

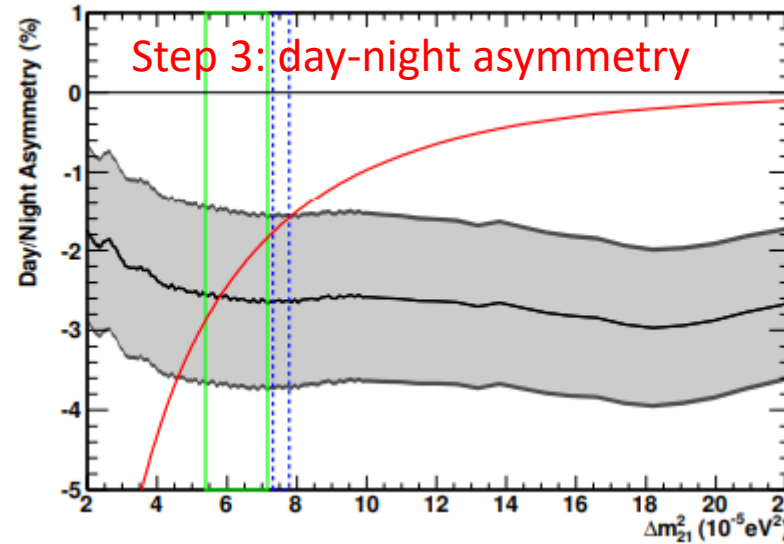
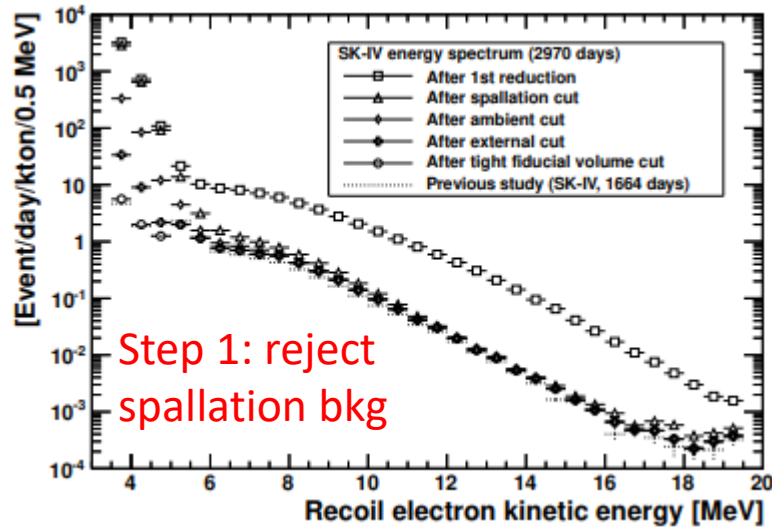
New solar neutrino analyses with argon detectors

Dan Pershey – Florida State University

Oct 17, 2024



How current solar neutrino analysis works

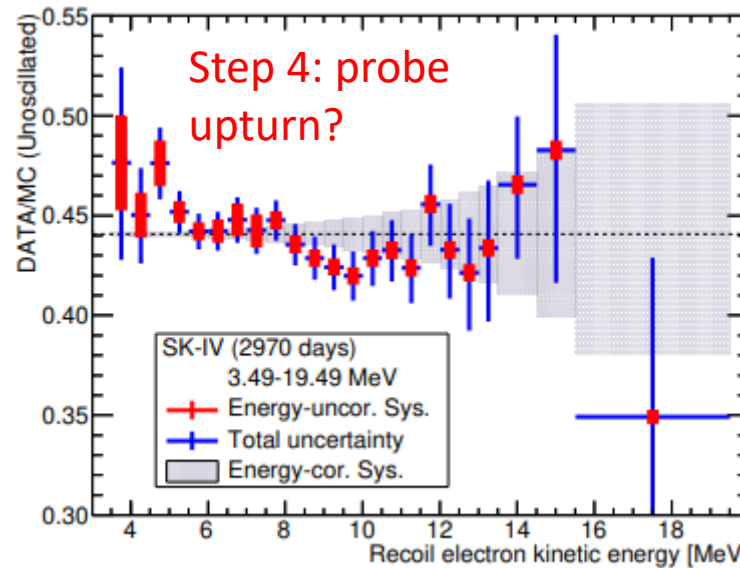
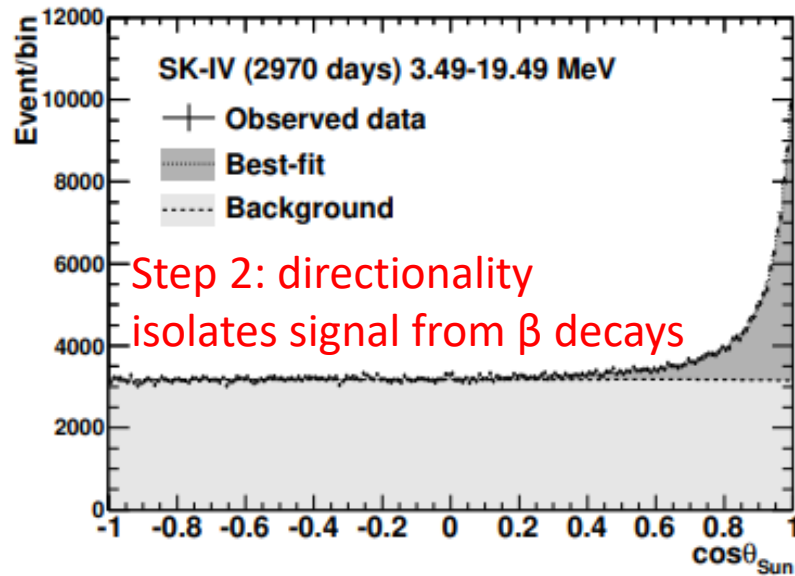


Electron scatter signal isolated by spallation cuts and direction

Difficult to understand oscillations that depend on E_ν with E_e spectra

ES channel also sensitive to all neutrino flavors, but with different weights, also smears out characteristics of ν_e flux

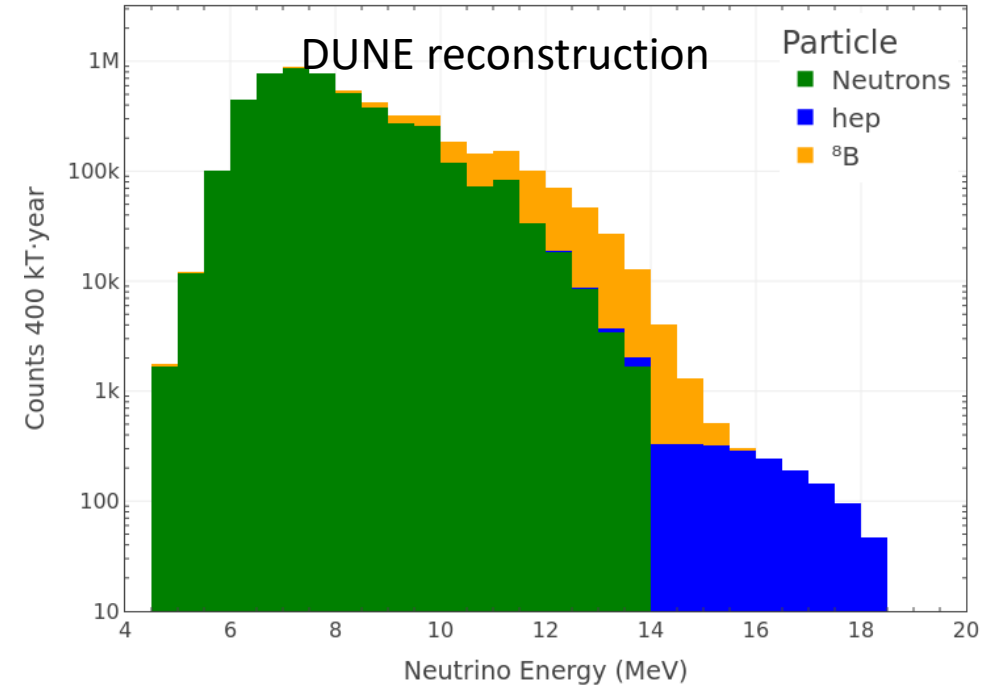
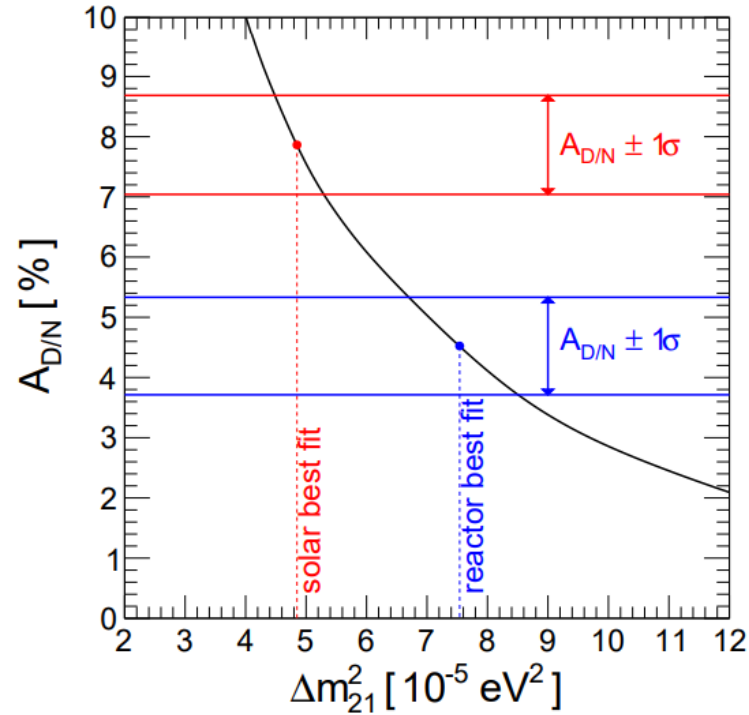
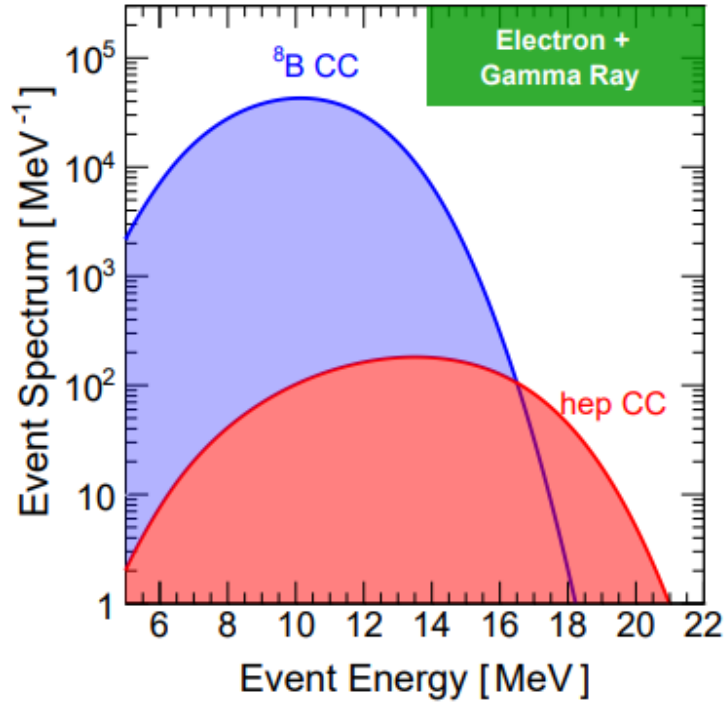
Desire for next generation experiments to precisely understand day-night diffs and probe the solar upturn



Finding the solar upturn



The power of argon



- ❑ Large CC xsec – good energy resolution + day-night asymmetry
- ❑ TPC technology – good tracking and calorimetry
- ❑ High light yield – scintillation gives additional analysis stream
 - Q+L -> improved energy reconstruction + quenching ID (α rejection)

Problem: intrinsic argon backgrounds

- ❑ Environmental radioactive noble nuclei contaminate liquid argon
- ❑ Most decays are not signal-like and can be rejected
 - Charge quenching rejects α 's
 - Low- Q value β 's below energy ROI BiPo's
 - Fiducialization removes plate-out
- ❑ Perform light-only (almost) analysis
 - Poisson energy smearing: LY = 180 PE/MeV
 - Assume 1 m vertex resolution
 - α rejection uses charge quenching

Source	Nucleus	Rejection	Activity
^{39}Ar	^{39}Ar	$Q_\beta = 0.565 \text{ MeV}$	—
^{42}Ar	^{42}Ar	$Q_\beta = 0.600 \text{ MeV}$	—
^{42}Ar	^{42}K	—	0.040 mBq/kg
^{85}Kr	^{85}Kr	$Q_\beta = 0.687 \text{ MeV}$	—
^{222}Rn	^{222}Rn	α	—
^{222}Rn	^{218}Po	α	—
^{222}Rn	^{214}Pb	—	0.35 mBq/kg
^{222}Rn	^{214}Bi	BiPo	—
^{222}Rn	^{214}Po	α	—
^{222}Rn	^{210}Tl	—	73 nBq/kg
^{222}Rn	^{210}Pb	plate-out	—
^{222}Rn	^{210}Bi	plate-out	—
^{222}Rn	^{210}Po	plate-out	—
^{220}Rn	^{220}Rn	α	—
^{220}Rn	^{216}Po	α	—
^{220}Rn	^{212}Pb	$Q_\beta = 0.569 \text{ MeV}$	—
^{220}Rn	^{212}Bi	BiPo	—
^{220}Rn	^{212}Po	α	—
^{220}Rn	^{208}Tl	α	0.056 mBq/kg

Solution: the golden channel

□ Story as old as the first detection of the neutrino – find a signal channel with time-correlated energy deposits

- Explore forest of deexcitation gammas – there is one with long, 480 ns lifetime
- Decay happens in 65% of Fermi transitions and 25% of total solar interactions

□ Golden channel – prompt+delayed scintillation time profile

- $E_1 = E_\nu - 2.73 \text{ MeV}$
- $E_2 = 1.65 \text{ MeV}$

□ High performance PDS gives good E / vtx resolution for correlation

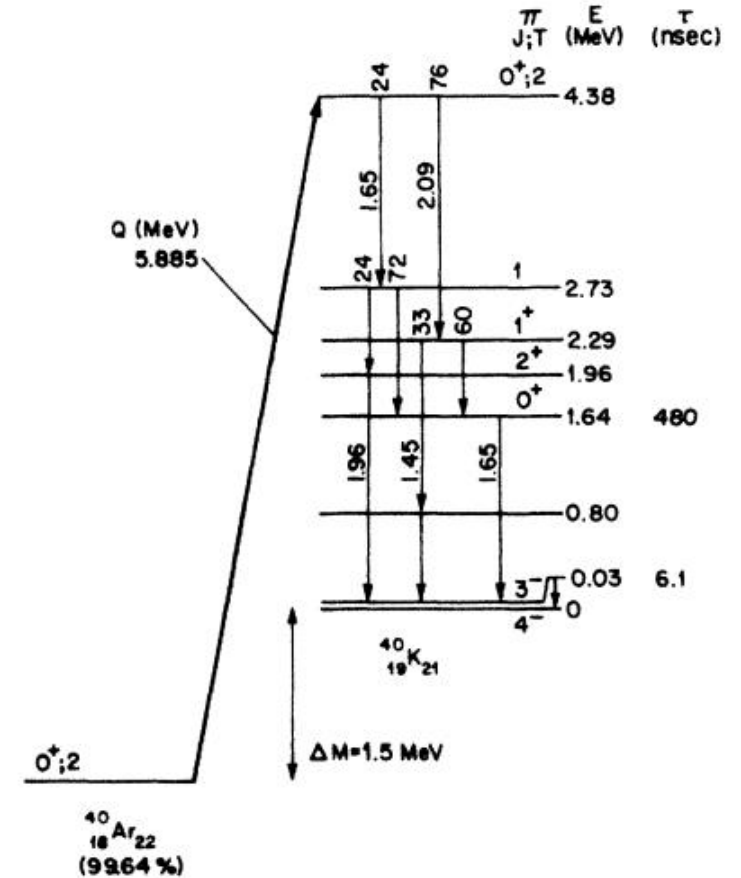
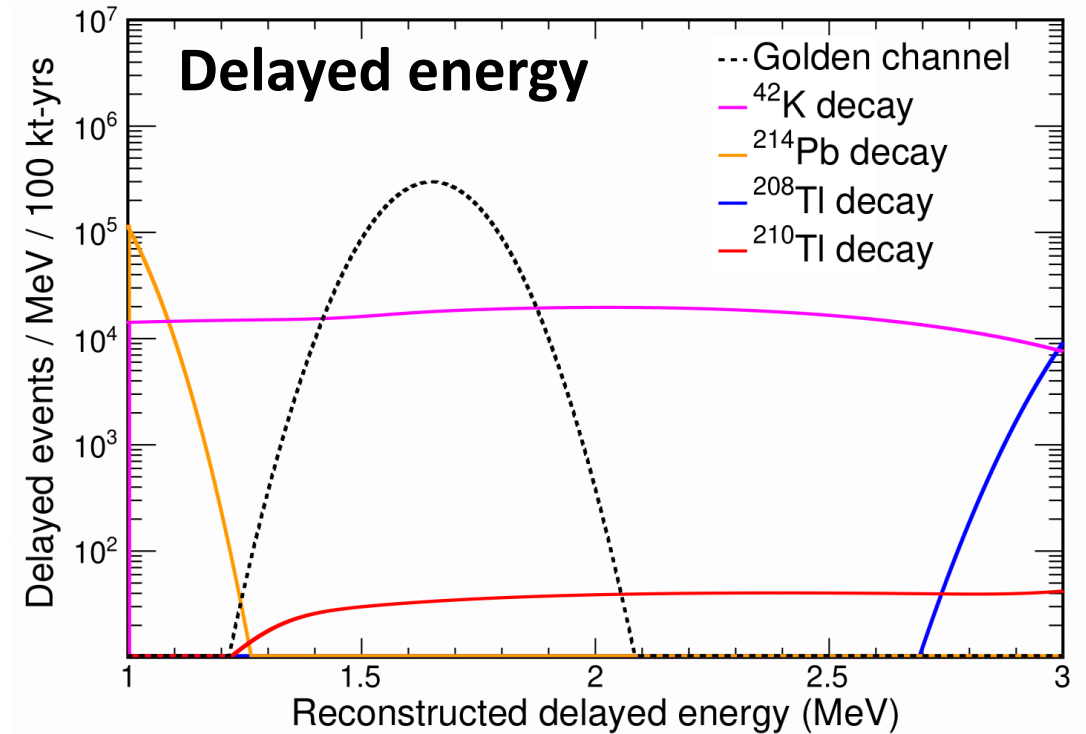
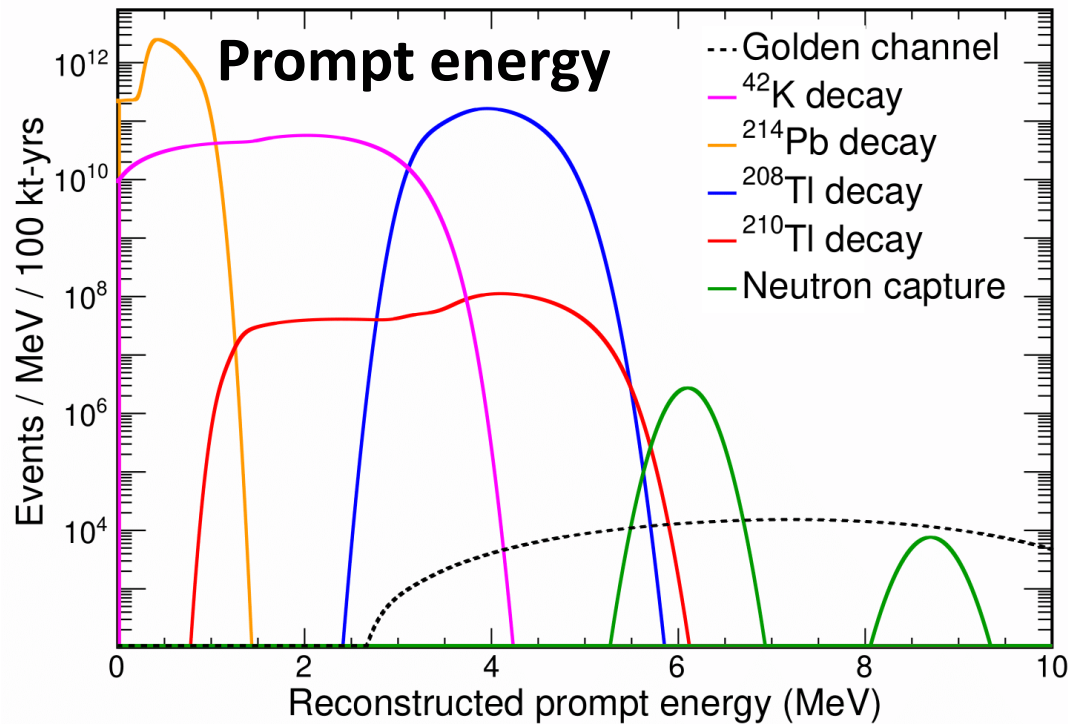


FIG. 1. Level scheme of ^{40}Ar - ^{40}K relevant to ν_e capture in argon.

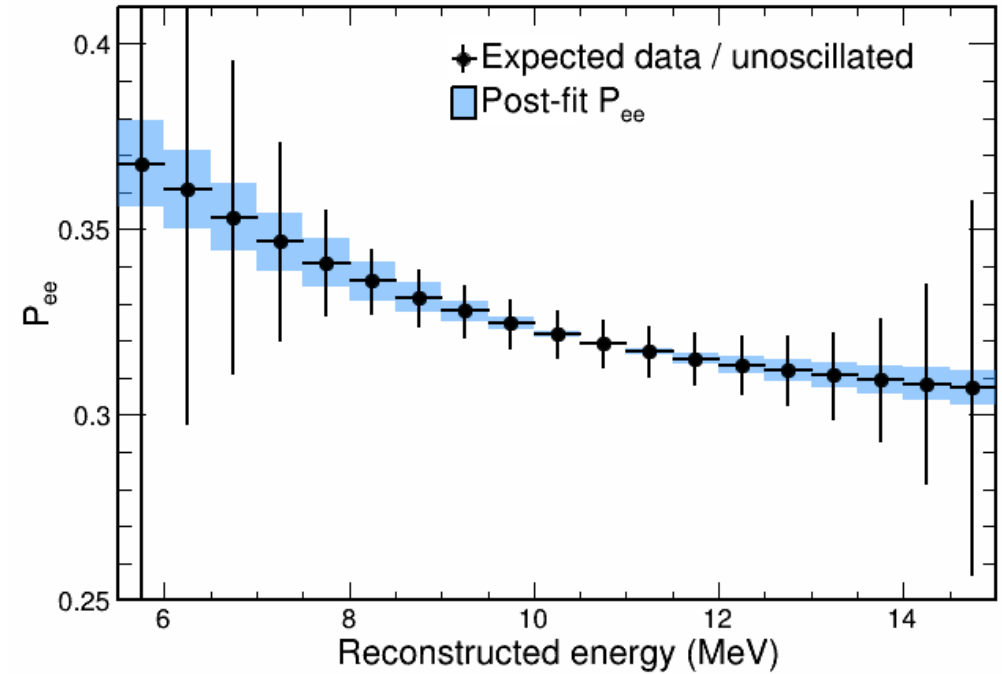
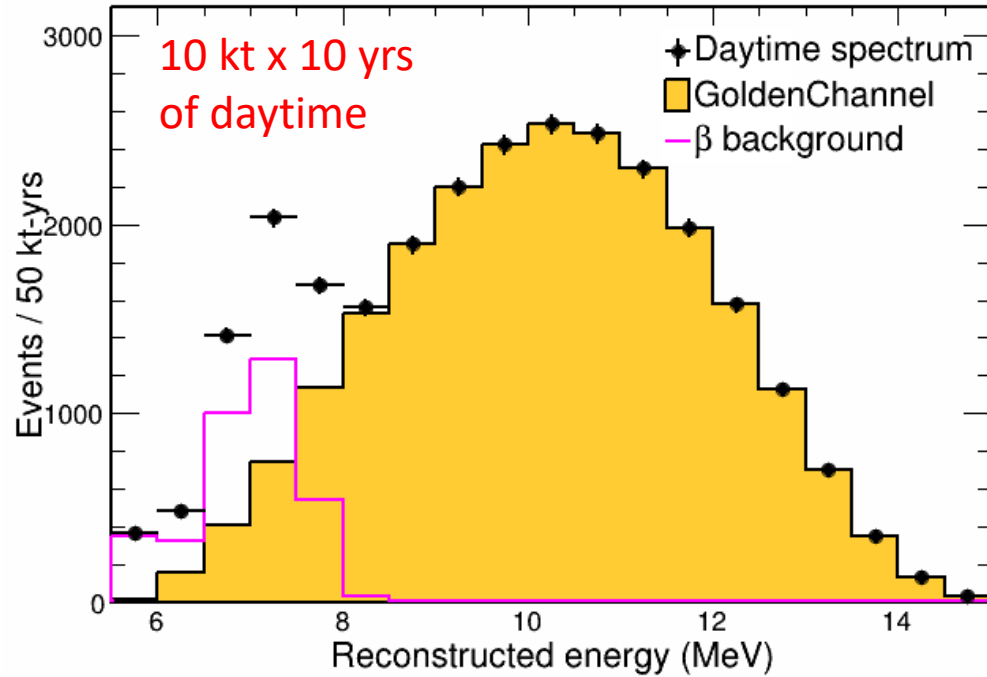
Selected background rate for DUNE



□ $O(1e12 \text{ bkg}) / 100 \text{ kt-yrs}$ in prompt ROI (dominated by ^{208}Tl decay)

□ $O(3e-8)$ bkg rejection factor after requiring delayed coincidence

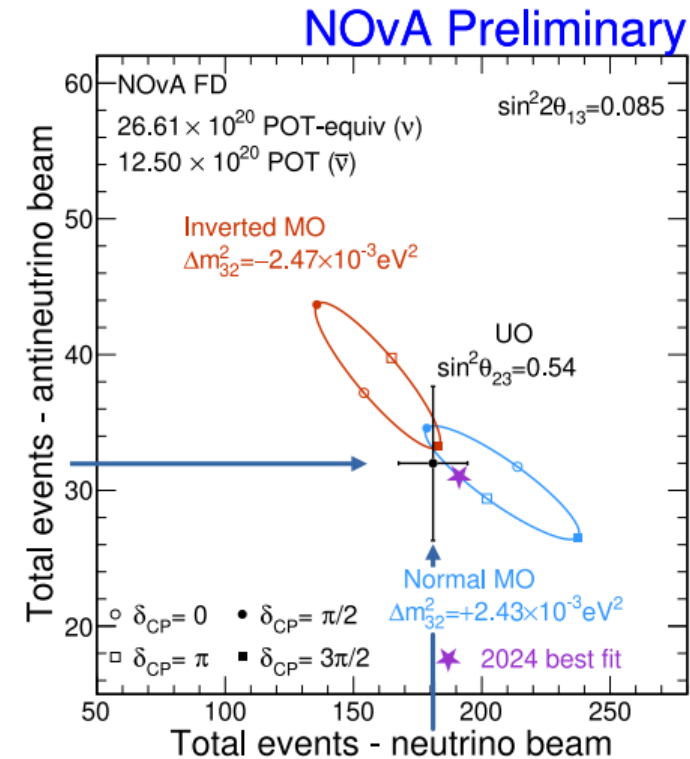
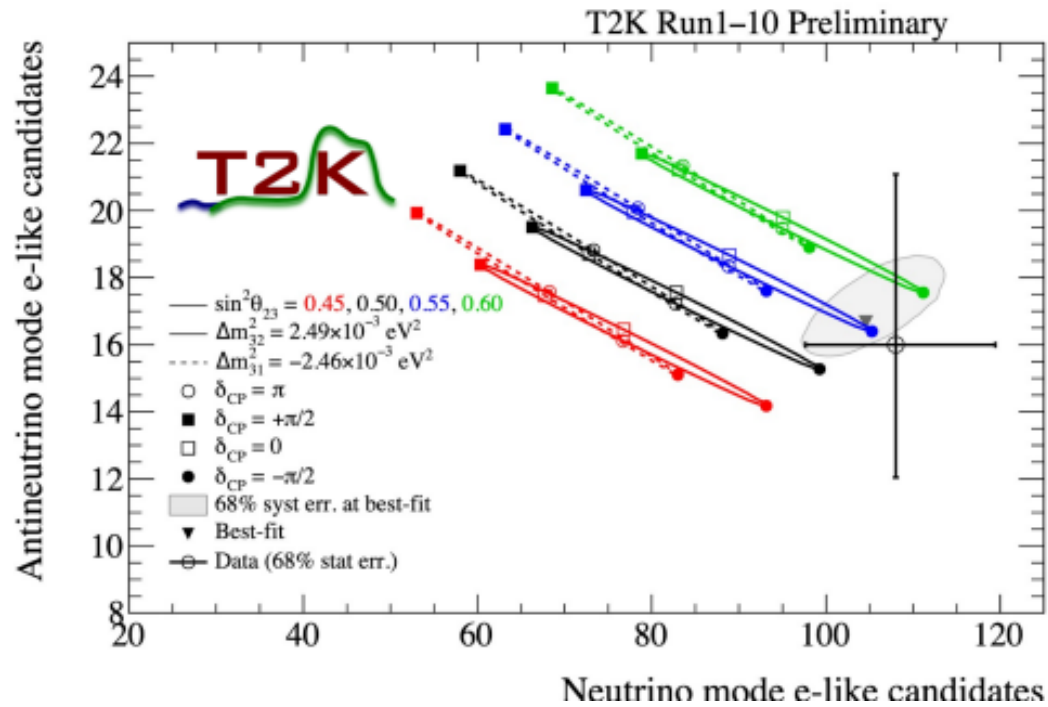
Golden channel analysis



- Select only night-time events to mitigate complication of wiggles
 - 2370/yr signal + 355/yr bkg
 - Upturn visible down to 6 MeV + bkg free above 8 MeV
- Shape-only fit: upturn visible at 4σ and Δm^2_{21} measured at 15%
- Allows night-time studies to complement -> extra physics with solar ν

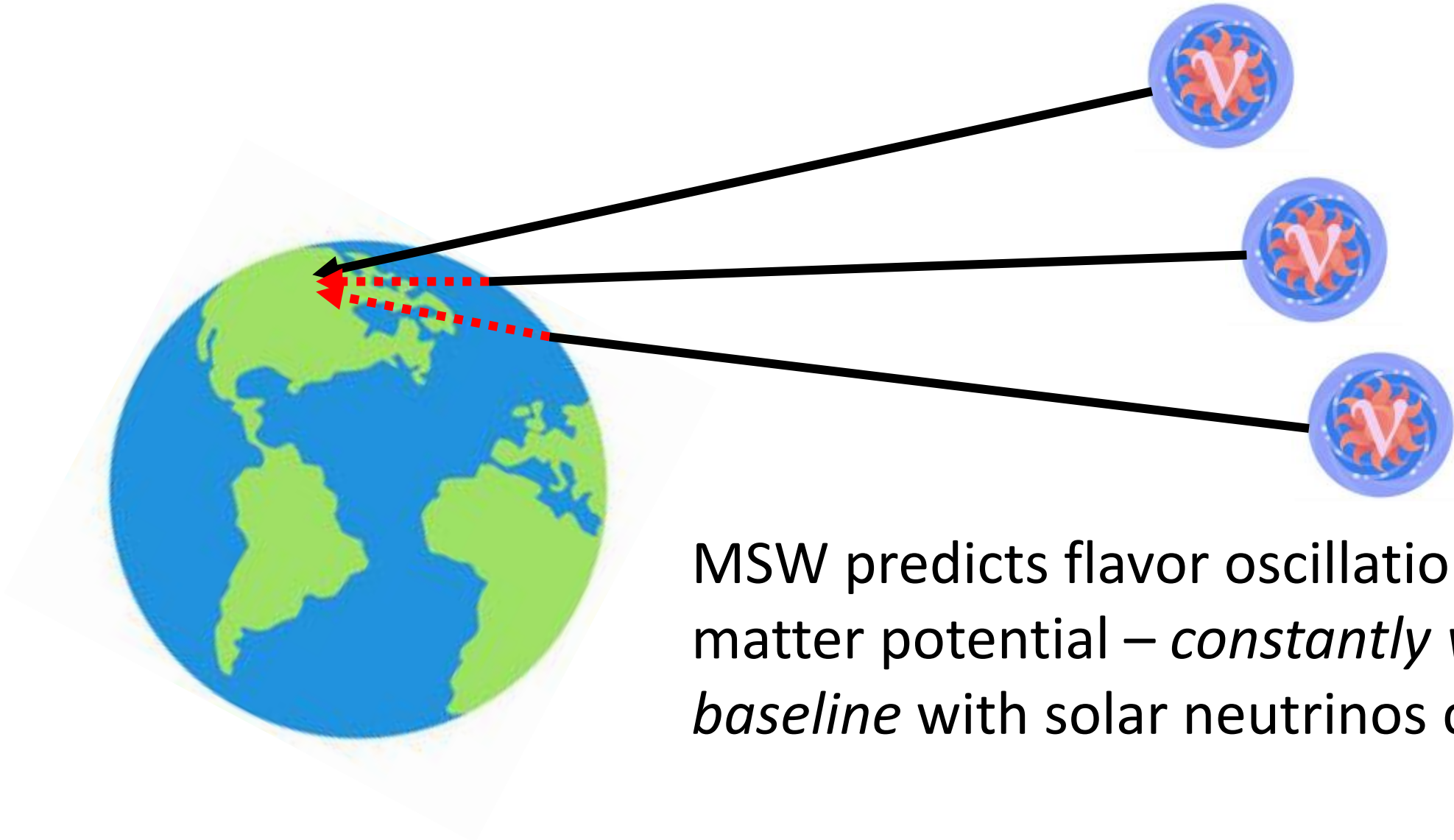
Twilight neutrinos

Testing the MSW effect



- Long baseline results rely on matter potential to break degeneracy
- Matter effects calculated with MSW model – is it valid?
- NOvA and T2K may now point to non-standard behavior (stats low!)

Testing MSW with twilight neutrinos



MSW predicts flavor oscillations in matter potential – *constantly varying baseline* with solar neutrinos can test

Review of mixing nomenclature in the sun

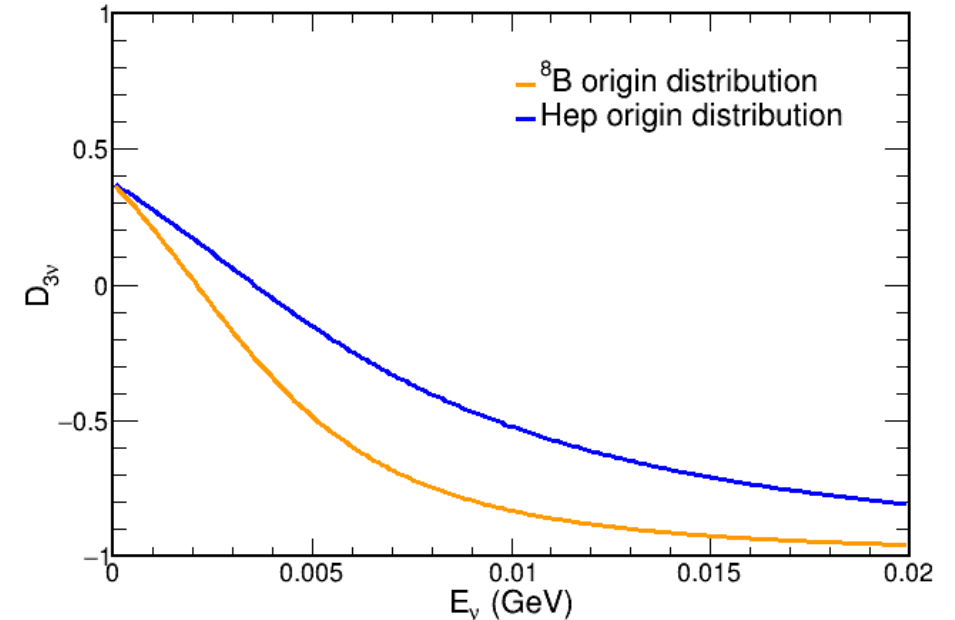
☐ Neutrinos arrive at Earth as a mix of mass eigenstates

- $\approx \cos \theta_{13} \left(\sqrt{\frac{1+\cos 2\theta_{\text{mat}}}{2}} |v_1\rangle + \sqrt{\frac{1-\cos 2\theta_{\text{mat}}}{2}} |v_2\rangle \right) + \sin \theta_{13} |v_3\rangle$

- $\cos 2\theta_{\text{mat}} = \frac{K \cos 2\theta_{12} - V_{CC}}{\sqrt{(K \cos 2\theta_{12} - V_{CC})^2 + (K \sin 2\theta_{12})^2}}$ $K \equiv \frac{\Delta m_{21}^2}{2E_\nu}$

- $D_{3\nu} = \int \cos(2\theta_{\text{mat}}(r_{\text{prod}})) P(r_{\text{prod}}) dr_{\text{prod}}$

[Blennow, Ohlsson, Snellman, PRD 69 073006 \(2004\)](#)



$D_{3\nu} = -1$ for $|v_e\rangle_{\text{mat}} = |v_2\rangle$ never fully realized
Observed signals in experiments depend on the local electron density at production

Study overview

- In high-energy limit, MSW resonance completely activated, solar neutrinos arrive as $\approx |\nu_2\rangle$
- As they enter Earth, matter effects change flavor content of $|\nu_2\rangle$ which induces oscillations – the wiggles
 - Night-time excess of event rate observed

$$\square P_{ee}^n - P_{ee}^d \cong \frac{V_{CC} E_\nu}{\Delta m_{21}^2} \cos^6 \theta_{13} D_{3\nu} \sin^2 2\theta_{12} \sin^2 \frac{\Delta m_{21}^2 L}{4E_\nu}$$

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$$V_{CC} = \sqrt{2} G_F N_e$$

Matter-induced mixing $\sin^2 2\theta_m$

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Technically, mass splitting depends on matter density only matters in matter-dominated limit

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$$V_{CC} = \sqrt{2} G_F N_e$$

Matter-induced mixing $\sin^2 2\theta_m$

Typical oscillation kinematic factor

□ Goal – determine if DUNE solar neutrinos give a precision test of MSW model of oscillation frequency and amplitude

□ Limited by E res, short-baseline “twilight” neutrinos carry sensitivity

Comparison to long-baseline

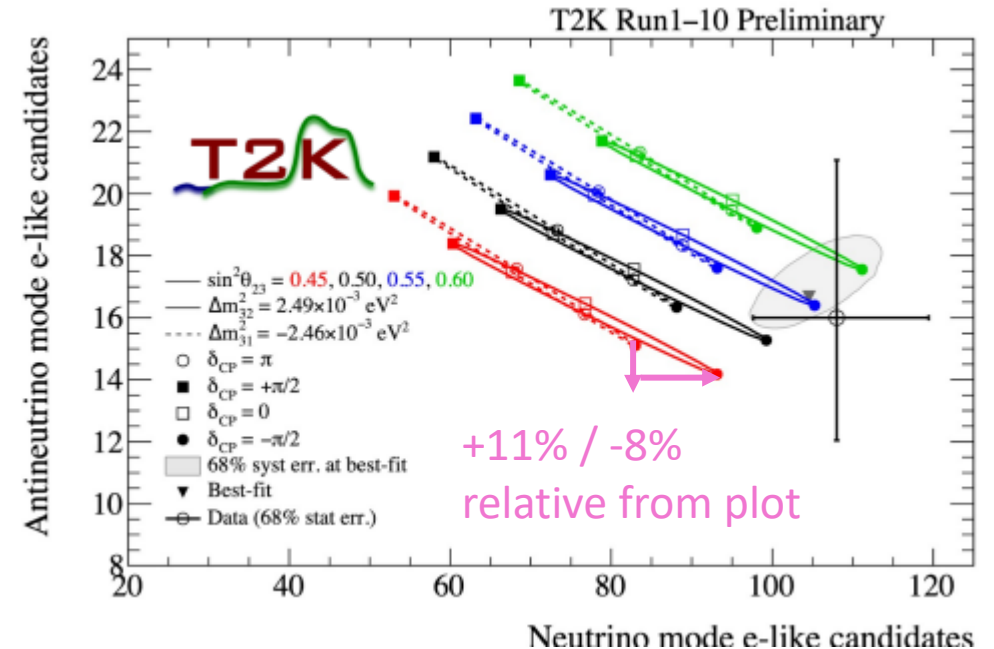
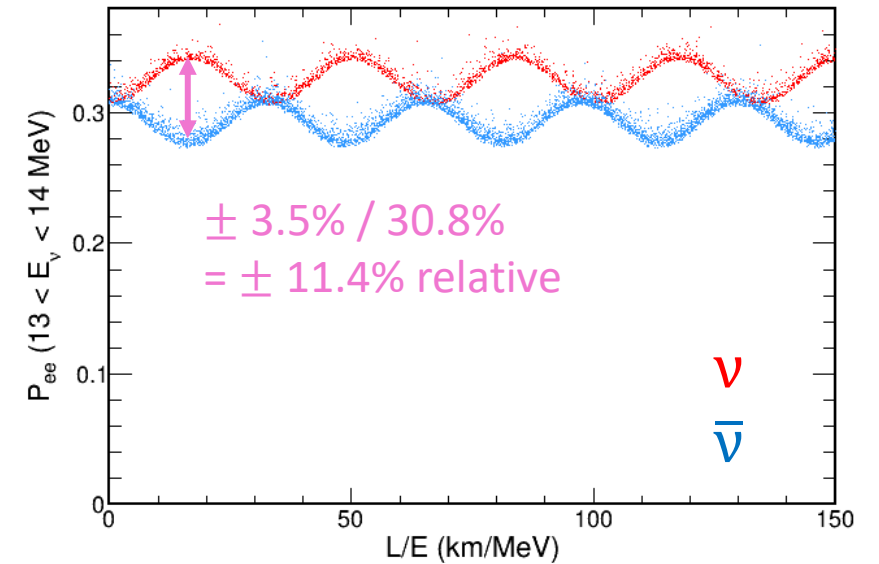
Night excess scale

$$A = \frac{V_{CC} E_\nu}{\Delta m_{ij}^2}$$

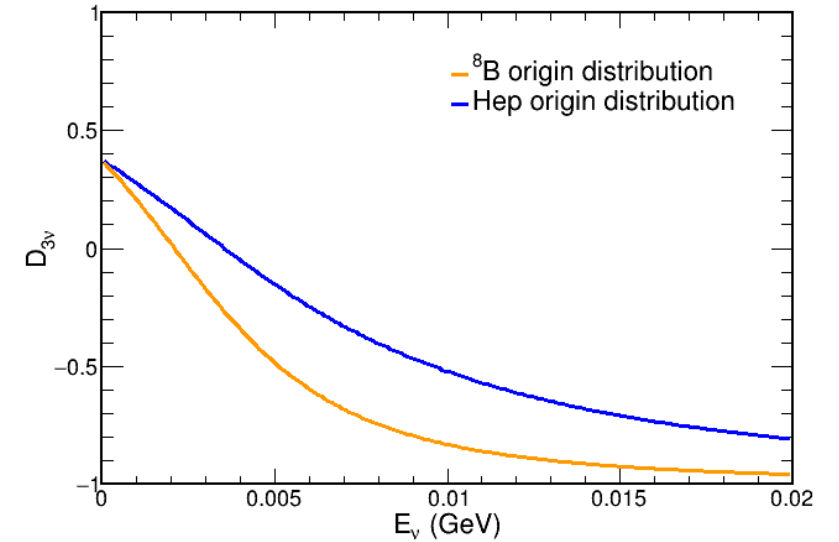
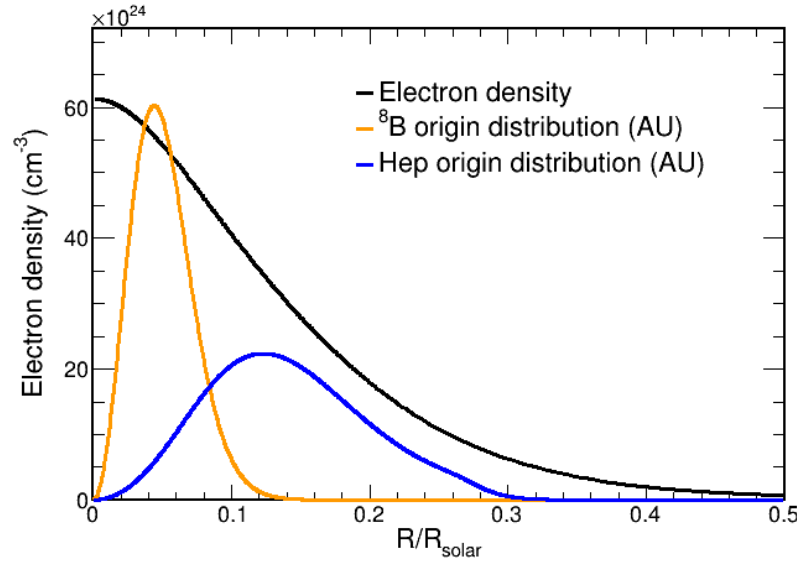
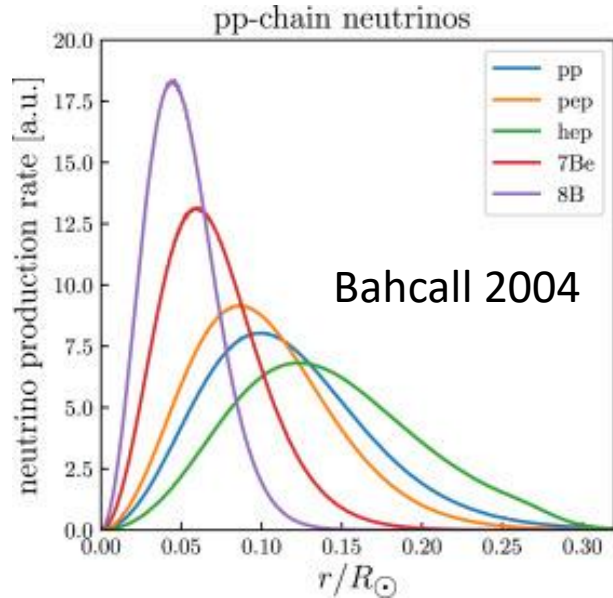
$$V_{CC} \approx 1.2e-13 \text{ eV @ } 3 \text{ gm/cm}^3$$

	V_{CC}	E_ν	Δm_{ij}^2	A
^8B Solar	1.3e-13 eV	14 MeV	7.5e-5 eV ²	0.027
DUNE LBL	1.2e-13 eV	2.8 GeV	2.4e-3 eV ²	0.14
T2(H)K LBL	1.2e-13 eV	0.6 GeV	2.4e-3 eV ²	0.030

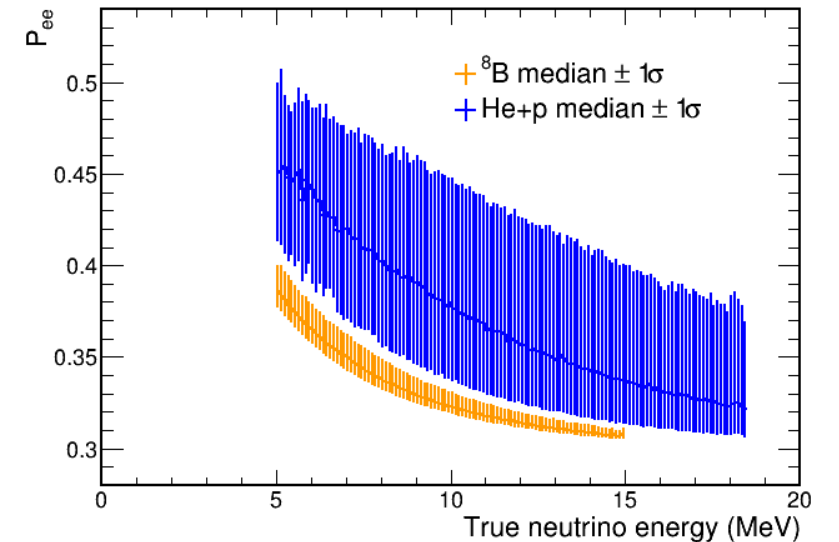
Solar neutrinos have approximately same degree of matter asymmetry as T2(H)K for long-baseline



Regions that support fusion complicates picture



- Distribution of production densities
→ distribution $|\nu_2\rangle$ of content →
distribution of critical E_{ν} for upturn
- This is a precision oscillation study,
incorporate this effect



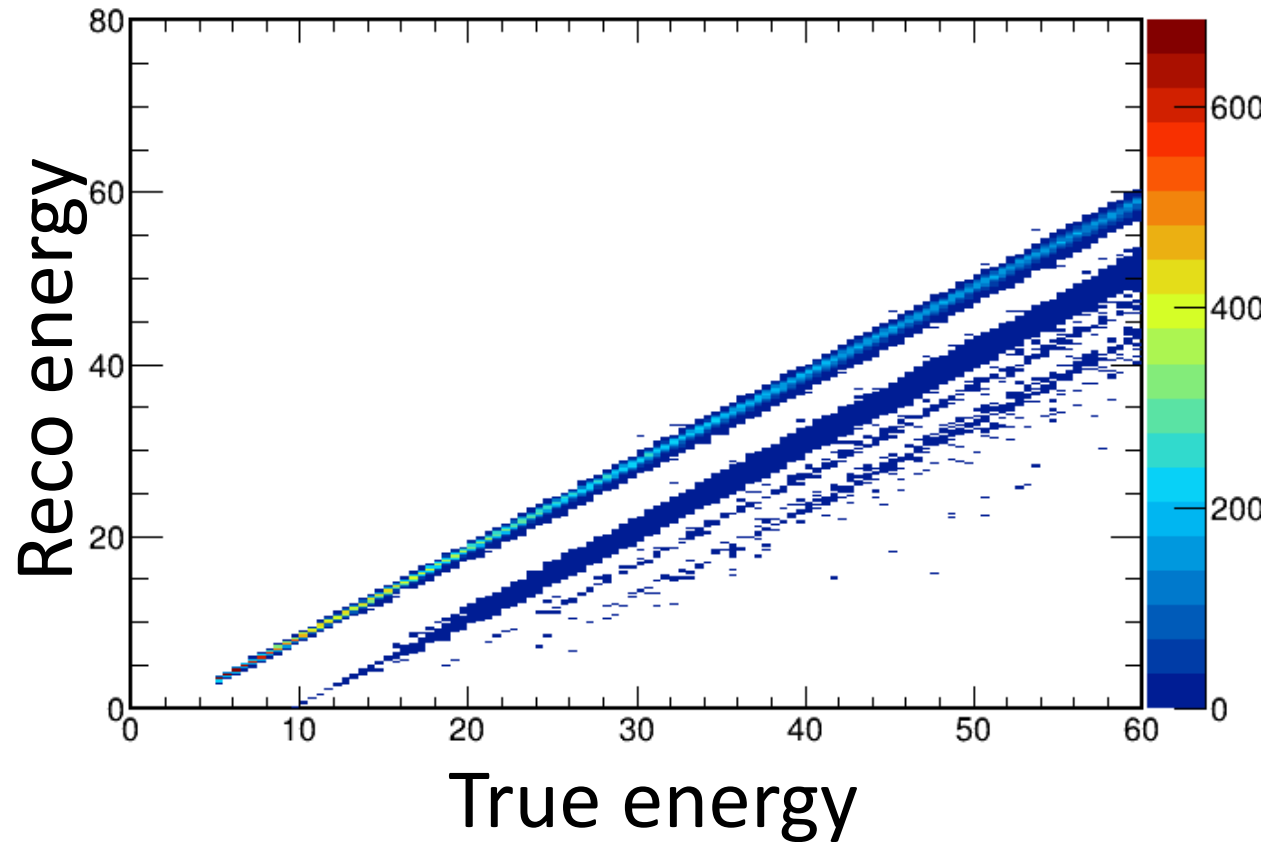
Simulating a solar neutrino sample

□ Focus on ^8B neutrinos

1. Pull electron density from spectrum radial distribution
2. Pull random nadir angle from expected yearly distribution
3. Pull random neutrino energy uniformly over flux energy range
4. Calculate $P_{ee}(r_{\text{prod}}, \cos \eta, E_\nu)$
5. Apply event weight = $P_{ee}(r_{\text{prod}}, \cos \eta, E_\nu)$
6. Continue until spectrum integral is:
 - 20.6k events / 17 kt-yrs x (10 x 17 kt-yrs) – Sergio's selected event rate

Detector resolution assumptions

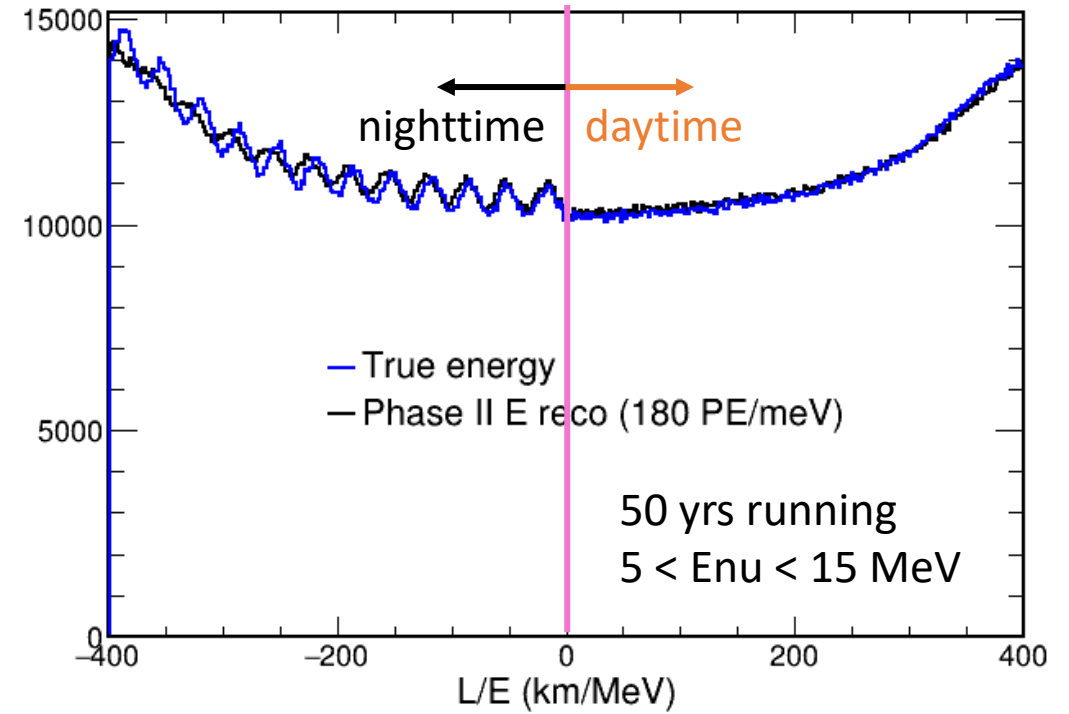
- Energy resolution matters! Use phase II detector response from Wei Shi's study
 - 75 keV for Q threshold
 - 180 PE/MeV LY
- Looking at phase I detector res from Sergio is on the docket



Predicting twilight neutrinos

Expected observed solar neutrino rate (5 – 15 MeV) as a function of baseline through Earth

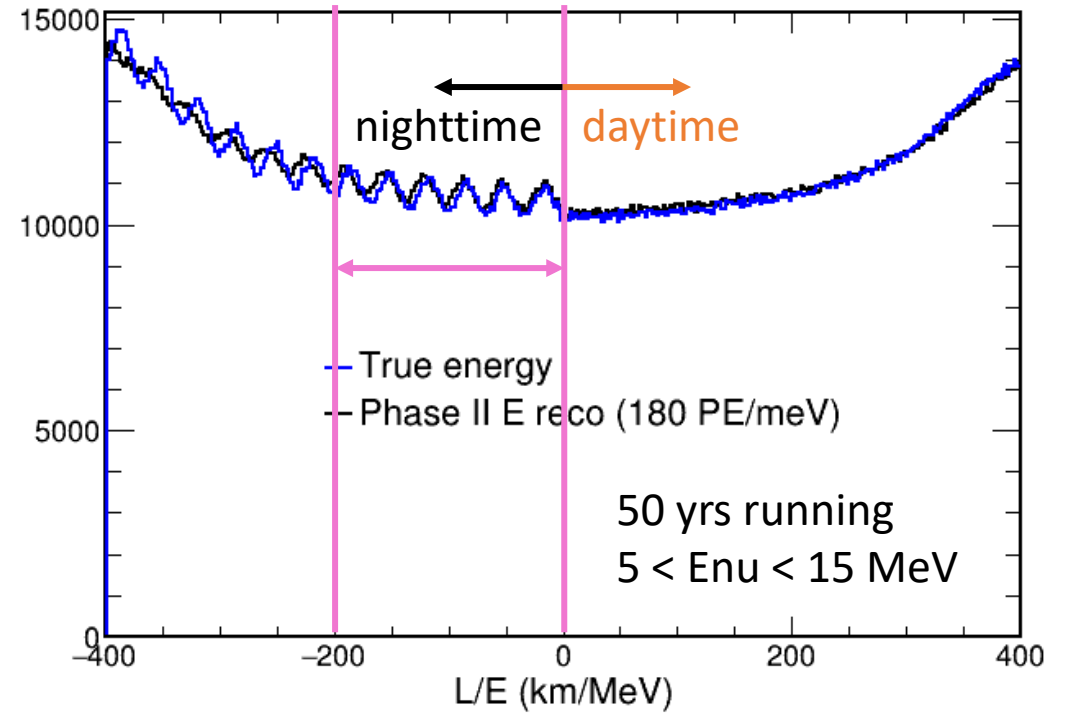
- $$P_{ee}^n - P_{ee}^d \cong C \sin^2 \frac{\Delta m_{21}^2 L}{4E_\nu}$$



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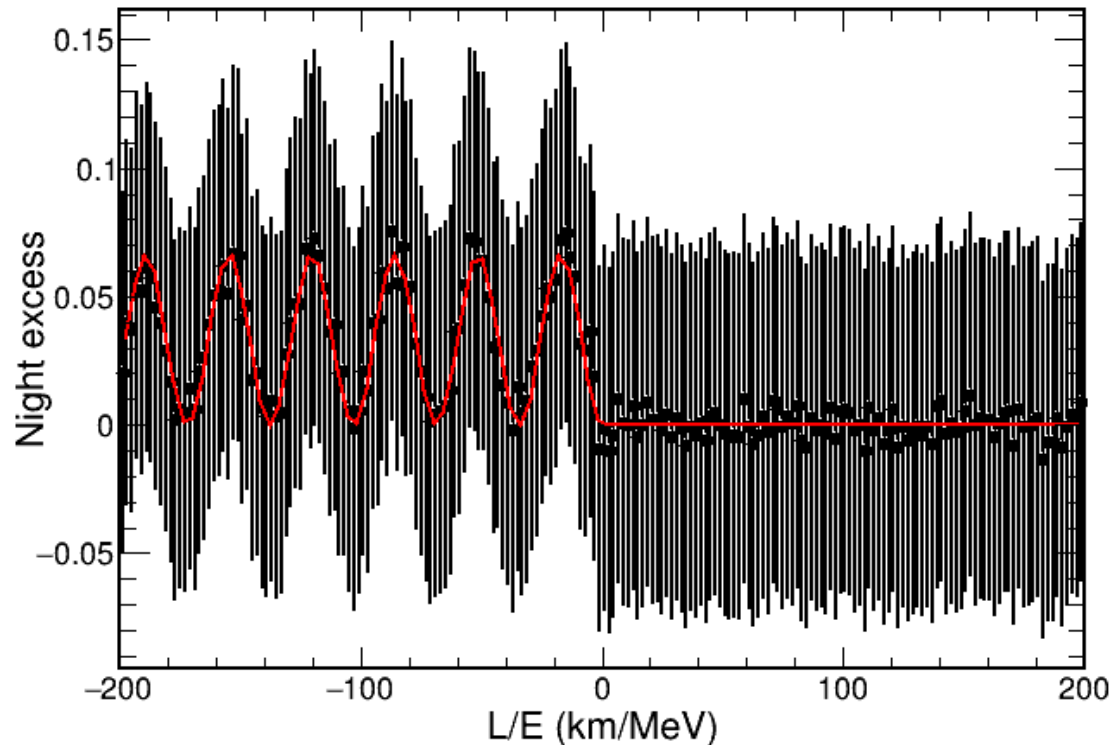


For $L/E < \sim -200$ km/MeV, energy resolution smear out oscillation behavior even for phase II resolution

- How well can we determine amplitude / frequency of oscillation for $L/E < 200$

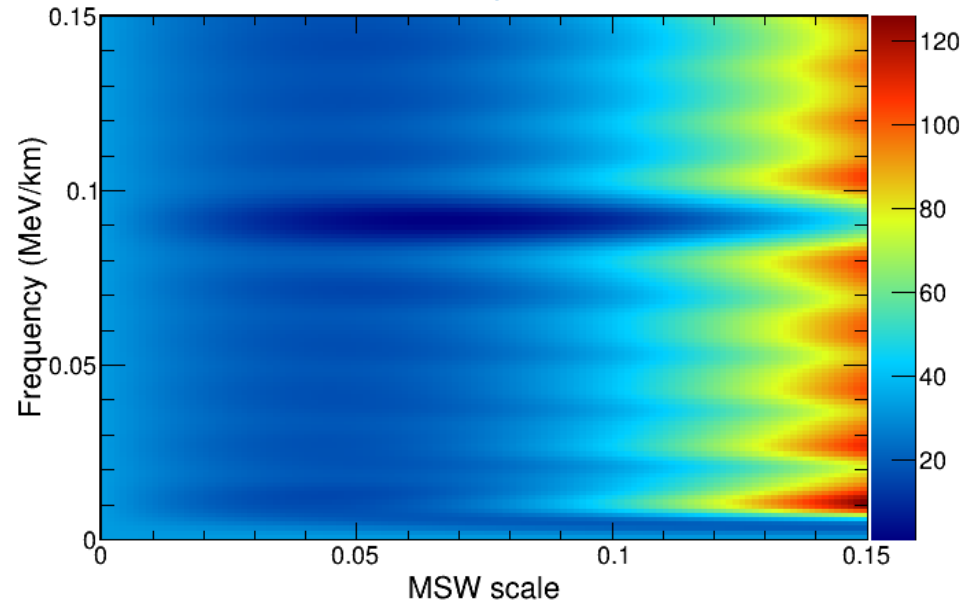
Mock data – 10 yrs of phase II

$$5 < E_{\text{reco}} < 15 \text{ MeV}$$

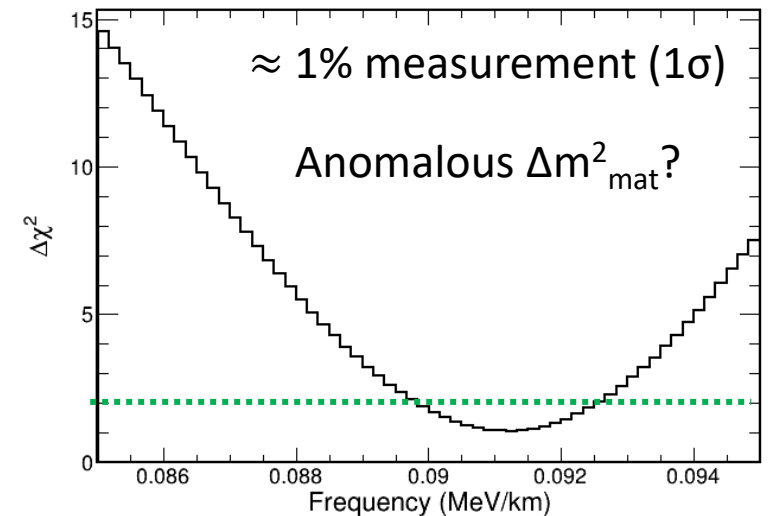
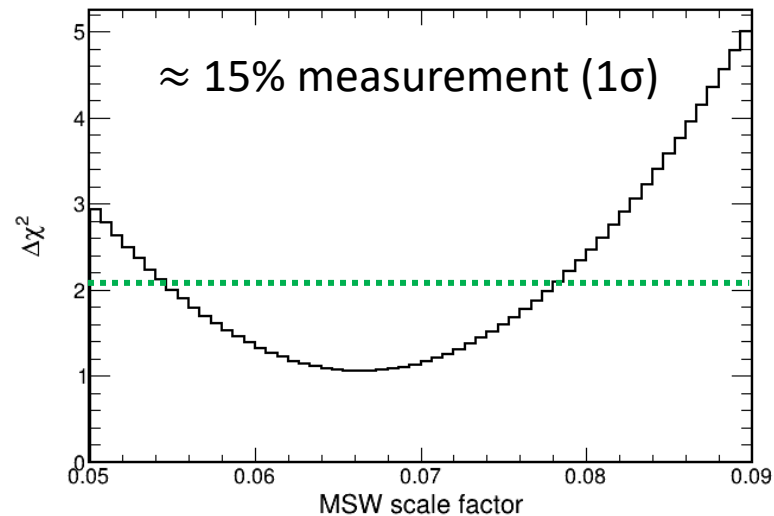
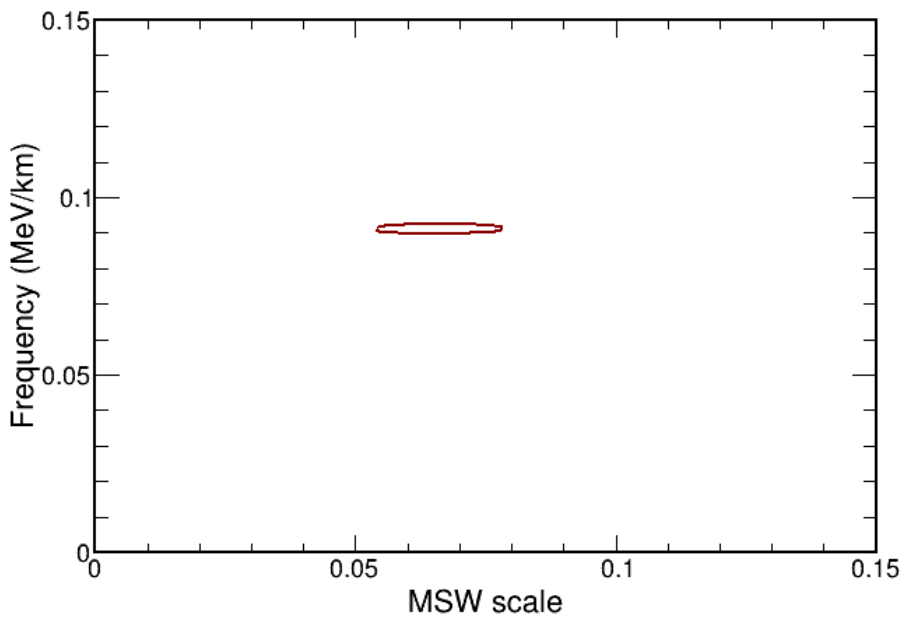


- ❑ Correct for nadir distribution shape, fit to polynomial: $p_0 + p_2x^2 + p_4x^4$ and apply to baseline prediction for night
- ❑ Normalize by day-time prediction to determine the night excess
- ❑ Select $L/E < 200 \text{ km/MeV}$
 - With $E \approx 10 \text{ MeV} \rightarrow L < 2000 \text{ km}$
 $\rightarrow < 40 \text{ min after sunset / before dawn}$
- ❑ Six oscillation maxima are observed

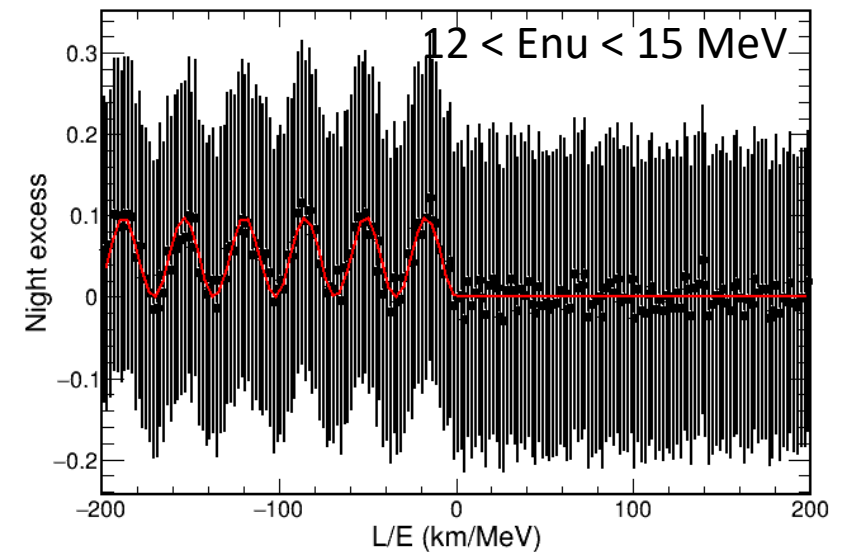
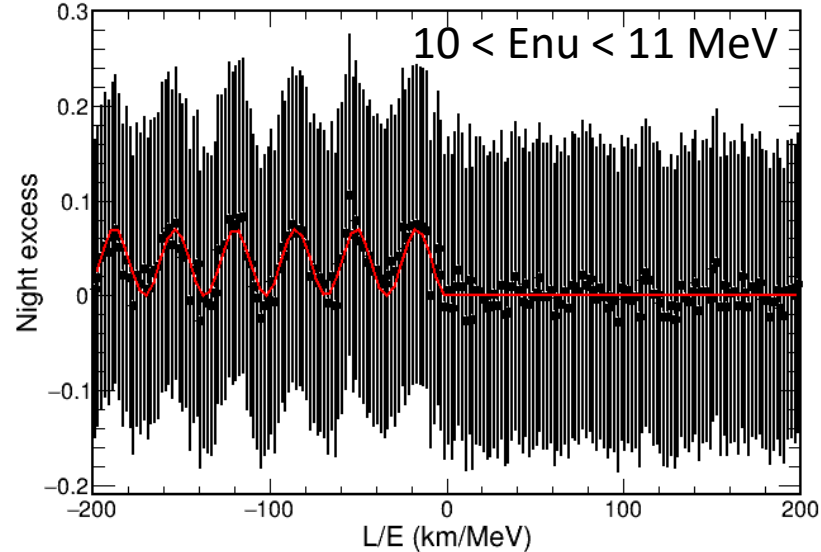
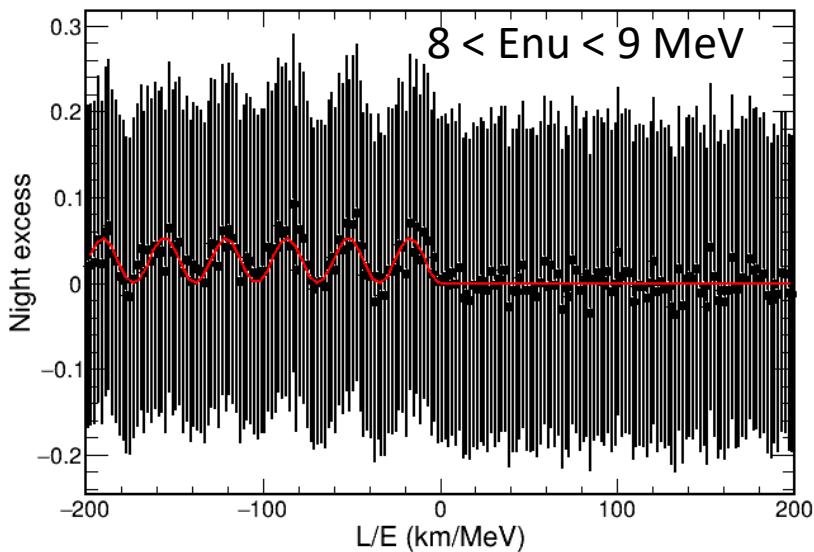
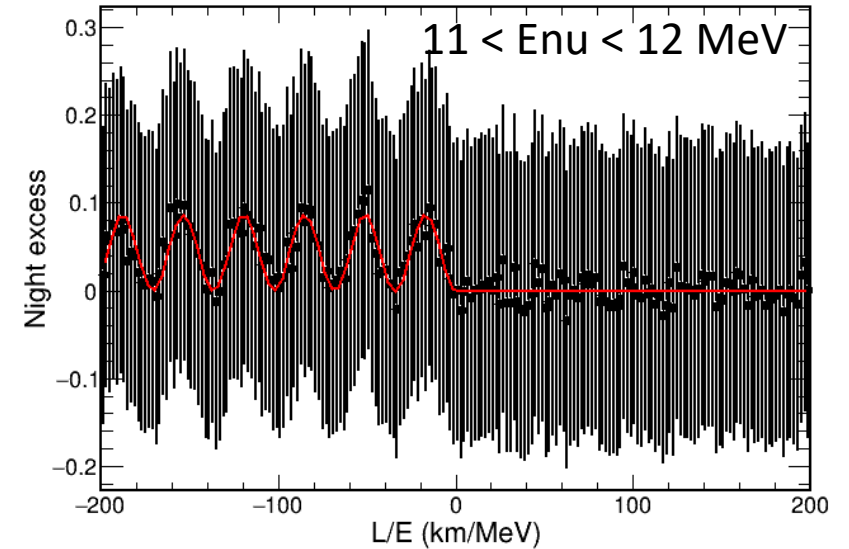
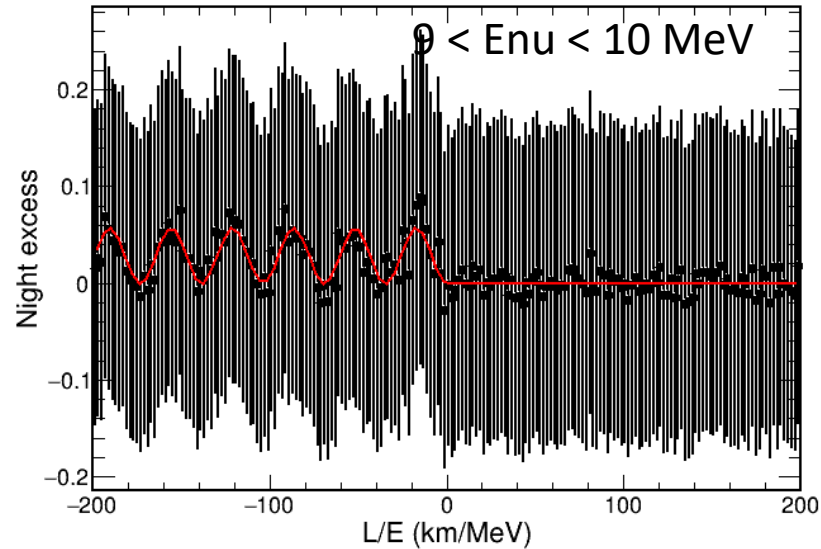
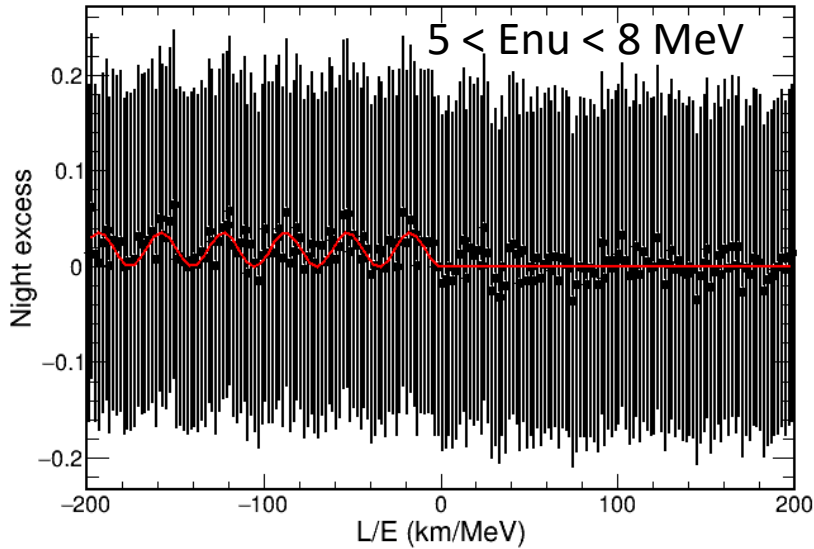
First attempt at oscillation contours



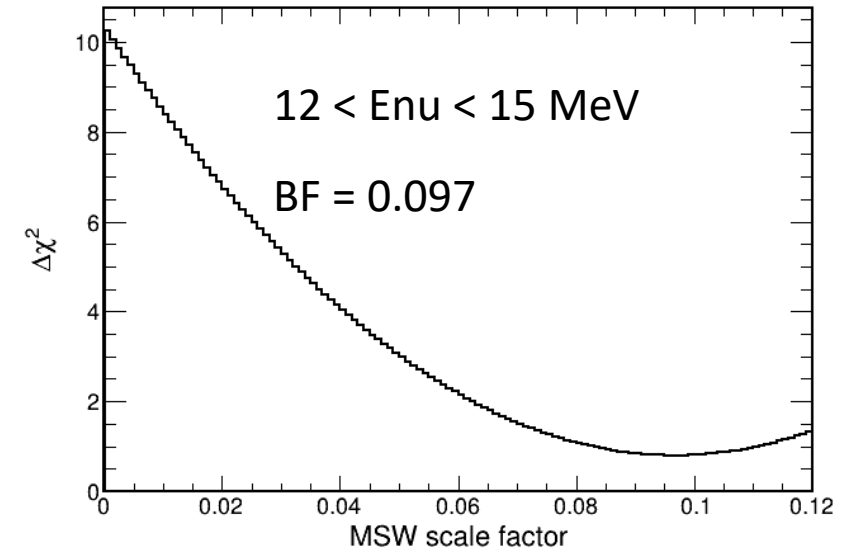
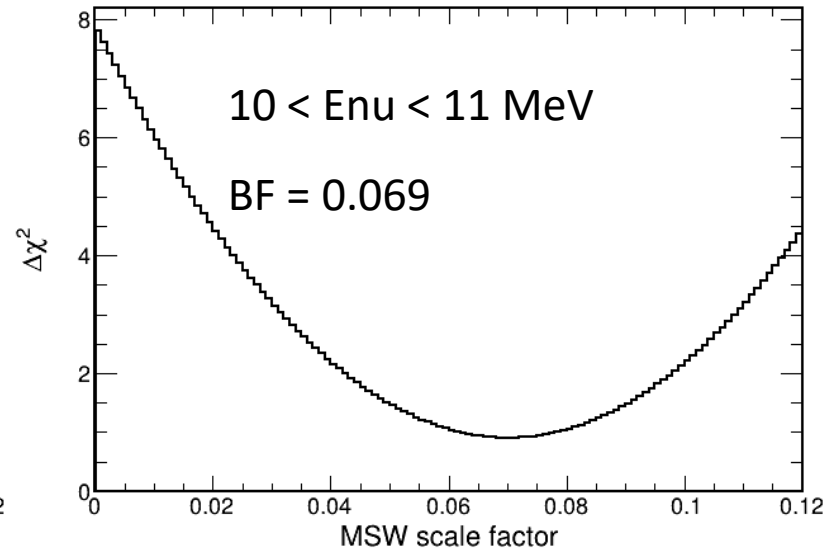
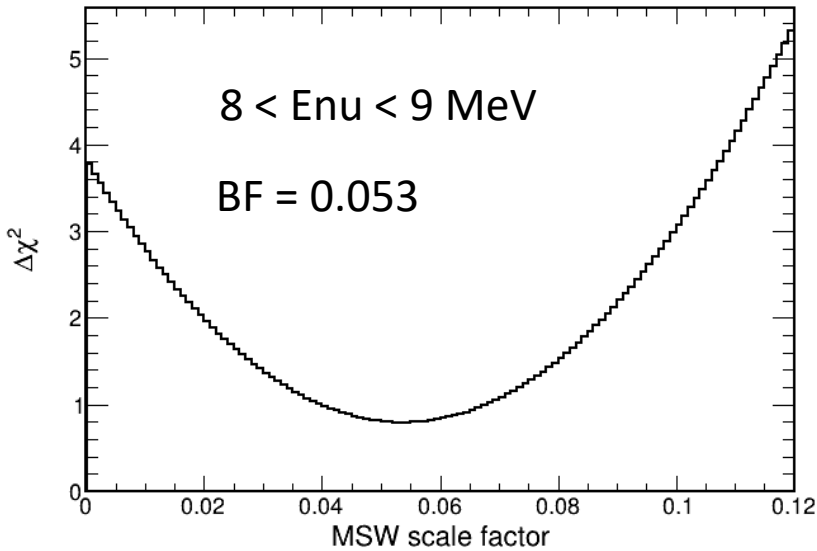
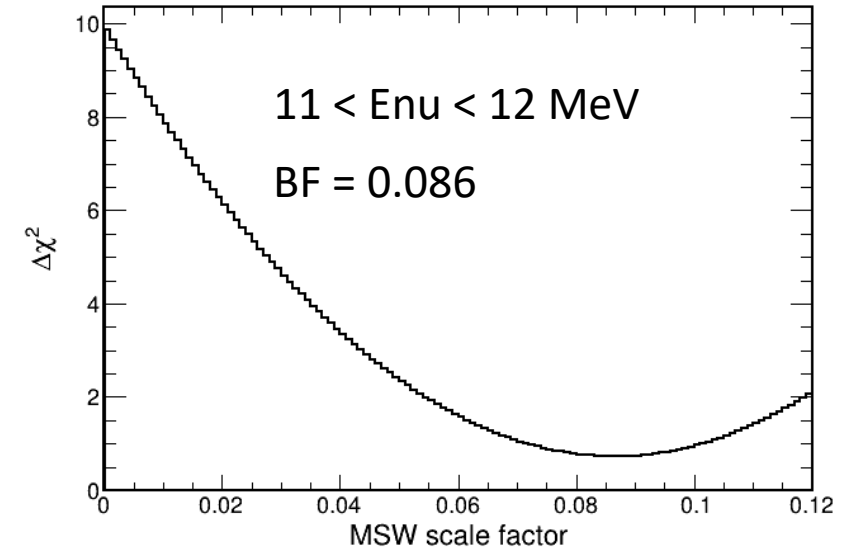
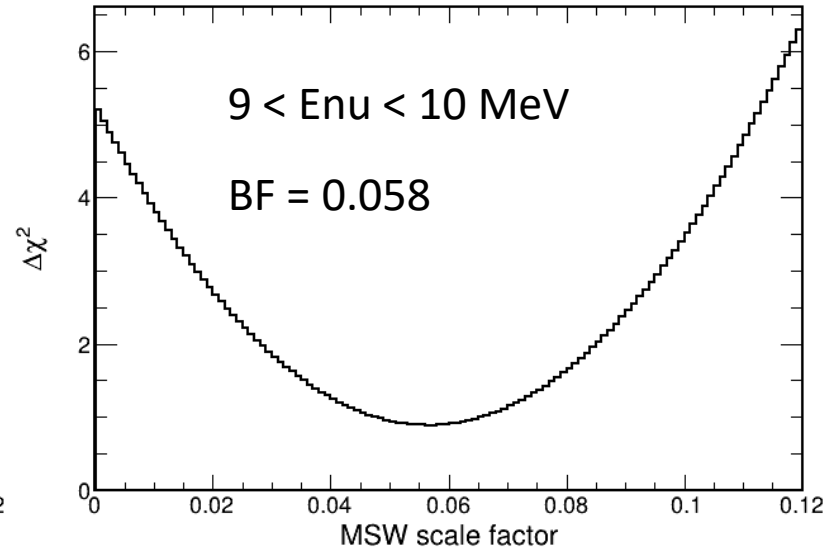
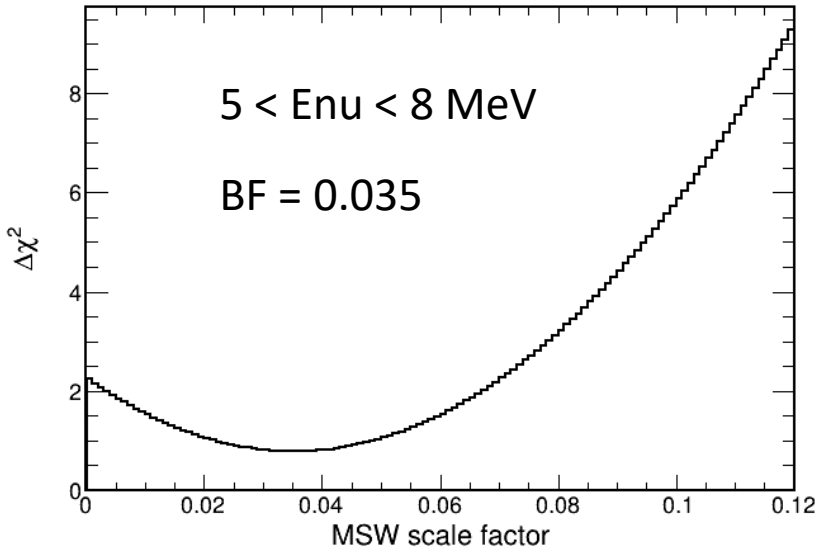
- 10 yrs of phase II solar neutrino data
- Great 1σ measurement of Δm^2
 - Can test adjustment due to matter effects
- Scale factor less-well determined, can be improved with late night neutrinos



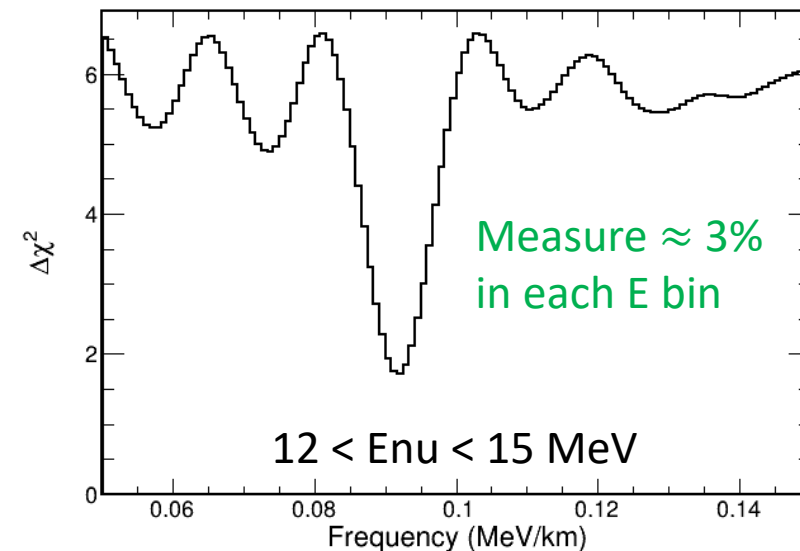
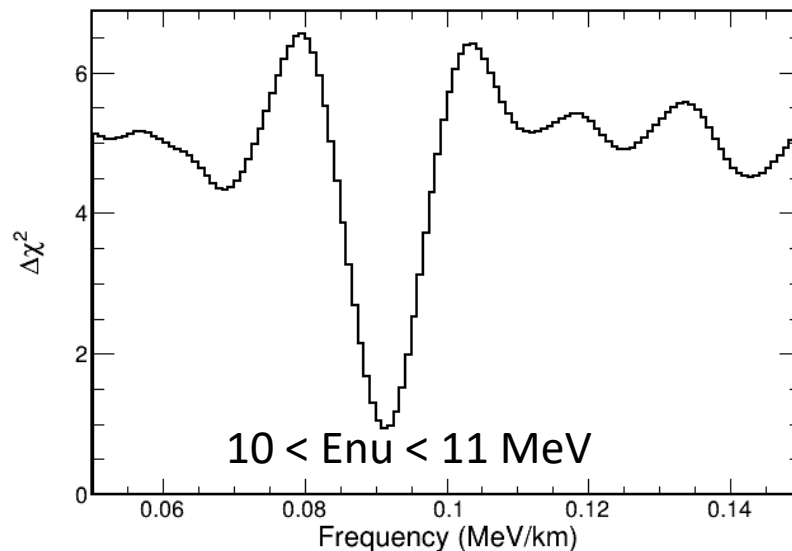
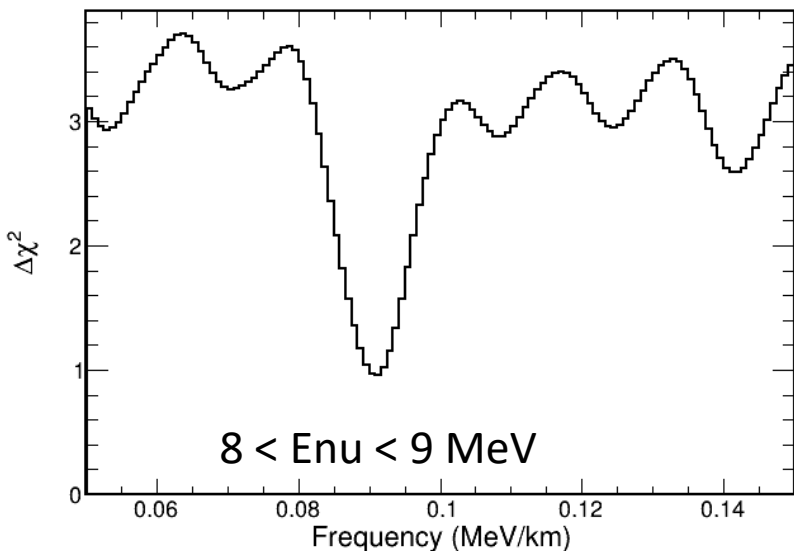
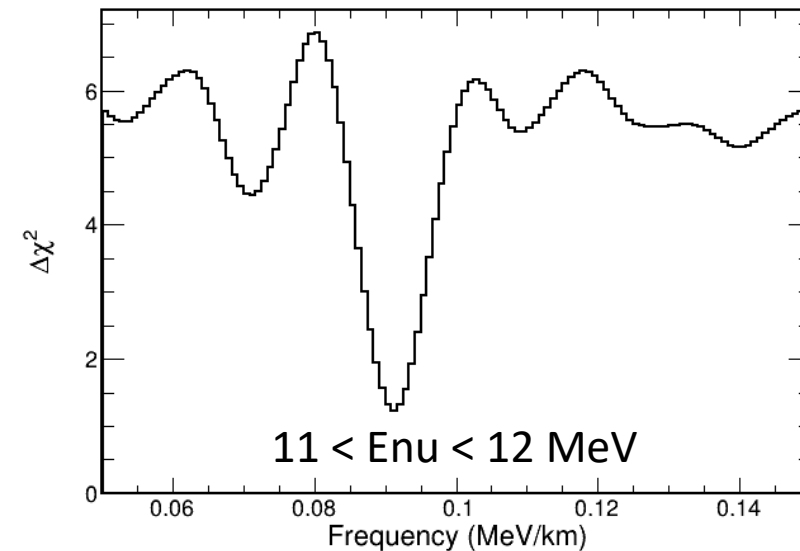
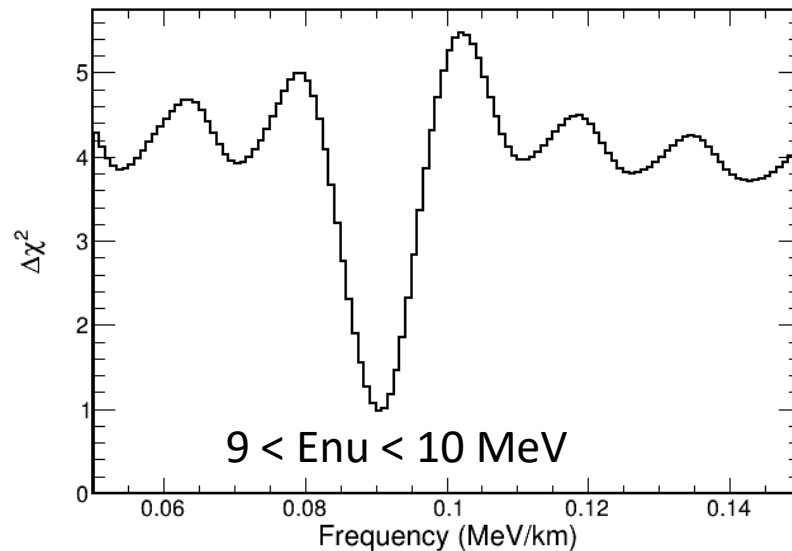
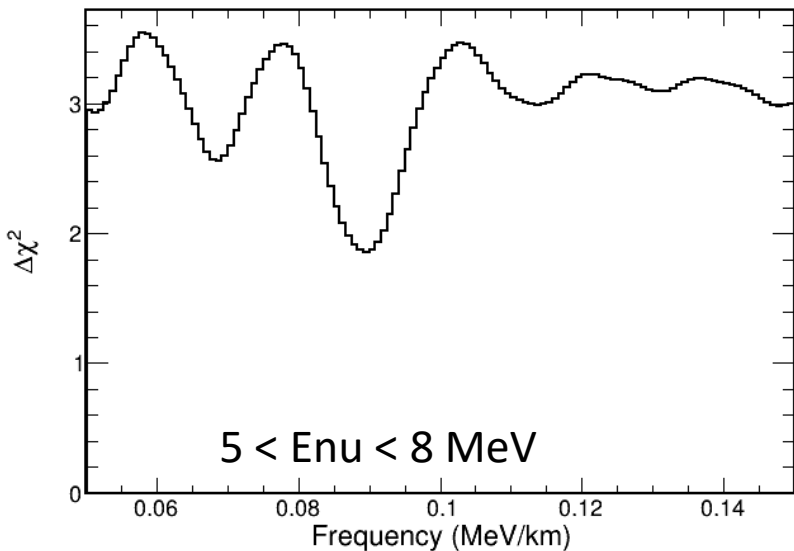
Spectra in energy bins



Sensitivity in energy bins – MSW scale



Sensitivity in energy bins – frequency



Impressions and next steps

- ❑ New physics available if we have 1 module with excellent resolution
 - Phase II owns this physics!
- ❑ Golden analysis channel allows analysis of the solar upturn giving independent measurement of Δm^2
 - Good resolution tests energy dependence of upturn
- ❑ Twilight neutrinos also Δm^2 (1% at 1σ) – maybe test departure of Δm^2 from vacuum value?
 - With JUNO, multiple probes overconstrain the solar neutrino picture allowing new physics tests
- ❑ Sensitivity to MSW scale less good (15% at 1σ) – resolution-averaged neutrinos from late at night can also measure this parameter
- ❑ Very basic studies – need more study with realistic detector models