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Analytical and Numerical Studies of Dark Current in Radiofrequency Structures for Wakefield Acceleration

Gaurab Rijal¹ Xueying Lu^{1,2}, Michael Shapiro¹, John Power² ¹Northern Illinois University, Dekalb, IL, USA ²Argonne National Laboratory, Lemont, IL, USA









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OUTLINE

- Motivation
- Short Pulse Acceleration

15th International Particle Accelerator Conference, Nashville, TN JACoW Publishing ISBN: 978-3-95450-247-9 ISSN: 2673-5490 doi: 10.18429/JACoW-IPAC2024-TUPR1 DARK CURRENT SIMULATIONS IN ACCELERATING STRUCTURES **OPERATING WITH SHORT RF PULSES**

G. Rijal¹, M. Shapiro¹, J. Power², and X. Lu^{1,2*}

¹Northern Illinois University, DeKalb, IL, USA ²Argonne National Laboratory, Lemont, IL, USA

- Short pulse \rightarrow High gradient \rightarrow Compact Accelerators
- Understanding RF breakdown phenomena with short RF pulses
- Dark Current Modeling
 - Two breakdown processes under investigation:
 - Field emission
 - Multipacting
 - Two approaches: Analytical modeling + PIC simulations

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Conclusions

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C.S. Department of Energy laboratory







MOTIVATION

 Recent experimental demonstration of high accelerating gradient in RF structures powered by short RF pulses at AWA

 Close to 400 MV/m achieved on the surface of an X-band photocathode with ~9 ns RF pulses

 $_{\odot}$ RF breakdown with nanosecond-long pulses has not been carefully characterized $_{\odot}$ Previous studies mostly cover a range of ~100 ns to ~1 μs

o Goals:

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- Understanding of RF breakdown physics on various time scales
- Experimental strategies to mitigate breakdown using short pulses









SHORT-PULSE ACCELERATION

Evidence of the benefits of short-pulse acceleration from:





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DARK CURRENT STUDIES

Dark current

Unwanted current in accelerators that limits the gradient

Consequences

- Energy loss and beam instabilities
- Emittance growth
- Secondary radiation

Sources of dark current

- Field emission
- Secondary electron emission
- Multipacting









THEORY OF FIELD EMISSION

- Field emission of electrons
 - Fowler-Nordheim field emission model
 - Solution of Schrödinger equation by considering an electron in conduction band in presence of triangular potential.





MULTIPACTING

- Multipacting resonance
 - Synchronization between electron motion and RF field
 - \rightarrow Exponential growth of dark current
- Two conditions for electron multipactor
 - Resonant electron motion
 - Average secondary electron yield (SEY) > 1
- Secondary electron emission
 - SEE yield depends on incident electron energy
 - Multipacting often plays a larger role at lower gradients













ANALYTICAL MODELING OF MULTIPACTOR



ANALYTICAL RESULTS OF MULTIPACTOR



DARK CURRENT SIMULATIONS

- Short-Pulse photocathode
 - X band, 1.5 cell RF gun
 - Operating in π mode, 11.7 GHz
 - Strongly over-coupled



Normalized longitudinal electric field distribution in the X-band photocathode cavities at 11.7 GHz













Surface I









DEPENDENCE ON PULSE SHAPE



The number of multipactor electrons with surface I sidewall with 3-3-3 ns pulse





The number of multipactor electrons with surface I sidewall with 8-2-2 ns pulse





FIELD EMISSION SIMULATIONS





Field emitted electrons from the iris

• Dark electrons blocked on the path

Field emitted electrons from photocathode surface At Field 120 MV/m and beta 5





CONCLUSIONS

- Short-pulse acceleration is promising
 - An 11.7 GHz traveling wave gun powered by short pulses recently demonstrated at AWA with close to 400 MV/m on the cathode
- Modeling and experimental characterization of RF breakdown phenomena with short nanosecond RF pulses
 - Multipacting simulations
 - Comparisons of analytical and numerical dark current results
 - Multipactor current growth depends on the pulse shape
 - Field emission simulations
- Future plans: Study on plasma related processes in the cavity, improved diagnostics





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GAC²⁴





full cell sidewall





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ANALYTICAL MODELING OF MULTIPACTOR

The motion equations are:

$$egin{aligned} &mrac{dec v}{dt} = -eec E - e\mu_0ec v imes ec H \ &ec V imes ec H = -ec V_z H_ heta + ec Z V_r H_ heta \ &mrac{dv_z}{dt} = -eE_z - e\mu_0 v_r H_ heta \ &mrac{dv_r}{dt} = -eE_r + e\mu_0 v_z H_ heta \end{aligned}$$



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$$\sigma = \sigma_m \left[w \exp(1 - w) \right]^k$$

where
$$w = \frac{W_i}{W_m}$$
, $k = \begin{cases} 0.56 & \text{if } w < 1\\ 0.25 & \text{if } w > 1 \end{cases}$
 $W_m = W_{\max} \left(1 + \frac{\theta_i^2}{2\pi}\right)$, $W_{\max} = 200 \text{ eV}$
 $\sigma_m = \sigma_{\max} \left(1 + \frac{\theta_i^2}{2\pi}\right)$, $\sigma_{\max} = 2.0$

 ✓ Vaughan's secondary emission model is implemented to track the secondary emitted particle and SEY.





BACKUP SLIDES







