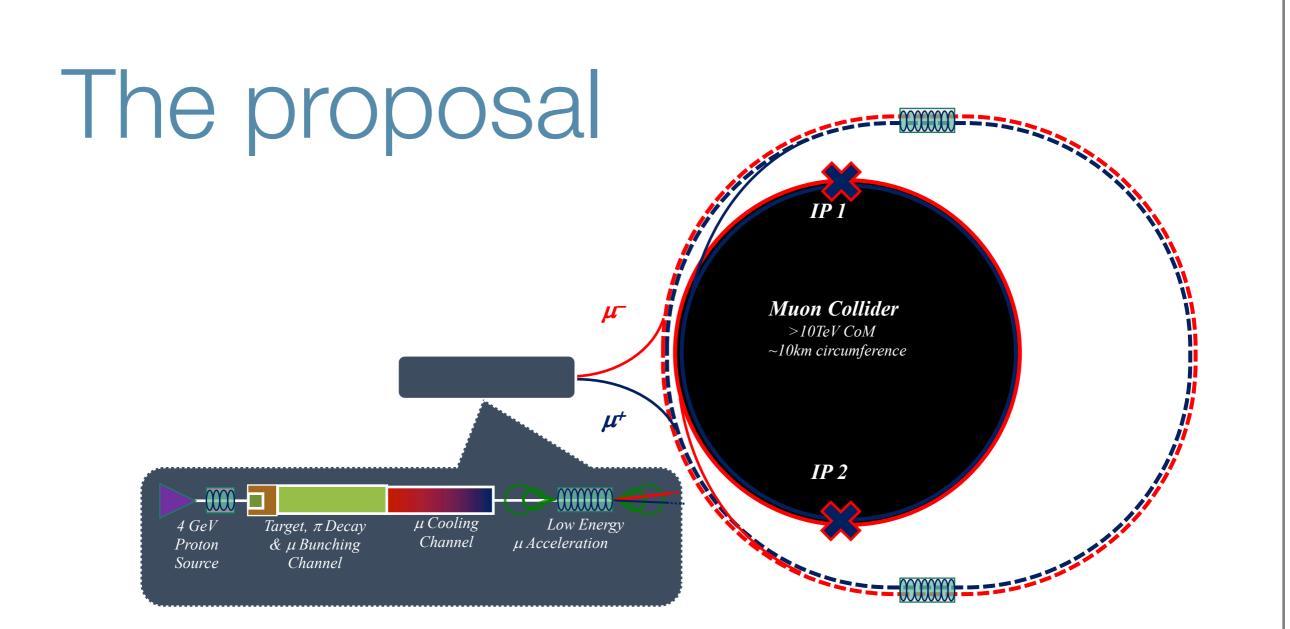
Physics potential of Muon Collider

LianTao Wang University of Chicago

Muons in Minneapolis workshop, Dec 6, U. Minnesota



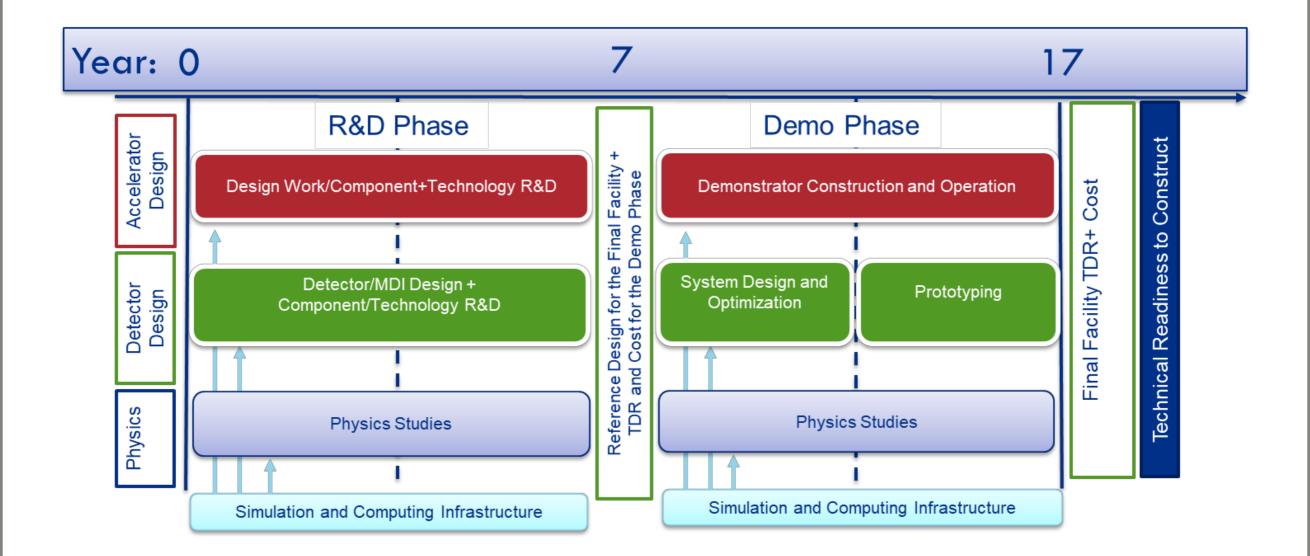
Running scenarios $E_{\rm CM} = 3$ TeV, $\mathscr{L} = 1$ ab^{-1} $\mathscr{L} \propto E_{\rm CM}^2$ $E_{\rm CM} = 10$ TeV, $\mathscr{L} = 10$ ab^{-1}

Ready to build?

Concepts		WFA MuC Spp	c FCC-hh
Collider Con	Collider-in-Sea	ReLIC (≤3 TeV) Multi-TeV ILC CCC (Nb ₃ Sn) (TeV)	FCC-eh CLIC TeV ILC (Nb)
Technical Maturity	 Low maturity conceptual development. Proof-of-principle R&D required. Concepts not ready for 	 Emerging accelerator concepts requiring significant basic R&D and design effort to bring to maturity. 	 prior R&D and design efforts. Critical project risks have been identified and sub-system focused
Funding Approach T	 facility consideration. Funding for basic R&D required. Availability of "generic' accelerator test facility access often necessary 	• Some large-ticket demonstrators are generally	• Funding approach typically cransitions to "project-style" efforts with significant dedicated
			M. Palmer

Before the decision to build it, needs more R&D.

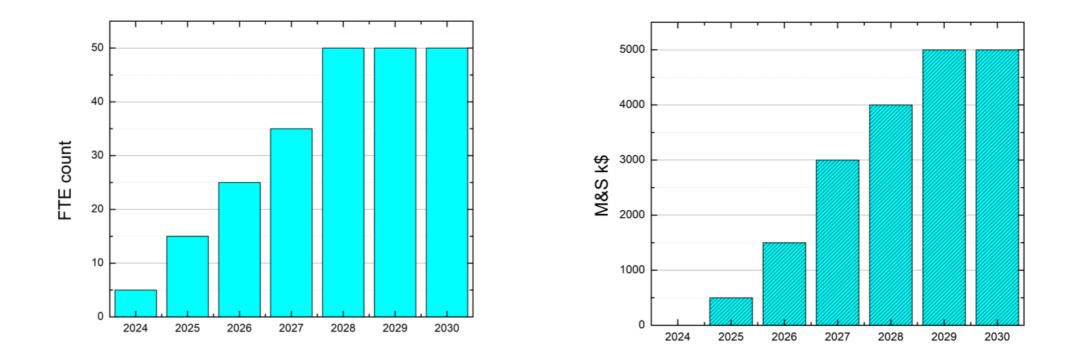
A plan.



Ready to build in early 2040s.



Need support

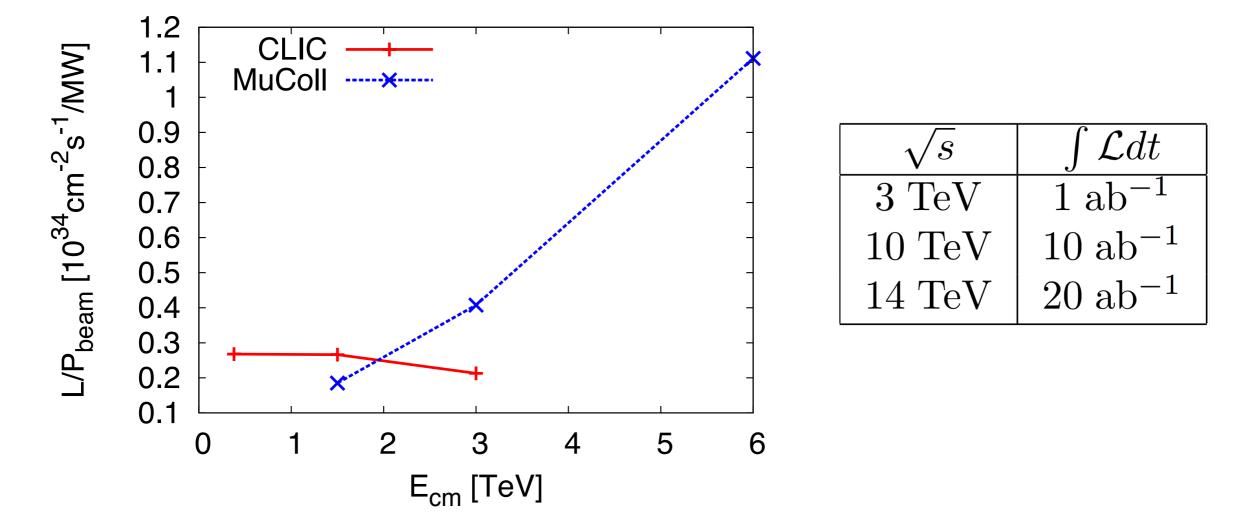


International Muon Collider Collaboration established after European strategy update in 2020.

US support (money, people) will be also crucial A first step: P5 report (~ 24 hours from now)

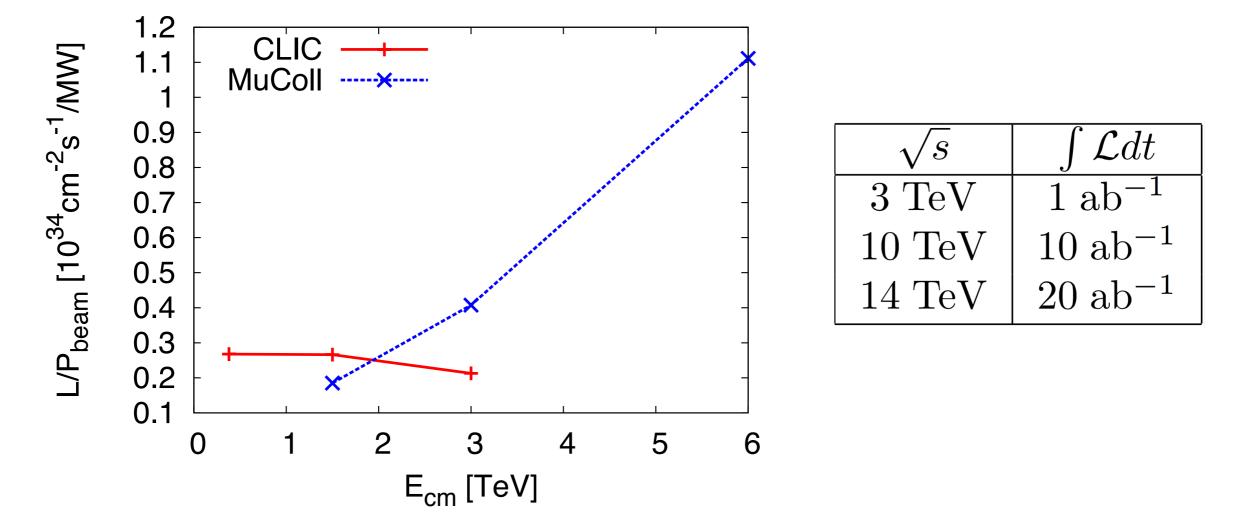
Why muon collider?

The obvious: higher energy, shorter distances



All Carifield

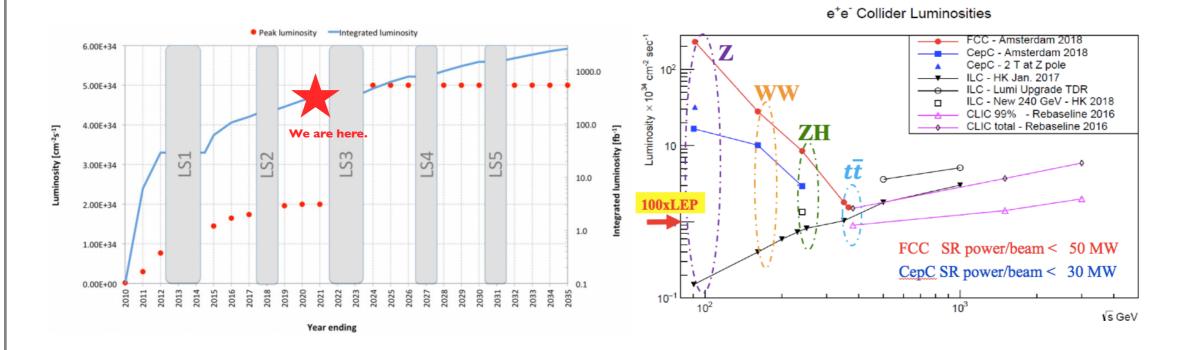
The obvious: higher energy, shorter distances



Good enough? Yes (for most of us).

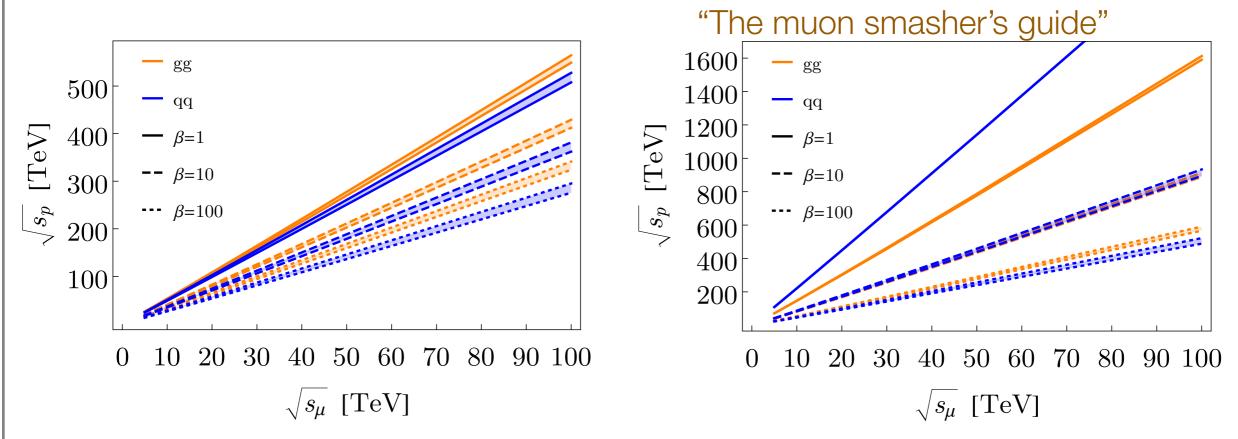
Still why much collider at these energies?

The coming decades



- * Main "near term" targets: precision, rare processes.
- Muon collider, going beyond these options, such as LHC and the low energy Higgs factories.

Comparison with 100-ish TeV pp collider Such as FCC-hh or SPPC



- * Naively, 100 TeV pp \approx 10+ TeV muon collider.
- Lepton collider "cleaner". Good for precision, search in difficult channels.

- * Are there guaranteed discoveries?
 - * No.

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- * It is a significant step beyond our current reach.
 - * A lot of interesting physics to cover.

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 - * No.
- * It is a significant step beyond our current reach.
 - * A lot of interesting physics to cover.
- * Higher energy is of course better for physics.
 - * Limited by resource.

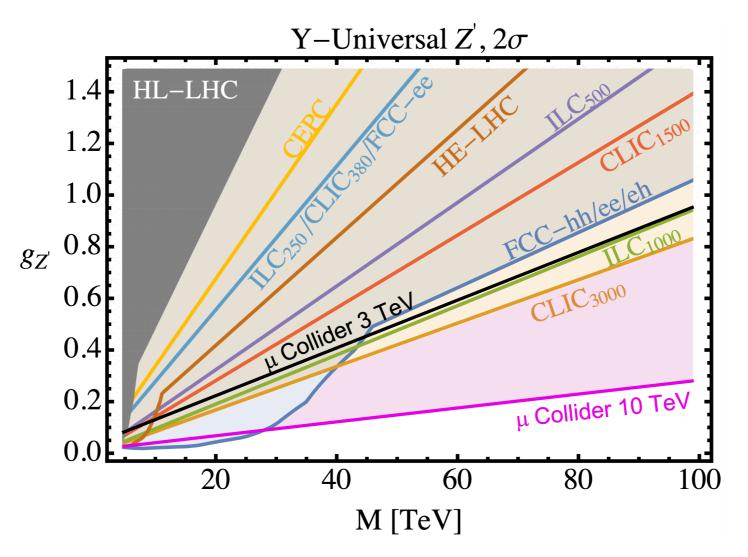
- * Are there guaranteed discoveries?
 - * No.

Rest of the talk

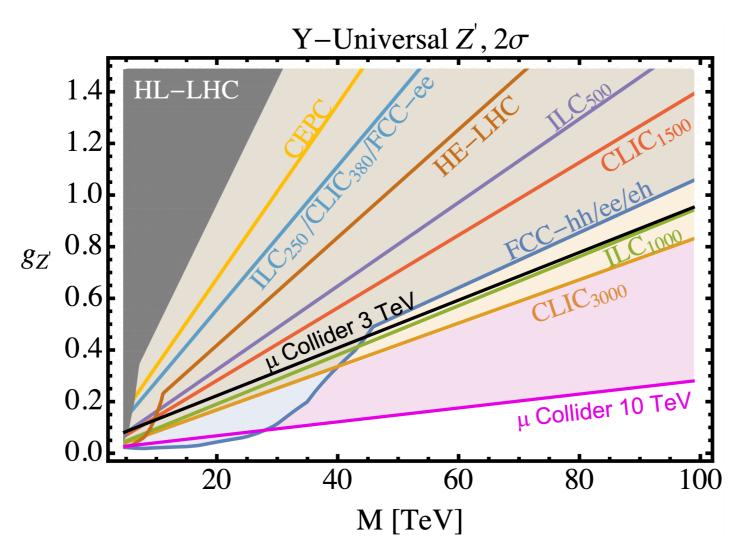
It is a significant step beyond our current reach.

- * A lot of interesting physics to cover.
- * Higher energy is of course better for physics.
 - * Limited by resource.

Obvious: big step in NP searches

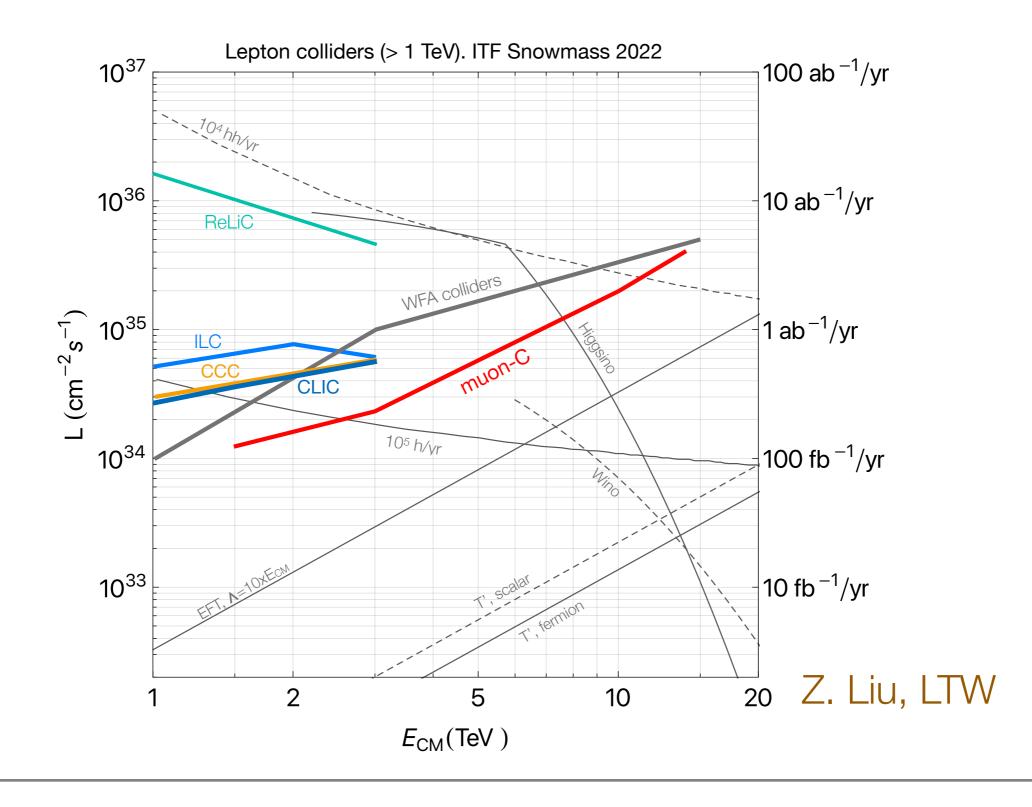


Obvious: big step in NP searches



What are the interesting questions it can help answer?

Basic physics output



Physics program at a muon collider

- * Higgs and electroweak.
- * New physics at higher energies.
 - Dark matter
 - * Flavor, CP
 - *

References:

Muon smasher's guide, IMCC input to Snowmass, Muon collider forum report Snowmass BSM working group report

Higgs and electroweak

- Obvious: important to understand the Higgs (more broadly, electroweak physics) better.
 - * Origin of the weak scale.
 - * Nature of EW symmetry breaking.

Higgs and electroweak

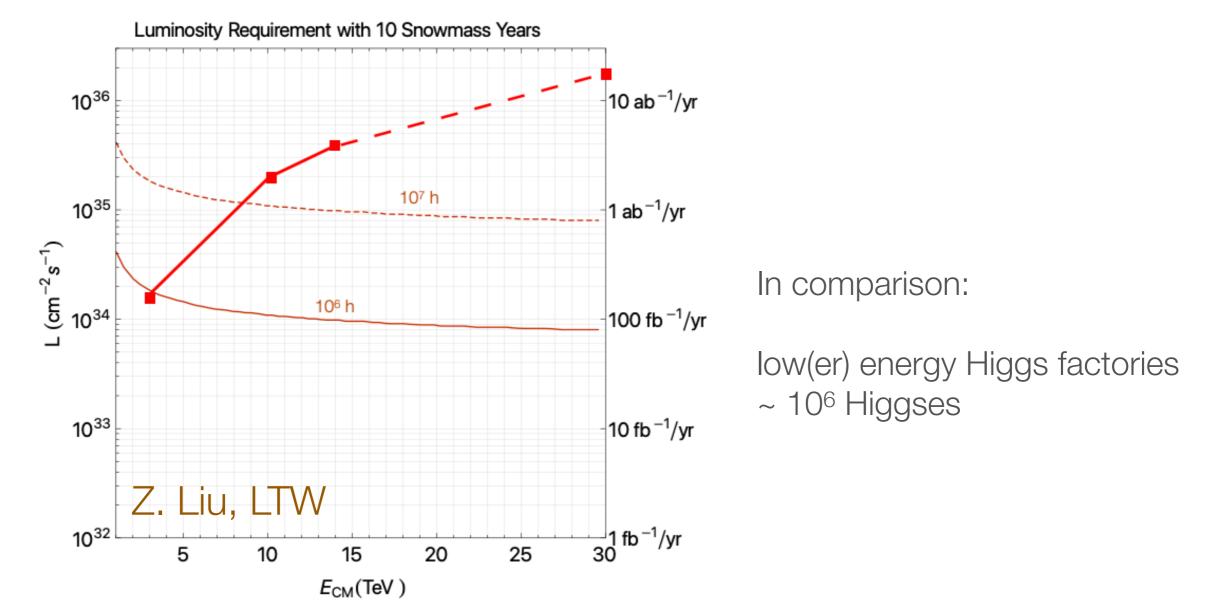
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Higgs and electroweak

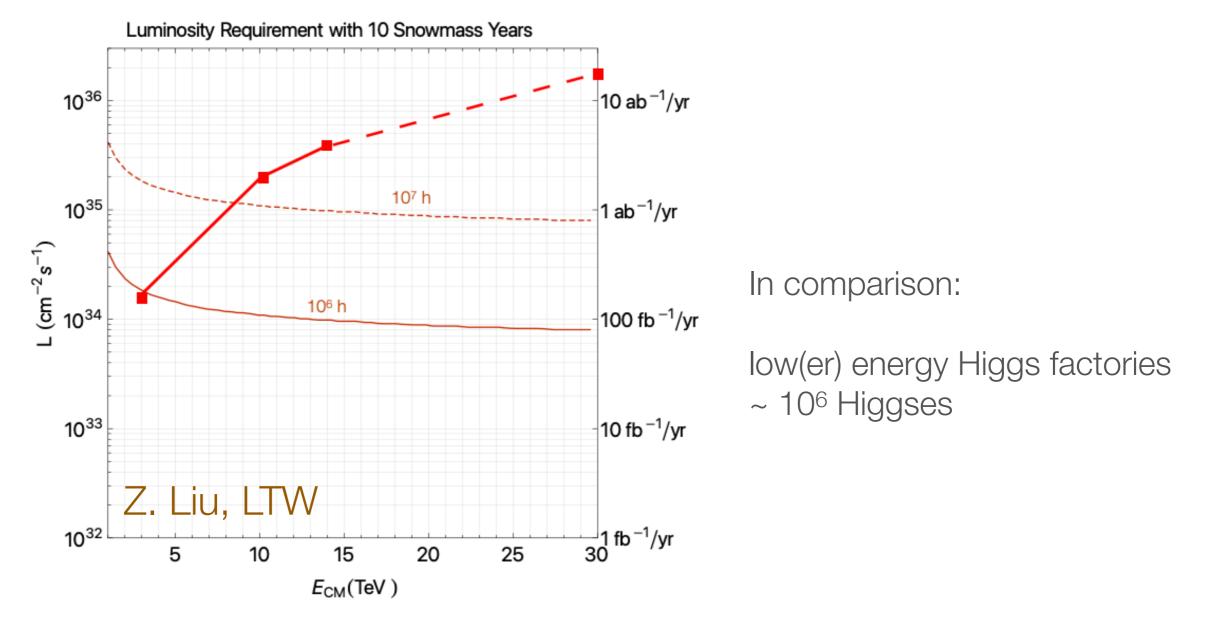
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 - * Nature of EW symmetry breaking.
- * Not so obvious: what does muon collider bring to the table?

It will reach energies much beyond the electroweak scale!

MC as Higgs factory

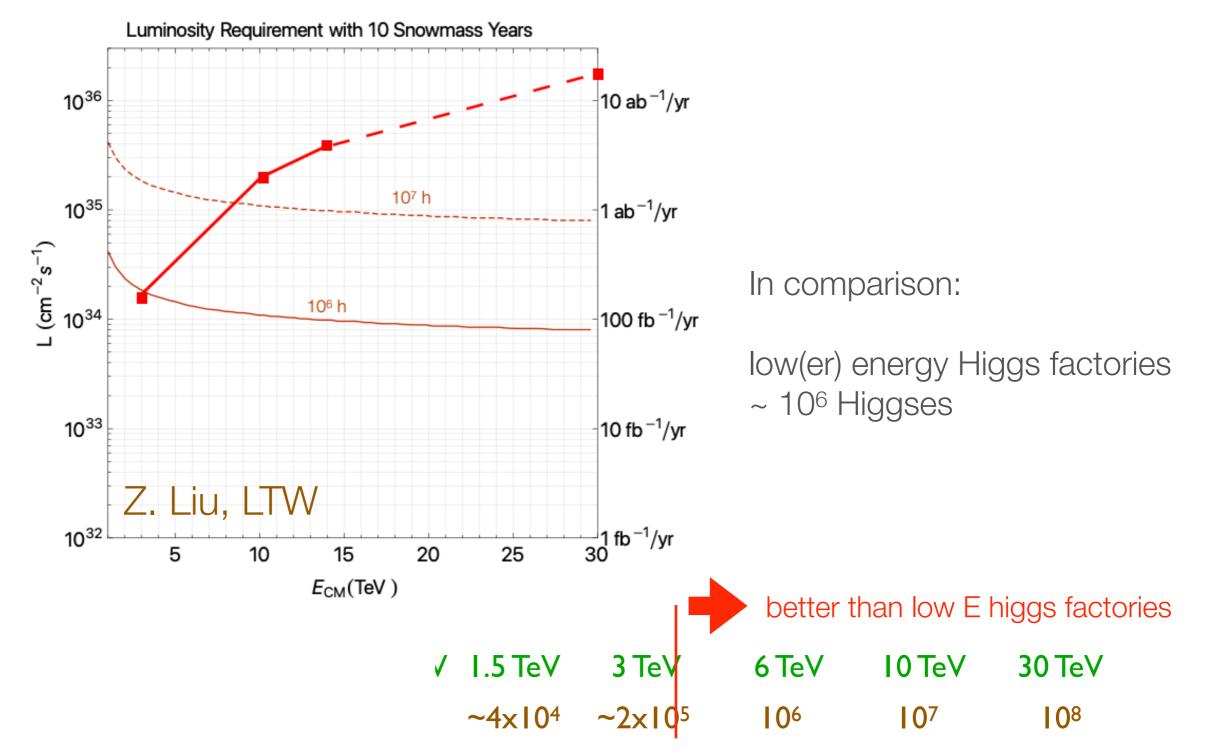


MC as Higgs factory



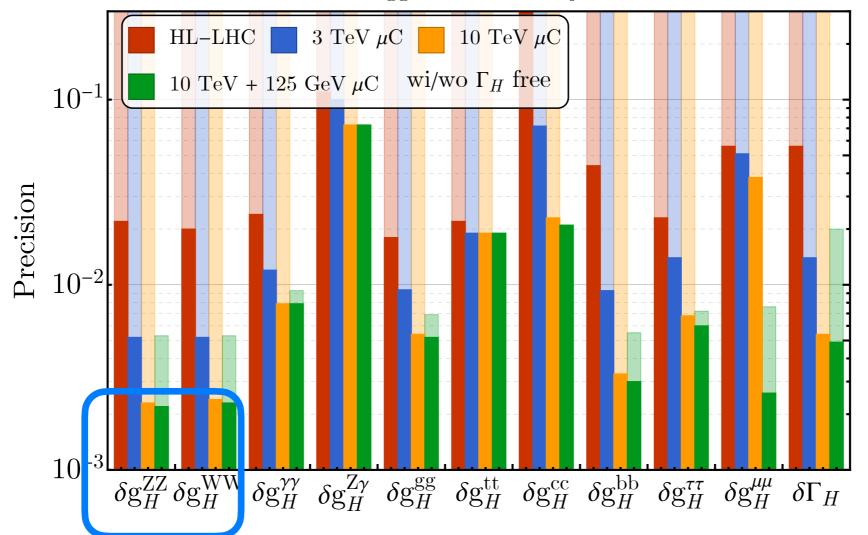
✓ I.5 TeV 3 TeV 6 TeV 10 TeV 30 TeV
 ~4×10⁴ ~2×10⁵
 10⁶
 10⁷
 10⁸

MC as Higgs factory



Higgs precision

Muon Collider Higgs Precision Projections (SMEFT)



0.1% level or better measurement possible at higher energies. A factor of 10 better than the HL-LHC Comparable to e+e- Higgs factories

Introduction

The discovery of the Higgs boson at the CERN Large Hadron Collider (LHC) opens a new venue in particle physics. On the one hand, the existence of the Higgs boson completes the particle spectrum in the Standard Model (SM) and provides a self-consistent mechanism in uantum field theory for mass generation of elementary particles. On the other hand, the SM loes not address the underlying mechanism for the electroweak symmetry breaking (EWSB) nd thus fails to understand the stability of the weak scale with respect to the Planck scale. In rder to gain for ther insight for those fundamental questions, it is of high prio199 ab study the liggs boson properties to high precision in the hope to identify hints for new physics beyond Higgs-Higgs-Gauge Boson couplings.^{MS ad}

In the SM, the Higgs sector is constructed from a complex scalar doublet Φ . After he EWSB, the neutral real Winpowent is the Higgs boston excitation H and the other three legrees of freedom become the longitudinal components of the massive gauge bosons. As such, study. Here it is made implicit that $\kappa_{V_r} = \kappa_W = \kappa_Z$. tudying the Higgs-gauge boson couplings would be the most direct probe to the underlying been verified to a good accuracy by precision EW measurement nechanism of the electroweak symmetry breaking. After the EWSB, the Higgs sector can bush to be more general and will not be assuming a correlate A fully consistent and theoretically-sound framework would arameterized as (EFT), by augmenting the SM Lagrangian with higher dimension

$$\mathcal{L} \supset \left(M_W^2 W_\mu^+ W^{-\mu} + \frac{1}{2} M_Z^2 Z_\mu Z^\mu \right) \left(\kappa_V \frac{2H}{v} + \kappa_V_2 \frac{H^2}{v^2} \right) - \frac{m_H^2}{2v} \left(\kappa_3 H^3 + \frac{1}{4v} \kappa_4 H^4 \right), \text{ out the heavier states [3]. While a systematic account for the effective of the composed and beyond the scope of the composed of the composed and beyond the scope of the composed of the composed and beyond the scope of the composed of the composed and beyond the scope of the composed of the composed and beyond the scope of the composed of the composed and beyond the scope of the composed and the scope of the composed and beyond the scope of the composed and beyond the scope of the composed and the$$

 $\sim c_H \frac{1}{\Lambda^2}$

consider the following two operators for the purpose of illustration **Higgs-Gauge B** $\mathcal{O}_{H} = \frac{c_{H}}{2\Lambda^{2}}\partial_{\mu}(\Phi^{\dagger}\Phi)\partial^{\mu}(\Phi^{\dagger}\Phi)$, $\mathcal{O}_{6} = -\frac{c_{6}\lambda}{\Lambda^{2}}(\Phi^{\dagger}\Phi)$ -Gauge Boson Reuplings fields for the vacuation of the Higgs field v = 246 GeV is the vacuation expectation value of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation of the Higgs field v = 246 GeV is the vacuation expectation expectation of the Higgs field v = 246 GeV is the vacuation expectation expectation of the Higgs field v = 246 GeV is the vacuation expectation expectation of the Higgs field v = 246 GeV is the vacuation expectation expec ouplings at tree-level. This " κ -scheme" is a conversion phenomenological parameterization f deviations from the SM expectations, which is suitable for the exploratory nature of the state where new physics sets in, and λ is the constant of the state front of $(H^{\dagger}H)^2$ term in the SM Higgs potential. At the dimensi

operators that are most relevant for our study. An additional open be removed by a suitable field-redefinition [5]. The resulting shi Ч. are^{1} $\Delta \kappa_V = -\frac{c_H}{2} \frac{v^2}{\Lambda^2}, \qquad \lambda \kappa_V = -\frac{c_H}{2} \frac{v^2}{\Lambda^2}, \qquad \Delta \kappa_{V2} = -2c_H \frac{v^2}{\partial_{\mu}(H^{\dagger}H)\partial^{\mu}(H^{\dagger}H)}, \qquad \mathcal{O}_H = \frac{\partial_{\mu}(H^{\dagger}H)\partial^{\mu}(H^{\dagger}H)\partial^{\mu}(H^{\dagger}H)}{\mathcal{O}_H} = \frac{\partial_{\mu}(H^{\dagger}H)\partial^{\mu}(H^{\dagger}H)\partial^{\mu}(H^{\dagger}H)}{\mathcal{O}_H}$

We see that deviations in the $V \psi H$ and $VVHH (V \neq W h^{\pm}, Z)$

out the heavier states [3]. While a systematic account for the effe

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Introduction

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We see that deviations in the $V \psi H$ and $VVHH (V \neq W h^{\pm}, Z)$

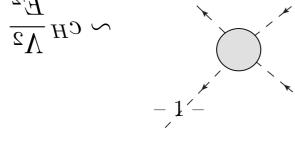
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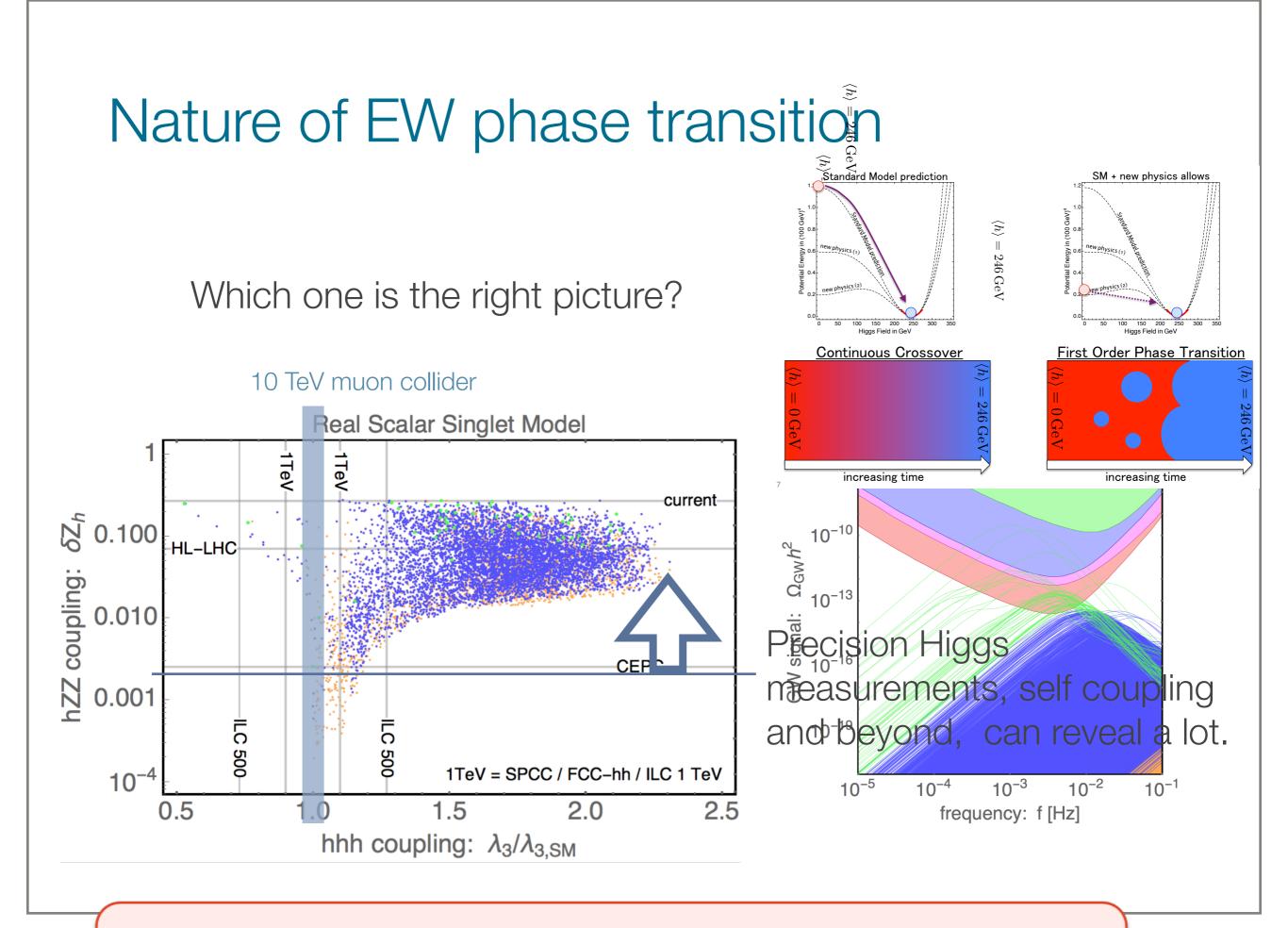
 are^{1}

$$\mathcal{L} \supset \left(M_W^2 W_{\mu}^+ W^{-\mu} + \frac{1}{2} M_Z^2 Z_{\mu} Z^{\mu} \right) \left(\kappa_V \frac{2H}{v} + \kappa_{V_2} \frac{H^2}{v^2} \right) - \frac{m_H^2}{2v} \left(\kappa_3 H^3 + \frac{1}{4v} \kappa_4 H^4 \right)$$

$$\frac{(H^{\dagger} H)^{\mu} 6(H^{\dagger} H)_{\mu} 6}{(H^{\dagger} H)^{\mu} 6(H^{\dagger} H)_{\mu} 6} = H^0$$

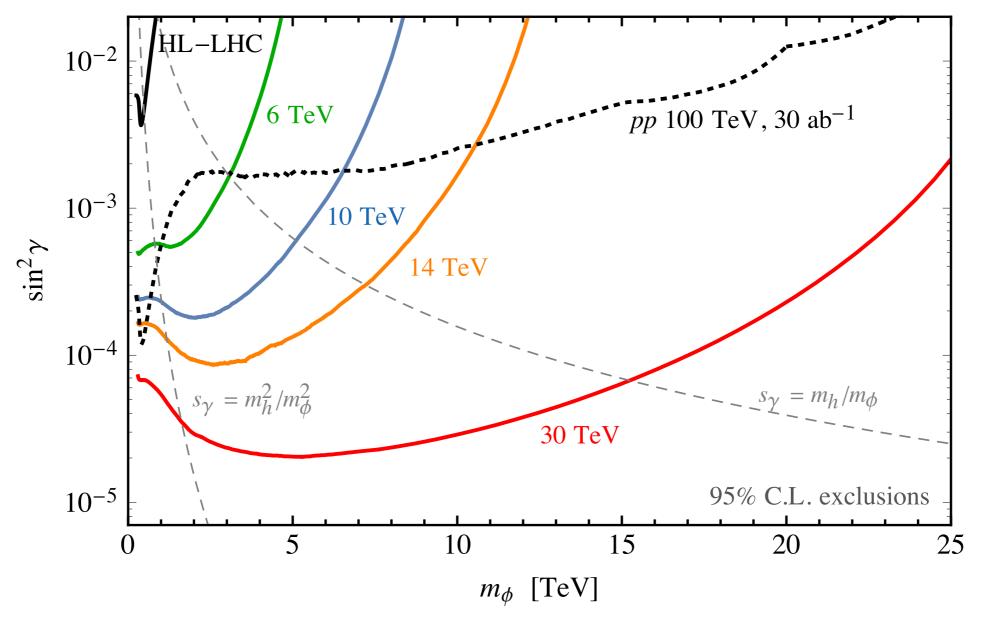
where v = 246 GeV is the vaction expectation value of the Higgs field and the expectation value of the exploratory nature of the expectation of the Higgs potential. At the dimension of the Higgs potential.





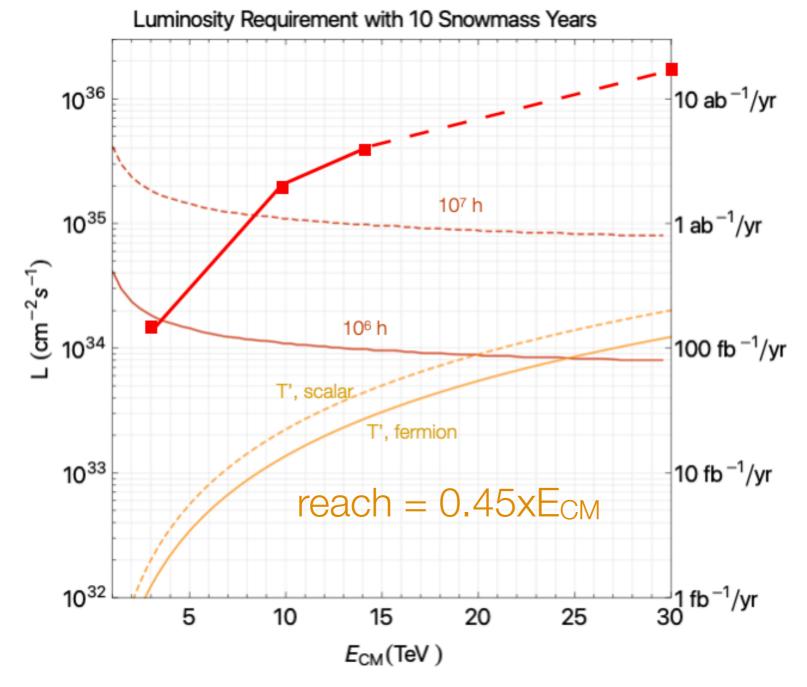
Higgs's friends

D. Buttazzo, D. Redigolo, F. Sala, A. Tesi, 1807.04743



Singlet scalar mixing with the Higgs.

Top partner



Pair production

$$\mu^+\mu^- \to T'\tilde{T}'$$

Examples: SUSY stop Composite T

. . .

Spectacular signal, 10s event needed for discovery

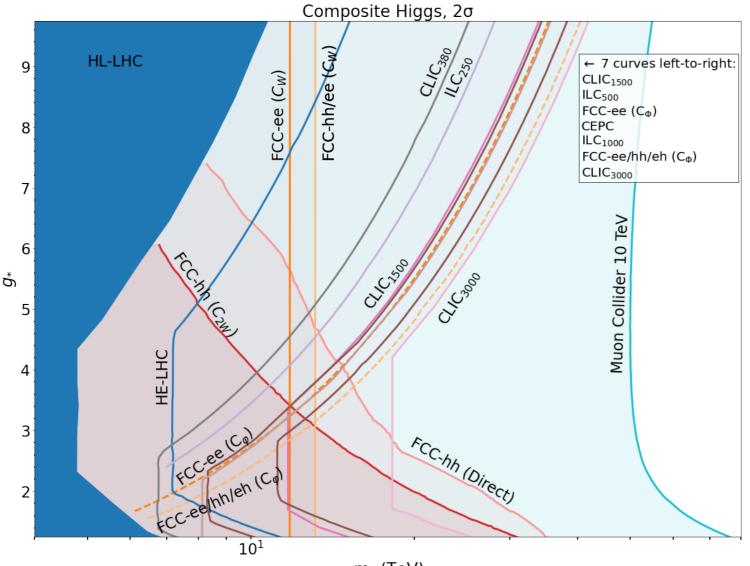
Energy⇒precision

- The effect of heavy new physics can be parameterized by higher dimensional (EFT) operators.
- * Their effect grows at higher energies.
 - * e.g. if new physics lead to dim-6 operators

$$\frac{\mathcal{O}^{(6)}}{\Lambda_{\rm NP}^2} \to d\sigma \propto E^2$$

Higher energy \Rightarrow better precision

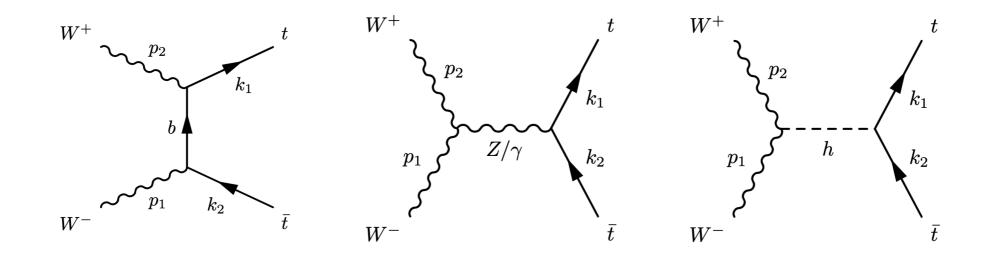
Composite Higgs



m∗ (TeV)

 $\frac{c_{\phi}}{\Lambda^2} \frac{1}{2} \partial_{\mu} \left(\phi^{\dagger} \phi \right) \partial^{\mu} \left(\phi^{\dagger} \phi \right)$

Top quark - Higgs coupling



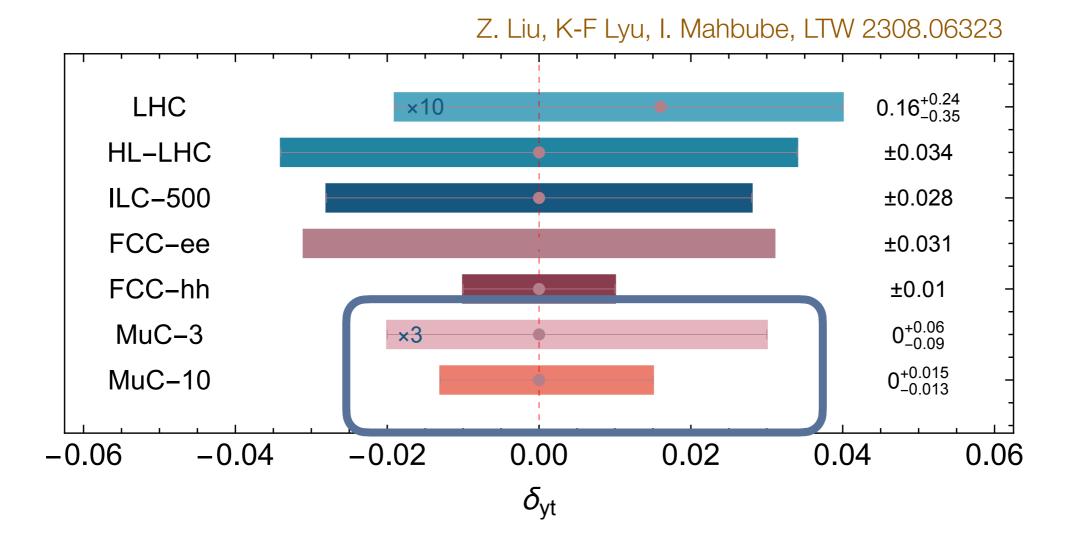
Largest Higgs coupling.

Plays an important role in weak scale dynamics. Sensitive to new physics.

e.g. NP effect =
$$\frac{1}{\Lambda^2} H^{\dagger} H H Q t \rightarrow \mathcal{M} \propto E$$

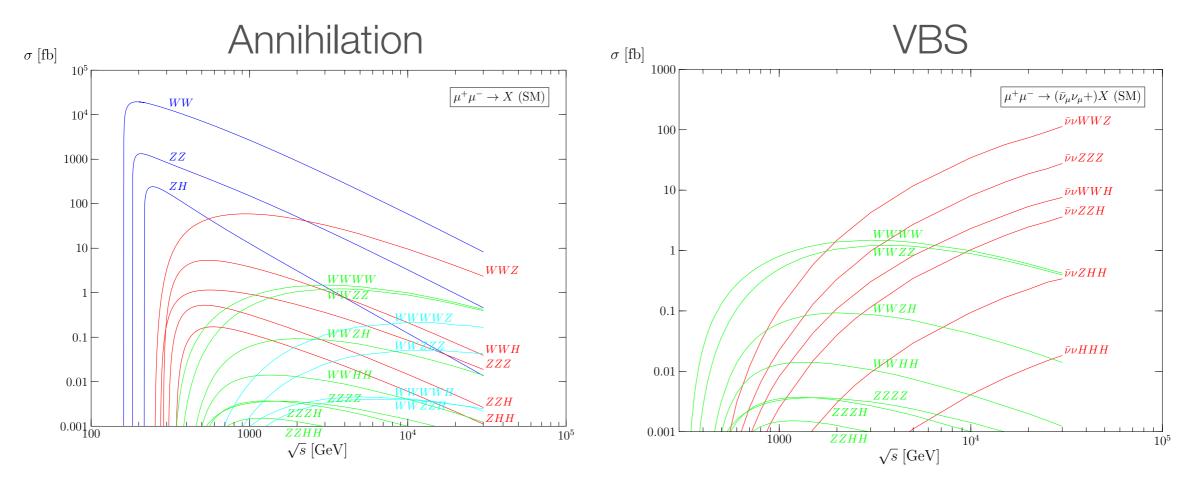
Larger effect at high energies!

Top quark - Higgs coupling



Higher energies at muon collider lead to better precision.

Multi-boson production



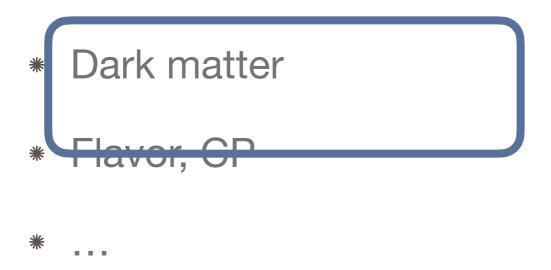
From W. Kilian, PittPACC workshop on muon Collider, 2020 Calculation: WHIZARD (\Rightarrow from CLICdp studies, cf. CERN YR / arXiv:1812.02093.)

W, Z, h "massless".

Can be sensitive to higher order NP effects.

Physics program at a muon collider

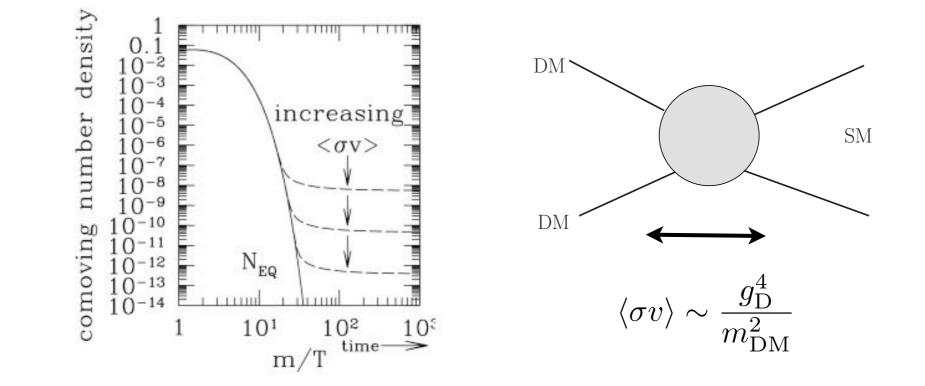
- * Higgs and electroweak.
- * New physics at higher energies.



References:

Muon smasher's guide, IMCC input to Snowmass, Muon collider forum report Snowmass BSM working group report

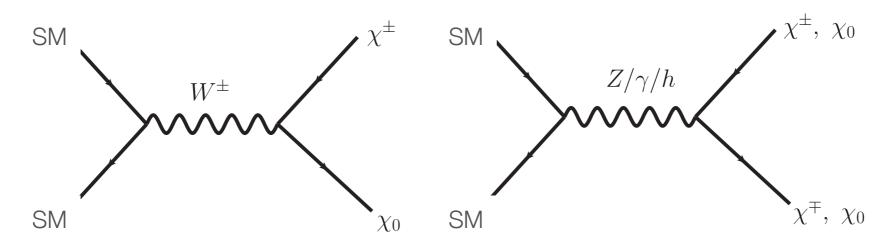
WIMP:



 Simple assumption: DM in thermal eq. with the SM in early universe

Simplest model: part of an EW multiplet

"Minimal dark matter", Cirelli, Fornengo and Strumia, hep-ph/0512090, 0903.3381



- * Simplicity: there is no additional new mediator.
 - * Mediated by W/Z/h. Very predictive.
- * In SUSY, there are two such examples
 - * Higgsino: doublet. Wino: triplet.

Thermal targets

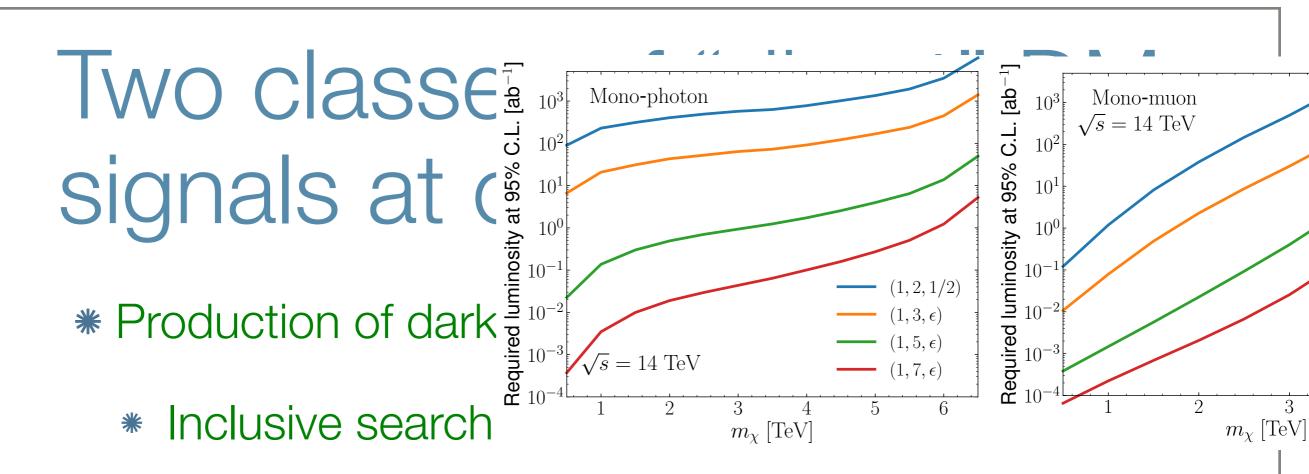
$\begin{array}{c} \text{Model} \\ (\text{color}, n, Y) \end{array}$		Therm. target
$(1,\!2,\!1/2)$	Dirac	1.1 TeV
(1,3,0)	Majorana	2.8 TeV
$(1,3,\epsilon)$	Dirac	2.0 TeV
(1,5,0)	Majorana	14 TeV
$(1,5,\epsilon)$	Dirac	6.6 TeV
(1,7,0)	Majorana	48.8 TeV
$(1,7,\epsilon)$	Dirac	16 TeV

Correct relic abundance ⇒ Thermal targets Reach up to thermal target ≈ complete coverage for WIMP candidate

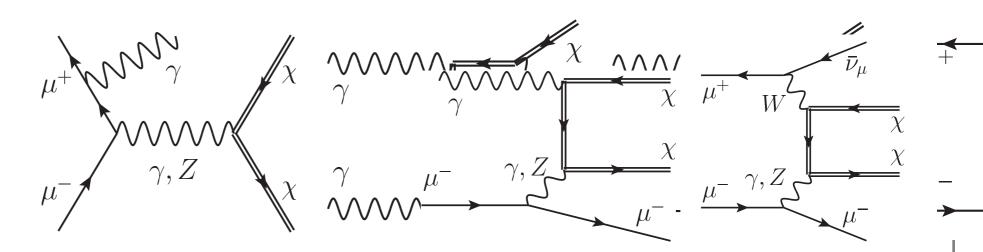
Mitridate, Redi, Smirnov, Strumia, 1702.01141

S. Bottaro, D. Buttazzo, M. Costa, R. Franceschini, P. Panci, D. Redigolo, L. Vittorio, 2107.09688

Way beyond LHC reach.



similar to mono-jet at hadron colliders

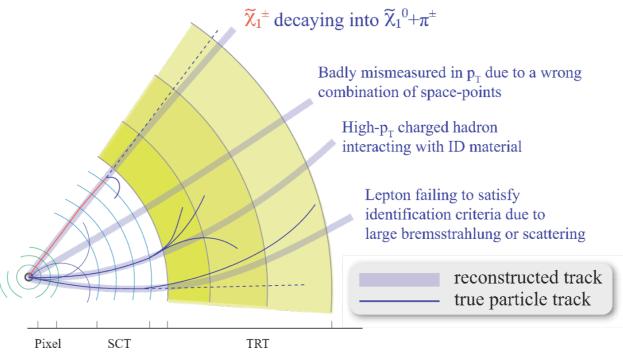


Examples:

Challenges: sizable background, systematics

Two classes of "direct" DM signals at colliders

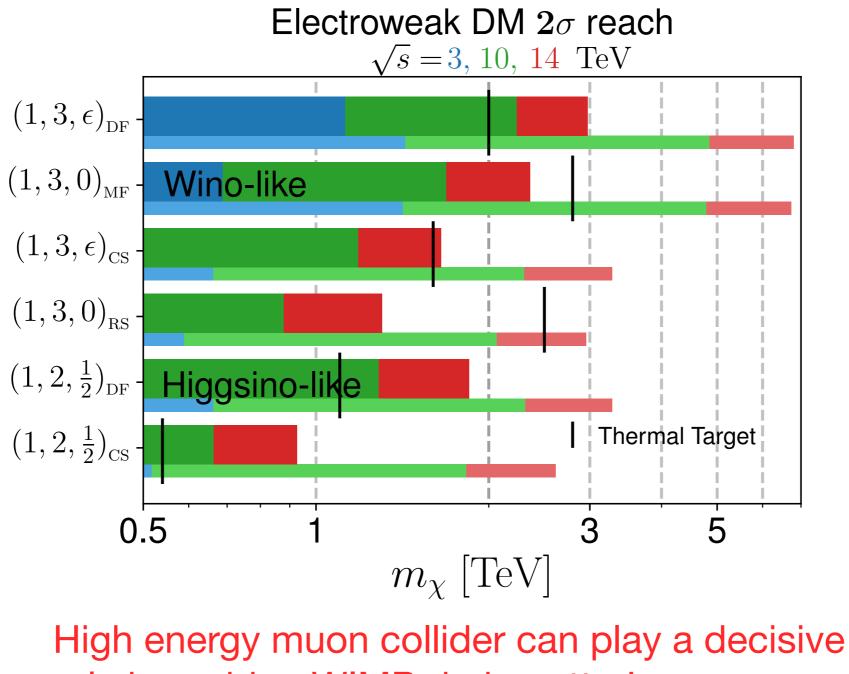
- Small EW induced mass splitting, charged member long-lived.
 - * Disappearing track



Challenge: detector need to be close, beam induced background

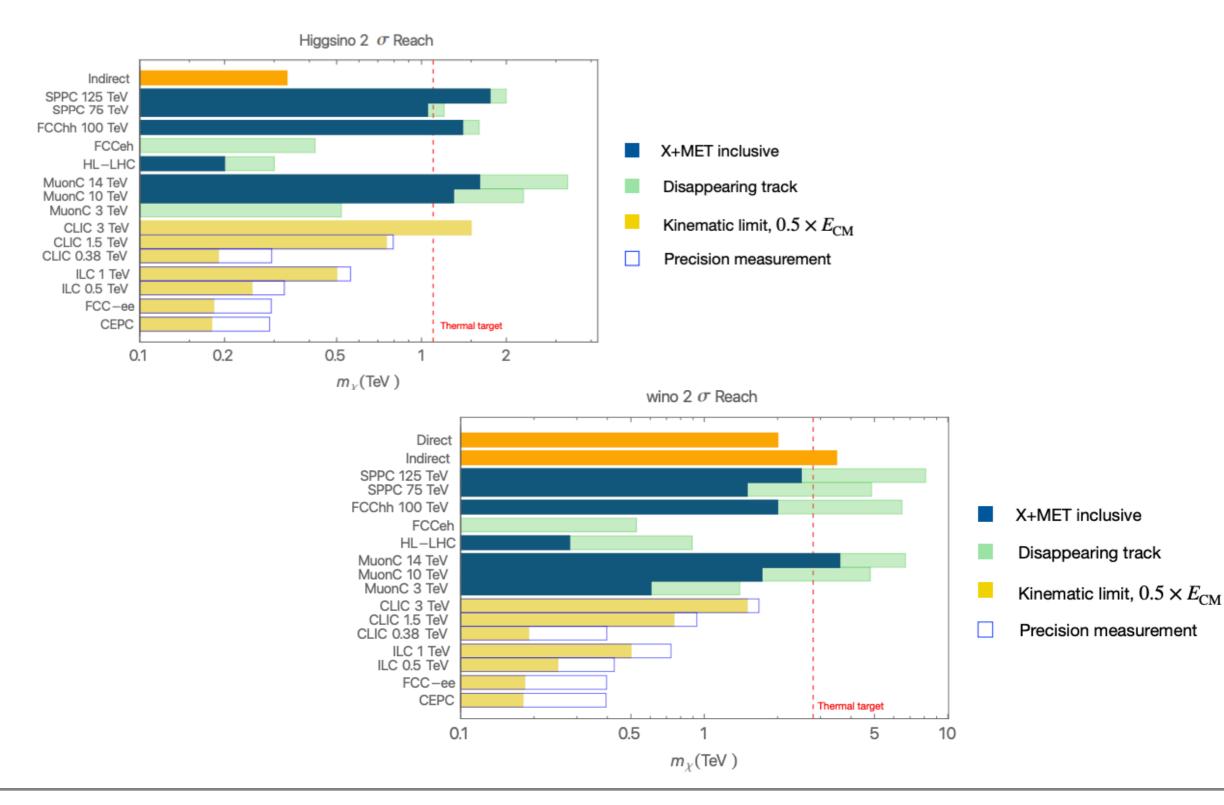
WIMP reach

T. Han, Z. Liu, X. Wang, LTW, 2009.11287, 2203.07351



role in probing WIMP dark matter!

Muon Collider vs others



Physics program at a muon collider

- * Higgs and electroweak.
- * New physics at higher energies.
 - Dark matter

Flavor, CP

References:

Muon smasher's guide, IMCC input to Snowmass, Muon collider forum report Snowmass BSM working group report

Flavor (CP)

- * What is the scale of new flavor/CP physics?
- Flavor (CP) measurements have consistently pushed this to be (far) beyond weak scale.
 - * e.g. Lepton flavor violation
 - * e.g. EDM
- High energy muon colliders offers new windows to probe them.

Lepton flavor violation

Exp limit: $BR(\mu \rightarrow 3e) < 10^{-12}$ Constraint: $\frac{c}{\Lambda^2}(e\Gamma\mu)(e\Gamma e), \Lambda > 2 \times 10^2 \text{ TeV}$ Exp limit: $BR(\tau \rightarrow 3\mu) < 2.1 \times 10^{-8}$

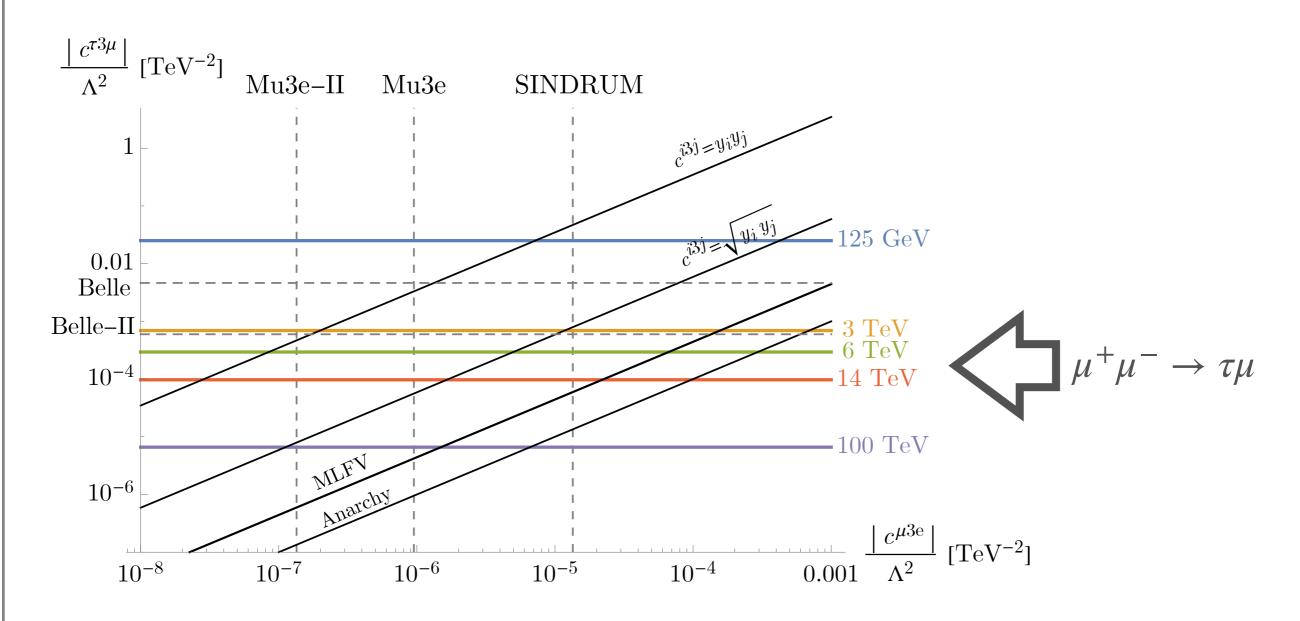
Constraint: $\frac{c}{\Lambda^2}(\mu\Gamma\tau)(\mu\Gamma\mu)$, $\Lambda > 10 \text{ TeV}$

Direct probe at muon colliders: μ^{-1}

$$\mu^+\mu^- \to \ell_i \ell_j$$

Probing lepton flavor violation

S. Homiller, Q. Lu, M. Reece, 2203.08825, Smasher's guide

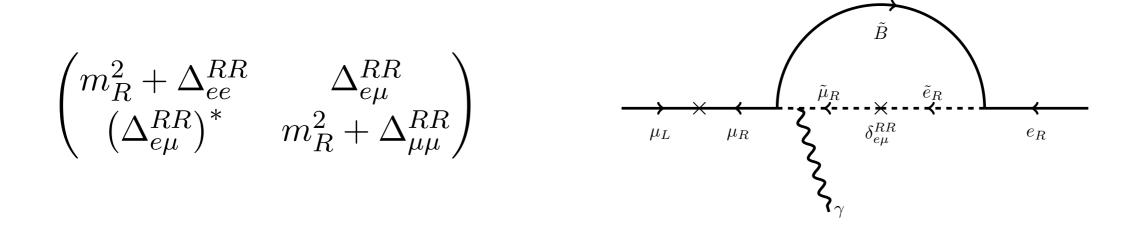


LFV at loop order

- * At loop order, the NP masses is lower.
 - * Opportunity to search for it directly at muon collider.

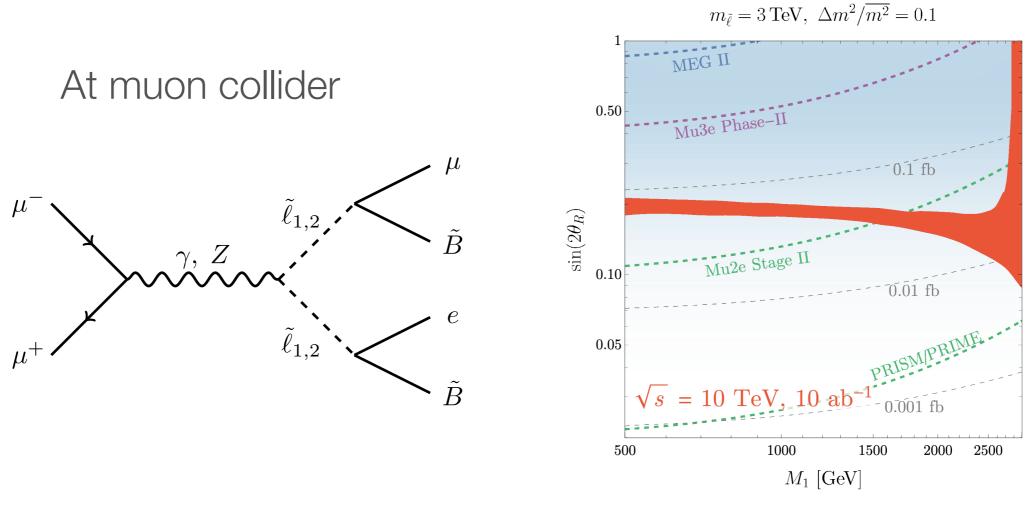
LFV at loop order

- * Supersymmetry.
- * Slepton flavor violation.



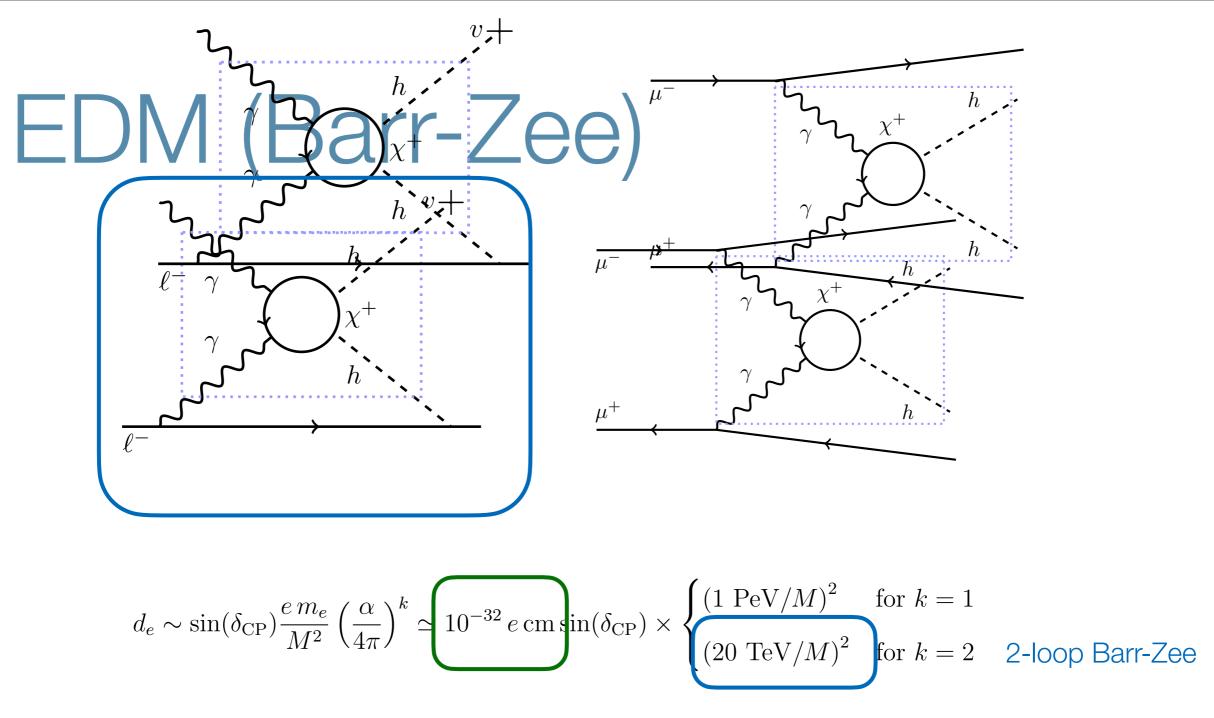
 $BR(\mu \rightarrow e\gamma) < 4.2 \times 10^{-13} \rightarrow m_{SUSY} > TeV$

LFV at loop order

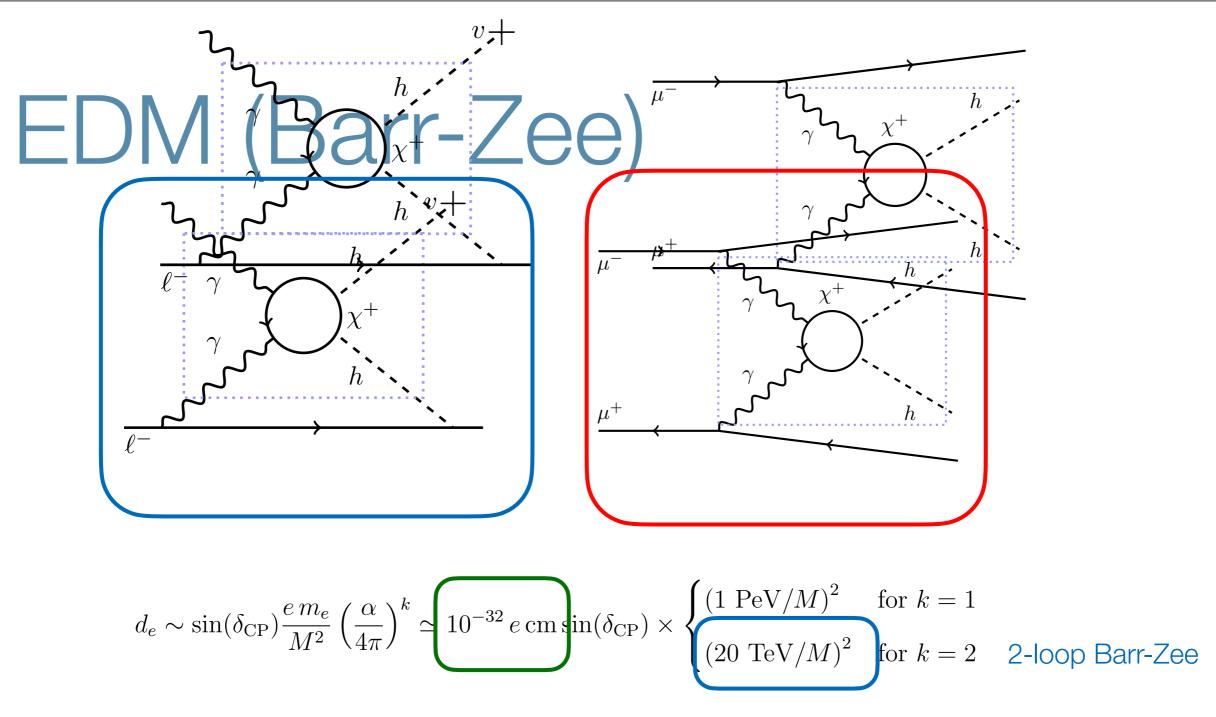


S. Homiller, Q. Lu, M. Reece, 2203.08825

10 TeV muon collider can have interesting reaches. Consinter entang, ton Row Menergy measurement.



Potential sensitivity of next generation exp.



Potential sensitivity of next generation exp.

Same process probed by 10(s) TeV muon collider!

Conclusion

- High energy muon collider holds promise of getting the next (10 TeV) high energy frontier.
- * A lot of interesting physics to cover.
- My hope: we will do solid R&D in the coming decades to make it into a mature project.

DM part of a EW multiplet

 $DM \in (1, n, Y)$ of $SU(3)_C \times SU(2)_L \times U(1)_Y$

- * n odd. Fermionic.
 - * n>7, Landau pole close to M_{DM} .
 - * After EWSB, mass splitting (minimally) generated at 1-loop.
 - * Choose Y=0. Lightest member electric neutral. Potential DM candidate.

DM part of a EW multiplet

 $DM \in (1, n, Y)$ of $SU(3)_C \times SU(2)_L \times U(1)_Y$

- * n even. Fermionic
 - Choose Y=(n-1)/2 ensures lightest member is neutral.
 - Direct detection rules out the minimal case due to tree level Z exchange.
 - Can be avoided to introduce a small splitting, δm > 10² keV, of the neutral states (for example, from a dim-5 operator). Not quite minimal (additional model dependence).
 - Famous example: Higgsino (1,2)_{1/2}

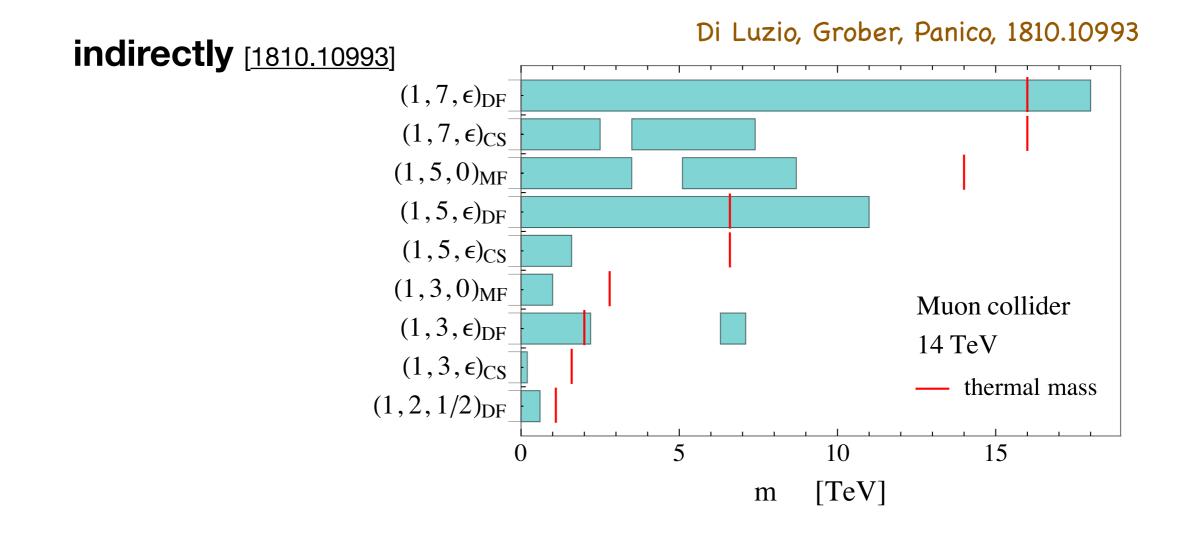
DM part of a EW multiplet

 $DM \in (1, n, Y)$ of $SU(3)_C \times SU(2)_L \times U(1)_Y$

Scalar (real and complex)

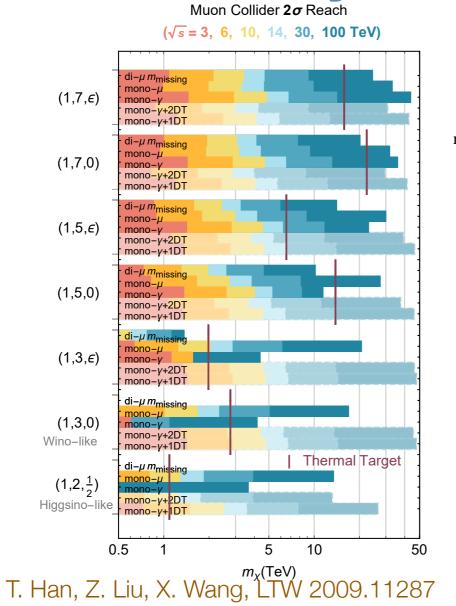
- Minimal case: mass splitting, stability discussion parallel to that of the fermionic multiplets.
- Addition couplings of the form H⁺ H X⁺ X. More parameters involved in a full analysis.
- * More focus on the fermion case (so far).

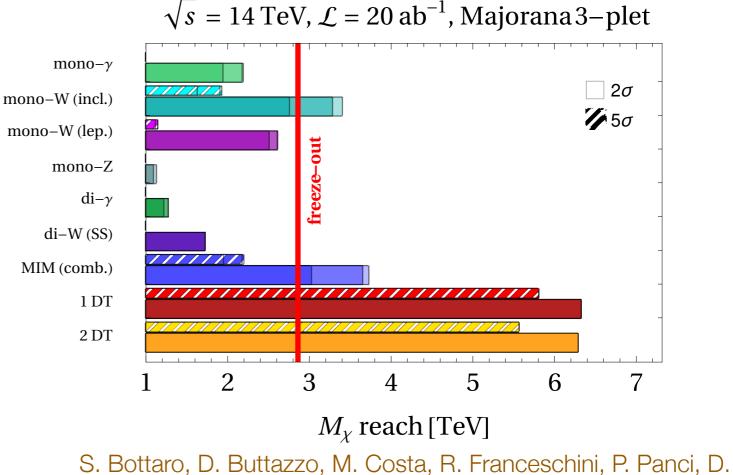
"indirect", from precision measurement



At loop level, modifying the $q\bar{q}(\text{ or } \ell^+\ell^-) \rightarrow f\bar{f}$ amplitude

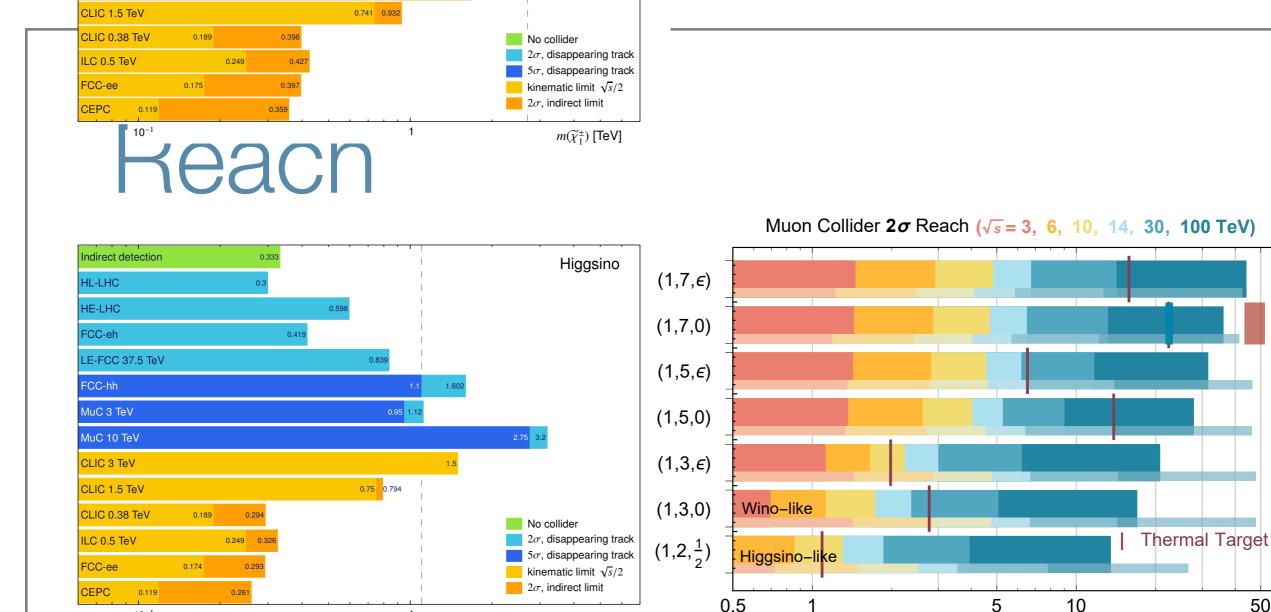
Reach by channel





Redigolo, L. Vittorio, 2107.09688

mono-X, more generic model independent. Interesting channels: muon-mu, mono-W. Disappearing track. Some model dependence. Important to have the right BIB estimates.



 $m(\widetilde{\chi}_1^{\pm})$ [TeV]

R. Capdevilla F. Meloni, J. Zurita, 2102.11292 T. Han, Z. Liu, X. Wang, LTW 2009.11287

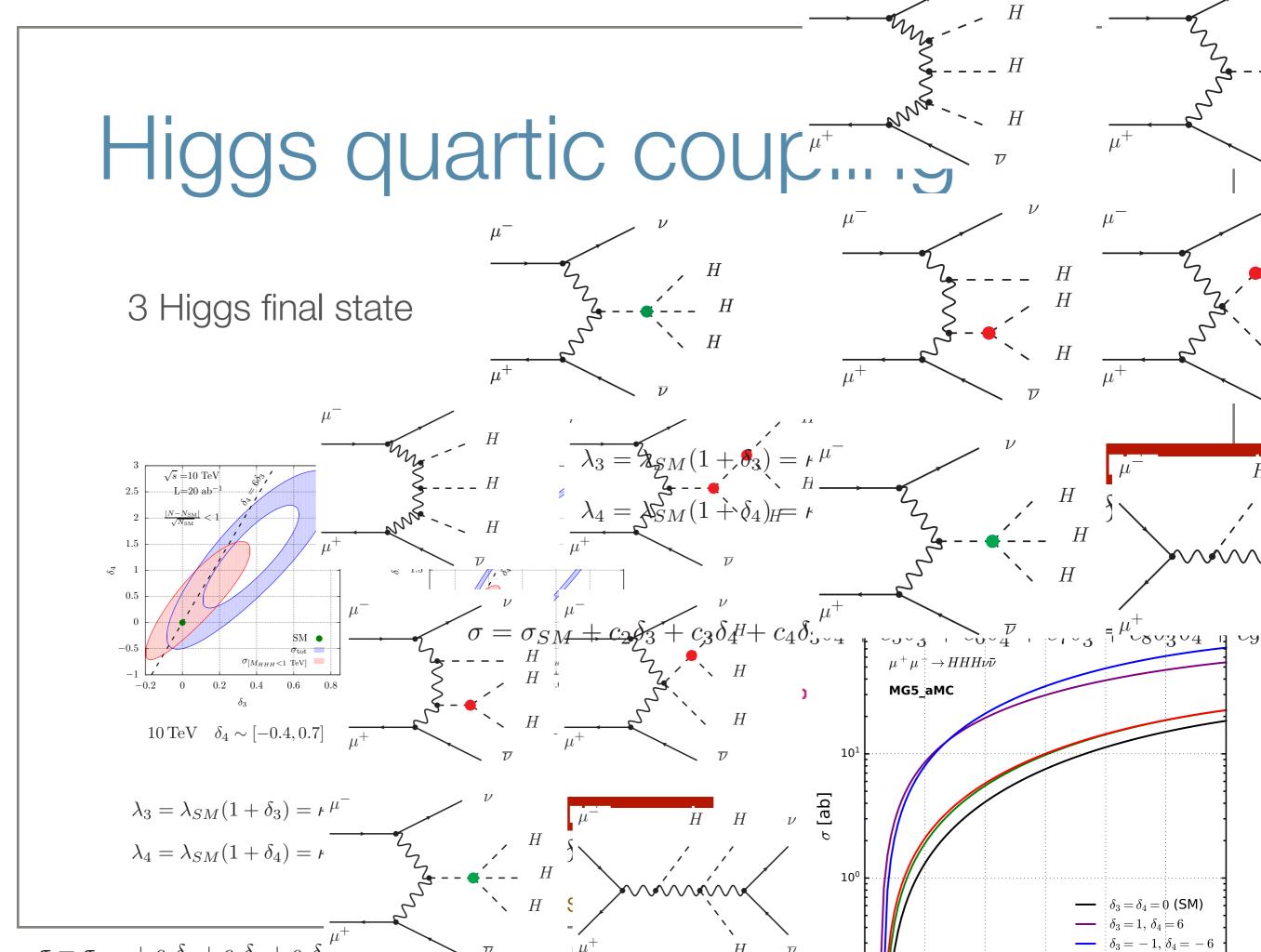
10⁻¹

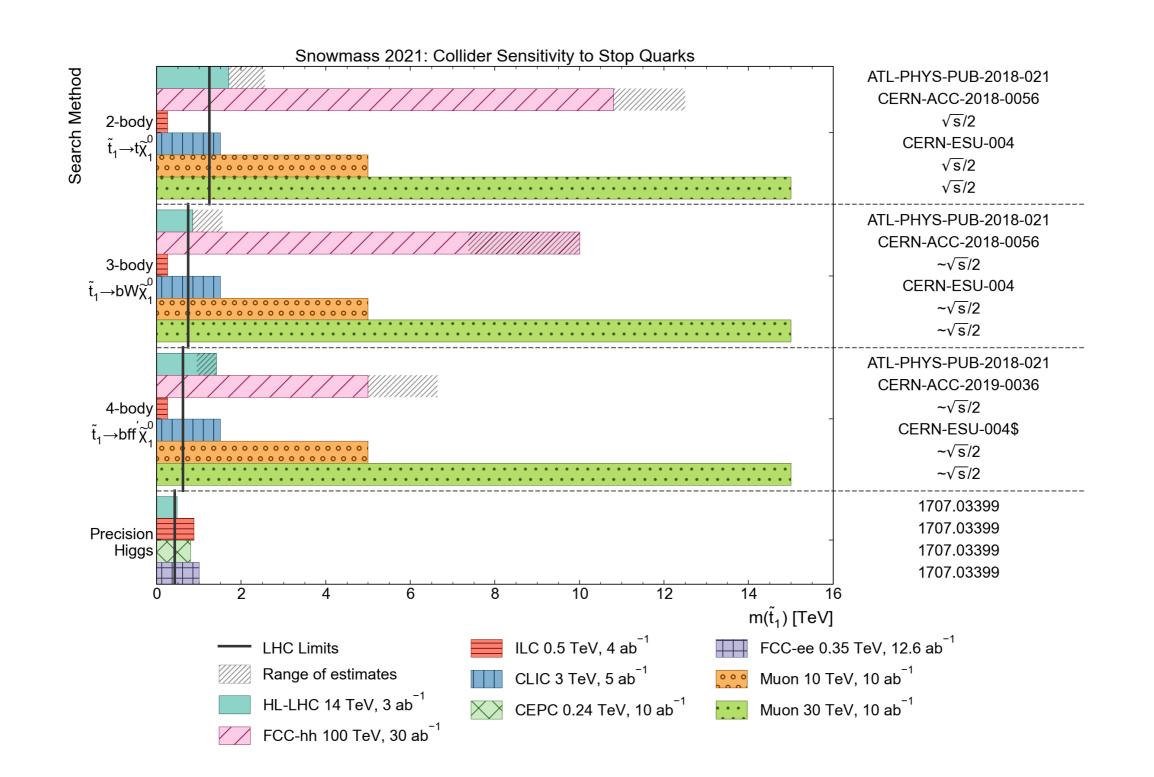
 $m_{\chi}(\text{TeV})$

50

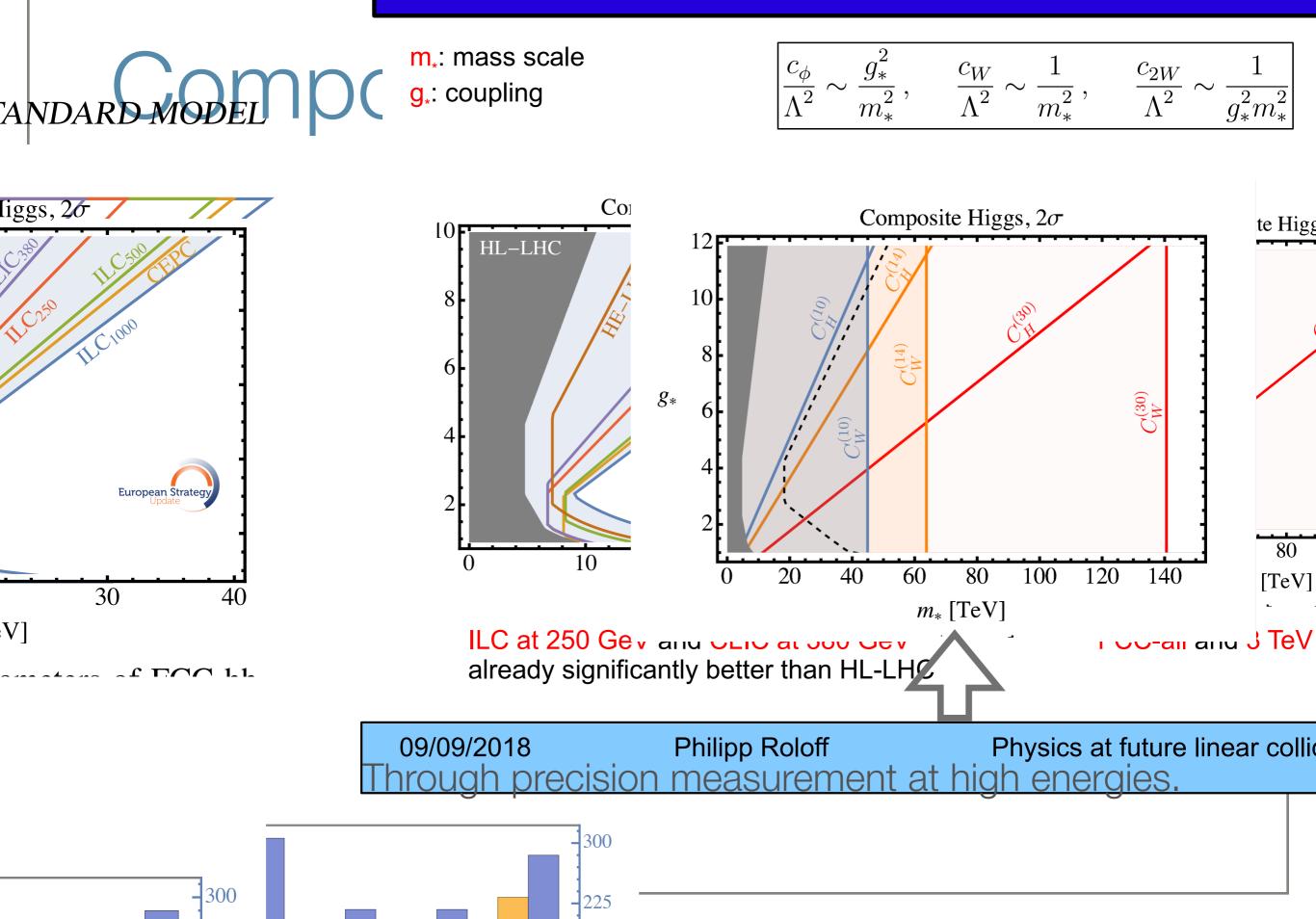
With inclusive signal: $E_{CM} \approx 14$ TeV enough to cover n≤3 multiplets. Higher energy needed to cover higher multiplets.

If we have disappearing track: potential to reach almost $m_{\chi} \approx 1/2 E_{CM}$





Composite Higgs



Blind spot for muon collider?

* New physics only charged under color.

- * e.g. Gluino
- * Muon blind flavor specific coupling.