Tau physics Intensity Frontier Workshop

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Enter a fill

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Tau physics: theory

What Yuval said:"We would like to invite you to give a talk at the
meeting. The idea of the talk will be to give an
overview off tau physics, and we hope that you
will put your personal interest into the talk and
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Tau physics: theory

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3 days before...



What I heard: Dear Martin, now you not only have to survive till April 26th without getting any sleep, but you also have to prepare an overview talk about tau physics...

M. González-Alonso

Tau physics: theory

- Introduction;
- □ Lepton universality;
- □ Leptonic decays: Michel parameters;



□ LFV decays;

- CPV in tau decays;
- Lots of things left out:
 LNV in tau decays, tau physics in hadron colliders, tau g-2, tau EDM, ...



Introduction

Discovered in 1975 by M. Perl et al. at SLAC-LBL: U particle!

1. INTRODUCTION

"Evidence for, and properties of, the new charged heavy lepton" (M. Perl @ XII Recontre de Moriond, 1977)

It is very likely that a mass $1.90 \pm .10 \text{ GeV/c}^2$ charged heavy lepton has been found in e⁺e⁻ annihilation. This paper summarizes the evidence for the new lepton and then presents what we know so far about its properties. The summary of the evidence is very brief; the original papers must be consulted for details on the event selection criteria, on the elimination of conventional explanations, and for the background calculations.

When we first found evidence for this new particle -- the eµ events in the data of the SLAC-LBL Magnetic Detector Collaboration¹ -- we called the particle U as a temporary name because its nature was unknown. Since there is now substantial evidence that it is a lepton, we wish to designate it by a lower case Greek letter. We use τ^{-} because it appears to be the third charged lepton to be found and $\tau \rho \tau \sigma v$ means third in Greek. We feel the old use of τ to designate the three pion decay mode of the K is now obsolete.

Introduction



Discovered in 1976 by M. Perl et al. at SLAC-LBL

- Mass: 1776.82(16) MeV (PDG2012)
- Lifetime: 2.096(10)x10⁻¹³ s. (PDG2012)
- (Lots of) semileptonic decays!
- Enormous progress in tau physics in since then! (CLEO, LEP, Babar, Belle, ...)
 - Early years: checking tau as a standard lepton;
 - **Better data: det. of fundamental parameters & QCD studies;**
 - Recently: tau = tool to search for NP (rare decays, final states in hadron colliders, ...)





Experiment	Number of τ Pairs
LEP	~3×10⁵
CLEO	~1×10 ⁷
BaBar	~5×10 ⁸
Belle	~9×10 ⁸

Sw. Banerjee (TAU2010)



Outline

- Introduction;
- Lepton universality;
- □ Leptonic decays: Michel parameters;
- □ (Inclusive) Hadronic tau decays within the SM (and beyond!);
- □ LFV decays;
- CPV in tau decays;
- Lots of things left out: LNV in tau decays, tau physics in hadron colliders, tau g-2, tau EDM, ...

• E-mu universality tested at 0.14% from tau leptonic BRs (0.28% achieved in Z decays!)



• What about the third family?



Summary: universality tested at the 0.1% level and good agreement...

- ... with too exceptions (\sim 3 sigma level anomalies):
 - Old one: W decay
 - New one: B decay

$$R(D^{(*)}) = \frac{\Gamma(\overline{B} \to D^{(*)}\tau^{-}\overline{\nu}_{\tau})}{\Gamma(\overline{B} \to D^{(*)}l^{-}\overline{\nu}_{l})}$$

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Tau physics

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$$R(D^{(*)}) = \frac{\Gamma(\overline{B} \to D^{(*)}\tau^{-}\overline{\nu}_{\tau})}{\Gamma(\overline{B} \to D^{(*)}l^{-}\overline{\nu}_{l})}$$

$$R_{\tau\ell}^W = \frac{2 \operatorname{BR} \left(W \to \tau \,\overline{\nu}_{\tau} \right)}{\operatorname{BR} \left(W \to e \,\overline{\nu}_e \right) + \operatorname{BR} \left(W \to \mu \,\overline{\nu}_{\mu} \right)} = 1.077(26) \qquad \text{SM result: } 0.999 \dots \text{ [2.8 } \sigma\text{]}$$

Winter 2005 - LEP Preliminary



L3

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Winter 2005 - LEP Preliminary





Why there aren't one thousand papers* in the arXiv explaining this anomaly?



EFT with [U(2)xU(1)]⁵ flavor symmetry:

$$\mathcal{L}_{eff.} = \mathcal{L}_{SM} + \frac{1}{\Lambda^2} \sum \alpha_i \mathcal{O}_i \longrightarrow \Gamma(W \to \tau \overline{\nu}_{\tau}) = \frac{G_F M_W^3}{6\sqrt{2}\pi} \left(1 - w_{\tau}^2\right)^2 \left(1 + \frac{w_{\tau}^2}{2}\right) \left(1 + 4\alpha_{hL}^3\right)$$
[Buchmüller & Wyler'86,
Leung et al.'86]
$$\begin{bmatrix}O_{hL}^3 = i \left(h^{\dagger} D_{\mu} \tau^i h\right) \left(\overline{L} \gamma^{\mu} \tau^i L\right) + \text{h.c.}\end{bmatrix}$$

$$w_{\tau} = m_{\tau}/M_W$$

$$w_{\tau} = m_{\tau}/M_W$$

$$\psi_{\tau} = m_{\tau}/M_W$$

$$w_{\tau} = m_{\tau}/M_W$$

$$w_{\tau} = r_{\tau}/M_W$$

0.02

Leptonic τ decays

0.01

 $2\alpha_{\rm hL}^3 = (g_\tau - g_e)/g_e$

3-4 TeV!!!

LEPII

0.03

0.04

-0.01

0.00

EFT with [U(2)xU(1)]⁵ flavor symmetry:

What if we have 2 (or more) operators?

All 17 operators that are $[U(2)xU(1)]^5$ inv. affecting the 3^{rd} family are not able to accommodate the discrepancy!

Relaxing the symmetry to U(2)⁵: 2 more operators... still no room!

$$O_{\tau B}^{t} = (\overline{L} \sigma^{\mu\nu} \tau) h B_{\mu\nu} + \text{h.c.} \qquad R_{\tau}^{V}$$
$$O_{\tau W}^{t} = (\overline{L} \sigma^{\mu\nu} \tau^{i} \tau) h W_{\mu\nu}^{i} + \text{h.c.} \qquad \text{LE}$$

 $R_{ au\ell}^W = 1.01 \pm 0.01$ LEPII result: 1.077(26) SM result: 0.999...

(Filipuzzi, MGA & Portolés, 2012)



EFT with [U(2)xU(1)]⁵ flavor symmetry:

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$$\begin{bmatrix} O_{hL}^3 = i \left(h^{\dagger} D_{\mu} \tau^i h \right) \left(\overline{L} \, \gamma^{\mu} \, \tau^i L \right) + \text{h.c.} \end{bmatrix}$$

$$w_{\tau} = m_{\tau} / M_W$$

Conclusion:

The loophole in the naïve thinking (cancellations) is really small...

Thus...

- highly non-trivial flavor structure;
- EFT assumptions are wrong (e.g. new light fields)
- Or it's just a statistical fluctuation (or systematic...);



Lepton universality in B decays

$$R(D^{(*)}) = \frac{\Gamma(\overline{B} \to D^{(*)}\tau^{-}\overline{\nu_{\tau}})}{\Gamma(\overline{B} \to D^{(*)}l^{-}\overline{\nu_{l}})}$$



- BaBar & Belle agree!
- Statistical errors dominate for both exp!
- Updates expected for both exp! (& SuperBs!!)
- Also tensions in $B \rightarrow \tau v$ (ameliorated by recent Belle result)

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$$R(D^{(*)}) = \frac{\Gamma(\overline{B} \to D^{(*)}\tau^{-}\overline{\nu_{\tau}})}{\Gamma(\overline{B} \to D^{(*)}l^{-}\overline{\nu_{l}})}$$



New physics?

- Type II 2HDM excluded [*BaBar'2012*];
- MSSM with MFV excluded [Altmannshofer et al'2012];
- It can be accommodated in other SM extensions: type III 2HDM [*Crivellin et al.'2012*], Aligned 2HDM [*Celis et al.'2012*], EFT approach (within/beyond MFV) [*Fajfer et al.'2012, Datta et al.'2012, Biancofiore et al, 2013*], ..."

□ Just SM3

- Recent 2+1 flavor lattice calculation : R(D)_{SM} = 0.316±0.012±0.007, [*FNAL/MILC*, 2012];
- Also: R(D)_{SM} = 0.31±0.02 [Becirevic et al.'2012];
 - \rightarrow Discrepancy below $2\sigma!$

BABAR	0.440 ± 0.071	0.332 ± 0.029	
SM	0.297 ± 0.017	0.252 ± 0.003	

R(D)

(^{*}Q) 8(0)

0.2

0.2

0.4

[*] Kamenik & Mescia'08; Fajfer et al.'12; BaBar'12

$t = tan\beta / m_{H^+}$ (GeV⁻¹)

BABAR 2HDM

Tau physics

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- Also: R(D)_{SM} = 0.31±0.02 [Becirevic et al.'2012];
 - → Discrepancy below 2σ ...

	R(D)	R(D*)
BABAR	0.440 ± 0.071	0.332 ± 0.029
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R(D)

(* 0.4

0.2

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BABAR 2HDM

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Tau leptonic decays

- **Observables:**
 - Total rates \rightarrow lepton universality tests;
 - Angular & energy distribution (Michel parameters):



$$\frac{d^2\Gamma}{dxd(\cos\theta)} \propto x^2 \left\{ (3-3x) + \frac{2}{3}\rho(4x-3) + \frac{6}{3}\eta_0 \frac{(1-x)}{x} + \mathcal{P}_{x}\xi \cos\theta \left[(1-x) + \frac{2}{3}\delta(4x-3) \right] \right\},$$
*more parameters if the electron/muon polarization is measured
$$x = \frac{E_e}{E_{\text{max}}}, \qquad x_0 = \frac{m_e}{E_{\text{max}}},$$

Theory side:

$$L_{eff} = -\frac{4G_F}{\sqrt{2}} \Big[\overline{e} \gamma_{\mu} (1 - \gamma_5) v_e \cdot \overline{v_{\tau}} \gamma_{\mu} (1 - \gamma_5) \tau + \sum_{i} \varepsilon_{i} \overline{e} \Gamma v_e \cdot \overline{v_{\tau}} \Gamma \tau \Big]$$

$$\rho = 3/4 + f(\varepsilon_i)$$

$$\eta = 0 + f(\varepsilon_i)$$

$$\xi = 1 + f(\varepsilon_i)$$

$$\delta = 3/4 + f(\varepsilon_i)$$

Agreement SM/exp

$$\Rightarrow$$
 bounds over epsilons.

In the 80s-90s, a few papers about the interest for specific NP models (Stahl'94, Pich & Silva'95, ...)

 $E_{\rm max}$

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Tau physics

Tau leptonic decays: an EFT approach



and muon! Tau leptonic decays: an EFT approach

$$L_{eff} \approx -\frac{4G_F}{\sqrt{2}} \left(1 + 2\operatorname{Re}\varepsilon_L\right) \left[\overline{e}\gamma_\mu (1 - \gamma_5)v_e \cdot \overline{v}_\tau \gamma_\mu (1 - \gamma_5)\tau + \varepsilon_R \overline{e}(1 - \gamma_5)v_e \cdot \overline{v}_\tau (1 + \gamma_5)\tau\right]$$

Deviations from (quark-)lepton universality

Remarkably simple result! Consequences:

- Only some Michel parameters are "interesting"
- 2) We can calculate the precision needed to compete with LEP [interesting for (super)-B factories]

 $\varepsilon_{R} = +2[\hat{\alpha}_{le}]_{2112}$

$$O_{le} = \left(\bar{l}\gamma^{\mu}l\right)(\bar{e}\gamma_{\mu}e),$$

Energy & angular distribution (Michel parameters!)

$$O_{le} = (\bar{l}\gamma^{\mu}l)(\bar{e}\gamma_{\mu}e), = \bar{v}_{\tau}\gamma_{\mu}(1-\gamma_{5})v_{e}\cdot\bar{e}\gamma^{\mu}(1+\gamma_{5})\tau + \tau\gamma_{\mu}(1-\gamma_{5})e\cdot\bar{e}\gamma^{\mu}(1+\gamma_{5})\tau$$
$$\bar{\tau}_{L}\tau_{R}\cdot\bar{e}_{R}e_{L}$$

(Cirigliano & MGA, 2010)

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The tau is the only lepton heavy enough to decay hadronically. Moreover, its mass is significantly larger than QCD scale!!! (key point) This allows precise analytical studies of inclusive observables (via OPE):

 α_S , V_{us} and low-energy parameters.

Hadronic tau decays

I'll focus on the following inclusive observables:

$$\Gamma(\tau \rightarrow v_{\tau} + hadrons_{S=0})$$

$$\Gamma(\tau \rightarrow v_{\tau} + hadrons_{S\neq0})$$

LEP measured with precision not only the total BRs but also the energy distribution of the hadronic system (spectral functions!!). This triggered enormous QCD activity!

Very active field: experimentally and theoretically;





• Extraction of α_s and V_{us} . The idea is simple:

$$R_{\tau} = \frac{\Gamma(\tau \rightarrow v_{\tau} + hadrons)}{\Gamma(\tau \rightarrow v_{\tau} e^{-} v_{e})} \approx N_{C}$$

$$R_{\tau} = R_{\tau}^{S=0} + R_{\tau}^{S\neq0} \approx N_{C} |V_{ud}|^{2} + N_{C} |V_{us}|^{2} \approx 2.85 + 0.15$$

Tau physics

QCD switch

$$O$$

 ON
 $OFFF$
 $(\alpha_S=0)$
 O





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$$R_{\tau} = \frac{\Gamma(\tau \to v_{\tau} + hadrons)}{\Gamma(\tau \to v_{\tau} \ e^{-} \ v_{e})} \approx N_{C} \quad [\text{exp: } \sim 3.628(9)]$$

$$R_{\tau} = R_{\tau}^{S=0} + R_{\tau}^{S\neq0} \approx N_{C} |V_{ud}|^{2} + N_{C} |V_{us}|^{2} \approx 2.85 + 0.15$$

$$[\text{exp: } \sim 3.467(8) + 0.161(3)]$$

Tau physics



Hadrons

(Braaten et al, 1992)

$$R_{\tau} = \frac{\Gamma(\tau^{-} \to v_{\tau} + \text{had})}{\Gamma(\tau^{-} \to v_{\tau} e^{-} \overline{v_{e}})} = 12\pi \int_{0}^{1} dx (1-x)^{2} \left[(1+2x) \operatorname{Im} \Pi^{(1)}(x m_{\tau}^{2}) + \operatorname{Im} \Pi^{(0)}(x m_{\tau}^{2}) \right] \qquad x = \frac{s}{m_{\tau}^{2}}$$

$$= 6\pi i \oint_{|x|=1} dx (1-x)^{2} \left[(1+2x) \Pi^{(0+1)}(x m_{\tau}^{2}) - 2x \Pi^{(0)}(x m_{\tau}^{2}) \right] \qquad x = \frac{s}{m_{\tau}^{2}}$$

$$= N_{C} S_{EW} \left(1 + \delta_{P} + \delta_{NP} \right) \qquad \Pi^{(r)}(s) = \sum_{D=2n} \frac{C_{D}^{(r)}(s,\mu) \langle O_{D}(\mu) \rangle}{(-s)^{D/2}}$$

$$S_{EW} = 1.0201 (3) \quad \text{Marciano-Sirlin, Braaten-Li, Erler}$$

$$\delta_{P} = a_{\tau} + 5.20 \ a_{\tau}^{2} + 26 \ a_{\tau}^{3} + 127 \ a_{\tau}^{4} + \dots \approx 20\%$$

$$\delta_{NP} = -0.0059 \pm 0.0014 \qquad a_{\tau} = a_{s}(m_{\tau})/\pi$$

$$(Baikov et al., '08)$$

Intense theoretical activity!!!

Davier et al'08, Maltman & Yavin'08, Beneke & Jamin'08, Maltman'09, Menke'09, Caprini & Fischer'09, Descotes-Genon & Malaescu'09, Cvetic et al'10, Abbas et al.'12, Beneke et al'12, Boito et al'12, ...

Exp: still using LEP data!
 A more precise measurement of the spectral functions would be welcome;



(Braaten et al, 1992)

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Tau physics

 $R_{\tau}^{S=0} \approx N_C \left| V_{ud} \right|^2 (1+\delta)$ $R_{\tau}^{S \neq 0} \approx N_C \left| V_{us} \right|^2 (1 + \delta) - \delta R_{\tau}$

SU(3) breaking correction! Calculable theoretically = $f(m_a)$ [Pich & Prades'98,'99] $\delta R_{\tau} = 0.239(30)$

[Gamiz'2013] [Maltman'2010: larger error]

$$|V_{us}|^{2} = \frac{R_{\tau}^{S \neq 0}}{\frac{R_{\tau}^{S=0}}{|V_{ud}|^{2}} - \delta R_{\tau}}$$

[Gamiz et al'03, '05, '08]

$$V_{us} = 0.2173 \pm 0.0020_{exp} \pm 0.0010_{th}$$

 $|V_{us}| = 0.2255(13) \text{ [from } K_{13}\text{]}$

$$R_{\tau,S}^{00} = 0.1612 (28)$$
$$R_{\tau,V+A}^{00} = 3.4671 (84)$$
$$|V_{ud}| = 0.97425 (22)$$

NOTE: BRs from B-factories are smaller than previous world averages!



K₁₃ decays, FlaviaNet 2010 0.2254 ± 0.0013 K₁₂ decays, FlaviaNet 2010 0.2252 ± 0.0013 CKM unitarity 0.2255 ± 0.0010 $\tau \rightarrow K\nu / \tau \rightarrow \pi\nu$, HFAG 2012 0.2229 ± 0.0021 $\tau \rightarrow K\nu$, HFAG 2012 0.2214 ± 0.0022 $\tau \rightarrow$ s inclusive, HFAG 2012

τ average, HFAG 2012 0.2202 ± 0.0015



[Interlude: CKM unitarity test]

 Apart from the importance per se (fundamental parameter!!!), it is important for NP searches: CKM unitarity tests!

 $|V_{ud}|^2 + |V_{us}|^2 + |V_{ub}|^2 = 0.9999 \pm 0.0006$

Shown in specific NP scenarios, and within an EFT approach (with MFV): [Barbieri et al. (1985), Marciano & Sirlin (1987), Hagiwara et al. (1995), Kurylov & Ramsey-Musolf (2002), Bauman, Erler & Ramsey-Musolf (2012) ...]

 $\Lambda_i^{eff} > 11 \text{TeV} (90\% \text{ CL})$

$$\Delta_{CKM} = 4 \left(-\hat{\alpha}_{\varphi l}^{(3)} + \hat{\alpha}_{\varphi q}^{(3)} - \hat{\alpha}_{lq}^{(3)} + \hat{\alpha}_{ll}^{(3)} \right) = -(1 \pm 6) \cdot 10^{-4}$$

[Cirigliano, MGA & Jenkins, 2010]

5x stronger than the constraint from LEP & other EWPO!!!

$$O_{ll}^{(3)} = \frac{1}{2} (\bar{l}\gamma^{\mu}\sigma^{a}l)(\bar{l}\gamma_{\mu}\sigma^{a}l)$$
$$O_{lq}^{(3)} = (\bar{l}\gamma^{\mu}\sigma^{a}l)(\bar{q}\gamma_{\mu}\sigma^{a}q)$$
$$O_{\varphi l}^{(3)} = i(h^{\dagger}D^{\mu}\sigma^{a}\varphi)(\bar{l}\gamma_{\mu}\sigma^{a}l) + \text{h.c.},$$
$$O_{\varphi q}^{(3)} = i(\varphi^{\dagger}D^{\mu}\sigma^{a}\varphi)(\bar{q}\gamma_{\mu}\sigma^{a}q) + \text{h.c.},$$

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Tau physics

[Interlude: CKM unitarity test]

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EFT with a more general NP flavor structure (beyond MFV)

$$V_{us}^{pheno}\Big|_{Kl3} \neq V_{us}^{pheno}\Big|_{\tau(excl)} \neq V_{us}^{pheno}\Big|_{\tau(incl)}$$

Is it possible to accommodate the current tensions without spoiling collider & EWPO?

- Identification of relevant operators;
- Bounds on those operators from EWPO;
- **c** Calculation of the effect of those operators on the V_{us} extraction;

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- Bounds on those operators from EWPO,

Calculation of the effect of those operators on the V_{us} extraction

Not trivial in the case of inclusive tau decays! [By-product: NP in the extraction of a_s ?]

$$\begin{aligned}
\mathbf{Im}(\mathbf{q}^{2}) & \mathbf{The QCDSR derivation} \\
\mathbf{involves new correlators!!} \\
\mathbf{s}_{0} & \mathbf{re}(\mathbf{q}^{2}) \\
\mathbf{ke}(\mathbf{q}^{2}) & \mathbf{ke}(\mathbf{q}^{2}) \\
\delta R_{\tau} &\equiv \frac{R_{\tau,nS}}{|V_{ud}|^{2}} - \frac{R_{\tau,S}}{|V_{us}|^{2}} \\
\end{aligned}$$

$$\begin{aligned}
R_{\tau} &= 6\pi i |V_{ud}|^{2} \oint_{|s|=m_{\tau}^{2}} \frac{ds}{m_{\tau}^{2}} \left(1 - \frac{s}{m_{\tau}^{2}}\right)^{2} \times \\
\times \left\{ |\kappa_{V}|^{2} \left[\left(1 + \frac{2s}{m_{\tau}^{2}}\right) \Pi_{VV}^{(1)}(s) + \Pi_{VV}^{(0)}(s) \right] \\
+ |\kappa_{A}|^{2} \left[\left(1 + \frac{2s}{m_{\tau}^{2}}\right) \Pi_{AA}^{(1)}(s) + \Pi_{AA}^{(0)}(s) \right] \\
+ 2\operatorname{Re}(\kappa_{V}\kappa_{S}^{*}) \underbrace{\Pi_{VS}(s)}_{m_{\tau}} + 2\operatorname{Re}(\kappa_{A}\kappa_{P}^{*}) \underbrace{\Pi_{AP}(s)}_{m_{\tau}} \\
+ 6 \times 2\operatorname{Re}(\kappa_{V}\kappa_{T}^{*}) \underbrace{\Pi_{VT}(s)}_{m_{\tau}} \right\} \times \left[1 - 2\tilde{v}_{L} \right],
\end{aligned}$$

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Lepton Flavor Violation

- □ LFV is an "accidental" symmetry in the SM (m_v=0) This accident does not occur in general in bSM theories.
- □ Neutrino have mass \rightarrow LFV! \rightarrow Charged LFV... but zero for all practical purposes.

$$\frac{\tau}{W_{\tau}} \underbrace{\nabla \tau}_{W_{\tau}} \underbrace{\nabla \mu}_{W_{\tau}} \underbrace{\mu}_{V_{\tau}} \underbrace{\Psi}_{V_{\tau}} \underbrace{\mu}_{V_{\tau}} \underbrace{\mu}_{V_{\tau}}$$

- LFV processes probe physics at scales Λ>>TeV.
 Moreover, realistic NP models often produce LFV at visible rates.
- Muon-LFV is more constrained experimentally... but we don't know the NP flavor structure!! Tau/muon LFV is model-dependent → model-diagnostic!! Several NP models find interesting tau/muon LFV correlations. This will become a key issue if MEG finds a NP signal.

Lepton Flavor Violation



Different tau-LFV processes are also model-diagnostic!

LFV in example models (Hitlin, IF2011)		$\tau \rightarrow \mu \gamma \ \tau \rightarrow \ell \ell \ell$	
SM + v mixing	Lee, Shrock, PRD 16 (1977) 1444 Cheng, Li, PRD 45 (1980) 1908	Undetectable	
SUSY Higgs	Dedes, Ellis, Raidal, PLB 549 (2002) 159 Brignole, Rossi, PLB 566 (2003) 517	10-10	10-7
SM + heavy Maj $v_{\rm R}$	Cvetic, Dib, Kim, Kim, PRD66 (2002) 034008	10-9	10-10
Non-universal Z'	Yue, Zhang, Liu, PLB 547 (2002) 252	10-9	10-8
SUSY SO(10)	Masiero, Vempati, Vives, NPB 649 (2003) 189 Fukuyama, Kikuchi, Okada, PRD 68 (2003) 033012	10-8	10-10
mSUGRA + Seesaw	Ellis, Gomez, Leontaris, Lola, Nanopoulos, EPJ C14 (2002) 319 Ellis, Hisano, Raidal, Shimizu, PRD 66 (2002) 115013	10-7	10-9

Lepton Flavor Violation



EIC would improve the limits on LFV (e/tau). [Gonderinger & Ramsey-Musolf, 2010]

CPV in tau decays;

- □ Tau physics offers a variety of CPV observables (rate, angular asymmetries, triple products, ...) that can be studied to search for new sources of CPV.
- □ Interesting recent result:

$$A_{\tau} = \frac{\Gamma(\tau^{+} \to \pi^{+}K_{S}\bar{\nu}_{\tau}) - \Gamma(\tau^{-} \to \pi^{-}K_{S}\nu_{\tau})}{\Gamma(\tau^{+} \to \pi^{+}K_{S}\bar{\nu}_{\tau}) + \Gamma(\tau^{-} \to \pi^{-}K_{S}\nu_{\tau})} = (-3.6 \pm 2.3 \pm 1.1) \cdot 10^{-3}$$

$$A_{\tau}^{SM} = (3.6 \pm 0.1) \cdot 10^{-3}$$

$$A_{\tau}^{SM} = (3.6 \pm 0.1) \cdot 10^{-3}$$

$$(\text{due to } K^{0} - \overline{K}^{0} \text{ mixing})$$

$$[Bigi \& Sanda, 2005]$$

$$[Grossman \& Nir, 2011]$$

- Belle has also searched for CPV in this decay mode through a difference in the angular distributions, finding a null result at the 0.2–0.3% level [*Belle, 2011*].
- □ NP explanations? [*Bigi, 2012*]
 - Charged Higgs?
 - W_L - W_R mixing?

LQ?



Tau physics in colliders

Excellent tool for New Physics searches: MSSM Higgses, Z' boson, ...

□ Largest Higgs-lepton coupling (4th Higgs BR)



Summary/Outlook

- $\begin{array}{c} & & & & & \\ & & &$
- Very exciting future in tau physics: Belle II, BESIII, LHCb, ...
- Several intriguing anomalies: Lepton universality (W and B decays), V_{us} extraction, CPV, ...
- Lot of theory work to do too:
 QCD analysis of inclusive hadronic tau decays (pQCD, DV, OPE, ...),
 New Physics studies, ...
- Crowning jewel: tau-LFV searches;
- Lots of (very interesting) topics not addressed in this talk: *low-energy QCD studies, (g-2)*_{tau/muon}, *EDM*_{tau}, *collider studies, …*









Backup slides

M. González-Alonso

Tau physics

NP flavor structure?

$$O_{ql}^{(3)} = i(\varphi^{\dagger}D^{\mu}\sigma^{a}\varphi)(\overline{l}\gamma_{\mu}\sigma^{a}l) + h.c.$$

$$\alpha_{ql}^{(3)}O_{ql}^{(3)} = \sum_{\alpha,\beta=1}^{3} [\alpha_{ql}^{(3)}]_{\alpha\beta} [O_{ql}^{(3)}]_{\alpha\beta} + h.c. = i(\varphi^{\dagger}D^{\mu}\sigma^{a}\varphi) \left((\overline{l_{e}} \quad \overline{l_{\mu}} \quad \overline{l_{\tau}})\gamma_{\mu}\sigma^{a} (\alpha_{21} \quad \alpha_{22} \quad \alpha_{23}) (\beta_{\mu} \quad \beta_{\mu}) + h.c. \right)$$
• Symmetry of the SM for m_{f1}=m_{f2}=m_{f3}=0 \rightarrow U(3)⁵

$$\alpha_{ql}^{(3)} = \begin{pmatrix} \alpha_{11} \quad \alpha_{12} \quad \alpha_{13} \\ \alpha_{21} \quad \alpha_{22} \quad \alpha_{23} \\ \alpha_{21} \quad \alpha_{22} \quad \alpha_{23} \\ \alpha_{31} \quad \alpha_{32} \quad \alpha_{33} \end{pmatrix} = \overline{\alpha}_{ql}^{(3)} \begin{pmatrix} 1 \quad 0 \quad 0 \\ 0 \quad 1 \quad 0 \\ 0 \quad 0 \quad 1 \end{pmatrix} \qquad \text{Lepton univ. holds by def.}$$

• A way of singularizing the 3rd family is relaxing the symmetry to $[U(2)xU(1)]^5$.

$$\alpha_{\varphi l}^{(3)} = \begin{pmatrix} \alpha_{11} & \alpha_{12} & \alpha_{13} \\ \alpha_{21} & \alpha_{22} & \alpha_{23} \\ \alpha_{31} & \alpha_{32} & \alpha_{33} \end{pmatrix} = \begin{pmatrix} \alpha_{h\ell}^3 & 0 & 0 \\ 0 & \alpha_{h\ell}^3 & 0 \\ 0 & 0 & \alpha_{hL}^3 \end{pmatrix}$$

• "Less symmetry": U(2)⁵ [SM symmetry for $m_{f1}=m_{f2}=0$, but $m_{f3}\neq 0$]

