140 \text{ GeV}$  in the Standard Model. This corresponds to a  $h\mu\mu$  coupling determination at about 15% accuracy assuming  $ht\bar{t}$ ,  $hb\bar{b}$  couplings SM-like. The observation becomes more promising in the MSSM for  $\tan\beta > 8$  and  $M_A < 130$ .



# CHARM-QUARK YUKAWA COUPLING

$$y_f(h + \frac{v}{\sqrt{2}})\bar{f}f \rightarrow y_c^{SM} = \frac{\sqrt{2}m_c}{v} \sim 4 \times 10^{-3}$$

In the  $\kappa$ -scheme for BSM:  $y_c = \kappa_c y_c^{SM}$

Measuring  $H\bar{c}c$  coupling at the LHC is very difficult:

- BR( $h \rightarrow cc$ )=2.9%
- Formidable QCD backgrounds
- c-tagging challenging
- ... ..

**(1).  $H \rightarrow cc$  at high  $p_T$**  The leading channel:

$$pp \rightarrow Vh, \quad \text{with } V \rightarrow \ell^+\ell^-, \ell^\pm\nu, h \rightarrow \bar{c}c$$

subject to huge backgrounds 2-jets, cc;  $h \rightarrow bb \dots$

At LHC Run 2:

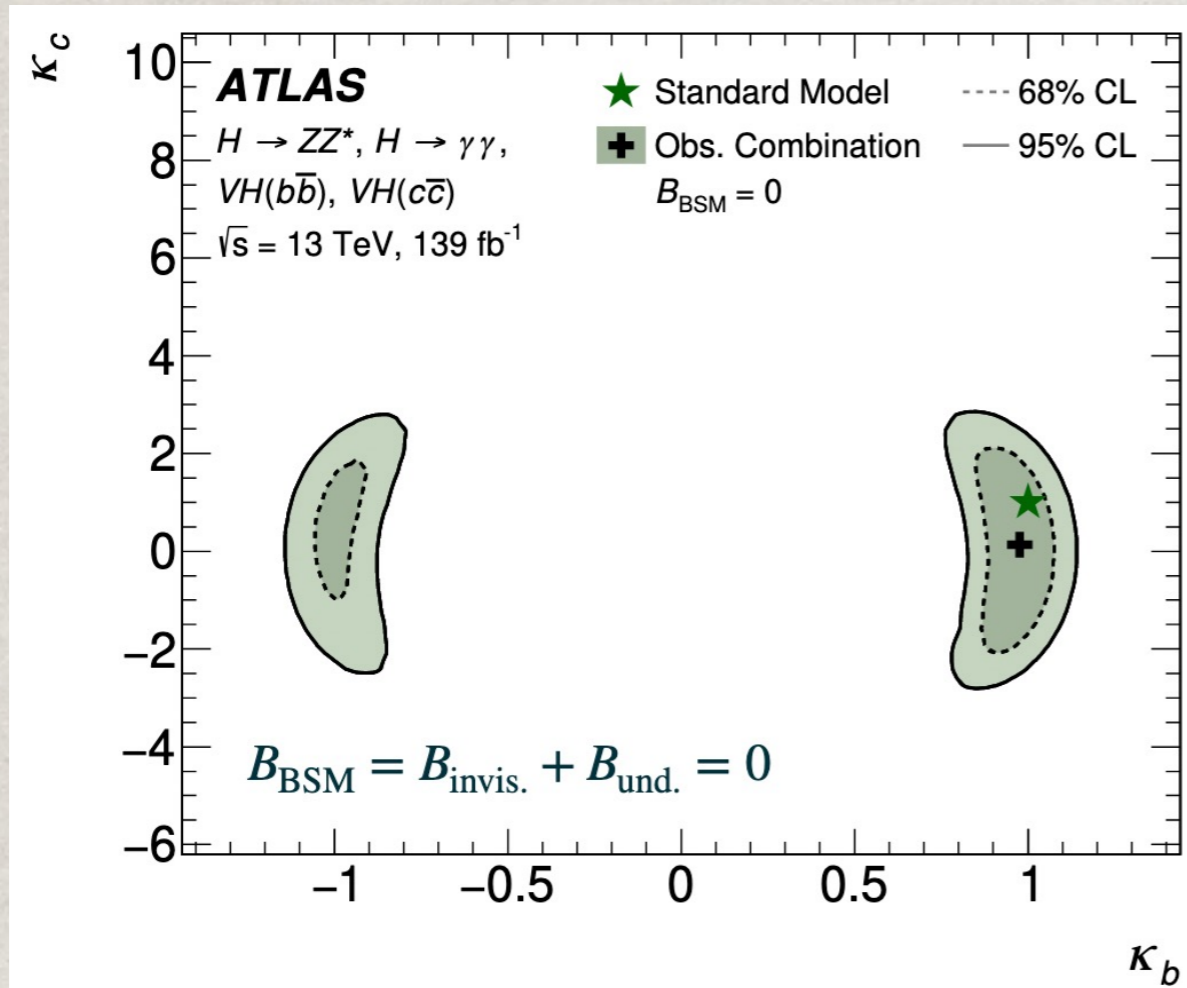
$$\kappa_c < 8.5@2\sigma \quad \text{ATLAS-CONF-2021-021}$$

Inclusive high  $p_T(h \rightarrow cc)$ :

$$\kappa_c < 6.9@2\sigma \quad \text{CMS: arXiv:2211.14181}$$

Nick Smith, next talk.

# LHC Run2 fit:



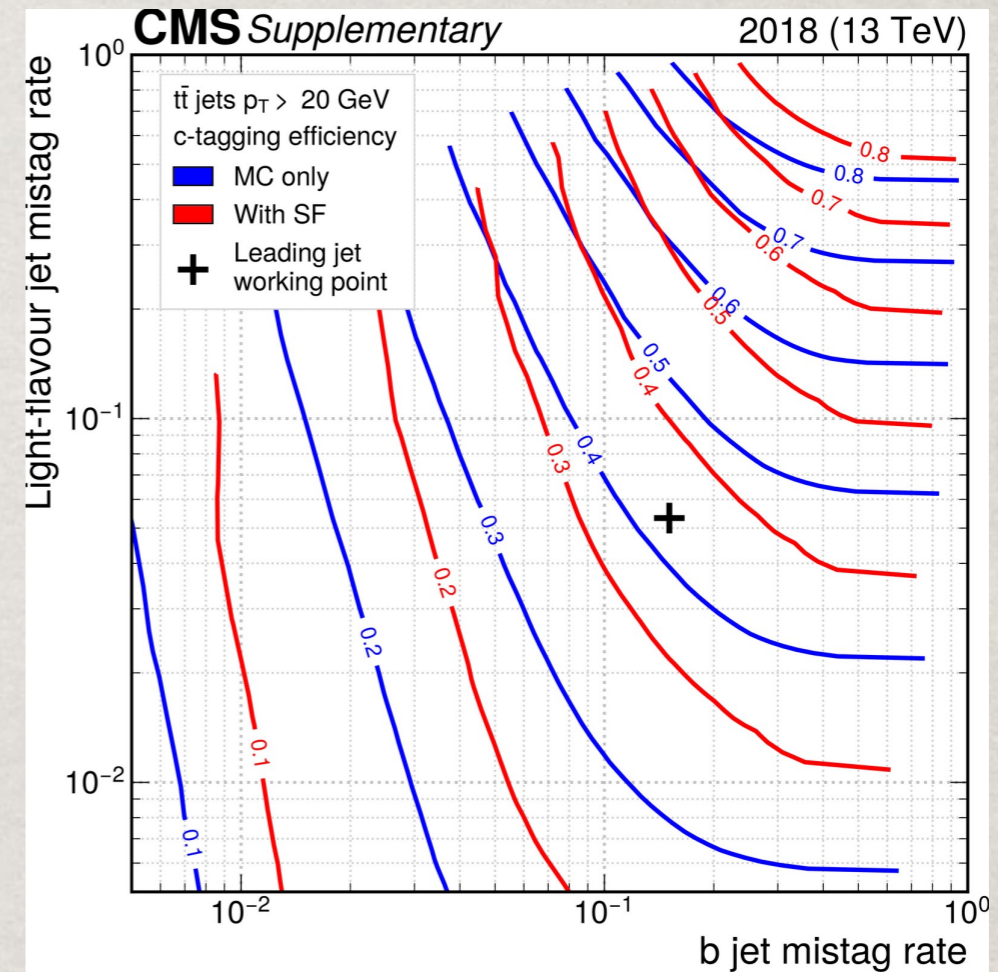
Submitted to JHEP [2207.08615]

At HL-LHC with  $3 \text{ ab}^{-1}$ :

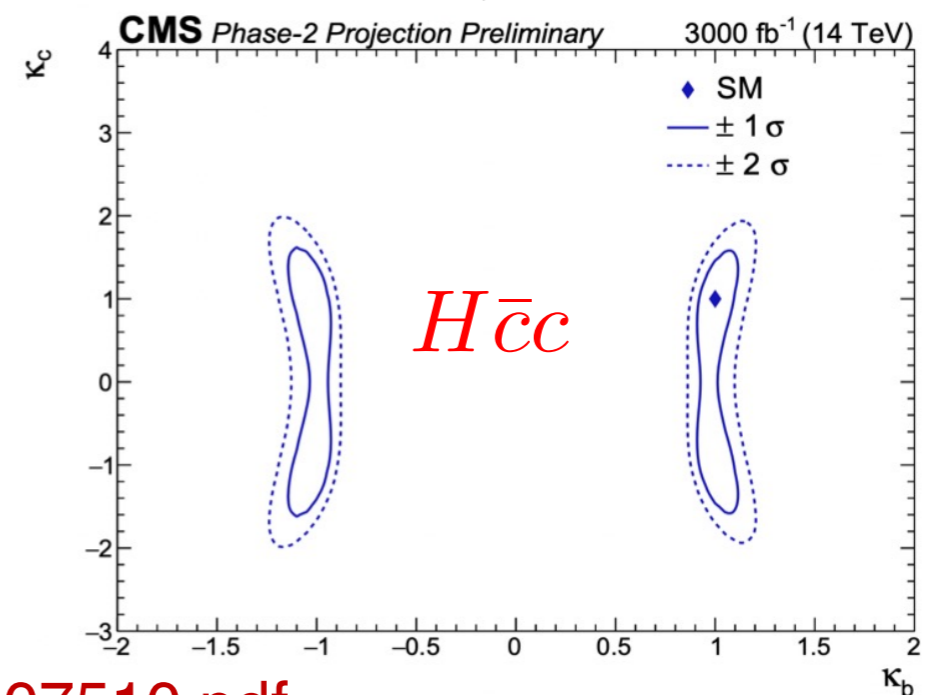
arXiv:1902.10229; ATLAS-CONF-2021-021

EF01/02 report: <https://arxiv.org/pdf/2209.07510.pdf>

Nick Smith, next talk.



CMS-HIG-21-008 ; CERN-EP-2022-081



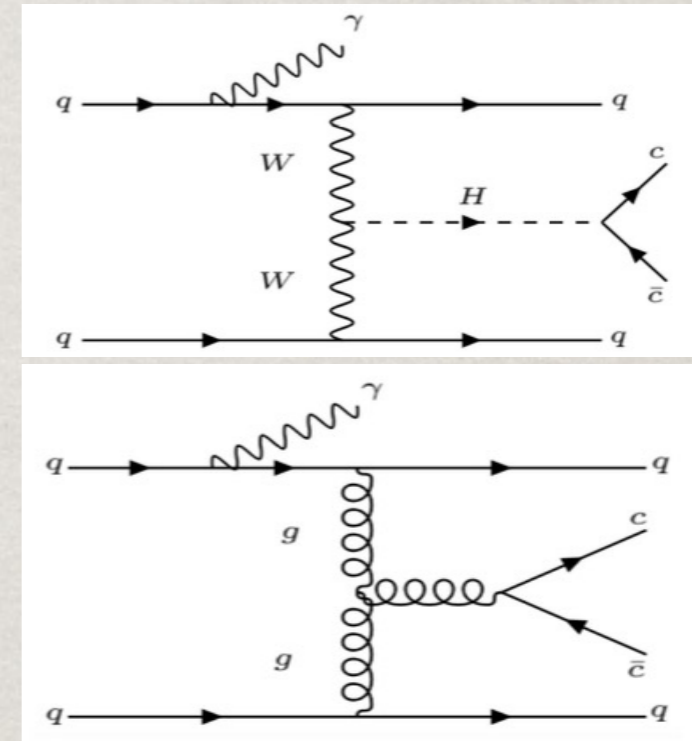
First extrapolation of  $\kappa_c$

Higgs production rate is high:  $\#H@LHC \sim 50 \text{ M}/\text{ab} !$   
 Need new ideas!

## (2). Higgs to charm + $\gamma$ in VBF

### A New Approach: VBF + $\gamma$

- Striking signatures and sizable signal events
- Additional photon results in lower rate
- Compensated by
  - Extra handle to trigger on
  - Suppression of gluon-rich background



$$p_T^j > 35 \text{ GeV}, \quad |\eta_j| < 5, \quad p_T^\gamma > 25 \text{ GeV}, \quad |\eta_\gamma| < 3.$$

	13 TeV	14 TeV	30 TeV	100 TeV
$\sigma_{\text{VBF}+\gamma}$ (pb)	0.024	0.027	0.099	0.43
$\sigma_{pp \rightarrow 4j+\gamma}$ (pb)	830	940	3700	21000

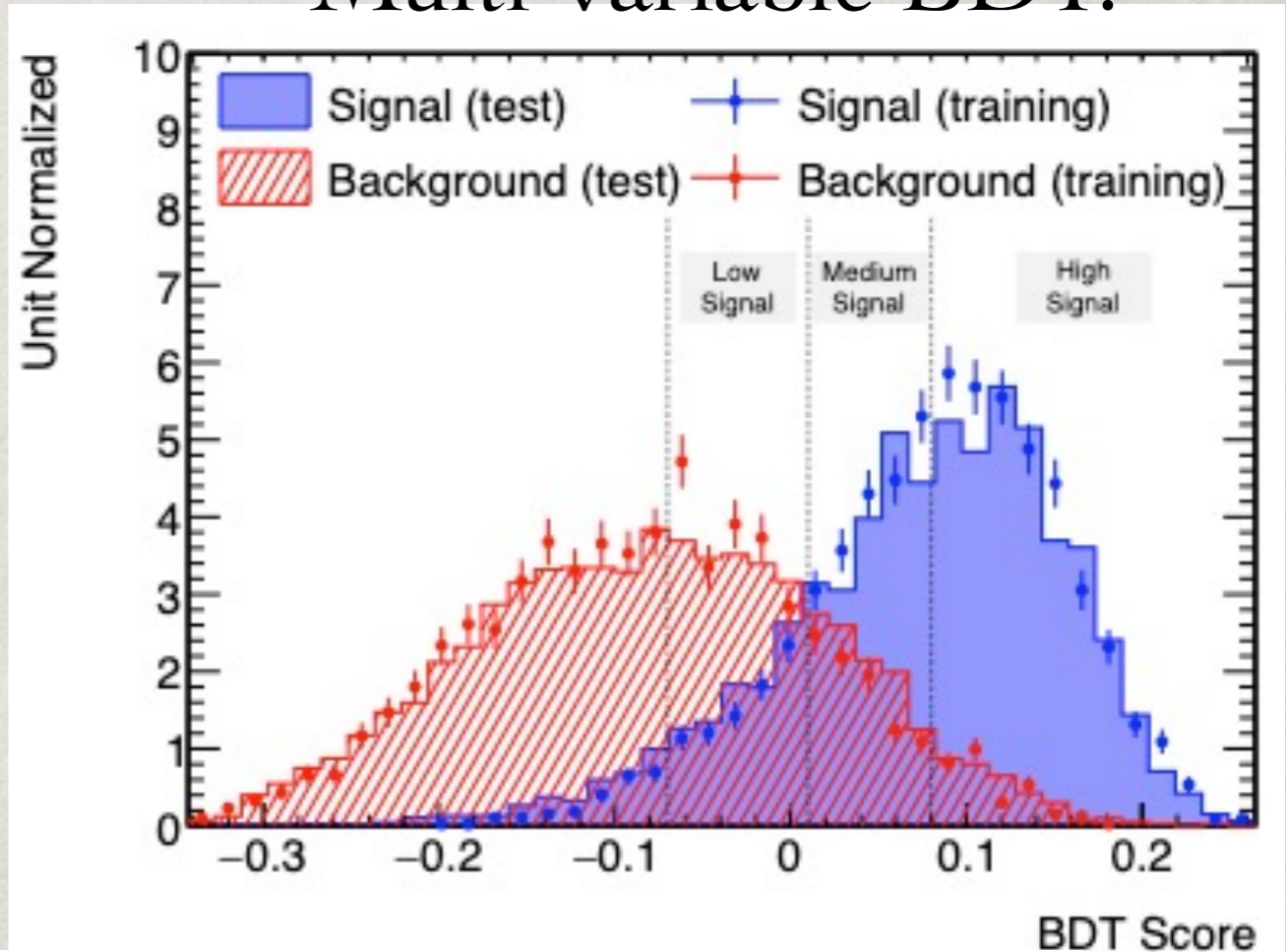
**Table 2:** Cross sections of signal and background at different center-of-mass energies, with the basic acceptance cuts in Sec. 2.1.

B. Carlson, TH, S.C.I. Leung, aXive:2105.08738

## (2). Higgs to charm+ $\gamma$ in VBF

Multi-variable BDT:

- Photon  $E_T^\gamma > 30$  GeV;
- At least four jets with  $p_T^j > 40$  GeV;
- At least one pair of jets with  $m_{jj} > 700$  GeV;
- At least one  $b$ -tagged jet with 77% efficiency.



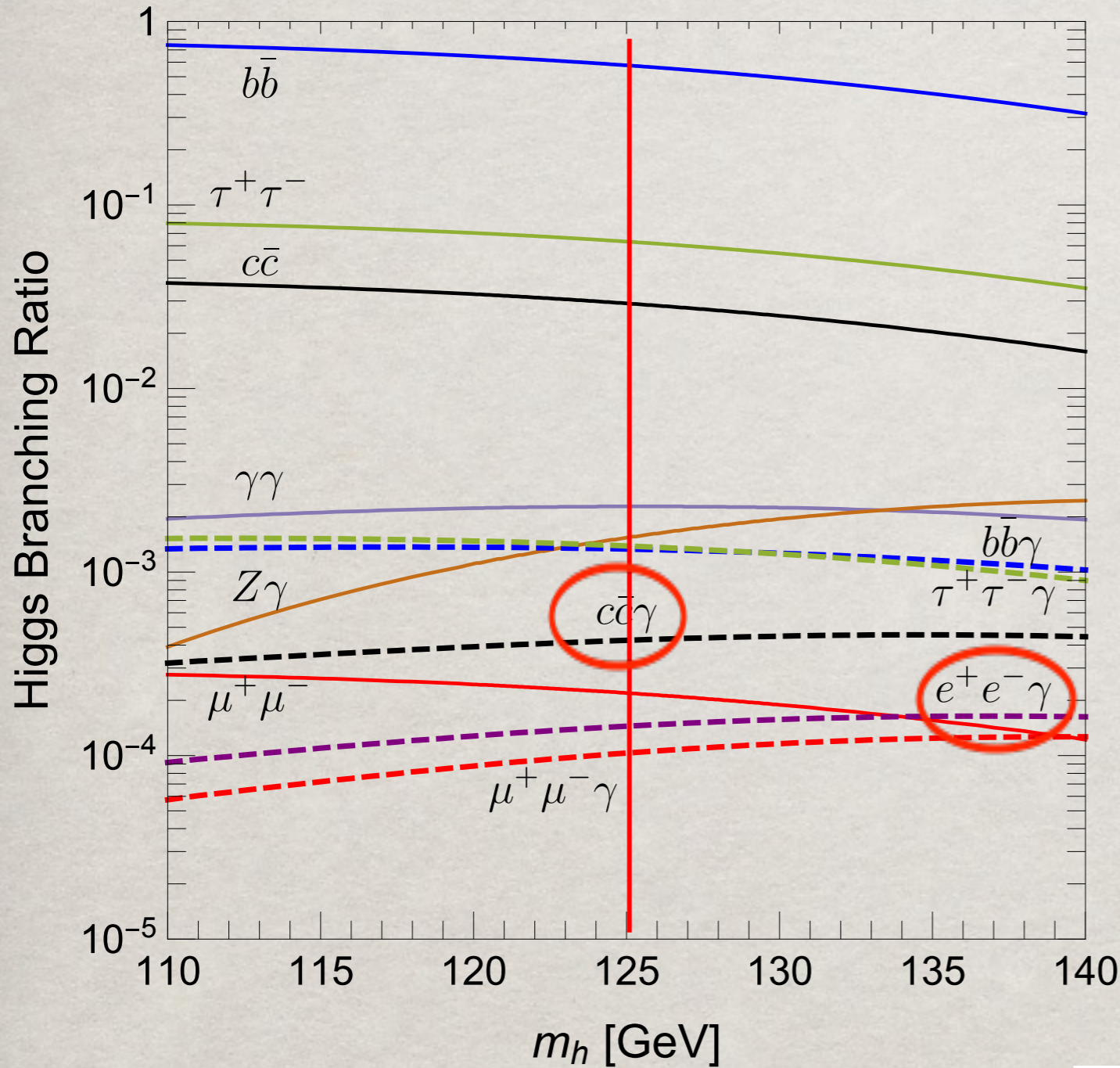
	LHC	Cut-based	BDT	$ZH$ [16, 17]	Fit [33]	$Hc$ [31]	$H \rightarrow c\bar{c}\gamma$ [41]
$\kappa_c$	36.1 fb <sup>-1</sup>	20	16	10	-	-	-
	3 ab <sup>-1</sup>	6.5	5.4	2.5	1.2	2.6 - 3.9	8.6

**Table 5:** The expected 95% CL<sub>s</sub> upper limit on the charm-Yukawa coupling from this analysis using 36.1 fb<sup>-1</sup> and 3 ab<sup>-1</sup> of data at 13 TeV, respectively, in comparison with other searches as quoted.

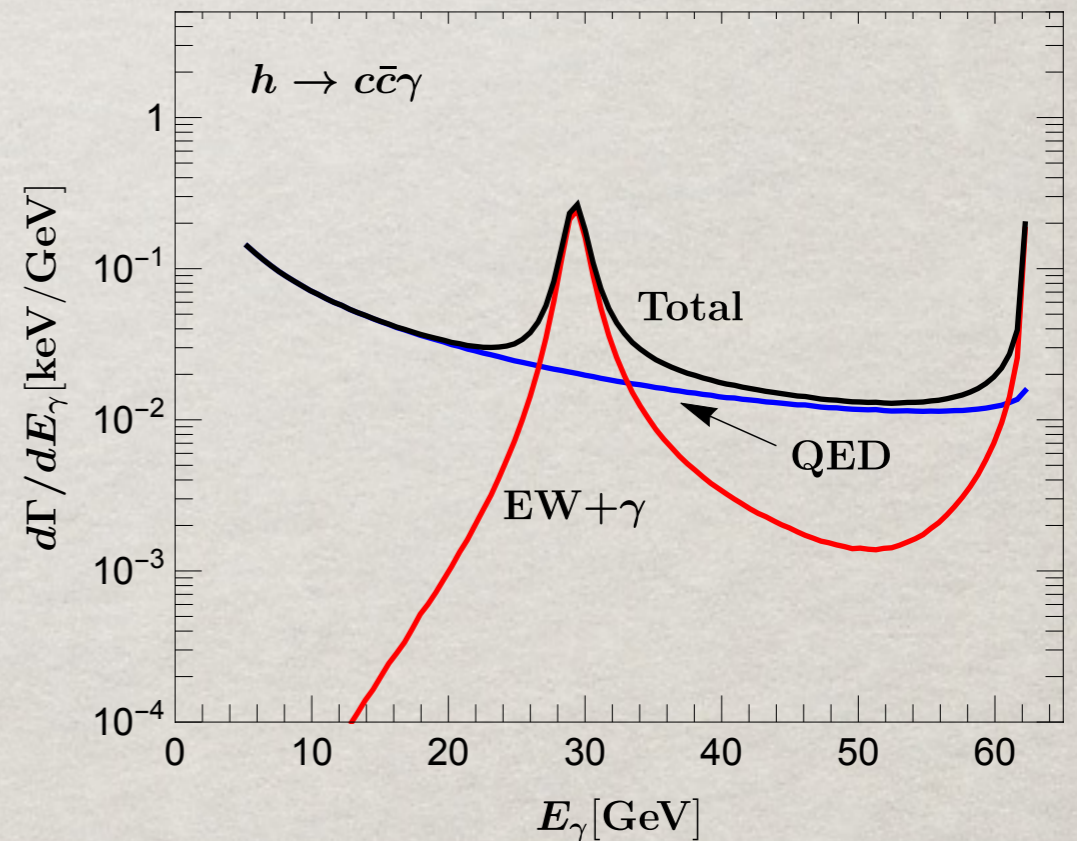
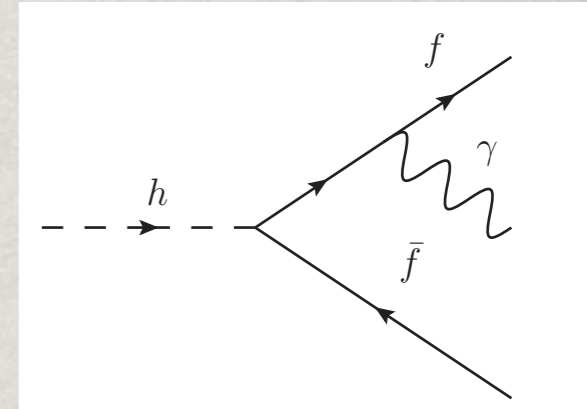
B. Carlson, TH, S.C.I. Leung, aXive:2105.08738;  
 [33]: J. de Blas et al., JHEP 01 (2020) 139, [1905.03764]

# (3). Higgs radiative decay

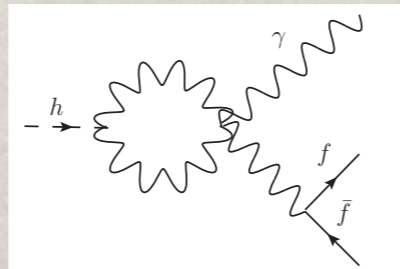
TH and X. Wang, arXiv:1704.00790



→ additional  $\gamma$  to trigger on:



Not  $\text{EW} + \gamma$ ,  
insensitive to  $\kappa_c$



$$E_\gamma = \frac{m_h}{2} \left(1 - \frac{m_Z^2}{m_h^2}\right) \approx 30 \text{ GeV, for } \gamma Z \text{ production,}$$

$$E_\gamma = \frac{m_h}{2} \left(1 - \frac{m_{\gamma^*}^2}{m_h^2}\right) \approx 63 \text{ GeV, for } \gamma\gamma^* \text{ production.}$$

# Results for $gg \rightarrow h \rightarrow c \bar{c} \gamma$ :

Operating Point	$\epsilon_c$	$\epsilon_b$	$\epsilon_j$
I	20%	33%	0.13%
II	30%	33%	1%
III	41%	50%	3.3%

c-tagging efficiency;  
b and jet mis-tagging rates

	Working Point	Signal (QED)	Background events	Background event rate [Hz]	$S/\sqrt{S+B}$ [ $10^{-2}$ ]
Level-1 (L1)	No Tag	-	-	$9.55 \times 10^3$	-
1 <i>c</i> -tag	I	269	$3.37 \times 10^8$	5.62	1.47
	II	349	$5.18 \times 10^8$	8.63	1.54
	III	401	$8.83 \times 10^8$	14.7	1.35
2 <i>c</i> -tags	I	29	$1.14 \times 10^7$	0.191	0.878
	II	66	$2.23 \times 10^7$	0.371	1.42
	III	126	$5.79 \times 10^7$	0.966	1.66

**Table 2:** Expected numbers of events of the signal and background, and event rates, in the range of  $100 < M_{jj\gamma} < 140$  GeV at the HL-LHC with  $\mathcal{L} = 3 \text{ ab}^{-1}$ . The first row gives the event rate at L1, with only the requirements in Sec. 2 applied. Systematic uncertainties are not accounted for in the significance calculation in the last column.

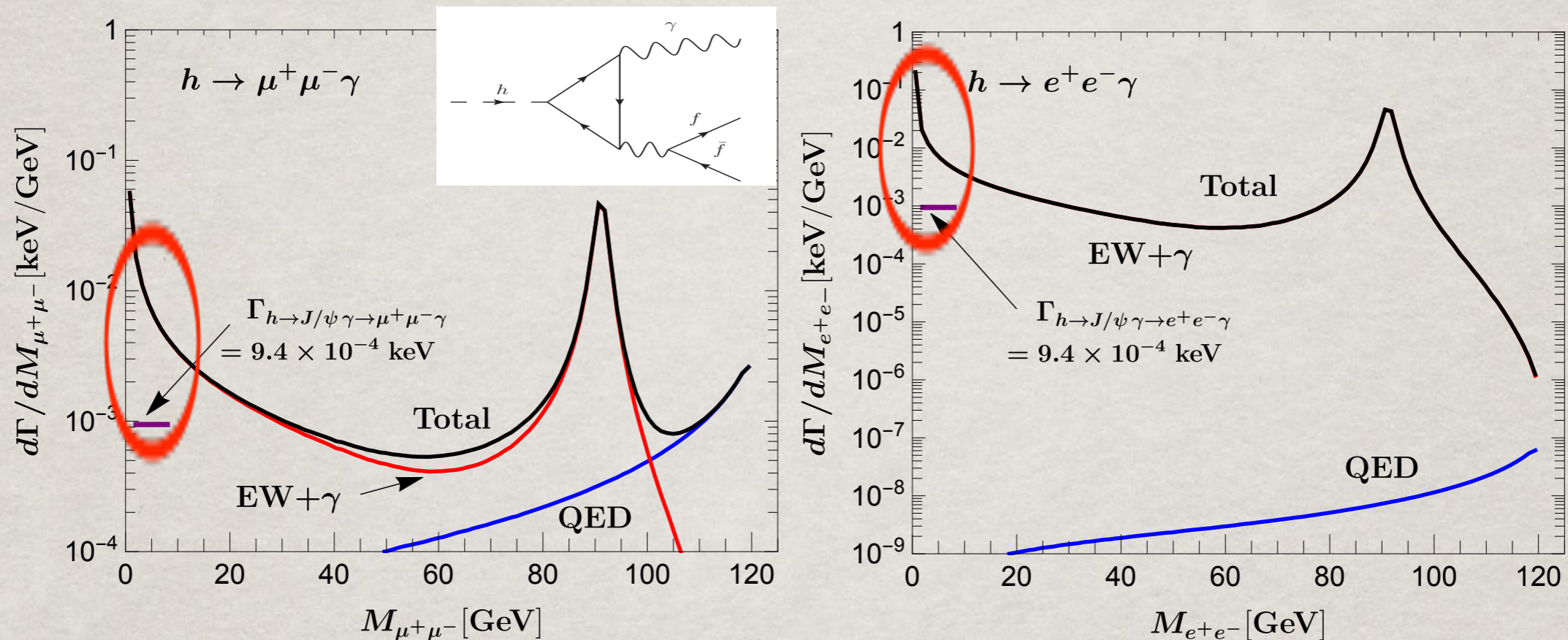
95%  $\text{CL}_s$  upper limit on the signal strength via BDT-based analysis to be

$$\mu < 89, 71, 70, \Rightarrow \kappa_c < 9.4, 8.4, 8.3.$$

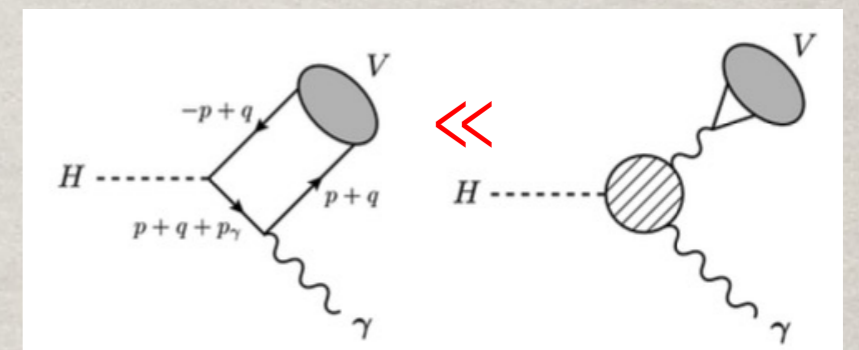
TH, B. Nachman, X. Wang, arXiv:1812.06992

# A remark on $gg \rightarrow h \rightarrow J/\psi \gamma \rightarrow \ell^+ \ell^- \gamma$ :

Bodwin, Petriello et al. (2013, 2014, 2017); Konig, Neubert (2015)



**$J/\psi \gamma$**  is more than an order of magnitude smaller;  
 Further, the dominant contribution is from the  
 “vector-meson dominance”  $\gamma^* \rightarrow J/\psi$ ,  
*insensitive* to the direct  $hcc$  Yukawa:



TH and X. Wang, arXiv:1704.00790

# (4). Higgs decay to charmonia: c-fragmentation

- $H \rightarrow J/\psi$  via charm-quark fragmentation:

$$H \rightarrow c + \bar{c} + J/\psi \text{ (or } \eta_c)$$

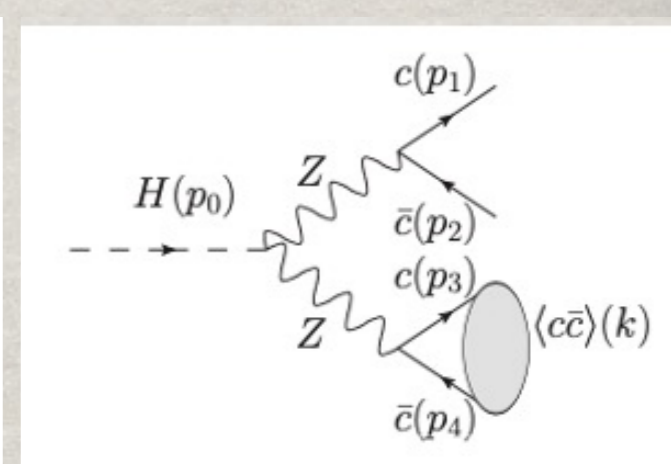
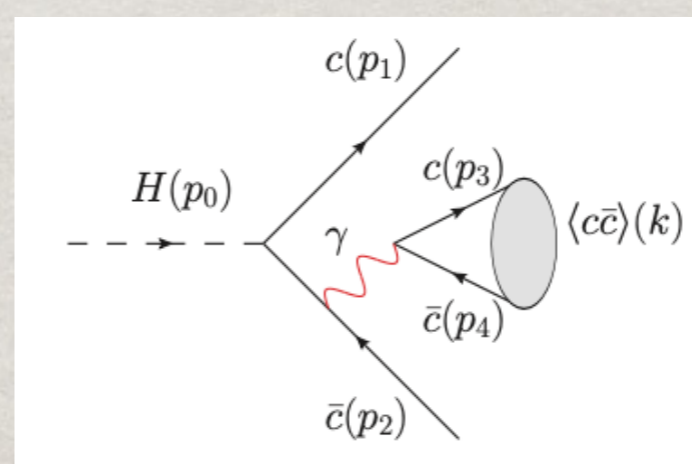
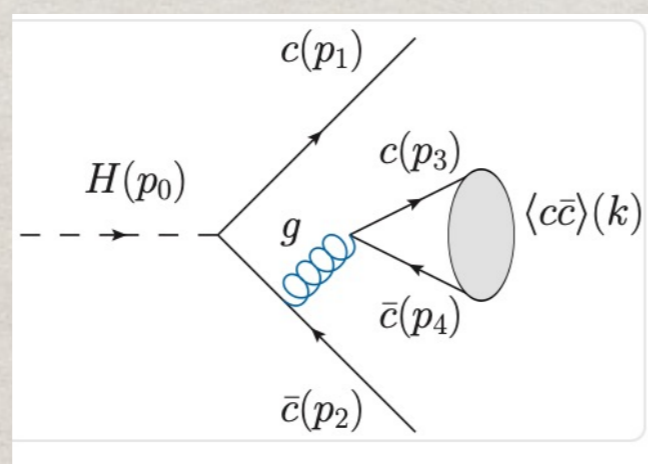
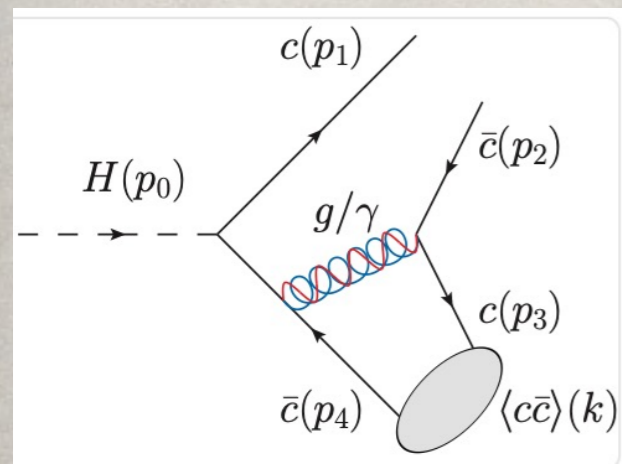
- Enhanced from the fragmentation
- Direct coupling to charm!

Color-singlet (CS)  
(leading)

Color-octet (CO)  
(sub-leading,  $1/2$  of CS)

QED  
(sub-leading,  $1/4$  of CS)

EW  
(sub-sub-leading)



TH, A. Leibovich, Y. Ma, X.Z. Tan: aXive:2202.08273



## Nonrelativistic QCD framework

$$\Gamma = \sum_{\mathbb{N}} \hat{\Gamma}_{\mathbb{N}}(H \rightarrow (Q\bar{Q})[\mathbb{N}] + X) \times \langle \mathcal{O}^h[\mathbb{N}] \rangle, \quad d\hat{\Gamma}_{\mathbb{N}} = \frac{1}{2m_H} \frac{|\mathcal{M}|^2}{\langle \mathcal{O}^{Q\bar{Q}} \rangle} d\Phi_3$$

## Long distance matrix element (LDME)

Related to the wave function at origin

$$\begin{aligned} \langle \mathcal{O}^{J/\psi}[{}^3S_1^{[1]}] \rangle &= \frac{3N_c}{2\pi} |R(0)|^2, & \langle \mathcal{O}^{\eta_c}[{}^1S_0^{[1]}] \rangle &= \frac{N_c}{2\pi} |R(0)|^2 \\ \langle \mathcal{O}^{Q\bar{Q}} \rangle &= 6N_c, \text{ for } {}^3S_1^{[1]}, & \langle \mathcal{O}^{Q\bar{Q}} \rangle &= 2N_c, \text{ for } {}^1S_0^{[1]} \end{aligned}$$

A quarkonium can also be produced through color-octet  $Q\bar{Q}$  Fock states

New states involved:  ${}^3S_1^{[8]}$ ,  ${}^1S_0^{[8]}$ ,  ${}^3P_J^{[8]}$ , and  ${}^1P_1^{[8]}$

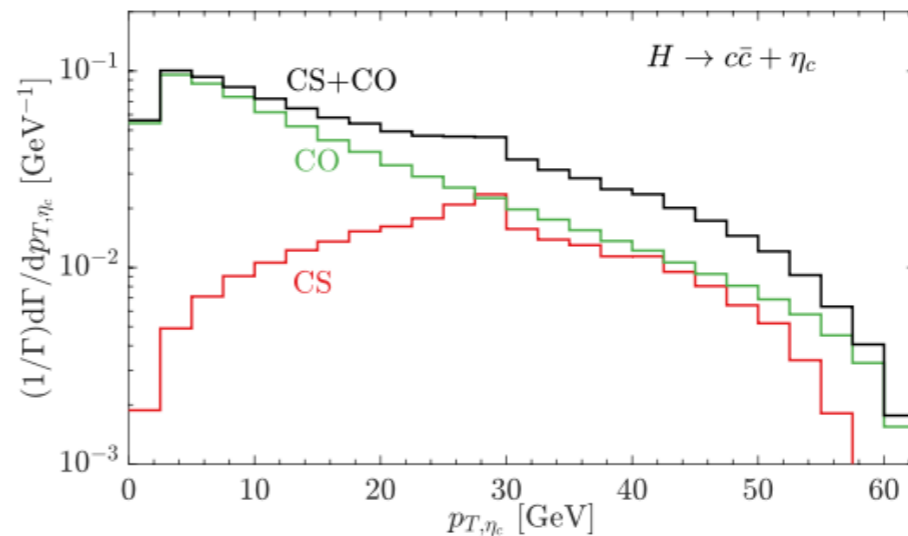
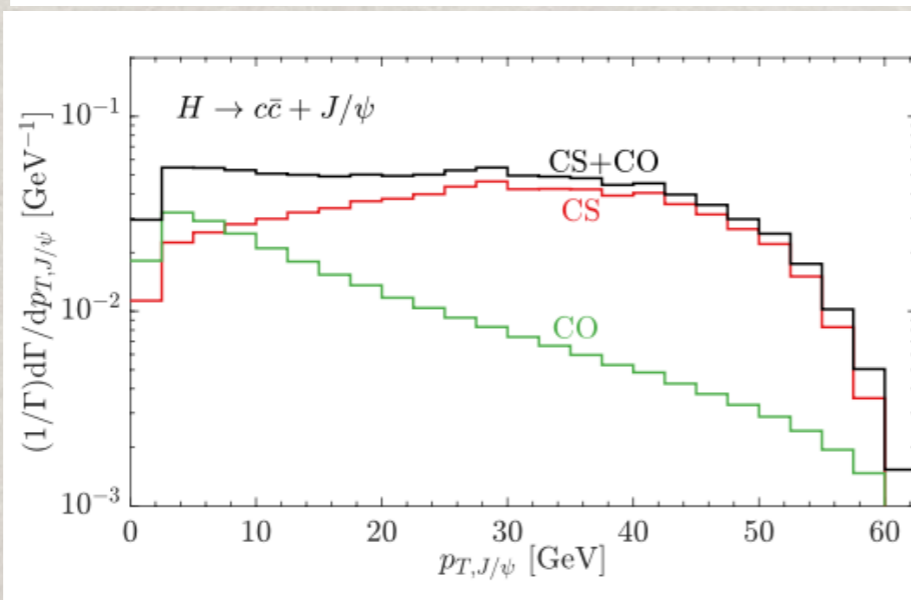
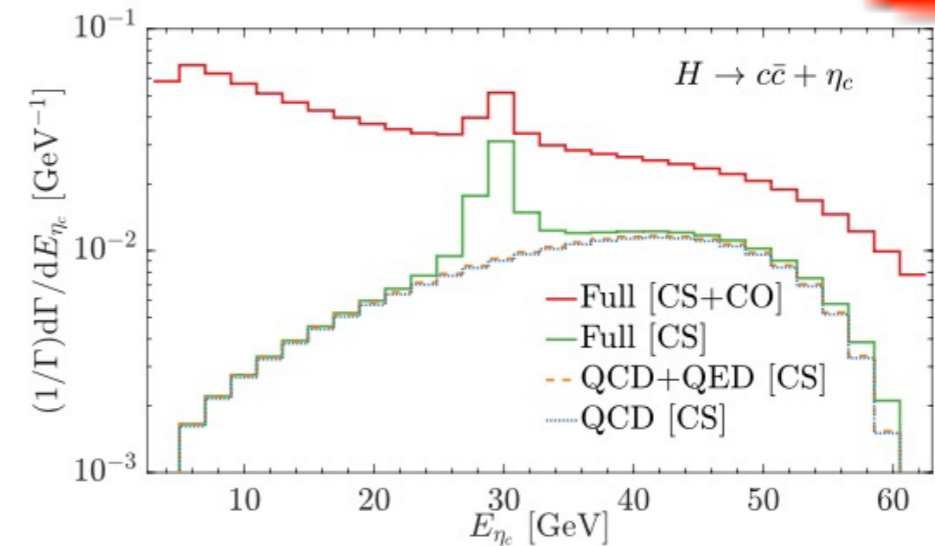
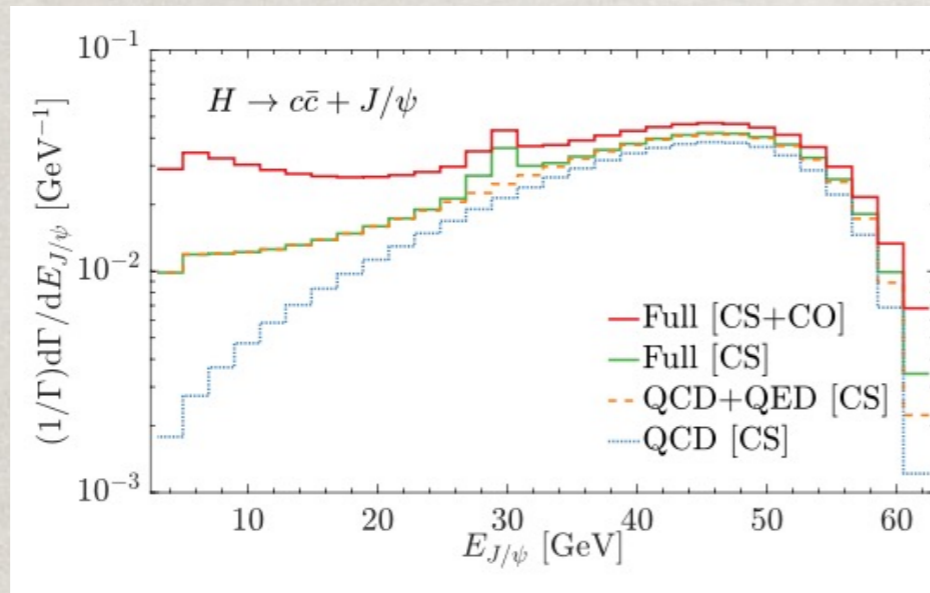
The LDMEs  $\langle \mathcal{O}^h[{}^{2S+1}L_J^{\text{color}}] \rangle$  need to be fitted from experimental data

Reference	$\langle \mathcal{O}^{J/\psi}[{}^1S_0^{[8]}] \rangle$	$\langle \mathcal{O}^{J/\psi}[{}^3S_1^{[8]}] \rangle$	$\langle \mathcal{O}^{J/\psi}[{}^3P_0^{[8]}] \rangle / m_c^2$
G. Bodwin,	$(9.9 \pm 2.2) \times 10^{-2}$	$(1.1 \pm 1.0) \times 10^{-2}$	$(4.89 \pm 4.44) \times 10^{-3}$
K.T. Chao,	$(8.9 \pm 0.98) \times 10^{-2}$	$(3.0 \pm 1.2) \times 10^{-3}$	$(5.6 \pm 2.1) \times 10^{-3}$
Y. Feng,	$(5.66 \pm 4.7) \times 10^{-2}$	$(1.77 \pm 0.58) \times 10^{-3}$	$(3.42 \pm 1.02) \times 10^{-3}$

# Our results: TH, A. Leibovich, Y. Ma, X.Z. Tan: aXive:2202.08273

## Decay width and branching fraction

	QCD [CS]	QCD+QED [CS]	Full [CS]	Full [CO]	Full [CS+CO]
$\Gamma(H \rightarrow c\bar{c} + J/\psi)$ (GeV)	$4.8 \times 10^{-8}$	$5.8 \times 10^{-8}$	$6.1 \times 10^{-8}$	$2.2 \times 10^{-8}$	$8.3 \times 10^{-8}$
$\text{BR}(H \rightarrow c\bar{c} + J/\psi)$	$1.2 \times 10^{-5}$	$1.4 \times 10^{-5}$	$1.5 \times 10^{-5}$	$5.3 \times 10^{-6}$	$2.0 \times 10^{-5}$
$\Gamma(H \rightarrow c\bar{c} + \eta_c)$ (GeV)	$4.9 \times 10^{-8}$	$5.1 \times 10^{-8}$	$6.3 \times 10^{-8}$	$1.8 \times 10^{-7}$	$2.4 \times 10^{-7}$
$\text{BR}(H \rightarrow c\bar{c} + \eta_c)$	$1.2 \times 10^{-5}$	$1.2 \times 10^{-5}$	$1.5 \times 10^{-5}$	$4.5 \times 10^{-5}$	$6.0 \times 10^{-5}$



Higgs production cross section at LHC  $\sigma_H \sim 50\text{M} / \text{ab}^{-1}$

Detection efficiency  $\varepsilon$  for the final state  $c\bar{c} + \ell^+\ell^-$

$\text{BR}(J/\psi \rightarrow \ell^+\ell^-) \sim 12\%$ ,  $\text{BR}(H \rightarrow J/\psi + c\bar{c}) \sim 2 \times 10^{-5}$

Event number

$$N = L\sigma_H \varepsilon \text{BR}(\ell^+\ell^-) \approx 24 \kappa_c^2 \times \frac{L}{\text{ab}^{-1}} \times \frac{\varepsilon}{20\%}$$

Considering the statistical error only  $\delta N \sim \sqrt{N}$  gives

$$\Delta\kappa_c \approx 10\% \times \left( \frac{L}{\text{ab}^{-1}} \times \frac{\varepsilon}{20\%} \right)^{-1/2}$$

**With  $\varepsilon \sim 20\%$ , we see  $\Delta\kappa_c \sim 6\%$  at ATLAS and CMS.** for  $3 \text{ab}^{-1}$

Formidable backgrounds:

Prompt  $J/\psi$  production  $\text{BR}(J/\psi \rightarrow \mu^+\mu^-) \times \sigma(pp \rightarrow J/\psi) \simeq 860 \text{ pb}$

Cross section of  $pp \rightarrow J/\psi + c\bar{c}$  at high  $p_T \sim 25 \text{ fb}$

→ At the end, should be better than  $J/\psi + \gamma: \kappa_c \sim 50$

→ May not beat  $W/Z+H \rightarrow W/Z+cc: \kappa_c \sim 3$

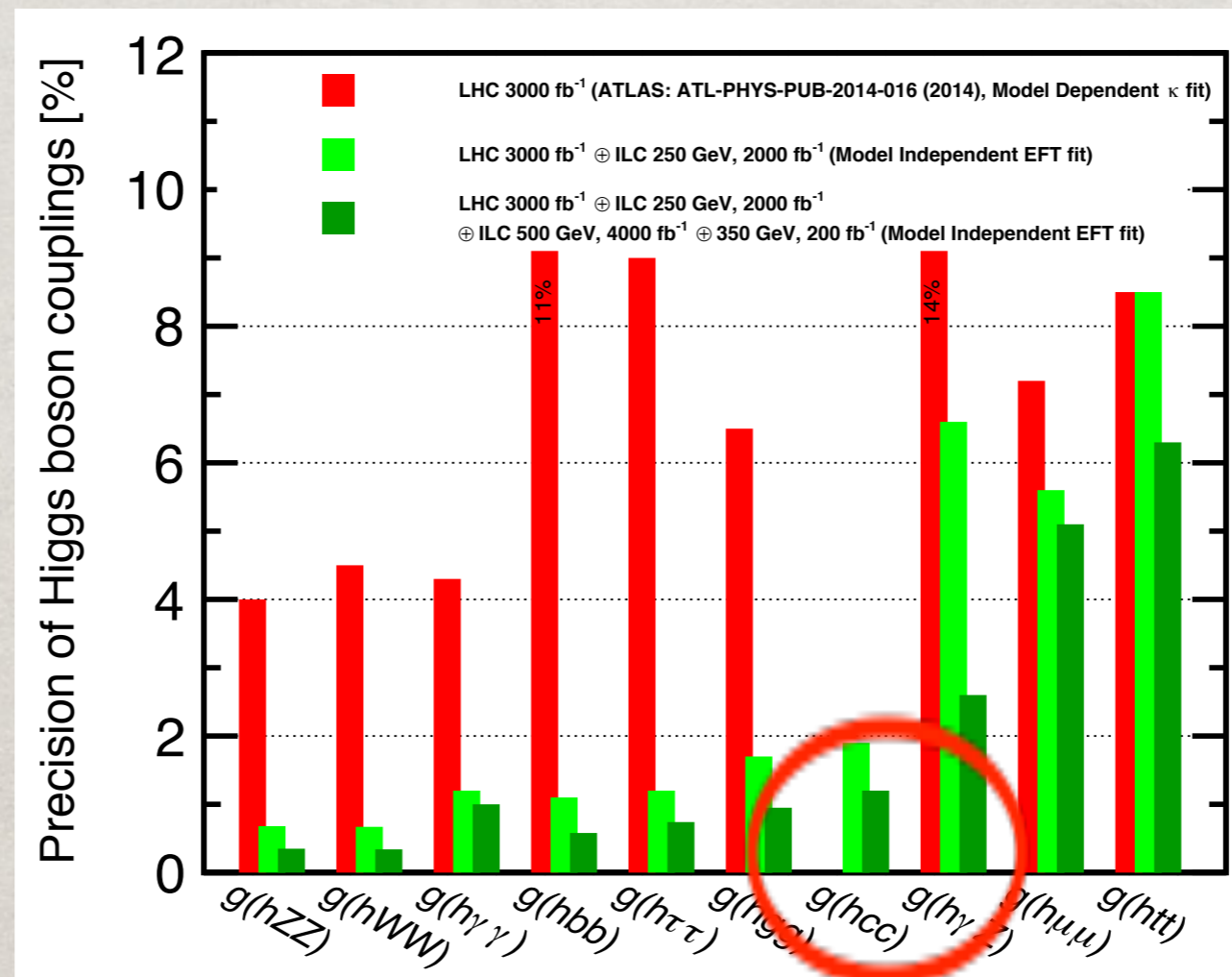
**Active study/simulation on-going!**

# SUMMARY ON HCC COUPLING

$\kappa_c$  sensitivity @ HL-LHC:

$gg \rightarrow h \rightarrow c \bar{c} \gamma$  : 8.3    Vh: hc kinematics: Global fit:  
 charm+ $\gamma$  in VBF : 5.4    2.5    3.0    1.5  
 $h \rightarrow J/\psi + c\bar{c}$ : hopeful    Need to combine all !

Snowmass report  $\rightarrow$  lepton colliders help!



**$Y_c$  measurement challenging & important!**



# HIGGS BOSON PROPERTIES:

**Run 1 + 2016 results:** 15 Feb 2020, PLB 805 (2020) 135425

$125.38 \pm 0.14 \text{ GeV}$  ( $\sim 1 \text{ ppm}$ )  $\Gamma_H = 4.0^{+1.3}_{-1.0} \text{ MeV}$

**Run 1 + Run 2:**

expect precision better than 100 MeV

**HL-LHC:**

Statistical uncertainty  $\sim 10 \text{ MeV}$

The game will be all about reducing systematics

**Latest updates for the well-established decays ( $>5\sigma$ ):**

$H \rightarrow \gamma\gamma$ :  $\mu = 1.12 \pm 0.09$  12 Mar 2021, arXiv:2103.06956 [Run 2]

$H \rightarrow ZZ$ :  $\mu = 0.94 \pm 0.11$  8 Mar 2021, arXiv:2103.04956 [Run 2]

$H \rightarrow WW$ :  $\mu_{\text{fid}} = 1.05 \pm 0.12$  4 July 2020, JHEP 03 (2021) 003 [Run 2]

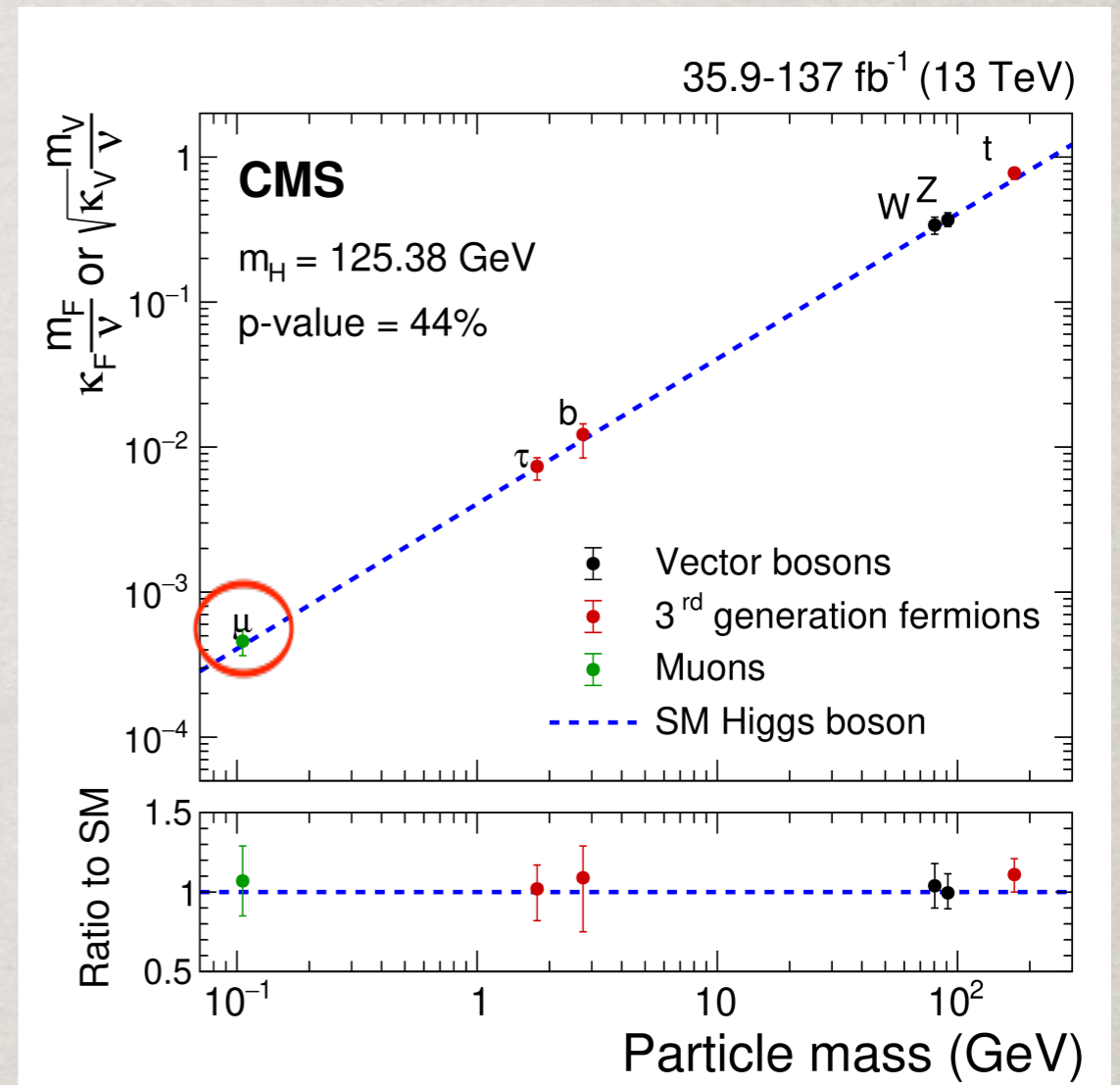
$H \rightarrow \tau\tau$ :  $\mu = 0.85 \pm 0.12$  31 July 2020, HIG-19-010 [Run 2]

$H \rightarrow bb$ :  $\mu = 1.04 \pm 0.20$  PRL 121 (2018) 121801 [2016+2017 data]

**Evidence for  $H \rightarrow \mu\mu$**  9 Sep 2020, JHEP 01 (2021) 148 [Run 2]

– Significance 3.0

– Signal strength  $\mu = 1.19 \pm 0.42$  **Korytov@NPHP21**



- Monumental achievements in theory & experiment!
- Higgs boson brings in new fundamental questions...
- A driver in precision & energy frontiers!

**The next target: the Higgs boson to charm!**

