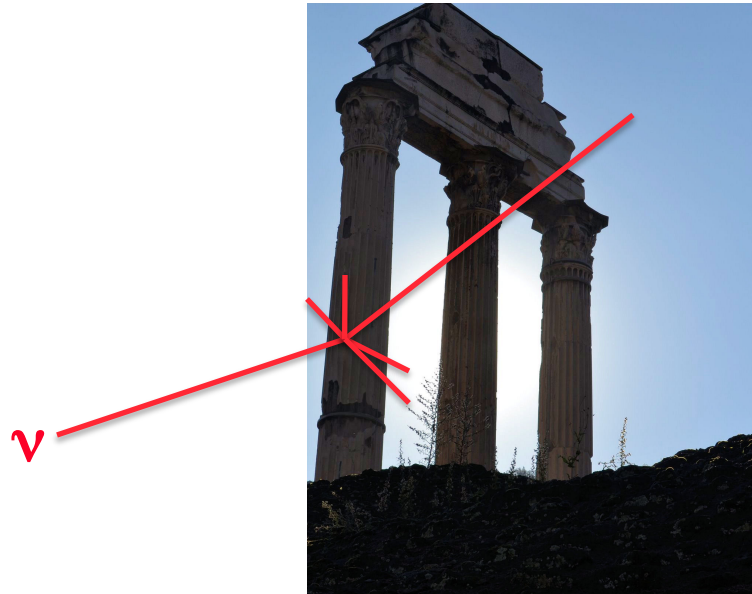


Neutrino – Nucleus Scattering Physics with nuSTORM

What can a dedicated nuSTORM neutrino-nucleus scattering physics program deliver?



Jorge G. Morfín
Fermilab

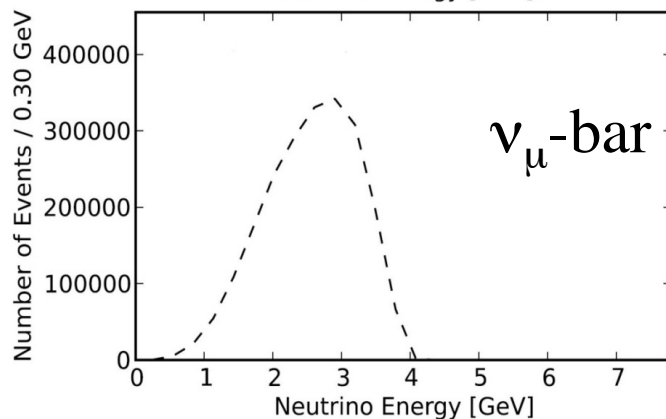
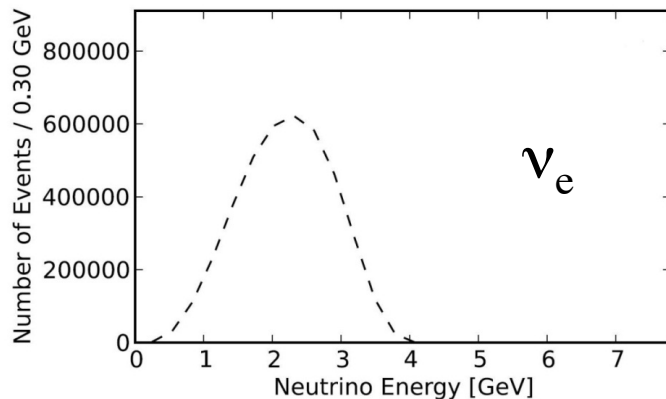
nuSTORM Workshop – Fermilab, November 2013

The nuSTORM Neutrino Beam: The Advantages



Mitchell YU

- ◆ The nuSTORM beam will provide a **very well-known** ($\delta \phi(E) \leq 1\%$) beam of ν and $\bar{\nu}$. Intensity of ν_μ good **not great** for 10 ton detector.
- ◆ A high-intensity source of **ν_e events** for experiments.



μ^+		μ^-	
Channel	N_{evts}	Channel	N_{evts}
$\bar{\nu}_\mu$ NC	844,793	$\bar{\nu}_e$ NC	709,576
ν_e NC	1,387,698	ν_μ NC	1,584,003
$\bar{\nu}_\mu$ CC	2,145,632	$\bar{\nu}_e$ CC	1,784,099
ν_e CC	3,960,421	ν_μ CC	4,626,480

event rates per 1E21 POT -
100 tons at 50m

3.8 GeV μ^+ stored, 226m straight, flux at 50m

nuSTORM Event Rates

- ◆ A beam of 700 kW yields order 6×10^{20} POT/year. The 1100 kW for LBNE could yield close to 10^{21} POT/year (depends on E of LBNE proton beam). **WE DON'T GET ALL THE PROTONS!**

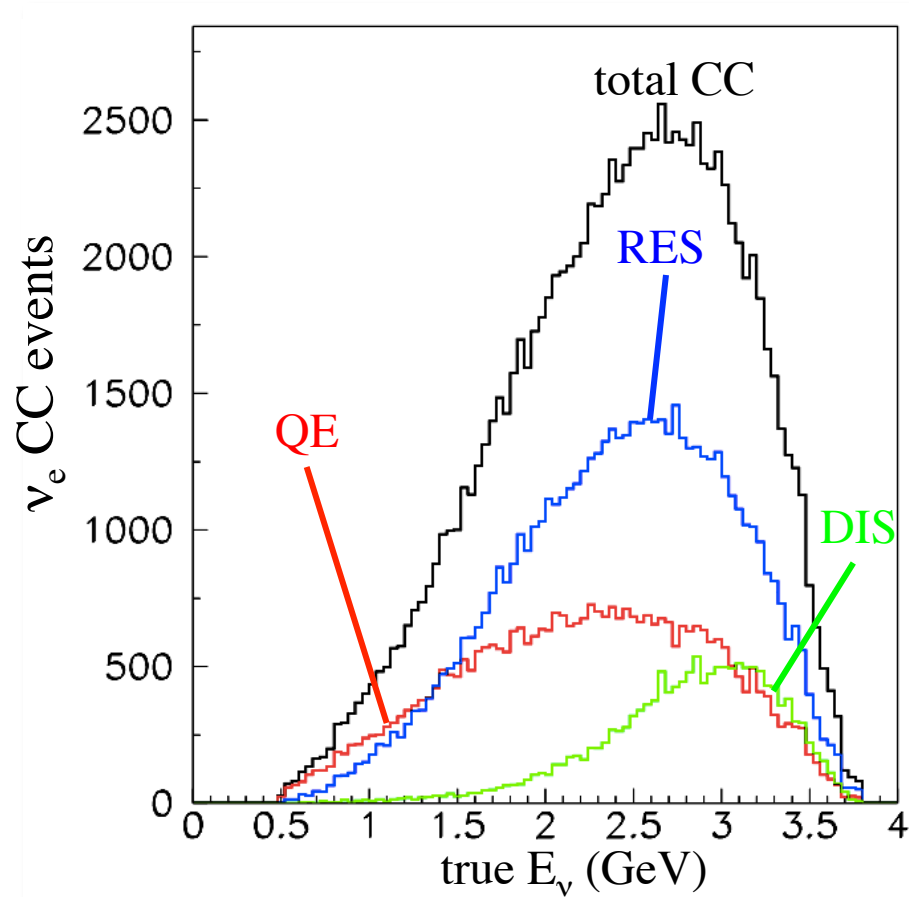
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ν_e CC	3,960,421	ν_μ CC	4,626,480

event rates per $1E21$ POT -
100 tons at 50m

- ◆ It would be preferable to have a near detector of order 30-50 ton fiducial volume.

ν Produced Events in a nuSTORM Near Detector

- ◆ ν produced by 3.8 GeV μ^+ beam.



**out of the CC
modes:**

- * 56% resonant
- * 32% QE
- * 12% DIS

- ◆ For $\bar{\nu}$ sample, 52% resonant, 40% QE, 8% DIS)

Why is Neutrino Nucleus Scattering Important?

What do we observe in our detectors?

- ◆ The events we observe in our detectors are convolutions of:

$$Y_{c\text{-like}}(E) \propto \phi(E' \geq E) \otimes \sigma_{c,d,e..}(E' \geq E) \otimes \text{Nuc}_{c,d,e.. \rightarrow c}(E' \geq E)$$

- ◆ $\phi(E)$ is the energy dependent neutrino flux that enters the detector. Currently, with traditional meson-decay-source neutrino beams, $\phi(E) \approx 10\%$ absolute and $\approx 7\%$ energy bin-to-bin accuracy. **Significant contribution to systematics.**
- ◆ $\sigma_{c,d,e..}(E' \geq E)$ is the measured or the Monte Carlo (model) energy dependent neutrino cross section off a **nucleon within a nucleus.**
- ◆ **$\text{Nuc}_{c,d,e.. \rightarrow c}(E' \geq E)$ – Nuclear Effects**
 - ◆ **Nuclear Effects** – a migration matrix that mixes produced/observed channels and energy
 - ◆ In general the interaction of a neutrino with energy E' creating initial channel d,e... can appear in our detector as energy E and channel c .
 - ◆ Particularly **fierce bias** when using the **QE hypothesis** to calculate E and Q^2 !
- ◆ $Y_{c\text{-like}}(E)$ is the event energy and channel / topology of the event observed in the detector. Appears to be channel c but may not have been channel c at interaction.

What do we observe in our detectors?

Further implications for Oscillation Experiments

- ◆ The events we observe in our detectors are convolutions of:

$$Y_{c\text{-like}}(E) \propto \phi(E' \geq E) \otimes \sigma_{c,d,e..}(E' \geq E) \otimes \text{Nuc}_{c,d,e.. \rightarrow c}(E' \geq E)$$

← effective $\sigma_c^A(E)$

- ◆ Experimentally, the convolution of initial cross section and nuclear effects are combined into an effective cross section $\sigma_c^A(E)$ that **depends on incoming neutrino energy spectrum and nuclear effects that populate the yield $Y_c^A(E)$.**
- ◆ This implies, for example, effective $\sigma_{\pi^+}^C(1 \text{ GeV})$ measured in the Booster beam **will be different** than the same effective $\sigma_{\pi^+}^C(1 \text{ GeV})$ observed in the higher energy NuMI beam due to, for example, more feed down from multi-pi events. **Can not simply plug in effective $\sigma_{\pi^+}^A$ from experiments in a different beam.**
- ◆ In a two-detector LBL oscillation experiment, neutrino flux entering the FD is different than the neutrino flux at the ND due to geometry and oscillations. **The $\sigma_c^A(E)$ effective that should be applied to expectations (Monte Carlo) at FD is NOT the same as that which we would measure at the ND. The ND results give us an excellent starting point for calculating the difference.**

How do we address this ν_μ scattering problem? What could a nuSTORM Scattering analysis add?

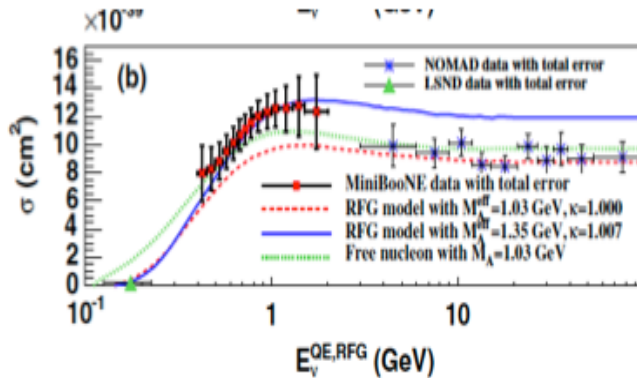
- ◆ Use the unique qualities of the nuSTORM beam meaning the fantastic knowledge of absolute and relative flux. And we can vary the energy distribution of this well-known beam.
- ◆ Combine with a **high-resolution near detector with multiple nuclear targets** to provide detailed studies of the final states including the vertex multiplicities and energy flow..
- ◆ **nuSTORM, providing a beam with knowledge of the flux to $\leq 1\%$, to such a near detector would provide an outstanding neutrino-nucleus scattering experiment addressing both electroweak and nuclear physics questions. It would allow us to measure, for example, $\sigma_{\pi^+}^A(E)$ for multiple A and various E.**

QE History

Very important for nuSTORM

- problem of low Q^2

- problem of axial mass



- problem of np-nh
(more complex nuclear effects:
SRC + MEC)

- M_A provides a convenient tool to describe exp. data in shape & normalization with Fermi Gas

Have they got the flux wrong?
nuSTORM could tell us? Distinction between Fermi Gas and Spectral Functions? With accurate $d\phi / dE$ – maybe nuSTORM could help.

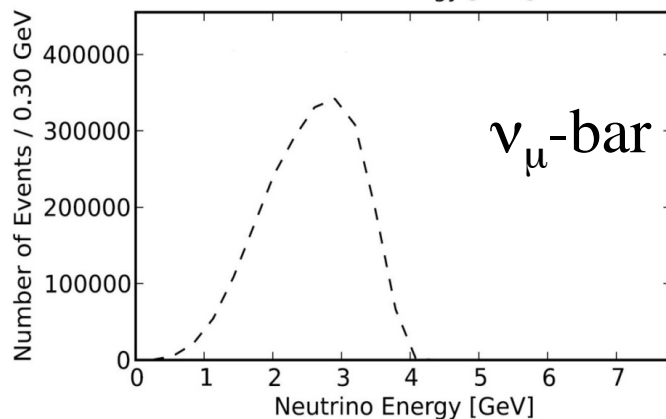
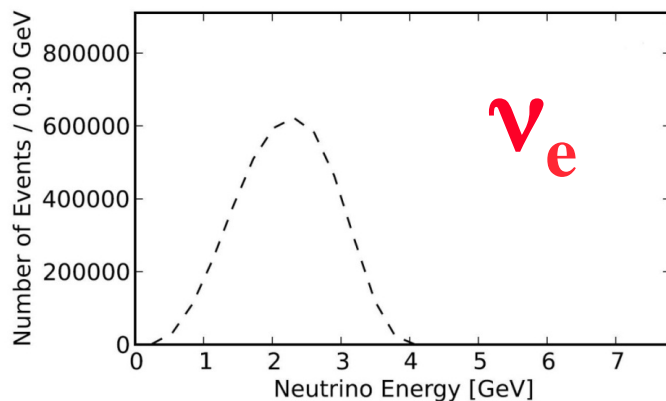
- However there is an alternative path involving a more sophisticated nuclear models. MEC implies extra tracks and energy at vertex,
nuSTORM could help. Time to retire impulse approximation with RFG!

The nuSTORM Neutrino Beam

$$\mu^+ \rightarrow \bar{\nu}_\mu + \nu_e + e^+$$

$$\mu^- \rightarrow \nu_\mu + \bar{\nu}_e + e^-$$

- ◆ The nuSTORM beam will provide a **very well-known** ($\delta \phi(E) \leq 1\%$) beam of ν and $\bar{\nu}$.
- ◆ **A high-intensity source of ν_e events for experiments.**



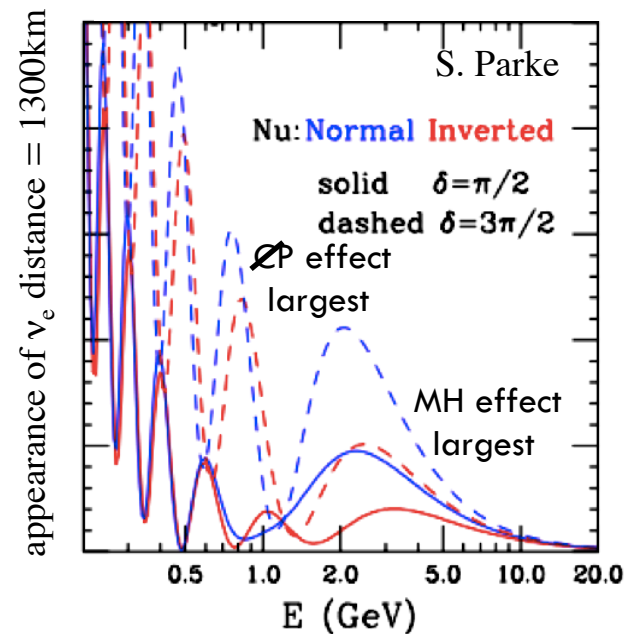
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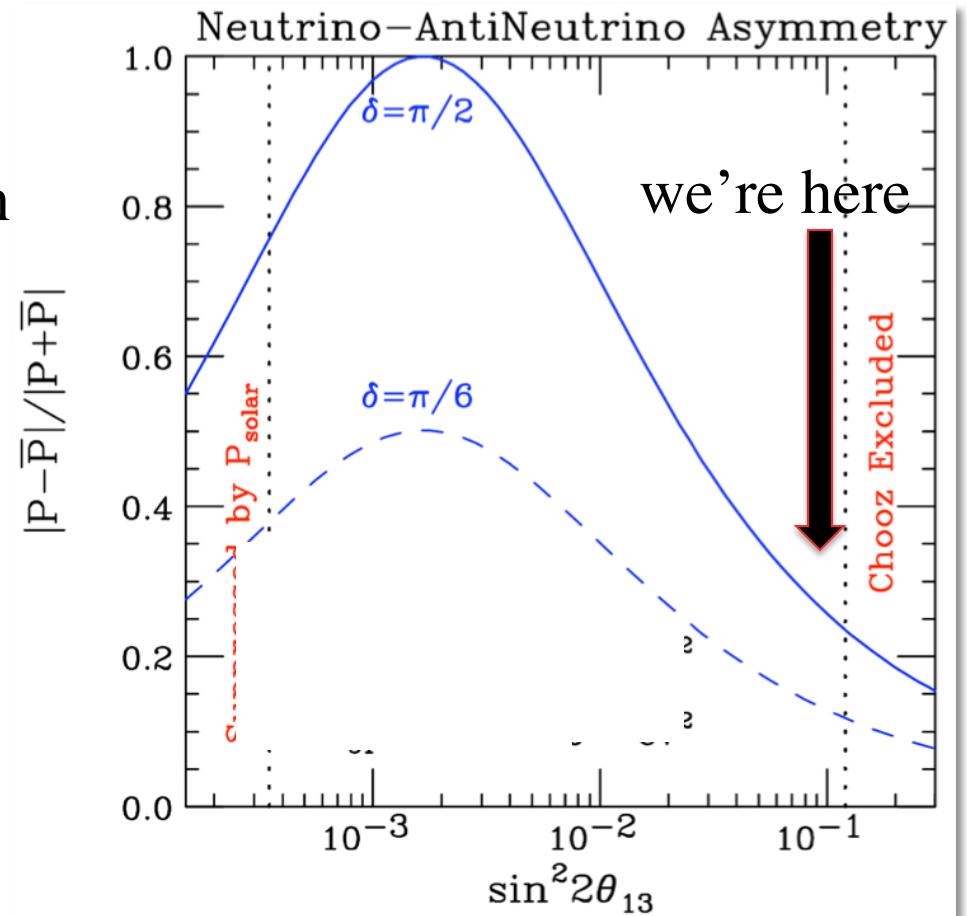
Why are ν_e Cross Sections Important?

- ◆ ν_e A – scattering results are interesting on their own.
- ◆ Recent determination of large θ_{13} has opened up possibilities of
 - ▼ Determining ν mass ordering.
 - ▼ Searching for CP-violation in the ν sector.
- ◆ To be sensitive to these effects, current/near-future long-baseline experiments will be looking for ν_μ to ν_e and $\bar{\nu}_\mu$ to $\bar{\nu}_e$ oscillations over a range of energies.
- ◆ These will no longer be only “counting” experiments but rather will depend on observing distortions in the far detectors neutrino energy spectrum in **both neutrino and anti-neutrino samples**.



Why are ν_e and $\bar{\nu}_e$ Cross Sections Important?

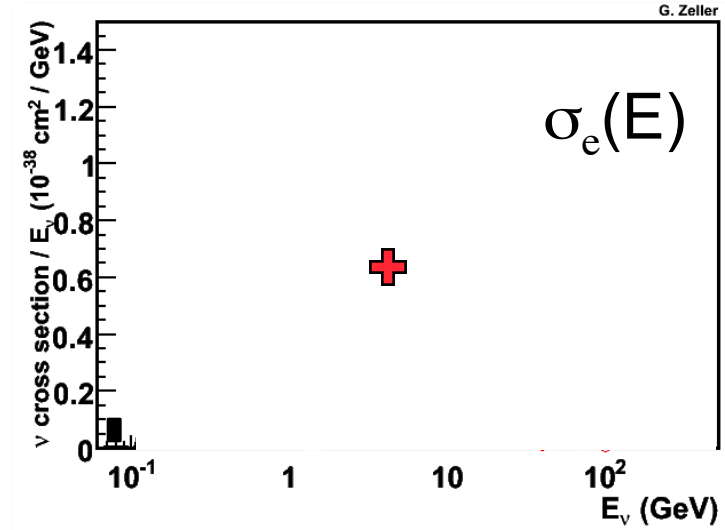
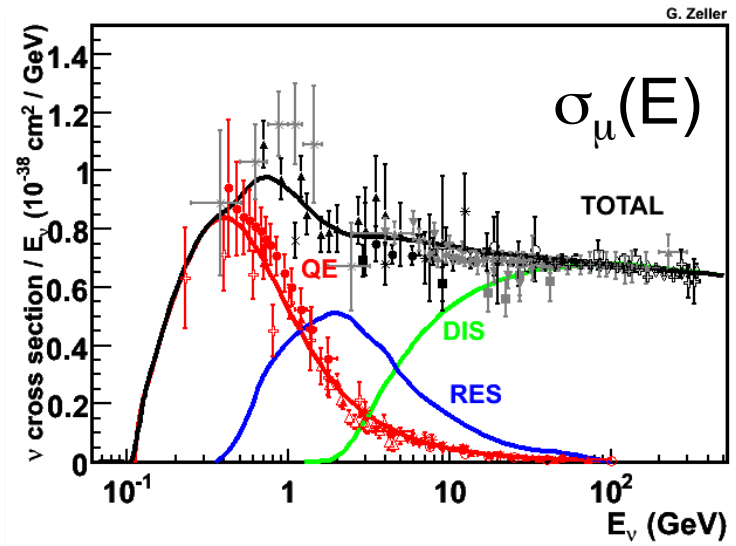
- ◆ Large θ_{13} means we could have reasonable statistics.
- ◆ However, as the now-well-known plot at right suggests, the asymmetry between ν and $\bar{\nu}$ will be small and the goal of constraining the range of δ will demand minimal systematic errors.
- ◆ One of these systematics will be our knowledge of ν_e and $\bar{\nu}_e$ cross sections in the relevant energy range.



(not including matter effects & backgrounds)
(S. Parke)
 11

How well do we know ν_e cross sections?

- ◆ **WE DON'T!** Need to measure the $\sigma_{\nu_e}(E)$ of multiple channels to fully predict a spectrum at a far detector for LBL experiments.



- ▼ **We infer them from $\sigma_{\nu_{\mu}}(E)$ results. The validity of this inference directly impacts the uncertainty of the measurements.**

One (+ ONE) to add to the collection...of one

- ◆ Gargamelle experiment published with around 200 electron and 60 positron events.
- ◆ Error bars on the order of 30%.

TOTAL CROSS SECTIONS FOR ν_e AND $\bar{\nu}_e$ INTERACTIONS
AND SEARCH FOR NEUTRINO OSCILLATIONS AND DECAY

Gargamelle Collaboration

J. BLIETSCHAU, H. DEDEN, F.J. HASERT, W. KRENZ, D. LANSKE, J.
MORFIN, M. POHL, K. SCHULTZE, H. SCHUMACHER, H. WEERTS and
L.C. WELCH

III. Physikalisches Institut der Technischen Hochschule, Aachen, Germany

Gargamelle: Nucl. Phys. B **133**: 205 (1978)

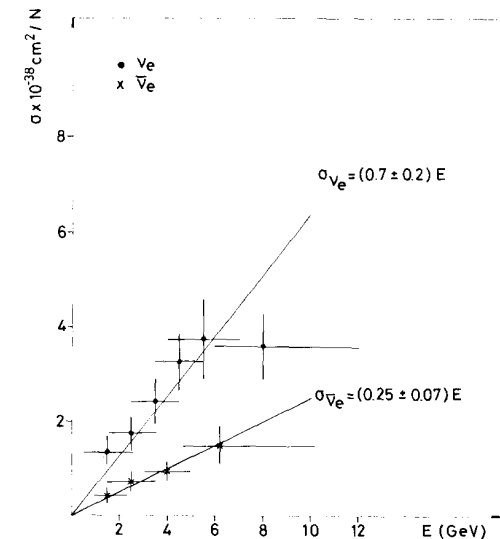


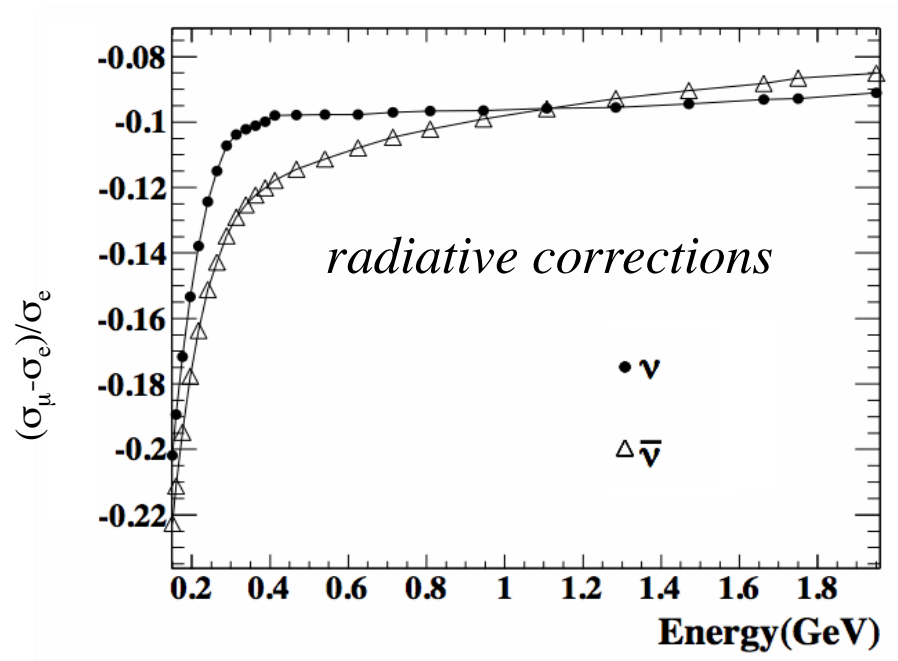
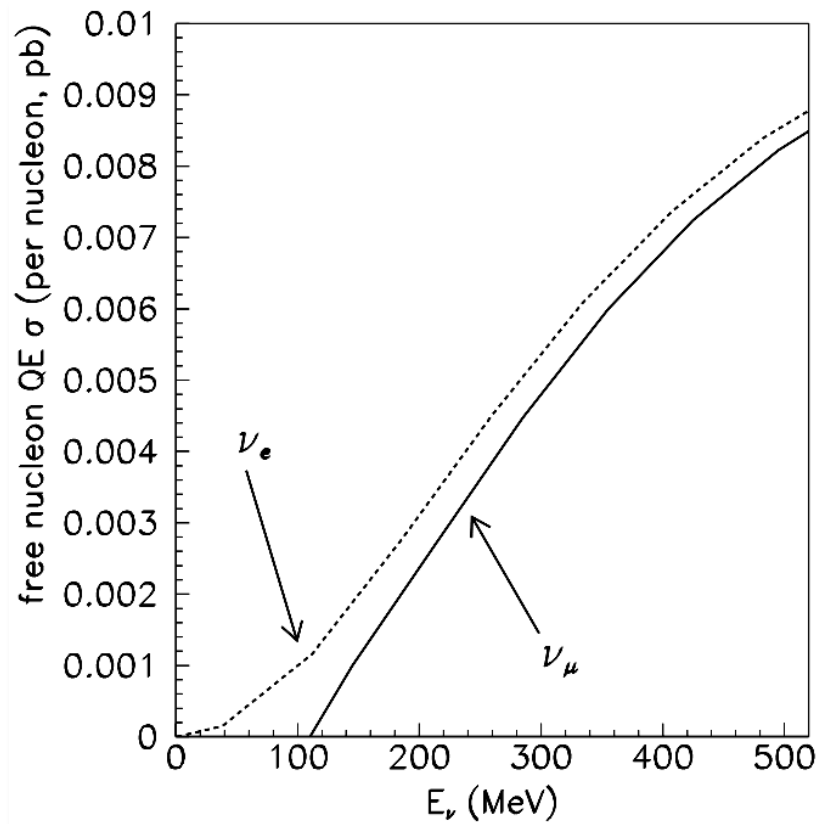
Fig. 2. Neutrino and antineutrino cross sections as a function of energy.

- ◆ MINERvA will have a sample of 65k ν_e CC events; 2.5k ν_e CCQE-like **produced** events in LE beam. In the ME beam probably factor of 5 higher statistics. Will be systematics limited \rightarrow FLUX.
- ◆ **Until then, we infer them from $\sigma_{\nu\mu}(E)$ results. The validity of this inference directly impacts the uncertainty of measurements.**

What are the Differences $\sigma_{\nu\mu}(E)$ and $\sigma_{\nu e}(E)$? Quasi-elastic Scattering

Day-McFarland study: Phys.Rev. D86 (2012) 053003

- ◆ QE scattering dominates at low energies (2nd oscillation maxima)
- ◆ Sources of possible differences and uncertainties - obvious:
 - ▼ Kinematic limits from μ / e mass difference.
 - ▼ Radiative Corrections. This may be an overestimate. Need full calculation.

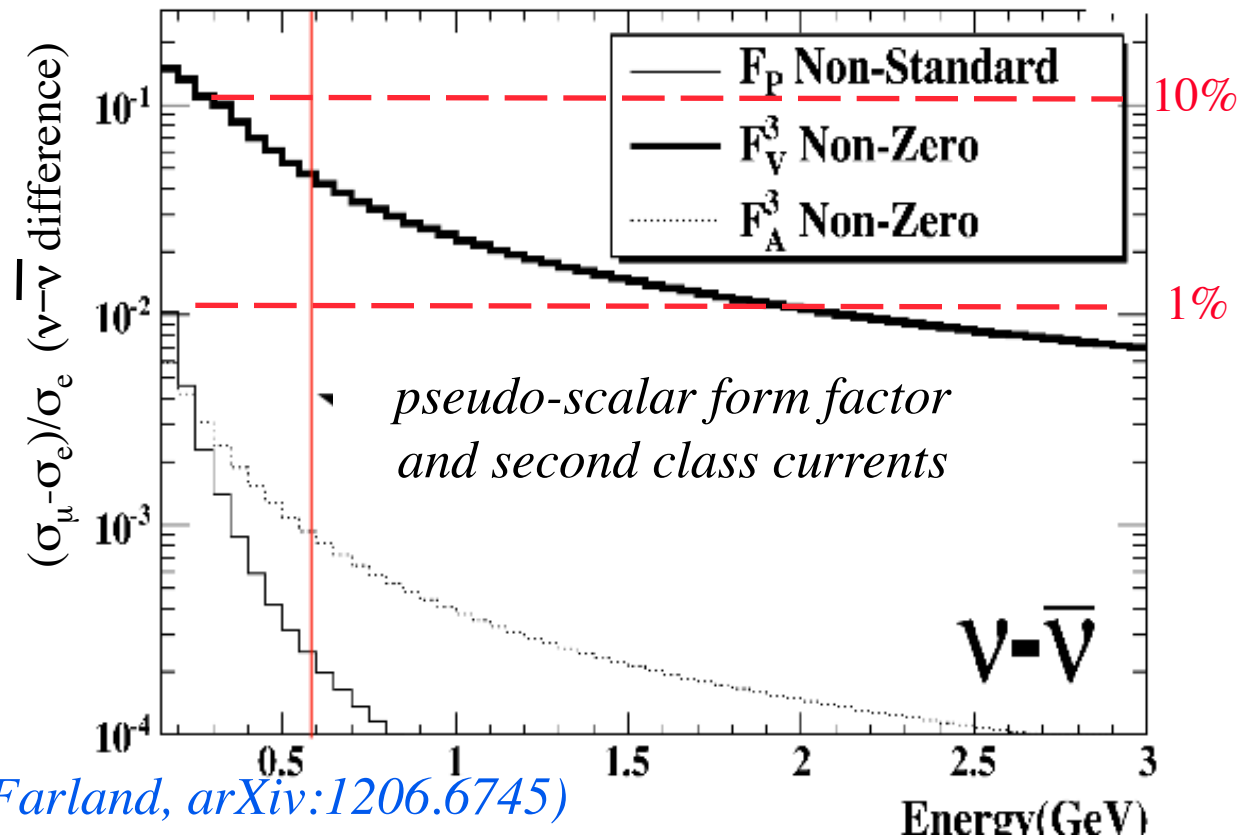


(M. Day, K. McFarland, arXiv:1206.6745)

What are the Differences $\sigma_{\nu\mu}(E)$ and $\sigma_{\nu e}(E)$? Quasi-elastic Scattering

Day-McFarland study: Phys.Rev. D86 (2012) 053003

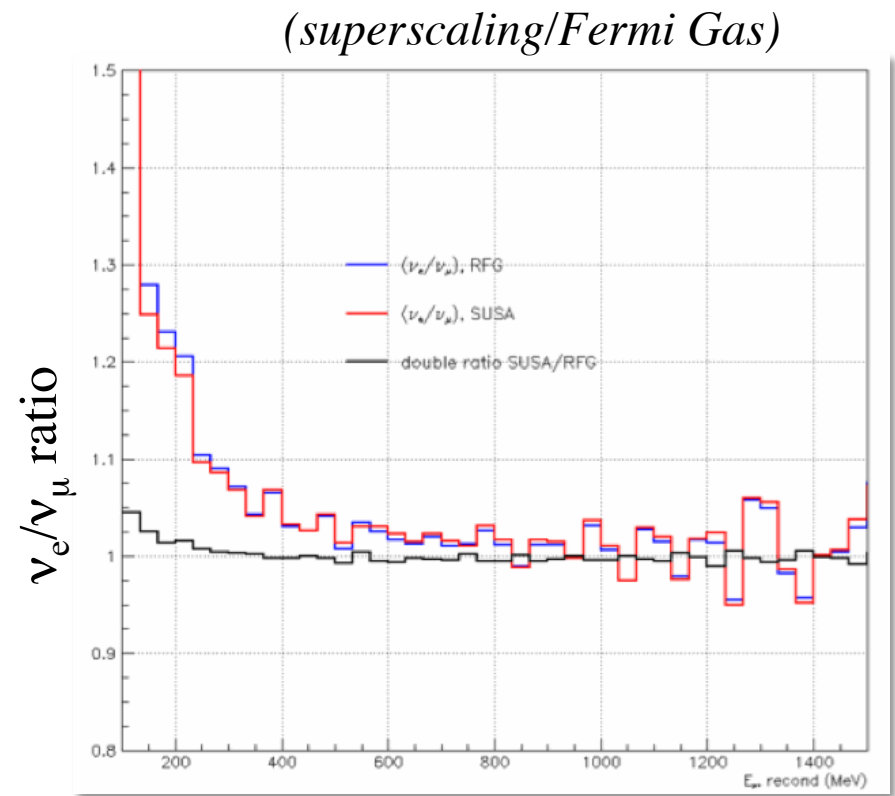
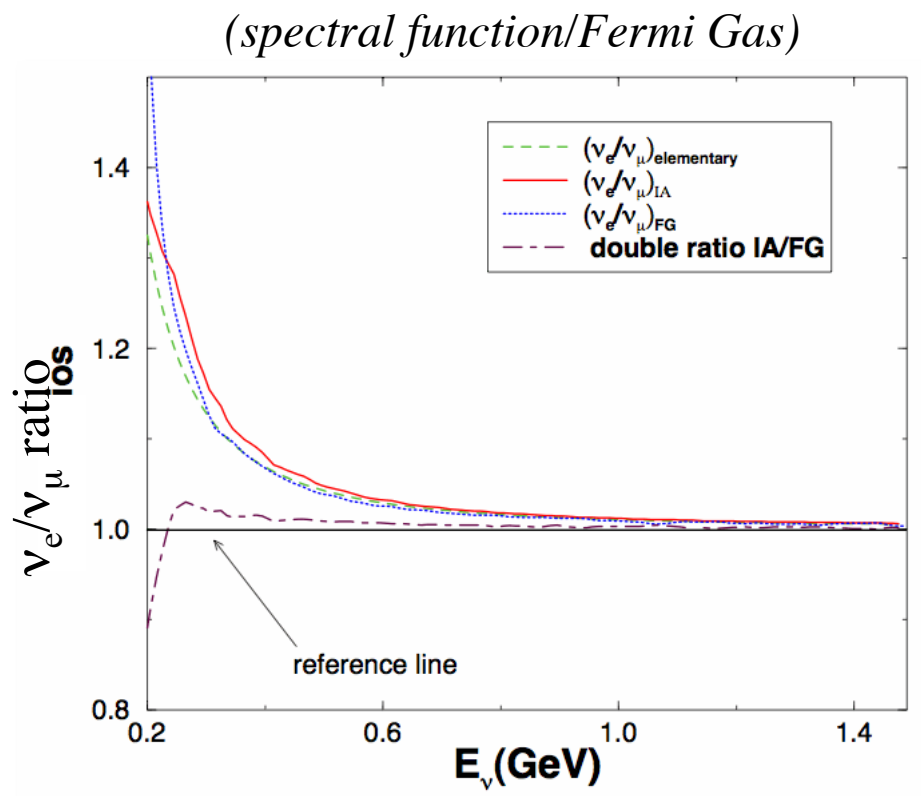
- ◆ Sources of possible differences: form factor uncertainties entering through lepton mass alterations - much more subtle:
 - ▼ Form factor contributions – both Axial and Pseudoscalar
 - ▼ Second class current contributions to vector and axial-vector form factors
- ◆ Possible contribution to CP uncertainties: effect on the FF could be different for ν and $\bar{\nu}$



(M. Day, K. McFarland, arXiv:1206.6745)

What are the Differences? Mainly QE Scattering Due to Nuclear Effects

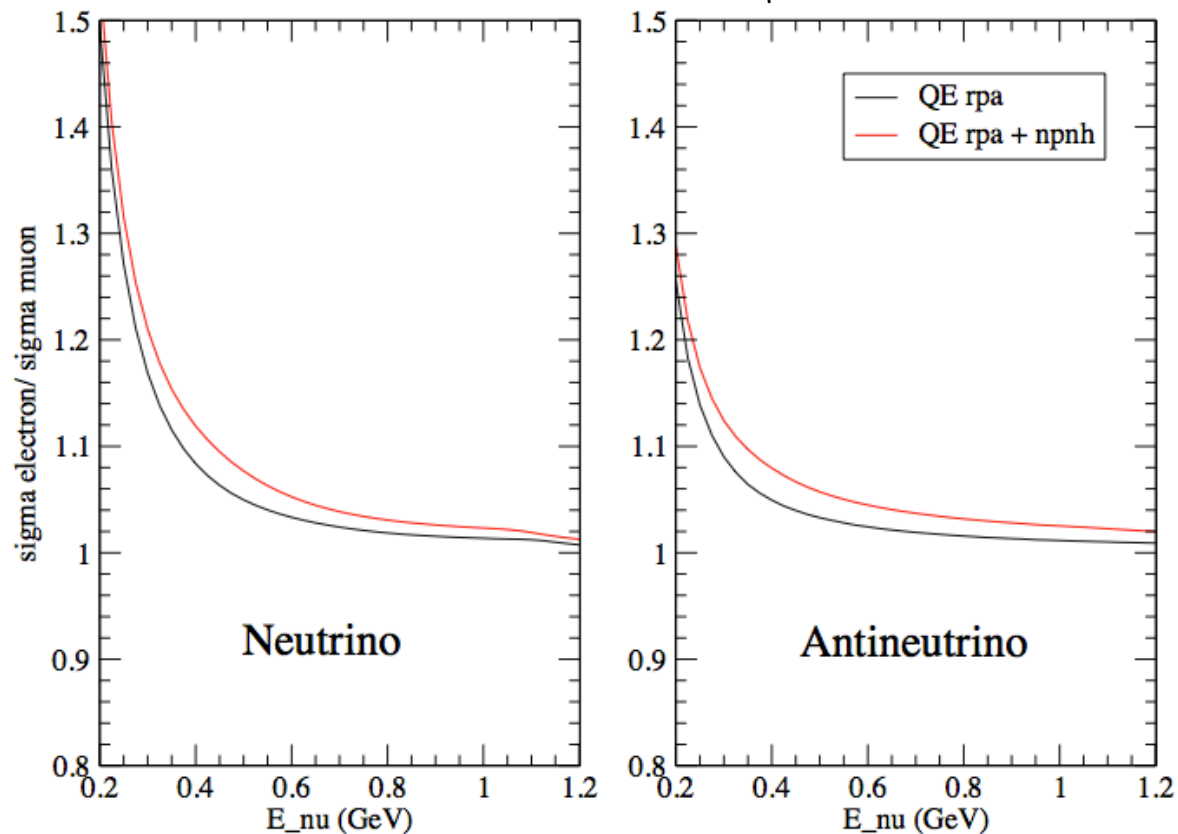
- ◆ For standard models, $\leq 5\%$ differences on ν_e/ν_μ ratio $E < 200$ MeV



What are the Differences? Mainly QE Scattering

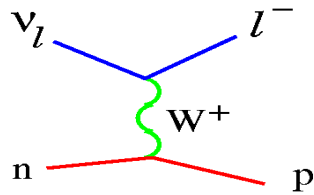
Meson-exchange Current Contributions – Marco Martini

- ◆ Hadronic part (nuclear response functions) is the same for ν_e or ν_μ cross section.
- ◆ However, the lepton tensor changes \rightarrow the relative weight of the nuclear responses in the several channels may change.
- ◆ The double ratio suggests the effect on the ν_e/ν_μ cross section ratio is $\leq 5\%$ (S. Zeller)



Nuclear Effects can Change the Energy Reconstruction for “QE” Events

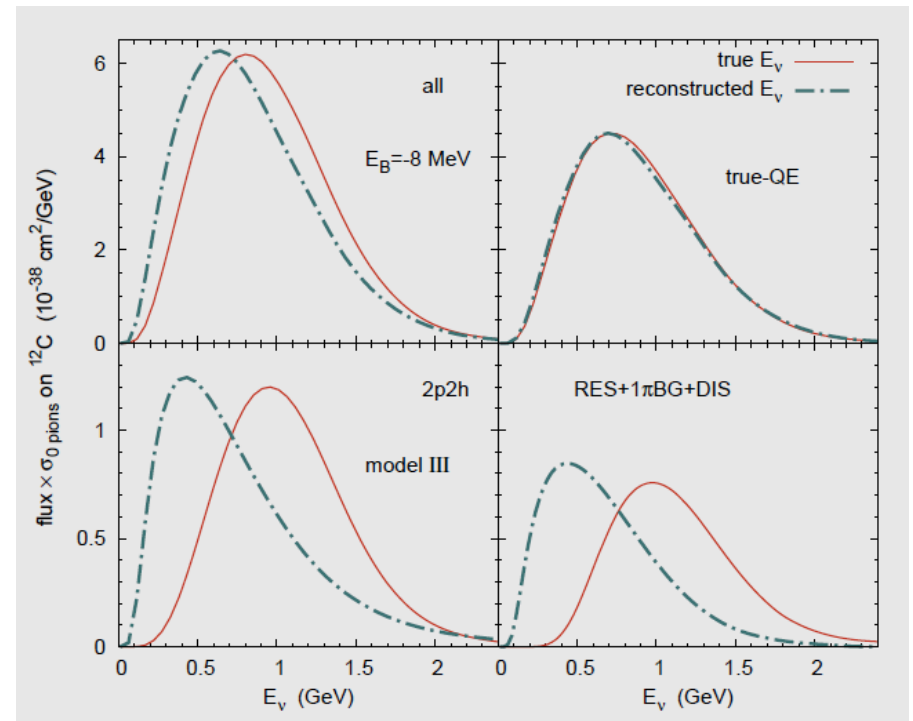
- In pure QE scattering on a nucleon at rest, the outgoing lepton can determine the neutrino energy:



$$E_\nu = \frac{2M_N E_\mu - m_\mu^2}{2(M_N - E_\mu + p_\mu \cos \theta_\mu)}$$

However, not on nuclei.

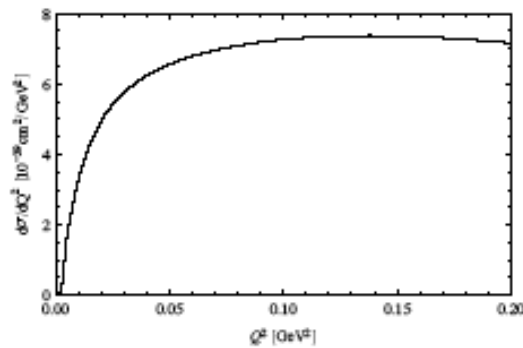
Reconstructed energy is shifted to lower values for all processes other than true QE off nucleon at rest



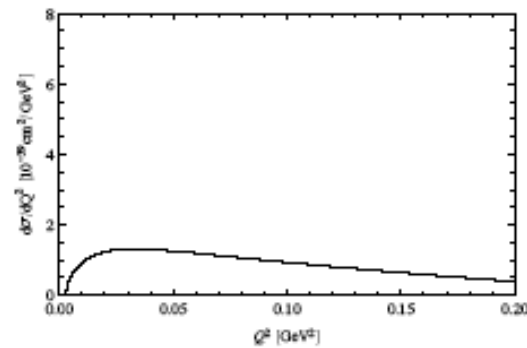
What are the Differences? Δ Production

Paschos – Schalla: arXiv:1209.4219

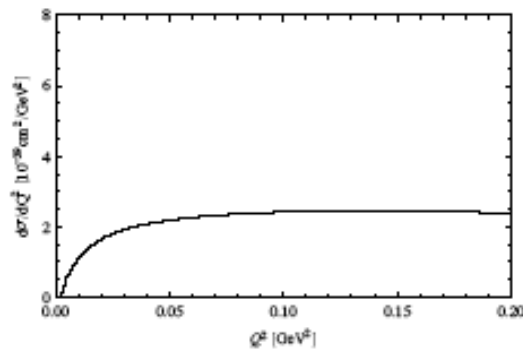
- ◆ Manny and his student have investigated ν_μ and $\bar{\nu}_\mu$ differences in Δ production in the low- Q ($Q^2 \approx m_\pi^2$) region where PCAC dominates the axial contribution.
- ◆ At $E = 1-2$ GeV, V part and V/A interference same size \rightarrow cancel for $\bar{\nu}$
- ◆ Use the Adler-Nussinov-Paschos model for nuclear corrections.



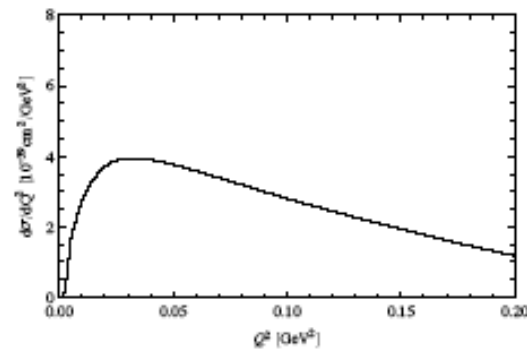
(a) $\nu_\mu p \rightarrow \mu^- X^{++}$



(b) $\bar{\nu}_\mu p \rightarrow \mu^+ X^0$



(c) $\nu_\mu n \rightarrow \mu^- X^+$

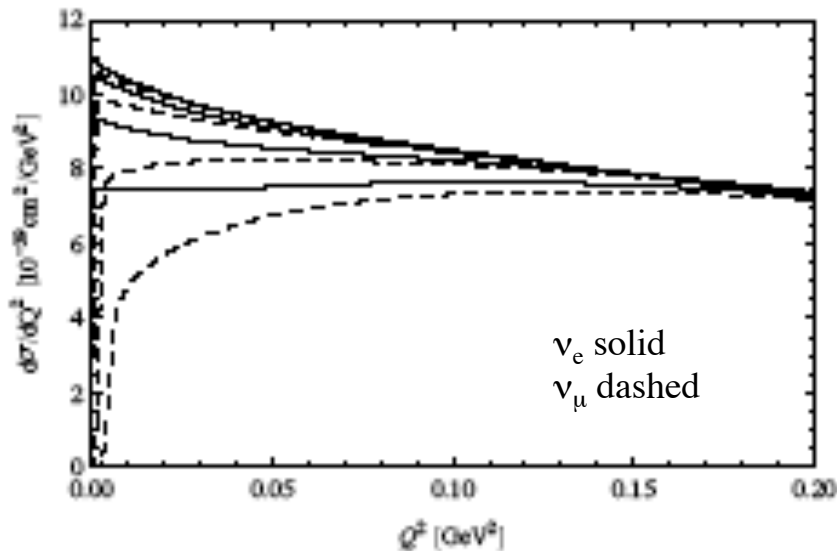


(d) $\bar{\nu}_\mu n \rightarrow \mu^+ X^-$

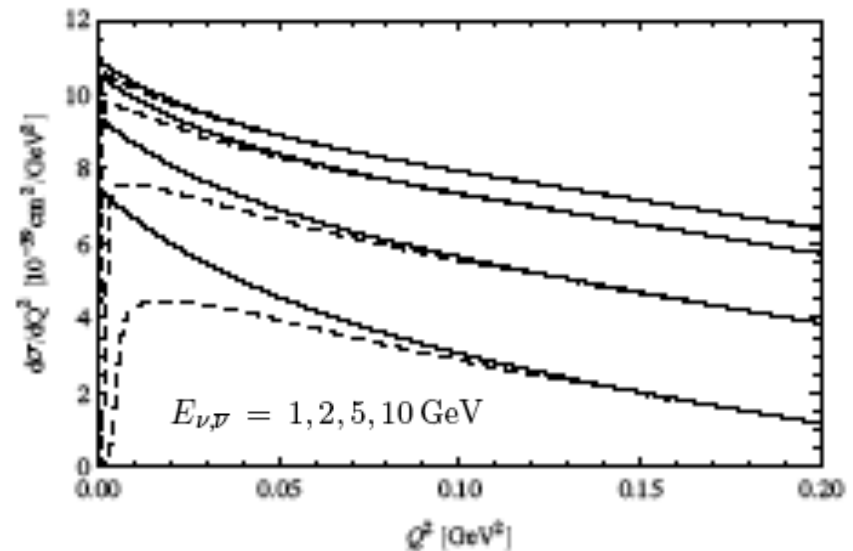
What are the Differences? Δ Production

Paschos – Schalla: arXiv:1209.4219

- ◆ Paschos-Schalla predicts the following differences in cross sections where only the lepton mass term contributions are shown and any differences in form factors are not yet included.



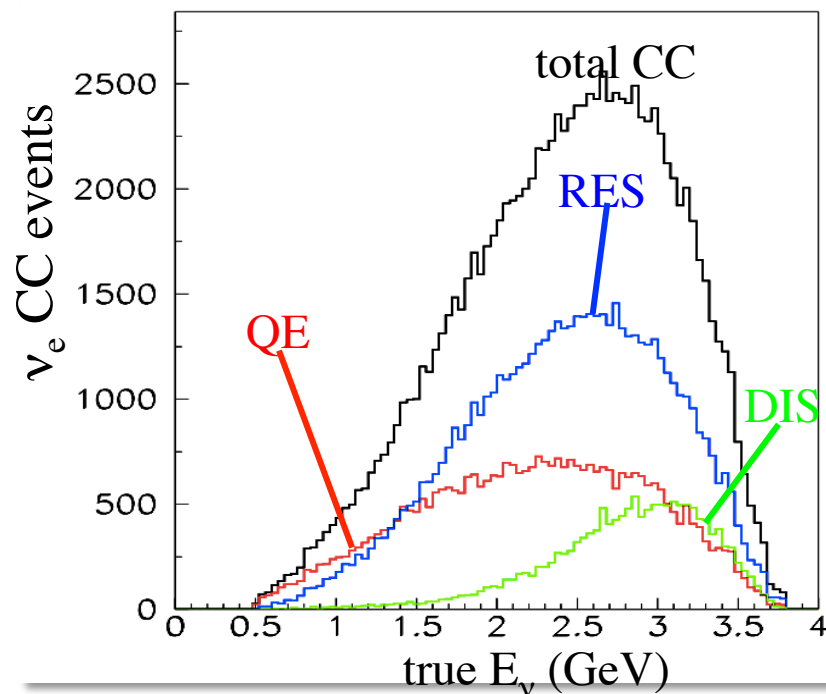
(a) $\nu_e p \rightarrow e^- \Delta^{++}$



(b) $\bar{\nu}_e n \rightarrow e^+ \Delta^-$

Can we Actually MEASURE these Differences in the 0.5 – 6 GeV region

- ◆ Need to measure the $\sigma_e(E)$ of multiple channels to predict spectrum at the far detector.
 - ▼ Want an intense source of ν_e events.
 - ▼ Would like to know the flux of ν_e (and ν_μ , by the way) to order 1%.



μ^+		μ^-	
Channel	N_{evts}	Channel	N_{evts}
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ν_e NC	1,387,698	ν_μ NC	1,584,003
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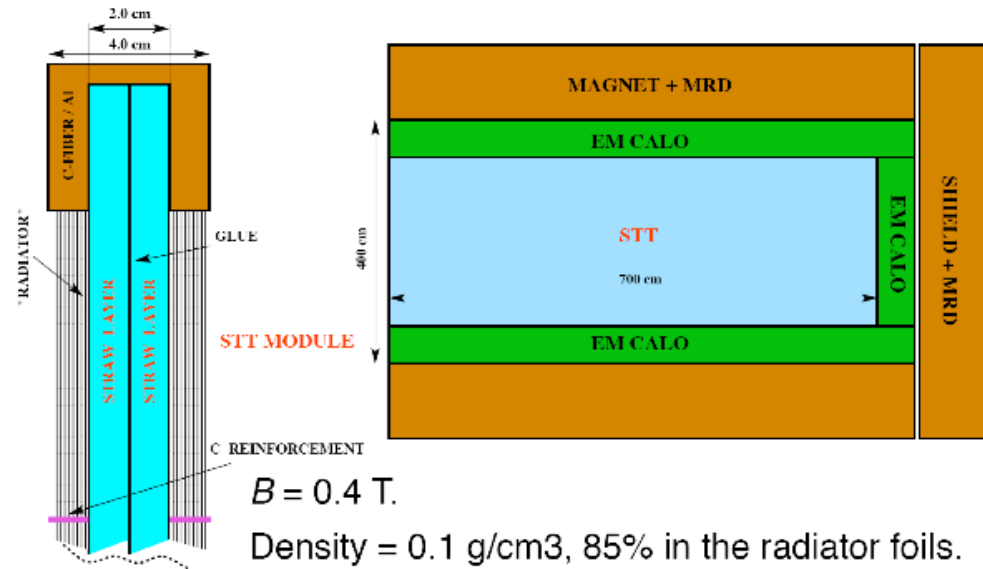
event rates per 1E21 POT -
100 tons at 50m

nuSTORM Near Detector(s)

Etam MESSOMO

1) HighRes -
High Resolution Straw-tube.
transition radiation
Magnetized Detector

Considered now by LBNE



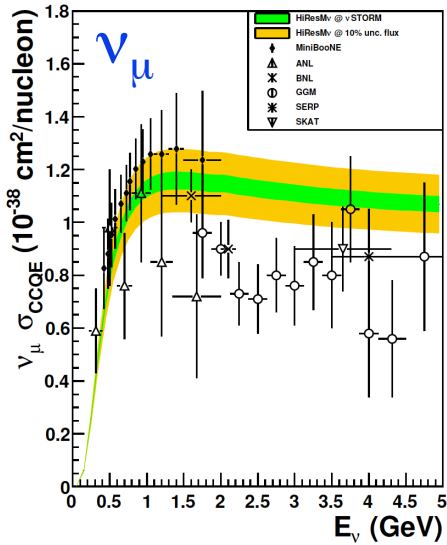
Transition Radiation \Rightarrow Electron ID $\Rightarrow \gamma$ (w. Kinematics)
 dE/dx \Rightarrow Proton, π , K ID
Magnet/Muon Detector $\Rightarrow \mu$

2) A 1-2 ton fiducial liquid hydrogen/deuterium track sensitive target upstream of HiRes. This could be a “bubble chamber”.

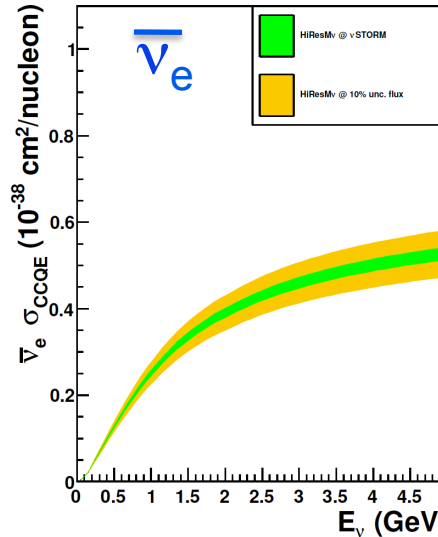
Scattering Measurements with nuSTORM + Near Detector

nuSTORM provides a well-known ($\delta \phi(E) \approx 1\%$) beam of ν and $\bar{\nu}$.

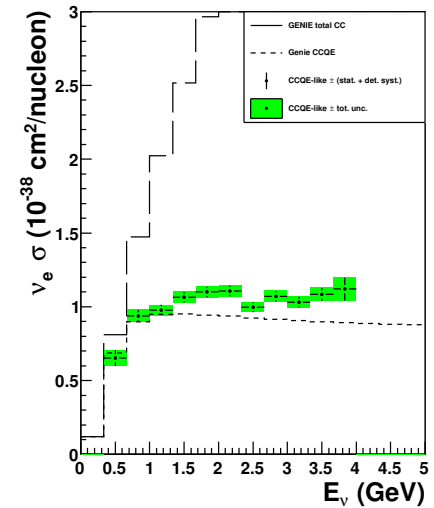
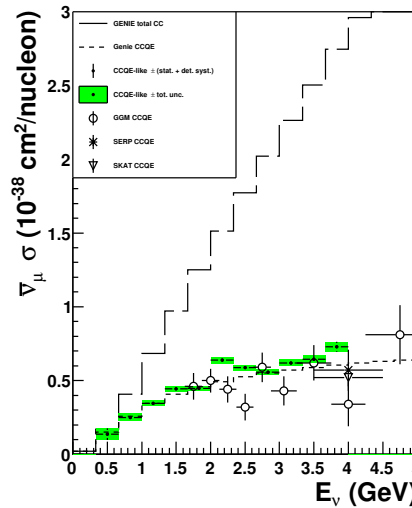
Ed Santos – Imperial College



HIRESM ν – systematics only



Total errors:
μBooNE LAr



Conclusions: What does nuSTORM bring to Neutrino-nucleus Interaction Physics?

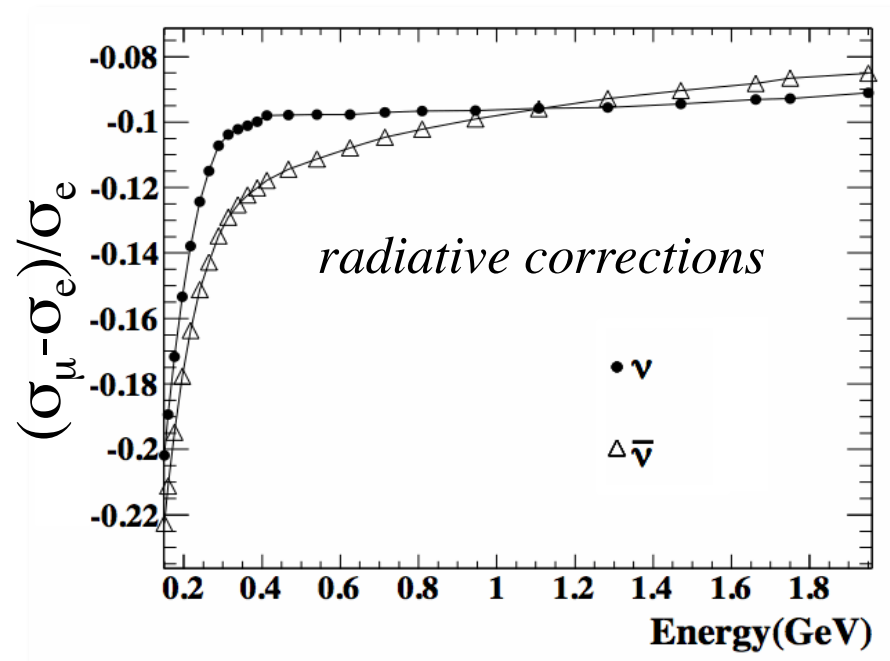
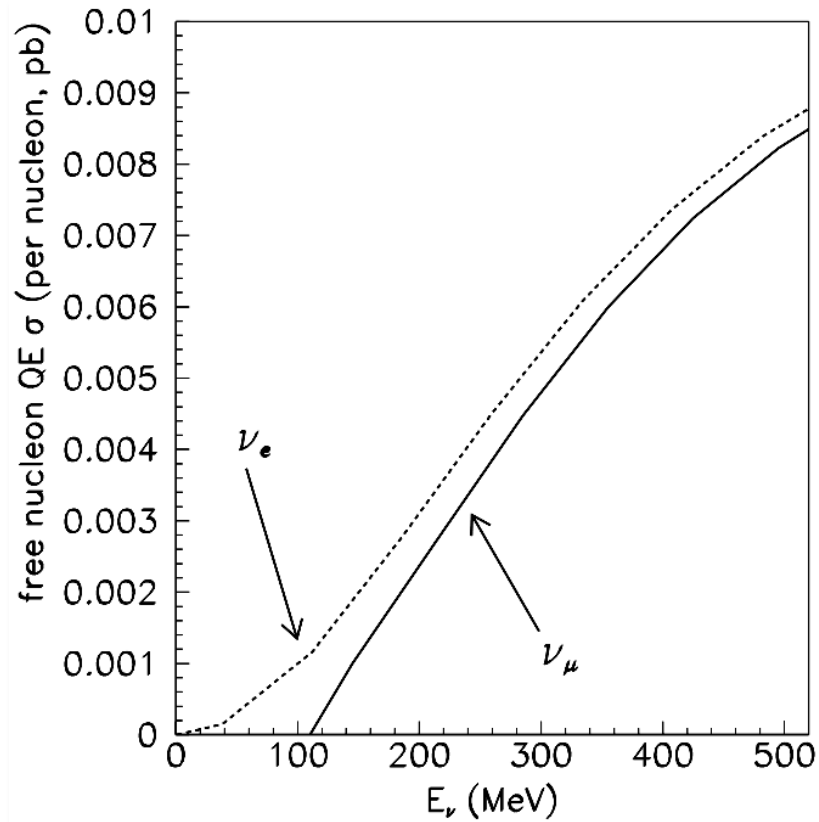
- ◆ Obvious benefit in measuring ν_e events.
- ◆ Use vastly improved nuSTORM neutrino flux uncertainty to measure neutrino cross sections and nuclear effects!
- ◆ There are many physics topics important in their own right and essential for precision neutrino oscillation experiments that will be awaiting the results of a high-resolution detector in the accurate nuSTORM flux!
- ◆ However, this is not the nuSTORM experiment that was approved. Adding at least a precision near detector and perhaps a H/D target.
- ◆ **This calls for a new collaboration using the nuSTORM facility. First meeting/workshop to establish this collaboration in November.**

BACKUP

What are the Differences $\sigma_{\nu\mu}(E)$ and $\sigma_{\nu e}(E)$? Quasi-elastic Scattering

Day-McFarland study: Phys.Rev. D86 (2012) 053003

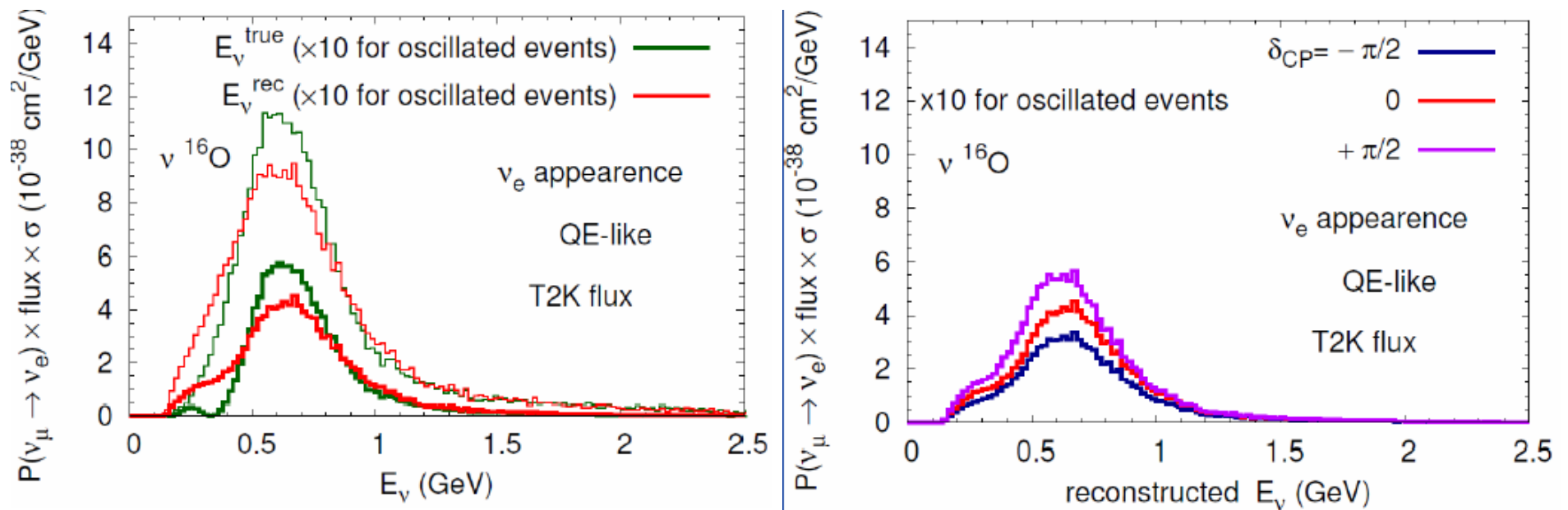
- ◆ QE scattering dominates at low energies (2nd oscillation maxima)
- ◆ Sources of possible differences and uncertainties - obvious:
 - ▼ Kinematic limits from μ / e mass difference.
 - ▼ Radiative Corrections. **This may be overestimated. Need full calculation.**



(M. Day, K. McFarland, arXiv:1206.6745)

Nuclear Effects and Oscillation Measurements

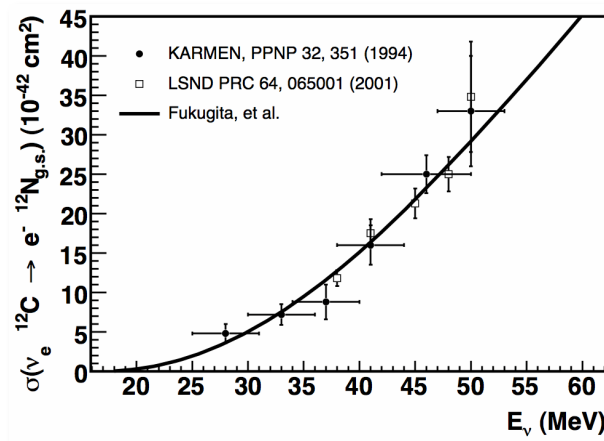
Ulrich Mosel using his Giessen Boltzmann-Uehling-Uhlenbeck (GiBUU) Transport Model looking at T2K



How well do we know cross sections: ν_e vs. ν_μ ?

Existing ν_e Cross Section Data

- ◆ What do we know about $\sigma_{\nu_e}(E)$? **Mostly very low energy results.**
 - ▼ Reactor neutrinos studying Inverse Beta Decay
 - ▼ Solar neutrino off deuterium (SNO)
 - ▼ Stopping π/μ decay neutrinos off higher A targets
 - ▼ See Formaggio and Zeller **Rev. Mod. Phys. 84, 1307–1341 (2012)**.
- ◆ One of few measurements of spectral shape of σ reflects the upper limit of most existing measurements, $E \leq 50$ MeV.



(Formaggio & Zeller,
Rev. Mod. Phys. 2012)

Addressing the lack of F_2 **Neutrino** Nuclear Effects Analyses

Nuclear PDFs from neutrino deep inelastic scattering

nCTEQ

**K. Kovarik (Karlsruhe) I. Schienbein (LPSC-Grenoble),
J.-Y. Yu (SMU), C. Keppel (Hampton/JeffersonLab)
J.G.M. (Fermilab), F. Olness (SMU), J.F. Owens (Florida State U)**

Also analyses by:

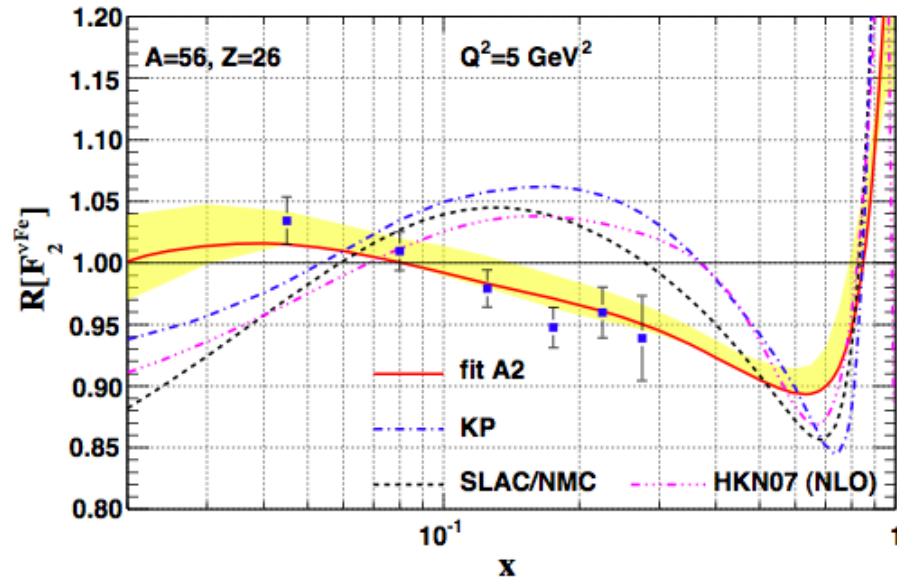
K. Eskola, V. Kolhinen and C. Salgado
and

D. de Florian, R. Sassot, P. Zurita and M. Stratmann

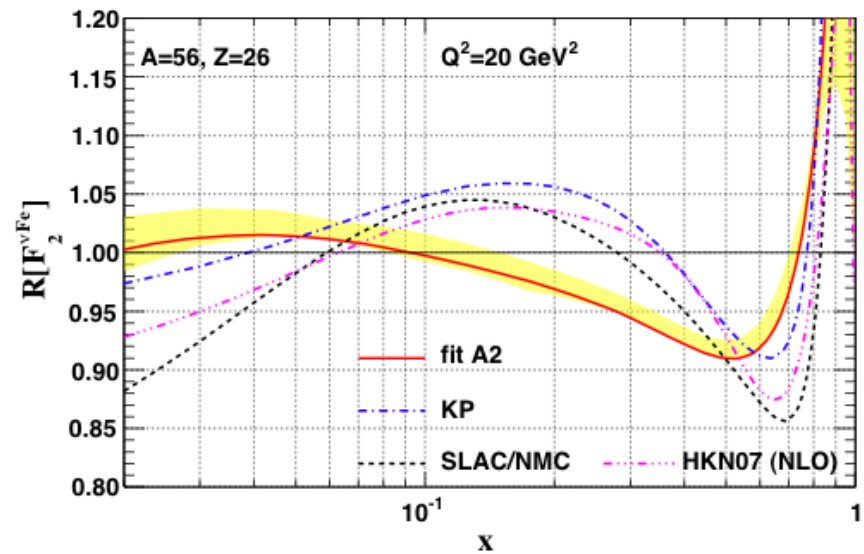
Significant Implications for Oscillation Experiments

- ◆ Can not simply plug in effective $\sigma_{\pi^+}^A$ from experiments in a significantly different beam.
- ◆ In a two-detector LBL oscillation experiment the neutrino flux entering the far detector is altered from the neutrino flux at the near detector due to geometry and oscillations.
- ◆ The $\sigma_c^A(E)$ effective that should be applied to expectations (Monte Carlo) at the far detector is NOT the same as that which we would measure at the near detector.
However, the near detector results give us an excellent starting point for calculating the difference.
- ◆ **The convoluted $\phi(E' \geq E)$ \otimes $\sigma(E)$ \otimes Nuc($E' \geq E$) systematics need to be correctly incorporated in determining the systematics of oscillation parameter measurements. Who is addressing this important consideration now?**

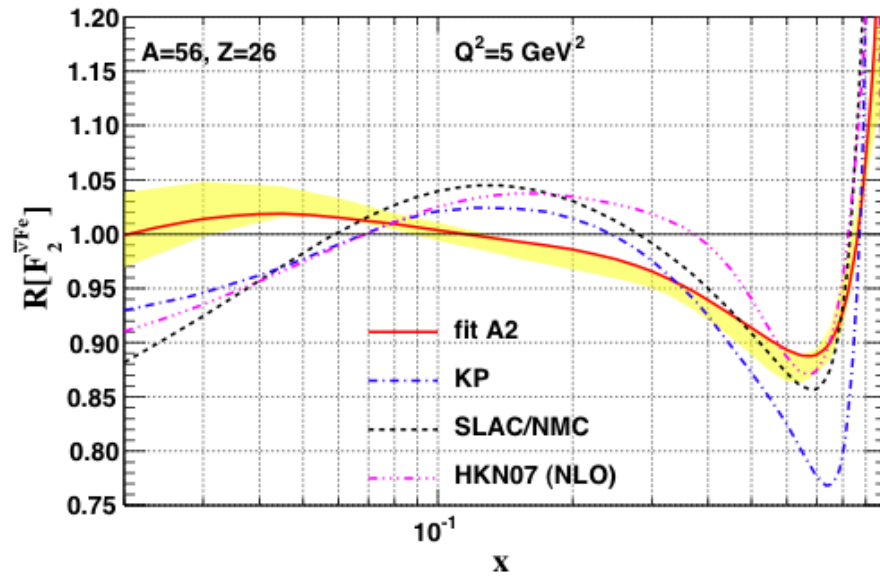
F_2 Structure Function Ratios: ν -Iron



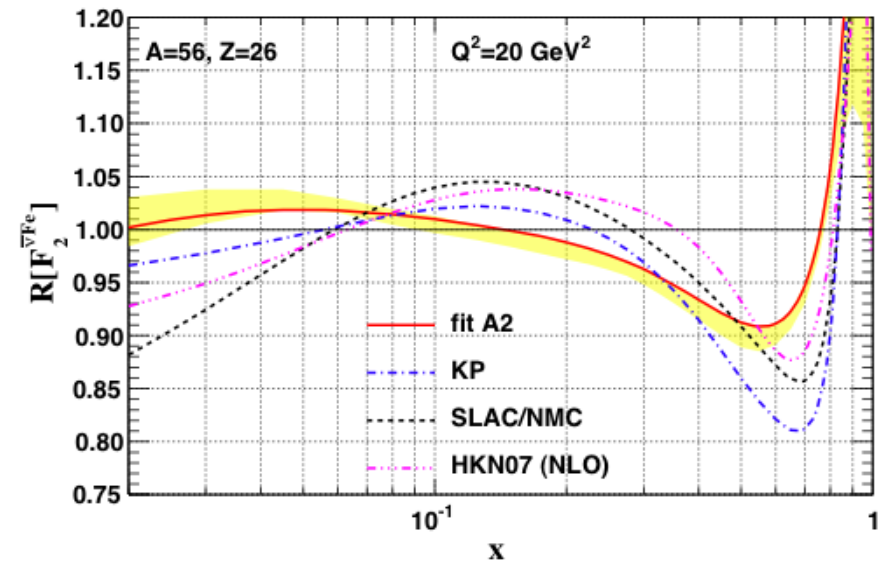
$$\frac{F_2(\nu + \text{Fe})}{F_2(\nu + [n+p])}$$



F_2 Structure Function Ratios: $\bar{\nu}$ -Iron

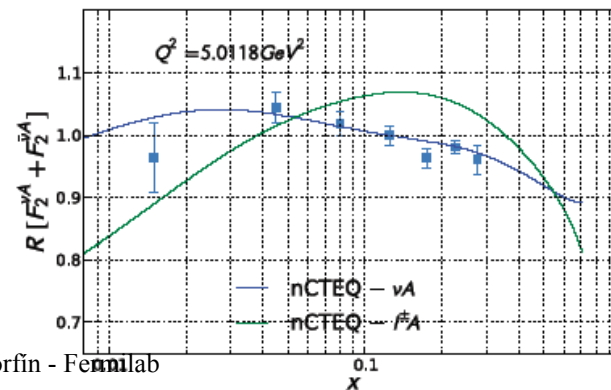
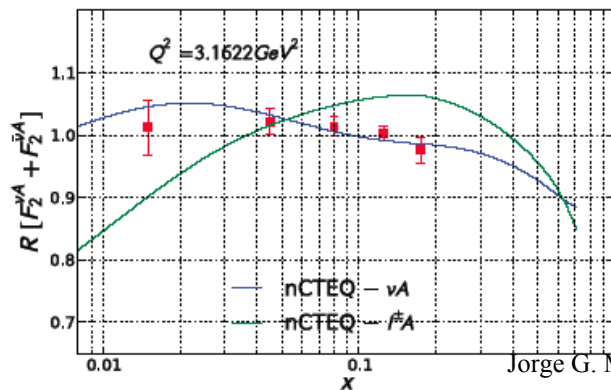
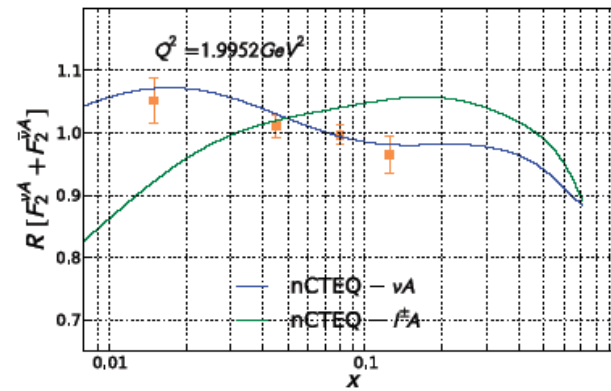
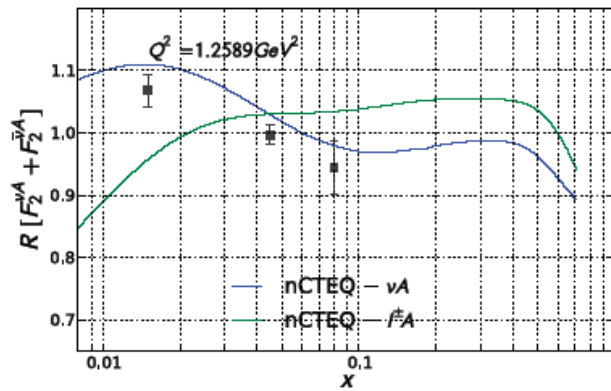


$$\frac{F_2(\nu + \text{Fe})}{F_2(\nu + [n+p])}$$



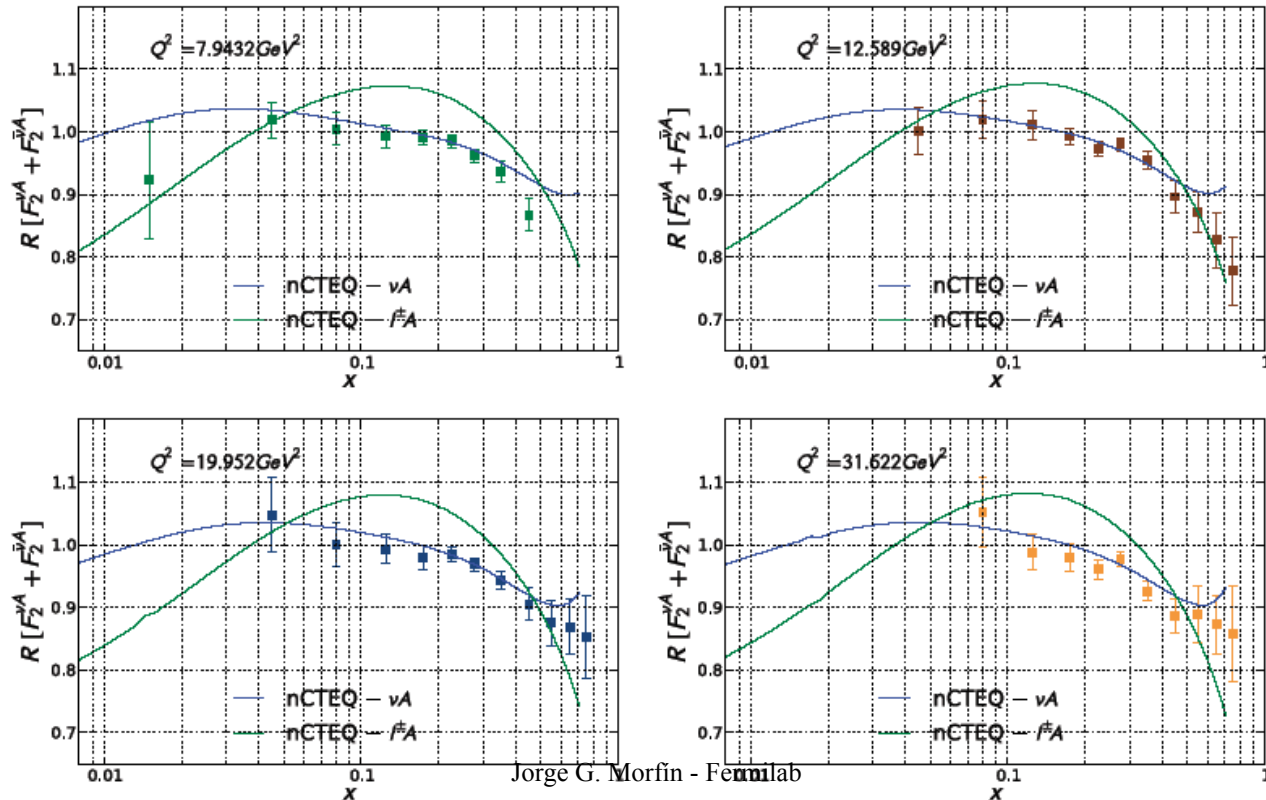
A More-Detailed Look at Differences

- ◆ NLO QCD calculation of $\frac{F_2^{\nu A} + F_2^{\bar{\nu} A}}{2}$ in the ACOT-VFN scheme
 - ▼ charge lepton fit undershoots low-x data & overshoots mid-x data
 - ▼ low- Q^2 and low-x data cause tension with the shadowing observed in charged lepton data



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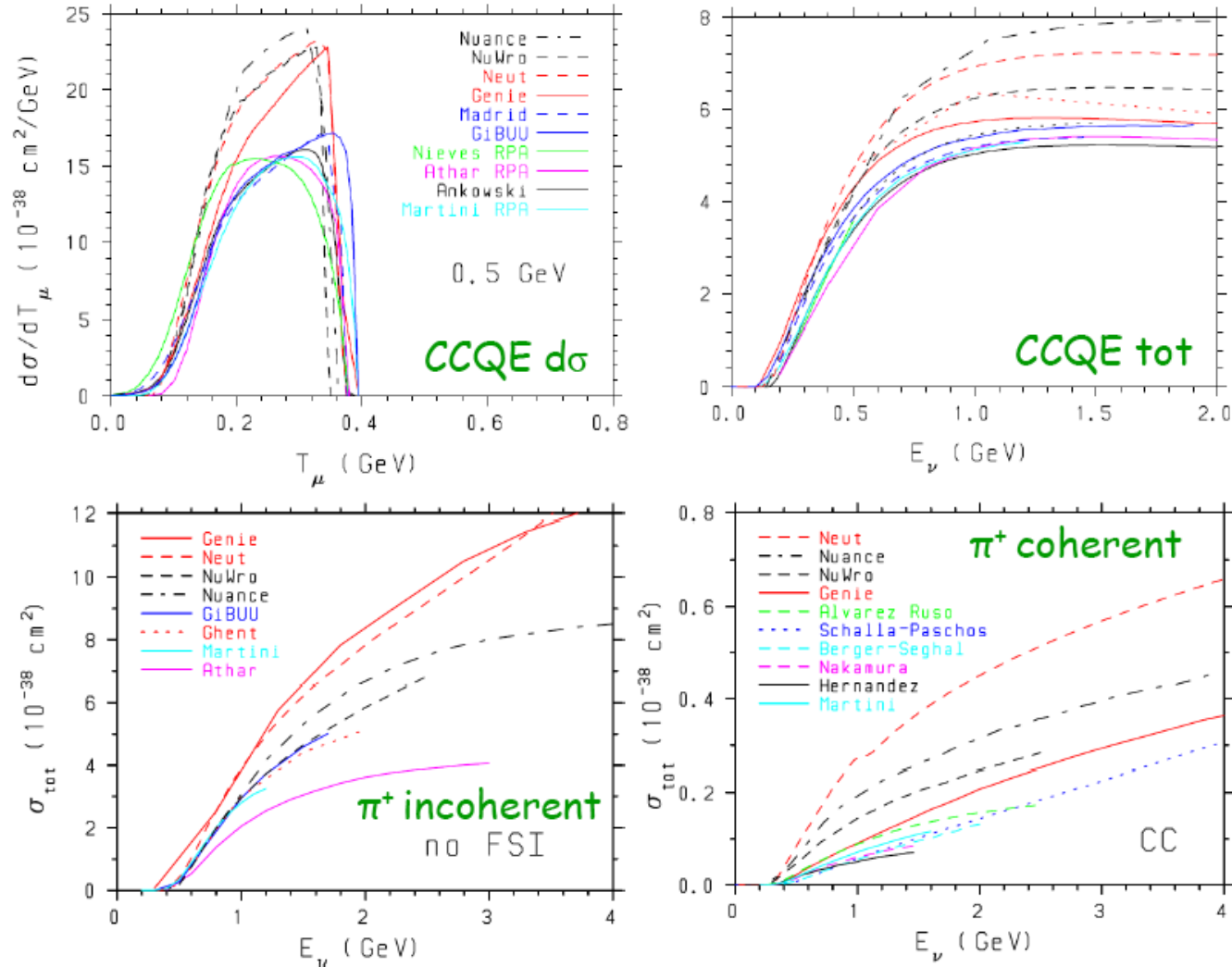


What are these Nuclear Effects $\text{Nuc}_{c,d,e.. \rightarrow c} (E' \geq E)$ in Neutrino Nucleus Interactions? (Partial List)

- ◆ Target nucleon in motion – classical Fermi gas model or the superior spectral functions (Benhar et al.)
- ◆ Multi-nucleon initial states: Short-range correlations, meson exchange currents.
- ◆ Form factors, structure functions, resonance widths, parton distribution functions and, consequently, cross sections are modified within the nuclear environment. (Butkevich / Kulagin, Tsushima et al., Kovarik et al.)
- ◆ Produced topologies are modified by final-state interactions modifying topologies and possibly reducing **detected** energy and **increasing** wrong-sign background.
 - ▼ Convolution of $\delta\sigma(n\pi)$ (x) formation zone uncertainties (x) π -charge-exchange/absorption probabilities and nuclear density uncertainties.
- ◆ **Systematics associated with each of these effects.**
 - ◆ Monte Carlos – like GENIE – try to include all these effects.
GENIE needs improvements! GENIE group needs additional help from the community.

How well off are we with ν_μ Cross sections: Range of Existing Model (MC) Predictions off C

NuInt09 – Steve Dytman



Example Model Uncertainties

Cross Section Model Uncertainties

Uncertainty	1 σ
M_A (Elastic Scattering)	$\pm 25\%$
F_A (Elastic scattering)	$\pm 30\%$
M_A (CCQE Scattering)	+25% -15%
CCQE Normalization	+20% -15%
CCQE Vector Form factor model	on/off
CC Resonance Normalization	$\pm 20\%$
M_A (Resonance Production)	$\pm 20\%$
M_V (Resonance Production)	$\pm 10\%$
1pi production from $\nu p / \bar{\nu} n$ non-resonant interactions	$\pm 50\%$
1pi production from $\nu n / \bar{\nu} p$ non-resonant interactions	$\pm 50\%$
2pi production from $\nu p / \bar{\nu} n$ non-resonant interactions	$\pm 50\%$
2pi production from $\nu n / \bar{\nu} p$ non-resonant interactions	$\pm 50\%$
Modtly Pauli blocking (CCQE) at low Q^2 (change PB momentum threshold)	$\pm 30\%$

Intranuclear Rescattering Uncertainties

Uncertainty	1 σ
Pion mean free path	$\pm 20\%$
Nucleon mean free path	$\pm 20\%$
Pion fates – absorption	$\pm 30\%$
Pion fates – charge exchange	$\pm 50\%$
Pion fates – Elastic	$\pm 10\%$
Pion fates – Inelastic	$\pm 40\%$
Pion fates – pion production	$\pm 20\%$
Nucleon fates – charge exchange	$\pm 50\%$
Nucleon fates – Elastic	$\pm 30\%$
Nucleon fates – Inelastic	$\pm 40\%$
Nucleon fates – absorption	$\pm 20\%$
Nucleon fates – pion production	$\pm 20\%$
AGKY hadronization model – x_T distribution	$\pm 20\%$
Delta decay angular distribution	On/off
Resonance decay branching ratio to photon	$\pm 50\%$

Hugh Gallagher

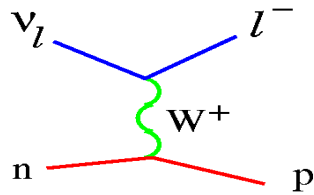
References: (1) www.genie-mc.org, (2) arXiv:0806.2119, (3) D. Bhattacharya, Ph. D Thesis (U. Pittsburgh) 2009.

Have these experiments really measured M_A ?

- ◆ Just as we have noted, we are observing an **effective $\sigma_c^A(E)$** in our detectors...
 - ▼ What has been measured is a parameter M_a^{eff}
 - ▼ It depends on the use of RFGM or Spectral Functions.
 - ▼ It depends on the nucleus used and...
 - ▼ It depends on the incoming flux.
 - ▼ It also depends on number of initial nucleons involved
- ◆ **Need nuSTORM with its accurate flux and series of nuclear targets with high-resolution detector(s).**
- ◆ Also need at least a good model for pion production which, through FSI, is the main background for QE.
- ◆ (QE) measurements calculating E and Q^2 via the muon are in trouble!

Nuclear Effects can Change the Energy Reconstruction for “QE” Events

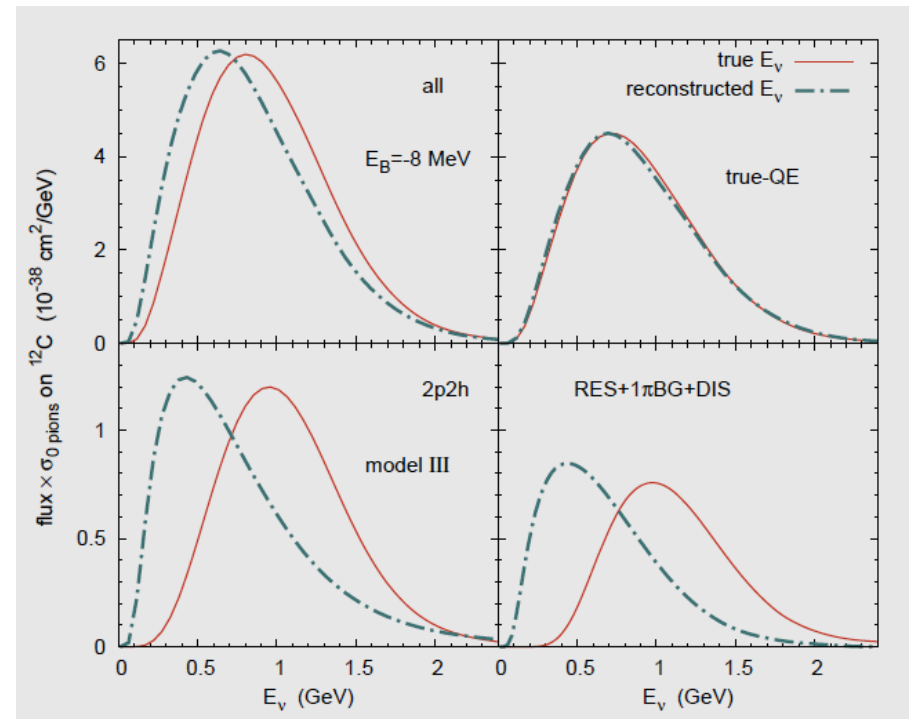
- In pure QE scattering on a nucleon at rest, the outgoing lepton can determine the neutrino energy:



$$E_\nu = \frac{2M_N E_\mu - m_\mu^2}{2(M_N - E_\mu + p_\mu \cos \theta_\mu)}$$

However, not on nuclei.

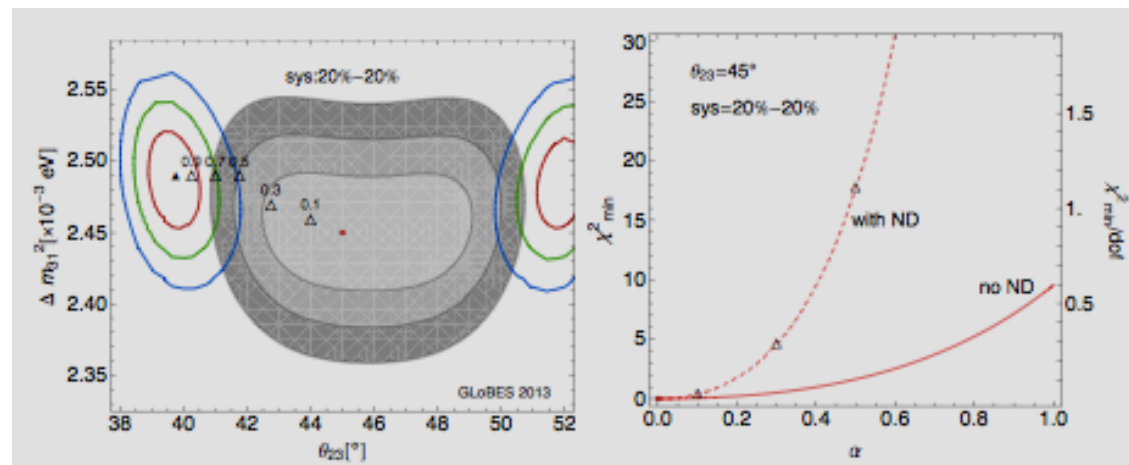
Reconstructed energy is shifted to lower values for all processes other than true QE off nucleon at rest



Detailed Study by P. Coloma and P. Huber

arXiv 1307.1243

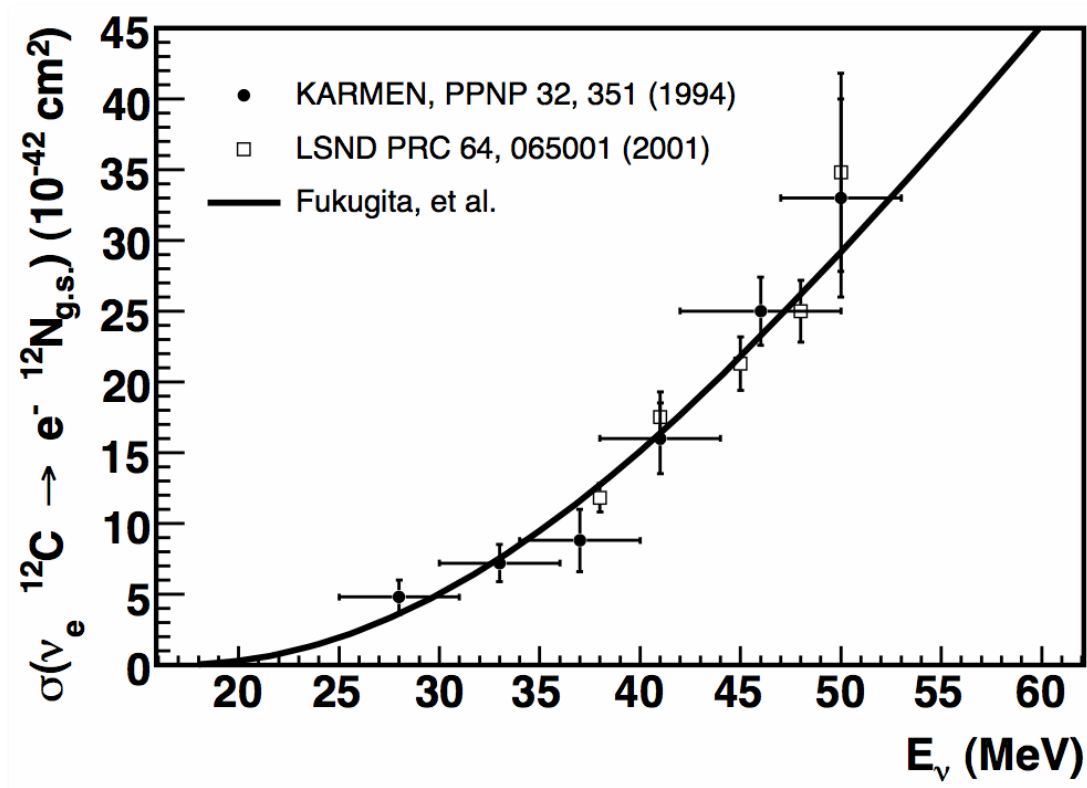
- ◆ Disappearance experiment using CC QE-like signal events. T2K – 5 years; 850 QE
- ◆ QE-like includes pion absorption and scattering off nucleon pairs. 1300 QE-like
- ◆ E_ν is reconstructed from the observed muon which gives a lower E_ν for non-QE.
- ◆ Give a quantitative estimate of this problem using: $N_i^{\text{test}}(\alpha) = \alpha \times N_i^{\text{QE}} + (1 - \alpha) \times N_i^{\text{QE-like}}$
- ◆ $\alpha = 1$ implies completely ignore nuclear effects while $\alpha = 0$ implies you know/model the nuclear effects completely.
- ◆ The importance of a near detector to help normalize the signal is obvious. However have not yet included different near and far incoming neutrino spectra.
- ◆ Even with ND, $\alpha = 0.3 \rightarrow 1 \sigma$ bias in parameters! **Need accurate nuclear model!**



Example of Existing Data: Carbon



- ◆ One of few measurements of spectral shape of σ reflects the upper limit of most existing measurements, $E \leq 50$ MeV.



(Formaggio & Zeller, Rev. Mod. Phys. 2012)