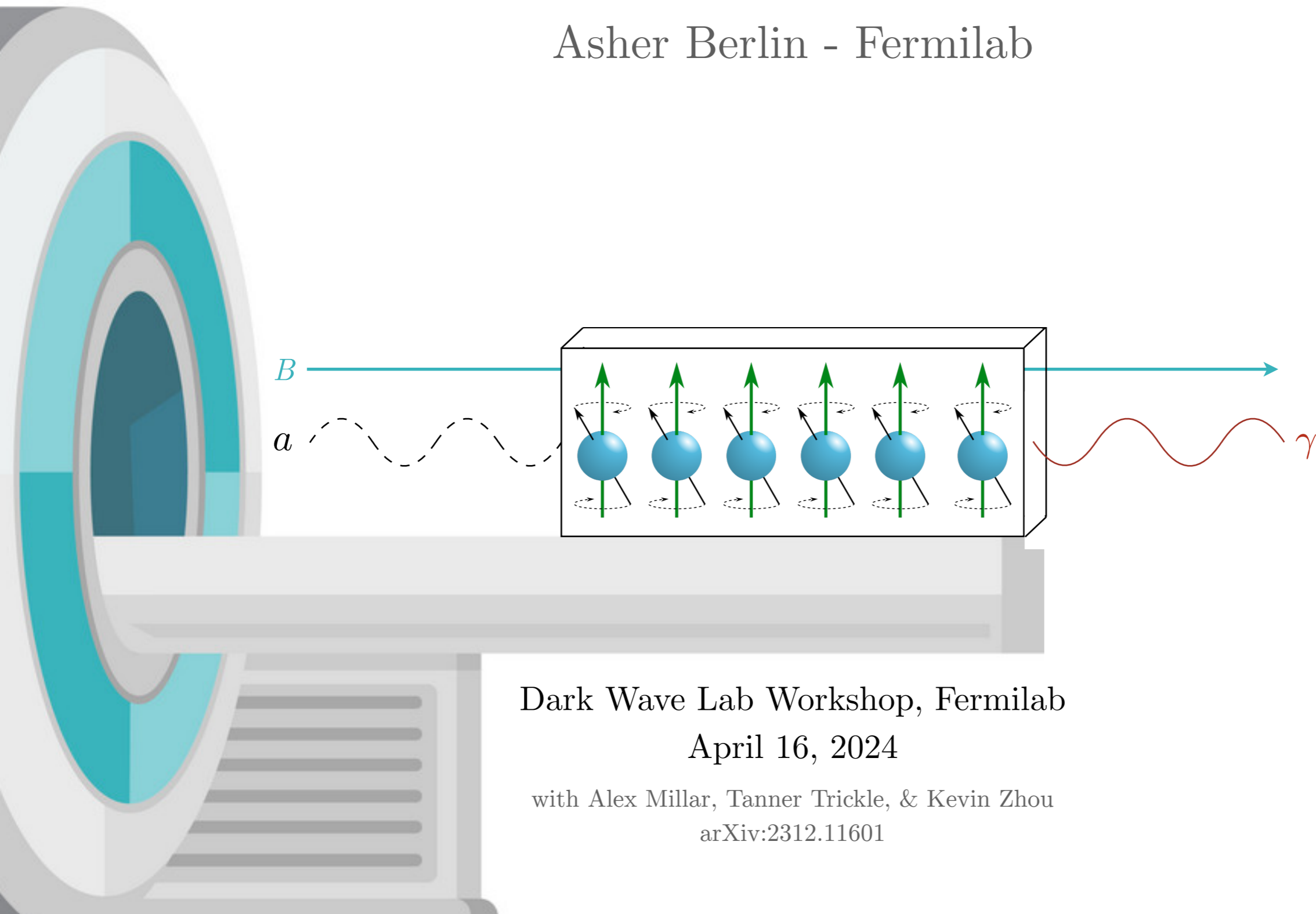


Magnetized Stacks for Electron-Coupled Axions

Asher Berlin - Fermilab



Dark Wave Lab Workshop, Fermilab

April 16, 2024

with Alex Millar, Tanner Trickle, & Kevin Zhou

arXiv:2312.11601

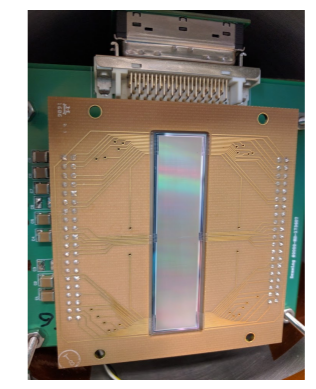
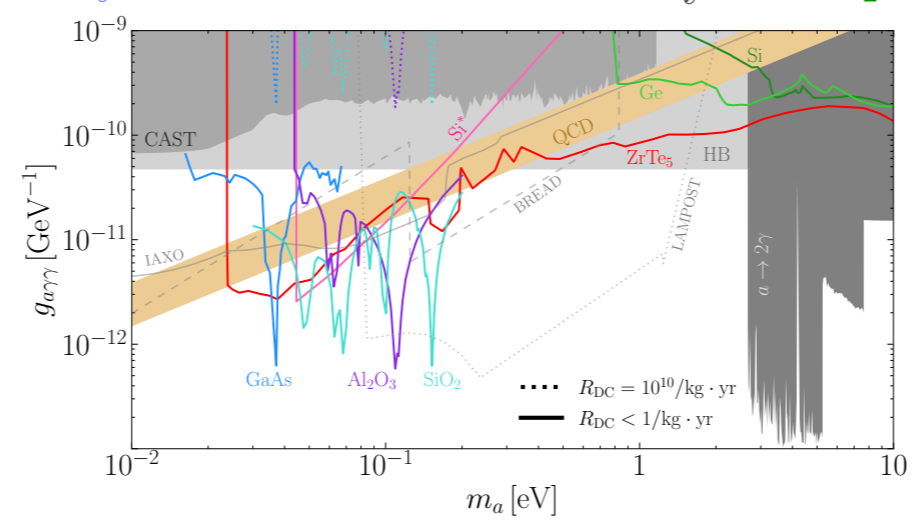
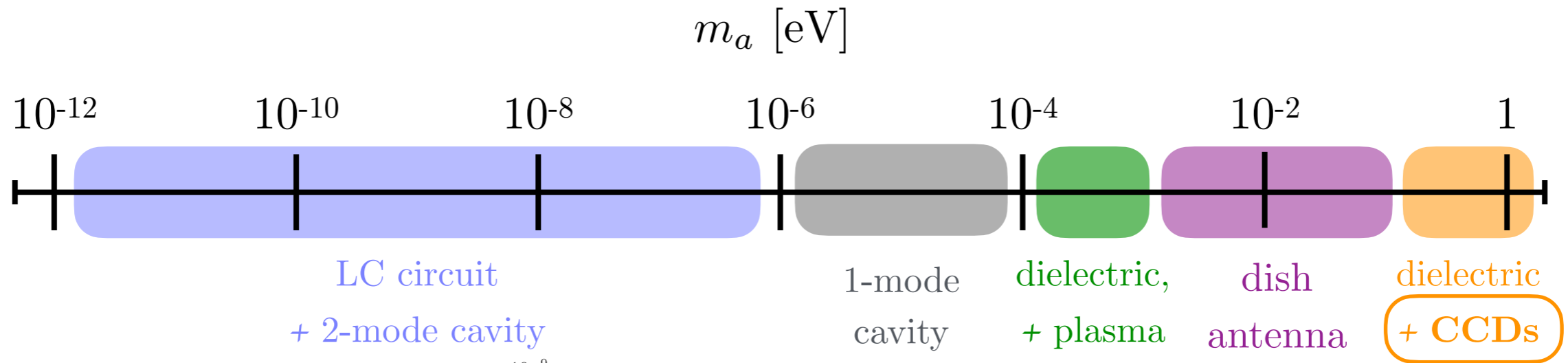
Axion Couplings

$$\mathcal{L} = \frac{\partial_\mu a}{f_a} \left(J_{\text{EM}}^\mu + J_{\text{QCD}}^\mu + J_{\text{spin}}^\mu + \dots \right)$$

Photon-Coupling

$$\mathcal{L} = \frac{\partial_\mu a}{f_a} \left(J_{\text{EM}}^\mu + J_{\text{QCD}}^\mu + J_{\text{spin}}^\mu + \dots \right)$$

Experiments that are potentially sensitive to canonical QCD axion

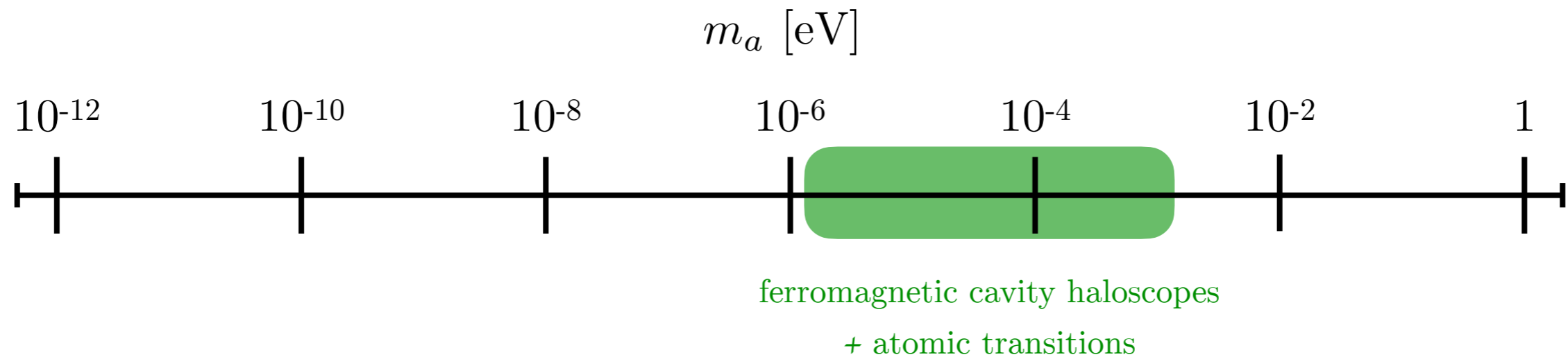


with Tanner Trickle
arXiv:2305.05681,
Phys. Rev. Lett.

Electron-Coupling

$$\mathcal{L} = \frac{\partial_\mu a}{f_a} \left(J_{\text{EM}}^\mu + J_{\text{QCD}}^\mu + J_{e\text{-spin}}^\mu + \dots \right)$$

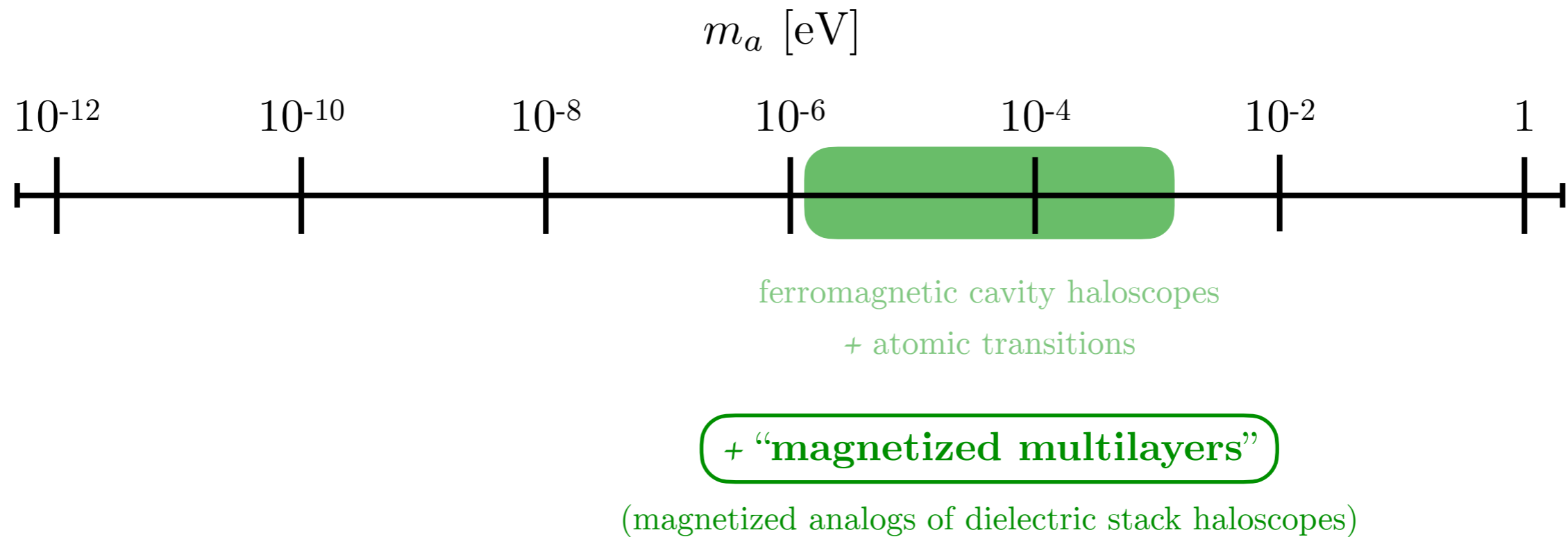
Experiments that are potentially sensitive to canonical QCD axion



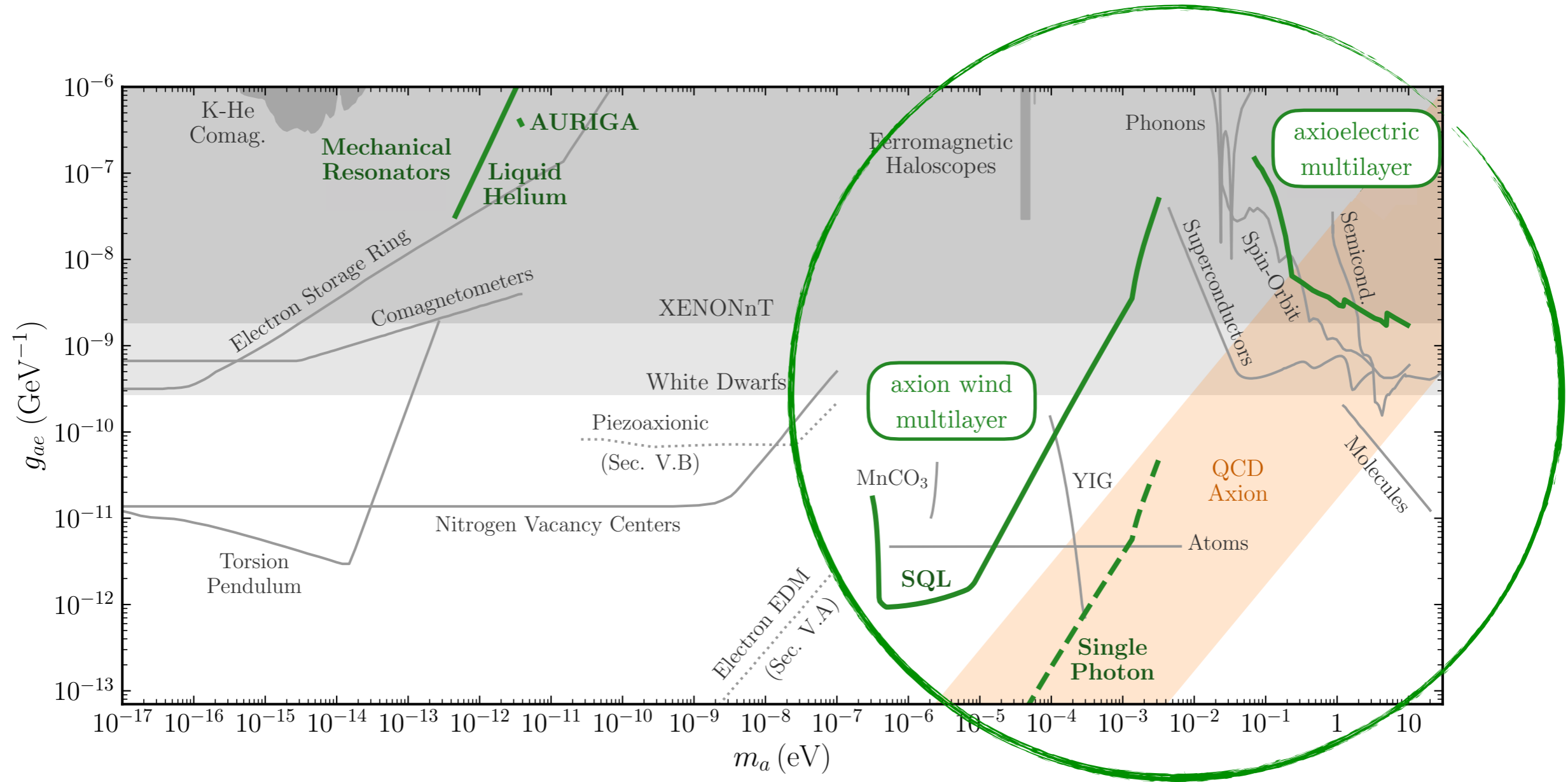
Electron-Coupling

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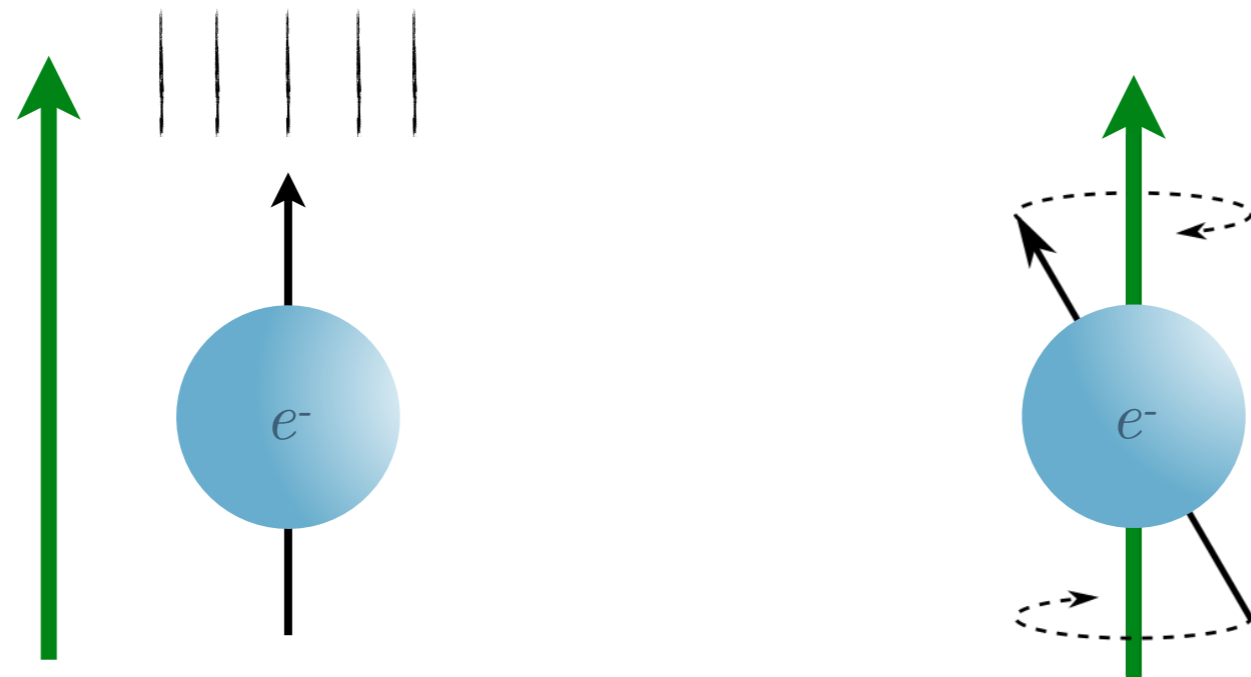
Experiments that are potentially sensitive to canonical QCD axion



Electron-Coupling



“magnetized multilayers”



I. Spin-Coupled Axions

Electron-Coupling

axion-electron interaction

$$L = \int d^3\mathbf{x} g_{ae} (\partial_\mu a) \langle \bar{e} \gamma^\mu \gamma^5 e \rangle$$

Electron-Coupling

axion-electron interaction

$$\begin{aligned}
 L &= \int d^3\mathbf{x} g_{ae} (\partial_\mu a) \langle \bar{e} \gamma^\mu \gamma^5 e \rangle \\
 &\simeq g_{ae} (\partial_\mu a) (\mathbf{v}_e \cdot \hat{\mathbf{s}}_e, \hat{\mathbf{s}}_e)^\mu \\
 &= g_{ae} (\nabla a) \cdot \hat{\mathbf{s}}_e + g_{ae} (\partial_t a) \hat{\mathbf{s}}_e \cdot \mathbf{v}_e
 \end{aligned}$$

Electron-Coupling

axion-electron interaction

$$\begin{aligned}
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 &= g_{ae} (\nabla a) \cdot \hat{\mathbf{s}}_e + g_{ae} (\partial_t a) \hat{\mathbf{s}}_e \cdot \mathbf{v}_e
 \end{aligned}$$

analogous to normal electromagnetism

$$L_{\text{EM}} = -\mu_B \mathbf{B} \cdot \hat{\mathbf{s}}_e - e \mathbf{A} \cdot \mathbf{v}_e$$

Electron-Coupling

axion-electron interaction

$$\begin{aligned}
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 \end{aligned}$$

analogous to normal electromagnetism

$$L_{\text{EM}} = -\mu_B \mathbf{B} \cdot \hat{\mathbf{s}}_e - e \mathbf{A} \cdot \mathbf{v}_e$$

$$\implies \boxed{\mathbf{B}_{\text{eff}} = -\frac{g_{ae}}{\mu_B} \nabla a, \quad \mathbf{A}_{\text{eff}} = -\frac{g_{ae}}{e} \partial_t a \hat{\mathbf{s}}_e}$$

Electron-Coupling

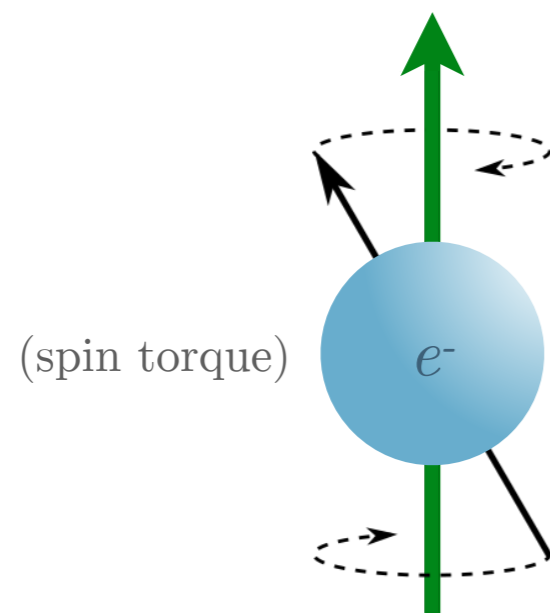
$$\mathbf{B}_{\text{eff}} = -\frac{g_{ae}}{\mu_B} \nabla a \quad , \quad \mathbf{A}_{\text{eff}} = -\frac{g_{ae}}{e} \partial_t a \hat{\mathbf{s}}_e$$



“Axion Wind”

Effective *magnetic* field that
torques spins along the dark matter wind

$$\boldsymbol{\tau} = g_{ae} \boldsymbol{\sigma}_e \times \nabla a \sim g_{ae} m_a a \boldsymbol{\sigma}_e \times \mathbf{v}_a$$



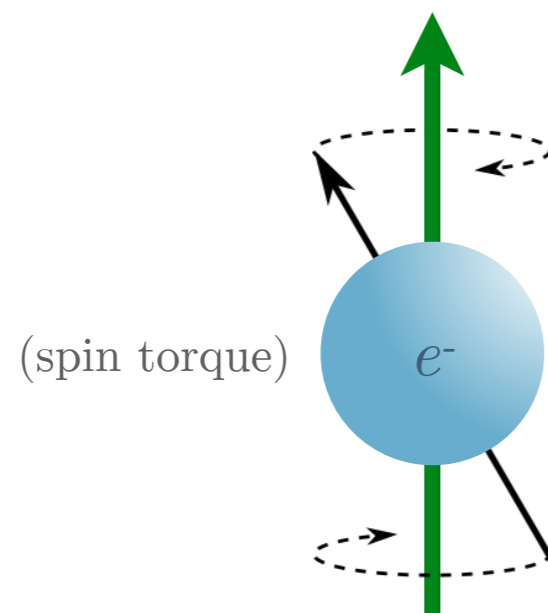
Electron-Coupling

$$\mathbf{B}_{\text{eff}} = -\frac{g_{ae}}{\mu_B} \nabla a \quad , \quad \mathbf{A}_{\text{eff}} = -\frac{g_{ae}}{e} \partial_t a \hat{\mathbf{s}}_e$$

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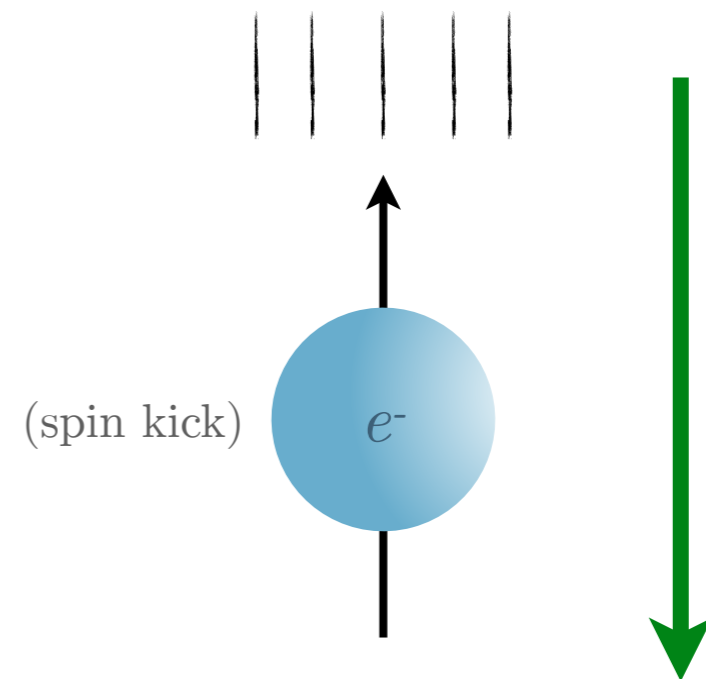
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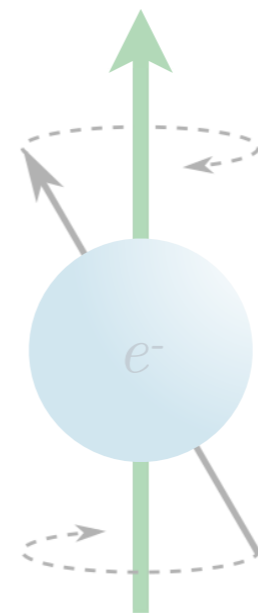
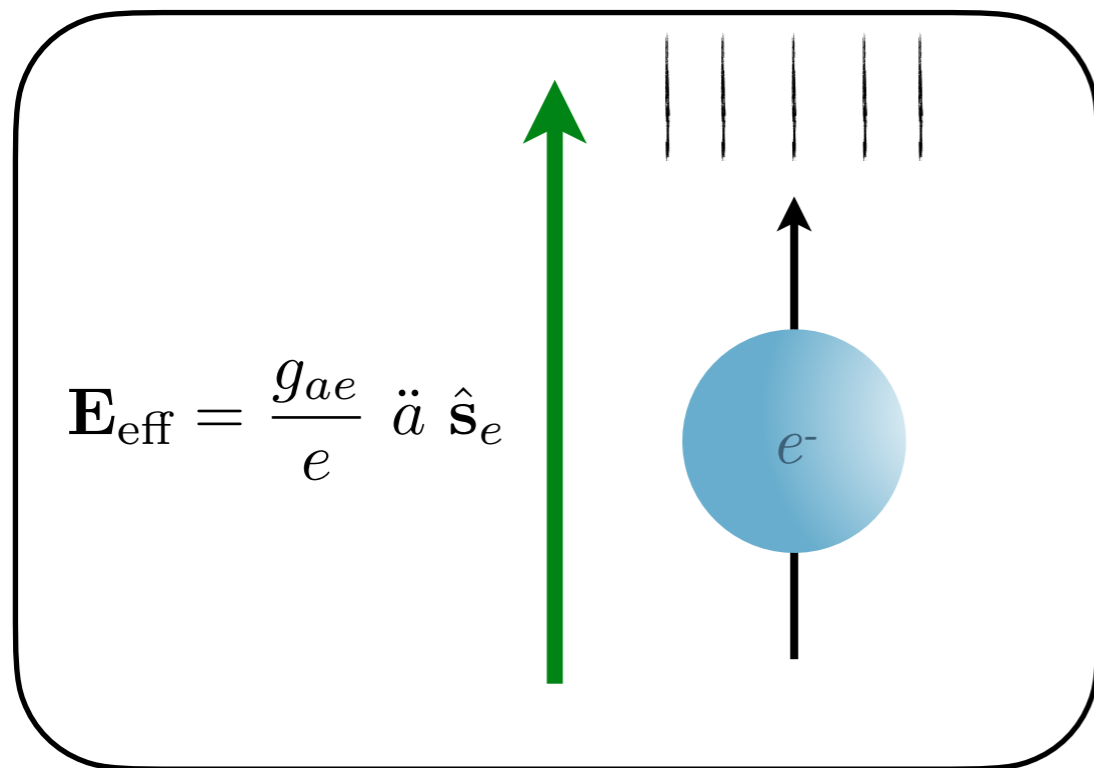


“Axioelectric”

Effective *electric* field that
pushes spins along the spin

$$\mathbf{F} = -g_{ae} \ddot{a} \hat{\mathbf{s}}_e \sim g_{ae} m_a^2 a \hat{\mathbf{s}}_e$$





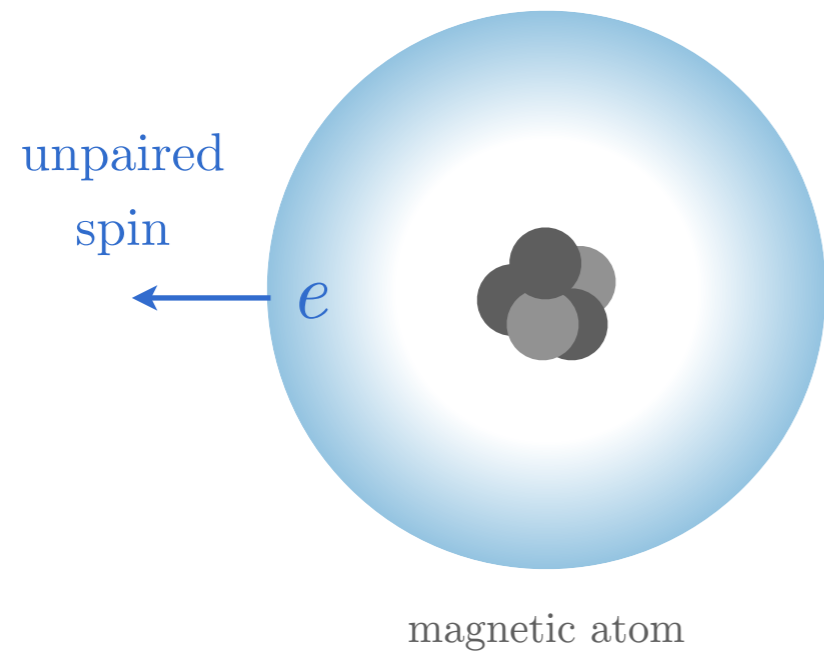
$$\mathbf{B}_{\text{eff}} = -\frac{g_{ae}}{\mu_B} \nabla a$$

II. Axioelectric Multilayer

Axioelectric

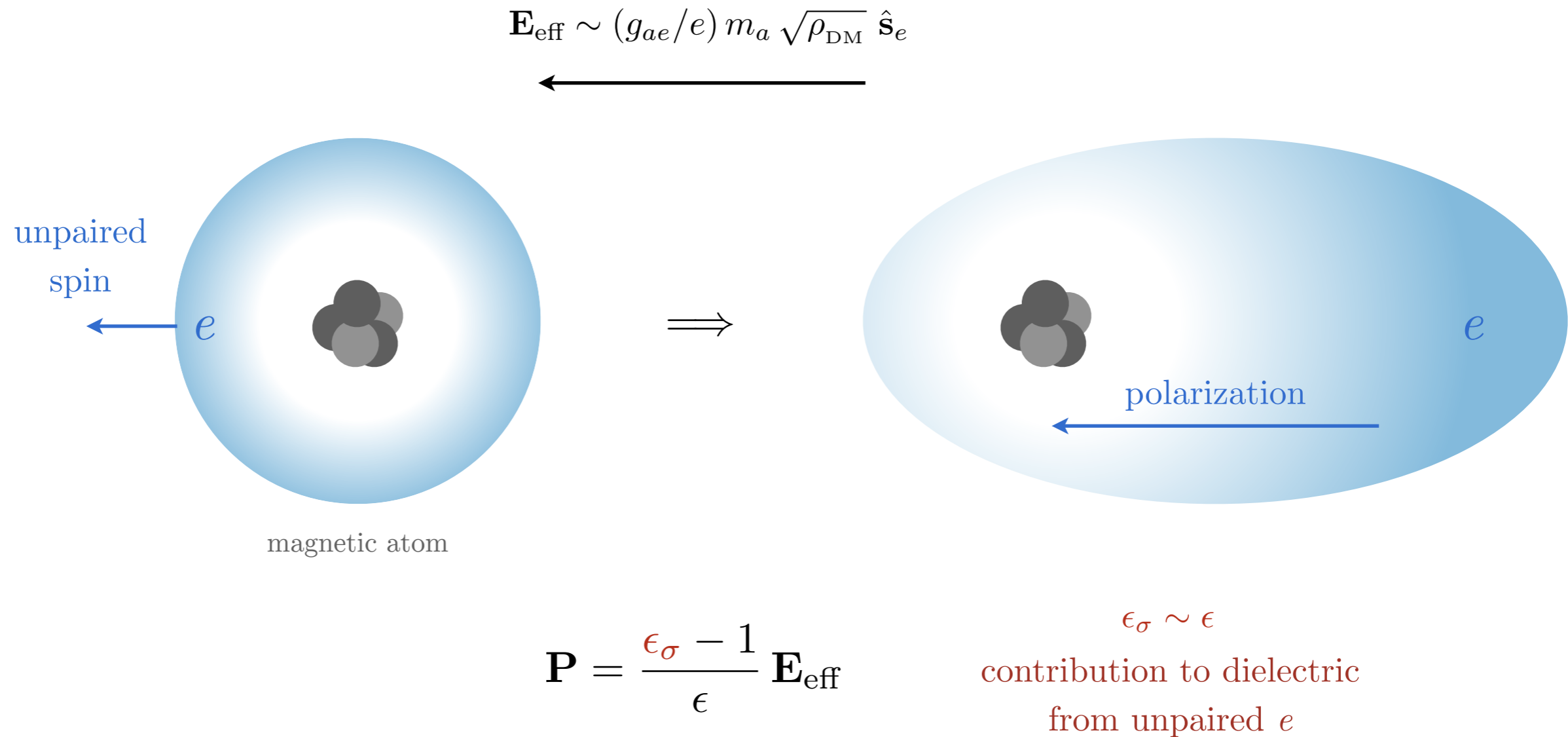
Axioelectric Effective Electric Field

$$\mathbf{E}_{\text{eff}} \sim (g_{ae}/e) m_a \sqrt{\rho_{\text{DM}}} \hat{\mathbf{s}}_e$$



Axioelectric

Axioelectric Effective Electric Field

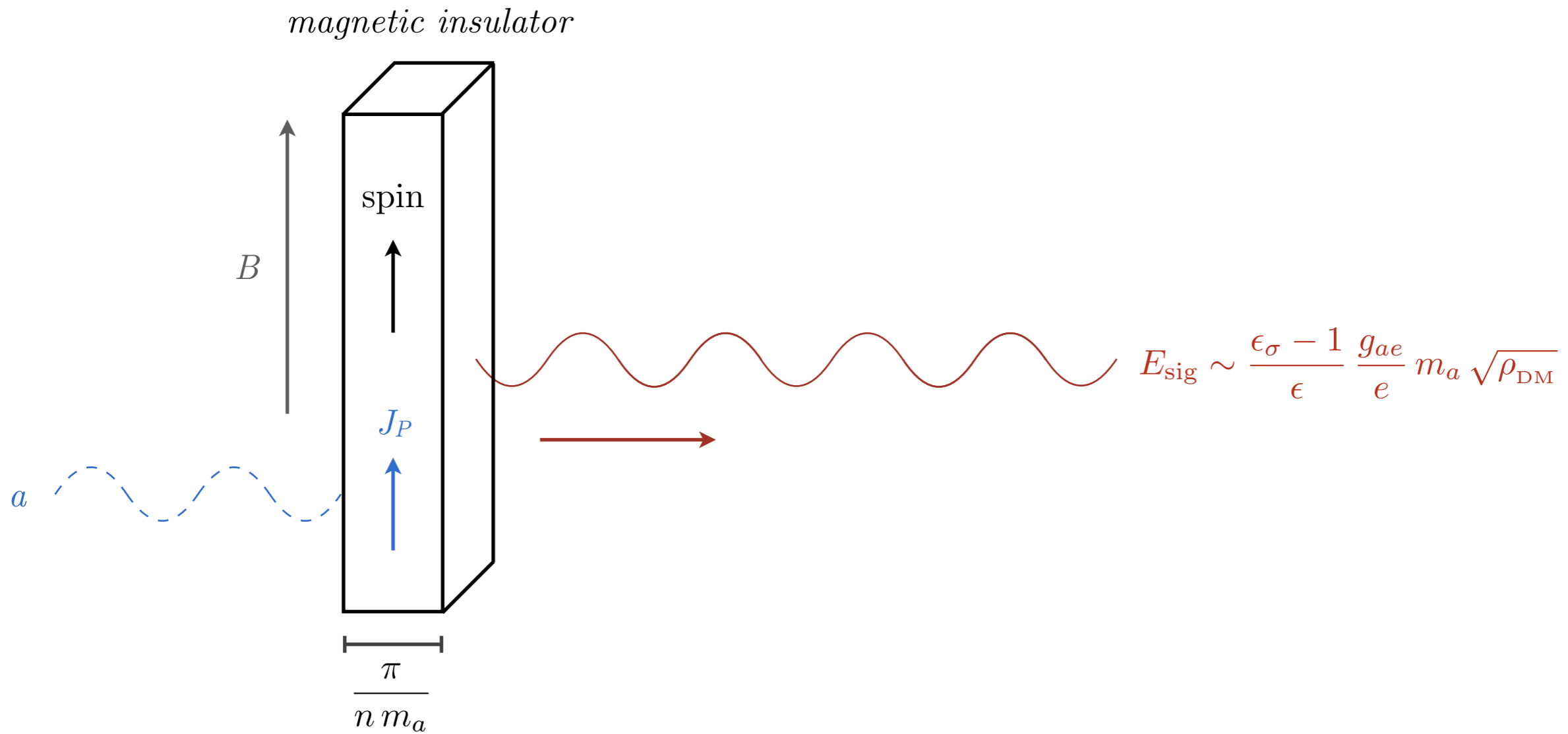


physical current “ J_P ” from the oscillating polarization \implies electromagnetic signal

$$\mathbf{J}_P = -\partial_t \mathbf{P}$$

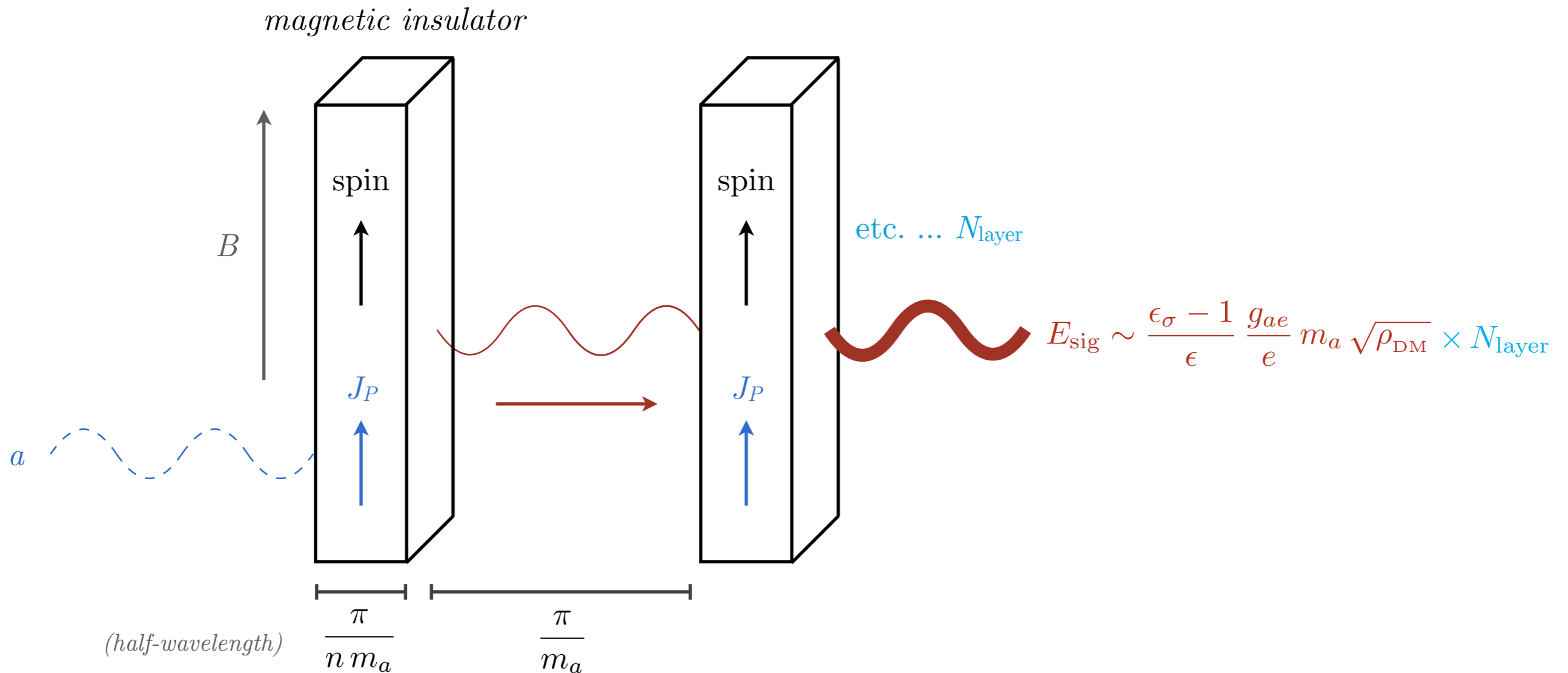
Axioelectric

Axioelectric Multilayer



Axioelectric

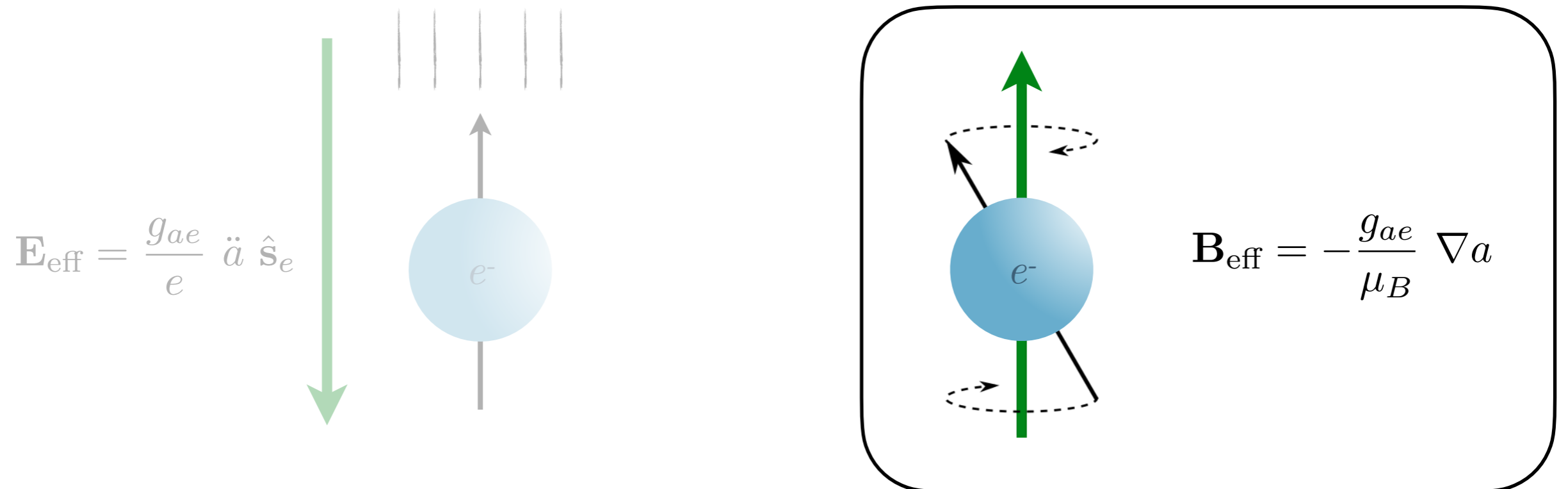
Axioelectric Multilayer



analogous to dielectric stacks MADMAX, MuDHI, and **LAMPOST**

$$g_{ae} \leftrightarrow g_{a\gamma\gamma} (e B / m_a^2)$$


(QCD axion sensitivity
at optical frequencies)

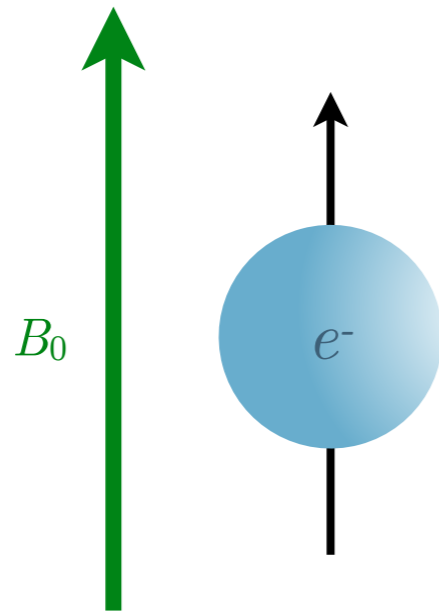


III. Axion Wind Multilayer

Axion Wind

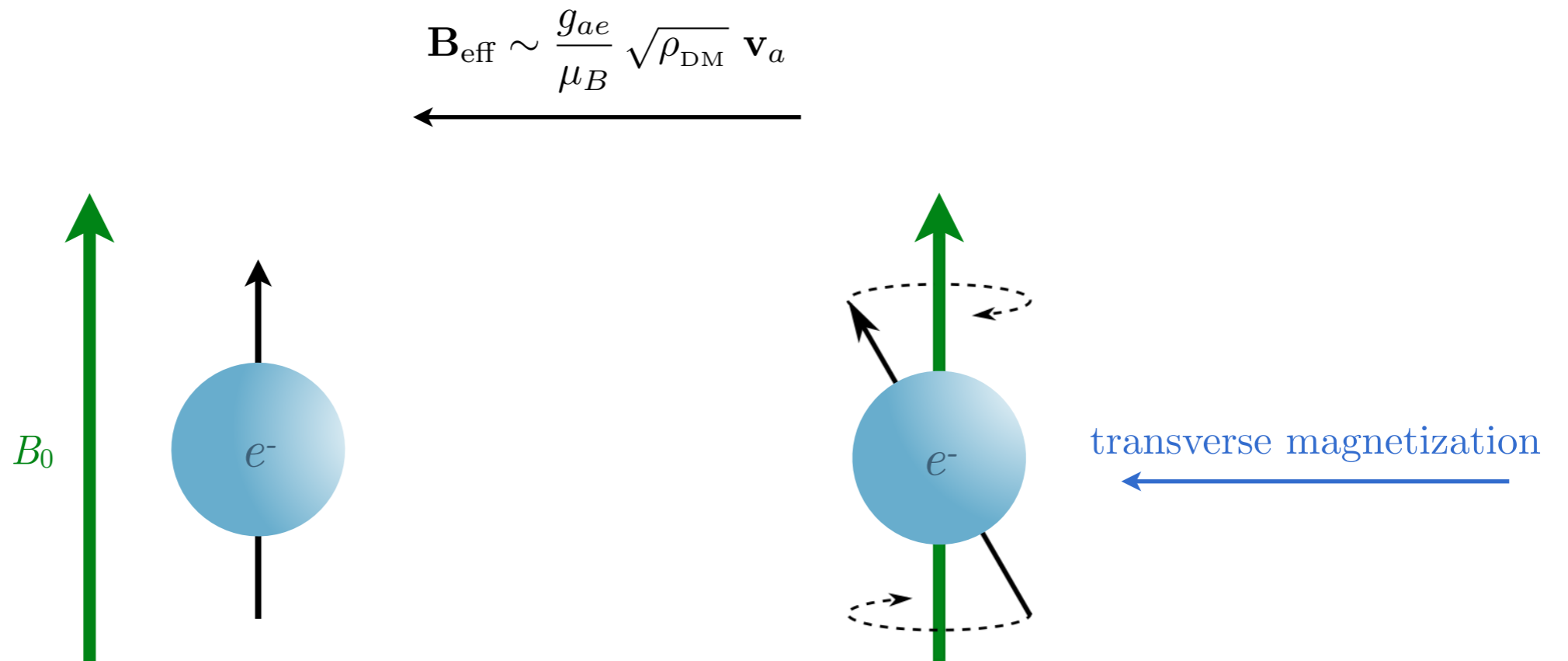
Axion Wind Effective Magnetic Field

$$\mathbf{B}_{\text{eff}} \sim \frac{g_{ae}}{\mu_B} \sqrt{\rho_{\text{DM}}} \mathbf{v}_a$$




Axion Wind

Axion Wind Effective Magnetic Field



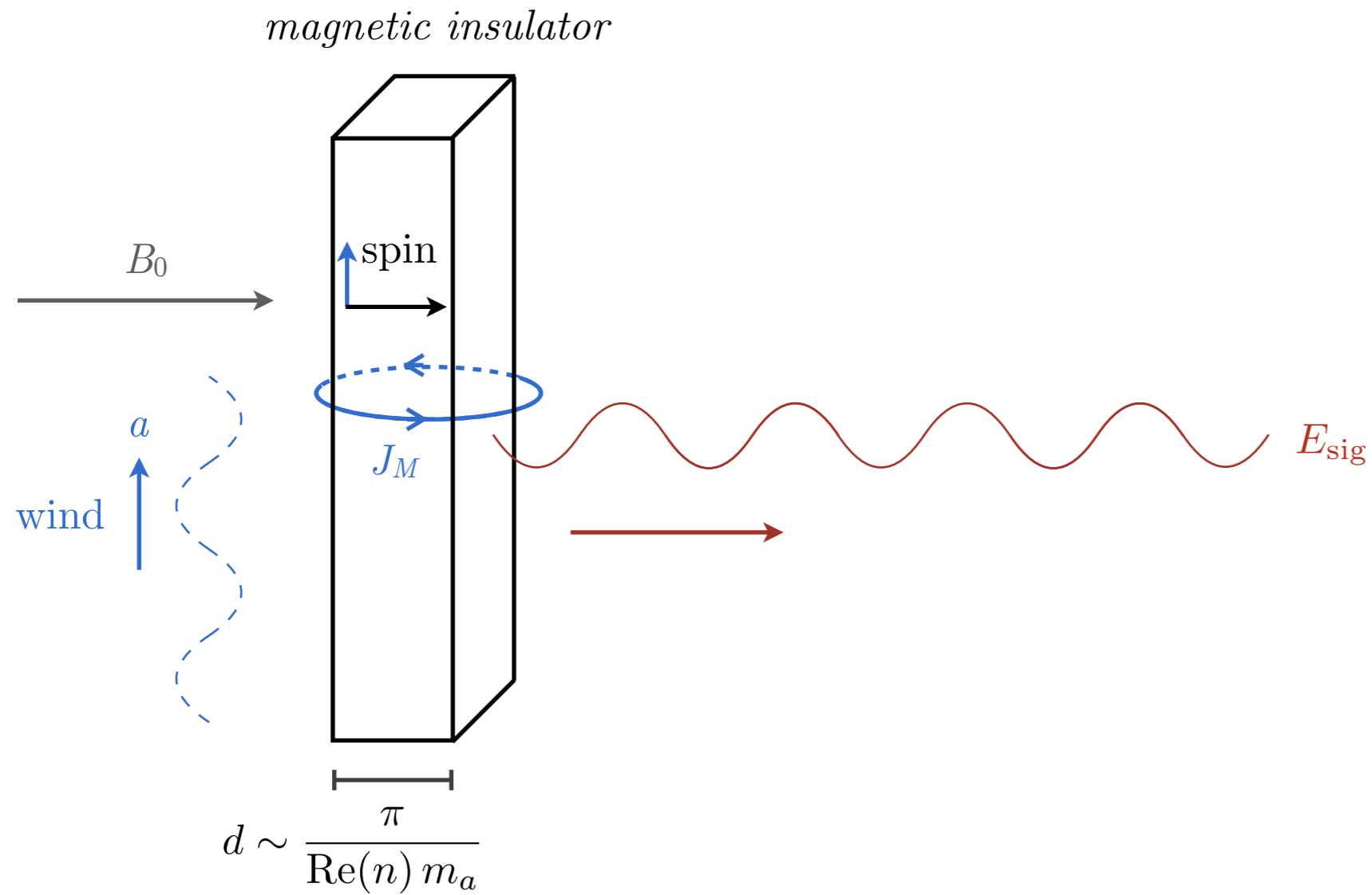
$$\mathbf{M}_{\perp} = (1 - \mu^{-1}) \mathbf{B}_{\text{eff}}$$

physical current “ J_M ” from the oscillating magnetization \implies electromagnetic signal

$$\mathbf{J}_M = \nabla \times \mathbf{M}$$

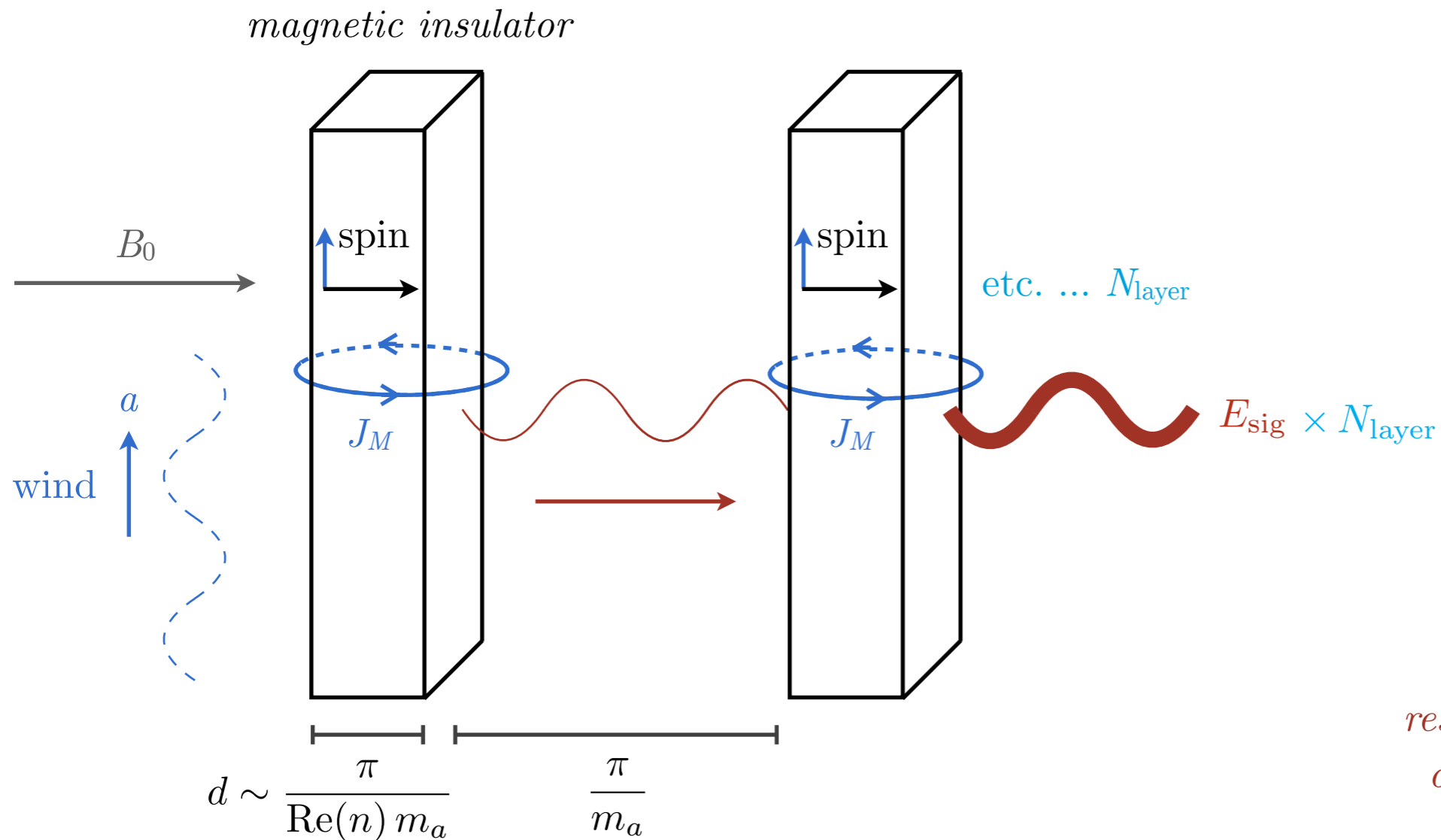
Axion Wind

Axion Wind Multilayer



Axion Wind

Axion Wind Multilayer

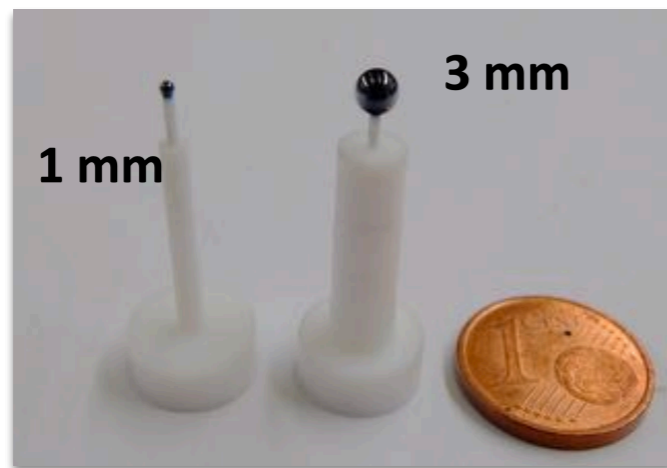


*response is strong function
of frequency and B-field
(e.g., FMR)*

$$\mu = \mu(m_a, B_0)$$

$$n = n(m_a, B_0)$$

$$E_{\text{sig}} \sim \frac{\mu - 1}{\mu - in \cot(n m_a d/2)} \frac{g_{ae}}{e} m_e v_a \sqrt{\rho_{\text{DM}}}$$

*Magnetic Insulators*YIG

Recent efforts have focused on small samples of high-purity YIG

$$Q_{\text{spin}} \sim 10^4$$

limited by small volume

(e.g., QUAX and others)

Ferrite ✓

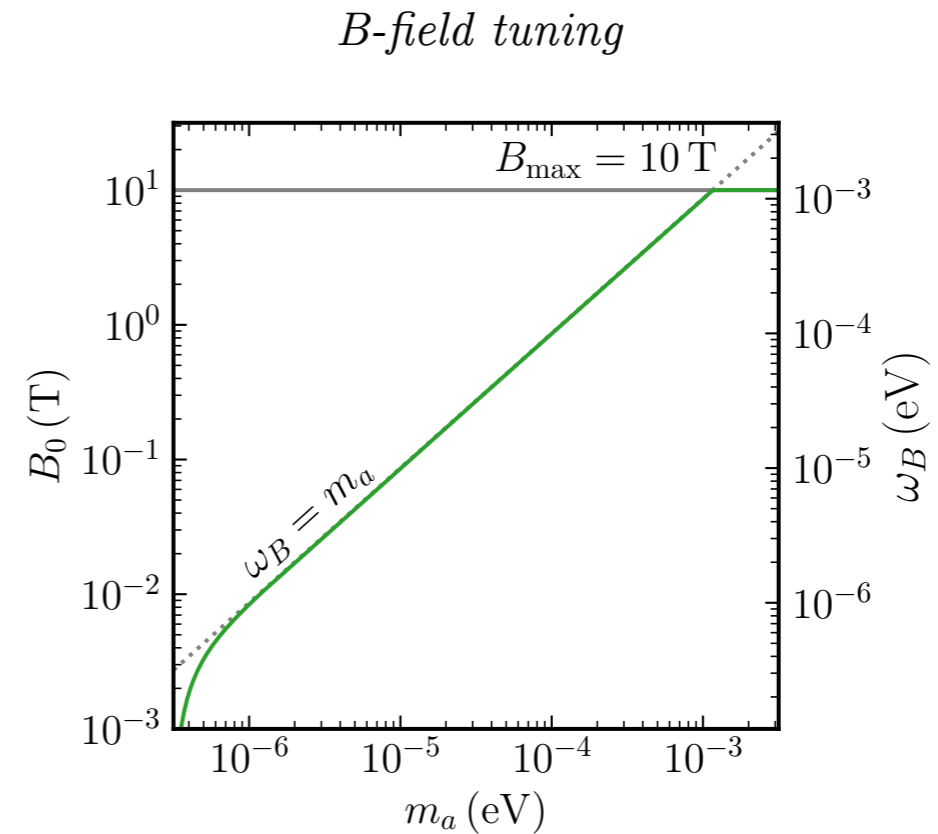
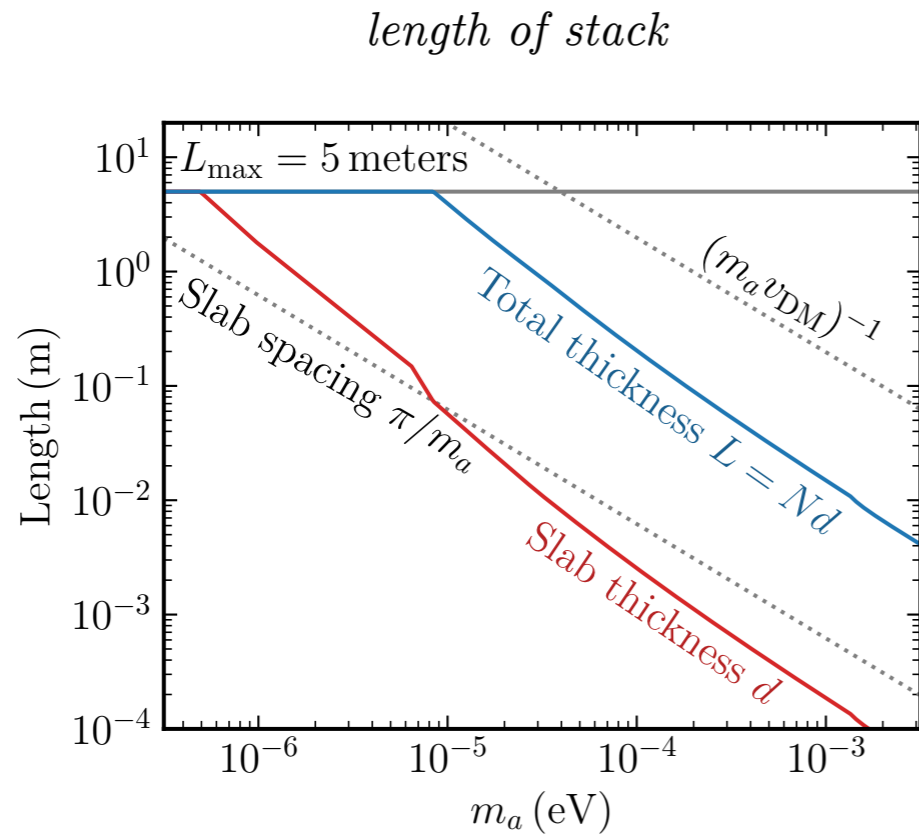
Cheaply mass produced

$$Q_{\text{spin}} \sim 10^2$$

enhanced by large volume

Axion Wind

Axion Wind Multilayer



$$\ell_{\text{screen}} \sim \frac{1}{m_a \text{Im}(n)}$$

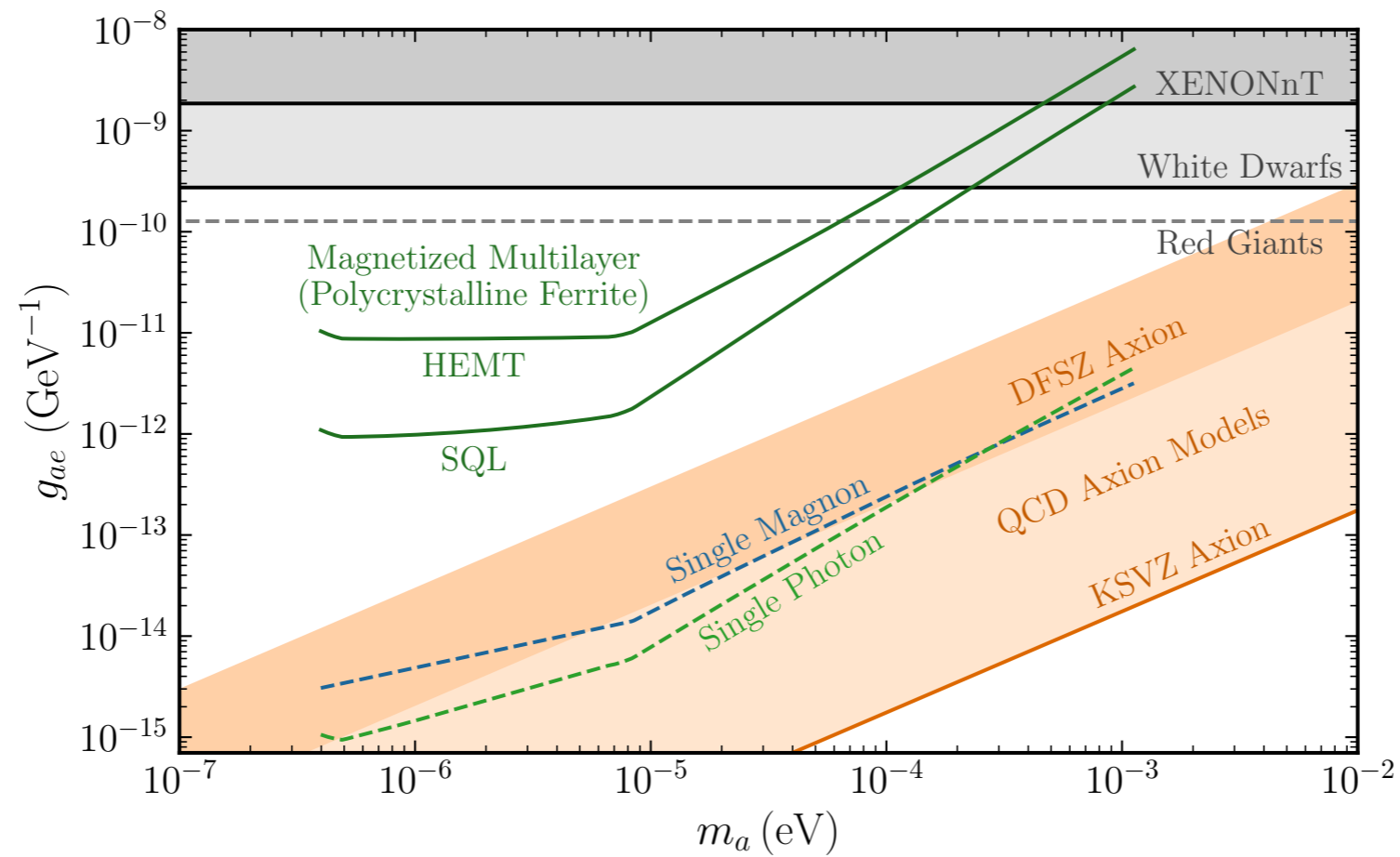
$$B_0 \sim \frac{m_a}{\mu_B} - \text{“small offset”} \quad (FMR)$$

$$m_a \gtrsim 10^{-6} \text{ eV} \quad \longleftrightarrow \quad \text{“MADMAX-like”} \quad \longleftrightarrow \quad m_a \lesssim 10^{-3} \text{ eV}$$

never requires B-field that is strong AND large

Axion Wind

Axion Wind Multilayer



large volume of
low- Q ferrite

axion absorption into magnon-polariton hybrid more likely than into magnon

Magnetic Stacks

$$L \sim g_{aff} \dot{a} \boldsymbol{\sigma} \cdot \mathbf{v}_f + g_{aff} \nabla a \cdot \boldsymbol{\sigma}$$

axioelectric *axion wind*

optical magnetic stack
(~ LAMPOST/MuDHI)

RF magnetic stack
(~ MADMAX)

Everything at leading order stems from these effects, expressed in different frames.

Magnetic Stacks

$$L \sim \underbrace{g_{aff} \dot{a} \boldsymbol{\sigma} \cdot \mathbf{v}_f}_{\text{axioelectric}} + \underbrace{g_{aff} \nabla a \cdot \boldsymbol{\sigma}}_{\text{axion wind}}$$

optical magnetic stack
(~ LAMPOST/MuDHI)

RF magnetic stack
(~ MADMAX)

Everything at leading order stems from these effects, expressed in different frames.

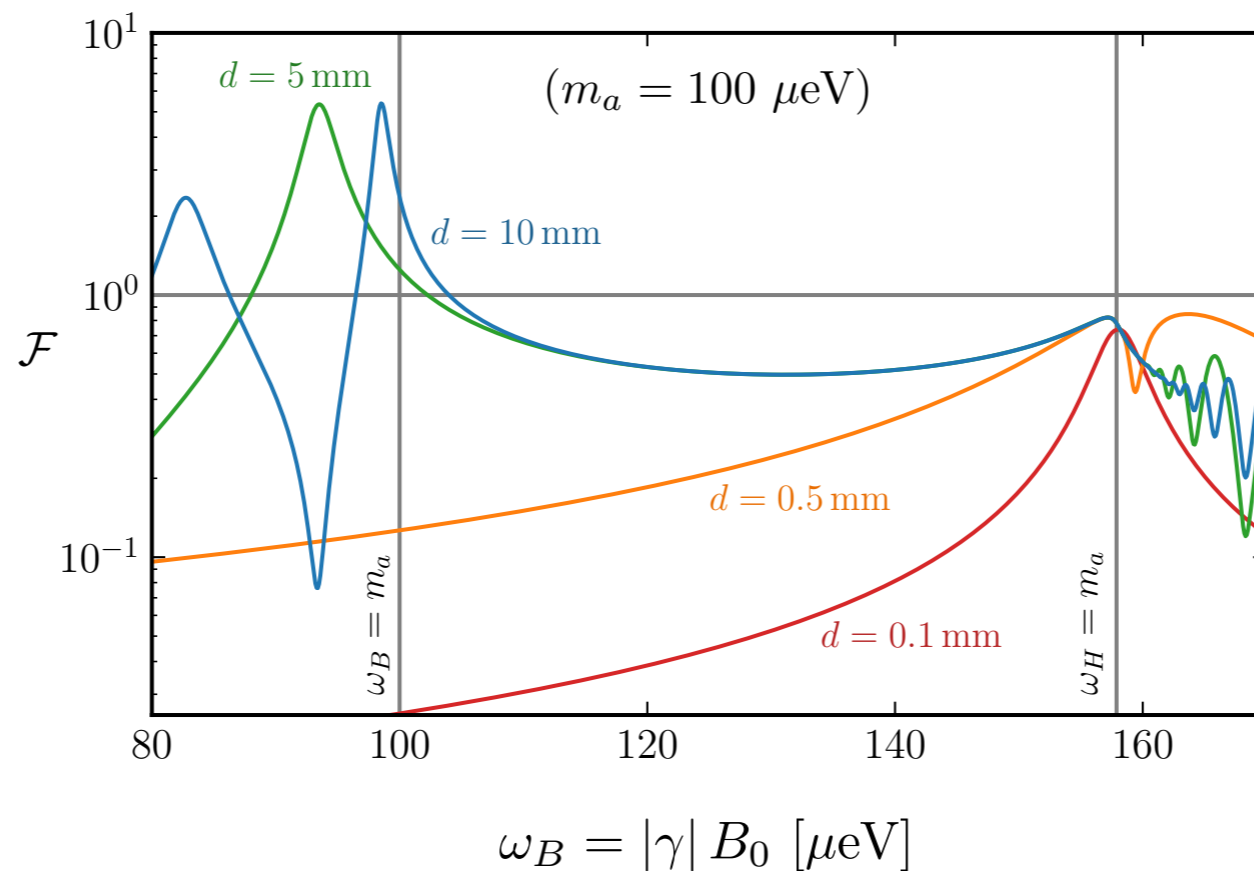
Other Nuggets

- Inclusive absorption rate $\sim \text{Im}[-1/\mu]$ (“magnetic energy loss function”)
- Bulk mechanical forces on spin-polarized sensors
- **No** leptonic EDMs proportional to the field value
- **No** axioelectronic energy shifts to leading order

Back Up Slides

Axion Wind

single layer optimization



$$\mathcal{F} = \left| \frac{\mu - 1}{\mu + in \cot(nm_a d/2)} \right|$$

$$\mu = \frac{m_a - \omega_B + i\omega_B/2Q}{m_a - \omega_H + i\omega_H/2Q}$$

- **Thin slab:** $\mu(\omega_H) \sim 2iQ(\omega_M/\omega_H)$, $\mathcal{F}_{\text{thin}}(\omega_H) \simeq Q_{\text{eff}} \omega_M d$, $Q_{\text{eff}}^{-1} \equiv Q^{-1} + Q_{\text{rad}}^{-1}$, $Q_{\text{rad}} = \frac{1}{\omega_M d}$
(spin response enhanced, but $\lambda \sim L_{\text{skin}}$ small \implies incoherent emission across slab \implies form factor < 1)

- **Thick slab:** $\mu(\omega_B) \sim (i/2Q)(\omega_B/\omega_M)$, $\mathcal{F}_{\text{thick}}(\omega_B) \simeq \frac{\sqrt{2Q\omega_M/(\omega_B\epsilon)}}{|\cot(n\omega_B d/2)|}$

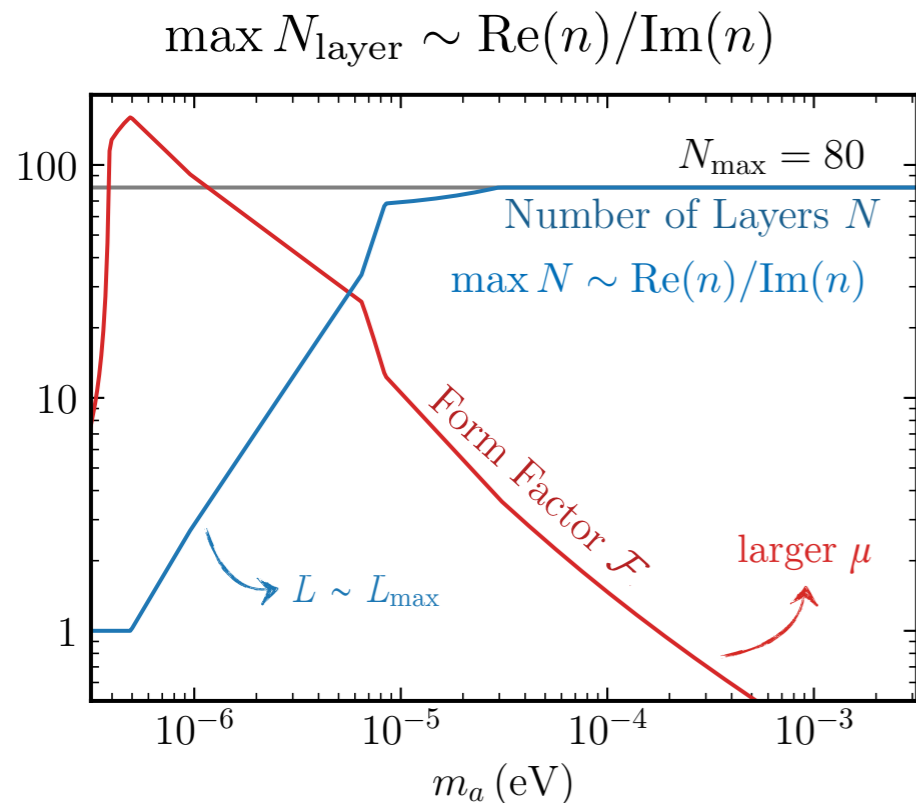
($\lambda \sim L_{\text{skin}}$ large \implies coherent emission across slab \implies form factor > 1)

$\omega_B < m_a \implies \mu$ small but real \implies slab \sim cavity with form factor peaking when $n m_a d \sim \pi$

Axion Wind

multi-layer optimization

$$\omega_B = m_a - \Delta\omega_B \quad (\Delta\omega_B \gg \omega_B/Q)$$



Parameters	Description	Variable	Value	
Material	Saturation magnetization	M_S	0.5 T	
	Magnetic quality factor	Q	10^2	
	Permittivity	ϵ	15	
Experimental	Slab area	A	1 m ²	
	e -fold scanning time	t_e	1 year	
	Maximum number of layers	N_{max}	80	
	Maximum total material thickness	L_{max}	5 m	
	Maximum applied B field	B_{max}	10 T	
Noise			HEMT	SQL
	Physical temperature	T	4 K	40 mK
	Amplifier noise temperature	T_{amp}	1 K $\left(\frac{m_a}{2\pi \times 4 \text{ GHz}}\right)$	m_a

sensitivity bandwidth: $N \frac{\Delta m_a}{m_a}, N \frac{\Delta \text{Re}(n)}{\text{Re}(n)} \lesssim 1 \implies \Delta m_a \sim \frac{\min(m_a, \Delta\omega_B)}{N} \sim \frac{\Delta\omega_B}{N}$

$$\Delta\omega_B \uparrow \implies \text{Re}(\mu) \uparrow + \text{Im}(\mu) \downarrow \implies \lambda \downarrow + L_{\text{skin}} \uparrow \implies N_{\text{layer}} \uparrow \quad (\text{compensated by smaller } \mathcal{F})$$