

High-energy neutrino and multi-messenger signatures from extreme astrophysical phenomena

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Pennsylvania State University

Fermilab Neutrino Seminar Series
Fermilab, Batavia, Illinois
March 21, 2024

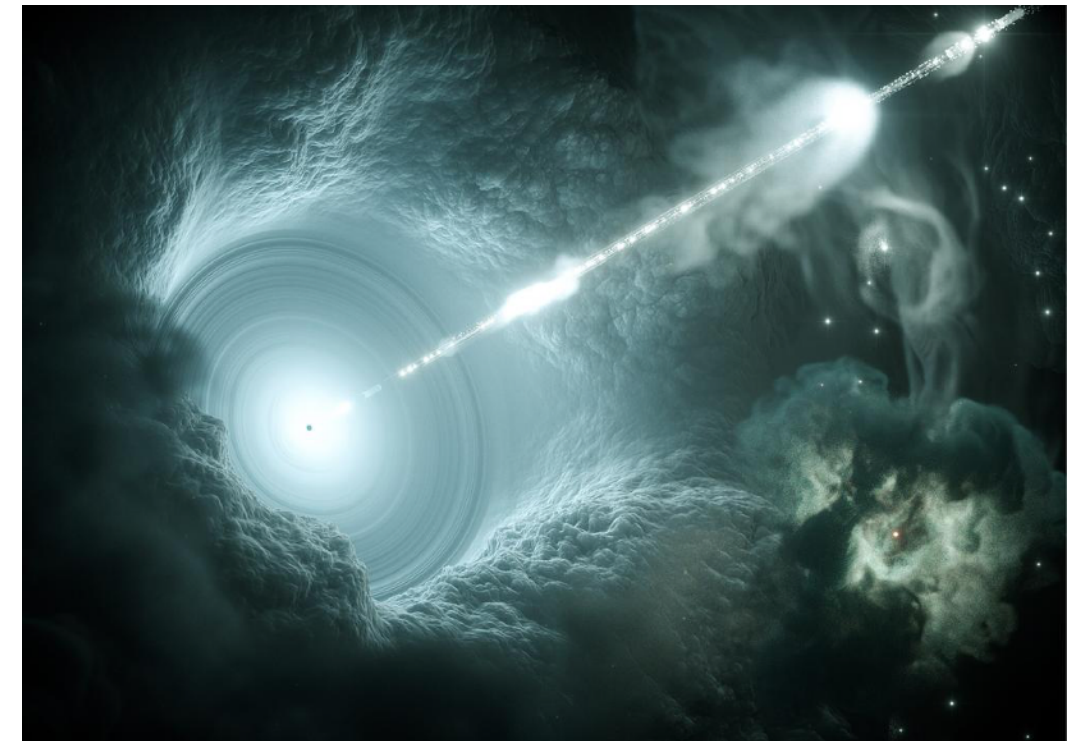
Prologue

New physics, understanding the fundamentals,....

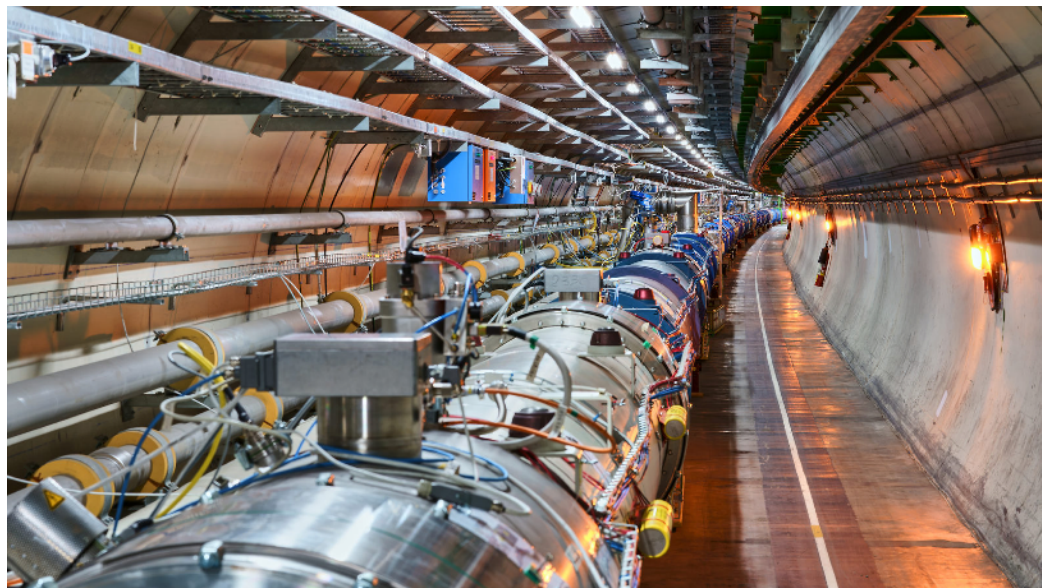
Man-made Accelerators

Cosmic Accelerators

Tevatron



LHC



Extreme astrophysical phenomena

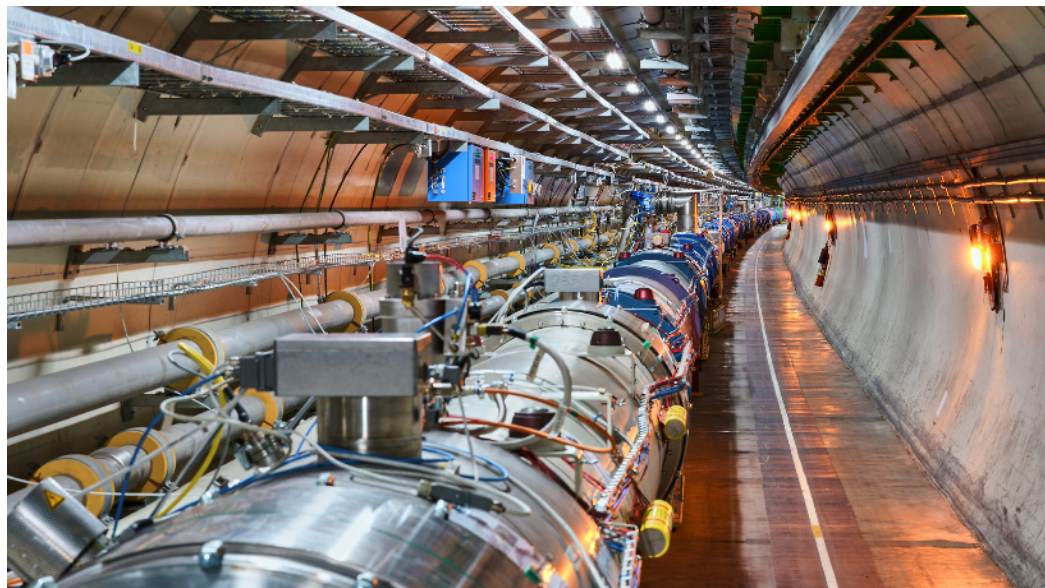
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New physics, understanding the fundamentals,....

Man-made Accelerators

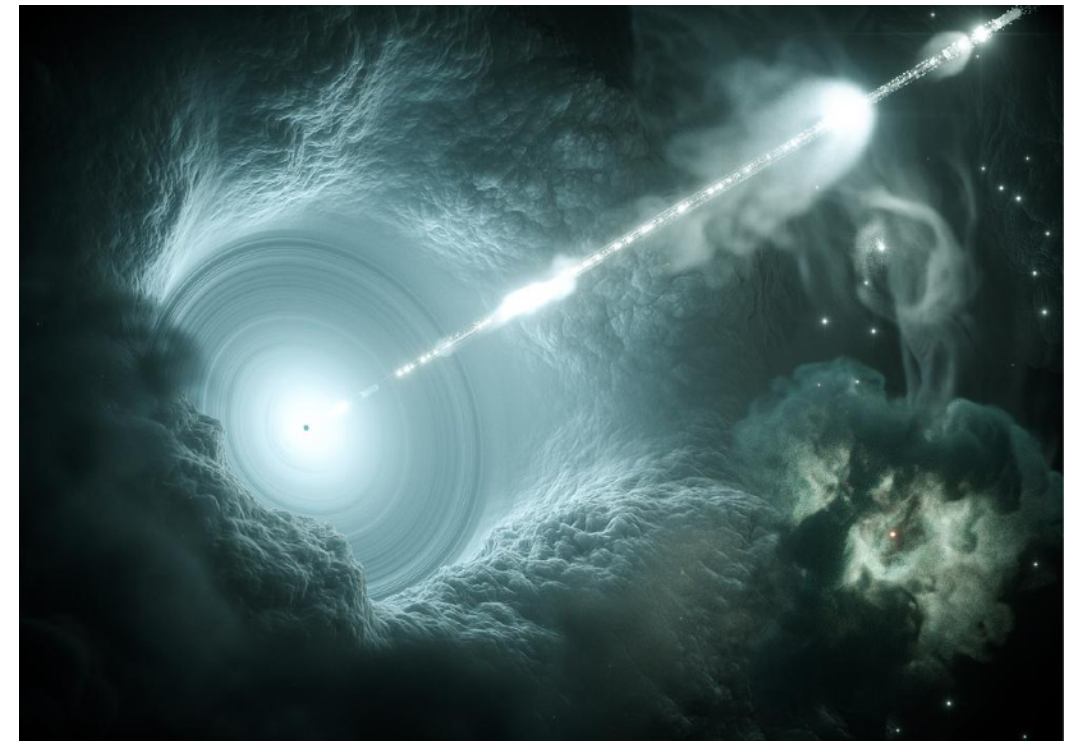


Tevatron



LHC

Cosmic Accelerators



Extreme astrophysical phenomena

The multi-messenger paradigm

Compact object mergers, TDEs, CCSNe,....

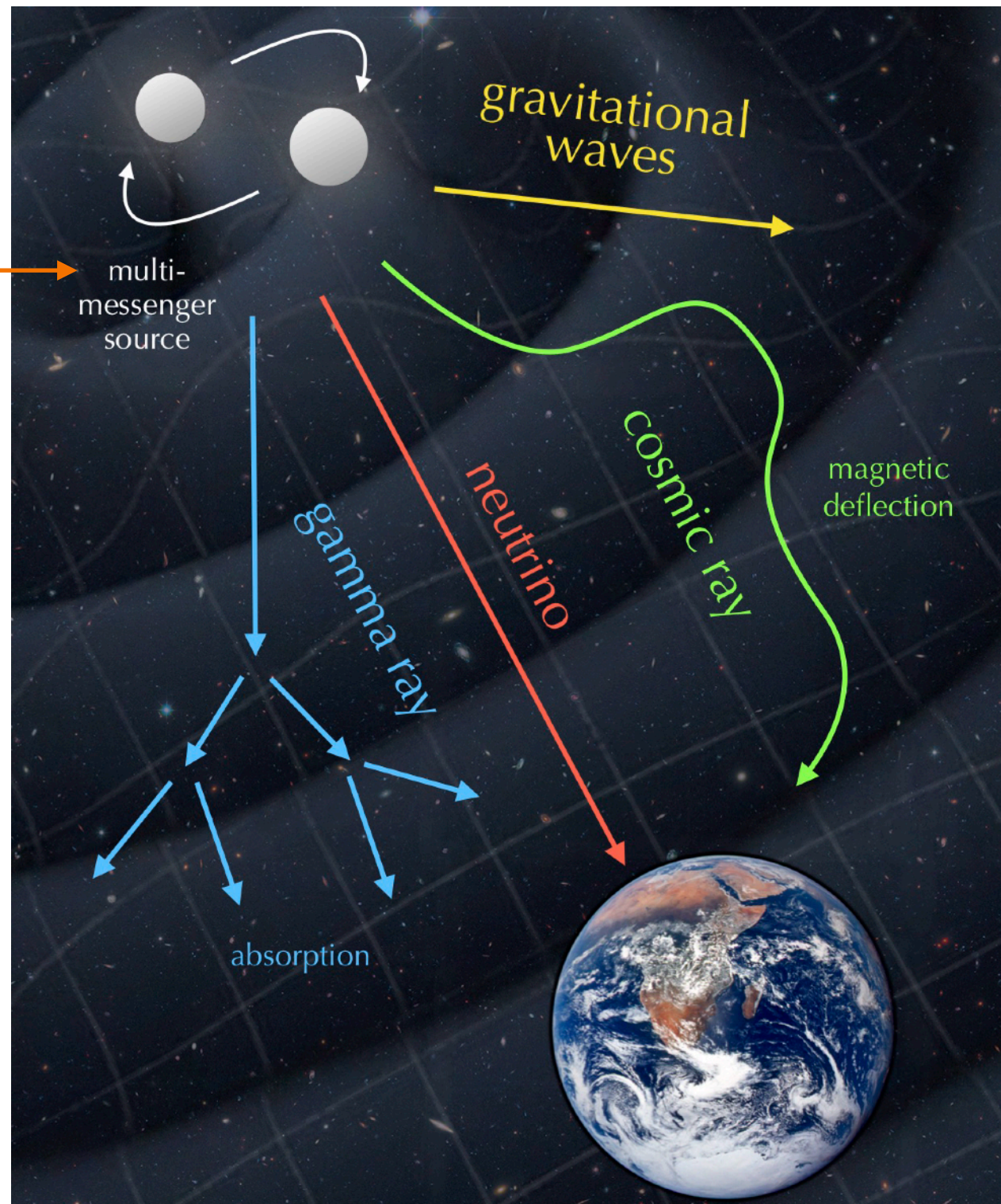
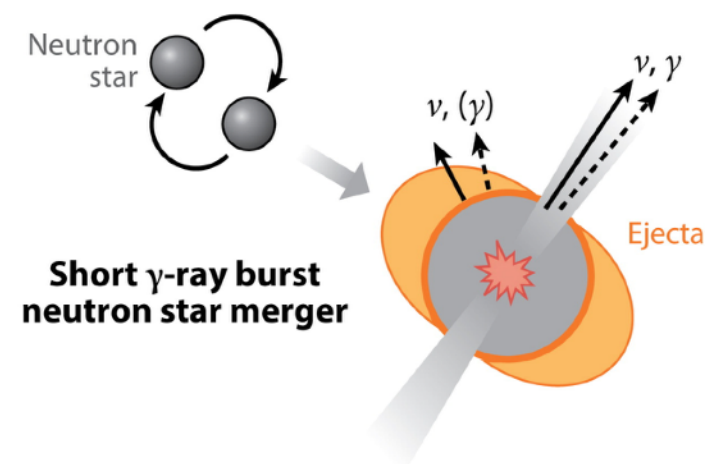
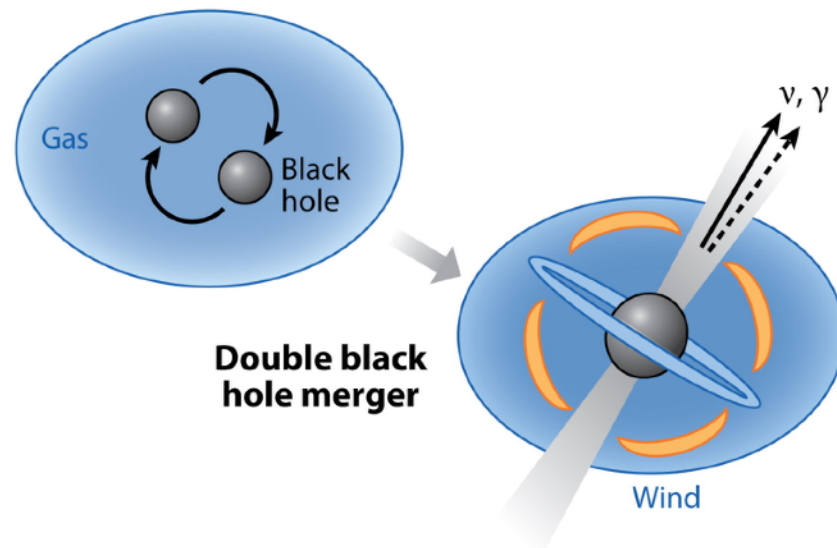
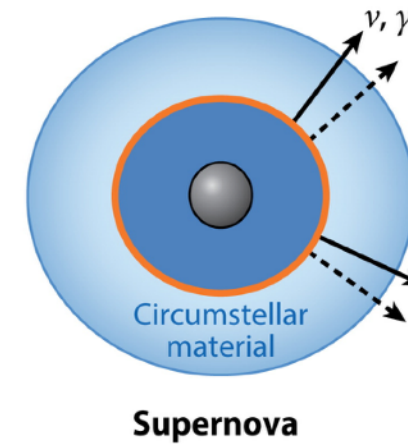
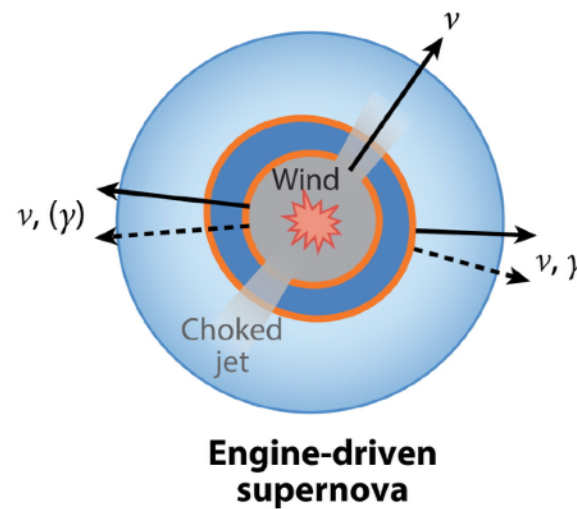
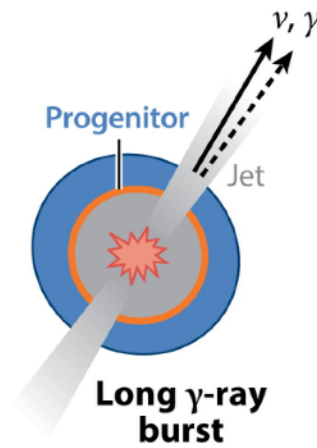
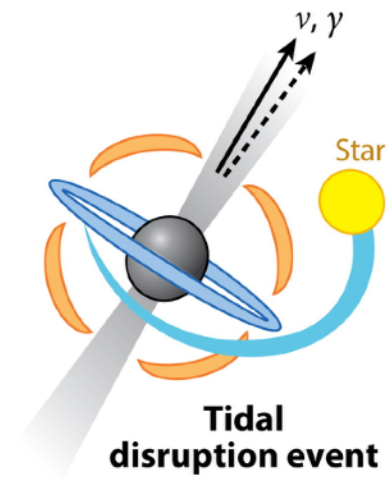
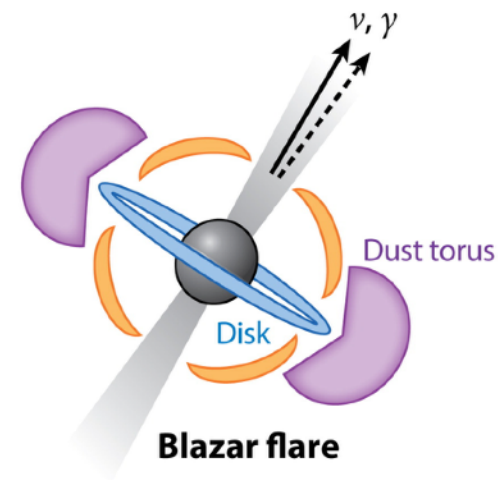


Image credits: <https://nbi.ku.dk/english/research/experimental-particle-physics/icecube/astroparticle-physics/>

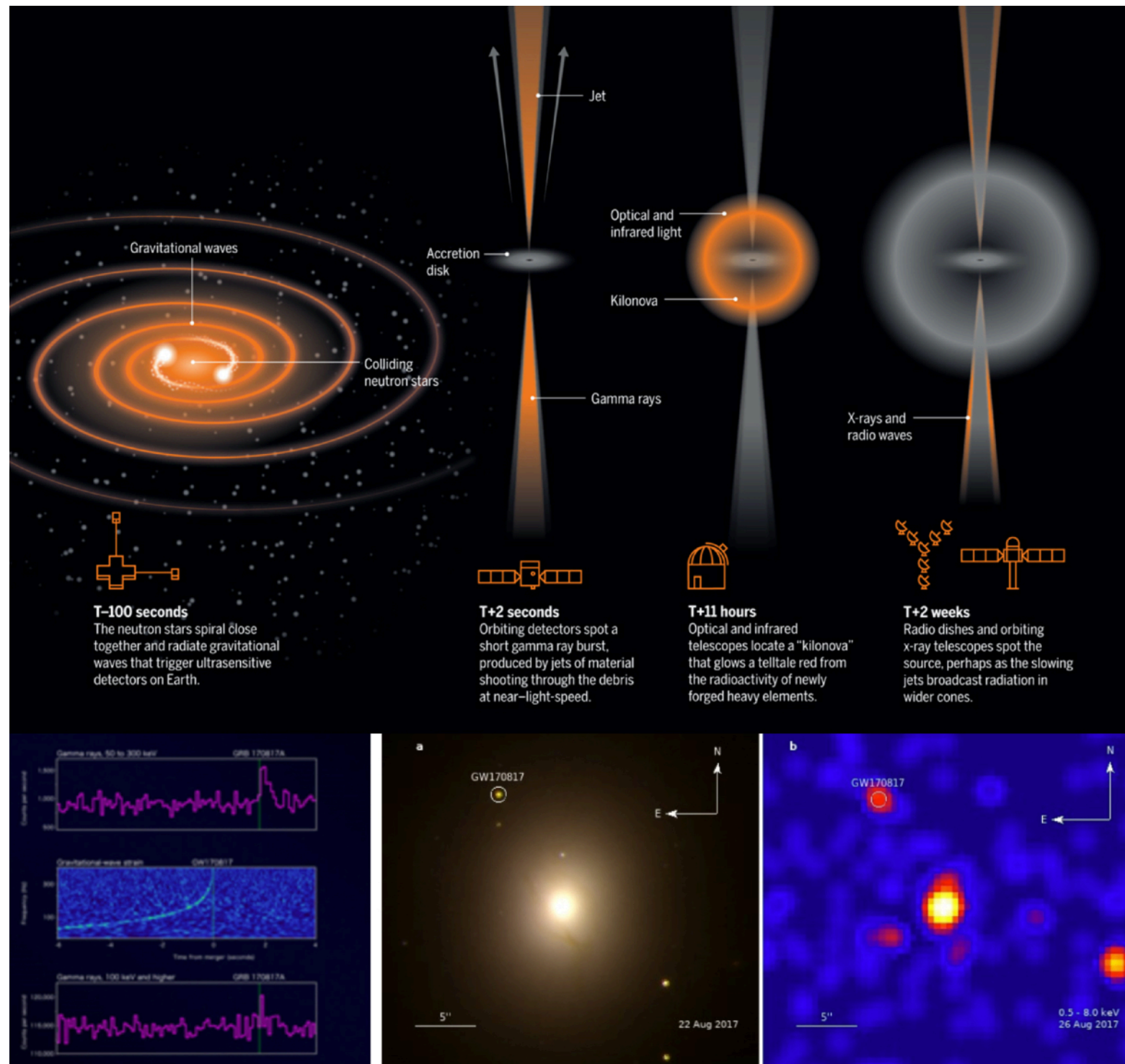
The high-energy multi-messenger transients

Extreme astrophysical phenomena



GW170817

~ 40 Mpc (NGC 4993)



No neutrinos :(

Gamma rays
(Fermi+Integral)

GW
(Adv. LIGO+Virgo)

X-rays
(Chandra)

Optical
(HST)

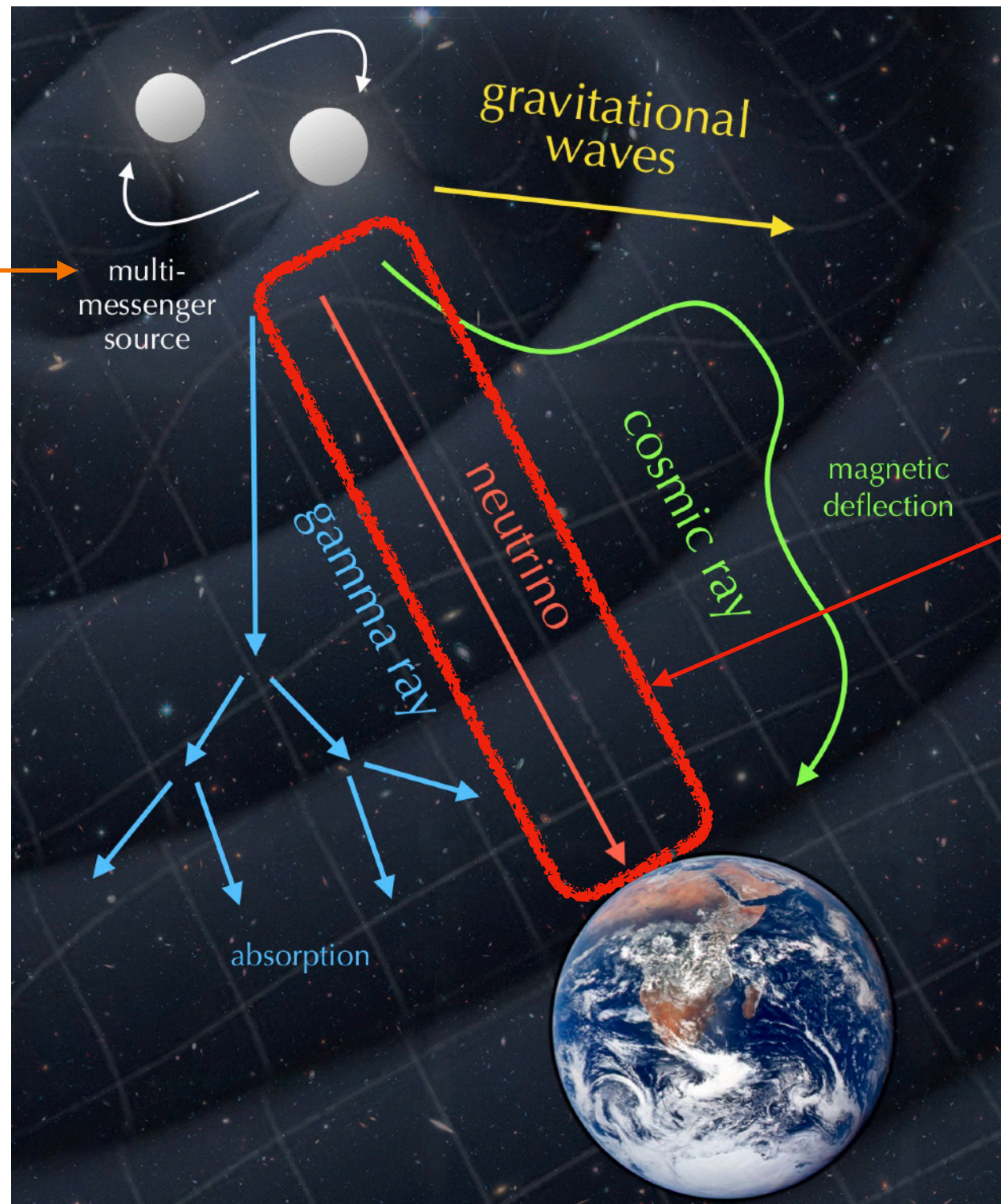
Image credits: <https://ahead.iaps.inaf.it>

Abbott et al. 2017, ApJ 848, L13

Troja, Piro, van Earthen et al., 2017, Nature, 551, 71 6

The multi-messenger paradigm

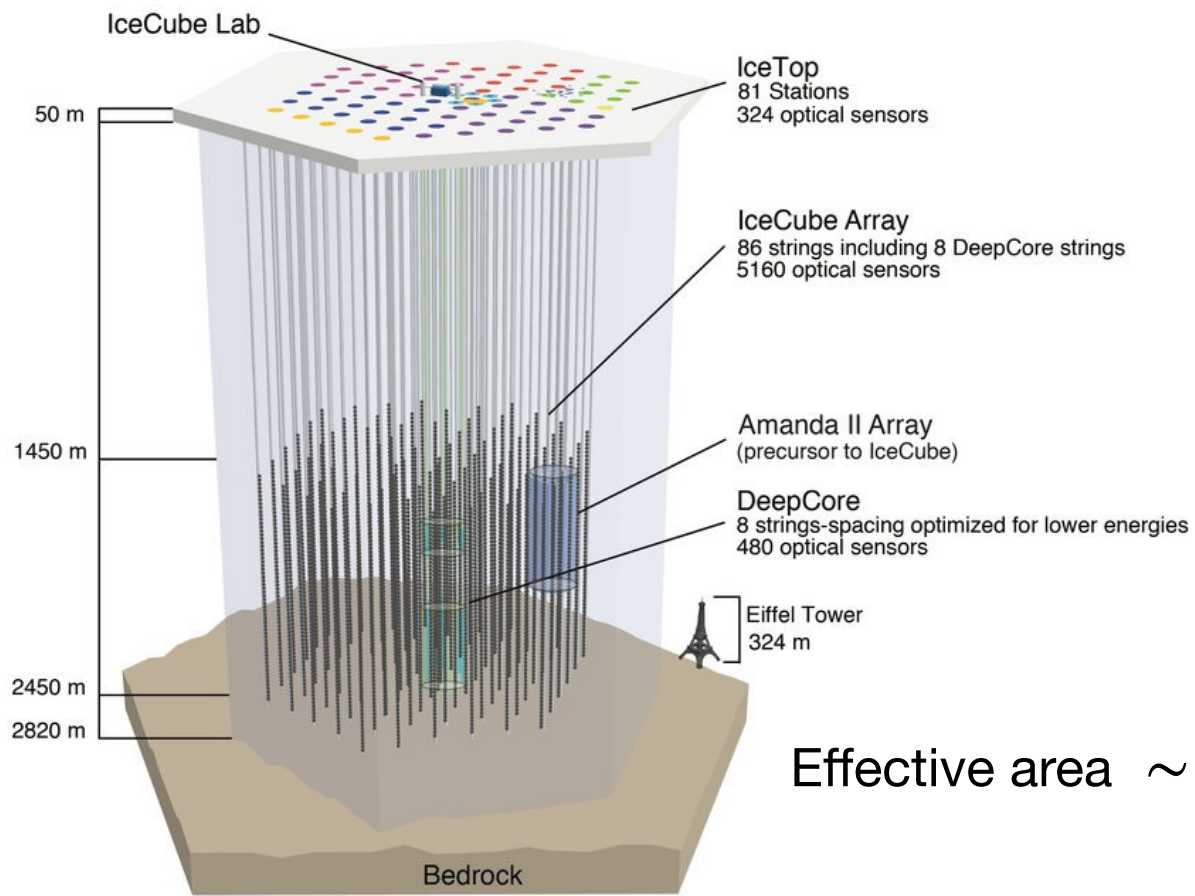
Compact object mergers, TDEs, CCSNe,....



High-energy neutrinos

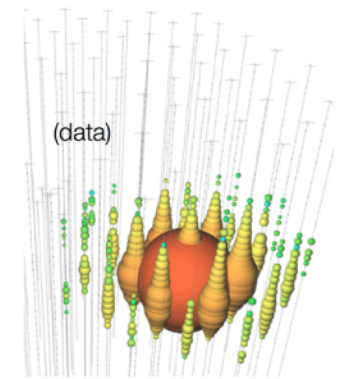
Image credits: <https://nbi.ku.dk/english/research/experimental-particle-physics/icecube/astroparticle-physics/>

High-energy neutrino detectors



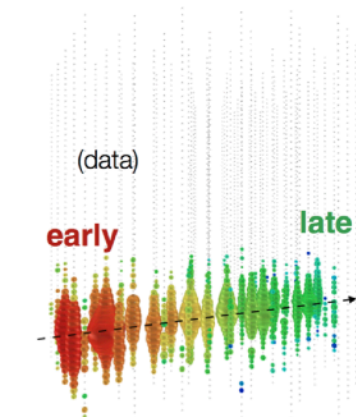
Effective area $\sim 1 \text{ km}^3$

Neutral-current / ν_e



Isolated energy deposition (cascade) with no track

Charged-current ν_μ



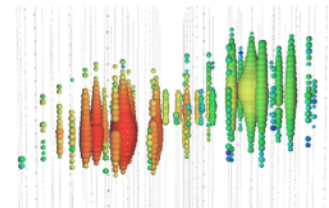
Up-going track

Charged-current ν_τ

IceCube observes seven astrophysical tau neutrino candidates

Posted on March 7, 2024 by Alisa King-Klemperer

(simulation)

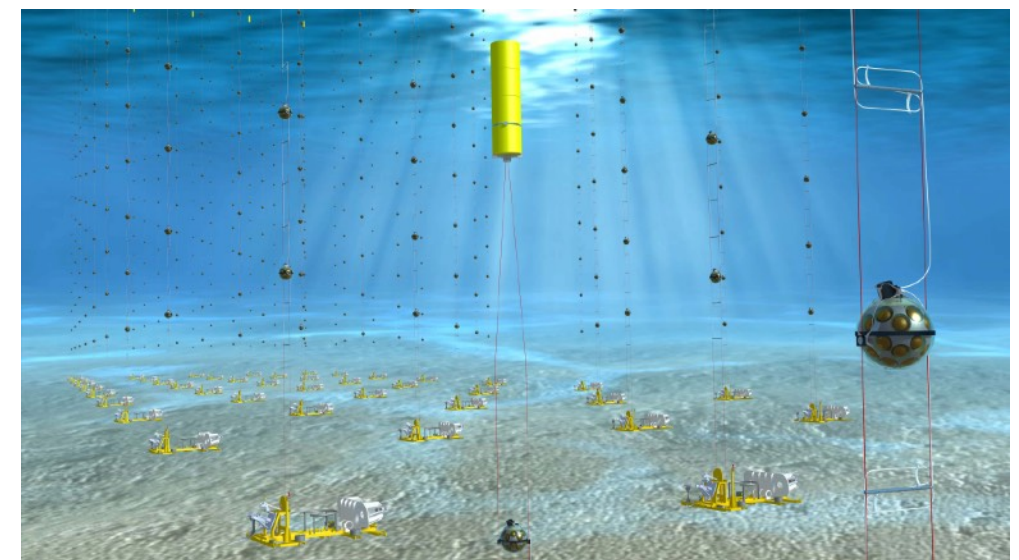


Double cascade

Baikal GVD

ANTARES

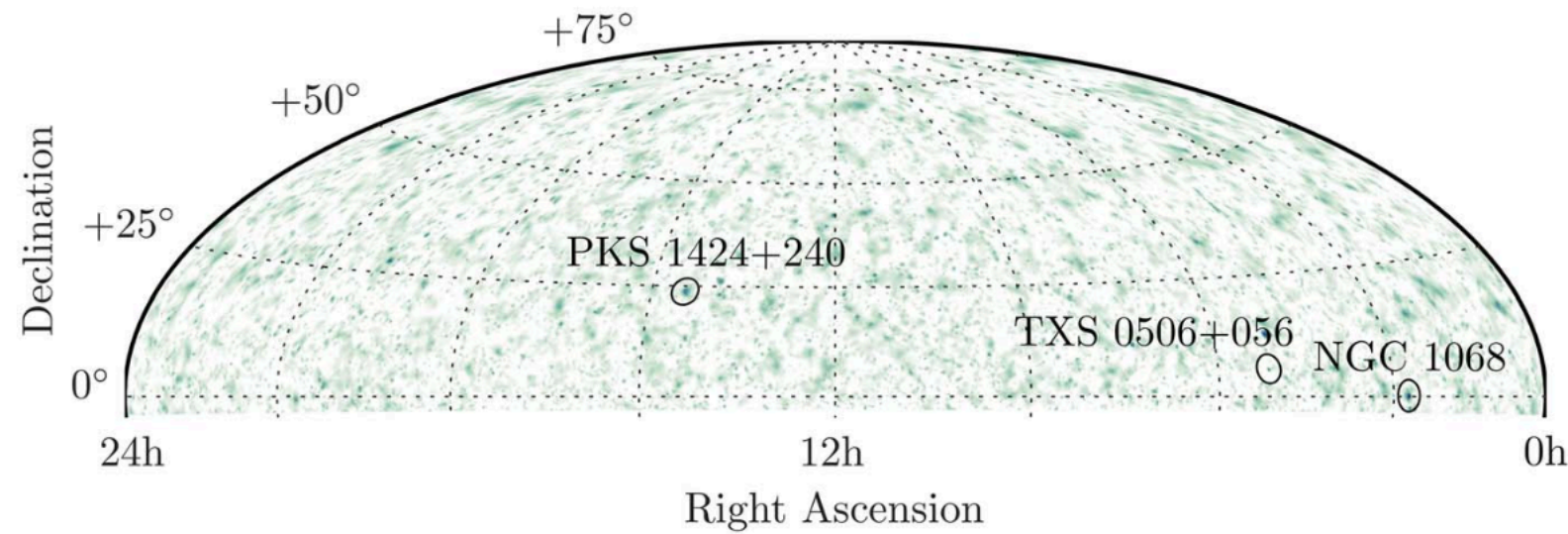
KM3NeT



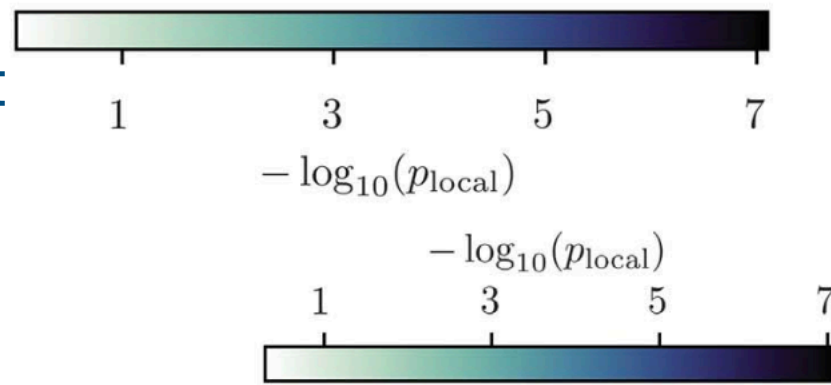
Future detectors: IceCube-Gen2, RNO-G, GRAND,.....

NGC 1068 (also TXS 0506+056)

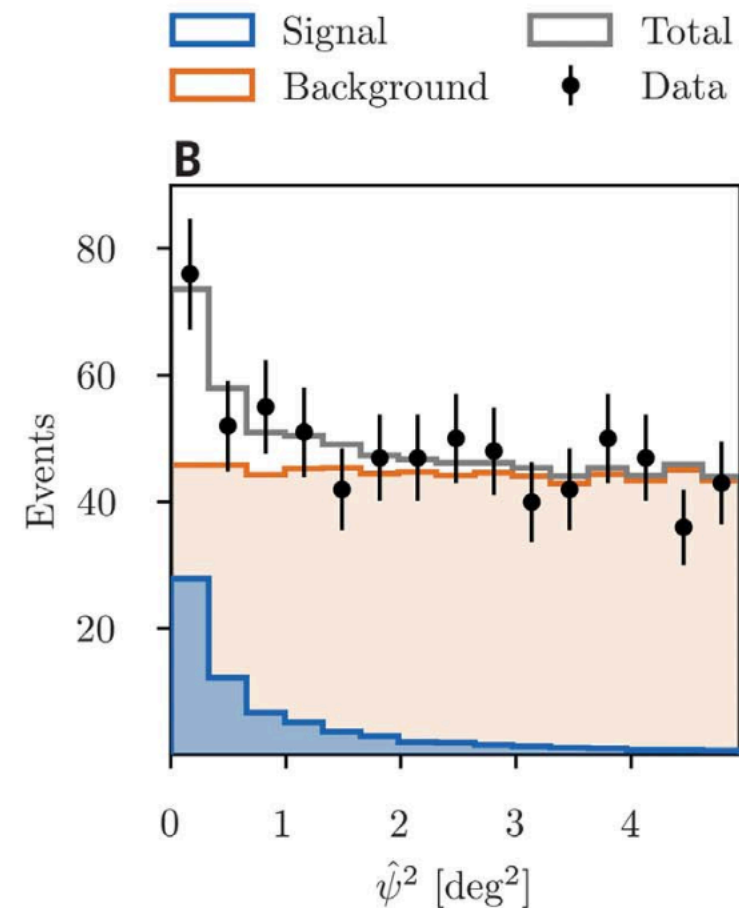
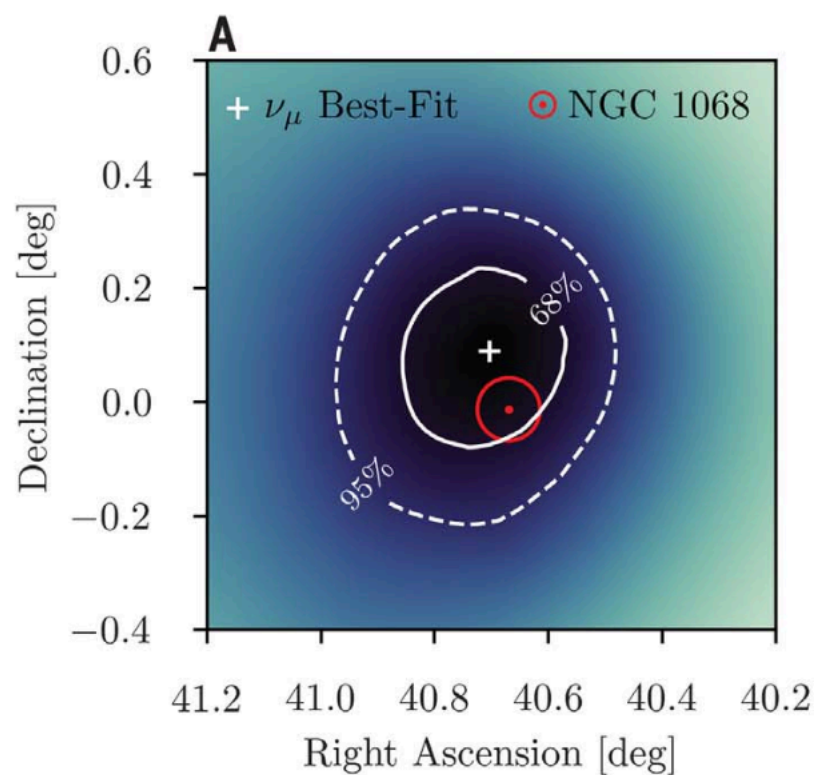
10 years of PS
data
(2011-2020)



$\sim 4.2\sigma$ w.r.t
110 known
gamma ray
sources



Source Name	Source Type	α [°]	δ [°]	\hat{n}_s	$\hat{\gamma}$	$-\log_{10} p_{\text{local}}$	$\Phi_{90\%}$
NGC 1068	SBG/AGN	40.67	-0.01	79	3.2	7.0 (5.2σ)	9.6
PKS 1424+240	BLL	216.76	23.80	77	3.5	4.0 (3.7σ)	11.4
TXS 0506+056	BLL/FSRQ	77.36	5.70	5	2.0	3.6 (3.5σ)	7.5

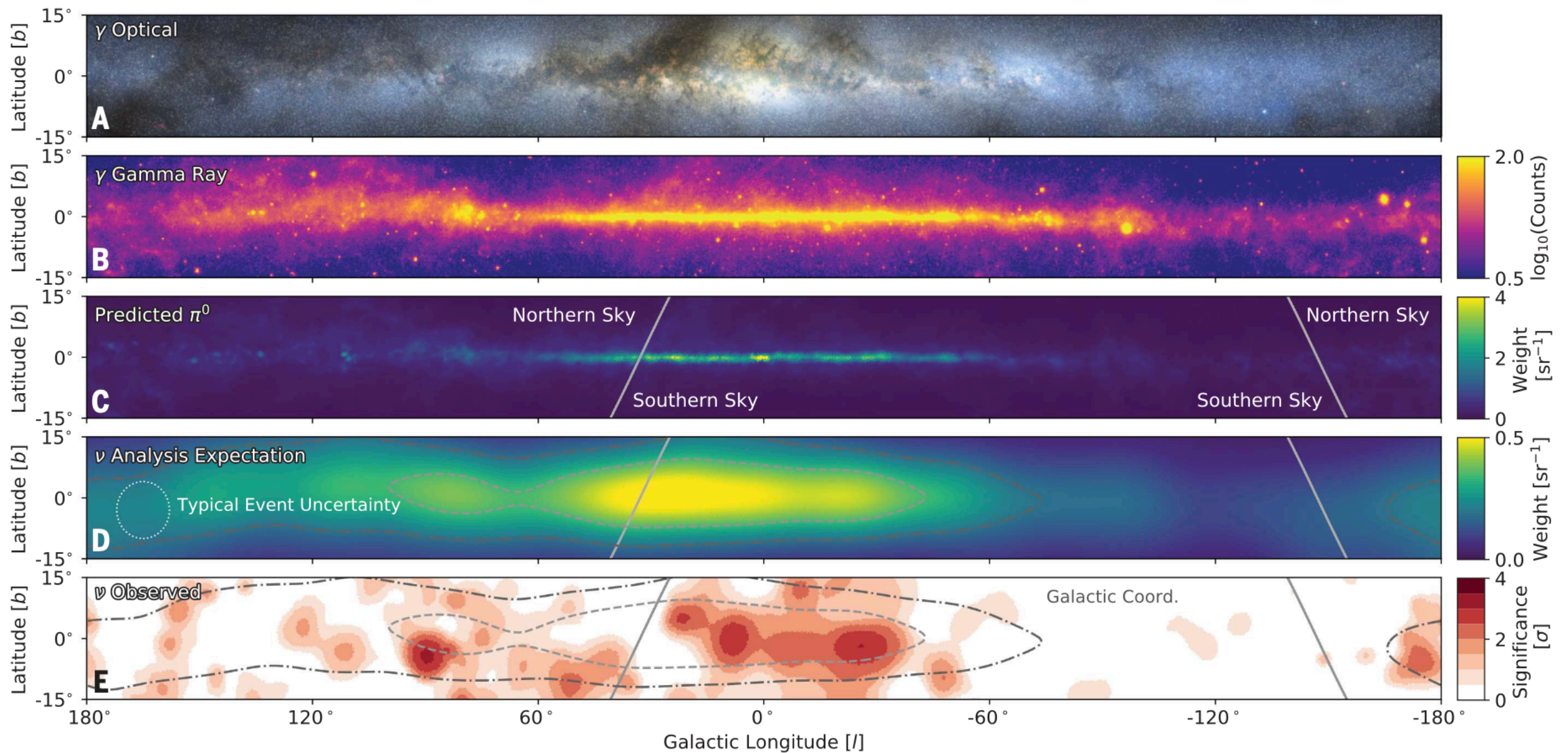


$\sim 79^{+22}_{-20}$
excess events

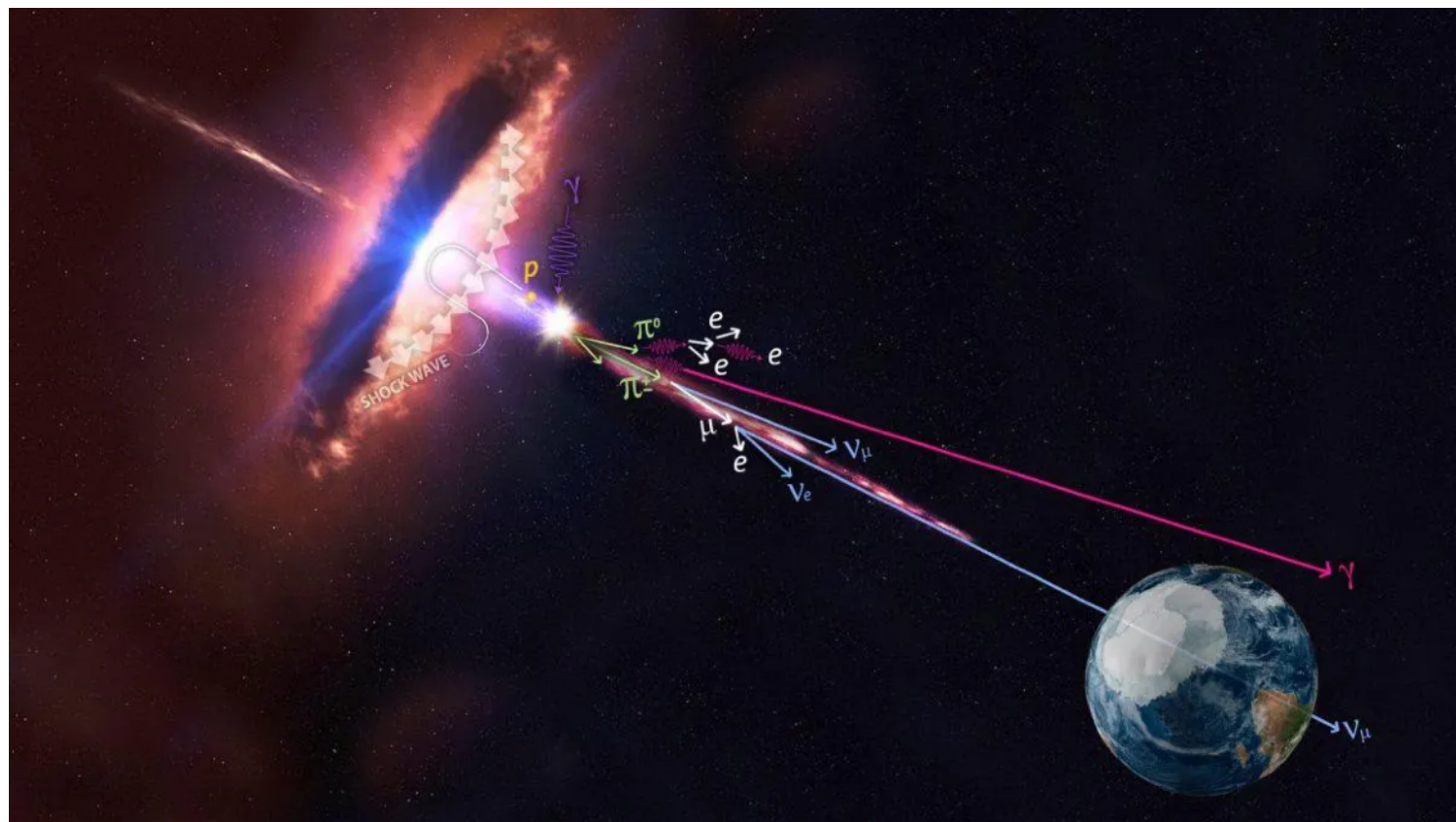
The Galactic plane

10 years of PS data
(2011-2020)

$\sim 4.5\sigma$ diffuse emission models
w.r.t background only hypothesis



High-energy neutrinos



$$p + p \rightarrow N\pi + X \quad p + \gamma \rightarrow N\pi + X$$

$$\pi^\pm \rightarrow \nu_\mu + \bar{\nu}_\mu + \nu_e(\text{or } \bar{\nu}_e) + e^\pm$$

$$\pi^0 \rightarrow \gamma + \gamma$$

Proton energy loss due to p-p interactions

Conditions for HE- ν production:

- Acceleration of ions (p and nuclei) to sufficiently high energies - Shocks, magnetic reconnection, stochastic acceleration aided by turbulence
- Rate of acceleration $>$ Rate of energy loss
- Significant density on target media - matter and radiation
- (a) and (b) \rightarrow production of charged mesons - pions that decay into neutrinos, charged leptons, and gamma-rays

$$t_{pp}^{-1} = n_N \kappa_{pp} \sigma_{pp} c$$

Nucleon density \rightarrow n_N Proton inelasticity \rightarrow κ_{pp} p-p cross-section \rightarrow σ_{pp}

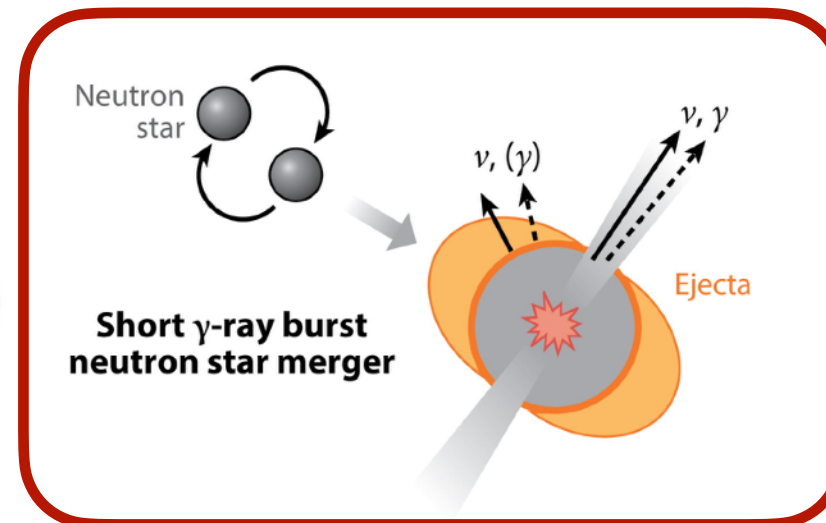
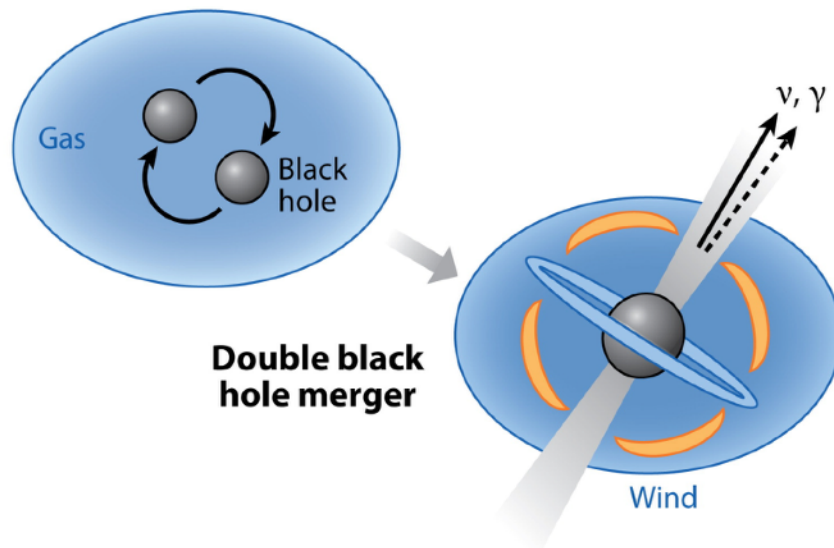
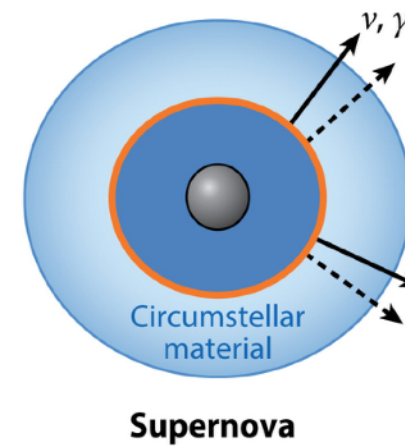
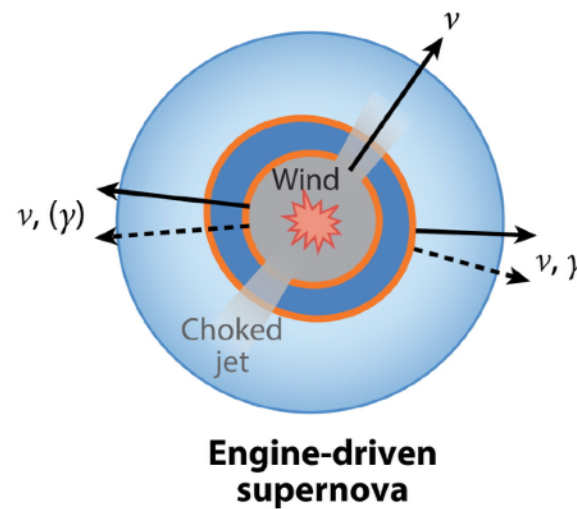
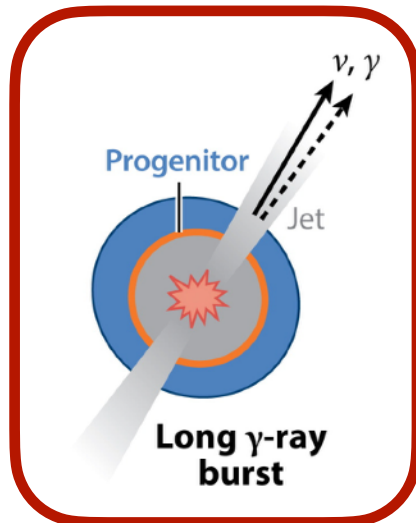
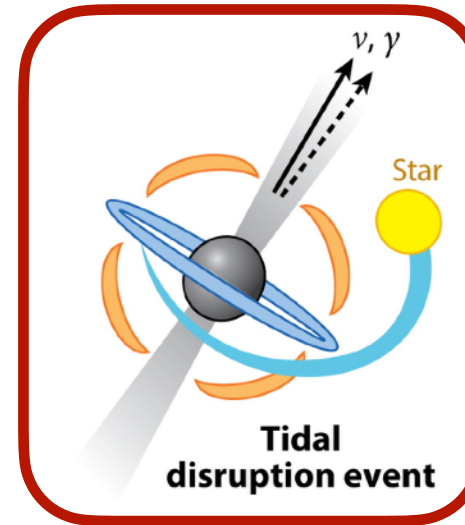
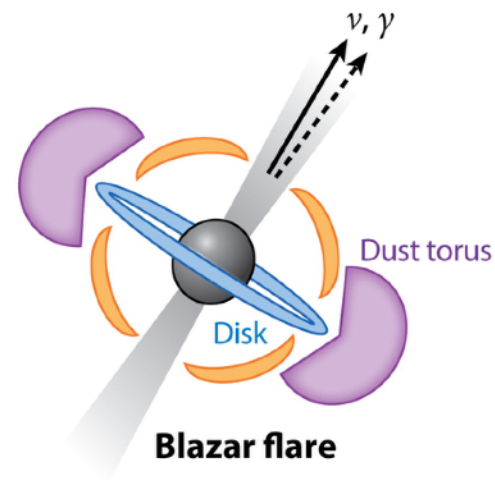
$$t_{p\gamma}^{-1}(\epsilon_p) = \frac{c}{2\gamma_p^2} \int_{\bar{\epsilon}_{th}}^{\infty} d\bar{\epsilon} \kappa_{p\gamma}(\bar{\epsilon}) \sigma_{p\gamma}(\bar{\epsilon}) \bar{\epsilon} \int_{\bar{\epsilon}/2\gamma_p}^{\infty} d\epsilon \epsilon^{-2} n_\epsilon$$

Proton energy \rightarrow ϵ_p Photon energy in proton rest frame \rightarrow $\bar{\epsilon}$ p- γ cross-section \rightarrow $\sigma_{p\gamma}$

Proton energy loss due to p-p interactions

The high-energy multi-messenger transients

Extreme astrophysical phenomena



Outline

Part 1: Can choked delayed jets explain the neutrino coincidences associated with TDEs?

Based on: [Multi-messenger signatures of delayed choked jets in tidal disruption events](#)
[MM](#), M. Bhattacharya, K. Murase
[\(submitted to MNRAS\)](#) ([arXiv: 2309.02275](#)).

Part 2: Hunting for neutrinos from BNS mergers at next-generation GW and neutrino detectors

Based on: [Gravitational wave triggered high energy neutrino searches from BNS mergers: prospects for next generation detectors](#)
[MM](#), S. S. Kimura, K. Murase
[Phys. Rev. D 109, 4, 043053 \(2024\)](#) ([arXiv: 2310.16875](#))

Part 3: Constraints from non-detection of neutrinos from the BOAT - GRB 221009A

Based on: [Neutrinos from the Brightest Gamma-Ray Burst?](#)
K. Murase, [MM](#), A. Kheirandish, S. S. Kimura, K. Fang
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Tidal disruption event (TDE)

The shredding apart of a star when it comes close to a SMBH, due to its tidal forces

Disruption starts



~ Half the debris is lost:
Unbound orbit

~ Half the debris falls back:
Bound orbit



Debris circularizes

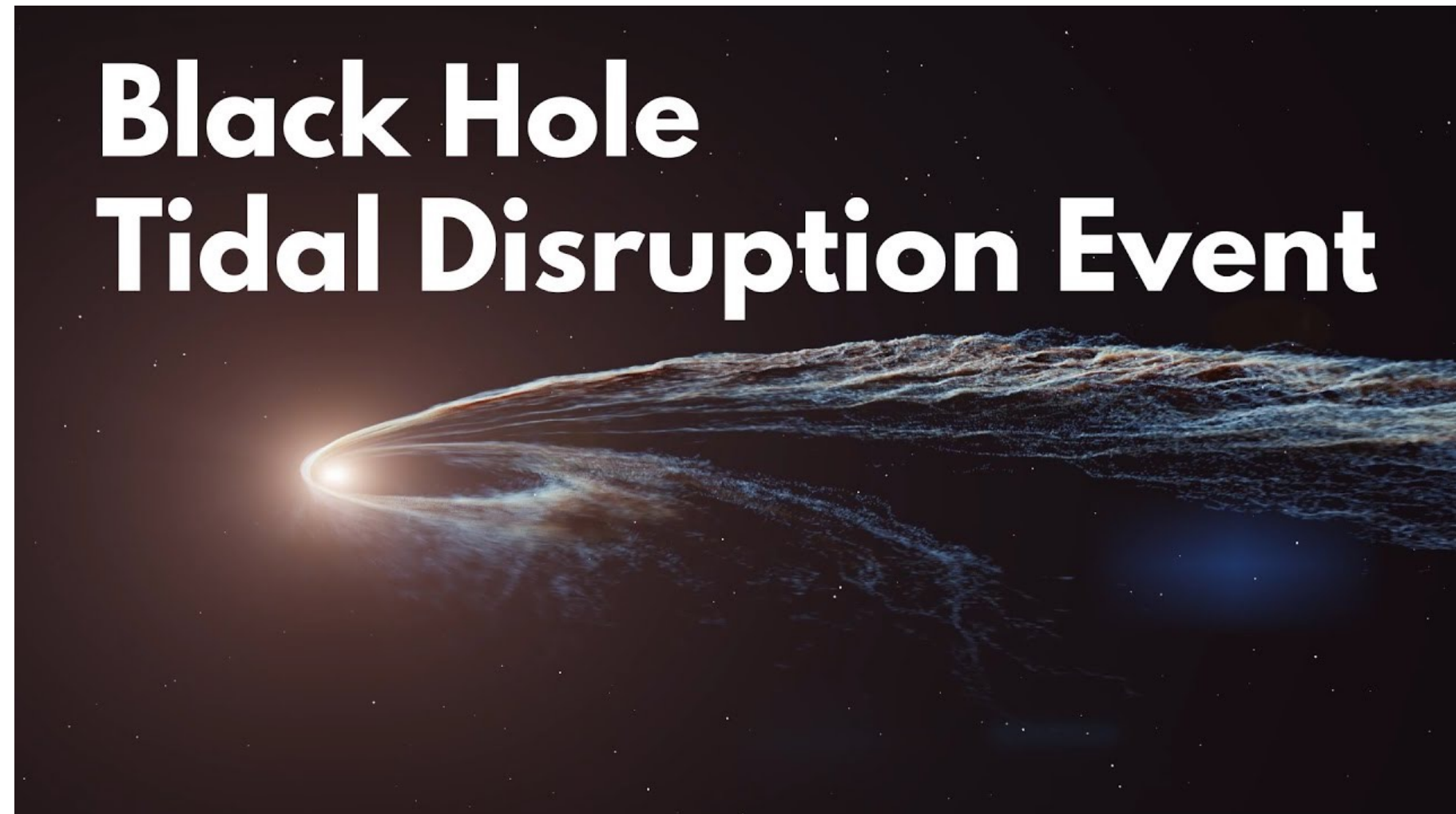


Part of the debris may form an
accretion disk



Winds, Outflows,
etc.

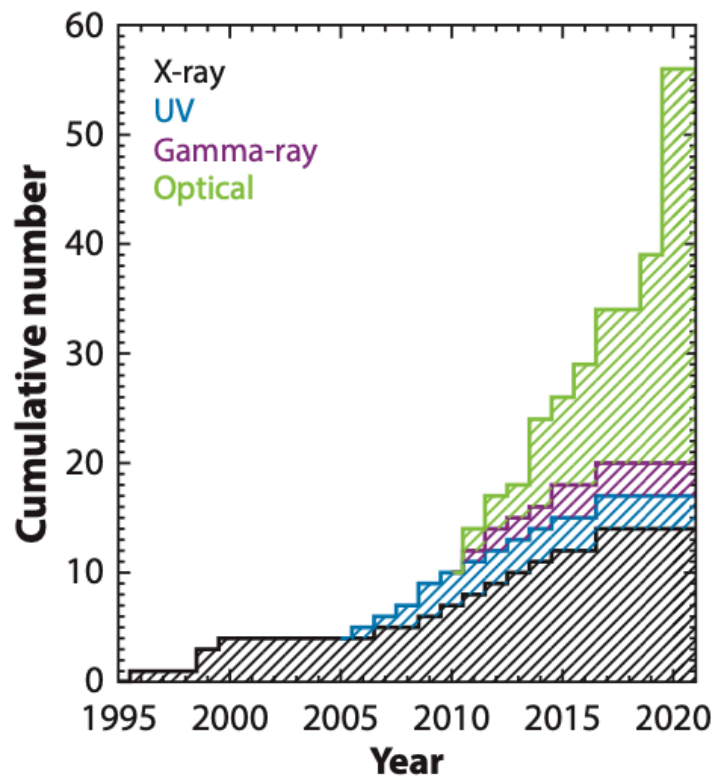
(Timescales are also uncertain)



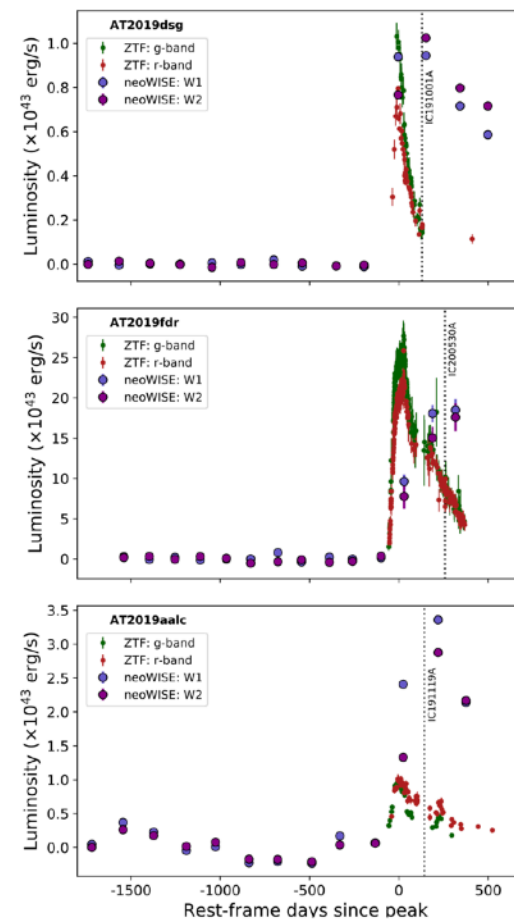
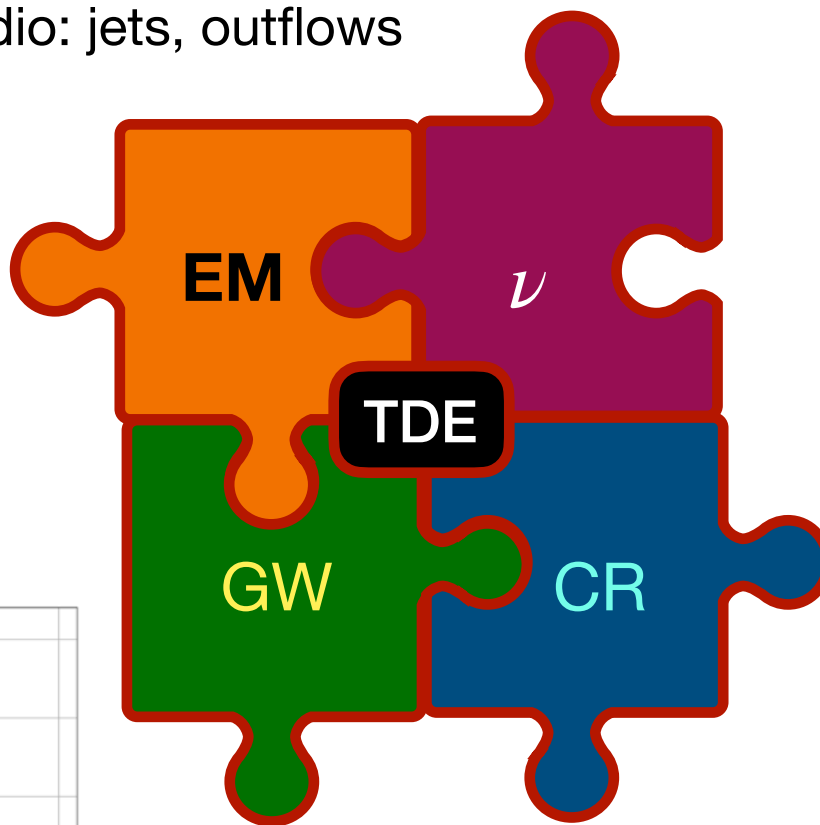
Credits: Science Communication Lab/DESY

TDEs: particle accelerators and multi-messenger zoo

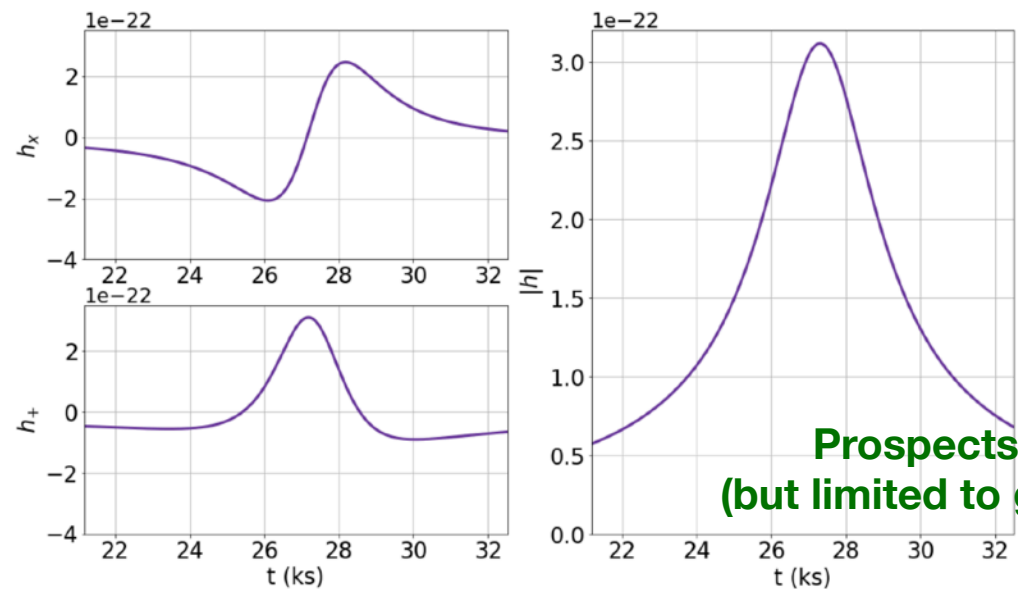
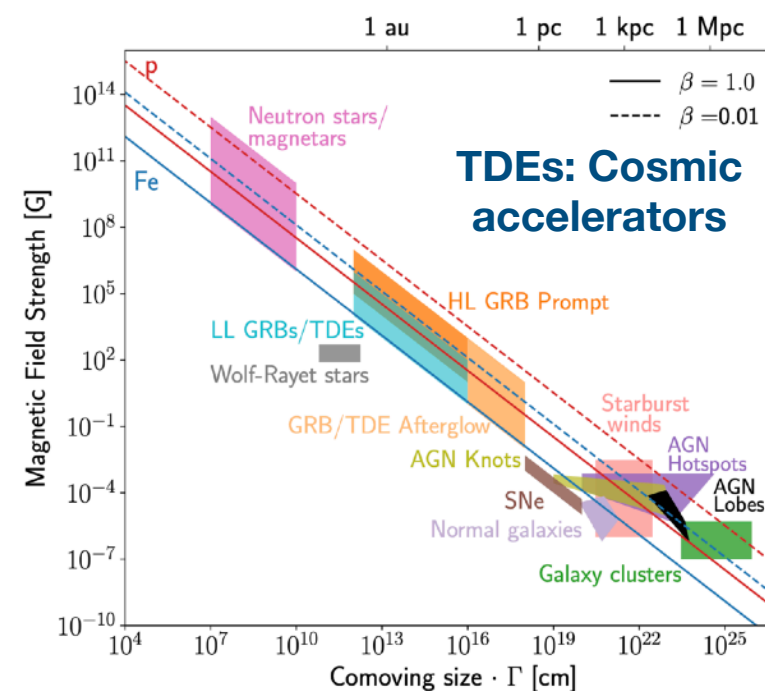
Observed



Optical-UV (blackbody): mass fallback $\sim t^{-5/3}$
X-rays: wide variability
Radio: jets, outflows



Associations



Prospects for LISA: IMBH-WD
(but limited to galactic scales ~ 10 kpc)

$$h \sim 2 \times 10^{-22} \beta (D/10\text{Mpc})^{-1} m_*^{4/3} r_*^{-1} M_{\text{BH},6}^{2/3}$$

$$f \sim 6 \times 10^{-4} \text{Hz} \beta^{3/2} m_*^{1/2} r_*^{-3/2}$$

($\beta = R_T/R_S$)

S. Gezari, *Annu. Rev. Astron. Astrophys.* 2021. 59:21–58

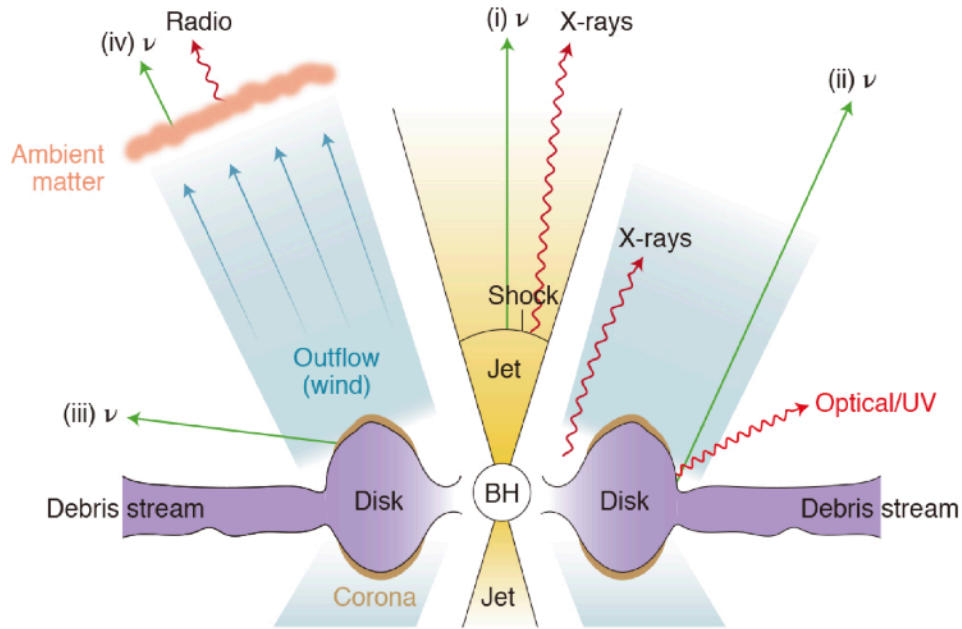
Van Velzen et al. (2021)

M. Toscani, G. Lodato, D.J. Price, D. Liptai, *MNRAS* (arXiv: 2111.05145)

Batista et al., *Front. Astron. Space Sci.* 6 (2019), 23

TDEs: high-energy neutrinos

▲ Soft X-ray TDEs *Hayasaki, Nat. Astar. (2021)*



▶ Optical/UV TDEs

(i) Relativistic jets

(Wang+16, Senno, Murase & Meszaros 17, Murase+ 20, Lunardini & Winter 17, 21)

(ii) Disk (RIAF - MAD)

(Hayasaki & Yamazaki 19, Murase+ 20)

(iii) Disk corona

(Murase+ 20)

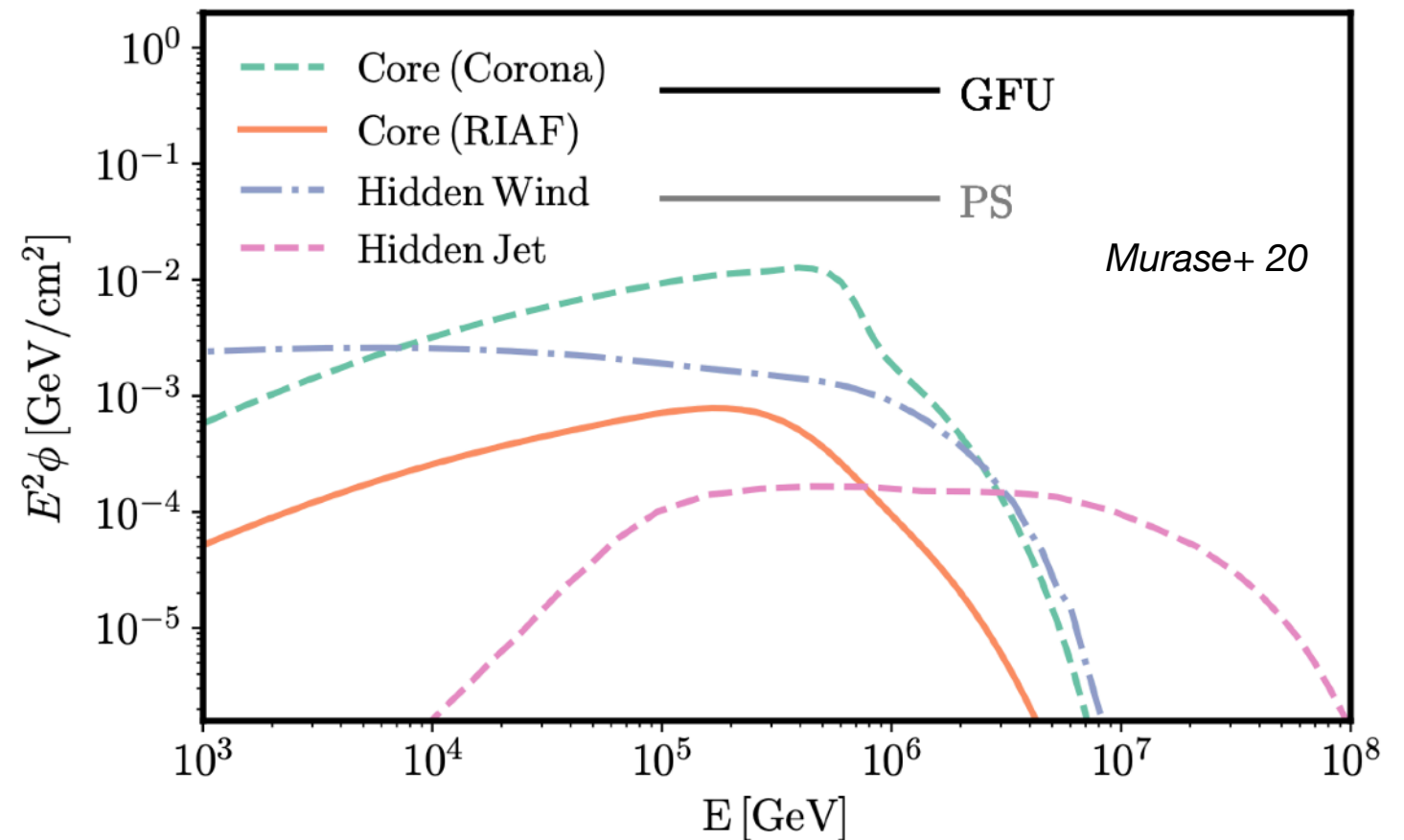
(iv) Wind/Outflow

(Murase+ 20, Wu+ 22, Winter & Lunardini 23)

TeV-PeV neutrinos

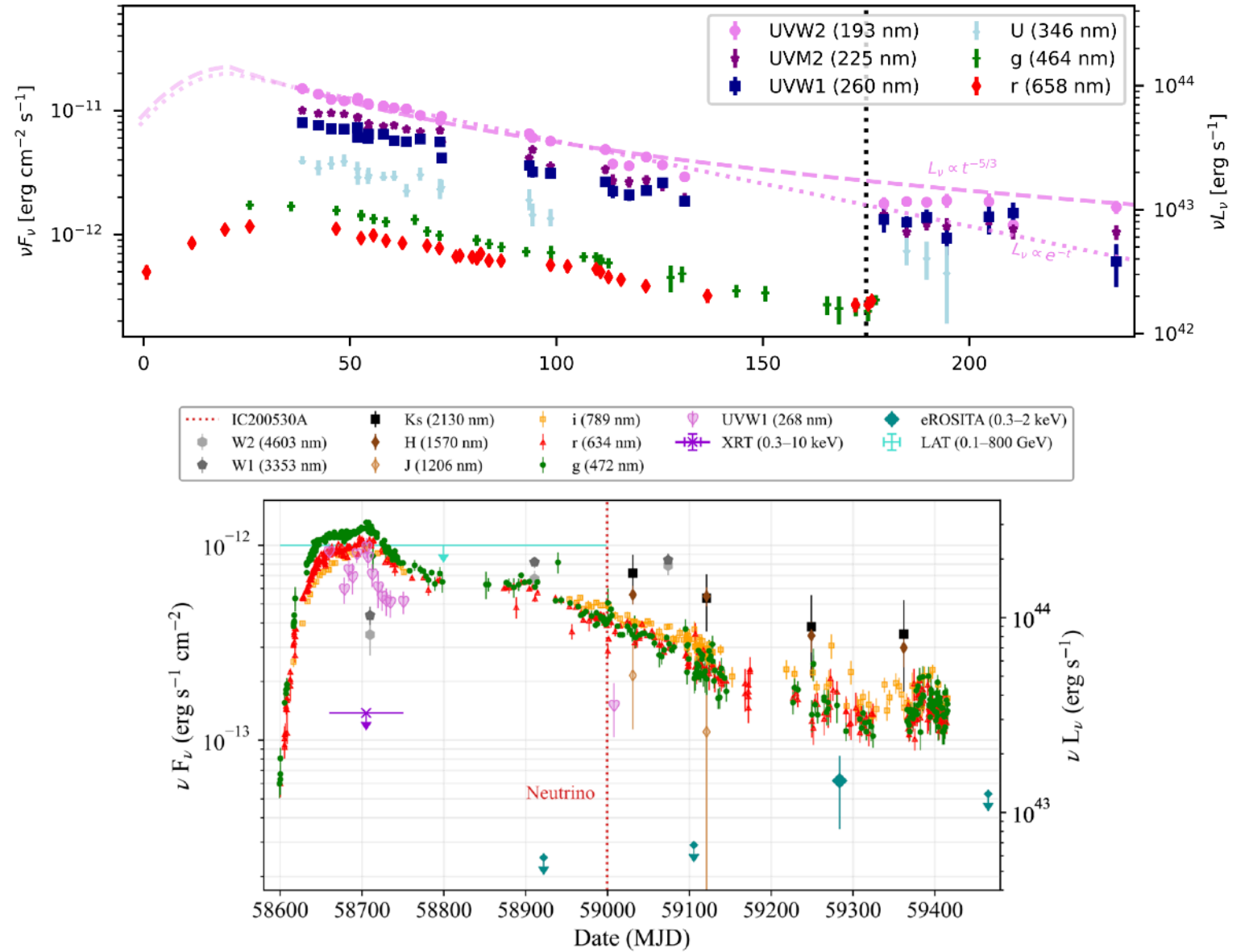
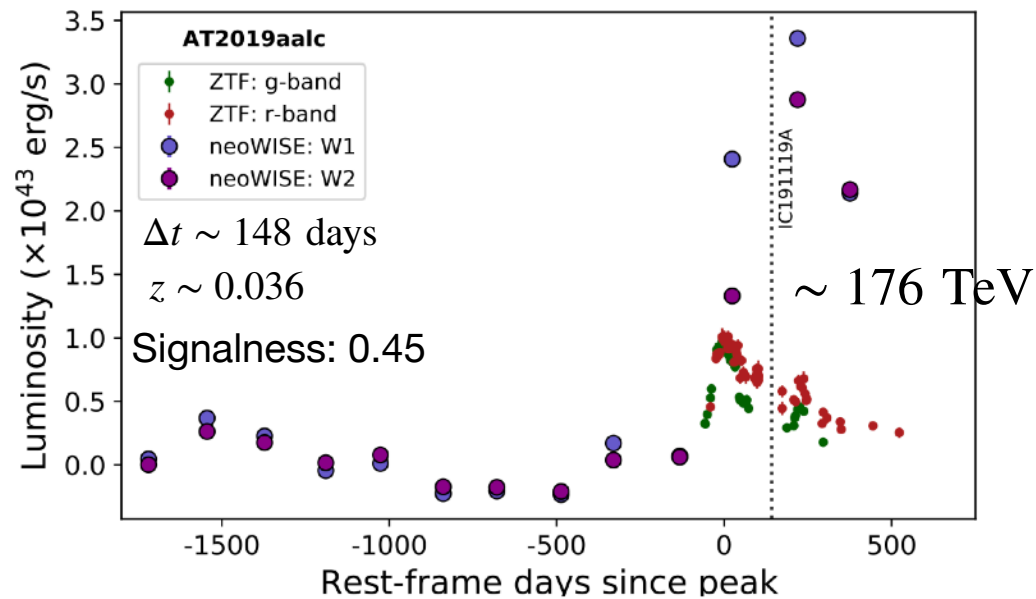
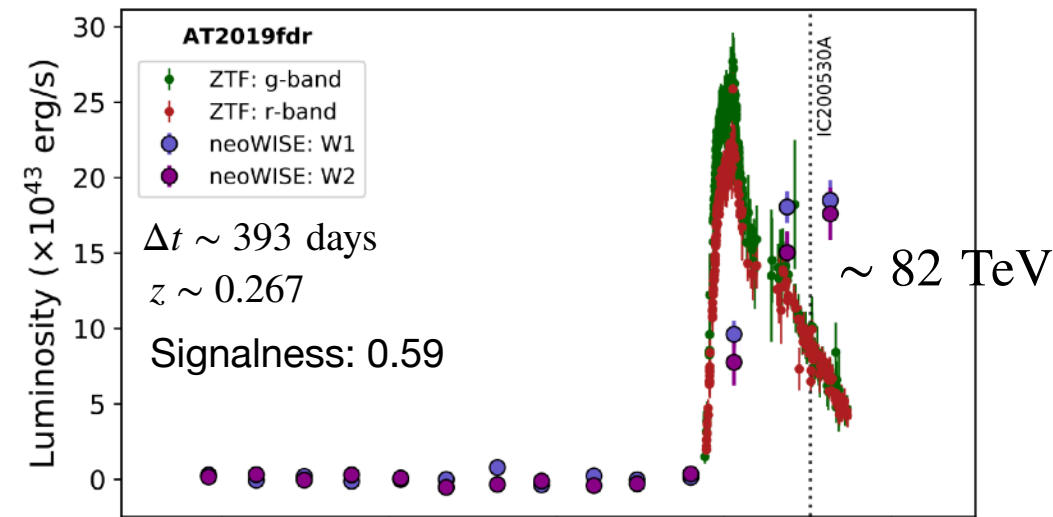
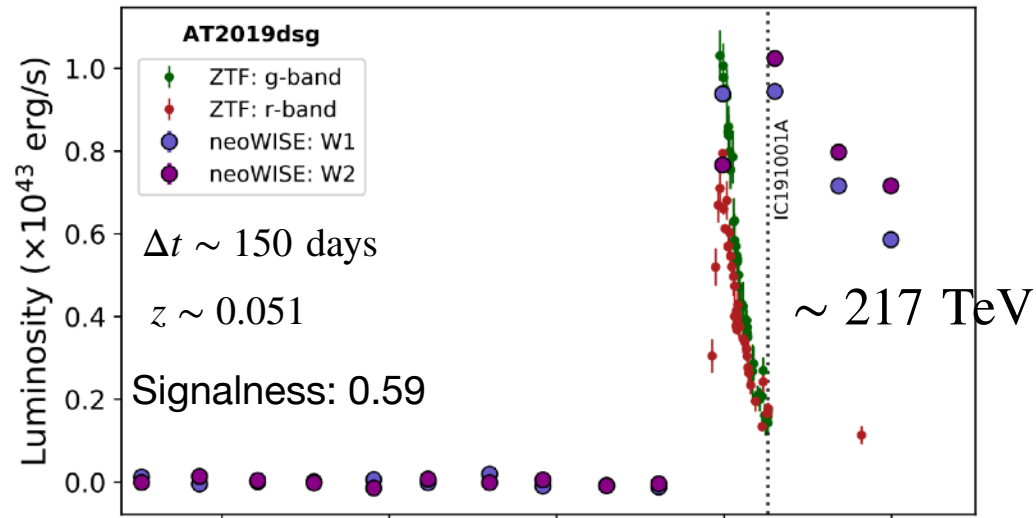
Various acceleration sites

Detectable at IceCube



Murase+ 20

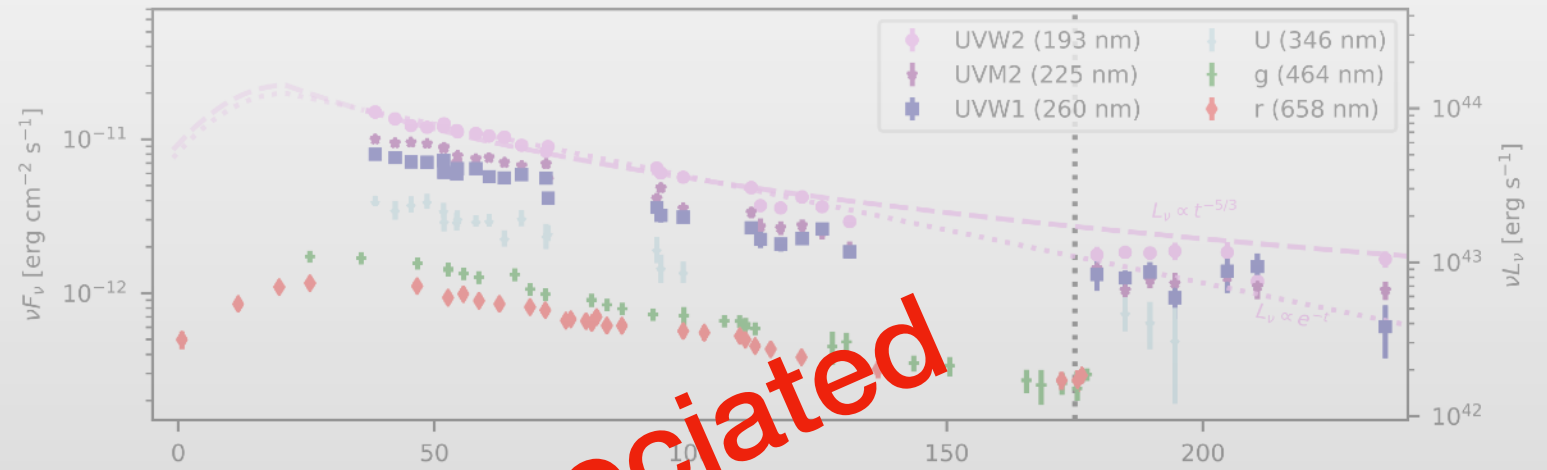
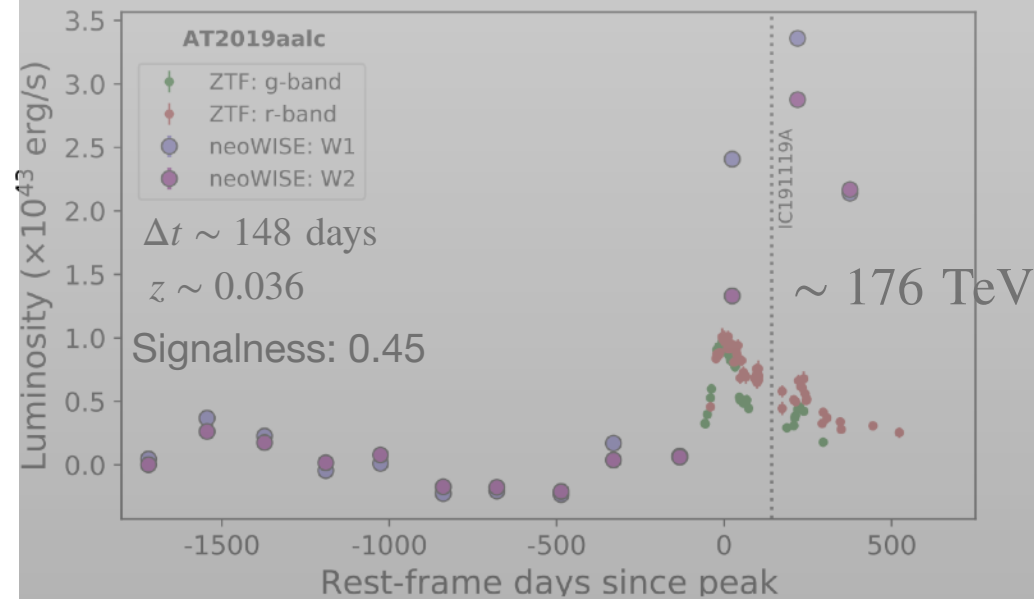
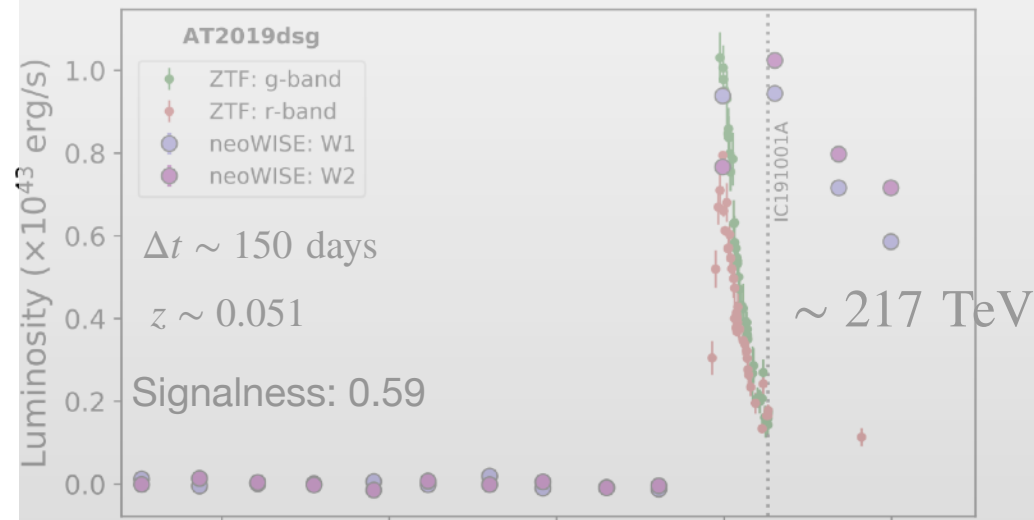
Observational aspects - ν Associations



Flare name	Neutrino name	z	Δt days	Δd deg	X-ray temp. keV	Fermi limit erg/s/cm ²	M_{BH} M_\odot	f_{Edd}
AT2019dsg	IC191001A	0.051	150	1.3	0.07 ± 0.01	$10^{-11.7}$	$10^{6.7}$	3.1
AT2019fdr	IC200530A	0.267	393	1.7	0.06 ± 0.03	$10^{-12.0}$	$10^{7.1}$	0.5
AT2019aalc	IC191119A	0.036	148	1.9	0.17 ± 0.01	$10^{-11.2}$	$10^{7.2}$	0.6

Stein et al. (2021), Nat. Astro.
 Reusch et al. (2022), PRL
 Van Velzen et al. (2021)

Observational aspects - ν Associations

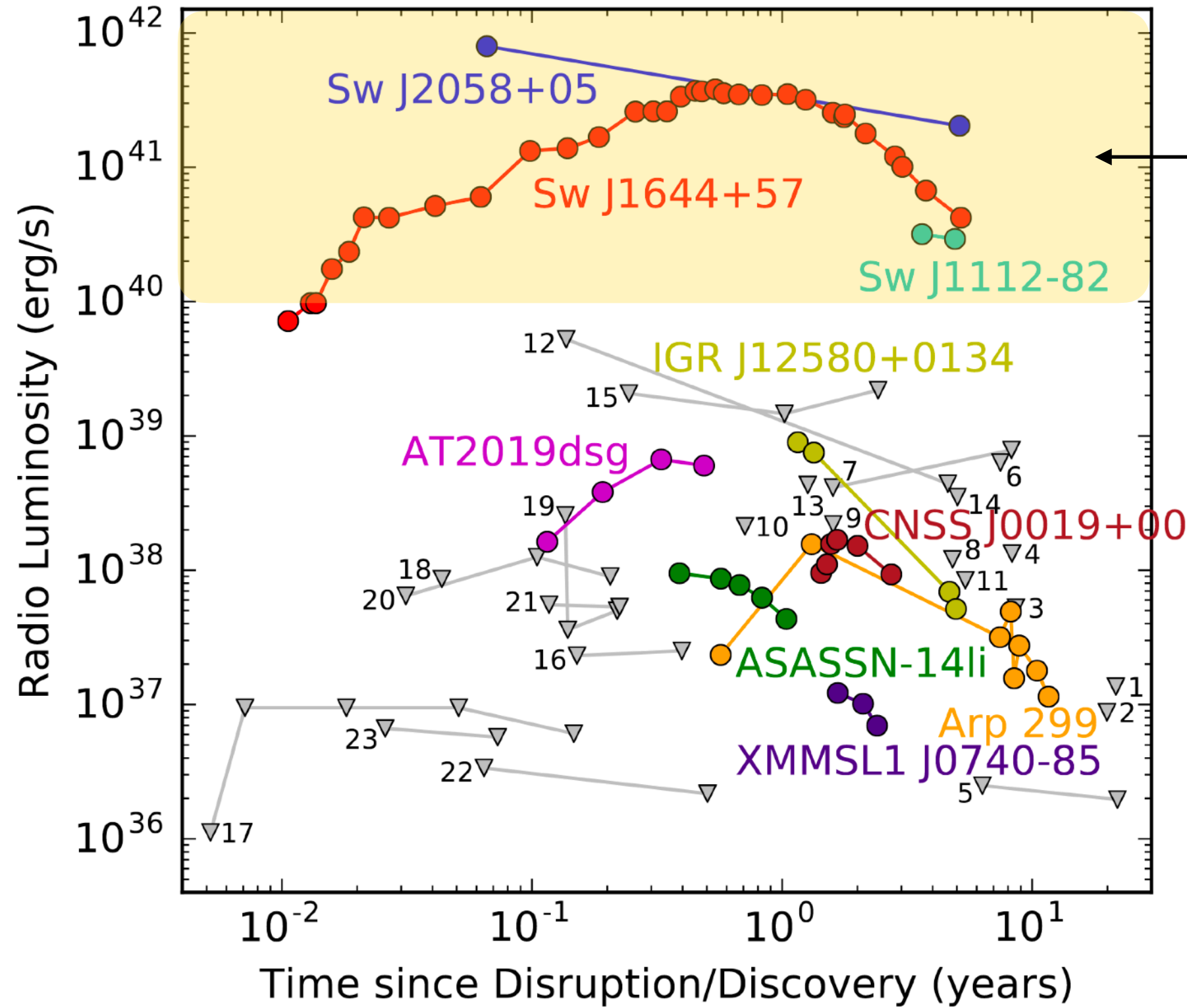


Delayed arrival of associated neutrino events

Flare name	Neutrino name	z	Δt days	Δd deg	X-ray temp. keV	Fermi limit erg/s/cm ²	M_{BH} M_{\odot}	f_{Edd}
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Van Velzen et al. (2021),

Jetted TDEs



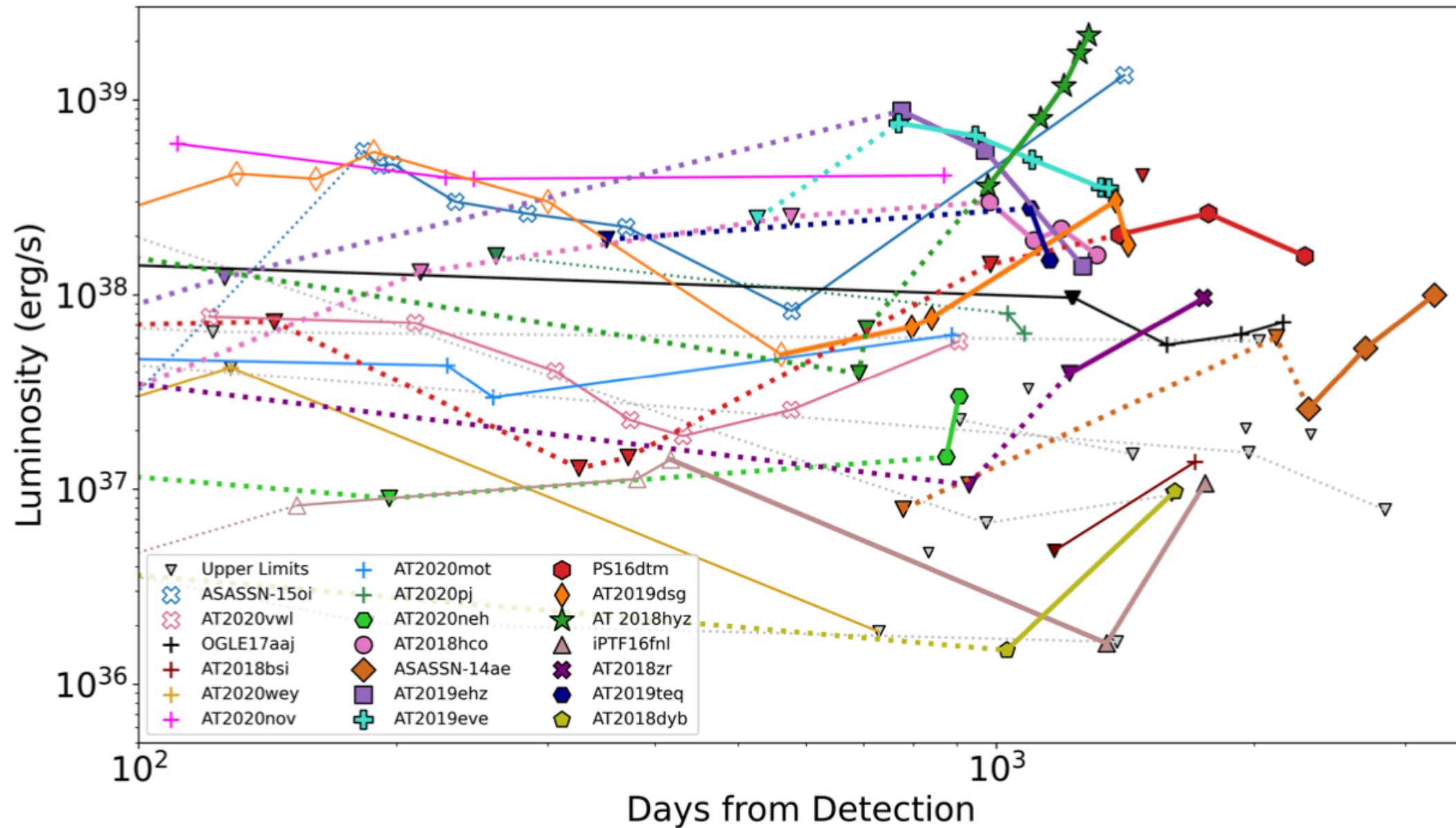
Radio loud
Jetted TDEs

~ 1% of TDEs
can be jetted

Non-relativistic
outflows

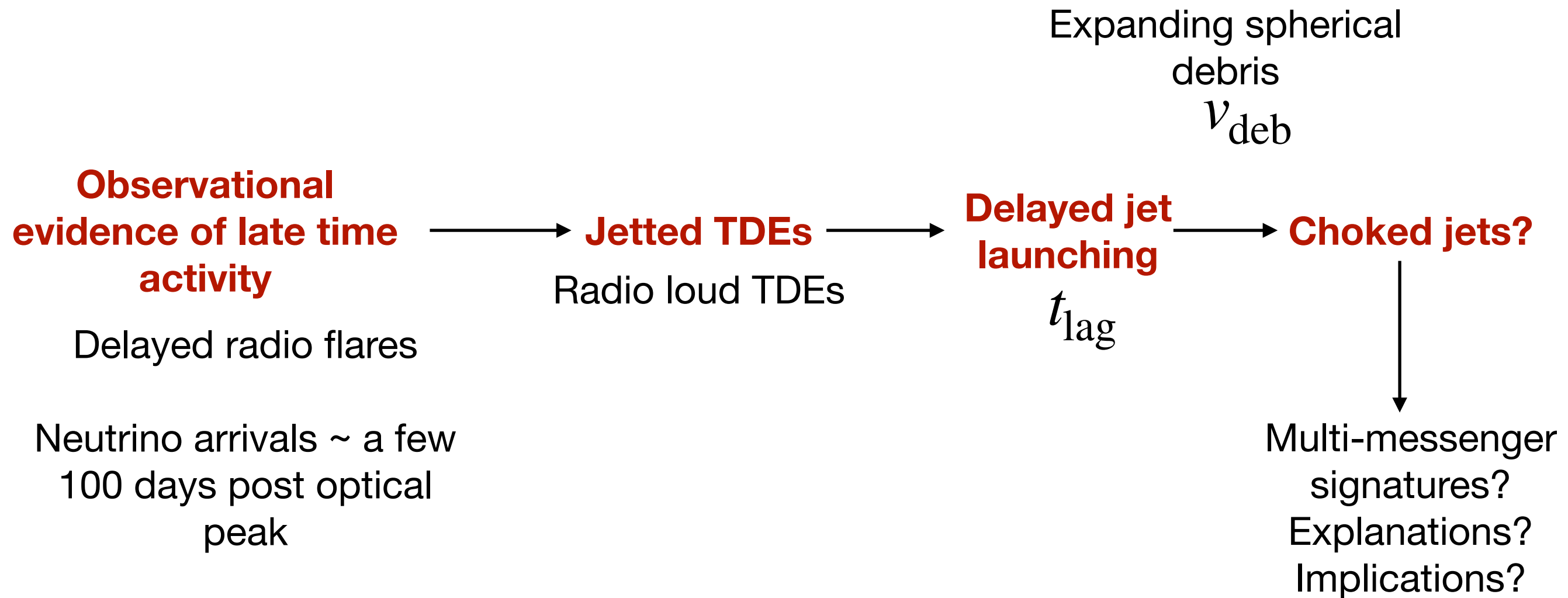
Observational aspects - EM radio flares

Delayed radio flares

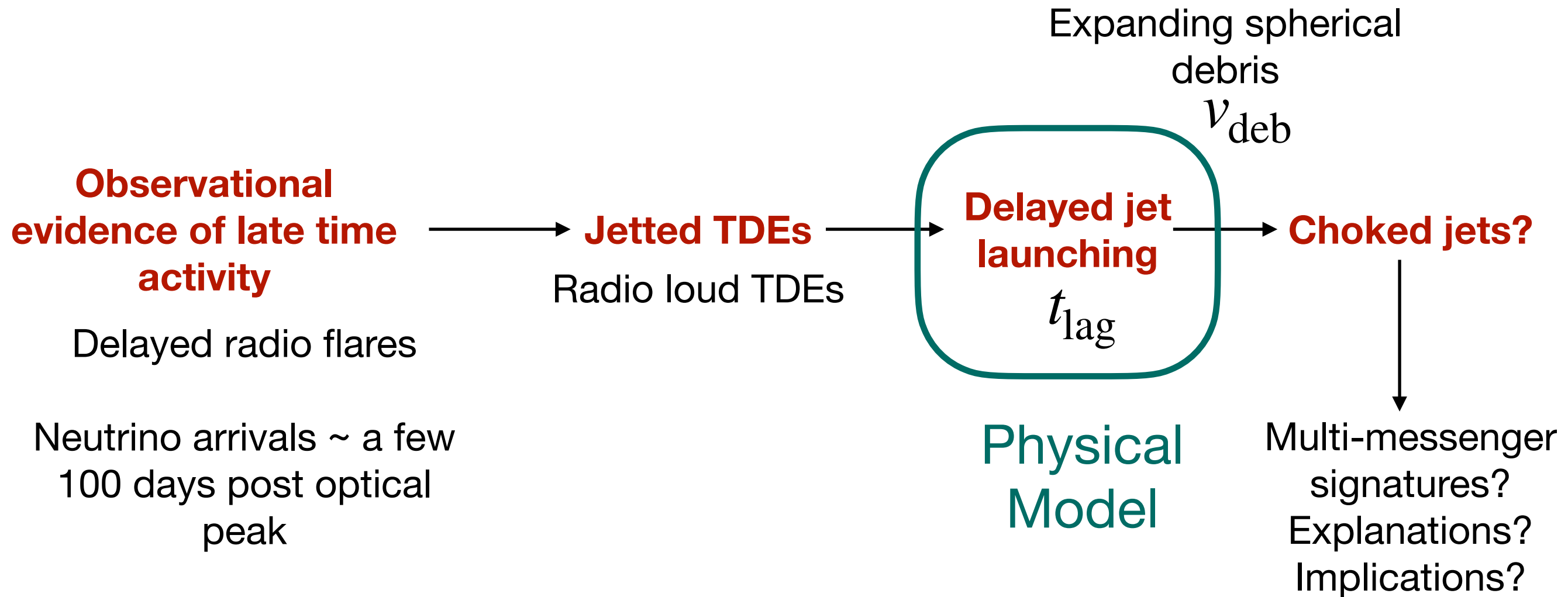


Delayed radio flares: Evidence for late time-activity

Motivations



Motivations



Physical Model: Expanding debris

Convention:

T : Time since TDE

$$T = t + t_{\text{lag}}$$

t_{lag} : Time lag associated with the launching of the jet

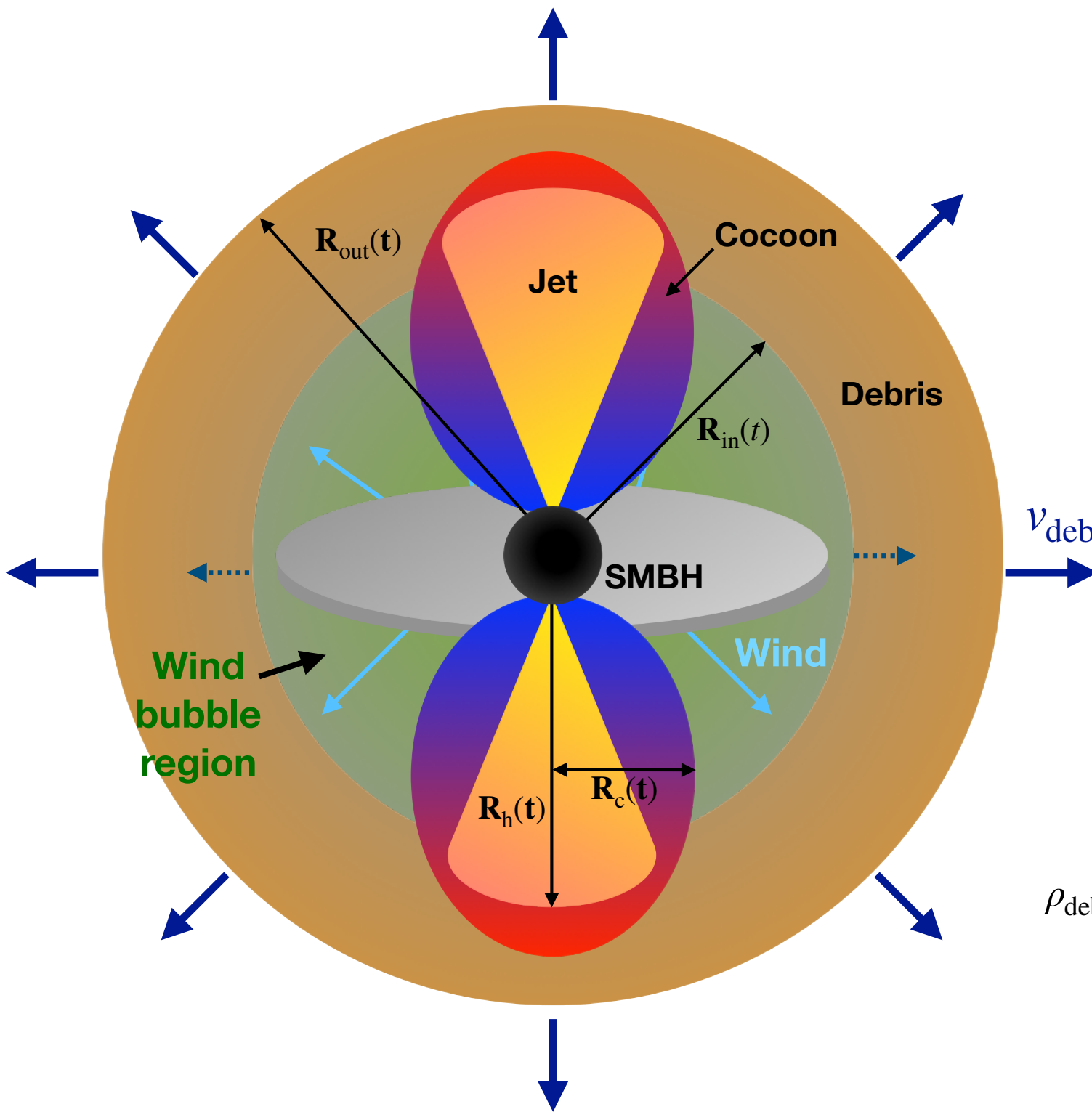
t : Time since the launch of the jet

$$R_{\text{out}}(T) = v_{\text{deb}} T$$

$$R_{\text{in}}(T) = \begin{cases} R_{\text{circ}}, & 0 < T \leq t_{\text{fb}} \\ R_{\text{circ}}(T/t_{\text{fb}}), & T > t_{\text{fb}} \end{cases}$$

$$\rho_{\text{deb}}(t, r) = \mathcal{N} \frac{M_{\text{deb}}}{4\pi R_{\text{out}}^3} \begin{cases} \left(r/R_{\text{out}} \right)^{-2}, & r \geq R_{\text{fb}} \\ \left(R_{\text{fb}}/R_{\text{out}} \right)^{-2} \left(r/R_{\text{fb}} \right)^{-\delta}, & r < R_{\text{fb}} \end{cases} \quad \delta = 1$$

$$R_{\text{fb}}(T) = \begin{cases} R_{\text{in}}(T = 0), & T < t_{\text{fb}} \\ R_{\text{in}}(T = 0) + v_{\text{deb}}(T - t_{\text{fb}}), & T \geq t_{\text{fb}} \end{cases}$$



$$t_{\text{coc}} < t < t_{\text{br}} \text{ or } t_{\text{coc}} < t < t_{\text{fin}}$$

Static and contracting envelopes have been considered

Physical Model: Expanding debris

Convention:

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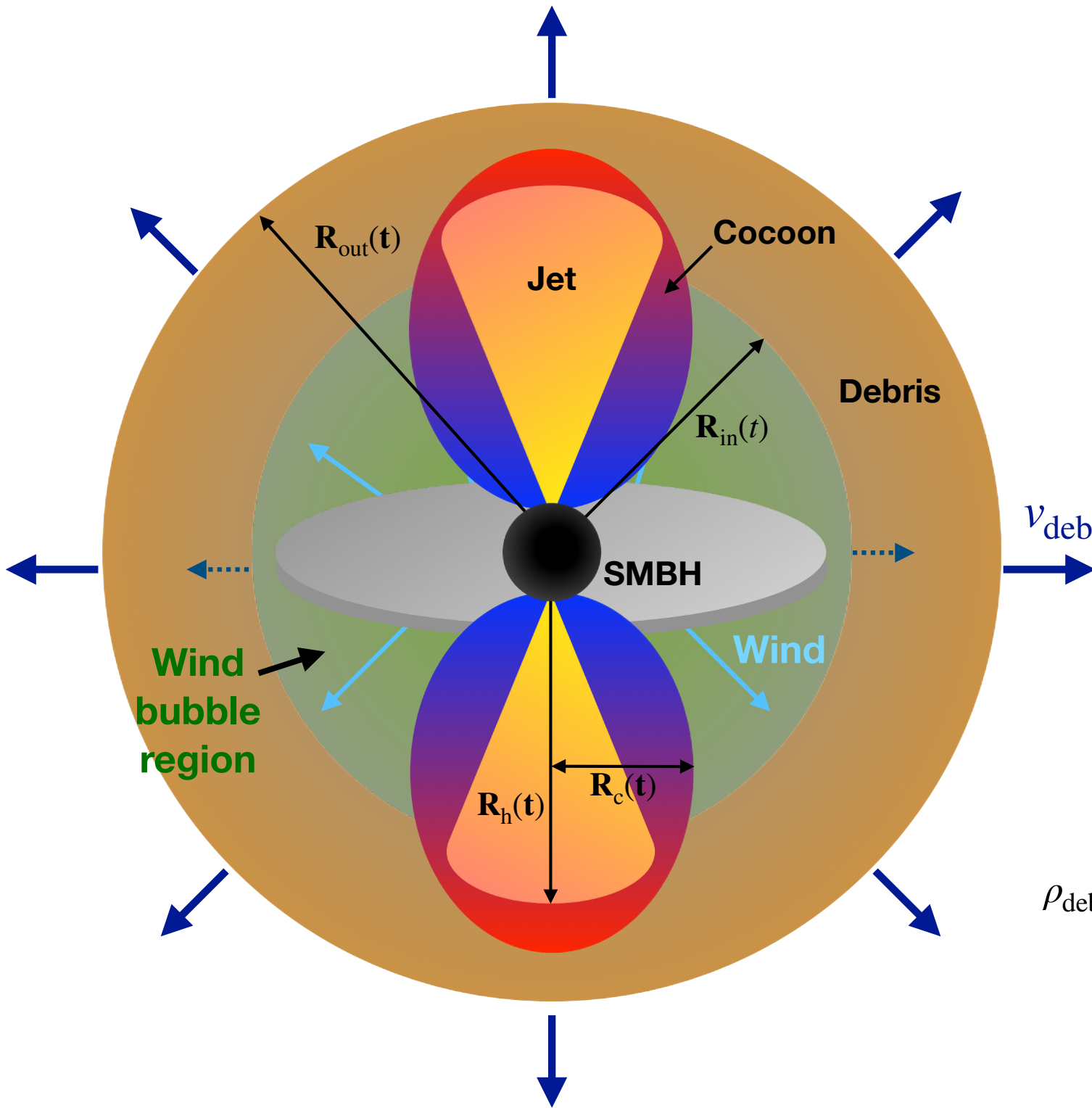
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$$\rho_{\text{deb}}(t, r) = \mathcal{N} \frac{M_{\text{deb}}}{4\pi R_{\text{out}}^3} \begin{cases} \left(r/R_{\text{out}} \right)^{-2}, & r \geq R_{\text{fb}} \\ \left(R_{\text{fb}}/R_{\text{out}} \right)^{-2} \left(r/R_{\text{fb}} \right)^{-\delta}, & r < R_{\text{fb}} \end{cases} \quad \delta = 1$$

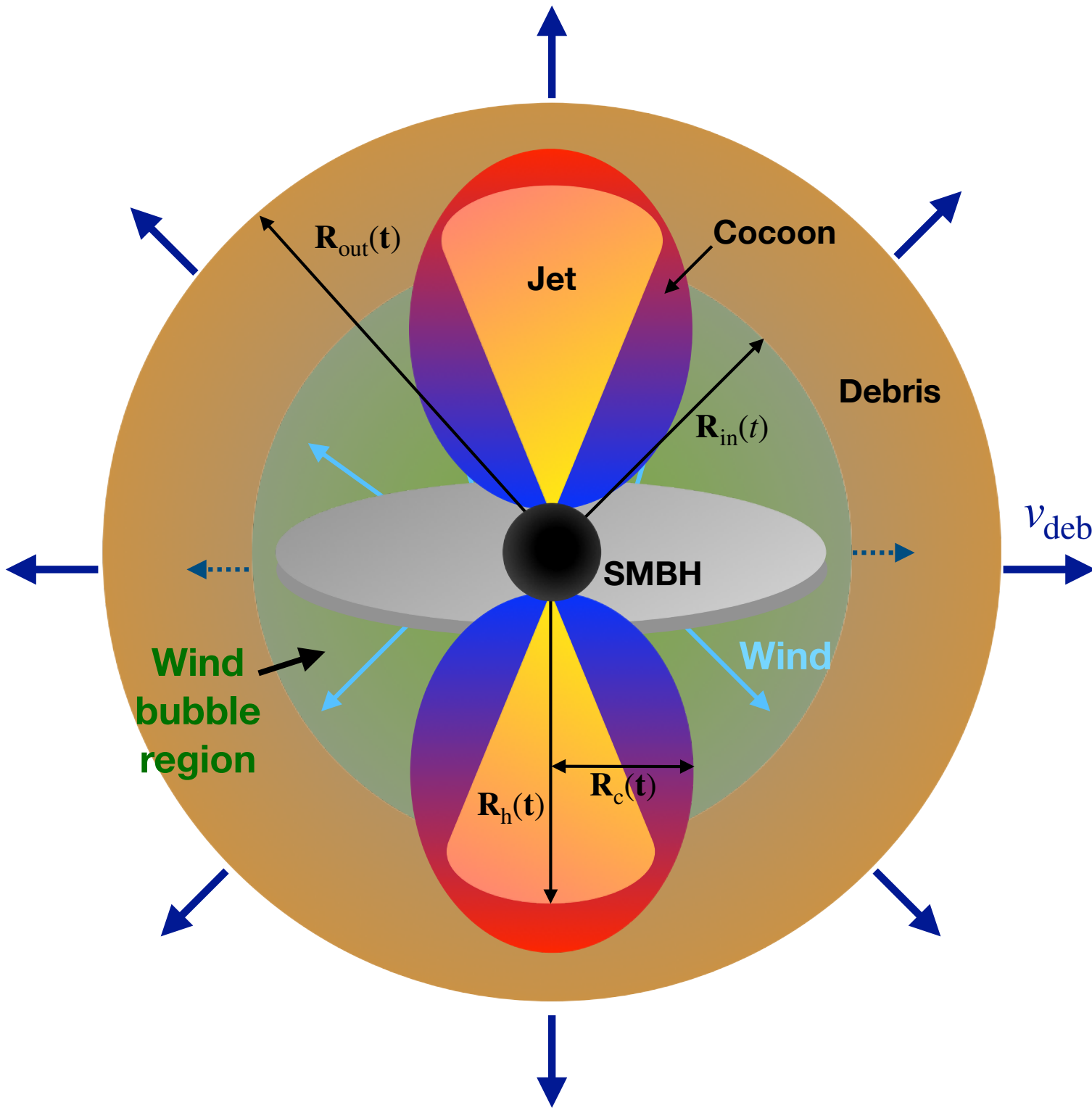
$$R_{\text{fb}}(T) = \begin{cases} R_{\text{in}}(T = 0), & T < t_{\text{fb}} \\ R_{\text{in}}(T = 0) + v_{\text{deb}}(T - t_{\text{fb}}), & T \geq t_{\text{fb}} \end{cases}$$



$$t_{\text{coc}} < t < t_{\text{br}} \text{ or } t_{\text{coc}} < t < t_{\text{fin}}$$

Static and contracting envelopes have been considered

Physical Model: Jet propagation in expanding debris



$$R_h(t = 0) = R_s = 2GM_{\text{BH}}/c^2$$

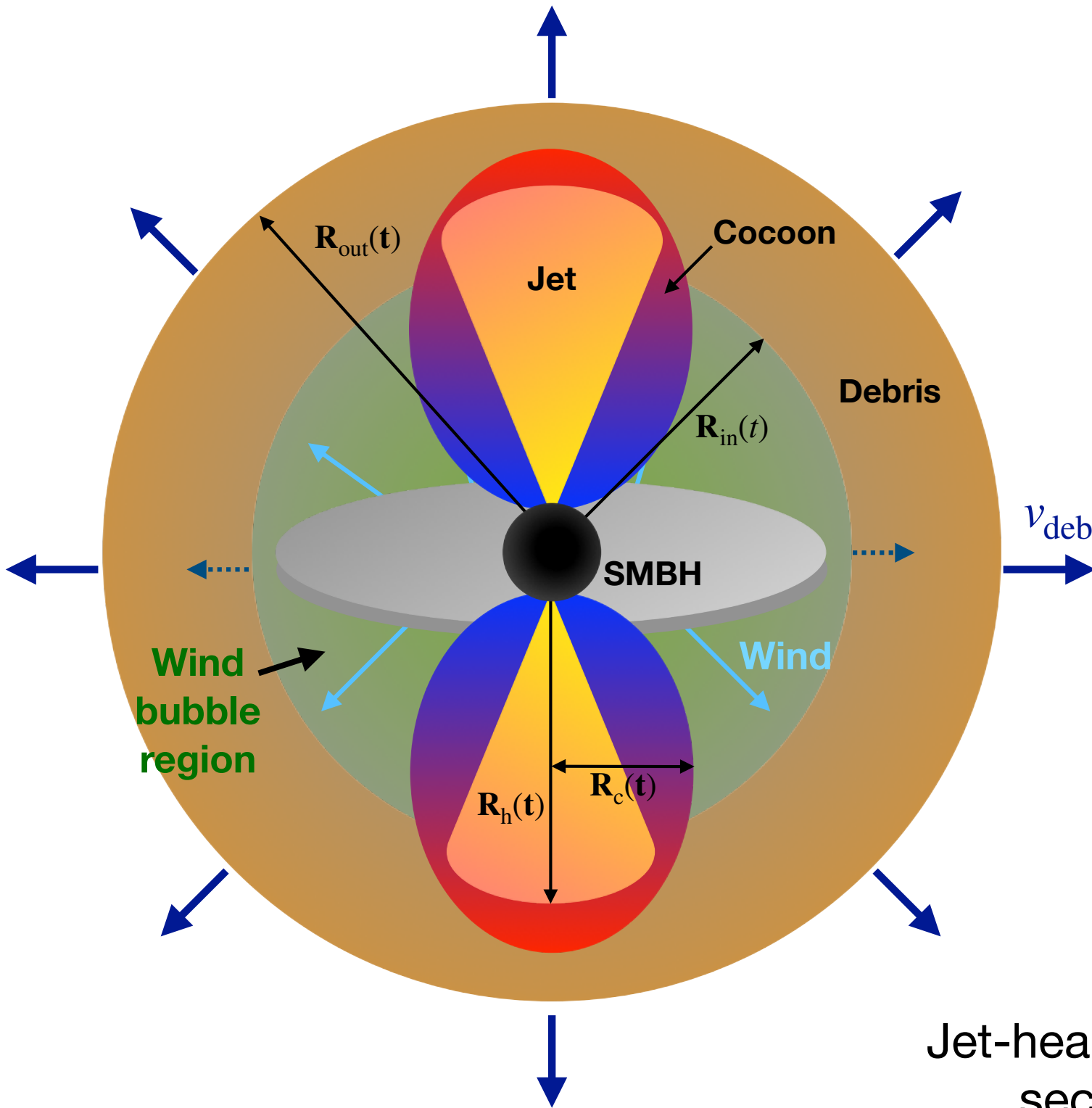
$$\dot{R}_h = c\beta_h$$

$$\beta_h = \frac{\beta_j - \beta_a}{1 + \tilde{L}_c^{-1/2}} + \beta_a$$

Ambient medium:
surrounding medium of
jet-head

$$t_{\text{coc}} < t < t_{\text{br}} \text{ Or } t_{\text{coc}} < t < t_{\text{fin}}$$

Physical Model: Jet propagation in expanding debris



$$t_{\text{coc}} < t < t_{\text{br}} \text{ Or } t_{\text{coc}} < t < t_{\text{fin}}$$

$$R_h(t_{\text{dur}} = 0) = R_s = 2GM_{\text{BH}}/c^2$$

$$\dot{R}_h = c\beta_h$$

$$\beta_h = \frac{\beta_j - \beta_a}{1 + \tilde{L}_c^{-1/2}} + \beta_a$$

Ratio of energy density between jet and ambient medium

$$\tilde{L}_c = N_s^2 \tilde{L}$$

Calibration factor

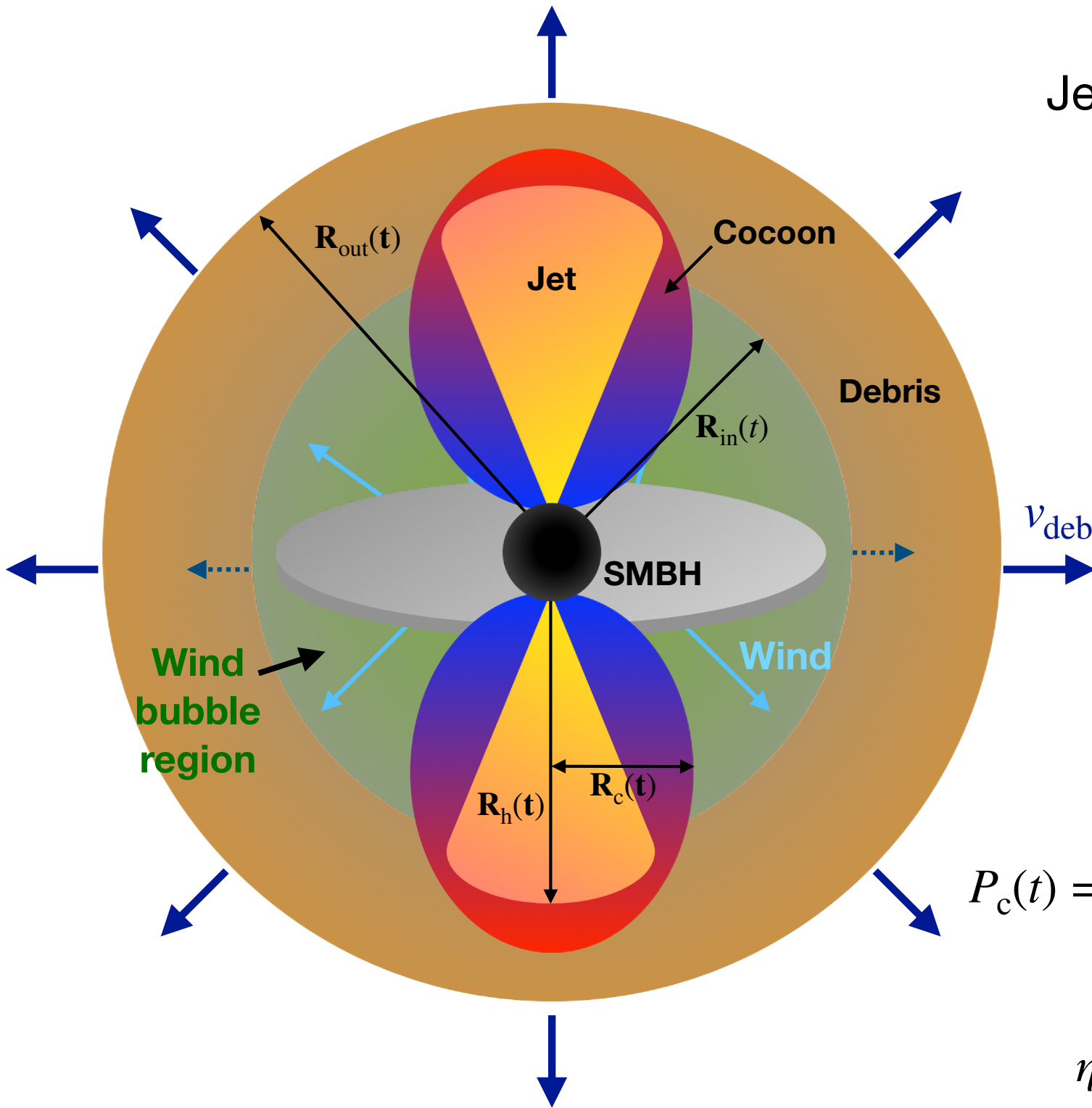
$$\tilde{L} = \frac{L_j}{\Sigma_j(t)\rho_a(t)c^3\Gamma_a^2}$$

Jet-head cross-section

$$\rho_a(t) = \rho_{\text{deb}}(t, r = R_h)$$

Density of ambient medium

Physical Model: Formation of cocoon and interaction



Jet head reaches the inner radius of the debris

$$R_c(t = t_{\text{coc}}) = R_j(t = t_{\text{coc}}) = R_h(t = t_{\text{coc}})\theta_0$$

$$\dot{R}_c = c\beta_c$$

$$\beta_c \approx \frac{1}{c} \sqrt{\frac{P_c}{\rho_a(t)}} + \frac{R_c(t)}{R_{\text{out}}(t)} \frac{v_w}{c}$$

Cocoon pressure

$$P_c(t) = \frac{E_c}{3V_c} = \frac{\eta}{4\pi R_c(t)^2 R_h(t)} \int_{t_c}^t d\tilde{t} L_j(\tilde{t}) (1 - \beta_h(\tilde{t}))$$

η : Fraction of jet energy deposited in cocoon

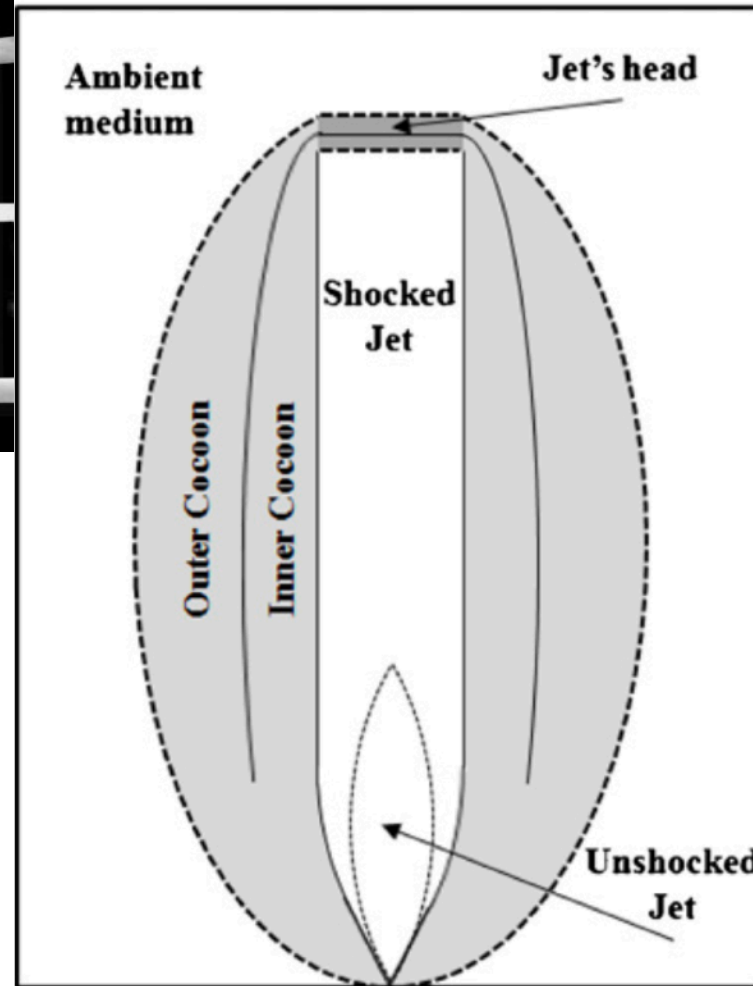
$$t_{\text{coc}} < t < t_{\text{br}} \text{ Or } t_{\text{coc}} < t < t_{\text{fin}}$$

Dynamics: To collimate or not collimate

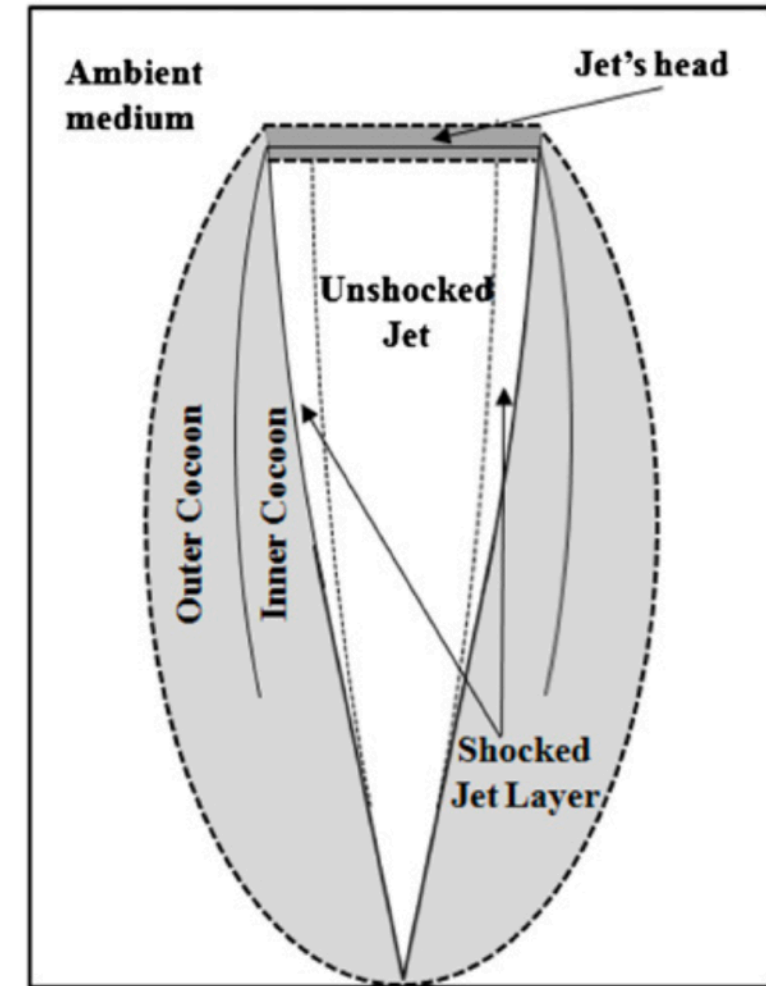


Jet pressure

Cocoon pressure
(Lateral)



Collimated



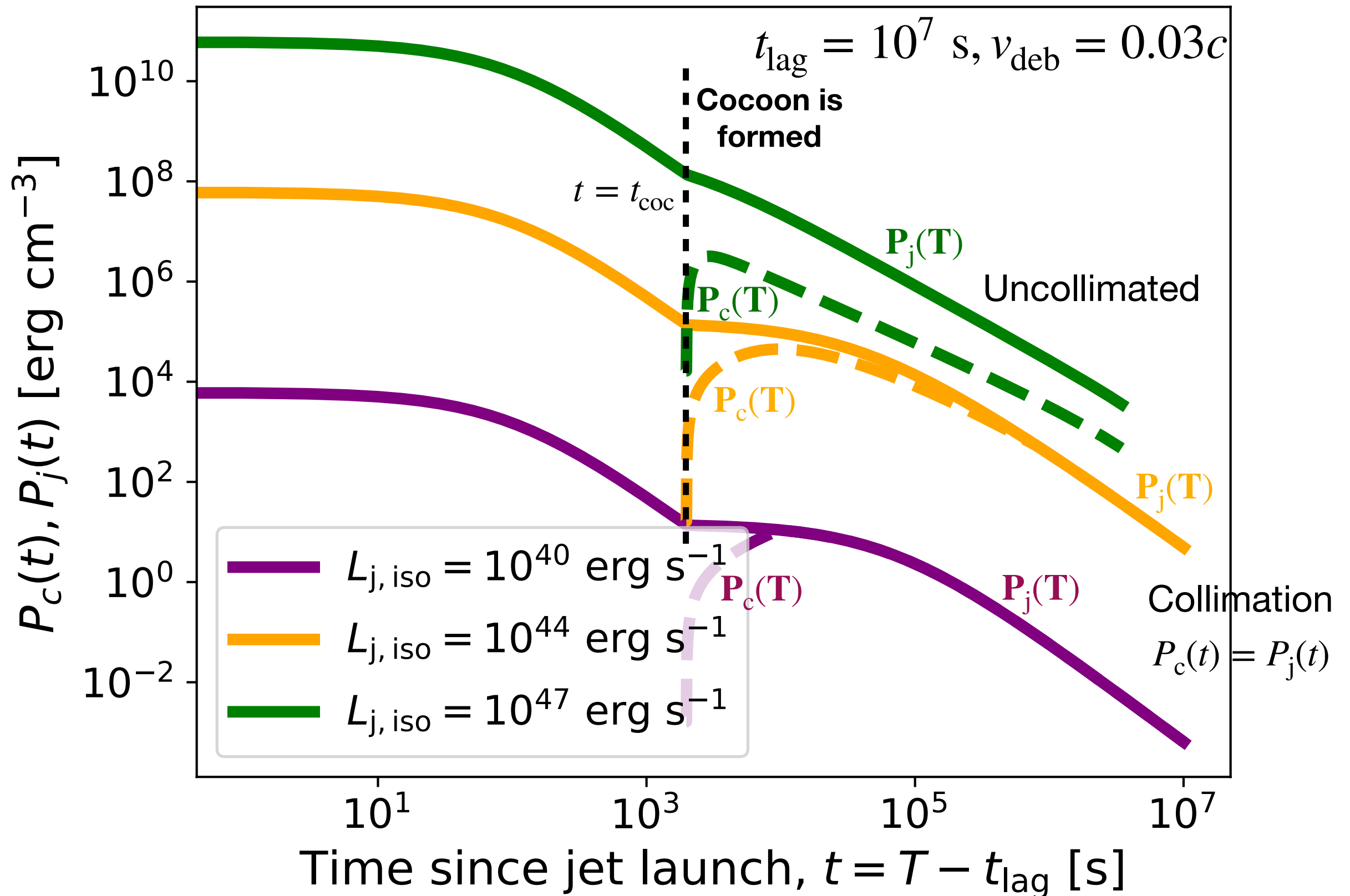
Uncollimated

Collimation criteria:

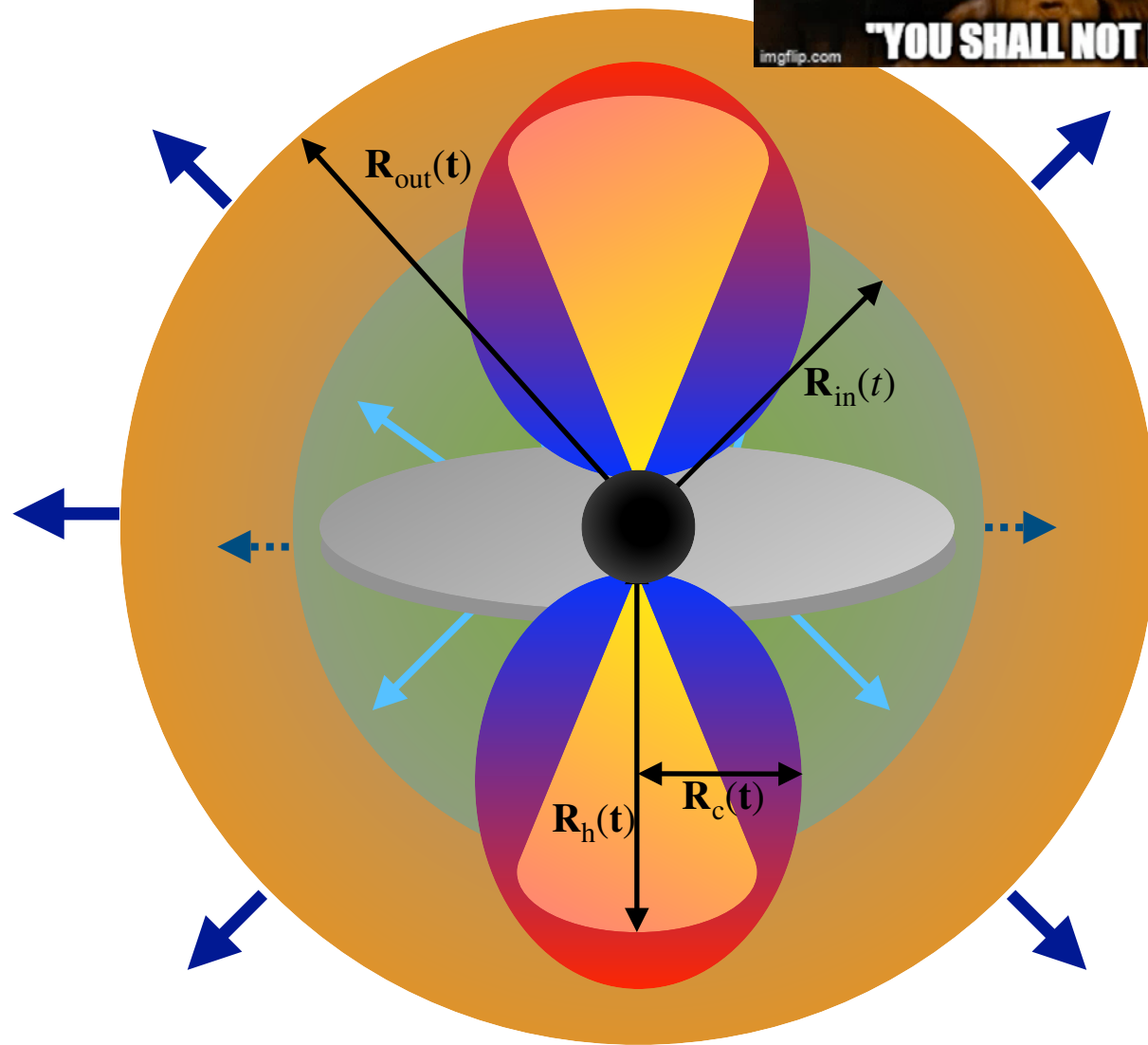
$$P_c(t) = P_j(t)$$

Bromberg et al (2011)

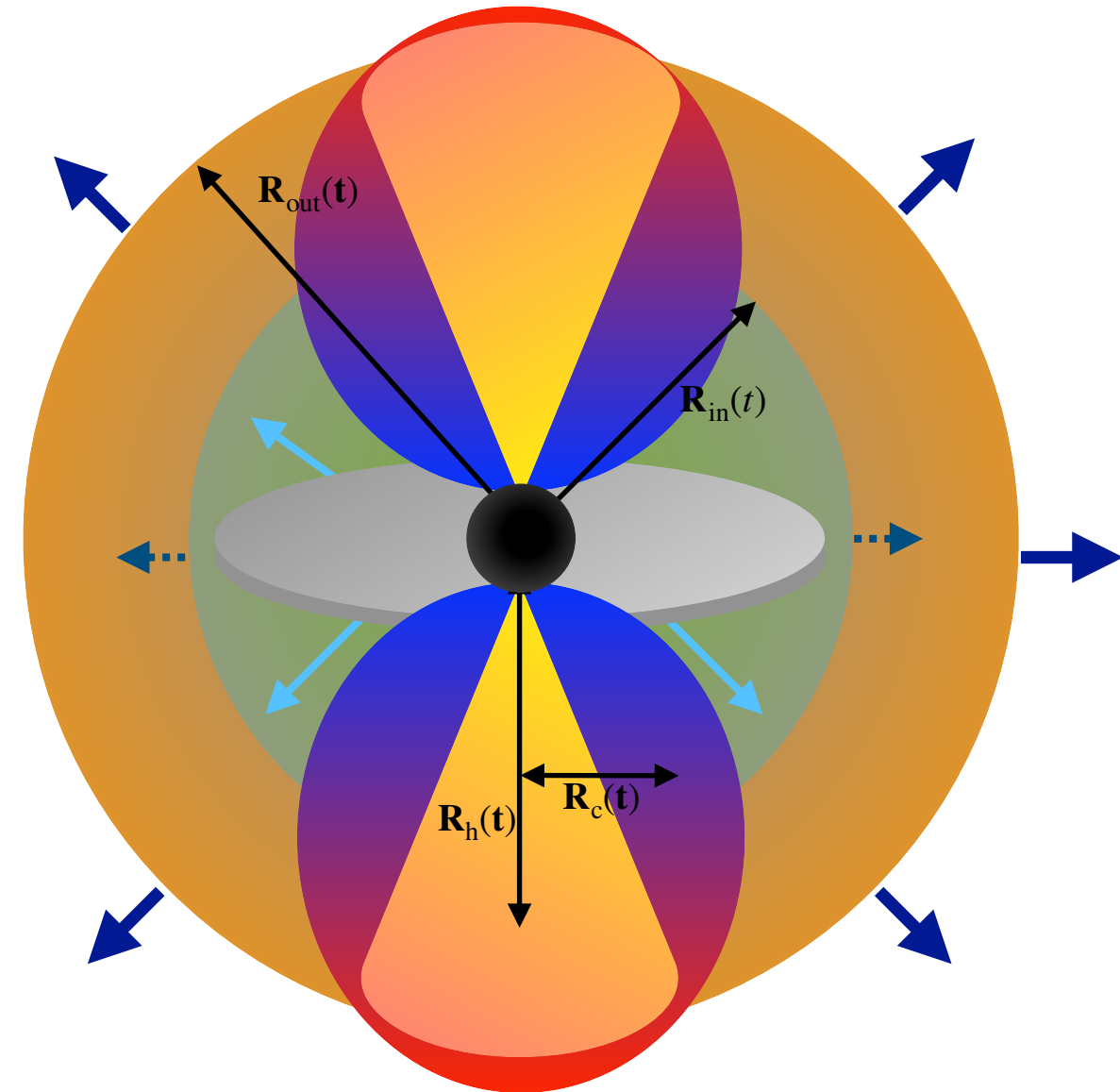
Dynamics: To collimate or not collimate



Dynamics: To choke or not to choke



Choked

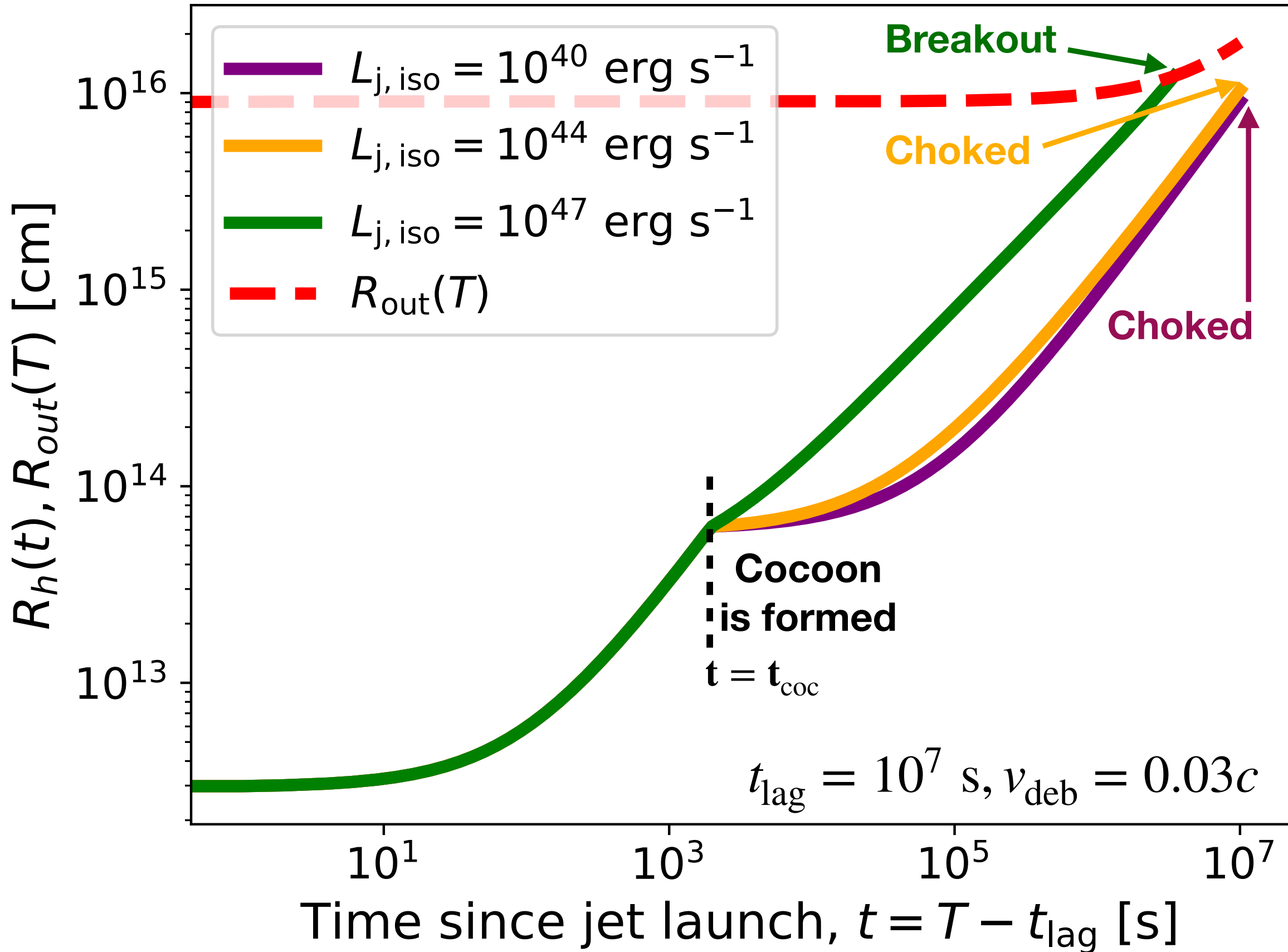


Breakout

Choking criteria:

$$R_h(t_{\text{dur}}) \leq R_{\text{out}}(T)$$

Dynamics: To choke or not to choke



Dynamics: Analytical estimate for choking

Choking criteria

$$R_h(t_{\text{dur}}) \leq R_{\text{out}}(t_{\text{fin}})$$

$$T = t_{\text{fin}} = t_{\text{dur}} + t_{\text{lag}}$$

Total evolution time

$$R_{\text{out}} \simeq 1.8 \times 10^{16} \text{ cm} \left(\frac{\beta_{\text{deb}}}{0.03} \right) \left(\frac{t_{\text{dur}}}{10^7 \text{ s}} \right) \left(\frac{\chi_{\text{lag}}}{2} \right) \quad \chi_{\text{lag}} = (1 + t_{\text{lag}}/t_{\text{dur}})$$

Assuming uncollimated jets

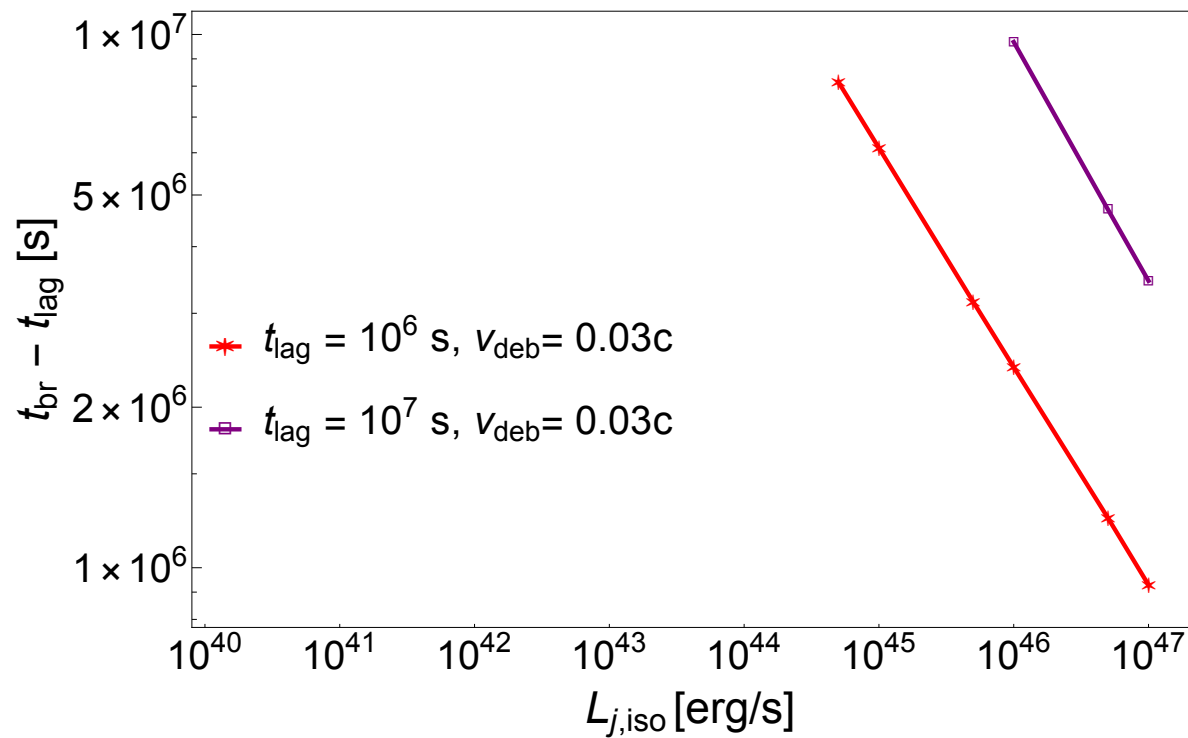
$$R_h \simeq 5.6 \times 10^{15} \text{ cm} \left(\frac{N_s}{0.35} \right)^{5/3} \left(\frac{L_{\text{j,iso}}}{10^{44} \text{ erg/s}} \right)^{1/3} \left(\frac{M_{\text{deb}}}{0.5 M_{\odot}} \right)^{-1/3} \left(\frac{\theta_0}{0.17} \right)^{-2/3} \left(\frac{\beta_{\text{deb}}}{0.03} \right)^{1/3} \left(\frac{t_{\text{dur}}}{10^7 \text{ s}} \right)^{4/3} \left(\frac{\chi_{\text{lag}}}{2} \right)^{1/3}$$

$$L_{\text{j,iso}} \lesssim 3.2 \times 10^{45} \text{ erg/s} \left(\frac{N_s}{0.35} \right)^{-5} \left(\frac{M_{\text{deb}}}{0.5 M_{\odot}} \right) \left(\frac{\theta_0}{0.17} \right)^2 \left(\frac{\beta_{\text{deb}}}{0.03} \right)^2 \left(\frac{t_{\text{dur}}}{10^7 \text{ s}} \right)^{-1} \left(\frac{\chi_{\text{lag}}}{2} \right)^2$$

Fairly good estimates

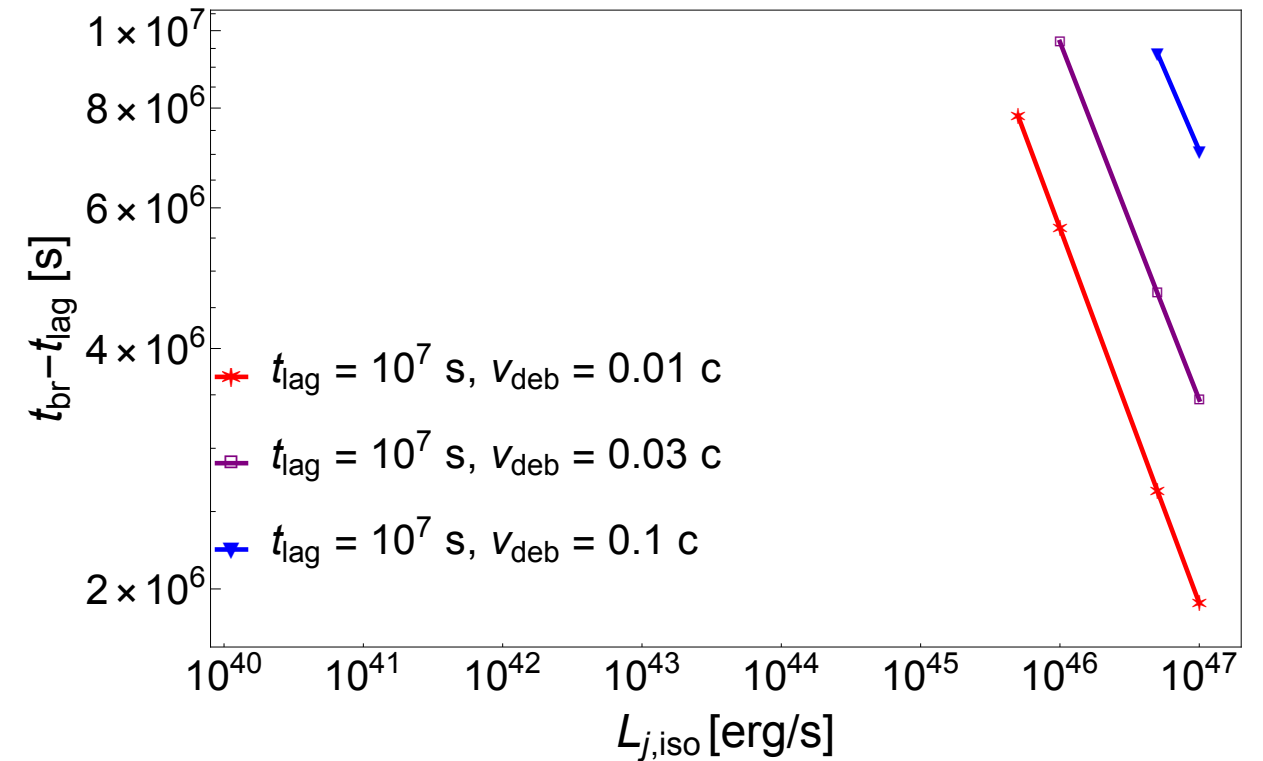
Dynamics: the land of choked jets

$$L_{j,\text{iso}} \lesssim 3.2 \times 10^{45} \text{ erg/s} \left(\frac{N_s}{0.35} \right)^{-5} \left(\frac{M_{\text{deb}}}{0.5 M_{\odot}} \right) \left(\frac{\theta_0}{0.17} \right)^2 \left(\frac{\beta_{\text{deb}}}{0.03} \right)^2 \left(\frac{t_{\text{dur}}}{10^7 \text{ s}} \right)^{-1} \left(\frac{\chi_{\text{lag}}}{2} \right)^2$$



$t_{\text{lag}} : \uparrow$ The debris has more time to expand

Jets require higher luminosity to breakout



$v_{\text{deb}} : \uparrow$

The debris expands with a higher velocity: extends to larger radii

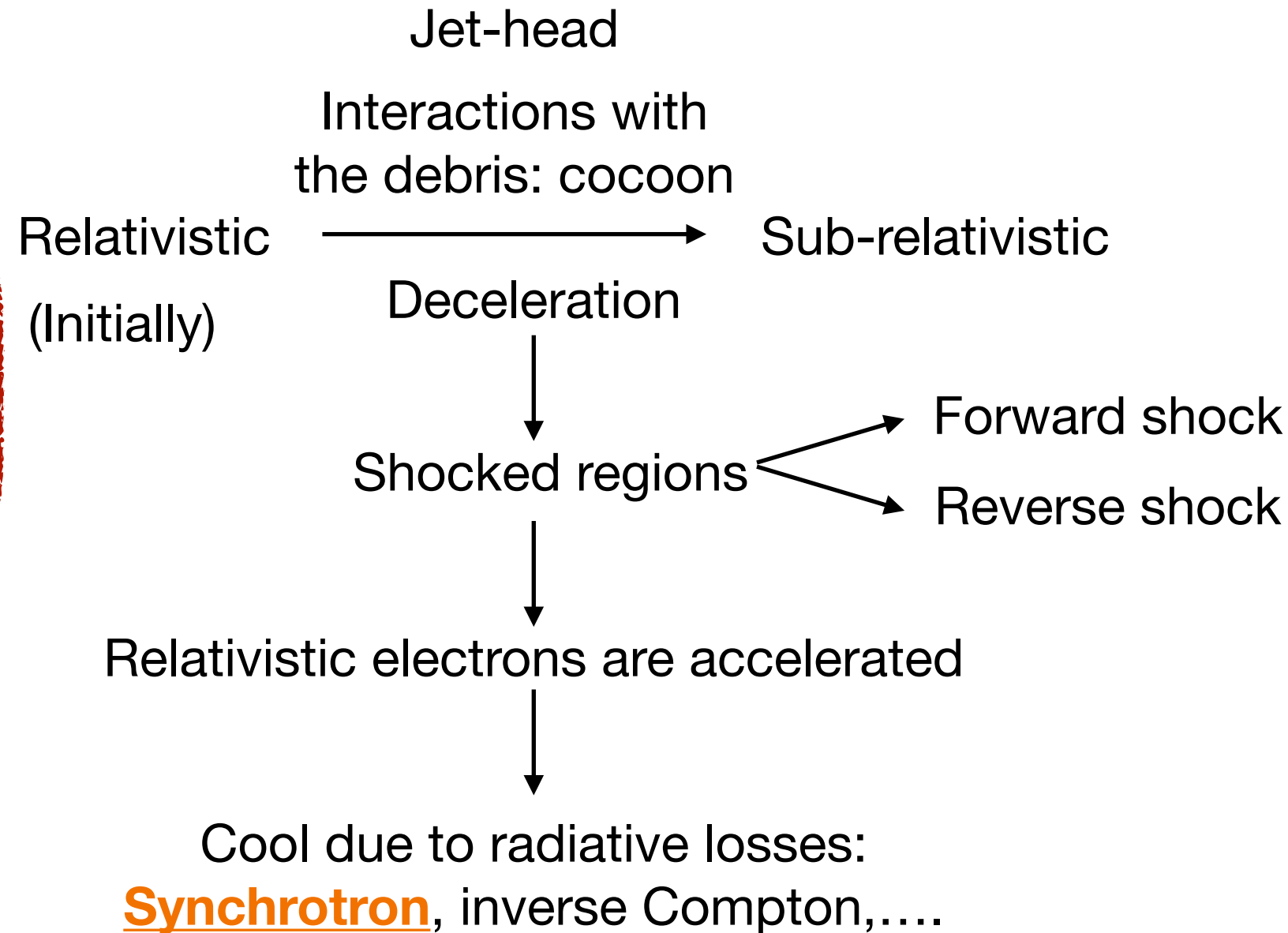
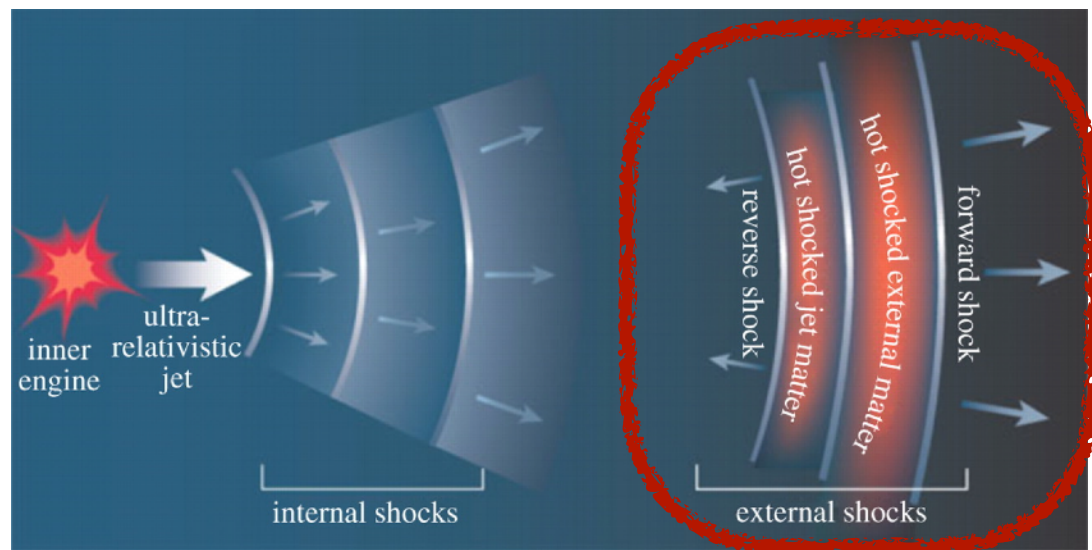
Jets require higher luminosity to breakout

Electromagnetic (EM) and Neutrino Signatures

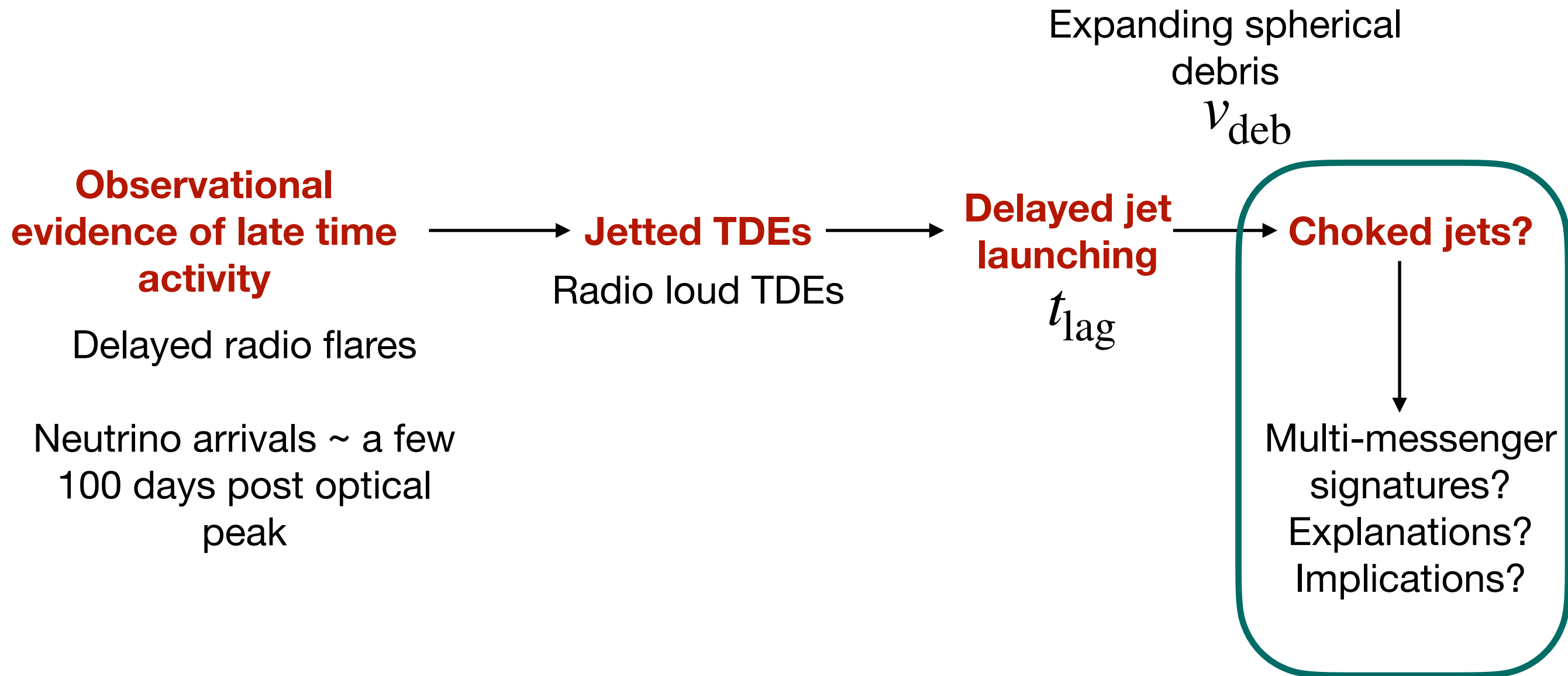
Signatures from delayed choked jets

Particle acceleration within jets:

- UHECRs
- High-energy neutrinos
- EM signatures

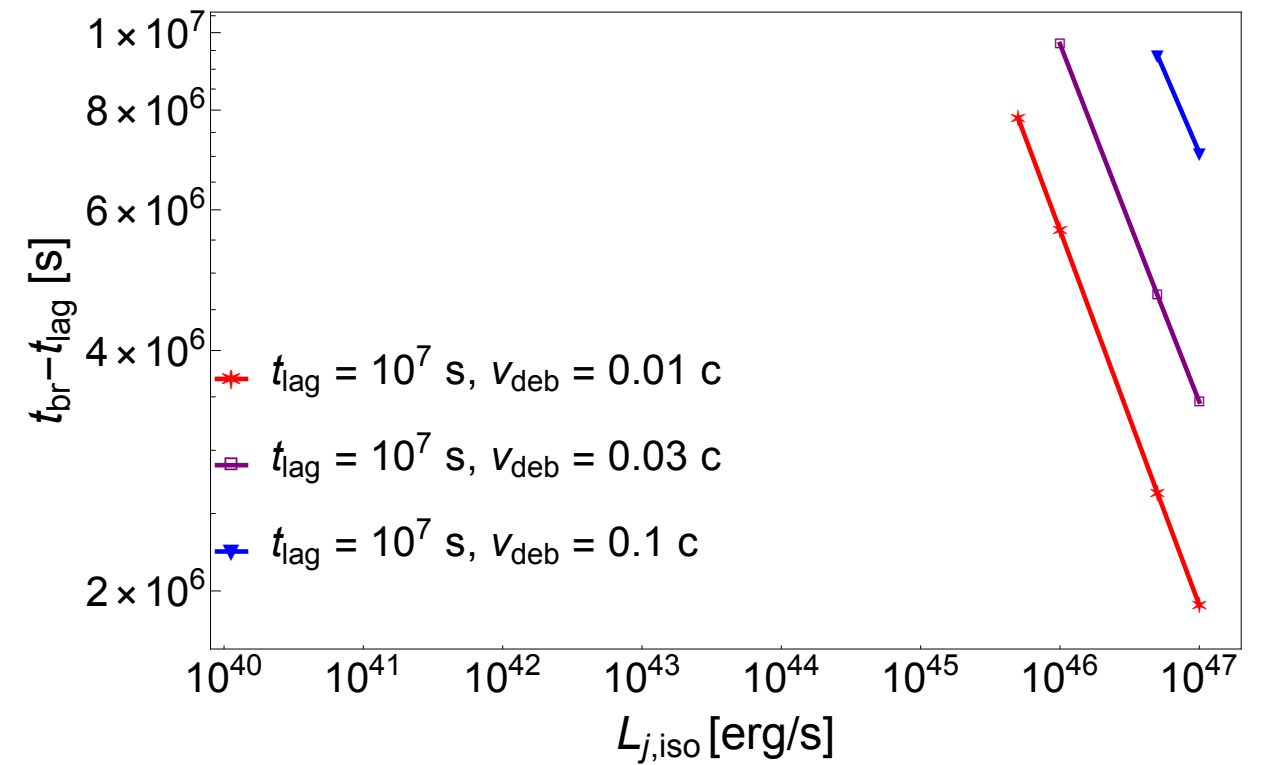
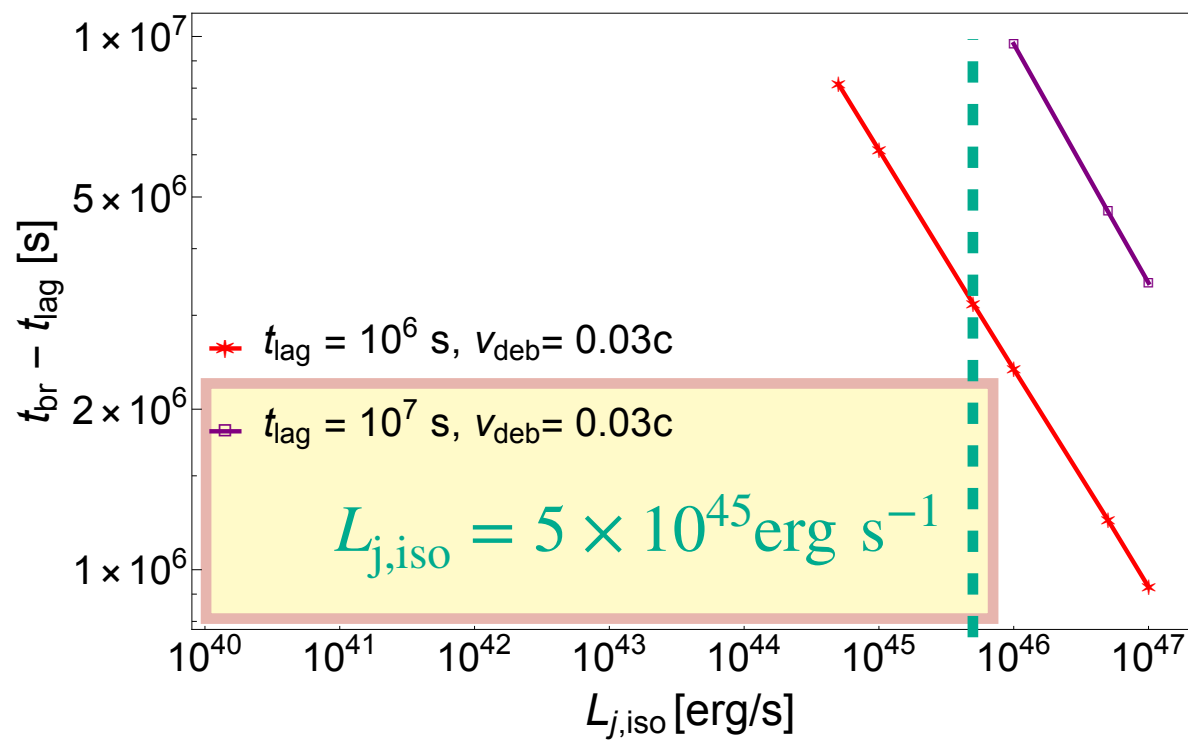


Motivations



Electromagnetic (EM) and Neutrino Signatures

$$L_{j,\text{iso}} \lesssim 3.2 \times 10^{45} \text{ erg/s} \left(\frac{N_s}{0.35} \right)^{-5} \left(\frac{M_{\text{deb}}}{0.5 M_{\odot}} \right) \left(\frac{\theta_0}{0.17} \right)^2 \left(\frac{\beta_{\text{deb}}}{0.03} \right)^2 \left(\frac{t_{\text{dur}}}{10^7 \text{ s}} \right)^{-1} \left(\frac{\chi_{\text{lag}}}{2} \right)^2$$

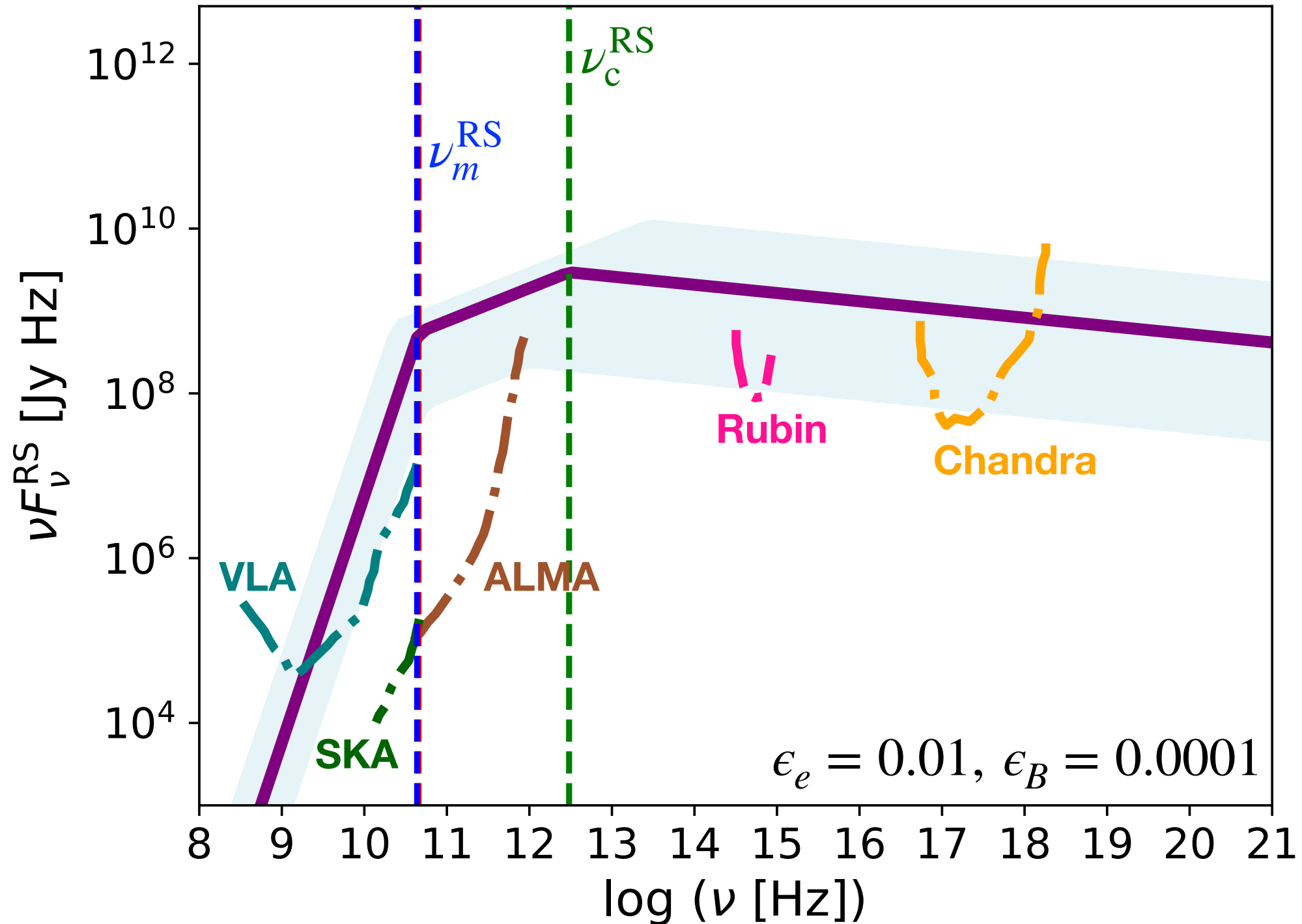


EM Signatures: Reverse Shock - Slow Cooling ($z = 0.05$)

$$F_{\text{syn,max}}^{\text{RS}} \simeq 37 \text{ mJy} (f_e/0.48) n_{2.53}^{\text{RS}} R_{h,16.21}^3 \Gamma_{0.70}^{\text{RS}} B^{\text{RS}} (1+z) d_{L,26.82}^{-2}$$

$$B^{\text{RS}} \simeq 0.32 \text{ G}$$

Reverse Shock (Slow cooling)

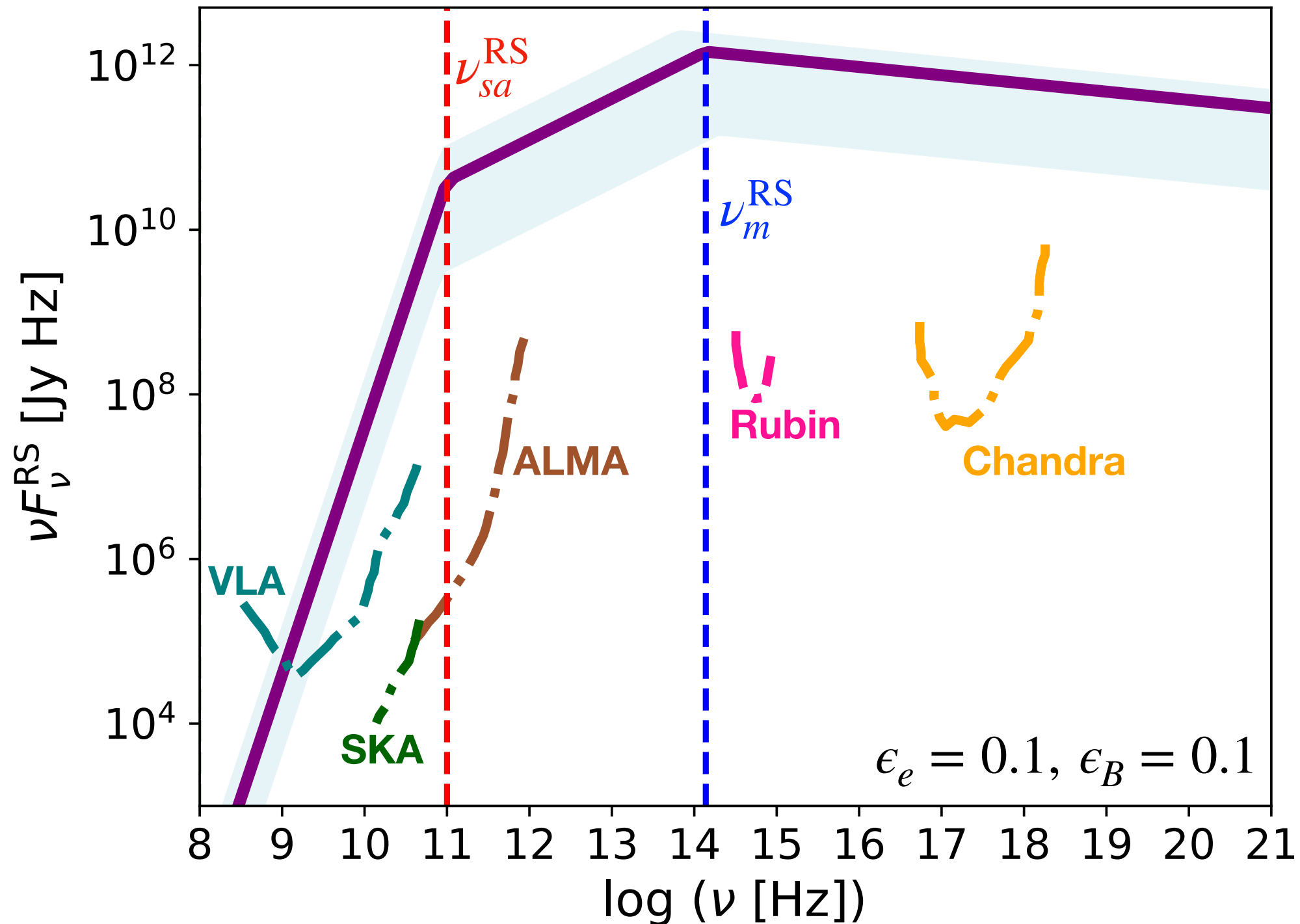


EM Signatures: Reverse Shock - Fast Cooling ($z = 0.05$)

$$F_{\text{syn,max}}^{\text{RS}} \simeq 37 \text{ mJy} (f_e/0.48) n_{2.53}^{\text{RS}} R_{h,16.21}^3 \Gamma_{0.70}^{\text{RS}} B^{\text{RS}} (1+z) d_{L,26.82}^{-2}$$

$$B^{\text{RS}} \simeq 10.25 \text{ G}$$

Reverse Shock (Fast cooling)

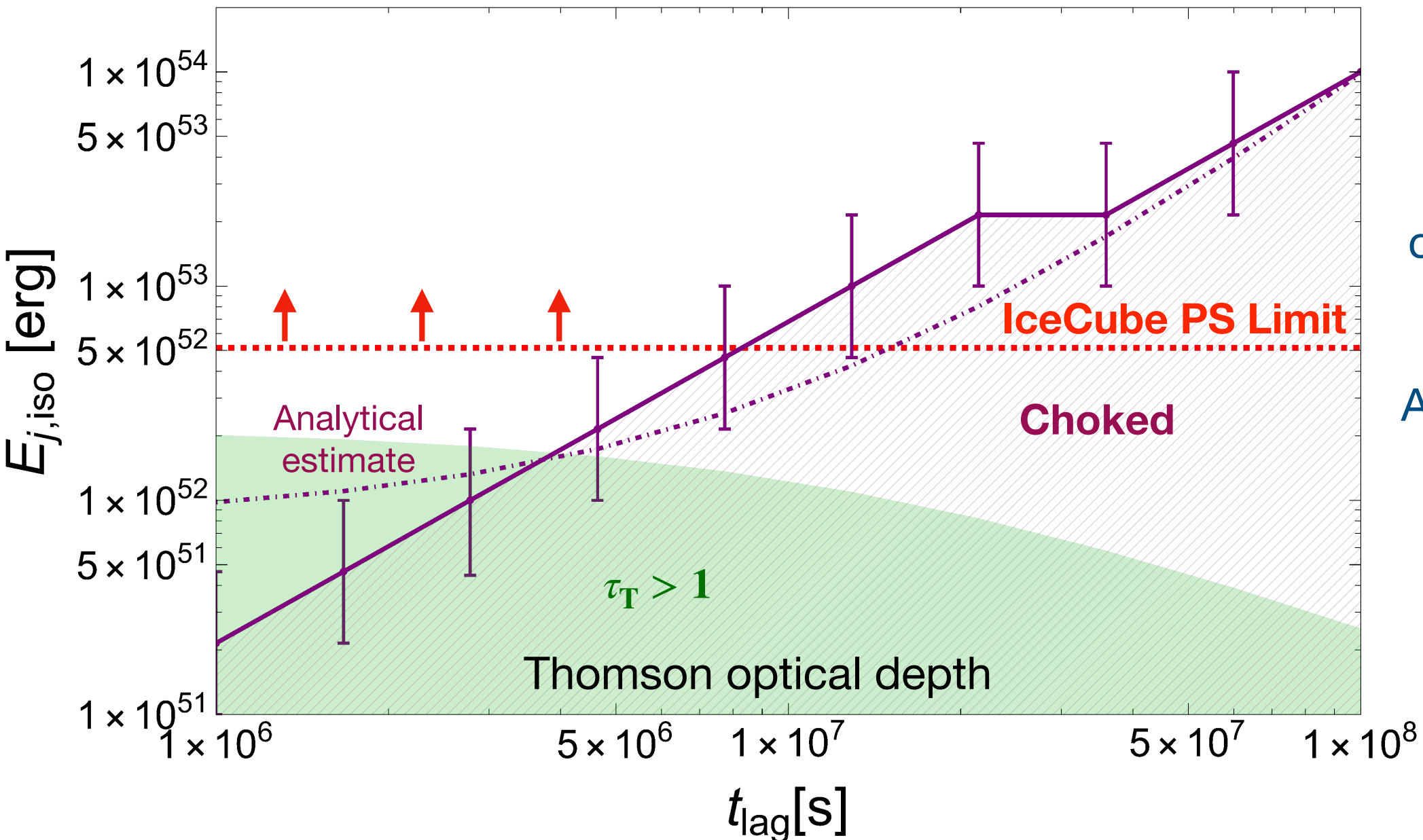


Neutrino Signatures: Choked jets ($t_{\text{lag}} = 10^7 \text{ s}, z = 0.05$)

What is the energy budget required for the jet to produce 1 neutrino event given the IceCube point source (PS) limit

$$E_\nu = 1\% \text{ of } E_j$$

AT2019dsg



Can explain the coincident neutrino observations

Also for AT2019aalc

AT2019fdr is still challenging

Outline

Part 1: Can choked delayed jets explain the neutrino coincidences associated with TDEs?

Based on: [Multi-messenger signatures of delayed choked jets in tidal disruption events](#)
MM, M. Bhattacharya, K. Murase
[\(submitted to MNRAS\)](#) (*arXiv: 2309.02275*).

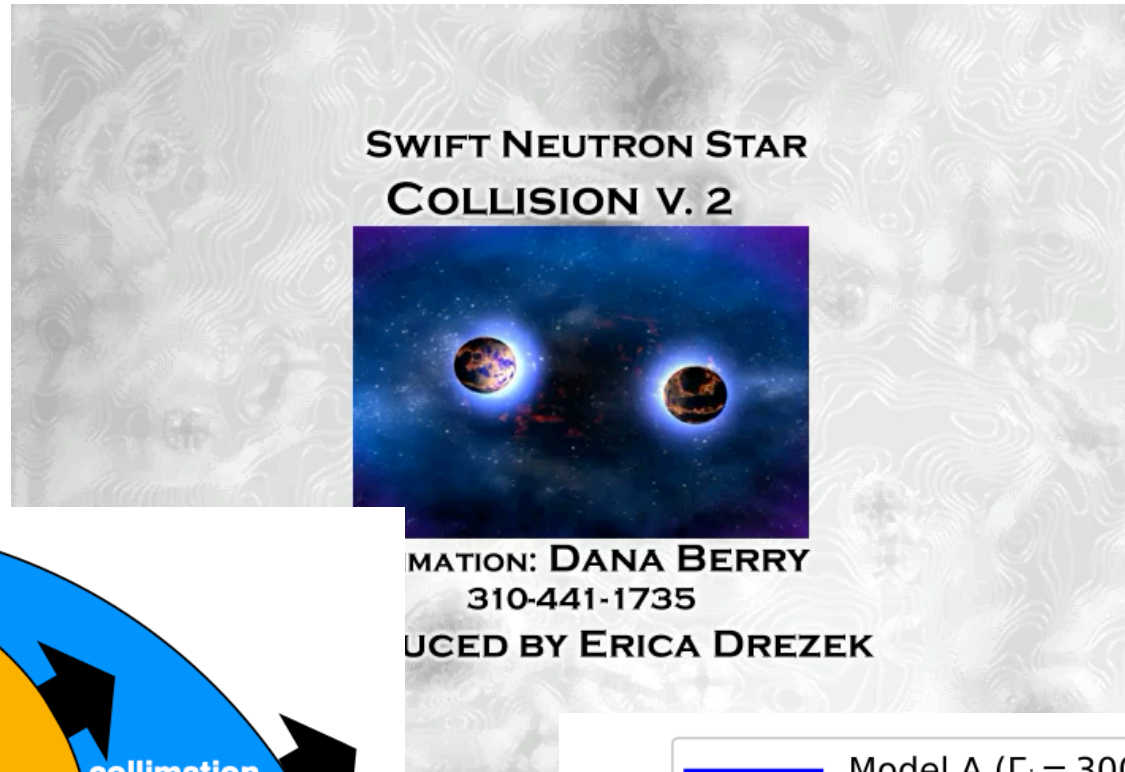
Part 2: Hunting for neutrinos from BNS mergers at next-generation GW and neutrino detectors

Based on: [Gravitational wave triggered high energy neutrino searches from BNS mergers: prospects for next generation detectors](#)
MM, S. S. Kimura, K. Murase
[Phys. Rev. D 109, 4, 043053 \(2024\)](#) (*arXiv: 2310.16875*)

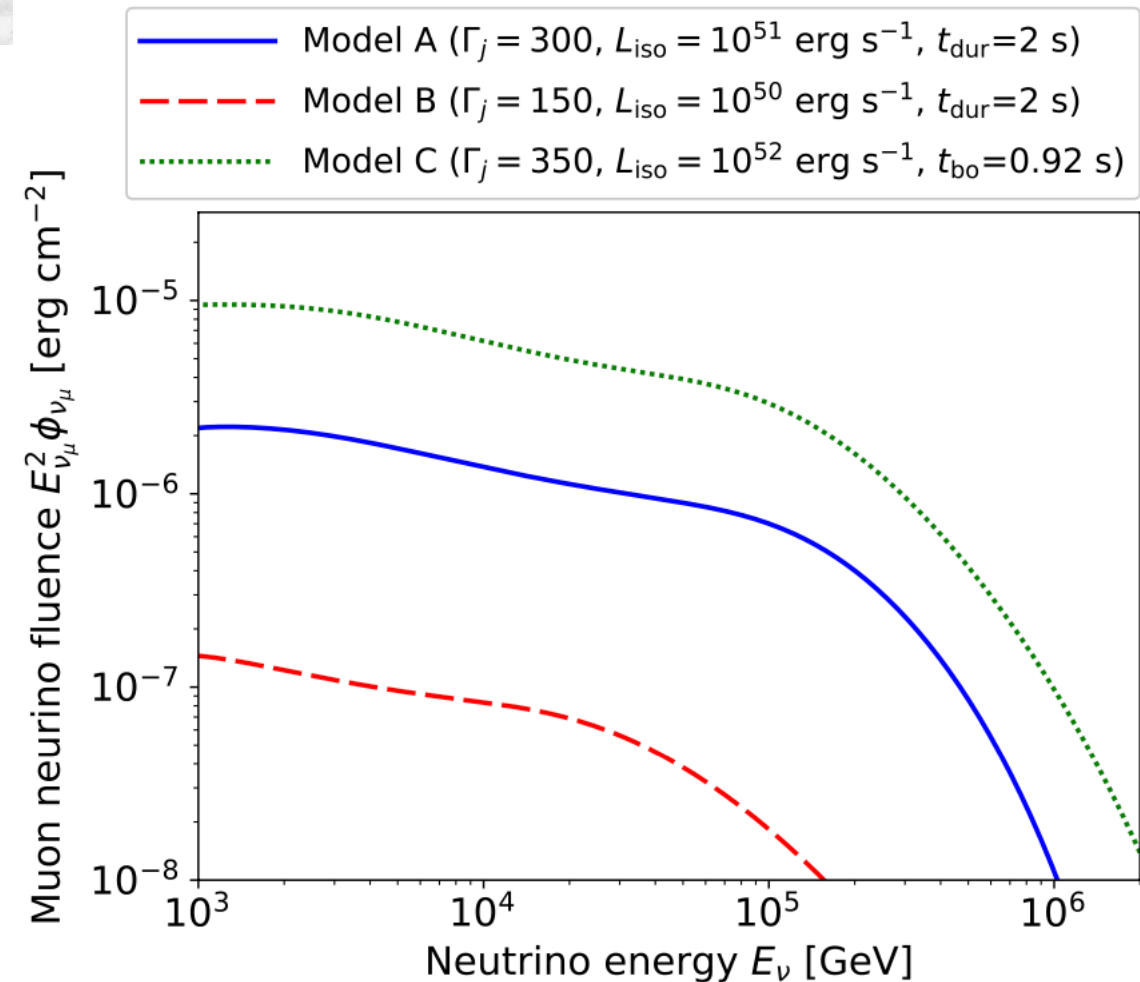
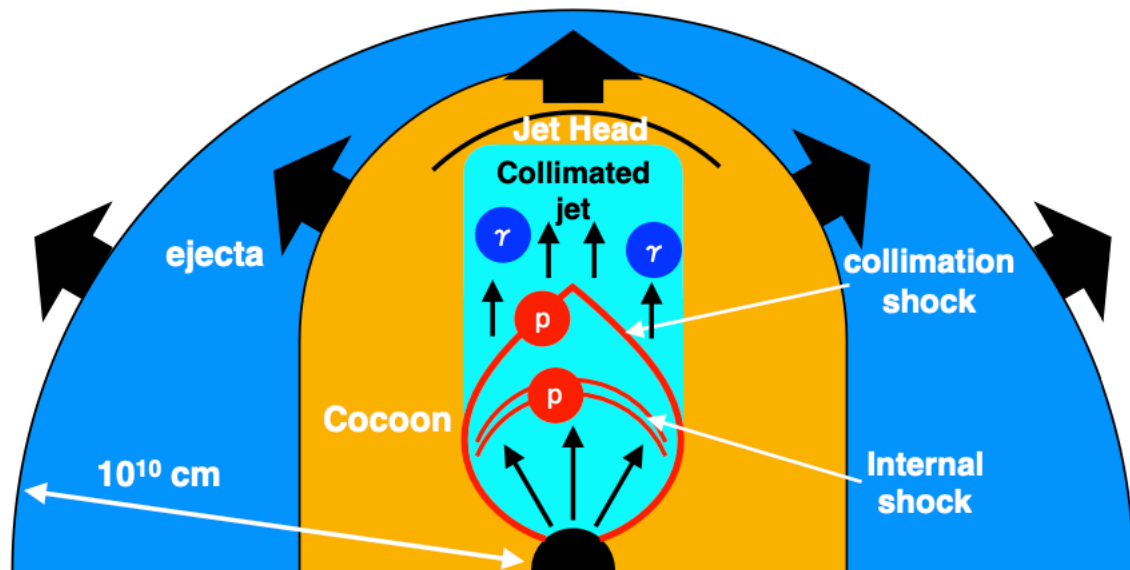
Part 3: Constraints from non-detection of neutrinos from the BOAT - GRB 221009A

Based on: [Neutrinos from the Brightest Gamma-Ray Burst?](#)
K. Murase, MM, A. Kheirandish, S. S. Kimura, K. Fang
[ApJ Letters 941 \(2022\) 1, L10](#) (*arXiv: 2210.15625*)

High energy neutrinos from BNS mergers



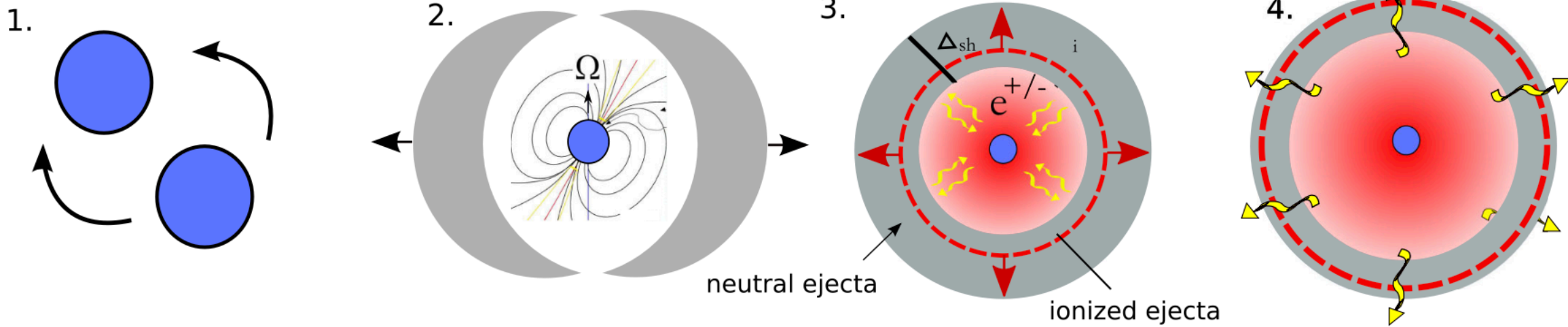
No detections yet :(



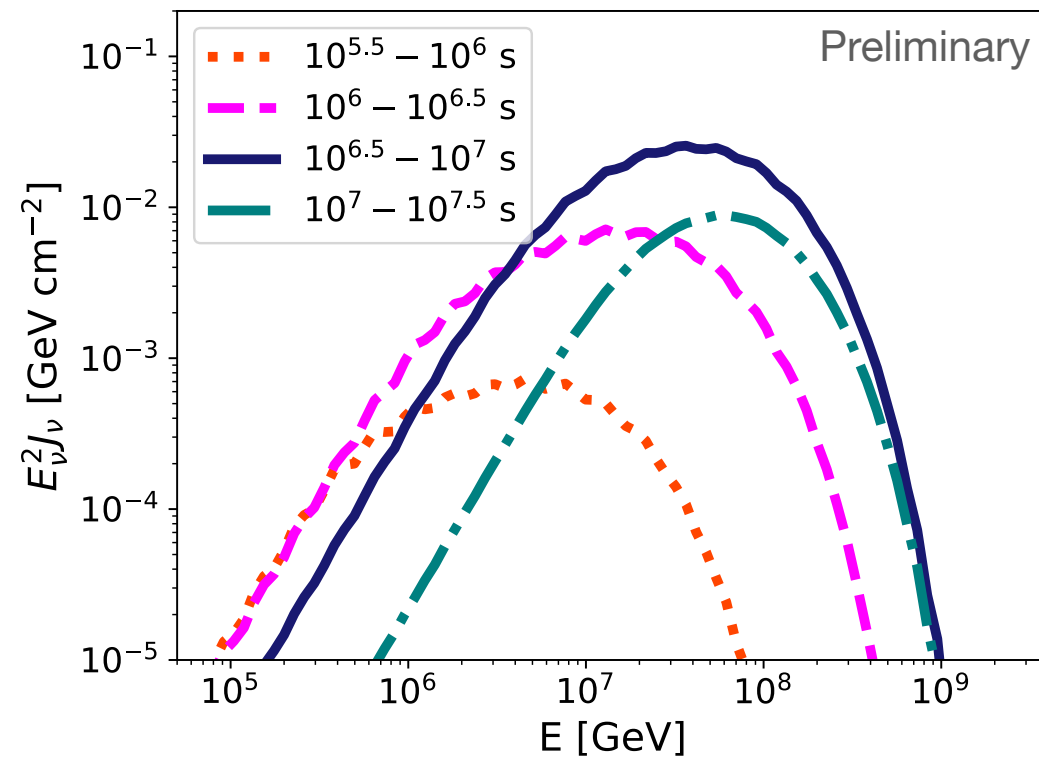
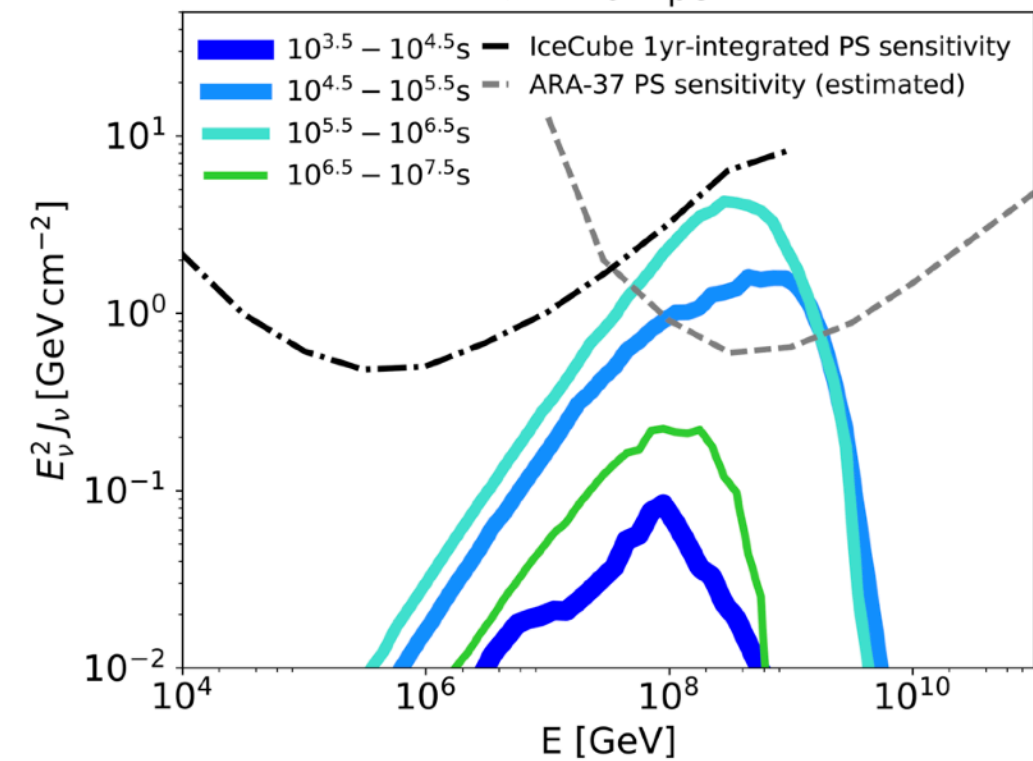
High energy neutrinos from BNS mergers



No detections yet :(

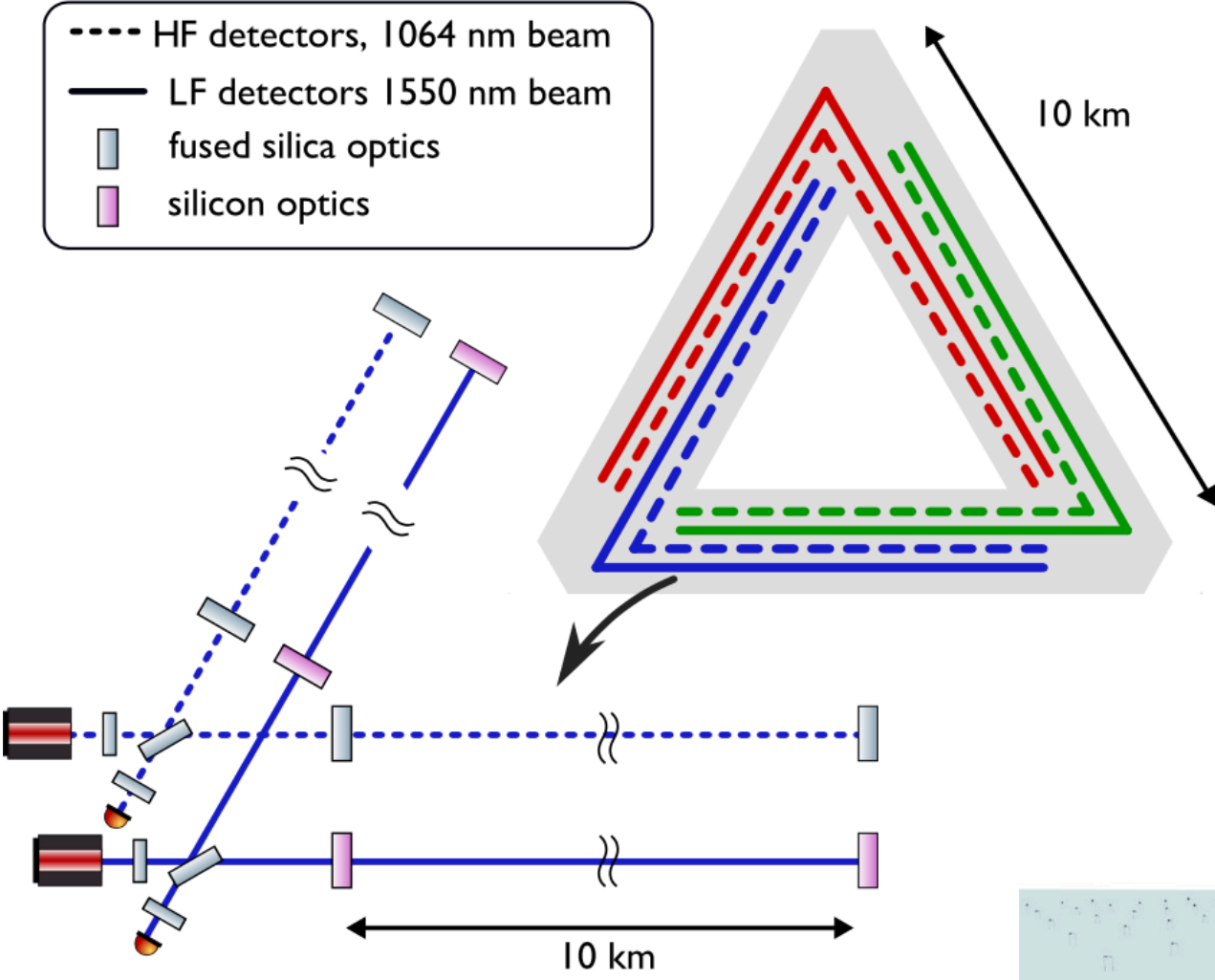


$D = 10$ Mpc

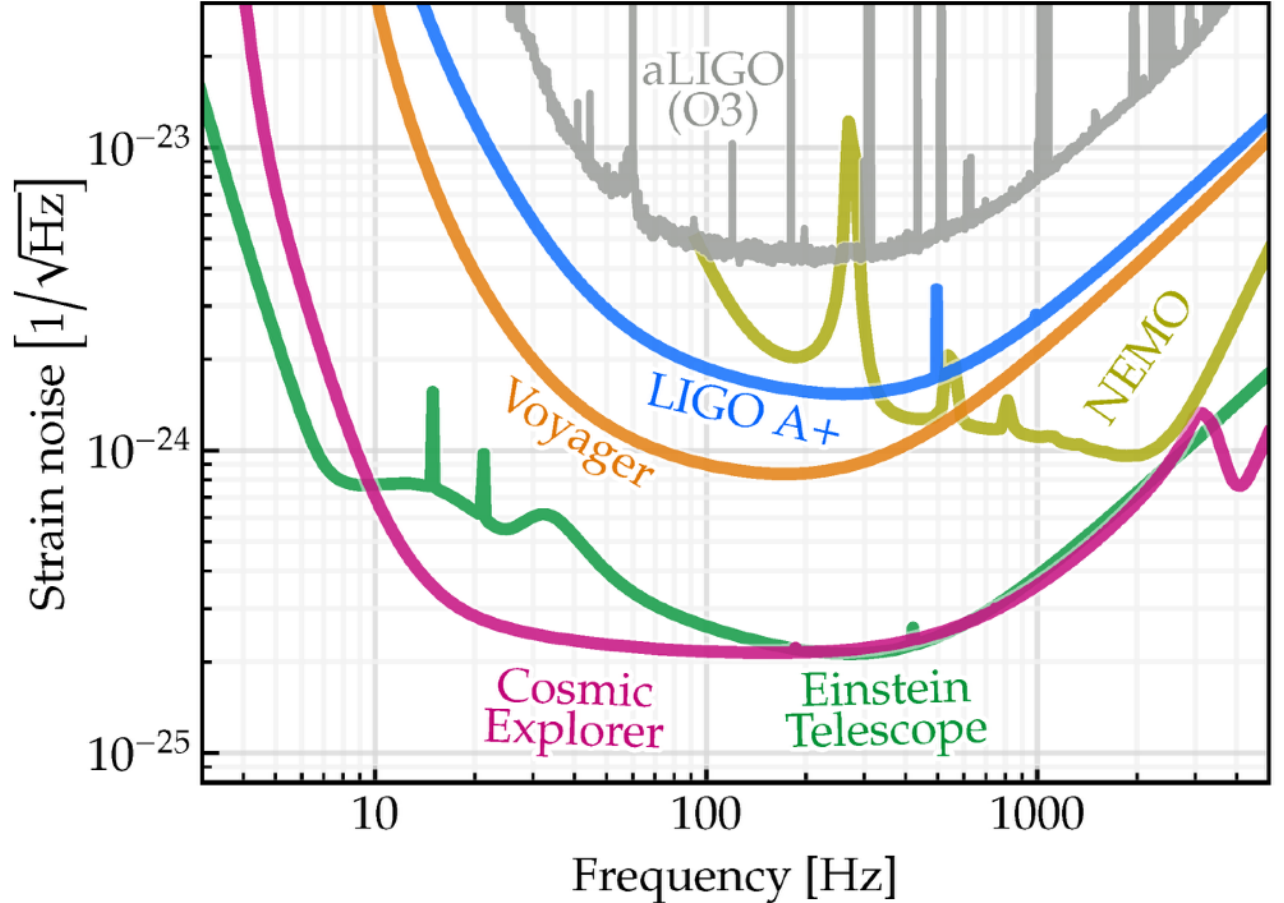


Fang and Metzger, *ApJ* (2017)
Metzger and Piro, *MNRAS* (2014)
MM and S.S. Kimura (in prep)

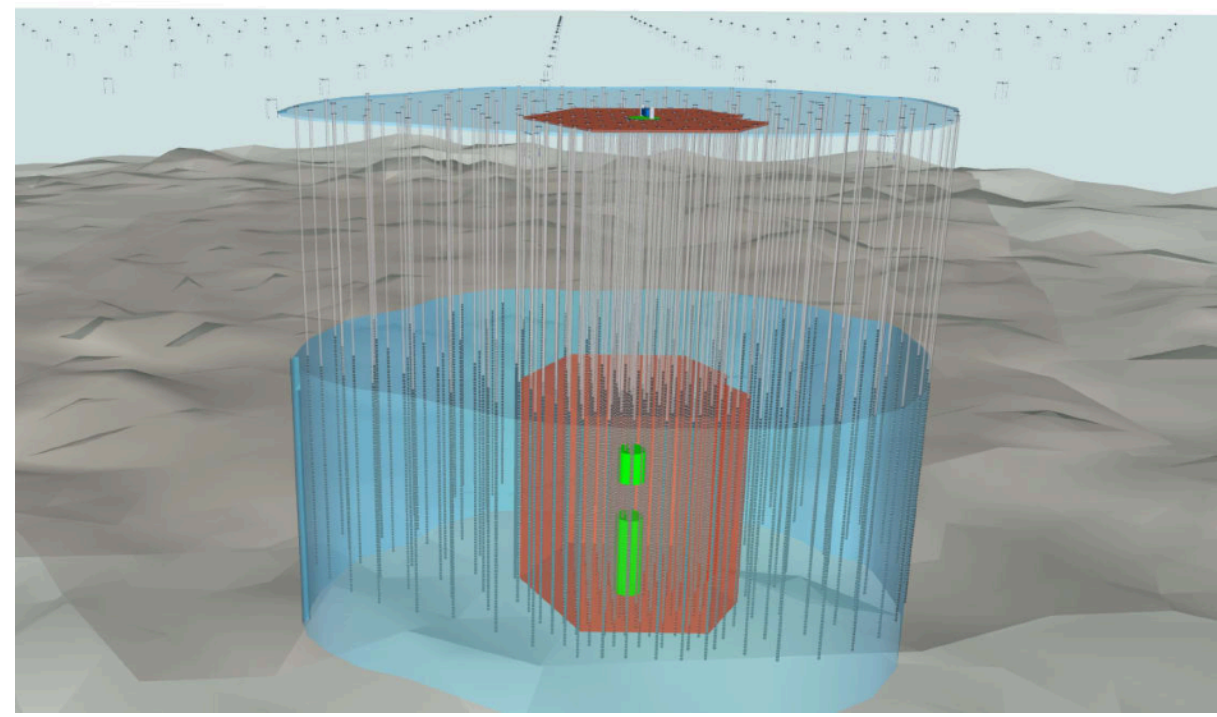
Next-generation GW and neutrino detectors



Einstein Telescope (ET)

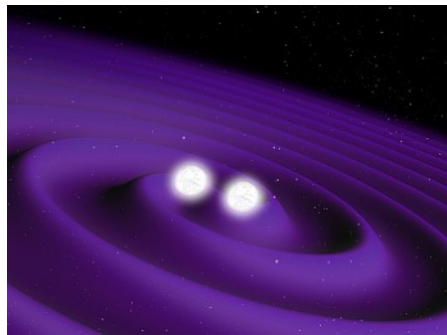


IceCube-Gen2



Evans et al., (2021)

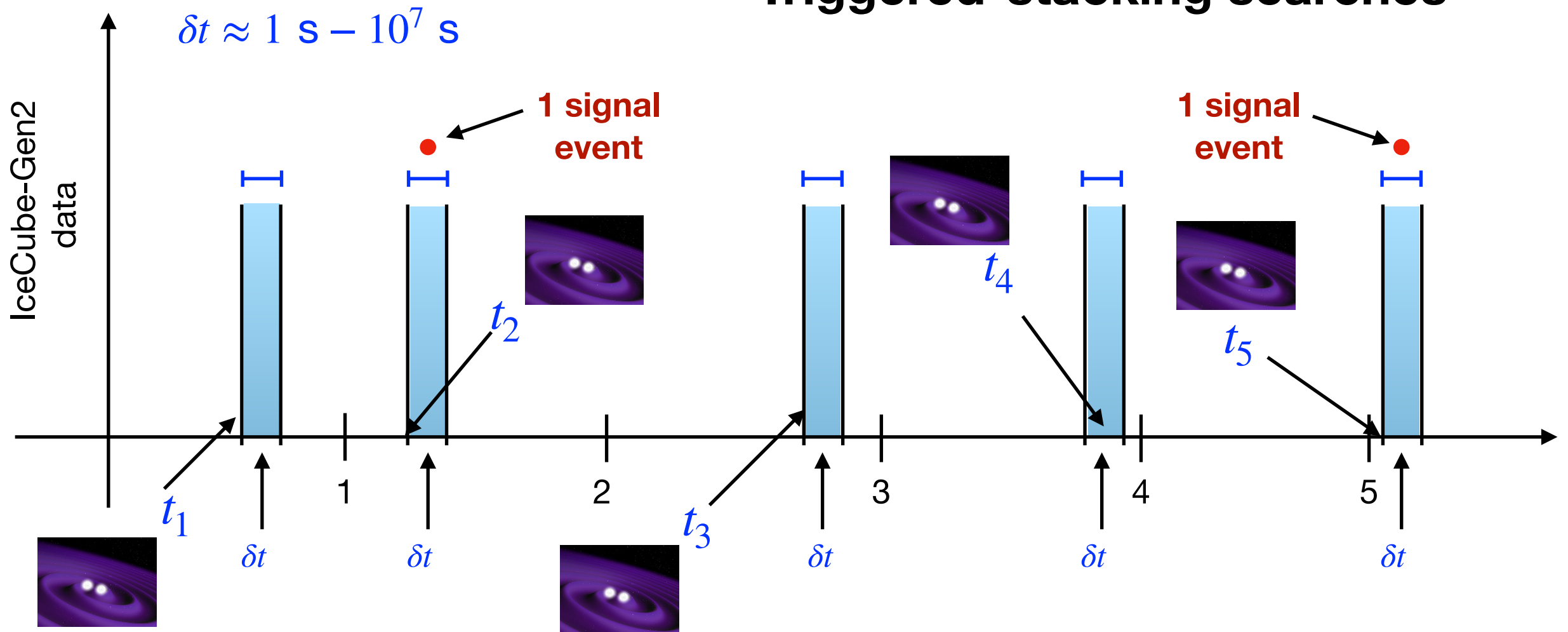
Detection strategy: triggered stacking search



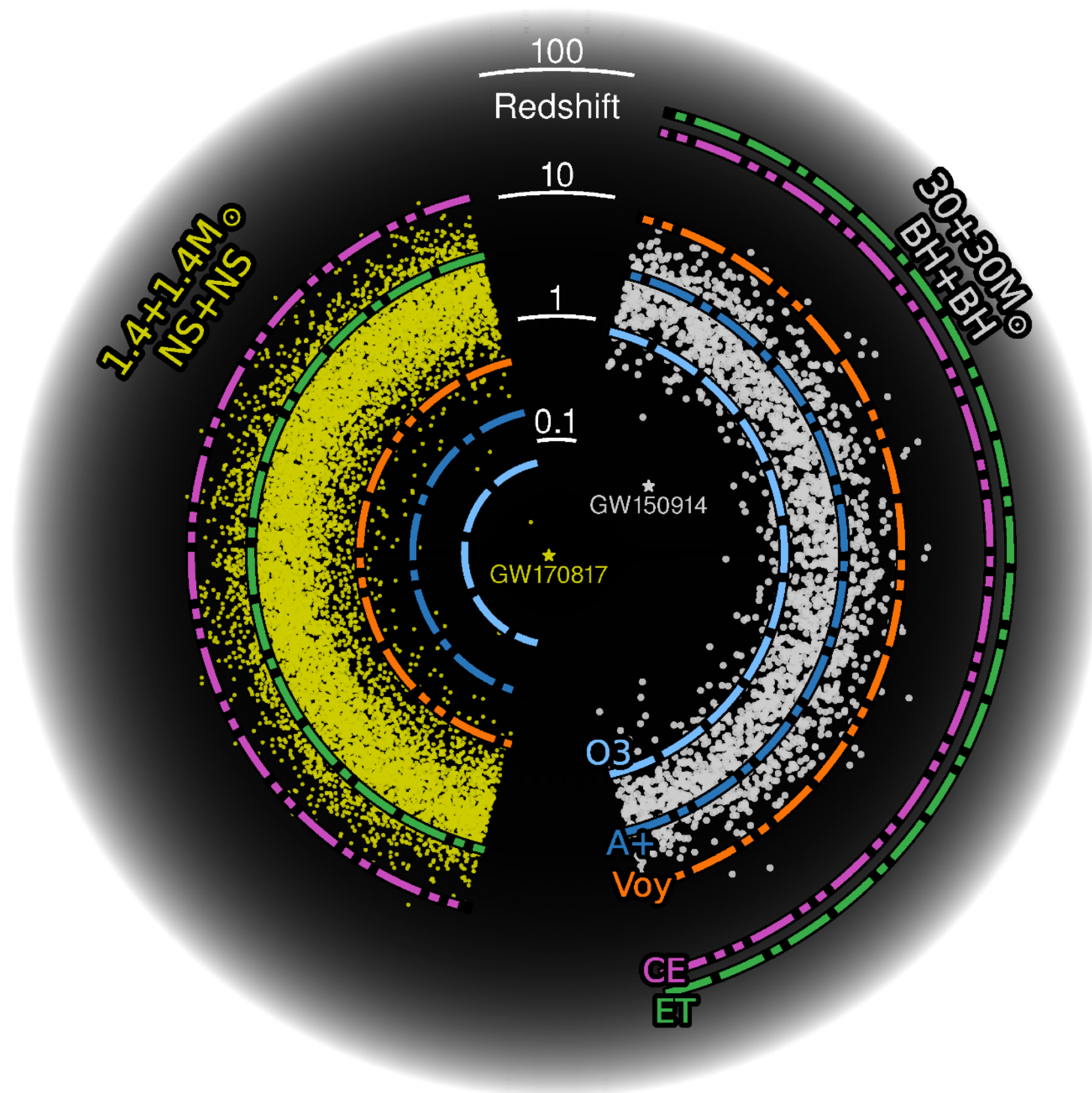
Trigger from next-gen GW detectors

Neutrinos in IceCube-Gen 2

Triggered-stacking searches

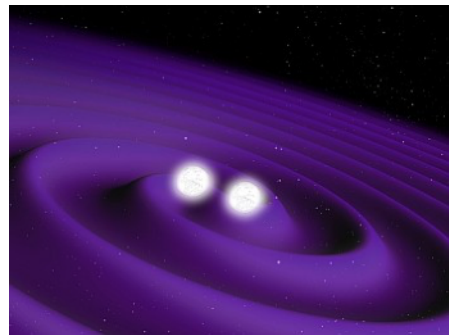


Next-generation GW detectors



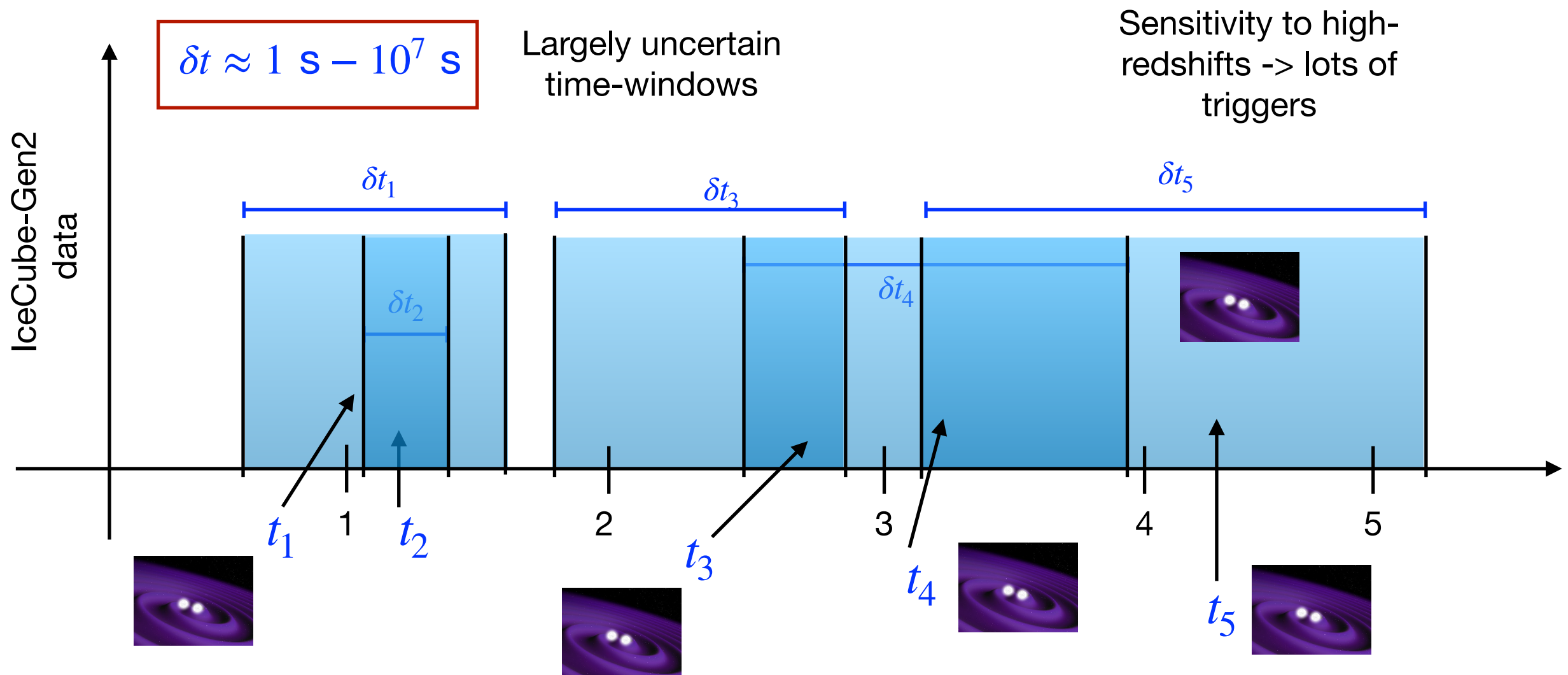
Sensitive to NS-NS
mergers from very
high redshifts

Impacts on triggered stacking searches



Trigger from next-gen GW detectors

Neutrinos in IceCube-Gen 2

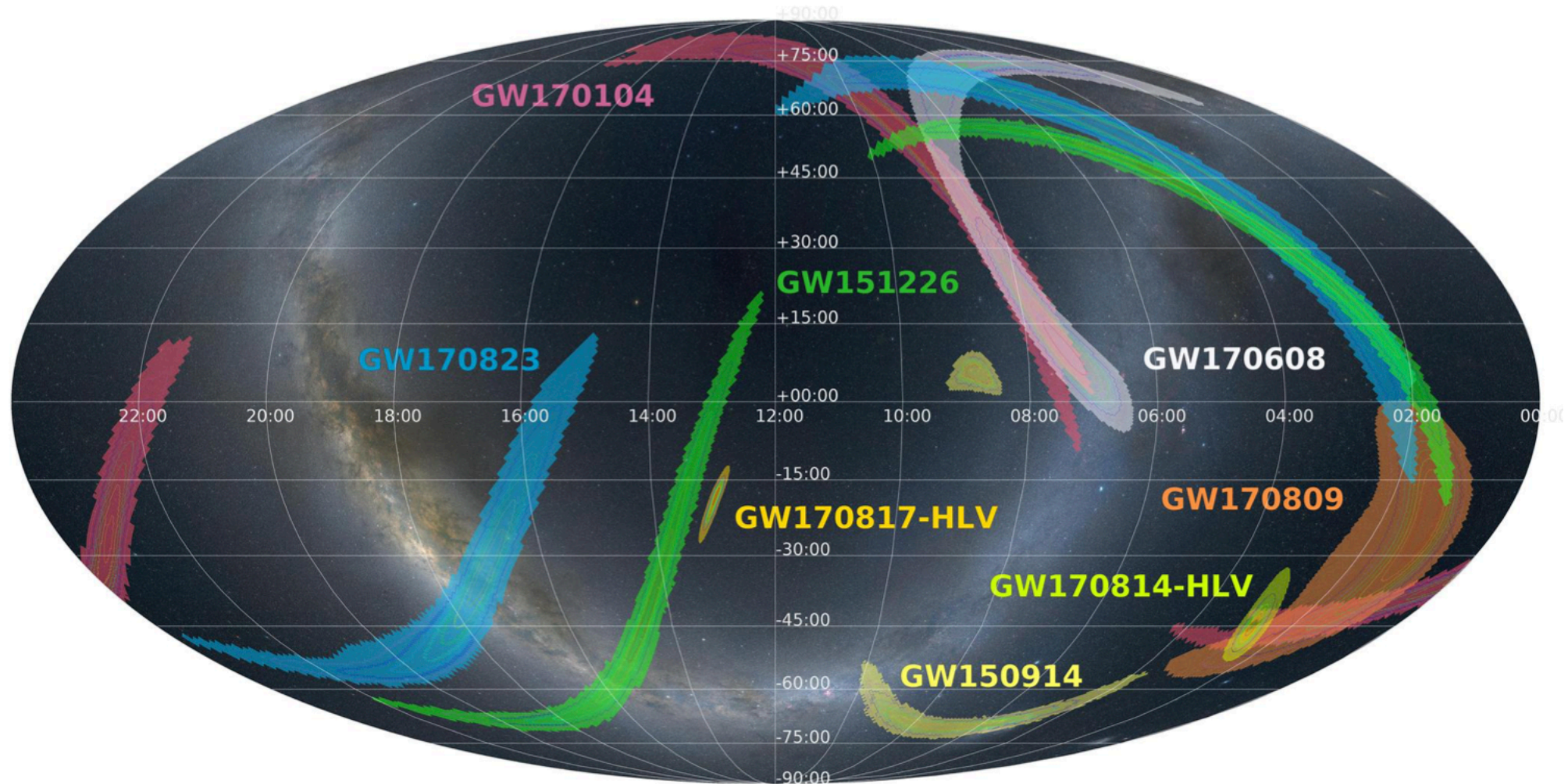


Spoils triggered stacking searches

How do we find meaningful triggers?

Motivations: How to obtain meaningful triggers?

Use the sky localization capabilities of the GW detectors....



Fraction of total sky area covered

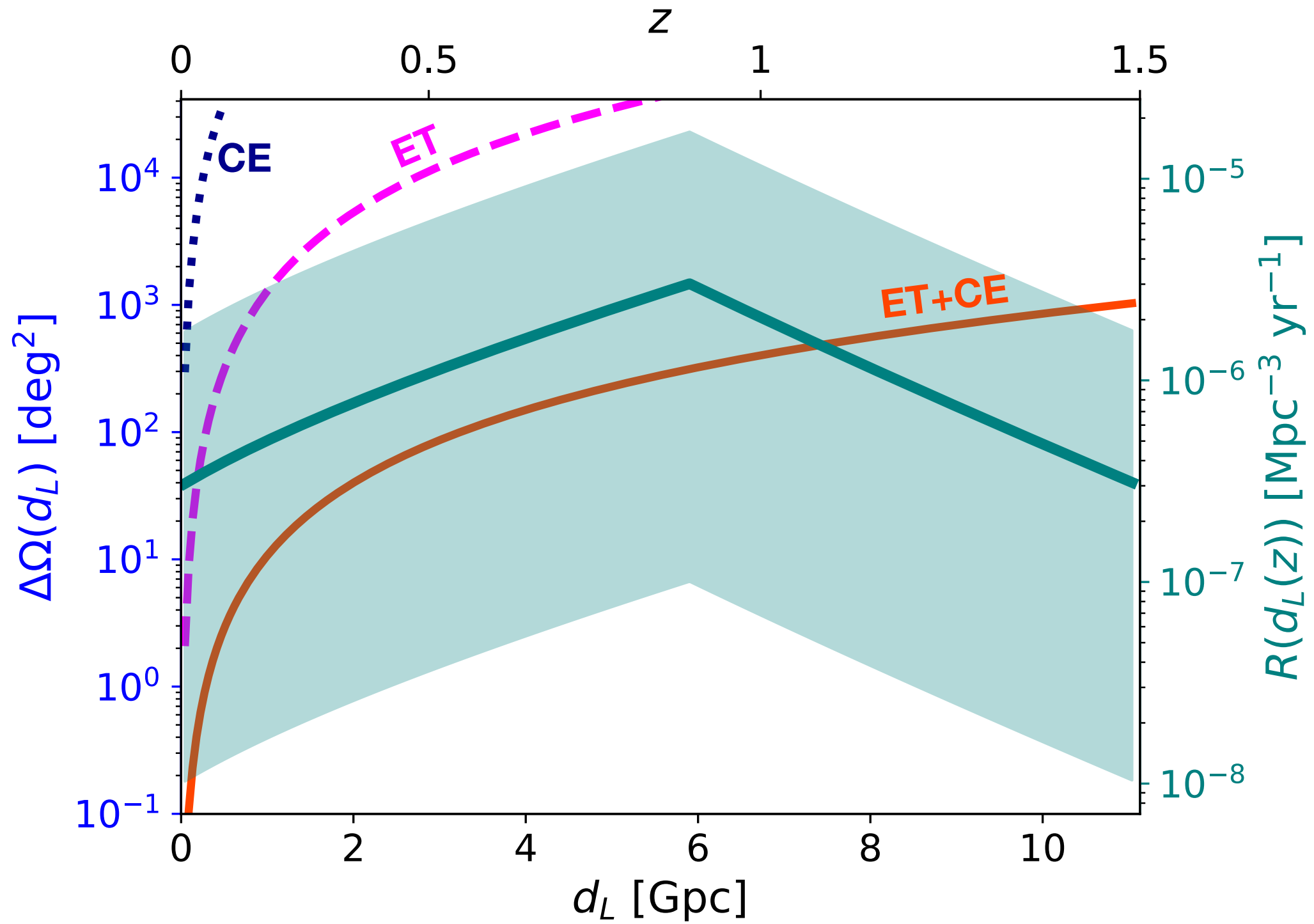


Set threshold: f_{th}



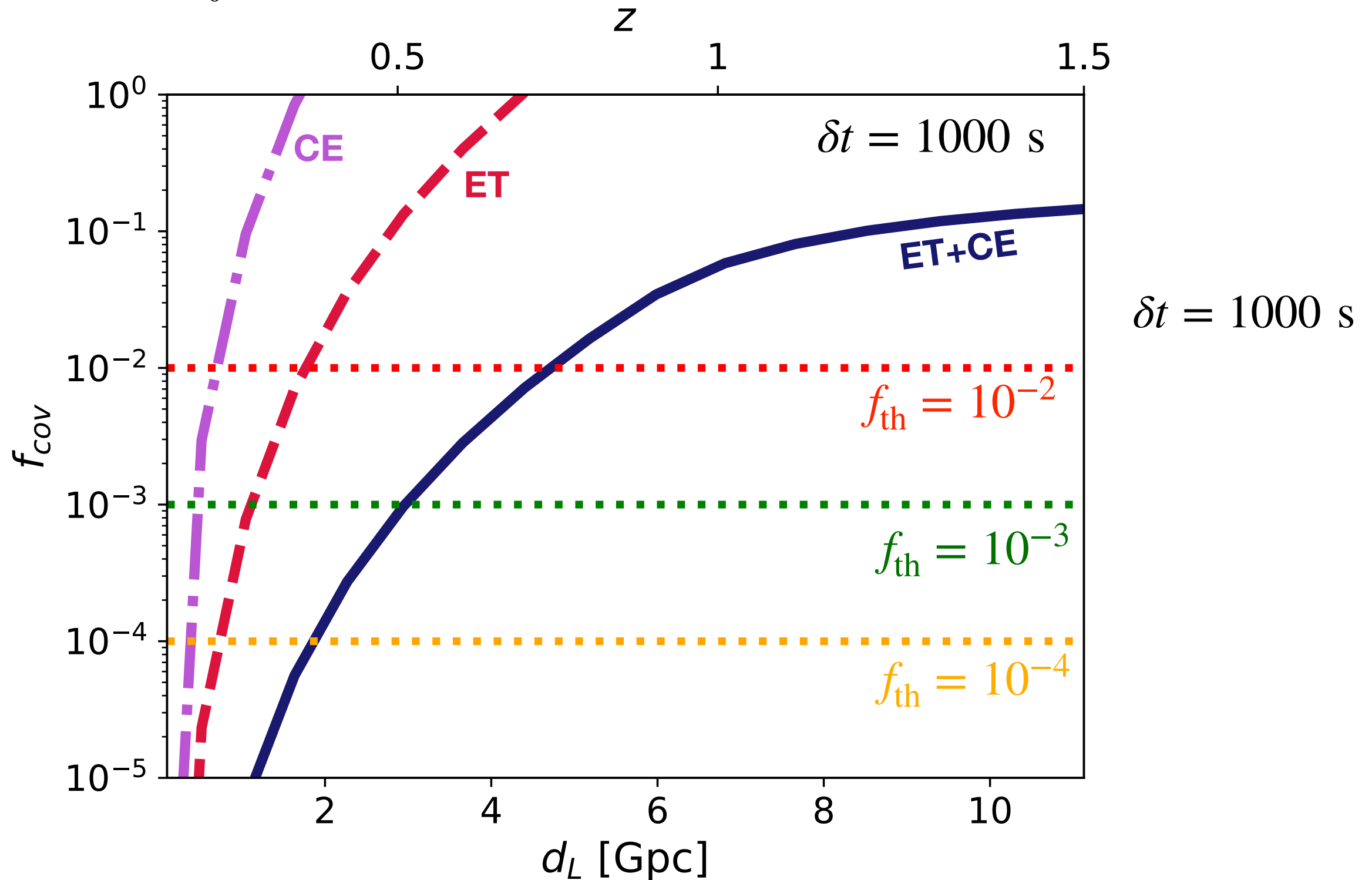
Obtain distance limits for GW detectors to collect meaningful triggers

Sky localization and BNS merger rate

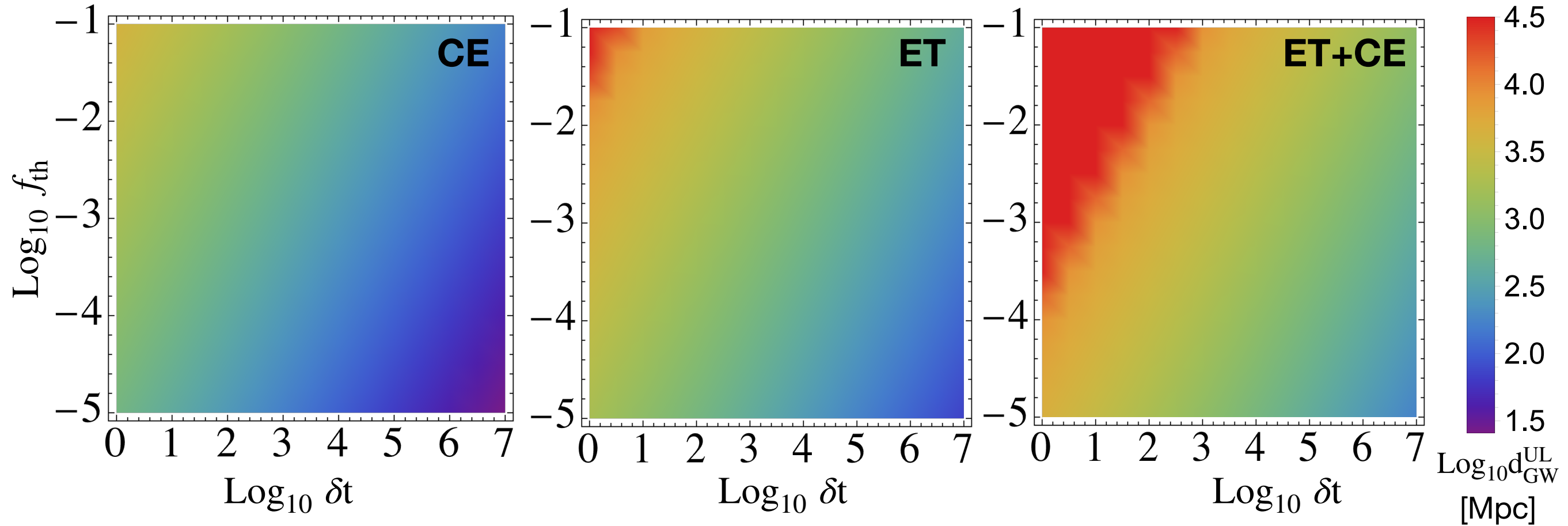


Distance limits for GW detectors

$$\int_0^{d_{\text{GW}}^{\text{lim}}} d(d_{\text{com}}) \frac{\Delta\Omega(d_L)}{4\pi} R(z) 4\pi d_{\text{com}}^2 \delta t = f_{\text{cov}}(d_{\text{GW}}^{\text{lim}})$$



Distance limits for GW detectors - $\delta t - f_{\text{th}}$ plane



High energy neutrinos from BNS mergers

Probability to detect more than one neutrino associated with GW signal in T_{op}

$$q(d_{GW}^{UL}, T_{op}) = 1 - \exp\left(-T_{op}I(d_{GW}^{UL})\right)$$

$$I(d_{GW}^{UL}) = 4\pi \int_0^{d_{GW}^{UL}} d(d_{com}) \frac{T_{op}}{(1+z)} R(z) d_{com}^2 P_{n \geq 1}(d_L)$$

Probability to detect more than one neutrino

$$d_{GW}^{UL} = \min(d_{GW}^{lim}, d_{GW}^{hor})$$

Depends on f_ν

Depends on δt

Assume a Poissonian probability

The event rate is calculated is convoluting the IceCube 10 years point source effective area with the muon neutrino flux

$$\phi_\nu(\mathcal{E}_\nu^{HE,iso}, E_\nu, d_L) = \frac{(1+z)}{4\pi d_L^2} \frac{\mathcal{E}_\nu^{HE,iso}}{\ln(\epsilon_\nu^{max}/\epsilon_\nu^{min})} E_\nu^{-2}$$

$$\mathcal{E}_\nu^{HE,iso} = \frac{\mathcal{E}_\nu^{HE,true}}{f_{bm}} = \left(\frac{f_\nu}{f_{bm}}\right) \mathcal{E}_{GW}$$

The flux is calculated assuming a $dN_\nu/dE_\nu \propto E_\nu^{-2}$ spectrum.

$$\mathcal{E}_\nu^{HE,true} = f_\nu \mathcal{E}_{GW}$$

$$\mathcal{E}_{GW} \sim \alpha \mathcal{E}^{tot}$$

$$\alpha \sim 1\%$$

Results - varying f_ν and δt

f_ν

10^{-5}

5×10^{-5}

Fiducial Parameters:

$$f_\nu = 2.5 \times 10^{-5}$$

$$\delta t = 1000 \text{ s}$$

$$E^{\text{tot}} \sim 5 \times 10^{54} \text{ erg}$$

Motivated by
physical models

1 s

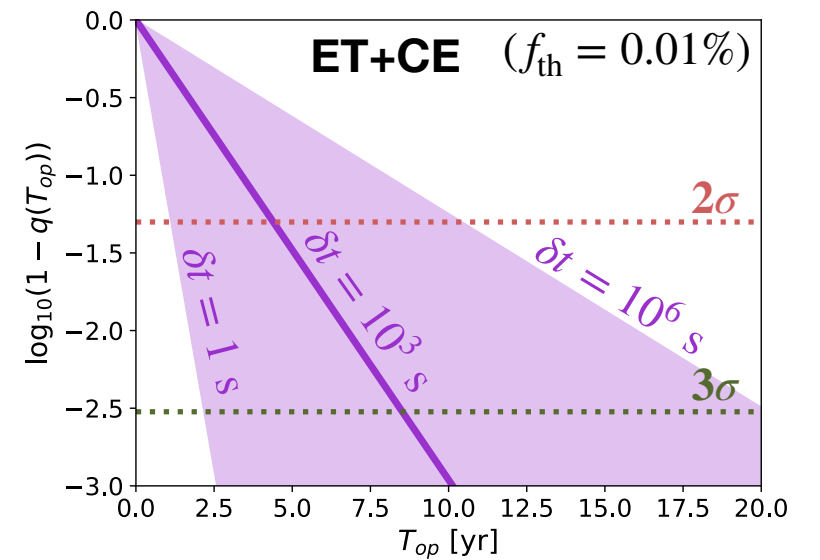
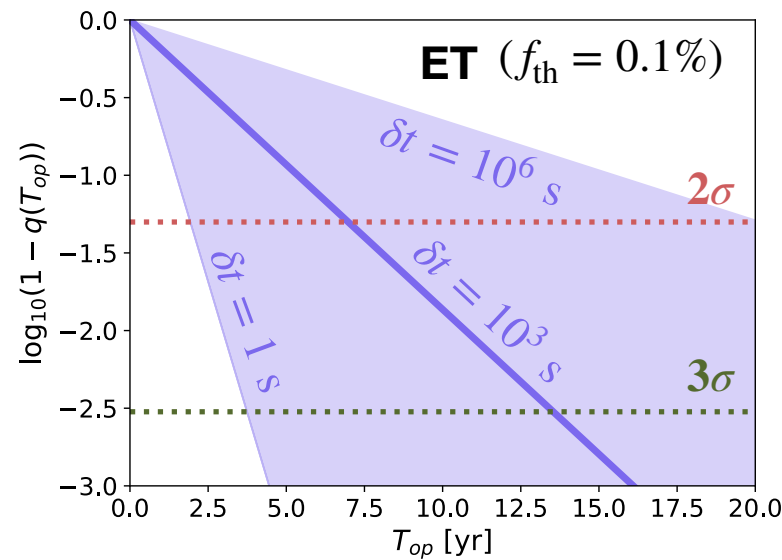
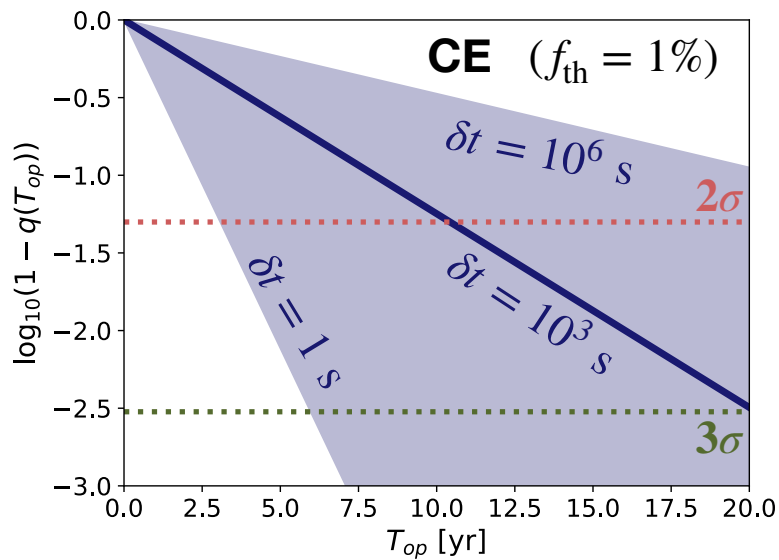
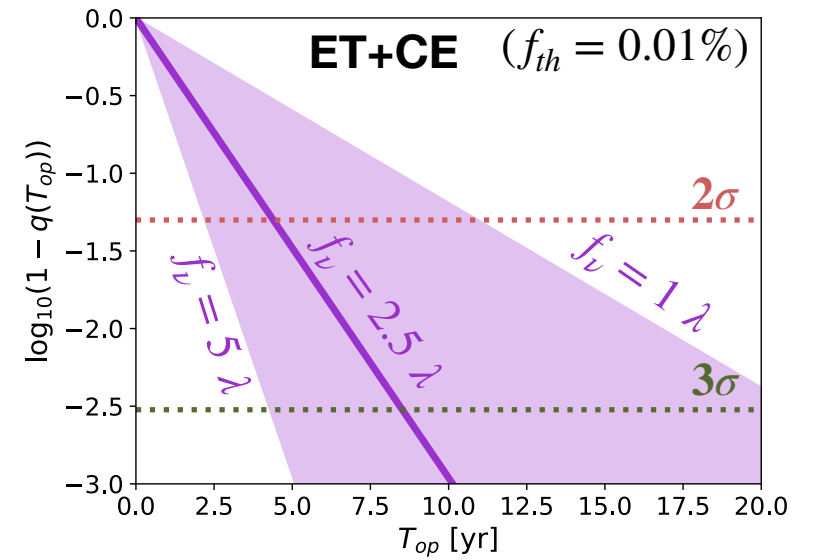
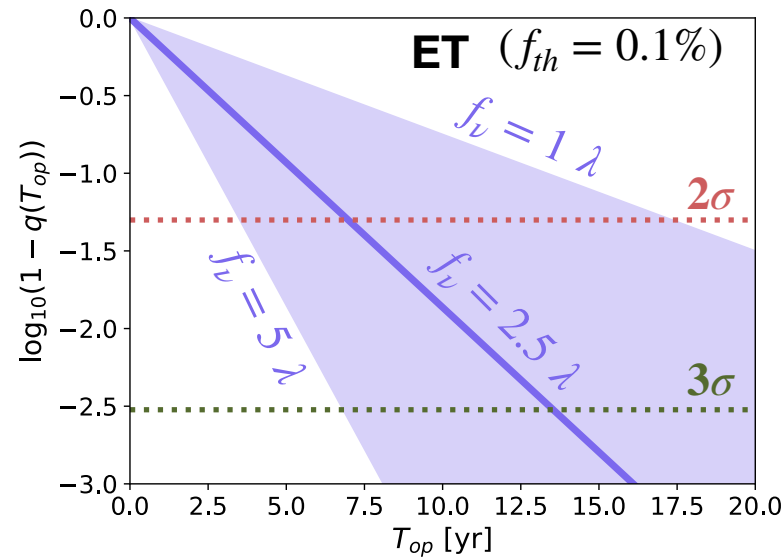
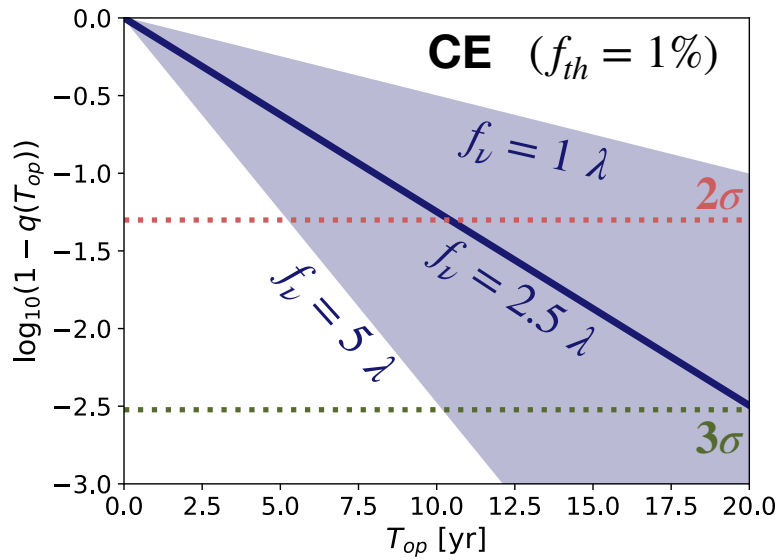
10^6 s

δt

Results - varying f_ν and δt

$$f_\nu$$

$$\delta t = 1000 \text{ s}$$

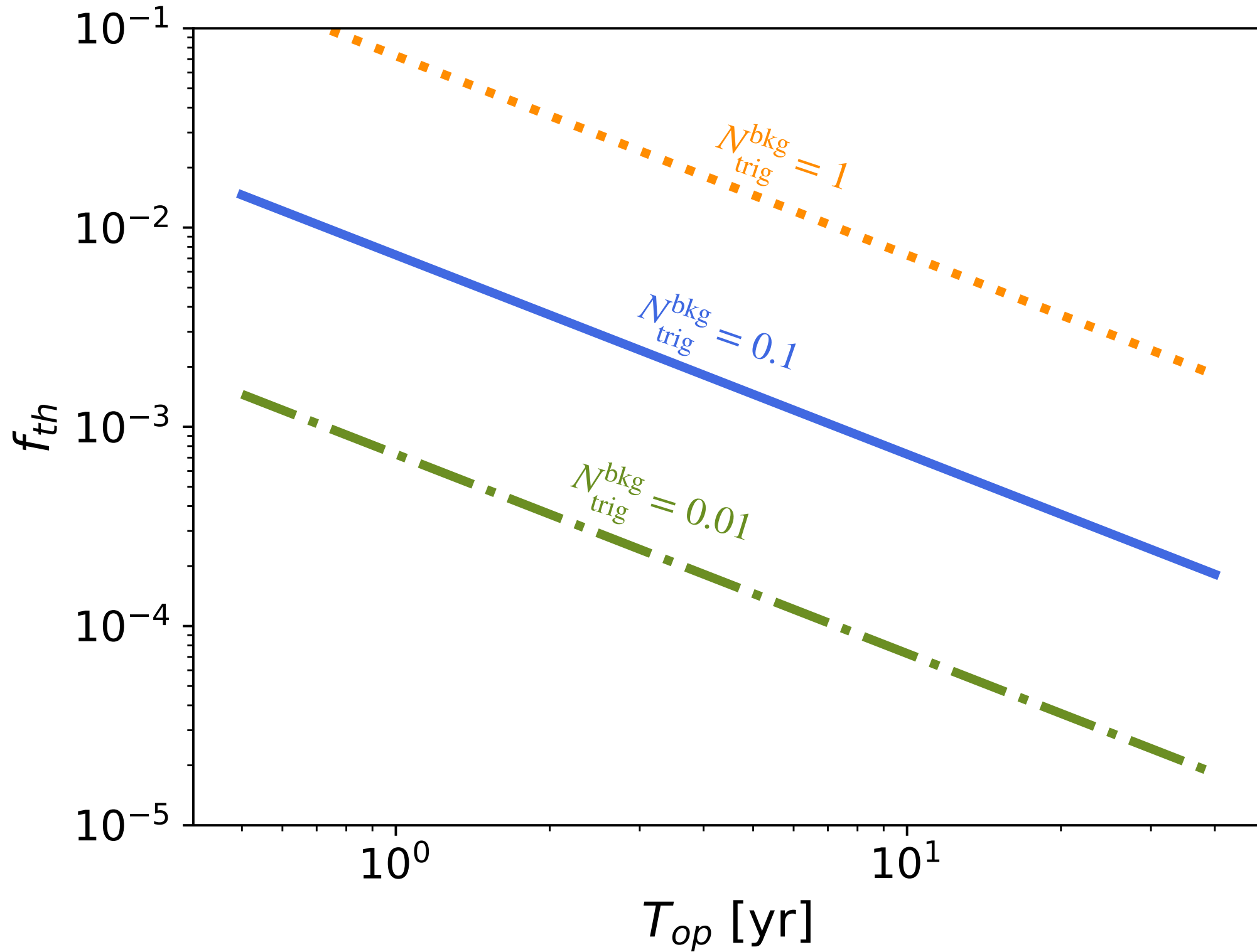


$$f_\nu = 2.5 \lambda$$

$$\delta t$$

$$\lambda = 10^{-5}$$

Backgrounds



Takeaways

Part 1: Multi-messenger signatures of choked delayed jets in TDEs

Late time activity associated with the SMBH from observations:

- Delayed radio flares
- Coincident neutrino detections: arrival after ~ 150 days, ~ 393 days, and 148 days post the optical peaks for AT2019dsg, AT2019fdr and AT2019aalc, respectively

Possibility of choked delayed jets

- Spherical debris envelope surrounding the SMBH, expanding outwards possibly driven by wind.
- Jet-cocoon interactions: collimation and choking - Higher delay times and debris velocities help with choking

Electromagnetic and neutrino signatures

- Synchrotron radiation from delayed choked jets: Reverse shock: slow and fast cooling cases
- Optical and X-ray observatories: good prospects, radio observations seem likely as well.
- Can explain the coincident observations by IceCube - AT2019dsg and AT2019aalc with this scenario of choked delayed jets.

Part 2: GW triggered searches for high energy neutrinos from BNS mergers: prospects for next-gen

- ET+CE can give coincident neutrino events or 3σ level constraints on the parameter space, due to extremely good sky localization capabilities over a timescale of ~ 20 years even for the less optimistic scenarios.
- ET can lead to 2σ constraints owing to its good sky localization capabilities over a time scale of 20 - 30 years for the less optimistic cases. CE has comparatively poor sky localization and hence may be good for coincident detections or 2σ -level constraints over reasonable time scales for optimistic parameters.
- Our analysis can constrain f_ν for a population of BNS sources: understanding emissions from BNS mergers
- Model independent analysis can help constrain models: neutrinos from choked jet scenarios and hence provide insights regarding GRB jets, neutrino emission sites and mechanisms.

Thank You!

Backup

What are TDEs?

The shredding apart of a star when it comes close to a SMBH, due to its tidal forces

Disruption starts



~ Half the debris is lost: Unbound orbit
~ Half the debris falls back: Bound orbit



Debris circularizes



Part of the debris may form an accretion disk

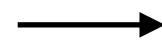
(Timescales are also uncertain)

Tidal disruption radius: $R_T = R_* \frac{M_{BH}}{M_*}$

Fallback time: $t_{fb} = 2\pi \sqrt{a_{min}^3 / (GM_{BH})}$

Semi-major axis: $a_{min} \approx R_T^2 / (2R_*)$

Circularization radius: $R_{circ} = 2R_T$

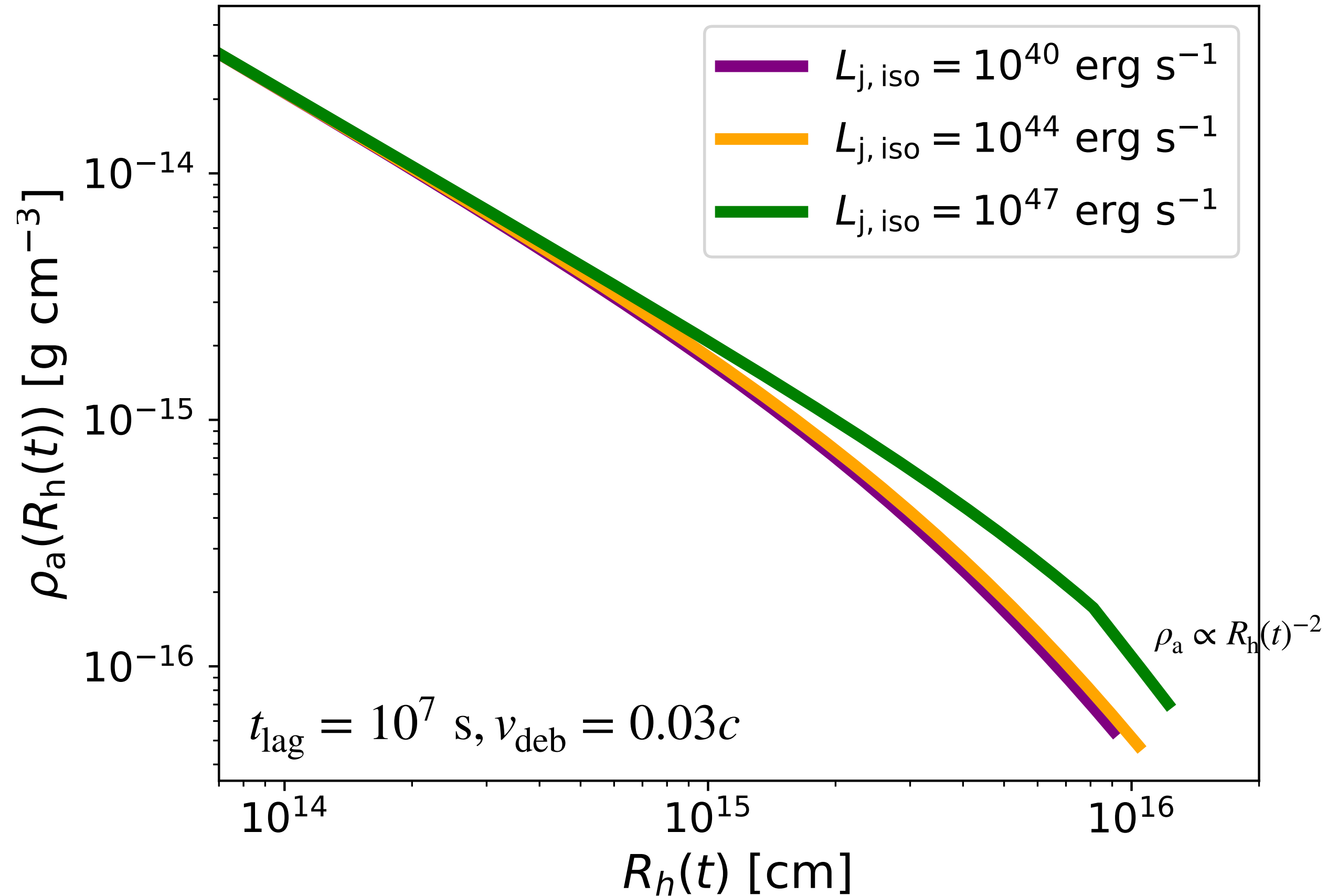


Winds, Outflows,
etc.

The neutrino associations

Property	AT2019dsg	AT2019fdr	AT2019aalc
TDE	yes	strong candidate	candidate
Peak bol. luminosity	$3.5 \times 10^{44} \text{ erg s}^{-1}$	$1.3 \times 10^{45} \text{ erg s}^{-1}$	–
SMBH Mass	$10^6 - 10^{6.7} M_{\odot}$	$10^{7.55} M_{\odot}$	$10^{7.2} M_{\odot}$
Radio	evolving	not evolving	archival det.
UV	very bright	bright	–
X-ray	early, soft spectrum	late, soft spectrum	soft spectrum
Dust echo strength	92.2	39.2	15.7
ν delay	5 months	10 months	5 months
ν production	possible	possible	possible
ν energy	217 TeV	82 TeV	176 TeV
ν 90% uncertainty box	25.5 sq. deg	25.2 sq. deg	61.2 sq. deg
ν signalness	0.59	0.59	0.45

Ambient medium density (ρ_a)



Dynamics: the land of choked jets

$L_{j,\text{iso}}$ (in erg/s)	$t_{\text{lag}} = 10^6 \text{ s}$	$t_{\text{lag}} = 10^7 \text{ s}$	$t_{\text{lag}} = 10^8 \text{ s}$
$10^{40} - 10^{43}$		Collimation; No breakout	
5×10^{43}	Collimation; No breakout	Collimation; No breakout	No Collimation; No breakout
10^{44}	Collimation; No breakout	Collimation; No breakout	No collimation; No breakout
5×10^{44}	Collimation; Breakout	No collimation; No breakout	No collimation; No breakout
10^{45}	Collimation; Breakout	No collimation; No breakout	No collimation; No breakout
5×10^{45}	Collimation; Breakout	No collimation; No breakout	No collimation; No breakout
10^{46}	Collimation; Breakout	No collimation; Breakout	No collimation; No breakout
5×10^{46}	No collimation; Breakout	No collimation; Breakout	No collimation; No breakout
10^{47}	No collimation; Breakout	No collimation; Breakout	No collimation; No breakout

$L_{j,\text{iso}}$ (in erg/s)	$v_{\text{deb}} = 0.01c$	$v_{\text{deb}} = 0.03c$	$v_{\text{deb}} = 0.1c$
$10^{40} - 10^{44}$		Collimation; No breakout	
$5 \times 10^{44} - 10^{45}$		No collimation; No breakout	
5×10^{45}	No collimation; Breakout	No collimation; No breakout	No collimation; No breakout
10^{46}	No collimation; Breakout	No collimation; Breakout	No collimation; No breakout
$5 \times 10^{46} - 10^{47}$		No collimation; Breakout	

Electromagnetic (EM) signatures

$$\nu_{\alpha}^{\text{ES}} = \frac{3}{4\pi} \frac{eB^{\text{ES}}}{m_e c} \frac{\Gamma^{\text{ES}}}{(1+z)} (\gamma_{\alpha}^{\text{ES}})^2$$

ES: External shock can be Forward or Reverse shock region

α : Can be injection frequency (m) or cooling frequency (c)

B: Magnetic field strength in the region

Γ : Bulk Lorentz factor in the shocked region

γ : Lorentz factor associated with the electrons

The absorption frequency ν_{sa} is given by setting the synchrotron self-absorption optical depth to 1

$$B^{\text{ES}} = \left[32\pi\epsilon_B \Gamma^{\text{ES}} (\Gamma^{\text{ES}} - 1) n^{\text{ES}} m_p c^2 \right]^{1/2}$$

↑
Fraction of electron energy converted
to magnetic field energy