Precision measurements in $t\bar{t}$ final states

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Introduction

Cross sections of 830 pb (920 pb) at 13 TeV (13.6 TeV) \rightarrow about 100M t \bar{t} pairs in Run 2



(Differential) cross section measurements:

- in all decay channels; at parton and particle level; including high $p_{\rm T}$ (boosted) top quarks
- used to extract m_t, charge asymmetry, PDFs ...
- provides interface to theory: test against new SM or BSM calculations
- but relies on unfolding

Direct extraction of parameters from data distribution:

- *m*t, top quark Yukawa coupling ...
- avoids unfolding, easier to constrain uncertainties from control regions.
- but depends directly on theory predictions

Full spectrum differential $t\bar{t}$ cross sections measurements in e/μ +jets events

137 fb⁻¹, 13 TeV, CMS: *Phys. Rev. D* 104, 092013

Analysis uses combination of resolved and boosted reconstructions:

- resolved: 1 isolated e/ μ and at least 4 jets (2 categories with tight and relaxed b-tagging requirement)
- boosted t_1 : 1 non-isolated lepton within $\Delta R < 0.6$ of a b jet. Selection using neural network (NN) based on isolation and kinematic variables.
- boosted t_h : 1 anti- k_T (size=0.8) jet with $p_T > 400$ GeV. Identified using NN with information based on subjets (H_{NN} fitted for background subtraction).
- combined fit to extract cross sections in resolved and boosted categories simultaneously



BHRL: boosted $t_{\rm h},$ resolved $t_{\rm l};$ BHBL: boosted $t_{\rm h}$ and $t_{\rm l}$

Uncertainties



- depending on the phase space region the precision is limited by systematic or statistical uncertainty.
- systematic dominated by experimental uncertainties like jet energy calibration and b-tagging.
- particle-level measurements reduce dependency on theoretical predictions (avoids extrapolation to full phase space)





- *p*_T well described by NNLO calculation.
- trend of harder spectrum in NLO calculations disappears above 600 GeV

Comparison of measurements to various predictions using χ^2 -tests uncertainties in measurements and predictions are take into account.



Most of the predictions are in good agreement with the measurement—with a few exceptions:

- m(tt) vs. p_T(t_h) and p_T(tt) vs. p_T(t_h) shows largest disagreements. but theory uncertainties might be underestimated using fully correlated scale variations
- at particle level additional jets vs. kinematic observable are not well described. *depends strongly on PS tuning*

Use cases

[M.V. Garzelli, S.-O. Moch, S. Zenaiev] Measurements compared to NNLO calculations (MATRIX) to extract m_t or PDFs:



JHEP 04 (2022) 079: calculation of $t\bar{t}$ NNLO+PS with POWHEG MINNLO_{PS} (direct comparison of particle level measurements to NNLO calculation)



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 $\mathrm{t}\bar{\mathrm{t}}$ precision measurements

Measurement of the $t\bar{t}$ charge asymmetry with highly Lorentz-boosted top quarks

138 fb⁻¹, 13 TeV, Acc. by Phys. Lett. B

Asymmetry only expected for $q\bar{q}$ induced $t\bar{t}$ production. At high $m(t\bar{t}) > 750 \text{ GeV}$ the contribution of valence quarks increases. Better sensitivity to BSM contributions expected.

Select e/μ +jets with high $p_{\rm T}$ top quarks

- t_h : one jet (size 0.8) containing all top decay products, one jet (0.4) + a jet (0.8) with W boson decay products, or fully resolved
- t₁: non-isolated lepton ($p_T(\mu) > 55$, $p_T(e) > 80 \text{ GeV}$) + jet with $\Delta R(lj) < 0.4$ or p_T relative to jet > 25 GeV



Unfold to full and fiducial phase space: (migration matrix A, unfolded cross sections μ):

$$\mathcal{L}_{k} = \prod_{j=1}^{N_{\text{reco}}} P\left(n_{j}; \sum_{i=1}^{N_{\text{gen}}} A_{ji}(\vec{\delta_{u}}) \mu_{i} + b_{j}\right) N(\vec{\delta_{u}})$$



Obtain asymmetry:

$$A_{\rm C} = \frac{N(\Delta|y|>0) - N(\Delta|y|<0)}{N(\Delta|y|>0) + N(\Delta|y|<0)}$$



Results are compatible with SM expectation

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Differential $\mathrm{t}\bar{\mathrm{t}}$ cross sections in allhadronic events: boosted

36 fb⁻¹, 13 TeV, PRD 103 (2021) 052008

- 2 jets (size: 0.8) $p_{\rm T} >$ 400 GeV, soft-drop mass: 120 220 GeV, b-tagged
- NN based on n-subjettiness used to define signal and control region:



- CR: NN < 0.8 normalized using $m_{\rm t}$ -fit
- Unregularized unfolding (matrix inversion)
 - good agreements of shapes between predictions and measurements in boosted regime



Measurement of multi-differential $\mathrm{t}\bar{\mathrm{t}}$ cross sections in dilepton events

138 fb⁻¹, 13 TeV, CMS: *CMS-TOP-20-006*

- selection: ee, $e\mu$, $\mu\mu$ at least 2 jets, at least 1 b jet.
- find analytic solutions for neutrino momenta; use solution with lowest $m(t\bar{t})$. Repeat procedure 100 times varying objects within their resolutions.



- uncertainty dominated by experimental uncertainties
- NNLO calculation improves description of $p_{\rm T}$ spectrum

Interpretation of double-differential $\mathrm{t}\bar{\mathrm{t}}$ measurements



36 fb⁻¹, 13 TeV, CMS: EPJC 80 (2020) 658



- m(tt̄) vs |y(tt̄)| shown to be sensitive to high x-gluon PDFs.
- these measurements are/will be included in new PDF fits.

Extraction of $\mathrm{t}\bar{\mathrm{t}}$ polarization and spin correlation

36 fb⁻¹, 13 TeV, *Phys. Rev. D* 100 (2019) 072002

• in helicity-frame the differential cross section of particles from tt decay can be parameterized by the polarization vector B and spin correlation matrix C:

$$\frac{d^4\sigma}{d\Omega d\bar{\Omega}} \propto 1 + \kappa \vec{B} \cdot \Omega + \bar{\kappa} \vec{B} \cdot \bar{\Omega} + \kappa \bar{\kappa} \Omega \cdot C \cdot \bar{\Omega}$$

• use dilepton events with same reconstruction as in differential cross section measurement.





- from a set of unfolded distributions extract all parameters of B and C
- measured polarization and spin correlation in agreement with SM expectations



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$t\bar{t}$ precision measurements

Extraction $m_{\rm t}^{\rm pole}$ from kinematic distributions

36 fb⁻¹, 13 TeV, *CMS-TOP-21-008 acc. JHEP*

- dilepton events use same reconstruction as in differential cross section measurement.
- sensitive variable $\rho = 2m_0/m_{t\bar{t}+jet}$ presence of additional jet increases sensitivity on m_t^{pole}
- resolution improved based on regression using NN (input 10 kinematic parameters mostly direct calculations of ρ and invariant object masses)





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Pred. / Data

 $t\bar{t}$ precision measurements

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Direct measurement of $m_{\rm t}$ (MC parameter) in e/μ +jets events

36 fb⁻¹, 13 TeV, CMS-TOP-20-008 Sub. to Eur. Phys. J. C

- select events with e/μ and >= 4jets
- perform kinematic fitting with constraints of two equal top quark masses and W mass
- goodness of fit also used to determine best parton-jet assignment
- up to 5 distributions used to extract m_t







- ullet best fitted value $m_{
 m t}=171.77\pm0.37\,{
 m GeV}$
- leading uncertainty is final-state PS scales 0.21 GeV in contrast to previous measurements uncorrelated scales for different splittings used
- modelling very important for this measurement

Direct measurement of top quark Yukawa coupling in dilepton events

137 fb⁻¹, 13 TeV, Phys. Rev. D 102 (2020) 092013



Calculations of EW corrections to $m(t\bar{t})$ and $|\Delta y_{t\bar{t}}|$ show significant dependency on the Y_t



- select events with ee, $\mu\mu$, $e\mu$ plus 2 b jets
- to avoid systematic unc. in $p_{\rm T}^{\rm miss}$ reconstruction, $m_{\rm blbl}$ and $\Delta y_{\rm bl}$ are used





- best fitted value $Y_t = 1.16^{+0.24}_{-0.35}$
- leading uncertainties
 - jet energy corrections
 - uncertainty in EW correction (in particular combination NLO(NNLO) QCD + EW)
 - parton shower modelling

Conclusion





Measurement of differential cross sections:

- performed in all channels
- even in low populated phase space regions
- reaches great precision: uncertainties mostly dominated by experimental sources
- to some degree dealing with theory uncertainties can be postponed (in particular for fiducial measurements)
- further used to extract important parameters: $m_{\rm t}$, $t\bar{t}$ spin correlations, charge asymmetry, PDFs

Direct extractions of parameters:

- $\bullet\,$ can provide very precise measurements $m_{\rm t},\,Y_{\rm t}$ if experimental and theoretical uncertainties can be constrained
- depends on precise understanding of theory and its uncertainties.

Backup

Perform χ^2 -fit to extract the cross sections:



$$\chi^{2} = \sum_{\ell} \sum_{y} \sum_{r} (\mathbf{m}_{yr\ell} - \mathbf{b}_{yr\ell}(\nu_{\alpha}) - M_{yr\ell}(\nu_{\alpha})\boldsymbol{\sigma})^{T} \boldsymbol{C}_{yr\ell}^{-1} (\mathbf{m}_{yr\ell} - \mathbf{b}_{yr\ell}(\nu_{\alpha}) - M_{yr\ell}(\nu_{\alpha})\boldsymbol{\sigma}) + \kappa(\nu_{\alpha})$$

- **m** measured distribution of events with the covariance matrix *C*; per year *y*, reconstruction *r* method (res. 1t1l, res. 2t, boosted), and lepton channels ℓ (18 categories)
- σ vector of cross sections (free parameters of interest).
- ν_{α} nuisances representing the uncertainties. These are constrained in $\kappa(\nu_{\alpha})$ taking into account year-by-year correlations.
- $M(\nu_{\alpha})$ response matrices that map σ to the $t\bar{t}$ event yields at detector level.
- $\mathbf{b}(\nu_{\alpha})$ non $t\bar{t}$ background.
- No regularization condition is used



- worst description of $p_{\rm T}$ in events without additional jet.

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Double-differential measurement of $p_{\rm T}(t\bar{t})$ vs. $p_{\rm T}(t_{\rm h})$



- $p_{\rm T}({
m t_h})$ prediction too hard in low $p_{\rm T}({
m t\bar{t}})$ -bin. Agreement better at high $p_{\rm T}({
m t\bar{t}})$.

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 $t\bar{t}$ precision measurements

Inclusive Cross Section

Obtain inclusive cross section from integration of differential distributions:



• Result of best expected measurement $(m(t\bar{t})vs\cos(\theta^*))$, also best observed:

$$\sigma_{e/\mu + \text{jets}} = 227.6 \pm 6.8 \,\text{pb.}$$

• With a branching fraction of $28.77 \pm 0.32\%$ to e/μ +jets:

 $\sigma_{t\bar{t}} = 791 \pm 25 \,\text{pb} \,(\pm 1(\text{stat}) \pm 21(\text{sys}) \pm 14(\text{lumi}) \,\text{pb})$

With 3.2% uncertainty, one of the most precise $t\bar{t}$ cross section measurements. Dominant uncertainties: jet energy, lepton identification, NNLO reweighting of NLO MC