Neutrino NSI from Ultralight Scalars

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Scalar Neutrino Non-standard Interactions (SNSI)

[\[2401.02107\]](https://arxiv.org/pdf/2401.02107)

$$
\mathcal{L_S} = \bar{\nu} \left(i \gamma^\mu \partial_\mu - m_\nu \right) \nu - (y_\nu)_{\alpha \beta} \bar{\nu}_\alpha \nu_\beta \phi - y_f \bar{f} f \phi - \frac{1}{2} \left(\partial_\mu \phi \right)^2 - \frac{m_\phi^2}{2} \phi^2
$$

• We consider a light (< keV) to ultralight (<< eV) scalar ϕ

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- We consider a light (< keV) to ultralight (<< eV) scalar ϕ
- Phenomenology signatures:
	- \circ Scattering in detector: light recoils $\alpha \frac{y_f(y_\nu)_{\alpha\beta}}{q^2-m_\phi^2}$ $m_\phi << q \rightarrow \frac{y_f(y_\nu)_{\alpha\beta}}{q^2}$
	- \circ Oscillation: in the forward, coherent scattering in-medium: $\Rightarrow \frac{y_f(y_\nu)_{\alpha\beta}}{m^2}$

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- We consider a light (< keV) to ultralight (<< eV) scalar ϕ
- Phenomenology signatures:
	- o Scattering in detector: light recoils $\propto \frac{y_f(y_\nu)_{\alpha\beta}}{q^2-m_\phi^2}$ $m_\phi << q \rightarrow \frac{y_f(y_\nu)_{\alpha\beta}}{q^2}$
	- \circ Oscillation: in the forward, coherent scattering in-medium: $\frac{y_f(y_\nu)_{\alpha\beta}}{m^2}$
- matter effect in **DUNE long-baseline beam neutrino oscillations**

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Scalar or Vector NSI?

$$
\mathcal{L}^6_{\mathcal{S}} \supset \frac{(y_\nu)_{\alpha\beta} y_f}{m_\phi^2} (\bar{f}f)(\bar{\nu}_\alpha \nu_\beta)
$$

See e.g. Ge, Parke 1812.08376 Smirnov, Xu 1909.07505 5

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What about the long-ranged scalar potential?

Consider long-ranged ϕ potential from matter sources inducing a neutrino mass correction:

$$
\Rightarrow \frac{n_f y_{\alpha\beta} y_f}{m_{\phi}^2} \quad \text{for} \quad m_f \text{L}_{\text{max}}(R_{\theta} e_{\gamma} + 1)
$$

Recovers the short-range description

Cutoff for DUNE beam depth: $m_{\phi} \sim 10^{-14}$ eV

Wise, Zhang 1803.00591 Smirnov, Xu 1909.07505 Babu, Chauhan, Dev 1912.13488 Agarwalla, Bustamante, Singh, Swain 2404.02775

How to parameterize the physics?

Flavor Basis: Normalize to one of the mass splittings

$$
\delta \mathbb{M}_{\alpha\beta} \equiv \sqrt{|\Delta m^2_{31}|} \begin{pmatrix} \eta_{ee} & \eta_{e\mu} e^{i\phi_{e\mu}} & \eta_{e\tau} e^{i\phi_{e\tau}} \\ \eta_{e\mu} e^{-i\phi_{e\mu}} & \eta_{\mu\mu} & \eta_{\mu\tau} e^{i\phi_{\mu\tau}} \\ \eta_{e\tau} e^{-i\phi_{e\tau}} & \eta_{\mu\tau} e^{-i\phi_{\mu\tau}} & \eta_{\tau\tau} \end{pmatrix} \quad \eta_{\alpha\beta} = \frac{n_f y_f y_{\alpha\beta}}{\sqrt{|\Delta m^2_{31}|} m_\phi^2}
$$

$$
H\supset{\mathbb{M}}^{\dagger}\cdot\delta{\mathbb{M}}\supset m_1\times\eta\hspace{0.4cm}\text{[modulo PMNS elements]}
$$

- **Depends on choice for** *m***₁**
- One choice: fixing Δm_{21}^2 and Δm_{31}^2 to measured values, and specify a choice of m_i to constrain η 's

Light SNSI Oscillations at DUNE Far Detector and δ_{cp} Extraction

[2401.02107]

- We consider $L \sim 1300$ km baseline oscillations from the DUNE beam neutrinos
- 4 component analysis: electron (anti)neutrino appearance, muon (anti)neutrino disappearance
- 5% flux normalization uncertainty

See also: Moon Moon's talk, [2309.12249], [2210.00109] **9** and the same of the

Resolving the CP Phase at DUNE

- Test for $\delta_{\rm CP}$ at DUNE with scalar NSI (SNSI) in the long-baseline oscillations
- We marginalize over the all SNSI (η magnitudes and phases)
- We expect degeneracies to show up in the measurement just like vector NSI

Dependence on the lightest neutrino mass

- Generically get the best upper limit (U.L.) when $m_{_1}$ is larger
	- Enhances the $M₀$ δM term
- To evaluate the best possible sensitivity, take the largest $m₁$ value

Other Constraints on Non-standard Oscillations (light scalars) [1912.13488] Babu, Chauhan, Dev

In Medium Mass: Our Sun

Long-ranged scalar potential induces in -medium mass correction \sim γ $\langle \phi \rangle$ in the **sun**

In Medium Mass: Supernova

Likewise, the long ranged potential in supernova

Additional constraints from, e.g.

- BBN
- \triangle Neff and CMB
- Scattering experiments (high masses)

Long-ranged mass correction should not exceed solar data : $\Delta m_{Sun}^{}$ = $n^{\rm Sun}_{f}$ $f y_f y_v / m_{\phi}^2$ < 7.4 meV

Long-ranged mass correction should not modify neutrino free streaming :

 $\Delta m_{_{SN}}$ = $n^{\rm SN}$ $_{f}$. $f y_f y_v / m_\phi^2 < T_v \sim 5$ MeV

Seeing the full picture: translating to neutrino/electron Yukawas

- Marginalize over all η 's and δ_{CD}
- Take the best constraint on $\eta_{\alpha\beta}$
(and the associated Yukawa) after marginalizing all others

…and the neutrino/nucleon Yukawas

- Marginalize over all η ⁵s and δ_{CP}
- \bullet Take the best constraint on $\eta_{\alpha\beta}$
(and the associated Yukawa) after marginalizing all others

Key Points

- Light-mediator scalar NSI modify the effective mass matrix in the propagation Hamiltonian
- Degeneracies with the absolute neutrino masses
- Degeneracies with oscillation parameters, e.g. δ_{CP}
- However DUNE appears to only be sensitive to SNSI already ruled out by more powerful long-ranged effects : $\Delta m_{\rm SN}$ and $\Delta m_{\rm Sun}$

Backup Deck

Scalar or Vector NSI?

$$
\bar{\nu}_{\alpha} \left[\gamma^{0} (i\partial_{0} - V_{\alpha\beta} - \frac{n_{f} g_{f} g_{\alpha\beta}}{m_{V}^{2}}) - i\gamma \cdot \nabla - M_{\alpha\beta} \right] \nu_{\beta} = 0
$$
\n**Propagation Hamiltonian**

Realized as a **potential energy shift**

$$
H_{\alpha\beta} = \frac{1}{2E_{\nu}} \mathbb{M}^{\dagger}_{\alpha\beta} \mathbb{M}_{\alpha\beta} + (V_{\text{CC}} + V_{\alpha\beta}^{\text{NSI}})
$$

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Many-parameters vs. one-at-a-time

All floating, 1-D marginals

One-at-a-time

Multi-dimension al parameter scan

MultiNest

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Multi-dimension al parameter scan

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$$

$$
\delta \mathbb{M}_{ij} = U^{\dagger} \delta \mathbb{M}_{\alpha\beta} U = \begin{pmatrix} \epsilon_{11} & \epsilon_{12} & \epsilon_{13} \\ \epsilon_{12}^* & \epsilon_{22} & \epsilon_{23} \\ \epsilon_{13}^* & \epsilon_{23}^* & \epsilon_{33} \end{pmatrix} \qquad \begin{array}{c} \epsilon_{11} \to -\epsilon_{11} - 2m_1 \\ \epsilon_{22} \to -\epsilon_{22} - 2m_2 \\ \epsilon_{33} \to -\epsilon_{33} - 2m_3 \\ \epsilon_{ij} \to -\epsilon_{ij}; \quad i \neq j \end{array}
$$

Physical Degrees of Freedom

$$
\mathbb{M}^2_{\text{eff}} \equiv \begin{pmatrix} m_1^2 + \mu_{11} & \mu_{12} & \mu_{13} \\ \mu_{12}^* & m_2^2 + \mu_{22} & \mu_{23} \\ \mu_{13}^* & \mu_{23}^* & m_3^2 + \mu_{33} \end{pmatrix}
$$

$$
\rightarrow \begin{pmatrix} 0 & \mu_{12} & \mu_{13} \\ \mu_{12}^* & \Delta m_{12}^2 - \mu_{11} + \mu_{22} & \mu_{23} \\ \mu_{13}^* & \mu_{23}^* & \Delta m_{13}^2 - \mu_{11} + \mu_{33} \end{pmatrix}
$$

- **● Depends on choice for** *m1*
- One choice: fixing Δm_{21}^2 and Δm_{31}^2 to measured values, and specify a choice of m_i to constrain η 's

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$$

[9-1=8 physical oscillation parameters; subtract constant times identity matrix]

$$
\mu_{11} = 2m_1\epsilon_{11} + |\epsilon_{11}|^2 + |\epsilon_{12}|^2 + |\epsilon_{13}|^2,
$$

\n
$$
\mu_{22} = 2m_2\epsilon_{22} + |\epsilon_{22}|^2 + |\epsilon_{12}|^2 + |\epsilon_{23}|^2,
$$

\n
$$
\mu_{33} = 2m_3\epsilon_{33} + |\epsilon_{33}|^2 + |\epsilon_{13}|^2 + |\epsilon_{23}|^2,
$$

\n
$$
\mu_{12} = (m_1 + m_2 + \epsilon_{11} + \epsilon_{22})\epsilon_{12} + \epsilon_{13}\epsilon_{23}^*,
$$

\n
$$
\mu_{13} = (m_1 + m_3 + \epsilon_{11} + \epsilon_{33})\epsilon_{13} + \epsilon_{12}\epsilon_{23}^*,
$$

\n
$$
\mu_{23} = (m_2 + m_3 + \epsilon_{22} + \epsilon_{33})\epsilon_{23} + \epsilon_{13}\epsilon_{12}^*.
$$

See also: Denton, Giarnetti, Meloni [2210.00109] 22

DUNE Spectra

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