Assessing MAGIS capabilities to measure the properties of GW signals from compact binaries

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What is special about AI GW detectors?

- Single baseline
 - ⇒ different antenna functions
- Two timescales (atom free-fall time, laser travel time)
 ⇒ different shape of the sensitivity curve
- Terrestrial or mid-Earth orbit

 \Rightarrow detector re-orientation on hour-timescales

- Small ratio of GW frequency⁻¹/GW wavelength to effective "baseline"
 - \Rightarrow excellent sky localization of sources

Quick reminder of sky-localization capabilities

Three options:

- High signal-to-noise ratio measurement of the waveform in many spatial directions (very hard)
- Time-of-flight between multiple detectors
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Ultimately, these both lead to the diffraction limit:

$$\sigma_{\rm V\Omega} \propto \lambda_{\rm GW} \, / \, L_{\rm baseline}$$

$$\sigma_{\Omega}^{\rm LIGO} \sim 10^{\circ} \left(\frac{1\,\rm kHz}{f_{\rm GW}}\right) ; \quad \sigma_{\Omega}^{\rm AI} \sim 0.3^{\circ} \left(\frac{1\,\rm Hz}{f_{\rm GW}}\right) ; \quad \sigma_{\Omega}^{\rm LISA} \sim 30^{\circ} \left(\frac{10\,\rm mHz}{f_{\rm GW}}\right)$$

Leading-order inspiral waveforms (pol. basis)

• Frequency evolution ("chirp") $\frac{df_{\rm GW}}{dt} = \frac{96}{5}\pi^{8/3}\mathcal{M}_c f_{\rm GW}^{11/3}$

with the (detector frame) "chirp mass" $\mathcal{M}_c = \frac{(m_1 m_2)^{3/3}}{(m_1 + m_2)^{1/5}}$

• Polarization basis waveforms: $\underline{\underline{h}}(t) = h_+(t)\underline{\underline{e}}_+ + h_{\times}(t)\underline{\underline{e}}_{\times}$

$$h_{+/\times}(t) = \frac{2\mathcal{M}_c^{5/3} \left[\pi f_{\rm GW}(t)\right]^{2/3}}{d_L} \begin{cases} \left[1 + \left(\hat{\boldsymbol{L}} \cdot \hat{\boldsymbol{n}}\right)^2\right] \cos \Phi_{\rm GW}(t) \\ 2\left(\hat{\boldsymbol{L}} \cdot \hat{\boldsymbol{n}}\right) \sin \Phi_{\rm GW}'(t) \end{cases}$$
$$\cos \iota \equiv \hat{\boldsymbol{L}} \cdot \hat{\boldsymbol{n}} \end{cases}$$

Leading-order inspiral waveforms (antenna func.)

• The response of a single-baseline detector is then

$$h(t) = \hat{l}(t)\underline{h}(t)\hat{l}(t) = F_{+}(t)h_{+}(t) + F_{\times}(t)h_{\times}(t)$$
$$F_{+/\times}(t) \equiv \hat{l}(t) \cdot \underline{e}_{+/\times} \cdot \hat{l}(t)$$

Parameters describing the binary

Parameter		Benchmark value
Chirp Mass	$\mathcal{M}_c = (m_1 m_2)^{3/5} / (m_1 + m_2)^{1/5}$	
Mass ratio	$q=m_1/m_2$	1.15
Luminosity distance	d_L	
Binary phase at $f_{\rm GW}^{\rm ref}$	Φ_0	0
Time of merger	t_c	solar equinox
Inclination angle	ι	45°
Polarization angle	ψ	60°
Right ascension	lpha	60°
Declination	δ	$6.6^{\circ} (30^{\circ} \text{ above ecliptic})$

Overview of the code

Fisher-matrix approach to forecast parameter reconstruction

- Computationally "cheap"
- Only reliable in high signal-to-noise ratio limit

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one-sided detector noise PSD

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- Define the inner product: $\langle h, g \rangle = 4 \operatorname{Re} \left[\int_{0}^{\infty} df \, \frac{\widetilde{h}^{*}(f) \widetilde{g}(f)}{S_{n}(f)} \right]$
- Signal-to-noise ratio $\rho^2 = \langle h, h \rangle$ and Fisher information matrix $\Gamma_{ij} = \left\langle \frac{\partial h}{\partial \theta_i}, \frac{\partial h}{\partial \theta_j} \right\rangle$

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Fisher-matrix approach to forecast parameter reconstruction

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- Calculate (3.5/3.0 PN) waveforms in time-domain (good for accelerated det., but computationally expensive)
- FFT to frequency domain (requires choosing window functions... current implementation uses Planck window)
- Some subtleties due to the accelerated detector frame (what do we mean by "detector frame" masses? What is the clock? ...)
- Include priors on periodic parameters (α , δ , ι , ψ)

Detector benchmarks

• (Network of) terrestrial MAGIS-1 km detectors



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- MAGIS-space in mid-Earth orbit
 - 20,000 km radius / 7.8 h period orbit
 - 3,600 km separation between the satellites
 - \circ Consider orbits inclined {0°, 45°, 90°} to ecliptic
 - For comparison, study also heliocentric orbits at {0.5, 1, 2} AU

Results: SNR in $\mathcal{M}_c - d_L$ plane



Heuristics of parameter reconstruction



Results: $\sigma(d_I)$ in $\mathcal{M}_c - d_I$ plane



Results: $\sigma(\sqrt{\Omega_n})$ in $\mathcal{M}_c - d_L$ plane









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Space detector. SNR in skymap



Space detector. SNR in skymap



Space detectors. SNR vs $\sigma(d_L)$ in skymap





Space detectors. SNR vs $\sigma(d_I)$ in skymap



Space detectors. SNR vs $\sigma(\sqrt{\Omega_n})$ in skymap



Space detector - the role of Earth



Space-detector. SNR in L-space



Space-detector. $\sigma(d_L)$ in *L*-space



Space-detector. $\sigma(\sqrt{\Omega_n})$ in *L*-space



Summary and comments

- Flexible code allows one to explore the capabilities of MAGIS to reconstruct the parameters of GW signals, i.e. to better understand the science case of MAGIS
- Evaluate possible sites for terrestrial MAGIS / different orbits for MAGIS-space
- The mid-band offers unique possibilities for the sky-localization of GW sources
 - Allows for an early warning system for electromagnetic follow-up!
- The mid-band is quite poor at measuring the luminosity distance
 - Combine with terrestrial laser-interferometers?

ICRS: Heliocentric equatorial coordinate system

