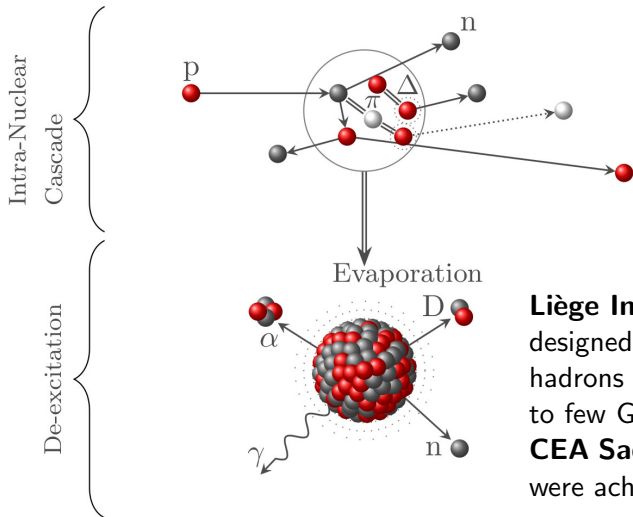


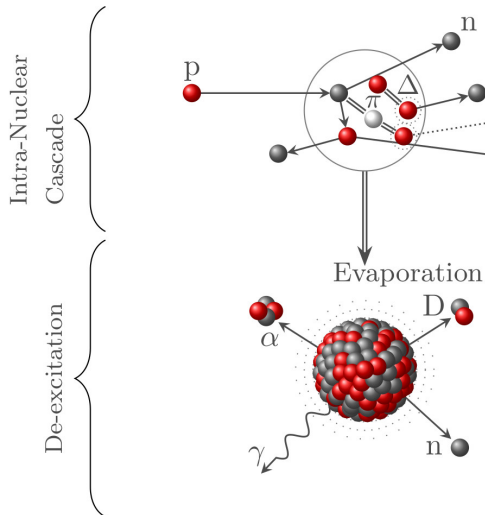
# Role of deexcitation in the final-state interactions of protons in neutrino-nucleus interactions

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**Liège Intra Nuclear Cascade:** is a code initially designed to simulate nuclear reactions induced by hadrons and light nuclei in the few tens of MeV to few GeV range. Its development is ongoing at **CEA Saclay**, and the results presented in this talk were achieved there.



**Projectiles:** baryons (nucleons,  $\Lambda$ ,  $\Sigma$ ), mesons (pions and Kaons) or light nuclei ( $A \leq 18$ ). **No neutrinos** yet! We use neutrino vertex from **NuWro**

**Versatile tool:** has been implemented in GEANT4, GENIE, and other

**De-excitation:** ABLA, SMM, GEMINI

We will use **ABLA**, since it is proved to work for the **light nuclei** (Phys. J. Plus 130, 153 (2015))

**Neutrino simulation results:**

PRD 106, 3 (2022), PRD 108, 11 (2023)

## Potential

Each nucleon in the nucleus has its **position and momentum** and moves **freely** in a square potential well. Nuclear model is essentially **classical**, with some additional ingredients to mimic quantum effects.

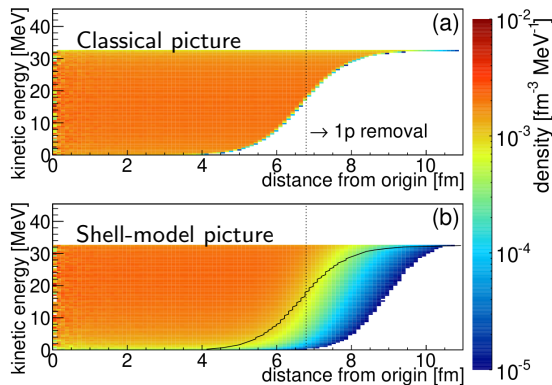
## Pauli Blocking

- **strict**: blocked is  $p < p_{Fermi}$
- **statistical**: count only nearby nucleons
- strict for the first event and statistical for the subsequent ones

## Events inside cascade

- decay/collision
- reflection/transmission with probability to **leave the nucleus as a nuclear cluster**

Space-kinetic-energy density of protons in  $^{208}\text{Pb}$



Phys.Rev.C 91, 034602 (2021)

## Potential

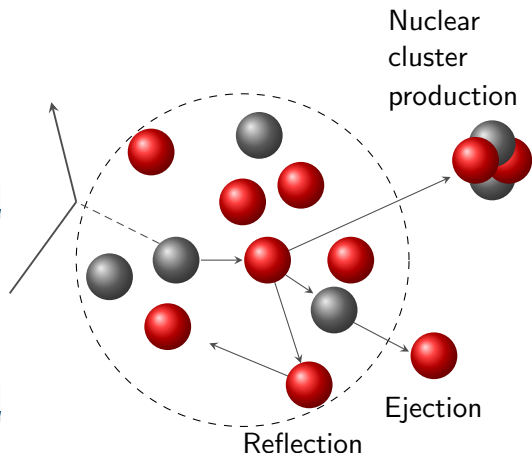
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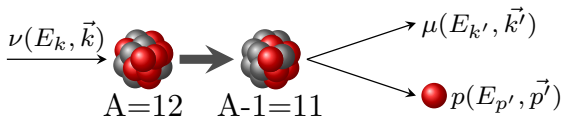
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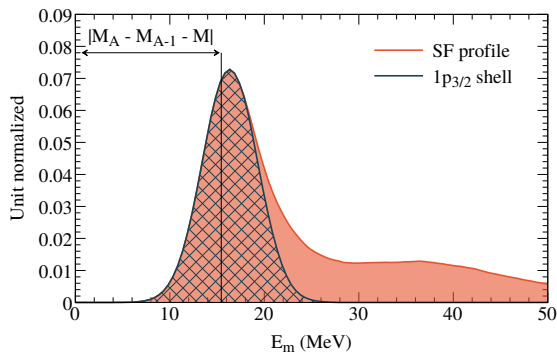


Experimental definition:

$$E_x^{\text{exp}} = E_{\text{missing}} - \underbrace{(M_A - M_{A-1} - M)}_{\text{constant}}$$

- A constant shift of missing energy by  $\sim 15.4$  MeV leads to **non-physical, negative values**
- We use experimental data (J. Phys. G: Nucl. Part. Phys. 16 507 (1999)) to simulate discrete levels

$M_{A-1}$  is the rest mass of the  $A - 1$  nucleus  
 $M_A$  is the rest mass of the initial  $A$  nucleus  
 $M$  is the rest mass of the target nucleon  
 $E_{\text{missing}}$  is the missing energy  
 For interaction on carbon,  
 $M_A - M_{A-1} - M = 15.4$  MeV



We assume all strength below the peak comes from the symmetric  $1p_{3/2}$  shell

Experimental data used for the carbon spectral function

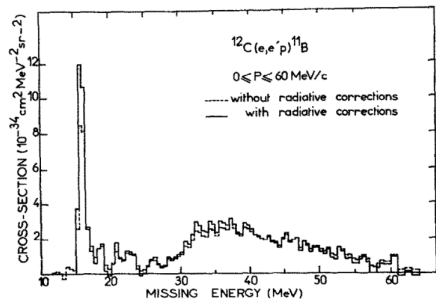


Fig. 7. Energy spectrum of the  $^{12}\text{C}(e, e'p)^{11}\text{B}$  reaction before and after the radiative corrections.

Experimental data used to simulate the discrete levels

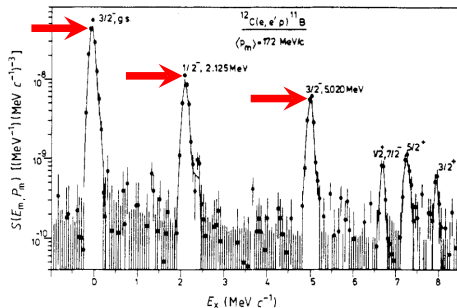
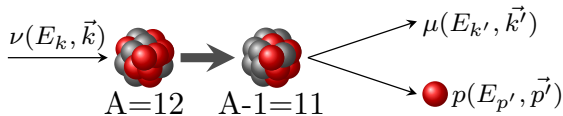


Figure 22. Excitation-energy spectrum of  $^{11}\text{B}$  observed in the reaction  $^{12}\text{C}(e, e'p)$ . Both negative and positive-parity final states are shown.

Nuclear Physics A262 (1976) 461-492

J.Phys.G: Nucl.Part.Phys. 16 507 (1990)



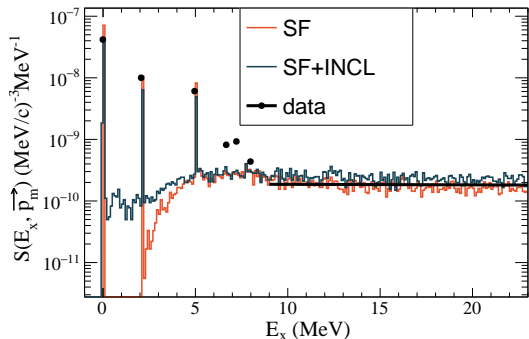
$M_R^*$  is the mass of the excited remnant  
 $M_R$  is the rest mass of the remnant  
 $T_R$  is the kinetic energy of the excited remnant  
 $p_m$  is the missing momentum

For the continuous spectrum part, we can calculate excitation energy as:

$$E_x = M_R^* - M_R, \text{ where:}$$

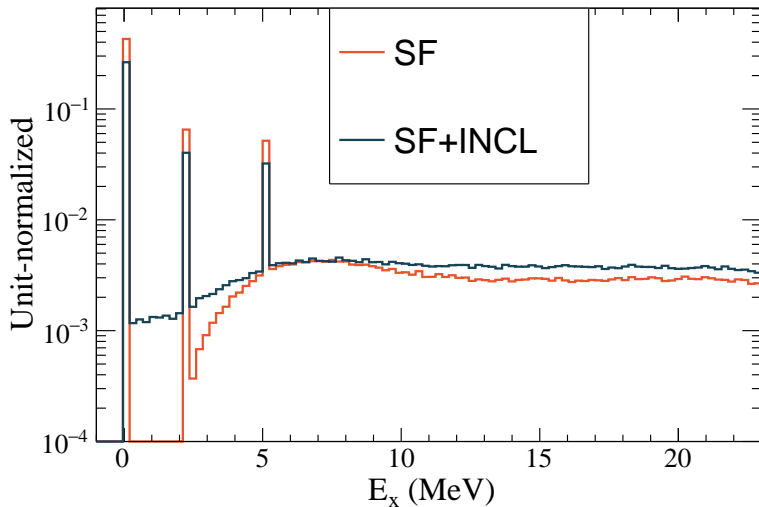
$$M_R^* = \sqrt{(E_k + M_A - E_{k'} - E_{p'})^2 - |\vec{p}_m|^2}$$

Otherwise, we model **3 discrete peaks** with strength of 79%, 12%, and 9% (**p-shell**)



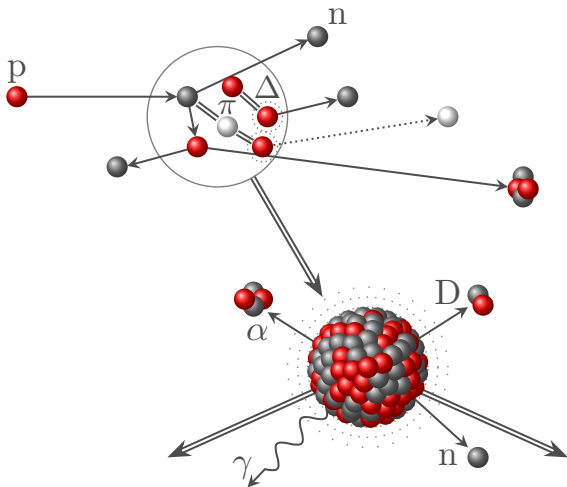
Excitation energy of the nuclear remnant after neutrino interactions for  $169.5 < p_m < 174.5$  MeV





Final distribution of the excitation energy

We use INCL+ABLA to simulate **neutrino CCQE** events with the **T2K** energy spectrum. We compare the obtained results to **NuWro**.

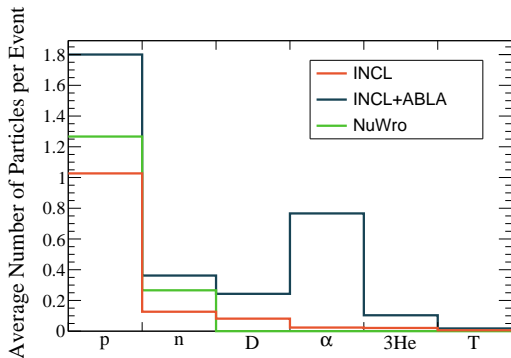


In the paper [Phys.Rev.D 106, 3 \(2022\)](#) we show the **nuclear cluster production for the first time** in FSI.

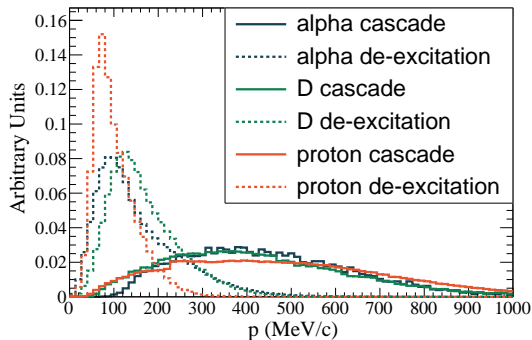
In the paper [Phys.Rev.D 108, 11 \(2023\)](#) we study the impact of the subsequent **de-excitation modelling**, that predicts **more nuclear clusters**.

**Important!** Gamma production during de-excitation is just a **small fraction** of possible scenarios

ABLA features a **massive production** of particles with **low momentum**.

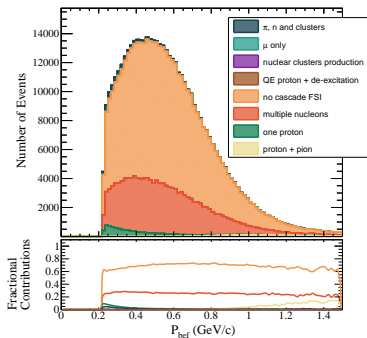


Average number of particles per event produced by INCL, INCL+ABLA, and NuWro.

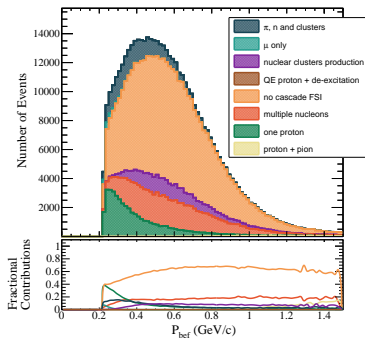


Momentum of nuclear clusters produced during the cascade and de-excitation

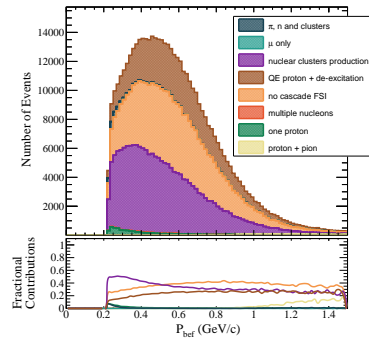
- Large fraction of "no FSI" events (i.e. proton untouched) is now **feature production of other particles** (and nuclear clusters) in the final state due to de-excitation
- Events with **only nucleon production** now feature **nuclear cluster production**



NuWro

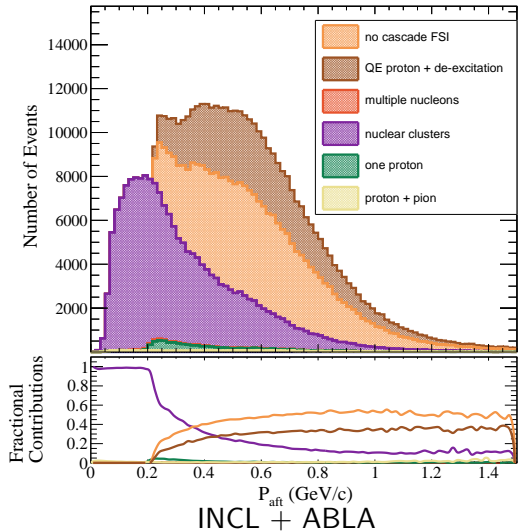
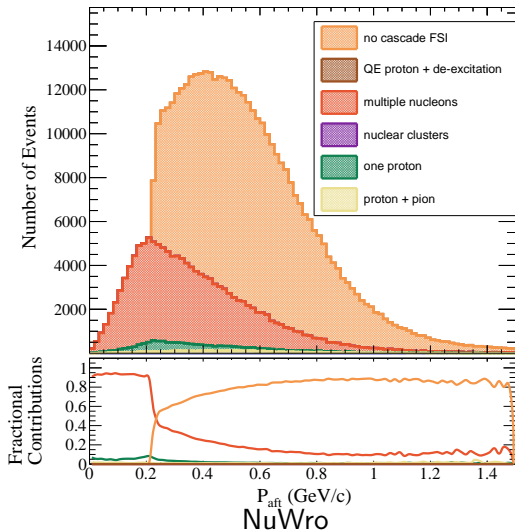


INCL



INCL + ABLA

INCL+ABLA simulation features **massive difference** in nucleon kinematics in comparison to NuWro

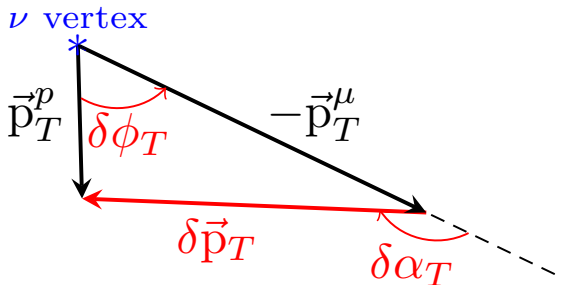


We use **Single Transverse Variables (STV)** that allow to disentangle different effects for better FSI estimation. STV are **observable** and **measurable**.




sensitive to FSI: 
$$\delta\alpha_T = \arccos \frac{-\vec{k}'_T \cdot \delta\vec{p}'_T}{k'_T \cdot \delta p'_T}$$




sensitive to Fermi Motion:

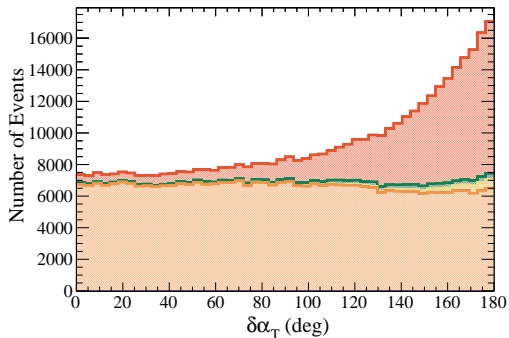
$$\delta\vec{p}_T = \vec{p}_T^{\vec{p}} + \vec{p}_T^{\vec{\mu}} = \vec{p}_T^{\vec{n}}$$



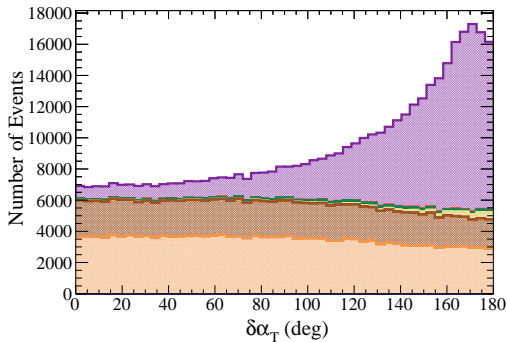
High  $\delta\alpha_T$  strongly depends on FSI and is affected by de-excitation and Pauli blocking

-  no cascade FSI
-  QE proton + de-excitation
-  multiple nucleons

-  nuclear clusters production
-  one proton
-  proton + pion



NuWro

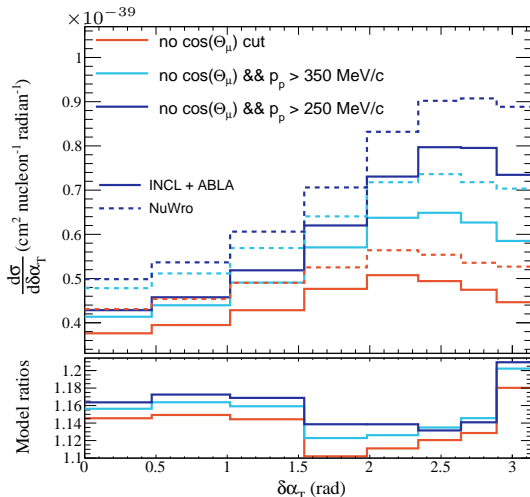
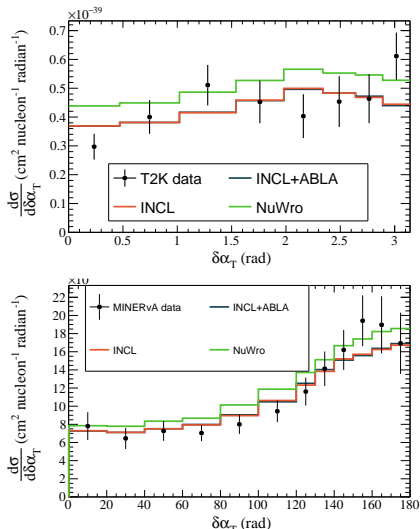


INCL+ABLA

Current detector threshold in ND280 and MINERvA scintillators is **too large**, so we **cannot see the difference** between INCL and

Lower threshold provides better sensitivity to distinguish models

NuWro





Using  $\mu + p$  is **better** than using muon only, but here we show that we gain even **higher precision** by using all subleading particles

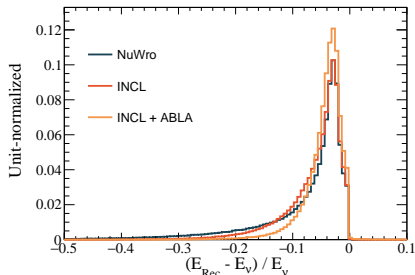
proton only:

$$E_{rec} = E_{\mu} + T_p$$

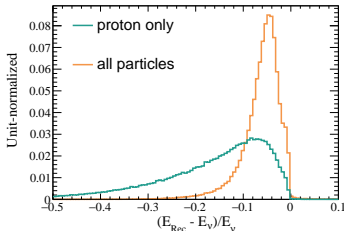


all particles (including clusters)

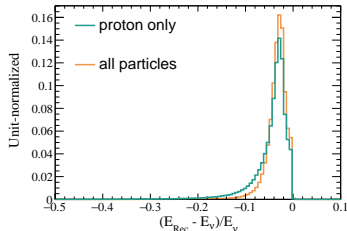
$$E_{rec} = E_{\mu} + \sum_i T_i$$



"all particles" reconstruction

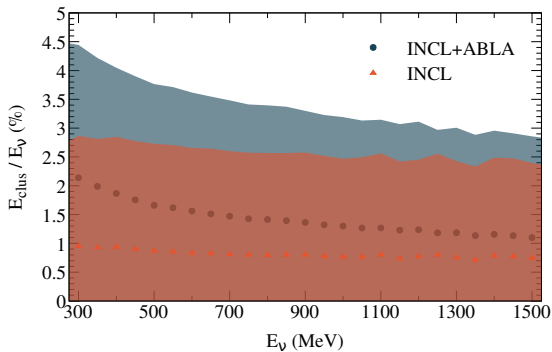


INCL+ABLA cascade FSI

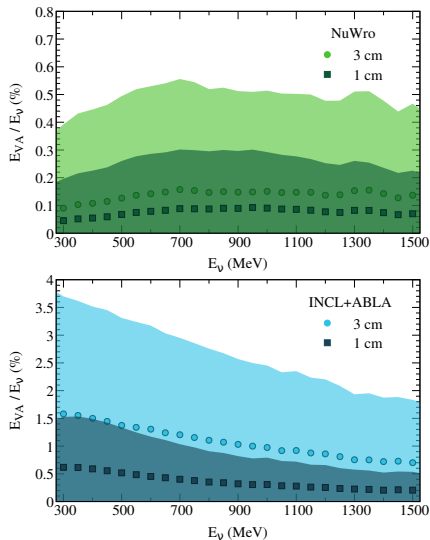


INCL+ABLA no cascade FSI

The **actual fraction** of neutrino energy going to the kinetic energy of the subleading hadrons is **non-negligible**.



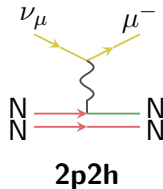
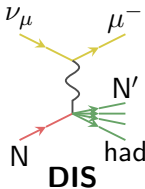
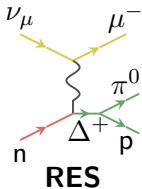
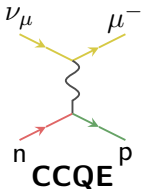
What can be actually seen in the detector (Birks quenching applied):



- the effort is ongoing to integrate **CCQE** reaction in INCL with subsequent implementation in **Geant4**
- "transparent events" are **no always** transparent: nuclear clusters may be produced
- INCL+ABLA simulation features **important difference** in nucleon kinematics in comparison to NuWro (and the other similar generator used in neutrino scattering)
- An essential novelty of this study is the **simulation of nuclear cluster production** during cascade and de-excitation. It is important for the understanding of the **vertex activity** and calorimetric method of  $\nu$  **energy reconstruction**

For **precise neutrino energy reconstruction** (e.g. "calorimetric method") is important to include **vertex activity** ( $\sim 1 - 2\%$ ), and to have proper model of it to correct for detector quenching. Large portion of VA comes from the **de-excitation**.

# BACK UP

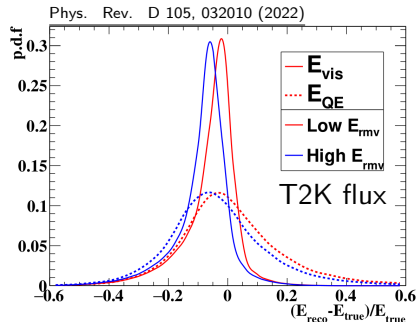


Energy reconstruction using only muon kinematics  
(works well for **quasi-elastic reaction**):

$$E_{\nu}^{QE} = \frac{m_p^2 - (m_n - E_B)^2 - m_{\mu}^2 + 2(m_n - E_B)E_{\mu}}{2((m_n - E_B) - E_{\mu} + p_{\mu} \cos \theta_{\mu})}$$

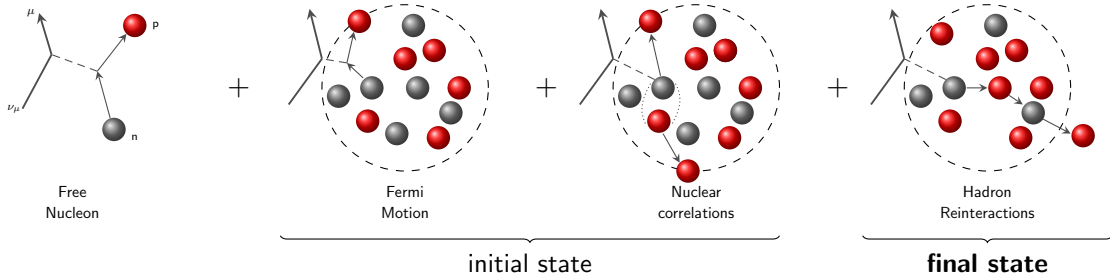
Energy reconstruction using **muon and kinetic energy of the nucleon**:

$$E_{\nu}^{vis} = E_{\mu} + T_N$$

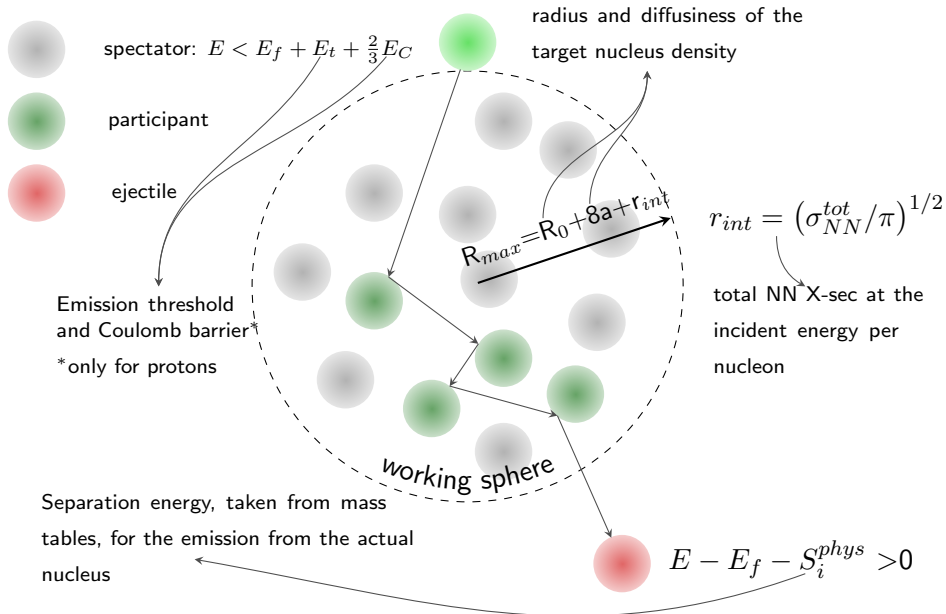


$E_{\nu}^{vis}$ , dashed line — QE formula  
solid line —  $\mu + N$  formula

$\mu + N$  formula gives us more **opportunities**, but also it creates more **challenges** for modelling and we need to **understand better nuclear effects** also on neutrons and protons.

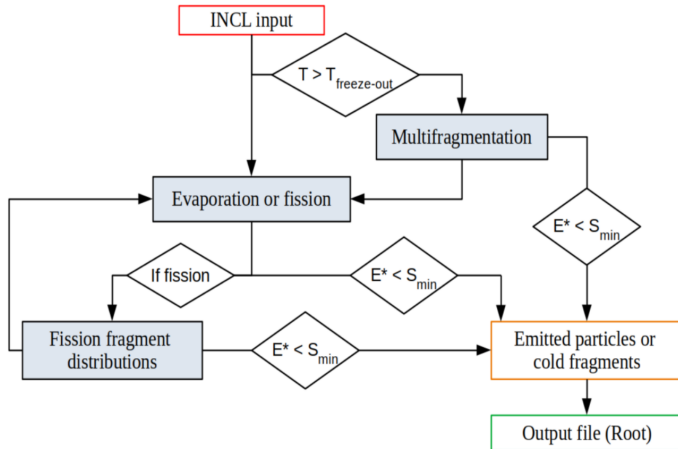


We will focus on **CCQE**  $\nu$  reaction channel and the **Final State Interactions (FSI)** that are described by **cascade models** and on the nuclear **excitation energy**.



The ablation model **ABLA** describes the **de-excitation** of an excited nuclear system through the emission of  $\gamma$ -rays, neutrons, light-charged particles, and intermediate-mass fragments (or fission in case of hot and heavy remnants).

Phys. Rev. C 105, 014623 (2022)



$$T_{freeze-out} = \max \left[ 5.5, 9.33e^{(-2.82 \times 10^{-3} A_{rem})} \right]$$

$S_{min}$ —minimum particle separation energy