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# Challenges of Plasma Wakefield Accelerator Beams as Drivers for Free-Electron Lasers

Alex H. Lumpkin Accelerator Directorate

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# **Challenging Path to Compact X-ray FELS**

- The PWFA community is targeting compact x-ray FELS, but..
- We consider the seeded free-electron laser (FEL) tests at 276 nm at HZDR in Germany and the self-amplified spontaneous emission (SASE) FEL test at 27 nm in China.
- Note: To obtain exponential gain in an X-ray SASE FEL, one will need significantly better beam quality than these cases.
- The basic slippage of e-beam behind seed laser or FEL pulse is considered for these two experiments.
- I compare the Pierce parameter at LCLS-1 at 1.5 angstroms to plasma wakefield accelerator (PWFA) case.
- Summary of rf linac-driven FEL parameters is also provided for comparison. Include ~1 angstrom cases at LCLS-1 and SACLA.

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# Slippage of e-beam Relative to Seed Laser and/or FEL

- The e-beam slips behind in z one resonant wavelength,  $\lambda_1$ , in traveling one undulator period.
- The total slippage at 276 nm for 97 periods is 26.9 μm.
  dz= N<sub>u</sub> λ<sub>1</sub> = 26.9 μm → dt=80 fs compared to original ~30-fs
  e-beam pulse where N<sub>u</sub> is total period number. Gain issue.
- The chicane is used to lengthen the 30-fs e-beam pulse to 0.9 ps and is matched to the seed laser pulse of about 1 ps. Timing/synchronization argument, but this step also makes slippage small compared to the pulse length.
- For Chinese 27-nm case, slippage is 13 fs at  $L_u$ =4.5 m, but would be more limiting at  $L_u$ =20 m and 63-fs slippage. Laser bunch length of 25 fs (FWHM) so expect similar for e-beam.

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#### **French-German test at HZDR**

- This was a seeded FEL test with a single 2-m long undulator.
- Gain detected at a red-shifted wavelength 276 nm from the seed wavelength of 270 nm. If not red shifted, hard to detect.
- Moreover, difficult to select only the FEL radiation for a user.



Fig. 1 | Experimental layout. The LPA is driven by the DRACO laser (for more details on the DRACO footprint, see ref. <sup>55</sup>). The electron beam generated in the LPA is first characterized using a removable electron spectrometer and then sent through a triplet of quadrupoles (QUAPEVAs) for beam transport to the undulator and FEL radiation generation. ICTs, integrated current transformers. Non-labelled elements: dipoles, red blocks; optical lenses, blue disks; mirrors, grey circled

black disks. **a**, Particle-in-cell simulation rendering of the accelerating structure driven by the laser pulse (red); the electron cavity sheet formed from the plasma medium (light blue) is in purple and the accelerated electron bunch in green. **b**–**d**, Electron-beam transverse distribution measured at the LPA exit (**b**), the undulator entrance (**c**) and the undulator exit (**d**).

Labat, M. et al., Seeded free-electron laser driven by a compact laser plasma accelerator. Nat. Phys. 17, 150–156 (2022).

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#### Seeded FEL Results with Red-shifted Output

- Investigators report radiation gain at a 6-nm longer wavelength than the seed laser.
- This is not user friendly since seed laser energy is larger.





intensity and displayed in logarithmic (dB) scale. In **a,b,c(iv**), the distributions are displayed in a linear scale. Simulation parameters (electron-beam parameters given at the source point):  $E_e$  = 188.8 MeV, charge = 150 pC,  $\sigma_z$  = 2 µm (r.m.s.), normalized emittance  $\epsilon_{x,y}$  = (1.5; 1.0) mm mrad, divergence  $\sigma_{x',y'}$  = (1.5; 1.0) mrad (r.m.s.),  $\sigma_e$  = 5% (r.m.s.),  $R_{56}$  = -1.8 mm, QUAPEVA 2 strength detuned by -2%,  $E_{seed}$  = 0.5 µJ,  $\lambda_{seed}$  = 269 nm,  $\Delta \lambda_{seed}$  = 3.9 nm (FWHM) and  $\Delta T_{seed}$  = 1.0 ps (FWHM).

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# **Pierce Parameter for SASE FEL**

- Dimensionless scaling parameter, p, determines the main characteristics of high gain FEL systems and enables order of magnitude estimates. (p.97 of Kim, Huang, Lindberg Book.)
- $\rho = [1/8\pi I/I_A (K(JJ)/1 + K^2/2)^2 \gamma \lambda_1^2/2\pi \sigma_x^2]^{1/3}$  where I is peak current,  $I_A$  is the Alfven Current, K is undulator parameter,  $\gamma$  is Lorentz factor,  $\lambda_1$  is the resonant fundamental wavelength,  $\sigma_x$ is the transverse e-beam size.
- Gain length ~  $\lambda_u/4\pi\rho$
- Saturation power ~ ρ x (e-beam power)
- Saturation length  $L_{sat} \sim \lambda_u / \rho$
- Transverse mode size  $\sigma_r \sim [\lambda_1 \ \lambda_u \ /16 \ \pi^2 \rho]^{1/2}$
- LCLS-1, ρ= 4 x 10<sup>-4</sup>

(Kim et al. Book.)

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• Chinese case  $\rho = 5 \times 10^{-3}$  reported.

#### **Chinese SASE FEL Results at 27 nm**

• Demonstration of gain~100 after 3rd undulator





beams from the LWFA for 20 consecutive shots. **c**, **d**, Measured transverse profiles of the electron beam at the entrance (**c**) and exit (**d**) of the undulators. The scale bars are normalized.

Wang, W. et al. Free-electron lasing at 27 nanometres based on a laser wakefield accelerator. Nature 595, 516–520 (2021).



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### **Chinese SASE FEL results initially reported in 2021**

- Output at 27 nm with e-beam energy of 490 MeV: Q=30 pC.
- SASE FEL exponential gain of 100 reported after 4.5 m.
- ρ~5 x 10<sup>-3</sup> estimated implied gain length of 0.23 m, but effectively ~ 1m obtained. Saturation needs 20 gain lengths.
- FEL performance is best electron-beam diagnostic.



Fig. 2 | Measurement of undulator radiation. a, Measured transverse radiation pattern of a typical pulse on the X-ray CCD camera located 12 m downstream from the gas target. The scale bar is normalized. b, Shot-to-shot radiation energy over 270 pulses. c, d, Measured radiation spectra (c) and the corresponding electron-beam energy spectra (d) detected by the second spectrometer located at the exit of the undulator. e, f, Image (e) and count profile (f) of the interference pattern generated when radiation propagates through two 10-µm slits with a slit separation of 40 µm.

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# Summary Table of SASE FELs Driven by rf Linac Beams

Table 1. Representative parameters of some FEL experiments and facilities.\* Somewhat dated from 2017. New SC rf linacs now.

Parameter	UCLA /LANL	VISA	LEUTL	TTF	FERMI	LCLS-1	SACLA
γmc <sup>2</sup> (GeV)	0.018	0.071	0.22	0.233	1.2	13.6	6.14
σ <sub>γ</sub> /γ (%)	0.25	0.18	0.2	0.15	0.015	0.01	0.01
I (kA)	0.17	0.25	0.2	0.4	0.8	3	3
ε <sub>xn</sub> (μm)	4	2	7	6	0.8	0.4	0.85
$2\pi\beta_x$ (m)	1.2	1.8	9	6	50	140	185
λ <sub>u</sub> (cm)	2.05	1.8	3.3	2.7	3.5	3	1.5
К	1.2	1.04	3.1	1.2	~1	3.5	1.36
L <sub>u</sub> (m)	2	4	20	27	20	110	100
λ <sub>1</sub> (nm)	12000	800	385	100	4	0.15	0.1

\*Kim, Huang, Lindberg, "Synchrotron Radiation and FELs" book, p 225.

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# SUMMARY

- The two discussed PWFA cases show promise, but also reveal challenges on e-beam quality: Charge, energy spread, pulse length, transverse emittance, beam-parameter fluctuations.
- Pierce parameter can be a guide for high gain FEL scaling.
- Slippage of e-beam within undulator should be checked.
- FEL performance can be electron-beam-quality diagnostic.
- The step to x-ray wavelengths carries serious challenges.
- SASE FEL has spiky spectra due to starting from noise which will also have the PWFA beam energy fluctuations.
- Solutions to enhance microbunching at short wavelengths would help. Laser Seeding at harmonics, long. compression of e-beam, laser undulator (?), High-gain harmonic generation...
- Single-shot microbunching diagnostics at visible wavelengths: LaBerge talk Wed. AM, also recent Nature Photonics article

#### **Beam-based Plasma-wakefield Acceleration: Rome, Italy**

- Electron Source is PC rf gun plus three rf accelerating sections giving 88 MeV acceleration of driver and witness bunches.
- Witness is accelerated by wakefields in a capillary to 93 MeV and directed into 6 undulators. Exponential gain was reported for FEL at 826  $\pm$ 9 nm. L<sub>a</sub> ~1.2 m



Fig. 1 | Experimental setup. The driver (D) and witness (W) electron bunches are produced by the photo-injector and their temporal separation is continuously monitored with a non-intercepting EOS diagnostics. The bunches are focused by a triplet of PMQs in a 3-cm-long capillary containing the plasma produced by ionizing hydrogen gas with a high-voltage discharge. The accelerated witness is extracted by a second triplet of PMQs and transported using six electromagnetic quadrupoles. A dipole spectrometer is used to measure its energy with a scintillator screen installed on a 14° beamline. The FEL beamline consists of six planar undulators with tunable gaps and five quadrupoles in between to transport the beam. The emitted FEL radiation is collected by an in-vacuum metallic mirror and measured with an imaging spectrometer equipped with a diffraction grating and a cooled intensified camera (iCCD).

R. Pompili et al., Nature, (2022).



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### Plasma Wakefield Acceleration of the Witness Bunch (20 pC)

• Wakefield accelerates witness by ~5 MeV.



Fig. 2 | Witness acceleration in plasma. Snapshots of the driver (D) and witness (W) spectrum with plasma turned off (a) and on (b). In a, the RF deflector is turned on to vertically separate the two bunches. In b, the decelerated driver energy spectrum is obtained by merging the images obtained with different currents of the spectrometer. c, Energy (top) and spread (bottom) distributions of 500 consecutive shots of the accelerated

witness. **d**, Numerical simulations. The top plot shows a snapshot of the two bunches moving through the plasma background. The white dashed line shows the axial accelerating field along the co-moving coordinate  $\xi$ . The plasma density is reported by the colour bar, in units of cm<sup>-3</sup>. The evolution of the average energy (blue) and energy spread (red) along the capillary longitudinal coordinate *z* is reported in the bottom plot.



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