

# FCC-ee: Physics & Detectors



Christoph Paus, MIT

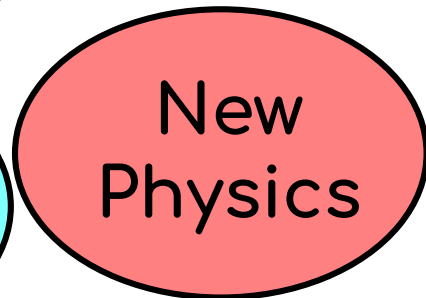
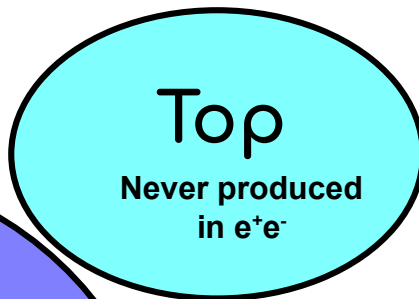
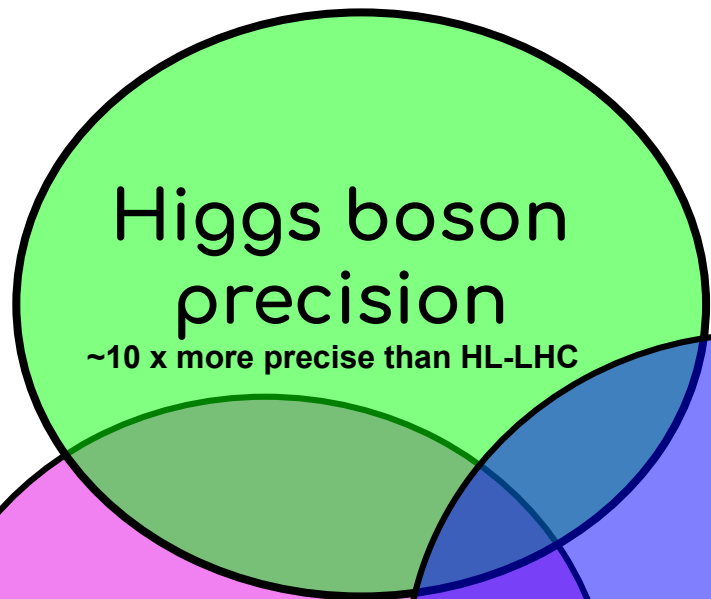
July 11, 2024  
57<sup>th</sup> Users Meeting, Fermilab

# Reporting from June 2024 FCC Week in San Francisco

<https://indico.cern.ch/event/1298458/timetable/>



# *Far away vision*



'Low mass' and high luminosity



Kind of 'Intensity frontier'



# FCC-ee Specs

Parameter	Z	WW	H (ZH)	ttbar
beam energy [GeV]	45	80	120	182.5
beam current [mA]	1280	135	26.7	5.0
number bunches/beam	10000	880	248	36
bunch intensity [ $10^{11}$ ]	2.43	2.91	2.04	2.64
SR energy loss / turn [GeV]	0.0391	0.37	1.869	10.0
total RF voltage 400/800 MHz [GV]	0.120/0	1.0/0	2.08/0	4.0/7.25
long. damping time [turns]	1170	216	64.5	18.5
horizontal beta* [m]	0.1	0.2	0.3	1
vertical beta* [mm]	0.8	1	1	1.6
horizontal geometric emittance [nm]	0.71	2.17	0.64	1.49
vertical geom. emittance [pm]	1.42	4.34	1.29	2.98
horizontal rms IP spot size [ $\mu\text{m}$ ]	8	21	14	39
vertical rms IP spot size [nm]	34	66	36	69
luminosity per IP [ $10^{34} \text{ cm}^{-2}\text{s}^{-1}$ ]	182	19.4	7.3	1.33
total integrated luminosity / year [ $\text{ab}^{-1}/\text{yr}$ ] 4 IPs	87	9.3	3.5	0.65
beam lifetime (rad Bhabha + BS+lattice)	8	18	6	10

From Fabiola's talk

Currently assessing technical feasibility of changing operation sequence (e.g. starting at ZH energy)

4 years  
 $5 \times 10^{12}$  Z  
 $\text{LEP} \times 10^5$

2 years  
 $> 10^8$  WW  
 $\text{LEP} \times 10^4$

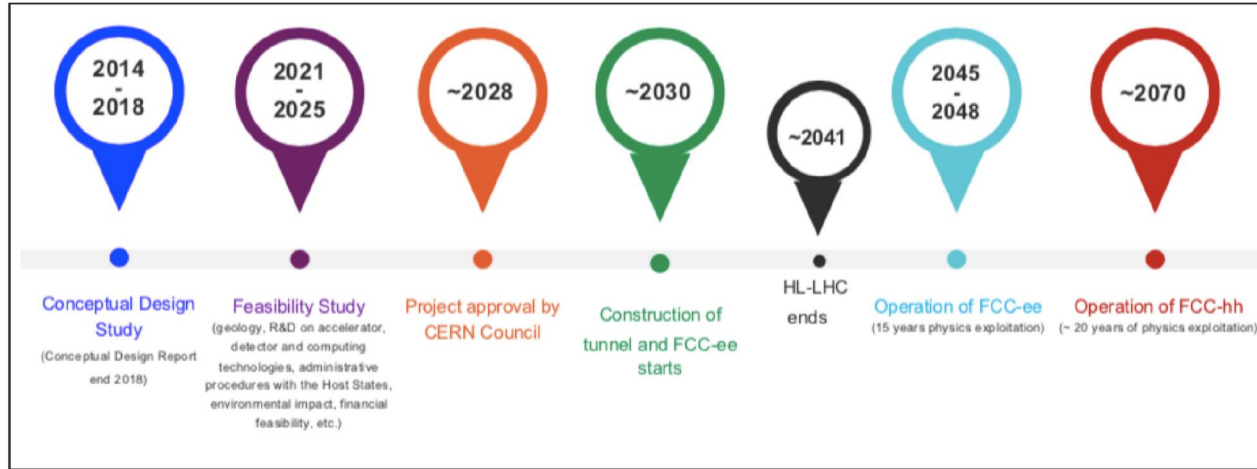
3 years  
 $2 \times 10^6$  H

5 years  
 $2 \times 10^6$  tt pairs

- $\times 10$ -50 improvements on all EW observables
- up to  $\times 10$  improvement on Higgs coupling (model-indep.) measurements over HL-LHC
- $\times 10$  Belle II statistics for b, c,  $\tau$
- indirect discovery potential up to  $\sim 70$  TeV
- direct discovery potential for feebly-interacting particles over 5-100 GeV mass range

Up to 4 interaction points  $\rightarrow$  robustness, statistics, possibility of specialised detectors to maximise physics output

# FCC-ee Schedule



From Fabiola's talk

**1<sup>st</sup> stage collider FCC-ee:**

electron-positron collisions 90-360 GeV:  
electroweak and Higgs factory

**2<sup>nd</sup> stage collider FCC-hh:**

proton-proton collisions at ~ 100 TeV

“Realistic” schedule taking into account:

- past experience in building colliders at CERN
- the various steps of approval process: ESPP update, CERN Council decision
- HL-LHC will run until ~ 2041

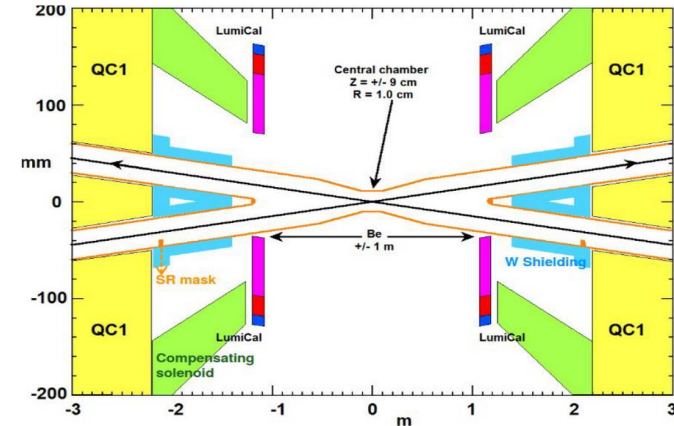
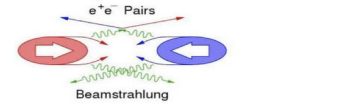
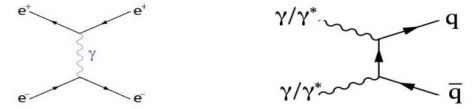
→ ANY future collider at CERN cannot start physics operation before ~ 2045 (but construction will proceed in parallel to HL-LHC operation)

**Care should be taken when comparing to other proposed facilities, for which in most cases only the (optimistic) technical schedule is shown.** In particular, studies related to **territorial implementation** (surface sites, roads, connection to water and electricity, environmental impact, admin procedures, etc.), which for FCC are being carried out in the framework of the Feasibility Studies, **take years.**

# Basic Detector Requirements

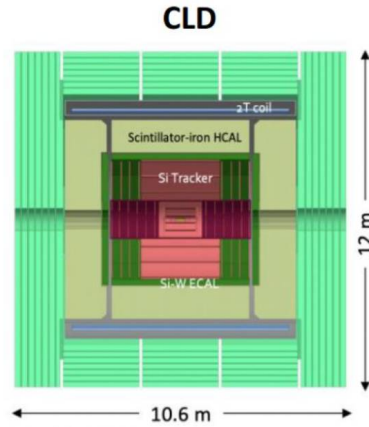
## Detector requirements - general considerations

- Requirements for Higgs and above have been studied to some extent by LC:
  - we want a detector that is able to withstand a **large dynamic range**:
    - in energy ( $\sqrt{s} = 90 - 365 \text{ GeV}$ )
    - in luminosity ( $L = 10^{34} - 10^{36} \text{ cm}^2/\text{s}$ )
- most of the **machine induced limitations are imposed by the Z pole run**:
  - large collision rates  $\sim 33 \text{ MHz}$  and continuous beams
    - no power pulsing possible
  - large event rates  $\sim 100 \text{ kHz}$ 
    - **fast detector response / triggerless** design challenging (but rewarding)
    - **high occupancy** in the inner layers/forward region (Bhabha scattering/ $\gamma\gamma$  hadrons)
  - beamstrahlung
- **complex MDI**: last focusing quadrupole is  $\sim 2.2\text{m}$  from the IP
  - magnetic field limited to  $B = 2\text{T}$  at the Z peak (to avoid disrupting vertical emittance/inst. Lumi via SR)
    - **limits the achievable track momentum resolution**
  - “anti”-solenoid
    - limits the acceptance to  $\sim 100 \text{ mrad}$

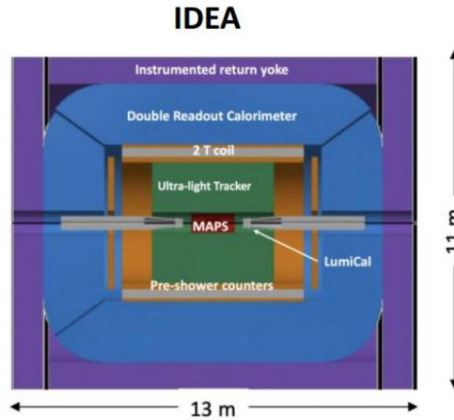


# Basic Detector Designs

## Detector Benchmarks



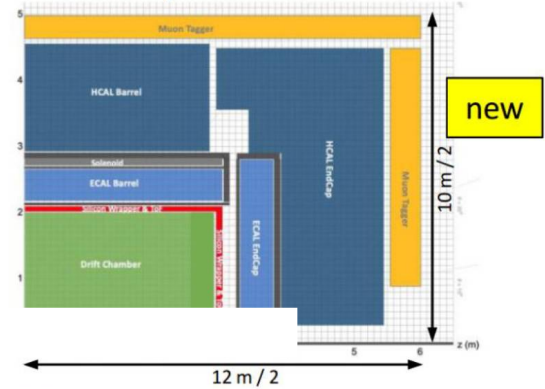
- Well established design
  - ILC -> CLIC detector -> CLD
- Full Si vtx + tracker;
- CALICE-like calorimetry;
- Large coil, muon system
- Engineering still needed for operation with continuous beam (no power pulsing)
  - Cooling of Si-sensors & calorimeters
- Possible detector optimizations
  - $\sigma_p/p$ ,  $\sigma_E/E$
  - PID ( $\mathcal{O}(10\text{ ps})$  timing and/or RICH)?



- A bit less established design
  - But still ~15y history
- Si vtx detector; ultra light drift chamber w powerful PID; compact, light coil;
- Monolithic dual readout calorimeter;
  - Possibly augmented by crystal ECAL
- Muon system
- Very active community
  - Prototype designs, test beam campaigns, ...

## ALLEGRO

**Noble Liquid ECAL based**



- A design in its infancy
- Si vtx det., ultra light drift chamber (or Si)
- High granularity Noble Liquid ECAL as core
  - Pb/W+LAr (or denser W+LKr)
- CALICE-like or TileCal-like HCAL;
- Coil inside same cryostat as LAr, outside ECAL
- Muon system.
- Very active Noble Liquid R&D team
  - Readout electrodes, feed-throughs, electronics, light cryostat, ...
  - Software & performance studies



# Computing / Software are crucial ingredient

## Assumptions and baseline needs<sup>1</sup>



- Integrated luminosities
  - Nominal: {90, 12, 5, 0.2, 1.5} ab<sup>-1</sup> at  $\sqrt{s} = \{91.2, 160, 240, 350, 365\}$  GeV
  - # of evts:  $3 \times 10^{12}$  visible Z decays,  $10^8$  WW events,  $10^6$  ZH events,  $10^6$  tt events

### ● Baseline event sizes / processing time for hadronic evts at Z

- DELPHES: 7.5 kB/evt, 0.4 s/evt
  - Full stat sample sizes: 30 PB,  $\approx 10^{10}$  s/core  $\approx 0.5$  MHS06<sup>2</sup>
- Full sim: CLD reference: 1 - 2 MB/evt, 10 s/evt
  - Full stat sample sizes: 3 EB,  $\approx 3 \cdot 10^{13}$  s/core  $\approx 10$ -15 MHS06<sup>2</sup> / detector<sup>3</sup>

(HL-)LHC is similar in scope.

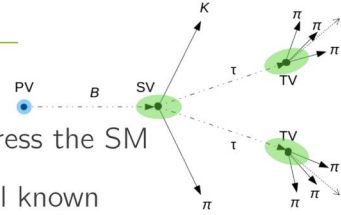
1. See also: GG, C Helsens: [EPJ Plus \(2022\) 137:30](#)
2. If done over a year, assuming similar number per each detector benchmark
  - a. CERN Openstack Core = 10-15 HEPSpec06 (HS06)
  - b. CERN OpenStack node used for tests: 16 cores, 32 GB RAM
3. Not applying to DELPHES, because in principle one sample can be re-adapted to other detector concepts

# Basic Measurements

Marina Nogueira, Ang Li, Michele Selvaggi,  
Lars Röhrig, Fabrizio Palla, Nicola De Filippis



# B Physics to benchmark vertex



## Vertex resolution

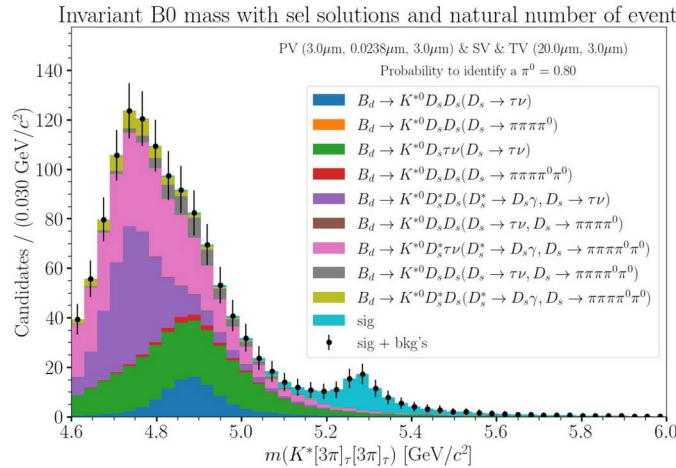
- Requires  $5 \rightarrow 3 \mu\text{m}$

## And BTW

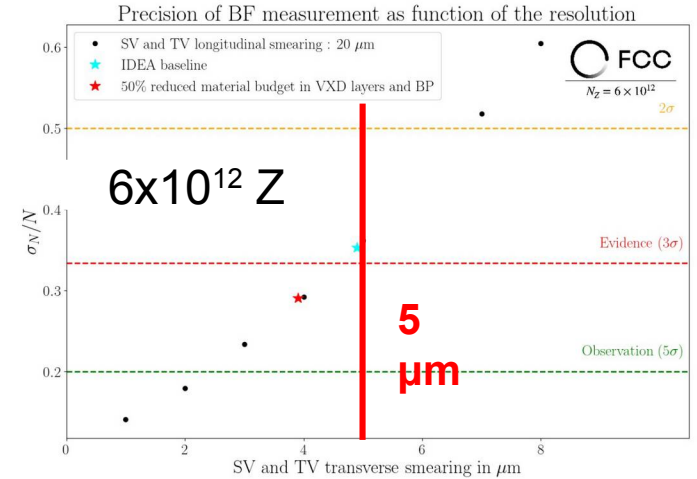
- FCC-ee has ~20 times more b and tau pairs than Belle II
- And the b/tau pairs are boosted

Vertex requirements: setting the stage with  $b \rightarrow s\tau^+\tau^-$

- EW penguin transitions of b quark in the SM very rare  $\rightarrow$  good laboratory to stress the SM
  - Third generation transitions in  $B^0 \rightarrow K^*\tau^+\tau^-$  couplings experimentally less well known
- $\rightarrow$  Feasibility depends on neutrino reconstruction  $\checkmark \rightarrow$  depends on **vertex precision**



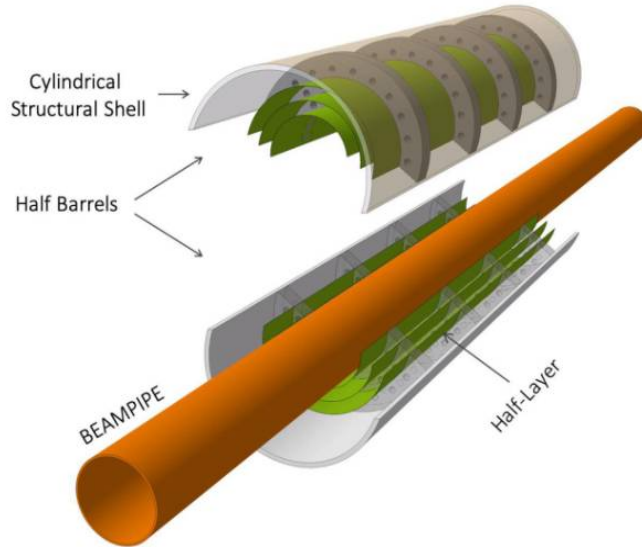
$\rightarrow$  Transverse vertex resolution of  $\mathcal{O}(5 \mu\text{m})$  required (limited by the **material budget of the beampipe**)



# Benchmark Vertex measurements

## MAPS

- Come to the workshop at CERN July 1/2



- A mini-workshop on vertex detector technologies (including system integration and mechanical aspects) will be held at CERN on July 1 and 2, with a lot of discussions:

<https://indico.cern.ch/event/1417976/>

***Lightweight layout using an ALICE ITS3 inspired design***

*(~0.05 %  $X/X_0$  material budget per layer – 5 times less than the Mid-Term one)*

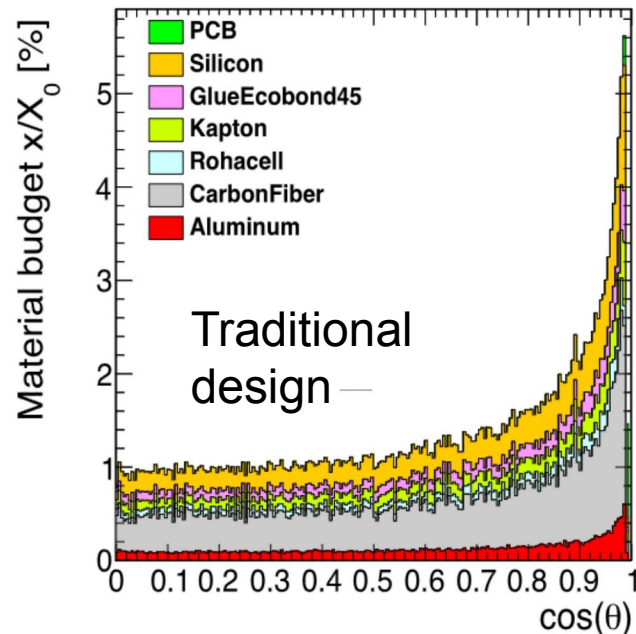
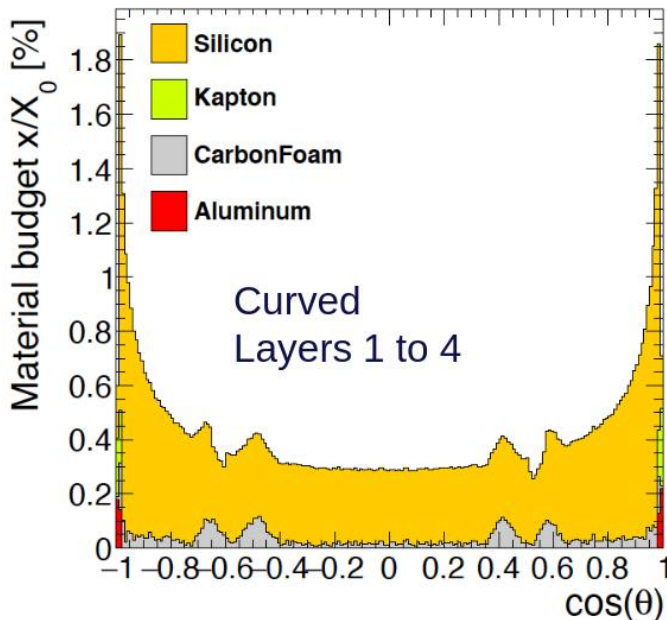
*After fruitful discussions with C. Gargiulo, A. Junique, G. Aglieri Rinella, W. Snoeys*

# Benchmark Vertex measurements

## MAPS

- Workshop at CERN July 1-2 was a great event to hear what people are doing ([link](#)).

## Material budget inner vertex



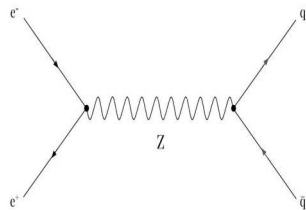
**~ 5  $\mu\text{m}$**  resolution, but want to further improve

EW

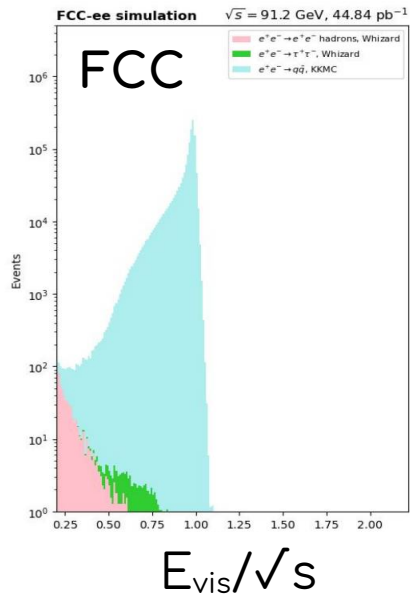
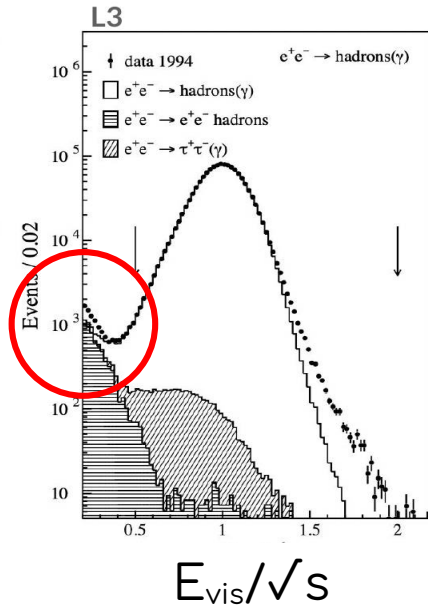
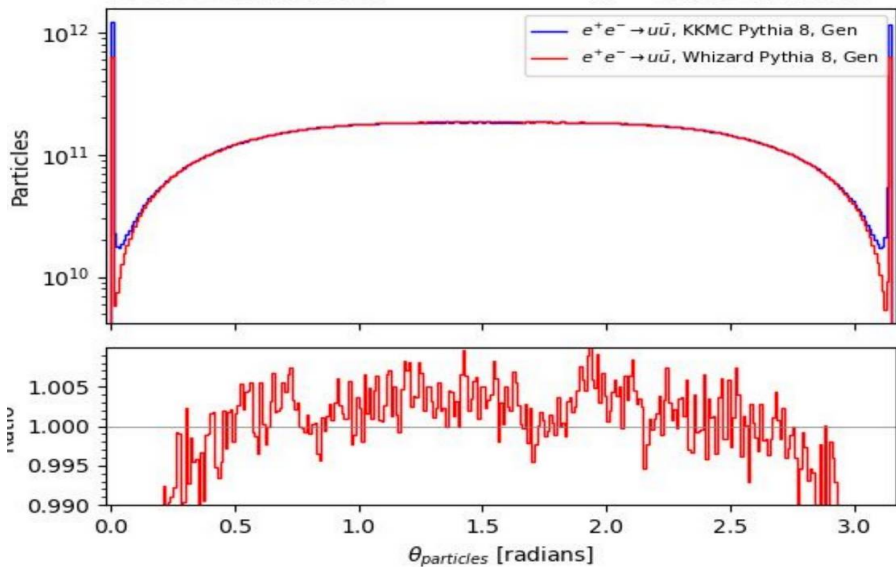
# Most copious process at FCC

$$e^+e^- \rightarrow qq$$

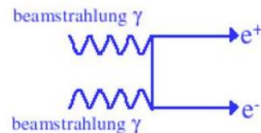
- Most basic MC needs work
- Acceptance not precise
- Beam and two photon background poorly modelled



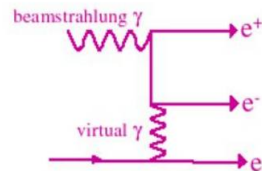
FCC-ee simulation  $\sqrt{s} = 91.2 \text{ GeV}, 75 \text{ ab}^{-1}$



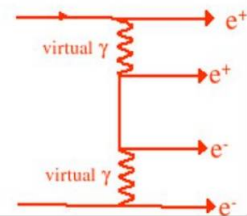
Breit-Wheeler



Bethe-Heitler



Landau-Lifshitz



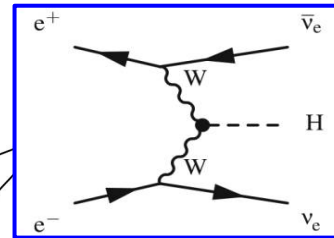
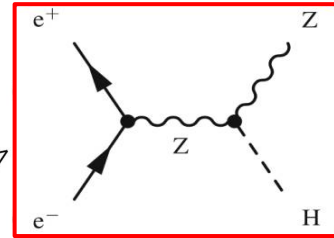
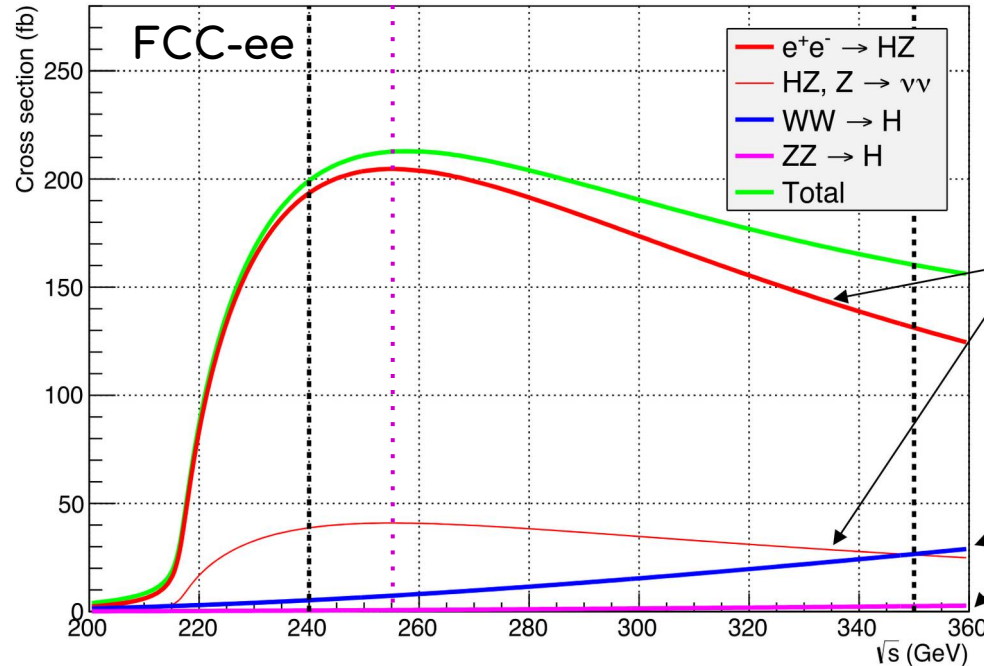
# Higgs Physics at $e^+e^-$ Colliders

ZH Threshold turns on at  $91 + 125 \text{ GeV} = 216 \text{ GeV}$  reaches a maximum at around  $255 \text{ GeV}$

Vector boson fusion rises steadily, but is small

FCC-ee:  
most Higgses at  $240 \text{ GeV}$   
considering lumi profile

Unpolarized cross sections



# Higgs Physics at $e^+e^-$ Colliders

## Leading strategy

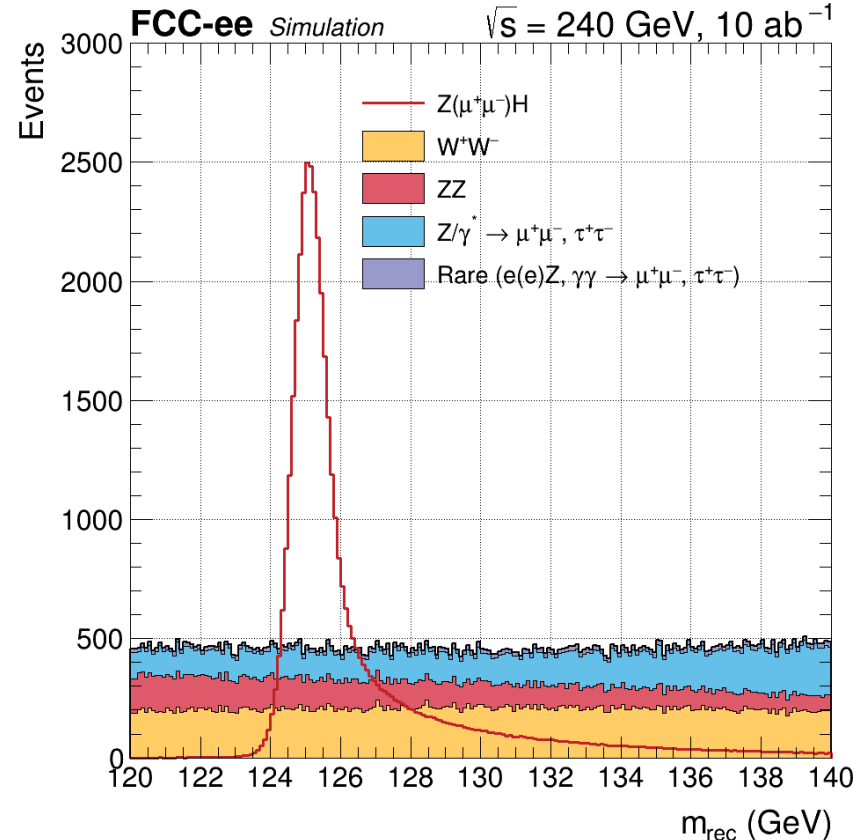
- Tag the Z boson (leptons or jets)
- Recoil mass peaks sharply at Higgs mass

$$\begin{aligned} m_{recoil}^2 &= (\sqrt{s} - E_{ff})^2 - p_{ff}^2 \\ &= s + m_Z^2 - 2E_{ff}\sqrt{s} \approx m_H^2 \end{aligned}$$

- Direct Higgs reconstruction not required, **model independent**  $\sigma_{ZH}$  measurement
- Dominant background: WW, ZZ and Z/ $\gamma^*$

## Challenges

- Detectors: resolution, tracking, vertexing, timing, angular
- Flavour tagging for Higgs couplings
- Jet reconstruction algorithms



This plot does not work at hadron colliders.



# Basic Higgs Properties

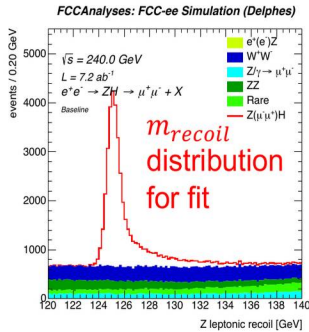
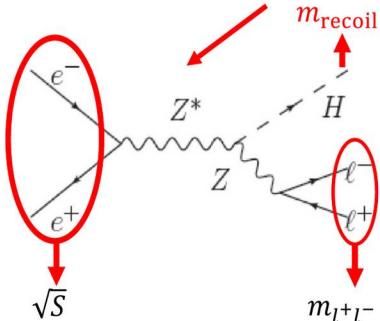
## Higgs mass

- Basic SM parameter
- Not a limiting factor for radiative corrections
- Essential for producing Higgs directly  $e^+e^- \rightarrow H$
- Widths 4.1 MeV

## Higgs mass

- ❖ Current best from LHC  $\delta m_H \sim 100 \text{ MeV}$
- ❖ At FCC-ee, Higgs mass will reach **MeV level accuracy**, ( $\Gamma_H \sim 4.1 \text{ MeV}$ )
- ❖ Electron and Muons final states:  $e^+e^- \rightarrow ZH \rightarrow l^+l^- + XX$ , ( $Z \rightarrow \mu^+\mu^-, e^+e^-$ )
- ❖  $M_{recoil}$  from the Z production without measuring the Higgs production final state

$$m_{recoil}^2 = (\sqrt{s} - E_{l\bar{l}})^2 - p_{l\bar{l}}^2 = s - 2E_{l\bar{l}}\sqrt{s} + m_{l\bar{l}}^2$$

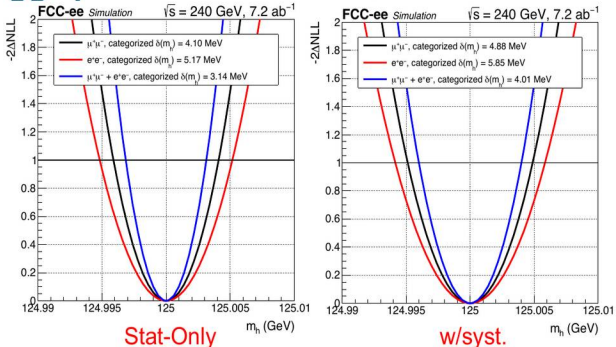


### Higgs mass, Fit with analytic shape

- Signal Shape: 2 Crystal-Ball with Gaussian core
- Backgrounds modelled as polynomial (3rd order)
- Signal and background injected in Combine,  $m_H$  as POI

$\sqrt{s} = 240 \text{ GeV}$   
 $L = 7.2 \text{ ab}^{-1}$

Gregorio Bernardi  
 Jan Eysermans  
 Ang Li  
 DOI [10.17181](https://doi.org/10.17181)



### Uncertainty Stat-Only, and w/ systematics:

➤ Higgs mass: **3.1 MeV  $\rightarrow$  4.0 MeV**

### Dominant Syst. Unc. :

**Centre-of-mass with  $\sim 2 \text{ MeV}$**



# Basic Higgs Properties

## Higgs mass

- Basic SM parameter
- Not a limiting factor for radiative corrections
- Essential for producing Higgs directly  $e^+e^- \rightarrow H$
- Width 4.1 MeV

## Higgs Mass – Detector Requirements

Extended studies performed regarding detector/accelerator effects on the Higgs mass

→ Looking at impact on  $m_H$  uncertainty stat. (stat.+syst.) in MeV

### Nominal configuration

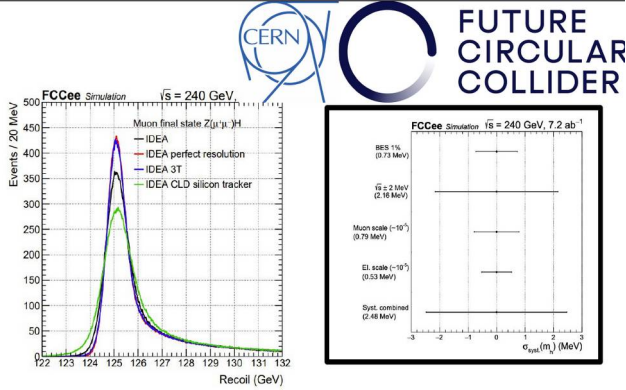
Crystal ECAL to Dual Readout

Nominal 2 T → field 3 T

IDEA drift chamber → CLD Si tracker

Impact of Beam Energy Spread uncertainties

Perfect (=gen-level) momentum resolution



Fit configuration	$\mu^+\mu^-$ channel	$e^+e^-$ channel	combination
Nominal	4.10 (4.88)	5.17 (5.85)	3.14 (4.01)
Inclusive	4.84 (5.53)	6.16 (6.73)	3.75 (4.50)
Degradation electron resolution (*)	4.10 (4.88)	5.98 (6.49)	3.32 (4.11)
Magnetic field 3T	3.38 (4.28)	4.30 (5.00)	2.60 (3.54)
CLD 2T (silicon tracker)	5.51 (6.07)	6.20 (6.70)	4.01 (4.66)
BES 6% uncertainty	4.10 (5.01)	5.17 (6.10)	3.14 (4.09)
Disable BES	2.27 (3.42)	3.11 (4.04)	1.80 (2.99)
Ideal resolution	2.89 (3.95)	3.89 (4.56)	2.39 (3.33)
Freeze backgrounds	4.10 (4.88)	5.17 (5.85)	3.14 (4.00)
Remove backgrounds	3.37 (4.34)	3.85 (4.80)	2.49 (3.56)

# Momentum resolution

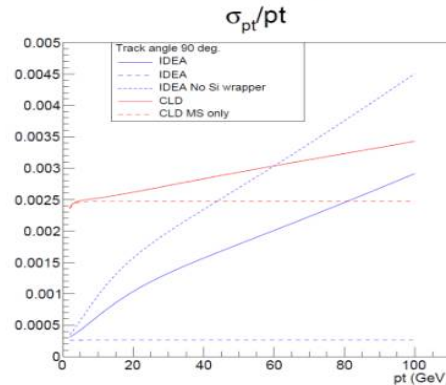
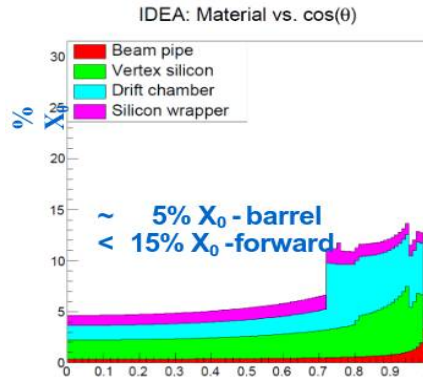
## Requirements on track momentum resolution

### Minimal material matters

- Drift chamber is ultra light
- Silicon detectors could be as light?
- Larger radius improves resolution ...
- Higher magnetic field improves resolution: 2T to 3T improves momentum 50% and mass by 14%

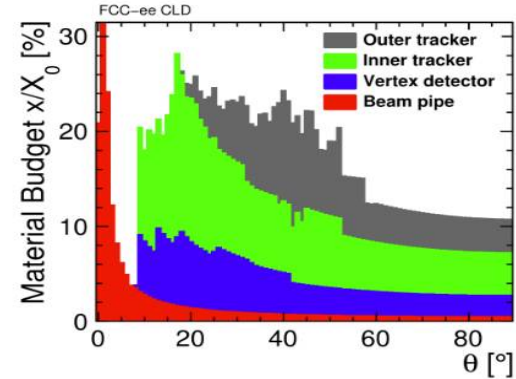
The IDEA Drift Chamber is designed to cope with transparency

- a unique-volume, high granularity, fully stereo, low-mass cylindrical
- gas: He 90% -  $iC_4H_{10}$  10%
- inner radius 0.35m, outer radius 2m
- length  $L = 4m$



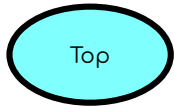
The CLD silicon tracker is made of:

- six barrel layers, at radii ranging between 12.7 cm and 2.1 m, and of eleven disks.
- the material budget for the tracker modules is estimated to be 1.1 – 2.1% of a radiation length per layer



For 10 GeV (50 GeV)  $\mu$  emitted at an angle of  $90^\circ$  w.r.t the detector axis, the  $p_T$  resolution is

- about 0.05 % (0.15%) with the very light IDEA DCH
- about 0.25% (0.3%) with the CLD full silicon tracker, being dominated by the effect of MS



# Basic Top Properties



## Top mass

- Basic SM parameter
- Top never directly produced at lepton collider
- Hadron colliders have problematic definition of mass
- Theoretically much cleaner access at lepton colliders

## Top Threshold

### Current run plan at the top threshold

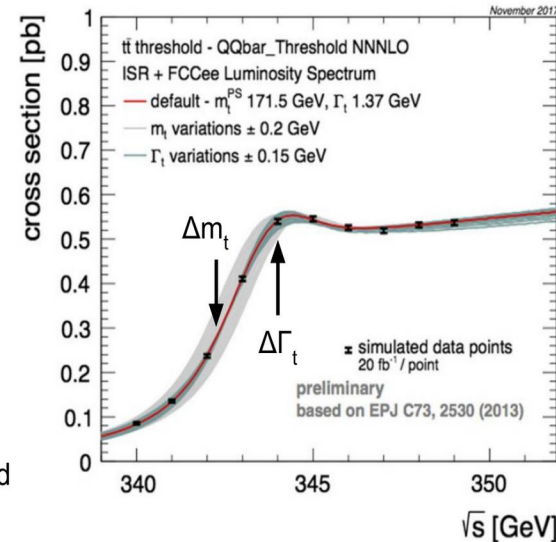
- 1 year threshold scan 340–350 GeV: total  $\sim 1.4 \text{ ab}^{-1}$
- 4 years at 365 GeV: total  $\sim 2.3 \text{ ab}^{-1}$

### Threshold scan to extract the Top mass and width (similar as WW)

- Relative large uncertainty on top mass ( $\pm 0.5 \text{ GeV}$  from HL-LHC)
- Need to constrain shape in optimal way
- Possible to constrain backgrounds (below) and  $t\bar{t}H$  (above)
- Multipoint scan in 5 GeV window [340, 345], each  $\sim 25 \text{ /fb}$  to be studied

### At 365 GeV, with $2.3 \text{ ab}^{-1}$

- Top properties
- Higgs properties ( $ee \rightarrow \nu\nu H$ ): total cross-section, couplings, width



→  $\Delta m_t \text{ (stat)} \sim 17 \text{ MeV}$   
 →  $\Delta \Gamma_t \text{ (stat)} \sim 45 \text{ MeV}$

*Higgs Couplings  
beyond the third  
generation fermions*

David d'Enterria, Francis Petriello, Loukas Gouskos, Michele Selvaggi,  
Daniel Elvira, Xunwu Zuo



# Why measure Higgs couplings?

## BSM O(1TeV): Impact on H-couplings

Model	$b\bar{b}$	$c\bar{c}$	$gg$	$WW$	$\tau\tau$	$ZZ$	$\gamma\gamma$	$\mu\mu$
MSSM [40]	+4.8	-0.8	-0.8	-0.2	+0.4	-0.5	+0.1	+0.3
Type II 2HD [42]	+10.1	-0.2	-0.2	0.0	+9.8	0.0	+0.1	+9.8
Type X 2HD [42]	-0.2	-0.2	-0.2	0.0	+7.8	0.0	0.0	+7.8
Type Y 2HD [42]	+10.1	-0.2	-0.2	0.0	-0.2	0.0	0.1	-0.2
Composite Higgs [44]	-6.4	-6.4	-6.4	-2.1	-6.4	-2.1	-2.1	-6.4
Little Higgs w. T-parity [45]	0.0	0.0	-6.1	-2.5	0.0	-2.5	-1.5	0.0
Little Higgs w. T-parity [46]	-7.8	-4.6	-3.5	-1.5	-7.8	-1.5	-1.0	-7.8
Higgs-Radion [47]	-1.5	-1.5	+10.	-1.5	-1.5	-1.5	-1.0	-1.5
Higgs Singlet [48]	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5	-3.5

[1708.08912](#)

$$\frac{v^2}{\Lambda^2} \sim \frac{6\%}{\Lambda^2(\text{TeV})}$$

e.g.  $\Lambda=1$  (5)TeV  $\rightarrow$   $\sim 5$  (0.1)%

### ■ HL-LHC:

◆ Direct searches: O(5) TeV

◆ H-couplings:

- Bosons/ 3<sup>rd</sup>-Gen fermions @ few %
- 2<sup>nd</sup> Gen fermions: maybe evidence of  $H \rightarrow cc$
- Self-coupling  $\sim 50\%$

### ■ Future $e^+e^-$ collider:

◆ Measure H-couplings at O(0.1)% level

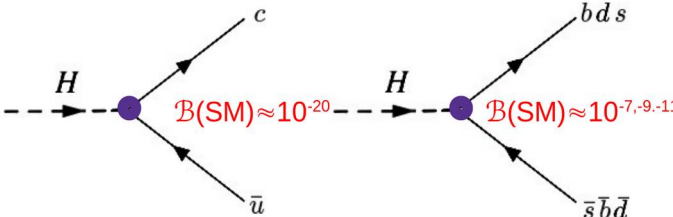
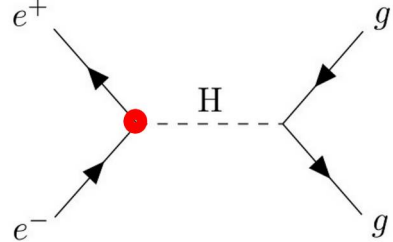
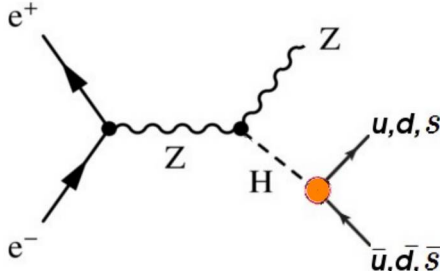
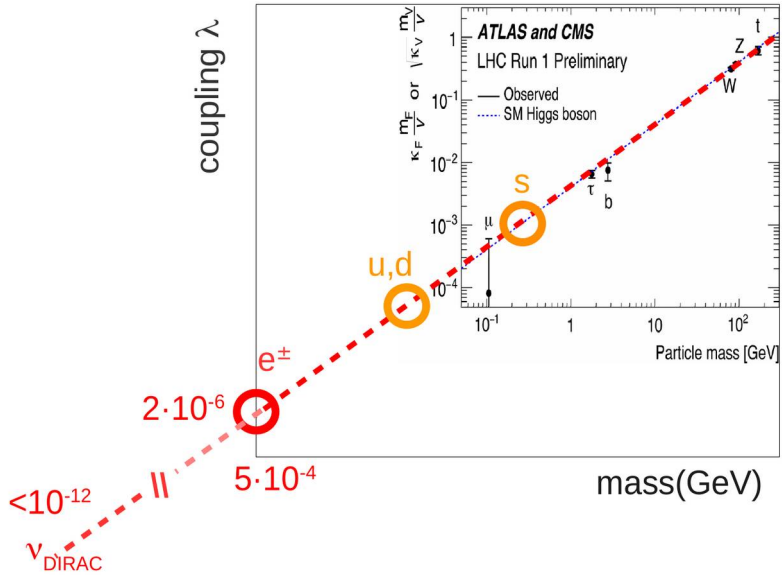


# Higgs to first generation

## Unresolved Higgs decay modes

- Muons are only second generation seen
- Lighter fermions are very difficult

■ Do the lightest fermions (u,d,s,e) acquire their masses through their Higgs (Yukawa) couplings?



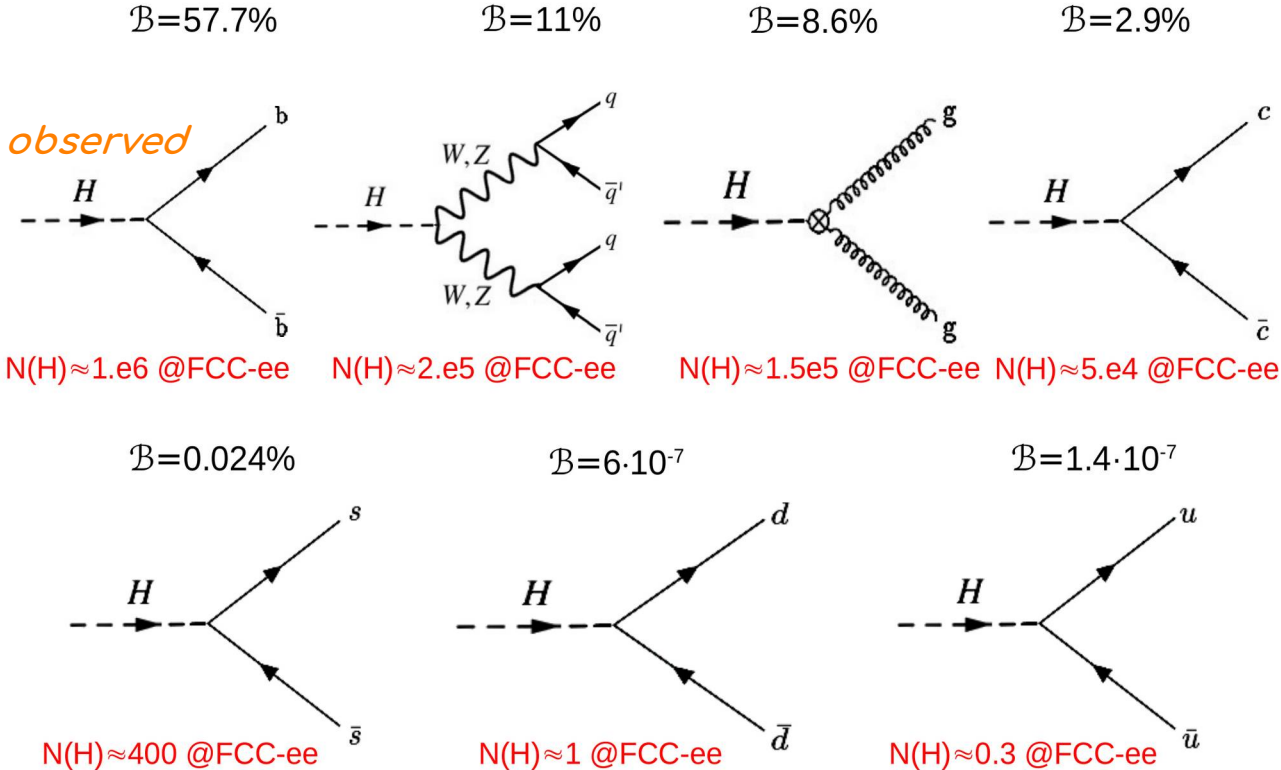
■ Does the Higgs boson mediate  $H \rightarrow qq'$  FCNCs at tree level?

# Hadronic final states dominate

## Why is this important?

- At LHC those are often hopeless – background
- FCC-ee offers cleaner environment, more handles and data calibration

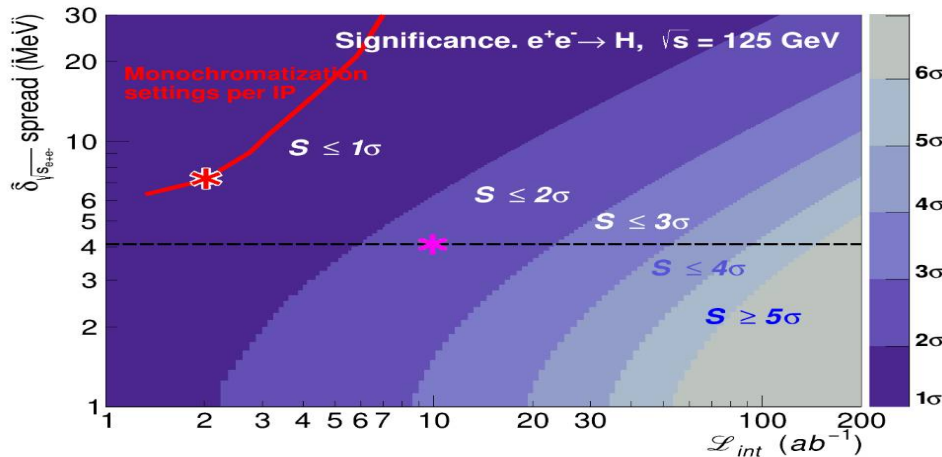
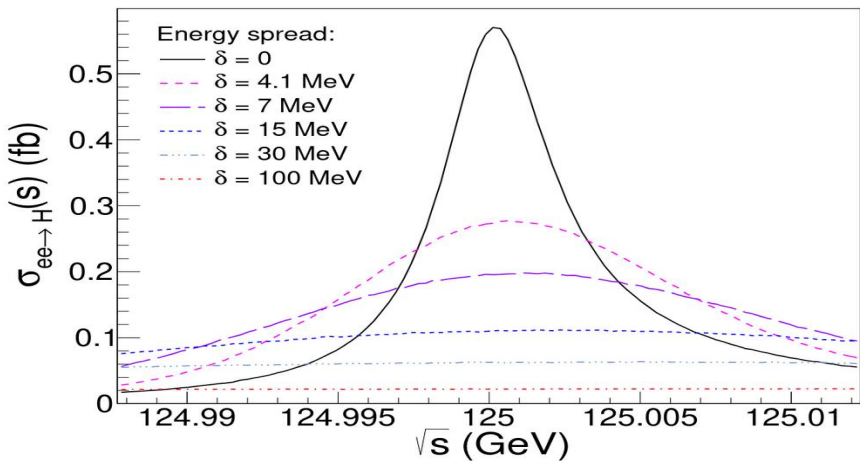
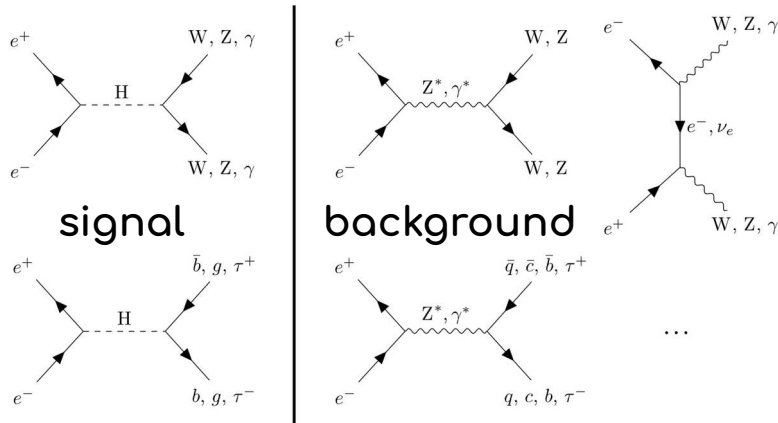
■ 80% of the Higgs decays are **fully hadronic**. Mostly measurable at FCC-ee!



# Higgs Electron Yukawa Coupling

Measure  $e^+e^- \rightarrow H \rightarrow e^+e^-$ : how?

- $\Gamma_H$  is 4.1 MeV, measure  $m_H$  at MeV level
- Dial collider  $E_{CM}$  to  $m_H$ , *precisely!*
- Monochromatize energy:  $\sim 4$  MeV spread
- Signal is tiny and background is very large
- **1.3 std significance per IP and per year**



# Higgs to gluon gluon

## Gluon tagging

- Major progress in tagging makes it feasible
- H→gg has no continuum background
- But can we distinguish well enough between u,d,s and gluons?

■ No e<sup>+</sup>e<sup>-</sup> background can generate 2 true gluon jets !

■ Analysis performances assumed:  
 2 gluon-tagged jets (with 70% effic. each)  
 u,d,s mistagging rate: ~1%  
 Challenging, but not impossible (see next)  
 Retains 50% of σ(H→gg) = 24 ab signal

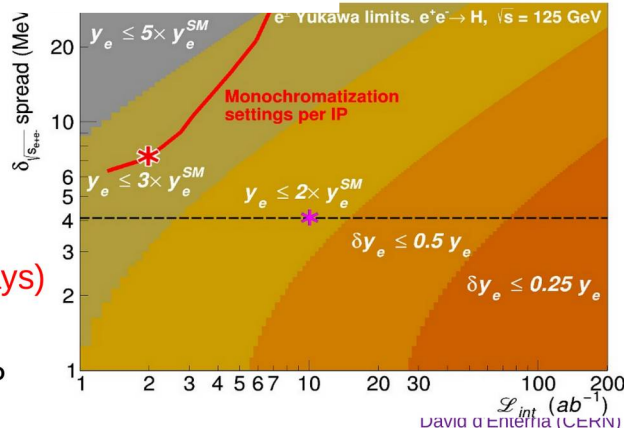
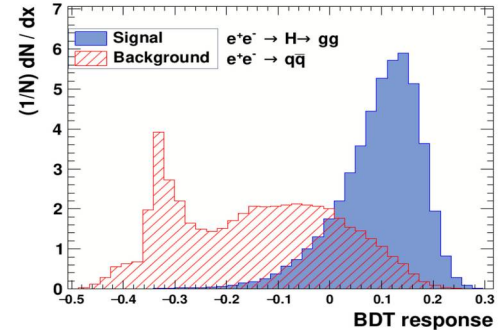
■ BDT MVA result (removing jet vars. potentially already used in g-uds discrimination):

Signal reduction ~50%  
 Backgd. reduction: x17

■ For  $\mathcal{L}_{int} = 10 \text{ ab}^{-1}$  :  
 $S/\sqrt{B} = 55/\sqrt{2500} \approx 1.1$   
 Significance  $\approx 1.1\sigma$  (1.3σ, other decays)

With current best monochromatization:  
 $y_e < 2.5 \times y_{e,SM}$  (95% CL) per year & per IP

[DdE,Poldaru/Wojcik arXiv:2107.02686]





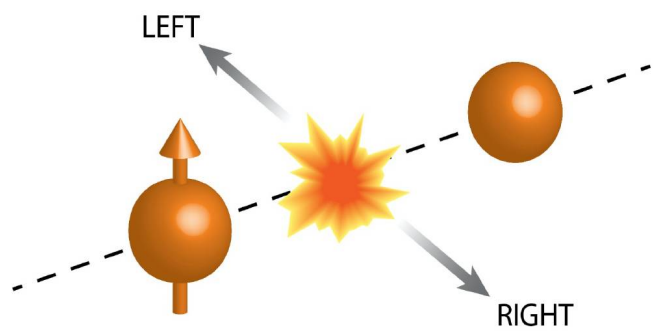
# $e^+e^- \rightarrow$ Higgs with polarized beams

## Beam polarization

- Transverse is more obvious
- 80% not unreasonable
- Longitudinal much less clear
- Needs polarimeter: expensive
- Can work: 30%?

## Transverse spin asymmetries

- The idea is to use transverse spin asymmetries to increase the sensitivity to the electron Yukawa coupling. We consider the following observables in our study.



$$A = \frac{N}{D}$$

Electron polarized, positron unpolarized (SP <sup>0</sup> ):	$N = \frac{1}{2}(\sigma^{+0} - \sigma^{-0})$
	$D = \frac{1}{2}(\sigma^{+0} + \sigma^{-0})$
Electron transversely polarized, positron longitudinally polarized (DP):	$N = \frac{1}{4}(\sigma^{++} - \sigma^{+-} - \sigma^{-+} + \sigma^{--})$
	$D = \frac{1}{4}(\sigma^{++} + \sigma^{+-} + \sigma^{-+} + \sigma^{--})$
Electron transversely polarized, positron longitudinally polarized (SP <sup>+</sup> ):	$N = \frac{1}{2}(\sigma^{++} - \sigma^{-+})$
	$D = \frac{1}{2}(\sigma^{++} + \sigma^{-+})$
Electron transversely polarized, positron longitudinally polarized (SP <sup>-</sup> ):	$N = \frac{1}{2}(\sigma^{+-} - \sigma^{--})$
	$D = \frac{1}{2}(\sigma^{+-} + \sigma^{--})$



# $e^+e^- \rightarrow$ Higgs with polarized beams

## Single Spin Asymmetry

## Theoretical structure of transverse SSAs

- Imaginary part in amplitude: interference
- Requires resonance (Higgs)

- The structure of transverse SSAs is dictated by the discrete symmetries of the SM.

Two key points:

$$S_T \cdot p_q = \beta_q \frac{\sqrt{s}}{2} \sin(\theta) \cos(\phi),$$

$$\epsilon(p_e, p_{\bar{e}}, p_q, S_T) = -\beta_e \beta_f \frac{s^{3/2}}{4} \sin(\theta) \sin(\phi)$$

$$S_T \cdot p_q \Rightarrow \text{P odd, } A_t \text{ even}$$

$$\epsilon(p_e, p_{\bar{e}}, p_q, S_T) \Rightarrow \text{P even, } A_t \text{ odd}$$

1. These two structures have different azimuthal dependence (orientation between final-state bottom quark and transverse spin direction); they can be separated by weighting the final-state phase-space integral

2. To get a structure odd under  $A_t$  we need an imaginary part in an amplitude. At tree-level this can only come when we are on a particle resonance

$$\frac{1}{s - M^2 + iM\Gamma}$$





# $e^+e^- \rightarrow$ Higgs with polarized beams

## Origin

## Application to the $ee \rightarrow bb$ process

- ZH interference
- Does not work for  $H \rightarrow gg$  !
- Term is proportional to mass!
- Azimuthal structure is different!

- Study the structure of the asymmetry numerator (DP in this example). Three diagrams contribute at tree-level: s-channel photon, Z-boson, and Higgs exchange.

$$N = \frac{1}{2s} \int dLIPS \left\{ \frac{R_{\gamma\gamma}}{s^2} + \frac{R_{ZZ}}{(s - M_Z^2)^2} + \frac{R_{\gamma Z}}{s(s - M_Z^2)} + \frac{R_{\gamma H}(s - M_H^2)}{s[(s - M_H^2)^2 + M_H^2\Gamma_H^2]} + \frac{R_{ZH}(s - M_H^2) + I_{ZH}M_H\Gamma_H}{(s - M_Z^2)[(s - M_H^2)^2 + M_H^2\Gamma_H^2]} \right\}$$

$$R_{\gamma\gamma} = 96e^4 Q_e^2 Q_q^2 m_e (S_T \cdot p_q)(t - u)$$

$$R_{ZZ} = 96m_e (S_p \cdot p_b) g_Z^4 g_{ve}^2 (g_{vq}^2 + g_{aq}^2)(t - u) + 192m_e (S_T \cdot p_q) g_Z^4 g_{ve} g_{ae} g_{vq} g_{aq} s$$

$$R_{\gamma Z} = 192e^2 g_Z^2 Q_e Q_q m_e (S_T \cdot p_b) g_{ve} g_{vq}(t - u) + 96e^2 g_Z^2 Q_e Q_u m_e (S_p \cdot p_q) g_{ae} g_{aq} s$$

$$R_{\gamma H} = -96e^2 Q_e Q_q y_e y_q (S_T \cdot p_q) m_q$$

$$R_{ZH} = -96g_Z^2 g_{ve} g_{vq} y_e y_q (S_T \cdot p_q) s$$

$$I_{ZH} = -192g_Z^2 g_{ae} g_{vq} y_e y_q m_q \epsilon(p_e, p_{\bar{e}}, p_q, S_T).$$

- Comes from the imaginary part of the Higgs propagator and is enhanced by a factor of  $M_H/\Gamma_H$ .
- All terms are suppressed **linearly** by the electron mass; this structure is directly proportional to the electron Yukawa couplings
- Can be isolated due to its different azimuthal structure, which follows from the discussion on the previous slide

# $e^+e^- \rightarrow \text{Higgs with polarized beams}$

to MeV from resonance invariant mass cut

$b\bar{b}$

Observable	$e^-e^+ \rightarrow b\bar{b}$
DP	0.27
SP <sup>0</sup>	0.19
SP <sup>+</sup>	0.11
SP <sup>-</sup>	0.38
Reference	0.11

Definite improvement using transverse polarization; further improvement if the second beam can be longitudinally polarized

Obtained using unpolarized cross section; in good agreement with  $S/\sqrt{B}=0.13$  in 2107.02686

Observable	$e^-e^+ \rightarrow b\bar{b}$
DP	0.41 (39%)
SP <sup>0</sup>	0.30 (33%)
SP <sup>+</sup>	0.17 (44%)
SP <sup>-</sup>	0.58 (39%)

Second column gives polar angle cut in terms of percentage of phase spaced removed

Observable	$e^-e^+ \rightarrow WW \rightarrow ll\nu\nu$
DP	0.45
SP <sup>0</sup>	0.80
SP <sup>+</sup>	2.9
SP <sup>-</sup>	0.33
Reference	0.45

Over a factor of 6 improvement if  $(P_T, P_L) = (80, 30)\%$  can be obtained

$W^+W^-$

Quoted are significance of the signal.

**Major improvements of up to factors of 6 possible for  $b\bar{b}$  and  $WW$**

Obtained using unpolarized cross section;  $S/\sqrt{B}=0.53$  in 2107.02686, likely due to use of BDT rather than simple cuts

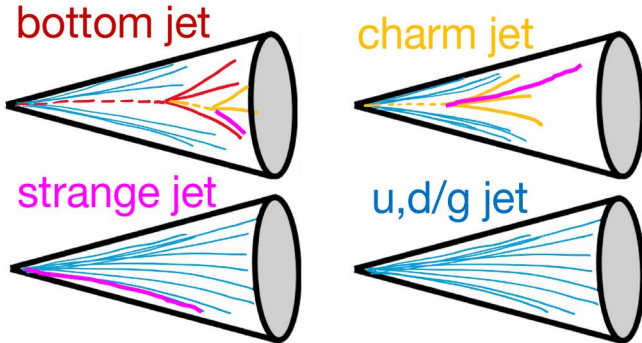
# Tagging Challenge: Higgs $\rightarrow$ $ss$

## Jet tagging

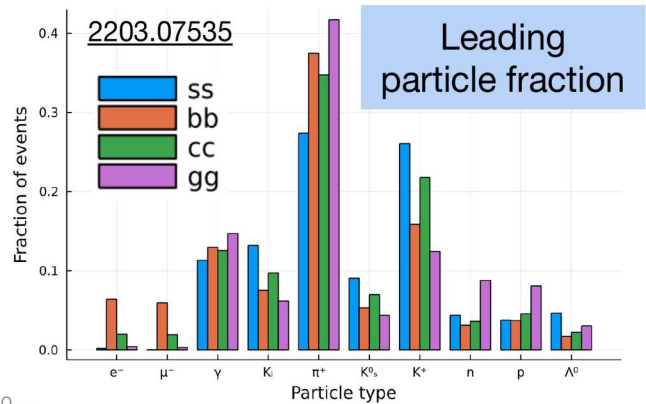
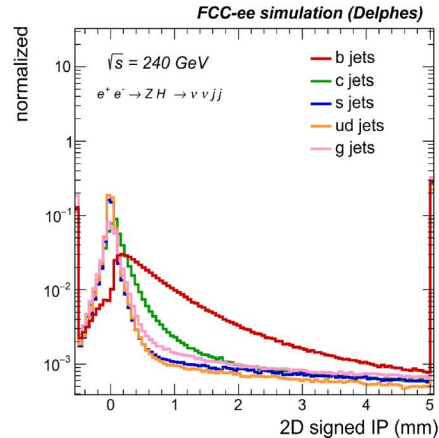
- Key to quarks and gluons
- Very different from LHC
- Huge Z and W boson decay samples to calibrate
- PID is crucial input
- Charm, strange gluon tagging works
- Seeing  $H \rightarrow ss$  is least obvious, but should be possible



### Basics



- **Bottom/charm tagging**
  - ◆ Large lifetime
  - ◆ Displaced vertices/tracks
  - ◆ Non-isolated  $e/\mu$
- **Strange tagging**
  - ◆ Enhanced Kaon fraction
  - ◆ Large momentum fraction



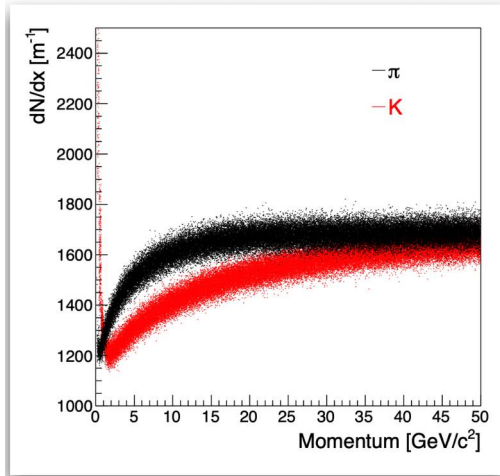
# Strange Tagging needs PID



## Handles for PID

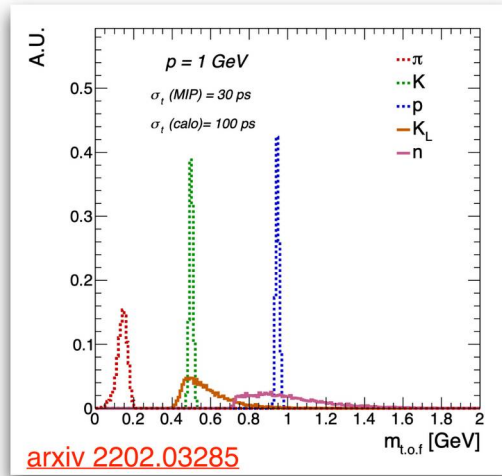
$dN/dx$  or  $dE/dx$ ,

- Ionization cluster count or energy per path length
- Good separation in wide momentum range
- “Blind” region around 1 GeV
- Currently assume 2% resolution



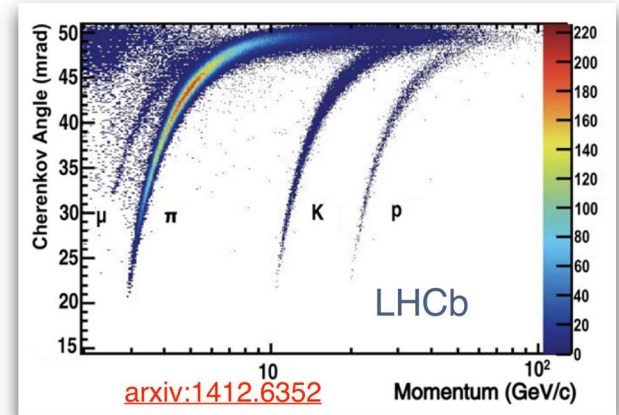
Time of flight

- Good separation at low momentum ( $\sim 1$  GeV)
- Requires  $\sim 100$  ps resolution to cover PID  $\sim 1$  GeV
- Current studies assume 30 ps resolution



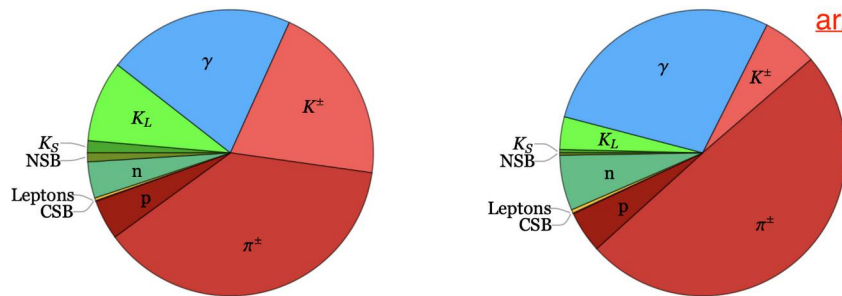
RICH

- Good separation in a wide momentum range
- Need enough radiation length for good PID



# Tagging Challenge: strange

## Strange jet tagging



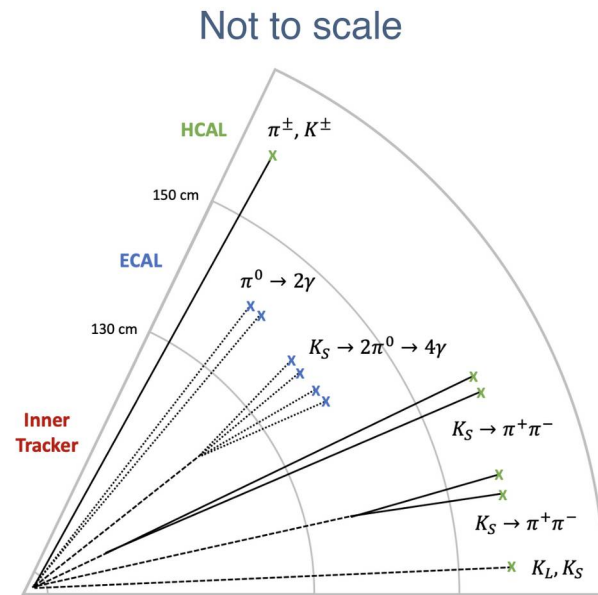
[arxiv:2003.09517](https://arxiv.org/abs/2003.09517)

Momentum fraction

Strange  $p_T = 45 \text{ GeV}$

Down  $p_T = 45 \text{ GeV}$

- Higher fraction of momentum carried by **kaons**
  - $K^+/\pi^+$  separation is the key
- Neutral kaons and s-baryons are long-lived
  - $c\tau(b/c) \approx 0.5 \text{ mm}$ ,  $c\tau(s) \approx 50 \text{ mm}$
  - Requirement on vertexing, see talk by L. Roerig



Not to scale

s-baryons  $\Lambda, \Sigma, \Xi$  have  $c\tau \approx 1 - 10 \text{ cm}$



# Tagging Challenge: Higgs $\rightarrow$ $ss$

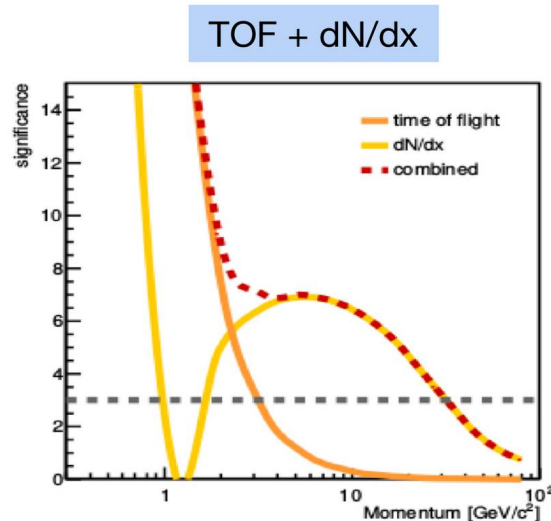
## Jet tagging

- Key to quarks and gluons
- Very different from LHC
- Huge Z and W boson decay samples to calibrate
- PID is crucial input
- Charm, strange gluon tagging works
- Seeing  $H \rightarrow ss$  is least obvious, but should be possible



## Strange tagging: Particle ID

- Big effort to design optimal PID detectors and algorithms to exploit their full potential [e.g., ECFA  $H \rightarrow ss$  team, [Wiki](#)]
  - ♦ IDEA detector:



Achieve  $3\sigma$   $\pi/K$  separation for up to  $\sim 30$  GeV momenta

### But:

We need to carefully assess impact of detector proposals to the full Higgs [and not only] physics program in general [more later]

# Tagging Challenge: Higgs → ss

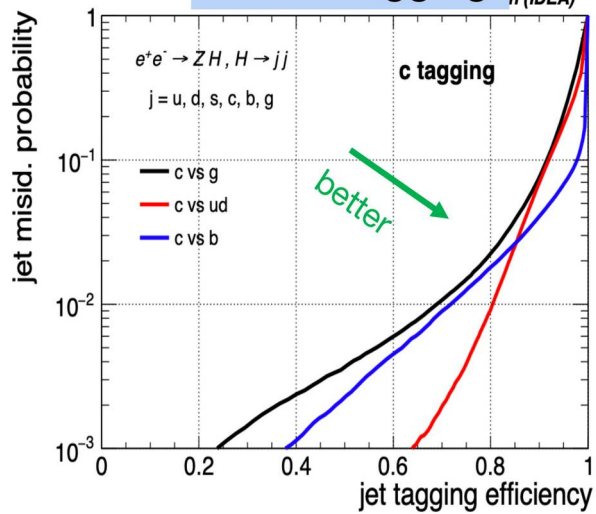
## Jet tagging



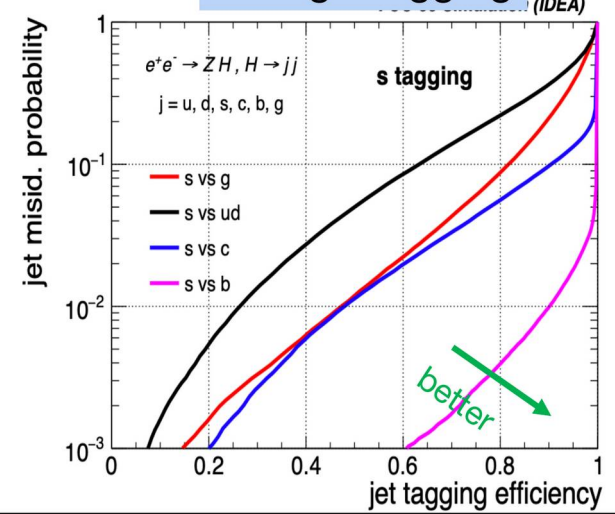
## Jet tagging: Performance

- Key to quarks and gluons
- Very different from LHC
- Huge Z and W boson decay samples to calibrate
- PID is crucial input
- Charm, strange gluon tagging works
- Seeing H → ss is least obvious, but should be possible

charm-tagging  $n$  (IDEA)



strange-tagging (IDEA)



Eff (c)	Mistag (g)	Mistag (ud)	Mistag (b)
90%	7%	7%	4%
80%	2%	0.8%	2%

Eff (s)	Mistag (g)	Mistag (ud)	Mistag (c)	Mistag (b)
90%	20%	40%	10%	1%
80%	9%	20%	6%	0.4%

# Tagging Challenge: Higgs $\rightarrow$ ss

## Jet tagging

- Key to quarks and gluons
- Very different from LHC
- Huge Z and W boson decay samples to calibrate
- PID is crucial input
- Charm, strange gluon tagging works
- Seeing H  $\rightarrow$  ss is least obvious, but should be possible

$E_{\text{CM}} = 240 \text{ GeV} [10.8 \text{ ab}^{-1}, 4 \text{ IP}]$

Decay mode	Z( $\rightarrow$ LL)H( $\rightarrow$ jj) [%]	Z( $\rightarrow$ vv)H( $\rightarrow$ jj) [%]	Z( $\rightarrow$ jj)H( $\rightarrow$ jj) [%]	Combination
H $\rightarrow$ bb	0.55	0.24	0.20	0.15
H $\rightarrow$ cc	3.35	1.77	2.38	1.20
H $\rightarrow$ ss	280	93	296	80
H $\rightarrow$ gg	1.86	0.75	1.63	0.65

$E_{\text{CM}} = 365 \text{ GeV} [2.3 \text{ ab}^{-1}, 4 \text{ IP}]$

Decay mode	Z( $\rightarrow$ LL)H( $\rightarrow$ jj) [%]	Z( $\rightarrow$ vv)H( $\rightarrow$ jj) [%]	Z( $\rightarrow$ jj)H( $\rightarrow$ jj) [%]	Combination
H $\rightarrow$ bb	1.23	0.68	0.52	0.39
H $\rightarrow$ cc	8.20	3.95	4.68	2.83
H $\rightarrow$ ss	1153	214	664	201
H $\rightarrow$ gg	4.24	2.51	4.15	1.92

# Calorimeter resolution matters ...

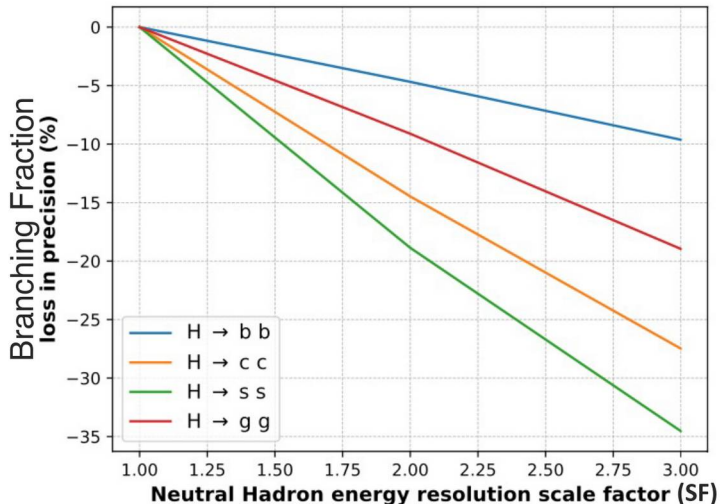
## Reconstruction of Higgs hadronic final states

(Case study 1)



Higgs  $\rightarrow$  2 jets signal ID in HZ events relies on the calorimeter (and vertex detector) performance: Mass resolution of Higgs and recoil system, flavor tagging efficiency

Study to measure impact of variation in neutral hadron resolution by a factor of 2 (3) with respect to the baseline) on  $H \rightarrow$  jet-jet, with jet = b, c, s, g, with  $Z \rightarrow$  lepton-lepton



Precision of  $H \rightarrow s\bar{s}$  degrades by 20% (35%)

- A bit larger than similar degradation in the number of ionization clusters per unit length ( $dN/dx$ ) – IDEA gas chamber ( $dN/dx$  provides particle ID)

The effect the  $H_{cc}$ ,  $H_{gg}$ ,  $H_{bb}$  couplings is smaller

- Increases as the s/b decreases

SF=1 ( dual readout calorimeter: **30%** $\sqrt{E}$  )

2 ( ATLAS type-calorimeter: **50%** $\sqrt{E}$  )

3 ( CMS-type calorimeter: **100%** $\sqrt{E}$  )

# Calorimeter technologies match

## Single particle, jet, and invariant mass resolution



Expected energy resolution for the different technologies: measurements when available, otherwise obtained from (DELPHI) simulation. Those values marked with “?” are estimates since neither measurement nor simulation exists

Detector technology (ECAL & HCAL)	E.m. energy res. stochastic term	E.m. energy res. constant term	ECAL & HCAL had. energy resolution (stoch. term for single had.)	ECAL & HCAL had. energy resolution (for 50 GeV jets)	Ultimate hadronic energy res. incl. PFlow (for 50 GeV jets)
Highly granular Si/W based ECAL & Scintillator based HCAL	15 – 17 % [12,20]	1 % [12,20]	45 – 50 % [20,45]	≈ 6 % ?	4 % [20]
Highly granular Noble liquid based ECAL & Scintillator based HCAL	8–10 % [24,27,46]	< 1 % [24,27,47]	≈ 40 % [27,28]	≈ 6 % ?	3–4 % ?
Dual-readout Fibre calorimeter	11 % [48]	< 1 % [48]	≈ 30 % [48]	4–5 % [49]	3–4 % ?
Hybrid crystal and Dual-readout calorimeter	3 % [30]	< 1 % [30]	≈ 26 % [30]	5–6 % [30,50]	3–4 % [50]
<a href="#">IDEA [48]</a> <a href="#">JINST 15 C06015</a>	<a href="#">CLD [20]</a> <a href="#">LCD-Note-2019-001</a>	<a href="#">Calos for FCC-hh [27]</a> <a href="#">CERN-FCC-PHYS-2019-0003</a>	<a href="#">Crystal Calos for FCs [30]</a> <a href="#">J. Instrum. 15, P11005–P11005 (2020)</a>		

Traditionally, the physics drivers for the “ultimate” ~3-4% PFlow jet energy resolution

- High efficiency for W/Z/H boson mass separation
- Separation of boosted objects (at higher energies)



# *SMEFT (at NLO)*

*Generically beyond SM*

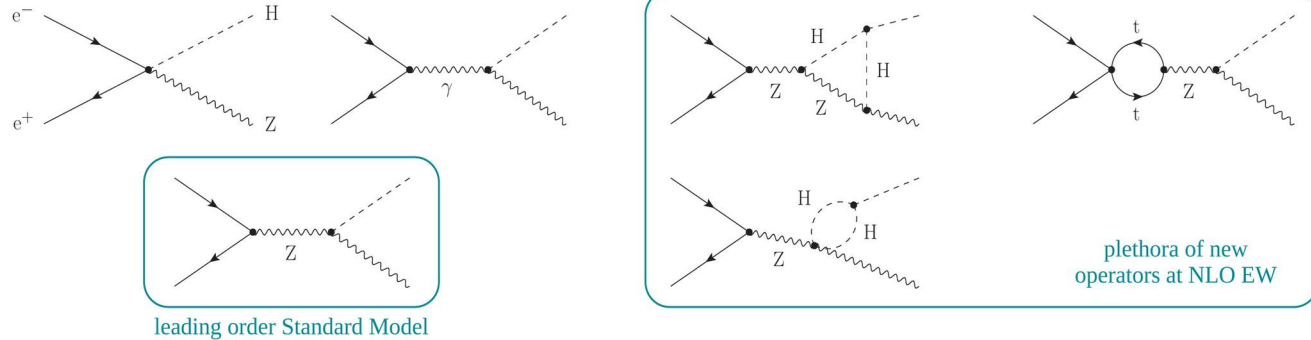
Konstantin Asteriadis

# Precision measurements and SMEFT

## SMEFT @ NLO

- Well studied on LEP data but LO only
- NLO opens up number of new operators
- Precision from FCC-ee is essential

### Higgstrahlung in the SM and SMEFT



- SM results available at NLO EW [Fleischer, Jegerlehner '83; Kniehl '92, Denner, Kublbeck, Mertig, Bohm '92; Bondarenko, Dydyshka, Kalinovskaya, Rumyantsev, Sadykov, Yermolchik '19]  
... many pieces known at NNLO accuracy [Sun, Feng, Jia, Sang '17; Gong, Li, Xu, Yang, Zhao '17; Song, Freitas '21; Chen, Guan, He, Li, Liu, Ma '22; Freitas, Song, Xie '23]
- SMEFT at LO extensively studied using LEP data  $\rightarrow$  precision of future lepton collider might allow the indirect study of operators not present at LO
- **Next step: SMEFT at NLO in the electro-weak expansion** (first studies published KA, Dawson, Giardino, Szafron, arXiv:2406.03557 ... more to come soon)

# Precision measurements and SMEFT

## SMEFT @ NLO

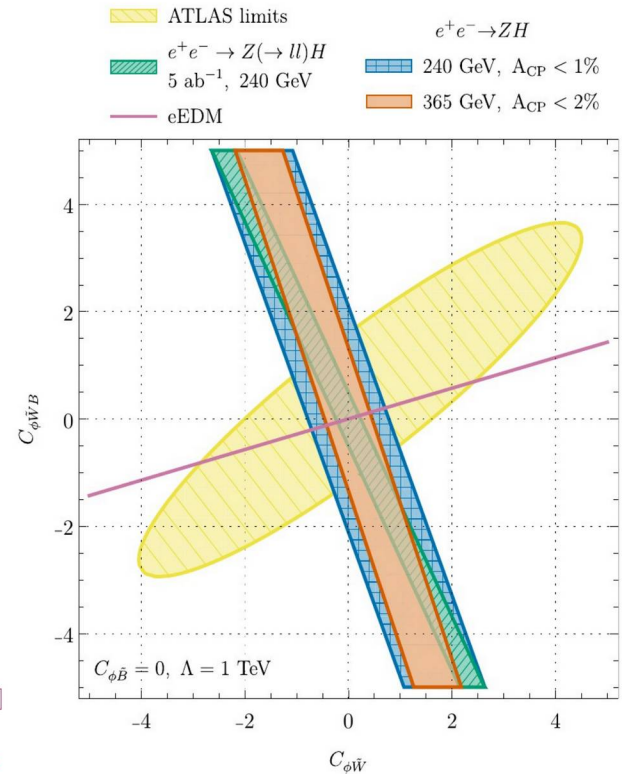
- Additional CP violating operators from NLO
- virtual corrections can develop imaginary contributions
- Instead of using full differential cross sections asymmetries should be sufficient
- LHC/FCC-ee complementary
- eEDM adds orthogonal very precise contribution

## CP Violation in Higgstrahlung

- Define CP violating asymmetry

$$A_{\text{CP}} = \frac{\sigma(\cos\theta < 0) - \sigma(\cos\theta > 0)}{\sigma_{\text{SM,NLO}}}$$

- Expected precision for the total cross section at FCC-ee might be as low as  $\sim 0.5\%$  at 240 GeV (365 GeV  $\sim 1\%$ )  
→ Assume half the precision
- Consider  $C_{\phi\tilde{W}}$  and  $C_{\phi\tilde{W}B}$  (other Wilson coefficients set to 0)
- Limits from  $H \rightarrow 4$  lepton decay at LHC [ATLAS, JHEP 05, 105 (2024)]
- Strong limits from electron electric dipole moment (eEDM) that also depends on SMEFT coefficients [ACME, Nature 562, 355 (2018)]
- Potential limits through angular observables [JHEP 03, 050 (2016)]



# Precision measurements and SMEFT

## Higgs Tri-linear and Top Quark Couplings

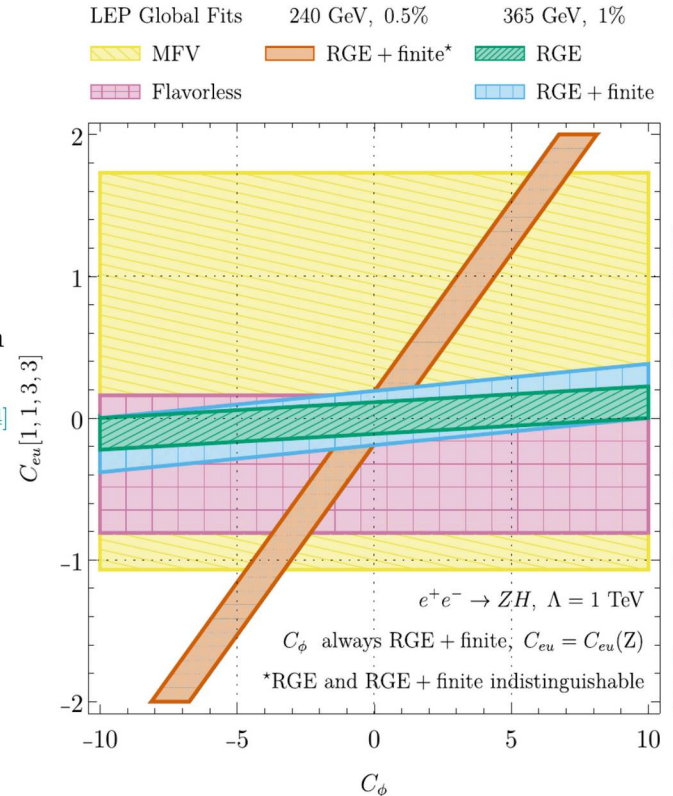
### SMEFT @ NLO

- Constraints on Higgs trilinear and electron-top coupling operators largely benefit from measurements at different  $E_{\text{CM}}$
- Different contributions have different dependence on the  $E_{\text{CM}}$

- Consider Higgs self-interaction  $C_\phi$  and electron-top 4-fermion operator  $C_{eu}[1, 1, 3, 3]$
- SMEFT Wilson coefficients are regulated in  $\overline{\text{MS}}$   $\rightarrow$  Scale dependent contributions  $\bar{\Delta}_i$  can be obtained from RGE evolution [Jenkins, Manohar, Trott '13 '14; Alonso, Jenkins, Manohar, Trott '14]

$$\frac{\sigma_{\text{NLO}}}{\sigma_{\text{SM,NLO}}} = 1 + \sum_i \frac{C_i(\mu)}{\Lambda^2} \left\{ \Delta_i + \bar{\Delta}_i \log \frac{\mu^2}{s} \right\}$$

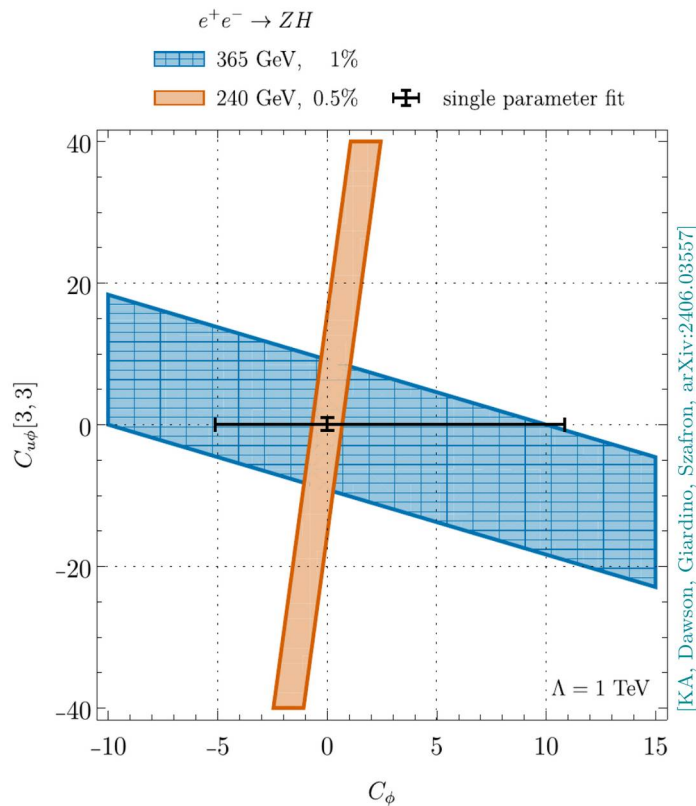
- Finite contributions  $\Delta_i$  only from exact higher order computations



# Precision measurements and SMEFT

## SMEFT @ NLO Higgs Tri-linear and Top Quark Couplings

- Systematic framework and evaluating all different operators
- SMEFT community is catching up and will join the fun
- Consider Higgs self-interaction  $C_\phi$  and anomalous top-Yukawa coupling  $C_{u\phi}[3,3]$
- Single parameter limits from global fit to LHC Higgs data [JHEP 04, 279 (2021)] and HH searches [ATLAS, arXiv:2404.05498]
- Measurement at two energy scales complementary





# *Beyond the Standard Model Physics*

Kevin Langhoff, Chris Verhaaren, Zeynep Demiragli

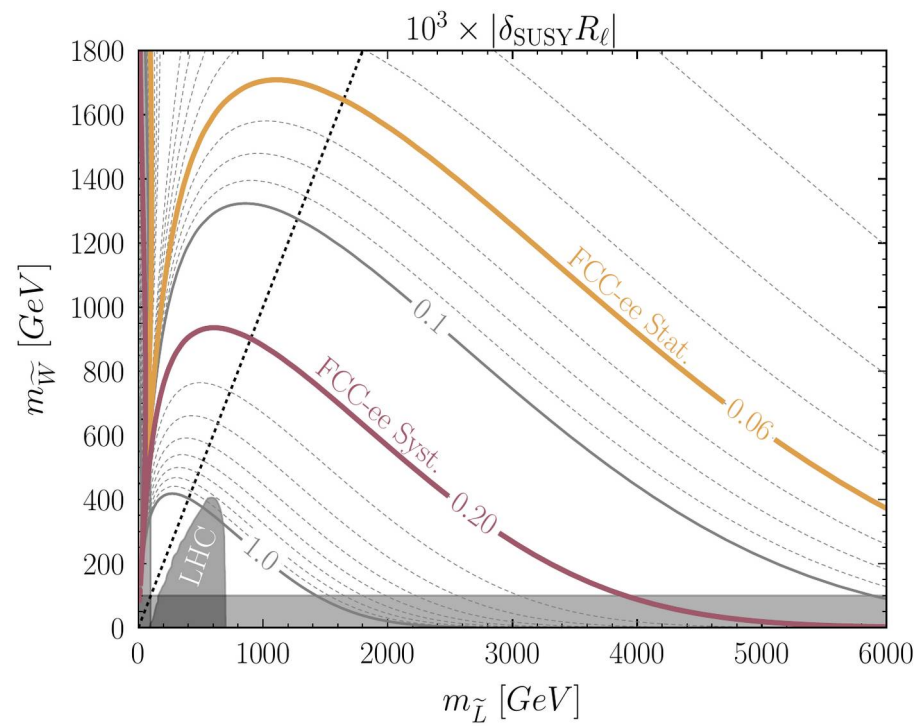
# SUSY has still open phase space

## SUSY

## Results

- For certain areas of the phase space there is still room
- Careful with older plots as the LHC might do better than indicated but ... there is room

$\cos(2\beta) = 0$



Wino + LH Slepton  
(Preliminary)

[1908.08215]

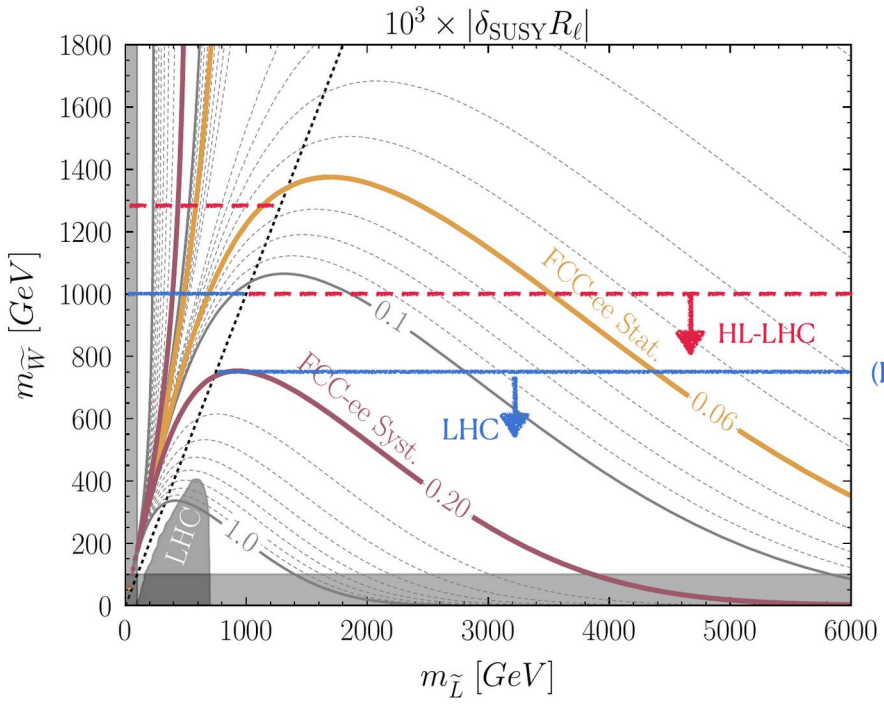
# SUSY has still open phase space

## SUSY

## Results

- LHC might do better than indicated but ... there is room
- Systematic uncertainties at FCC-ee are very important
- Theorists need to review the options, experimentalist the uncertainties

$\cos(2\beta) = -1$



Wino + LH Slepton (Preliminary)

(Depends on  $m_{\chi_1^0}$ !)

[1908.08215]

# Dark Sector: Axion-Like Particles

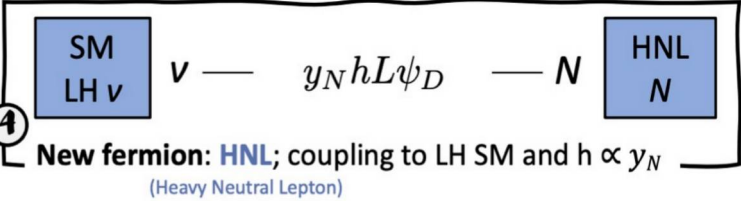
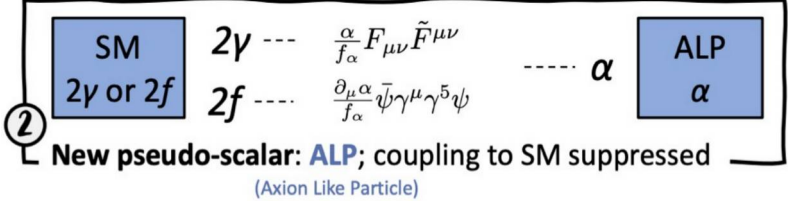
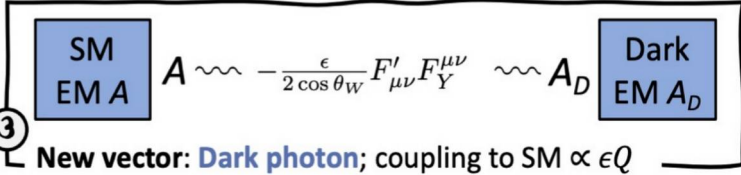
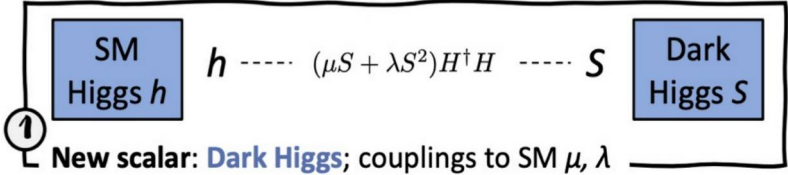


## Dark Sectors

### Personal Favorite Motivator: Dark Sector

While the dynamics of the dark sector could be complicated...  
to observe a dark sector, we need a portal interaction:

- Going to lower energy from LHC
- → use the event counts and look for lower masses: intensity frontier

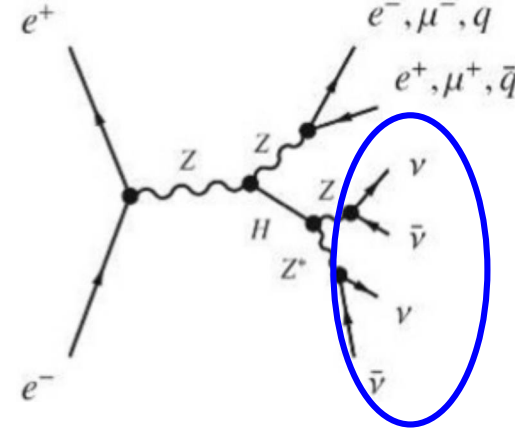


New Physics could be light and feebly interacting with SM

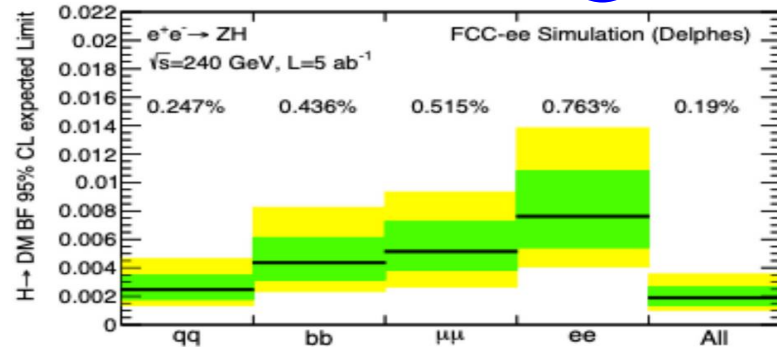
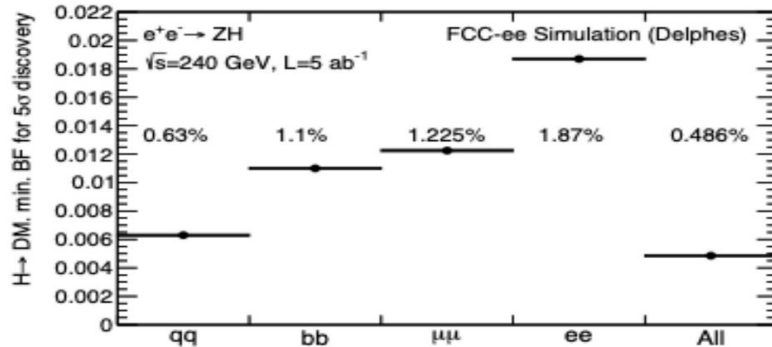
# Higgs Invisible Width

## Higgs boson: portal to dark world

- Use recoil and require nothing else
- Measure  $H \rightarrow ZZ \rightarrow \nu\nu\nu$
- Then remove as SM background
- SM precision 0.1%; NP at BF of 0.5%



Invisible decay products



Recent work ([FCC MIT Workshop](#)) compares CLD full sim and CLD & IDEA Delphes fast sim.

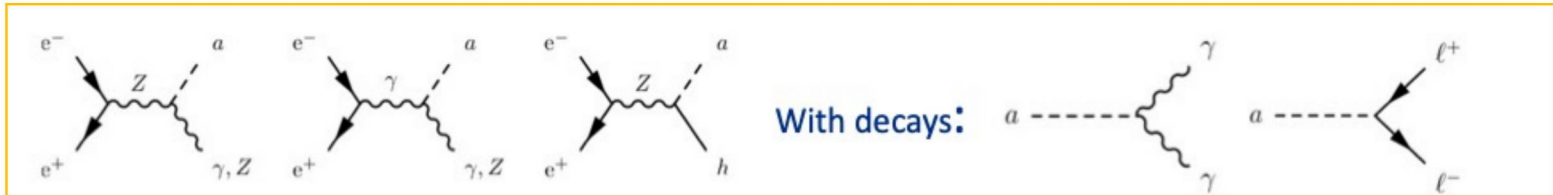
- Efficiency is ~identical for IDEA and CLD fast simulations!
- Electron eff is worse for full sim than for fast sim & Muon eff is very similar for full & fast sim.



# Dark Sector: Axion-Like Particles

## Case Study: Axion-like Particles

“Standard” approach:

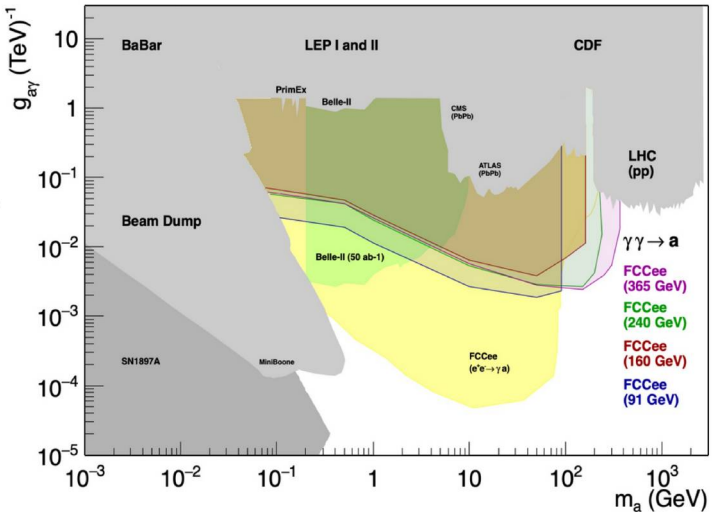


### Dark Sectors

- Going to lower energy from LHC
- → use the event counts and look for lower masses: intensity frontier

- $\gamma\gamma \rightarrow a$  extends current LHC limits for  $m_a = 5 - 350\text{GeV}$  by 2(O) magnitude
- $e^+ e^- \rightarrow Z \rightarrow \gamma a$  extends current LHC limits for  $m_a = 0.1 - 90\text{ GeV}$  by 3(O) magnitude

*For low ALP mass, sophisticated detectors & techniques are needed to isolate the overlapping photons*



# Can the Z be a portal?

## Rare Z decays

- Not only the Higgs can be a portal
- Z resonance holds potential for exotic decays

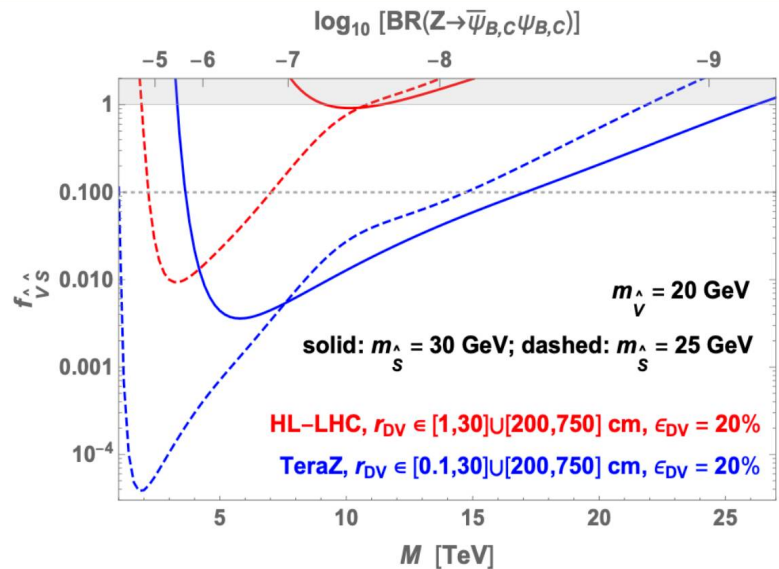
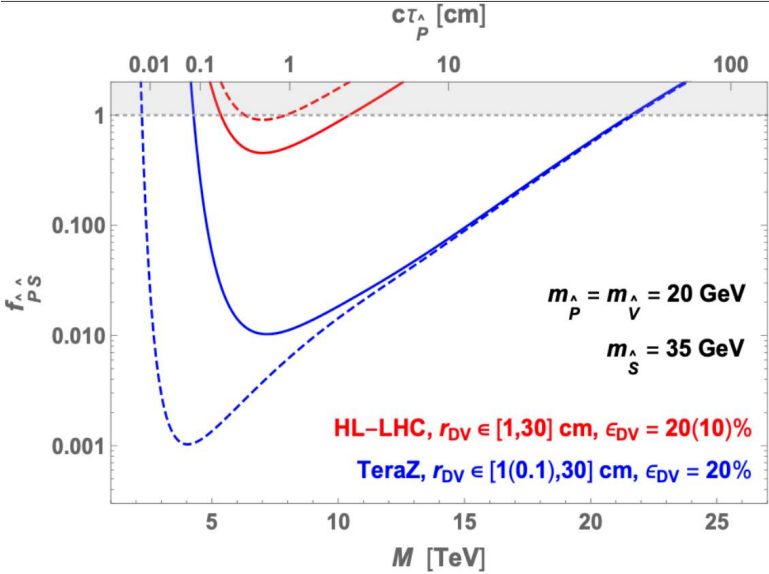
## Z-portal

Fraction of Z decays to hidden sector that are XY final state:  $f_{XY}$

Grey lines motivated benchmarks

See that FCC-ee has impressive reach

Cheng, Li, Salvioni, CV 1906.02198



# Sensitive to fractional charge?

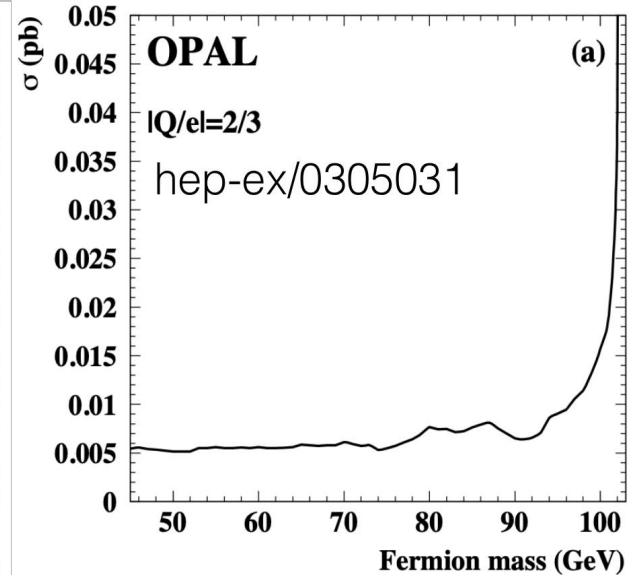
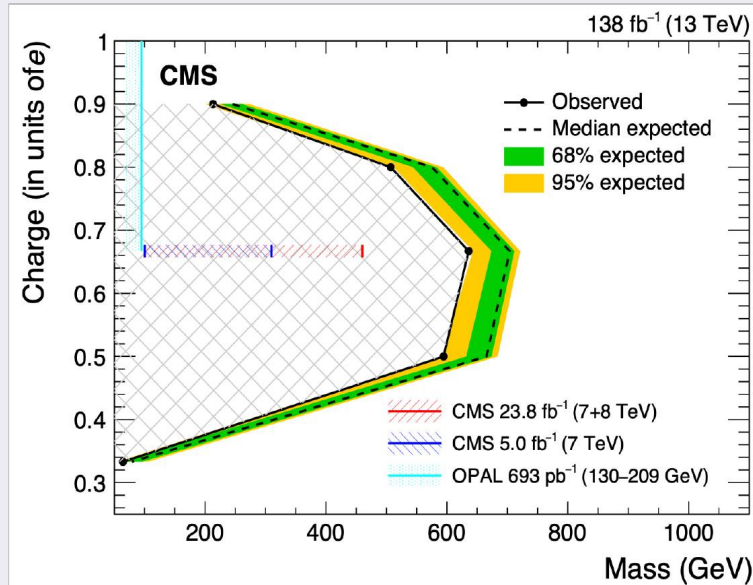
## GUT

- Recent paper motivate GUTs with possible  $e/6$  charges
- Can the FCC-ee make a contribution

## Fractionally Charge Particles

Not clear that the charge  $e/6$  target can be probed at the LHC

Can the FCC-ee make a comprehensive search/discovery up to the top-quark mass?



# Higgs and Glueballs?

## Higgs Physics

### Comments

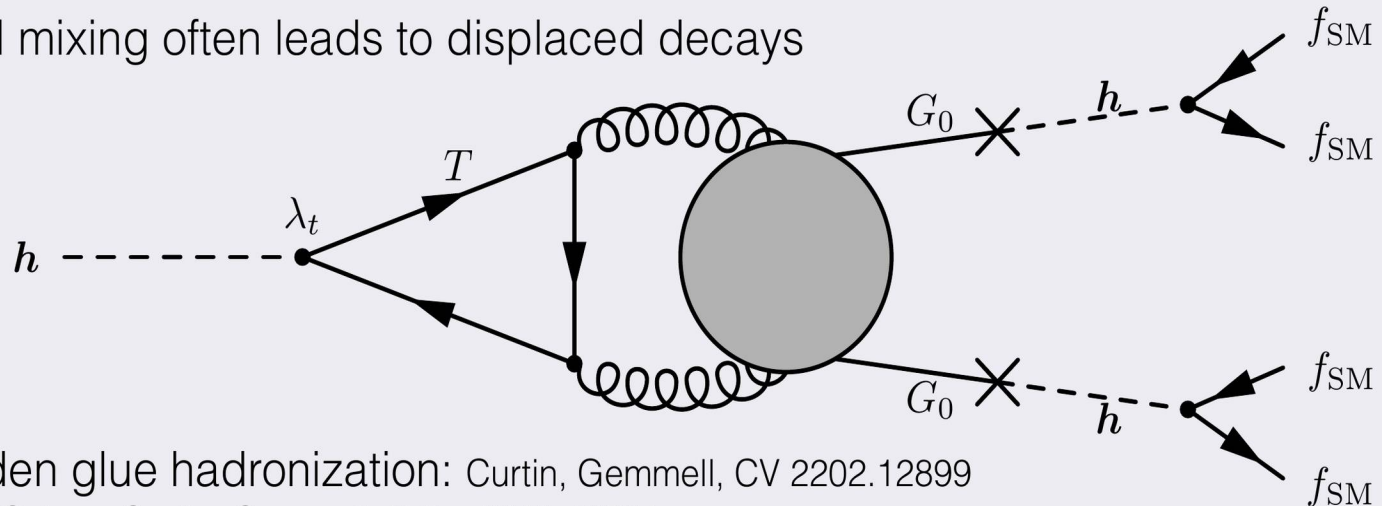
- Glueball searches are hard at the LHC
- The excess in  $b\bar{b}b\bar{b}$  final states will be hard to distinguish from the background

### Exotic Higgs Decays

Lightest hidden glueball mixes with the Higgs

$$h \rightarrow G_0 G_0 \rightarrow \bar{f} f \bar{f} f \quad (\text{Mostly to b-quarks})$$

Small mixing often leads to displaced decays



Hidden glue hadronization: Curtin, Gemmell, CV 2202.12899

Batz, Cohen, Curtin, Gemmell, Kribs 2310.13731

# Conclusions

## Status

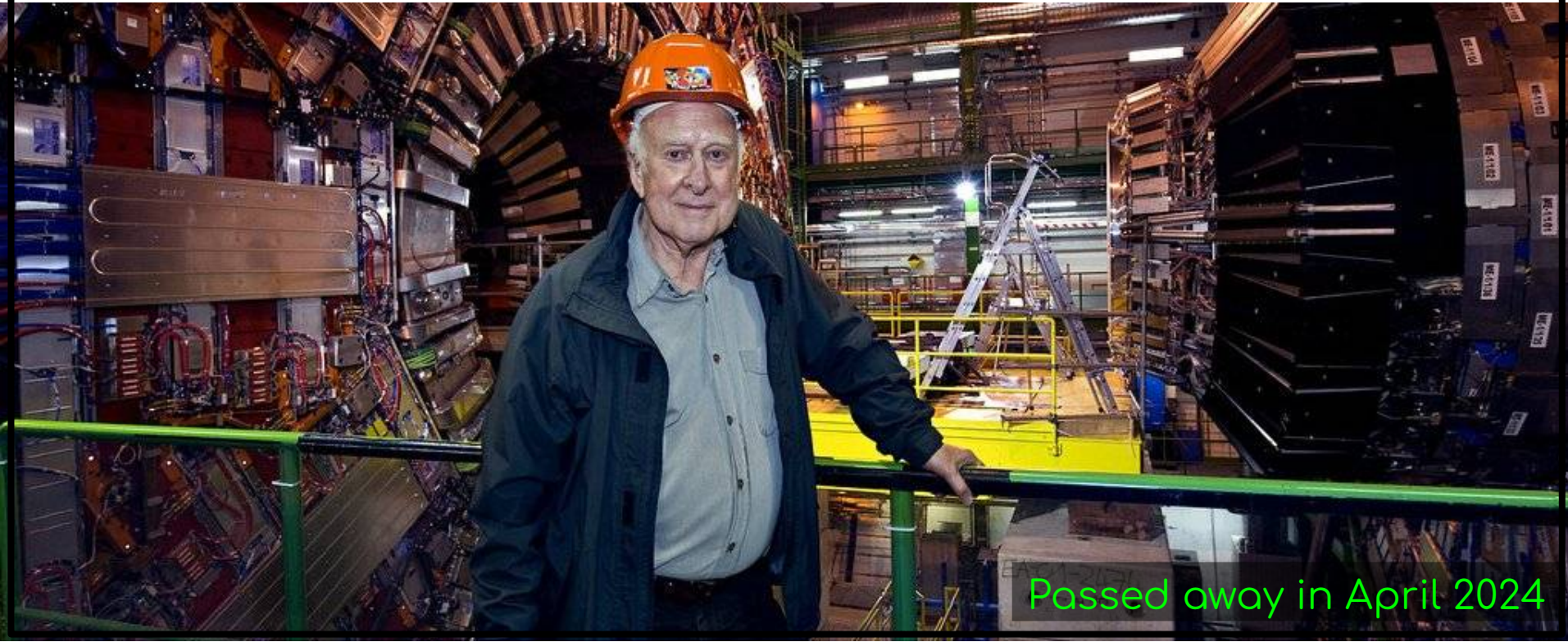
- FCC-ee produces  $\sim 2.2$ M Higgs bosons *in pristine conditions* and thus has a strong Higgs program (HL-LHC  $\sim 180$ M Higgses)
- There are extraordinary electroweak precision, flavor precision, top, and BSM programs
- Detector design ideas exist and match the requirements
- New ideas for even better solutions are being investigated

## Work that needs doing

- Work on systematics and the theory is essential
- New detector technology should be supported
- Detector integration is starting to move into focus



# R.I.P. Peter Higgs



Passed away in April 2024



*More*

# FCC-ee: Physics & Detectors

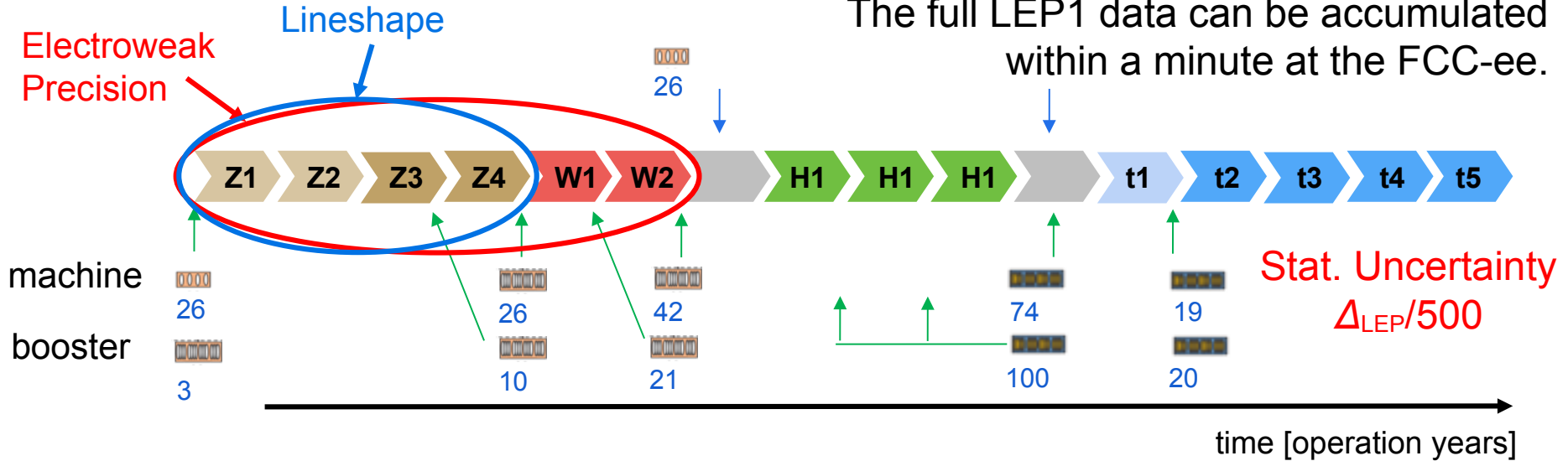
A panoramic view of the San Francisco skyline and the Golden Gate Bridge over the water. The bridge is in the foreground, stretching across the water towards the city. The city skyline is visible in the background, with various skyscrapers and buildings. The sky is clear and blue.

July 11, 2024  
57<sup>th</sup> Users Meeting, Fermilab

Christoph Paus, MIT

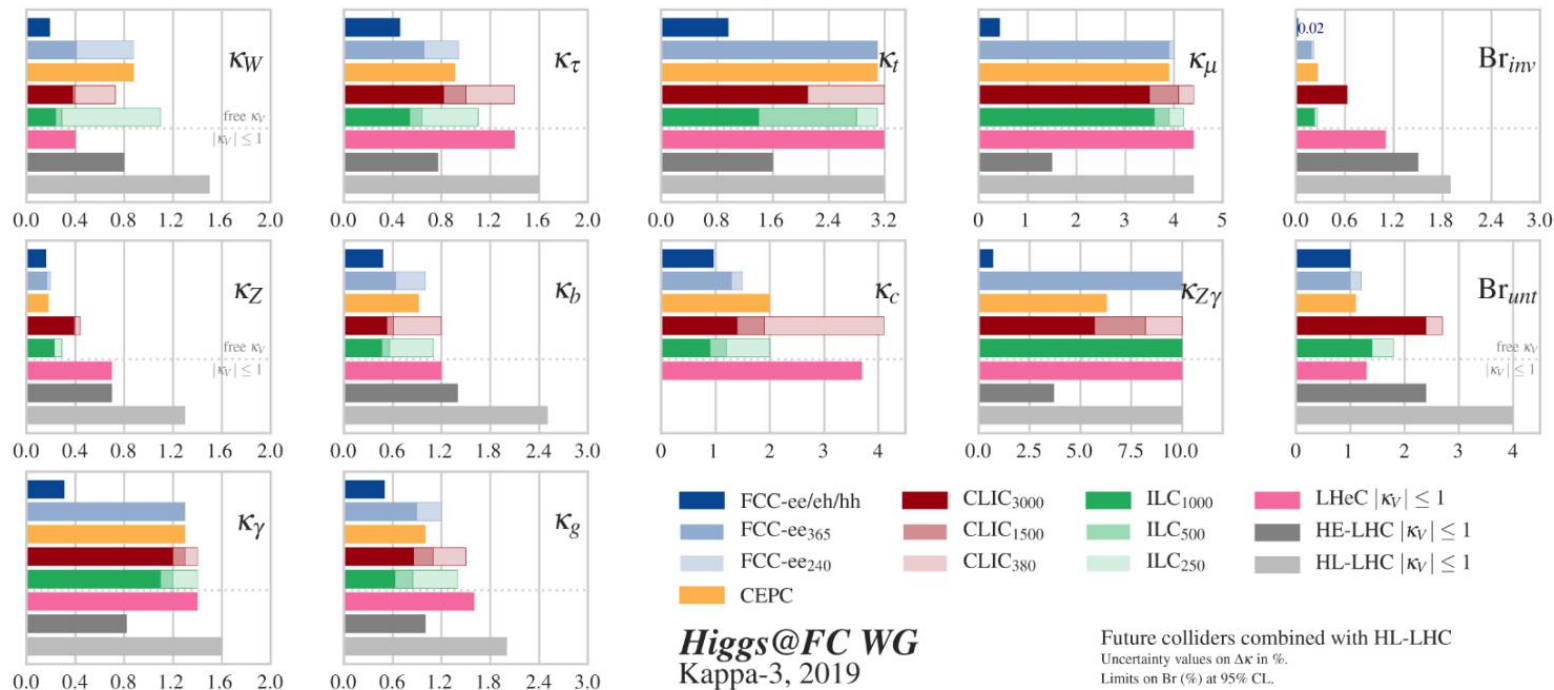
# FCC-ee Schedule

The full LEP1 data can be accumulated within a minute at the FCC-ee.



Working point	Z, years 1-2	Z, later	WW, years 1-2	WW, later	ZH	t $\bar{t}$	
$\sqrt{s}$ (GeV)	88, 91, 94		157, 163		240	340–350	365
Lumi/IP ( $10^{34} \text{ cm}^{-2} \text{ s}^{-1}$ )	70	140	10	20	5.0	0.75	1.20
Lumi/year ( $\text{ab}^{-1}$ )	34	68	4.8	9.6	2.4	0.36	0.58
Run time (year)	2	2	2	0	3	1	4
Number of events	$6 \cdot 10^{12}$ Z		$2.4 \cdot 10^8$ WW		$1.45 \cdot 10^6$ HZ + 45k WW $\rightarrow$ H	$1.9 \cdot 10^6$ t $\bar{t}$ +330k HZ +80k WW $\rightarrow$ H	

# Higgs Couplings Precision



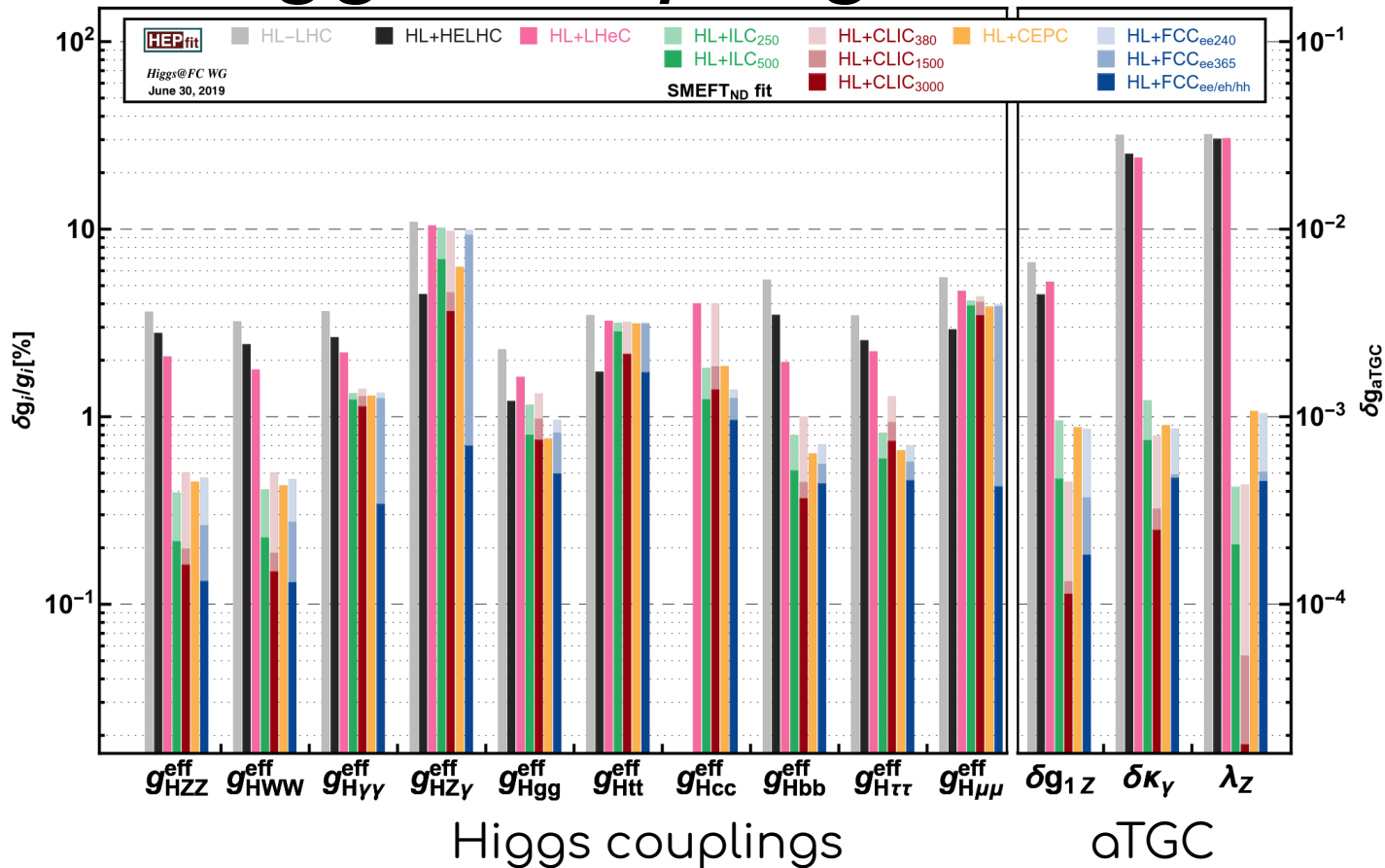
Not very dependent on the  $e^+e^-$  option

- Sensitivity to Higgs coupling is mostly around a percent
- Details of the uncertainties are dependent on the specific implementations

# Higgs Couplings Precision

Sensitivity to deviations for

- Different effective Higgs couplings
- and aTGC



# Lineshape Summary

Quantity	current	ILC250	ILC-GigaZ	FCC-ee	CEPC	CLIC380
$\Delta\alpha(m_Z)^{-1} (\times 10^3)$	17.8*	17.8*		3.8 (1.2)	17.8*	
$\Delta m_W$ (MeV)	12*	0.5 (2.4)		0.25 (0.3)	0.35 (0.3)	
$\Delta m_Z$ (MeV)	2.1*	0.7 (0.2)	0.2	0.004 (0.1)	0.005 (0.1)	2.1*
$\Delta m_H$ (MeV)	170*	14		2.5 (2)	5.9	78
$\Delta\Gamma_W$ (MeV)	42*	2		1.2 (0.3)	1.8 (0.9)	
$\Delta\Gamma_Z$ (MeV)	2.3*	1.5 (0.2)	0.12	0.004 (0.025)	0.005 (0.025)	2.3*
$\Delta\sigma_{\text{had}}^0$ (pb)	37*			0.035 (4)	0.05 (2)	37*
$\delta R_e (\times 10^3)$	2.4*	0.5 (1.0)	0.2 (0.5)	0.004 (0.3)	0.003 (0.2)	2.7
$\delta R_\mu (\times 10^3)$	1.6*	0.5 (1.0)	0.2 (0.2)	0.003 (0.05)	0.003 (0.1)	2.7
$\delta R_\tau (\times 10^3)$	2.2*	0.6 (1.0)	0.2 (0.4)	0.003 (0.1)	0.003 (0.1)	6
$\delta R_b (\times 10^3)$	3.0*	0.4 (1.0)	0.04 (0.7)	0.0014 (< 0.3)	0.005 (0.2)	1.8
$\delta R_c (\times 10^3)$	17*	0.6 (5.0)	0.2 (3.0)	0.015 (1.5)	0.02 (1)	5.6

Z pole run



# Lineshape Summary

WW threshold run

Quantity	current	ILC250	ILC-GigaZ	FCC-ee	CEPC	CLIC380
$\Delta\alpha(m_Z)^{-1} (\times 10^3)$	17.8*	17.8*		3.8 (1.2)	17.8*	
$\Delta m_W$ (MeV)	12*	0.5 (2.4)		0.25 (0.3)	0.35 (0.3)	
$\Delta m_Z$ (MeV)	2.1*	0.7 (0.2)	0.2	0.004 (0.1)	0.005 (0.1)	2.1*
$\Delta m_H$ (MeV)	170*	14		2.5 (2)	5.9	78
$\Delta\Gamma_W$ (MeV)	42*	2		1.2 (0.3)	1.8 (0.9)	
$\Delta\Gamma_Z$ (MeV)	2.3*	1.5 (0.2)	0.12	0.004 (0.025)	0.005 (0.025)	2.3*
$\Delta\sigma_{\text{had}}^0$ (pb)	37*			0.035 (4)	0.05 (2)	37*
$\delta R_e (\times 10^3)$	2.4*	0.5 (1.0)	0.2 (0.5)	0.004 (0.3)	0.003 (0.2)	2.7
$\delta R_\mu (\times 10^3)$	1.6*	0.5 (1.0)	0.2 (0.2)	0.003 (0.05)	0.003 (0.1)	2.7
$\delta R_\tau (\times 10^3)$	2.2*	0.6 (1.0)	0.2 (0.4)	0.003 (0.1)	0.003 (0.1)	6
$\delta R_b (\times 10^3)$	3.0*	0.4 (1.0)	0.04 (0.7)	0.0014 (< 0.3)	0.005 (0.2)	1.8
$\delta R_c (\times 10^3)$	17*	0.6 (5.0)	0.2 (3.0)	0.015 (1.5)	0.02 (1)	5.6

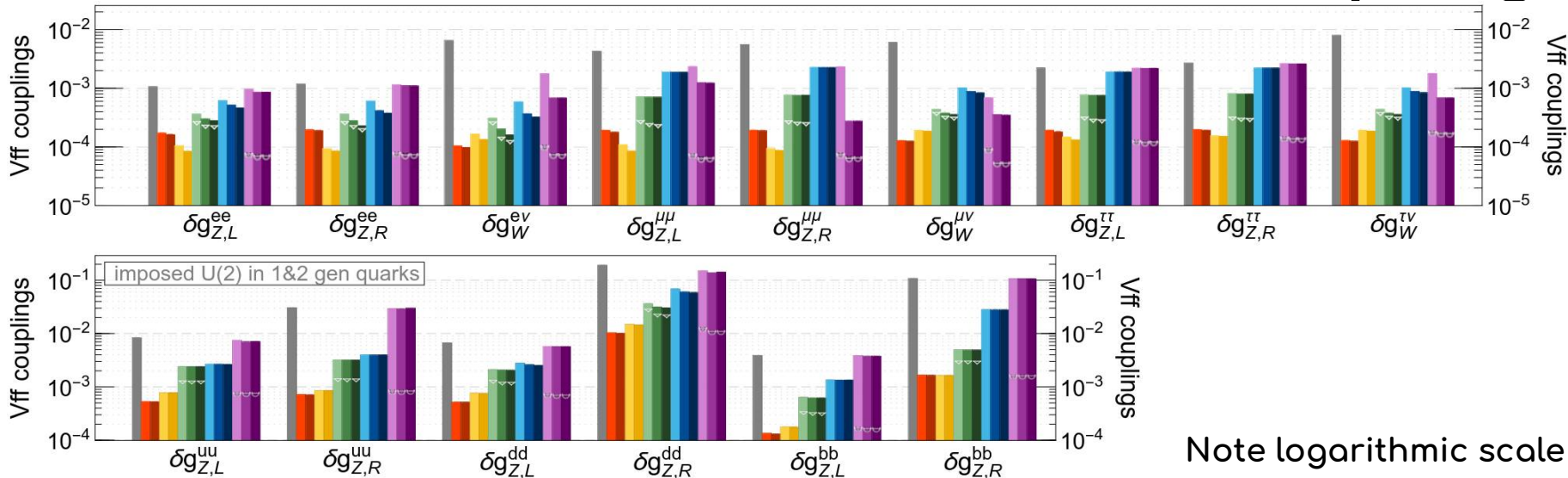
# Asymmetry Summary

Quantity	current	ILC250	ILC-GigaZ	FCC-ee	CEPC	CLIC380
$\Delta A_e (\times 10^5)$	190*	14 (4.5)	1.5 (8)	0.7 (2)	1.5	64
$\Delta A_\mu (\times 10^5)$	1500*	82 (4.5)	3 (8)	2.3 (2.2)	3.0 (1.8)	400
$\Delta A_\tau (\times 10^5)$	400*	86 (4.5)	3 (8)	0.5 (20)	1.2 (6.9)	570
$\Delta A_b (\times 10^5)$	2000*	53 (35)	9 (50)	2.4 (21)	3 (21)	380
$\Delta A_c (\times 10^5)$	2700*	140 (25)	20 (37)	20 (15)	6 (30)	200

## A few points to note

- Z pole running creates substantially improved precision for all 'LEP' measurements by close to 3 orders of magnitude (statistically speaking)
- Major work for experimental and theory community to bring that precision to bear

# Global Fit focus $W/Z$ couplings



Note logarithmic scale

As expected

- Precision on couplings of  $W$  and  $Z$  bosons to fermions is more competitive at circular collider