

Closing the Wedge with 300 and 3000 fb^{-1} at the LHC

Ian Lewis

Brookhaven National Lab

University of Washington
Snowmass Energy Frontier

June 30, 2013

Motivation

- We have discovered a Higgs boson:
 - Couplings to massive bosons.
 - Nonuniversal couplings to fermions.
 - Appears the new boson couples according to mass and is related to EWSB.
- Era of directly probing the mechanism of electroweak symmetry breaking (EWSB) has begun.
- Necessary to investigate if there is more to the EWSB sector.
- One popular beyond the standard model theory that addresses this is the Supersymmetry.
- For studies at the LHC energies it is useful to study the minimal low energy realization, the Minimal Supersymmetric Standard Model (MSSM).

MSSM Higgs Sector

- Due to the need for anomaly cancellation, the MSSM contains two Higgs doublets, H_1, H_2 .
- There are five physical Higgs bosons: h, H, A, H^\pm
- At tree level, Higgs sector described by two parameters, typically chosen to be $\tan\beta = v_1/v_2$ and M_A
- We will consider the scenario where the lightest Higgs boson, h , is identified as the Higgs discovered at $M_h = 125.5$ GeV.

MSSM Higgs Sector

- Due to the need for anomaly cancellation, the MSSM contains two Higgs doublets, H_1, H_2 .
- There are five physical Higgs bosons: h, H, A, H^\pm
- At tree level, Higgs sector described by two parameters, typically chosen to be $\tan\beta = v_1/v_2$ and M_A
- We will consider the scenario where the lightest Higgs boson, h , is identified as the Higgs discovered at $M_h = 125.5$ GeV.
- At tree level $M_h \leq M_Z$. Fortunately, stop loops provide considerable correction to the Higgs mass.
- The leading component to the Higgs mass from stop loops:

$$\varepsilon = \frac{3\bar{m}_t^4}{2\pi^2 v^2 \sin^2\beta} \left[\log \frac{M_S^2}{\bar{m}_t^2} + \frac{X_t^2}{M_S^2} \left(1 - \frac{X_t^2}{12M_S^2} \right) \right]$$

- Stop mixing parameter: $X_t = A_t - \mu \cot\beta$
- $M_S = \sqrt{\bar{m}_{\tilde{t}_1} \bar{m}_{\tilde{t}_2}}$ is the geometric mean of the stop masses.
- This correction is maximized when $X_t = \sqrt{6}M_S$.

Benchmarks

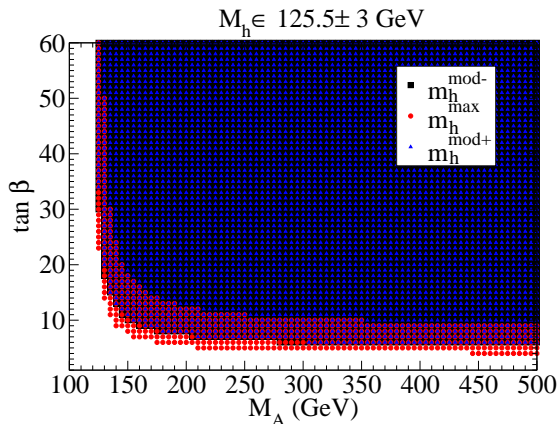
- Will consider benchmark points from [Carena, Heinemeyer, Stål, Wagner, Weiglein, 1302.7033](#).
- m_h^{\max} scenario chosen to maximize one-loop stop correction to Higgs mass
[Carena, Heinemeyer, Wagner, Weiglein, hep-ph/0202167](#)
 - $X_t^{\text{OS}} = 2M_{\text{SUSY}}$ (on-shell calculation)
 - $X_t^{\text{MS}} = \sqrt{6}M_{\text{SUSY}}$ ($\overline{\text{MS}}$ calculation)
- However, m_h^{\max} does its job too well... much of $\tan\beta - M_A$ region produces too large a Higgs mass.
- Slightly decrease stop mixing parameter.
- $m_h^{\text{mod}+}$:
 - $X_t^{\text{OS}} = 1.5M_{\text{SUSY}}$
 - $X_t^{\text{MS}} = 1.6M_{\text{SUSY}}$
- $m_h^{\text{mod}-}$:
 - $X_t^{\text{OS}} = -1.9M_{\text{SUSY}}$
 - $X_t^{\text{MS}} = -2.2M_{\text{SUSY}}$
- Third generation squark masses: $M_{\text{SUSY}} \equiv M_{\tilde{t}_L} = M_{\tilde{b}_L} = M_{\tilde{t}_R} = M_{\tilde{b}_R}$

Benchmarks

- Common parameters [Carena, Heinemeyer, Stål, Wagner, Weiglen, 1302.7033](#):

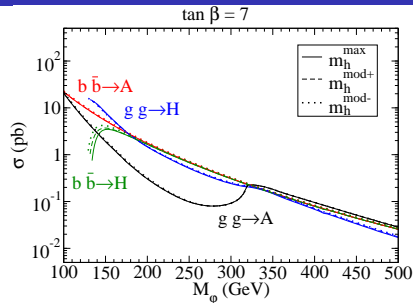
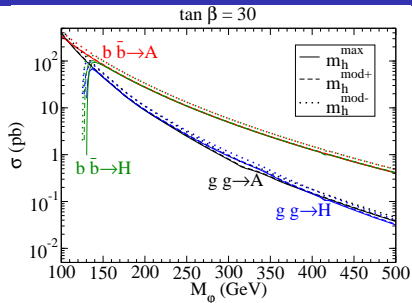
- $M_{\tilde{q}_{1,2}} = 3M_{\tilde{l}_{1,2}} = 1500 \text{ GeV}$
- $A_f = 0, \quad (f = c, s, u, d, \mu, e)$
- $m_t = 173.2 \text{ GeV}$
- $M_1 = \frac{5}{3} \frac{s_w^2}{c_w^2} M_2$
- $\mu = M_2 = 200 \text{ GeV}$
- $m_{\tilde{g}} = \frac{3}{2} M_{\tilde{l}_3} = 1500 \text{ GeV}$
- $A_b = A_\tau = A_t$

M_h range from benchmarks

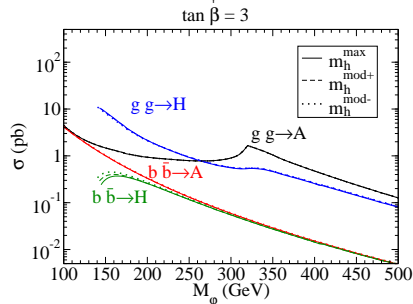


- Band in $\tan\beta - M_A$ plane for which M_h inside $125.5 \pm 3 \text{ GeV}$ range [Carena, Heinemeyer, Stål, Wagner, Weiglen, 1302.7033](#).
- Decoupling regime $M_A \gg M_Z$: $M_h^2 = M_Z^2 \cos^2 2\beta + \varepsilon \sin^2 \beta$
- FeynHiggs [Heinemeyer, Hollik, Weiglein, hep-ph/9812320](#) was used to produce this plot, and it is used extensively in the rest of the presentation.

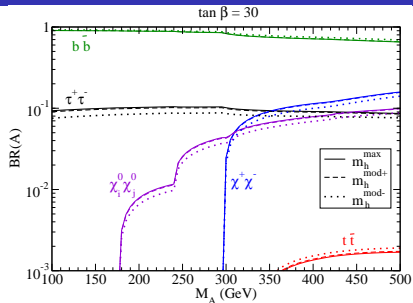
Production at 7 TeV LHC



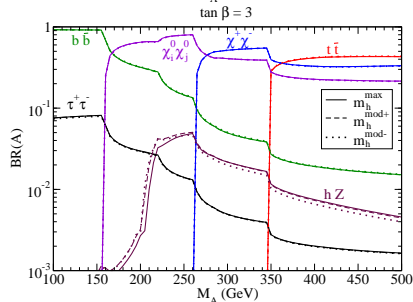
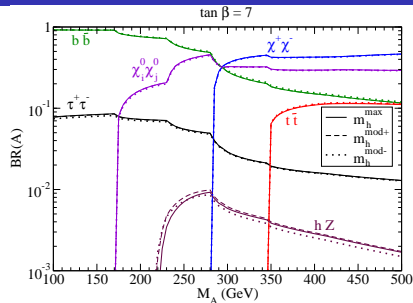
- Kink at $M_\phi \sim 320$ GeV from imaginary part in top loop.
- The $b\bar{b}$ cross section has been rescaled from the SM value produced by bbh@nnlo [Harlander, Kilgore, hep-ph/0304035](http://arxiv.org/abs/hep-ph/0304035), and the gg cross section is rescaled from the values given by the LHC Higgs Cross Section Working Group.



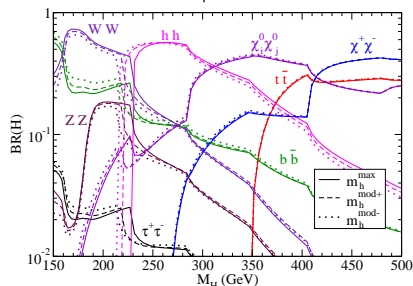
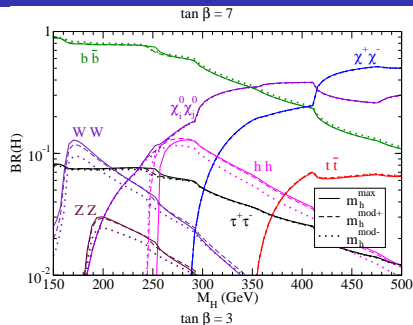
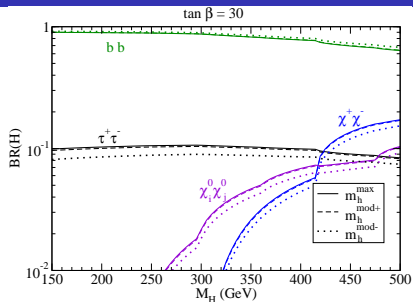
Decays of Pseudoscalar



- At high $\tan\beta$, $b\bar{b}$ and $\tau^+\tau^-$ dominate.
- At lower $\tan\beta$, neutralino and chargino decays much more important.

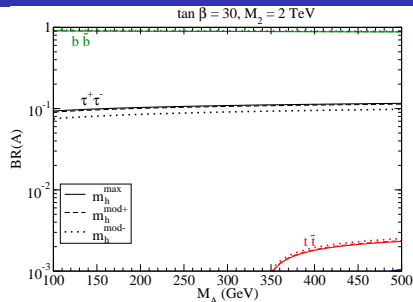
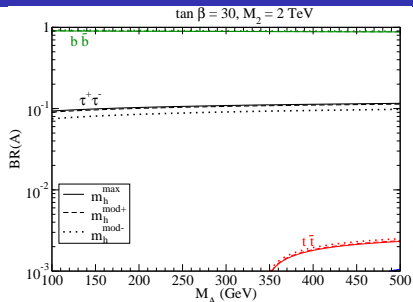


Decays of Heavy Scalar

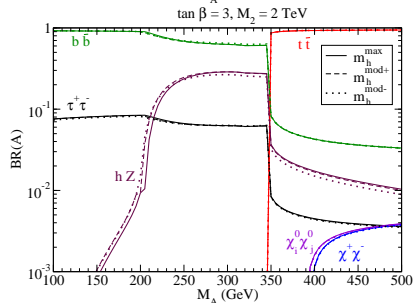


- At high tan β , $b\bar{b}$ and $\tau^+\tau^-$ dominate.
- More complicated at lower tan β , depends on mass range.

Decays of Pseudoscalar With $M_2 = 2$ TeV



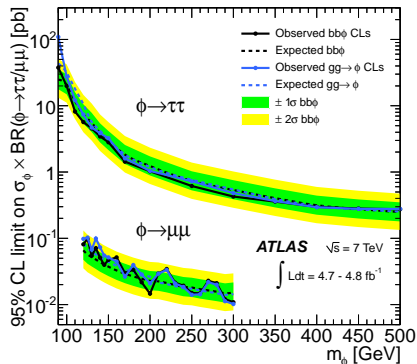
- $M_2 = 2$ TeV alters neutralino mass spectrum and branching ratios of heavy Higgs.
- $A \rightarrow \tau^+ \tau^-$ more viable for all tan β .
- $A \rightarrow hZ$ more important at low tan β
- Extends range for $H \rightarrow hh$



The Wedge

- Different search modes at LHC relevant for high and low $\tan\beta$.
- Since production cross sections and branching ratios similar in three benchmarks, will focus on m_h^{\max}
- Focus on $\Phi \rightarrow \tau^+ \tau^-$ searches, extending current LHC limits.
- Also investigate $A \rightarrow hZ$ to open lower $\tan\beta$ regime.
- All cross section bounds assumed to scale as $1/\sqrt{\text{Luminosity}}$

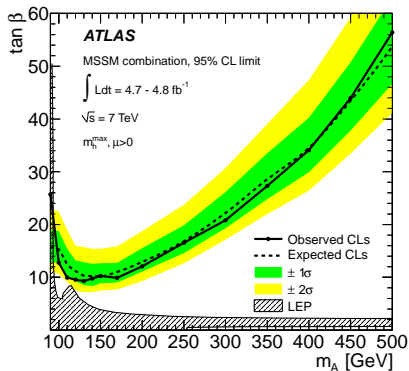
Current bounds on $\sigma(pp \rightarrow \phi) \times \text{BR}(\phi \rightarrow \tau^+\tau^-)$



JHEP02(2013)095

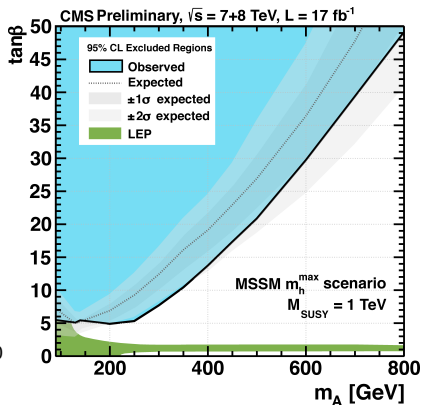
1211.6956

Current bounds



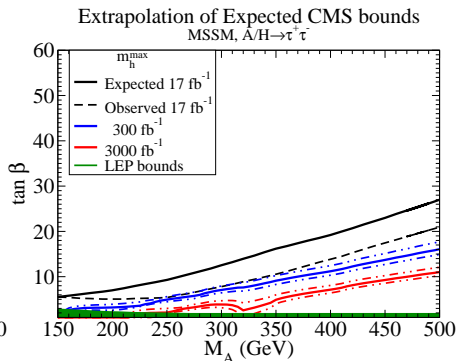
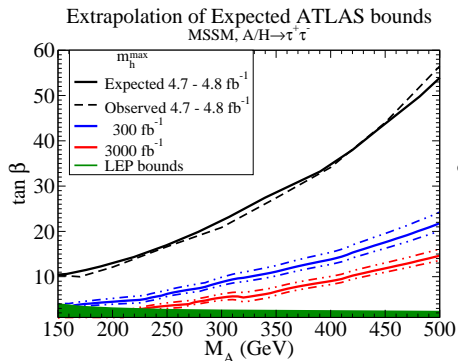
JHEP02(2013)095

1211.6956



CMS-PAS-HIG-12-050

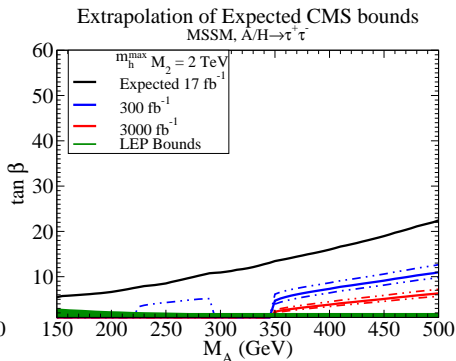
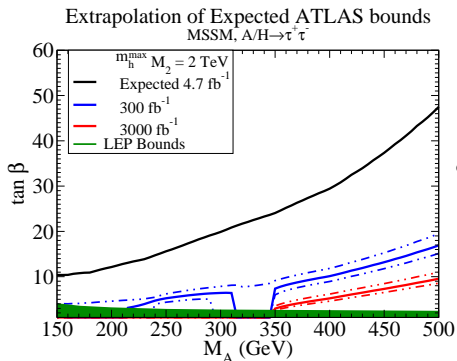
Projection of bounds to 300 and 3000 fb⁻¹



- Included conservative error bands of $\Delta\sigma(\Phi) \times \text{BR}(\Phi \rightarrow \tau^+\tau^-) = \pm 25\%$

Baglio, Djouadi, 1012.0530; Djouadi, Quevillon, 1304.1787

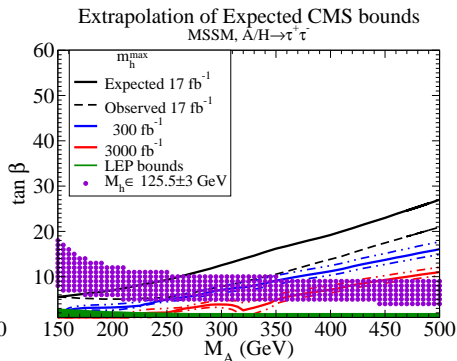
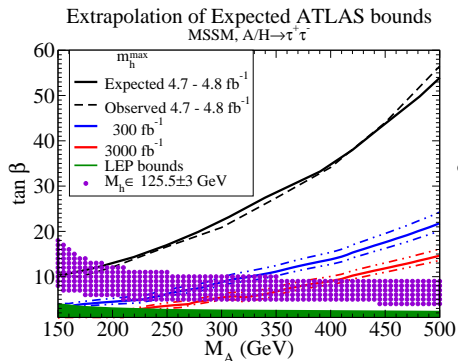
Projection of bounds to 300 and 3000 fb⁻¹ with $M_2 = 2$ TeV



- Included conservative error bands of $\Delta\sigma(\Phi) \times \text{BR}(\Phi \rightarrow \tau^+\tau^-) = \pm 25\%$

Baglio, Djouadi, 1012.0530; Djouadi, Quevillon, 1304.1787

Projection of bounds to 300 and 3000 fb⁻¹



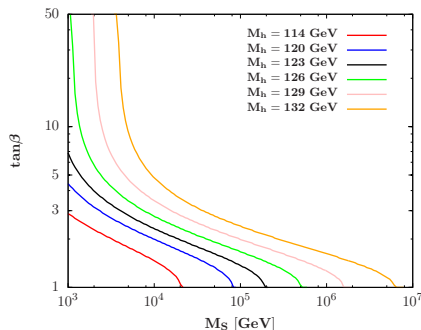
Bounds on M_A in m_h^{\max} scenario:

Luminosity	ATLAS	CMS
300 fb ⁻¹	210-240 GeV	230-260 GeV
3000 fb ⁻¹	260-290 GeV	290-360 GeV

Low tan β

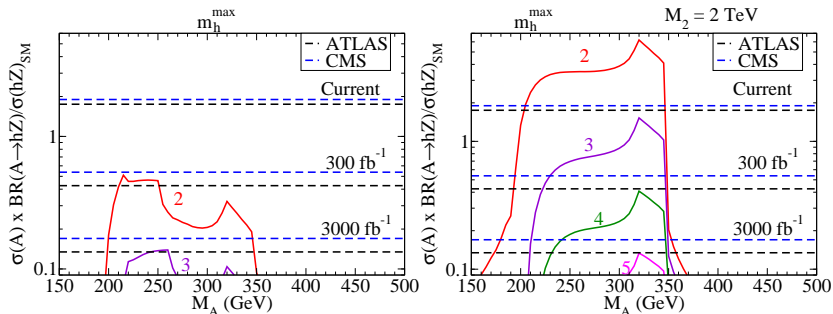
- Although benchmarks considered here do not reproduce correct Higgs mass at low tan β , it is possible to reproduce Higgs mass by increasing M_{SUSY} considerably

Djouadi, Quevillon, 1304.1787



- Search strategy largely depends on neutralino and chargino spectrum.
- Additional signals of additional interest are $A \rightarrow hZ$, $H \rightarrow hh$, and $H \rightarrow WW$.

Projection for $A \rightarrow hZ$



- For a 125 GeV Higgs, current bounds on associated productions are
 - ATLAS: $\sigma(VH) < 1.75 \times \sigma(VH)_{\text{SM}}$ with 17.7 fb⁻¹ [ATLAS-CONF-2012-161](#)
 - CMS: $\sigma(VH) < 1.9 \times \sigma(VH)_{\text{SM}}$ with 24 fb⁻¹ [CMS-PAS-HIG-13-012](#)
- No dedicated resonance search, so apply these upper bounds on $\sigma(A \rightarrow hZ)$
- With $200 \text{ GeV} \lesssim M_A \lesssim 250 \text{ GeV}$, sensitive to $\tan\beta \lesssim 2$ with 300 fb⁻¹ and $\tan\beta \lesssim 3$ with 3000 fb⁻¹

ILC

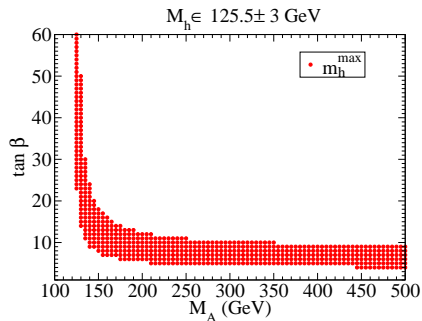
- Search for Higgs-strahlung, $e^+e^- \rightarrow Z^* \rightarrow Z\Phi$, or pair production $e^+e^- \rightarrow Z^* \rightarrow \Phi A$
- For $M_A \gtrsim 200$ GeV
 - $Z - Z - H$ and $Z - h - A$ couplings suppressed
 - $Z - Z - h$ and $Z - H - A$ coupling nearly maximal.
- For direct production, if have sufficient energy can have unsuppressed production of $e^+e^- \rightarrow Z^* \rightarrow HA$.
- Substantial decay channels such as $A/H \rightarrow b\bar{b}$, $A/H \rightarrow t\bar{t}$, $H \rightarrow hh \rightarrow 4b$, and $A \rightarrow Zh \rightarrow Zb\bar{b}$ easier to detect than LHC.
- Not as much uncertainty as LHC in production cross section.
- Should be able to close out wedge without as much ambiguity about theory errors.

Conclusions

- Higgs discovery just beginning of the exploration of EWSB.
- One popular alternative to the SM EWSB scenario is the MSSM.
- Contains 5 Higgs bosons: h, H, A, H^\pm
- We used current constraints from ATLAS and CMS on $\sigma(pp \rightarrow H/A \rightarrow \tau^+\tau^-)$ to explore high $\tan\beta$ regime and extrapolate bounds on the $M_A - \tan\beta$ plane at 300 and 3000 fb^{-1}
- Found with 300 fb^{-1}
 - ATLAS can exclude m_h^{max} benchmark for $M_A \lesssim 210 - 240 \text{ GeV}$
 - CMS can exclude m_h^{max} benchmark for $M_A \lesssim 230 - 260 \text{ GeV}$
- Found with 3000 fb^{-1}
 - ATLAS can exclude m_h^{max} benchmark for $M_A \lesssim 260 - 290 \text{ GeV}$
 - CMS can exclude m_h^{max} benchmark for $M_A \lesssim 290 - 360 \text{ GeV}$
- To examine low $\tan\beta$ regime, also placed bounds on $A \rightarrow Zh$ production via the bound on SM Vh production with $M_h = 125 \text{ GeV}$.
- In the m_h^{max} scenario with $200 \text{ GeV} \lesssim M_A \lesssim 250 \text{ GeV}$, LHC is sensitive to $\tan\beta \lesssim 2$ with 300 fb^{-1} and $\tan\beta \lesssim 3$ with 3000 fb^{-1}
- Increasing M_2 can substantially strengthen all bounds.

EXTRA SLIDES

M_h in the m_h^{\max} scheme



- Now know $M_h \sim 125.5$
- Band in $\tan \beta - M_A$ plane for which M_h inside 125.5 ± 3 GeV range.