

A fully plasma based electron injector for a linear collider or XFEL

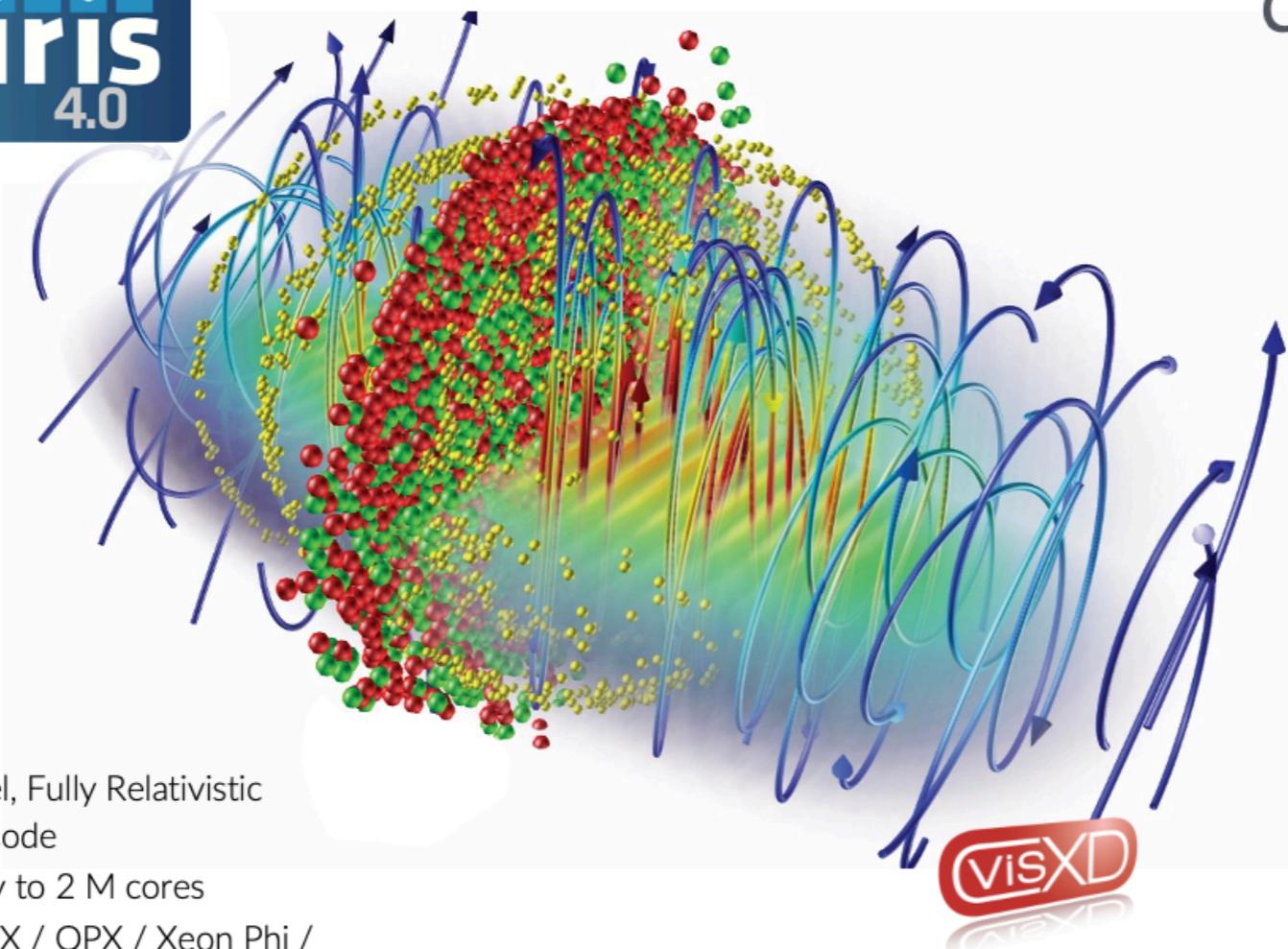
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Advanced Accelerators Concepts
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OSIRIS framework

- Massively Parallel, Fully Relativistic Particle-in-Cell Code
- Parallel scalability to 2 M cores
- Explicit SSE / AVX / QPX / Xeon Phi / CUDA support
- Extended physics/simulation models

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Committed to open science

Open-access model

- 40+ research groups worldwide are using OSIRIS
- 300+ publications in leading scientific journals
- Large developer and user community
- Detailed documentation and sample inputs files available

Using OSIRIS 4.0

- The code can be used freely by research institutions after signing an MoU
- Find out more at:
<http://epp.tecnico.ulisboa.pt/osiris>

Open-source version

- Search [osiris-code/osiris](https://github.com/osiris-code/osiris) on GitHub



Beam generation & optimal acceleration

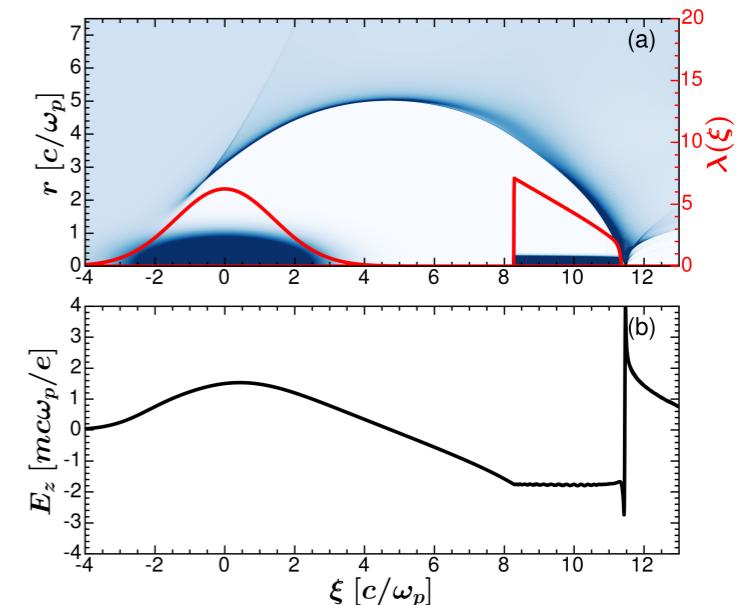
- PBA can enable high gradient acceleration (> 1 GV/cm) and high quality beam generation for next-gen XFEL and LCs.

$$Q \sim 1 \text{ nC}, \quad \epsilon_n \sim 100 \text{ nm}$$

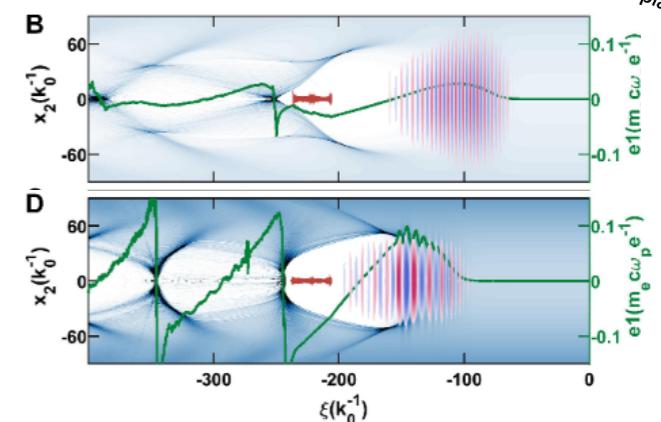
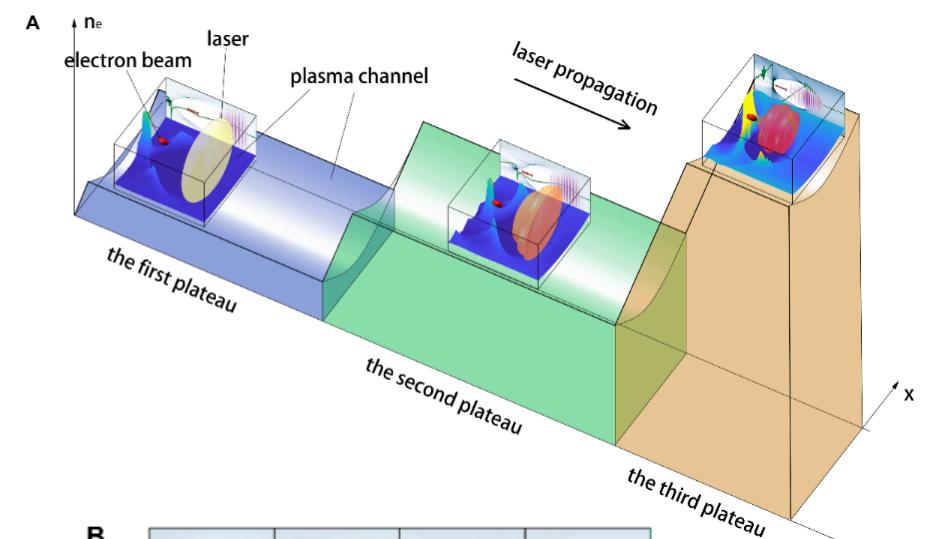
- Main challenge: low energy spread acceleration.
- Ideal loading difficult to realize with injection.
- Dynamic beam loading (DBL) schemes (e.g., tri-plateau plasma) aim to flatten average field

$$\langle d_\xi E_z \rangle_{acc} \approx 0$$

$$\xi = z - ct$$



T. N. Dalichaouch et al., Phys. Plasmas, 28, 063103 (2021).



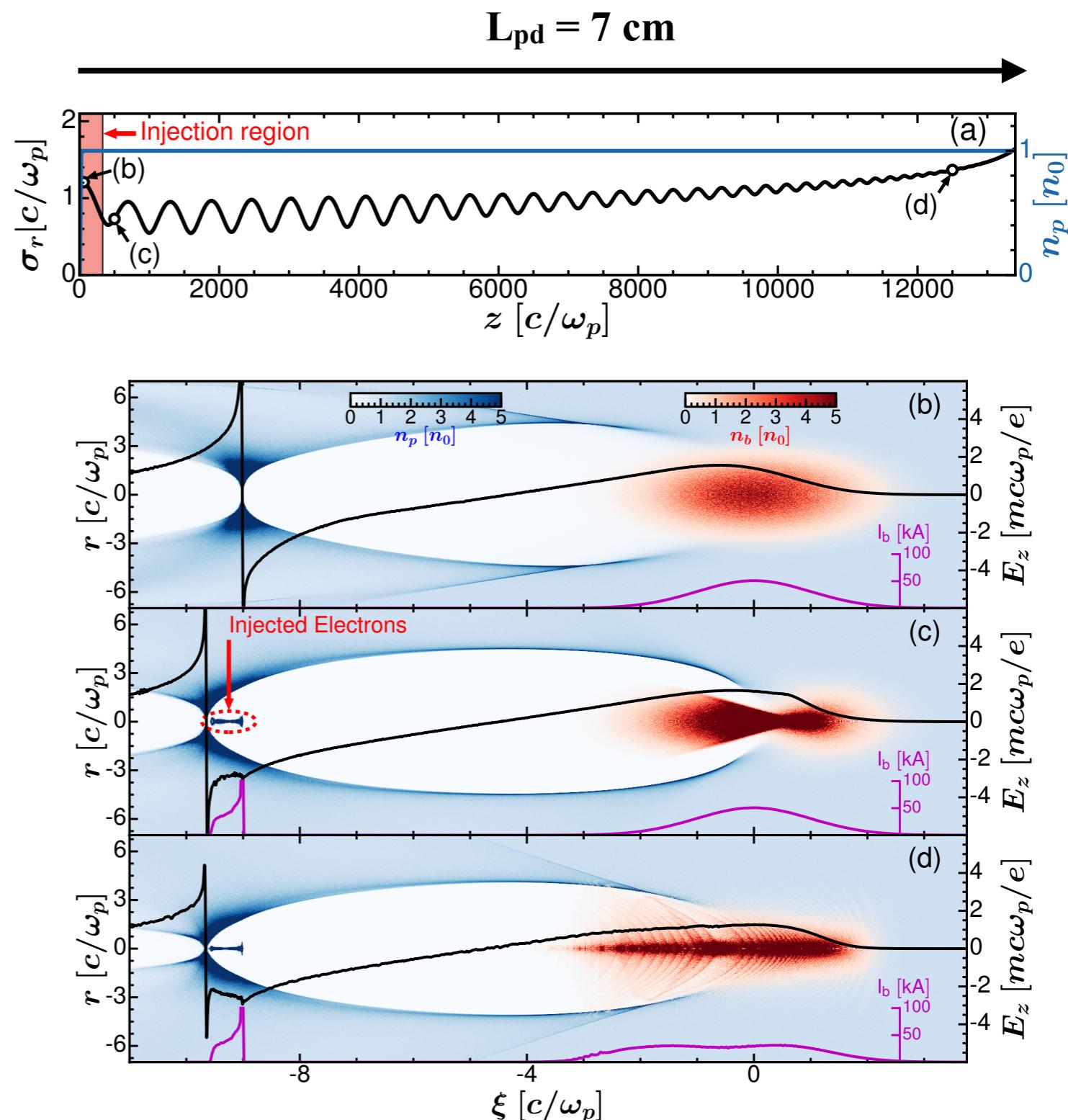
S. Liu, et al., Research, 0, 34133 (2024).

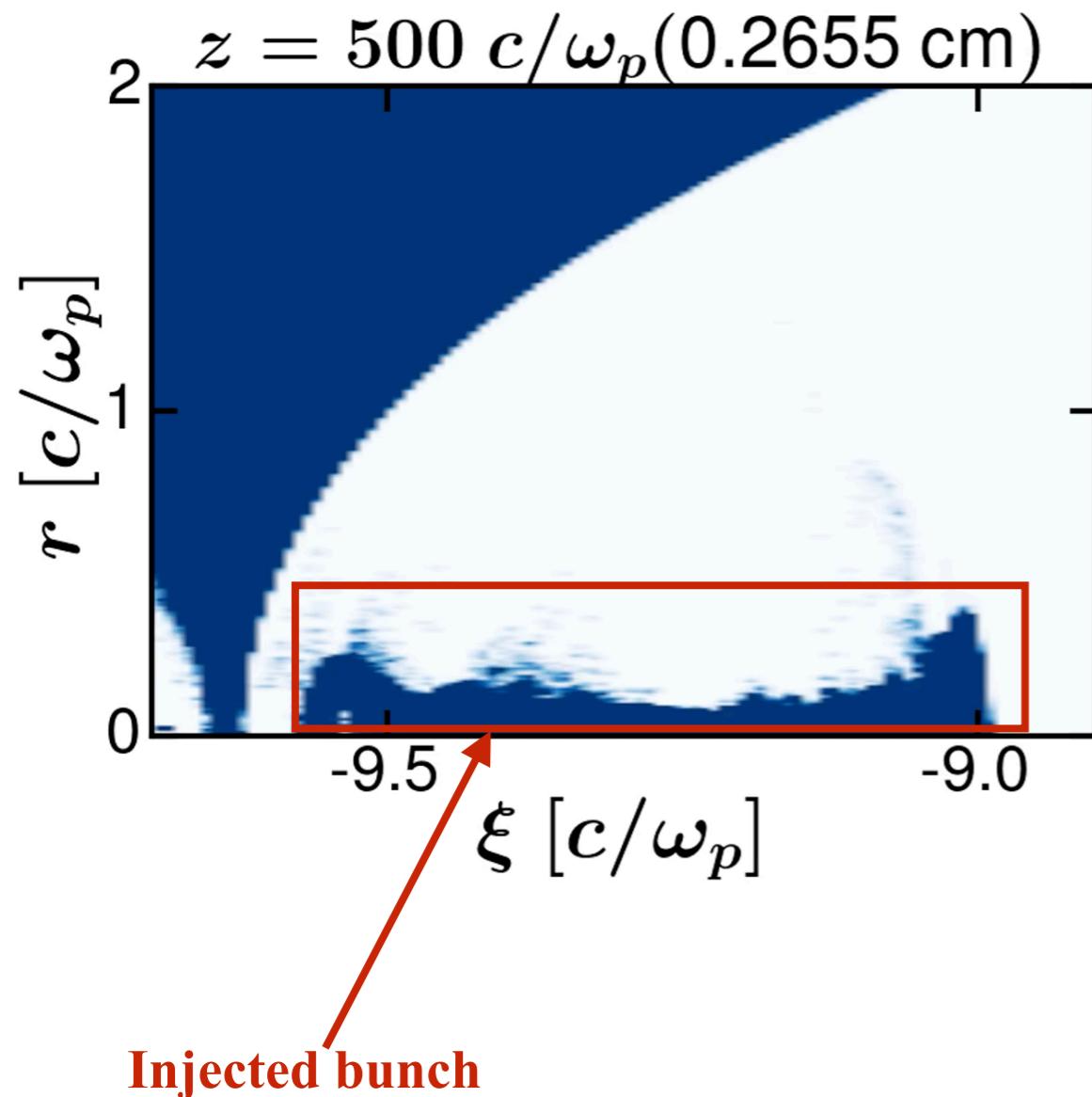
Proposed Scheme

- An unmatched e- driver (2.3 nC) self-focuses leading to injection (0.5 nC).
- DBL is then induced by wake evolution.
- Transitions from underloaded ($d_\xi E_z > 0$) to overloaded $d_\xi E_z < 0$
- DBL is dictated by pump depletion.

Drive/Plasma Parameters

$$I_{pk} \simeq 51 \text{ kA}, \quad \sigma_r \simeq 6.5 \text{ } \mu\text{m}, \quad \sigma_z \simeq 5.3 \text{ } \mu\text{m}, \\ \gamma_b = 10\text{GeV}, \quad \beta^* = 1.7 \text{ cm}, \quad n_p \simeq 10^{18} \text{ cm}^{-3}$$



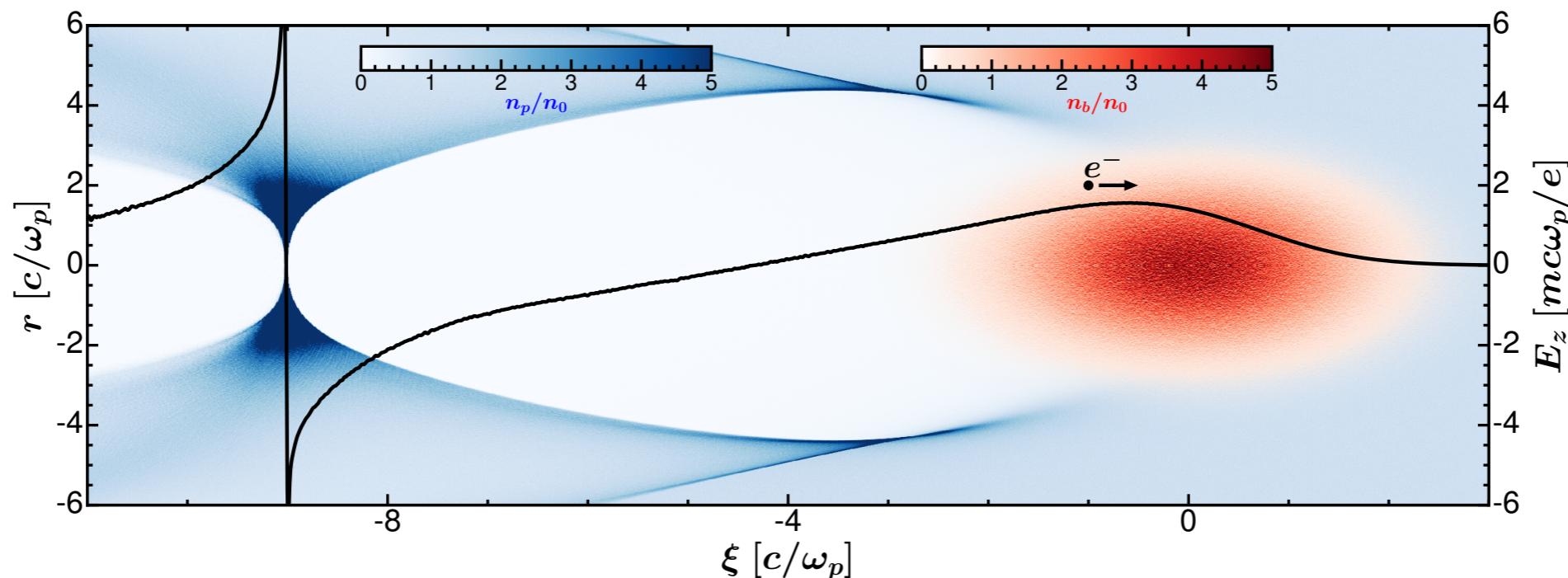


Nonlinear theory

- The loading of the wake depends on the bubble radius $r_b(\xi)$ and beam current $\lambda(\xi)$

$$\frac{dE_z}{d\xi} = \frac{1}{2} + \frac{1}{2} \left(\frac{dr_b}{d\xi} \right)^2 - \frac{\lambda(\xi)}{r_b^2}$$

- DBL induced by wake evolution: r_b decreases, space-charge force increases.
- $d_\xi E_z$ changes sign from positive to negative.



Transverse dynamics

$$x''(z) + \frac{\gamma'_b(z)}{\gamma_b} x'(z) + k_\beta(z)^2 x(z) = 0$$

$$x(z) \approx x_i \left[\frac{\gamma_{bi}}{\gamma_b(z)} \right]^{1/4} \cos(\phi)$$

Longitudinal dynamics

$$\frac{d\xi}{cdt} = \frac{v_z}{c} - 1 \approx -\frac{|x'|^2}{2}$$

$$\Delta\xi(z) \approx \frac{-k_p^2 r_i^2 z}{4\gamma_b(1 + \sqrt{\gamma_{bi}/\gamma_b})}$$

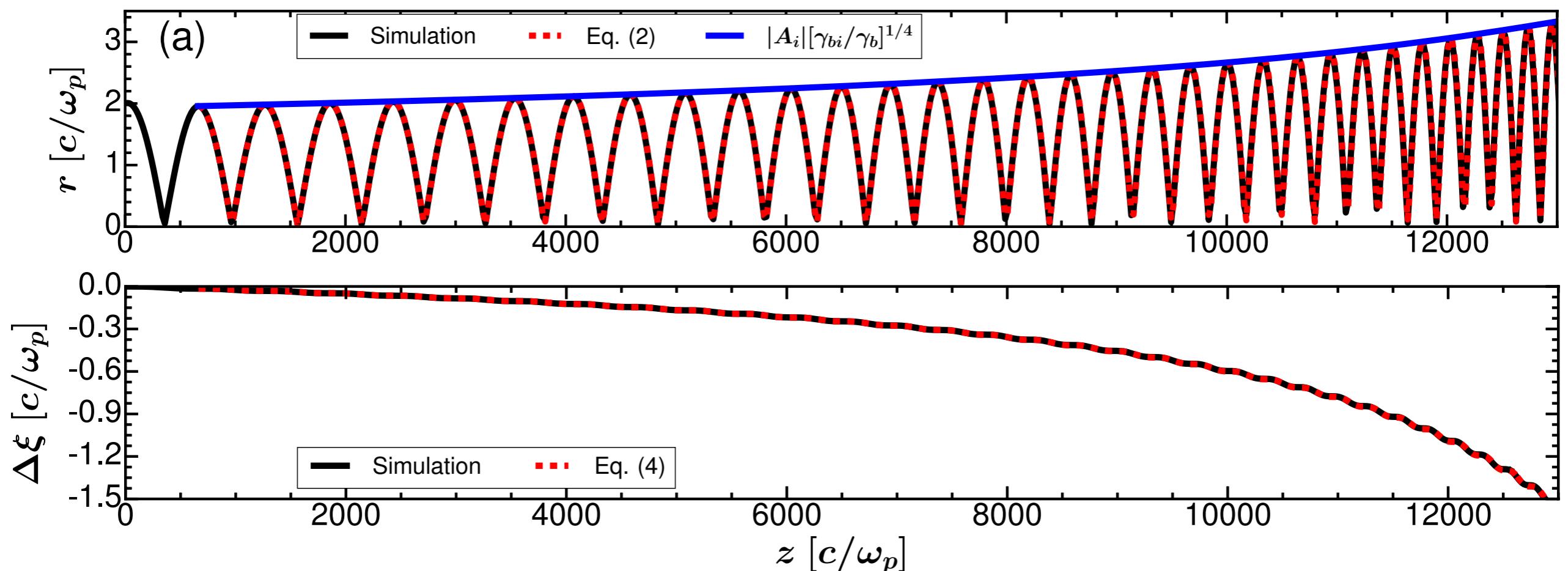
$$k_\beta(z) = \frac{k_p}{\sqrt{2\gamma_b(z)}} \quad \phi = \int_0^z k_\beta ds$$

**Trans/long dynamics
coupled to energy**

Example drive beam electron motion

$$x(z) \approx x_i \left[\frac{\gamma_{bi}}{\gamma_b(z)} \right]^{1/4} \cos(\phi) \quad (2)$$

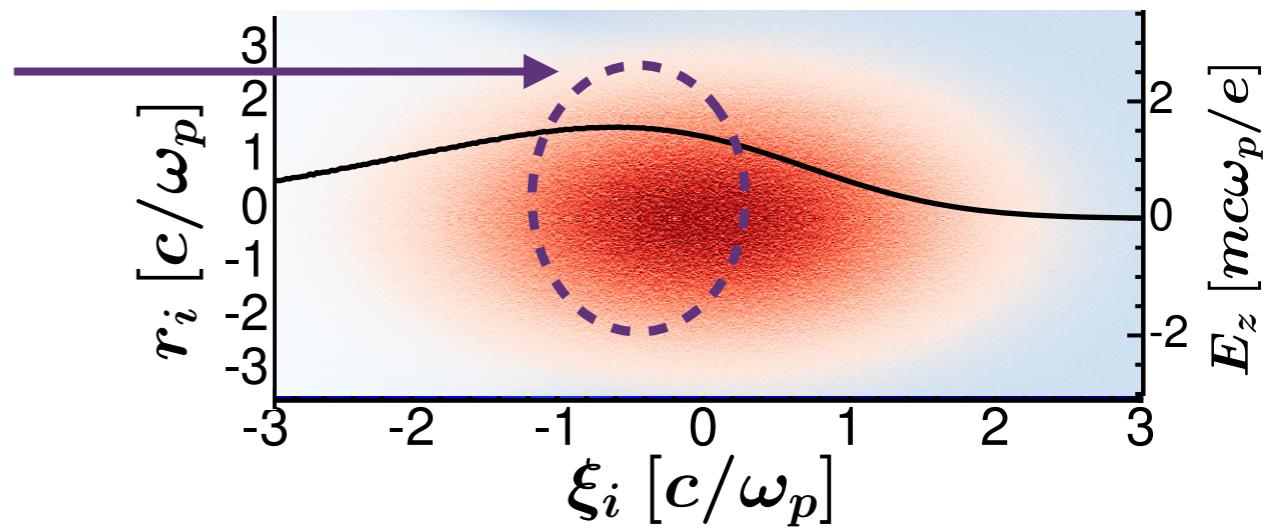
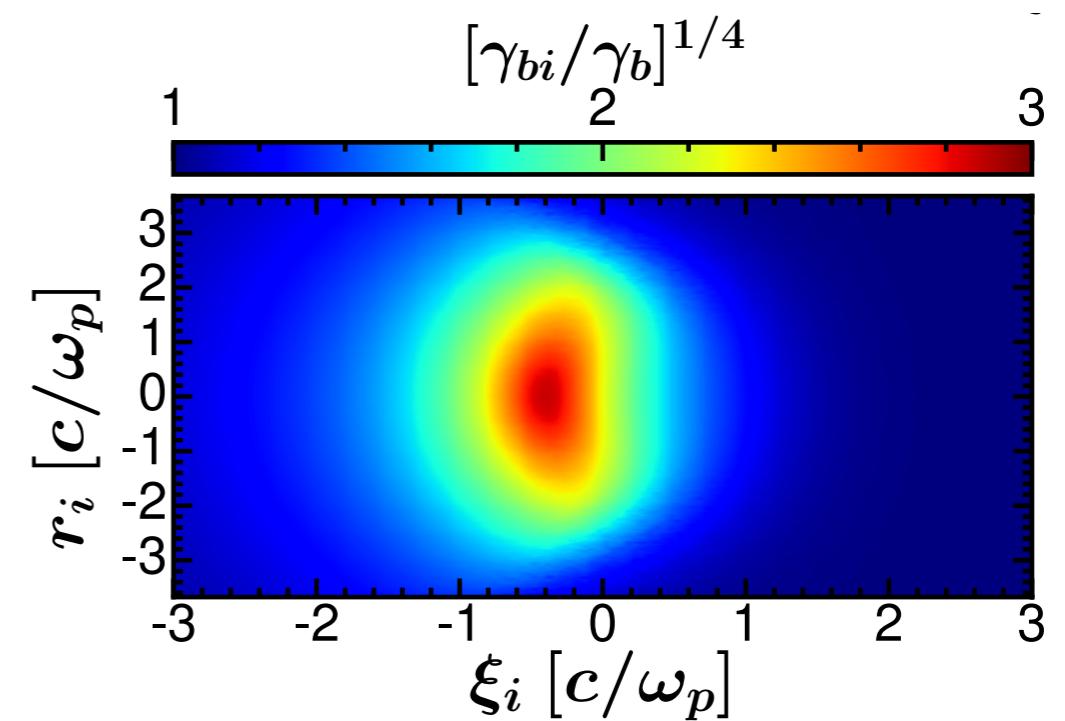
$$\frac{d\xi}{cdt} \approx -\frac{|x'|^2}{2} \quad (4)$$



- Betatron oscillations and pump depletion lead to defocusing+dephasing.

Which beam components drive defocusing?

- Energy factor $x(z) \sim [\gamma_{bi}/\gamma_b(z)]^{1/4}$ for drive beam electrons shown after **6.9 cm (13000 c/wp)**
- Spot size expansion driven by significant pump depletion (93%+ energy loss) behind the centroid
 $(-1 \lesssim k_p \xi_i \lesssim 0)$
- Strong defocusing where decelerating field is large.

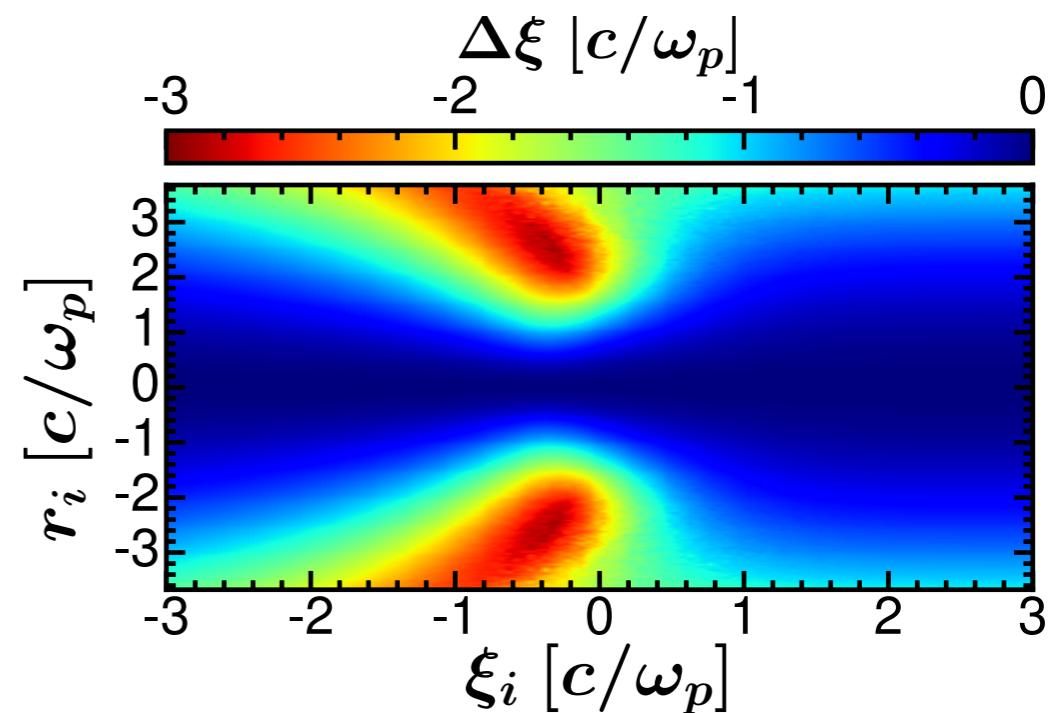


Which beam components drive dephasing?

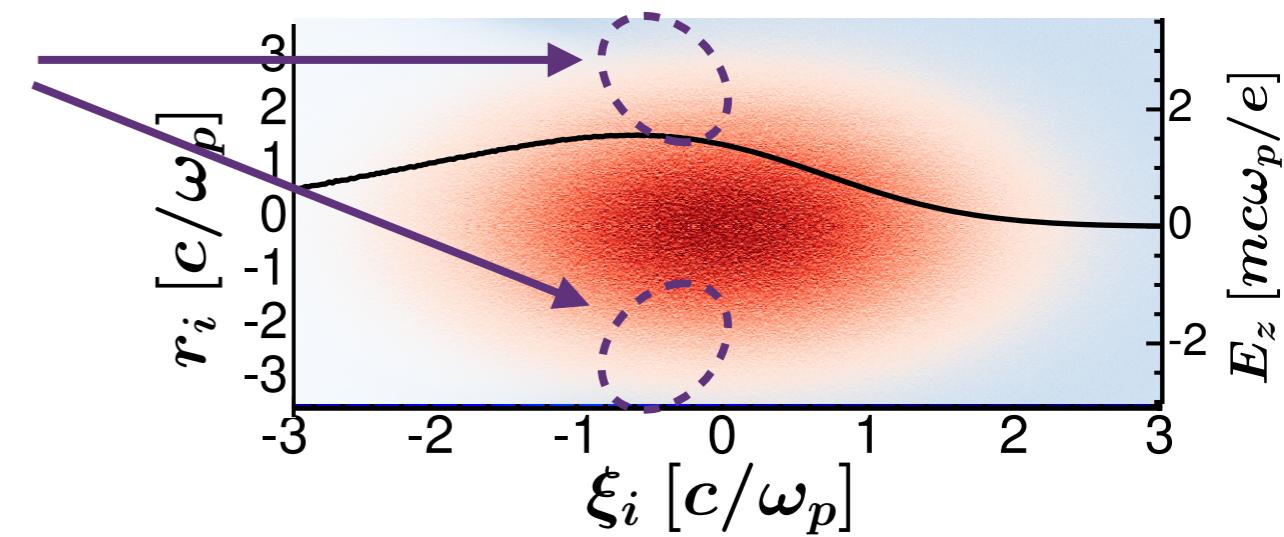
- Dephasing $\Delta\xi$ shown for driver beam electrons after $z = 6.9$ cm (13000 c/wp) in plasma.

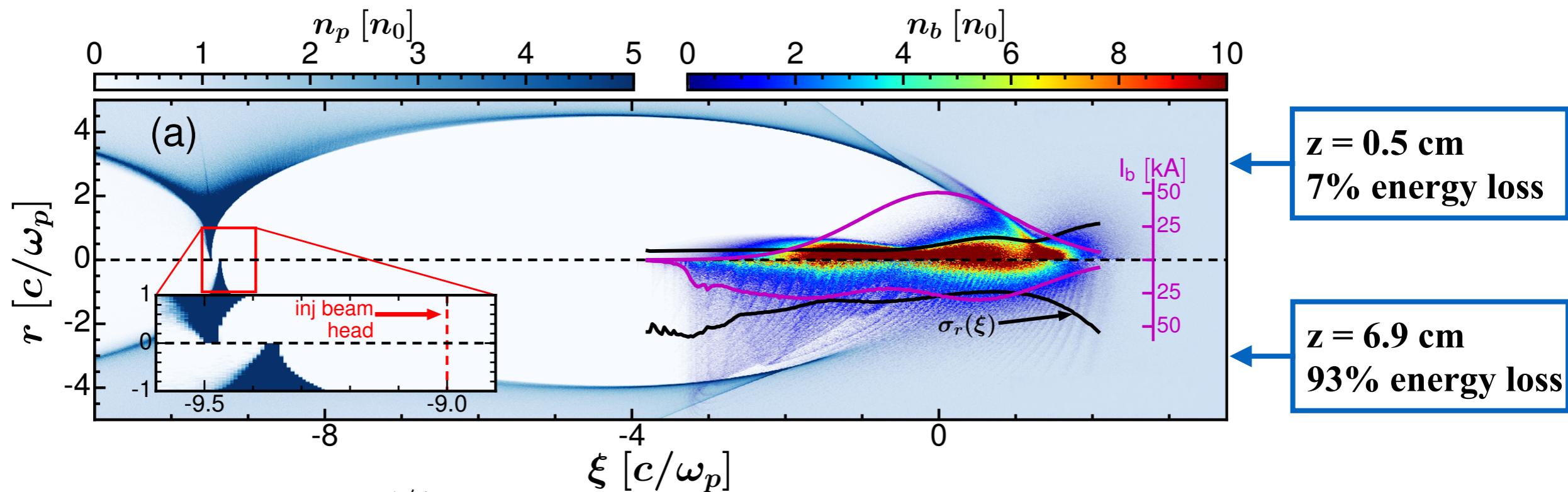
$$\Delta\xi(z) \approx \frac{-k_p^2 r_i^2 z}{4\gamma_b(1 + \sqrt{\gamma_{bi}/\gamma_b})}$$

$$\frac{d\xi}{cdt} \approx -\frac{|x'|^2}{2}$$



- Dephasing is most prominent for particles with high frequency, large betatron oscillations.

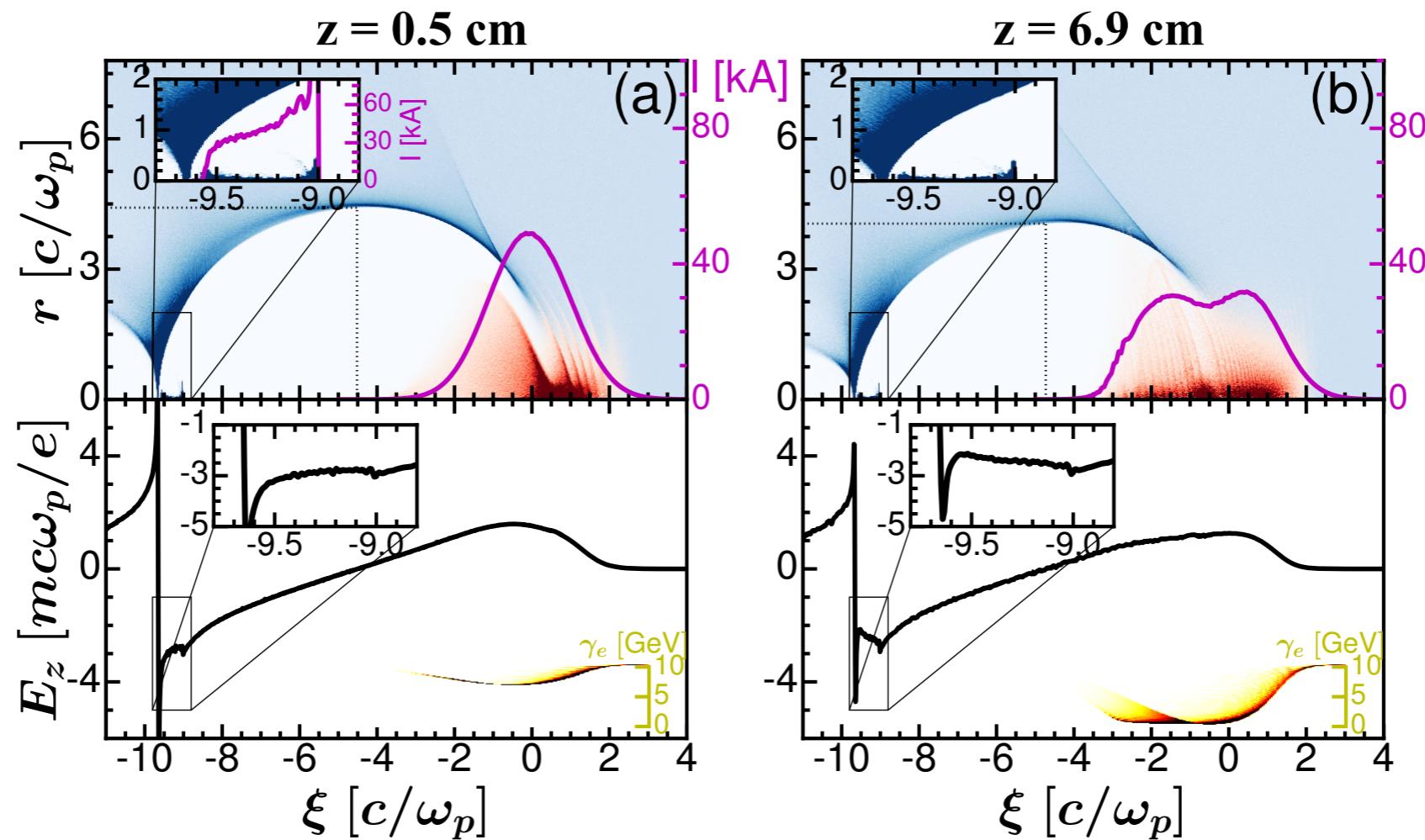




Wake evolution w/out trailing beam

- Spot size defocusing reduces the wake length.
- Dephasing reduces peak current and blowout radius $k_p r_m \approx 2\sqrt{\Lambda}$
- Smaller $r_b(\xi)$ along injected beam leads to stronger loading.

Loaded wake evolution over pump depletion



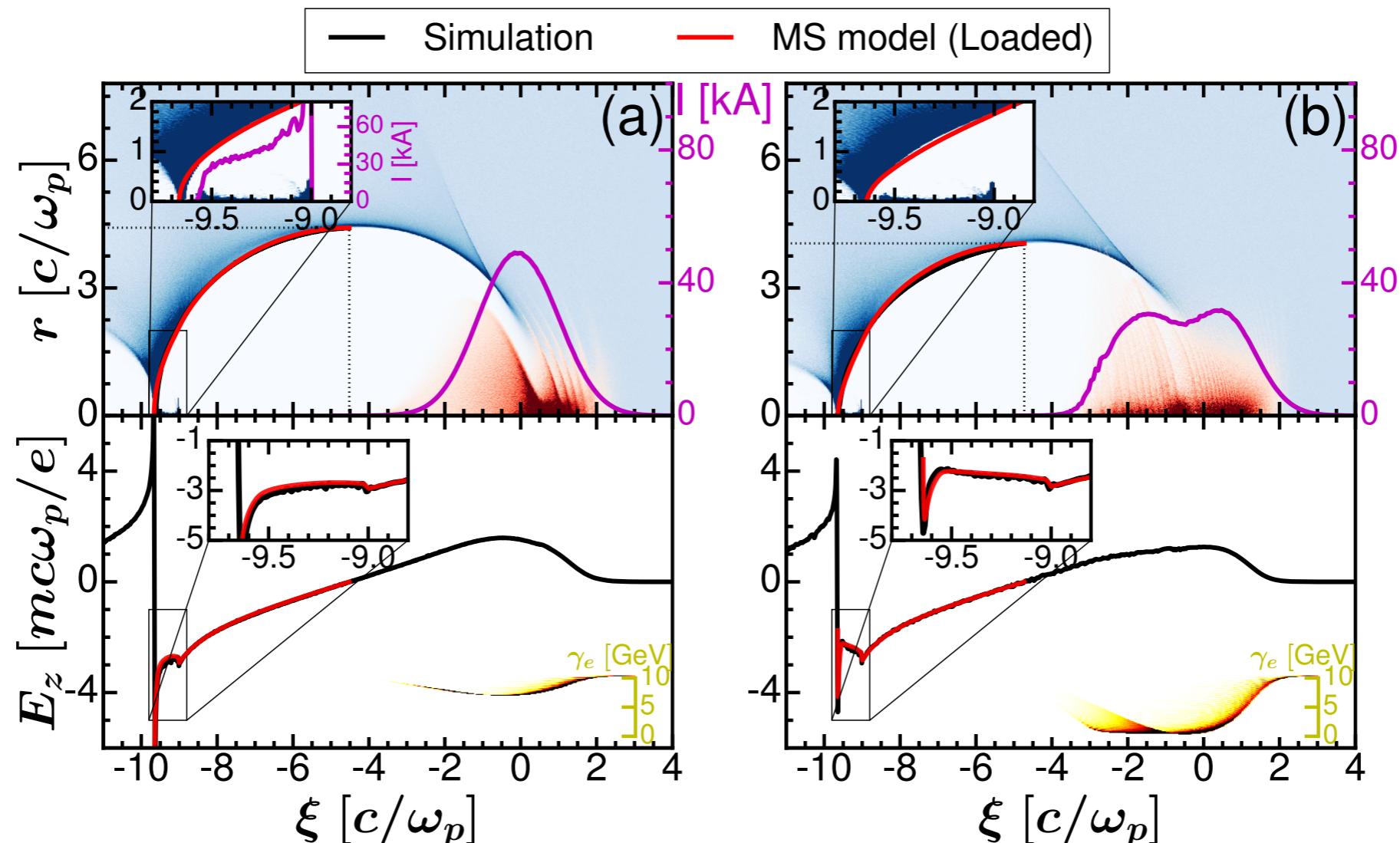
- Loaded wake remains fully expanded due to the space-charge force.
- Beam force $\lambda(\xi)/r_b$ increases ($r_b \downarrow$) and slope $dr_b/d\xi$ decreases ($|p_\perp| \downarrow, \psi \uparrow$).

$$\frac{dE_z}{d\xi} = \frac{1}{2} + \frac{1}{2} \left(\frac{dr_b}{d\xi} \right)^2 - \frac{\lambda(\xi)}{r_b^2}$$

$$\frac{dr_b}{d\xi} = \frac{-p_\perp}{1 + \psi}$$

- Beam term dominates so chirp becomes negative (overloaded).

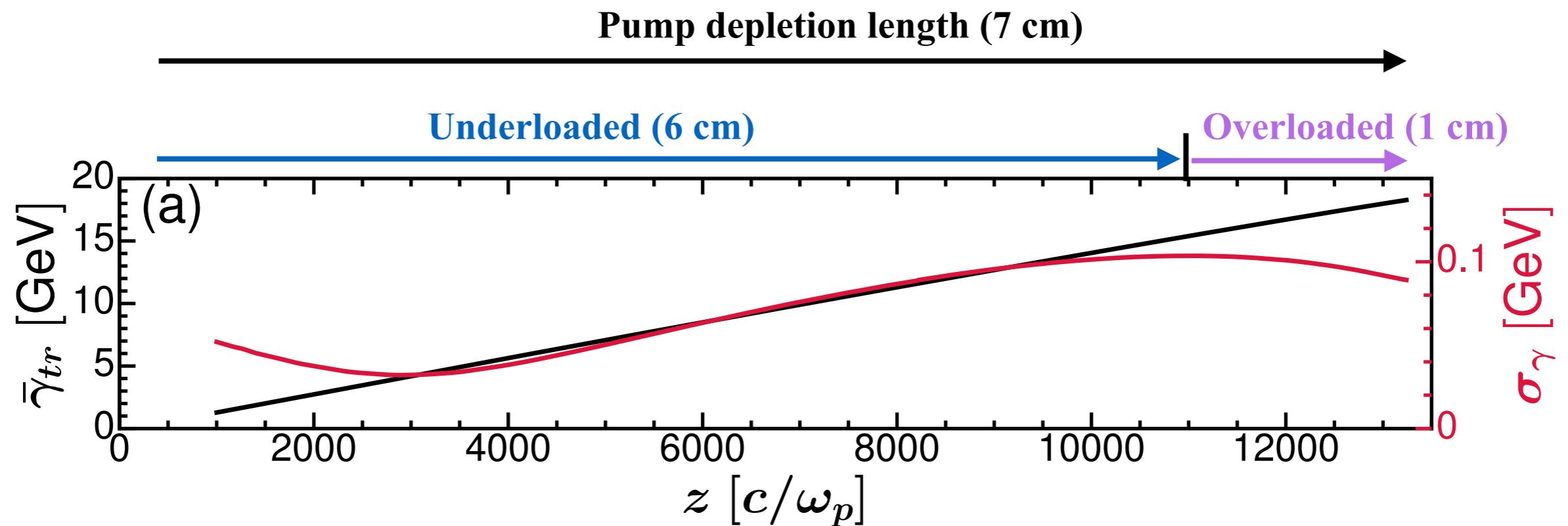
DBL effect verified by multi-sheath model



Multi-sheath model

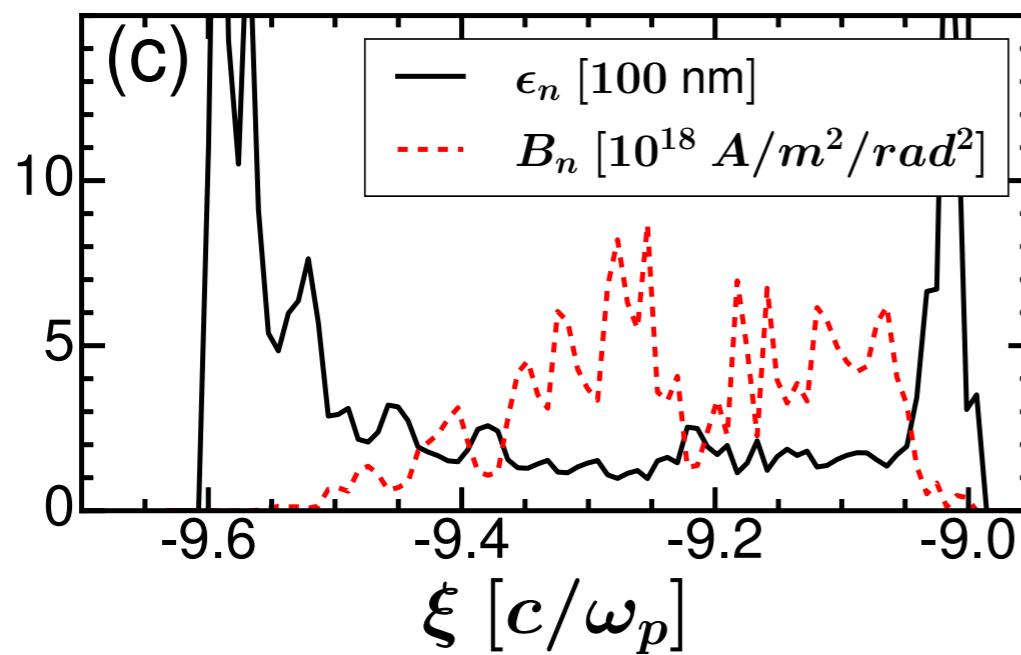
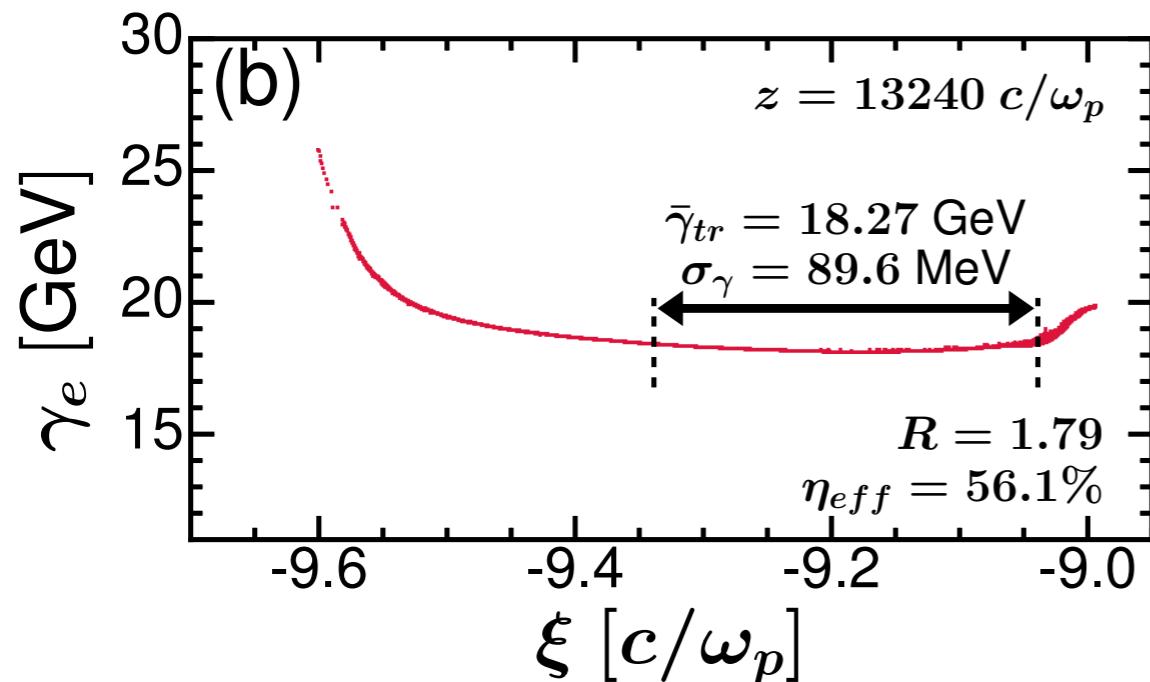
Verified this DBL effect using multi-sheath model for wake potential. Bubble radius and electric field obtained from 2nd order pde

$$\psi_0(\xi) = (1 + \beta')r_b^2/4, \quad E_z(\xi) = d\psi_0/d\xi$$



- Wake is underloaded until final 10-20% of pump depletion distance
- Energy spread deviates (sub 0.1 GeV) due to strong overloading.
- Normalized energy spread $\hat{\sigma}_\gamma = \sigma_\gamma / \langle \gamma \rangle \sim \langle d_\xi E_z \rangle$ is minimized

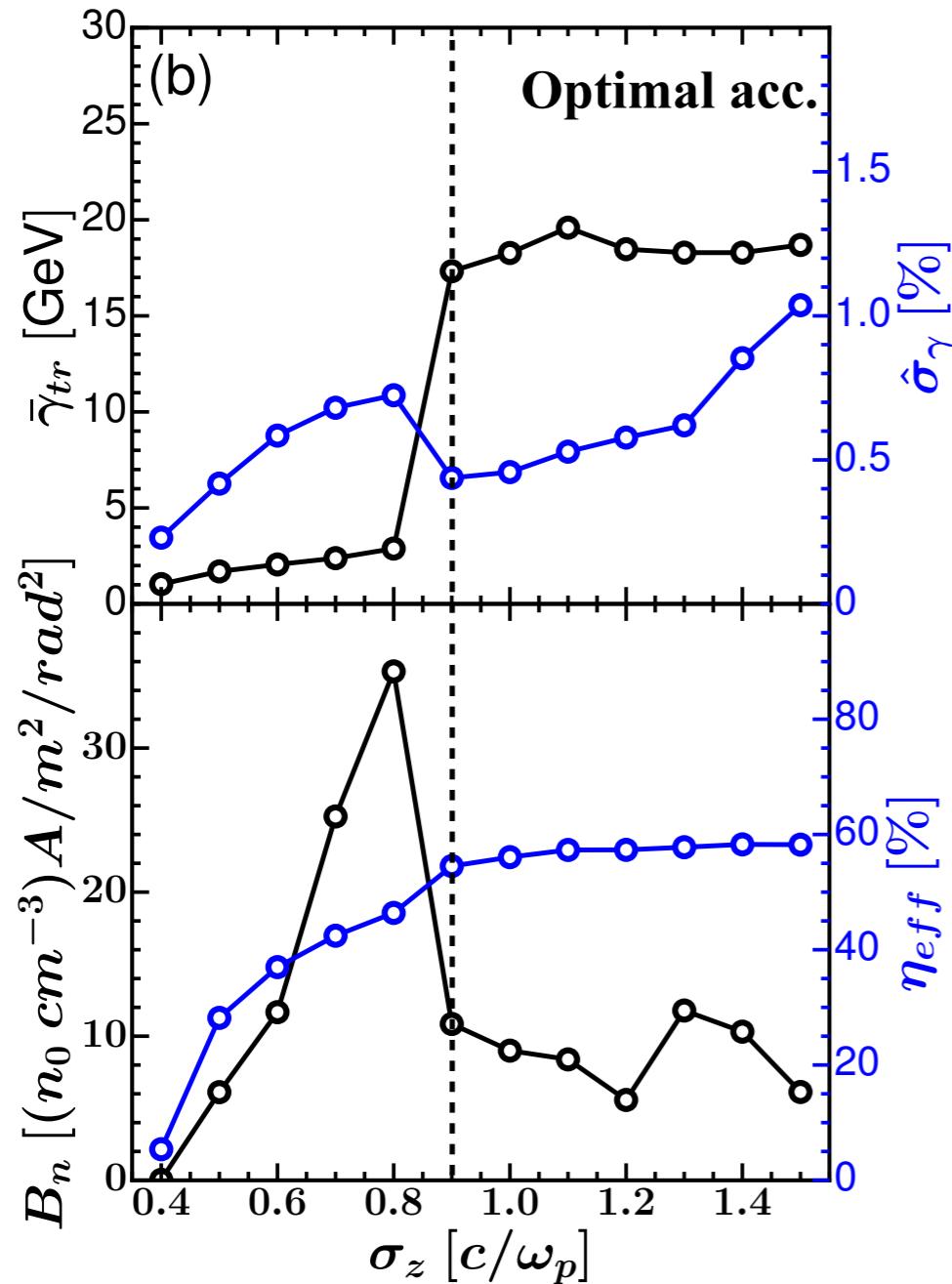
Energy and final beam slice parameters (7 cm)



- Beam extracted with **~18.3 GeV** and **~90 MeV (0.49%)** after 7 cm propagation distance.
- High efficiency acceleration **56%** and high transformer ratio **~1.8**.

Energy/Brightness Booster		
Drive	Trailing	
Energy	10 GeV	18.3 GeV
Emittance	$\sim 50 \mu\text{m}$	$\sim 100 \text{ nm}$
B_n [$\text{A/m}^2/\text{rad}^2$]	10^{14}	10^{19}

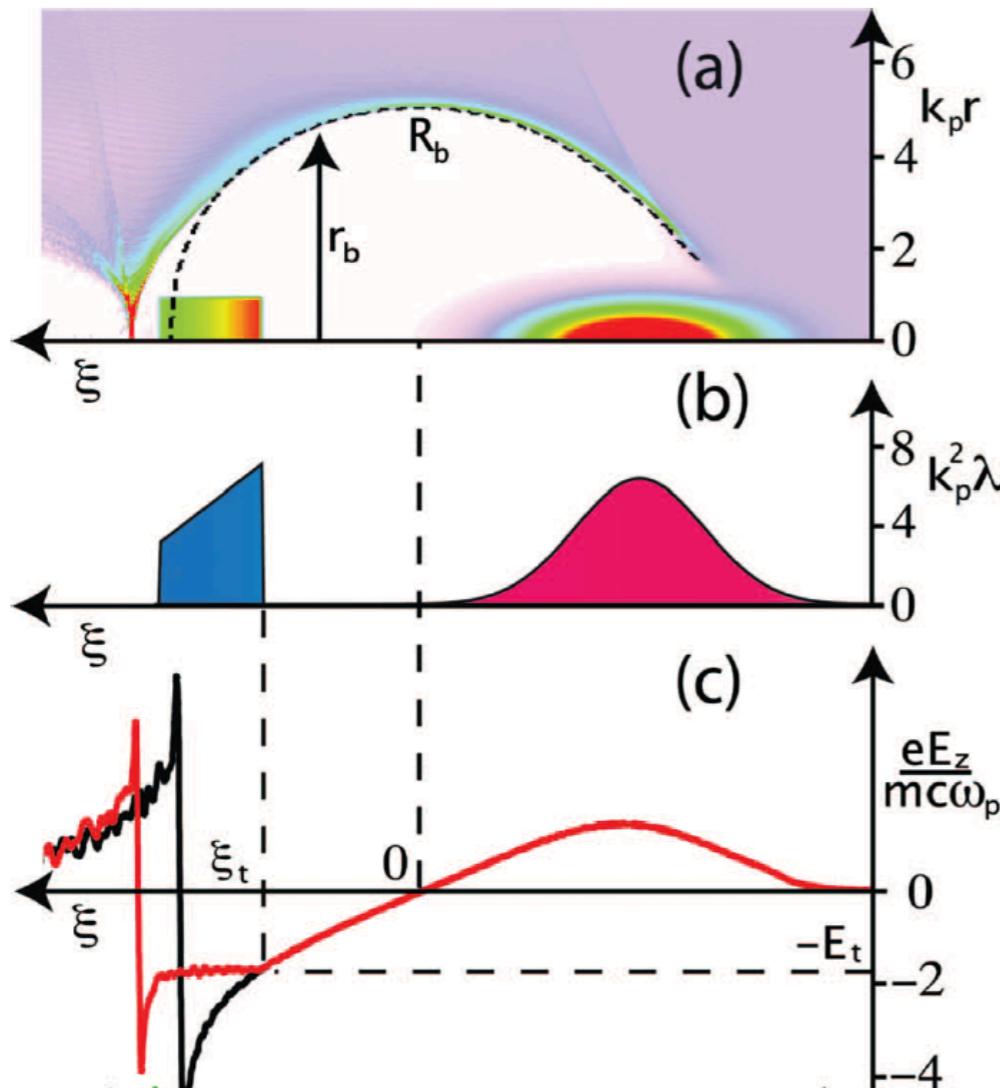
Longer drivers that generate optimal DBL



- Parameter scans show optimal DBL can be achieved by using longer drivers $\sigma_z \gtrsim 0.9 c/\omega_p$

- Can produce high quality multi-GeV beams:
 - 20 GeV energies,
 - sub-1% energy spreads.
 - High efficiency (> 50%)
 - High brightness ($10^{19} \text{ A/m}^2/\text{rad}^2$)

Thank you for your time!



Tzoufras et. al, Phys. Plasmas 16, 056705 (2009).

Basic idea & Limitations

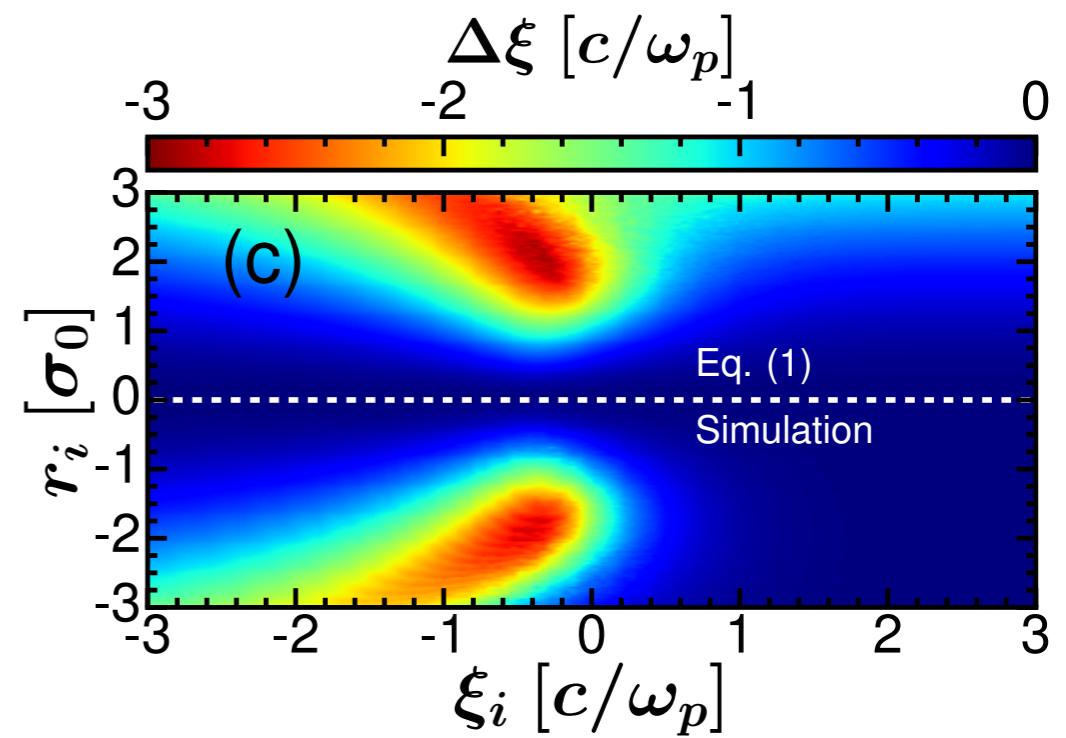
- The loading of the wake depends on the bubble radius $r_b(\xi)$ and beam current $\lambda(\xi)$

$$\frac{dE_z}{d\xi} = \frac{1}{2} + \frac{1}{2} \left(\frac{dr_b}{d\xi} \right)^2 - \frac{\lambda(\xi)}{r_b^2}$$

- Conventional picture: wake is “static” and trapezoidal beam flattens E_z .
- Does not work if wake evolves

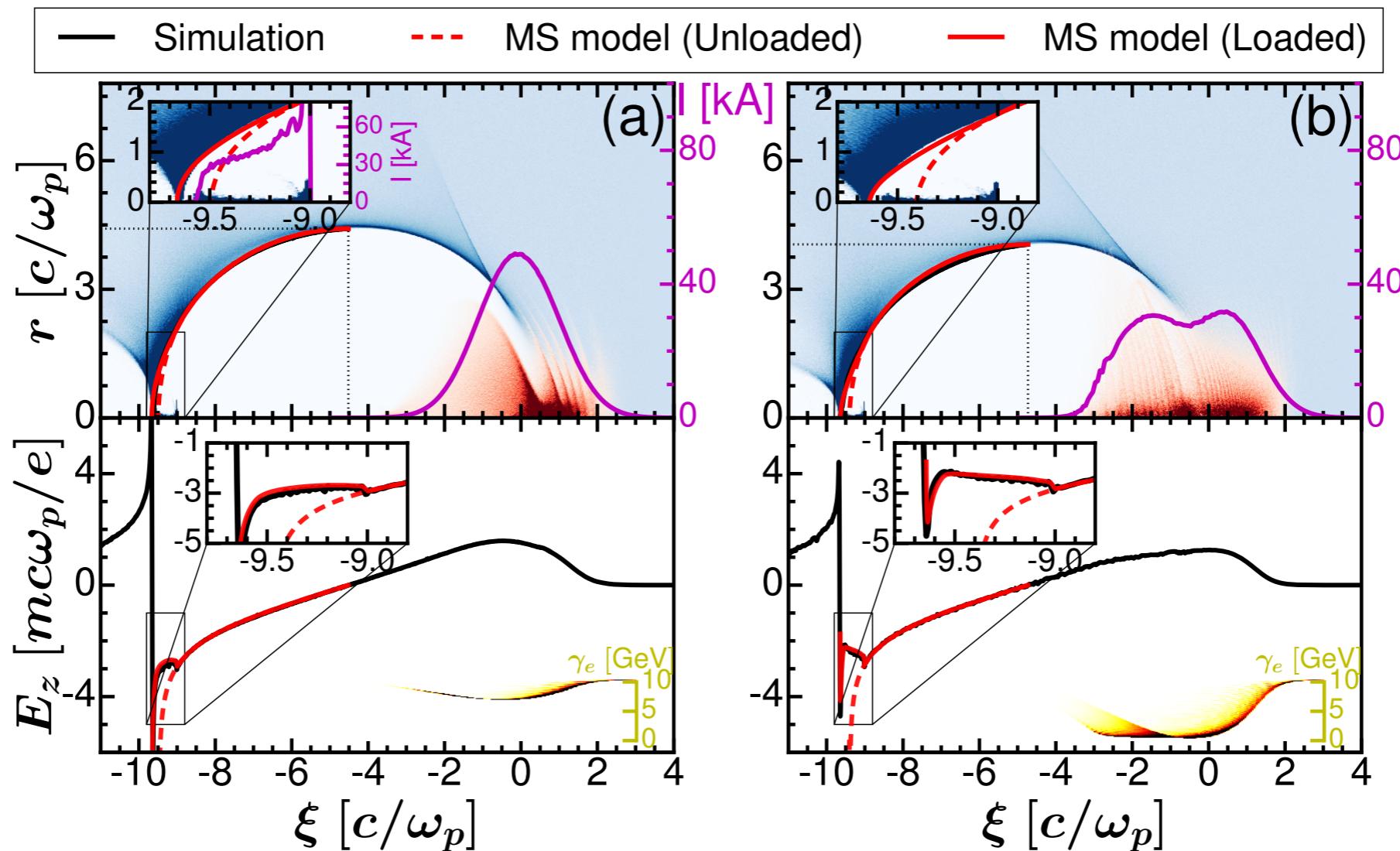
Dephasing: Simulation vs Theory

- Dephasing theory vs simulation shown for drive beam electrons after $z = 6.9$ cm (13000 c/wp) in plasma.
- Strong agreement along most of driver except at front.
- This is because Eq. (1) assumes linear focusing force which is not valid at beam head.



$$\Delta\xi(z) \approx \frac{-k_p^2 r_i^2 z}{4\gamma_b(1 + \sqrt{\gamma_{bi}/\gamma_b})} \quad (1)$$

Modeling unloaded wake



Multi-sheath model

Can calculate the bubble radius and electric field using the multi-sheath model for the wake potential

$$\psi_0(\xi) = (1 + \beta')r_b^2/4, \quad E_z(\xi) = d\psi_0/d\xi$$