

FLAVOR

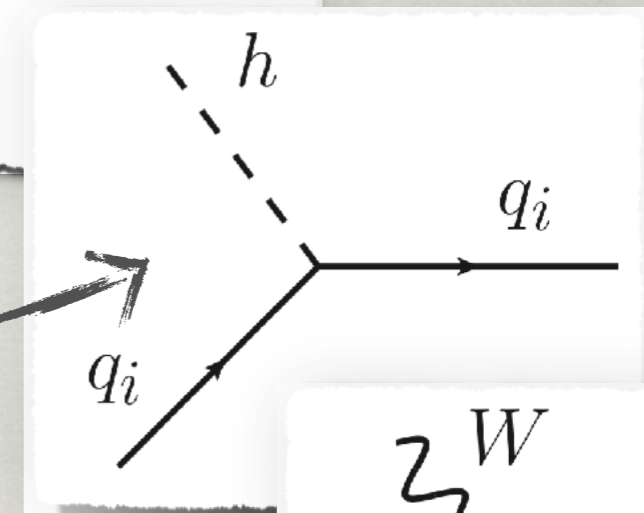
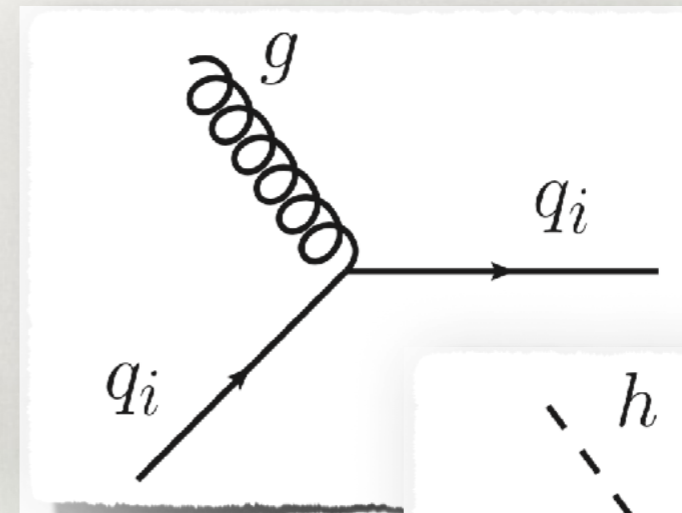
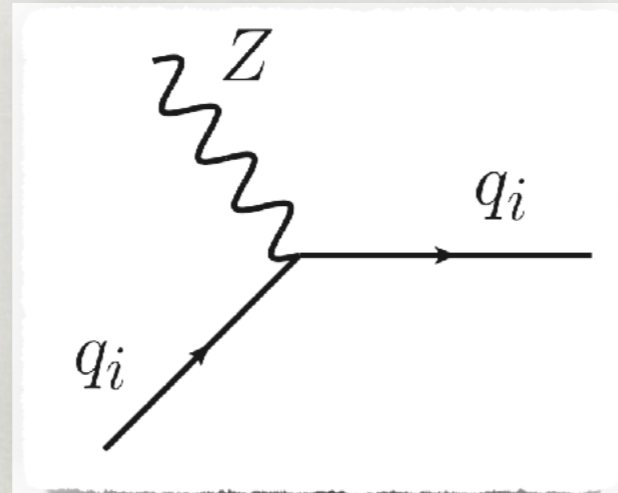
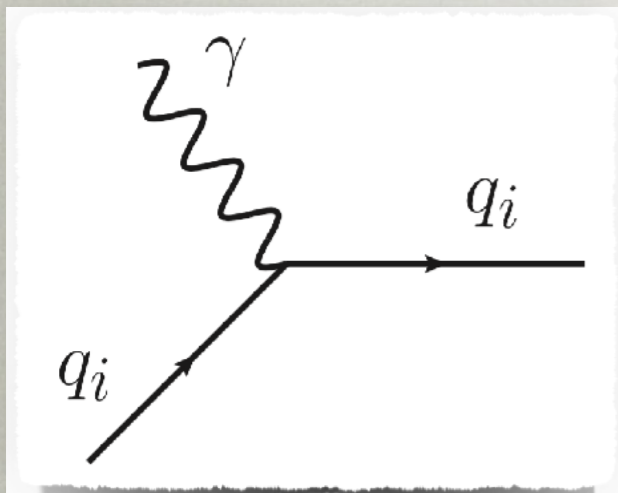
@MUON COLLIDER

JURE ZUPAN
U. OF CINCINNATI

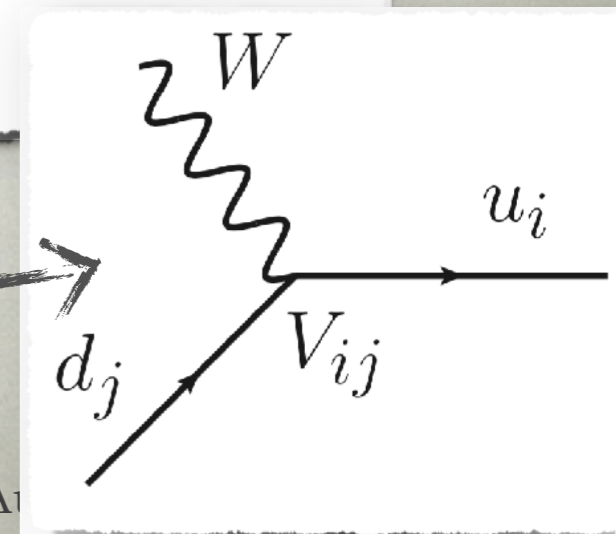
Inaugural US Muon Collider Meeting, Fermilab, Aug 8 2024

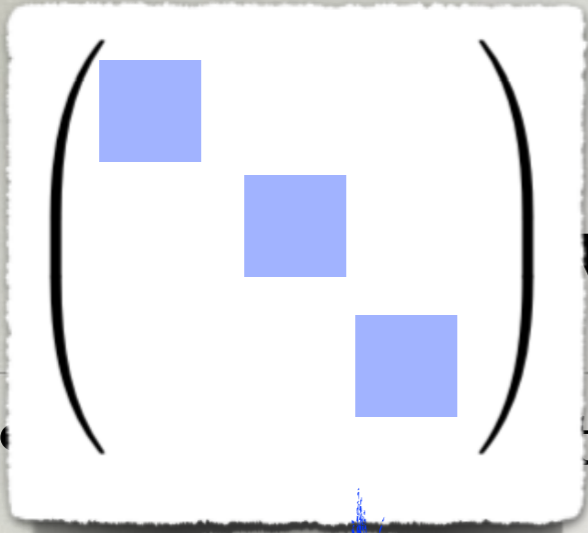
FLAVOR IN THE SM

- neutral currents are flavor conserving (at tree level)
 - photon, gluon, Z: have *flavor (generation) universal* interactions



- Higgs has *flavor diagonal* interactions
 - proportional to quark mass
- charged currents are *flavor changing*
 - W couplings are flavor changing

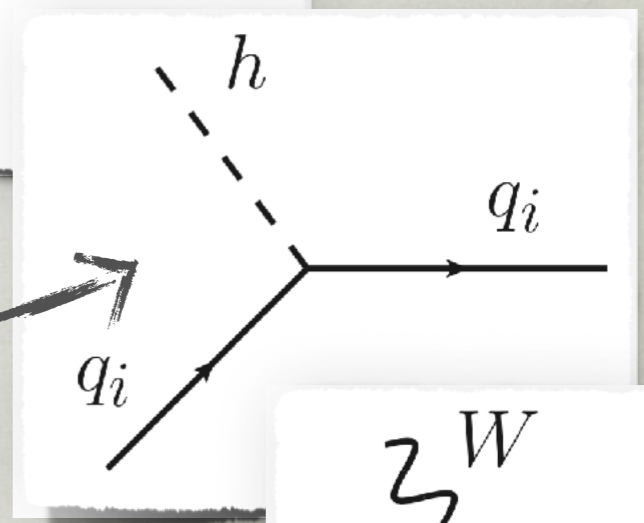
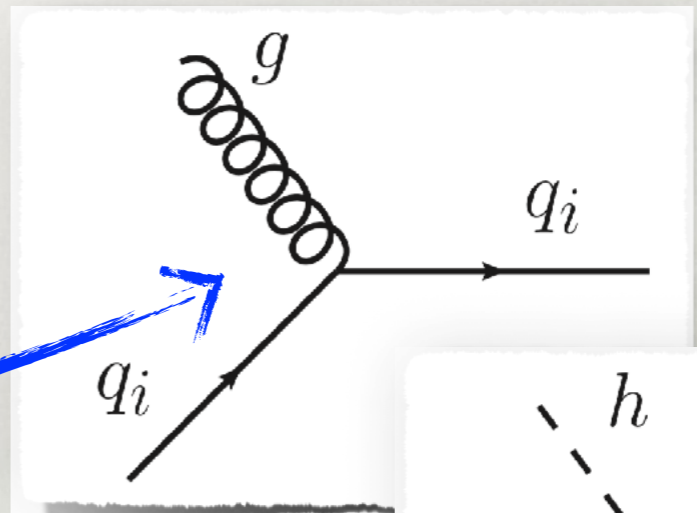
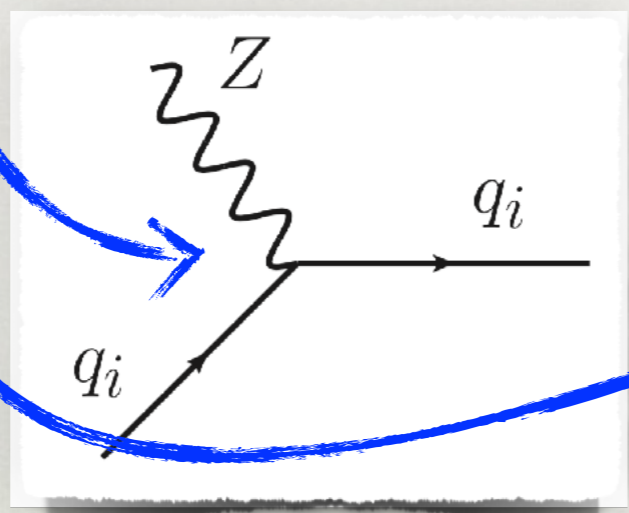
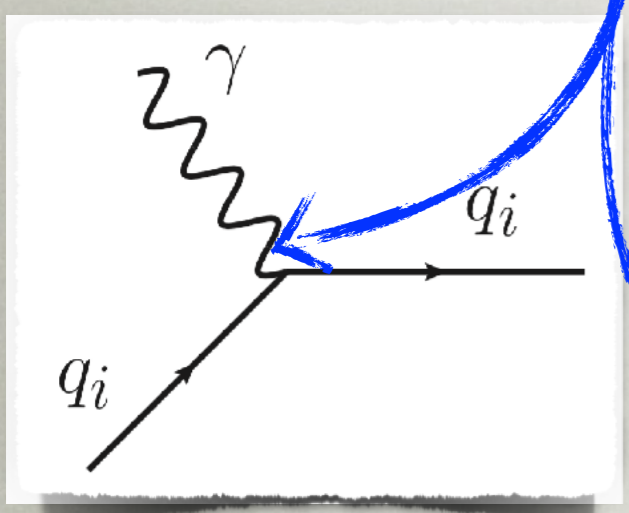




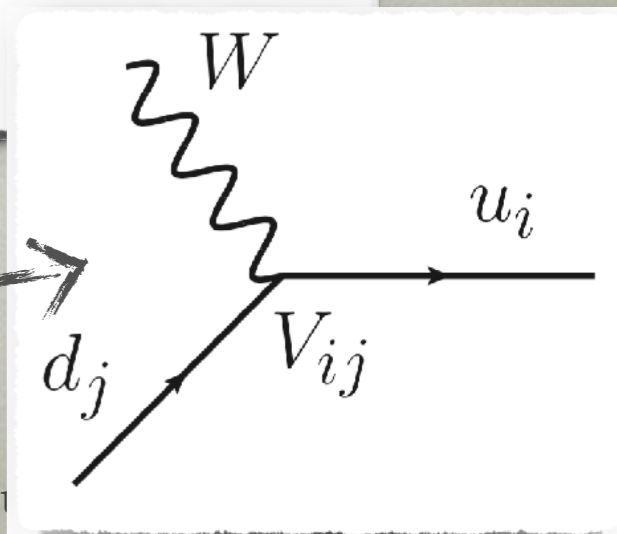
FLAVOR IN THE SM

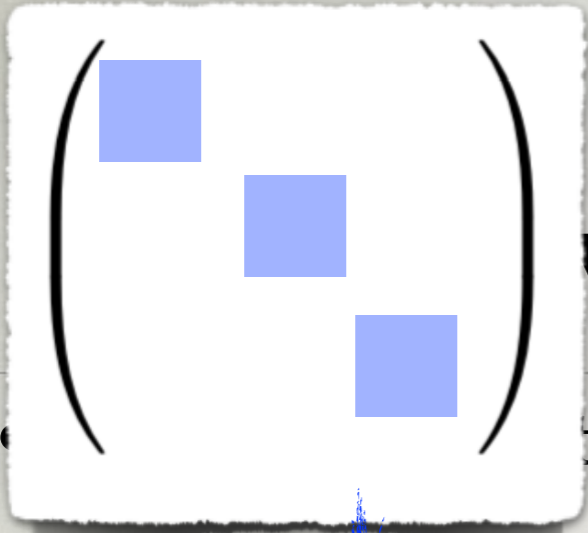
• γ, Z, g are flavor conserving (at tree level)

- photon, gluon, Z: have *flavor (generation) universal* interactions

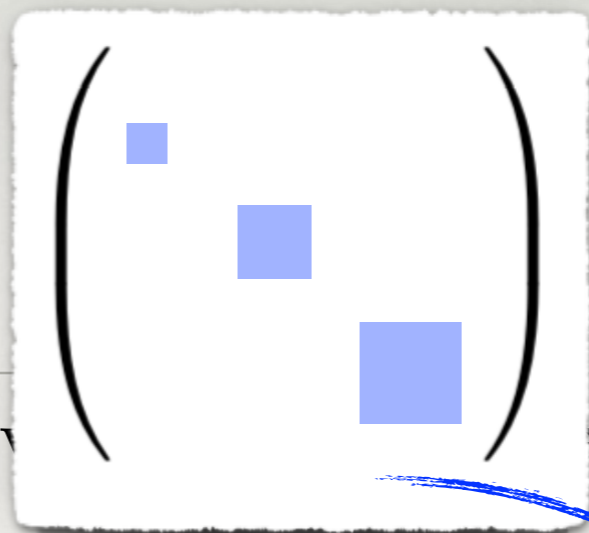


- Higgs has *flavor diagonal* interactions
 - proportional to quark mass
- charged currents are *flavor changing*
 - W couplings are flavor changing





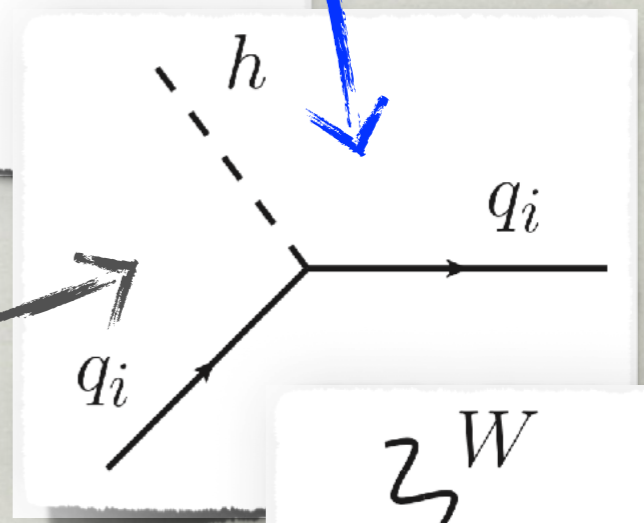
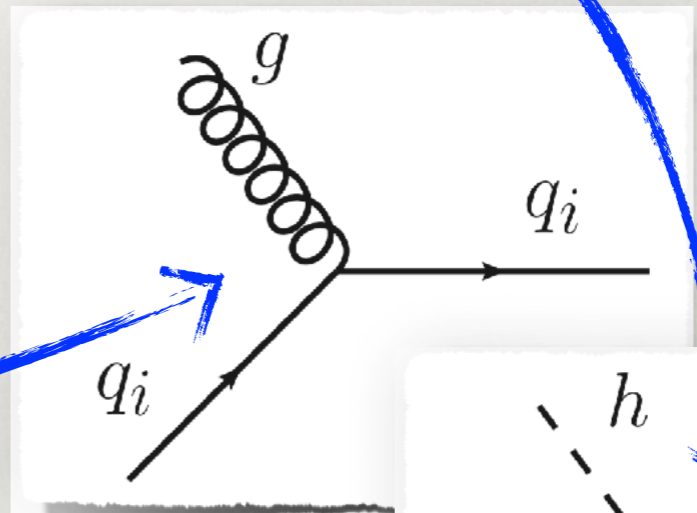
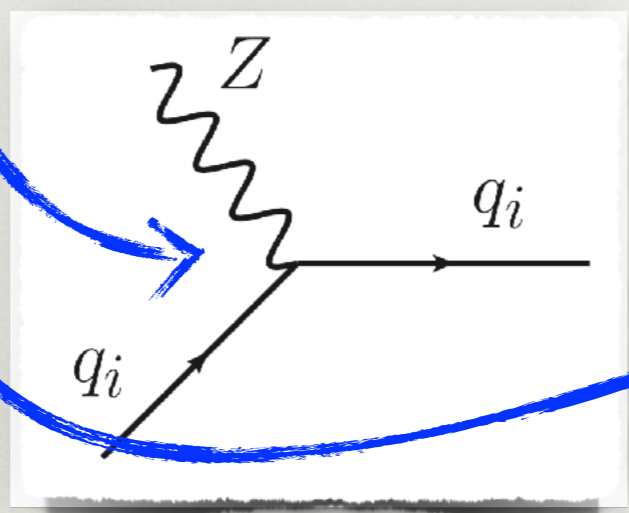
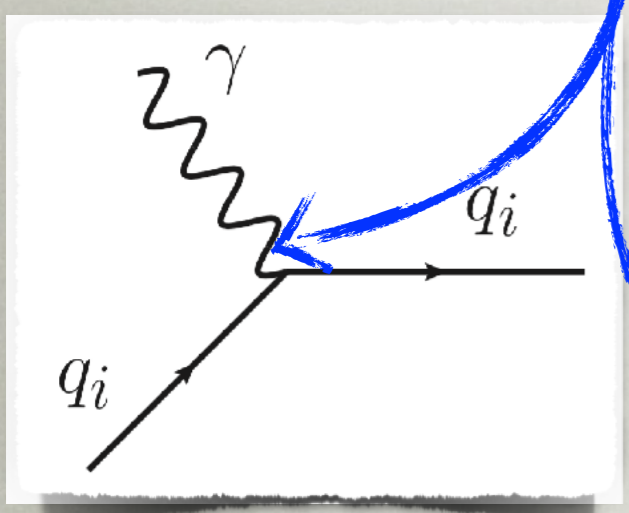
VOR



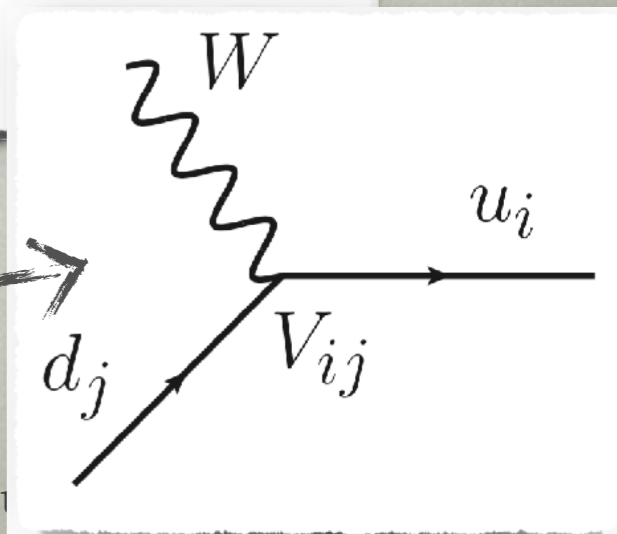
SM

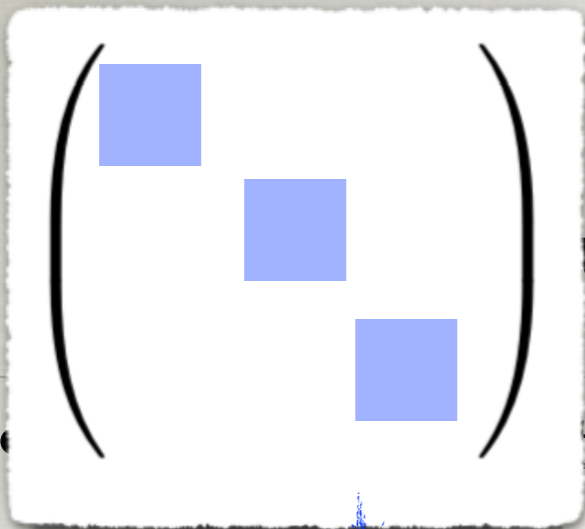
(...s are flavor diagonal (at tree level))

- photon, gluon, Z: have *flavor (generation) universal* interactions



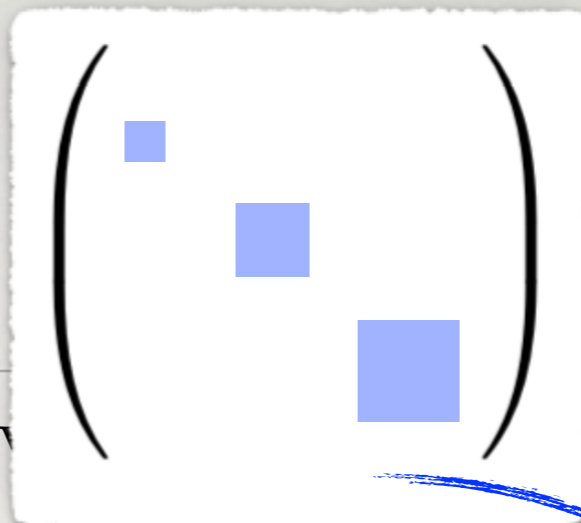
- Higgs has *flavor diagonal* interactions
 - proportional to quark mass
- charged currents are *flavor changing*
 - W couplings are flavor changing





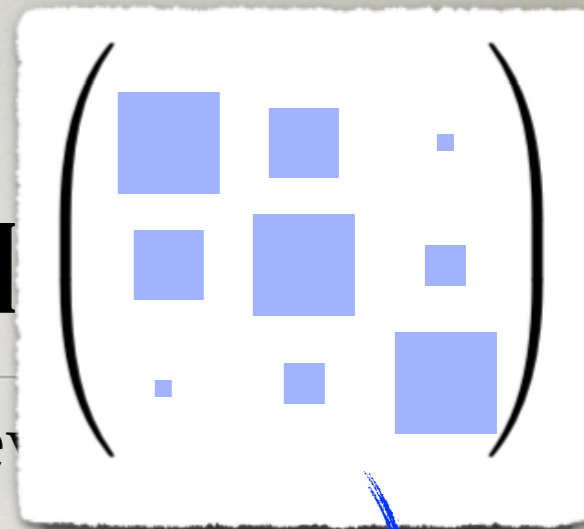
VOR

ts are flav

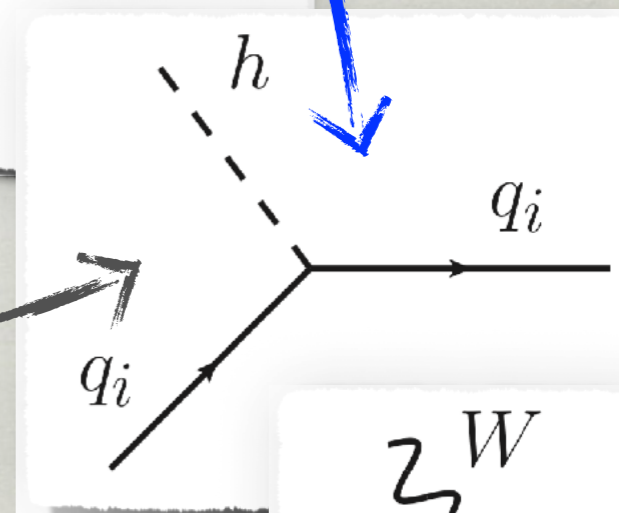
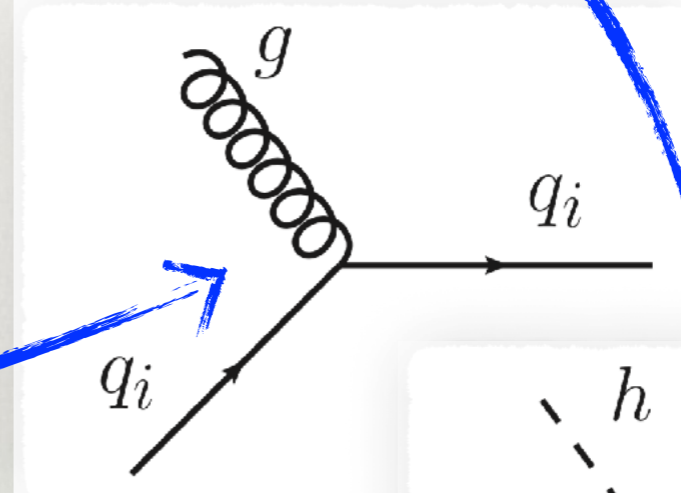
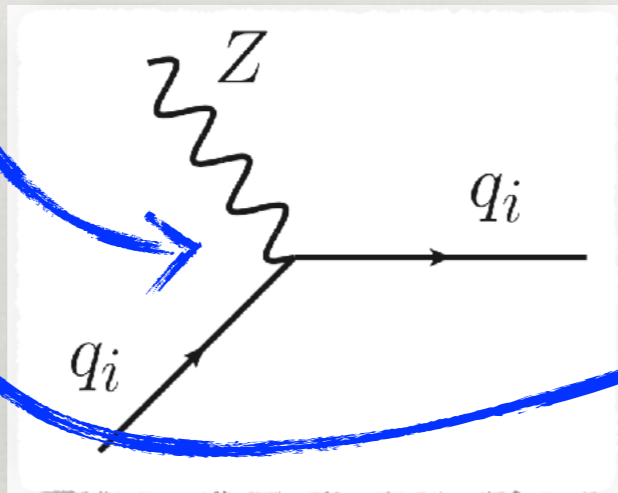
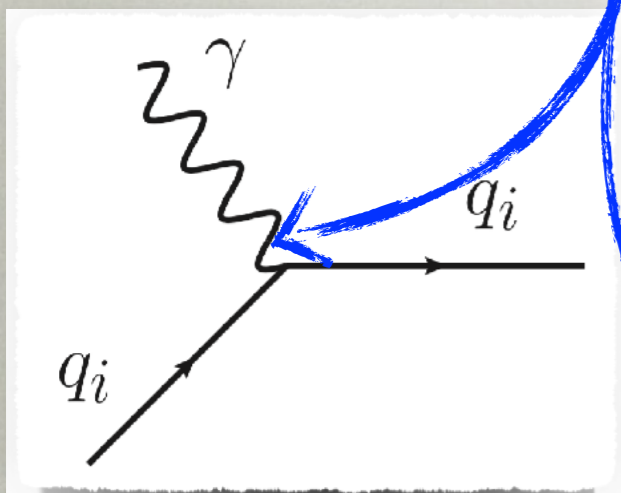


SM

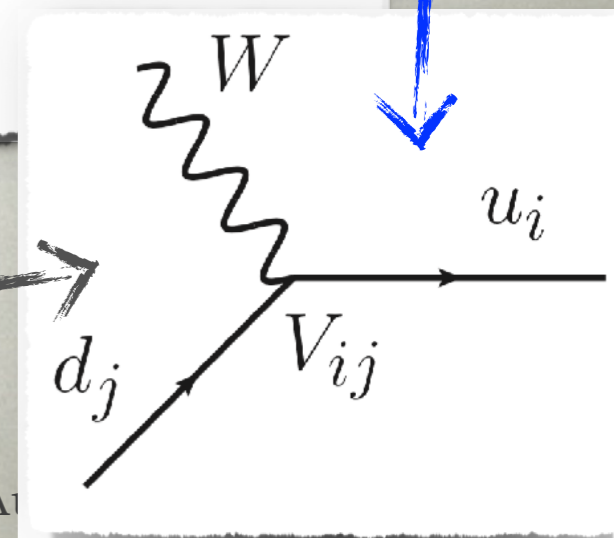
at tree lev



- photon, gluon, Z: have *flavor (generation) universal* interactions

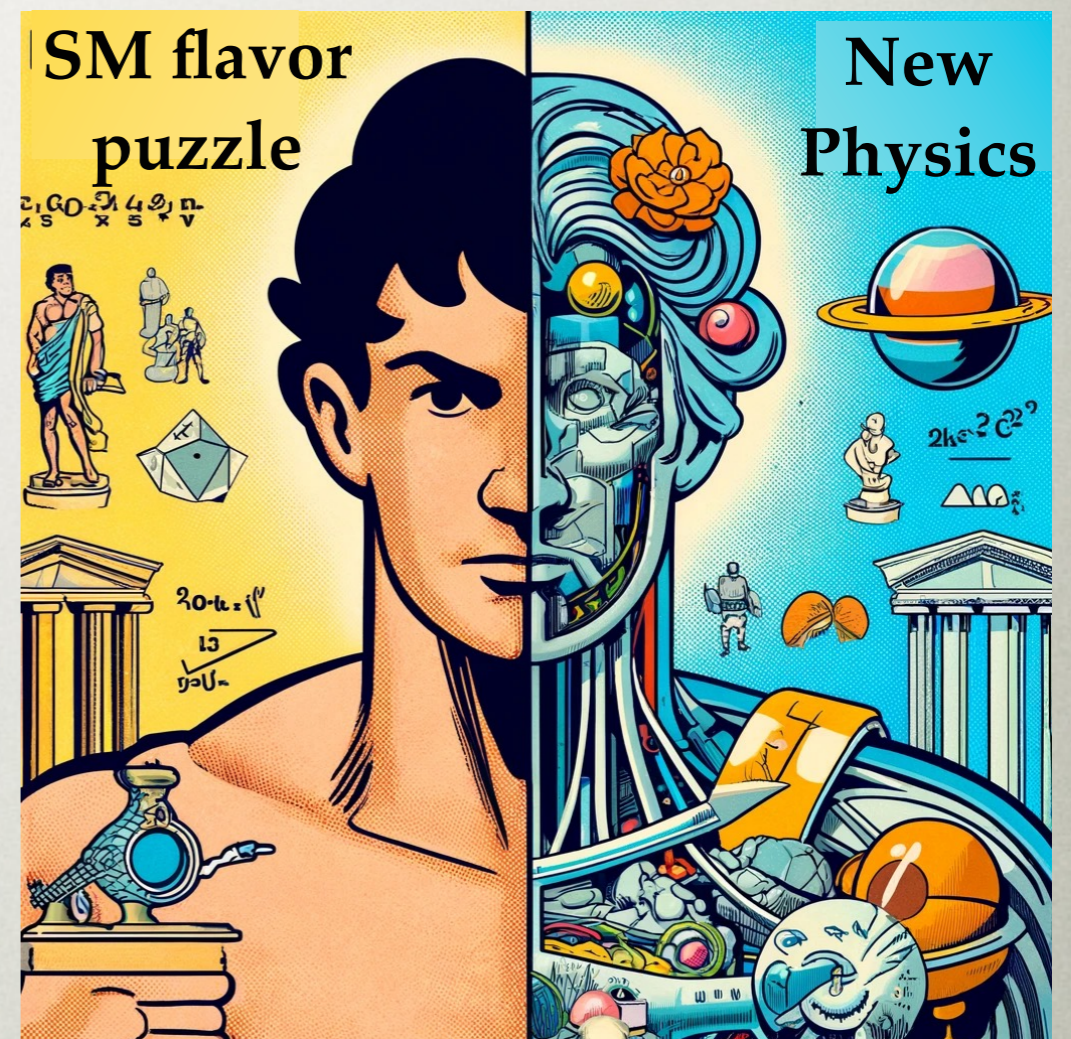


- Higgs has *flavor diagonal* interactions
 - proportional to quark mass
- charged currents are *flavor changing*
 - W couplings are flavor changing



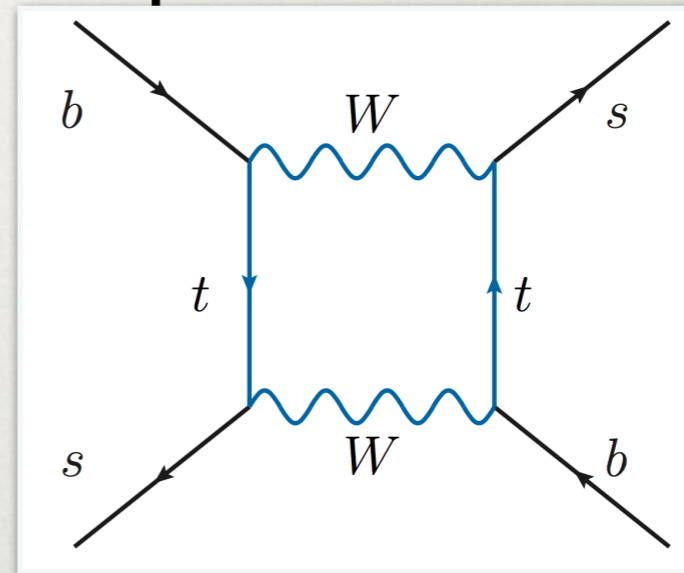
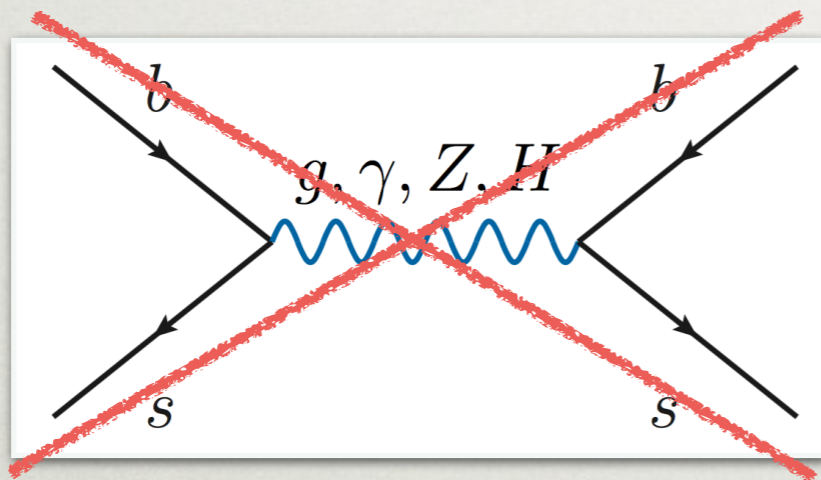
TWO FACES OF FLAVOR PHYSICS

- no flavor changing neutral currents in the SM
 - \Rightarrow flavor transitions sensitive probes of new physics
- why the observed structure of quark and lepton masses and mixings?
 - \Rightarrow flavor model building

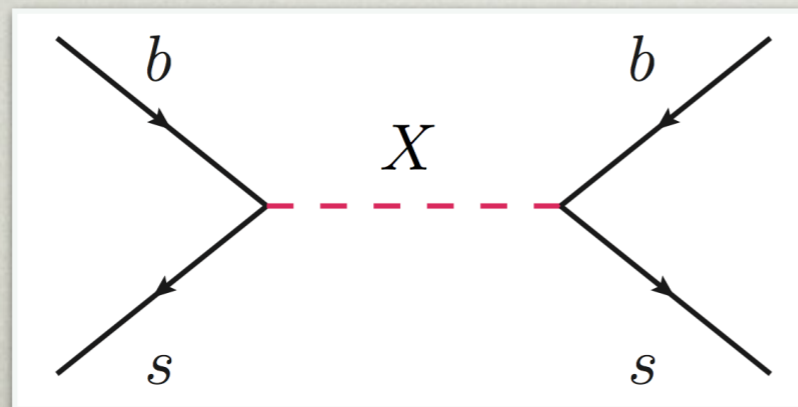


SEARCHING FOR OFF-SHELL NEW PHYSICS

- FCNC processes only at loop level in the SM



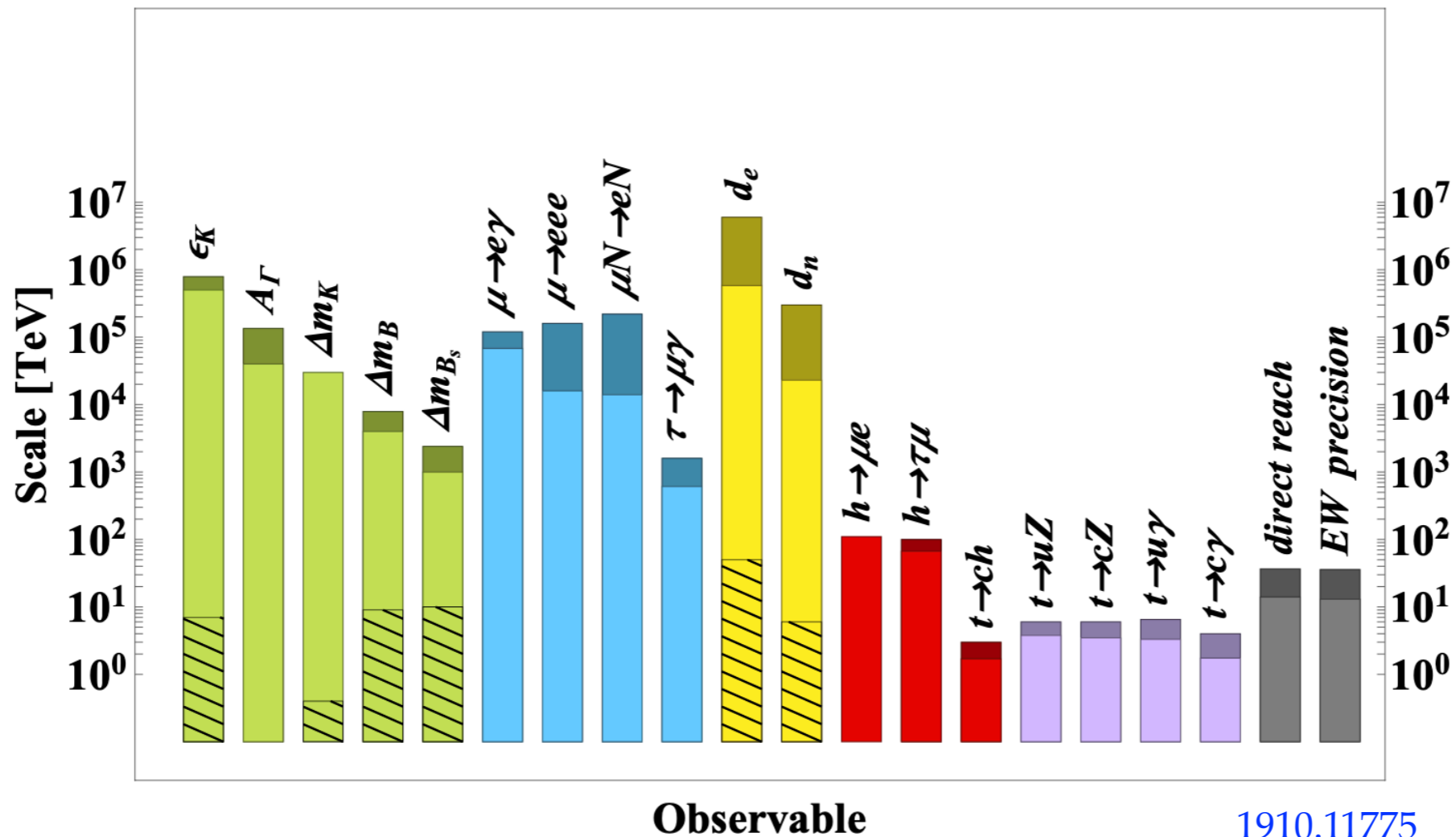
- can search for off-shell new physics



$$\delta C^{\text{NP}} \propto \frac{g_{sb}^2}{M_{\text{NP}}^2}$$

HEAVY NEW PHYSICS

- compare exp. and SM prediction
 - does it agree? \Rightarrow place bounds
 - for $g_{\text{NP}} \sim \mathcal{O}(1) \Rightarrow$ probe high scales



FLAVOR @ MUON COLLIDER

- searching for NP through precision probes
 - off-shell NP coupling to muons: $\mu^+ \mu^- \rightarrow \tau \mu$
 - Higgs as a probe of flavor
 - $h \rightarrow \tau \mu, \tau e, \mu e$
 - $h \rightarrow bs, cu$
- production of new states
 - flavor structure of these \Rightarrow learn about the solution to the SM flavor puzzle?

PRECISION FLAVOR
PROBES @MUC

$\tau \rightarrow 3\mu$

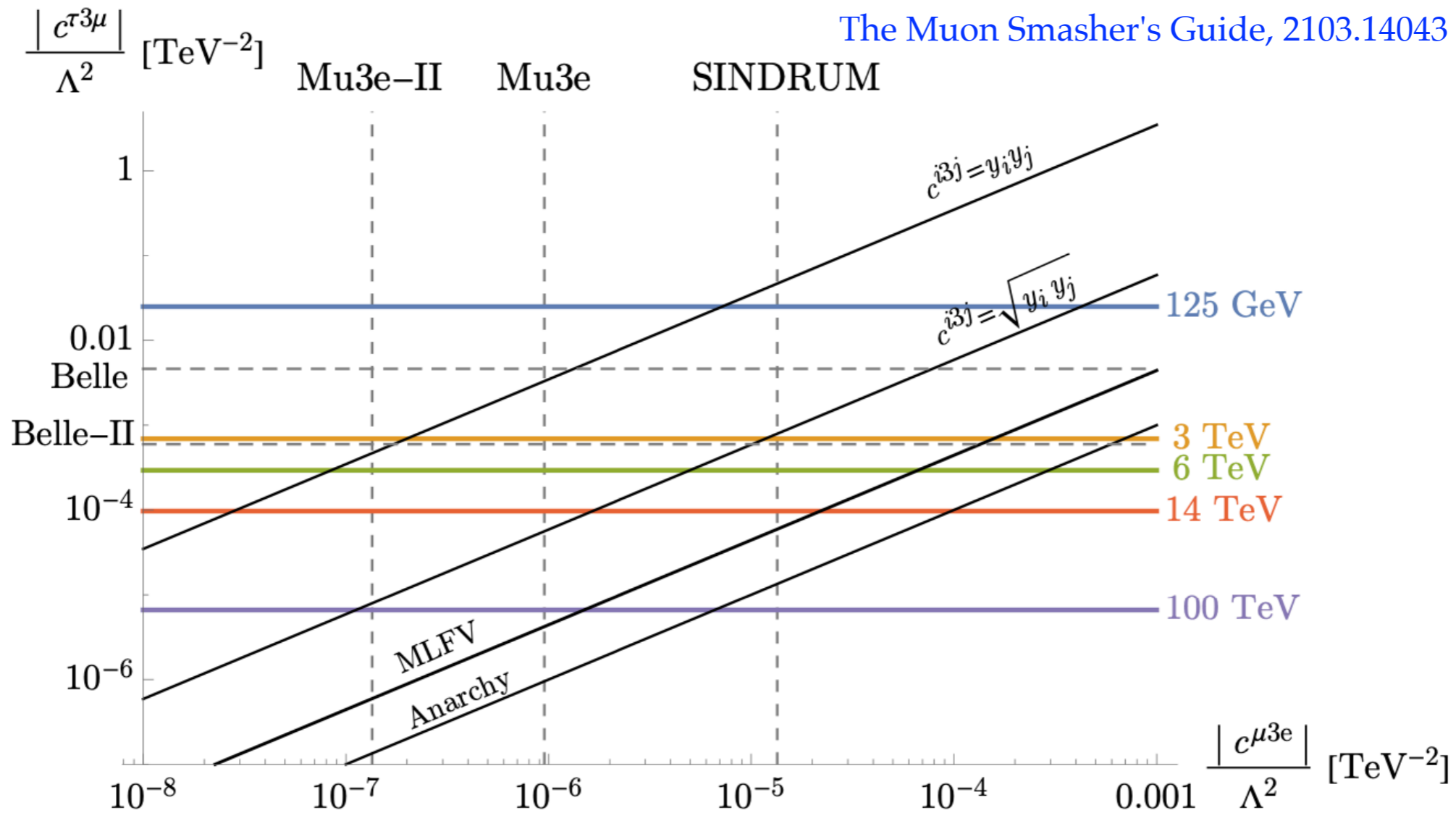
- assume NP is heavy \Rightarrow SMEFT
- focus on $\tau \rightarrow 3\mu$ couplings
 - dimension 6 operators

The Muon Smasher's Guide, 2103.14043

$$\mathcal{L} \supset V_{LL}^{\tau 3\mu} (\bar{\mu} \gamma^\mu P_L \mu) (\bar{\tau} \gamma_\mu P_L \mu) + V_{LR}^{\tau 3\mu} (\bar{\mu} \gamma^\mu P_L \mu) (\bar{\tau} \gamma_\mu P_R \mu) + (L \leftrightarrow R) + \text{h.c.},$$

$$V_{LL}^{\tau 3\mu} = V_{LR}^{\tau 3\mu} = V_{RL}^{\tau 3\mu} = V_{RR}^{\tau 3\mu} = \frac{c^{\tau 3\mu}}{\Lambda^2}$$

- indirect probe from (Belle)
 $Br(\tau \rightarrow 3\mu) < 2.1 \times 10^{-8} \Rightarrow \Lambda \gtrsim 15 \text{ TeV}$
- gain for MuC: $\sigma(\mu\mu \rightarrow \tau\mu) \propto s/\Lambda^4$
 - can probe indirectly much higher scales than Belle II



$\mathcal{L} \supset$

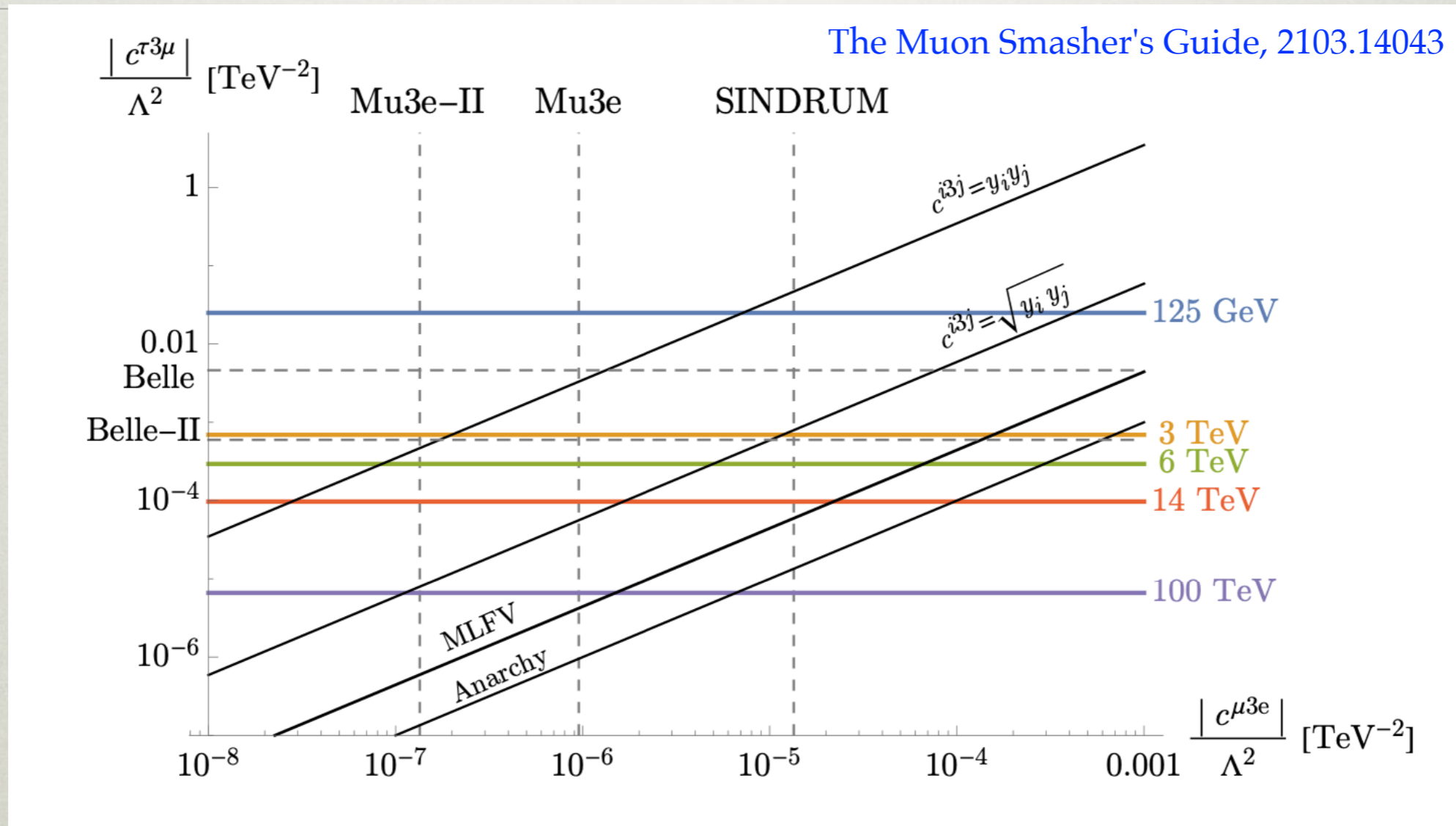
$$v_{LL} = v_{LR} = v_{RL} = v_{RR} = v_{\tau 3\mu} = \frac{c^{\tau 3\mu}}{\Lambda^2}$$

- indirect probe from (Belle)

$$Br(\tau \rightarrow 3\mu) < 2.1 \times 10^{-8} \Rightarrow \Lambda \gtrsim 15 \text{ TeV}$$
- gain for MuC: $\sigma(\mu\mu \rightarrow \tau\mu) \propto s/\Lambda^4$
 - can probe indirectly much higher scales than Belle II

$$\tau \rightarrow 3\mu$$

The Muon Smasher's Guide, 2103.14043



- for $\mu\mu \rightarrow \mu e$ bounds from $\mu \rightarrow 3e$ much more stringent ($\Lambda \gtrsim 270$ TeV)
 - relation between $\tau \rightarrow 3\mu$ and $\mu \rightarrow 3e$ model dependent
- similar analysis for $\mu\mu \rightarrow \tau e$

$$\mu\mu \rightarrow bs$$

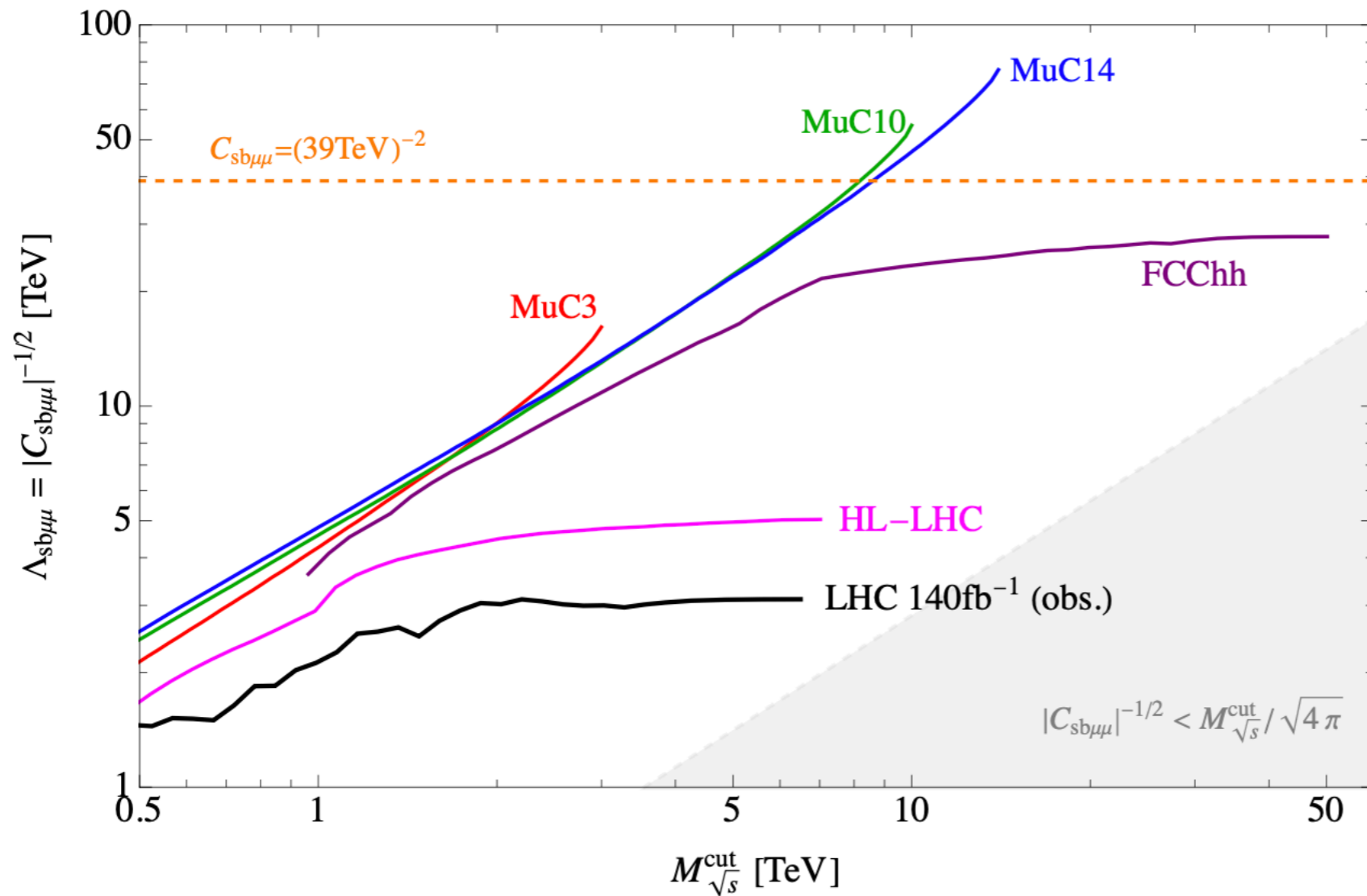
Azatov et al, 2205.13552

- the $\mu^+\mu^- \rightarrow \bar{b}s + \bar{s}b$ a probe of $B \rightarrow K^{(*)}\mu\mu$ anomalies

- assume that dim-6 EFT operator

$$(\bar{s}_L \gamma_\alpha b_L)(\bar{\mu}_L \gamma^\alpha \mu_L)$$

- MuC could probe interesting parameter space
- other flavors could be explored as well
 $\mu\mu \rightarrow tc, tu, cu, bd, (sd?)$



Azatov et al, 2205.13552

$\rightarrow K^{(*)}\mu\mu$

or

- MuC could probe interesting parameter space
- other flavors could be explored as well
 $\mu\mu \rightarrow tc, tu, cu, bd, (sd?)$

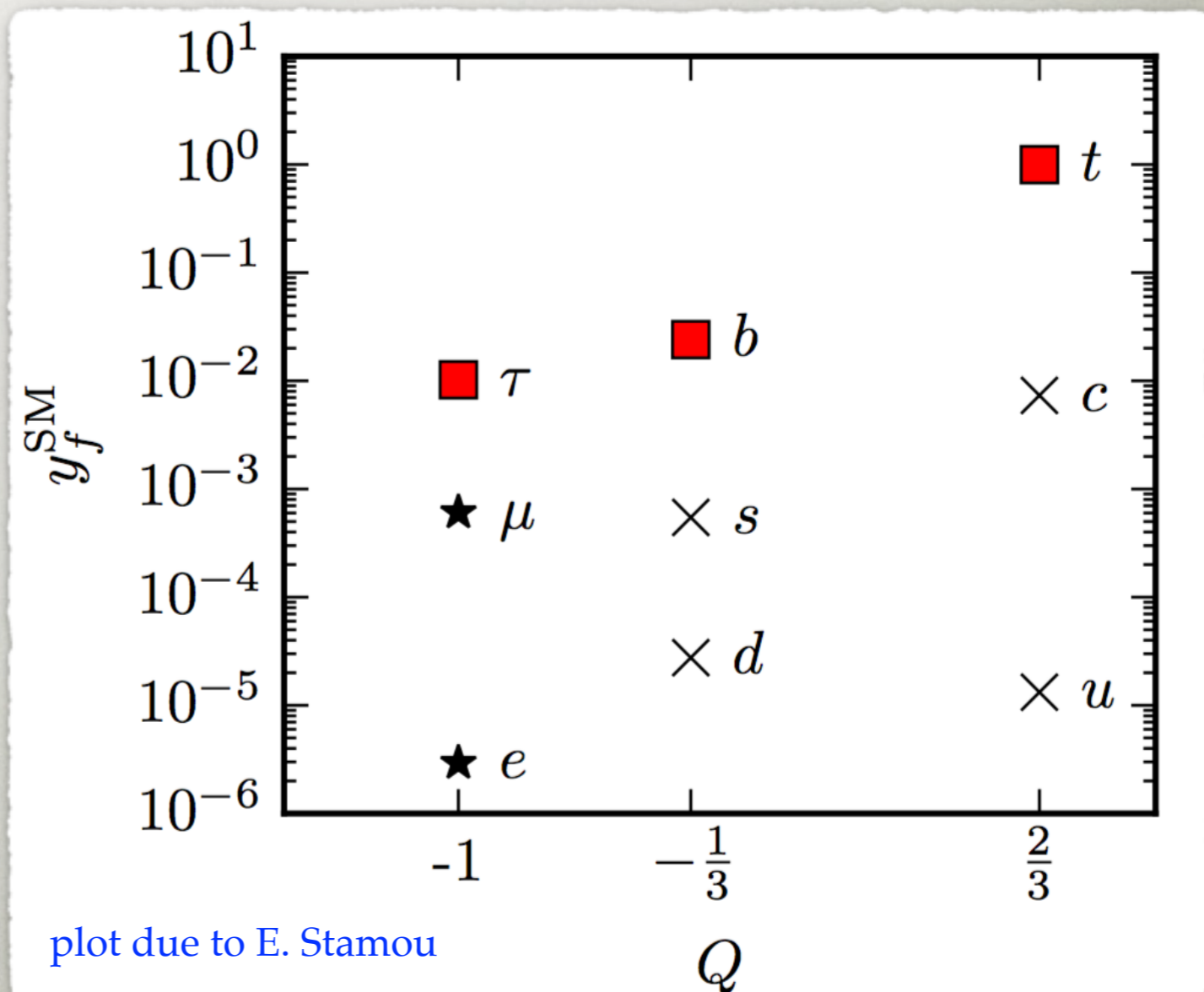
HIGGS AS A PROBE OF FLAVOR

HIGGS = NONTRIVIAL FLAVOR STRUCTURE

- generation of masses in the SM through the Higgs mechanism
- implies Higgs has hierarchical couplings to fermions
- in the SM

$$y_f = \sqrt{2}m_f/v$$

- we want to test this



TESTING THE FLAVOR OF THE HIGGS

Nir, 1605.00433; JZ 1903.05062

- several questions

- proportionality

$$y_{ii} \propto m_i$$

- factor of proportionality

$$y_{ii}/m_i = \sqrt{2}/v$$

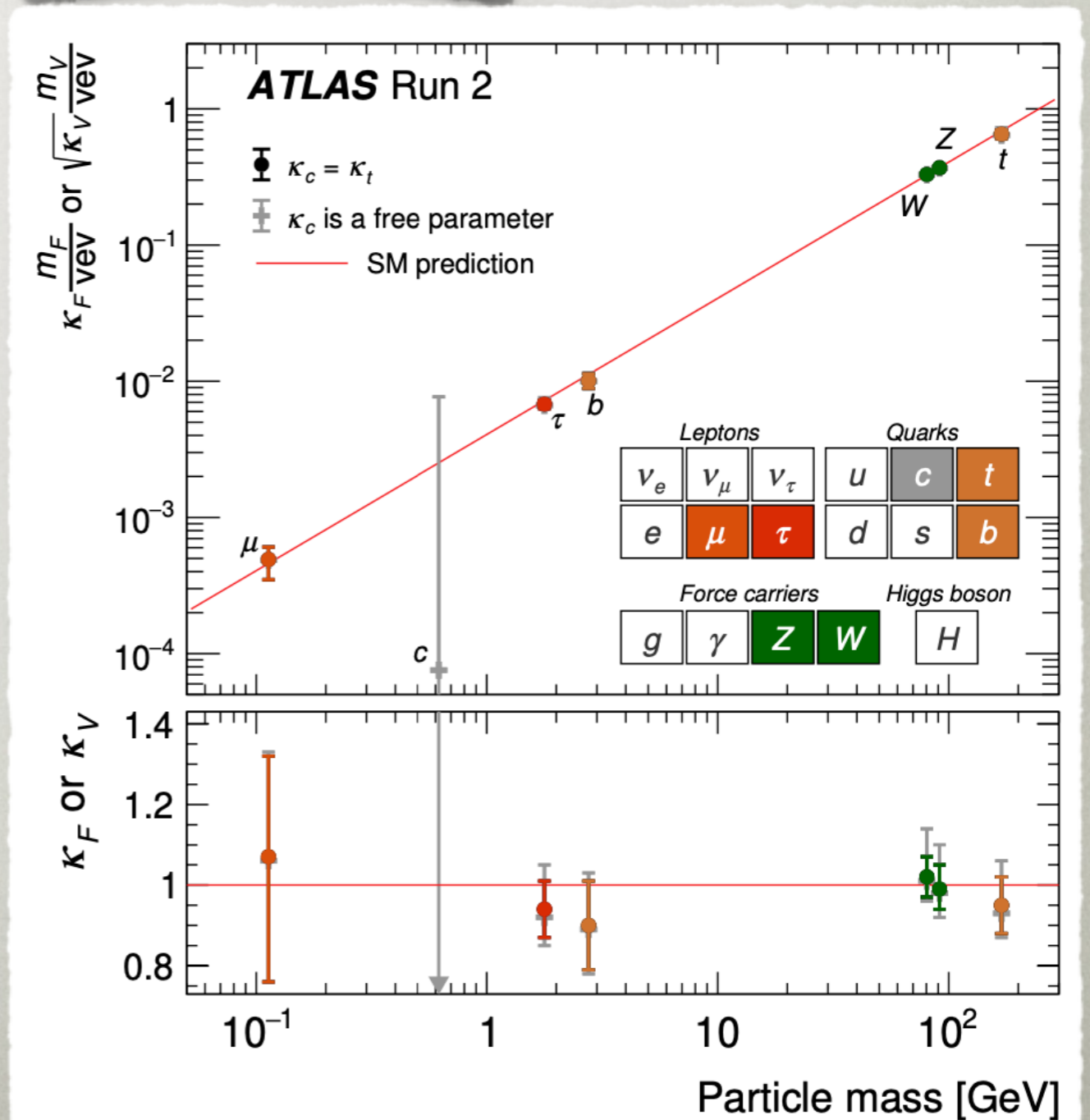
- diagonality

$$y_{ij} = 0, \quad i \neq j$$

- reality

$$\text{Im}(y_{ij}) = 0$$

$$y_f^{\text{SM}} = \sqrt{2}m_f/v$$



TESTING THE FLAVOR OF THE HIGGS

Nir, 1605.00433; JZ 1903.05062

- several questions

- proportionality

$$y_{ii} \propto m_i$$

- factor of proportionality

$$y_{ii}/m_i = \sqrt{2}/v$$

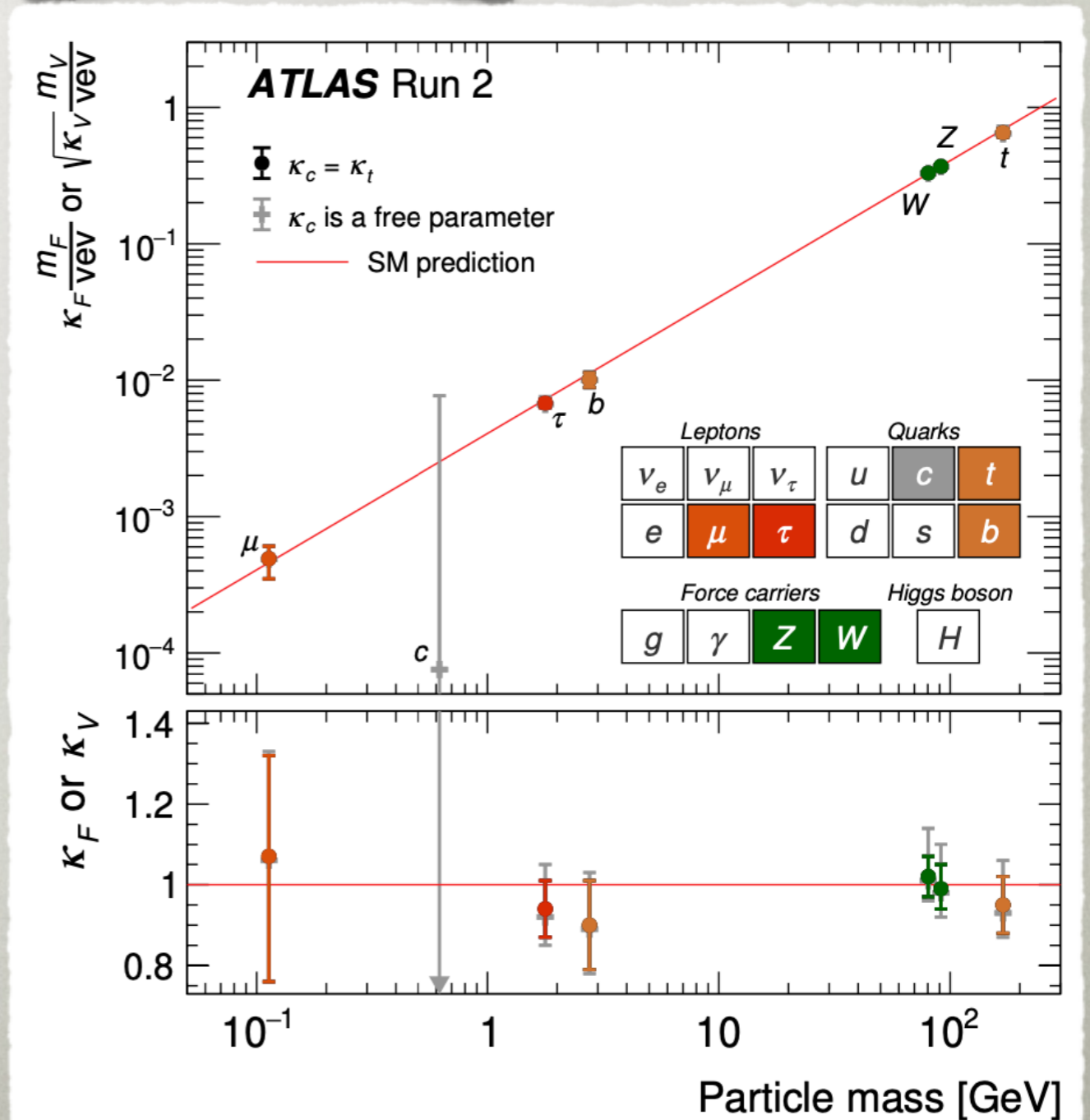
- diagonality

$$y_{ij} = 0, \quad i \neq j$$

- reality

$$\text{Im}(y_{ij}) = 0$$

$$y_f^{\text{SM}} = \sqrt{2}m_f/v$$



NEW PHYSICS CORRECTIONS TO HIGGS COUPLINGS

- in SMEFT, the Yukawas get corrected by higher dim. ops

$$\mathcal{L}_{SM} = - [\lambda_{ij} (\bar{f}_L^i f_R^j) H + h.c.]$$

$$\Delta\mathcal{L}_Y = -\frac{\lambda'_{ij}}{\Lambda^2} (\bar{f}_L^i f_R^j) H (H^\dagger H) + h.c. + \dots$$

- there could be other sources of EWSB
- NP in general misaligns mass and Yukawa matrices

$$\mathcal{L}_Y = -m_{ij} \bar{f}_L^i f_R^j - Y_{ij} (\bar{f}_L^i f_R^j) h + h.c. + \dots$$

$$m = \left(\lambda + \frac{v^2}{2\Lambda^2} \lambda' \right) \frac{v}{\sqrt{2}}$$

$$Y = \left(\lambda + 3 \frac{v^2}{2\Lambda^2} \lambda' \right) \frac{1}{\sqrt{2}}$$

$$m = v \begin{pmatrix} * + *' & * + *' & * + *' \\ * + *' & * + *' & * + *' \\ * + *' & * + *' & * + *' \end{pmatrix}$$

$$Y = \begin{pmatrix} * + *' & * + *' & * + *' \\ * + *' & * + *' & * + *' \\ * + *' & * + *' & * + *' \end{pmatrix}$$

NEW PHYSICS CORRECTIONS TO HIGGS COUPLINGS

- in SMEFT, the Yukawas get corrected by higher dim. ops

$$\mathcal{L}_{SM} = - [\lambda_{ij} (\bar{f}_L^i f_R^j) H + h.c.]$$

$$\Delta\mathcal{L}_Y = -\frac{\lambda'_{ij}}{\Lambda^2} (\bar{f}_L^i f_R^j) H (H^\dagger H) + h.c. + \dots$$

- there could be other sources of EWSB
- NP in general misaligns mass and Yukawa matrices

$$\mathcal{L}_Y = -m_{ij} \bar{f}_L^i f_R^j - Y_{ij} (\bar{f}_L^i f_R^j) h + h.c. + \dots$$

$$m = \left(\lambda + \frac{v^2}{2\Lambda^2} \lambda' \right) \frac{v}{\sqrt{2}}$$

$$Y = \left(\lambda + 3 \frac{v^2}{2\Lambda^2} \lambda' \right) \frac{1}{\sqrt{2}}$$

$$m = v \begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$$Y = \begin{pmatrix} * & * & * \\ * & * & * \\ * & * & * \end{pmatrix}$$

CPV AND FV HIGGS COUPLINGS TO SM FERMIONS

$$m = v \begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$$Y = \begin{pmatrix} * & * & * \\ * & * & * \\ * & * & * \end{pmatrix}$$

- NP contris in general lead to
 - flavor violating Higgs decays
 - CPV Higgs decays
- different models lead to different patterns of flavor diagonal and flavor violating Yukawas
- note: large Yukawas for light fermions require significant tuning in contributions to fermion masses

Harnik, Kopp, JZ, 1209.1397
Blackenburg, Ellis, Isidori, 1202.5704

OTHER CONSTRAINTS

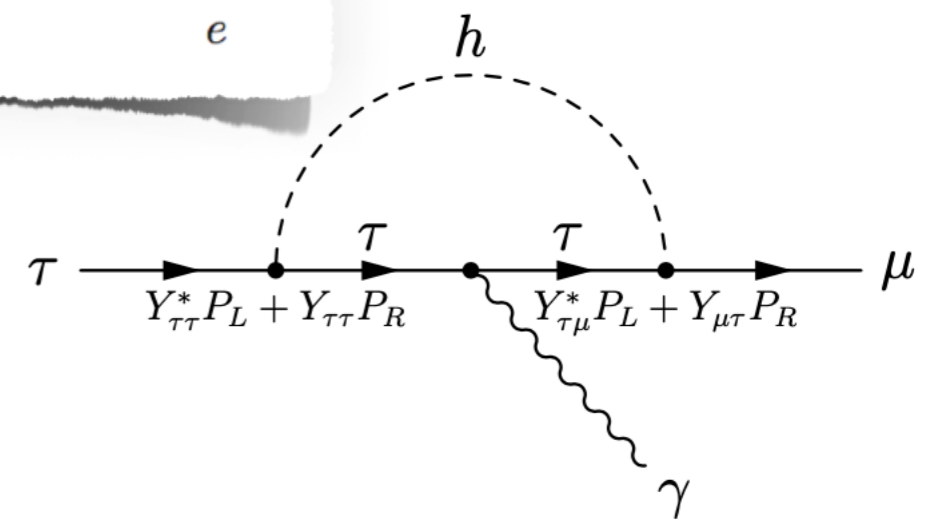
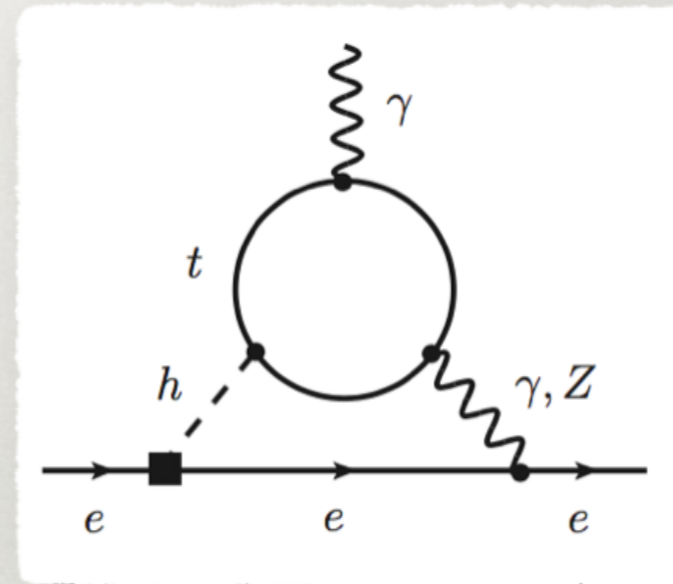
- both CPV and FV Higgs couplings face severe indirect constraints

- CPV from EDM
- FV from low energy FCNC processes

- wide open for HL-LHC & future colliders:

$$h \rightarrow \tau\mu, \tau e \text{ and } t \rightarrow hc, hu$$

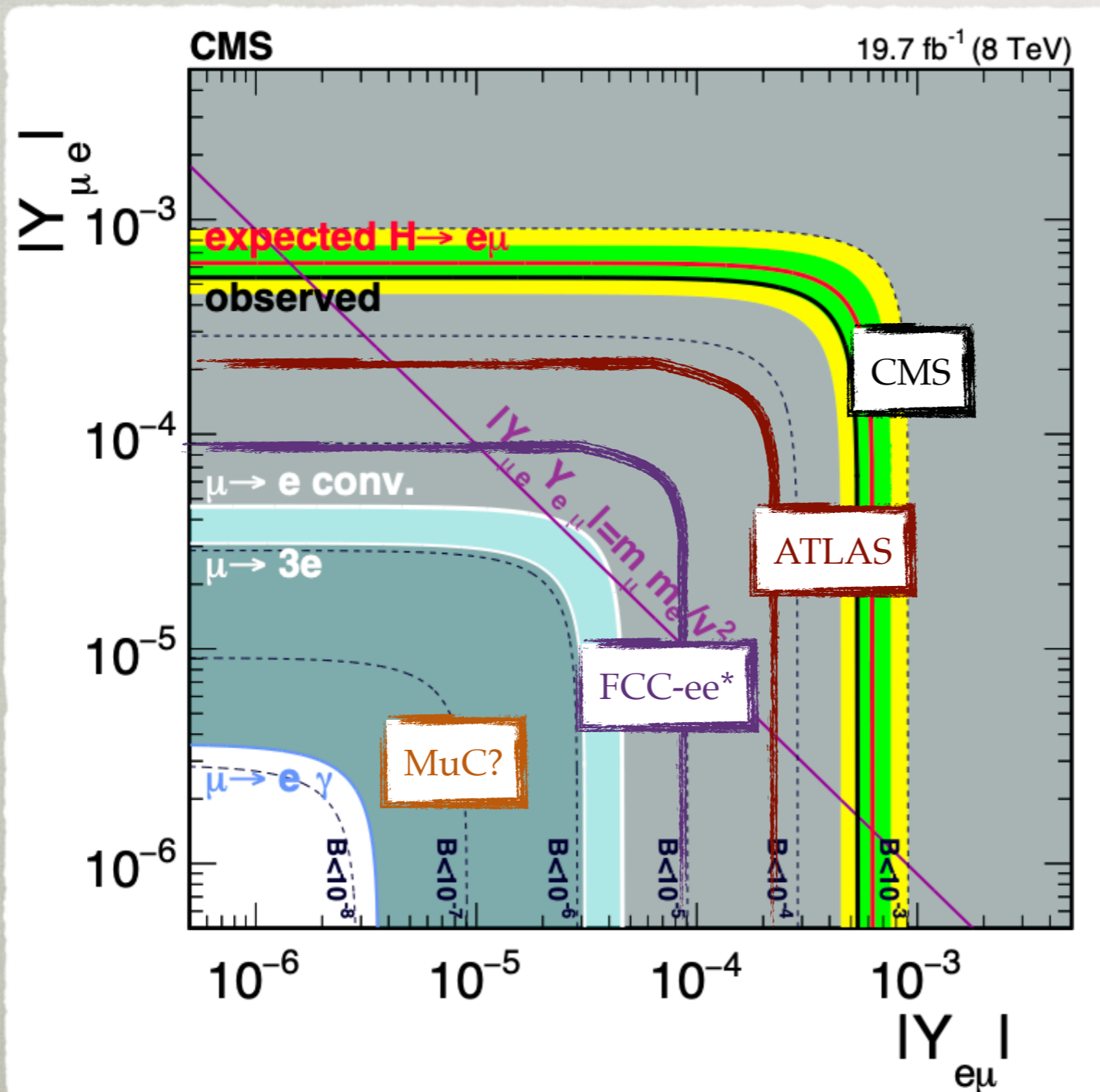
- $h \rightarrow bs, cu$ with some tuning



$Y_{\mu e}$ YUKAWA

ATLAS, 1909.10235

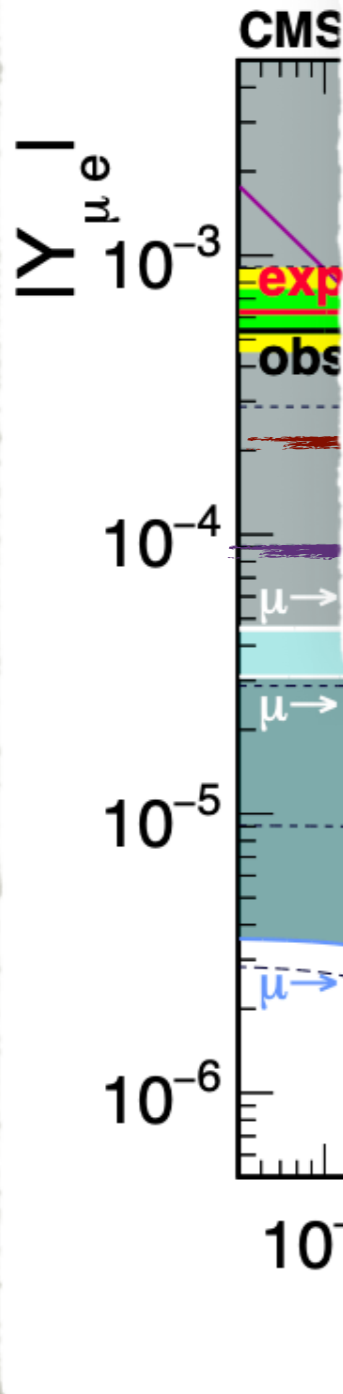
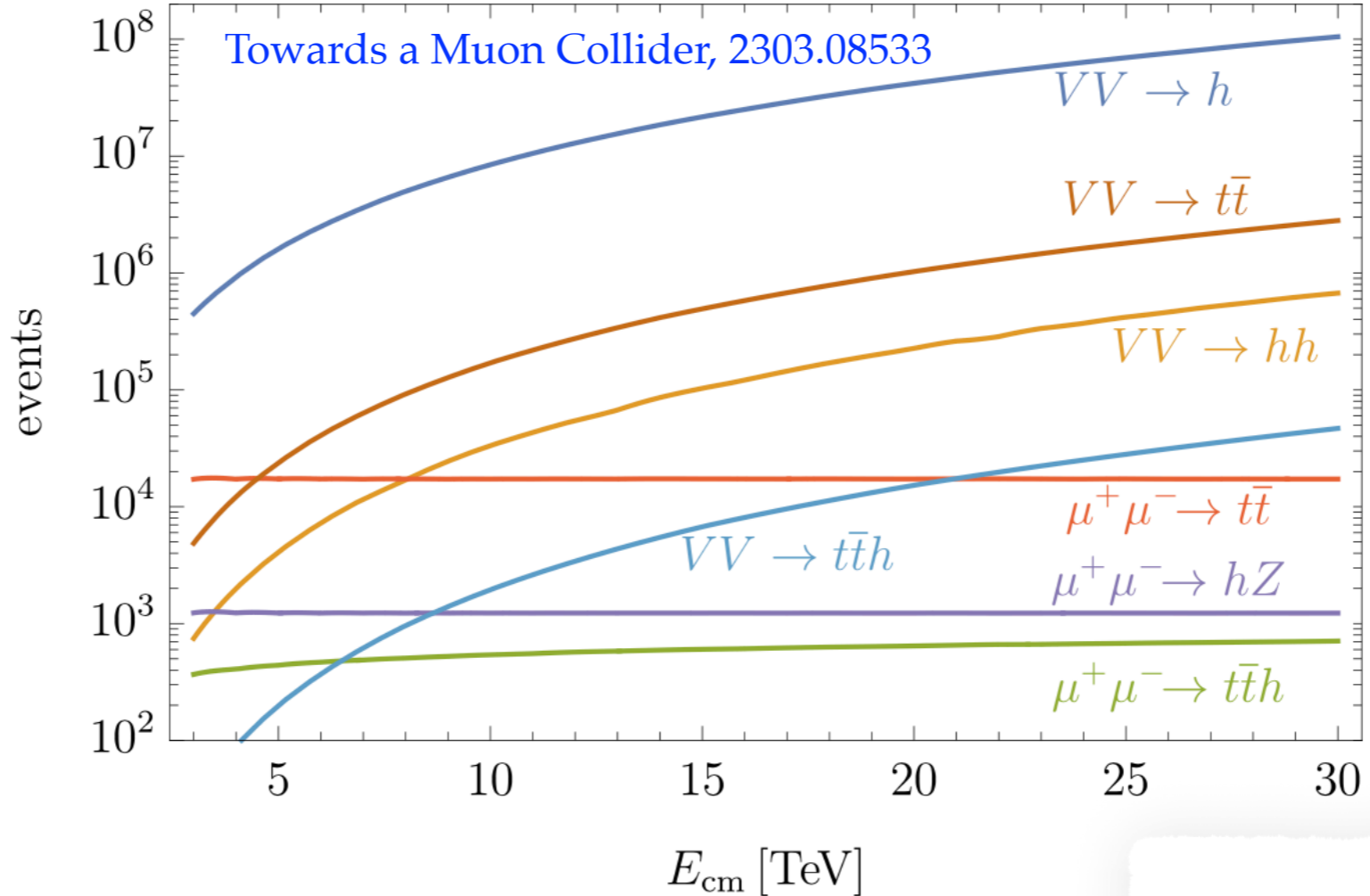
CMS, 1607.03561



*from Qin et al, 1711.07243

ATLAS, 1909.10235

CMS, 1607.03561

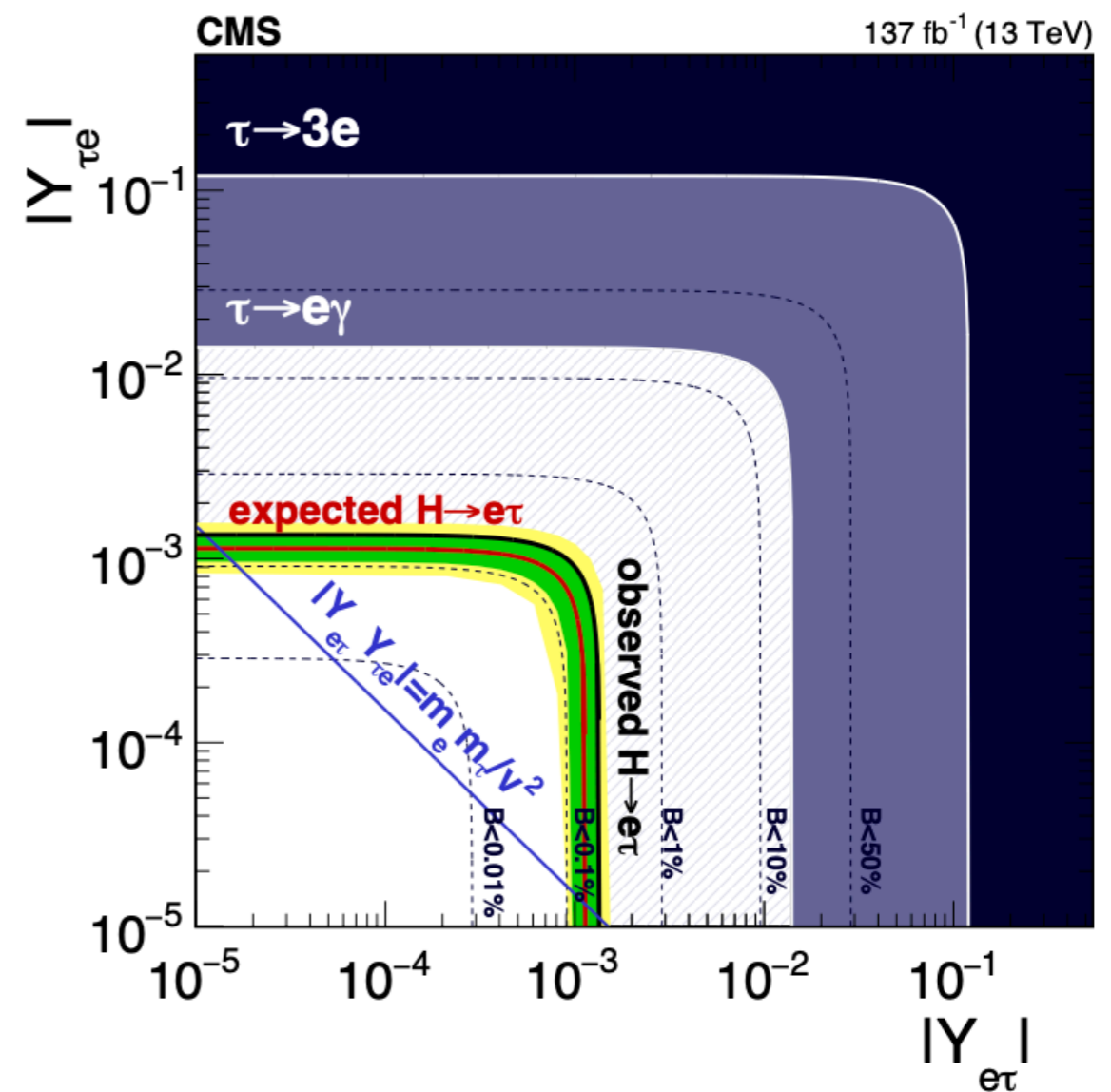
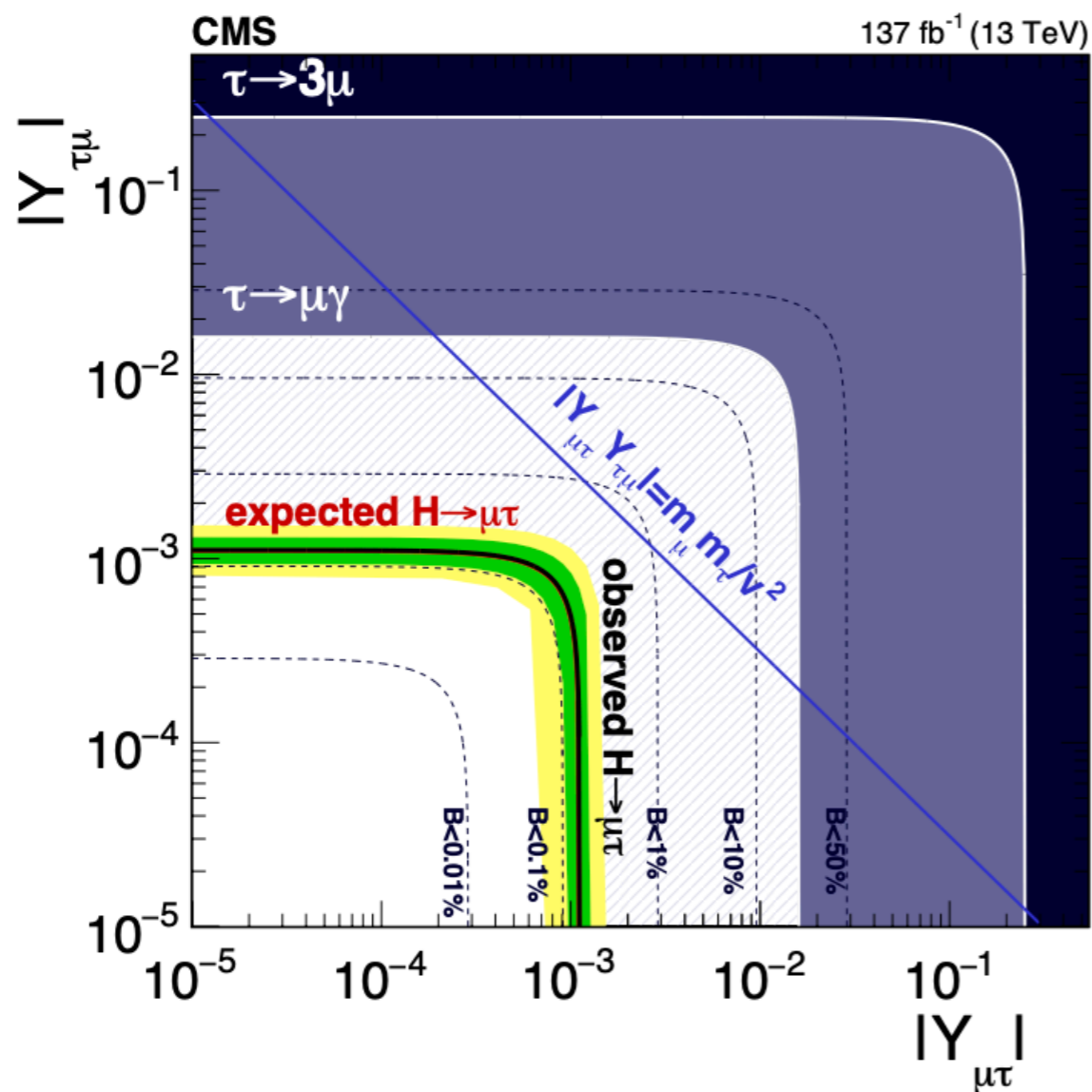


$$\mathcal{L}_{int} = 10 \text{ ab}^{-1} \left(\frac{E_{cm}}{10 \text{ TeV}} \right)^2$$

*from Qin et al, 1711.07243

$Y_{\tau\mu}, Y_{\tau e}$ YUKAWAS

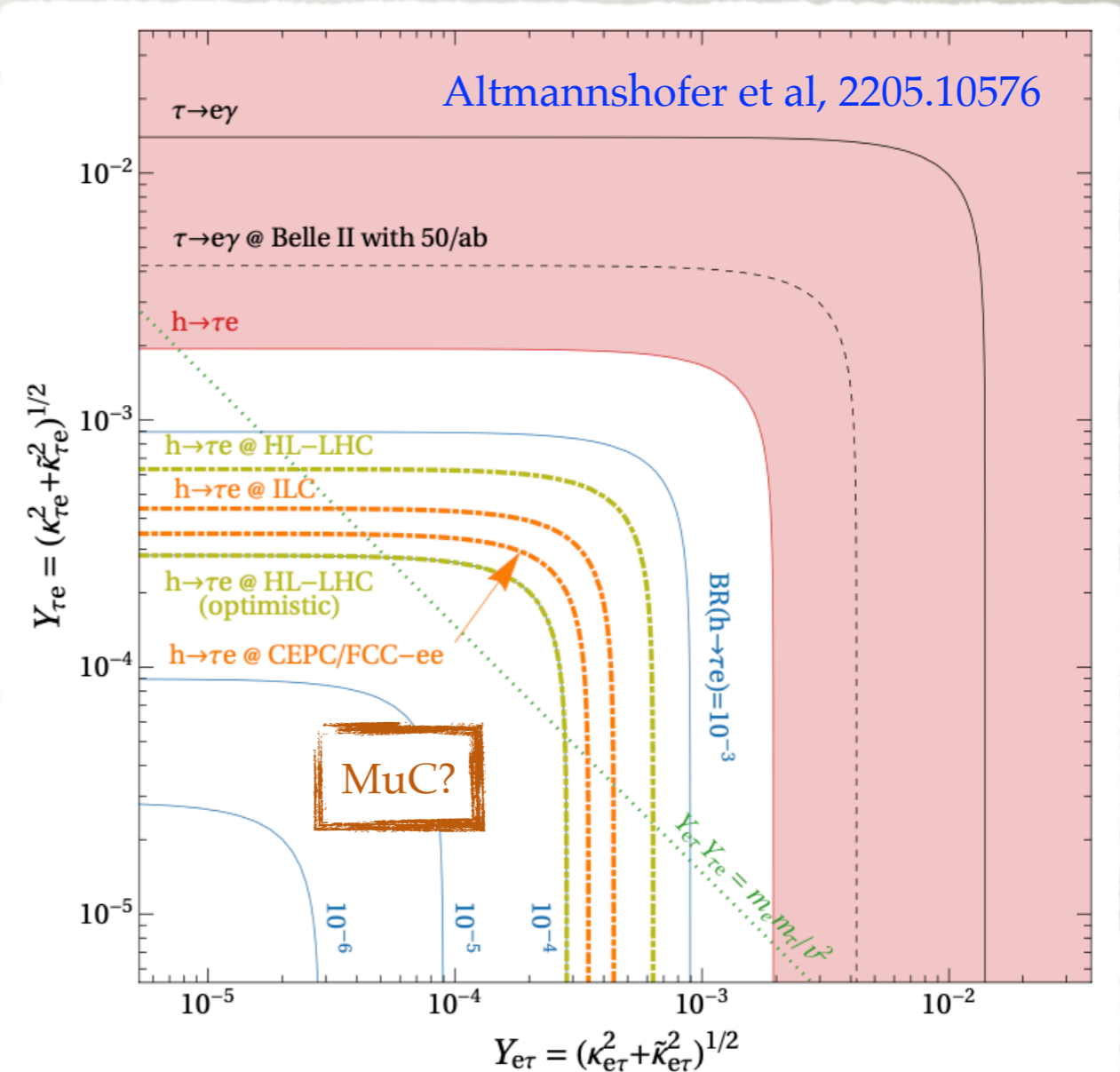
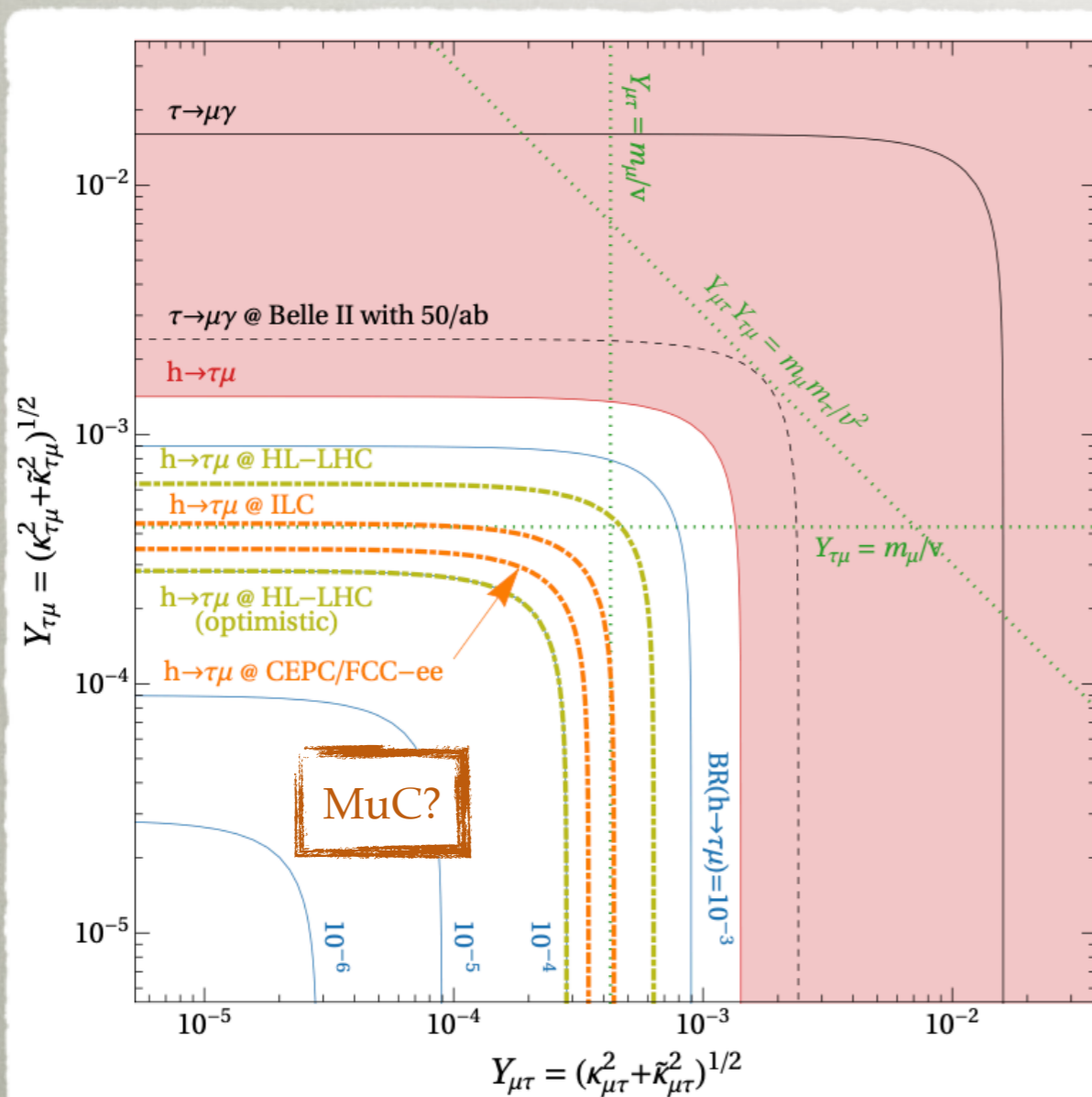
CMS, 2105.03007
ATLAS, 1907.06131



$Y_{\tau\mu}, Y_{\tau e}$ YUKAWAS - FUTURE

see also Qin et al, 1711.07243
Arroyo-Urena et al, 2002.04120

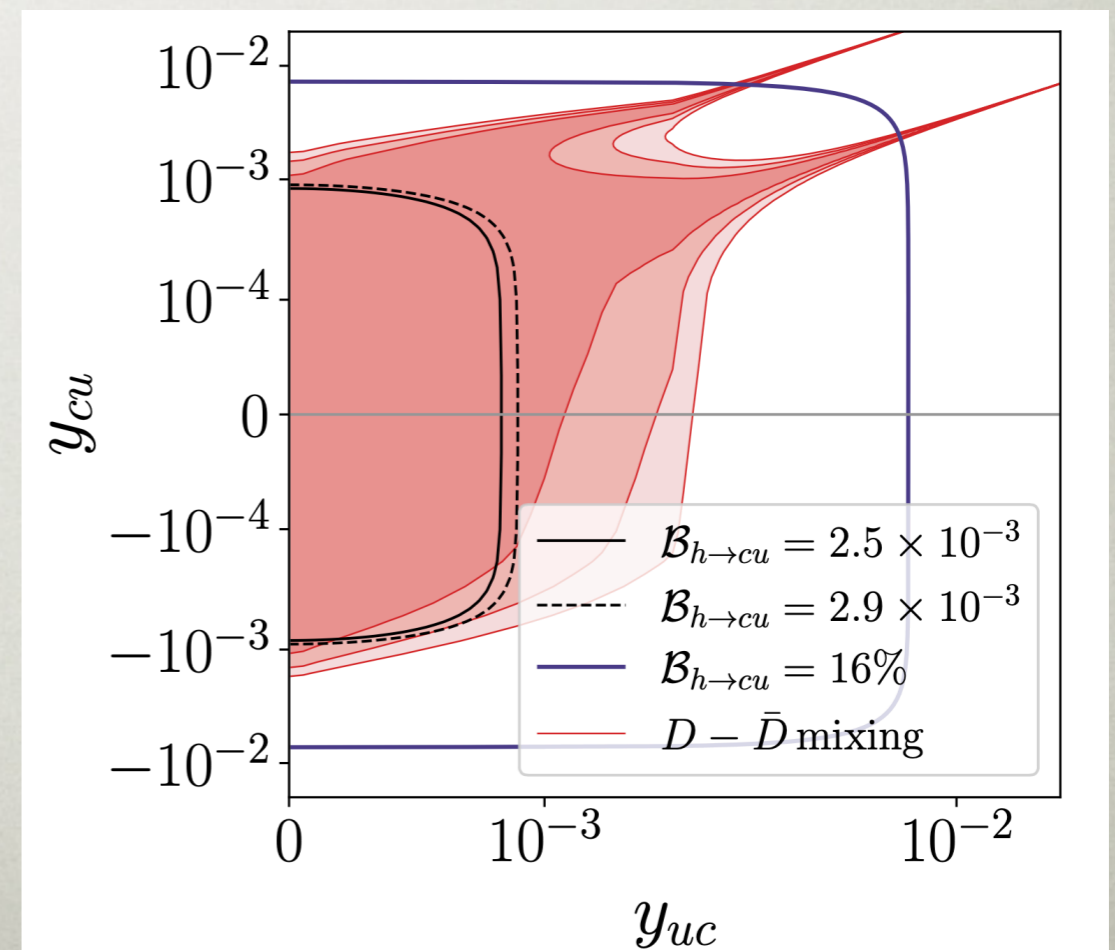
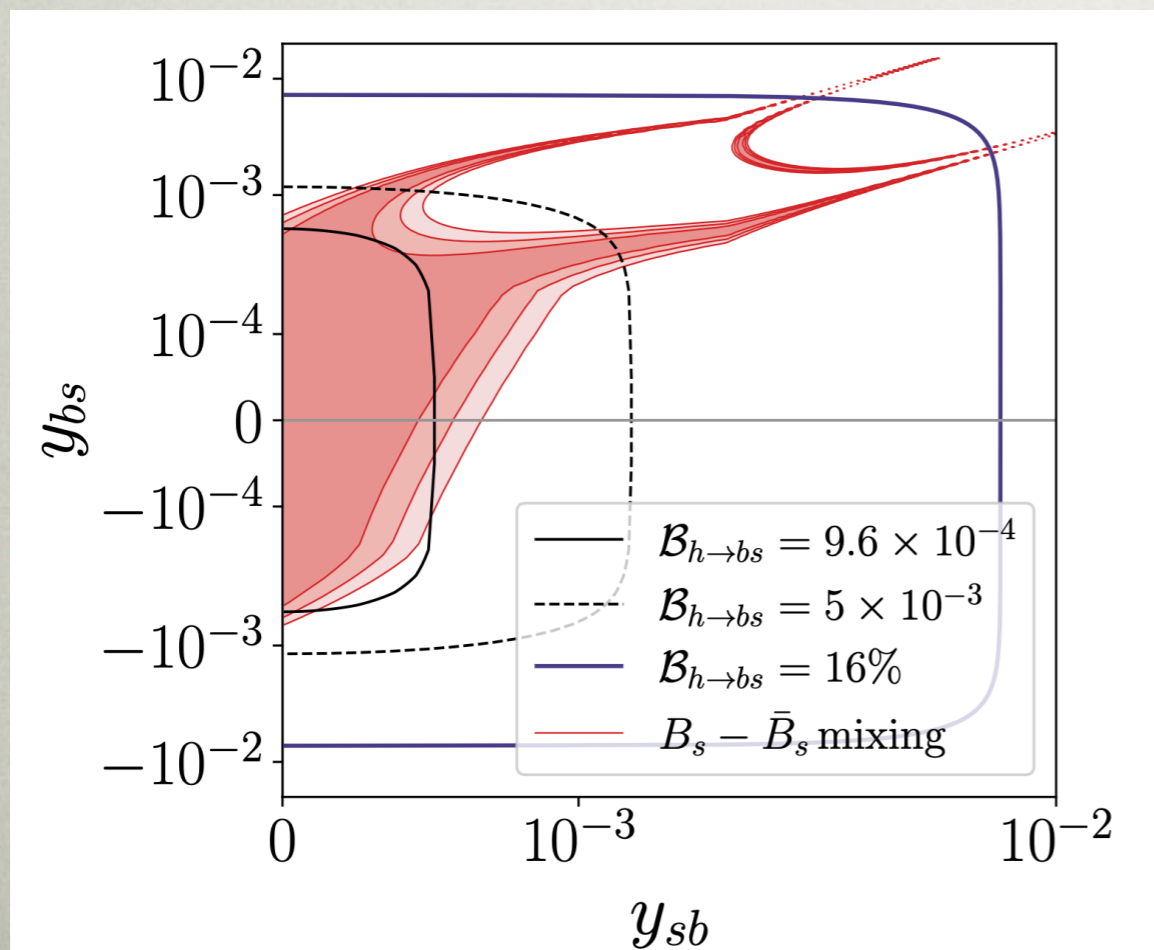
- $\text{Br} \lesssim \mathcal{O}(10^{-4}) \Rightarrow \Lambda \gtrsim 10 \text{ TeV} (C_{ij} = 1)$

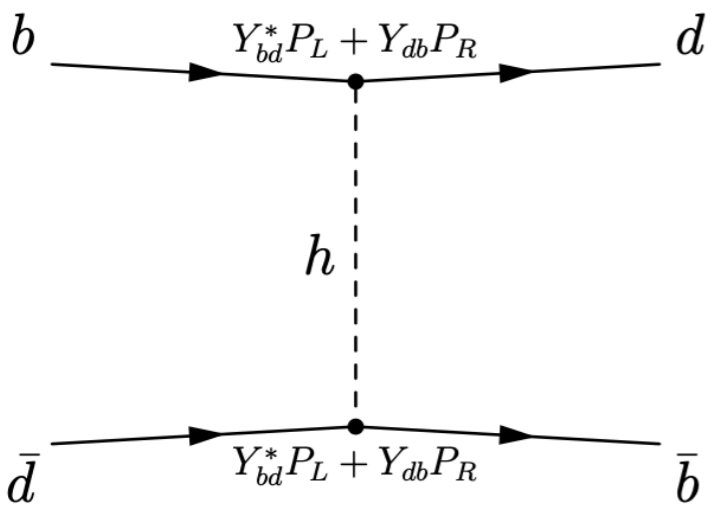


$h \rightarrow bs, cu$

- y_{bs}, y_{cu} bounded from meson mixing
- for $Br(h \rightarrow qq') \sim 10^{-3}$ direct searches become more important

Kamenik et al, 2306.17520



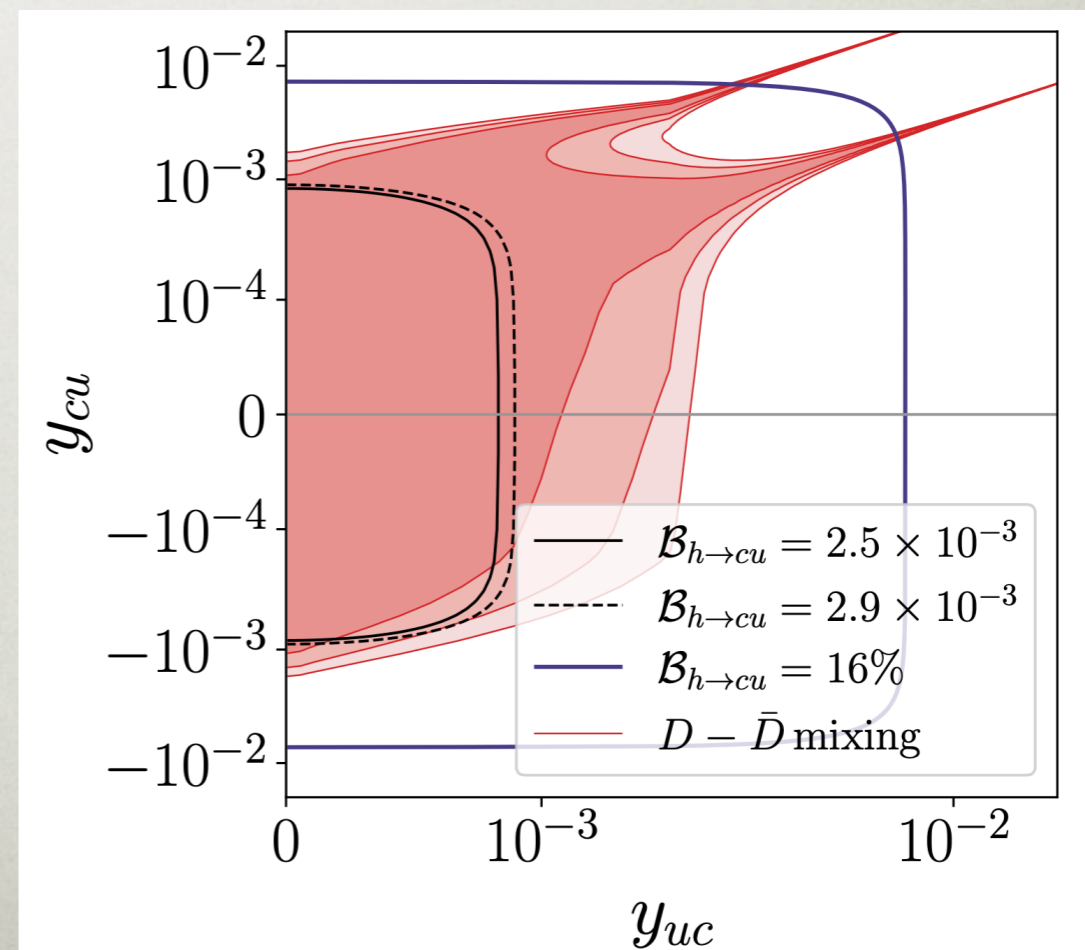
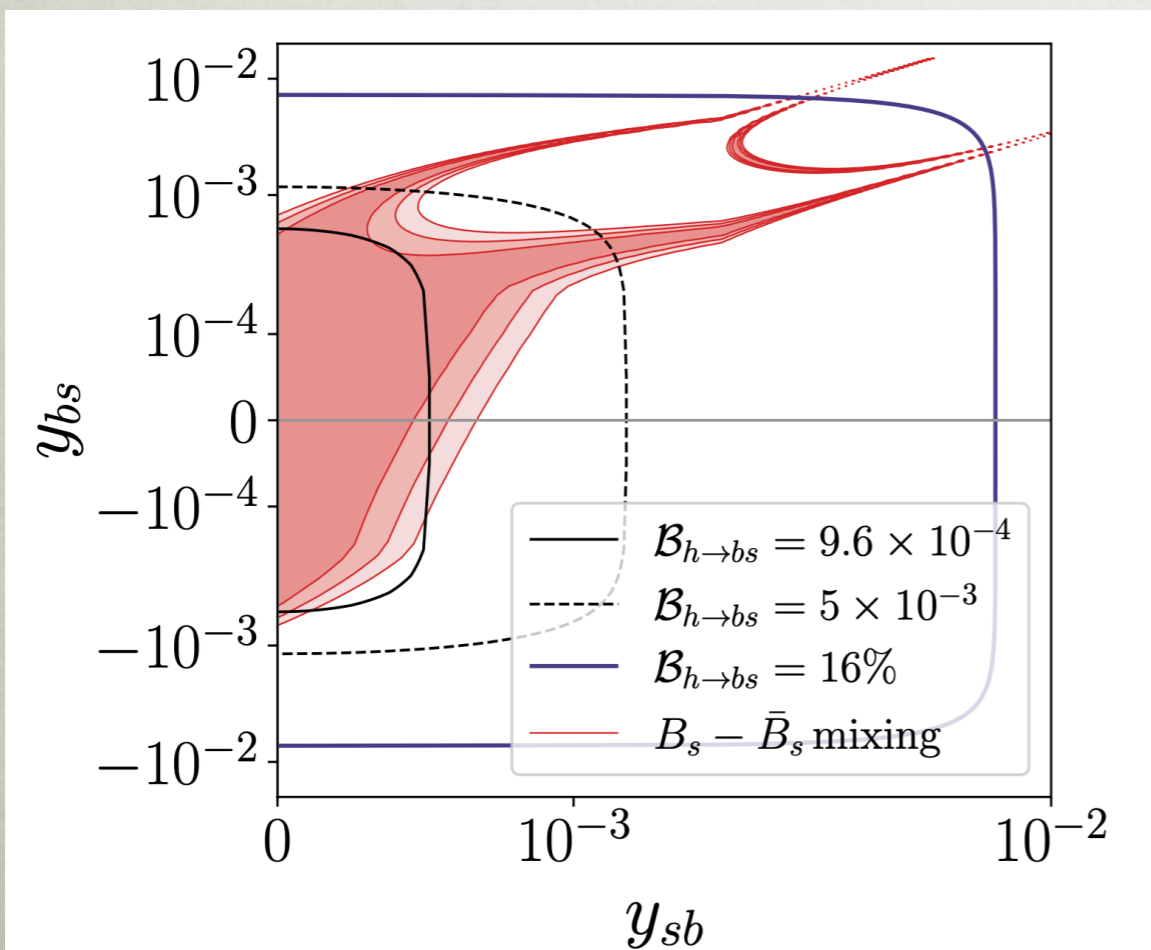


$h \rightarrow bs, cu$

bounded from meson mixing

- for $Br(h \rightarrow qq') \sim 10^{-3}$ direct searches become more important

Kamenik et al, 2306.17520

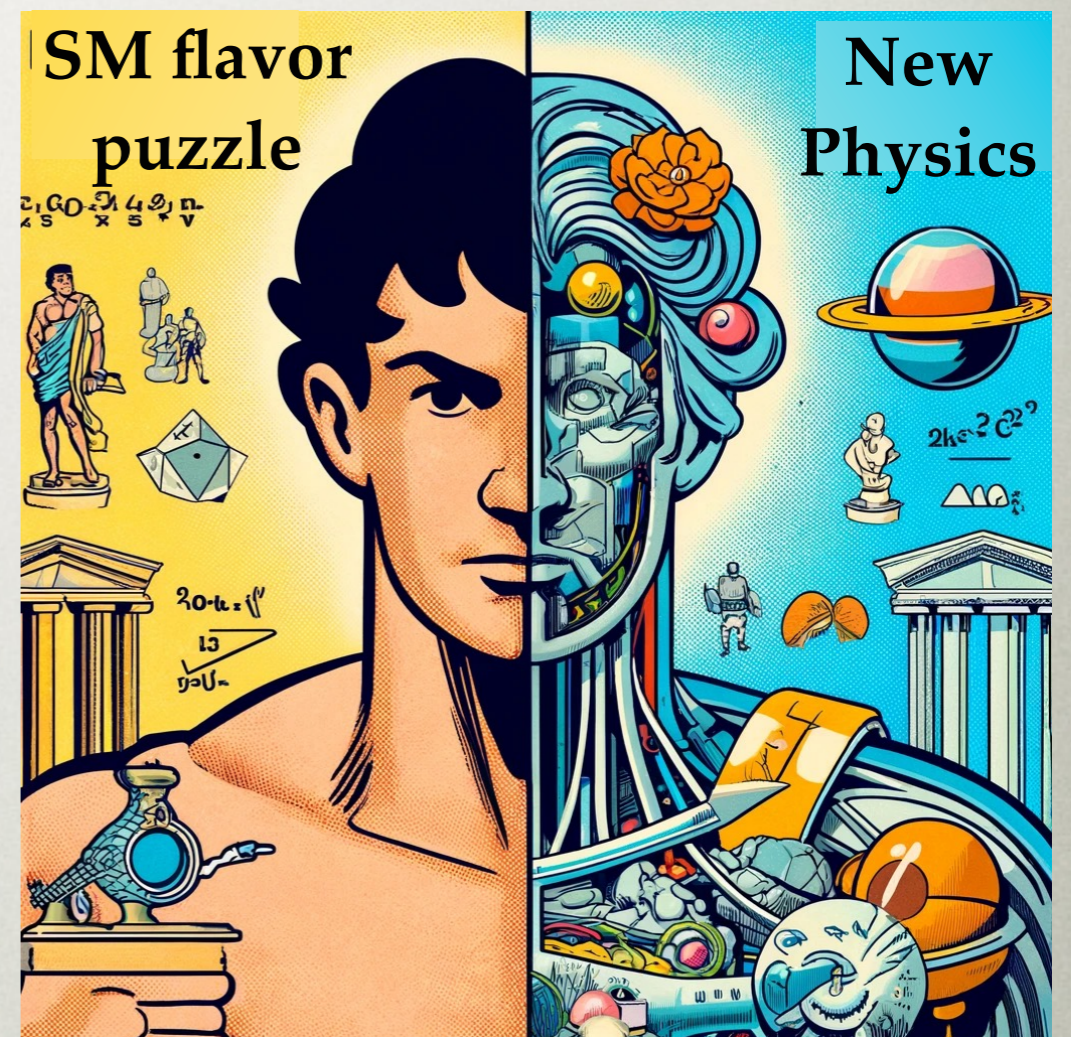


STANDARD MODEL FLAVOR PUZZLE

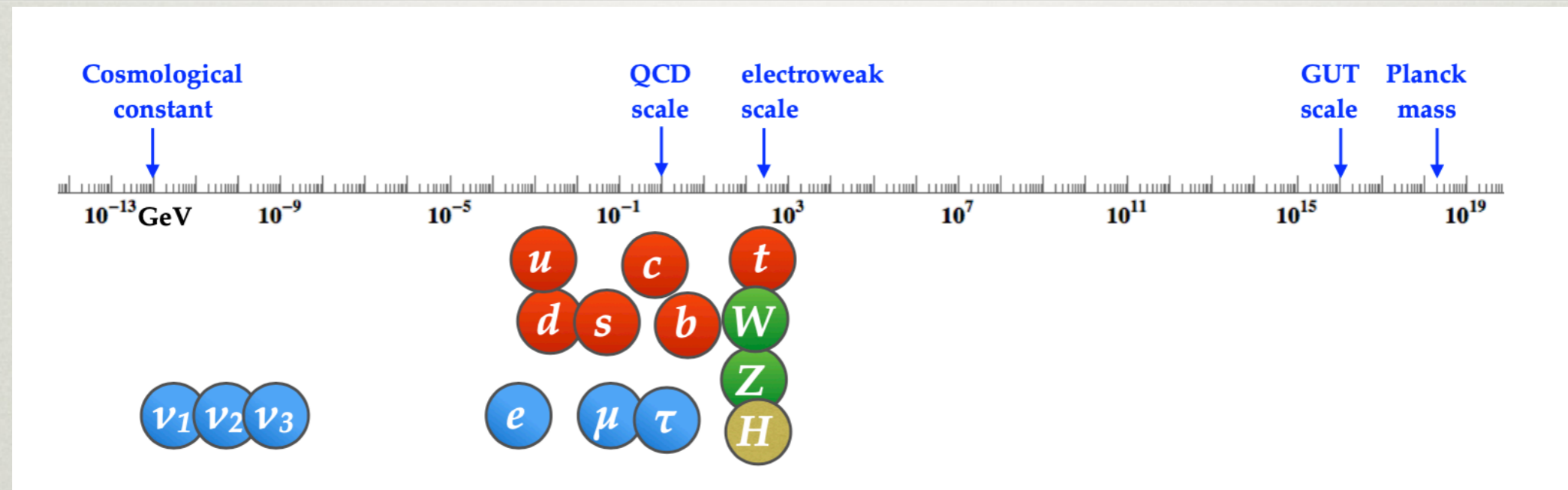
TWO FACES OF FLAVOR PHYSICS

- no flavor changing neutral currents in the SM
 - \Rightarrow flavor transitions sensitive probes of new physics

- why the observed structure of quark and lepton masses and mixings?
 - \Rightarrow flavor model building



FLAVOR MODEL BUILDING



- dynamical explanations
 - horizontal flavor symmetries
 - warped extra dimensions
 - partial compositeness
 - radiative fermion masses
- common to all: extra states, new sources of flavor violation
 - generically above the reach of MuC

SM FLAVOR PUZZLE @MUC

- if SM flavor puzzle has a dynamical explanation
- \Rightarrow same dynamics would control flavor couplings of NP states
 - for instance, if squarks are produced
 - flavor violating decays $\tilde{q}_i \rightarrow q_j \tilde{\chi}^0$ in principle possible
 - (squark mass degeneracy would avoid indirect constraints from FCNCs)
- the NP spectrum and couplings carry more information
 - e.g., for squarks, mass terms are $\tilde{q}_{Li}^* \tilde{q}_{Lj}$ and $\tilde{q}_{Ri}^* \tilde{q}_{Rj}$ instead of $\bar{q}_{Li} q_{Rj}$ for SM quarks

MORE FLAVOR @MUC

- other examples / possibilities
 - other LFV SMEFT operators through $\mu Z \rightarrow \tau h, \mu\gamma \rightarrow \tau h, \dots$ poster by H. Bagherian
 - CP violating $h \rightarrow \tau\tau, ht\bar{t}$ couplings
 - Y_{tc} coupling from $\mu^+\mu^- \rightarrow \bar{t}c + X$ or $\mu^+\mu^- \rightarrow \bar{t}ch + X$ similar to Y_{tt} in Liu et al, 2308.06323
 - probes of $(g - 2)_\mu$ UV physics see, e.g., Buttazzo, Paradisi, 2012.02769
 - V_{cb}, V_{cs} from W decays (so far analysis only for FCC-ee) Marzocca, Szevec, Tammaro, 2405.08880
 - exotic Z decays
 - flavor violating $Z \rightarrow \tau\mu, \tau e, \mu e$ decays
 - Z decays that could gain from boosts, e.g., $Z \rightarrow \tau(\tau \rightarrow \mu a)$
 - TeV muon beams to produce flavor violating ALPs in $\mu A \rightarrow \tau A a$ Batell et al, 2407.15942
 - $Y_{e\mu}$ from $e^-\mu^+ \rightarrow h$ thanks to Zhen Liu
 -

CONCLUSIONS

- two ways flavor may become important for physics at Muon Collider
 - probe off-shell new physics from precision flavor probes
 - example: higgs FV decays
 $h \rightarrow \tau\mu, h \rightarrow \tau e, h \rightarrow \mu e$
 $h \rightarrow bs, h \rightarrow cu$
 - flavor structure of new physics states produced on-shell
 - extra heavy Higgses, SUSY partners, ...

BACKUP SLIDES

EDM CONSTRAINTS

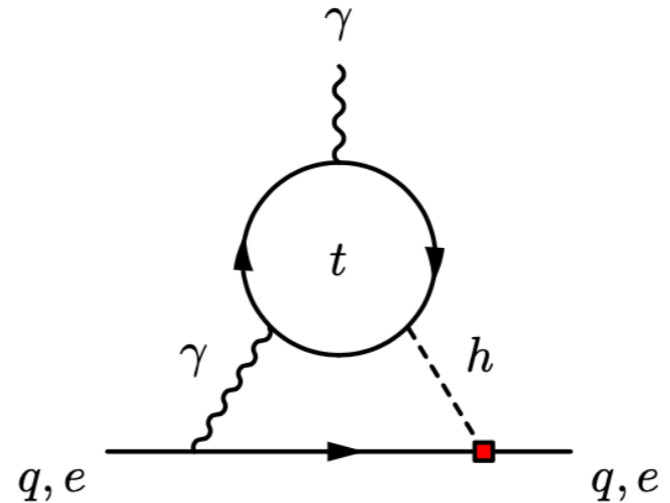
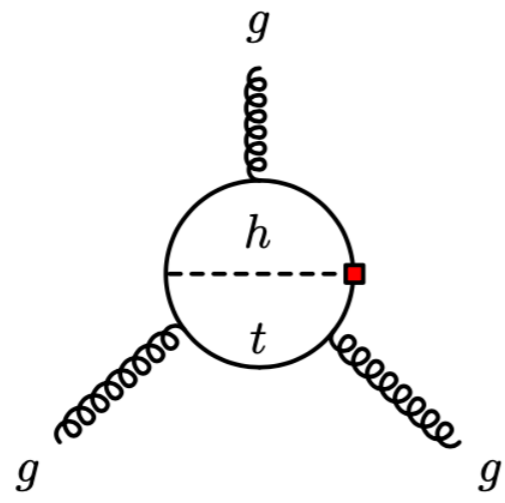
Brod et al, 2203.03736

- CPV dim. 6 ops contribute to EDMs through loops

$$\mathcal{L}_h = - \sum_{f=u,d,\ell} \frac{h}{\sqrt{2}} \bar{f}_L \left(y_f^{\text{SM}} - \frac{v^2}{\Lambda^2} C_{fH} \right) f_R + \text{h.c.},$$

$$C_{fH,ij} = C_{fH+,ij} + iC_{fH-,ij}$$

- 2-loop electroweak mixing \Rightarrow large logs \Rightarrow leading log
QCD resummation
- if single flavor dominates \Rightarrow in general stringent constraints
 - two extreme cases: top and muon
- if more than one flavor important
 - cancellations in EDMs possible but not in LHC constr.
 - the parameter space for CPV couplings opens up \Rightarrow FCC



Brod et al, 2203.03736

- CPV dim. 6 ops contribute to EDMs through loops

$$\mathcal{L}_h = - \sum_{f=u,d,\ell} \frac{h}{\sqrt{2}} \bar{f}_L \left(y_f^{\text{SM}} - \frac{v^2}{\Lambda^2} C_{fH} \right) f_R + \text{h.c.},$$

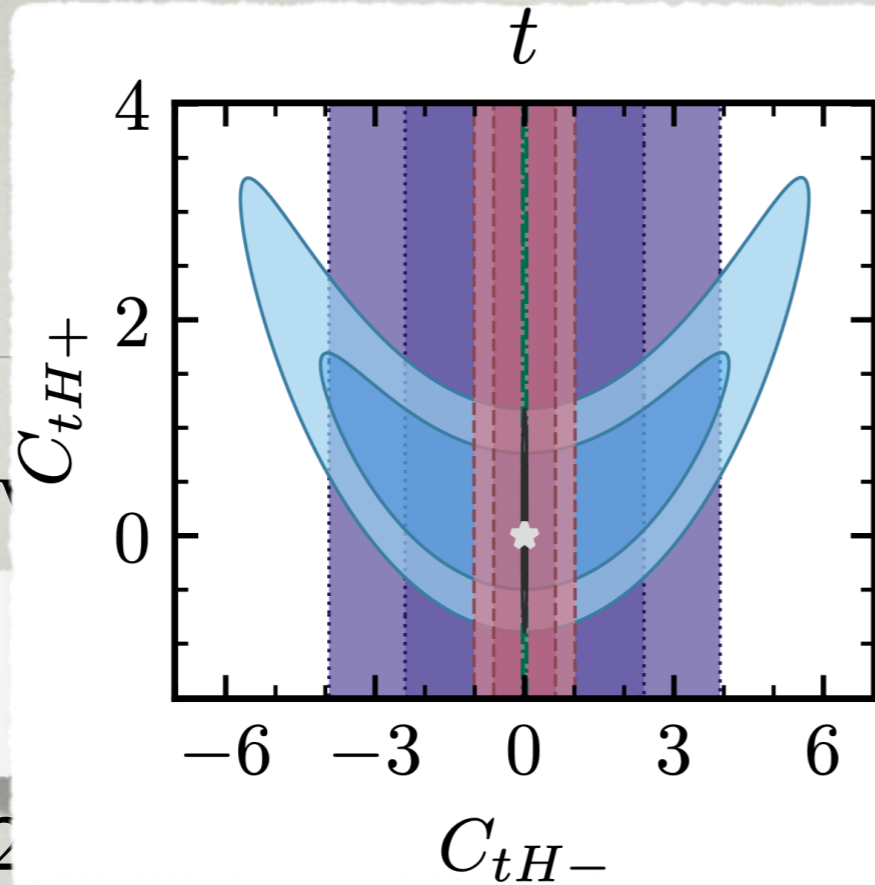
$$C_{fH,ij} = C_{fH+,ij} + iC_{fH-,ij}$$

- 2-loop electroweak mixing \Rightarrow large logs \Rightarrow leading log QCD resummation
- if single flavor dominates \Rightarrow in general stringent constraints
 - two extreme cases: top and muon
- if more than one flavor important
 - cancellations in EDMs possible but not in LHC constr.
 - the parameter space for CPV couplings opens up \Rightarrow FCC

CONSTRAINTS

Brod et al, 2203.03736

- CPV



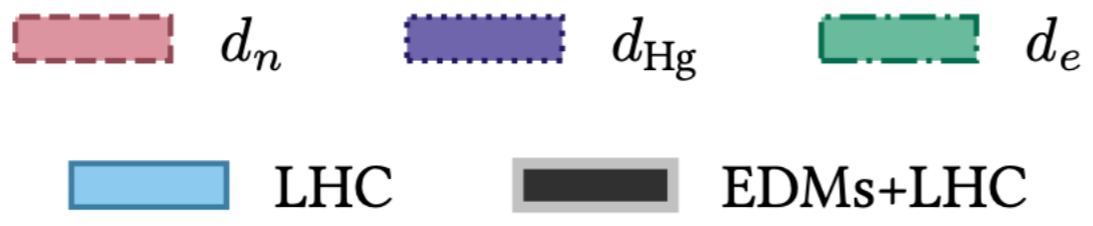
EDMs through loops

$$\left(f_R + \text{h.c.}, \right)$$

$$C_{fH,ij} = C_{fH+,ij} + iC_{fH-,ij}$$

\Rightarrow large logs \Rightarrow leading log

- if si



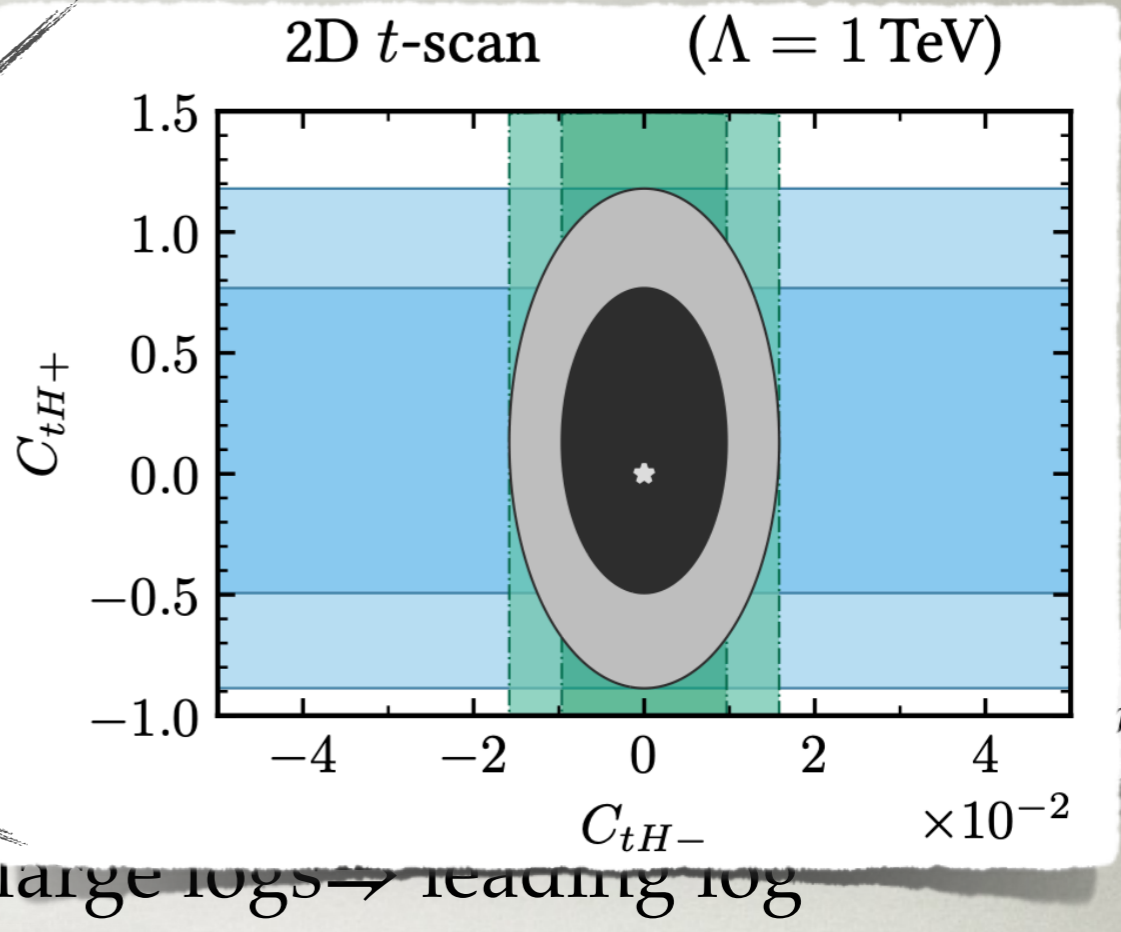
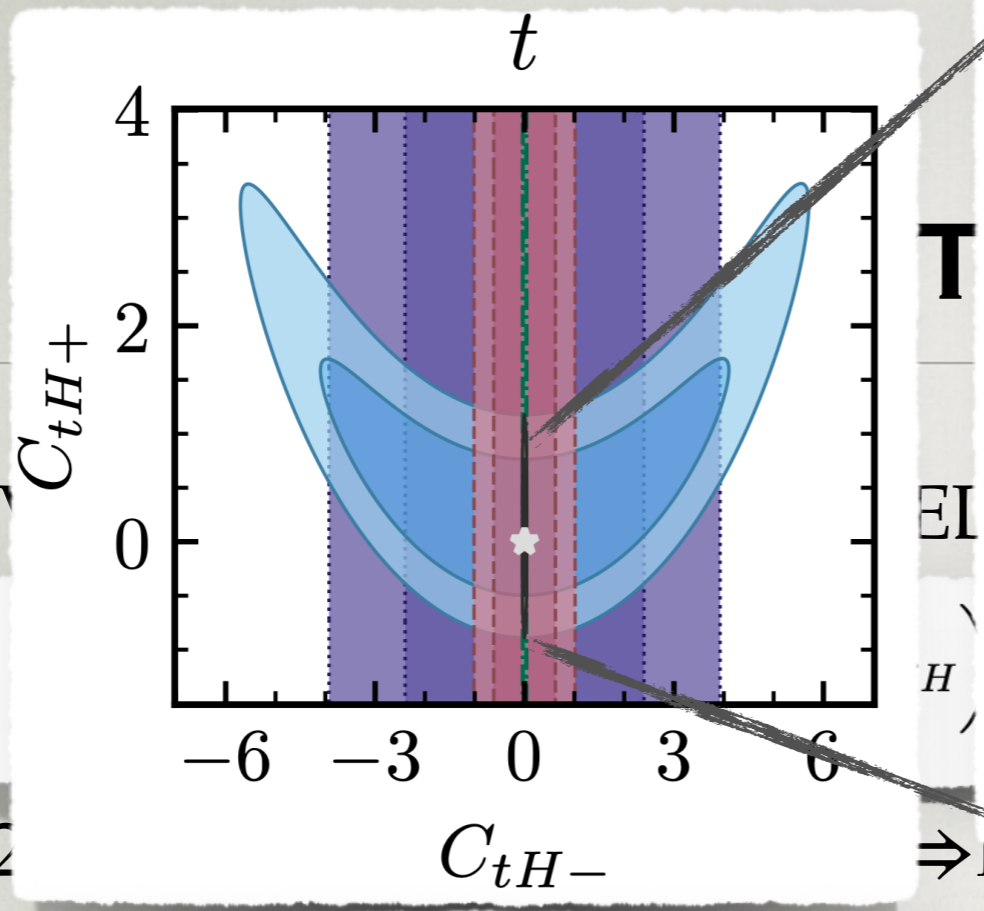
al stringent

- if more than one flavor important

- cancellations in EDMs possible but not in LHC constr.

- the parameter space for CPV couplings opens up \Rightarrow FCC

- CPV

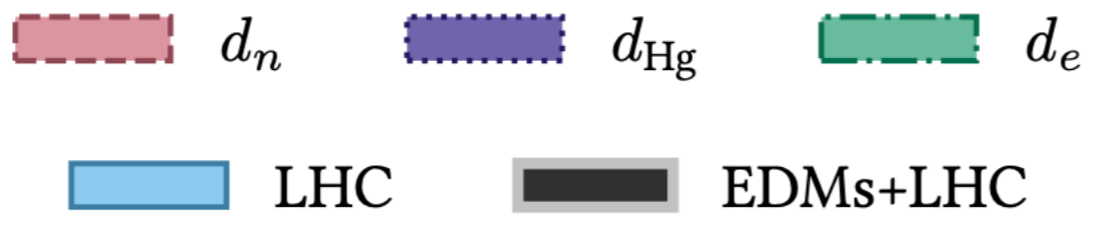


⇒ large logs ⇒ leading log

- 2

- if si

con



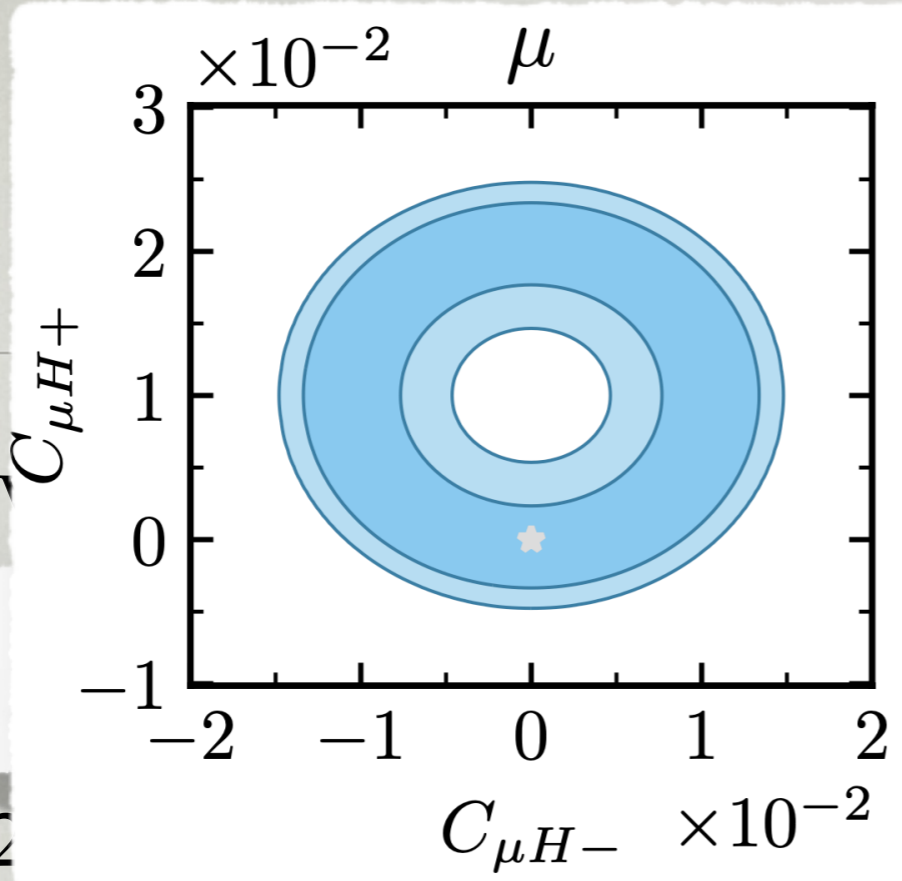
al stringent

- two extreme cases: top and muon
- if more than one flavor important
 - cancellations in EDMs possible but not in LHC constr.
 - the parameter space for CPV couplings opens up ⇒ FCC

CONSTRAINTS

Brod et al, 2203.03736

- CPV








EDMs through loops

$$\left(f_R + \text{h.c.}, \right)$$

$$C_{fH,ij} = C_{fH+,ij} + iC_{fH-,ij}$$

- \Rightarrow large logs \Rightarrow leading log

- if si

	d_n		d_{Hg}		d_e
	LHC		EDMs+LHC		

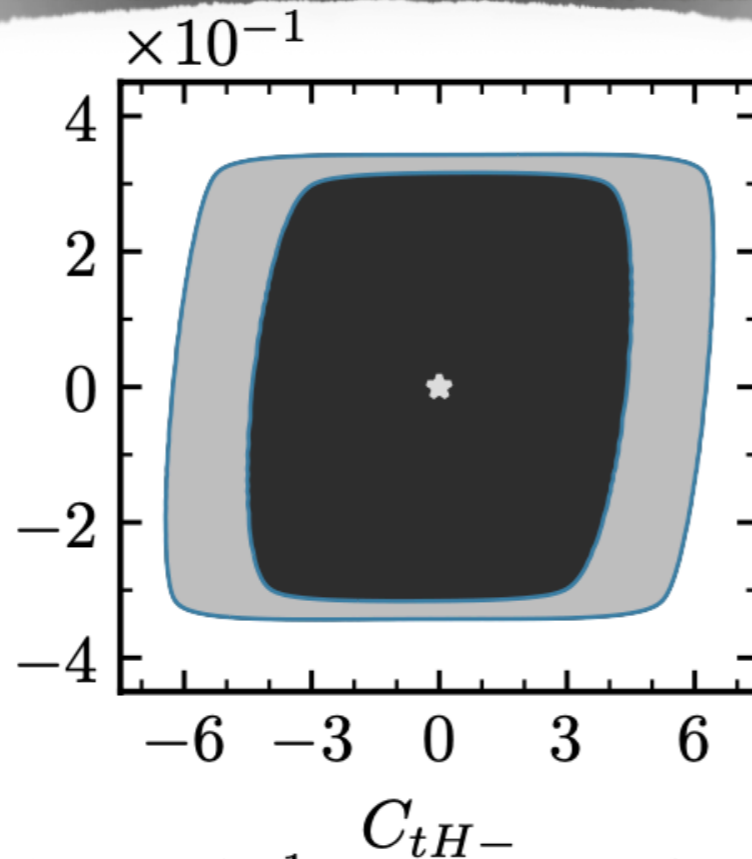
al stringent

- two extreme cases: top and muon
- if more than one flavor important
 - cancellations in EDMs possible but not in LHC constr.
 - the parameter space for CPV couplings opens up \Rightarrow FCC

CONSTRAINTS

Brod et al, 2203.03736

- CPV



EDMs through loops

$$(f_H) f_R + \text{h.c.},$$

$$C_{fH,ij} = C_{fH+,ij} + iC_{fH-,ij}$$

\Rightarrow large logs \Rightarrow leading log

- 2

- if si

d_n d_{Hg} d_e

al stringent

con

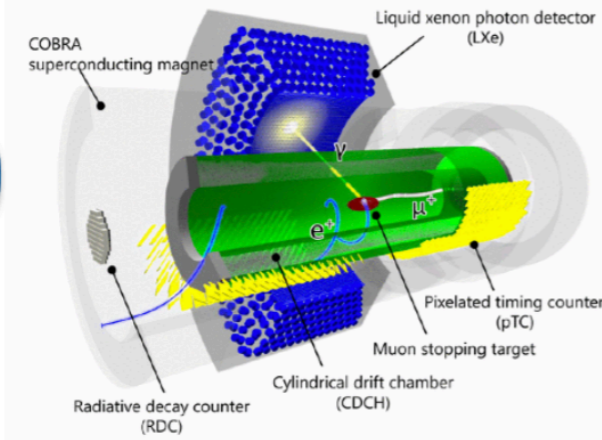
LHC EDMs+LHC

- two extreme cases: top and muon
- if more than one flavor important
 - cancellations in EDMs possible but not in LHC constr.
 - the parameter space for CPV couplings opens up \Rightarrow FCC

cLFV experiments in the world

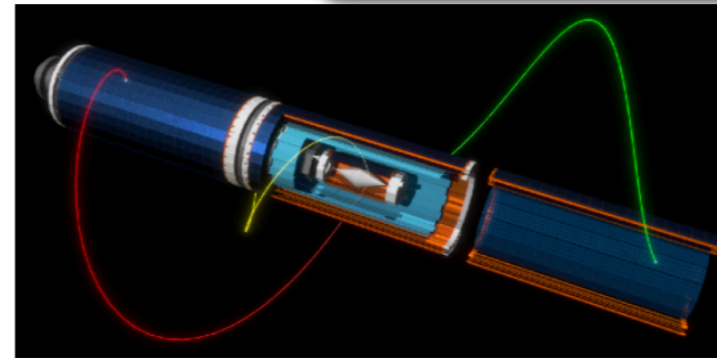
MEG II

$\mu^+ \rightarrow e^+ \gamma$



Mu3e

$\mu^+ \rightarrow e^+ e^+ e^-$



Coincidence measurement:
DC beam needed to minimize
backgrounds from accidental
coincidences

$BKG \propto (Rate)^2$

PSI



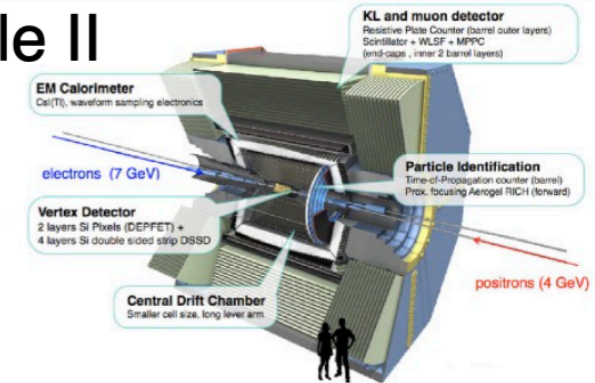
CERN

LHCb/ATLAS/CMS

$\tau \rightarrow 3\mu, \tau \rightarrow \mu\gamma$

KEK

Belle II

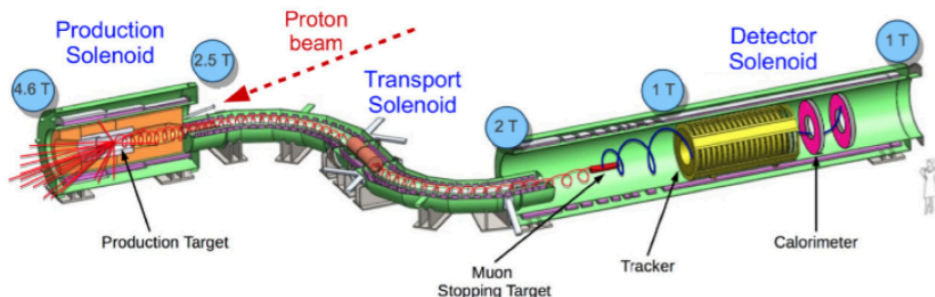


Fermilab

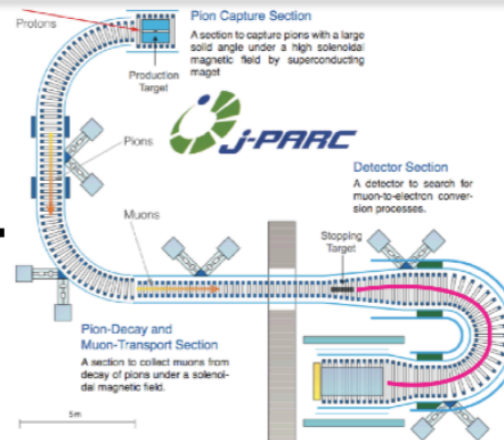
$\mu-N \rightarrow e-N$

J-PARC

Mu2e



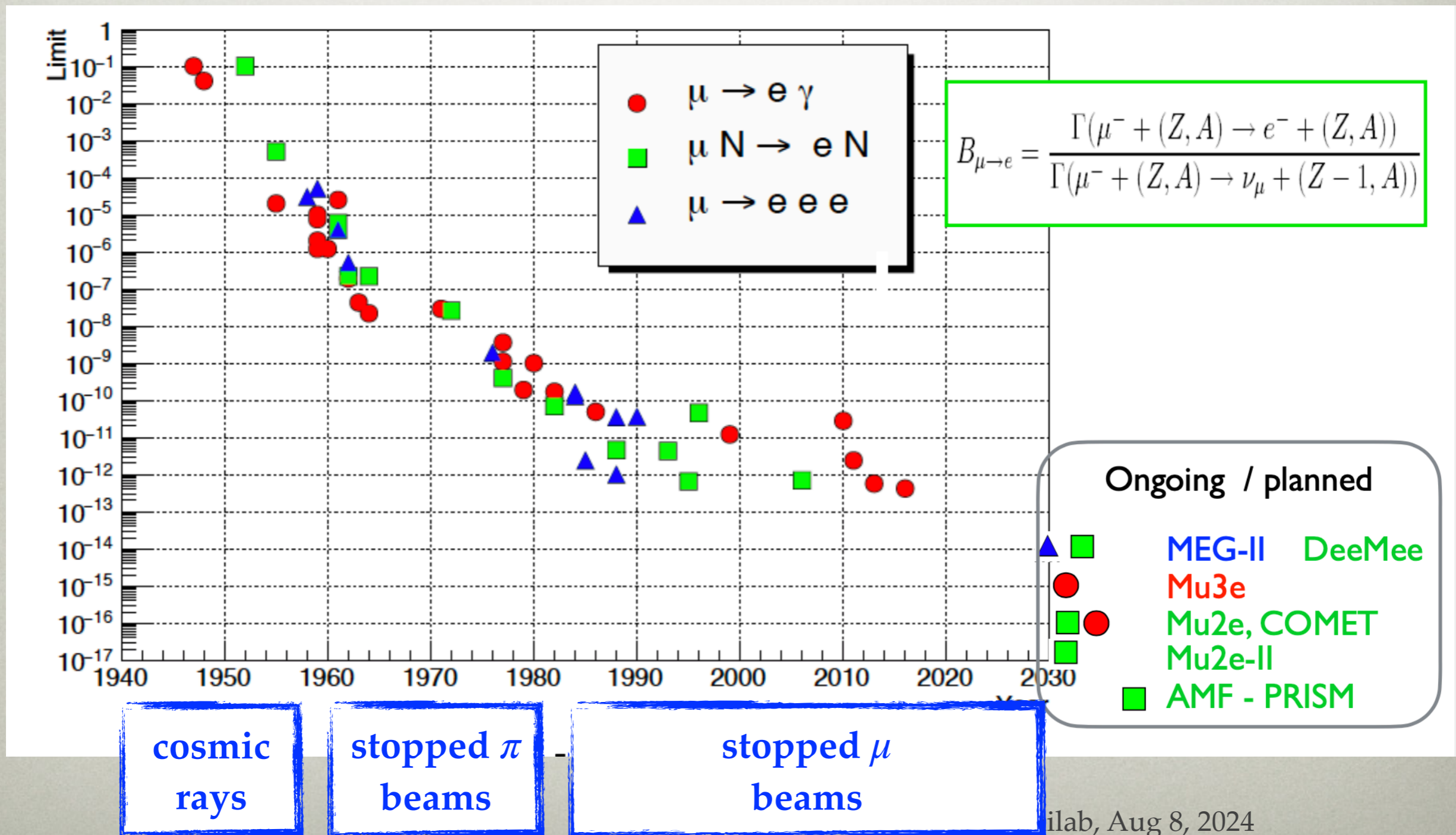
DeeMe,
COMET



Single e^- measurement:
pulsed beam needed
Many pion-induced
backgrounds after
proton pulse
wait it out with 26 ns
lifetime

EXPERIMENTAL PROGRESS

- steady experimental progress since 1940s



CPV AND FV HIGGS COUPLINGS TO SM FERMIONS

$$m = v \begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$$Y = \begin{pmatrix} * & * & * \\ * & * & * \\ * & * & * \end{pmatrix}$$

- NP contris in general lead to
 - flavor violating Higgs decays
 - CPV Higgs decays
- different models lead to different patterns of flavor diagonal and flavor violating Yukawas
- note: large Yukawas for light fermions require significant tuning in contributions to fermion masses

Harnik, Kopp, JZ, 1209.1397
Blackenburg, Ellis, Isidori, 1202.5704

CPV AND FV HIGGS

C

Model	$\kappa_{ct(tc)}/\kappa_t$	$\kappa_{ut(tu)}/\kappa_t$	$\kappa_{uc(cu)}/\kappa_t$
GL & GL2	$\epsilon(\epsilon^2)$	$\epsilon(\epsilon^2)$	ϵ^3
MFV	$\frac{\text{Re}(c_u m_b^2 V_{cb}^{(*)})}{\Lambda^2} \frac{\sqrt{2} m_{t(c)}}{v_W}$	$\frac{\text{Re}(c_u m_b^2 V_{ub}^{(*)})}{\Lambda^2} \frac{\sqrt{2} m_{t(u)}}{v_W}$	$\frac{\text{Re}(c_u m_b^2 V_{ub(cb)} V_{cb(ub)}^*)}{\Lambda^2} \frac{\sqrt{2} m_{c(u)}}{v_W}$
RS	$\sim \lambda^{(-)2} \frac{m_{t(c)}}{v_W} \bar{Y}^2 \frac{v_W^2}{m_{KK}^2}$	$\sim \lambda^{(-)3} \frac{m_{t(u)}}{v_W} \bar{Y}^2 \frac{v_W^2}{m_{KK}^2}$	$\sim \lambda^{(-)1} \frac{m_{c(u)}}{v_W} \bar{Y}^2 \frac{v_W^2}{m_{KK}^2}$
pNGB	$\mathcal{O}(y_*^2 \frac{m_t}{v_W} \frac{\lambda_{L(R),2} \lambda_{L(R),3} m_W^2}{M_*^2})$	$\mathcal{O}(y_*^2 \frac{m_t}{v_W} \frac{\lambda_{L(R),1} \lambda_{L(R),3} m_W^2}{M_*^2})$	$\mathcal{O}(y_*^2 \frac{m_c}{v_W} \frac{\lambda_{L(R),1} \lambda_{L(R),2} m_W^2}{M_*^2})$

Bishara et al, 1504.04022

- NP contris in general lead to
 - flavor violating Higgs decays
 - CPV Higgs decays
- different models lead to different patterns of flavor diagonal and flavor violating Yukawas
- note: large Yukawas for light fermions require significant tuning in contributions to fermion masses

Harnik, Kopp, JZ, 1209.1397

Blackenburg, Ellis, Isidori, 1202.5704

CPV AND FV HIGGS COUPLINGS TO SM FERMIONS

$$m = v \begin{pmatrix} * & 0 & 0 \\ 0 & * & 0 \\ 0 & 0 & * \end{pmatrix}$$

$$Y = \begin{pmatrix} * & * & * \\ * & * & * \\ * & * & * \end{pmatrix}$$

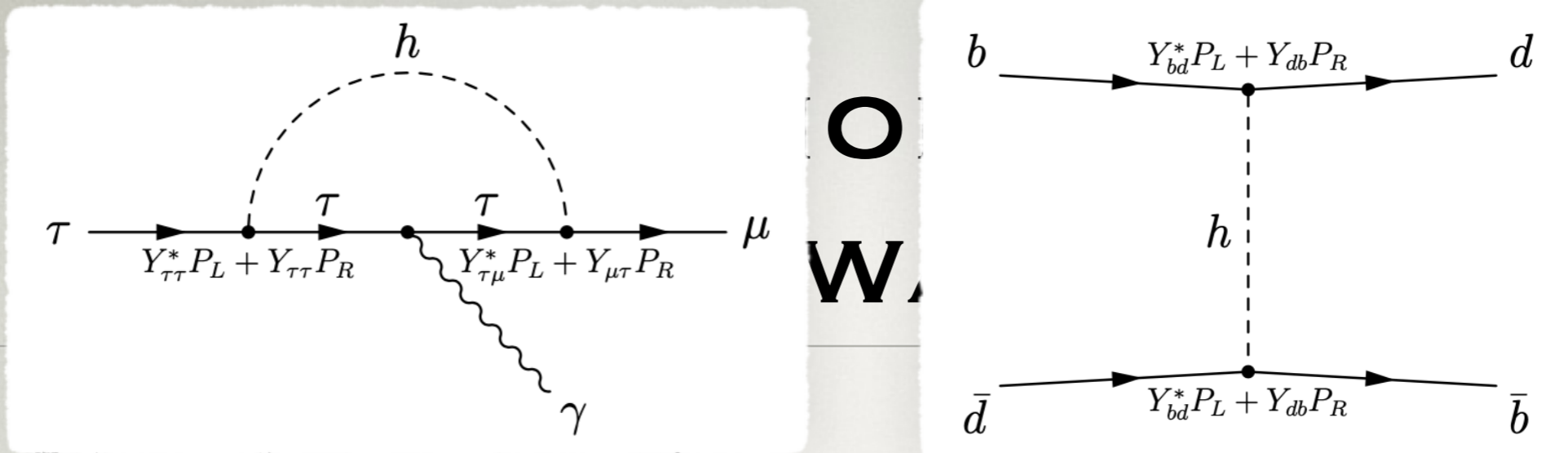
- NP contris in general lead to
 - flavor violating Higgs decays
 - CPV Higgs decays
- different models lead to different patterns of flavor diagonal and flavor violating Yukawas
- note: large Yukawas for light fermions require significant tuning in contributions to fermion masses

Harnik, Kopp, JZ, 1209.1397
Blackenburg, Ellis, Isidori, 1202.5704

FLAVOR VIOLATING YUKAWAS

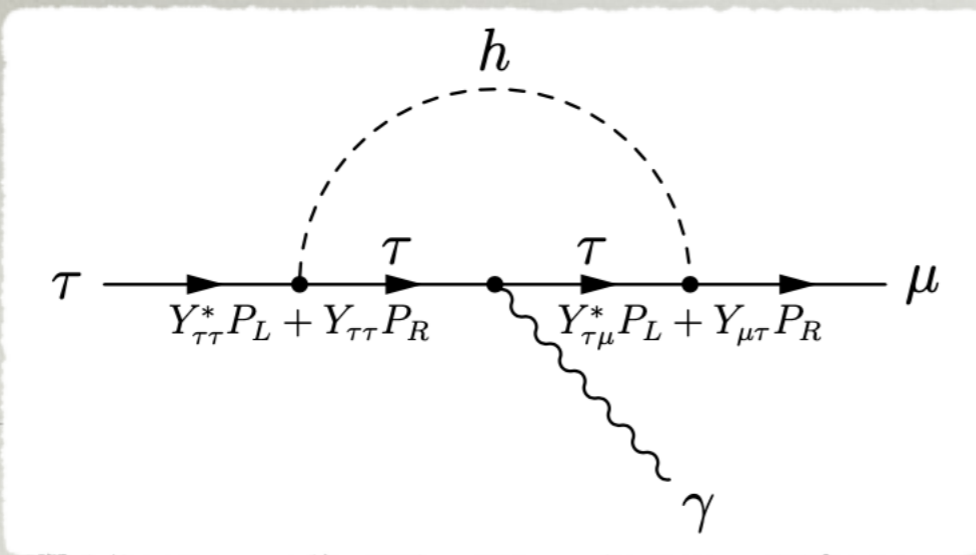
Harnik, Kopp, JZ, 1209.1397
Blackenburg, Ellis, Isidori, 1202.5704

- FV Yukawas induce other FCNCs
 - D, B, B_s, K mixing \Rightarrow constrain FV quark Yukawas
 - these very poor for top $\Rightarrow Y_{tq}$ Yukawas from LHC, HL-LHC, future coll.
 - $Y_{bs'}$, Y_{cu} maybe in reach of future coll
 - $\mu \rightarrow e\gamma, \mu \rightarrow 3e, \dots \Rightarrow$ stringent constraints on $Y_{\mu e'}$, less on $Y_{\tau e}, Y_{\tau\mu}$
- "model independent" \Rightarrow not easy to suppress in concrete models

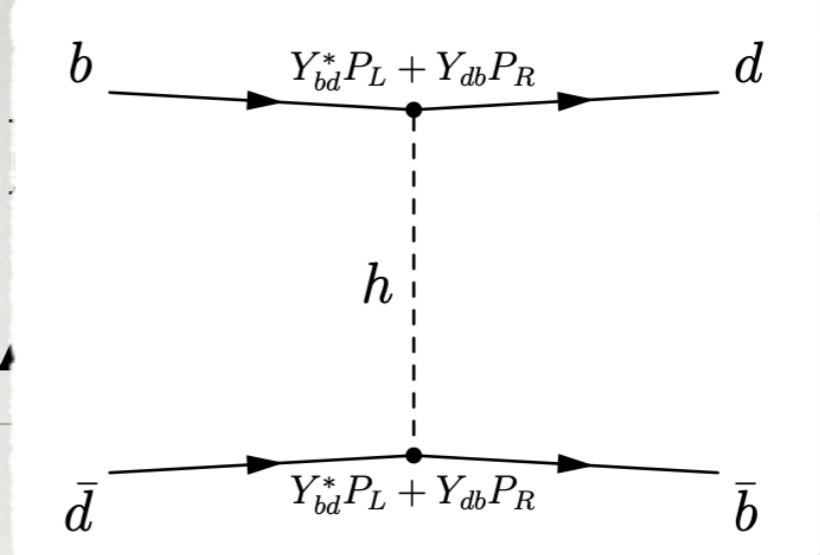


pp, JZ, 1209.1397
 sidori, 1202.5704

- FV Yukawas induce other FCNCs
 - D, B, B_s, K mixing \Rightarrow constrain FV quark Yukawas
 - these very poor for top $\Rightarrow Y_{tq}$ Yukawas from LHC, HL-LHC, future coll.
 - $Y_{bs'}$, Y_{cu} maybe in reach of future coll
 - $\mu \rightarrow e\gamma, \mu \rightarrow 3e, \dots \Rightarrow$ stringent constraints on $Y_{\mu e'}$ less on $Y_{\tau e}, Y_{\tau\mu}$
- "model independent" \Rightarrow not easy to suppress in concrete models



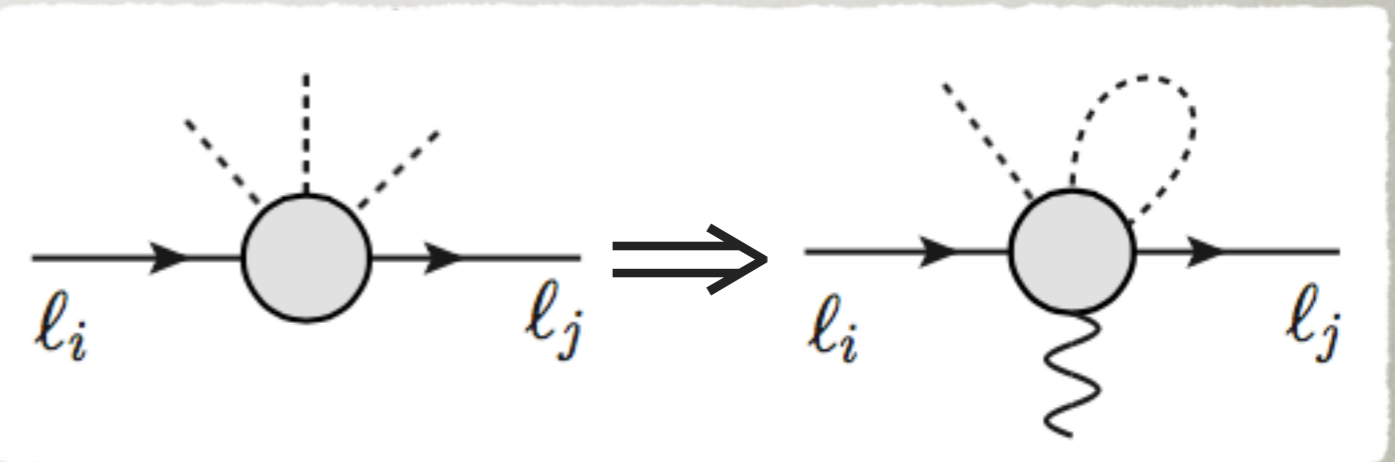
\mathcal{O}
 \mathcal{W}_A



pp, JZ, 1209.1397
sidori, 1202.5704

- FV Yukawas induce other FCNCs
 - D, B, B_s, K mixing \Rightarrow constrain FV quark Yukawas

- these very poor
 - LHC, HL-LHC,
 - $Y_{bs'}$, Y_{cu} maybe in



- $\mu \rightarrow e\gamma, \mu \rightarrow 3e, \dots \Rightarrow$ stringent constraints on $Y_{\mu e'}$ less on $Y_{\tau e}, Y_{\tau\mu}$

- "model independent" \Rightarrow not easy to suppress in concrete models