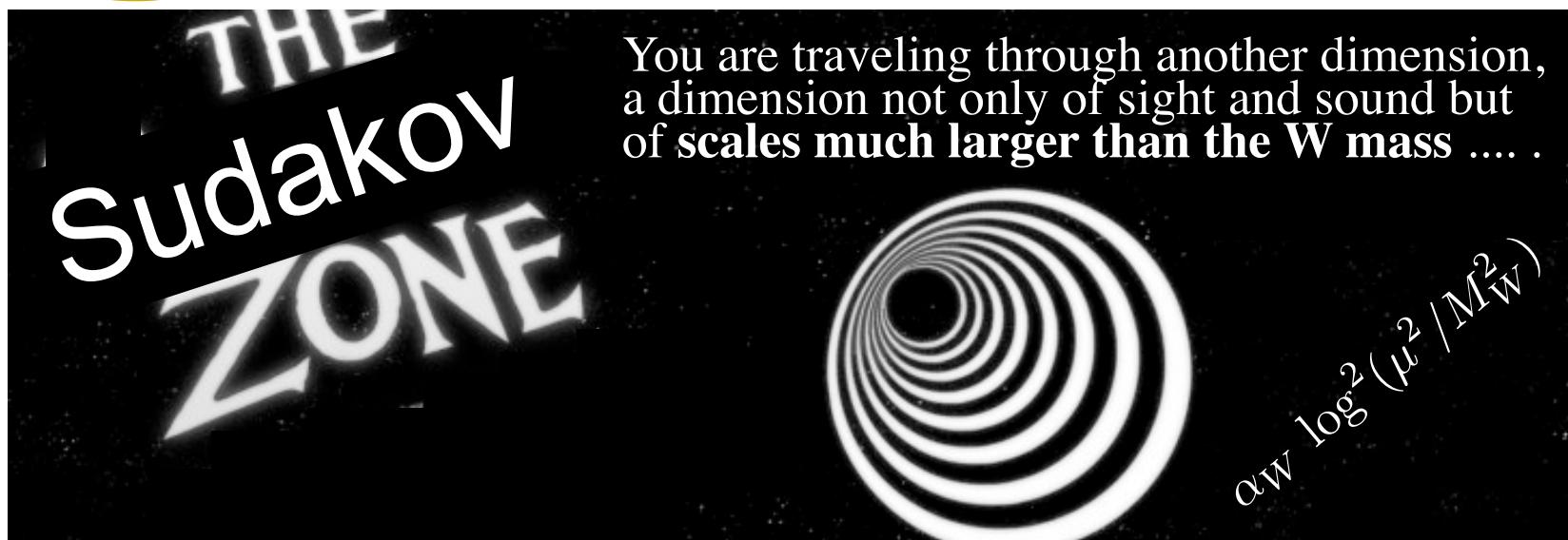




Electroweak Corrections at High Energies

Kalanand Mishra, *Fermilab*



Snowmass: Seattle Energy Frontier Workshop, July 2, 2013

Outline

- Introduction
 - What are EWK corrections, when/where they matter
- A survey of the most abundant processes at LHC for sensitivity to EWK corrections: at 8 TeV, 14 TeV, 33 TeV, and 100 TeV collider energies
 - inclusive jet, dijets
 - inclusive W/Z
 - W/Z+jets
 - Diboson (WW)
- Summary table

Definition

If experimental and PDF errors are < EWK corrections, then I'll call the measurement as being in the Sudakov zone. Typically need a few dozen events above 1 TeV to be in this zone.

When electroweak corrections are too strong

- EWK corrections become important in certain kinematic regimes
- At the \sim TeV scale, they significantly modify production cross sections through Sudakov logs of order $\alpha_W \log^2(\mu^2/M_W^2)$

$$= 25 * \alpha_W \quad @\mu=1\text{TeV}$$

This typically exceeds $\sim 10\%$ of LO when μ is of order 1 TeV

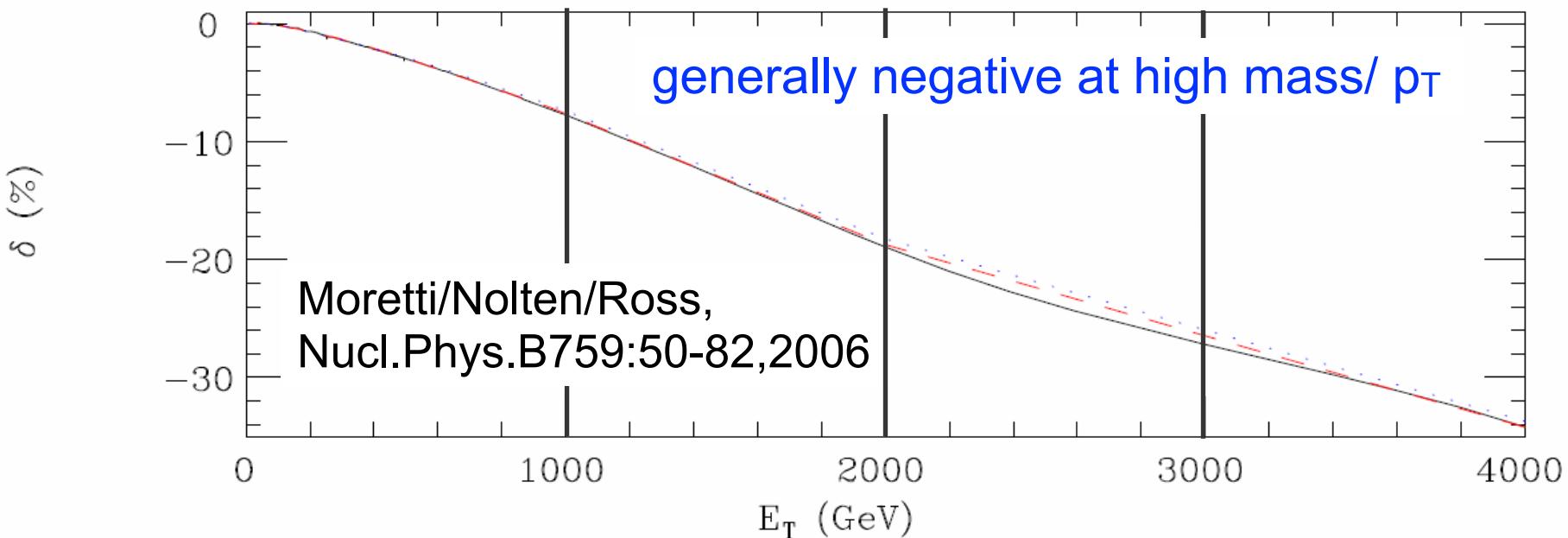
- We are now encountering LHC data where experimental precision at the TeV scale is of order 10%

Important not only for nominally electroweak processes, but for nominally QCD ones as well!

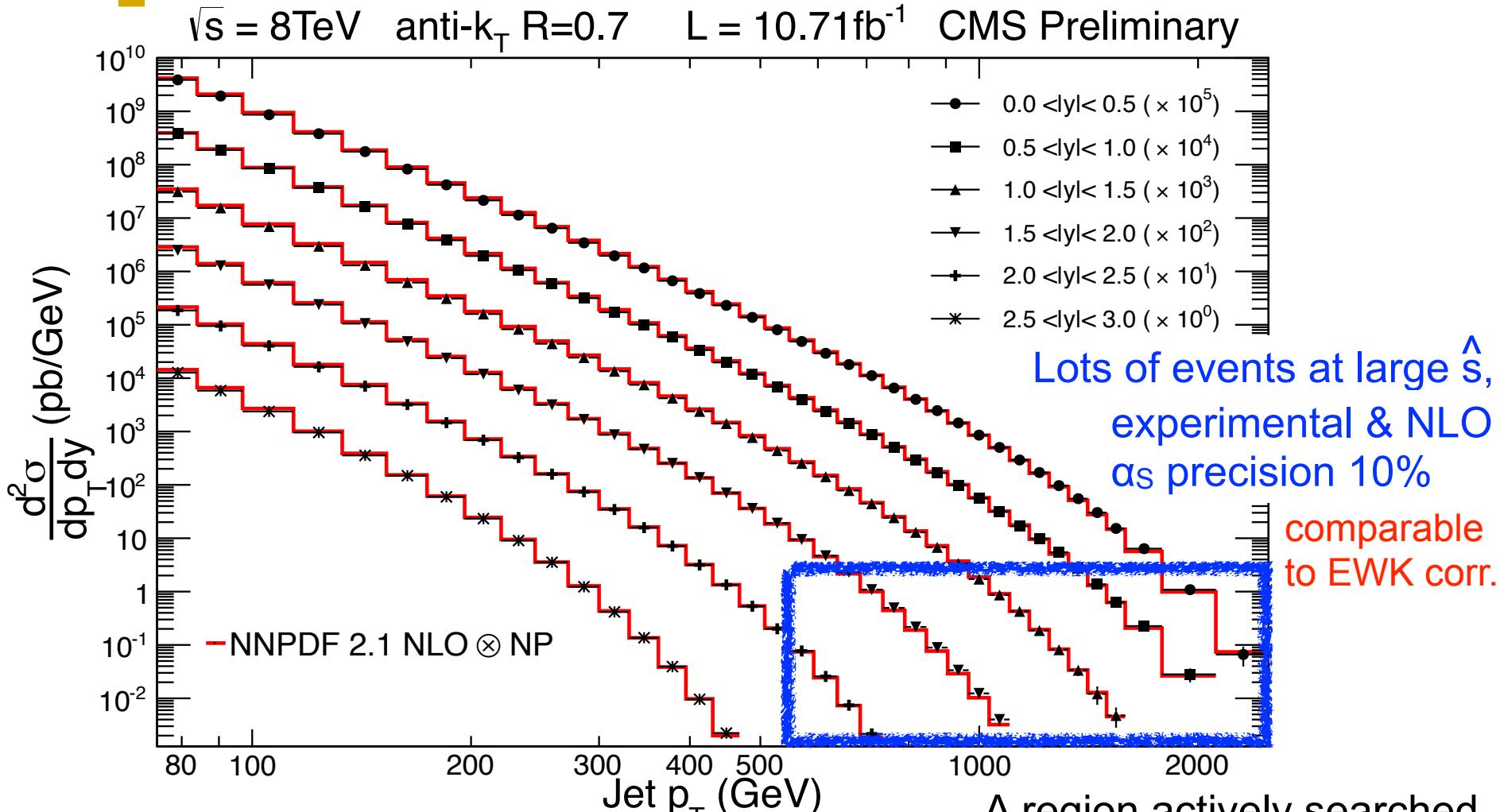
How big an effect we are talking about ?

Inclusive jet production at LHC

For inclusive jet production at 14 TeV, negative virtual contribution dominates over real emission at ~ 1 TeV and **cross section decreases from LO result by $\sim 20\%$ at 2TeV**



Are EWK corrections too strong for jets?

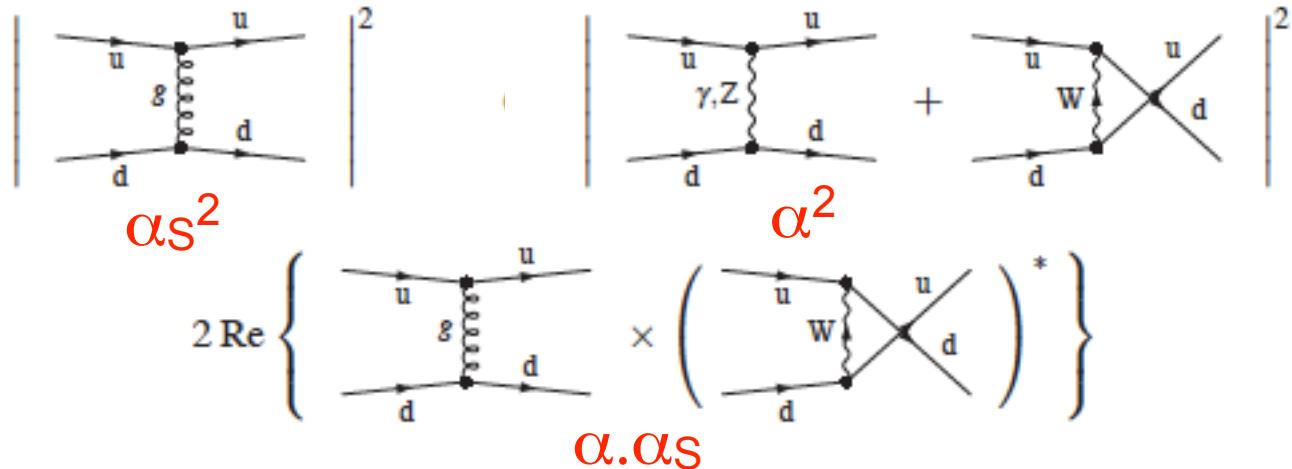


<https://cdsweb.cern.ch/record/1547589>

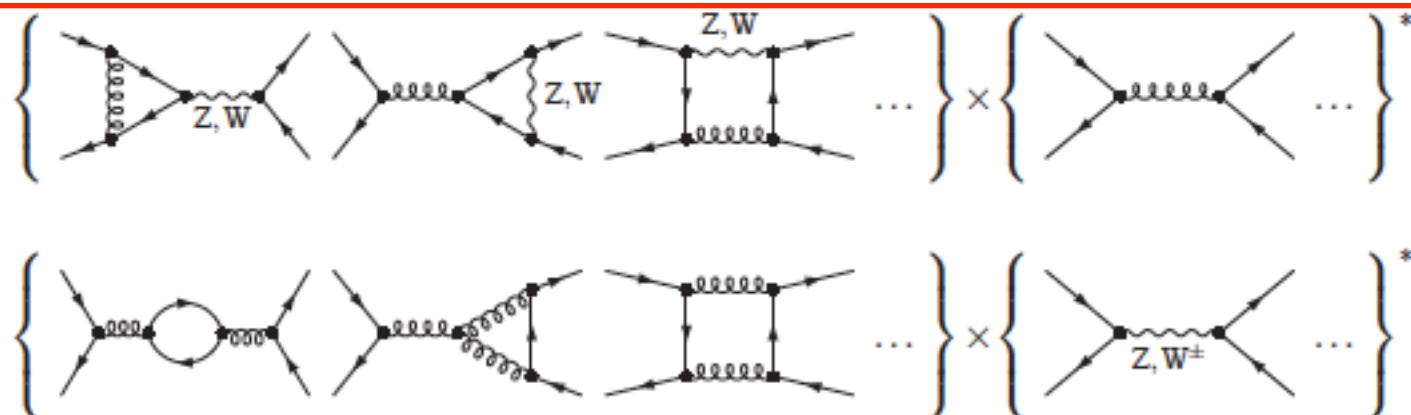
A classic example: EWK corr for dijet production

arXiv:1210.0438,
arXiv:1306.6298

Tree-level diagrams
(+ve contribution)

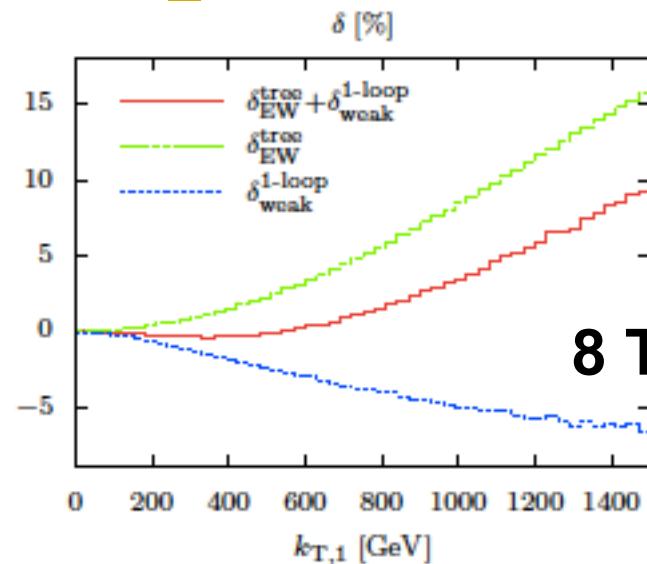


Loop diagrams
(-ve contribution,
virtual, Sudakov)

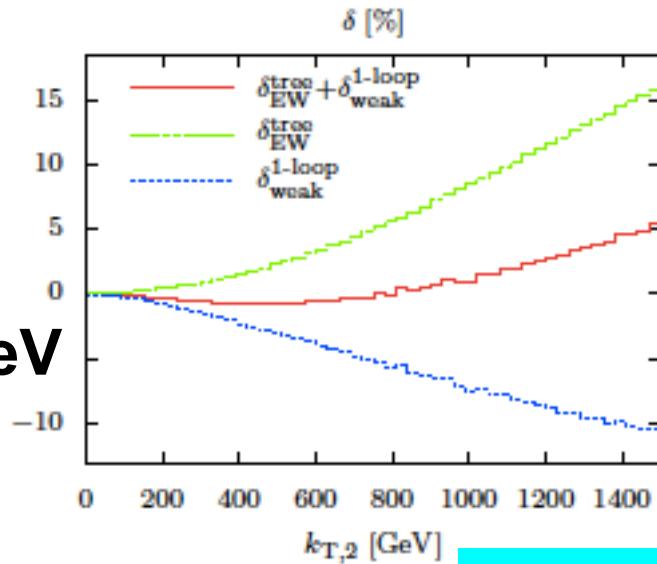


Although negligible for integrated xsec, Sudakov corrections can reach 10–20% for TeV jets in interesting kin regimes

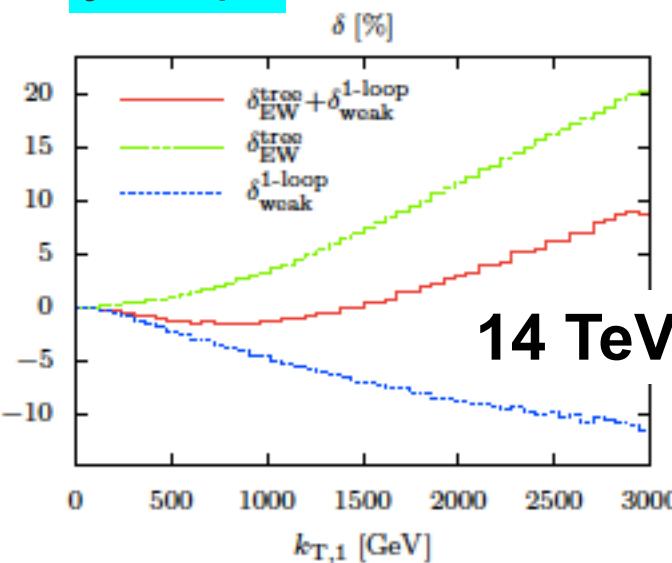
Dijet expectation at 8 TeV and 14 TeV



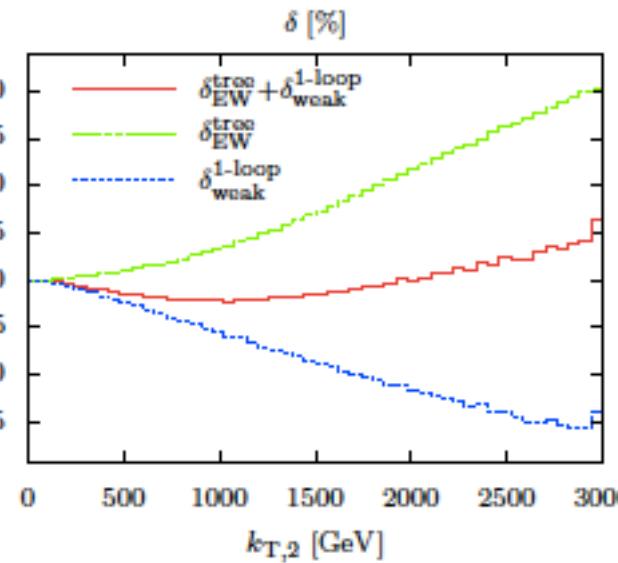
jet 1 p_T



jet 2 p_T



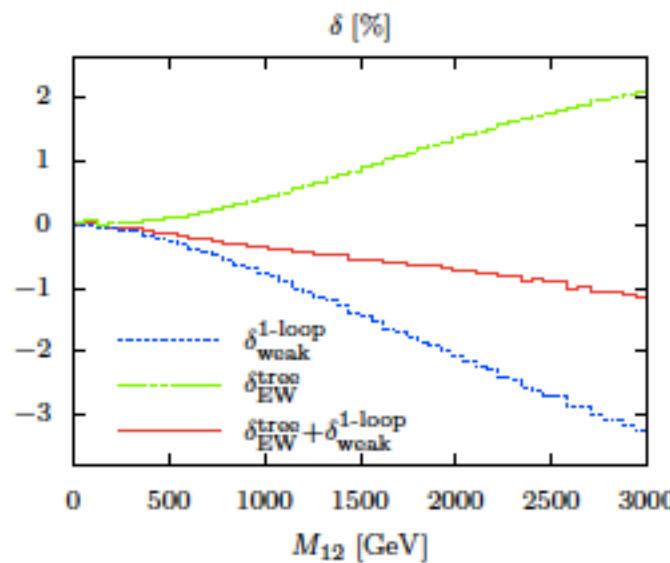
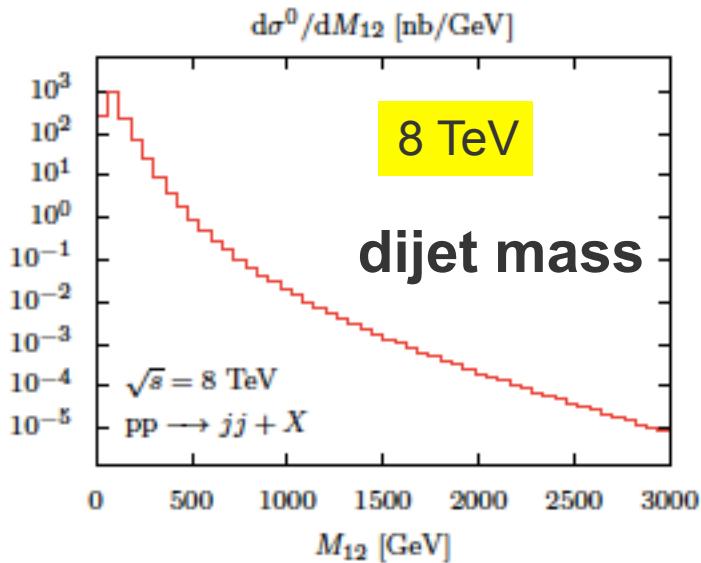
14 TeV



arXiv:1210.0438

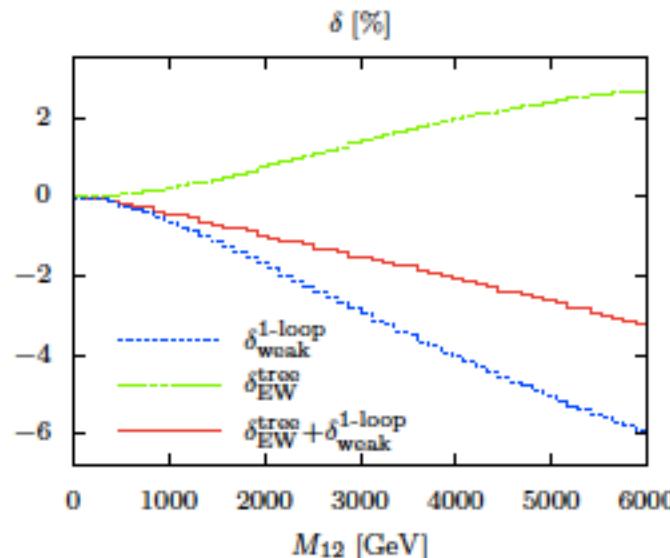
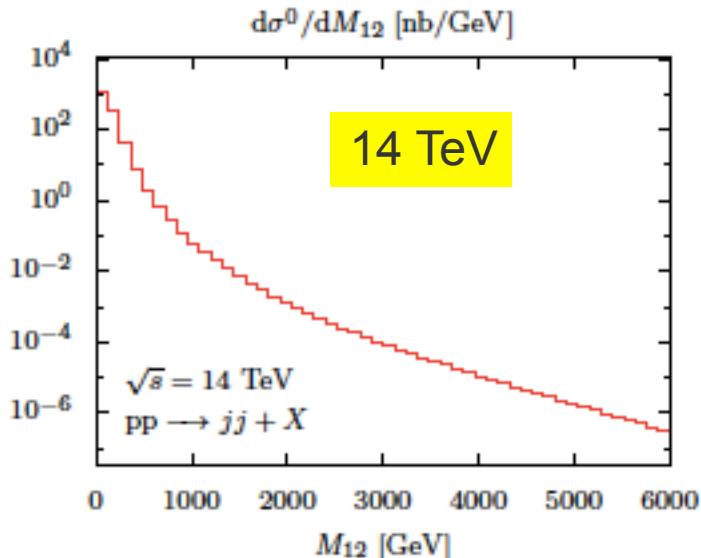
Sudakov virtual correction is large, although the net EWK correction is small

Dijet expectation at 8 TeV and 14 TeV



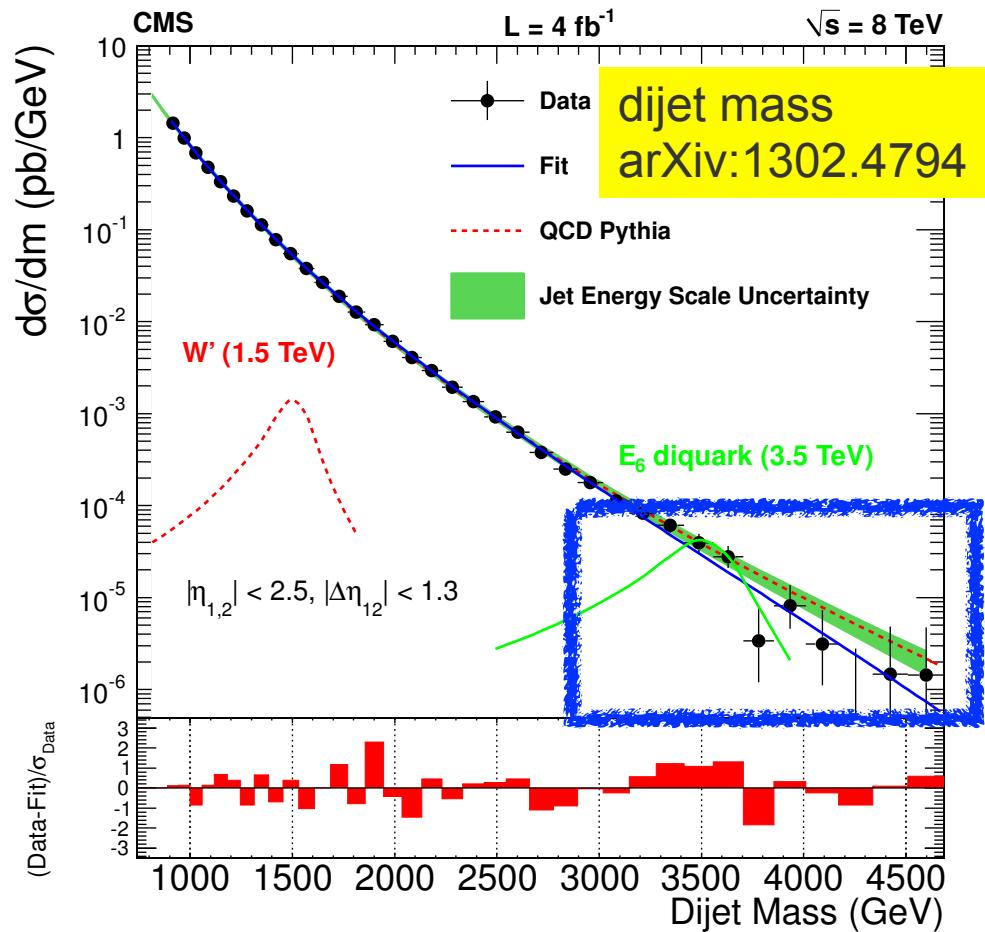
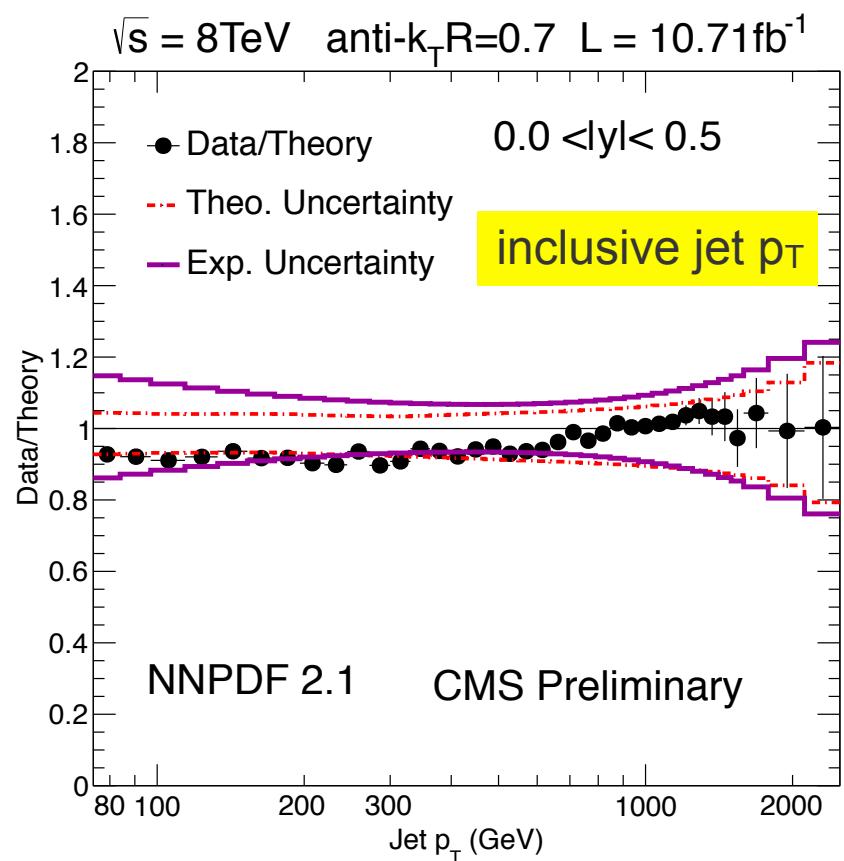
Typically we are interested in low y^* , where the correction can be even larger.

arXiv:1210.0438



See the paper for plots of individual y^* bins

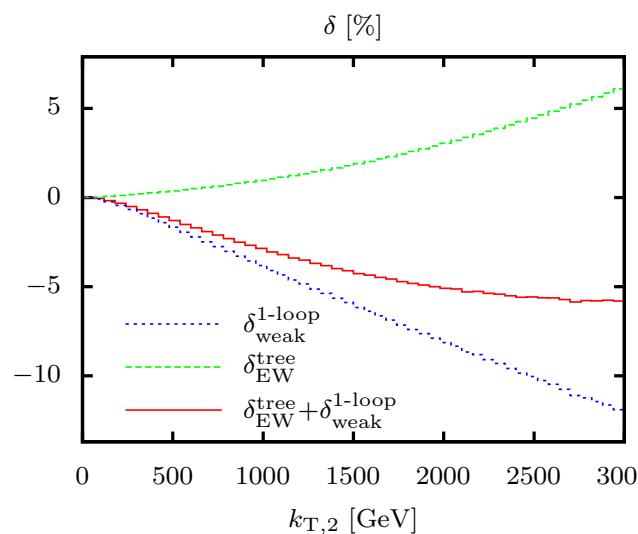
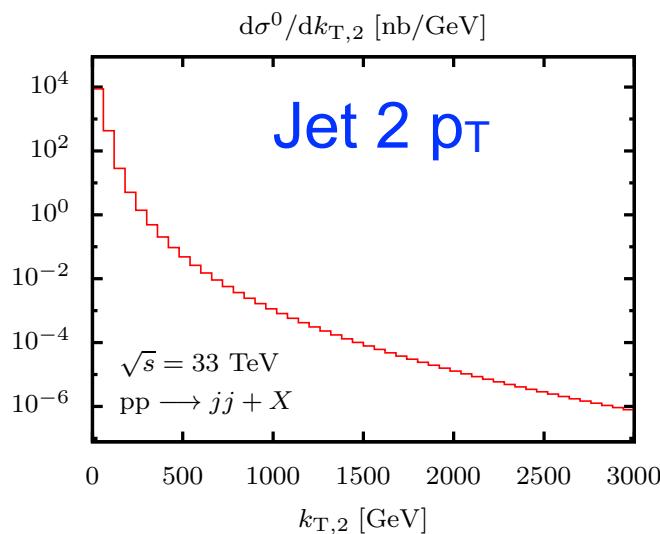
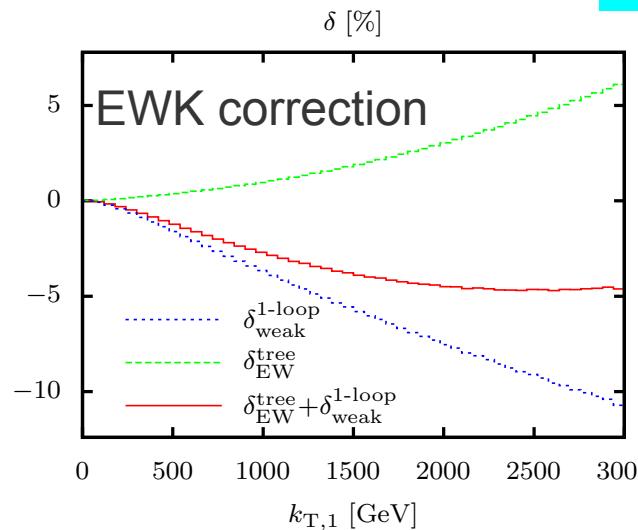
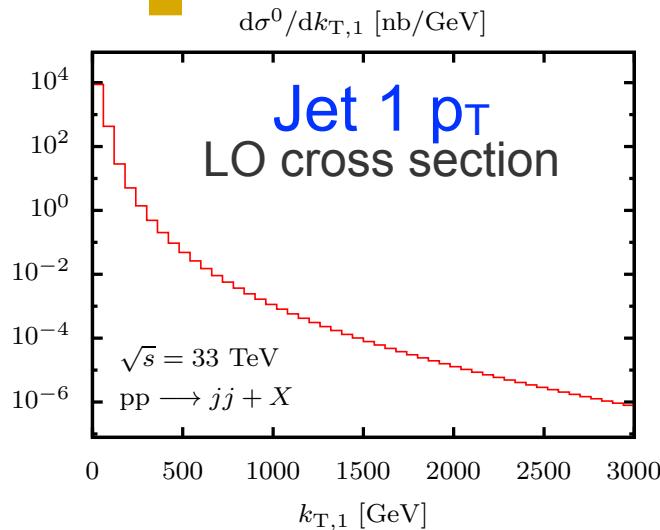
Too strong for jets? At today's precision, yes!



- Theory and experimental error for $\sim 2(4)$ TeV jets(dijets) is $\sim 10\%$
 - Comparable to the size of EWK corrections

So, what we expect at 33 TeV?

Provided by Alexander Huss & Stefan Dittmaier

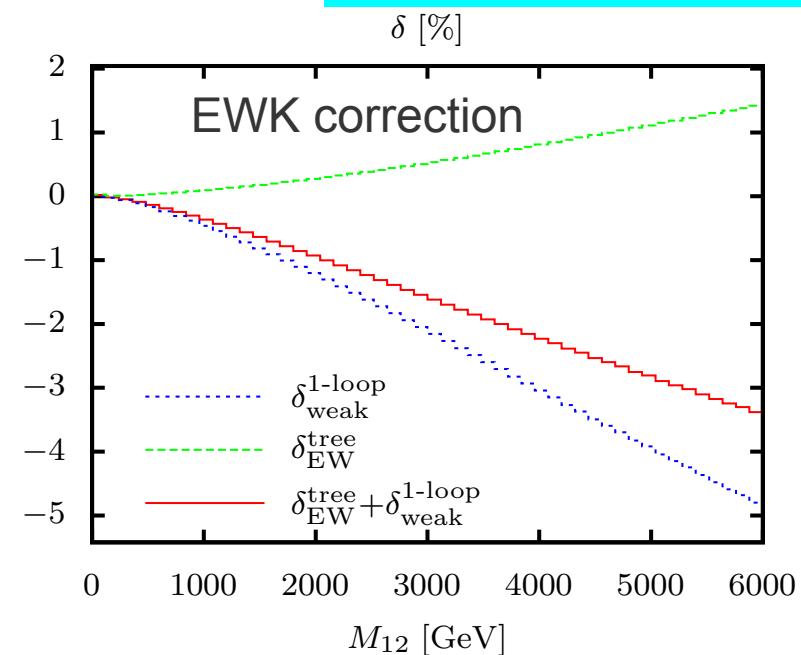
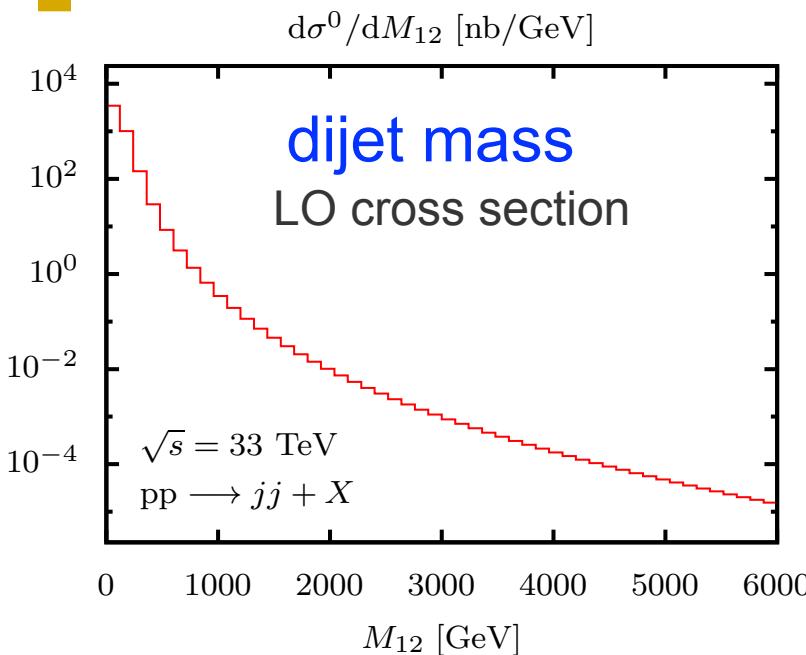


- Weak 1-loop corr not so dependent on collider energy in the case of p_T distribution for which the tails are dominated by the Sudakov regime.
- Tree-level corr are dependent on collider energy due to 4-quark subprocesses that gain in importance at higher E.

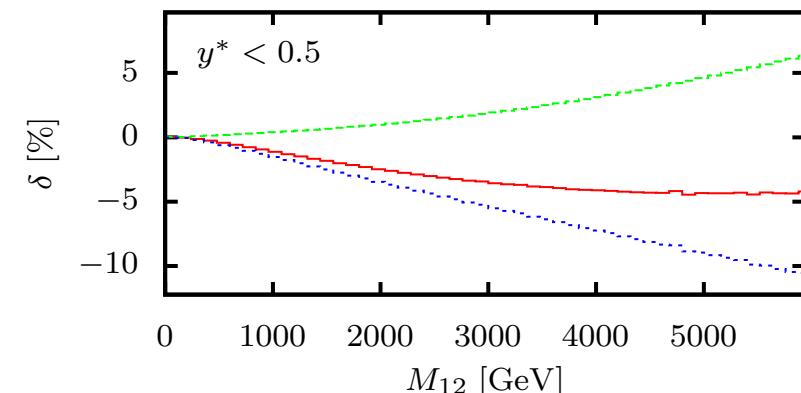
[arXiv:1306.6298](https://arxiv.org/abs/1306.6298)

So, what we expect at 33 TeV?

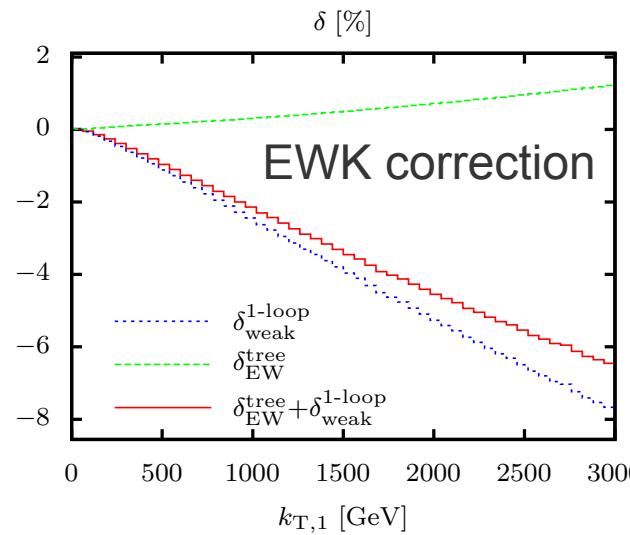
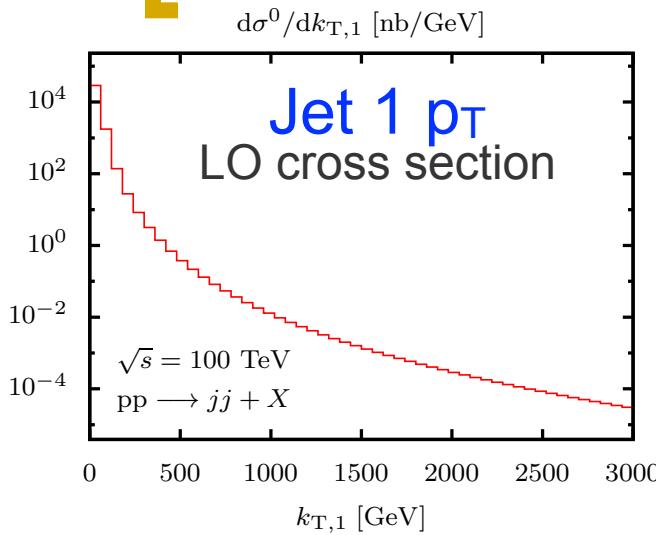
Provided by Alexander Huss & Stefan Dittmaier



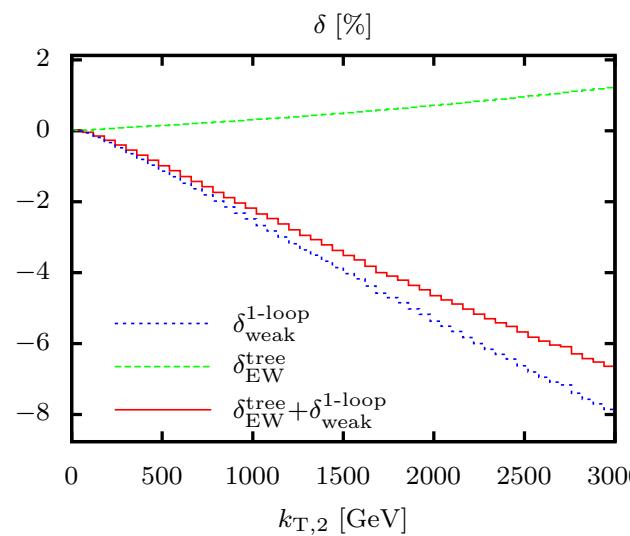
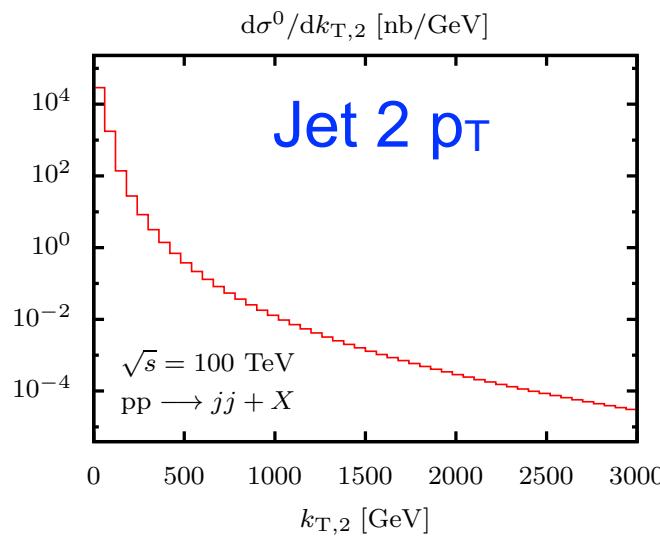
For M_{jj} the energy dependence is stronger, since not probing pure Sudakov regime (not solely described by the universal Sudakov logarithms). However, the lowest y^* bins (in Sudakov regime) are less dependent on the energy & show a similar behavior to the p_T distributions.



At 100 TeV



Provided by
Alexander Huss &
Stefan Dittmaier

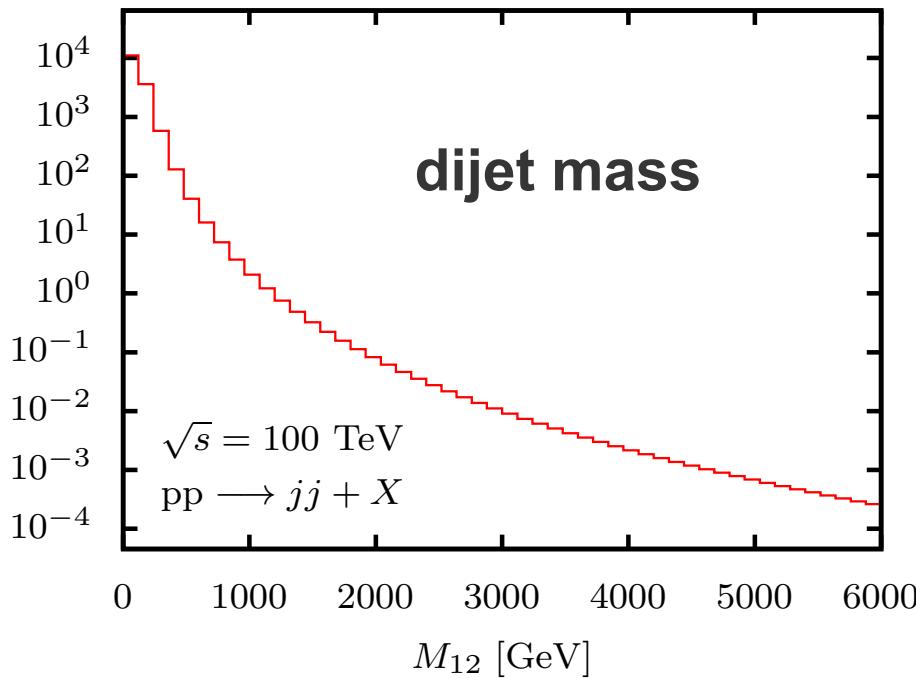


Very similar to
the 33 TeV result

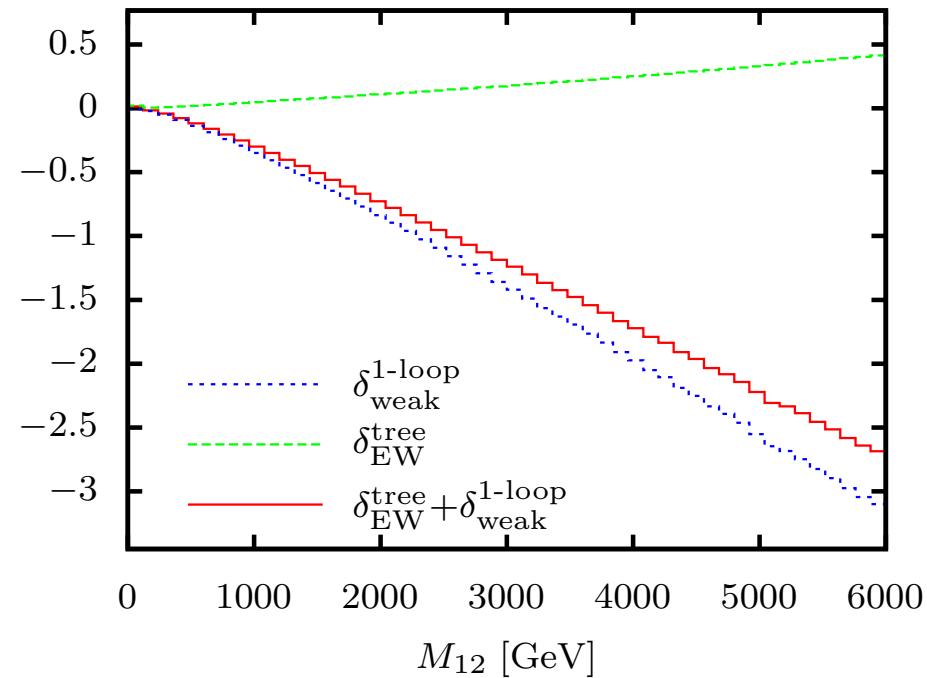
At 100 TeV

Provided by Alexander
Huss & Stefan Dittmaier

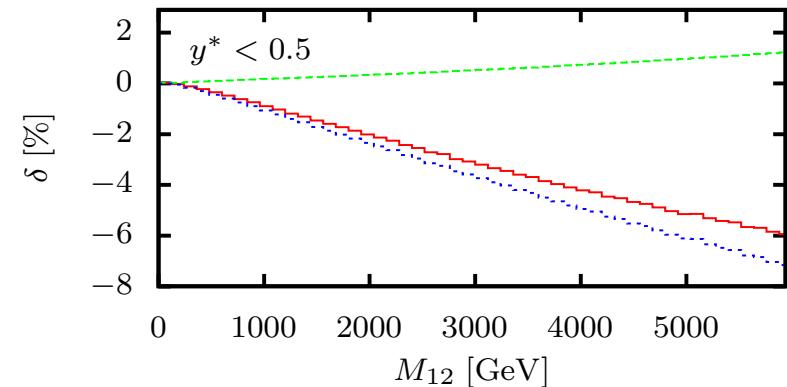
$d\sigma^0/dM_{12}$ [nb/GeV]



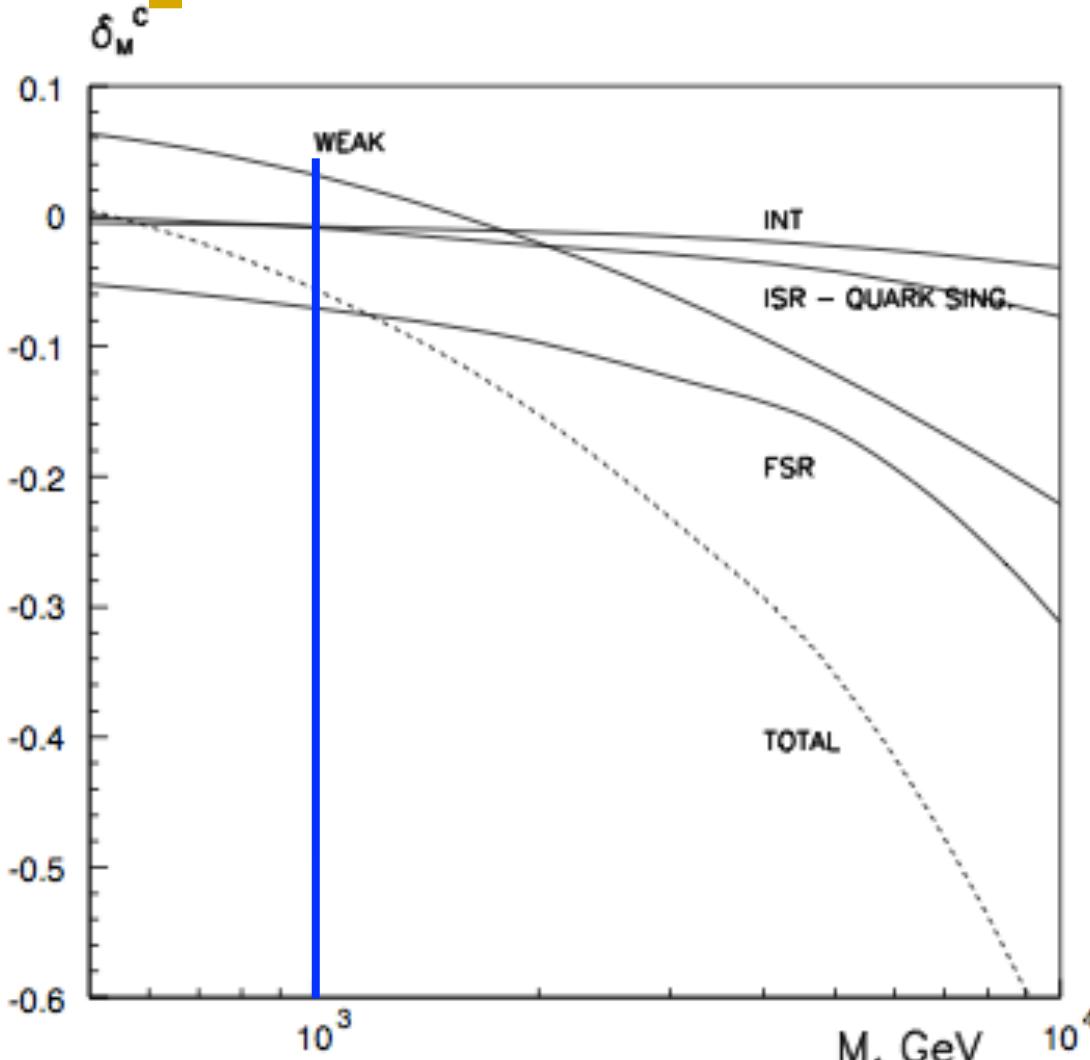
δ [%]



Very similar to the 33 TeV result



What about electroweak processes

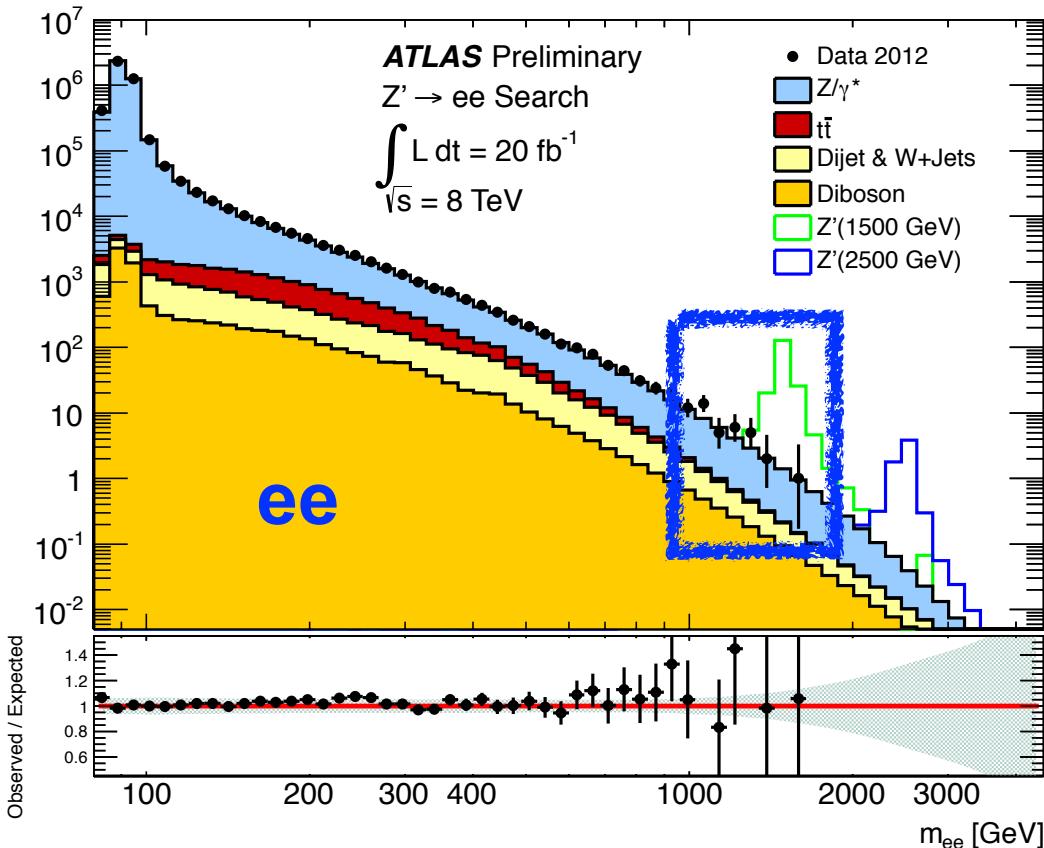


Almost everything we could measure in 2010 (jets, photons, W, Z) has entered the “Sudakov zone” in 2012.

As in the case of inclusive jet, the QCD NLO correction (+ve) compensates for EWK correction (-ve).

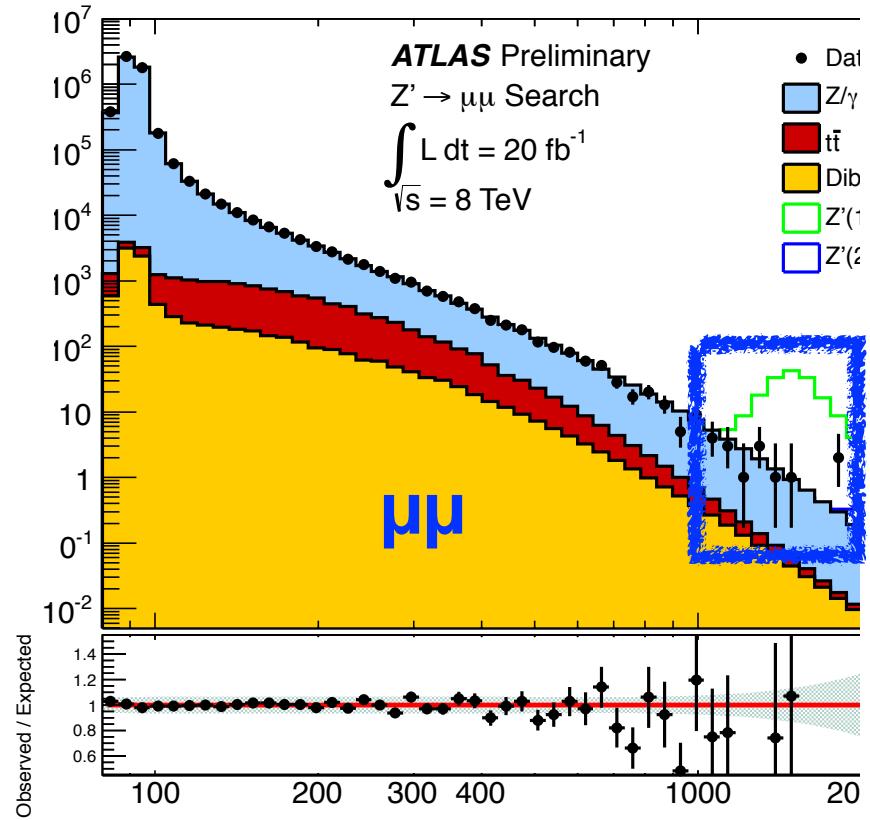
Zykunov, arXiv: hep-ph/0702-203

Too strong for Z? Yes, in some cases.

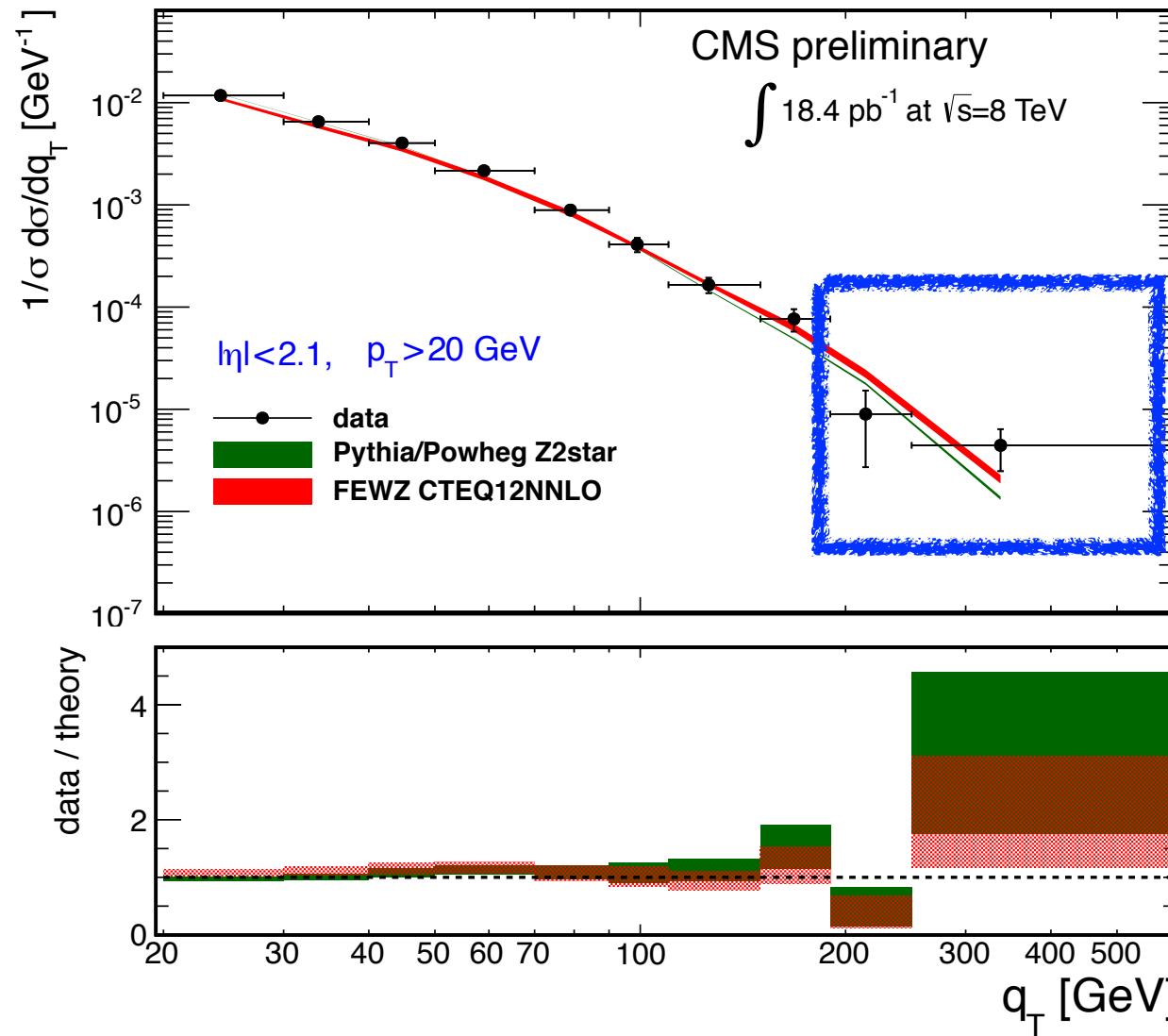


Dozens of Z events above 1 TeV!
Statistical precision $>1 \text{ TeV}$ is $\sim 10\%$.
So in the Sudakov zone.

<http://cds.cern.ch/record/1525524>



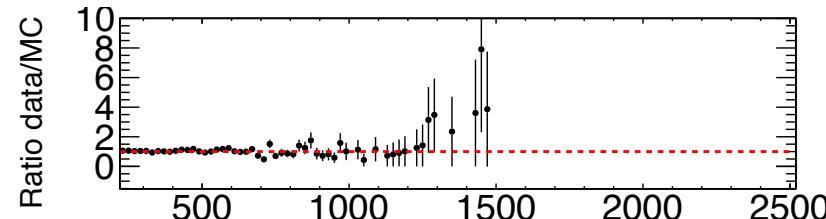
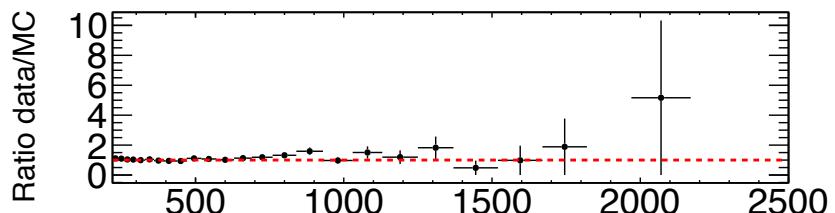
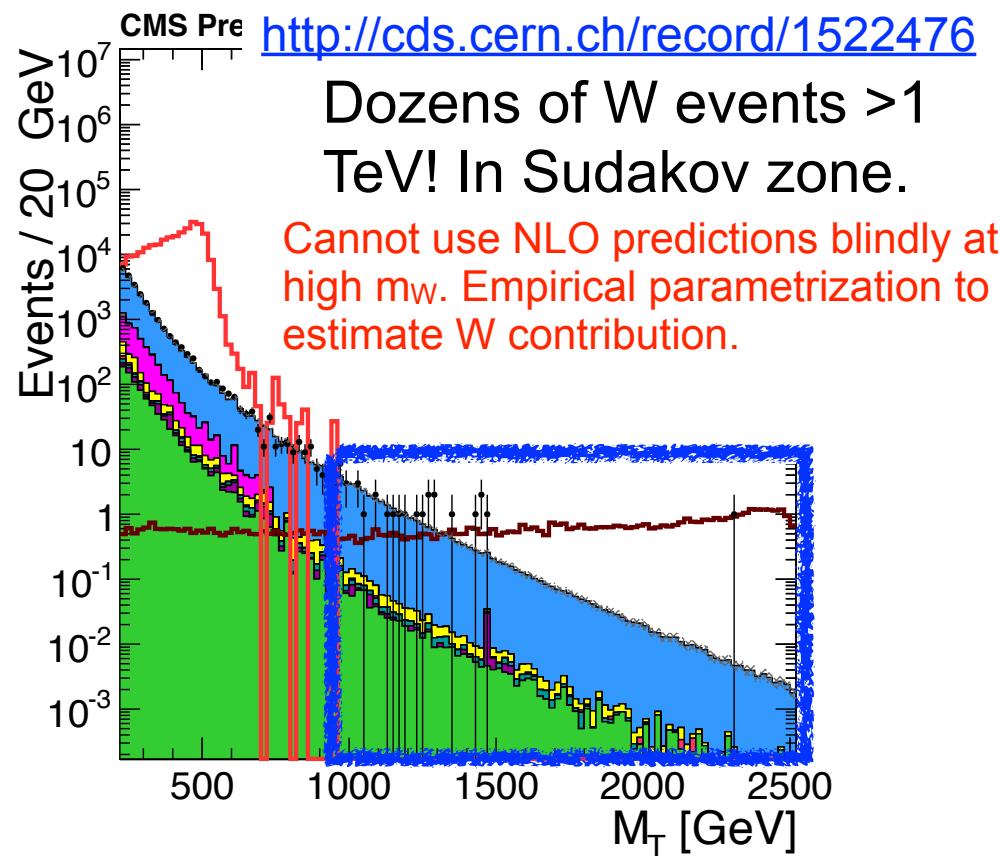
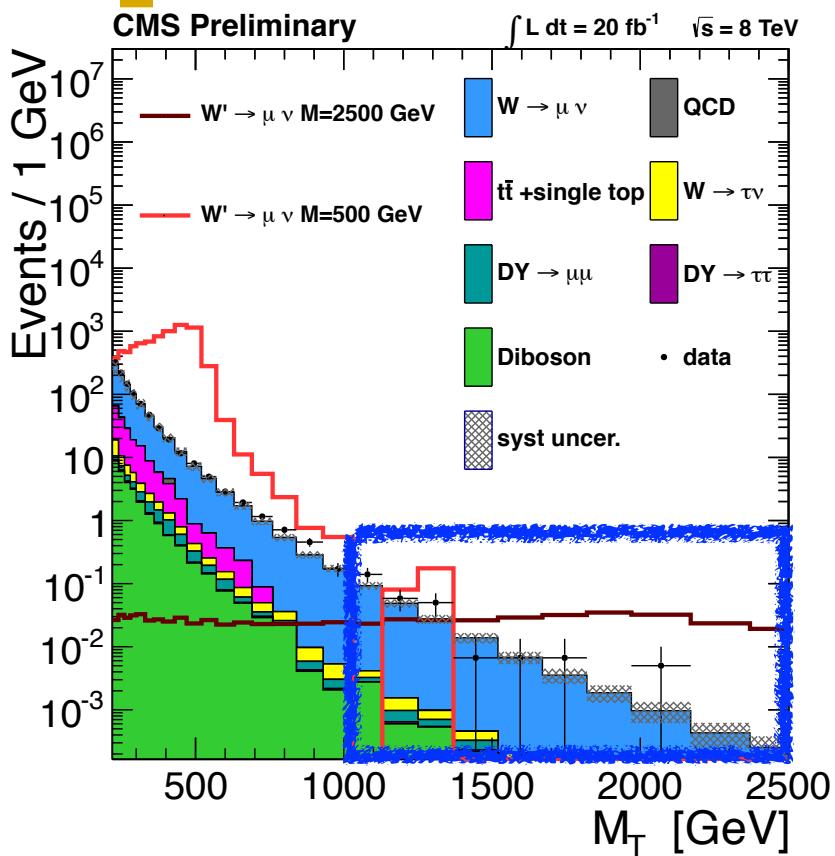
Z p_T distribution in 8 TeV LHC data



<https://cdsweb.cern.ch/record/1528579>

Expect $\sim 10\%$ negative correction in the last few bins. However, the effect gets diluted due to NLO QCD correction.

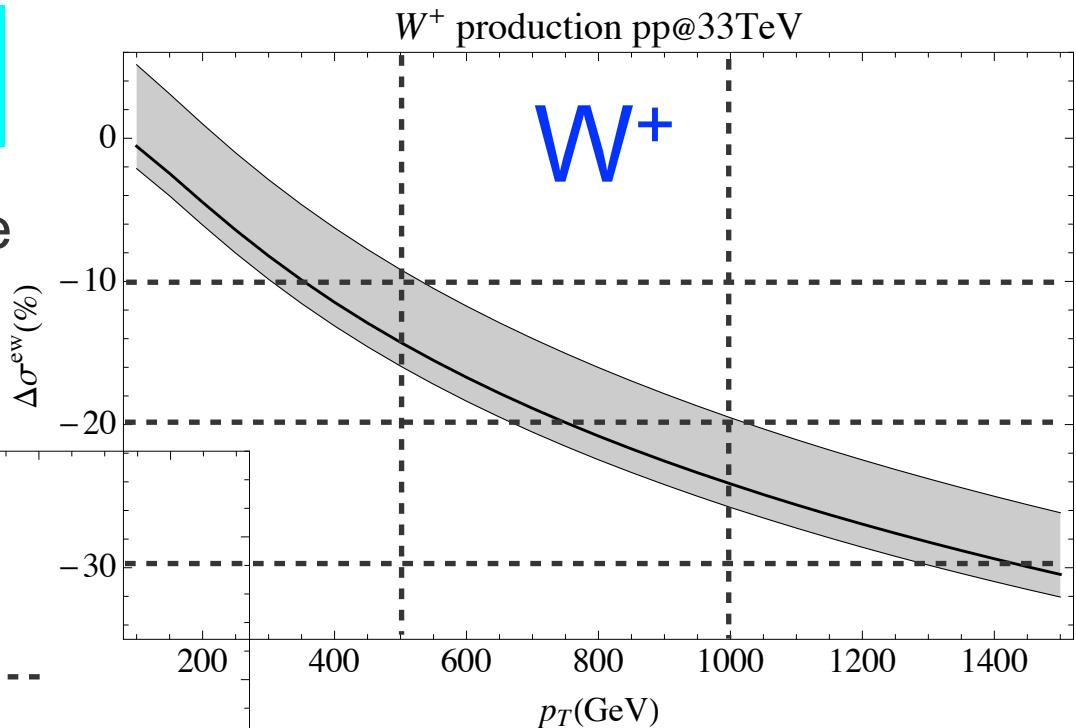
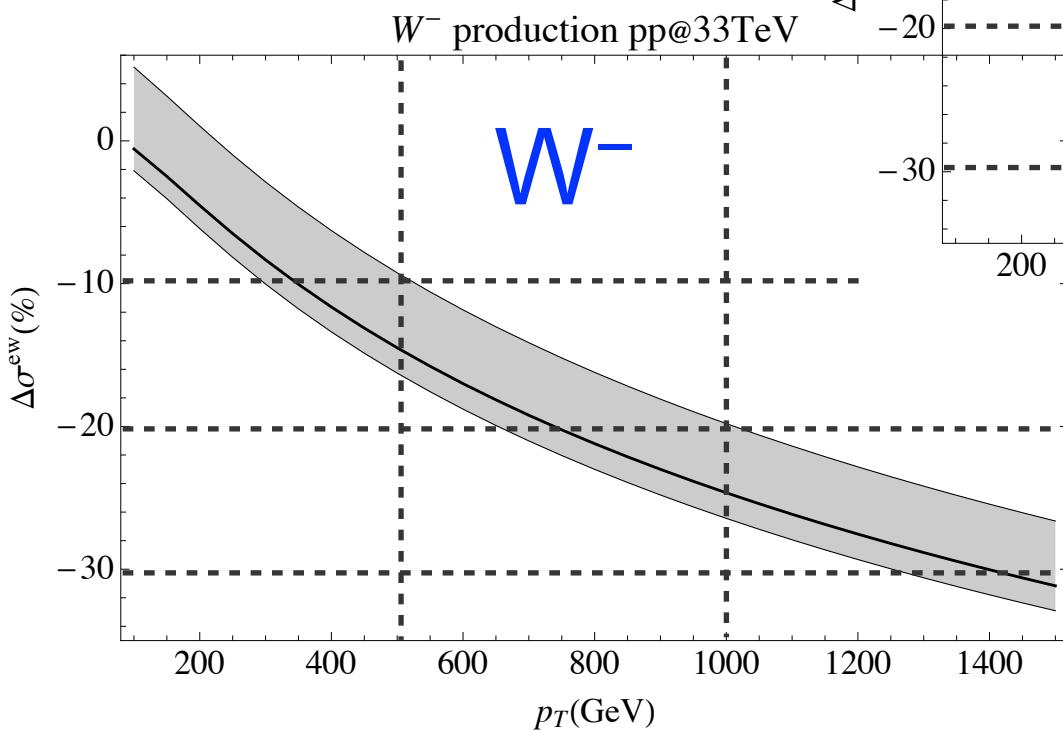
Too strong for W? Yes



Inclusive W^+ and W^- at 33 TeV

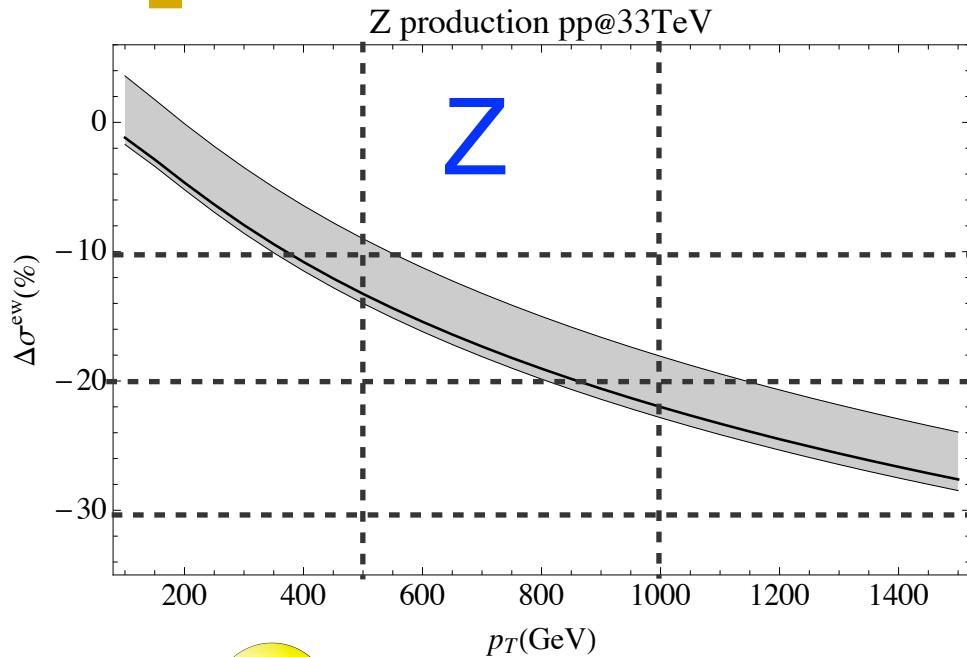
Provided by Xavier Garcia
Tormo & Thomas Becher

For 8 TeV & 14 TeV results see
[arXiv:1305.4202](https://arxiv.org/abs/1305.4202)



- About the same for W^+ & W^-
- Corrections \sim same as in the case of 8 TeV and 14 TeV

Inclusive Z and γ at 33 TeV



Bad

Getting increasingly more important b/c precision of experimental data & NLO as calculation finer than EWK corr.

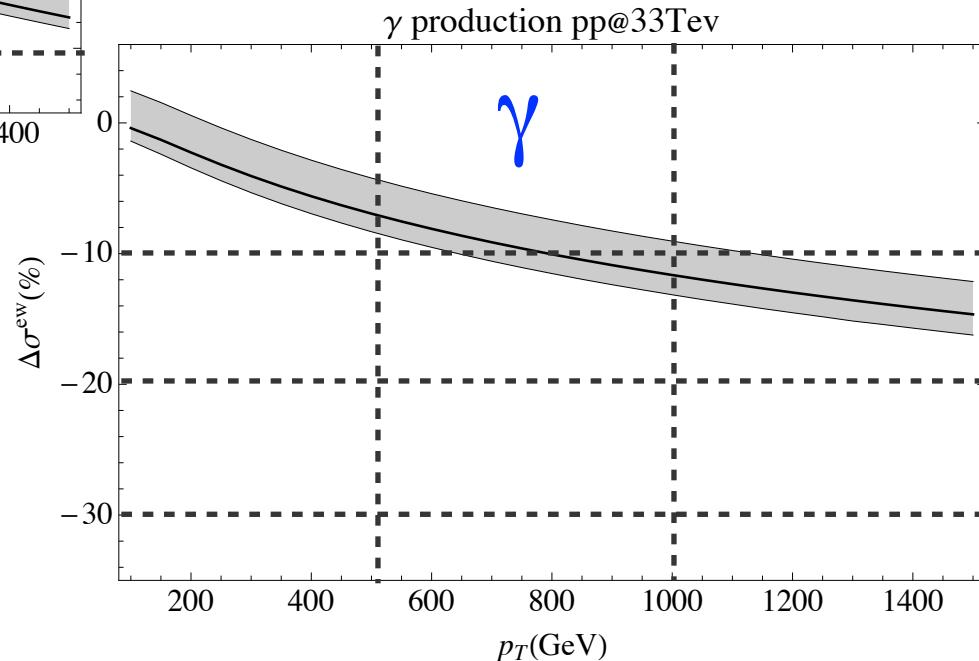
Coincidental cancellation of worryingly large magnitude.

Provided by Xavier Garcia Tormo & Thomas Becher



Good

EWK corrections almost independent of the collider energy.



Z+jets at 33 TeV as func. of $M_{\ell\ell}$

Provided by Tobias Kasprzik
in the format of arXiv:1103.0914

pp $\rightarrow l^+l^-$ jet + X at $\sqrt{s} = 33$ TeV

$M_{\ell\ell}$ / GeV	100 – ∞	200 – ∞	400 – ∞	1000 – ∞	2000 – ∞	4000 – ∞
$\sigma_{\text{Born}}^{\mu=M_Z}$ / pb	19.924(6)	1.6890(2)	0.28005(4)	0.022682(3)	0.0024968(4)	0.00015626(2)
$\sigma_{\text{Born}}^{\text{var}}$ / pb	19.849(6)	1.6482(2)	0.26618(4)	0.020604(3)	0.0021755(3)	0.00012945(2)
$\delta_{\text{EW}}^{\mu=M_Z}$ / %	-9.6(1)	-5.74(1)	-8.26(1)	-14.31(2)	-21.69(3)	-31.92(4)
$\delta_{\text{EW}}^{\text{rec}, \mu=M_Z}$ / %	-5.3(1)	-3.06(1)	-5.13(1)	-10.11(2)	-16.34(3)	-24.77(4)
$\delta_{\text{EW}}^{\text{var}}$ / %	-9.46(8)	-5.69(1)	-8.14(1)	-14.18(2)	-21.56(3)	-31.88(6)
$\delta_{\text{EW}}^{\text{rec, var}}$ / %	-5.05(7)	-2.94(1)	-4.93(1)	-9.93(2)	-16.14(3)	-24.75(6)
$\delta_{\text{QCD}}^{\mu=M_Z}$ / %	29.(1)	14.8(2)	-0.6(1)	-29.5(1)	-57.1(1)	-89.5(1)
$\delta_{\text{QCD}}^{\text{var}}$ / %	27.9(6)	15.9(2)	2.2(1)	-23.0(1)	-46.6(1)	-74.8(1)
$\delta_{\text{QCD, veto}}^{\mu=M_Z}$ / %	5.0(6)	-8.9(2)	-25.5(1)	-54.9(1)	-82.4(1)	-115.3(1)
$\delta_{\text{QCD, veto}}^{\text{var}}$ / %	6.1(4)	-7.2(2)	-20.8(2)	-45.8(1)	-69.8(1)	-98.1(1)
$\delta_{\gamma, \text{Born}}^{\mu=M_Z}$ / %	0.669(1)	2.097(5)	2.409(6)	2.168(6)	1.844(5)	1.610(3)
$\delta_{\gamma, \text{Born}}^{\text{var}}$ / %	0.710(1)	2.298(5)	2.721(7)	2.510(7)	2.135(5)	1.822(3)
$\sigma_{\text{full, veto}}^{\text{var}}$ / pb / %	19.32(9)	1.473(4)	0.1965(6)	0.00877(2)	0.000234(3)	-0.0000365(2)

Z+jets at 33 TeV as func. of jet p_T

Provided by Tobias Kasprzik
in the format of arXiv:1103.0914

$pp \rightarrow l^+l^- \text{ jet} + X$ at $\sqrt{s} = 33 \text{ TeV}$

$p_{T,\text{jet}}/\text{GeV}$	$50 - \infty$	$100 - \infty$	$200 - \infty$	$400 - \infty$	$1000 - \infty$	$2000 - \infty$
$\sigma_{\text{Born}}^{\mu=M_Z}/\text{pb}$	121.26(1)	33.737(4)	6.3314(7)	0.79429(9)	0.026920(2)	0.0010264(1)
$\sigma_{\text{Born}}^{\text{var}}/\text{pb}$	120.70(1)	32.654(4)	5.6403(7)	0.60157(6)	0.014491(1)	0.00038663(4)
$\delta_{\text{EW}}^{\mu=M_Z}/\%$	-4.68(1)	-5.47(1)	-9.25(1)	-16.15(2)	-29.25(5)	-41.8(2)
$\delta_{\text{EW}}^{\text{rec}, \mu=M_Z}/\%$	-3.29(3)	-4.48(2)	-8.50(4)	-15.42(3)	-28.59(8)	-41.0(1)
$\delta_{\text{EW}}^{\text{var}}/\%$	-4.61(1)	-5.31(1)	-8.92(1)	-15.65(2)	-28.48(5)	-40.5(3)
$\delta_{\text{EW}}^{\text{rec, var}}/\%$	-3.16(3)	-4.42(4)	-8.12(2)	-14.89(2)	-27.62(5)	-39.9(1)
$\delta_{\text{QCD}}^{\mu=M_Z}/\%$	56.1(2)	78.7(1)	113.0(1)	160.5(1)	240.4(2)	330.6(3)
$\delta_{\text{QCD}}^{\text{var}}/\%$	54.0(1)	77.4(1)	117.3(1)	186.0(1)	347.2(2)	609.0(4)
$\delta_{\text{QCD, veto}}^{\mu=M_Z}/\%$	13.1(1)	23.3(1)	29.7(1)	27.5(1)	-2.8(1)	-42.2(1)
$\delta_{\text{QCD, veto}}^{\text{var}}/\%$	12.9(1)	25.6(2)	38.8(1)	49.8(1)	52.6(1)	54.0(1)
$\delta_{\gamma, \text{Born}}^{\mu=M_Z}/\%$	0.1377(4)	0.1647(4)	0.1960(5)	0.2442(7)	0.341(1)	0.457(1)
$\delta_{\gamma, \text{Born}}^{\text{var}}/\%$	0.1489(4)	0.1923(5)	0.2562(7)	0.371(1)	0.661(2)	1.095(3)
$\sigma_{\text{full, veto}}^{\text{var}}/\text{pb}/\%$	130.9(2)	39.36(6)	7.342(9)	0.8091(7)	0.01808(2)	0.000443(1)

See backup slides for prediction as a function of jet p_T

Z+jets at 100 TeV as func. of M_{ll}

Provided by Tobias Kasprzik
in the format of arXiv:1103.0914

$pp \rightarrow l^+l^- \text{ jet} + X \text{ at } \sqrt{s} = 100 \text{ TeV}$

M_{ll}/GeV	100 – ∞	200 – ∞	400 – ∞	1000 – ∞	2000 – ∞	4000 – ∞
$\sigma_{\text{Born}}^{\mu=M_Z}/\text{pb}$	60.70(2)	5.3194(9)	0.9142(1)	0.08462(1)	0.012762(2)	0.0016229(2)
$\sigma_{\text{Born}}^{\text{var}}/\text{pb}$	61.76(2)	5.3548(9)	0.9016(1)	0.07980(1)	0.011500(2)	0.0013850(2)
$\delta_{\text{EW}}^{\mu=M_Z}/\%$	-9.4(1)	-5.85(1)	-8.41(1)	-14.28(2)	-21.05(3)	-29.97(5)
$\delta_{\text{EW}}^{\text{rec}, \mu=M_Z}/\%$	-5.0(1)	-3.23(1)	-5.29(1)	-10.42(3)	-16.42(3)	-24.34(5)
$\delta_{\text{EW}}^{\text{var}}/\%$	-9.5(2)	-5.76(1)	-8.31(1)	-14.08(2)	-20.85(3)	-29.87(5)
$\delta_{\text{EW}}^{\text{rec, var}}/\%$	-5.2(2)	-3.07(1)	-5.15(2)	-10.11(2)	-16.22(4)	-24.09(5)
$\delta_{\text{QCD}}^{\mu=M_Z}/\%$	36.(4)	20.8(3)	7.3(2)	-18.1(1)	-43.3(1)	-73.7(1)
$\delta_{\text{QCD}}^{\text{var}}/\%$	29.3(9)	18.4(3)	6.0(1)	-16.3(1)	-37.4(1)	-63.3(1)
$\delta_{\text{QCD, veto}}^{\mu=M_Z}/\%$	5.0(9)	-7.0(2)	-21.2(2)	-46.8(1)	-71.8(1)	-102.1(1)
$\delta_{\text{QCD, veto}}^{\text{var}}/\%$	3.8(7)	-7.4(3)	-19.6(2)	-41.7(1)	-63.1(1)	-88.5(2)
$\delta_{\gamma, \text{Born}}^{\mu=M_Z}/\%$	0.532(1)	1.679(4)	2.019(6)	1.988(6)	1.792(6)	1.511(4)
$\delta_{\gamma, \text{Born}}^{\text{var}}/\%$	0.568(1)	1.871(4)	2.363(7)	2.464(7)	2.259(7)	1.901(5)
$\sigma_{\text{full, veto}}^{\text{var}}/\text{pb}/\%$	58.6(4)	4.75(1)	0.672(1)	0.0373(1)	0.00211(2)	-0.000228(3)

Z+jets at 100 TeV as func. of jet p_T

Provided by Tobias Kasprzik
in the format of arXiv:1103.0914

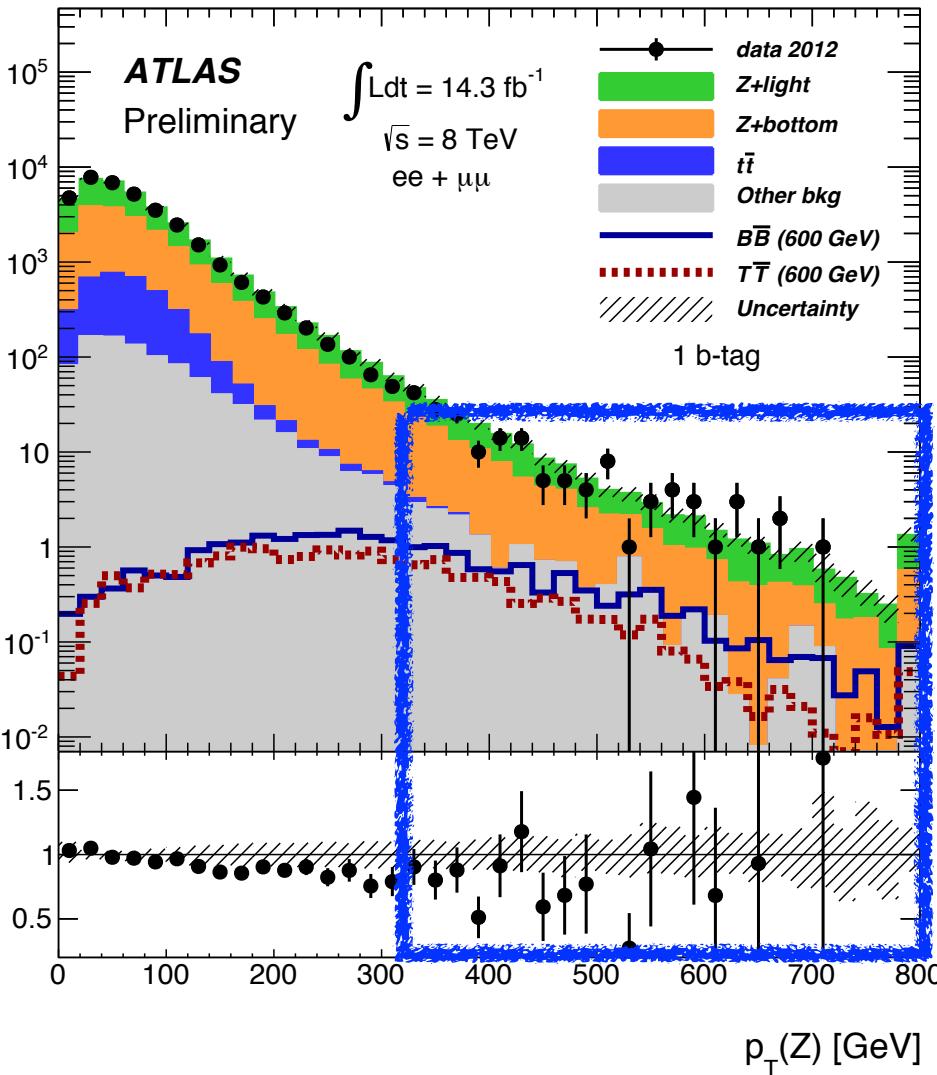
pp → l⁺l⁻ jet + X at $\sqrt{s} = 100 \text{ TeV}$

$p_{\text{T},\text{jet}}/\text{GeV}$	100 – ∞	200 – ∞	400 – ∞	800 – ∞	2000 – ∞	4000 – ∞
$\sigma_{\text{Born}}^{\mu=M_Z}/\text{pb}$	114.29(1)	23.772(3)	3.5452(4)	0.42003(4)	0.017238(1)	0.00094403(9)
$\sigma_{\text{Born}}^{\text{var}}/\text{pb}$	118.30(1)	23.762(3)	3.1922(3)	0.31583(3)	0.0091290(9)	0.00035205(3)
$\delta_{\text{EW}}^{\mu=M_Z}/\%$	-5.62(1)	-9.57(1)	-16.86(2)	-27.11(8)	-43.5(1)	-58.8(1)
$\delta_{\text{EW}}^{\text{rec},\mu=M_Z}/\%$	-4.65(3)	-8.72(2)	-16.08(2)	-26.29(4)	-43.15(7)	-58.5(2)
$\delta_{\text{EW}}^{\text{var}}/\%$	-5.50(1)	-9.29(1)	-16.38(3)	-26.36(4)	-43.2(2)	-57.5(1)
$\delta_{\text{EW}}^{\text{rec,var}}/\%$	-4.48(2)	-8.52(2)	-15.62(2)	-25.64(4)	-42.21(7)	-56.8(1)
$\delta_{\text{QCD}}^{\mu=M_Z}/\%$	97.4(2)	146.0(1)	215.2(2)	288.7(2)	378.0(3)	472.6(5)
$\delta_{\text{QCD}}^{\text{var}}/\%$	85.4(2)	130.0(2)	201.7(1)	298.8(2)	487.9(3)	769.0(7)
$\delta_{\text{QCD,veto}}^{\mu=M_Z}/\%$	35.7(2)	54.2(1)	66.7(1)	61.3(1)	13.2(2)	-43.1(1)
$\delta_{\text{QCD,veto}}^{\text{var}}/\%$	29.6(2)	47.4(1)	65.5(1)	76.4(3)	65.6(2)	51.5(1)
$\delta_{\gamma,\text{Born}}^{\mu=M_Z}/\%$	0.1218(3)	0.1400(4)	0.1681(5)	0.2114(7)	0.291(1)	0.382(1)
$\delta_{\gamma,\text{Born}}^{\text{var}}/\%$	0.1407(3)	0.1799(5)	0.2482(7)	0.365(1)	0.630(2)	1.006(5)
$\sigma_{\text{full,veto}}^{\text{var}}/\text{pb}/\%$	147.0(2)	32.86(4)	4.767(4)	0.475(1)	0.01124(2)	0.0003343(7)

See backup slides for prediction as a function of jet p_T

Distribution in 8 TeV LHC data

Entries / 20 GeV



ATLAS analysis of $Z + \text{jets}$ data with one b-tagged jet

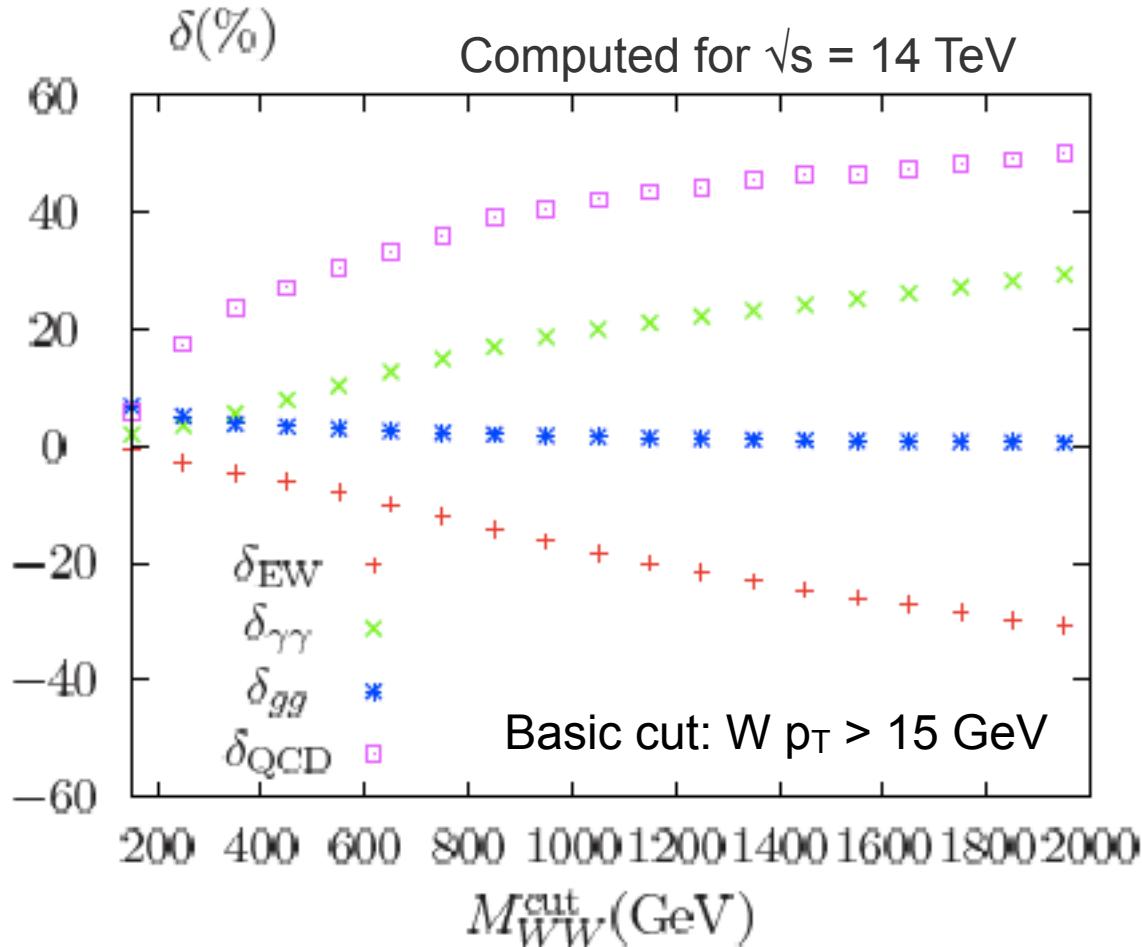
<http://cds.cern.ch/record/1557773>

In Sudakov zone

Additional experimental plot
in backup

WW production: a curious case

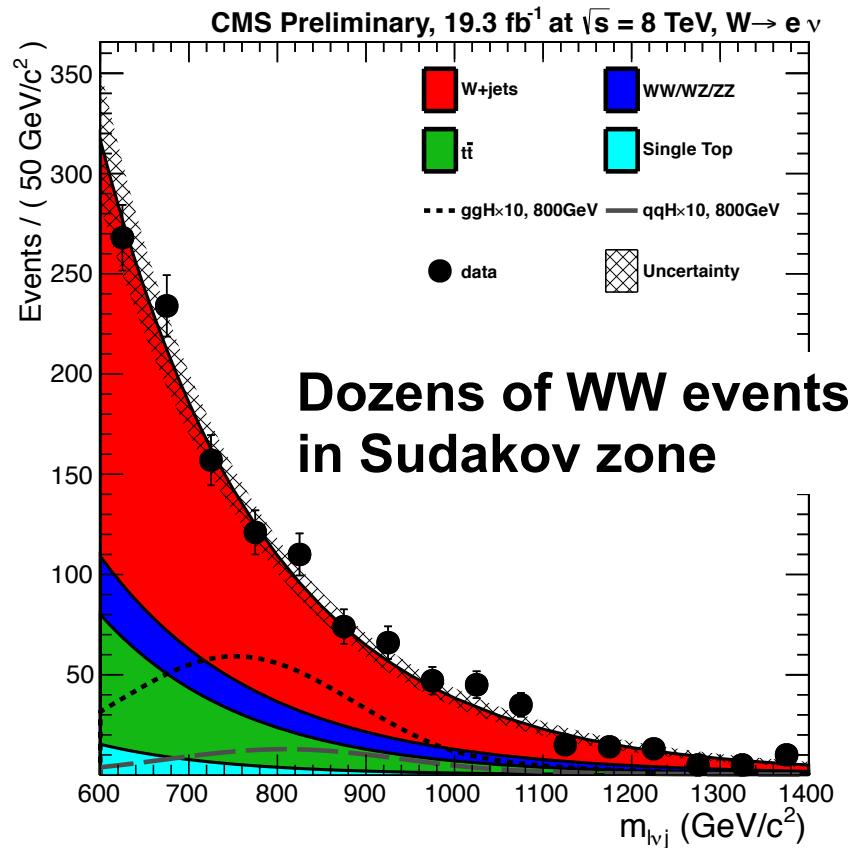
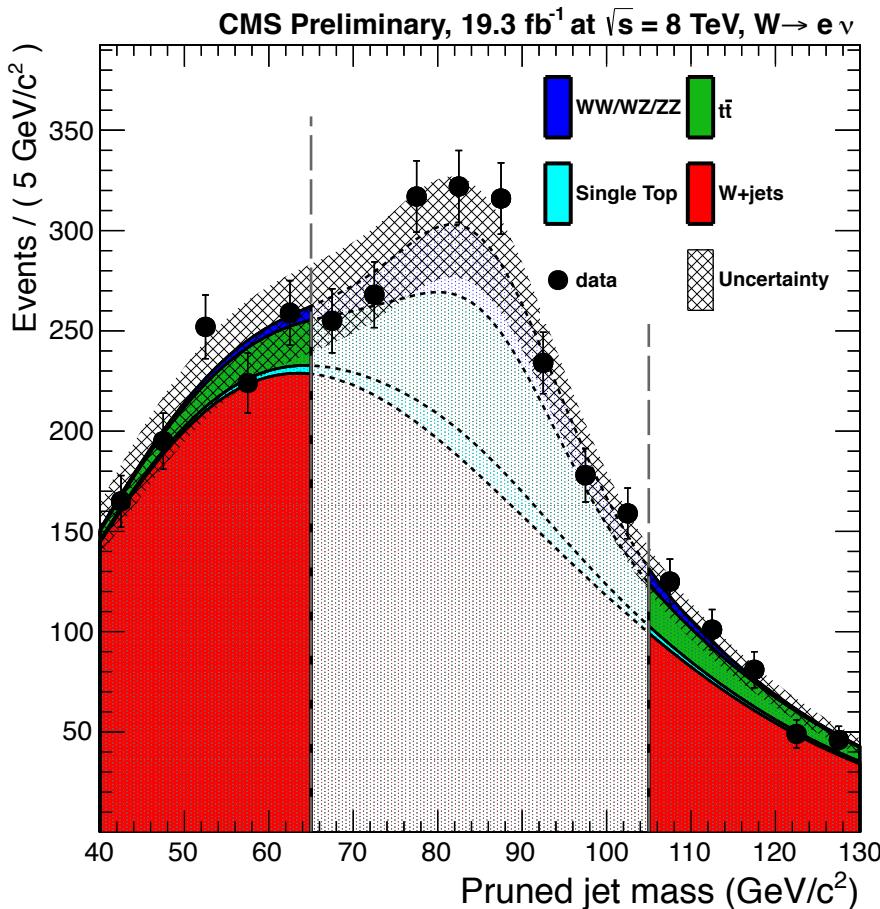
See talk by T. Kasprzik at ICHEP 2012



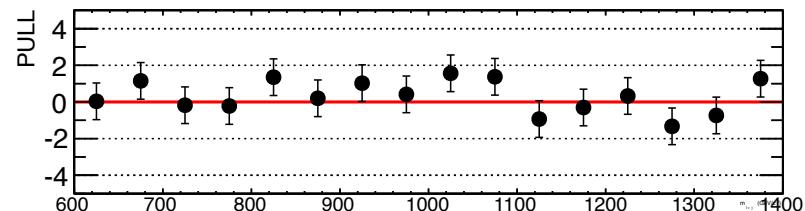
EWK and QCD corrections partially cancel out, resulting in net positive enhancement at high WW invariant mass

Again, coincidental cancellation of worryingly large magnitude.

Too strong for WW? Yes, in some cases.

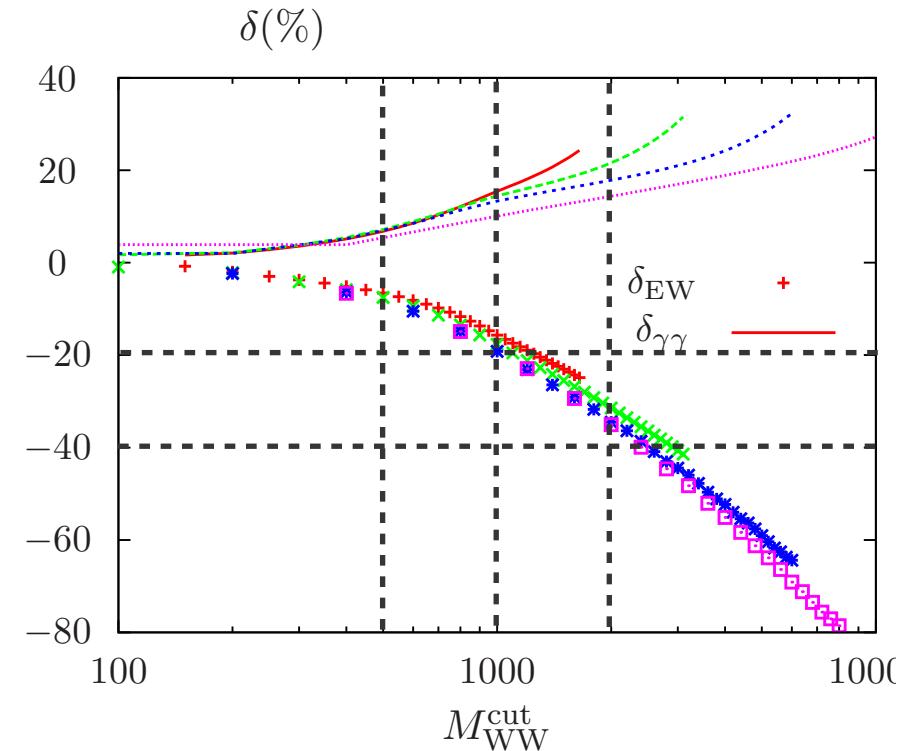
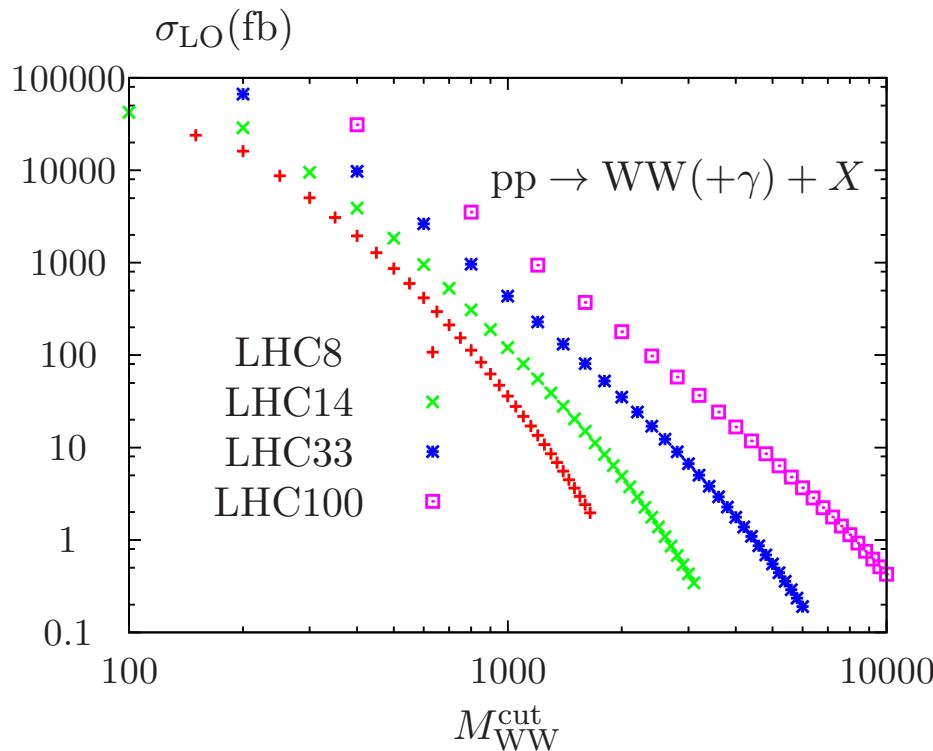


Additional experimental plot in backup



WW at 33 TeV and 100 TeV

Provided by Tobias Kasprzak

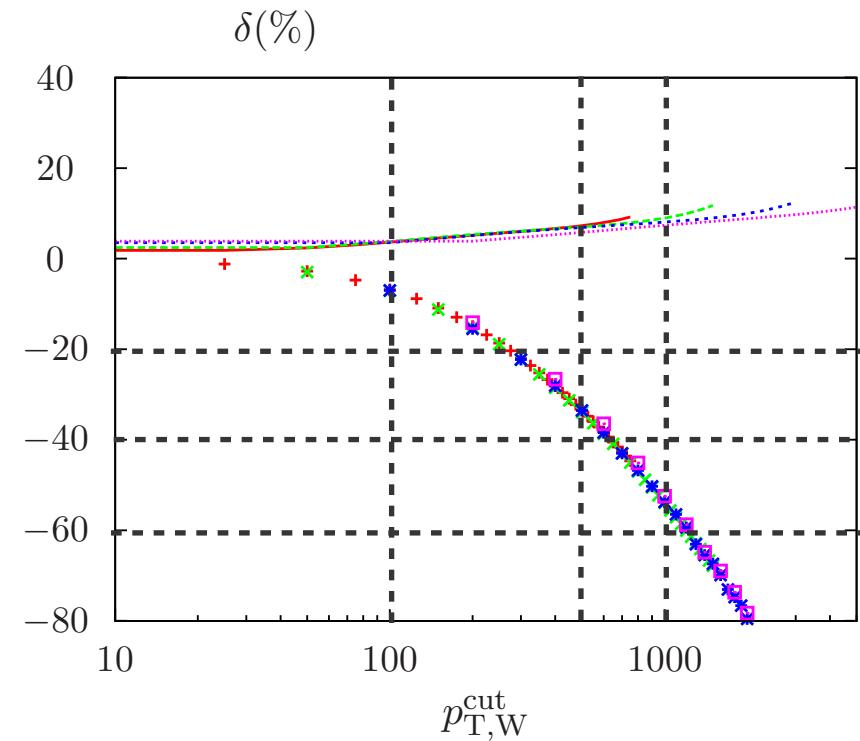
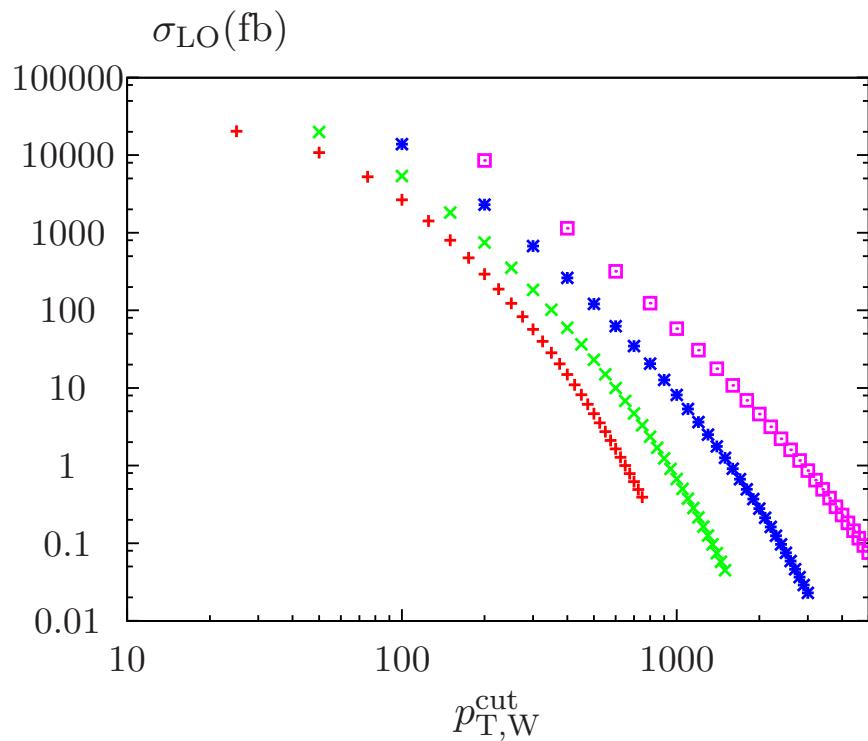


The EWK corrections by themselves do not depend much on the collider energy, but $\gamma\gamma \rightarrow WW$ increases at a much slower rate at high collider energies. So, the sum of the two is more -ve. Need to watch out for NLO QCD correction to understand the net cancellation effect.

WW at 33 TeV and 100 TeV

Provided by Tobias Kasprzik

Also computed as a function of $p_{T,W}$ cut



Summary: Quick recap of Sudakov zone survey

Are we in the Sudakov zone ?

Process	Now	14 TeV	33/100 TeV
Inclusive jet, dijet	Yes	Yes	Yes
Inclusive W/Z tail	Yes	Yes	Yes
$V\gamma (\ell\ell\gamma, \ell\nu\gamma)$ tail	No	Yes	Yes
$V+jets$ tail	Yes	Yes	Yes
WW leptonic	Close	Yes	Yes
WZ,ZZ leptonic	No	No	Yes
VV semi-leptonic	Yes	Yes	Yes

Previous report and additional survey of 7 TeV LHC data at:
<http://conference.ippp.dur.ac.uk/conferenceDisplay.py?confId=331>

Summary

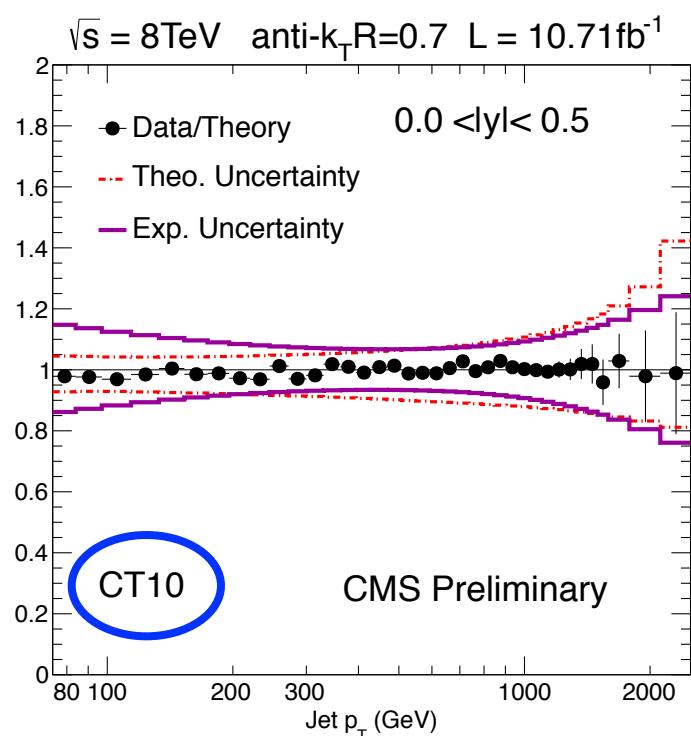
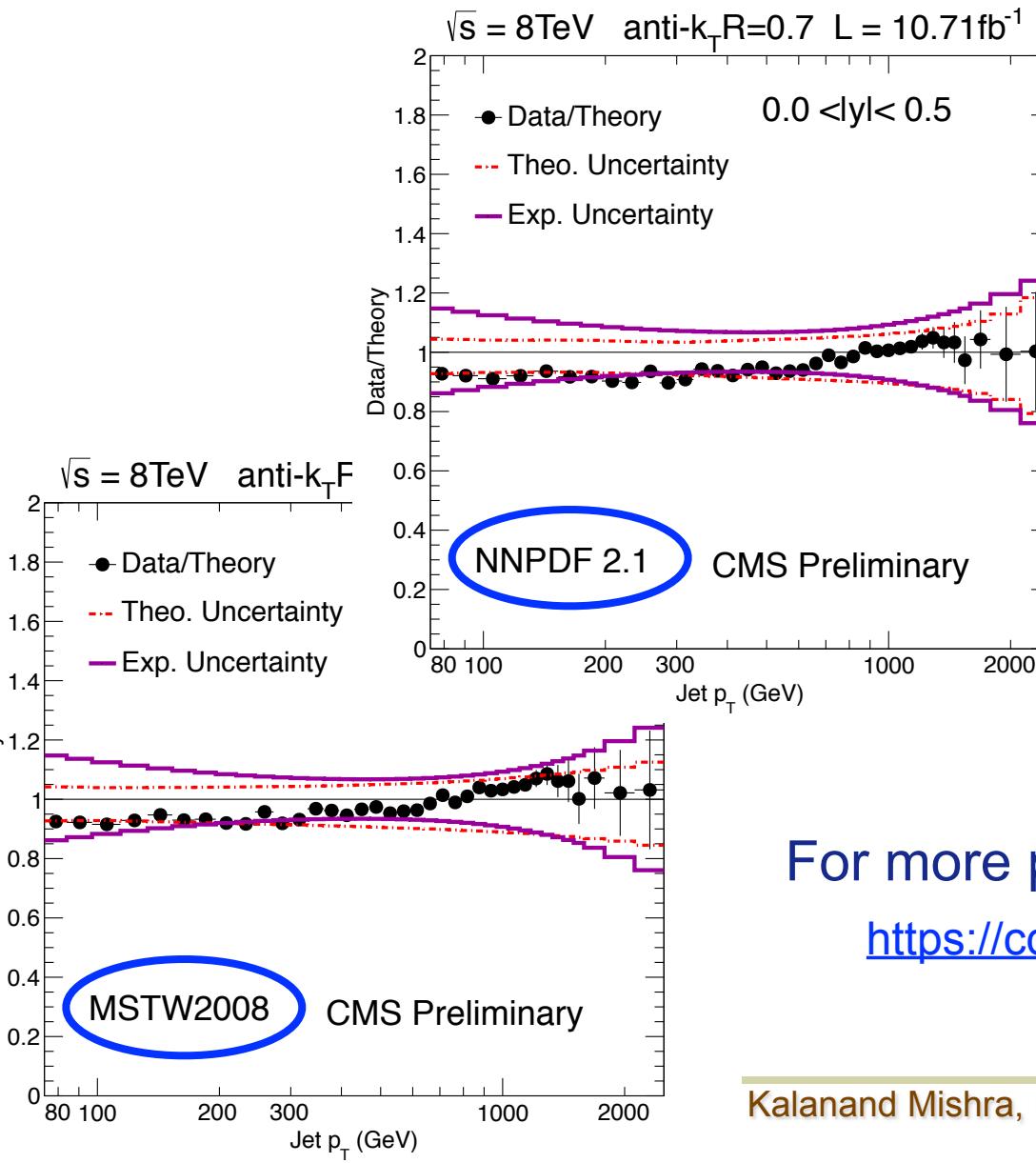
- ✓ The size of EWK corrections for many processes at LHC are approaching the current precision of the measurement
 - Important at high mass/ p_T tails

- ✓ Extrapolations made using current LHC data and dedicated calculations for 33 TeV & 100 TeV collider energies suggest that
 - Corrections are large in kin regime accessible at higher energies
 - Important for almost all abundant (background) processes
 - Do not change dramatically with collider energy

Plan to write up this report & include in the Snowmass white paper

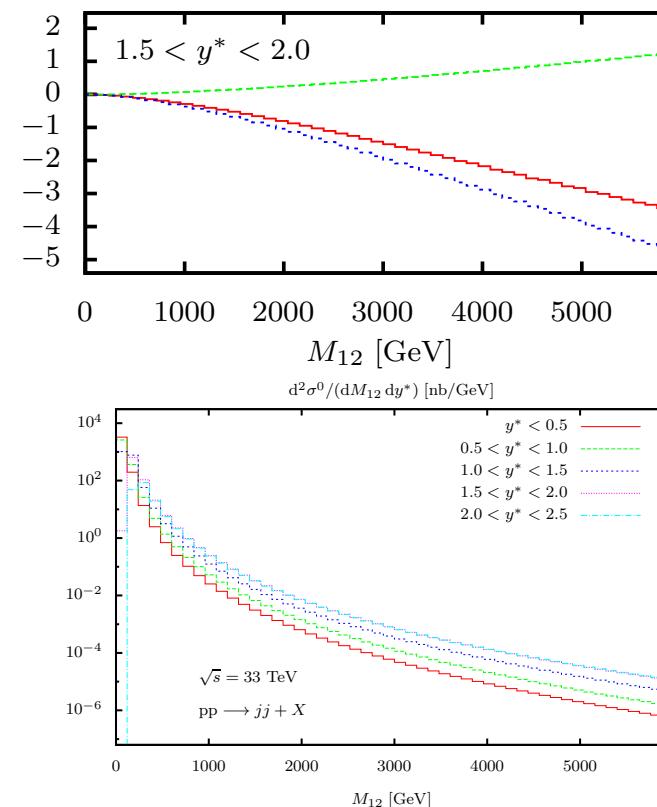
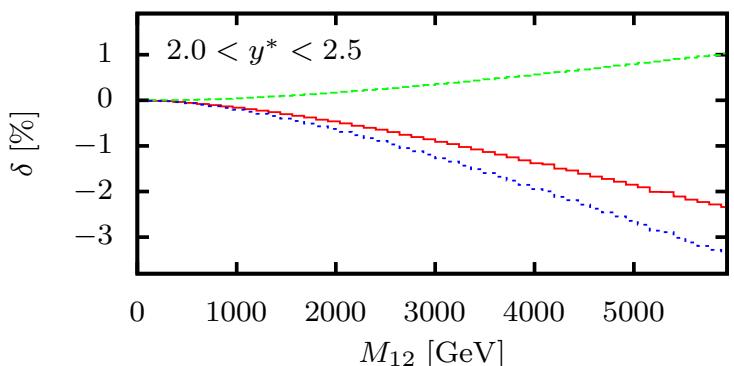
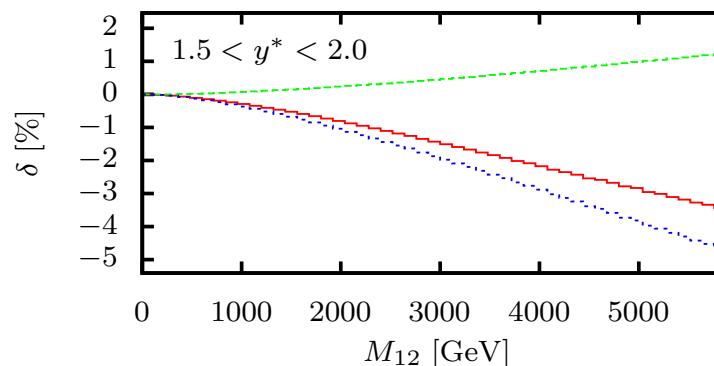
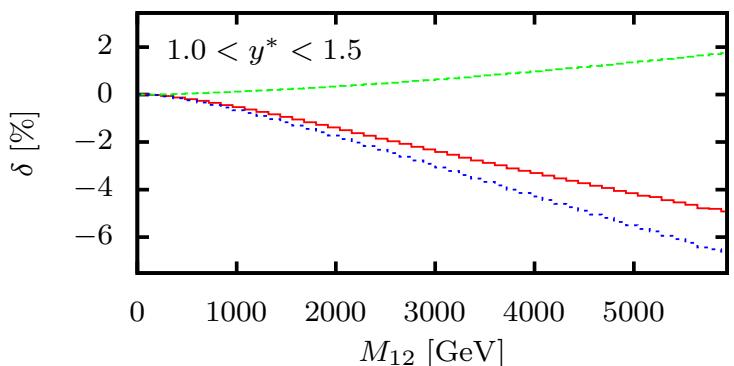
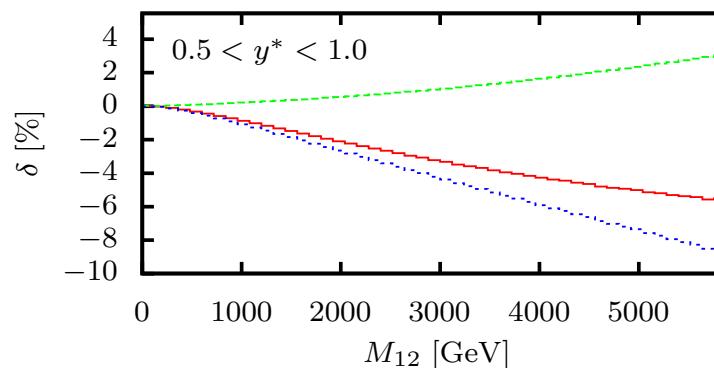
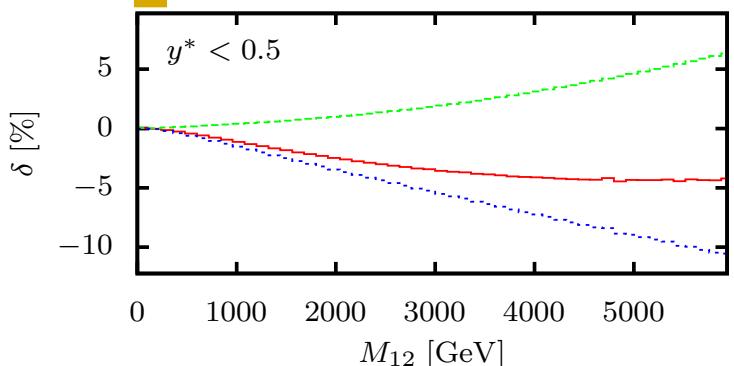
BACKUP SLIDES

Aside: PDF dependence of QCD (α_s) uncertainty

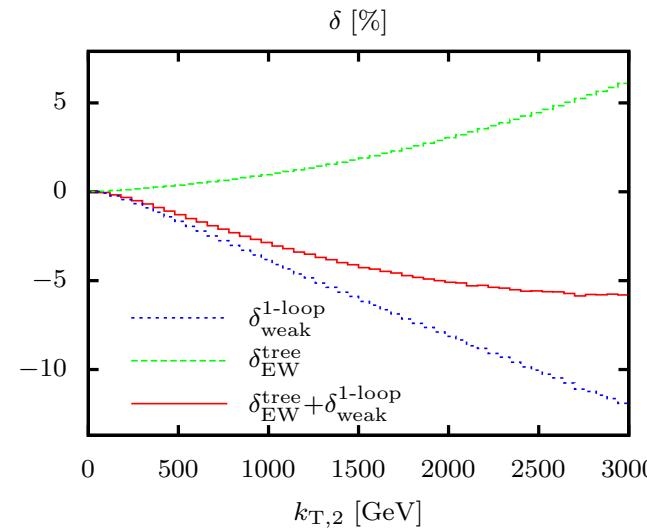
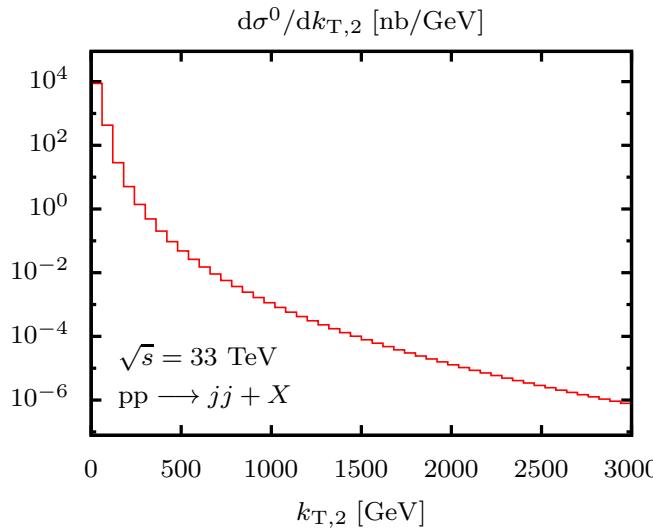
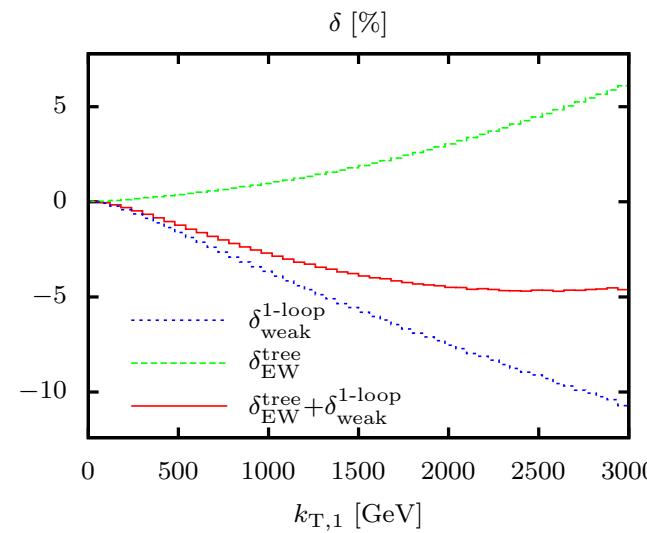
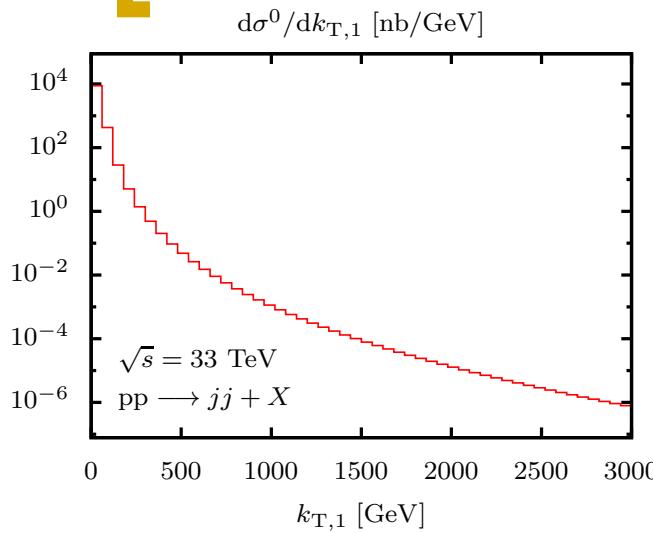


For more plots (other rapidity bins) see
<https://cdsweb.cern.ch/record/1547589>

Dijet @33 TeV: Dissecting into rapidity regions



Dijet @33 TeV: differential in jet p_T



Z+jets at 33 TeV: as a function of lepton p_T

pp → l^+l^- jet + X at $\sqrt{s} = 33$ TeV

$p_{T,l+}$ / GeV	50 – ∞	100 – ∞	200 – ∞	400 – ∞	1000 – ∞	2000 – ∞
$\sigma_{\text{Born}}^{\mu=M_Z}$ / pb	138.38(2)	12.641(2)	1.7316(3)	0.19986(3)	0.006439(1)	0.00023702(5)
$\sigma_{\text{Born}}^{\text{var}}$ / pb	138.01(2)	12.011(2)	1.4856(2)	0.14747(2)	0.0037329(7)	0.00011702(2)
$\delta_{\text{EW}}^{\mu=M_Z}$ / %	-6.99(1)	-9.26(2)	-13.57(7)	-19.96(7)	-31.0(2)	-37.1(3)
$\delta_{\text{EW}}^{\text{rec}, \mu=M_Z}$ / %	-3.97(9)	-6.11(4)	-10.14(5)	-16.12(9)	-25.8(2)	-31.7(2)
$\delta_{\text{EW}}^{\text{var}}$ / %	-7.00(1)	-9.08(2)	-13.18(3)	-19.11(6)	-26.7(3)	-31.(2)
$\delta_{\text{EW}}^{\text{rec, var}}$ / %	-4.12(8)	-5.82(5)	-9.65(4)	-15.07(7)	-21.8(3)	-23.3(7)
$\delta_{\text{QCD}}^{\mu=M_Z}$ / %	32.3(2)	56.3(2)	49.6(2)	28.4(3)	-21.4(2)	-67.9(3)
$\delta_{\text{QCD}}^{\text{var}}$ / %	30.5(2)	53.2(2)	51.7(2)	42.0(2)	15.9(2)	-15.1(7)
$\delta_{\text{QCD, veto}}^{\mu=M_Z}$ / %	7.9(2)	22.5(3)	17.3(1)	-1.9(2)	-46.4(2)	-89.9(3)
$\delta_{\text{QCD, veto}}^{\text{var}}$ / %	7.7(2)	23.4(2)	25.5(3)	19.8(1)	-1.(1) × 10 ¹	-35.4(6)
$\delta_{\gamma, \text{Born}}^{\mu=M_Z}$ / %	0.1204(3)	0.3068(8)	0.426(1)	0.525(1)	0.683(3)	0.817(3)
$\delta_{\gamma, \text{Born}}^{\text{var}}$ / %	0.1284(4)	0.355(1)	0.559(1)	0.792(2)	1.200(4)	1.498(7)
$\sigma_{\text{full, veto}}^{\text{var}}$ / pb / %	139.2(3)	13.78(2)	1.677(4)	0.1497(3)	0.0025(7)	0.000041(2)

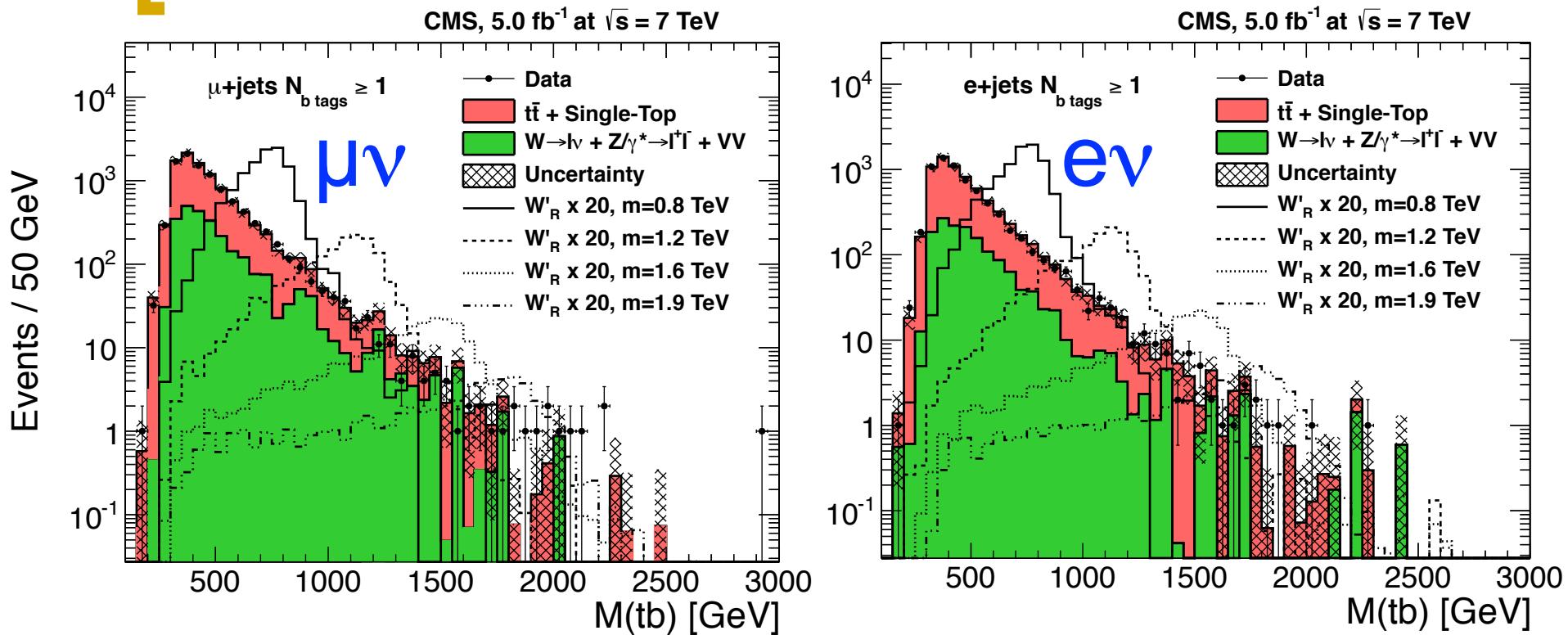
Z+jets at 100 TeV: as a function of lepton p_T

pp → l^+l^- jet + X at $\sqrt{s} = 100$ TeV

p_{T,l^+} / GeV	100 – ∞	200 – ∞	400 – ∞	800 – ∞	2000 – ∞	4000 – ∞
$\sigma_{\text{Born}}^{\mu=M_Z}$ / pb	44.075(8)	6.865(1)	0.9684(1)	0.11379(2)	0.0045639(8)	0.00024532(5)
$\sigma_{\text{Born}}^{\text{var}}$ / pb	45.107(9)	6.643(1)	0.8362(1)	0.08329(1)	0.0025868(5)	0.00011752(3)
$\delta_{\text{EW}}^{\mu=M_Z}$ / %	-9.68(2)	-14.56(3)	-22.5(3)	-31.1(1)	-44.5(3)	-51.8(8)
$\delta_{\text{EW}}^{\text{rec}, \mu=M_Z}$ / %	-6.53(5)	-11.17(4)	-18.17(6)	-26.8(1)	-39.8(2)	-49.(1)
$\delta_{\text{EW}}^{\text{var}}$ / %	-9.54(2)	-14.22(3)	-21.20(6)	-29.6(2)	-41.8(8)	-41.8(6)
$\delta_{\text{EW}}^{\text{rec, var}}$ / %	-6.32(6)	-10.81(4)	-17.38(6)	-25.1(1)	-39.(1)	-35.4(8)
$\delta_{\text{QCD}}^{\mu=M_Z}$ / %	73.6(3)	77.4(2)	68.2(2)	41.1(2)	-20.5(6)	-75.2(8)
$\delta_{\text{QCD}}^{\text{var}}$ / %	60.1(4)	62.5(2)	58.2(2)	44.9(2)	12.(2)	-24.4(4)
$\delta_{\text{QCD, veto}}^{\mu=M_Z}$ / %	34.7(2)	40.5(2)	33.4(2)	10.0(3)	-49.1(4)	-99.6(4)
$\delta_{\text{QCD, veto}}^{\text{var}}$ / %	26.9(2)	33.2(4)	32.7(1)	23.4(3)	-7.5(3)	-45.1(4)
$\delta_{\gamma, \text{Born}}^{\mu=M_Z}$ / %	0.2264(7)	0.3036(9)	0.368(1)	0.441(1)	0.564(3)	0.667(3)
$\delta_{\gamma, \text{Born}}^{\text{var}}$ / %	0.2562(7)	0.385(1)	0.538(1)	0.743(3)	1.093(5)	1.32(1)
$\delta_{\gamma, \text{NLO}}^{\mu=M_Z}$ / %	0.2264(7)	0.3036(9)	0.368(1)	0.441(1)	0.564(3)	0.667(3)
$\delta_{\gamma, \text{NLO}}^{\text{var}}$ / %	0.2562(7)	0.385(1)	0.538(1)	0.743(3)	1.093(5)	1.32(1)
$\sigma_{\text{full, veto}}^{\text{var}}$ / pb / %	53.1(1)	7.93(2)	0.937(1)	0.0788(3)	0.00134(2)	0.0000169(9)

[W+jets with at least 1 b-tag]

arXiv:1208.0956



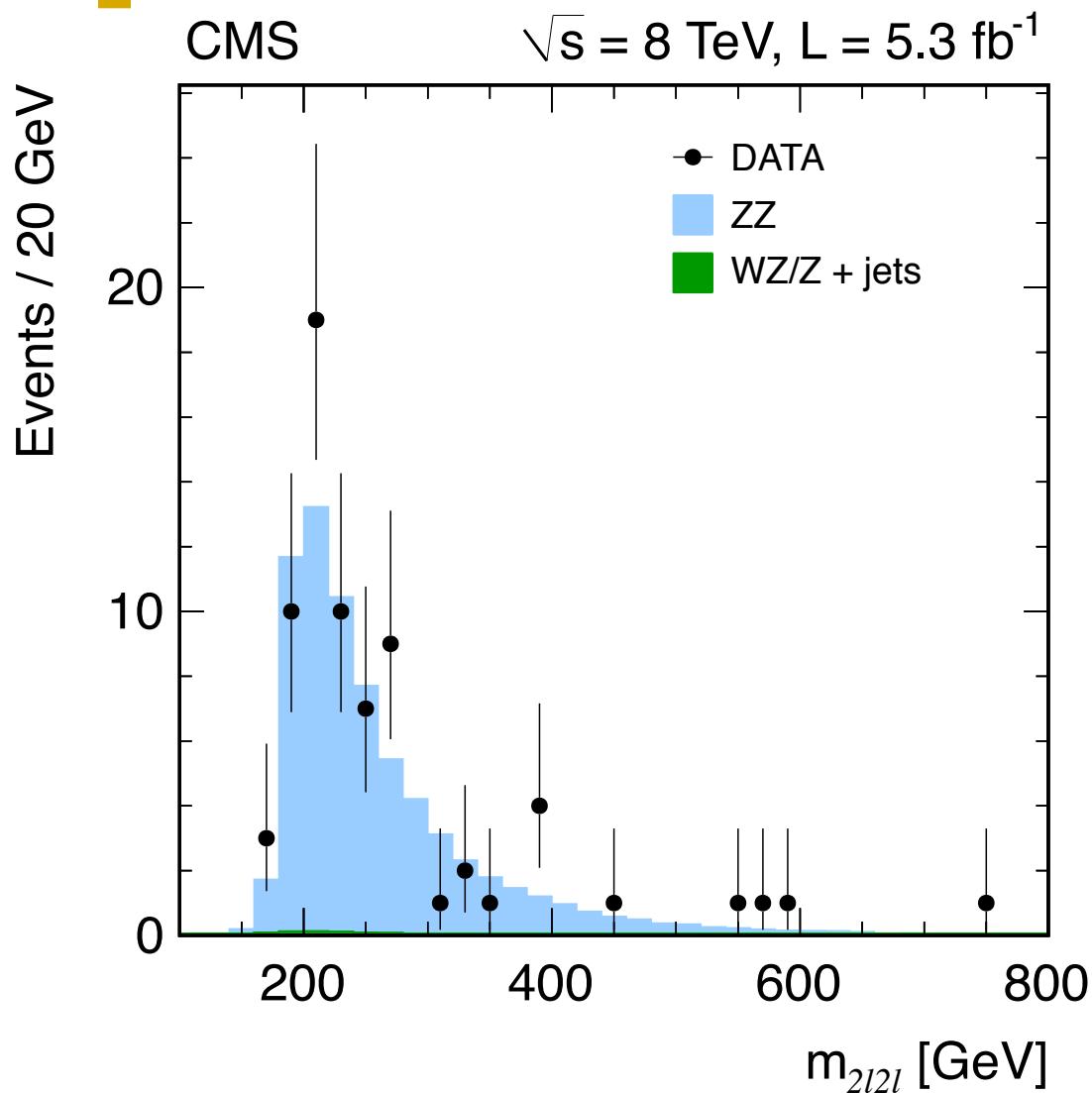
Dozens of events above 1 TeV

- Better than 5% precision expected in 8 TeV data

Already in Sudakov zone!

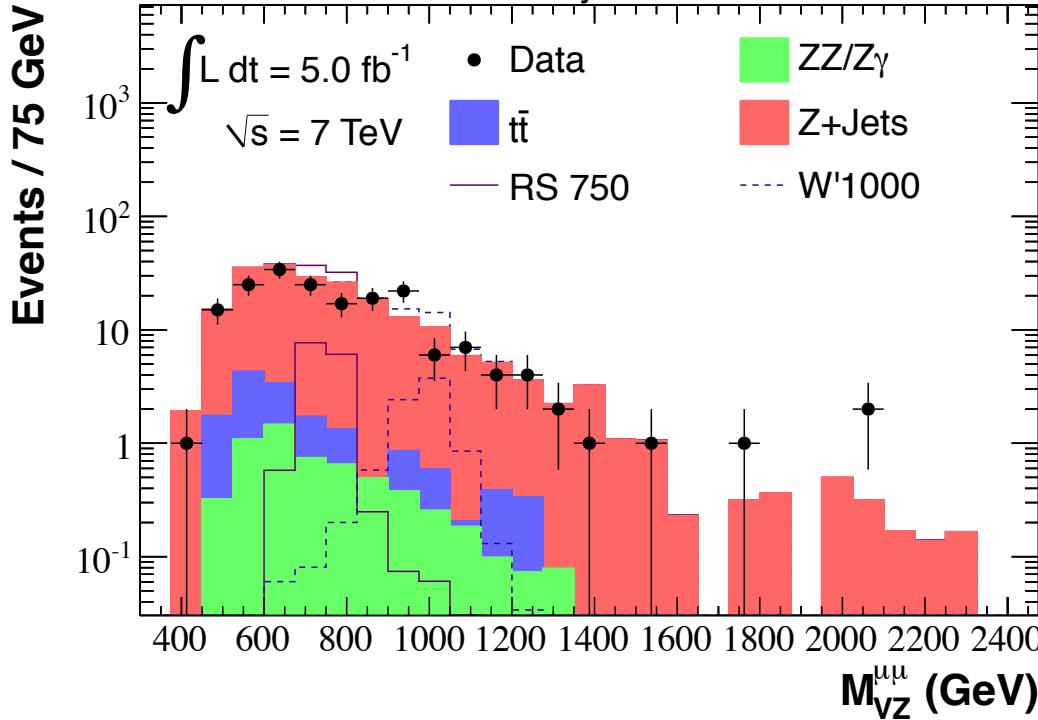
Validate W+jets shape using data: compare 0 b-tagged sample in data with W+jets MC; take difference as systematics.

[ZZ (leptonic) invariant mass in 8 TeV data



Diboson at 7 TeV: WZ/ZZ $\rightarrow\ell\ell+j$ (boosted)

<http://cdsweb.cern.ch/record/1444879>



- ◆ The other boson (W or Z) decays hadronically into a single (merged) jet
 - anti-kT 0.7 jet
 - highly boosted: $p_T > 250$ GeV
- ◆ Plot invariant mass of the VZ system.

A dozen events beyond 1 TeV. Will be in the Sudakov zone in 2012.

Fit the smoothly falling spectrum to set limit.