Neutrino-induced meson productions in resonance region

arXiv:1506.03403

Satoshi Nakamura

Osaka University, Japan

Collaborators : H. Kamano (RCNP, Osaka Univ.), T. Sato (Osaka Univ.)

Contents

★ Introduction νN scattering in resonance region

★ Dynamical coupled-channels (DCC) model

Analysis of $\gamma N,\ \pi N \rightarrow \pi N,\ \eta N,\ K\Lambda,\ K\Sigma$ data

and electron scattering data

Extension to $\nu N \rightarrow l^- X$ ($X = \pi N,\ \pi\pi N,\ \eta N,\ K\Lambda,\ K\Sigma$)

★ Results for $\nu N \rightarrow l^- X$

Introduction

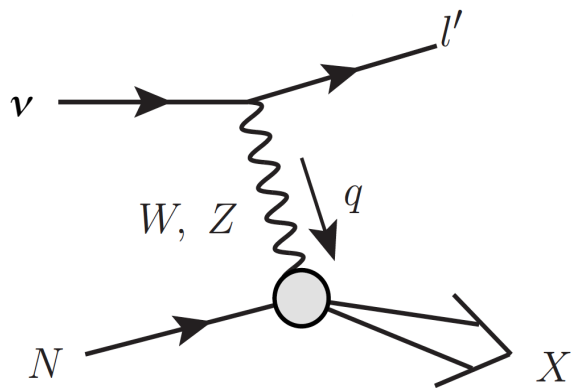
Neutrino-nucleus scattering for ν-oscillation experiments

Next-generation exp. ➔ leptonic $\not{C}\not{P}$, mass hierarchy

ν−nucleus scattering needs to be understood more precisely

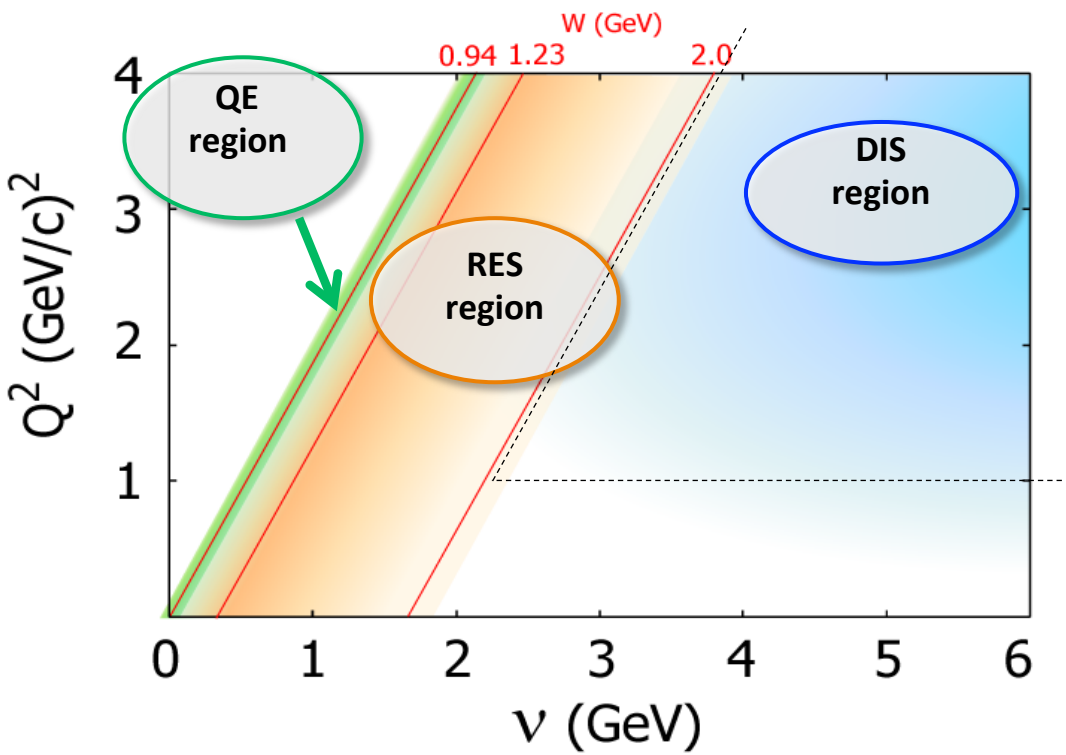
All ν-oscillation experiments measure ν-flux through ν-nucleus interaction

$Q^2 = -q^2$
$\nu = q^0$



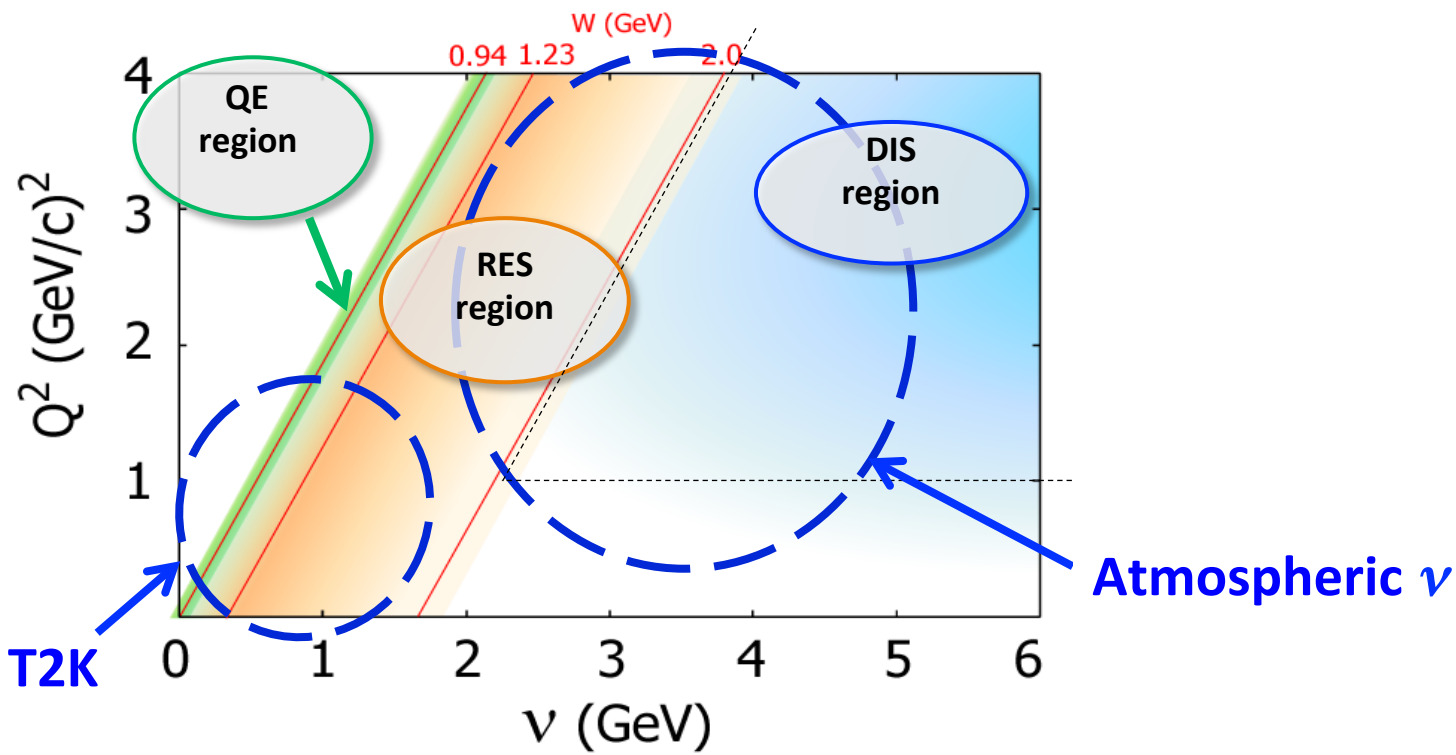
Neutrino-nucleus scattering for ν-oscillation experiments

Next-generation exp. ➔ leptonic $\not{C}P$, mass hierarchy



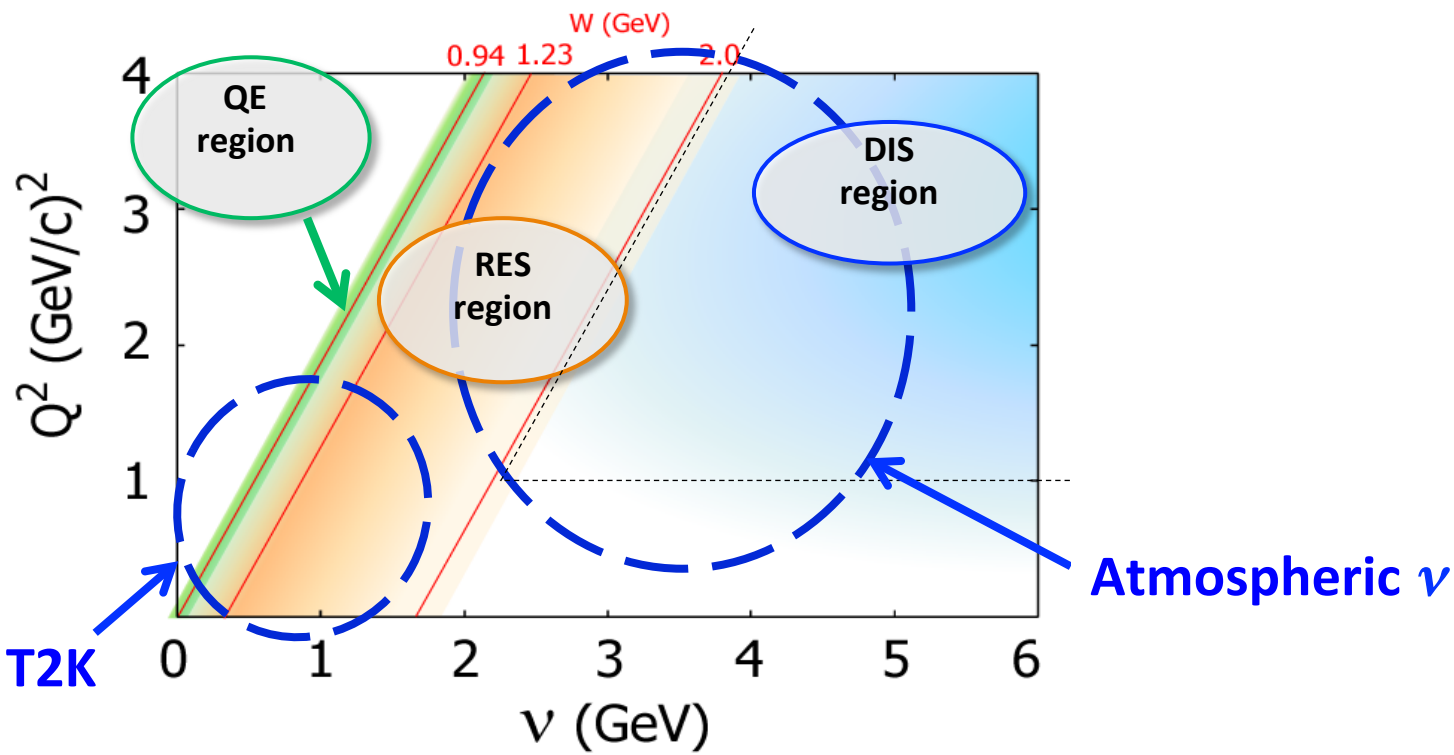
Neutrino-nucleus scattering for ν-oscillation experiments

Next-generation exp. ➔ leptonic $\not{C}\not{P}$, mass hierarchy



Neutrino-nucleus scattering for ν-oscillation experiments

Next-generation exp. ➔ leptonic $\not{C}P$, mass hierarchy

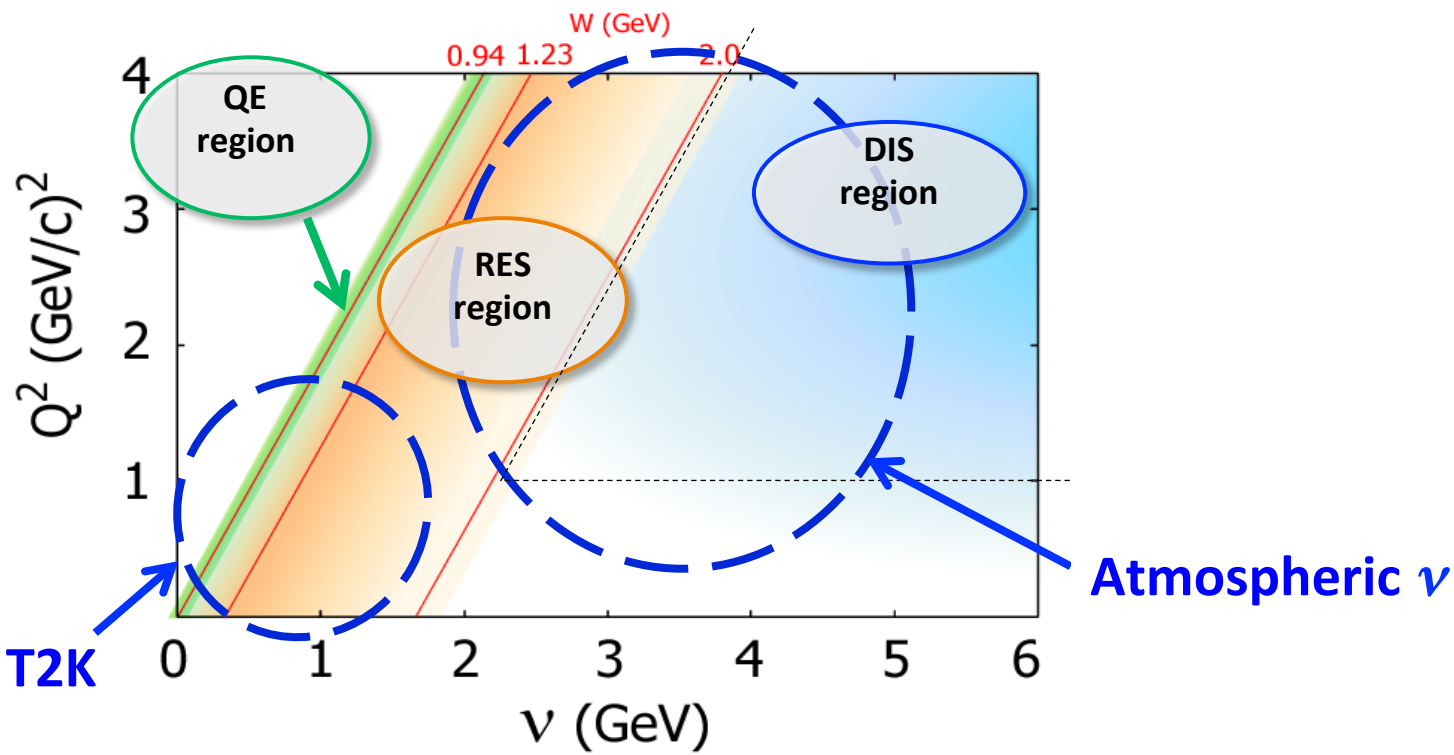


Wide kinematical region with different characteristic

➔ Combination of different expertise is necessary

Neutrino-nucleus scattering for ν-oscillation experiments

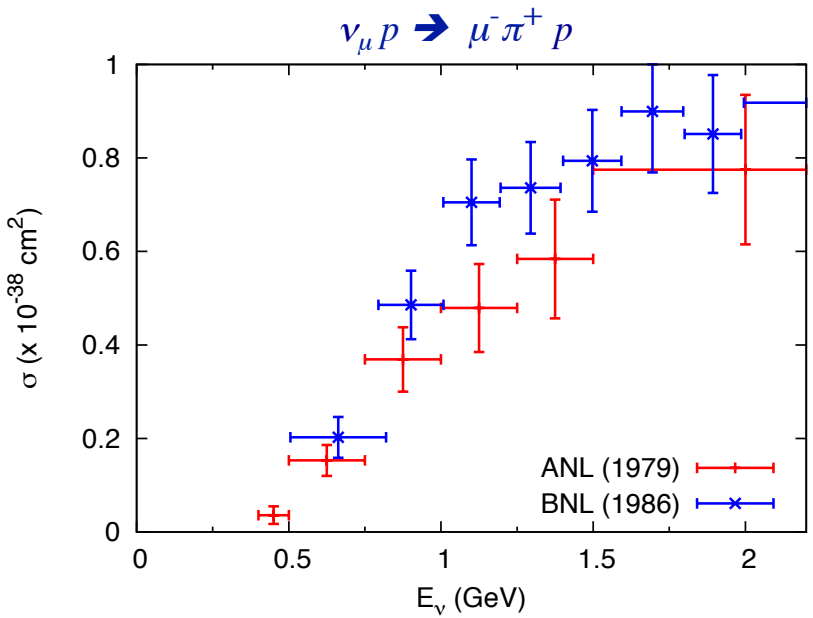
Next-generation exp. ➔ leptonic $\not{C}\not{P}$, mass hierarchy

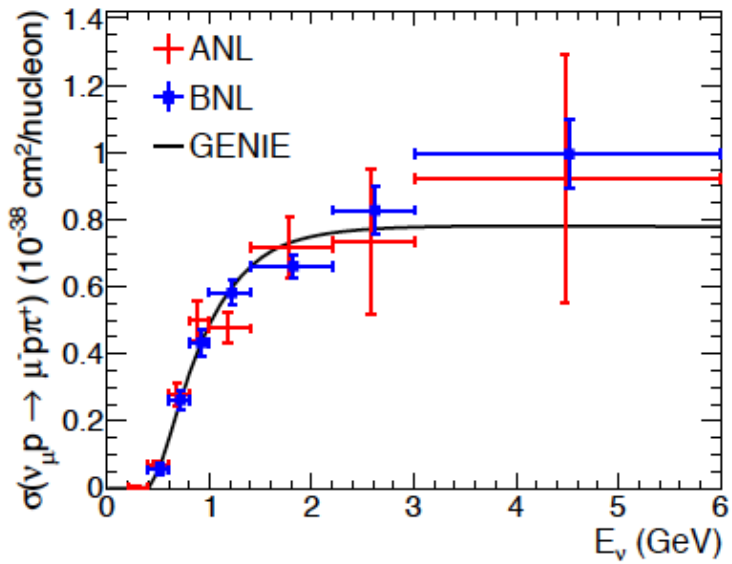


Collaboration at J-PARC Branch of KEK Theory Center

http://j-parc-th.kek.jp/html/English/e-index.html

Neutrino interaction data in resonance region





- Data to fix nucleon axial current ($g_{AN\Delta}$)
- Discrepancy between BNL & ANL data

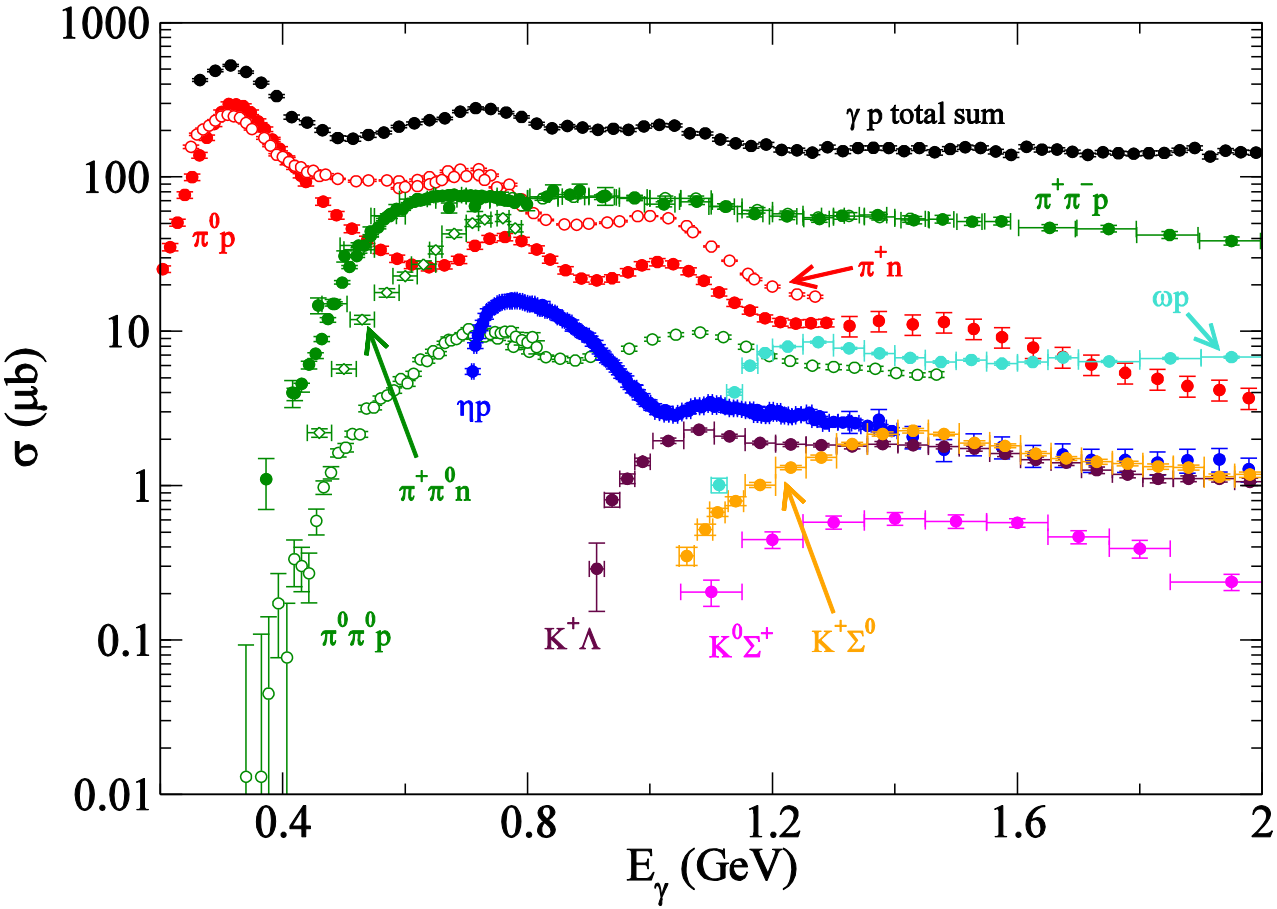
- Recent reanalysis of original data

  → discrepancy resolved (!?)

  PRD 90, 112017 (2014)

Resonance region (single nucleon)

$\gamma N \rightarrow X$



Multi-channel reaction

- 2π production is comparable to 1π
- η, K productions (ν case: background of proton decay exp.)

GOAL : Develop νN-interaction model in resonance region

Problems in previous models

- Channel-couplings required by unitarity is missing
- Important 2 π production model is missing

Our strategy to overcome the problems…

We develop a dynamical coupled-channels model

★ Dynamical coupled-channels (DCC) model for γN, πN ➔ πN, $\pi\pi N$, ηN, $K\Lambda$, $K\Sigma$

★ Extension to νN ➔ $l^- X$ ($X = \pi N$, $\pi\pi N$, ηN, $K\Lambda$, $K\Sigma$)

Dynamical Coupled-Channels model for meson productions

DCC (Dynamical Coupled-Channel) model

Matsuyama et al., Phys. Rep. **439**, 193 (2007)

Kamano et al., PRC 88, 035209 (2013)

Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_c V_{ac} G_c T_{cb}$$

$$\{a, b, c\} = \pi N,\ \eta N,\ \pi\pi N\ , \pi\Delta, \sigma N, \rho N\ ,\ K\Lambda,\ K\Sigma$$

By solving the LS equation, coupled-channel unitarity is fully taken into account

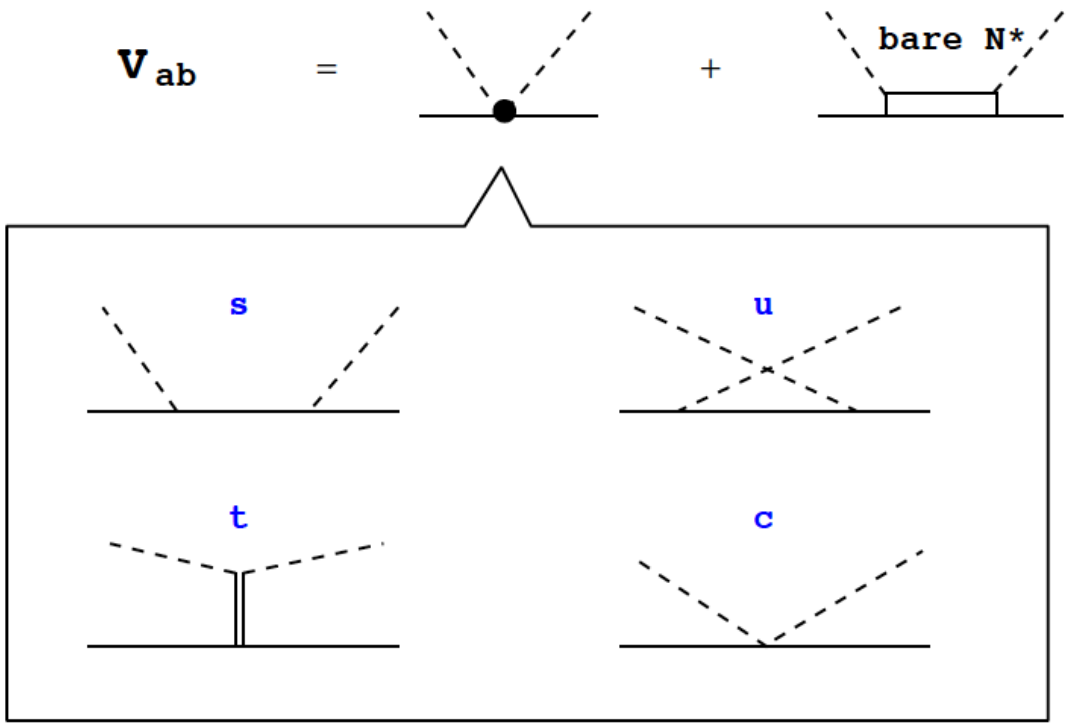
DCC (Dynamical Coupled-Channel) model

Matsuyama et al., Phys. Rep. **439**, 193 (2007)

Kamano et al., PRC 88, 035209 (2013)

Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_c V_{ac} G_c T_{cb}$$



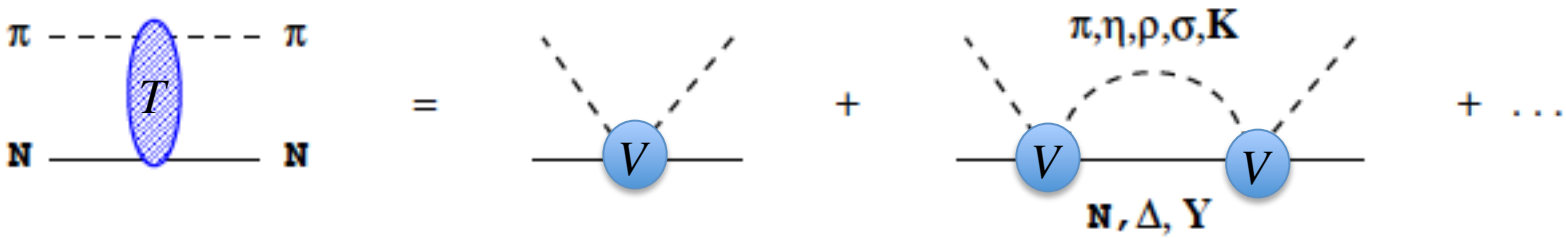
DCC (Dynamical Coupled-Channel) model

Matsuyama et al., Phys. Rep. **439**, 193 (2007)

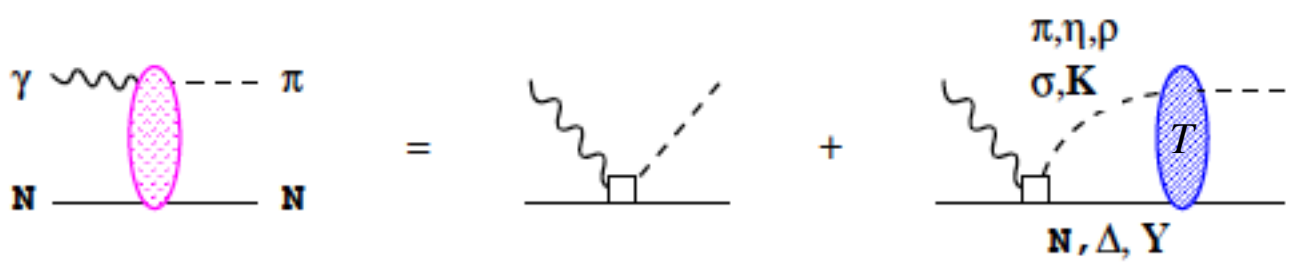
Kamano et al., PRC 88, 035209 (2013)

Coupled-channel Lippmann-Schwinger equation for meson-baryon scattering

$$T_{ab} = V_{ab} + \sum_c V_{ac}\, G_c\, T_{cb}$$



In addition, γN, $W^{\pm}N$, ZN channels are included perturbatively



Relation between neutrino and electron (photon) interactions

Charged-current (CC) interaction (e.g. $\nu_\mu + n \rightarrow \mu^- + p$)

$$L^{cc} = \frac{G_F V_{ud}}{\sqrt{2}} [J^{cc}_\lambda \ell^\lambda_{cc} + h.c.]$$

$$J^{cc}_\lambda = V_\lambda - A_\lambda$$

$$\ell^\lambda_{cc} = \bar{\psi}_\mu \gamma^\lambda (1-\gamma_5) \psi_\nu$$

Electromagnetic interaction (e.g. $\gamma^{(*)} + p \rightarrow p$)

$$L^{em} = e\, J^{em}_\lambda A^\lambda_{em}$$

$$J^{em}_\lambda = V_\lambda + V^{IS}_\lambda$$

V and V^{IS} in J^{em} can be separately determined by analyzing photon (Q^2=0) and electron reaction ($Q^2 \neq 0$) data on both proton and neutron targets, because:

$$< p | V_\lambda | p > = - < n | V_\lambda | n >$$

$$< p | V^{IS}_\lambda | p > = < n | V^{IS}_\lambda | n >$$

Matrix element for the weak vector current is obtained from analyzing electromagnetic processes

$$< p | V_\lambda | n > = 2 < p | V_\lambda | p >$$

DCC model for axial current

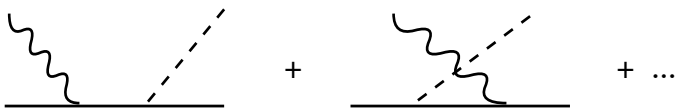
Because neutrino reaction data are scarce, axial current cannot be determined phenomenologically
→ Chiral symmetry and PCAC (partially conserved axial current) are guiding principle

PCAC relation $< X' | q \cdot A | X > \sim i f_\pi < X' | T | \pi X >$

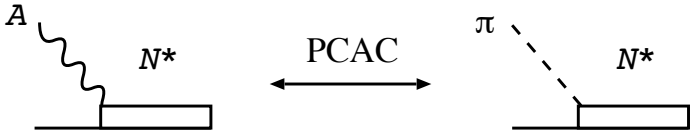
Q^2=0

non-resonant mechanisms

$$\partial_\mu \pi \rightarrow f_\pi A_\mu^{external}$$



resonant mechanisms



Interference among resonances and background can be uniquely fixed within DCC model

DCC model for axial current

$Q^2 \neq 0$ $F_A(Q^2)$: axial form factors

non-resonant mechanisms $$F_A(Q^2) = \left(\frac{1}{1 + Q^2 / M_A^2} \right)^2$$ $M_A = 1.02$ GeV

resonant mechanisms $$F_A(Q^2) = \left(\frac{1}{1 + Q^2 / M_A^2} \right)^2$$

More neutrino data are necessary to fix axial form factors for ANN^*

Neutrino cross sections will be predicted with this axial current

DCC analysis of $\gamma N, \pi N \rightarrow \pi N, \eta N, K\Lambda, K\Sigma$

and electron scattering data

DCC analysis of meson production data

Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Fully combined analysis of $\gamma N, \pi N$ ➔ $\pi N, \eta N, K\Lambda, K\Sigma$ data

$d\sigma / d\Omega$ and polarization observables (W ≤ 2.1 GeV)

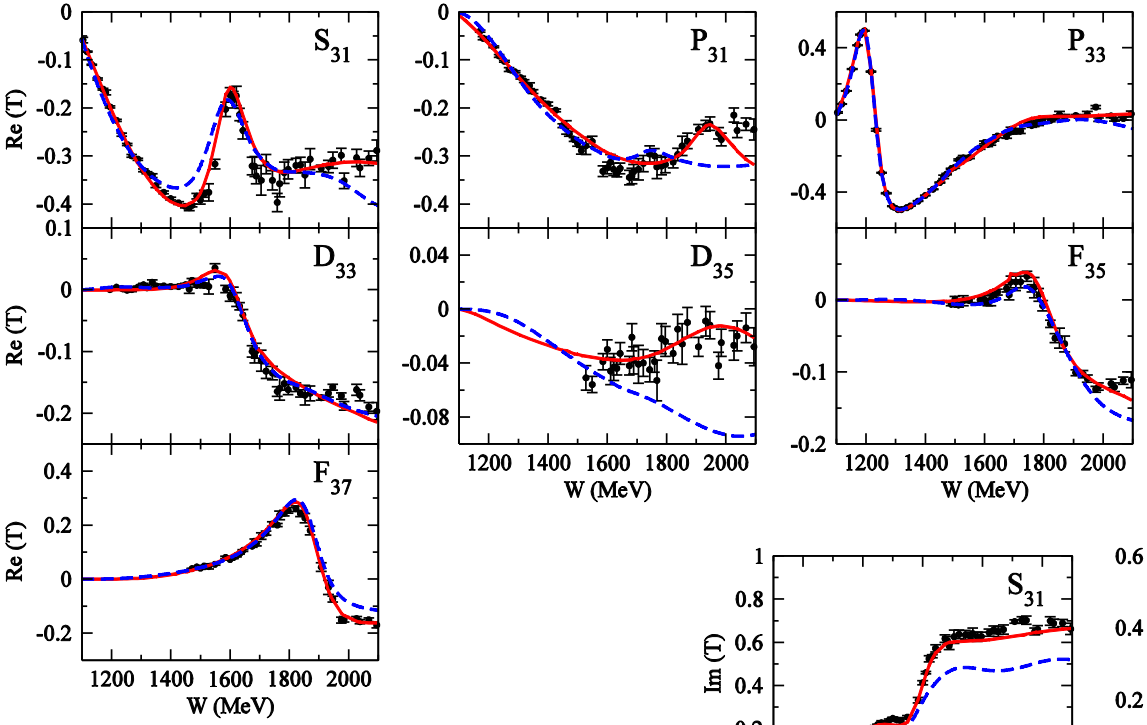
~ 23,000 data points are fitted

by adjusting parameters (N^* mass, N^* ➔ MB couplings, cutoffs)

Data for electron scattering on proton and neutron are analyzed by adjusting

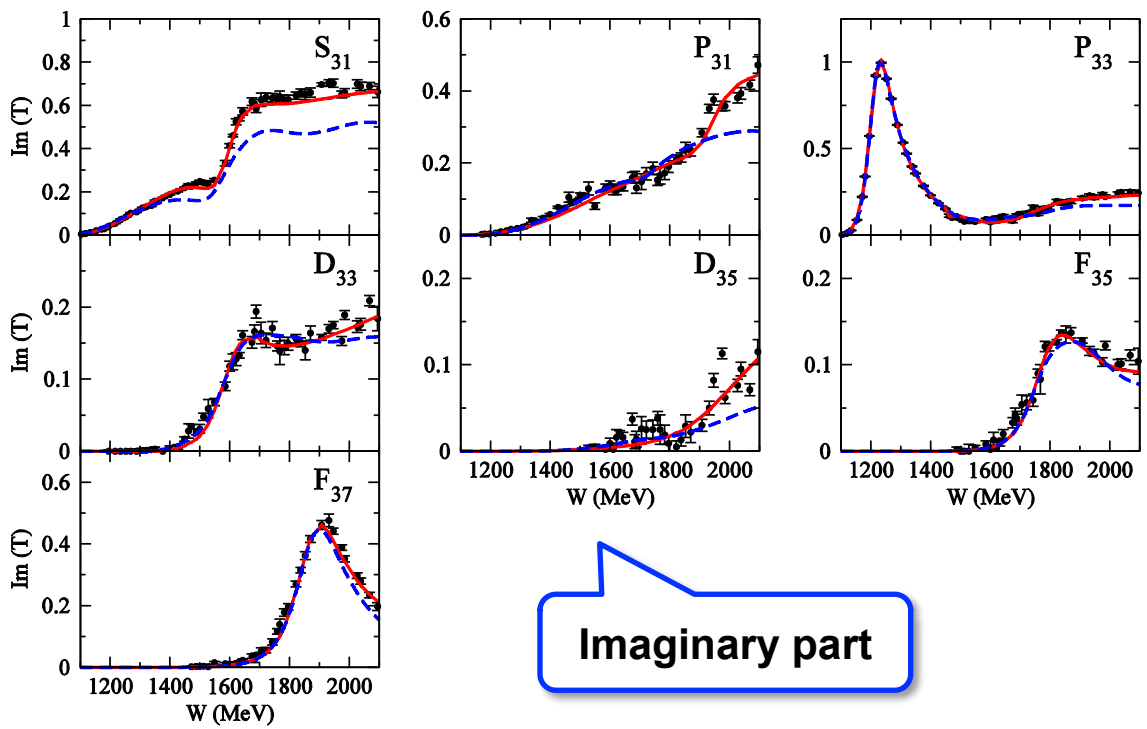
$\gamma^* N$ ➔ N^* coupling strength at different Q^2 values ($Q^2 \leq 3\ (\mathrm{GeV}/c)^2$)

Partial wave amplitudes of π N scattering



Real part

$$I = \frac{3}{2}$$



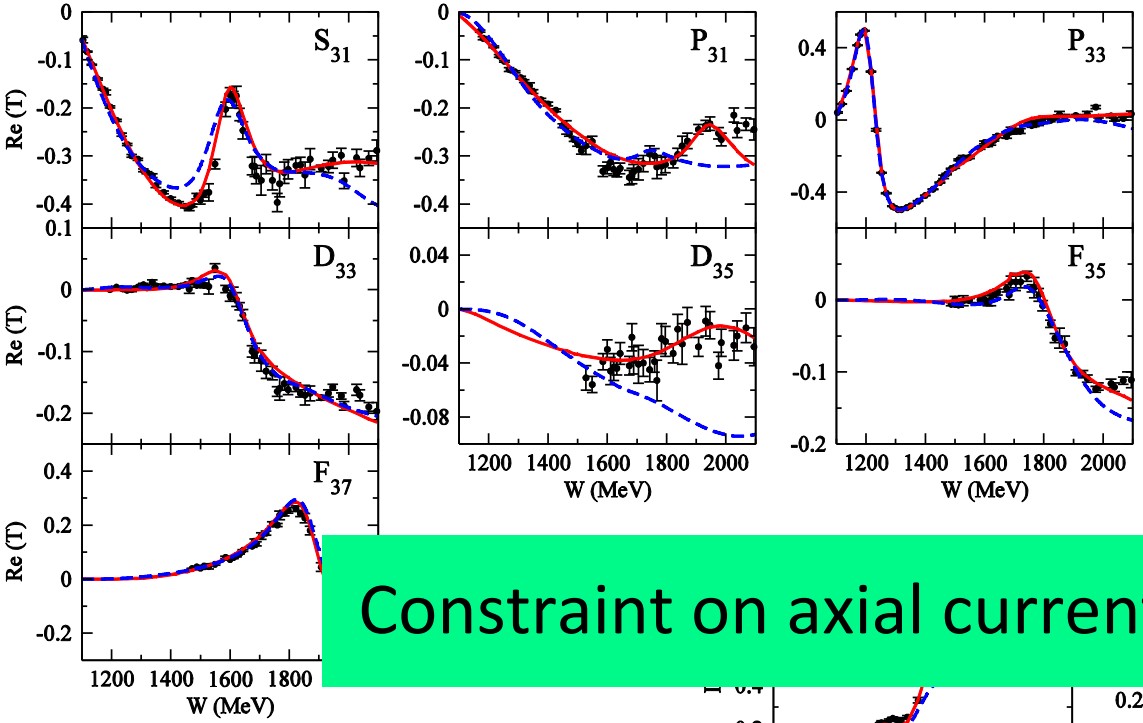
Imaginary part

Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Previous model
(fitted to πN → πN data only)
[PRC76 065201 (2007)]

Data: SAID πN amplitude

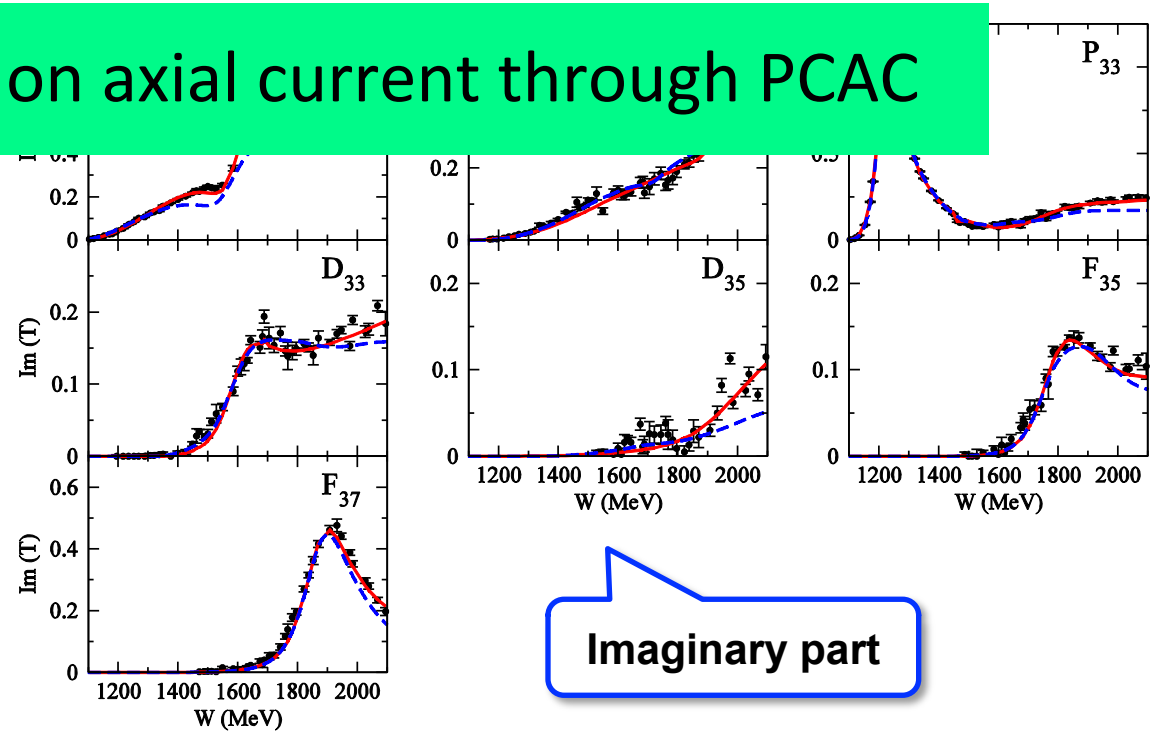
Partial wave amplitudes of π N scattering



Real part

$$I = \frac{3}{2}$$

Constraint on axial current through PCAC



Kamano, Nakamura, Lee, Sato, PRC 88 (2013)

Previous model
(fitted to πN → πN data only)
[PRC76 065201 (2007)]

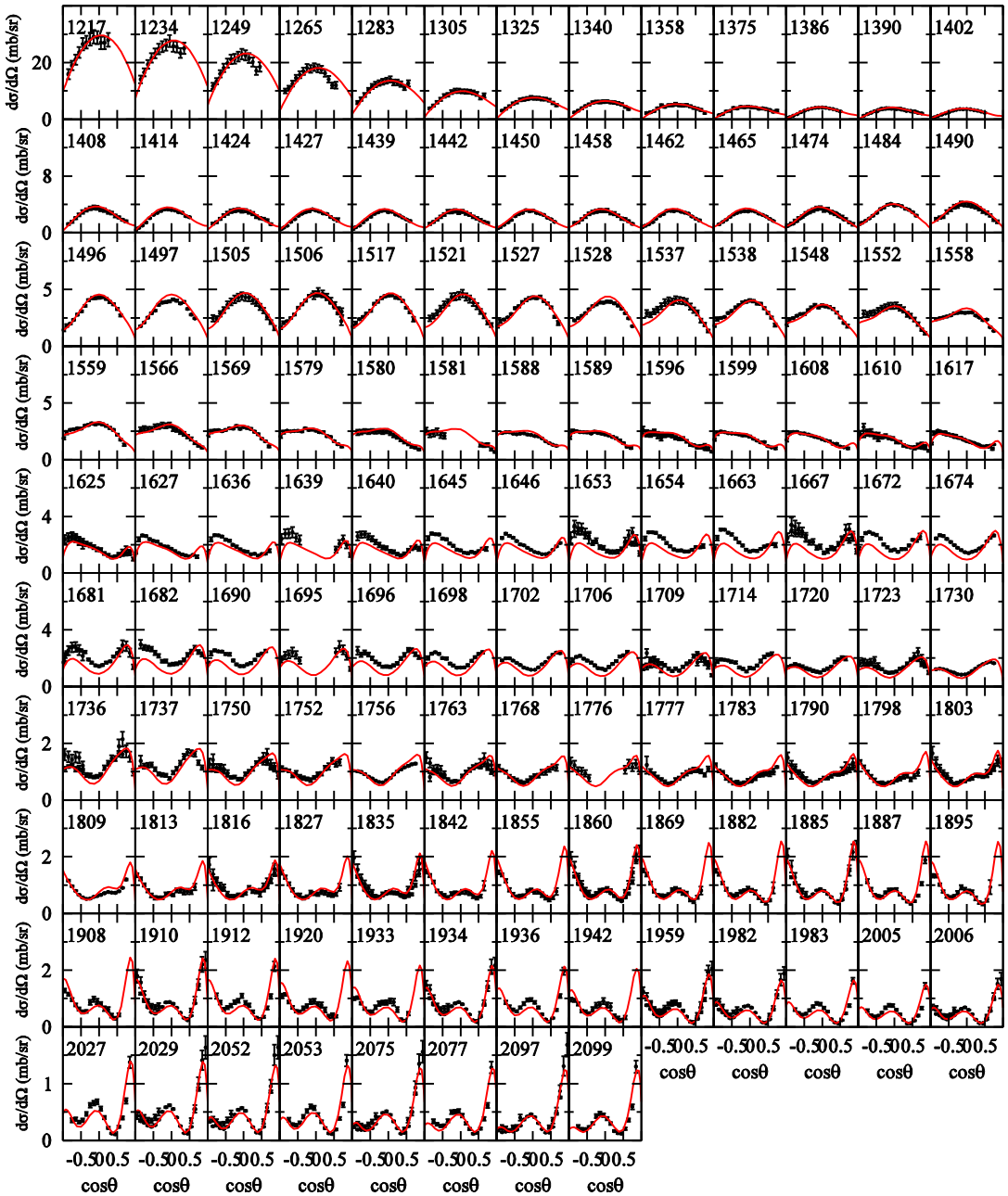
Data: SAID πN amplitude

Imaginary part

$\gamma p \rightarrow \pi^0 p$

$d\sigma/d\Omega$ for W < 2.1 GeV

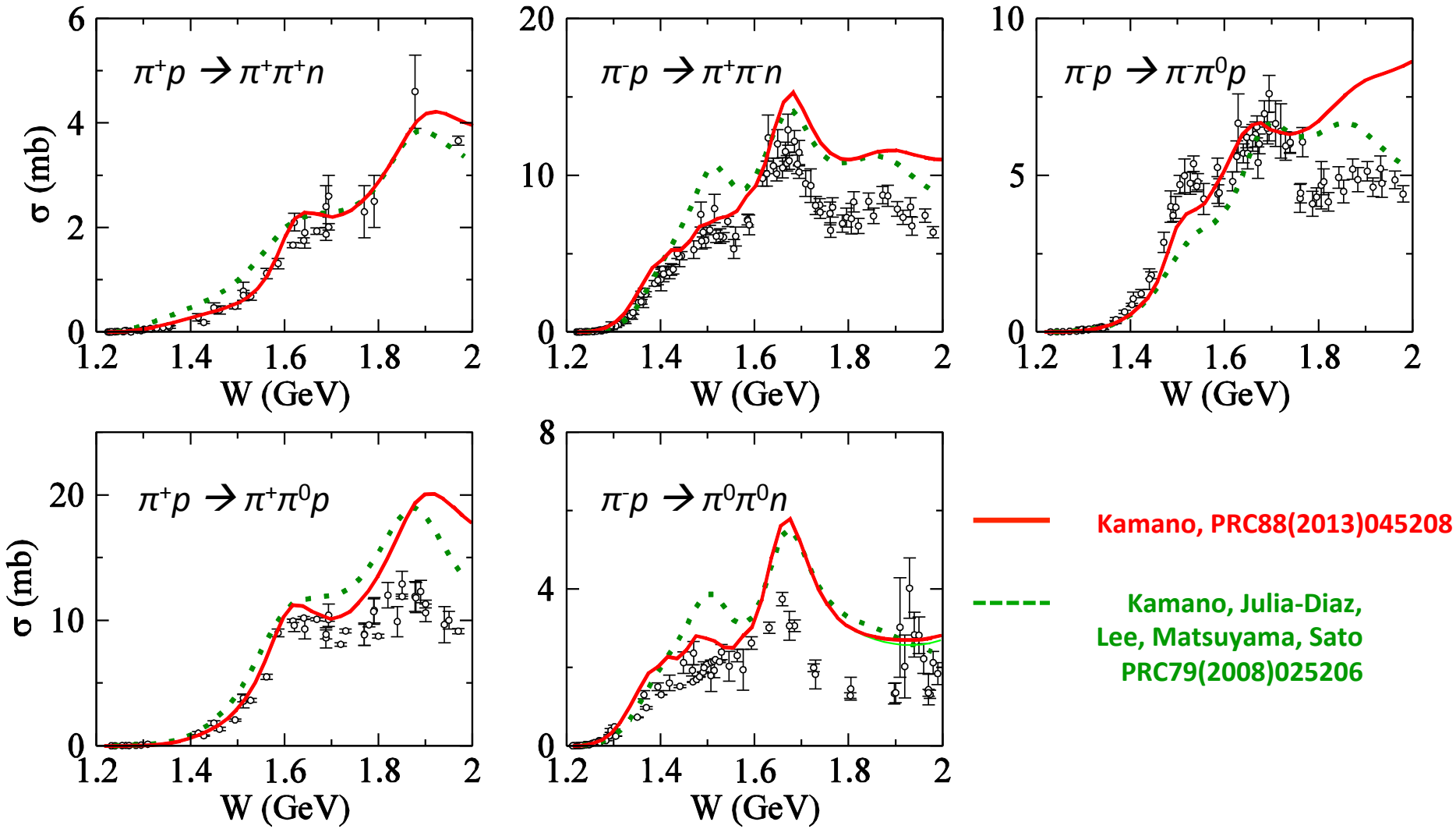
Kamano, Nakamura, Lee, Sato, PRC 88 (2013)



Vector current (Q^2=0) for 1π

Production is well-tested by data

Predicted πN → ππN total cross sections with our DCC model

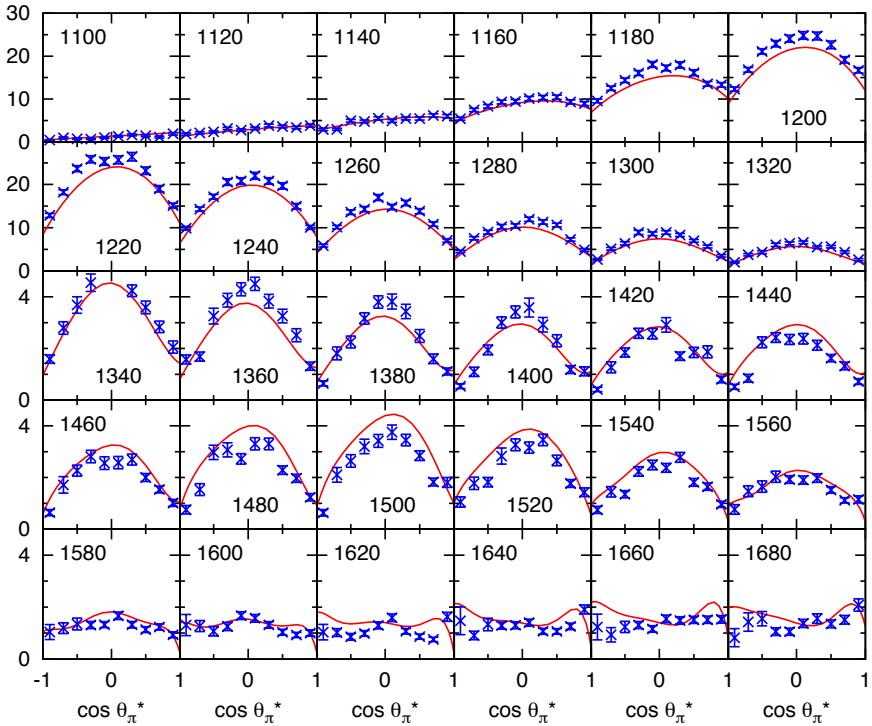


Single π production in electron-proton scattering

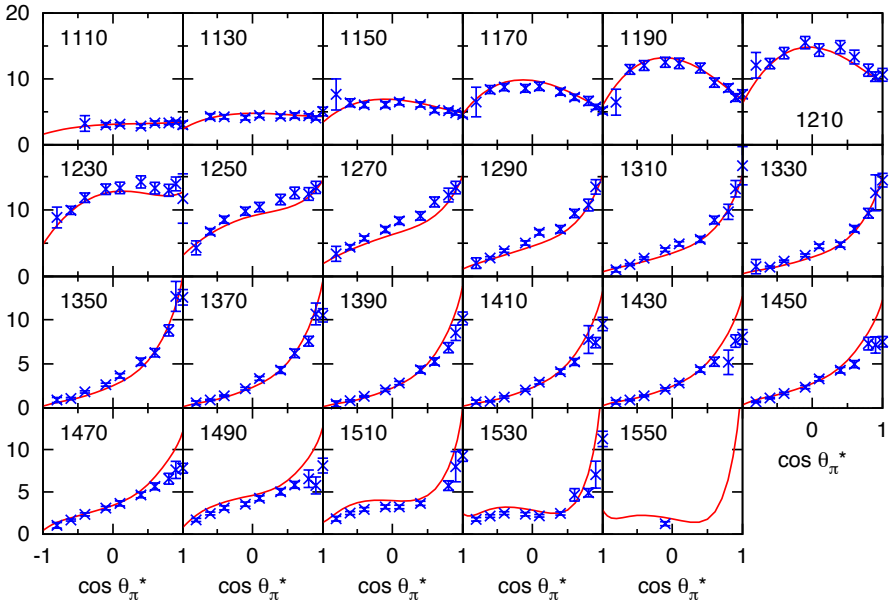
Purpose : Determine Q^2–dependence of vector coupling of p-N^* : $VpN^*(Q^2)$

$\sigma_T + \varepsilon\, \sigma_L$ for Q^2=0.40 (GeV/c)^2 and W=1.1 – 1.68 GeV

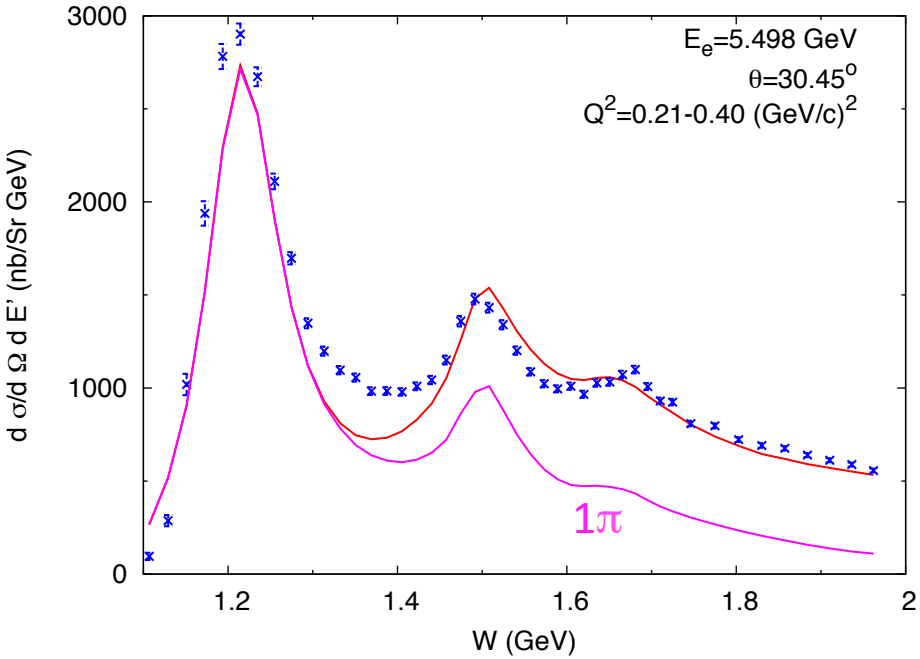
$p(e,e'\pi^0)p$

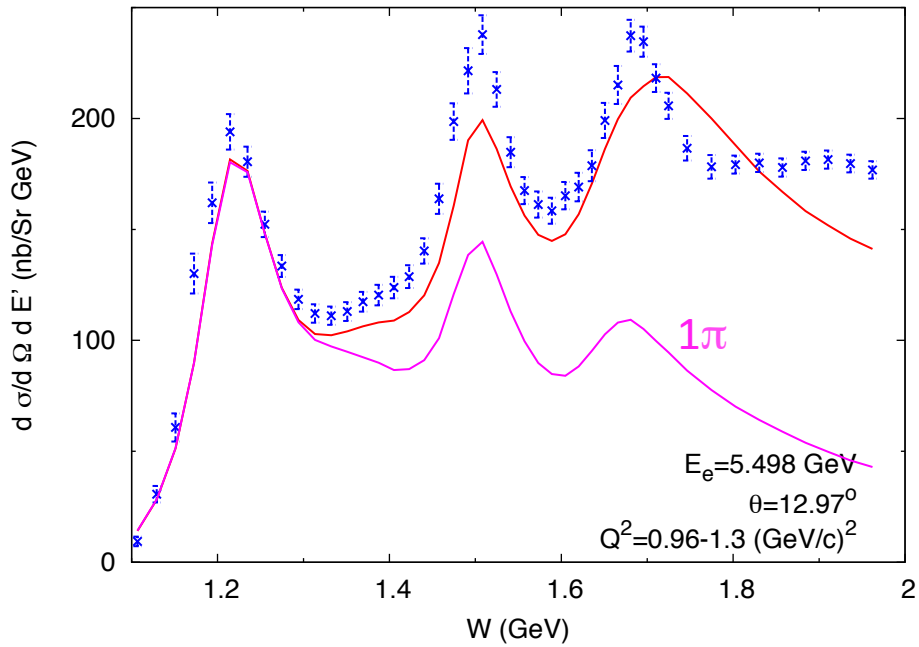


$p(e,e'\pi^+)n$



Inclusive electron-proton scattering





Data: JLab E00-002 (preliminary)

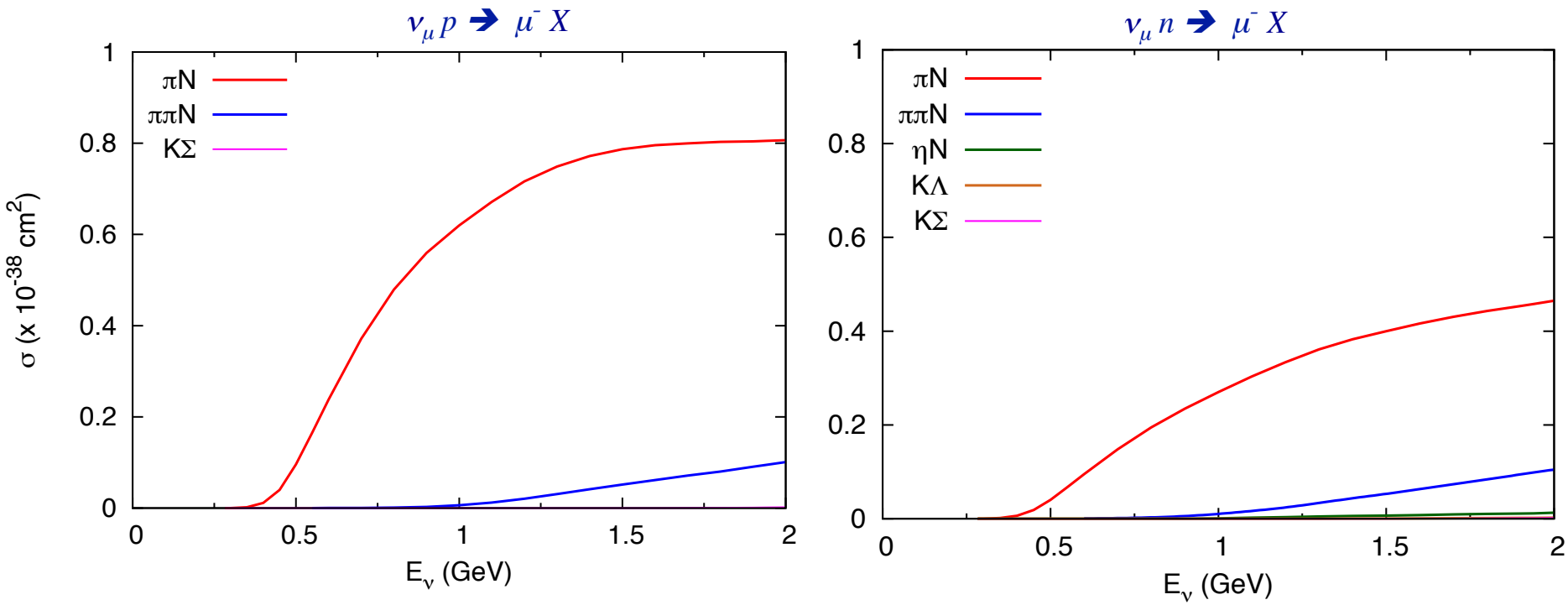
- Reasonable fit to data for application to neutrino interactions
- Important 2π contributions for high W region

Similar analysis of electron-neutron scattering data has also been done

DCC vector currents has been tested by data for whole kinematical region relevant to neutrino interactions of $E_\nu \leq 2$ GeV

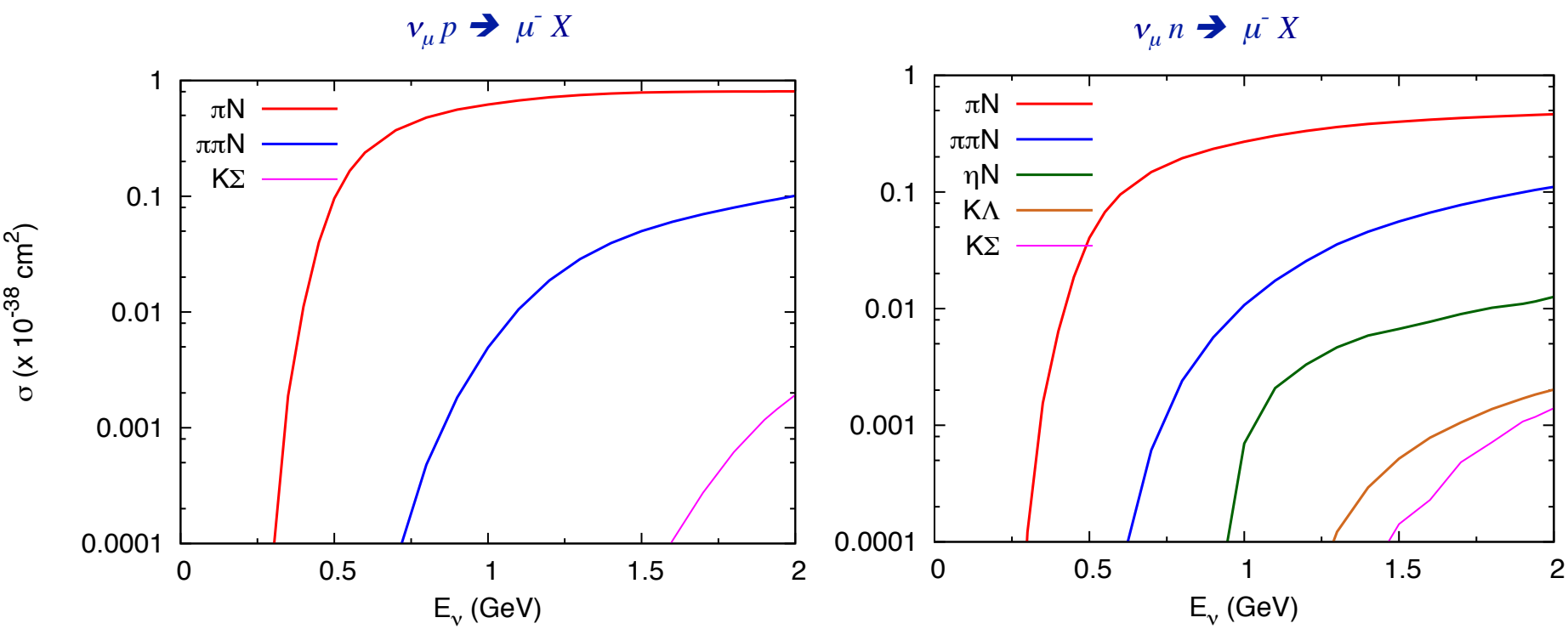
Neutrino Results

Cross section for $\nu_\mu N \rightarrow \mu^- X$



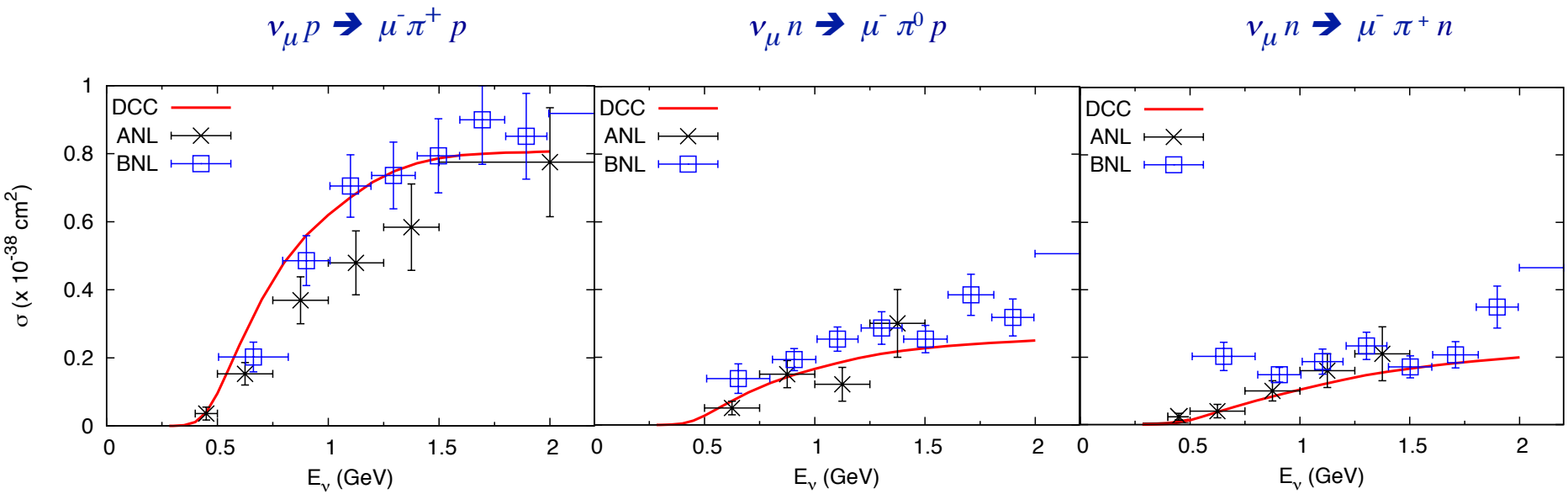
- πN & $\pi\pi N$ are main channels in few-GeV region
- DCC model gives predictions for all final states
- ηN, KY cross sections are $10^{-1} - 10^{-2}$ smaller

Cross section for $\nu_\mu N \rightarrow \mu^- X$



- πN & $\pi\pi N$ are main channels in few-GeV region
- DCC model gives predictions for all final states
- ηN, KY cross sections are $10^{-1} - 10^{-2}$ smaller

Comparison with single pion data



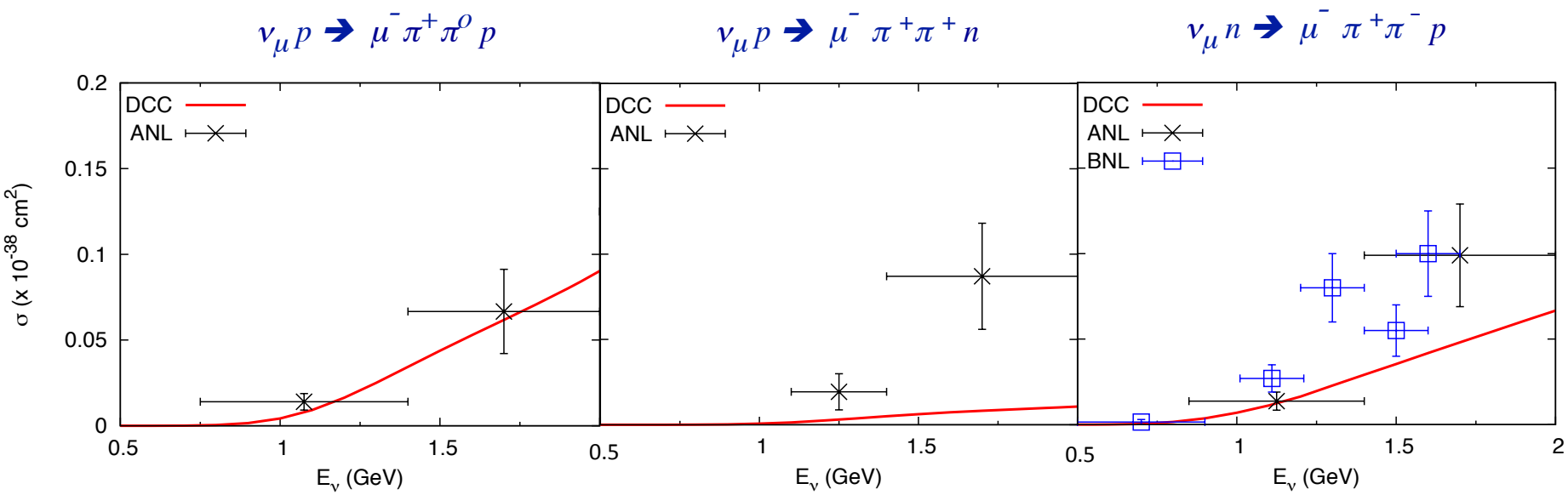
DCC model prediction is consistent with data

ANL Data : PRD **19**, 2521 (1979)
BNL Data : PRD **34**, 2554 (1986)

- DCC model has flexibility to fit data ($ANN^*(Q^2)$)
- Data should be analyzed with nuclear effects

  (Wu et al. , arXiv:1412:2415)

Comparison with double pion data

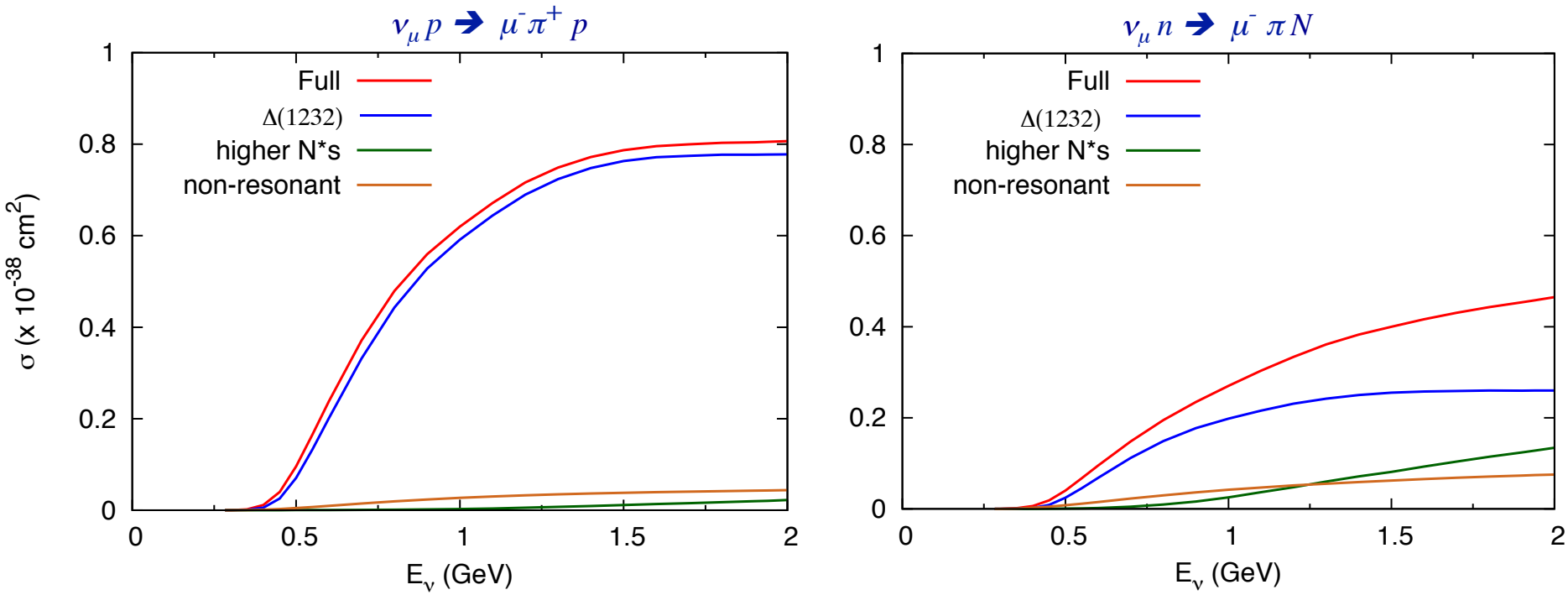


Fairly good DCC predication

ANL Data : PRD **28**, 2714 (1983)
BNL Data : PRD **34**, 2554 (1986)

First dynamical model for 2 π production in resonance region

Mechanisms for $\nu_\mu N \rightarrow \mu^- \pi N$

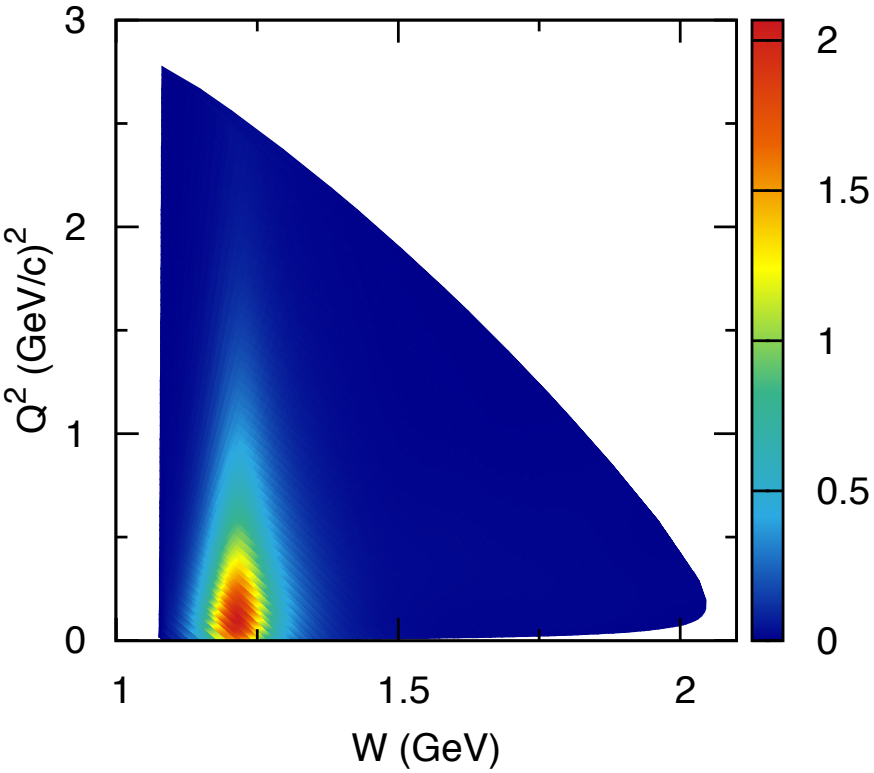


- $\Delta(1232)$ dominates for $\nu_\mu p \rightarrow \mu^- \pi^+ p$ $(I=3/2)$ for $E_\nu \leq 2$ GeV
- Non-resonant mechanisms contribute significantly
- Higher N^*s becomes important towards $E_\nu \approx 2$ GeV for $\nu_\mu n \rightarrow \mu^- \pi N$

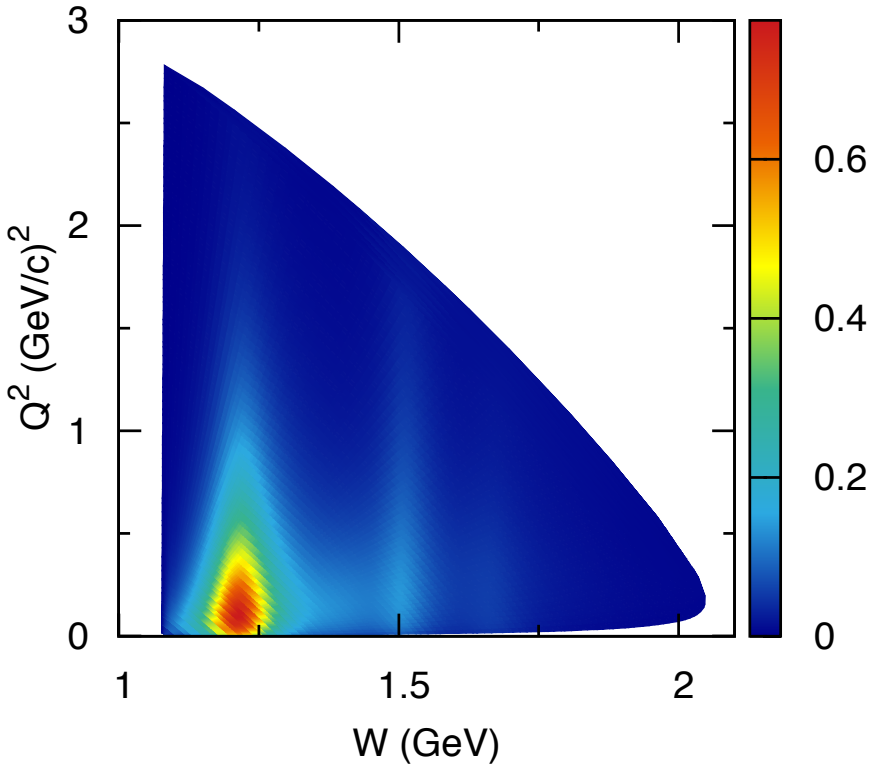
$d\sigma / dW\, dQ^2$ ($\times 10^{-38}$ cm^2 / GeV^2)

E_ν = 2 GeV

$\nu_\mu p \rightarrow \mu^- \pi^+ p$



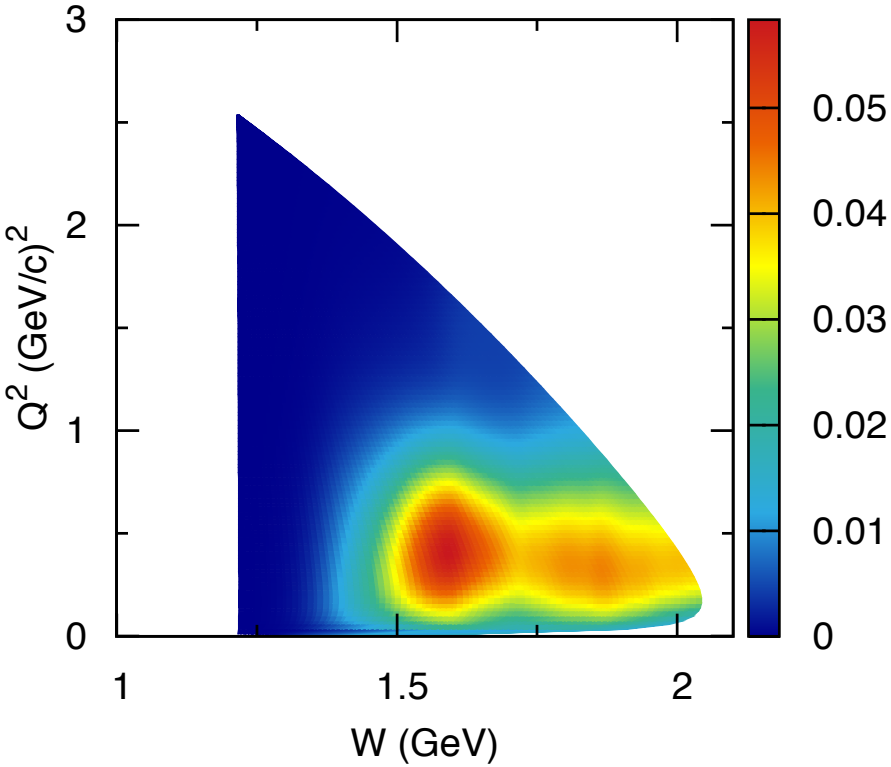
$\nu_\mu n \rightarrow \mu^- \pi N$



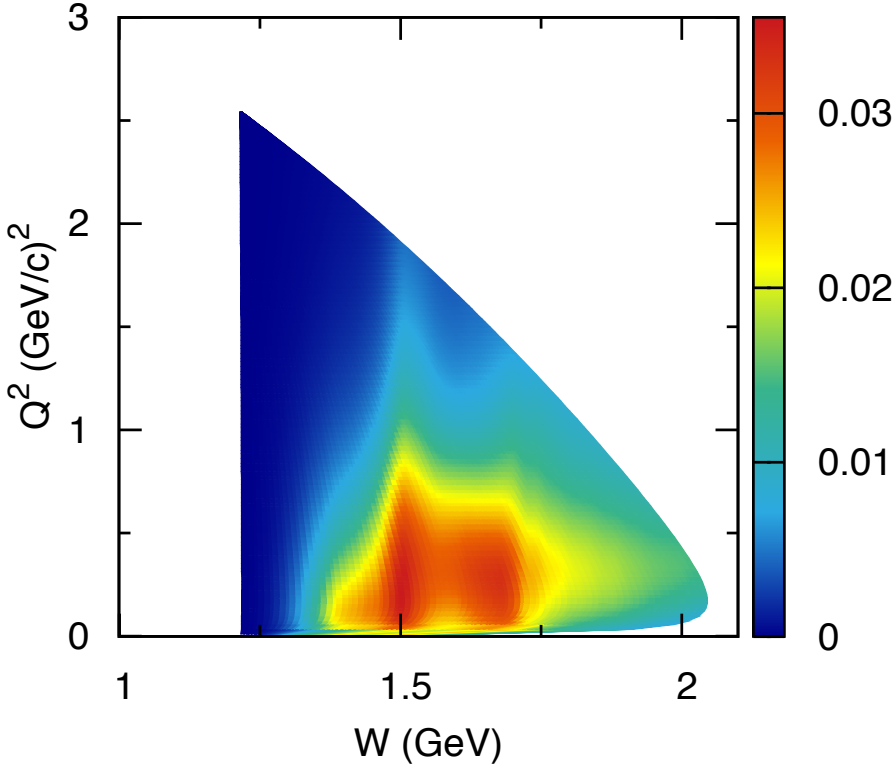
$d\sigma / dW\, dQ^2$ ($\times 10^{-38}$ cm^2 / GeV^2)

E_ν = 2 GeV

$\nu_\mu p \rightarrow \mu^- \pi^+ \pi^0 p$



$\nu_\mu n \rightarrow \mu^- \pi^+ \pi^- p$



Conclusion

Development of DCC model for νN interaction in resonance region

Start with DCC model for γN, πN ➔ πN, $\pi\pi N$, ηN, $K\Lambda$, $K\Sigma$

➔ extension of vector current to $Q^2 \neq 0$ region, isospin separation

through analysis of e^-- p & e^--'n' data for $W \leq 2$ GeV , $Q^2 \leq 3$ (GeV/c)^2

➔ Development of axial current for νN interaction; PCAC is maintained

Conclusion

- πN & $\pi\pi N$ are main channels in few-GeV region
- DCC model prediction is consistent with BNL data
- Δ, N^*s, non-resonant are all important in few-GeV region (for $\nu_\mu n$ ➔ $\mu^- X$)

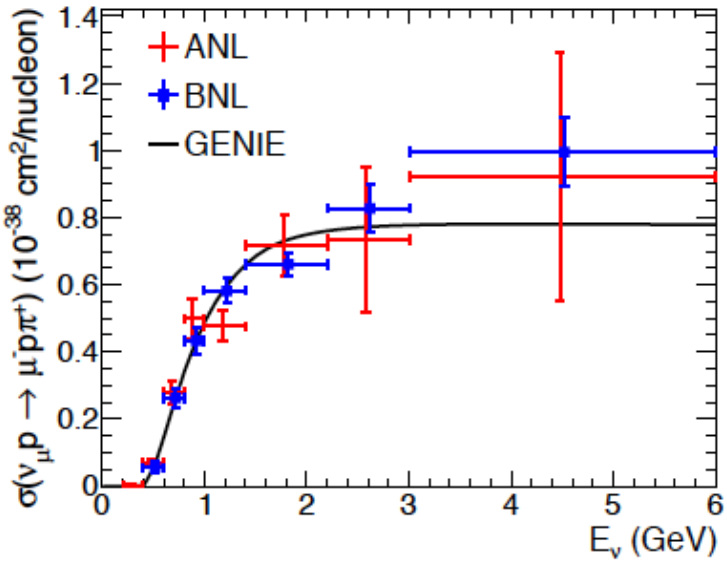
➔ essential to understand interference pattern among them

➔ DCC model can do this; consistency between π interaction and axial current

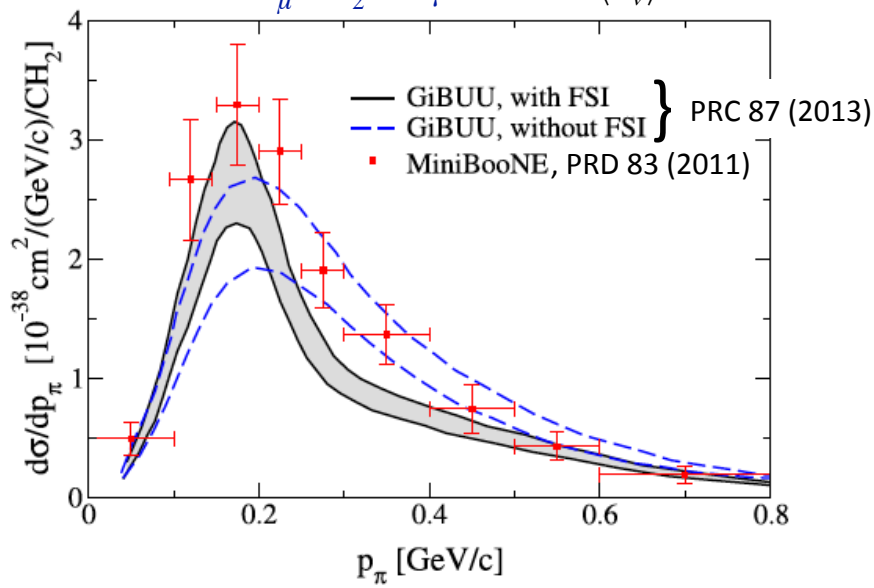
BACKUP

Neutrino interaction data in resonance region

$\nu_\mu p \rightarrow \mu^- \pi^+ p$



$\nu_\mu CH_2 \rightarrow \mu^- \pi^0 X$ $\langle E_\nu \rangle = 0.8$ GeV



- Data to fix nucleon axial current ($g_{AN\Delta}$)
- Discrepancy between BNL & ANL data
- Recent reanalysis of original data

  → discrepancy resolved (!?)

  PRD 90, 112017 (2014)

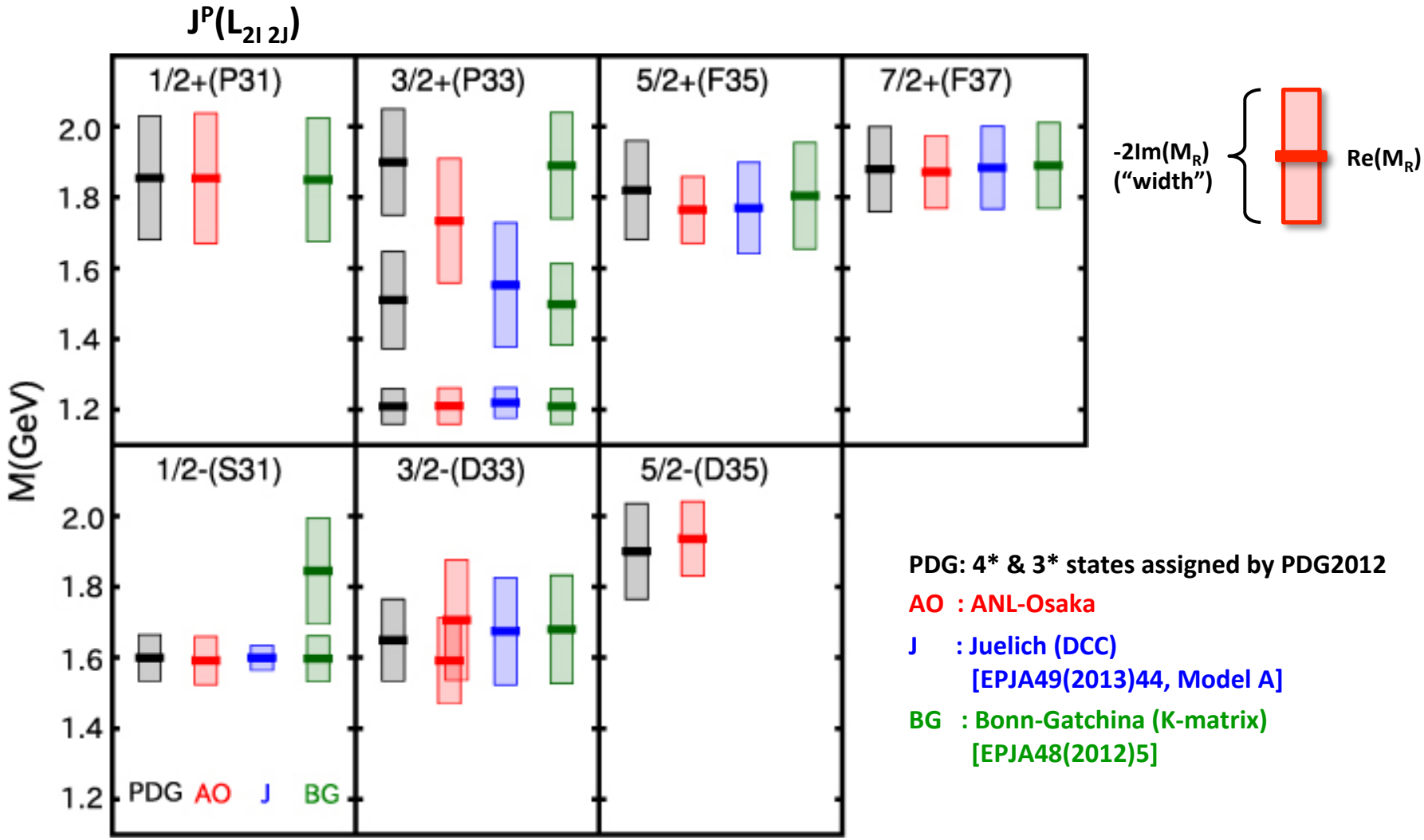
- Final state interaction (FSI) changes charge, momentum, number of π
- Cross section shape is worse described with FSI
- MINERνA data (arXiv:1406.6415) favor FSI

  $\langle E_\nu \rangle = 4.0$ GeV

More data are coming → better understanding of neutrino-nucleus interaction

“Δ” resonances (I=3/2)

Kamano, Nakamura, Lee, Sato, PRC 88 (2013)



PDG: 4* & 3* states assigned by PDG2012

AO : ANL-Osaka

J : Juelich (DCC) [EPJA49(2013)44, Model A]

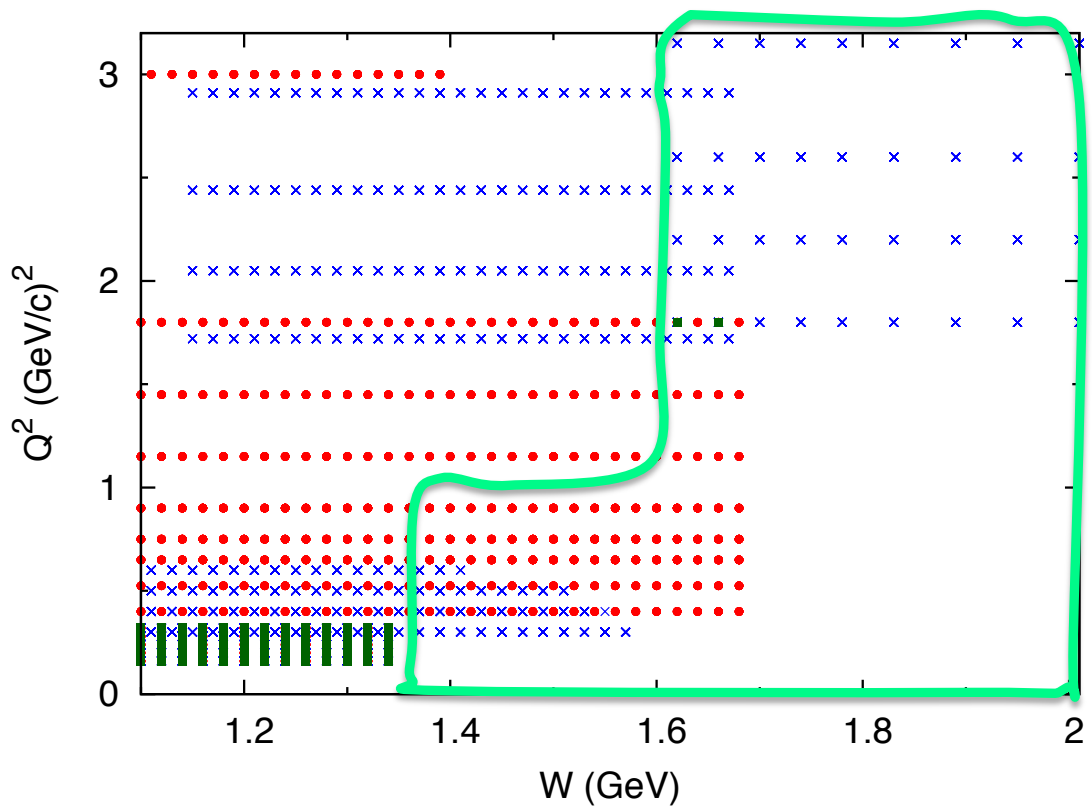
BG : Bonn-Gatchina (K-matrix) [EPJA48(2012)5]

Analysis of electron-proton scattering data

Purpose : Determine Q^2–dependence of vector coupling of p-N^* : $VpN^*(Q^2)$

Data : * 1π electroproduction

* Empirical inclusive inelastic structure functions σ_T, σ_L ← Christy et al, PRC 81 (2010)



Database

- $p(e,e'\pi^0)p$
- $p(e,e'\pi^+)n$
- both

region where inclusive σ_T & σ_L are fitted

Analysis of electron-'neutron' scattering data

Purpose : Vector coupling of neutron-N^* and its Q^2–dependence : $VnN^*(Q^2)$ $(I=1/2)$

$I=3/2$ part has been fixed by proton target data

Data : * 1π photoproduction (Q^2=0)

* Empirical inclusive inelastic structure functions σ_T, σ_L ($Q^2 \neq 0$)

← Christy and Bosted, PRC 77 (2010), 81 (2010)

Done

DCC vector currents has been tested by data for whole kinematical region relevant to neutrino interactions of $E_\nu \leq 2$ GeV

Formalism

Cross section for $\nu N \rightarrow l\, X$ ($X = \pi N, \pi\pi N, \eta N, K\Lambda, K\Sigma$)

$\theta \rightarrow 0$

$$\frac{d\sigma}{dE_\ell d\Omega_\ell} = \frac{G_F^2}{2\pi^2} E_\ell^2 \left(\cancel{2W_1 \sin^2\frac{\theta}{2}} + W_2 \cos^2\frac{\theta}{2} \pm \cancel{W_3 \frac{E_\nu + E_\ell}{m_N} \sin^2\frac{\theta}{2}} \right)$$

$Q^2 \rightarrow 0$

$$W_2 = \frac{Q^2}{\vec{q}^2} \sum \left[\cancel{\frac{1}{2}\left(|\langle J^x \rangle|^2 + |\langle J^y \rangle|^2\right)} + \frac{Q^2}{\vec{q}_c^2} \left| \left\langle \cancel{J^0} + \frac{\omega_c}{Q^2} q \cdot J \right\rangle \right|^2 \right]$$

CVC & PCAC

$$\langle q \cdot J \rangle = \langle q \cdot V \rangle - \langle q \cdot A \rangle = i f_\pi m_\pi^2 \langle \hat{\pi} \rangle$$

LSZ & smoothness

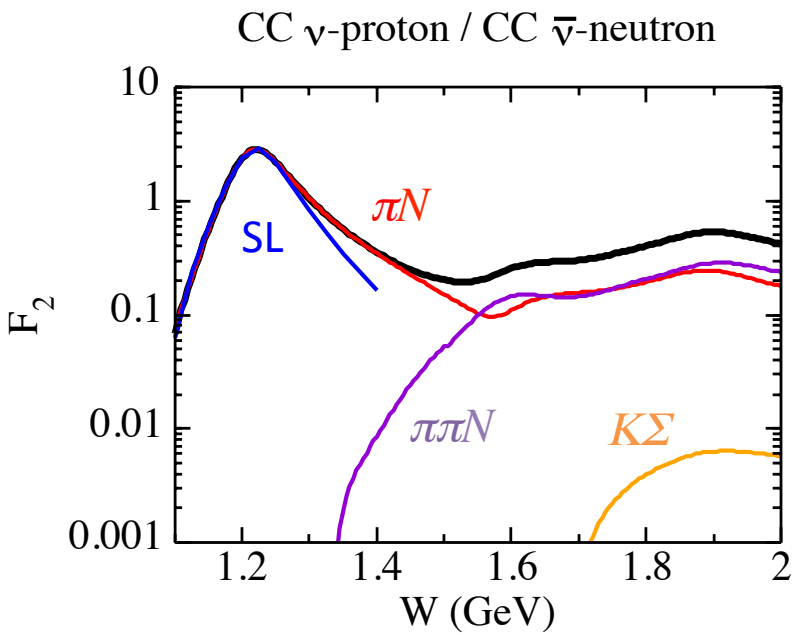
$$\langle X | \hat{\pi} | N \rangle = \frac{\sqrt{2\omega_c}}{m_\pi^2} \mathcal{T}_{\pi N \rightarrow X}(0) \sim \frac{\sqrt{2\omega_c}}{m_\pi^2} \mathcal{T}_{\pi N \rightarrow X}(m_\pi^2)$$

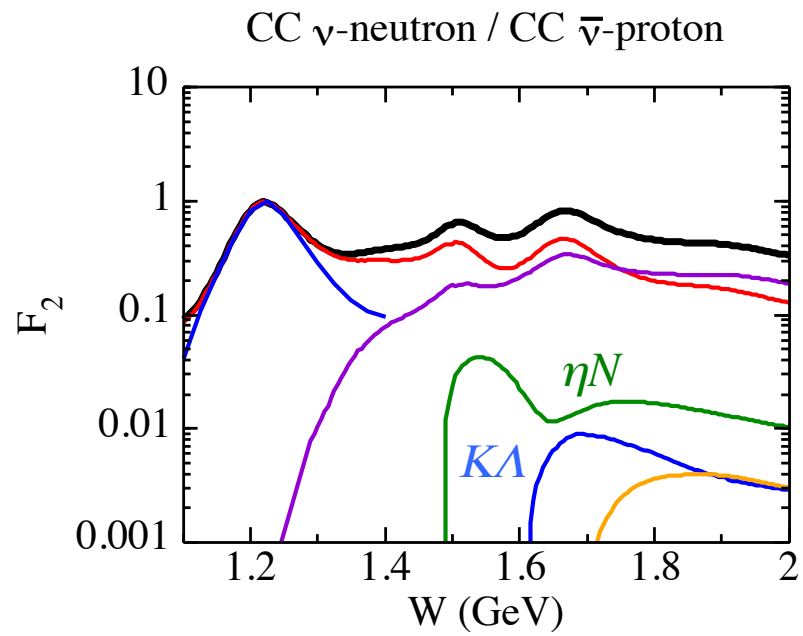
Finally

$$F_2 \equiv \omega W_2 = \frac{2 f_\pi^2}{\pi} \sigma_{\pi N \rightarrow X}$$

$\sigma_{\pi N \rightarrow X}$ is from our DCC model

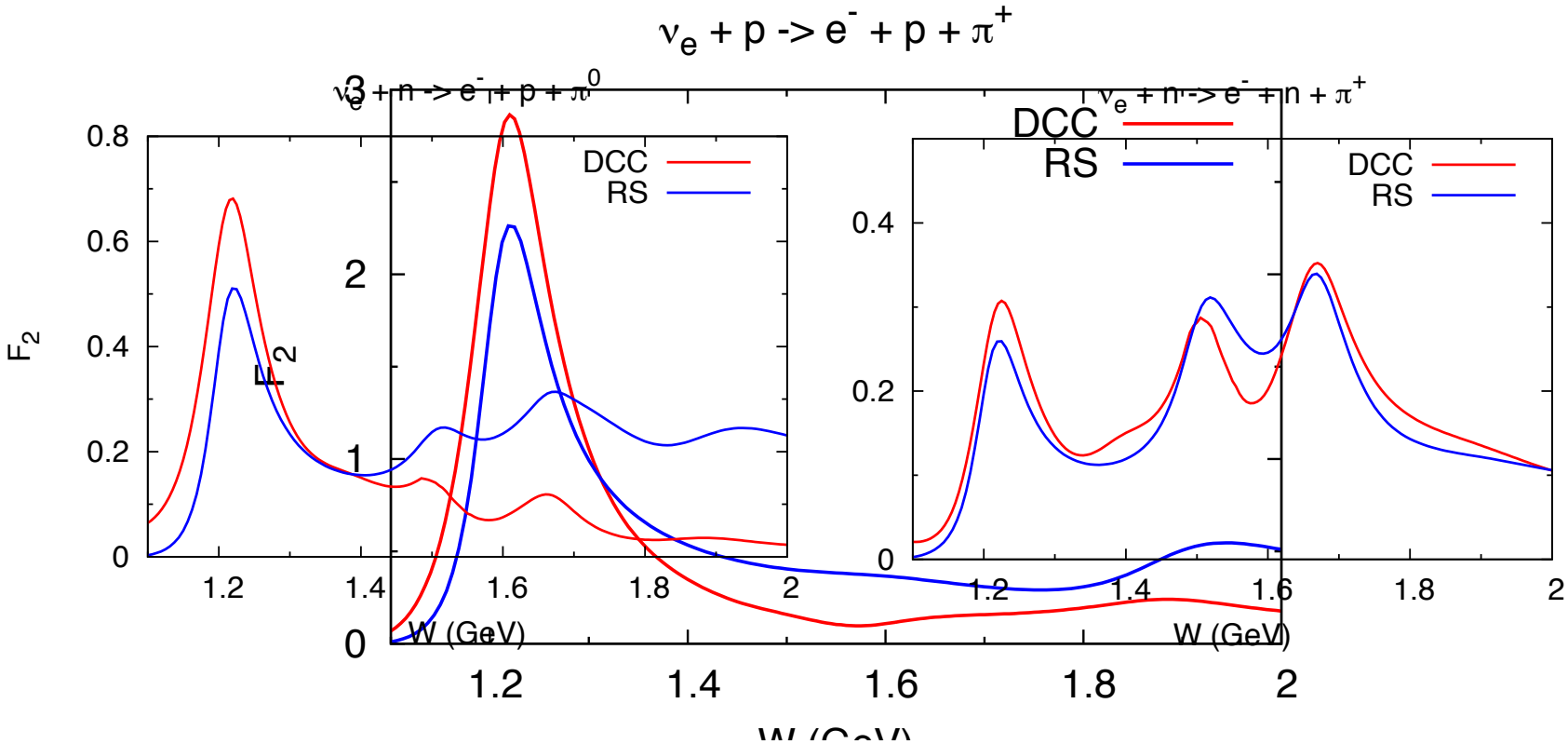
Results





- Prediction based on model well tested by data (first νN ➔ $\pi\pi N$)
- πN dominates for $W \leq 1.5$ GeV
- $\pi\pi N$ becomes comparable to πN for $W \geq 1.5$ GeV
- Smaller contribution from ηN and KY $O(10^{-1})$ - $O(10^{-2})$
- Agreement with SL (no PCAC) in Δ region

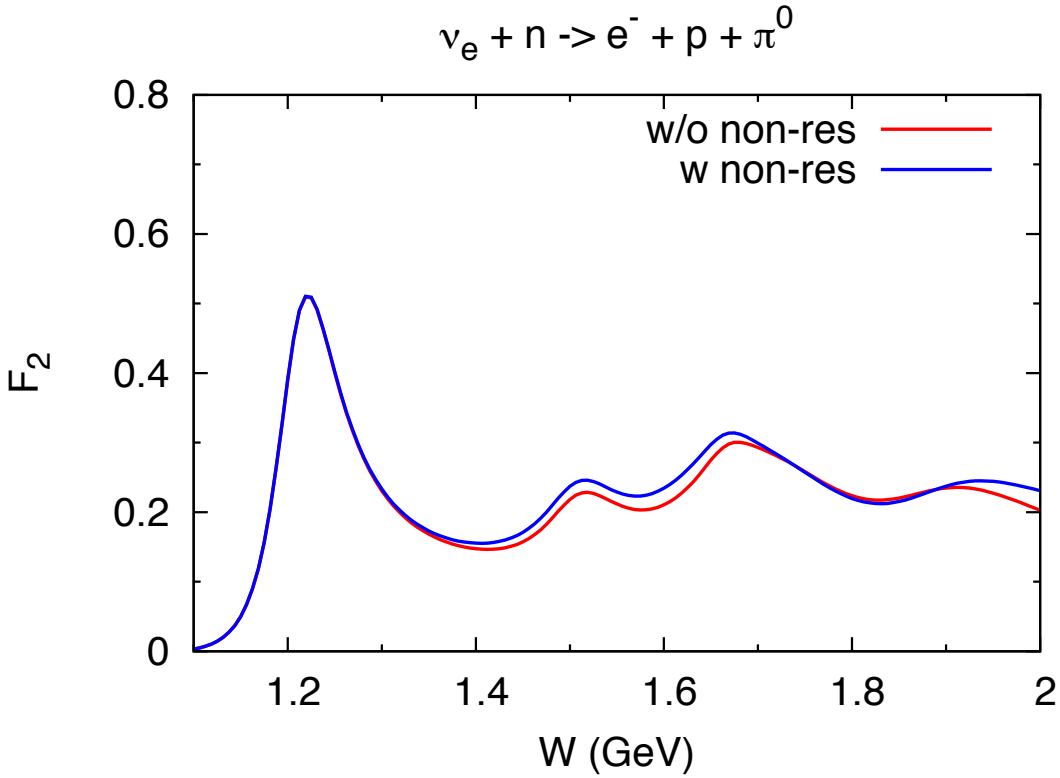
Comparison with Rein-Sehgal model



Comparison in whole kinematical region will be done

after axial current model is developed

F_2 from RS model

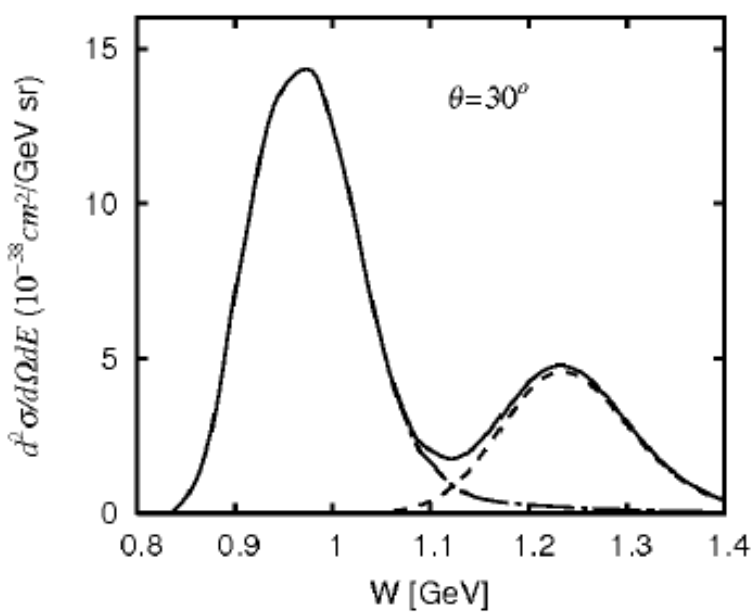


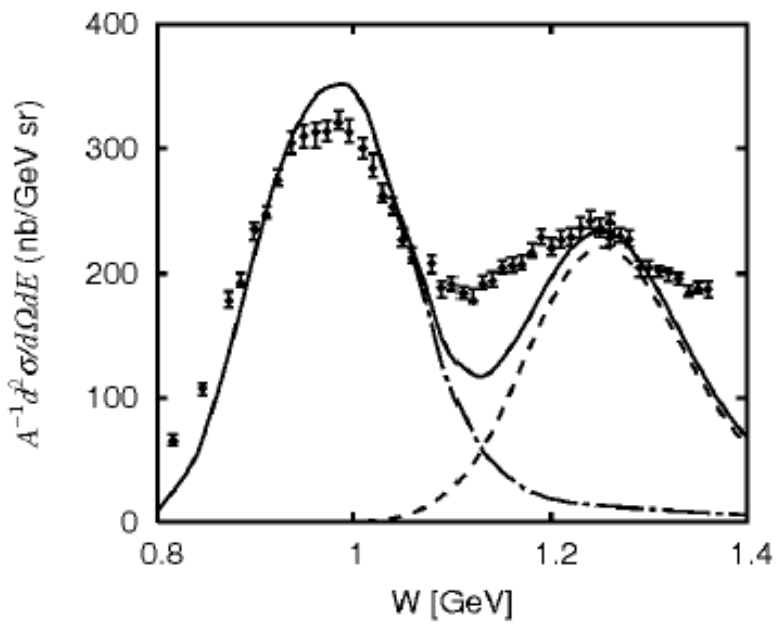
SL model applied to ν−nucleus scattering

1 π production

$\nu_e + {}^{12}\mathrm{C} \rightarrow e^- + X$ ($E_\nu = 1$ GeV)

$e^- + {}^{12}\mathrm{C} \rightarrow e^- + X$ ($E_e = 1.1$ GeV)



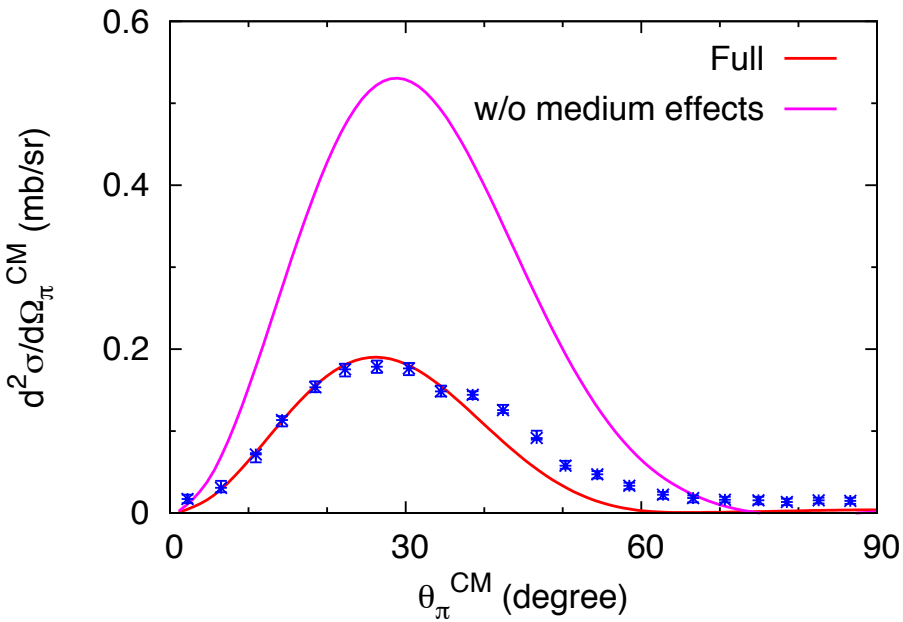


Szczerbinska et al. (2007)

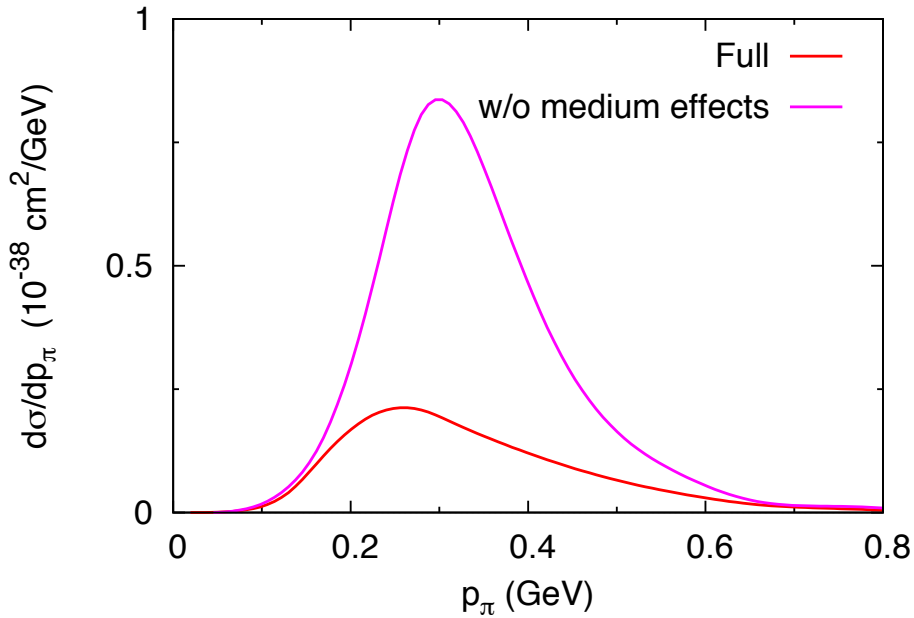
SL model applied to ν−nucleus scattering

coherent π production

$$\gamma + {}^{12}\mathrm{C} \rightarrow \pi^0 + {}^{12}\mathrm{C}$$



$$\nu_\mu + {}^{12}\mathrm{C} \rightarrow \mu^- + \pi^0 + {}^{12}\mathrm{C}$$



Nakamura et al. (2010)

DCC (Dynamical Coupled-Channel) model

Matsuyama et al., Phys. Rep. **439**, 193 (2007)

Coupled-channel Lippmann-Schwinger equation

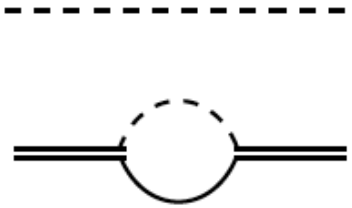
$$T_{ab} = V_{ab} + \sum_c V_{ac}\, G_c\, T_{cb}$$

$\mathbf{V_{ab}}$ =

+ **Z**

$G\,c$ =

for stable channels

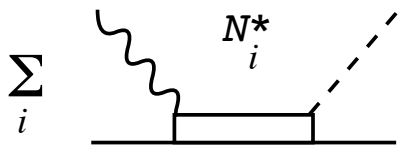


for unstable channels

Previous models for ν–induced 1π production in resonance region

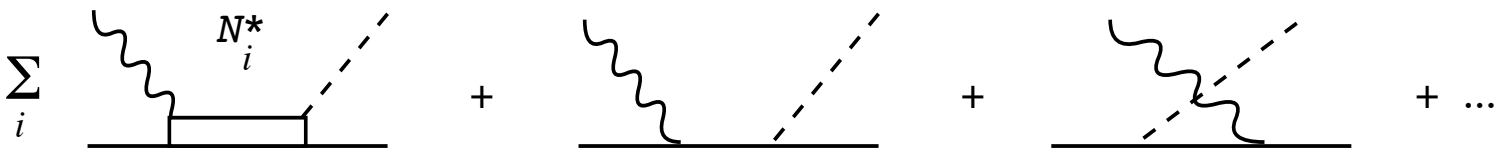
resonant only

Rein et al. (1981), (1987) ; Lalalulich et al. (2005), (2006)



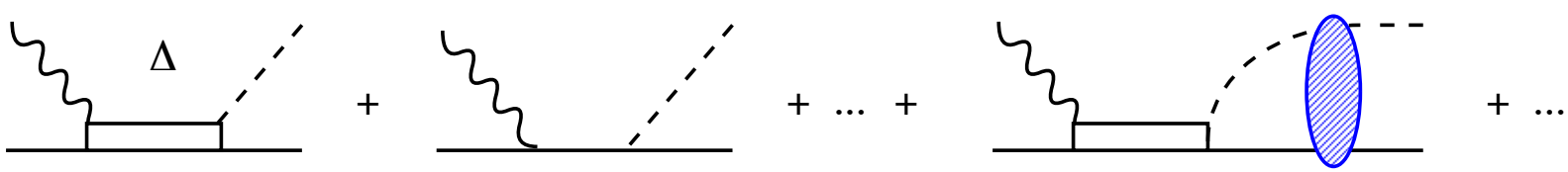
+ non-resonant (tree-level)

Hernandez et al. (2007), (2010) ; Lalakulich et al. (2010)



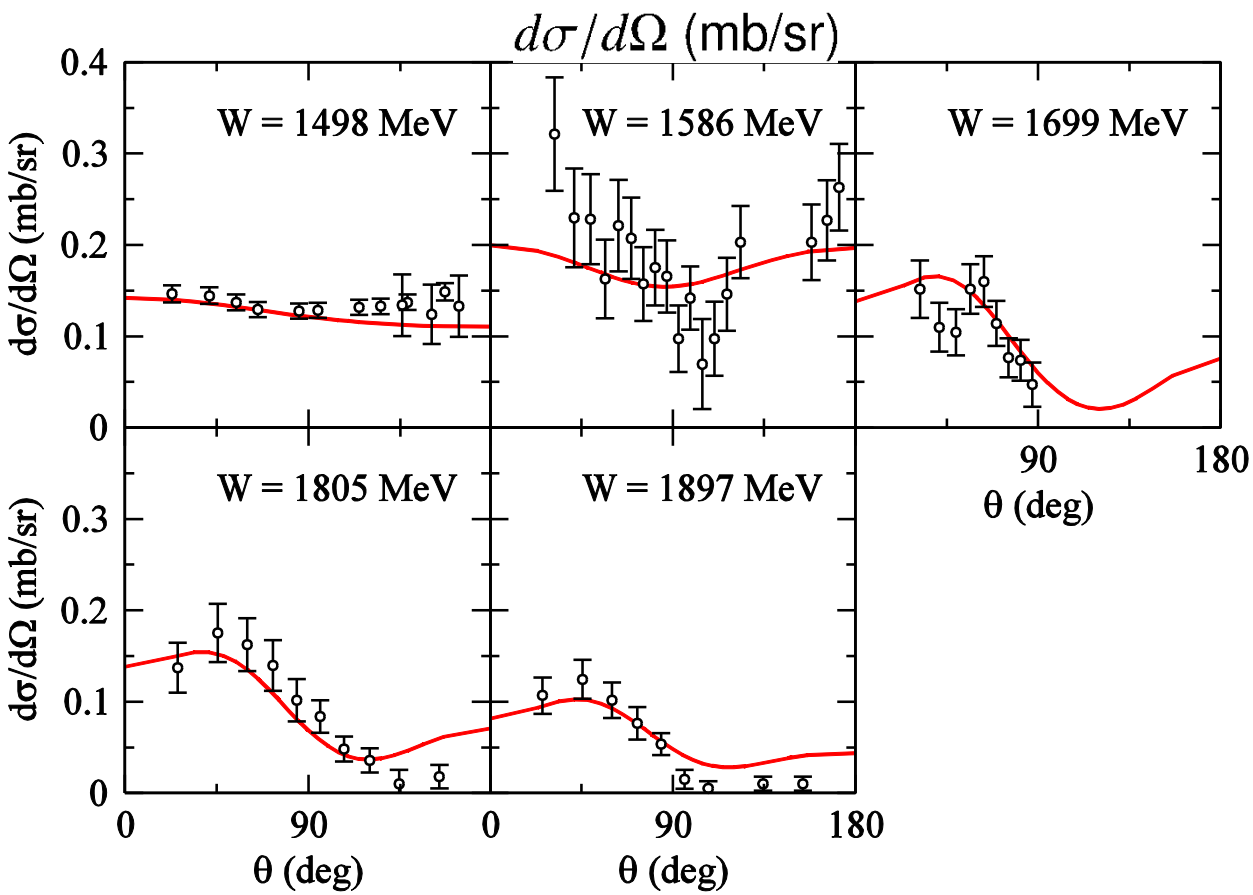
+ rescattering (πN unitarity)

Sato, Lee (2003), (2005)

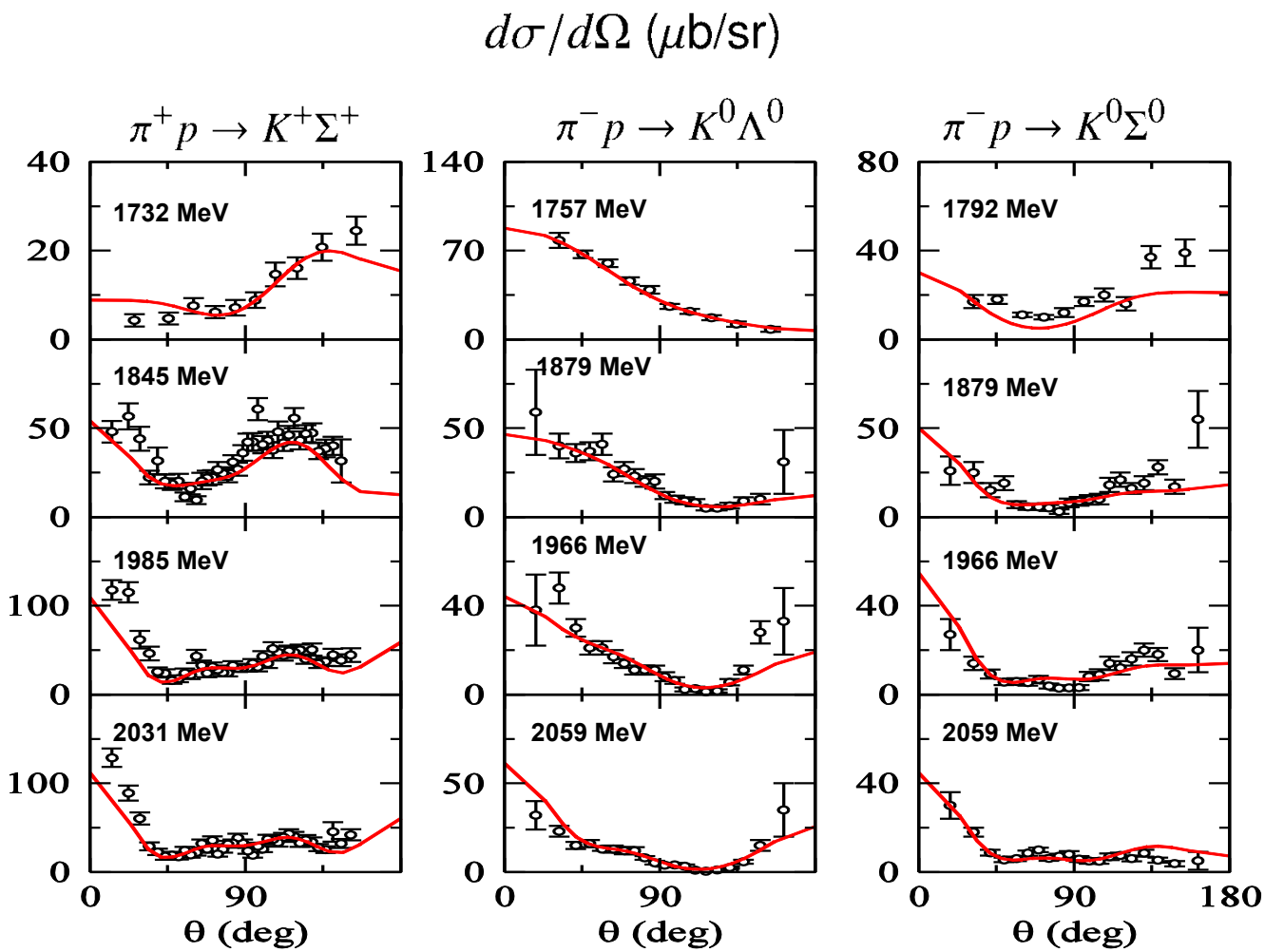


Eta production reactions

$$\pi^- p \rightarrow \eta n$$



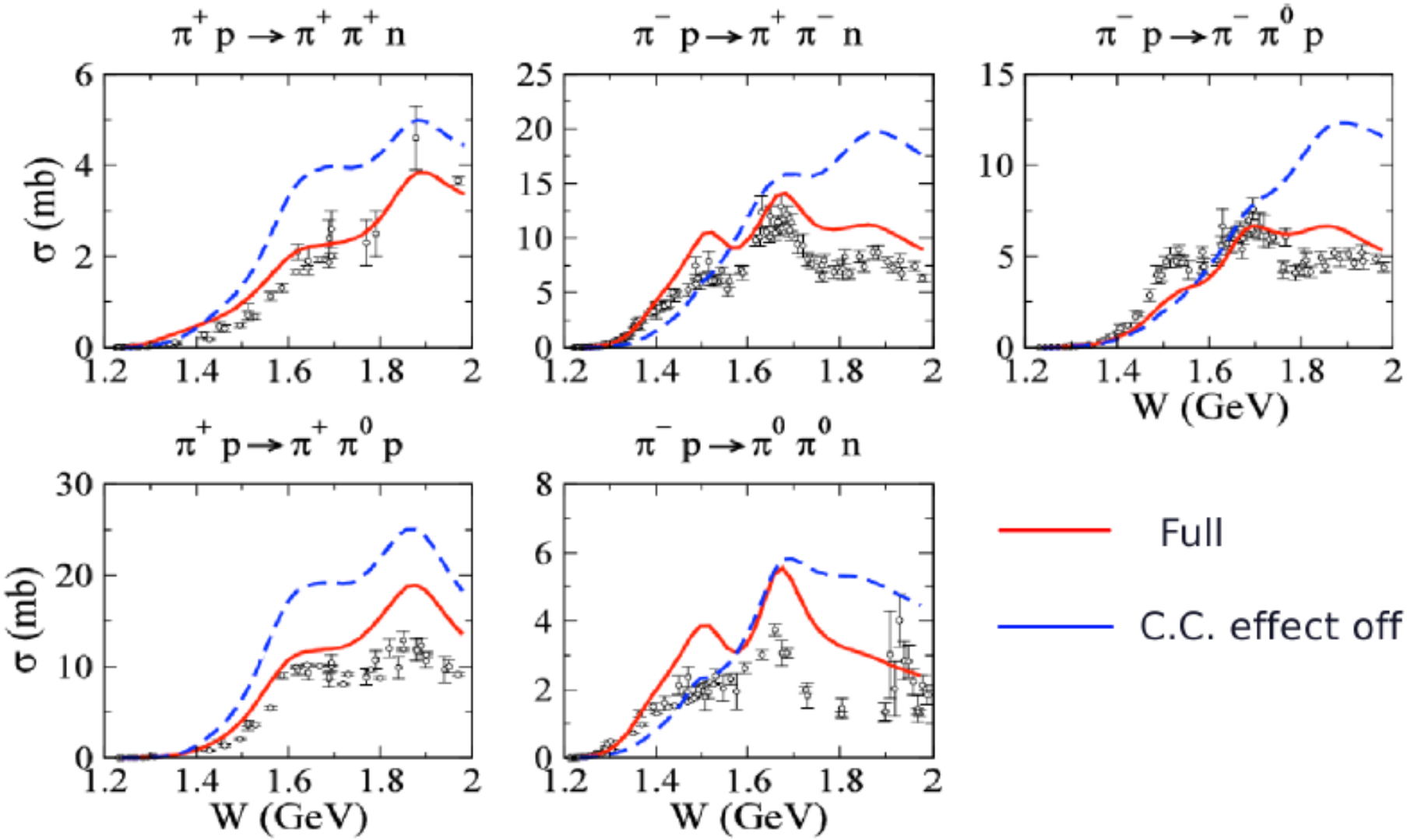
KY production reactions

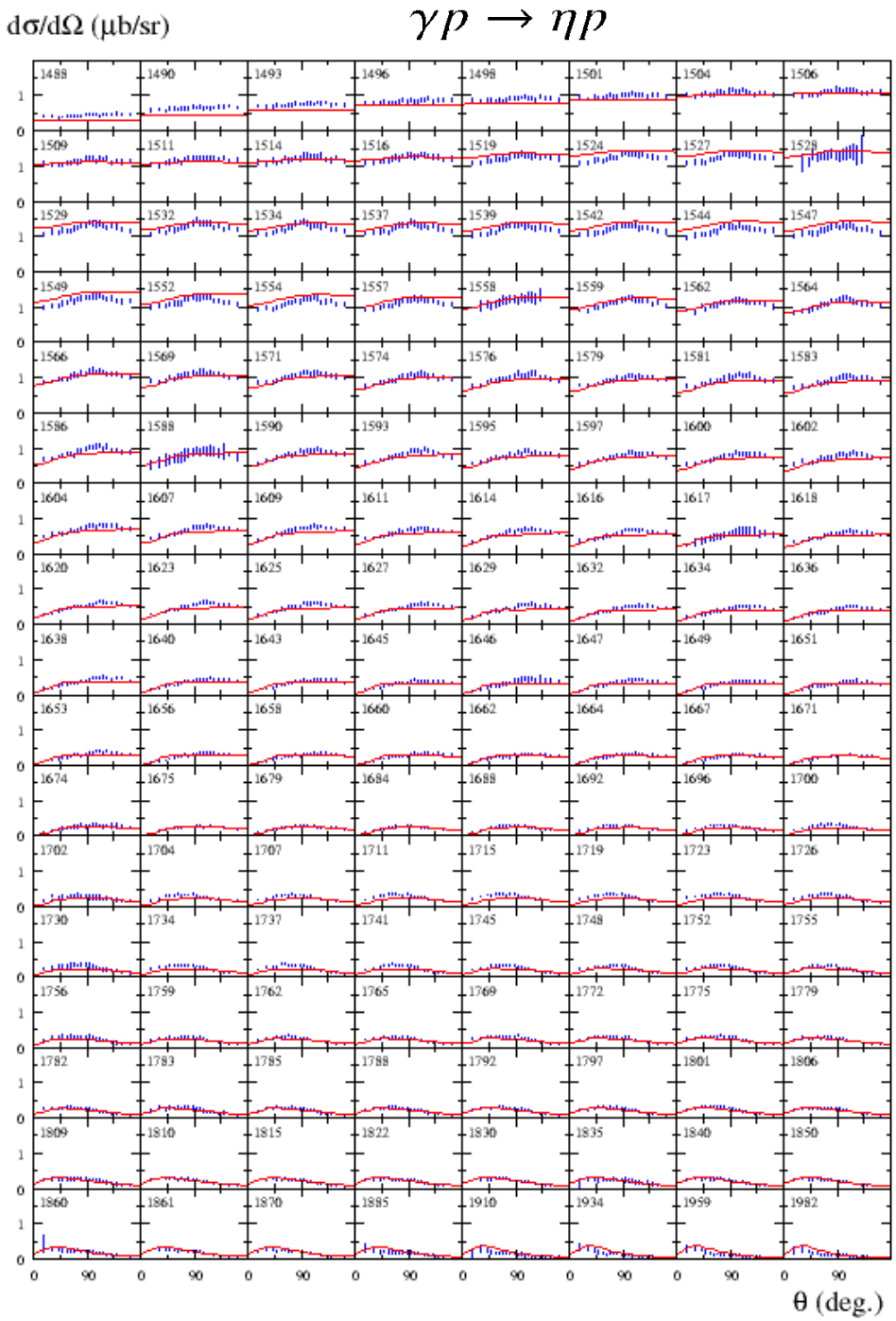


$\pi N \to \pi\pi N$

(parameters had been fitted to $\pi N \to \pi N$)

Kamano, Julia-Diaz, Lee, Matsuyama, Sato, PRC79 025206 (2009)





Kamano, Nakamura, Lee, Sato, arXiv:1305.4351

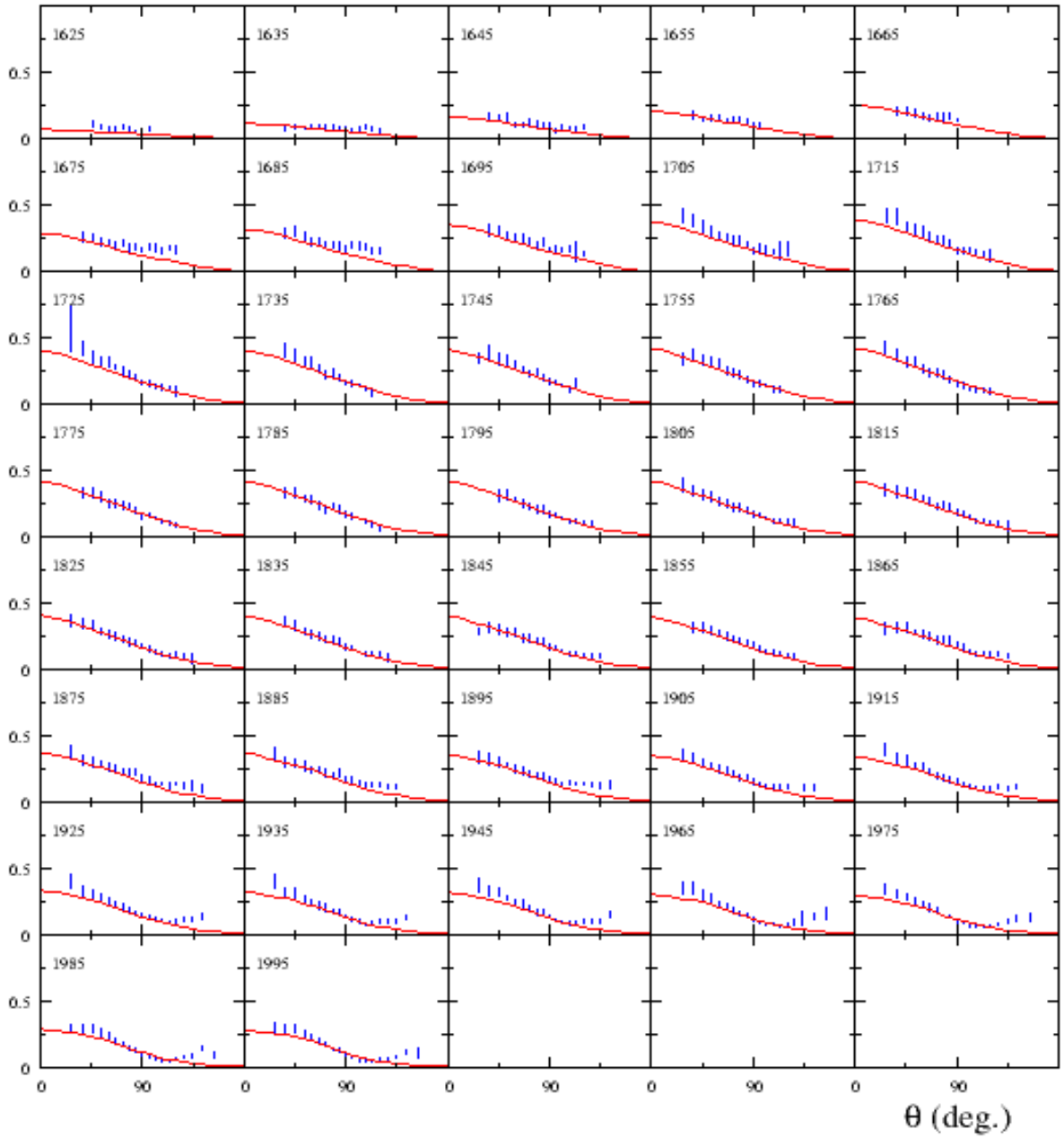
Vector current (Q^2=0) for η

Production is well-tested by data

$d\sigma/d\Omega$ (μb/sr)

$$\gamma p \rightarrow K^{+}\Lambda$$

Kamano, Nakamura, Lee, Sato, arXiv:1305.4351



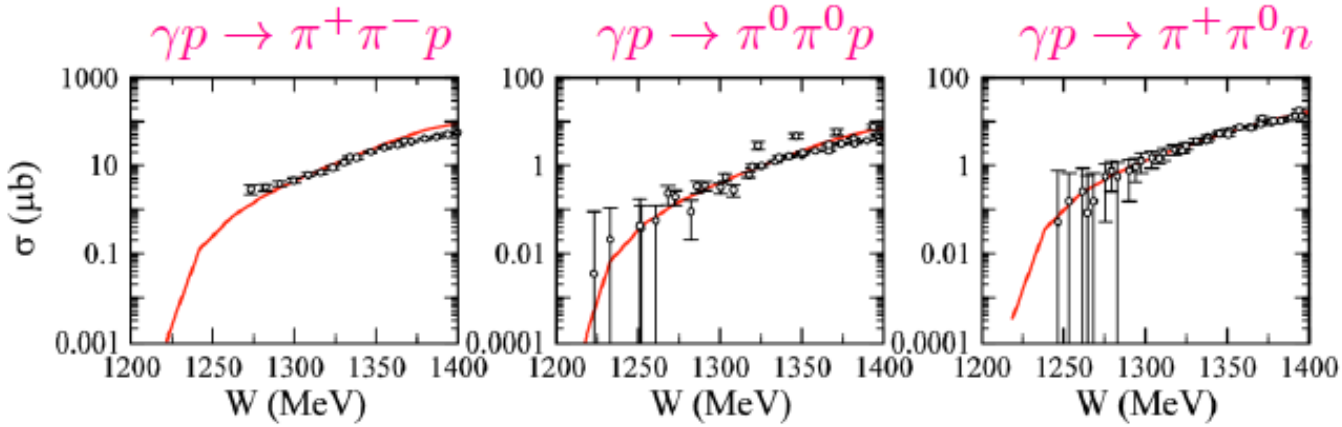
Vector current (Q^2=0) for K

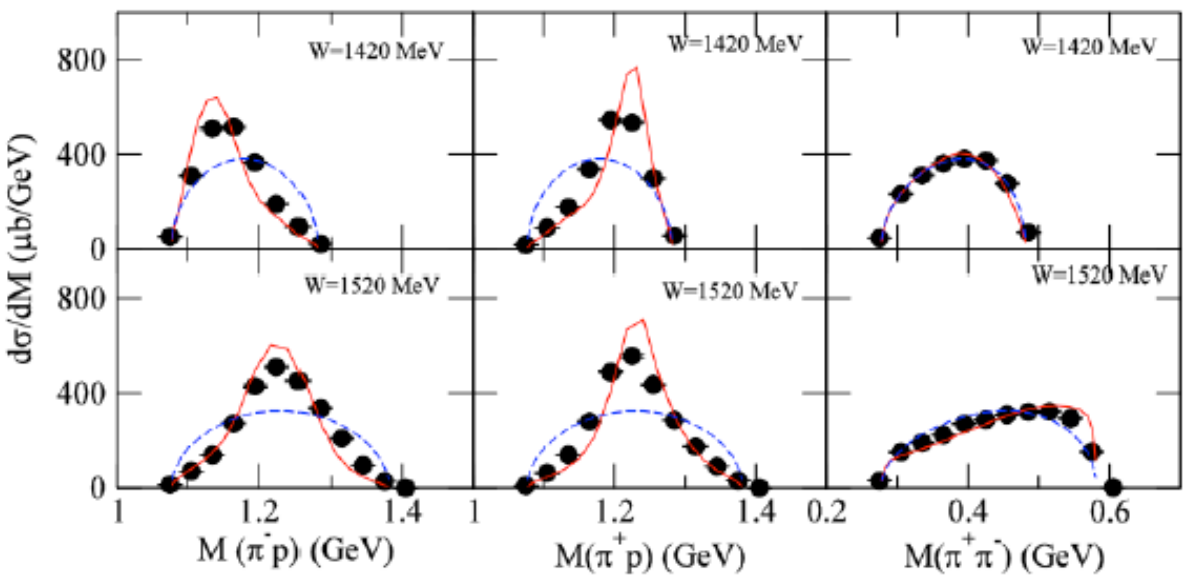
Production is well-tested by data

$\gamma N \to \pi\pi N$

(parameters had been fitted to $\pi N, \gamma N \to \pi N$)

Kamano, Julia-Diaz, Lee, Matsuyama, Sato, PRC80 065203 (2009)



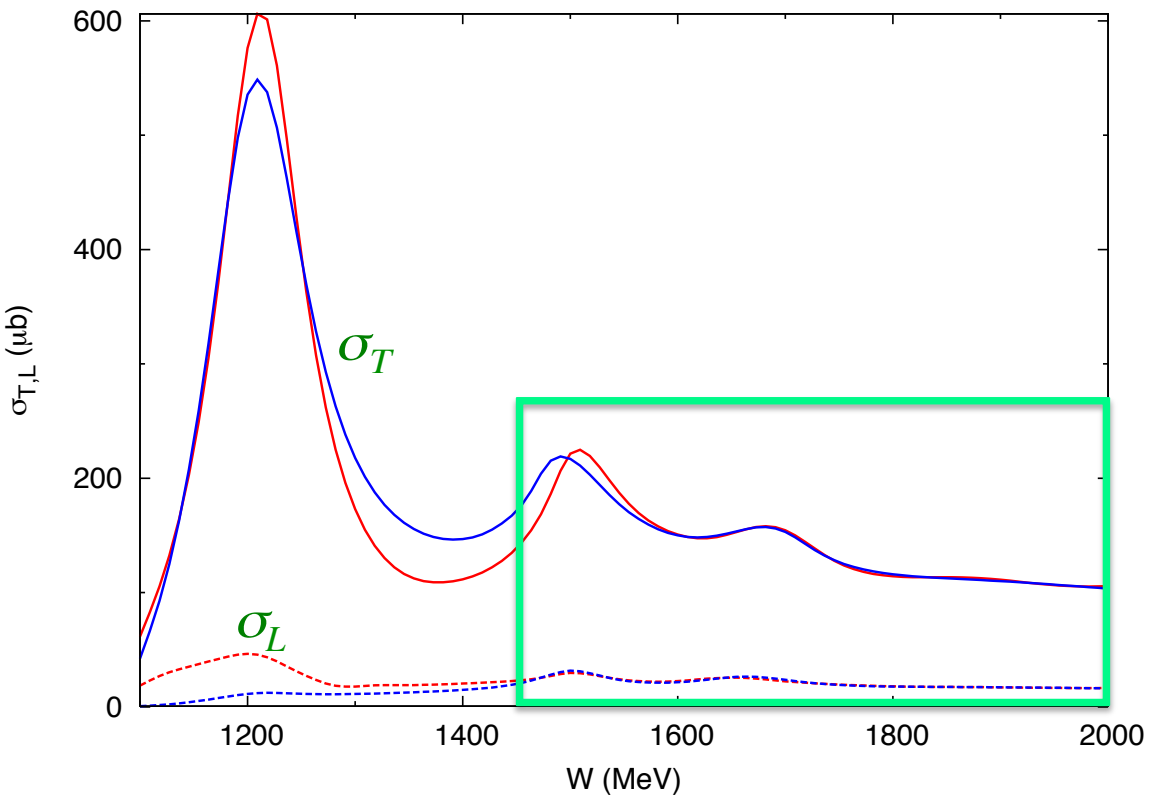


* Good description near threshold
* Good shape of invariant mass distribution
* Total cross sections overestimate data for $W \geq 1.5$ GeV

Analysis result

Q^2=0.16 (GeV/c)^2

σ_T & σ_L (inclusive inelastic)



DCC

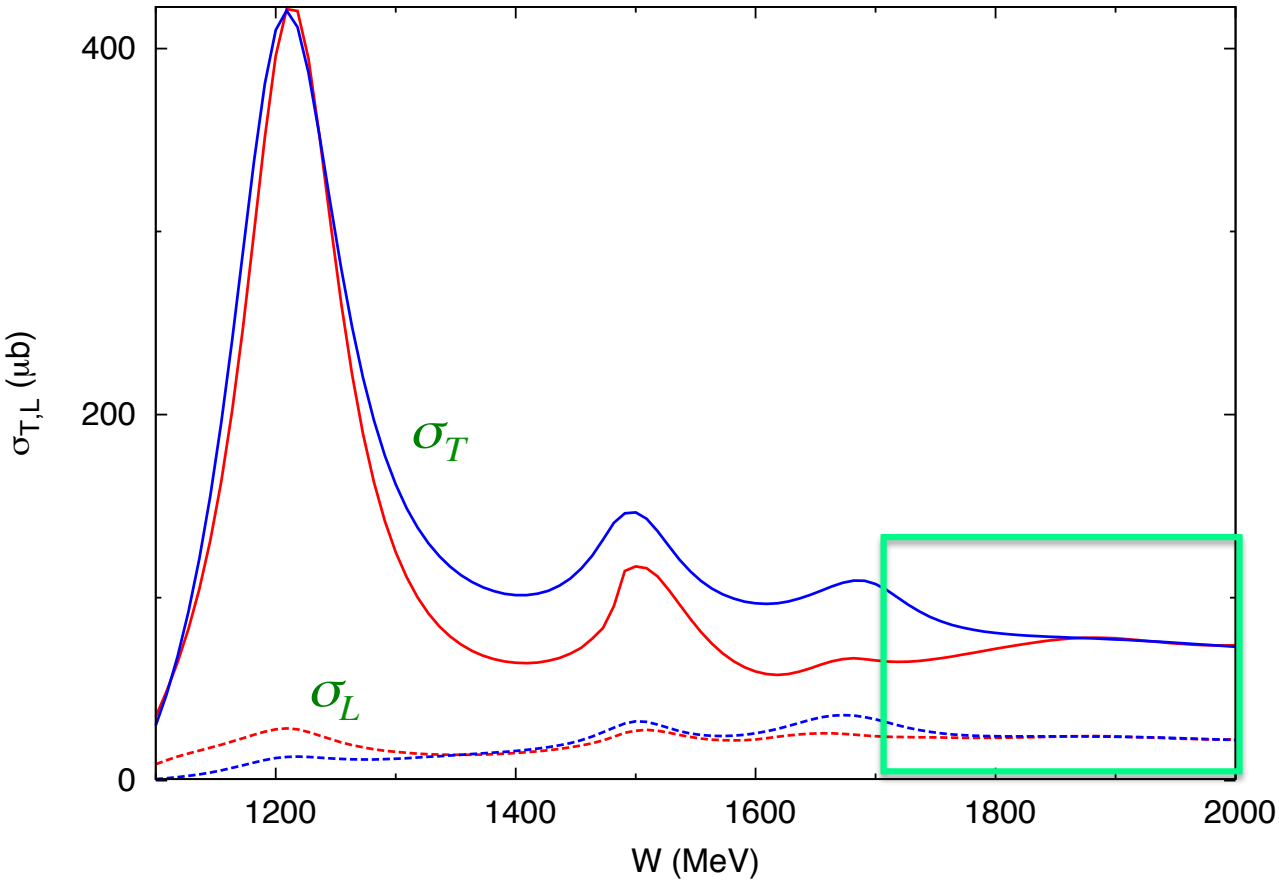
Christy et al PRC 81

region where inclusive σ_T & σ_L are fitted

Analysis result

Q^2=0.40 (GeV/c)^2

σ_T & σ_L (inclusive inelastic)



DCC

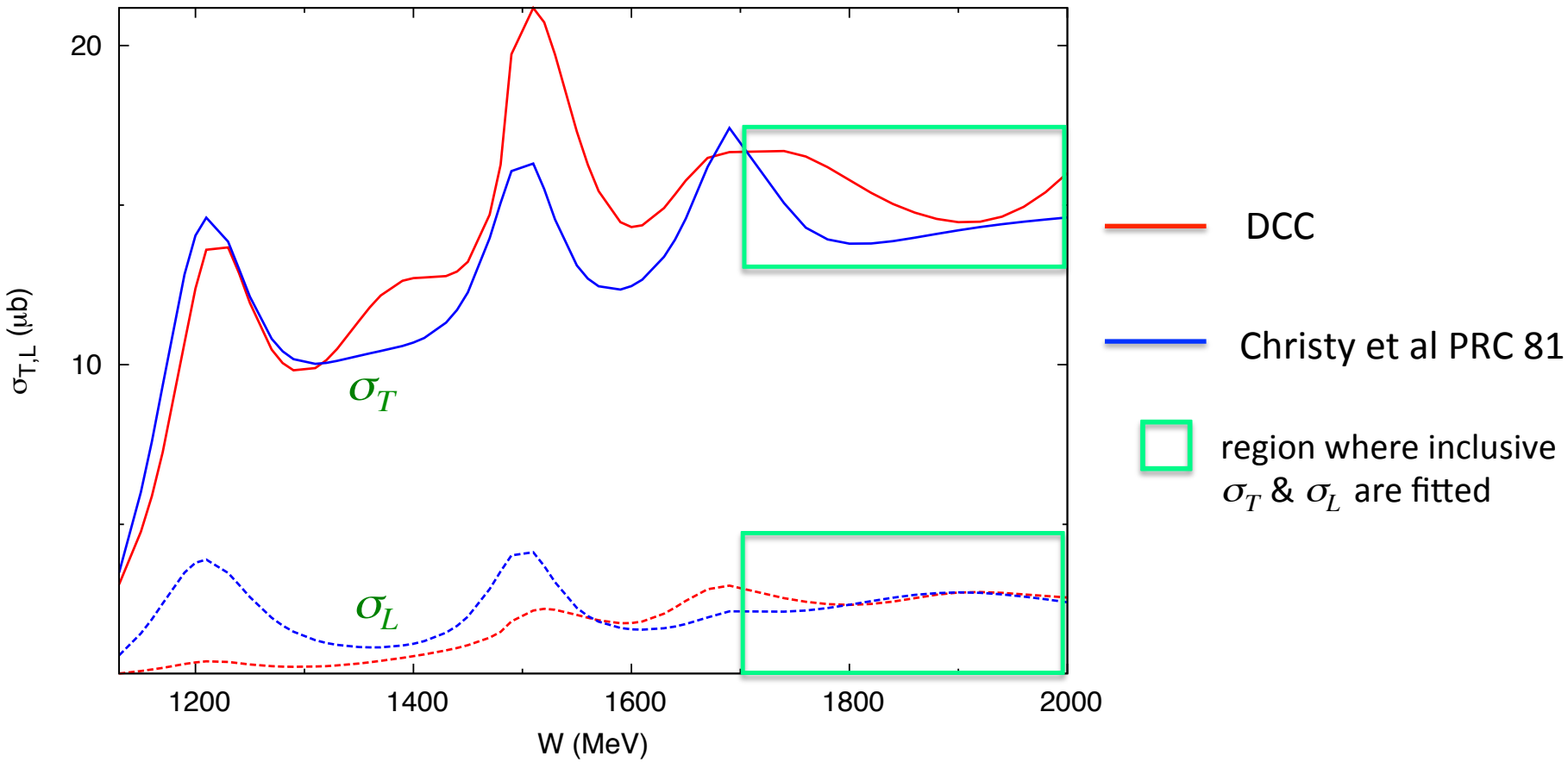
Christy et al PRC 81

region where inclusive σ_T & σ_L are fitted

Analysis result

Q^2=2.95 (GeV/c)^2

σ_T & σ_L (inclusive inelastic)

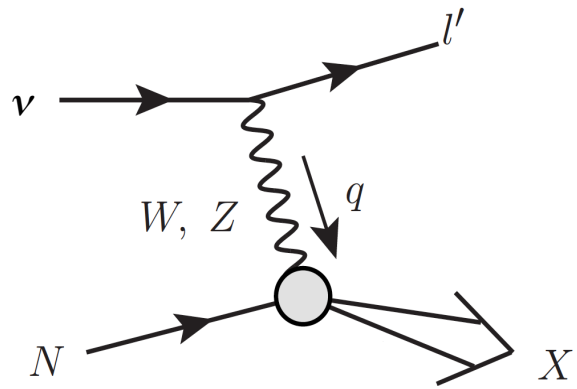


DCC model for neutrino interaction

$\nu N \rightarrow l\, X$ ($X = \pi N, \pi\pi N, \eta N, K\Lambda, K\Sigma$)

at forward limit $Q^2=0$

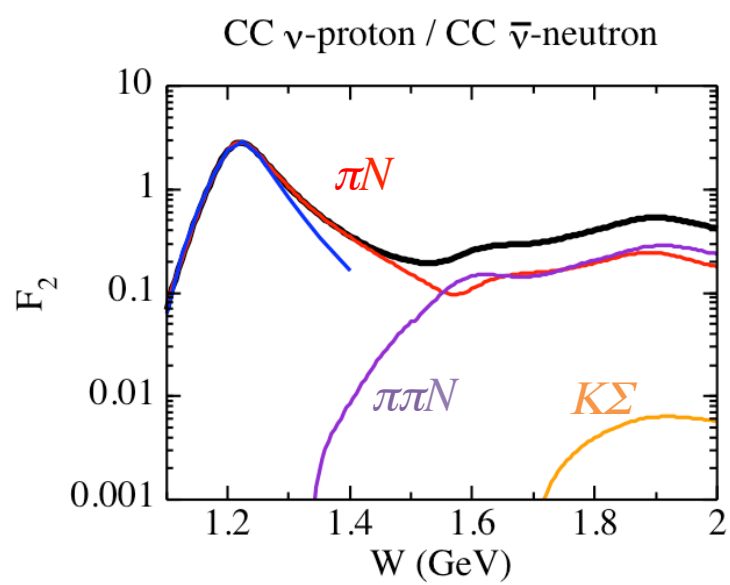
Kamano, Nakamura, Lee, Sato, PRD 86 (2012)



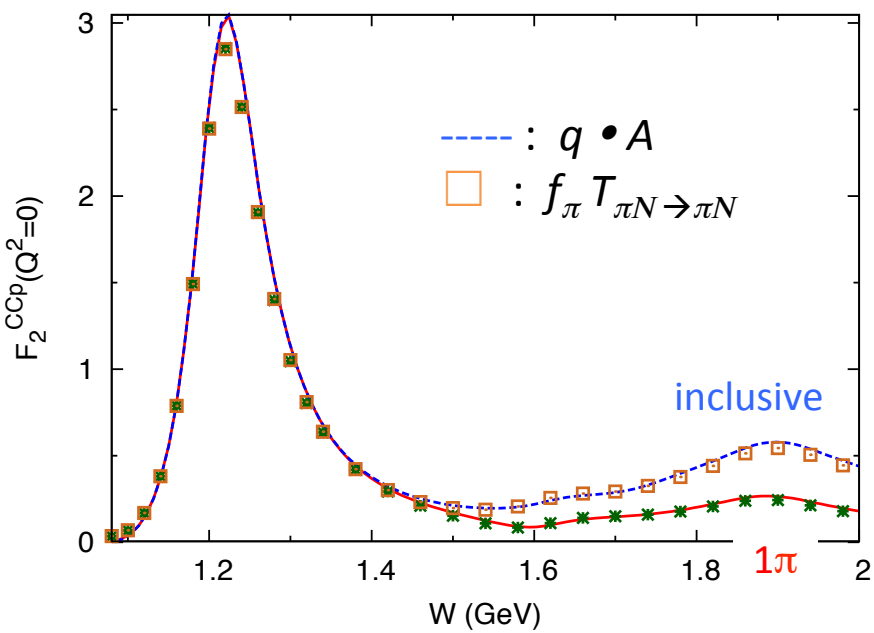
$$\frac{d\sigma}{dE_\ell d\Omega_\ell} = \frac{G_F^2}{2\pi^2} E_\ell^2\, W_2$$

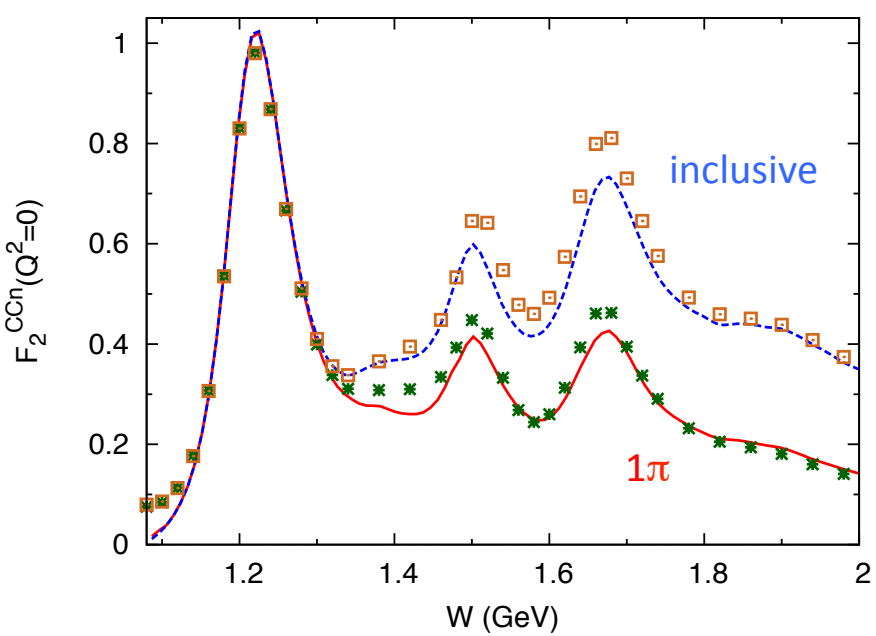
via PCAC $$F_2 \equiv \omega W_2 = \frac{2 f_\pi^2}{\pi} \sigma_{\pi N \rightarrow X}$$

$\sigma_{\pi N \rightarrow X}$ is from our DCC model



$F_2(Q^2=0)$ from DCC model and PCAC



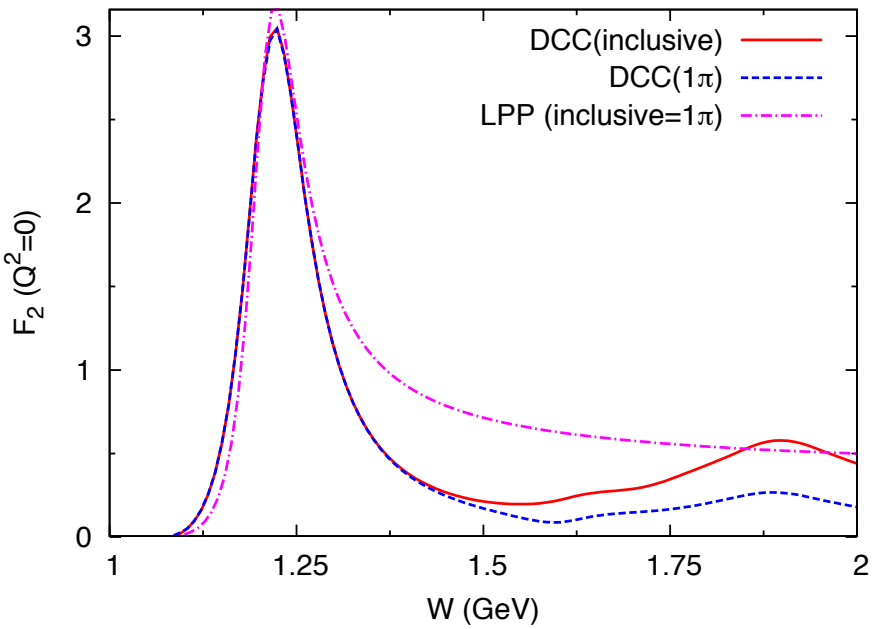


DCC model keeps good consistency with PCAC

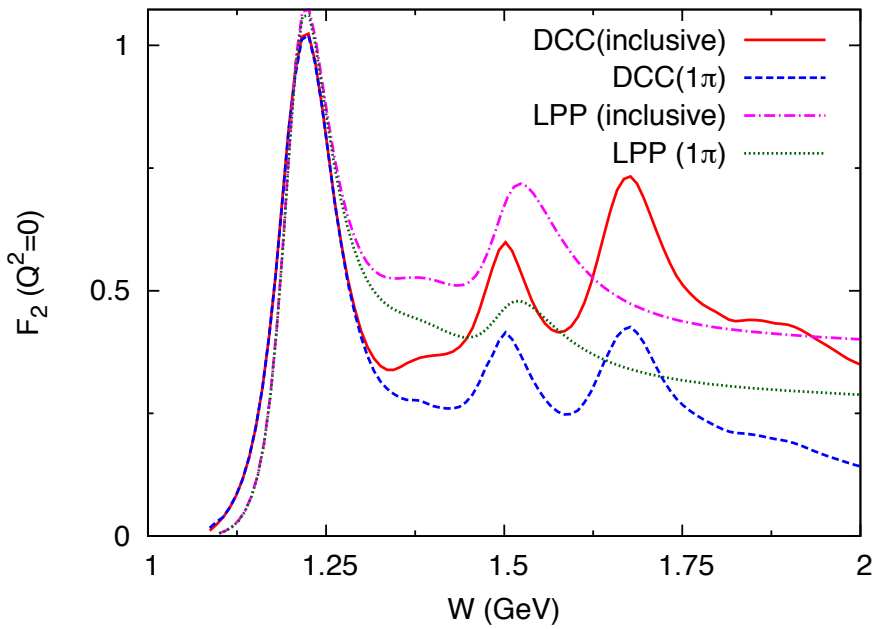
Comparison with LPP model

LPP model : Lalakulich et al, PRD 74 (2006)

CCνp



CCνn



- Large difference beyond $\Delta(1232)$ region
- Importance of consistency between axial-current and πN interaction

Future development

- Axial form factor

  more neutrino data is ideal

  $\pi N \rightarrow \rho N$ (t-ch π) (maybe possible at J-PARC)

- $\pi\pi N$ channel

  $\pi N \rightarrow \pi\pi N$ experiment (J-PARC, K. Hicks et al.)

  $\gamma N \rightarrow \pi\pi N$ experiment (ELPH, JLab)

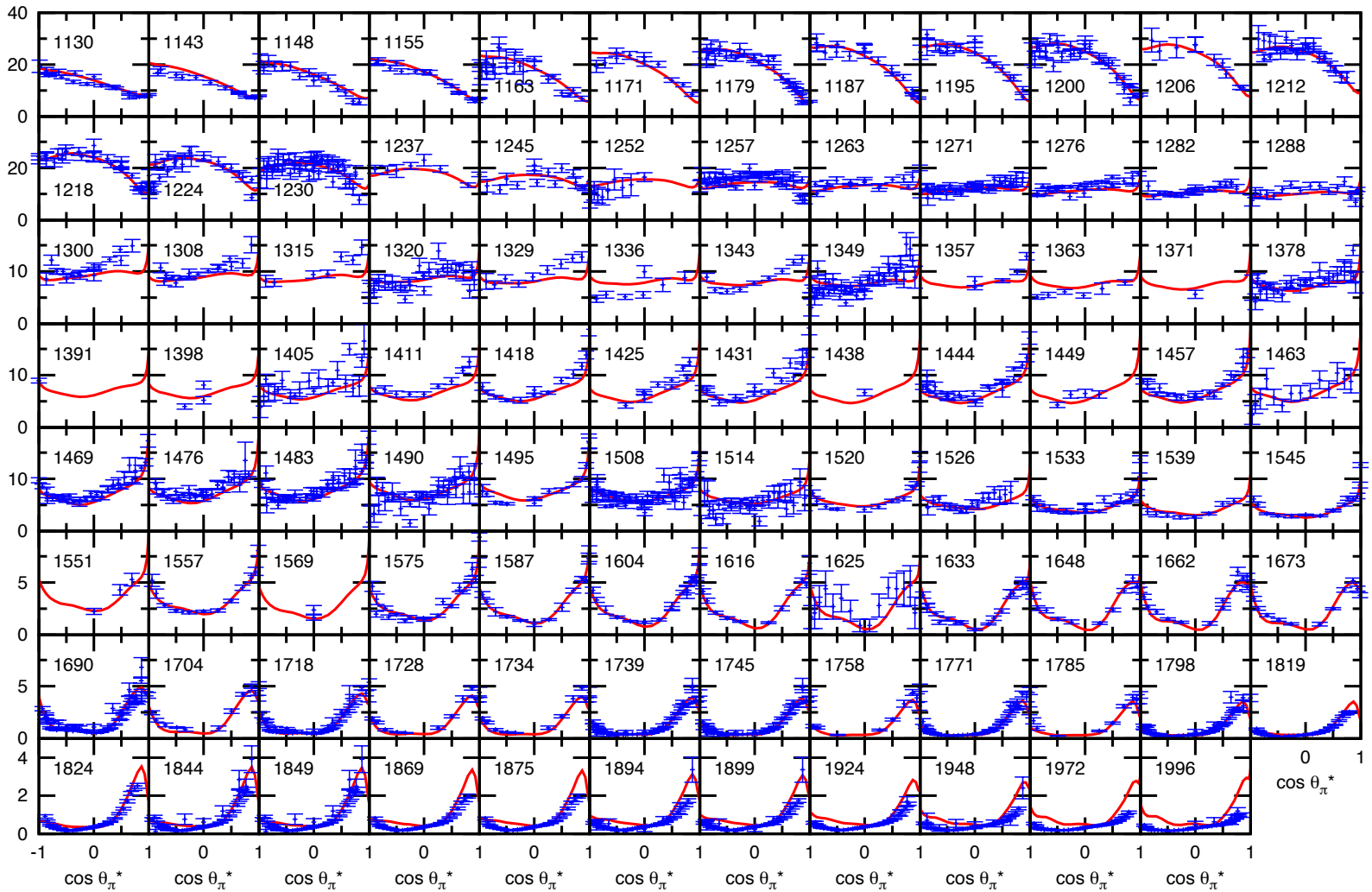
- Nuclear models

  deuteron model , Δ-hole type model

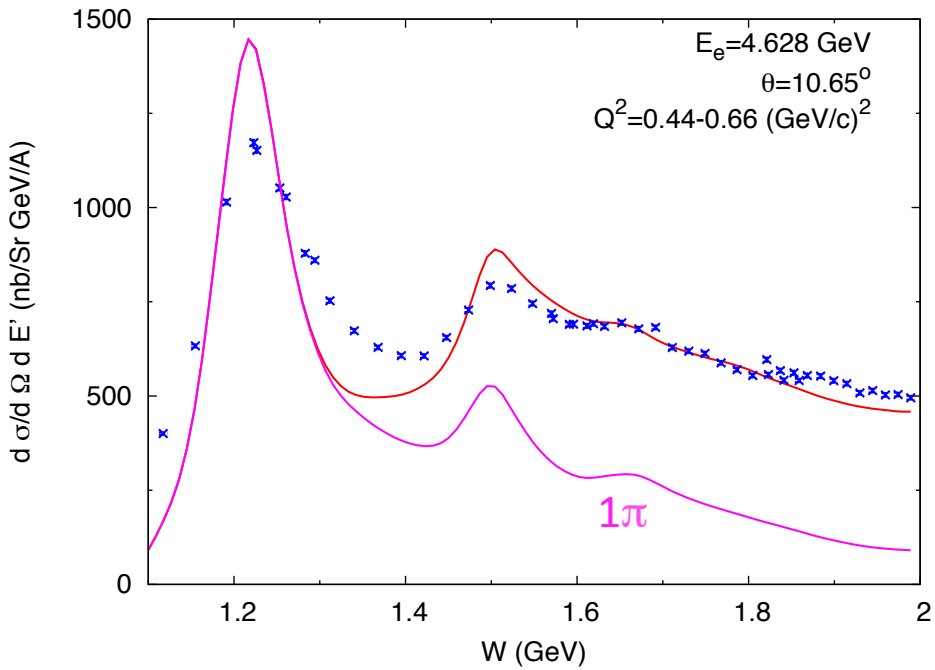
Analysis result (single π)

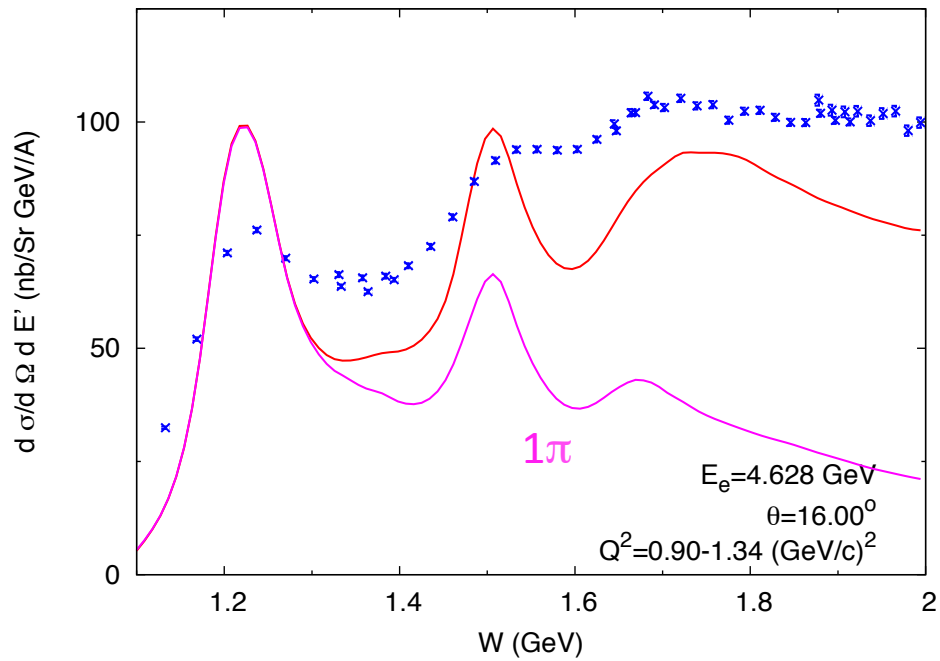
Q^2=0

$d\sigma / d\Omega$ $(\gamma n \rightarrow \pi^- p)$ for W=1.1 – 2.0 GeV



Analysis result (inclusive e^--d)





Data: NP Proc. Suppl. 159, 163 (2006)

- Our calculation : $[\ \sigma(e^-\text{-}p) + \sigma(e^-\text{-}n)\ ] / 2$
- Too sharp resonant peaks → fermi motion smearing, other nuclear effects needed
- Reasonable starting point for application to neutrino interactions

For application to neutrino interactions

Analysis of electron scattering data

➔ $VpN^*(Q^2)$ & $VnN^*(Q^2)$ fixed for several Q^2 values

➔ Parameterize $VpN^*(Q^2)$ & $VnN^*(Q^2)$ with simple analytic function of Q^2

$I=3/2$	:	$VpN^*(Q^2) = VnN^*(Q^2)$	➔ CC, NC
$I=1/2$ isovector part	:	$(VpN^*(Q^2) - VnN^*(Q^2)) / 2$	➔ CC, NC
$I=1/2$ isoscalar part	:	$(VpN^*(Q^2) + VnN^*(Q^2)) / 2$	➔ NC

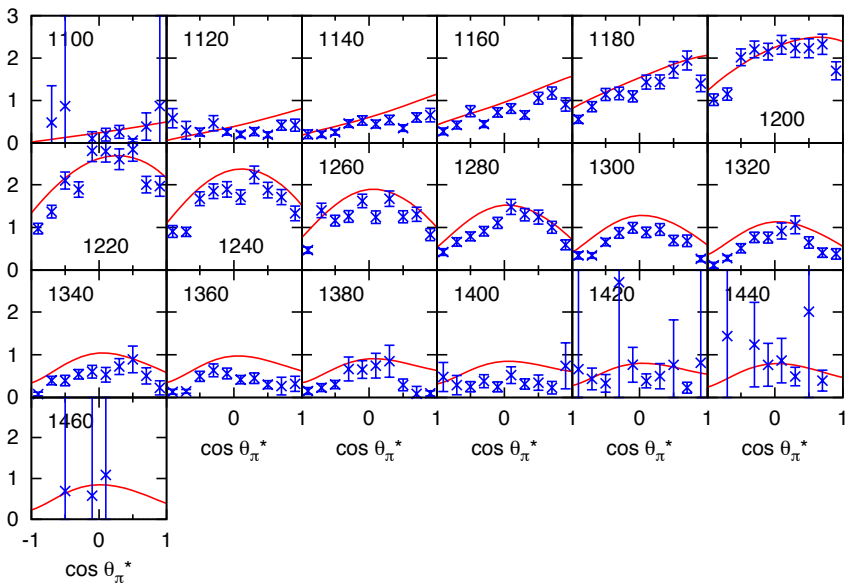
DCC vector currents has been tested by data for whole kinematical region relevant to neutrino interactions of $E_\nu \leq 2$ GeV

Analysis result (single π)

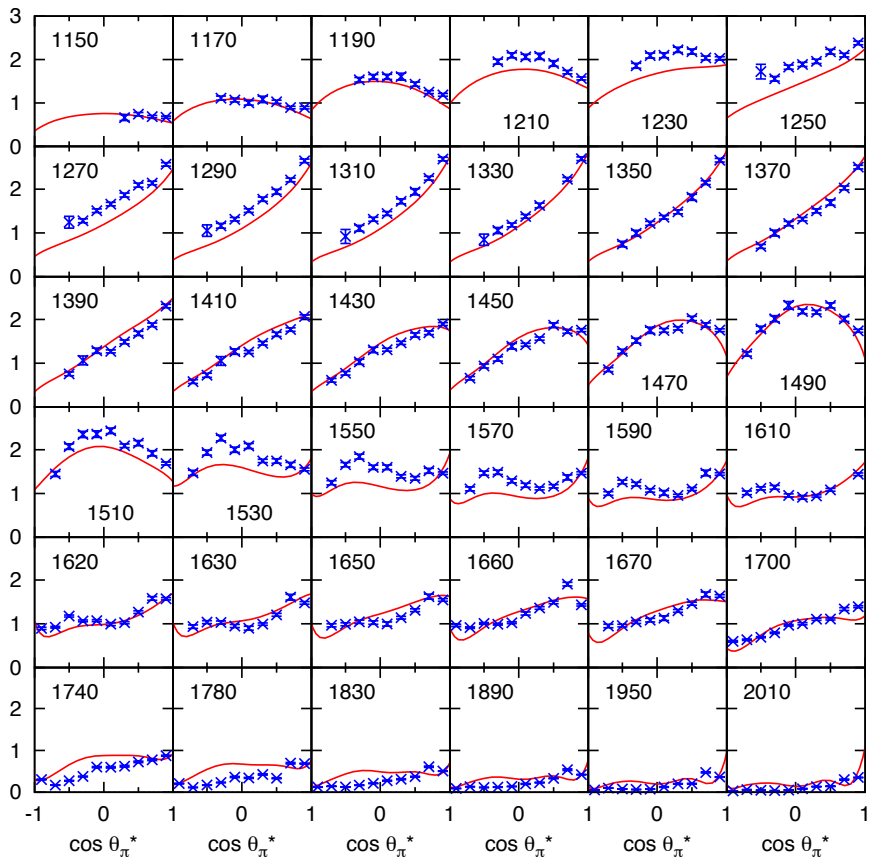
Q^2=1.76 (GeV/c)^2

$\sigma_T + \varepsilon\, \sigma_L$ for W=1.10 – 2.01 GeV

$p(e,e'\pi^0)p$



$p(e,e'\pi^+)n$

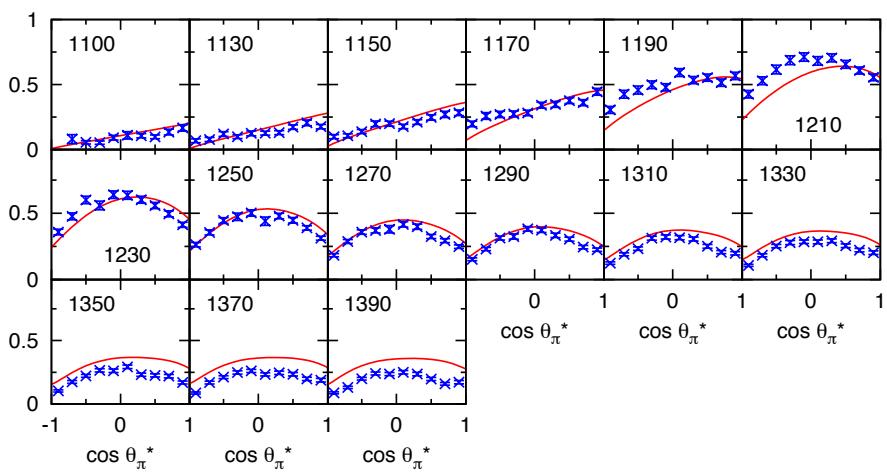
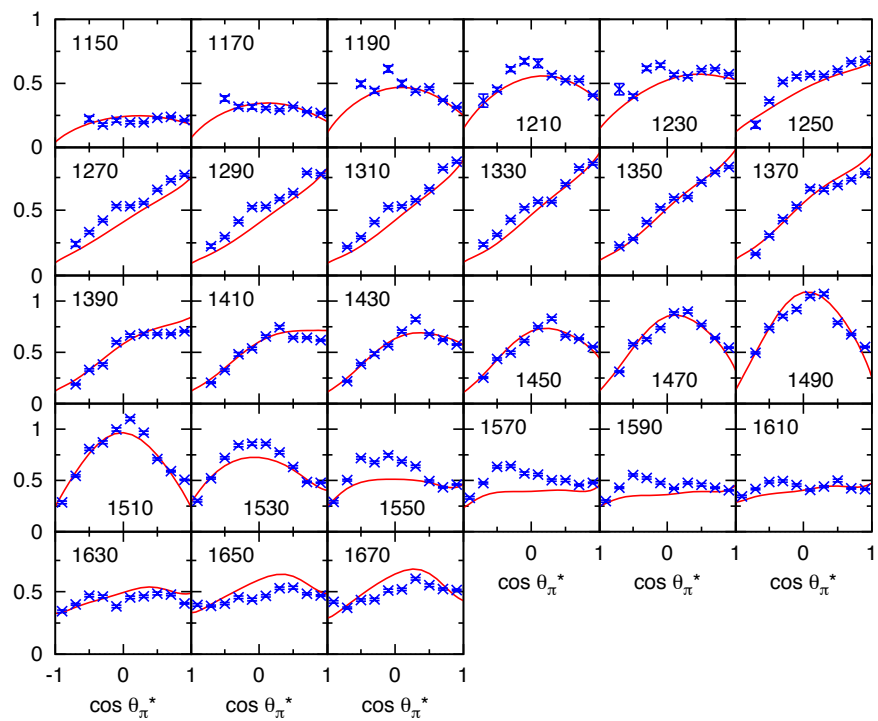


Analysis result (single π)

Q^2=2.91-3.00 (GeV/c)^2

$\sigma_T + \varepsilon\, \sigma_L$ for W=1.10 – 1.67 GeV

$p(e,e'\pi^0)p$



$p(e,e'\pi^+)n$