

T violation at a future neutrino factory

R. Kitano¹, J. Sato², S. Sugama²
¹KEK, ²Yokohama National Univ.

Ref. arXiv: 2407.05807 [hep-ph]

16-21 September 2024
NuFact 2024 @ Argonne National Laboratory

Contents

1. Introduction
2. CP and T violation in neutrino oscillations
3. Result
4. Summary

Motivation

- Testing T violation in lepton sector has not been achieved
- CP and T violation measurements as a non-trivial check of the CPT theorem in QFT

Introduction

Framework of standard 3 flavor ν oscillation in matter

$$i \frac{d}{dt} \begin{pmatrix} \nu_e(\bar{\nu}_e) \\ \nu_\mu(\bar{\nu}_\mu) \\ \nu_\tau(\bar{\nu}_\tau) \end{pmatrix} = [U \text{diag}(0, \Delta E_{21}, \Delta E_{31}) U^\dagger + \text{diag}(\pm A, 0, 0)] \begin{pmatrix} \nu_e(\bar{\nu}_e) \\ \nu_\mu(\bar{\nu}_\mu) \\ \nu_\tau(\bar{\nu}_\tau) \end{pmatrix}$$

+ for ν , - for $\bar{\nu}$

Diagonalized by \tilde{U}
and eigen values \tilde{E}_j

$$= \tilde{U}^{(\pm)} \text{diag}(\tilde{E}_1^{(\pm)}, \tilde{E}_2^{(\pm)}, \tilde{E}_3^{(\pm)}) \tilde{U}^{(\pm)\dagger} \begin{pmatrix} \nu_e(\bar{\nu}_e) \\ \nu_\mu(\bar{\nu}_\mu) \\ \nu_\tau(\bar{\nu}_\tau) \end{pmatrix}$$

$$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_{23} & s_{23} \\ 0 & -s_{23} & c_{23} \end{pmatrix} \begin{pmatrix} c_{13} & 0 & s_{13} e^{-i\delta} \\ 0 & 1 & 0 \\ -s_{13} e^{i\delta} & 0 & c_{13} \end{pmatrix} \begin{pmatrix} c_{12} & s_{12} & 0 \\ -s_{12} & c_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix}$$

CP phase

Introduction

CP violation : if $\delta = 0$, this difference does not zero
→ **Matter effect may mimic CP violation**

$$\begin{aligned} & P(\nu_\alpha \rightarrow \nu_\beta) - P(\bar{\nu}_\alpha \rightarrow \bar{\nu}_\beta) \\ &= -4 \sum_{j>k} \left\{ \text{Re} \left[\tilde{U}^{(+)}_{\beta j} \tilde{U}^{(+)*}_{\alpha j} \tilde{U}^{(+)*}_{\beta k} \tilde{U}^{(+)}_{\alpha k} \right] \sin^2 \left(\frac{\Delta \tilde{E}_{jk}^{(+)} L}{2} \right) \right. \\ &\quad \left. - \text{Re} \left[\tilde{U}^{(-)}_{\beta j} \tilde{U}^{(-)*}_{\alpha j} \tilde{U}^{(-)*}_{\beta k} \tilde{U}^{(-)}_{\alpha k} \right] \sin^2 \left(\frac{\Delta \tilde{E}_{jk}^{(-)} L}{2} \right) \right\} \\ &\quad + 2 \sum_{j>k} \left\{ \text{Im} \left[\tilde{U}^{(+)}_{\beta j} \tilde{U}^{(+)*}_{\alpha j} \tilde{U}^{(+)*}_{\beta k} \tilde{U}^{(+)}_{\alpha k} \right] \sin \left(\Delta \tilde{E}_{jk}^{(+)} L \right) \right. \\ &\quad \left. - \text{Im} \left[\tilde{U}^{(-)}_{\beta j} \tilde{U}^{(-)*}_{\alpha j} \tilde{U}^{(-)*}_{\beta k} \tilde{U}^{(-)}_{\alpha k} \right] \sin \left(\Delta \tilde{E}_{jk}^{(-)} L \right) \right\} \end{aligned}$$

Introduction

T violation : if $\delta = 0$, this difference is exactly zero

→ **Matter effect is not important**

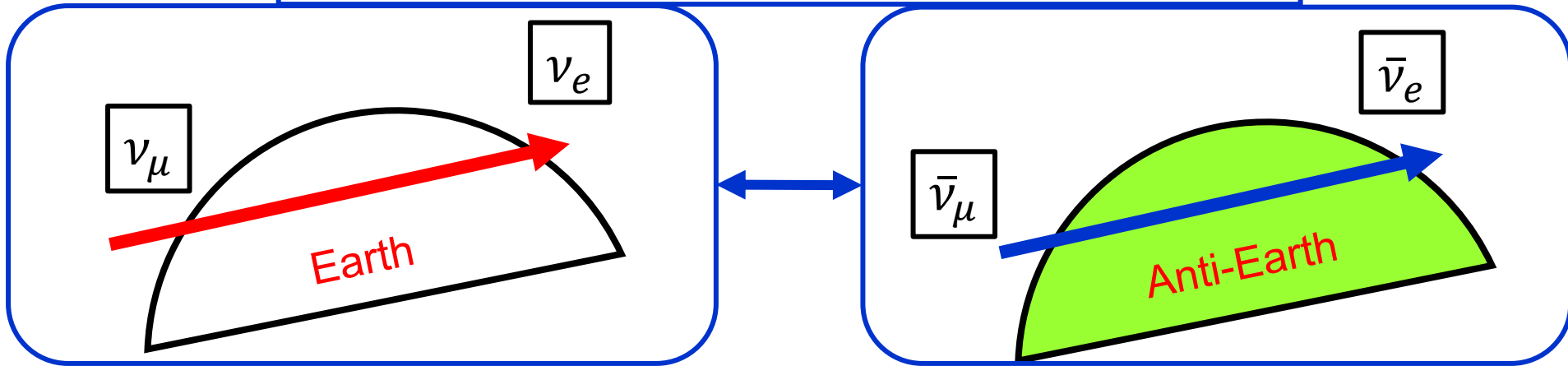
→ **Pure CP (T) violating effect**

$$\begin{aligned} & P(\nu_\alpha \rightarrow \nu_\beta) - P(\nu_\beta \rightarrow \nu_\alpha) \\ &= -16\text{Im} \left[\tilde{U}^{(+)}_{\beta 2} \tilde{U}^{(+)*}_{\alpha 2} \tilde{U}^{(+)*}_{\beta 1} \tilde{U}^{(+)}_{\alpha 1} \right] \\ &\times \sin^2 \left(\frac{\Delta \tilde{E}_{31}^{(+)} L}{2} \right) \sin^2 \left(\frac{\Delta \tilde{E}_{32}^{(+)} L}{2} \right) \sin^2 \left(\frac{\Delta \tilde{E}_{21}^{(+)} L}{2} \right) \end{aligned}$$

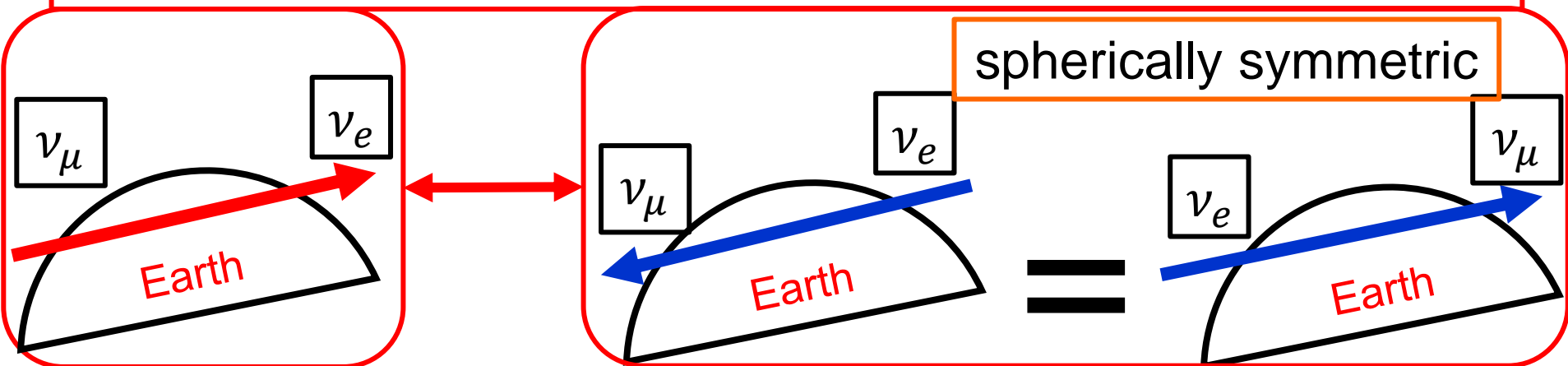
$\propto \sin \delta$

CP violation and T violation in matter

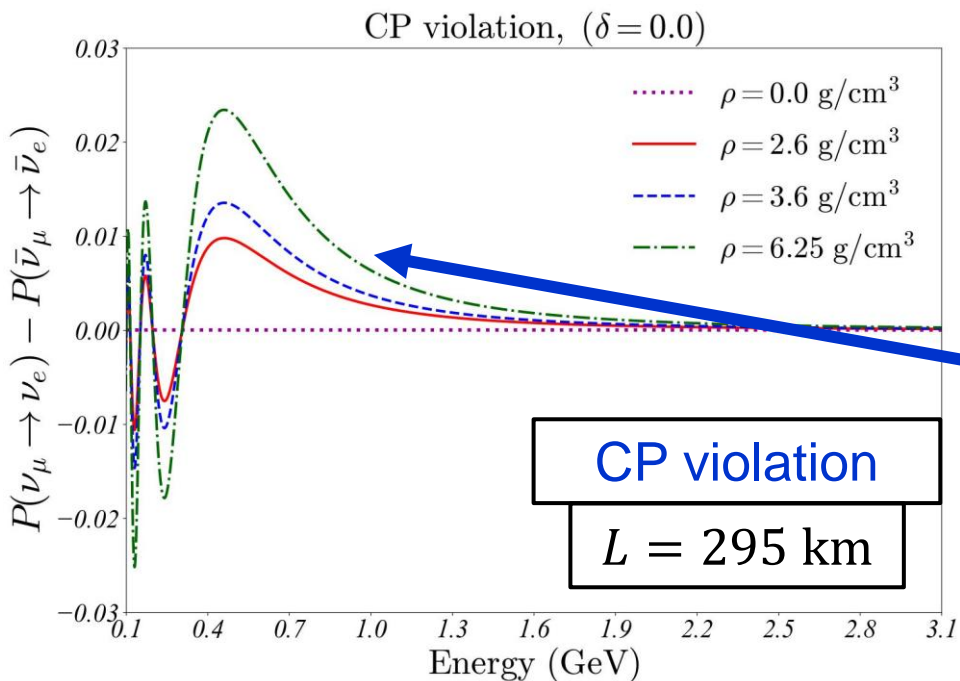
Ideal CP-test : need the anti-Earth



Ideal T-test : need ν beam from opposite direction



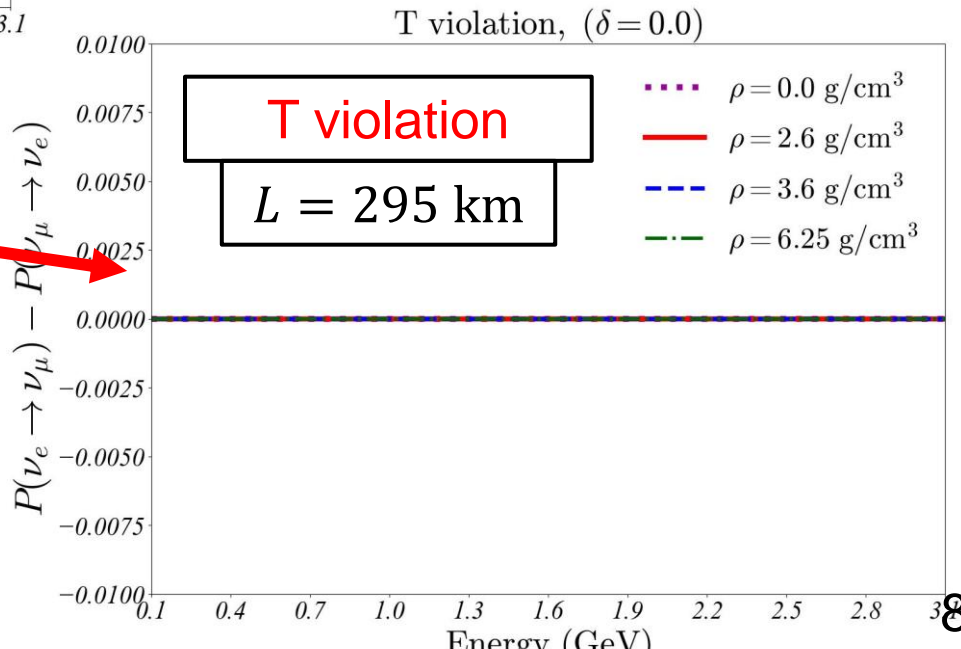
CP violation and T violation in matter



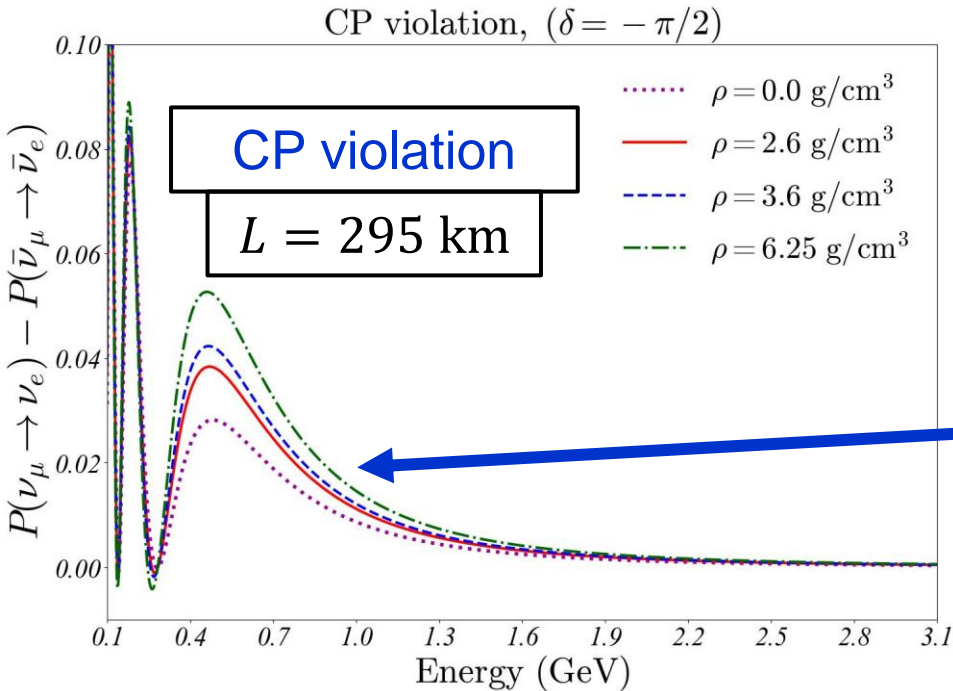
Assumption
 $\delta = 0^\circ$
(No CP violation)
 Normal Ordering

Fake CP violation exists
 As long as the matter density is non-zero, i.e. not in a vacuum.

No fake CP (T) violation,
 regardless of matter effects.



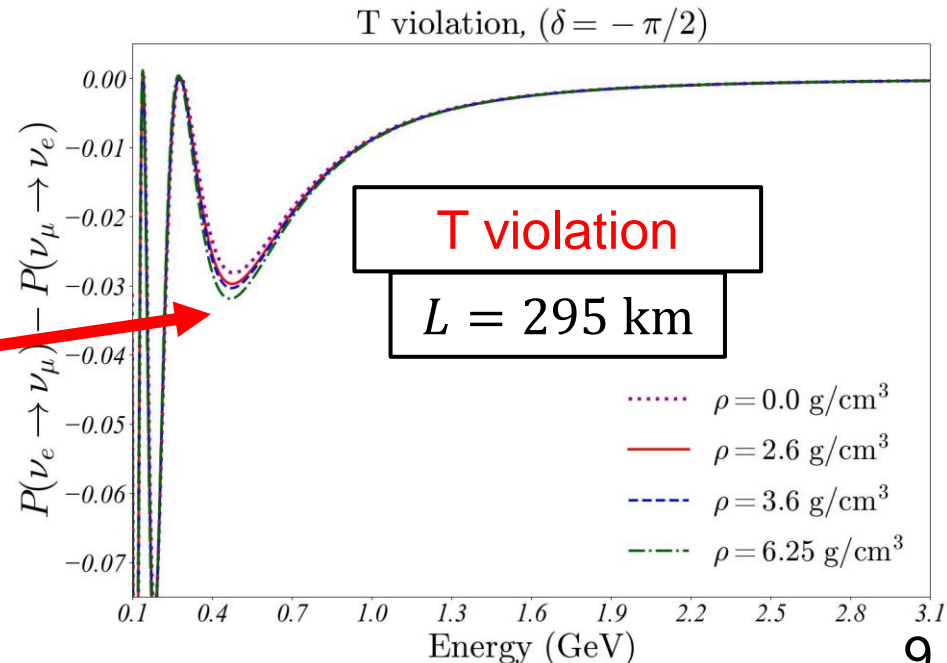
CP violation and T violation in matter



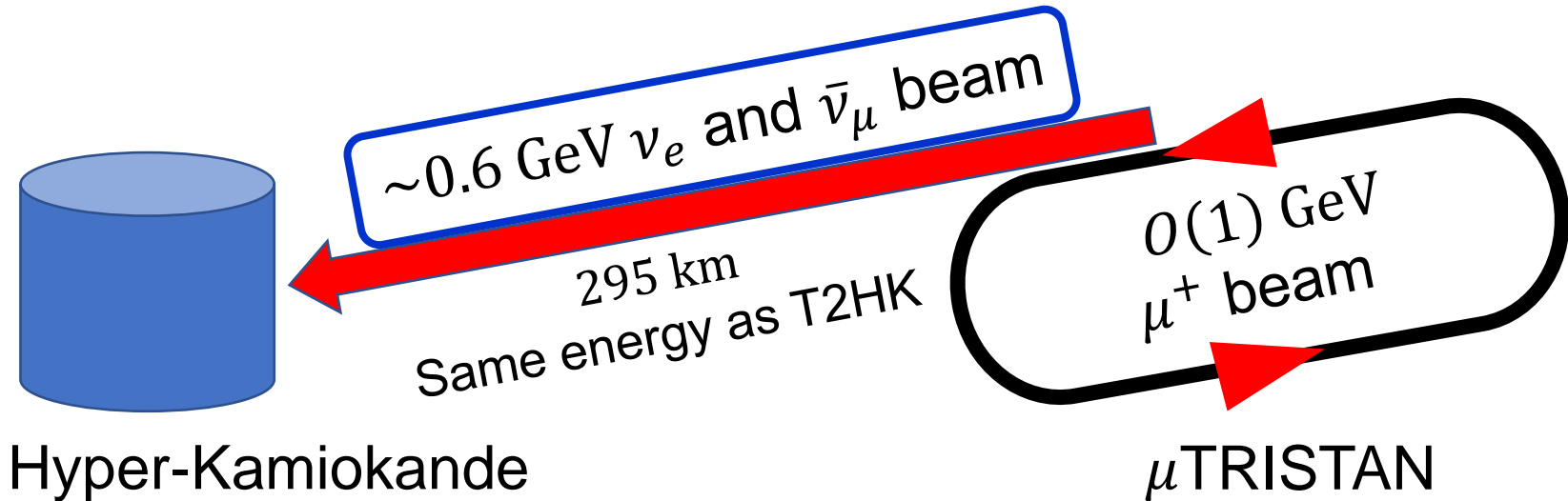
Assumption
 $\delta = 270^\circ$
(Maximally CP violation)
Normal Ordering

Dependence on matter density is **large**

Dependence on matter density is **tiny**



Future neutrino factory



$P(\nu_e \rightarrow \nu_\mu)$ from $\mu\text{TRISTAN}$

$P(\nu_\mu \rightarrow \nu_e)$ from T2HK

combine

T violation $P(\nu_e \rightarrow \nu_\mu) - P(\nu_\mu \rightarrow \nu_e)$

Robust test of CP phase free from matter effects!

Probability of $\nu_e \rightarrow \nu_\mu$

At the HK, in principle, ν_μ and $\bar{\nu}_\mu$ are distinguished by neutron tagging method.

We can define the oscillation probability $P(\nu_e \rightarrow \nu_\mu)$ as

Perfect charge identification

$$P(\nu_e \rightarrow \nu_\mu) = \frac{N_{\text{far}}^{\nu_e \rightarrow \nu_\mu}}{N_{\text{near}}^{\nu_e \rightarrow \nu_e}}$$

No charge identification

$$P(\nu_e \rightarrow \nu_\mu) = \frac{\left(N_{\text{far}}^{\nu_e \rightarrow \nu_\mu} + N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} \right) - N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} \Big|_{\text{T2HK}}}{N_{\text{near}}^{\nu_e \rightarrow \nu_e}}$$

Statistical analysis

Definition of χ_{TV}^2

$$\chi_{\text{TV}}^2 \equiv \sum_j \frac{[P_j^{\text{TV}}(\delta_0, \rho_0) - P_j^{\text{TV}}(\delta^{\text{test}}, \rho^{\text{test}})]^2}{(\Delta P_j^{\text{TV}})^2}$$

$$P_j^{\text{TV}}(\delta, \rho) \equiv P_j(\nu_e \rightarrow \nu_\mu) - P_j(\nu_\mu \rightarrow \nu_e) \Big|_{\text{T2HK}}$$

j runs over energy bins, ρ is matter density of the Earth

In this study, we consider only statistical error.

The ΔP_j^{TV} is obtained by modified oscillation probability

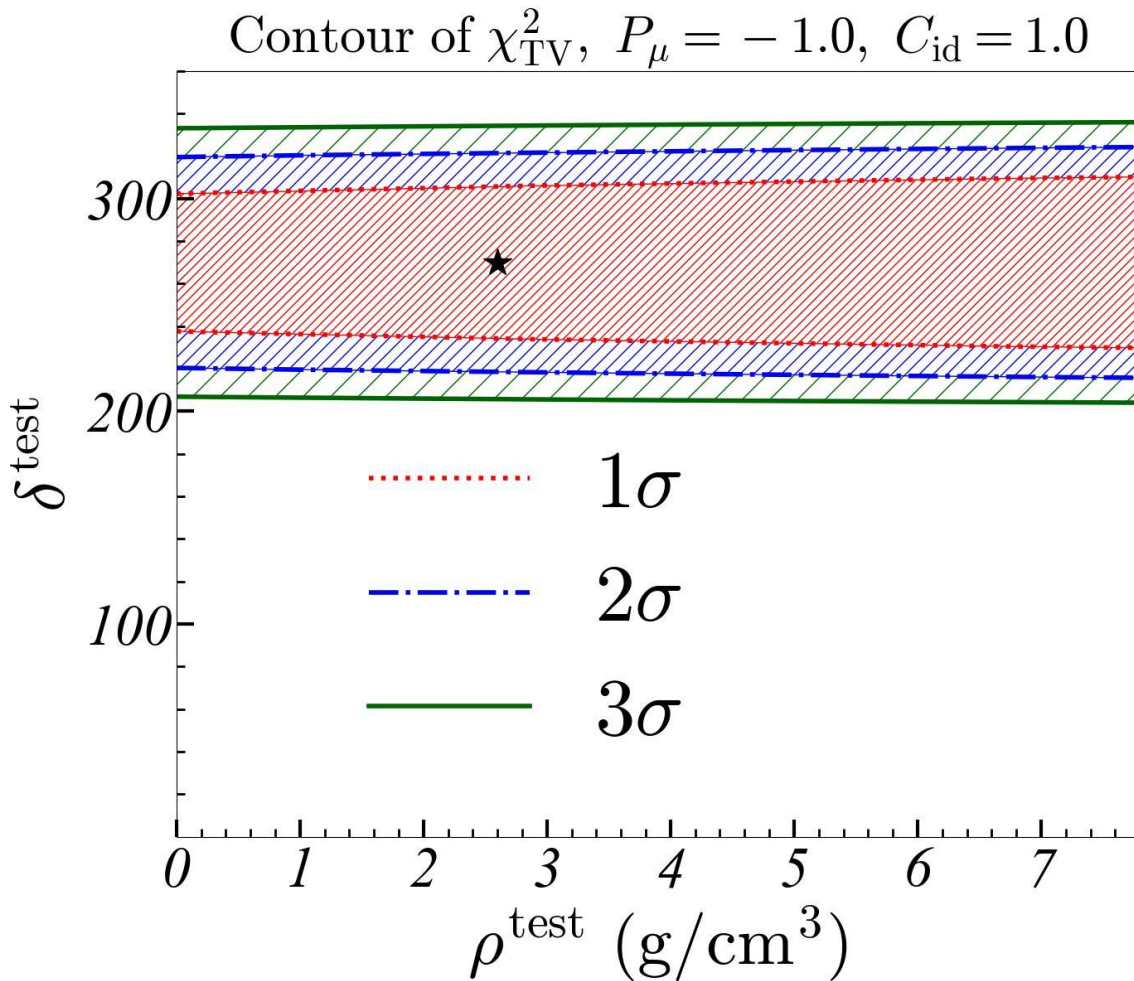
Reference values

$\Delta m_{21}^2/10^{-5} \text{ eV}$	$\Delta m_{31}^2/10^{-3} \text{ eV}$	θ_{12}	θ_{13}	θ_{23}
7.43	2.432	33.9°	8.49°	48.1°

The reference values are the arithmetic average of bfp in Particle Data Group 2022.

In this study, we consider only Normal Ordering.

Result of T violation



True value

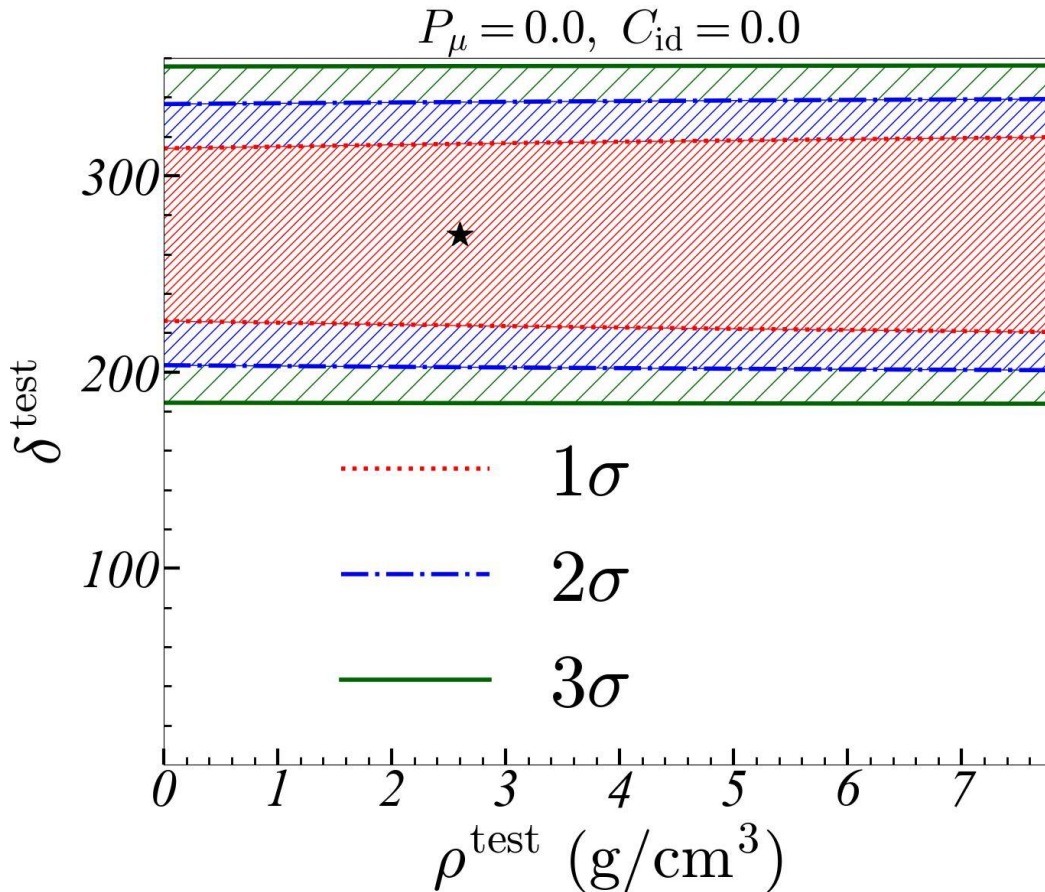
$$\delta_0 = 270^\circ$$

$$\rho_0 = 2.6 \text{ g}/\text{cm}^3$$

- χ_{TV}^2 only depends on δ .
- The confidence level as that for a 1 d.o.f., i.e., $\chi^2 = 1$ as 1σ .
- Perfect charge id and muon polarization, the δ is determined to be about $\pm 30^\circ$.

Free from matter density!

Result of T violation

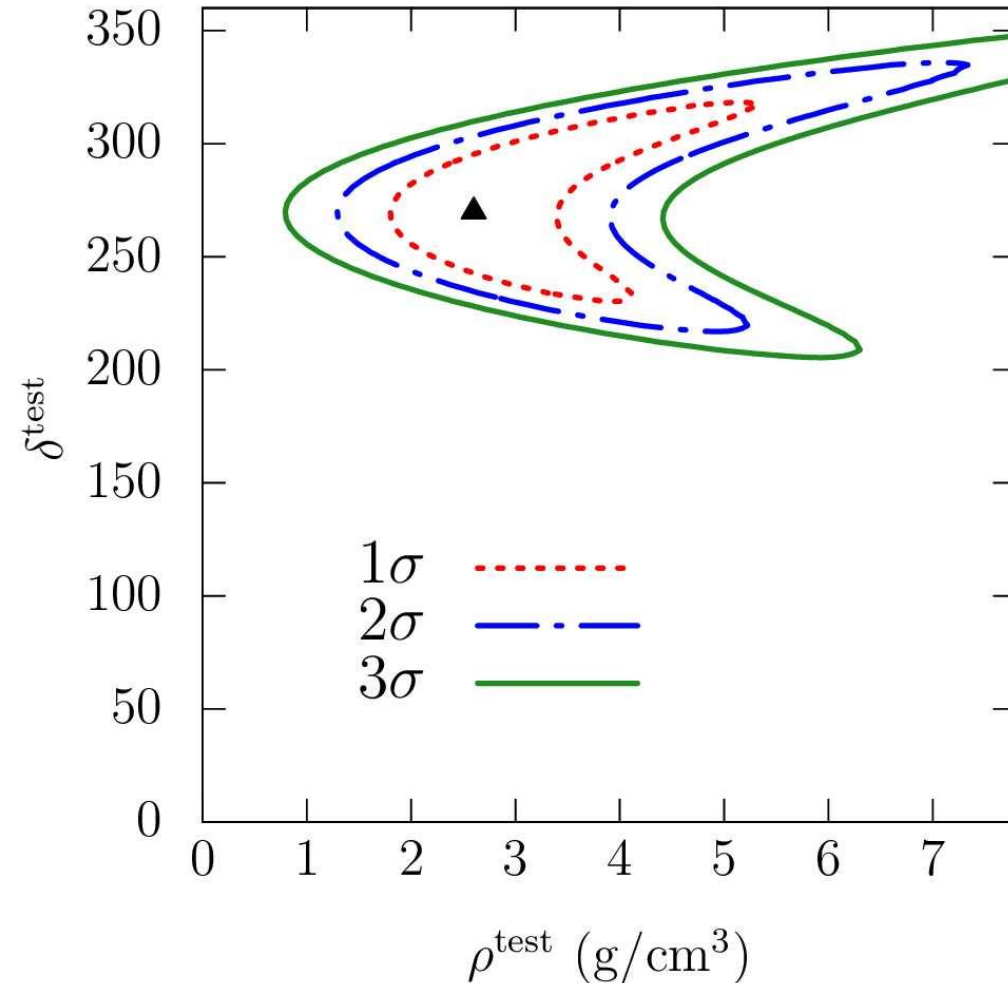


- χ_{TV}^2 only depends on δ .
- No charge id and unpolarized muon, $\delta = 0^\circ$ and 180° (CP (or T) conserving point), can be excluded at the level of 3σ .
- With no ability of distinguish ν_μ and $\bar{\nu}_\mu$, the sensitivities do not change significantly.

Free from matter density!

Comparing with CP violation

Contour of χ_{CP}^2



- Define χ_{CP}^2 in a similar way as T violation.
- **Non-trivial ρ dependence.**
- A good knowledge of the **matter density profile will be necessary.**
- T violation will be an **important additional information** for the measurement of the CP angle δ .

Summary

- **Testing T violation** in the lepton sector has not been achieved, so it is **important in particle physics**.
- We study the possibility of measuring T violation, by **combining $\nu_e \rightarrow \nu_\mu$ (μ TRISTAN \rightarrow HK) and $\nu_\mu \rightarrow \nu_e$ (T2HK)**.
- If nature has chosen $\delta = 270^\circ$, we can **exclude $\delta = 0^\circ, 180^\circ$ at more than 3σ** .
- Comparing with the χ^2 for T violation and CP violation, **T violation would not suffer from the uncertainty in the matter density profile of the earth**.
- In this study, we only consider the statistical error. A more complete analysis will be necessary to establish the feasibility.

Back up

Background subtraction

At the HK, in principle, ν_μ and $\bar{\nu}_\mu$ are distinguished by neutron tagging method.

We can define the oscillation probability $P(\nu_e \rightarrow \nu_\mu)$ as

$$P(\nu_e \rightarrow \nu_\mu) = \frac{1}{\kappa} \frac{\left(\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu} + (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} \right) - (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu} \Big|_{\text{T2HK}}}{\kappa N_{\text{near}}^{\nu_e \rightarrow \nu_e}}$$

$$\kappa \equiv \frac{1 + C_{\text{id}}}{2}, \quad C_{\text{id}}: \text{charge identification efficiency}$$

In the case of $C_{\text{id}} = 0.0$, we just do not perform the charge identification analysis and simply add the background events, and subtract the estimated amount by using the T2HK data, i.e. we take $\kappa = 1$ and $1 - \kappa = 1$.

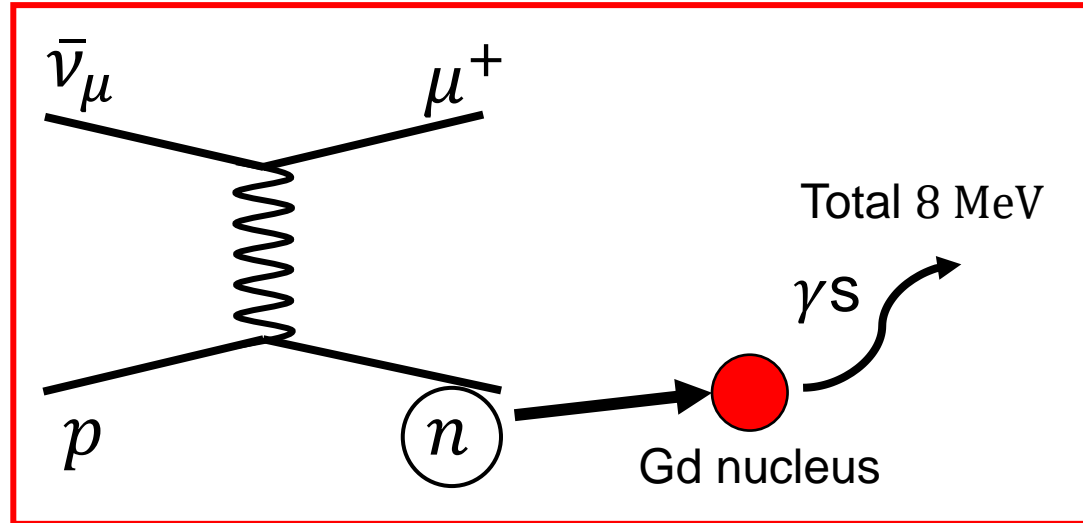
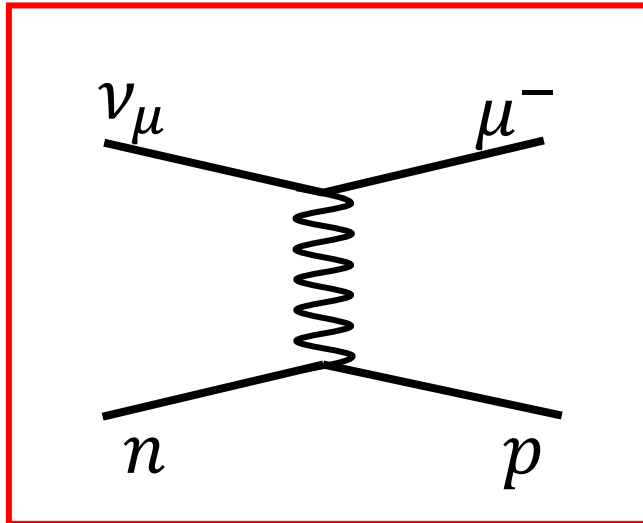
Statistical error ΔP_j^{TV}

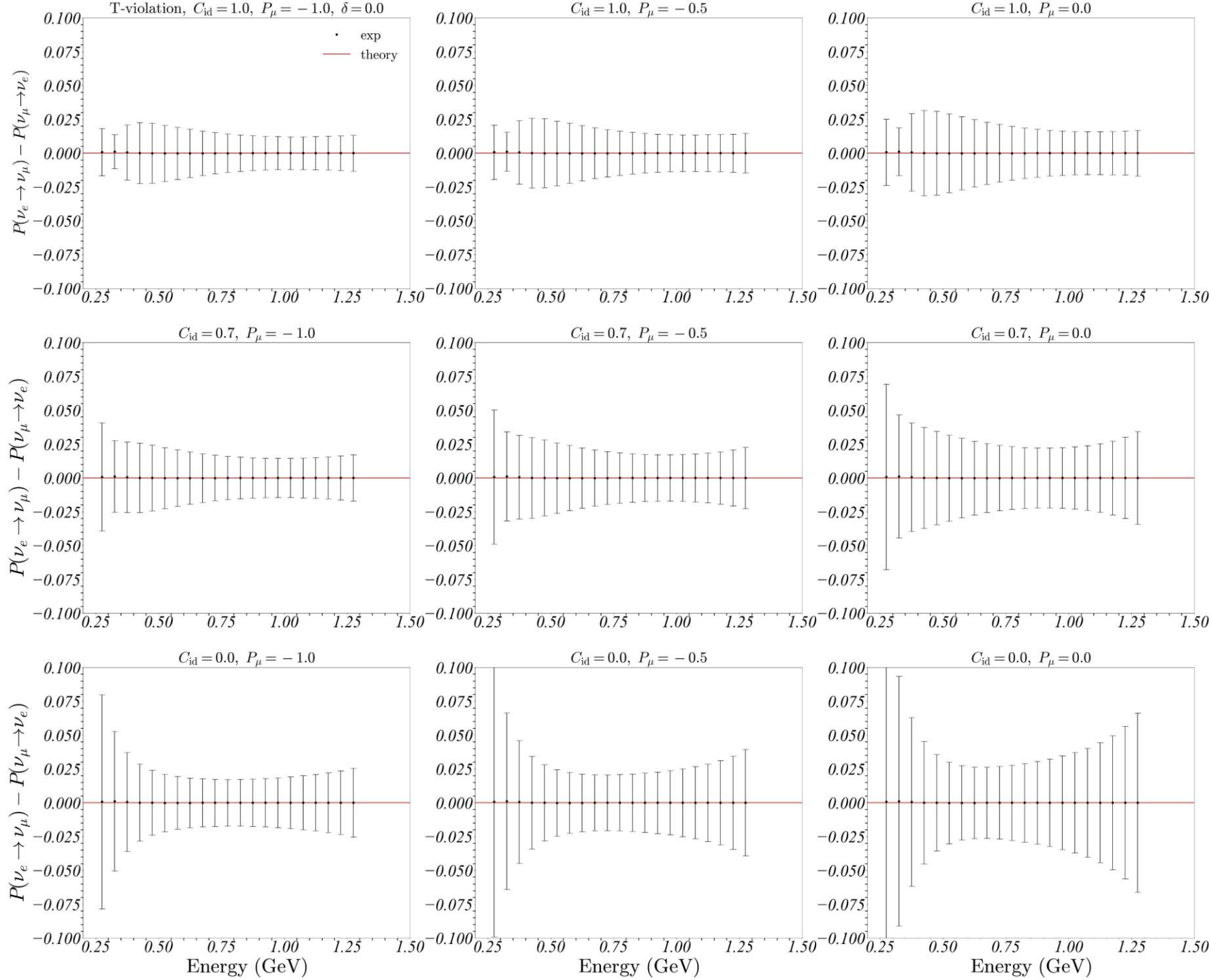
$$\begin{aligned}
 (\Delta P_j^{\text{TV}})^2 &= \left(\Delta P_j^{\nu_e \rightarrow \nu_\mu}\right)^2 + \left(\Delta P_j^{\nu_\mu \rightarrow \nu_e}\right)^2 \\
 &= \left(P_j^{\nu_e \rightarrow \nu_\mu}\right)^2 \left[\left(\frac{\sqrt{\left(\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu} + (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu}\right)}}{\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu}} \right)_j^2 \right. \\
 &\quad \left. + \left(P_j^{\nu_\mu \rightarrow \nu_e}\right)^2 \left[\left(\frac{\sqrt{N_{\text{far}}^{\nu_\mu \rightarrow \nu_e}}}{N_{\text{far}}^{\nu_\mu \rightarrow \nu_e}} \right)_j^2 + \left(\frac{\sqrt{N_{\text{near}}^{\nu_\mu \rightarrow \nu_e}}}{N_{\text{near}}^{\nu_\mu \rightarrow \nu_e}} \right)_j^2 \right] \right] \\
 &\sim \left(P_j^{\nu_e \rightarrow \nu_\mu}\right)^2 \left(\frac{\sqrt{\left(\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu} + (1 - \kappa) N_{\text{far}}^{\bar{\nu}_\mu \rightarrow \bar{\nu}_\mu}\right)}}{\kappa N_{\text{far}}^{\nu_e \rightarrow \nu_\mu}} \right)_j^2 + \left(P_j^{\nu_\mu \rightarrow \nu_e}\right)^2 \left(\frac{\sqrt{N_{\text{far}}^{\nu_\mu \rightarrow \nu_e}}}{N_{\text{far}}^{\nu_\mu \rightarrow \nu_e}} \right)_j^2
 \end{aligned}$$

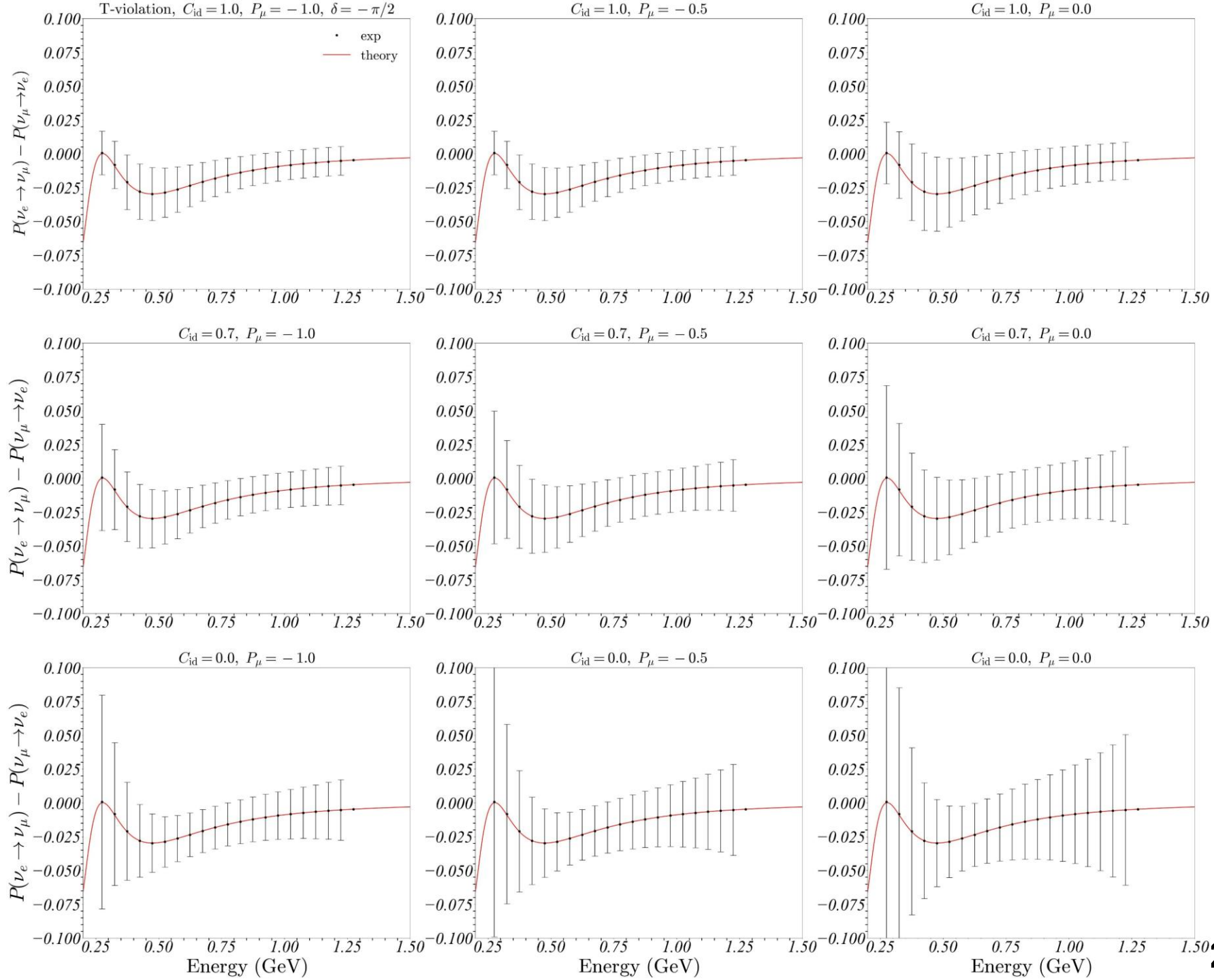
Neutron tagging

SK-Gd : efficiency $\sim 70\%$
Hyper-K : efficiency $\geq 70\%$

R. Akutsu. Ph.D thesis, Tokyo University, 2019.

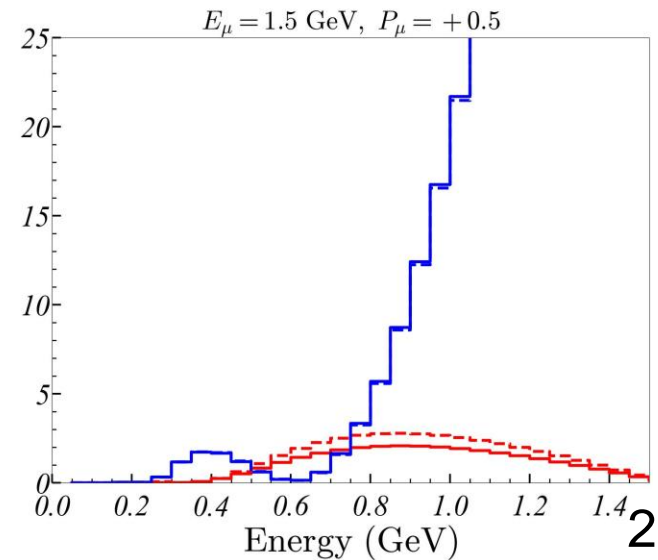
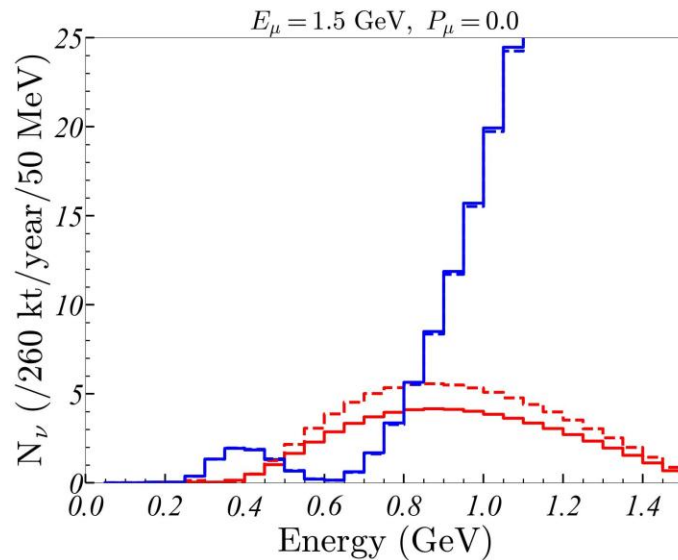
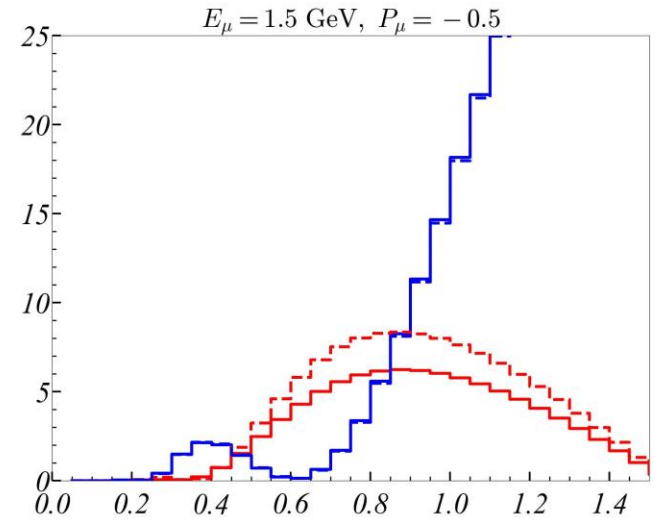
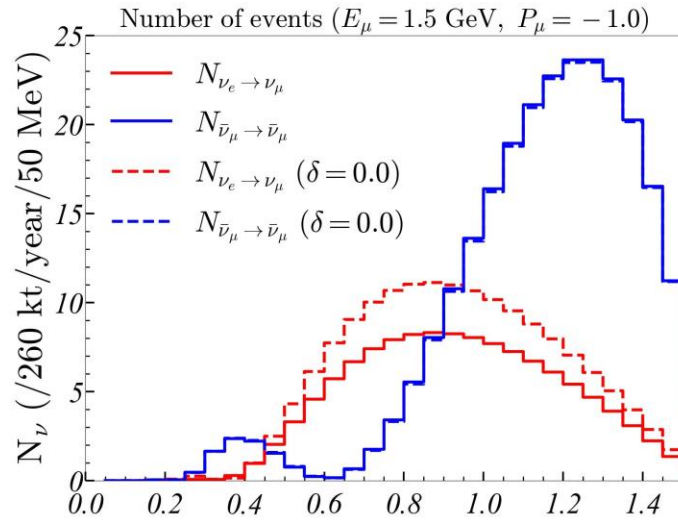




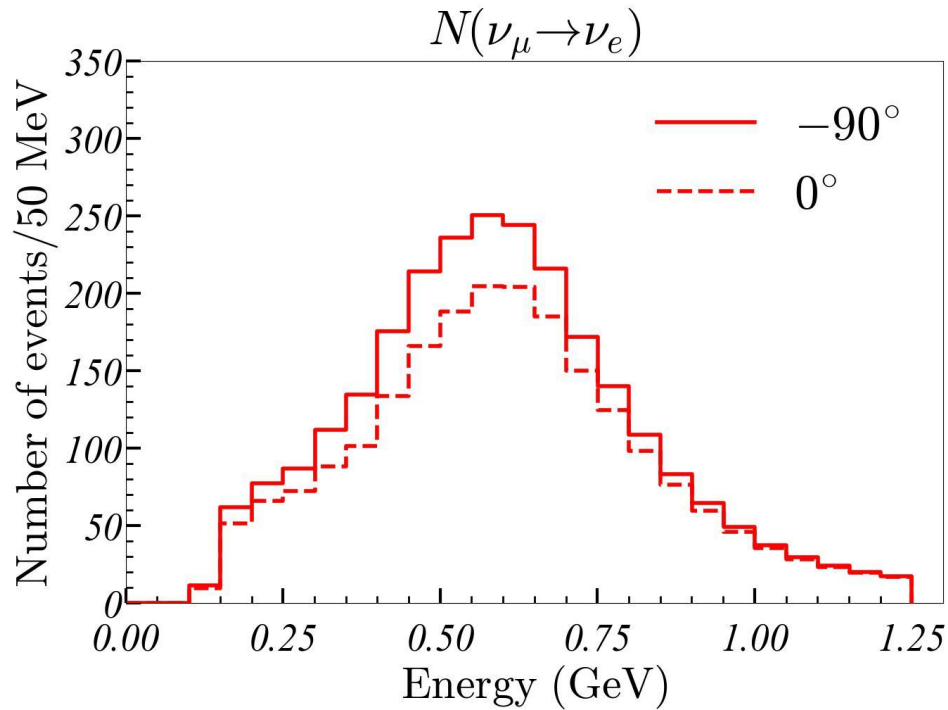


Number of events

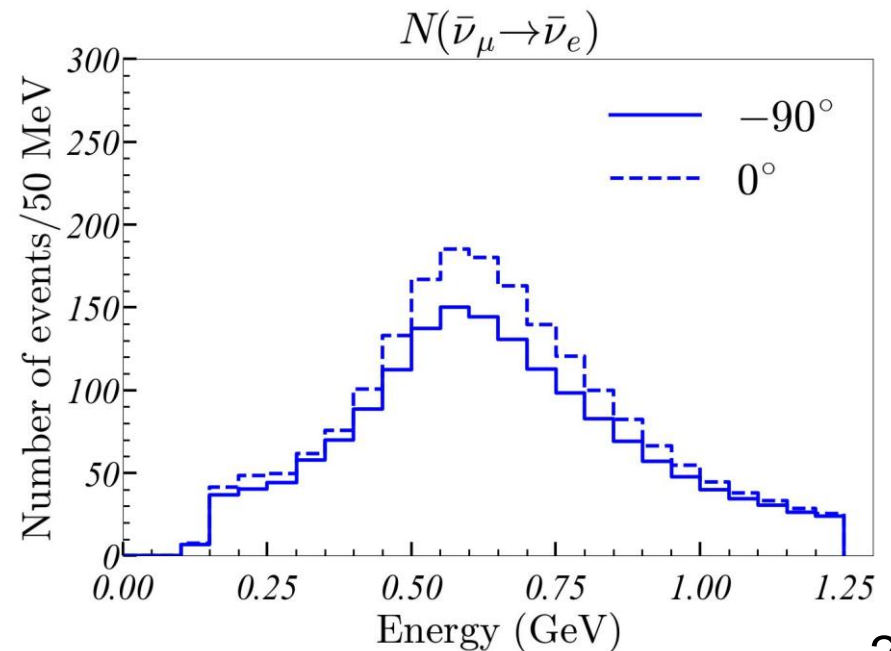
μ TRISTAN
 $N_\mu \sim 10^{21}$ muons/year



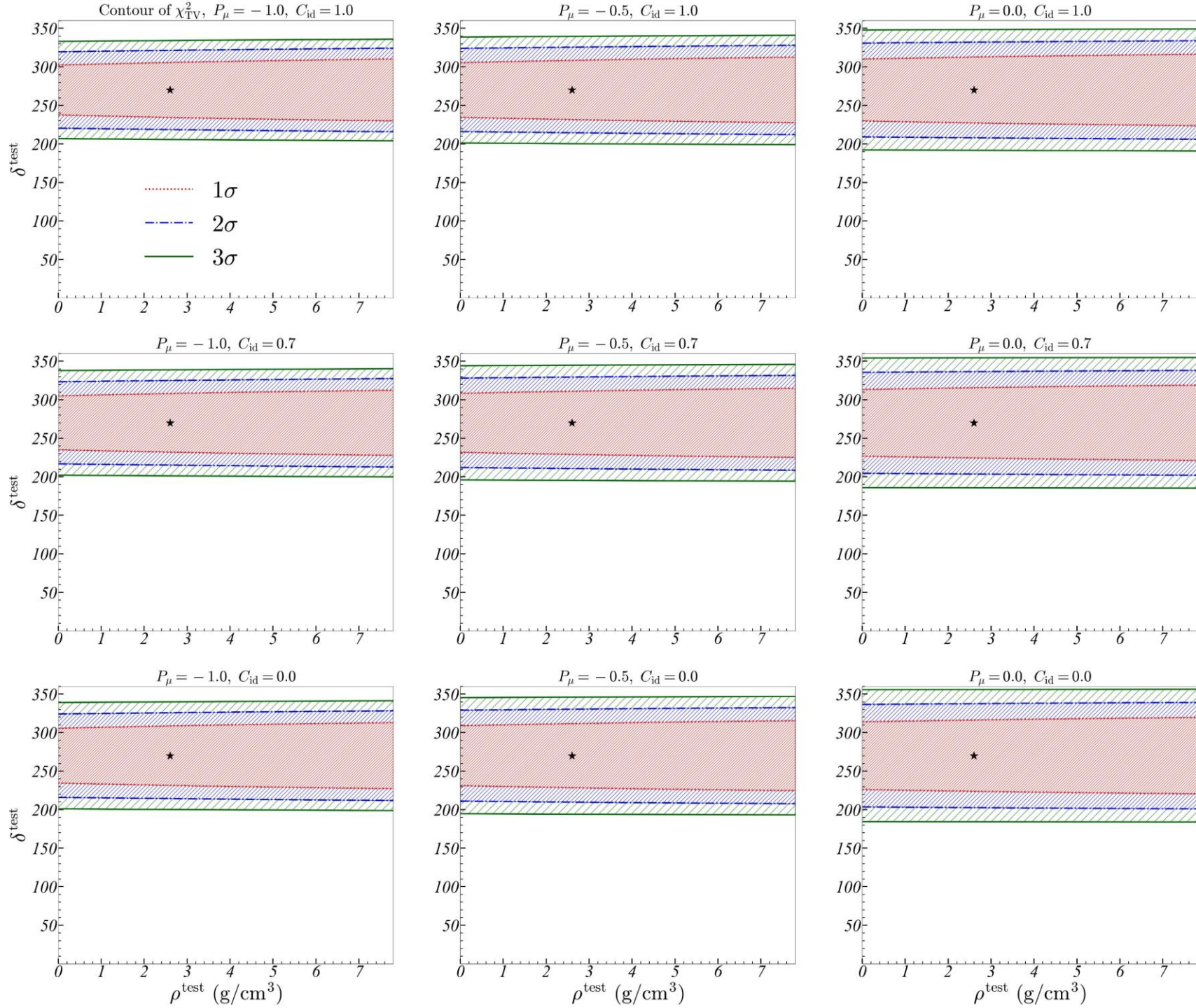
Number of events

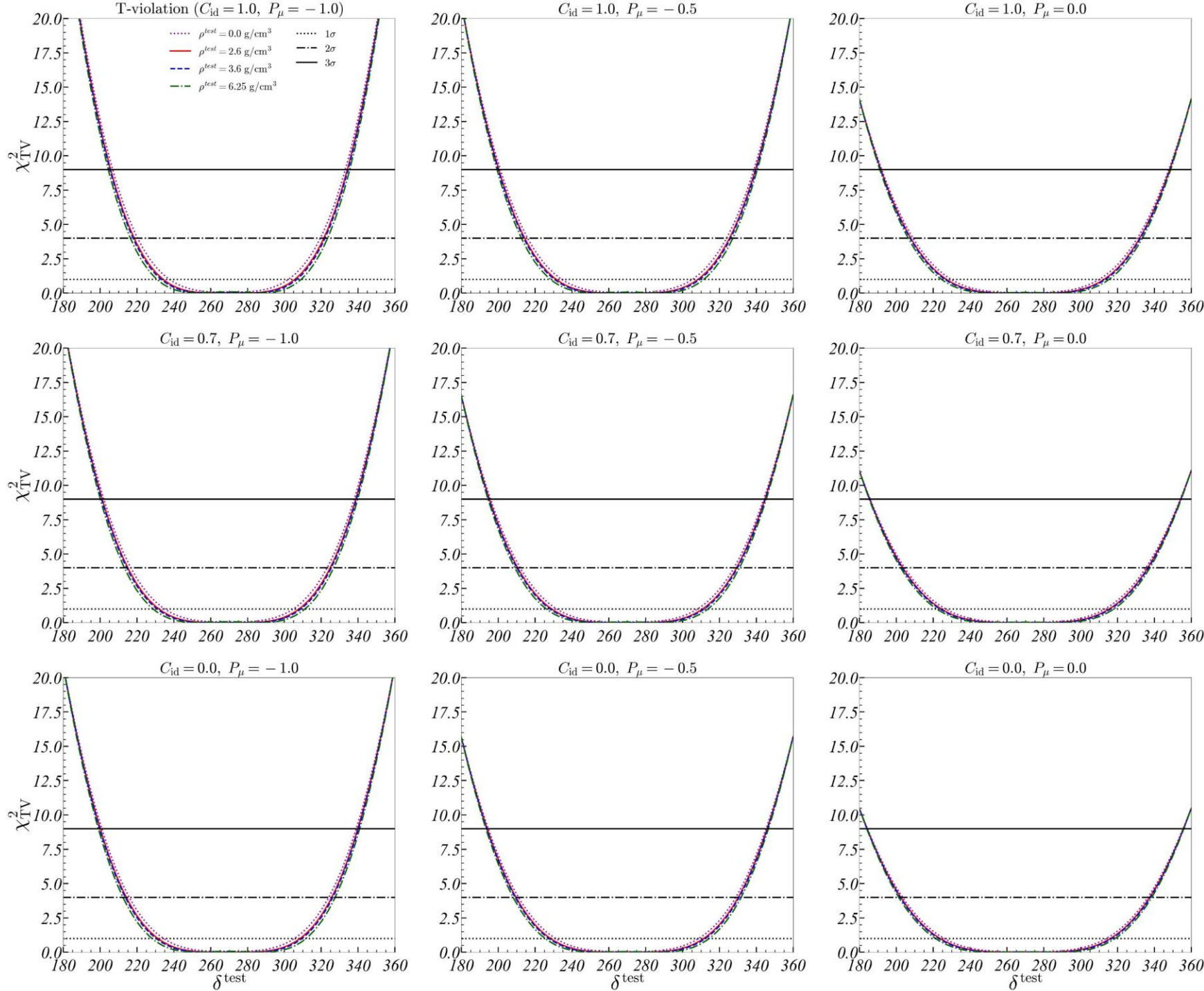


T2HK
 $\sim 10^{21}$ POT



K. Abe et al. [Hyper-Kamiokande Proto-], 2015





Possible CPT test?

Although we do not try in this study, one would be able to perform a similar analysis for **CPT violation**,

$$P_j^{CPT} = P_j(\nu_e \rightarrow \nu_\mu) - P_j(\bar{\nu}_\mu \rightarrow \bar{\nu}_e) \Big|_{T2HK} = P_j^{TV} + P_j^{CP}$$

Under our assumptions, this quantity would not measure anything. The analysis of this kind will be a quite important fundamental test of symmetry in physical laws of the Universe. Nevertheless, this will provide us with a quite important test of our underlying assumptions such as the three-neutrino scheme as well as quantum field theory.