

# Beam Dynamics for Crab Cavities in the APS Upgrade

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#### **Outline**

- Why mention APS Upgrade at this workshop?
- Deflecting cavity scheme description
- Challenges in beam dynamics of crab cavities

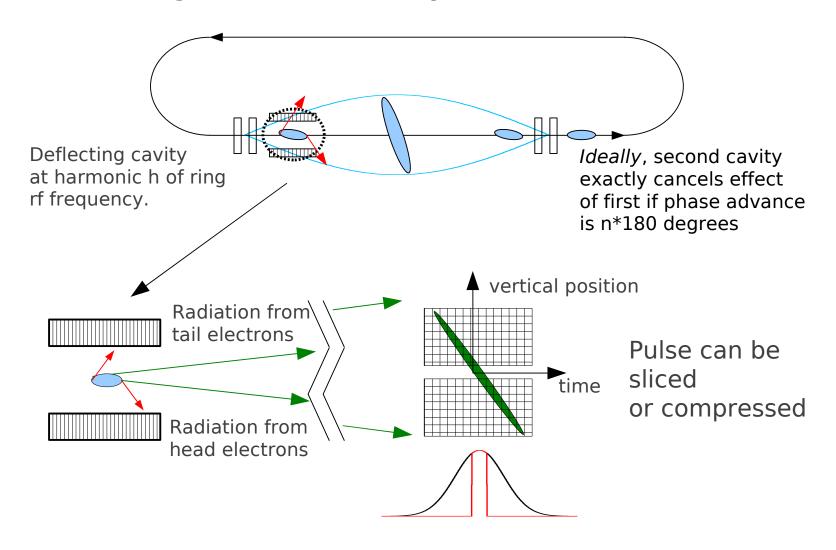
## Applications of deflecting cavities in storage rings

- Two major applications for deflecting cavities:
  - Restoring head-on collisions in crab crossing in colliders
    - Suppresses synchro-betatron resonances excited by crab crossing
  - Generating short X-ray pulses in light sources
    - Allows to take advantage of small vertical beam size to generate temporally short pulses
- Some beam dynamics issues are similar:
  - Additional impedance
  - Cavity generated beam noise
- Some are different
  - Beam-beam related effects in colliders
  - Coupling increase and related nonlinear dynamics complications in light sources
- Major difference is deflection plane: vertical for light sources and horizontal for colliders





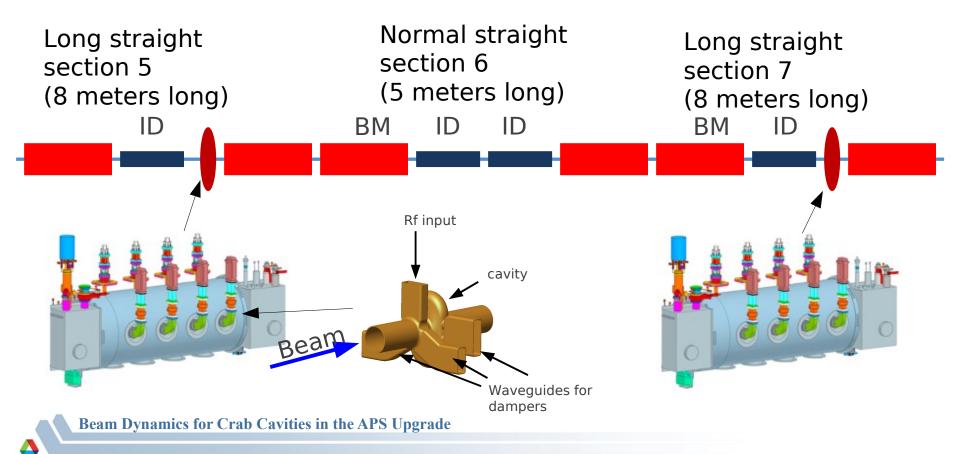
# Deflecting cavities concept<sup>1</sup>



<sup>&</sup>lt;sup>1</sup>A. Zholents et al., NIM A 425, 385 (1999).

## Short-Pulse X-ray source

- Few picosecond x-ray pulses by applying a local (y,y')-z correlation ("chirp") bump to stored beam
- Superconducting radio-frequency deflecting cavities operated in continuous-wave mode
- Up to 4 ID and 2 BM beam lines, operation in 24 singlets mode

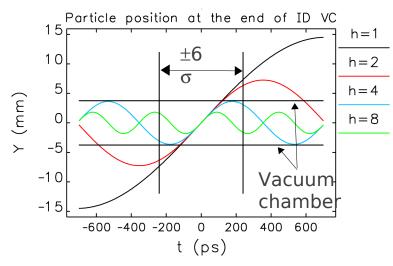


## Choice of parameters

 To obtain rms pulse length of 1 ps (2 ps FWHM), the deflecting voltage amplitude times harmonic has to be (assuming no changes to SR optics):

$$h \ V \approx \frac{E}{\sigma_s \ f_{rf}} \sqrt{\frac{\beta_{id}}{\beta_{rf}}} \sqrt{\frac{\epsilon_y}{\beta_{id}} + \frac{\lambda_{rad}}{2 \ L_u}} \approx 15 \ MV$$

- Cavities will share straight sections with insertion devices which means there will be narrow-gap vacuum chamber
- Large vertical beam size inside narrow-gap VC puts lower limit on frequency due to lifetime, h > 4
- Chosen deflecting voltage parameters: next to them



#### Effect of cavities on the beam

- Less than total kick cancellation at the second cavity could lead to beam emittance increase and to orbit distortion
- Nonlinear beam dynamics is affected
- Cavities introduce additional impedance, and therefore can affect single-bunch and multi-bunch instabilities

#### Effect on emittance

- In a real machine, many effects could lead to emittance degradation
  - Various errors and imperfections are first things coming to mind
- However, even in a perfect machine the emittance can increase many ways
  - Path length dependence on the particle energy leads to incomplete kick canceling in the second cavity
  - Betatron phase advance dependence on energy (chromaticity) leads to closed bump condition breaking
  - Sextupoles between cavities introduce nonlinearities that generate betatron phase advance dependence on amplitude and linear coupling between horizontal and vertical planes

#### Momentum compaction

- This effect comes from the path length difference between the cavities for particles with different energy
- This effect is present even if there are no errors and nonlinearities
- For a particle with energy deviation  $\delta_i$ , the time of flight differential  $\Delta t_i = \alpha_c \delta_i T_0$
- Additional kick after the second cavity is  $\Delta y_i' = \frac{-V \omega \Delta t_i}{E}$

which gives emittance increase of

$$\frac{\Delta \epsilon_{y}}{\epsilon_{v}} = \frac{\sqrt{\sigma_{y'}^{2} + \sigma_{\Delta y'}^{2}}}{\sigma_{v'}} - 1$$

 For APS case, it gives about 0.3% increase of emittance in a single turn which gives negligible effect on overall emittance increase



## **Chromaticity**

- The second cavity is placed at  $n\pi$  phase advance to cancel the kick of the first cavity
- If there is non-zero chromaticity  $\xi_y$  between the cavities, the phase advance of a particle with  $\delta_i$  is changed by  $-2\pi\xi_y\delta_i$  which leads to a particle position change at the second cavity  $y_2 = \beta \ y'_1 \sin(2\pi \ \xi_y \delta_i)$
- The rms value of the residual amplitude is

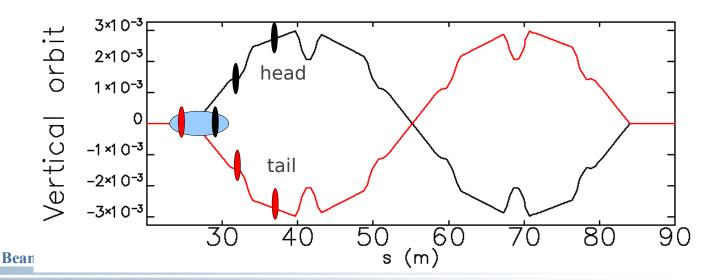
$$\sigma_{y_2} = 2\pi \xi_y \beta \frac{V \omega}{E} \sigma_{\delta} \sigma_t$$

- For APS parameters with uncompensated chromaticity (no sextupoles in these two sectors), this works out to be over 50% of the nominal vertical beam size of 11 μm
- To avoid this emittance increase, sextupoles are required between the cavities

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#### Sextupole nonlinearities

- Introduces amplitude-dependent focusing
  - for particles going off-axis the kick cancellation at the second cavity is not perfect
- Introduces transverse coupling
  - deflecting cavities generate large vertical trajectories in sextupoles
  - Vertical trajectory in sextupoles creates coupling between large horizontal and small vertical emittances



#### Beam dynamics simulation methods

- We use tracking to simulate beam dynamics
- We use parallel elegant<sup>1</sup> typically utilizing 10-50 CPU cores
- Accelerating cavities are required to simulate synchrotron motion
- Synchrotron radiation is essential: to damp initial cavity effects
  - Tracking is done for 10k turns about 4 damping times
- Deflecting cavity is simulated as TM-like mode, deflection is radius independent resulting from combination of TM- and TE-like field<sup>2</sup>

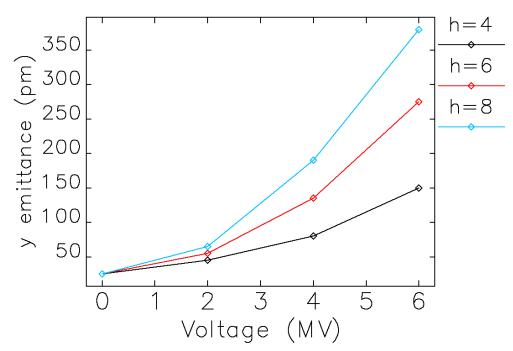
<sup>&</sup>lt;sup>2</sup>M. Nagl, tesla.desy.de/fla/publications/talks/seminar/FLA-seminar\_230904.pdf



<sup>&</sup>lt;sup>1</sup>Y. Wang et al., AIP 877, 241 (2006).

# Initial results of the deflecting cavity application

 Right away, we have found significant blow-up of vertical emittance due to increased coupling. This can be fixed by adjusting sextupole gradients in the two sectors, but creates a major problem

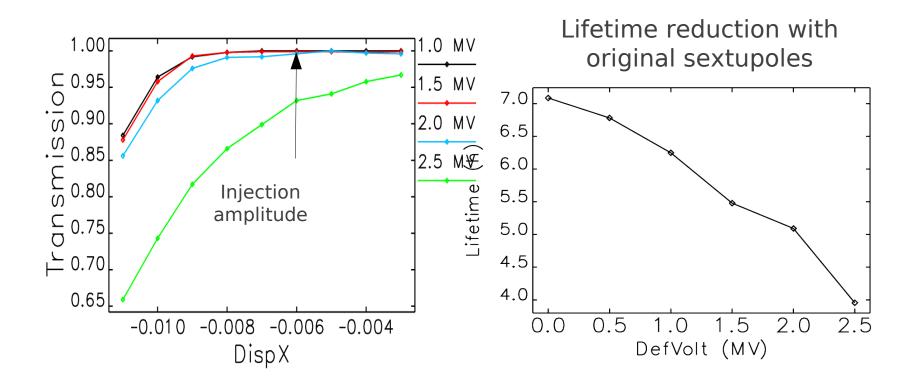


### Nonlinear dynamics challenge in general

- Light sources tend to minimize their beam emittance to the level where Dynamic Aperture (DA) and lifetime are barely enough for operation
- Many sextupole families are utilized to achieve workable DA and lifetime, i.e. for symmetric optics without deflecting cavities.
- A local sextupole adjustment that minimizes vertical emittance growth will violates the earlier sextupole optimization of the whole ring
- Even small reduction of DA and lifetime can be crucial
- Further investigations requires including the deflecting cavity effects on nonlinear dynamics
- The cavity effects are defined by large vertical trajectories between deflecting cavities:
  - Physical acceptance is decreased
  - Additional linear and nonlinear coupling is introduced



### Injection and lifetime with deflecting cavities



Reduction due to vertical physical aperture

Reduction due to a skew sextupole resonance with original sextupole distribution



### Sextupole optimization with deflecting cavities

- Sextupoles between the cavities are needed to compensate for natural chromaticity
- At the same time large vertical trajectories in sextupoles lead to vertical emittance increase and nonlinear dynamics deterioration
- Optimization of sextupoles between cavities allows to solve each problem separately
- Now we need to satisfy everything at the same time
- The best way to do it is to use multi-objective optimization, and do it as a part of overall lattice design

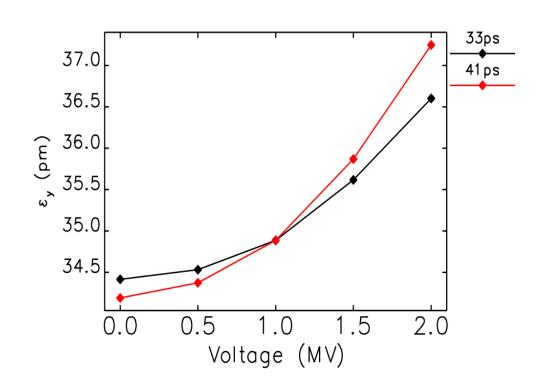
## Sextupole optimization (2)

- The optimization is done using a genetic optimizer
- Every optimizer evaluation consists of
  - Linear optics design (if required)
  - Interior sextupoles optimization for vertical emittance blowup minimization
  - Exterior sextupole optimization for DA/LMA
- The penalty functions are vertical emittance increase, DA area, and lifetime
- It is very CPU-hungry process, it requires parallel computations, but it gives satisfactory results
  - We are able to achieve satisfactory dynamic aperture and lifetime without any increase of vertical emittance
- DA/LMA evaluation with cavities on is not included in optimization yet



# Vertical emittance after global sextupole optimization

- Particles are tracked for 10k turns (several damping times)
- Sextupoles were optimized for extreme case of 50-ps-long bunch and 4MV
- Vertical emittance growth below 10% is achieved
- Two bunch lengths corresponding to two different operating conditions are shown



#### Deflecting voltage tolerances

- The voltage could vary in amplitude and phase, and variations at both cavities could follow each other (common-mode) or not (differential-mode)
- Common-mode variations affect the beam only between the cavities
  - Important for colliders
  - Not as important for light sources because the beam size between cavities is greatly increased
- Differential-mode variations affect the beam everywhere
  - Give very tight tolerances for light sources due to small vertical beam sizes
- Will not talk about common-mode tolerances



#### Differential mode tolerances

• When the voltage waveform in the second cavity does not exactly follow the first cavity, the resulting effect of two cavities on the beam is non-zero:

$$V\sin(\omega t) - (V + \Delta V)\sin(\omega t + \Delta \phi) \approx V\cos(\omega t)\sin(\Delta \phi) - \Delta V\sin(\omega \phi)$$

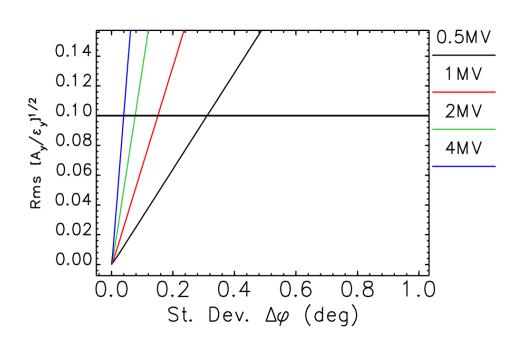
- The first term provides a net orbit kick because its value is non-zero at the center of the bunch (t=0)
- The second term generates beam tilt outside of the deflecting cavities and affects projected beam sizes
- The effect can be treated as a single source orbit distortion and a single deflecting cavity with voltage  $\Delta V$ .

#### **Tolerances: Orbit**

- Want to keep orbit variation under some fraction of nominal beam emittance (total APS beam motion budget in terms of beam motion invariant is 1% of beam emittance)
- Using APS parameters, we get:

 $\Delta \phi$  < 0.08 deg or 80 fs

 This is quite a tight tolerance for rf phase



#### **Tolerances: Emittance**

- Various errors affect the outside beam sizes
  - Differential deflecting voltage
  - Vertical betatron phase advance not equal to  $N^*\pi$
  - Beta function mismatch
  - Cavity and magnet roll
- All these errors except differential deflecting voltage are static
  - Beta function error can be compensated by changing relative voltage of second cavity
  - Phase advance error can be compensated by changing relative voltage of first and second sets of cells in second cavity
  - Cavity roll is found to be a weak effect<sup>1</sup>
  - Magnet roll can be corrected with additional skew quadrupoles
- We will only look at effect of differential voltage errors

<sup>1</sup>M. Borland, PRSTAB 8, 074001 (2005).

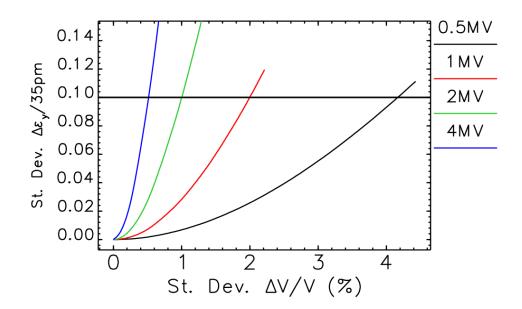


## Tolerances: Emittance (3)

• If we require that the beam size increase does not exceed 10% of the total beam size, for APS parameters we get:

$$\frac{\Delta V}{V} < 0.01$$

Realistic tracking simulations of the emittance sensitivity to the voltage errors show good agreement:



#### Collective effects

- Can be separated into short- and long-range effects
- Long-range effects generate multi-bunch instabilities
- Short-range wake fields limit single bunch current

## Cavity impedance requirements

 Initial estimates of largest allowable resonator impedances (assuming high Qs) for bunch train stability were given to rf designers

Shunt impedance	Limit
Longitudinal	
$(R_s f_{HOM})$ for one monopole HOM/LOM	$0.44~\mathrm{M}\Omega\text{-GHz}$
$R_s$ for one monopole HOM/LOM at 2 GHz	$0.22~\mathrm{M}\Omega$
Transverse	
$R_t$ for one x-plane HOM/LOM	$1.3~\mathrm{M}\Omega/\mathrm{m}$
$R_t$ for one y-plane HOM/LOM	$3.9~\mathrm{M}\Omega/\mathrm{m}$

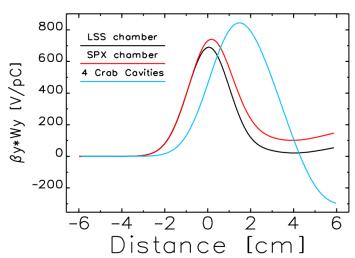
- Dampers were designed that produced very low Qs and shunt impedances
- Monte Carlo simulations of the damped-Q HOM resonators (with randomized frequency) verifies stable beam conditions

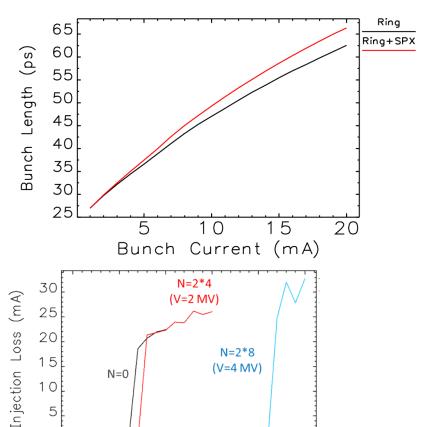
## Collective effects (2)

Short-range wake fields could limit single bunch current

Additional impedance comes from cavities and vacuum

chamber transitions





25

Stored Current (mA)

30

40

#### Cavity alignment requirements

- Cavity misalignment has several effects:
  - Beam-induced power generation due to transverse misalignment could damage the rf components
  - Beam arrival jitter combined with transverse offset leads to rf phase noise
  - Cavity roll can affect beam emittance
- Beam orbit can only be steered through "average" cavity center but cavity-to-cavity misalignment cannot be compensated
- Realistically achievable alignment is taken into account

	Cryomodule	Cavity inside
	alignment	cryomodule
$\Delta X$	$\pm 500~\mu\mathrm{m}$	$\pm 500~\mu\mathrm{m}$
$\Delta Y$	$\pm 200~\mu\mathrm{m}$	$\pm 200~\mu\mathrm{m}$
$\Delta Z$	$\pm 1000~\mu\mathrm{m}$	$\pm 1000~\mu\mathrm{m}$
Yaw	$\pm 10~\mathrm{mrad}$	$\pm 10~\mathrm{mrad}$
Pitch	$\pm 10~\mathrm{mrad}$	$\pm 10~\mathrm{mrad}$
Roll	$\pm 10~\mathrm{mrad}$	$\pm 10~\mathrm{mrad}$

Here X is horizontal, Y is vertical, and Z is longitudinal directions.

#### Conclusions

- Deflecting cavities could affect single particle beam dynamics through nonlinearities on large trajectories between the cavities
  - Sextupoles and nonlinearities of the deflecting fields could limit momentum and dynamics aperture
  - Sextupoles could greatly increase transverse coupling
  - Sextupole distribution solved by genetic algorithm and massive tracking on computer cluster
- Cavities could increase beam emittance and generate beam motion through rf noise in cavities
  - Leads to engineering tolerances
- Cavities introduce additional impedance, and therefore can affect single-bunch and multi-bunch instabilities
  - Approach the same way as other rf cavities, i.e. dampers, careful design of tapers, feedback systems.

