

String hadronization in PYTHIA8:

from e^+e^- to pp and beyond

Javira Altmann - Monash University, visiting University of Oxford

- **Confinement** in High-Energy Collisions
- String **Hadronisation** → Modelling in PYTHIA (QCD Colour Reconnections)
- String **Junctions**
- Strings from vacuum → small systems → heavy ion collisions

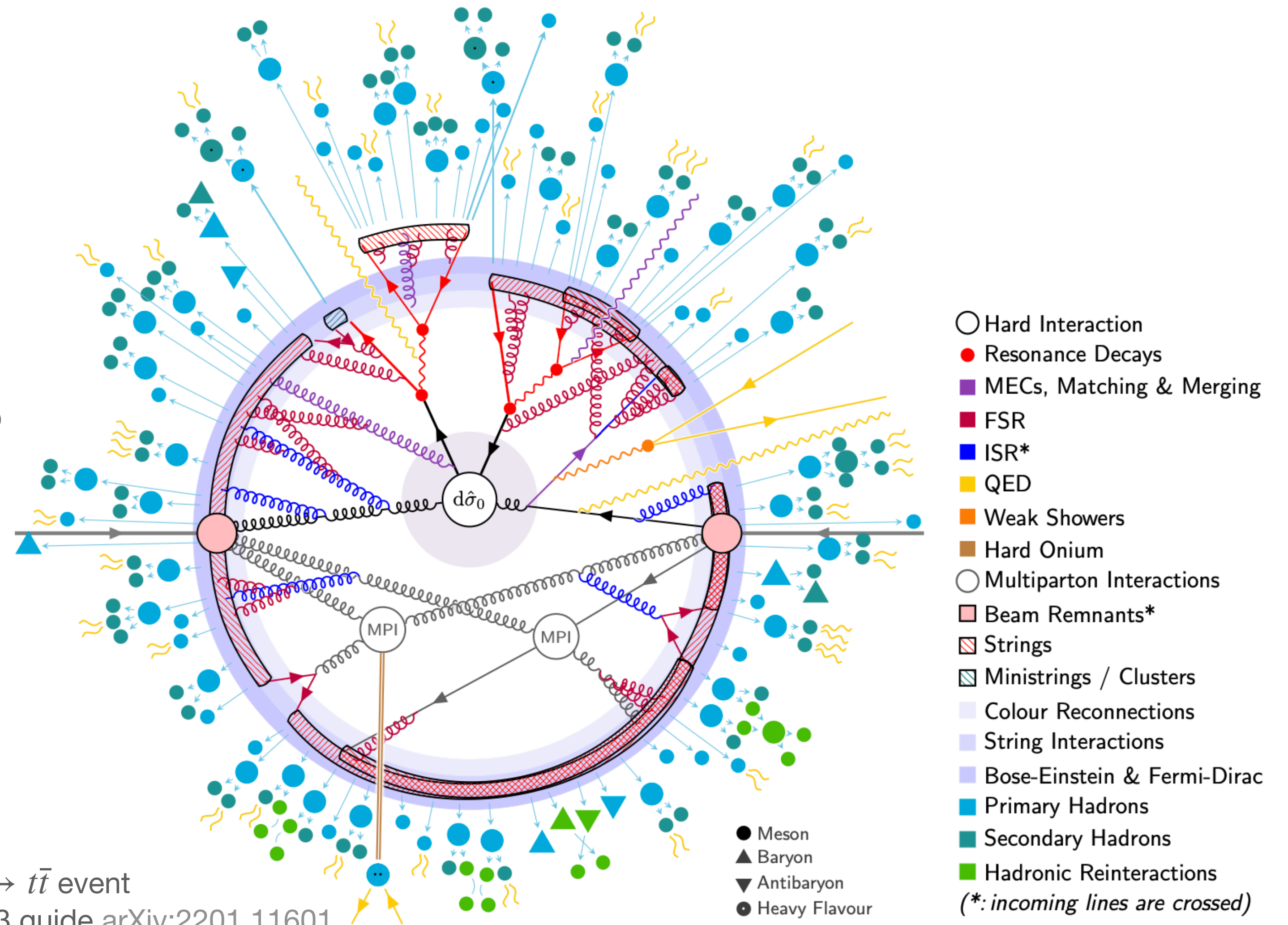
Confinement in high energy collisions

Consider “hard” processes with large momentum transfers $Q^2 \gg \Lambda_{QCD}^2$

At wavelengths $\sim r_{proton} \sim 1/\Lambda_{QCD}$

Need a dynamical process to ensure partons (**quarks and gluons**) become **confined** within hadrons

i.e. **non-perturbative parton \rightarrow hadron map**



Example of $pp \rightarrow t\bar{t}$ event
From PYTHIA 8.3 guide arXiv:2201.11601

Colour neutralisation

Require colour neutralisation:

- The point of confinement is that partons are **coloured** → a physical model needs two or more partons to create **colour neutral** objects

What does this **confinement field** look like?

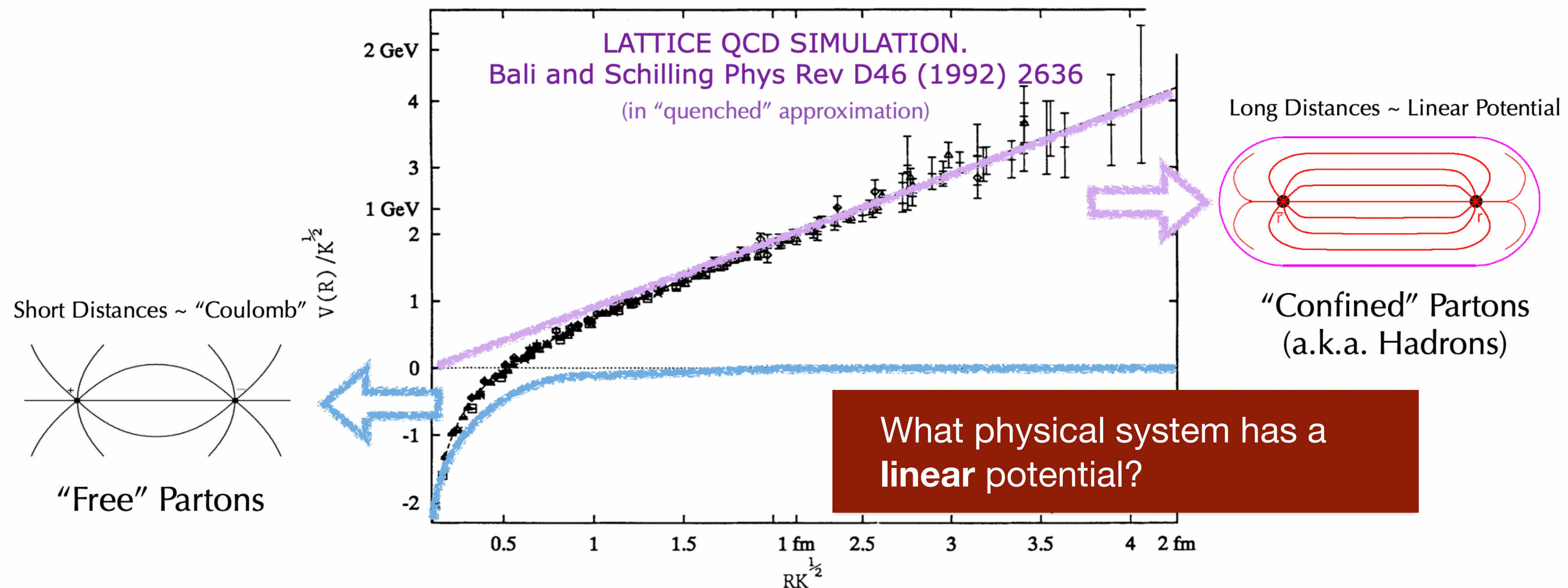
Colour neutralisation

Require colour neutralisation:

- The point of confinement is that partons are **coloured** → a physical model needs two or more partons to create **colour neutral** objects

Lattice QCD **“Cornell potential”** $V(r) = -\frac{a}{r} + \kappa r$ with $\kappa \sim 1 \text{ GeV/fm}$

shows us the potential energy of a colour singlet $q\bar{q}$ at separation distance r



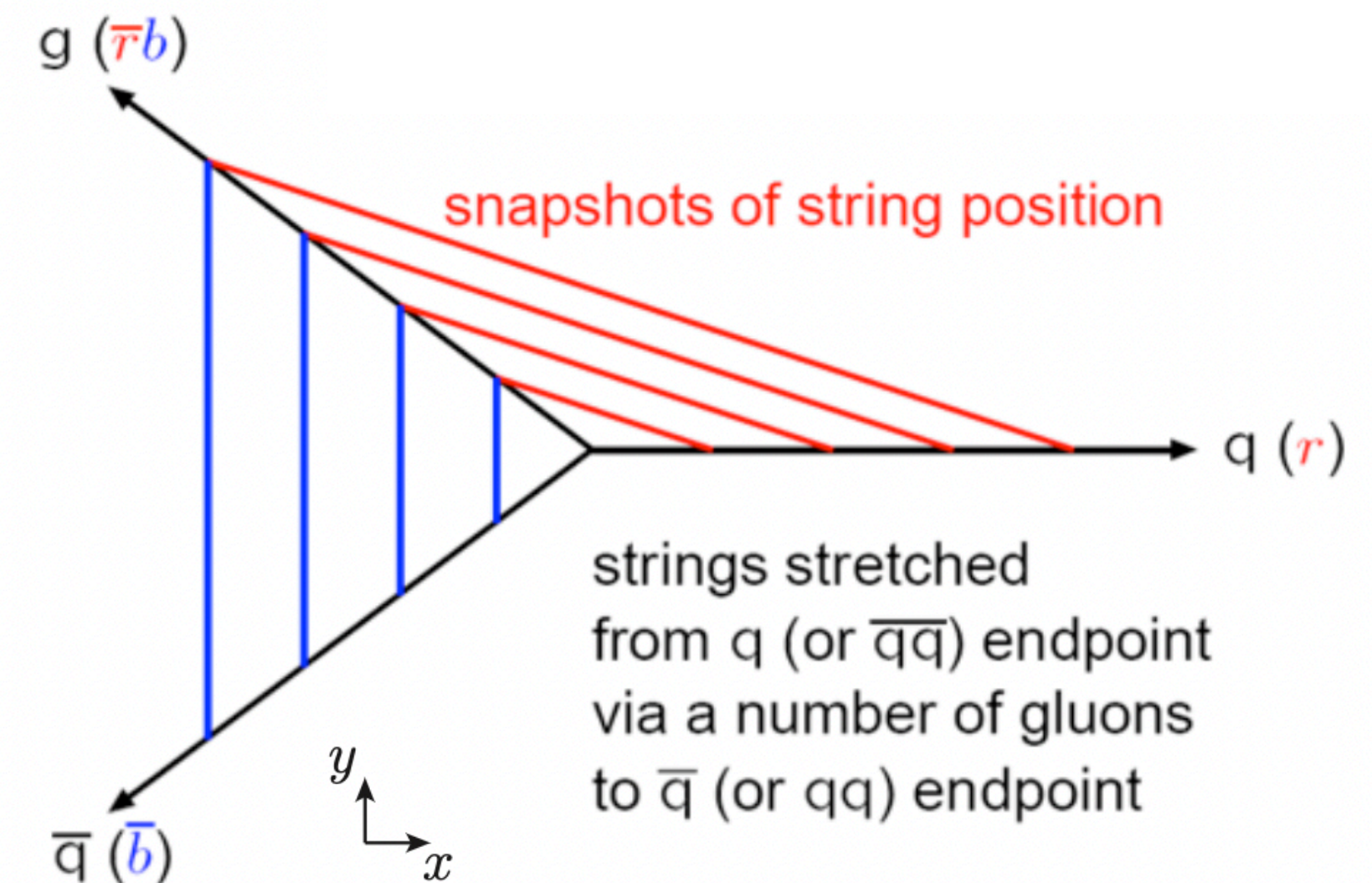
Lund String Model

Lund String Model:

Model the **confining field** between colour charges as a **string**

Collapse the colour field into a **narrow flux tube** (relativistic 1+1 dimensional world sheet) with uniform energy density

$$\kappa \sim 1 \text{ GeV/fm}$$



Example of a “dipole” string

Lund String Model

Lund String Model:

Model the **confining field** between colour charges as a **string**

Collapse the colour field into a **narrow flux tube** (relativistic 1+1 dimensional world sheet) with uniform energy density

$$\kappa \sim 1 \text{ GeV/fm}$$

Quarks / antiquarks

(anti)triplet \rightarrow carry (anti)**colour**

\rightarrow connected via a string to an anticolour charge

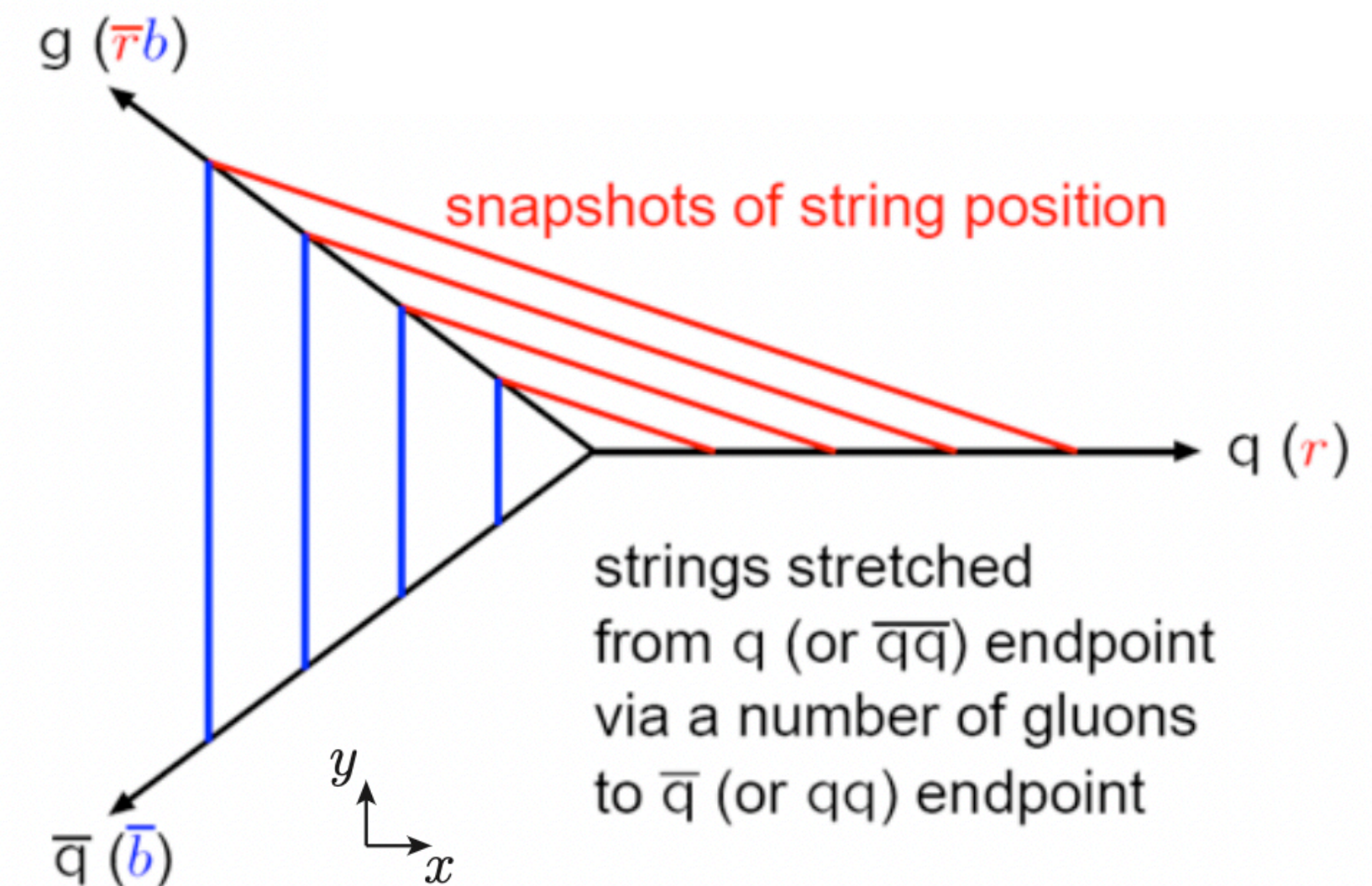
\rightarrow string **endpoints**

Gluons

Octet \rightarrow carry a **colour** and an **anticolour**

\rightarrow connected via a string to both a colour and an anticolour charge

\rightarrow transverse excitations on the **string** (“kinks”)



Example of a “dipole” string

Lund String Model

Lund String Model:

Model the **confining field** between colour charges as a **string**

Collapse the colour field into a **narrow flux tube** sheet) with uniform energy density

$$\kappa \sim 1 \text{ GeV/fm}$$

Quarks / antiquarks

(anti)triplet \rightarrow carry (anti)**colour**

\rightarrow connected via a string to an anticolour charge

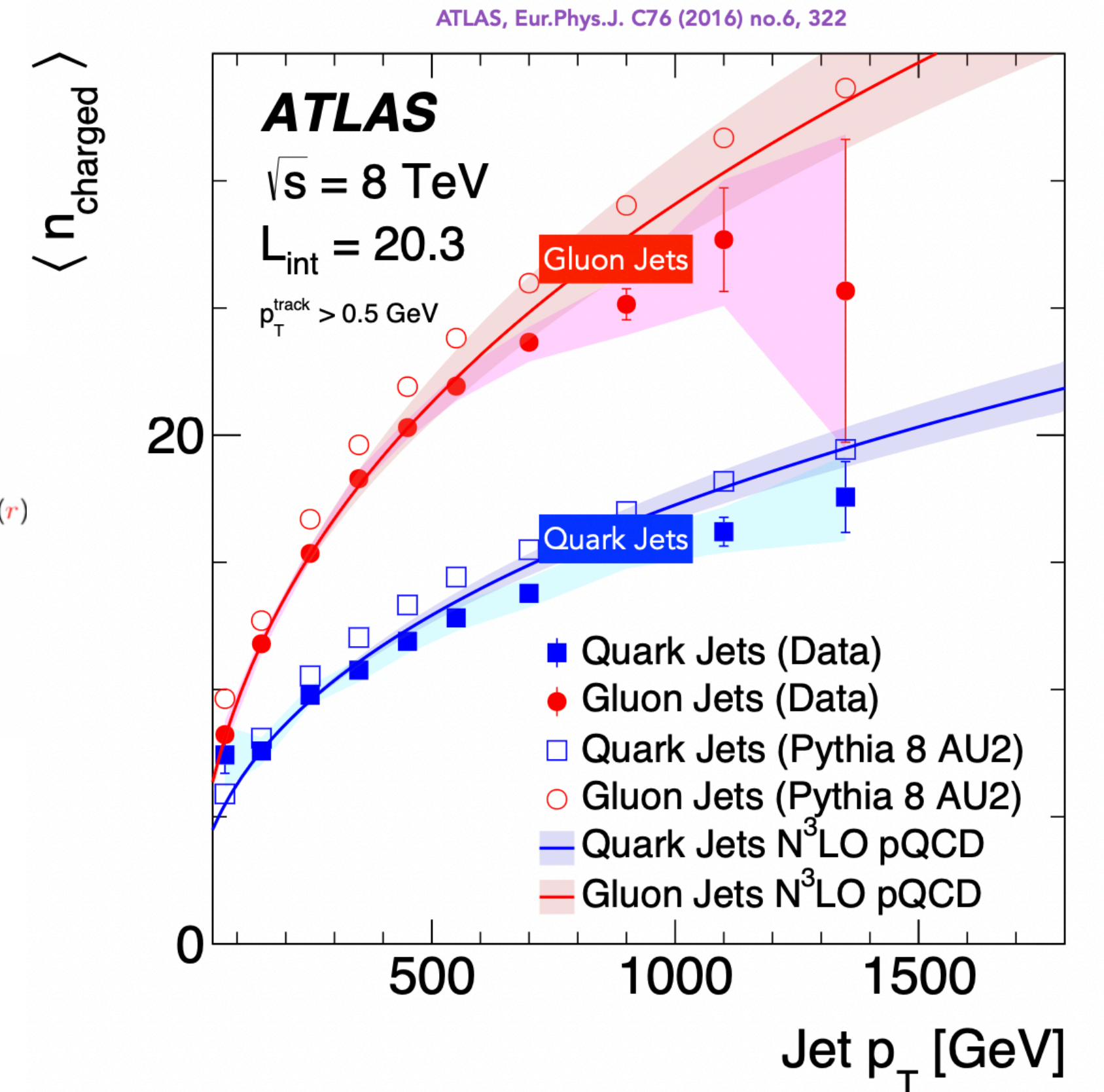
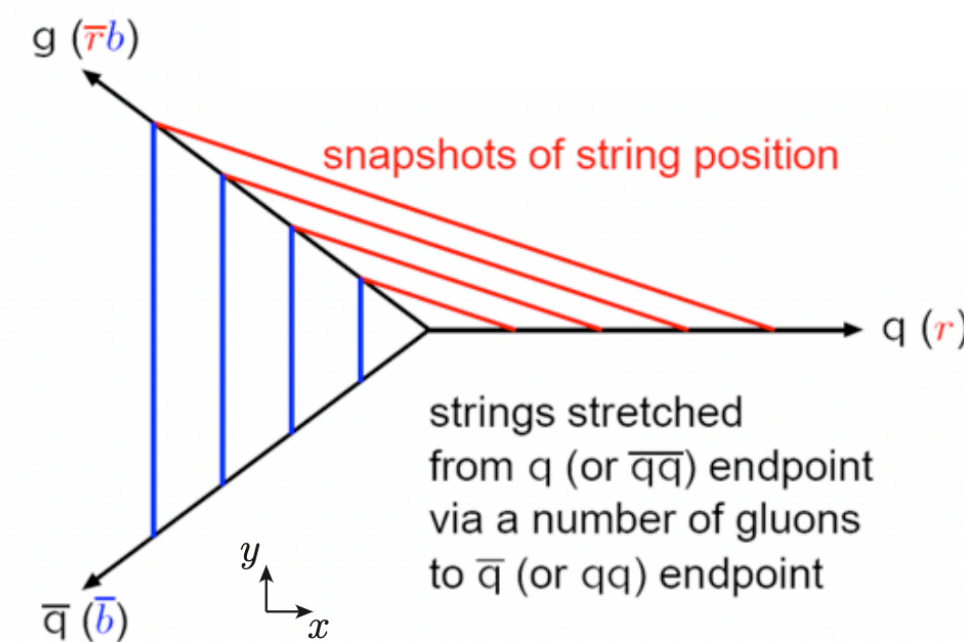
\rightarrow string **endpoints**

Gluons

Octet \rightarrow carry a **colour** and an **anticolour**

\rightarrow connected via a string to both a colour and an anticolour charge

\rightarrow transverse excitations on the **string ("kinks")**



Signatures of gluon-kinks have been seen
 Factor ~ 2 more particles in gluon jets

Lund String Model

Lund String Model:

Model the **confining field** between colour charges as a **string**

Collapse the colour field into a **narrow flux tube** (sheet) with uniform energy density

$$\kappa \sim 1 \text{ GeV/fm}$$



Quarks / antiquarks

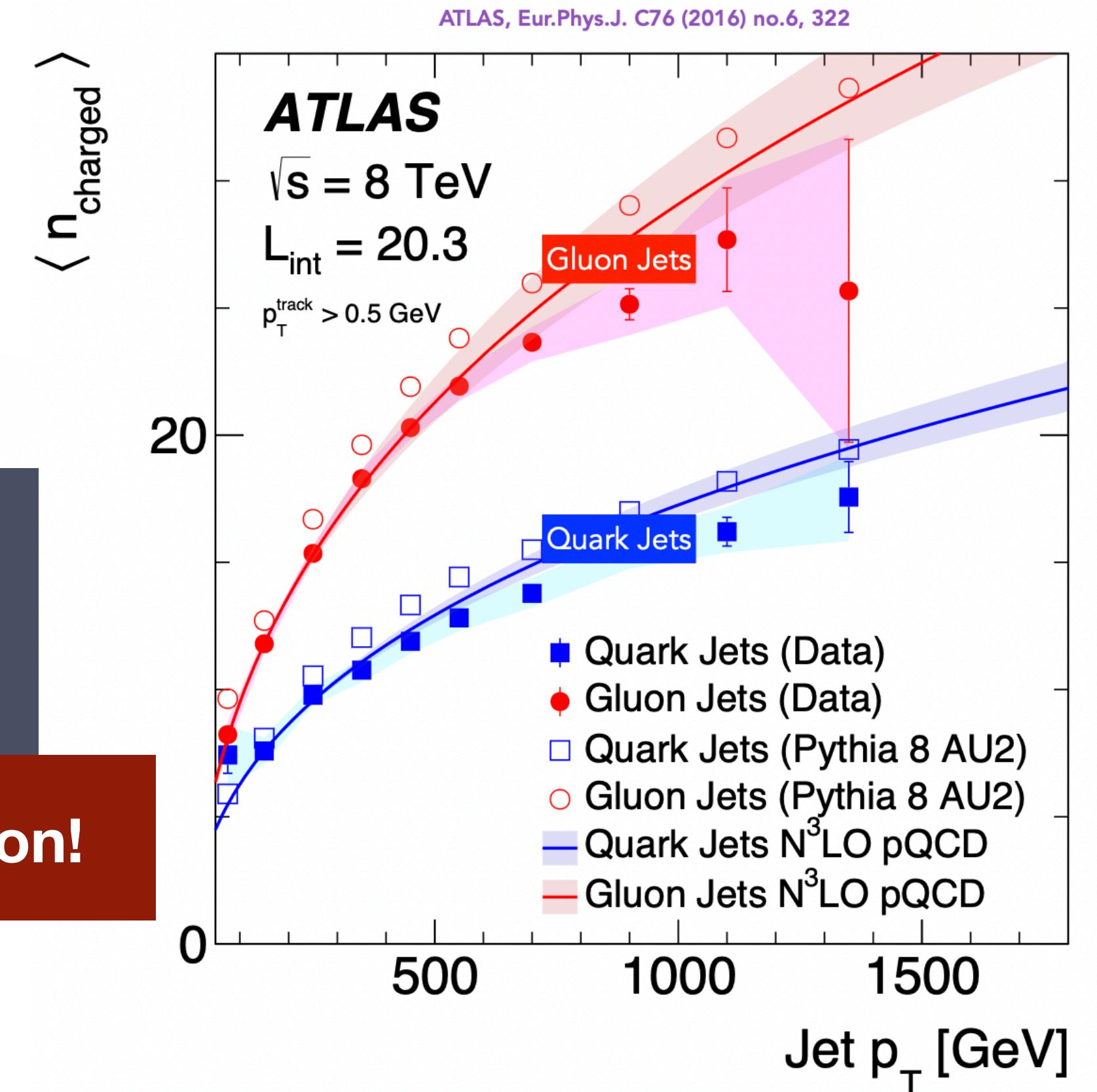
- (anti)triplet → carry (anti)colour
- connected via a string
- string **endpoints**

How does this map partons onto hadrons in high-energy collisions?

String fragmentation!

Gluons

- Octet → carry a **colour** and an **anticolour**
- connected via a string to both a colour and an anticolour charge
- transverse excitations on the **string ("kinks")**



Signatures of gluon-kinks have been seen
 Factor ~ 2 more particles in gluon jets

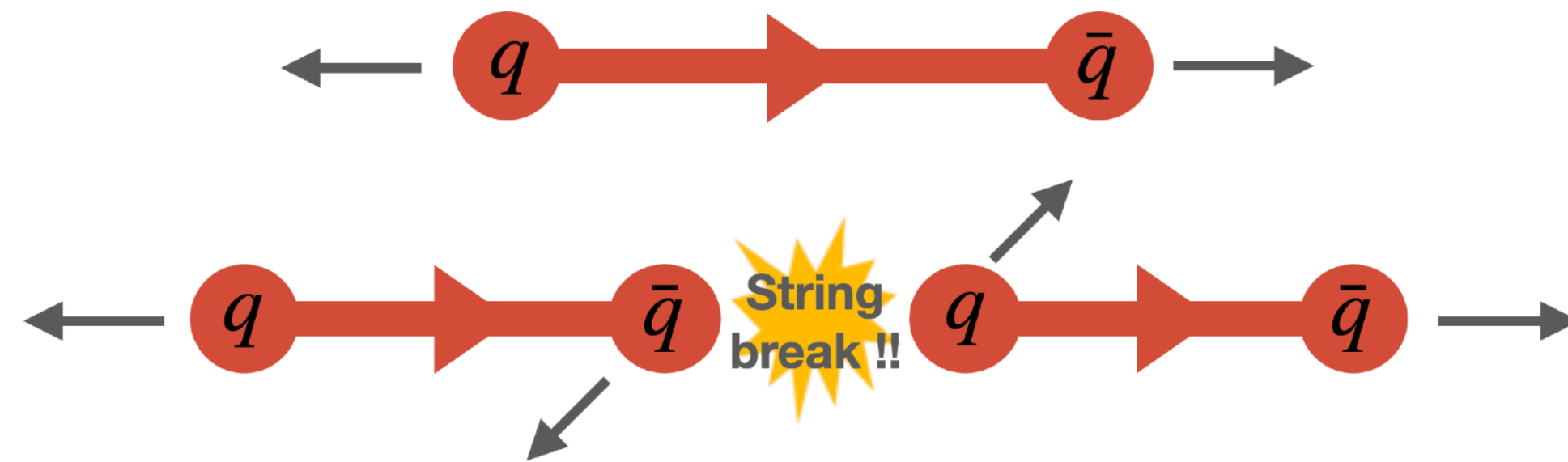
Partons → Hadrons

Hadronisation:

Partons move apart and stretch the string → **string breaks**

These happen at **non-perturbative** scales, can't use $P_{g \rightarrow q\bar{q}}(z)$

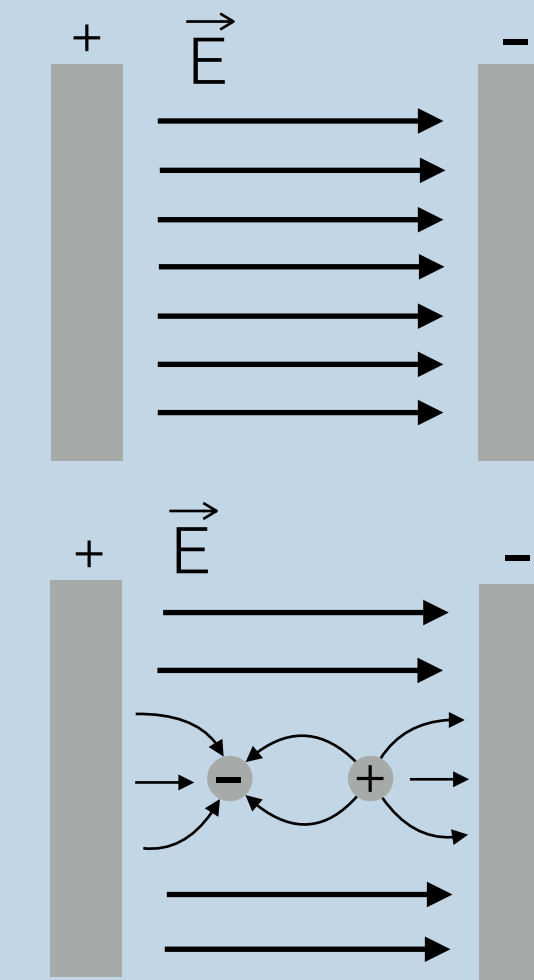
Instead use the **Schwinger mechanism**



Schwinger → Gaussian p_{\perp} spectrum and heavy flavour suppression $\text{Prob}(u:d:s) \approx 1 : 1 : 0.2$

Heavy quarks are only produced from hard processes
→ must be **string endpoints**

Schwinger mechanism QED



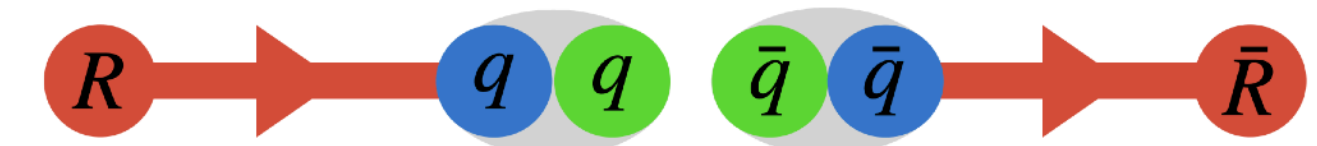
Non-perturbative creation of e^+e^- pairs in a string electric field

Probability from tunnelling factor

$$\mathcal{P} \propto \exp\left(\frac{-m^2 - p_{\perp}^2}{\kappa/\pi}\right)$$

Gaussian suppression of high $m_{\perp} = \sqrt{m_q^2 + p_{\perp}^2}$

Baryons formed from beam remnants or **diquark-antidiquark** pair creation



Partons \rightarrow Hadrons

Hadronisation:

Schwinger \rightarrow Gaussian p_{\perp} spectrum and heavy flavour suppression $\text{Prob}(u:d:s) \approx 1 : 1 : 0.2$

String breaks are **causally disconnected**

\rightarrow can fragment off hadrons from either end of the string

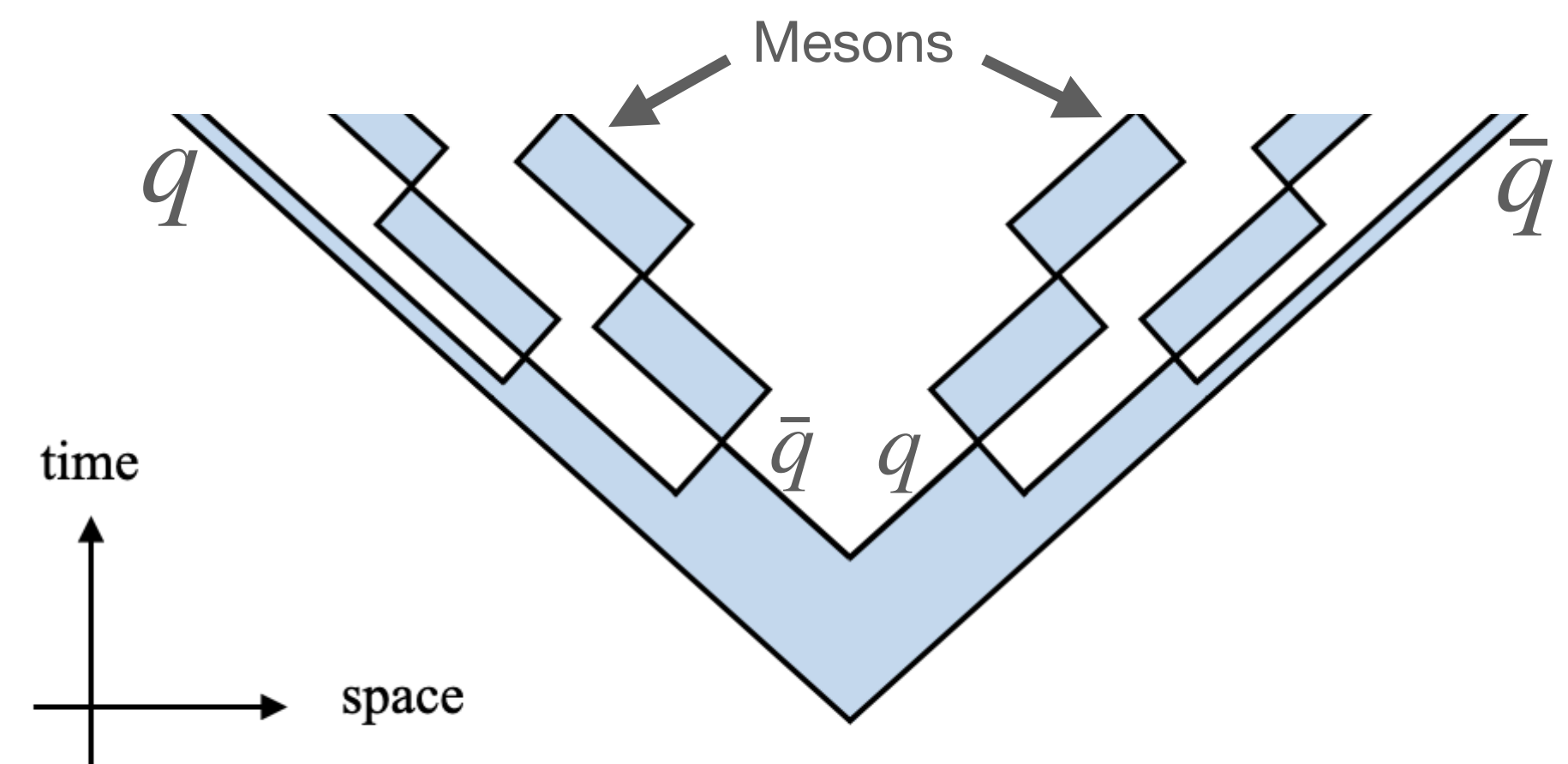
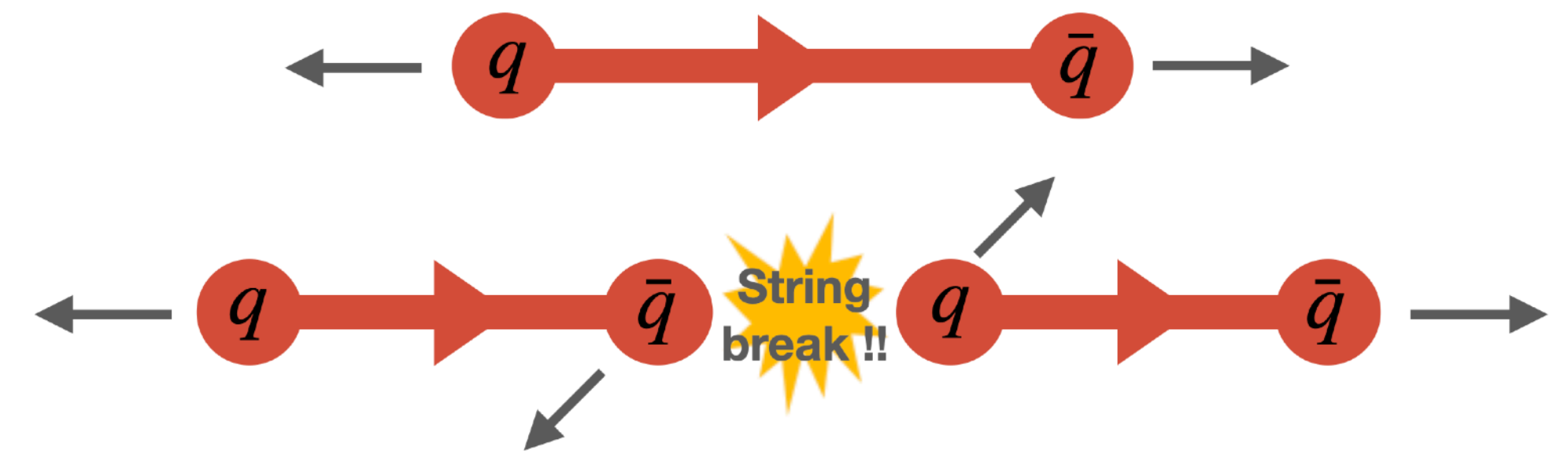
Probability distribution for the **fraction of quark**

momenta, z , the hadron will take is parametrised by the

Lund Symmetric Fragmentation Function

$$f(z) \propto \frac{1}{z} (1-z)^a \exp\left(\frac{-b(m_h^2 + p_{\perp h}^2)}{z}\right)$$

Free tuneable parameters a and b



Partons → Hadrons

Hadronisation:

Schwinger → Gaussian p_{\perp} spectrum and heavy flavour suppression $\text{Prob}(u:d:s) \approx 1 : 1 : 0.2$

String breaks are **causally disconnected**

→ can fragment off the top of a string with a different flavour

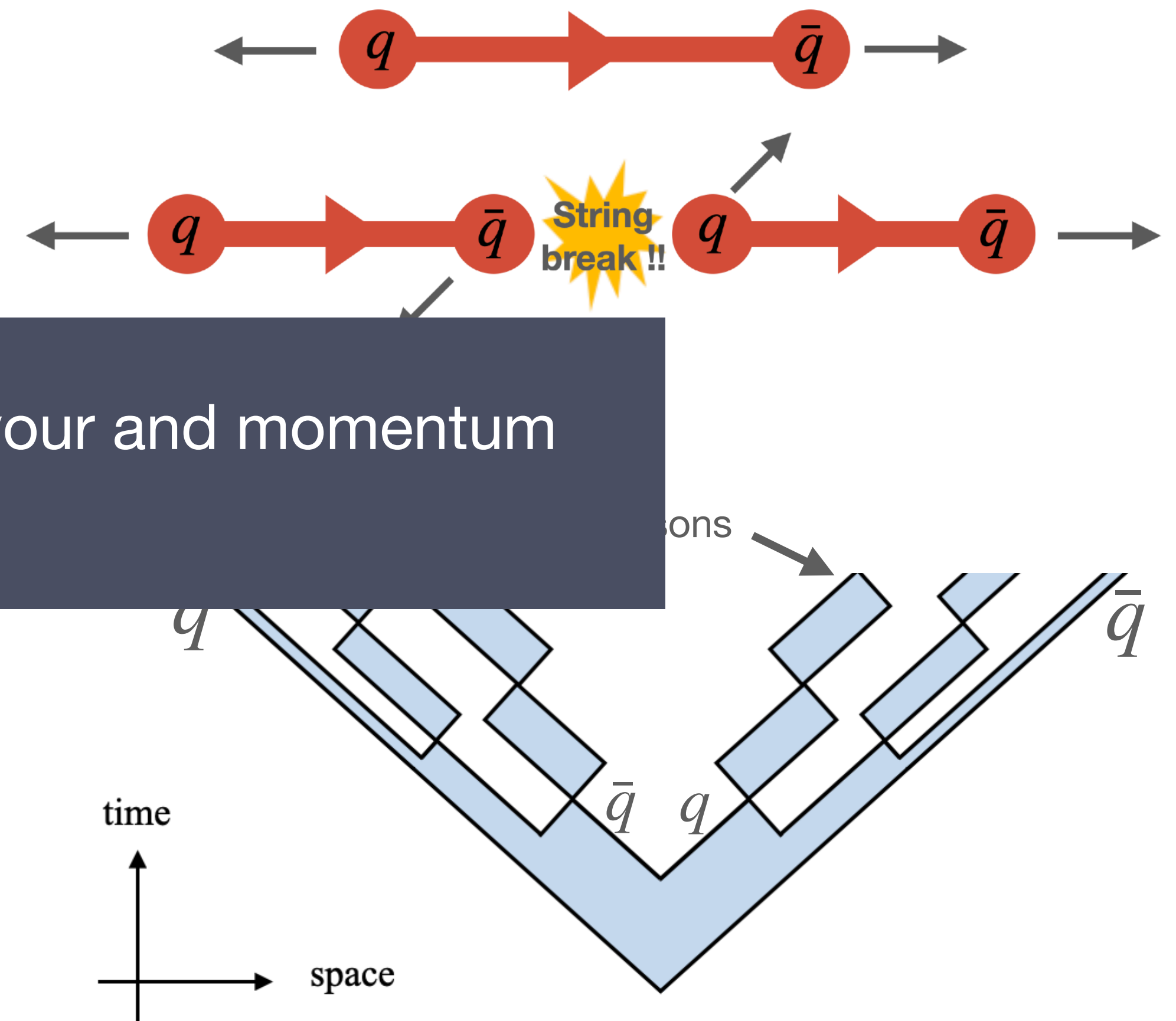
Probability of
momenta, z

So far we have notion of hadron flavour and momentum
What about colour?

Lund Symmetric Fragmentation Function

$$f(z) \propto \frac{1}{z} (1-z)^a \exp\left(\frac{-b(m_h^2 + p_{\perp h}^2)}{z}\right)$$

Free tuneable parameters a and b



Modelling Colour

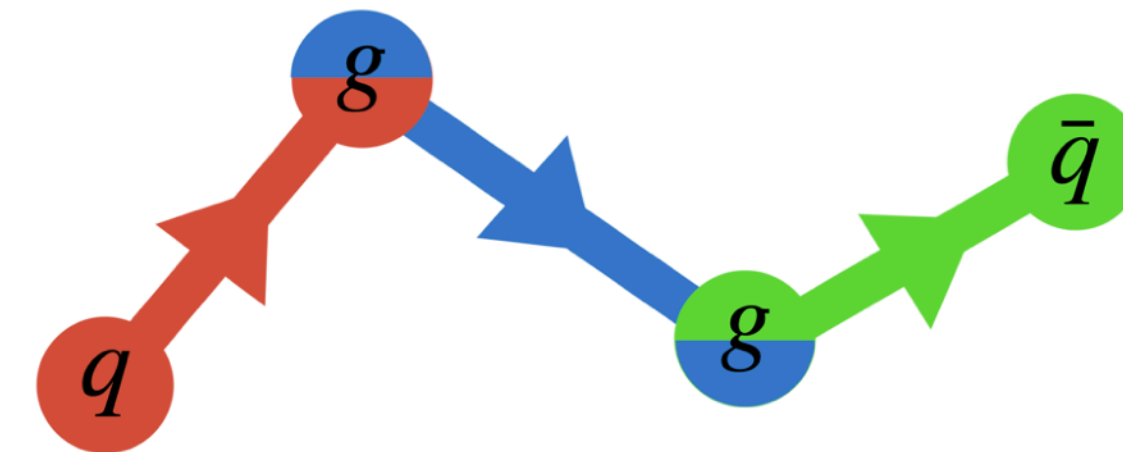
Leading Colour limit:

Starting point for Monte Carlo event generators $N_C \rightarrow \infty$

- Each **colour is unique** \rightarrow only one way to make colour singlets
- Only **dipole** strings
- Used by PYTHIA in the default (Monash 2013) tune

In e^+e^- collisions :

- Corrections suppressed by $1/N_C^2 \sim 10\%$
- Not much overlap in phase space



e.g. a dipole string configuration which make use of the **colour-anticolour** singlet state

Modelling Colour

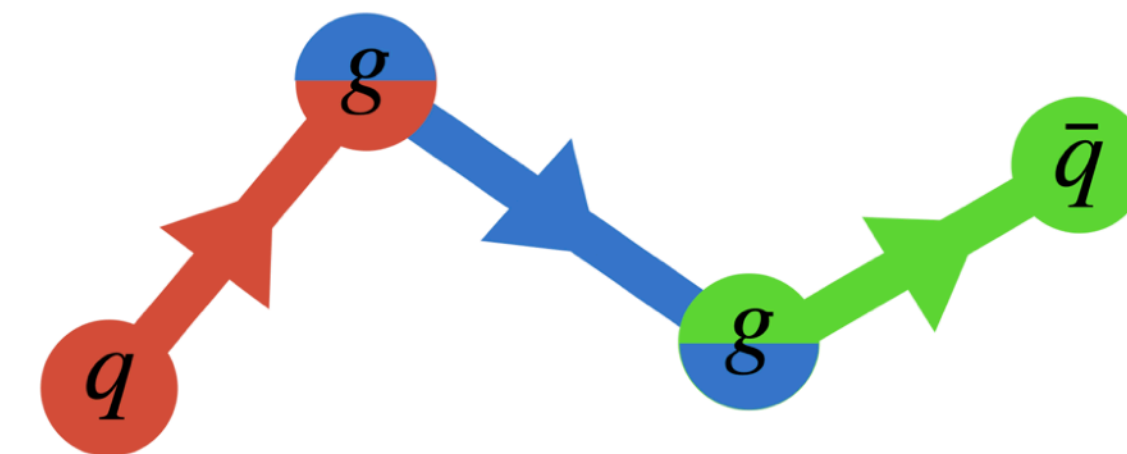
Leading Colour limit:

Starting point for Monte Carlo event generators $N_C \rightarrow \infty$

- Each **colour is unique** \rightarrow only one way to make colour singlets
- Only **dipole** strings
- Used by PYTHIA in the default (Monash 2013) tune

In e^+e^- collisions :

- Corrections suppressed by $1/N_C^2 \sim 10\%$
- Not much overlap in phase space



e.g. a dipole string configuration which make use of the **colour-anticolour** singlet state

But high-energy pp collisions involve **very many** coloured partons with significant **phase space overlaps**

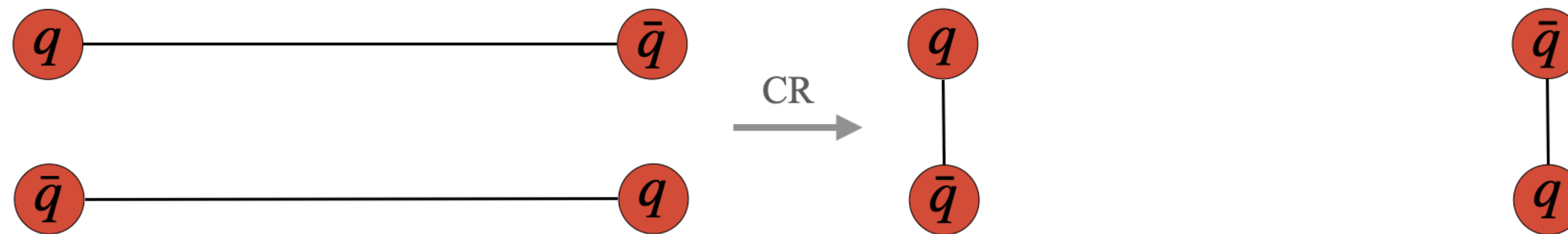
QCD Colour Reconnection (CR) model

QCD Colour Reconnections

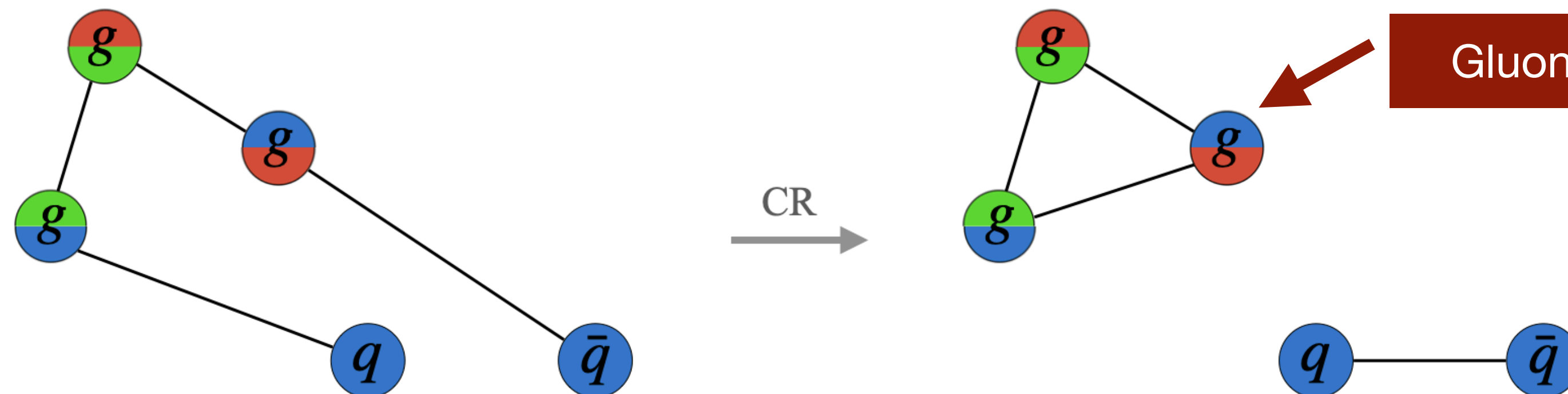
Stochastically restores colour-space ambiguities according to **SU(3) algebra**

➤ Allows for reconnections to **minimise string lengths**

Colour - anticolour singlet state



Dipole reconnection



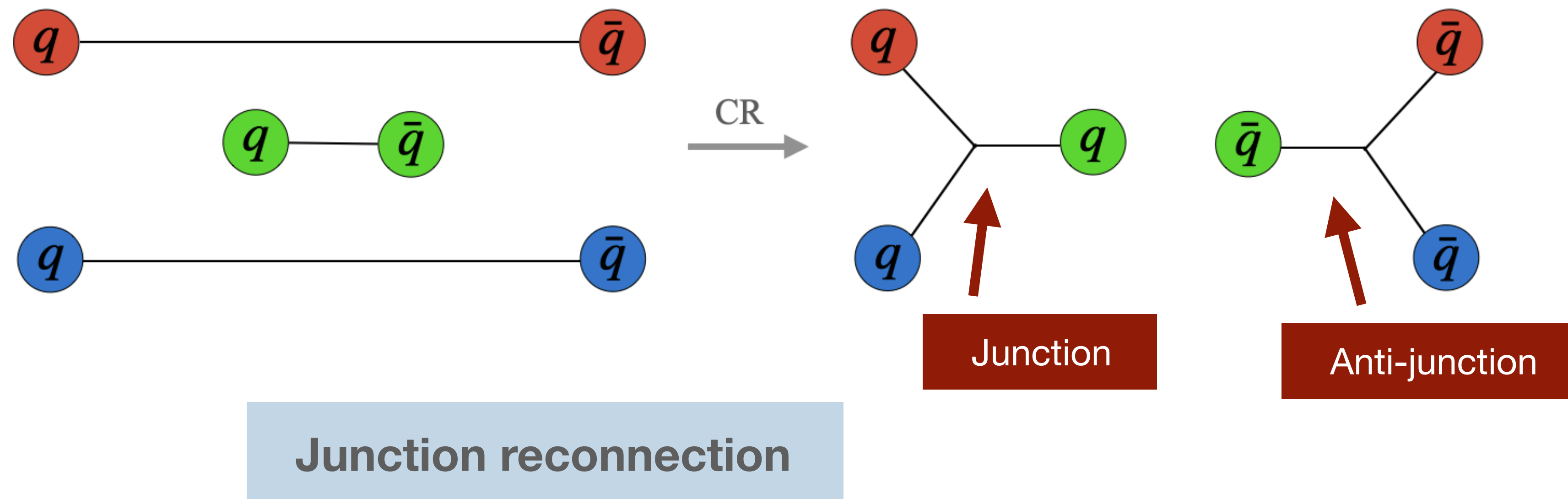
Gluon loop formation

QCD Colour Reconnections

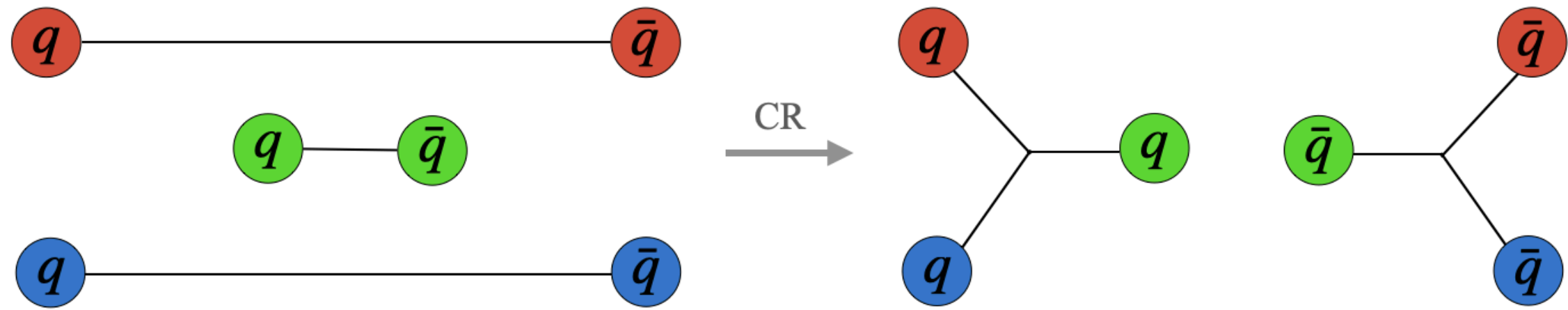
Stochastically restores colour-space ambiguities according to **SU(3) algebra**

- Allows for reconnections to **minimise string lengths**

What about the **red-green-blue** colour singlet state?

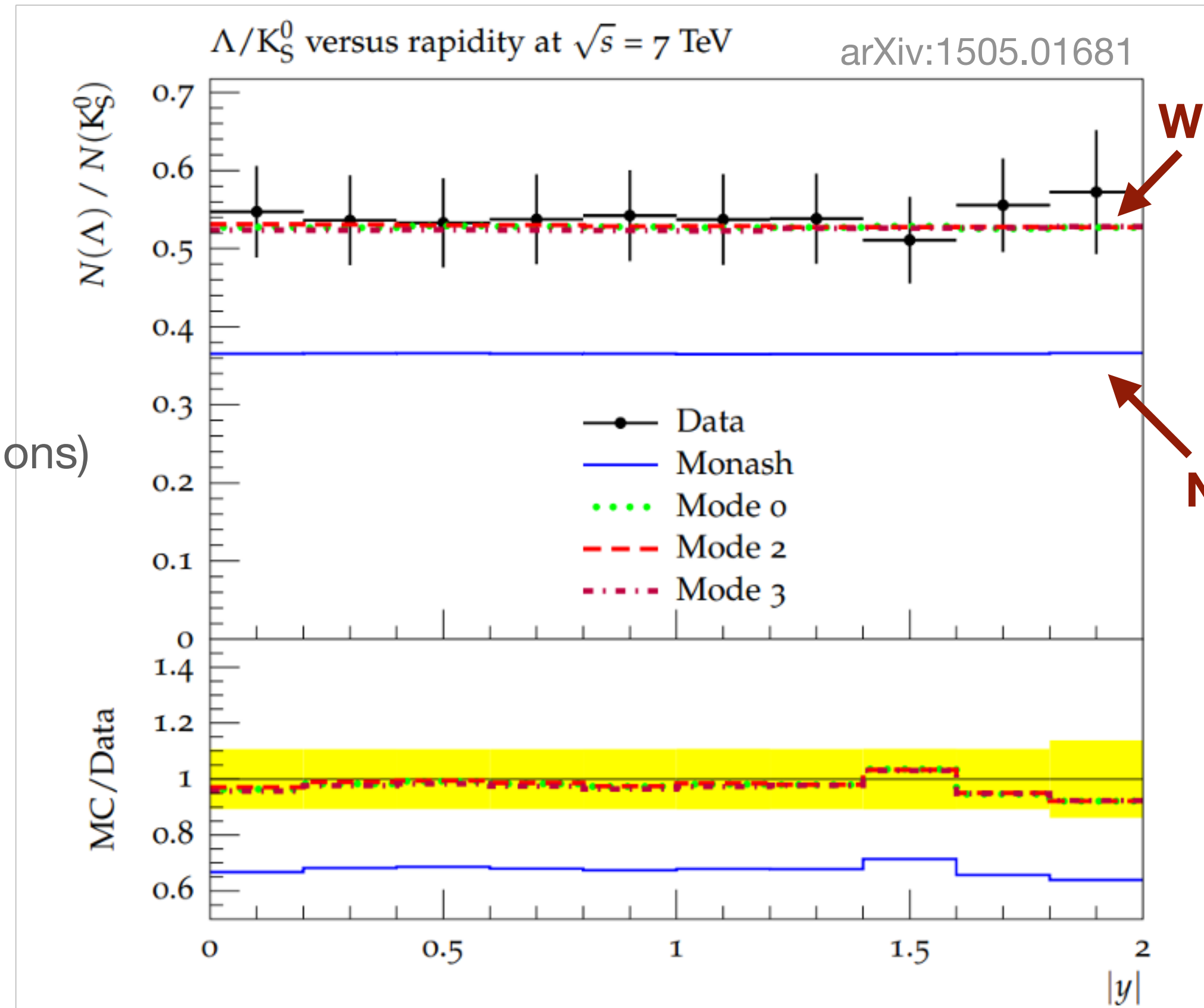
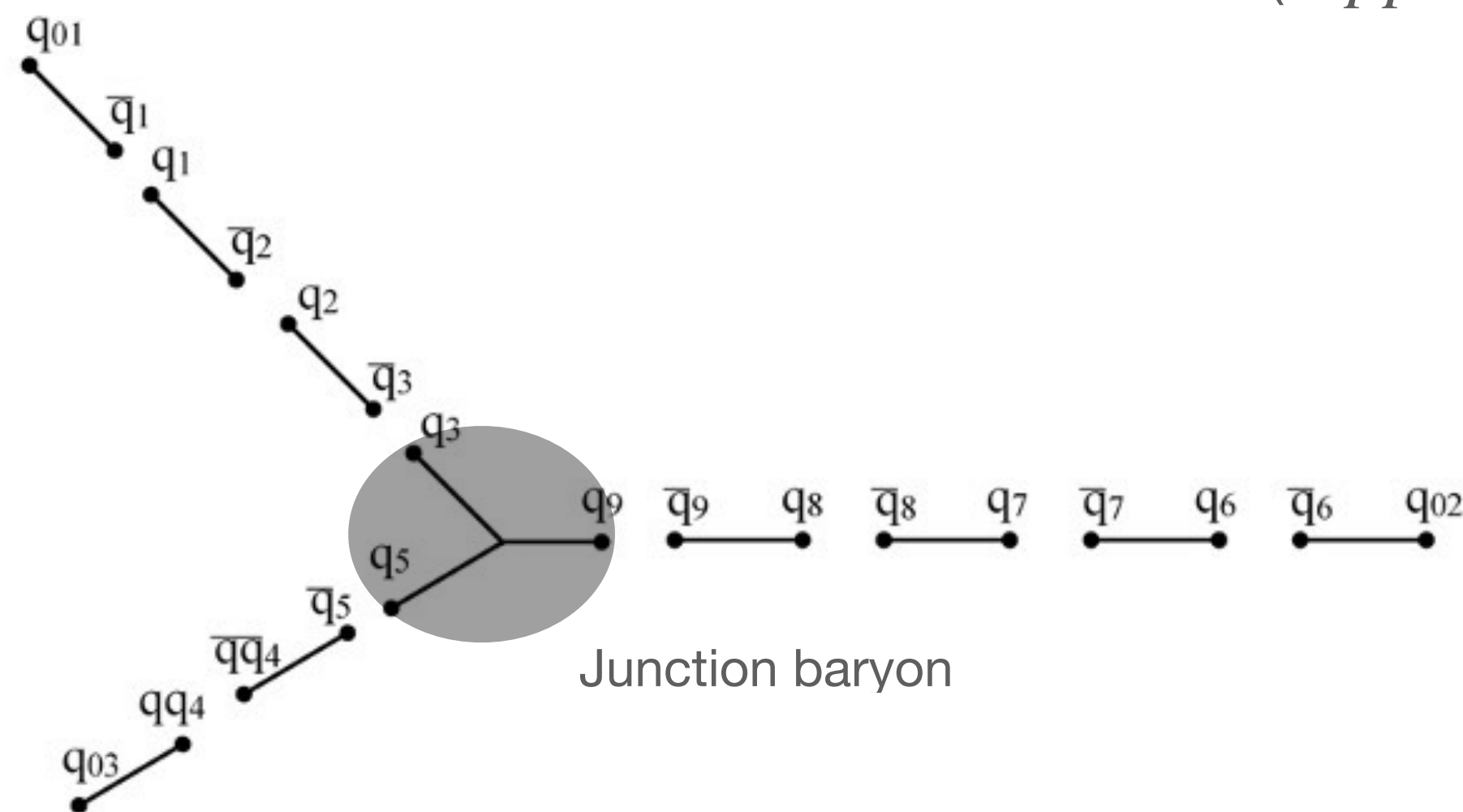


Junctions

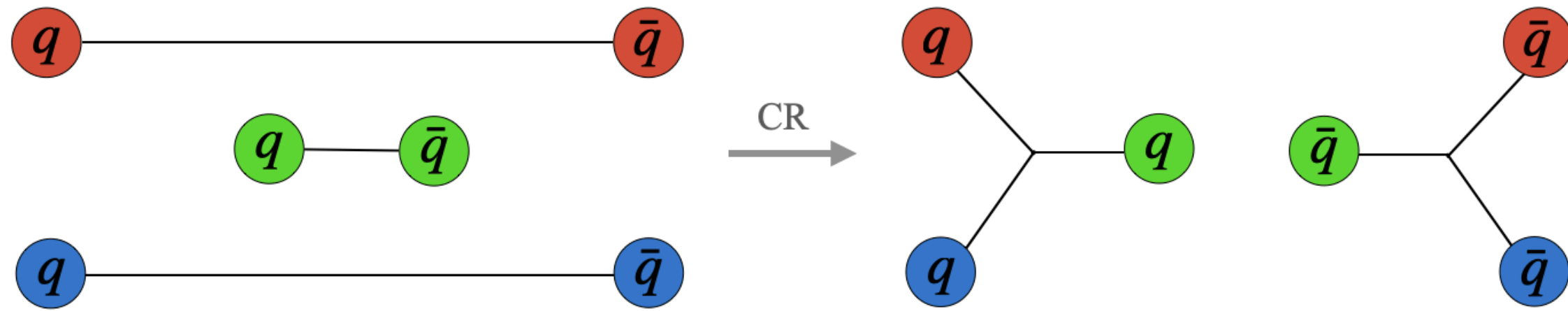


Mechanism for **baryon production**

➤ ~40% of baryons are from junctions in PYTHIA
(in pp collisions)



Junctions



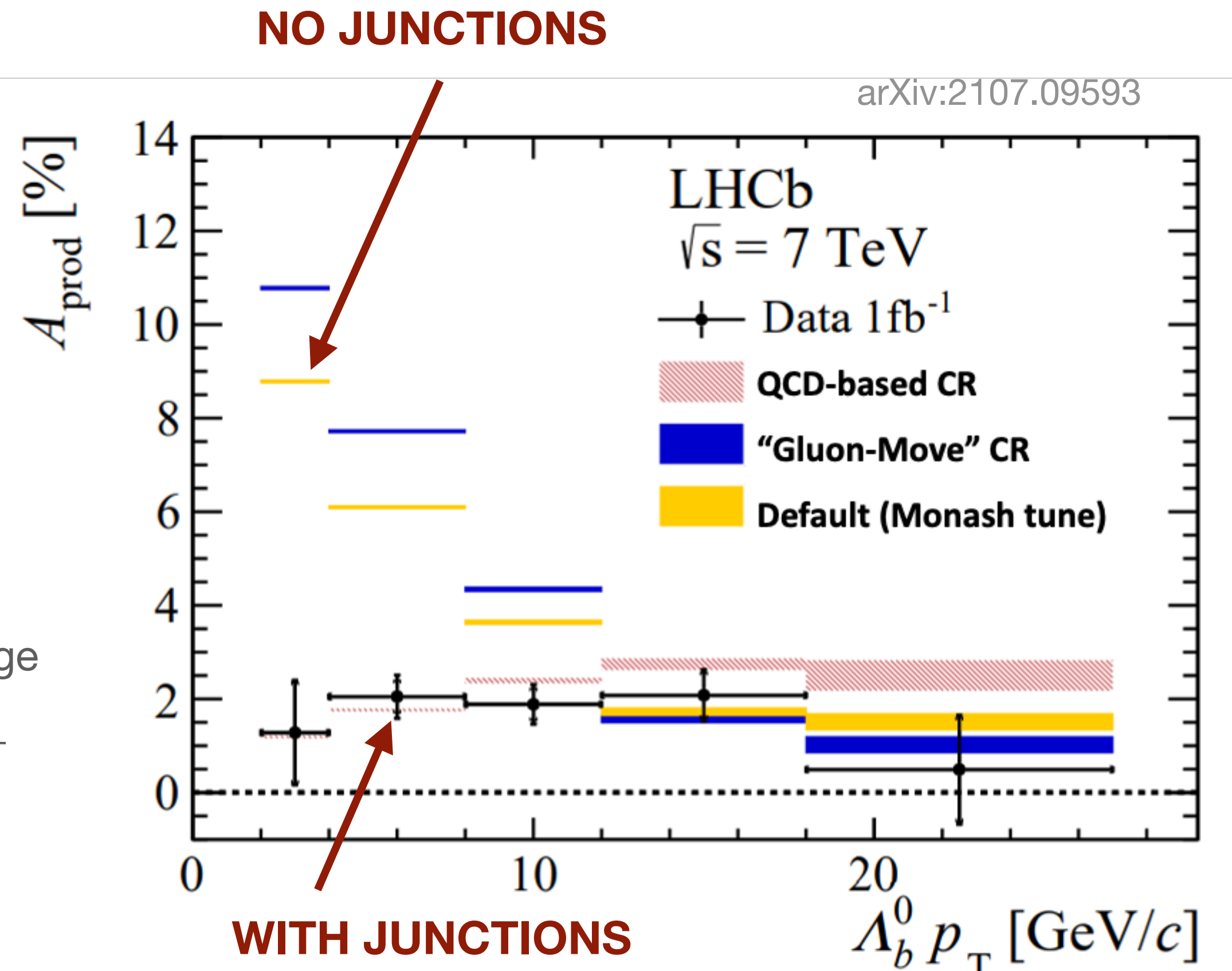
Mechanism for **baryon production**

- ~40% of baryons are from junctions in PYTHIA

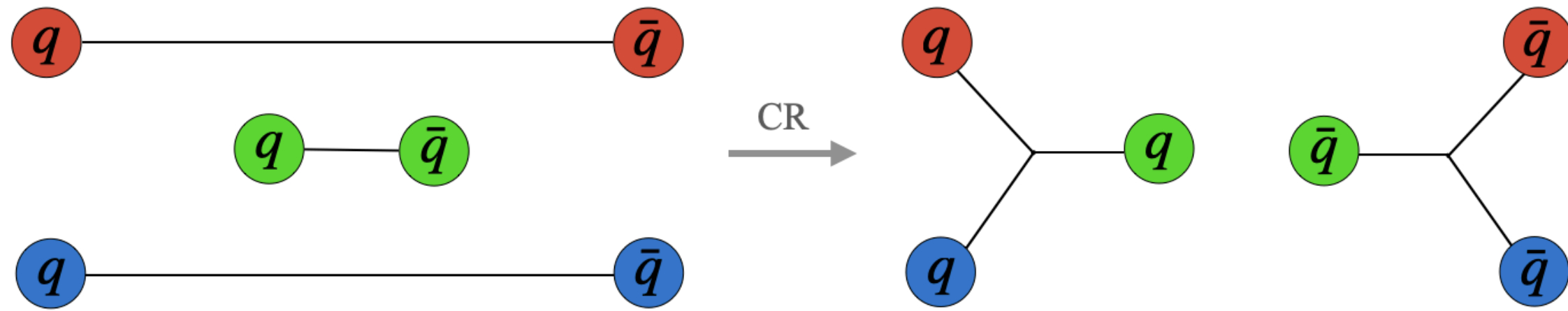
Asymmetries

- Equal amount of junctions and anti junctions are formed

Junctions typically **form between jets** → as jets are likely to have large opening angles due to available phase space, **junction sits at low p_{\perp}**



Junctions



Mechanism for baryon production

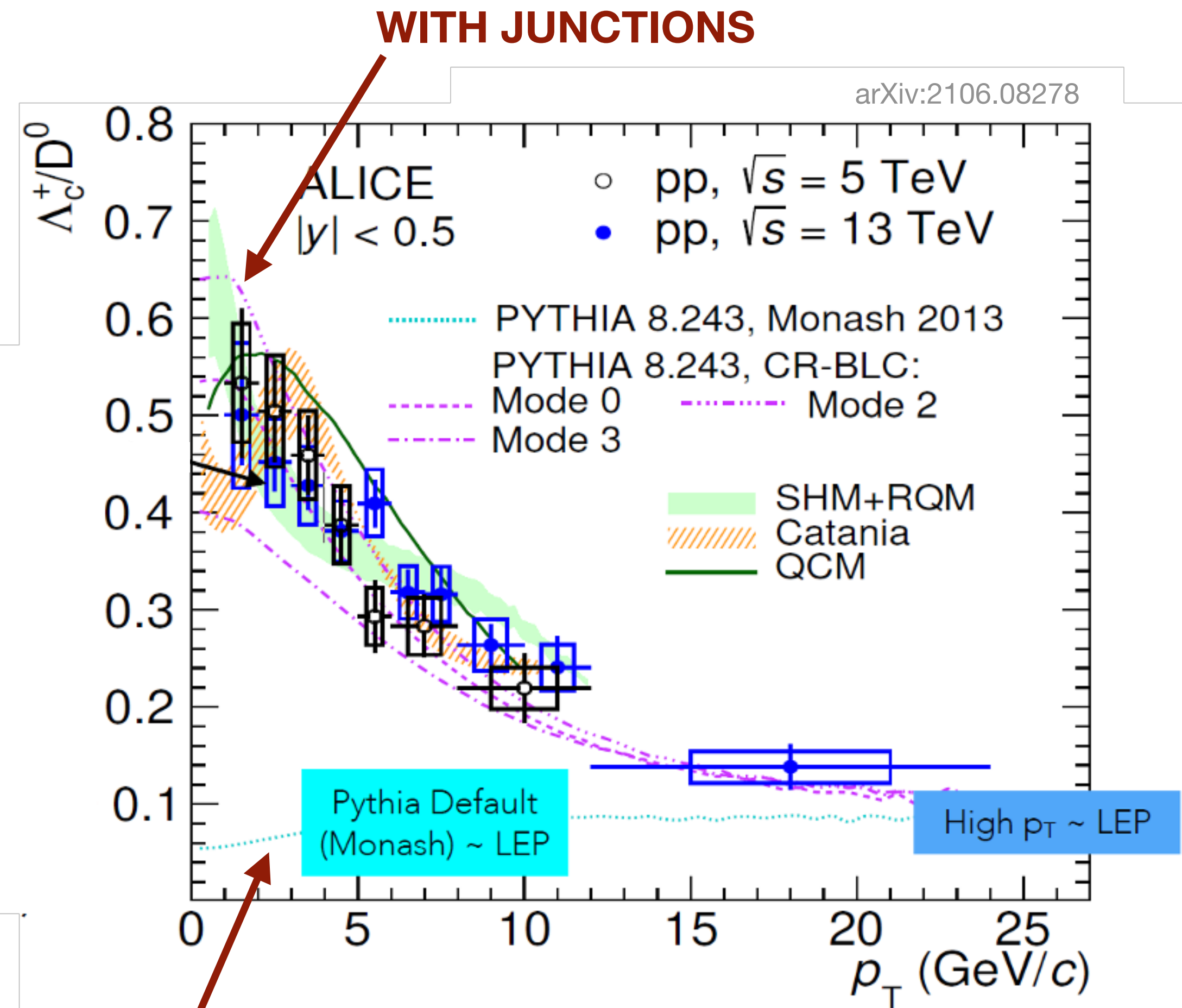
- ~40% of baryons are from junctions in PYTHIA

Asymmetries

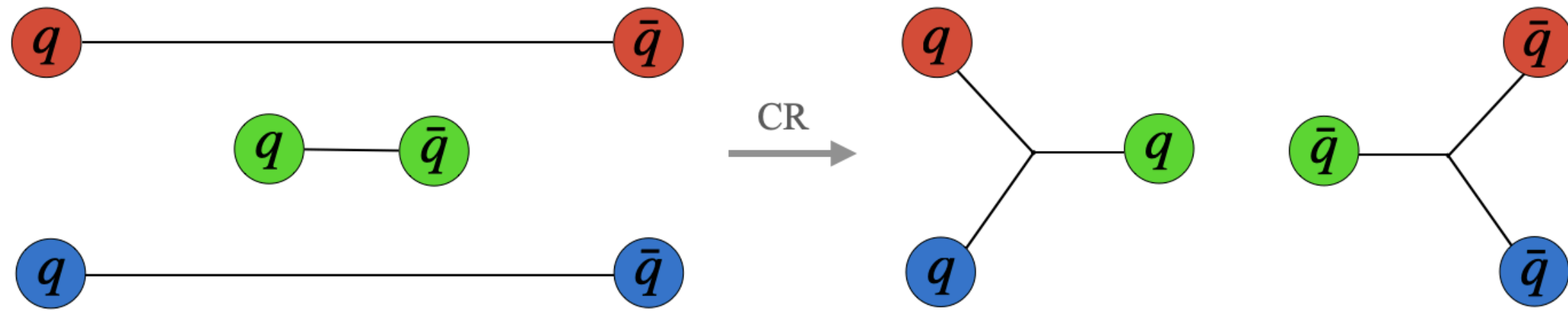
- Equal amount of junctions and anti junctions are formed
- Junctions typically **form between jets** → as jets are likely to have large opening angles due to available phase space, **junction sits at low p_{\perp}**

Heavy flavour baryons

- **~70% of heavy baryons** are from junctions in PYTHIA



Junctions



Mechanism for baryon production

- ~40% of baryons are from junctions in PYTHIA

Asymmetries

- Equal amount of junctions and anti junctions are formed

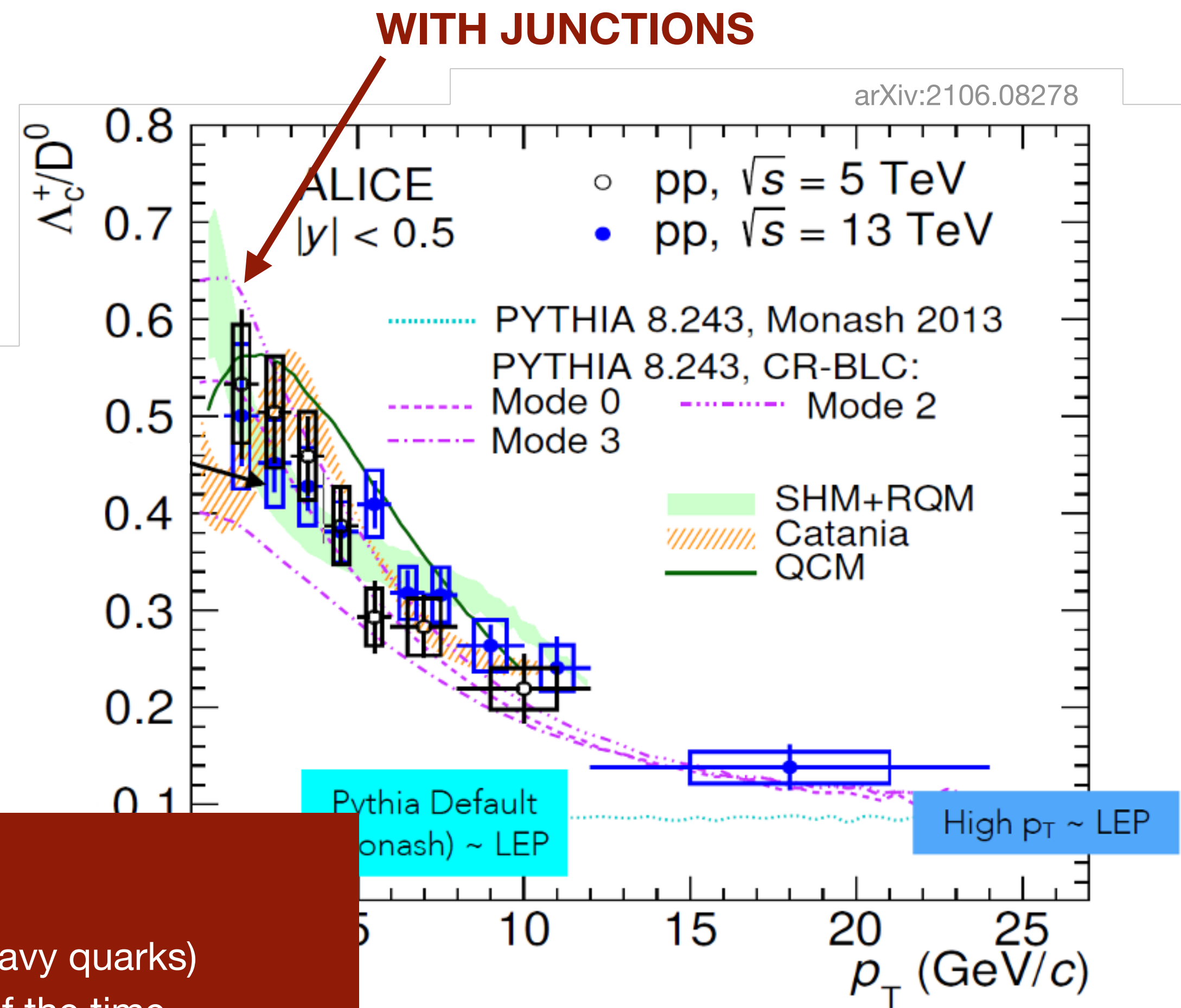
Junctions typically **form between jets** → as jets are likely to have large opening angles due to available phase space, **junction sits at low p_{\perp}**

Heavy flavour baryons

- ~70% of heavy flavour baryons

Current implementation

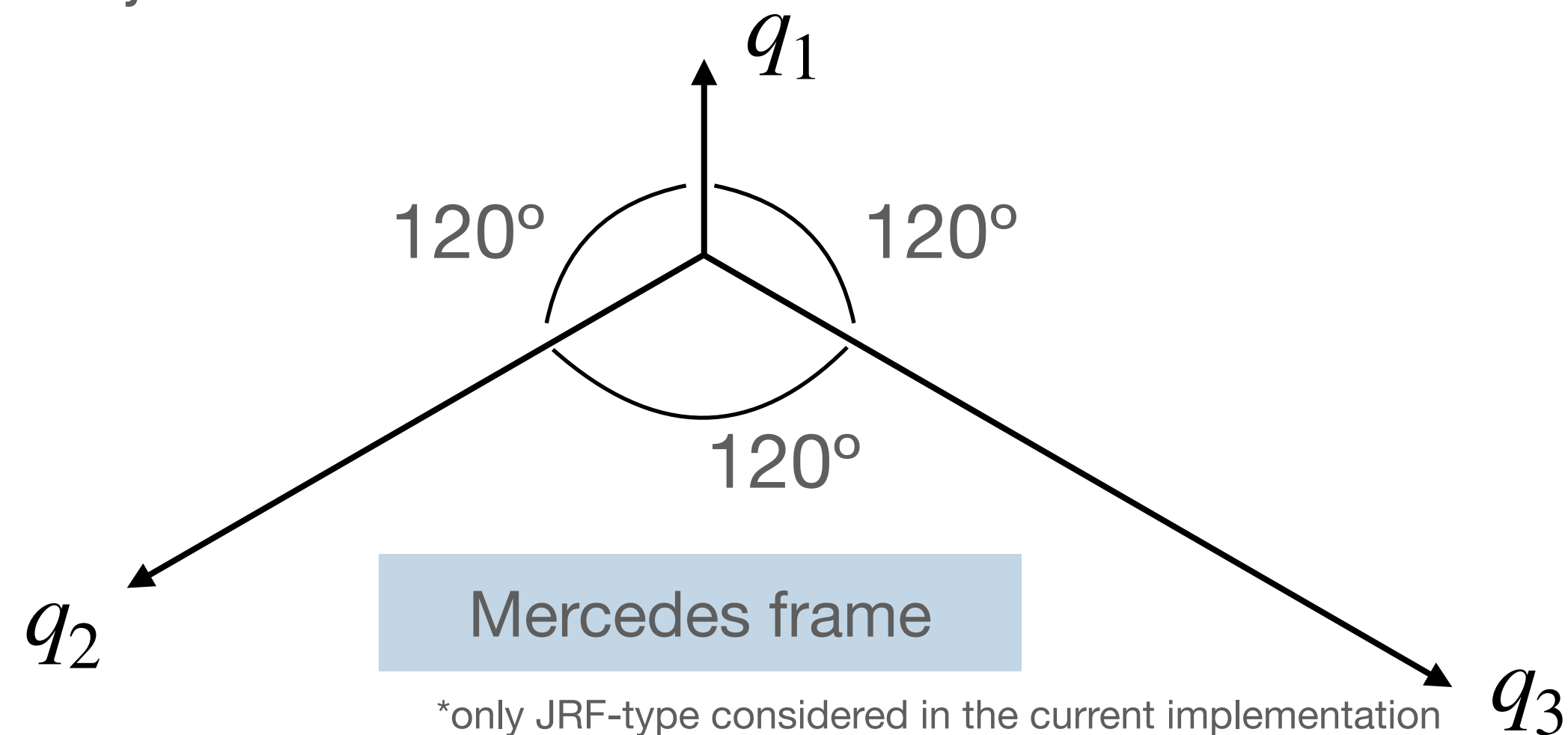
- Runs into cases with no solution (particularly for heavy quarks)
- Relies on convergence procedure that fails ~10% of the time



Junction Rest Frame

What is the junction rest frame?

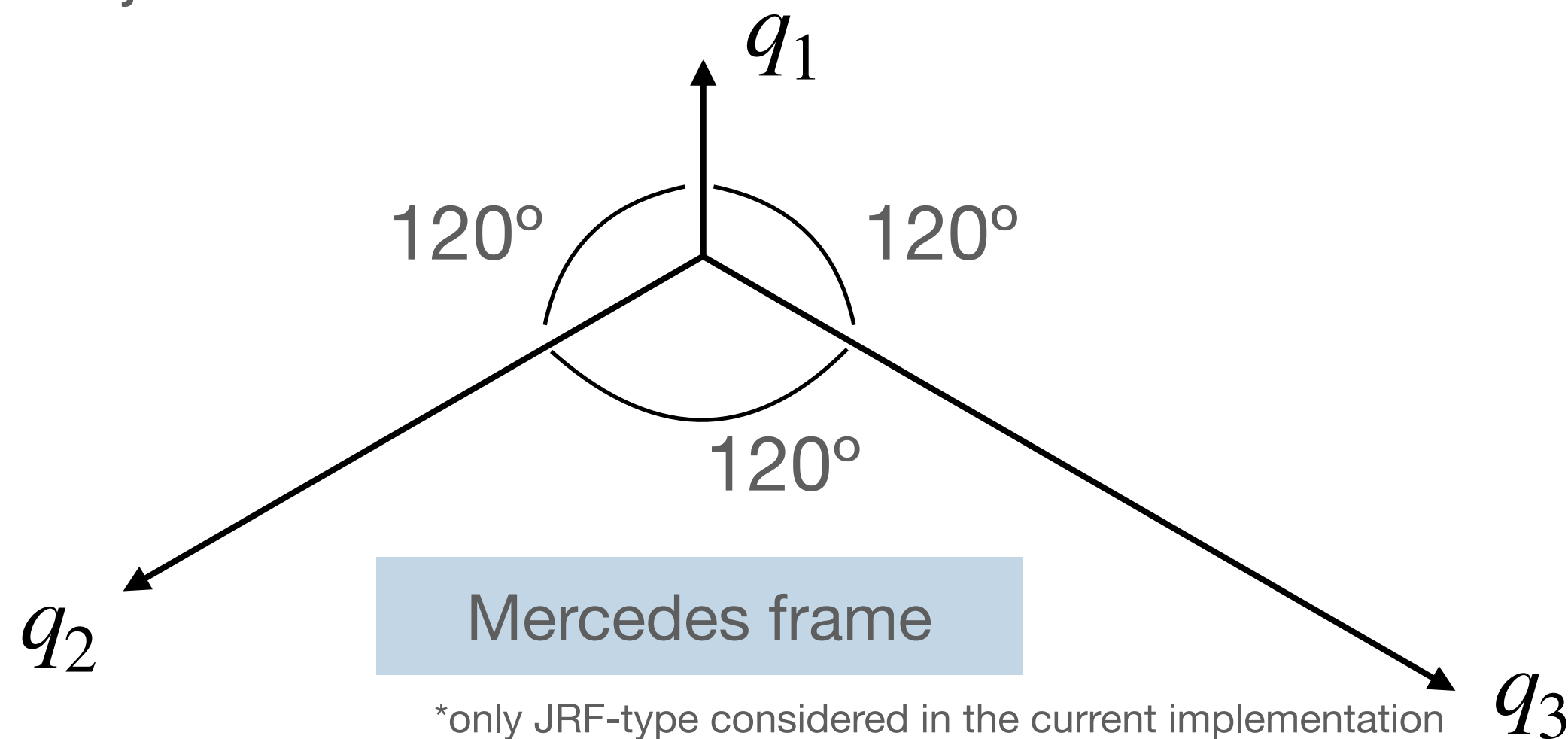
- If the momenta of the junction legs are at 120° angles
- the pull in each direction on the junction is equal
- junction is at rest



Junction Rest Frame

What is the **junction rest frame**?

- If the momenta of the junction legs are at 120° angles
- the pull in each direction on the junction is equal
- junction is at rest

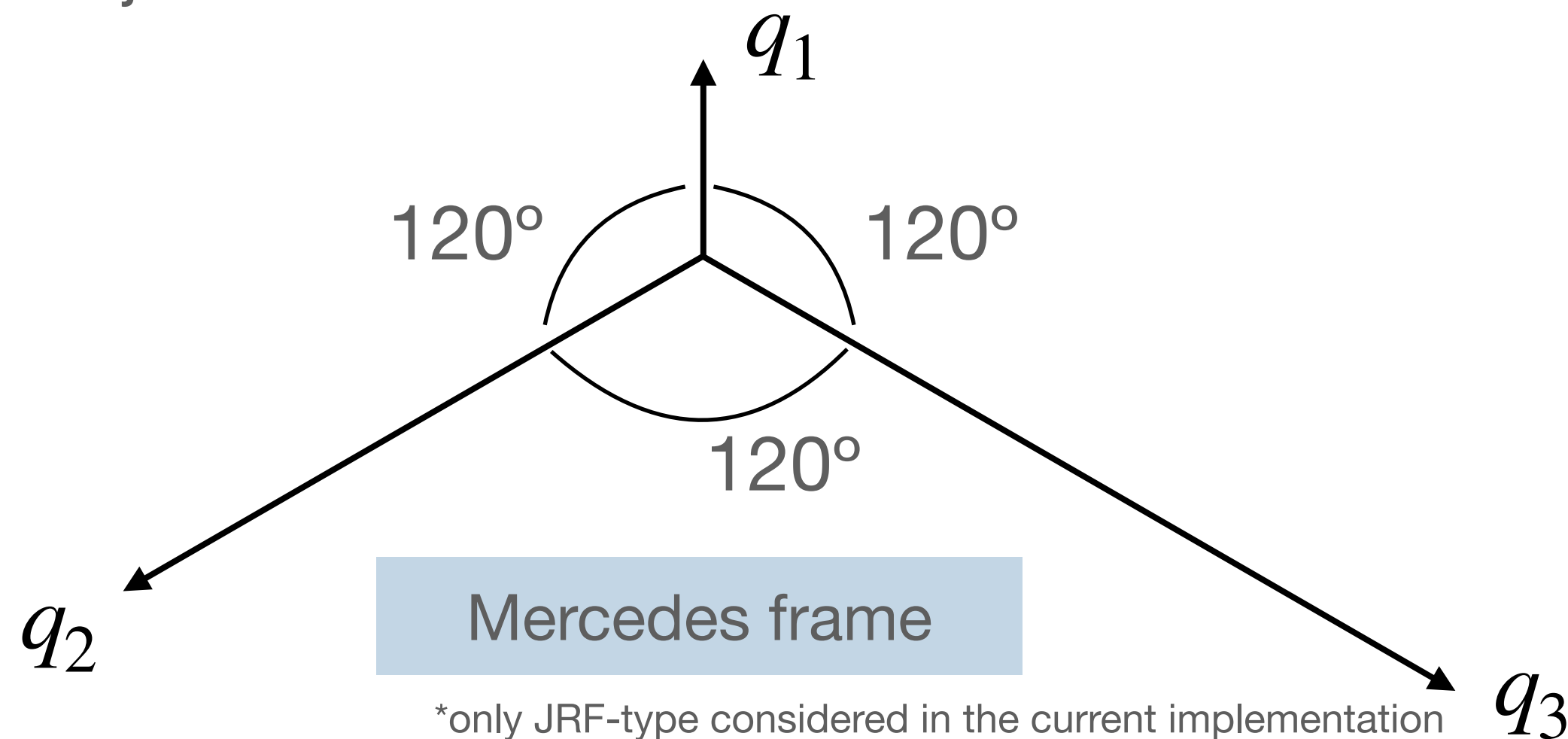


Does a boost to the mercedes frame always exist?

Junction Rest Frame

What is the junction rest frame?

- If the momenta of the junction legs are at 120° angles
- the pull in each direction on the junction is equal
- junction is at rest

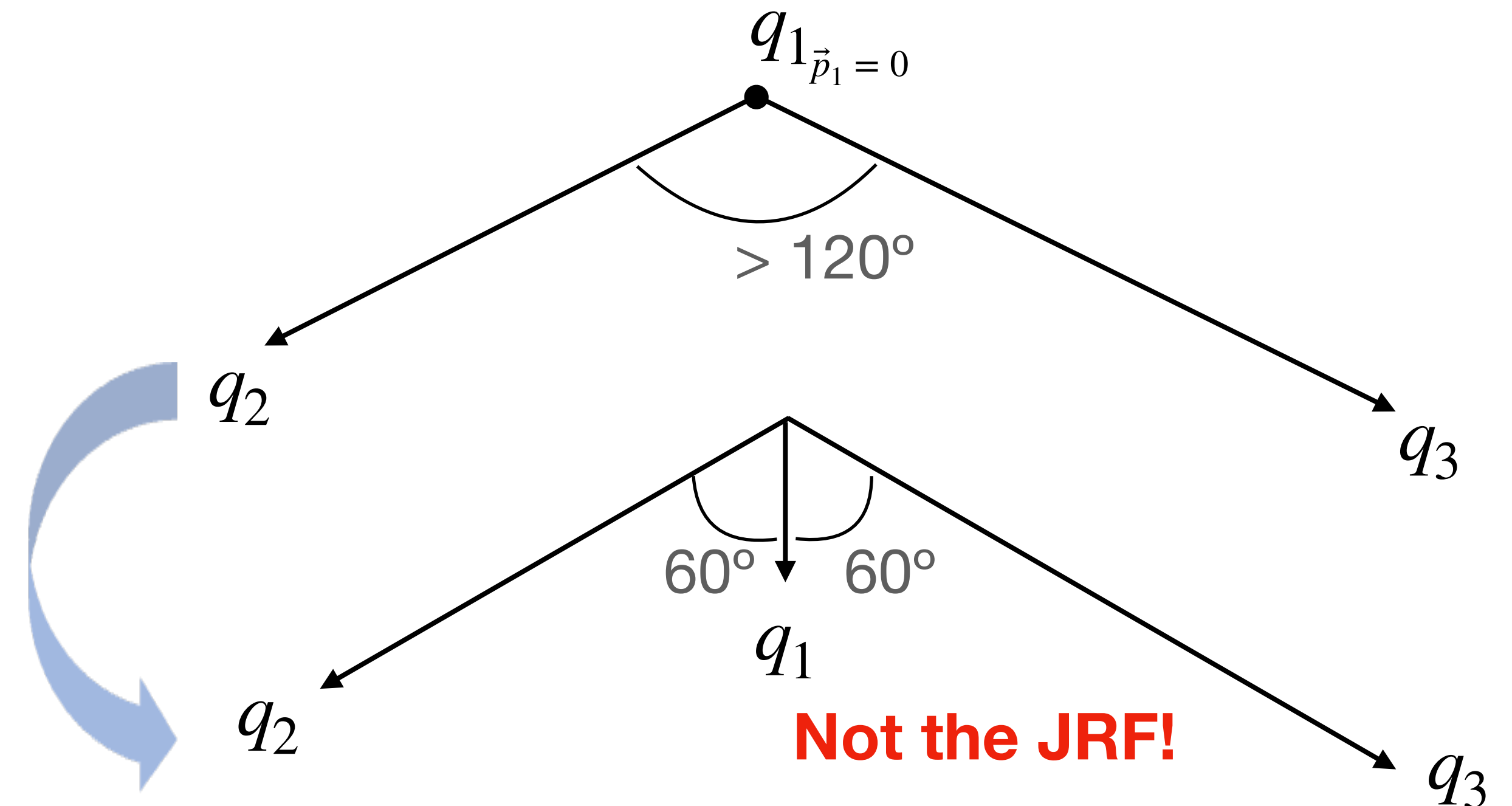


Does a boost to the mercedes frame always exist?

Consider the following:

In the **rest frame of one of the partons**, and the angle between the other two partons is **greater than 120°**

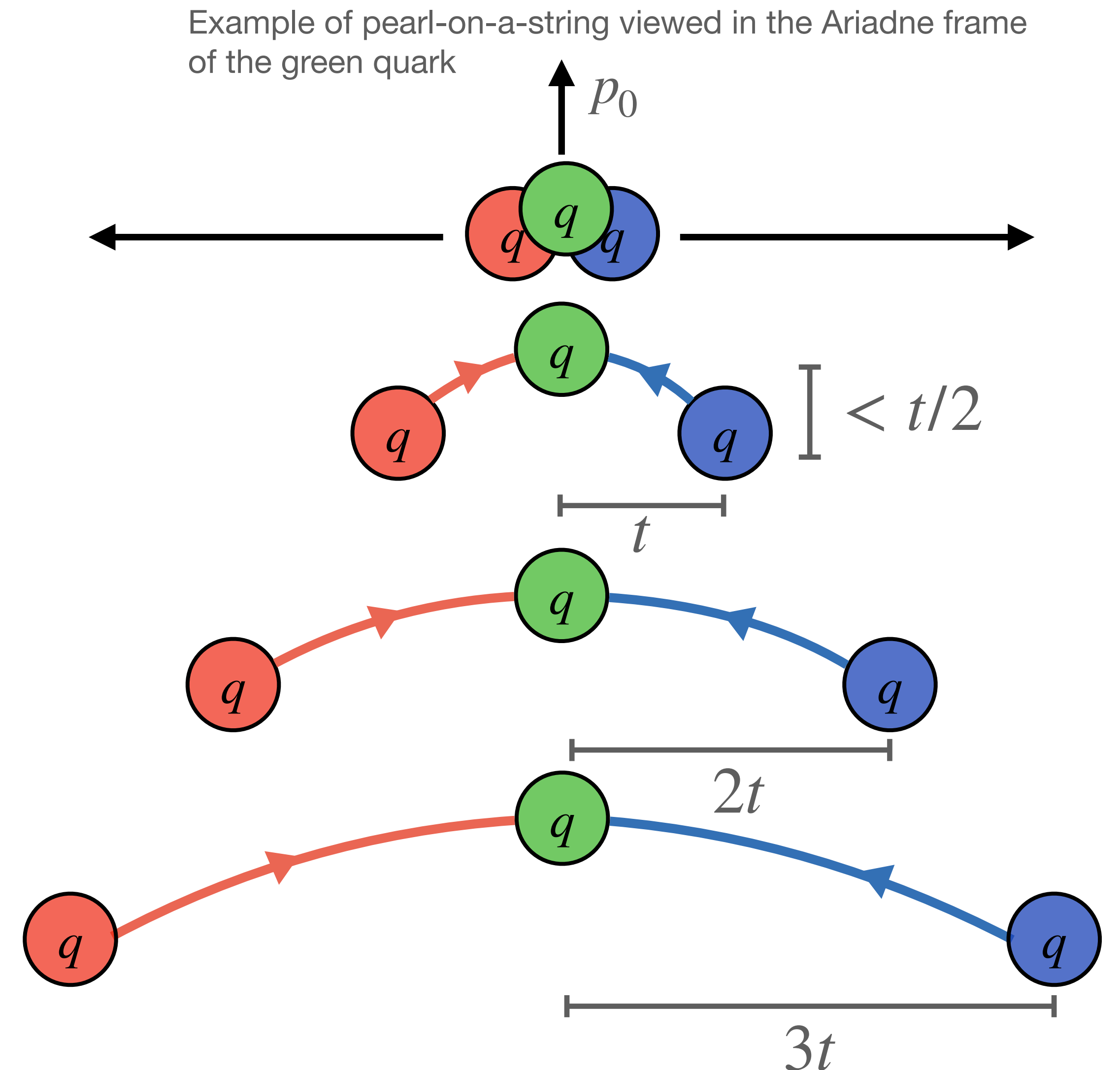
*no special consideration for these cases in current implementation



Pearl-on-a-string

The **junction gets “stuck”** to the soft quark, which we call a **pearl-on-a-string**

- More likely to occur for junctions with heavy flavour endpoints

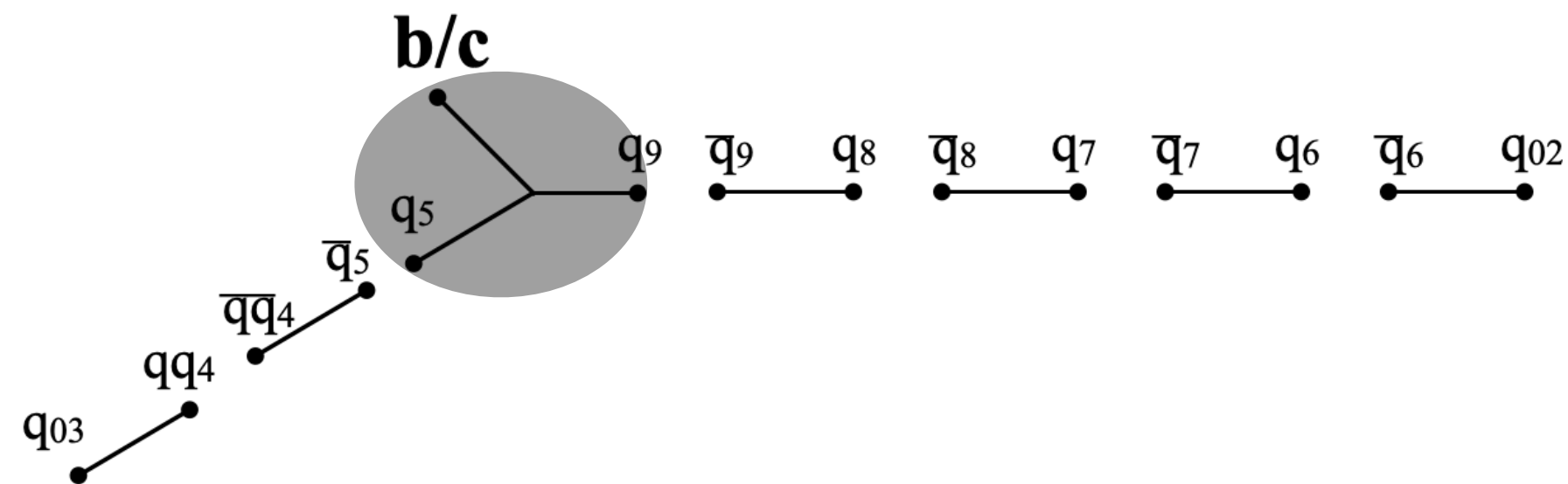


Pearl-on-a-string

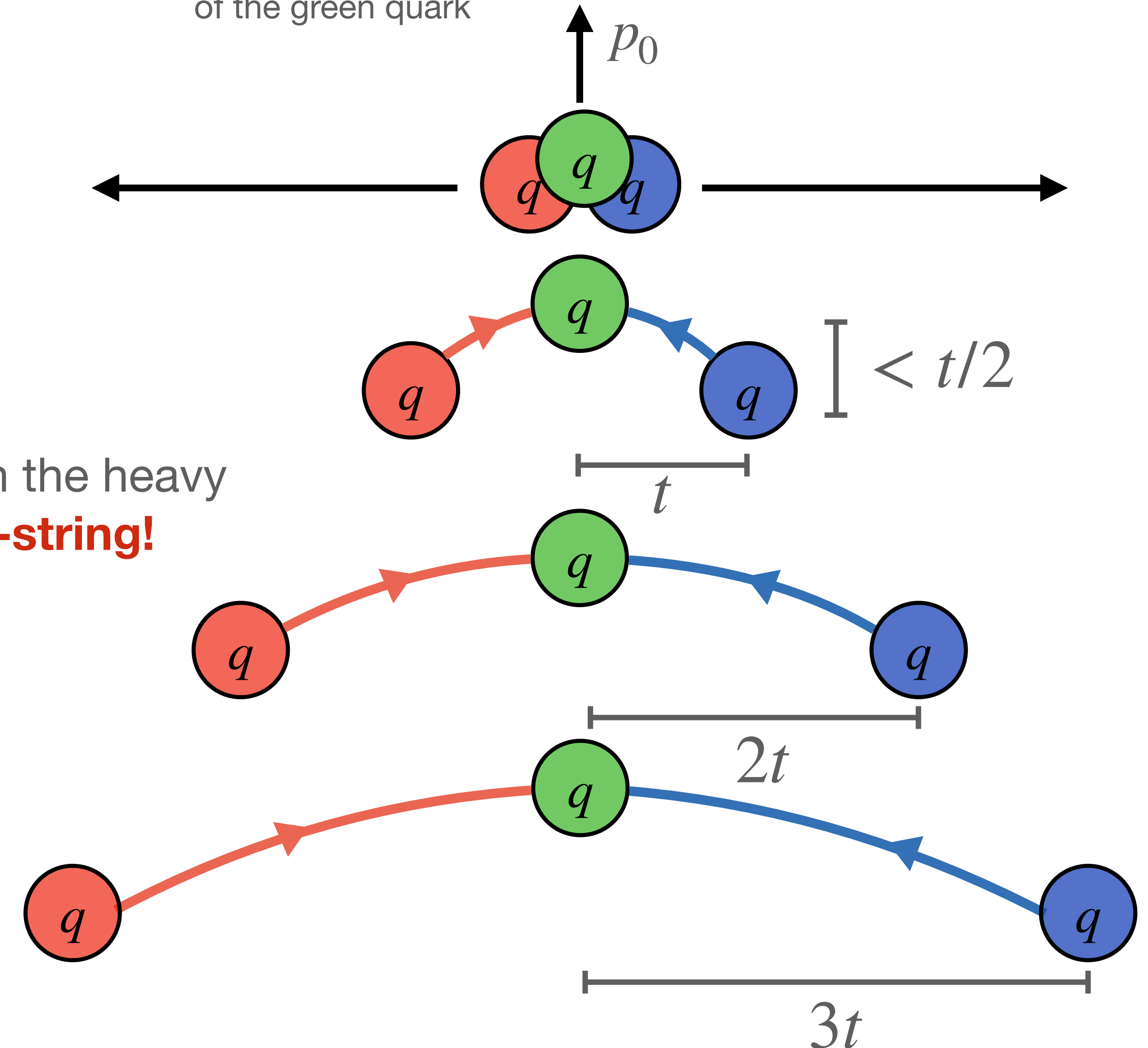
The **junction gets “stuck”** to the soft quark, which we call a **pearl-on-a-string**

- More likely to occur for junctions with heavy flavour endpoints

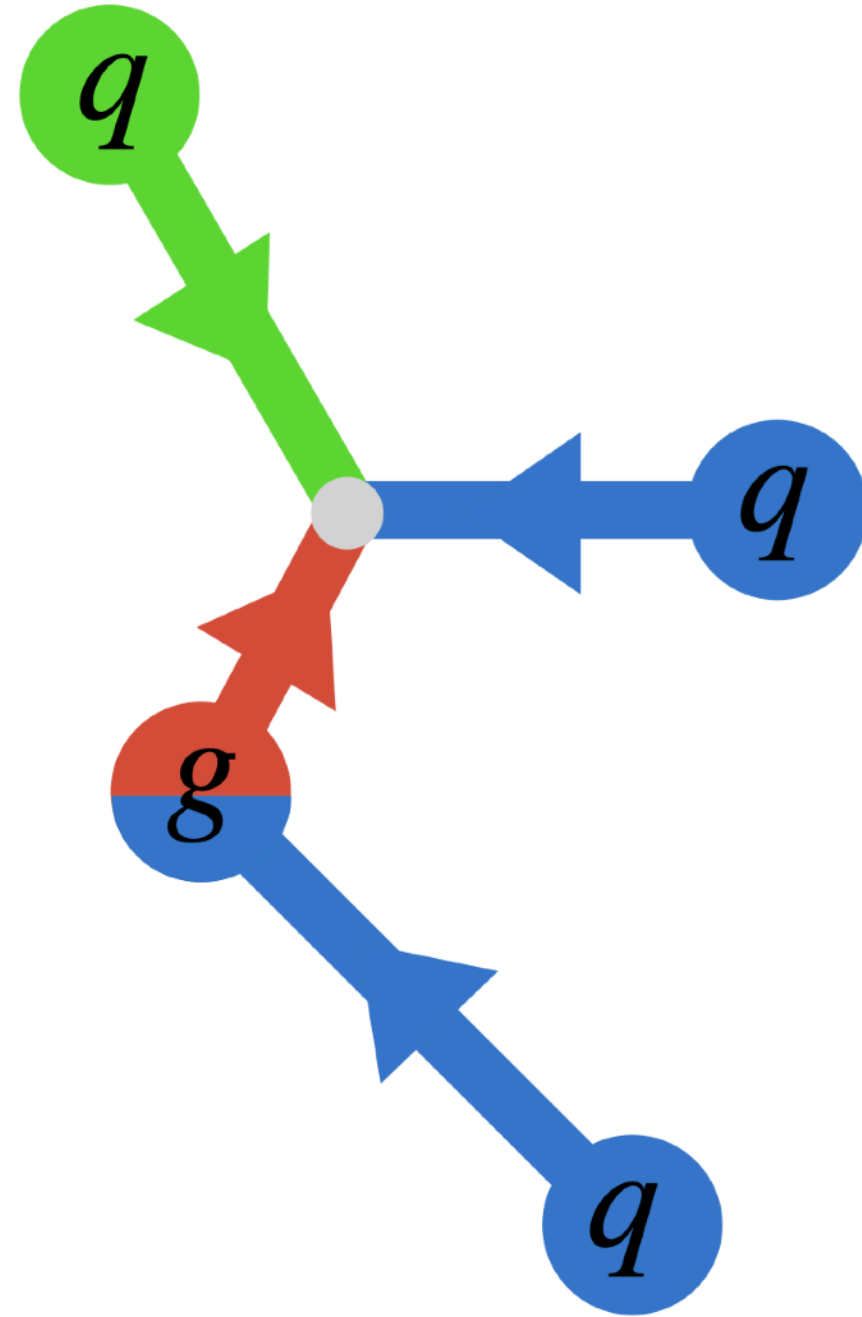
For a junction to make a **heavy baryon**, the junction leg with the heavy quark can't fragment (*i.e.* a “soft” junction leg) = **pearl-on-a-string!**



Example of pearl-on-a-string viewed in the Ariadne frame of the green quark



Updates to averaging



Use an “**average**” JRF

Current procedure assumes the **average is the mercedes frame**

- Uses **energy weighted sum** of momenta on each junction leg
- Relies on convergence **procedure that fails ~10%** of cases

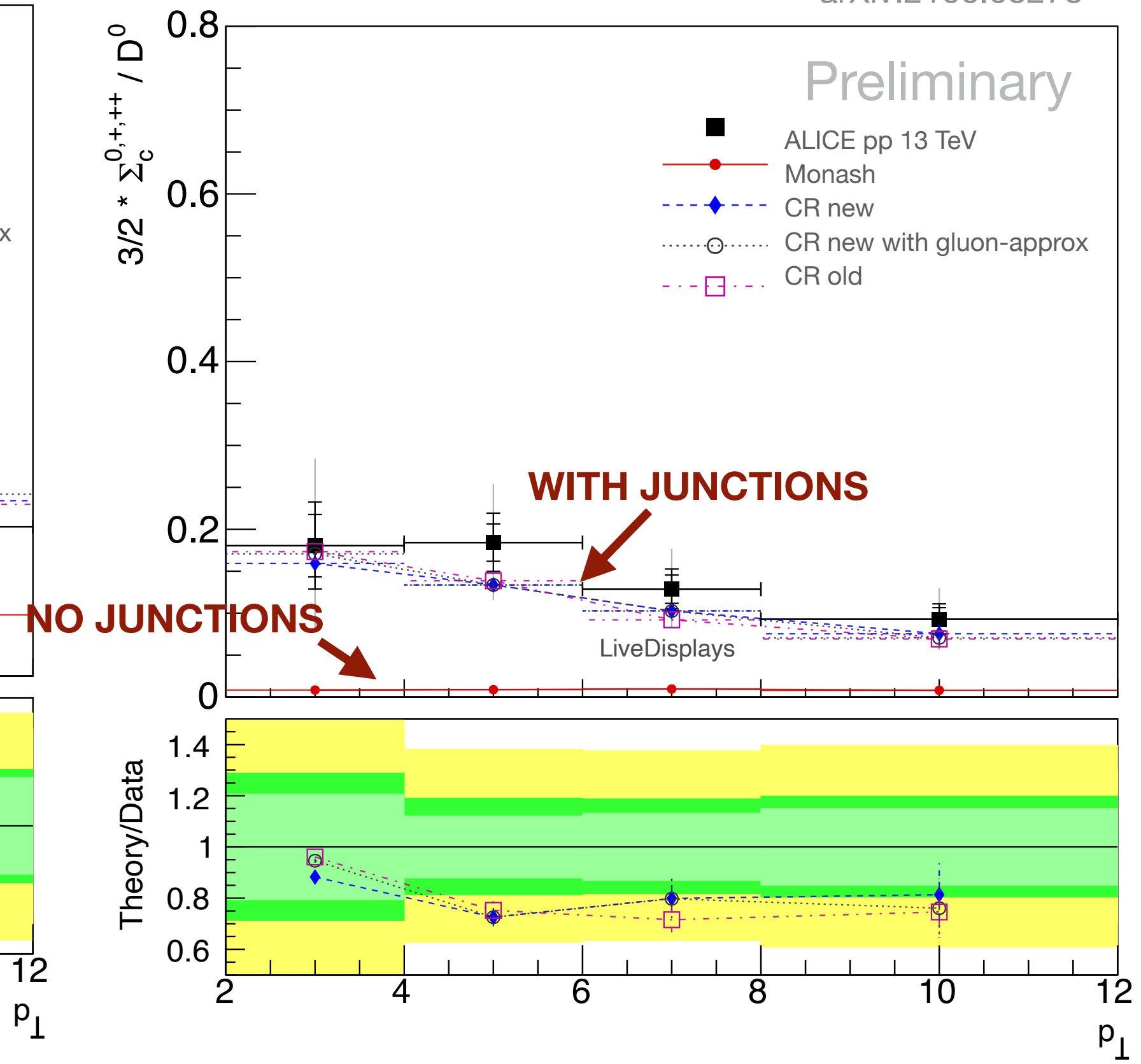
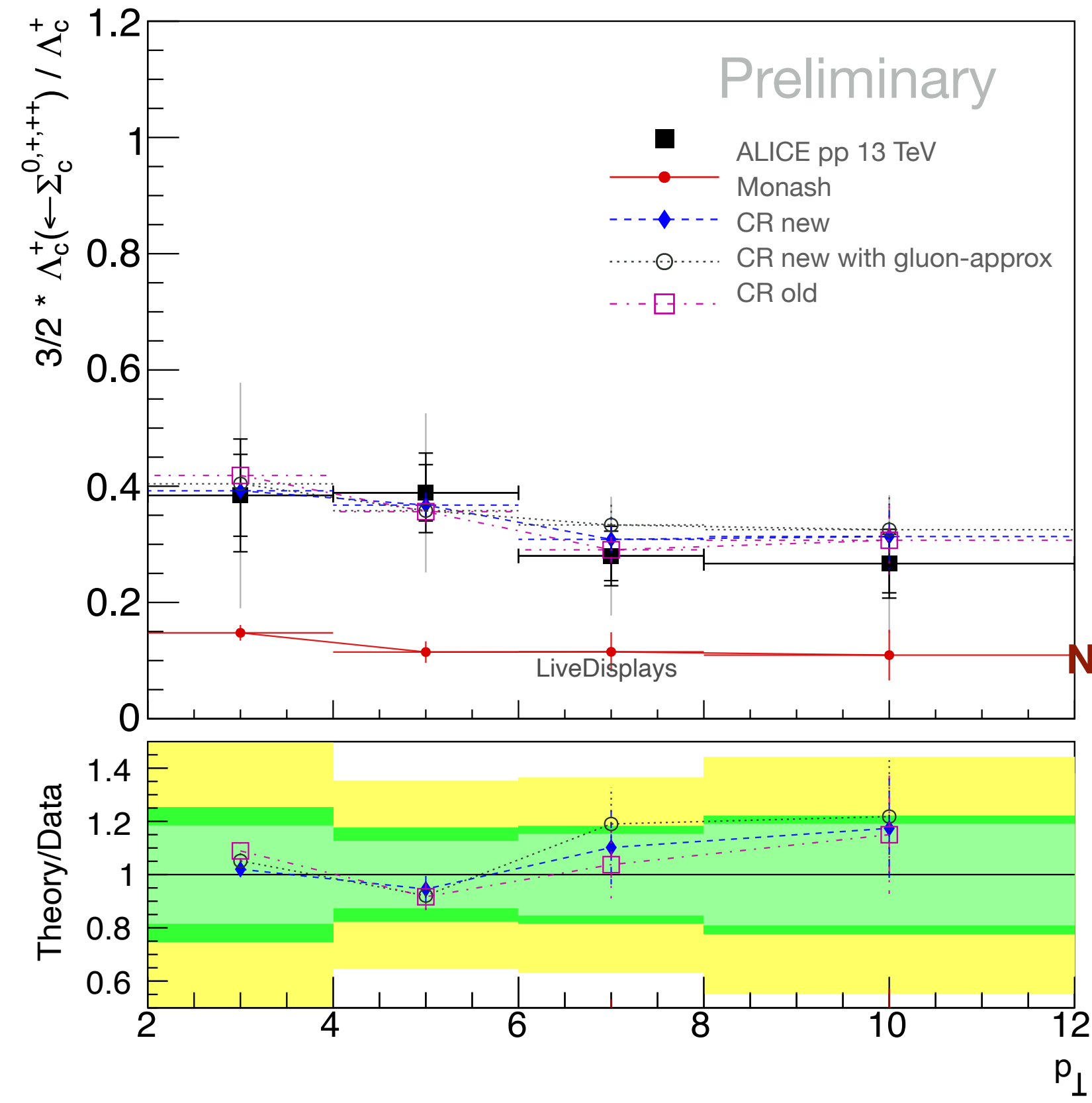
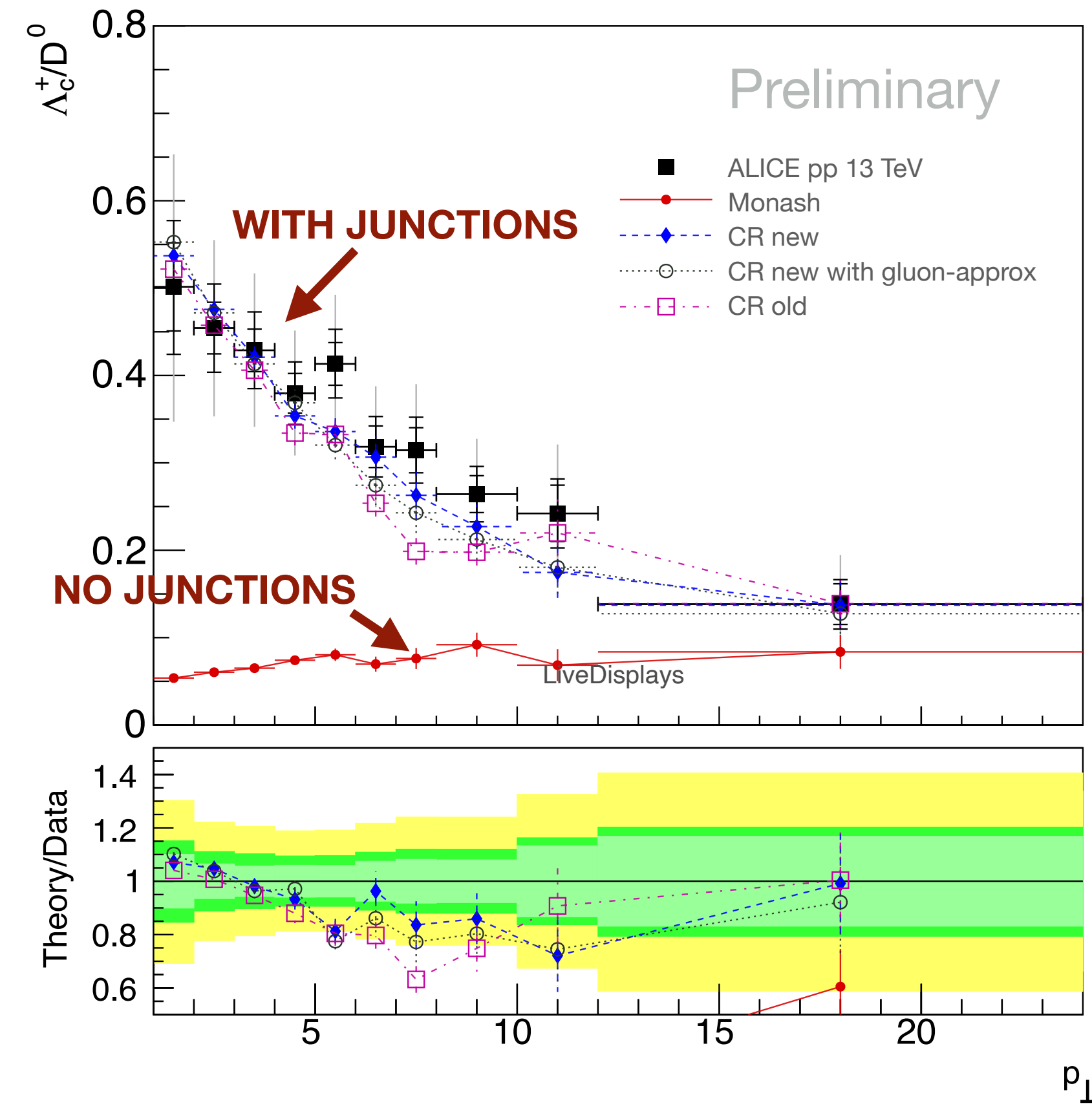
New treatment:

- Considers pull on junction over time and average over junction motion
- Includes pearl-on-a-string
- Allow endpoint oscillations
- No reliance on convergence

- Early time JRF defined by the first parton on each leg
 - Use smallest leg momentum as a measure of effective time for the JRF
- When softest parton has lost its momentum, the next parton dominates the pull

Junctions

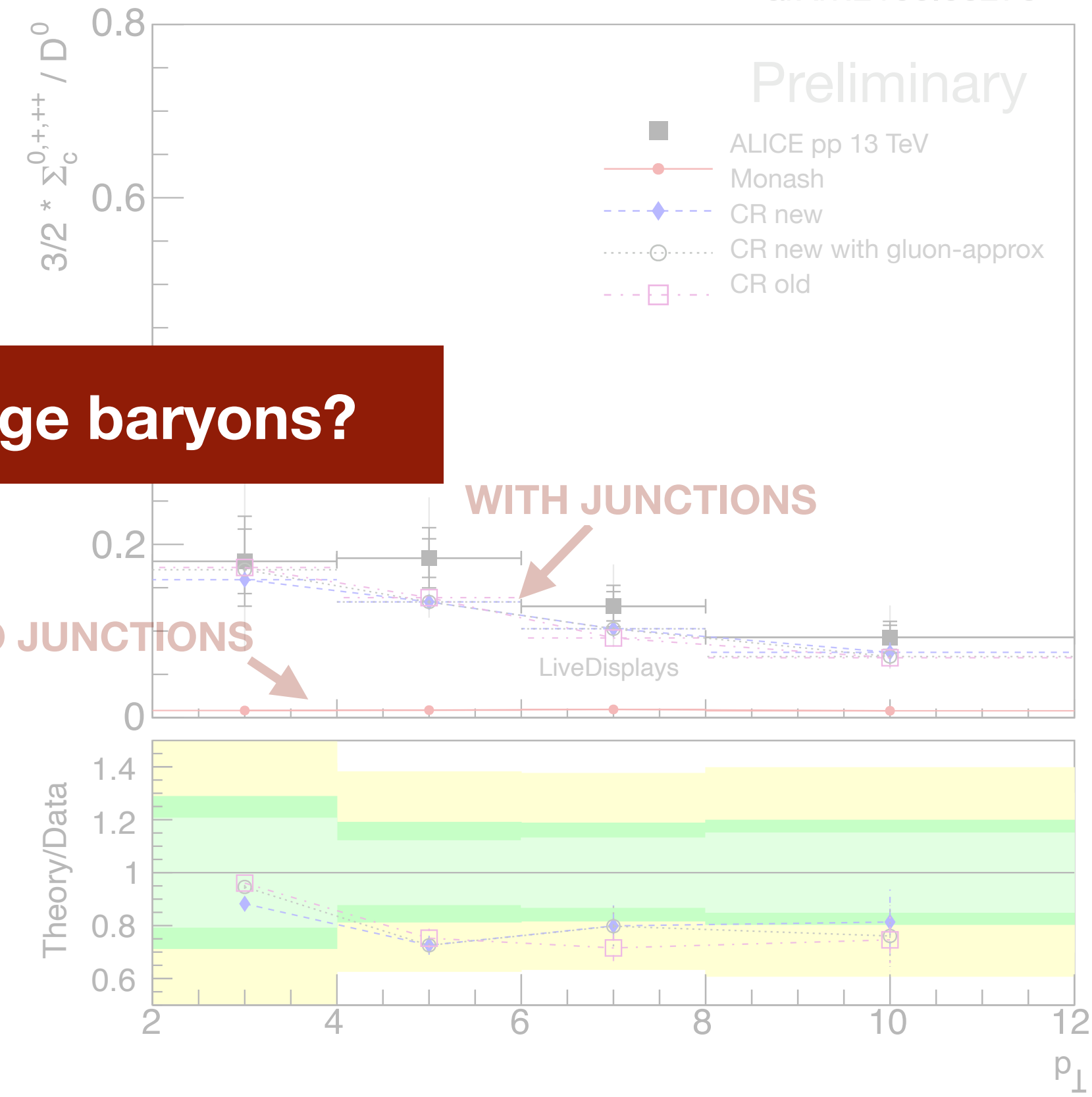
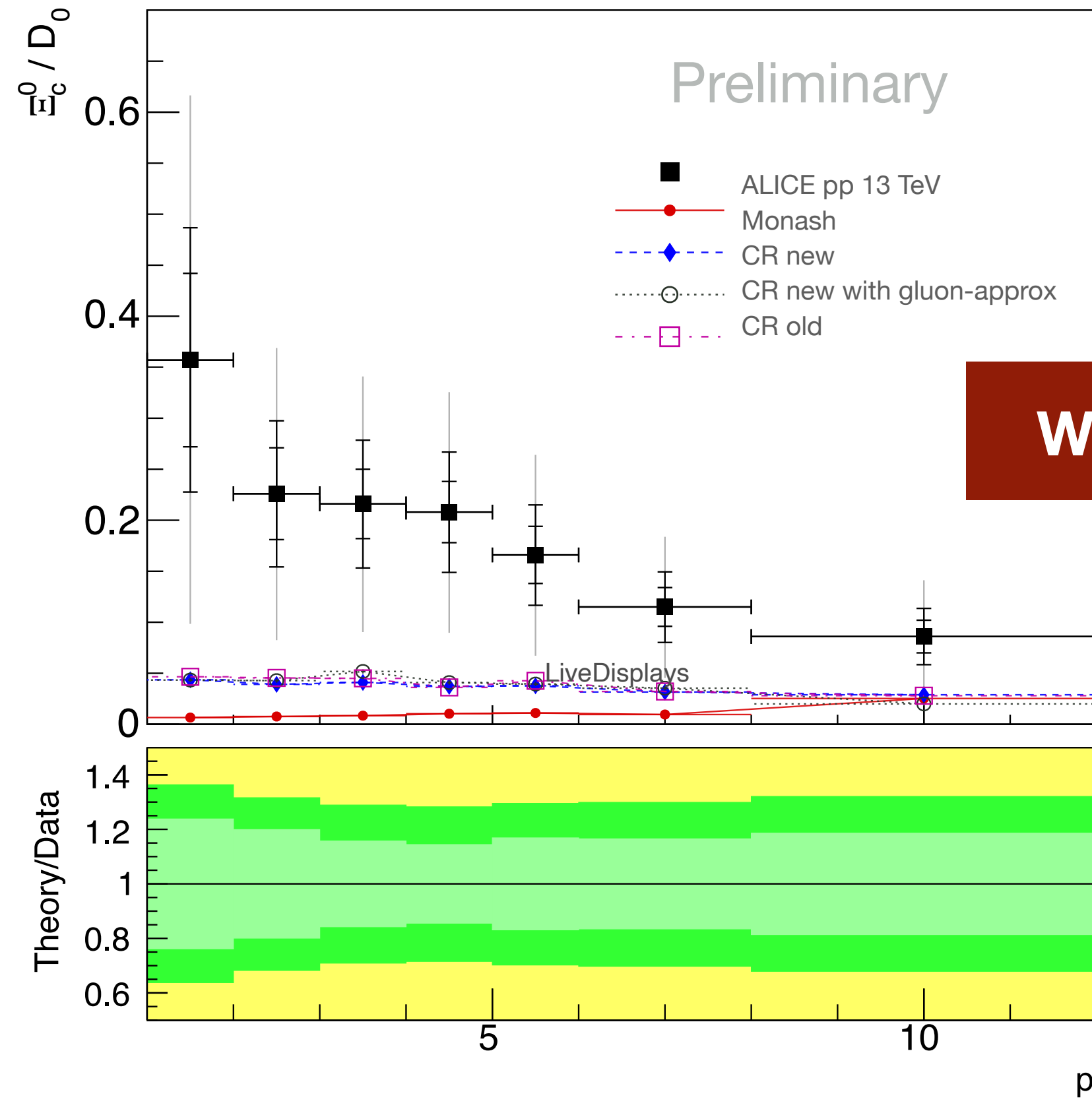
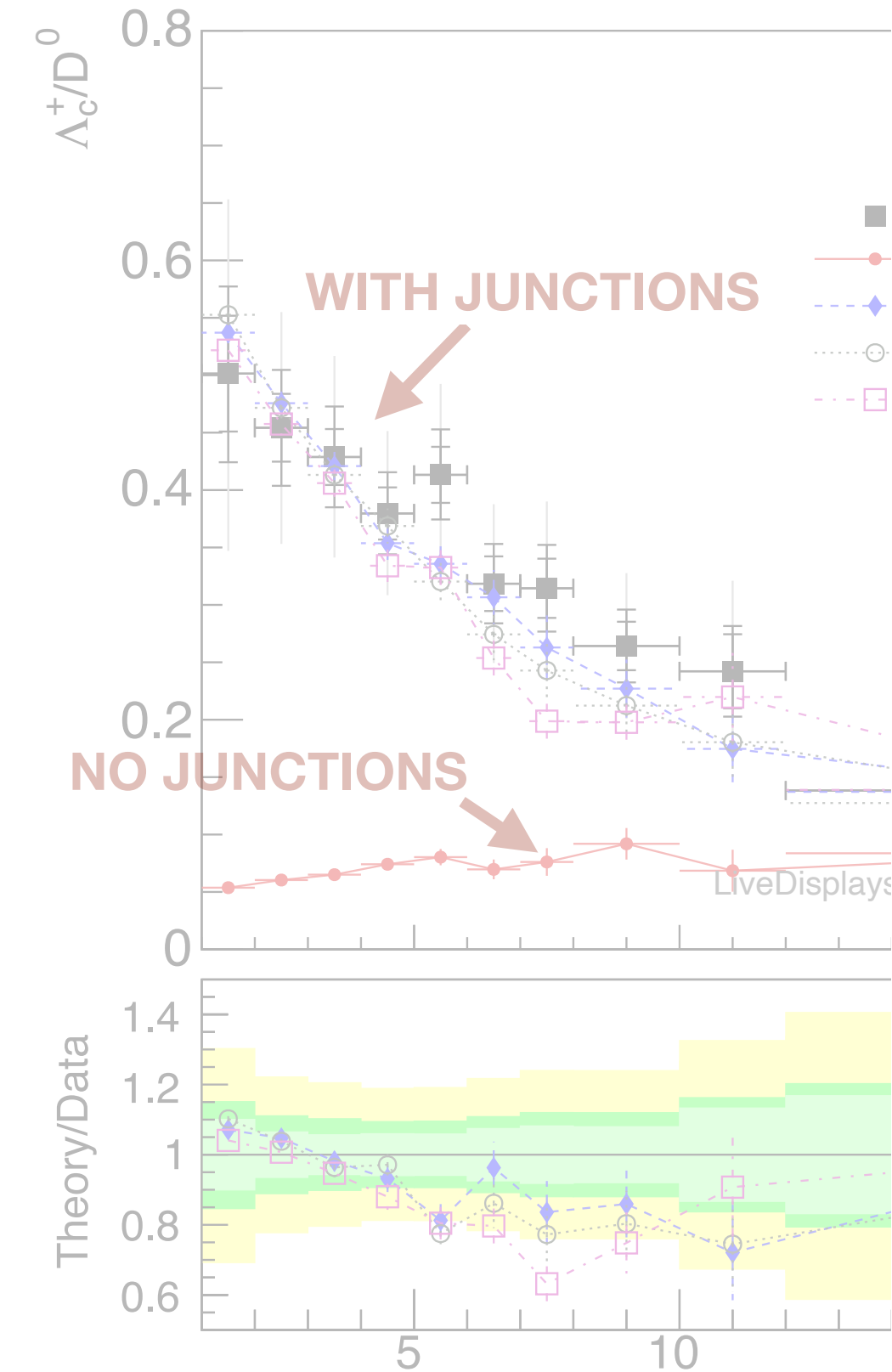
arXiv:2106.08278



Junctions

arXiv:2105.05187

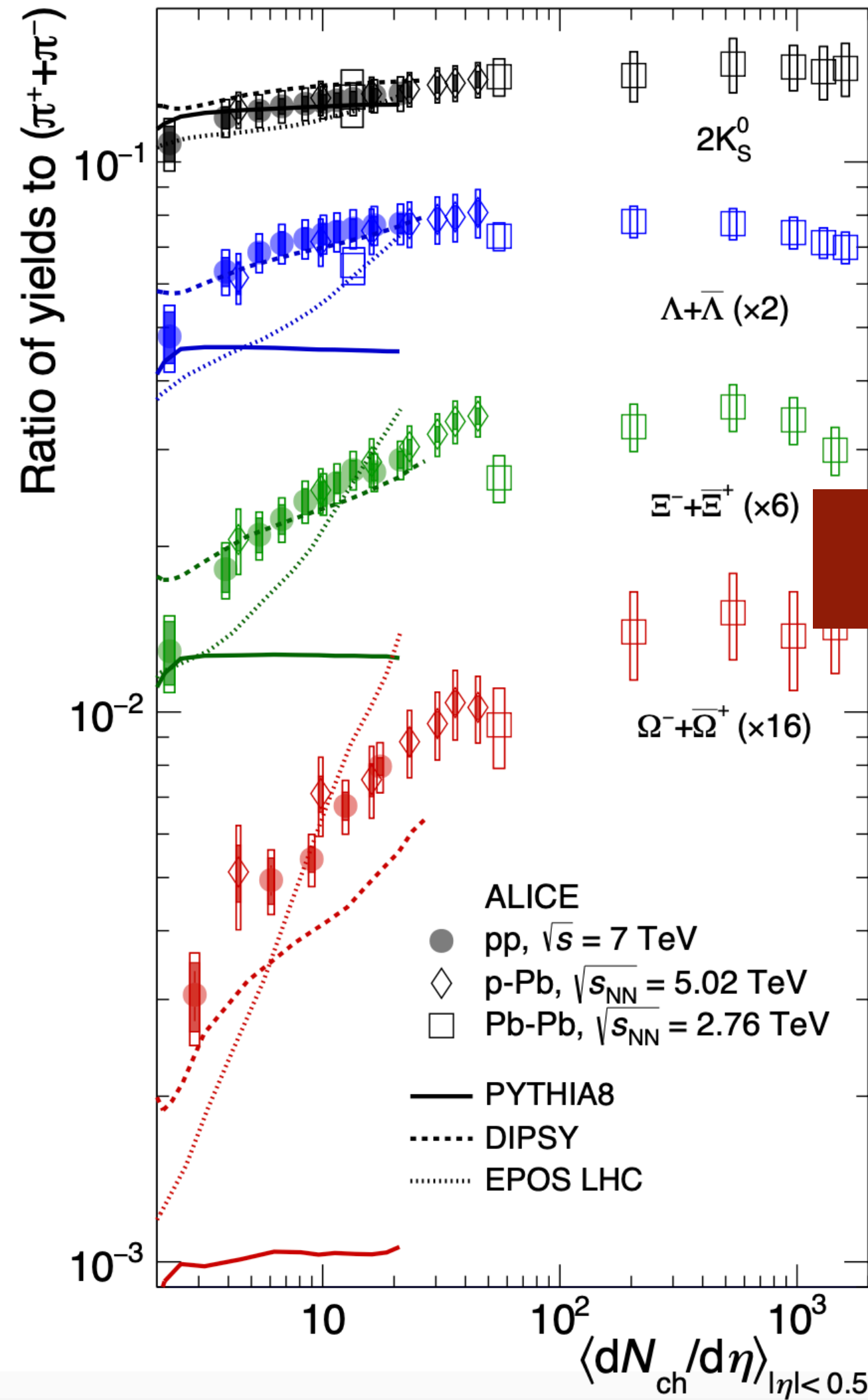
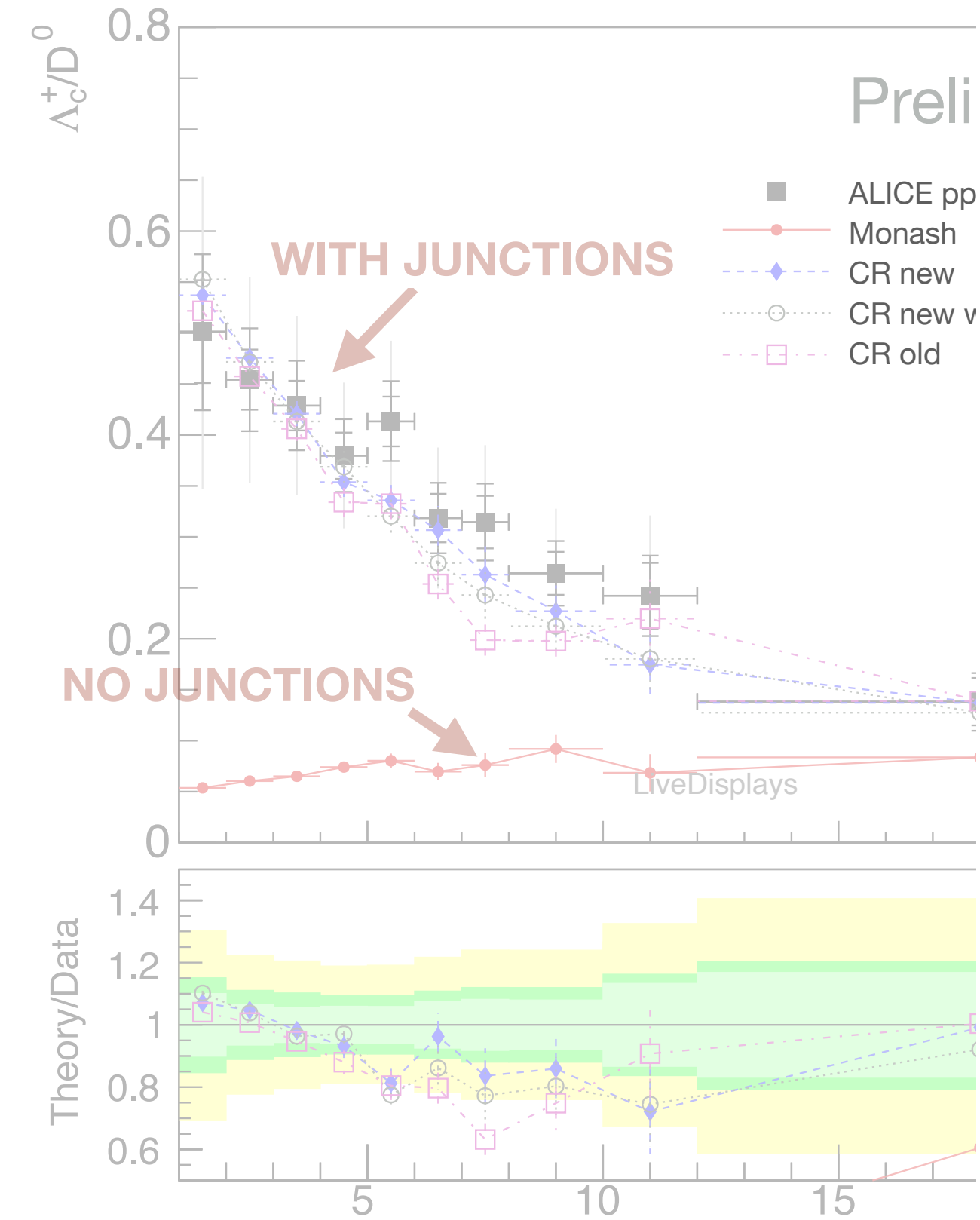
arXiv:2106.08278



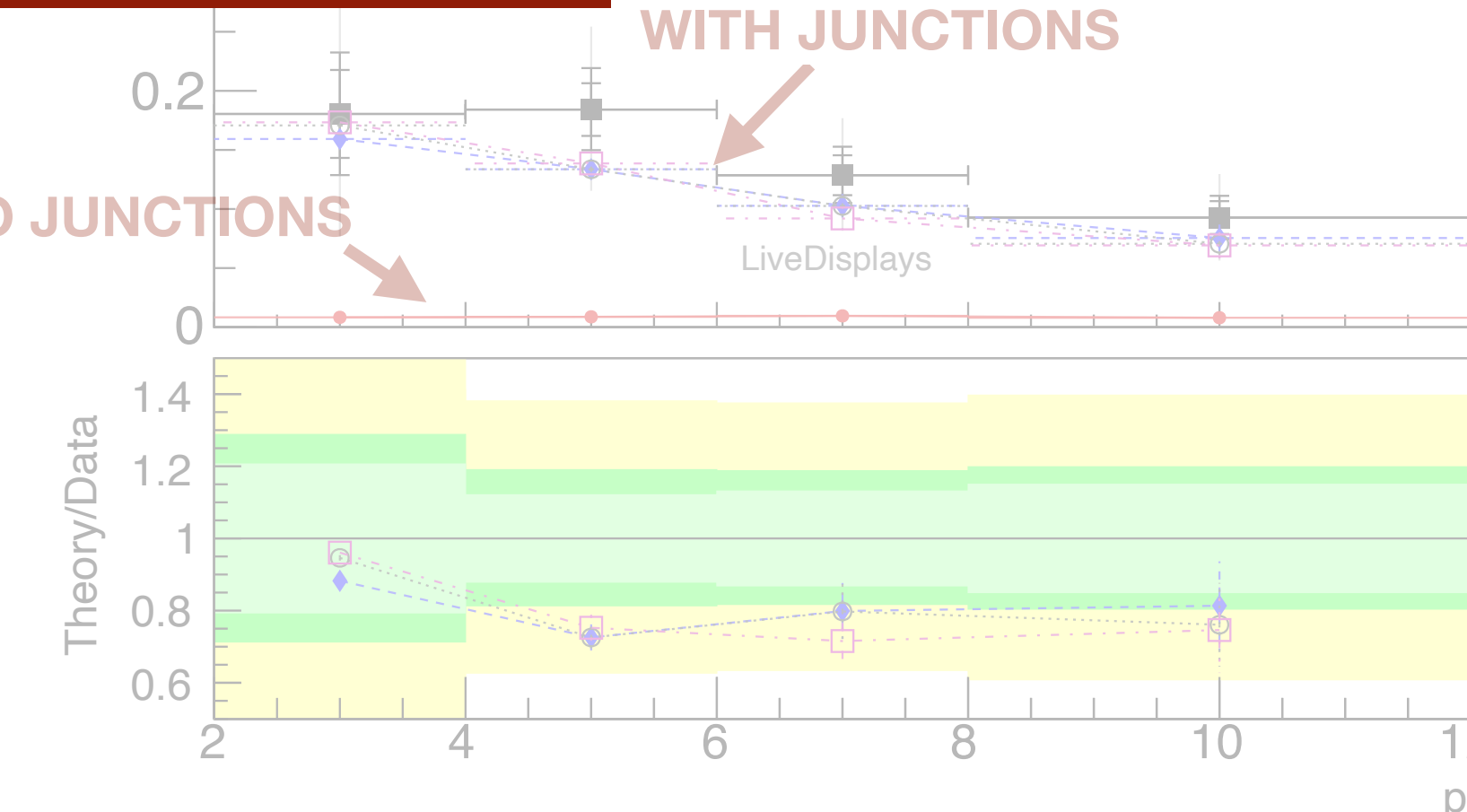
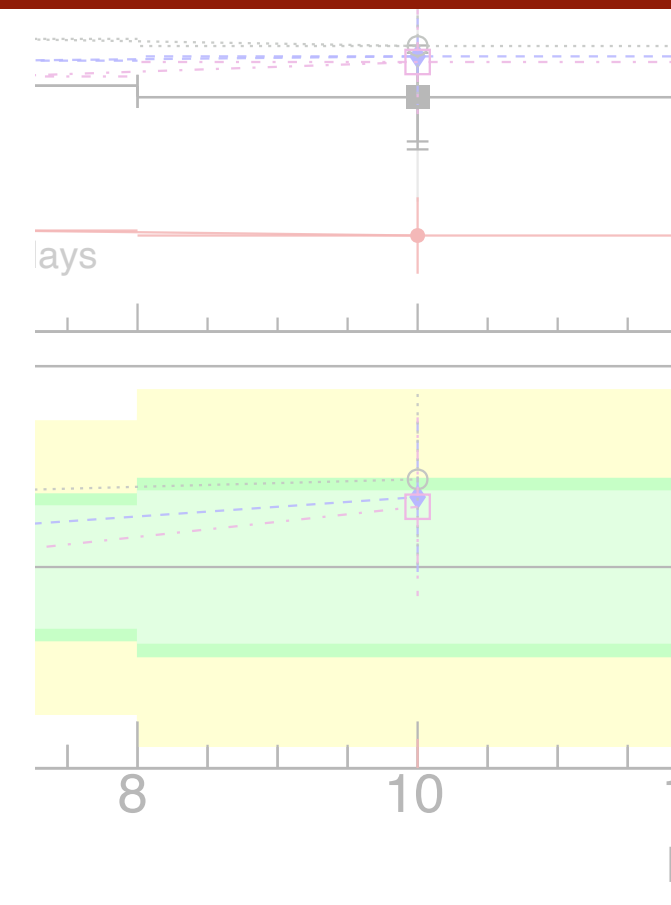
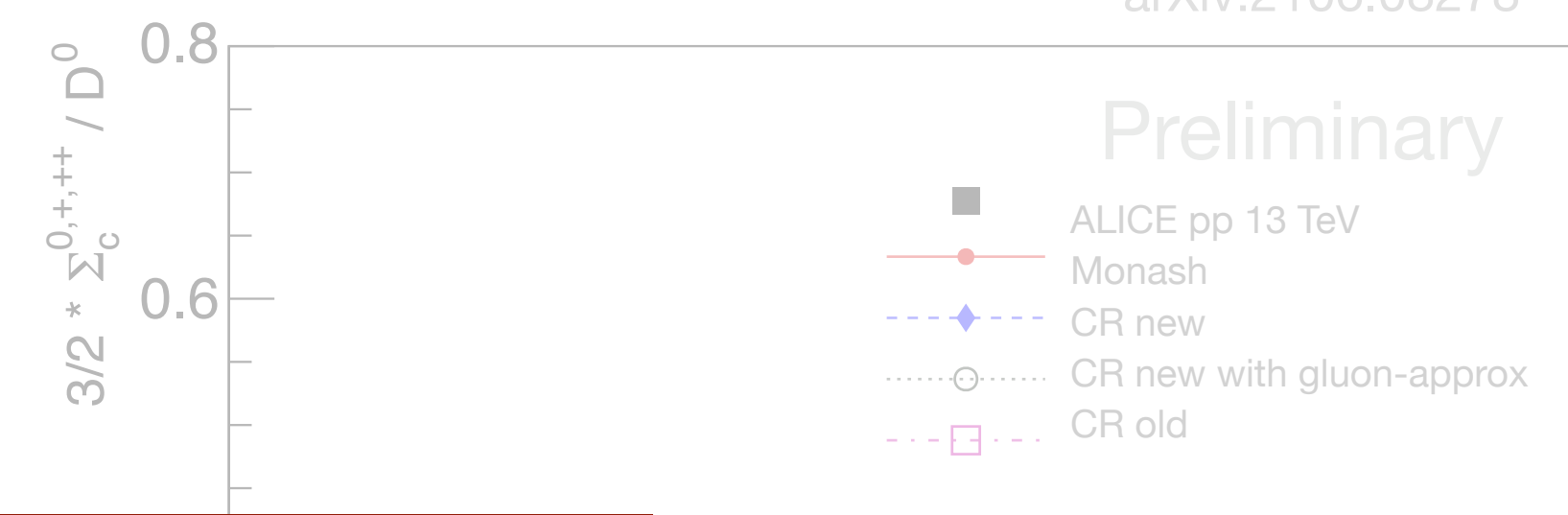
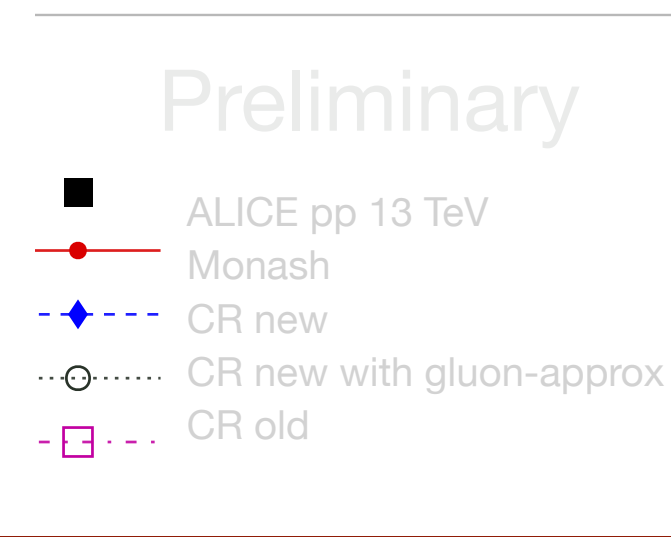
What about strange baryons?

Junctions

arXiv:2106.08278



What about strange baryons?



Clear observations of strangeness enhancement with respect to charged multiplicity [e.g. ALICE Nature Phys. 13, 535 (2017)]

Colour Reconnections

Recent brief review on CR arXiv:2405.19137

Starting point for Monte Carlo is **leading colour** $N_C \rightarrow \infty$ i.e. unique colour singlet configurations determined by colour tracing in hard processes

CR restores missing colour correlations from $SU(3)$ assuming **string “length” minimisation**

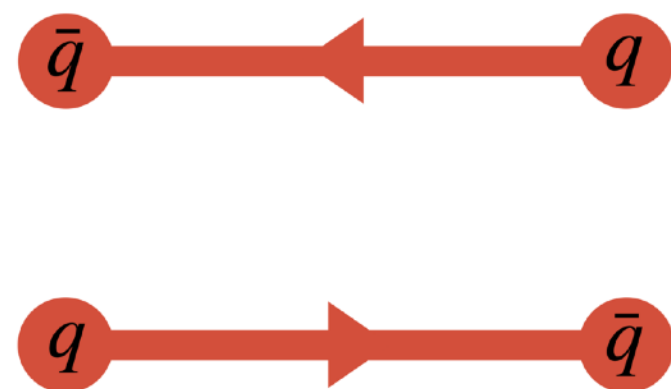
Aims to **stochastically restore** these colour correlations using $SU(3)$ algebra

$$3 \otimes \bar{3} = 8 \oplus 1 \text{ (colour-anticolour)}$$

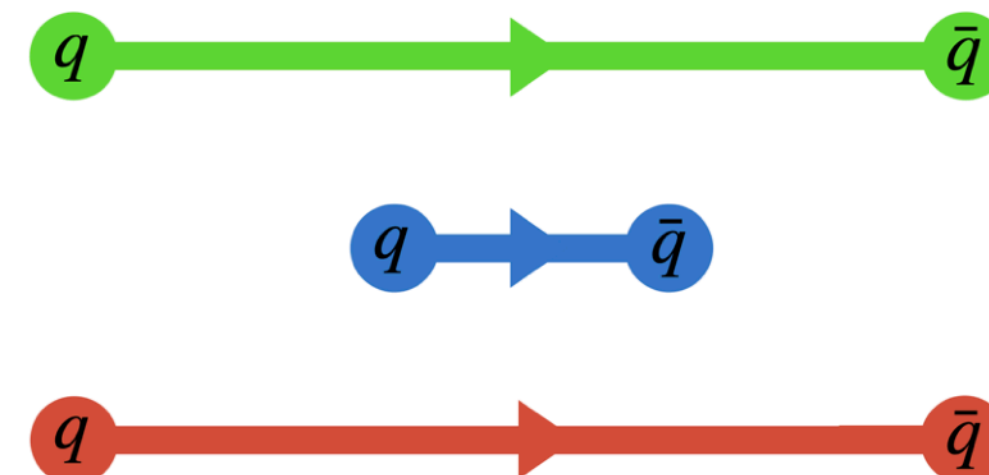
$$3 \otimes 3 = 6 \oplus \bar{3} \text{ (colour-colour)}$$

“string length” is **not** a spatial measure but measure of approx how many hadrons a string can make e.g. rapidity-type measure or invariant mass of the dipole

Dipole-type reconnection: colour-anticolour



Junction-type reconnection: red-green-blue



Colour Reconnections

Recent brief review on CR arXiv:2405.19137

Starting point for Monte Carlo is **leading colour** $N_C \rightarrow \infty$ i.e. unique colour singlet configurations determined by colour tracing in hard processes

CR restores missing colour correlations from $SU(3)$ assuming **string “length” minimisation**

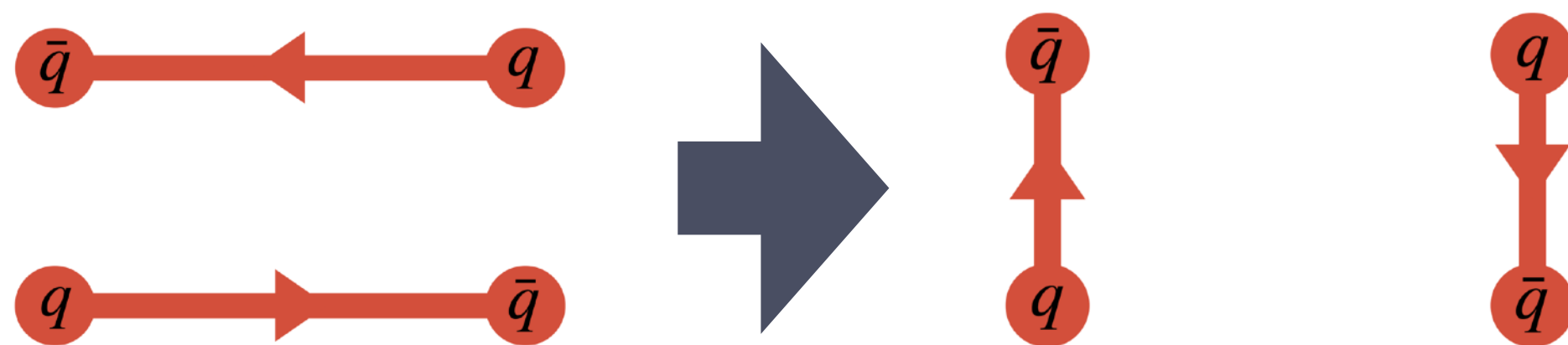
Aims to **stochastically restore** these colour correlations using $SU(3)$ algebra

$$3 \otimes \bar{3} = 8 \oplus 1 \text{ (colour-anticolour)}$$

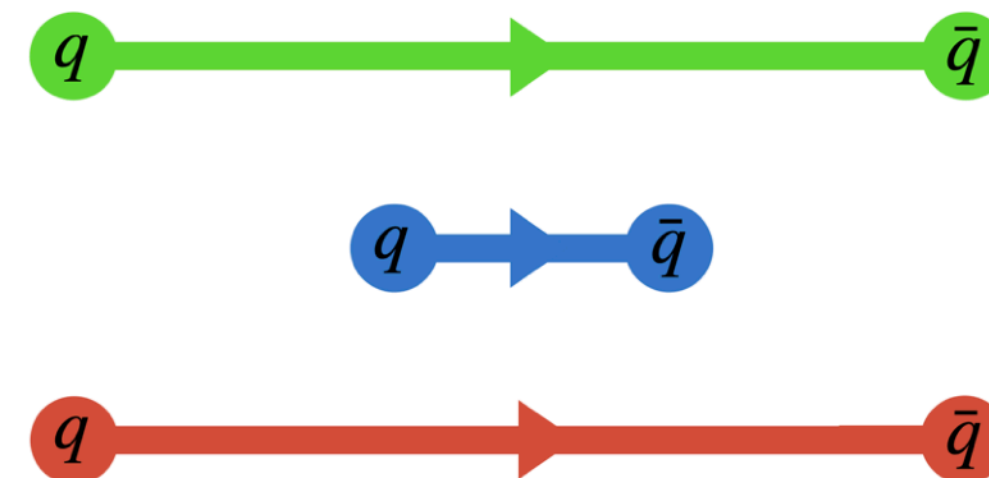
$$3 \otimes 3 = 6 \oplus \bar{3} \text{ (colour-colour)}$$

“string length” is **not** a spatial measure but measure of approx how many hadrons a string can make e.g. rapidity-type measure or invariant mass of the dipole

Dipole-type reconnection: colour-anticolour



Junction-type reconnection: red-green-blue



Independently hadronising MPI does not result in **increasing** $\langle p_{\perp} \rangle$ with multiplicity

Colour Reconnections

Recent brief review on CR arXiv:2405.19137

Starting point for Monte Carlo is **leading colour** $N_C \rightarrow \infty$ i.e. unique colour singlet configurations determined by colour tracing in hard processes

CR restores missing colour correlations from $SU(3)$ assuming **string “length” minimisation**

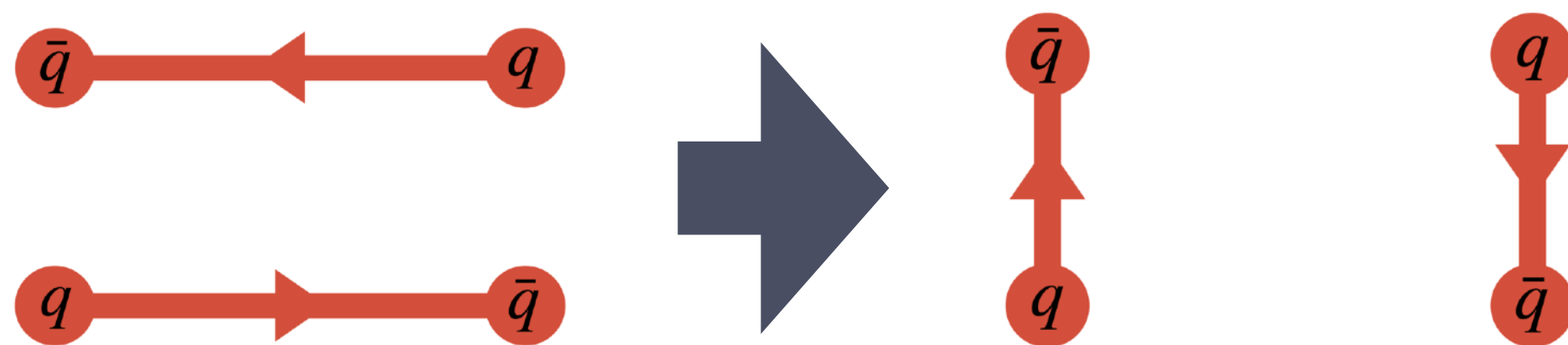
Aims to **stochastically restore** these colour correlations using $SU(3)$ algebra

$$3 \otimes \bar{3} = 8 \oplus 1 \text{ (colour-anticolour)}$$

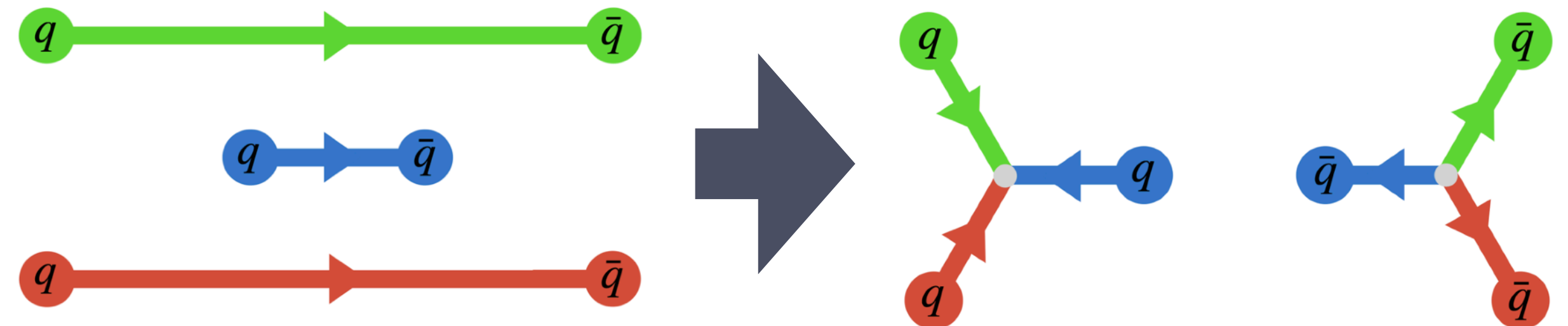
$$3 \otimes 3 = 6 \oplus \bar{3} \text{ (colour-colour)}$$

“string length” is **not** a spatial measure but measure of approx how many hadrons a string can make e.g. rapidity-type measure or invariant mass of the dipole

Dipole-type reconnection: colour-anticolour



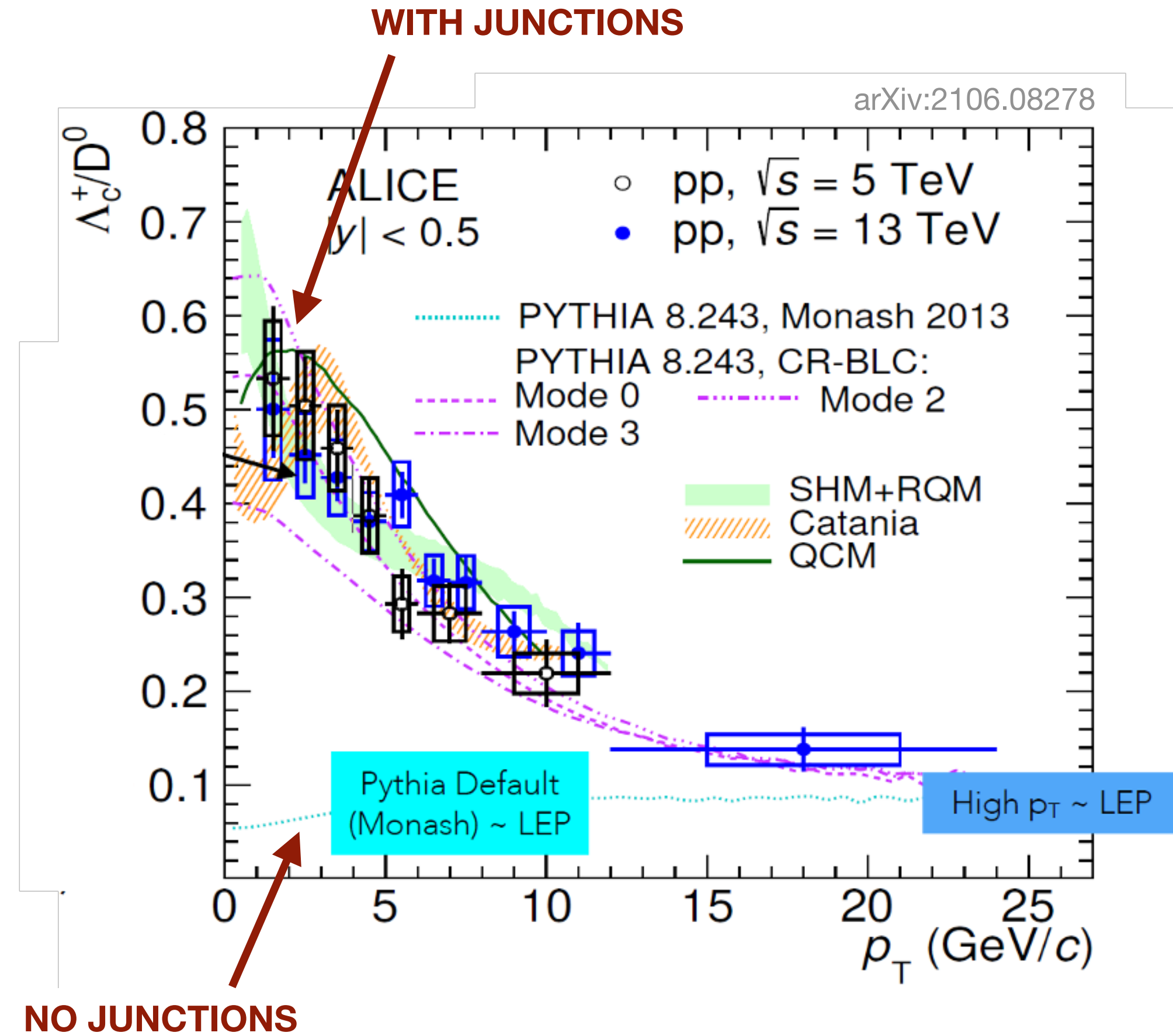
Junction-type reconnection: red-green-blue



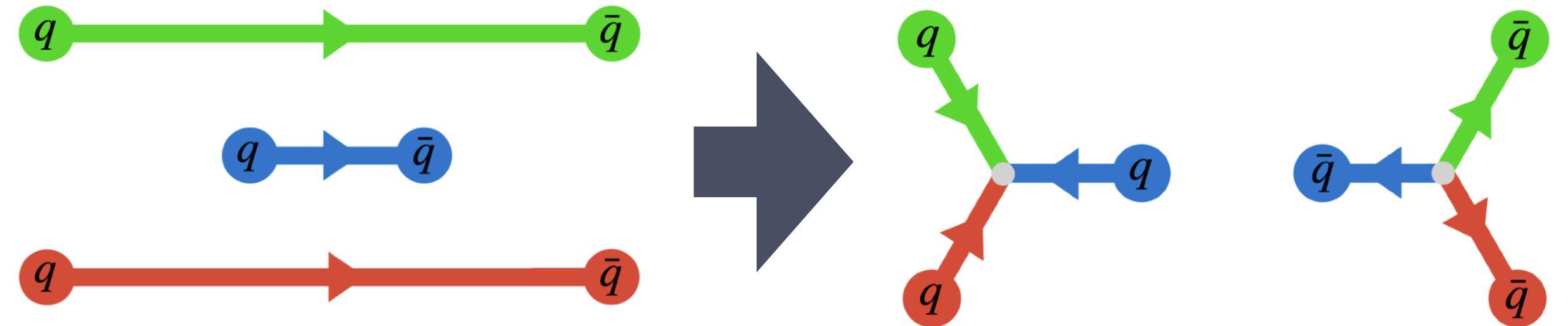
Independently hadronising MPI does not result in **increasing** $\langle p_{\perp} \rangle$ with multiplicity

Colour Reconnections

Recent brief review on CR arXiv:2405.19137



Junction-type reconnection: red-green-blue

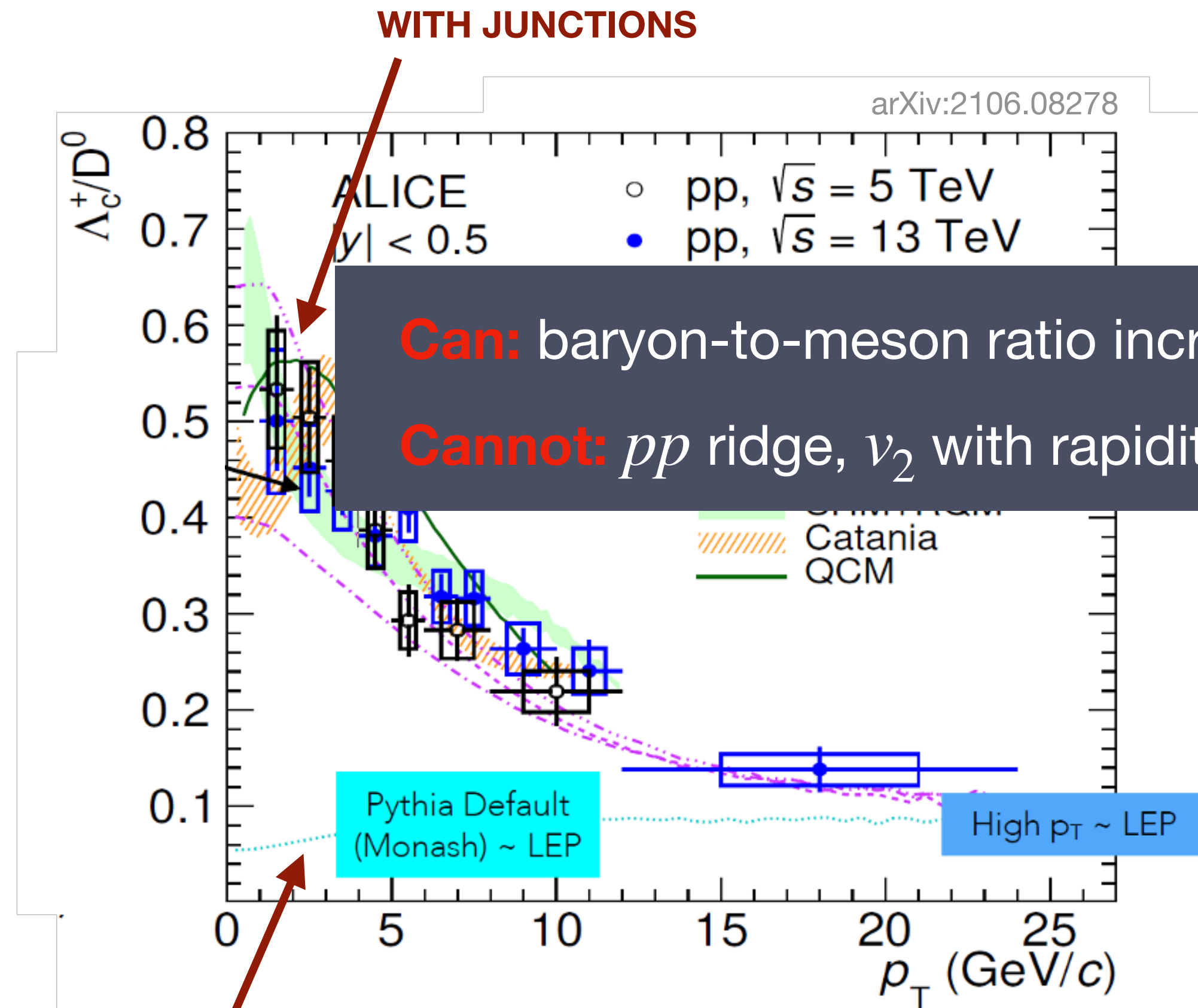


Independently hadronising MPI does not result in **increasing $\langle p_{\perp} \rangle$ with multiplicity**

Junctions result in baryons \rightarrow increase in **baryon-to-meson ratio**

Colour Reconnections

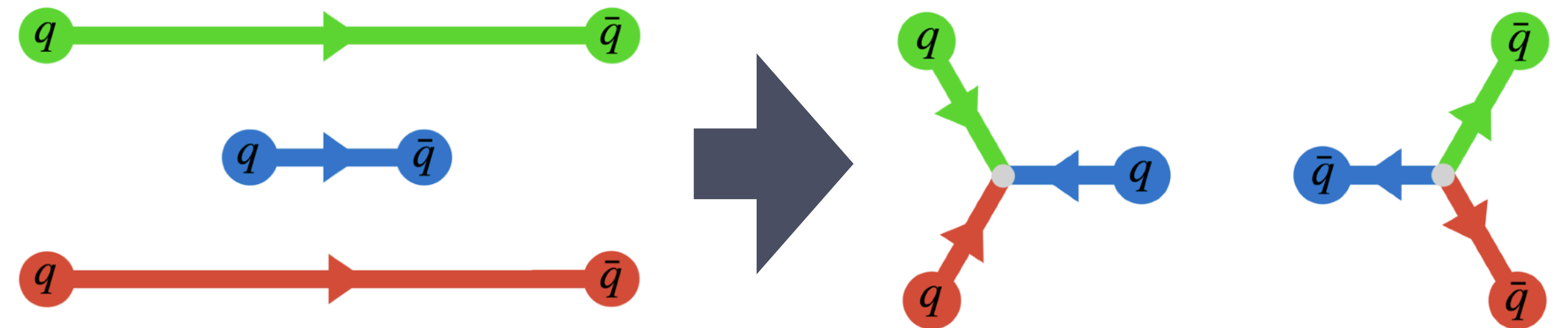
Recent brief review on CR arXiv:2405.19137



Can: baryon-to-meson ratio increase, $\langle p_{\perp} \rangle$ increase with multiplicity, some flow-like effects

Cannot: pp ridge, v_2 with rapidity gap between particles of interest

Junction-type reconnection: red-green-blue



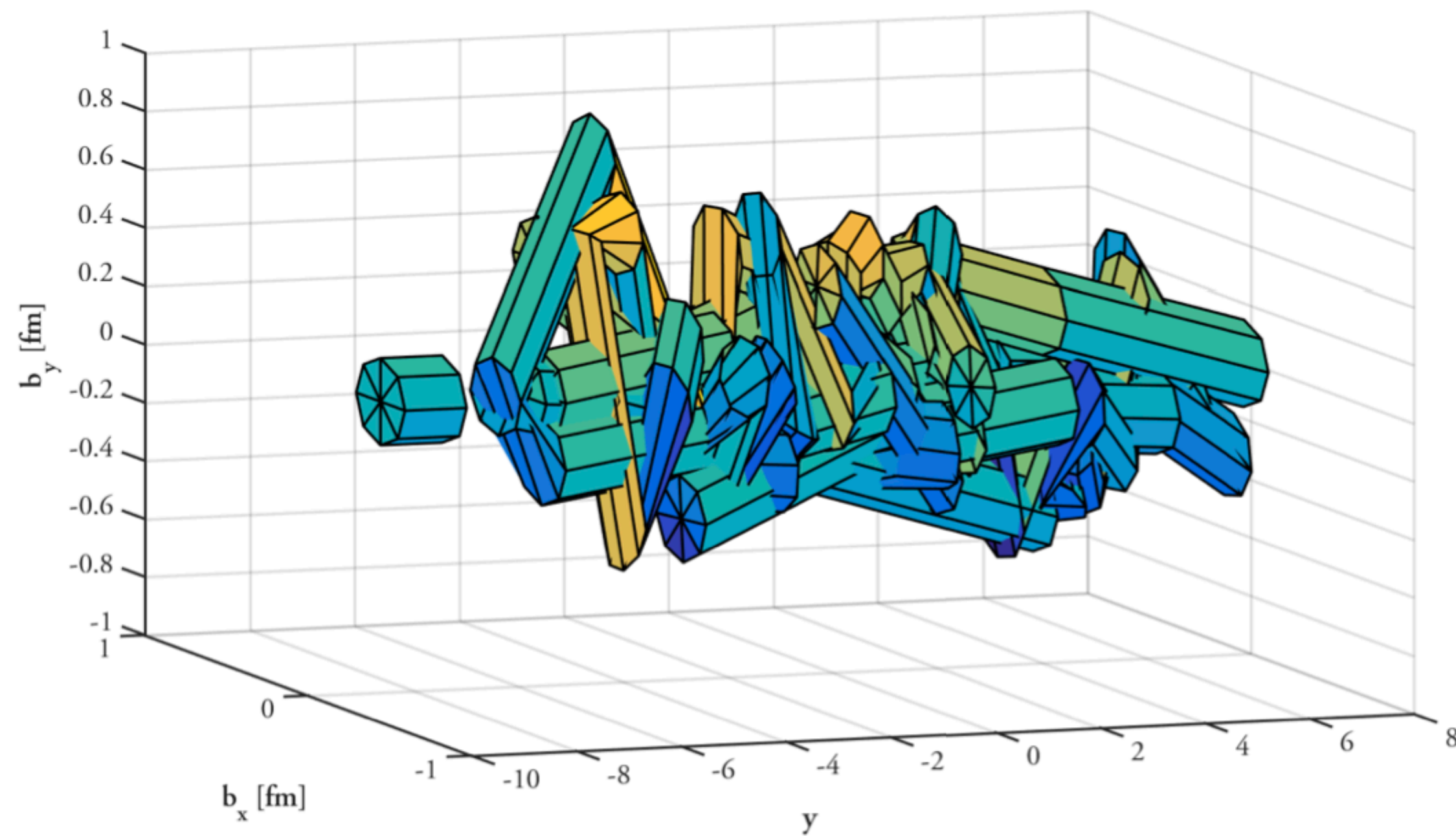
Independently hadronising MPI does not result in **increasing $\langle p_{\perp} \rangle$ with multiplicity**

Junctions result in baryons \rightarrow increase in **baryon-to-meson ratio**

Shoving

arXiv:2010.07595

After the string has had time after its initial creation to expand to its full **transverse size**, strings will start **“shoving”**



e.g. $\sqrt{s} = 7$ TeV collision example

*uses string radius of 0.2 fm for illustration purposes but in reality can be much larger

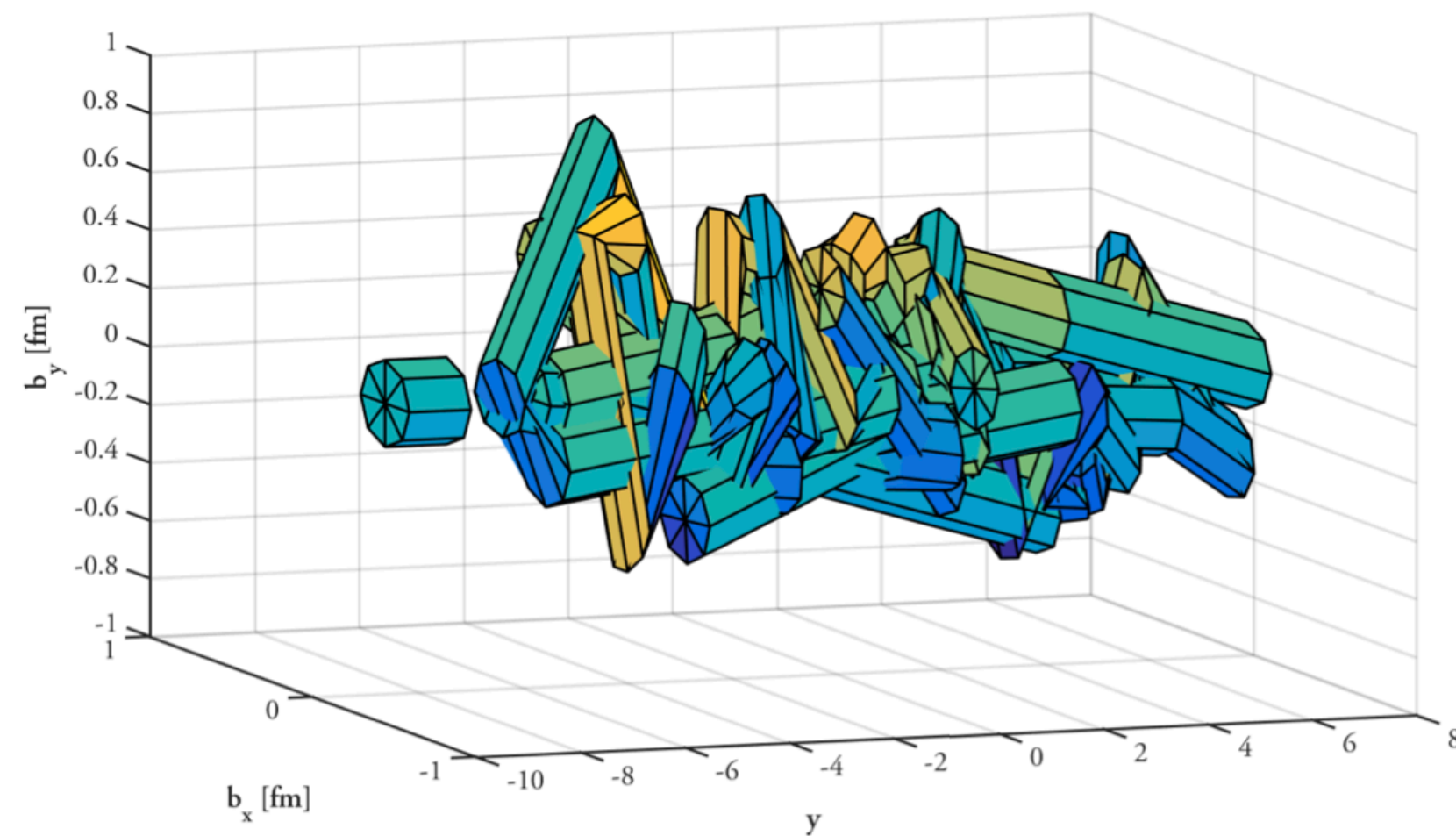
Shoving

After the string has had time after its initial creation to expand to its full **transverse size**, strings will start **“shoving”**

CR has already occurred with string minimisation choosing singlet configurations

→ only octet states would likely be near one another

→ only **repulsion** left



e.g. $\sqrt{s} = 7$ TeV collision example

*uses string radius of 0.2 fm for illustration purposes but in reality can be much larger

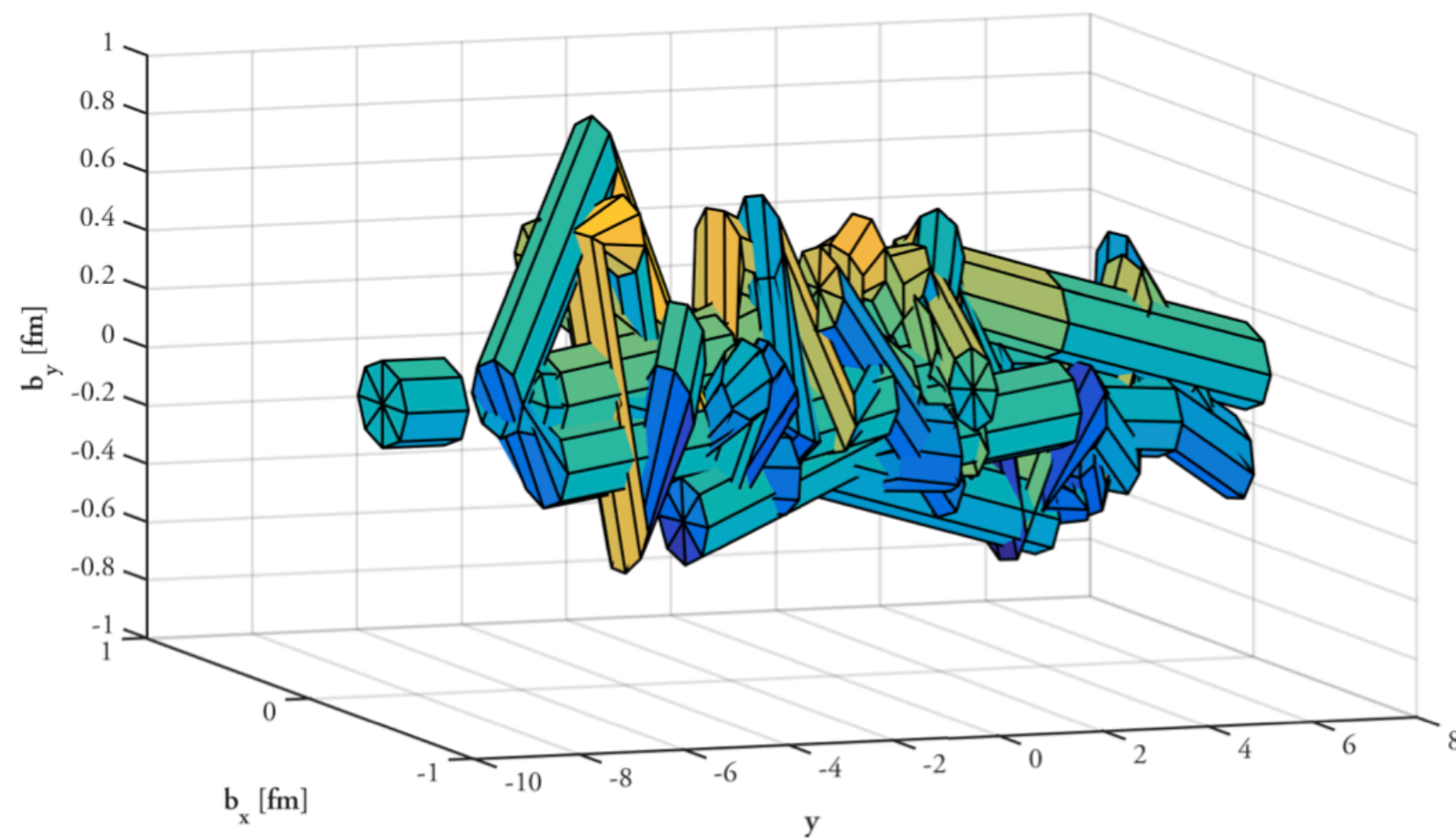
Shoving

After the string has had time after its initial creation to expand to its full **transverse size**, strings will start **“shoving”**

CR has already occurred with string minimisation choosing singlet configurations

→ only octet states would likely be near one another

→ only **repulsion** left



e.g. $\sqrt{s} = 7$ TeV collision example

*uses string radius of 0.2 fm for illustration purposes but in reality can be much larger

Force calculable from the field $E = N \exp(-\rho^2/2R^2)$

Energy per unit length of two strings

overlapping
$$\int d^2\rho \frac{(E_1 + E_2)^2}{2}$$

Force between two strings transversely separated

by d_{\perp} is then
$$f(d_{\perp}) = \frac{g\kappa d_{\perp}}{R^2} \exp\left(-\frac{d_{\perp}^2}{4R^2}\right)$$

ρ is the radius in cylindrical coordinates

R is the equilibrium radius

N is a normalization factor, determined by letting the energy in the field correspond to a **fraction g of the total string tension.**

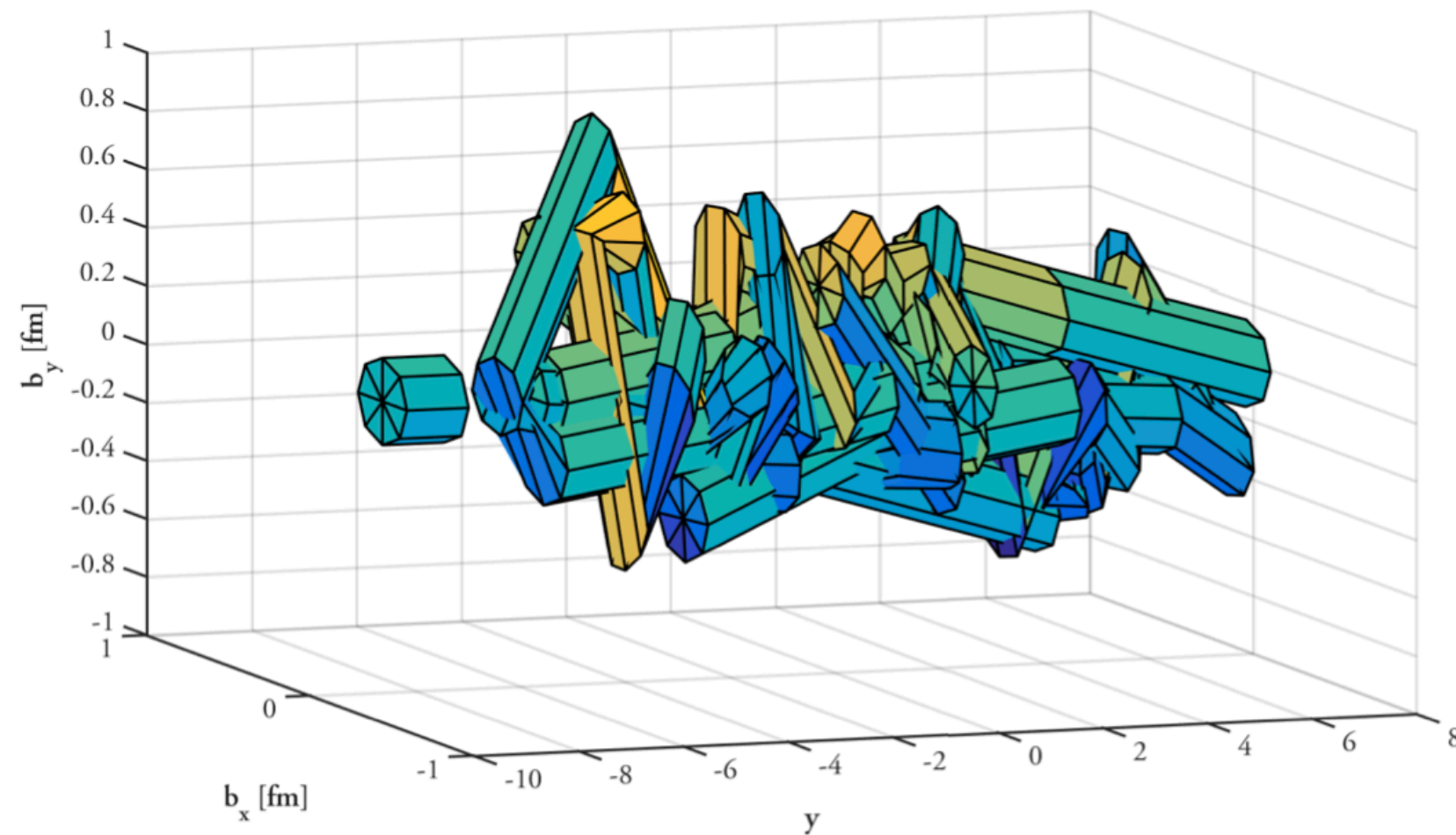
Shoving

After the string has had time after its initial creation to expand to its full **transverse size**, strings will start **“shoving”**

CR has already occurred with string minimisation choosing singlet configurations

→ only octet states would likely be near one another

→ only **repulsion** left



e.g. $\sqrt{s} = 7$ TeV collision example

*uses string radius of 0.2 fm for illustration purposes but in reality can be much larger

Force calculable from the field $E = N \exp(-\rho^2/2R^2)$

Energy per unit length of two strings

overlapping $\int d^2\rho \frac{(E_1 + E_2)^2}{2}$

Force between two strings transversely separated

by d_{\perp} is then $f(d_{\perp}) = \frac{g\kappa d_{\perp}}{R^2} \exp\left(-\frac{d_{\perp}^2}{4R^2}\right)$

ρ is the radius in cylindrical coordinates

R is the equilibrium radius

N is a normalization factor, determined by letting the energy in the field correspond to a **fraction g of the total string tension.**

Monte Carlo implementation details

- Use parallel dogbone frame
- Ordered in p_{\perp} in similar spirit to parton shower ordering

Requires space-time picture of strings

Shoving in pp

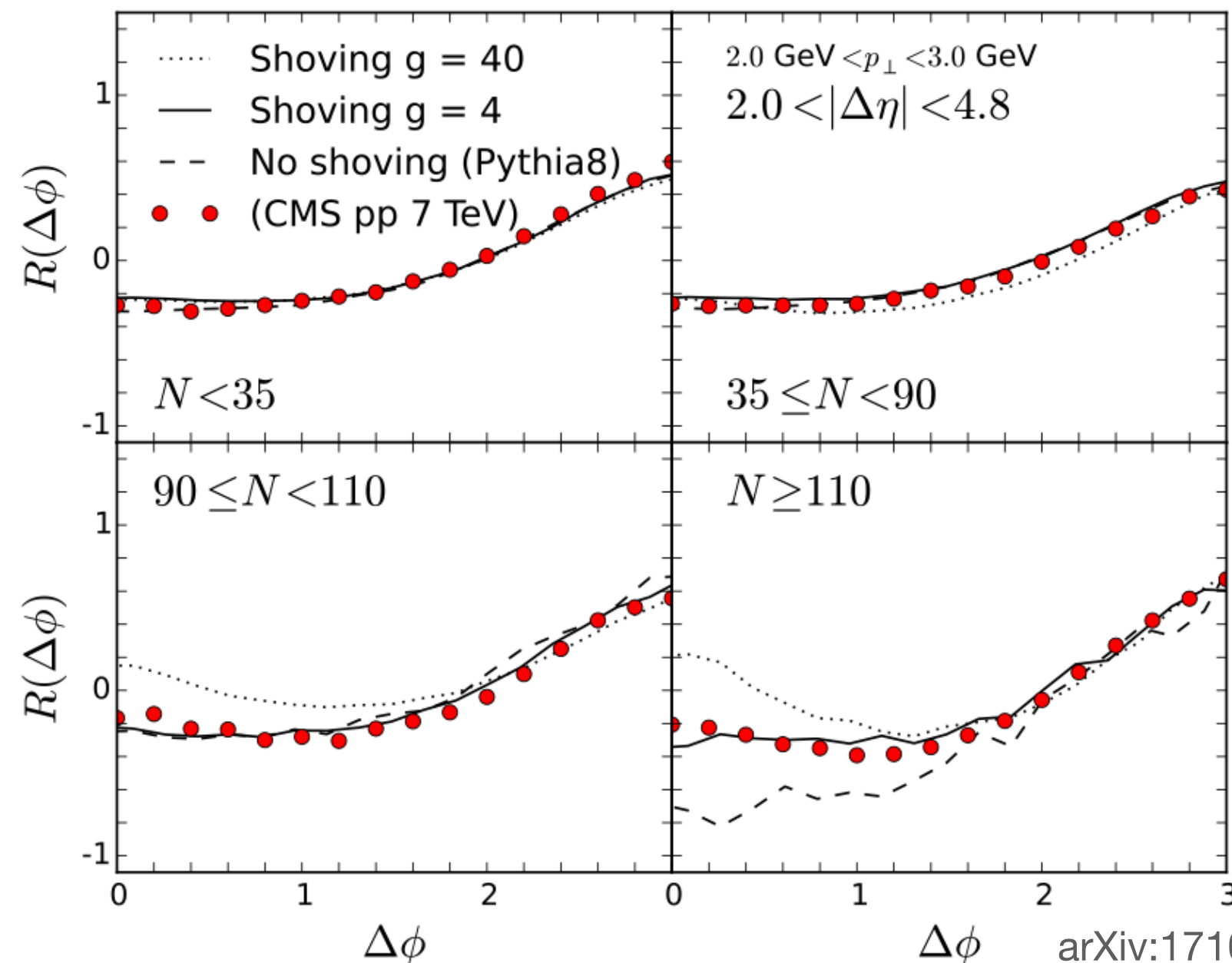
After the string has had time after its initial creation to expand to its full **transverse size**, strings will start **“shoving”**

CR has already occurred with string minimisation choosing singlet configurations

→ only octet states would likely be near one another

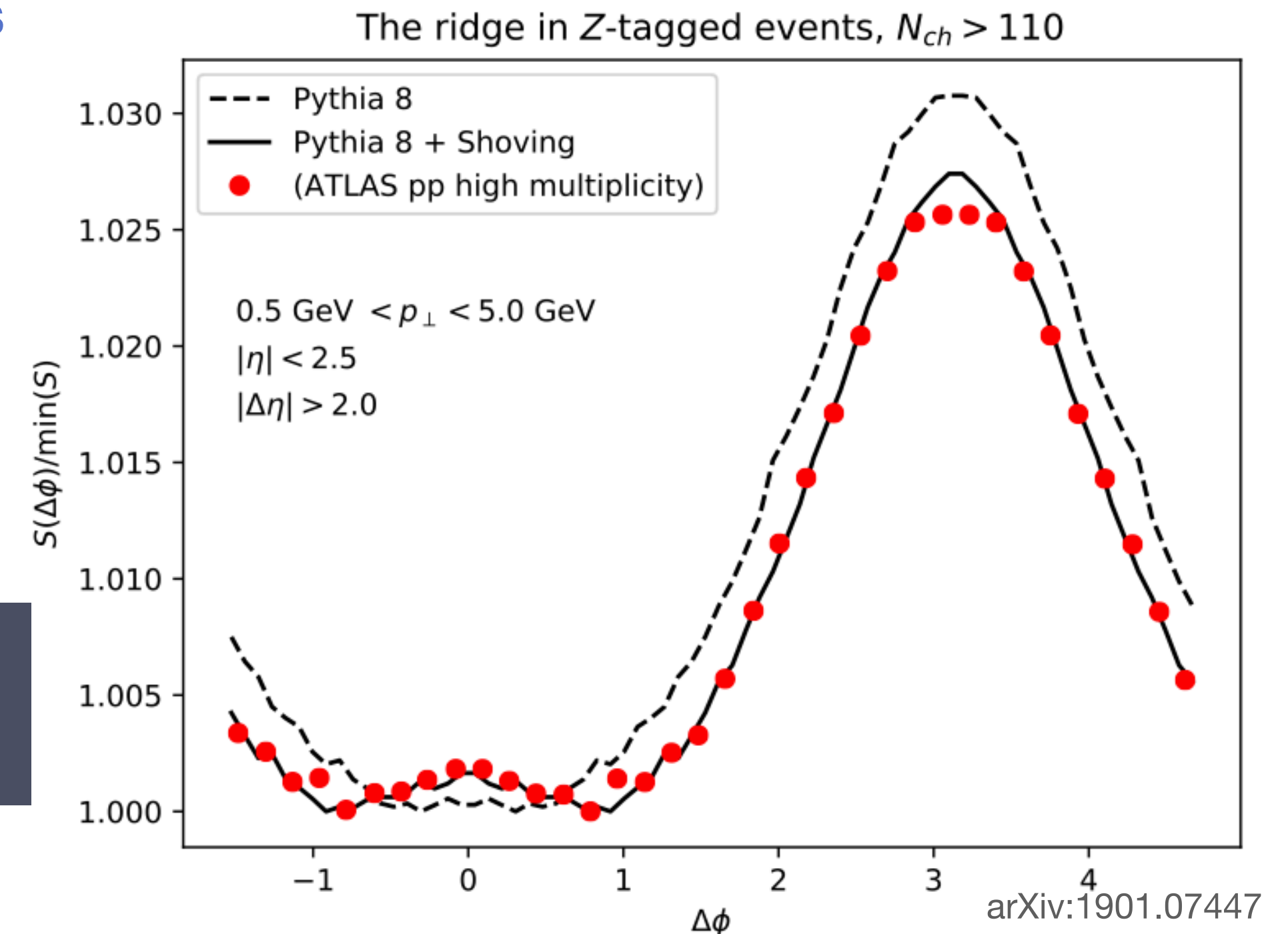
→ only **repulsion** left

Generate anisotropic flow as a response to the **spatial initial conditions**



Z-tagged results - m_Z scale
as the largest hard scale of the collision, which could alter the distribution in p_{\perp} of MPIs w.r.t minimum bias

*model of shoving used here is old implementation that manifests the “shove” as soft gluons



*note g parameter differs by normalisation factor from g in equation

*results from old implementation of shoving w.r.t to the beam axis rather than the dog-bone implementation

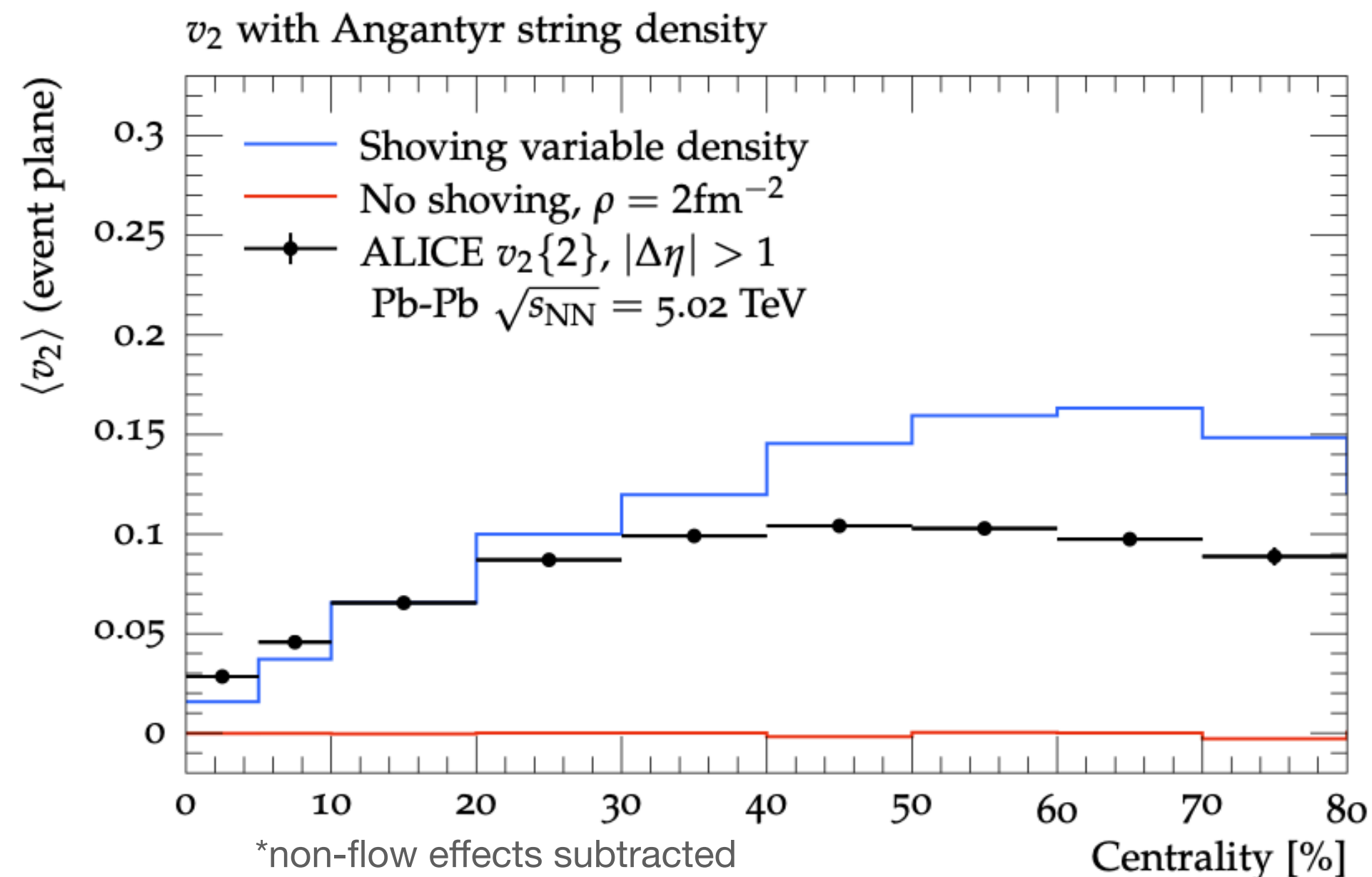
Shoving in AA

First look at **toy case**

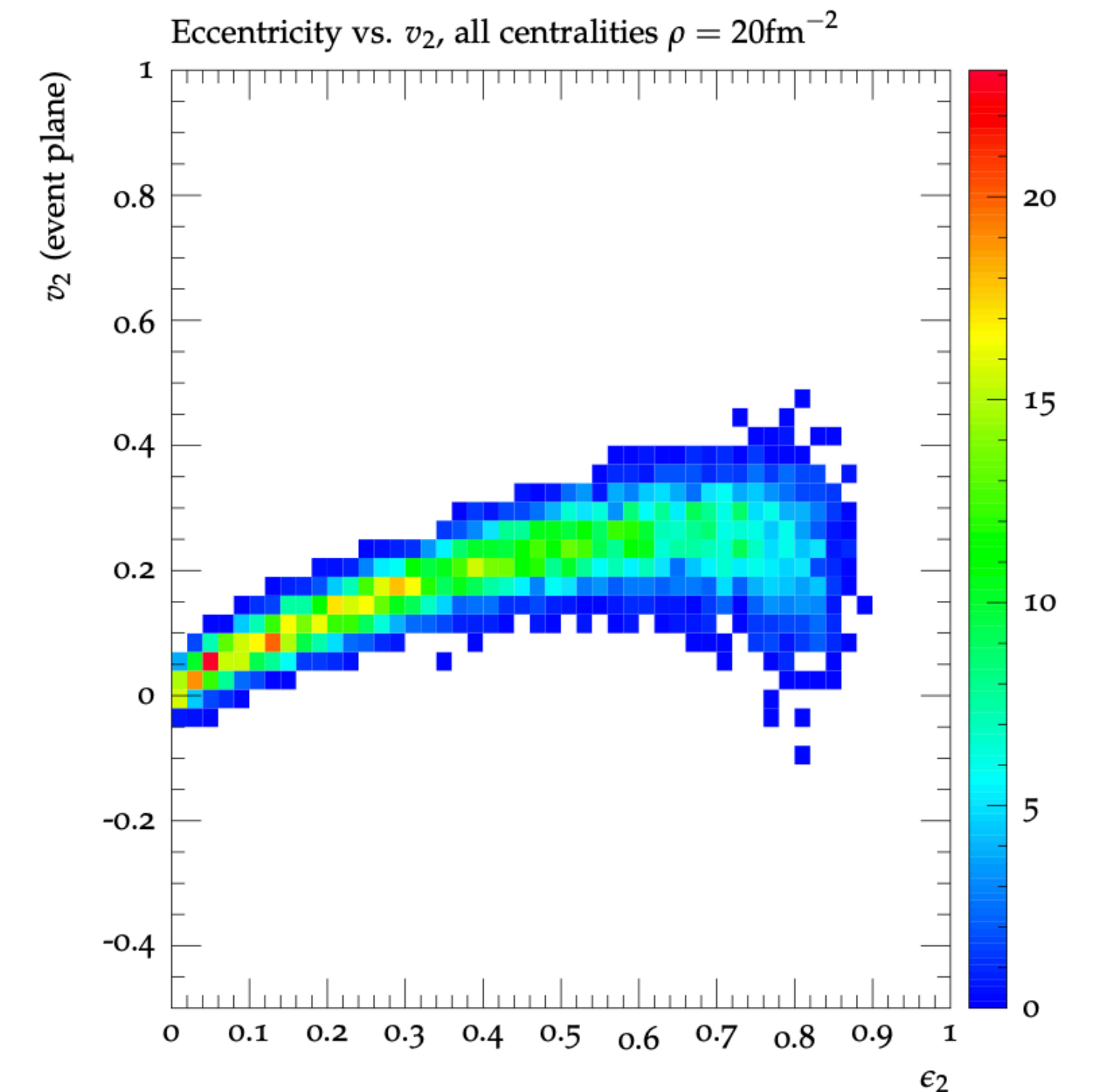
- Multiplicity generated by a single string well known (approx one hadron per unit of rapidity)
- System of **straight strings (no gluon kinks)** that corresponds to the multiplicity of AA collisions in a given centrality interval

Not perfect agreement however is only a toy model and uses **same parameters as pp collision systems**

source of flow can be the same across collision systems!!!



Correlation between initial state ϵ_2 and final state v_2 is linear in hydrodynamic deconfined QGP phase - **similarly with shoving** → **hydrodynamic behaviour is not limited to deconfined systems**



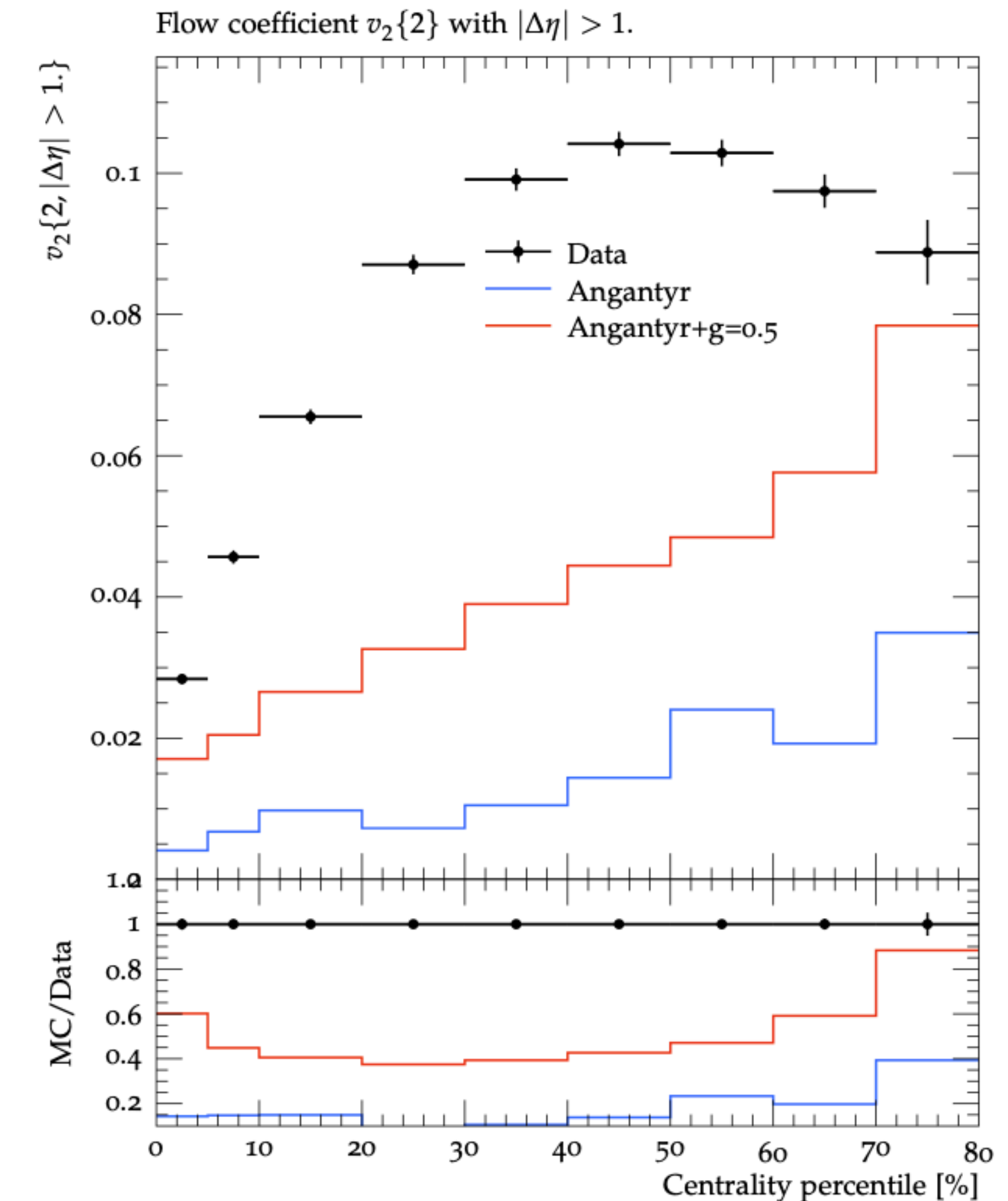
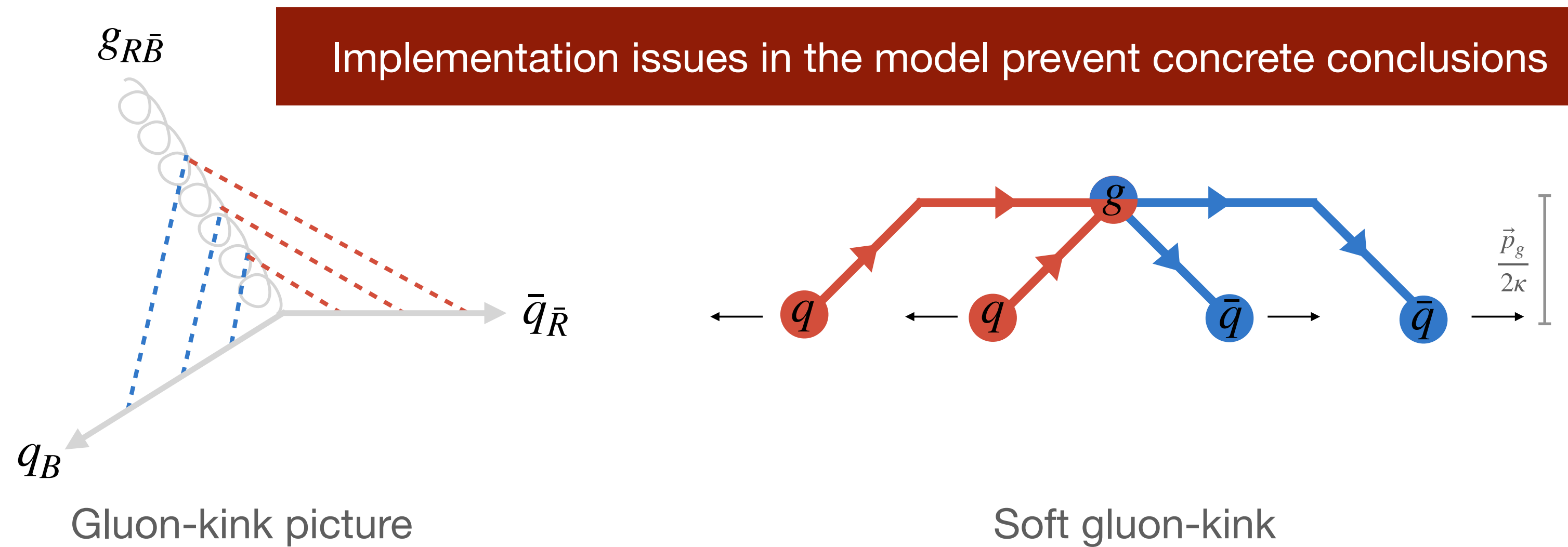
Shoving in AA

Full Pb-Pb collision in Angantyr

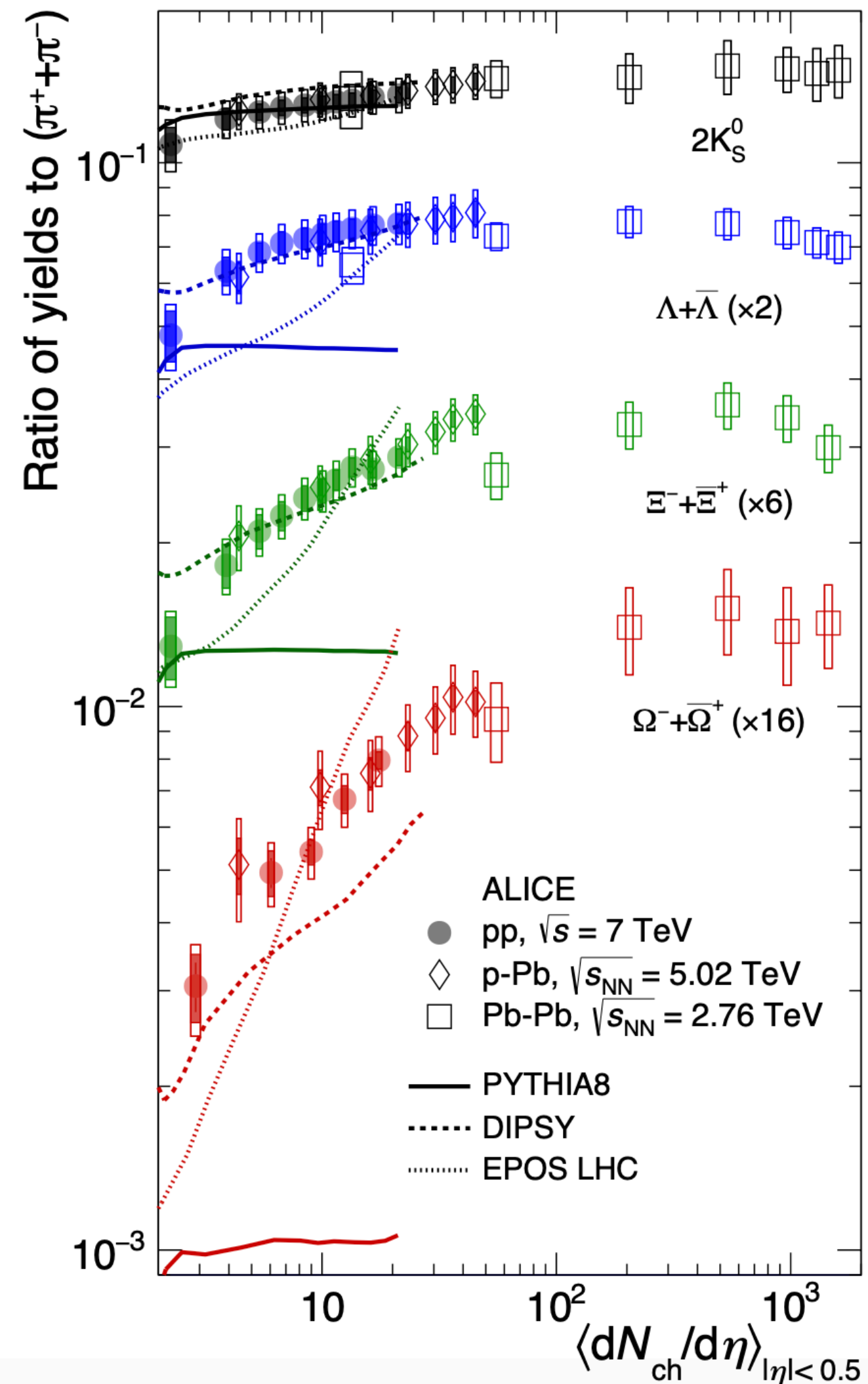
➤ Implementation issues

Many **soft gluons** → **short interaction time** for shoving mechanism as the mechanism does not consider the region formed from soft gluons → insufficient level of shoving

➤ Trend is in the correct direction but insufficient, also lacks curved shape

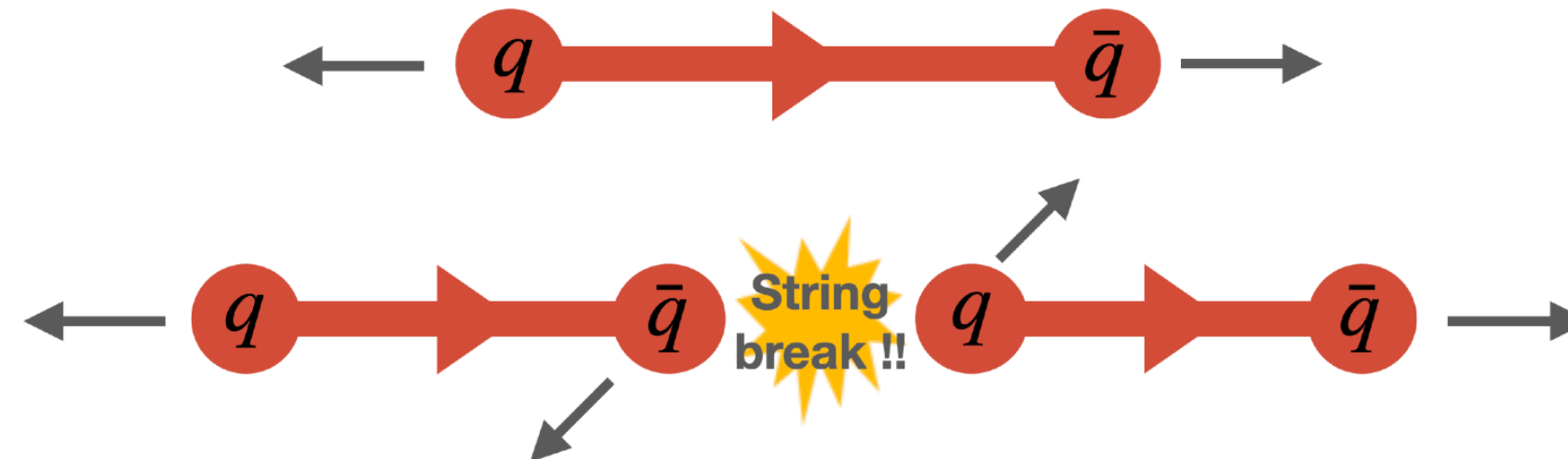


Strangeness Enhancement

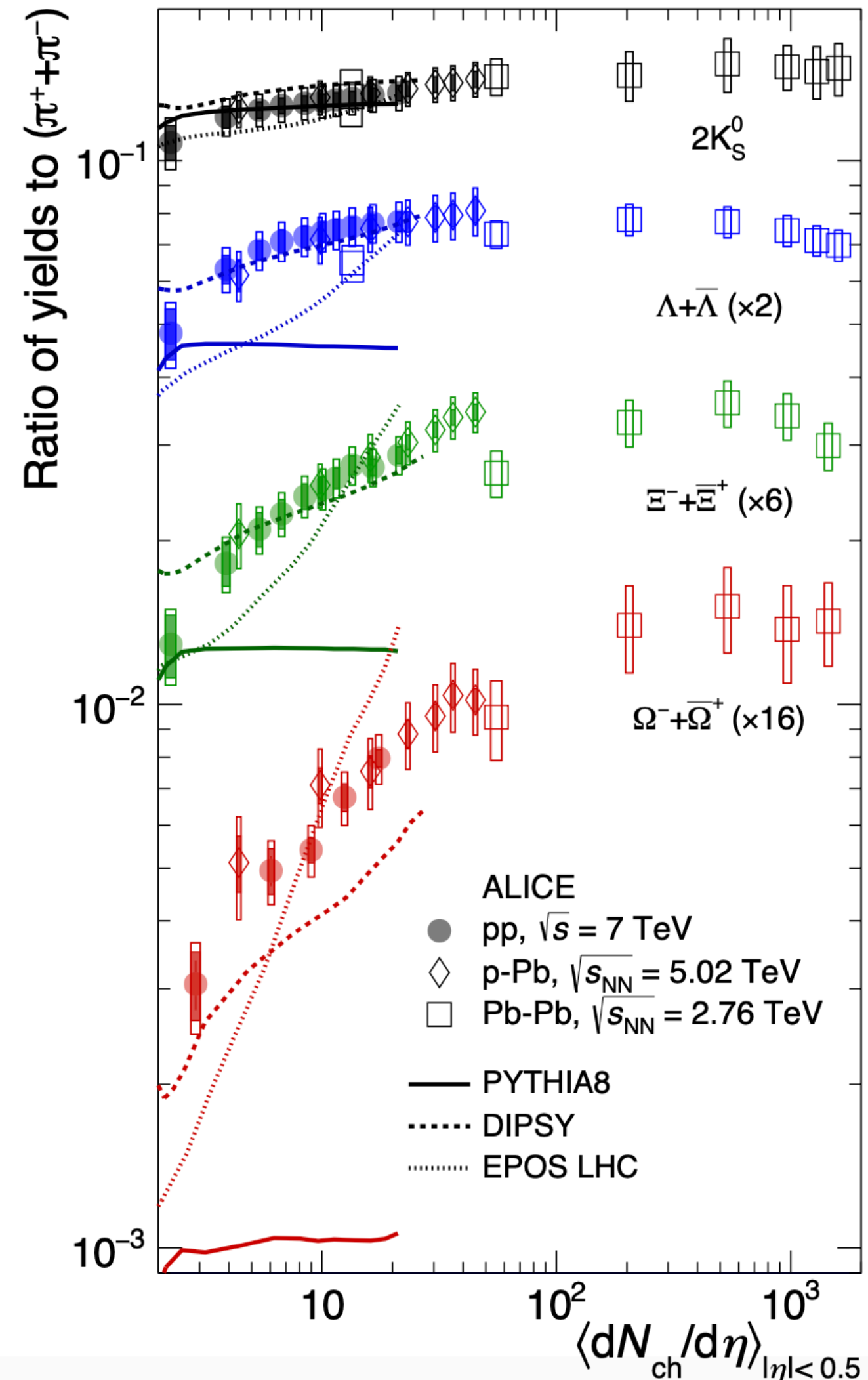


Strange production in the string picture

Use **Schwinger mechanism** to model **tunnelling of quark-antiquark pairs** created by string breaks

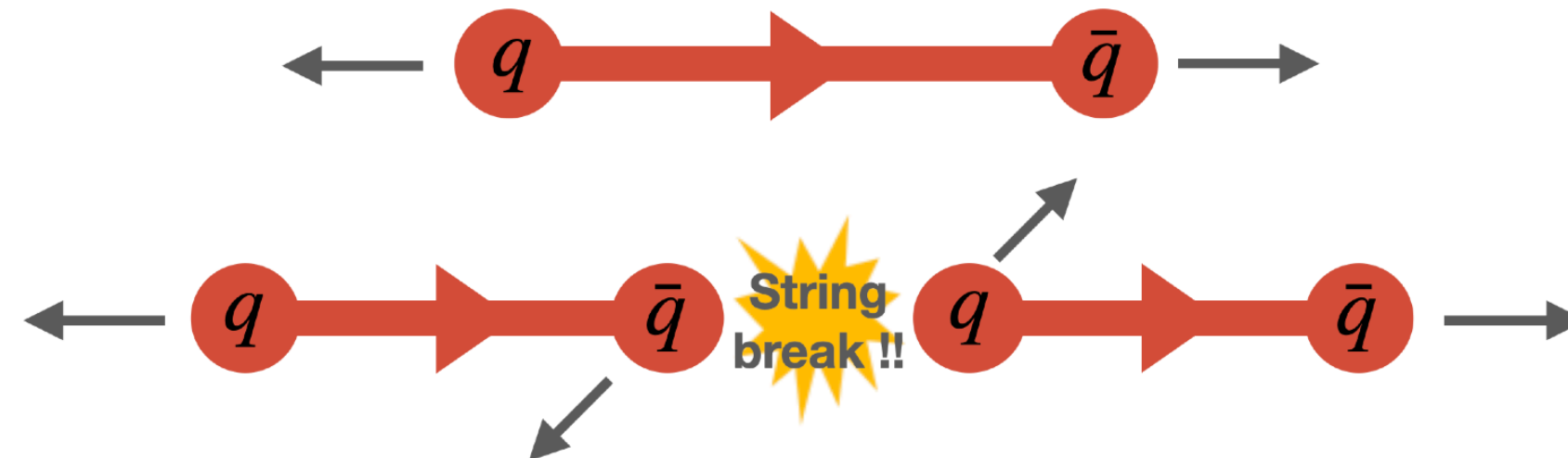


Strangeness Enhancement

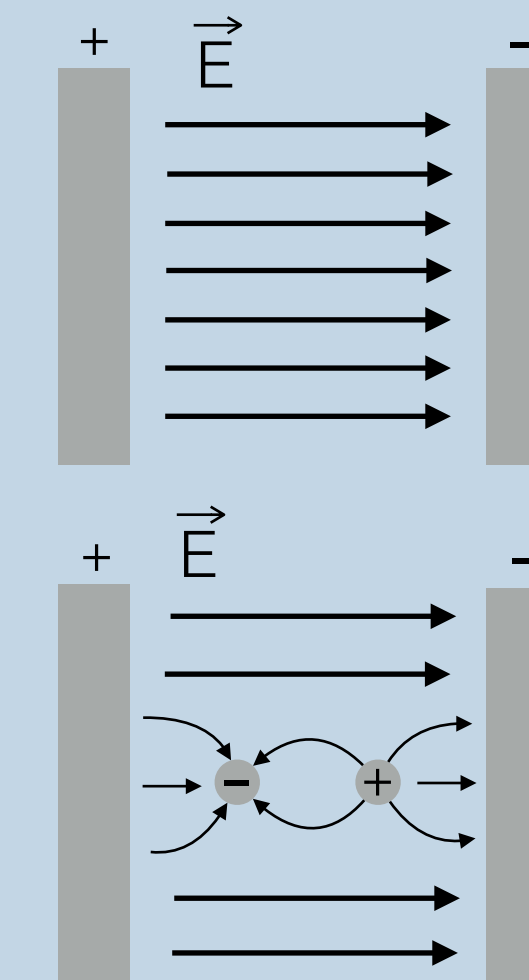


Strange production in the string picture

Use **Schwinger mechanism** to model **tunnelling of quark-antiquark pairs** created by string breaks



Schwinger mechanism QED



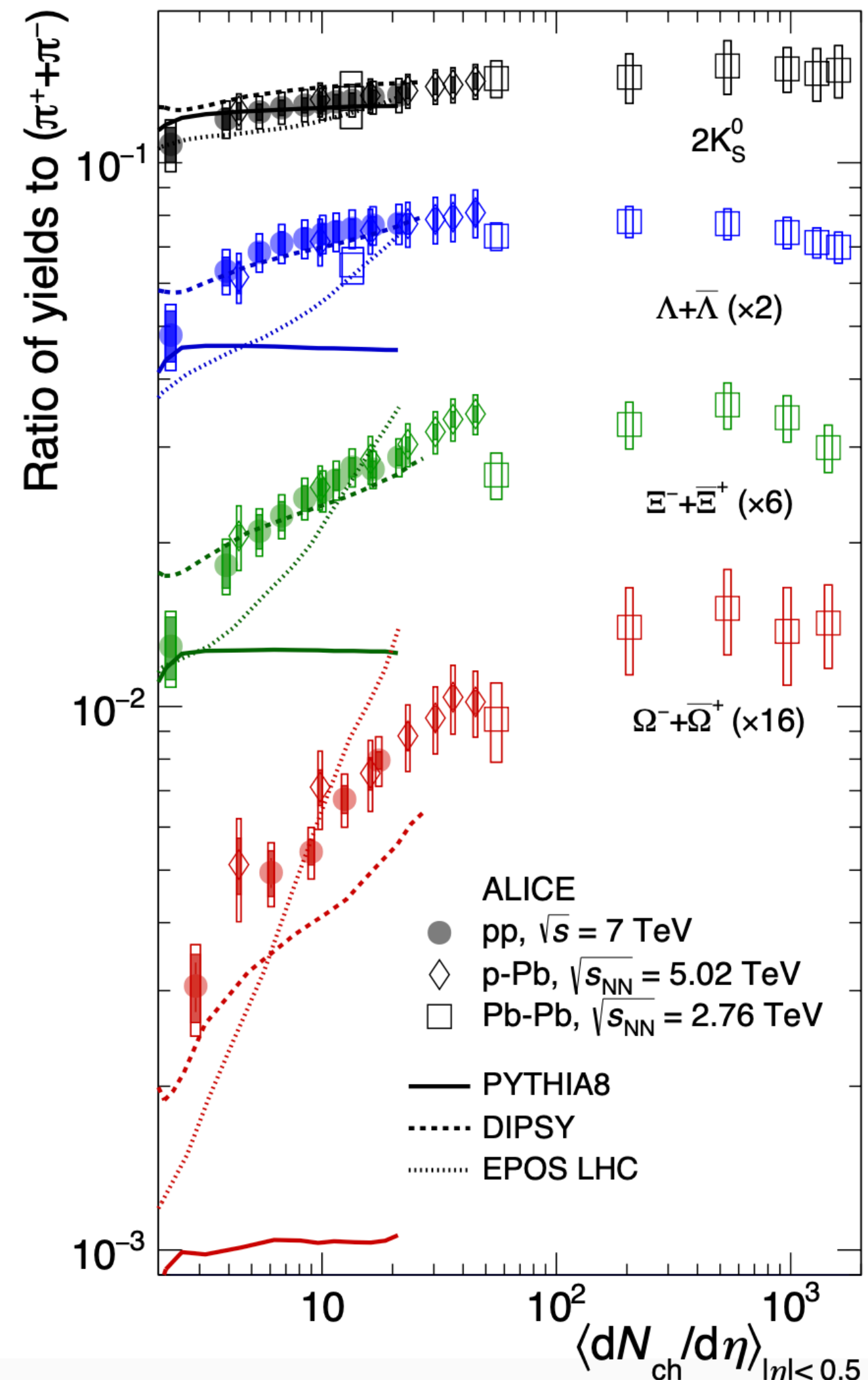
Non-perturbative creation of e^+e^- pairs in a string electric field

Probability from tunnelling factor

$$\mathcal{P} \propto \exp\left(\frac{-m^2 - p_{\perp}^2}{\kappa/\pi}\right)$$

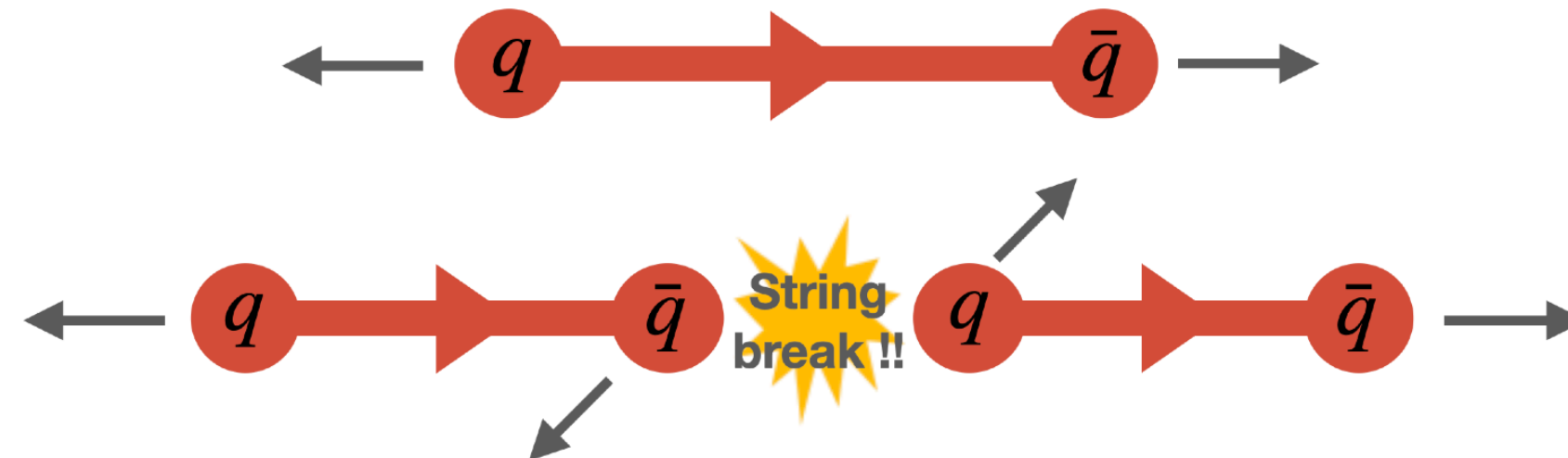
$\kappa = \text{string tension}$

Strangeness Enhancement



Strange production in the string picture

Use **Schwinger mechanism** to model **tunnelling of quark-antiquark pairs** created by string breaks

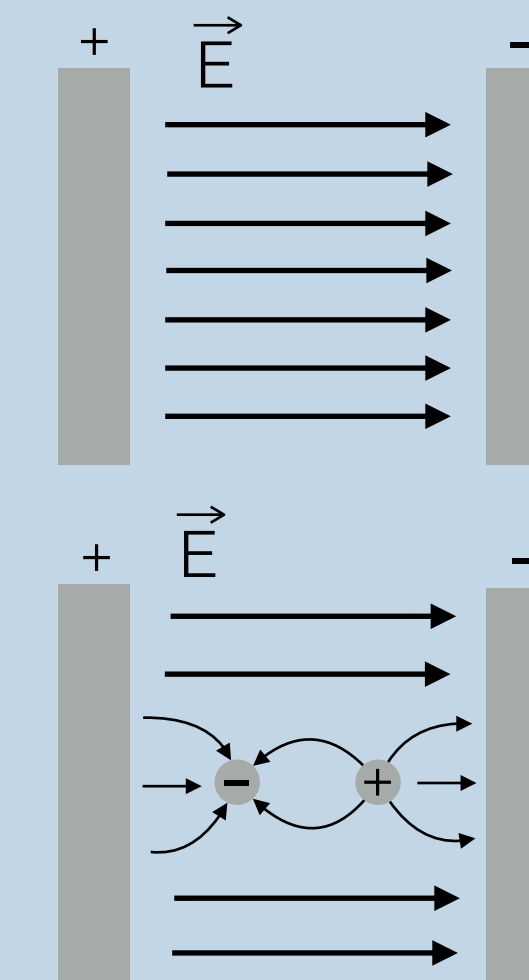


Schwinger → **Gaussian p_{\perp} spectrum** and heavy flavour suppression

Prob(u:d:s) \approx 1 : 1 : 0.2

Heavy quarks (charm and bottom) are only produced from hard processes → must be **string endpoints**

Schwinger mechanism QED



Non-perturbative creation of e^+e^- pairs in a string electric field

Probability from tunnelling factor

$$\mathcal{P} \propto \exp\left(\frac{-m^2 - p_{\perp}^2}{\kappa/\pi}\right)$$

$\kappa = \text{string tension}$

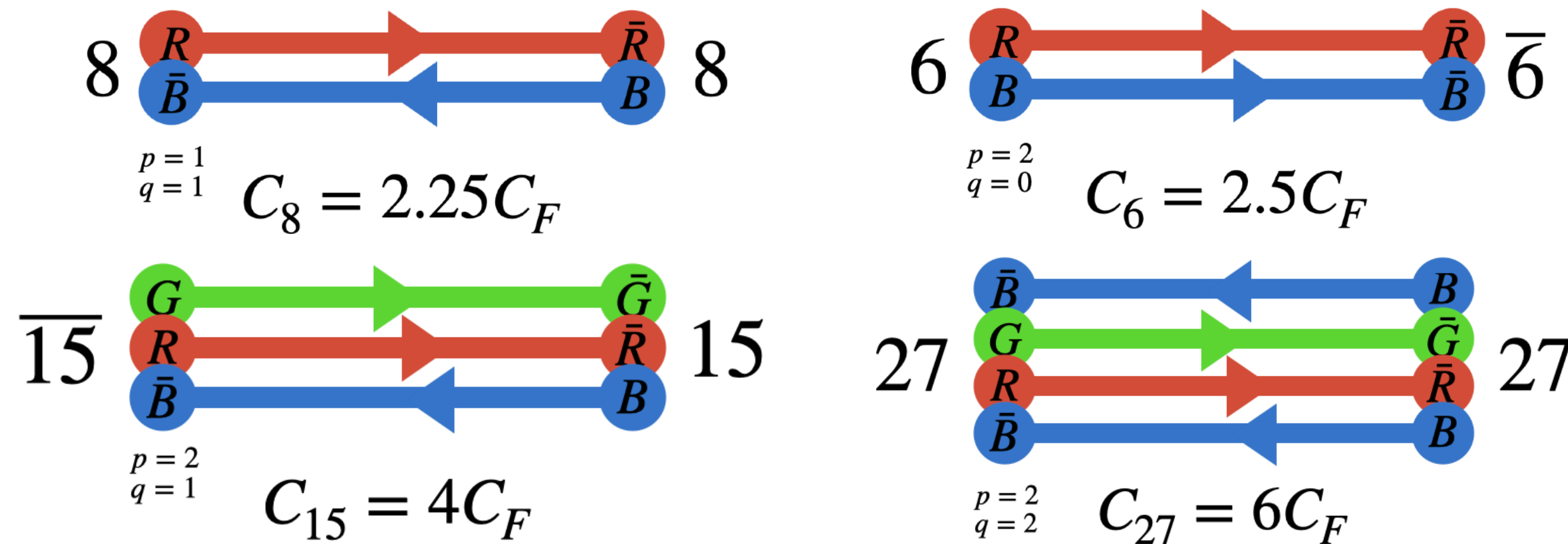
Strangeness Enhancement

Rope hadronisation

arXiv:1412.6259

After shoving, if strings are still overlapping \rightarrow form a rope

Enhance string tension for higher multiplets



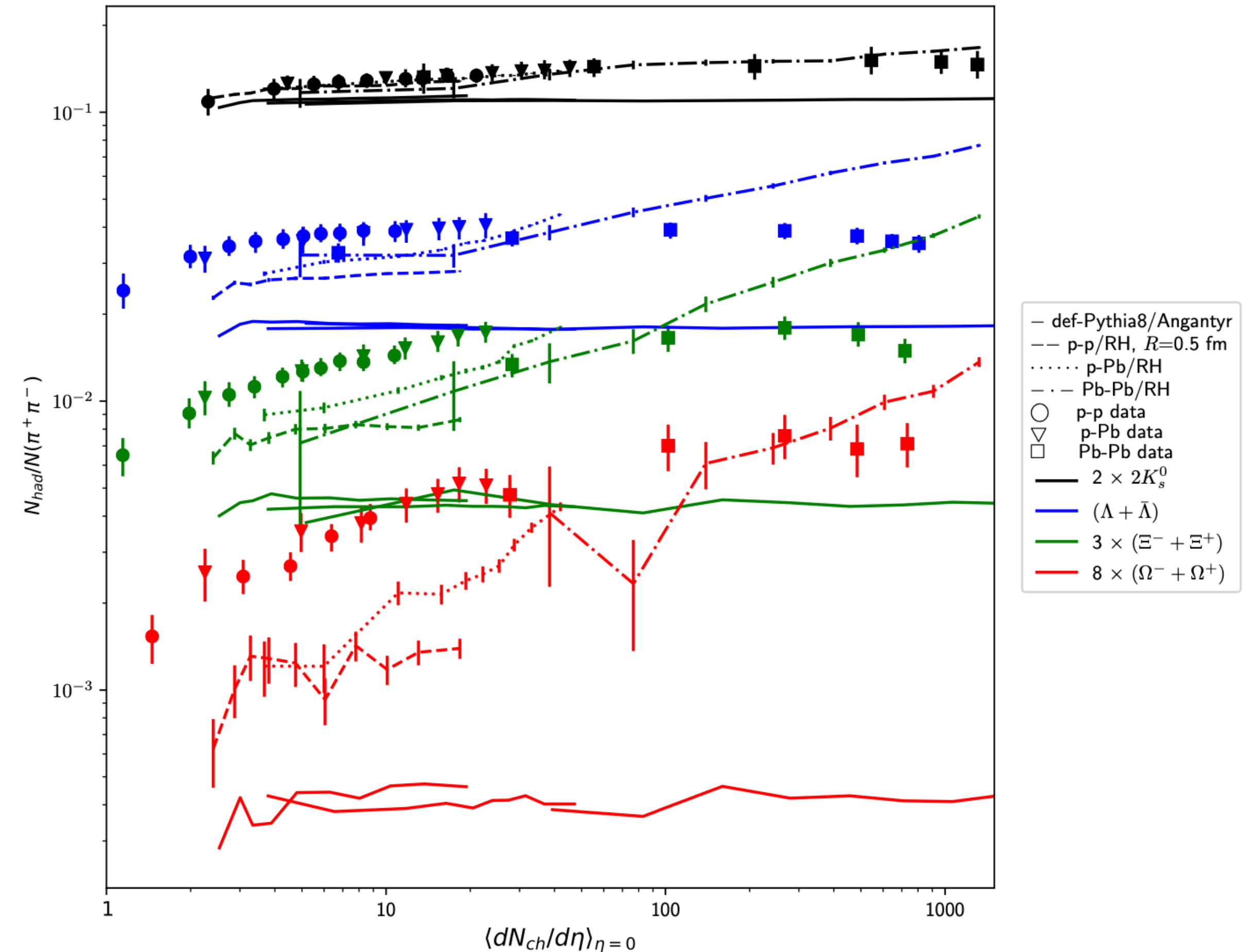
Dense string environments

\rightarrow Casimir scaling of **effective string tension**

\rightarrow Higher probability of strange quarks

Higher $\kappa \rightarrow$ lower strangeness suppression

$N/N(\pi^+\pi^-)$ vs. $\langle dN_{ch}/d\eta \rangle$ for p-p 7 TeV, p-Pb 5.02 TeV and Pb-Pb 2.76 TeV



Strangeness Enhancement

Rope hadronisation

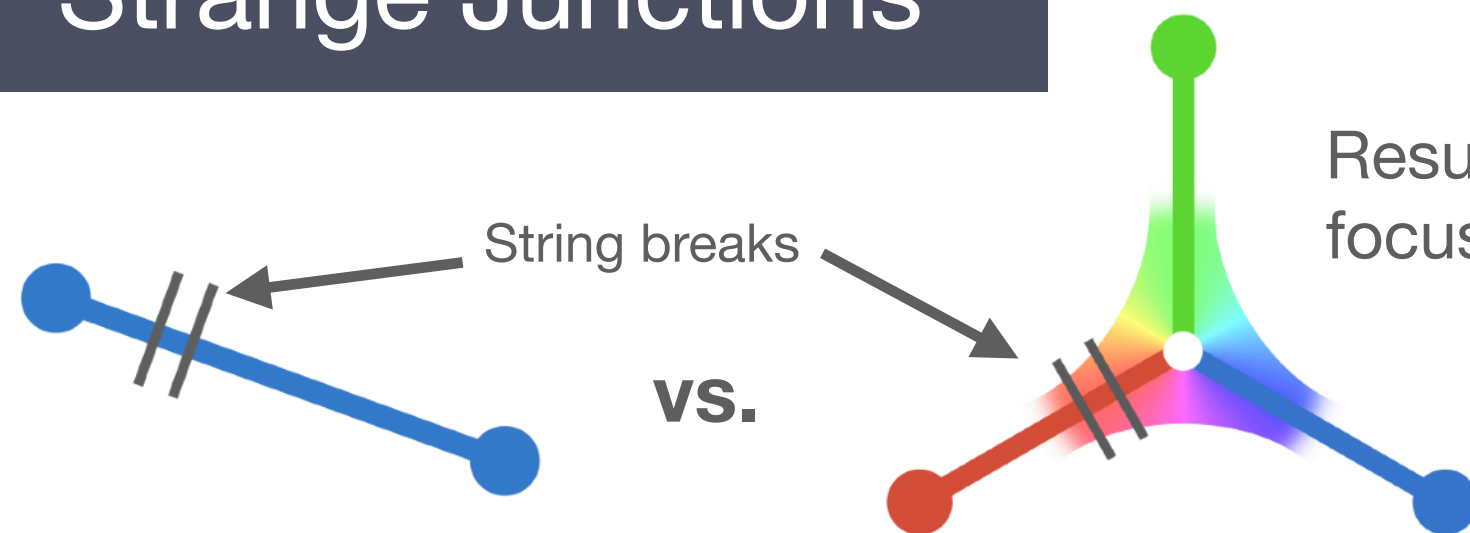
arXiv:1412.6259

After shoving, if strings are still overlapping \rightarrow form a rope
 Enhance string tension for higher multiplets

Close-packing

Simpler implementation of ropes fully in momentum space
 \rightarrow not in conjunction with shoving

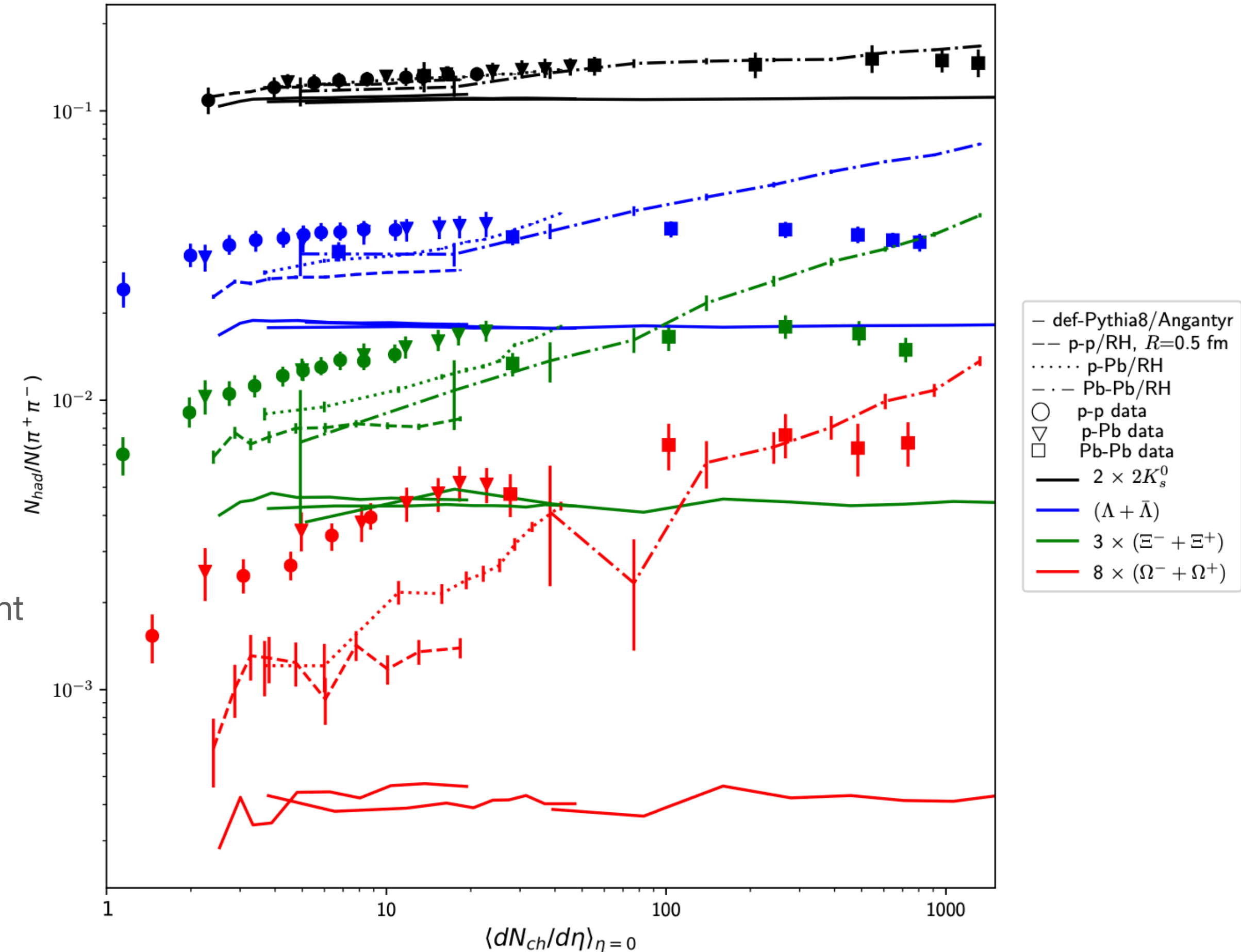
Strange Junctions



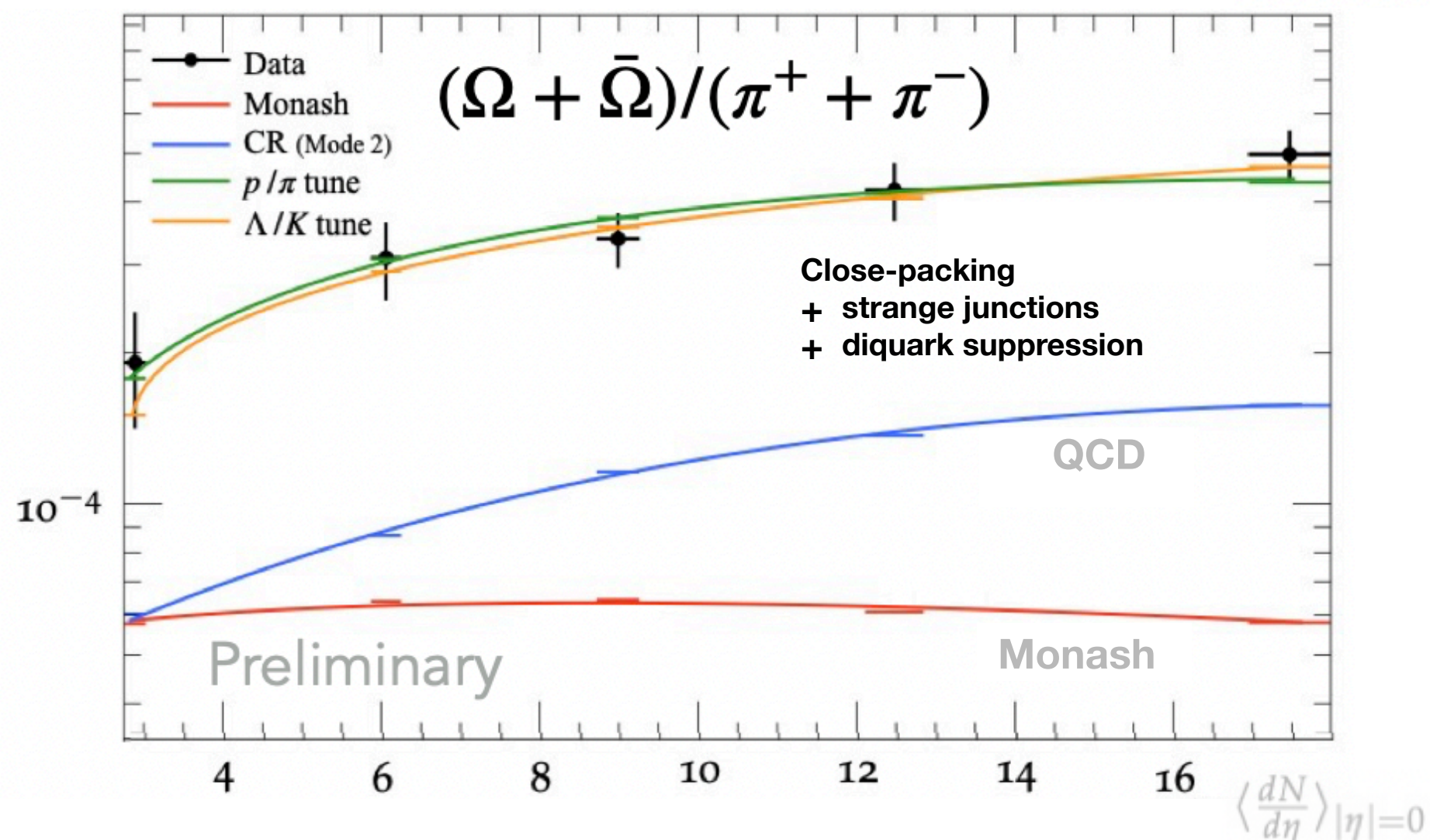
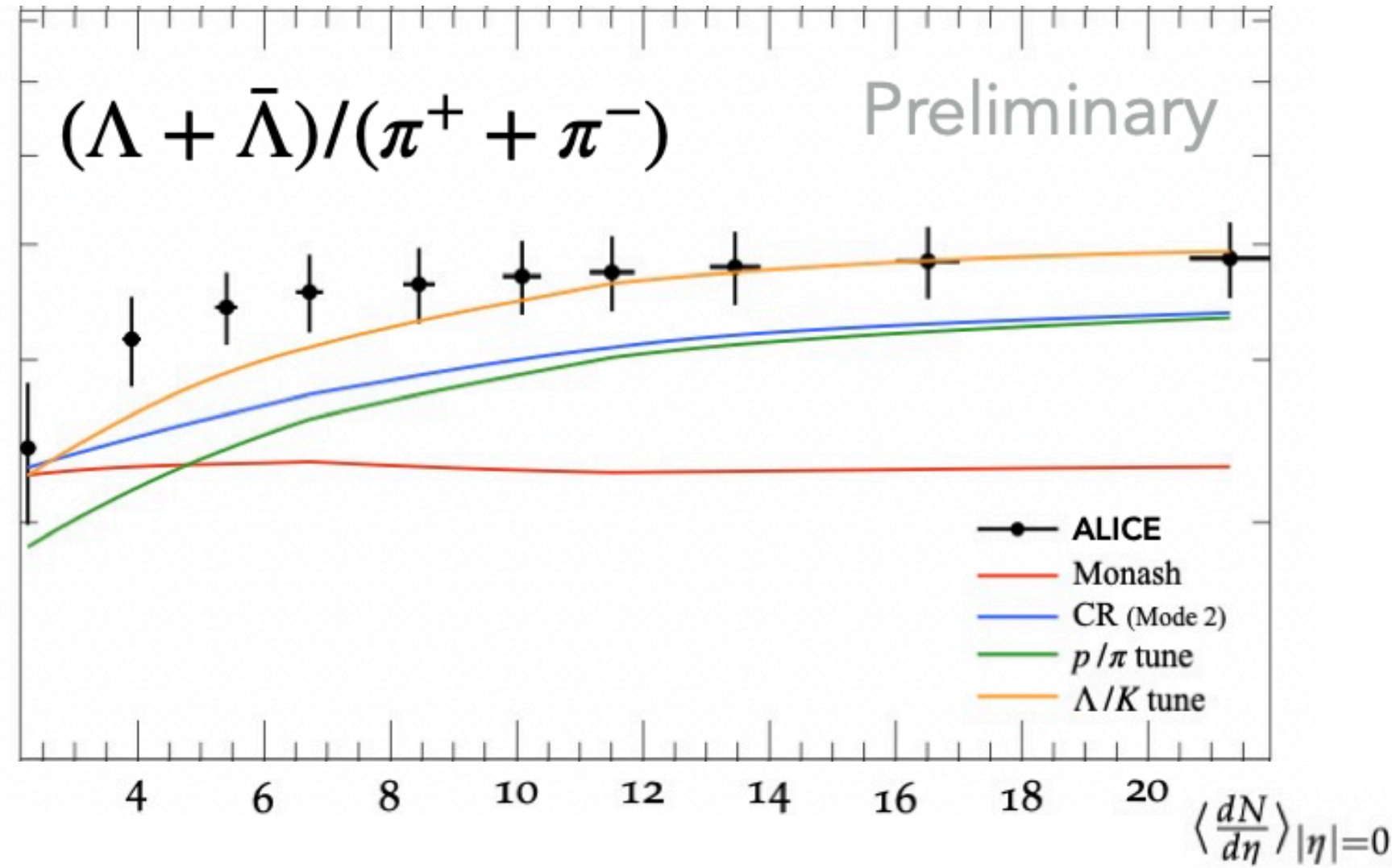
Results in strangeness enhancement
 focused in baryon sector

String tension could be different from the vacuum
 case compared to near a junction

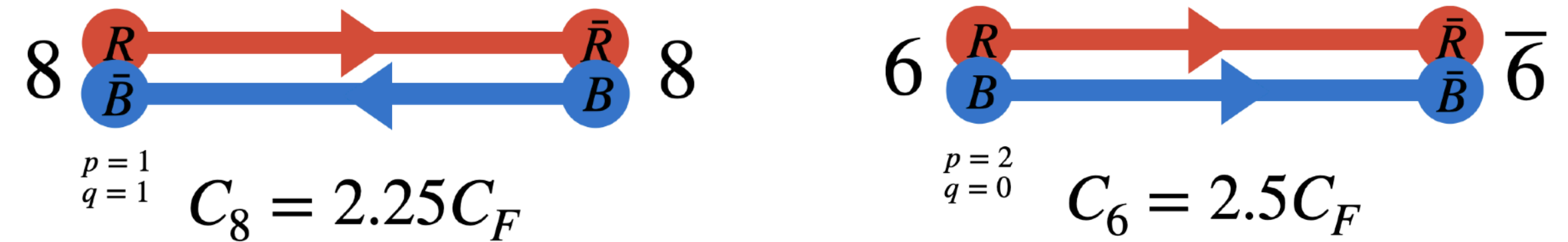
$N/N(\pi^+ \pi^-)$ vs. $\langle dN_{ch}/d\eta \rangle$ for p-p 7 TeV, p-Pb 5.02 TeV and Pb-Pb 2.76 TeV



Strangeness Enhancement



Close-packing + Ropes

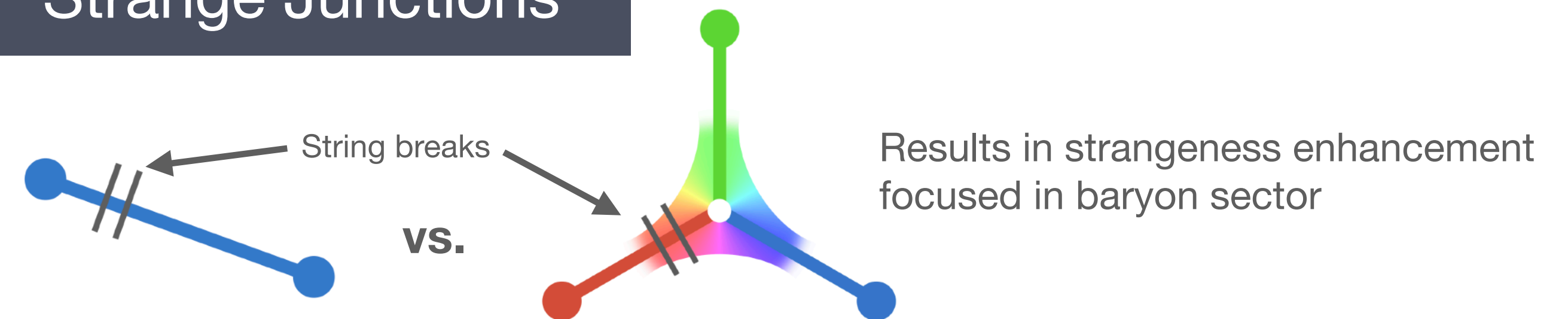


Dense string environments

→ Casimir scaling of **effective string tension**

→ Higher probability of strange quarks

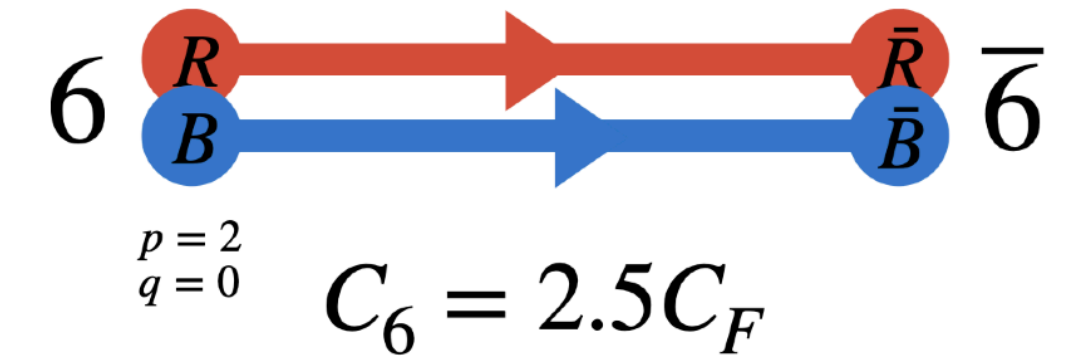
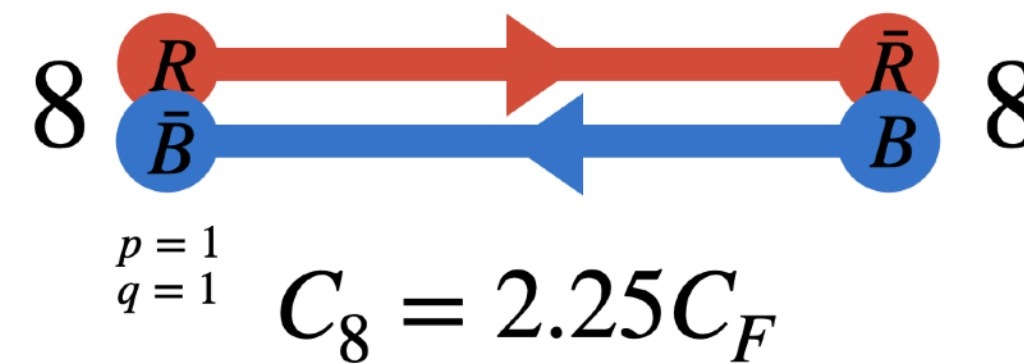
Strange Junctions



String tension could be different from the vacuum case compared to near a junction

Strangeness Enhancement

Close-packing + Ropes

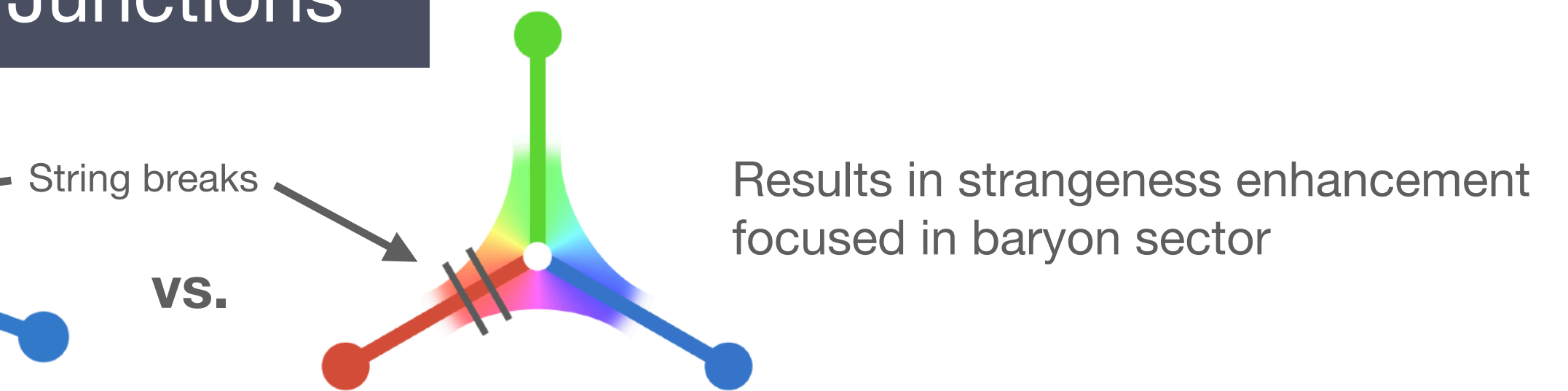


Dense string environments

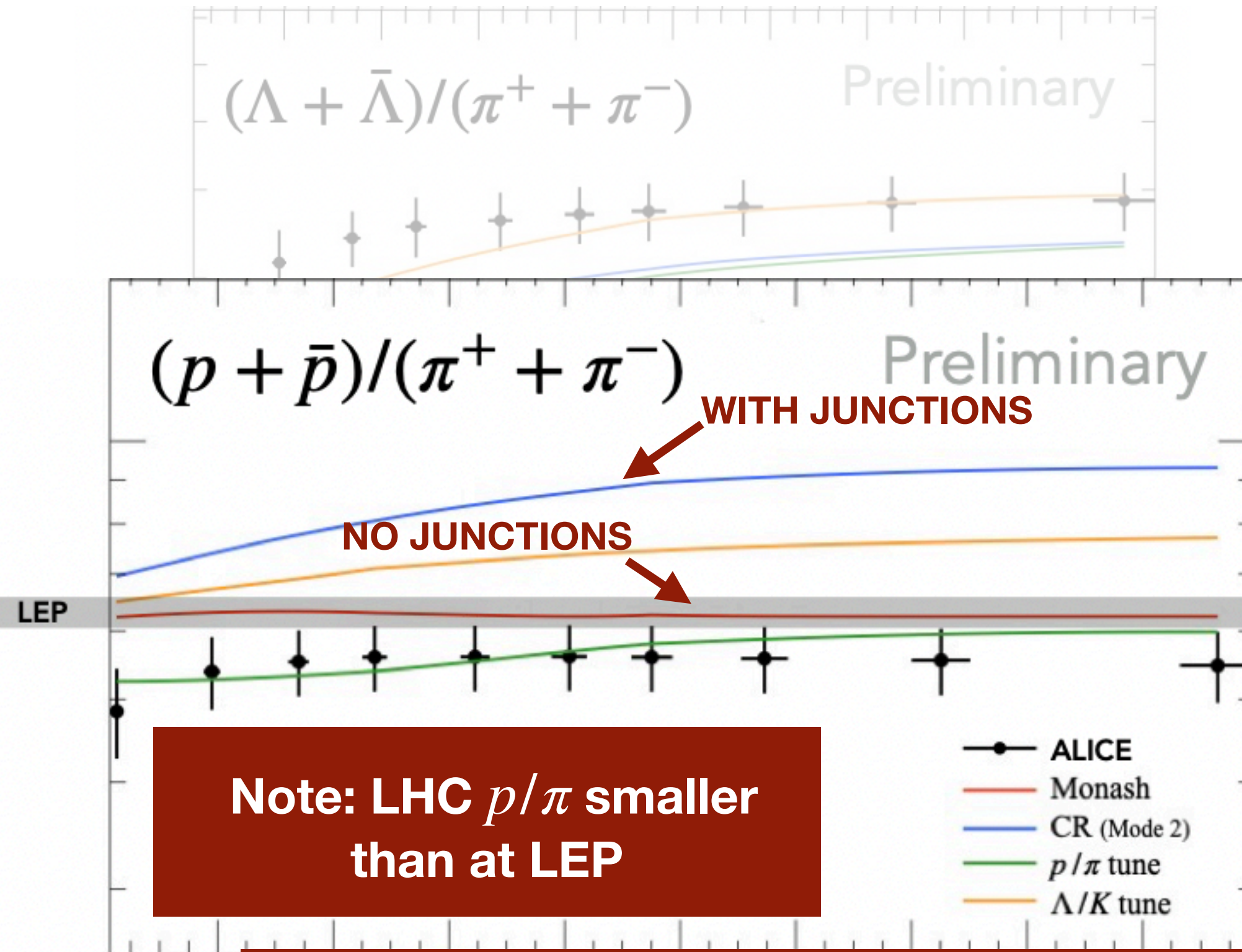
→ Casimir scaling of **effective string tension**

→ Higher probability of strange quarks

Strange Junctions

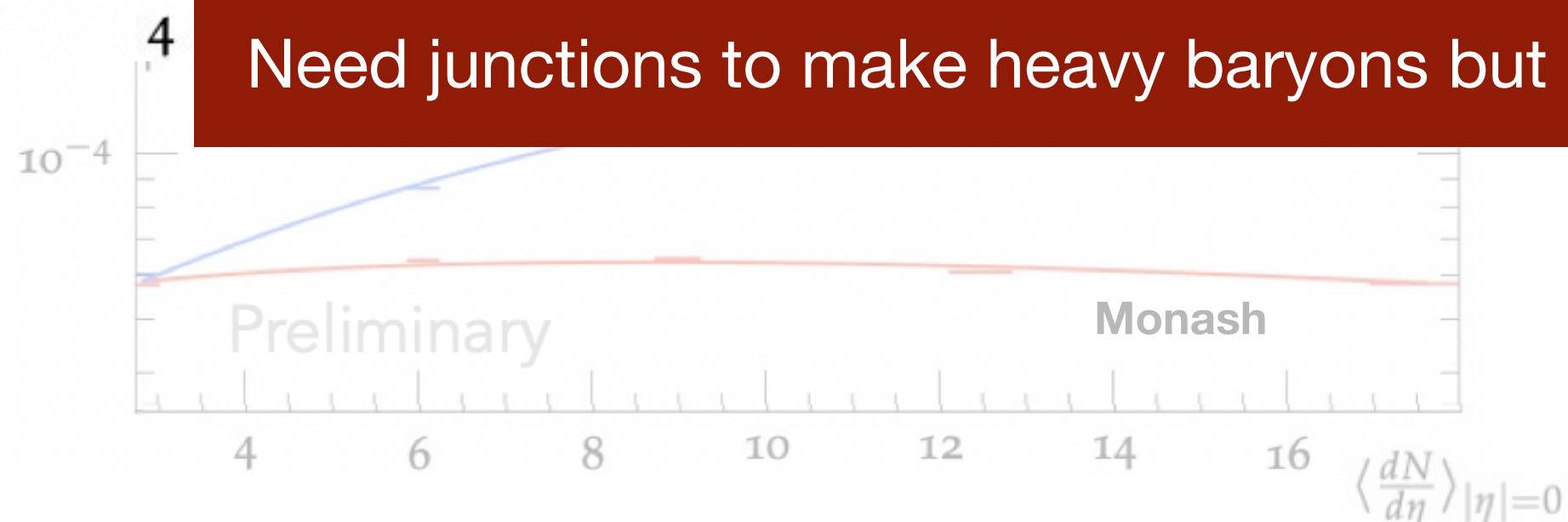


String tension could be different from the vacuum case compared to near a junction



Note: LHC p/π smaller than at LEP

Need junctions to make heavy baryons but need less protons?



Strangeness Enhancement

Rope hadronisation

arXiv:1412.6259

After shoving, if strings are still overlapping \rightarrow form a rope

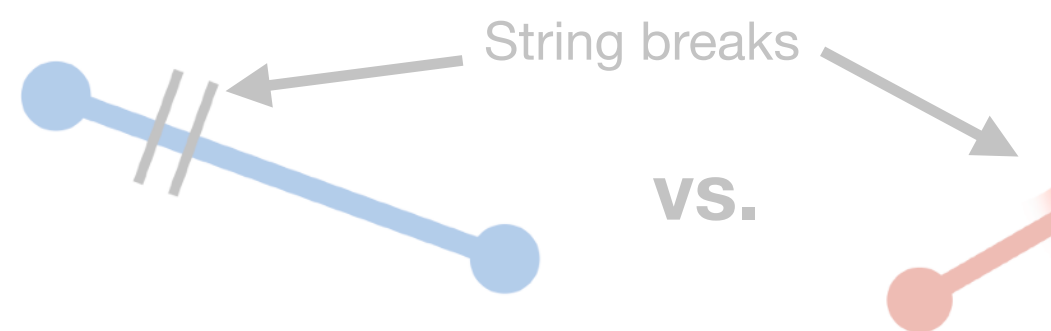
Enhance string tension for

What about the non-strange baryon-to-meson ratio p/π ?

Close-packing

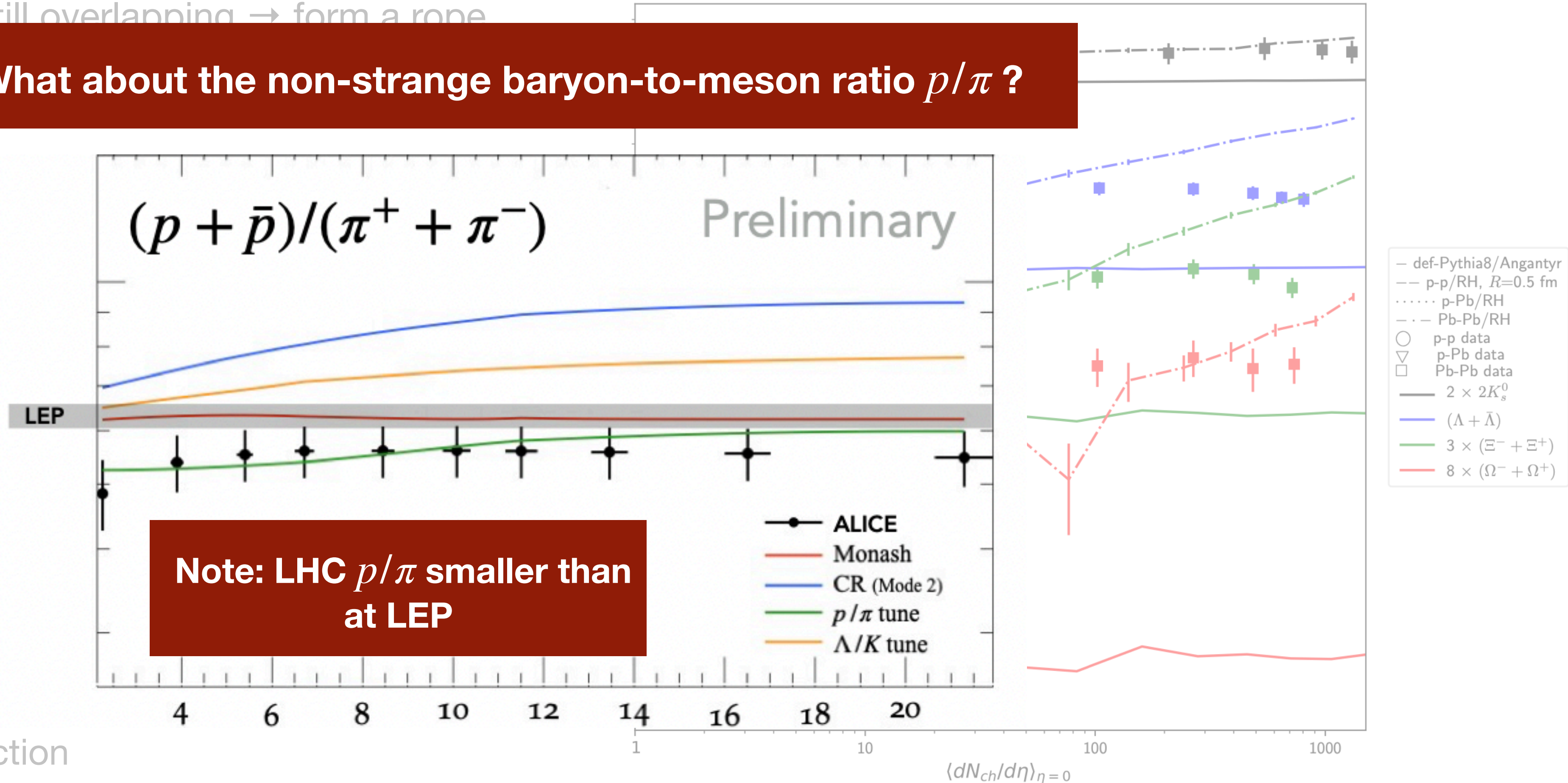
Simpler implementation of rope
 \rightarrow not in conjunction with shoving

Strange Junctions



String tension could be different case compared to near a junction

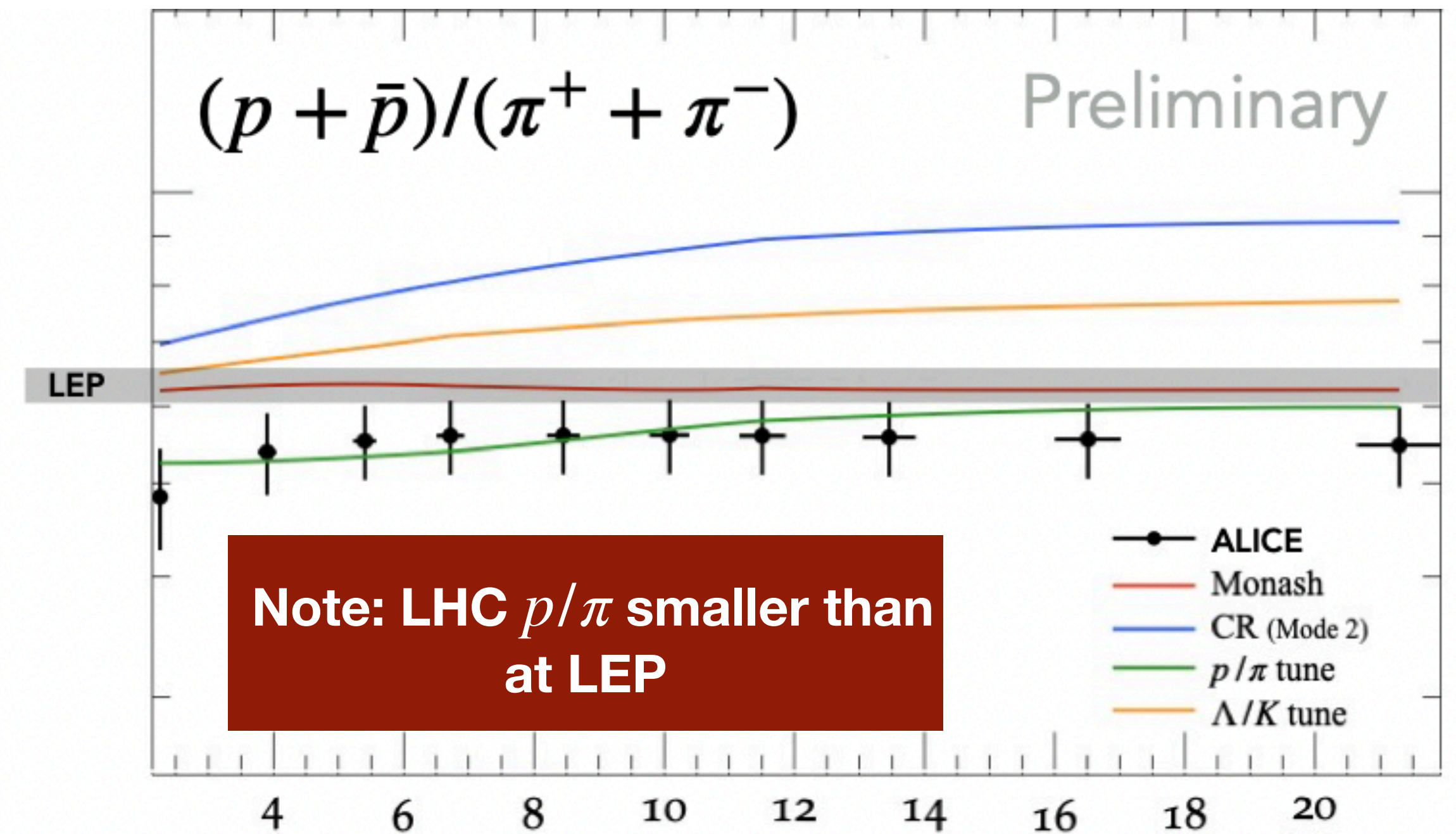
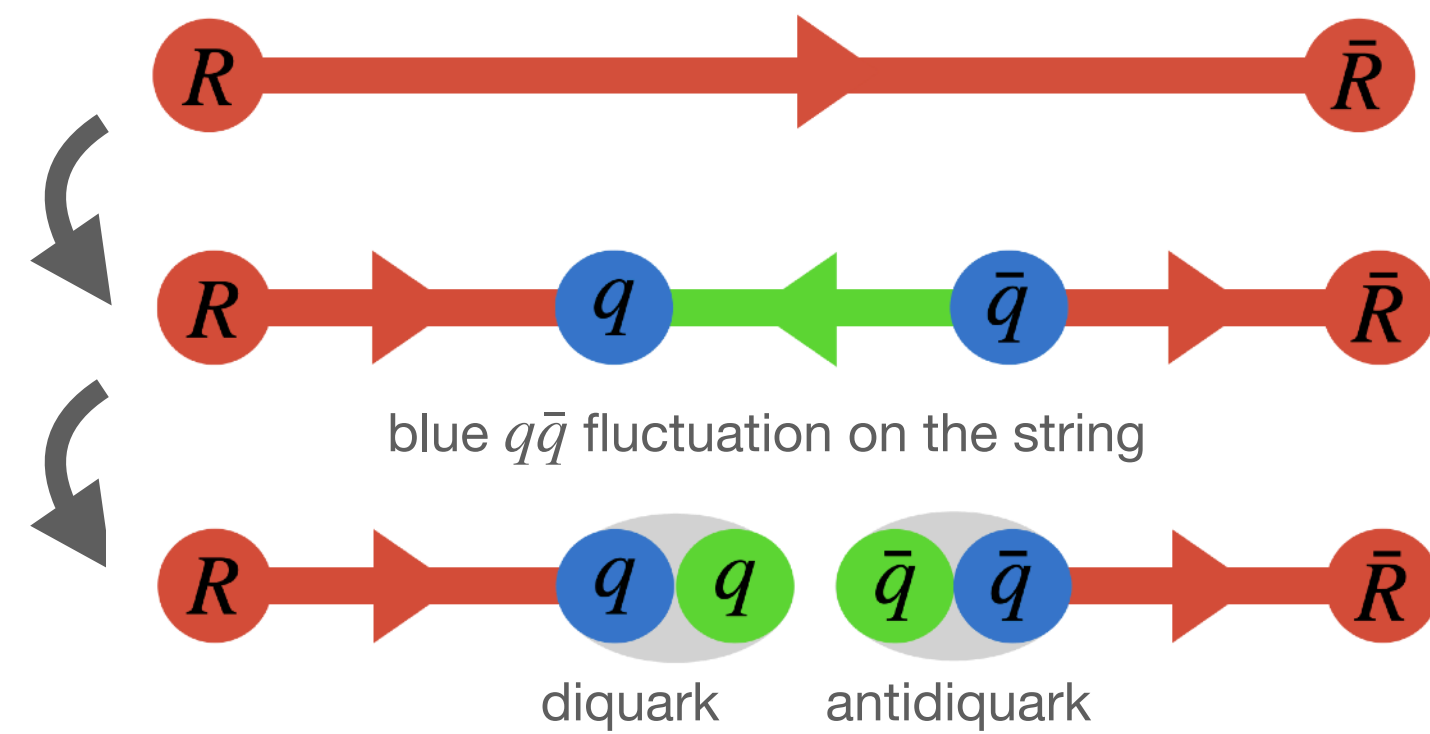
$N/N(\pi^+ \pi^-)$ vs. $\langle dN_{ch}/d\eta \rangle$ for p-p 7 TeV, p-Pb 5.02 TeV and Pb-Pb 2.76 TeV



Proton problem

Popcorn mechanism for diquark production

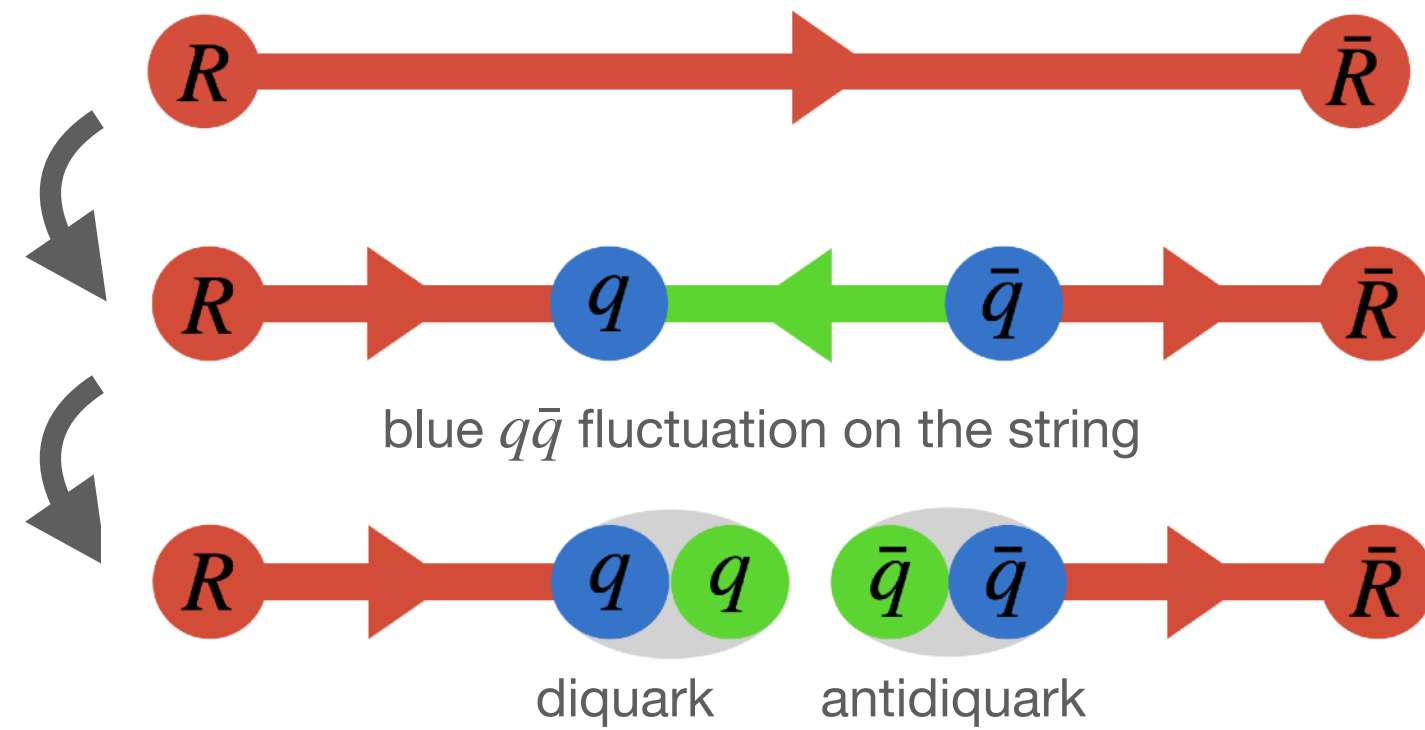
Diquark formation via **successive colour fluctuations** — popcorn mechanism



Proton problem

Popcorn mechanism for diquark production

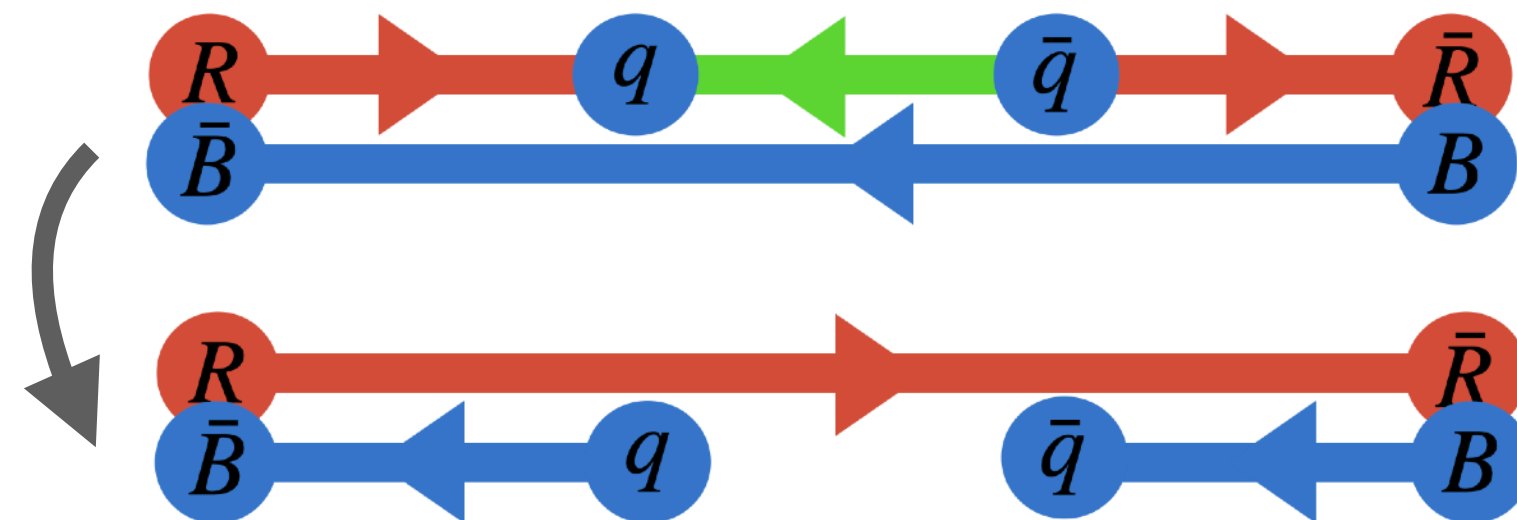
Diquark formation via **successive colour fluctuations** — popcorn mechanism



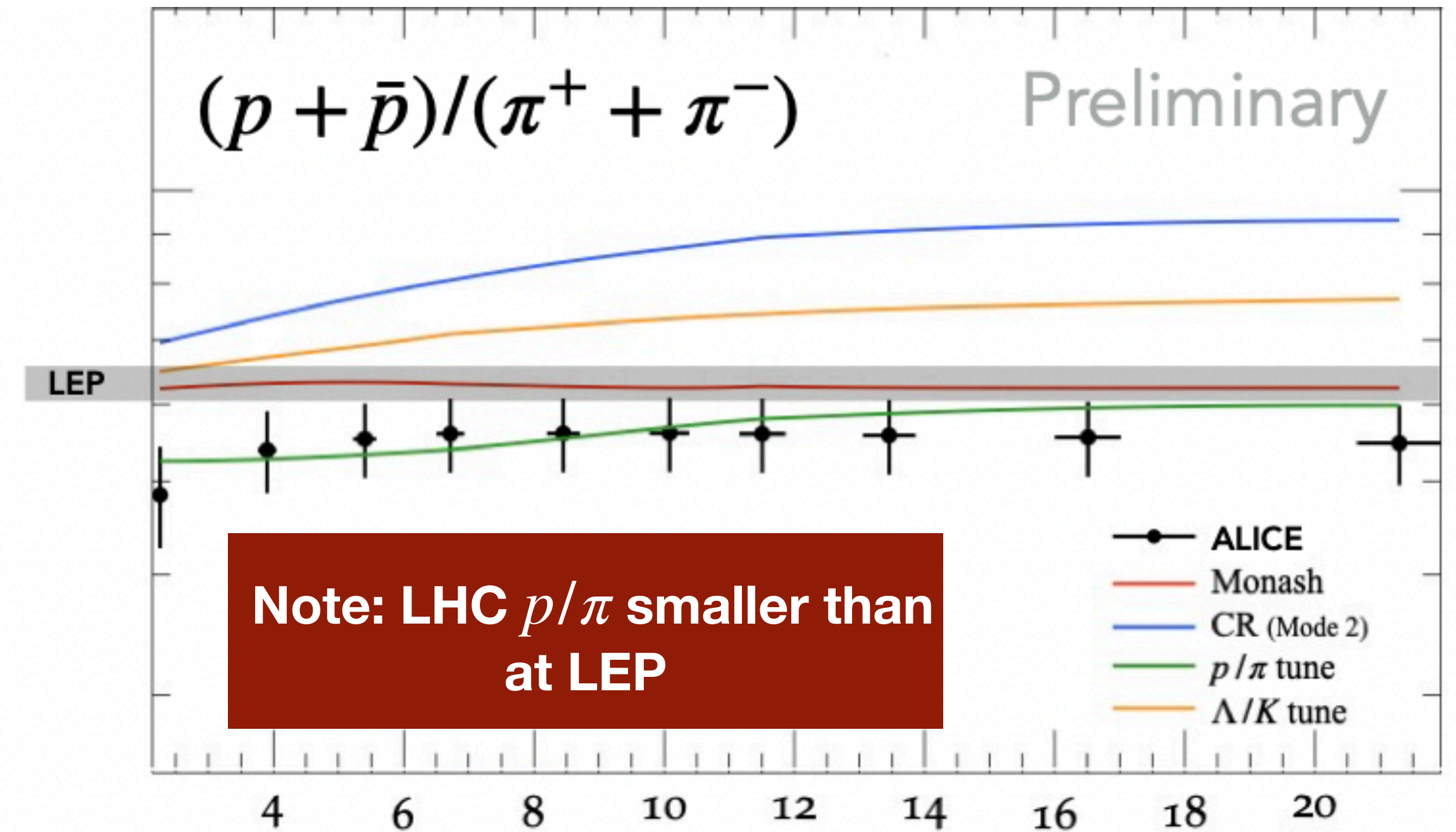
Popcorn destructive interference

NEW

What if there's a blue string nearby?

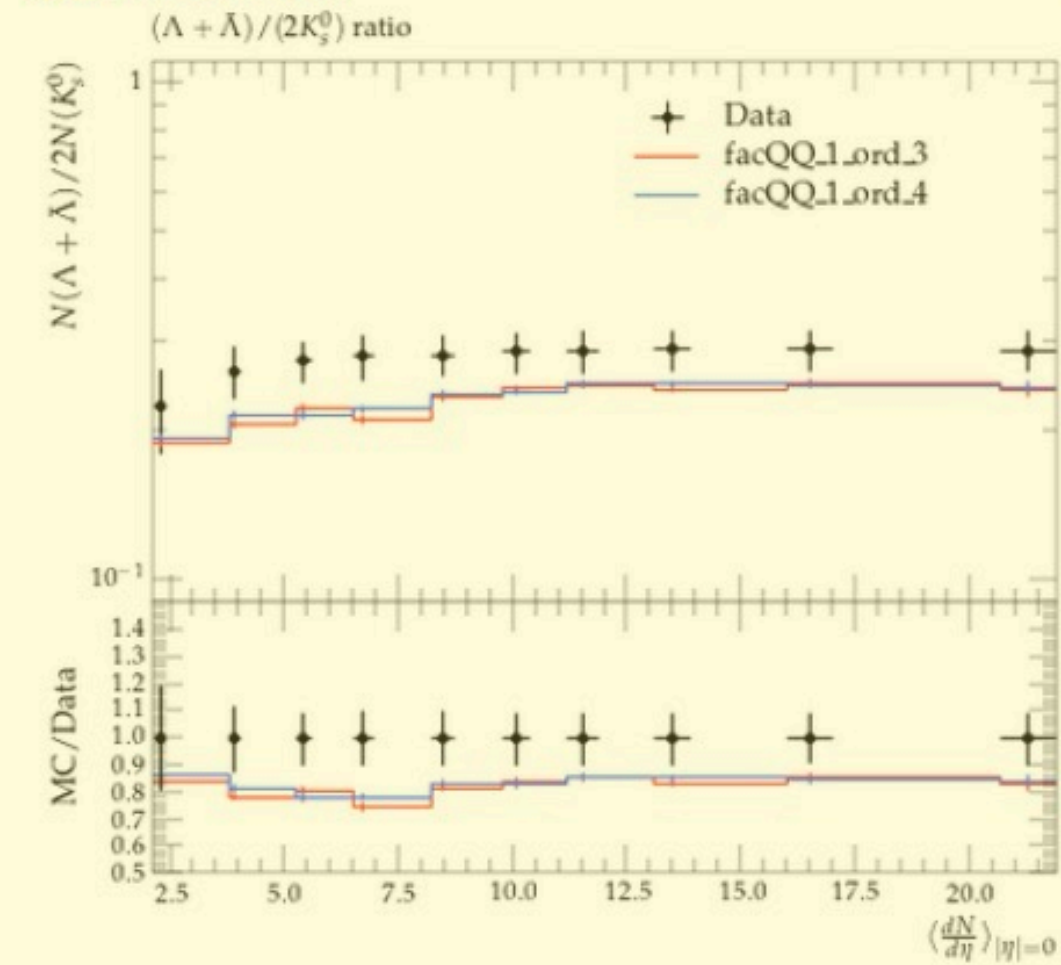


blue $q\bar{q}$ fluctuation breaks nearby blue string, preventing diquark formation

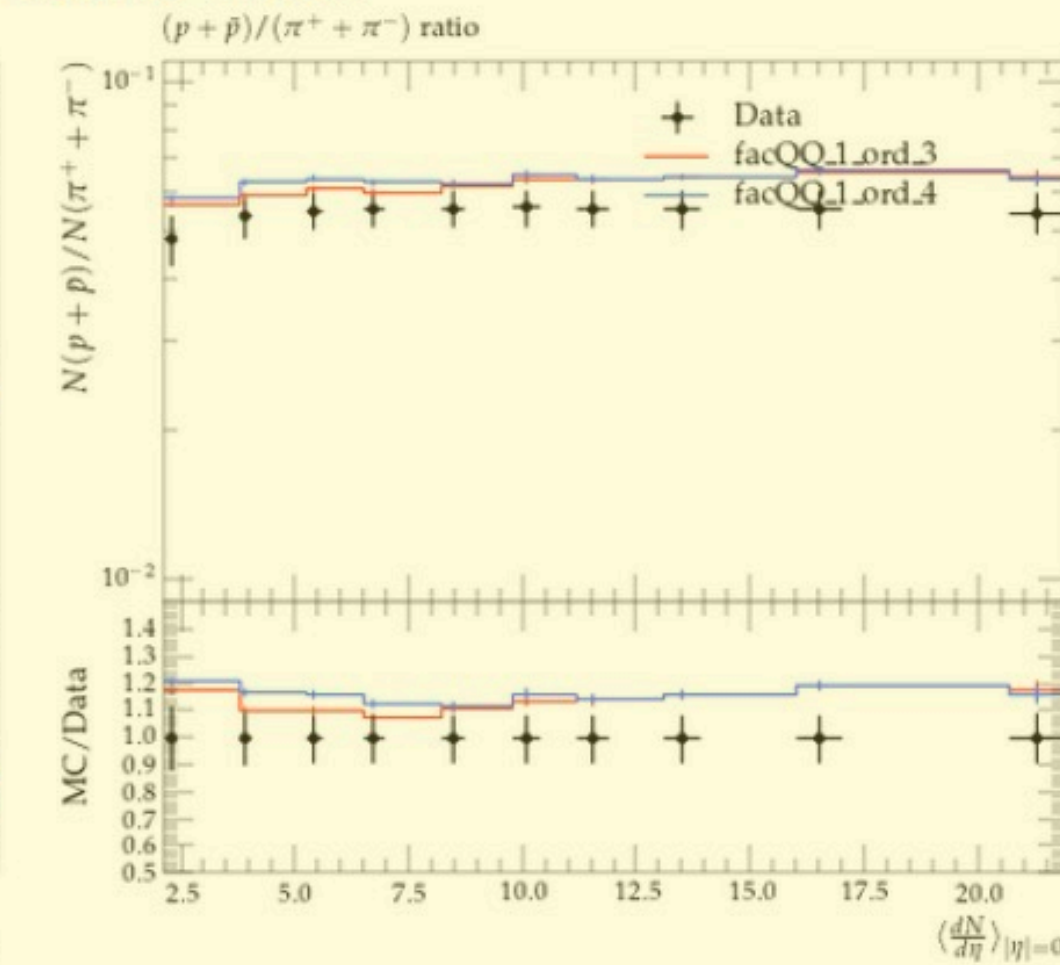


Results — ongoing

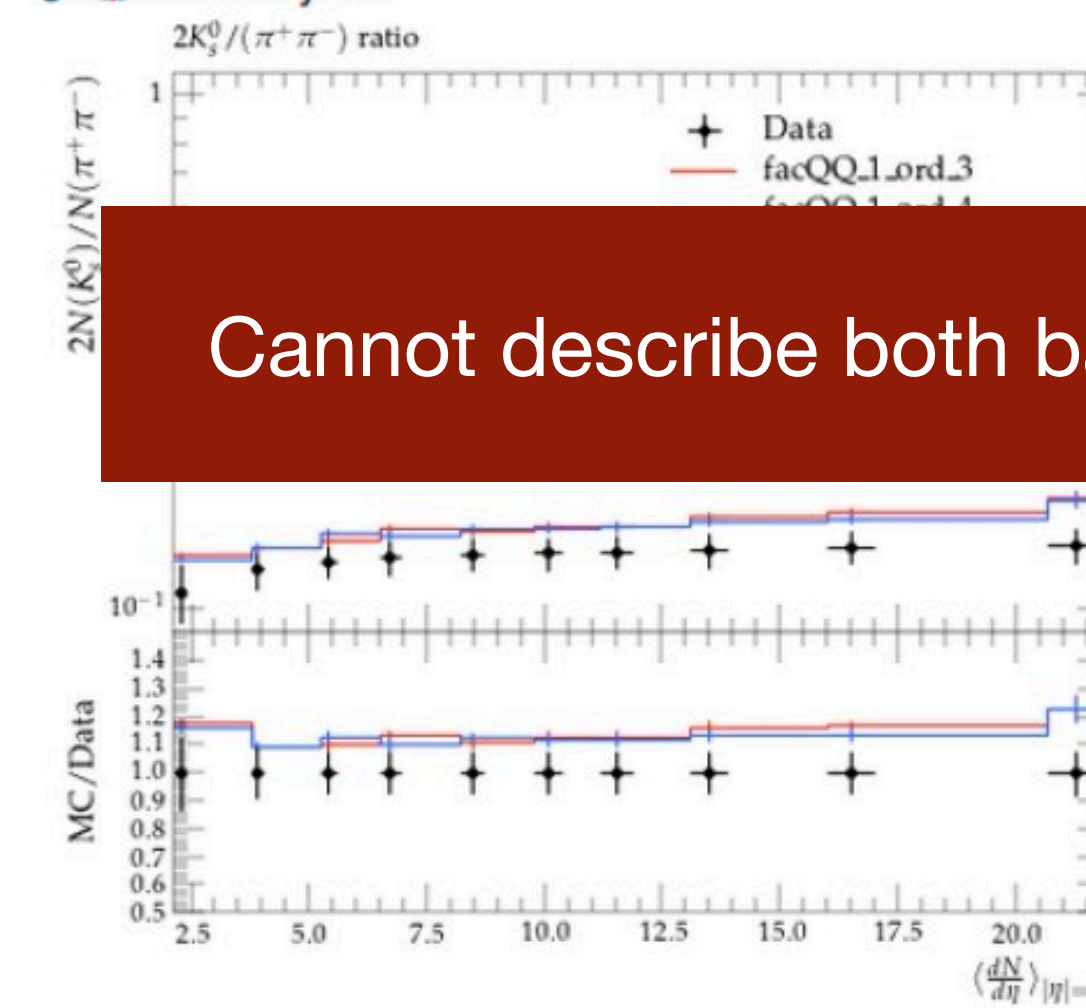
d46-x01-y01:



d47-x01-y01:



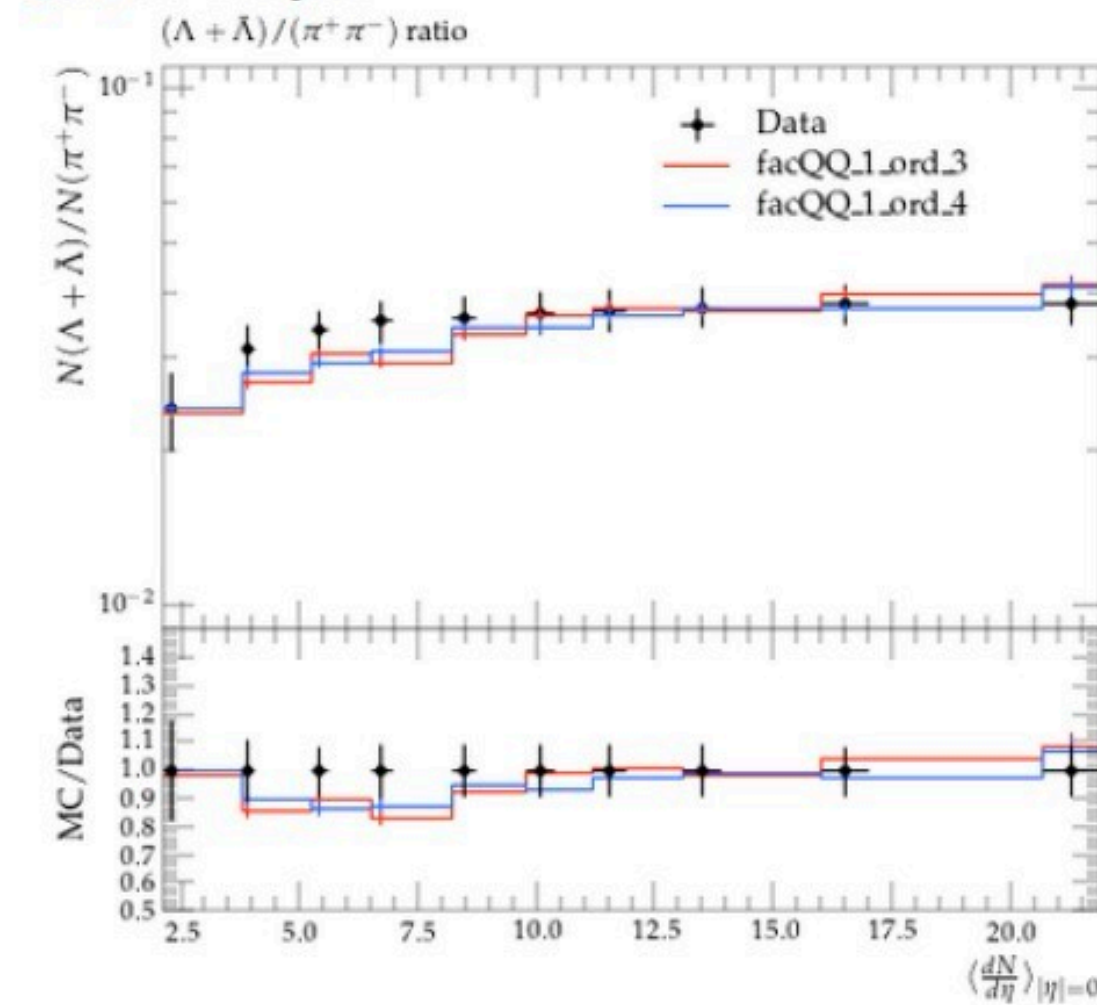
d36-x01-y01:



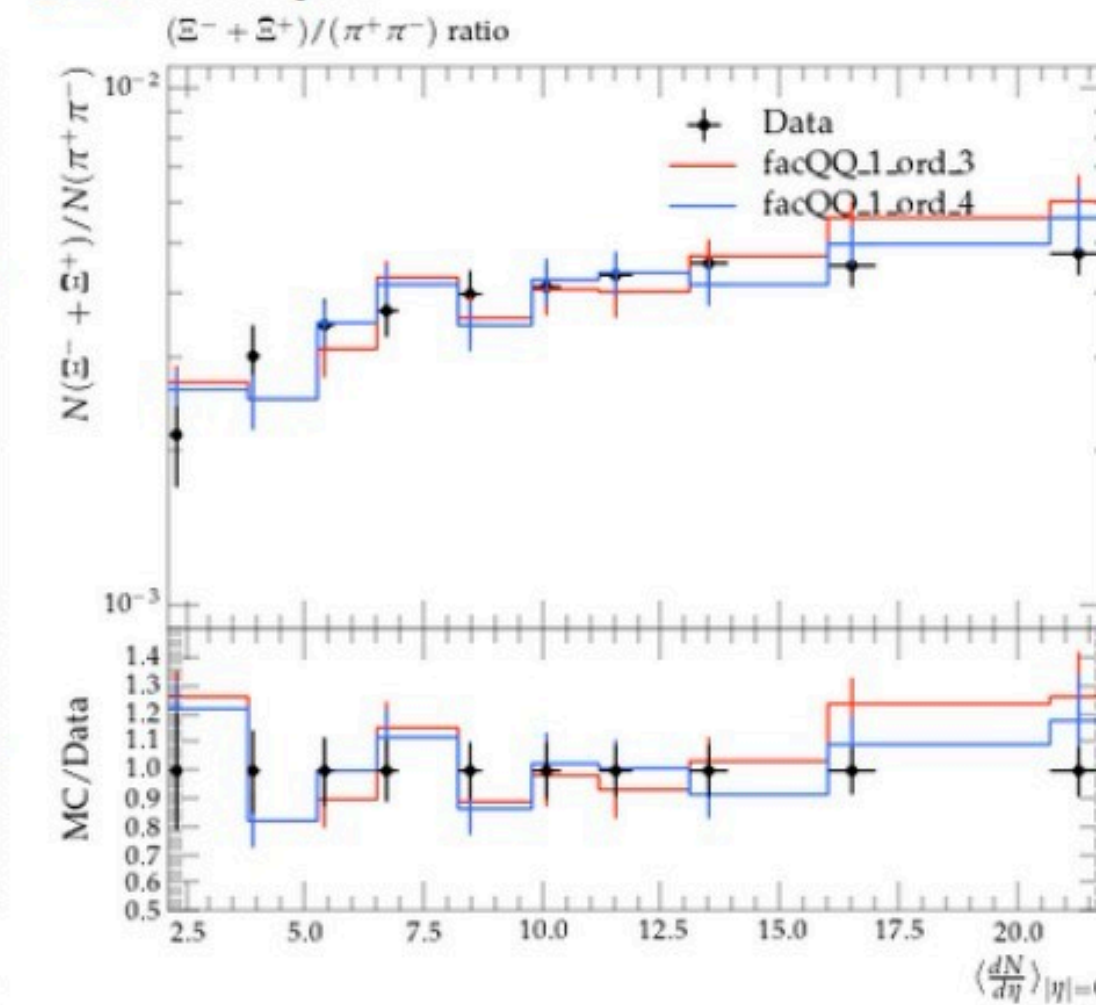
Cannot describe both baryon-to-meson ratios simultaneously

Taken from slide by Lorenzo Bernadinis: masters student currently undertaking tuning project with the model

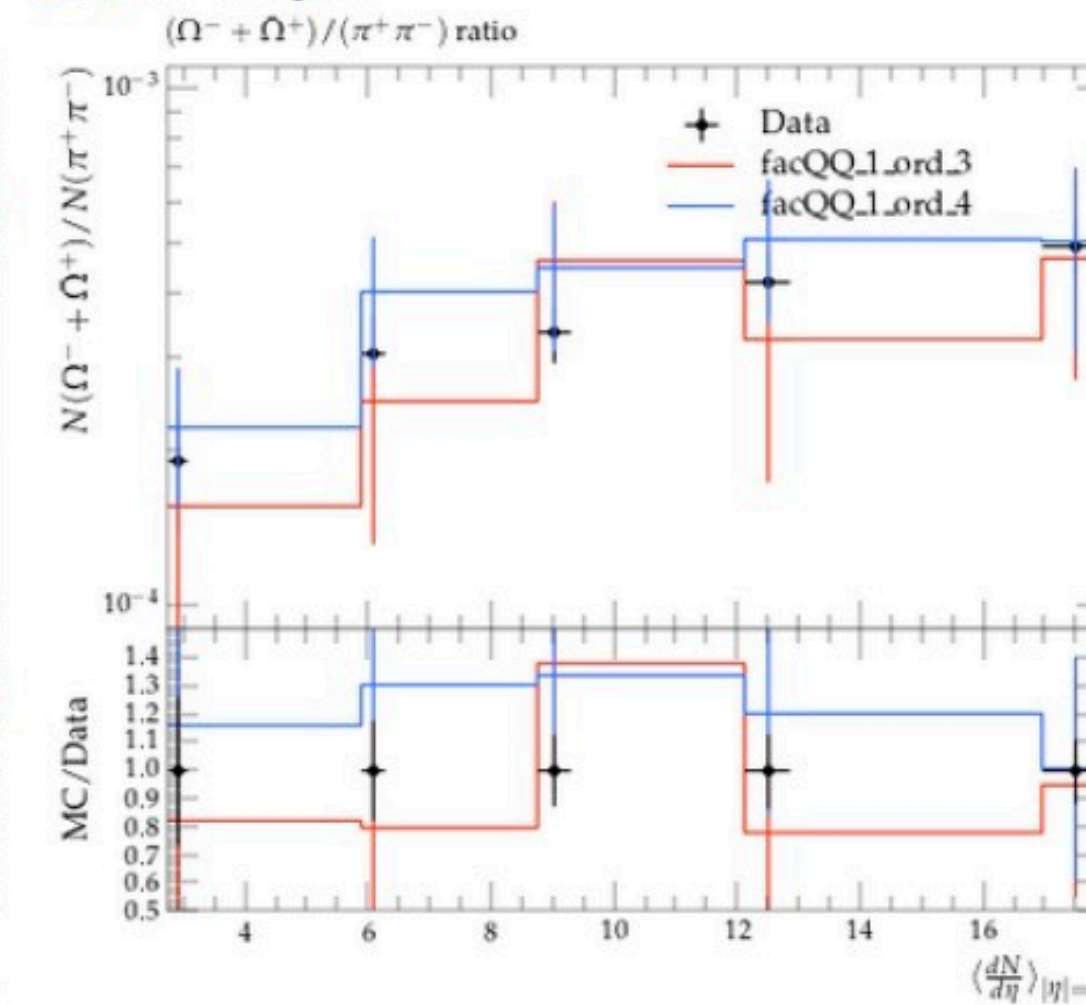
d37-x01-y01:



d38-x01-y01:



d39-x01-y01:



Summary

Evidence that **collective effects can arise from non-QGP** sources

CR restores $SU(3)$ colour correlations

→ baryons-to-meson ratio enhancement, $\langle p_{\perp} \rangle$ increase with multiplicity, some flow-like

Angantyr allows for pA and AA using strings instead of QGP

→ multiplicity distributions for AA

Shoving string interactions before hadronisation

→ near-sided ridge in pp , some v_2 with full description hindered by implementation technicality issues

Ropes

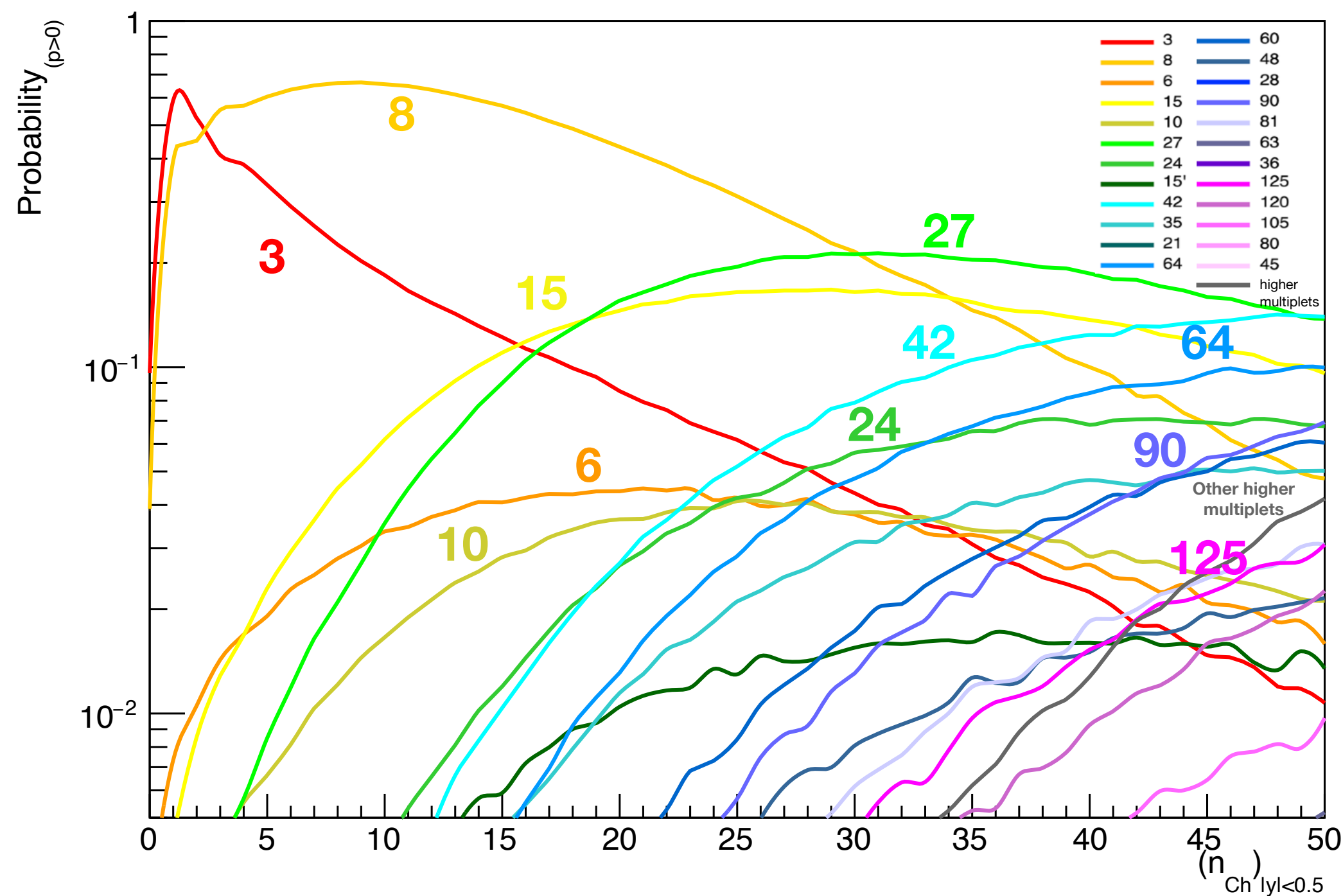
→ strangeness enhancement

Unmentioned: jet quenching, hadron rescattering

Future studies: shoving **considering regions formed by soft gluons**, reexamination of results given **updates to CR in Angantyr** (previous modelling only included CR within each nucleon-nucleon collision, now CR is allowed between nucleon-nucleon collisions)

Vacuum → High multiplicities

Multiplets ($y=0$, pp 7 TeV)



Clear observations of strangeness enhancement with respect to charged multiplicity [e.g. ALICE Nature Phys. 13, 535 (2017)]

Protons are composite

- lots of quarks and gluons inside
- multiple parton-parton interactions
- lots of colour charges

Strangeness enhancement with charged multiplicity suggests **higher multiplicity** string systems act **different to the vacuum case**

Number of fundamental and anti-fundamental flux lines at central rapidity in pp collisions give us **effective multiplet representation**

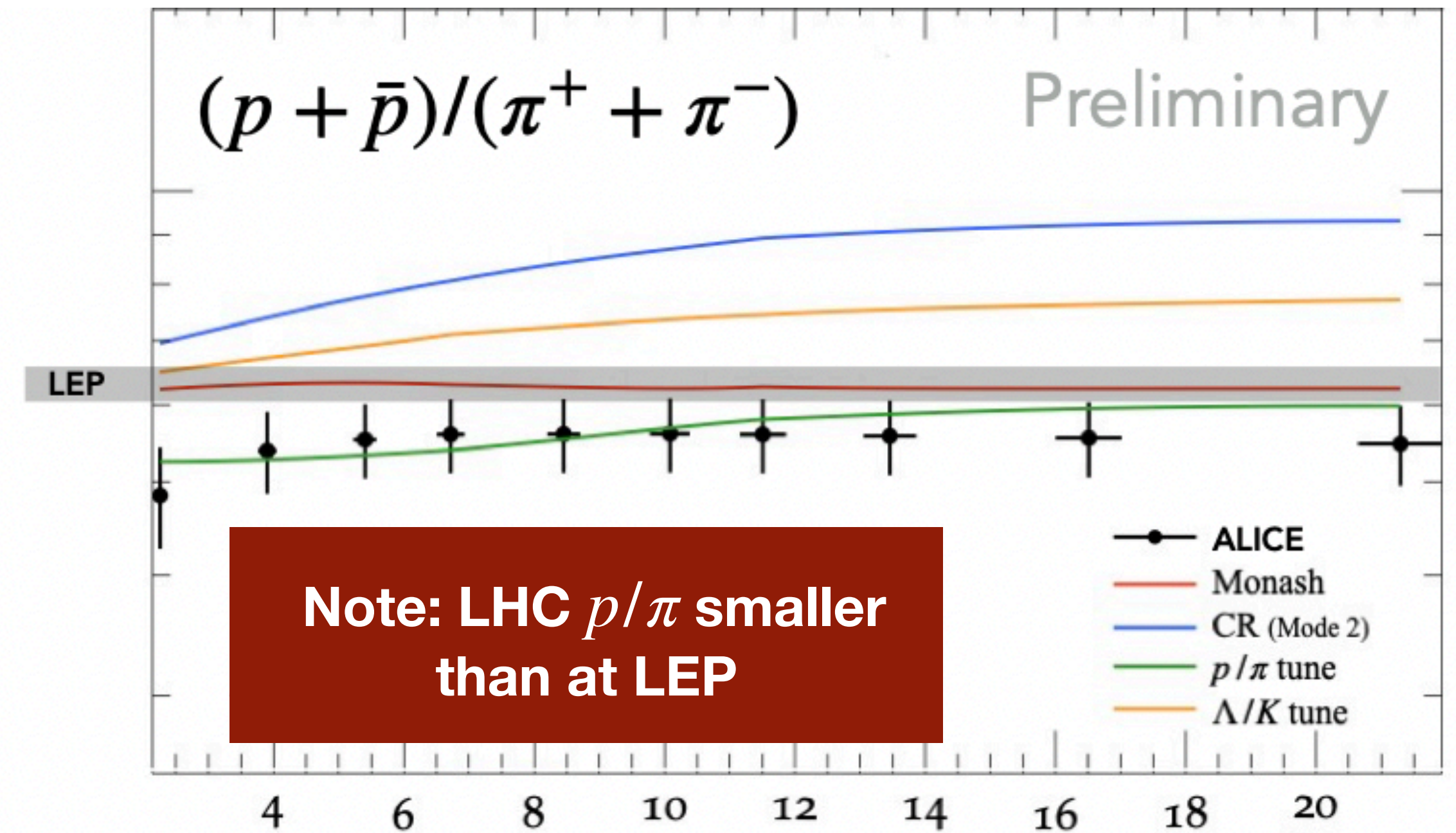
Reach higher than simple quark-antiquark triplet string

Popcorn mechanism

Popcorn Mechanism

arXiv:hep-ph/9606454

Diquark formation via **successive colour fluctuations**

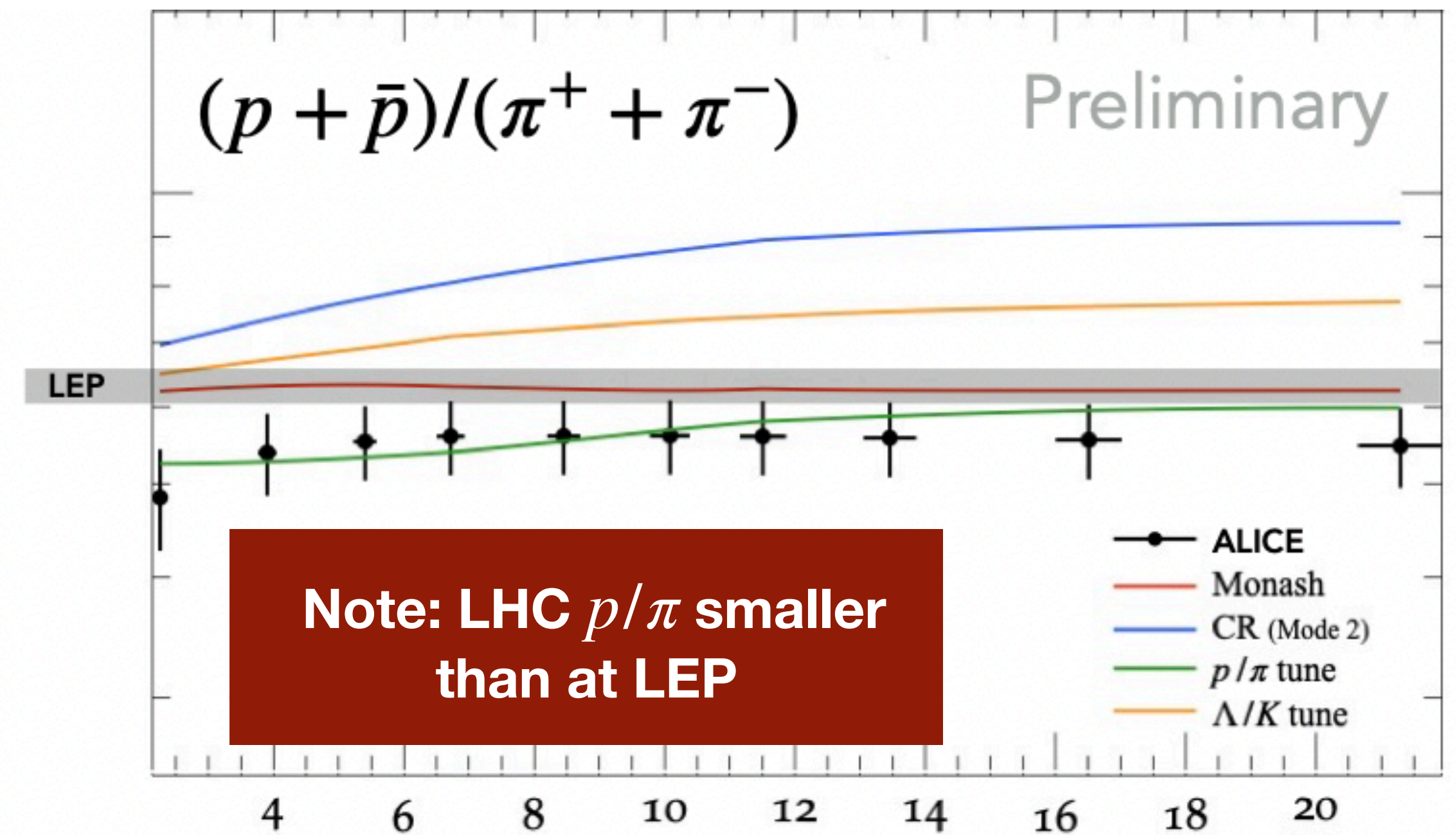
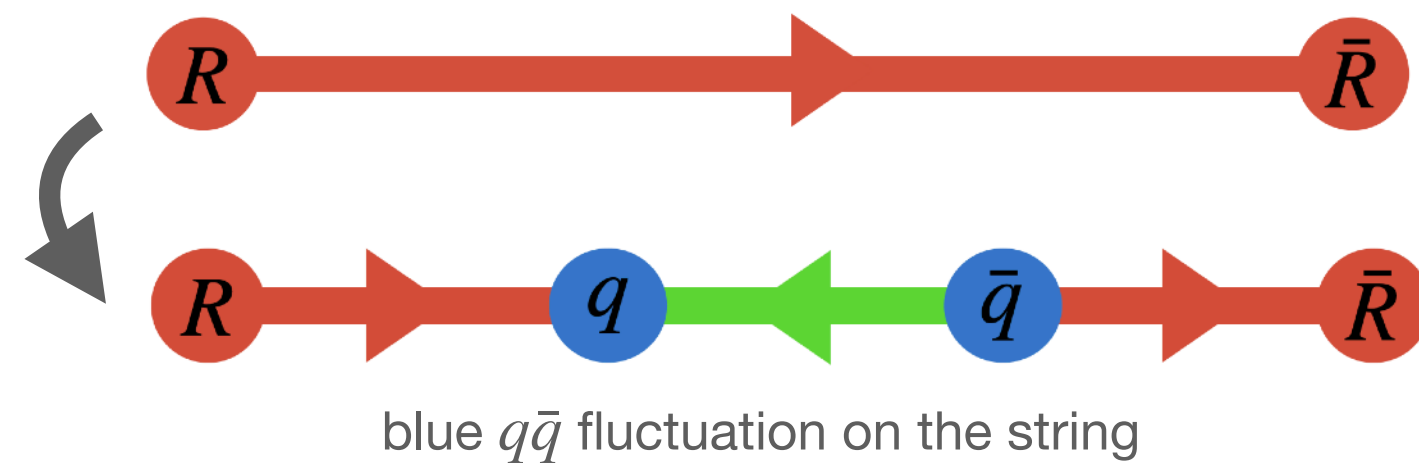


Popcorn mechanism

Popcorn Mechanism

arXiv:hep-ph/9606454

Diquark formation via **successive colour fluctuations**

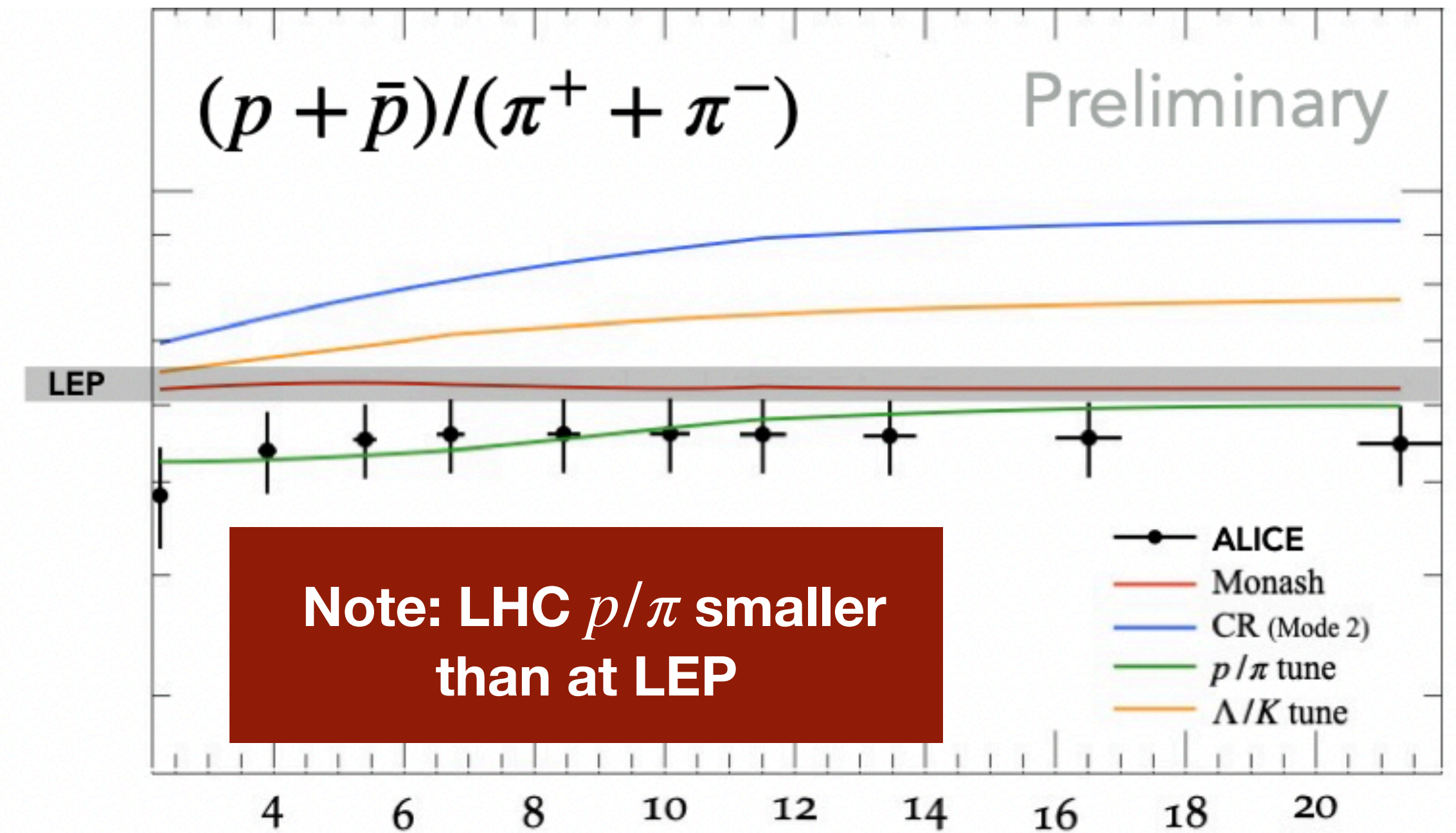
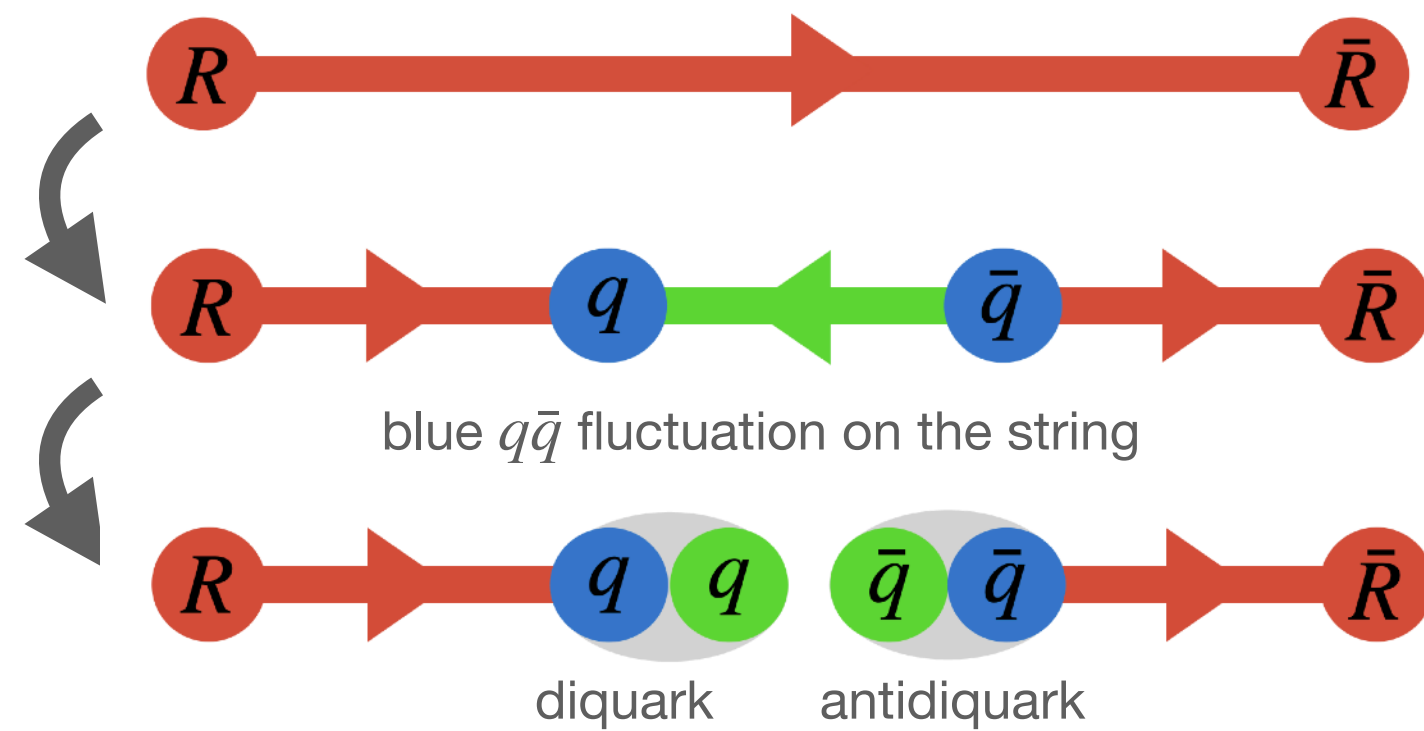


Popcorn mechanism

Popcorn Mechanism

arXiv:hep-ph/9606454

Diquark formation via **successive colour fluctuations**

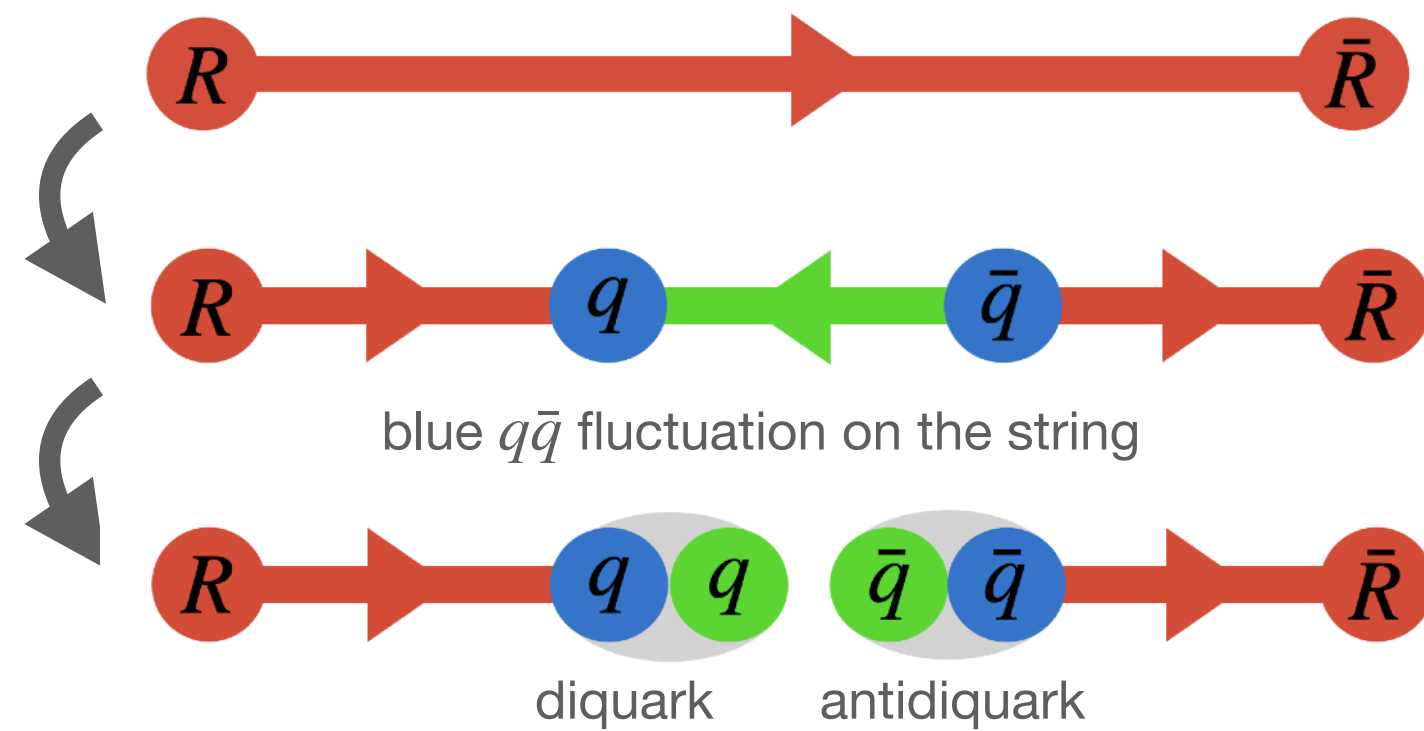


Popcorn mechanism

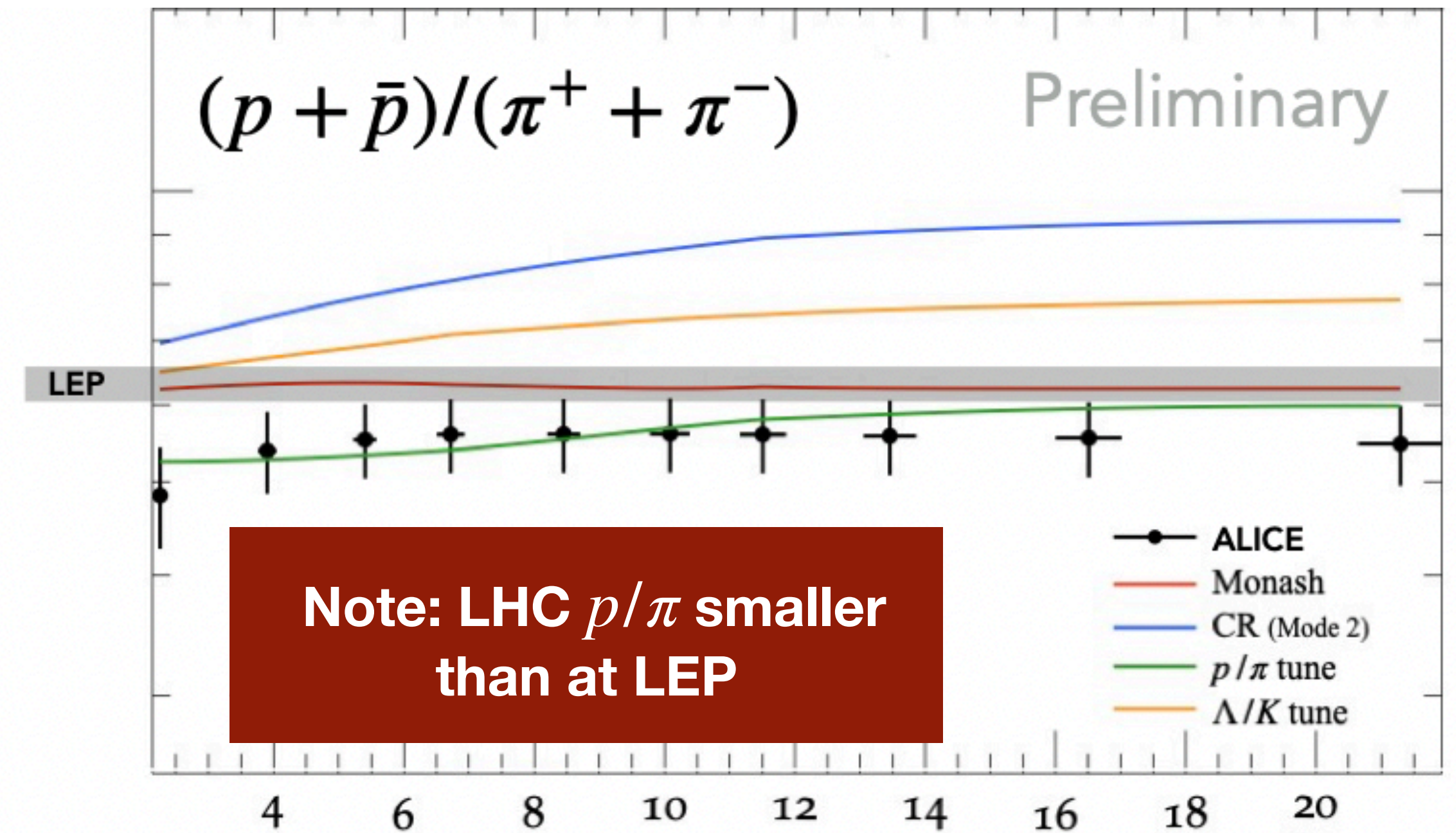
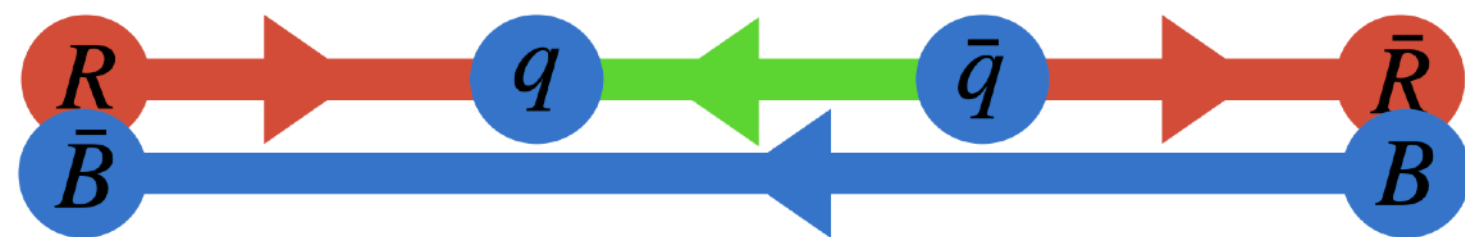
Popcorn Mechanism

arXiv:hep-ph/9606454

Diquark formation via **successive colour fluctuations**



What if there's a blue string nearby?

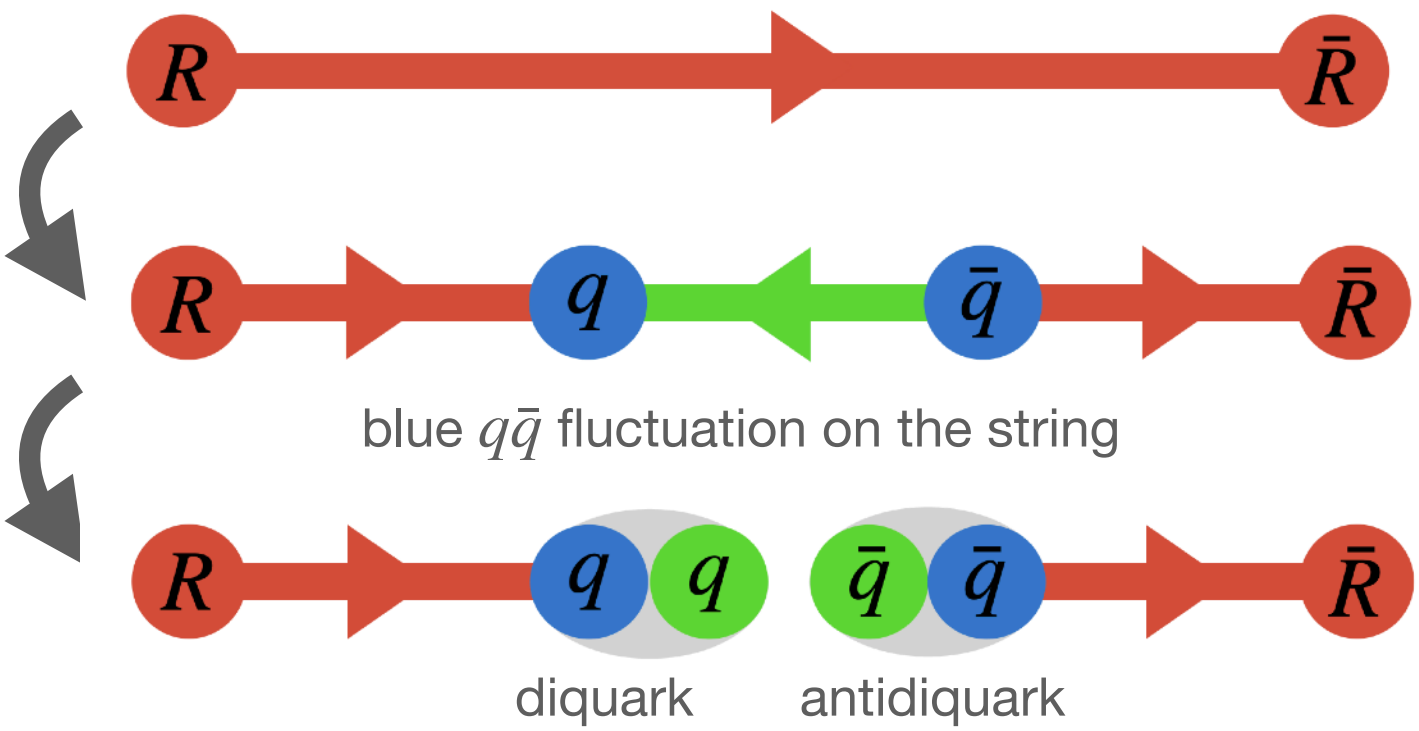


Popcorn mechanism

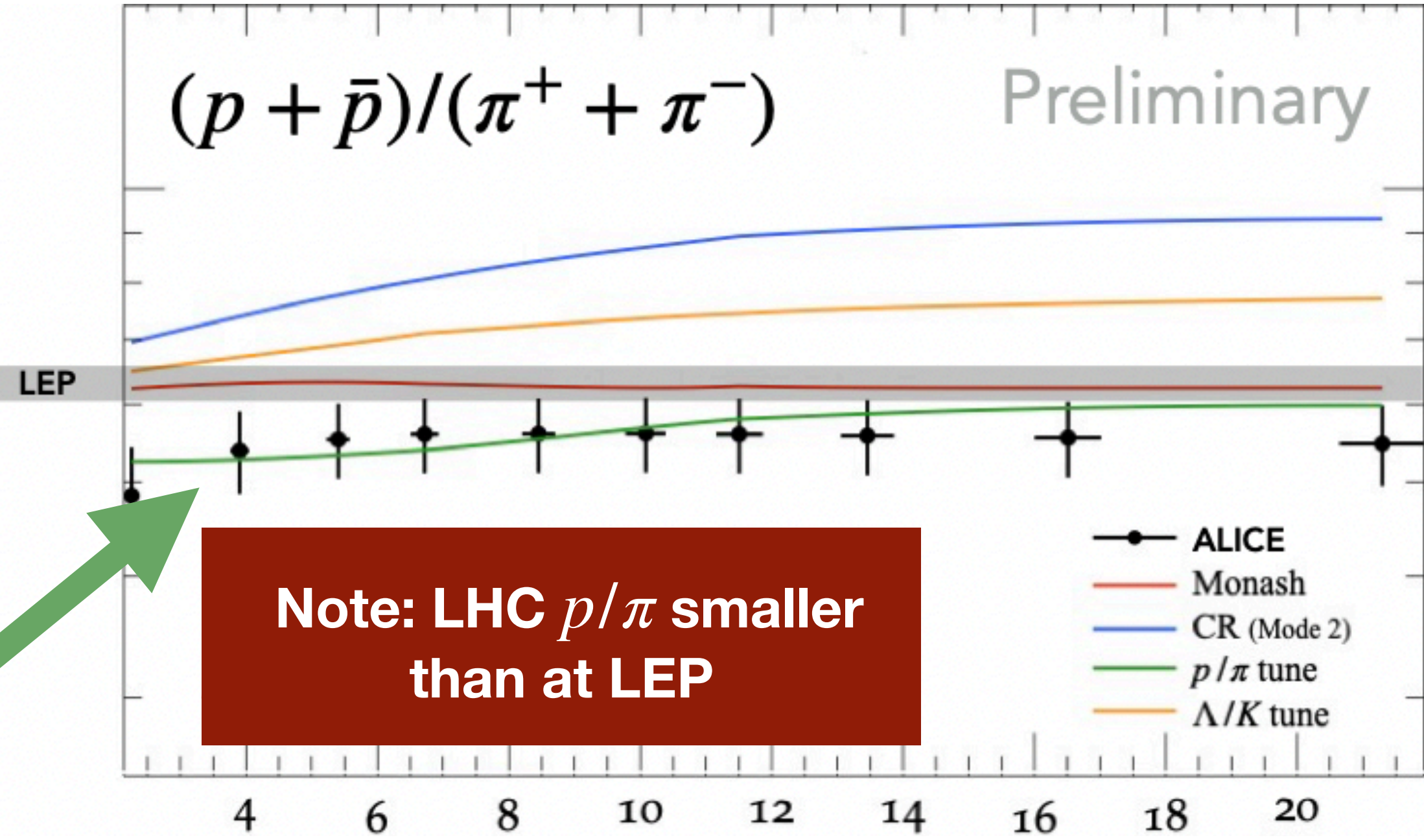
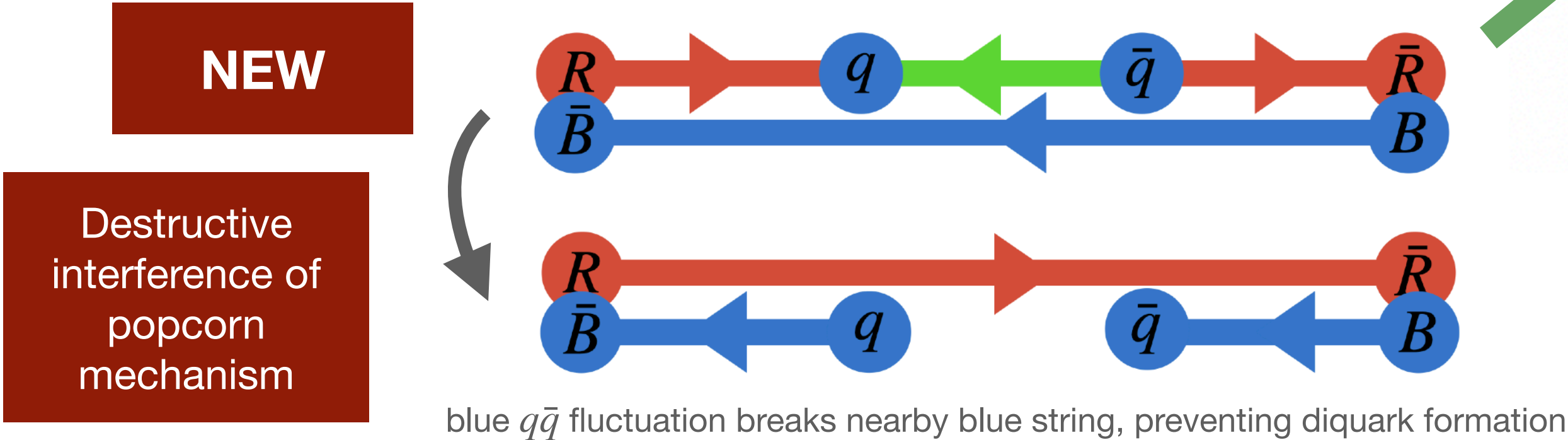
Popcorn Mechanism

arXiv:hep-ph/9606454

Diquark formation via **successive colour fluctuations**



What if there's a blue string nearby?



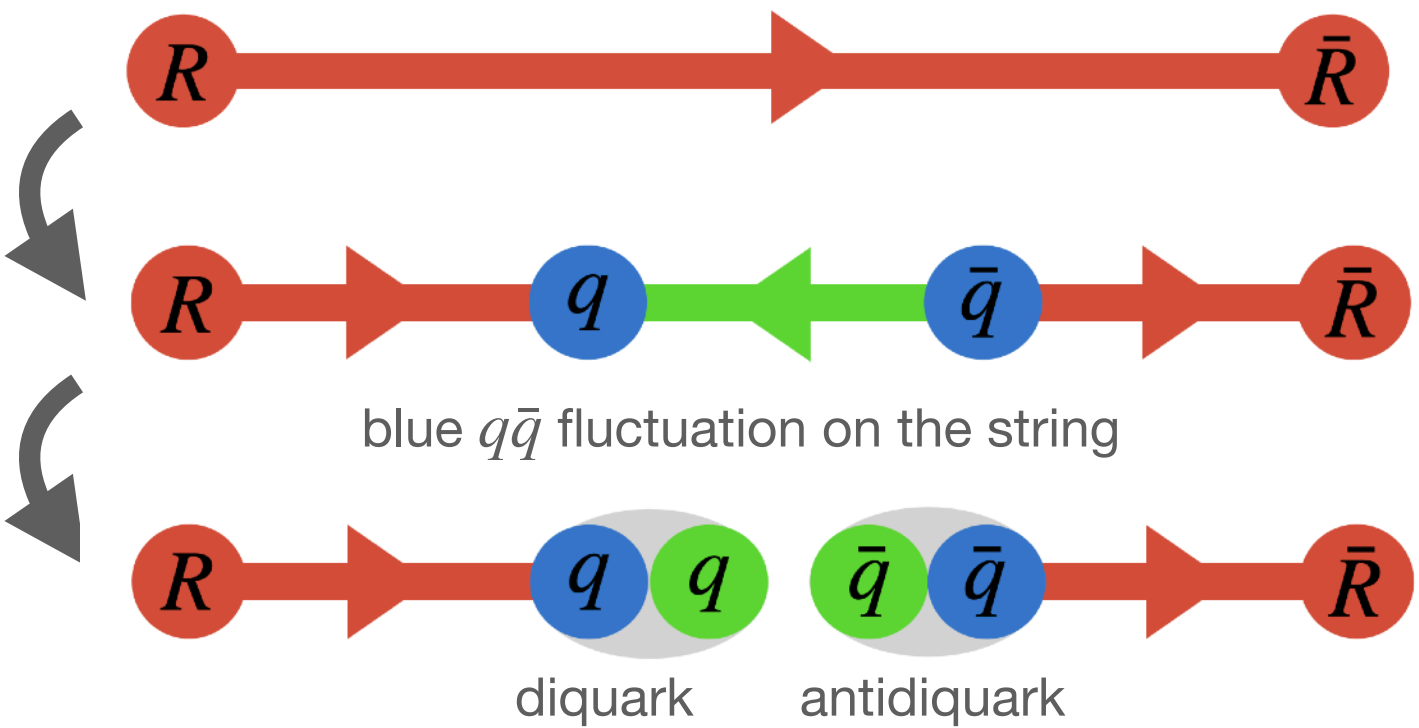
Only basic model implemented thus far, further improvements on the modelling still happening!

Popcorn mechanism

Popcorn Mechanism

arXiv:hep-ph/9606454

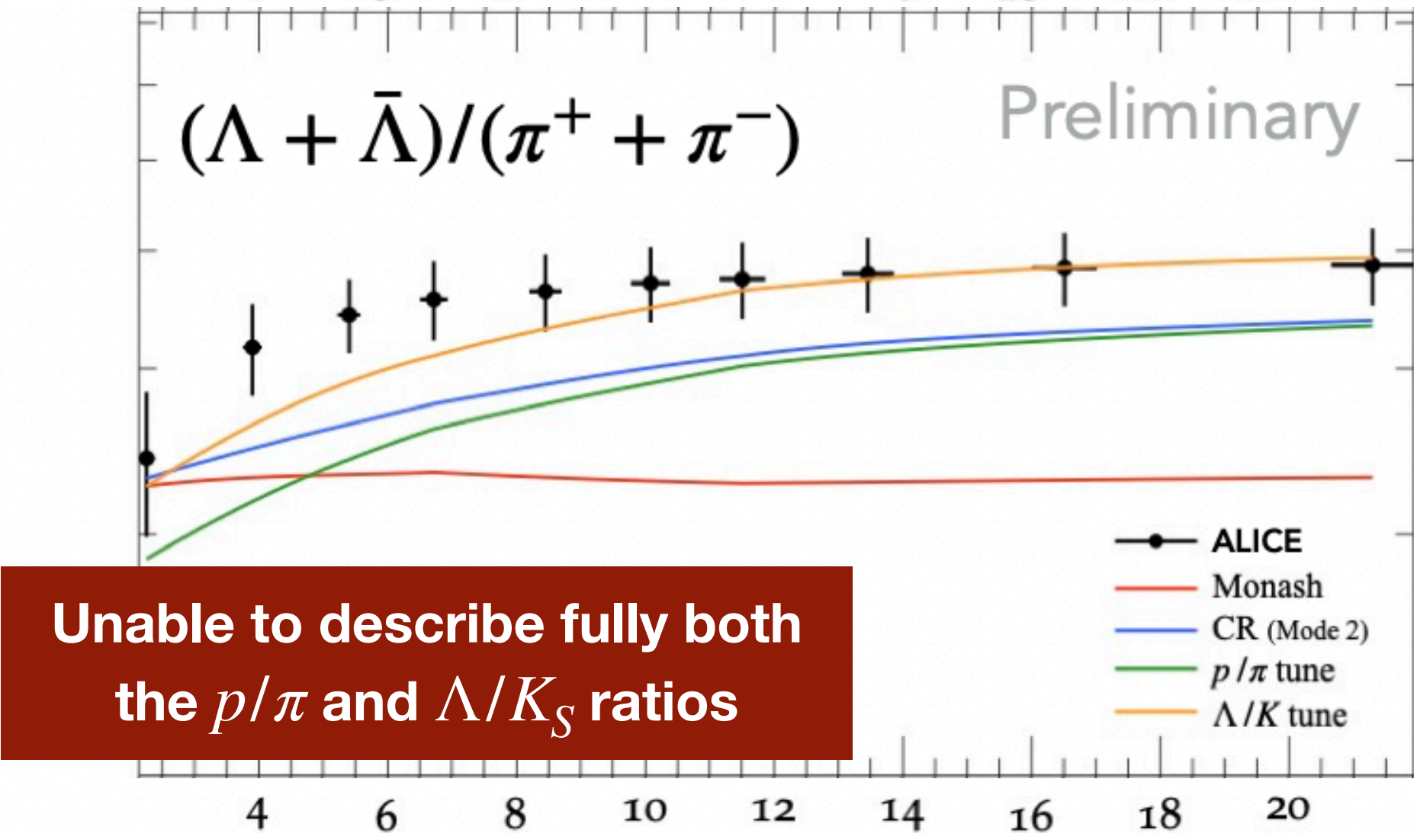
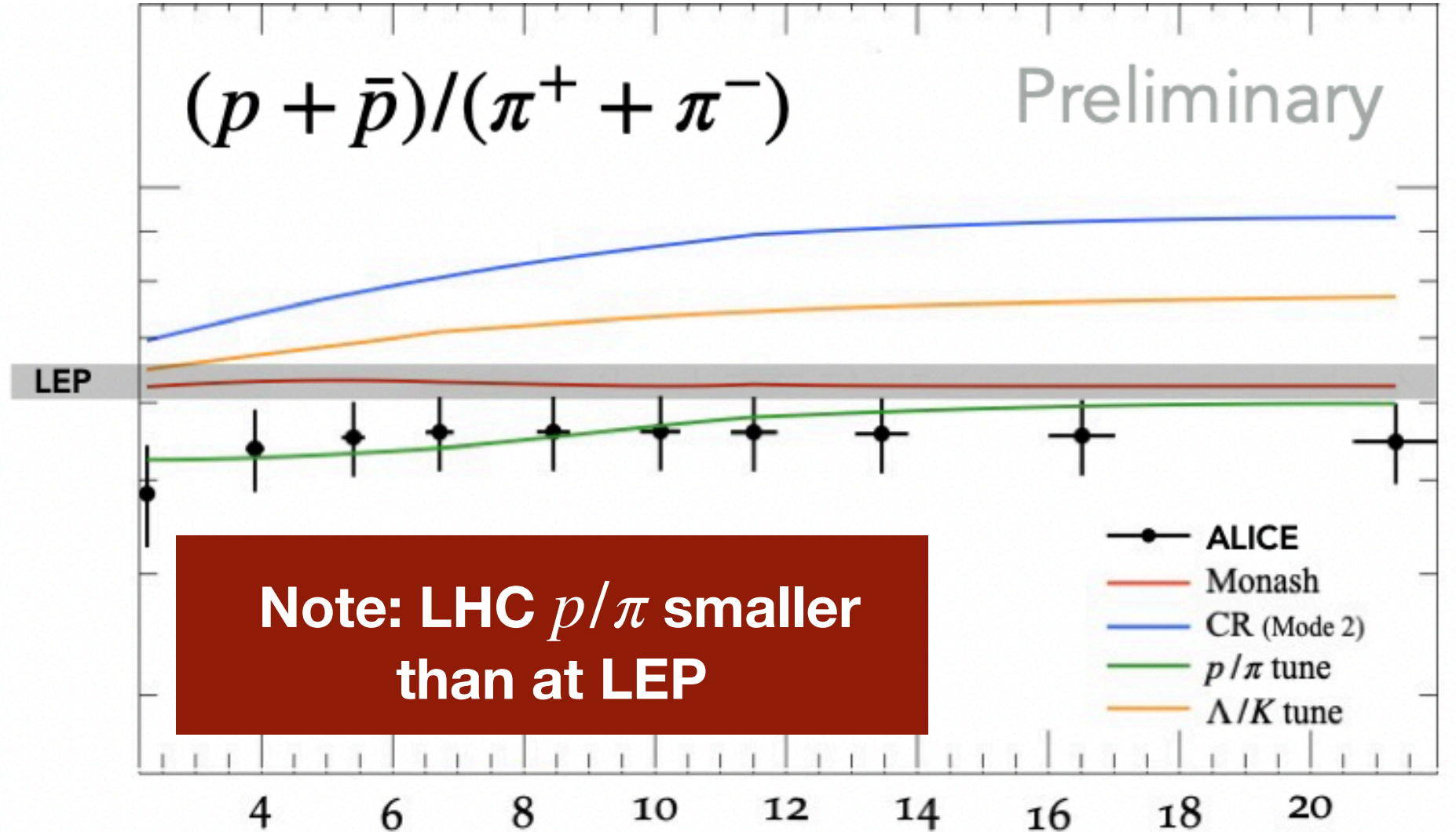
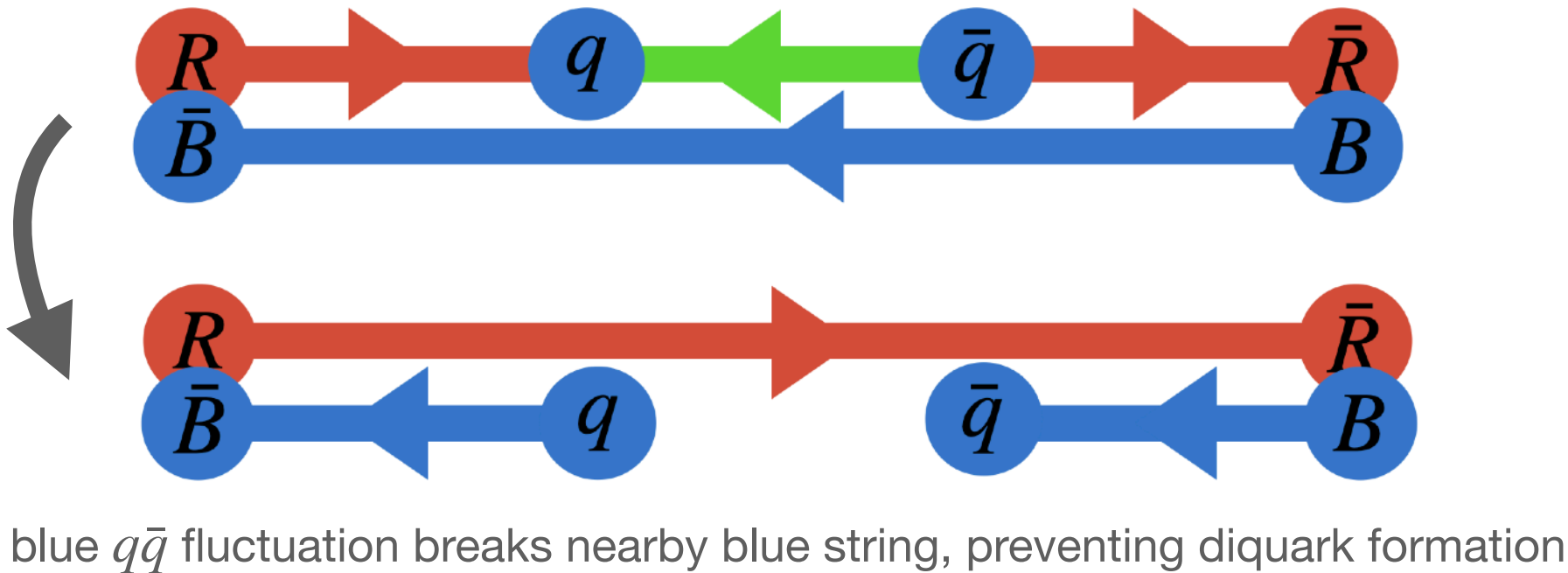
Diquark formation via **successive colour fluctuations**



What if there's a blue string nearby?

NEW

Destructive interference of popcorn mechanism



Vacuum → Small Systems → Heavy Ion

String model has well described e^+e^- **systems** (i.e. cases with not many strings), and we've explored **high multiplicity small systems**, but what about **heavy ion systems**?

Do we still have strings? Do we have QGP? Is it a mix of both, or is there a smooth transition between the strings and QGP?

Angantyr uses PYTHIA as its base to do pA and AA collisions, using only strings (**no QGP** formation)

How far can we push the string model?

Collective effects of strings can describe features that are typically described as signature of QGP

- Near-sided ridge → string shoving
- v_2 → string repulsion?
- Strangeness enhancement → ropes/close-packing

Thank you for listening!
