

# 100 TeV PDFs for ~~the LHC~~

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Michigan State University

## References:

- J. Rojo: kickoff meeting for FCC at CERN, Feb. 2014
- Snowmass QCD writeup: arXiv:1310.5189
- Les Houches session 1 writeup: [www.pa.msu.edu/~huston/leshouches\\_2013/LH13.pdf](http://www.pa.msu.edu/~huston/leshouches_2013/LH13.pdf) (to be posted on archive this week)

## Workshop on Physics at a 100 TeV Collider

April 23-25, 2014, SLAC



Workshop Topics  
PDFs and Generators  
Detector Challenges  
SM at 100 TeV  
Physics Reach  
BSM Spectroscopy

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[www.slac.stanford.edu/th/100TeV.html](http://www.slac.stanford.edu/th/100TeV.html)

MSU 24 – Stanford 20



# Introduction

- Cross sections at any hadron-hadron collider depend on PDFs determined by global fits to collider data, but (more importantly, at least so far) from DIS data (HERA and fixed target) and from fixed target Drell-Yan data
- For hadron collider predictions, global fit data taken at smaller  $Q^2$  is evolved to higher  $Q^2$  values using DGLAP evolution, at LO/NLO/NNLO

$$\sigma = \sum_{a,b} \int_0^1 dx_1 f_{a/A}(x_1, \mu_F^2) \int_0^1 dx_2 f_{b/B}(x_2, \mu_F^2) \left\{ \int d\hat{\sigma}_{ab}^{LO}(\alpha_s) \Theta_{\text{obs}}^{(m)} + \alpha_s(\mu_R^2) \left[ \int (d\hat{\sigma}_{ab}^V(\alpha_s, \mu_R^2) + d\hat{\sigma}_{ab}^C(\alpha_s, \mu_F^2)) \Theta_{\text{obs}}^{(m)} + \int d\hat{\sigma}_{ab}^R(\alpha_s) \Theta_{\text{obs}}^{(m+1)} \right] \right\} + \dots$$

- In addition to the PDFs themselves, it is often useful to define a PDF luminosity

$$\frac{dL_{ij}}{d\hat{s} dy} = \frac{1}{s} \frac{1}{1 + \delta_{ij}} [f_i(x_1, \mu) f_j(x_2, \mu) + (1 \leftrightarrow 2)] . \quad \text{...or integrated over } y$$

# Some history: PDF4LHC

- In 2010, we carried out an exercise to which all PDF groups were invited to participate
- A comparison of NLO predictions for benchmark cross sections at the LHC (7 TeV) using MCFM with prescribed input files
- Benchmarks included
  - ◆  $W/Z$  production/rapidity distributions
  - ◆  $t\bar{t}$  production
  - ◆ Higgs production through  $gg$  fusion
    - ▲ masses of 120, 180 and 240 GeV
- PDFs used include CTEQ6.6, MSTW08, NNPDF2.0, HERAPDF1.0, ABKM09, GJR08
- Results in Higgs YR1 and YR2

## The PDF4LHC Working Group Interim Report

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All of the benchmark processes were to be calculated with the following settings:

1. at NLO in the  $\overline{MS}$  scheme
2. all calculation done in a the 5-flavor quark ZM-VFNS scheme, though each group uses a different treatment of heavy quarks
3. at a center-of-mass energy of 7 TeV
4. for the central value predictions, and for  $\pm 68\%$  and  $\pm 90\%$  c.l. PDF uncertainties
5. with and without the  $\alpha_s$  uncertainties, with the prescription for combining the PDF and  $\alpha_s$  errors to be specified
6. repeating the calculation with a central value of  $\alpha_s(m_Z)$  of 0.119.

arXiv:1101.0536v1 [hep-ph] 3 Jan 2011



# Followup in 2013

- Study of NNLO PDFs from 5 PDF groups (no new updates for JR)
  - ◆ drawing from what Graeme Watt had done at NNLO, but now including CT10 NNLO, and NNPDF2.3 NNLO
    - ▲ HERAPDF has upgraded to HERAPDF1.5; ABM09->ABM11
  - ◆ using a common values of  $\alpha_s$  (0.118) as a baseline; varying in range from 0.117 to 0.119)
  - ◆ including a detailed comparisons to LHC data which have provided detailed correlated systematic error information, keeping track of required systematic error shifts, normalizations, etc
    - ▲ ATLAS 2010 W/Z rapidity distributions
    - ▲ ATLAS 2010 inclusive jet cross section data
    - ▲ CMS 2011 W lepton asymmetry
    - ▲ LHCb 2010 W lepton rapidity distributions in forward region
- The effort was led by Juan Rojo and Pavel Nadolsky and has resulted in an independent publication
- The results from this paper will be utilized in a subsequent PDF4LHC document(s)
- ...and are now in YR3

# Benchmark paper

- Not officially a PDF4LHC document but will be used as input to future recommendations
- Comparisons only at NNLO, but NLO comparisons available at <http://nnpdf.hepforge.org/html/pdfbench/catalog>

arXiv:1211.5142v2 [hep-ph] 5 Apr 2013

CERN-PH-TH/2012-263  
Edinburgh 2012/21  
SMU-HEP-12-16  
LCTS/2012-26  
IFUM-1003-FT

## Parton distribution benchmarking with LHC data

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## Abstract:

We present a detailed comparison of the most recent sets of NNLO PDFs from the ABM, CT, HERAPDF, MSTW and NNPDF collaborations. We compare parton distributions at low and high scales and parton luminosities relevant for LHC phenomenology. We study the PDF dependence of LHC benchmark inclusive cross sections and differential distributions for electroweak boson and jet production in the cases in which the experimental covariance matrix is available. We quantify the agreement between data and theory by computing the  $\chi^2$  for each data set with all the various PDFs. PDF com-

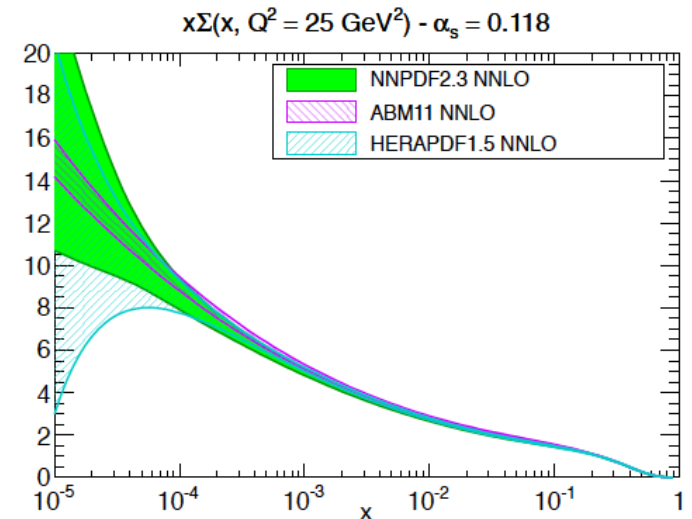
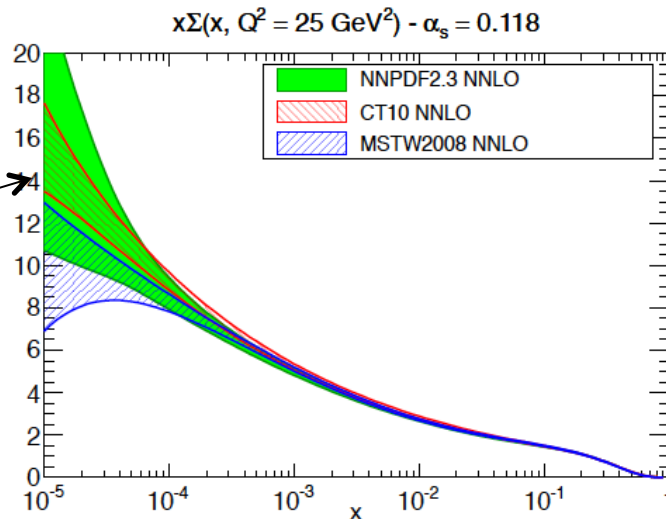
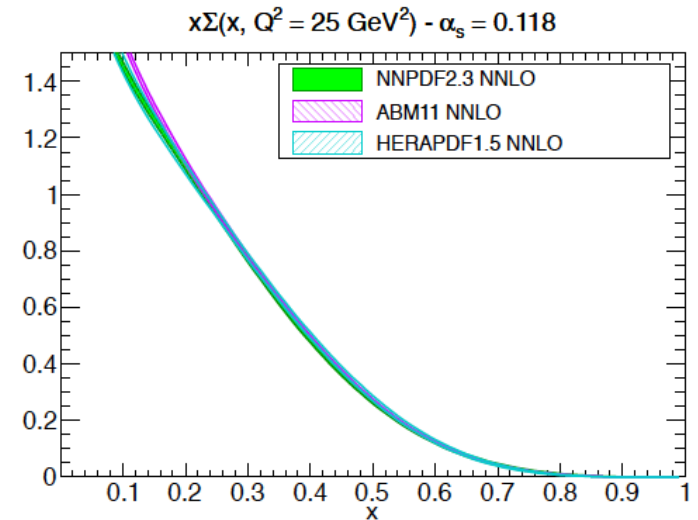
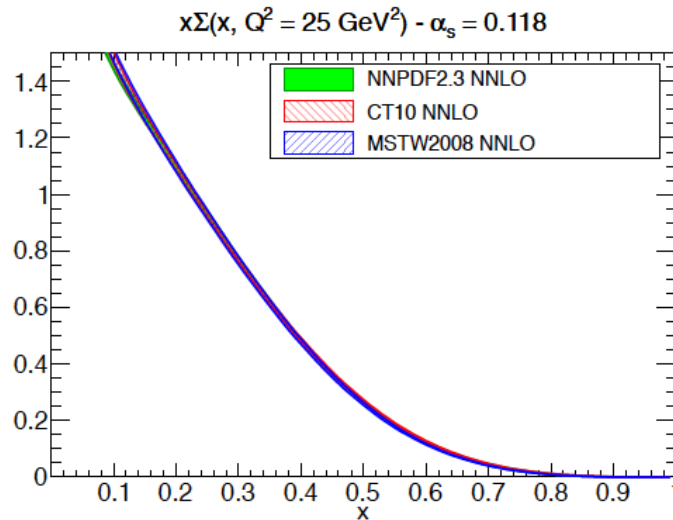
# PDF comparisons

## quark singlet PDFs

...results for other values of  $\alpha_s$  and at NLO available on the HEPFORGE website

good agreement for all sets for quark singlet distribution

note the blowup of uncertainties for low  $x$



# Comparison of PDFs

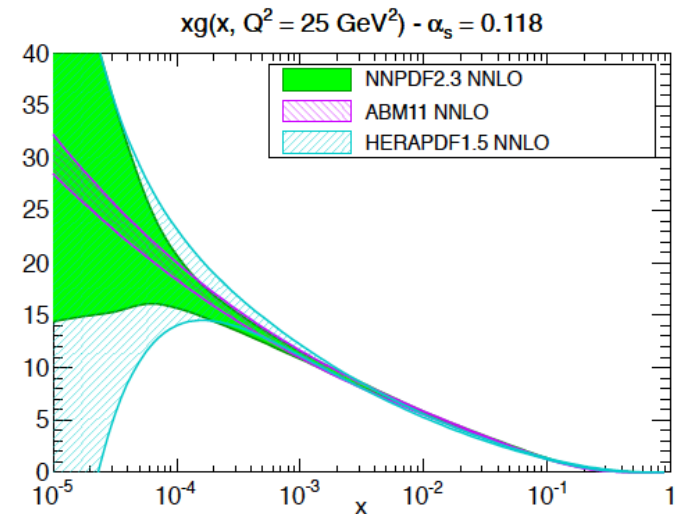
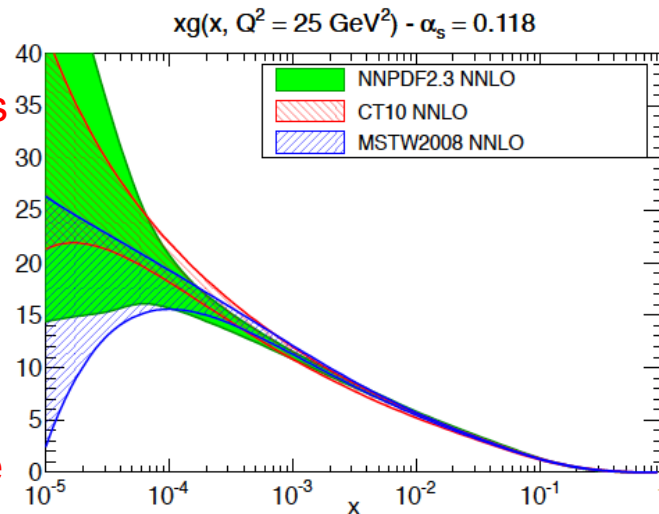
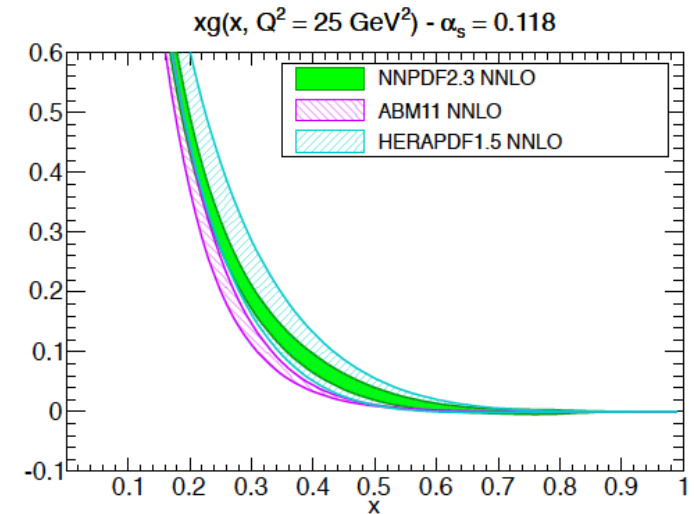
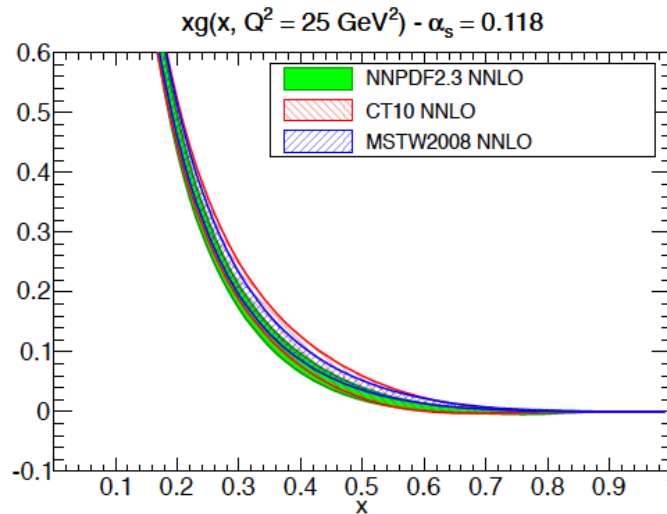
CT10, MSTW08  
and NNPDF2.3  
gluon distributions  
all in reasonable,  
but not perfect,  
agreement

The 1-sigma  
uncertainty  
bands overlap  
for all values of  
 $x$

again, uncertainties  
blow up for small  $x$

HERAPDF  
uncertainties  
somewhat larger  
at low  $x$ ; noticeably  
larger at high  $x$  due  
to lack of collider  
jet data

gluon PDF





# PDF luminosities

gluon-gluon and  
gluon-quark  
luminosities in  
reasonable, but  
again not perfect,  
agreement  
for CT10,  
MSTW08 and  
NNPDF2.3 for full  
range of invariant  
masses

HERAPDF1.5  
uncertainties larger in  
general

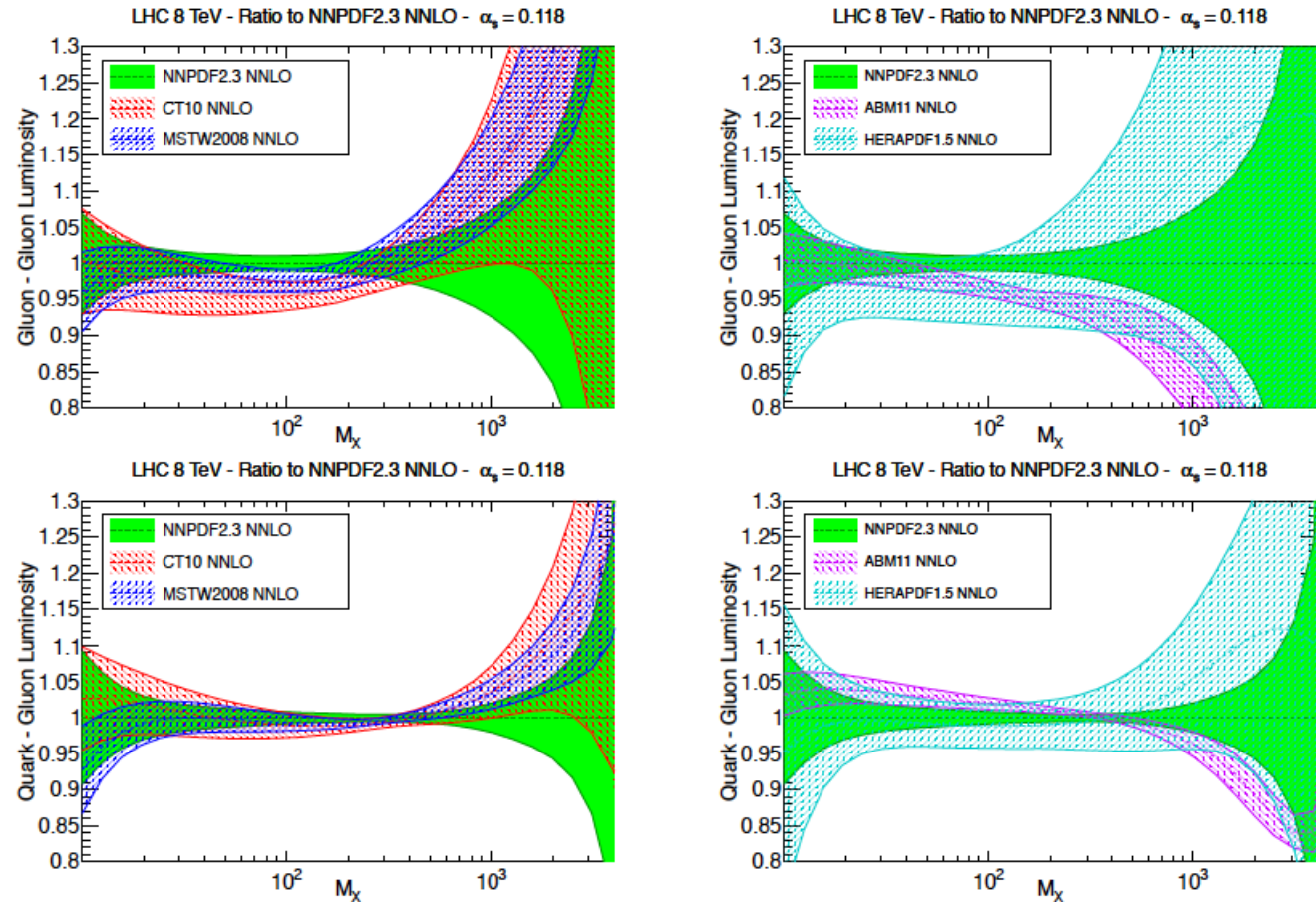


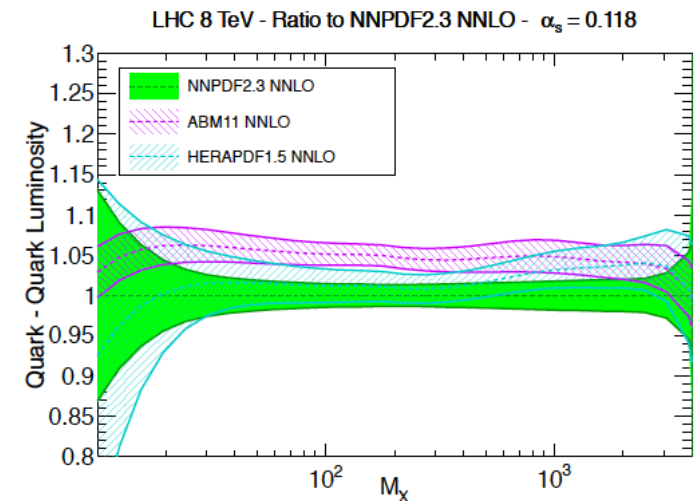
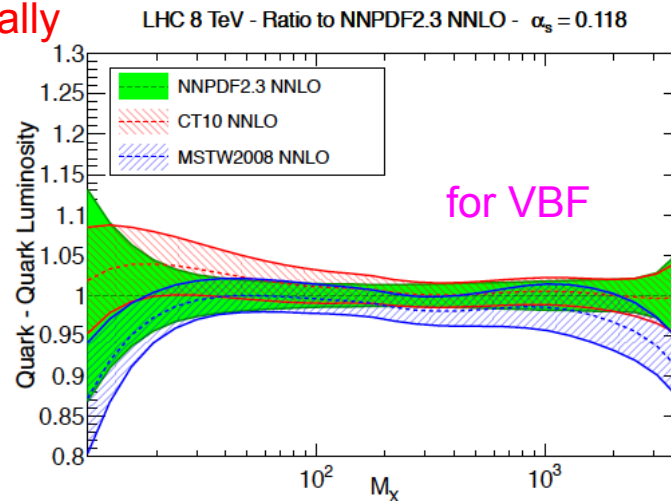
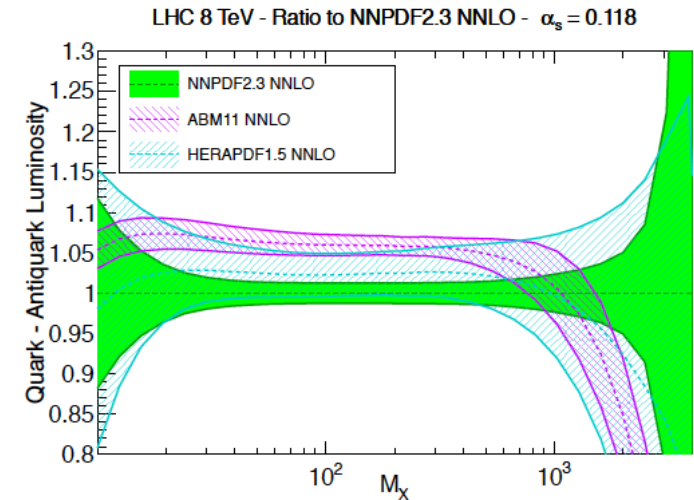
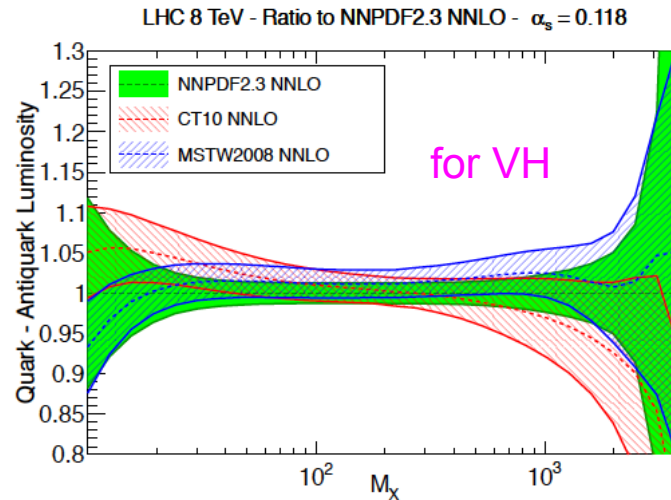
Figure 6: The gluon-gluon (upper plots) and quark-gluon (lower plots) luminosities, Eq. (2), for the production of a final state of invariant mass  $M_X$  (in GeV) at LHC 8 TeV. The left plots show the comparison between NNPDF2.3, CT10 and MSTW08, while in the right plots we compare NNPDF2.3, HERAPDF1.5 and MSTW08. All luminosities are computed at a common value of  $\alpha_s = 0.118$ .

# PDF luminosities

quark-quark and quark-antiquark

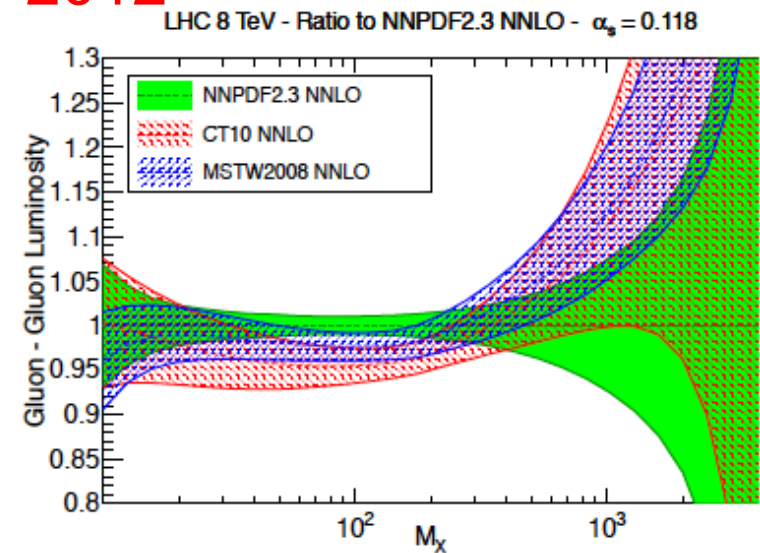
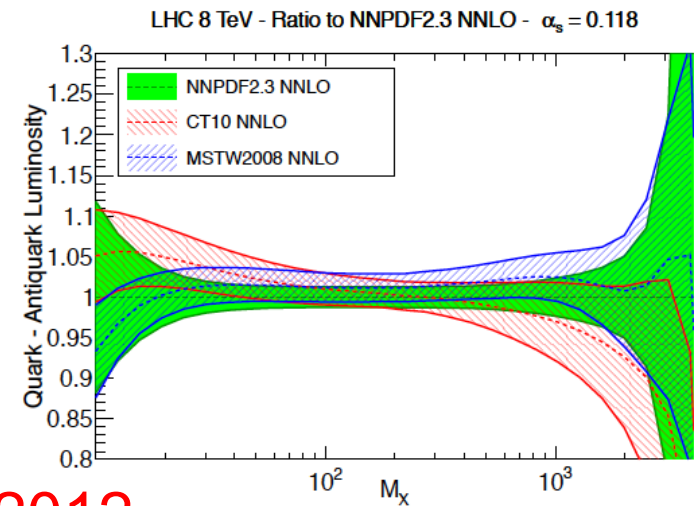
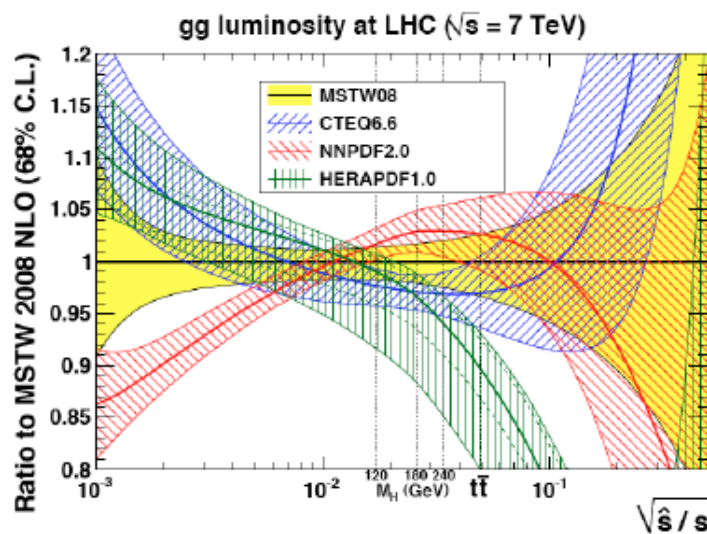
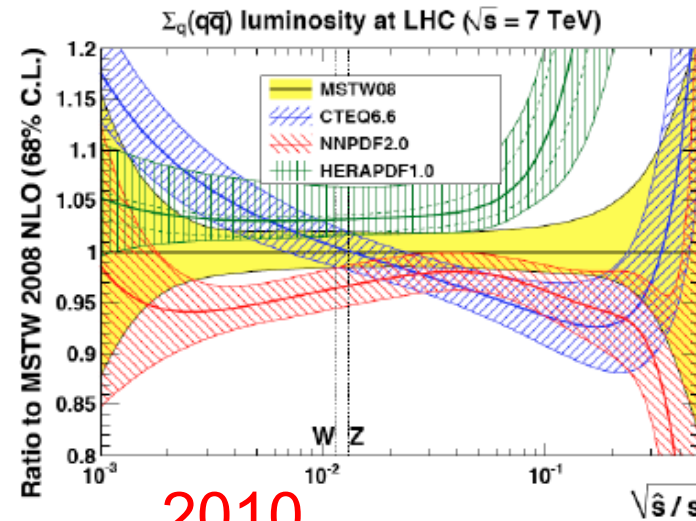
quark-antiquark luminosities for CT10, MSTW08 and NNPDF2.3 overlap almost 100% in W/Z range

ABM11 systematically larger at small mass, then falls off more rapidly at high mass



# Uncertainties have improved

- ...with additional data and in going from NLO to NNLO



# Compare relative luminosity uncertainties

good agreement in size of uncertainties between the 3 global PDFs

larger uncertainties of HERAPDF1.5 apparent

ABM11 uncertainties smaller at high mass

note the uncertainties starting to blow up at low mass; low mass  $x$  values become moderate mass  $x$  values at 100 TeV

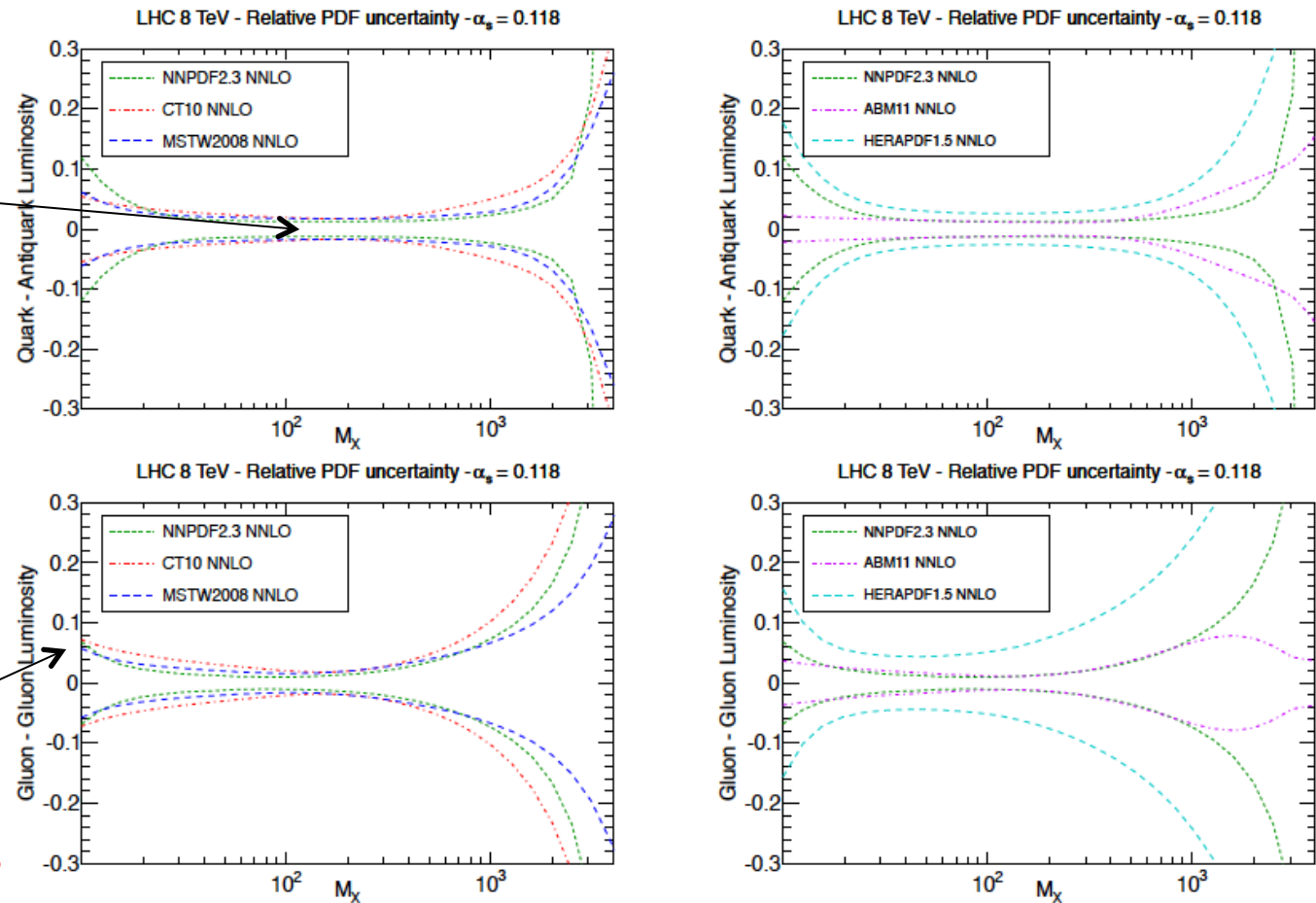
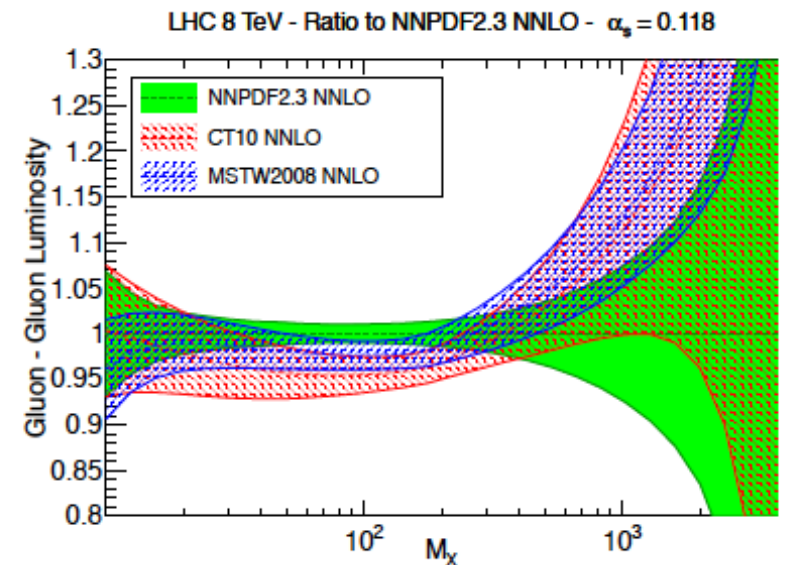
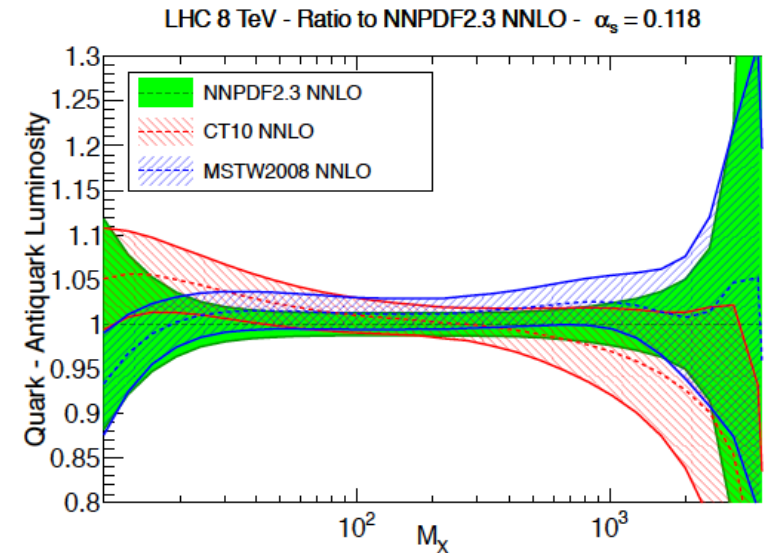


Figure 8: The relative PDF uncertainties in the quark-antiquark luminosity (upper plots) and in the gluon-gluon luminosity (lower plots), for the production of a final state of invariant mass  $M_X$  (in GeV) at the LHC 8 TeV. All luminosities are computed at a common value of  $\alpha_s = 0.118$ .



# NNLO PDF uncertainties

- Nice convergence for qQ PDF luminosities in range of W/Z masses (at 8 TeV)
  - ♦ but not so for lower masses
- Also not so for gg PDF luminosities around 125 GeV at 8 TeV
  - ♦ better overlap, but with larger uncertainties, at low mass
  - ♦ PDF+as error dominant theory error
- Project started at Les Houches
  - ♦ understand differences in central luminosity value from CT10, MSTW08, NNPDF2.3 and HERAPDF1.5
  - ♦ progress report in Les Houches
  - ♦ meetings continuing



# $\alpha_s(m_Z)$

- Right now the Higgs Cross Section Working Group is using a mean value for  $\alpha_s(m_Z)$  of 0.118 with 90% CL error of 0.002 (68%CL error of 0.012), or an inflation of the world average uncertainties; the  $\alpha_s$  error is added in quadrature with the PDF error
- The world average is dominated by lattice results
- Are the lattice results are robust enough, so that an uncertainty of 0.012 (at 68% CL) may be an overestimate? Will the uncertainty in  $\alpha_s$  be a non-issue at the time of any 100 TeV collider\*

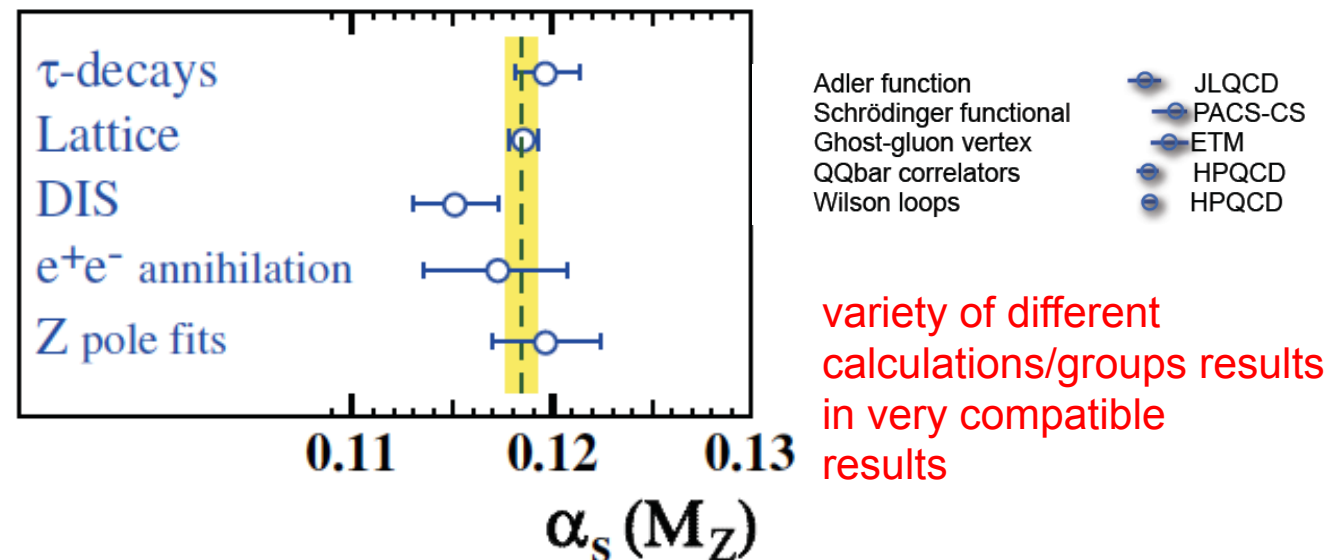


Figure 1-1. Summary of values of  $\alpha_s(M_Z^2)$  obtained for various sub-classes of measurements. The world average value of  $\alpha_s(M_Z^2) = 0.1184 \pm 0.0007$  is indicated by the dashed line and the shaded band. Figure taken from [1].

\*will I be alive to worry about the issue?

# On to 100 TeV

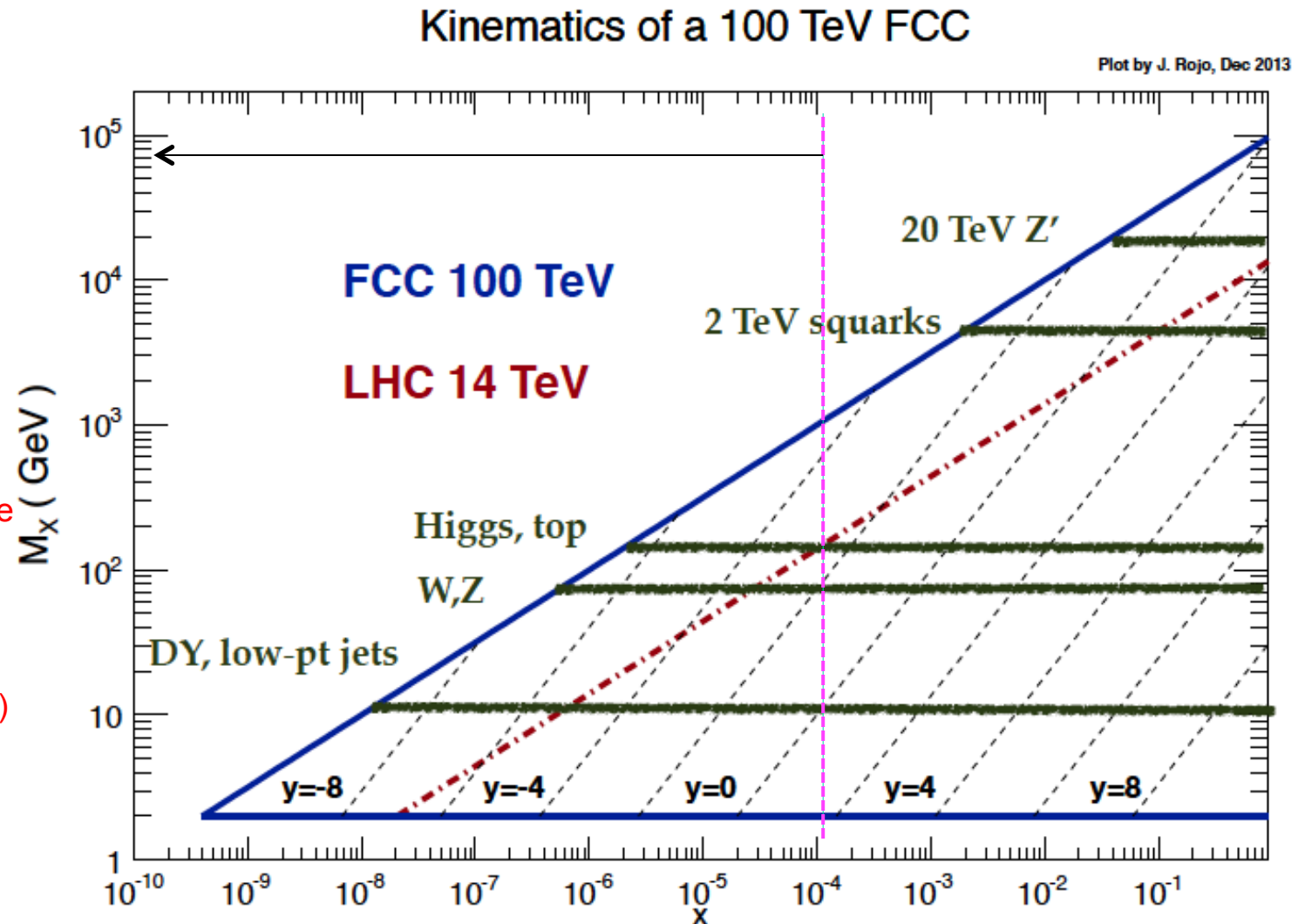
will access  
smaller  $x$ ,  
larger  $Q^2$

currently  
have  
no  
constraints  
on PDFs  
for  $x$   
values below  
 $1E-4$

we don't know where  
at low  $x$ , BFKL  
effects start to  
become important

poor constraints (still)  
as well for  
high  $x$  PDFs

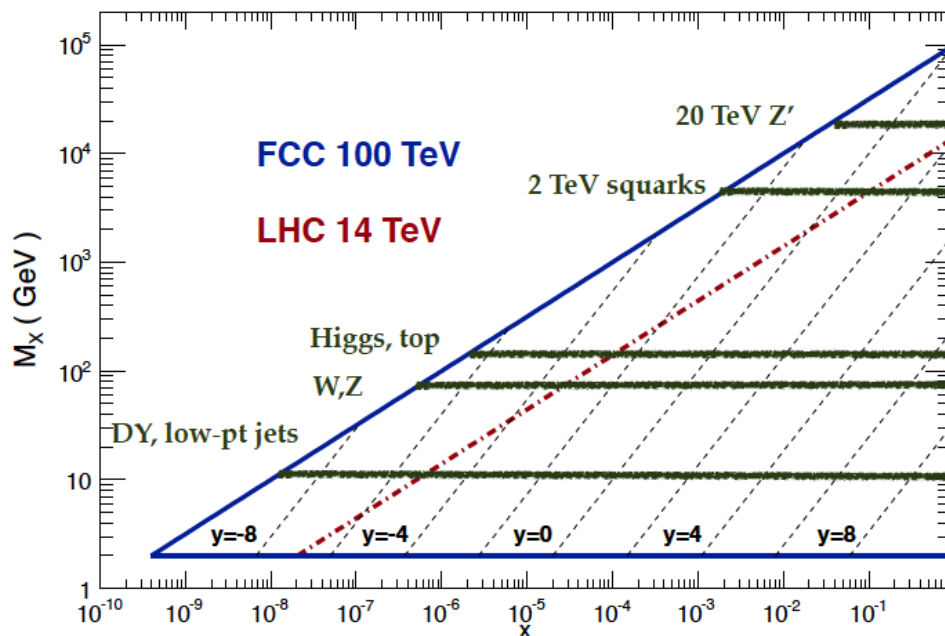
at high masses  
( $Q^2$ ), rely on  
DLAP evolution; we know at large  $Q^2$ ,  
EW effects also become important



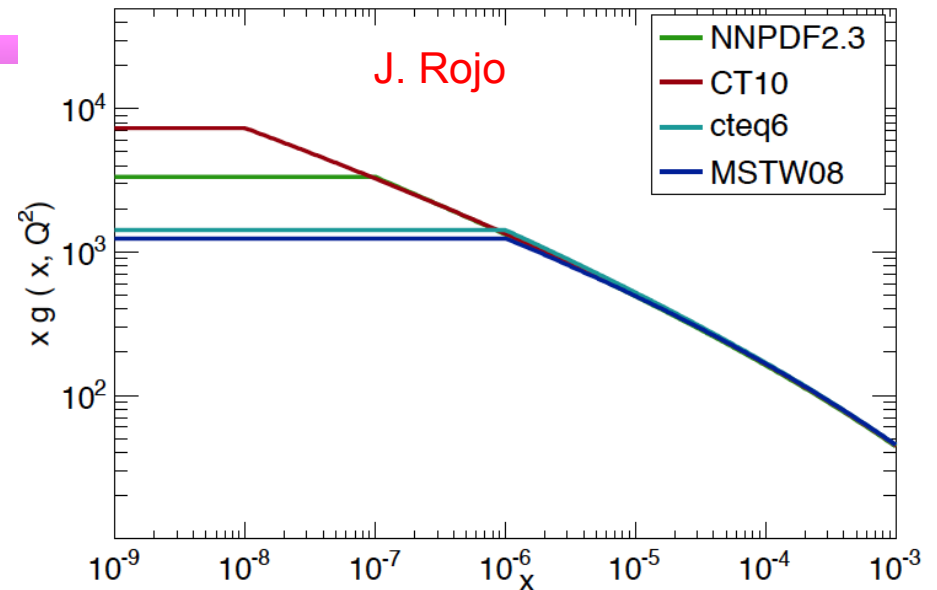
- Information for PDFs below  $x$  value of  $1E-4$  very sparse
- Most current PDFs cut off at some low  $x$  value
- Can extrapolate, but it is just that, extrapolation, perhaps based on some Regge arguments

Kinematics of a 100 TeV FCC

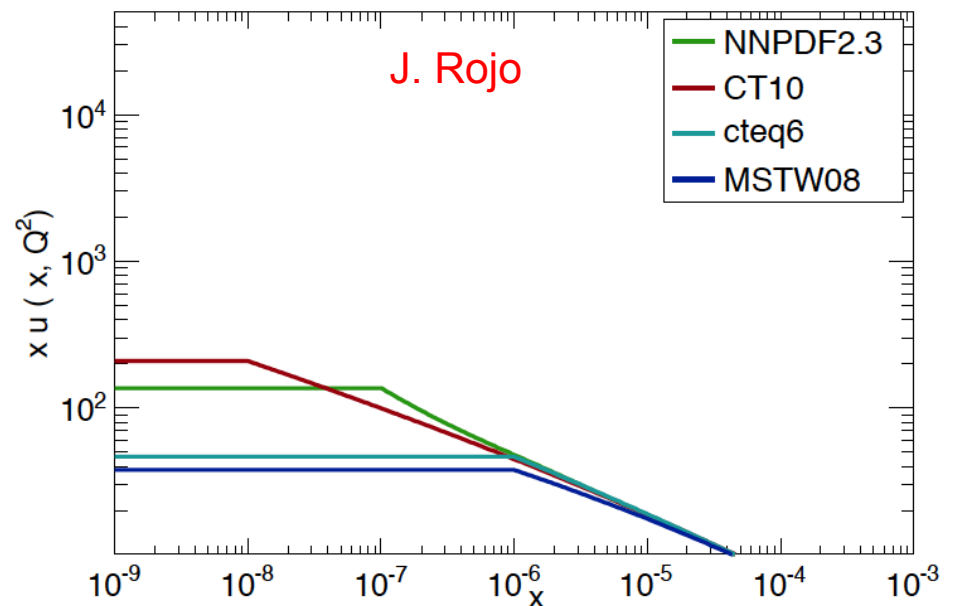
Plot by J. Rojo, Dec 2013



Small- $x$  NNLO PDFs for FCC studies



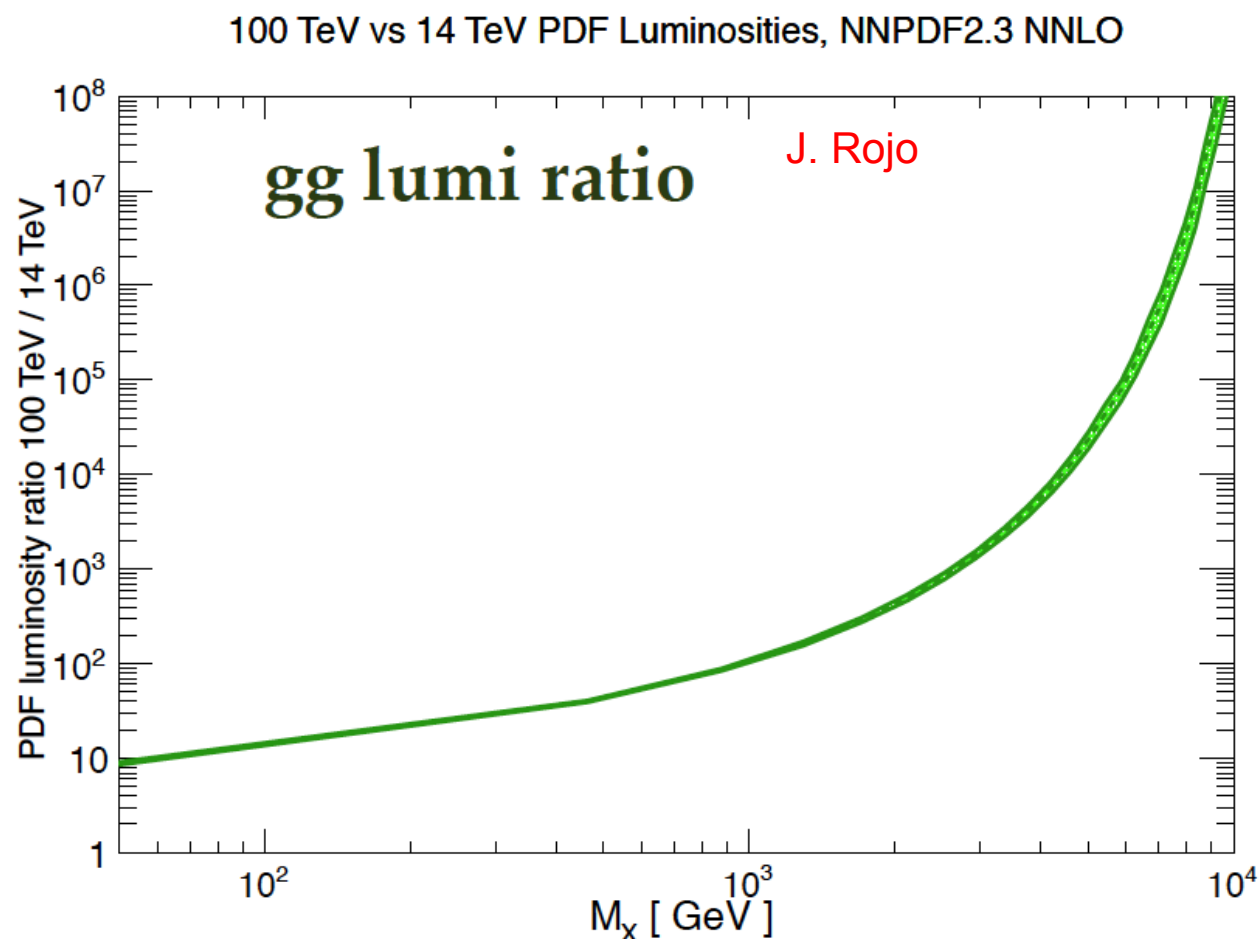
Small- $x$  NNLO PDFs for FCC studies





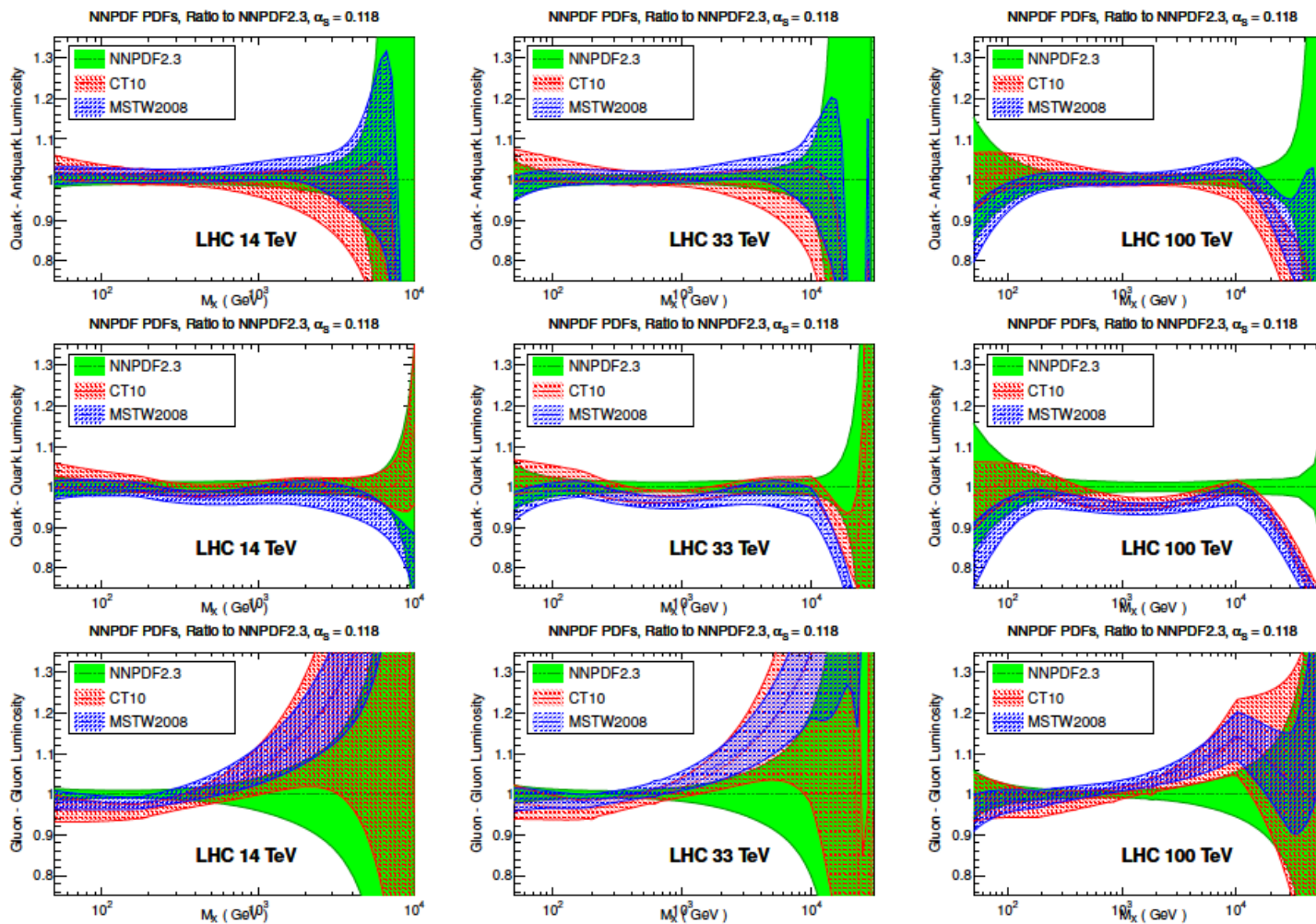
# PDF luminosities at 100 TeV

- gg luminosity ratio at order of 100 at TeV scale 1E8 at 10 TeV scale
- Similar increases for other PDFs



# PDFs at higher energies: as part of the Snowmass exercise

PDFs are HERA/fixed target dominated for  $x \sim 0.05-0.1$ ; LHC data at 14 TeV offers opportunity for shrinking uncertainties in new physics search range



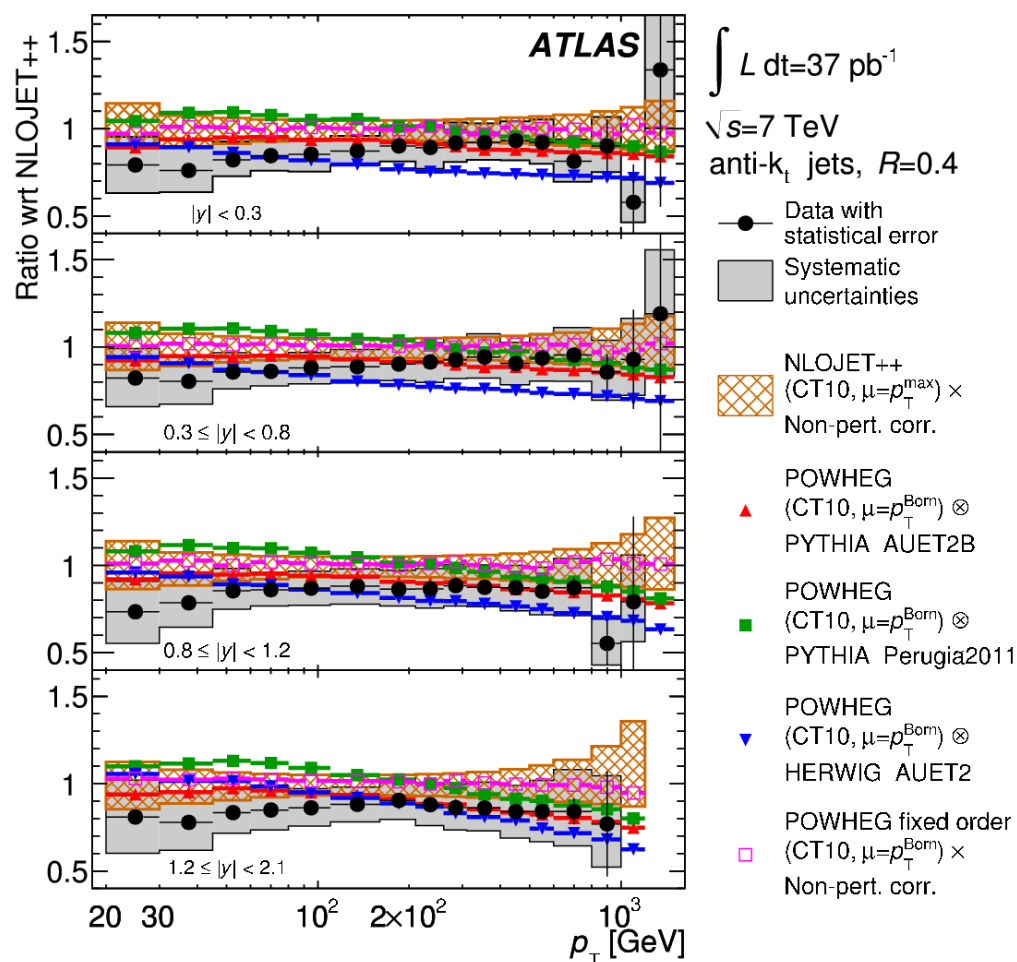
high masses  
always a  
problem, with  
current uncer-  
tainties

low masses  
become a  
problem at  
very high  
energy colliders

# LHC data in global PDF fits

- LHC data will become increasingly important in global fits
- Not just inclusive jet data but for processes such as inclusive photon production, Drell-Yan, W/Z rapidity,  $t\bar{t}$  mass and rapidity
- For any process to be used in a global PDF fit, correlated systematic errors must be provided
- 2010 inclusive jet data from ATLAS provides no discrimination
- Data from 2011/2012, with increased statistics and improved systematics may
- **BUT**, Note that LHC data is competing against HERA data where two experiments have been combined and statistical and systematic errors are a few percent
  - ♦ may be difficult to compete in the precision physics range a la  $gg \rightarrow \text{Higgs}$

- 2010 ATLAS data lies below NLOJET++ prediction using CT10 at high  $p_T/y$
- difference if Powheg used instead of fixed order? extra radiation?



# New PDF4LHC exercise

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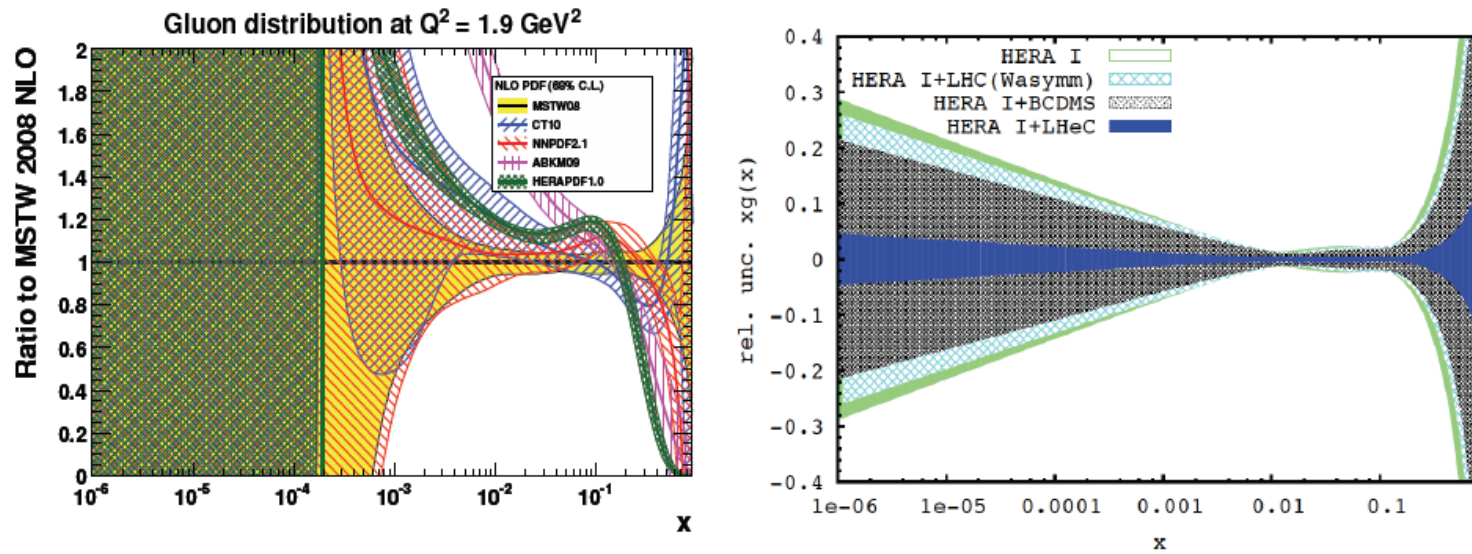
- Lay out a coherent coordinated plan for QCD(+EW) measurements, among ATLAS, CMS and LHCb, that can reduce PDF systematics using LHC data
  - ◆ again systematic errors will be very important
- Wiki is now up, PDF4LHC meeting in May

[https://twiki.cern.ch/twiki/bin/view/  
PDF4LHC/WebHome](https://twiki.cern.ch/twiki/bin/view/PDF4LHC/WebHome)



# LHeC

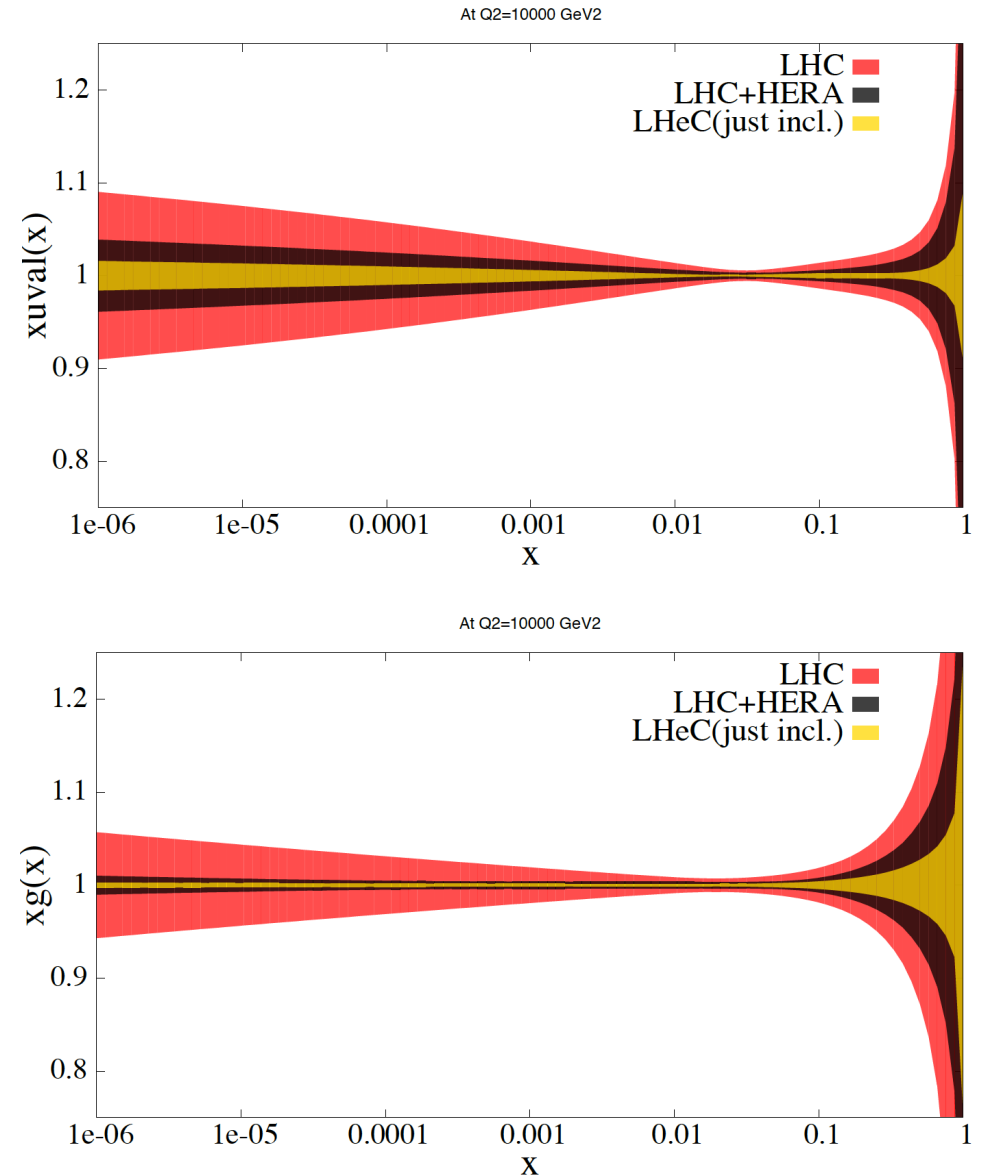
- Further improvement might come from an LHeC



**Figure 1-5.** *Uncertainty of the gluon distribution at  $Q^2 = 1.9 \text{ GeV}^2$  as a function of Bjorken  $x$ , see text. The LHeC PDF set, corresponding to the inner blue error band, is available on LHAPDF.*

# Snowmass exercise with LHC data

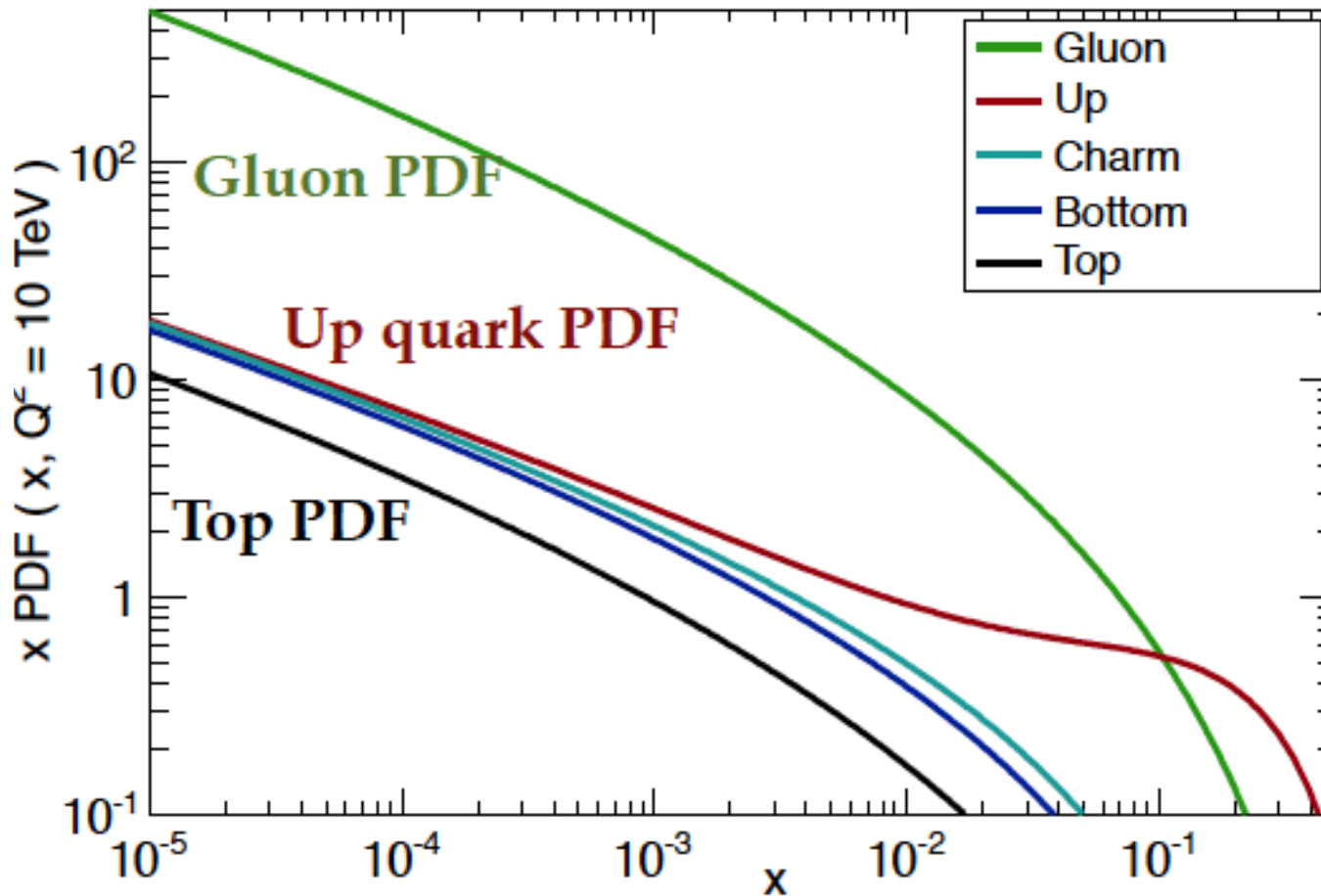
- Use current LHC data in global PDF fits, find no great restraint
  - ◆ impact comes from inclusion of HERA data
- With  $100 \text{ fb}^{-1}$ , will have precision measurements of DY production from 60 to 1500 GeV, with systematic errors half of the current values, stat errors 5% at high mass
  - ◆ Phase 1 ( $300 \text{ fb}^{-1}$ ) and phase 2 ( $3000 \text{ fb}^{-1}$ ) will provide strong improvement in PDF uncertainties at high mass (BSM search region)



# Top quark PDFs

- At very high  $Q^2$ , top mass becomes small, and top PDFs may need to be taken into account

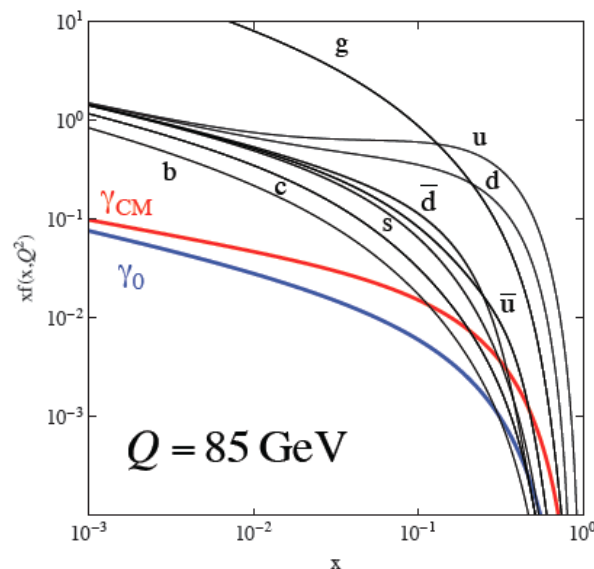
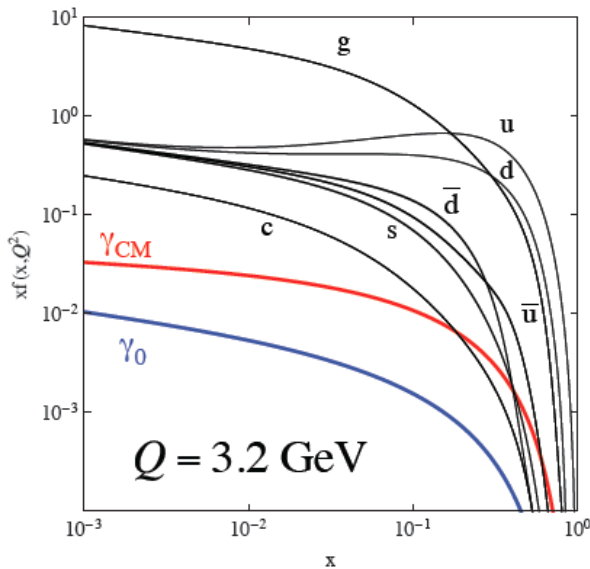
NNPDF2.3 NNLO  $N_F = 6$



see talk of  
Ismail Ahmed

# Photon PDFs

- Photon PDFs: photon PDFs can be larger than antiquark distributions at high  $x$ ; the LHC is a  $\gamma\gamma$  collider; even more true of a 100 TeV collider
- NNPDF has developed photon PDFs + QED corrections (in addition to MRST2004QED)
- CT10 in progress (see talk of C. Schmidt at DIS2014)
  - ♦ fitting to photon production in DIS



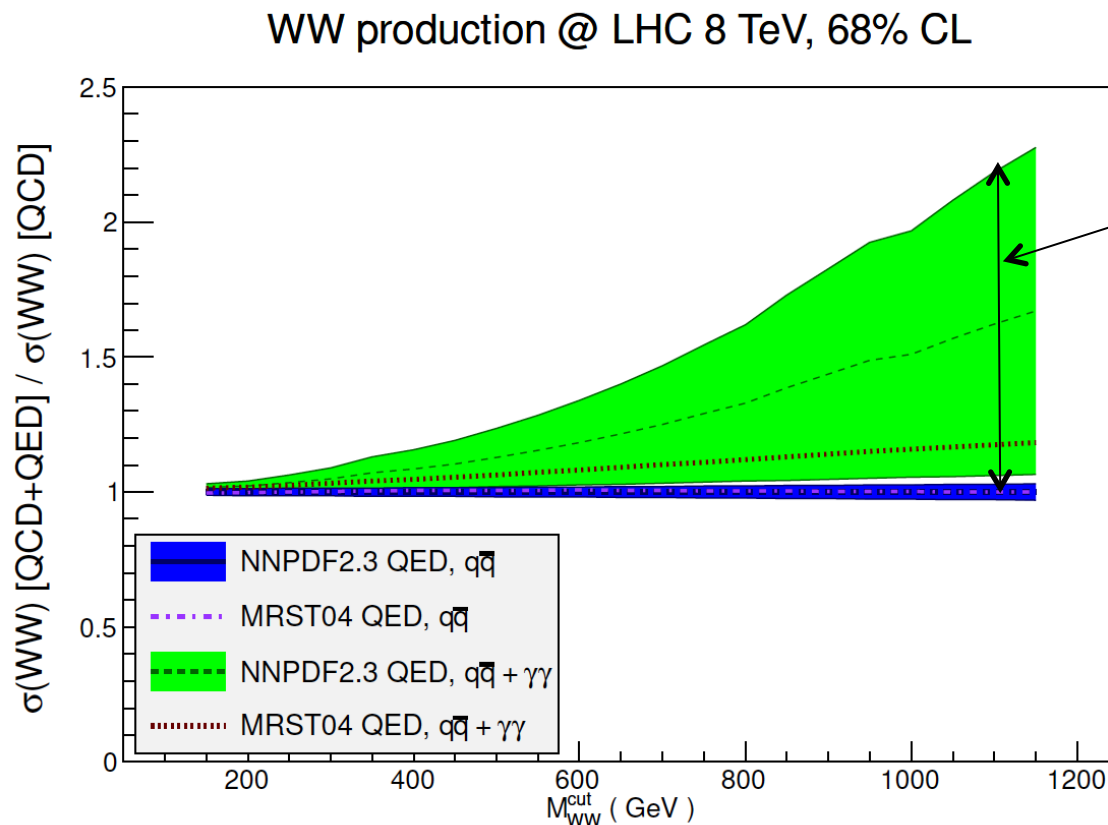
$\gamma$  momentum fraction:

$p^\gamma(Q)$	$\gamma(x, Q_0) = 0$	$\gamma(x, Q_0)_{\text{CM}}$
$Q = 3.2 \text{ GeV}$	0.05%	0.34%
$Q = 85 \text{ GeV}$	0.22%	0.51%

allow for non-perturbative component of photon at  $Q_0$

# WW production and the photon PDF

- photon-induced WW production can contribute significantly at high mass
- ...and understanding high mass WW production will be important in the next run
- a better understanding of the photon PDF is thus crucial
  - ◆ first steps taken with LHC DY data

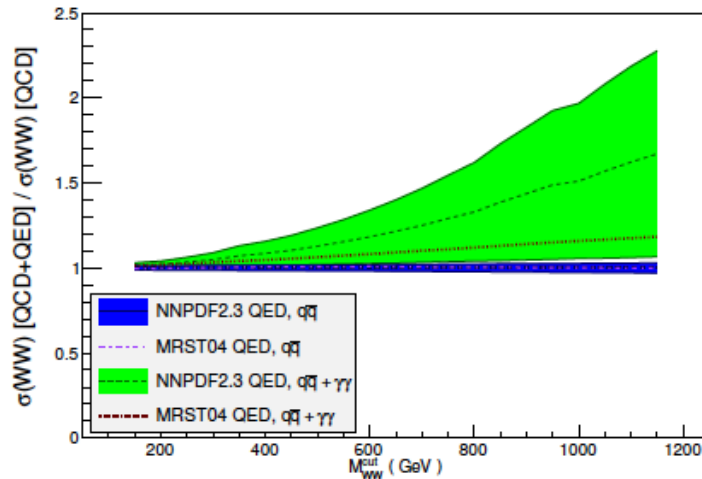


with currently a  
very large  
uncertainty due to  
lack of  
knowledge of the  
photon PDF

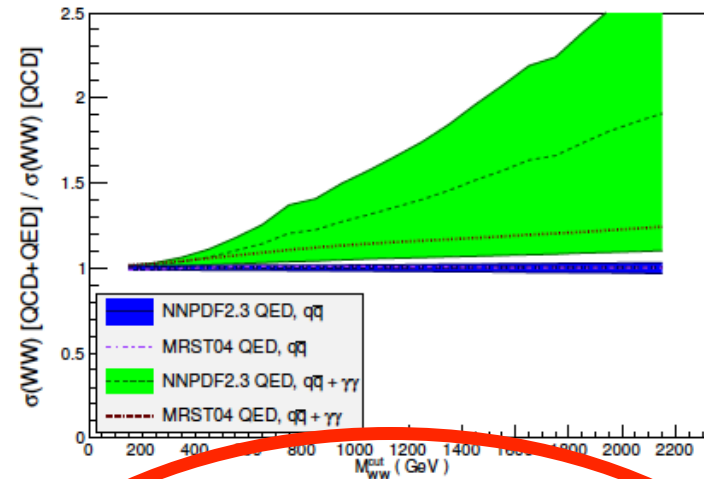


# QED corrections

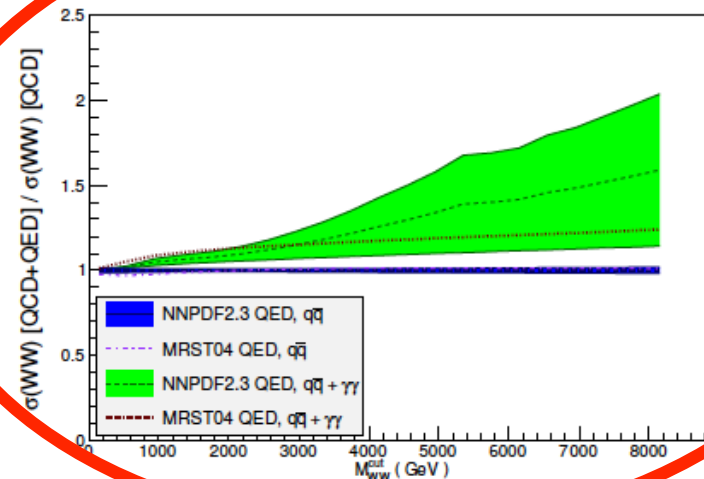
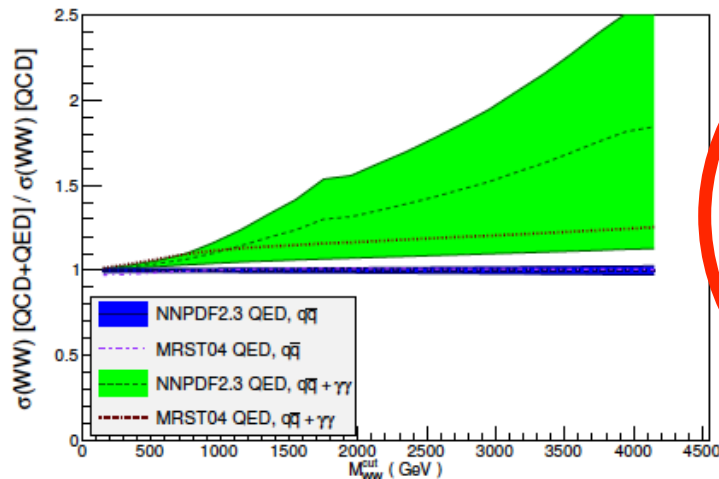
- Photon PDFs will become important as energies increase for processes such as  $\gamma\gamma \rightarrow WW$



WW production @ LHC 33 TeV, 68% CL



WW production @ LHC 100 TeV, 68% CL



# EW corrections

---

- At high  $Q^2$ , logs of  $\alpha \ln(Q^2/m_W^2)$  become large; EW corrections become as large as higher order QCD corrections
- Need EW evolution for PDFs
  - ◆ W and Z PDFs
  - ◆ Ciafaloni and Comelli, 2002, 2005
- ...in Les Houches proceedings, a *dictionary* for QCD+EW corrections has been provided by Stefan Dittmaier

# Meta-PDFs:arXiv:1401.0013

- Take NNLO PDFs

<i>NNLO</i>	<i>Initial scale</i>	$a_s$	<i>Error type</i>	<i>Error sets</i>
<i>CT10</i>	<i>1.3</i>	<i>0.118</i>	<i>Hessian</i>	<i>50</i>
<i>MSTW'08</i>	<i>1.0</i>	<i>0.1171</i>	<i>Hessian</i>	<i>40</i>
<i>NNPDF2.3</i>	<i>1.414</i>	<i>0.118</i>	<i>MC</i>	<i>100</i>

- Choose a meta-parametrization of PDFs at initial scale of 8 GeV (away from thresholds) for 9 PDF flavors (66 parameters in total)

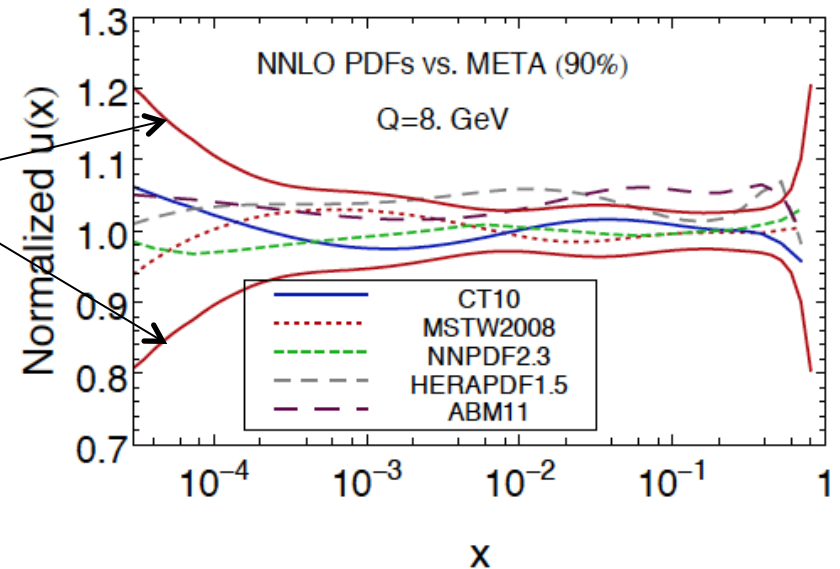
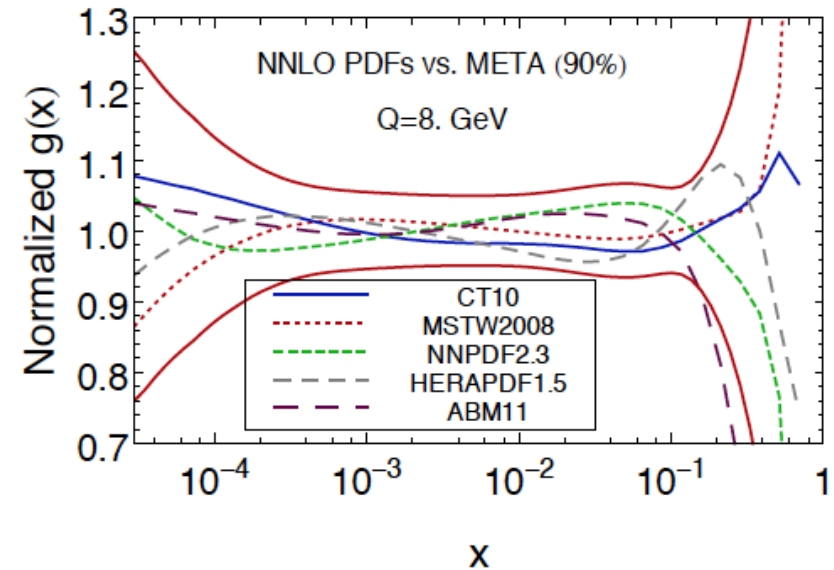
$$f(x, Q_0; \{a\}) = e^{a_1} x^{a_2} (1 - x)^{a_3} e^{\sum_{i \geq 4} a_i [T_{i-3}(y(x)) - 1]}$$

- Generate MC replicas for all 3 groups and merge with equal weights, finding meta parameters for each of the replicas by fitting PDFs in x ranges probed at LHC
- Construct 50 eigenvectors using Hessian method
- These 50 eigenvectors provide a very good representation of the PDF uncertainties for all of the 3 PDF error families above

# meta-PDFs

- The meta-PDFs provide both an average of the chosen PDFs, as well as a good estimation of the total PDF uncertainty

meta-PDF uncertainty band



# Higgs observables

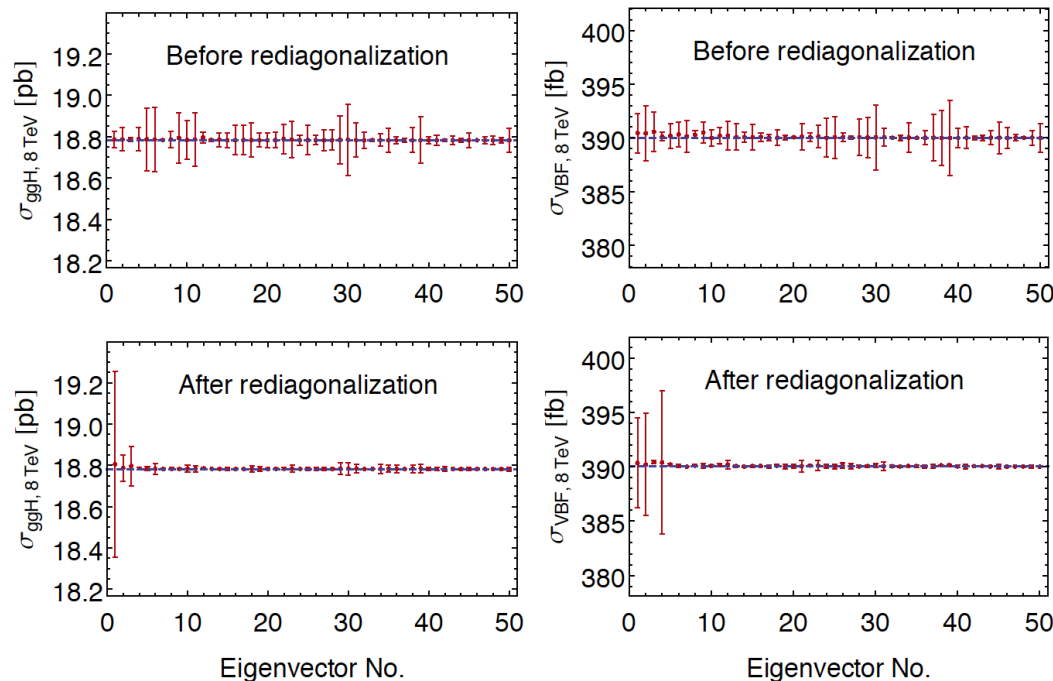
- Select global set of Higgs cross sections at 8 and 14 TeV (46 observables in total; more can be easily added if there is motivation)

production channel	$\sigma(inc.)$	$\sigma( y_H  > 1)$	$\sigma(p_{T,H} > m_H)$	scales
$gg \rightarrow H$	iHixs1.3 [32] at NNLO	MCFM6.3 [33] at LO	—	$m_H$
$b\bar{b} \rightarrow H$	iHixs at NNLO	—	—	$m_H$
VBF	VBFNLO2.6 [34] at NLO	same	same	$m_W$
$HZ$	VHNNLO1.2 [35] at NNLO	CompHEP4.5 [36] at LO	CompHEP at LO	$m_Z + m_H$
$HW^\pm$	VHNNLO at NNLO	—	—	$m_W + m_H$
$HW^+$	CompHEP at LO	same	same	$m_W + m_H$
$HW^-$	CompHEP at LO	same	same	$m_W + m_H$
$H + 1jet$	MCFM at LO	same	same	$m_H$
$Ht\bar{t}$	MCFM at LO	CompHEP at LO	CompHEP at LO	$2m_t + m_H$
$HH$	Hpair [37] at NLO	—	—	$2m_H$



# Data set diagonalization (arXiv:0904.2424)

- There are 50 eigenvectors, but can re-diagonalize the Hessian matrix to pick out directions important for the Higgs observables listed on previous page; with rotation of basis, 50 eigenvectors become 6



J. Gao,  
J. Huston  
P. Nadolsky  
(in progress)

**It's possible to define a few eigenvectors which completely encompass the PDF and  $\alpha_s$  uncertainties for CT10, MSTW08 and NNPDF2.3 for Higgs production for 8-14 TeV; no reason this cannot be expanded to 100 TeV**

# arXiv:1004.4624

- Treat  $\alpha_s$  input as another eigenvector;  $\alpha_s$  and PDF uncertainties can be added in quadrature ( $\alpha_s(m_Z)=0.118\pm0.0012$ )
- So 7 eigenvectors to represent all PDF+ $\alpha_s$  uncertainty

LHC	$\Delta\alpha_s(M_Z)$	GGH inc.	GGH 0j exc.	GGH 1j exc.	GGH 2j inc.	VBF inc.
LHC 8 TeV	+1 $\sigma$	2.2%	1.6%	3.0%	4.8%	-0.23%
	-1 $\sigma$	-2.2%	-1.6%	-2.8%	-4.8%	0.11%
LHC 14 TeV	+1 $\sigma$	2.1%	1.4%	2.6%	4.5%	0.05%
	-1 $\sigma$	-2.0%	-1.4%	-2.5%	-4.4%	-0.09%

❖ using PDF  $\alpha_s$  series of the META PDFs

# NNLO QCD+NLO EW wishlist

Process	known	desired	details
H	$d\sigma$ @ NNLO QCD $d\sigma$ @ NLO EW finite quark mass effects @ NLO	$d\sigma$ @ NNNLO QCD + NLO EW MC@NNLO finite quark mass effects @ NNLO	H branching ratios and couplings
H + j	$d\sigma$ @ NNLO QCD (g only) $d\sigma$ @ NLO EW finite quark mass effects @ LO	$d\sigma$ @ NNLO QCD + NLO EW finite quark mass effects @ NLO	H $p_T$
H + 2j	$\sigma_{\text{tot}}(\text{VBF})$ @ NNLO(DIS) QCD $d\sigma(\text{gg})$ @ NLO QCD $d\sigma(\text{VBF})$ @ NLO EW	$d\sigma$ @ NNLO QCD + NLO EW	H couplings
H + V	$d\sigma$ @ NNLO QCD $d\sigma$ @ NLO EW	with $H \rightarrow b\bar{b}$ @ same accuracy	H couplings
t $\bar{t}$ H	$d\sigma(\text{stable tops})$ @ NLO QCD	$d\sigma(\text{top decays})$ @ NLO QCD + NLO EW	top Yukawa coupling
HH	$d\sigma$ @ LO QCD (full $m_t$ dependence) $d\sigma$ @ NLO QCD (infinite $m_t$ limit)	$d\sigma$ @ NLO QCD (full $m_t$ dependence) $d\sigma$ @ NNLO QCD (infinite $m_t$ limit)	Higgs self coupling

- LO  $\equiv \mathcal{O}(1)$ ,
- NLO QCD  $\equiv \mathcal{O}(\alpha_s)$ ,
- NNLO QCD  $\equiv \mathcal{O}(\alpha_s^2)$ ,
- NLO EW  $\equiv \mathcal{O}(\alpha)$ ,
- NNNLO QCD  $\equiv \mathcal{O}(\alpha_s^3)$ ,
- NNLO QCD+EW  $\equiv \mathcal{O}(\alpha_s\alpha)$

Table 1: Wishlist part 1 – Higgs ( $V = W, Z$ )

S. Dittmaier, N. Glover, J. Huston

In the writeup, we justified the requested precision based on current/extrapolated experimental errors

# NNLO QCD + NLO EWK wishlist

Process	known	desired	details
$t\bar{t}$	$\sigma_{\text{tot}}$ @ NNLO QCD $d\sigma(\text{top decays})$ @ NLO QCD $d\sigma(\text{stable tops})$ @ NLO EW	$d\sigma(\text{top decays})$ @ NNLO QCD + NLO EW	precision top/QCD, gluon PDF, effect of extra radiation at high rapidity, top asymmetries
$t\bar{t} + j$	$d\sigma(\text{NWA top decays})$ @ NLO QCD	$d\sigma(\text{NWA top decays})$ @ NNLO QCD + NLO EW	precision top/QCD top asymmetries
single-top	$d\sigma(\text{NWA top decays})$ @ NLO QCD	$d\sigma(\text{NWA top decays})$ @ NNLO QCD (t channel)	precision top/QCD, $V_{tb}$
dijet	$d\sigma$ @ NNLO QCD (g only) $d\sigma$ @ NLO weak	$d\sigma$ @ NNLO QCD + NLO EW	Obs.: incl. jets, dijet mass → PDF fits (gluon at high x) → $\alpha_s$ CMS <a href="http://arxiv.org/abs/1212.6660">http://arxiv.org/abs/1212.6660</a>
3j	$d\sigma$ @ NLO QCD	$d\sigma$ @ NNLO QCD + NLO EW	Obs.: $R3/2$ or similar → $\alpha_s$ at high scales dom. uncertainty: scales CMS <a href="http://arxiv.org/abs/1304.7498">http://arxiv.org/abs/1304.7498</a>
$\gamma + j$	$d\sigma$ @ NLO QCD $d\sigma$ @ NLO EW	$d\sigma$ @ NNLO QCD +NLO EW	gluon PDF $\gamma + b$ for bottom PDF

Table 2: Wishlist part 2 – jets and heavy quarks

S. Dittmaier, N. Glover, J. Huston

# NNLO QCD + NLO EWK wishlist

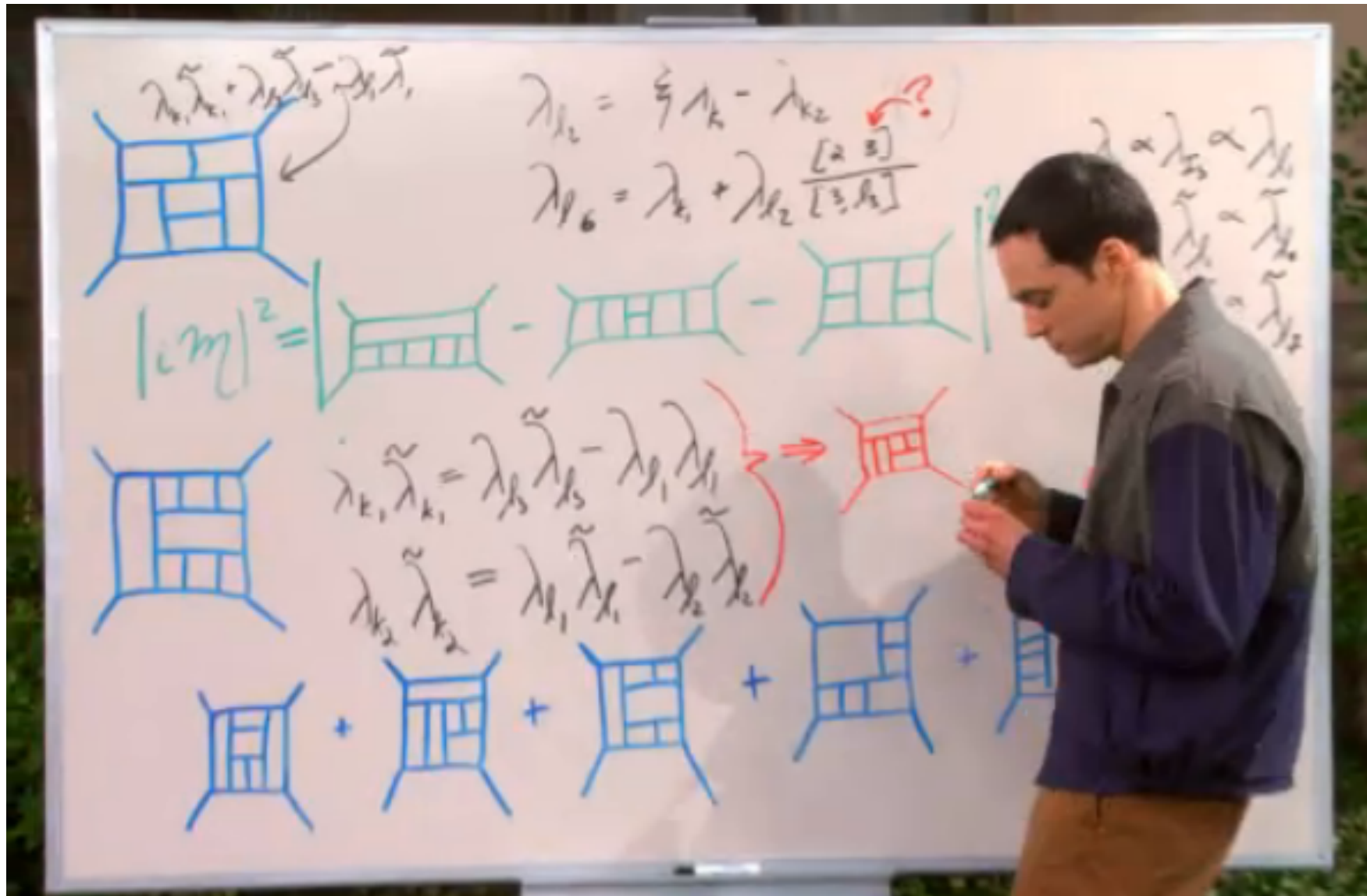
S. Dittmaier,  
N. Glover,  
J. Huston

Process	known	desired	details
V	$d\sigma(\text{lept. V decay}) @ \text{NNLO QCD}$ $d\sigma(\text{lept. V decay}) @ \text{NLO EW}$	$d\sigma(\text{lept. V decay})$ @ NNNLO QCD + NLO EW MC@NNLO	precision EW, PDFs
V + j	$d\sigma(\text{lept. V decay}) @ \text{NLO QCD}$ $d\sigma(\text{lept. V decay}) @ \text{NLO EW}$	$d\sigma(\text{lept. V decay})$ @ NNLO QCD + NLO EW	Z + j for gluon PDF W + c for strange PDF
V + jj	$d\sigma(\text{lept. V decay}) @ \text{NLO QCD}$	$d\sigma(\text{lept. V decay})$ @ NNLO QCD + NLO EW	study of systematics of H + jj final state
VV'	$d\sigma(\text{V decays}) @ \text{NLO QCD}$ $d\sigma(\text{stable V}) @ \text{NLO EW}$	$d\sigma(\text{V decays})$ @ NNLO QCD + NLO EW	off-shell leptonic decays TGCs
gg → VV	$d\sigma(\text{V decays}) @ \text{LO QCD}$	$d\sigma(\text{V decays})$ @ NLO QCD	bkg. to $H \rightarrow VV$ TGCs
V $\gamma$	$d\sigma(\text{V decay}) @ \text{NLO QCD}$ $d\sigma(\text{PA, V decay}) @ \text{NLO EW}$	$d\sigma(\text{V decay})$ @ NNLO QCD + NLO EW	TGCs
Vb $\bar{b}$	$d\sigma(\text{lept. V decay}) @ \text{NLO QCD}$ massive b	$d\sigma(\text{lept. V decay}) @ \text{NNLO QCD}$ massless b	bkg. for $VH \rightarrow b\bar{b}$
VV' $\gamma$	$d\sigma(\text{V decays}) @ \text{NLO QCD}$	$d\sigma(\text{V decays})$ @ NLO QCD + NLO EW	QGCs
VV'V''	$d\sigma(\text{V decays}) @ \text{NLO QCD}$	$d\sigma(\text{V decays})$ @ NLO QCD + NLO EW	QGCs, EWSB
VV' + j	$d\sigma(\text{V decays}) @ \text{NLO QCD}$	$d\sigma(\text{V decays})$ @ NLO QCD + NLO EW	bkg. to H, BSM searches
VV' + jj	$d\sigma(\text{V decays}) @ \text{NLO QCD}$	$d\sigma(\text{V decays})$ @ NLO QCD + NLO EW	QGCs, EWSB
$\gamma\gamma$	$d\sigma @ \text{NNLO QCD}$		bkg to $H \rightarrow \gamma\gamma$

Table 3: Wishlist part 3 – EW gauge bosons ( $V = W, Z$ )

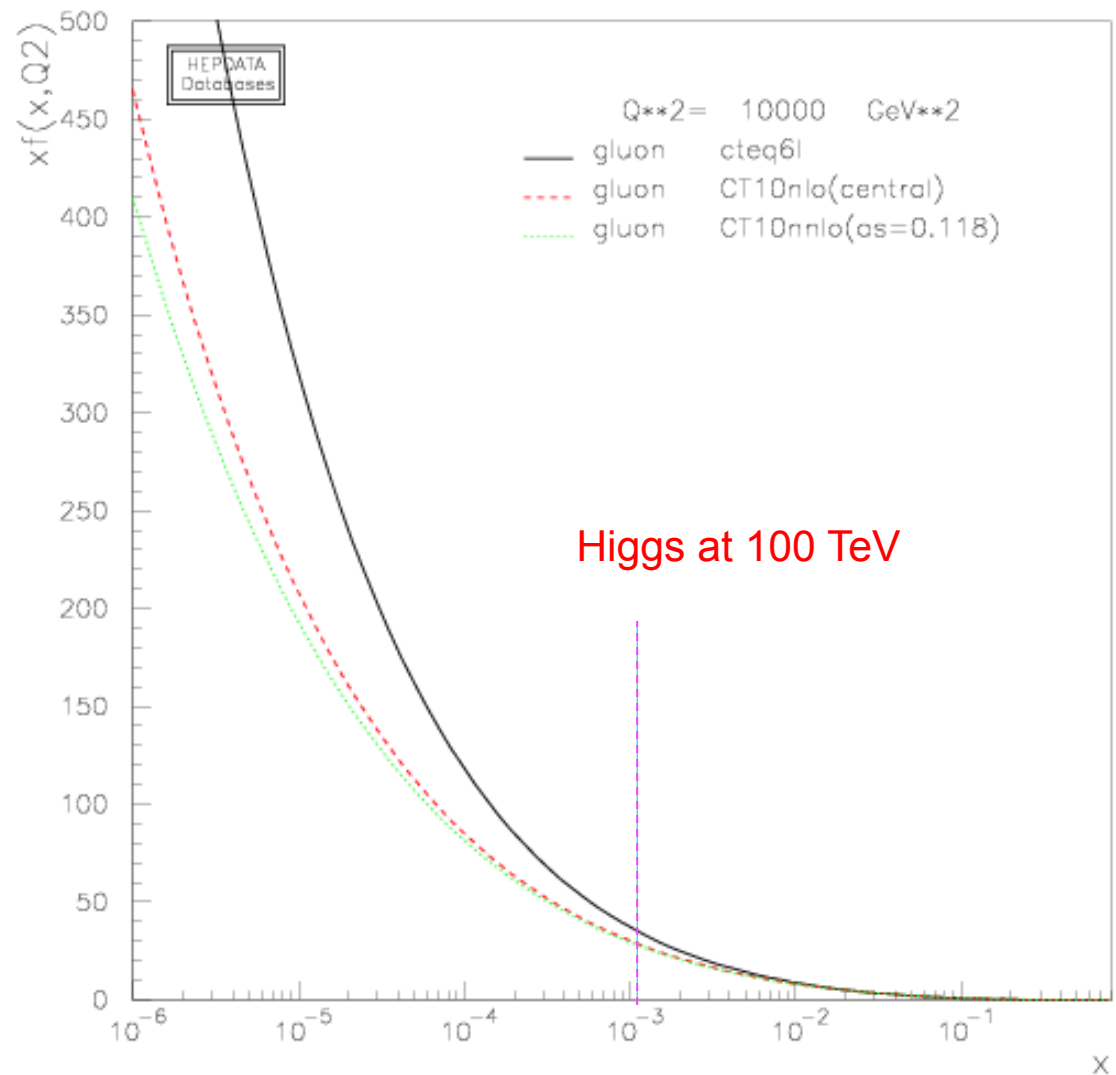


# The frontier



# Will we need N<sup>3</sup>LO PDFs for 100 TeV?

- There's a big change in the gluon distribution in going from LO to NLO
- Much smaller change from NLO to NNLO
- In Higgs kinematic region, scale uncertainties will dominate over PDF order effects
  - ◆ Forte, Isgro and Vita, arXiv: 1312.6688
- Maybe for precision physics at smaller  $x$ ?



# Summary

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- A 100 TeV pp collider will be exploring new kinematic regions in  $x$  and  $Q^2$ , where current knowledge is just extrapolation
- The 100 TeV data will be useful in determining PDFs in these new kinematic regions
- ...but there is also the opportunity to partially explore these kinematic regions in advance using LHC data and a possible LHeC
- Meanwhile, theorists have enough to keep themselves out of trouble with the new high precision wishlist, perhaps until the 100 TeV collider arrives

# Wu Ki Tung Award for Early Career Research on QCD

- See information at [http://tigger.uic.edu/~varelas/tung\\_award/](http://tigger.uic.edu/~varelas/tung_award/)
- Contribute at <https://www.givingto.msu.edu/gift/?sid=1480>
- **MSU will match any donations**



# Extras



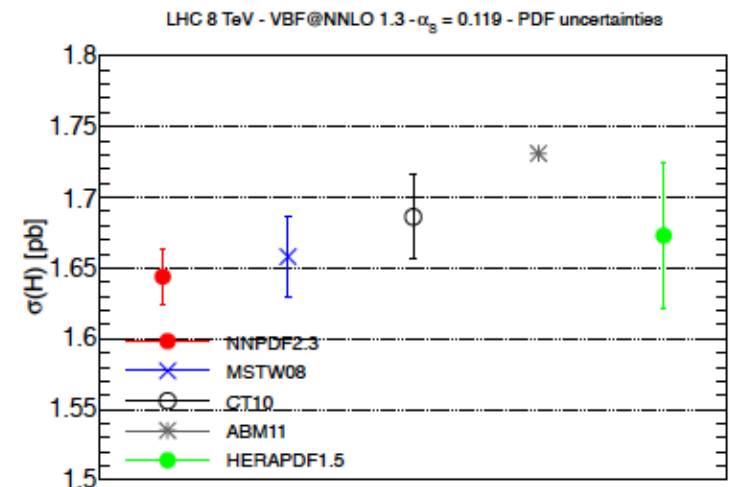
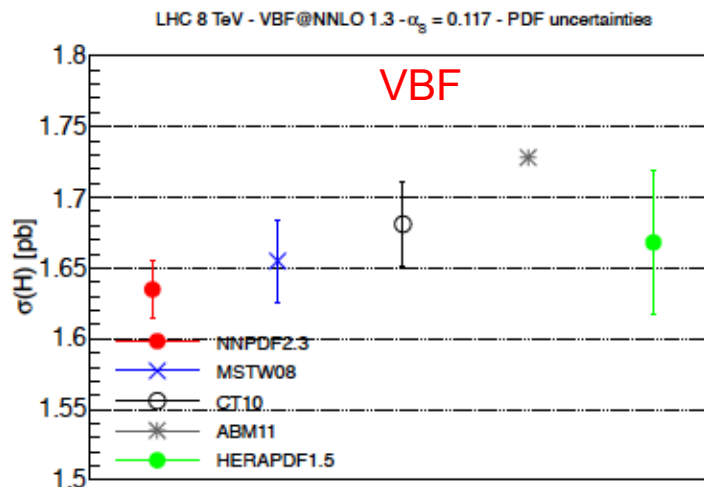
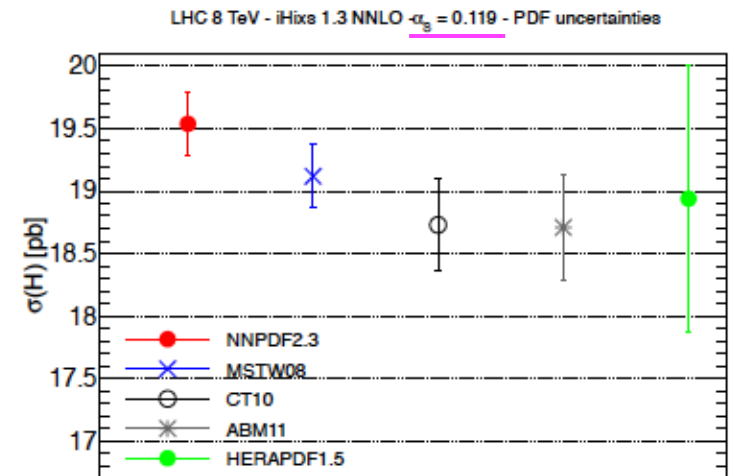
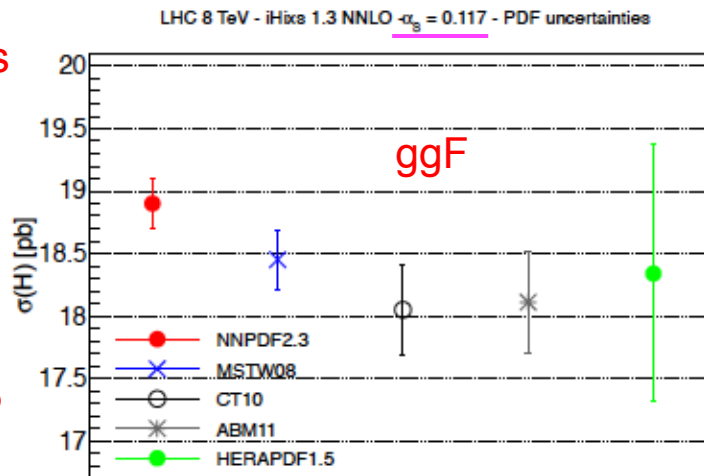


# 8 TeV Higgs cross section predictions

cross sections  
calculated at  
NNLO  
using a scale  
of  $m_H$

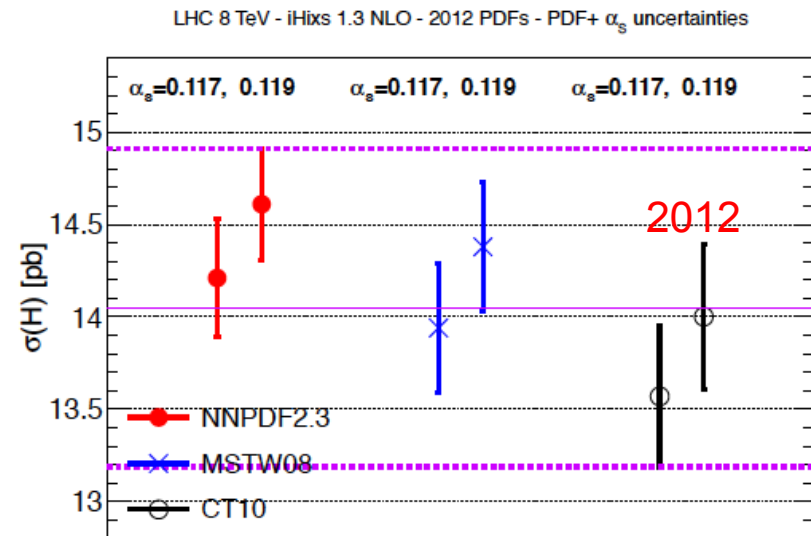
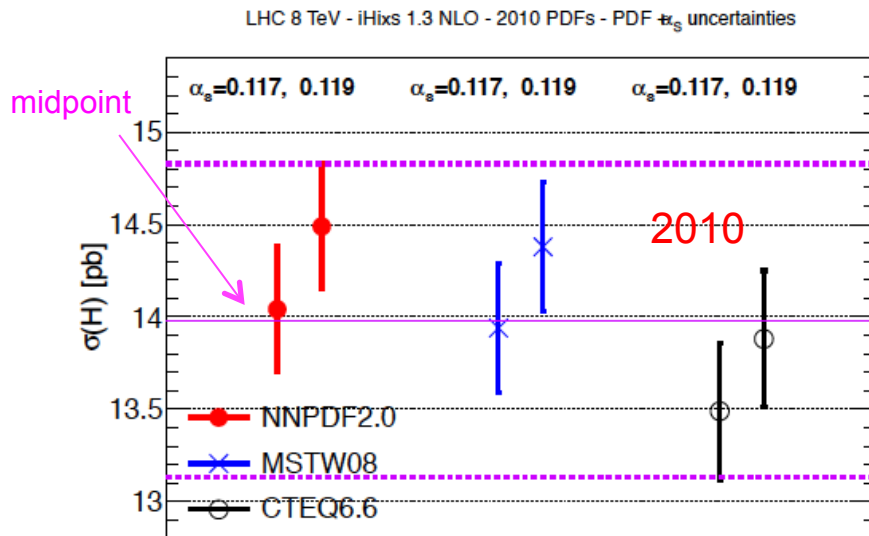
ABM11 and  
HERAPDF1.5  
predictions  
within  
error  
envelope

NB: ABM11  
cross section  
would be  
lower if  
native value  
of  $\alpha_s$  (0.1134)  
used



# Revisit prescriptions (for 8 TeV cross sections for gg fusion)

$$\sigma_H^{\text{NLO}} = 13.98 \pm 0.85 \text{ pb}, \quad (\pm 6.1\% \text{ "PDF} + \alpha_s") \rightarrow \sigma_H^{\text{NLO}} = 14.05 \pm 0.86 \text{ pb}, \quad (\pm 6.1\% \text{ "PDF} + \alpha_s").$$

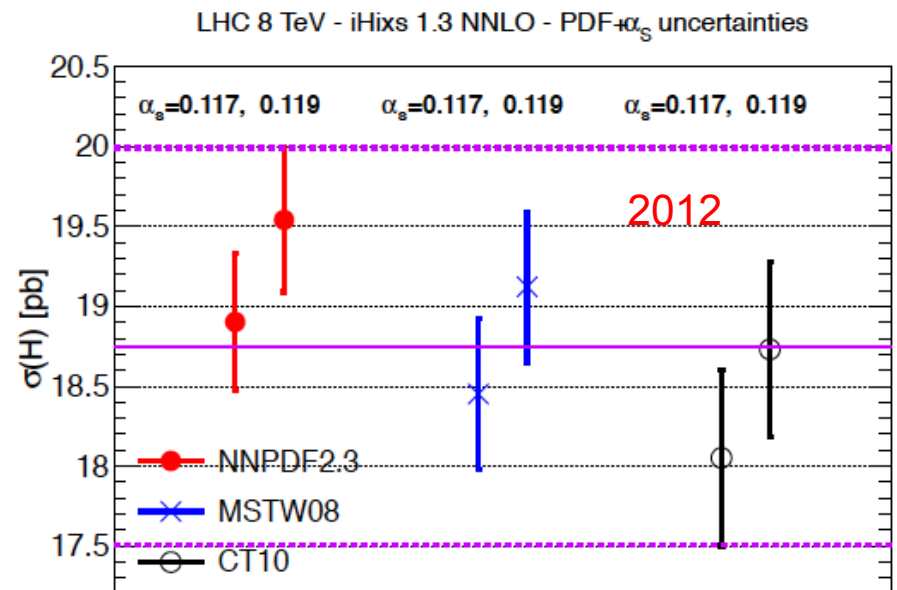


# Revisit prescriptions (for 8 TeV cross sections for gg fusion)

## 2012 NNLO result

$$\sigma_H^{NNLO} = 18.75 \pm 1.24 \text{ pb}, \quad (6.6\% \text{ "PDF} + \alpha_s\text{"}).$$

Compare to MSTW08 NNLO value of  
18.45 pb  
(2010 prescription)

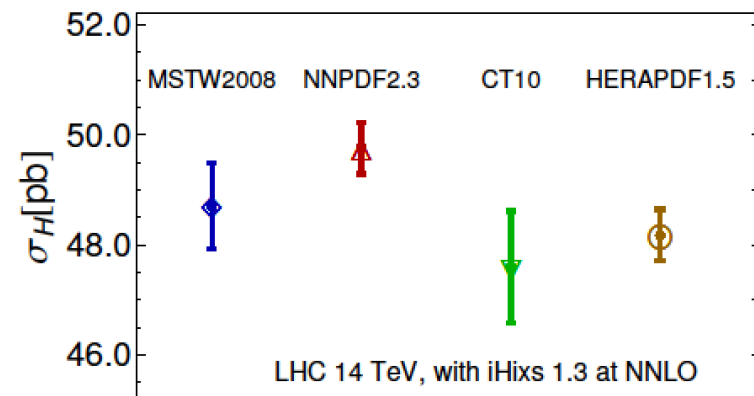
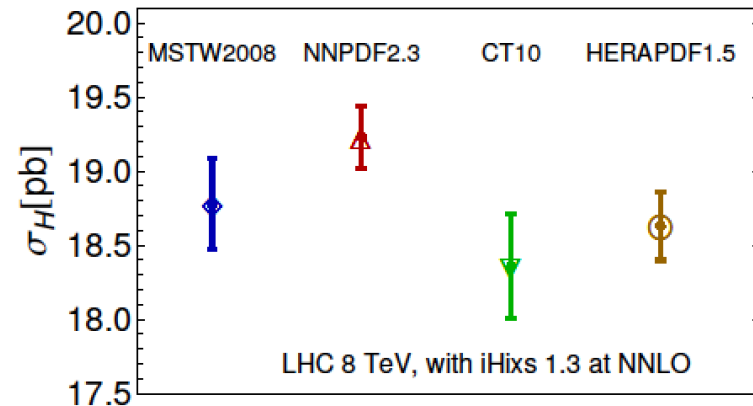


HXSWG 8 TeV NNLO cross section  
NNLO+NNLL

$$\sigma_H^{NNLO} = 19.52 \pm 1.41 \text{ pb}, \quad (\pm 7.2\% \text{ "PDF} + \alpha_s\text{"}).$$

# Les Houches study

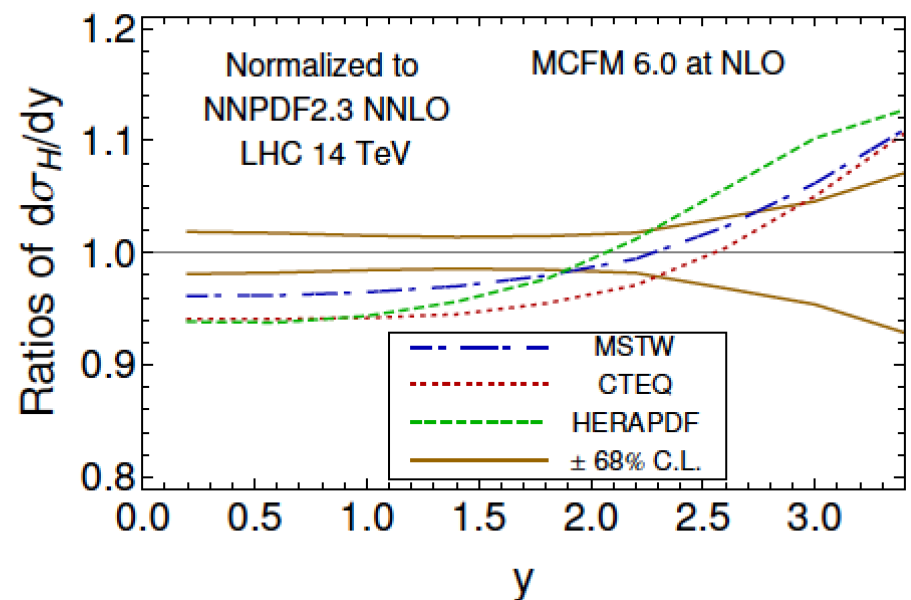
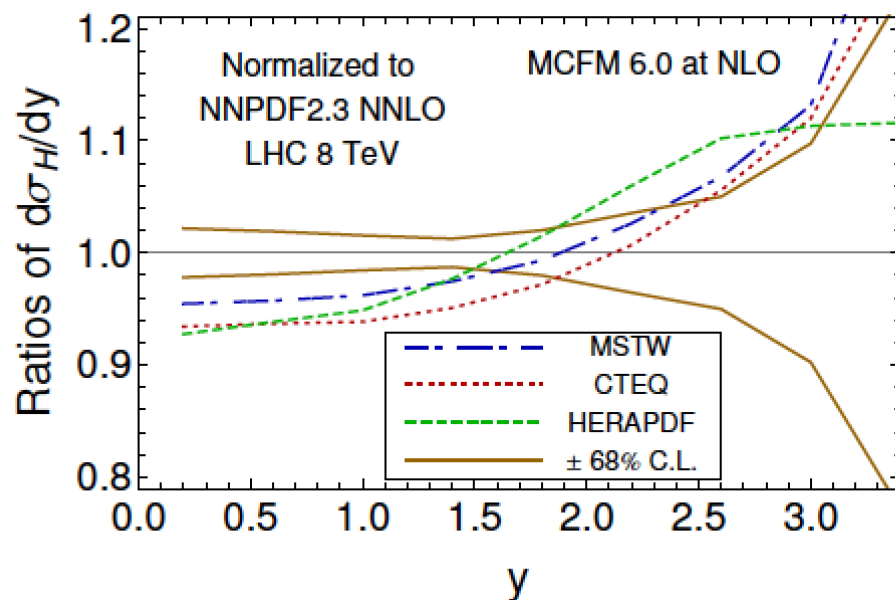
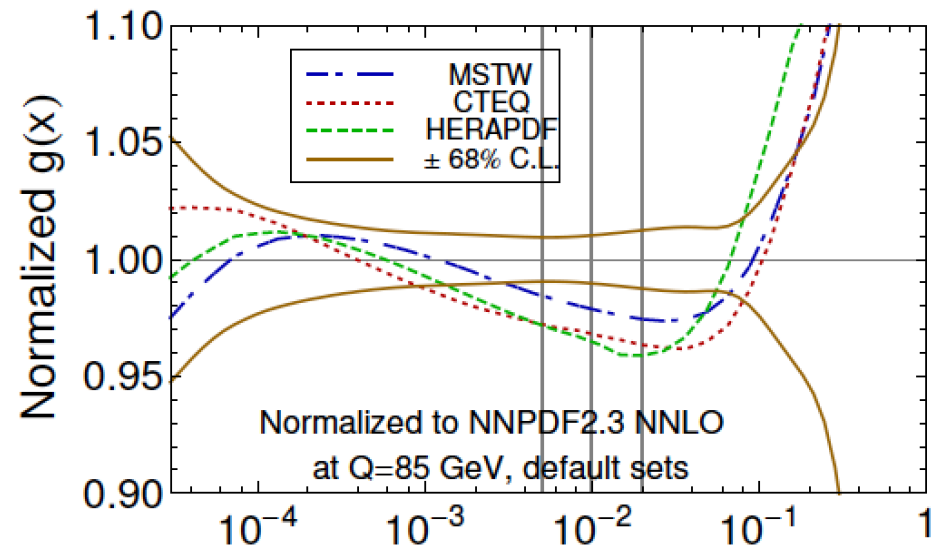
- The study used the PDFs that use a variable flavor number scheme, i.e. the 4 shown on the curves to the right
- The discrepancies present at 8 TeV persist at 14 TeV with the same pattern



$\sigma_H$ [pb]	CT10	MSTW'08	NNPDF2.3	HERAPDF 1.5
LHC 8 TeV	$18.36 \pm 0.35$	$18.78 \pm 0.31$	$19.23 \pm 0.21$	$18.63 \pm 0.23$
LHC 14 TeV	$47.60 \pm 1.02$	$48.71 \pm 0.77$	$49.76 \pm 0.47$	$48.18 \pm 0.47$

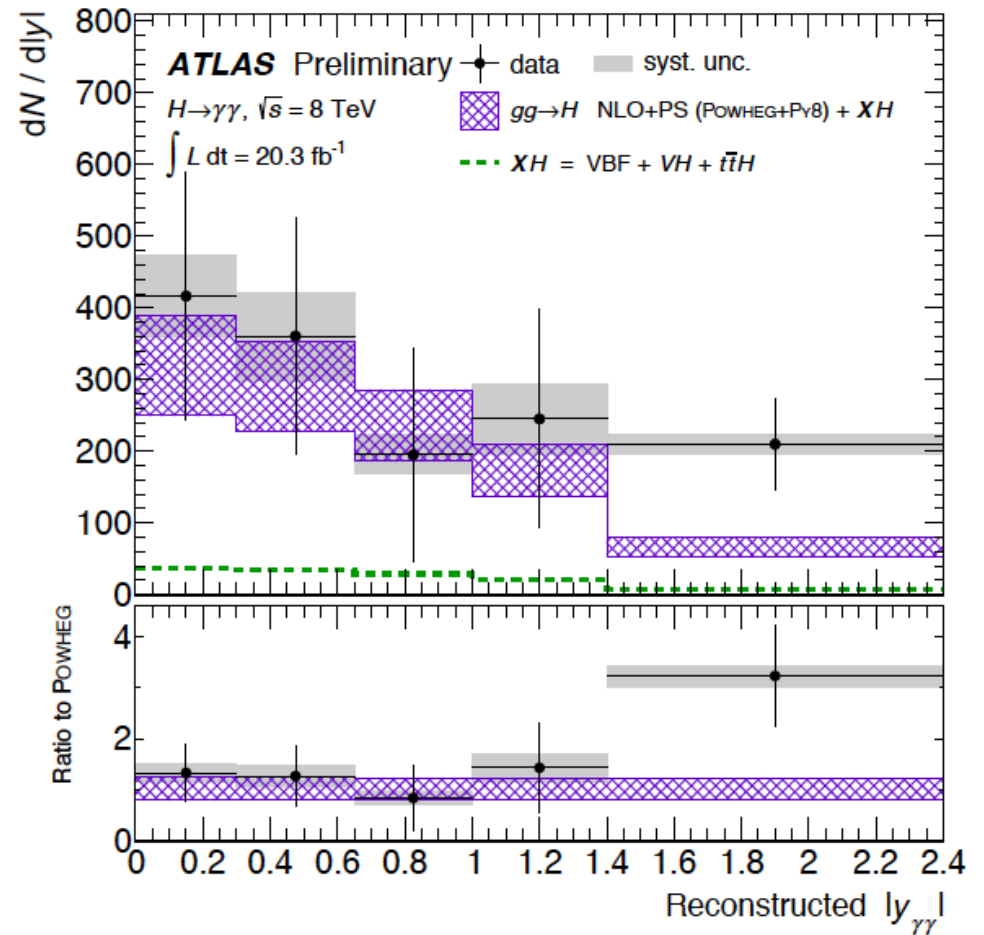
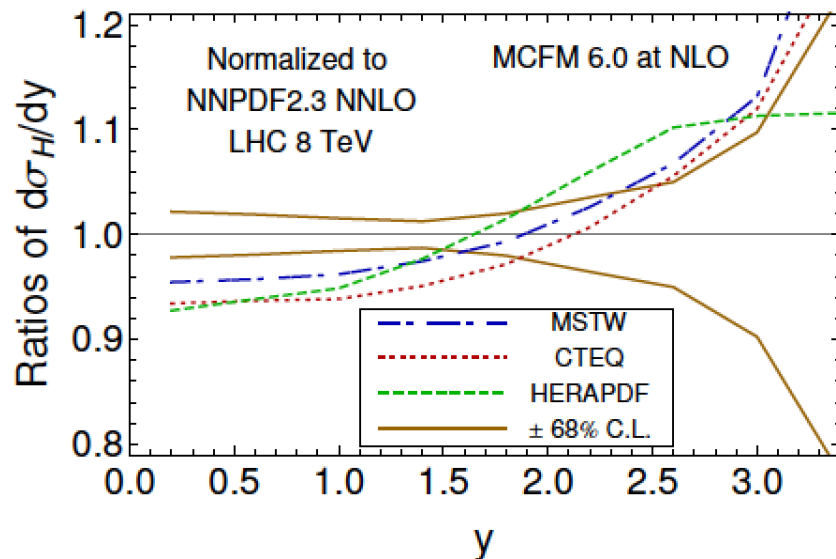
# Gluon distributions and Higgs $y$

- On the right some obvious shape differences as a function of  $x$  can be seen
- These shape differences translate into different predictions for Higgs rapidity as seen below



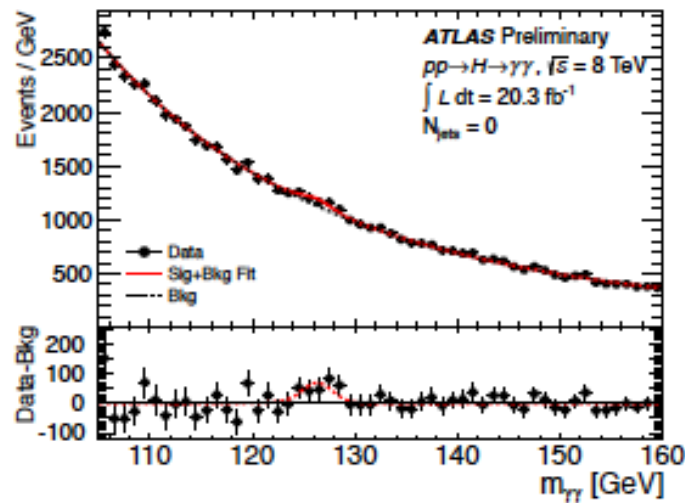
# Gluon distributions and Higgs $y$

- These shape differences translate into different predictions for Higgs rapidity
- Alas, we're not quite there for using the Higgs rapidity distribution to tune PDFs

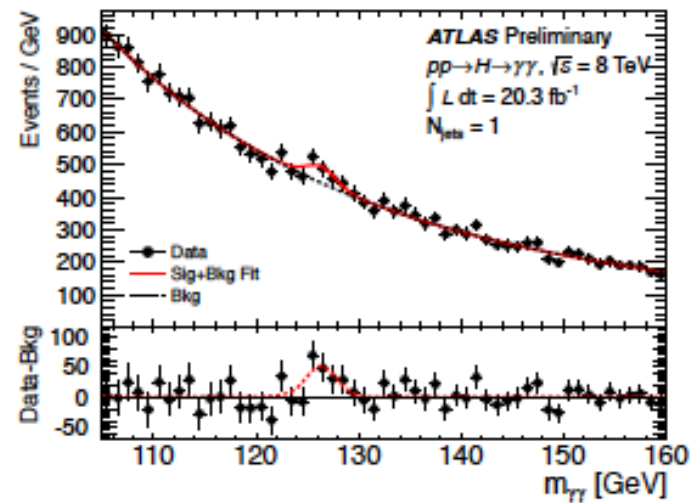




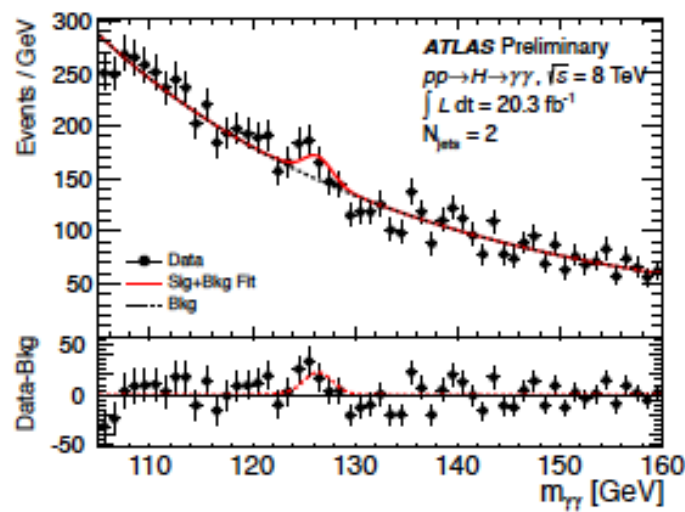
# Aside



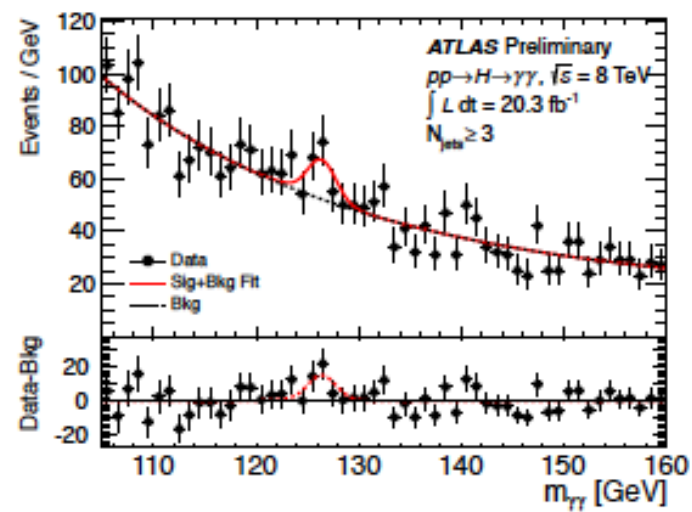
(a)  $N_{\text{jets}} = 0$



(b)  $N_{\text{jets}} = 1$



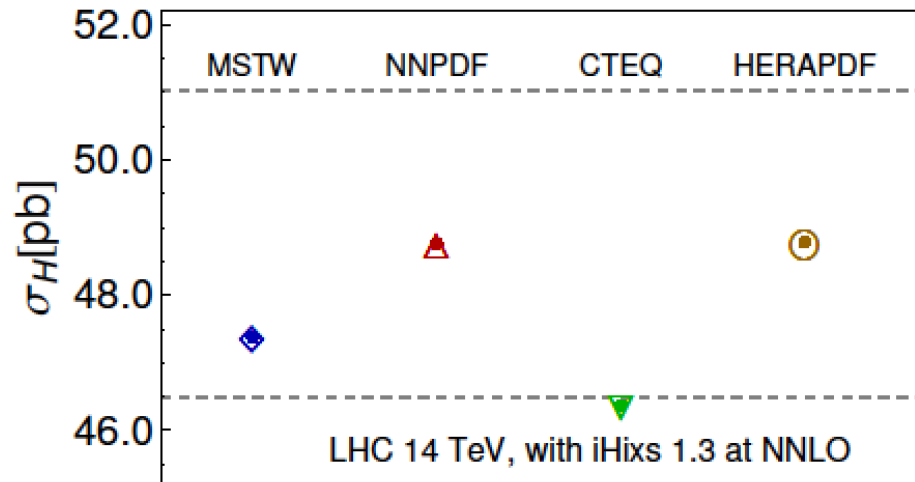
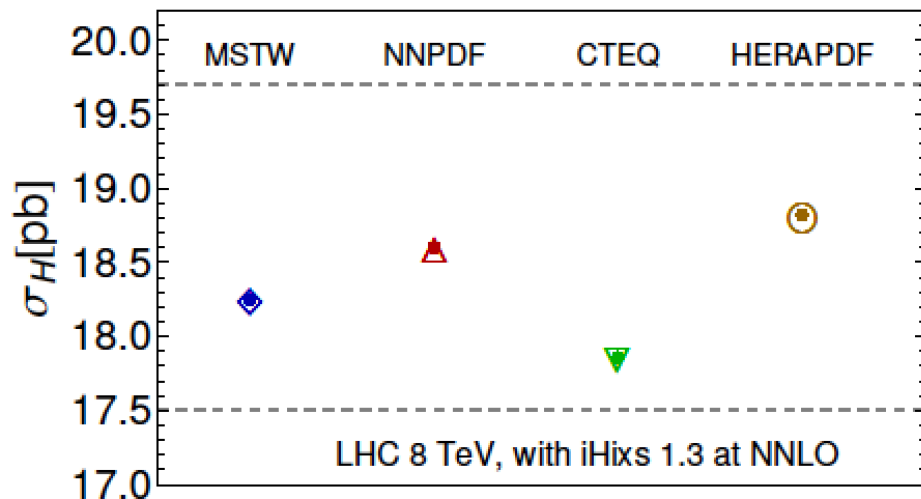
(c)  $N_{\text{jets}} = 2$



(d)  $N_{\text{jets}} \geq 3$

# Results of fits to HERA1 only

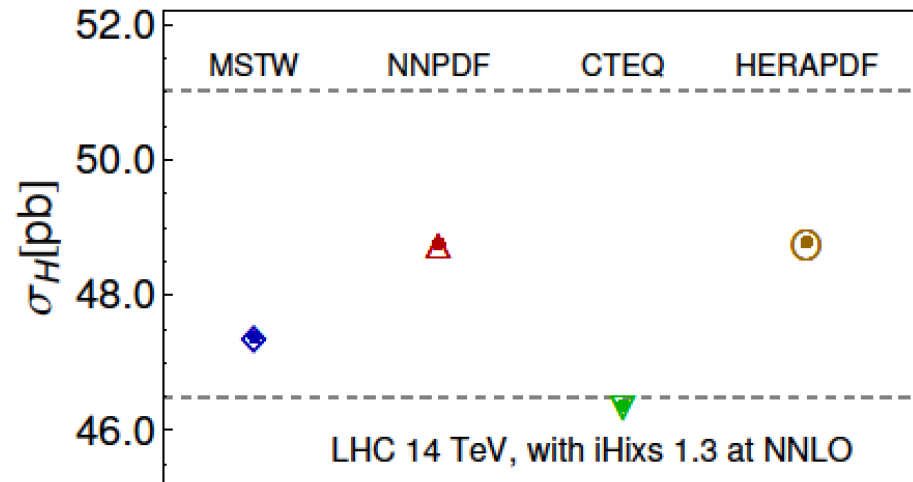
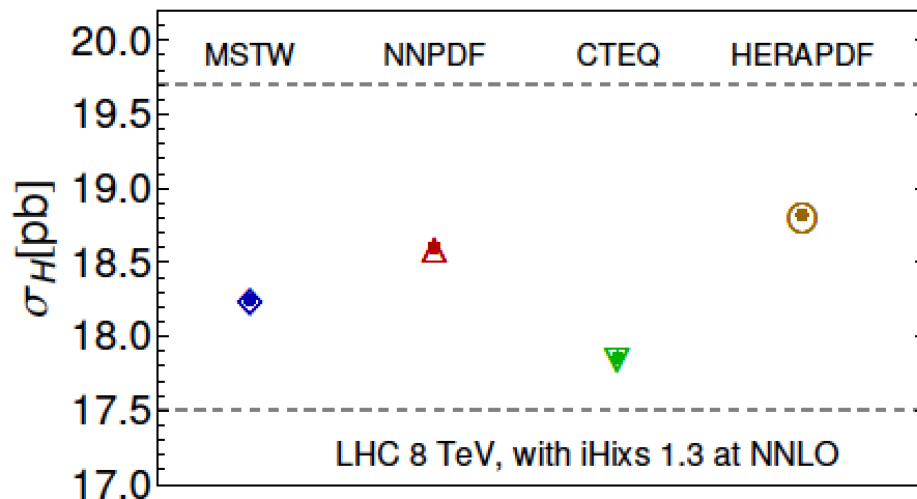
- Square 1: benchmark comparison of NNLO neutral-current DIS structure functions for four fitting codes with same toy PDF->good agreement
- Then, fit to one well-defined (and important) dataset: HERA Run 1
- Result: all predictions for Higgs cross sections (except HERA1) decrease
- All predictions within (expanded) PDF uncertainty of NNPDF2.3
- ...but hierarchy remains the same
- ...why?



# Results of fits to HERA1 only

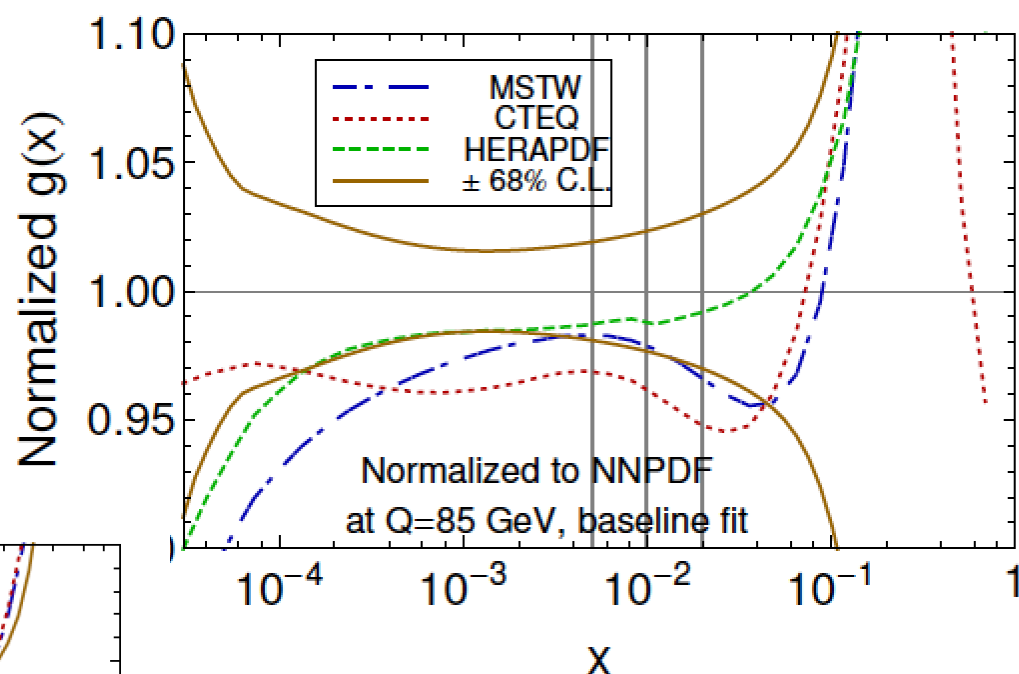
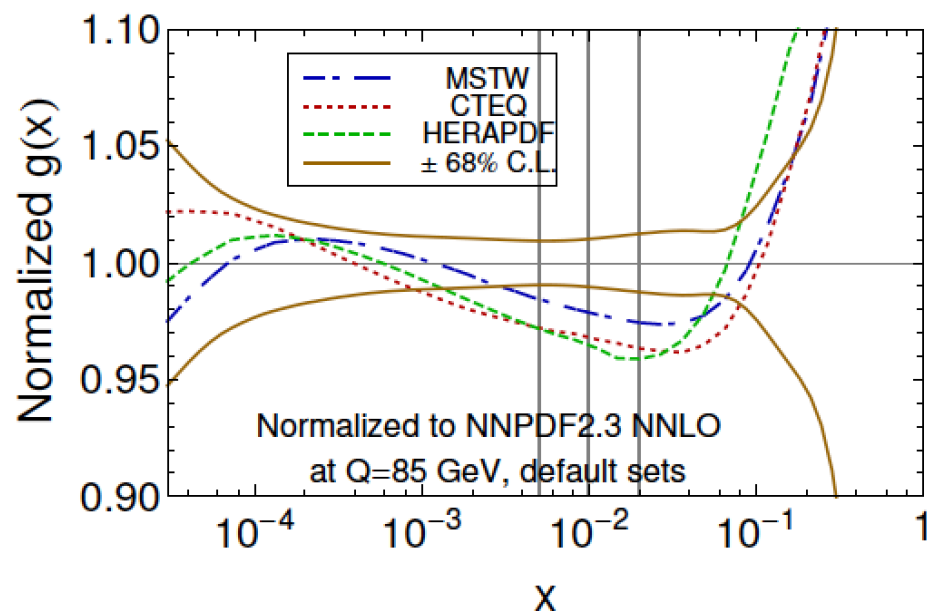
- All predictions (except HERA1) decrease

566 data points	CTEQ	MSTW	NNPDF	HERAPDF
$\chi^2$	521.8	514.8	548.5	535.0
lum. shift	-0.19	0.27	0.16	0.18
max. shift	1.64	1.51	1.82	1.81
$\sigma_H$ [pb], 8 TeV	17.86	18.25	$18.60 \pm 1.10$	18.82
$\sigma_H$ [pb], 14 TeV	46.37	47.38	$48.76 \pm 2.26$	48.78



# HERA1 benchmark PDFs

- More similarity in shape at low  $x$ , but still big differences at high  $x$  (but within expanded uncertainties)



# Gluon-gluon luminosities

- ...as a function of mass
- Again, as noted before, all predictions are within the expanded uncertainties

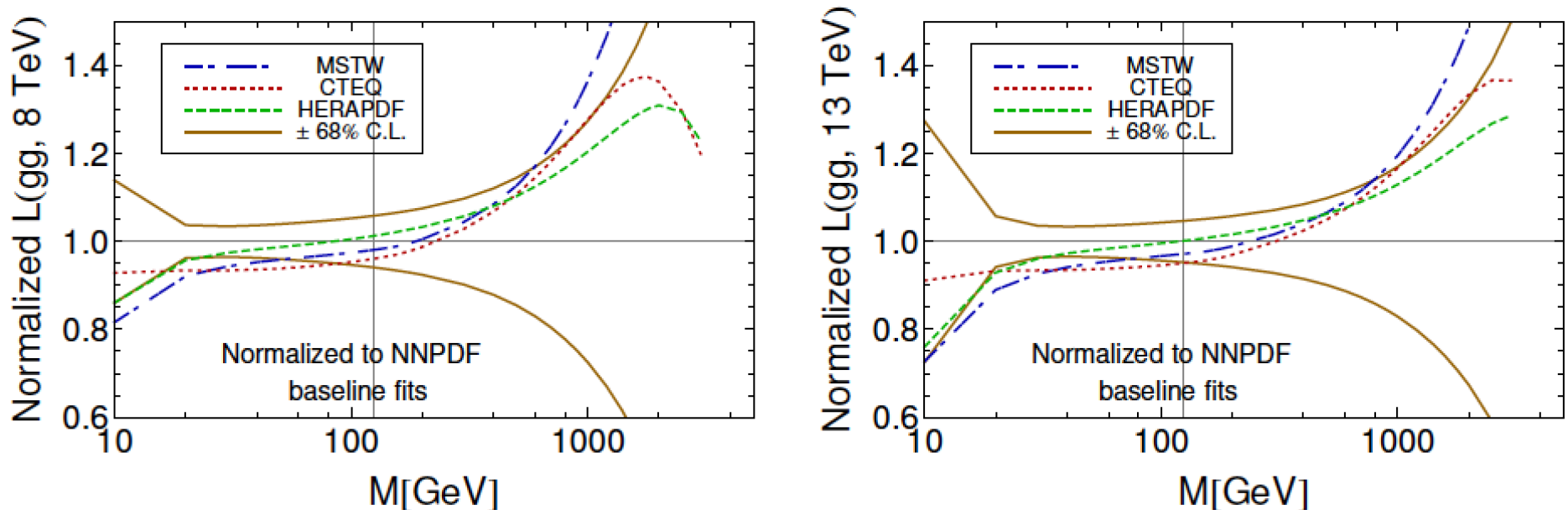


Fig. IV.14: Comparison of the gluon-gluon parton luminosity as a function of invariant mass at the LHC 8 and 13 TeV from the HERA-1-only NNLO fits, normalized to the NNPDF central prediction. The

# Summary (of study)

---

- Lots of other detail in the Les Houches writeup
  - ◆ exact definitions of  $\chi^2$
  - ◆ checks of parametrisation, scale choices, heavy quark schemes, ...
- HERA-1 only fits prefer smaller Higgs cross sections
- Predictions using HERA-1 follow same pattern as with full global data sets
- Next step (post-Les Houches): add additional data sets into comparisons sequentially, ensuring all groups use exactly the same data points, uncertainties, definition of the systematic uncertainties, etc
- Compare the impact of LHC data sensitive to the gluon PDF



- ...but consider the 2012 inclusive jet measurement from CMS (8 TeV) where CT10 seems to provide a good description
- ...with much higher statistics and improved systematics
- Errors aren't public yet so don't know the impact on global PDF fits

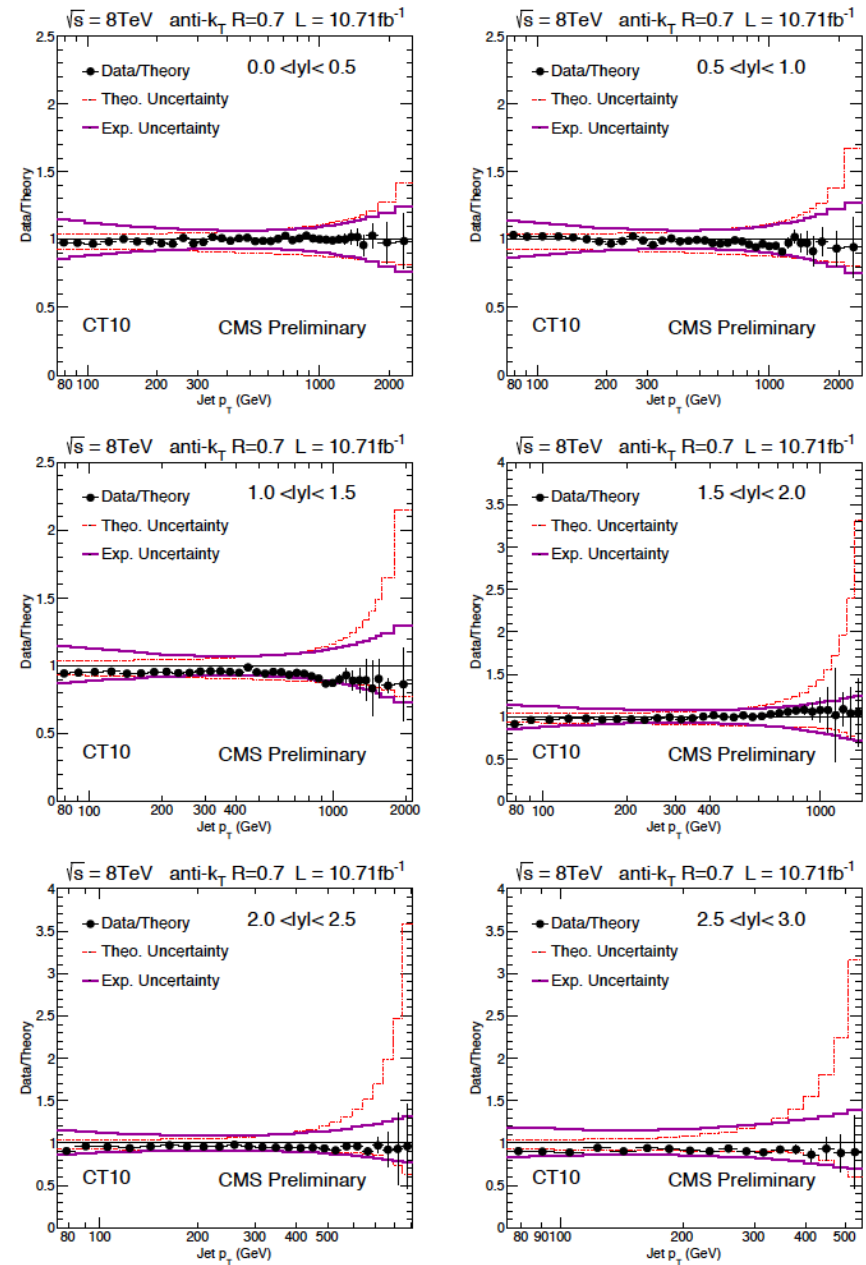


Figure 3: Ratio of data over theory at NLO times NP correction for the CT10 PDF set. For comparison the total theoretical (band enclosed by dashed red lines) and the total experimental systematic uncertainty (band enclosed by full magenta lines) are shown as well. The error bars correspond to the statistical uncertainty of the data.

- ...whereas NNPDF2.3 (or MSTW08) seems to be below the data at high  $p_T$

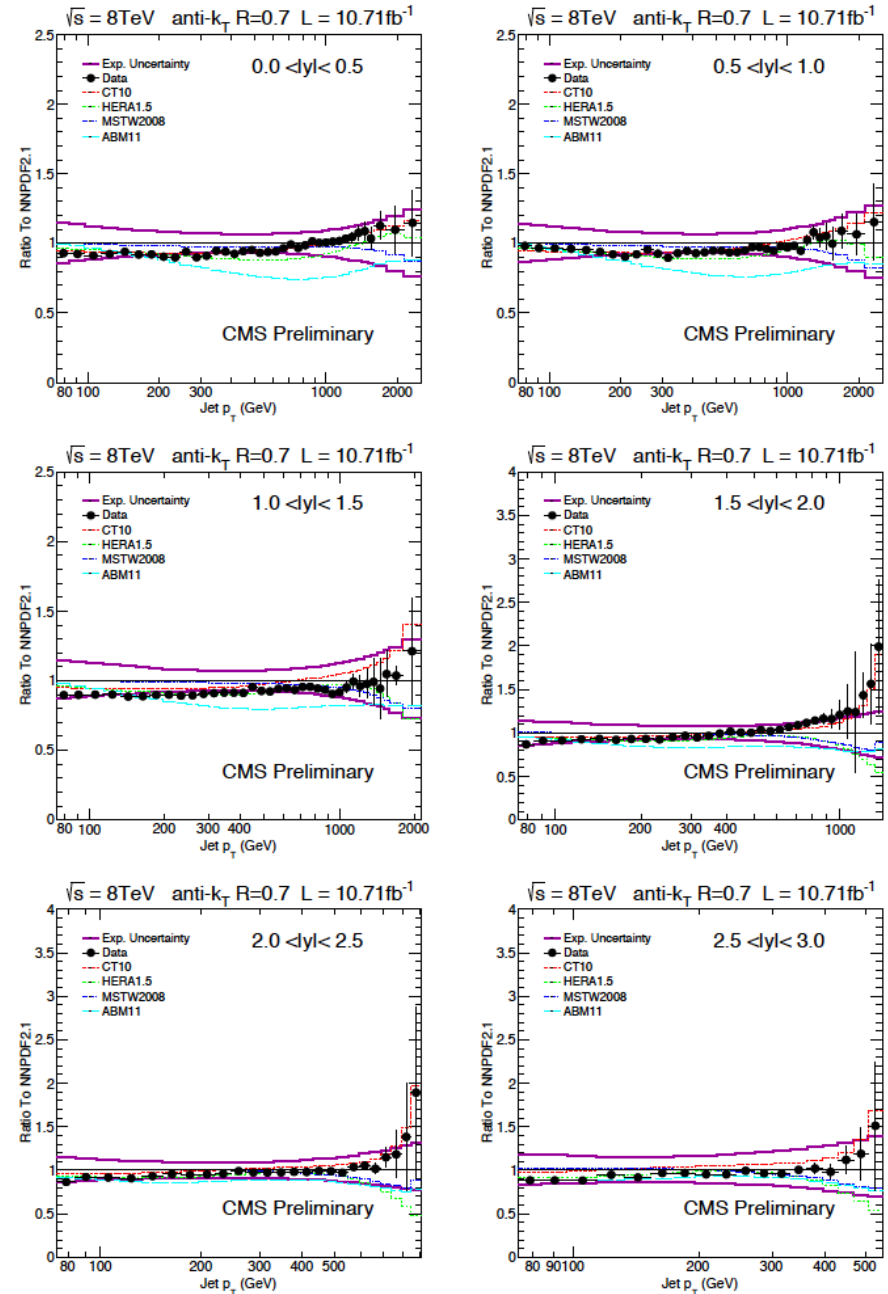


Figure 7: Ratio of data over theory at NLO times NP correction for the NNPDF2.1 PDF set. For comparison predictions employing four other PDF sets are shown in addition to the total experimental systematic uncertainty (band enclosed by full magenta lines). The error bars correspond to the statistical uncertainty of the data.

# Fits of the fits: META PDFs

PDFs from different groups have different physics inputs. But if we only focus on the phenomenological studies at the LHC with the limited  $x$  and  $Q$  ranges, the idea of META PDFs is reasonable and also feasible.

Procedure (for LHC):

- 1, selecting a specific  $x$ - $Q$  range, and a parameterization form to describe all the PDFs at an initial scale above the bottom quark mass;
- 2, check that the fitted PDFs can well represent the original PDFs at the  $x$ - $Q$  range studied;
- 3, choosing a scheme to combine the PDF measurements of different groups in the new PDF parameter space;

Benefits:

- 1, A nature way to compare and combine the LHC predictions from different PDF groups independent of the process, works similarly as the PDF4LHC prescriptions but directly in the PDF parameter space;
- 2, Especially desirable for including results from large number of PDF groups, in this case also minimizing numerical computation efforts for massive NNLO calculations

**It's possible to define a few eigenvectors which completely encompass the PDF and  $\alpha_s$  uncertainties for CT10, MSTW08 and NNPDF2.3 for Higgs production for 8-14 TeV; no reason this cannot be expanded to 100 TeV**

Jun Gao, Pavel Nadolsky, JH

# PDF+ $\alpha_s$ uncertainties

VBF exc., 14 TeV, LO	-0.43	-0.49	-0.3	0.09	0.09	0.06	0.92	0.92	-0.39	-0.42	-0.33	0.02	0.02	0.	1.
	<b>-0.44</b>	<b>-0.5</b>	<b>-0.33</b>	<b>0.09</b>	<b>0.09</b>	<b>0.09</b>	<b>0.93</b>	<b>0.93</b>	<b>-0.4</b>	<b>-0.44</b>	<b>-0.35</b>	<b>0.02</b>	<b>0.02</b>	<b>0.</b>	<b>1.</b>
VBF inc., 14 TeV, LO	-0.43	-0.49	-0.3	0.09	0.09	0.06	0.92	0.92	-0.39	-0.42	-0.33	0.02	0.02	0.	
	<b>-0.44</b>	<b>-0.5</b>	<b>-0.33</b>	<b>0.09</b>	<b>0.09</b>	<b>0.09</b>	<b>0.93</b>	<b>0.93</b>	<b>-0.4</b>	<b>-0.44</b>	<b>-0.35</b>	<b>0.02</b>	<b>0.02</b>	<b>0.</b>	
GGH 2j full mass, 14 TeV, LO	0.45	0.25	0.72	0.96	0.96	0.96	-0.04	-0.04	0.31	0.08	0.47	0.99	0.99		
	<b>0.42</b>	<b>0.22</b>	<b>0.71</b>	<b>0.98</b>	<b>0.98</b>	<b>0.98</b>	<b>-0.05</b>	<b>-0.05</b>	<b>0.28</b>	<b>0.05</b>	<b>0.46</b>	<b>0.99</b>	<b>0.99</b>		
GGH 2j exc., 14 TeV, LO	0.43	0.22	0.71	0.97	0.97	0.97	-0.01	-0.01	0.29	0.07	0.46	0.99			
	<b>0.44</b>	<b>0.23</b>	<b>0.72</b>	<b>0.98</b>	<b>0.98</b>	<b>0.98</b>	<b>-0.02</b>	<b>-0.02</b>	<b>0.29</b>	<b>0.07</b>	<b>0.48</b>	<b>0.99</b>			
GGH 2j inc., 14 TeV, LO	0.43	0.22	0.71	0.97	0.97	0.97	-0.01	-0.01	0.29	0.07	0.46				
	<b>0.44</b>	<b>0.23</b>	<b>0.72</b>	<b>0.98</b>	<b>0.98</b>	<b>0.98</b>	<b>-0.02</b>	<b>-0.02</b>	<b>0.29</b>	<b>0.07</b>	<b>0.48</b>				
GGH 1j exc., 14 TeV, NLO	0.98	0.94	0.93	0.3	0.3	0.3	-0.34	-0.34	0.97	0.89					
	<b>0.98</b>	<b>0.94</b>	<b>0.94</b>	<b>0.33</b>	<b>0.33</b>	<b>0.33</b>	<b>-0.34</b>	<b>-0.34</b>	<b>0.97</b>	<b>0.9</b>					
GGH 0j exc., 14 TeV, NLO	0.91	0.96	0.7	-0.07	-0.07	-0.07	-0.4	-0.4	0.97						
	<b>0.92</b>	<b>0.97</b>	<b>0.73</b>	<b>-0.08</b>	<b>-0.08</b>	<b>-0.08</b>	<b>-0.4</b>	<b>-0.4</b>	<b>0.97</b>						
GGH inc., 14 TeV, NNLO	0.97	0.97	0.84	0.14	0.14	0.14	-0.38	-0.38							
	<b>0.97</b>	<b>0.98</b>	<b>0.87</b>	<b>0.14</b>	<b>0.14</b>	<b>0.14</b>	<b>-0.39</b>	<b>-0.39</b>							
VBF exc., 8 TeV, LO	-0.41	-0.44	-0.31	0.06	0.06	0.04	1.								
	<b>-0.41</b>	<b>-0.45</b>	<b>-0.33</b>	<b>0.05</b>	<b>0.05</b>	<b>0.05</b>	<b>0.99</b>								
VBF inc., 8 TeV, LO	-0.41	-0.44	-0.31	0.06	0.06	0.04									
	<b>-0.41</b>	<b>-0.45</b>	<b>-0.33</b>	<b>0.05</b>	<b>0.05</b>	<b>0.05</b>									
GGH 2j full mass, 8 TeV, LO	0.27	0.06	0.57	0.99	0.99										
	<b>0.29</b>	<b>0.08</b>	<b>0.6</b>	<b>0.99</b>	<b>0.99</b>										
GGH 2j exc., 8 TeV, LO	0.27	0.06	0.57	0.99											
	<b>0.29</b>	<b>0.08</b>	<b>0.6</b>	<b>0.99</b>											
GGH 2j inc., 8 TeV, LO	0.27	0.06	0.57												
	<b>0.29</b>	<b>0.08</b>	<b>0.6</b>												
GGH 1j exc., 8 TeV, NLO	0.93	0.83													
	<b>0.93</b>	<b>0.83</b>													
GGH 0j exc., 8 TeV, NLO	0.97														
	<b>0.97</b>														
GGH inc., 8 TeV, NNLO	3.3%	3.2%	3.6%	6.9%	6.9%	7.7%	2.4%	2.4%	3.3%	3.2%	3.4%	5.7%	5.7%	5.8%	2.1%
	<b>3.3%</b>	<b>3.2%</b>	<b>3.5%</b>	<b>6.8%</b>	<b>6.8%</b>	<b>6.8%</b>	<b>2.4%</b>	<b>2.4%</b>	<b>3.3%</b>	<b>3.2%</b>	<b>3.4%</b>	<b>5.7%</b>	<b>5.7%</b>	<b>5.8%</b>	<b>2.1%</b>

Correlation table for Higgs cross sections

Red indicates  $|\cos(\phi)| > 0.7$

Numbers in Italic–bold (plain) for 6 eigenvectors (full set 50 eig.)

META PDF 1.0 (LHCH)

VBF-like cut applied for 2 or more jets final states

jet (anti- $k_T$ , 0.4) selection with  $|y| < 4.5$  and  $p_T > 30$  GeV

including  $\alpha_s$  uncertainty

PDF+ $\alpha_s$  uncertainties: total unc. in last row; correlations in grids

# Nota bene

- For the PDFs to be fully NNLO, we need to use NNLO matrix elements for inclusive jet production, crucial to the determination of the high  $x$  gluon
- So far, we have them for the  $gg$  channel

- ◆ corrections are sizeable; I would expect them to be smaller for the  $gq$  and  $qQ$  channels, following the Dixon conjecture

Casimir for biggest color representation final state can be in

Simplistic rule

$$C_{i1} + C_{i2} - C_{f,\max}$$

L. Dixon

Casimir color factors for initial state

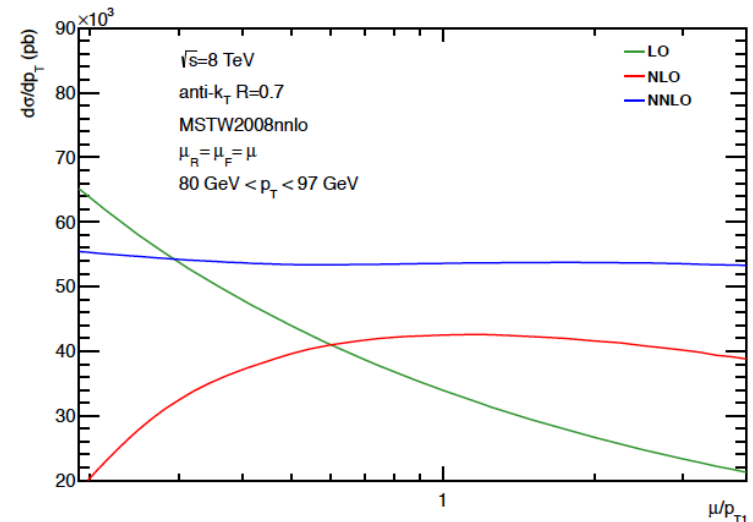
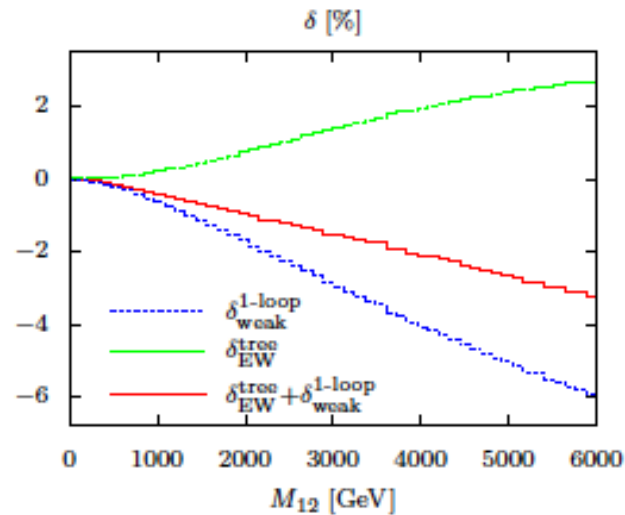
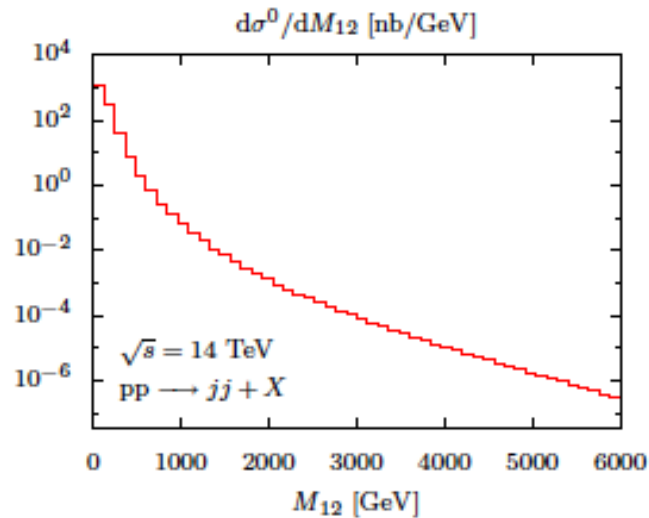


FIG. 2: Scale dependence of the inclusive jet cross section for  $pp$  collisions at  $\sqrt{s} = 8$  TeV for the anti- $k_T$  algorithm with  $R = 0.7$  and with  $|y| < 4.4$  and  $80 \text{ GeV} < p_T < 97 \text{ GeV}$  at NNLO (blue), NLO (red) and LO (green).

We know that NLO describes jet sections for  $R=0.6$  and  $R=0.7$  better than for  $R=0.4$  and  $R=0.5$ ; need extra gluon that's in NNLO?

Completion of NNLO this year?

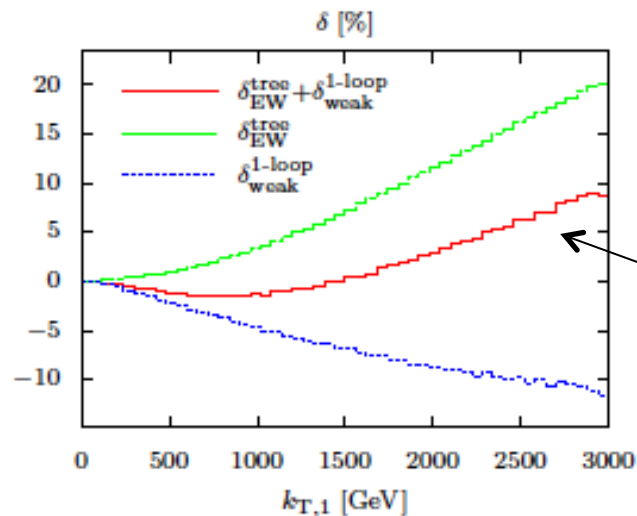
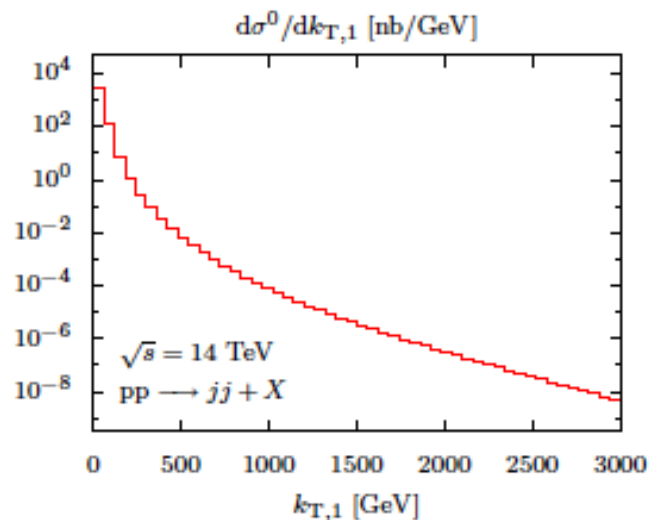
## Weak corrections to dijet production – numerical results



S.D., Huss, Speckner '12

### Weak corrections

- small for integrated XS
- growing in distributions for larger scales



Cancellations between tree and loop corrections (cut-sensitive!)

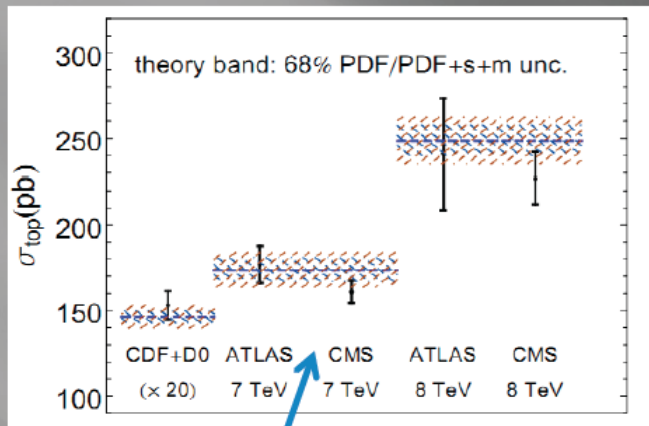
total EW impact starts to become noticeable at high  $p_T$



# Meta-PDFs

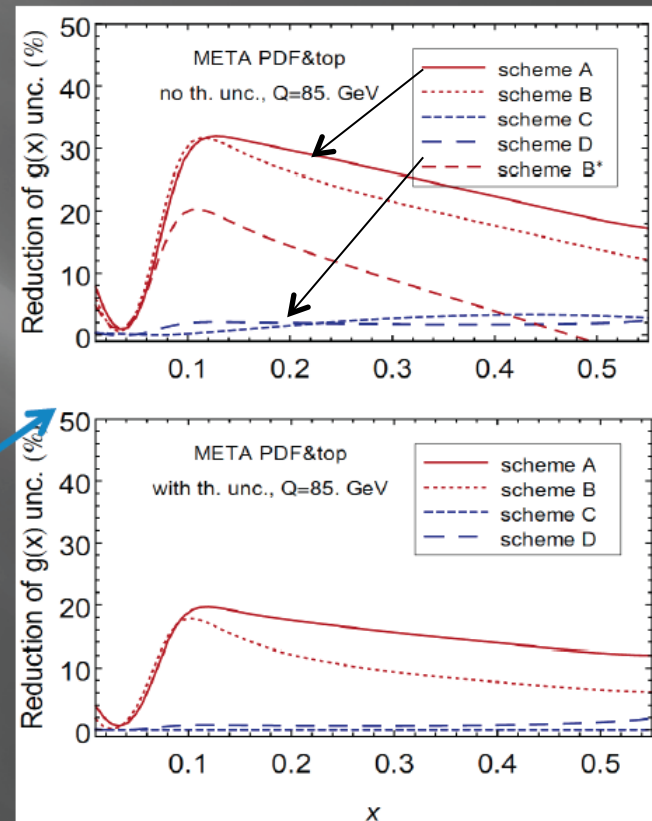
## Examples: top quark data

We perform a similar study as in (1303.7215, M. Czakon, et al.) using the measurements of top quark pair inclusive rate to constraint the gluon PDFs.



Comparison of META predictions with data before reweighting

Reduction of the gluon PDF uncertainties under different schemes with and without including theoretical uncertainties.



PRELIMINARY

effect of tolerance on impact of new data in global fits needs to be better understood

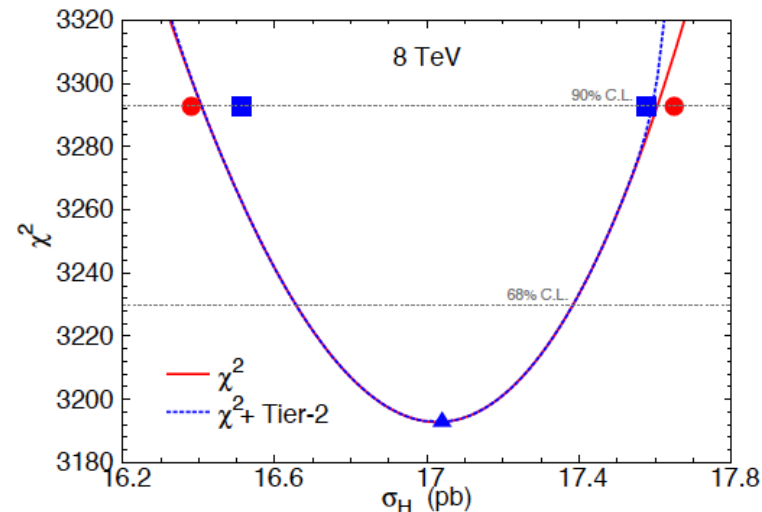
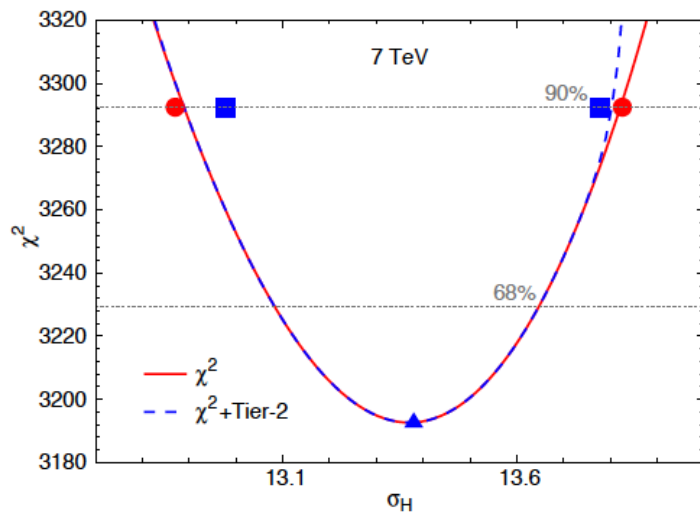
CTEQ/MSTW may be different than NNPDF?

investigate for Les Houches Writeup

use-cases for META-PDFS or equivalent

# Scaling issues: 90%CL->68%CL

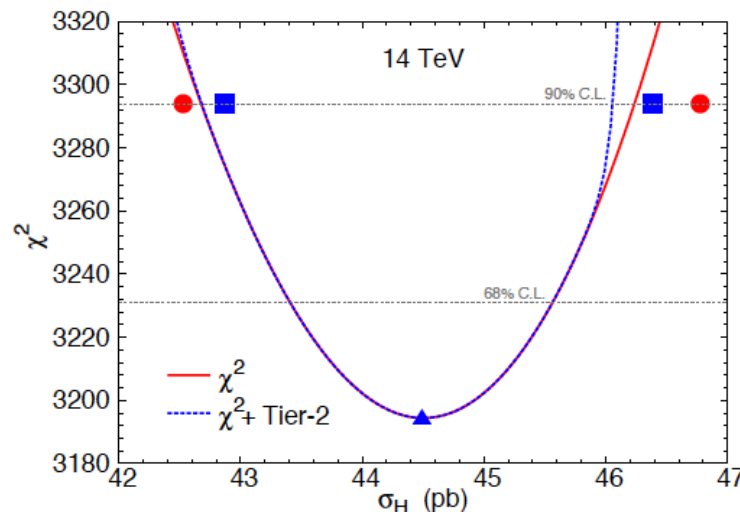
- New CT paper dealing with PDF and  $\alpha_s$  uncertainties for  $gg \rightarrow \text{Higgs}$  production, comparing Hessian and Lagrange Multiplier Techniques



LM technique not dependent on assumption of quadratic  $\chi^2$  behavior, so more robust than Hessian

Tier 2 penalty prevents the fit to any one experiment from degrading too Much

all predictions at NNLO using  $\mu = m_H$



curves are LM calculations of global fit  $\chi^2$  vs Higgs  $\sigma$  with (blue) and without (red) 'Tier 2 penalty'

The blue (red) points are the Hessian determination of the of the PDF uncertainty with (without) the Tier 2 penalty

# PDF+ $\alpha_s$ uncertainties

- LM estimates of PDF(+ $\alpha_s$ ) uncertainties slightly larger than Hessian determinations, but close, especially for the combined PDF+ $\alpha_s$  errors

Method	90% CL			68% CL		
	7 TeV	8 TeV	14 TeV	7 TeV	8 TeV	14 TeV
LM (PDF-only)	+3.2/-3.7	+3.2/-3.7	+3.5/-4.1	+2.0/-2.2	+2.0/-2.3	+2.2/-2.4
Hessian (PDF-only)	+3.0/-3.0	+3.2/-3.1	+4.3/-3.6	+1.8/-1.8	+1.9/-1.9	+2.6/-2.2
LM (PDF + $\alpha_s$ )	+4.8/-5.0	+4.6/-4.6	+5.2/-5.2	+2.9/-3.2	+2.8/-2.9	+3.4/-3.2
Hessian (PDF + $\alpha_s$ )	+4.7/-4.6	+4.8/-4.7	+5.4/-5.0	+2.9/-2.8	+2.9/-2.8	+3.3/-3.0

- The 68% CL errors agree with the naïve scaling factor of 1.645

# Comparisons to 2011 data

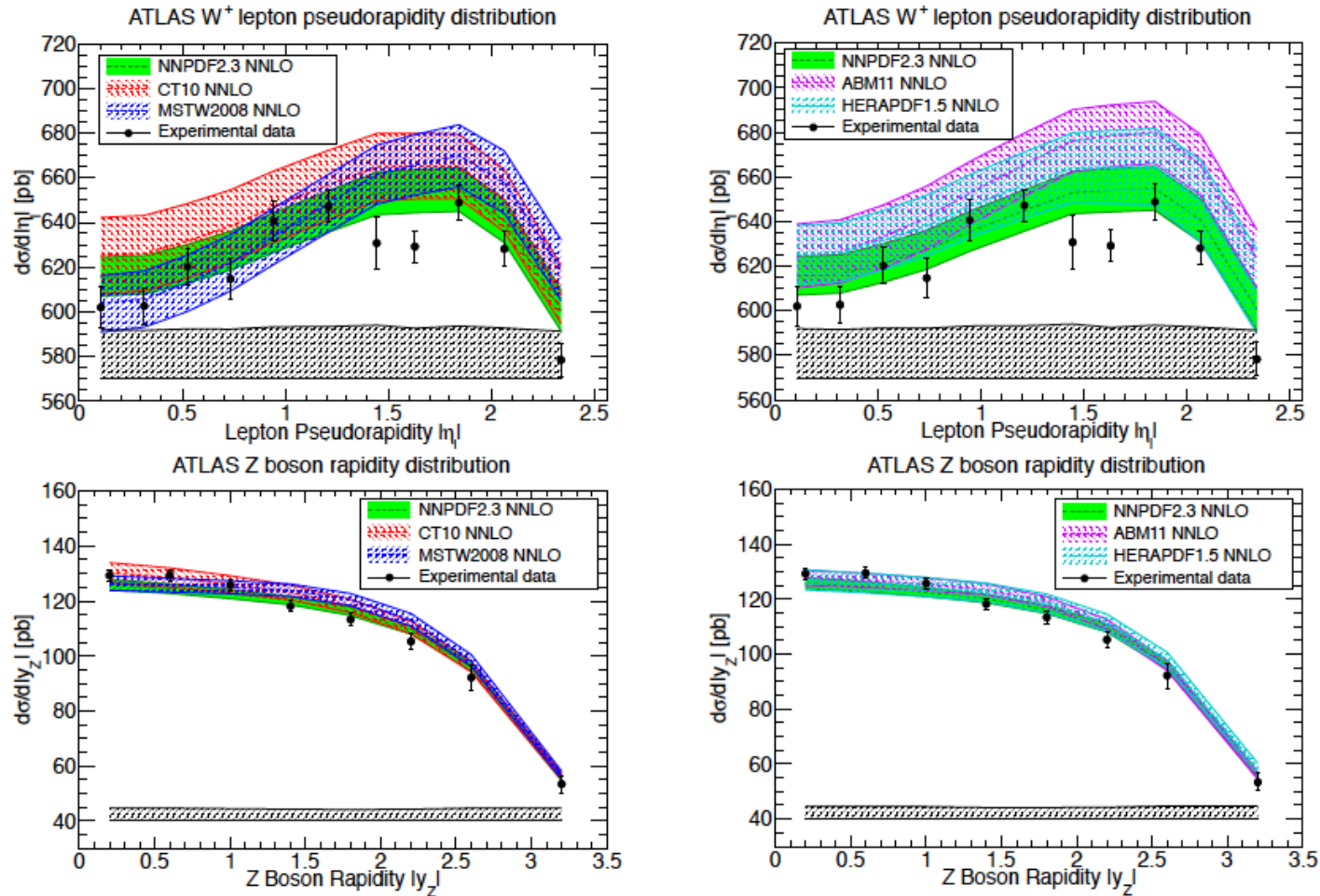


Figure 12: Comparison of the ATLAS electroweak vector boson production data with the NNPDF2.3, CT10 and MSTW2008 predictions with  $\alpha_s = 0.118$ . The error bars correspond to statistical uncertainties, while the band in the bottom of the plot indicates the correlated systematics (including normalization errors).

# Comparisons to 2011 data

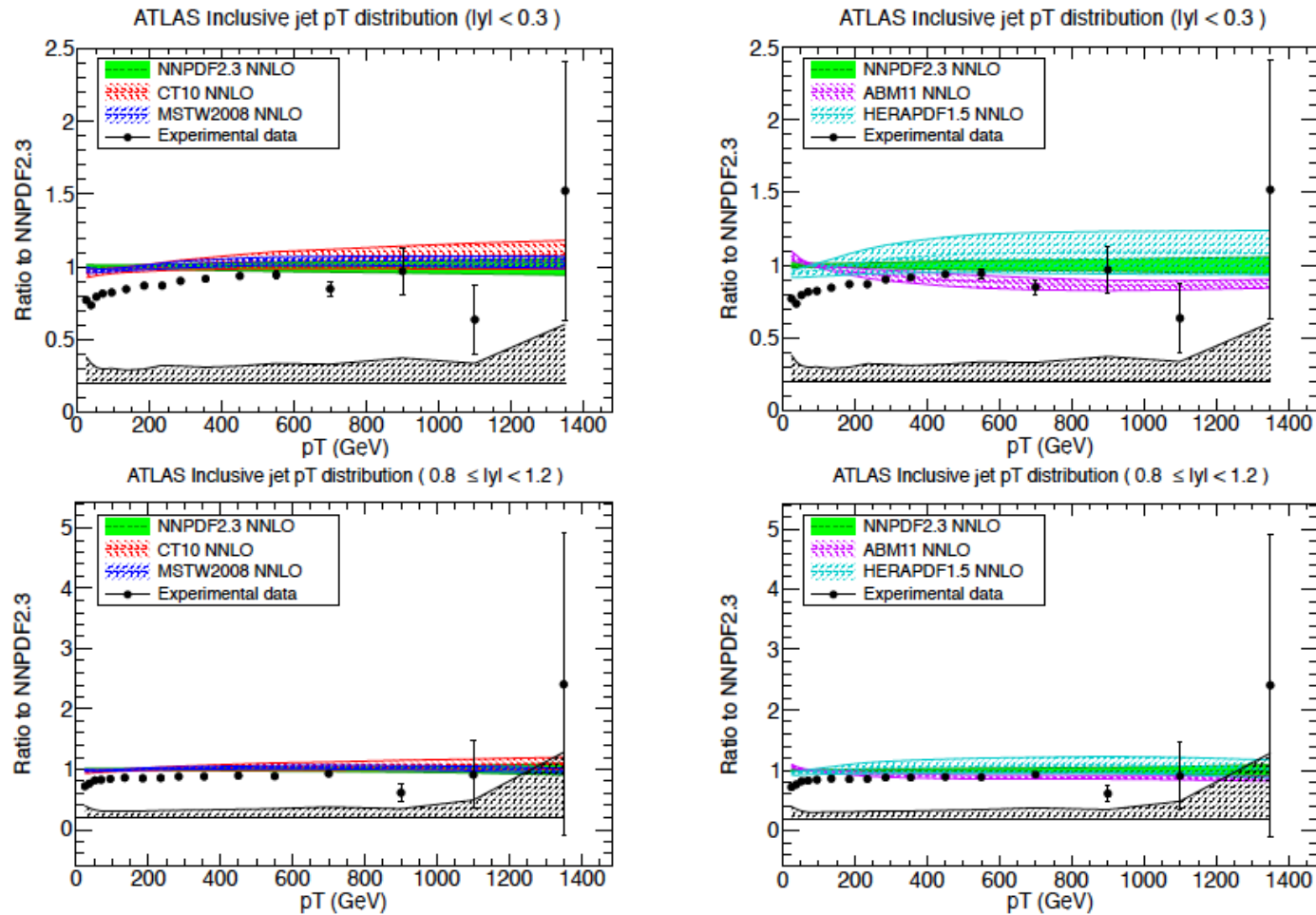


Figure 14: Comparison of the ATLAS  $R = 0.4$  inclusive jet production data from the 2010 dataset with the NNPDF2.3, CT10 and MSTW2008 NNLO PDF sets and  $\alpha_S = 0.118$ . The error bars correspond to statistical uncertainties, while the band in the bottom of the plot indicates the correlated systematics (including normalization errors)

# Comparisons to 2011 data

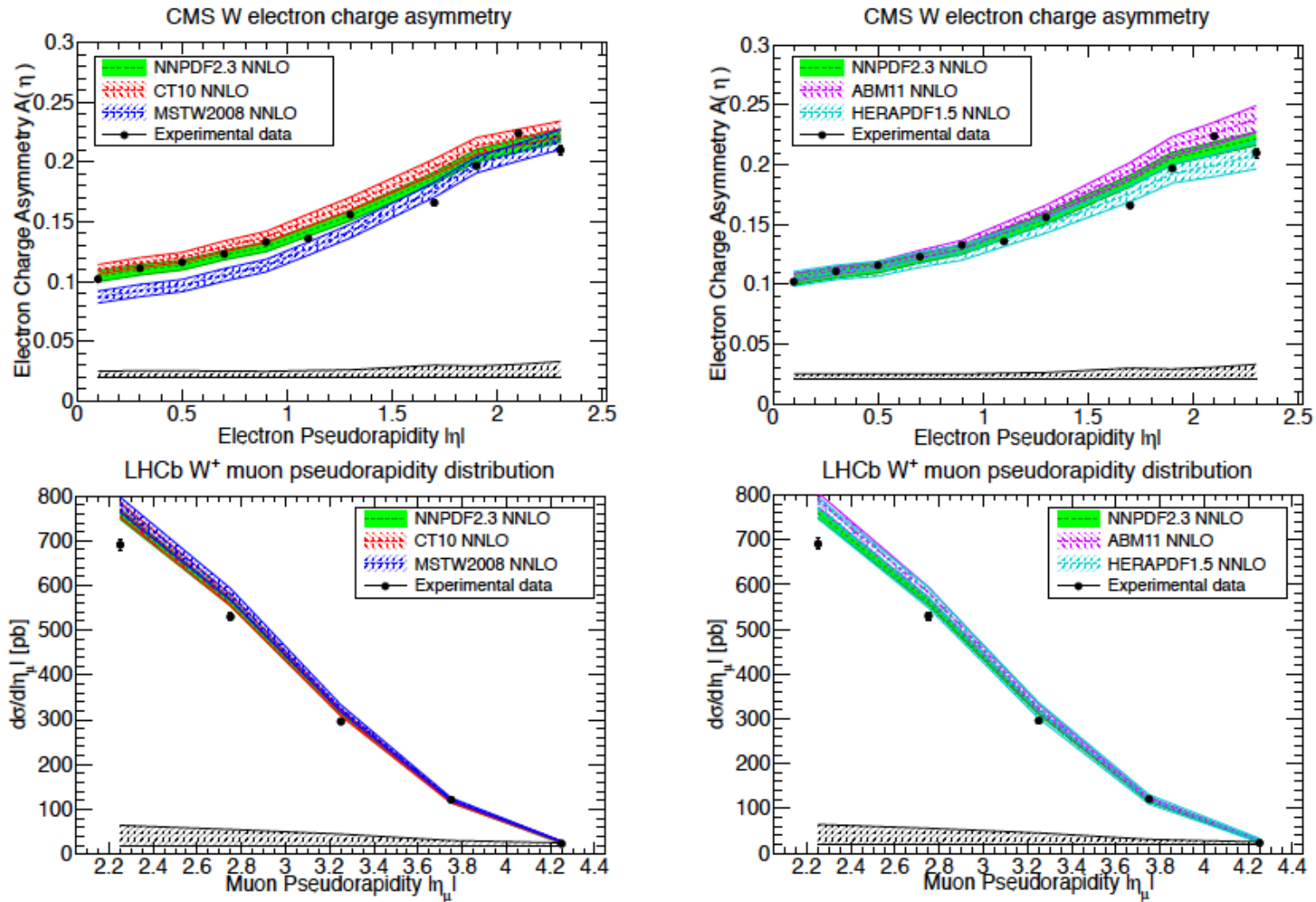


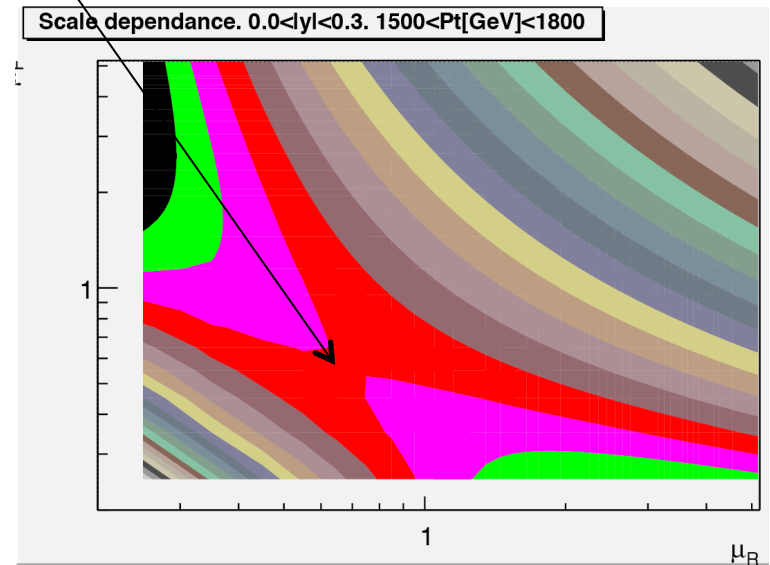
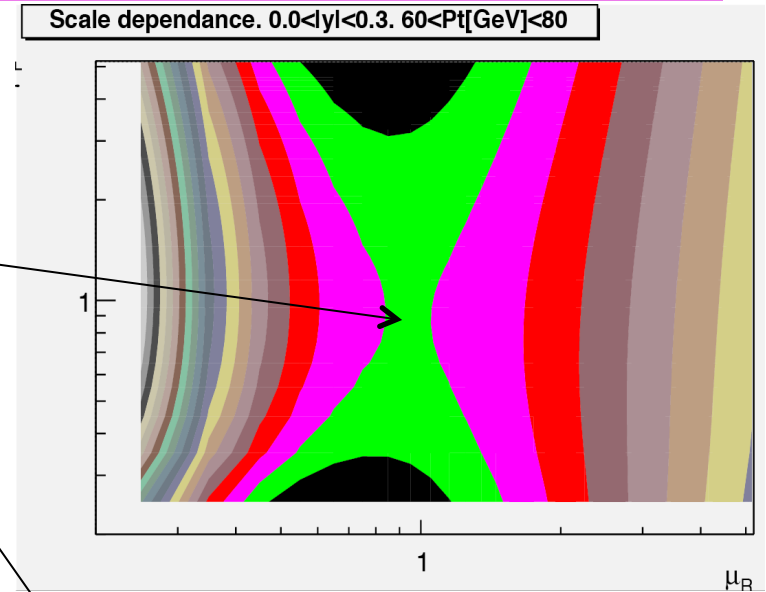
Figure 13: Same as Fig. 12 for CMS and LHCb  $W$  production.



# Aside: Scale choices

- Take inclusive jet production at the LHC
- Canonical scale choice at the LHC is  $\mu_r = \mu_f = 1.0 * p_T$ 
  - ♦ CDF used  $0.5p_T$
  - ♦ CTEQ6.6/CT10 used this scale for determination of PDFs
  - ♦ new CT PDFs use  $p_T$
- Close to saddle point for low  $p_T$
- But saddle point moves down for higher  $p_T$  (and the saddle region rotates)
- Our typical scale choices don't work for all LHC kinematics; more extreme movements for some of measured cross sections
- Rather than look for some magic formula, we should try to understand what is going on on the kinematic/scale point-of-view

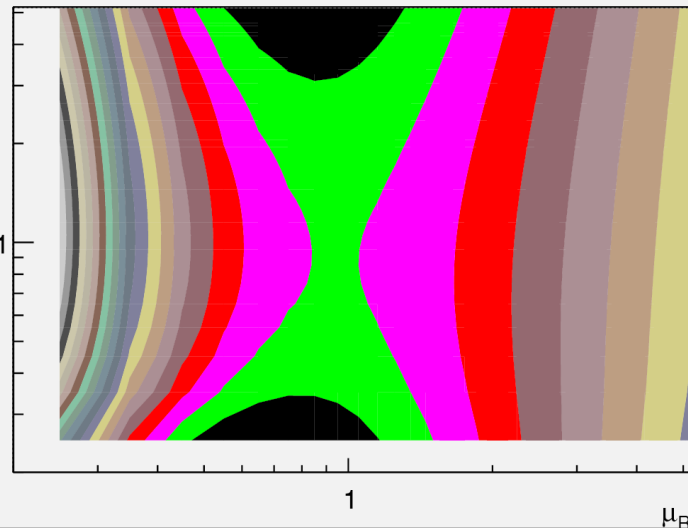
$R=0.4$   
antikT



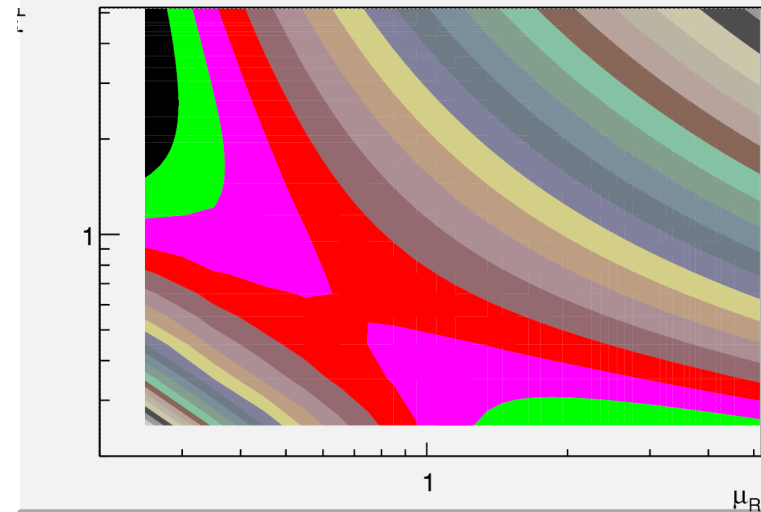
# Scale dependence also depends on jet size;

R=0.4  
antikT

Scale dependence.  $0.0 < |y| < 0.3$ .  $60 < P_t [\text{GeV}] < 80$

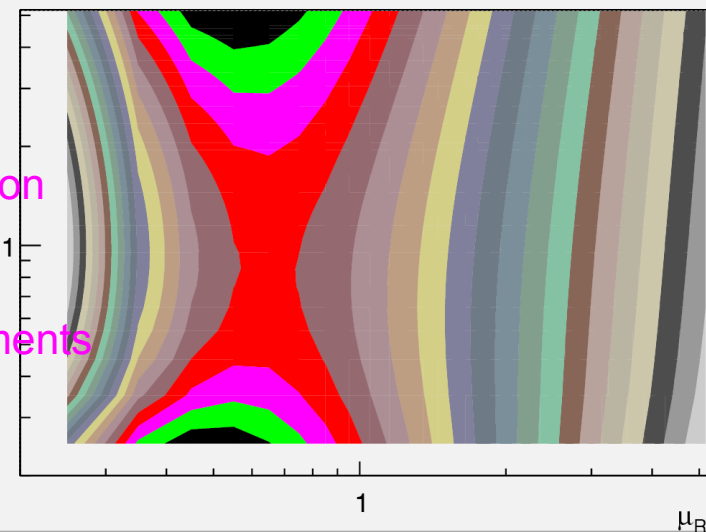


Scale dependence.  $0.0 < |y| < 0.3$ .  $1500 < P_t [\text{GeV}] < 1800$

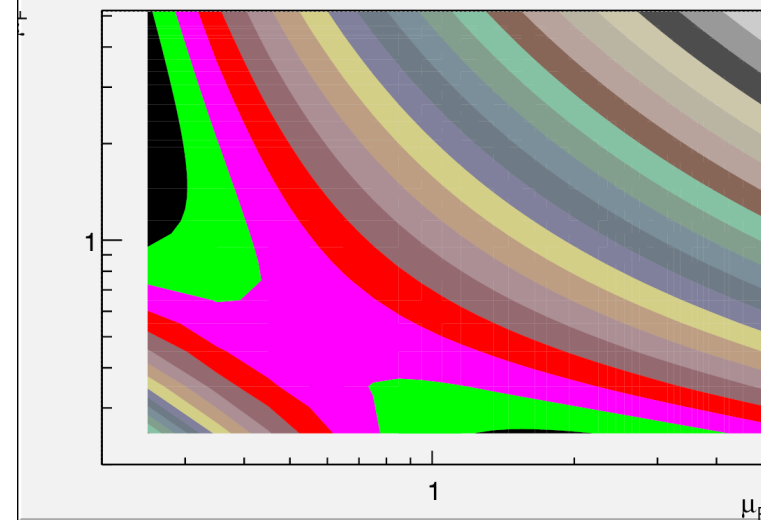


R=0.6  
antikT

Scale dependence.  $0.0 < |y| < 0.3$ .  $60 < P_t [\text{GeV}] < 80$



Scale dependence.  $0.0 < |y| < 0.3$ .  $1500 < P_t [\text{GeV}] < 1800$



NB: Tevatron  
inclusive  
jet  
measurements  
with  
R=0.7

# Calculation of $\chi^2$

Given the knowledge of the statistical, systematic and normalization uncertainties for a given experiment, we define the experimental covariance matrix used to quantify the data/theory quality as follows:

$$(\text{cov})_{IJ} = \left( \sum_{l=1}^{N_c} \sigma_{I,l} \sigma_{J,l} + \delta_{IJ} \sigma_{I,s}^2 \right) F_I F_J + \left( \sum_{n=1}^{N_a} \sigma_{I,n} \sigma_{J,n} + \sum_{n=1}^{N_r} \sigma_{I,n} \sigma_{J,n} \right) F_I F_J \quad (2)$$

where  $I$  and  $J$  run over the experimental points,  $F_I$  and  $F_J$  are the measured central values for the observables  $I$  and  $J$ . The uncertainties, given as relative values, are:  $\sigma_{I,l}$ , the  $N_c$  correlated systematic uncertainties;  $\sigma_{I,n}$ , the  $N_a$  ( $N_r$ ) absolute (relative) normalization uncertainties;  $\sigma_{I,s}$  the statistical uncertainties (which includes uncorrelated systematic uncertainties). Note that Eq. (2) cannot be used in an actual PDF fit since it is affected by the D'Agostini bias for the treatment of normalization errors [21], but it is suitable to compare predictions from different PDF sets.

Other definitions of the covariance matrix rather than Eq. (2) will lead to somewhat different results, as well as different treatments of systematic and luminosity uncertainties, can lead to somewhat different results. We will study in the appendix the impact of different definitions of the covariance matrix in the context of the ATLAS 2010 inclusive jet measurements.

# Which $\chi^2$ ?

- There are a number of  $\chi^2$  values being quoted that can differ greatly depending on the details of the definition

PDF	Code	$\chi^2$ definition			
		Eq. (A1), $\sigma_k = D_k$	Eq. (A4), $\sigma_k = D_k$	Eq. (A1), $\sigma_k = T_k(\text{CT10})$	Eq. (A1), $\sigma_k = T_k(\text{NN2.3})$
CT10	FNLO	0.95	0.95	0.55	0.60
CT10	MEKS1	1.11	1.11	0.67	0.71
CT10	MEKS2	1.00	1.00	0.65	0.68
NN2.3	FNLO	0.86	0.87	0.60	0.57
NN2.3	MEKS1	1.11	1.12	0.80	0.82
NN2.3	MEKS2	0.90	0.90	0.65	0.62
NN2.3	APPLGRID	1.00	1.00	0.64	0.58

Table II:  $\chi^2/N_{pt}$  values for the ATLAS inclusive jet production data ( $\sqrt{s} = 7$  TeV,  $R = 0.4$ ) obtained with various NLO PDFs, computer codes, and definitions of the  $\chi^2$  function. The cross sections are computed at NLO using FASTNLO (FNLO), MEKS with  $\mu_{F,R}$  equal to the individual jet  $p_T$  (MEKS1) or  $p_T$  of the hardest jet (MEKS2), and APPLGRID. The correlation matrix is obtained from the raw experimental matrix as the percentage of the central experimental value (columns 1 and 2), CT10 theoretical prediction (column 3) and NNPDF2.3 theoretical prediction (column 4).

$$\chi^2(\{a\}, \{\lambda\}) = \sum_{k=1}^{N_{pt}} \frac{1}{s_k^2} \left( D_k - T_k(\{a\}) - \sum_{\alpha=1}^{N_\lambda} \beta_{k\alpha} \lambda_\alpha \right)^2 + \sum_{\alpha=1}^{N_\lambda} \lambda_\alpha^2, \quad (\text{A1})$$

$$\tilde{\chi}^2(\{a\}, \{\lambda_0(a)\}) = \sum_{i,j=1}^{N_{pt}} (D_i - T_i) C_{ij}^{-1} (D_j - T_j) \quad C_{ij}^{-1} = \left[ \frac{\delta_{ij}}{s_i^2} - \sum_{\alpha,\beta=1}^{N_\lambda} \frac{\beta_{i\alpha}}{s_i^2} \mathcal{A}_{\alpha\beta}^{-1} \frac{\beta_{j\beta}}{s_j^2} \right]$$