

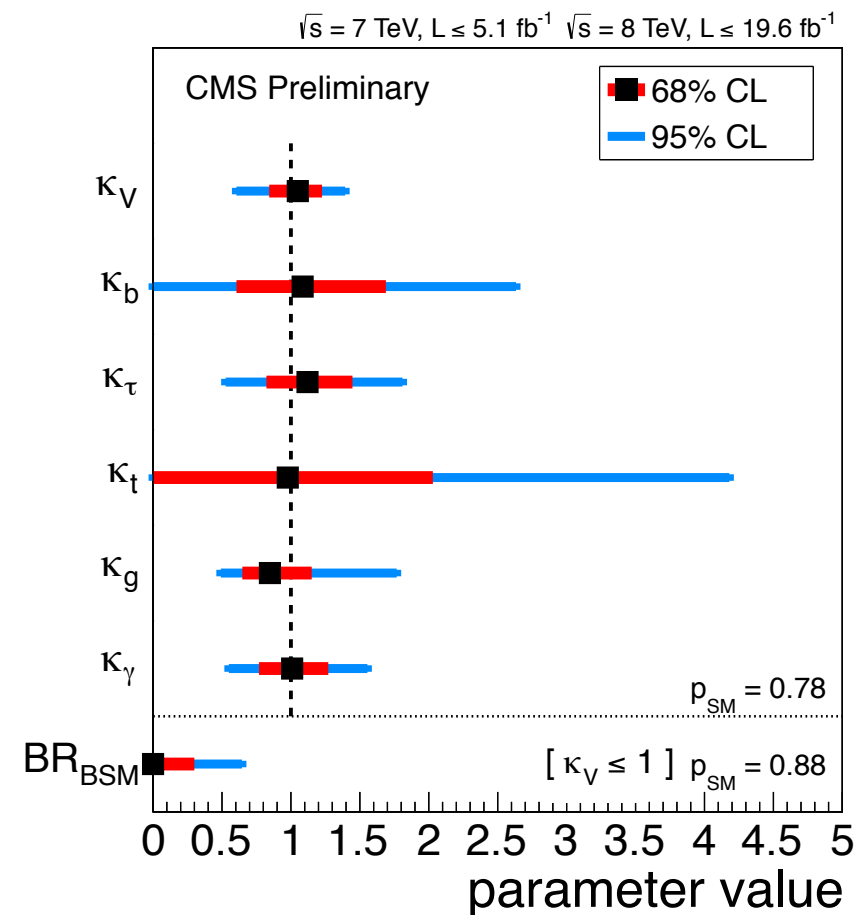
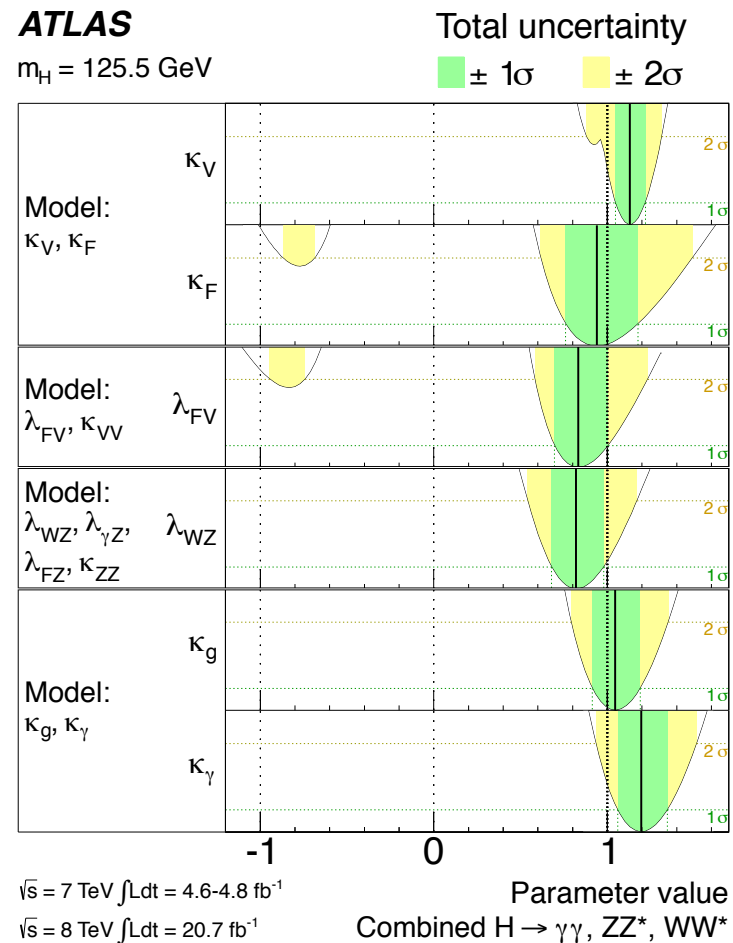
# Higgs Couplings and Electroweak Phase Transition

Maxim Perelstein, Cornell

BSM Higgs Workshop, LPC/Fermilab, November 3, 2014

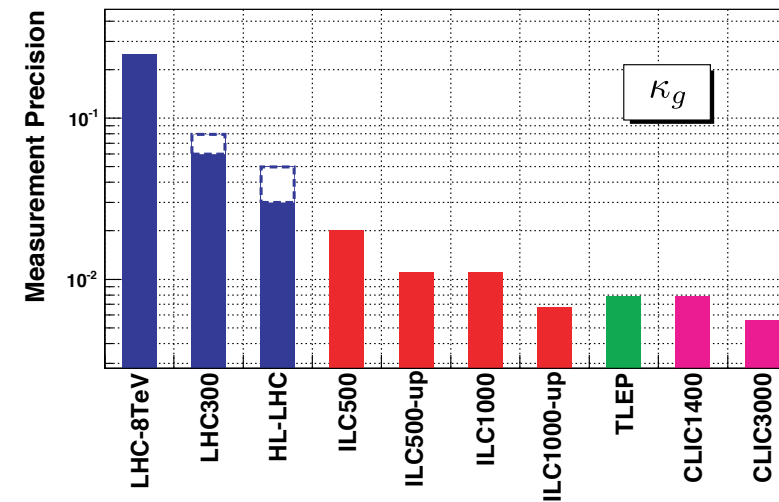
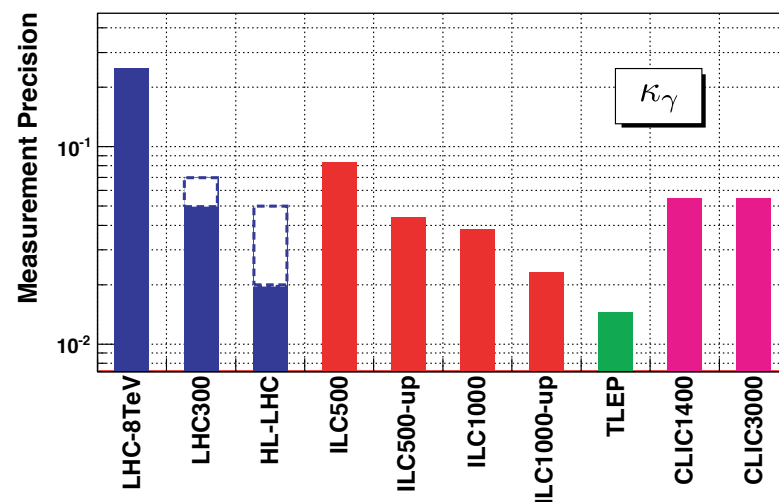
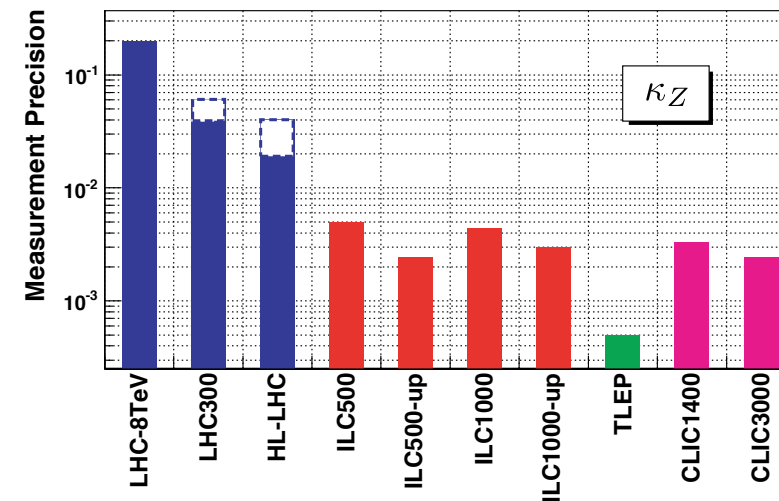
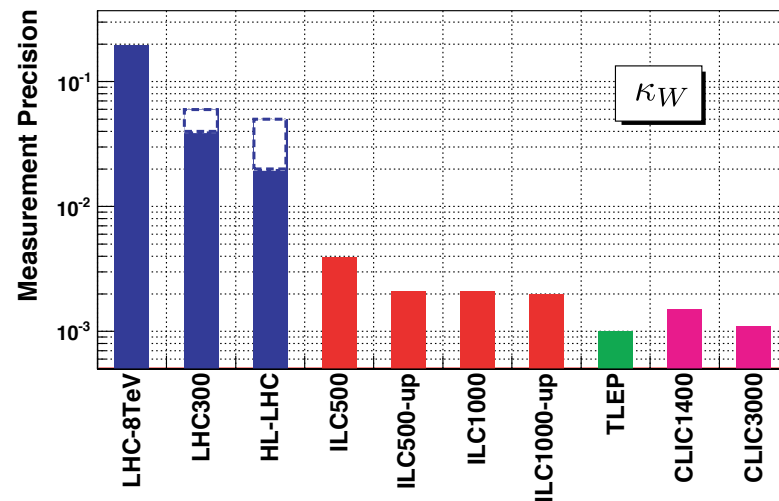
[Based on: Andrey Katz, MP, 1401.1827, JHEP]

# Higgs: Discovery to Precision



Current status: **~10-20%** precision on  
 some couplings (**V, g, gamma**)

# Future: Higgs Precision Program



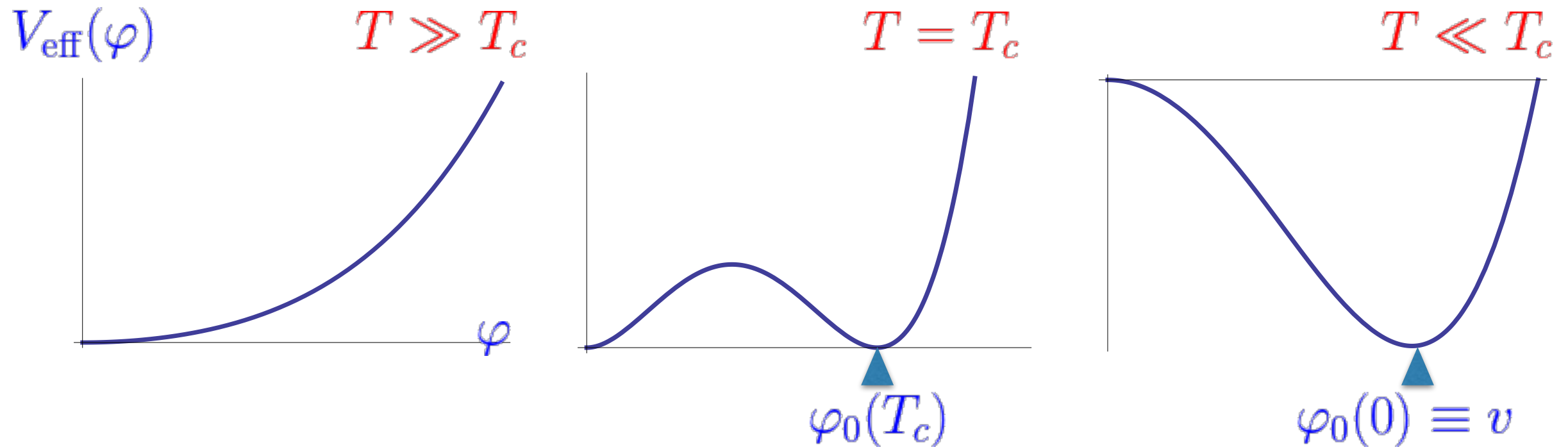
Snowmass Higgs Report

- 2-5% precision achievable at the HL-LHC
- 0.1% precision on  $V$ , 1% on  $g$  and  $\gamma$  at  $e^+e^-$  Higgs factories
- At or almost at precision electroweak levels!

# Electroweak Phase Transition

- At high temperatures (e.g. in the early Universe), electroweak symmetry is restored:  $m_{\text{eff}}^2 = -m^2 + cT^2 > 0 \Rightarrow \langle \varphi \rangle = 0$
- Electroweak Phase Transition into the current broken-EW phase occurred about  $10^{-10}$  sec after the Big Bang ( $T \sim 100$  GeV )
- Baryon asymmetry may have been produced during this phase transition - “electroweak baryogenesis”
- A strongly first-order transition is required for successful EWBG
- In the SM, transition is second order; BSM physics at the weak scale can modify dynamics, inducing a 1st order transition

# First-Order EWPT in Cartoons



- “Transition strength”  $\sim$  entropy release  $\xi = \varphi_0(T_c)/T_c$
- Numerical studies: EW Baryogenesis possible if  $\xi \geq 0.9$
- Otherwise, sphaleron washout of the baryon number

# HC and EWPT

- No possibility of producing “plasma” with restored EW symmetry (T-RHIC?) so no direct experimental probe
  - However, hard to induce large modifications of the **finite-T** potential without also modifying **T=0** Higgs potential and couplings
  - Can precise measurements of Higgs couplings conclusively probe the nature of EWPT?
  - Two basic mechanisms for first-order EWPT: **tree-level mixing** with other scalars; and **loop-induced** corrections (the famous  $T^3$  term)
  - We focused on **loop-y models** since they seem harder to probe\*
- [\* a study of tree-y models is now in progress...]

# HC and EWPT: Setup

- The cubic term at high-T is induced by loops of scalars, not fermions
- Add a single complex scalar  $\Phi$ , with  $V_\Phi = m_0^2|\Phi|^2 + \kappa|\Phi|^2|H|^2 + \eta|\Phi|^4$ .
- One-loop corrections to the potential at both T=0 and finite-T are well known:

$$V_1(\varphi) = \frac{g_i(-1)^{F_i}}{64\pi^2} \left[ m_i^4(\varphi) \log \frac{m_i^2(\varphi)}{m_i^2(v)} - \frac{3}{2} m_i^4(\varphi) + 2m_i^2(\varphi)m_i^2(v) \right];$$

$$V_T(\varphi; T) = \frac{g_i T^4 (-1)^{F_i}}{2\pi^2} \int_0^\infty dx x^2 \log \left[ 1 - (-1)^{F_i} \exp \left( \sqrt{x^2 + \frac{m_i^2(\varphi)}{T^2}} \right) \right],$$

- The key object is the Higgs-dependent  $\Phi$  mass! But recall:

$$\mathcal{L}_{h\gamma\gamma} = \frac{2\alpha}{9\pi v} C_\gamma h F_{\mu\nu} F^{\mu\nu}, \quad \mathcal{L}_{hgg} = \frac{\alpha_s}{12\pi v} C_g h G_{\mu\nu} G^{\mu\nu}$$

$$C_\gamma = 1 + \frac{3}{8} \sum_f^{\text{Dirac fermions}} N_{c,f} Q_f^2 \frac{\partial \ln m_f^2(v)}{\partial \ln v} + \frac{3}{32} \sum_s^{\text{scalars}} N_{c,s} Q_s^2 \frac{\partial \ln m_s^2(v)}{\partial \ln v}$$

$$C_g = 1 + \sum_f^{\text{Dirac fermions}} C(r_f) \frac{\partial \ln m_f^2(v)}{\partial \ln v} + \frac{1}{4} \sum_s^{\text{scalars}} C(r_s) \frac{\partial \ln m_s^2(v)}{\partial \ln v},$$

- Expect **direct correlation** between the size of the cubic coupling induced at finite-T and non-SM contributions to  $hgg$  and  $h\gamma\gamma$  (unless  $\Phi$  is color and EM-neutral)

# Analytic Example

- A special case can be studied analytically\*:  $m_0 = 0$

- High-temperature expansion of the thermal potential:

$$V_{\text{eff}}(\varphi; T) = V_0(\varphi) + V_T(\varphi; T) \approx \frac{1}{2} \left( -\mu^2 + \frac{g_\Phi \kappa T^2}{24} \right) \varphi^2 - \frac{g_\Phi \kappa^{3/2} T}{24\sqrt{2}\pi} \varphi^3 + \frac{\lambda}{4} \varphi^4.$$

- Location of the broken-symmetry minimum at finite T:  $\frac{\partial V_{\text{eff}}}{\partial \varphi} = 0 \Rightarrow \varphi_0(T)$

- Critical temperature:  $V_{\text{eff}}(0, T_c) = V_{\text{eff}}(\varphi_0(T_c), T_c)$

- Solve together:  $T_c^2 = \frac{24\mu^2}{g_\Phi \kappa \left( 1 - \frac{g_\Phi \kappa^2}{24\pi^2 \lambda} \right)}, \quad \varphi_+(T_c) = \frac{g_\Phi \kappa^{3/2} T_c}{12\sqrt{2}\pi \lambda}.$

- Strongly 1-st order if  $\kappa > 3.6 g_\Phi^{-2/3}$

- Gluon-Higgs coupling:  $R_g = \frac{1}{8} \frac{\kappa v^2}{m_0^2 + \frac{\kappa v^2}{2}} \approx 25\%$



# Numerical Studies

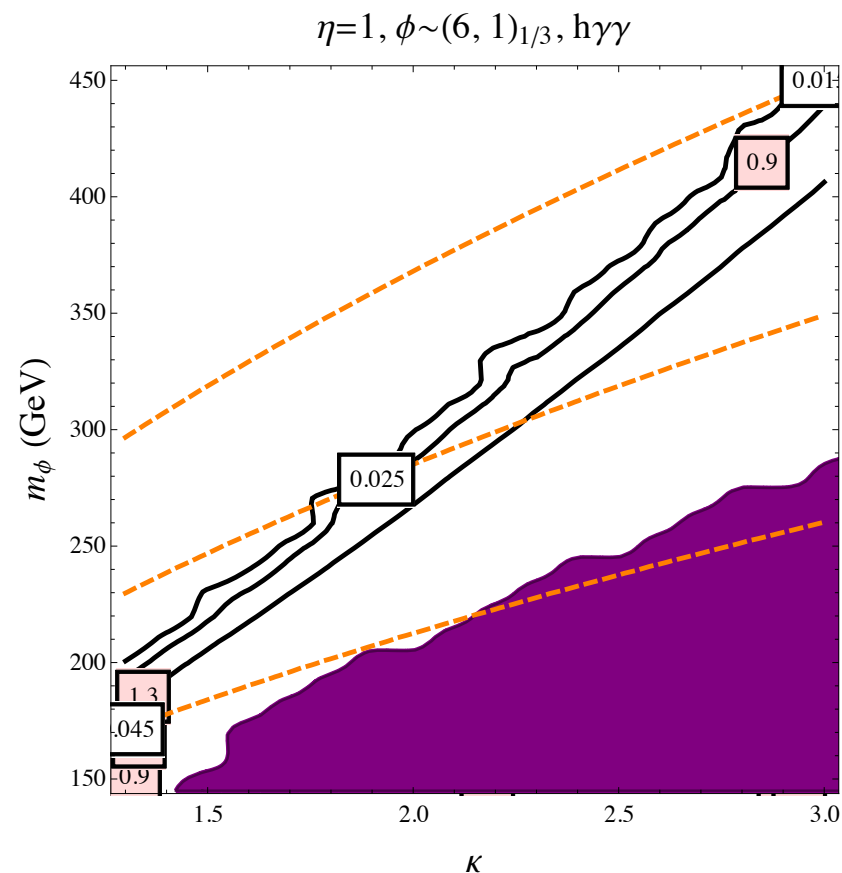
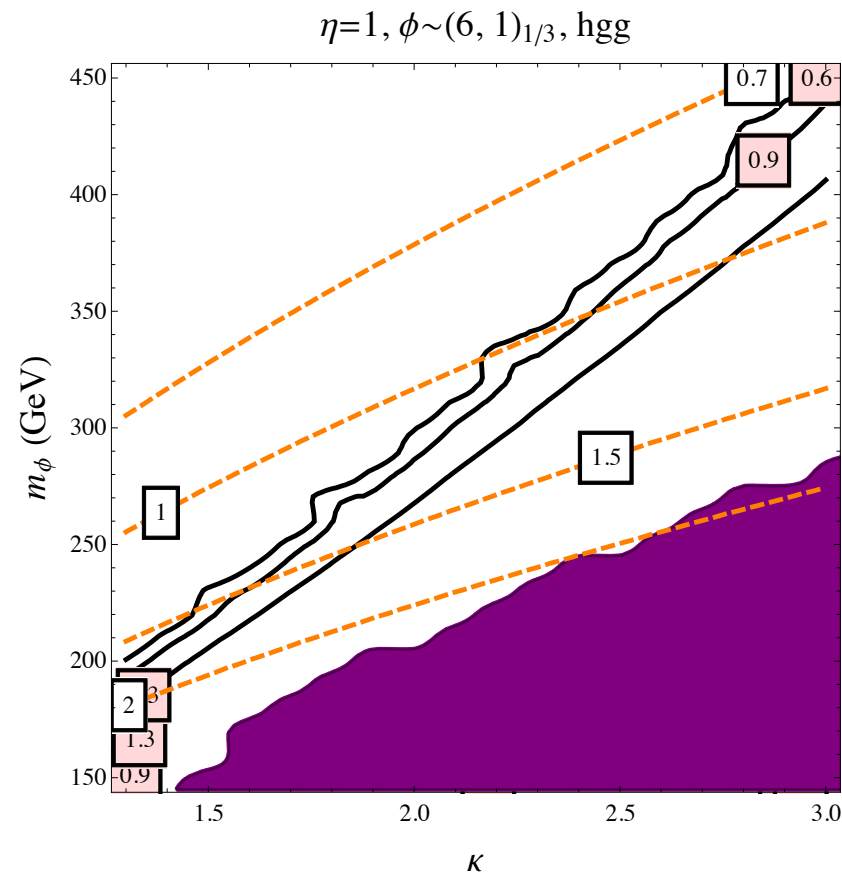
- In general, no analytic solution for critical T and order parameter - solve numerically
- Numerical code also includes SM contributions, “daisy resummations” etc.
- Analyzed a few toy models, representative of the range of possibilities for quantum numbers of the  $\Phi$  field

Model	$(SU(3), SU(2))_{U(1)}$	$g_\Phi$	$C_3$	$C_2$	$\frac{\Pi_W}{g^2 T^2}$	$\frac{\Pi_B}{g'^2 T^2}$	$\frac{\Delta\Pi_h}{\kappa T^2}$
“RH stop”	$(\bar{3}, 1)_{-2/3}$	6	4/3	0	11/6	107/54	1/4
Exotic triplet	$(3, 1)_{-4/3}$	6	4/3	0	11/6	131/54	1/4
Exotic sextet	$(\bar{6}, 1)_{8/3}$	12	10/3	0	11/6	227/54	1/2
“LH stau”	$(1, 2)_{-1/2}$	4	0	3/4	2	23/12	1/6
“RH stau”	$(1, 1)_1$	2	0	0	11/6	13/6	1/12
Singlet	$(1, 1)_0$	2	0	0	11/6	11/6	1/12

Table 1. Benchmark models studied in this paper.

[\* we treat  $\kappa$  as a free parameter, unlike SUSY]

# Results: “Sextet”

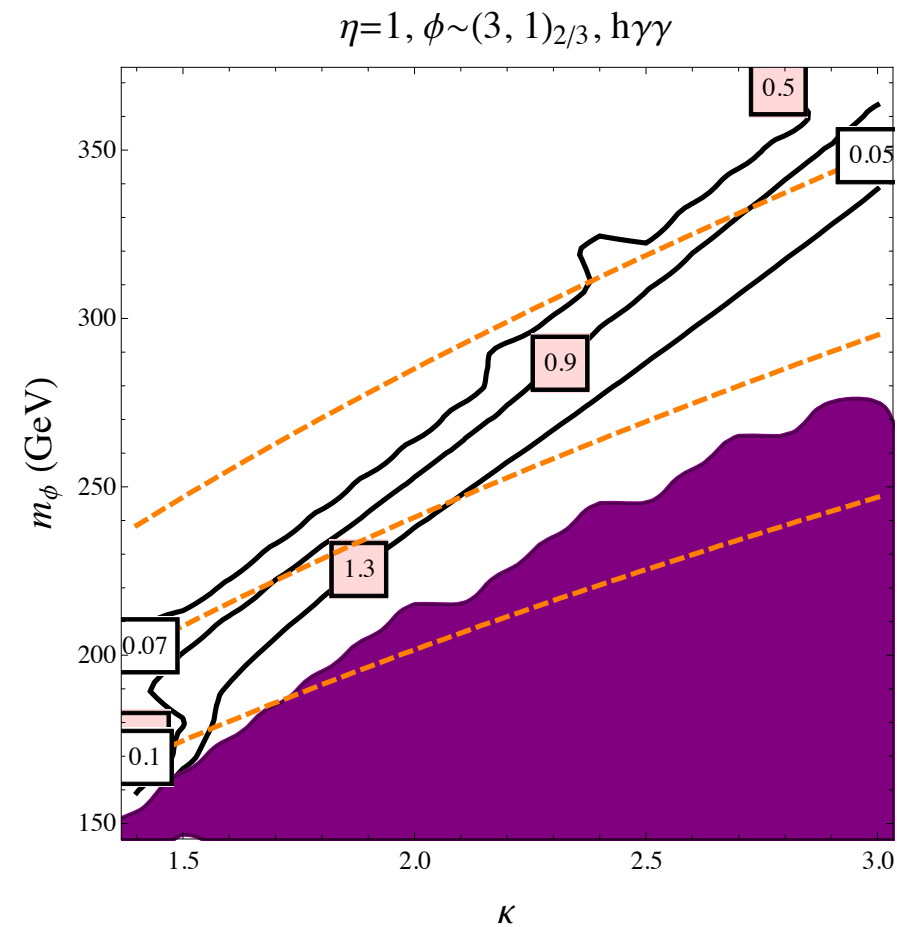
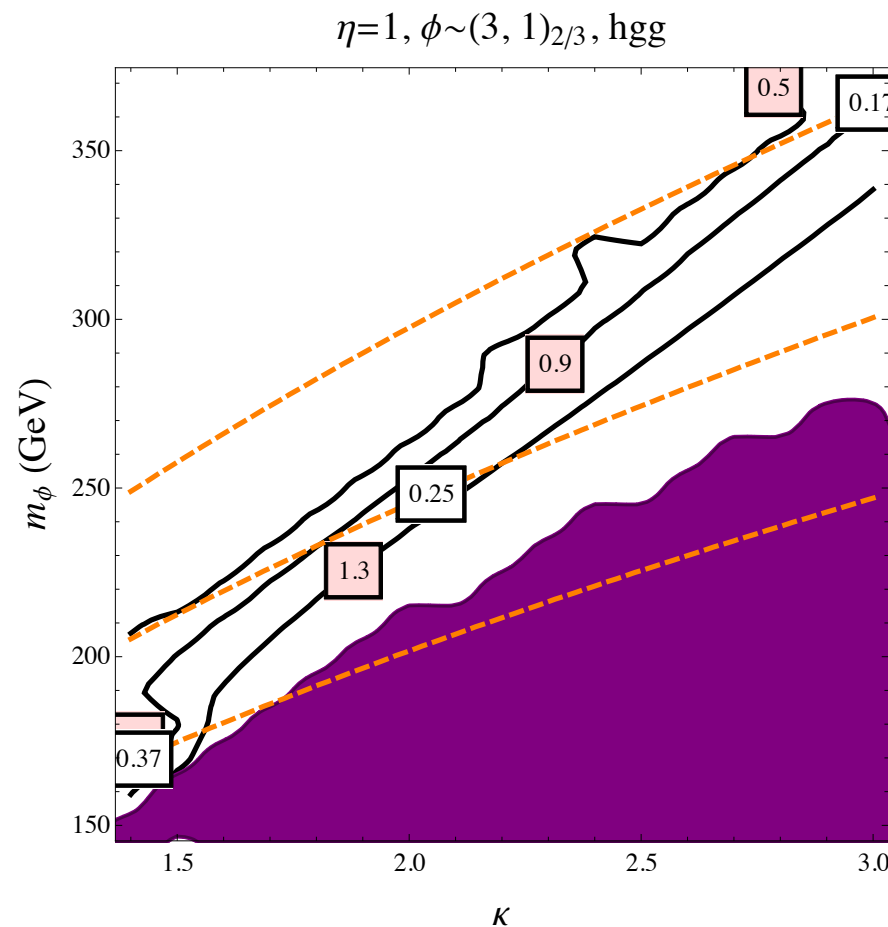


ATLAS:  $R_g = 1.08 \pm 0.14$ ,  
 $R_\gamma = 1.23^{+0.16}_{-0.13}$ .  $\rightarrow$  ruled out!\*

[\* usual caveat: SM total width assumed]

**NOTE:** Our sextet can decay to 4 jets  $\rightarrow$  no direct search!

# Results: “RH Stop”



NOT ruled out if  $\kappa > \kappa_{\text{MSSM}}$ !

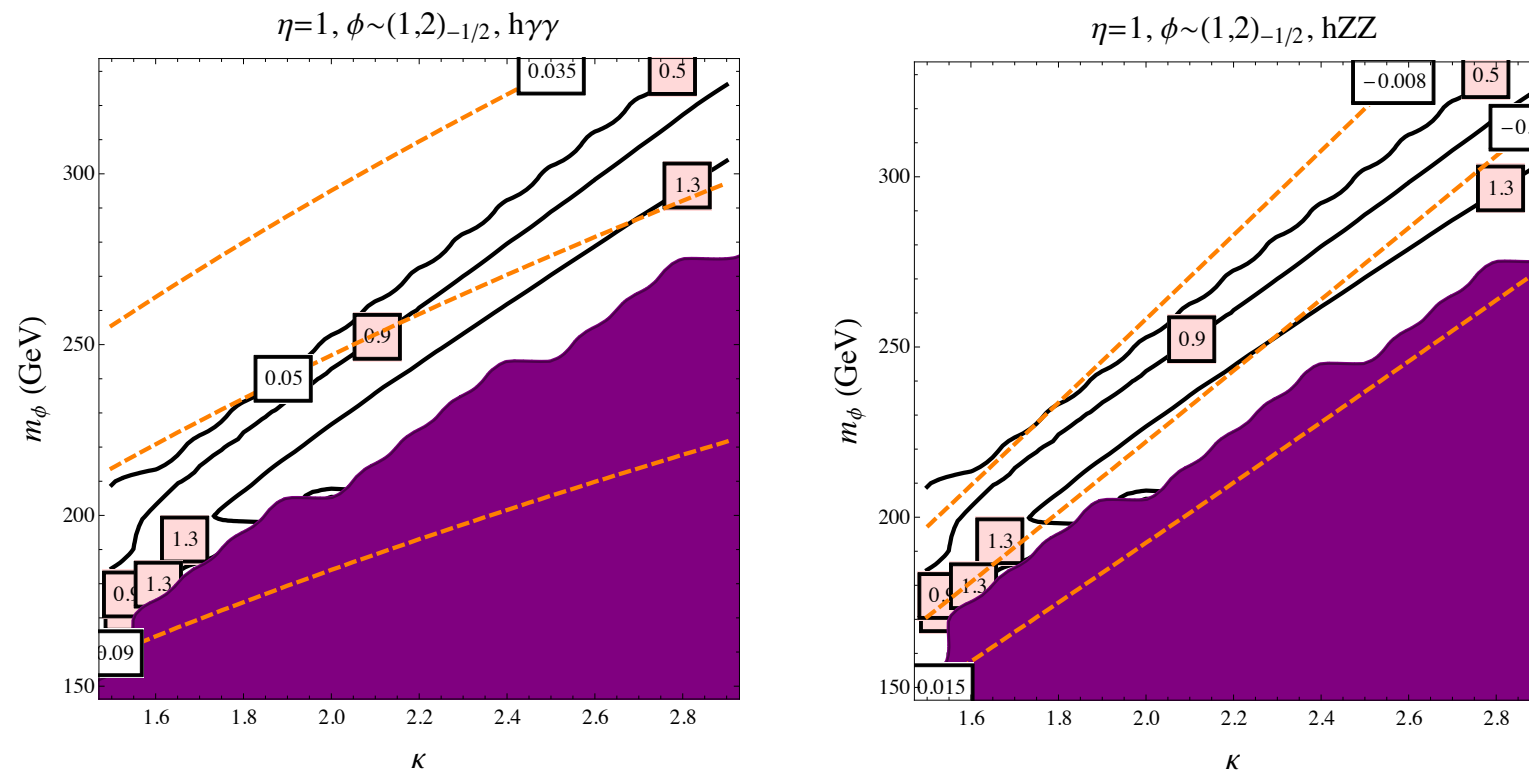
MIN deviation  $\sim 17\%$ , probed at 3-sigma at LHC-14

NOTE: The “RH stop” can decay to 2 jets or be “stealthy/ compressed” → avoid direct searches!

# Higgs and a Singlet

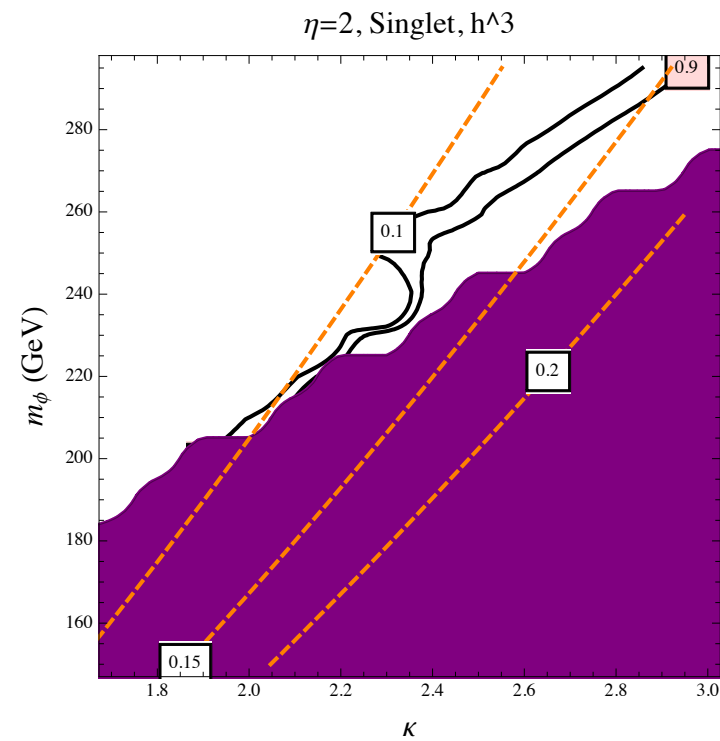
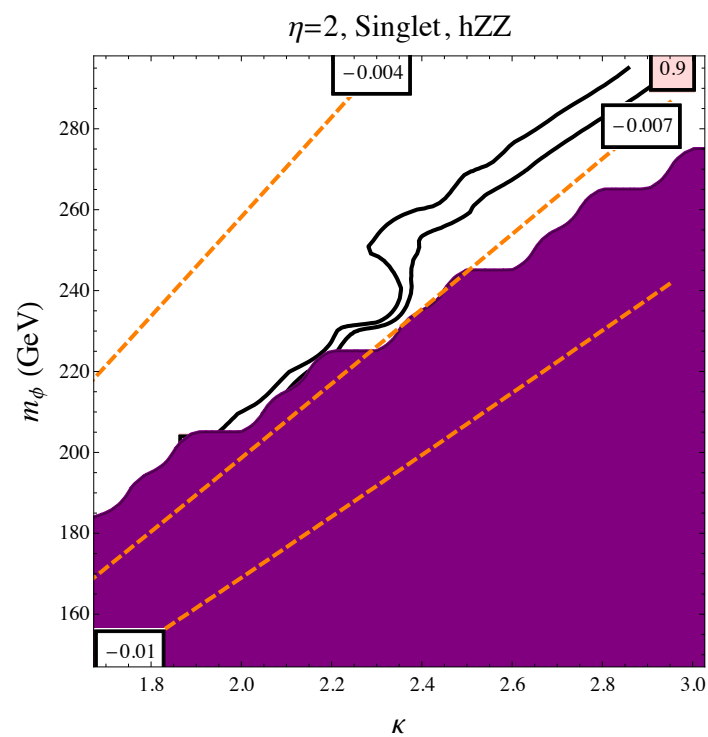
- $\Phi$  does not need to have SM gauge interactions to drive a first-order EWPT
- Obviously this scenario would not produce any deviation in  $hgg$  or  $h\gamma\gamma$
- However, it does predict a (small) deviation in  $hZZ$  coupling [Craig, Englert, McCullough, 1305.5251]
- Consider  $m_\Phi \gg v$ , integrate it out  $\rightarrow$  a dim-6 operator:  $\frac{\kappa^2}{16\pi^2 m_\Phi^2} (\partial_\mu |H|^2)^2$
- After Higgs gets a vev:  $\frac{\kappa^2}{16\pi^2 m_\Phi^2} (\partial_\mu h)^2$
- Canonically normalized Higgs  $\rightarrow$  shift in  $hZZ$  coupling
- Effect is small, but  $hZZ$  coupling can be determined very precisely from Higgsstrahlung cross section:  $\sim 0.25\%$  ILC,  $\sim 0.05\%$  “TLEP” [Snowmass Higgs report]

# Results: “LH Stau”



**hZZ: MIN** deviation **0.8%**, probed at 3-sigma at ILC

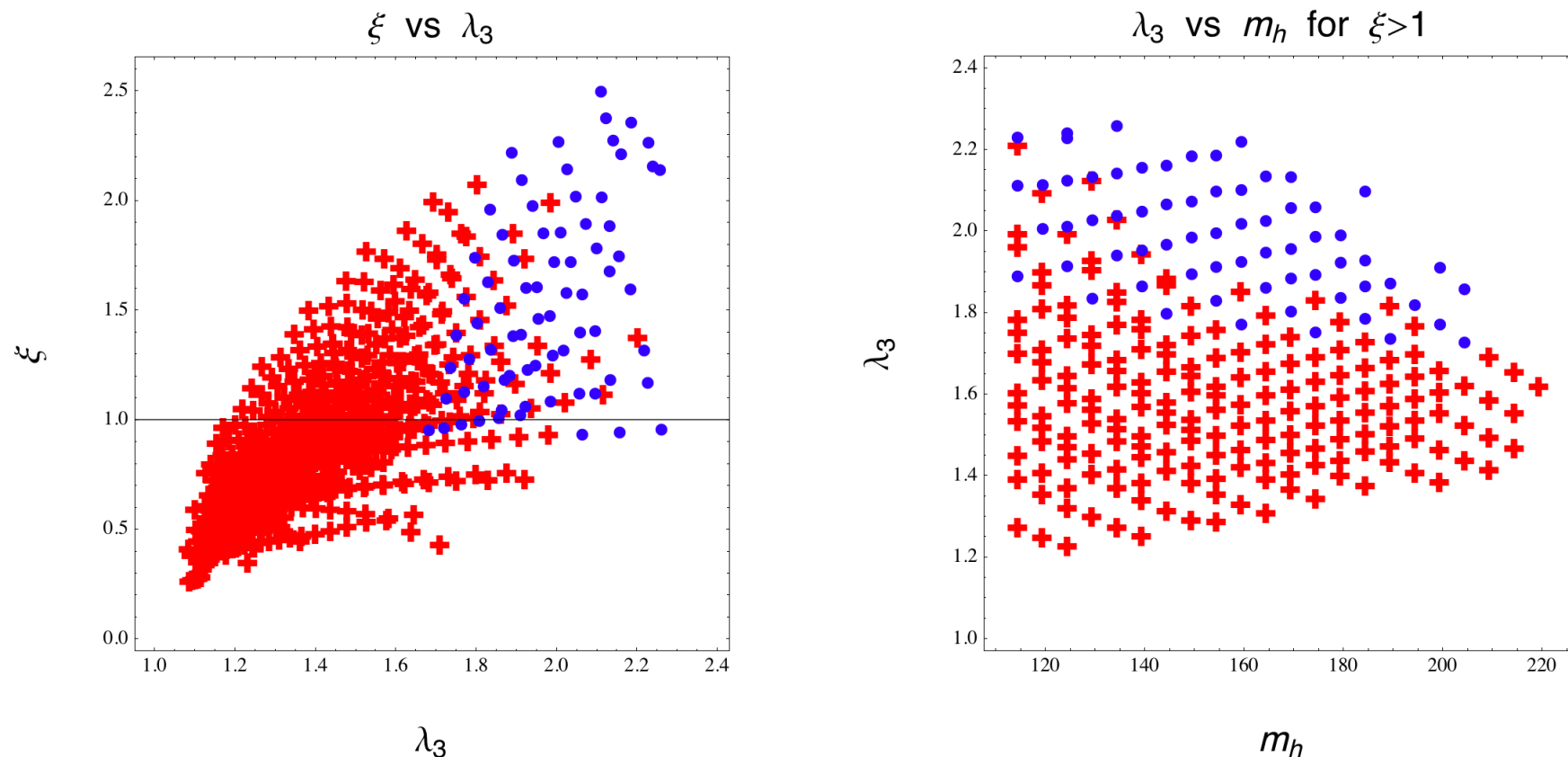
# Results: Singlet



**hZZ: MIN** deviation **0.5%**, probed at  $\sim 2$ -sigma at ILC,  
10-sigma at “TLEP”

# Higgs Self-Coupling

[Noble, MP, 0711.3018]



same correlation for Higgs self-coupling: deviations of **20% or more** in a broad range of models with first-order EWPT

Measure it at the ILC-1 TeV? a 100 TeV collider?

# Conclusions: EWPT

- Strongly first-order EWPT, and with it Electroweak Baryogenesis, remains a **viable possibility** in a general BSM context
- We focused on the models where first-order EWPT is induced by **loops** of a BSM scalar, with various SM quantum numbers
- In the case of colored scalar, **LHC-14** measurement of *hgg* will be able to conclusively probe the full parameter space with a 1-st order EWPT
- For non-colored scalars, **e<sup>+</sup>e<sup>-</sup> Higgs factories** will be necessary
- **Higgs factory may** be able to conclusively probe the full parameter space with 1-st order EWPT in **all models**, even if induced by a SM-singlet scalar