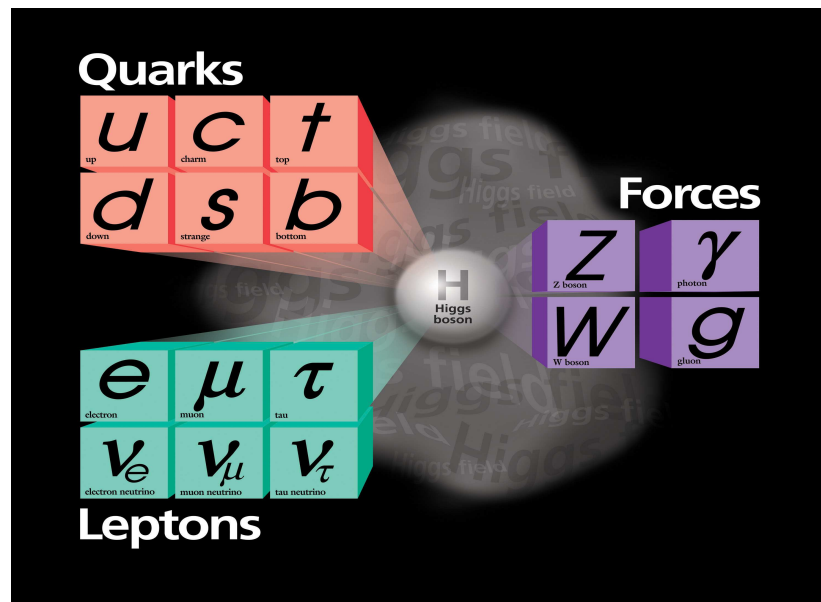


LHC tests of new theories

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+ ?

June 11, 2015 – Fermilab Users Meeting

The LHC is probing the laws of nature at the shortest distances accessible by humans so far.

Run I has confirmed many aspects of the Standard Model, and measured $M_h = 125.09 \pm 0.24 \text{ GeV}$ (ATLAS + CMS, 1503.07589).

We do not know what the LHC Run II will find...

The observed elementary particles are

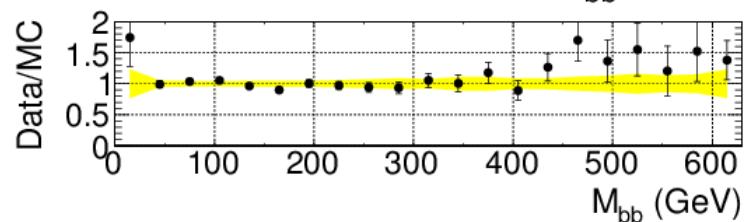
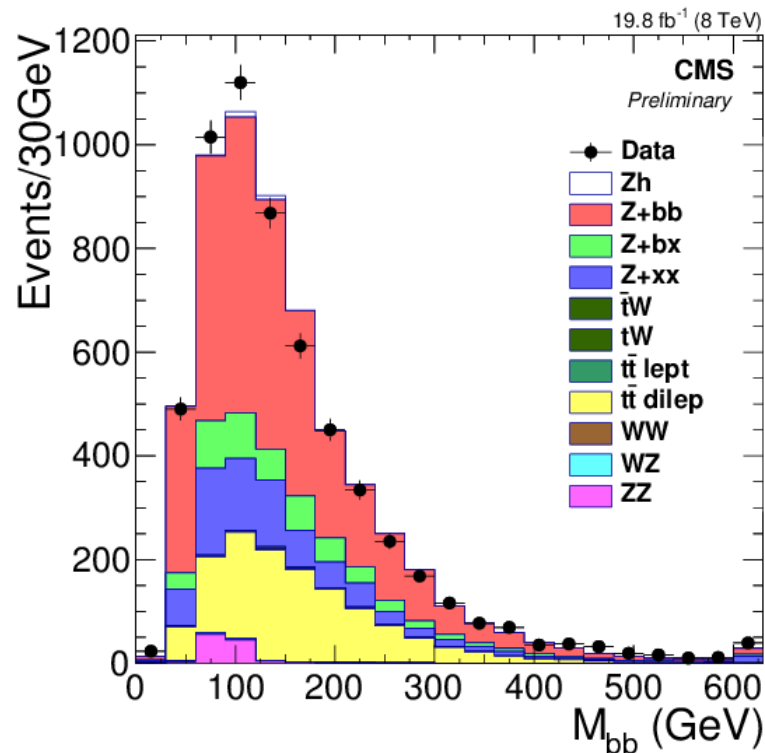
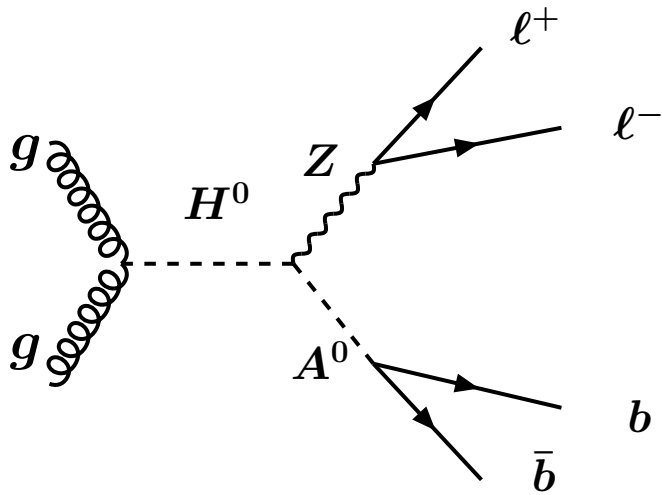
- Scalar (spin 0): Higgs boson
- Chiral fermions (spin 1/2): 6 quarks and 6 leptons (left-handed doublets and right-handed singlets)
- Gauge bosons (spin 1): γ , W , Z , g .

Simple assumption about new particles: heavier versions of the ones discovered so far!?

SM Higgs doublet has 4 degrees of freedom: h^0 , longitudinal Z, W^\pm

A second Higgs doublet – heavier spin-0 particles: H^0, A^0, H^\pm .

(required in the MSSM, composite Higgs models, etc.)



Other processes:

$$pp \rightarrow t\bar{b}H^- \rightarrow t\bar{b}\tau^- \nu, \dots$$

New fermions

All Standard Model fermions are chiral: their masses are not gauge invariant, and arise from the Higgs coupling.

Important legacy of Run I:

a 4th generation of chiral quarks and leptons has been ruled out by direct searches at the LHC.

Example:

$$pp \rightarrow t'\bar{t}' \rightarrow (W^+b)(W^-\bar{b}) \quad \text{with} \quad B(t' \rightarrow Wb) = 100\%$$

new limit: $m_{t'} > 770 \text{ GeV}$ (ATLAS 1505.04306)

already requires a nonperturbative coupling ...

Vectorlike (*i.e.* non-chiral) fermions – a new form of matter.

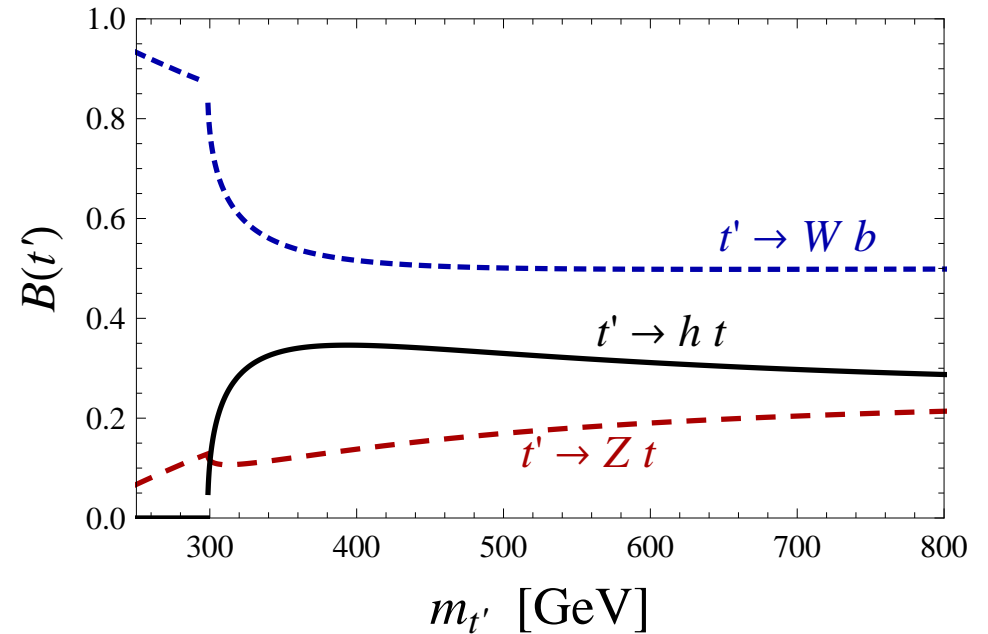
Masses allowed by $SU(3)_c \times SU(2)_W \times U(1)_Y$ gauge symmetry,
 \Rightarrow naturally heavier than the t quark.

Unlike chiral fermions, the vectorlike fermions have
a decoupling limit:

$m \gg v_H \approx 174 \text{ GeV} \longrightarrow$ Standard Model is recovered.

A vectorlike quark that transforms as $(3,1,+2/3)$ under $SU(3)_c \times SU(2)_W \times U(1)_Y$ would mix with the top quark.

Branching fractions of t' :
 ($\sin \theta_L = 0.1$)



For $m_{t'} \gg m_t$:

decays into longitudinal W and $Z \rightarrow B(ht) = B(Zt) = B(Wb)/2$.

QCD production of $t'\bar{t}'$ pairs, followed by t' decays, leads to various final states:

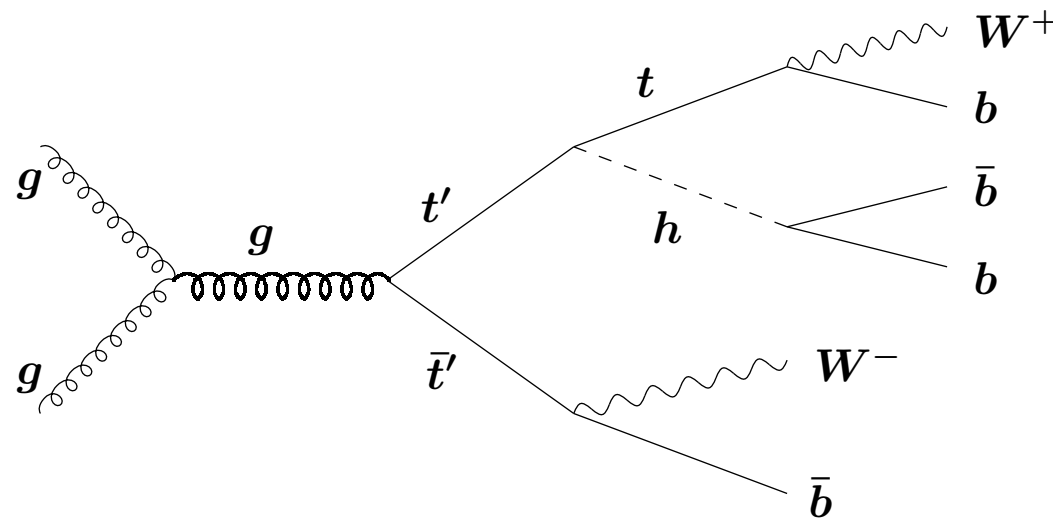
$(W^+b)(W^-\bar{b})$ usual “ t' search”

$(Zt)(W^-\bar{b})$ or $(Z\bar{t})(W^+b)$

$(ht)(W^-\bar{b})$ or $(h\bar{t})(W^+b)$, with $h \rightarrow b\bar{b}$ or $h \rightarrow W^+W^-$

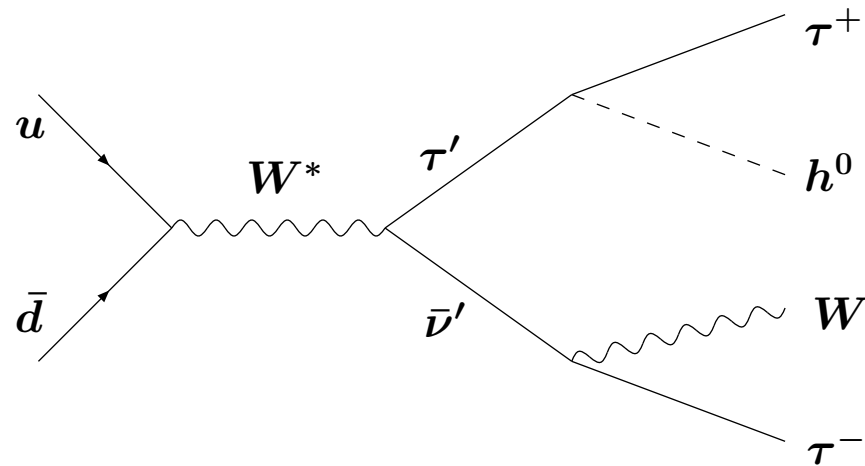
...

Example:



ATLAS and CMS limits on vectorlike quarks: $m \gtrsim 700$ GeV

Vectorlike leptons



Lower limits on vectorlike lepton masses ~ 270 GeV (1312.5329)

(stronger limits if the decays are into e or μ instead of τ)

Run II will be sensitive to significantly higher masses in these and other electroweak processes.

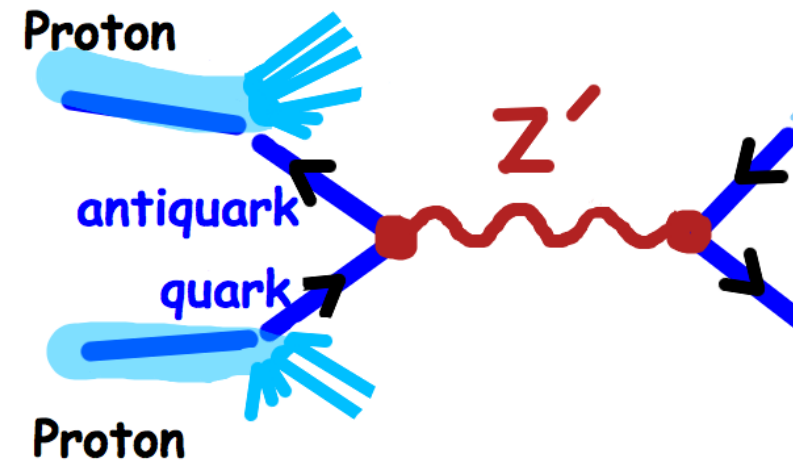
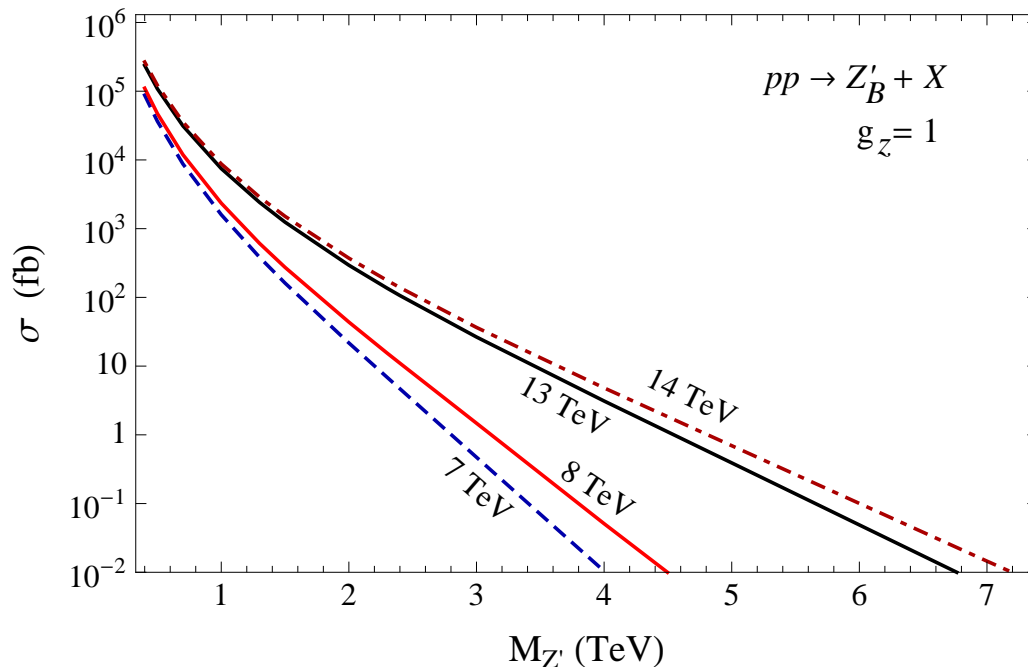
New gauge bosons

Hypothetical heavy particle of spin 1 and charge 0: Z' boson.

If couplings to quarks are flavor-independent:

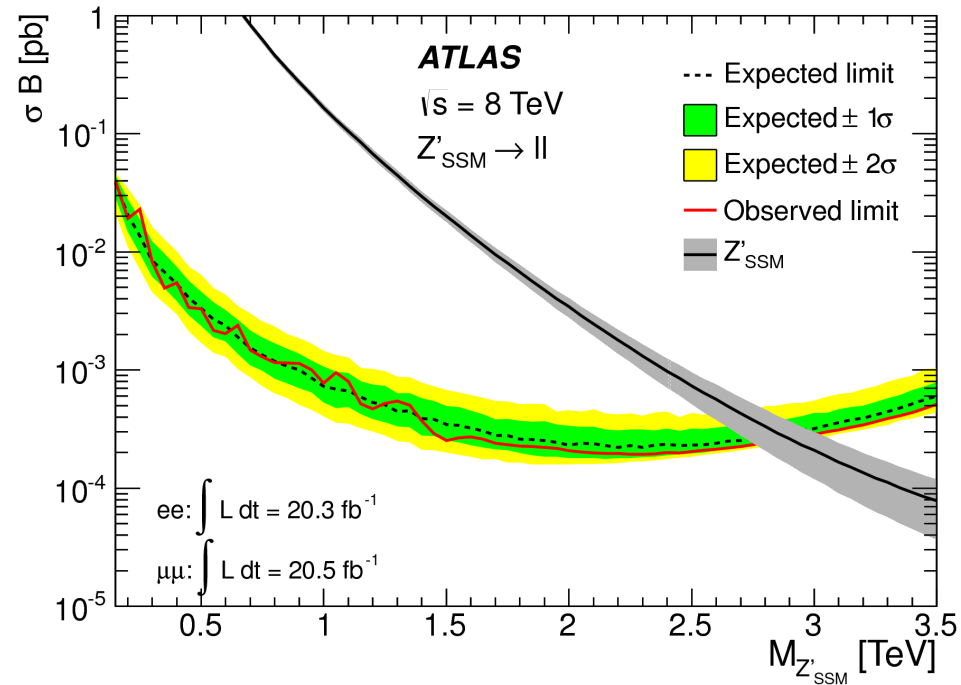
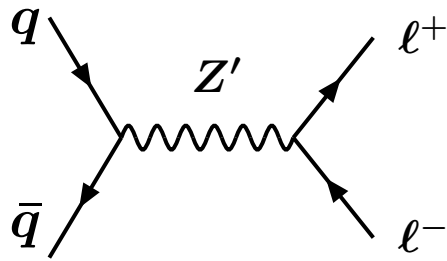
$$\mathcal{L}_q = \frac{g_z}{2} Z'_\mu \sum_q \left(\frac{1}{3} \bar{q}_L \gamma^\mu q_L + \frac{1}{3} \bar{q}_R \gamma^\mu q_R \right)$$

then s -channel production
of Z'_B at the LHC for $g_z = 1$:



Cross section scales as g_z^2
if $g_z \lesssim O(2)$.

If Z' couples to leptons, then stringent limits on g_z from dilepton resonance searches ($g_z < 0.1$ for $M_{Z'} \approx 2$ TeV):



If Z' couples only to quarks (“leptophobic”), then it decays to quark-antiquark pairs \rightarrow limits on g_z are loose (large background)

$g_z < 1.3$ for $M_{Z'} \approx 2$ TeV (Felix Yu, 1306.2629)

Spin-1 fields are well behaved in the UV provided that they are bound states (not discussed here) or gauge bosons.

Z' is associated with a new gauge symmetry.

Simple choice: $SU(3)_c \times SU(2)_W \times U(1)_Y \times U(1)_B$

Theoretical requirements:

- **$U(1)_B$ must be spontaneously broken.**

Simple choice: a new scalar field ϕ acquires a VEV.

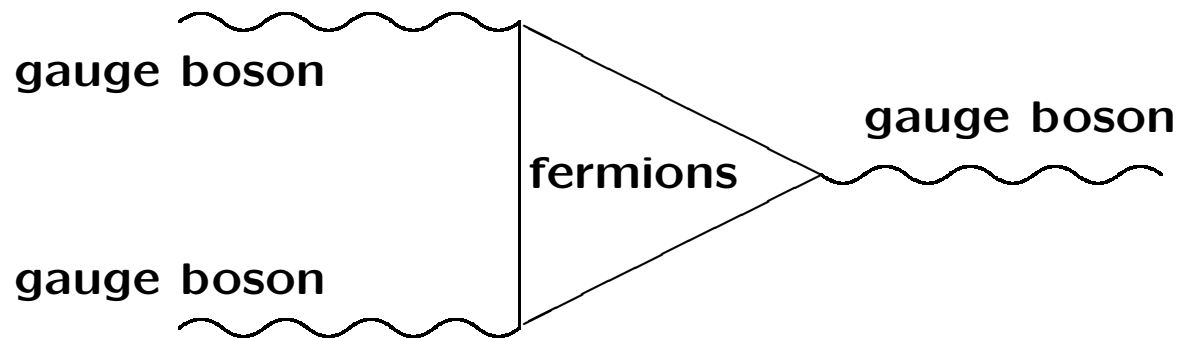
- **All $U(1)_B$ gauge anomalies must cancel.**

Gauge anomaly cancellation

W. Bardeen, 1969, ...

Gauge symmetries may be broken by quantum effects.

Cure: sums over fermion triangle diagrams must vanish.



Standard Model – anomalies cancel within each fermion generation:

$$[SU(3)_c]^2 U(1)_Y: \quad 2(1/6) + (-2/3) + (1/3) = 0$$

$$[SU(2)_W]^2 U(1)_Y: \quad 3(1/6) + (-1/2) = 0$$

$$[U(1)_Y]^3: \quad 3 \left[2(1/6)^3 + (-2/3)^3 + (1/3)^3 \right] + 2(-1/2)^3 + (-1)^3 = 0$$

... (u_L, d_L) u_R d_R (ν_L, e_L) e_R

Any leptophobic Z' that couples to quarks requires new charged fermions to cancel the anomalies.

(or to mix with the SM quarks - Fox et al, 1104.4127)

\Rightarrow The new fermions (“anomalons”) must be vectorlike with respect to $SU(3)_c \times SU(2)_W \times U(1)_Y$, and chiral with respect to the new gauge group.

New fields carrying $U(1)_B$ charge in the minimal leptophobic model:

field	spin	$SU(3)_c$	$SU(2)_W$	$U(1)_Y$	$U(1)_B$
L_L L_R	1/2	1	2	-1/2	-1 +2
E_L E_R	1/2	1	1	-1	+2 -1
N_L N_R	1/2	1	1	0	+2 -1
ϕ	0	1	1	0	+3

Mass mixing of the two charged anomalous (E and L^e) $\rightarrow E_D, E_S$

Two neutral anomalous (N and L^ν), mixing $\rightarrow N_D, N_S$

The decays of the four anomalon physical states depend on their mass ordering.

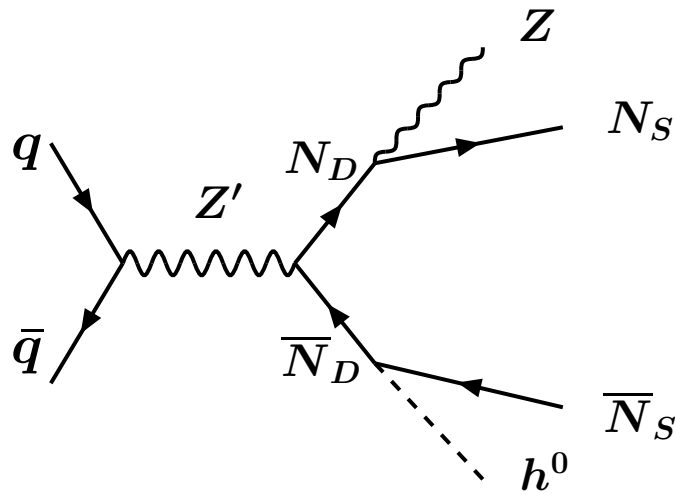
$U(1)_B$ symmetry is spontaneously broken down to a Z_3 subgroup

The anomalons have Z_3 charge $+1$

\Rightarrow lightest anomalon is stable (in the minimal model),
can be a dark matter component if it is N_S .

For $m_{E_S} > m_{E_D} > m_{N_D} > m_{N_S}$:

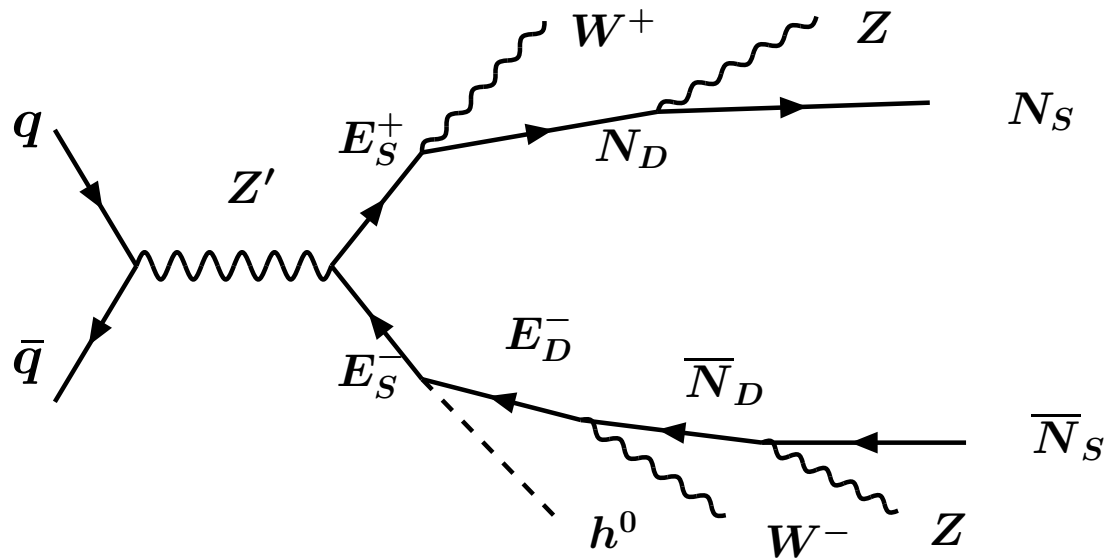
N_D anomalon has 2 decay modes: $N_S h^0$ and $N_S Z$.



E_D anomalon can decay into $N_D W$ or $N_S W$.

E_S has 3 main decay modes: $N_D W$ and $E_D h^0$, and $E_D Z$.

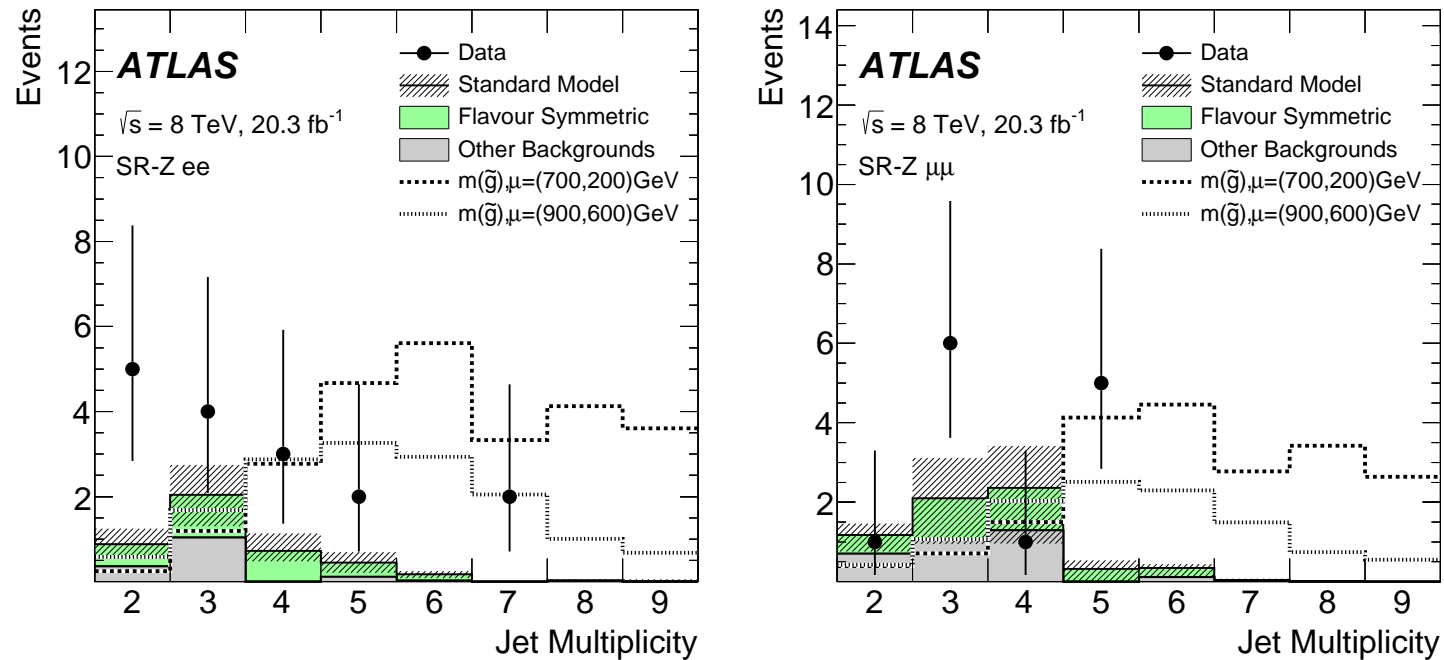
Long cascade decays:



Leptophobic Z' leads to final states with leptons!

Signatures with a $Z \rightarrow \ell^+ \ell^-$, missing energy and jets.

ATLAS 1503.03290 – events with $H_T > 600$ GeV (3σ excess)



Jet multiplicity is more consistent with the Z' interpretation than with supersymmetry.

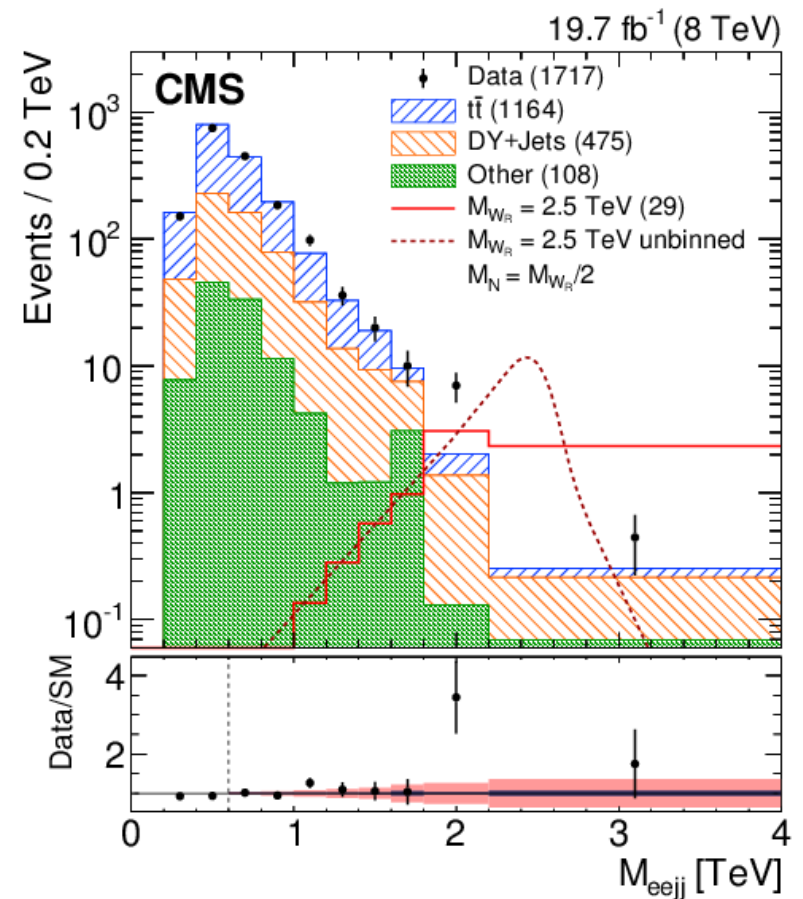
W' boson: hypothetical particle of spin 1 and charge ± 1 , requires an extra $SU(2)$ gauge group.

CMS search for $W'_R \rightarrow eN_R \rightarrow e^+e^-jj$

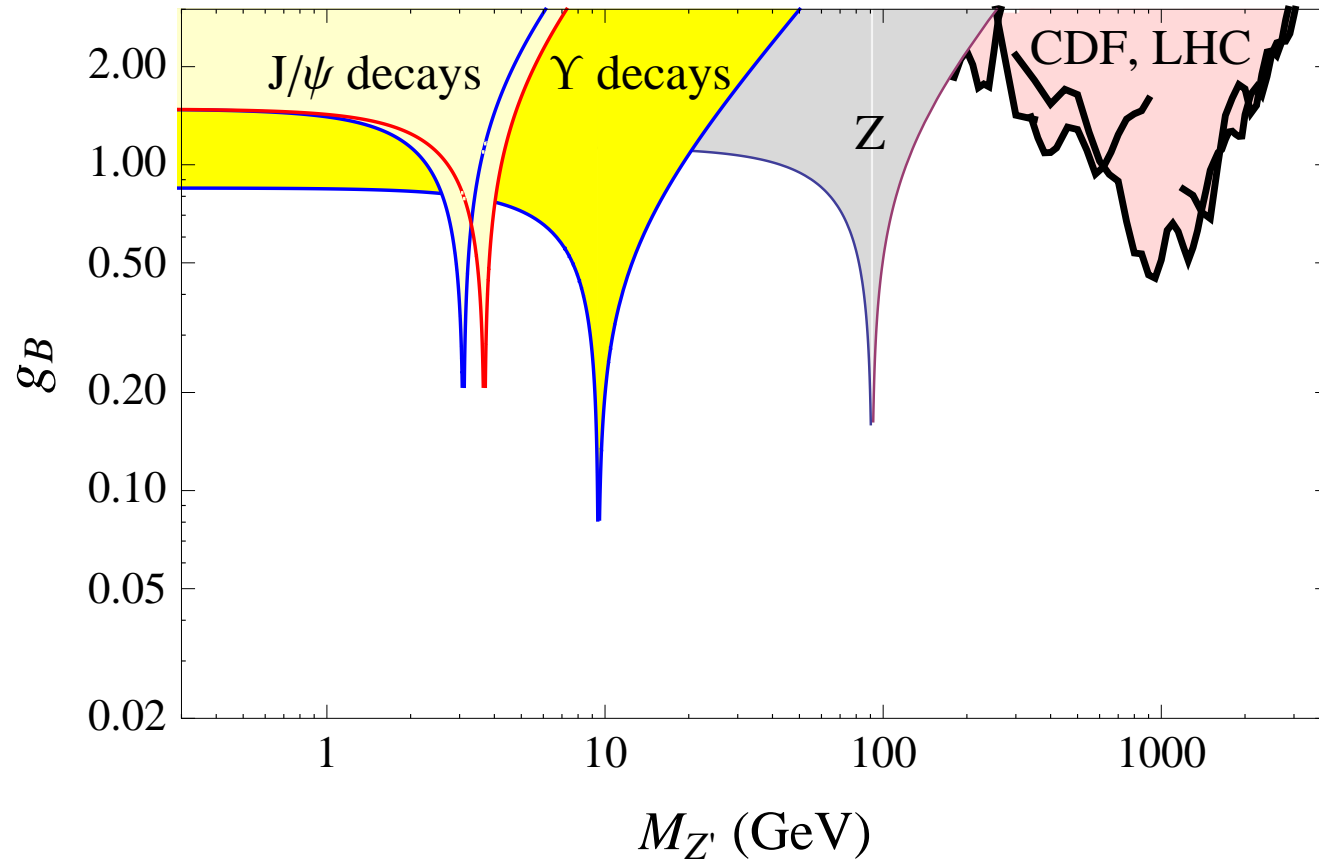
2.8σ local excess in
the 1.8-2.2 TeV bin

(1407.3683)

*Lepton universality violation
is natural in these models.*

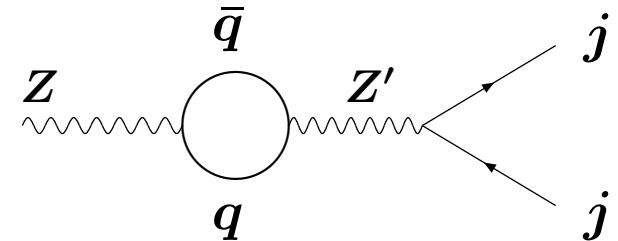


Very light leptophobic Z' bosons



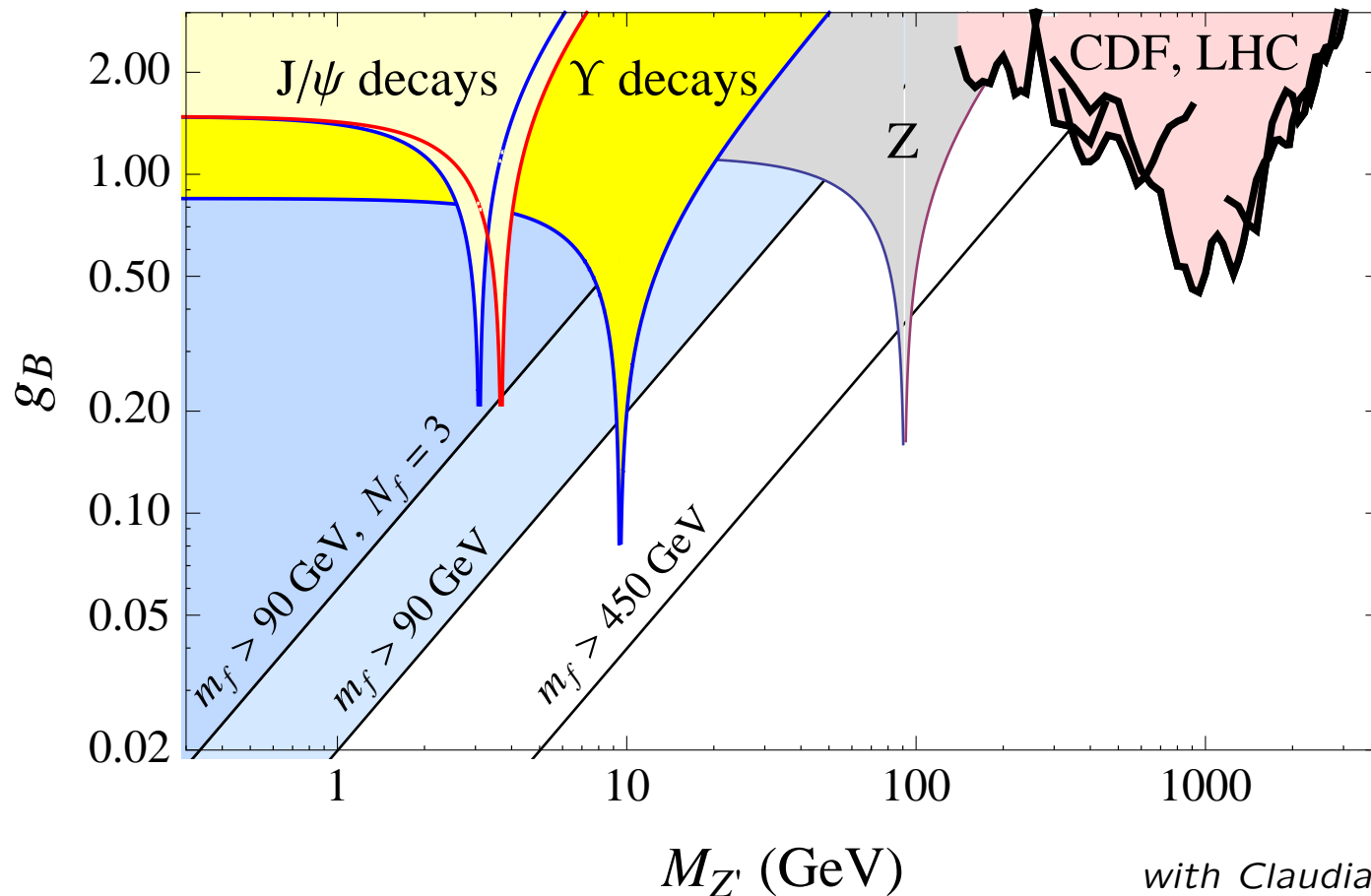
ARGUS experiment, 1986: limit on $\Upsilon \rightarrow Z'^* \rightarrow jj$

Small $Z - Z'$ mixing induced at one loop:



Z' requires a new gauge symmetry: $SU(3)_c \times SU(2)_W \times U(1)_Y \times U(1)_B$

$U(1)_Y [U(1)_B]^2$ anomaly cancelation \Rightarrow charged vectorlike fermions
 \Rightarrow colliders limits on fermion masses lead to a limit on g_B :



with Claudia Fruguele:
PRL 113, 061801 (2014)

Conclusions

- **The Run II of the LHC will push the boundaries of knowledge.**
Synergy with low-energy experiments, including neutrinos, dark matter, flavor processes, ... (many theoretical connections)
- **Heavier “versions” of the SM particles would not be copies:**
 - New gauge bosons can have cascade decays.**
 - New fermions are vectorlike.**
 - New Higgs particles produce different signals.**
- **Much more exotic phenomena are possible ...**