

A Global Jet Finding Algorithm



Yang Bai

University of Wisconsin-Madison

MC4BSM Workshop@LPC

May 20, 2015

In collaboration with:



Zhenyu Han

University of Oregon



Ran Lu

Univ. of Wisconsin-Madison

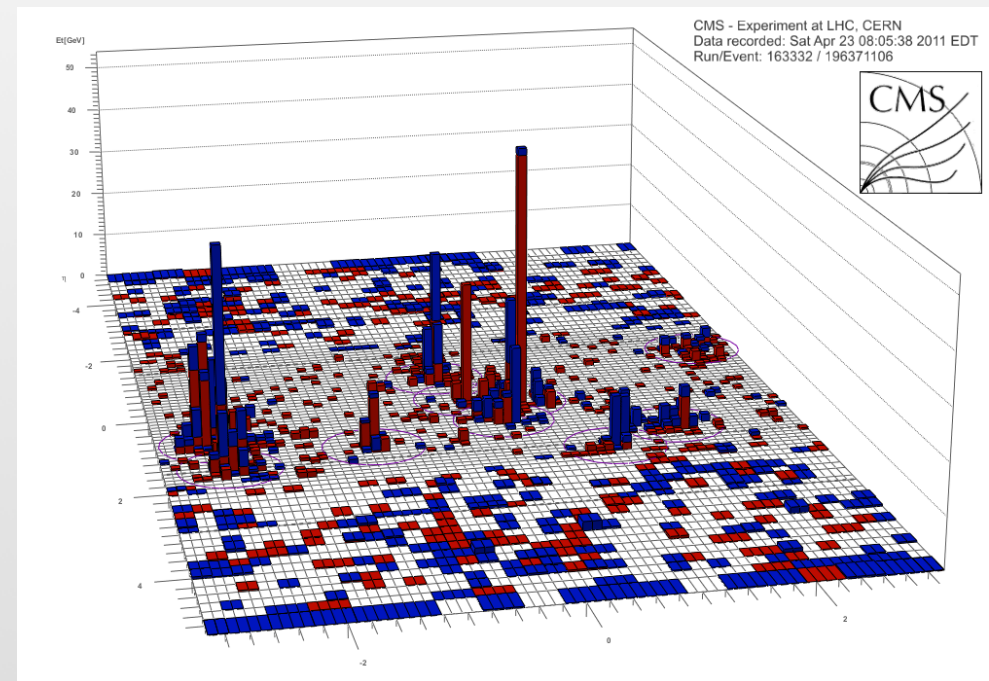
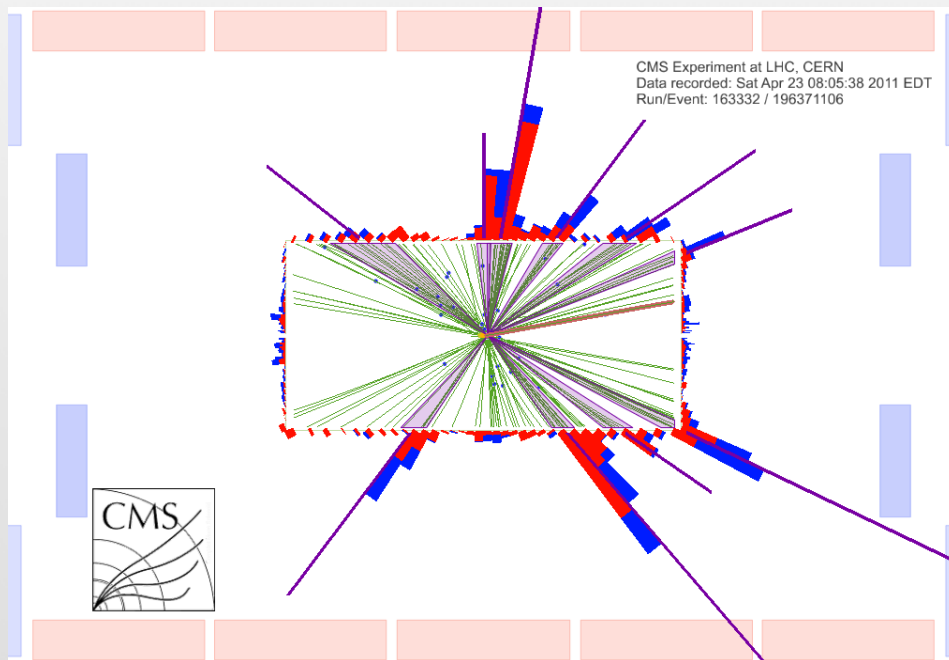
arxiv:1411.3705 and work in progress

Outline

- Motivation: physics beyond the Standard Model
- Review of jet finding algorithms
- Introduction of a global definition
- Application to QCD-like jets
- Extension to boosted two-prong jets
- Conclusion

What is a jet?

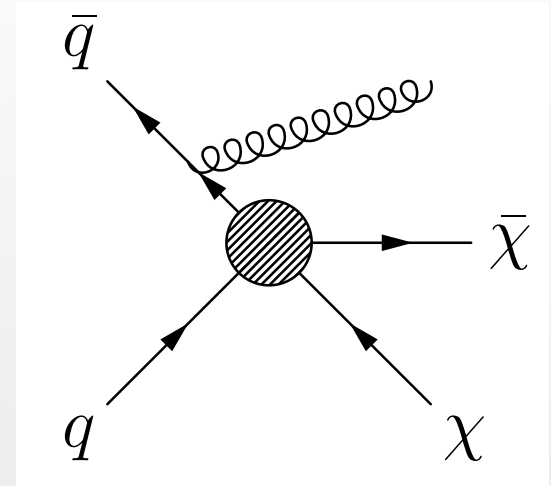
An ensemble of particles in detectors can be called a jet



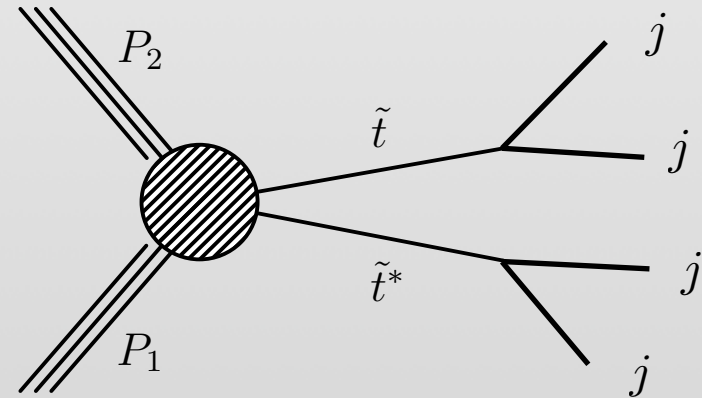
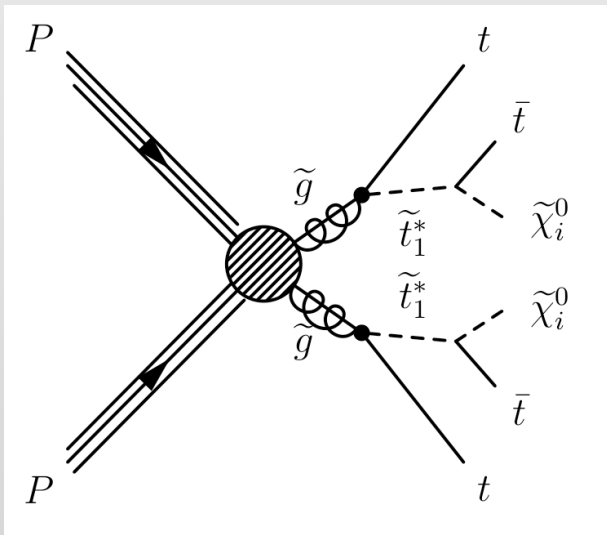
Jet-finding algorithm: how to group particles together?

Jets in BSM

- Mono-jet plus MET events as the dark matter signature



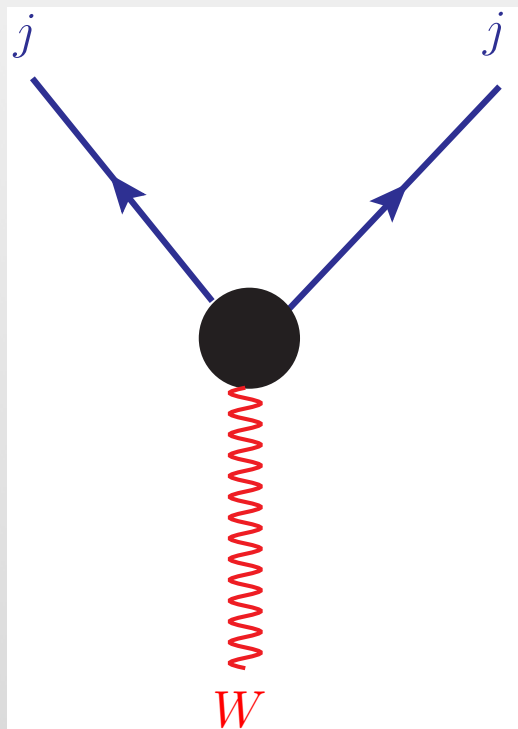
- Multi-jets plus MET for RPC SUSY or without MET for RPV SUSY



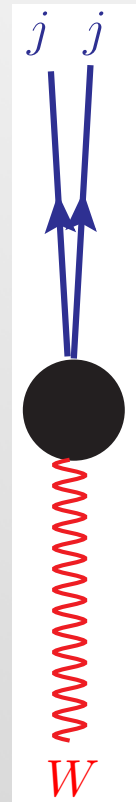
Fat-jet object

Search for a few TeV resonance decaying into $t, W, Z, h \dots$

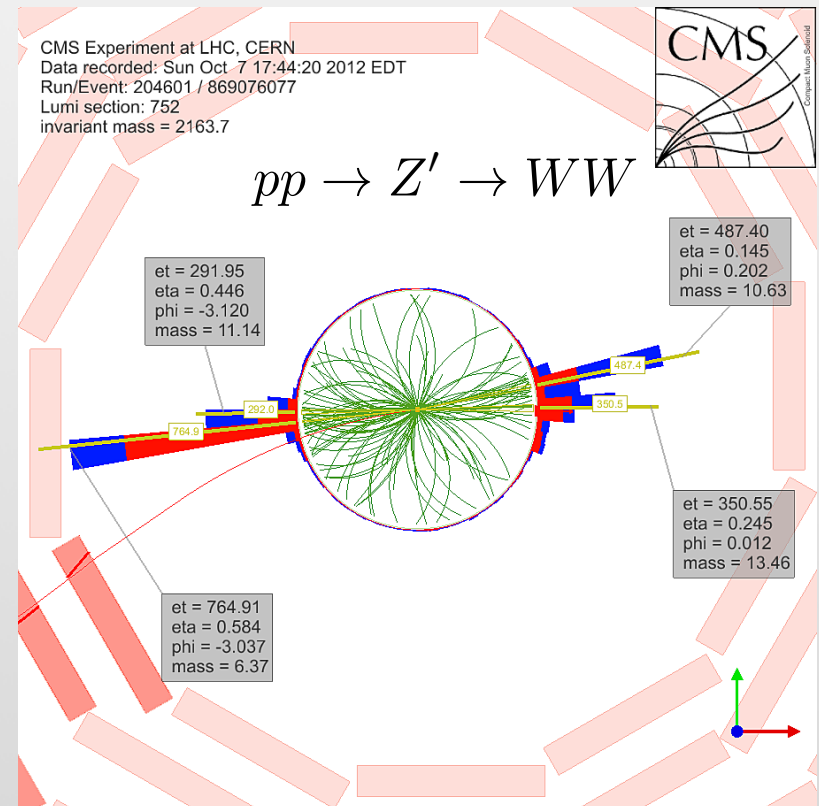
- Boosted top quark, W/Z , Higgs



$$p_T(W) \sim M_W$$



$$p_T(W) \gg M_W$$



Jet substructure

A jet may not be just a parton and it could have an internal structure

Many new objects: (incomplete list)

- ...; Butterworth, Cox, Forshaw, [WW scattering](#), [hep-ph/0201098](#)
- Butterworth, Davison, Rubin, Salam, [boosted Higgs](#), [0802.2470](#)
- Thaler and Wang, [boosted top](#), [0806.0023](#)
- Kaplan, Rehermann, Schwartz, Tweedie, [boosted top](#), [0806.0848](#)
- Almeida, Lee, Perez, Sterman, Sung, [boosted top](#), [0807.0234](#);...

Many new variables or procedures:

- mass drop, N-subjettiness, pull, dipolarity, without trees, ...
- Jet grooming: filtering, trimming, pruning ...

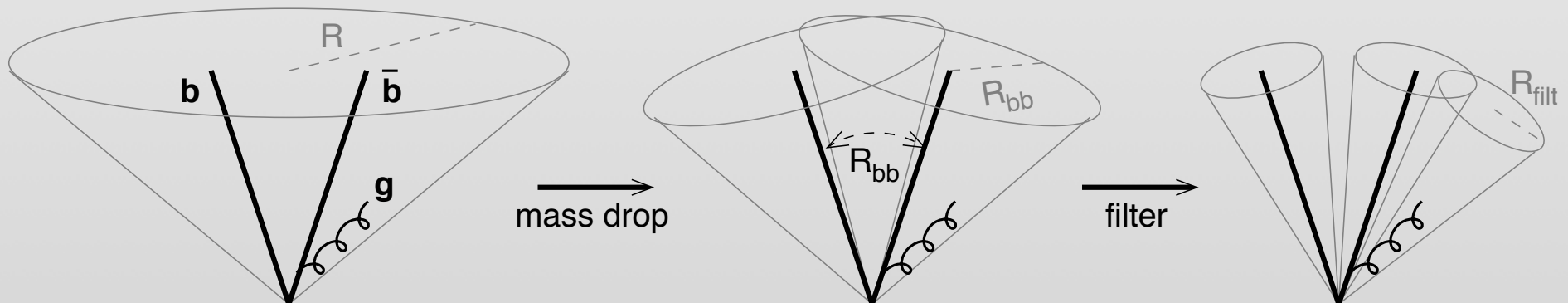
Jet substructure: an example

Boosted Higgs for measuring the $h \rightarrow b\bar{b}$ decay

Two steps:

Butterworth et.al., 0802.2470

- (1) start from a jet-finding algorithm (C/A) to cover a wider area
 - (2) **mass-drop**: (the QCD quark is massless) some subset of particles inside a Higgs-jet can have a much smaller mass.
- Filter**: (reduce underlying events) introduce a finer angular scale



Our motivation

Can we combine this two-step procedure into a single one?

- **Hope:** keep more hard process information and less underlying event contamination
- **Method:** define a new jet-finding algorithm suitable for a boosted heavy object

To proceed, let's start with traditional jet-finding algorithms for QCD jets

A brief review of jet-finding algorithms

★ Cone algorithm

- Started by Stermann and Weinberg in 70's
- CDF SearchCone, Mid point, SIScone ...
- Used at UA1, Tevatron

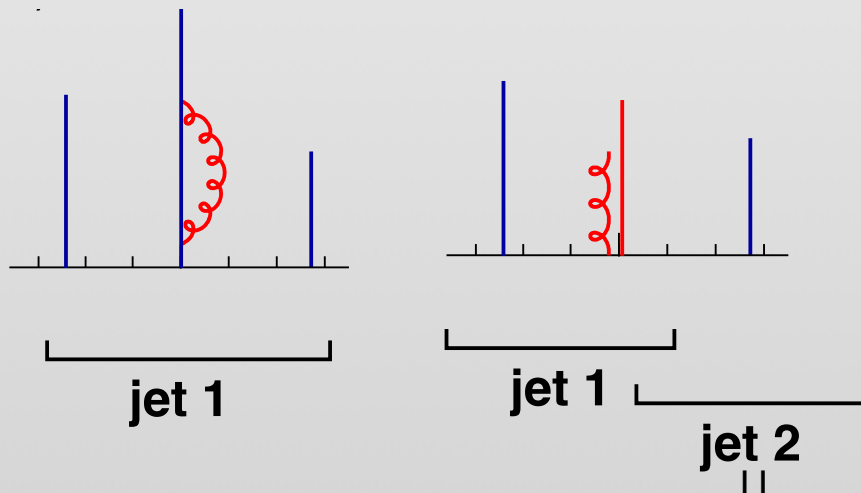
★ Sequential recombination algorithm

- Started by the JADE collaboration in 80's
- k_t , Cambridge/Aachen, anti- k_t
- Extensively used at the LHC

Cone algorithm

Iterative process:

- choose particle with highest transverse momentum as the seed particle
- draw a cone of radius R around the seed particle
- sum the momenta of all particles in the cone as the jet axis
- if the jet axis does not agree with the original one, continue; otherwise find a stable cone and stop



Colinear safety?



SISCone
(split-merge)

Anti- k_t algorithm

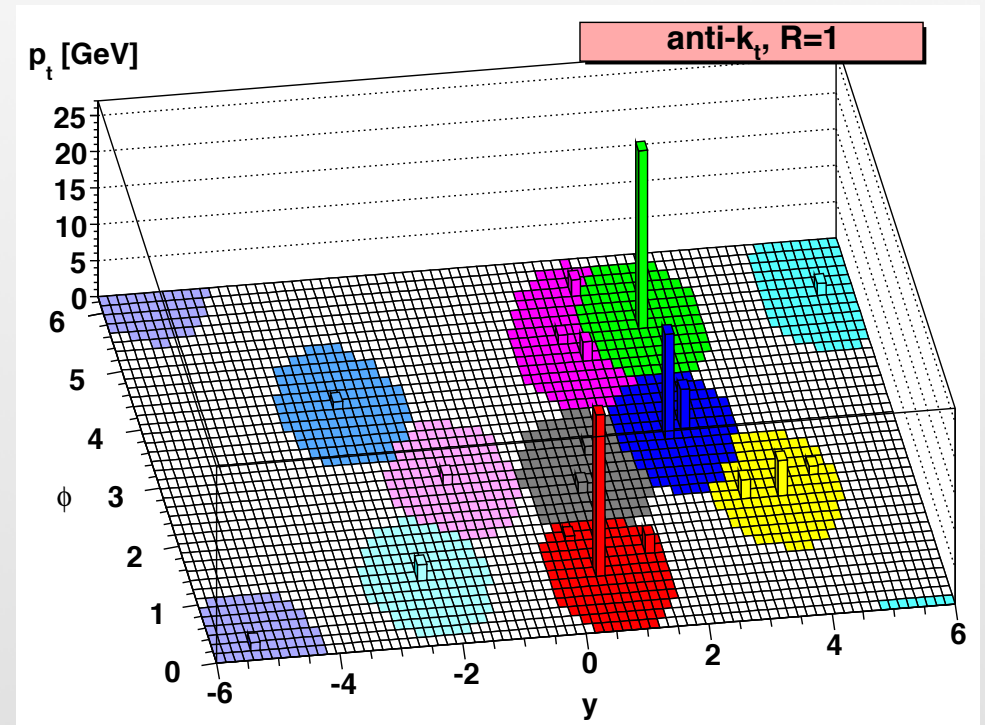
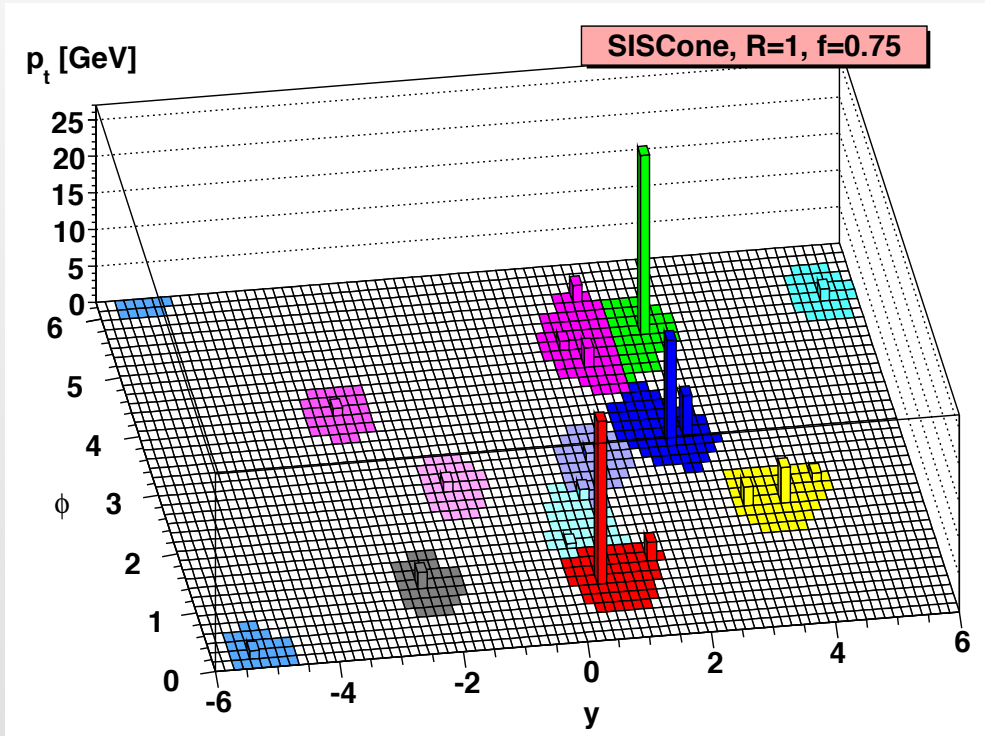
$$d_{ij} = \min(p_{ti}^{-2}, p_{tj}^{-2}) \frac{\Delta R_{ij}^2}{R^2} \quad d_{iB} = p_{ti}^{-2} \quad \Delta R_{ij}^2 = (y_i - y_j)^2 + (\phi_i - \phi_j)^2$$

Iterative process:

- Find the minimum of the d_{ij} and d_{iB}
- If it is a d_{ij} , recombine i and j into a single new particle, and repeat
- otherwise, if it is a d_{iB} , declare i to be a jet, and remove it from the list of particles
- stop when no particles remain

Infrared and collinear safe !

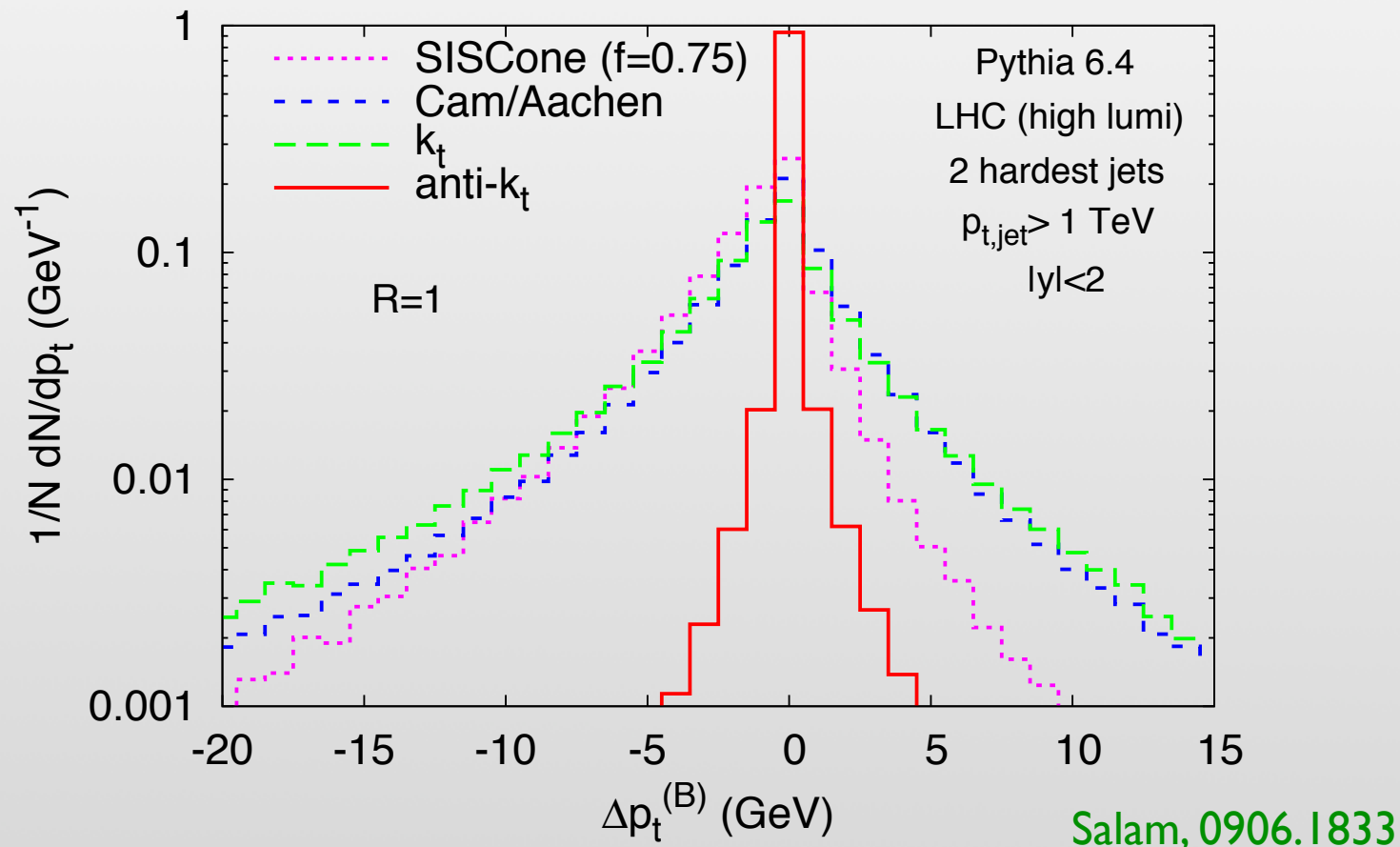
Behaviors of different algorithms



Salam, 0906.1833

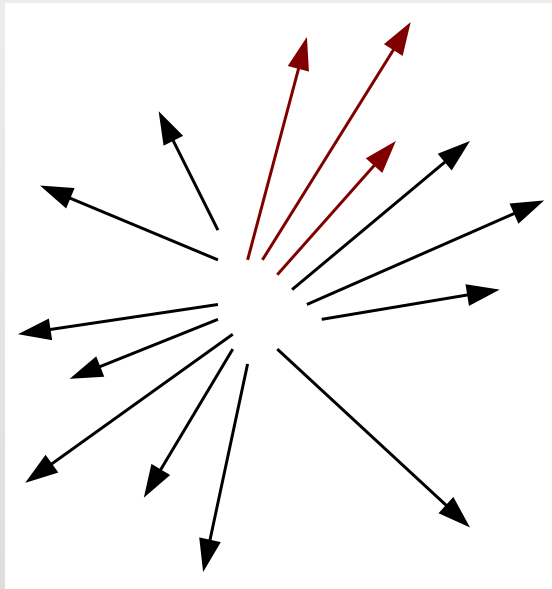
Quantify the goodness of algorithms

Back-reaction: how much adding soft background particles changes the original particles in a jet

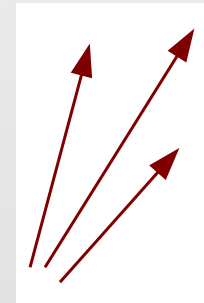


Can one has a more intuitive way to define a jet-finding algorithm?

events with
N particles

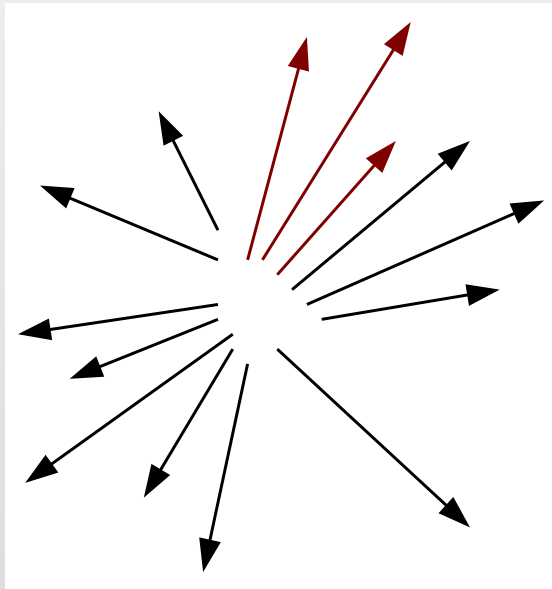


a jet with
subset particles



Can one has a more intuitive way to define a jet-finding algorithm?

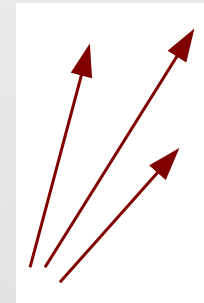
events with
N particles



A jet
function

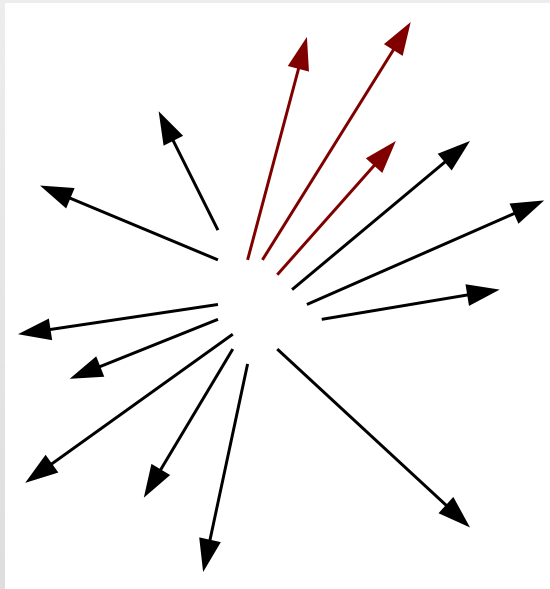


a jet with
subset particles



Can one has a more intuitive way to define a jet-finding algorithm?

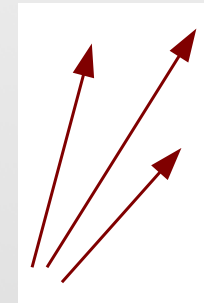
events with
N particles



A jet
function



a jet with
subset particles



Look for a simple jet definition function

Start with a QCD jet

- ★ QCD partons are massless
- ★ The jet function should
 - prefer increasing jet energy
 - disfavor increasing jet mass

Start with a QCD jet

- ★ QCD partons are massless
- ★ The jet function should
 - prefer increasing jet energy
 - disfavor increasing jet mass
- ★ The simple option at a **lepton collider**:

$$J(P^\mu) = E - \beta \frac{m^2}{E} \quad [\text{H. Georgi, 1408.1161}]$$

Start with a QCD jet

- ★ QCD partons are massless
- ★ The jet function should
 - prefer increasing jet energy
 - disfavor increasing jet mass
- ★ The simple option at a **lepton collider**:

$$J(P^\mu) = E - \beta \frac{m^2}{E} \quad [\text{H. Georgi, 1408.1161}]$$

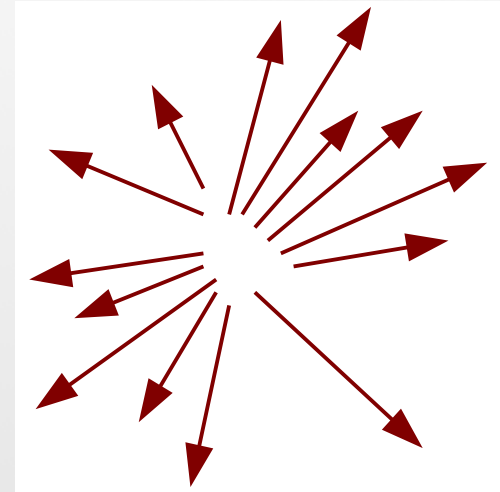
For N particles and 2^N possibilities, find the one maximizing this jet function. One does this iteratively to find all jets in one event.

Special cases

$$J(P^\mu) = E - \beta \frac{m^2}{E}$$

- $\beta = 0 : J = E$

include all particles in one jet

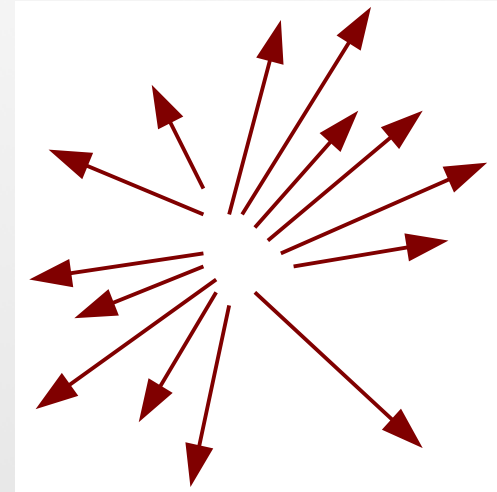


Special cases

$$J(P^\mu) = E - \beta \frac{m^2}{E}$$

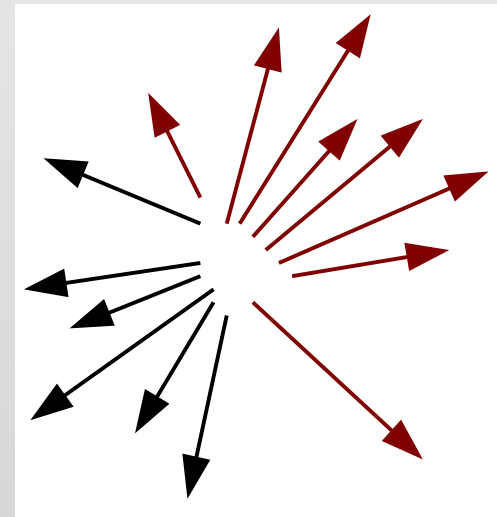
- $\beta = 0 : J = E$

include all particles in one jet



- $\beta = 1 : J = |\vec{P}|$

hemisphere way for two jets



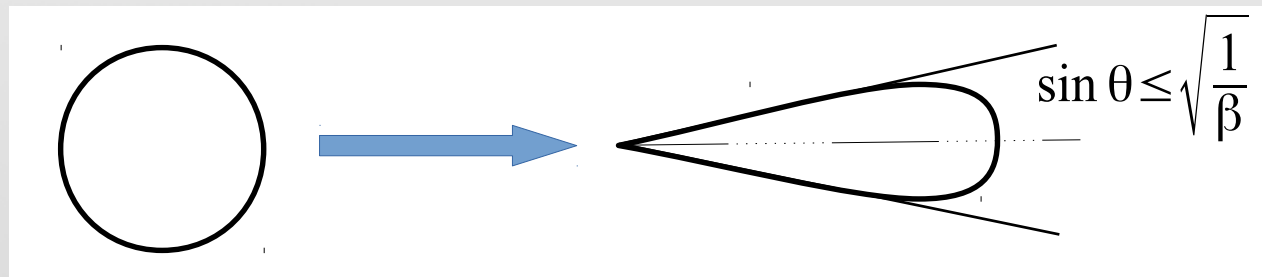
General cases

- ★ A group of particles will have a boost factor from its rest frame and has a jet function bigger than a soft particle

$$J = \frac{(E^2 - \beta m^2)}{E} = (\gamma^2 - \beta) \frac{m^2}{E} \geq 0$$

$$\gamma \geq \sqrt{\beta}$$

- ★ Relativistic beaming effect



- ★ The particles are inside a jet cone

A larger value of β means a smaller cone size

Extension to hadron colliders

★ The center-of-mass frame is likely to be highly boosted in the beam direction

★ The simplest way to extend the jet definition is

$$J_{E_T}(P_J^\mu) \equiv E_T - \beta \frac{m^2}{E_T}$$

★ One could also try other powers

$$J_{E_T}(P_J^\mu) \equiv E_T^\alpha (1 - \beta \frac{m^2}{E_T^2})$$

★ Does this new function has a similar cone geometry?

Try an “easier” function

★ For $\alpha = 2$,

$$J_{E_T^2} = E_T^2 - \beta m^2 = E^2 - P_z^2 - \beta m^2$$

★ Requiring $J_{E_T^2}(P_J^\mu) > J_{E_T^2}(P_J^\mu - p_j^\mu)$, the boundary satisfies

$$\frac{1}{|\underline{p}||P|} \left(P_x \underline{p}_x + P_y \underline{p}_y + \left(1 - \frac{1}{\beta}\right) P_z \underline{p}_z \right) = \frac{1}{v} \left(1 - \frac{1}{\beta}\right)$$

Try an “easier” function

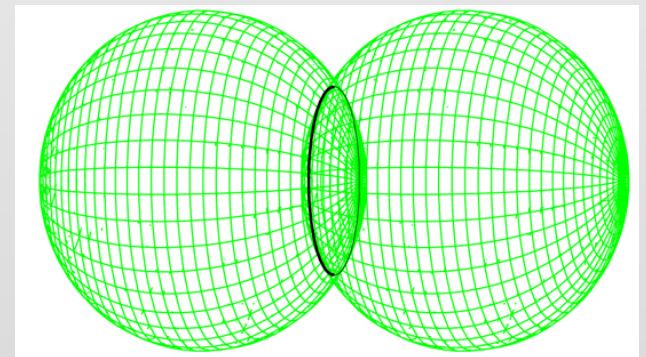
★ For $\alpha = 2$,

$$J_{E_T^2} = E_T^2 - \beta m^2 = E^2 - P_z^2 - \beta m^2$$

★ Requiring $J_{E_T^2}(P_J^\mu) > J_{E_T^2}(P_J^\mu - p_j^\mu)$, the boundary satisfies

$$\frac{1}{|\underline{p}||P|} \left(P_x \underline{p}_x + P_y \underline{p}_y + \left(1 - \frac{1}{\beta}\right) P_z \underline{p}_z \right) = \frac{1}{v} \left(1 - \frac{1}{\beta}\right)$$

$$\begin{cases} p_x^2 + p_y^2 + p_z^2 = C_1(P) \\ (p_x - P_x)^2 + (p_y - P_y)^2 + \left(p_z - \left(1 - \frac{1}{\beta}\right) P_z\right)^2 = C_2(P) \end{cases}$$



★ Can be interpreted as intersection of two spheres

Still a cone jet

★ For a general α , the boundary is

$$\frac{1}{|p||P|} (P_x p_x + P_y p_y + \kappa P_z p_z) = \frac{\kappa}{v}$$

$$\kappa = 1 - \frac{\alpha}{2\beta} + \frac{\alpha-2}{2} \frac{m^2}{E_T^2}$$

the center is shifted from the jet momentum towards the central region

$$\vec{\hat{P}}_c = \frac{1}{\sqrt{1 - (1 - \kappa^2) \hat{P}_J^z{}^2}} (\hat{P}_J^x, \hat{P}_J^y, \kappa \hat{P}_J^z) \quad \kappa < 1$$

particles belong to the jet is within a cone from the center

$$z_c \geq \frac{\kappa}{v_J \sqrt{1 - (1 - \kappa^2) \cos^2 \theta_J}}$$

Still a cone jet

★ For a general α , the boundary is

$$\frac{1}{|p||P|} (P_x p_x + P_y p_y + \kappa P_z p_z) = \frac{\kappa}{v}$$

$$\kappa = 1 - \frac{\alpha}{2\beta} + \frac{\alpha-2}{2} \frac{m^2}{E_T^2}$$

the center is shifted from the jet momentum towards the central region

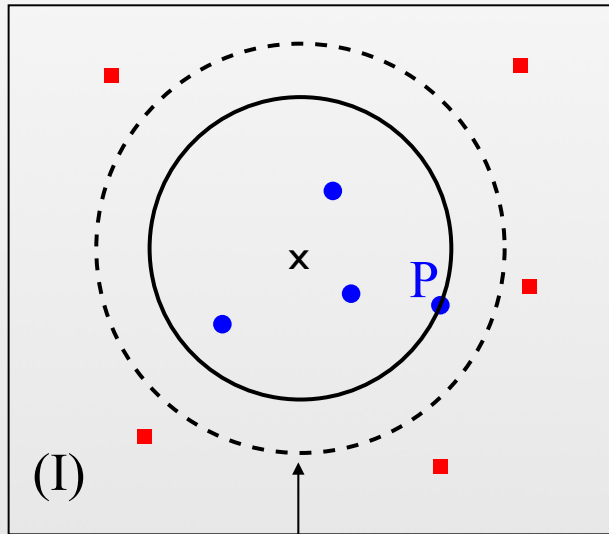
$$\vec{\hat{P}}_c = \frac{1}{\sqrt{1 - (1 - \kappa^2) \hat{P}_J^z{}^2}} (\hat{P}_J^x, \hat{P}_J^y, \kappa \hat{P}_J^z) \quad \kappa < 1$$

particles belong to the jet is within a cone from the center

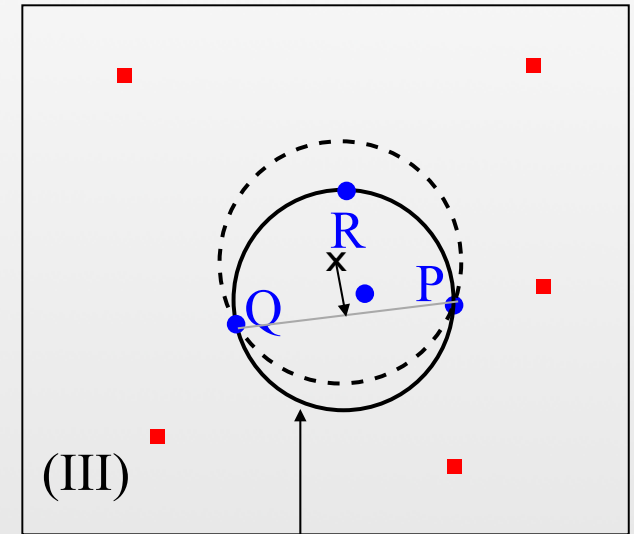
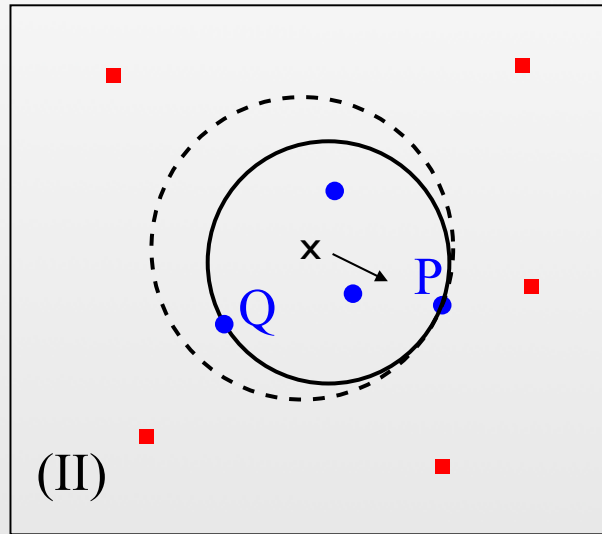
$$z_c \geq \frac{\kappa}{v_J \sqrt{1 - (1 - \kappa^2) \cos^2 \theta_J}}$$

★ The beam direction always stays away from the jet and does not need any special treatment

Cone identification

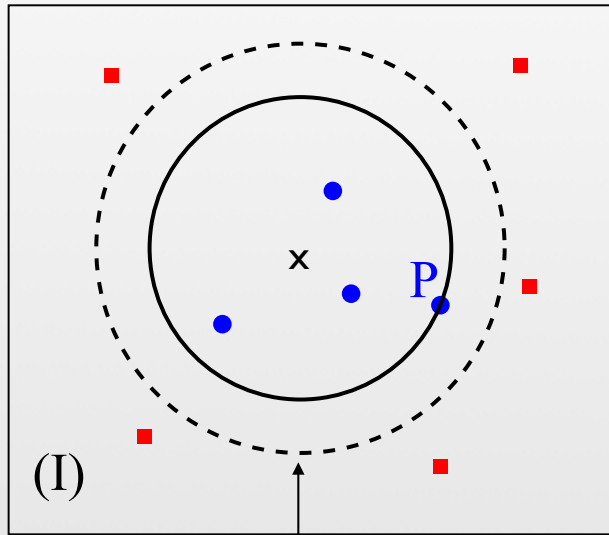


theoretical
boundary

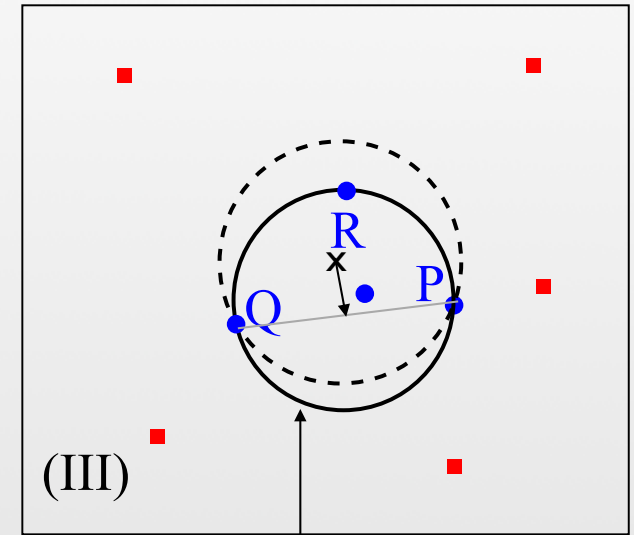
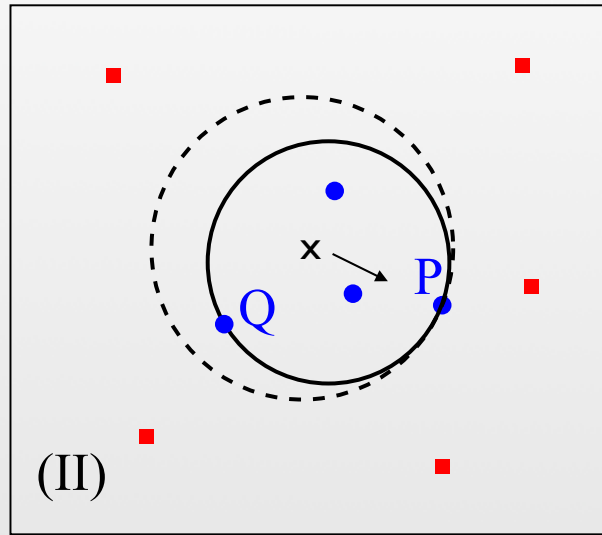


physical
boundary

Cone identification



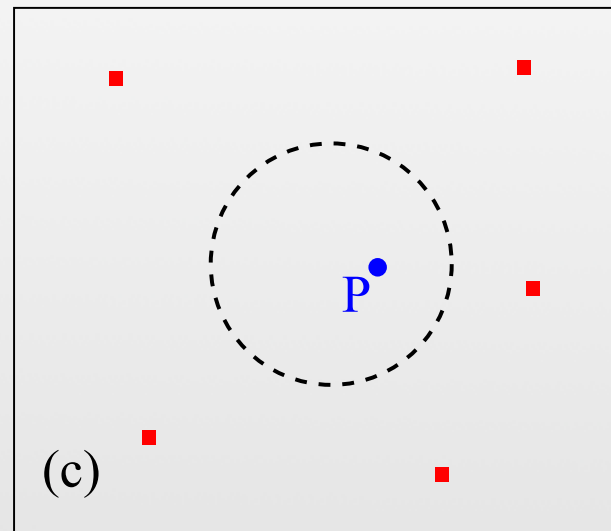
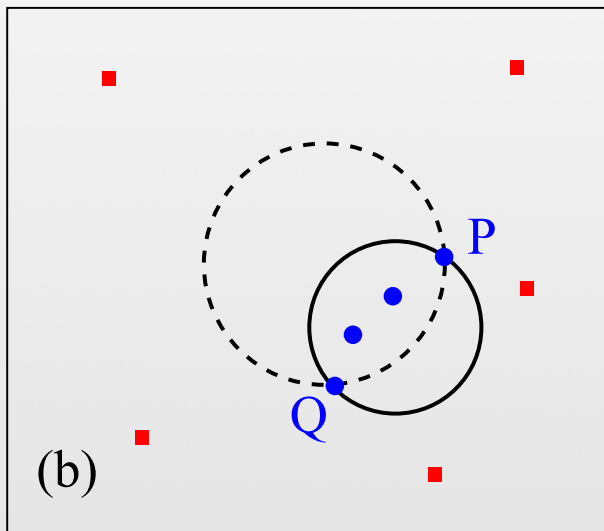
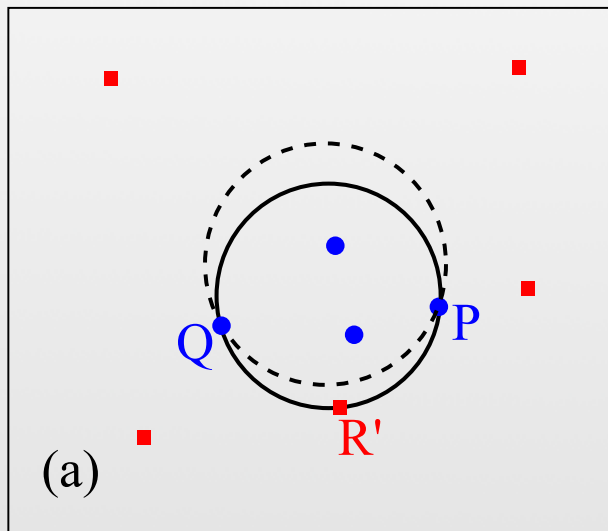
theoretical
boundary



physical
boundary

one can use three particles to identify a cone

Alternative boundaries



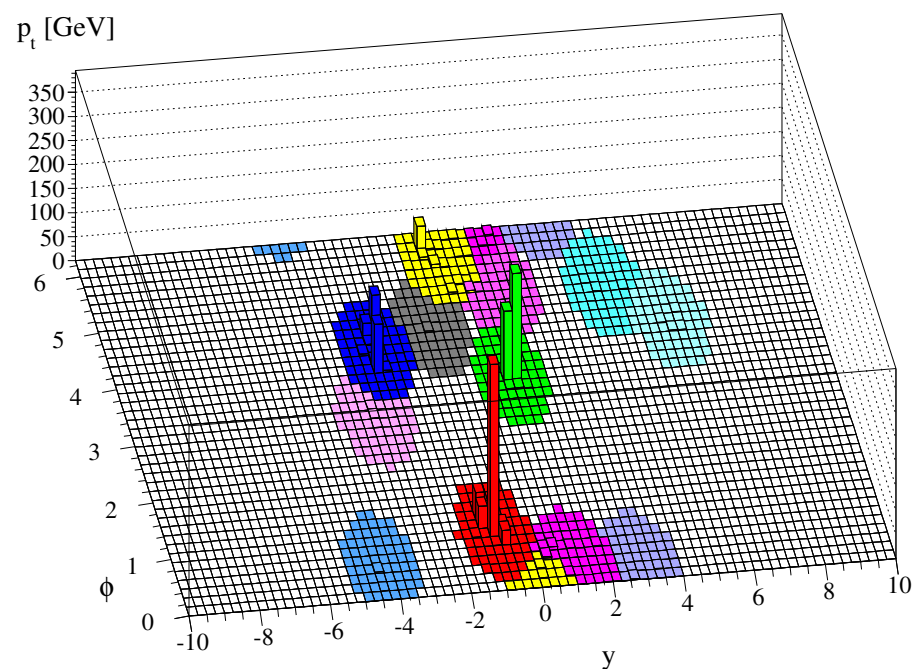
Numerical implementation

- ★ In general, we need to check all 2^N possible subsets of particles for a general function, which is not possible
- ★ Knowing the geometrical shape of jets, one only need to check all possible cones and choose the one maximizing the jet function — “global”
- ★ For each particle, one can also determine its fiducial region such that one only needs to check “ $n \ll N$ ” nearby particles as a neighbor
- ★ For each particle, the physically distinct cones is $O(n^3)$, the total operation time is $O(N n^3)$

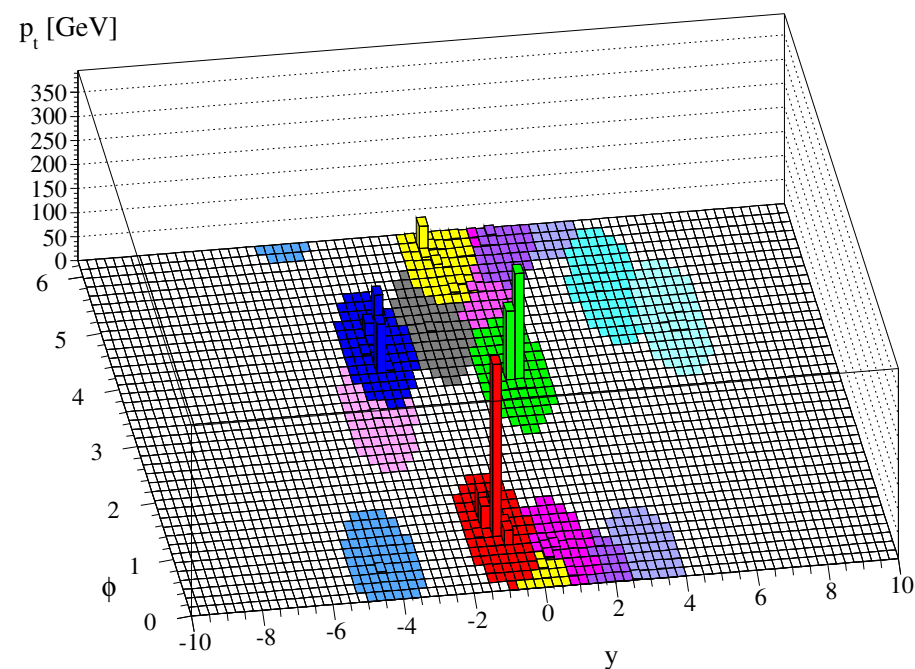
<https://github.com/LHCJet/JET>

Comparison: shape

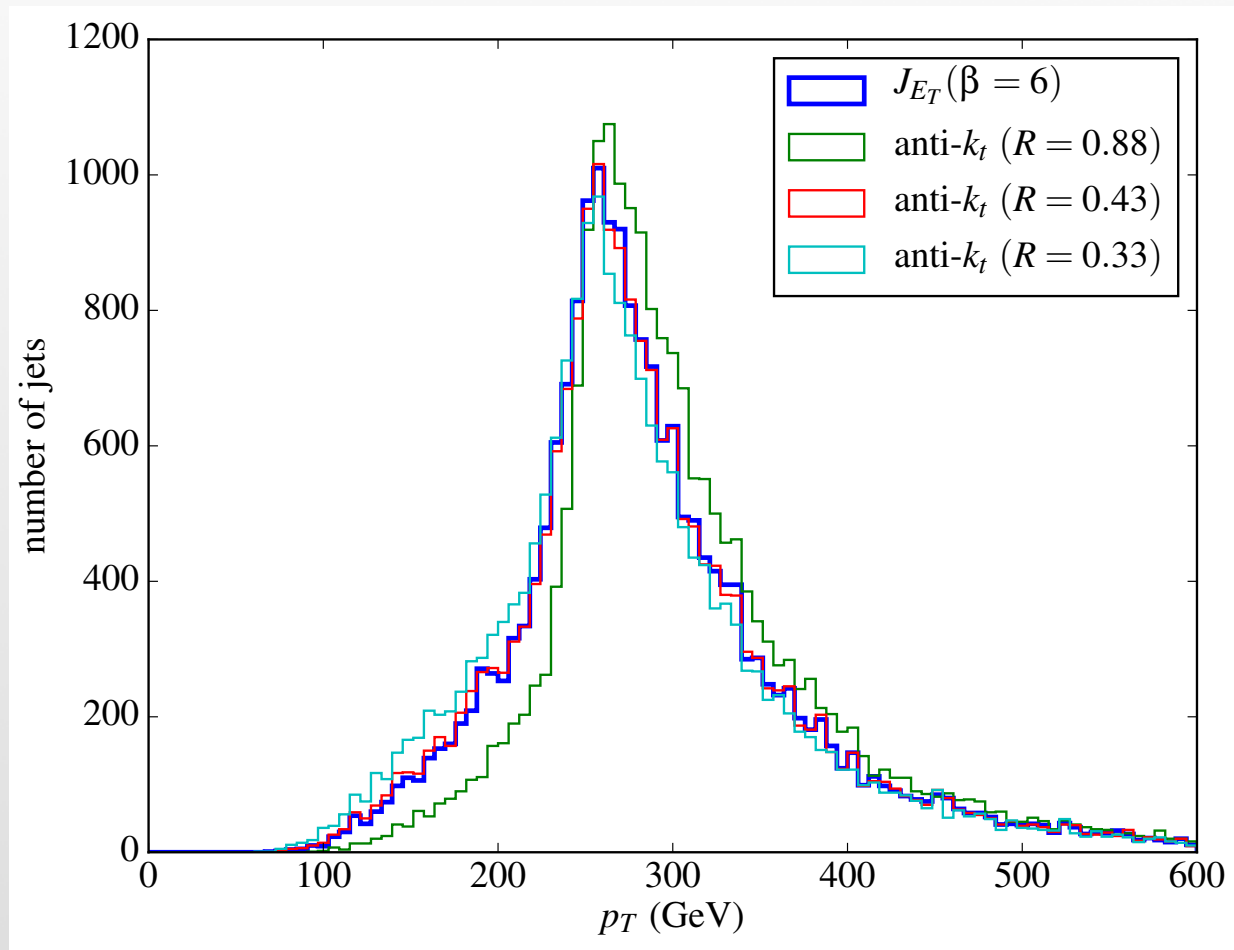
J_{E_T} with $\beta = 1.4$



anti- k_T with $R = 1.0$

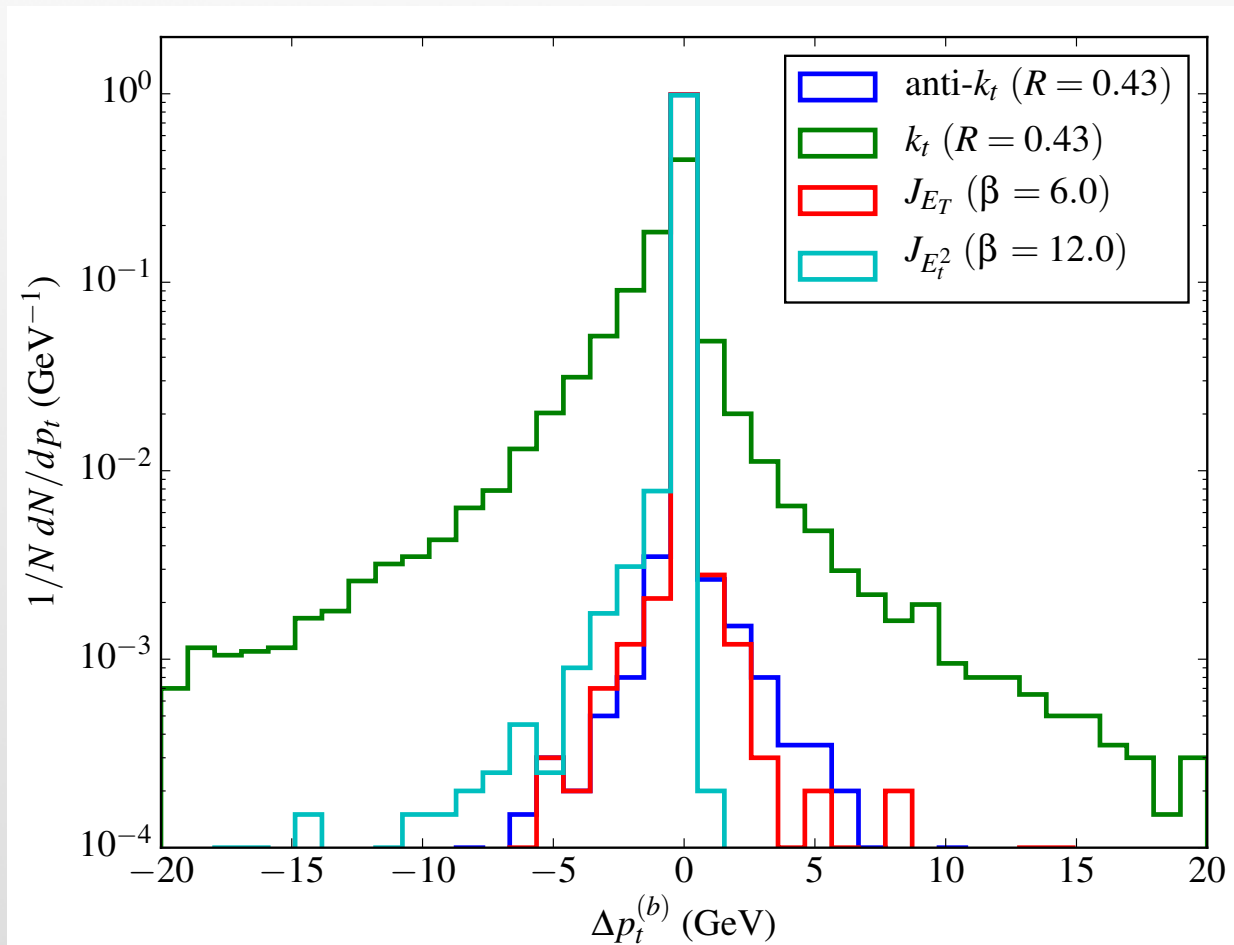


Comparison: size



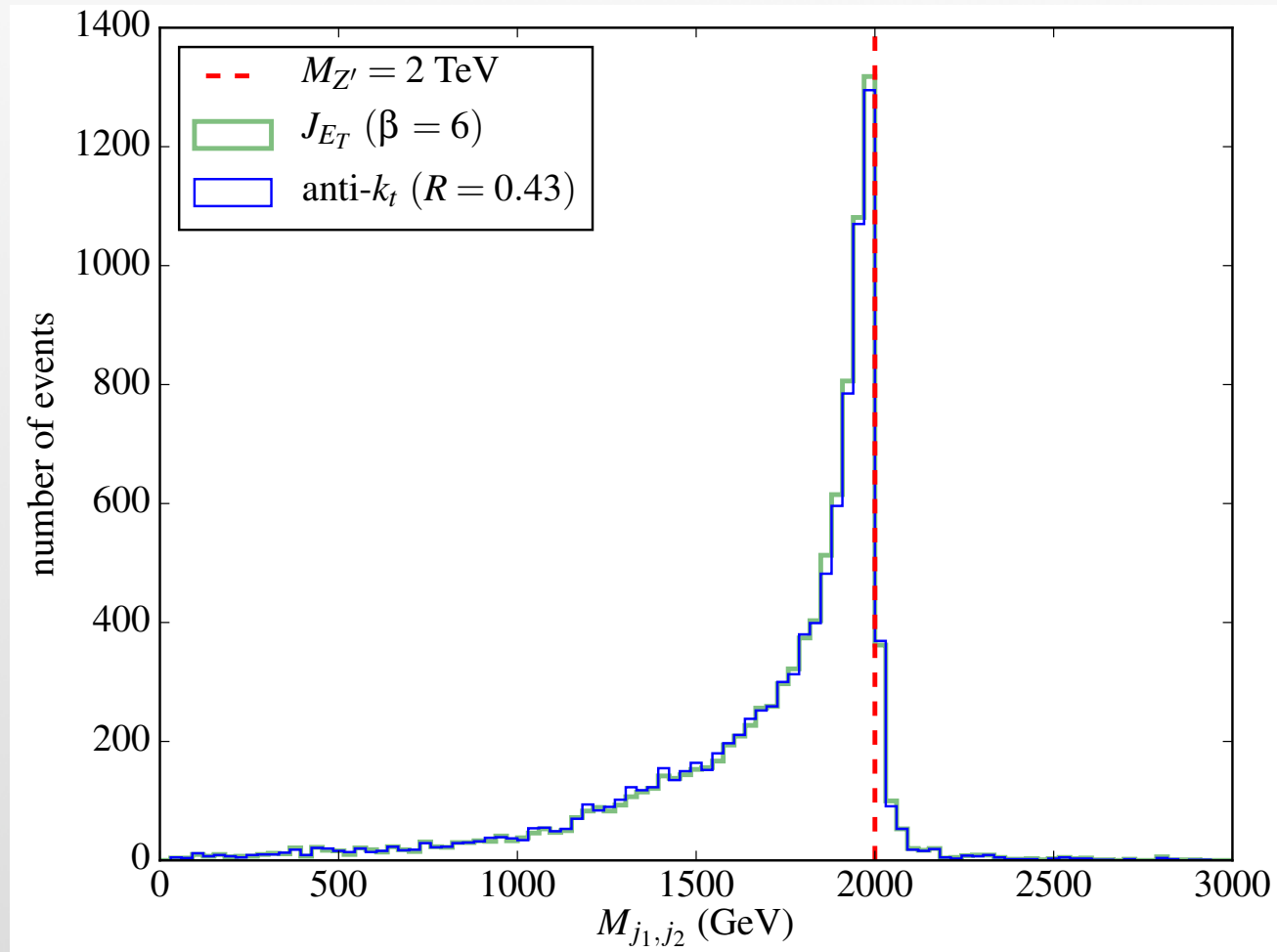
match anti-kt results very well for a QCD jet

Comparison: back-reaction



again, similar to the anti-kt results

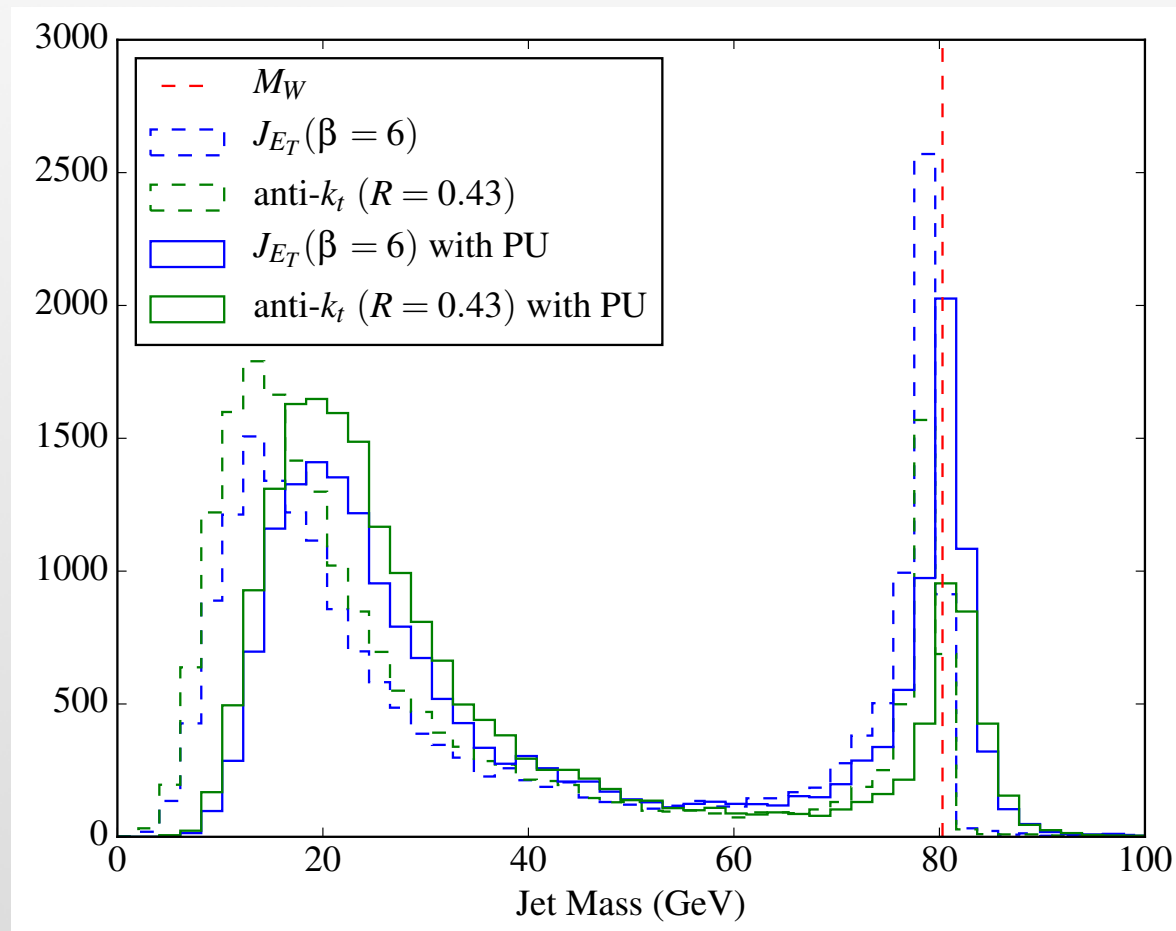
Comparison: dijet Z' mass



again, similar to the anti-kt results

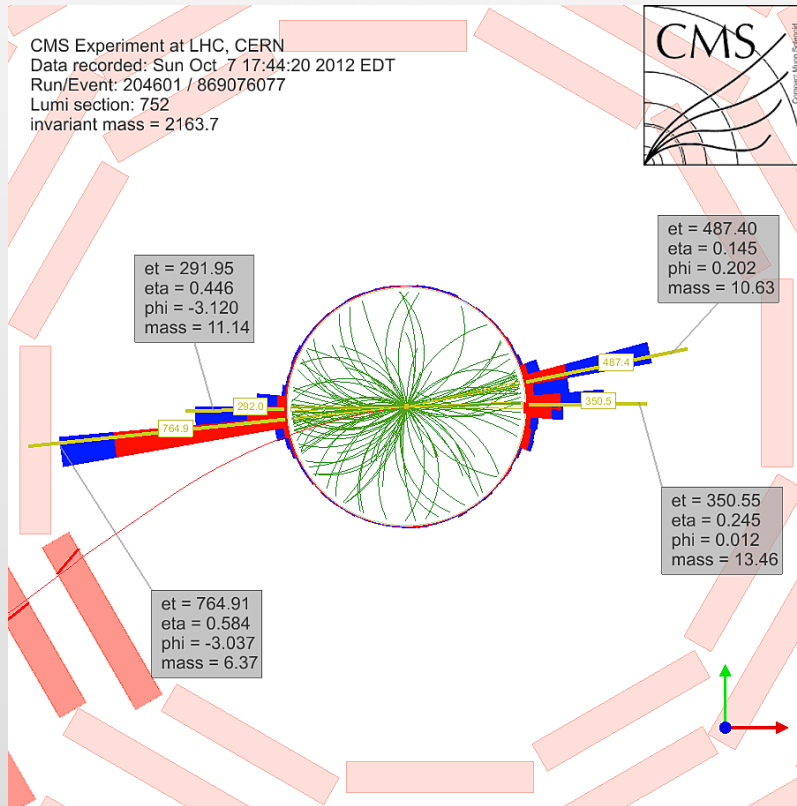
A naive comparison for W-jet

$$p_T(W) > 250 \text{ GeV}$$



our jet-finding algorithm is designed for QCD jets so far

Design a W-jet-finding function



- A boosted W-jet contains a two-prong structure
- Need to incorporate a jet shape in the function
- The existing part of J_{E_T} may be kept

Design a W-jet-finding function

$$J_{E_T}^W(P_J^\mu) = E_T^\alpha \left[1 - \beta \frac{m^2}{E_T^2} + \gamma \overline{H}_{2,J} \right]$$

- The new function need to prefer two-prong

Design a W-jet-finding function

$$J_{E_T}^W(P_J^\mu) = E_T^\alpha \left[1 - \beta \frac{m^2}{E_T^2} + \gamma \overline{H}_{2,J} \right]$$

- The new function need to prefer two-prong
- try the jet energy correlation functions:

$$\sum_{i \neq k} \frac{|\vec{p}_i| |\vec{p}_k|}{E_J^2} |\sin \varphi_{ik}|^a (1 - |\cos \varphi_{ik}|)^{1-a}$$

Banfi, Salam, Zanderighi,
hep-ph/0407286

$$\text{ECF}(N, \beta) = \sum_{i_1 < i_2 < \dots < i_N \in J} \left(\prod_{a=1}^N p_{T i_a} \right) \left(\prod_{b=1}^{N-1} \prod_{c=b+1}^N R_{i_b i_c} \right)^\beta$$

Larkoski, Salam, Thaler,
1305.0007

A working function

$$\overline{H}_{2,J} \equiv \frac{H_{2,J}}{E_T^2} = \frac{1}{E_T^2} \sum_{i,k} \frac{[m_J^2 p_i \cdot p_k - (P_J \cdot p_i)(P_J \cdot p_k)]^2}{m_J^2 (P_J \cdot p_i)(P_J \cdot p_k)}$$

- It is Lorentz invariant except the overall factor
- It becomes transparent in the jet rest frame

$$\overline{H}_{2,J} = \left(\sum_{i,k} \frac{|\vec{p}_i| |\vec{p}_k|}{E_T^2} \cos^2 \varphi_{ik} \right)_{\text{rest}} = \left(\sum_{i,k} \frac{(\vec{p}_i \cdot \vec{p}_k)^2}{E_T^2 |\vec{p}_i| |\vec{p}_k|} \right)_{\text{rest}}$$

- One can easily show that this function reaches its maximum for a two-prong structure

A working function

$$\overline{H}_{2,J} \equiv \frac{H_{2,J}}{E_T^2} = \frac{1}{E_T^2} \sum_{i,k} \frac{[m_J^2 p_i \cdot p_k - (P_J \cdot p_i)(P_J \cdot p_k)]^2}{m_J^2 (P_J \cdot p_i)(P_J \cdot p_k)}$$

- It is Lorentz invariant except the overall factor
- It becomes transparent in the jet rest frame

$$\overline{H}_{2,J} = \left(\sum_{i,k} \frac{|\vec{p}_i||\vec{p}_k|}{E_T^2} \cos^2 \varphi_{ik} \right)_{\text{rest}} = \left(\sum_{i,k} \frac{(\vec{p}_i \cdot \vec{p}_k)^2}{E_T^2 |\vec{p}_i||\vec{p}_k|} \right)_{\text{rest}}$$

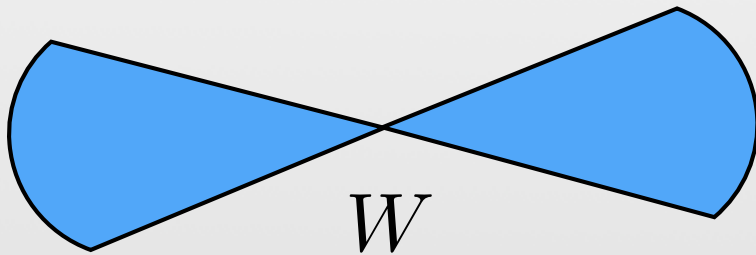
- One can easily show that this function reaches its maximum for a two-prong structure
- The function in rest frame is the **Fox-Wolfram moment**, introduced as an event shape at lepton colliders

Double-cone shape

$$J_{E_T}^W(P_J^\mu) = E_T^\alpha \left[1 - \beta \frac{m^2}{E_T^2} + \gamma \overline{H}_{2,J} \right]$$

in the lab frame

in the rest frame

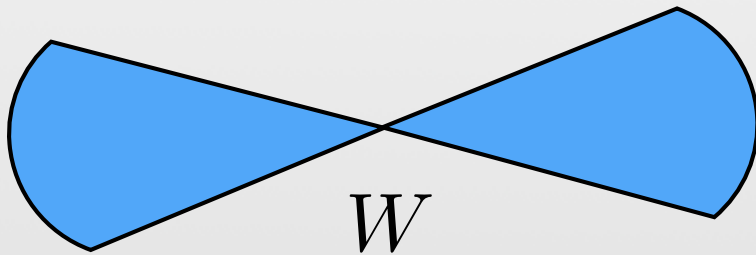


- a double-cone structure with the subjet size determined dynamically

Double-cone shape

$$J_{E_T}^W(P_J^\mu) = E_T^\alpha \left[1 - \beta \frac{m^2}{E_T^2} + \gamma \overline{H}_{2,J} \right]$$

in the rest frame

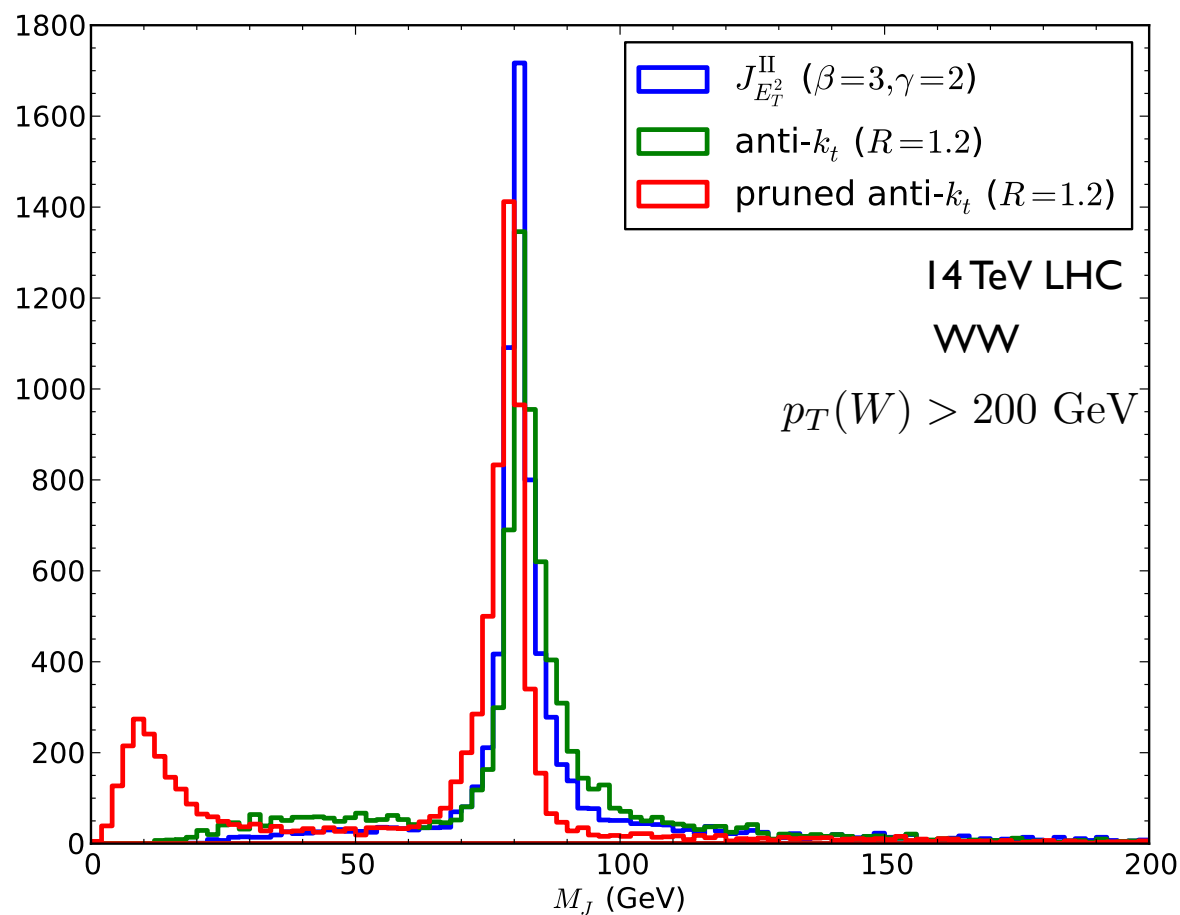


in the lab frame



- a double-cone structure with the subjet size determined dynamically
- $1/\sqrt{\beta}$ controls the subjet size and $1/(\beta - \gamma)$ controls the fat jet size

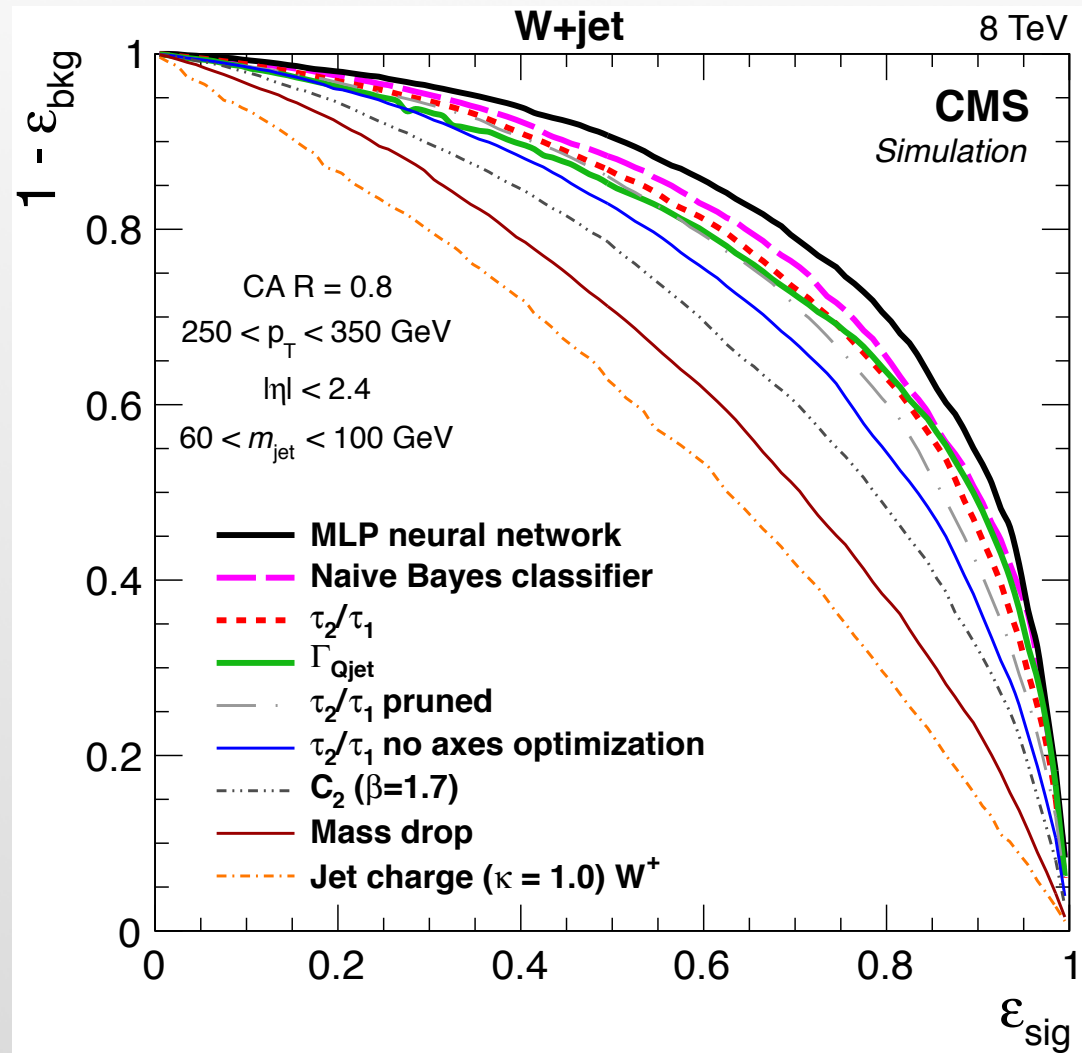
$J_{E_T^2}^W$ results



pruning jet: S. Ellis, Vermilion, Walsh; 0912.0033

no pile-up included yet

Variables used in CMS

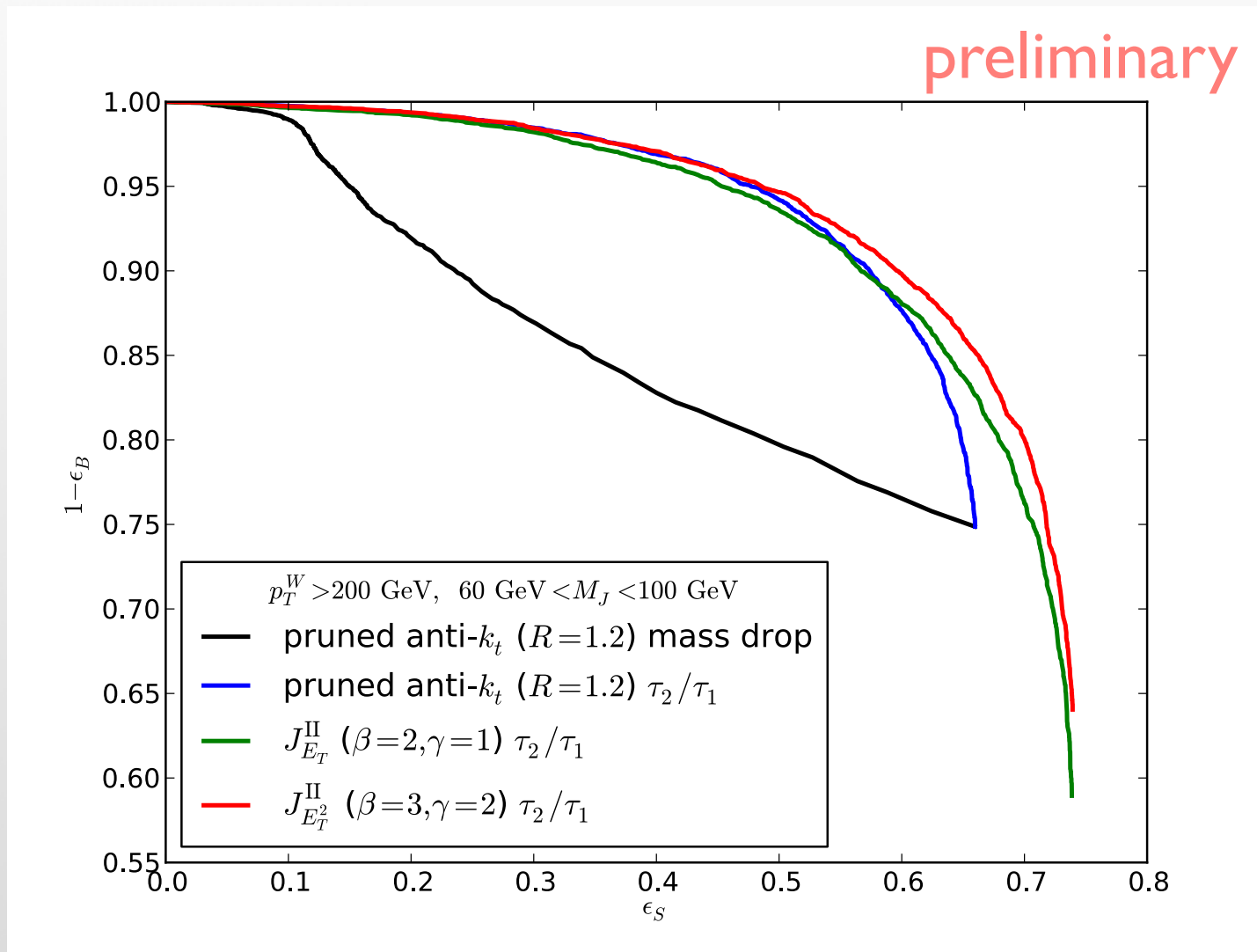


CMS; 1410.4227

N-subjettiness: Thaler and Tilburg; 1011.2268

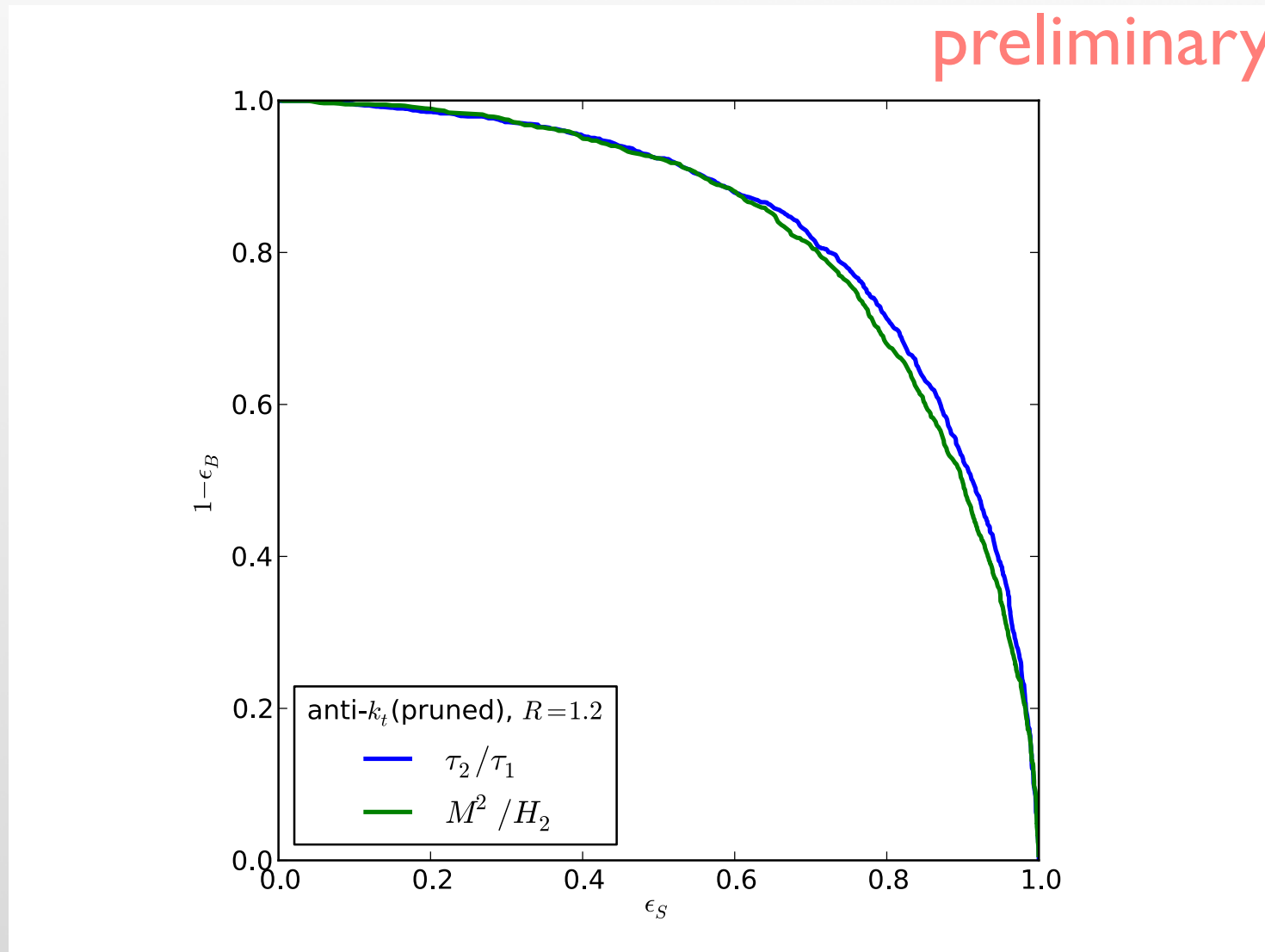
Q-jets: Ellis, Hornig, Roy, Krohn, Schwartz; 1201.1914

Performance w. Jet-sub. Variables



A better jet-finding algorithm makes some improvement

Byproduct: A New Event Shape Variable



Conclusions

- ★ A global jet-finding algorithm for maximizing a jet function works for a QCD jet
- ★ Our preliminary results show that our W -jet function can tag a W -jet very well
- ★ We are finalizing the numerical code with a trade-off between finding a global maximum and running speed
- ★ Other jet functions to tag top quark, black-hole multi-jets and new conformal gauge sector signatures are also interesting to explore

Thanks

Real proof for a cone jet

- ★ Check the angular distance of a soft particle from the jet momentum

$$z = \cos \theta = \frac{p_x P_x + p_y P_y + p_z P_z}{|p| |P|}$$

- ★ For a soft particle j belongs to the jet:

$$J(P) > J(P - p_j)$$

$$1 - \beta(1 - v_\alpha^2) > 1 - r_j - \beta \frac{1 - v_\alpha^2 - 2r_j(1 - z v_\alpha)}{1 - r_j}$$

$$z > \frac{\beta(1 + v_\alpha^2) - (1 - r_j)}{2\beta v_\alpha} > \frac{\beta(1 + v_\alpha^2) - 1}{2\beta v_\alpha} = \frac{1}{v_\alpha} \left(1 - \frac{1}{2\beta} \left(1 + \beta \frac{m^2}{E^2} \right) \right)$$

Real proof for a cone jet

★ For a soft particle j belongs to the jet:

$$z > \frac{\beta(1+v_\alpha^2) - (1 - r_j)}{2\beta v_\alpha} > \frac{\beta(1+v_\alpha^2) - 1}{2\beta v_\alpha} = \frac{1}{v_\alpha} \left(1 - \frac{1}{2\beta} \left(1 + \beta \frac{m^2}{E^2} \right) \right)$$

★ For a soft particle k not belongs to the jet:

$$z < \frac{\beta(1+v_\alpha^2) - (1 + r_k)}{2\beta v_\alpha} < \frac{\beta(1+v_\alpha^2) - 1}{2\beta v_\alpha} = \frac{1}{v_\alpha} \left(1 - \frac{1}{2\beta} \left(1 + \beta \frac{m^2}{E^2} \right) \right)$$

★ So, a cone-like boundary for individual jets

★ Soft particles are on the boundary; very IR safe