

# SN neutrino theory: lightning overview

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Fermilab



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intensity frontier experiment



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  - Create a ball of matter so dense ( $10^{14}$  g/cm<sup>3</sup>, nuclear densities) that it would be opaque even for neutrinos. Measure its cooling properties as a function of time.



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  - Create a dense neutrino gas ( $10^8$ – $10^{10}$  moles of neutrinos/cm<sup>3</sup>). Let this system expand. Measure the resulting collective flavor oscillation dynamics.



# This experiment is carried out in a core-collapse supernova!

- Inner  $\sim 1.4 M_{\odot}$  of material collapses to a super-dense object a few tens of km across
- 10% of the rest mass of the collapsed core is emitted in  $10^{58}$  neutrinos in a burst lasting  $\delta t \sim$  seconds
  - Neutrino diffusion time scale
- At  $\sim 100$  km, the number density of streaming neutrinos is
  - $\sim 10^{58}/4\pi r^2 c \delta t \sim 10^{32} \text{ cm}^{-3}$
  - Comparable to the number density of matter



# Evolution of the explosion is reflected in neutrinos

- Neutronization burst, accretion and cooling phases can all be seen in neutrinos
- Importantly, different for different progenitor masses

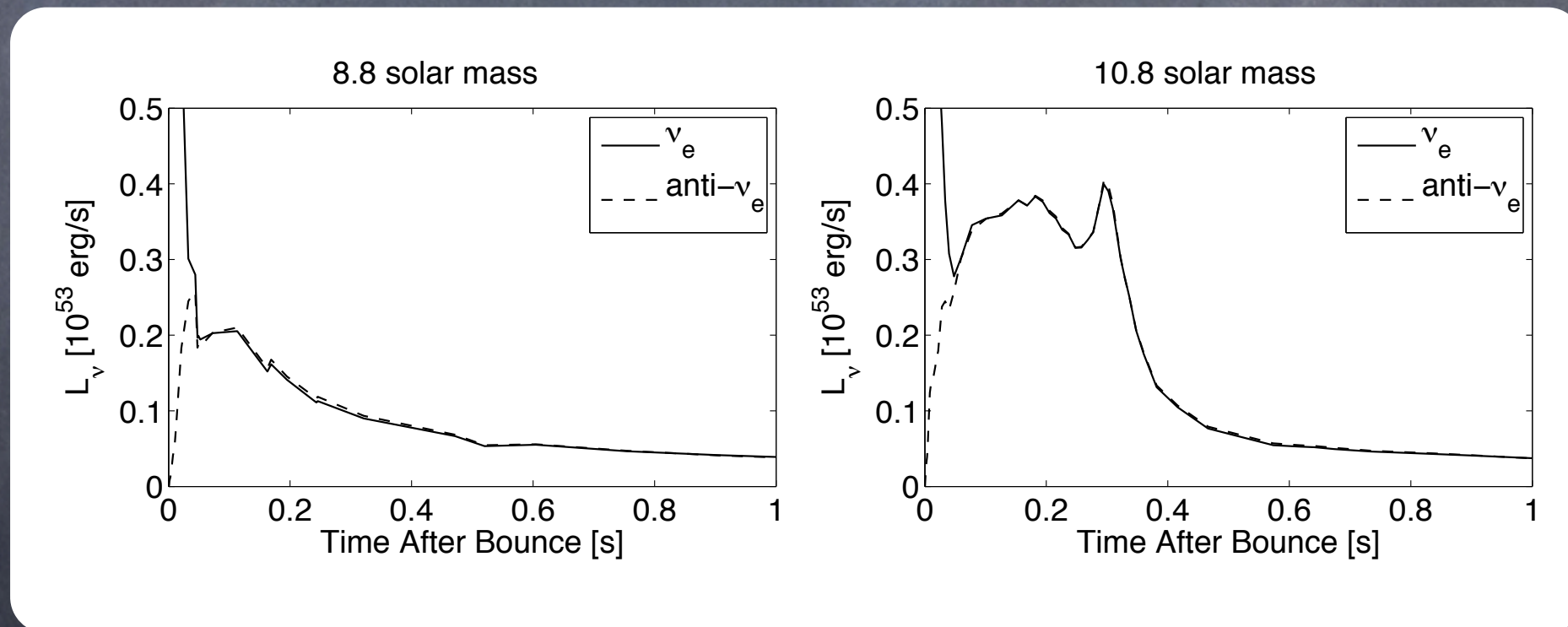


Fig from Fischer, Whitehouse, Mezzacappa, Thielemann, Liebendörfer, arXiv:0908.1871



# Measure each of the phases

- The Neutronization burst provides information about the onset of the explosion, shock breakout through the neutrinosphere; also, a useful sharp time structure
- During the Accretion stage the shock stalls at a few hundred km; we need to know when and how it is reenergized
  - 50-year question in SN theory!
- Cooling stage ends with the formation of a neutron star or a black hole. The signal is sensitive to new physics contributions to cooling (light hidden sector!). Monitor how the shock travels out and the turbulent bubble behind expands.
  - May be possible thanks to neutrino oscillations!



# The richest and most challenging neutrino oscillations problem known

- Possible matter effect in the Earth
- “Solar” MSW in the outer envelope of the progenitor
- “Atmospheric” MSW in the outer envelope of the progenitor
- Turbulent region behind the shock
- Collective oscillations near the neutrino-sphere
- This is schematic, the order of some of these ingredients could be interchanged, depending on the progenitor mass, stage of the explosion



# Earth effect

- The density of the Earth is close to resonant for the “solar” splitting and 20–40 MeV SN neutrinos
  - cf. the D/N effect in  $^8\text{B}$  solar neutrinos is expected at high energies
- Can help to distinguish between different mixing scenarios
- See, e.g.,
  - Smirnov, Spergel & Bahcall, PRD 1994
  - Lunardini & Smirnov, arXiv:hep-ph/0009356
  - Dighe, Kachelriess, Raffelt & Tomas, arXiv:hep-ph/0311172



# Sun: 2-state oscillations

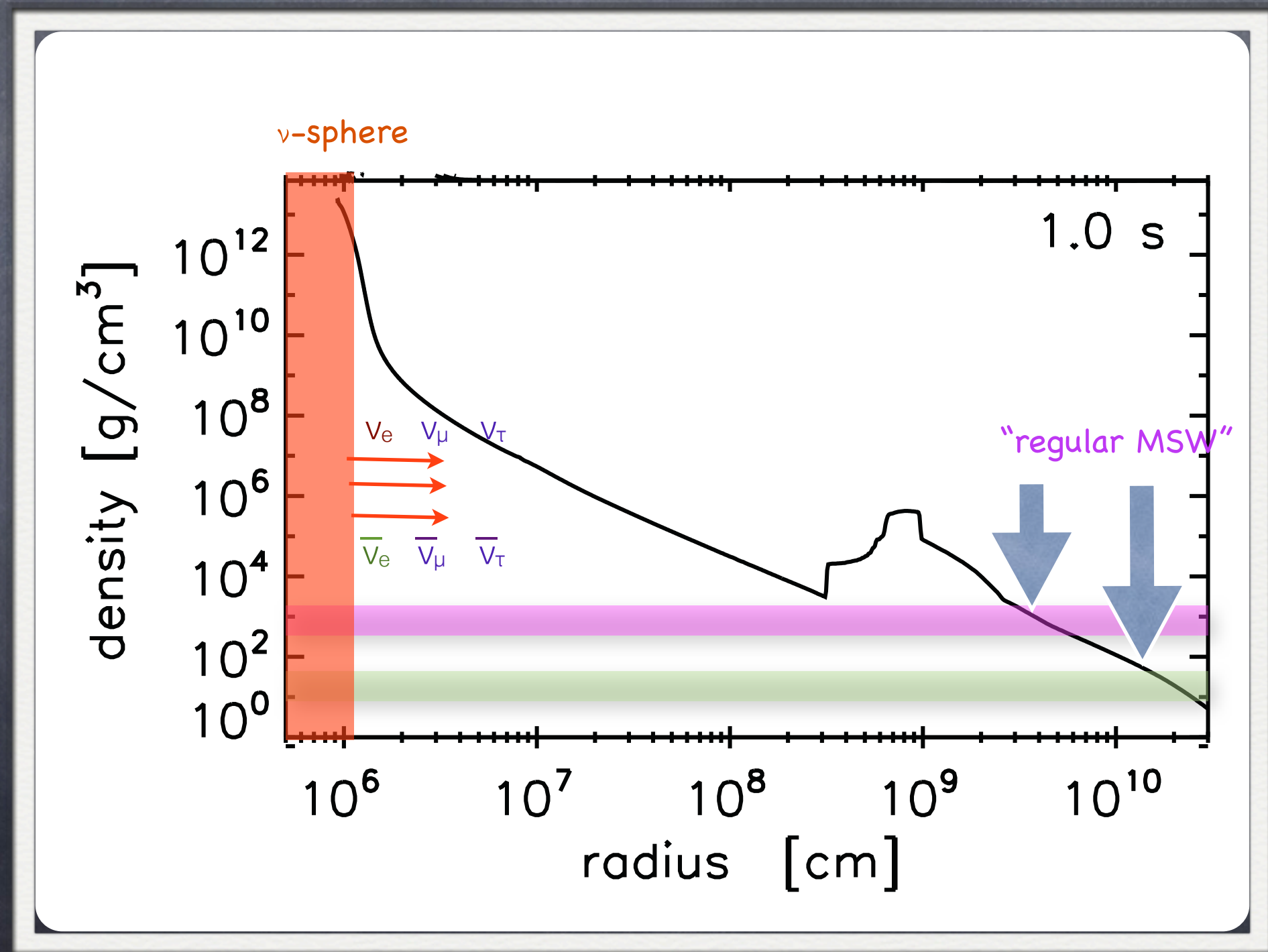
$$P_2(\nu_e \rightarrow \nu_e) = \sin^2 \theta \sin^2 \theta_\odot + \cos^2 \theta \cos^2 \theta_\odot$$



- The evolution is adiabatic (no level jumping), since  $l_{\text{osc}} \ll$  density scale height ( $|d \ln \rho / dr|^{-1}$ )
- Hint: for most of the Sun, the density scale height is  $R_{\text{sun}}/10$ , while  $l_{\text{osc}}$  is comparable to the width of Japan (KamLAND)



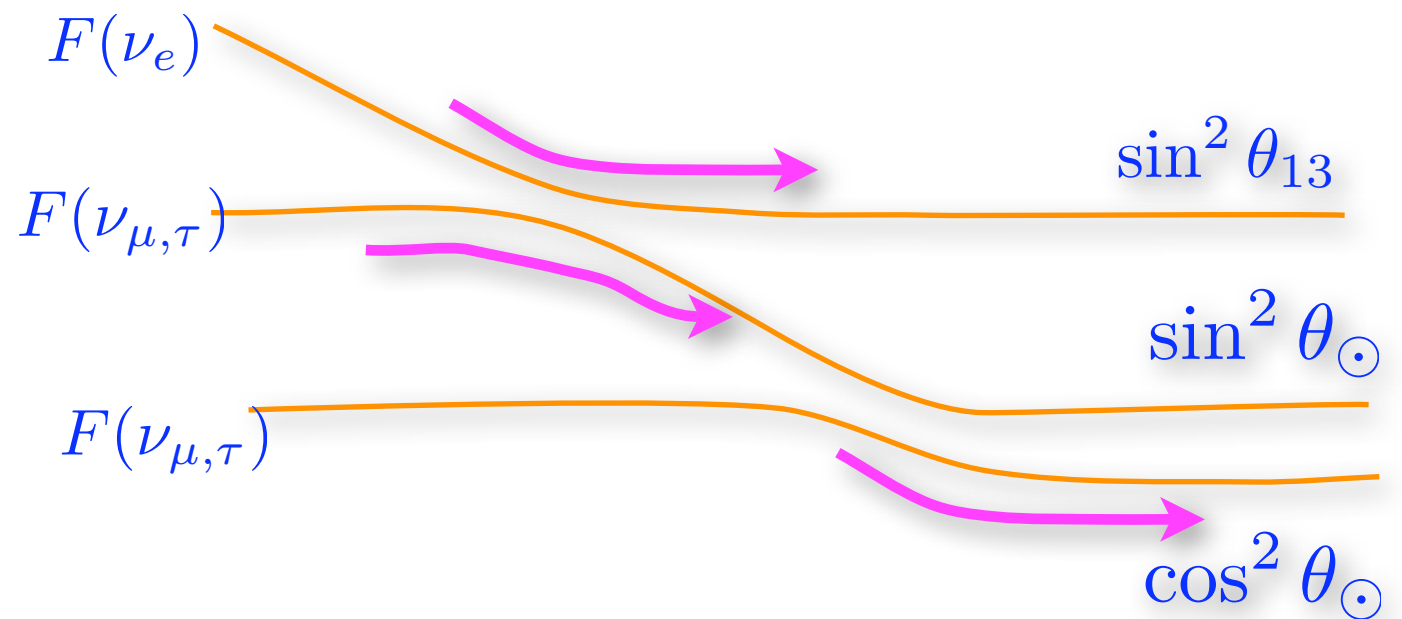
# SN $\nu$ oscillations: 2 MSW densities





# SN MSW transformations, schematics

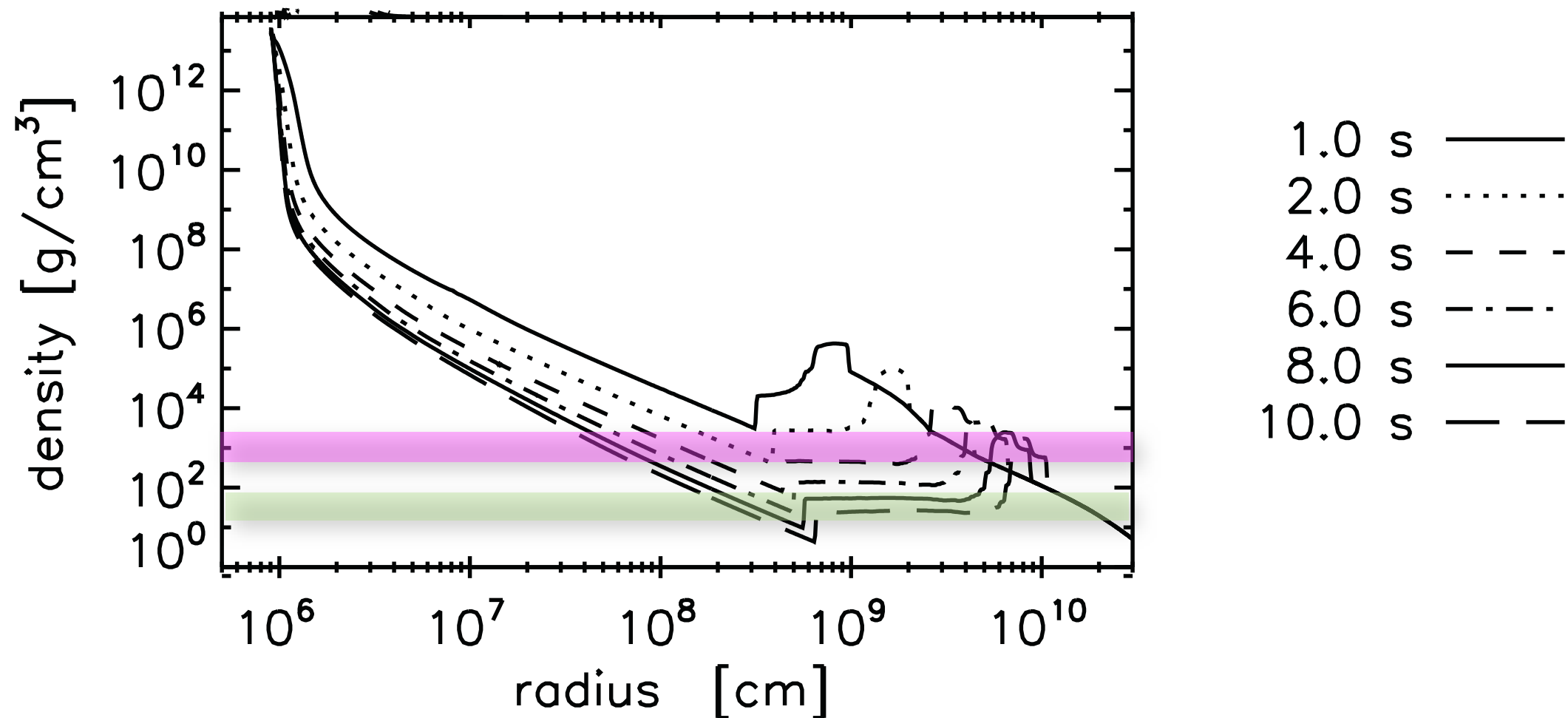
- Given the scale height in the progenitor, the evolution is very adiabatic
  - the adiabaticity of the atmospheric resonance is controlled by  $\theta_{13}$
- Prediction for the  $\nu_e$  signal during the **neutronization burst** is critically dependent on the sign of MH



For inverted hierarchy, the same happens in antineutrinos.



# Dynamical density profile



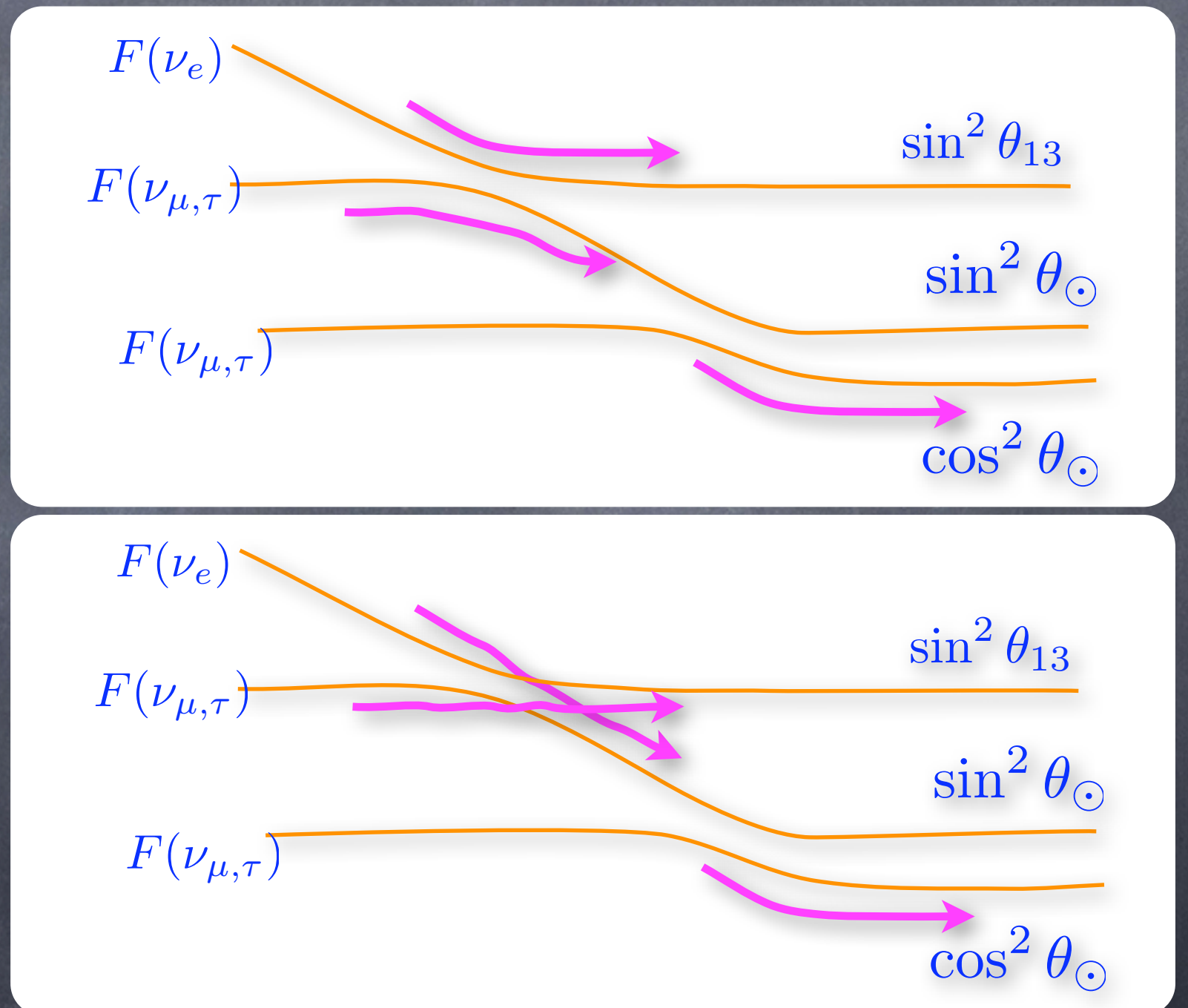
- Front shock reaches the regions where “atmospheric” and “solar” transformations happen, while neutrinos are being emitted

• See Schirato & Fuller (2002) [astro-ph/0205390](#)



# Moving shock and MSW transformations

- The shock is infinitely sharp from the neutrinos' point of view (photon mean free path).
- When it arrives at the resonance, the evolution becomes non-adiabatic.

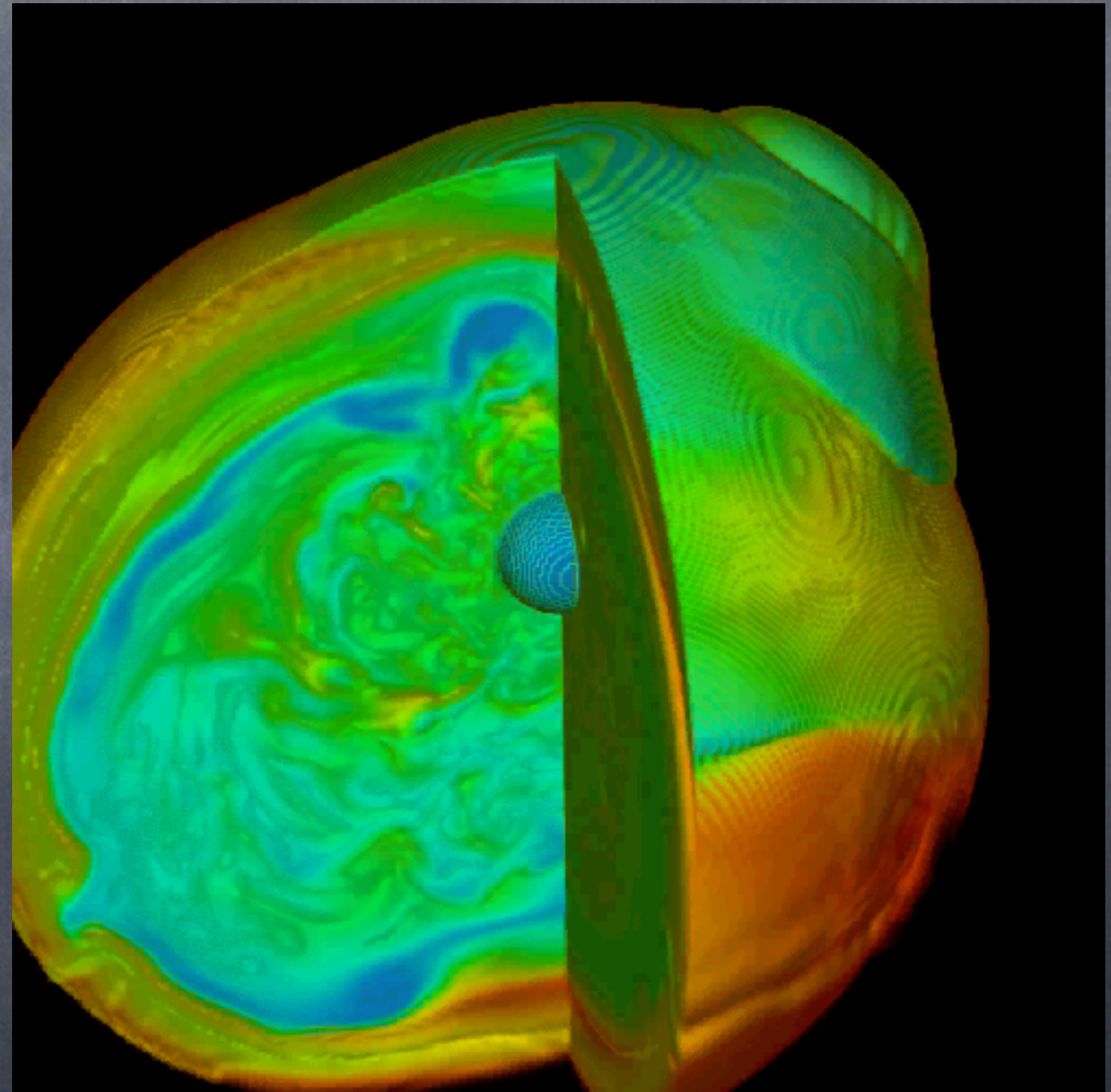


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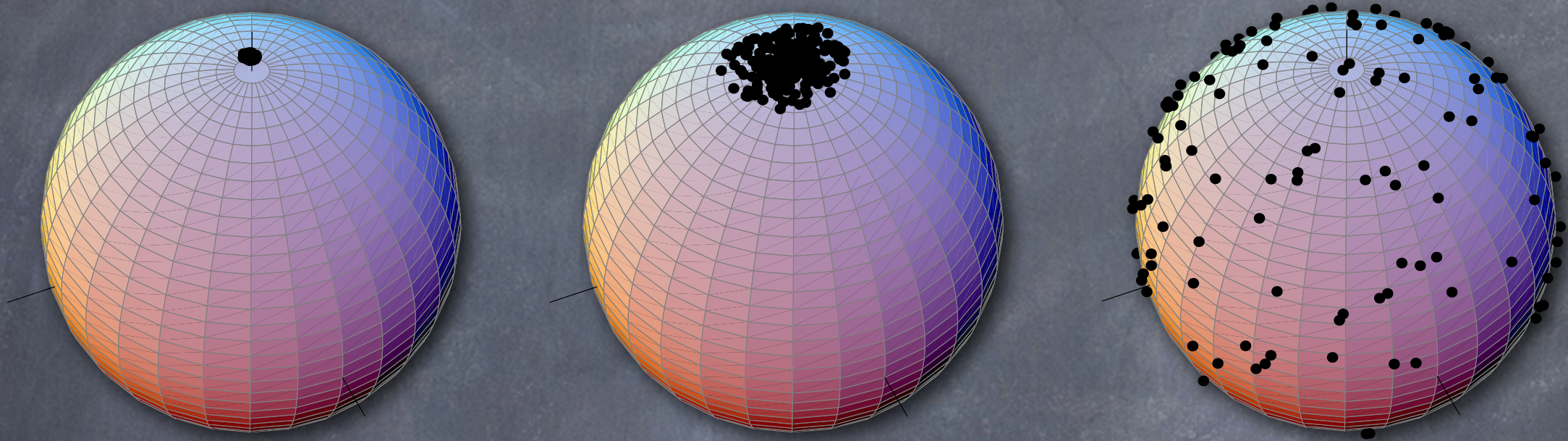
# 3D simulations show turbulence

- 3d simulations of the accretion shock instability  
Blondin, Mezzacappa, & DeMarino (2002)
- See <http://www.phy.ornl.gov/tsi/pages/simulations.html>
- extensive, well-developed turbulence behind the shock





# Turbulence makes neutrinos diffuse in the flavor space



- Need to estimate the rate of diffusion
  - Given large-scale fluctuations in published simulations (order 1) and the large measured value of  $\theta_{13}$ , observable signal expected a few seconds into the explosion



# Some technical details

- The level-jumping probability depends on fluctuations
  - relevant scales are small,  $O(10 \text{ km})$
  - take large-scale fluctuations from simulations, scale down with a Kolmogorov-like power law
  - contributions of different scales to the level-jumping probability are given by the following spectral integral

$$P \simeq \frac{G_F}{\sqrt{2}n'_0} \int dk C(k) G\left(\frac{k}{2\Delta \sin 2\theta}\right), \quad G(p) \simeq \frac{\Theta(p-1)}{p\sqrt{p^2-1}}.$$

for details, see Friedland & Gruzinov, [astro-ph/0607244](#)



# Neutrino “self-refraction”

- Neutrinos undergo flavor conversion in the background of other neutrinos
- The neutrino induced contribution depends on the flavor states of the background neutrinos

$$\sqrt{2}G_F \sum_{\vec{p}} n_i (1 - \cos \Theta_{\vec{p}\vec{q}}) |\psi_{\vec{p}}\rangle \langle \psi_{\vec{p}}|$$

- One has to evolve the neutrino ensemble as a whole
- Rich many-body physics, with many regimes

Fuller et al, Notzold & Raffelt 1988;  
Pantaleone 1992; ...  
Duan, Fuller, Qian, Carlson, 2006;  
+ hundreds more

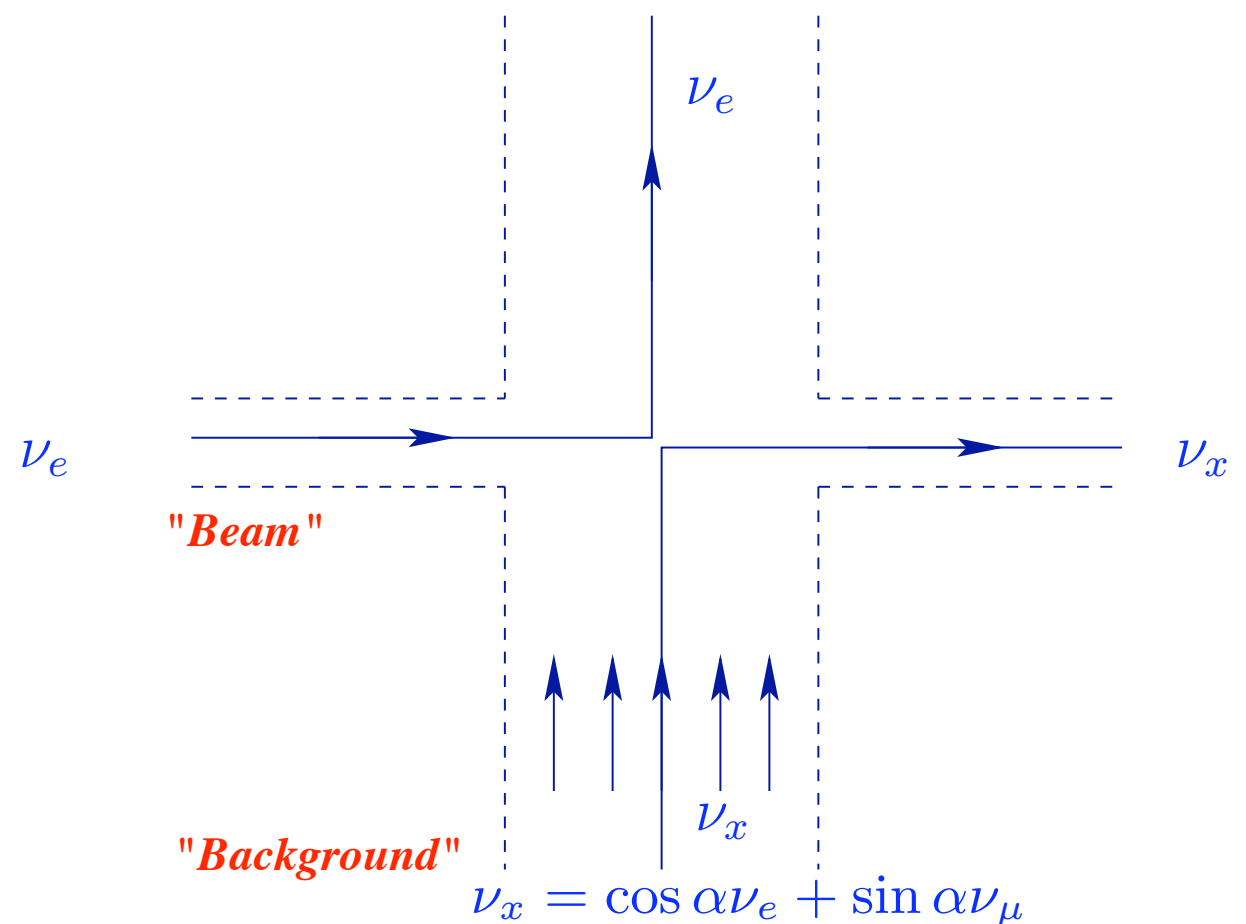
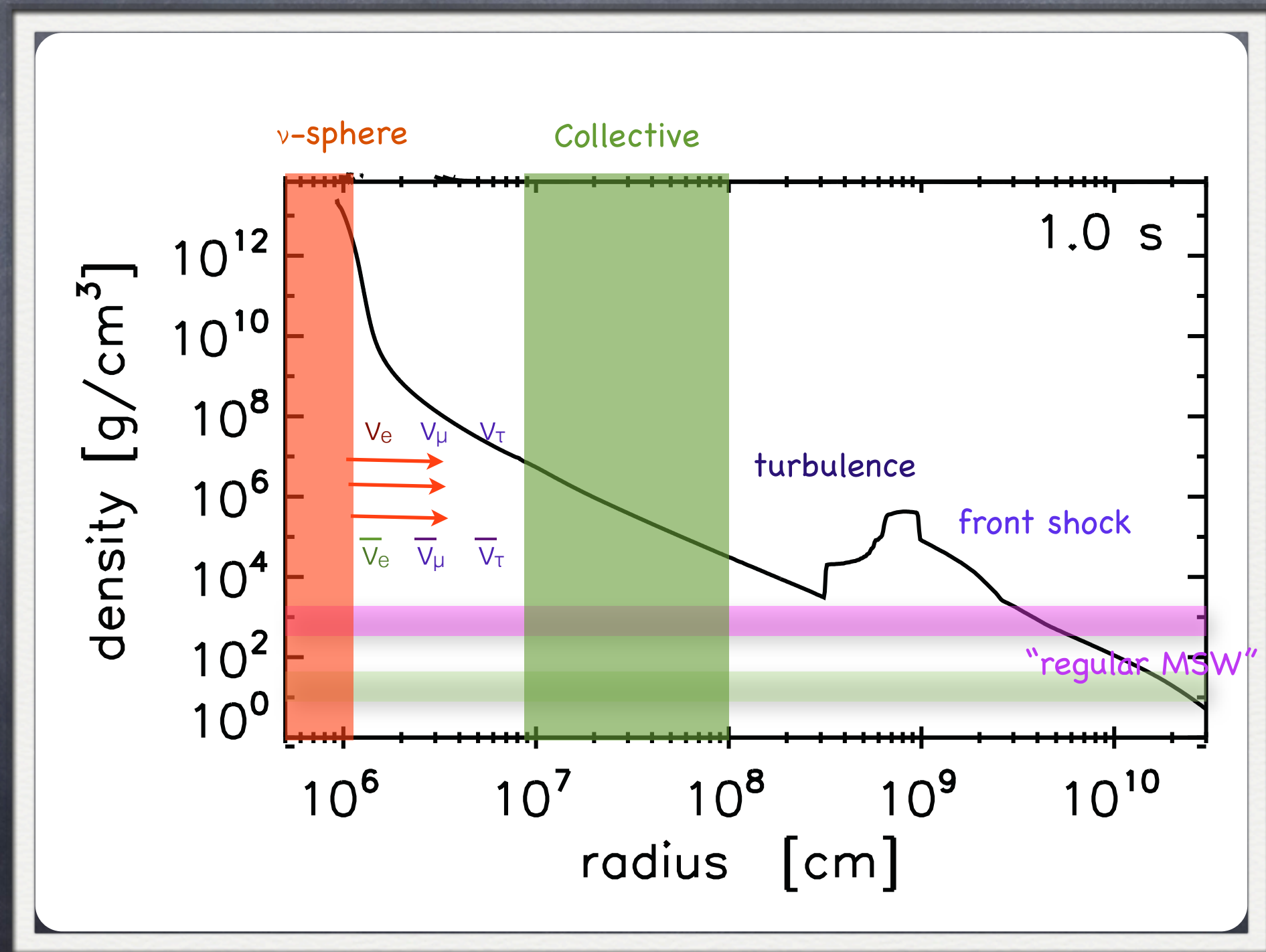


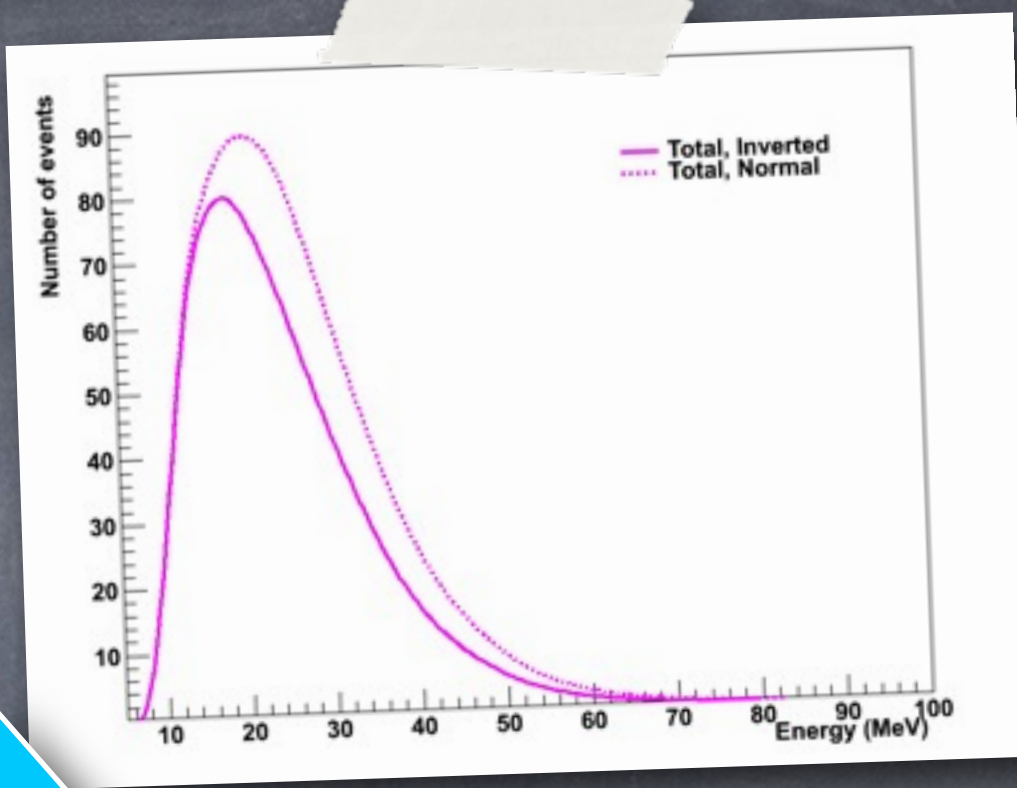
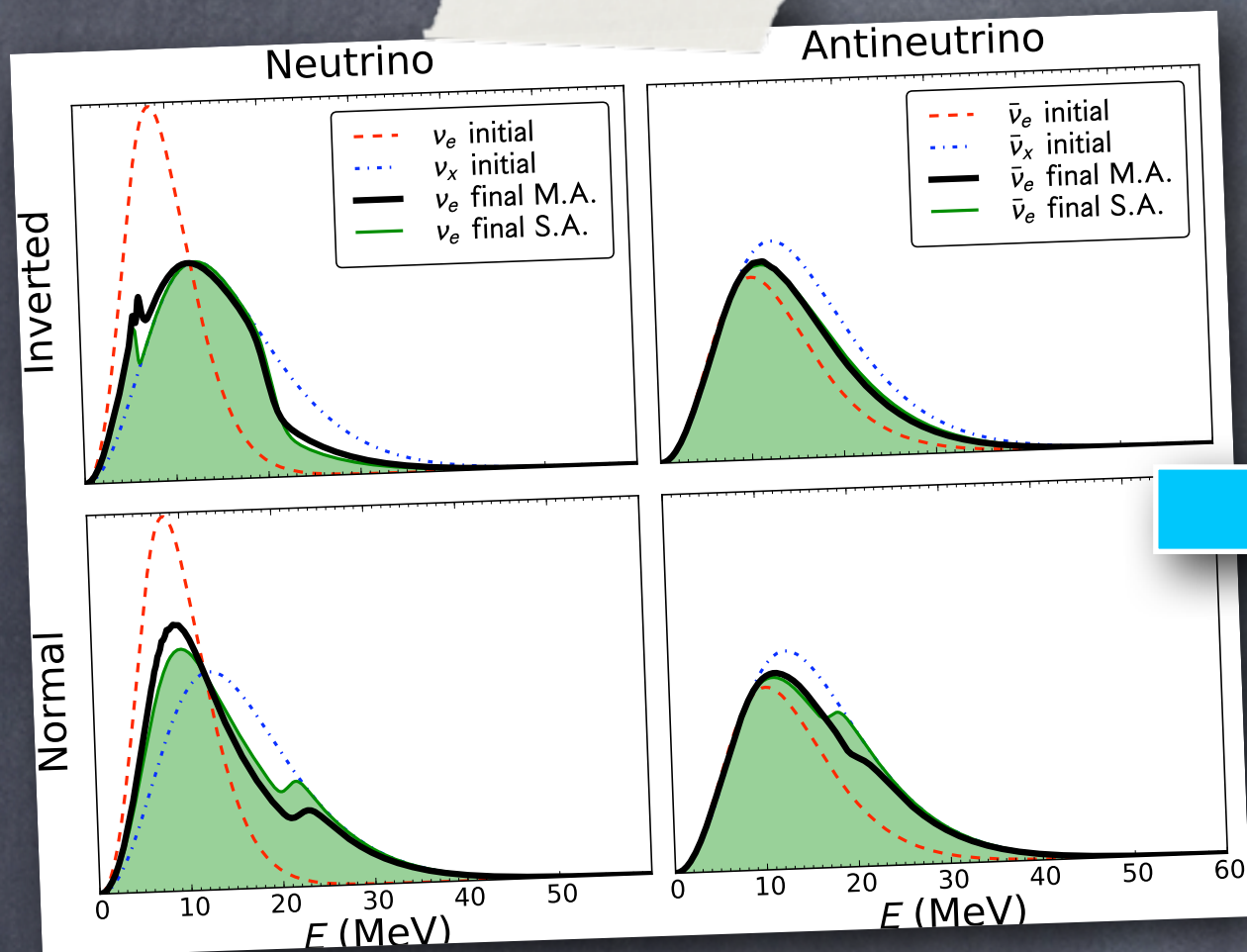
Figure from  
Friedland & Lunardini,  
Phys. Rev. D 68, 013007 (2003)



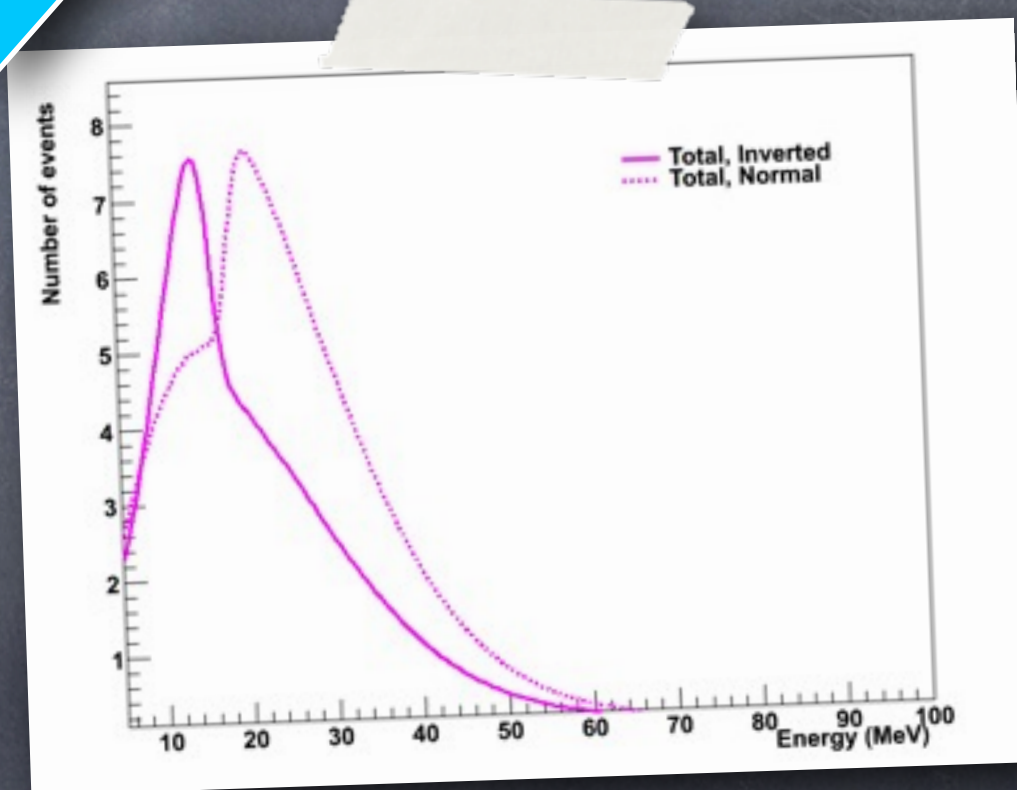
# SN $\nu$ : summary physics cartoon







WC



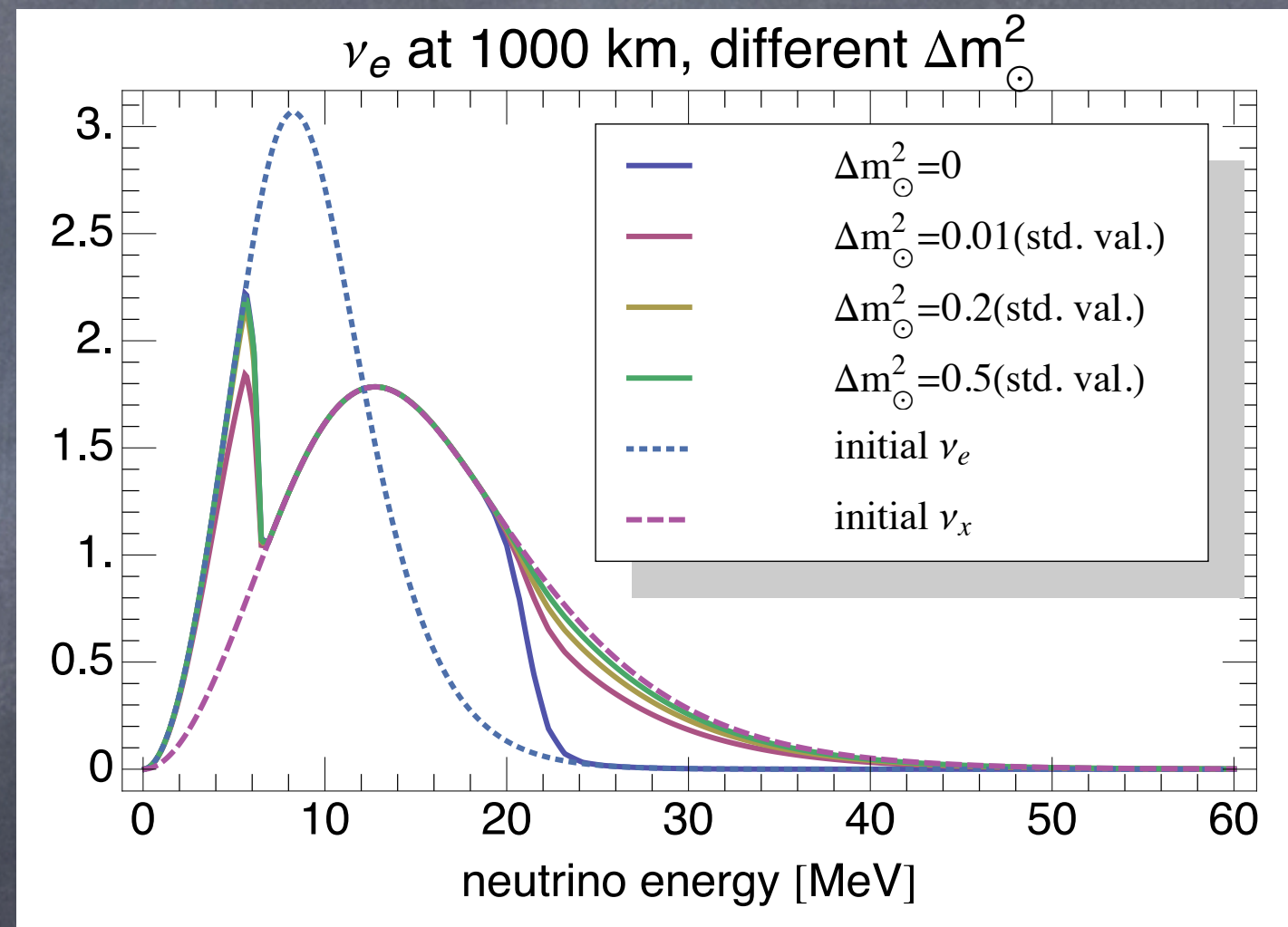
LAr

- \* spectra by Duan & Friedland, PRL 2011
- \* detector modeling by Kate Scholberg & co



# Collective oscillations must be done in 3 flavors

- Example where the solar mass splitting is turned on gradually
  - At  $\Delta m_{\odot}^2=0$ , 2-flavor result is reproduced
  - As soon as  $\Delta m_{\odot}^2 \neq 0$ , the answer is closer to the realistic  $\Delta m_{\odot}^2$  than to  $\Delta m_{\odot}^2=0$
- 2-flavor trajectory can be unstable in the 3-flavor space



For details, see Friedland, PRL (2010);  
also Dasgupta, Dighe, Raffelt, Smirnov, PRL (2009)

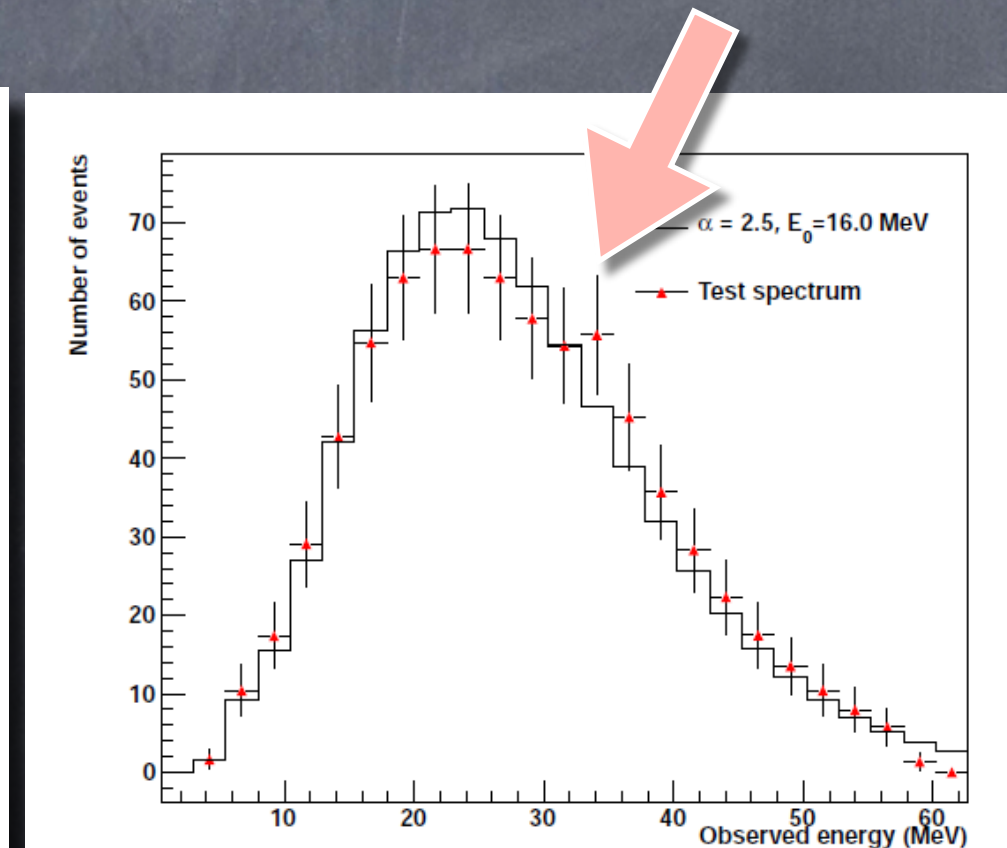
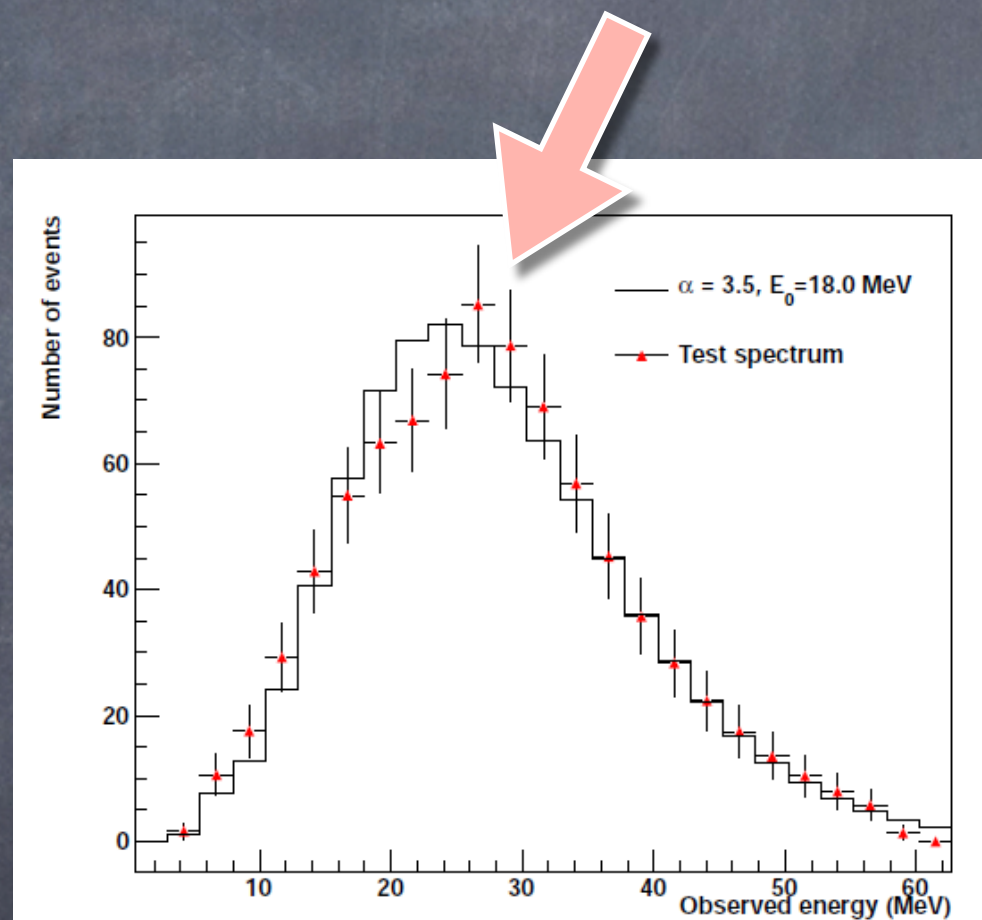


# What are we looking for?

## Smoking-gun features

Modeling  
multiangle  
collective +  
moving  
shock  
by A. F.

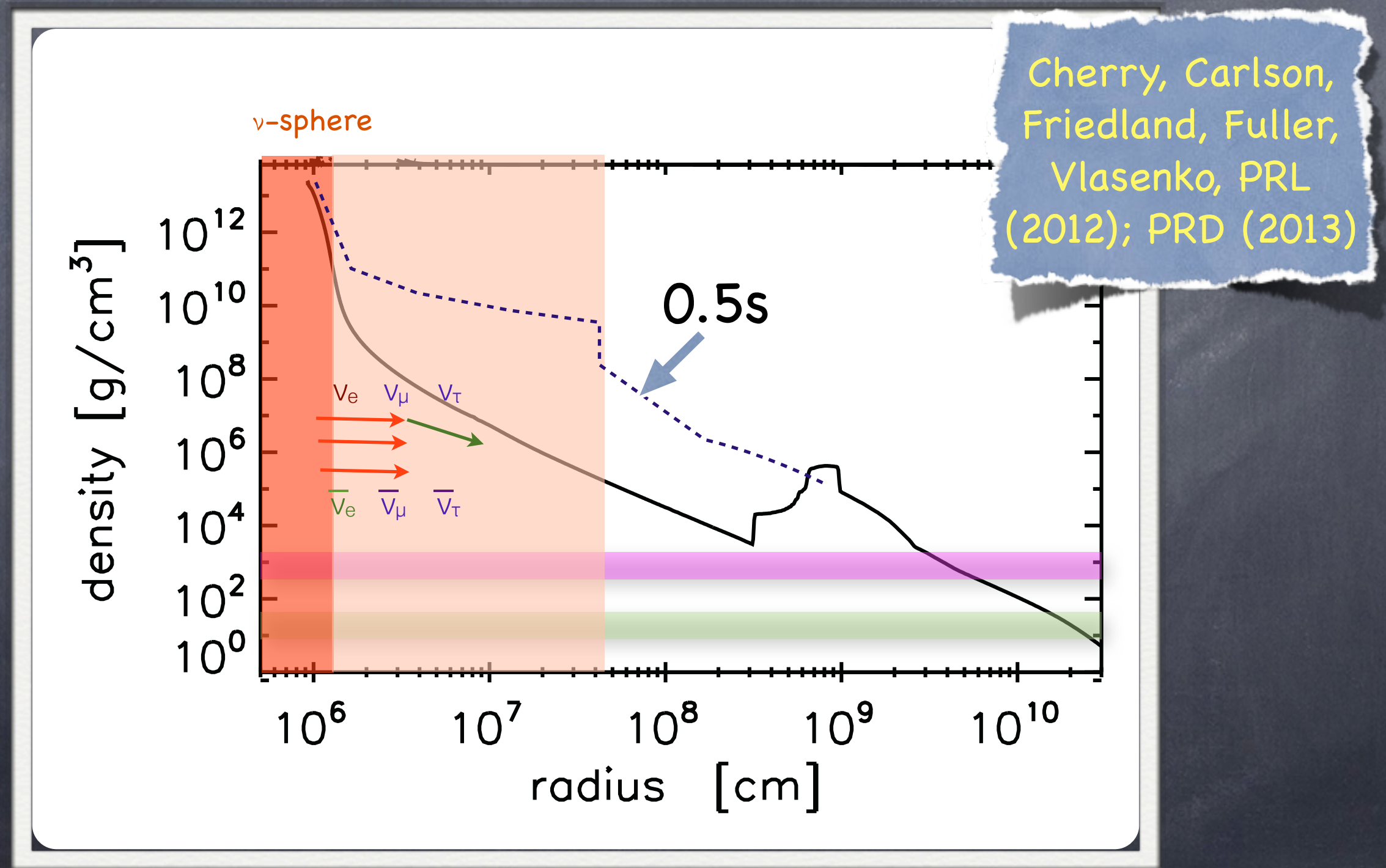
Detector  
model by K.  
Scholberg



- The neutrino spectrum is modulated, but not antineutrinos (simultaneously observed by SK/HK)



# Accretion phase: neutrinos scattering above $\nu$ -sphere?





# Much work is still to be done!

- The role of matter in collective oscillations
  - Do they always factorize?
- Dependence of collective transformations on luminosities and temperatures of different components
  - Transition from sharp spectral splits to decoherence
- Breaking of spherical symmetry
  - e.g., Raffelt, Sarikas de Sousa Seixas, PRL 111, 091101 (2013)
- Effects of nonstandard physics
  - e.g., de Gouvea and Shalgar, JCAP (2012, 2013)