

Dark Sector @ PIP2-BD

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Physics Opportunities at Beam Dump Facility in PIP-II and Beyond

Introduction

- New Physics: Dark matter(DM), neutrino masses and mixing, baryon abundance and various anomalies, $g-2$ of muon, MiniBooNE etc.

Are they all correlated? Is there a model?

- Where is the new physics scale?
- Many experiments are probing new physics scales: DM direct and indirect detections, LHC, neutrino experiments, beam dump experiments, rare decays, astrophysical observations etc.
- LHC is mostly probing scales above 1 GeV

Introduction

Investigation of scales below 1 GeV

- This region is difficult to search
- Anomalies, and puzzles can be addressed
- There are many new ideas

Models (Many ongoing activities):

Light mediators: scalar/pseudo-scalar, vector; sub-GeV DM

Low energy beam dump-based experiments, Forward physics facility at the LHC, Astrophysical observations, etc. can investigate low scale models

This talk will discuss: Exploration of various models at GeV scale proton beam dump-based stopped pion experiments:

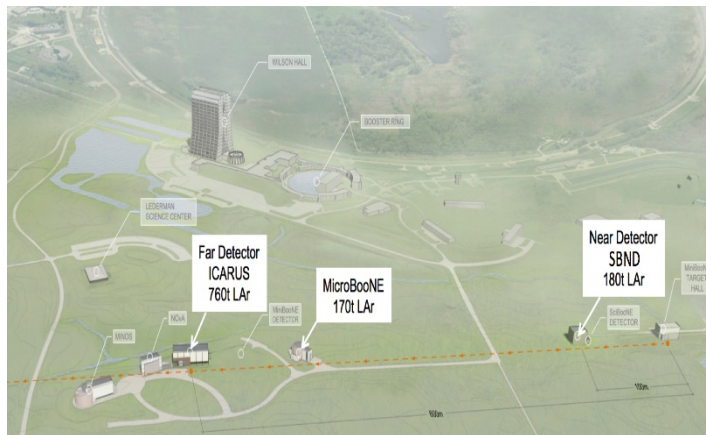
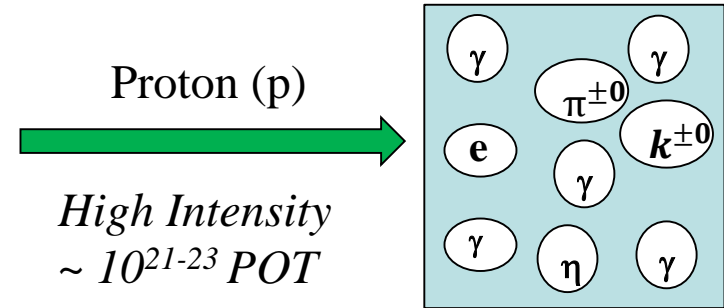
High intensity beam, various production possibilities, large detectors

Proton beam-based experiments

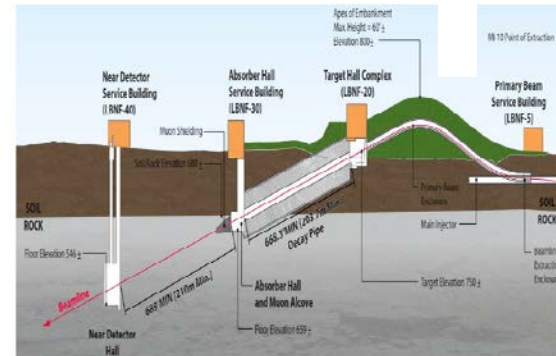
Beam dump-based (proton beam)

[ongoing]: 800 MeV-3 GeV: COHERENT (Oakridge), CCM (LANL), JSNS2(JPARC) Detectors, CsI, LAr, NaI, Ge

Fermilab SBN program: 120 GeV NUMI, 8 GeV BNB beams (ongoing)



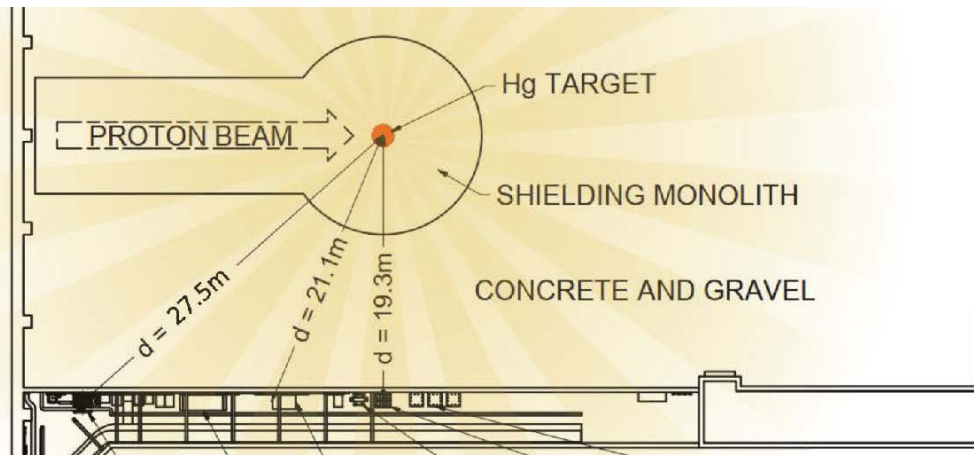
DUNE (120 GeV)



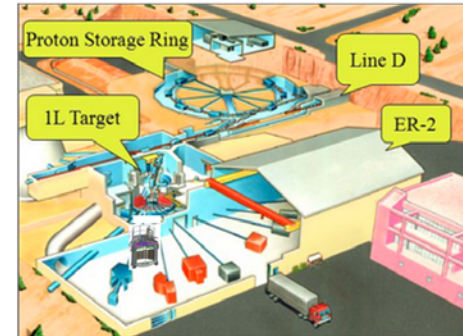
- FASER, FASER_v, SND are ongoing

How complementary are these searches?

Stopped pion experiments



800 MeV protons, 100kW, 275 nanosecond pulsed beam



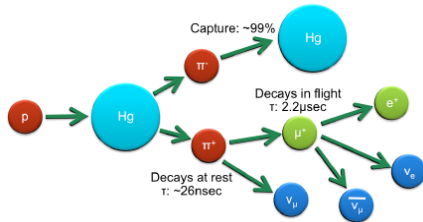
LANSCE-PSR-Lujan Target: Prolific source of charged/ neutral pions and photons that produce neutrinos and potential dark sector particles.

Creating a stopped-pion source with PIP-II: PIP2-BD

- PIP-II Accumulator Ring (PAR), Compact PIP-II Accumulator Ring (C-PAR), and Rapid Cycling Synchrotron Storage Ring (RCS-SR) are three accelerator scenarios we studied ahead of Snowmass 2022
- PAR and C-PAR are realizable in the timeframe of the start of the PIP-II accelerator and DUNE Phase I
- RCS-SR is a Booster Replacement scenario under ACE on the timescale of DUNE Phase II

Facility	Beam Energy (GeV)	Repetition Rate (Hz)	Pulse Length (s)	Beam Power (MW)
PAR	0.8	100	2×10^{-6}	0.1
C-PAR	1.2	100	2×10^{-8}	0.09
RCS-SR	2	120	2×10^{-6}	1.3

COHERENT

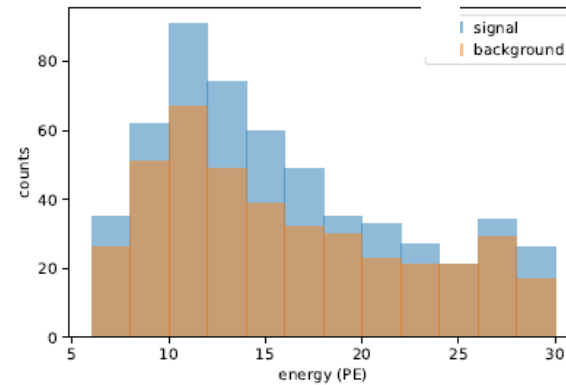
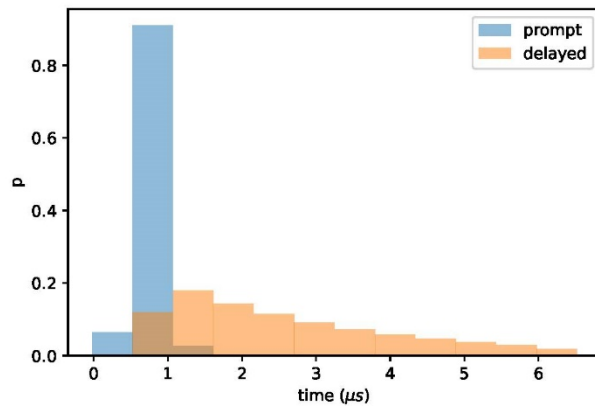


Proton beam hits a target Hg [COHERENT:
1 GeV]

$$\begin{aligned} \text{Prompt} \quad & \pi^+ \rightarrow \mu^+ + \nu_\mu \\ & \vdots \\ \text{Delayed} \quad & \mu^+ \rightarrow e^+ + \bar{\nu}_\mu + \nu_e \\ & \vdots \end{aligned}$$

CEvNS: Coherent elastic neutrino nucleus scattering, Nuclear recoil~ O(KeV)

COHERENT: CsI data:2018



COHERENT (2017) No CEvNS rejected at 6.7σ : CsI
(2020): 11.6σ

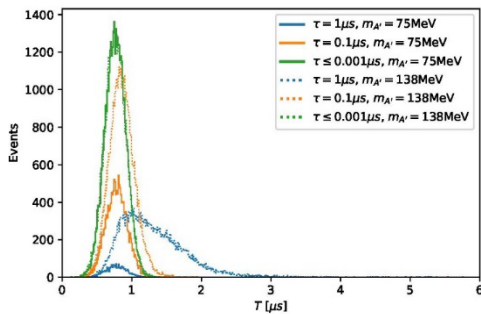
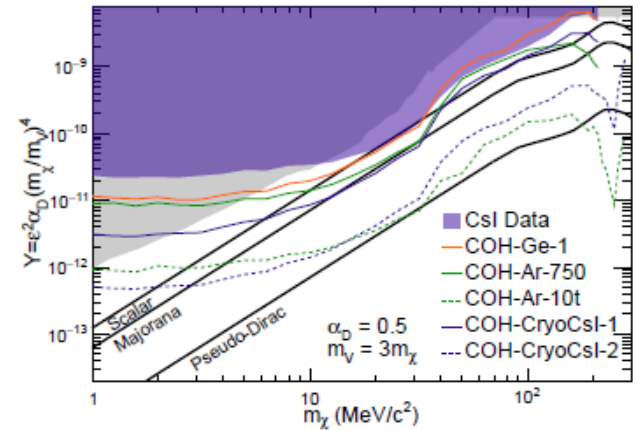
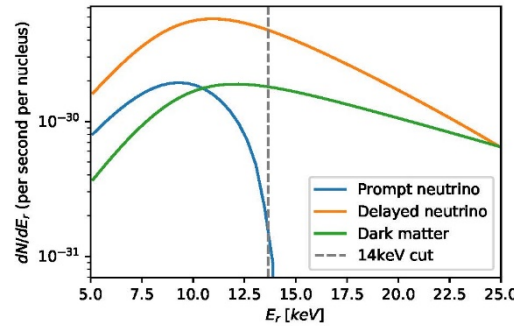
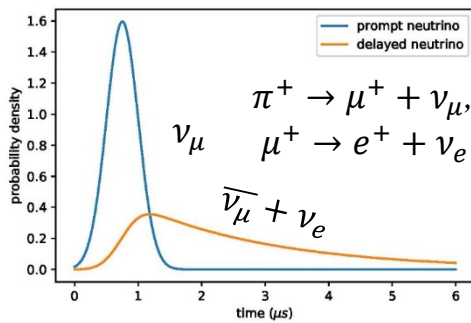
COHERENT (2020) No CEvNS rejected at 3.8σ : LAr

CENS-DM

The ongoing CE ν NS experiments, COHERENT, CCM etc. are probing light DM

Deniverville, Pospelov, Ritz, PRD, 2015, Ge, Shoemaker, JHEP, 2018, Dutta, Kim, Liao, Park, Shin, Strigari, PRL, 2020

- Both neutrinos, DM produce nuclear recoils: how to distinguish them?
- The timing and energy recoil measurement at COHERENT, CCM, JSNS² can be used



- For $t < 1.5\mu$ s, we have mostly prompt ν
- For $E_r > 14$ KeV (CsI) we remove the prompt ν
- ➔ We can remove the SM/NSI ν backgrounds
- ➔ Similar strategies apply to CCM, JSNS²

CCM: KeV to MeV

CCM recently showed sensitivities to energies from KeV to 100 MeV

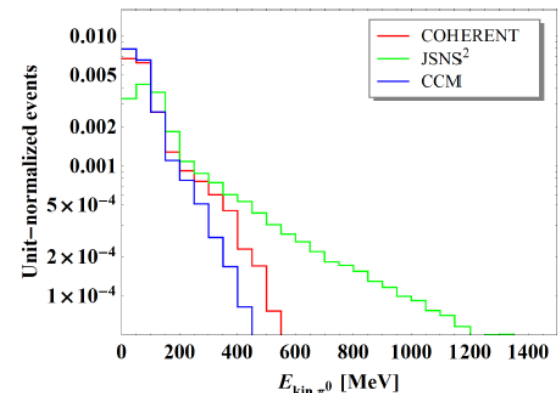
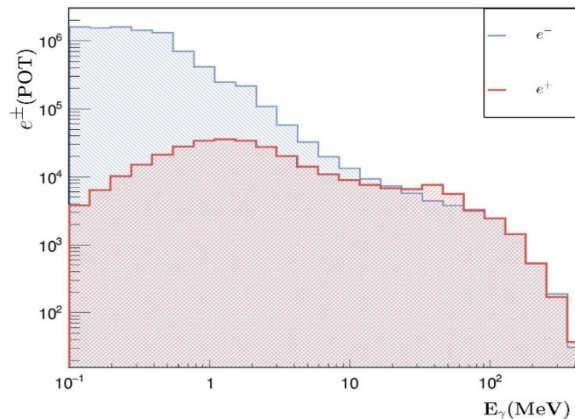
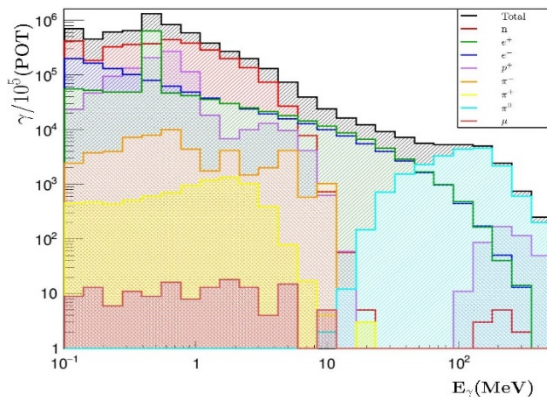
CCM, 2112.09979 (PRD)

This opens up various new possibilities:

Axion, Inelastic DM-nucleus scattering, Dark photons,
Light mediators: MiniBooNe anomaly, HNL, Mirror Neutrons etc.

New particles can be produced at the target using γ , e^\pm , pion flux

CCM/COHERENT: $0. \pi^+$ per proton: $10^{22/23}$ POT



CCM, 2112.09979

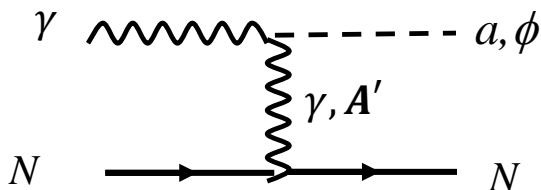
New physics at ν experiments

From γ :

A' : Vector
 ϕ =scalar
 a =pseudo-scalar

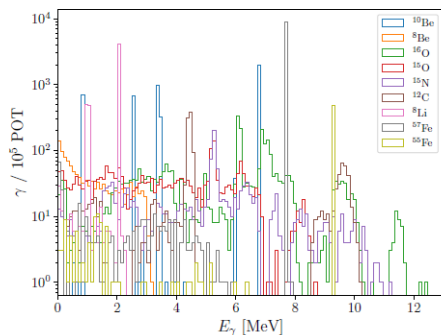


$$L \supset -\frac{\varepsilon}{4} F^{\mu\nu} F_{\mu\nu}^{(\prime)} - g_{a,\phi\gamma(Z')} \frac{(a, \phi)}{4} F^{\mu\nu} \tilde{F}_{\mu\nu}^{(\prime)}$$



Primakoff

Coherent scattering for γ exchange



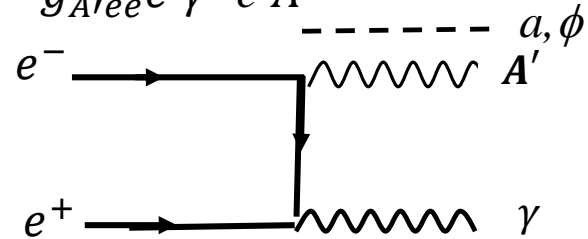
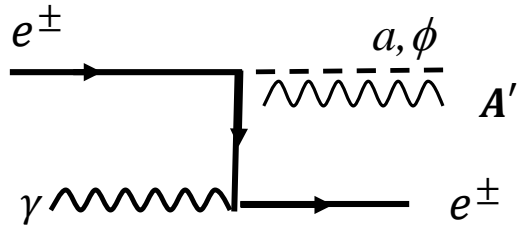
Deexcitation of N^* to N : produce a, s, A'

**L. Waites, A. Thompson, A. Bungau, J. Conrad, B. Dutta,
 W. Huang, D. Kim, M. Shaevitz, J. Spitz,
 arXiv:2207.13659**

New physics at ν experiments

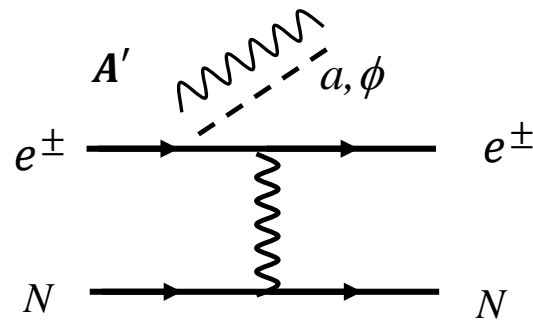
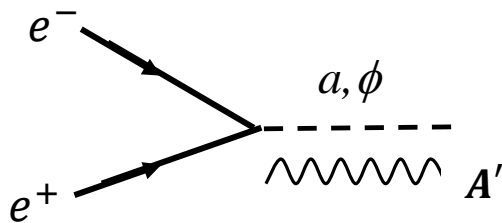
- From e^\pm : $L \supset -g_{\phi(a)ee} \bar{e} (i\gamma^5) e \phi(a) - g_{A'ee} \bar{e} \gamma^\mu e A'$

Compton



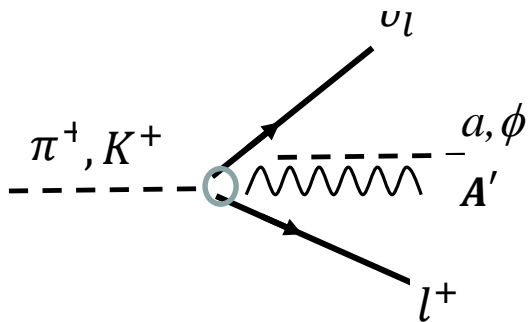
Associated

Resonance



Bremsstrahlung

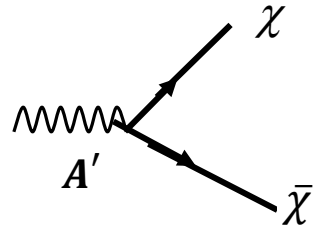
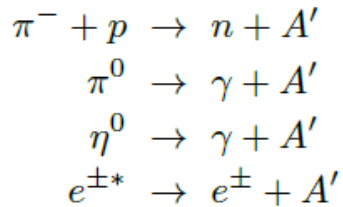
- Neutral meson decays $\eta^0, \pi^0 \rightarrow \gamma A'_\mu$
- From Charged meson decay: quarks and lepton couplings



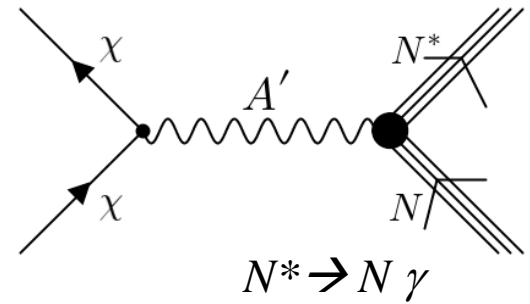
$$L \supset -g_{\phi(a)ff} \bar{f} (i\gamma^5) f \phi(a) - g_{A'ee} \bar{f} \gamma^\mu f A'_\mu$$

- Not helicity suppressed \rightarrow both electron and muon final states contribute
- Needs to include all the internal bremsstrahlung diagrams IB_i ($i=1,2,3$)

Example 1: MeV signal - DM



DM (χ) can be scalar/fermion



$$\begin{aligned} \frac{d\sigma_{\text{inel}}^{DM}}{dE_r} &= \frac{2e^2\epsilon^2 g_D^2 E'_\chi{}^2}{p_\chi p'_\chi (2m_N E_r + m_{A'}^2)^2} \frac{m_N}{2\pi} \frac{4\pi}{2J+1} \left\{ \sum_{J \geq 1, \text{spin}} \left[\frac{1}{2} (\vec{l} \cdot \vec{l}^* - l_3 l_3^*) \left(|\langle J_f | \hat{T}_J^{\text{mag}} | J_i \rangle|^2 + |\langle J_f | \hat{T}_J^{\text{el}} | J_i \rangle|^2 \right) \right] \right. \\ &+ \left. \sum_{J \geq 0, \text{spin}} \left[l_0 l_0^* |\langle J_f | \hat{M}_J | J_i \rangle|^2 + l_3 l_3^* |\langle J_f | \hat{L}_J | J_i \rangle|^2 - 2l_3 l_0^* \text{Re} \left(\langle J_f | \hat{L}_J | J_i \rangle \langle J_f | \hat{M}_J | J_i \rangle^* \right) \right] \right\} \end{aligned}$$

We use *Bigstick*: Shell Model code for this calculation

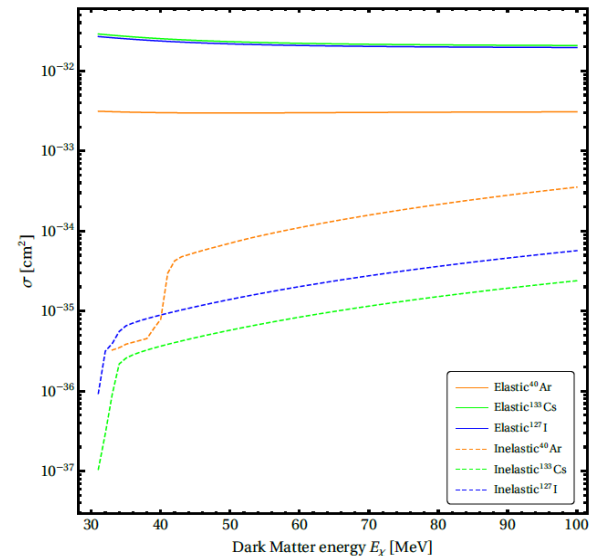
$$\begin{aligned} \frac{d\sigma_{\text{el}}^{DM}}{dE_r} &= \frac{e^2\epsilon^2 g_D^2 Z^2}{4\pi(E_\chi^2 - m_\chi^2)(2m_N E_r + m_{A'}^2)^2} F^2(q^2) \\ &\times \left[2E_\chi^2 m_N \left(1 - \frac{E_r}{E_\chi} - \frac{m_N E_r}{2E_\chi^2} \right) + E_r^2 m_N \right] \end{aligned}$$

For Ar40

$$\left(\frac{\text{Inelastic}}{\text{Elastic}} \right)_{\text{signal}} = 10^{-2} - 10^{-1}$$

$$\left(\frac{\text{Inelastic}}{\text{Elastic}} \right)_{\text{bkg}} = 10^{-4} - 10^{-3}$$

**B. Dutta, W. Huang,
J. Newstead, V. Pandey,
Phys.Rev.D 106 (2022) 11, 113006**

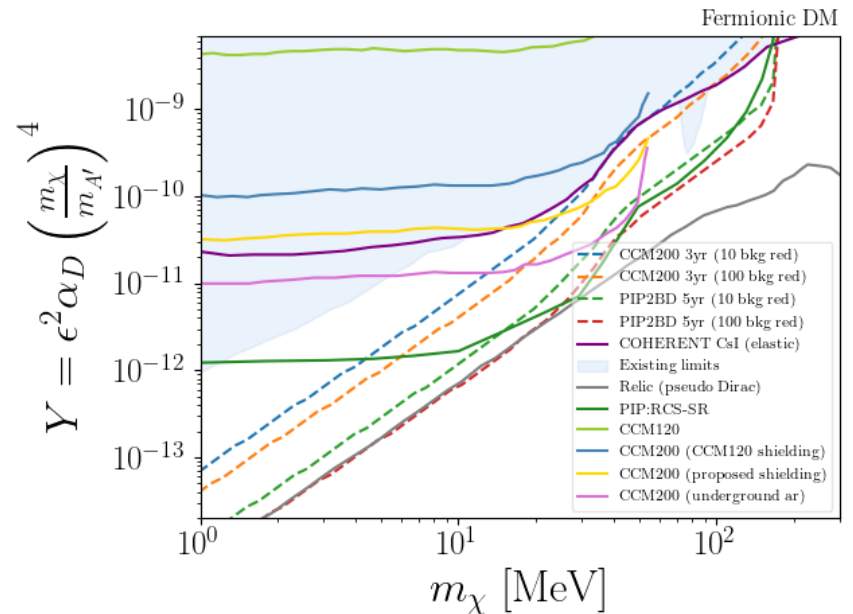
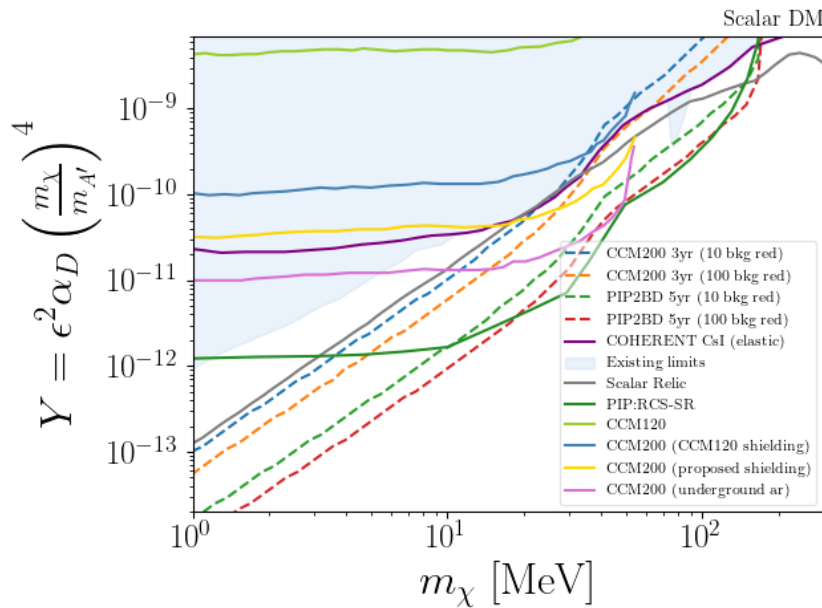


MeV signal - DM

$$g_D = \sqrt{2\pi}, \quad \frac{M_{A'}}{M_\chi} = 3 \quad \sigma_{\text{elastic}} \sim 1/(2m_N E_r + m_{A'}^2)^2$$

$$\approx 1/(2m_N E_r)^2 \text{ if } m_{A'}^2 < 2m_N E_r$$

B. Dutta, W. Huang, J. Newstead, arXiv:2302.10250

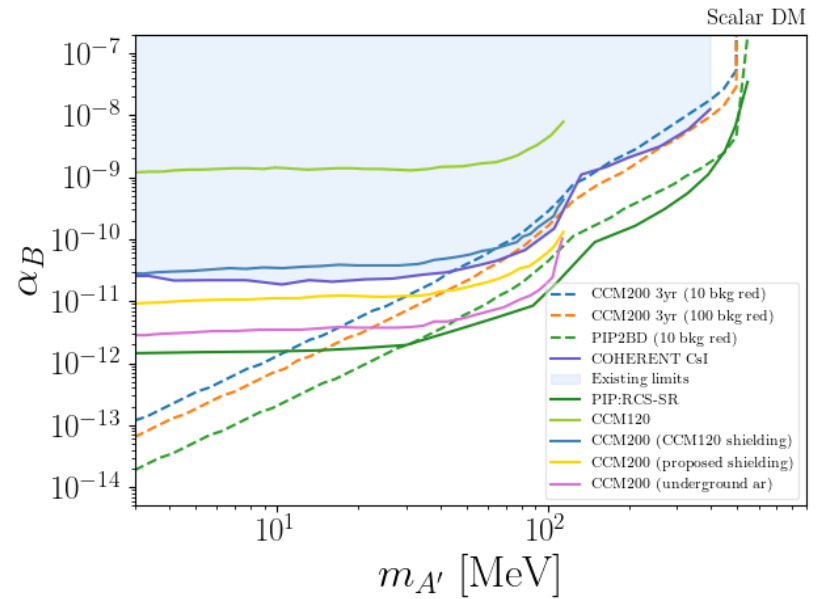
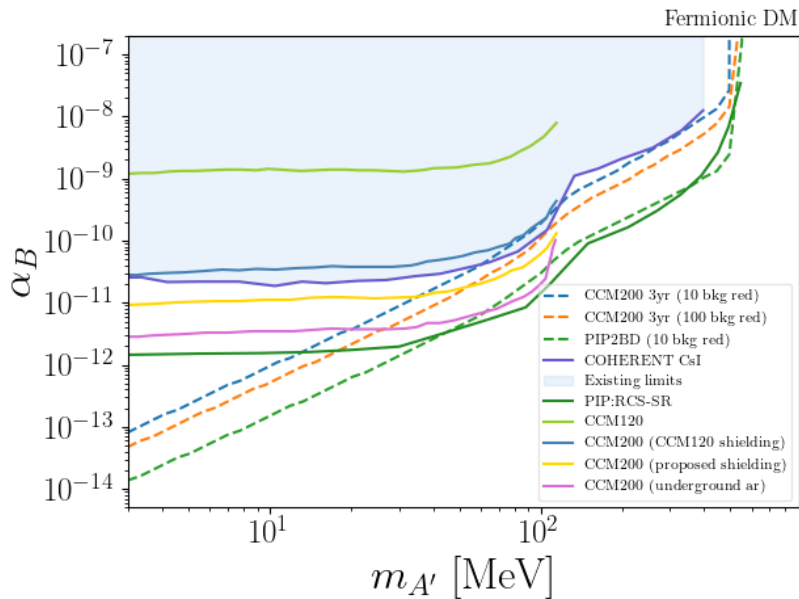


- We use CCM background measurements and projections
- $t < 200$ ns, prompt window reduces the neutrino background down to $O(1)$ events
- Rescale the shell model prediction to be consistent with the experiment, **W. Tornow et al., 2210.14316**
- A lower threshold detector will help to improve the sensitivity in the elastic channel
- Signal: Recoil (KeV), deexcitation photon

MeV signal – DM: Leptophobic

$U(1)_B$ Model

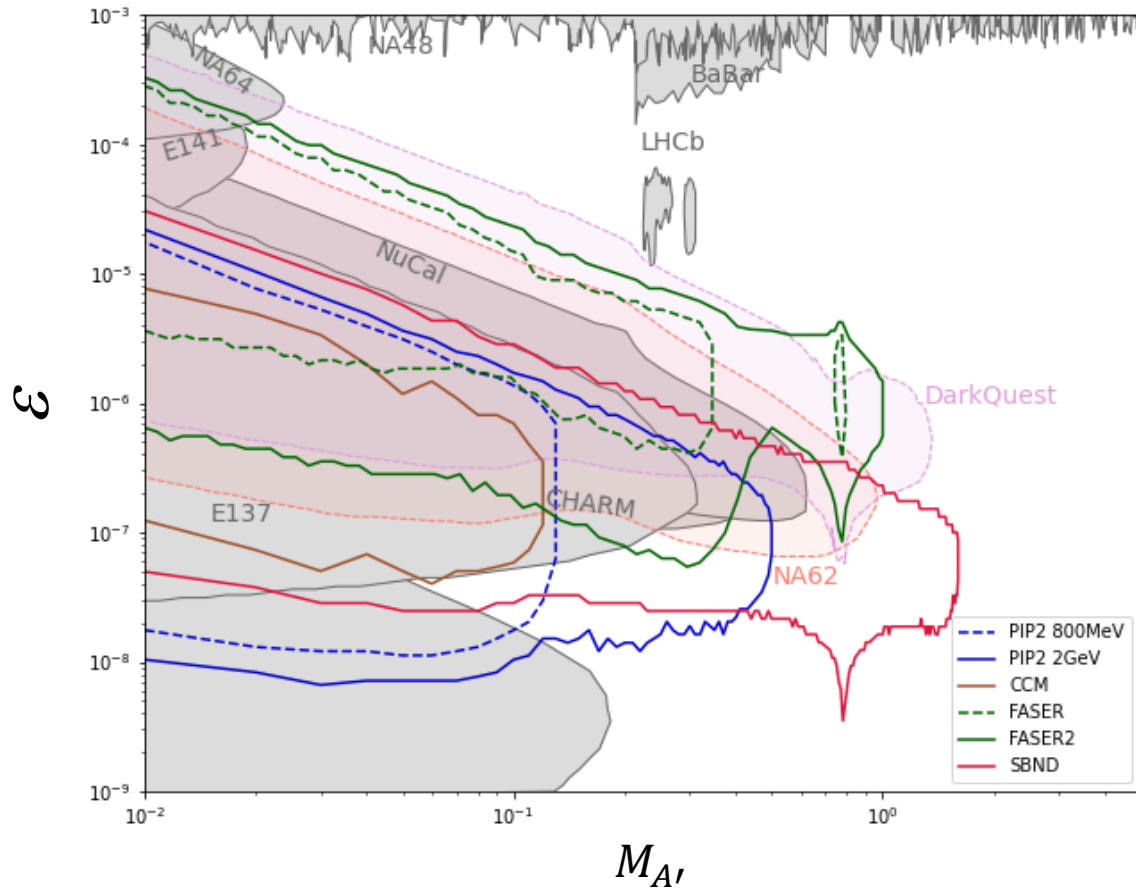
$$m_{A'}/m_\chi \approx 2$$



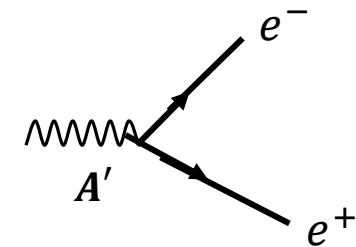
Can we implement inelastic scattering at SBN?

Example2: dark photon

Dark Photon



Decay final states



Scattering final states (ongoing)

Dutta, Karthikeyan, Kim, to appear

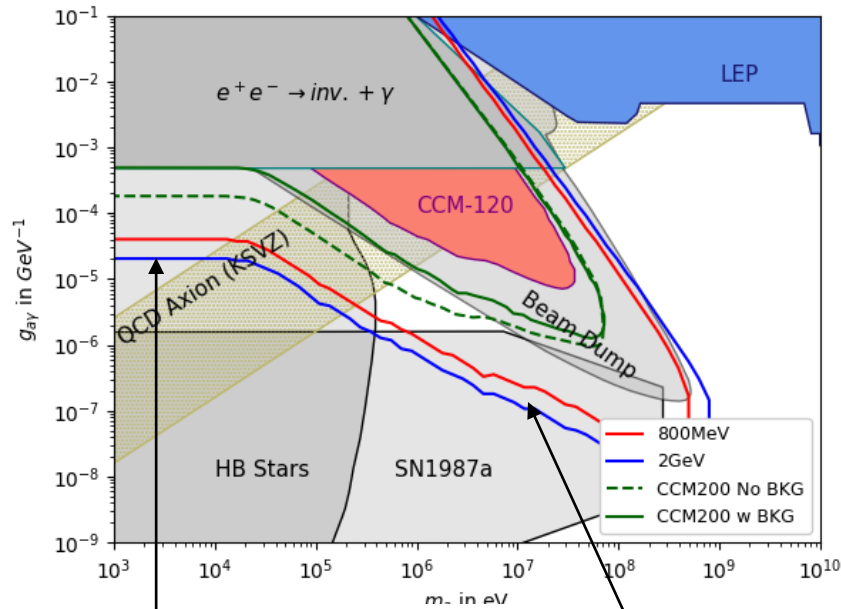
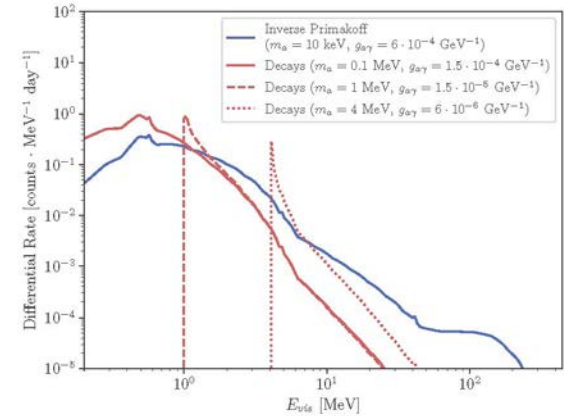
Productions involve: Charged and neutral meson decays, protons, electron/positron flux induces bremsstrahlung and various other production processes

Example 3: ALP

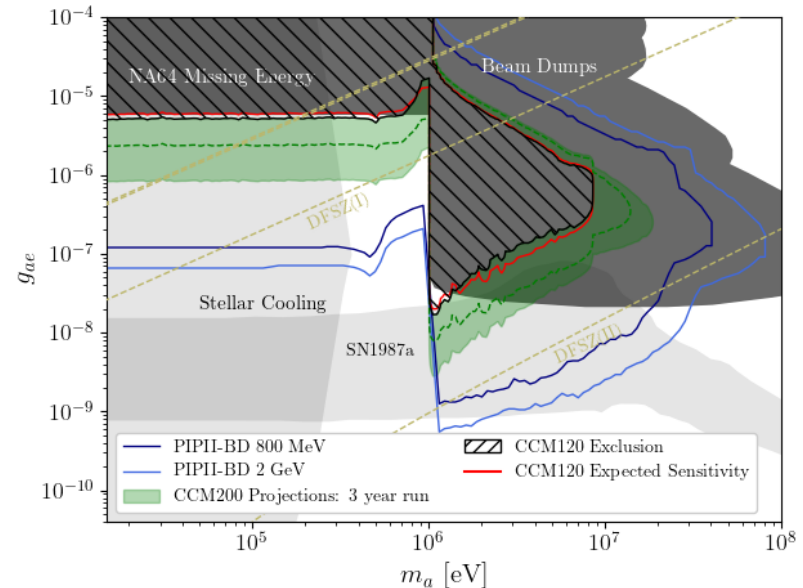
- γ , e^\pm fluxes, Primakoff/Compton productions, inverse Compton/Primakoff, decays final states are used by CCM(LANL) to explore new physics
- For g_{ae} coupling, the CCM engineering run data has started exploring new parameter space

$$L \supset -g_{a,\phi\gamma} \frac{a}{4} F^{\mu\nu} \tilde{F}_{\mu\nu}$$

CCM, 2112.09979



$$L \supset -g_{aee} \bar{e} (i\gamma^5) e a$$



Scattering

Decay

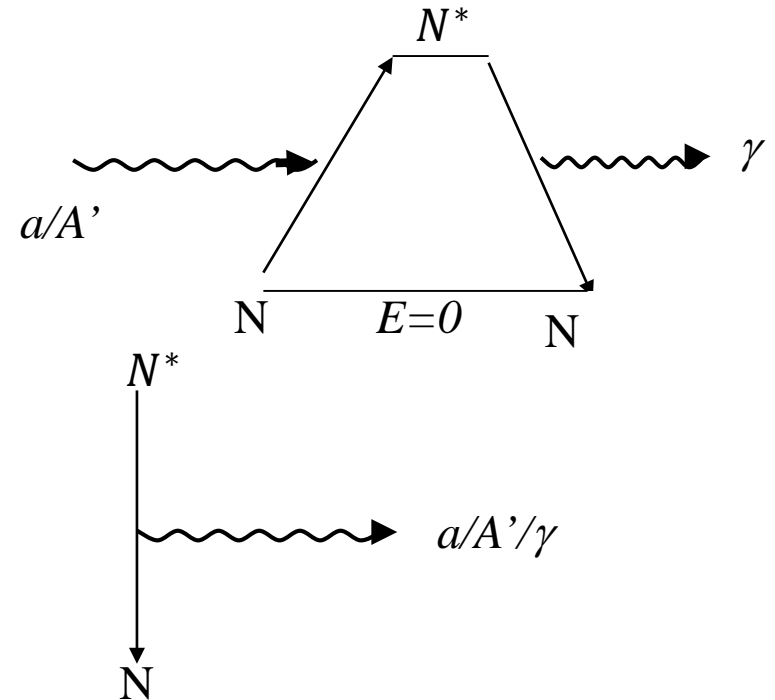
Example 3.1: Absorption

Additional detection and production modes:
 MeV Lines from inelastic absorption,
 e.g., axion absorption;

Axion productions from deexcitation
 lines at the target

Primakoff/ Compton production
 Inverse Primakoff/Compton detection

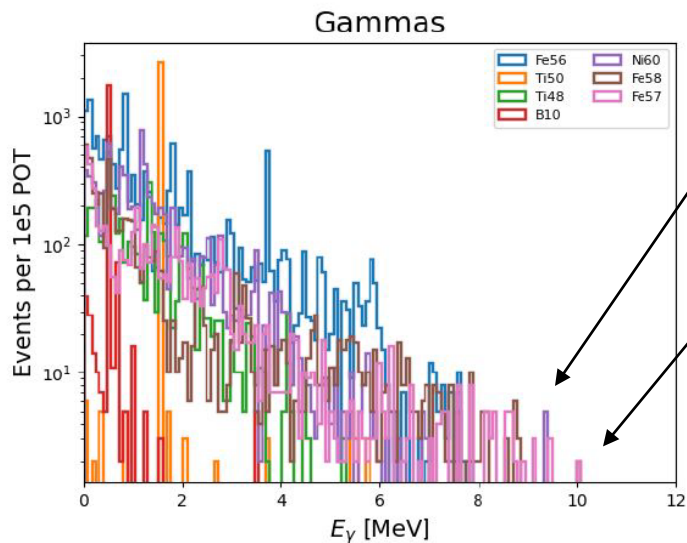
Deexcitation production
 Inelastic Absorption



$\mathcal{G}_{aee}, \mathcal{G}_{a\gamma\gamma}$

\mathcal{G}_{ann}

CCM/PIP2: ALP

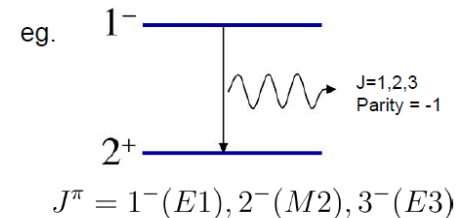
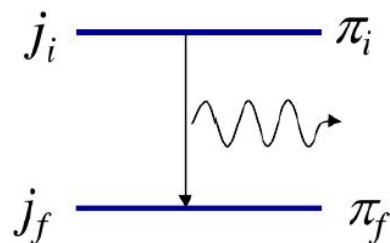


We use these lines to produce axions,
dark photon production

$$\pi = (-1)^I \text{ for E transitions}$$

$$\pi = (-1)^{I+1} \text{ for M transitions}$$

$$\text{multipolarity} = 2^I$$



I	π_γ	Radiation type	Label
1	-1	Electric Dipole radiation	E1
1	+1	Magnetic Dipole radiation	M1
2	-1	Magnetic Quadrupole radiation	M2
2	+1	Electric Quadrupole radiation	E2
3	-1	Electric Octapole radiation	E3
3	+1	Magnetic Octapole radiation	M3

$$|J_i - J_f| \leq J \leq J_i + J_f$$

$$\pi = \pi_i \pi_f$$

- These lines are also utilized at IsoDAR [60 MeV proton beam dump experiment with Be target]
- Can be utilized at PIP2 and DAMSA

CCM/PIP2: ALP

- $g_{ann}g_{\gamma\gamma}$

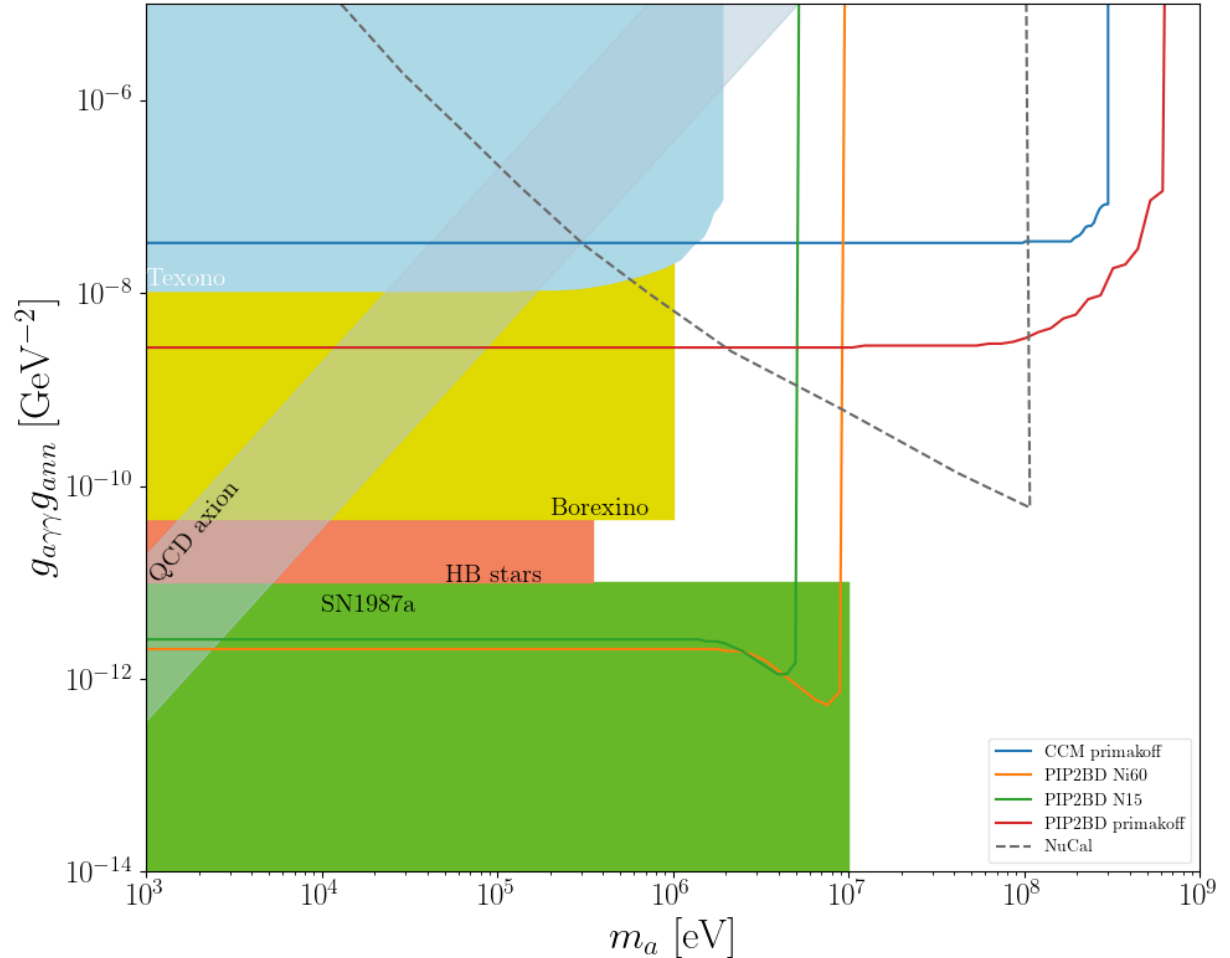
deexcitation line (target)+
inverse Primakoff/Compton
(detector);
Primakoff/Compton (Target)
+excitation (detector)

- g_{ann}

deexcitation line (target)+
+excitation (detector)

$$\mathcal{L}_{\text{int}} \supset \frac{\partial_\mu a}{f_a} \bar{N} \gamma^\mu \gamma^5 N$$

$$g_{ann} = \frac{1}{f_a}$$



Dutta, Newstead, Huang, to appear

- De-excitation line at the target is due to absorption:

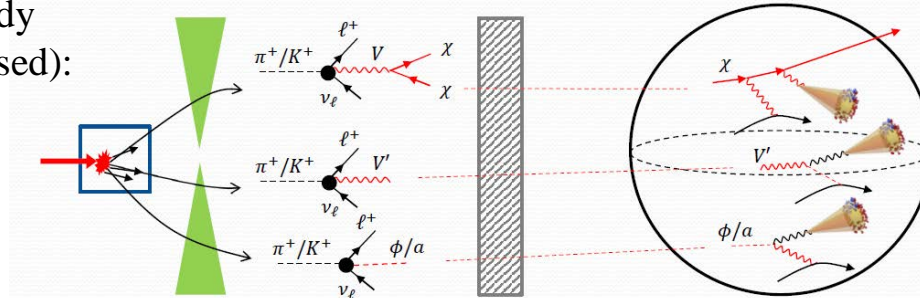
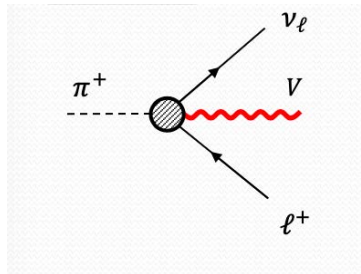
$$\sigma_{\text{abs}}(E_a, \Delta_2) = \frac{g_A^2 \pi}{6(2J+1)} g_{ann}^2 \delta(E_a - E_r - \Delta_2) p_a |\langle J_f || \sum_{i=1}^A \frac{1}{2} \hat{\sigma}_i \hat{\tau}_0 || J_i \rangle|^2$$

Example4: Dark Sector

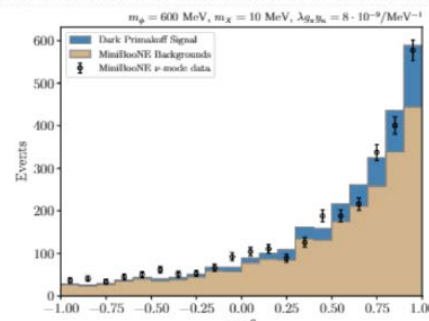
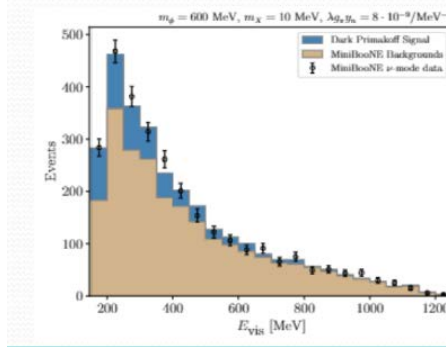
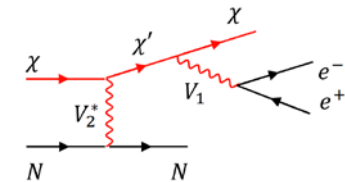
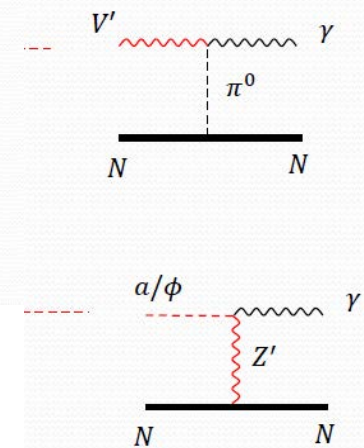
MiniBooNE anomaly: Dark sector

- For dark sector appearing from $\pi^0 \rightarrow V \gamma$ only: ruled out by MB dump

Use the charged pion 3 body
Decay (helicity unsuppressed):



Detection

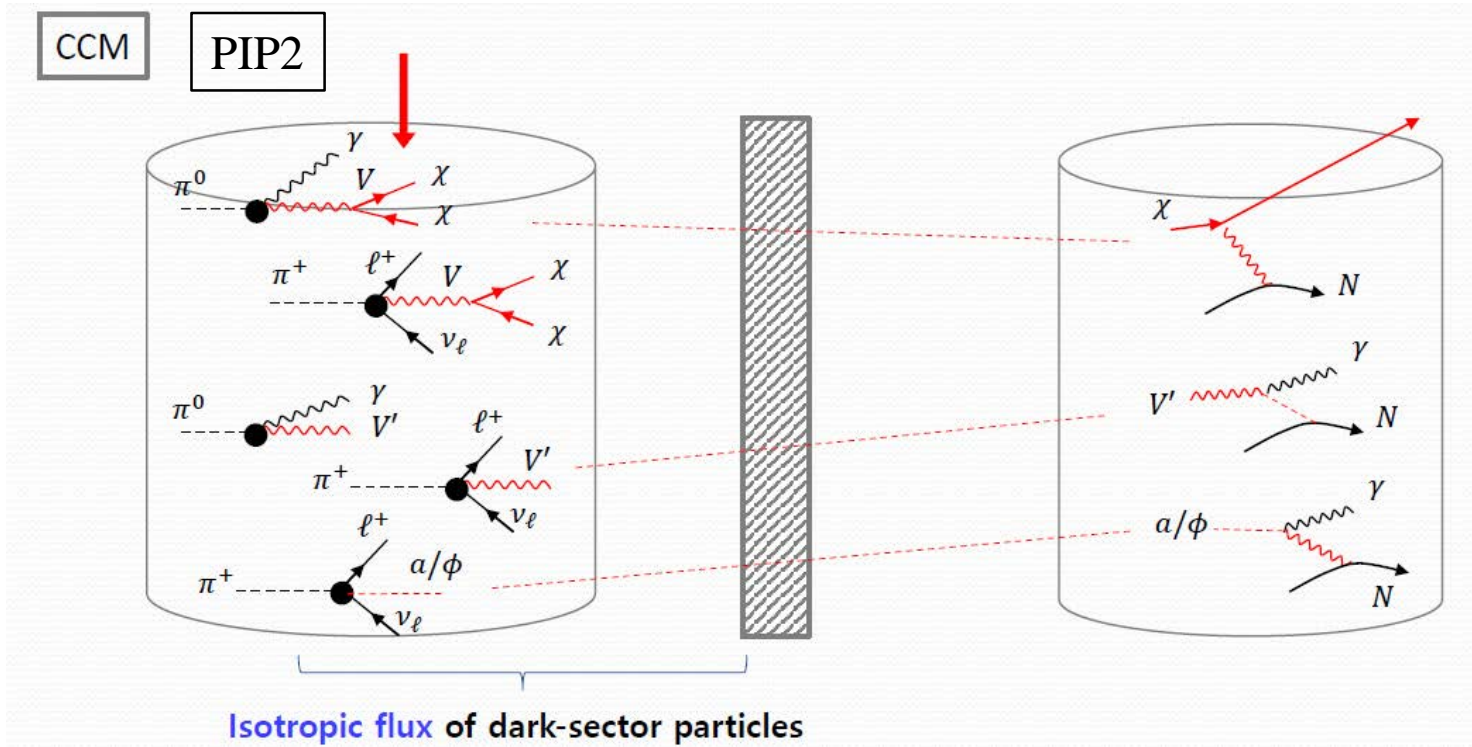


Can be probed at CCM, MicroBooNE, SBND, DUNE

Dutta, Kim, Remington, Thompson, Van de Water, PRL 129 (2022) 11, 111803

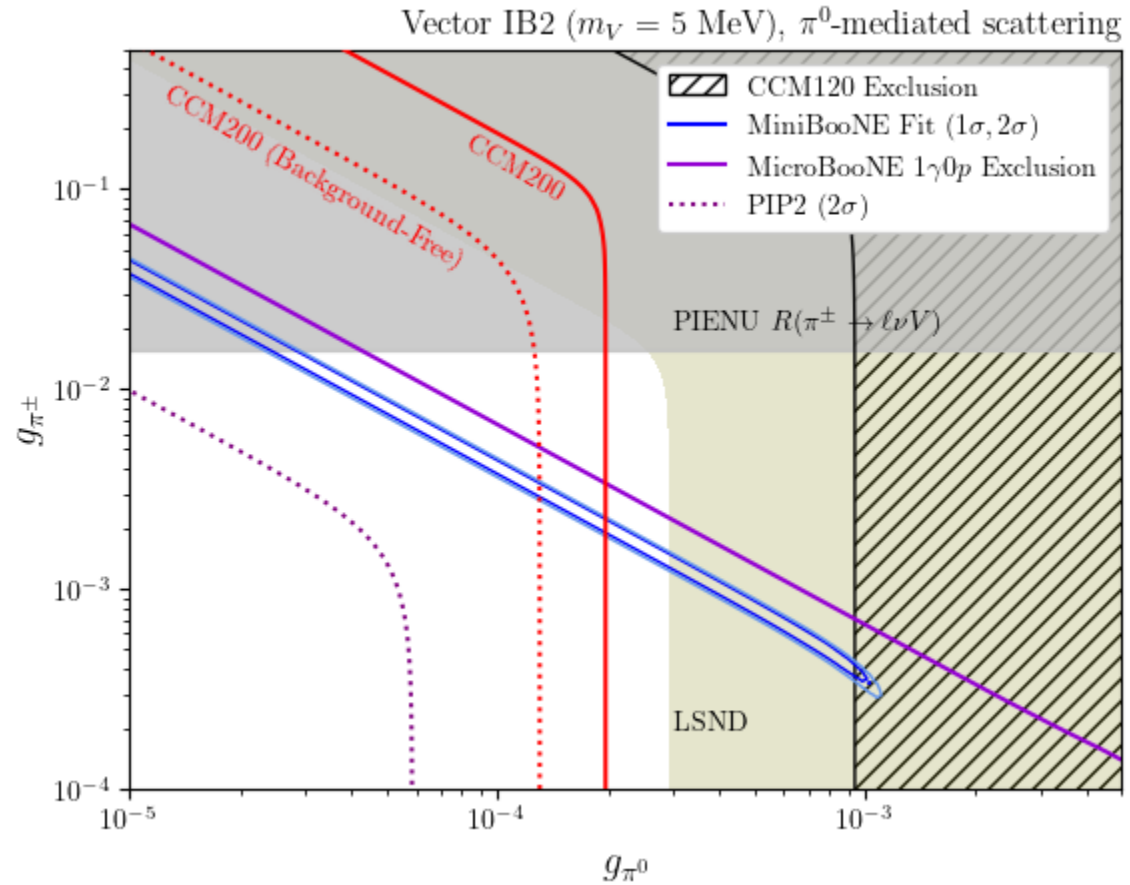
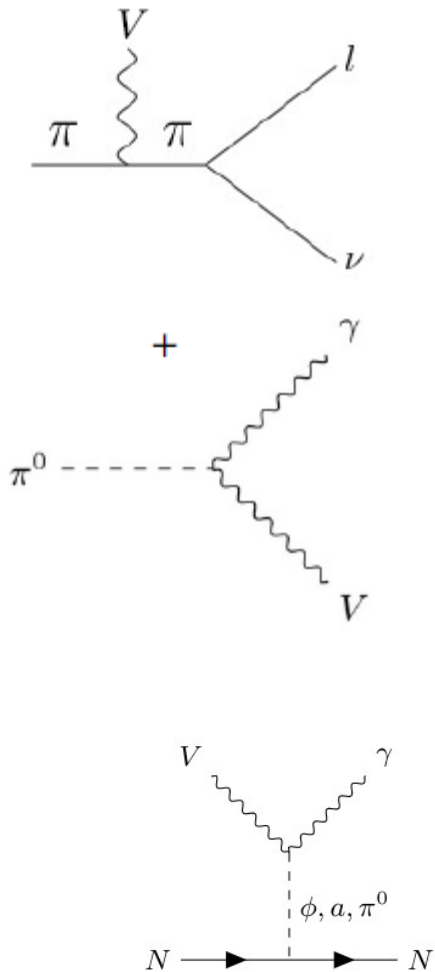
Dark Sector

MiniBooNE puzzle at CCM/PIP2



Productions via π^0 and π^+ can be investigated \rightarrow would help to discern light mediator models

Dark Sector



A. Thompson, CCM collaboration, to appear

Outlook

- Light mediator models can explain various anomalies and puzzles
- Many model possibilities
- $M(\text{new physics}) < \text{GeV}$ is not easy to probe, e.g., LHC, direct and indirect detection experiments mostly probe $M > \text{GeV}$
- Stopped pion based proton beam dump experiments have unique opportunities to probe the dark sector with MeV e, γ final states
neutrino-related background is small
- New physics can be searched using pion, γ, e^\pm , meson fluxes, Lower mass scales can be probed efficiently
- Ongoing CCM is already providing interesting results on dark sector model parameter space
- PIP2-BD can investigate a larger range of parameter space of various models