

beta-SRF Highlights

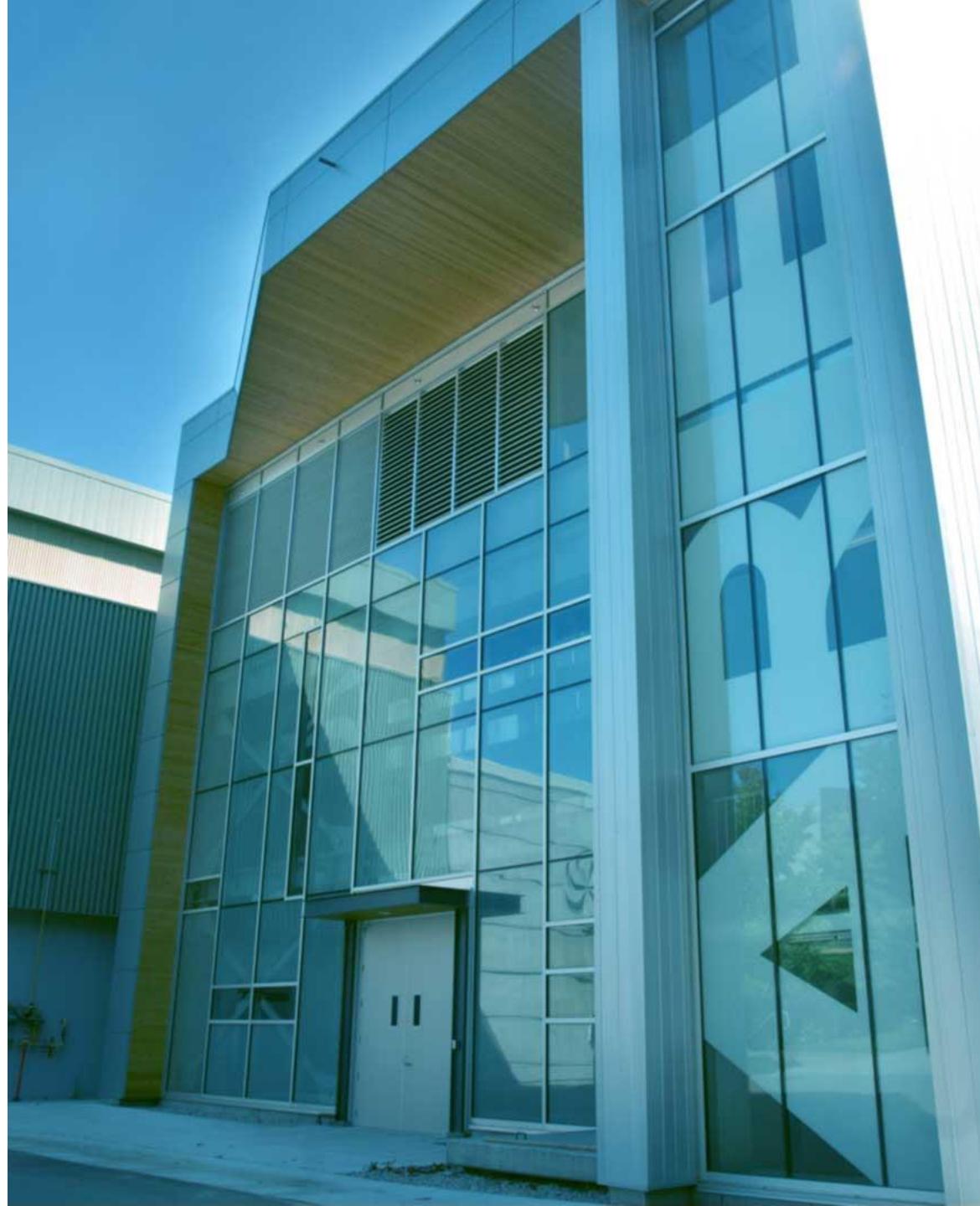
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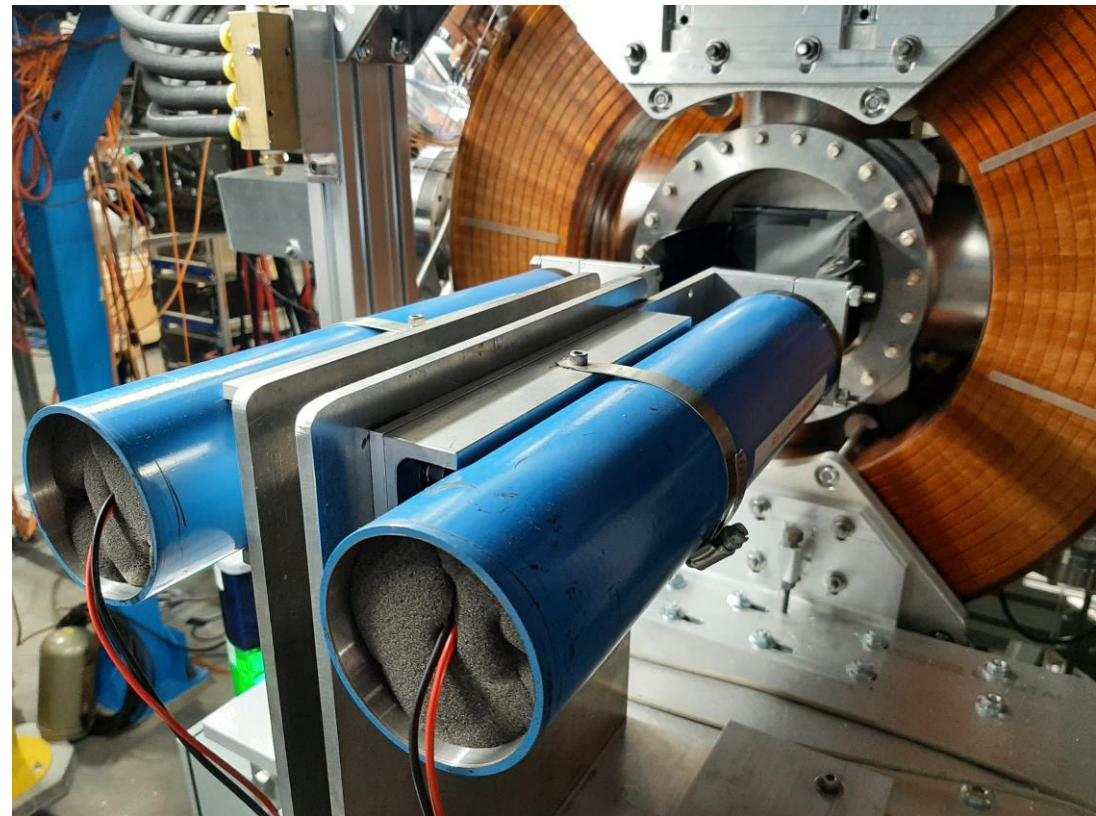
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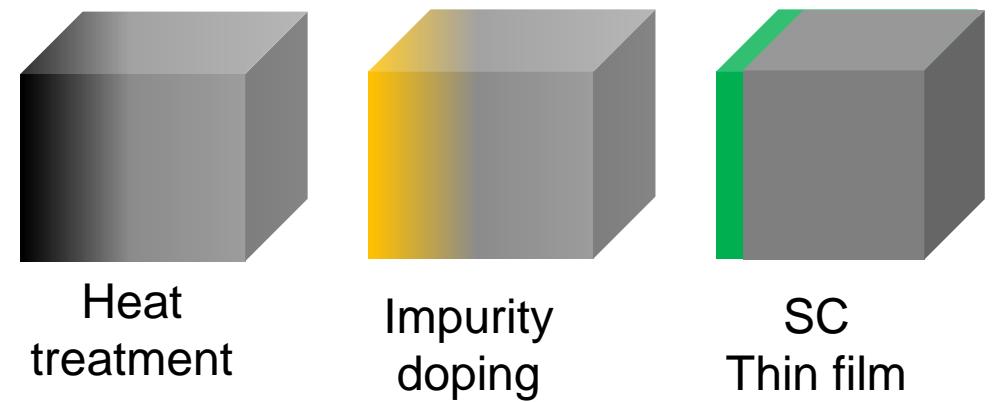
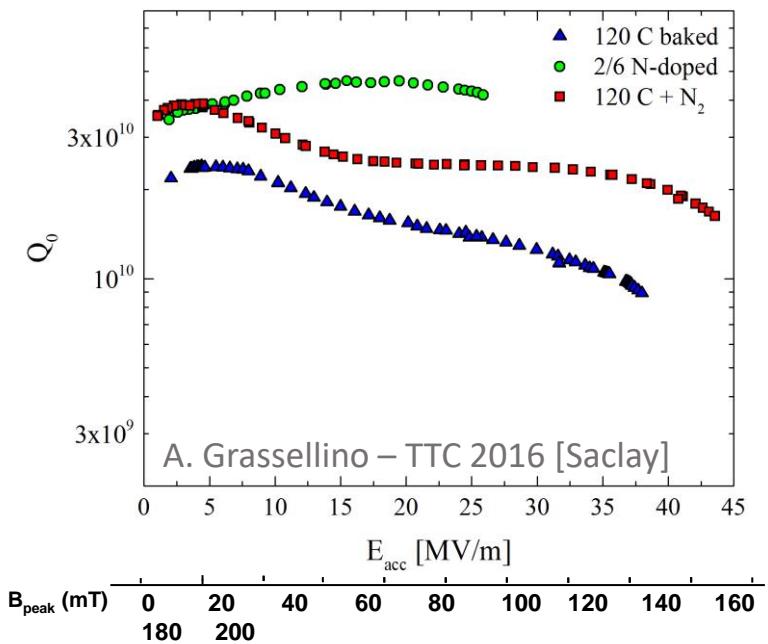
Outline

- Introduction/Motivation
- Beta-NMR Technique
- Results Highlights
- Summary & future outlook



SRF Cavities Performance

- Ultimate Goal:
 - **high Q** → **cheaper operation**
 - **high gradient** → **shorter LINAC**
- Underlying mechanisms ?
 - Macroscopic performance **very sensitive** to surface treatment
 - Nonlinear field dependence (**Q-slope**, **B_{quench}**)
- Empirical solutions:
 - Impurity engineering (baking, doping)
 - Thin film overlayer(s)



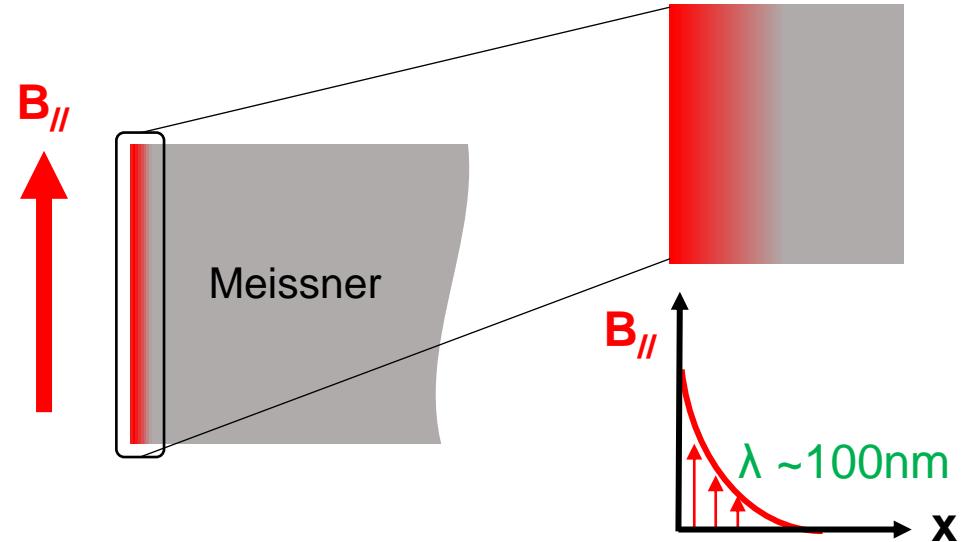
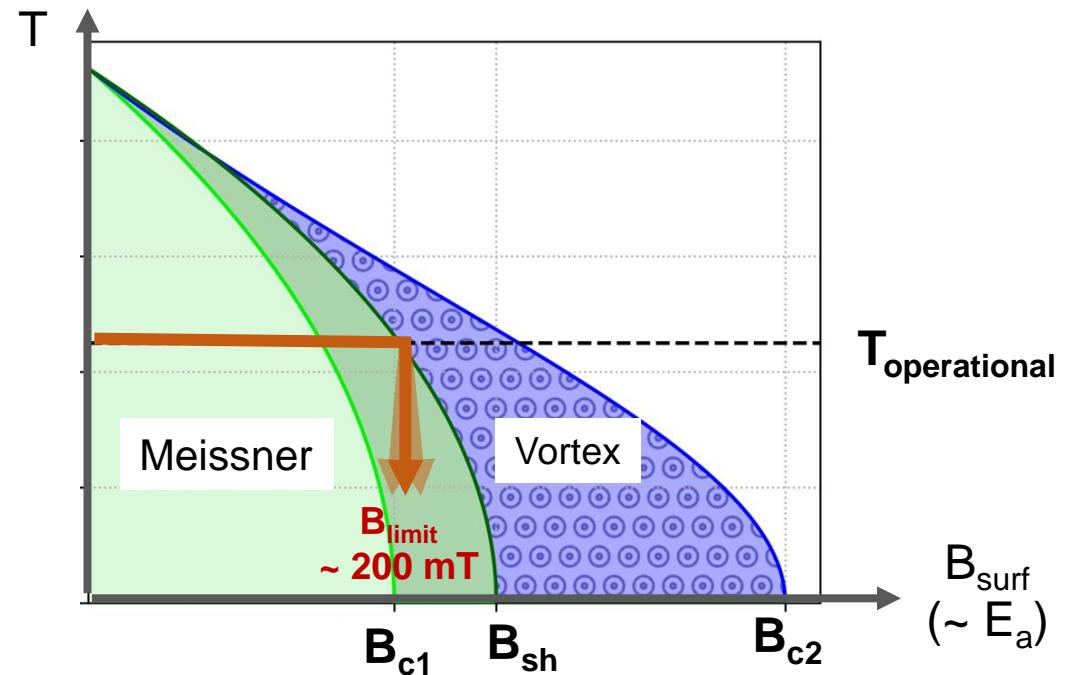
Nanometric Subsurface Role

- Achieving ideal performance

- **Accel. Gradient limit:**
 - Sustain Meissner phase to the limiting field (B_{limit})
- Screened magnetic fields within λ

- How do the subsurface variations & modifications affect

- Meissner limiting field: B_{limit} ?
- Screening response: λ vs. treatment ?
- Field dependence: λ vs. B_{app} ?



SRF Samples Characterization

- Requirements:

- local field within the London layer
- depth-resolved (within ~ 100 nm): $B(x)$
- up to $B_{\text{limit}} \sim 200 \text{ mT}$: $B(x)$ vs B_{applied}

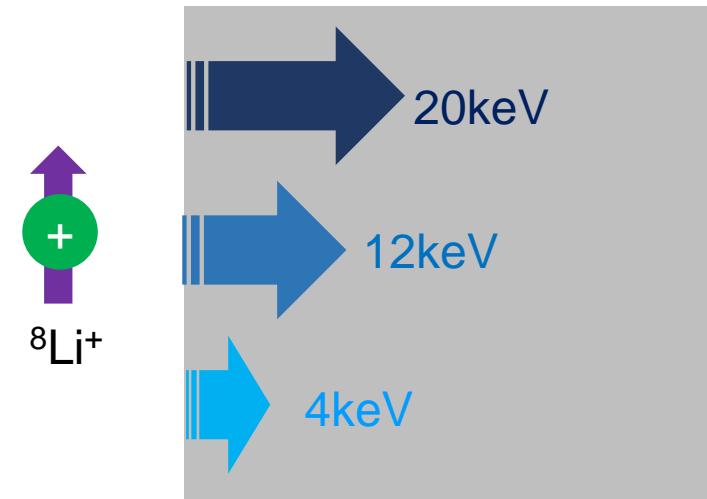
- Radioactive spin-polarized ions

- ion implantation $E \rightarrow \text{nm-scale, depth-resolved}$
- asymmetric radioactive decay → direct monitor of spin-polarization
- local field → evolution of spin-polarization

- Two facilities at TRIUMF

- HE- μ SR: bulk probe (100 μm)
- β -NMR: nm depth resolved (0 – 100 nm)

upgrade

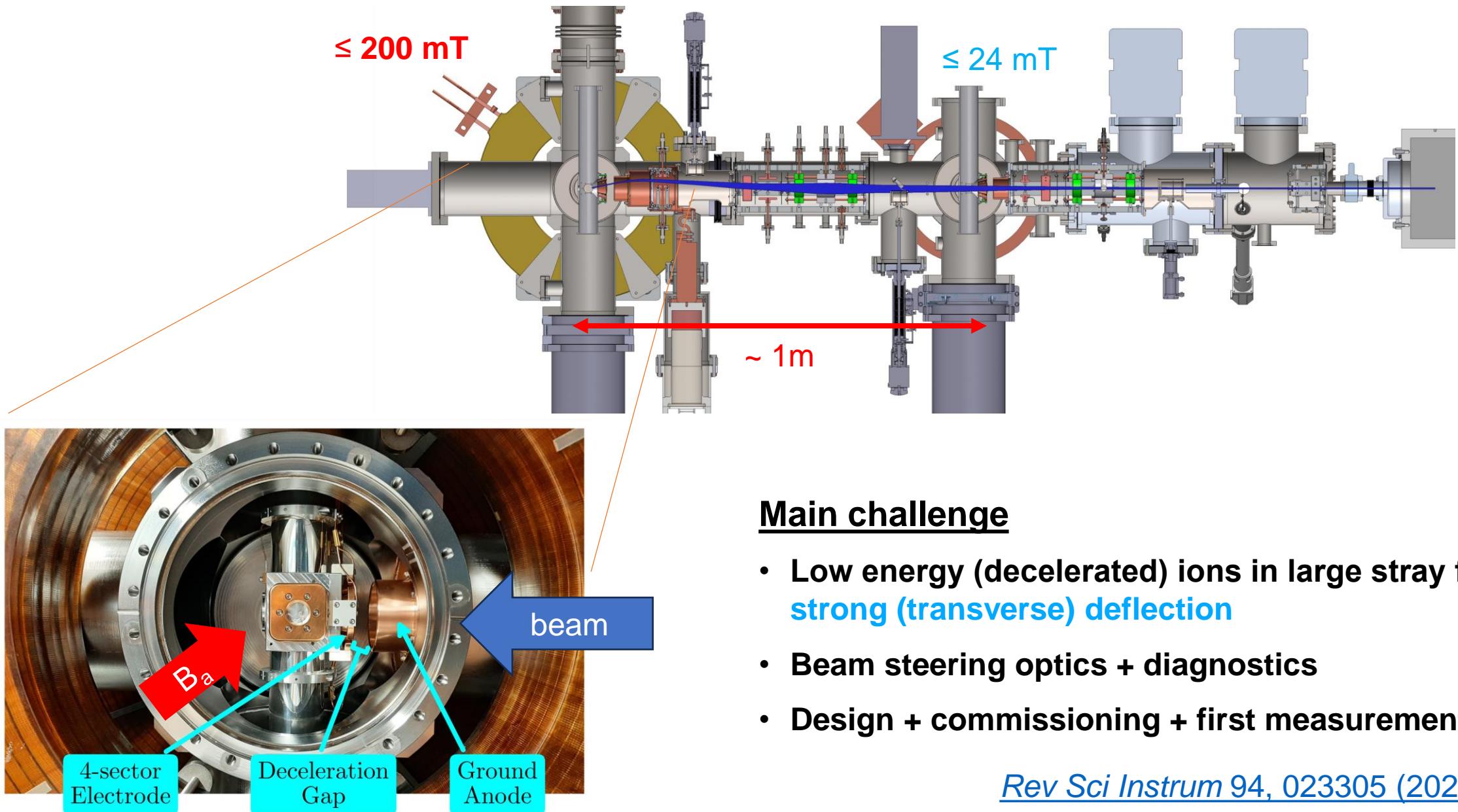


β -SRF beamline:

- high-parallel fields (200 mT)
- nm-scale depth resolved
- local field measurements

β SRF beamline

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Main challenge

- Low energy (decelerated) ions in large stray fields → strong (transverse) deflection
- Beam steering optics + diagnostics
- Design + commissioning + first measurements:

[Rev Sci Instrum 94, 023305 \(2023\)](#)

The beta-SRF Experiments

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SAMPLES

Two Nb samples measured: RRR
Niobium

- “**Baseline**”:
 - 1400 °C annealing for 4 hours + BCP
- Custom treated with mid-T bake (“**Oxygen doped**”):
 - “Baseline” + **400°C for 3 hours**
- **Field screening with applied fields 100 → 200 mT**



Sample ladder

LOCAL FIELD EXTRACTION

A) Relaxation Rate

B) Average Field

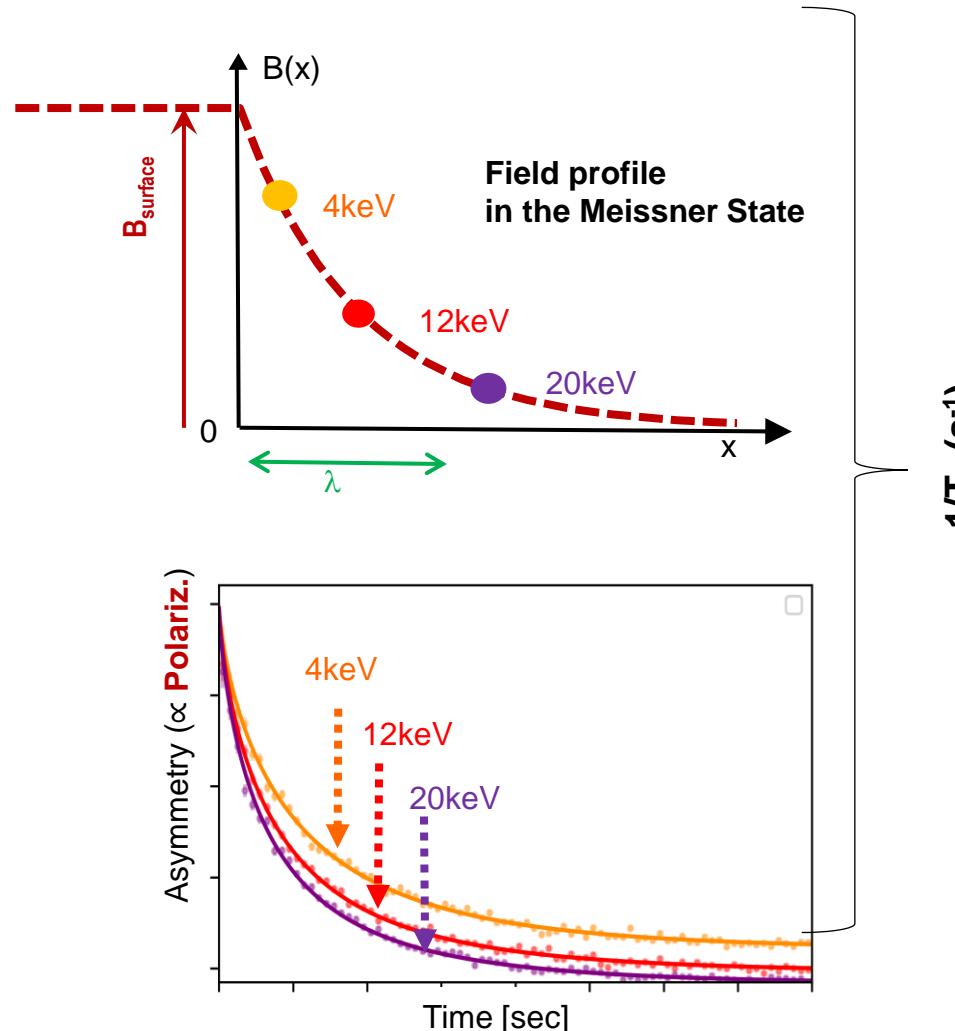
C) Local Field

A) Spin-lattice relaxation rate

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1. Depth-resolved:

- Varies **implantation energies**
- Energy \sim depth



2. Local field sensitivity:

- Probes spins depolarizes in sample \rightarrow direct monitor β -decay
- Char. depol. time = **SLR rate $1/T_1$**
- **Local B-field “slows-down” relaxation**

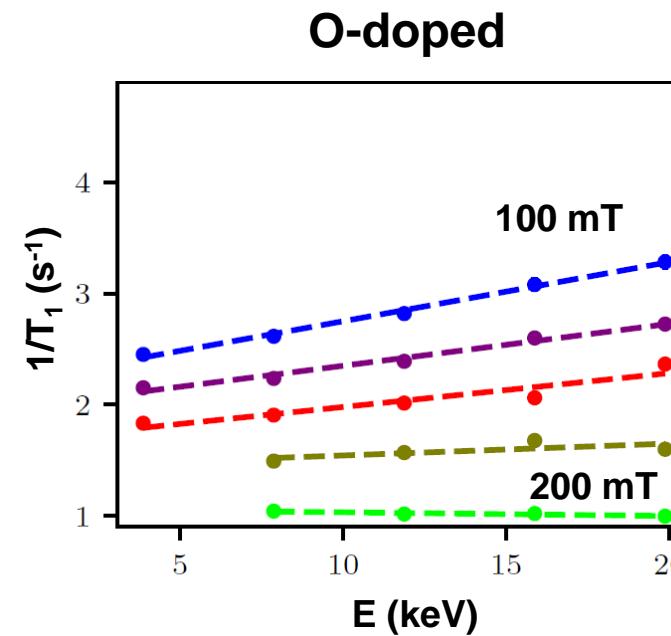
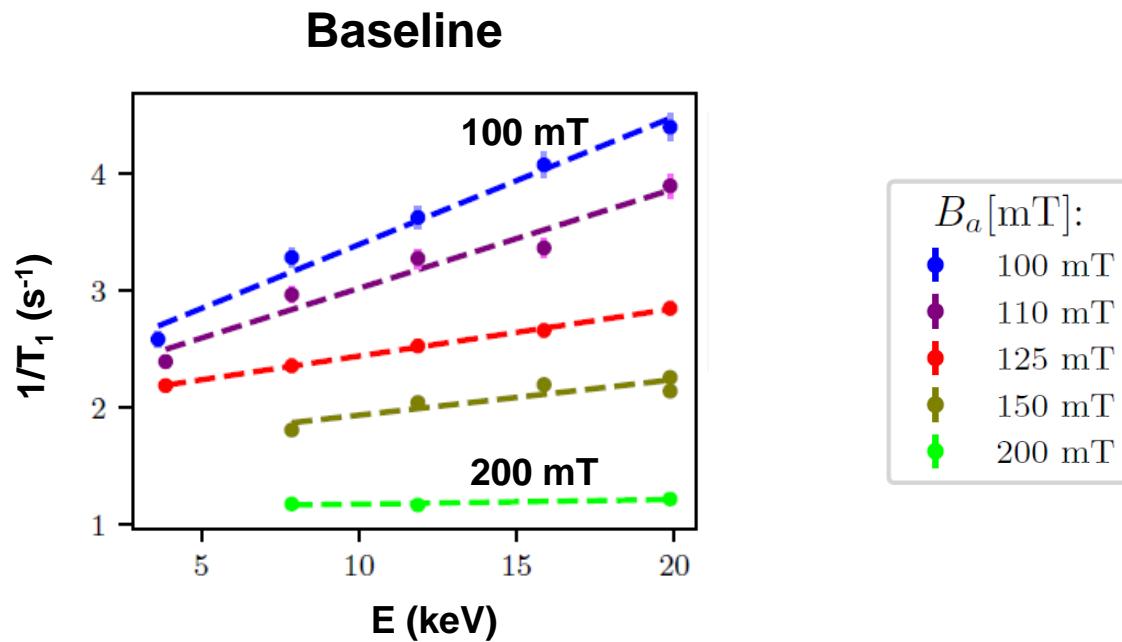
$$\frac{1}{T_1} = \frac{a}{b + B^2}$$

- **Slope = screening**

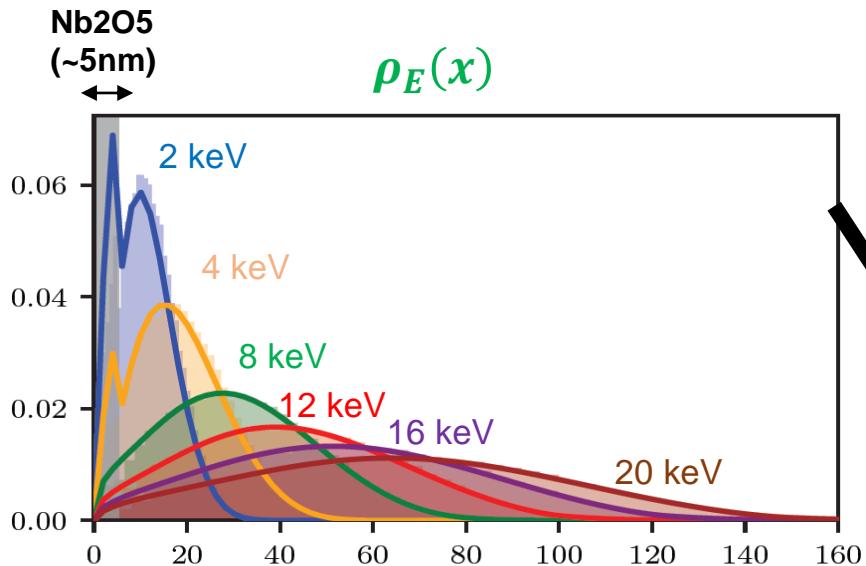
A) Spin-lattice relaxation rate data

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3. Increase B_{applied} :



B) Mapping $1/T_1$ to Average Field $\langle B \rangle_E$



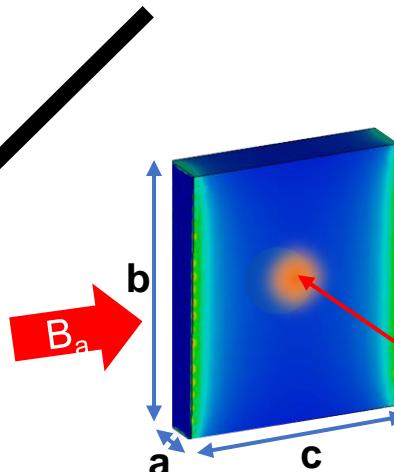
$$B(x) = B_{surf} \exp(-(x - \tilde{d})/\tilde{\Lambda})$$

$\tilde{\Lambda}$: screening penetration depth
 \tilde{d} : dead layer

B_{surf} (Meissner)

- Baseline: 8.1%
- O-doped: 4.3%

beamspot



$\langle B \rangle_E$

Due to stopping distribution:

- Measured = average $\langle 1/T_1 \rangle$
- $B \rightarrow \langle B \rangle_E = \int \rho_E(x) B(x) dx$

$$\left\langle \frac{1}{T_1} \right\rangle (E) = \frac{a}{b + \langle B \rangle_E^2}$$

CALCULATED

VS.

$$\left\langle \frac{1}{T_1} \right\rangle (E)$$

MEASURED

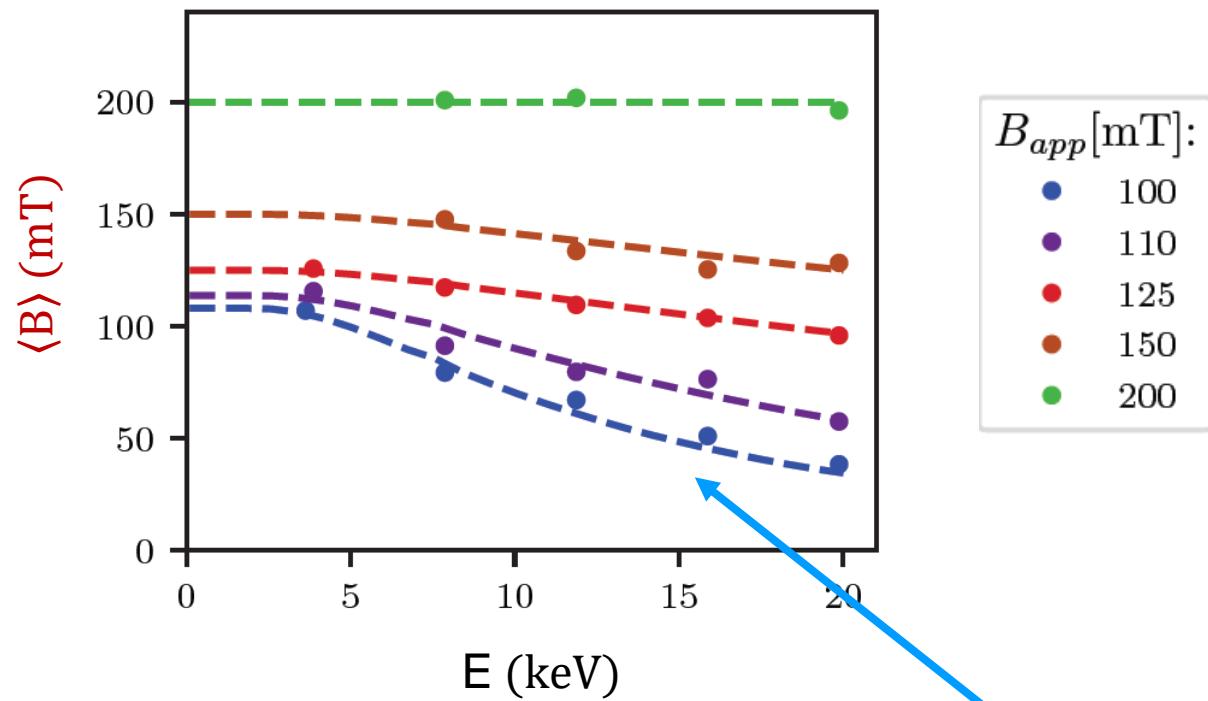
Lorentzian Mapping:
 Applied to other fields – extract $\tilde{\Lambda}(B_a)$, $B_{surf}(B_a)$

B) Average Field Results

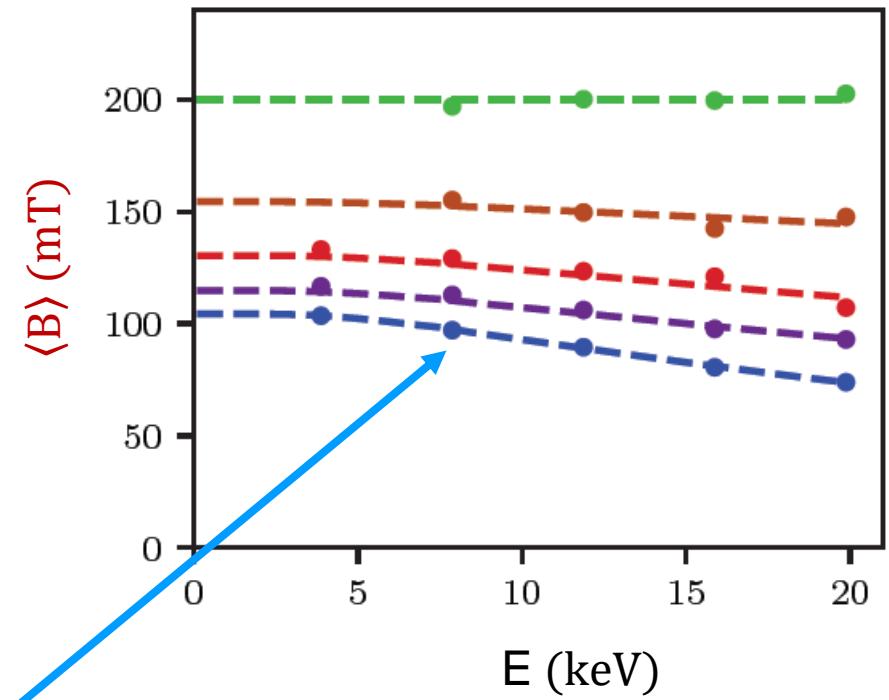
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$$\left\langle \frac{1}{T_1} \right\rangle (E) = \frac{a}{b + \langle B \rangle_E^2}$$

Baseline

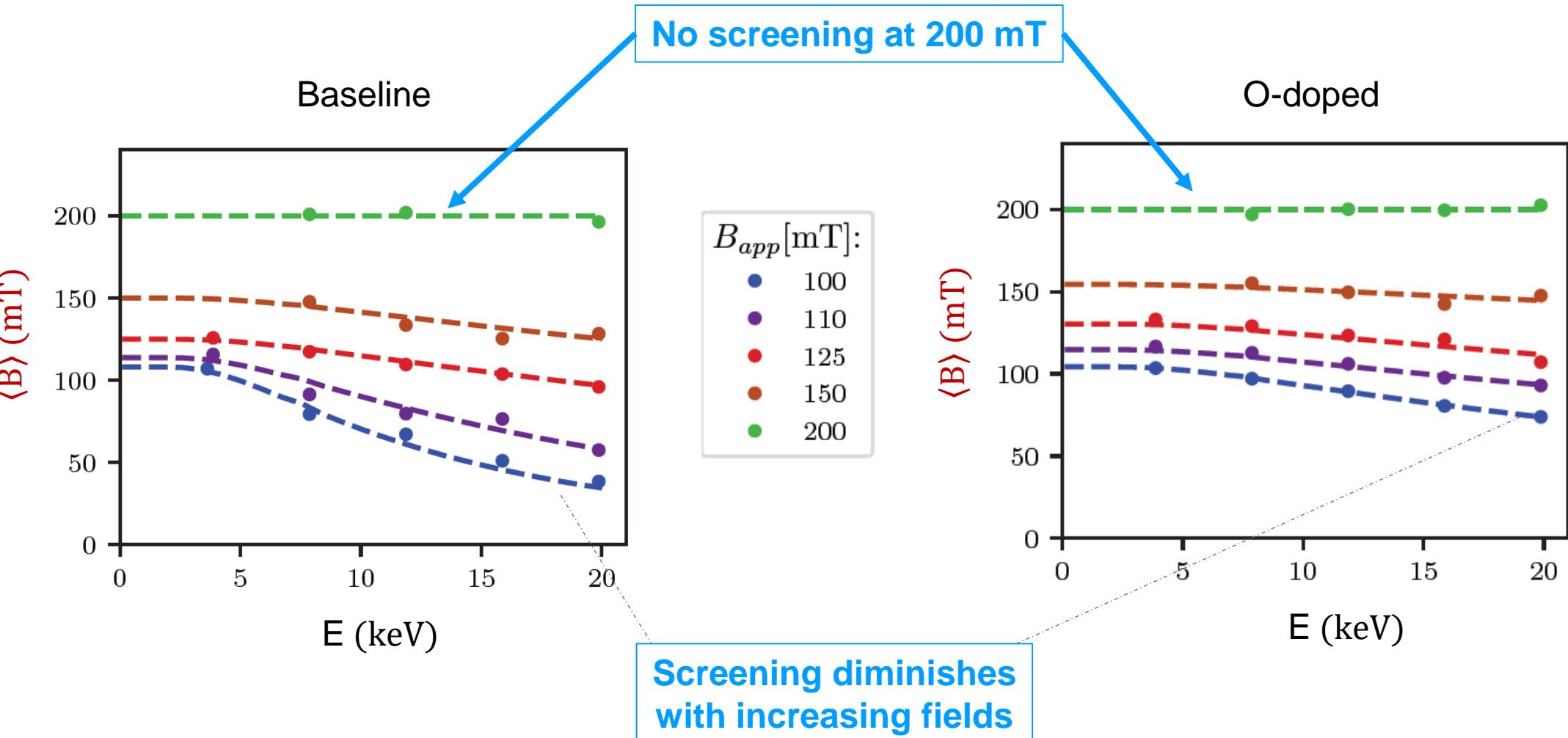


O-doped



Different screening
between the two samples

B) Average Field Results



C) Local Field Analysis

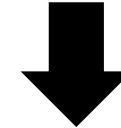
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More refined analysis:

- Directly average $\langle 1/T_1 \rangle_E$ instead of $\langle B \rangle_E$
 - individual ion relaxes with rate $1/T_1(x)$
 - No E-mapping: $\langle 1/T_1 \rangle_E \rightarrow \langle B \rangle_E$
- Map $B(x)$ to $1/T_1(x)$: more accurate modified Lorentzian \mathcal{L}^* (different probe vs. Nb nuclear spins)
- Same $B(x)$ model

$$B(x) = B_{\text{surf}} \exp(-(x - \tilde{d})/\tilde{\Lambda})$$

$$\frac{1}{T_1}(x) = \mathcal{L}^* [B(x)]$$



100 mT, 200 mT

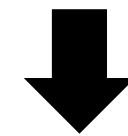
$$\left\langle \frac{1}{T_1} \right\rangle_E = \int \rho_E(x) \frac{1}{T_1}(x) dx$$

Calculated

VS.

$$\left\langle \frac{1}{T_1} \right\rangle(E)$$

Measured

Best \mathcal{L}^* 

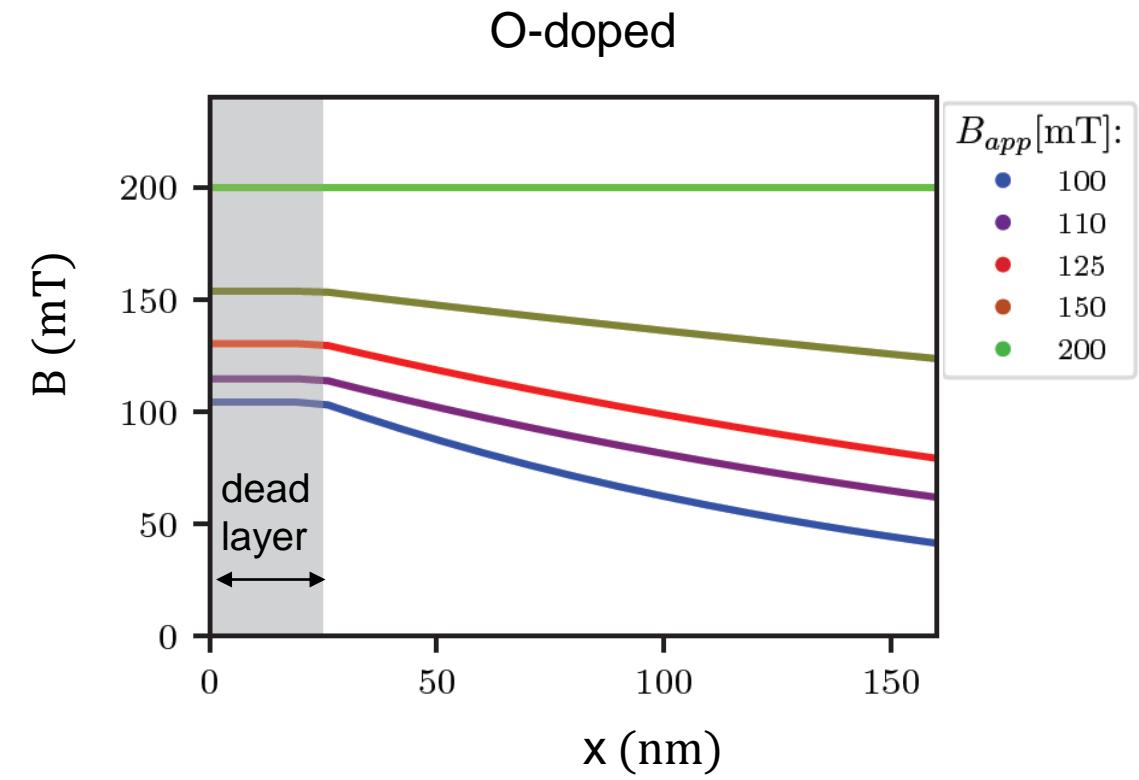
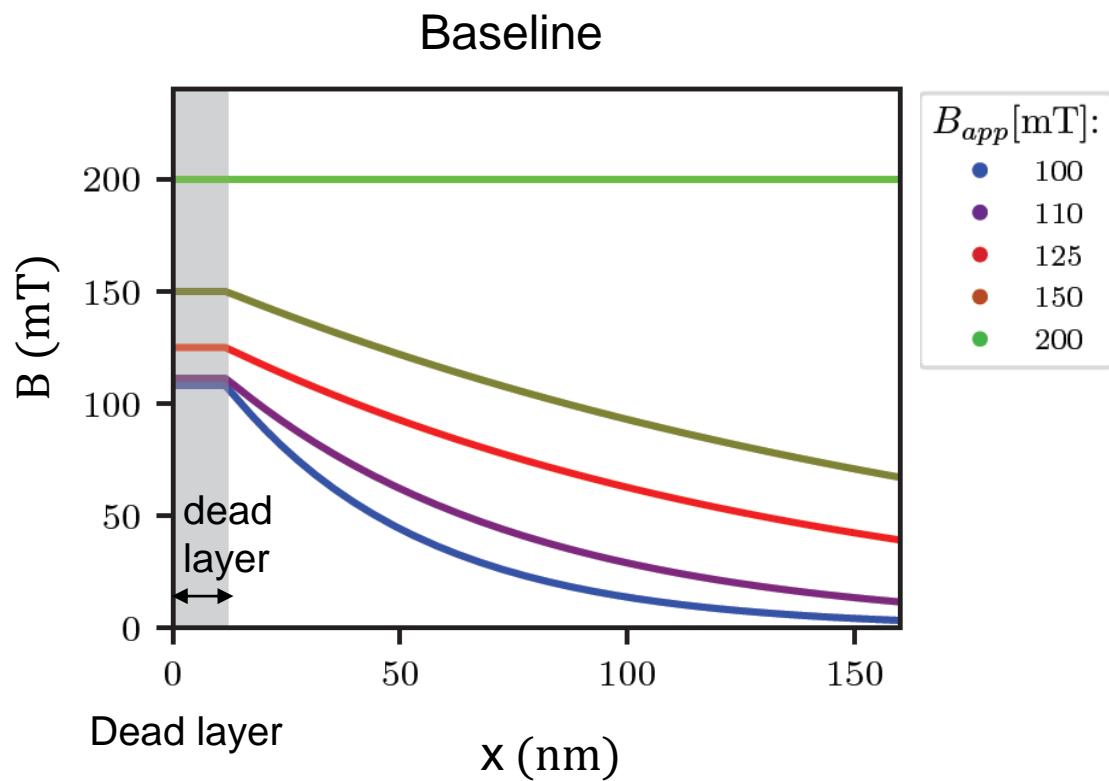
Using best-fit \mathcal{L}^* , extract:

- $\tilde{\Lambda}(B_{\text{app}})$: screening penetration depth
- $B_{\text{surf}}(B_{\text{app}})$: enhanced field at surface

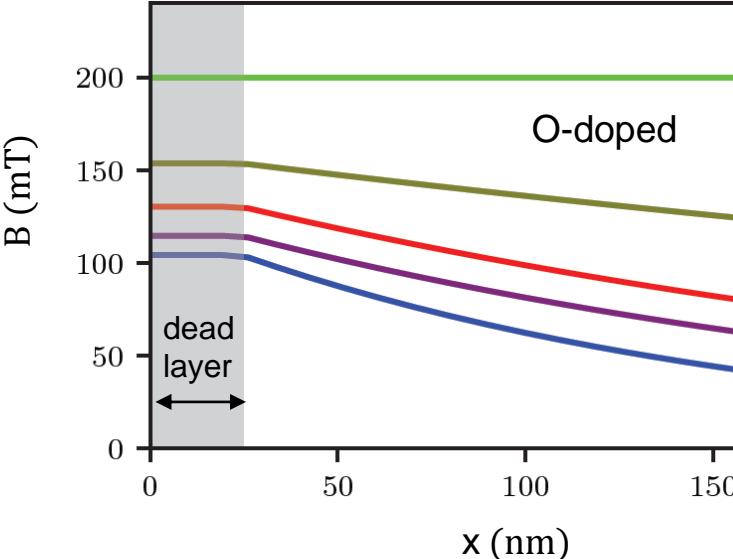
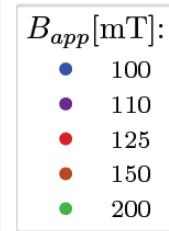
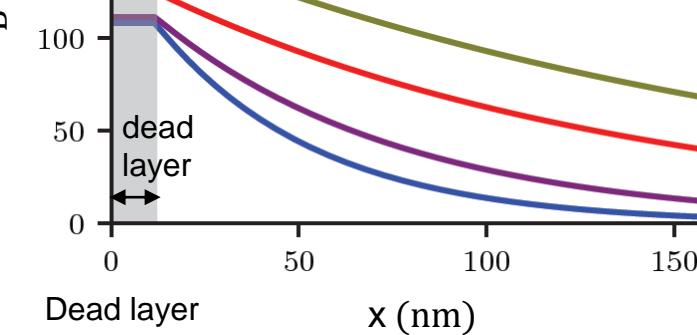
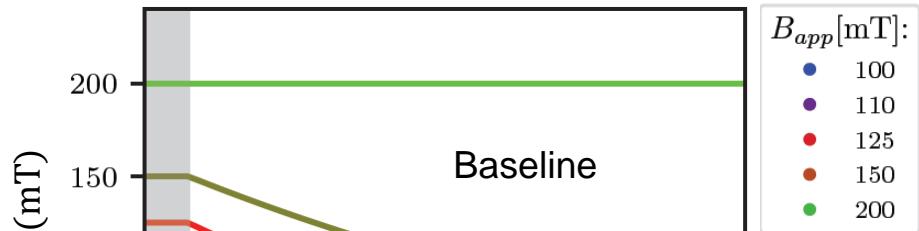
C) Local Field Analysis Results

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$$\left\langle \frac{\mathbf{1}}{\mathbf{T}_1} \right\rangle_E = \int \rho_E(x) \mathcal{L}^* [\mathbf{B}(x)] dx$$

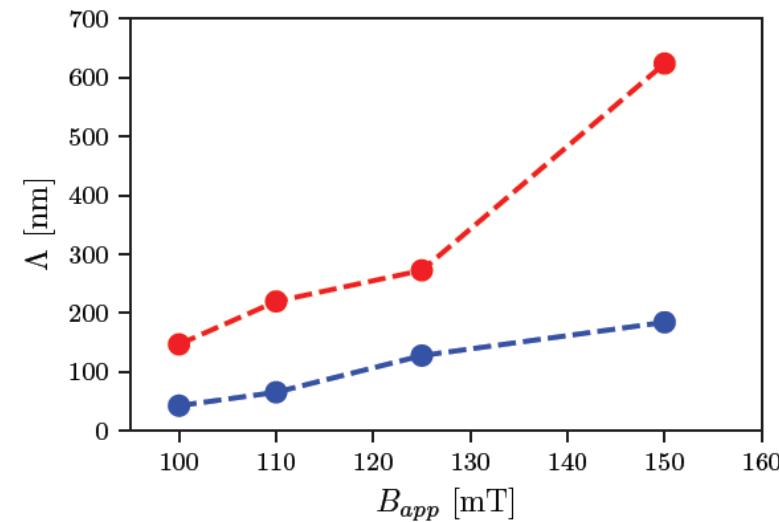


Results: Field-dependent Screening Penetration Depth (Λ)



B_{app}	Λ (Baseline)	Λ (O-doped)
(mT)	(nm)	(nm)
100	43	147
110	66	220
125	128	273
150	184	624

$$\lambda(0 K) = \lambda_L \sqrt{1 + \frac{\xi_0}{l}}$$



$\Lambda = \langle \lambda, n_v(x) \rangle$

- λ : Meissner
- $n_v(x)$: Vortex

Results Highlights: Meissner Region

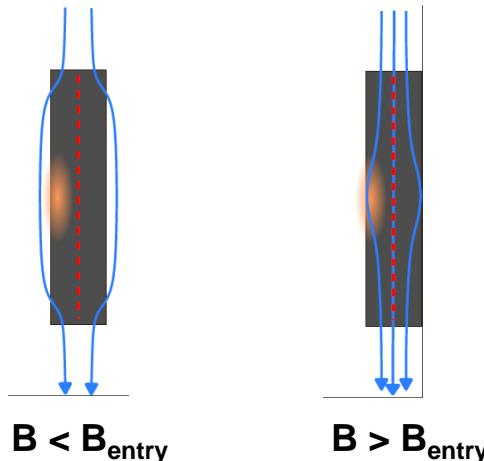
Flux entry detection via

- ❖ reduced surface enhancement B_{surf}
- ❖ Define χ [0,1], ~ flux free volume fraction :
 - $\chi = 1$ (Meissner)
 - $\chi < 1$ (mixed state)
 - $\chi = 0$ (vortex state)

Brandt's formulation:

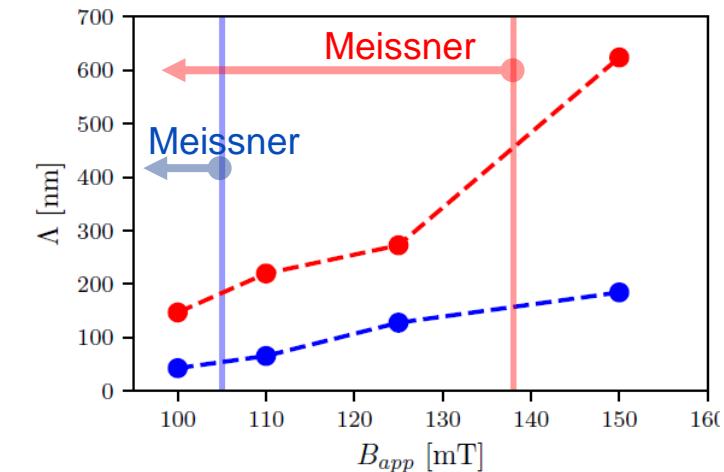
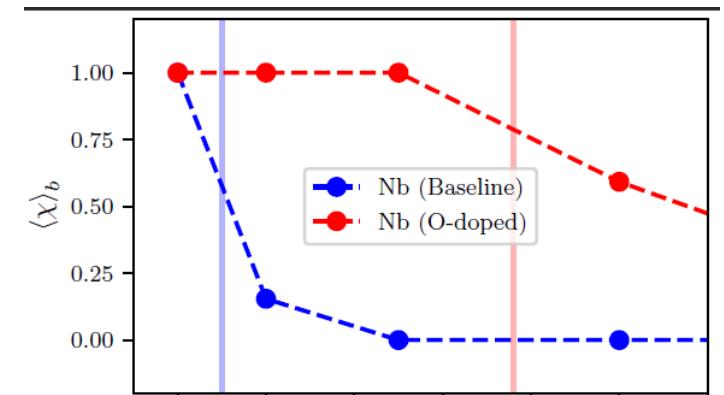
$$B_{entry} = B_{c1} \times \tanh(\sqrt{0.36 * c/a})$$

- Baseline: 121 mT
- O-doped: 127 mT



B_{app} (mT)	χ	SPD (nm)
100	1	43
110	0.15	66
125	0	128
150	0	184
200	0	∞

B_{app} (mT)	χ	SPD (nm)
100	1	147
110	1	220
125	1	273
150	0.59	624
200	0	∞



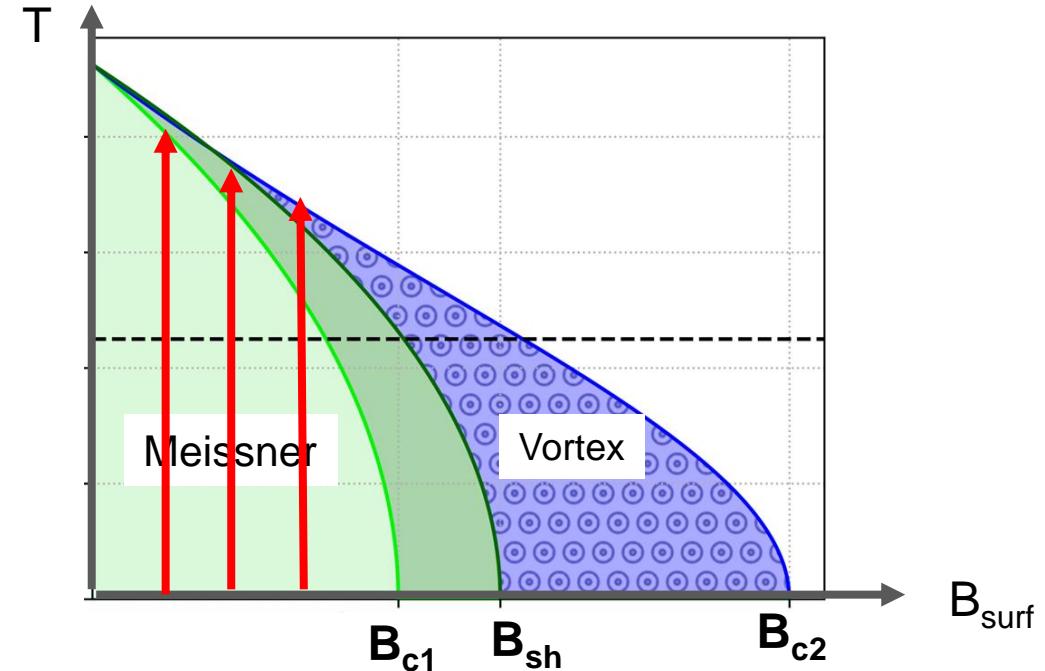
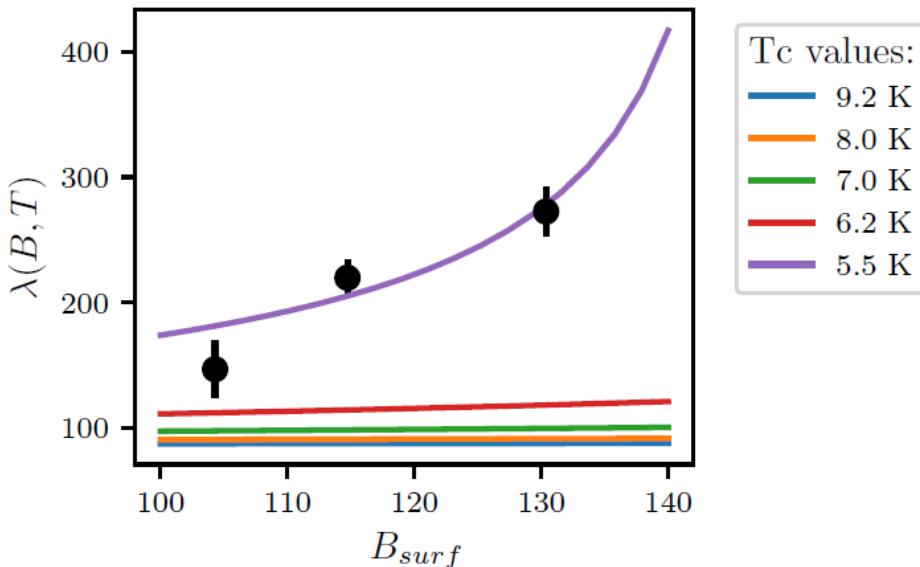
Interesting Features: Strong Field Dependence (O-doped)

- Meissner region @ $B = 100\text{-}125$ mT:
- Field dependence due to T_c

- Two-fluid model: λ diverges at T_c

$$\lambda(T) = \frac{\lambda(0)}{\sqrt{1-(T/T_c)^4}},$$

$$T_c(B) = T_c(B=0) \sqrt{\frac{1-(B/B_{c2})}{1+(B/B_{c2})}}$$



Variation of T_c due to high O-concentration (O-doping) ?

- T_c varies with O-concentration

[C.C. Koch, et al. Phys. Rev. B 9, 888 (1974)]

- Best-fit value of T_c , but seems very low

Interesting Features (2): Nonlinear Meissner Effect

- Nonlinear Meissner Effect: additional increase due to **quadratic field dependence**

$$\lambda(T, B) = \lambda(T, 0) \left[1 + \beta \left(\frac{B_{surf}}{B_c} \right)^2 \right]$$

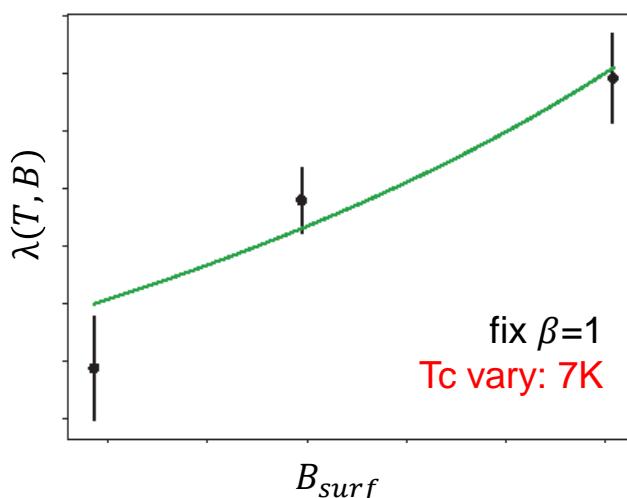
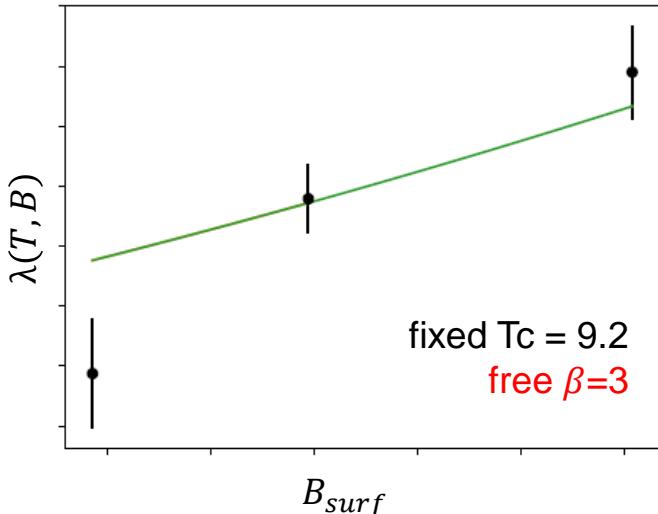
- from Ginzburg-Landau (GL) theory:

$$\beta = \frac{\kappa(\kappa+2^{3/2})}{8(\kappa+2^{1/2})^2}$$
 [J. Makita, et al. Phys. Rev. Research 4, 013156 (2022)]

- predicts $\beta \sim 0.12$. Best-fit at $\beta \sim 3$, **prefactor 25x larger than GL**

- combined NLME + reduced Tc:**

- if β_{max} bounded to $\beta = 1 + \text{vary } T_c$, best fit for $T_c \sim 7\text{K}$ due to e.g., O-concentration ~ 2 [at%] [C.C. Koch, et al. Phys. Rev. B 9, 888 (1974)]



- Localized vortex nucleation ?

Summary & Future Outlook

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Demonstrates β -SRF ability for:

- Clear differences in penetration depth between different treatment
- Clear evolution from Meissner state into mixed state
- Strong field dependence of penetration depth
- Proof of principle + potential to shed new light of SRF materials
- Details → manuscript in preparation

Recent studies:

1. SIS multilayer: NbTiN/AlN/Nb:

Md. Asaduzzaman + T. Junginger (Y. Kalboussi + T. Proslie, CEA Saclay)

- ❖ Measured at perp. field spectrometer (4.1 T) → study vortex state
- ❖ Characterize suitable measurement conditions for β -SRF beamline ($B_{\parallel} \leq 200$ mT)

2. Nb thin film + Nb oxide (Qubit): Fermilab SQMS + TRIUMF SRF/CMMS

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