

\mathcal{T}^{jet} resummation in Higgs production at $\text{NLL}' + \text{NLO}$

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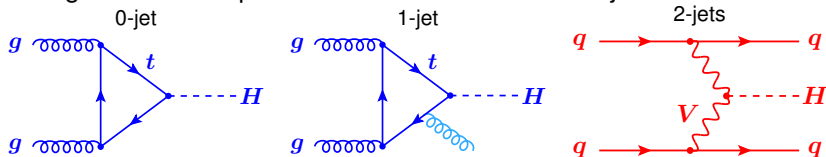
In collaboration with Maximillian Stahlhofen and Frank Tackmann
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Introduction

- ▶ Jet binning is important :

Distinguish different processes based on number of jets in the final state.



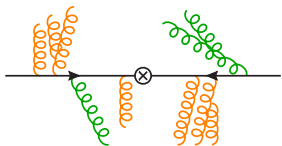
Separate different backgrounds : eg in $H \rightarrow W W^*$ any hard jets with $p_T^{\text{jet}} > p_T^{\text{cut}} \sim 20 - 30 \text{ GeV}$ are vetoed.

- ▶ Any type of such exclusive measurements or vetoes induce large Sudakov double logs of the veto scale p_T^{cut} : $\alpha_s^n \log^m [p_T^{\text{cut}}/m_H]$.
- ▶ Eg: $gg \rightarrow H + 0 \text{ jet}$

$$\sigma_0(p_T^{\text{cut}}) \propto \sigma_B \left(1 - 2 \frac{\alpha_s C_A}{\pi} \log^2 \frac{p_T^{\text{cut}}}{m_H} + \dots \right)$$

- ▶ These large logarithms must be resummed to all orders to obtain reliable precision predictions.

Jet veto Observables



Global Veto

restricts Σ of all emissions

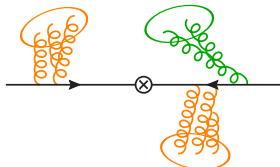
"beam broadening"

$$E_T = \sum_i |\vec{p}_{Ti}|$$

"beam thrust"

$$\mathcal{T} = \sum_i |\vec{p}_{Ti}| e^{-|y_i - Y|}$$

(y_i and Y are the jet and Higgs rapidity respectively.)



Local Veto

restrict (local chunks of)
individual emissions

"jet p_T "

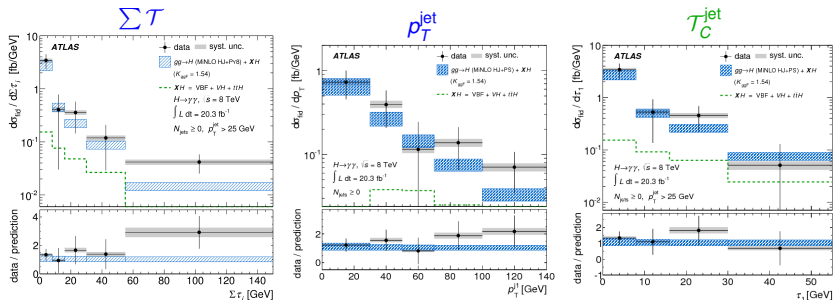
$$\vec{p}_T^{jet} = \max_{j \in jets} |\vec{p}_{Tj}|$$

"Rapidity-Dependent"

$$\mathcal{T}_f^{jet} = \max_{j \in jets} |p_{Tj}| f(y_j)$$

Jet veto Observables

- ▶ Fiducial differential cross section measurements exist for different jet veto observables.



- ▶ Resummation of various jet veto observables has been done :

- 1 Beam Thrust resummation: Berger, Marcantonini, Stewart, Tackmann, Waalewijn
- 2 p_T^{jet} resummation :
 H + 0-jet: Banfi, Monni, Salam, Zanderighi [1203.5773, 1206.4998]
 H + 0-jet: Becher, Neubert, Rothen [1205.3806, 1307.0025]
 H + 0-jet: Stewart, Tackmann, Walsh, Zuberi [1307.1808]
- 3 In this talk : $\mathcal{T}_{B(C)}^{\text{jet}}$ resummation

Rapidity dependent Jet Vetoes $\mathcal{T}_{B,C}$

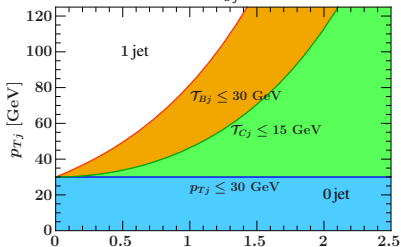
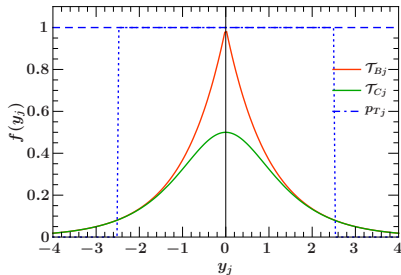
$$p_T^{\text{jet}} = \max_{j \in \text{jets}} \{ |\vec{p}_{Tj}| \theta(|y_j| < y^{\text{cut}}) \}, \quad \mathcal{T}_{fj} = |\mathbf{p}_{Tj}| f(y_j) \rightarrow \mathcal{T}_f^{\text{jet}} = \max_{j \in \text{jets}} \mathcal{T}_{fj} < \mathcal{T}^{\text{cut}}$$

$$\mathcal{T}_{Bcm}^{\text{jet}} : f(y_j) = e^{-|y_j|}$$

$$\mathcal{T}_B^{\text{jet}} : f(y_j) = e^{-|y_j - Y|}$$

$$\mathcal{T}_{Ccm}^{\text{jet}} : f(y_j) = \frac{1}{2 \cosh(y_j)}$$

$$\mathcal{T}_C^{\text{jet}} : f(y_j) = \frac{1}{2 \cosh(y_j - Y)}$$



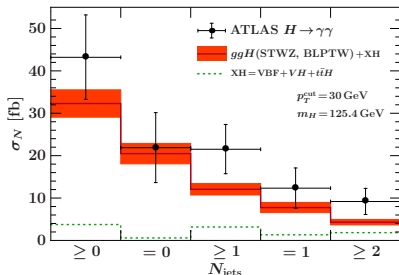
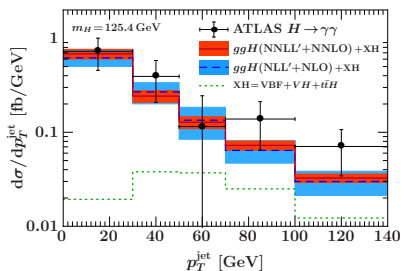
$$\text{0-jet bin: } \mathcal{T}_f^{\text{jet}} < \mathcal{T}^{\text{cut}}$$

$$\geq \text{1-jet bin: } \mathcal{T}_f^{\text{jet}} > \mathcal{T}^{\text{cut}}$$

Rapidity dependent jet vetoes $\mathcal{T}_{B,C}$

Motivation

- ▶ Jets can only be reconstructed down to some minimum p_T : limit on the jet veto cut. Low p_T jets hard to measure at forward rapidities.
- ▶ Rapidity dependent observables : A tight veto on central jets and veto constraint getting looser for forward jets.
- ▶ $\mathcal{T}_B^{\text{jet}}$ and $\mathcal{T}_C^{\text{jet}}$ are rapidity weighted p_T^{jet} , resumable to same level as p_T^{jet} .
- ▶ Provide complementary information in the exclusive jet bins.



[Stewart, Tackmann, Walsh, Zuberi], [Boughezal, Petriello, Liu, Walsh, Tackmann]

Soft-collinear Factorization and Resummation

The cross section in SCET for $gg \rightarrow H$:

$$d\sigma_0 = |C_{ggH}|^2 \langle p_a p_b | O_{ggH}^\dagger M^{\text{veto}} O_{ggH} | p_a p_b \rangle$$

$$O_{ggH} \sim H O_a O_s O_b \sim H B_{n_a}^\mu T [Y_{n_a}^\dagger Y_{n_b}] B_{n_b, \mu}$$

Measurement function M^{veto}

- ▶ Implements phase space constraints due to jet veto.
- ▶ Soft-collinear factorization implies

$$M^{\text{veto}} = M_a^{\text{veto}} \times M_b^{\text{veto}} \times M_s^{\text{veto}} + \mathcal{O}(R^2)$$

- ▶ Matrix elements factorize into independent soft and collinear matrix elements

$$B_{a,b}(\mu) = \langle p_a | O_{a,b}^\dagger M_{a,b} O_{a,b} | p_b \rangle$$

$$S_{a,b}(\mu) = \langle 0 | O_s^\dagger M_s O_s | 0 \rangle$$

Jet algorithm effects in local vetoes

- ▶ Measurement function for $\mathcal{T}_f^{\text{jet}} < \mathcal{T}^{\text{cut}}$: product of measurement functions on individual jets.

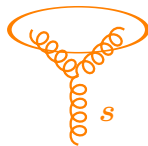
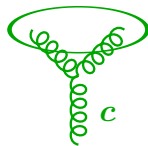
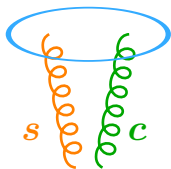
$$M_f^{\text{jet}}(\mathcal{T}^{\text{cut}}) = \theta(\mathcal{T}_f^{\text{jet}} < \mathcal{T}^{\text{cut}}) = \prod_{j \in J(R)} \theta(\mathcal{T}_{fj} < \mathcal{T}^{\text{cut}})$$

- ▶ Measurement factorized into soft and collinear components:

$$M_f^{\text{jet}} = (M_{fa}^{\text{jet}} + \Delta M_{fa}^{\text{jet}})(M_{fb}^{\text{jet}} + \Delta M_{fb}^{\text{jet}})(M_{fs}^{\text{jet}} + \Delta M_{fs}^{\text{jet}}) + \delta M_f^{\text{jet}}$$

δM_f^{jet} : $O(R^n)$ clustering of soft and collinear components into the same jet. Relevant for large $R \sim 1$.

ΔM_f^{jet} : $O(\ln^n R)$ clustering within each sector. Relevant for small $R \ll 1$



- ▶ Typical jet radius used in experiments ATLAS: 0.4, CMS: 0.5

Soft collinear Factorization

The full H+0-jet cross section differential in the Higgs rapidity Y and with $\mathcal{T}_{B,C}^{\text{jet}} < \mathcal{T}^{\text{cut}}$

$$\frac{d\sigma_0}{dY} = \sigma_B H_{gg}(m_t, m_H^2, \mu) B_g(m_H \mathcal{T}^{\text{cut}}, x_a, R, \mu) B_g(m_H \mathcal{T}^{\text{cut}}, x_b, R, \mu) \times \\ S_{gg}^{B,C}(\mathcal{T}^{\text{cut}}, R, \mu) + \frac{d\sigma_0^{\text{Rsub}}}{dY}(\mathcal{T}_{B,C}^{\text{jet}} < \mathcal{T}^{\text{cut}}, R) + \frac{d\sigma_0^{\text{non}}}{dY}(\mathcal{T}_{B,C}^{\text{jet}} < \mathcal{T}^{\text{cut}})$$

For $\mathcal{T}_{Bcm,Ccm}^{\text{jet}} < \mathcal{T}^{\text{cut}}$

$$\frac{d\sigma_0}{dY} = \sigma_B H_{gg}(m_t, m_H^2, \mu) B_g(m_H \mathcal{T}^{\text{cut}} e^Y, x_a, R, \mu) B_g(m_H \mathcal{T}^{\text{cut}} e^{-Y}, x_b, R, \mu) \times \\ S_{gg}^{B,C}(\mathcal{T}^{\text{cut}}, R, \mu) + \frac{d\sigma_0^{\text{Rsub}}}{dY}(\mathcal{T}_{Bcm,Ccm}^{\text{jet}} < \mathcal{T}^{\text{cut}}, R) + \frac{d\sigma_0^{\text{non}}}{dY}(\mathcal{T}_{Bcm,Ccm}^{\text{jet}} < \mathcal{T}^{\text{cut}})$$

Resummation

- ▶ Only difference between $\mathcal{T}_B^{\text{jet}}$ and $\mathcal{T}_C^{\text{jet}}$ to all orders: different **soft functions**.
- ▶ **Beam functions** same: Describe collinear ISR i.e. emissions with forward rapidities where $\mathcal{T}_B^{\text{jet}}$ and $\mathcal{T}_C^{\text{jet}}$ measurements are equal.
- ▶ Logarithms are split apart and resummed using RGE

$$\ln^2 \frac{\mathcal{T}^{\text{cut}}}{m_H} = 2 \ln^2 \frac{m_H}{\mu} - \ln^2 \frac{\mathcal{T}^{\text{cut}} m_H}{\mu^2} + 2 \ln^2 \frac{\mathcal{T}^{\text{cut}}}{\mu}$$

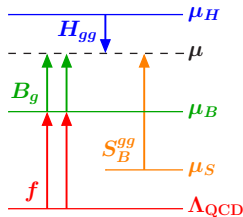
- ▶ Natural scales for **Hard**, **beam** and **soft** functions:

$$\mu_H \simeq m_H, \mu_B \simeq \sqrt{\mathcal{T}^{\text{cut}} m_H}, \mu_S \simeq \mathcal{T}^{\text{cut}}$$

- ▶ **Hard function** contains: $\ln^2 [-m_H^2/\mu_H^2]$

- ▶ **Beam functions**: $\ln^2 [m_H \mathcal{T}^{\text{cut}}/\mu_B^2]$

- ▶ **Soft function**: $\ln^2 [\mathcal{T}^{\text{cut}}/\mu_S]$



Resummation

- ▶ Resummation Structure:

$$\ln \sigma_0(\mathcal{T}^{\text{cut}}) \sim \sum_n \alpha_s^n \ln \left[\frac{\mathcal{T}^{\text{cut}}}{m_H} \right]^{n+1} (1 + \alpha_s + \alpha_s^2 + \dots) \sim \text{LL} + \text{NLL} + \text{NNLL} + \dots$$

- ▶ Including the RGE running the resummed cross section for \mathcal{T}^{jet}

$$\sigma_0(\mathcal{T}^{\text{cut}}) = H_{ggH}(m_H, \mu_H) U_H(m_H, \mu_H, \mu) \times B_g(x_a, R, \mathcal{T}^{\text{cut}}, \mu_B) \\ B_g(x_b, R, \mathcal{T}^{\text{cut}}, \mu_B) U_B(m_H, \mu_B, \mu)^2 \times S_{gg}^{B,C}(R, \mathcal{T}^{\text{cut}}, \mu_S) U_S(\mu_S, \mu)$$

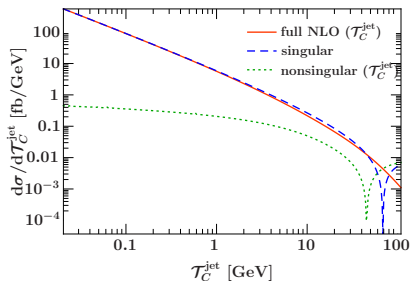
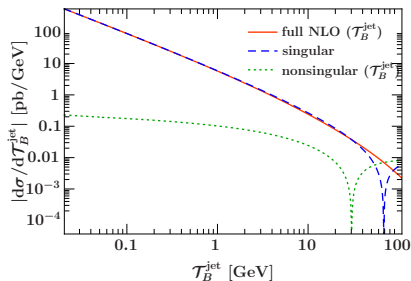
Log counting:	Fixed-order corrections		Resummation input		
	matching	nonsingular	$\gamma_{H,B,S}^\mu$	Γ_{cusp}	β
NLL	1	-	1-loop	2-loop	2-loop
NLL' + NLO	NLO	NLO	1-loop	2-loop	2-loop
NNLL' + NNLO	NNLO	NNLO	2-loop	3-loop	3-loop

Nonsingular Contribution

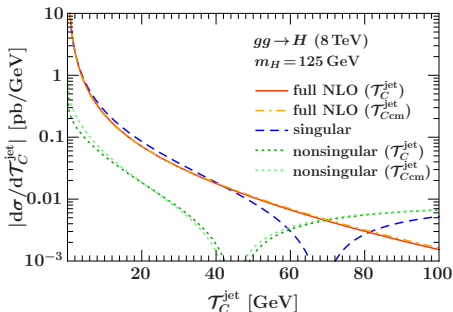
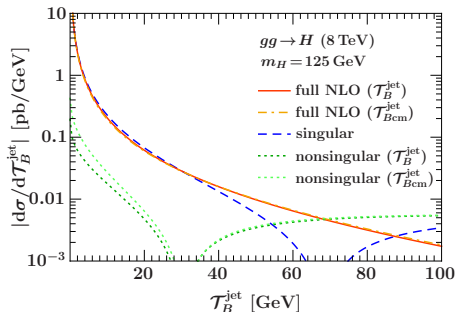
- ▶ To compute nonsingular corrections for different jet veto observables:

$$\frac{d\sigma_0^{\text{nonsing}}}{d\mathcal{T}_f^{\text{jet}} dY} = \frac{d\sigma_0^{\text{FO}}}{d\mathcal{T}_f^{\text{jet}} dY} - \frac{d\sigma_0^{\text{resum}}}{d\mathcal{T}_f^{\text{jet}} dY} (\mu_H = \mu_B = \mu_S = \mu_{\text{FO}})$$

- ▶ $d\sigma_0^{\text{FO}}$: Computed full FO cross section differential in Y and $\mathcal{T}_f^{\text{jet}}$. Implemented non-trivial measurement functions for $\mathcal{T}_f^{\text{jet}}$.
- ▶ Validation plots:



Resummation scales



- ▶ **Resummation region:** Logs are resummed using canonical scaling

$$|\mu_H| \sim m_H$$

$$\mu_S \sim \mathcal{T}^{\text{cut}}$$

$$\mu_B \sim \sqrt{m_H \mathcal{T}^{\text{cut}}}$$

- ▶ **FO region:** Resummation is turned off to get the right FO cross section at large \mathcal{T}^{cut}

$$\mu_B, \mu_S \rightarrow \mu_{\text{FO}} \sim m_H$$

- ▶ **Transition region:** Profiles for μ_B, μ_S provide smooth transition from resummation to fixed-order region.

Perturbative Uncertainties in Jet Binning

$$\begin{aligned}\sigma_{\text{total}} &= \int_0^{\mathcal{T}^{\text{cut}}} d\mathcal{T}^{\text{jet}} \frac{d\sigma}{d\mathcal{T}^{\text{jet}}} + \int_{\mathcal{T}^{\text{cut}}}^{\infty} d\mathcal{T}^{\text{jet}} \frac{d\sigma}{d\mathcal{T}^{\text{jet}}} \\ &\equiv \sigma_0(\mathcal{T}^{\text{cut}}) + \sigma_{\geq 1}(\mathcal{T}^{\text{cut}})\end{aligned}$$

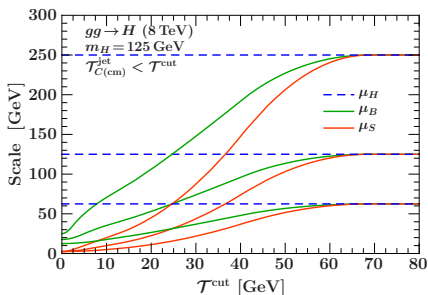
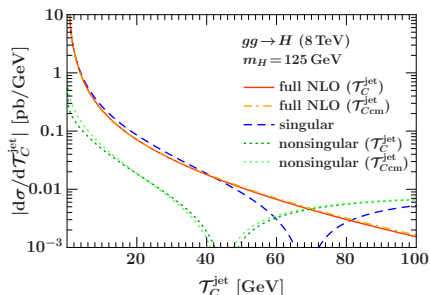
- Uncertainties in the jet binning can be described by fully correlated (yield) and fully anticorrelated (migration) components of a covariance matrix $\{\sigma_0, \sigma_{\geq 1}\}$ [Stewart,Tackmann: 1107.2117, SG,Tackmann: 1302.4437]

$$\mathbf{C} = \begin{pmatrix} (\Delta_0^y)^2 & \Delta_0^y \Delta_{\geq 1}^y \\ \Delta_0^y \Delta_{\geq 1}^y & (\Delta_{\geq 1}^y)^2 \end{pmatrix} + \begin{pmatrix} \Delta_{\text{cut}}^2 & -\Delta_{\text{cut}}^2 \\ -\Delta_{\text{cut}}^2 & \Delta_{\text{cut}}^2 \end{pmatrix}$$

- **yield uncertainty** $\Delta^y \equiv \Delta_{FO}$: Fixed-order scale uncertainty in σ_{total}
- **migration uncertainty** $\Delta_{\text{cut}} \equiv \Delta_{\text{resum}}$: Resummation uncertainty induced by the binning cut and drops out in the sum $\sigma_0 + \sigma_{\geq 1}$.

$$\Delta_0^2(\mathcal{T}^{\text{cut}}) = \Delta_{FO}^2(\mathcal{T}^{\text{cut}}) + \Delta_{\text{resum}}^2(\mathcal{T}^{\text{cut}})$$

Resummation Scales and Fixed-Order Uncertainty



► Central profiles:

$$\mu_H = -i\mu_{\text{FO}}, \quad \mu_{\text{ns}} = \mu_{\text{FO}},$$

$$\mu_S(\mathcal{T}^{\text{cut}}) = \mu_{\text{FO}} f_{\text{run}}(\mathcal{T}^{\text{cut}}/m_H),$$

$$\mu_B(\mathcal{T}^{\text{cut}}) = \sqrt{\mu_S(\mathcal{T}^{\text{cut}})\mu_{\text{FO}}} = \mu_{\text{FO}} \sqrt{f_{\text{run}}(\mathcal{T}^{\text{cut}}/m_H)}$$

► Fixed Order Uncertainty Δ_{FO} :

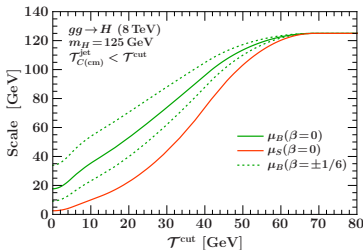
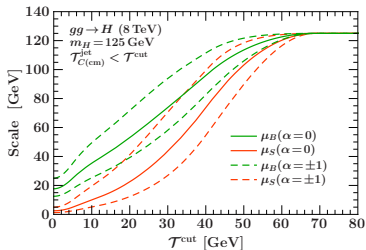
- Maximum of collective variation of all scales by a factor of 2 keeping scale ratios fixed.
- Reproduces the inclusive cross section uncertainty for large \mathcal{T}^{cut} .

Resummation Uncertainty

- ▶ Varying argument of logs estimates their size and missing higher log terms.
- ▶ These variations are smoothly turned off at large \mathcal{T}^{cut} where the resummation turns off.

$$\mu_S^{\text{vary}}(X, \alpha) = f_{\text{vary}}^\alpha(X) \mu_S(X) = \mu_{\text{FO}} f_{\text{vary}}^\alpha(X) f_{\text{run}}(X)$$

$$\mu_B^{\text{vary}}(X, \alpha, \beta) = \mu_S^{\text{vary}}(X, \alpha)^{1/2-\beta} \mu_{\text{FO}}^{1/2+\beta} = \mu_{\text{FO}} [f_{\text{vary}}^\alpha(X) f_{\text{run}}(X)]^{1/2-\beta}$$



Resummation Uncertainties

Choosing α and β for μ_S and μ_B variation (for SCET-I type observables)

- ▶ Retain natural scale hierarchy: $\mu_{\text{FO}} \sim \mu_H \gg \mu_B = \sqrt{\mu_S \mu_{\text{FO}}} \gg \mu_S$.
- ▶ A systematic way to vary μ_B and μ_S without double counting.

$$\mu_S^{\text{vary}}(X, \alpha) = \mu_{\text{FO}} f_{\text{vary}}^\alpha(X) f_{\text{run}}(X), \quad \mu_B^{\text{vary}}(X, \alpha, \beta) = \mu_{\text{FO}} [f_{\text{vary}}^\alpha(X) f_{\text{run}}(X)]^{1/2-\beta}$$

- ▶ Independent scale ratios entering the resummed logs are

$$\frac{\mu_B^2}{\mu_H^2} \sim \frac{\mathcal{T}^{\text{cut}}}{m_H}, \quad \frac{\mu_S^2}{\mu_B^2} \sim \frac{\mathcal{T}^{\text{cut}}}{m_H}$$

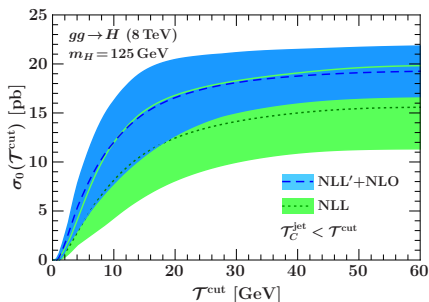
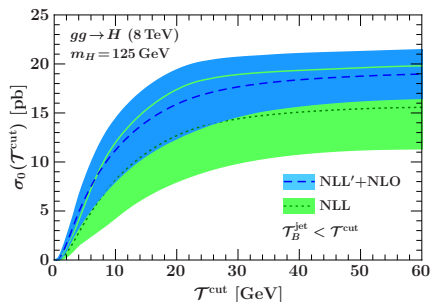
- ▶ α variation for both μ_S and μ_B : Equal changes in the log ratios.

$$\ln \frac{\mu_B^2}{\mu_H^2} \rightarrow \ln f_{\text{vary}}^\alpha + \ln \frac{\mu_B^2}{\mu_H^2}, \quad \ln \frac{\mu_S^2}{\mu_B^2} \rightarrow \ln f_{\text{vary}}^\alpha + \ln \frac{\mu_S^2}{\mu_B^2}.$$

- ▶ β variation: Equal magnitude opposite sign changes in the log ratios.

$$\ln \frac{\mu_B^2}{\mu_H^2} \rightarrow (1 - 2\beta) \ln \frac{\mu_B^2}{\mu_H^2}, \quad \ln \frac{\mu_S^2}{\mu_B^2} \rightarrow (1 + 2\beta) \ln \frac{\mu_S^2}{\mu_B^2}.$$

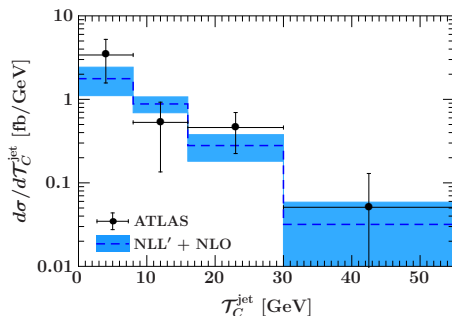
NLL'+NLO results for H+0-jet



- ▶ Good convergence and reduced uncertainties going from NLL to NLL'+NLO.
- ▶ At $\mathcal{T}^{\text{cut}} \sim 25$ GeV, 20% uncertainty for NLL'+NLO, comparable to the precision for p_T^{jet} .
- ▶ Expect significant reduction in uncertainties going to NNLL'+NNLO.

Comparison to ATLAS data

- ▶ The resummed results for different jet veto observables can be directly compared to the ATLAS measurements.
- ▶ Comparison of our NLL'+NLO predictions in bins of $\mathcal{T}_C^{\text{jet}}$ with the recent ATLAS measurement in $H \rightarrow \gamma\gamma$ channel.
- ▶ Several correction factors implemented given in ATLAS pub:1407.4222 for $H \rightarrow \gamma\gamma$ branching ratio, photon isolation efficiency, diphoton kinematic acceptance.



- ▶ **Resummation of jet veto logs**: important for accurate cross section predictions.
- ▶ **Rapidity dependent jet vetoes**:
 - 1 A natural and efficient central jet veto, Resummable to same level of precision as p_T^{jet} .
 - 2 Provide complementary way to divide phase space into exclusive jet bins.
→ Motivation to measure rapidity-dependent jet-vetoes in other processes.
- ▶ Determined **NLL'+NLO** resummed results for $\mathcal{T}_{Bcm}^{\text{jet}}$, $\mathcal{T}_B^{\text{jet}}$, $\mathcal{T}_{Ccm}^{\text{jet}}$ and $\mathcal{T}_C^{\text{jet}}$ with robust uncertainty estimates. Comparison with ATLAS data!
- ▶ Next step is **resummation to NNLL'+NNLO order** :
 - 1 Beam and soft functions at 2 loops + **In R^2 clustering corrections (W.I.P)**
 - 2 Extension to full NNLL'+NNLO.

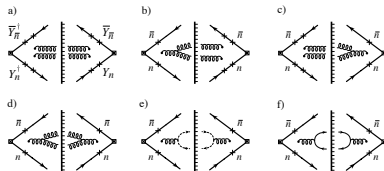
Outlook: Clustering correction to soft function

- ▶ Dijet hemisphere soft function

$$\Delta S^{(2)}(k_1, k_2) = \int \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} A_j(k_1, k_2) \Delta M^{\text{jet}(2)}(k_1, k_2) C(k_1) C(k_2)$$

$$\int \frac{d^d k_1}{(2\pi)^d} \frac{d^d k_2}{(2\pi)^d} = \int_0^\infty dk_1^+ dk_2^+ dk_1^- dk_2^- (k_1^+ k_1^- k_2^+ k_2^-)^{-\epsilon} \int_0^\pi d\Delta\phi \sin^{-2\epsilon} \Delta\phi$$

- ▶ For the clustering corrections only non-abelian matrix elements considered.
[Hornig, Lee, Stewart, Walsh, Zuberi '11]



- ▶ For a \mathcal{T}_B veto, the full measurement function can be written as
[following Tackmann, Walsh, Zuberi '12]

$$M_i^{\text{jet}}(\mathcal{T}^{\text{cut}}) = M_i(\mathcal{T}^{\text{cut}}) + \Delta M_i^{\text{jet}}(\mathcal{T}^{\text{cut}})$$

Outlook: Clustering correction to soft function

Inclusive \mathcal{T}_B measurement

$$M_i(\mathcal{T}^{\text{cut}}) = \theta(\sum_j \mathcal{T}_j < \mathcal{T}^{\text{cut}})$$

Measurement function for the clustering correction relative to M_i

$$\begin{aligned}\Delta M^{\text{jet}}(\mathcal{T}^{\text{cut}}) &= 2[\theta(\Delta R_{12} < \Delta R)\theta(\mathcal{T}_B^{\text{jet}} < \mathcal{T}^{\text{cut}}) + \theta(\Delta R_{12} > R)\theta(\mathcal{T}_{B1} < \mathcal{T}^{\text{cut}})\theta(\mathcal{T}_{B2} < \mathcal{T}^{\text{cut}}) \\ &\quad - \theta(\mathcal{T}_{B1} + \mathcal{T}_{B2} < \mathcal{T}^{\text{cut}})] \\ &= 2\theta(\Delta R_{12} > R)[\theta(\mathcal{T}_{B1} < \mathcal{T}^{\text{cut}})\theta(\mathcal{T}_{B2} < \mathcal{T}^{\text{cut}}) - \theta(\mathcal{T}_{B1} + \mathcal{T}_{B2} < \mathcal{T}^{\text{cut}})]\end{aligned}$$

The variables we use are

$$\begin{aligned}y_t &= \frac{1}{2}(y_1 + y_2) & \Delta y &= y_1 - y_2 & z &= \frac{\mathcal{T}_{B1}}{\mathcal{T}_{B1} + \mathcal{T}_{B2}} = \frac{k_1^+}{k_1^+ + k_2^+} \\ \mathcal{T}_T &= \mathcal{T}_{B1} + \mathcal{T}_{B2} & \cos \Delta\phi &= \frac{k_1^\perp \cdot k_2^\perp}{|k_1^\perp| |k_2^\perp|} = \frac{1/2(k_1^+ k_2^- + k_2^+ k_1^-) - k_1 \cdot k_2}{\sqrt{k_1^+ k_1^- k_2^+ k_2^-}}\end{aligned}$$

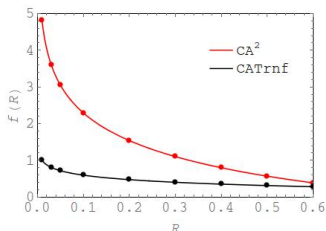
$$\Delta S^{(2)}(\mathcal{T}^{\text{cut}}) = \frac{8}{(16\pi^2)^2} \int \frac{\mathcal{T}_T^{-1-4\epsilon}}{z(1-z)} \mathcal{T}_T^4 z^2 (1-z)^2 (z^2(1-z)^2)^{-\epsilon} e^{-4\epsilon y_t} A_j \Delta M^{\text{jet}(2)} e^{4y_t} d\mathcal{T}_T dz dy_t d\Delta y$$

$$\begin{aligned}\Delta S^{(2)}(\mathcal{T}^{\text{cut}}) &= \frac{-8}{(16\pi^2)^2} \frac{1}{4\epsilon} \left(\frac{\mu}{\mathcal{T}^{\text{cut}}}\right)^{4\epsilon} \int_0^1 dz \int_{-\infty}^{\infty} d\Delta y \int_0^\pi d\Delta\phi \sin^{-2\epsilon} \Delta\phi \frac{\log[\max(z, 1-z)]}{z(1-z)} (z(1-z))^{-2\epsilon} \\ &\quad \theta(\Delta R_{12} > \Delta R) \mathcal{T}_T^4 z^2 (1-z)^2 A_j\end{aligned}$$

Outlook: Clustering correction to soft function

- ▶ Remaining integrals finite and result of the form: $a \log R + b + O(R)$.
- ▶ To extract $\log R$, express A_i in the small $\Delta R = \sqrt{\Delta y^2 + \Delta \phi^2}$ limit as A_i^R .
Write $A = (A - A^R) + A^R$
- ▶ To determine a : Analytic integration of A^R .
Extract b numerically by integrating $A - A^R$.

$$\Delta S^{(2)}(\mathcal{T}^{\text{cut}}) = \left(\frac{\alpha_s}{\pi}\right)^2 \left(\frac{\mu}{\mathcal{T}^{\text{cut}}}\right)^{4\epsilon} \left[\frac{1}{4\epsilon} \left[C_A^2 \left\{ \frac{1}{18} (131 - 12\pi^2 - 132 \log 2) \log R - 0.937 + 0.652 R^2 + \dots \right\} \right. \right. \\ \left. \left. + C_A T_R n_f \left\{ \frac{-1}{9} (23 - 24 \log 2) \log R + 0.747 + 0.019 R^2 + \dots \right\} \right] + C_A^2 \{-4.254 \log R \right. \\ \left. + 1.096 \log^2 R + 1.713 + O(R^2)\} + C_A T_R n_f \{-0.451 \log R + 0.177 \log^2 R + 0.184 + O(R^2)\} \right]$$

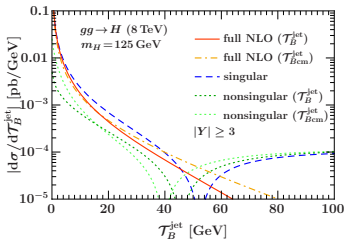
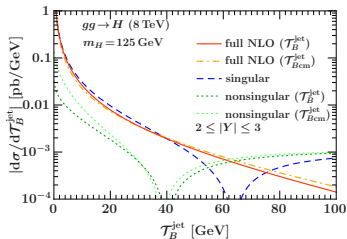


Analogous but more tedious calculation for $\mathcal{T}_C^{\text{jet}}$ gives the same anom. dim. as $\mathcal{T}_B^{\text{jet}}$

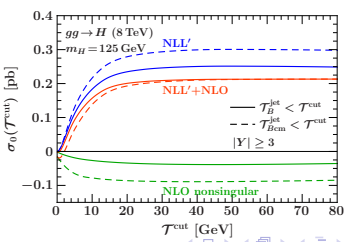
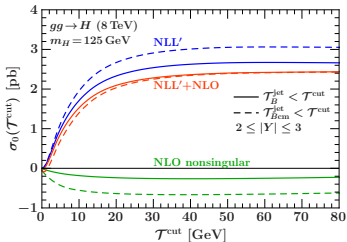
BackUp slides

NLL'+NLO results for H+0-jet

- ▶ Differential singular, nonsingular and full fixed order cross sections for $\mathcal{T}_B^{\text{jet}}$ in different Y bins.



- ▶ Cumulant NLL' and nonsingular for $\mathcal{T}_B^{\text{jet}}$ in different Y bins.



Soft function and Resummation uncertainty

- ▶ Soft function for $\mathcal{T}_B^{\text{jet}}$ and $\mathcal{T}_C^{\text{jet}}$:

$$S_{gg}^C(\mathcal{T}^{\text{cut}}, \mu) = 1 + \frac{\alpha_s(\mu) C_A}{\pi} \left(-2 \ln^2 \frac{\mathcal{T}^{\text{cut}}}{\mu} + \frac{\pi^2}{4} \right)$$

$$S_{gg}^B(\mathcal{T}^{\text{cut}}, \mu) = 1 + \frac{\alpha_s(\mu) C_A}{\pi} \left(-2 \ln^2 \frac{\mathcal{T}^{\text{cut}}}{\mu} + \frac{\pi^2}{12} \right),$$

$$f_{\text{run}}(x) = \begin{cases} x_0 \left(1 + \frac{x}{4x_0} \right) & x \leq 2x_0 \\ x & 2x_0 \leq x \leq x_1 \\ x + \frac{(2-x_2-x_3)(x-x_1)^2}{2(x_2-x_1)(x_3-x_1)} & x_1 \leq x \leq x_2 \\ 1 - \frac{(2-x_1-x_2)(x-x_3)^2}{2(x_3-x_1)(x_3-x_2)} & x_2 \leq x \leq x_3 \\ 1 & x_3 \leq x \end{cases}$$

$$f_{\text{vary}}(x) = \begin{cases} 2(1 - x^2/x_3^2) & 0 \leq x \leq x_3/2 \\ 1 + 2(1 - x/x_3)^2 & x_3/2 \leq x \leq x_3 \\ 1 & x_3 \leq x \end{cases},$$

Clustering Correction

- ▶ The variables we use are

$$y_1 = \frac{1}{2} \log \frac{k_1^-}{k_1^+}$$

$$y_2 = \frac{1}{2} \log \frac{k_2^-}{k_2^+}$$

$$y_t = \frac{1}{2}(y_1 + y_2) \quad \Delta y = y_1 - y_2 \quad z = \frac{\mathcal{T}_{B1}}{\mathcal{T}_{B1} + \mathcal{T}_{B2}} = \frac{k_1^+}{k_1^+ + k_2^+}$$

$$\mathcal{T}_T = \mathcal{T}_{B1} + \mathcal{T}_{B2} \quad \cos \Delta\phi = \frac{k_1^\perp \cdot k_2^\perp}{|k_1^\perp| |k_2^\perp|} = \frac{1/2(k_1^+ k_2^- + k_2^+ k_1^-) - k_1 \cdot k_2}{\sqrt{k_1^+ k_1^- k_2^+ k_2^-}}$$

$$A_f^R(\mathcal{T}^{cut}) = \frac{4g^4 C_A \mathcal{T}_R n_f \mu^{4\epsilon} e^{-4y_t}}{\mathcal{T}_T^4 z^2 (1-z)^2} \frac{4}{\Delta R^4} \left[\frac{z(1-z)}{2} (\Delta R^2 - 4z(1-z)\Delta y^2) \right]$$

$$A_A^R(\mathcal{T}^{cut}) = \frac{4g^4 C_A^2 \mu^{4\epsilon} e^{-4y_t}}{\mathcal{T}_T^4 z^2 (1-z)^2} \frac{2}{\Delta R^2} \left[(1-z+z^2) - z(1-z) + \frac{2z^2(1-z)^2 \Delta y^2 (1-\epsilon)}{\Delta R^2} \right]$$

$$A_f(\mathcal{T}^{cut}) = \frac{4g^4 C_A \mathcal{T}_R n_f \mu^{4\epsilon}}{\mathcal{T}_T^4 z^2 (1-z)^2} \left[\frac{1}{(\cosh \Delta y - \cos \Delta\phi)(\cosh \Delta y - (1-2z) \sinh \Delta y)} \right]^2 \frac{z(1-z)}{2} \left[1 - 2 \cos \Delta\phi \cosh \Delta y + 2(1-2z) \cos \Delta\phi \sinh \Delta y + (\cosh \Delta y - (1-2z) \sinh \Delta y)^2 \right]$$

$$A_A(\mathcal{T}^{cut}) = \frac{4g^4 C_A^2 \mu^{4\epsilon}}{\mathcal{T}_T^4 z^2 (1-z)^2} \frac{1}{(\cosh \Delta y - \cos \Delta\phi)(\cosh \Delta y - (1-2z) \sinh \Delta y)} \left[(1-z+z^2) \cos \Delta\phi \cosh \Delta y - (1-2z) \cos \Delta y \sinh \Delta\phi - z(1-z) + \frac{(1-\epsilon)z^2(1-z)^2 \sinh \Delta y^2}{(\cosh \Delta y - \cos \Delta\phi)(\cosh \Delta y - (1-2z) \sinh \Delta y)} \right]$$