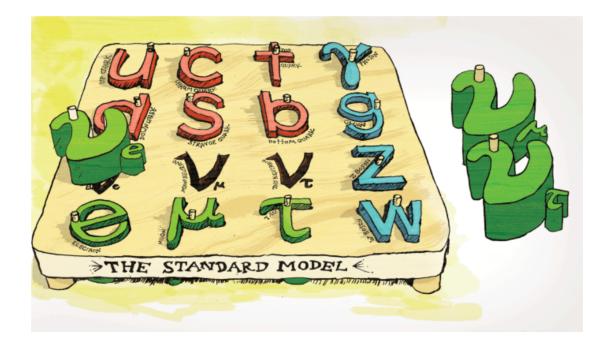
Neutrino Oscillations: Present, Future (and What for?)



André de Gouvêa – Northwestern University Seminar at MAP Meeting – SLAC December 3–7, 2014



International Committee for Future Accelerators

Sponsored by the Particles and Fields Commission of IUPAP

ICFA Neutrino Panel

News

Initial report from the ICFA Neutrino Panel (May 27, 2014)

Mission

To promote international cooperation in the development of the accelerator-based neutrino-oscillation program and to promote international collaboration in the development of a neutrino factory as a future intense source of neutrinos for particle physics experiments.

Panel

- <u>Membership</u>
- Email the panel
- Terms of Reference
- <u>Meetings</u>

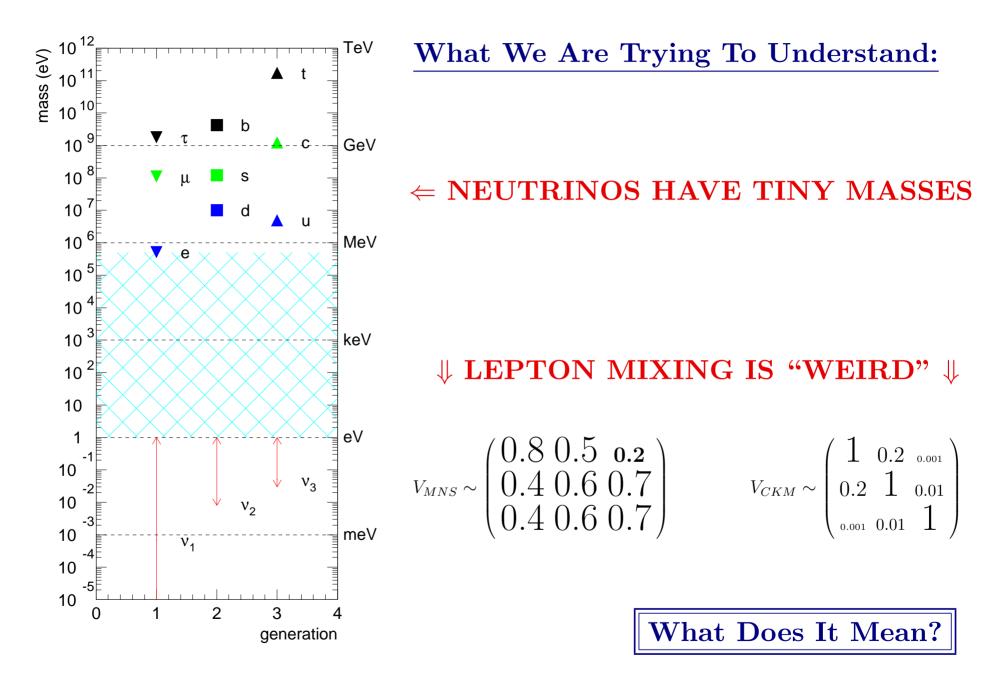
First Report:

http://arxiv.org/pdf/1405.7052.pdf

Communication

Mailing List

A mailing list has been set up to facilitate contact between the Panel and the international neutrino community and to December 3, 2014 allow discussion and information sharing within the community. Individuals may sign up to the list by visiting the list on the JiscMail site <u>here</u>.



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Future ν **s**: Oscillations

<u>Neutrino Masses</u>: Only^{*} "Palpable" Evidence of Physics Beyond the Standard Model

The SM we all learned in school predicts that neutrinos are strictly massless. Hence, massive neutrinos imply that the the SM is incomplete and needs to be replaced/modified.

Furthermore, the SM has to be replaced by something qualitatively different.

- What is the physics behind electroweak symmetry breaking? (Higgs \checkmark).
- What is the dark matter? (not in SM).
- Why is there more matter than antimatter in the Universe? (not in SM).
- Why does the Universe appear to be accelerating? Why does it appear that the Universe underwent rapid acceleration in the past? (not in SM).

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^{*} There is only a handful of questions our model for fundamental physics cannot explain (these are personal. Feel free to complain).

What is the New Standard Model? $[\nu SM]$

The short answer is – WE DON'T KNOW. Not enough available info!

\bigcirc

Equivalently, there are several completely different ways of addressing neutrino masses. The key issue is to understand what else the ν SM candidates can do. [are they falsifiable?, are they "simple"?, do they address other outstanding problems in physics?, etc]

We need more experimental input.

Neutrino Masses, Higgs Mechanism, and New Mass Scale of Nature

The LHC has revealed that the minimum SM prescription for electroweak symmetry breaking — the one Higgs double model — is at least approximately correct. What does that have to do with neutrinos?

The tiny neutrino masses point to three different possibilities.

- 1. Neutrinos talk to the Higgs boson very, very **weakly**;
- 2. Neutrinos talk to a **different Higgs** boson there is a new source of electroweak symmetry breaking!;
- 3. Neutrino masses are small because there is **another source of mass** out there a new energy scale indirectly responsible for the tiny neutrino masses, a la the seesaw mechanism.

We are going to need a lot of experimental information from all areas of particle physics in order to figure out what is really going on!

Piecing the Neutrino Mass Puzzle

Understanding the origin of neutrino masses and exploring the new physics in the lepton sector will require unique **theoretical** and **experimental** efforts ...

- understanding the fate of lepton-number. Neutrinoless double beta decay!
- A comprehensive long baseline neutrino program. (On-going T2K and $NO\nu A$. LBNF and HyperK next steps towards the ultimate "superbeam" experiment.)
- The next-step is to develop a qualitatively better neutrino beam e.g. muon storage rings (neutrino factories).
- Different baselines and detector technologies a must for both over-constraining the system and looking for new phenomena.
- Probes of neutrino properties, including neutrino scattering experiments.
- Precision measurements of charged-lepton properties (g 2, edm) and searches for rare processes $(\mu \rightarrow e\text{-conversion the best bet at the moment})$.
- Collider experiments. The LHC and beyond may end up revealing the new physics behind small neutrino masses.
- Neutrino properties affect, in a significant way, the history of the universe (Cosmology). Will we learn about neutrinos from cosmology, or about cosmology from neutrinos?

HOWEVER...

We have only ever objectively "seen" neutrino masses in long-baseline oscillation experiments. It is the clearest way forward!

Does this mean we will reveal the origin of neutrino masses with oscillation experiments? We don't know, and we won't know until we try!

[henceforth, I will concentrate on neutrino oscillations]

Something Funny Happened on the Way to the 21st Century ν Flavor Oscillations

Neutrino oscillation experiments have revealed that neutrinos change flavor after propagating a finite distance. The rate of change depends on the neutrino energy E_{ν} and the baseline L. The evidence is overwhelming.

- $\nu_{\mu} \rightarrow \nu_{\tau}$ and $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{\tau}$ atmospheric and accelerator experiments;
- $\nu_e \rightarrow \nu_{\mu,\tau}$ solar experiments;
- $\bar{\nu}_e \rightarrow \bar{\nu}_{other}$ reactor experiments;
- $\nu_{\mu} \rightarrow \nu_{\text{other}}$ and $\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{\text{other}}$ atmospheric and accelerator expts;
- $\nu_{\mu} \rightarrow \nu_{e}$ accelerator experiments.

The simplest and **only satisfactory** explanation of **all** this data is that neutrinos have distinct masses, and mix.

A Realistic, Reasonable, and Simple Paradigm:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{e\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

Definition of neutrino mass eigenstates (who are ν_1, ν_2, ν_3 ?):

- $m_1^2 < m_2^2$ $\Delta m_{13}^2 < 0$ Inverted Mass Hierarchy
- $m_2^2 m_1^2 \ll |m_3^2 m_{1,2}^2|$ $\Delta m_{13}^2 > 0$ Normal Mass Hierarchy

$$\tan^2 \theta_{12} \equiv \frac{|U_{e2}|^2}{|U_{e1}|^2}; \quad \tan^2 \theta_{23} \equiv \frac{|U_{\mu3}|^2}{|U_{\tau3}|^2}; \quad U_{e3} \equiv \sin \theta_{13} e^{-i\delta}$$

[For a detailed discussion see e.g. AdG, Jenkins, PRD78, 053003 (2008)]

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- Future νs : Oscillations

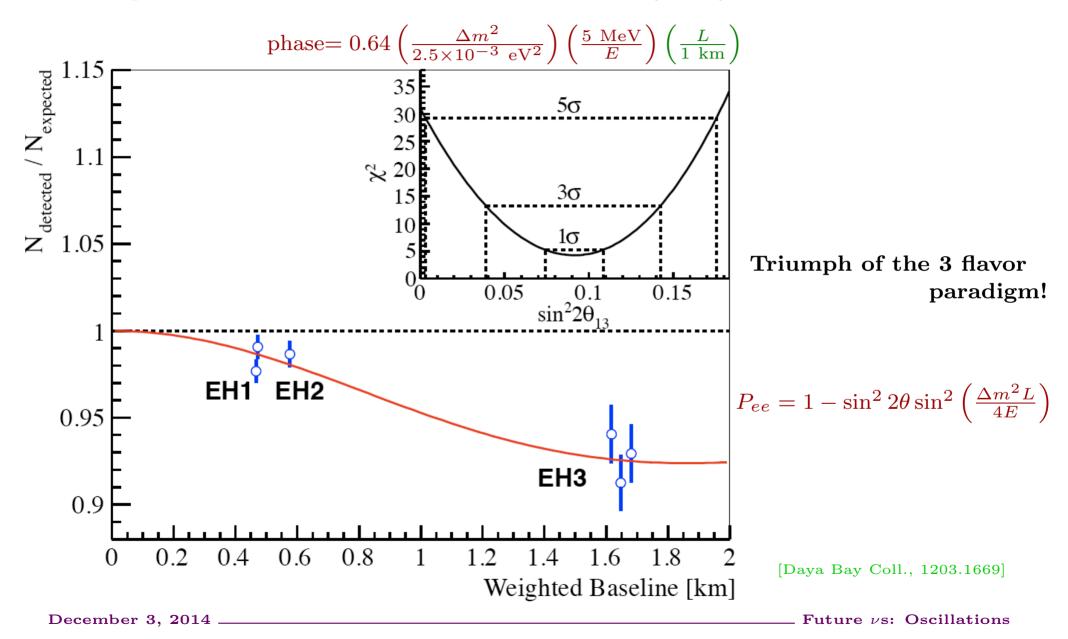
NuFIT 2.0 (2014)

	Normal Ordering $(\Delta \chi^2 = 0.97)$		Inverted Ordering (best fit)		Any Ordering
	bfp $\pm 1\sigma$	3σ range	bfp $\pm 1\sigma$	3σ range	3σ range
$\sin^2 heta_{12}$	$0.304\substack{+0.013\\-0.012}$	$0.270 \rightarrow 0.344$	$0.304\substack{+0.013\\-0.012}$	$0.270 \rightarrow 0.344$	$0.270 \rightarrow 0.344$
$ heta_{12}/^{\circ}$	$33.48_{-0.75}^{+0.78}$	$31.29 \rightarrow 35.91$	$33.48^{+0.78}_{-0.75}$	$31.29 \rightarrow 35.91$	$31.29 \rightarrow 35.91$
$\sin^2 heta_{23}$	$0.452_{-0.028}^{+0.052}$	$0.382 \rightarrow 0.643$	$0.579^{+0.025}_{-0.037}$	$0.389 \rightarrow 0.644$	$0.385 \rightarrow 0.644$
$ heta_{23}/^{\circ}$	$42.3^{+3.0}_{-1.6}$	$38.2 \rightarrow 53.3$	$49.5^{+1.5}_{-2.2}$	$38.6 \rightarrow 53.3$	$38.3 \rightarrow 53.3$
$\sin^2 heta_{13}$	$0.0218\substack{+0.0010\\-0.0010}$	$0.0186 \rightarrow 0.0250$	$0.0219\substack{+0.0011\\-0.0010}$	$0.0188 \rightarrow 0.0251$	$0.0188 \rightarrow 0.0251$
$ heta_{13}/^{\circ}$	$8.50^{+0.20}_{-0.21}$	$7.85 \rightarrow 9.10$	$8.51_{-0.21}^{+0.20}$	7.87 ightarrow 9.11	$7.87 \rightarrow 9.11$
$\delta_{ m CP}/^{\circ}$	306^{+39}_{-70}	$0 \rightarrow 360$	254_{-62}^{+63}	$0 \rightarrow 360$	$0 \rightarrow 360$
$\frac{\Delta m_{21}^2}{10^{-5} \ {\rm eV}^2}$	$7.50_{-0.17}^{+0.19}$	$7.02 \rightarrow 8.09$	$7.50_{-0.17}^{+0.19}$	$7.02 \rightarrow 8.09$	$7.02 \rightarrow 8.09$
$\frac{\Delta m_{3\ell}^2}{10^{-3} \text{ eV}^2}$	$+2.457^{+0.047}_{-0.047}$	$+2.317 \rightarrow +2.607$	$-2.449^{+0.048}_{-0.047}$	$-2.590 \rightarrow -2.307$	$ \begin{bmatrix} +2.325 \to +2.599 \\ -2.590 \to -2.307 \end{bmatrix} $

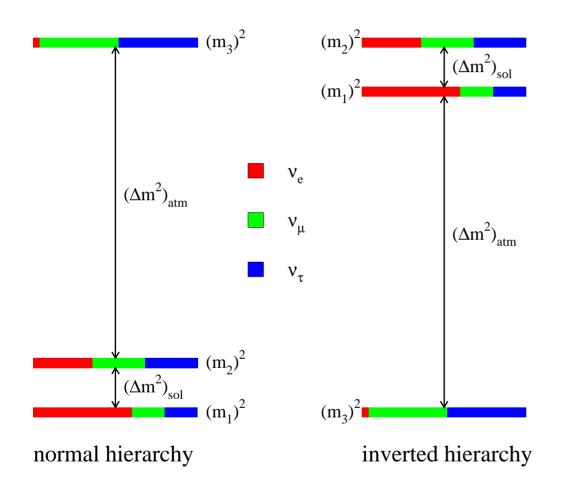
[Gonzalez-Garcia, Maltoni, Schwetz, 1409.5439, http://www.nu-fit.org]

*Modulo the Short-Baseline Anomalies, to be discussed later.

"Atmospheric Oscillations" in the Electron Sector: Daya Bay, RENO, Double Chooz

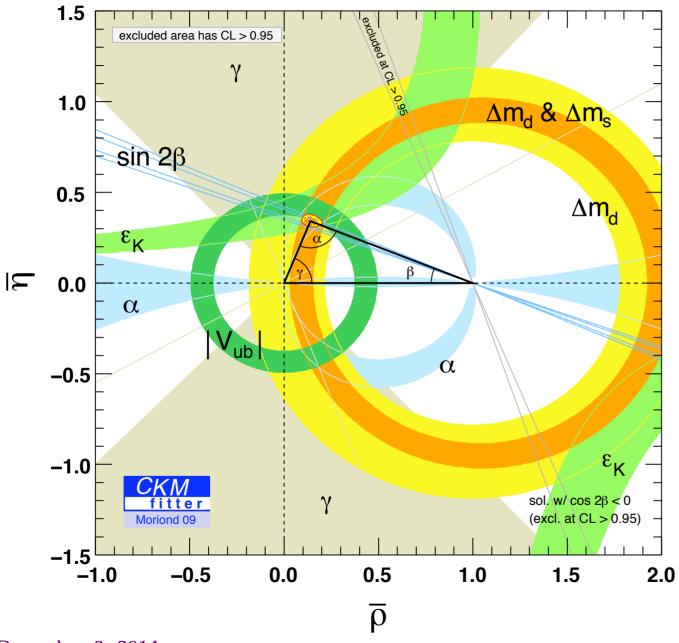


What We Know We Don't Know: Missing Oscillation Parameters



- What is the ν_e component of ν_3 ? $(\theta_{13} \neq 0!)$
- Is CP-invariance violated in neutrino oscillations? $(\delta \neq 0, \pi?)$
- Is ν_3 mostly ν_{μ} or ν_{τ} ? $(\theta_{23} > \pi/4, \theta_{23} < \pi/4, \text{ or } \theta_{23} = \pi/4?)$
- What is the neutrino mass hierarchy? $(\Delta m_{13}^2 > 0?)$
- ⇒ All of the above can "only" be addressed with new neutrino oscillation experiments

Ultimate Goal: Not Measure Parameters but Test the Formalism (Over-Constrain Parameter Space)



What we ultimately want to achieve:

We need to do <u>this</u> in the lepton sector!

Future ν **s**: Oscillations

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$$\left(\begin{array}{c}\nu_{e}\\\nu_{\mu}\\\nu_{\tau}\end{array}\right) = \left(\begin{array}{ccc}U_{e1}&U_{e2}&U_{e3}\\U_{\mu1}&U_{\mu2}&U_{\mu3}\\U_{\tau1}&U_{\tau2}&U_{\tau3}\end{array}\right) \left(\begin{array}{c}\nu_{1}\\\nu_{2}\\\nu_{3}\end{array}\right)$$

What we have **really measured** (very roughly):

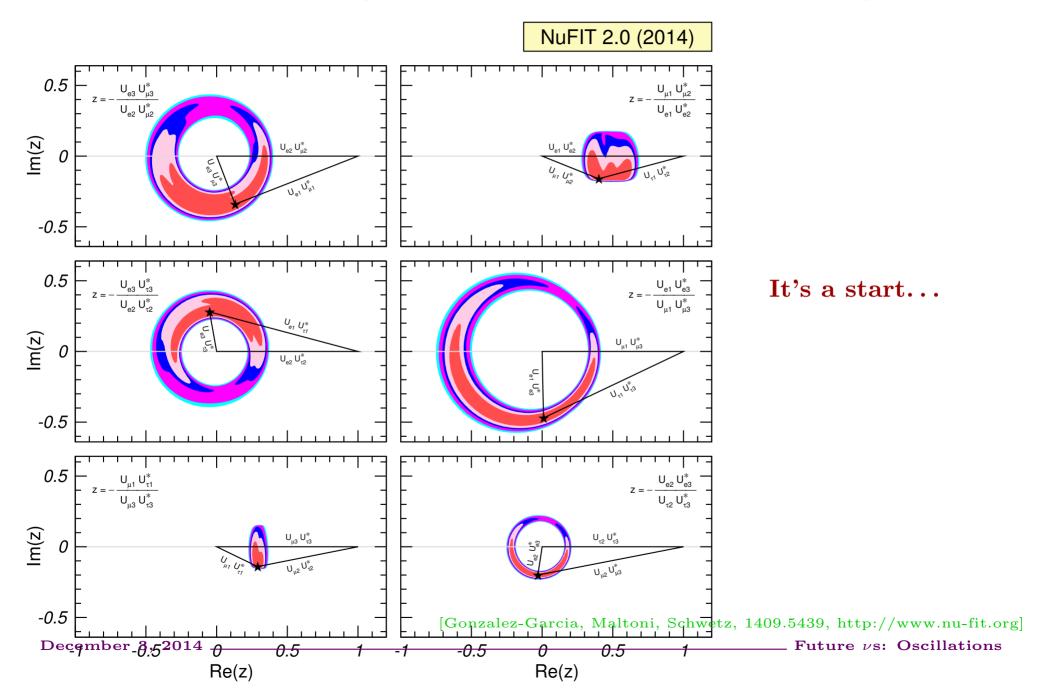
- Two mass-squared differences, at several percent level many probes;
- $|U_{e2}|^2$ solar data;
- $|U_{\mu 2}|^2 + |U_{\tau 2}|^2 \text{solar data};$
- $|U_{e2}|^2 |U_{e1}|^2 \text{KamLAND};$
- $|U_{\mu3}|^2(1-|U_{\mu3}|^2)$ atmospheric data, K2K, MINOS;
- $|U_{e3}|^2(1-|U_{e3}|^2)$ Double Chooz, Daya Bay, RENO;
- $|U_{e3}|^2 |U_{\mu3}|^2$ (upper bound \rightarrow evidence) MINOS, T2K.

We still have a ways to go!

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Future νs : Oscillations

Where We Are (?) [This is Not a Proper Comparison Yet!]



What New New Physics Might We Stumble Upon?

Large deviations from the three-flavor paradigm are allowed, including (not exhaustive list)

- There may be **more than three neutrino mass eigenstates**. These would manifest themselves in the form of new oscillation lengths or apparent non-Unitarity of the mixing matrix.
- Neutrinos may participate in **new**, weaker-than-weak interactions, which lead to nonstandard matter effects. [Concrete models are constrained by other experimental probes (usually charged-leptons) but there are scenarios, at the "existence-proof" level, that can only be constrained by neutrino oscillations.]
- Neutrino propagation may deviate from standard expectations. E.g., the **neutrinos might decay**, which leads to new decay lengths and potentially new mixing phenomena. "Macroscopic quantum interference" also very sensitive to **Lorentz invariance violation**, **new quantum mechanical decoherence** from hypothetical "neutrino-vacuum" interactions. Tiny neutrino masses also allow stringent **tests of the CPT-theorem**. (e.g., do fermions and antifermions have exactly the same mass?)

Golden Opportunity to Understand Matter versus Antimatter?

The SM with massive Majorana neutrinos accommodates **five** irreducible CP-invariance violating phases.

- One is the phase in the CKM phase. We have measured it, it is large, and we don't understand its value. At all.
- One is the θ_{QCD} term ($\theta G \tilde{G}$). We don't know its value but it is constrained to be very small. We don't know why. There are some good ideas, yet to be confirmed by experiment.
- Three are in the neutrino sector. One can be measured via neutrino oscillations. Instantaneous 50% increase on the amount of information!

We don't know much about CP-invariance violation. Is it really fair to presume that CP-invariance is generically violated in the neutrino sector solely based on the fact that it is violated in the quark sector? Why? Cautionary tale: "Mixing angles are small"

CP-invariance Violation in Neutrino Oscillations

The most promising approach to studying CP-violation in the leptonic sector seems to be to compare $P(\nu_{\mu} \rightarrow \nu_{e})$ versus $P(\bar{\nu}_{\mu} \rightarrow \bar{\nu}_{e})$.

The amplitude for $\nu_{\mu} \rightarrow \nu_{e}$ transitions can be written as

$$A_{\mu e} = U_{e2}^* U_{\mu 2} \left(e^{i\Delta_{12}} - 1 \right) + U_{e3}^* U_{\mu 3} \left(e^{i\Delta_{13}} - 1 \right)$$

where $\Delta_{1i} = \frac{\Delta m_{1i}^2 L}{2E}, i = 2, 3.$

The amplitude for the CP-conjugate process can be written as

$$\bar{A}_{\mu e} = U_{e2} U_{\mu 2}^* \left(e^{i\Delta_{12}} - 1 \right) + U_{e3} U_{\mu 3}^* \left(e^{i\Delta_{13}} - 1 \right).$$

[I assume the unitarity of $U, U_{e1}U_{\mu 1}^* = -U_{e2}U_{\mu 2}^* - U_{e3}U_{\mu 3}^*$]

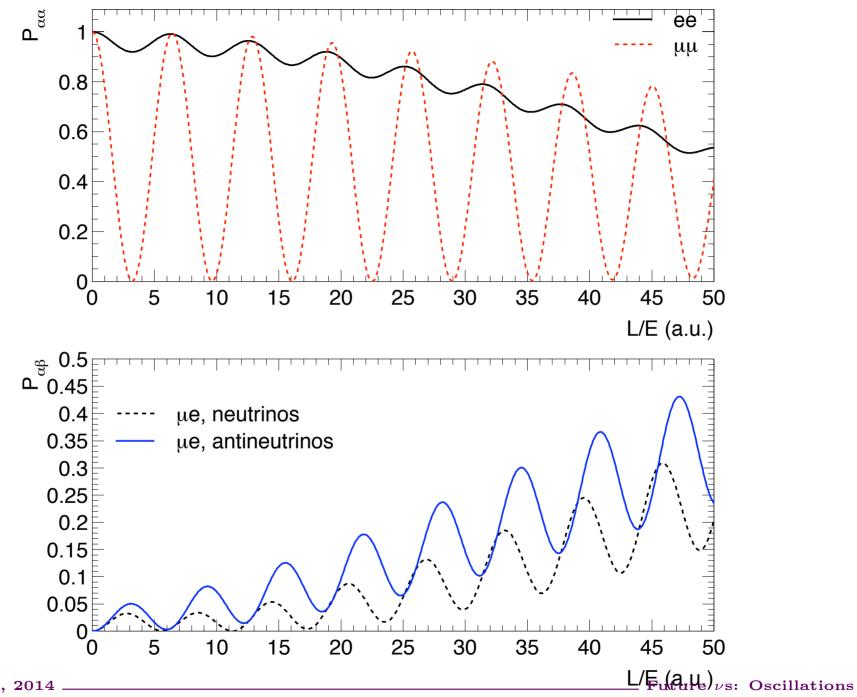
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Future νs : Oscillations

In general, $|A|^2 \neq |\bar{A}|^2$ (CP-invariance violated) as long as:

- Nontrivial "Weak" Phases: $\arg(U_{ei}^*U_{\mu i}) \to \delta \neq 0, \pi$;
- Nontrivial "Strong" Phases: $\Delta_{12}, \Delta_{13} \rightarrow L \neq 0$;
- Because of Unitarity, we need all $|U_{\alpha i}| \neq 0 \rightarrow$ three generations.

All of these can be satisfied, with a little luck: we needed $|U_{e3}| \neq 0$.



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How Precisely Do We Need to Measure Lepton Mixing Anyway?

In order to test the formalism, i.e., to look for new new physics we need to measure the same thing in different ways very precisely. "As precisely as possible."

This is clearly not a satisfactory answer, but honest answers are model dependent, and depend on the question.

Neutrino oscillations themselves, instead, offer some interesting goal posts a.k.a. **small parameters**. These include the ratio of the mass-squared differences and the smallest of the mixing parameters

$$\frac{\Delta m_{12}^2}{|\Delta m_{13}^2|} \simeq 0.03; \quad |U_{e3}| \simeq 0.15.$$

"Flavor models" often predict that mixing parameters are non-trivially related. In order to test that, we need to measure all mixing parameters with similar precision. Except for δ , $\sin^2 \theta_{23}$ has, currently the largest uncertainty (relative and absolute), while $\sin^2 \theta_{12}$ and $\sin^2 \theta_{13}$ are known at the 5% level.

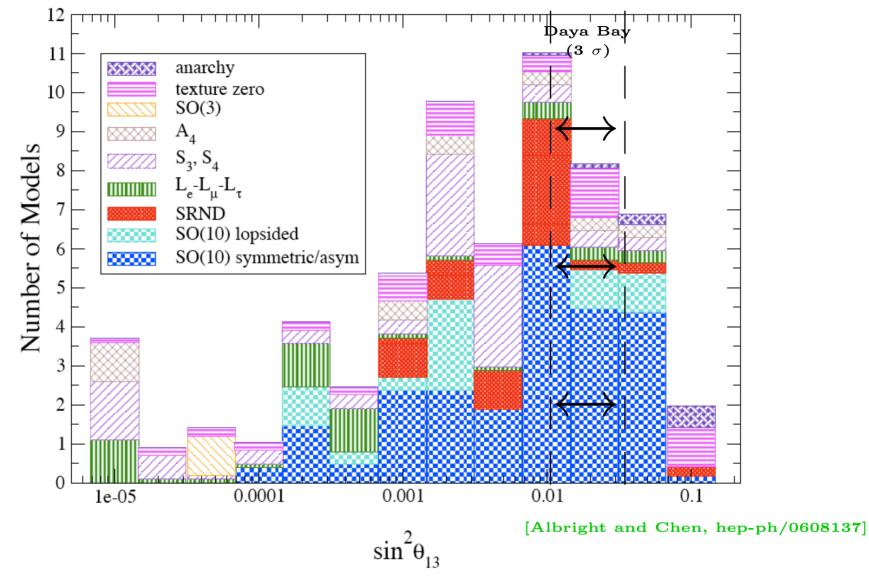
[toy example and some comments on flavor \rightarrow]

Understanding Fermion Mixing

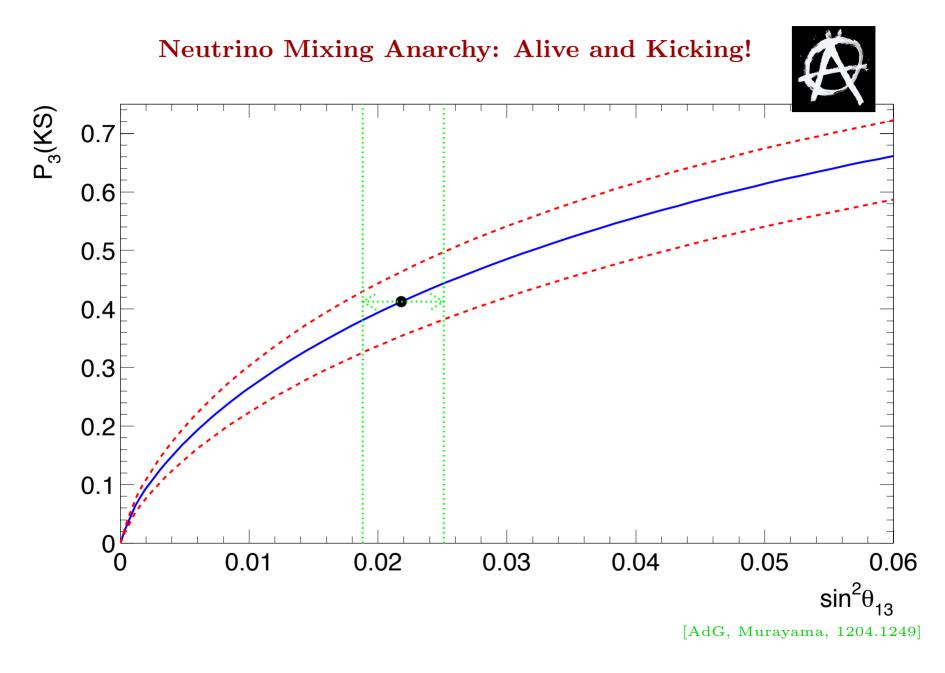
One of the puzzling phenomena uncovered by the neutrino data is the fact that Neutrino Mixing is Strange. What does this mean? It means that lepton mixing is very different from quark mixing:

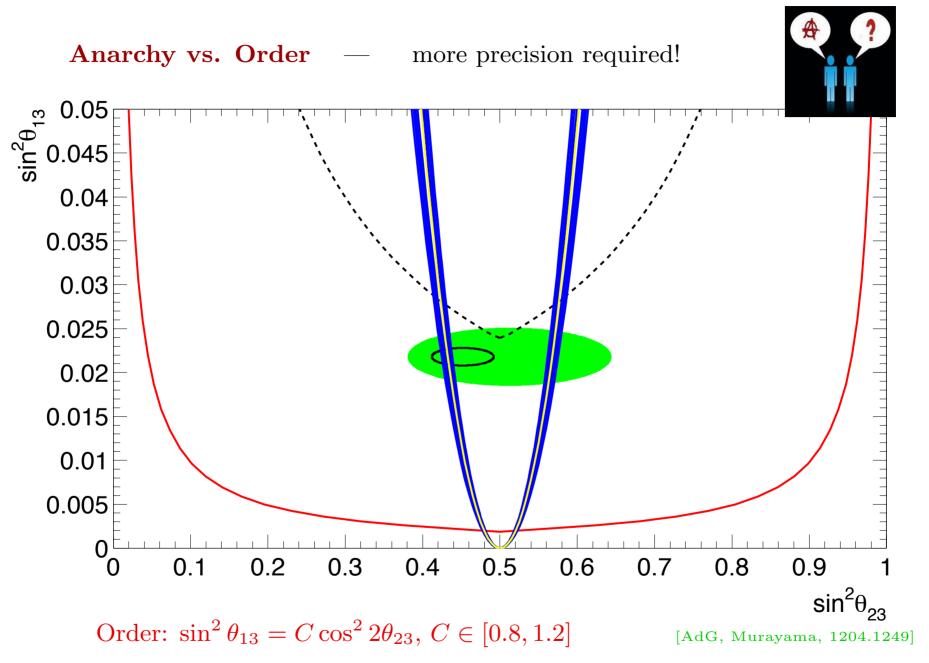
 $[|(V_{MNS})_{e3}| < 0.2]$

They certainly look VERY different, but which one would you label as "strange"?



"Left-Over" Predictions: δ , mass-hierarchy, $\cos 2\theta_{23}$





Long-Baseline Experiments, Present and Future (Not Exhaustive!)

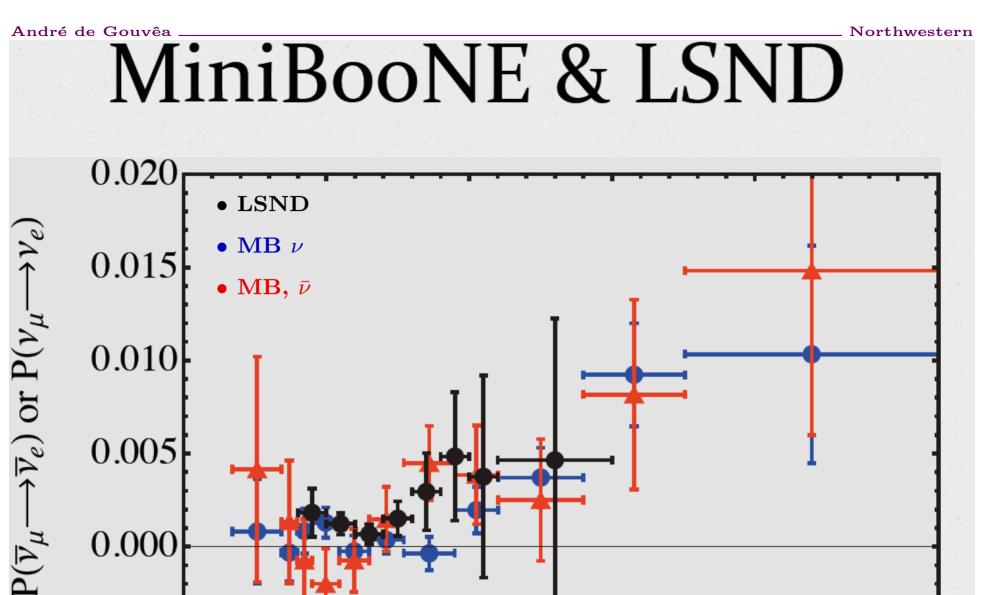
- [NOW] T2K (Japan), NOνA (USA) ν_μ → ν_e appearance, ν_μ disappearance – precision measurements of "atmospheric parameters" (Δm²₁₃, sin² θ₂₃). Pursue mass hierarchy via matter effects. Nontrivial tests of paradigm. First step towards CP-invariance violation.
- [~2020] JUNO (China) $\bar{\nu}_e$ disappearance precision measurements of "solar parameters" (Δm_{12}^2 , $\sin^2 \theta_{12}$). Pursue the mass hierarchy via precision oscillations..
- [~2020] PINGU (South Pole) atmospheric neutrinos pursue mass hierarchy via matter effects.
- [~2025] HyperK (Japan), LBNF (USA) Second (real opportunity for discovery!) step towards CP-invariance violation. More nontrivial tests of the paradigm. Ultimate "super-beam" experiments.
- [>2030(?)] Neutrino Factories Ultimate neutrino oscillation experiment. Test paradigm, precision measurements, solidify CP-violation discovery or improve sensitivity significantly.

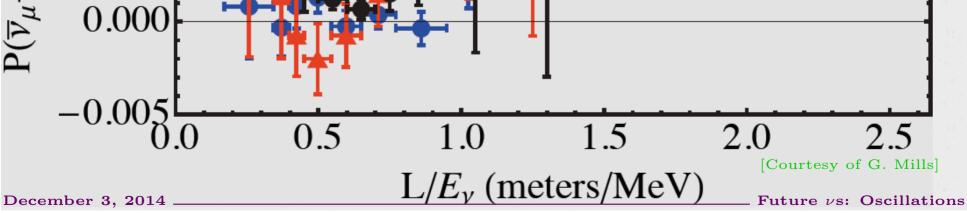
Not all is well(?): The Short Baseline Anomalies

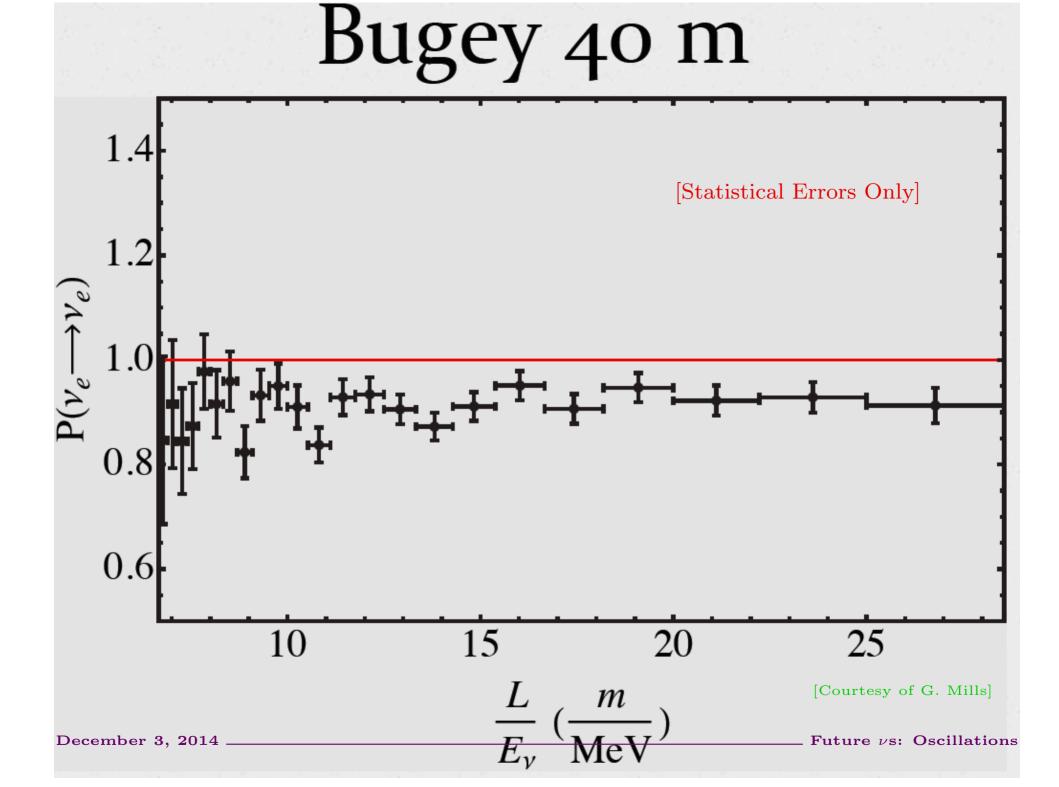
Different data sets, sensitive to L/E values small enough that the known oscillation frequencies do not have "time" to operate, point to unexpected neutrino behavior. These include

- $\nu_{\mu} \rightarrow \nu_{e}$ appearance LSND, MiniBooNE;
- $\nu_e \rightarrow \nu_{other}$ disappearance radioactive sources;
- $\bar{\nu}_e \rightarrow \bar{\nu}_{other}$ disappearance reactor experiments.

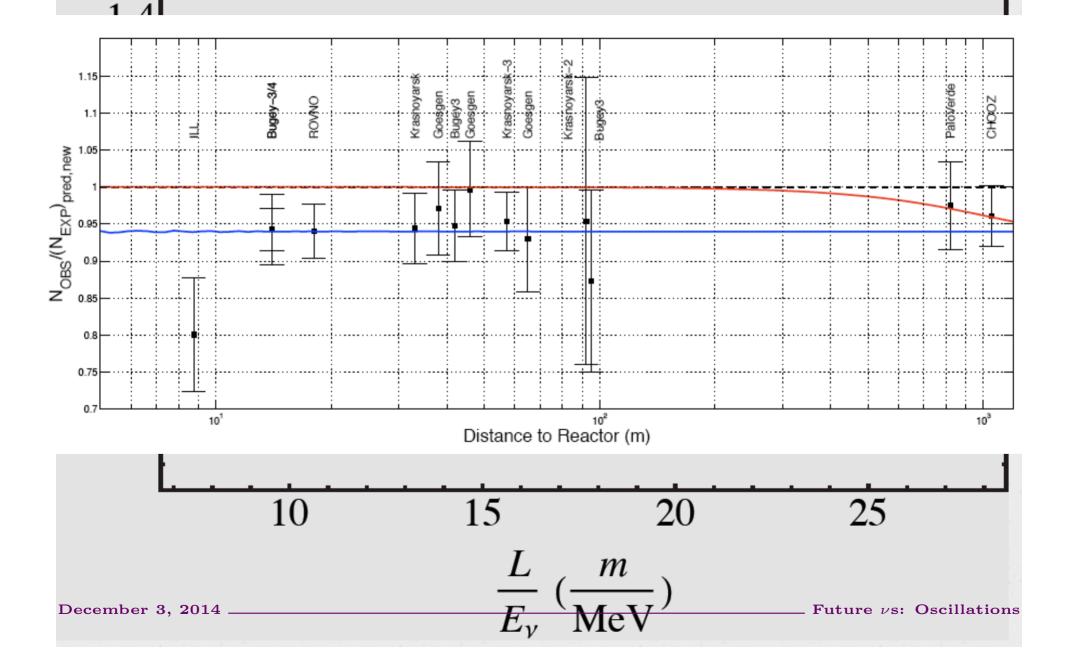
None are entirely convincing, either individually or combined. However, there may be something very very interesting going on here...







Bugey 40 m



What is Going on Here?

- Are these "anomalies" related?
- Is this neutrino oscillations, other new physics, or something else?
- Are these related to the origin of neutrino masses and lepton mixing?
- How do clear this up **definitively**?

Need new clever experiments, of the short-baseline type! [In practice, we address the oscillation interpretation of the anomalies, but we don't know what else to do...]

Observable wish list:

- ν_{μ} disappearance (and antineutrino);
- ν_e disappearance (and antineutrino);
- $\nu_{\mu} \leftrightarrow \nu_{e}$ appearance;
- $\nu_{\mu,e} \rightarrow \nu_{\tau}$ appearance.

If the oscillation interpretation of the short-baseline anomalies turns out to be correct ...

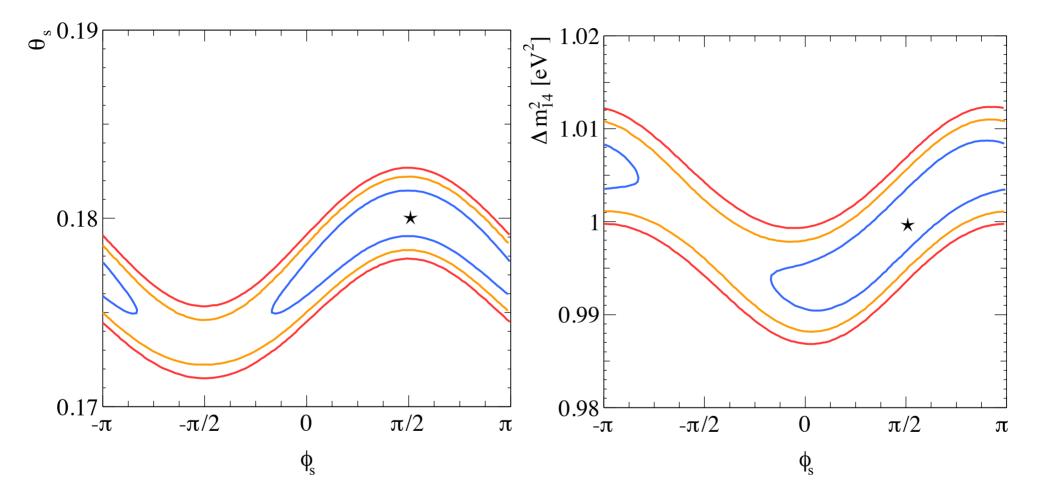
- We would have found new particle(s)!!!!!! [cannot overemphasize this!]
- Lots of Questions! What is it? Who ordered that? Is it related to the origin of neutrino masses? Is it related to dark matter?
- Lots of Work to do! Discovery, beyond reasonable doubt, will be followed by a panacea of new oscillation experiments. If, for example, there were one extra neutrino state the 4 × 4 mixing matrix would require three more mixing angles and three more CP-odd phases. Incredibly challenging. For example, two of the three CP-odd parameters, to zeroth order, can only be "seen" in tau-appearance.

For example, if the new neutrino states are the "right-handed neutrinos" from the standard seesaw, independent from the short-baseline anomalies (for an inverted mass hierarchy, $m_4 = 1 \text{ eV}(\ll m_5)$)...

[AdG, Huang, 1110.6122]

- ν_e disappearance with an associated effective mixing angle $\sin^2 2\vartheta_{ee} > 0.02$. An interesting new proposal to closely expose the Daya Bay detectors to a strong β -emitting source would be sensitive to $\sin^2 2\vartheta_{ee} > 0.04$;
- ν_{μ} disappearance with an associated effective mixing angle $\sin^2 2\vartheta_{\mu\mu} > 0.07$, very close to the most recent MINOS lower bound;
- $\nu_{\mu} \leftrightarrow \nu_{e}$ transitions with an associated effective mixing angle $\sin^{2} \vartheta_{e\mu} > 0.0004;$
- $\nu_{\mu} \leftrightarrow \nu_{\tau}$ transitions with an associated effective mixing angle $\sin^2 \vartheta_{\mu\tau} > 0.001$. A $\nu_{\mu} \rightarrow \nu_{\tau}$ appearance search sensitive to probabilities larger than 0.1% for a mass-squared difference of 1 eV² would definitively rule out $m_4 = 1$ eV if the neutrino mass hierarchy is inverted.

E.g., CPV in 3+1 Scenarios. ν STORM+, $\nu_e \rightarrow \nu_{\mu}$ at the "optimal" baseline...



[AdG, Kelly, Kobach, arXiv, to appear]

The Role of Muon Storage Rings

Qualitatively better beams – better than the pion-decay-in-flight beams – are necessary in order to

- seriously test the standard neutrino paradigm;
- more precisely measurements of oscillation parameters;
- precisely measure neutrino cross-sections (around 1%);
- carry out a serious "sterile neutrinos" short-baseline program.

Very few candidates on the market (muon-storage ring, pion-decay-at-rest). The most versatile one, by far, are the neutrino beams from muon-storage rings. It is hard to imagine a high-precision, thorough, long-term neutrino-oscillation campaign – a la B-Factories – without neutrino beams from muon storage rings.

In Conclusion

The venerable Standard Model sprung a leak in the end of the last century: neutrinos are not massless! (and we are still trying to patch it)

- 1. We know very little about the new physics uncovered by neutrino oscillations.
- 2. neutrino masses are very small we don't know why, but we think it means something important.
- 3. **neutrino mixing is "weird"** we don't know why, but we think it means something important.

- 4. We need more experimental input These will come from a rich, diverse experimental program which relies heavily on the existence of underground facilities capable of hosting large detectors (double-beta decay, precision neutrino oscillations, supernova neutrinos, nucleon decay). Also "required"
 - Powerful, well-characterized neutrino beams;
 - Precision studies of charged-lepton lepton properties and processes;
 - High energy collider experiments (the LHC will do for now).
- 5. There is plenty of **room for surprises**, as neutrinos are potentially very deep probes of all sorts of physical phenomena. Remember that neutrino oscillations are "quantum interference devices" potentially very sensitive to whatever else may be out there.

Backup Slides

Future ν s: Oscillations

High-energy seesaw has no other observable consequences, except, perhaps, ...

Baryogenesis via Leptogenesis

One of the most basic questions we are allowed to ask (with any real hope of getting an answer) is whether the observed baryon asymmetry of the Universe can be obtained from a baryon–antibaryon symmetric initial condition plus well understood dynamics. [Baryogenesis]

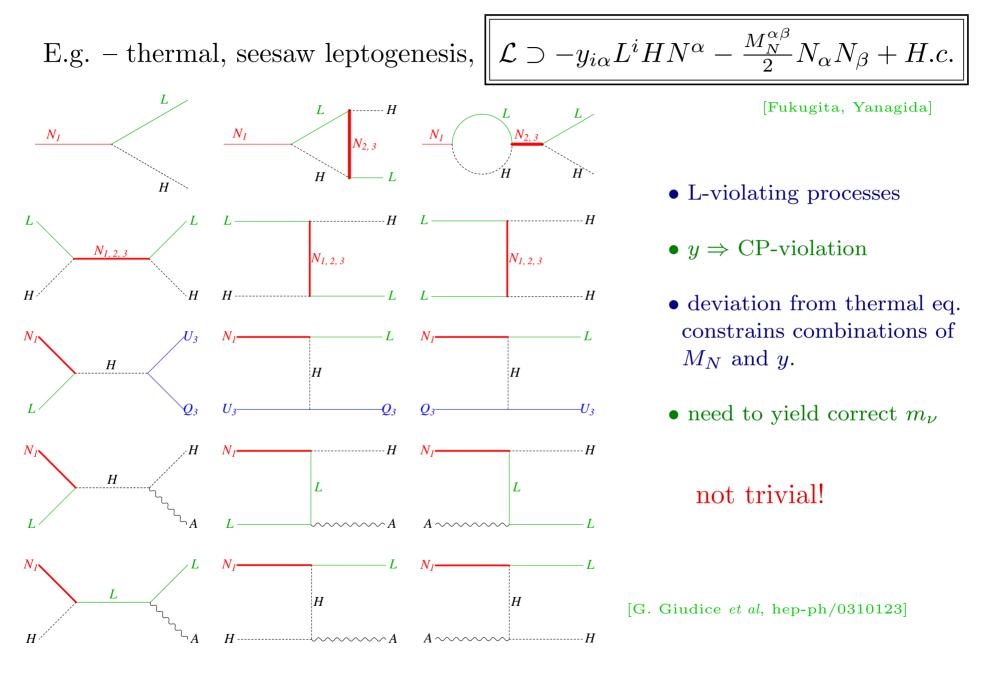
This isn't just for aesthetic reasons. If the early Universe undergoes a period of inflation, baryogenesis is required, as inflation would wipe out any pre-existing baryon asymmetry.

It turns out that massive neutrinos can help solve this puzzle!

In the old SM, (electroweak) baryogenesis does not work – not enough CP-invariance violation, Higgs boson too light.

Neutrinos help by providing all the necessary ingredients for successful baryogenesis via leptogenesis.

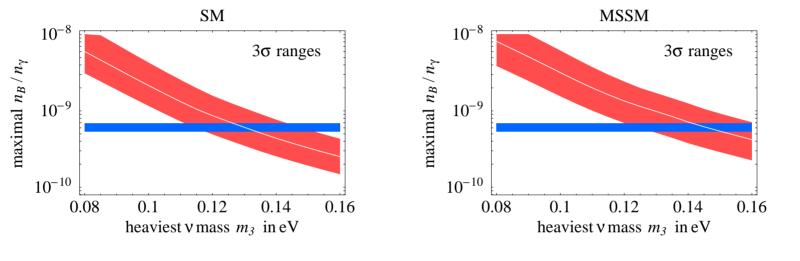
- Violation of lepton number, which later on is transformed into baryon number by nonperturbative, finite temperature electroweak effects (in one version of the ν SM, lepton number is broken at a high energy scale M).
- Violation of C-invariance and CP-invariance (weak interactions, plus new CP-odd phases).
- Deviation from thermal equilibrium (depending on the strength of the relevant interactions).



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Future ν **s**: **Oscillations**

E.g. – thermal, seesaw leptogenesis,
$$\|\mathcal{L} \supset -y_{i\alpha}L^iHN^{\alpha} - \frac{M_N^{\alpha\beta}}{2}N_{\alpha}N_{\beta} + H.c.$$



[G. Giudice et al, hep-ph/0310123]

It did not have to work – but it does MSSM picture does not quite work – gravitino problem (there are ways around it, of course...)

Relationship to Low Energy Observables?

In general ... no. This is very easy to understand. The baryon asymmetry depends on the (high energy) physics responsible for lepton-number violation. Neutrino masses are a (small) consequence of this physics, albeit the only observable one at the low-energy experiments we can perform nowadays.

see-saw: y, M_N have more physical parameters than $m_{\nu} = y^{t} M_N^{-1} y$.

There could be a relationship, but it requires that we know more about the high energy Lagrangian (model depent). The day will come when we have enough evidence to refute leptogenesis (or strongly suspect that it is correct) - but more information of the kind I mentioned earlier is really necessary (charged-lepton flavor violation, collider data on EWSB, lepton-number violation, etc). And that is not all! Neutrinos are unique probes of several different physics phenomena from vastly different scales, including...

- Dark Matter;
- Weak Interactions;
- Nucleons;
- Nuclei;
- the Earth;
- the Sun;
- Supernova explosions;
- The Origin of Ultra-High Energy Cosmic Rays;
- The Universe.



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