

Stepan Stepanyan JLAB Intensity Frontier Workshop 25-27 April 2013, ANL





HPS at JLAB

HPS experiment at JLAB will search for A'

- in the scattering of high energy (1.1 GeV, 2.2 GeV, and 6.6 GeV), high intensity (~500 nA) electron beams on tungsten target (0.125% r.l.)
- in the mass range from 20 MeV to 1000 MeV
- for couplings ε² > 10⁻⁷ with bump hunt and ε² <5x10⁻⁸ with displaced decay vertex search (*unique to HPS*)
- in the decay modes to e^+e^- and $\mu^+\mu^-$ (*unique to HPS*)

HPS will use a large acceptance forward spectrometer in experimental Hall-B at JLAB







Fixed target experiments at JLAB

Jefferson Lab - Precision and intensity frontier!



 $E_{max} = 12 \text{ GeV} (2.2 \text{ GeV/pass})$ $I_{max} = 100 \mu A$ P = 85% Simultaneous delivery of CW beam to 3 Halls











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- The HPS detector based on a 3-magnet chicane, with the second dipole as the analyzing magnet. It will detect and identify electrons and muons produced at angles θ >15 mr
- Detector package includes: 6-layer Silicon Vertex Tracker (SVT) installed inside the analyzing magnet vacuum chamber, Electromagnetic Calorimeter (ECal) and the muon system installed downstream of the analyzing magnet
- To avoid "wall of flame", crated by Multiple Coulomb scattered beam particles and radiative secondaries, the detectors will be split into two identical parts, installed above and below the "dead zone" (beam plane)





HPS apparatus - SVT

Precise measurements of momentum and production vertex of charged particles

- Will be installed in the vacuum inside the analyzing magnet
- First layer is located at 10 cm from the target, the silicon in the first layer is only 0.5 mm from the center of the beam
- First 3-layers are retractable
- Silicon will be actively cooled to remove heat and retard radiation damage
- The sensors have 60(30) μm readout(sense) pitch (hit position resolution ~6 μm)
- The sensors are read out continuously at 40 MHz using the APV25 chip

Layer number	1	2	3	4	5	6
nominal z , from target (cm)	10	20	30	50	70	90
Stereo Angle (mrad)	100	100	100	50	50	50
Bend-plane resolution (μm)	≈ 60	≈ 60	≈ 60	$\approx \! 120$	pprox 120	pprox 120
Non-bend resolution (μm)	≈ 6	≈ 6	≈ 6	≈ 6	≈ 6	≈ 6
Number of sensors	4	4	4	8	8	8
Nominal dead zone in $y \pmod{mm}$	± 1.5	± 3.0	± 4.5	± 7.5	± 10.5	± 13.5
Module power consumption (W)	6.9	6.9	6.9	13.8	13.8	13.8







HPS apparatus - ECal

Electron triggering and electron identification

- Lead-tungstate calorimeter with 442 16 cm long crystals (1.3x.1.3 cm² cross section) with APD readout (Hamamatsu S8664-55)
- In each half, crystals are arranged in rectangular formation in 5 layers, 4 layers have 46 crystals and one (closest to the beam) has 37
- Modules are located inside the thermal enclosure with temperature stability <1°C
- Readout and trigger are based on JLAB FADC250
- Pulse height, spatial and timing information from each crystal are available for the trigger decision
- Expected energy resolution $\sigma/E\approx4.5\%/\sqrt{E}$







HPS apparatus – Muon system

Muon trigger and muon identification.

- Four double layer (XY) scintillator hodoscopes sandwiched between Fe absorbers (30/15/15/15 cm)
- Optimized for π/μ rejection in momentum range 1 GeV to 4 GeV
- Readout and trigger are based on JLAB FADC250



The expected low background and high detection efficiency make the di-muon final state an attractive complement to the $e_{-}e_{-}$ final state







Electron beam parameters

- The HPS will use up to 500nA, 1.1 GeV, 2.2 GeV, and 6.6 GeV electron beams incident on a thin tungsten (W) target.
- The vertex resolution will benefit from a small beam size (< 50 μm) in the non-bend plane, Y direction, while the momentum measurement will not benefit from small beam sizes in the X direction
- Asymmetric beam profile is desirable (a small beam sizes in both dimensions will overheat the target foil)

The beamline optimizations performed for the 12 GeV CEBAF machine, including the proposed changes for Hall-B operations, demonstrated that required beam parameters are achievable



Parameter	Ree	Unit		
Е	1100	2200	6600	MeV
$\delta E/E$		$< 10^{-4}$		
Current	< 200	< 400	< 500	nA
Current Instability		< 5		%
σ_x		$\mu { m m}$		
σ_y		$\mu { m m}$		
Position Stability		$\mu { m m}$		
Divergence		$\mu \mathrm{rad}$		
Beam Halo $(> 5\sigma_Y)$		$< 10^{-5}$		

The same optics optimization program was proven to work well for 6 GeV CEBAF







HPS test run

The main goal - validate critical assumptions made in our simulations for rates and occupancies

- Large fraction of trigger rates and the tracker occupancies come from multiple Coulomb scattered electrons
- Correct simulation of the electromagnetic background is crucial for the design of the experiment
- Two simulation tools, GEANT4 and EGS5 is gave markedly different results in the rate estimates
- GEANT4 that was used for simulations of the background and trigger rates in the original proposal gives x2 higher rates than EGS5



Other goal of the test run was to demonstrate the feasibility of the proposed apparatus and data acquisition systems









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Performance of the test run apparatus



ECal occupances







Test Run Results: EGS5 is correct

- Multiple Coulomb scattering of beam electrons is the main contributor to the detector occupancies and determines the limits of sensitivity of the experiment
- In test run with photon beam, the angular distribution of the pair produced electrons and positrons emerging from the converter has been studied to validate simulations





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ECal

SVT



Target 0.25% r. l.

Electron beam

Detector performance studies

- Simulation of the detector response in GEANT4 design geometry, realistic energy deposition and pulse formation in SVT, Ecal, and muon detectors
- For trigger simulations, cluster finding algorithm and the trigger logic used in trigger FPGAs are applied to the simulated FADC signal time evolution
- EGS5 was used to simulate electromagnetic backgrounds generated in the target

Limiting factors for luminosity are: SVT Layer-1 occupancy and rates in ECal modules







HPS experimental reach







True muonium with HPS

- HPS experiment has the potential to discover "true muonium", a bound state of a $\mu + \mu -$
- The (μ+μ-) atom is hydrogen-like, and so has a set of excited states characterized by a principal quantum number *n*, with the binding energy of these states is E = -1407 eV/n²
- The (μ+μ-) "atom" can be produced by an electron beam incident on a target, a similar way as the A'
- HPS will discover the 1S, 2S, and 2P true muonium bound states with its proposed run plan
- Search will require a vertex cut at about 1.5 cm to reject almost all QED background events, then look for a resonance at 2m_u

$$N_{\left(\mu^{+}\mu^{-}\right)} = 200 \left(\frac{I}{450nA}\right) \left(\frac{t}{1month}\right)$$

In two weeks of 6.6 GeV run HPS should see ~15 true muonium events

A. Banburski and P. Schuster, Phys. Rev. D 86, 093007 (2012)





Summary

- HPS is one of the three experiments proposed at JLAB to search for hidden sector photons
- It is the only experiment so far that has capability to reach couplings of ϵ^2 <5x10⁻⁸ and search in the $\mu^+\mu^-$ decay mode
- Since the first presentation at JLAB PAC37 (2011), HPS succeeded to:
 - be funded and build a test setup (within 9 months)
 - run a test with photon beam (May 2012) and demonstrate feasibility of the proposed detector design
 - validated critical assumptions made in the simulation
- PAC39 liked the results of the test run, rated HPS physics with "A", conditionally approved full HPS, C1, and urged JLAB for physics running
- In addition of heavy photons, HPS is well suited to discover a bound state of a $\mu+\mu-$, a $(\mu+\mu-)$ atom
- HPS is looking forward to commission the new apparatus and run the first phase of the experiment in 2014-2015, if schedules at JLAB admit





Backups







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HPS run strategy

- The first phase will include:
 - a commissioning run end of 2014, 3 weeks with beam, which includes data taking at 1.1 GeV and 2.2 GeV beam energies
 - extensive data taking in 2015, with runs at 2.2 and 6.6 GeV (roughly 4 weeks each)
- The second phase can use remaining beam time any time after §
- purple-dashed: 1 week of 1.1 GeV
- blue-dashed: 1 week of 2.2 GeV
- blue-solid: 3 weeks of 2.2 GeV
- dark-green: 2 weeks of 6.6 GeV, detecting A' →e⁺e⁻,
- light-green: 2 weeks of 6.6 GeV, detecting A' →µ⁺µ⁻
- red: the statistical combination of all of the above
- green-shaded: 3 months each of 2.2 GeV and 6.6 GeV







HPS time-line

- Heavy Photon Search proposal was first presented to JLAB PAC 37 (January 2011), PAC endorsed the test run, and conditionally (C2) approved the full experiment
- The test run detector was installed in Hall-B for parasitic running with photon beams on April 19, 2012. Dedicated data taking on last shift of CEBAF 6 GeV operations
- The JLAB PAC39 (June 2012) graded HPS physics with an "A", approved a commissioning run with electrons (concurred PAC37 decision), and granted C1 approval for the full HPS experiment.
- □ The total requested beam time for the experiment is 180 days
- HPS experiment will be reviewed by DOE/HEP (July, 2013) and if funded, will be ready to take data in fall of 2014. If JLAB 12 GeV schedule permits, production data taking can take place in 2014-2015





Where and how to search for dark photons

Both "naturalness" arguments and fits to astrophysical data suggest $\alpha'/\alpha \equiv \epsilon^2 \sim 10^{-4} - 10^{-10}$ and $m_{\Delta'} \sim MeV - GeV$



A' can be electroproduced the same way as radiative tridents in the fixed target experiment (*J.D. Bjorken, R. Essig, P. Schuster, and N. Toro, Phys. Rev. D80, 2009, 075018*)







Sensitivity with bump hunt

The rate of the A' signal relates to the radiative trident cross-section within a small mass window of width δm as

$$\frac{d\sigma(e^{-}Z \to e^{-}Z(A' \to l^{+}l^{-}))}{d\sigma(e^{-}Z \to e^{-}Z(\gamma^{*} \to l^{+}l^{-}))} = \left(\frac{3\pi\varepsilon^{2}}{2N_{eff}\alpha}\right)\left(\frac{m_{A'}}{\delta m}\right)$$

Mass resolution is an important ingredient for the sensitivity of the experiment

Within the acceptance and signal region for the HPS experiment, the Bethe-Heitler reaction dominates the trident rate by 4:1





$$\frac{S}{\sqrt{B}}\Big|_{bin} = \left(\frac{N_{rad}}{N_{tot}}\right) \sqrt{N_{bin}} \left(\frac{3\pi\varepsilon^2}{2N_{eff}\alpha}\right) \left(\frac{m_{A'}}{\delta m}\right) \eta_{bin}$$

 N_{rad}/N_{tot} , the fraction of radiative events among all QED trident events in the search region is determined by simulation





Displaced decay vertex search

A search for resonances that decay with cm-scale displaced vertices opens up sensitivity to much smaller couplings



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Prompt γ cτ=3.5mm

γ ct=14mm

γ ct=35mm