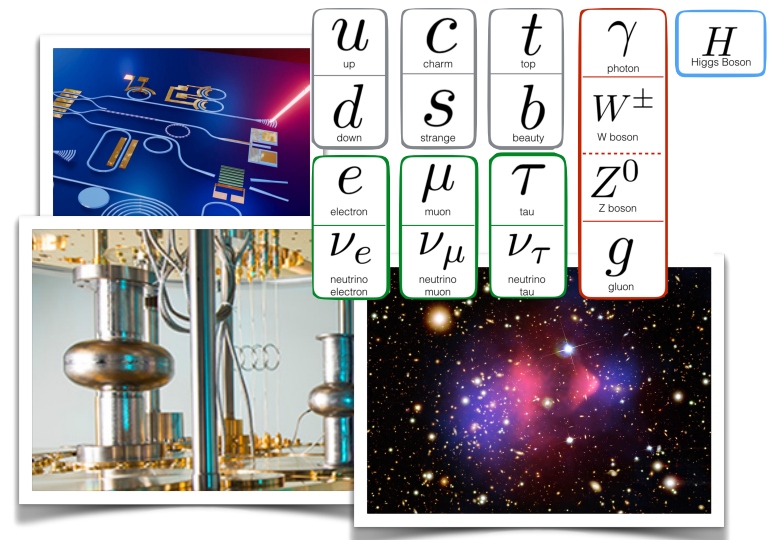


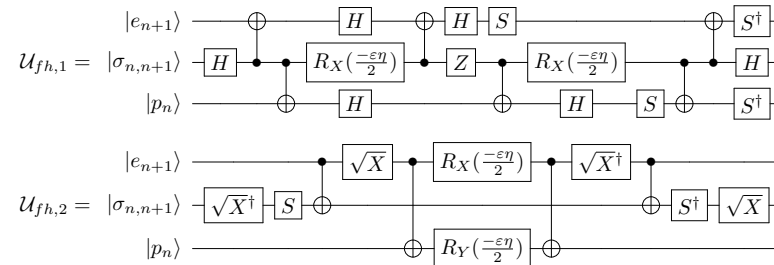
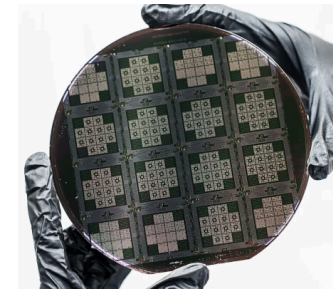
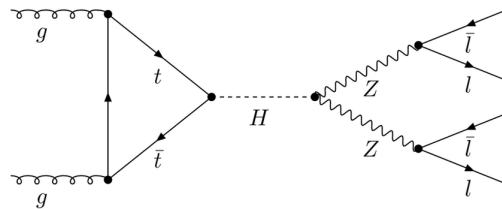
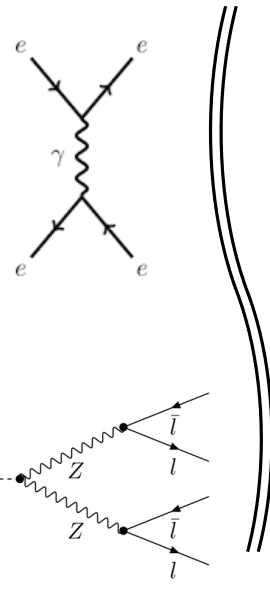
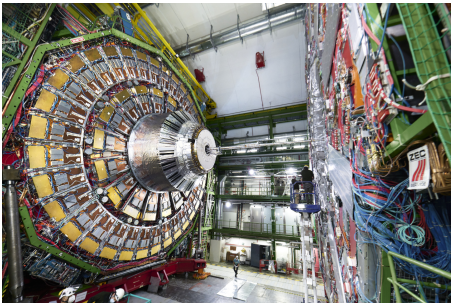
Quantum Sensing for Fundamental Physics

Roni Harnik,
Fermilab Quantum Theory Department
and SQMS



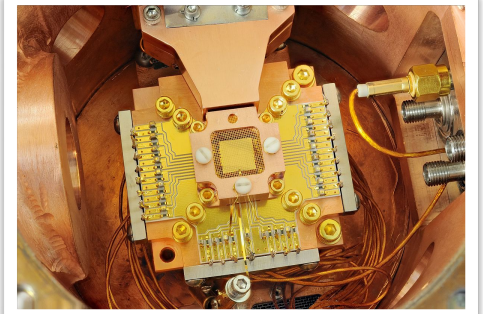
HEP - Quantum Interface

- Particle physics was always inherently quantum. Duh.
- A new field of Quantum physics is rapidly emerging, QIS.
- The interface is growing, but still small -

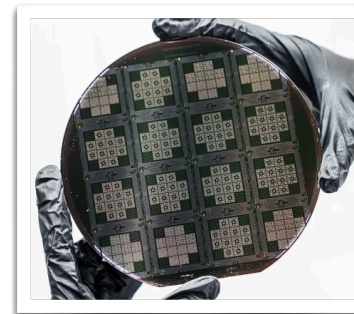


QIS spans many technologies.
 Very sensitive (and cool) devices!
 The goal is to manipulate information in quantum ways.

Ion traps



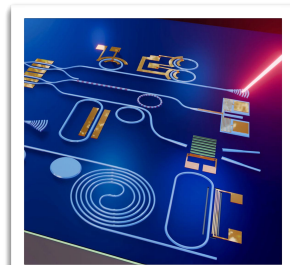
SC qubits



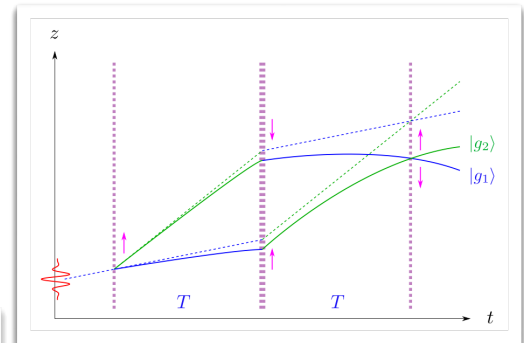
SC cavities



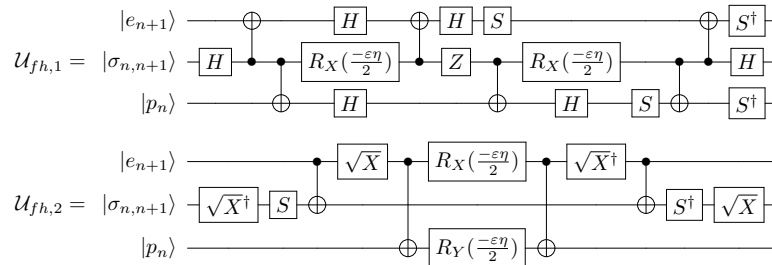
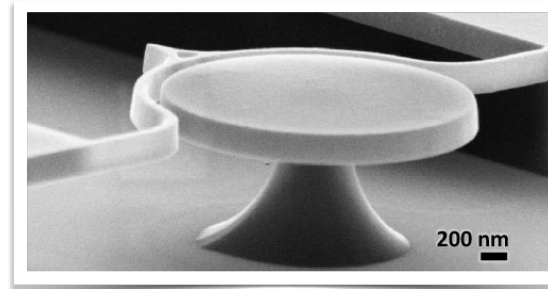
Integrated photonics



Atom interferometers



Optomechanical sensors



Quantum algorithms

Takeaway and Plan

Takeaway: we can bridge the gap b/w these fields

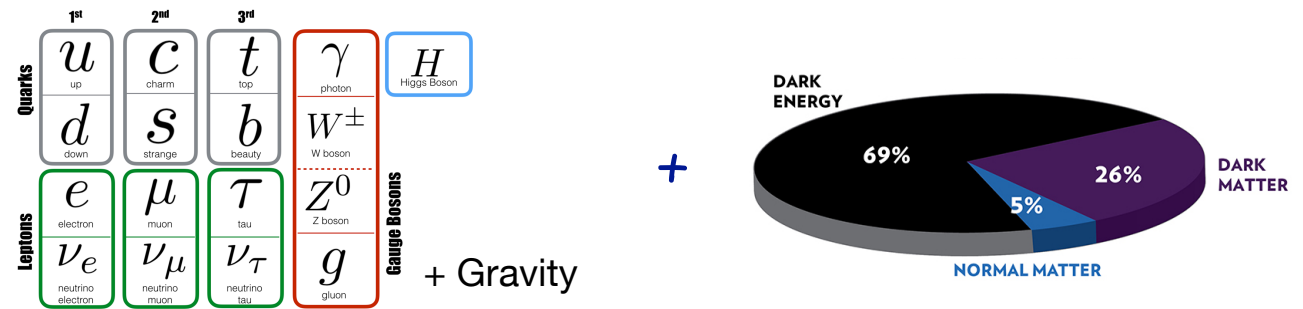
- HEP, QIS share a common foundation - Quantum Field Theory. These fields are not as distant as we sometimes think.
- We have new opportunities to explore Nature with quantum technology.

Plan:

- Talk about quantum devices in HEP language. (EFT)
 - Examples.
- Talk about BSM models in a QIS language.
 - Examples
- Examples of searches for BSM with sensors

Particle physics and its mysteries:

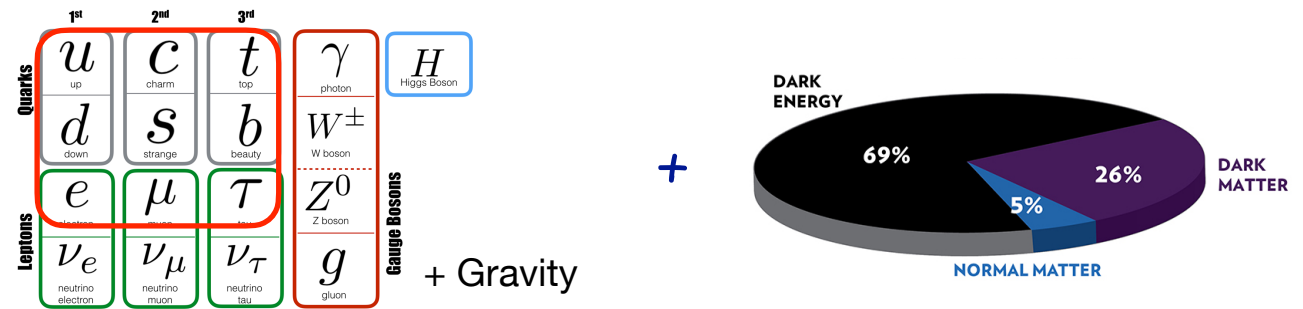
- The Standard model(s): a menu of particles & interactions + an energy budget,



+ a whole lot of mysteries.

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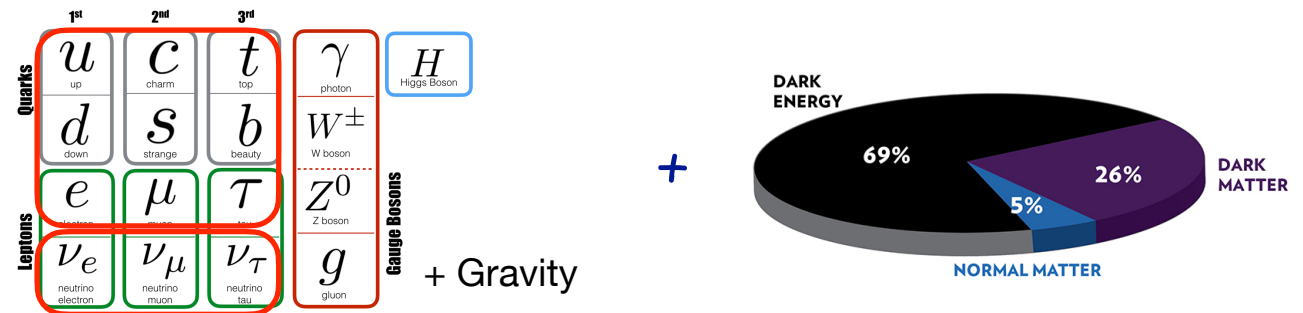
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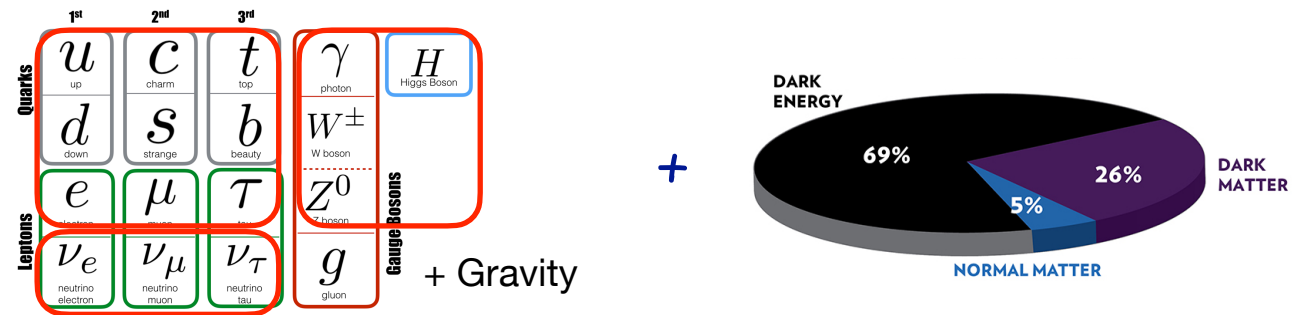
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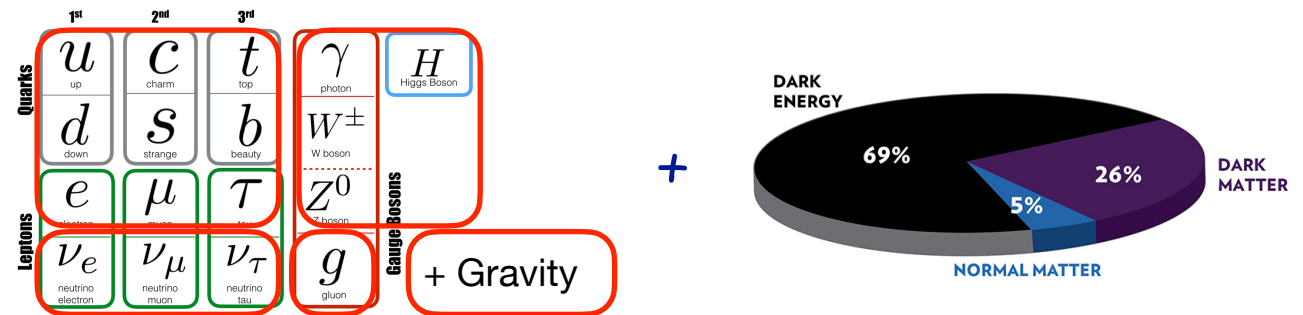
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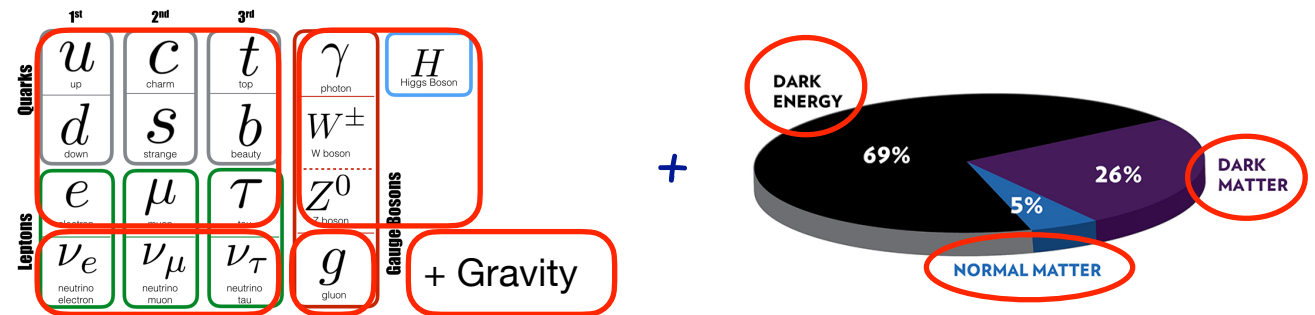
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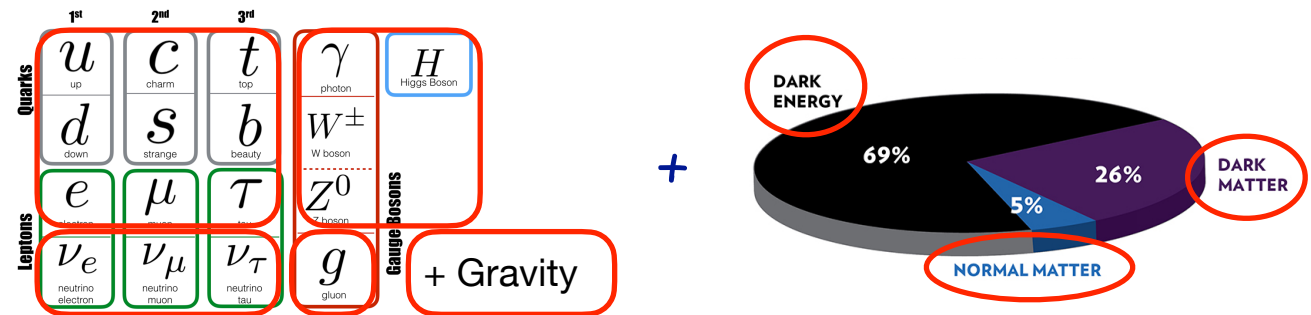
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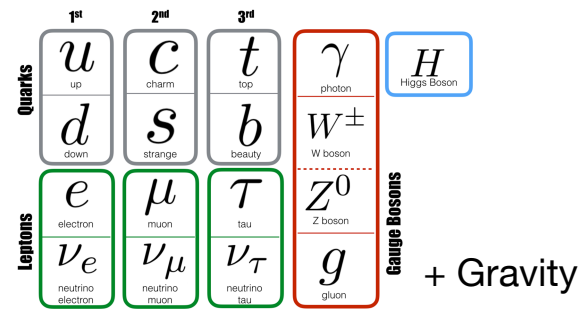
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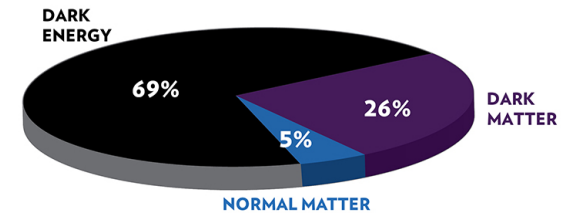
There is more. **BSM**. More fields! We'll get back to that!

Particle physics and its mysteries:

- The Standard model(s): a menu of particles & interactions + an energy budget,



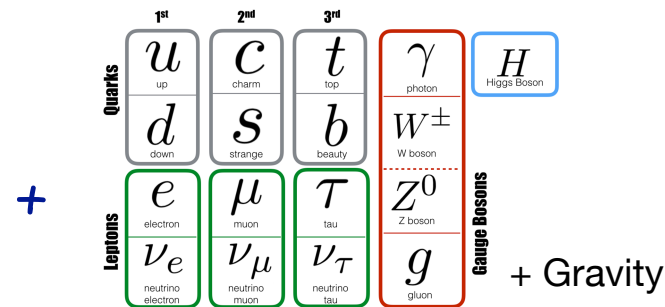
+



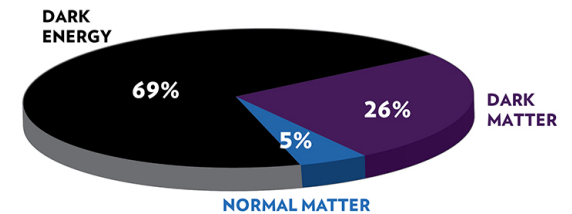
Particle physics and its mysteries:

- The Standard model(s): a menu of particles & interactions + an energy budget,

QFT:
 $|\psi(x)\rangle = \psi(x) |\Omega\rangle$
 $\mathcal{L} = \dots$



+



Every particle is a Field

QFT is continuum of interacting fields. All frequencies.



Quantum Field Theory

- At the heart of QFT is a mode expansion.

We get to pick the modes. Something like -

$$\phi(x_\mu) = \int \frac{d^3k}{(2\pi)^3} \frac{1}{\sqrt{2\omega}} \left(a_{\vec{k}} u_{\vec{k}}(\vec{x}) e^{i\omega t} + a_{\vec{k}}^\dagger u_{\vec{k}}^*(\vec{x}) e^{-i\omega t} \right)$$

Quantize: a, a^\dagger are operators.

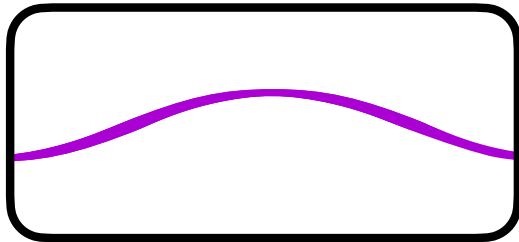
$$\text{Satisfy: } [a_{\vec{k}}, a_{\vec{k}'}^\dagger] = \delta_{\vec{k}\vec{k}'}$$

- This is sometimes referred to as "second quantization". For HEP its first!

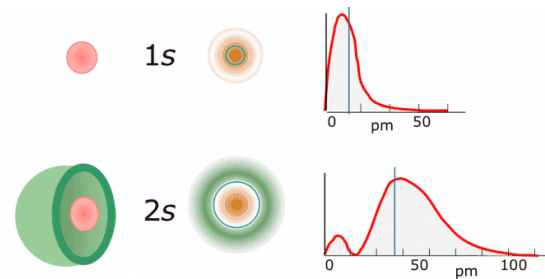
Quantum Fields in Small Devices

- In this big Universe, fields sometimes get localized to a finite regions. Either “naturally” or in a lab.

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or



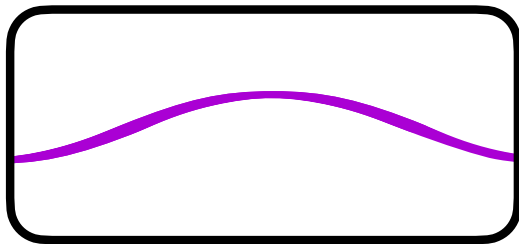
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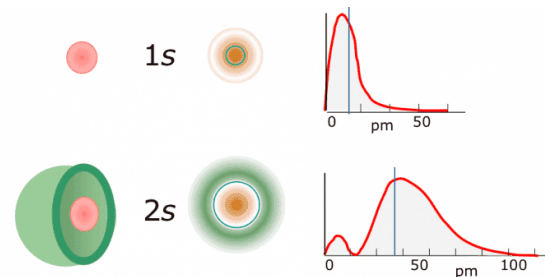
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Only a discretum satisfies boundary conditions.

$$+ \sum_j \frac{1}{\sqrt{2\omega}} \left(a_j u_j(\vec{x}) e^{i\omega t} + a_j^\dagger u_j^*(\vec{x}) (e^{-i\omega t}) \right)$$



or

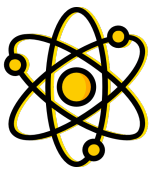


Quantum Fields in Small Devices

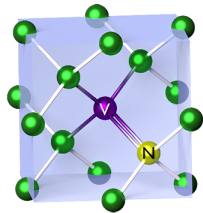
- Consider the low energy EFT of the discretum. Often in terms of a , a^\dagger

$$\phi_j(x_\mu) = \frac{1}{\sqrt{2\omega}} \left(a_j u_j(\vec{x}) e^{i\omega t} + a_j^\dagger u_j^*(\vec{x}) (e^{-i\omega t}) \right)$$

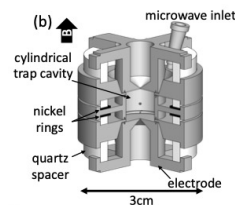
- In these EFTs, modes separate from the continuum, Quantum Mechanics shines:



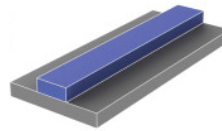
Atoms



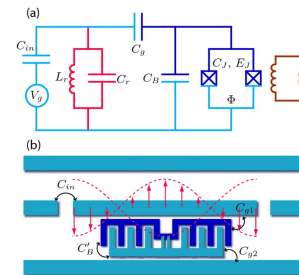
Defects



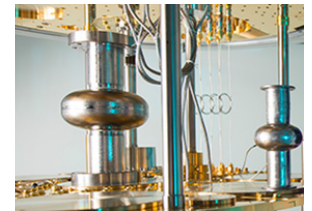
Artificial Atoms
(particle in trap)



Optical
waveguide



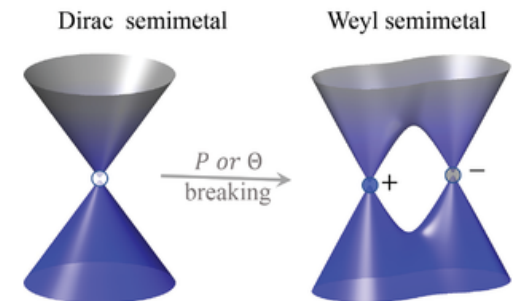
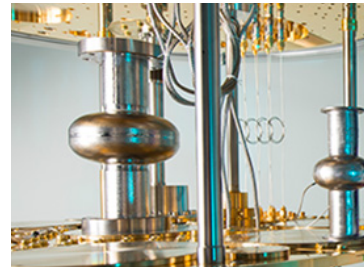
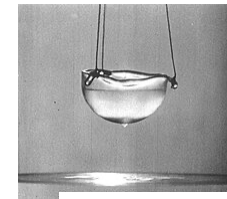
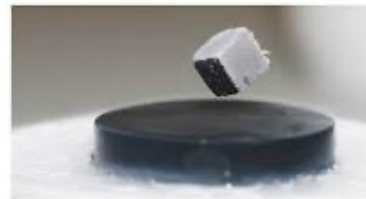
Superconducting
circuits



Electromagnetic
Cavities

New phases

- BTW: interesting quantum effects sometimes happen even without boundary conditions
- New phases, gaps, band structures, nontrivial dispersion... :
 - Superconductors
 - Superfluids
 - Semiconductors
 - Semi-metals
 - Spin glasses
 - Topological phase
 - ...



Of course, these phases are sometimes used to impose boundary conditions, and ...

New phases

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- New phases, gaps, band structures, nontrivial dispersion... :
 - Superconductors
 - Superfluids
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 - Topological phase
 - ...

In these phases, the convenient basis for the QFT mode expansion, the low energy EFT, involves collective excitations.

$$\phi_j(x_\mu) = \frac{1}{\sqrt{2\omega}} \left(a_j u_j(\vec{x}) e^{i\omega t} + a_j^\dagger u_j^*(\vec{x}) e^{-i\omega t} \right)$$

Apologies for over-simplifying a field that is not mine..

Of course, these phases are sometimes used to impose boundary conditions, and ...

Examples

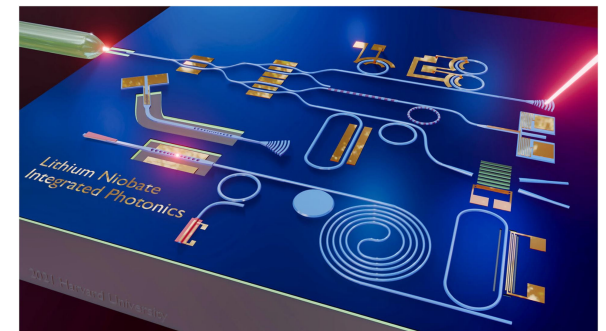
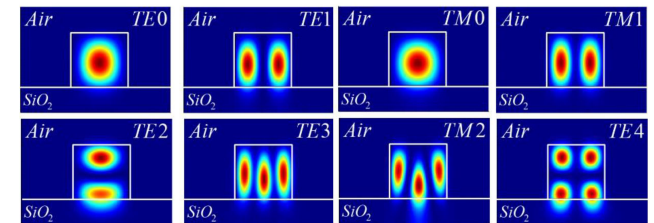
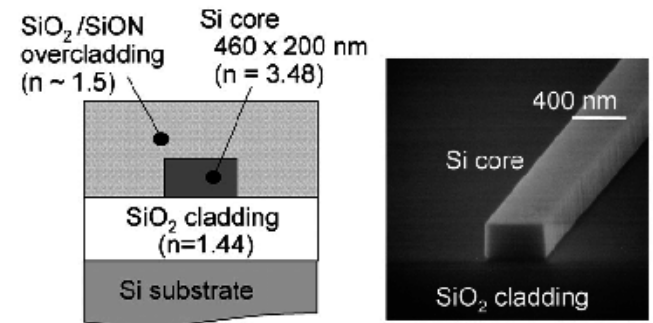
Optical Devices (e.g. integrated optics)

Superconducting circuits and cavities

Optical Devices

- Optics is the low energy EFT of light in matter.
- We can control the dispersion relation: $k = nw$.
Useful for localization.
- A waveguide admits a 1D EFT w/ modes quantized in transverse direction.
- Transverse wave function affects longitudinal dispersion relation (a la KK modes!)

Linear Optics: $H = E^2 + B^2 = \sum \hbar\omega(a^\dagger a + 1/2)$



"Integrated photonics"

Nonlinear Devices

- Like any EFT, in a quantum device there is a UV cutoff.
- We can add higher dim operators. For example, in optics

Dim-6:
$$H_{\text{SPDC}} = \int_{\text{crystal}} d^3\vec{x} \left(\chi_{jkl}^{(2)} E_j E_k E_l \right)$$

Dim-8:
$$H_{4\text{-wave}} = \int_{\text{crystal}} d^3\vec{x} \left(\chi_{jklm}^{(3)} E_j E_k E_l E_m \right)$$

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We can estimate χ 's in naive dimensional analysis:

When the field is set to that in an atom, we set (Dim-4 ~ Dim-6 ~ Dim-8):

$$E_{\text{atom}} \sim e/4\pi a_0^2 \quad \chi^{(2)} \sim \frac{\sqrt{4\pi}}{\alpha^{5/2} m_e^2} \quad \chi^{(3)} \sim \frac{4\pi}{\alpha^5 m_e^4} \quad \left(\text{by comparison, in vacuum } \begin{matrix} \chi^{(2)} = 0 \\ \chi^{(3)} = \frac{2\alpha^2}{45m_e^4} \end{matrix} \right)$$

Superconducting Cavities & Circuits

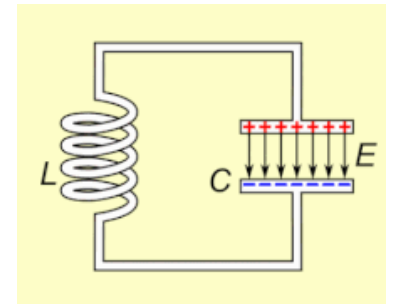
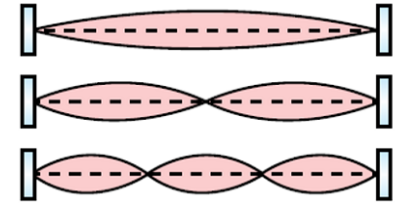
- Cavities: Light in a box. A discretum of states.
- Separation from the continuum is parametrized by Q .

$Q \sim 10^{10}$ is now routine. (Thank you accelerators!)

- LC Circuits: quantized current/flux.
- Control frequency with L & C , ($\omega^2 = 1/LC$).

Both are harmonic. Equally spaced levels.

$$H_{\text{mode}} = \hbar\omega(a^\dagger a + \frac{1}{2})$$

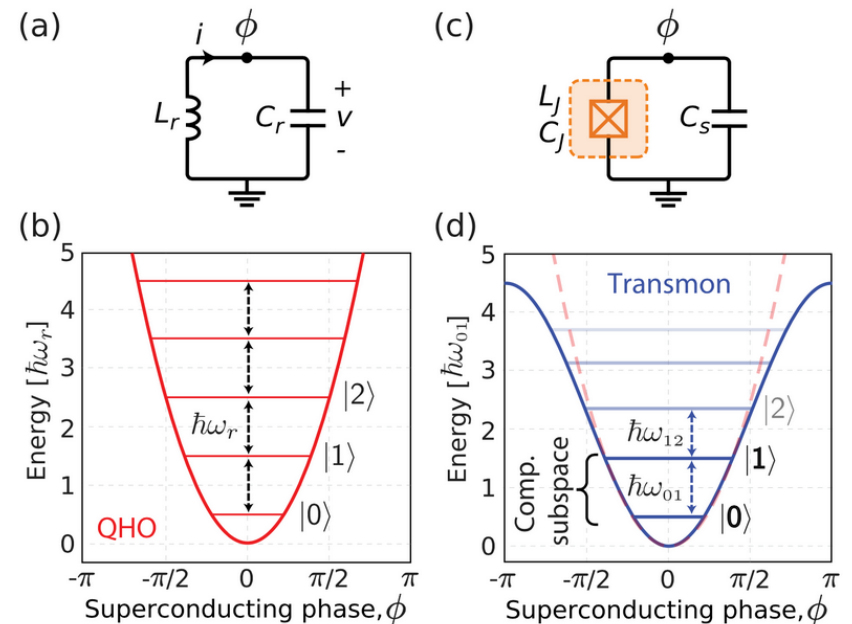


Nonlinear Devices

- Here too, we can arrange for a UV cutoff with higher dim operators,
- e.g. making L a function of $a^\dagger a$.

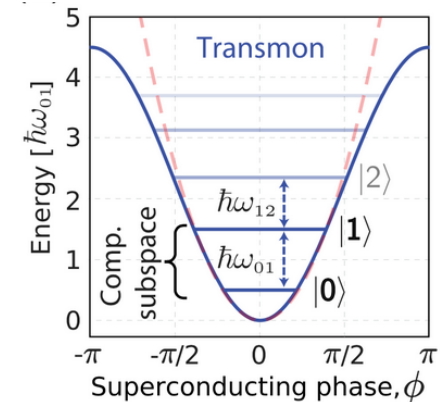
$$H = \hbar\omega(a^\dagger a + 1/2) + \kappa(a^\dagger a)^2$$

- Level spacing is nonuniform.
- Generically, $\kappa < \omega$. (“Naturalness” :-)
- This allows for control of individual levels of a given mode!

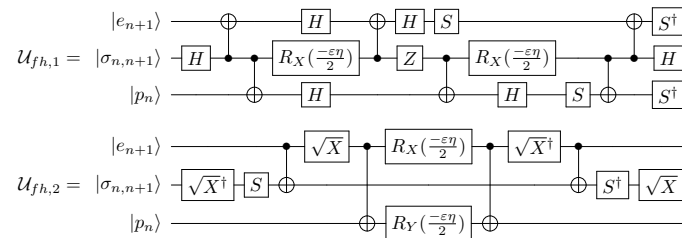


Quantum Control

- Within the device EFT, we can control the quantum state.
- Pulses of light can induce oscillation between states
- Occupation number, say, can store information & be read out.

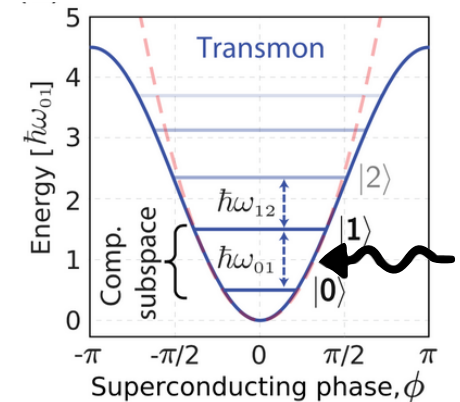


Qubit: $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle$

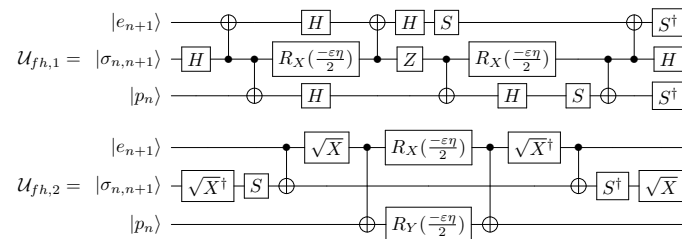


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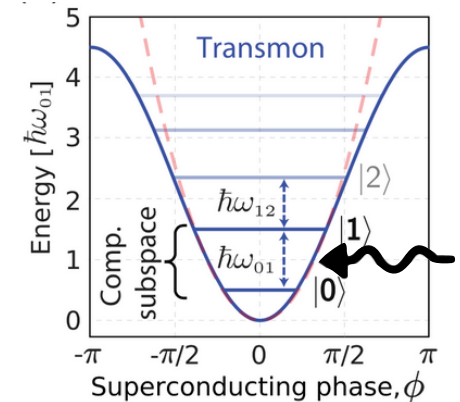


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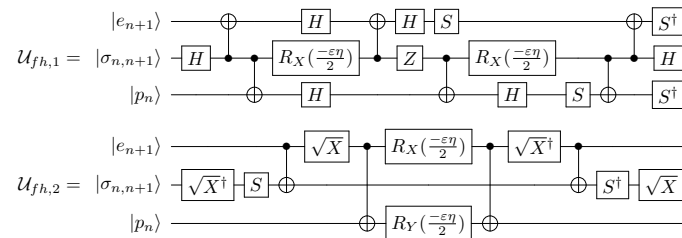


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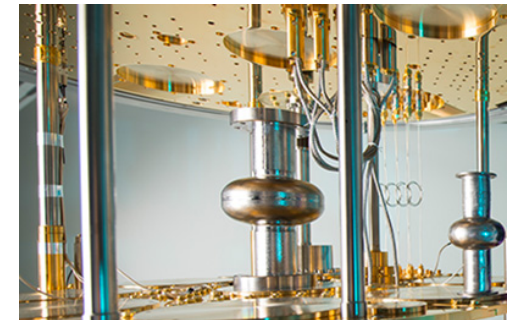
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Or Qudit: $|\psi\rangle = \alpha|0\rangle + \beta|1\rangle + \gamma|2\rangle \dots$



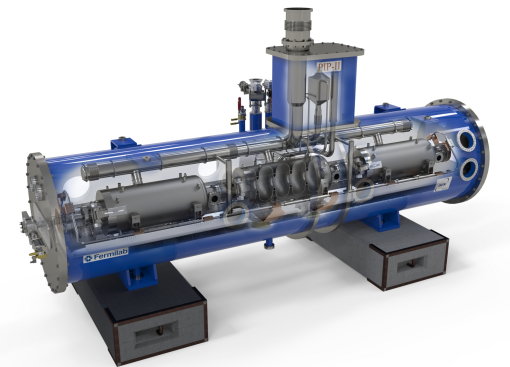
By the way ...



SQMS



- Superconducting devices are a central player in QIS.
- A limiting factor in their performance is loss (finite Q) & decoherence.
e.g. prepare a qubit, $\alpha|0\rangle + \beta|1\rangle$, after a while $|0\rangle \rightarrow |1\rangle$.
For modern qubits, a while \sim msec
- Fermilab, and other labs, had already developed ultra-high- Q cavities to accelerate particles! Lifetimes of $O(\text{sec})$!
- The SQMS center is using this technology and materials know-how for QIS.
- Achieving record coherence times in 2D and 3D systems, and understanding decoherence effects.

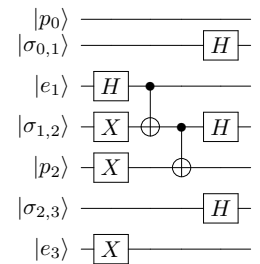
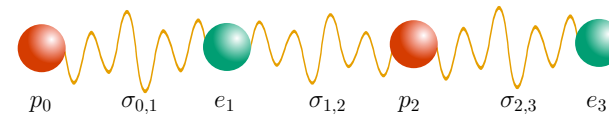


)

Quantum Simulation



- Famously, control of quantum states change the game. Quantum computers! (Shor's algorithm, etc).
- For HEP, quantum computers may tackle open computational challenges:
- Real-time simulation of HEP (e.g. of hadronization) on classical computers converge prohibitively slowly due to rapidly rotating phases (sign problem).
- But, even if quantum computers were fault tolerant today, we are not ready!
- Identifying efficient mappings of a lattice QFT to a QC, and the algorithm for simulation is an active area of research!

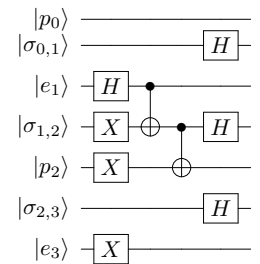
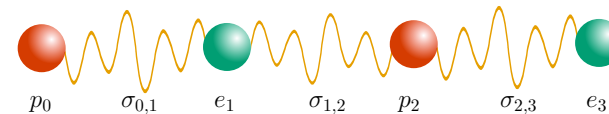


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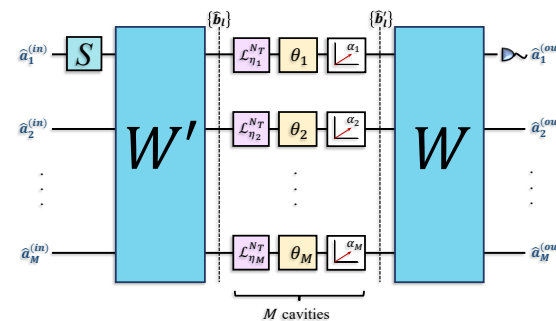
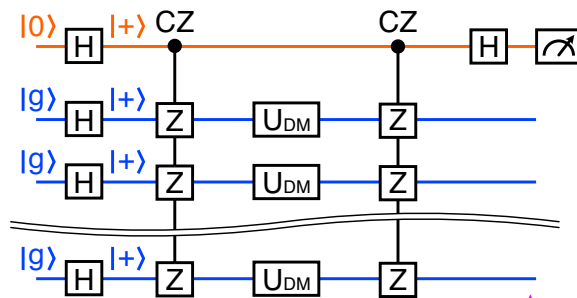
But not the focus of my talk.



Quantum Sensing

- The isolation of modes, and the ability to control them enables feeble effects to lead to dramatic consequences:
 - Appearance of mode occupation (Haloscope, light shining through wall)
 - Removal of mode occupation/phase transitions (TES, Nanowires: SC to normal)
 - Time evolution of ultra sensitive (entangled?) states

Chen et al, 2311.10413
Ito et al 2311.11632



Distributed squeezing:
Brady, RH, et al, *PRX Quantum* 3 (2022) 3, 030333

BSM - for Quantum Mechanics

New Physics → *New Fields*

	1 st	2 nd	3 rd	
Quarks	u up	C charm	t top	γ photon
	d down	S strange	b beauty	W^\pm W boson
Leptons	e electron	μ muon	τ tau	Z^0 Z boson
	ν_e neutrino electron	ν_μ neutrino muon	ν_τ neutrino tau	g gluon
				H Higgs Boson

+ Something new.

Ok. For concreteness,

(and because QIS is often about controlling light)

lets assume the new field couples to photons.

Linear or nonlinear?

Dark Photons - a Linear Extension

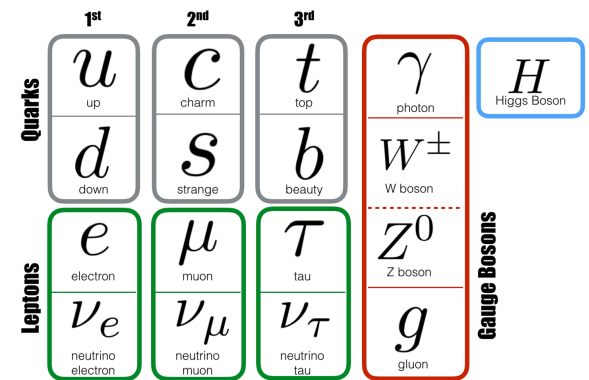
- If something mixes linearly with the photon, it must have the same quantum numbers:

- The dark Photon effective Hamiltonian:

$$\mathcal{H} \supset \mathcal{H}_{\text{QED}} + \varepsilon \vec{E} \cdot \vec{E}' + \vec{B} \cdot \vec{B}' \quad \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~}$$

(and dark photon also has a mass, and a longitudinal polarization!)

(OFC, a dark photon, if it exists, would teach us profound lessons!
New force of nature. Grand Unification, etc.)



Dark Photons - a Linear Extension

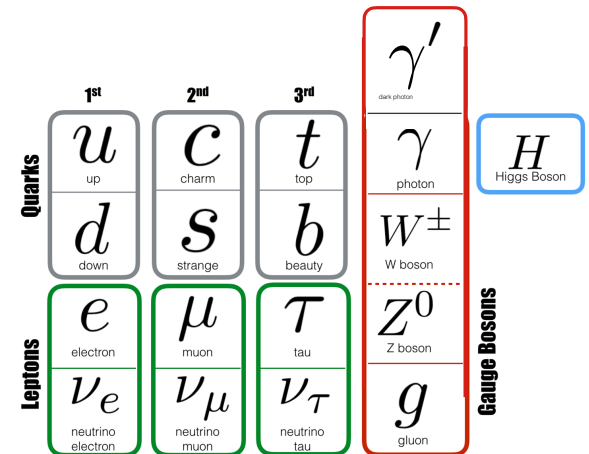
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□ The dark Photon effective Hamiltonian:

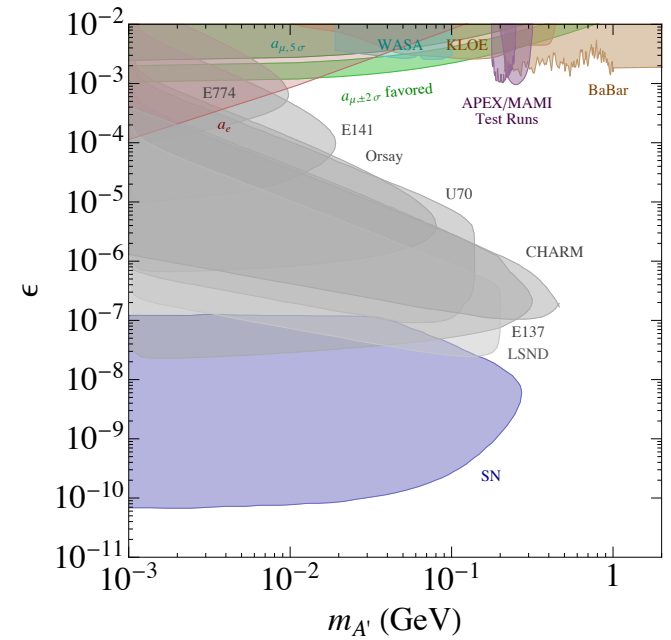
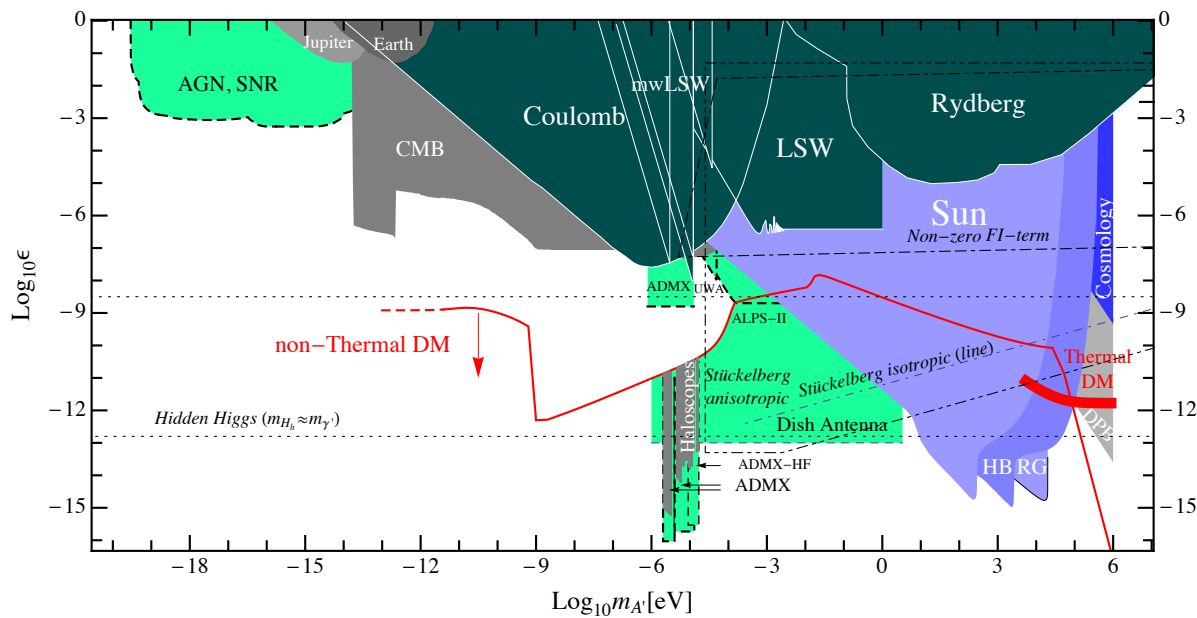
$$\mathcal{H} \supset \mathcal{H}_{\text{QED}} + \varepsilon \vec{E} \cdot \vec{E}' + \vec{B} \cdot \vec{B}' \quad \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~} \text{~~~~~}$$

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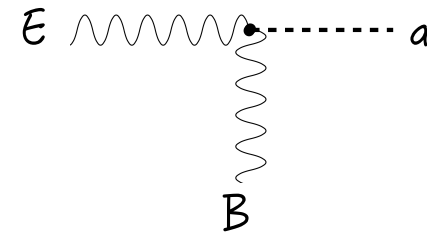


Pecci and Quinn (77)

Axions - A nonlinear extension of QED

- A nonlinear interaction, naively, would involve 2 photons & 1 new field.

$$\mathcal{L} \supset \frac{a}{f} F^{\mu\nu} \tilde{F}_{\mu\nu} = \frac{a}{f} \vec{E} \cdot \vec{B}$$



- Axion phenomenology:

- Axion mixing w/ photons polarized along background B field.
- Axion can be absorbed by photon \rightarrow up conversion.
- Axion exchange \rightarrow photon nonlinearity in vacuum.

[Sikivie]

[Berlin et al (2019)]

[Gao, RH (2020)]

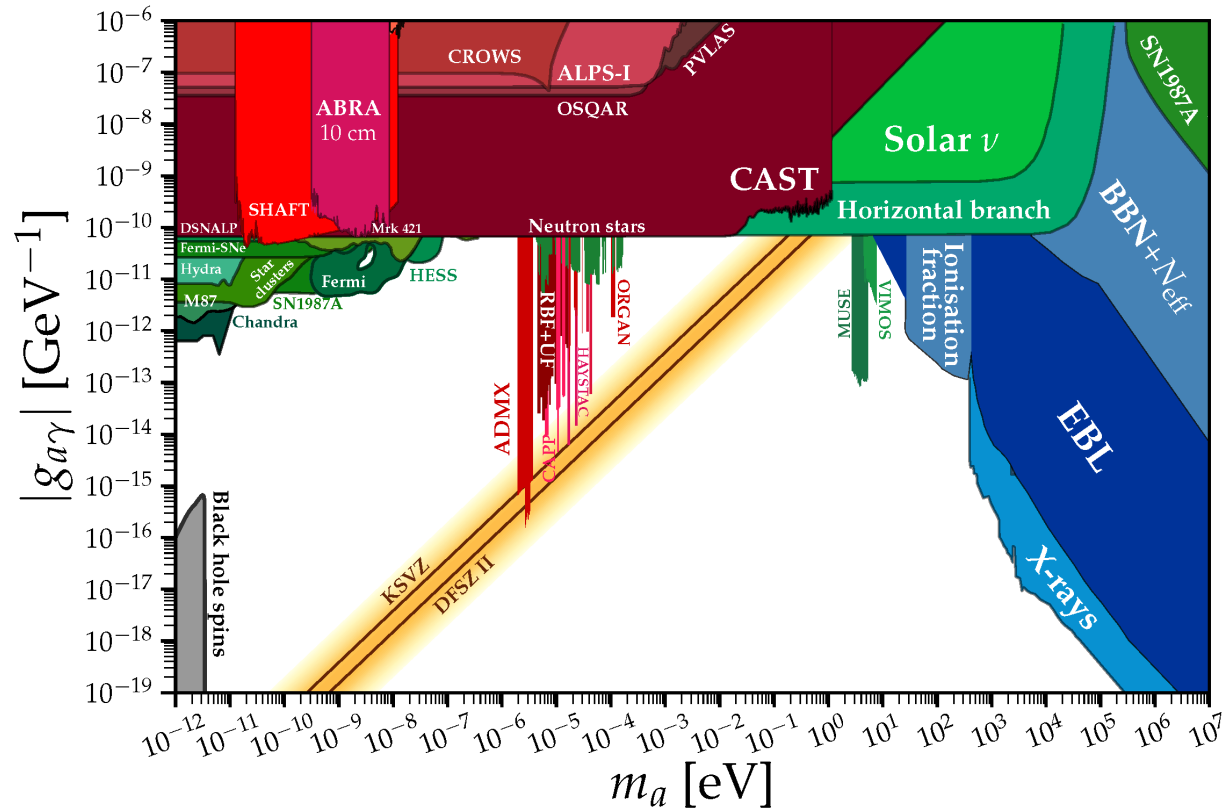
[PVLAS]

[e.g. Bogorad, Hook, Kahn, Soreq (2020)]

(Of course, the discovery of an axion will be a profound insight! Strong CP. etc.)

Pecci and Quinn (77)

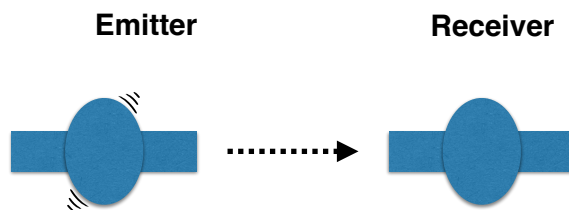
Axions - A nonlinear extension of QED



New Particle vs Dark Matter Searches

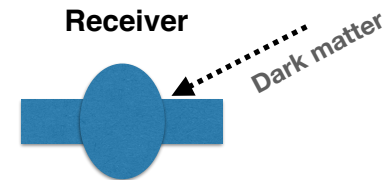
- We can test two distinct hypotheses:
 - Axion/Dark photons are new fields we can source in a lab.
 - They also make up the (cold) dark matter of the Universe.

Lab search



e.g. Light Shining through wall.
Frequency is known.

DM search



Frequency is set by DM mass
and is unknown. Need to scan!

Examples

Cavities

Optics

Qubits

Atoms

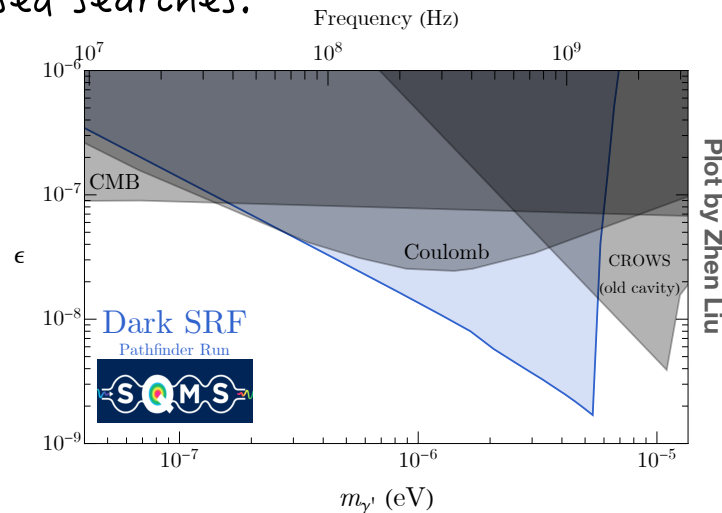
I probably have more example than we have time for...

Dark SRF: cavity-based search for the Dark Photon

Look for linear model with linear devices:
A light-shining-through-wall experiment.



Phase 1: Pathfinder run in LHe.
Demonstrated enormous potential for SRF based searches.



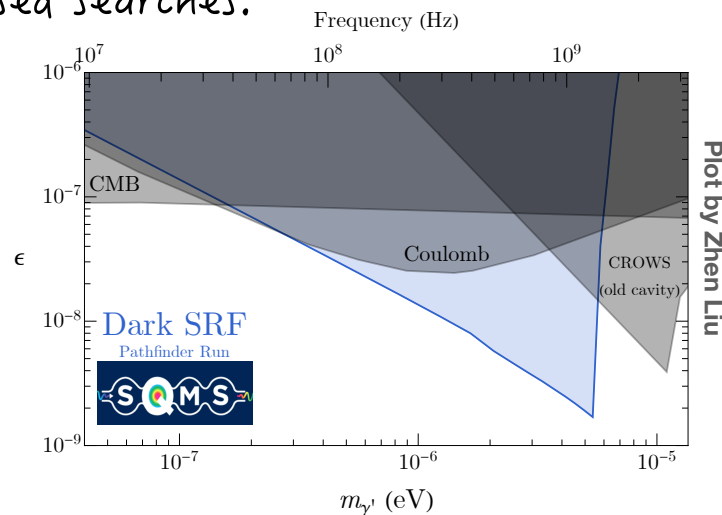
Romanenko et al., *Phys.Rev.Lett.* 130 (2023) 26, 261801

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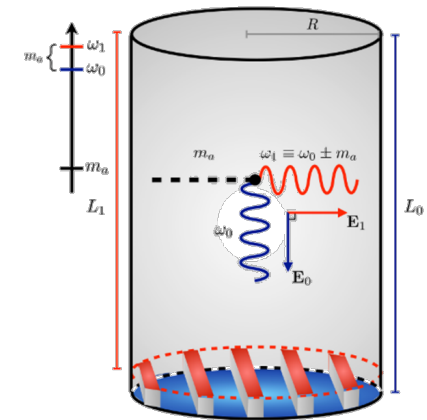
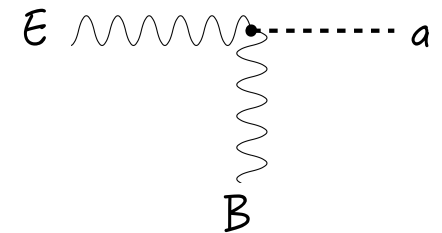
Phase 2: in DR,
receiver at \sim mk, in
quantum regime.
Improved frequency
stability. Photon
counting and/or
phase sensitive
readout.

Increased reach.

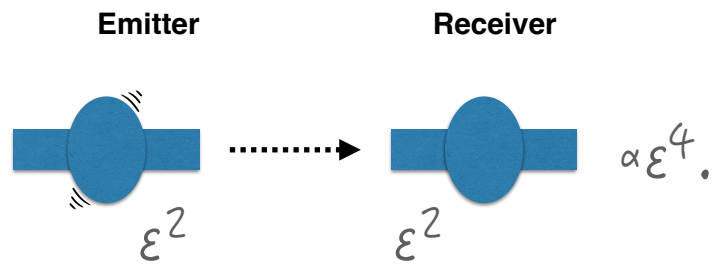


Axion searches with Cavities

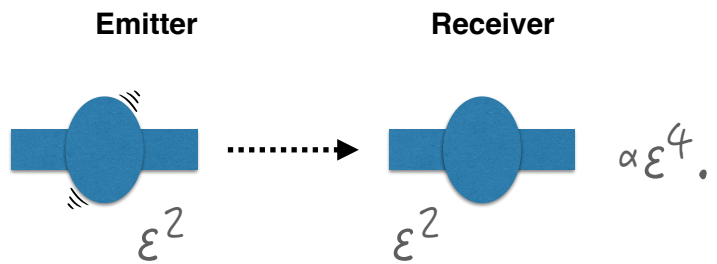
- We can look for nonlinear effects.
 - DM search: Turn on a B field*, or up conversion of pump mode photon into a nearby signal mode. [Sikivie] [Berlin et al (2019)]
 - Lab search: look for nonlinearity between modes (e.g. 2 pumps, one signal. [e.g. Bogorad, Hook, Kahn, Soreq (2020)]
- Condensed matter connection: In many cases, the inherent nonlinearity of SC, nonlinear-Meinser. [J. A. Sauls, Prog. Theor. Exp. Phys. 2022, 033103 (2022)]
- A double cavity setup likely needed. [Gao, RH (2020)]



* Turning on a strong B field in a SC cavity and keeping high Q is an open research challenge.

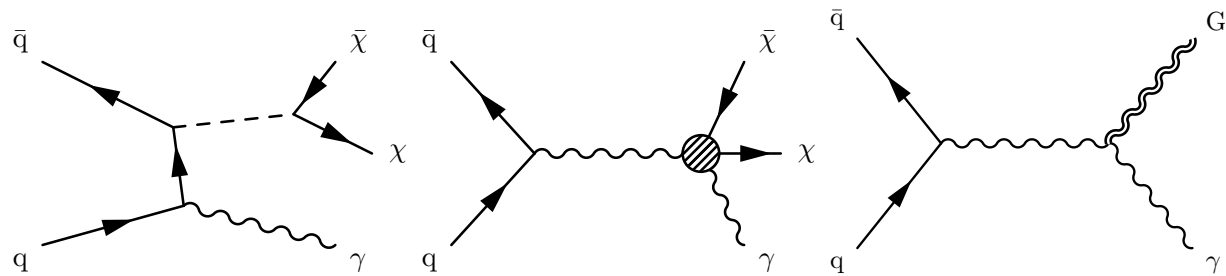


Can we do better?



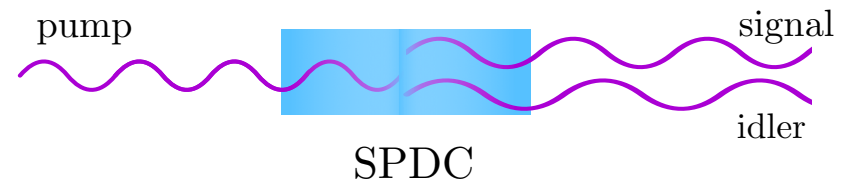
Can we do better?

CMS, ATLAS:

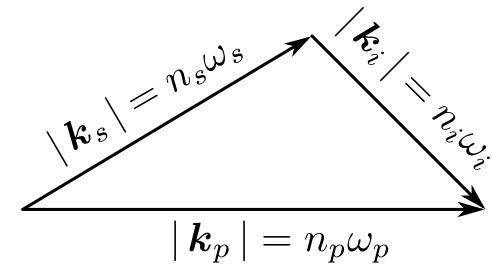


Nonlinear Optics with Dark States

- SPDC: a workhorse in quantum optics.
- Pump \rightarrow signal + idler (a "decay")
- A particle physicist may think such a decay has not phase space.
- Recall: pump, signal and idler can have different n 's!



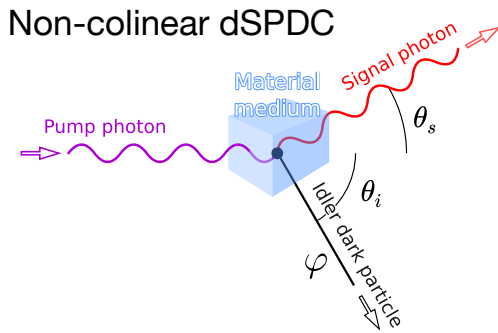
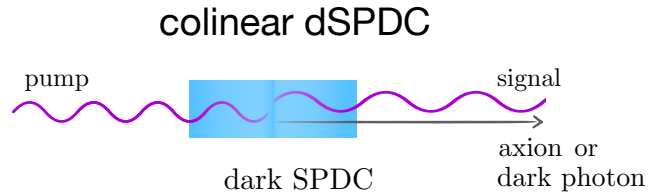
$$\omega_p = \omega_s + \omega_i \quad \text{and}$$



- Presence of idler is inferred. Might as well be invisible!

Nonlinear Optics with Dark States

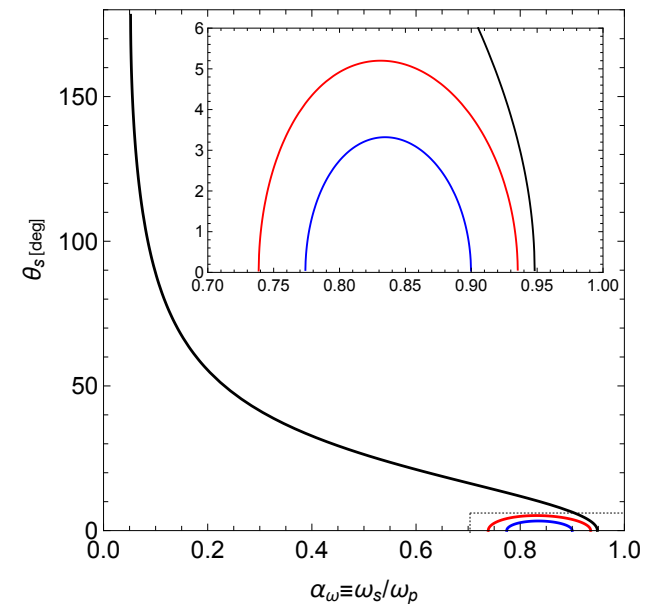
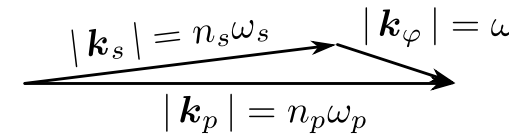
- Dark SPDC: Pump \rightarrow signal + axion or dark photon. Rate $\propto \epsilon^2$ (vs ϵ^4 for LSW)



Experimental design in progress with Niel Sinclair

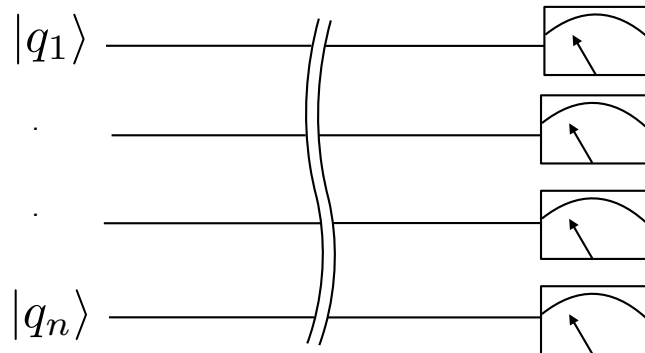
Note: the axion or dark photon have index of refraction of 1 (and a mass).

dSPDC has significantly different phase matching conditions. Different kinematics.



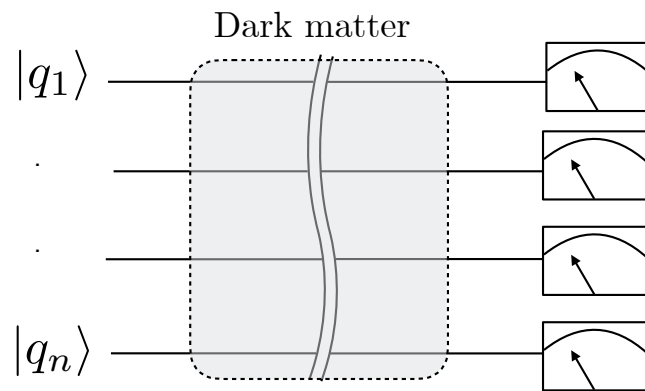
Quantum Advantage with Qubits

- Can we get quantum speed-up in DM detection? Consider n qubits interacting with DM:



Quantum Advantage with Qubits

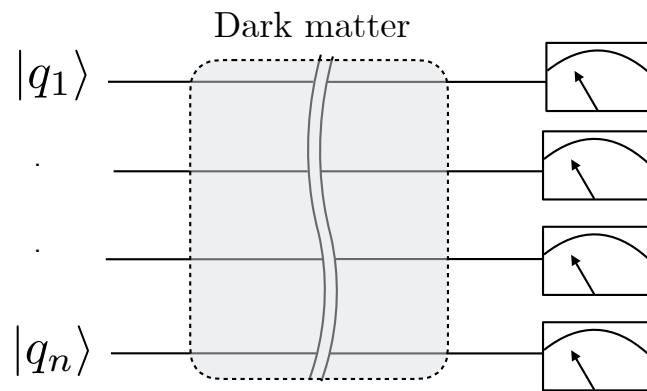
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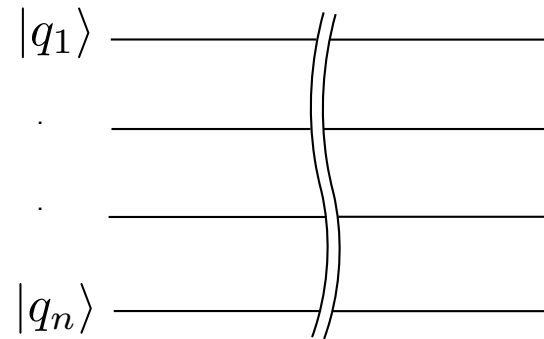
Rate scales as n .

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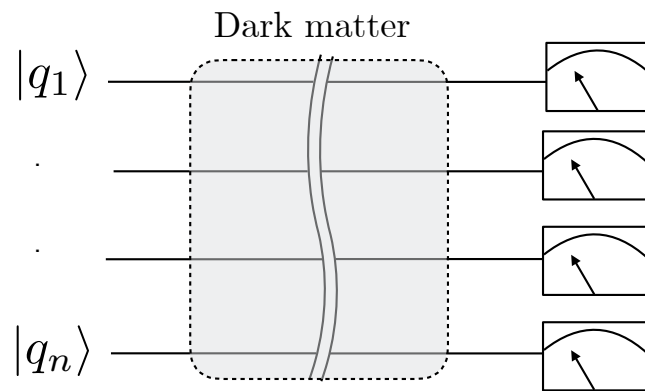


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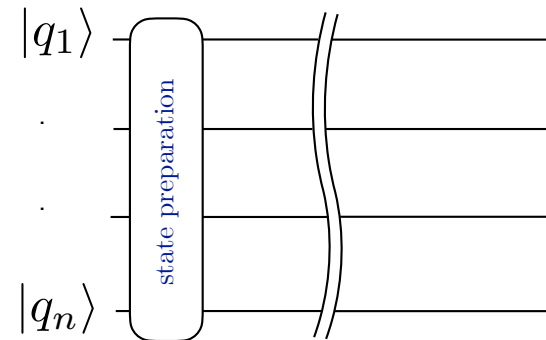


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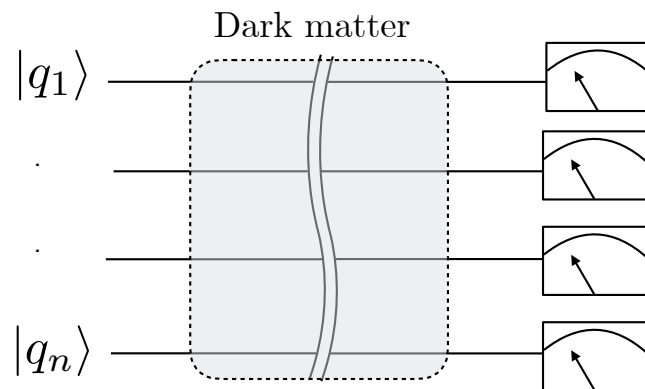


Rate scales as n .

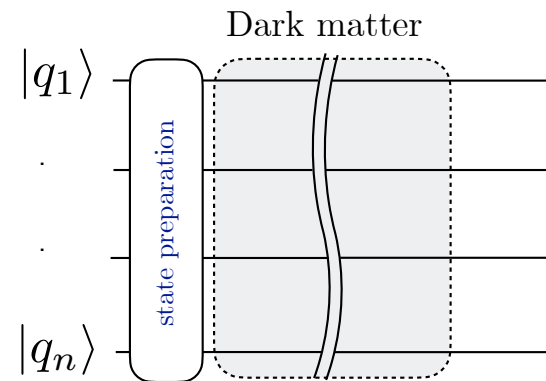


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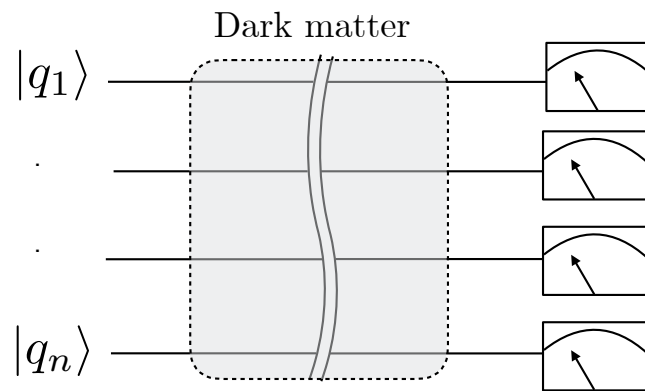


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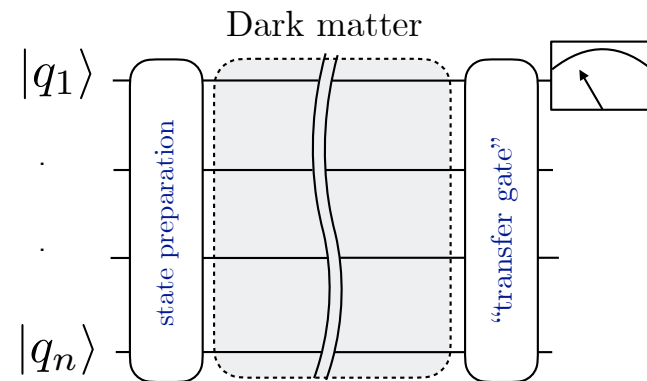


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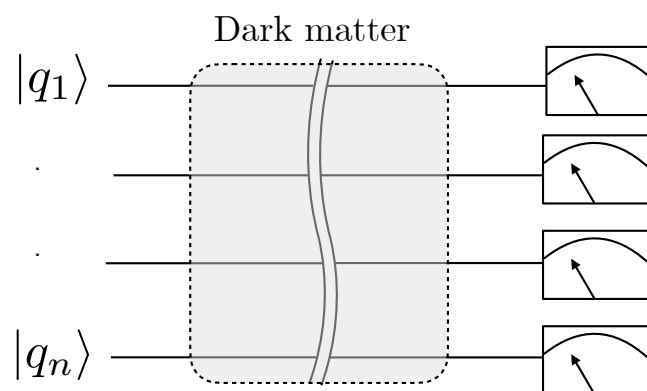


Rate scales as n .

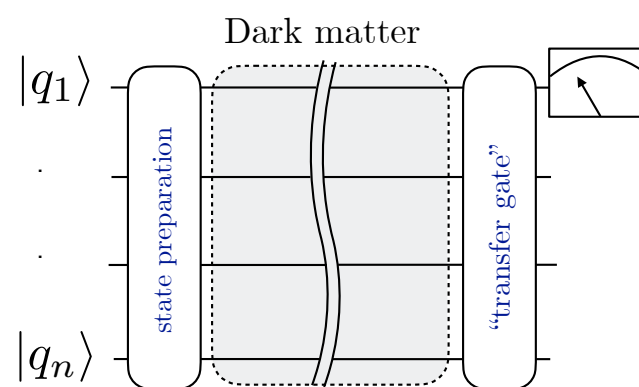


Quantum Advantage with Qubits

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Rate scales as n .



Which qubit interacted?

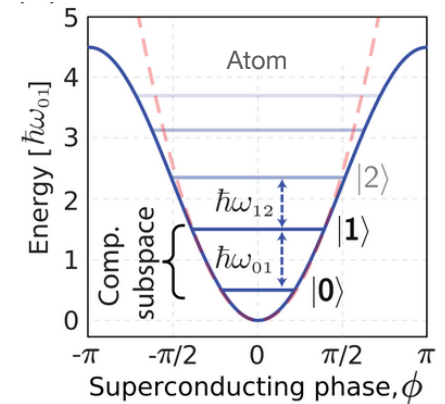
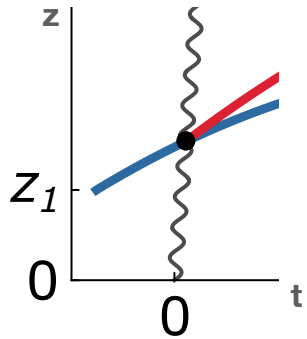
Amplitude scales as n .

Rate scales as n^2 !!

- A variety of algorithms for sensing (e.g. quantum phase estimation, work in progress).

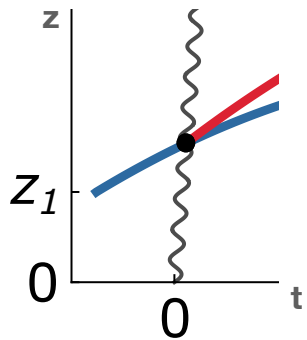
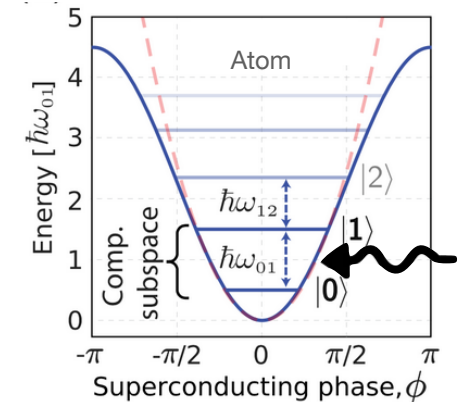
Quantum Control in Free Space

- Imagine we take a qubit in free space and induce a transition:
- We give the atom some energy, but also momentum!



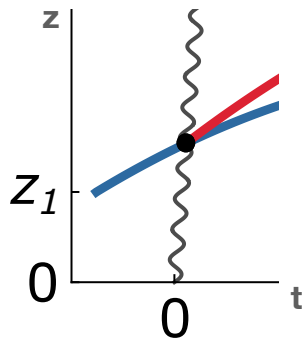
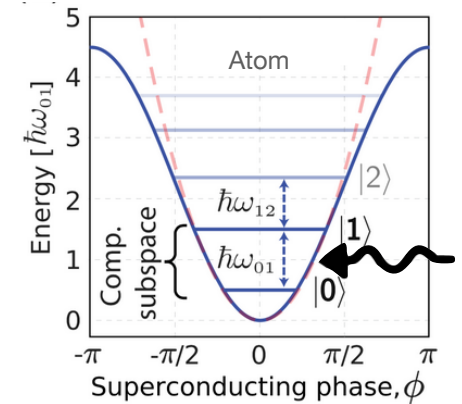
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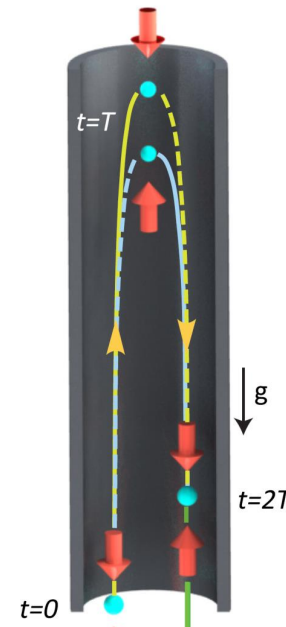


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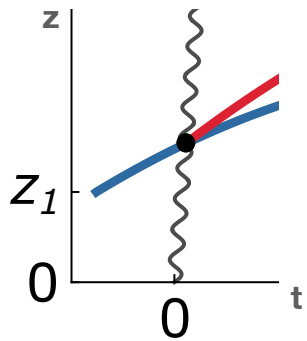
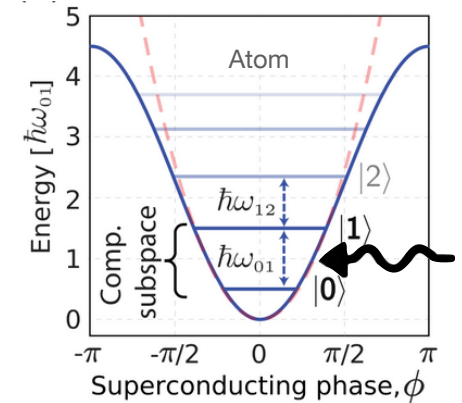


Wait ...



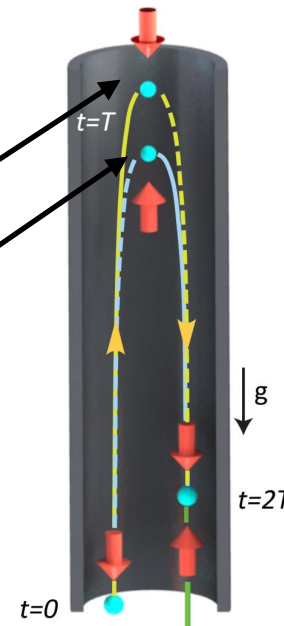
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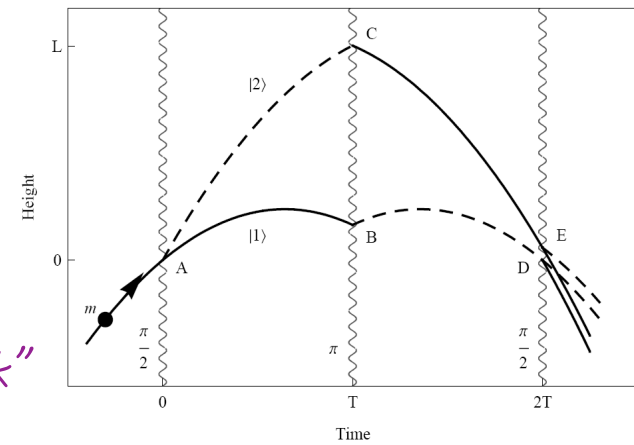
$$|\text{atom}\rangle = |\text{here}\rangle + |\text{there}\rangle$$



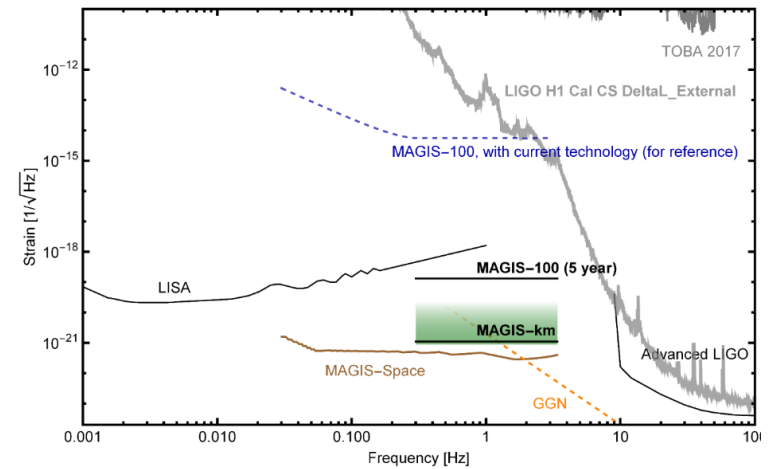
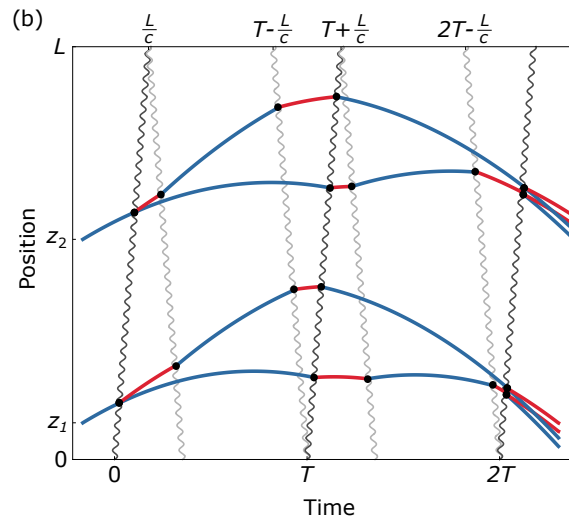
MAGIS 100

□ Now we can do cool things! Interferometry:

"A freely falling clock"



□ MAGIS 100:



MAGIS-100 is under construction in Fermila's MINOS shaft. :-)

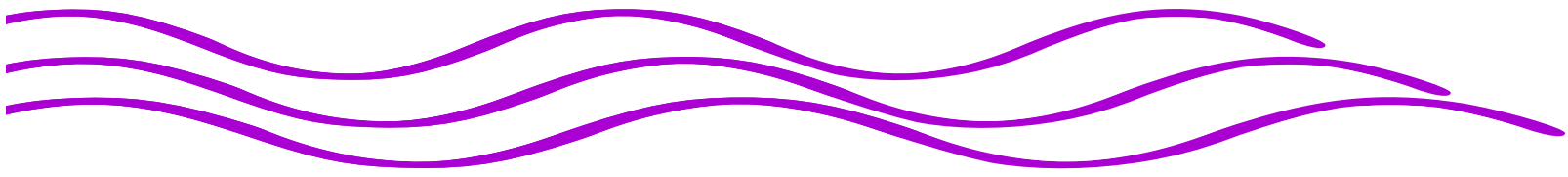
Summary

- Quantum Field theory extends into today's quantum devices!
- The interface of HEP and QIS is growing!
 - Quantum simulation of HEP (exciting, but not today's topic)
 - Quantum sensing
- Quantum sensors can probe new hypotheses in HEP

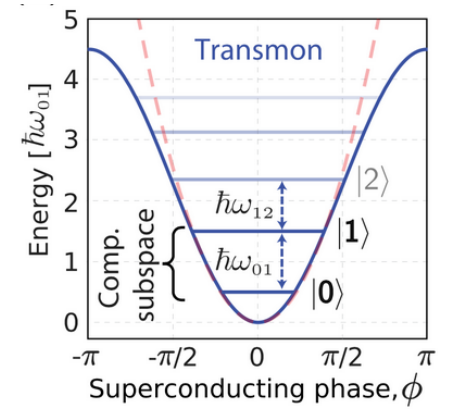


 **Fermilab**


SUPERCONDUCTING QUANTUM
MATERIALS & SYSTEMS CENTER



Deleted Scenes



Rates:

$$\Gamma_{\text{SPDC}} \sim \frac{P_p \chi_{\text{eff}}^{(2)2} \omega_s \omega_i L}{\pi n_p n_s n_i A_{\text{eff}}}$$

Motivates long crystals too.

Rates:

$$\Gamma_{\text{SPDC}} \sim \frac{P_p \chi_{\text{eff}}^{(2)2} \omega_s \omega_i L}{\pi n_p n_s n_i A_{\text{eff}}}$$

Motivates long crystals too.

$$\Gamma_{\text{dSPDC}}^{(A'_L)} \sim \epsilon^2 \frac{m_{A'}^2}{\omega_{A'}^2} \frac{P_p \chi_{A'_L}^{(2)2} \omega_s \omega_{A'} L}{n_p n_s A_{\text{eff}}}$$

$$\Gamma_{\text{dSPDC}}^{(\text{axion})} \sim \frac{P_p g_{a\gamma\gamma}^2 \omega_s L}{\omega_{\text{axion}} n_p n_s A_{\text{eff}}}$$

$$N_{\text{events}}^{(A'_L)} \sim 10^{21} \left(\epsilon^2 \frac{m_{A'}^2}{\omega_{A'}^2} \right) \left(\frac{P_p}{\text{Watt}} \right) \left(\frac{L}{\text{m}} \right) \left(\frac{t_{\text{int}}}{\text{year}} \right)$$

$$N_{\text{events}}^{(\text{axion})} \sim 40 \left(\frac{g_{a\gamma}}{10^{-6} \text{ GeV}^{-1}} \right)^2 \left(\frac{P_p}{\text{Watt}} \right) \left(\frac{L}{\text{m}} \right) \left(\frac{t_{\text{int}}}{\text{year}} \right)$$

	Dark Photon ($m_{A'} = 0.1 \text{ eV}$)	Axion-like particle ($m_a = 0.1 \text{ eV}$)
Current lab limit	$\epsilon < 3 \times 10^{-7}$	$g_{a\gamma} < 10^{-6} \text{ GeV}^{-1}$
Example dSPDC setup	$P_p = 1 \text{ W}$ $L = 1 \text{ cm}$ $\Gamma = 10/\text{day}$	$P_p = 1 \text{ kW}$ $L = 10 \text{ m}$ $\Gamma = 10/\text{day}$
Current Solar limit	$\epsilon < 10^{-10}$	$g_{a\gamma} < 10^{-10} \text{ GeV}^{-1}$
Example dSPDC setup	$P_p = 1 \text{ W}$ $L = 10 \text{ m}$ $\Gamma = 10/\text{year}$	$P_p = 100 \text{ kW}$ $L = 100 \text{ m}$ $\Gamma = 10/\text{year}$

Gravity Waves - A nonlinear extension

- A gravity wave also interacts w/ two photons

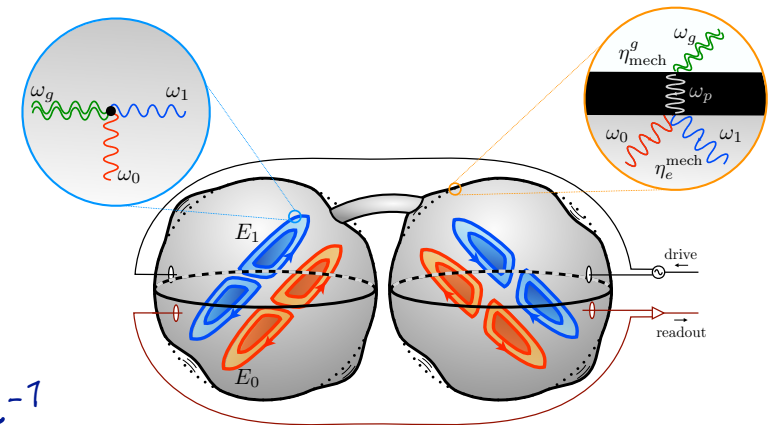
$$\mathcal{L} \supset F^{\mu\nu}F_{\mu\nu} \sim h(\mathbf{B} \cdot \mathbf{B} + \mathbf{E} \cdot \mathbf{E}) + \dots$$

- But often more important:

$$H = \hbar\omega(a^\dagger a + 1/2) \quad \text{with} \quad \omega \sim (1+h)L^{-1}$$

- Axion-like phenomenology:

- GW mixing w/ photons in background B field.
- GW can be absorbed by photon \rightarrow up conversion.

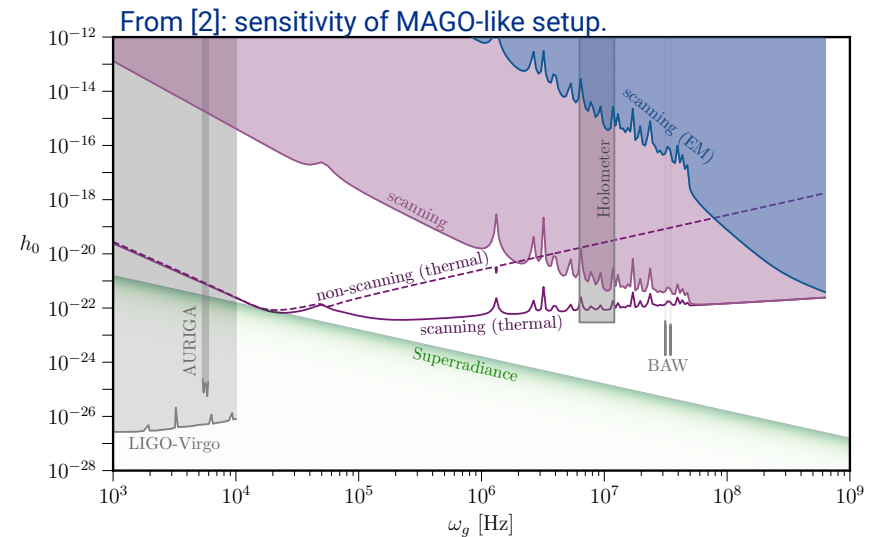
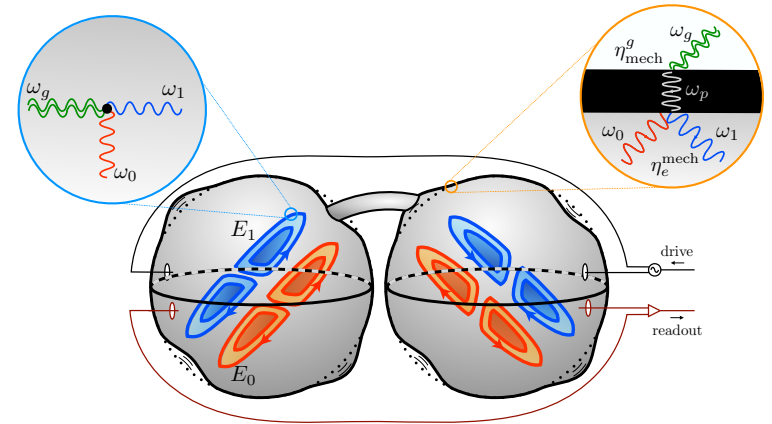


Gravitational waves

- Photon up-conversion due to GW.
- Current axion experiments have sensitivity to GHz Gravity waves [1].
- A dedicated cavity experiment, e.g. MAGO, has significant reach at MHz [2].
- MAGO traveled from INFN to DESY to Fermilab for testing
- **A Fermilab KEK collaboration to design new dedicated broadband cavity.**



MAGO (INFN)

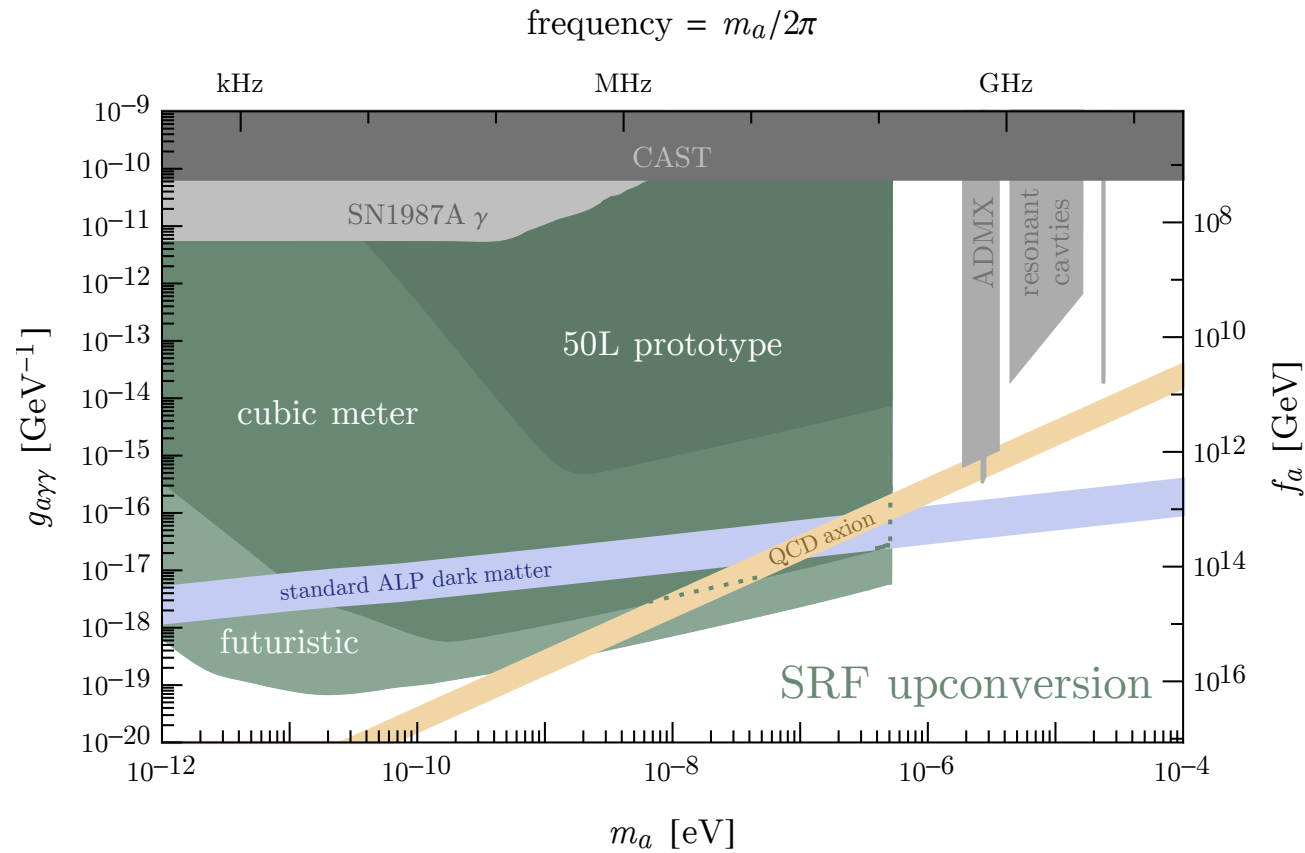


[1] *Phys.Rev.D* 105 (2022) 11, 116011

[2] *Phys.Rev.D* 108 (2023) 8, 084058

Multimode searches

Berlin, et al., JHEP, DOI:10.1007/JHEP07 (2020) 088
 Berlin, et al., arXiv:2203.12714, Snowmass WP (2022)
 Giaccone, et al., arXiv:2207.11346 (2022)

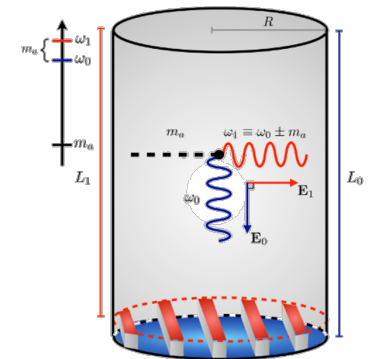


Snowmass name:

SRF-m³

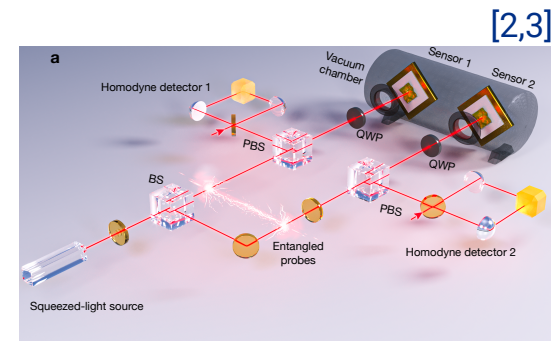
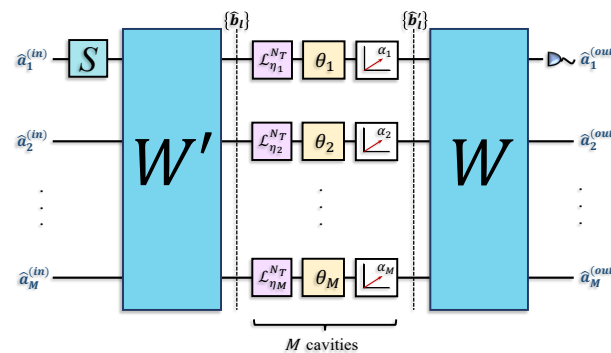
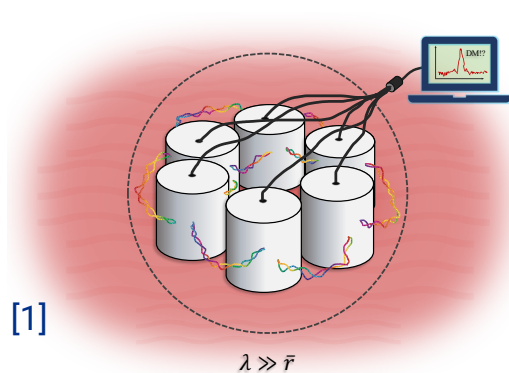
Asher's proposal:

SuperRAD



Dark Matter - Quantum Sensor Network

- A network of cavities can be used to enhance the sensitivity to dark matter.
- How should we distribute quantum resources in the network?
- **A distributed quantum sensing protocol for DM searches allows for enhanced scan rate for DM.**



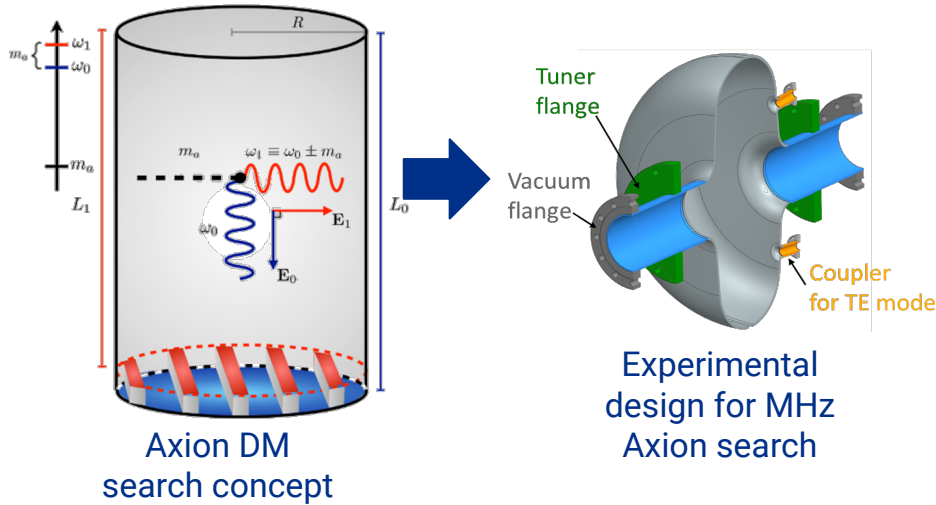
A new collaboration of HEP and QIS experts from across SQMS:

[1] *PRX Quantum* 3 (2022) 3, 030333

[2] 2210.07291 [quant-ph]

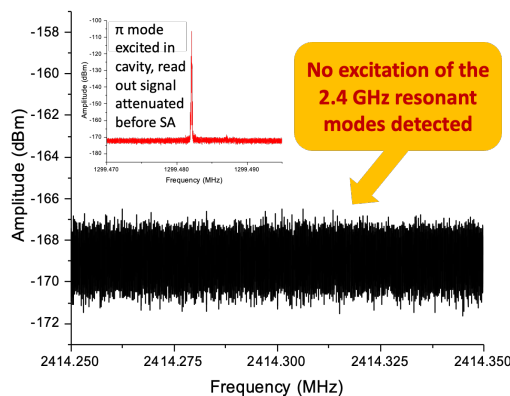
[3] 2210.16180 [quant-ph] (in review by Nature Photonics)

Multimode searches

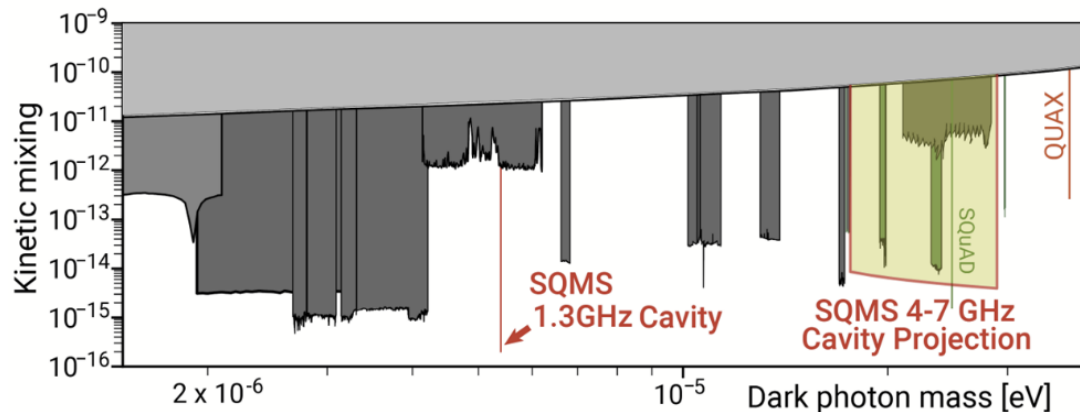
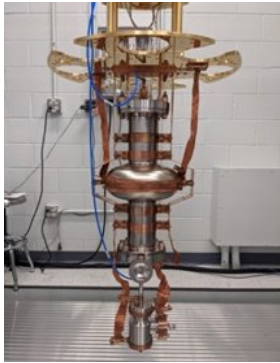


Bogorad, et al., PRL, DOI:10.1103/PhysRevLett.123.021801
 Berlin, et al., JHEP, DOI:10.1007/JHEP07 (2020) 088
 Gao & Harnik, JHEP, DOI:10.1007/JHEP07 (2021) 053
 Berlin, et al., arXiv:2203.12714, Snowmass WP (2022)
 Sauls, PTEP, DOI:10.1093/ptep/ptac034 (2022)
 Giaccone, et al., arXiv:2207.11346 (2022)

- **Axion DM** search based on the heterodyne detection scheme: cavity design is finalized, contract for cavity fabrication placed (cavity arrival: Fall 2023)
- In preparation for search:
 - Working on RF experimental set up and read out system
 - Addressing experimental challenges such as passive dampening of vibrations in LHe facility
 - Multimode feasibility study

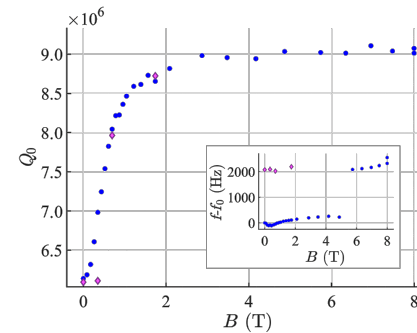
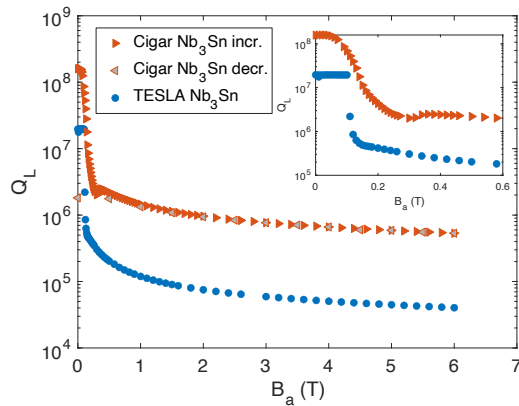
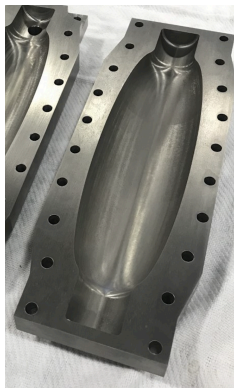


Ultrahigh Q for Dark Matter



Cervantes et al.,
arxiv:2208.03183, in
review in Phys. Rev. Lett.

No B-Field:
 $Q > 10^{10}$



With B-Field:
 $Q \geq 10^{5-7}$

Hybrid copper-dielectric cavity (INFN): Di Vora et al.,
PhysRevApplied.17.054013

Superconducting Nb₃Sn cavity (FNAL): Posen et al.,
arxiv:22014.10733, in review in Phys Rev Applied

Single Particle Qubit

- The most precise theory-experiment comparison in physics:

Electron magnetic moment $(g-2)_e$:

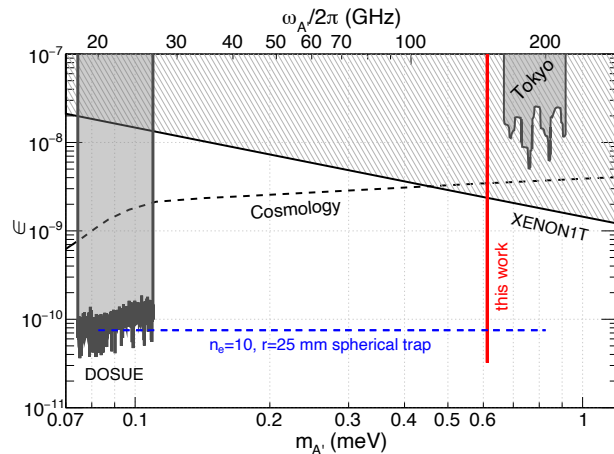
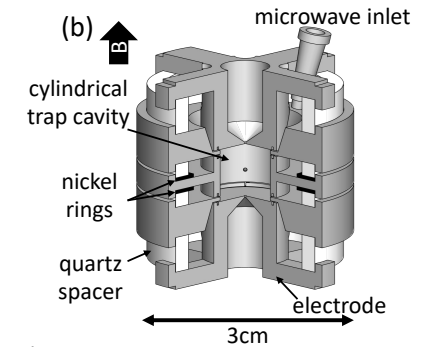
The quantum state of a single electron in a trap is monitored via a QND measurement.

$$-\frac{\mu}{\mu_B} = \frac{g}{2} = 1.001\,159\,652\,180\,59(13) \quad [0.13 \text{ ppt}]$$

[*Phys. Rev. Lett.* 130, 071801 \(2023\)](#)

Editors choice!

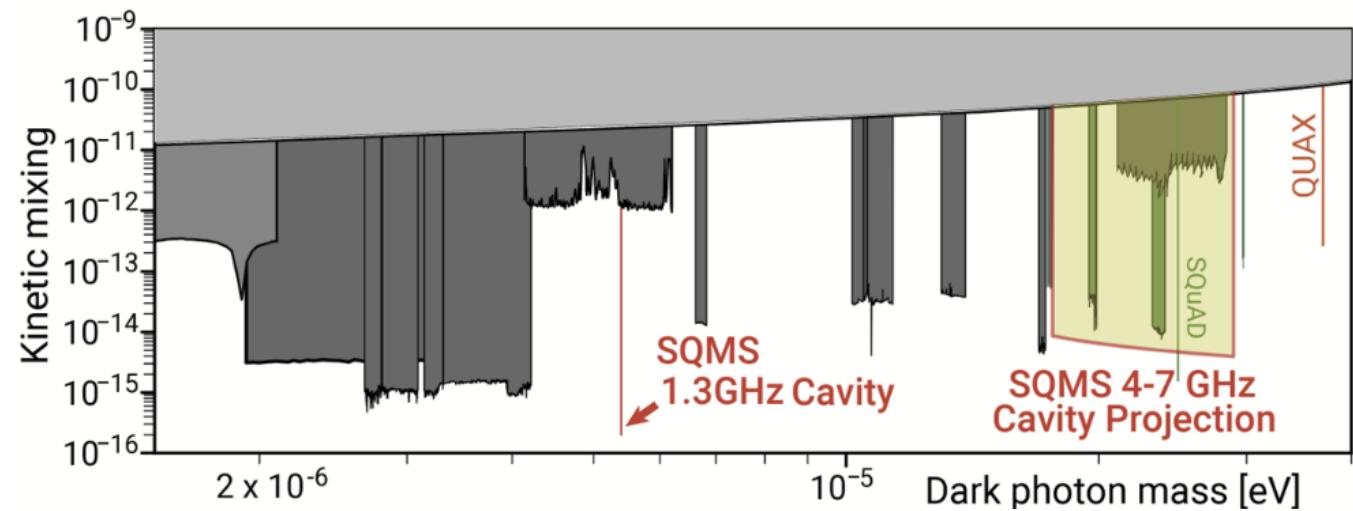
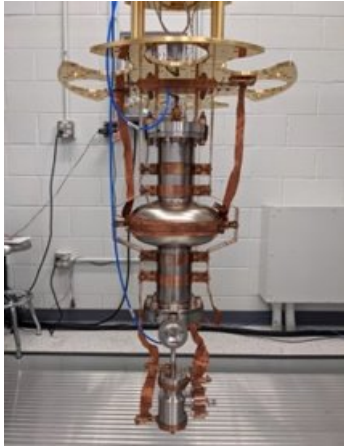
- SQMS joined the effort, contributed to understanding loss sources.



- SQMS bonus:** We also found that a single-electron qubit is a sensitive DM search in a challenging frequency range!
- Theory + proof-of-concept!

Phys.Rev.Lett. 129 (2022) 26, 261801
(a new NU-Stanford-Fermilab collaboration)

Ultrahigh Q for Dark Matter



Cervantes et al., arxiv:2208.03183, in review in Phys. Rev. Lett.

DPDM search with 1.3 GHz cavity with $Q_L \approx 10^{10}$.

Deepest exclusion to wavelike DPDM by an order of magnitude.

Next steps:

- Tunable DPDM search from 4-7 GHz. (“low hanging fruit”)
- Implement photon counting to subvert SQL noise limit.