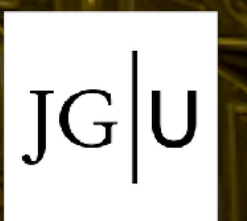
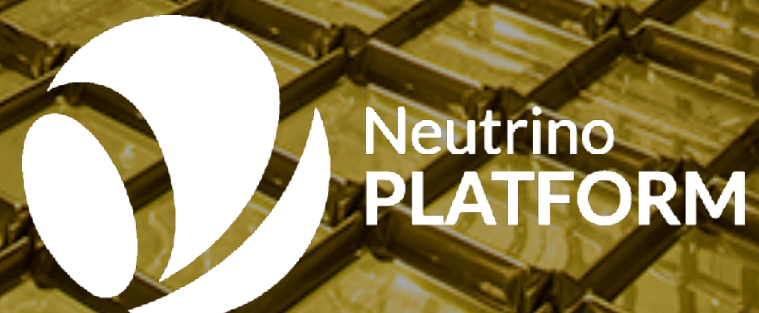


# Impact of Neutrino Interaction Uncertainties on BSM Searches

Joachim Kopp (CERN & JGU Mainz)  
NuInt 2024 • São Paulo • April 2024

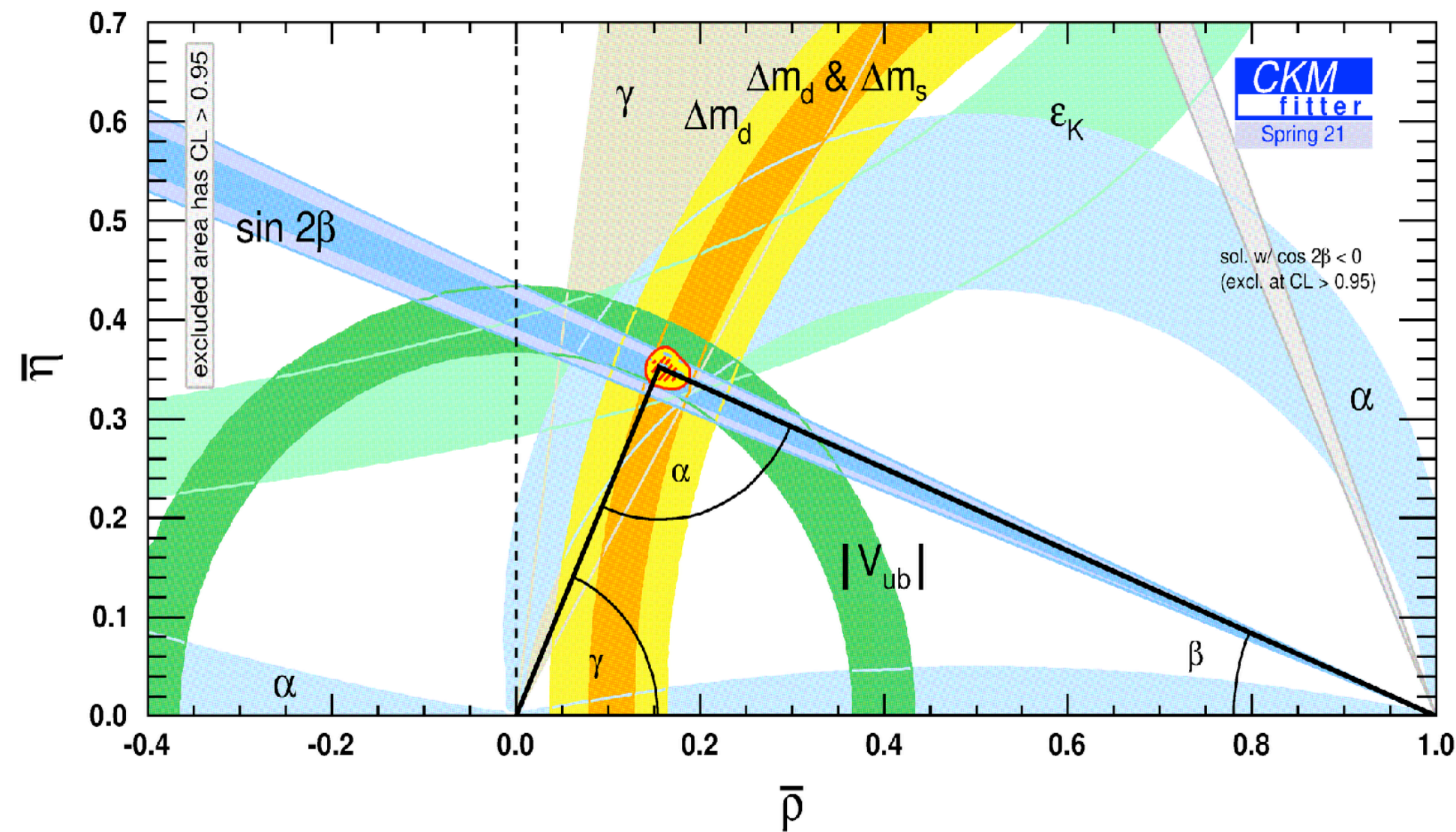


JOHANNES GUTENBERG  
UNIVERSITÄT MAINZ

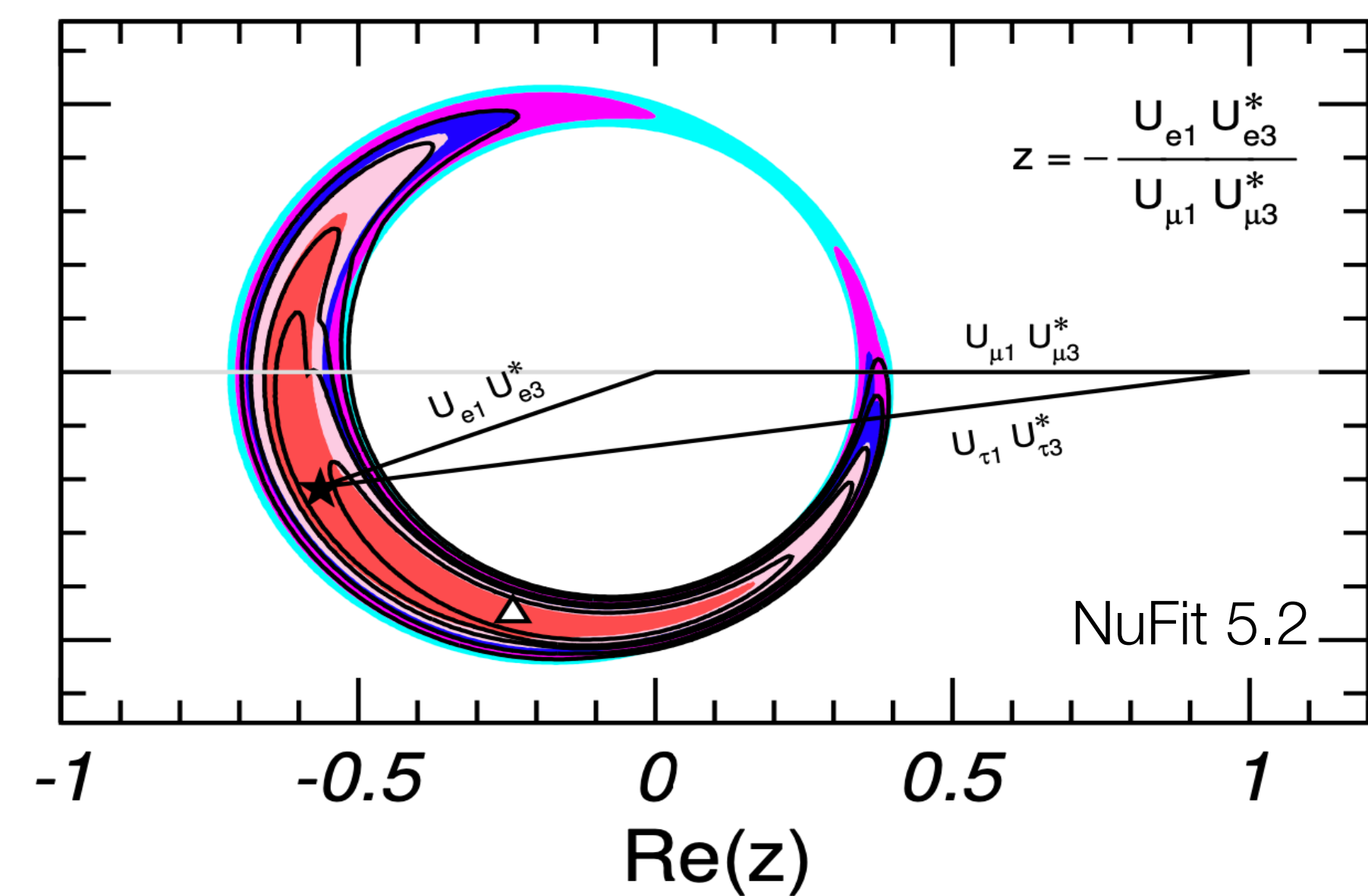


# Precision Neutrino Physics

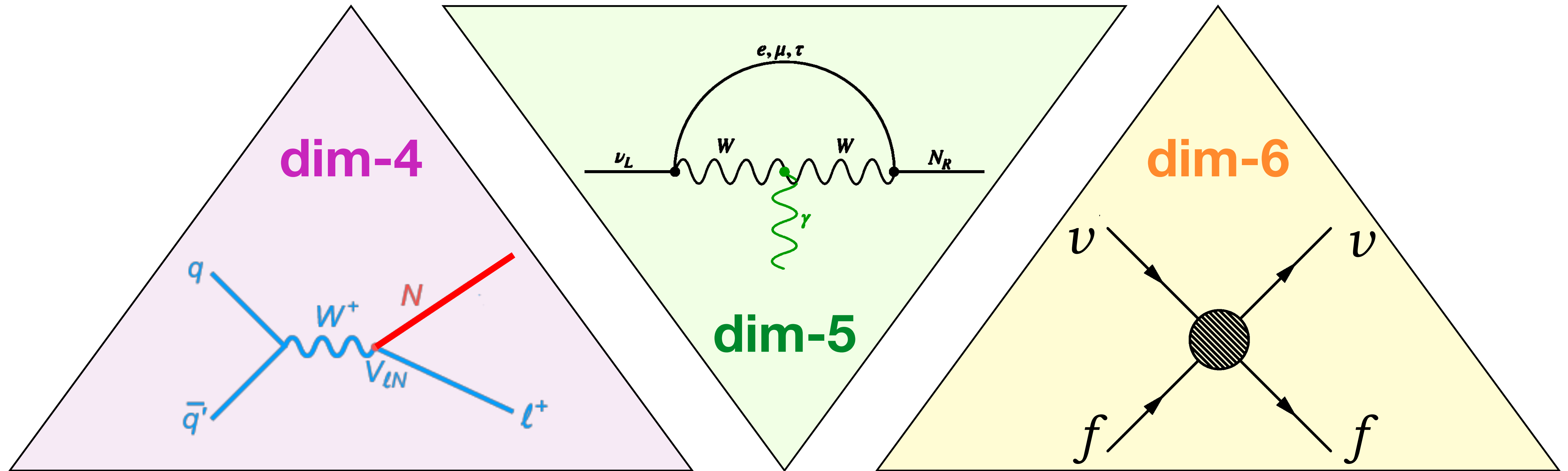
## Quarks



## Leptons

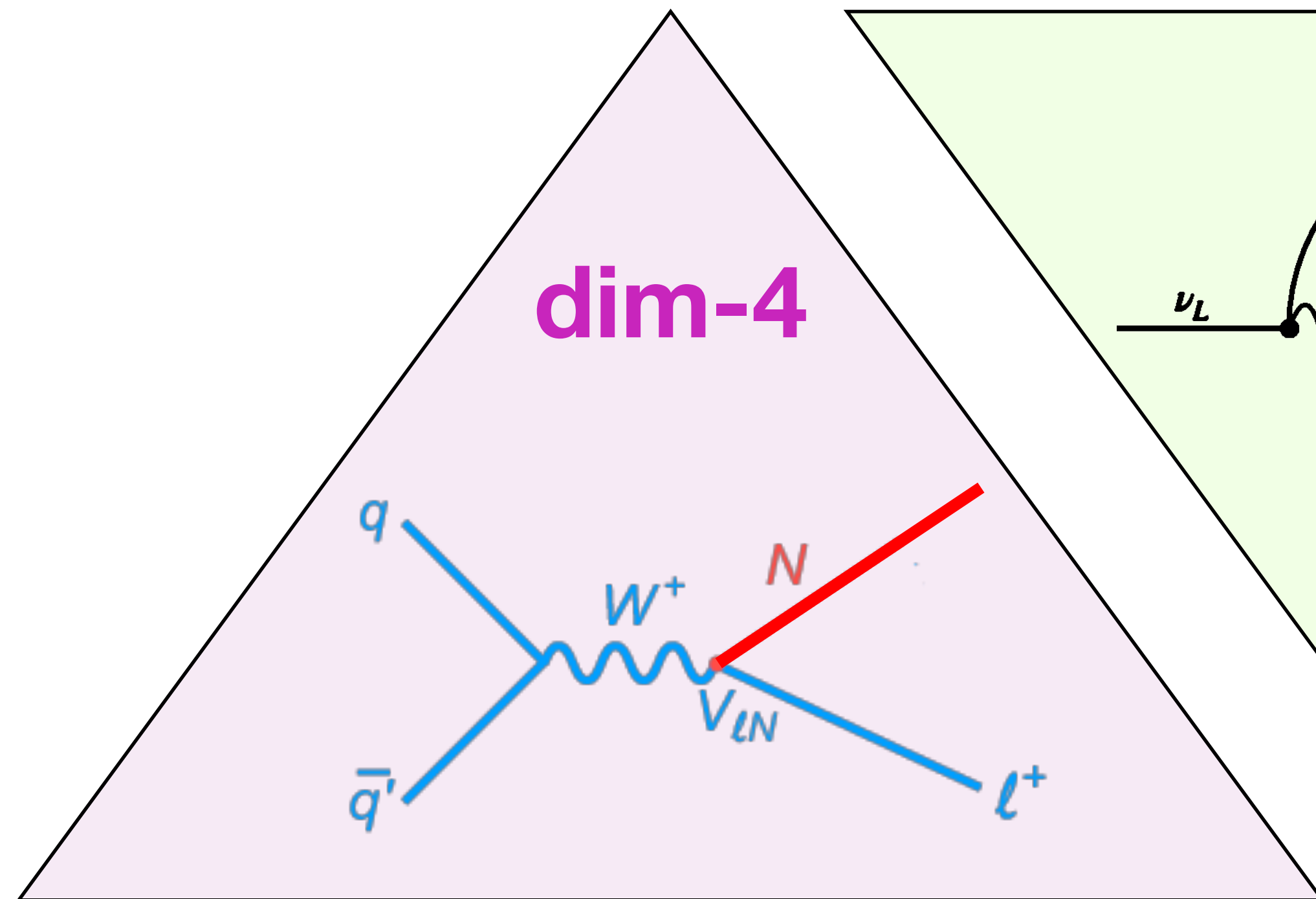


# Neutrino Physics Beyond the Standard Model

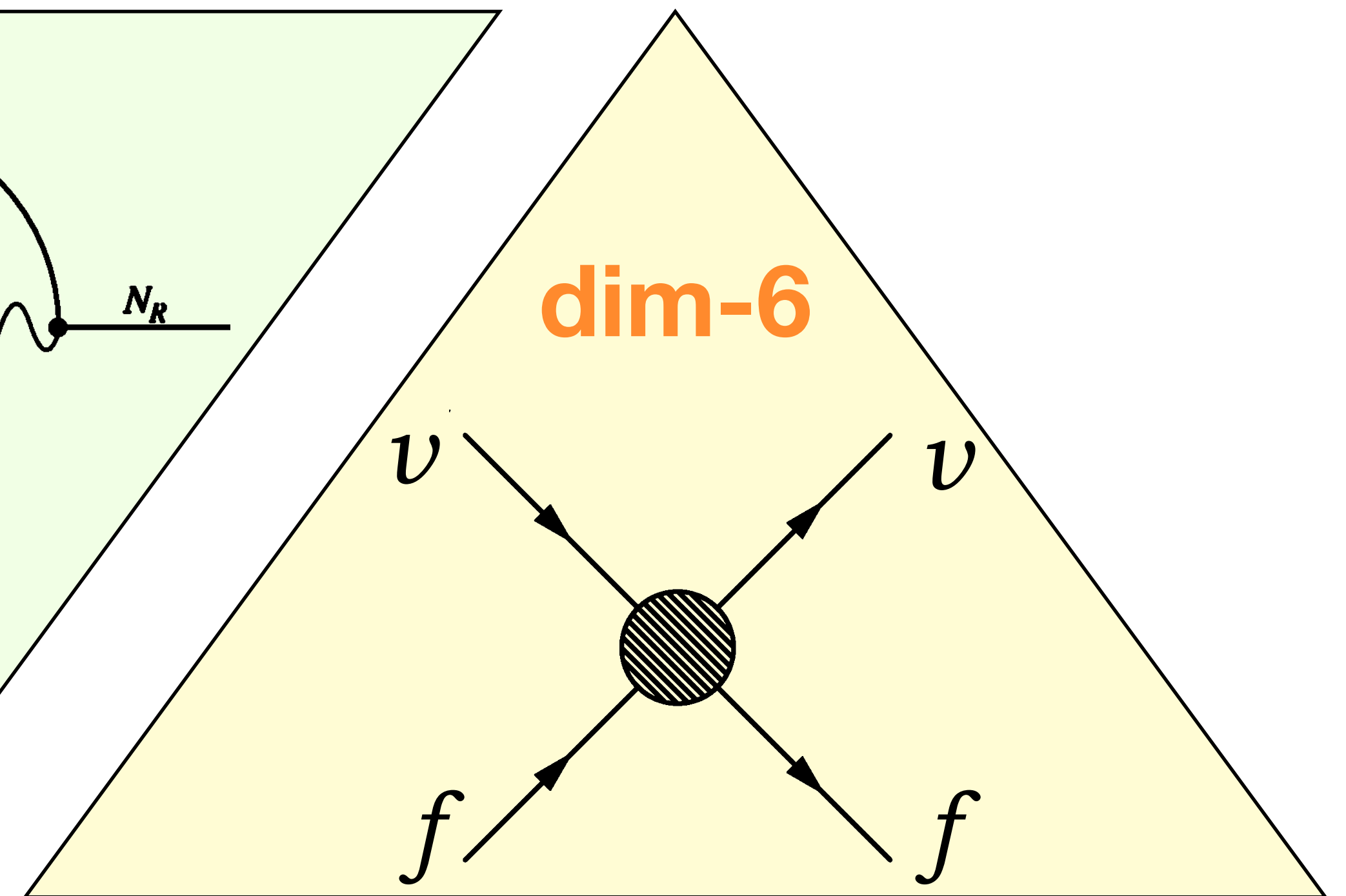
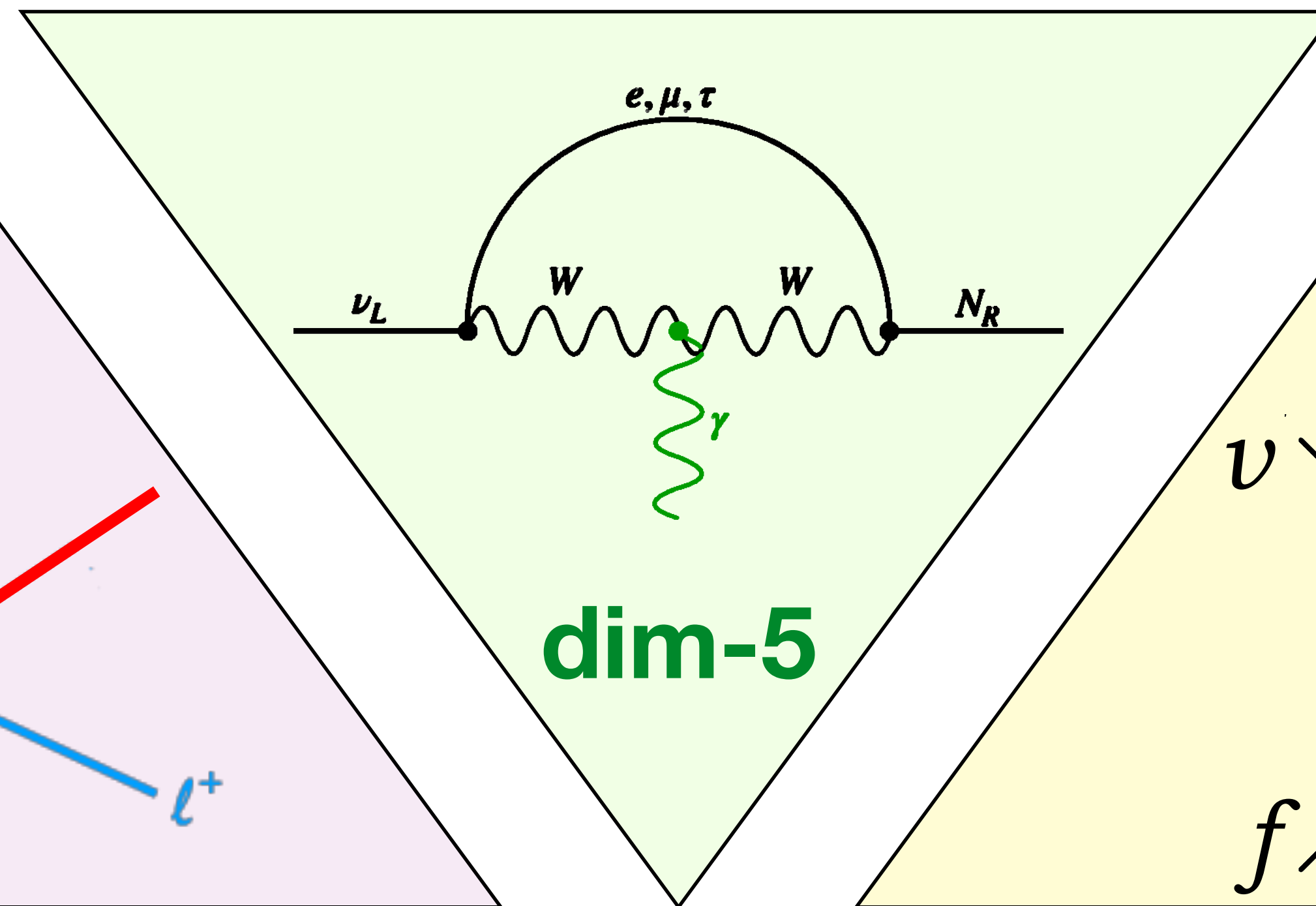


# Neutrino Physics Beyond the Standard Model

e.g. neutrino magnetic moments



e.g. sterile neutrinos



e.g. non-standard interactions

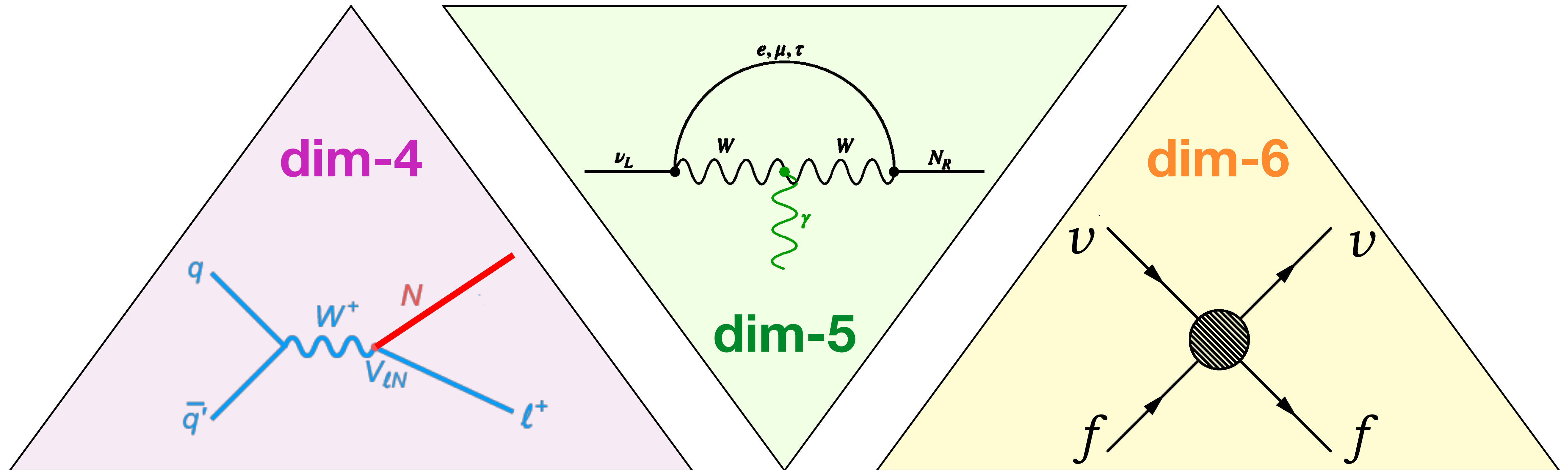


“With great precision comes great responsibility.”

*Tim Linden, WIN 2021*

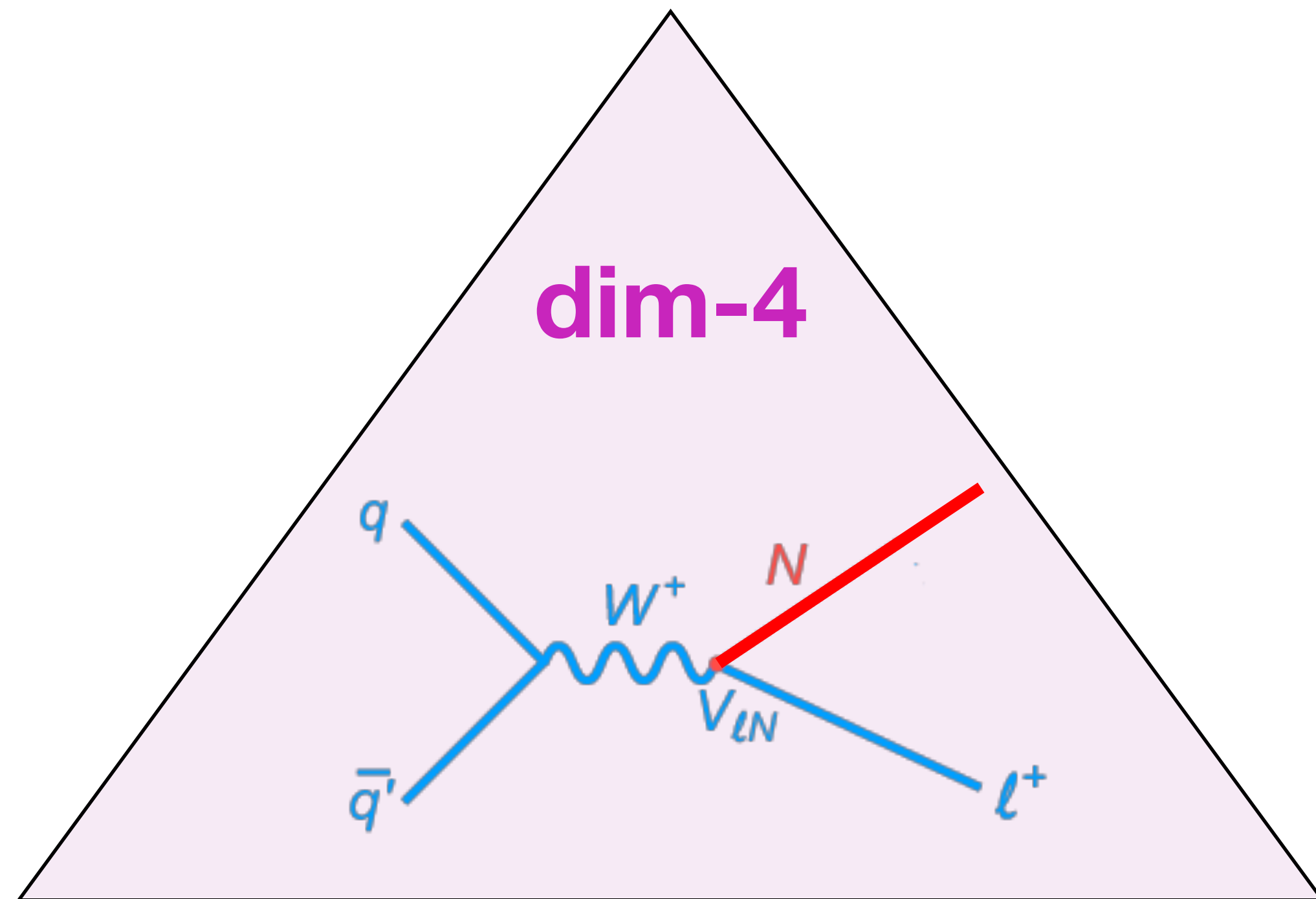


# Neutrino Physics Beyond the Standard Model



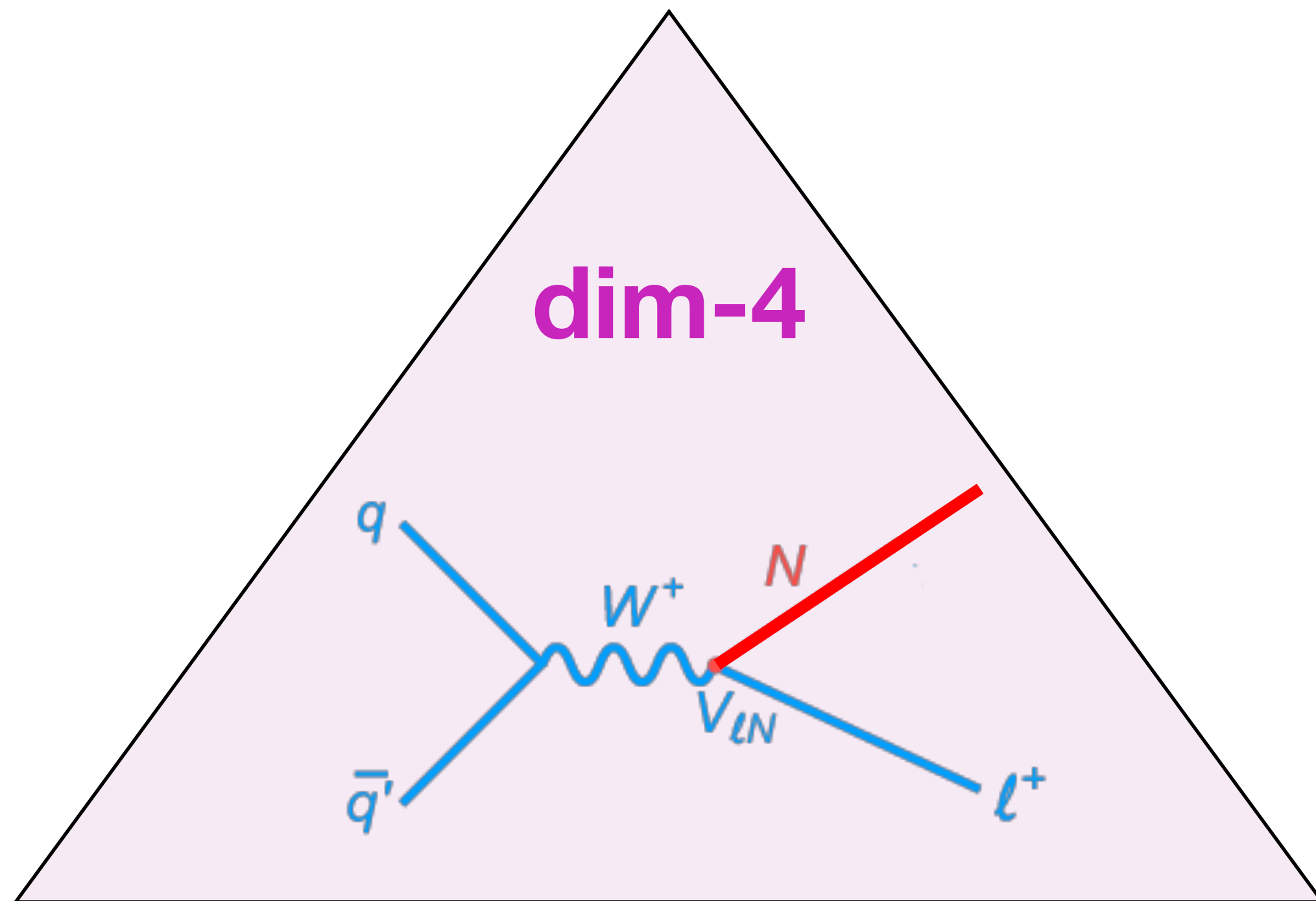


# Sterile Neutrinos = new, uncharged fermions



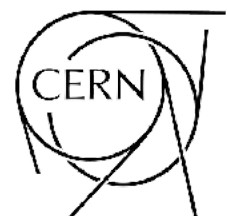


# Sterile Neutrinos = new, uncharged fermions



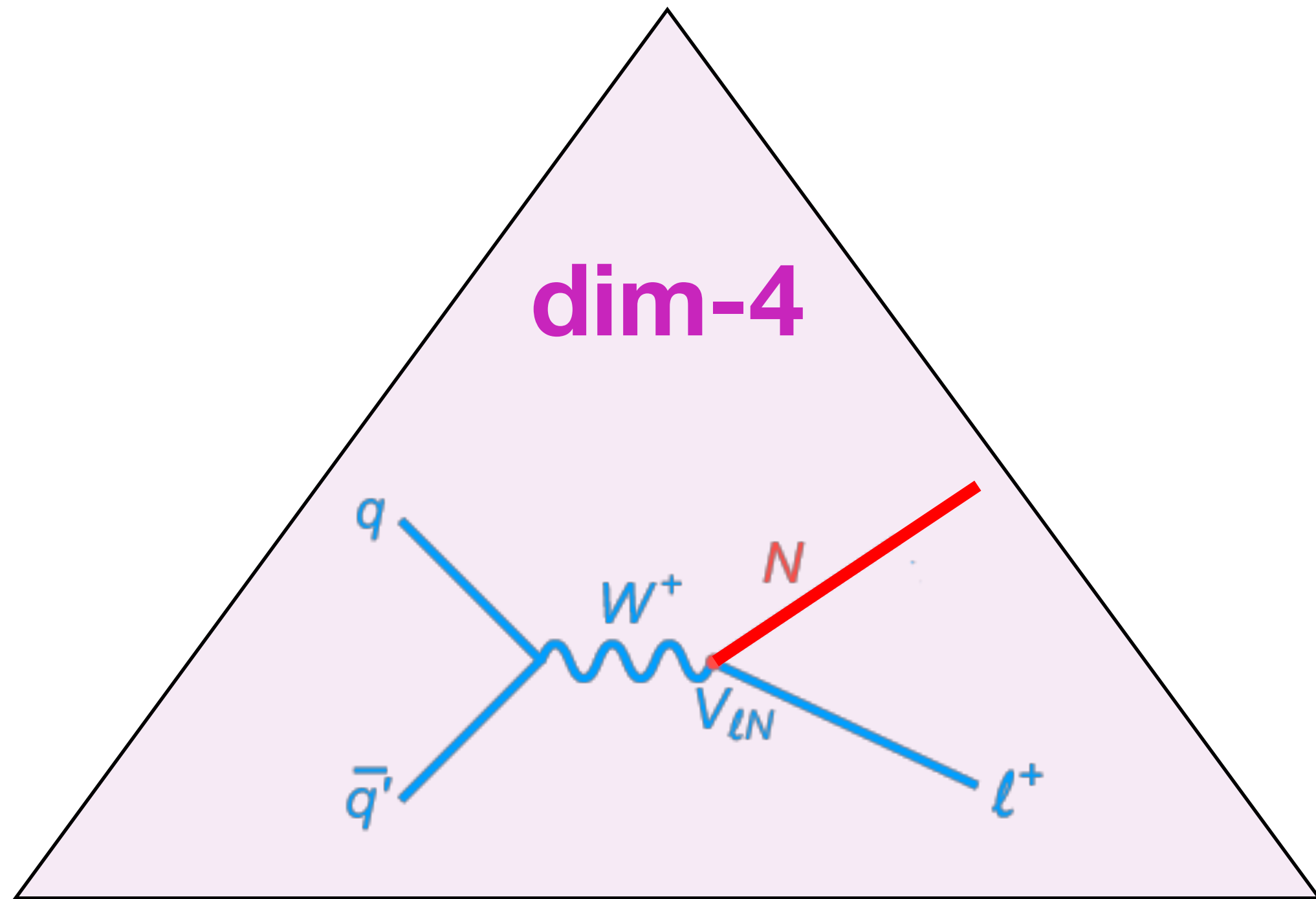
## Standard Model of Elementary Particles

three generations of matter (fermions)						
	I	II	III			
mass	$\approx 2.4 \text{ MeV}/c^2$	$\approx 1.275 \text{ GeV}/c^2$	$\approx 172.44 \text{ GeV}/c^2$	0	$\approx 125.09 \text{ GeV}/c^2$	
charge	$2/3$	$2/3$	$2/3$	0	0	
spin	$1/2$	$1/2$	$1/2$	1	0	
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon	<b>H</b> Higgs	
<b>QUARKS</b>	$\approx 4.8 \text{ MeV}/c^2$ $-1/3$ $1/2$ <b>d</b> down	$\approx 95 \text{ MeV}/c^2$ $-1/3$ $1/2$ <b>s</b> strange	$\approx 4.18 \text{ GeV}/c^2$ $-1/3$ $1/2$ <b>b</b> bottom	0 0 1 <b>γ</b> photon	<b>SCALAR BOSONS</b>	
	$\approx 0.511 \text{ MeV}/c^2$ $-1$ $1/2$ <b>e</b> electron	$\approx 105.67 \text{ MeV}/c^2$ $-1$ $1/2$ <b>μ</b> muon	$\approx 1.7768 \text{ GeV}/c^2$ $-1$ $1/2$ <b>τ</b> tau	$\approx 91.19 \text{ GeV}/c^2$ 0 1 <b>Z</b> Z boson		
<b>LEPTONS</b>	$< 2.2 \text{ eV}/c^2$ 0 $1/2$ <b>ν<sub>e</sub></b> electron neutrino	$< 1.7 \text{ MeV}/c^2$ 0 $1/2$ <b>ν<sub>μ</sub></b> muon neutrino	$< 15.5 \text{ MeV}/c^2$ 0 $1/2$ <b>ν<sub>τ</sub></b> tau neutrino	$\approx 80.39 \text{ GeV}/c^2$ $\pm 1$ 1 <b>W</b> W boson		
				<b>GAUGE BOSONS</b>		





# Sterile Neutrinos = new, uncharged fermions



## Standard Model of Elementary Particles

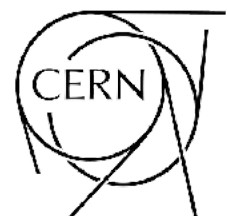
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spin	$1/2$	$1/2$	$1/2$	1		0
	<b>u</b> up	<b>c</b> charm	<b>t</b> top	<b>g</b> gluon		<b>H</b> Higgs
	<b>d</b> down	<b>s</b> strange	<b>b</b> bottom	<b><math>\gamma</math></b> photon		
	<b>e</b> electron	<b><math>\mu</math></b> muon	<b><math>\tau</math></b> tau	<b>Z</b> Z boson		
	<b><math>\nu_e</math></b> electron neutrino	<b><math>\nu_\mu</math></b> muon neutrino	<b><math>\nu_\tau</math></b> tau neutrino	<b>W</b> W boson		
	<b><math>\nu_s</math></b> sterile neutrino					

**QUARKS** (left side of the table)

**LEPTONS** (left side of the table)

**GAUGE BOSONS** (right side of the table)

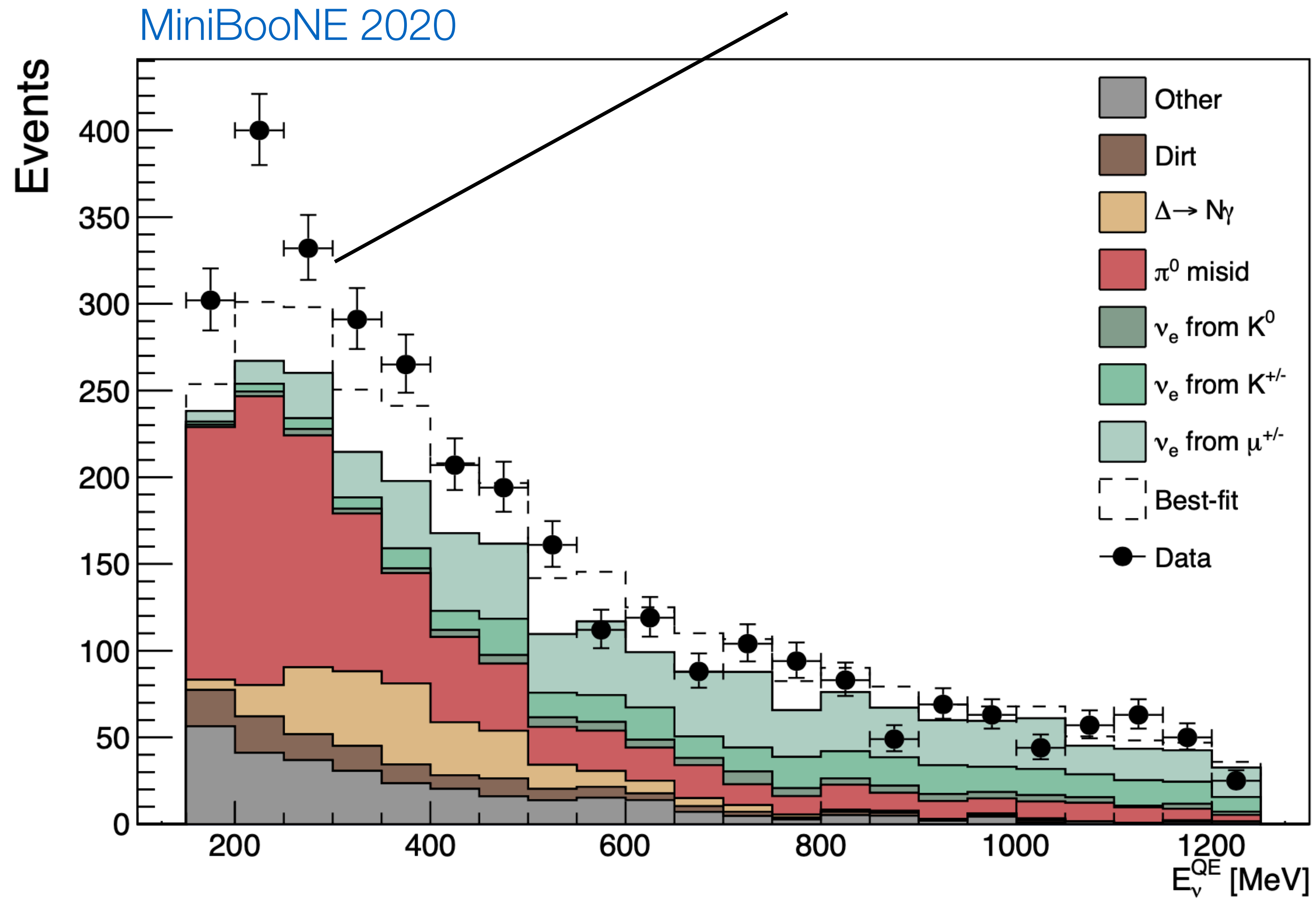
**SCALAR BOSONS** (right side of the table)





# MiniBooNE

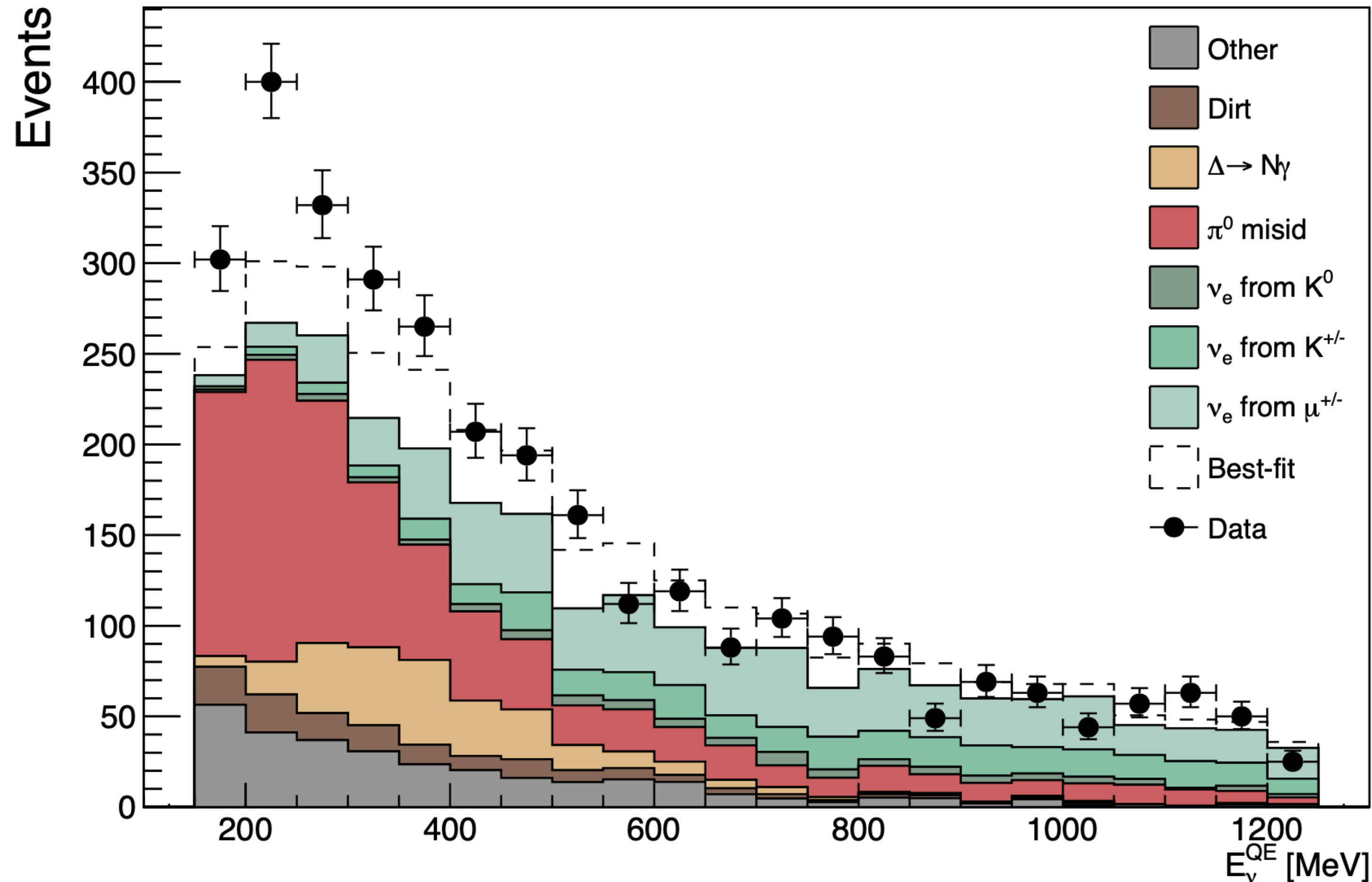
4.8  $\sigma$  excess of  $\nu_e$  in a  $\nu_\mu$  beam





# MiniBooNE

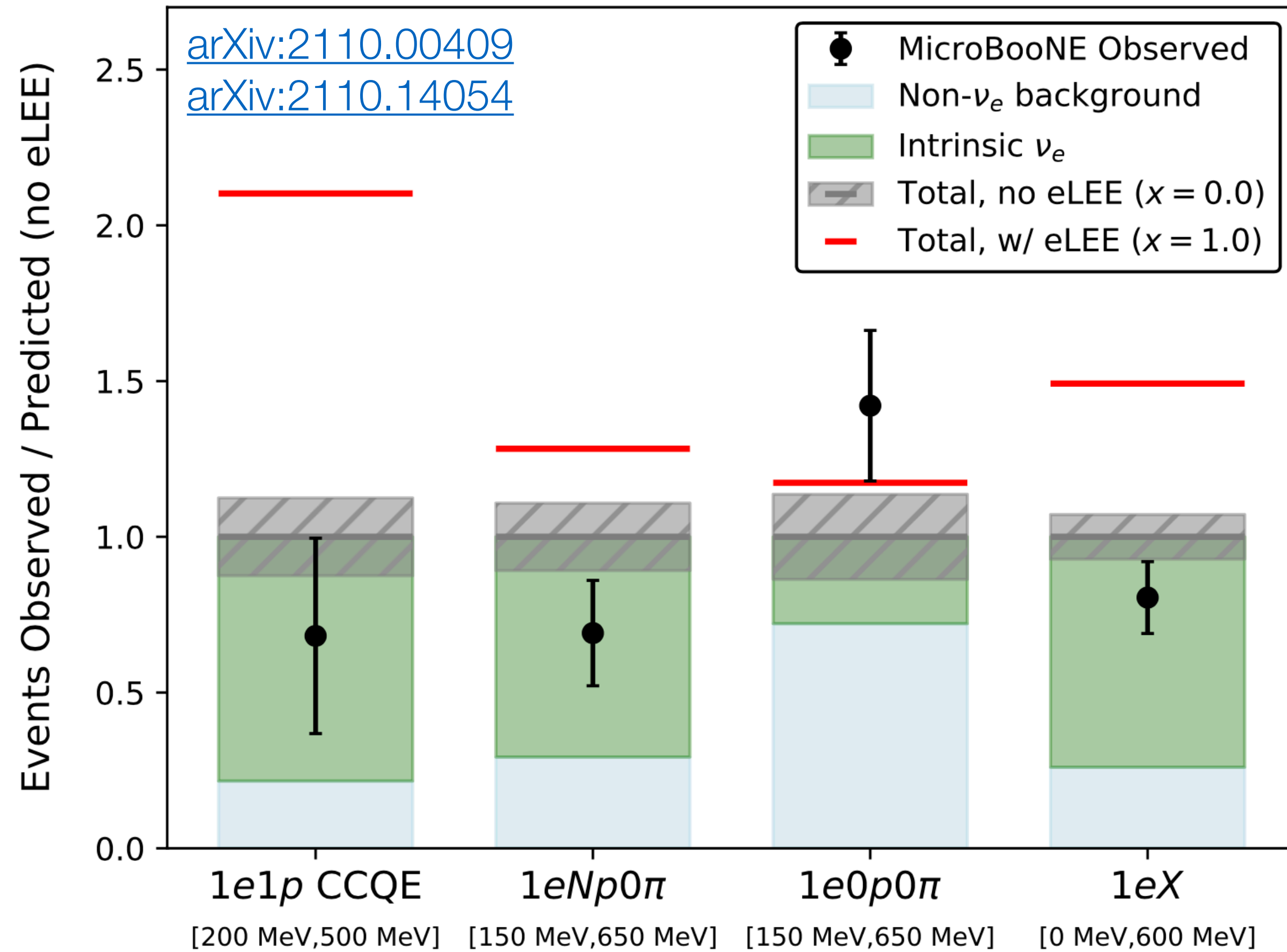
MiniBooNE 2020



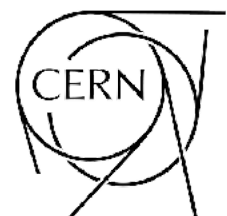
- baseline too short for std. oscillations
- but could be explained by **eV-scale sterile neutrino** (“ $\nu_{\mu} \rightleftharpoons \nu_s \rightleftharpoons \nu_e$ ”)
- $\sim$  consistent with **other anomalies**
- but inconsistent with **null searches**



# MicroBooNE

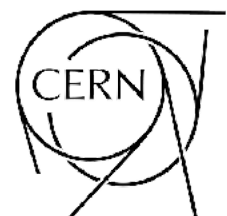
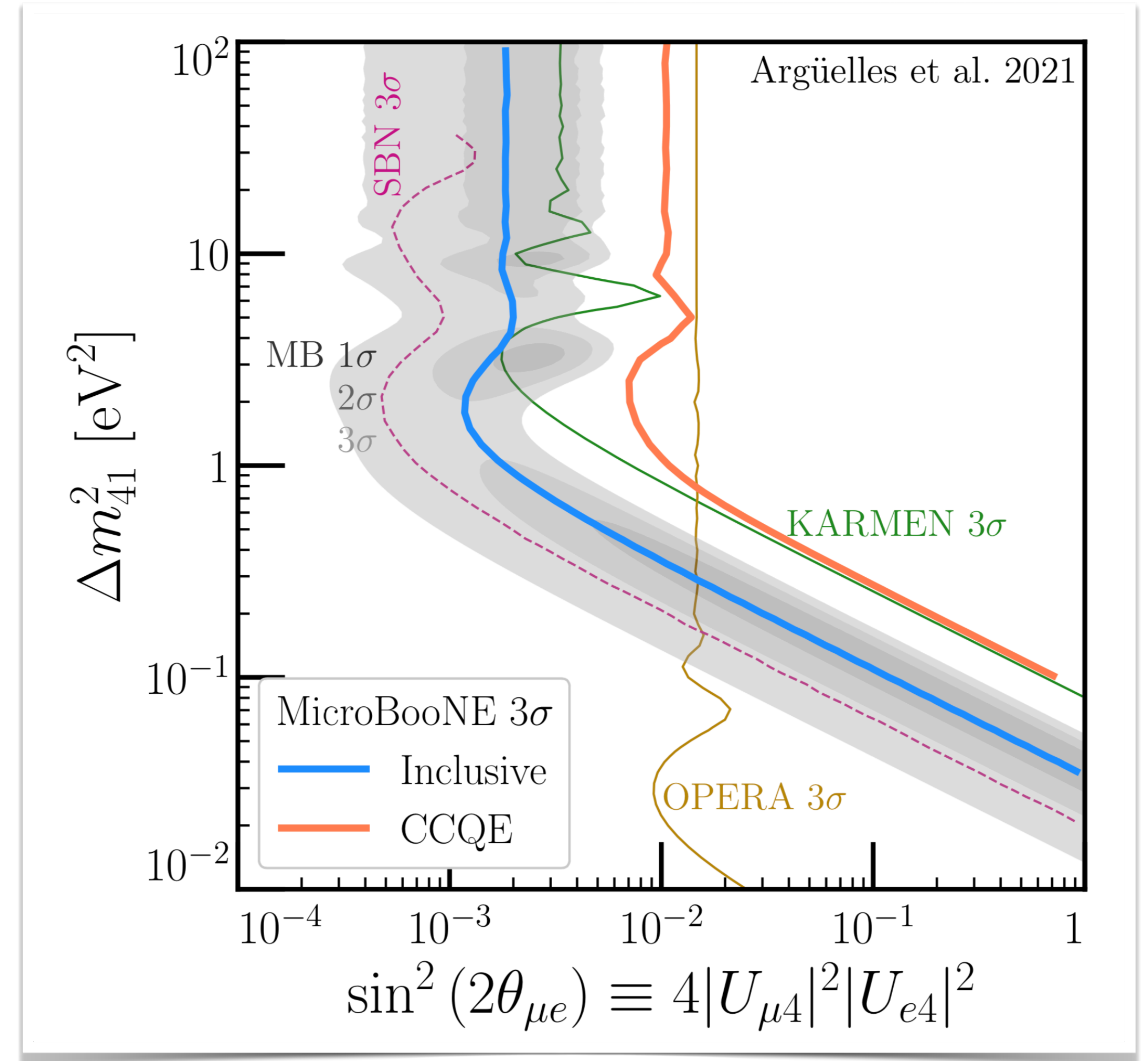
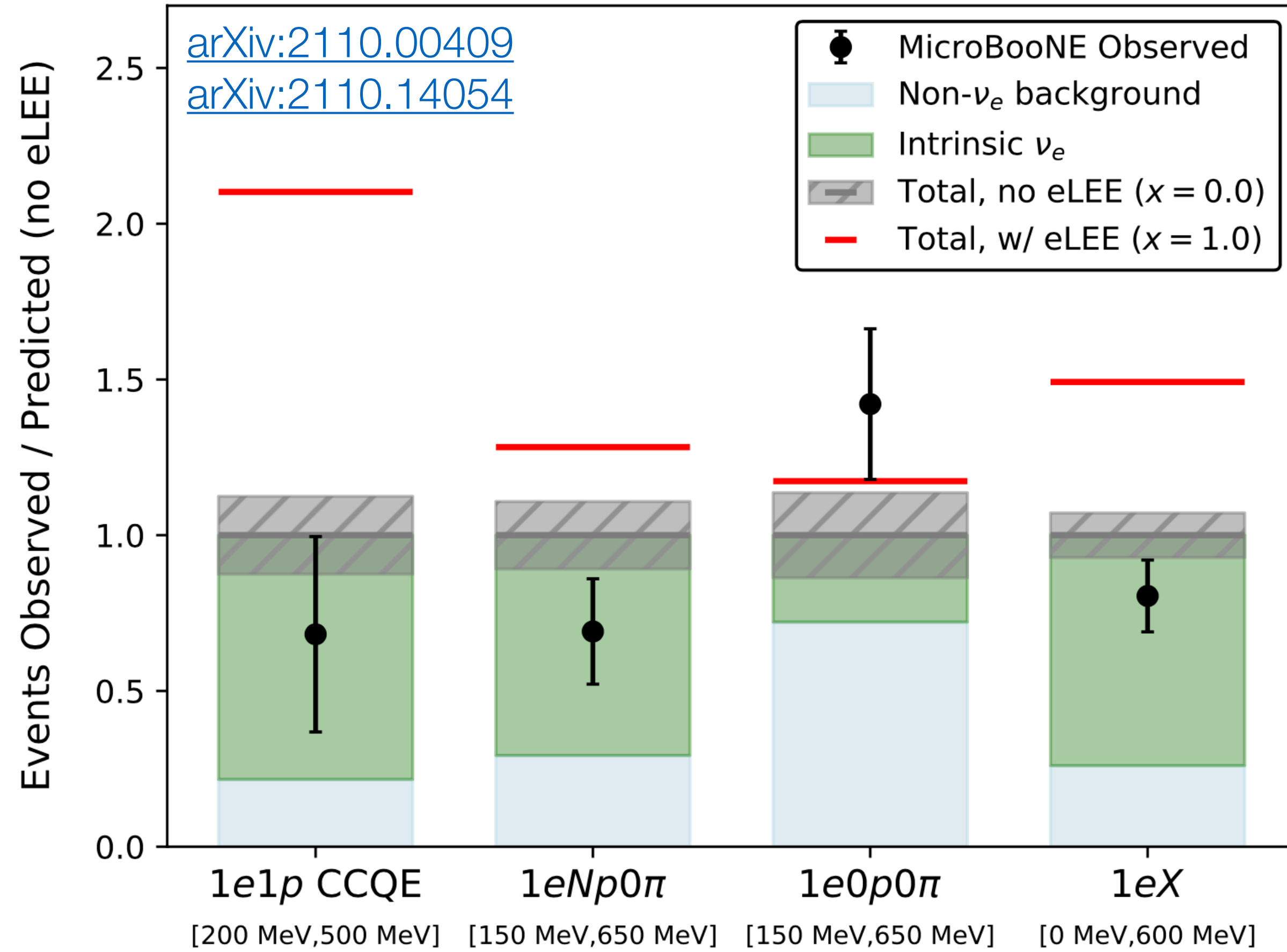


- LAr TPC  $\Rightarrow$  superior event reconstruction
- no excess seen so far (but still consistent with MiniBooNE)





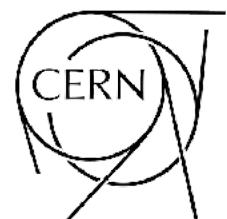
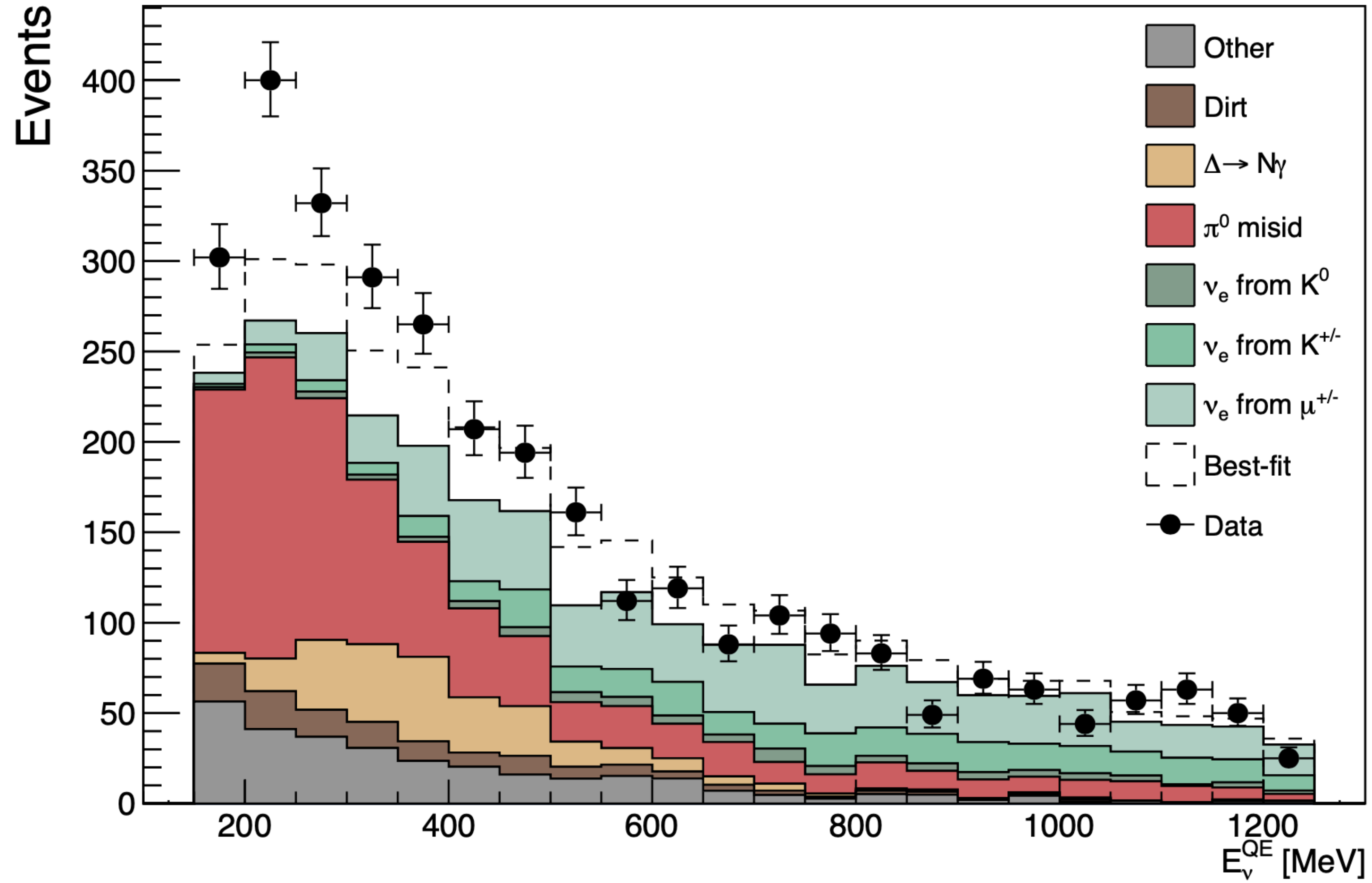
# MicroBooNE





# MiniBooNE

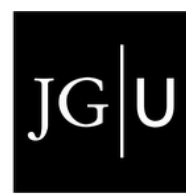
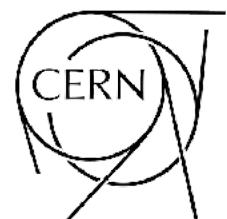
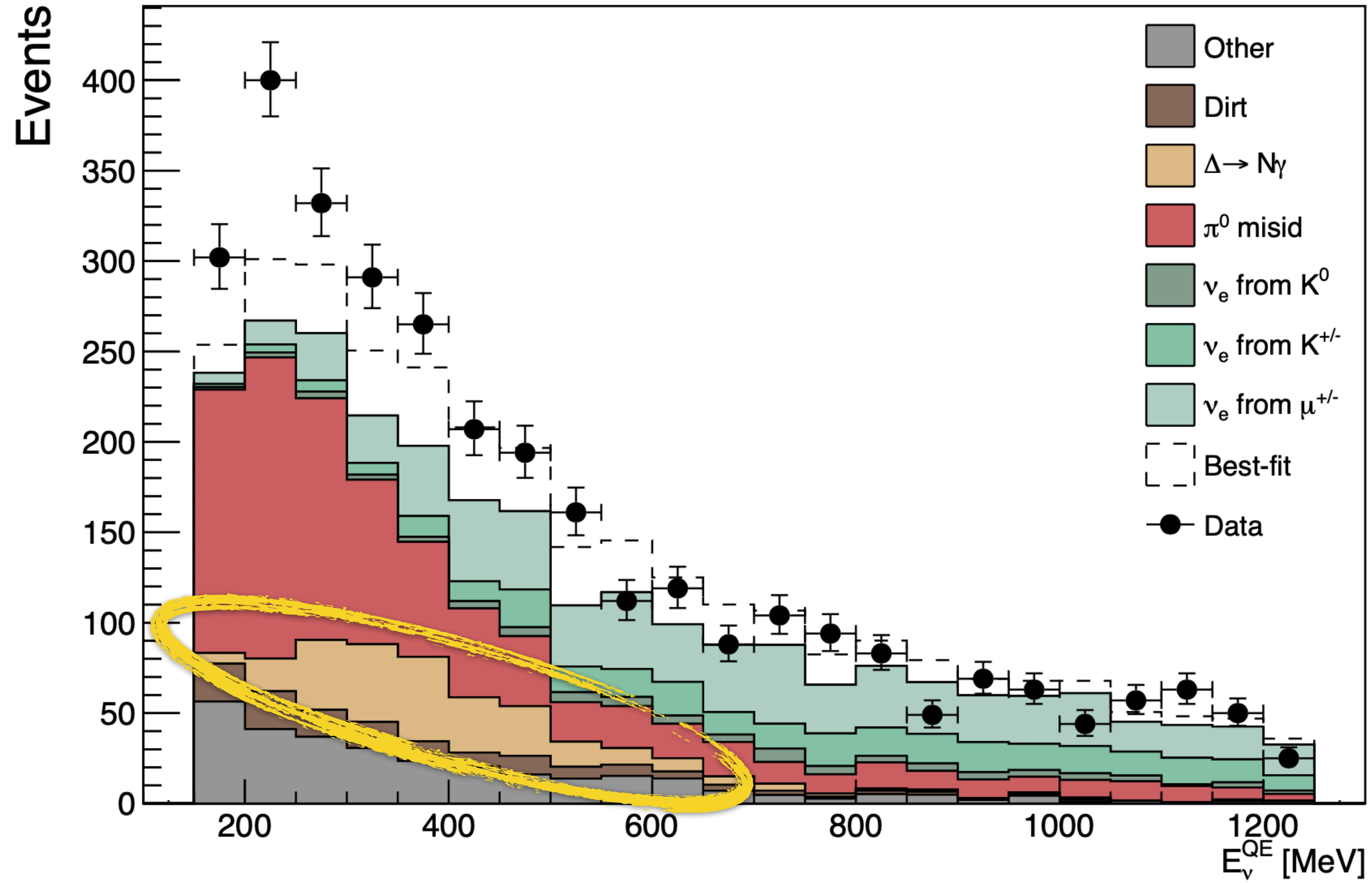
MiniBooNE 2020



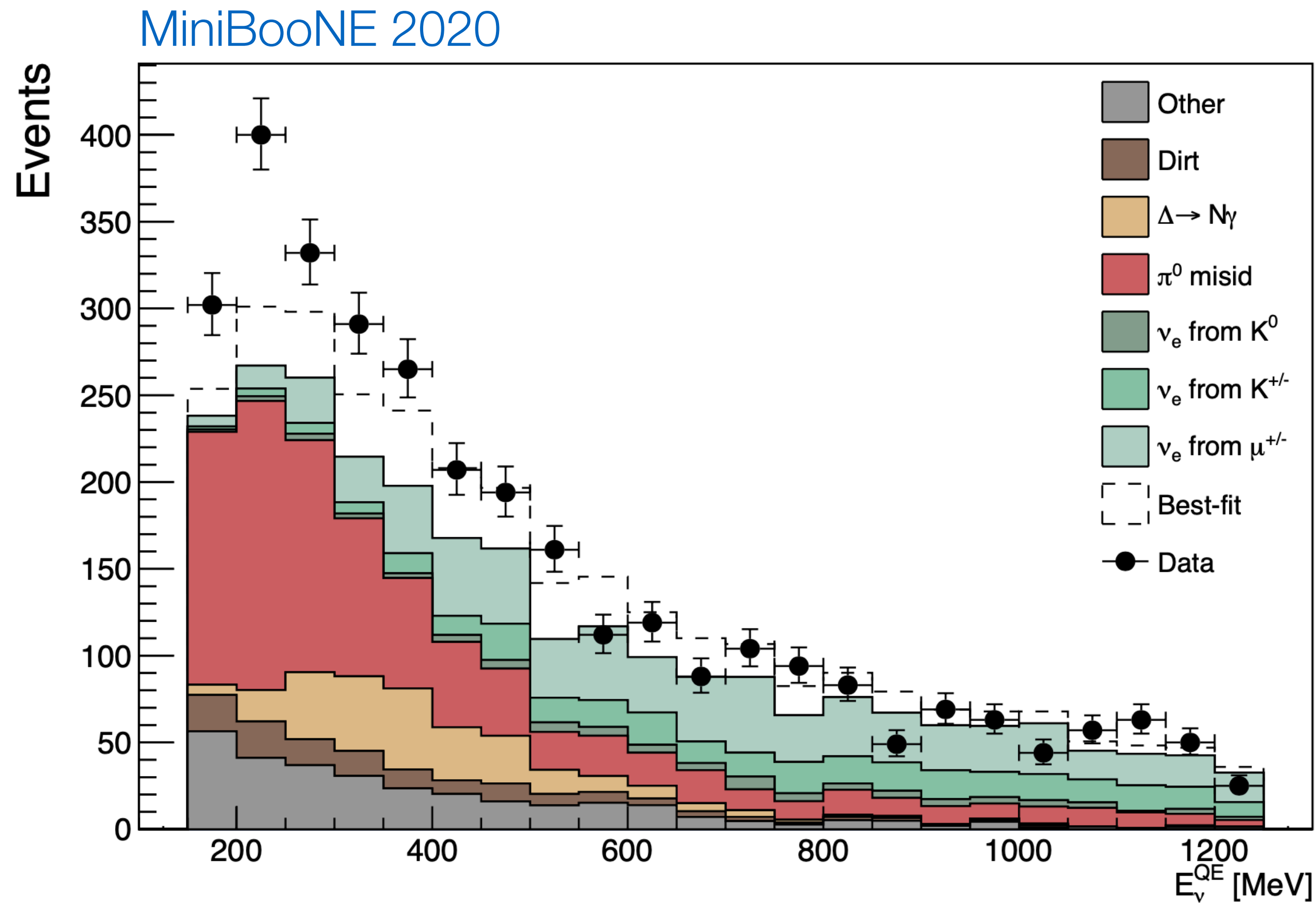


# MiniBooNE

MiniBooNE 2020



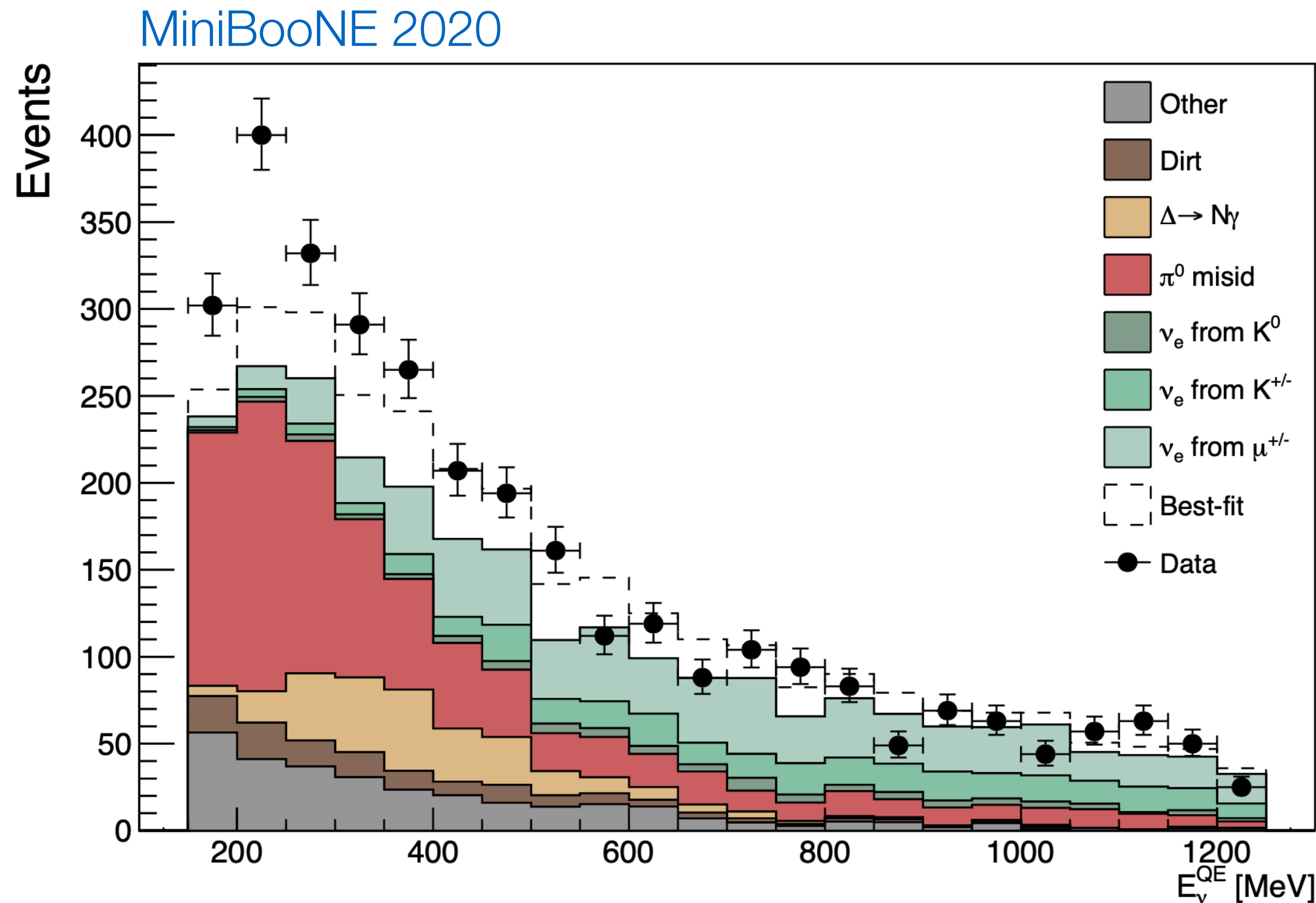
# $\Delta \rightarrow N \gamma$ Background



- NC interaction:  
 $\nu + N \rightarrow \nu + \Delta(1232)$
- Most  $\Delta(1232)$  decay to  $\pi + N$
- But rare decay exists to  $\gamma + N$
- MiniBooNE cannot distinguish single- $\gamma$  background from CC  $\nu_e$  signal



# $\Delta \rightarrow N \gamma$ Background

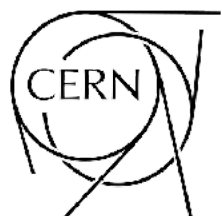


- $\Delta$  production rate can be estimated from  $\Delta \rightarrow \pi N$
- Pions may be **absorbed** on their way out of the nucleus
  - may **excite another  $\Delta(1232)$** 
    - ▮  $\Delta \rightarrow \gamma N$  enhanced
  - or may be **absorbed**
    - ▮ control region suppressed

Ioannian [1909.08571](#)  
Giunti Ioannian Ranucci [1912.01524](#)

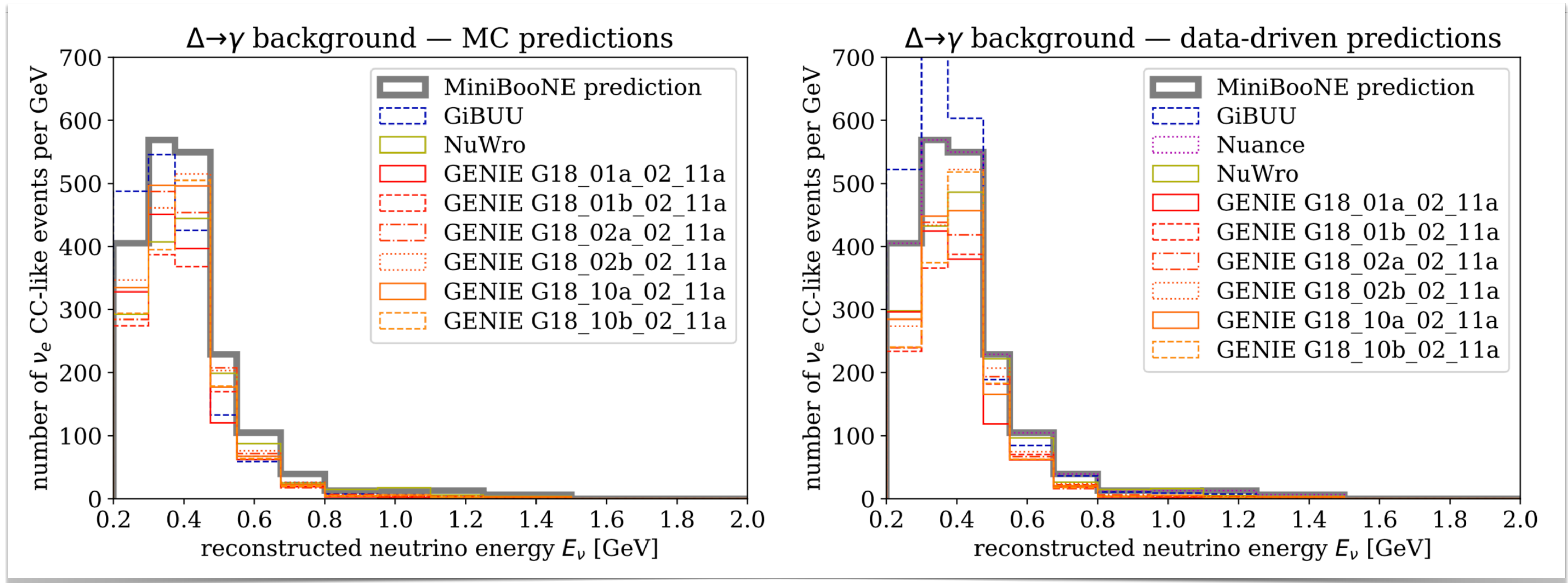
(These effects **have been taken into account** by MiniBooNE)

MiniBooNE, [arXiv:2006.16883](#)



# $\Delta \rightarrow N \gamma$ – Comparison of Generators

Brdar JK, arXiv:2109.08157



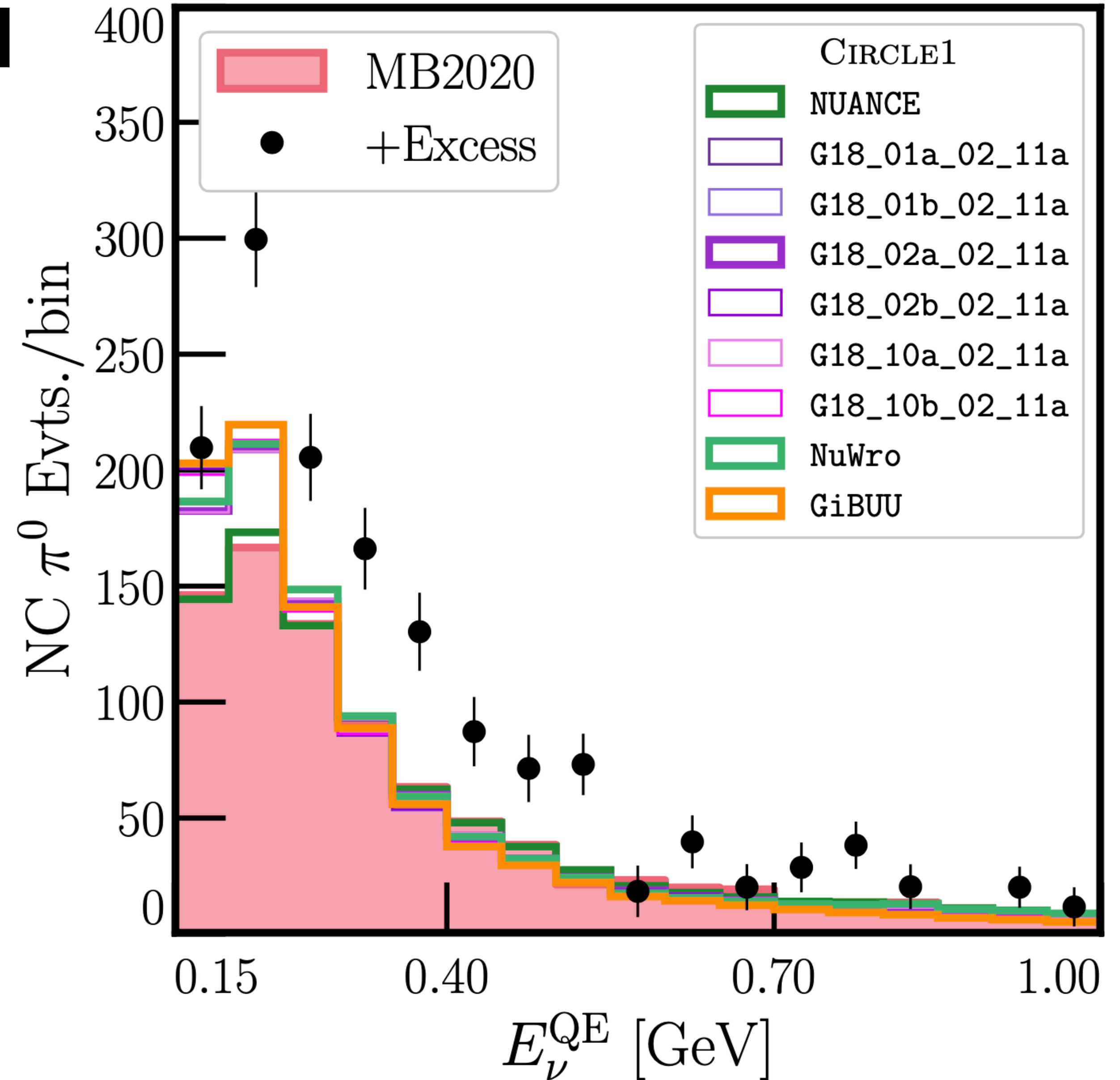
- NUANCE prediction matched to MB prediction, others rescaled in the same way
- using our [own implementation of radiative resonance decays](#) in GiBUU, NuWro, NUANCE



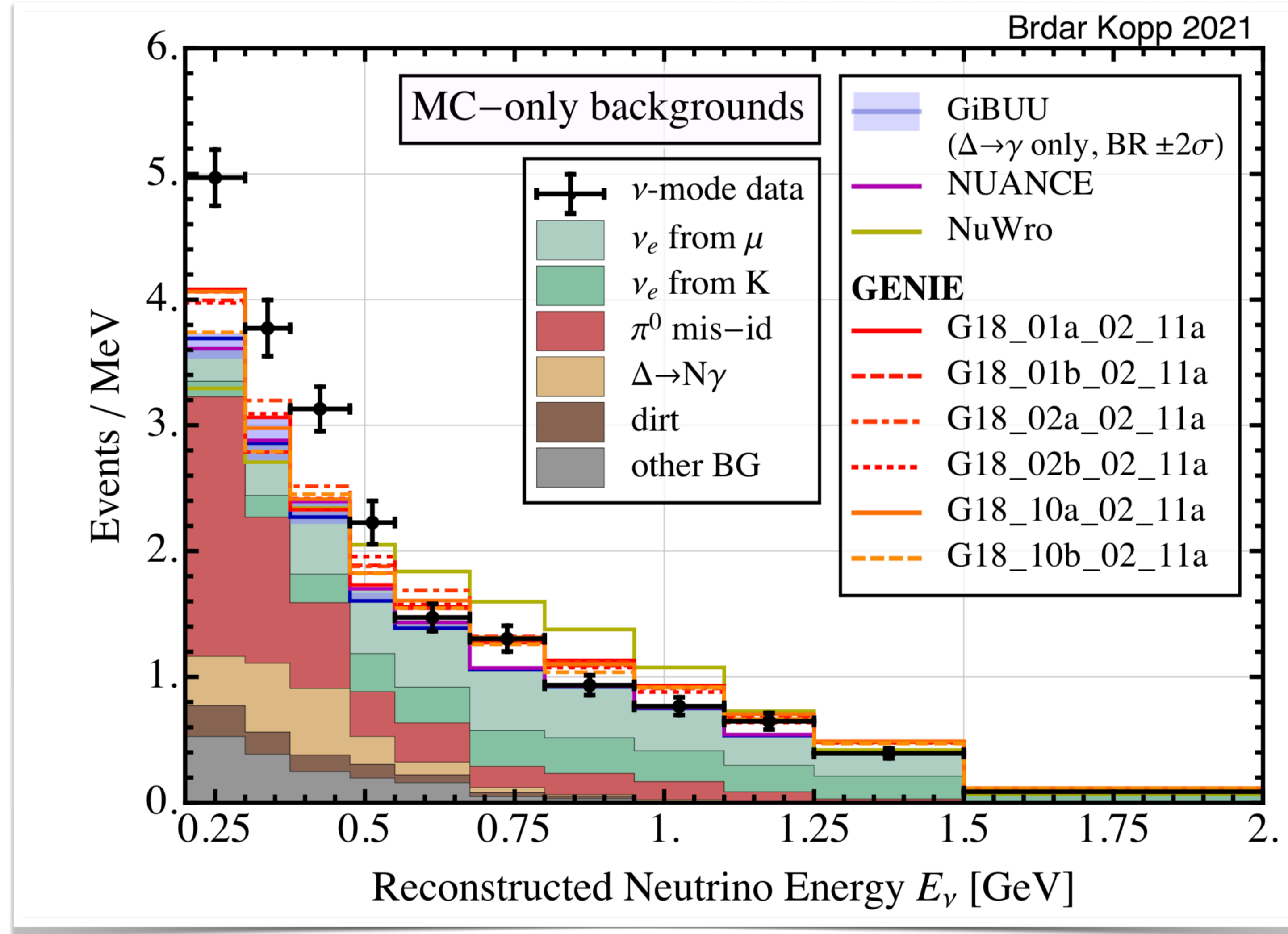
# MiniBooNE $\pi^0$ Background

- Attempt to reproduce MiniBooNE's  $\pi^0$  mis-ID probability
- using our [own implementation of radiative resonance decays](#) in GiBUU, NuWro, NUANCE

Kelly JK, arXiv:2210.08021



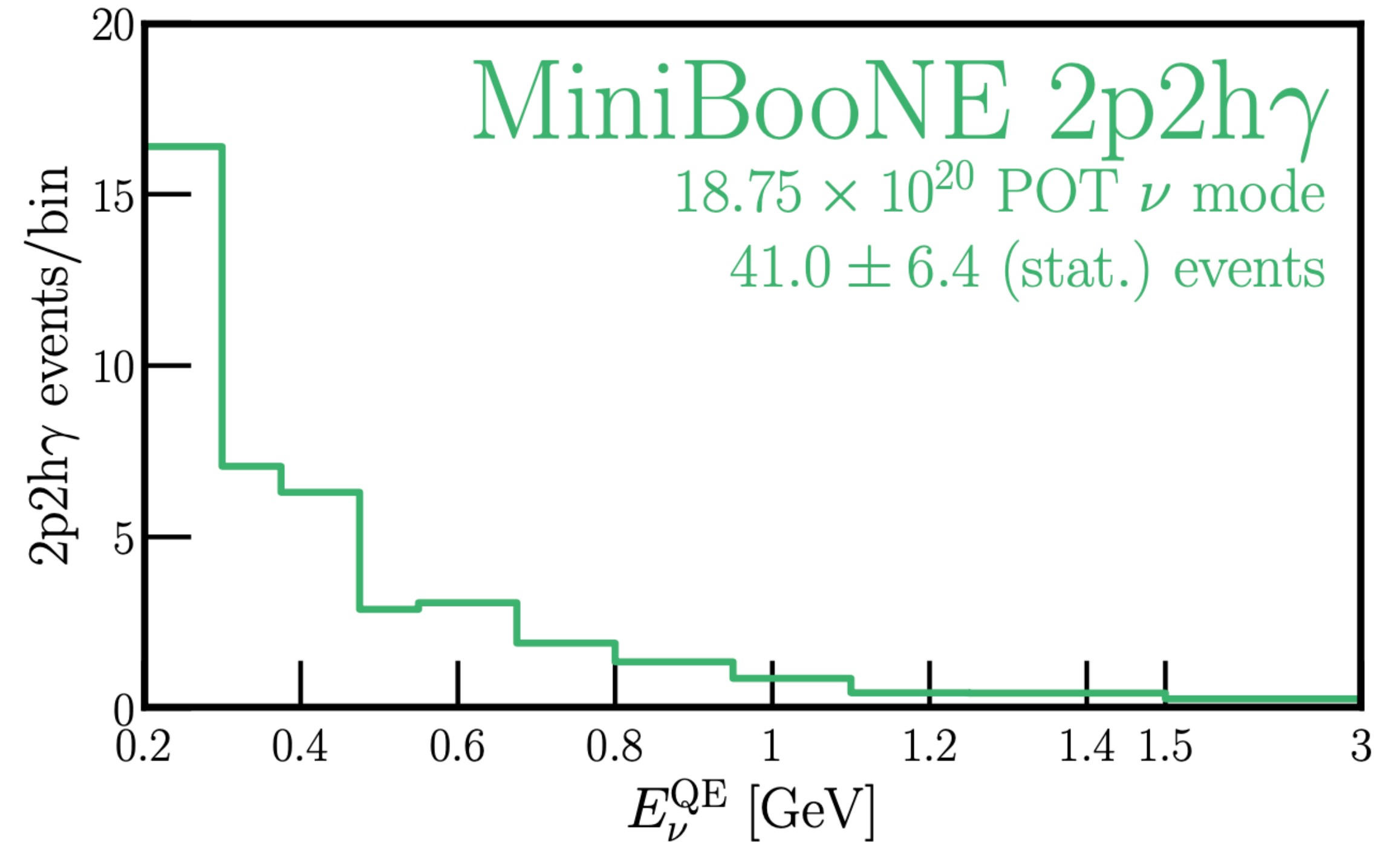
# All Backgrounds – Comparison of Generators





# New Processes? – 2p2h+ $\gamma$

- intuition: neutrino scatters on 2-nucleon system
  - ▢ radiated  $\gamma$  sees both nucleons coherently
  - ▢ enhanced by factor 2
- very naïve implementation (elastic scattering on 2-nucleon system)
- event rate falls far short of explaining the anomaly



Kelly JK, arXiv:2210.08021



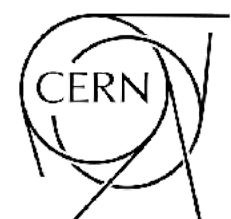
# My Bet for MiniBooNE

an  
“*Altarelli Cocktail*”



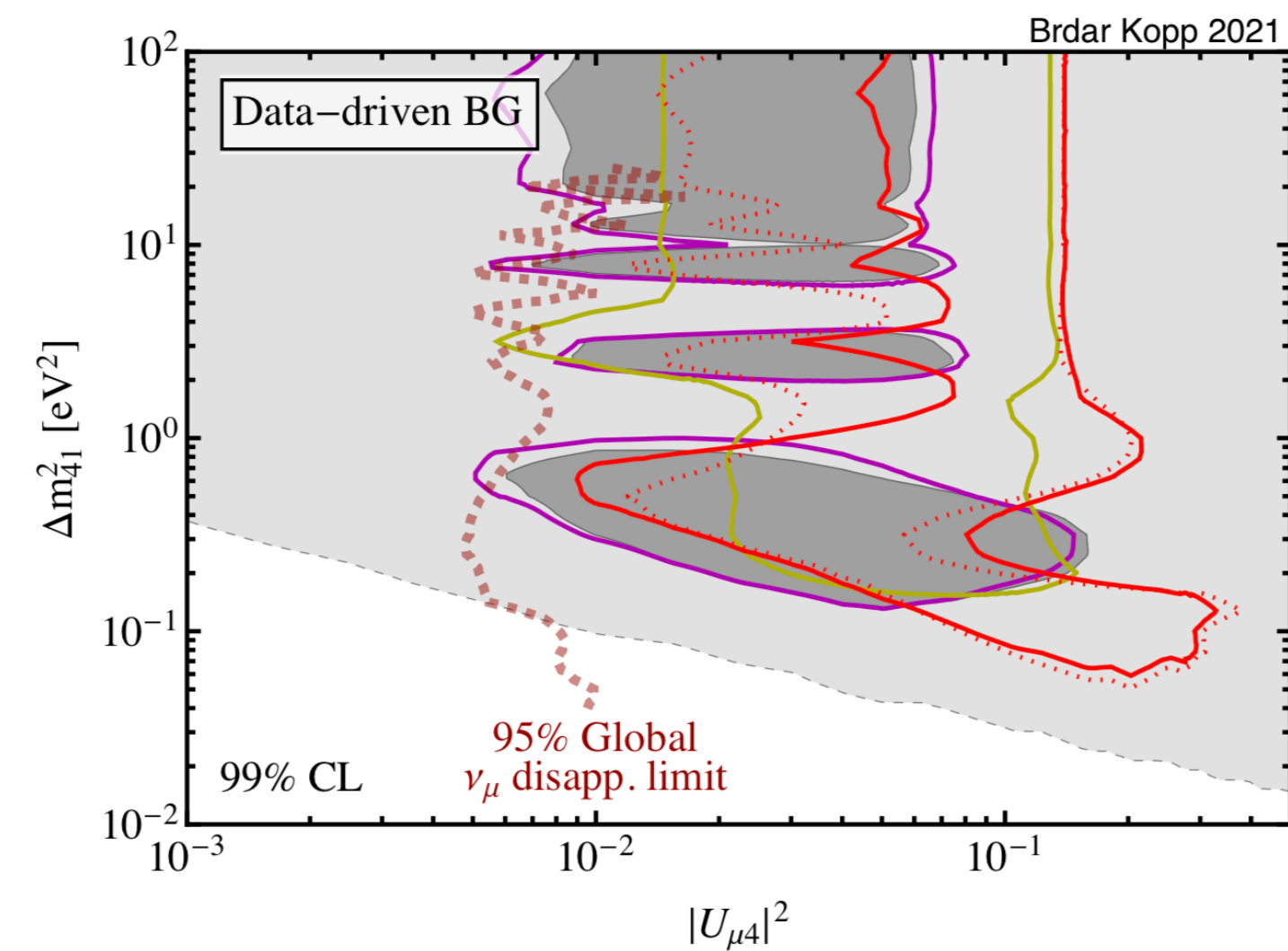
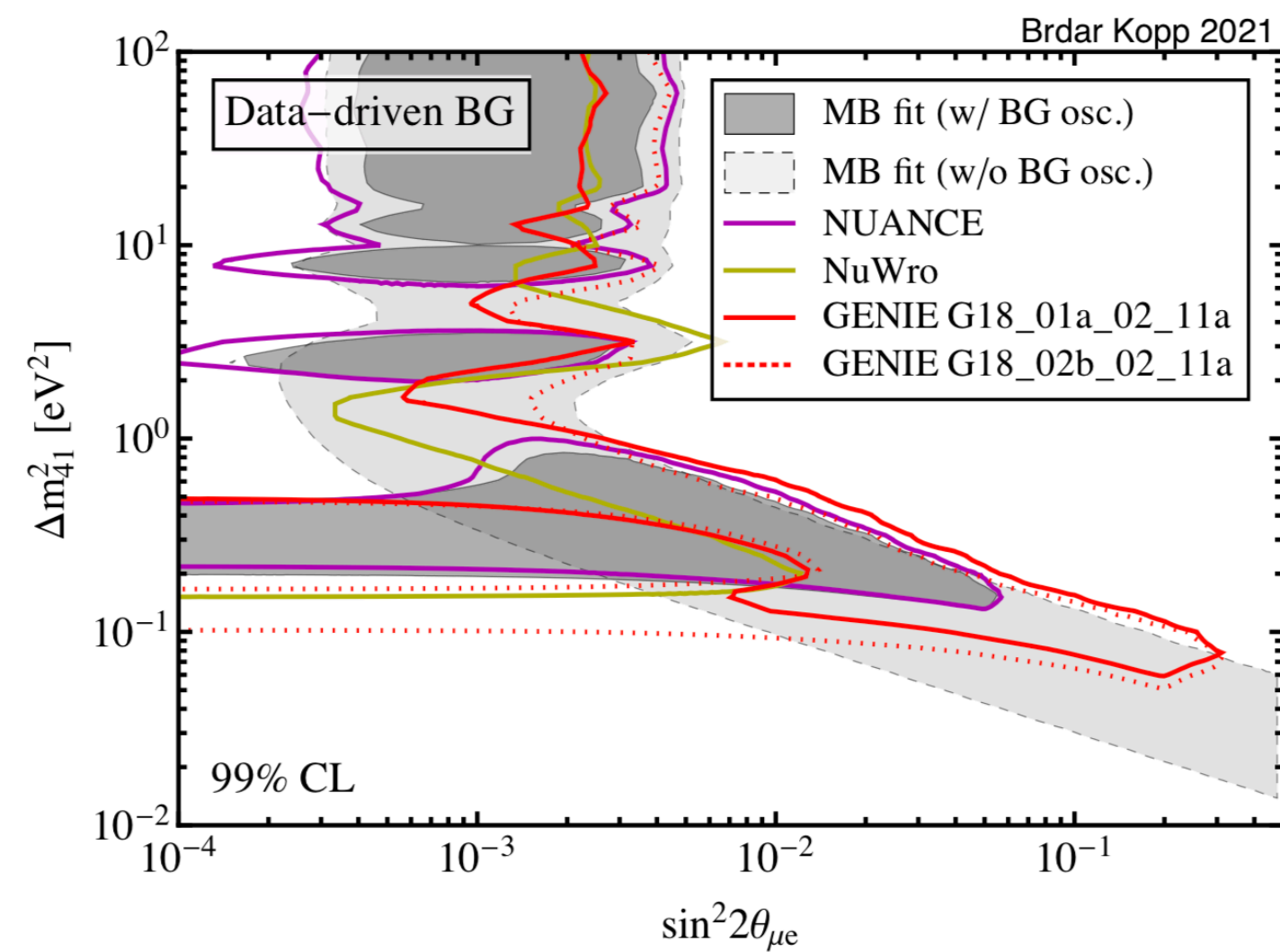
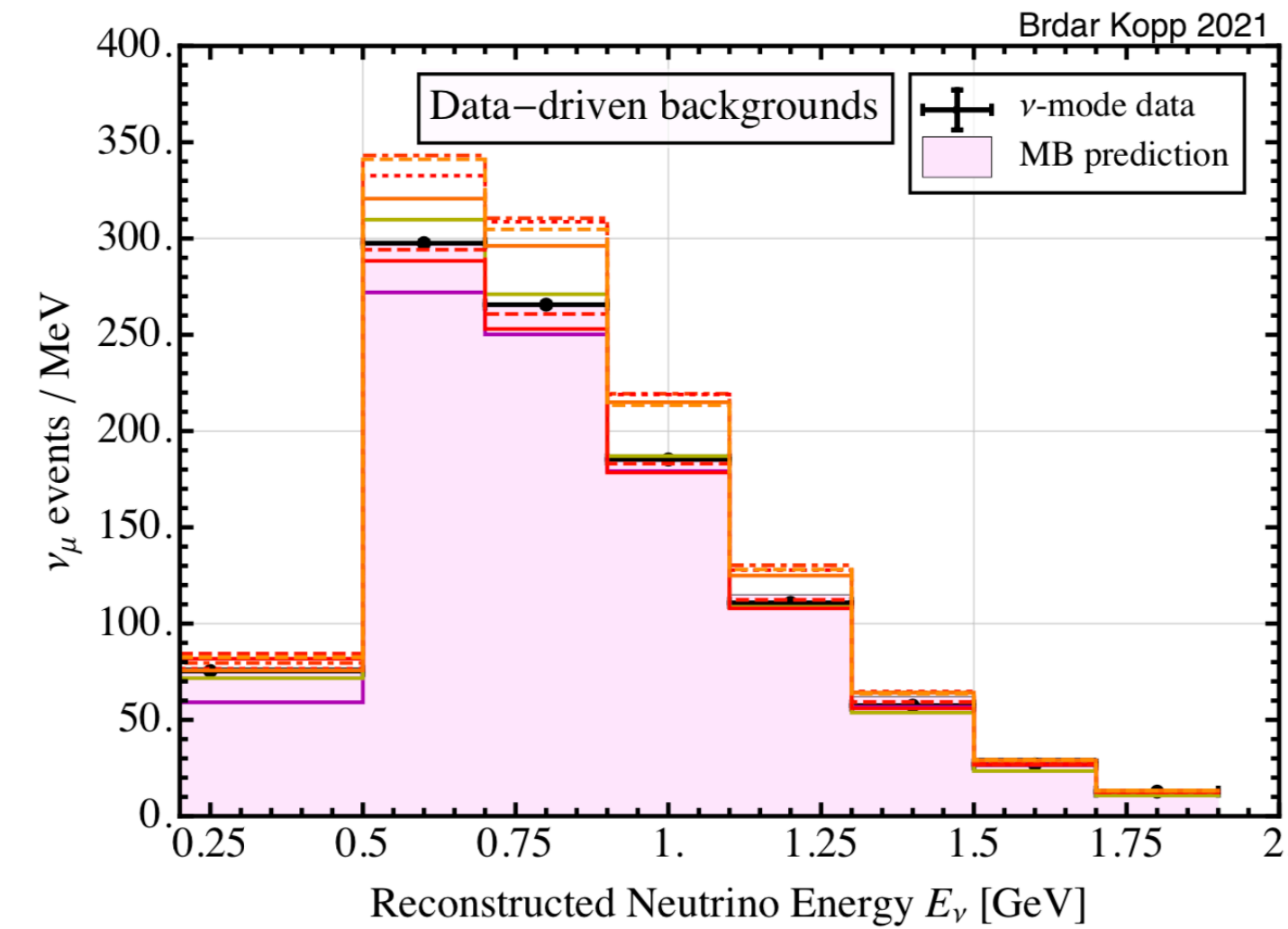
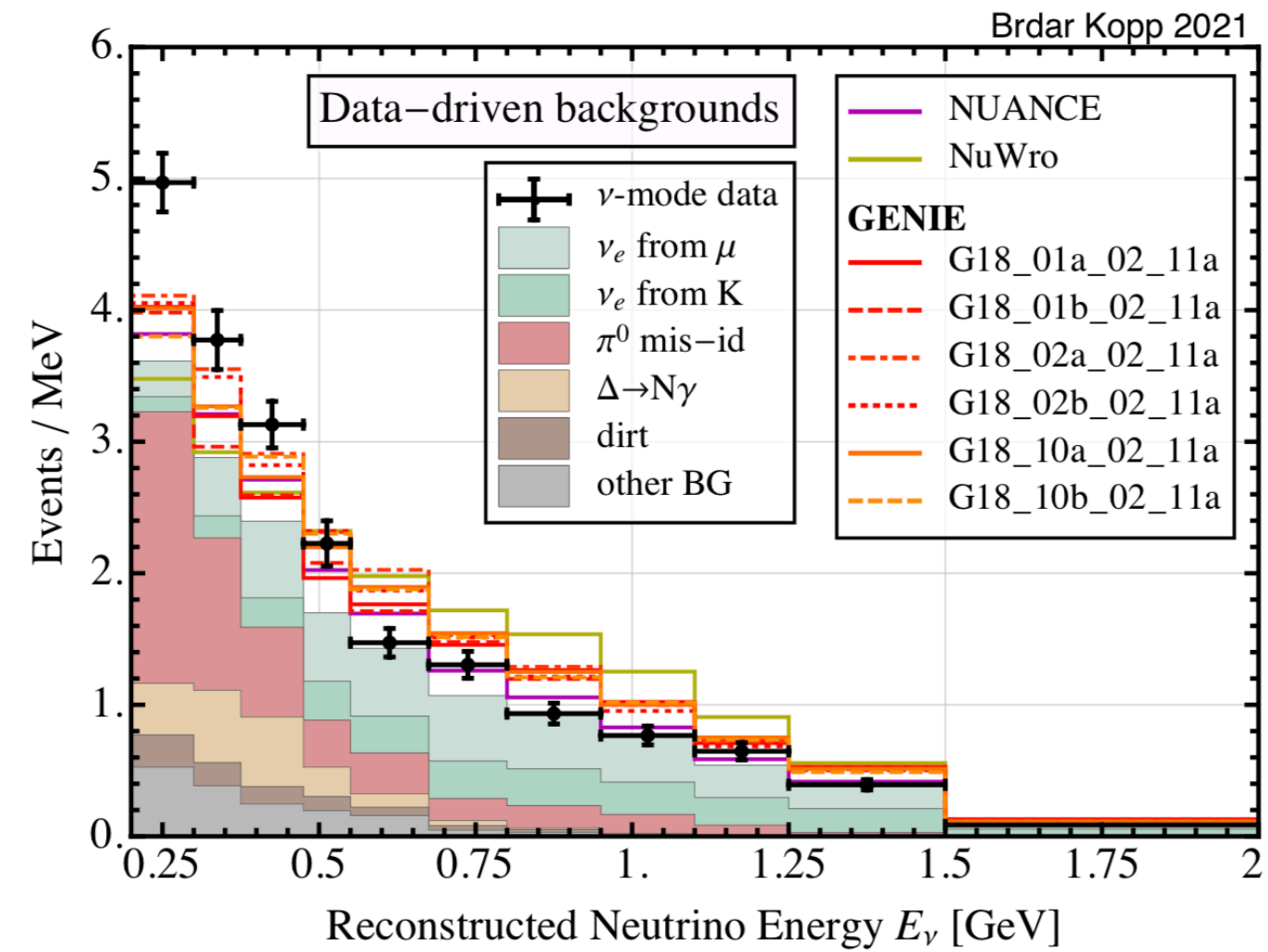
Image: ChatGPT

(who refused to draw an Italian physicist w/o a beard)



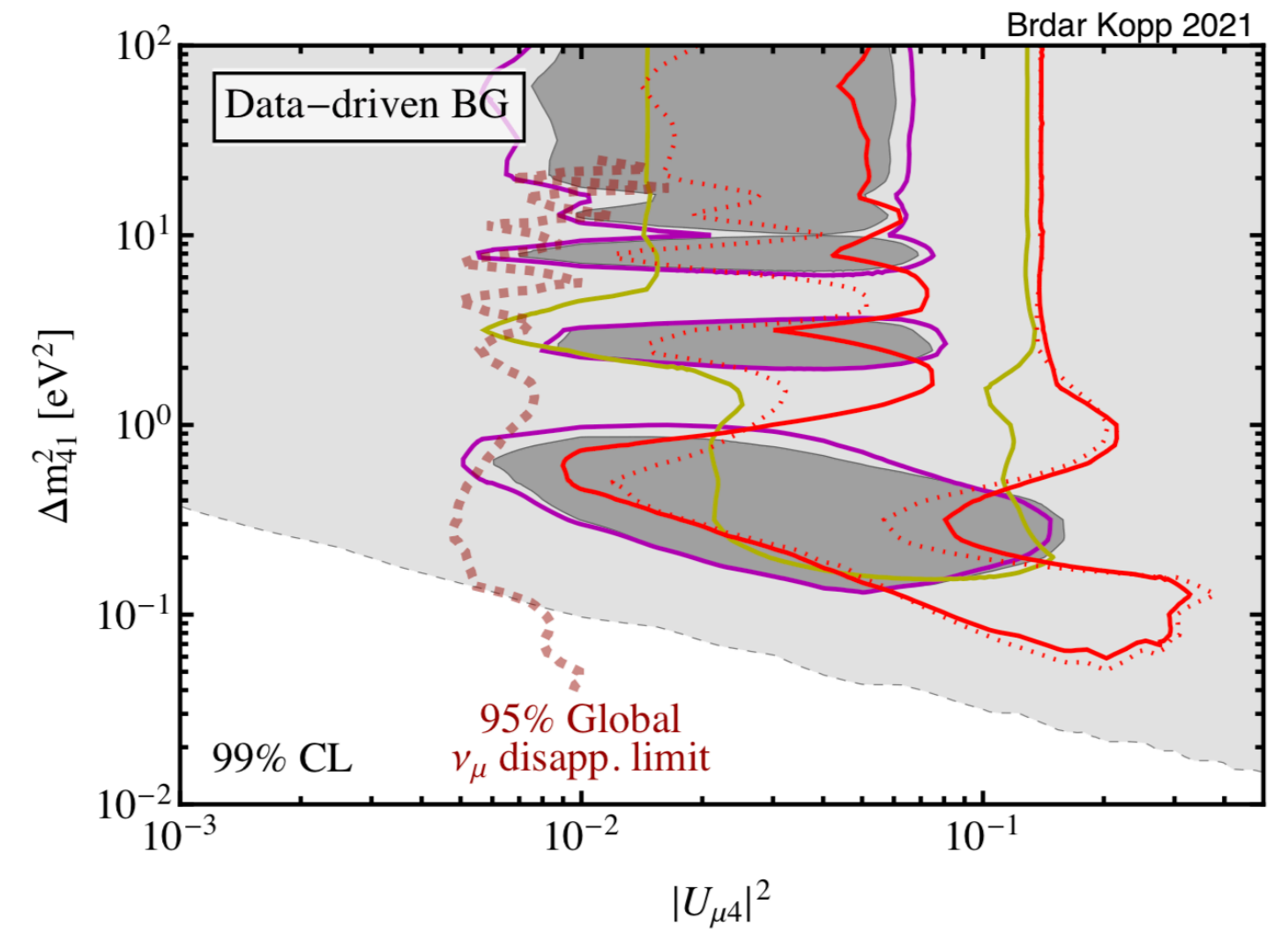
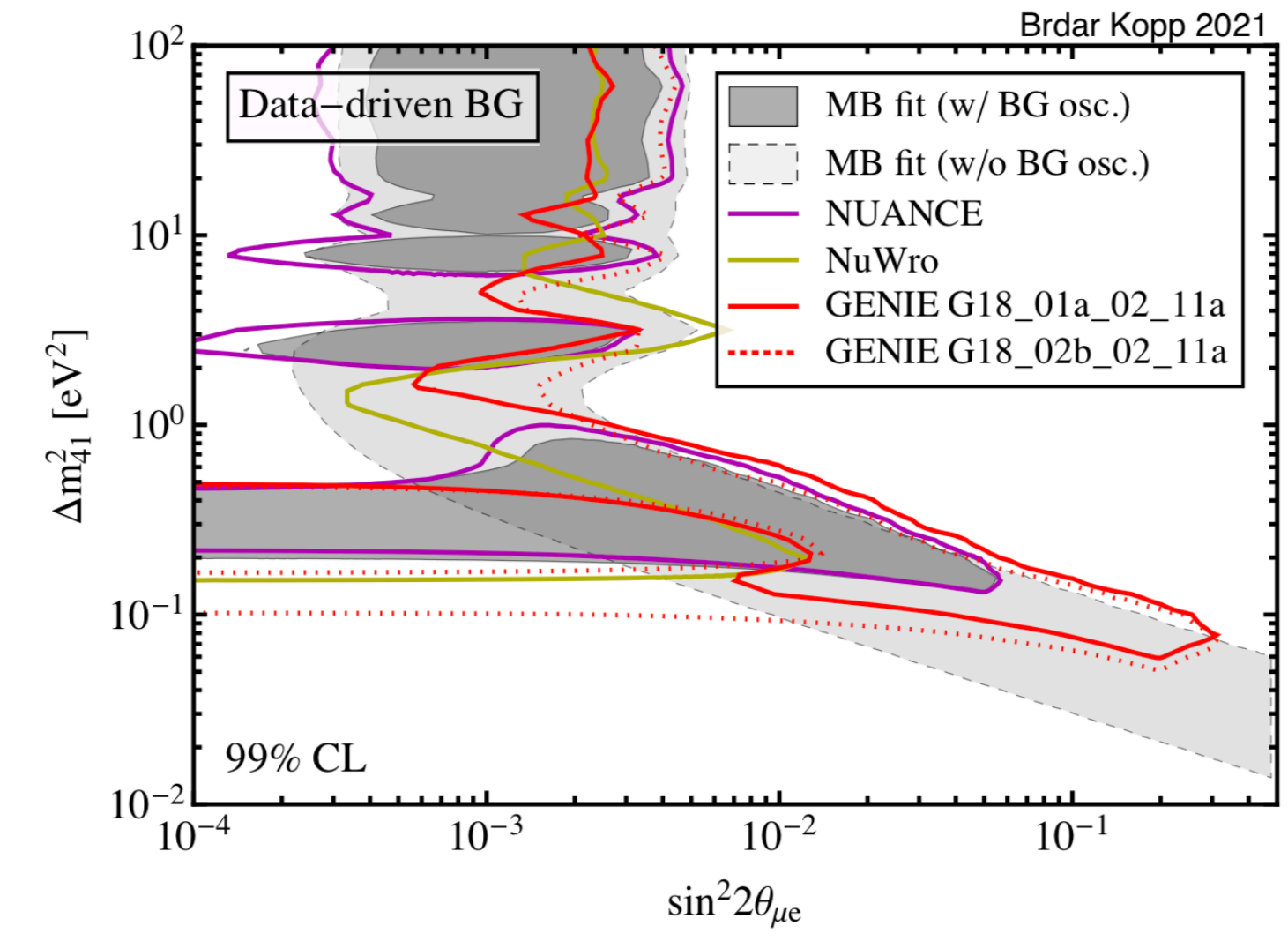
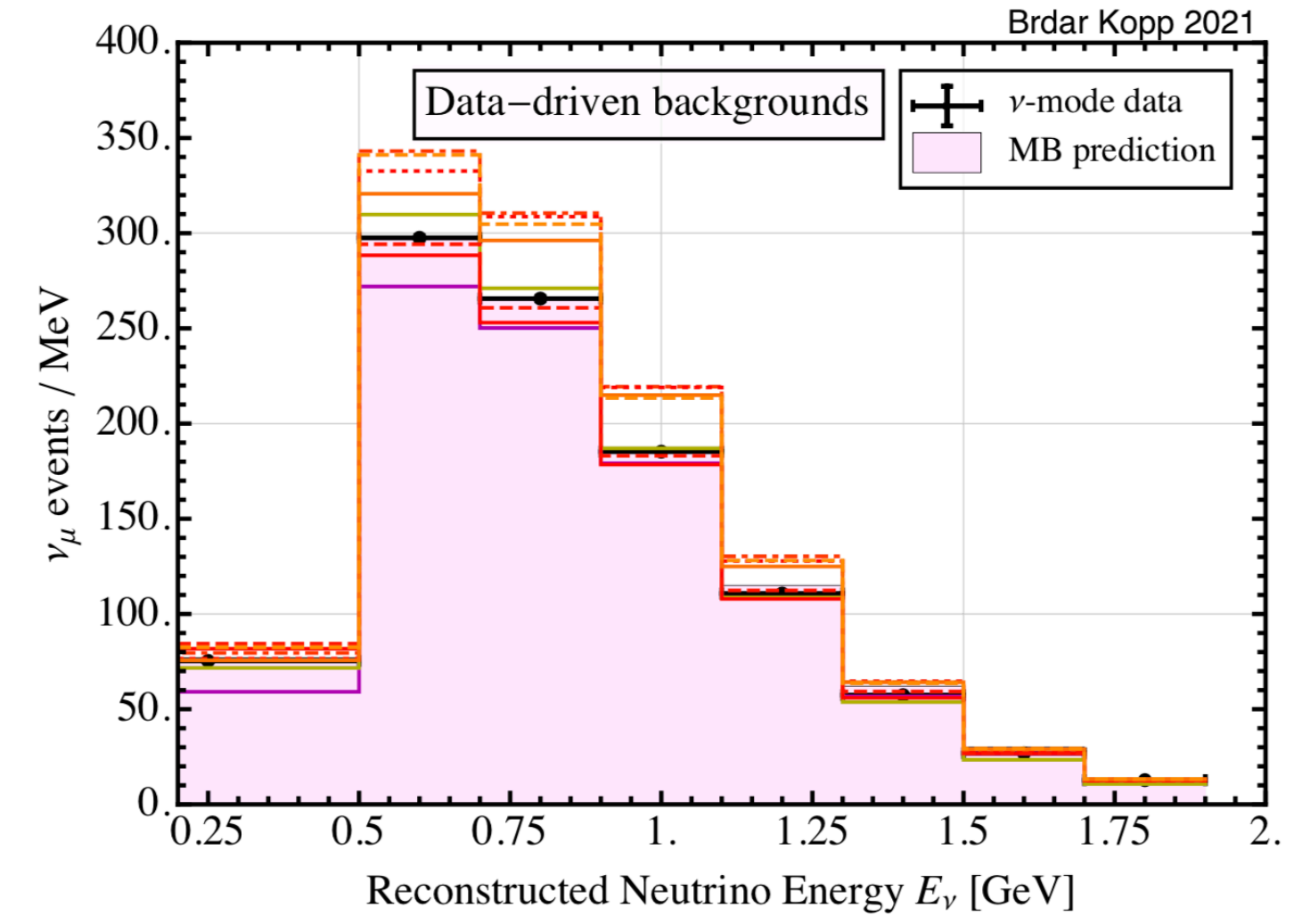
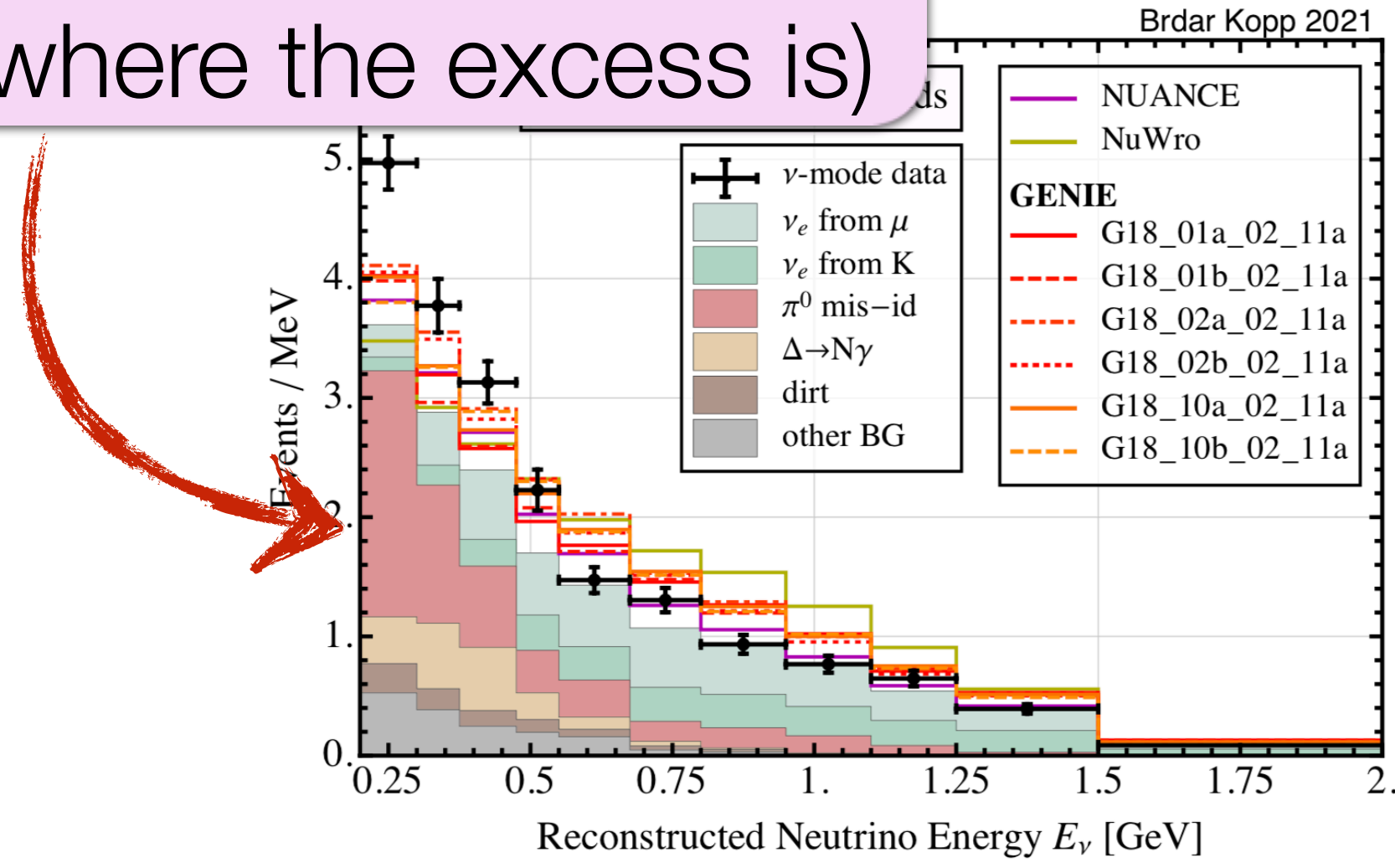


# 3+1 Models in MiniBooNE – Comparison of Generators



# 3+1 Models in MiniBooNE – Comparison of Generators

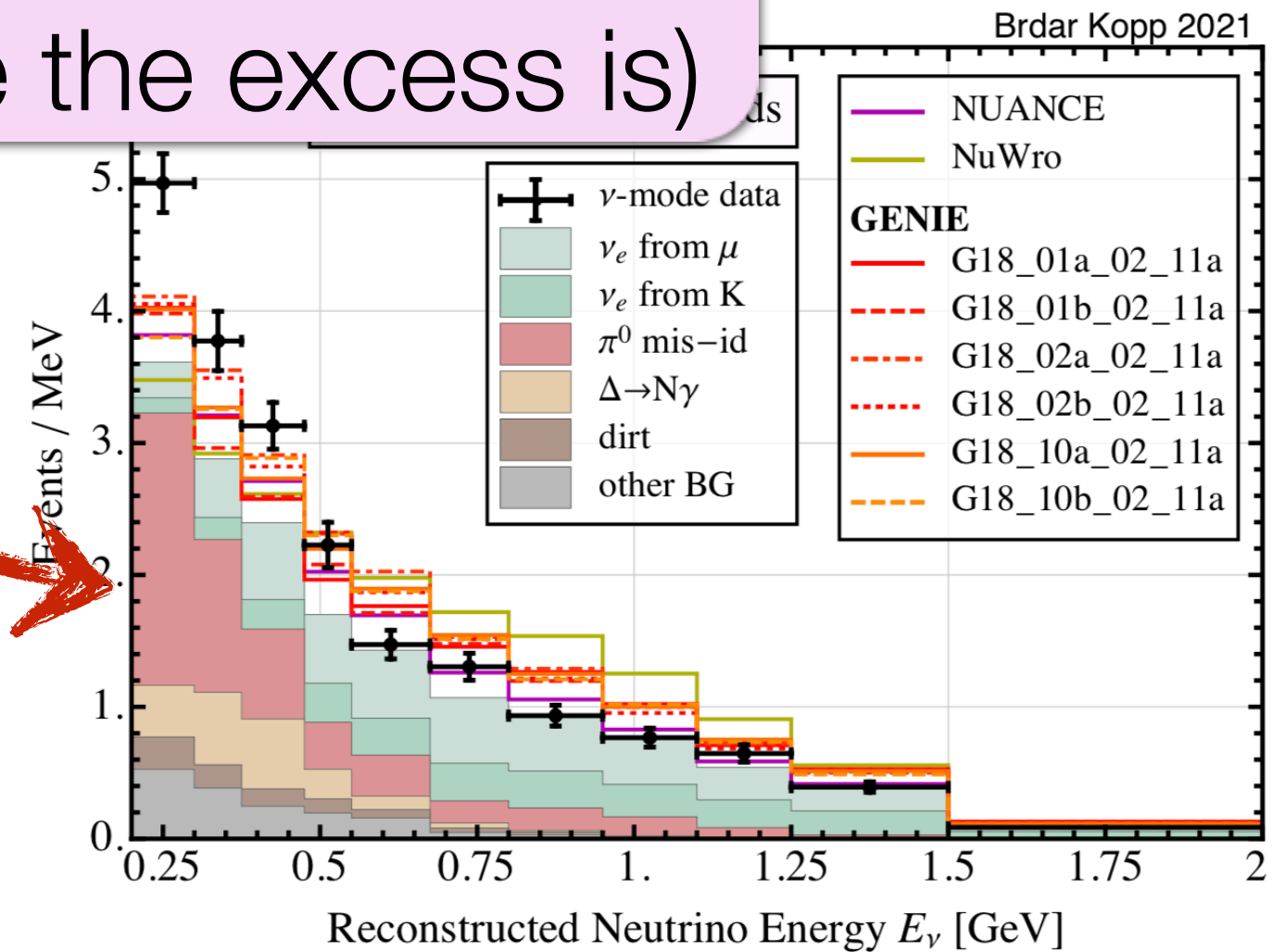
$\nu_e$  spectrum  
(where the excess is)



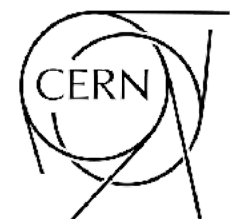
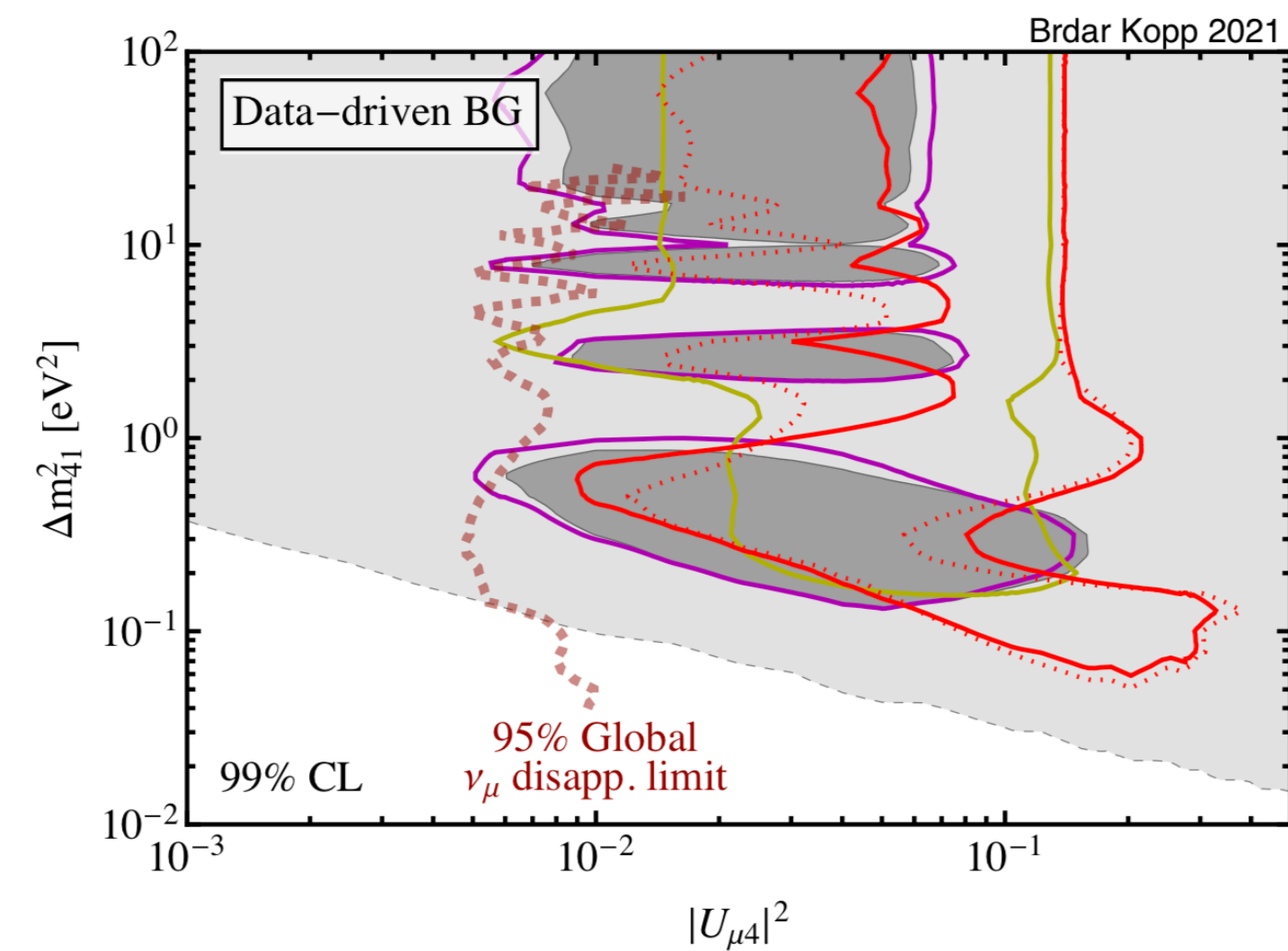
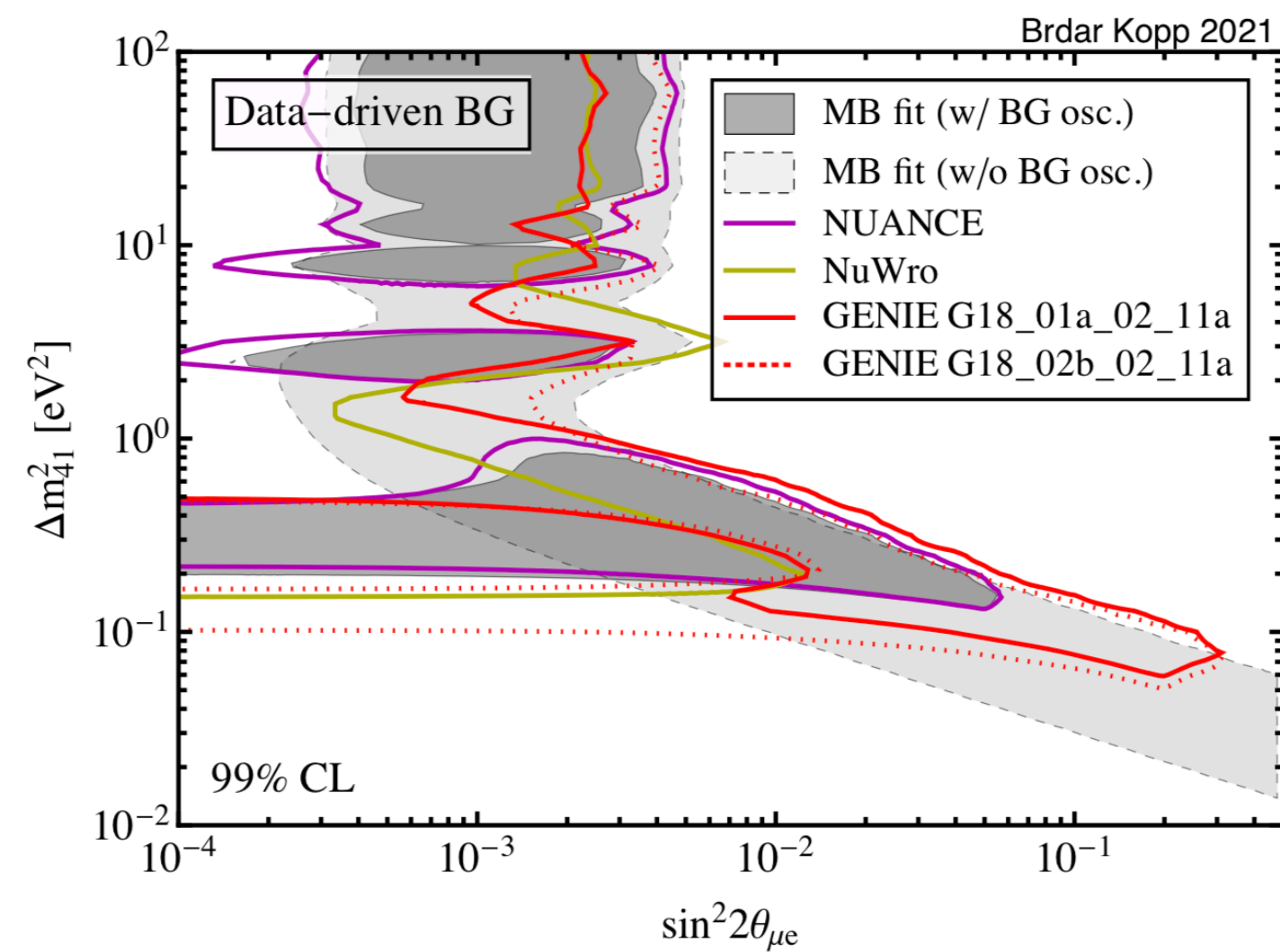
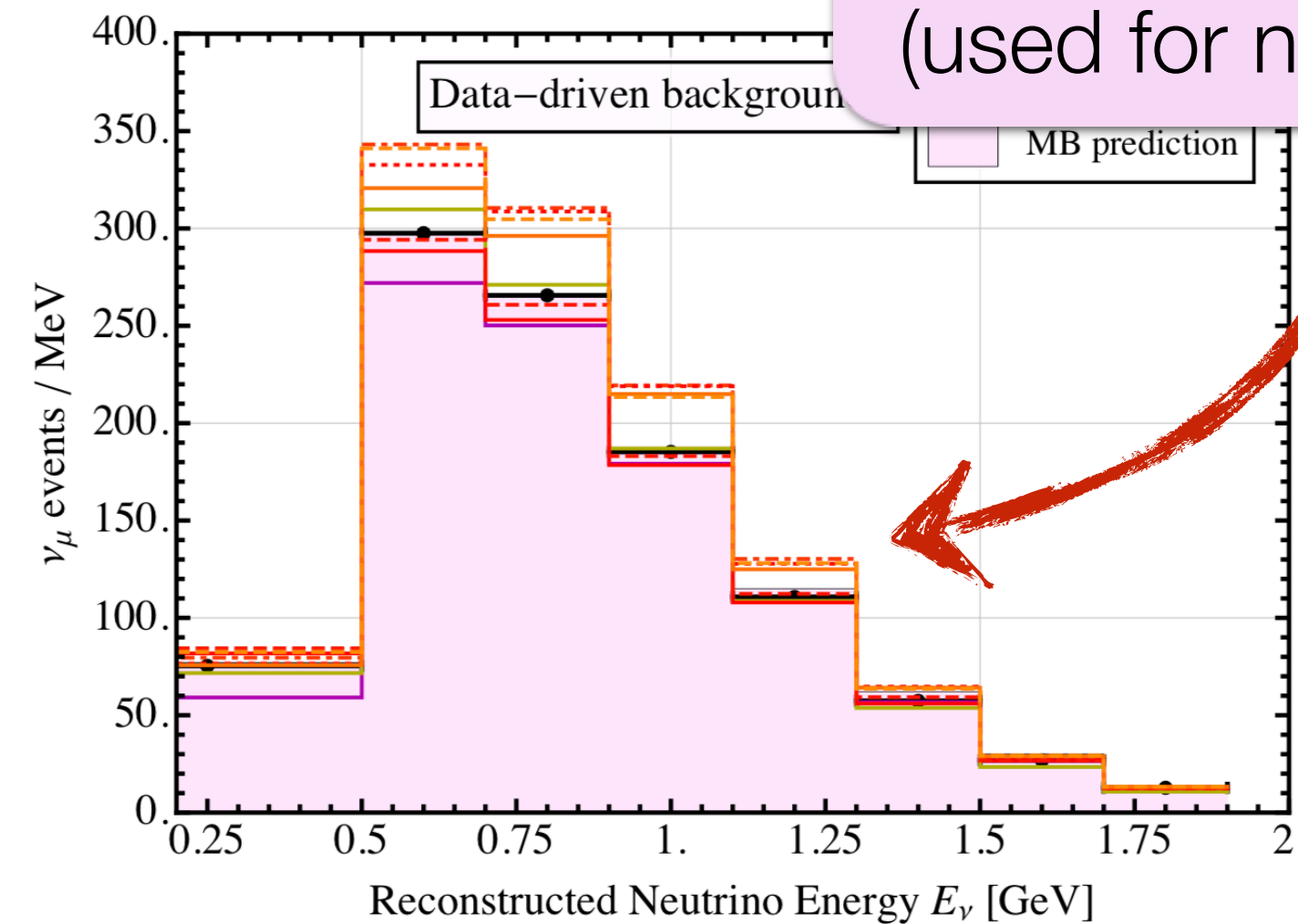


# 3+1 Models in MiniBooNE – Comparison of Generators

$\nu_e$  spectrum  
(where the excess is)

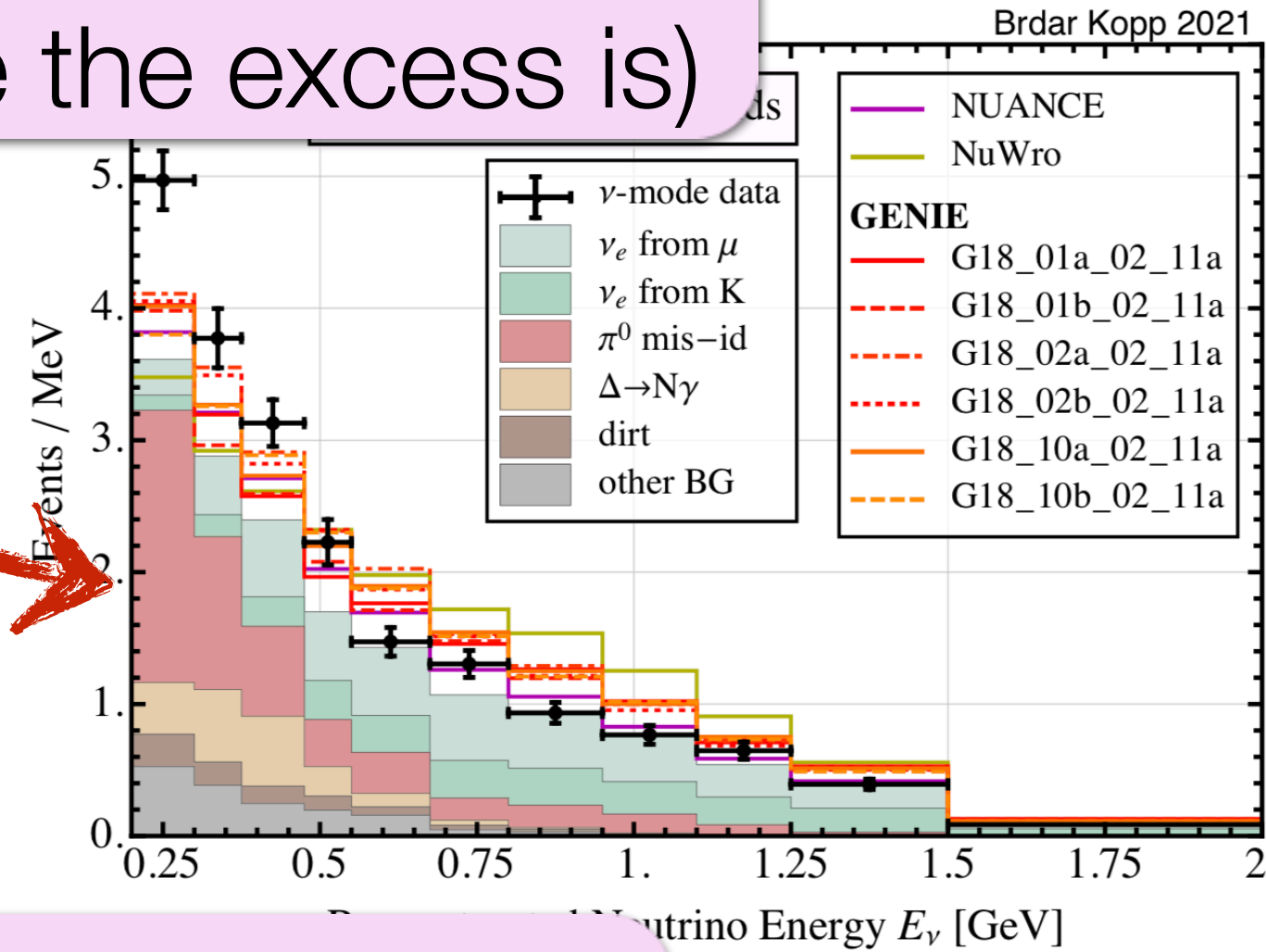


$\nu_\mu$  spectrum  
(used for normalization)

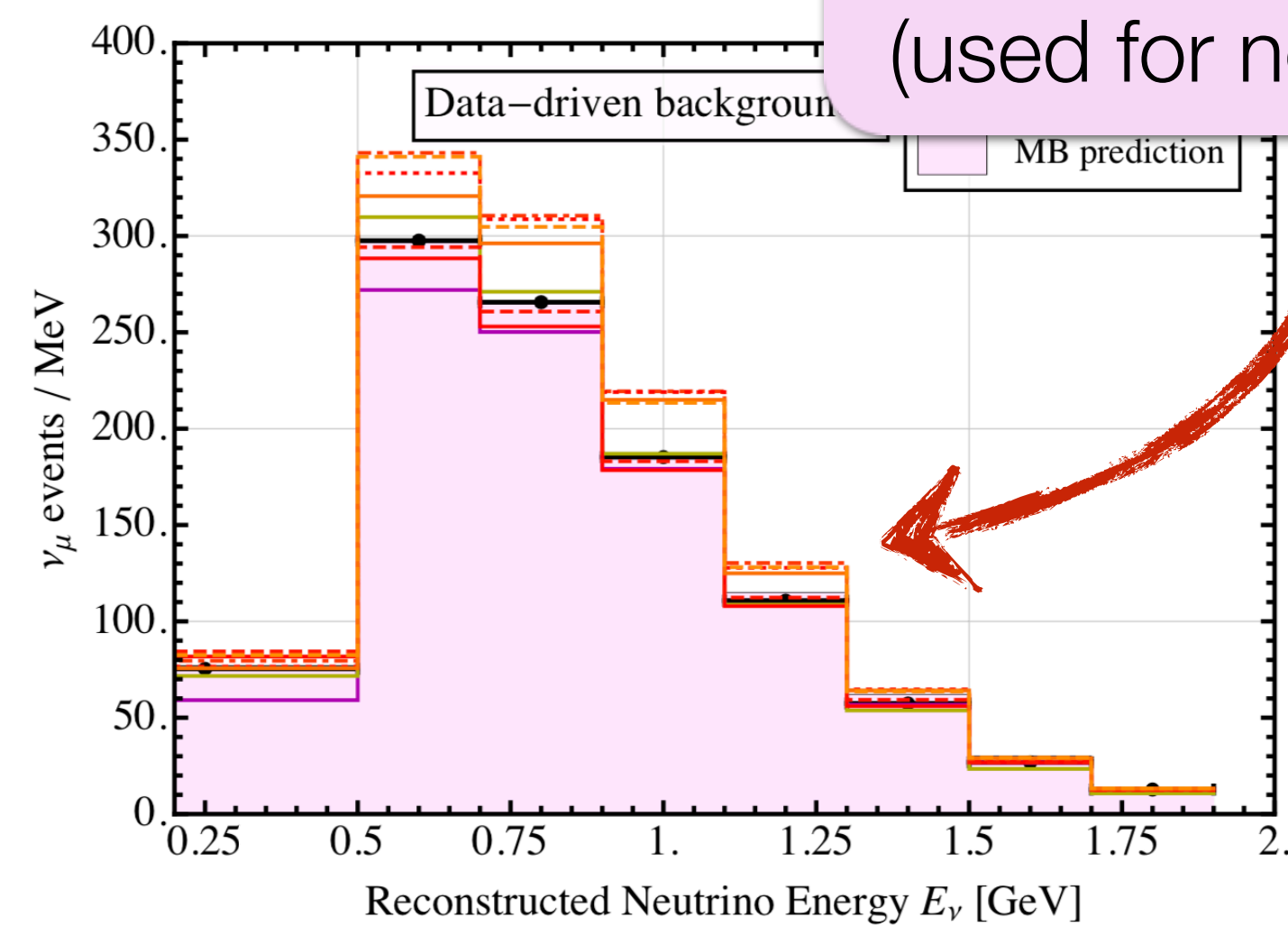


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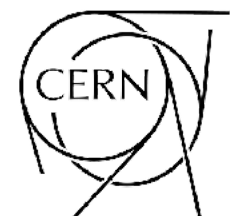
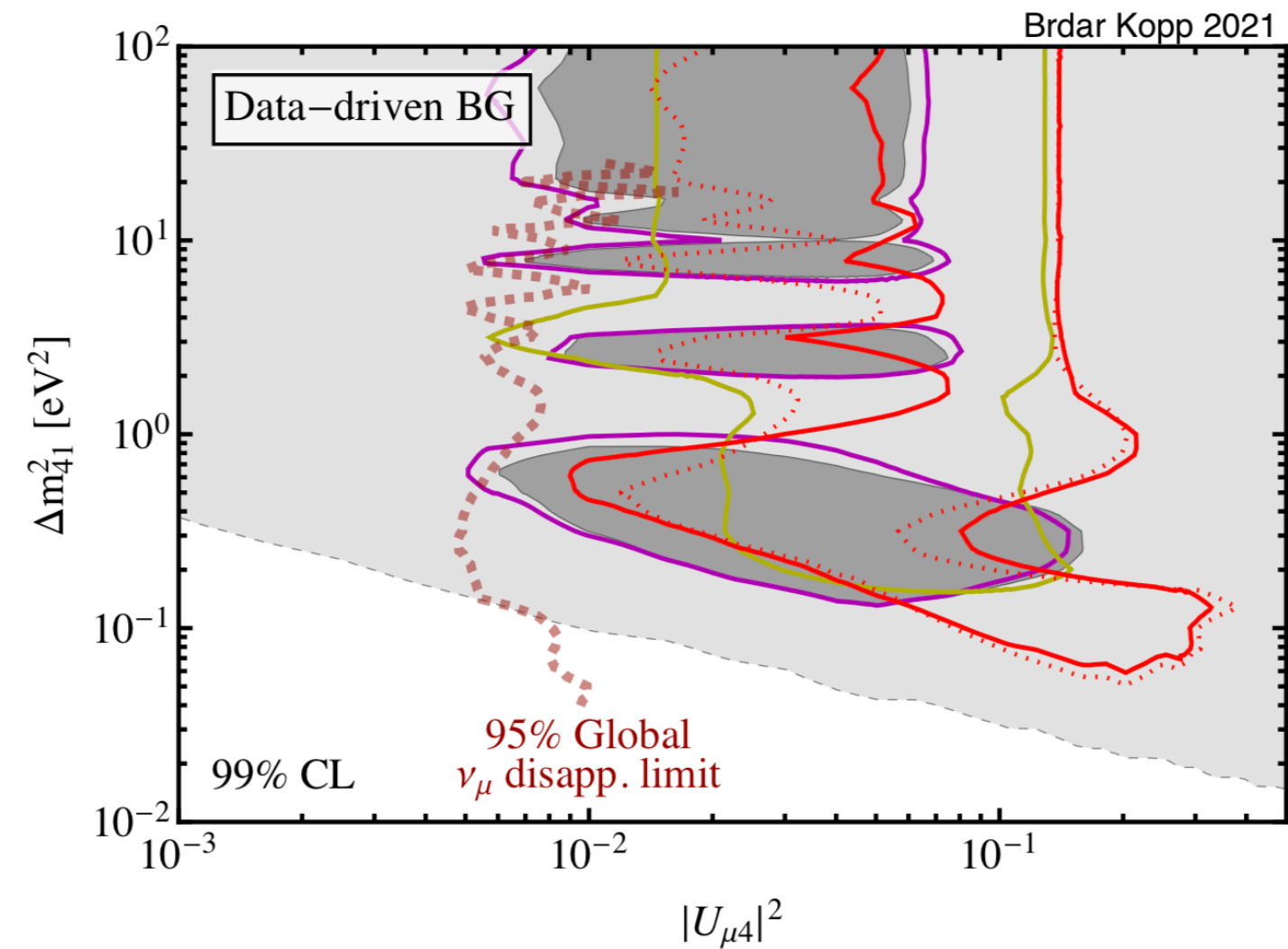
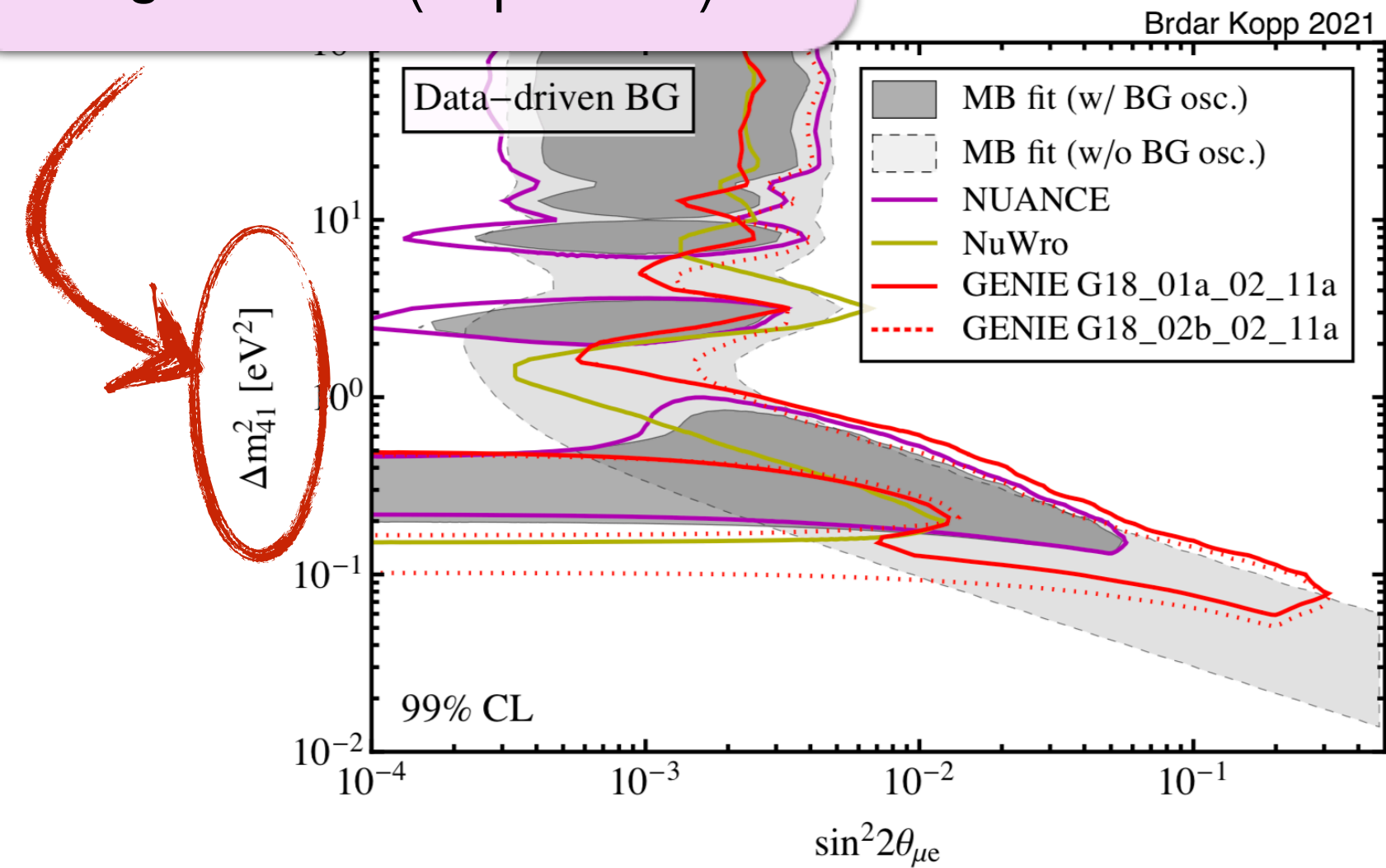
$\nu_e$  spectrum  
(where the excess is)



$\nu_\mu$  spectrum  
(used for normalization)



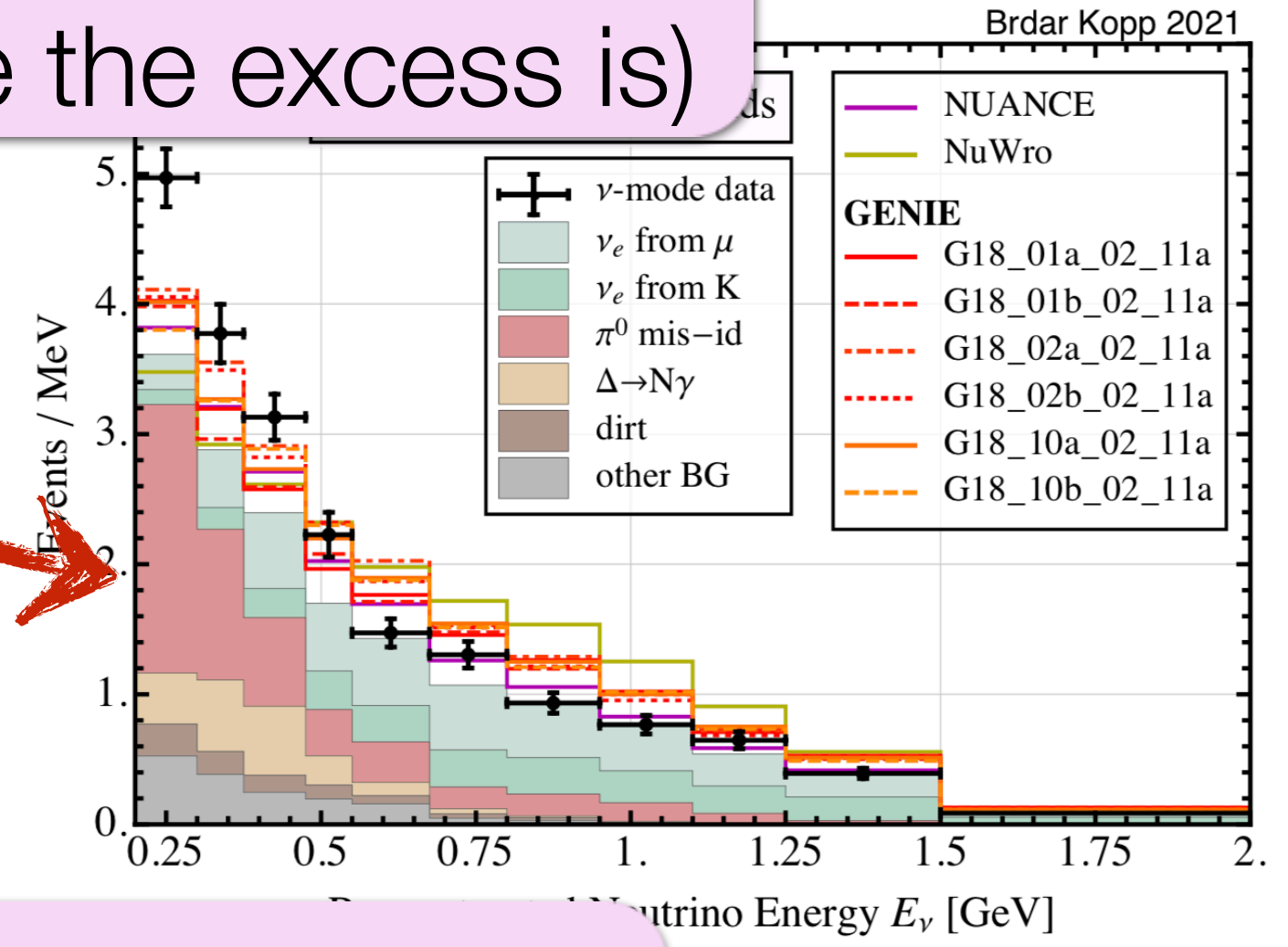
$\nu_s$  mass (squared)



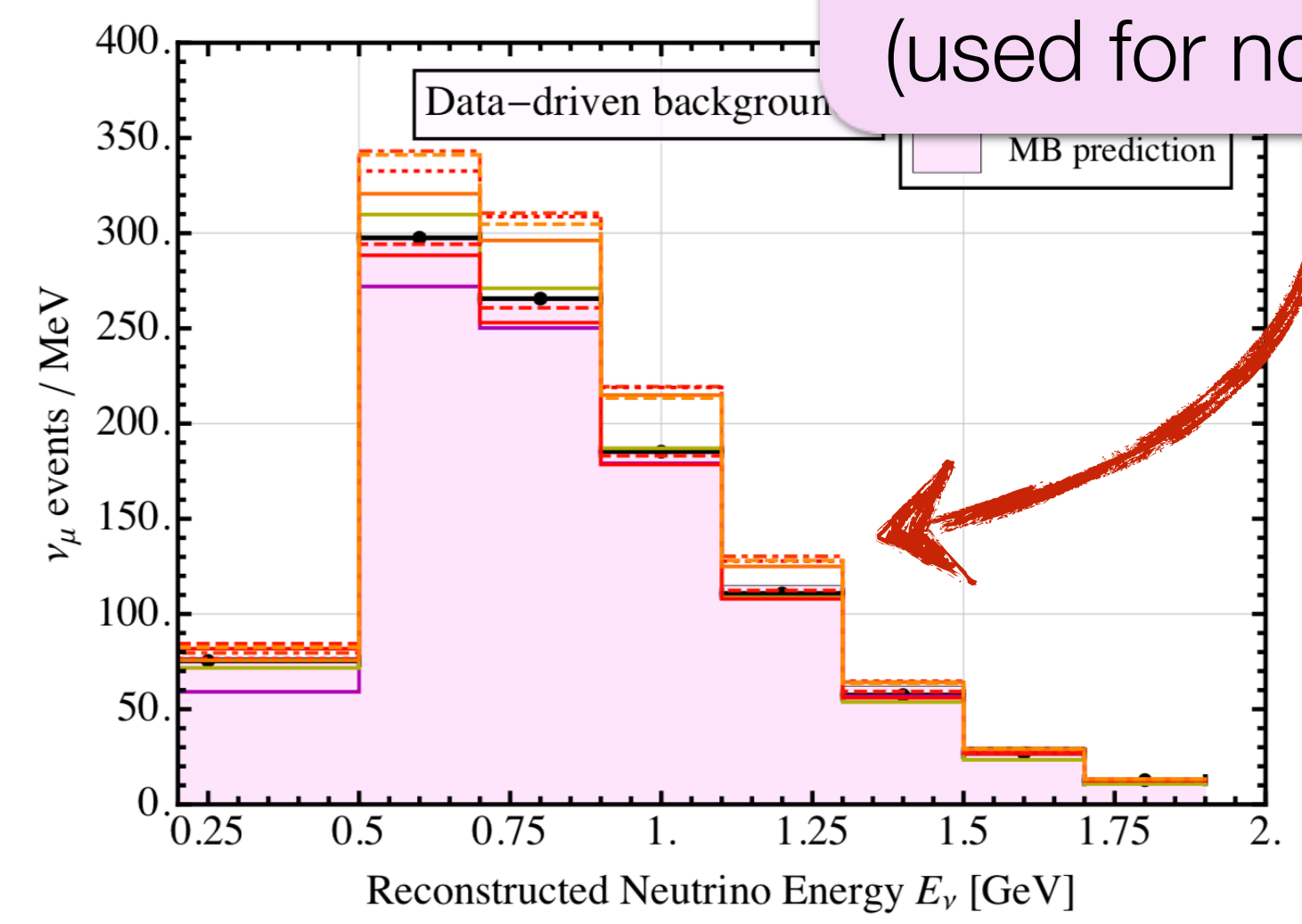


# 3+1 Models in MiniBooNE – Comparison of Generators

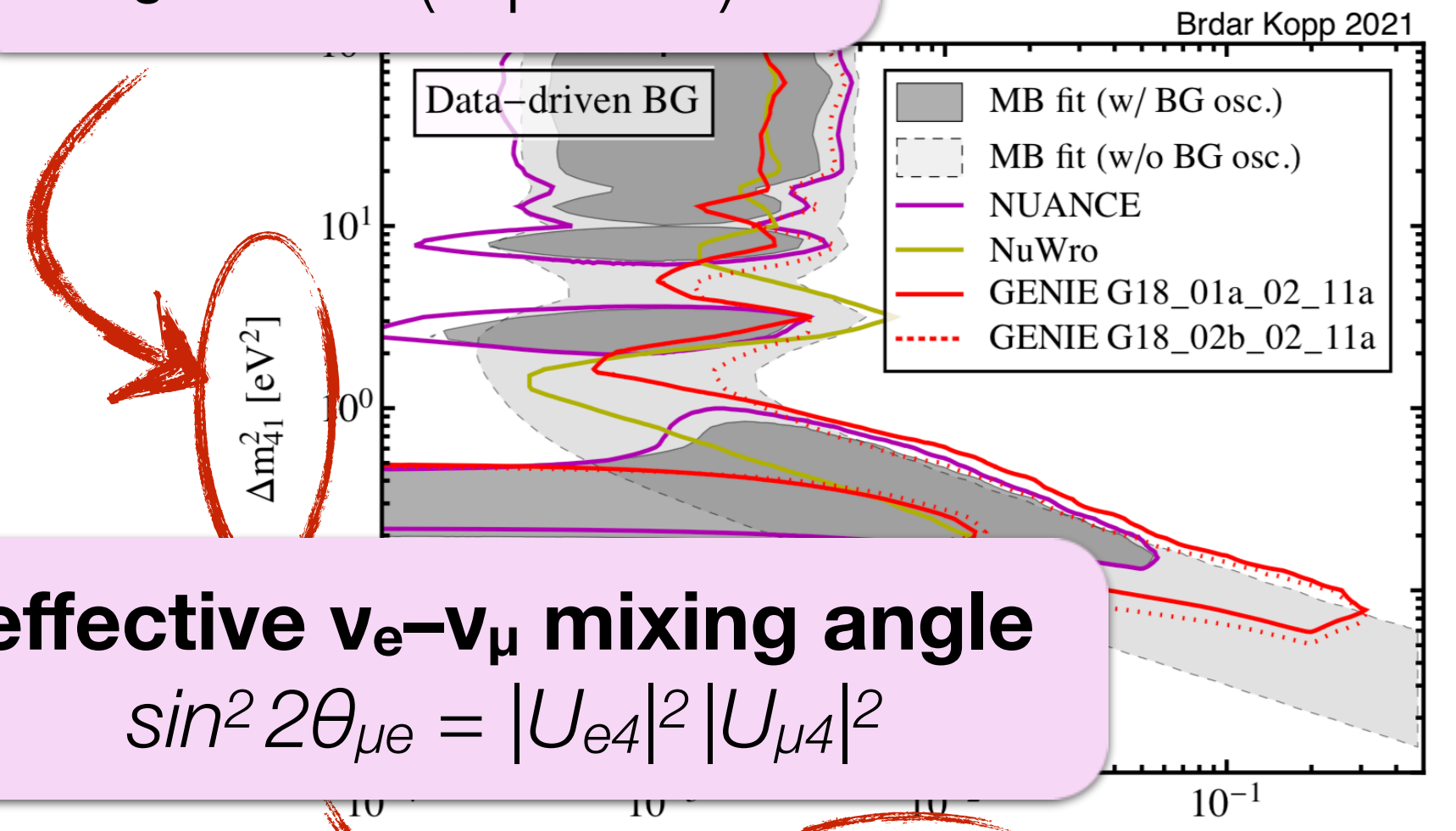
**$\nu_e$  spectrum**  
(where the excess is)



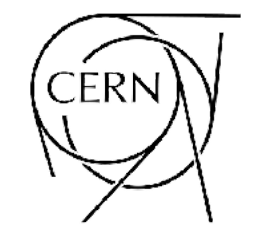
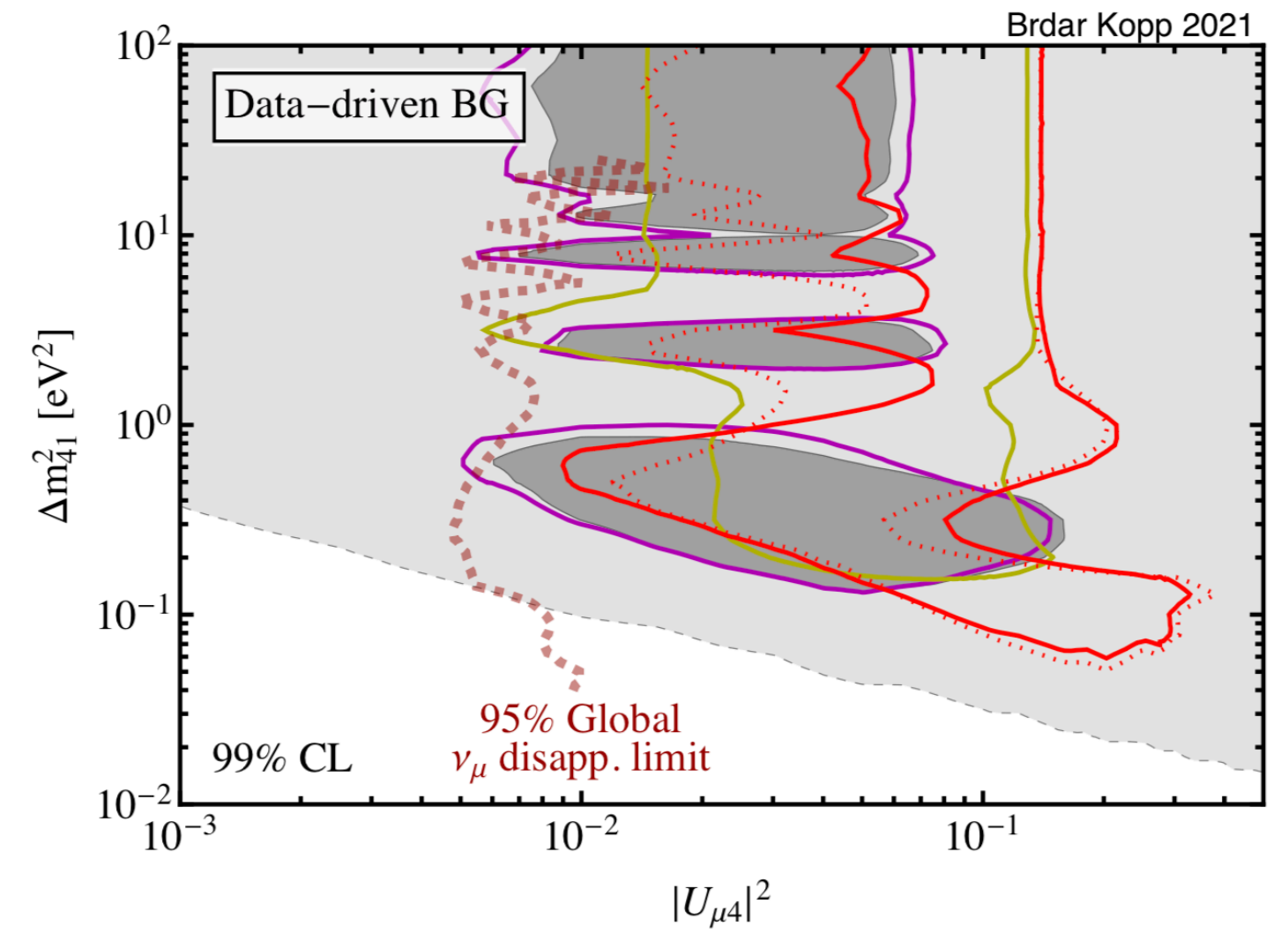
**$\nu_\mu$  spectrum**  
(used for normalization)



**$\nu_s$  mass (squared)**

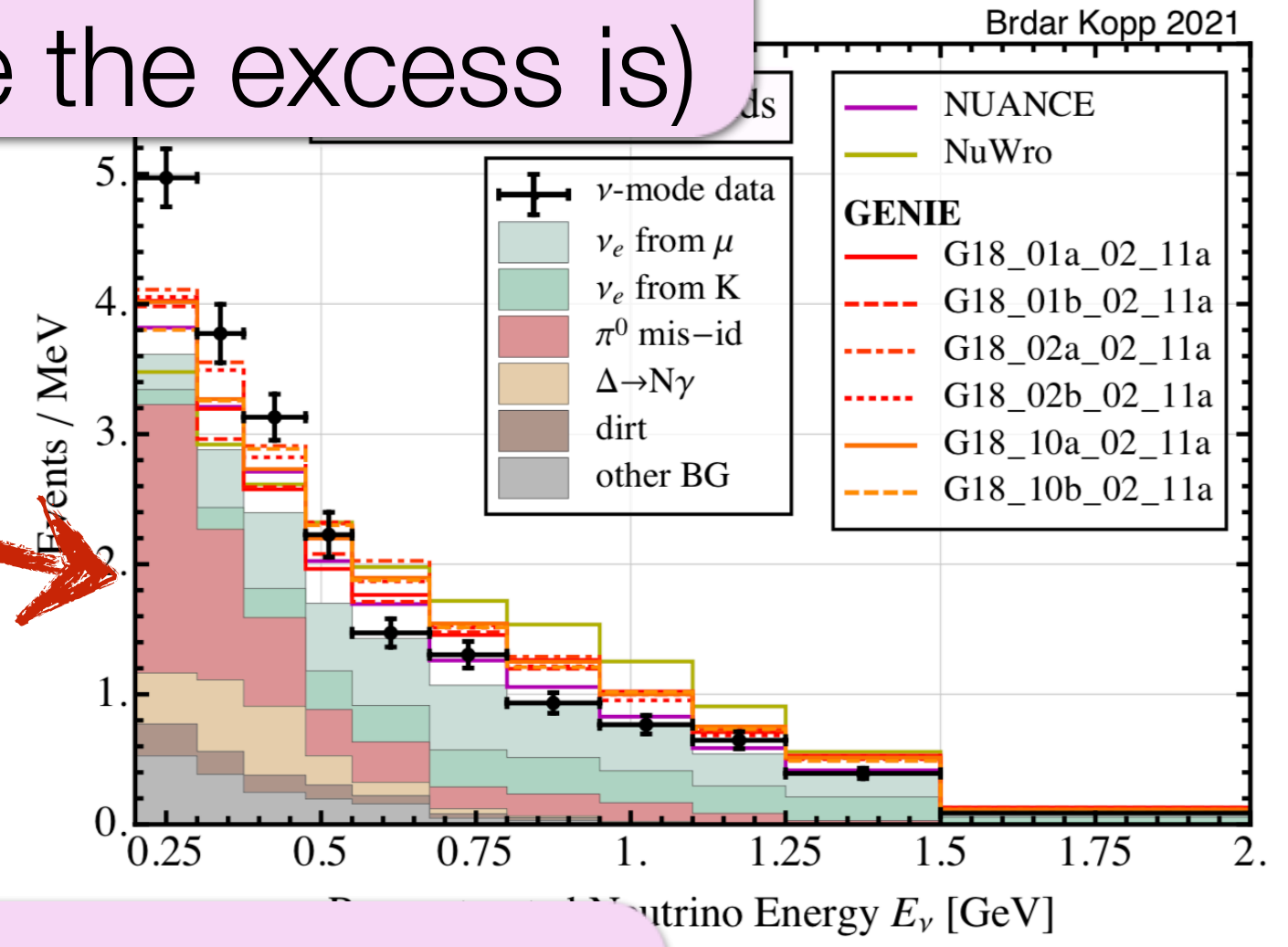


**effective  $\nu_e$ - $\nu_\mu$  mixing angle**  
 $\sin^2 2\theta_{\mu e} = |U_{e4}|^2 |U_{\mu 4}|^2$

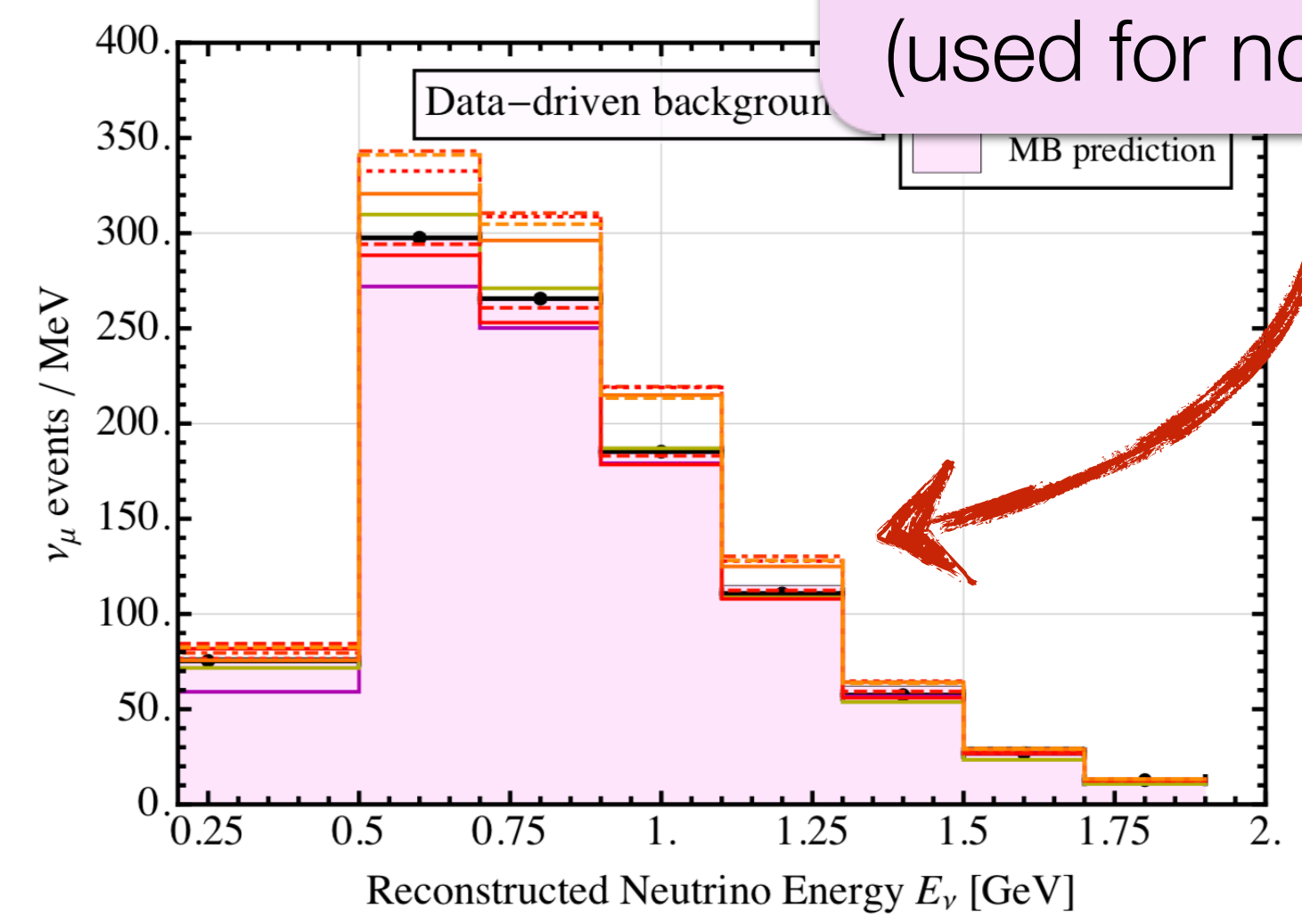


# 3+1 Models in MiniBooNE – Comparison of Generators

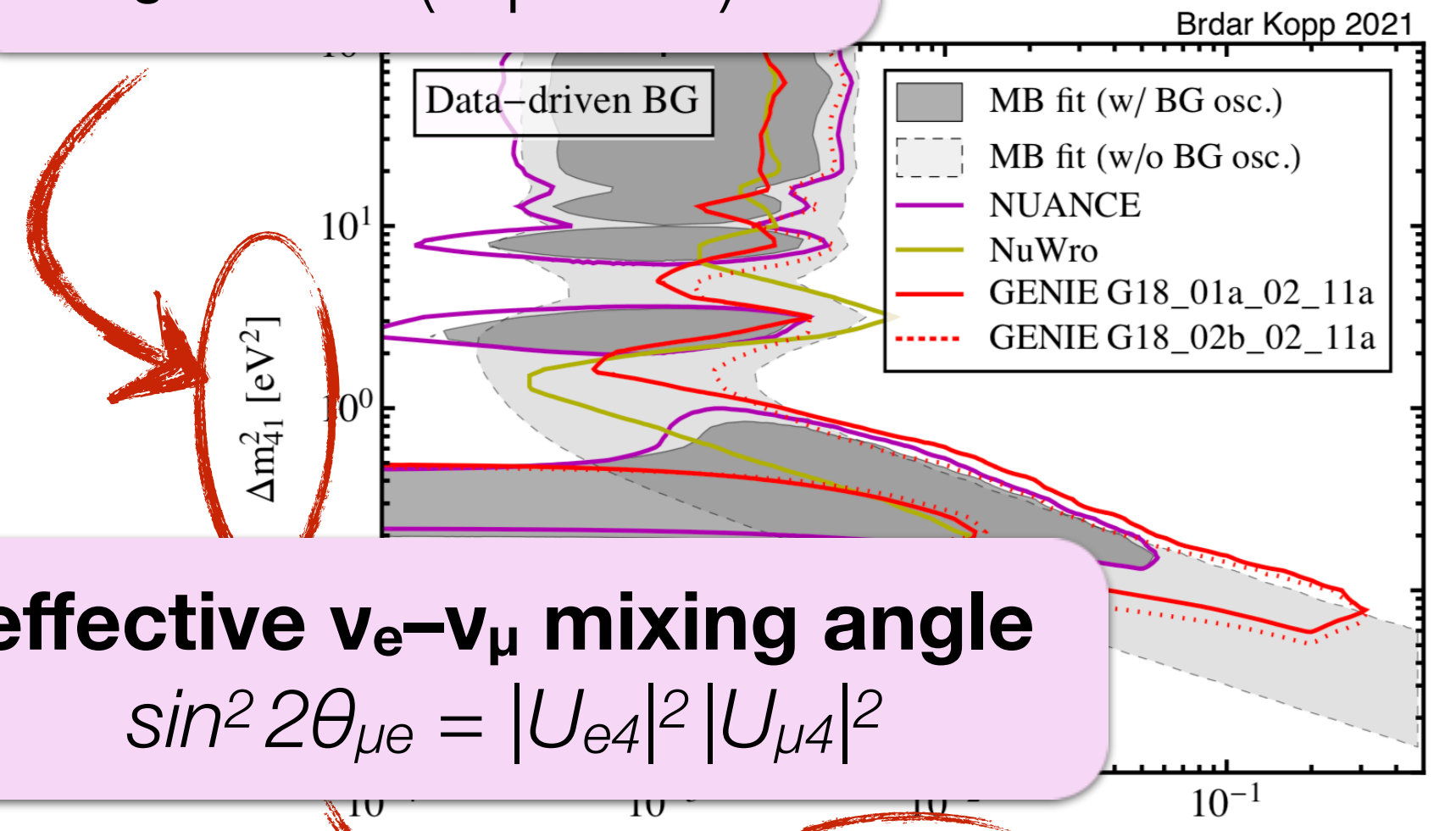
**$\nu_e$  spectrum**  
(where the excess is)



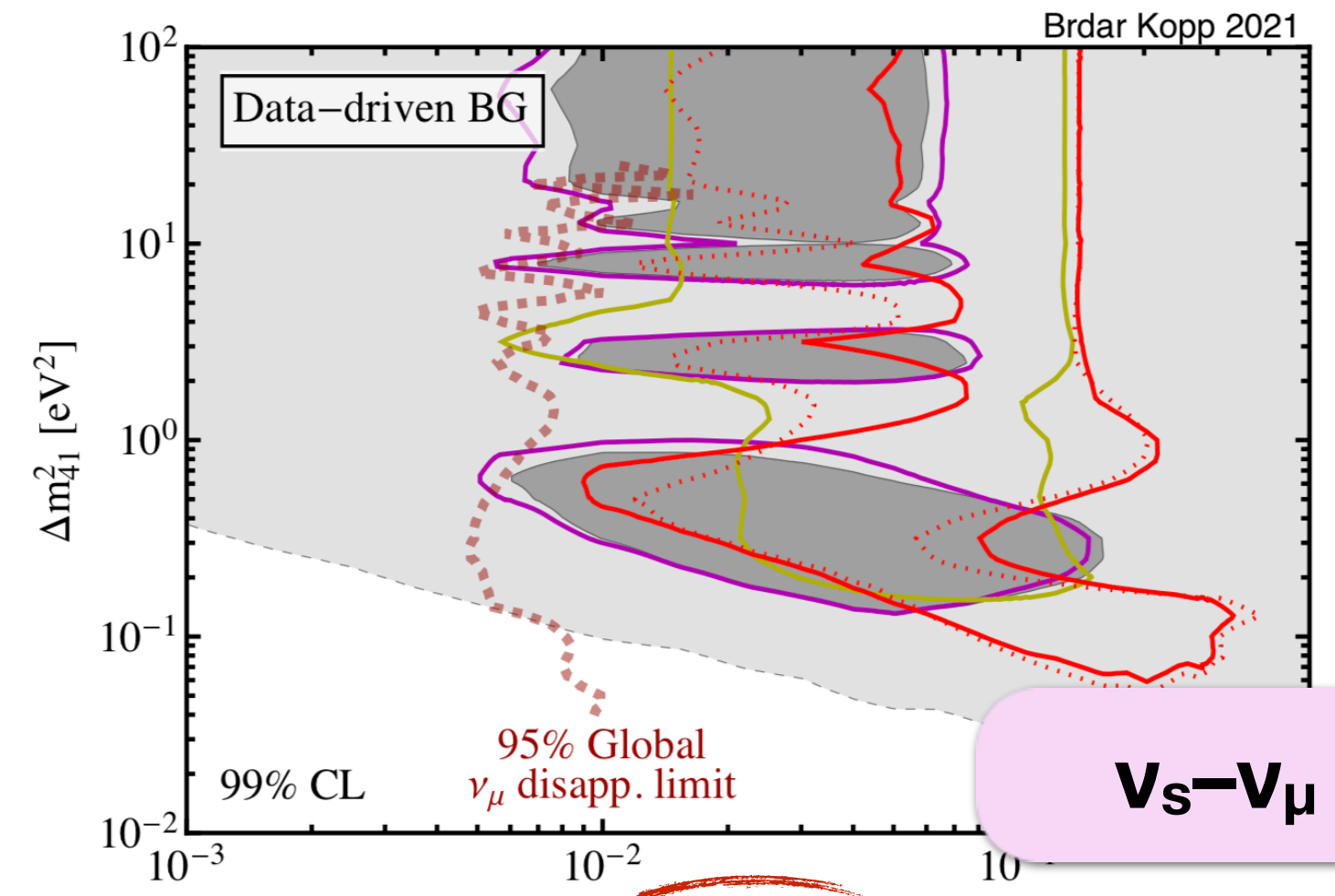
**$\nu_\mu$  spectrum**  
(used for normalization)



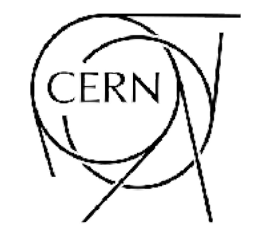
**$\nu_s$  mass (squared)**



**effective  $\nu_e$ - $\nu_\mu$  mixing angle**  
 $\sin^2 2\theta_{\mu e} = |U_{e4}|^2 |U_{\mu 4}|^2$



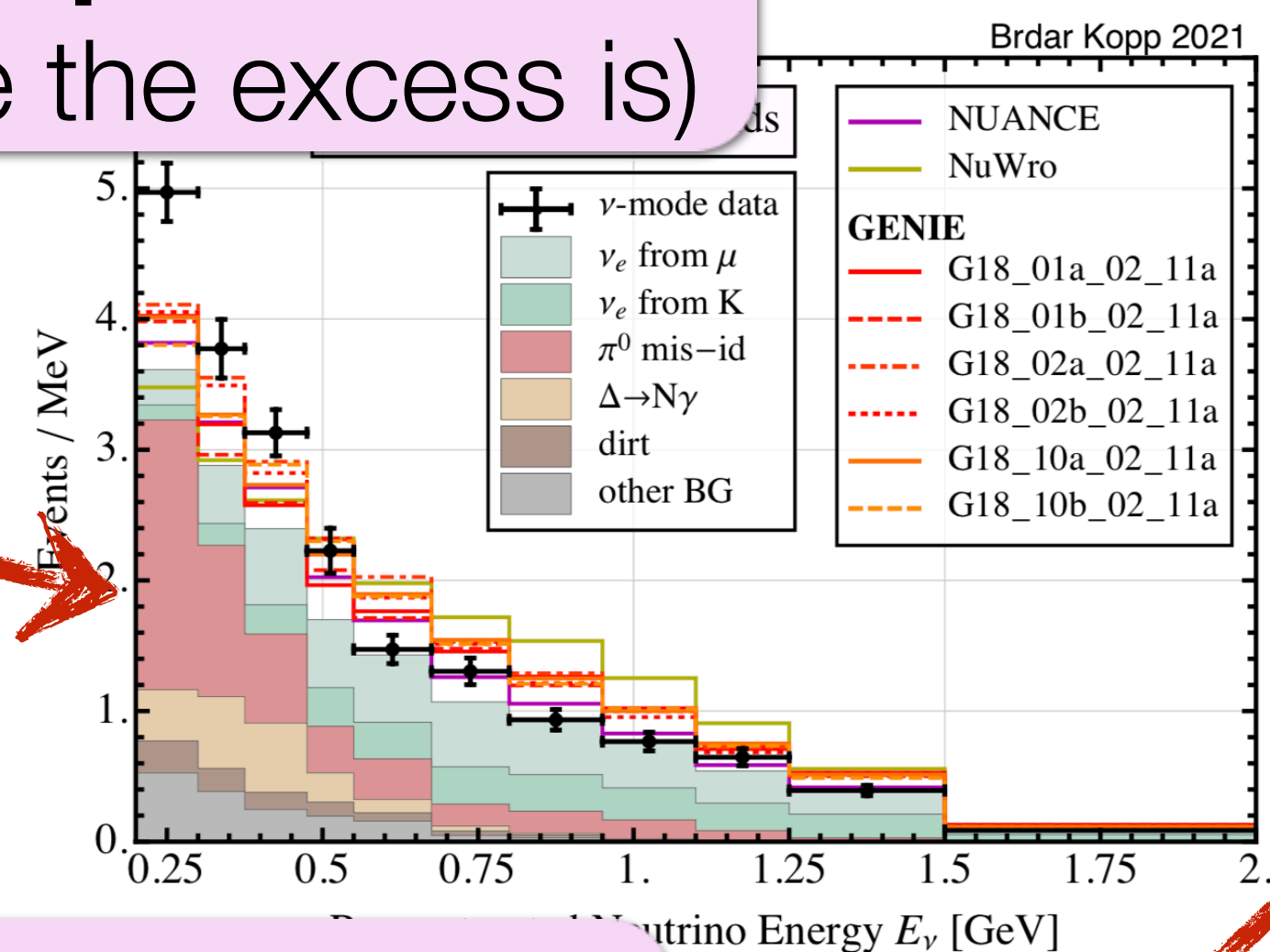
**$\nu_s$ - $\nu_\mu$  mixing**



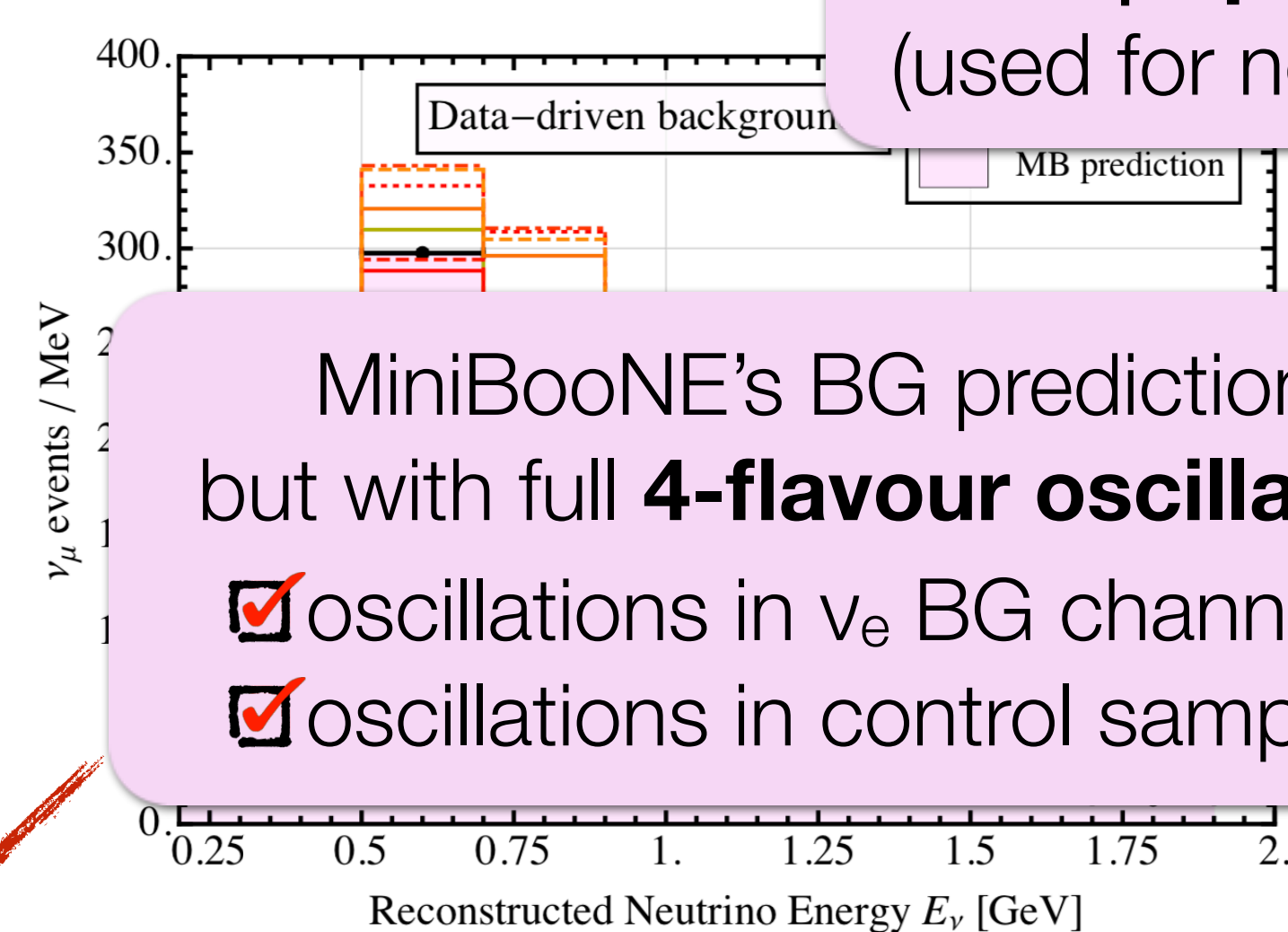


# 3+1 Models in MiniBooNE – Comparison of Generators

**$\nu_e$  spectrum**  
(where the excess is)



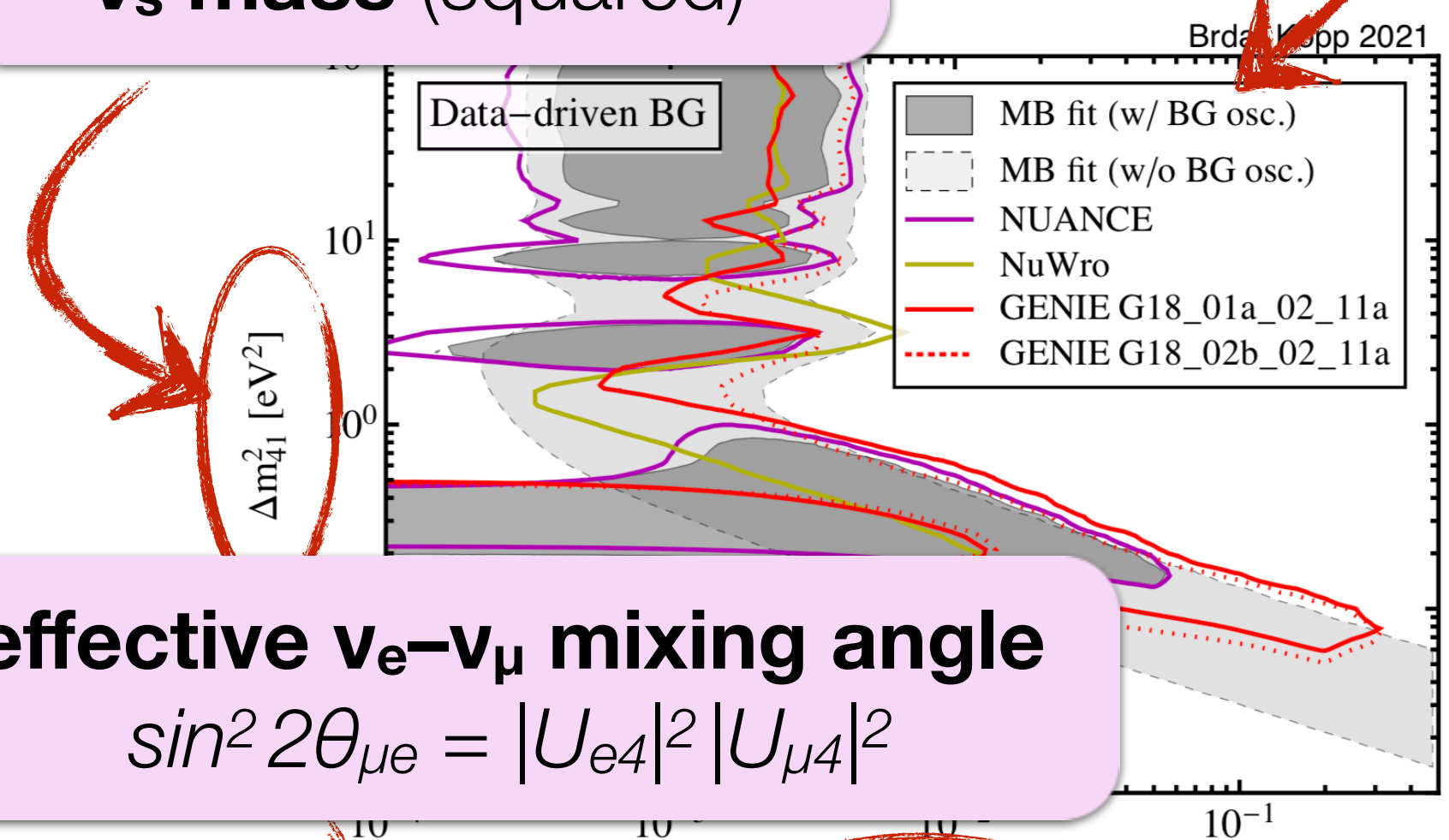
**$\nu_\mu$  spectrum**  
(used for normalization)



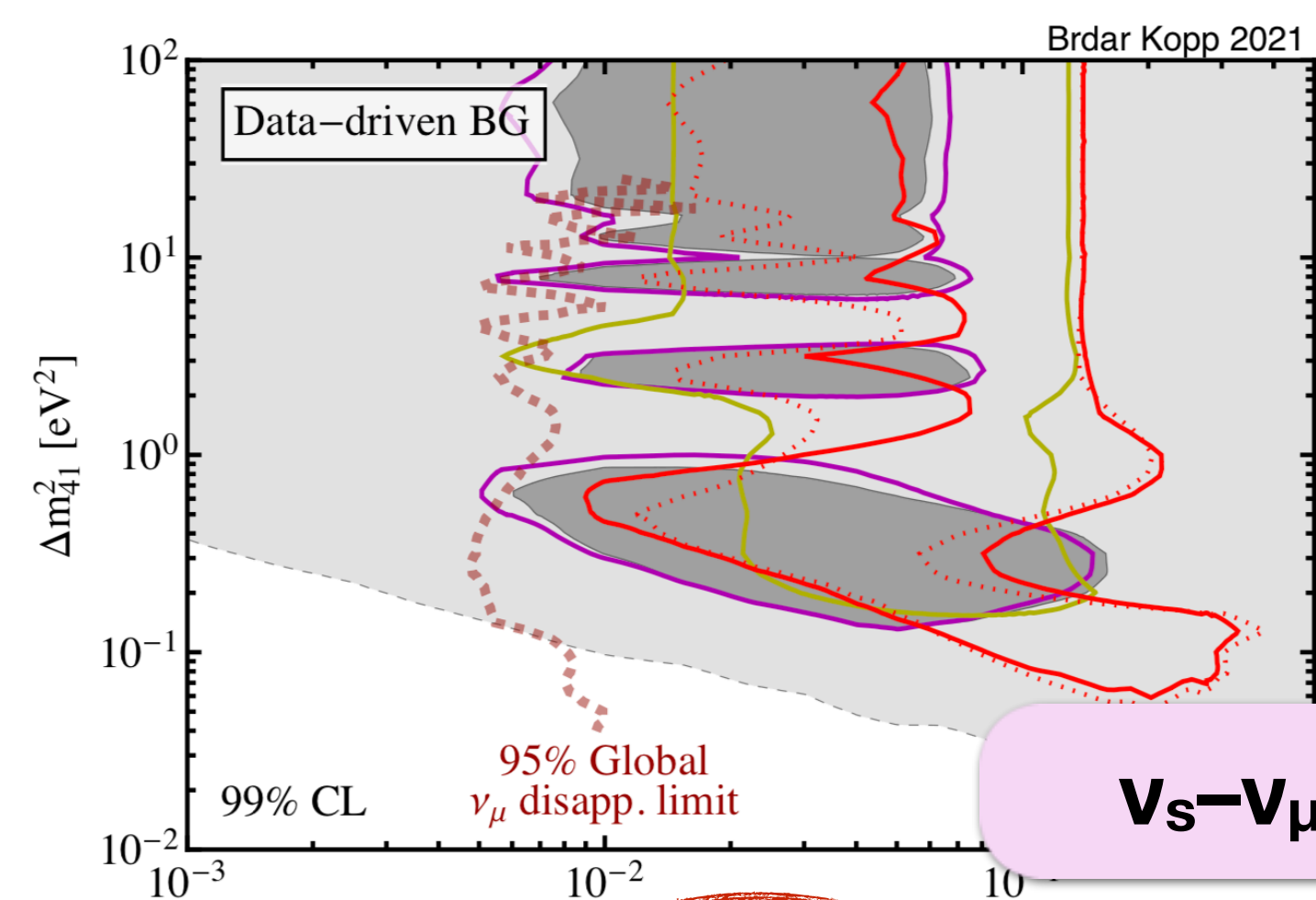
MiniBooNE's BG predictions, but with full **4-flavour oscillations**

- oscillations in  $\nu_e$  BG channels
- oscillations in control sample

**$\nu_s$  mass (squared)**



**effective  $\nu_e$ - $\nu_\mu$  mixing angle**  
 $\sin^2 2\theta_{\mu e} = |U_{e4}|^2 |U_{\mu4}|^2$

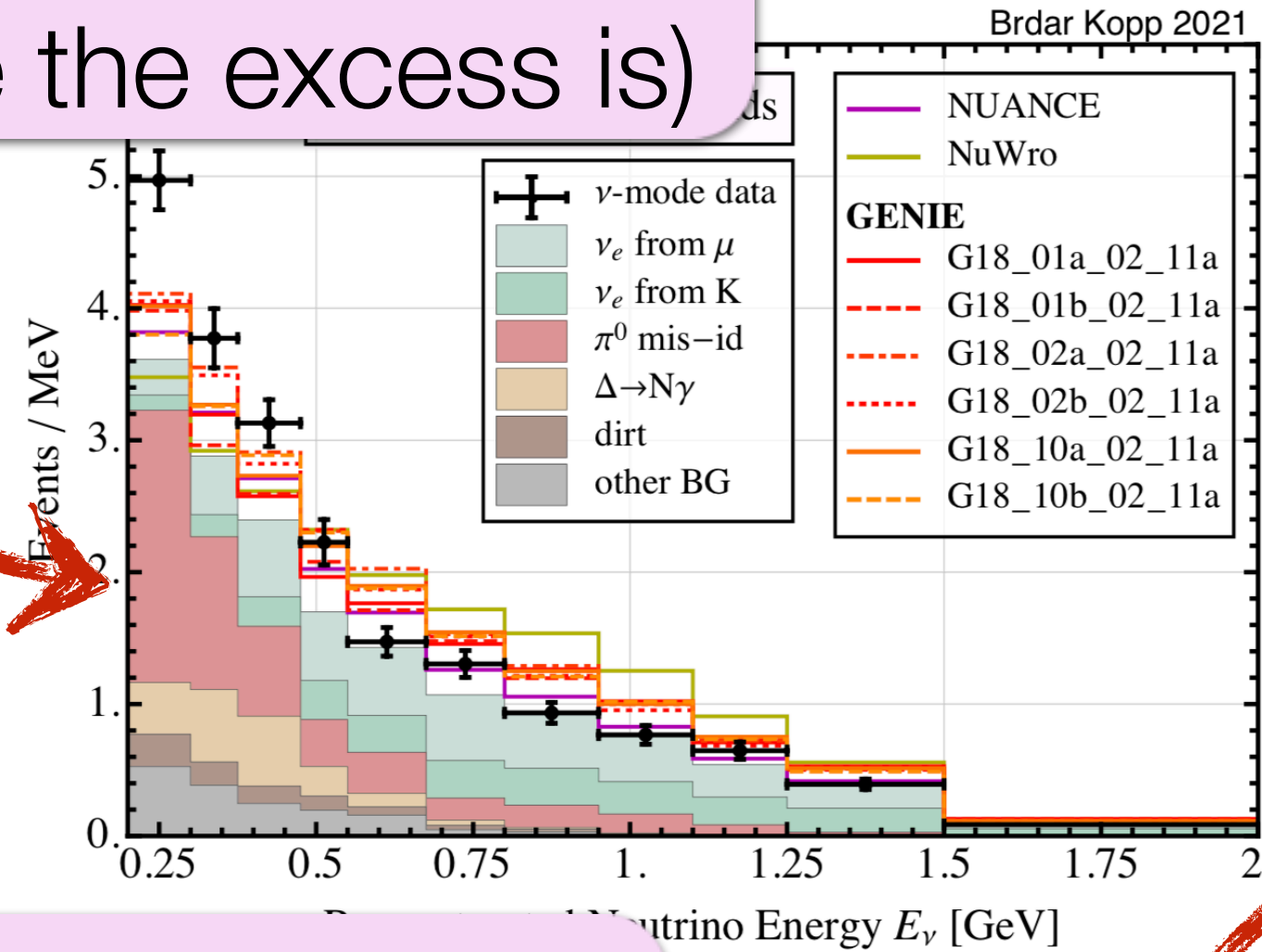


**$\nu_s$ - $\nu_\mu$  mixing**

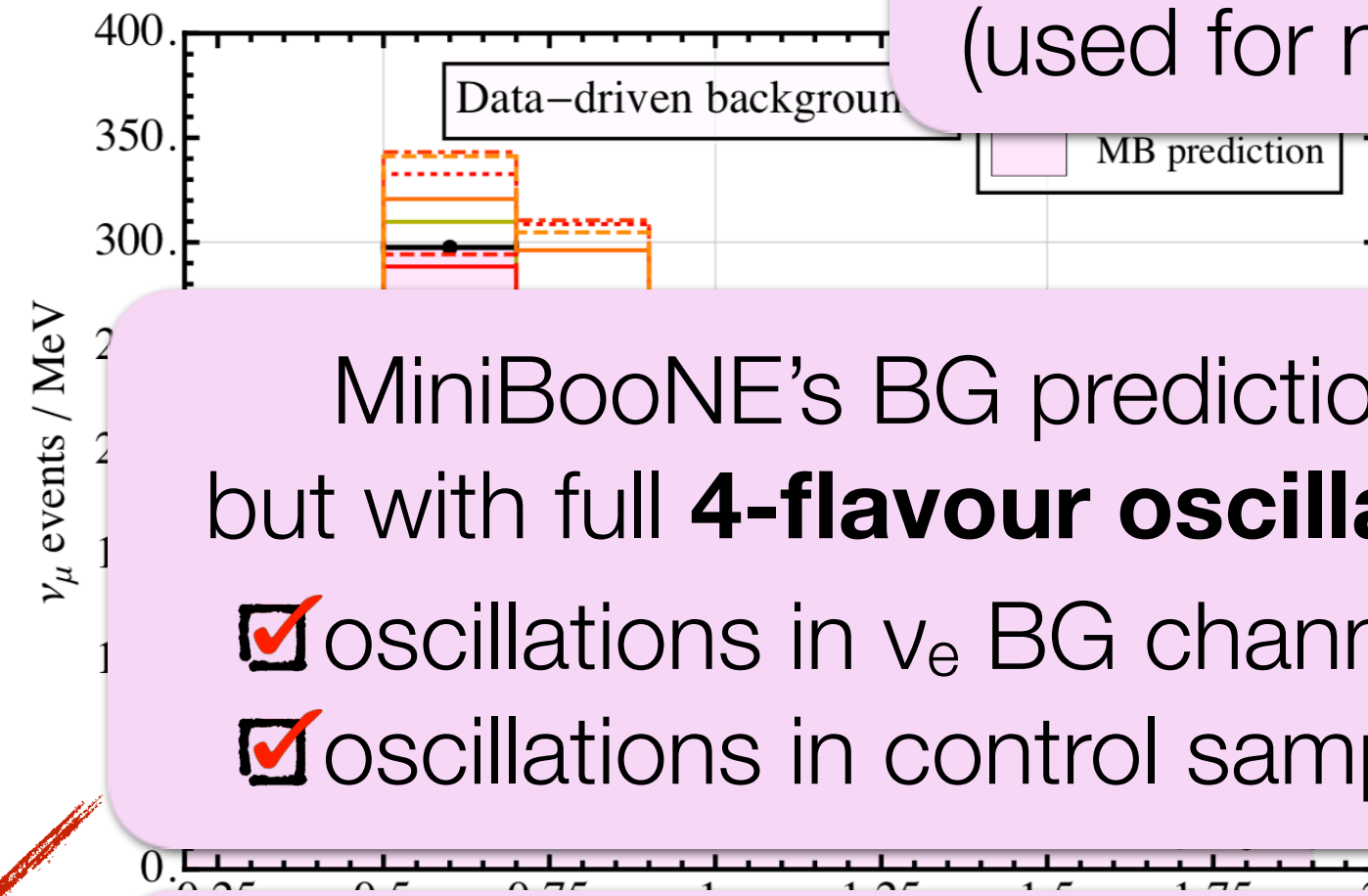


# 3+1 Models in MiniBooNE – Comparison of Generators

$\nu_e$  spectrum  
(where the excess is)



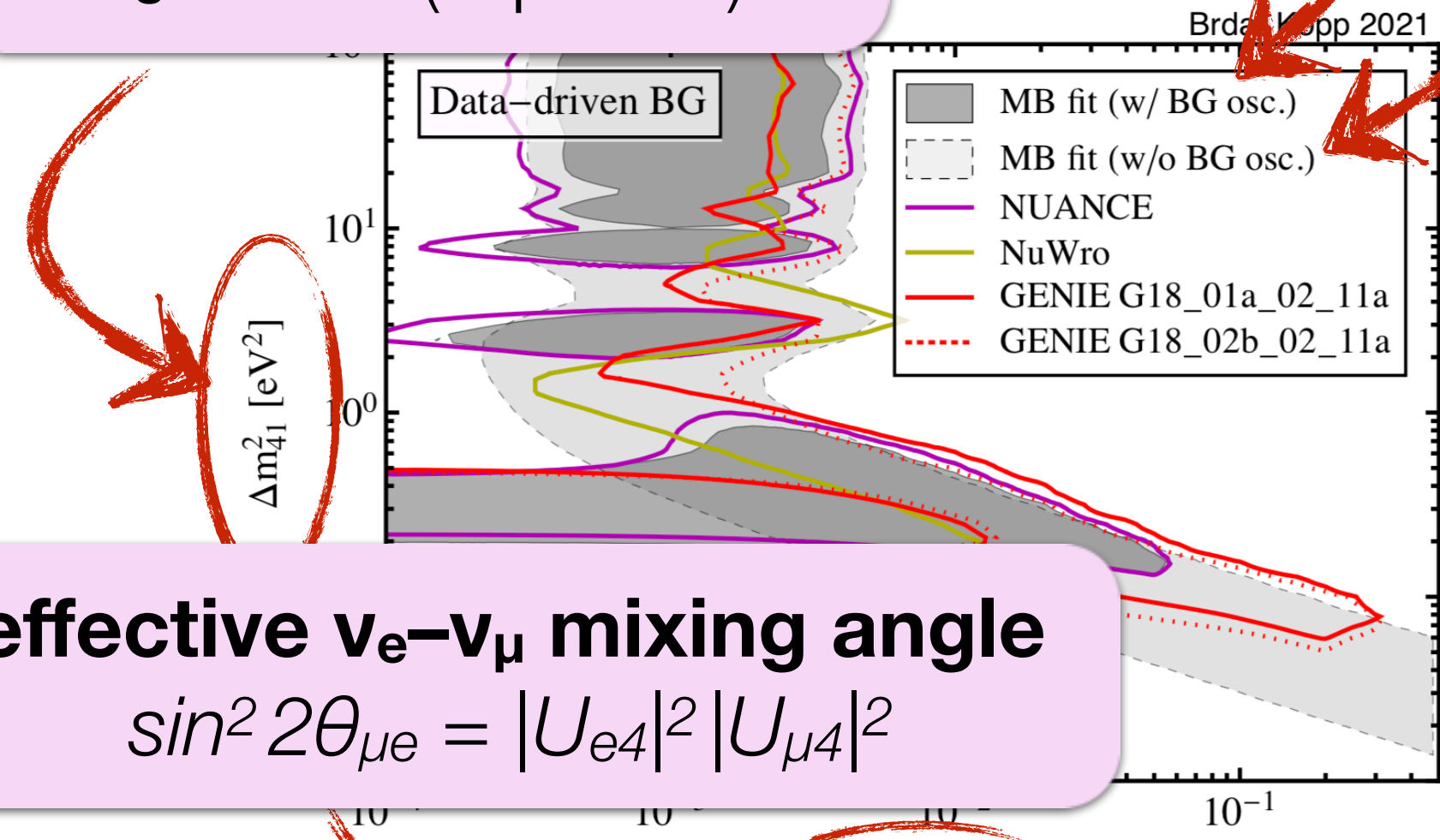
$\nu_\mu$  spectrum  
(used for normalization)



MiniBooNE's BG predictions, but with full **4-flavour oscillations**

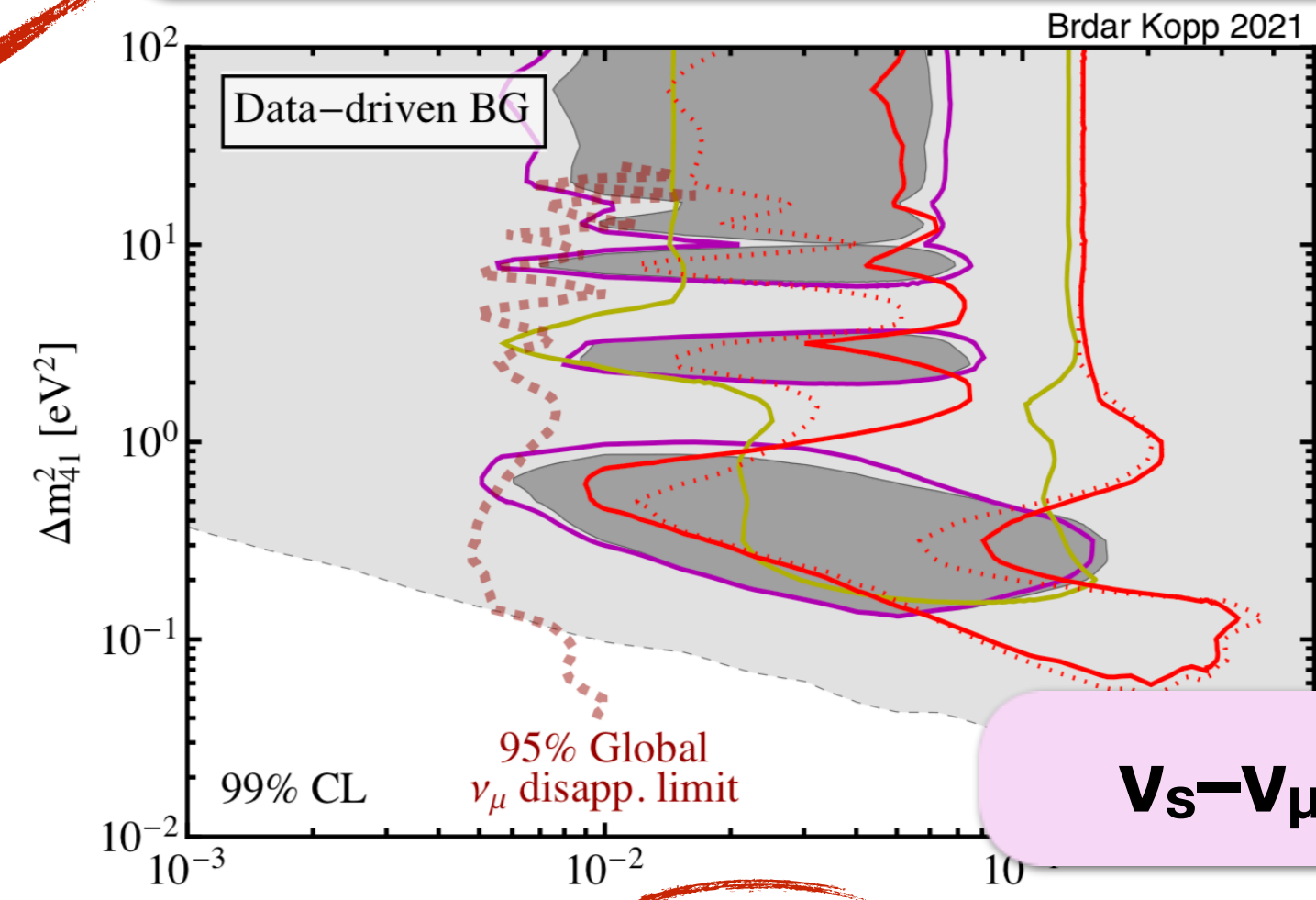
- oscillations in  $\nu_e$  BG channels
- oscillations in control sample

$\nu_s$  mass (squared)

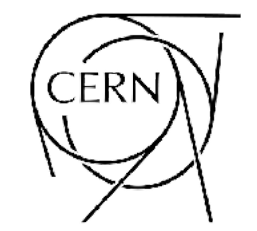


effective  $\nu_e$ - $\nu_\mu$  mixing angle  
 $\sin^2 2\theta_{\mu e} = |U_{e4}|^2 |U_{\mu 4}|^2$

MiniBooNE's fit (2-flavour oscillations)



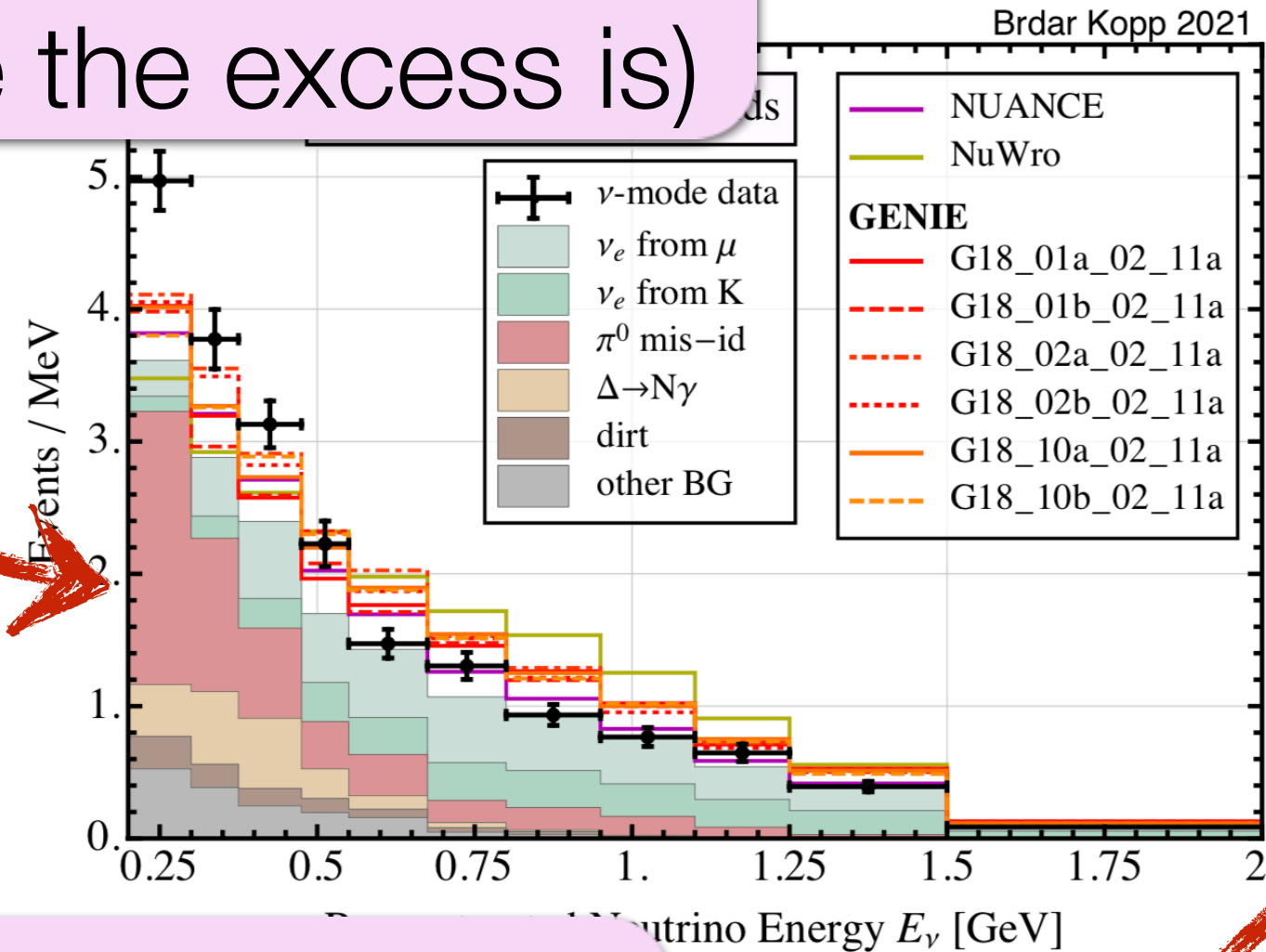
$\nu_s$ - $\nu_\mu$  mixing



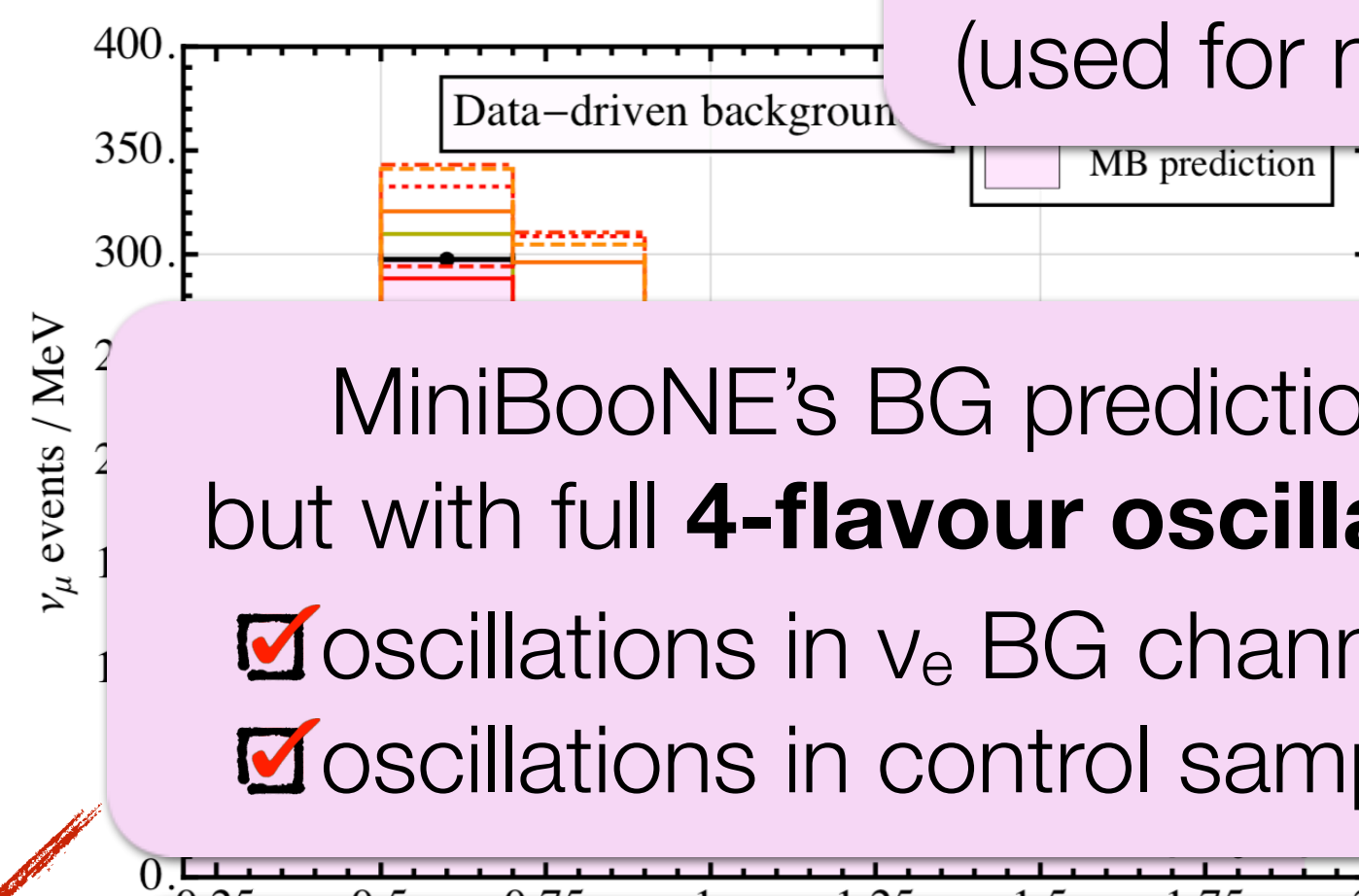


# 3+1 Models in MiniBooNE – Comparison of Generators

**$\nu_e$  spectrum**  
(where the excess is)



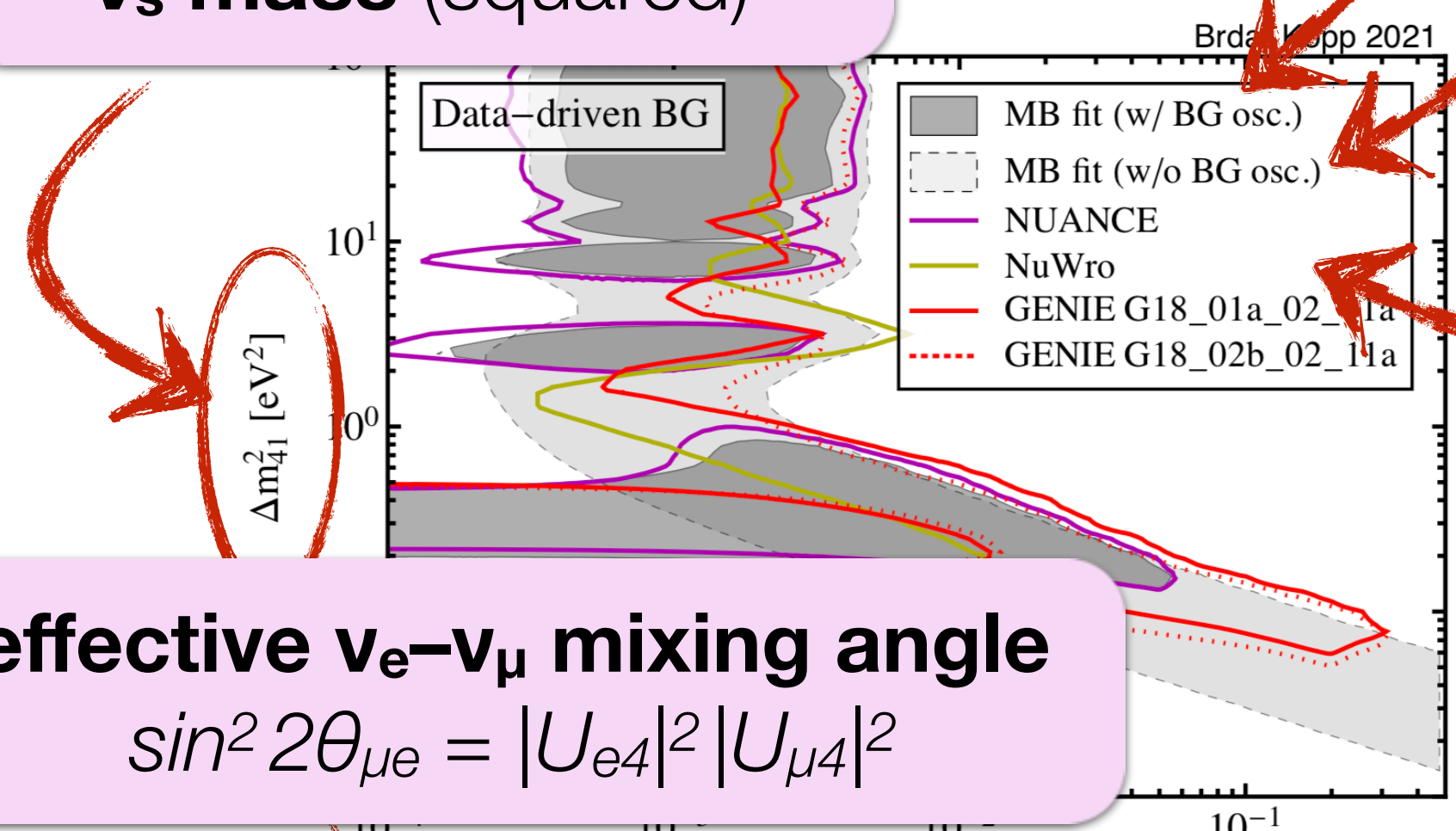
**$\nu_\mu$  spectrum**  
(used for normalization)



MiniBooNE's BG predictions, but with full **4-flavour oscillations**

- oscillations in  $\nu_e$  BG channels
- oscillations in control sample

**$\nu_s$  mass (squared)**



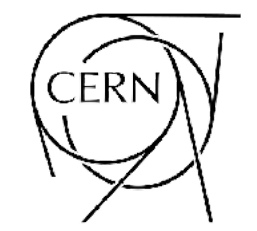
**effective  $\nu_e$ - $\nu_\mu$  mixing angle**  
 $\sin^2 2\theta_{\mu e} = |U_{e4}|^2 |U_{\mu 4}|^2$

**MiniBooNE's fit (2-flavour oscillations)**

using our own BG predictions:

- NUANCE:  $4\sigma$
- NuWro:  $3.1\sigma$
- G18\_01a\_02\_11a:  $4.4\sigma$
- G18\_02b\_02\_11a:  $3.7\sigma$

**$\nu_s$ - $\nu_\mu$  mixing**

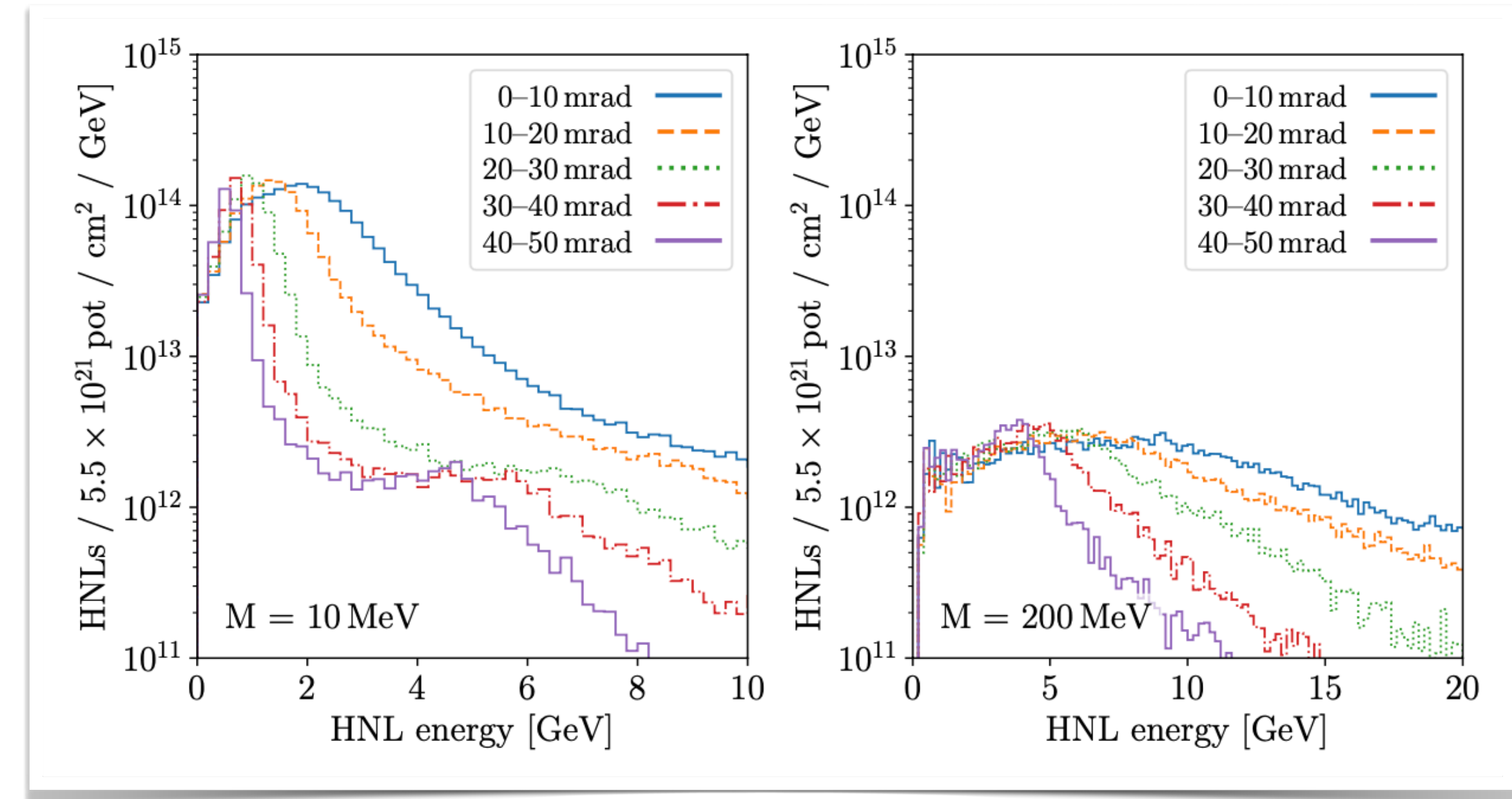


# Heavy Neutral Leptons

- Neutrino portal leads to mixing between  $\nu$  and heavy sterile neutrino  $N$

$$\mathcal{L} \supset y \bar{L} (i\sigma^2 H^*) N$$

- any process that makes  $\nu$  in the SM can also make  $N$  (suppressed by a mixing angle)
- production in meson decays



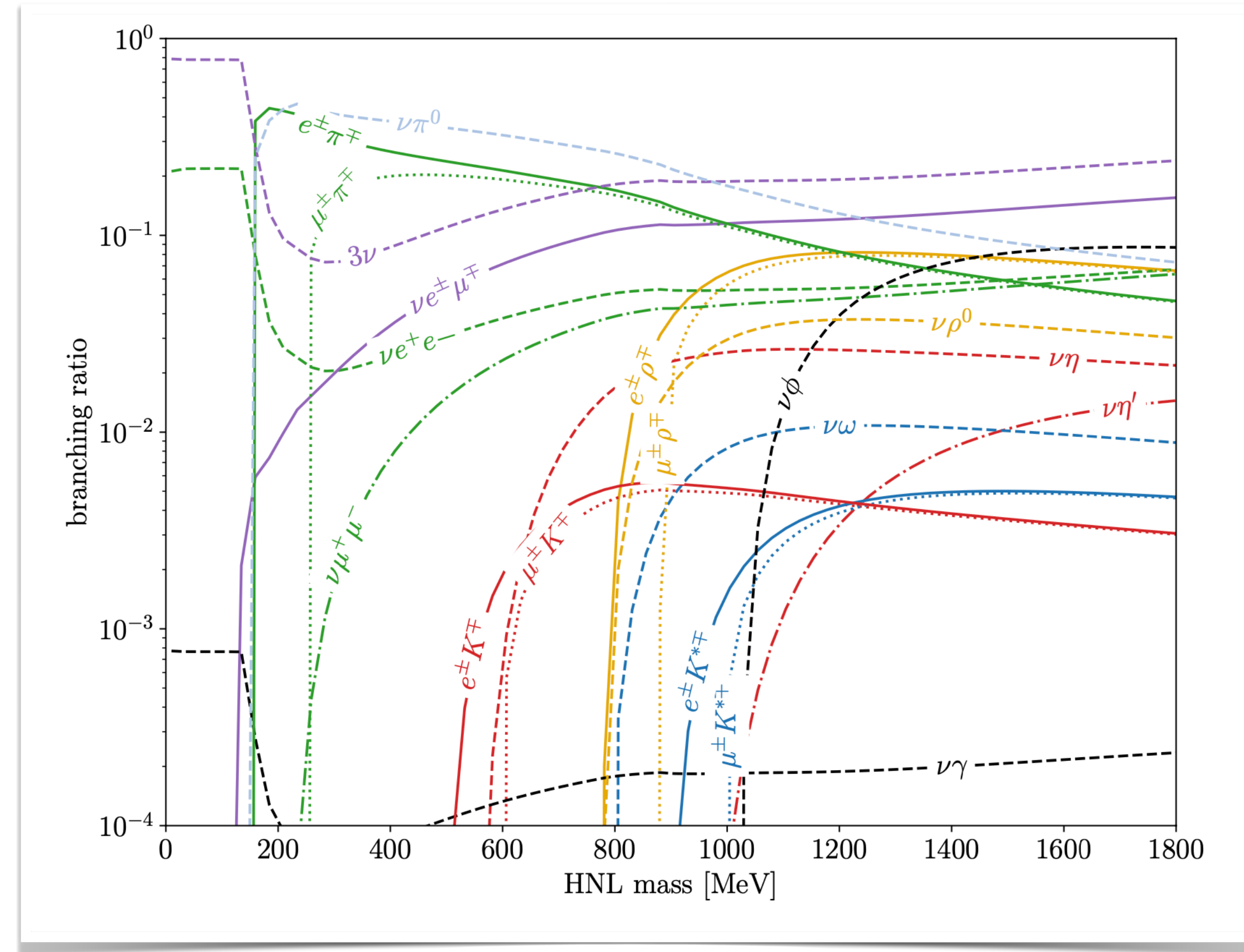


# Heavy Neutral Leptons

- Neutrino portal leads to mixing between  $\nu$  and heavy sterile neutrino  $N$

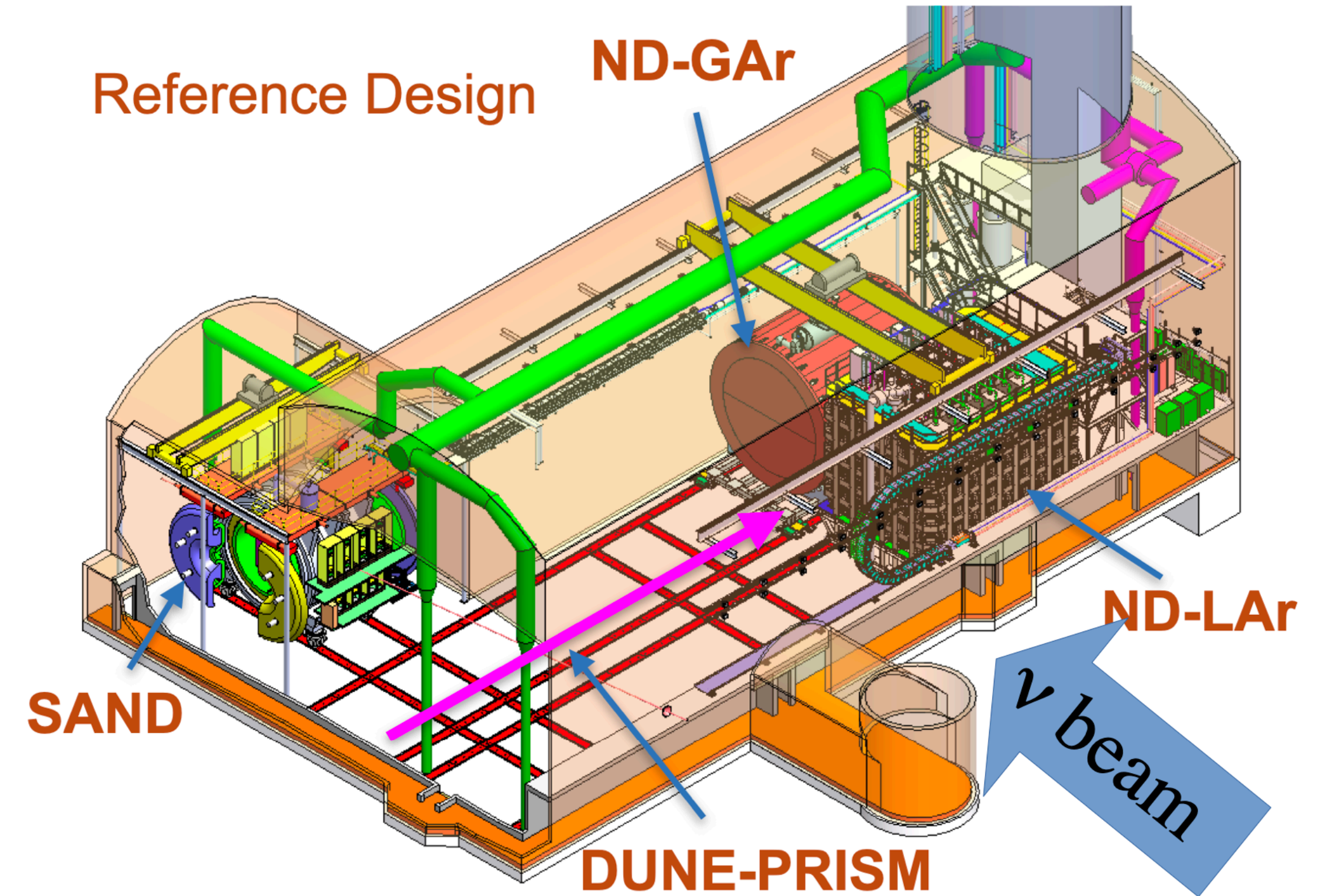
$$\mathcal{L} \supset y \bar{L} (i\sigma^2 H^*) N$$

- any process that makes  $\nu$  in the SM can also make  $N$  (suppressed by a mixing angle)
- production in meson decays
- many 2-body and 3-body decay channels
- background estimates require detailed understanding of **exclusive final states**



# DUNE-PRISM

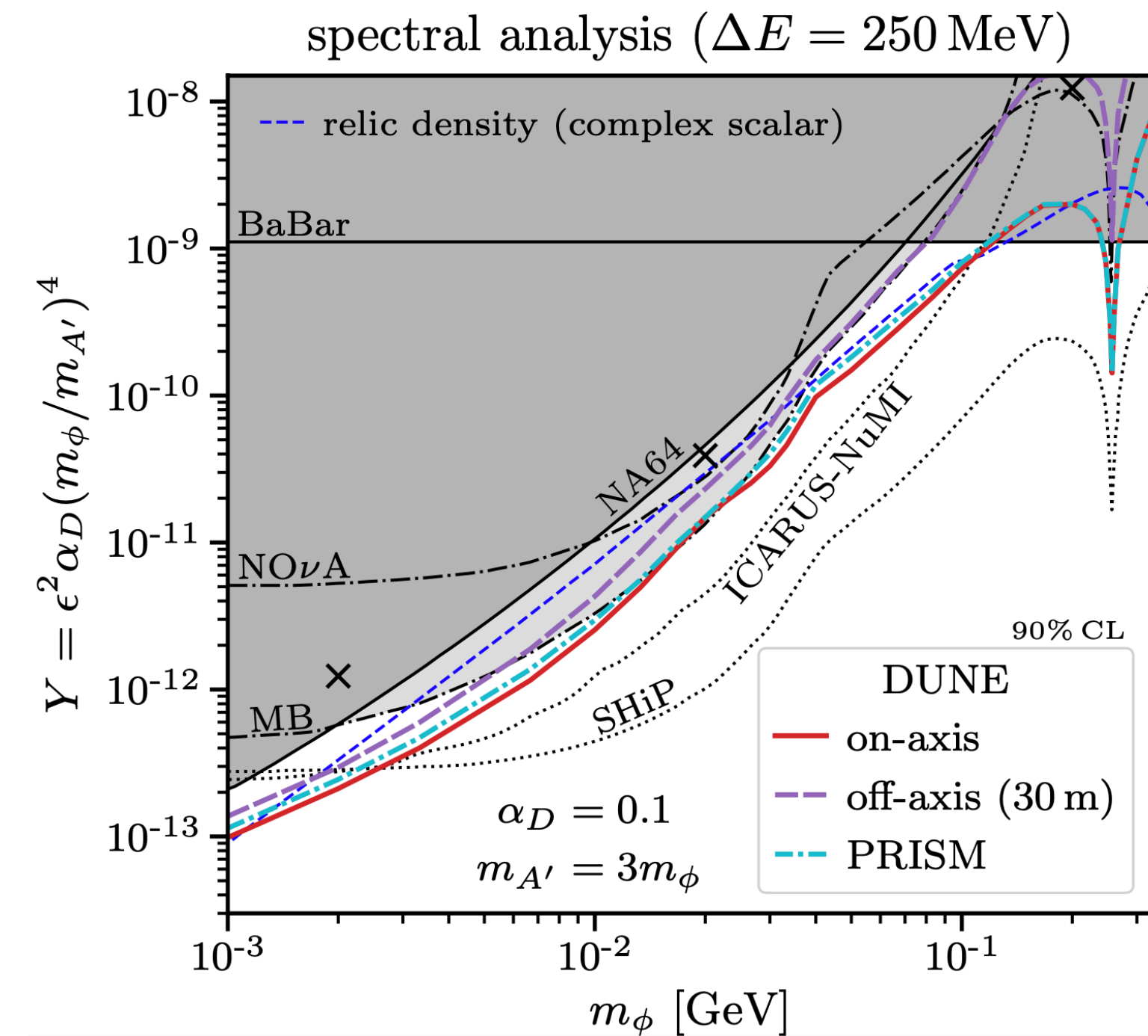
- Movable near detectors measure event rates at different off-axis angles
  - ▣ disentangle flux and  $x$ -sec uncertainties
- Interesting opportunities for BSM searches
  - heavy particle less boosted
  - backgrounds lower off-axis



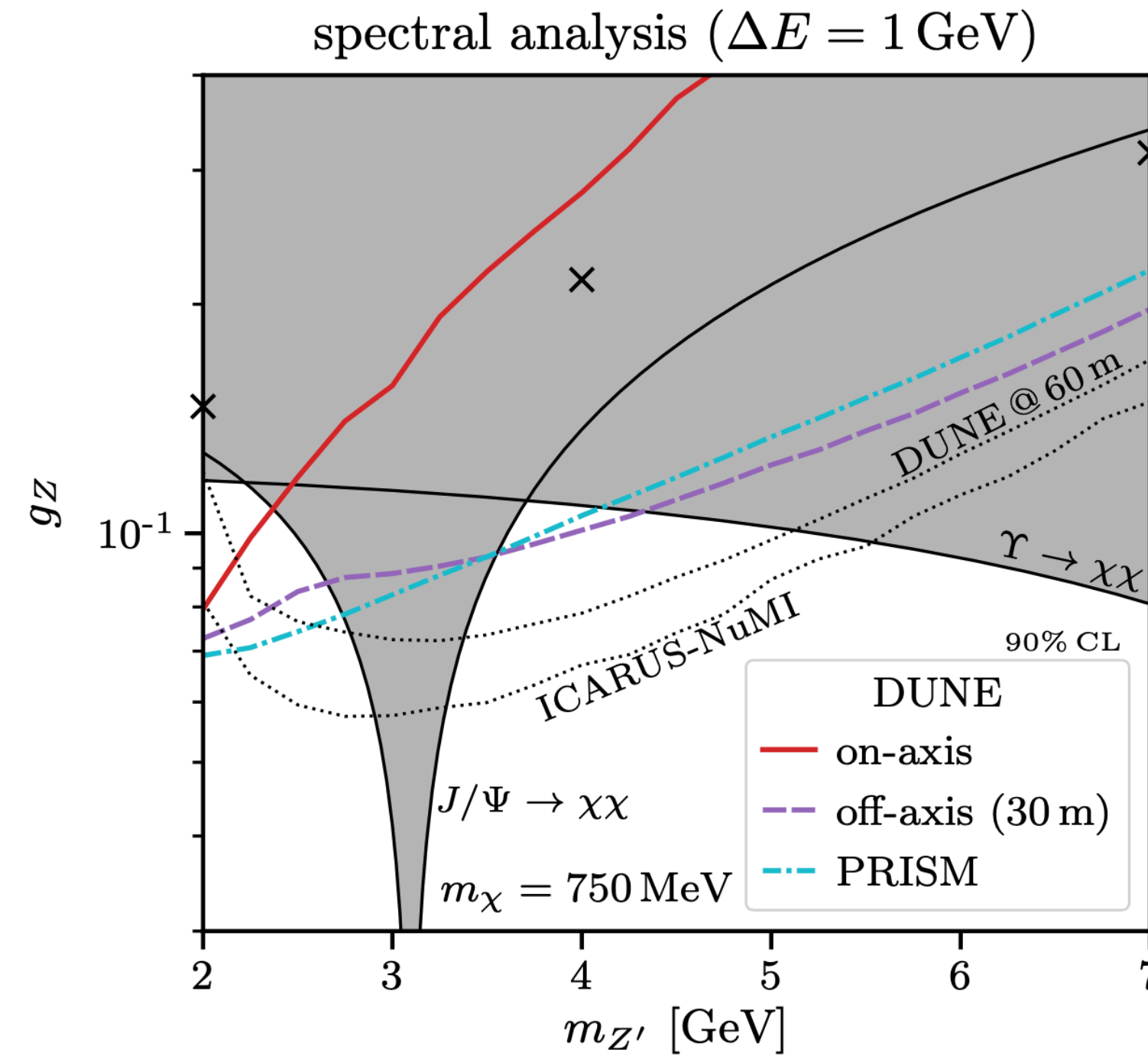


# DUNE-PRISM

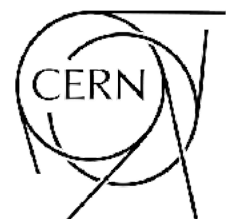
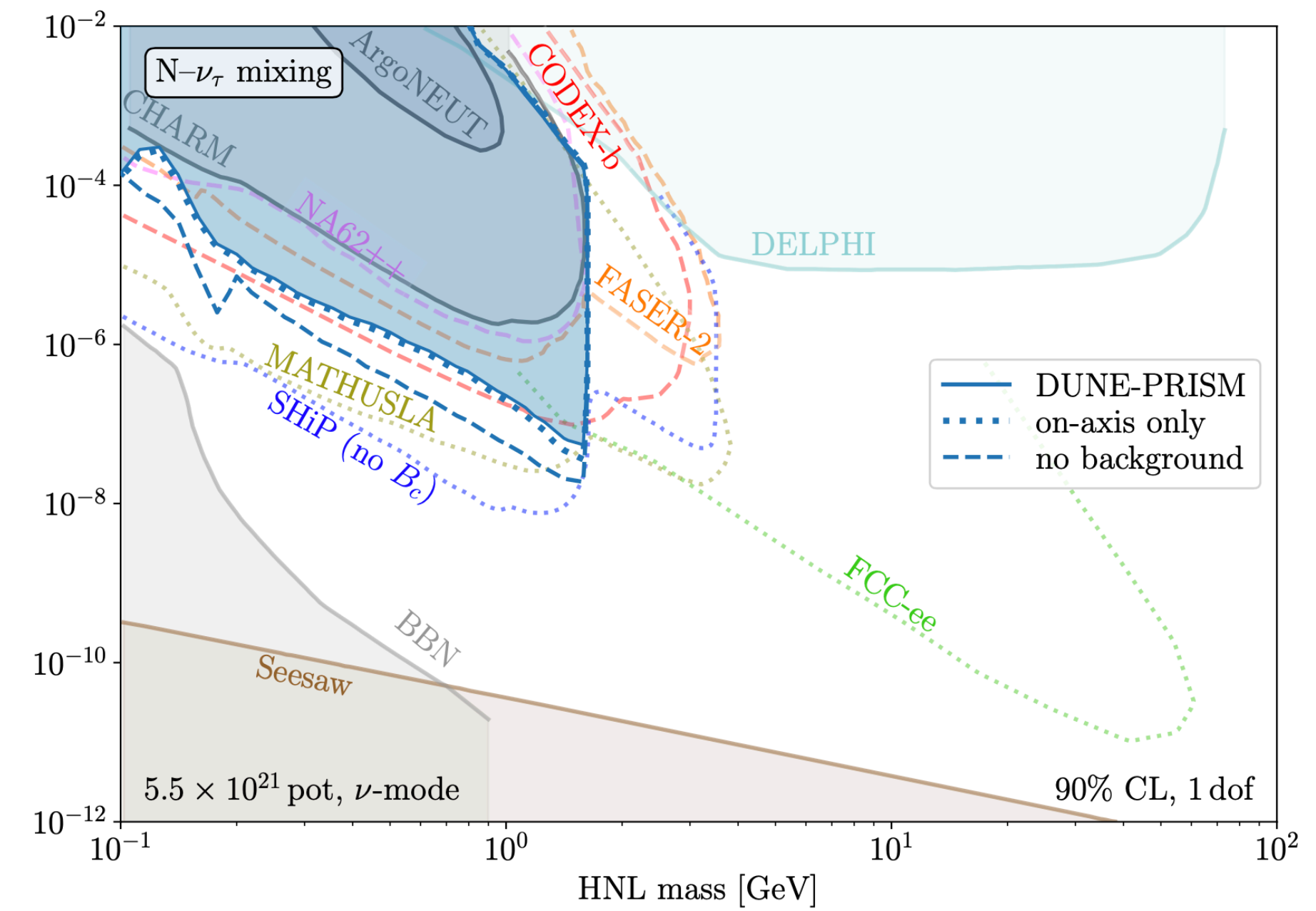
## A'-mediated DM



## leptophobic DM

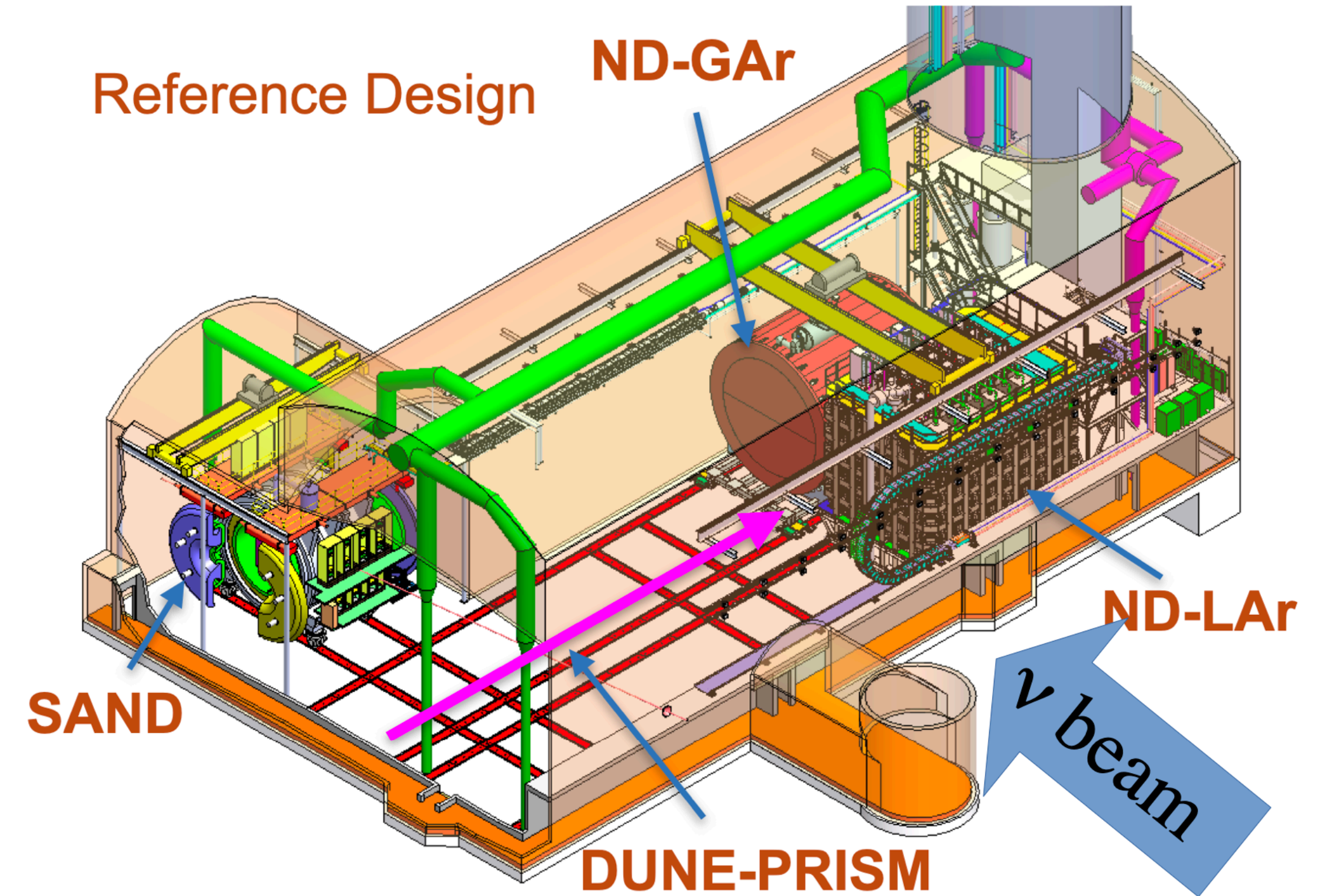


## heavy neutral leptons



# DUNE-PRISM

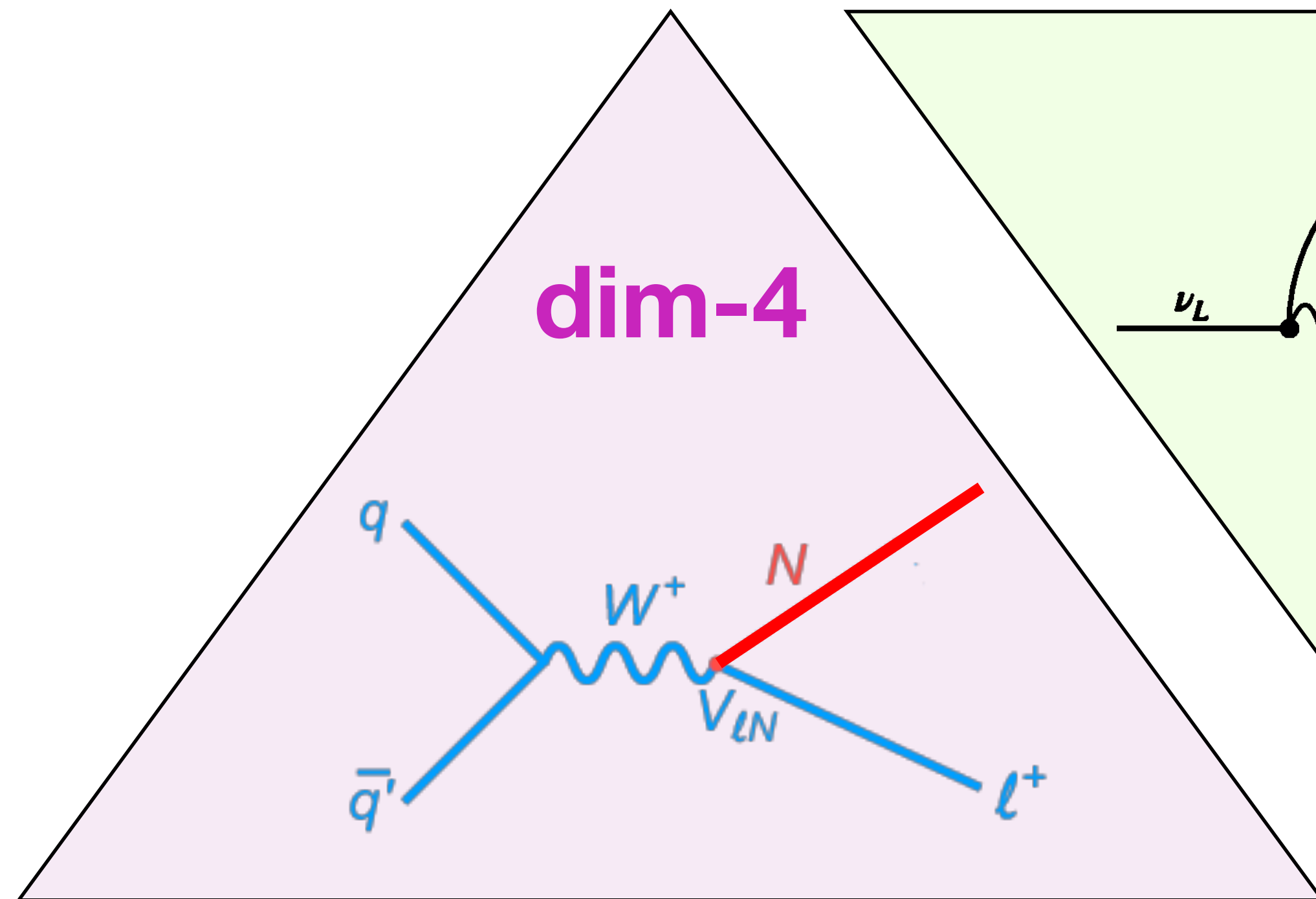
- Movable near detectors measure event rates at different off-axis angles
  - disentangle flux and  $x$ -sec uncertainties
- Interesting opportunities for BSM searches
  - heavy particle less boosted
  - backgrounds lower off-axis
- BSM searches are typically robust w.r.t. running strategy



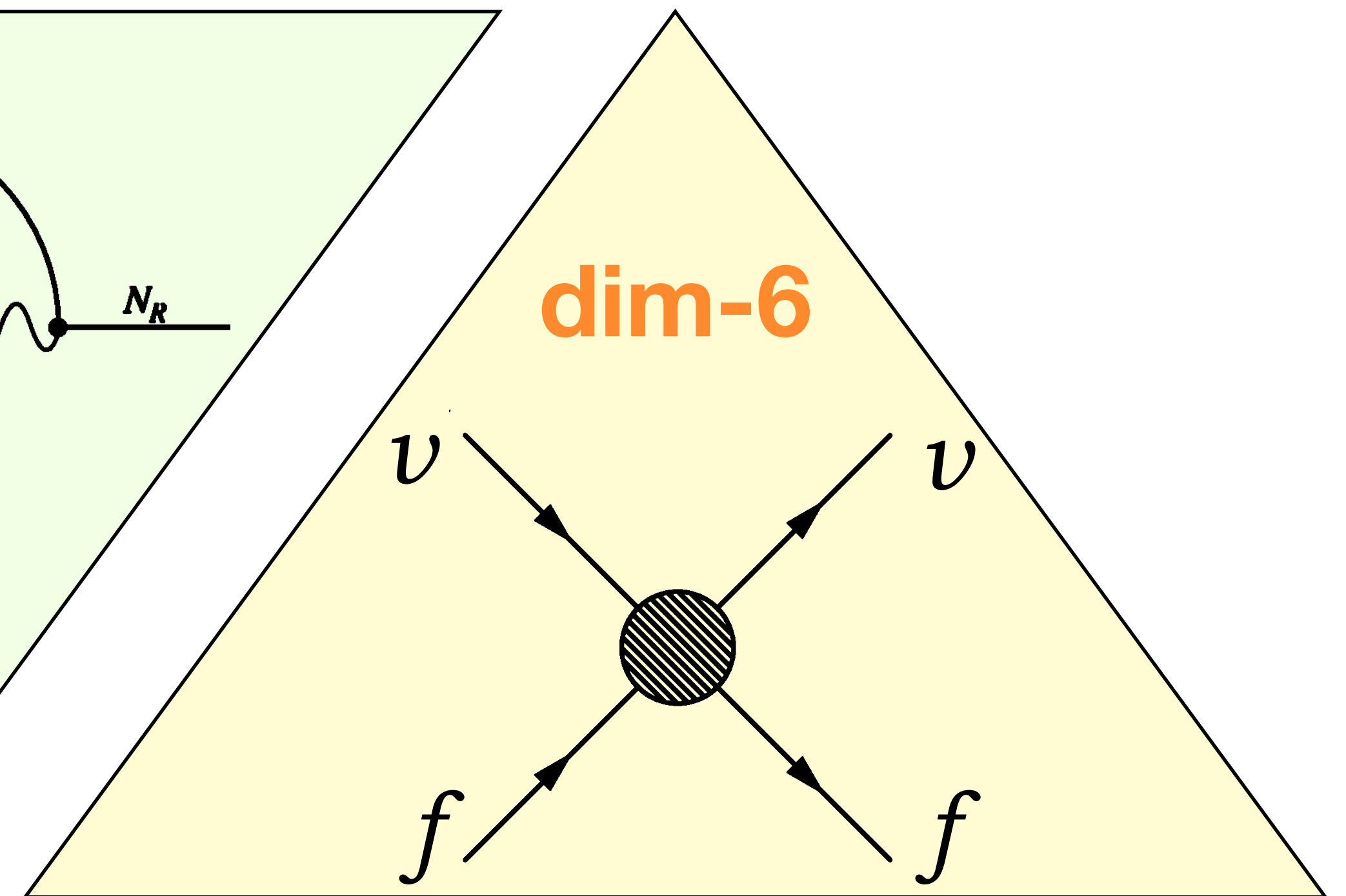
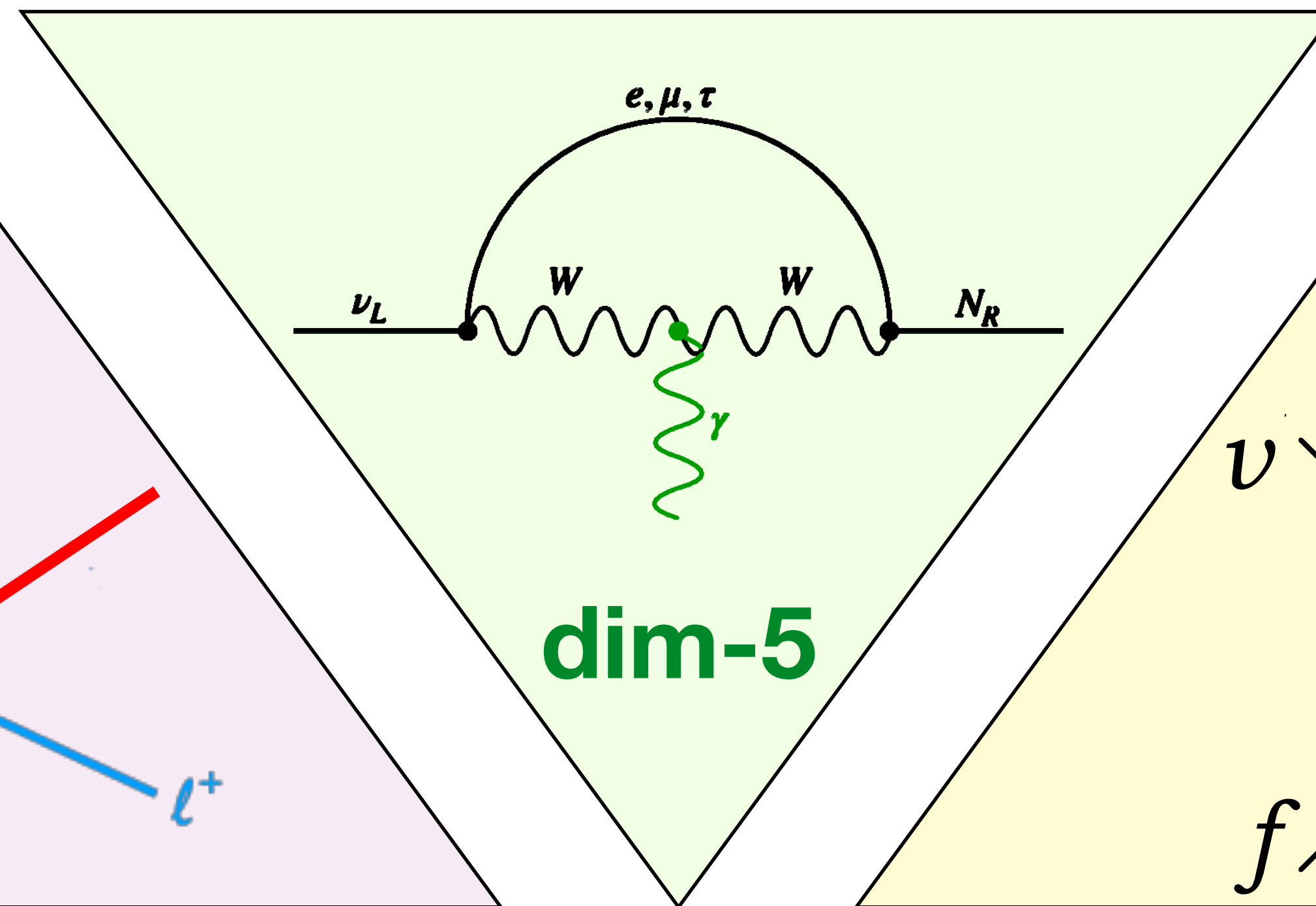


# Neutrino Physics Beyond the Standard Model

e.g. neutrino magnetic moments

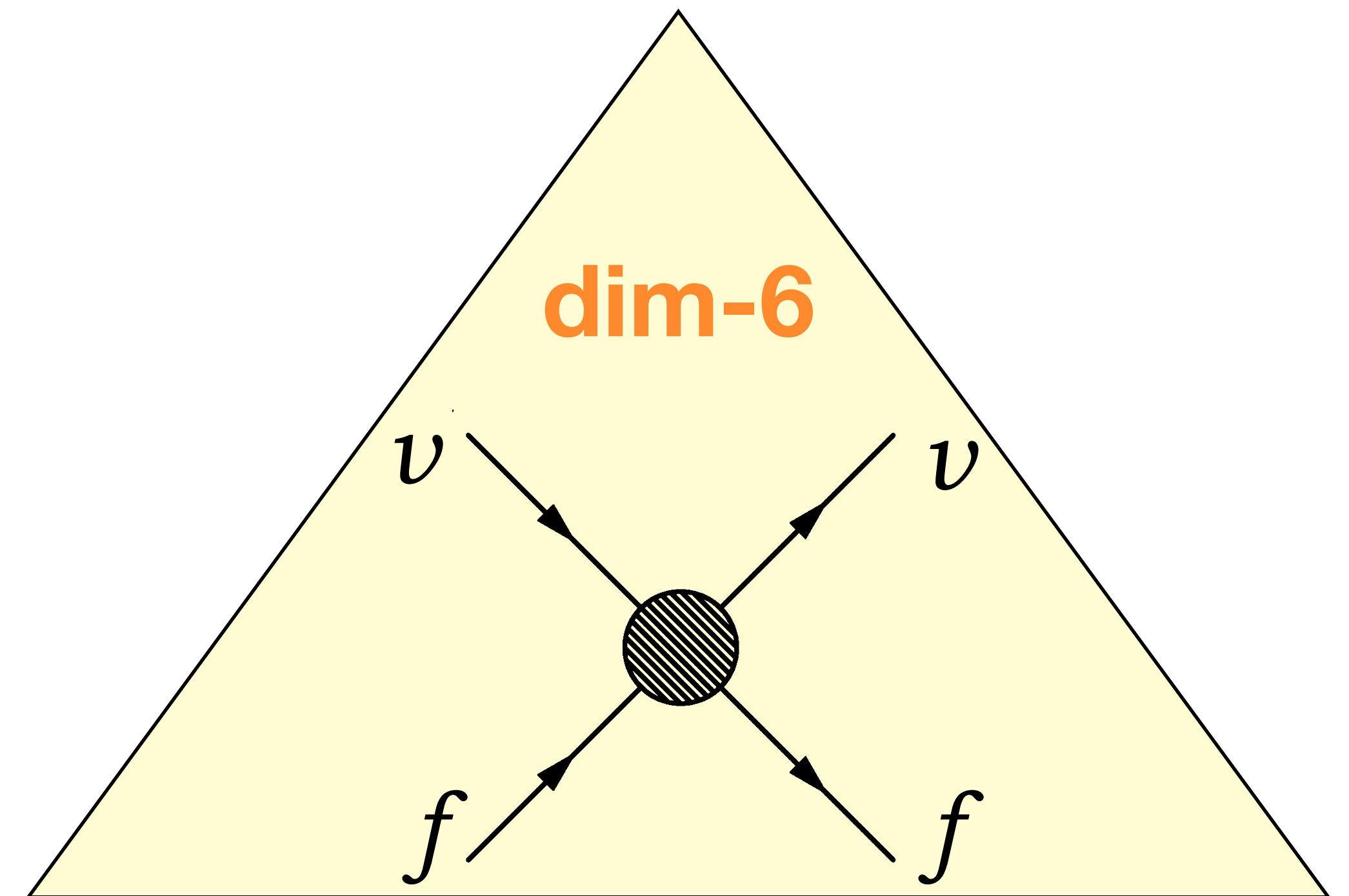


e.g. sterile neutrinos



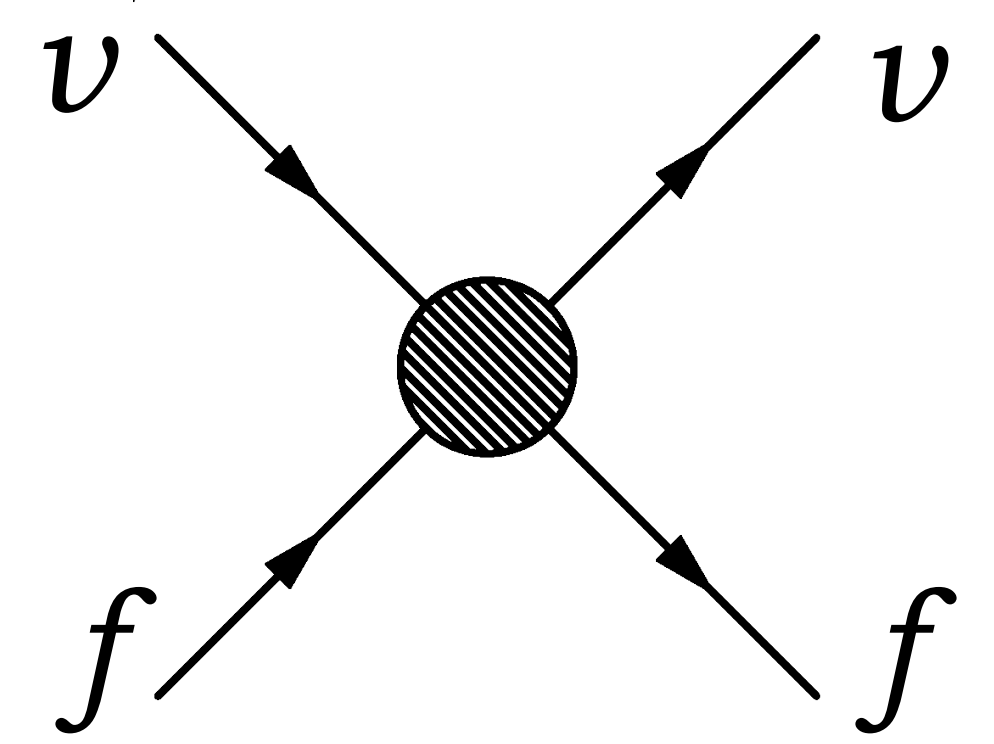
e.g. non-standard interactions

# Effective Field Theories for Neutrinos



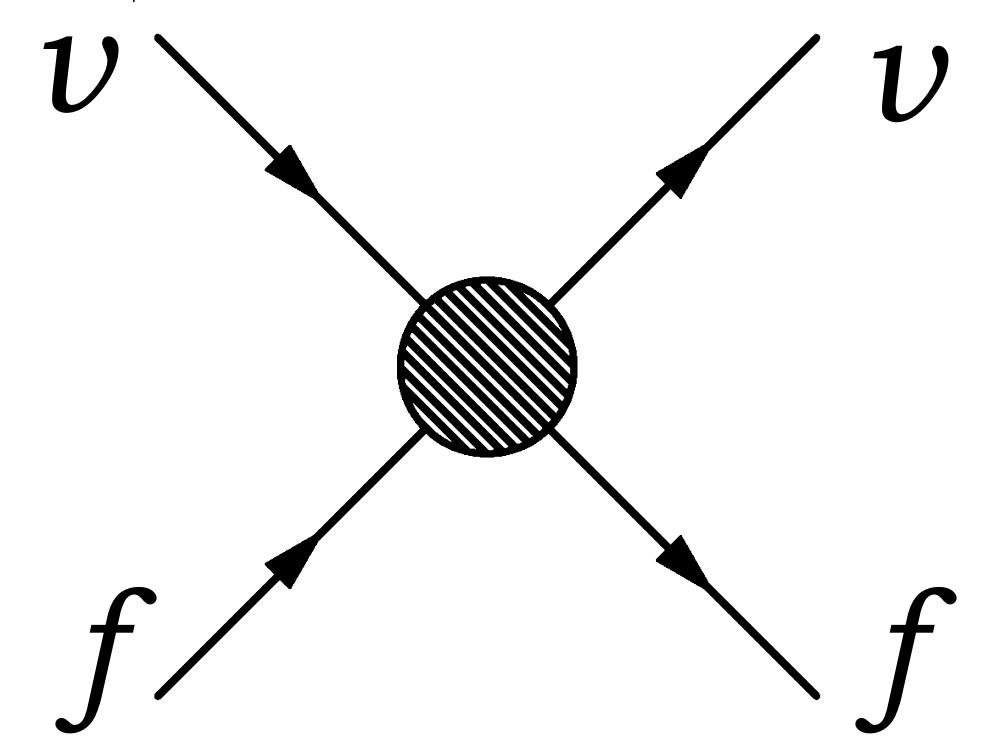


# Effective Field Theories for Neutrinos



$$\begin{aligned}
 \mathcal{L}_{\text{WEFT}} \supset & -\frac{2V_{ud}}{v^2} \left\{ [\mathbf{1} + \epsilon_L]_{\alpha\beta} (\bar{q}_u \gamma^\mu P_L q_d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) + [\epsilon_R]_{\alpha\beta} (\bar{q}_u \gamma^\mu P_R q_d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) \right. \\
 & + \frac{1}{2} [\epsilon_S]_{\alpha\beta} (\bar{q}_u q_d) (\bar{\ell}_\alpha P_L \nu_\beta) - \frac{1}{2} [\epsilon_P]_{\alpha\beta} (\bar{q}_u \gamma_5 q_d) (\bar{\ell}_\alpha P_L \nu_\beta) \\
 & \left. + \frac{1}{4} [\epsilon_T]_{\alpha\beta} (\bar{q}_u \sigma^{\mu\nu} P_L q_d) (\bar{\ell}_\alpha \sigma_{\mu\nu} P_L \nu_\beta) + \text{h.c.} \right\}.
 \end{aligned}$$

# Effective Field Theories for Neutrinos

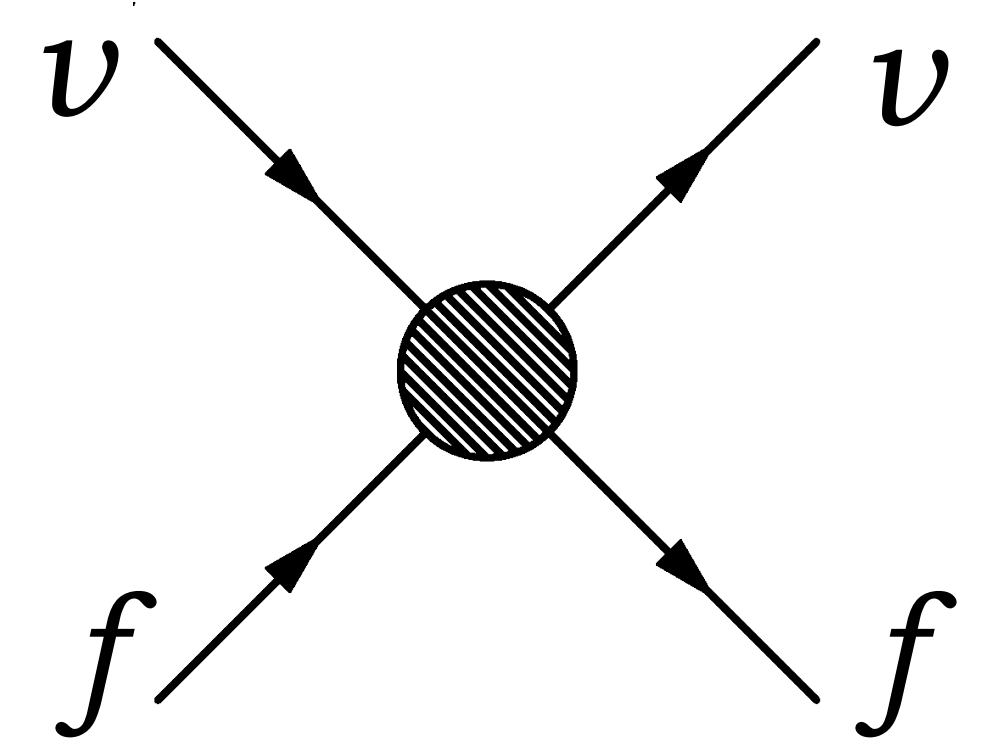


**dimensionless coefficients**  
 (interaction strength  
 relative to SM weak interactions)

$$\mathcal{L}_{\text{WEFT}} \supset -\frac{2V_{ud}}{v^2} \left\{ [\mathbf{1} + \epsilon_L]_{\alpha\beta} (\bar{q}_u \gamma^\mu P_L q_d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) + [\epsilon_R]_{\alpha\beta} (\bar{q}_u \gamma^\mu P_R q_d) (\bar{\ell}_\alpha \gamma_\mu P_L \nu_\beta) \right. \\
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 \left. + \frac{1}{4} [\epsilon_T]_{\alpha\beta} (\bar{q}_u \sigma^{\mu\nu} P_L q_d) (\bar{\ell}_\alpha \sigma_{\mu\nu} P_L \nu_\beta) + \text{h.c.} \right\}.$$



# Effective Field Theories for Neutrinos



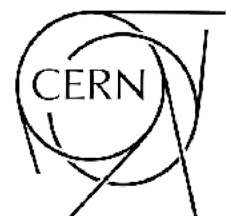
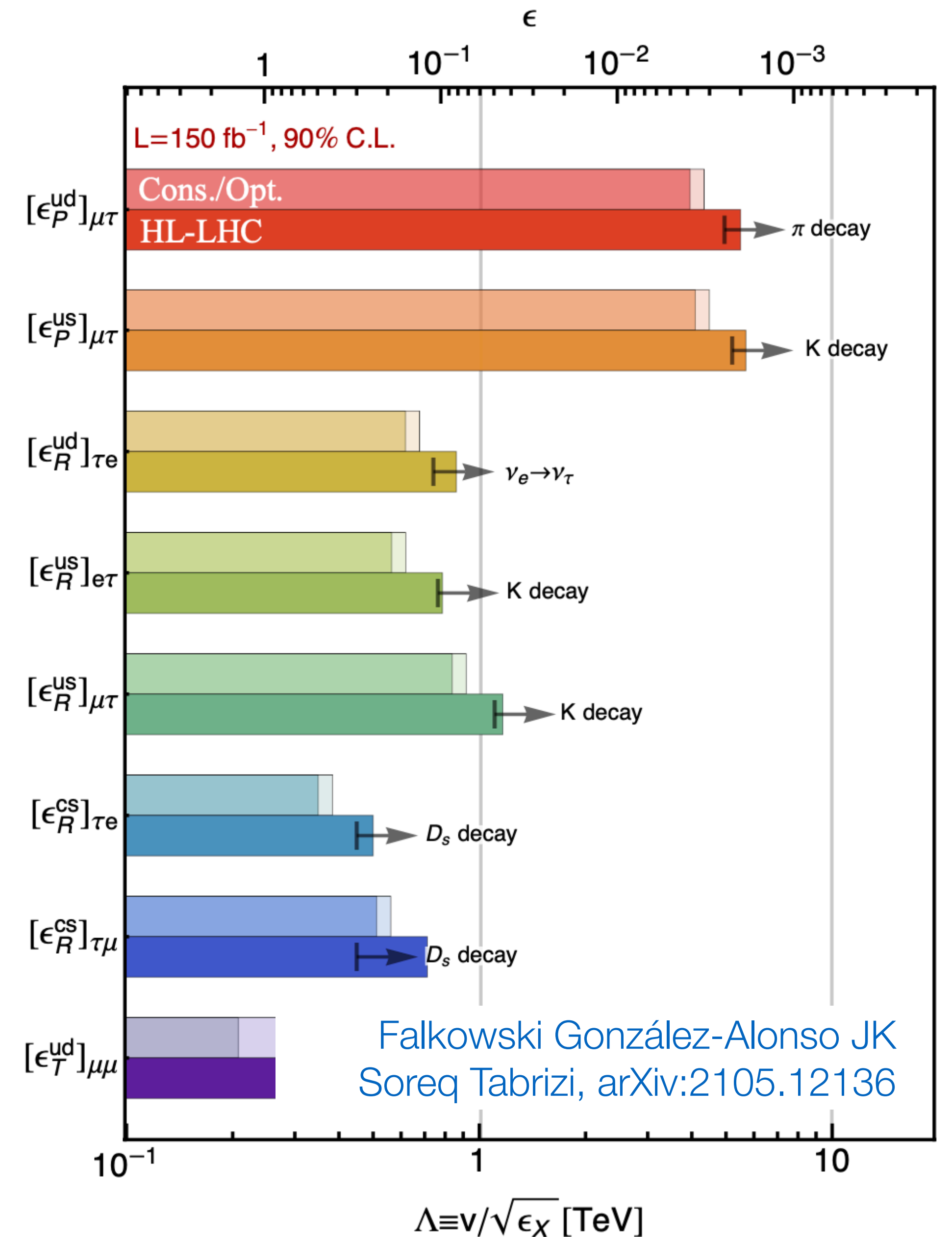
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**dim-6 operators**  
with different Lorentz structures

# New Interactions in DIS

- straightforward –  $\nu$  scattering described in terms of PDFs
- though flux uncertainties remain a problem



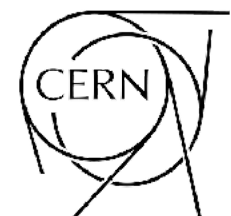
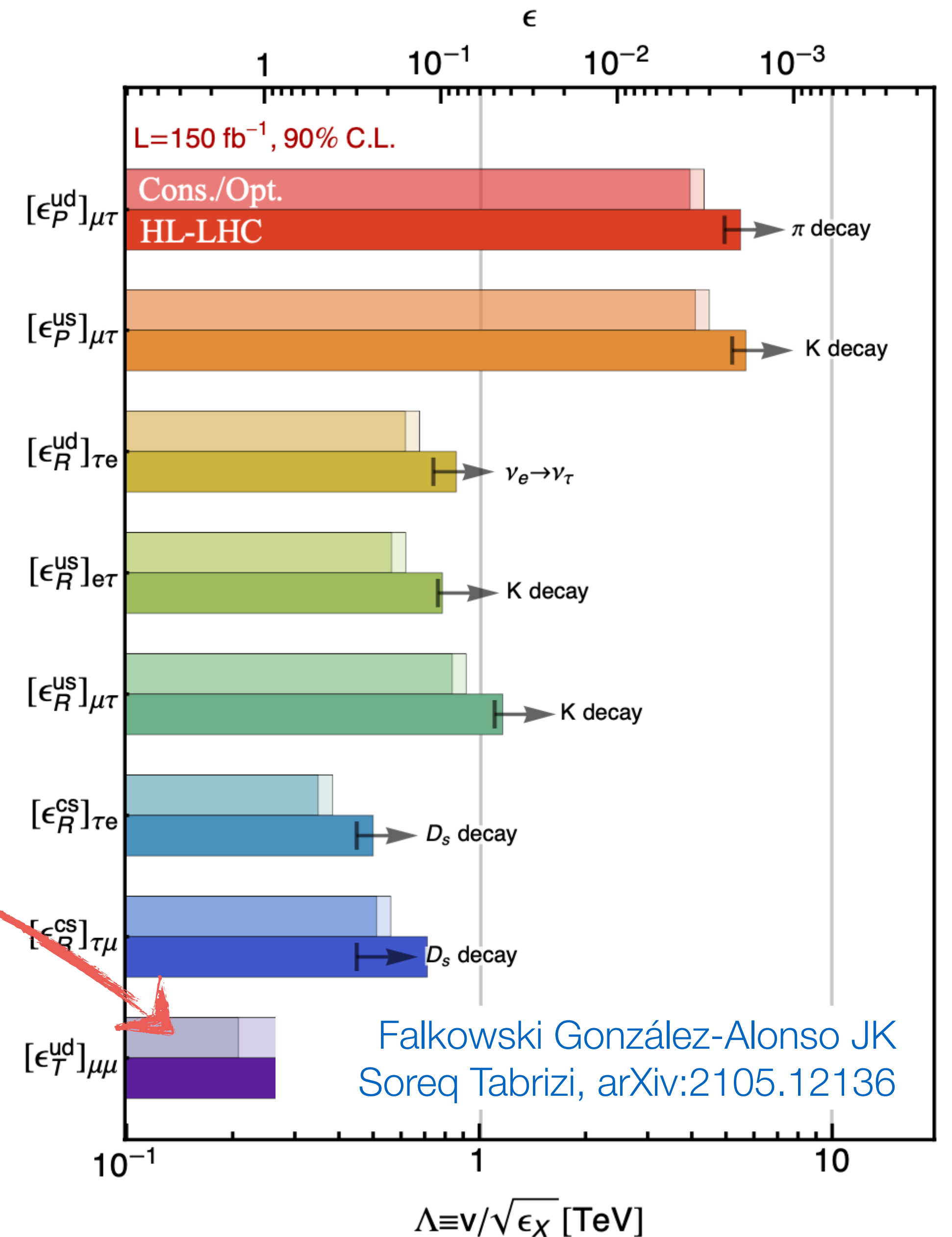


# New Interactions in DIS

- straightforward –  $\nu$  scattering described in terms of PDFs
- though flux uncertainties remain a problem

## conservative scenario

flux uncertainties 30% / 40% / 50%  
for  $\nu_e / \nu_\mu / \nu_\tau$



# New Interactions in DIS

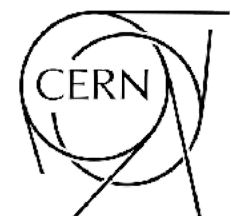
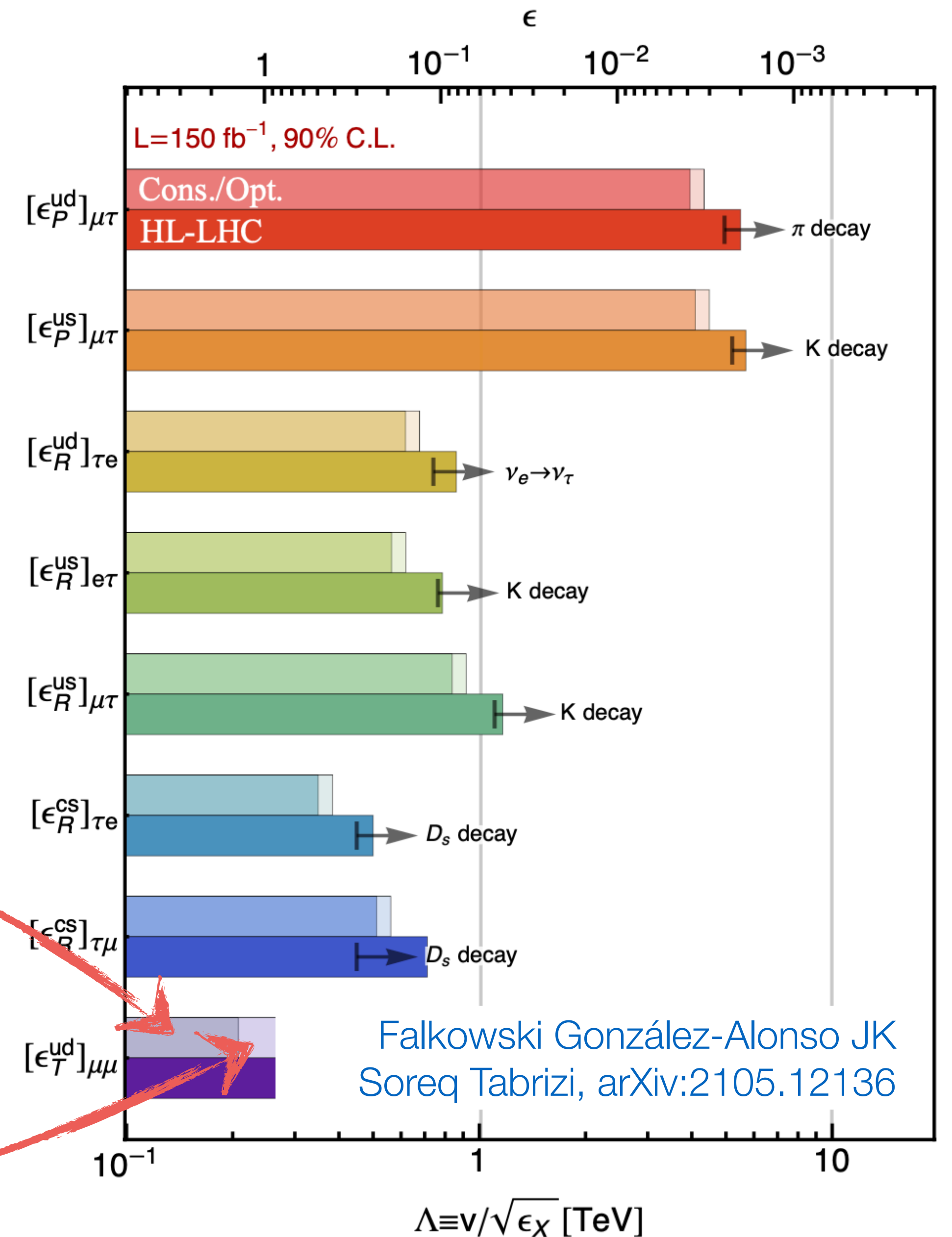
- straightforward –  $\nu$  scattering described in terms of PDFs
- though flux uncertainties remain a problem

## conservative scenario

flux uncertainties 30% / 40% / 50%  
for  $\nu_e / \nu_\mu / \nu_\tau$

## optimistic scenario

flux uncertainties 5% / 10% / 15%  
for  $\nu_e / \nu_\mu / \nu_\tau$





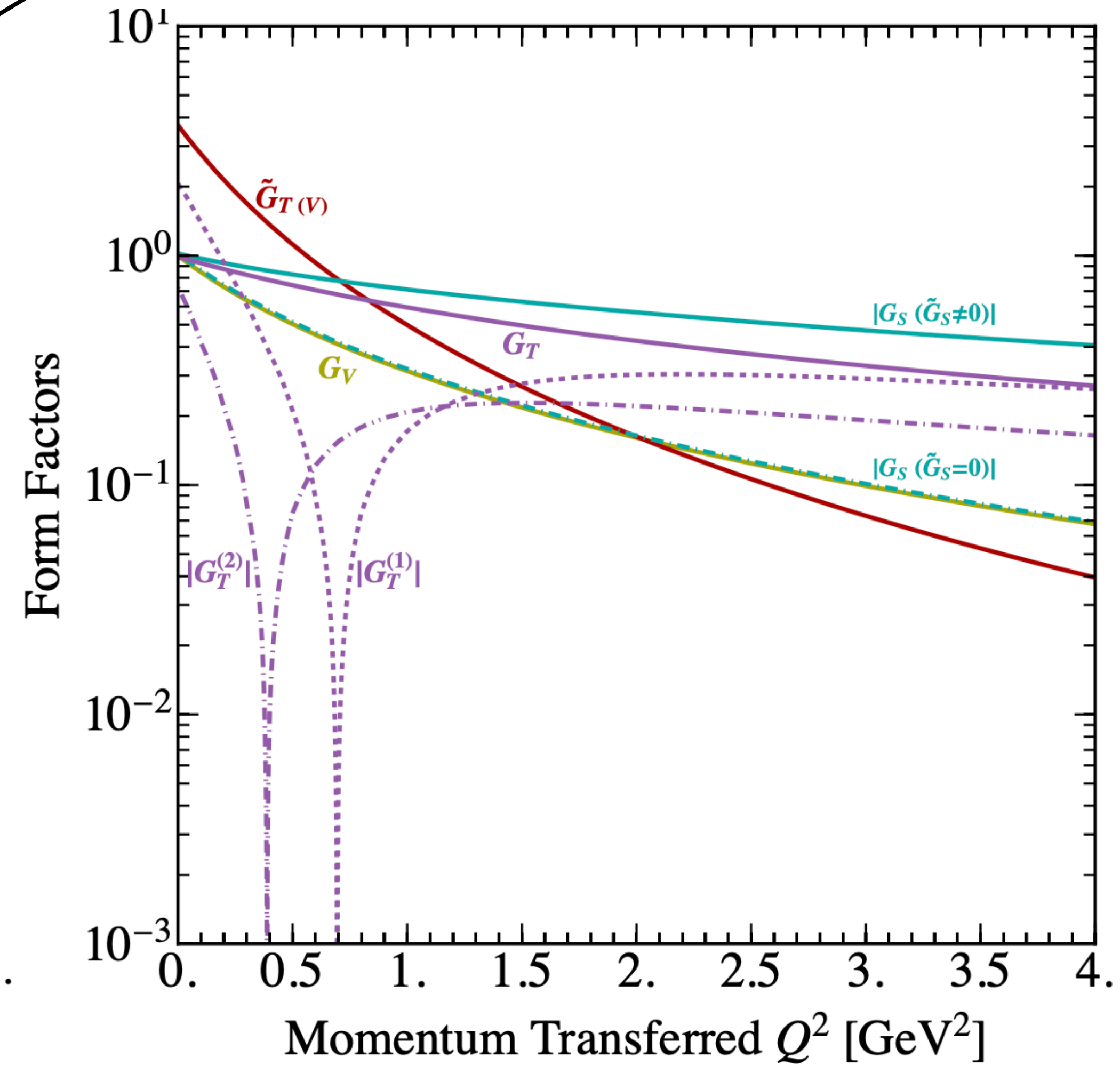
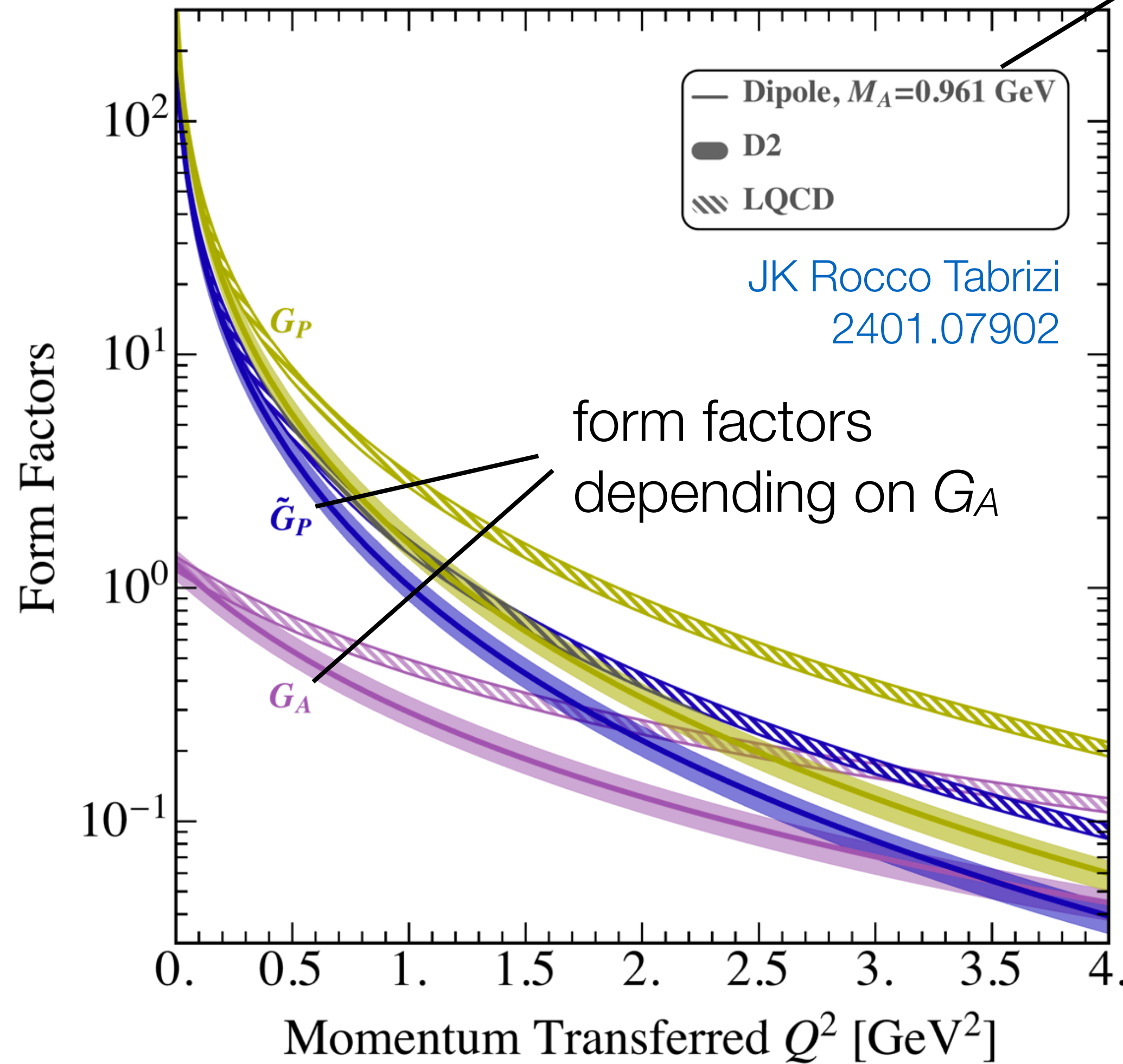
# New Interactions in QES

- For V–A: same uncertainties as in the SM
- For other Lorentz structures: significant nucleon form factor uncertainties

$$\begin{aligned}
 \langle p(p_p) | \bar{q}_u \gamma_\mu q_d | n(p_n) \rangle &= \bar{u}_p(p_p) \left[ G_V(Q^2) \gamma_\mu + i \frac{\tilde{G}_{T(V)}(Q^2)}{2M_N} \sigma_{\mu\nu} q^\nu - \frac{\tilde{G}_S(Q^2)}{2M_N} q_\mu \right] u_n(p_n), \\
 \langle p(p_p) | \bar{q}_u \gamma_\mu \gamma_5 q_d | n(p_n) \rangle &= \bar{u}_p(p_p) \left[ G_A(Q^2) \gamma_\mu \gamma_5 + i \frac{\tilde{G}_{T(A)}(Q^2)}{2M_N} \sigma_{\mu\nu} q^\nu \gamma_5 - \frac{\tilde{G}_P(Q^2)}{2M_N} q_\mu \gamma_5 \right] u_n(p_n), \\
 \langle p(p_p) | \bar{q}_u q_d | n(p_n) \rangle &= G_S(Q^2) \bar{u}_p(p_p) u_n(p_n), \\
 \langle p(p_p) | \bar{q}_u \gamma_5 q_d | n(p_n) \rangle &= G_P(Q^2) \bar{u}_p(p_p) \gamma_5 u_n(p_n), \\
 \langle p(p_p) | \bar{q}_u \sigma_{\mu\nu} q_d | n(p_n) \rangle &= \bar{u}_p(p_p) \left[ G_T(Q^2) \sigma_{\mu\nu} - \frac{i}{M_N} G_T^{(1)}(Q^2) (q_\mu \gamma_\nu - q_\nu \gamma_\mu) \right. \\
 &\quad \left. - \frac{i}{M_N^2} G_T^{(2)}(Q^2) (q_\mu P_\nu - q_\nu P_\mu) - \frac{i}{M_N} G_T^{(3)}(Q^2) (\gamma_\mu \not{q} \gamma_\nu - \gamma_\nu \not{q} \gamma_\mu) \right] u_n(p_n),
 \end{aligned}$$

# New Interactions in QES

three different estimates for the axial form factor





# New Interactions in QES

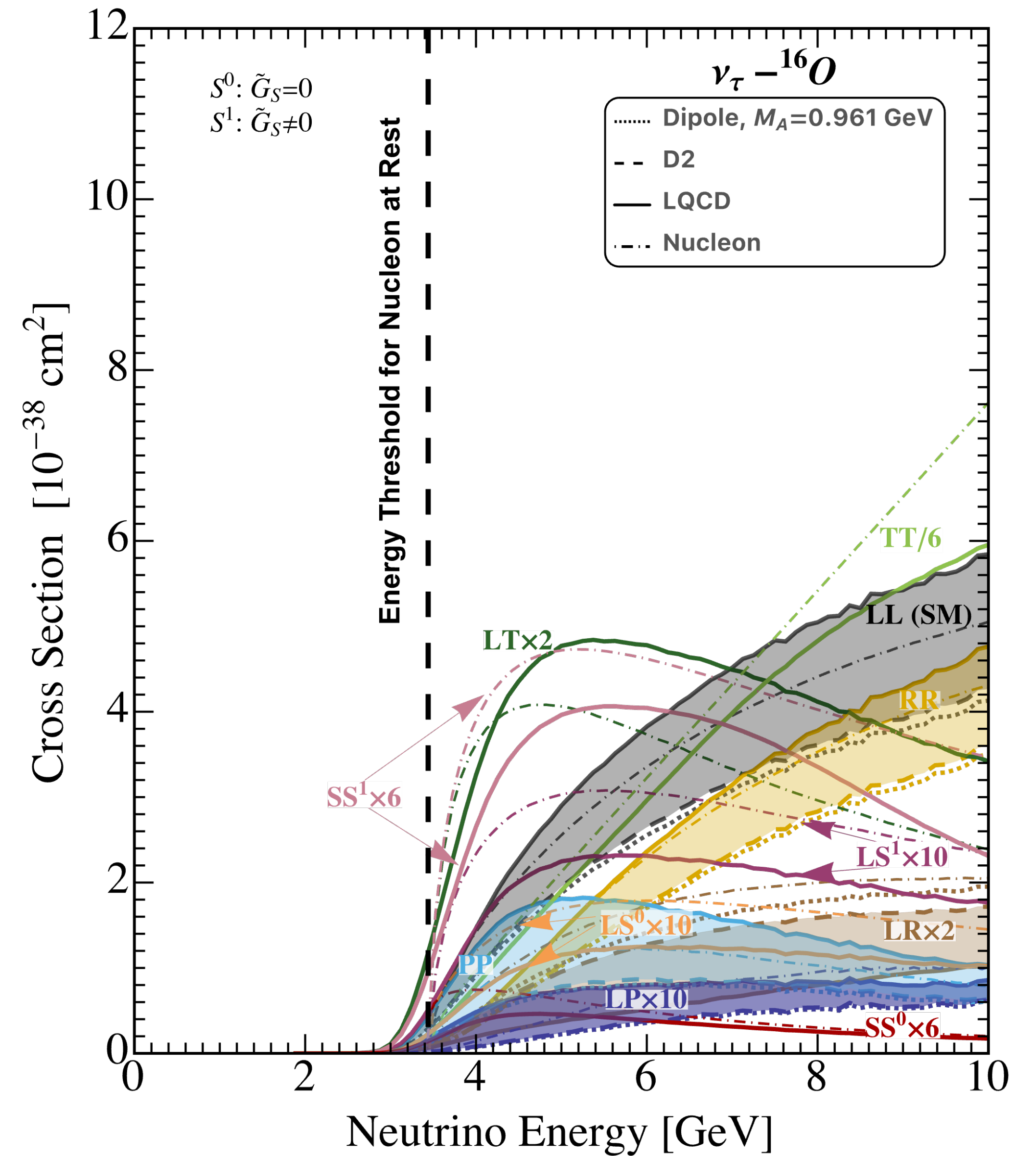
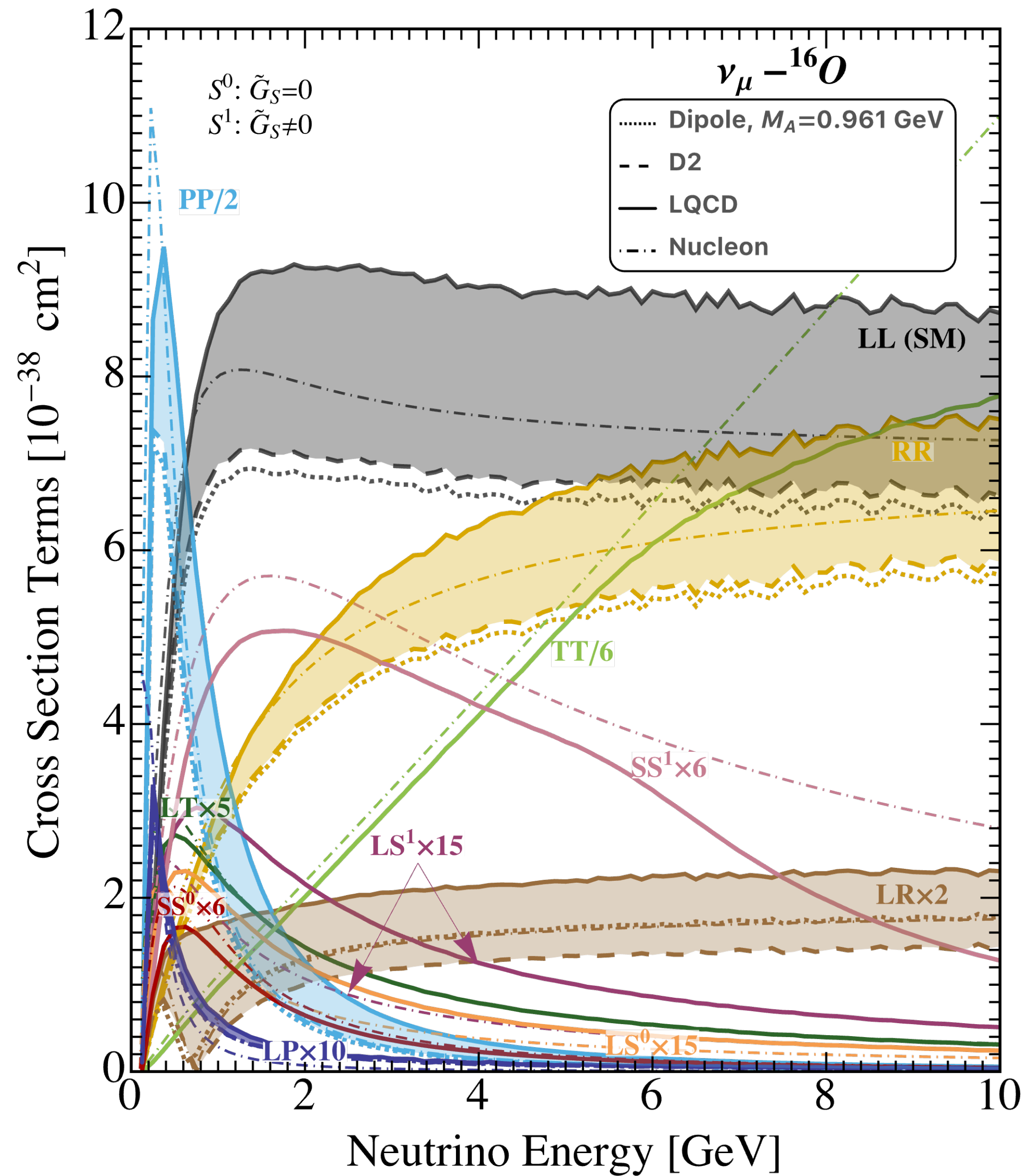
- in addition: **nuclear effects**, parameterized via a spectral function

$$\frac{1}{2} \sum_{\text{spin}} \mathcal{A}_{X,\alpha} \mathcal{A}_{Y,\alpha}^* = \frac{1}{2} \int \frac{d^3 p_N}{(2\pi)^3} P_h(\mathbf{p}_N, E^*) \frac{m_N}{e(\mathbf{p}_N)} \frac{m_N}{e(\mathbf{q} + \mathbf{p}_N)} \\ \times \sum_{\text{spin}} \sum_N A_{X,\alpha} A_{Y,\alpha}^* \delta(\tilde{\omega} + e(\mathbf{p}_N) - e(\mathbf{q} + \mathbf{p}_N)),$$

- cross-section decomposed into contributions with different Lorentz structures

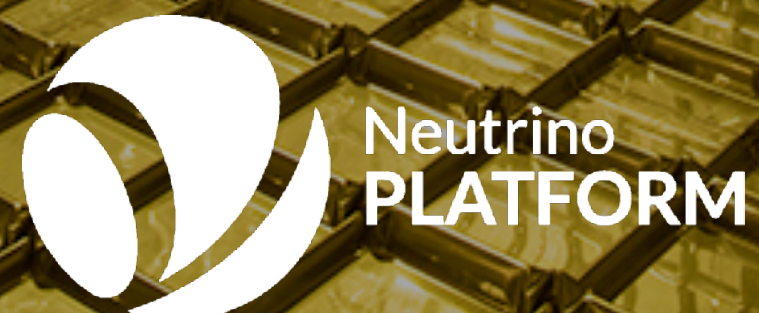
$$\frac{d\sigma_{\alpha\beta}}{dQ^2} = \frac{d\hat{\sigma}_{LL,\alpha}}{dQ^2} \delta_{\alpha\beta} + \sum_X \left( [\epsilon_X]_{\alpha\beta} \frac{d\hat{\sigma}_{LX,\alpha}}{dQ^2} \delta_{\alpha\beta} + h.c. \right) + \sum_{X,Y,\beta} [\epsilon_X]_{\alpha\beta} [\epsilon_Y]_{\alpha\beta}^* \frac{d\hat{\sigma}_{XY,\alpha}}{dQ^2}$$

# New Interactions in QES





**Thank You!**



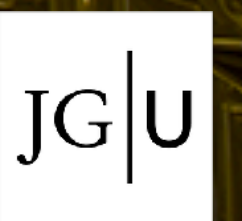
JOHANNES GUTENBERG  
UNIVERSITÄT MAINZ







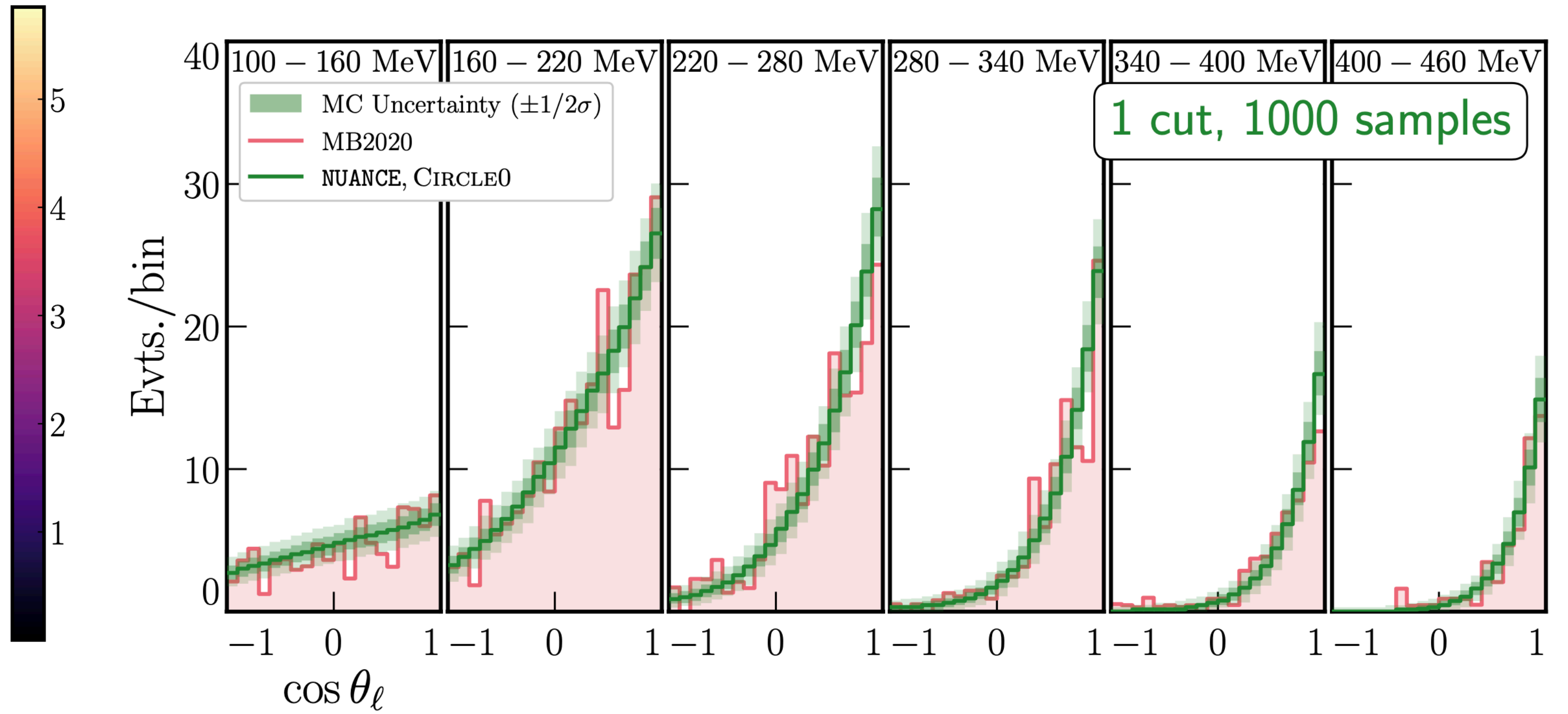
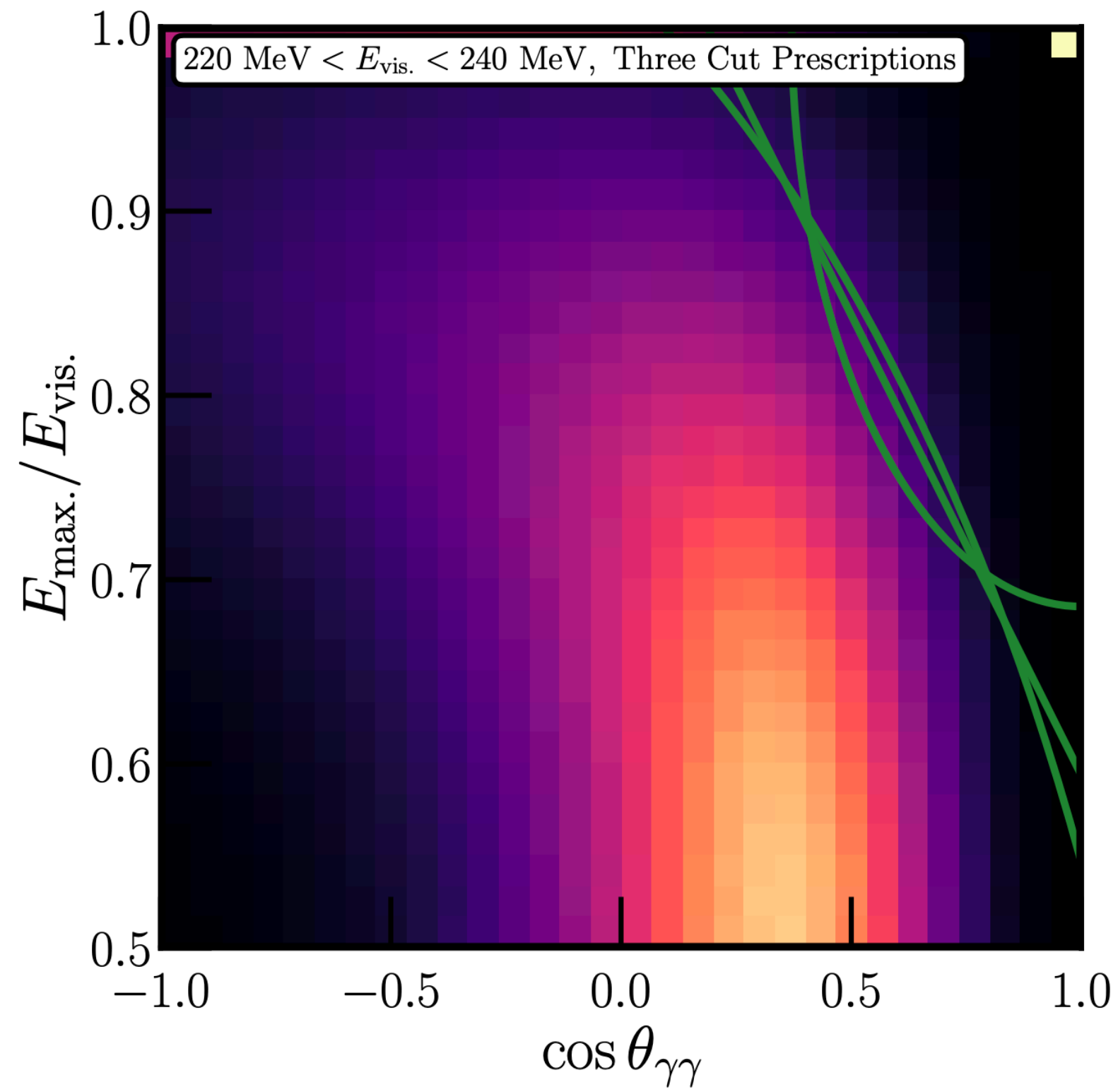
# Bonus Slides



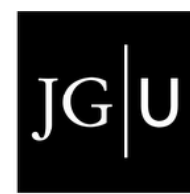
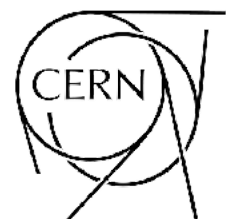
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UNIVERSITÄT MAINZ



# Modeling $\pi^0$ Mis-ID in MiniBooNE

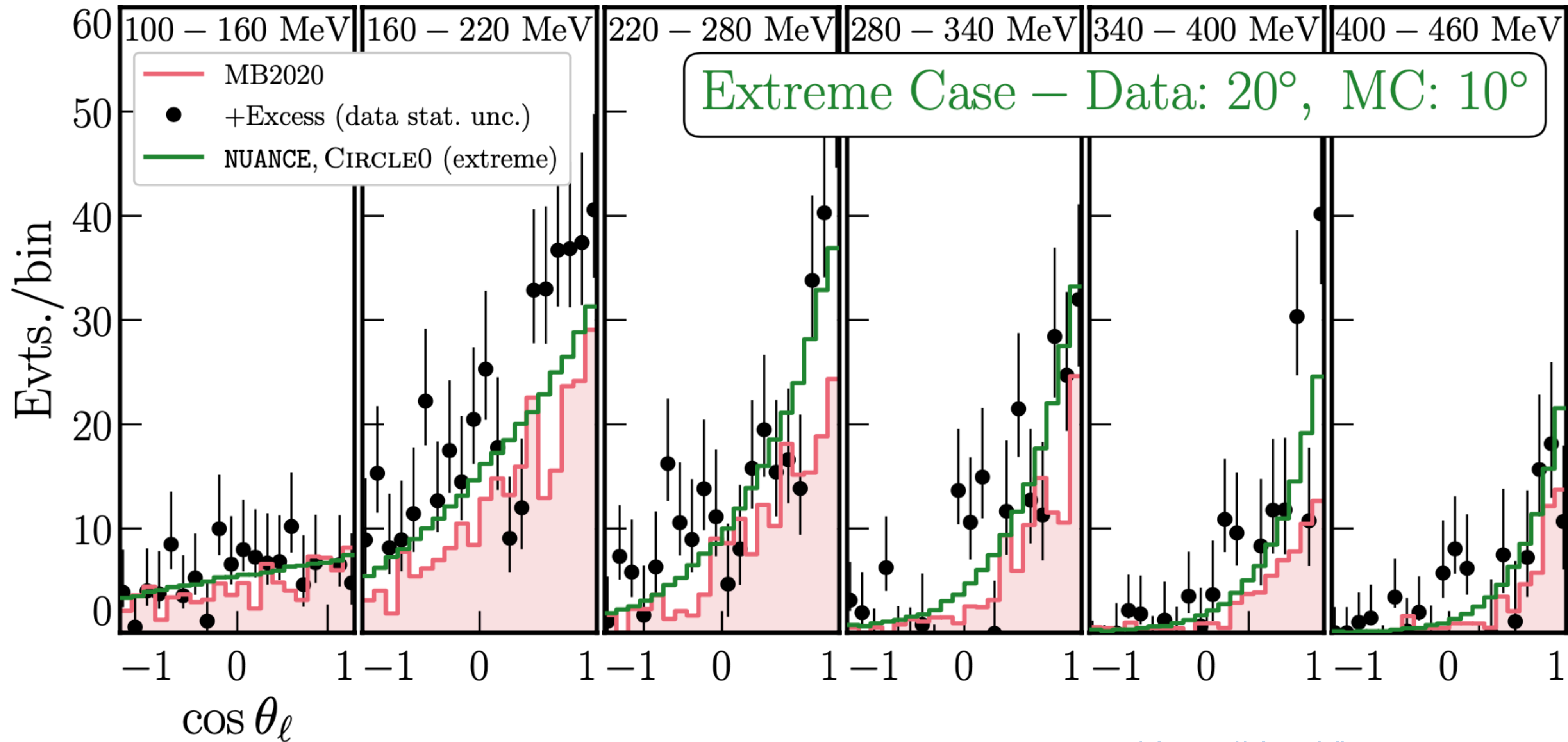


Kelly JK, arXiv:2210.08021





# Can mis-modeling of $\pi^0$ mis-ID explain the anomaly?



Kelly JK, arXiv:2210.08021

