

Physics of EDMs

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Review: M. Pospelov & AR, Ann. Phys. 318, 119 (2005) [hep-ph/0504231]

D. McKeen, M. Pospelov & AR, arXiv:1208.4597, 1303.1172

Physics of EDMs

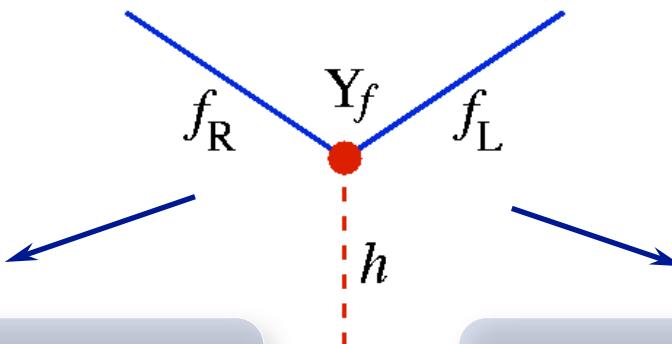
Recent FNAL workshop

The logo features a red sphere with a blue base, representing a particle. A vertical axis passes through it, with a purple arrow labeled 'S' pointing upwards. A horizontal arrow labeled ' \vec{d} ' points along the axis. Below the sphere, there is a small circular motion arrow. To the left of the sphere, three equations are listed: $t \rightarrow -t$, $\vec{s} \rightarrow -\vec{s}$, and $\vec{d} \rightarrow \vec{d}$. Below these, the equation $\vec{d} = d \frac{\vec{S}}{S}$ is given, followed by $d \rightarrow -d \rightarrow 0$.

**Winter Workshop on Electric Dipole
Moments (EDMs13)**

13-15 February 2013 *Fermi National
Accelerator Laboratory*
US/Central timezone

CP Violation in the Standard Model



$$\sin(\delta_{\text{KM}}) \propto \text{Arg Det}[Y_u Y_u^\dagger, Y_d Y_d^\dagger]$$

$$\bar{\theta}_{\text{QCD}} \sim \text{Arg Det}[Y_u Y_d]$$

(in a convenient basis)

$$\delta_{\text{KM}} \sim \mathcal{O}(1)$$

$$\theta < 10^{-10} !$$

Explains CP-violation in K and B meson mixing and decays

Constrained experimentally
(strong CP problem)

Do we anticipate other CP-odd sources ?

- Required by baryogenesis (Sakharov conditions)
- Generic with extra degrees of freedom

Experimental EDM Limits

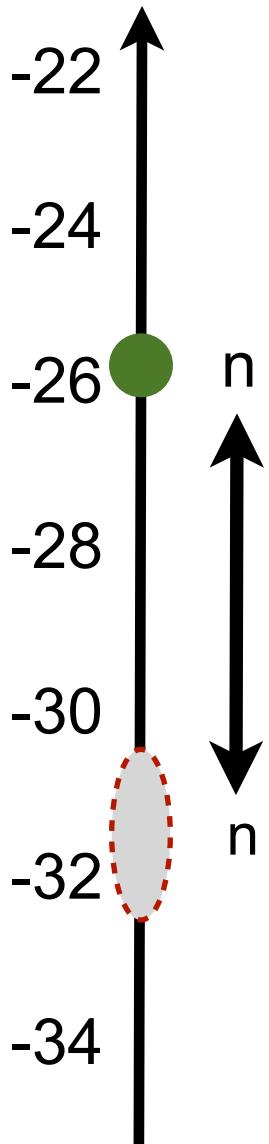
$$H = -d\vec{E} \cdot \frac{\vec{S}}{S}$$

- EDMs are powerful (amplitude-level) probes for new flavour-diagonal (T,P) violating sources
- Best current limits from neutrons, para- and dia-magnetic atoms and molecules

Neutron EDM	$ d_n < 3 \times 10^{-26} e \text{ cm}$	[Baker et al. '06]
Thallium EDM (paramagnetic)	$ d_{Tl} < 9 \times 10^{-25} e \text{ cm}$	[Regan et al. '02]
YbF “EDM” (paramagnetic)	$ “d_{YbF}” < 1.4 \times 10^{-21} e \text{ cm}$	[Hudson et al. '11]
Mercury EDM (diamagnetic)	$ d_{Hg} < 3 \times 10^{-29} e \text{ cm}$	[Griffith et al. '09]

SM (CKM) background

$\log(d [e \text{ cm}])$



SM (CKM) contribution is (at least)
4-5 orders of magnitude below the
current neutron sensitivity, and
even lower for the atomic EDMs

⇒ negligible CKM background

[Khriplovich & Zhitnitsky '82;
McKellar et al '87;
Mannel & Uraltsev '12]

Schematic new physics sensitivity

Negligible SM (CKM) contribution \Rightarrow EDMs are rather special in having clean (dim-6) amplitude-level sensitivity to new physics

- Generic corrections to the SM:

$$\Delta\Gamma, \Delta\sigma \sim \text{Re}(A_{\text{SM}} A_{\text{NP}}^*) + \delta_{\text{SM}}^{\text{error}}$$

linear dependence on A_{NP} , but sensitivity limited by SM precision

- Tests of SM forbidden/suppressed processes:

- often quadratic in the amplitude, e.g. LFV

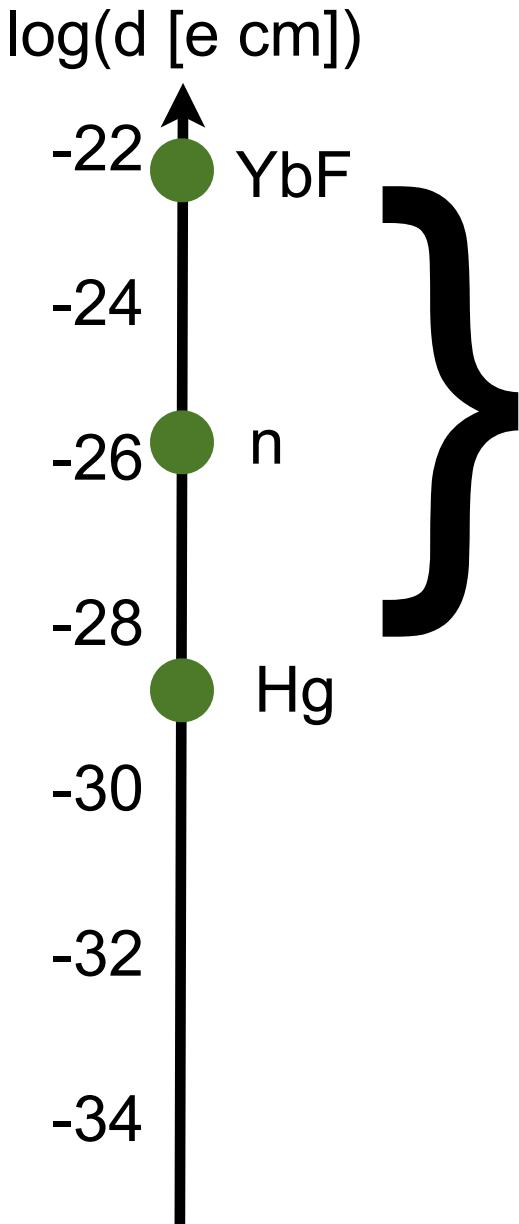
$$\Delta\Gamma \sim |A_{\text{NP}}|^2$$

- EDMs are linear!

$$\Delta d \sim A_{\text{NP}}$$

\Rightarrow Better reward for improvements in experimental sensitivity

Schematic view of the bounds



Difference of more than 6 orders of magnitude, but the sensitivity to underlying CP-odd sources is actually very similar...

CP-odd operator expansion (at $\sim 1\text{GeV}$)

(Flavor-diagonal) CP-violating operators at $\sim 1\text{GeV}$

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{c_n}{\Lambda^{d-4}} \mathcal{O}_d^{(n)}$$

$$\mathcal{L}_{\text{dim } 4} \supset \bar{\theta} \alpha_s G \tilde{G}$$

$$\bar{\theta} = \theta_0 - \text{ArgDet}(M_u M_d) \equiv \theta_0 - \theta_q$$

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$$\mathcal{L}_{\text{dim } 4} \supset \bar{\theta} \alpha_s G \tilde{G}$$

$$d_i \sim c Y_i \frac{v}{\Lambda^2}$$

$$\mathcal{L}^{\text{"dim 6"}} \supset \sum_{q=u,d,s} \left(d_q \bar{q} F \sigma \gamma_5 q + \tilde{d}_q \bar{q} G \sigma \gamma_5 q \right) + \sum_{l=e,\mu} d_l \bar{l} F \sigma \gamma_5 l$$

$$\mathcal{L}_{\text{dim } 6} \supset w g_s^3 G G \tilde{G} + \sum_{q,\Gamma} C_{qq} (\bar{q} \Gamma q)_{LL} (\bar{q} \Gamma q)_{RR}$$

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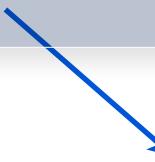
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$$C_{ij} \sim c Y_i Y_j \frac{v^2}{\Lambda^4}$$

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(Flavor-diagonal) CP-violating operators at $\sim 1\text{GeV}$

$$\mathcal{L}_{\text{eff}} = \sum_n \frac{c_n}{\Lambda^{d-4}} \mathcal{O}_d^{(n)}$$

$$d_N \bar{N} F \sigma \gamma_5 N + \bar{g}_{\pi NN}^{(1)} \pi^0 \bar{N} N + \dots$$

$$\mathcal{L}_{\text{dim } 4} \supset \bar{\theta} \alpha_s G \tilde{G}$$

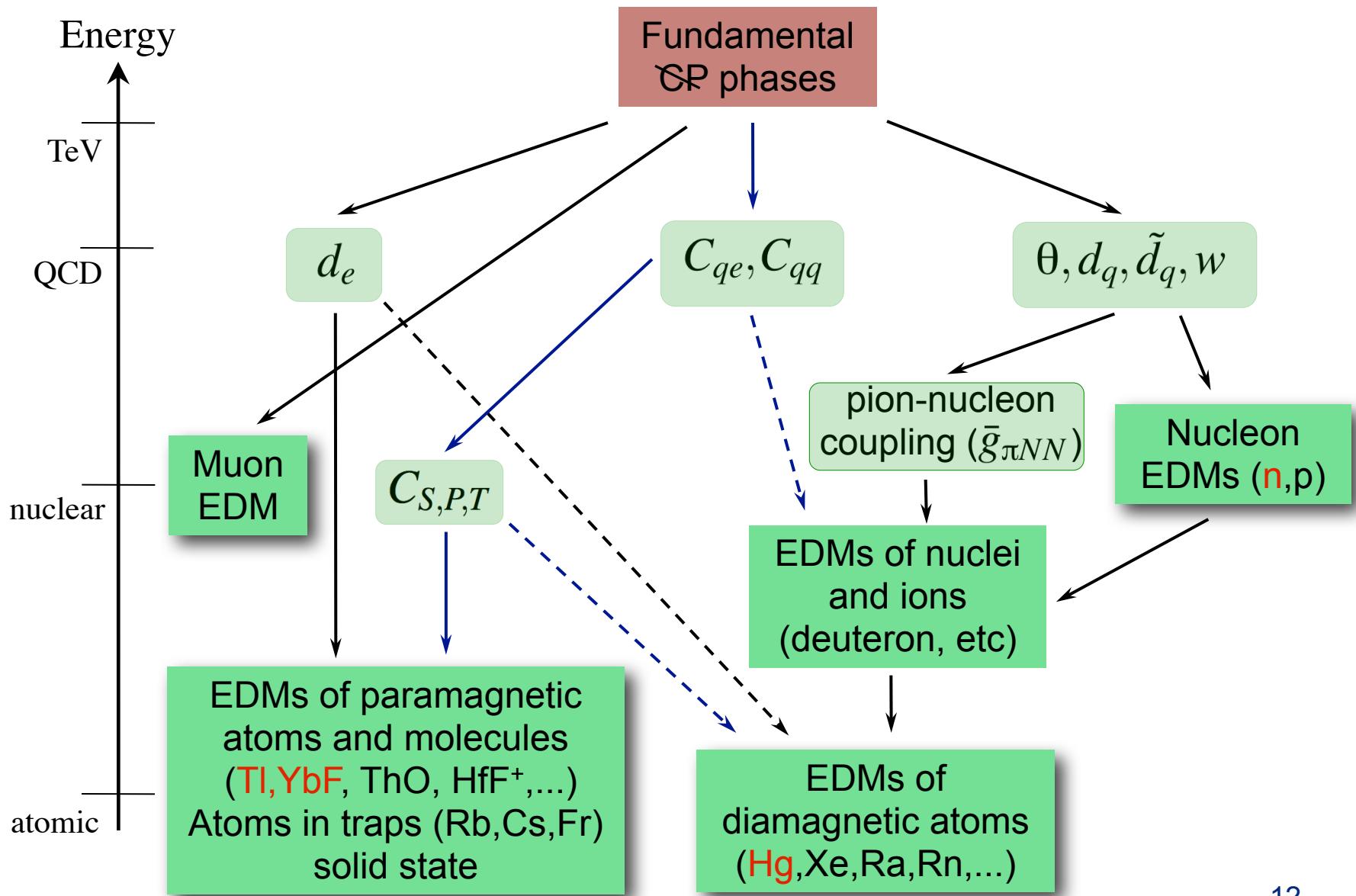
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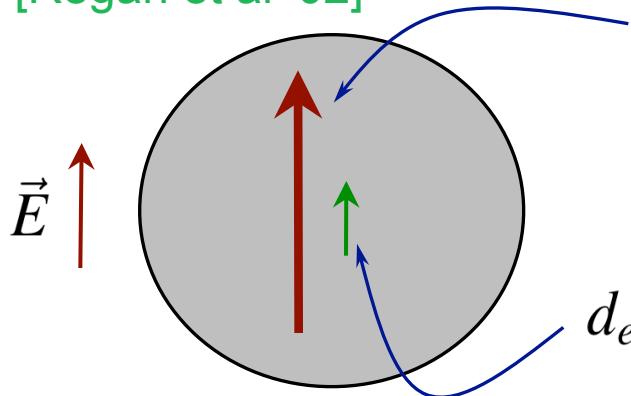
$$C_S \bar{N} N \bar{e} i \gamma_5 e + \dots$$

EFT hierarchy



Paramagnetic EDMs - “Schiff enhancement”

Atoms (e.g. Tl [Berkeley])
[Regan et al '02]



(relativistic violation of Schiff screening)

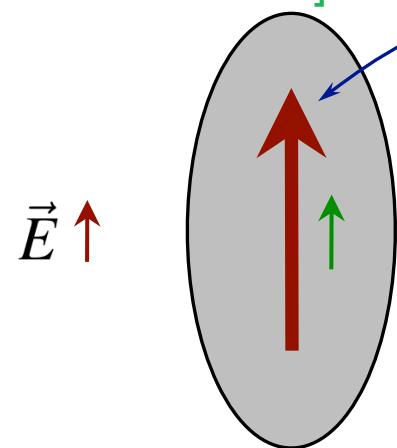
$$\alpha^2 Z^3 \vec{E}$$

[Salpeter '58;
Sandars '65]

$$d_{Tl} \sim -20\alpha^2 Z^3 d_e + \mathcal{O}(C_S)$$

≈ 585 [Liu & Kelly '92]

Polar molecules (e.g. YbF [Imperial]) [also ThO [Harvard/Yale]]
[Hudson et al '11]



Nonlinear function of E_{ext}

$$\Delta E_{\text{YbF}} \sim \mathcal{E}_{\text{eff}}(E_{\text{ext}}) d_e + \mathcal{O}(C_S)$$

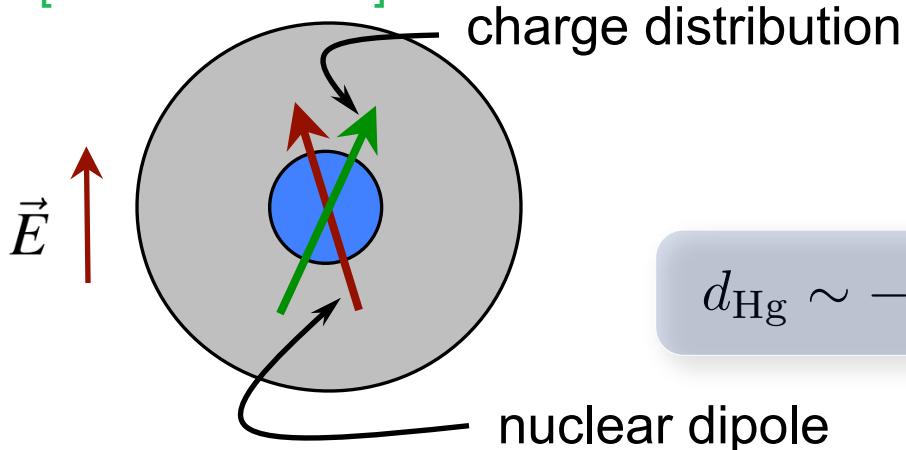
$$\text{“}d_{\text{YbF}}\text{”} \sim 10\alpha^2 Z^3 \frac{M_{\text{mol}}}{m_e} d_e + \mathcal{O}(C_S)$$

[Sushkov &
Flambaum, '78]

Diamagnetic EDMs - “Schiff suppression”

Atoms (e.g. Hg [Washington]) (finite size violation of Schiff screening)

[Griffith et al '09]



$$d_{\text{Hg}} \sim -3 \times 10^{-17} S[e \text{ fm}^2] + \mathcal{O}(d_e, C_{qe}, C_{qq})$$

$$d_{\text{Hg}} \sim 10Z^2(R_N/R_A)^2 d_{\text{nuc}}$$

$$O(10^{-3})$$

[Flambaum et al '86;
Dzuba et al. '02]

Schiff moment [Schiff '63]

$$S = S(\bar{g}_{\pi NN}^{(i)}, d_N, \dots)$$

[Flambaum et al. '86;
Dmitriev & Senkov '03;
de Jesus & Engel '05;
Ban et al '10]

• Octopole enhancements (e.g. Ra, Rn)

[R. Holt, Z.-T. Lu, et al,
T. Chupp et al]

- Schiff moment $O(100-1000)$ larger than Hg

[Flambaum et al.]

Nuclear EDMs - avoiding Schiff screening

- Neutron EDM via UCN bottles [...., PSI, Sussex/ILL, ORNL, Osaka/TRIUMF]

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- Nuclear EDMs (e.g. p,D, ${}^3\text{He}$,...) in storage rings [BNL, FNAL?
COSY/Julich]

$$d_{Hg} \sim 10Z^2(R_N/R_A)^2 d_{\text{nuc}}$$



O(10^{-3}) suppression could be avoided with a direct measurement of the nuclear EDM.

Nuclear EDMs - avoiding Schiff screening

- Neutron EDM via UCN bottles [...., PSI, Sussex/ILL, ORNL, Osaka/TRIUMF]
- Nuclear EDMs (e.g. p,D, ${}^3\text{He}$,...) in storage rings [BNL, FNAL?
COSY/Julich]

- [Best current limit on muon EDM]

$$|d_\mu| < 2 \times 10^{-19} \text{ ecm} \quad [\text{Muon g-2 Collab, BNL '09}]$$

- e.g. proton - similar sensitivity to the neutron ($d \leftrightarrow u$)

- e.g. deuteron

$$\begin{aligned} d_D &= (d_n + d_p)(\bar{\theta}, d_q, \tilde{d}_q) + d_D^{\pi NN}(\bar{\theta}, \tilde{d}_q) \\ &\approx -2 \times 10^{-14} \bar{g}_{\pi NN}^{(1)}(\bar{\theta}, \tilde{d}_q) \text{ e cm} + \mathcal{O}(\bar{g}_{\pi NN}^{(0)}) \\ &\approx -5e(\tilde{d}_d - \tilde{d}_u) + \dots \end{aligned}$$

via η - π mixing

- extended to other light nuclei (e.g. ${}^3\text{H}$, ${}^3\text{He}$) in recent work
[Stetcu et al '08, de Vries et al '11]

Dependences

1. TI EDM & YbF “EDM” (paramagnetic) (atomic/molecular)

$$d_{Tl} \sim -585d_e - 2e \sum_{q=d,s,b} C_{qe}/m_q$$

$$\text{“}d_{YbF}\text{”} \sim 1.4 \times 10^6 d_e + \mathcal{O}(C_S)$$

[Liu & Kelly '92; Khatsymovsky et al. '86]

Nonlinear function of E_{ext} [Kozlov et al. 94-98]

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2. neutron/proton EDM (chiralPT, NDA, QCD sum rules, ...) $\Rightarrow |\theta| < 10^{-10}$

$$d_n(\bar{\theta}) \sim 3 \times 10^{-16} \bar{\theta} \text{ ecm}$$

$$d_p(\bar{\theta}) \sim -4 \times 10^{-16} \bar{\theta} \text{ ecm}$$

$$d_n^{(PQ)} \sim (0.4 \pm 0.2)[4d_d - d_u + 2.7e(\tilde{d}_d + 0.5\tilde{d}_d) + \dots] + \mathcal{O}(d_s, w, C_{qq})$$

$$d_p^{(PQ)} \sim (0.4 \pm 0.2)[4d_u - d_d - 5.3e(\tilde{d}_u + 0.13\tilde{d}_d) + \dots] + \mathcal{O}(d_s, w, C_{qq})$$

[Pospelov
& AR
'99,'00]

NB: precision limited by: sum rules analysis, s-quark content, nucleon coupling,...

[Hisano
et al '12]

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[Pospelov & AR '99,'00]

NB: precision limited by: sum rules analysis, s-quark content, nucleon coupling,...

[Hisano et al '12]

3. Hg EDM (diamagnetic) (atomic+nuclear+QCD)

$$d_{Hg} \sim 10^{-3} d_{\text{nuc}} \sim -3 \times 10^{-17} S(\bar{g}_{\pi NN}^{(0,1,2)}, d_n, d_p) [\text{e fm}^3] + \mathcal{O}(d_e, C_{qq})$$

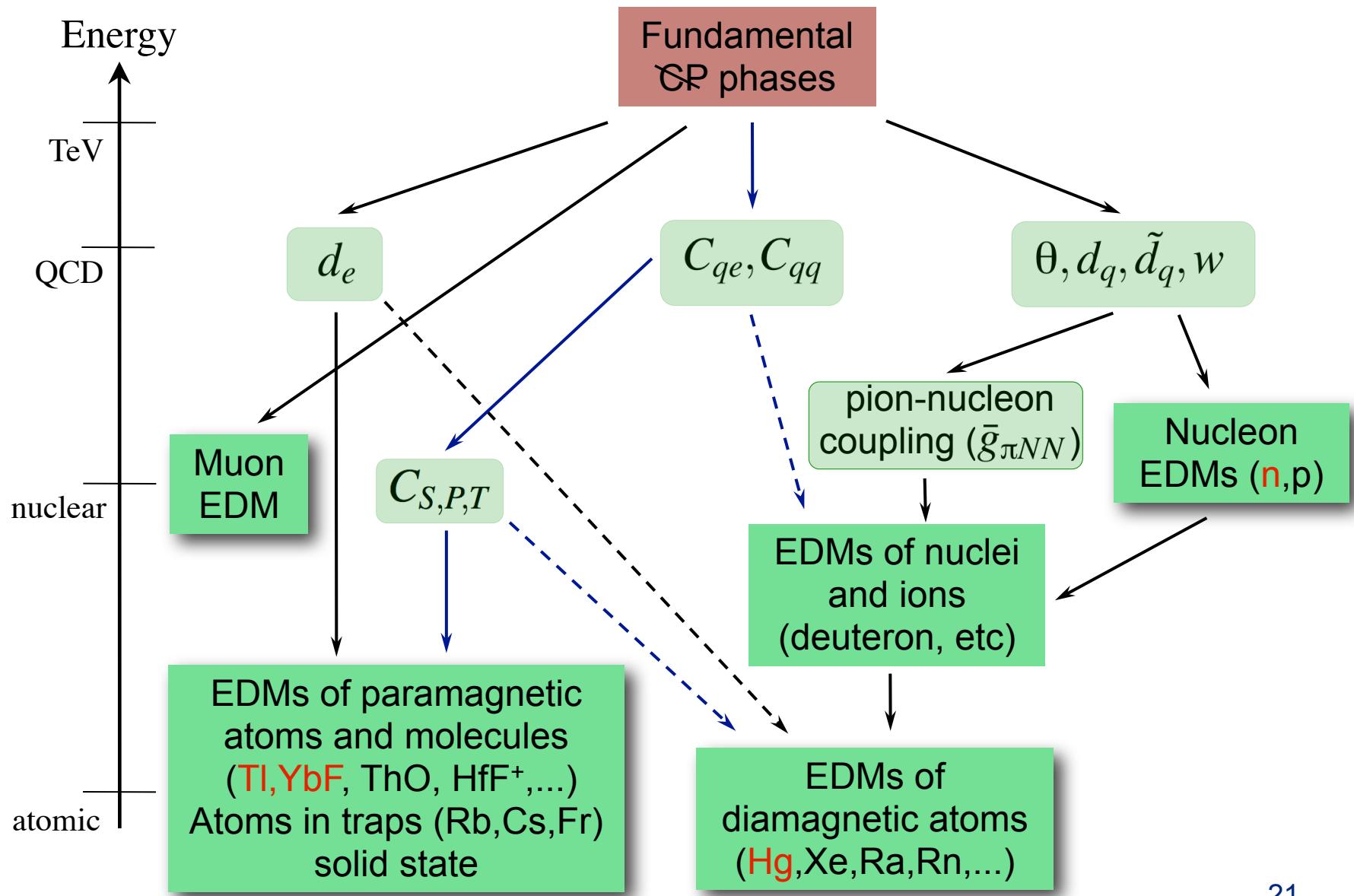


[Dzuba et al. '02; Flambaum et al. '86; Dmitriev & Senkov '03; de Jesus & Engel '05; Ban et al '10]

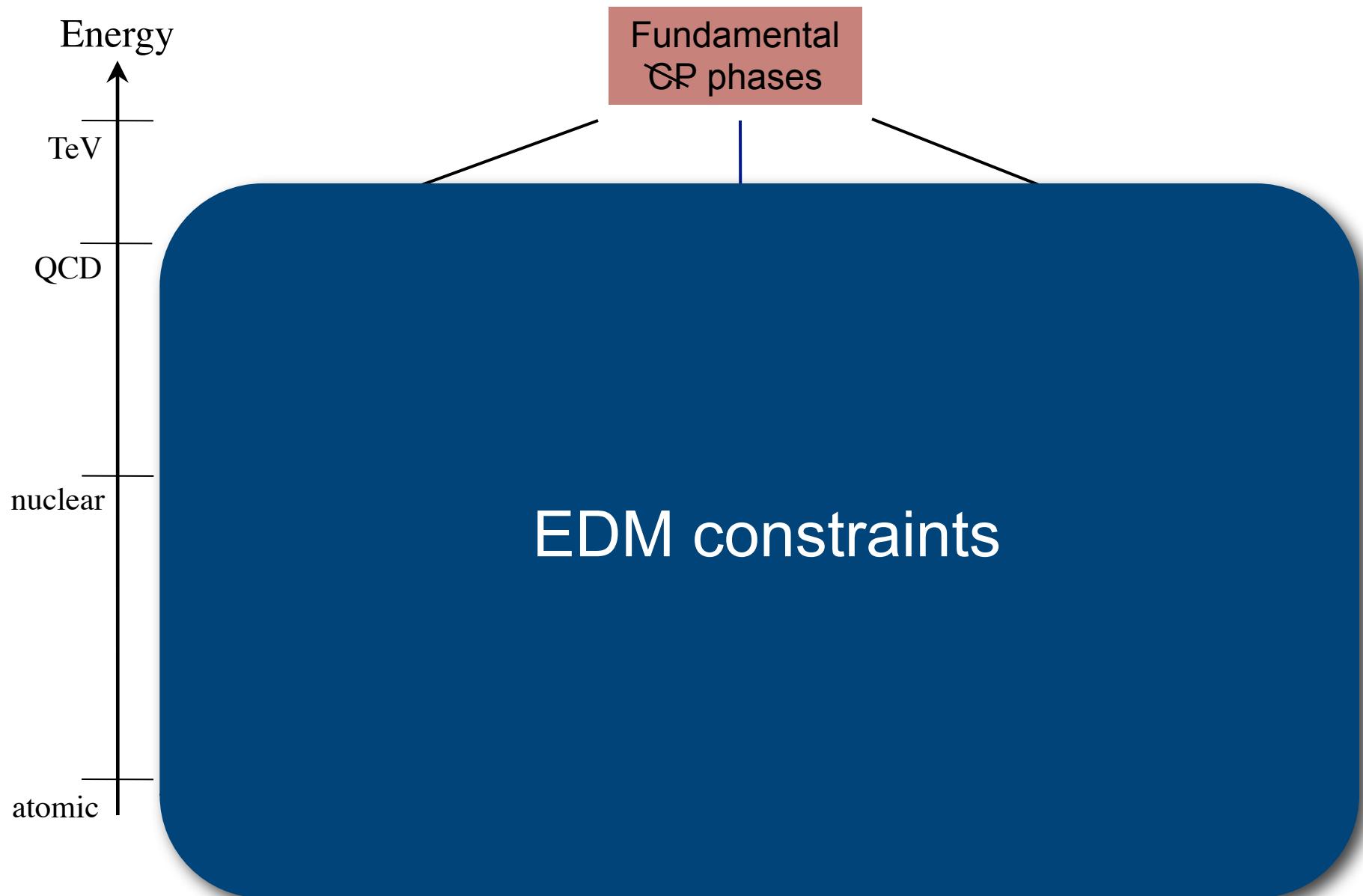
NB: Precision limited by: (i) $S(\bar{g}_{\pi NN}^{(0)}, \bar{g}_{\pi NN}^{(1)}, \bar{g}_{\pi NN}^{(2)})$ [Ban et al '10]

(ii) $\bar{g}_{\pi NN}^{(1)}(\tilde{d}_q) \sim (1-6)(\tilde{d}_u - \tilde{d}_d) + O(\tilde{d}_u + \tilde{d}_d, \tilde{d}_s, w)$ [Pospelov '01]

EFT hierarchy



Constraints on CP-violation



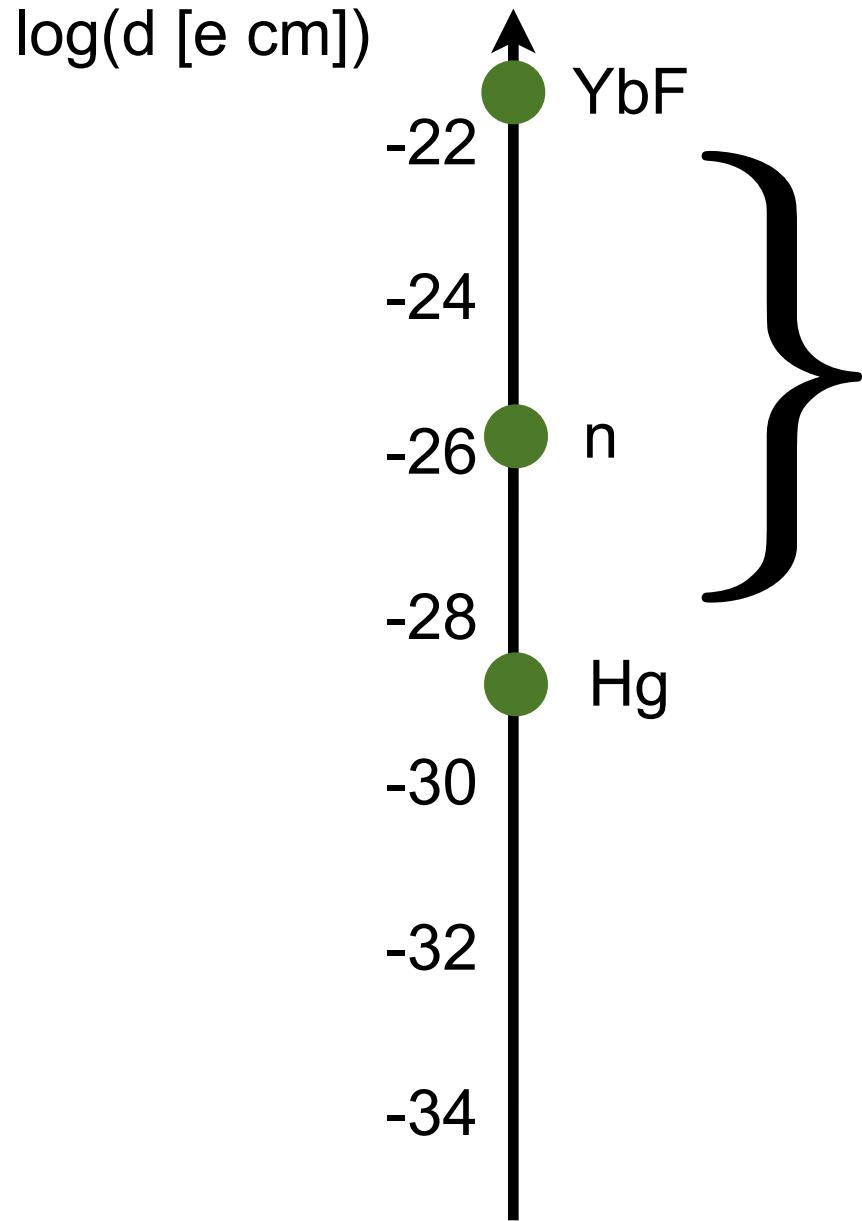
Resulting Bounds on fermion EDMs & CEDMs

YbF EDM [$\pm 20\%$]	$ d_e + \mathcal{O}(C_{eq}) < 1.1 \times 10^{-27} e \text{ cm}$
TI EDM [$\pm 20\%$]	$\left d_e + e(26 \text{ MeV})^2 \left(3 \frac{C_{ed}}{m_d} + 11 \frac{C_{es}}{m_s} + 5 \frac{C_{eb}}{m_b} \right) \right < 1.6 \times 10^{-27} e \text{ cm}$
Neutron EDM [$\pm 50\%?$]	$ e(\tilde{d}_d + 0.5\tilde{d}_u) + 1.3(d_d - 0.25d_u) + \mathcal{O}(\tilde{d}_s, w, C_{qq}) < 2 \times 10^{-26} e \text{ cm}$
Hg EDM [$\pm O(\text{few})?$]	$e \tilde{d}_d - \tilde{d}_u + \mathcal{O}(d_e, \tilde{d}_s, C_{qq}, C_{qe}) < 6 \times 10^{-27} e \text{ cm}$

Generic scaling: $d_f \sim (\text{couplings}) \times \frac{m_f}{\Lambda_{CP}^2}$

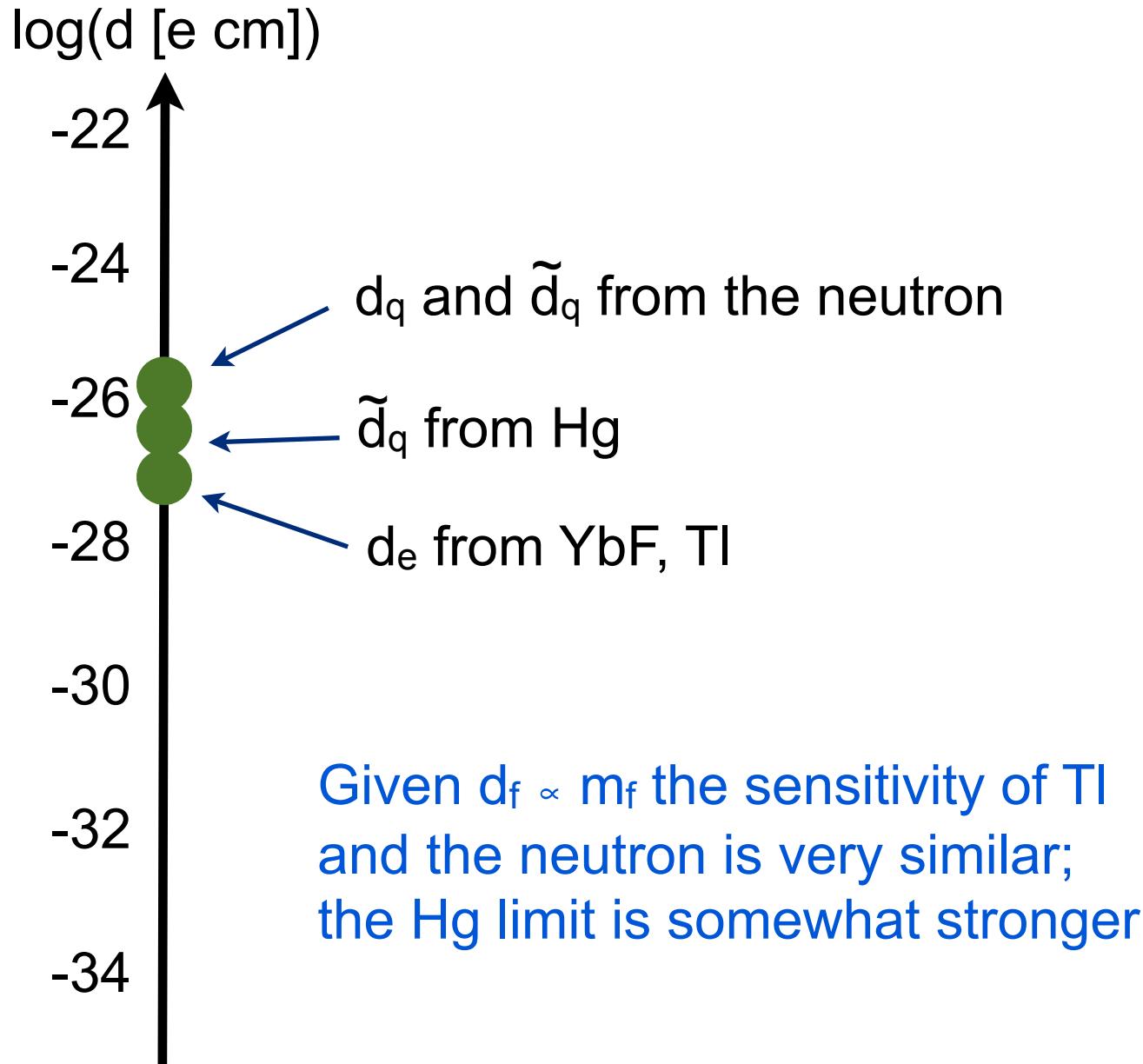
See also recent compilation of limits: [Engel, Ramsey-Musolf, van Kolck '13]

Summary of the bounds

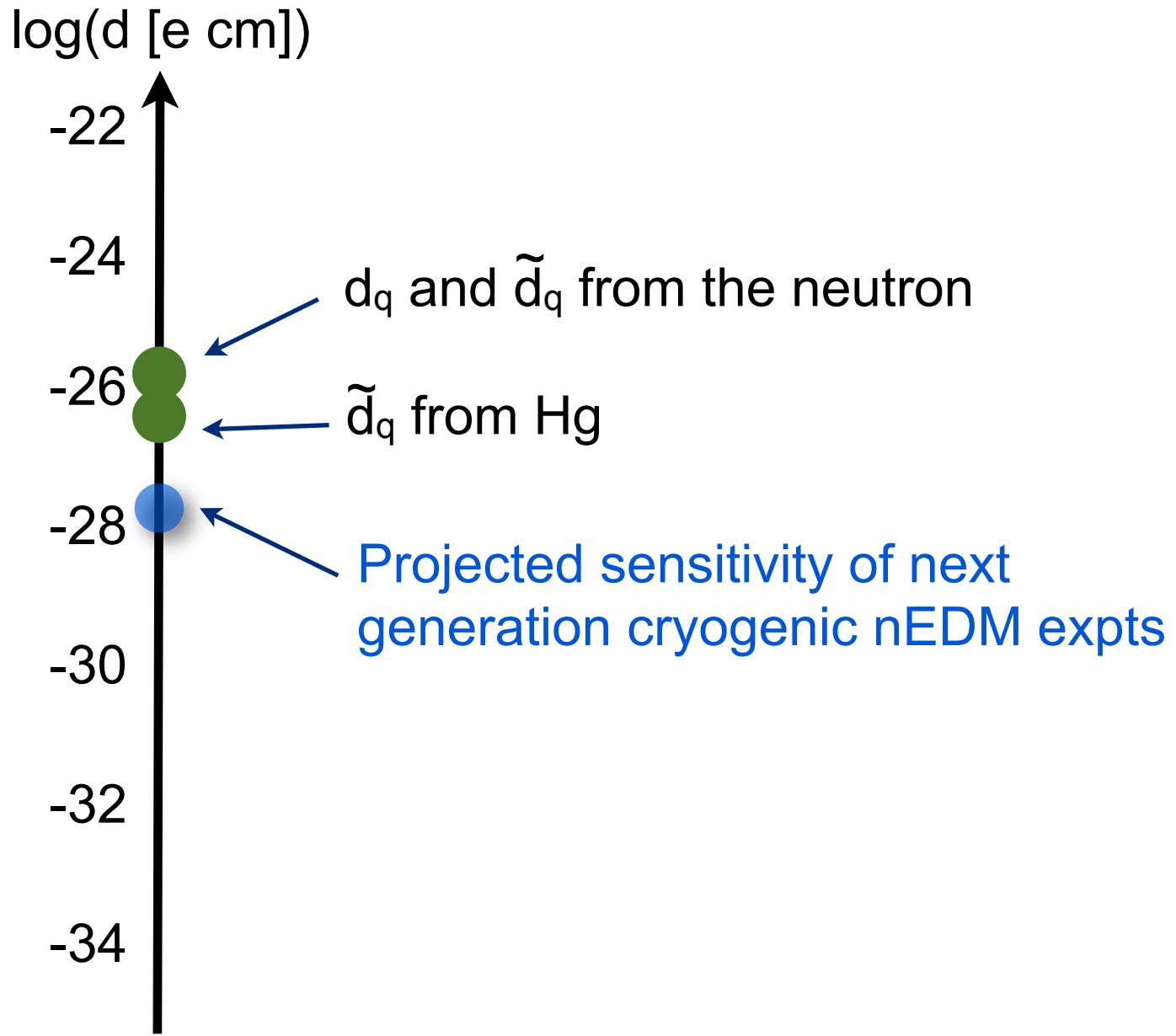


Difference of more than 6 orders of magnitude, but the sensitivity to many underlying CP-odd sources is similar...

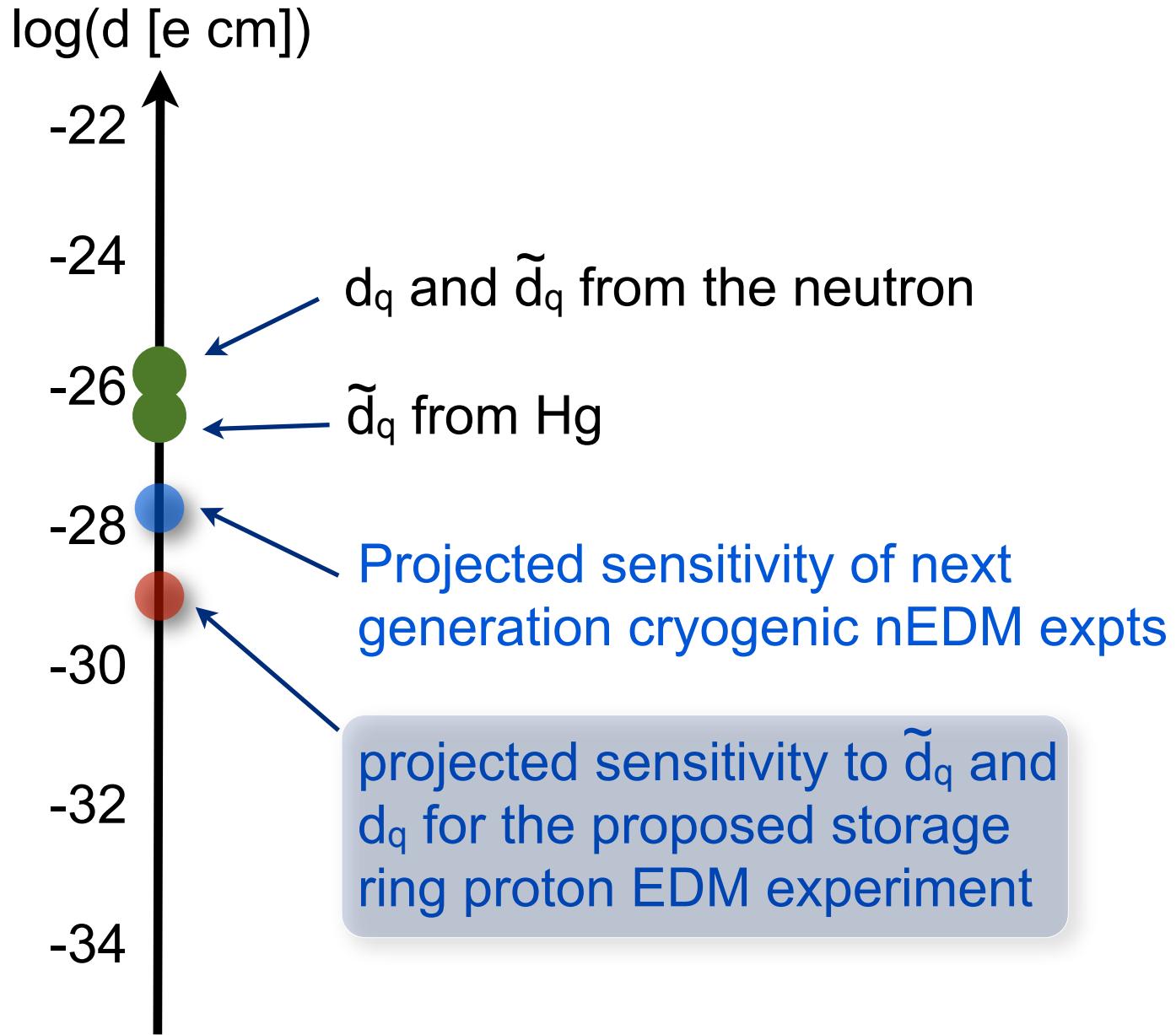
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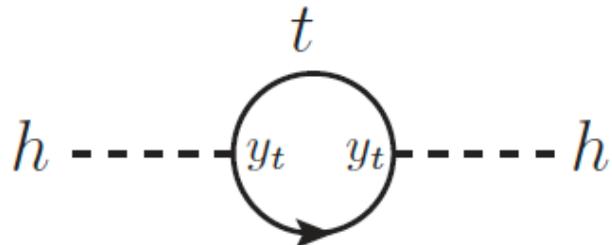


Summary of the bounds



LHC-era tests of CP-violating new physics

Expectation of new EW-scale physics is (or was) primarily associated with stabilising the Higgs sector...

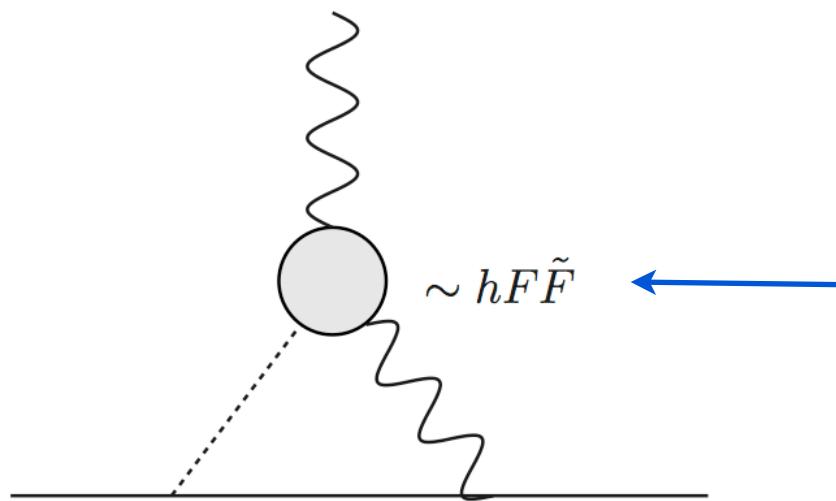


- This predominantly suggests new physics coupling strongly to the Higgs, 3rd generation, ...
 ⇒ EDMs at 2-loops

E.g. - CP-odd Higgs couplings

- Hints in 2012 that $\text{Br}(h \rightarrow \gamma\gamma) > \text{Br}_{\text{SM}}$ (still present in ATLAS data)
- EDMs significantly constrain any CP-odd contribution to $h \rightarrow \gamma\gamma$

$$\Delta\mathcal{L} = \frac{1}{e^2 \tilde{\Lambda}^2} H^\dagger H \left(a_h g_1^2 B_{\mu\nu} \tilde{B}^{\mu\nu} + b_h g_2^2 W_{\mu\nu} \tilde{W}^{\mu\nu} \right) \rightarrow \frac{\tilde{c}_h v}{\tilde{\Lambda}^2} h F_{\mu\nu} \tilde{F}^{\mu\nu} + \dots$$



This interaction corrects the Higgs width, but also generates 2-loop EDMs!

$$\frac{\Gamma_{\gamma\gamma}}{\Gamma_{\gamma\gamma}^{\text{SM}}} \simeq 1 + \left| \tilde{c}_h \frac{v^2}{\tilde{\Lambda}^2} \frac{8\pi}{\alpha A_{\text{SM}}} \right|^2$$

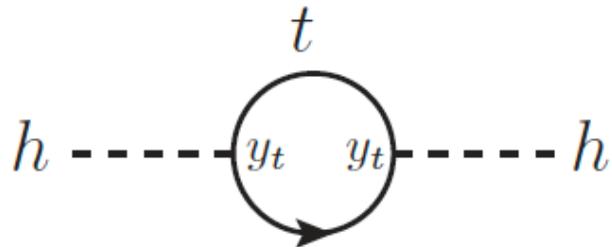
Current limit on d_e limits the shift of $\text{Br}(h \rightarrow \gamma\gamma)/\text{Br}_{\text{SM}}$ to $O(10^{-4})$!

[McKeen, Pospelov & AR '12]

[Harnik et al '12; Fan & Reece '13]

LHC-era tests of CP-violating new physics

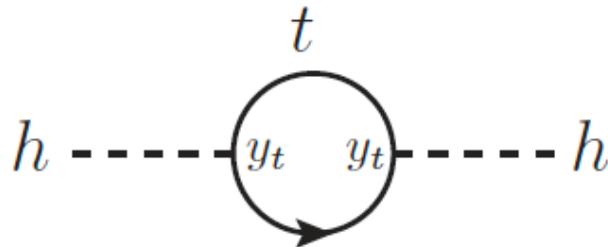
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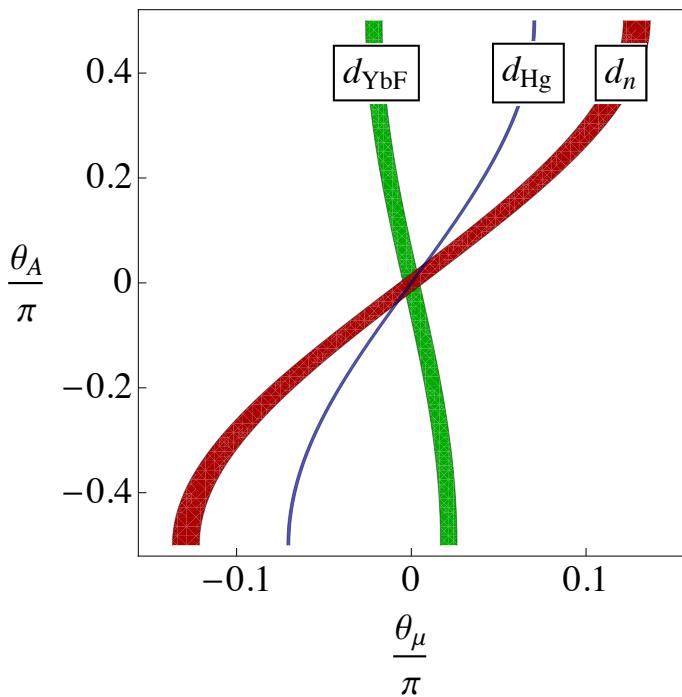


- This predominantly suggests new physics coupling strongly to the Higgs, 3rd generation, ...
 ⇒ EDMs at 2-loops
- SUSY provides new physics with strong coupling to 1st generation
 ⇒ EDMs at 1-loop! → SUSY CP problem!

E.g. - SUSY CP Problem (given LHC constraints)

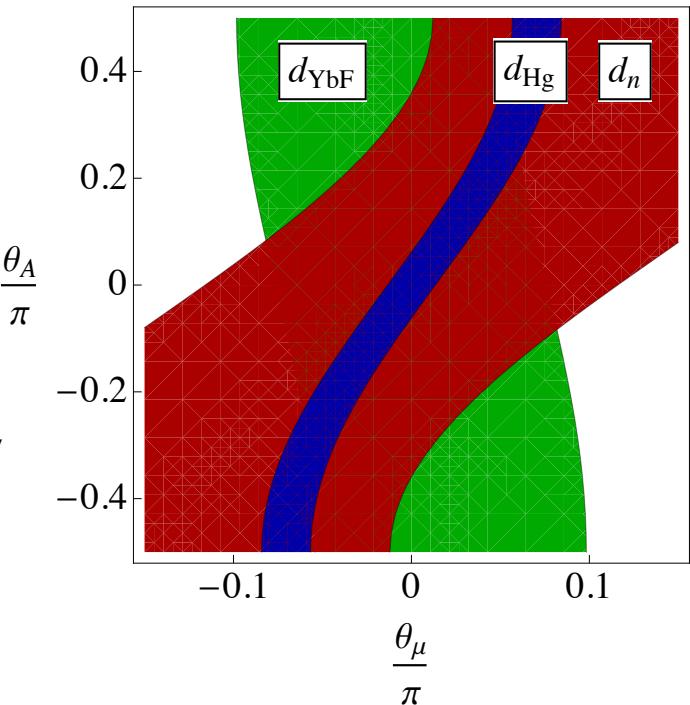
(pre-LHC)

$$M_{susy} = 500 \text{ GeV}$$



(now)

$$M_{susy} = 2 \text{ TeV}$$

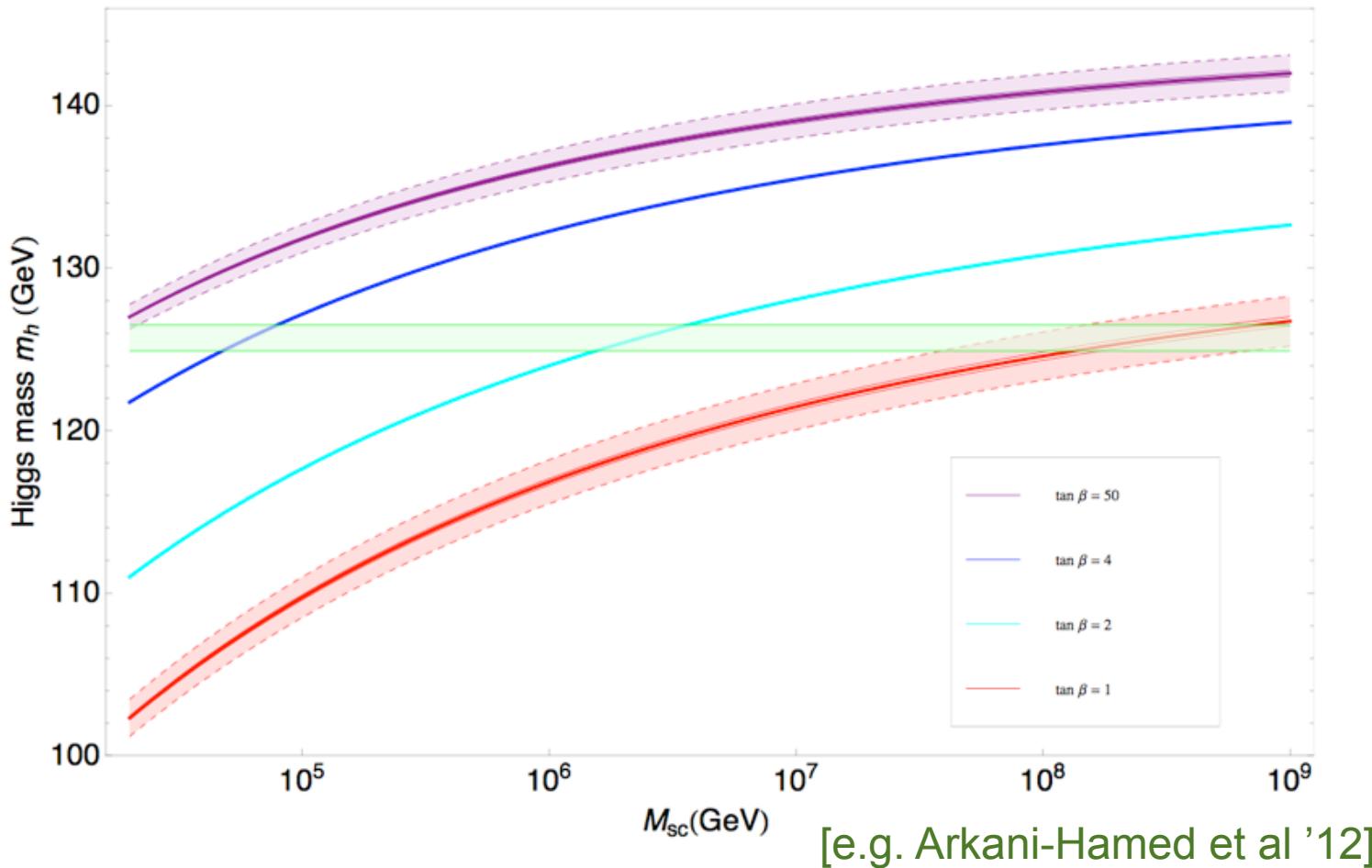


1st gen squarks
excluded by direct
searches at ~ 1 TeV

EDMs have for many years required (tuned) $O(10^{-3})$ CP-odd phases for generic weak-scale SUSY. The LHC appears to have resolved this by pushing mass limits on 1st generation sfermions above a TeV

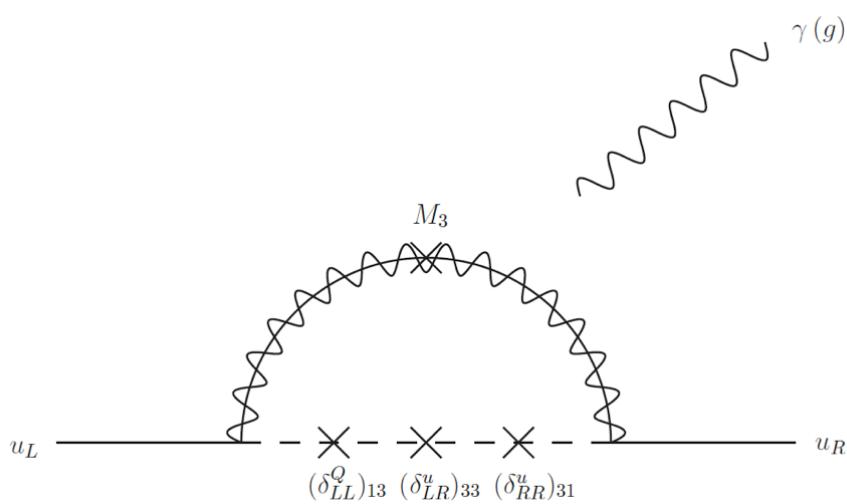
E.g. - PeV-scale SUSY sensitivity

- Within minimal SUSY, $m_h \gg m_Z$ points to PeV-scale s-partners
(\Rightarrow no soln to hierarchy problem)

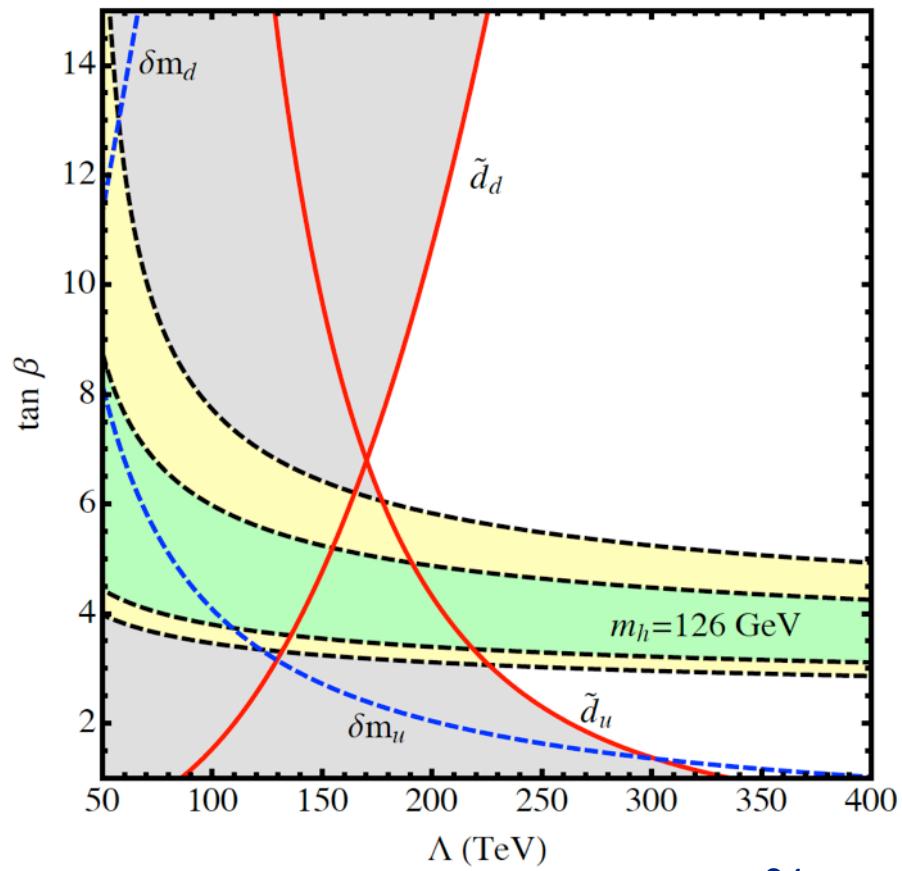


E.g. - PeV-scale SUSY sensitivity

- Within minimal SUSY, $m_h \gg m_Z$ points to PeV-scale s-partners
- The PeV scale allows a generic flavour structure and, with TeV gauginos, hadronic EDMs are one of the few observables able to probe this scale via log-enhanced quark CEDMs



[McKeen, Pospelov & AR '13]



Concluding Remarks

EDMs form an important class of flavour-diagonal CP-odd observables, testing/constraining new physics (motivated by the need for baryogenesis)

- Disentangling multiple CP-odd operators at 1 GeV requires multiple observables
- Storage rings are a novel new technology, which suggest a route around the Schiff screening that “hides” nuclear EDMs



recent applications

- Useful interplay between EDM constraints and precision tests of CP-odd Higgs couplings
- The SUSY CP problem hinted at by (1-loop) EDMs for more than 20 years has been “confirmed” by the LHC, with no s-partners seen near the weak scale (thus far).

Extra slides

Neutron/Proton EDM Precision

- Issues affecting precision, and tests:
 - numerical coefficients are consistent with NDA, NQM (for d_q), and the chiral log (for θ)
 - another test for $d_n(d_q)$ via (LQCD) nucleon tensor charge
[e.g. Falk et al '99]
$$\langle N | \frac{1}{2} d_q \bar{q} \tilde{F} \sigma q | N \rangle = \frac{1}{2} d_q \tilde{F}^{\mu\nu} \langle N | \sigma_{\mu\nu} | N \rangle = \frac{1}{2} g_T^q d_q \bar{N} \tilde{F} \sigma N$$
$$\implies d_n(d_q) = g_T^d(1 \text{ GeV}) d_d + g_T^u(1 \text{ GeV}) d_u \sim 1.1 d_d - 0.25 d_u$$

- sum-rules fixes ($d_n \sim \langle qq \rangle / \lambda^2$), so the normalization of the coupling matters

$$\lambda \sim 0.025 \text{ GeV}^3$$

from analysis of CP-even sum rules
for m_n , sigma term, etc (or lattice
result for tensor charge above)

[Pospelov & AR '99, '00]

$$\lambda \sim 0.044 \pm 0.01 \text{ GeV}^3$$

from LQCD [Aoki et al '08] run down
from 2 GeV, *BUT* $\langle qq \rangle$ is also
larger with LQCD values for m_q

[Hisano et al '12, Fuyuto et al '12]

- higher order dependence on s-quark EDM?